

SM precision measurements at the LHC

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on behalf of the
ATLAS, CMS, LHCb Collaborations

Moriond EW
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Introduction

- Exceptional detector performance, new analysis methodologies, and unprecedented higher-order theoretical calculations have made SM precision measurements a flagship of the LHC physics programme

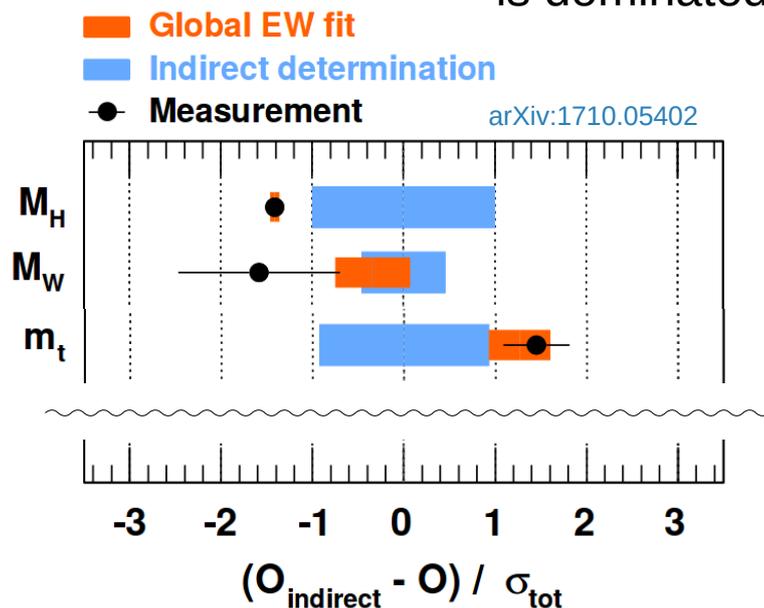
Presenting today a selection of recent SM measurements at the LHC

- ATLAS W-boson angular coefficients [arXiv:2509.13759](#)
- CMS boosted m_W [SMP-24-012](#)
- LHCb m_W from unfolded cross sections [arXiv:2509.18817](#)
- ATLAS Inclusive dijets cross sections [arXiv:2512.19073](#)
- CMS α_s from inclusive jet production [PLB 868 \(2025\) 139651](#)



Motivation for W measurements

The global fit of the electroweak observables is dominated by the measurement of m_W



	Measurement	SM Prediction (*)
m_H	125.1 ± 0.1	104.5 ± 21.7
m_t	172.56 ± 0.31	174.4 ± 2.1
m_W	80.369 ± 0.013	80.356 ± 0.006

(*) PDG 2025

The measurements of the Higgs and top-quark masses are currently more precise than their indirect determination from the global fit of the electroweak observables



Improving precision will not increase sensitivity to new physics

Indirect determination of m_W (± 6 MeV) is more precise than the experimental measurement

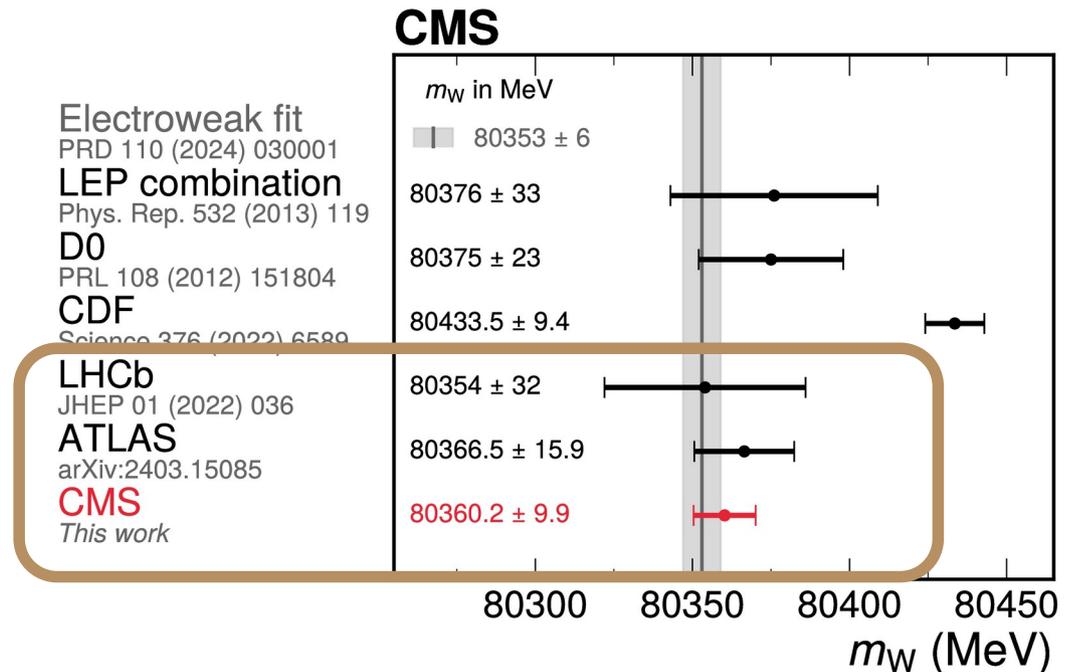
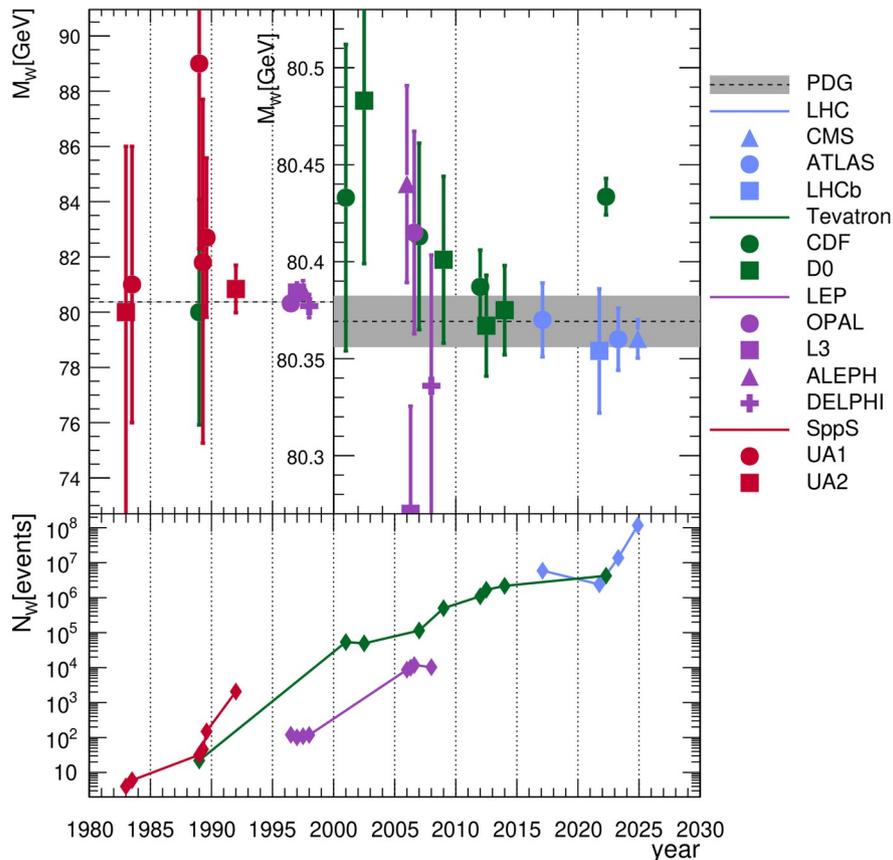
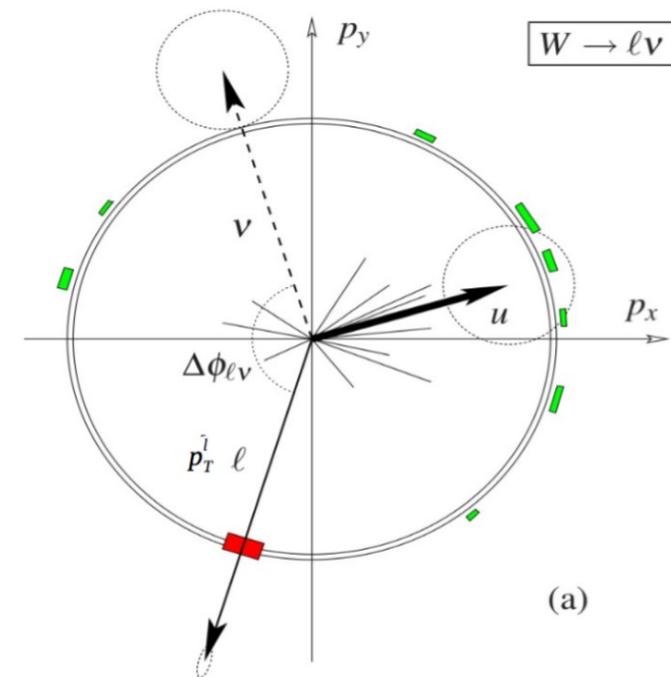


Call for $\delta m_W^{\text{exp}} \sim 5$ MeV

The W mass is nowadays the crucial measurement to improve the sensitivity of the global EW fits to new physics

W-boson measurements at the LHC

- Measurements of $W \rightarrow \ell\nu$ final states at the LHC are complicated by neutrinos which are undetected. Measurements rely either on the indirect reconstruction of the neutrino p_T from momentum balance, or on the kinematic properties of the charged lepton (electron or muon)
- The measurement of the W-boson mass is a milestone of the LHC physics programme. It requires a precise modelling of the W-boson production and decay process, which benefits from auxiliary measurements of W-boson final states



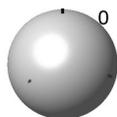
W-boson angular variables

- The $W \rightarrow \ell\nu$ cross section can be expressed through the factorisation of the production dynamic and the decay kinematic properties of the dilepton system

$$\frac{d^3\sigma^{U+L}}{dp_T dy dm} \left(1 + \cos^2\theta + \sum_{i=0}^7 A_i(y, p_T, m) P_i(\cos\theta, \phi) \right)$$

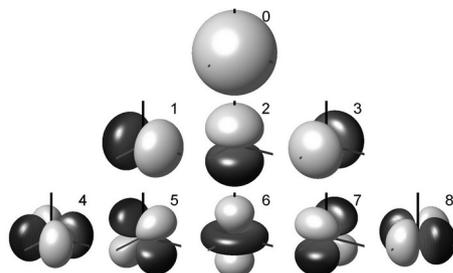
- Decomposition of $(\cos\theta, \phi)$ into 9 helicity cross sections \rightarrow basis of spherical harmonics

Why 9?



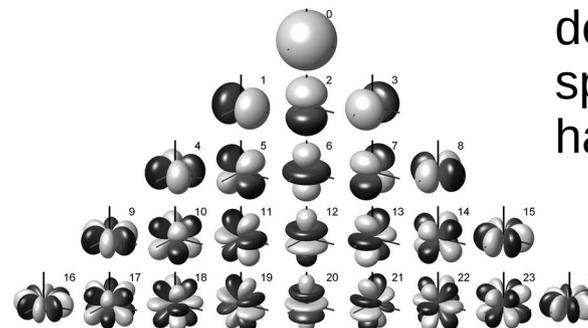
Spin 0 (Higgs)

$$\begin{aligned} \ell &= 0 \\ n &= 1 \end{aligned}$$



Spin 1 (W, Z, γ^*)

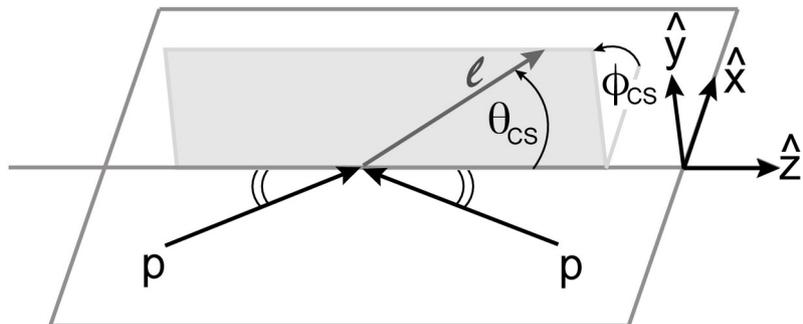
$$\begin{aligned} \ell &\leq 2 \\ n &= 1+3+5 = 9 \end{aligned}$$



Spin 2

$$\begin{aligned} \ell &\leq 4 \\ n &= 1+3+5+7+9 = 25 \end{aligned}$$

ℓ denotes the degree of the spherical harmonics



- Coefficients defined in the Collins-Soper frame

W-boson angular coefficients

- 0.3 fb⁻¹ of low pile-up data for better recoil resolution
- Exploit the angular variables decomposition to perform a simultaneous p_T measurement of
 - Angular coefficients
 - Unpolarised full-lepton phase space cross sections

$$\frac{d^2\sigma}{dp_T dy} \left(1 + \cos^2\theta + \sum A_i(p_T, y) P_i(\cos\theta, \phi) \right)$$

$$P_0(\cos\theta, \phi) = \frac{1}{2}(1 - 3\cos^2\theta)$$

$$P_1(\cos\theta, \phi) = \sin 2\theta \cos \phi$$

$$P_2(\cos\theta, \phi) = \frac{1}{2} \sin^2\theta \cos 2\phi$$

$$P_3(\cos\theta, \phi) = \sin\theta \cos\phi$$

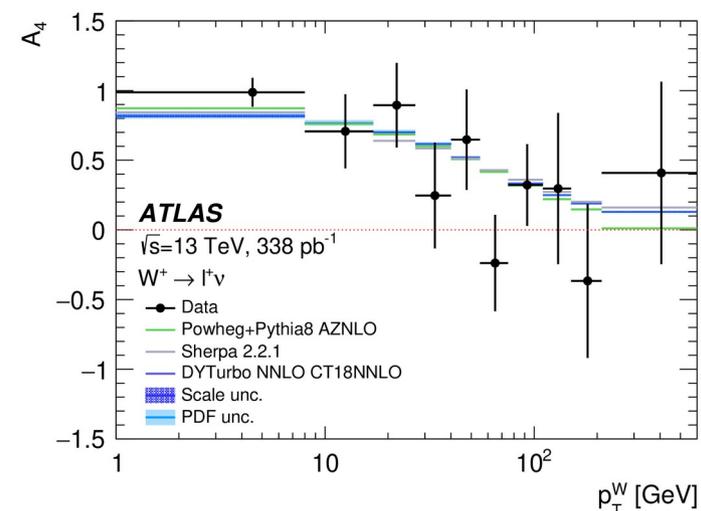
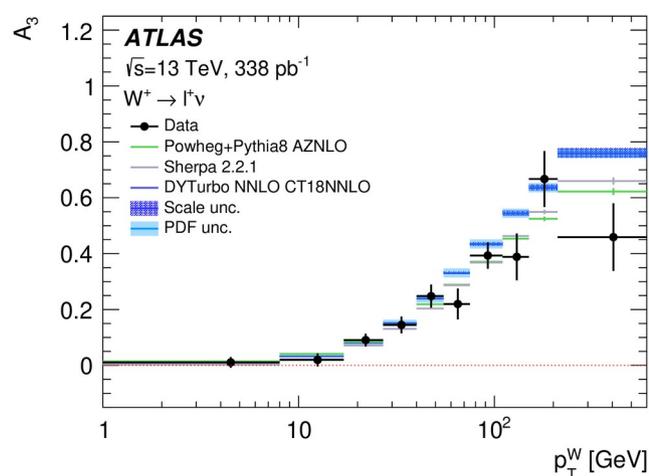
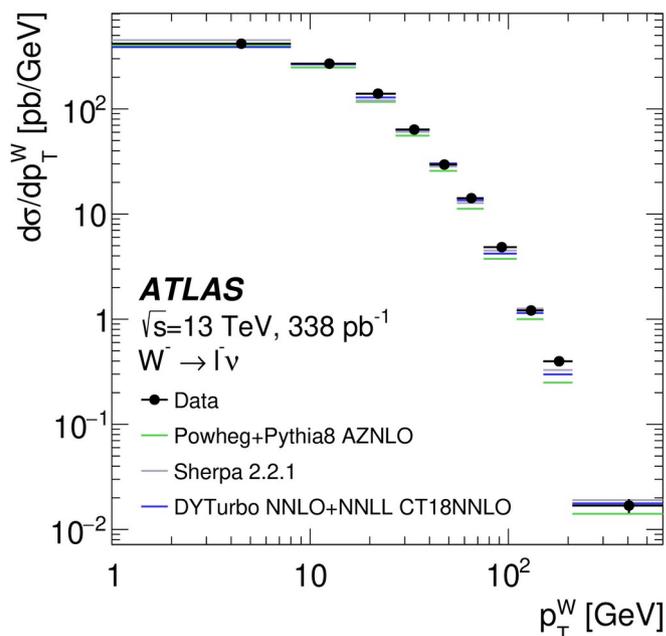
$$P_4(\cos\theta, \phi) = \cos\theta$$

$$P_5(\cos\theta, \phi) = \sin^2\theta \sin 2\phi$$

$$P_6(\cos\theta, \phi) = \sin 2\theta \sin \phi$$

$$P_7(\cos\theta, \phi) = \sin\theta \sin\phi$$

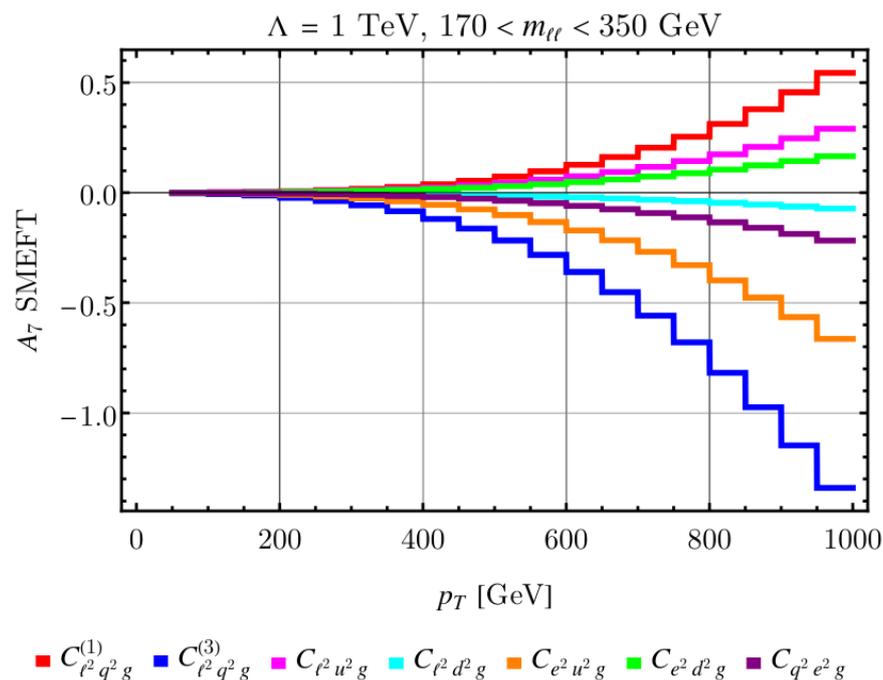
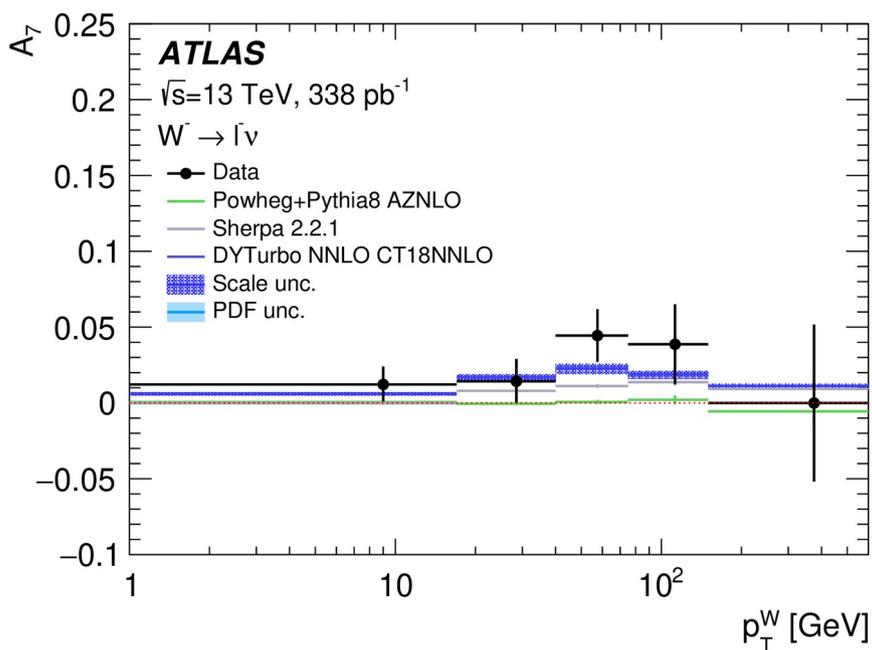
- Transverse momentum spectrum and polar-angle asymmetries A₃, A₄ are crucial ingredients for the modelling of W production needed for the m_W measurement



W-boson T-odd asymmetries

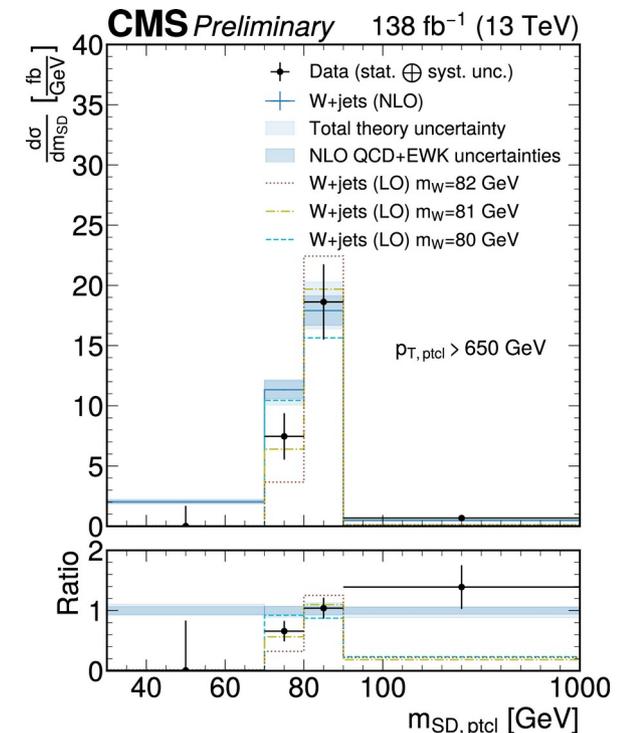
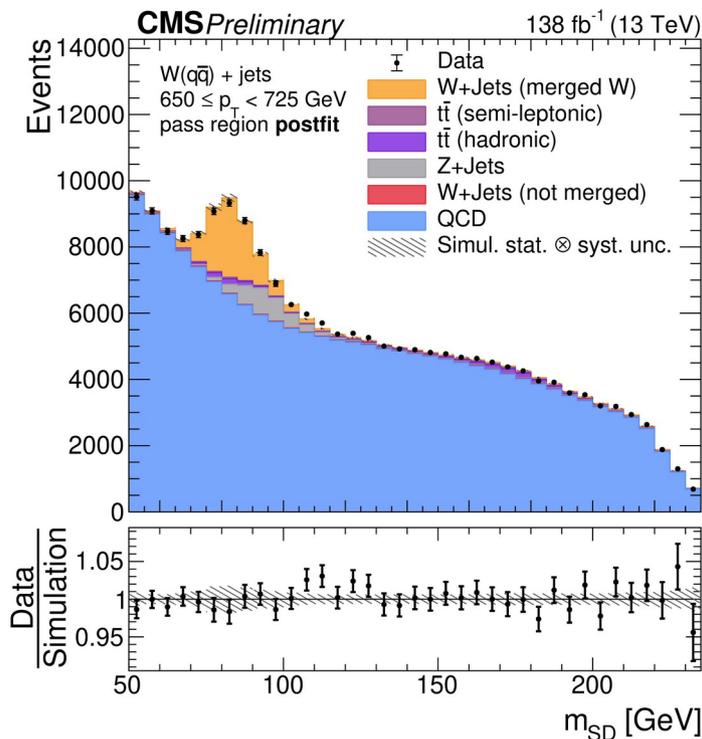
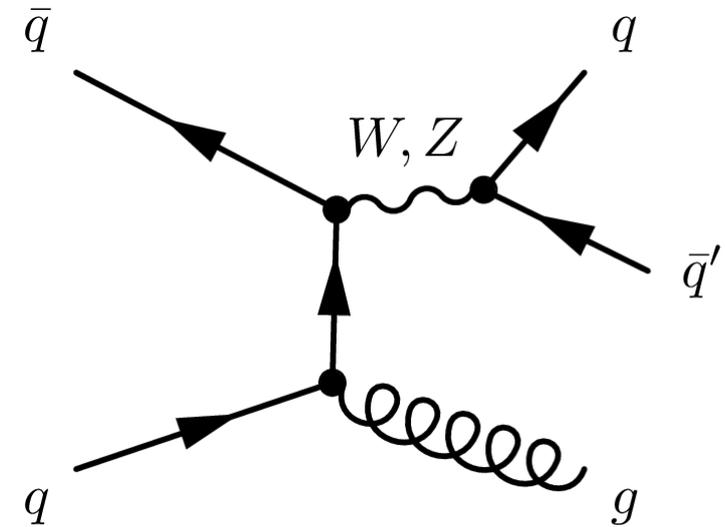
- $A_{5,6,7}$ coefficients are naive T-odd because ϕ flips sign under naive time reversal
- Sizeable non zero T-odd asymmetries in W production predicted by QCD starting at $O(\alpha_s^2)$ [PRL 52 \(1984\) 1076](#), [arXiv:1407.1016](#)
- Could be used to probe dimension-8 SMEFT [arXiv:2511.19617](#)
- Interesting hint of positive A_7 coefficient motivates further dedicated measurements

$$\begin{aligned}
 P_0(\cos \theta, \phi) &= \frac{1}{2}(1 - 3 \cos^2 \theta) \\
 P_1(\cos \theta, \phi) &= \sin 2\theta \cos \phi \\
 P_2(\cos \theta, \phi) &= \frac{1}{2} \sin^2 \theta \cos 2\phi \\
 P_3(\cos \theta, \phi) &= \sin \theta \cos \phi \\
 P_4(\cos \theta, \phi) &= \cos \theta \\
 P_5(\cos \theta, \phi) &= \sin^2 \theta \sin 2\phi \\
 P_6(\cos \theta, \phi) &= \sin 2\theta \sin \phi \\
 P_7(\cos \theta, \phi) &= \sin \theta \sin \phi
 \end{aligned}$$



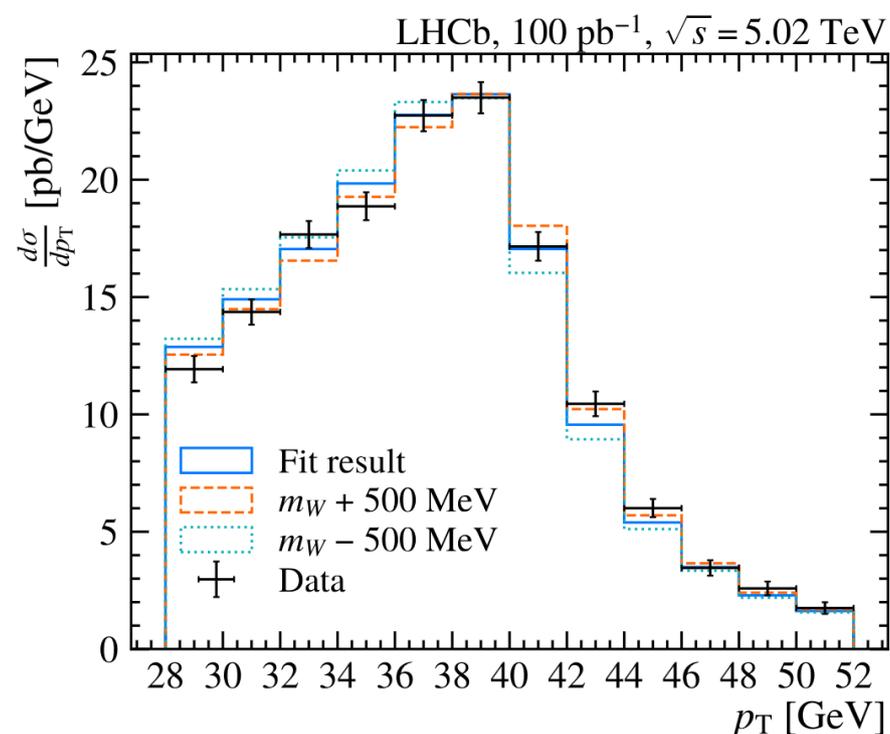
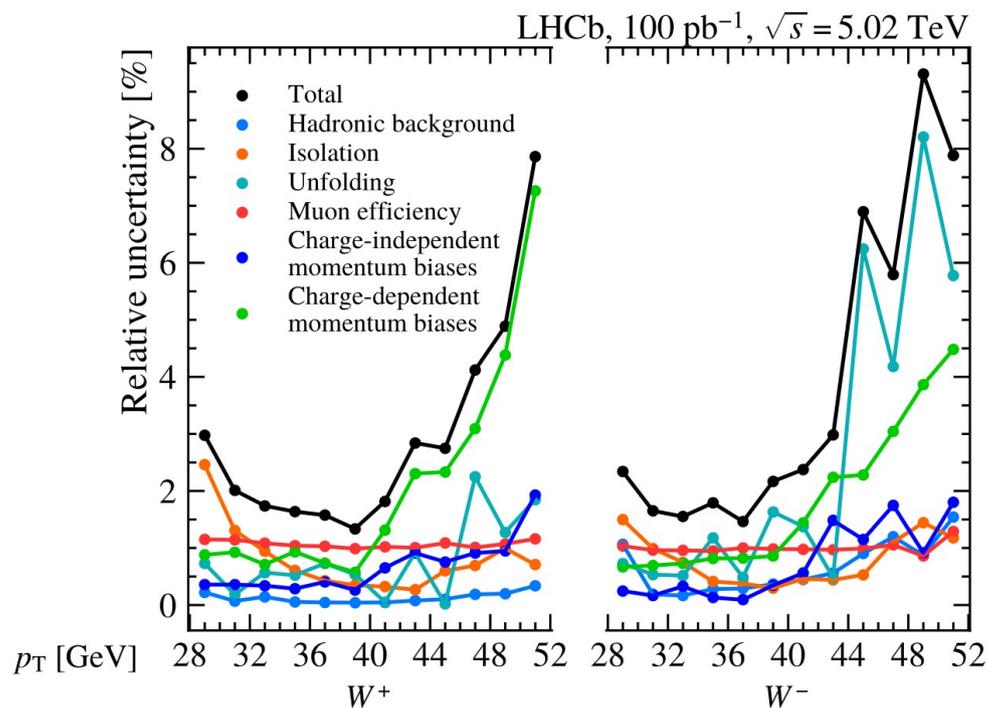
Boosted m_W

- Exploit hadronic decay of the W-boson to fully reconstruct the invariant mass
- Highly boosted W bosons, reconstructed as a large-radius jet, jet mass measured using the soft drop algorithm
- Challenge: Large background and worse energy resolution of jets compared to electrons and muons
- $m_W = 80.77 \pm 0.57 \text{ GeV}$



m_W from unfolded cross sections

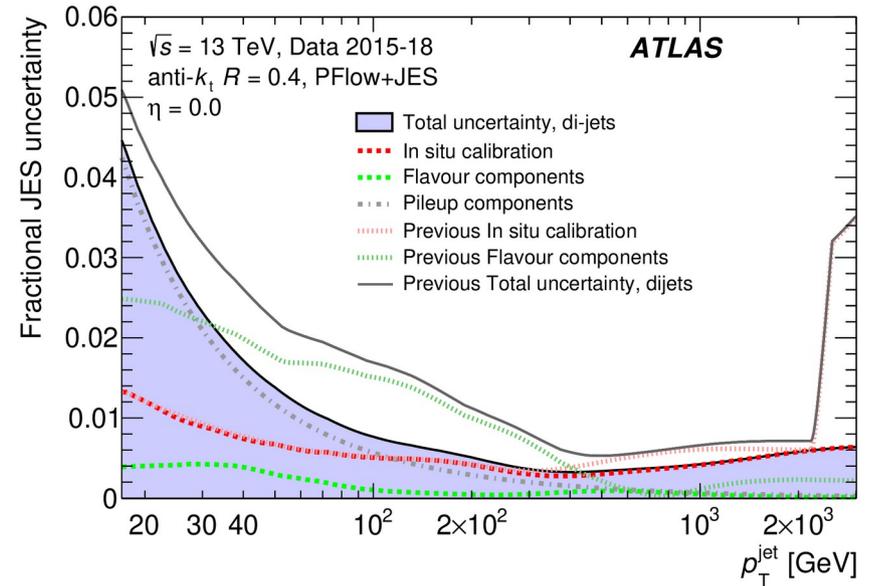
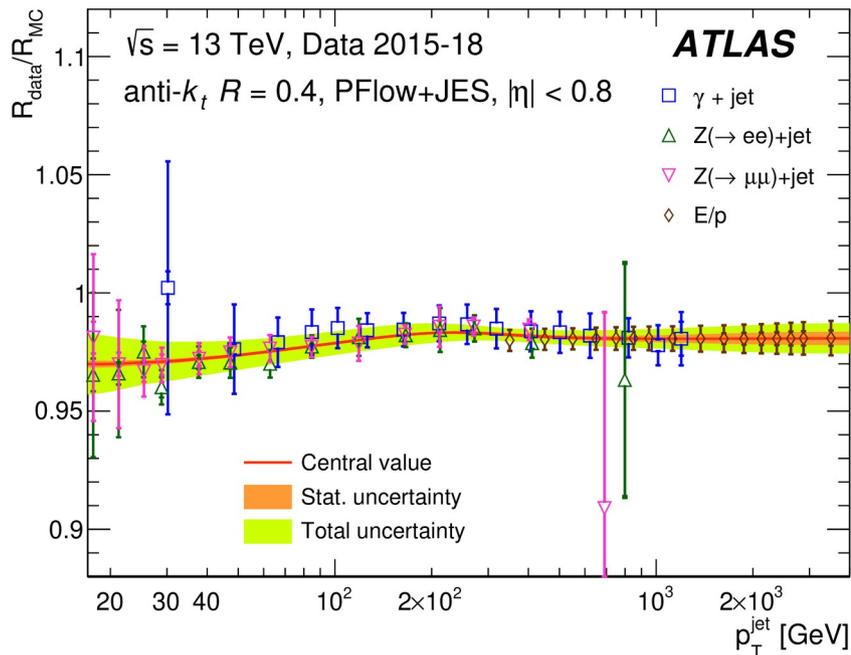
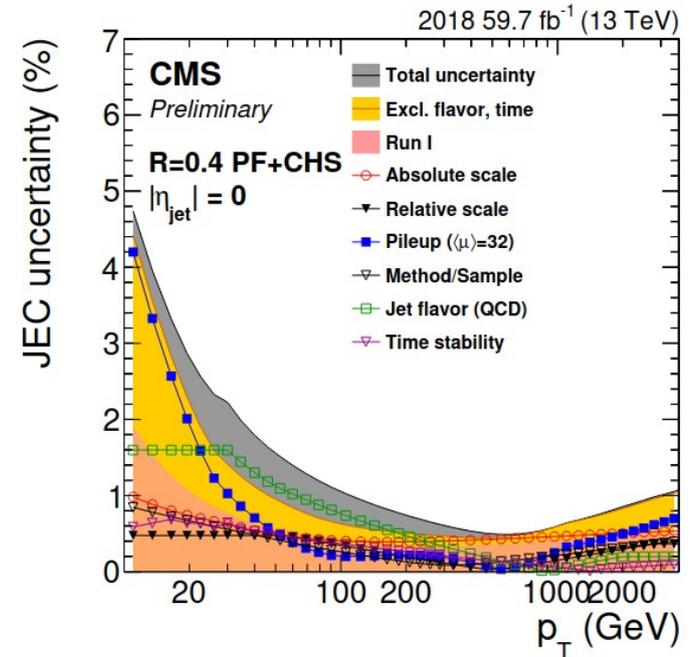
- LHCb measurement of m_W in the forward region from unfolded p_T muon cross section
- Performed on a statistically limited dataset as a proof of principle
- $M_W = 80.369 \pm 0.130(\text{exp.}) \pm 0.33(\text{th.}) \text{ GeV}$
- Interesting new technique, will allow improving existing measurements of m_W when more accurate theory predictions become available



Jet energy scale

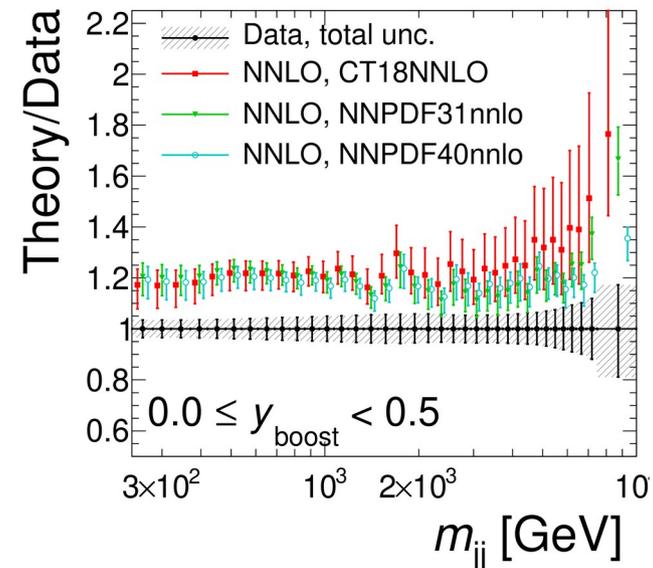
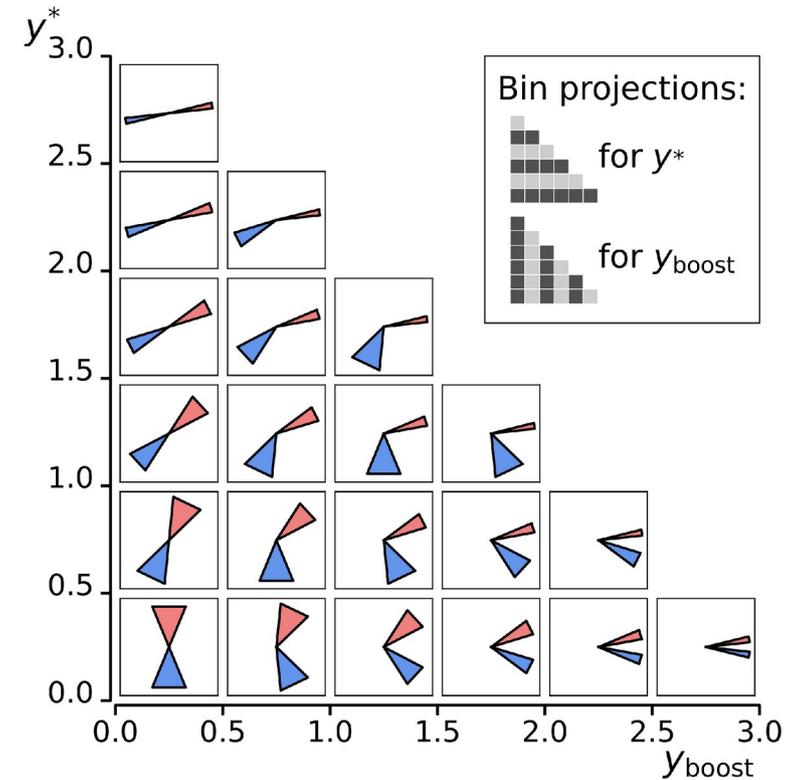
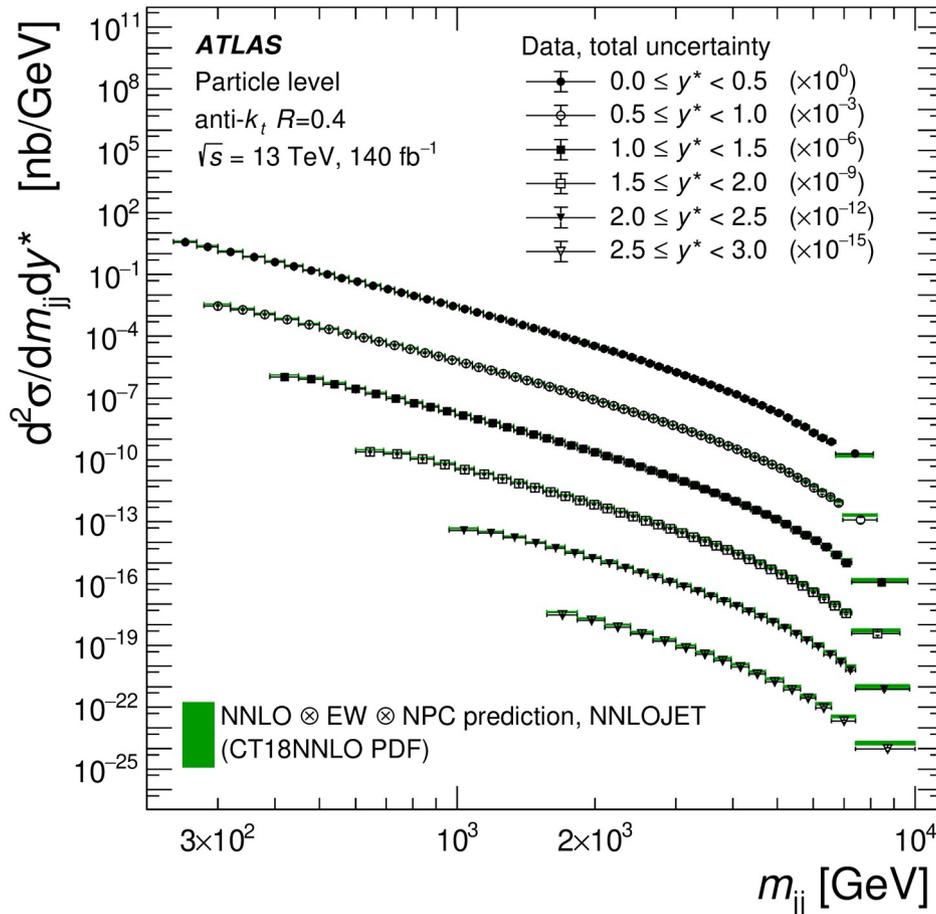
- ATLAS and CMS have very precise jet energy scale calibrations
- ATLAS recently made a step forward combining single-particle response and Z/ γ + jet p_T balance, reaching 0.3% at 300 GeV and 0.6% at 4 TeV

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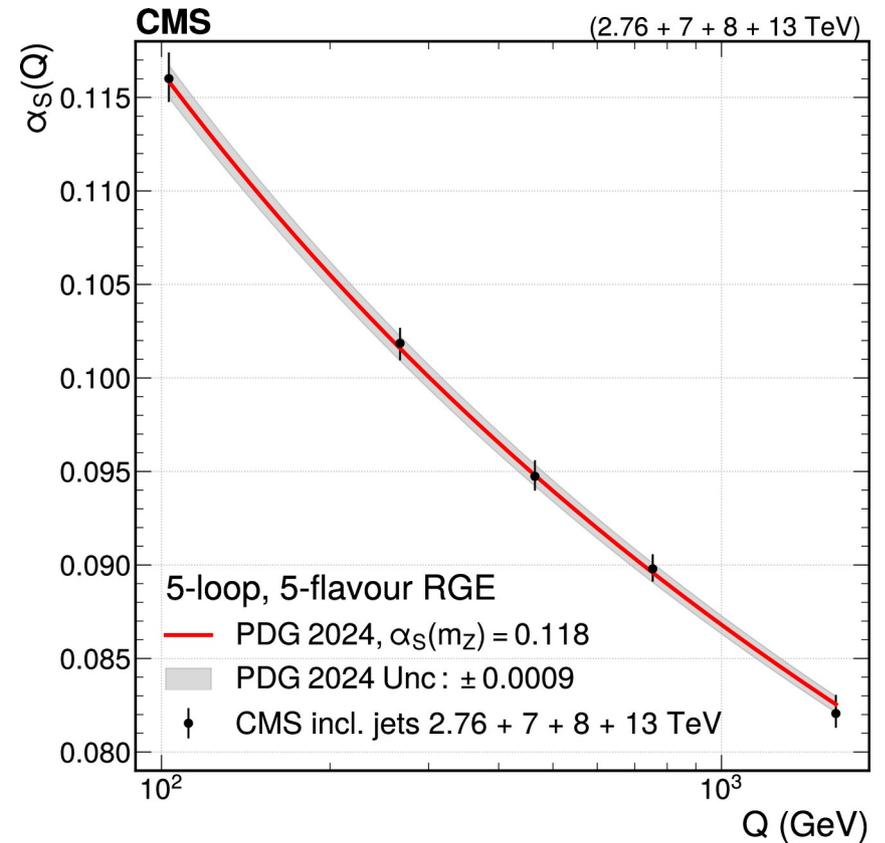
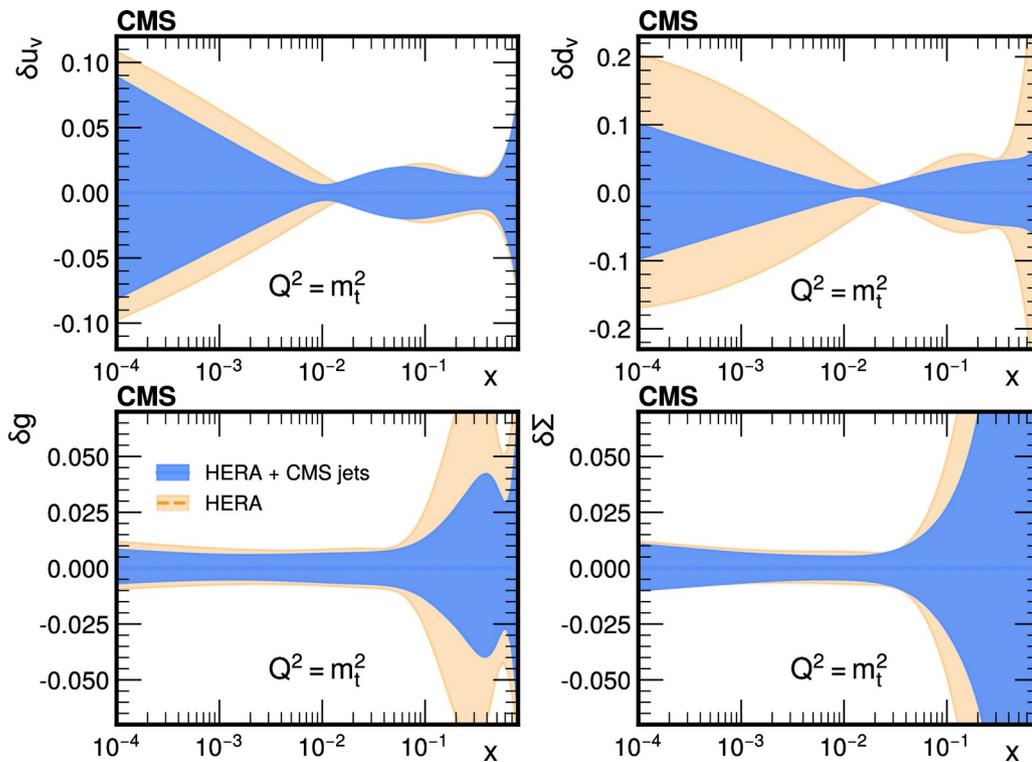
ATLAS inclusive dijets

- Measurement of inclusive dijets cross sections as functions of m_{jj} , y^* , y_{boost}
- Rely on very precise JES calibration
- Provide constraints on PDFs, and can be used for determination of the strong coupling



CMS $\alpha_s(m_Z)$ from jet production

- Significant reduction of the gluon PDF uncertainty at high x
- Tested RGE running of α_s up to the TeV scale



Summary

- SM precision measurements at the LHC are enabled by a tremendous advance of experimental techniques, physics performance, and development of theory tools beyond expectation
- They provide stringent tests of the consistency of the SM, and offer validation and improvement of theory predictions, which is crucial to increase the sensitivity of BSM searches

