

Mauro Mezzetto  
*Istituto Nazionale di Fisica Nucleare, Sezione di Padova*

## **“Sensitivities comparison EU/US/Japan”**

- Introduction
- Ingredients for a comparison
  - Signal Efficiency
  - Backgrounds
  - Systematic Errors
  - Degeneracies
- Updated comparison charts

# What future experiments will have to measure

- 1 Leptonic CP violation
  - 2 Mass Hierarchy
  - 3  $\theta_{13}$
  - 4  $\pi/4 - \theta_{23}$  and the octant
- Plus new physics searches through non standard neutrino interactions (NSI) and/or the unitarity of the mixing matrix

The above order following the prescriptions of the CERN strategy group

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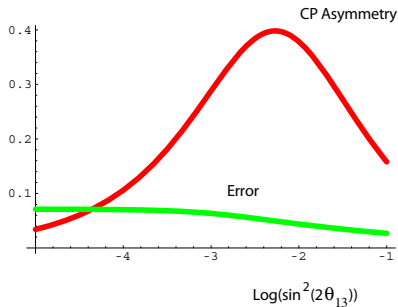
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Neutrino mass model builders would probably order as 3,2,1,4

# Measuring Leptonic CP violation

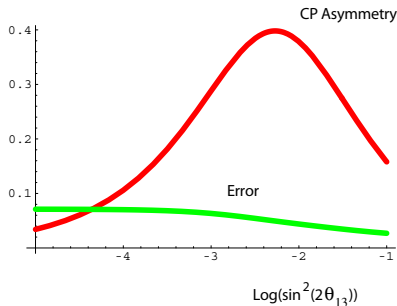
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LCPV asymmetry at the first oscillation maximum,  $\delta = 1$ , Error curve: dependence of the statistical+systematic (2%) computed for a beta beam the fixed energy  $E_{\nu} = 0.4$  GeV,  $L = 130$  km.

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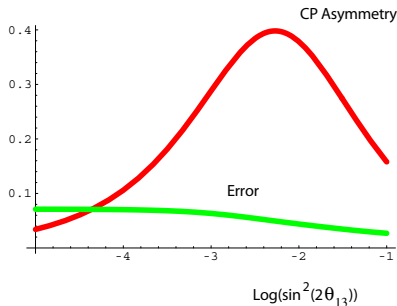


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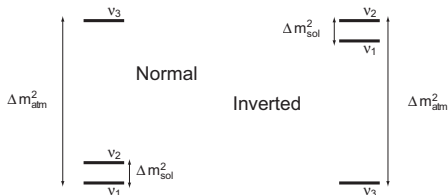


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- **The detection of such asymmetry is an evidence of Leptonic CP violation only in absence of competitive processes (i.e. matter effects, see following slides)  $\Rightarrow$  "short" Long Baseline experiments**
- Statistics and systematics play different roles at different values of  $\theta_{13} \Rightarrow$  impossible to optimize the experiment without a prior knowledge of  $\theta_{13}$
- Contrary to the common belief, the highest values of  $\theta_{13}$  are not the easiest condition for LCPV discovery

# Measuring mass hierarchy

An internal degree of freedom of neutrino masses is the sign of  $\Delta m_{31}^2$ :  $\text{sign}(\Delta m_{31}^2)$ .



This parameter decides how mass eigenstates are coupled to flavor eigenstates with important consequences to direct neutrino mass and double beta decay experiments.

# Neutrino Oscillations in Matter

$$\begin{aligned}P_{\theta_{13}} &= \sin^2(2\theta_{13})\sin^2\theta_{23}^2 \sin^2((\hat{A} - 1)\hat{\Delta})/(\hat{A} - 1)^2; \\p_{\sin \delta} &= \alpha \sin(2\theta_{13})\zeta \sin \delta \sin(L\hat{\Delta}) \sin(\hat{A}\hat{\Delta}) \sin((1 - \hat{A})\hat{\Delta})/((1 - \hat{A})\hat{A}); \\p_{\cos \delta} &= \alpha \sin(2\theta_{13})\zeta \cos \delta \cos \hat{\Delta} \sin(\hat{A}\hat{\Delta}) \sin(1 - \hat{A}\hat{\Delta})/((1 - \hat{A})\hat{A}); \\p_{\text{solar}} &= \alpha^2 \cos^2\theta_{23}^2 \sin^2 2\theta_{12} \sin^2(\hat{A}\hat{\Delta})/\hat{A}^2;\end{aligned}$$

$$\alpha = \text{Abs}(\Delta m_{21}^2/\Delta m_{31}^2); \quad \hat{\Delta} = \frac{L\Delta m_{31}^2}{4E} \quad \zeta = \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{23}$$

$$\hat{A} = \pm a/\Delta m_{31}^2; \quad a = 7.6 \cdot 10^{-5} \rho \cdot E_\nu (\text{GeV}) \quad \rho = \text{matter density (g cm}^{-3}\text{)}$$

The  $\hat{A}$  term changes sign with  $\text{sign}(\Delta m_{13}^2)$

**Matter effects require long “long baselines”**



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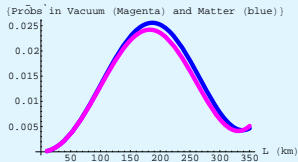
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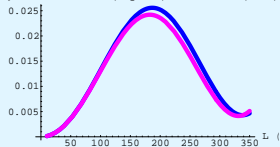
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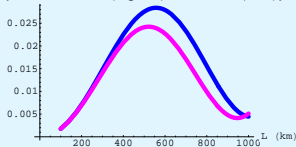
**Matter effects require long “long baselines”**

$$E_\nu = 0.35 \text{ GeV } L \simeq 130 \text{ km} \quad E_\nu = 1 \text{ GeV } L \simeq 500 \text{ km}$$

{Pröbs in Vacuum (Magenta) and Matter (blue)}



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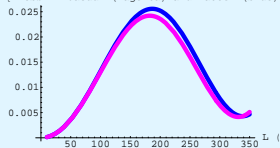
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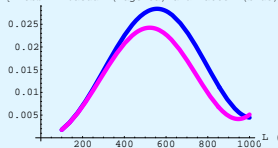
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$$E_\nu = 0.35\text{GeV } L \simeq 130 \text{ km} \quad E_\nu = 1\text{GeV } L \simeq 500 \text{ km} \quad E_\nu = 3\text{GeV } L \simeq 1500 \text{ km}$$

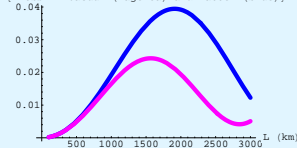
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# What to compare

**Sensitivity:** The highest value of a parameter (say  $\theta_{13}$ ) that can be excluded at a given CL in absence of a signal.

True value = 0; fit value  $\neq$  0.

**Discovery potential:** The smallest value of a parameter that can provide a signal that can't be fitted with a null value at a given CL.

True value  $\neq$  0; fit value = 0.

**Precision:** Fractional value of one std over the parameter (i.e.

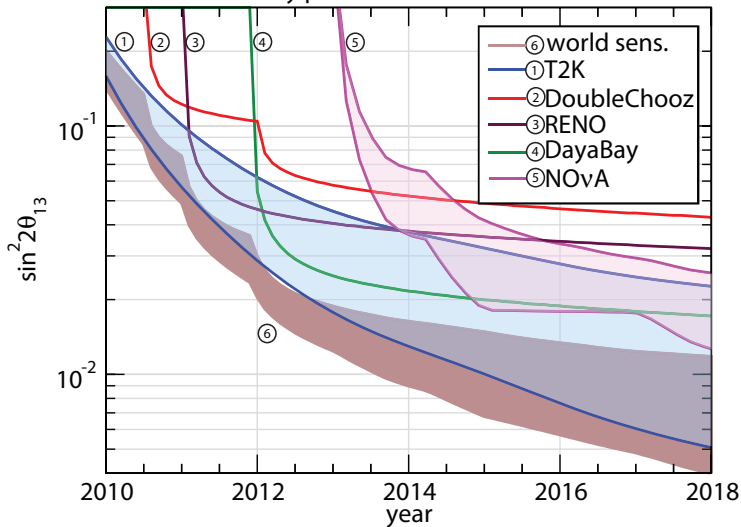
$$\frac{\sigma(\theta_{13})}{\theta_{13}}(\theta_{13}))$$

In the following we will talk about sensitivity

# Status after this generation of LBL experiments: $\theta_{13}$

From M.M. and T. Schwetz, arXiv:1003.5800

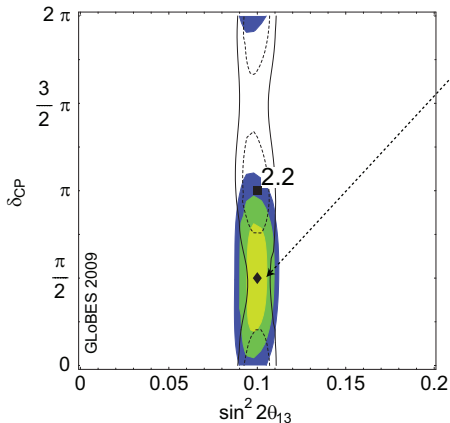
Discovery potential at  $3\sigma$  for NH



# Status after this generation of LBL experiments: CPV

From P. Huber et al., JHEP 0911:044,2009.

T2K + NOvA+Reactors  
after the nominal run

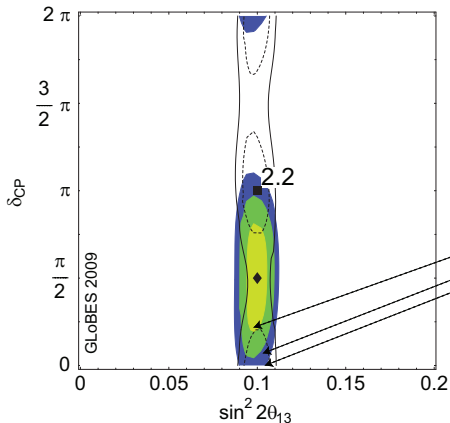


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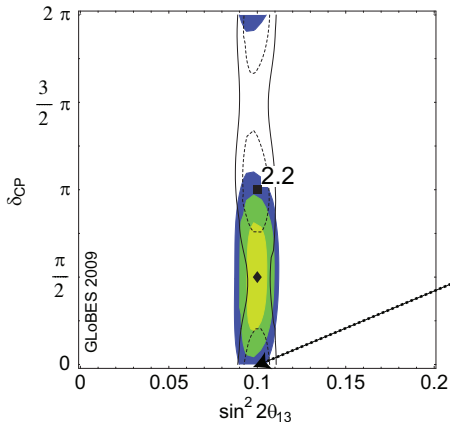
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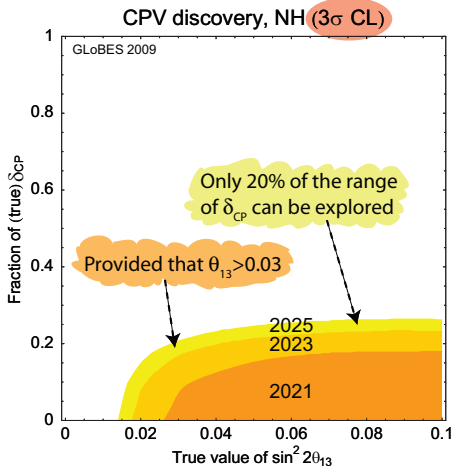
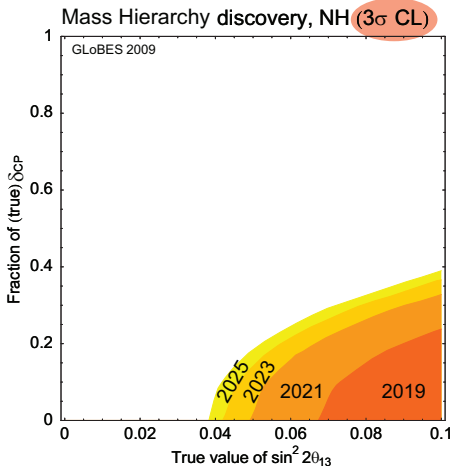
3) Null CP is compatible with data already at  $2\sigma$



# Status after accelerator upgrades

From P. Huber et al., JHEP 0911:044,2009.

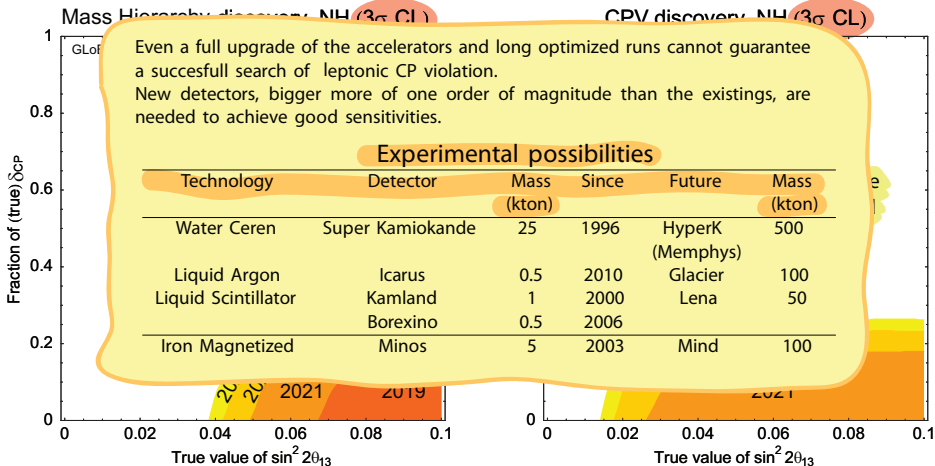
Prediction of sensitivity including a **fully optimized global run** (antineutrinos in T2K and NO $\nu$ A) and **full upgrade of the accelerators**: 1.6 MW at J-PARC and 2.4 MW at FNAL (Project-X)



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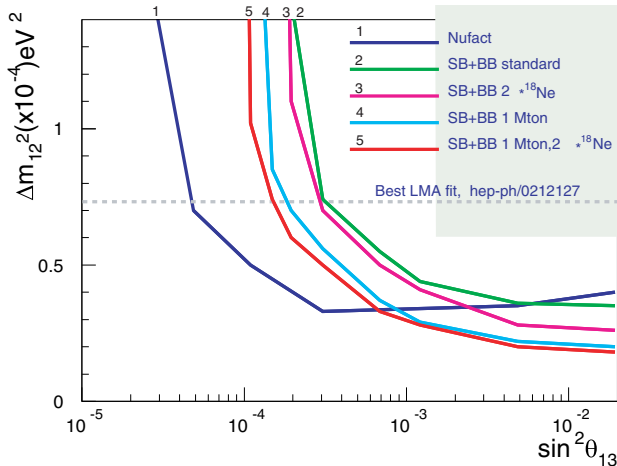
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# Back to 2002

Apollonio et al. hep-ph/0210192 (see also MM hep-ex/0302007)



Soon after SNO and Kamland gave us the good news: LMA is the right solution and  $\sin^2 2\theta_{12}$  is bigger than  $7 \cdot 10^{-5} \text{eV}^2$

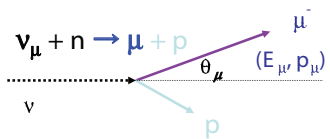
## Things to take care about

- Energy reconstruction
- Background simulation
- Systematic errors
- Degeneracies

# Energy reconstruction for beam neutrinos

The quasi elastic case

Select single ring events and assume they are Quasi Elastic



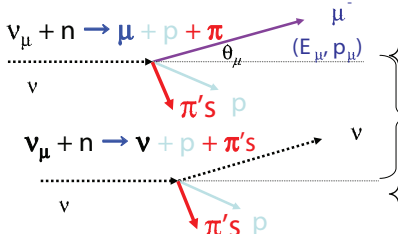
◇ **CC QE**

◇ can reconstruct  $E_\nu \leftarrow (\theta_\mu, p_\mu)$

$$E_\nu^{\text{rec}} = \frac{m_N E_\mu - m_\mu^2/2}{m_N - E_\mu + p_\mu \cos \theta_\mu}$$

$\delta E \sim 60 \text{ MeV} \quad \delta E/E \sim 10\%$

Single ring non Quasi Elastic are badly measured

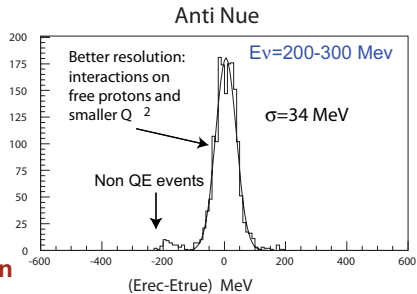
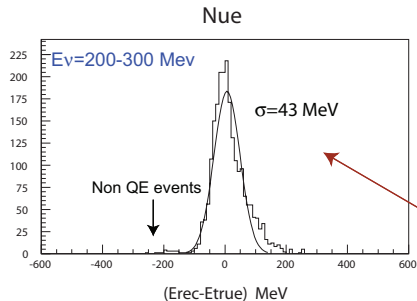


◇ bkg. for  $E_\nu$  measurement

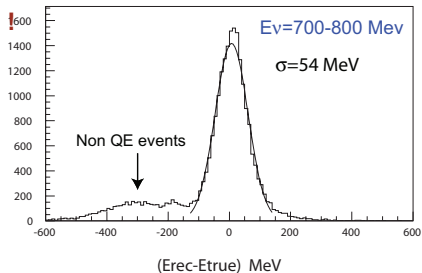
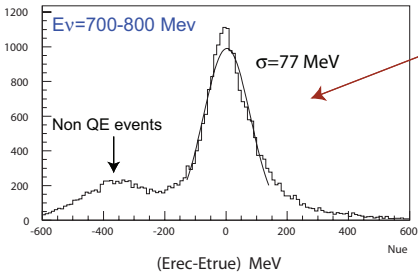
High energy part

◇ bkg. for e-appearance

# Goodness of energy reconstruction

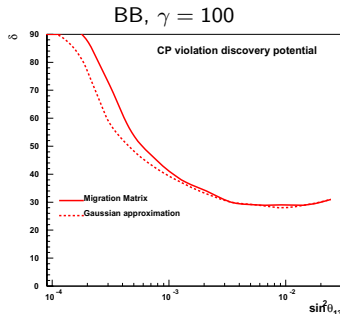
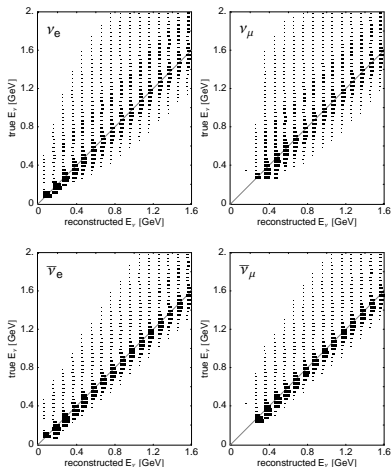


**better resolution at lower energies!**



# Migration Matrixes

A gaussian assumption for energy resolution is a too crude approximation



N.B. DIS event reconstruction requires to precisely measure the hadronic shower  $\rightarrow$  introducing again non-gaussianity of the energy resolution

# Backgrounds

The ultimate limit for CPV and  $\theta_{13}$  sensitivity are background rates and systematic errors.

In a Super Beam setup backgrounds come mainly from intrinsic  $\nu_e$  that look like signal.

For Beta Beams and Nufacts where the intrinsic backgrounds are absent and the background rate small, it is mandatory to compute them as accurately as possible, exploiting a full simulation and a full reconstruction.

Better if migration matrixes are built (most general case) but their energy spectrum is enough for a fair sensitivity computation. (Energy shape is important because it could be different from oscillated events, reducing the impact of the background rate, as in the case of beta beams).

Quite often anyway background rates are just “guessed”: typical case NC backgrounds assumed as a  $10^{-4}$  fraction of CC interactions.

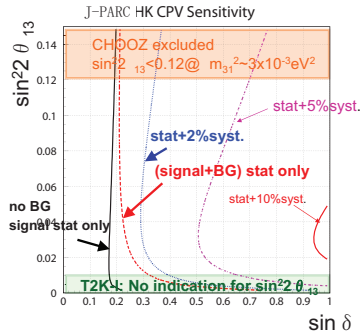
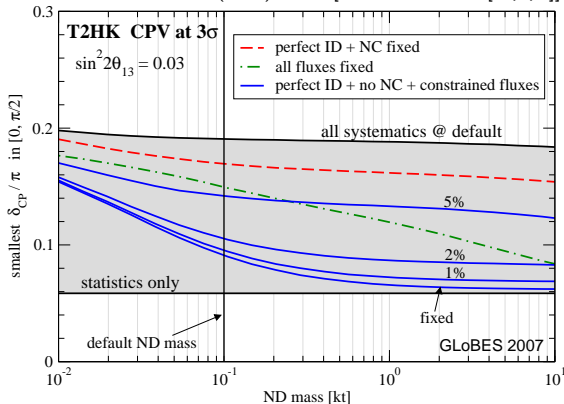


# Systematic errors

They could completely destroy leptonic CP violation.

Default value are often 5% for SuperBeams, 2% for Neutrino Factory and 2-5% for Beta Beams.

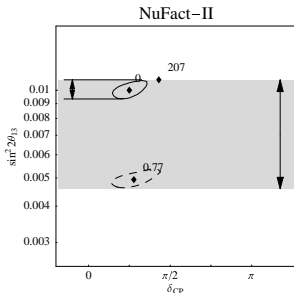
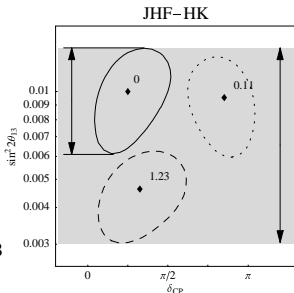
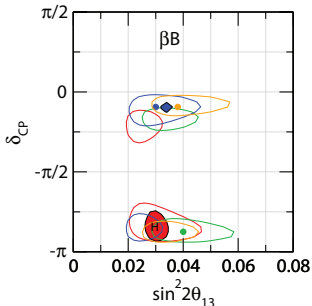
Are they realistic goals? From P. Huber, MM, T. Schwetz, JHEP **0803** (2008) 021 [arXiv:0711.2950 [hep-ph]].



# How to extract $\theta_{13}$ and $\delta_{CP}$

## The problem is not that simple

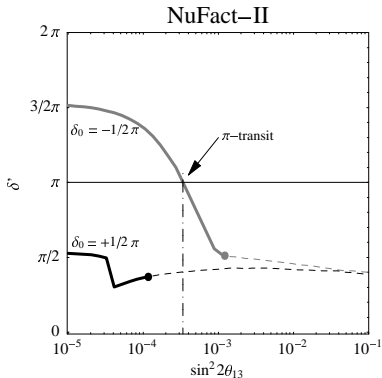
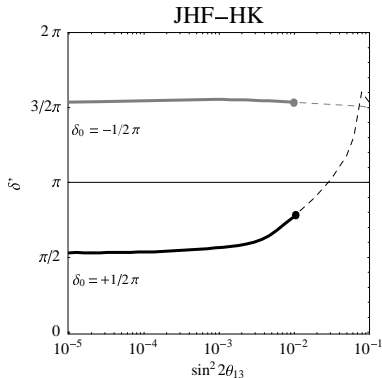
- The 3  $\nu$  oscillation formula contains all the mixing matrix parameters and  $\Delta m^2$ . The parameters already measured do have errors that will influence the extraction of the unknown parameters.
- Several parameters still unknown:  $\theta_{13}$ ,  $\delta_{CP}$ ,  $\text{sign}(\Delta m^2)$  (hierarchy), the octant of  $\theta_{23}$ . Different combinations of the above unknowns can fit the same data:  $\Rightarrow$  **The eightfold degeneracy**



# The $\pi$ -transit problem

From Huber, Lindner, Winter, Nucl. Phys B645:3-48, 2002

The  $\text{sign}(\Delta m_{13}^2)$  degenerate solution could show up at  $\delta_{\text{CP}} = \pi$  destroying any CPV sensitivity. Its position is function of  $\theta_{13}$  and depends from the baseline.



## To solve degeneracies:

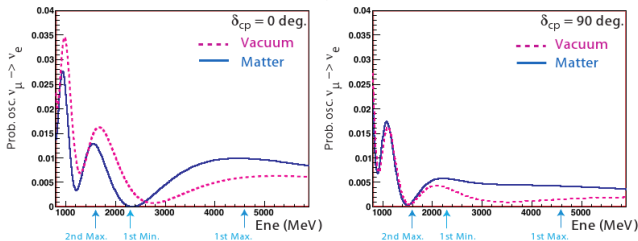
- A single experiment, single channel, can't get rid of degeneracies.
- The combination of different channels in the same detector can solve degeneracies, i.e. first a second oscillation maxima measurement in LBNE or at Okinoshima.
- Different signals in the same detector can also do the job, i.e. beta beams and atmospheric neutrinos.
- A third possibility is to combine the information of different detectors along the same neutrino beam, as exploited by several proposed neutrino factory configurations.
- Of course the combination of the above combinations can also measure all the unknown parameters: can we define an optimal strategy?

# The second oscillation maximum

Studied for DUSEL-LBNE, T2O (T2OK ?), PS2-Laguna

- The functional form of genuine CP and  $\text{sign}(\Delta m_{13}^2)$  terms (both bringing to  $\nu \neq \bar{\nu}$ ) are different at the first and the second oscillation maxima.
- This gives a powerful handle to solve degeneracies.
- It is very challenging on the detector side, covering the 0.5-10 GeV energy range with good background rejection.
- Furthermore the longer distance brings to better sensitivities about mass hierarchy.
- The (high) price to pay is a factor 9 reduction on neutrino fluxes (the second maximum is three times more distant from the first), bringing to reduced sensitivities on  $\theta_{13}$  and CPV.

From A. Rubbia, arXiv:1003.1921



# The synergy with atmospheric neutrinos

**P. Huber et al., Phys. Rev. D 71, 053006 (2005):** Combining Long Baseline data with atmospheric neutrinos (that come for free in the megaton detector):

- Degeneracies can be canceled, allowing for better performances in  $\theta_{13}$  and LCPV searches
- The neutrino mass hierarchy can be measured
- The  $\theta_{23}$  octant can be determined.

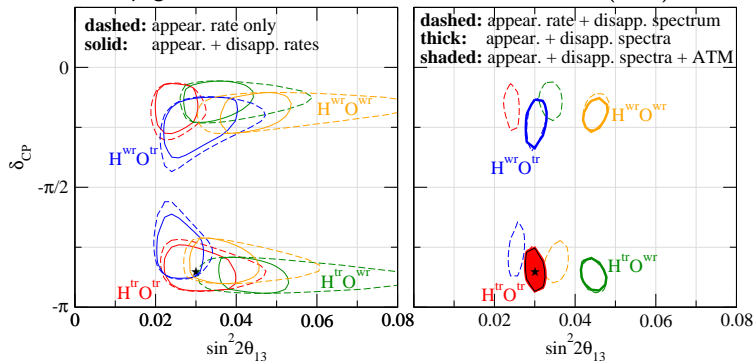
The main reasons are:

- **Octant** e-like events in the Sub-GeV data is  $\propto \cos^2 \theta_{23}$
- **Sign** e-like events in the Multi-GeV data, thanks to matter effects, especially for zenith angles corresponding to neutrino trajectories crossing the mantle and core where a resonantly enhancement occurs.

**NOTE:** LBL and atmospheric neutrinos are a true synergy. They add to each other much more than a simple gain in statistics. Atmospheric neutrinos alone could not measure the hierarchy, the octant,  $\theta_{13}$  and LCPV. While the Beta Beam at short baselines could not measure the hierarchy as well as the octant.

# Degeneracy removal: SPL

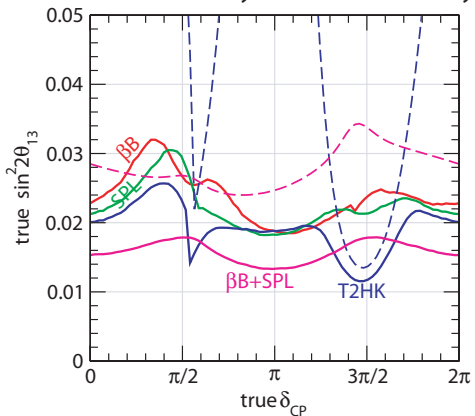
J.E.Campagne, M.Maltoni, M.M., T.Schwetz, JHEP **0704** (2007) 003



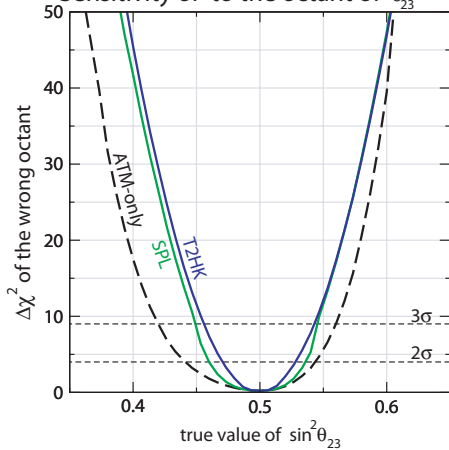
Resolving degeneracies in SPL by successively using the appearance rate measurement, disappearance channel rate and spectrum, spectral information in the appearance channel, and atmospheric neutrinos. Allowed regions in  $\sin^2 2\theta_{13}$  and  $\delta_{CP}$  are shown at 95% CL, and  $H^{tr/wr}(O^{tr/wr})$  refers to solutions with the true/wrong mass hierarchy (octant of  $\theta_{23}$ ). The true parameter values are  $\delta_{CP} = -0.85\pi$ ,  $\sin^2 2\theta_{13} = 0.03$ ,  $\sin^2 \theta_{23} = 0.6$ . The running time is  $(2\nu + 8\bar{\nu})$  yrs.

# Beta Beam plus atmo: determining mass hierarchy and the octant

2 $\sigma$  sensitivity to normal hierarchy



Sensitivity of to the octant of  $\theta_{23}$





# Sensitivity Comparison

Based to arXiv:1005.3146, the EuroNu midterm physics report

**WBB:** Fermilab to Duse, 1 MW for  $\nu$  running, proton energy: 120 GeV, 2 MW for  $\bar{\nu}$  running (5+5 yr), 100 kton liquid argon detector, according to Barger et al, Phys. Rev. D76:053005, 2007 (hep-ph/0703029). This setup is different from the proposed LBNE experiment.

**T2KK:** J-Parc  $\nu$  beam running at 4 MW. 270 kton WC detector at Kamioka (295 km) and 270 kton WC detector in Korea (1050 km), Barger et al, Phys. Rev. D76:053005, 2007 (hep-ph/0703029).

**PS2-Slanic** CERN-PS2 SuperBeam fired to 100 kton LAr detector at Slanic, as computed by A. Rubbia, arXiv:1003.1921.

**SPL:** Neutrino beam from CERN-SPL running at 3.5 GeV, 4 MW. 440 kton WC detector at Frejus (130 km). Campagne et al. JHEP **0704** (2007) 003 (hep-ph/0603172).

**Beta Beam  $\gamma = 100$**  Eurisol Beta Beam to Frejus (440 kton WC detector). Campagne et al. JHEP **0704** (2007) 003 (hep-ph/0603172).

**Beta Beam + SPL** The combination of the above two.

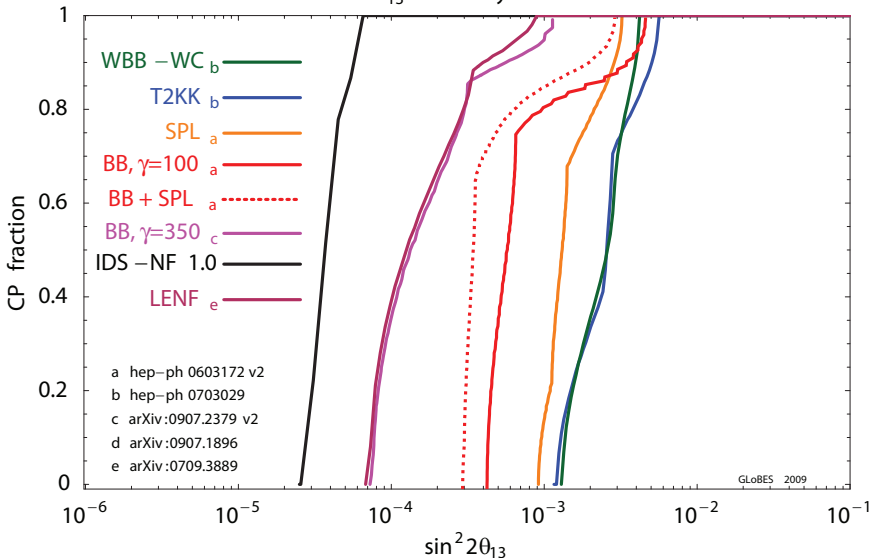
**Beta Beam  $\gamma = 350$**  Beta Beam at  $\gamma = 350$ , running  ${}^6\text{He}$  and  ${}^{18}\text{Ne}$  at the same decay rates as the Eurosol Beta Beam. WC detector of 500 kton at Canfranc (650 km). S. Choubey et al., JHEP 0912:020,2009 (arXiv:0907.2379)

**Low Energy Neutrino Factory (LENF)** Neutrino Factory running at 4.12 GeV delivering  $10^{21}$  muon decays/year for each sign, 30 kton No $\nu$ a like detector, fully magnetized (!) at 1480 km (Fermilab-Henderson mine). A. Bross et al, Phys.Rev.D77:093012,2008. (arXiv:0709.3889)

**IDS 1.0 Neutrino Factory** 25 GeV neutrino factory delivering  $0.5 \cdot 10^{21}$  muon decays/year for each sign, a 50 kton iron magnetized detector and a 10 kton Emulsion Cloud Chamber, at 4000 km and a 50 kton iron magnetized detector at 7500 km.

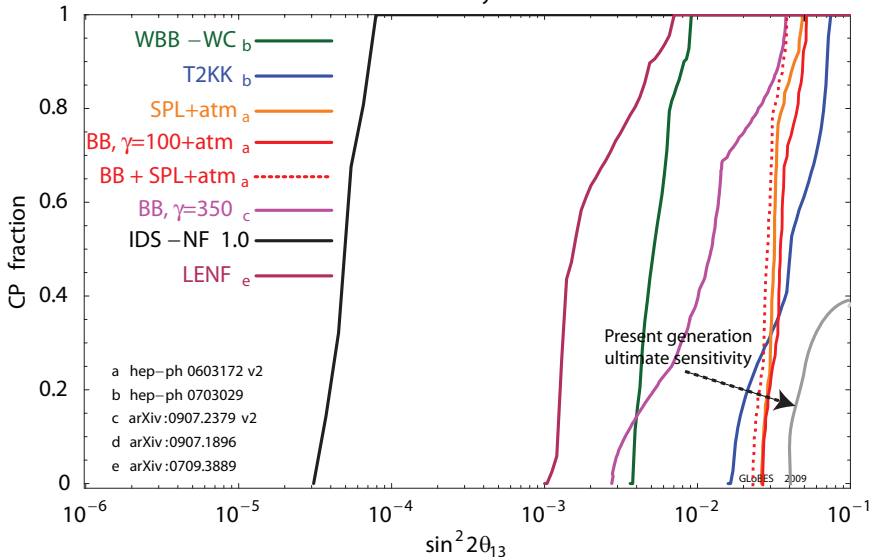
# Sensitivity Comparison: $\theta_{13}$

Elaborated from arXiv:1005.3146  
 $\sin^2 2\theta_{13}$  discovery at  $3\sigma$  CL



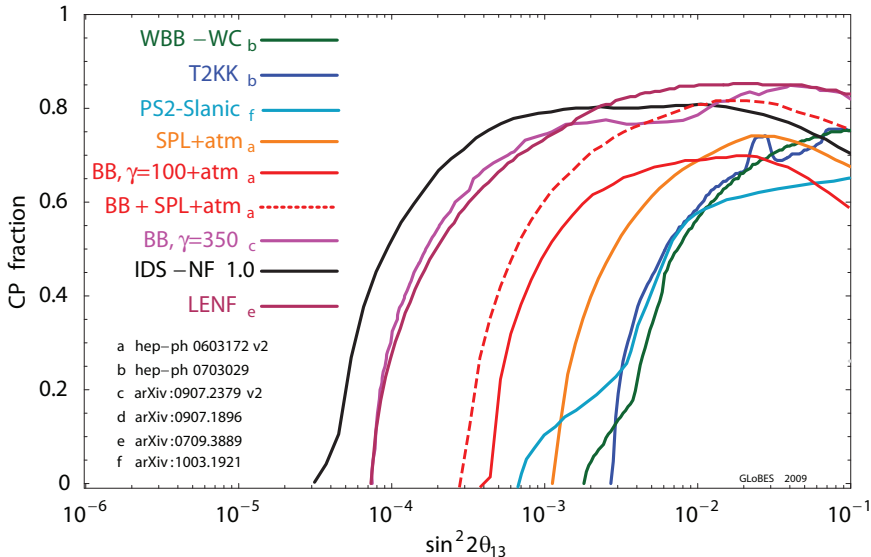
# Sensitivity Comparison: $\text{sign}(\Delta m_{13}^2)$

Elaborated from arXiv:1005.3146  
Mass hierarchy at  $3\sigma$  CL



# Sensitivity Comparison: LCPV

Elaborated from arXiv:1005.3146  
CP violation at  $3\sigma$  CL



# Conclusions

In absence of any information about  $\theta_{13}$  it is impossible to define an optimal strategy.

I'm personally skeptical about optimal strategies and I'm more in favor of realistic, pragmatic approaches. Keeping in mind that the goal for the future is to discover Leptonic CP Violation.

Difficult exercise to understand what is realistic at CERN.

If  $\theta_{13}$  is not within the reach of T2K and the reactors we have to plan Beta Beams or Neutrino Factories. This doesn't exclude a first stage with a Super Beam.