



Long-lived HNLs at Spallation Sources

IRN Neutrino meeting, 2th June, 2026
IJCLab Orsay, France

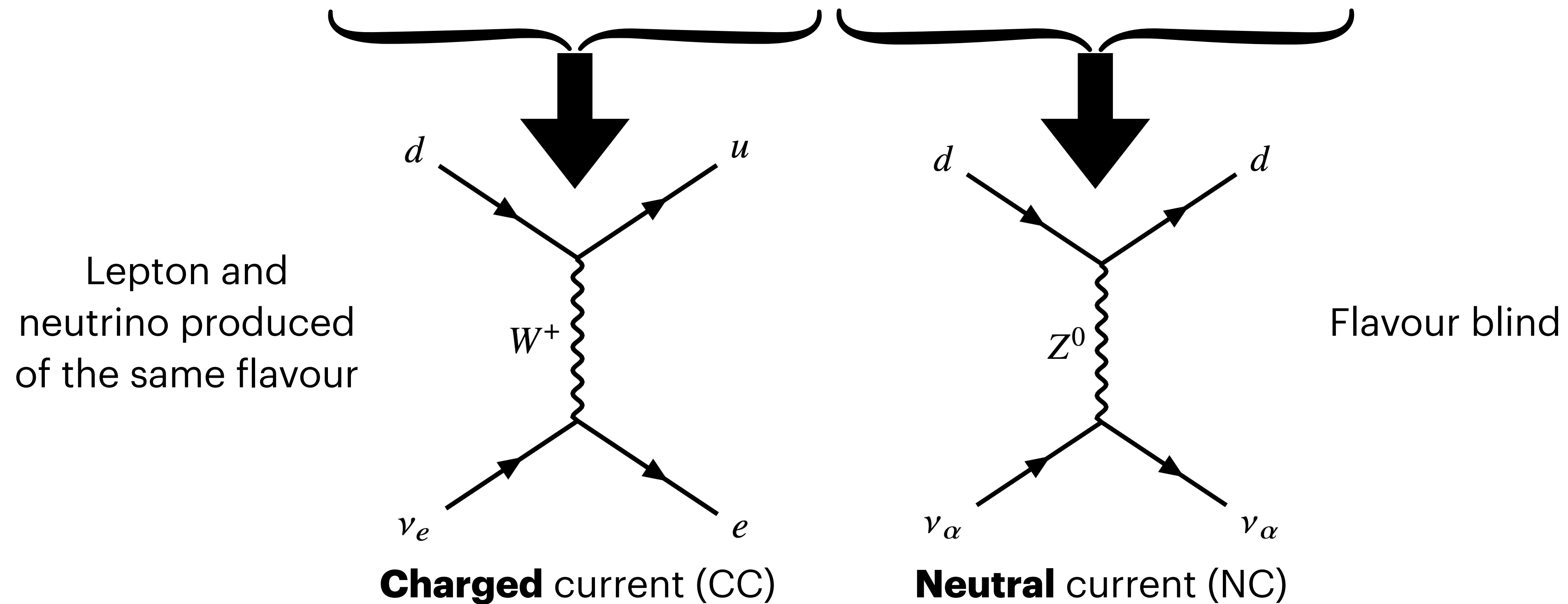
Based on arXiv:2606.XXXXX in collaboration with Valentina De Romeri,
Alfredo Galindo-Uribarri, Víctor Martín Lozano and Gonzalo Sánchez García

Ana Martín-Galán
(IFIC, Valencia - UV/CSIC)

Neutrinos in the Standard Model

Neutrinos in **SM** are **massless** and **weakly interacting**:

$$\mathcal{L}_{SM} \subset -\frac{g}{\sqrt{2}} \sum_{\alpha} \bar{\nu}_{\alpha L} \gamma^{\mu} \ell_{\alpha L} W_{\mu} - \frac{g}{2 \cos \theta_W} \sum_{\alpha} \bar{\nu}_{\alpha L} \gamma^{\mu} \ell_{\alpha L} Z_{\mu} + h.c.$$



Nonetheless, we know at least two of three light neutrino are **massive** \longrightarrow Physics beyond SM!

Heavy neutral leptons

One of the simplest extensions of SM is to **add** (at least) **two fermionic singlets** or **HNLs** that mix with the SM neutrinos:

$$\nu_\ell = U_{\ell N} N + \sum_{i=1}^3 U_{\ell i} \nu_i$$

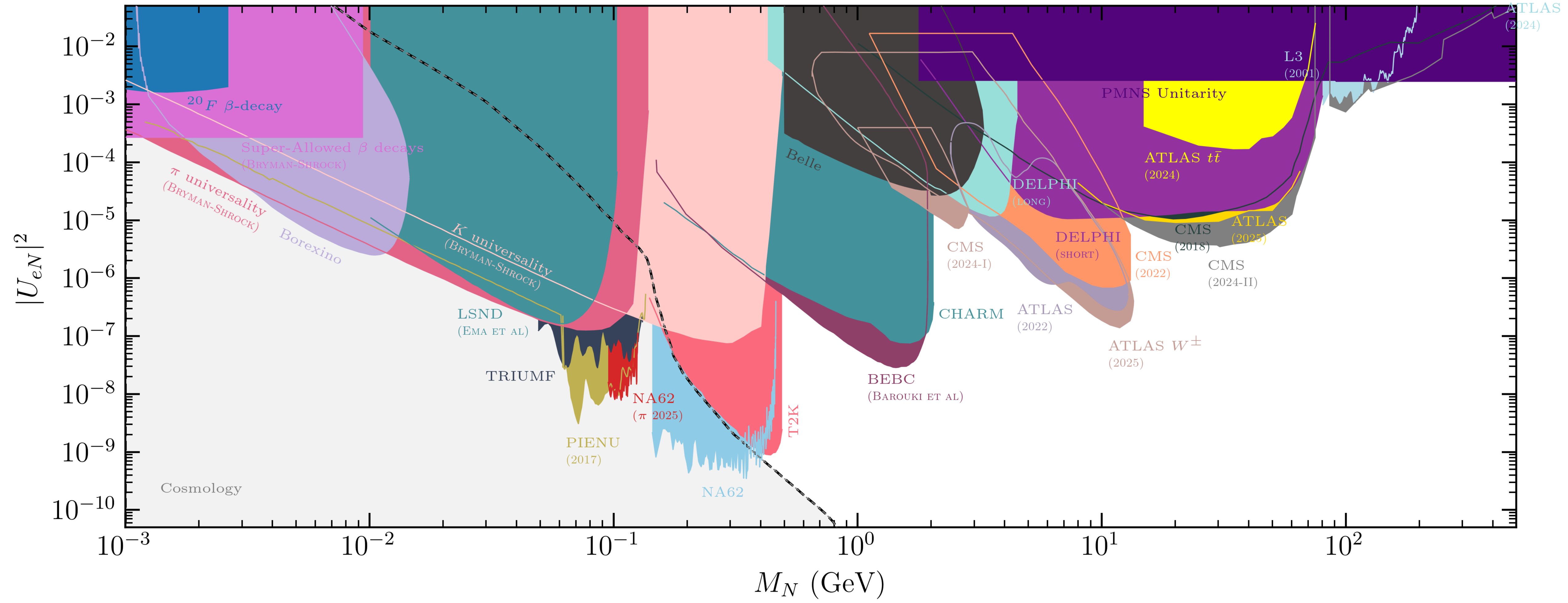
In the **minimal HNL Lagrangian** the charged and neutral current interactions with the SM leptons are defined as

$$\mathcal{L} \supseteq -\frac{g}{\sqrt{2}} U_{\ell N}^* \bar{\ell} W_\mu^- \gamma^\mu P_L N - \frac{g}{2 \cos \theta_W} U_{\ell N}^* \bar{\nu}_\ell Z_\mu \gamma^\mu P_L N + \text{h.c.}$$

$$U_{\ell N} \ll 1 \longrightarrow \text{Long-Lived Particle}$$

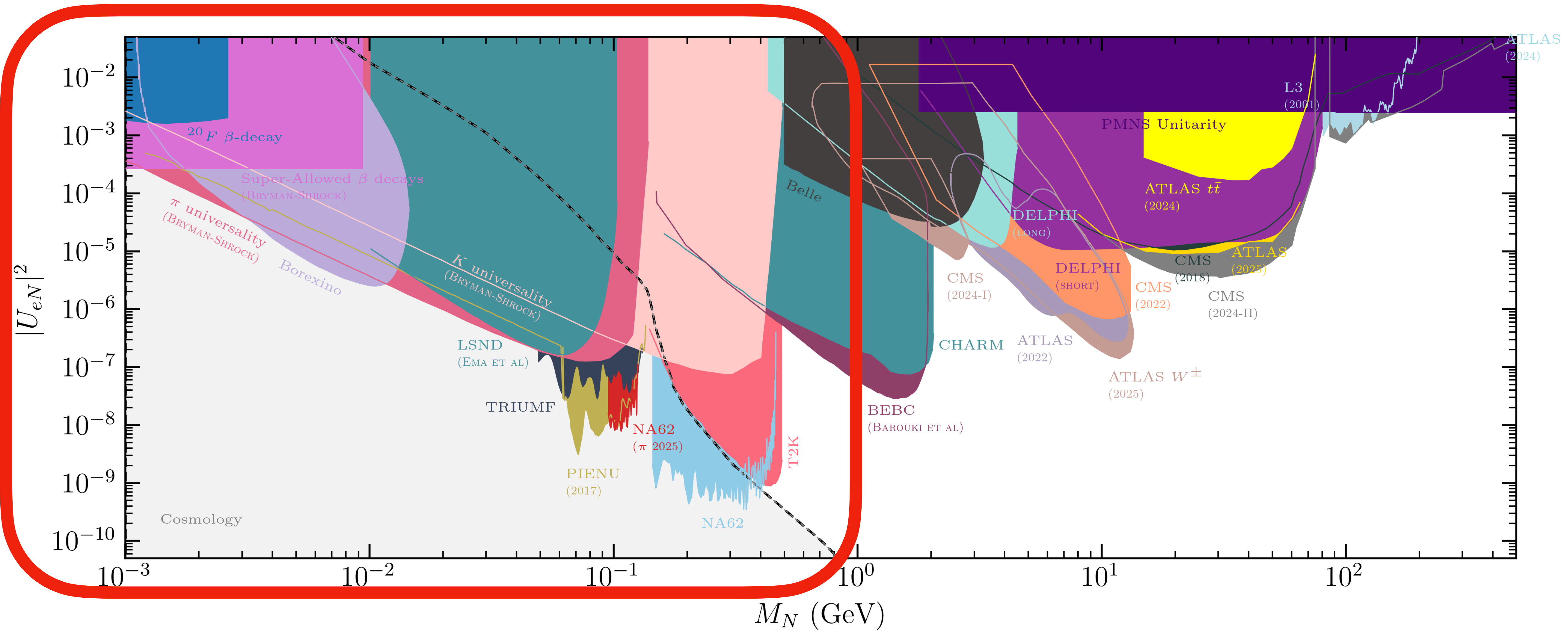
Our case: **only one of the HNLs is** kinematically accessible, which **mixes** with the light states. We assume them to be **pseudo-Dirac** with Majorana mass term negligible, $m_N \lesssim 1 \text{ GeV}$.

Current limits on HNLs



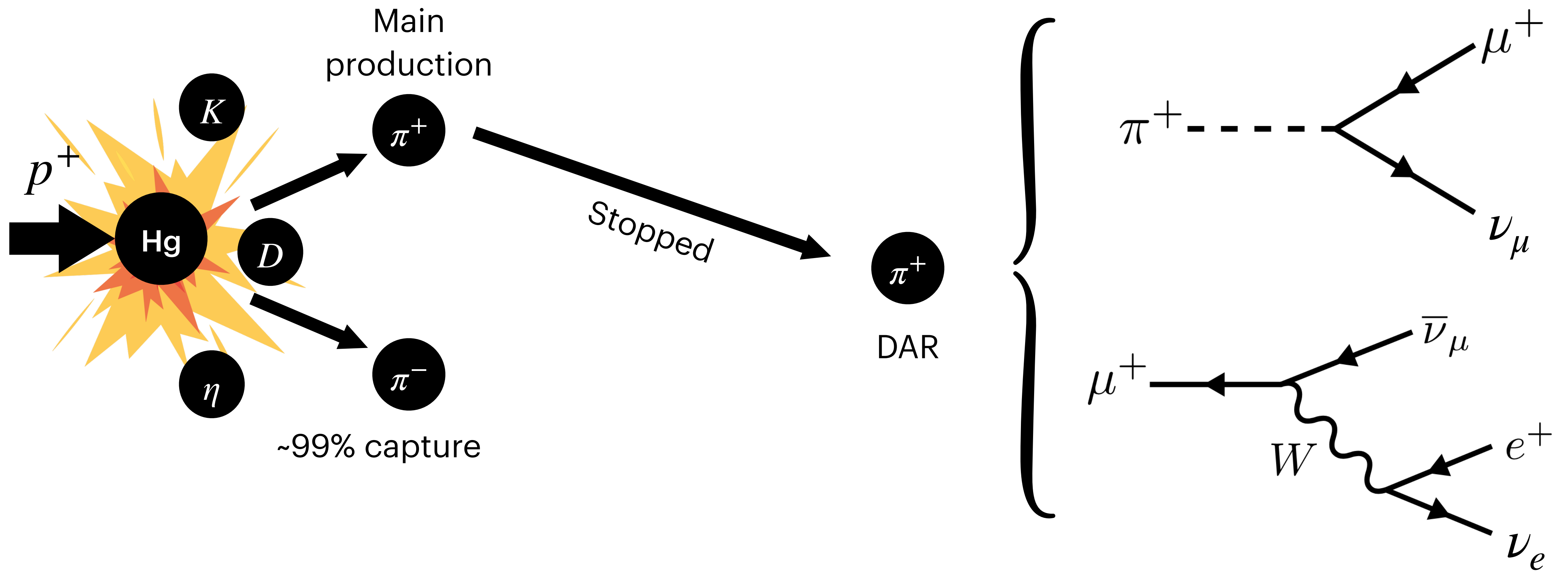
Current limits on HNLs

Parameter space **potentially accessible** via **spallation sources!**



SM neutrinos from stopped π at spallation sources

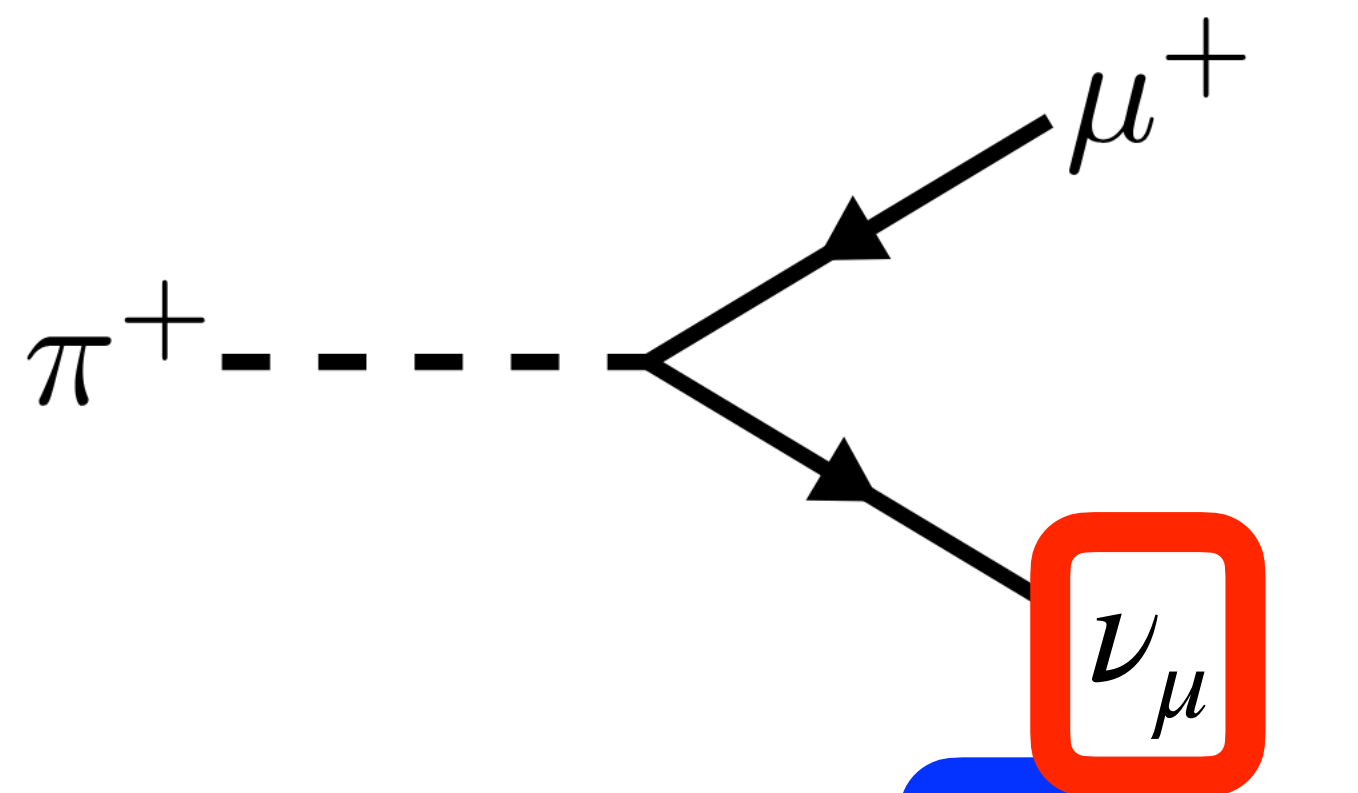
Pions as byproduct in spallation sources. **Three SM neutrino fluxes** produced from these pions.



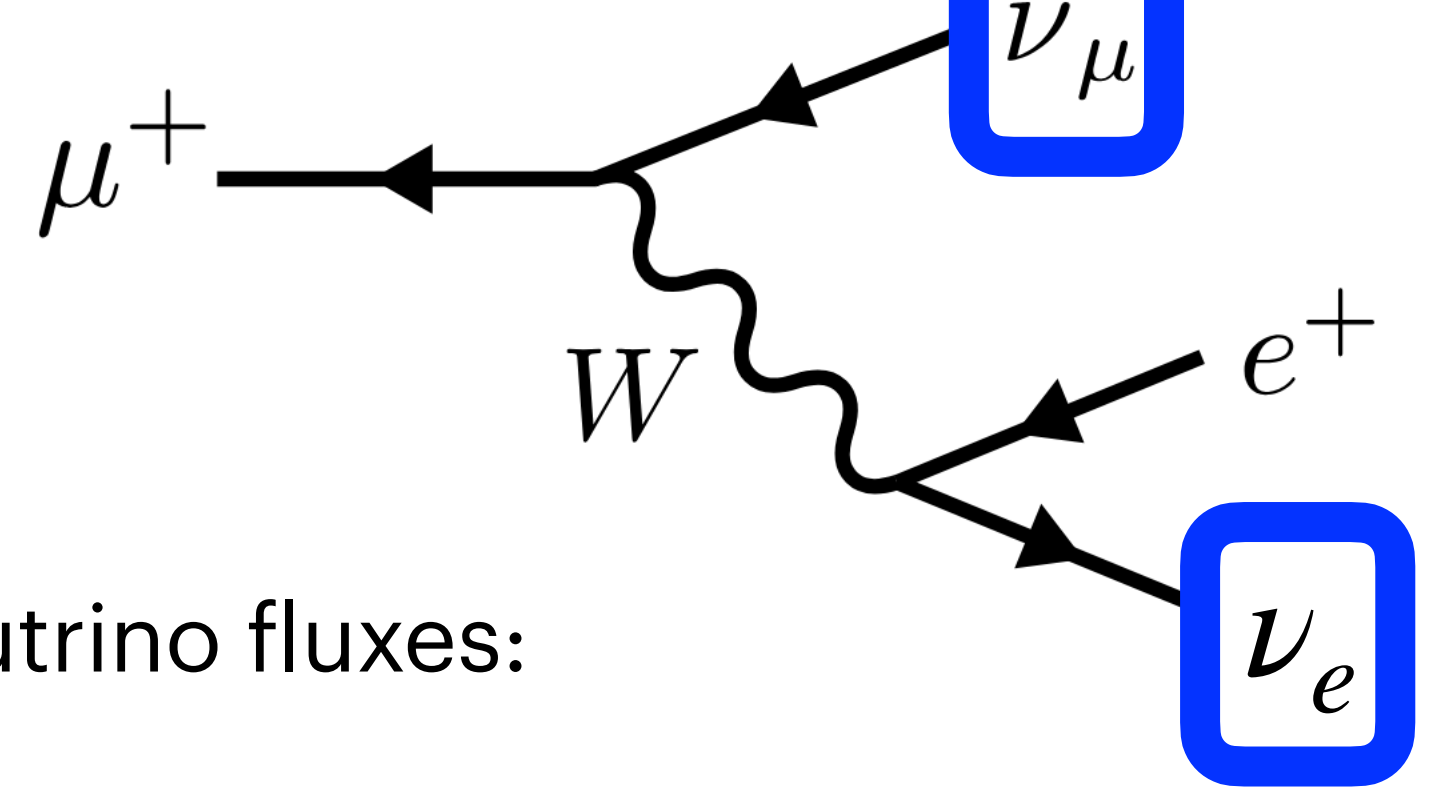
Timing of the SM neutrino signal

Pulsed beam and **arriving time of the neutrinos** help us to **distinguish** between **contributions** and **backgrounds**.

$\tau = 26 \text{ ns}$

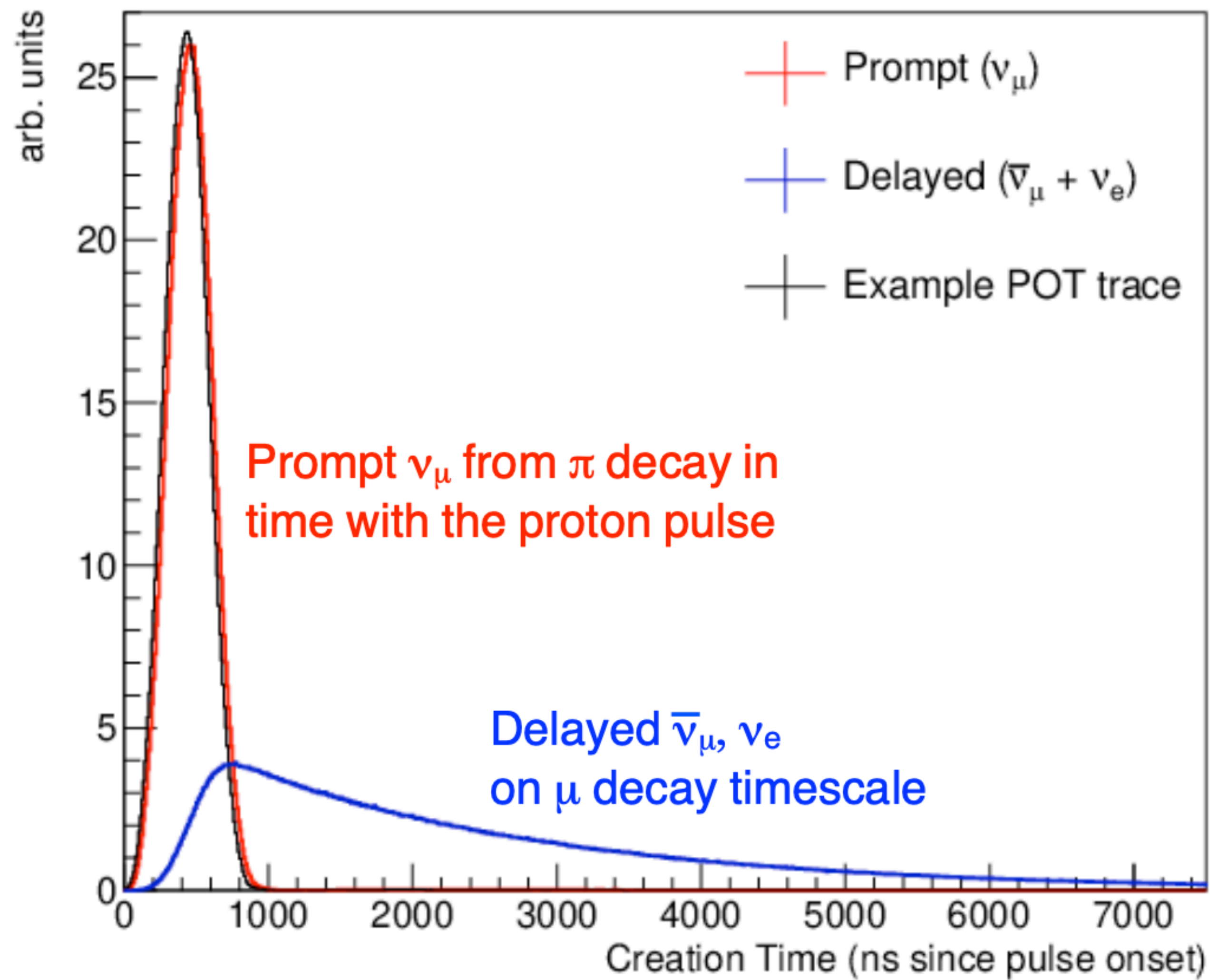


$\tau = 2.2 \mu\text{s}$



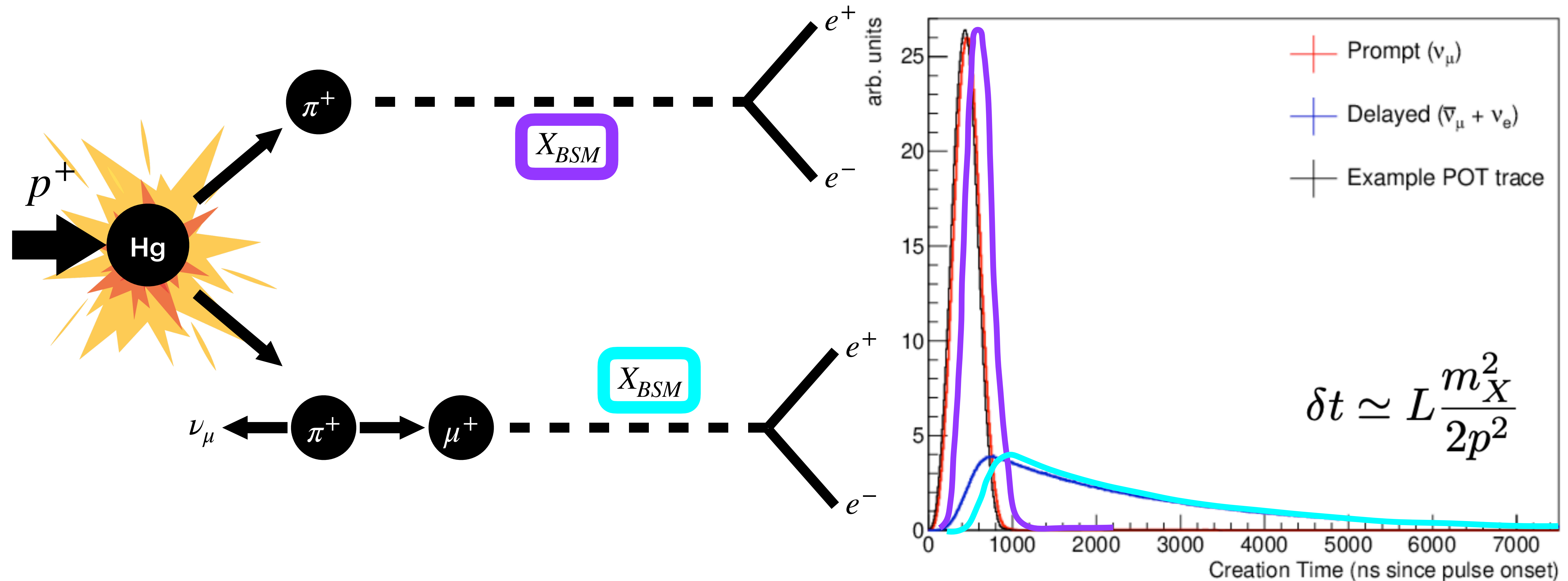
Three neutrino fluxes:

- **1 prompt**
- **2 delayed**



Spallation sources and physics BSM

Heavy new physics would give us a **delayed signal** with respect to the SM neutrino flux.



Spallation sources we consider

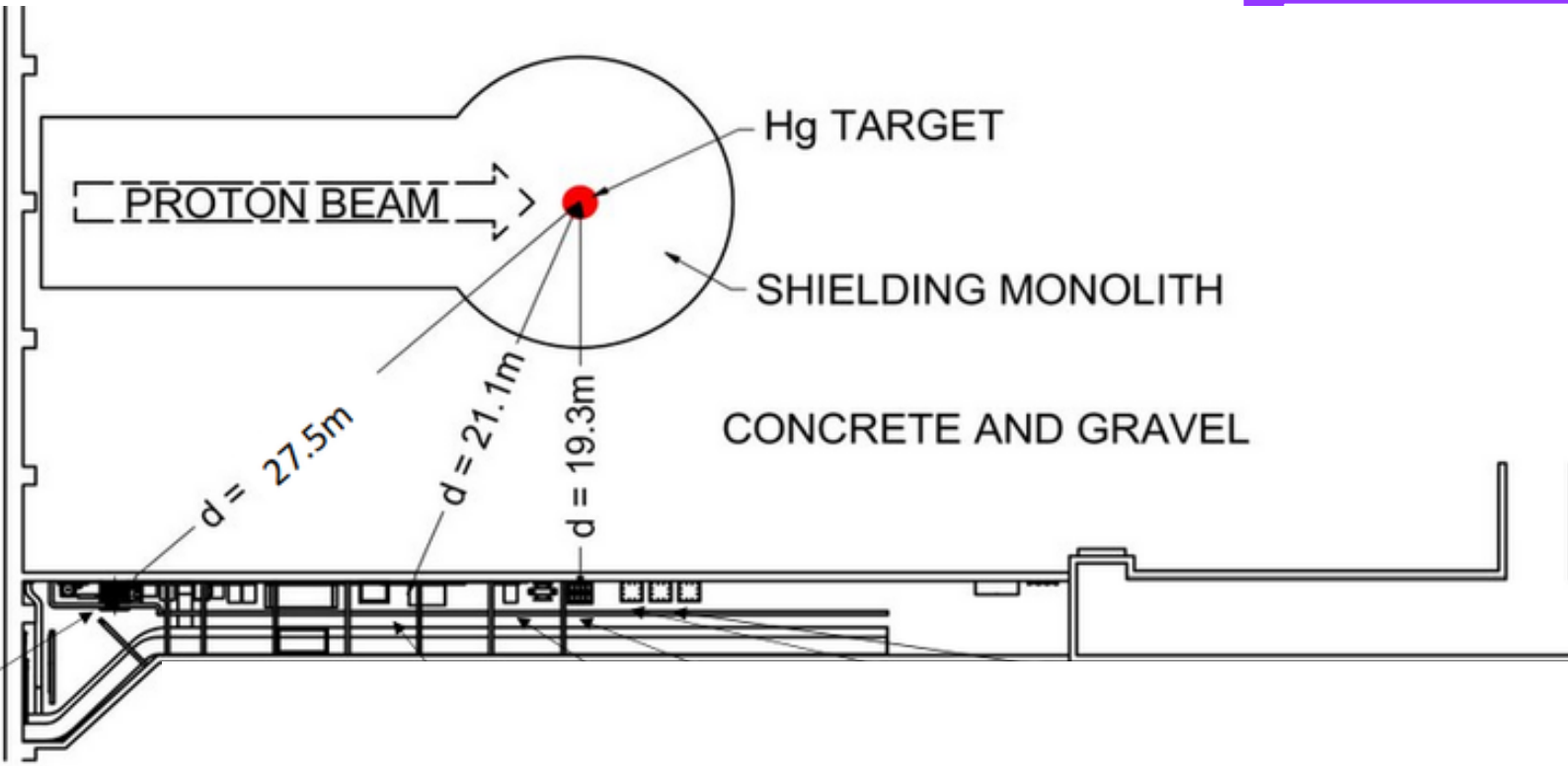


Spallation sources we consider



Beam characteristics:

- Mercury target
- (1.0 → 1.7 → 2.0) MW
- (1 → 1.3) GeV p^+
- (9 % → 11%) π DAR/POT



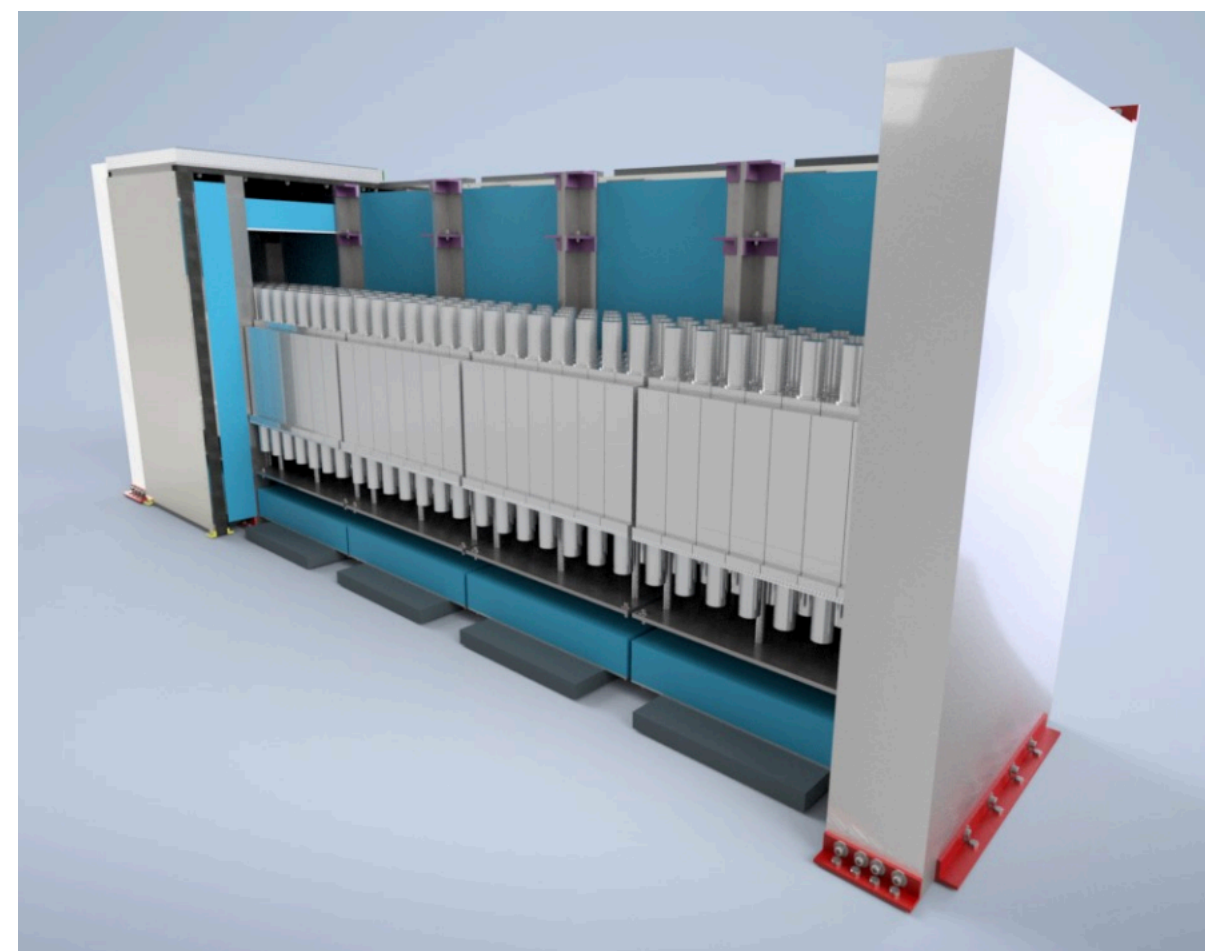
In this talk I will focus on SNS!



COHERENT experiment



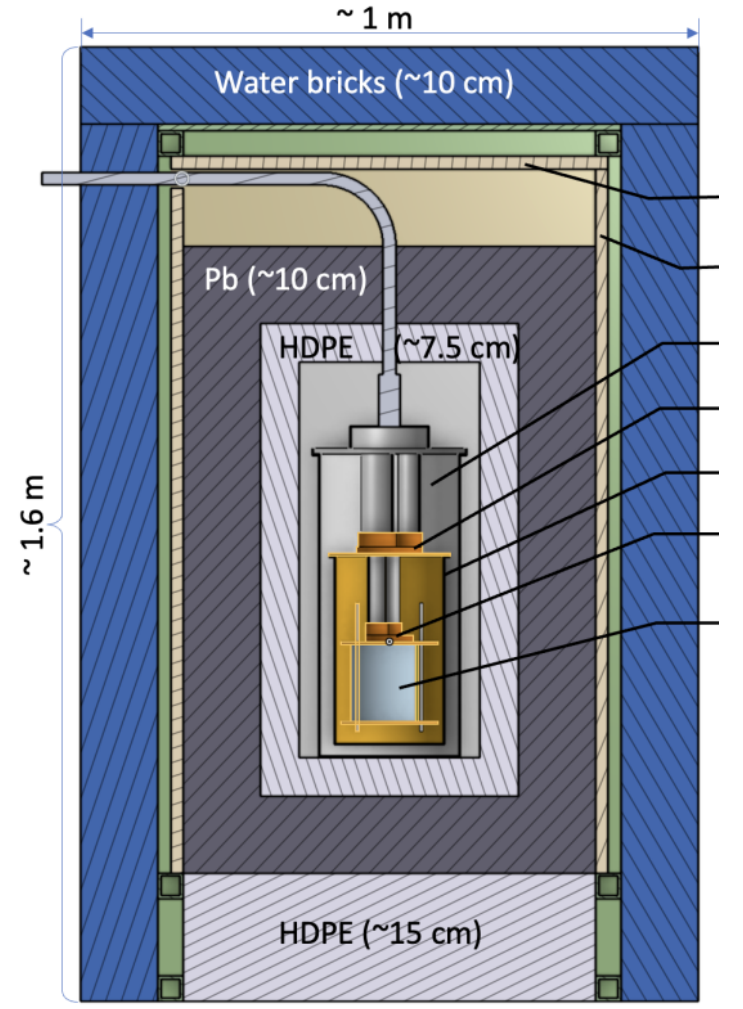
COHERENT was built with CE ν NS detection in mind.
 Nevertheless, BSM searches **could** be performed.



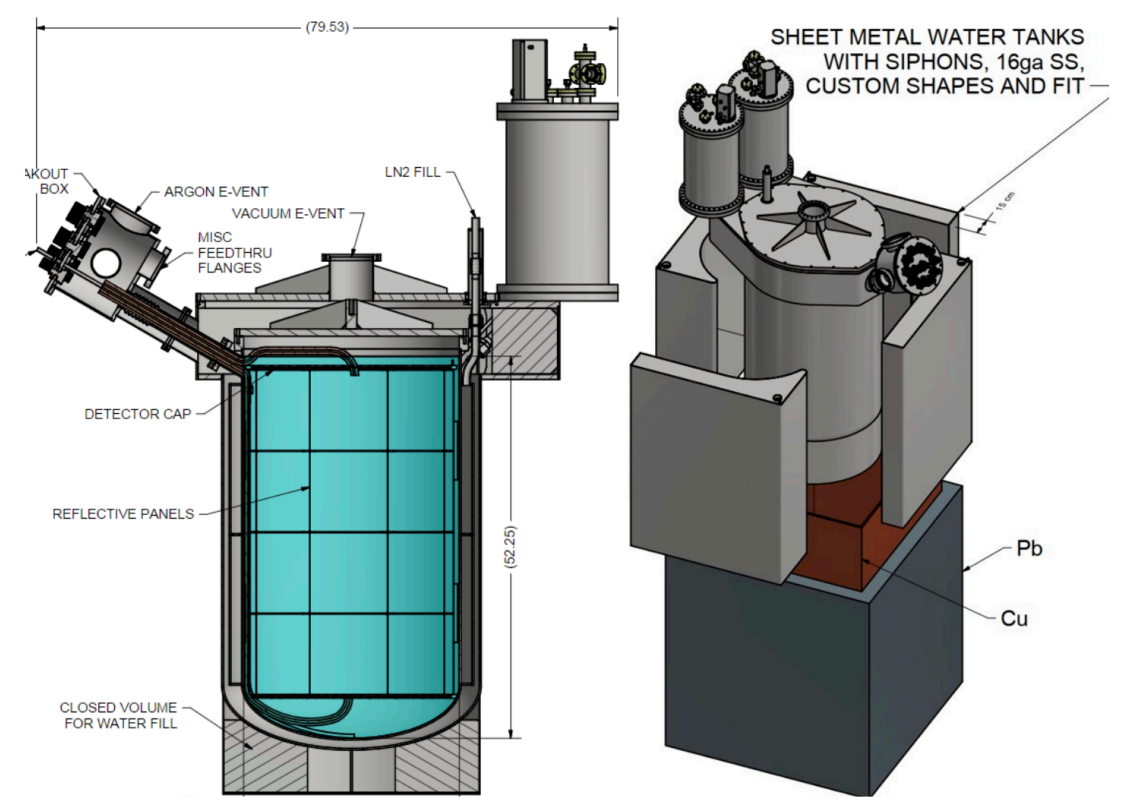
NaIvETE



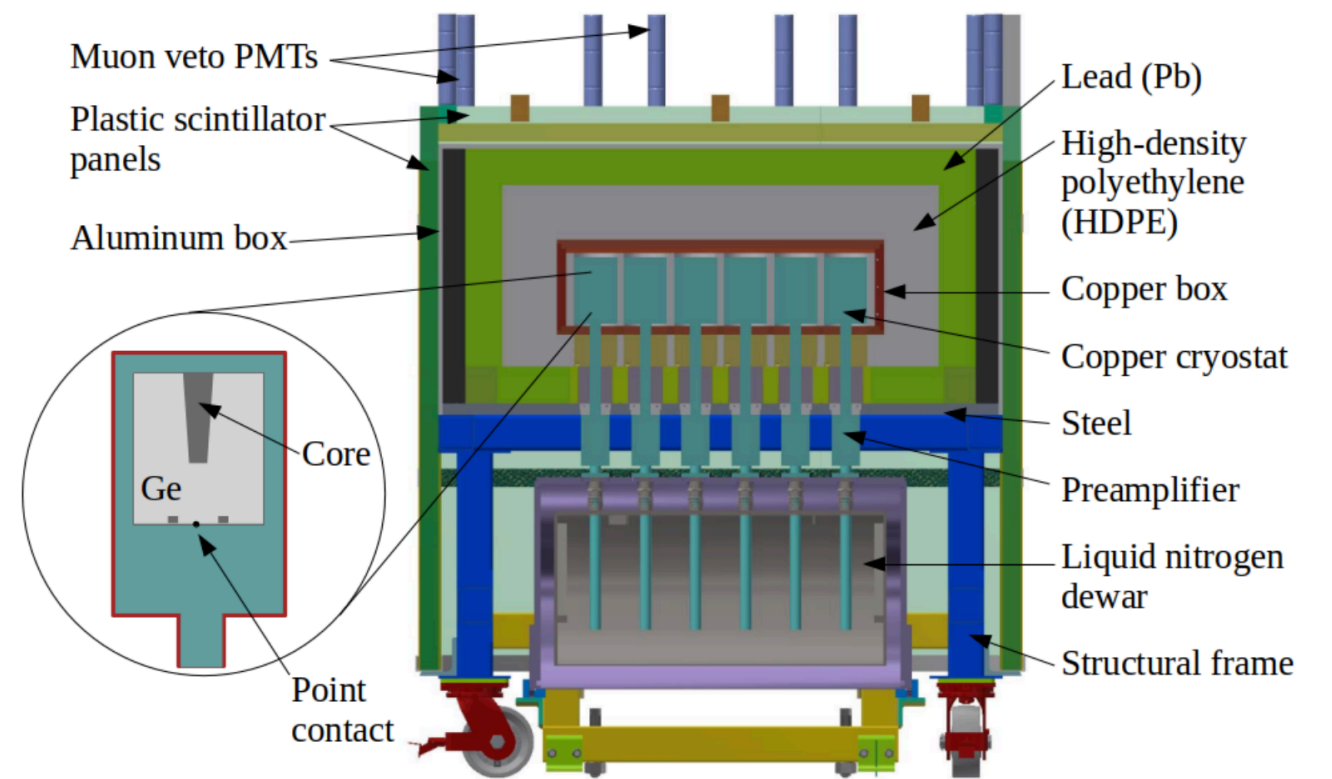
D₂O+H₂O



Cryo-CsI



LAr



Ge-mini

- A priori* good features for LLP detection:
- Volume
 - Flux
 - Distance
 - Backgrounds

From running data

Setup considered

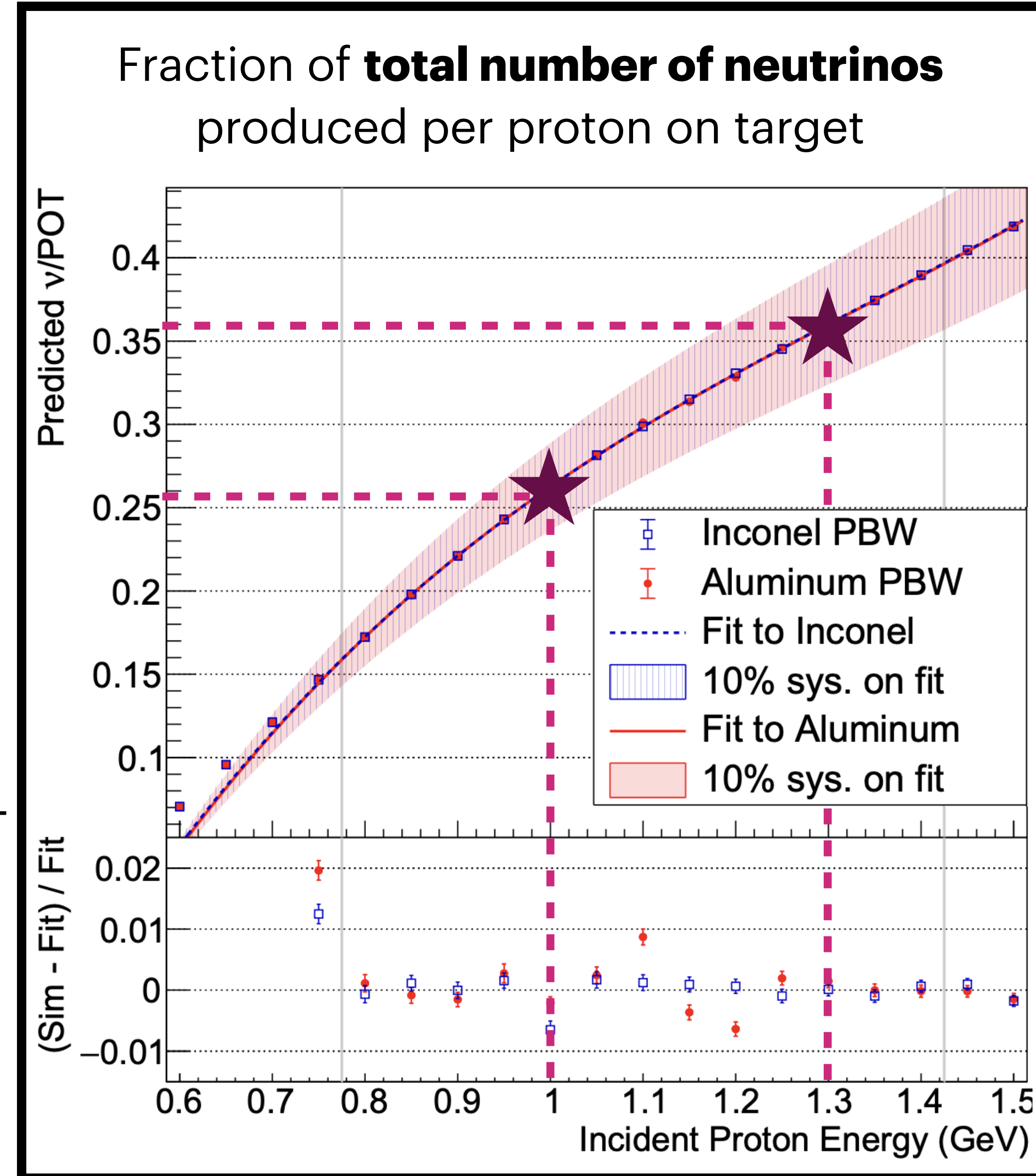
COHERENT Collaboration [2602.15652]
COHERENT Collaboration [2603.17951]

Detector	Total N^{POT}	c_{π^+}	Mass [Kg]
Ge	4.68×10^{22}	0.09	8.53
NaI v1	1.728×10^{23}	0.11	2400
NaI v2			3500
LAr			476
D2O+H2O			1000
cryo-CsI-1	Preliminary setup		10
cryo-CsI-2			100
LAr-750			750
LAr-10t			1000

Running or near future

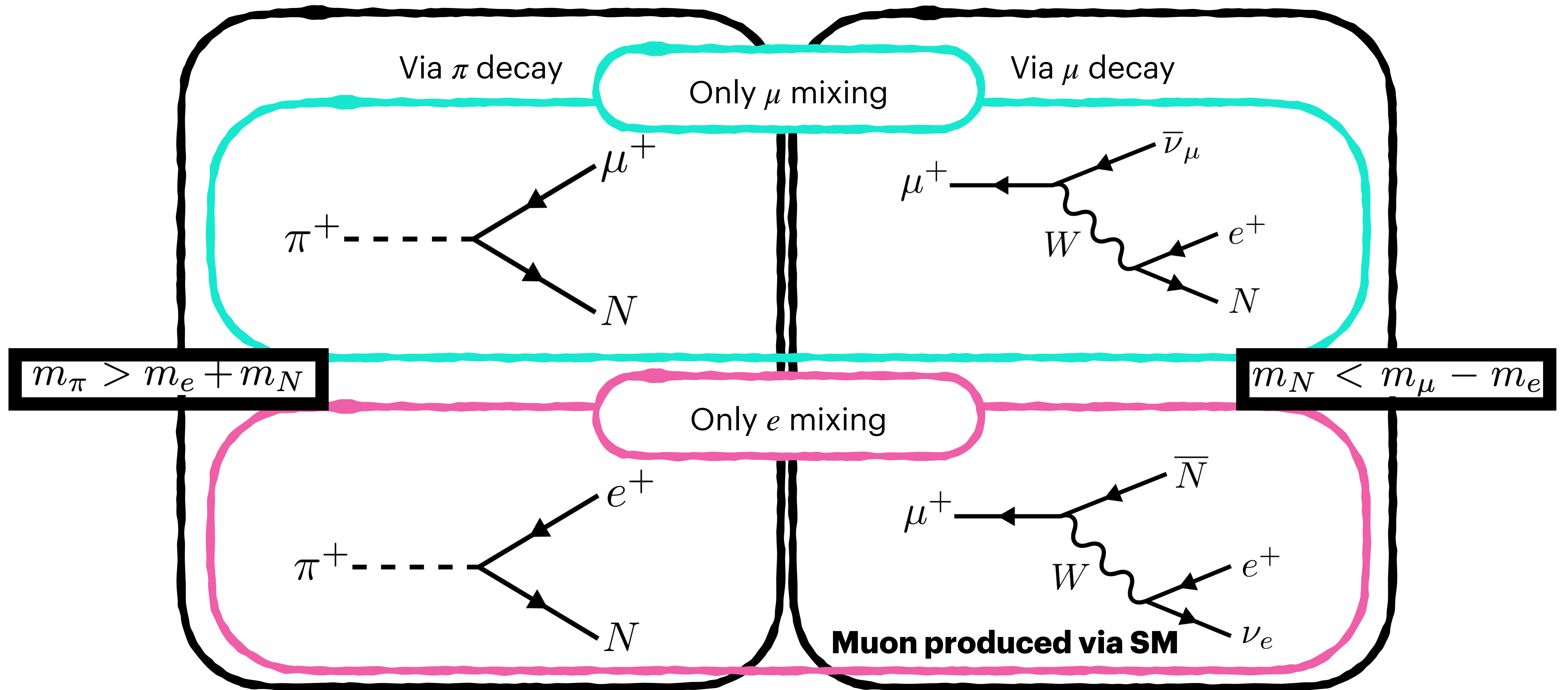
Proposals for future experiments

$$N^{POT} \propto \frac{P[\text{MW}]}{E[\text{GeV}]} t[\text{s}]$$



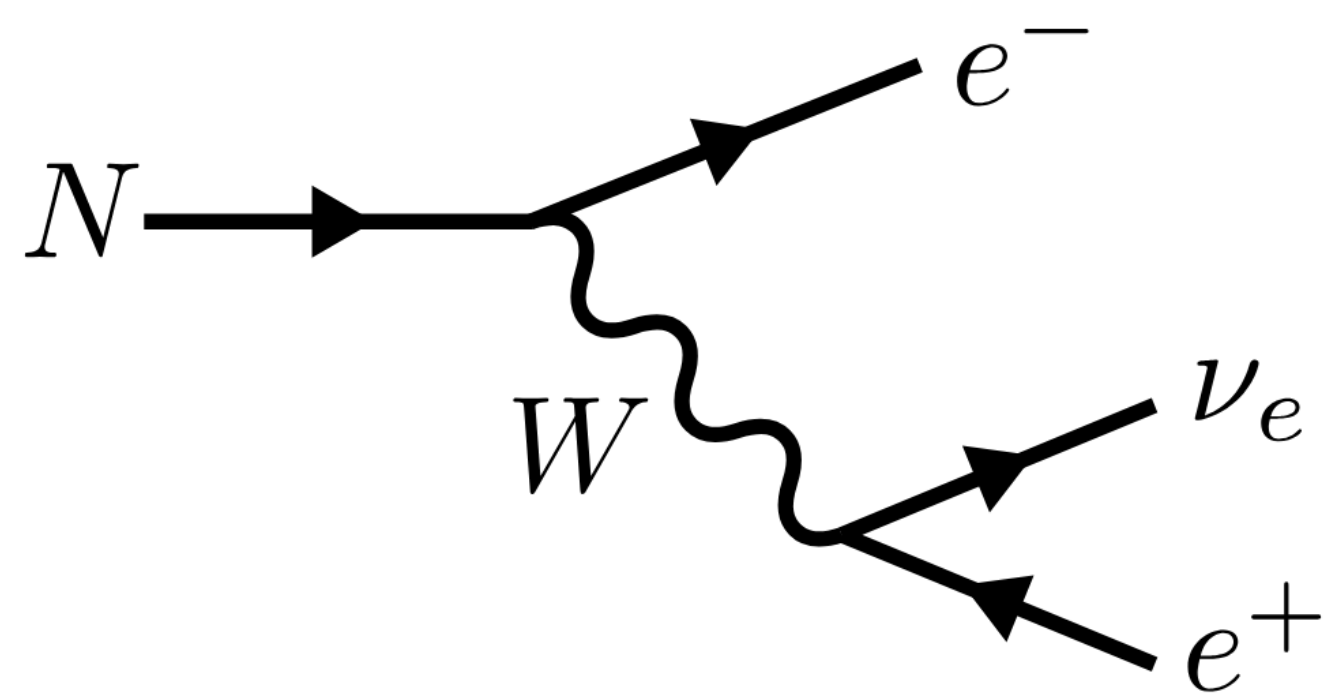
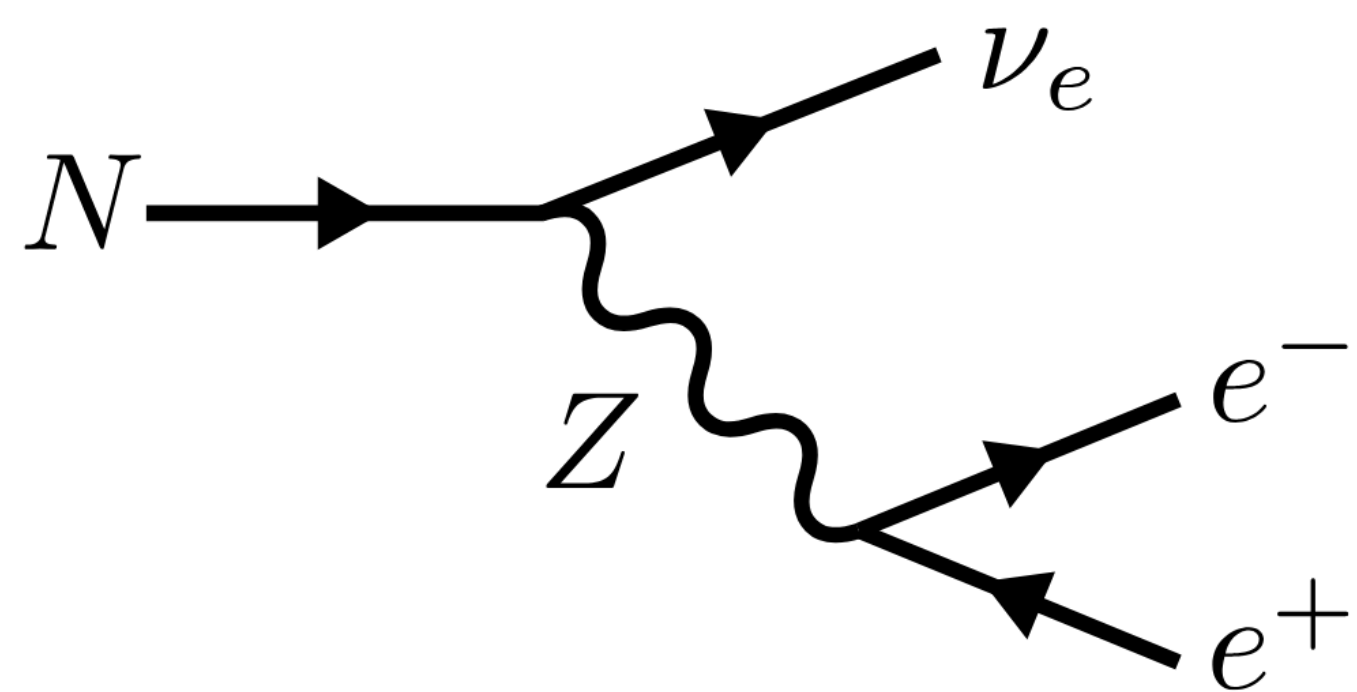
Production of HNLs at SNS

One mixing activated at the time. τ kinematically forbidden.

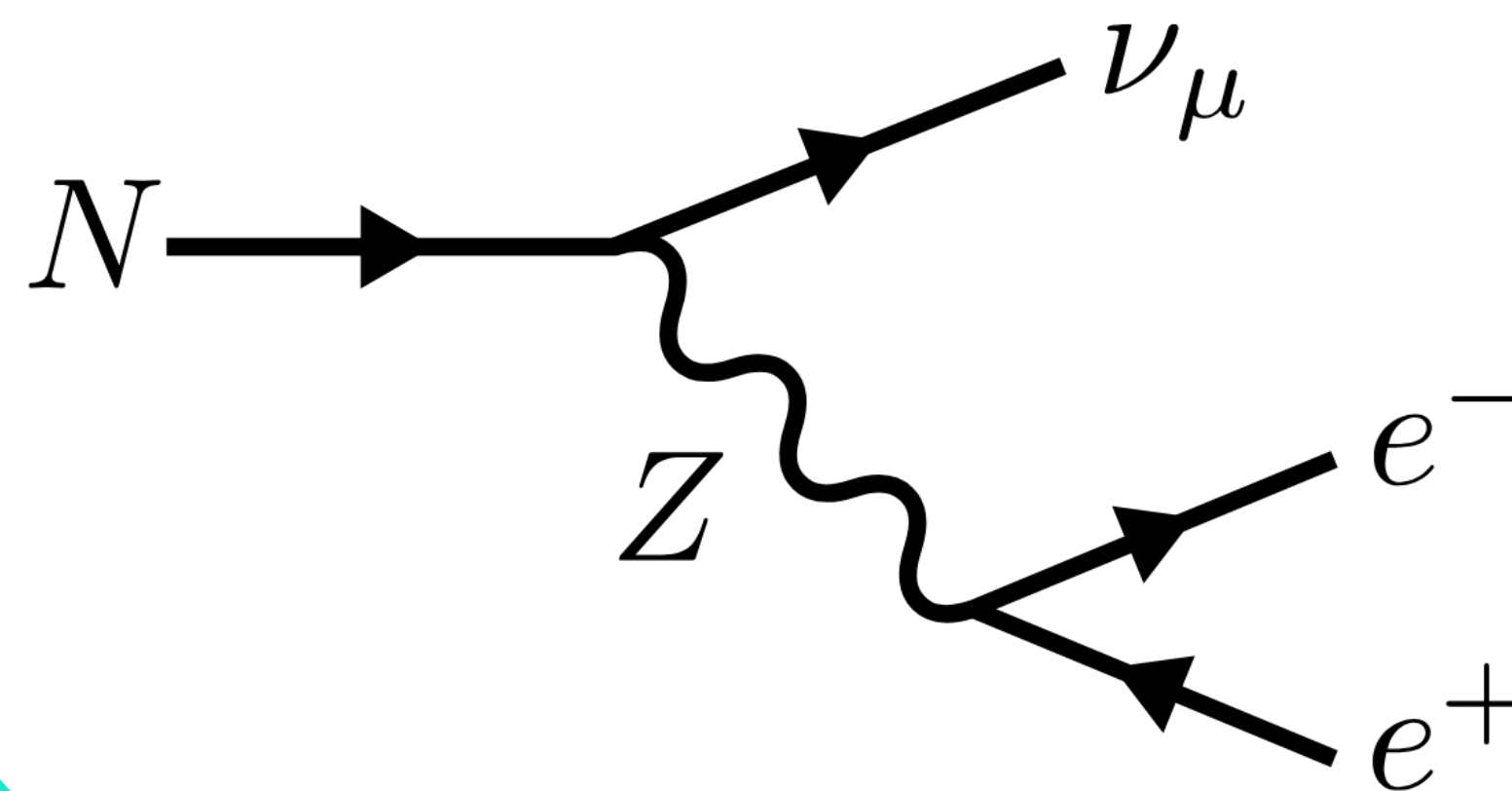


Decay of HNLs

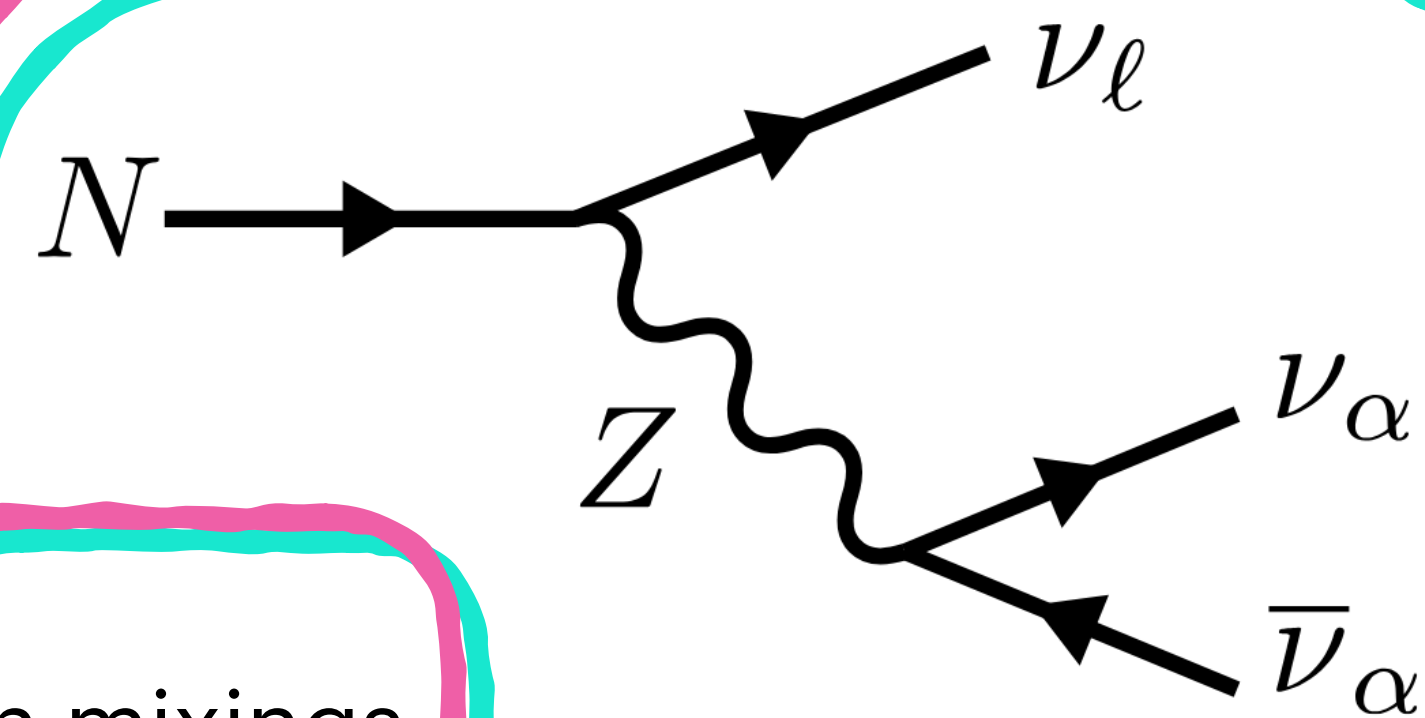
Only e mixing



Only μ mixing



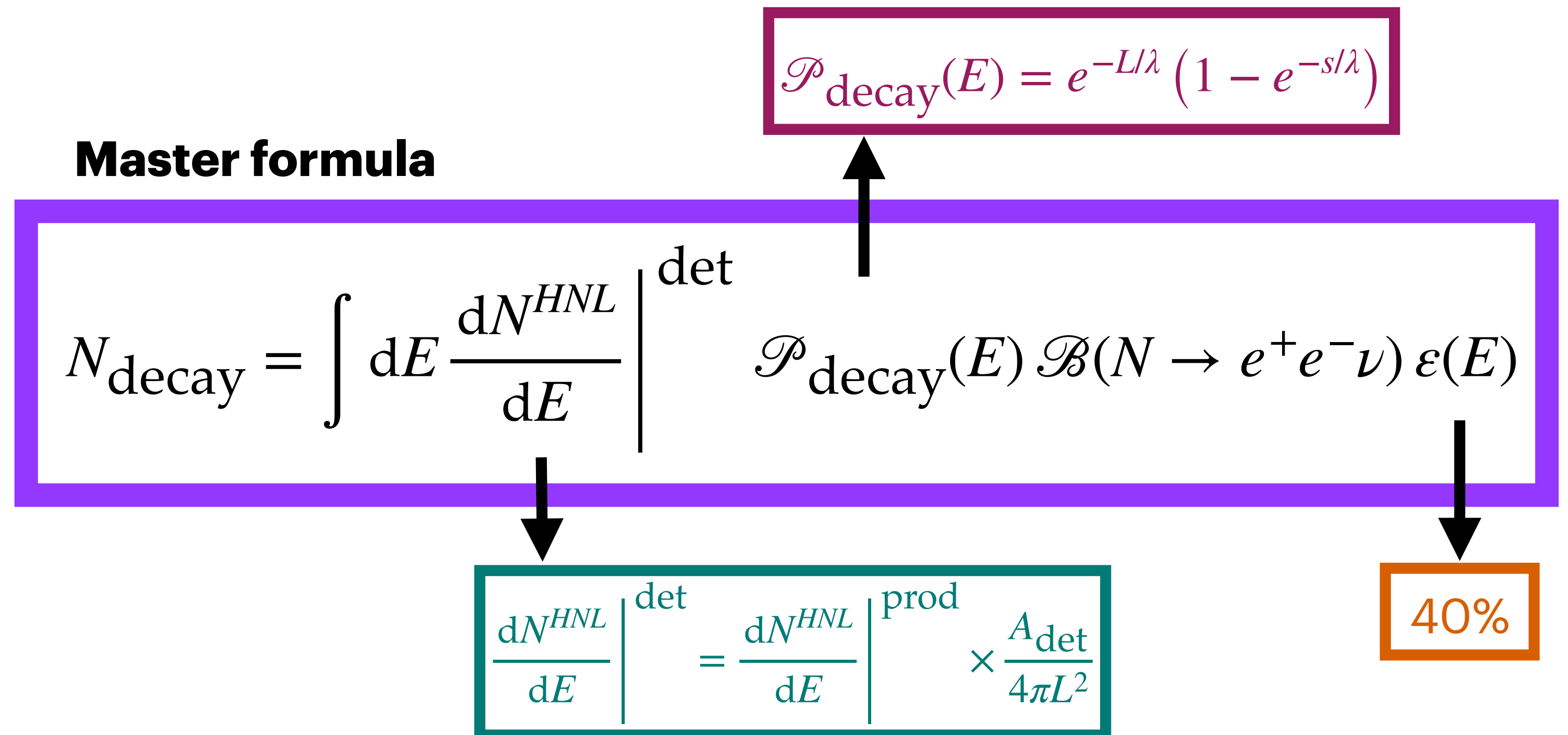
Both mixings



Detection pipeline

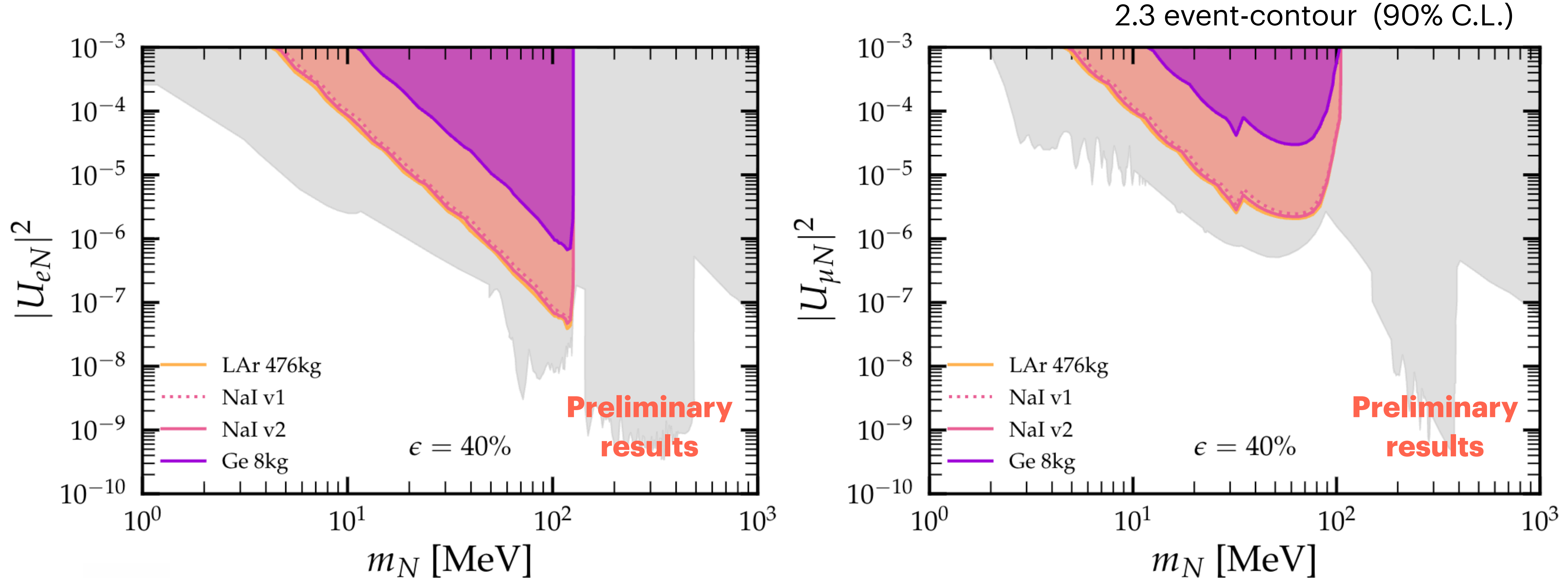
We seek a signal of a e^+e^- **pair**.

- Probability peaks for large L , given the long decay lengths for this parameter space.
- Isotropic flux decreases rapidly with distance L .
- Backgrounds considered through flat efficiency.



Even if decay lengths are large, **a shorter baseline** would increase the number of detected events.

Current sensitivities at COHERENT

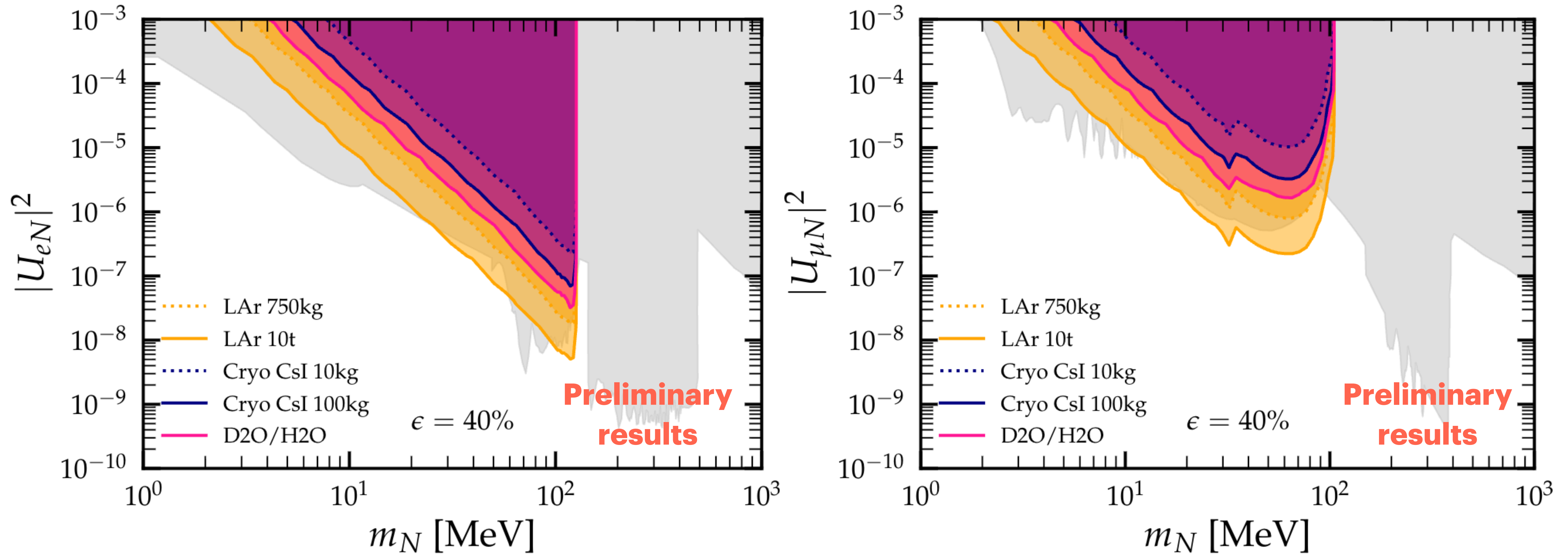


Reminder: this experimental setup was not originally built for this kind of process!

The **current** experimental setup **does not allow for improvement** of the current limits. They are still **competitive** with current limits at certain mass range and with conservative efficiency.

Future sensitivities at COHERENT

2.3 event-contour (90% C.L.)



The **proposed** experimental setup will **allow for improvements** to the current limits **by almost an order of magnitude** at certain masses of the parameter space.

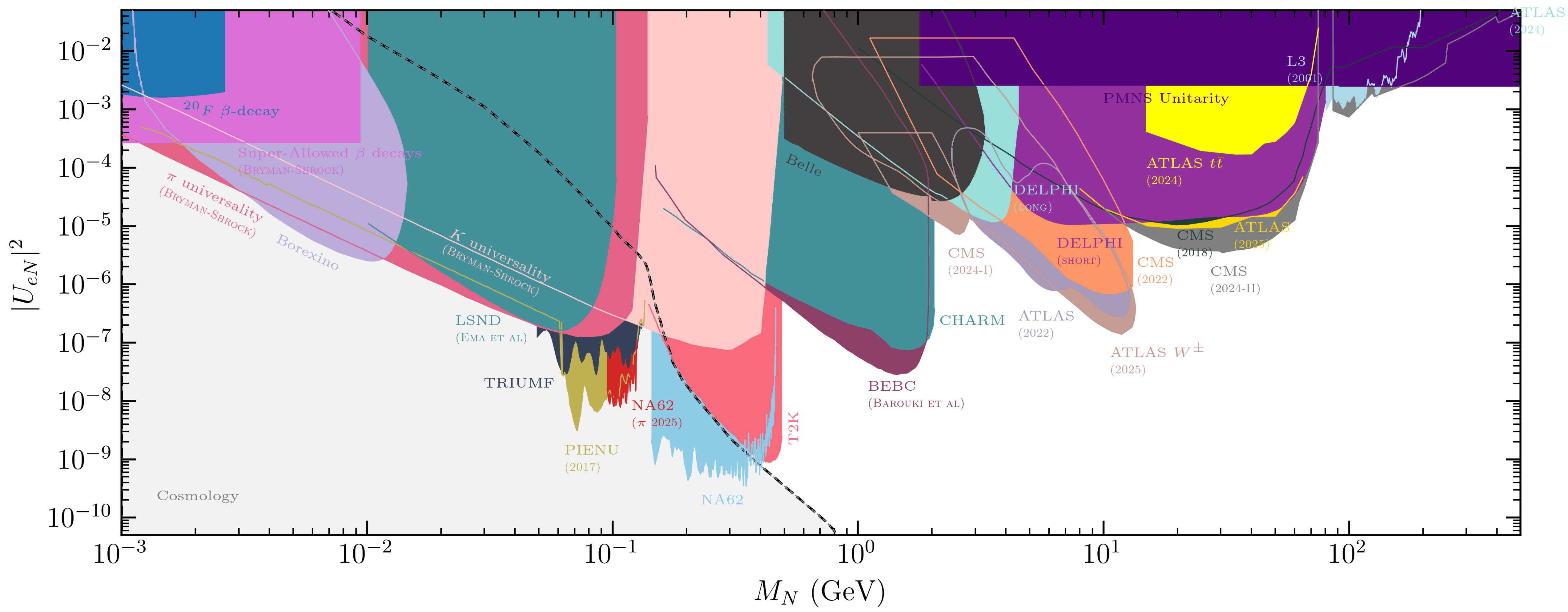
Conclusions

- ◆ HNLs are good candidates for **neutrino masses** explanation.
- ◆ HNLs could be **potentially produced** at Spallation Sources.
- ◆ Focus on SNS and COHERENT experiment:
 - ◆ Current experiments with **limited** capability for improving limits, but **competitive**.
 - ◆ Future proposals **potentially improve** the current limits by one order of magnitude at certain mass scales.
- ◆ **Similar analysis** will be done in other facilities like J-PARC, CSNS and the future ESS.
- ◆ Flat efficiency considered. **Further analysis of backgrounds needed.**

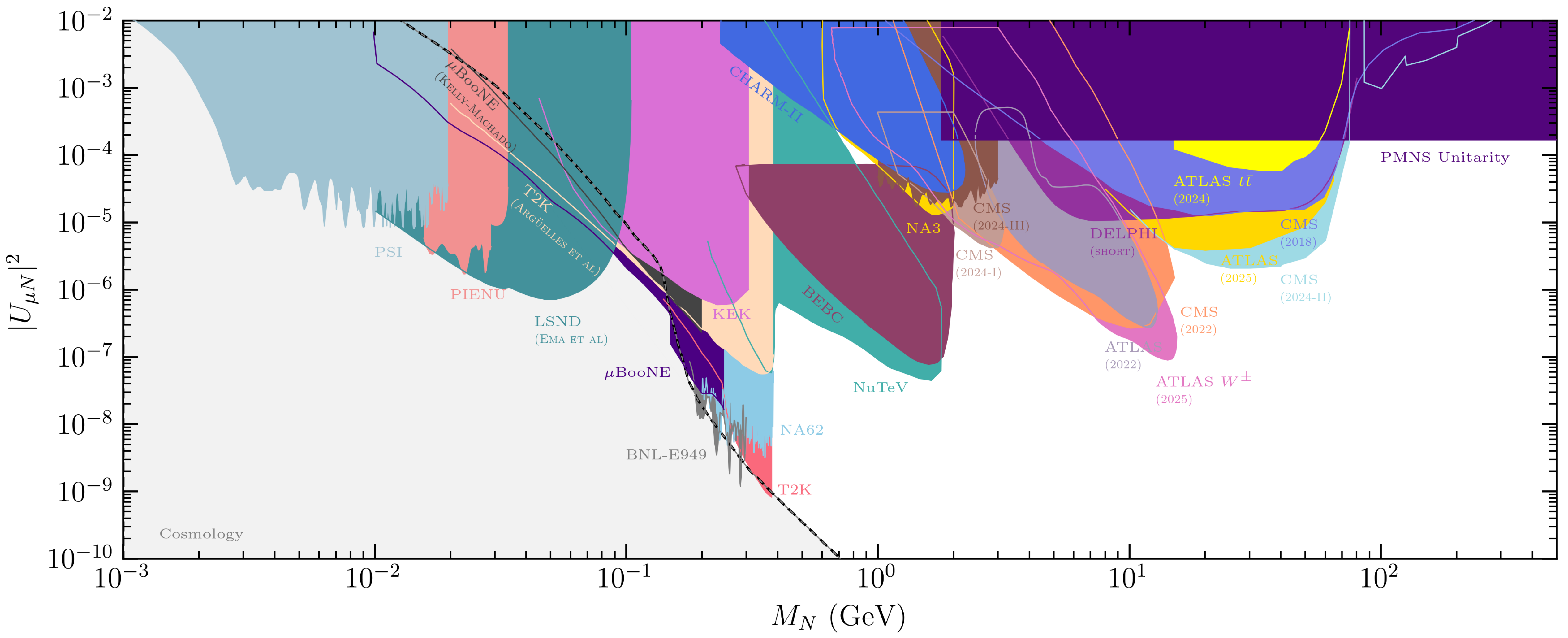
Given COHERENT **experimental baselines and volumes, dedicated searches** of HNLs would be **feasible**.

Backup

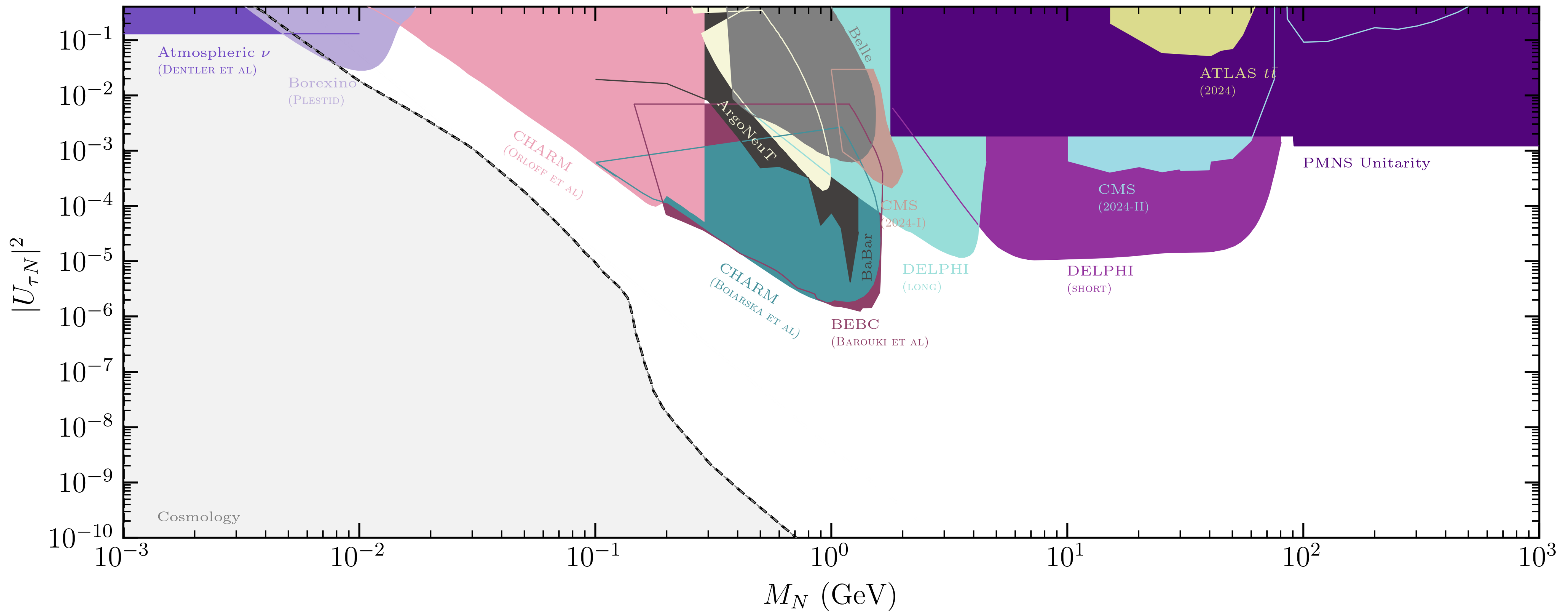
HNL limits: electron mixing case



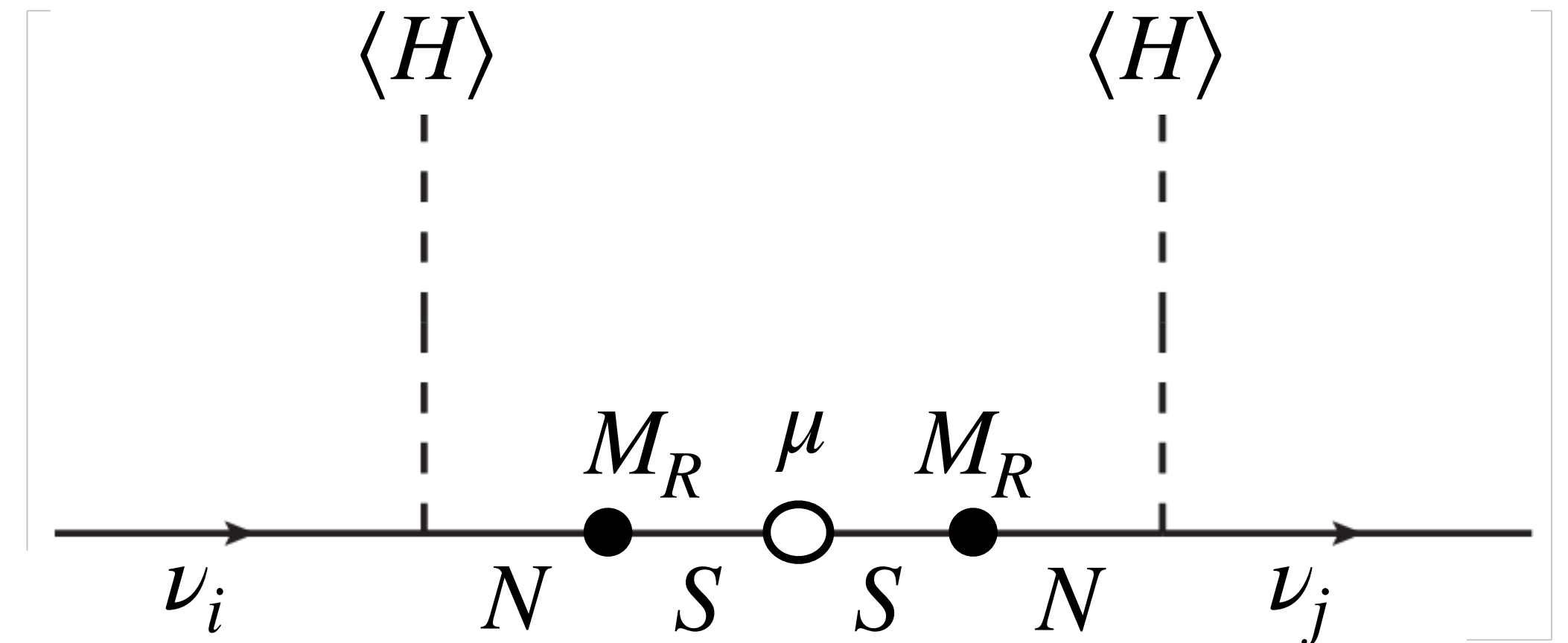
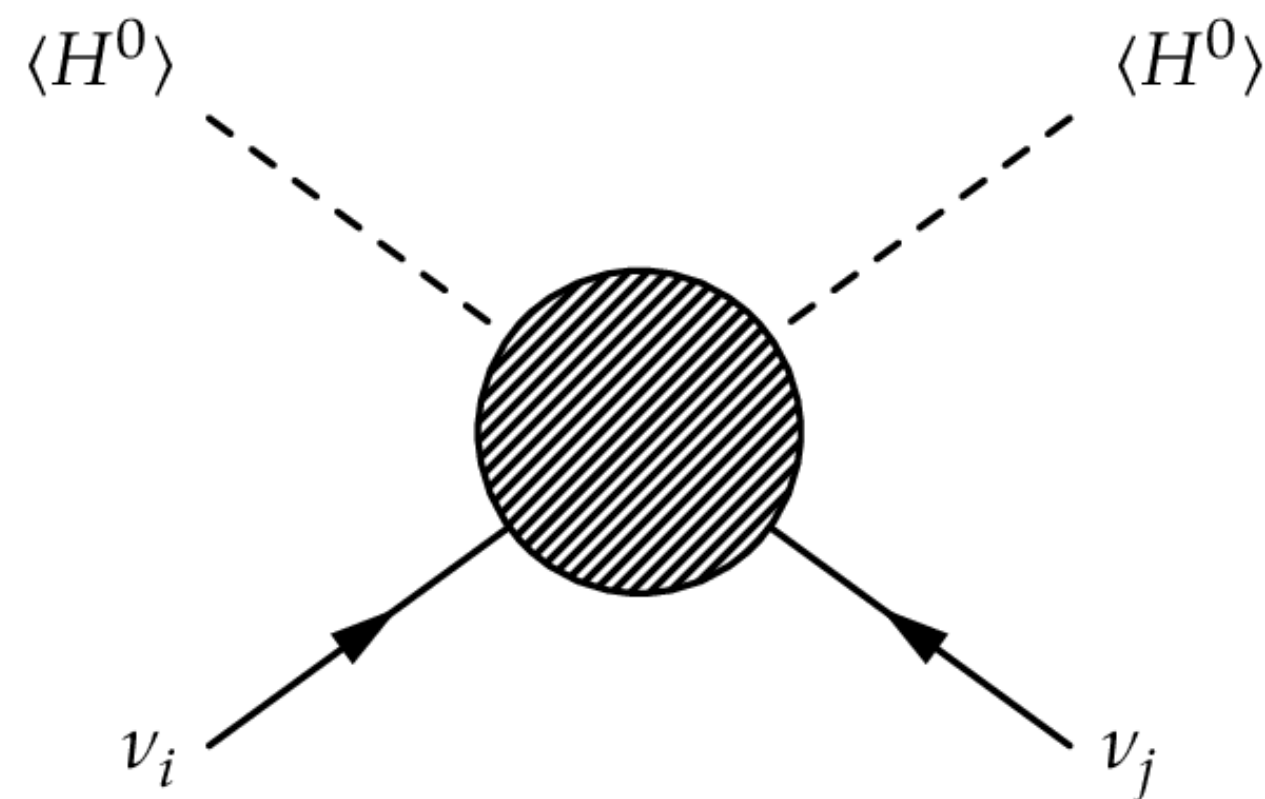
HNL limits: muon mixing case



HNL limits: tau mixing case



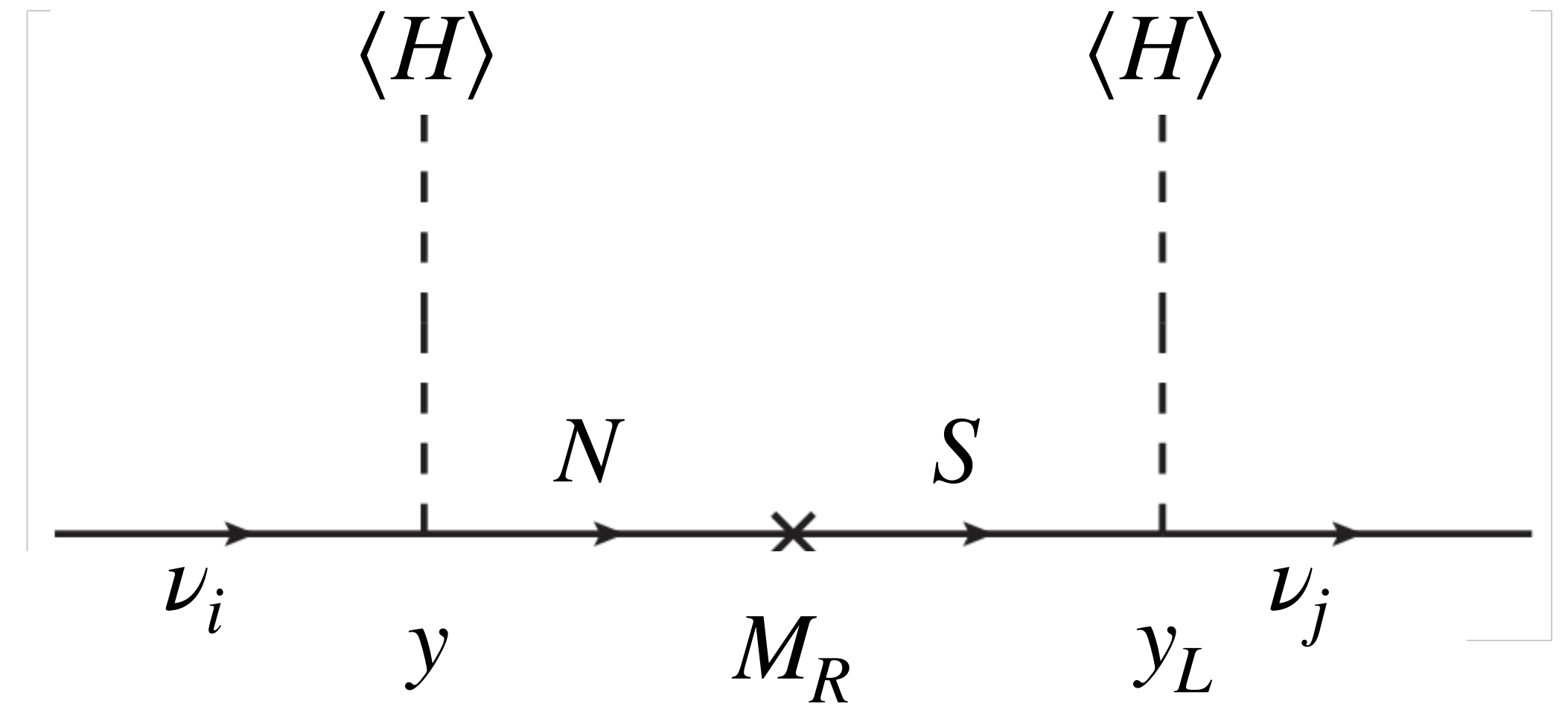
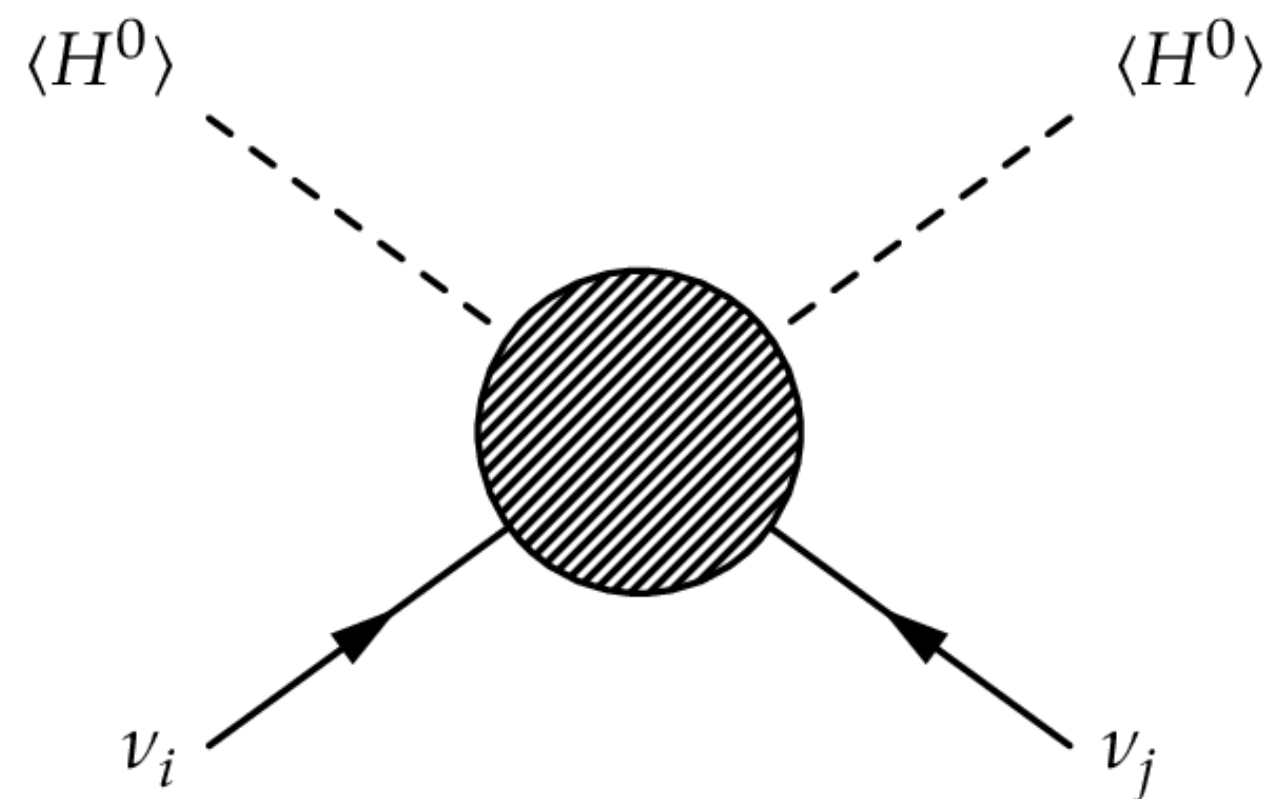
Inverted Seesaw



$$\mathcal{M}_{ISS} = \begin{matrix} & \nu_L & N & S \\ \begin{pmatrix} 0 & m_D^T & 0 \\ m_D^T & 0 & M_R \\ 0 & M_R^T & \mu \end{pmatrix} & \left\{ \begin{array}{l} m_\nu \approx \frac{m_D}{M_R} \mu \frac{m_D}{M_R} \\ V_{\nu N} \approx \frac{m_D}{M_R} \end{array} \right. \end{matrix}$$

- ◆ Small neutrino mass by small μ .
- ◆ Large mixing automatically allowed.

Linear Seesaw



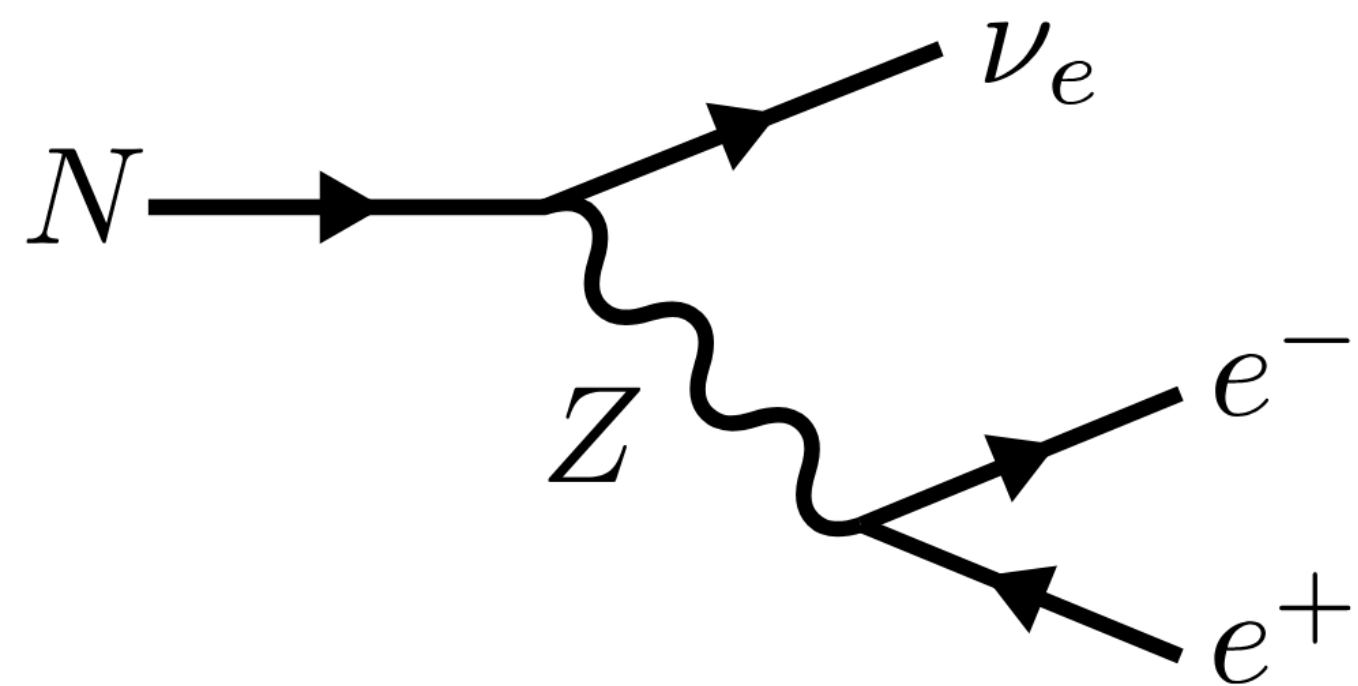
ν_L N S

$$\mathcal{M}_{LSS} = \begin{pmatrix} 0 & m_D^T & M_L^T \\ m_D^T & 0 & M_R \\ M_L^T & M_R^T & 0 \end{pmatrix} \begin{cases} m_\nu \approx (M_L M_R^{-1} m_D) + (\dots)^T \\ V_{\nu N} \approx \frac{m_D}{M_R} \end{cases}$$

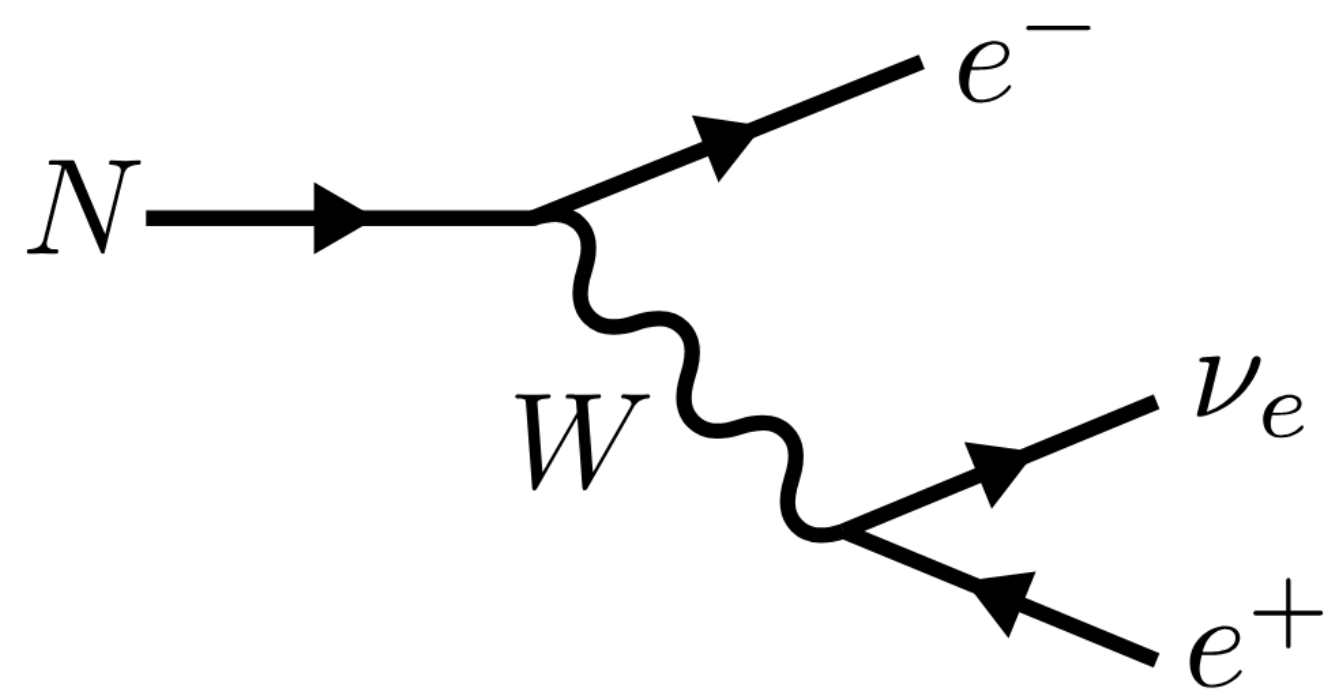
- ◆ Small neutrino mass by small M_L .
- ◆ Large mixing automatically allowed.

Decay width of HNLs

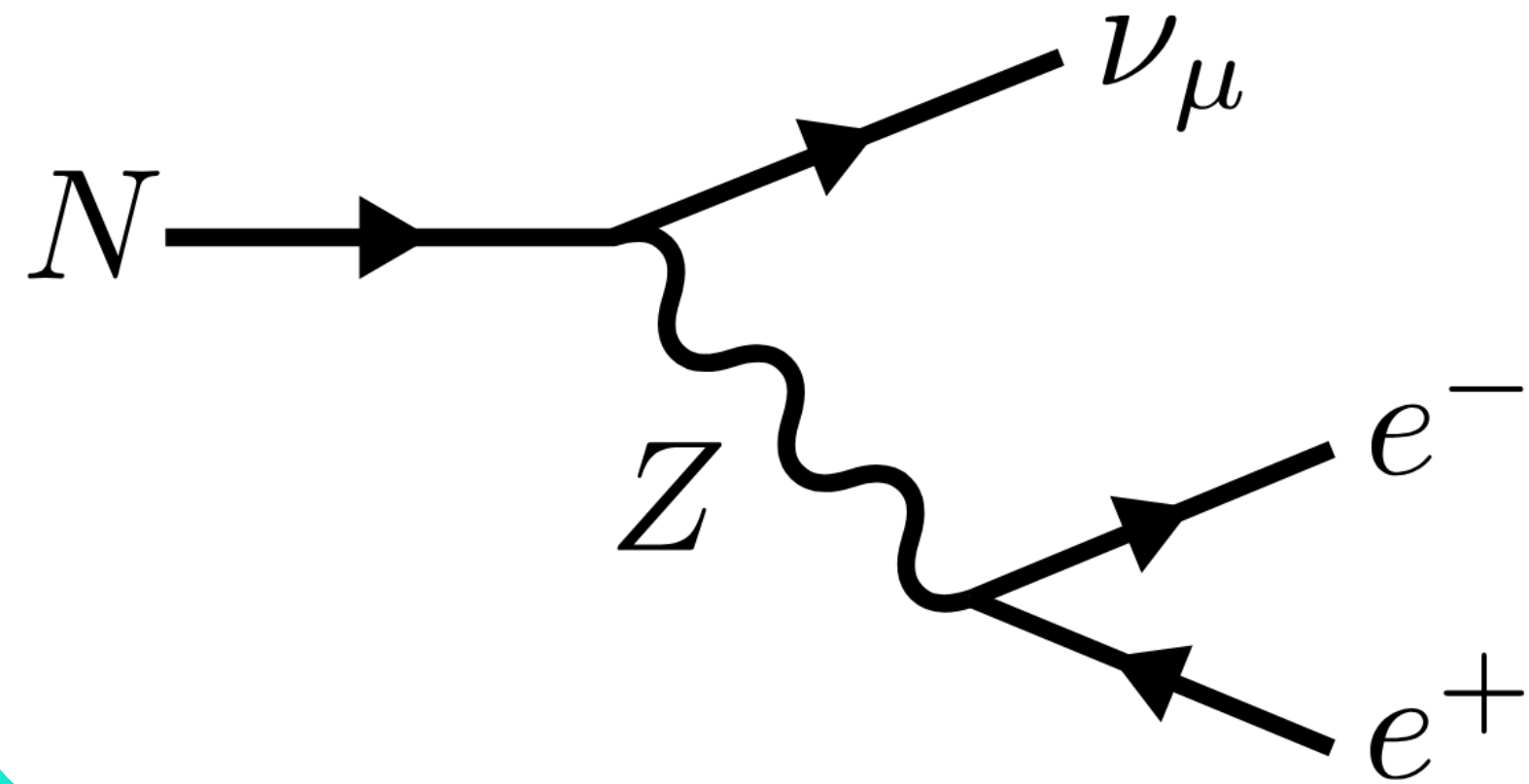
Only e mixing



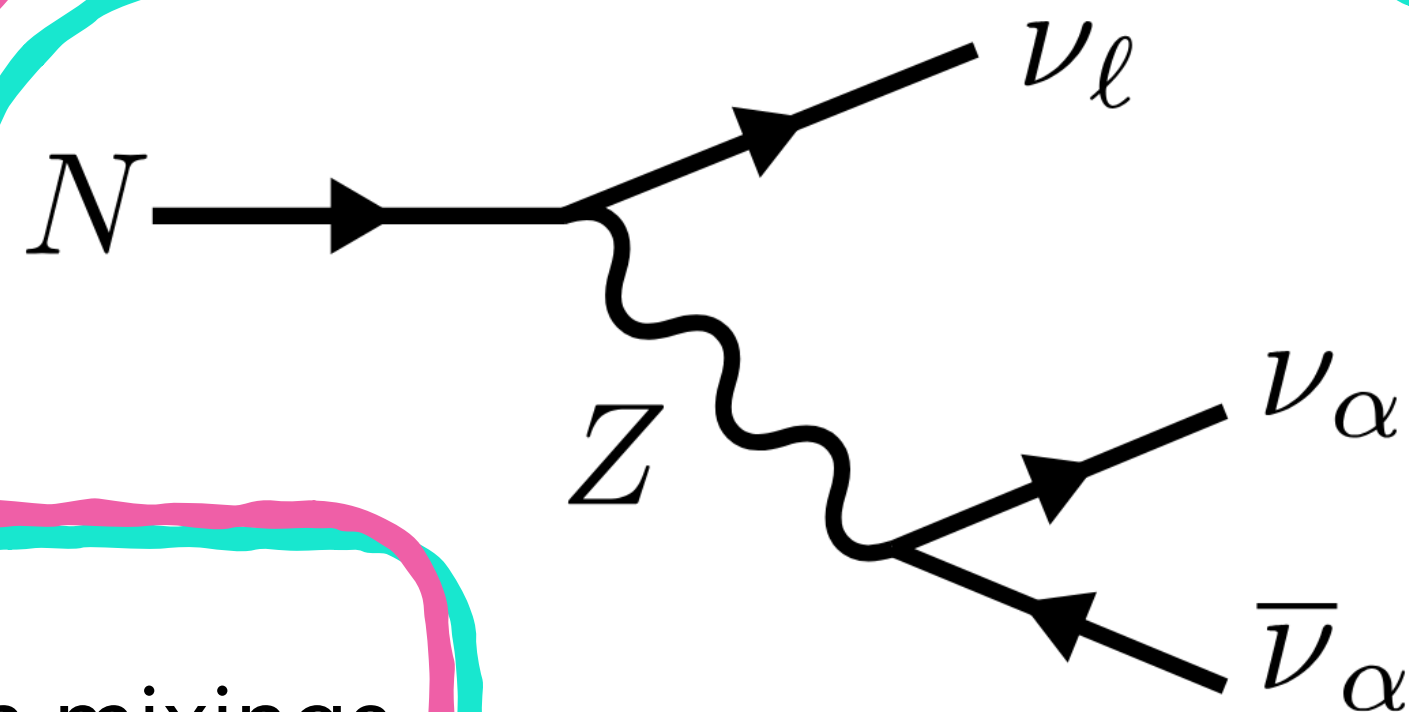
$$\Gamma(N \rightarrow \nu_e e^- e^+) = \frac{G_F^2 |U_{eN}|^2}{768\pi^3} m_N^5 \times (1 + 4 \sin^2 \theta_W + 8 \sin^4 \theta_W)$$



Only μ mixing



$$\Gamma(N \rightarrow \nu_\mu e^- e^+) = \frac{G_F^2 |U_{\mu N}|^2}{768\pi^3} m_N^5 \times (1 - 4 \sin^2 \theta_W + 8 \sin^4 \theta_W)$$



$$\Gamma(N \rightarrow \nu_l \nu_\alpha \bar{\nu}_\alpha) = \frac{G_F^2 |U_{lN}|^2}{768\pi^3} m_N^5$$

Both mixings

CE ν NS cross section

Neutrino cross sections

