

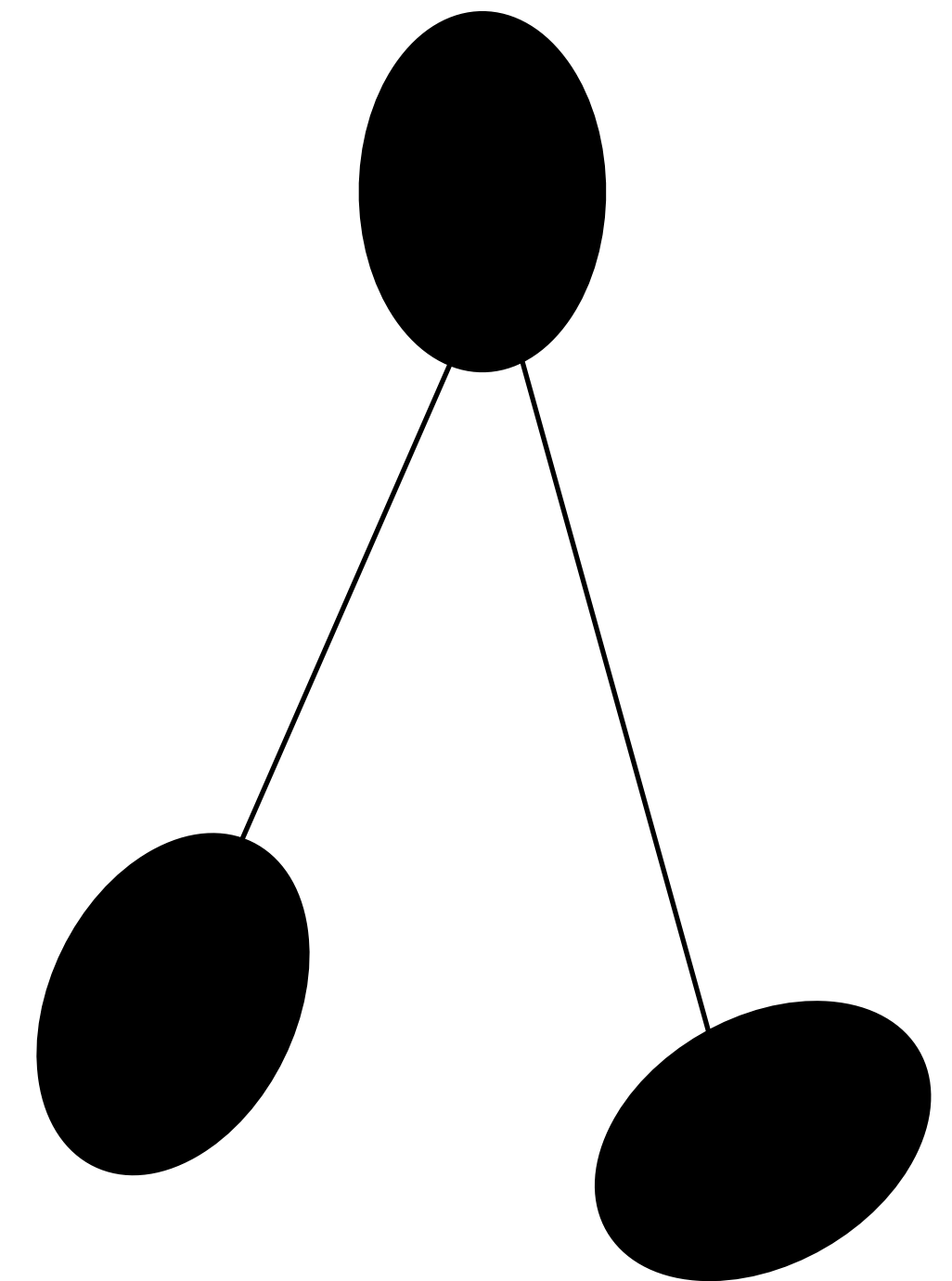


The Non-Gaussian Universe @ Heraklion, Crete

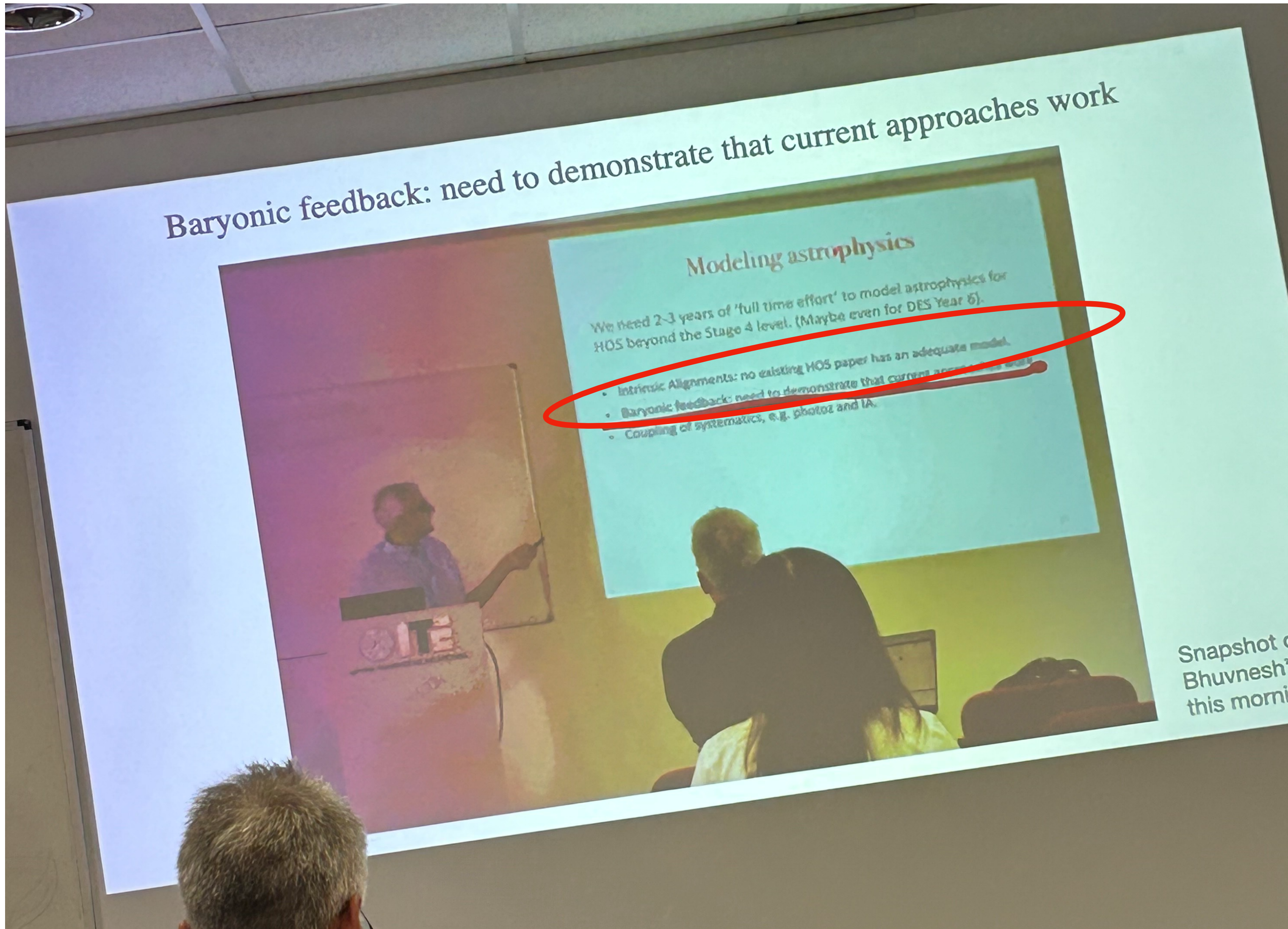
Three-point intrinsic galaxy alignments in FLAMINGO

Casper Vedder, PhD @ Leiden University

Based on [2601.17914](#) with Thomas Bakx, Elisa Chisari, Henk Hoekstra



Modeling astrophysics: IA

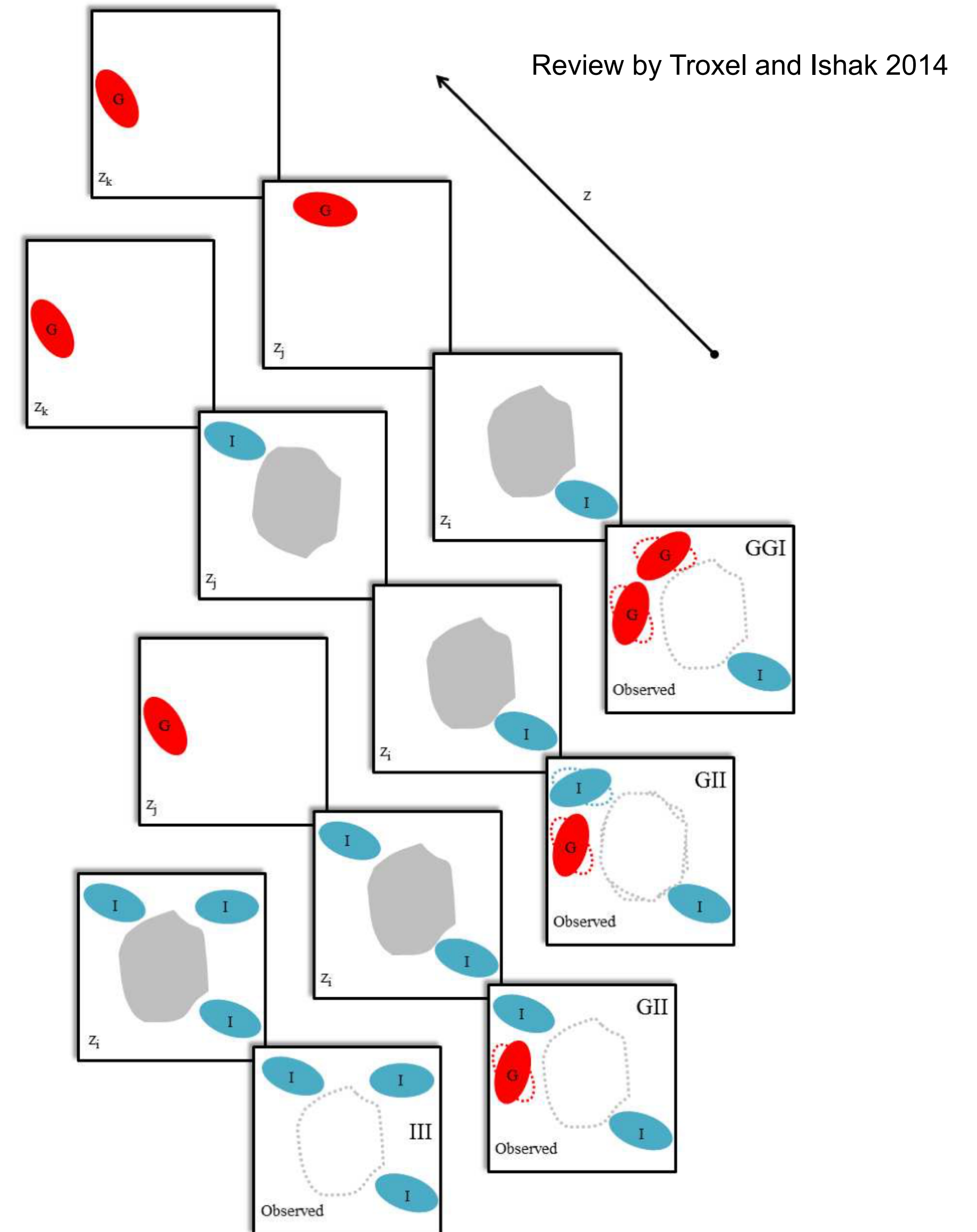


Snapshot of Bhuvnesh's talk in Giovanni's talk yesterday

Intrinsic Alignments

Challenges for (higher order) lensing*

- Tighter constraints give more stringent demands for modeling
- We specifically target non-linear information

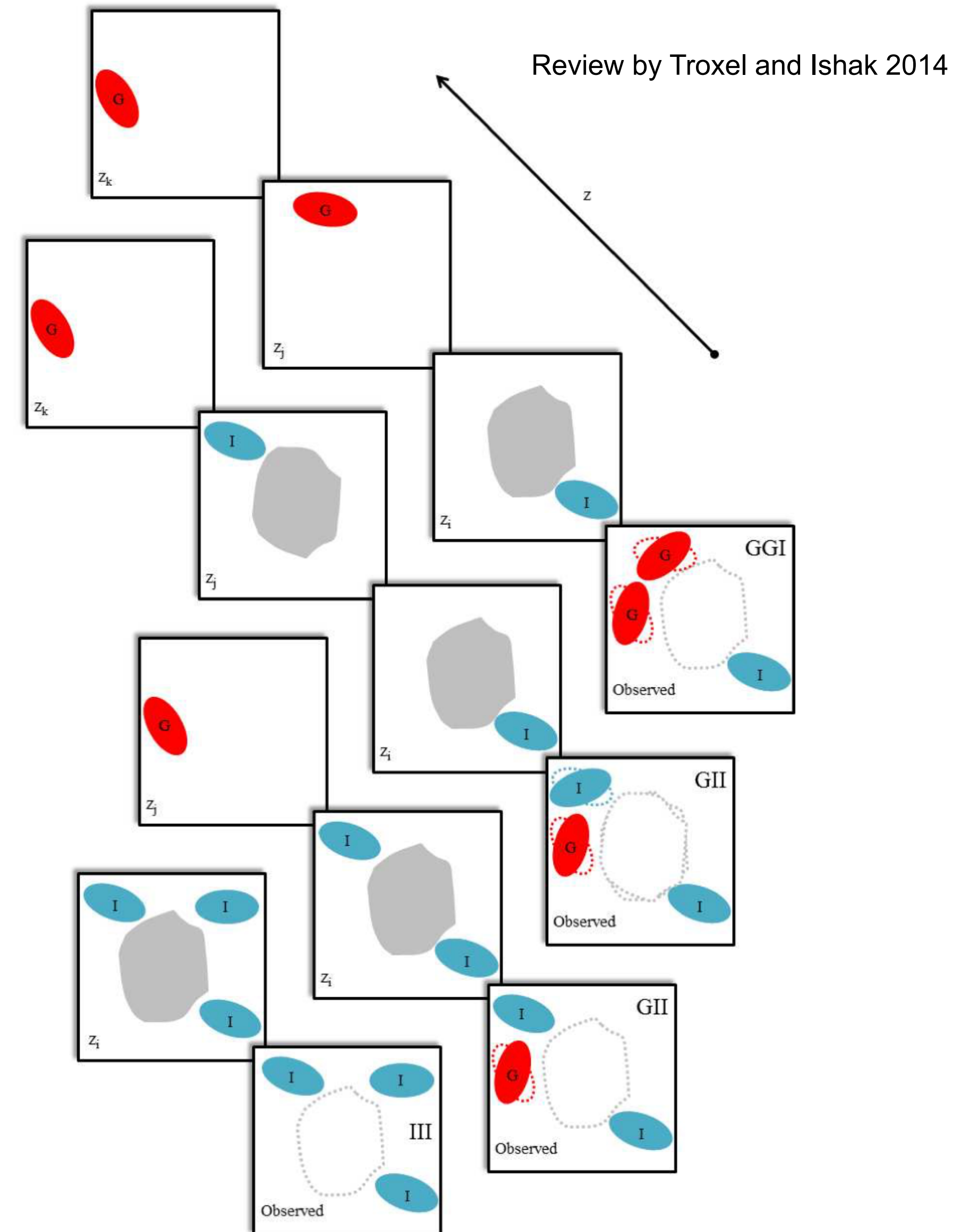


*but many more applications of IA!

Intrinsic Alignments

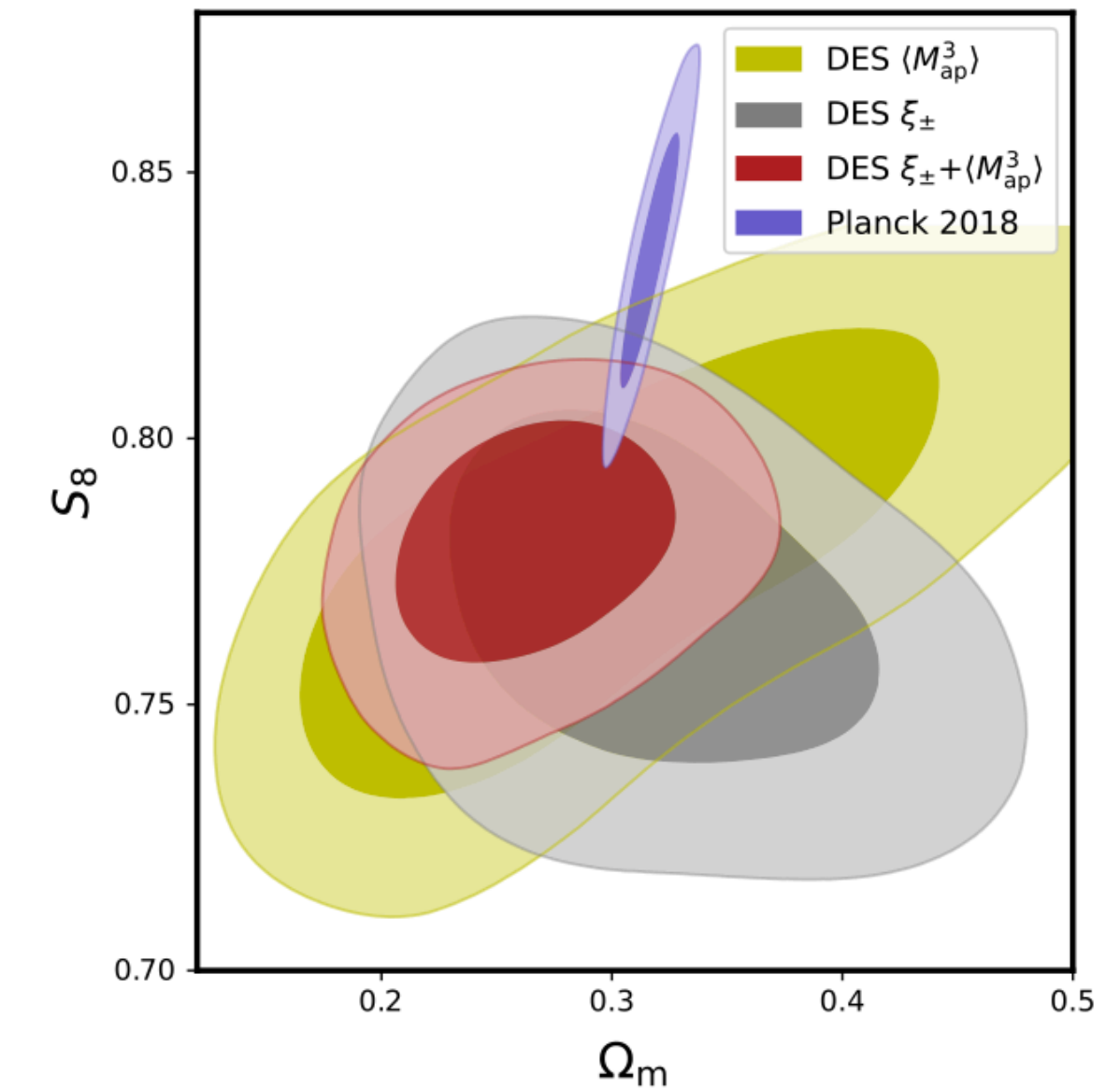
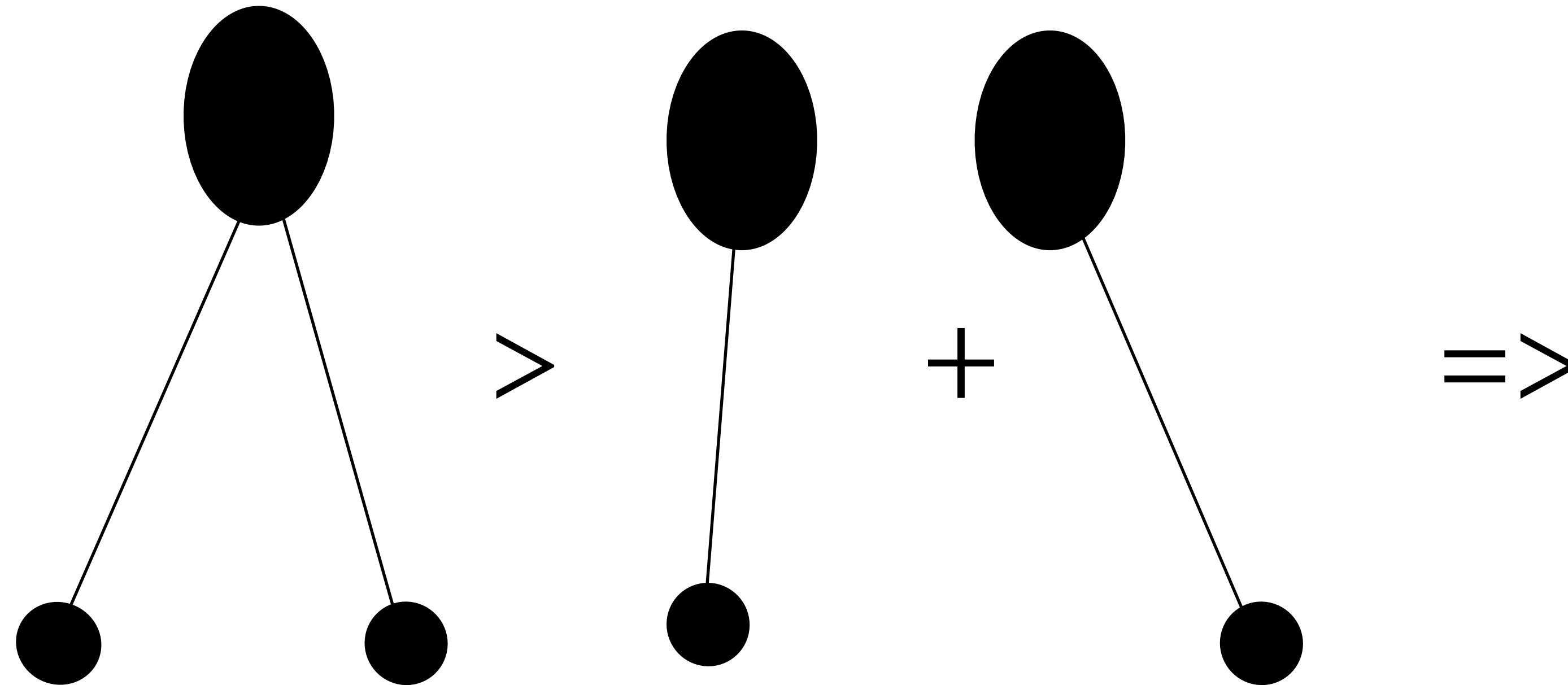
Challenges for (higher order) lensing*

- Tighter constraints give more stringent demands for modeling
- We specifically target non-linear information
 - Models from 2-pt. might not suffice
 - Important to test our models and assumptions



*but many more applications of IA!

Three-point statistics and cosmic shear

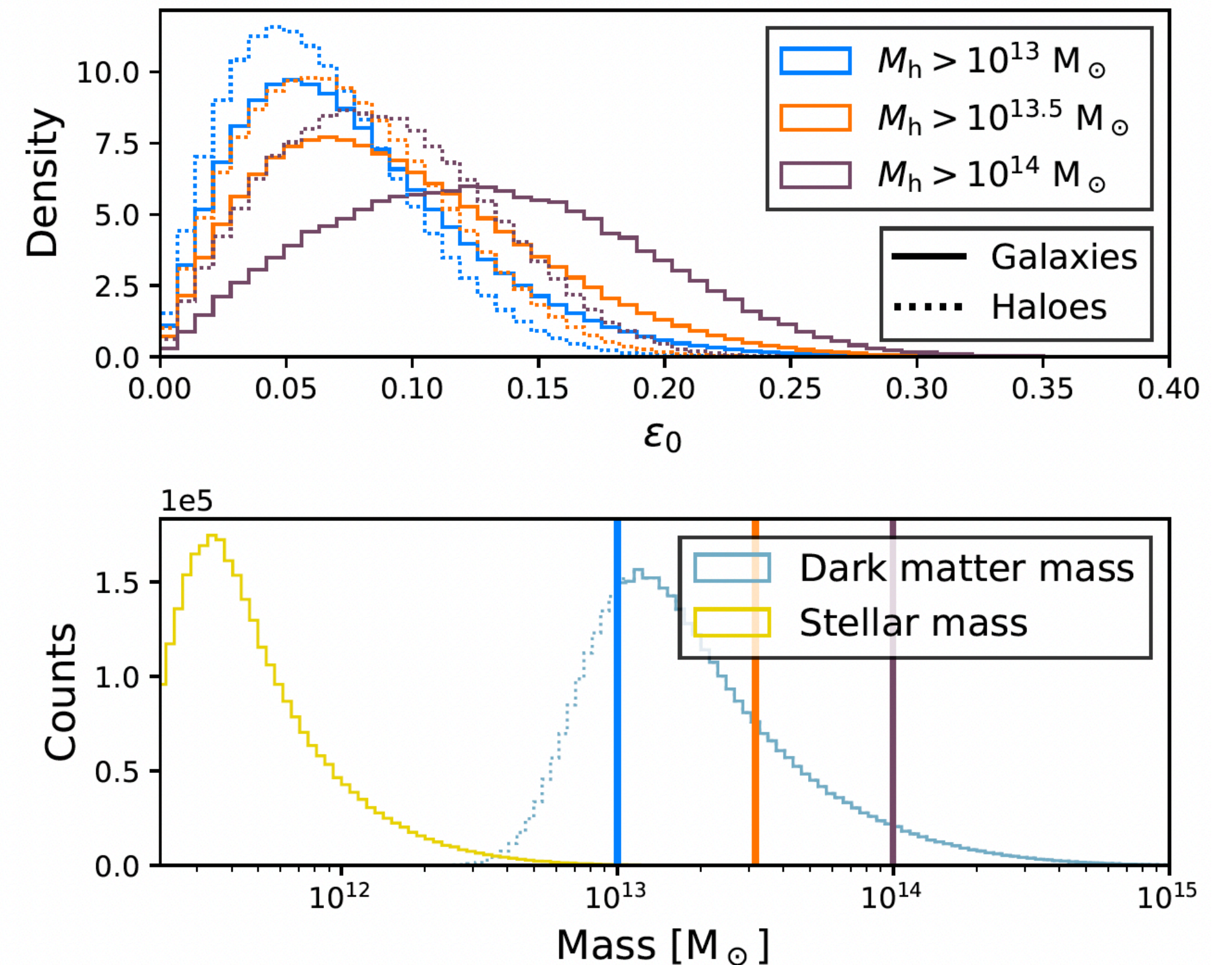


DES y3: Gomes et al, 2025, [2508.14018](https://arxiv.org/abs/2508.14018)

- Benefit of 3-pt: **analytical modeling** on large scales (EFT) + **some intuition**
- Obtained directly from a shape catalogue
- Many insights and challenges will carry over, so natural bridge to higher order stats!

FLAMINGO

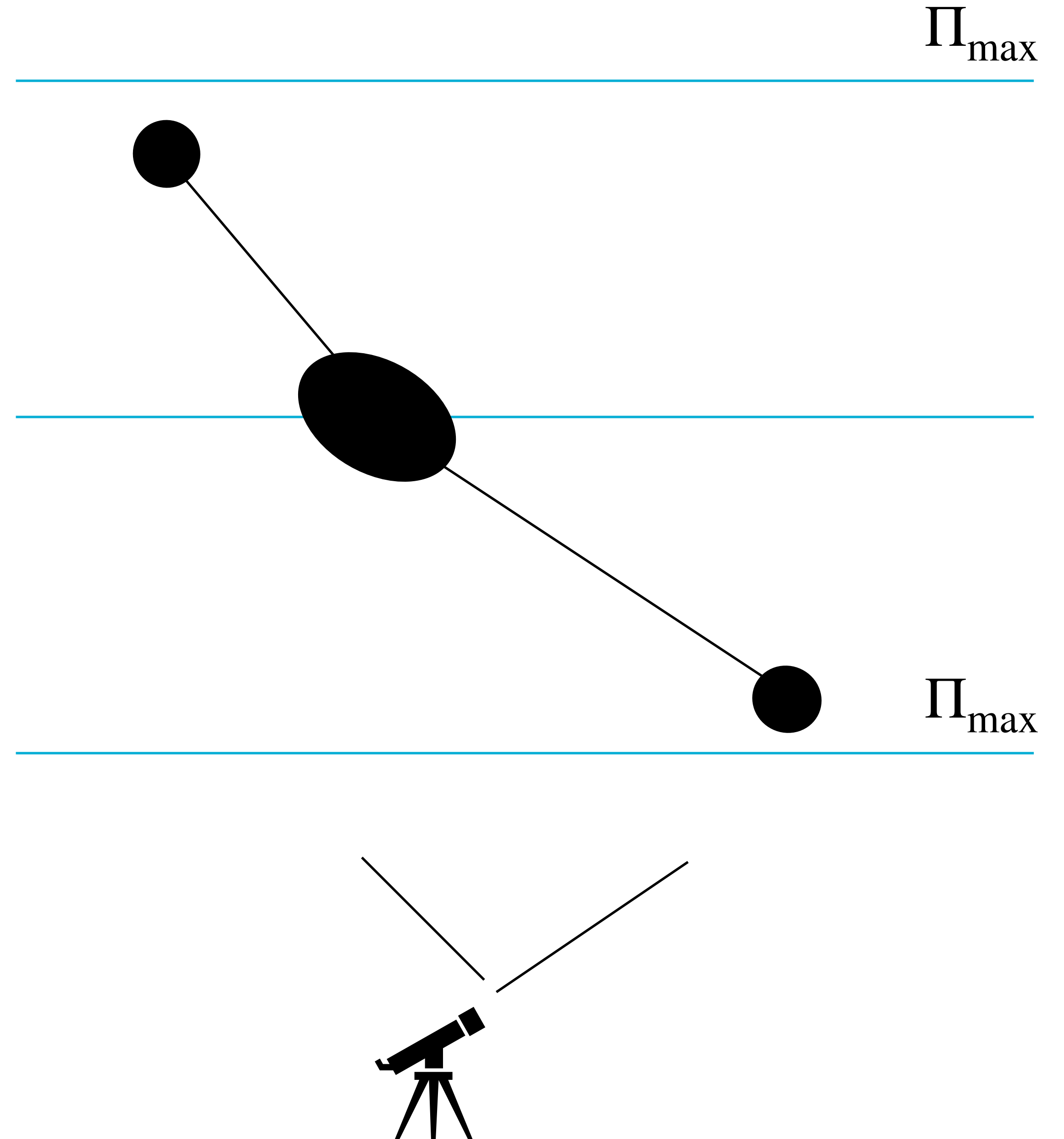
- One of the largest hydro sims to date
 - $(2.8 \text{ Gpc})^3$, 'm9 resolution'
- Construct a galaxy catalogue of about 3 mln LRG's
- Obtain shapes via inertia tensors and project into (ϵ_1, ϵ_2)
- Measure and model 3PCF of IA on this sample



3PCF of IA

- **Natural components:** complex valued correlation functions
- Isolate and project tracers within **LOS distance of < 20 Mpc** from one specific galaxy
- Measured efficiently with a multipole algorithm (Porth et al, 2024), using *TreeCorr* (Jarvis et al, 2004)
- g_{I} , g_{II} , g_{III}

$$g_{\text{I}} = \frac{DDS - 2DRS + RRS}{RRR}$$

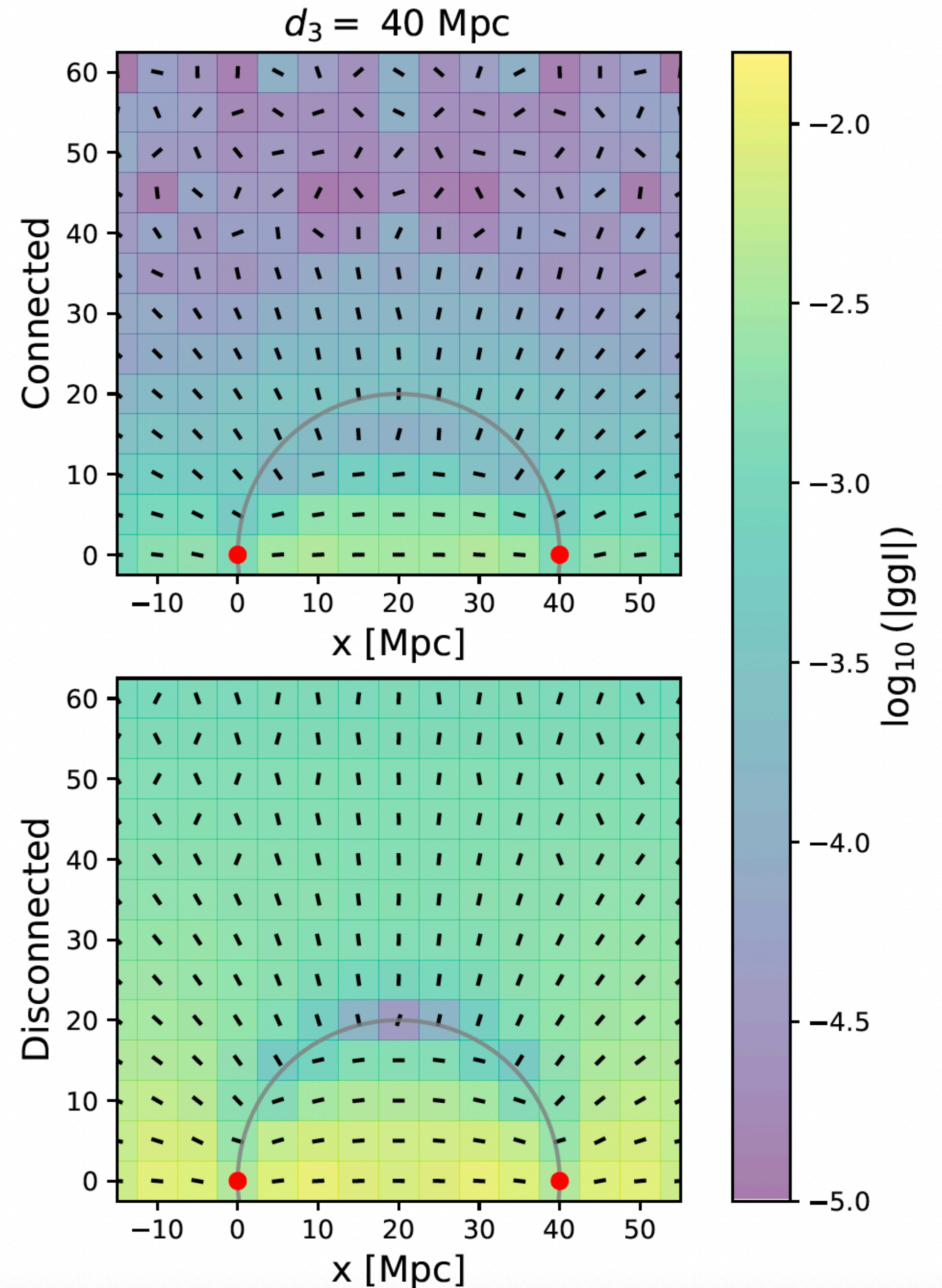
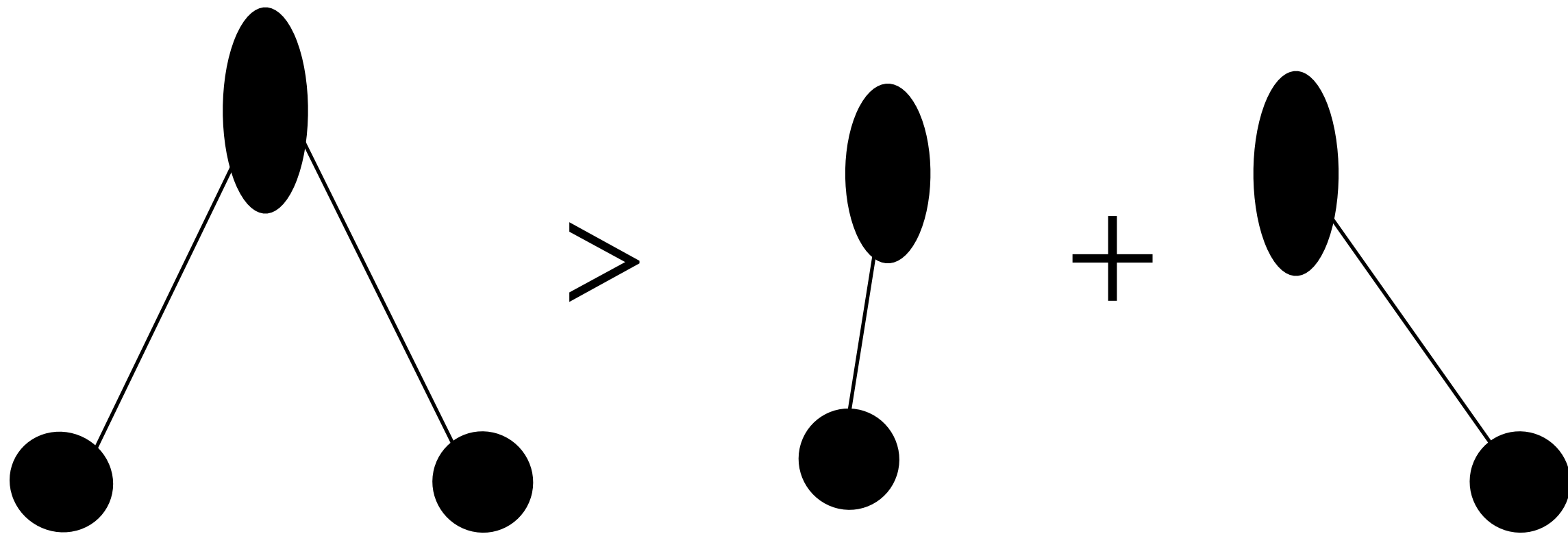


3PCF of IA

Example

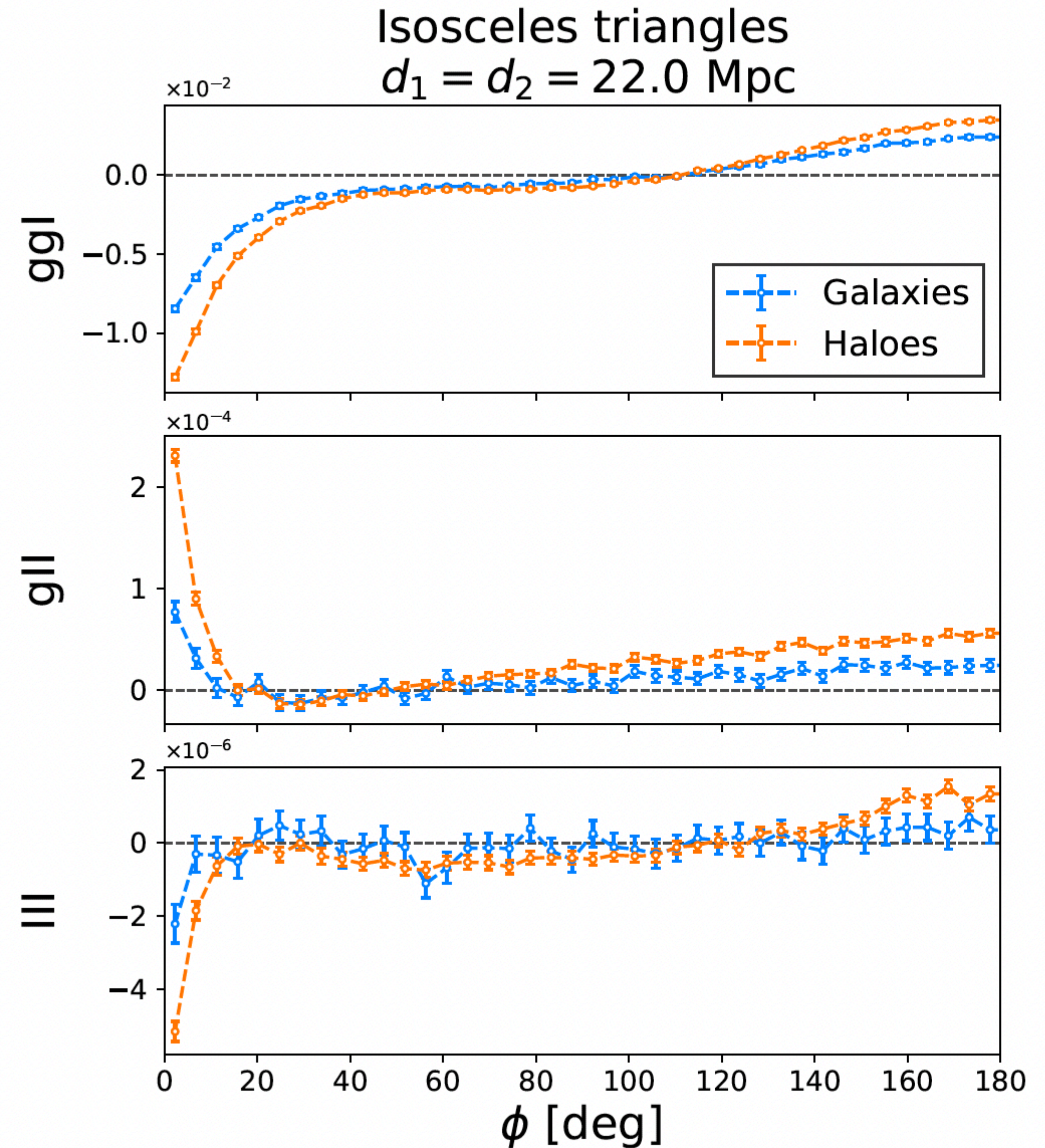
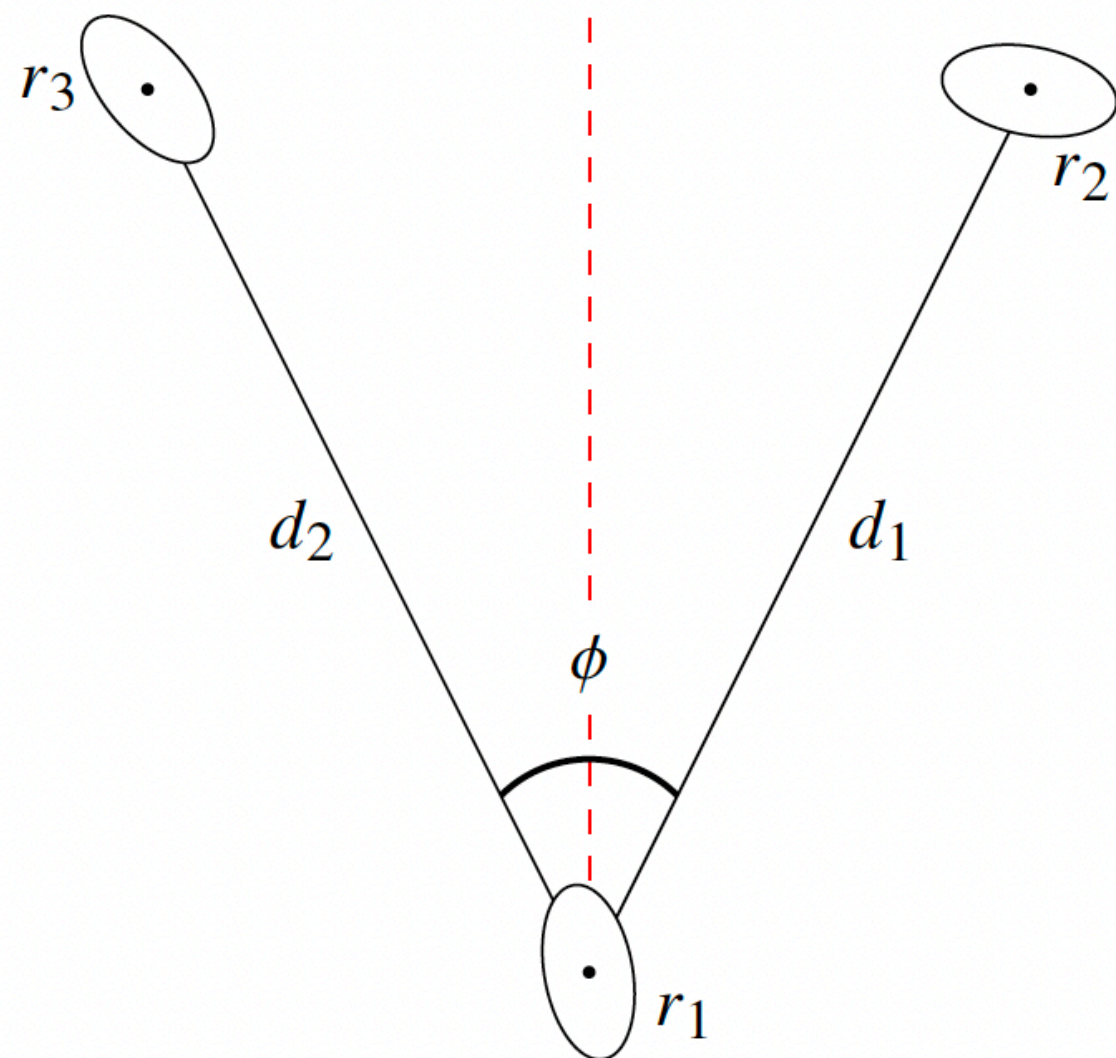
- “Excess alignment around galaxy pairs”

- $$ggI = \frac{DDS - 2DRS + RRS}{RRR}$$



3PCF of IA

- **Tangential components** of 3PCF
- Most signal in collinear triangles

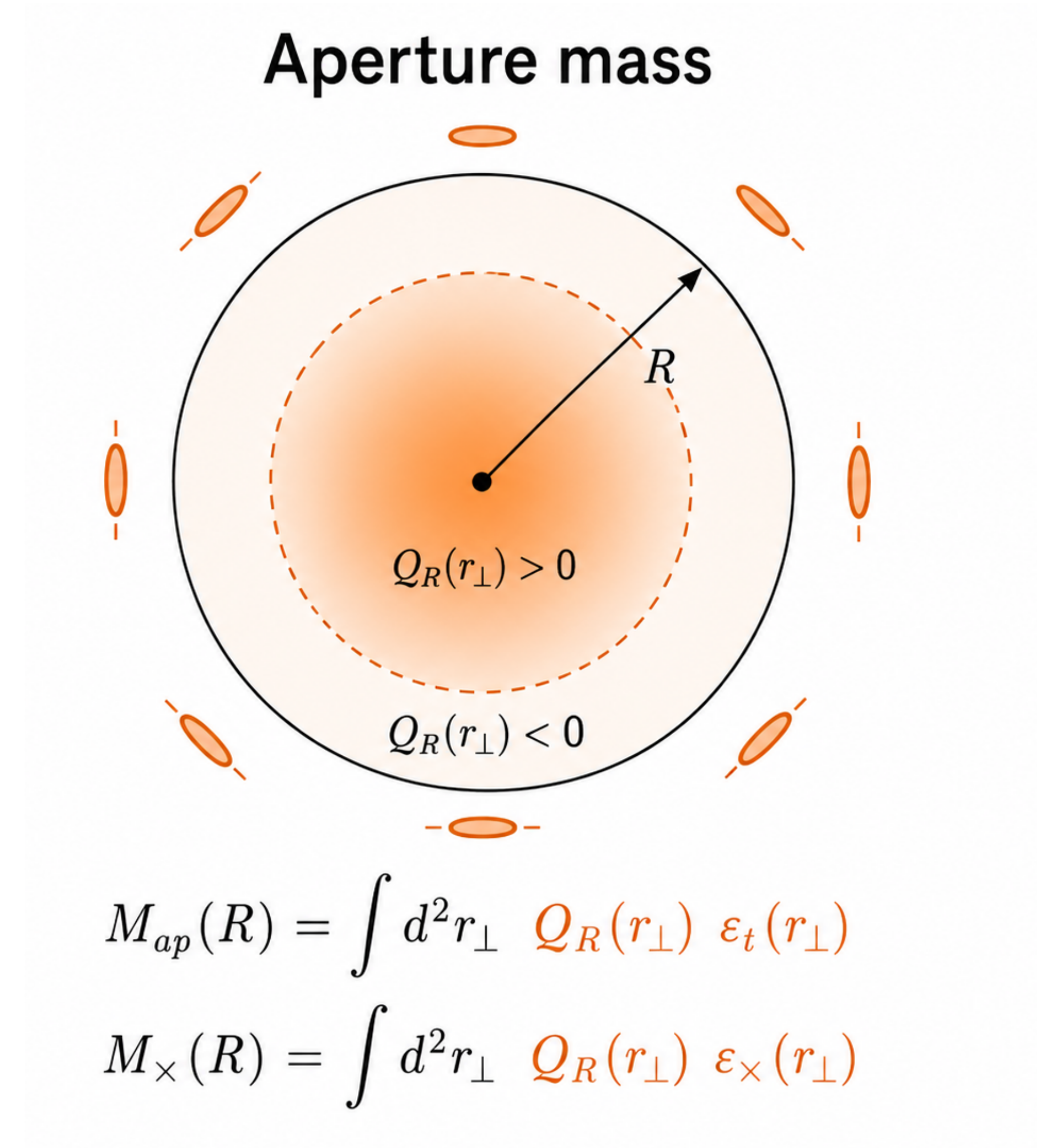


Aperture mass statistics

Practical compression

- 3PCF is rich in information, but somewhat unwieldy, apertures provide several benefits
 - Easier to model
 - **E/B separation**
 - Randoms subtraction automatically due to **compensated filter**
 - Relatively local probe of the bispectrum
- Integral transform of 3PCF, and easy to connect to the bispectrum

$$\bullet \langle NNM_{\text{ap}} \rangle \sim \int 3\text{PCF} \sim \int B_{\delta\delta E}$$



Modeling the bispectrum

Shapes and the bias expansion*

- Bias expansion

See review by Desjacques, Jeong and Schmidt (2018), see Vlah et al (2020) for IA

$$g_{ij} = \sum_{\mathcal{O}} b_{\mathcal{O}} \mathcal{O}_{ij} = b_K K_{ij} + b_{\delta K} \delta K_{ij} \dots$$

- Bias parameters encode our ignorance of galaxy evolution and formation
- Tree-level bispectrum is then obtained via $\langle \delta \delta g_{ij} \rangle \sim \langle \delta^{(1)} \delta^{(2)} g_{ij}^{(1)} \rangle + \langle \delta^{(2)} \delta^{(1)} g_{ij}^{(1)} \rangle + \langle \delta^{(1)} \delta^{(1)} g_{ij}^{(2)} \rangle$
 - So no $\langle \delta^{(1)} \delta^{(1)} g_{ij}^{(1)} \rangle$ due to Wicks theorem
 - **Hence, always non-linear bias terms**
 - Therefore we have $B_{mmE} \neq b_K B_{mmm}$ even at tree level. Additional bias terms contribute on all scales. This is fundamentally different from the 2-point function!

*In Fourier space also add stochastic contributions!

IA - Shape bias expansion

For bispectrum: Schmitz et al (2018), Vlah et al (2020), Bakx et al (2025a, 2025b)

- 4 operators relevant at tree-level, which can all be relevant
- **Models we use are a subset of this tree-level EFT**
 - **NLA**: Assume only b_K , expand matter-field to all orders
 - **EFT - no VS/TATT**: velocity shear set to 0
 - **EFT - LLB**, using co-evolutions relations based on linear Lagrangian bias

Akitsu et al (2023)

$$\begin{aligned} b_{KK} &= -b_K, \\ b_{\delta K} &= \left(b_1 - \frac{2}{3}\right)b_K, \\ b_t &= \frac{5}{2}b_K. \end{aligned}$$

$$g_{ij} \approx b_K K_{ij} + b_{\delta K} \delta K_{ij} + b_{KK} \text{TF} \left(K^2 \right)_{ij} + b_t t_{ij} + \dots,$$

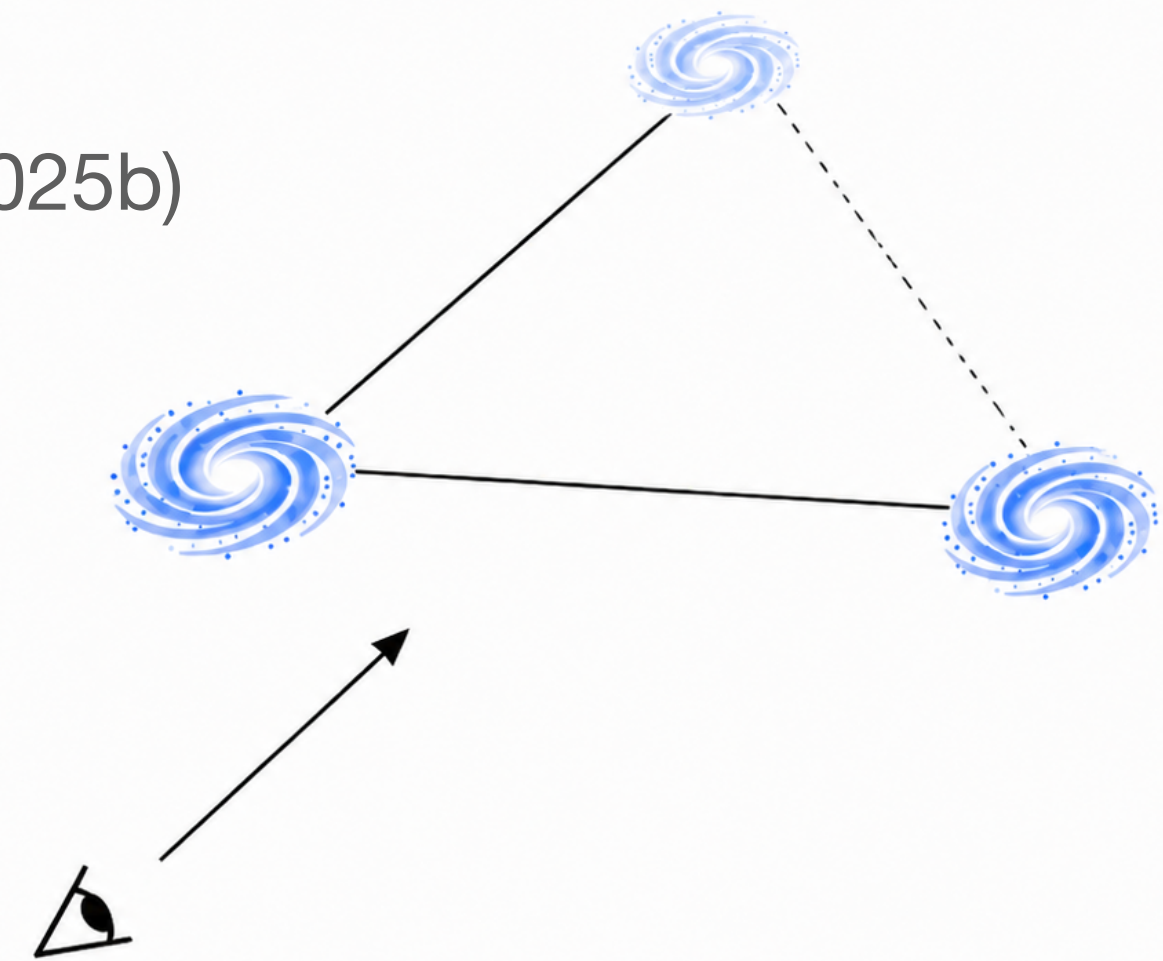
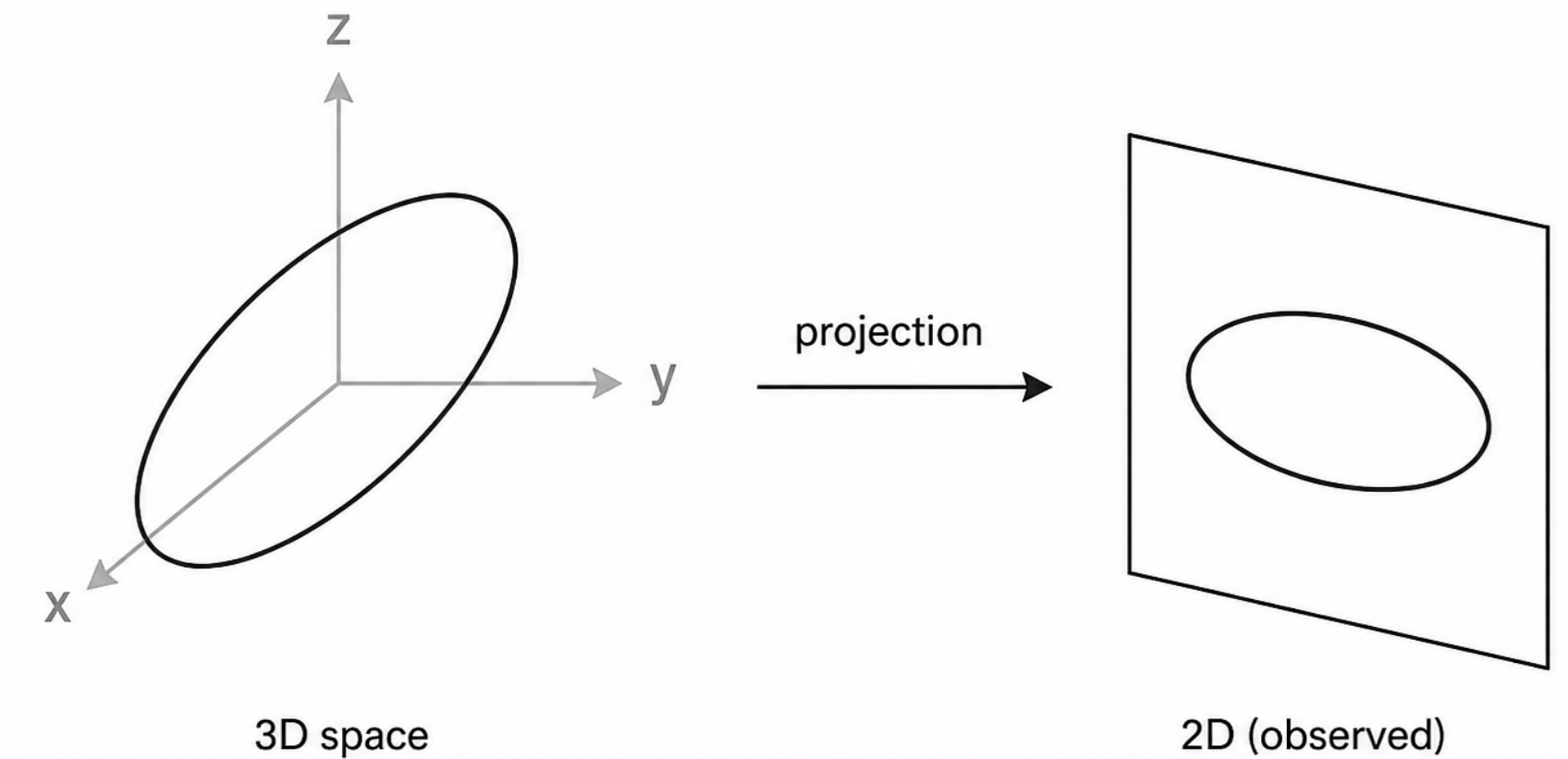
Projecting the shapes to the plane

- 3D shapes need to be projected to E and B modes

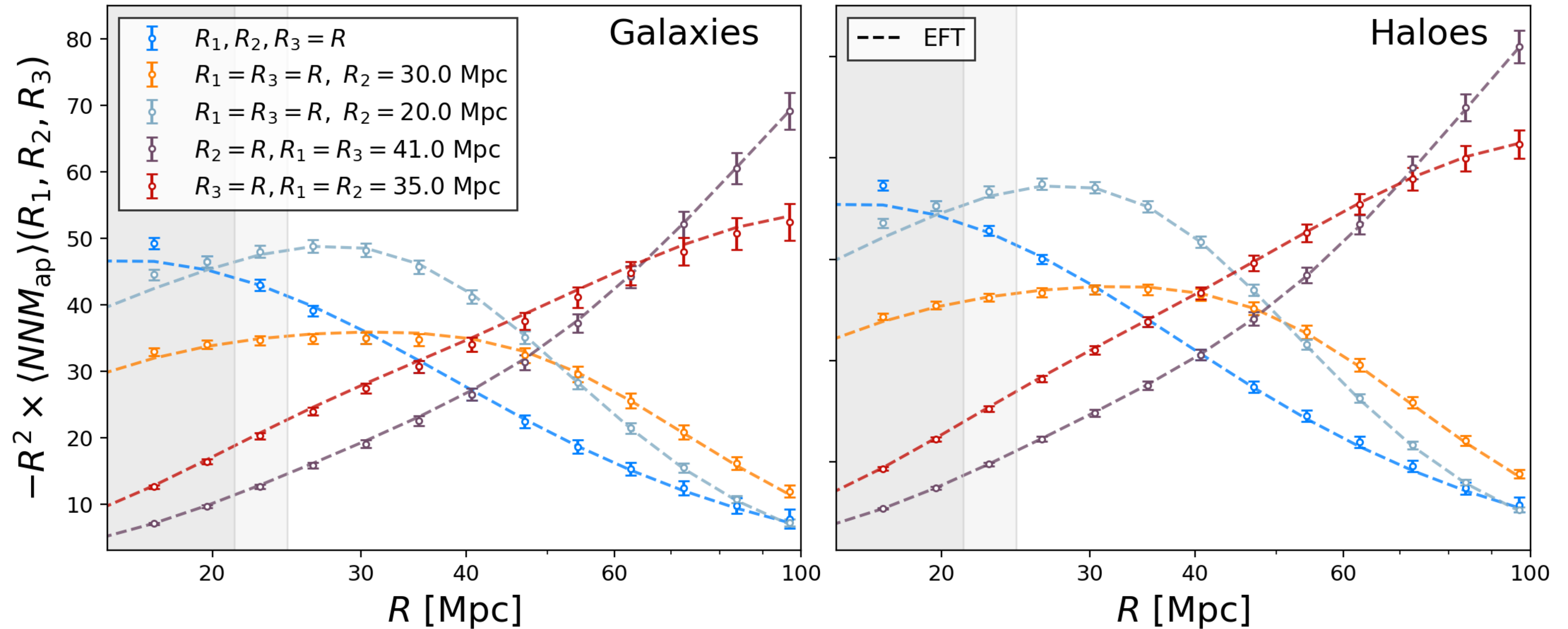
- $B_{\delta\delta g} \sim \langle \delta\delta g_{ij} \rangle \rightarrow B_{\delta\delta E} \sim \langle \delta\delta \epsilon_E \rangle \rightarrow \langle NNM_{\text{ap}} \rangle$

- Triangle depends on 5 variables, 3 k's and 2 angles for the line of sight dependence

Bakx et al (2025a, 2025b)



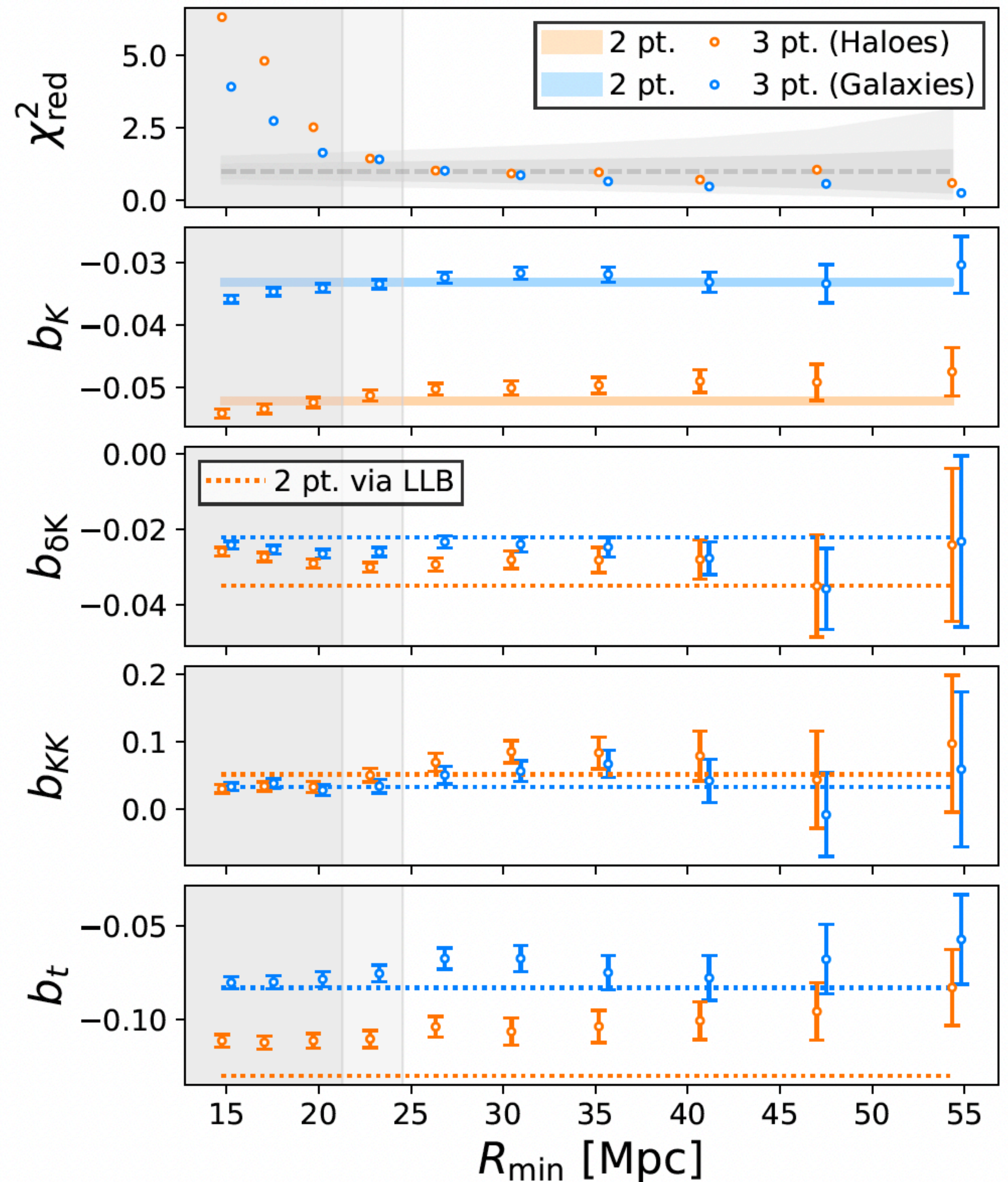
EFT vs data (1)



- Correlate shapes with matter instead of galaxies position to avoid galaxy biasing
- Unequal aperture scale add information and are needed to break degeneracies

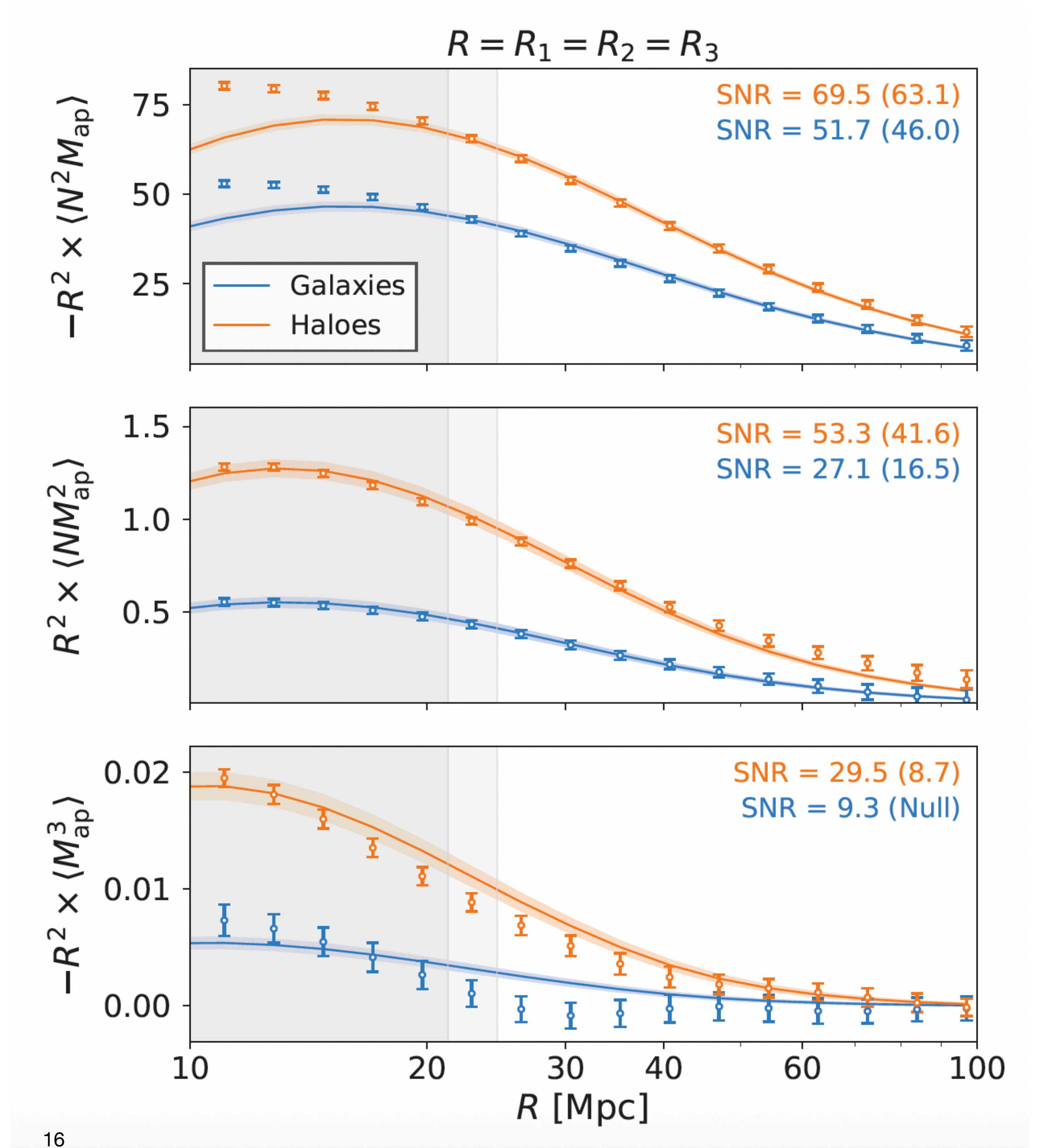
EFT vs data (2)

- All parameters are detected at several sigma
- EFT fits the data well until 22 Mpc
- Matches large scale 2 pt. b_K
- Co-evolution relations capture general trends



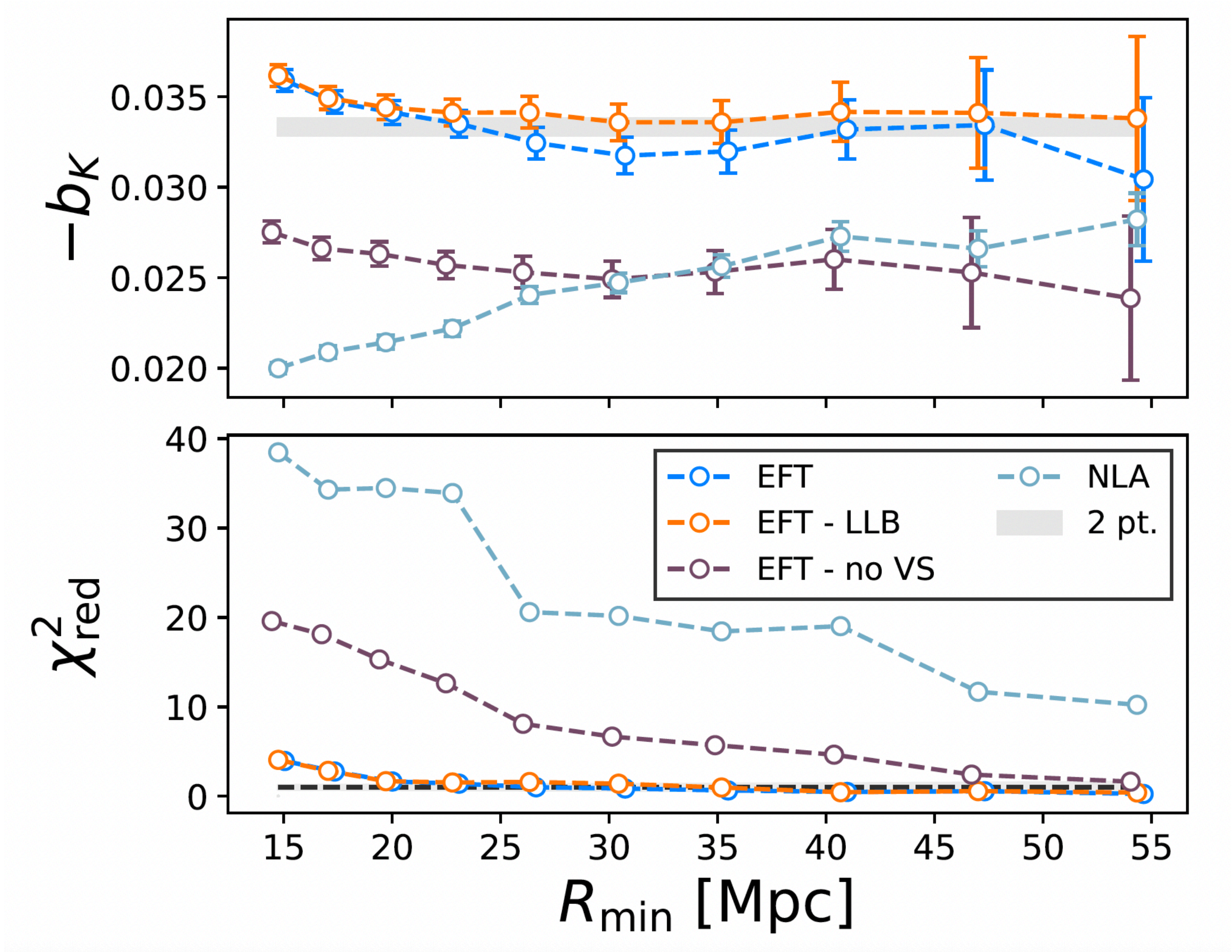
EFT vs data (3)

- Best fitting $\langle NNM_{\text{ap}} \rangle$ parameters consistent with $\langle NM_{\text{ap}}^2 \rangle$ and $\langle M_{\text{ap}}^3 \rangle$



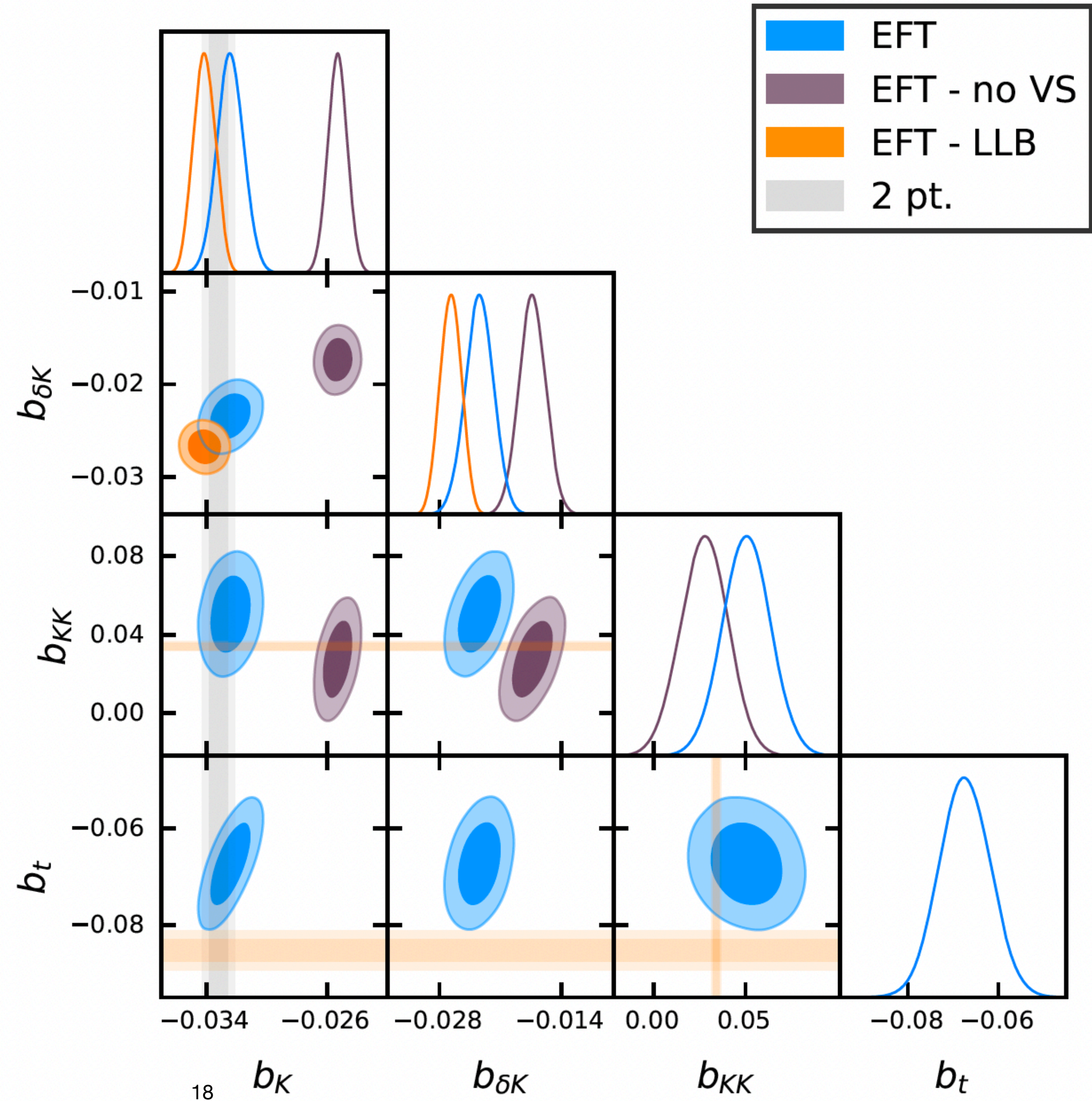
Model comparison

- All parameters relevant
 - **Also at large scales**
- Co-evolution gives a similar fit to data!

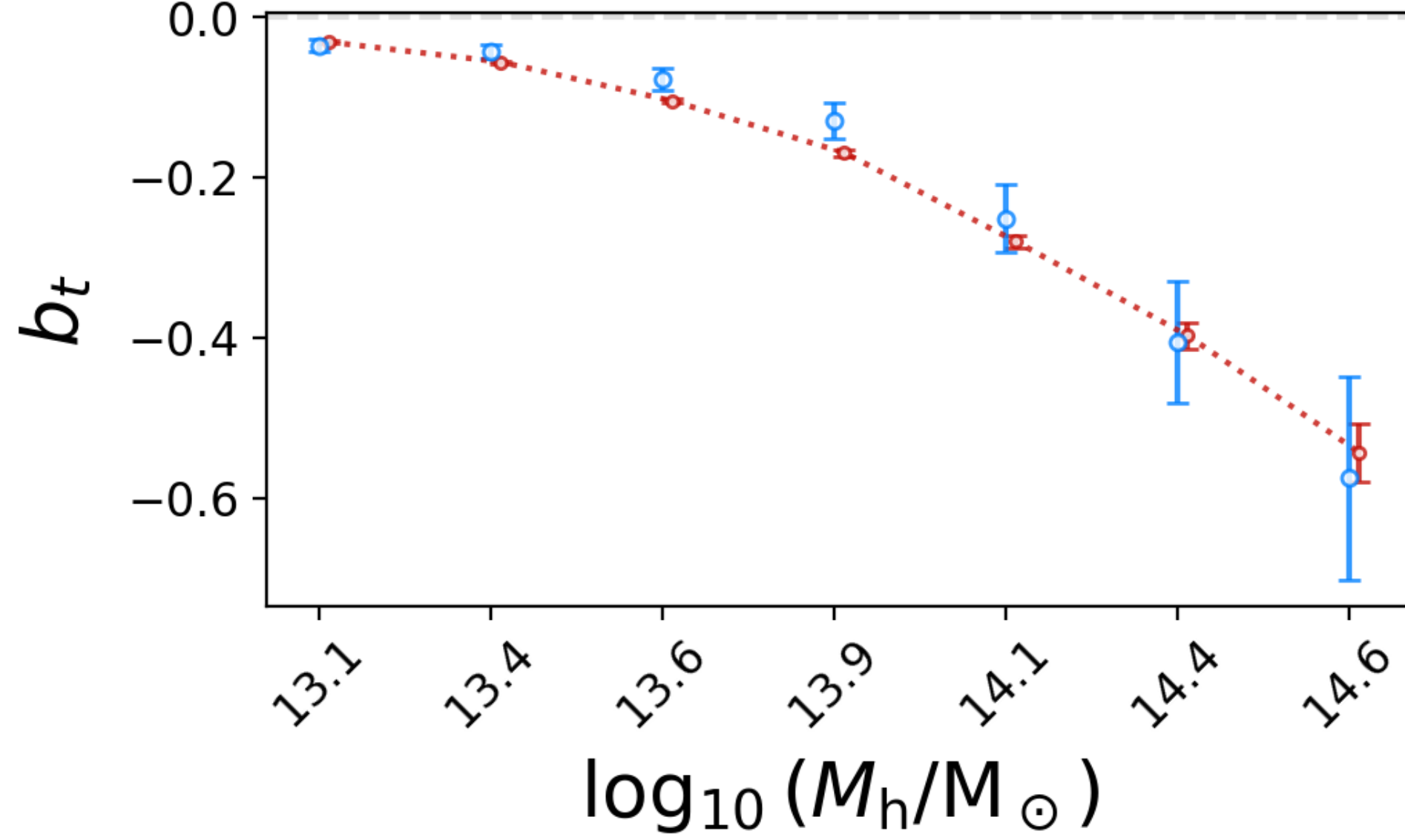
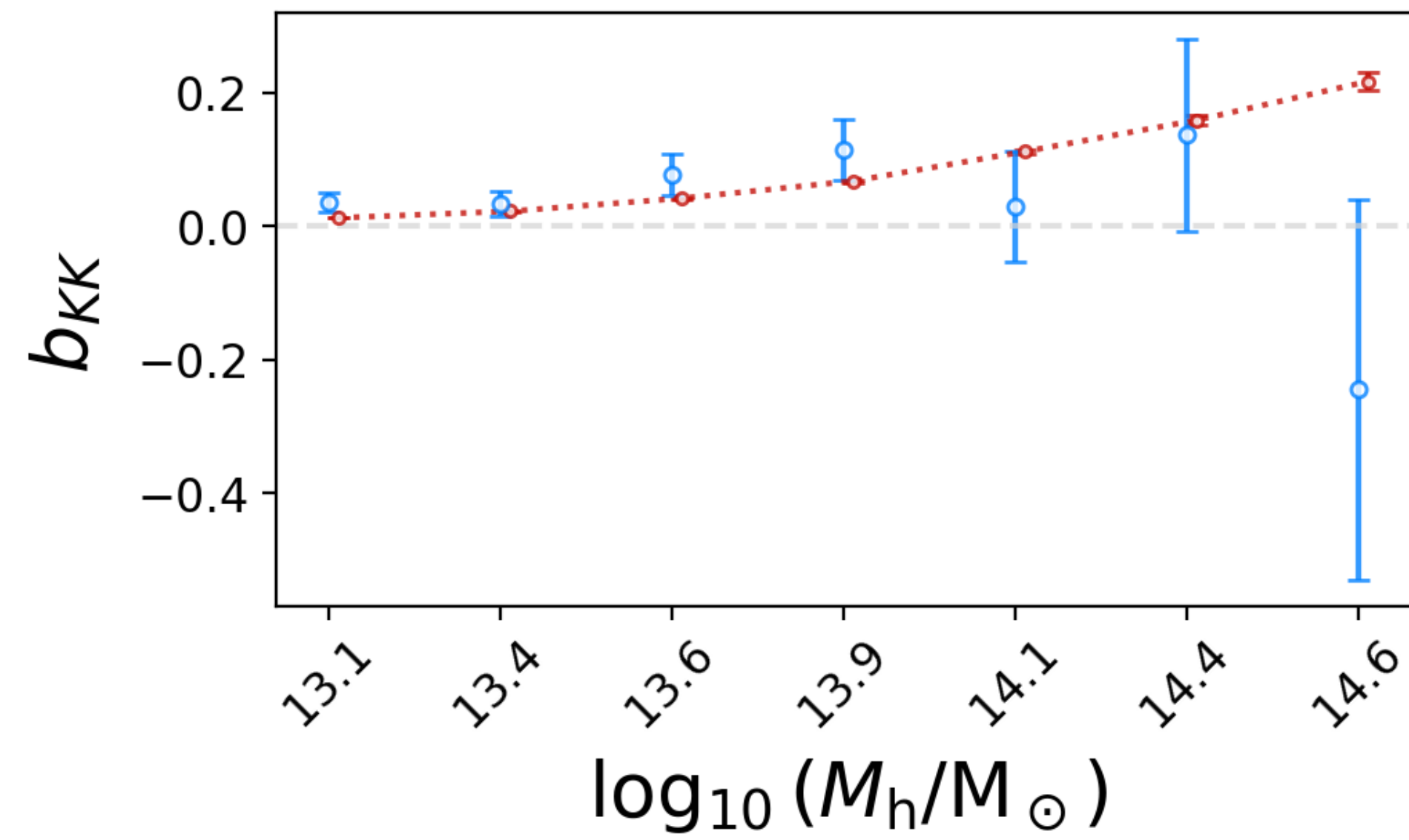
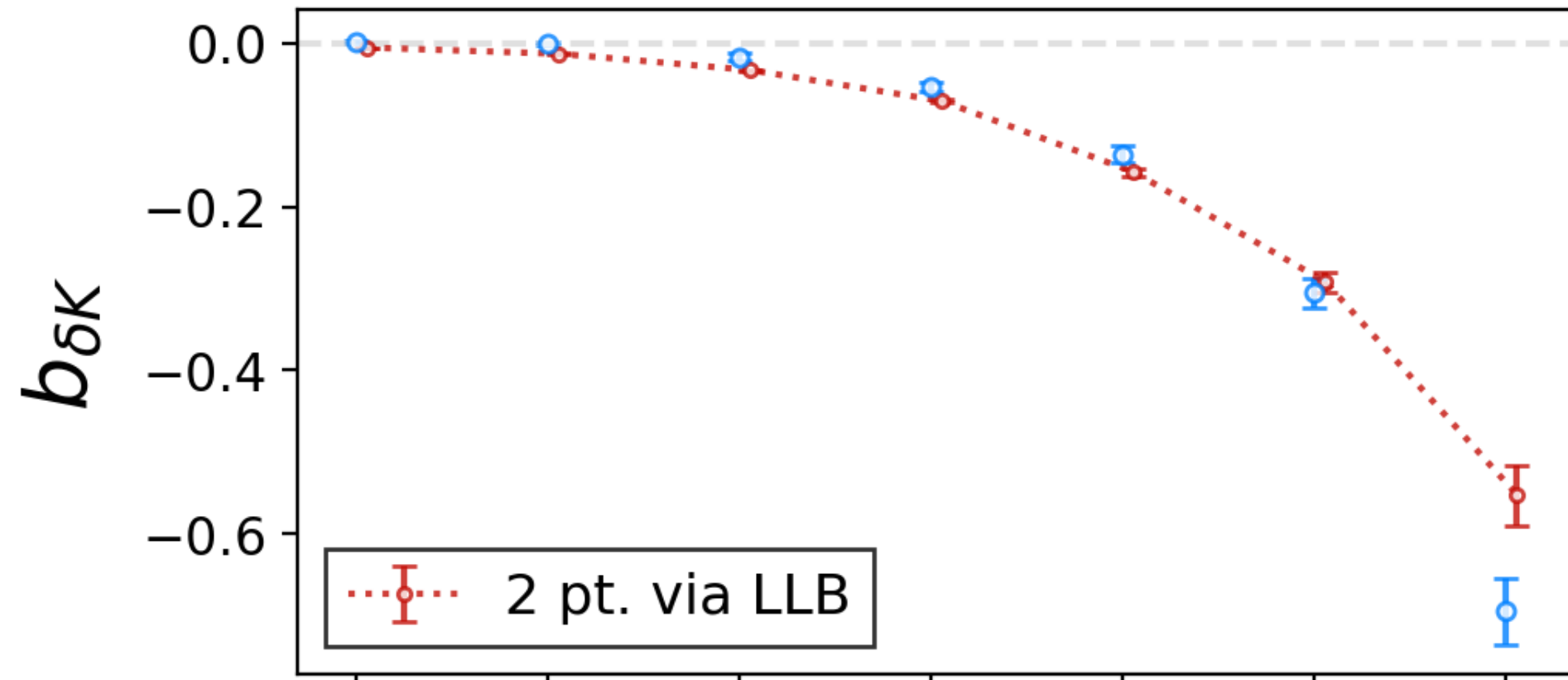
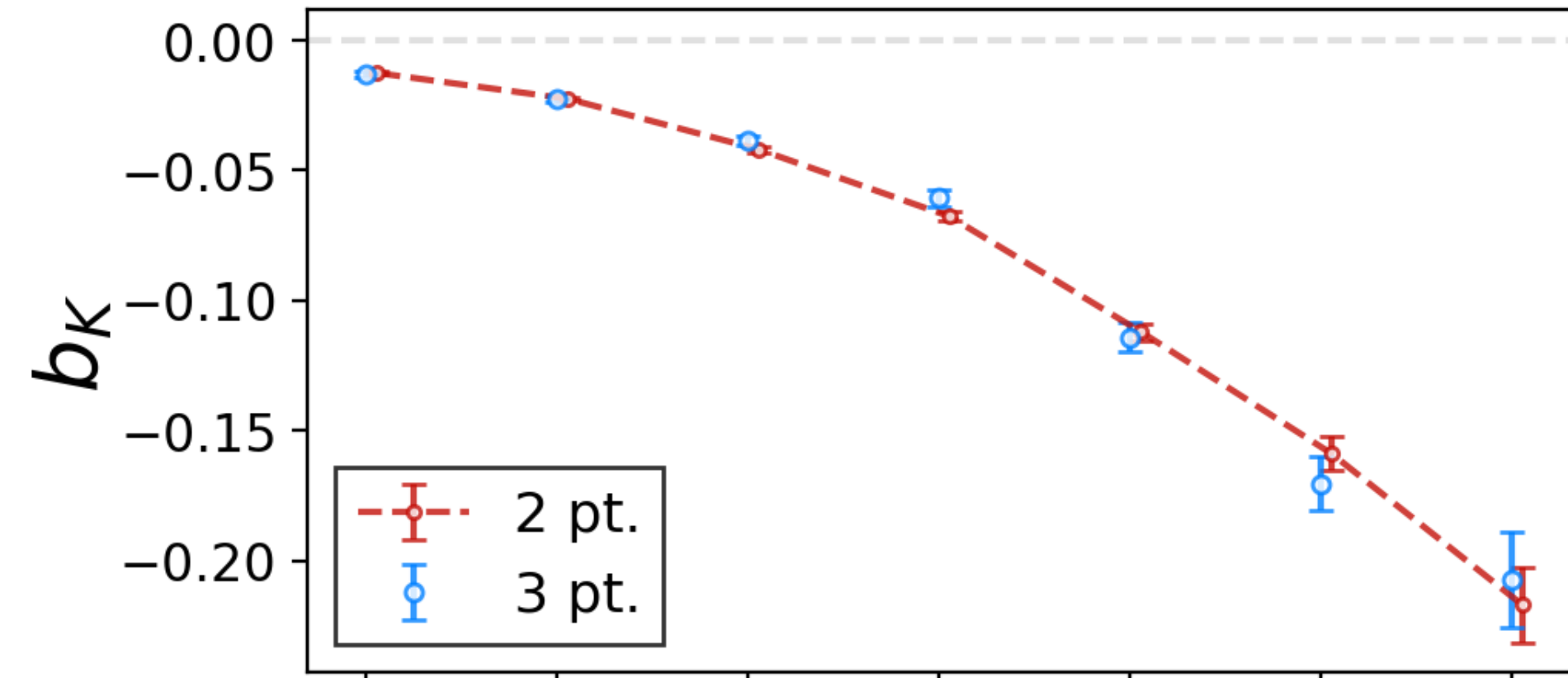


Contours

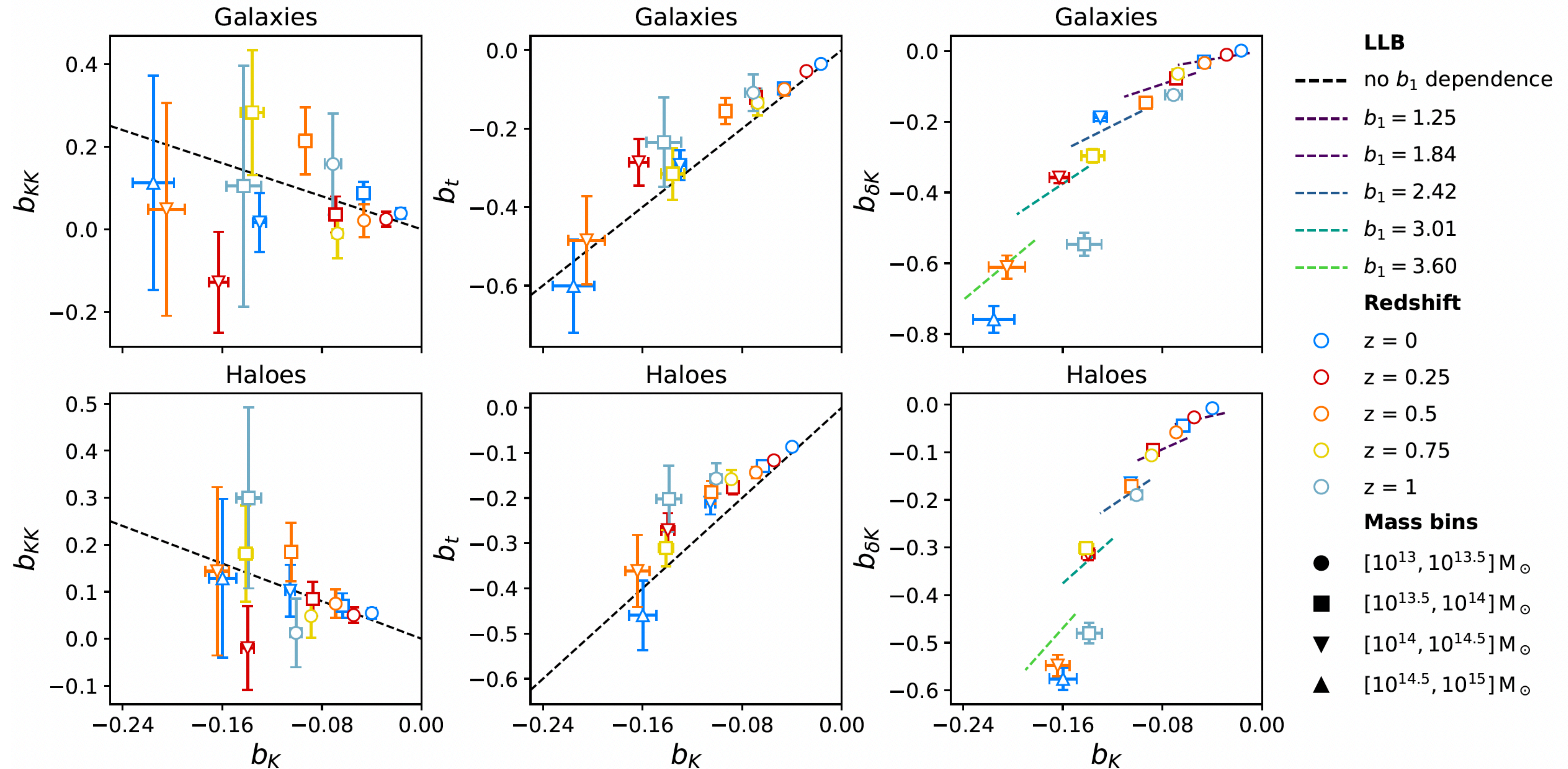
- Omission of velocity shear especially biasing to the results
- Forms a pair with b_K
 - if one goes up, the other goes down



Mass - scaling and co-evolution



Co-evolution relations across mass and redshift



Conclusions

- 3-pt IA of FLAMINGO galaxies can be accurately described using the EFT of IA, yielding consistent results
- Fixing parameters to co-evolutions reduces the parameter space without significant biases: likely a good option for lensing mitigations
- Neglecting parameters can lead to **significant biases** with respect to 2-pt.
 - Also a **cautionary tale** for higher order stats in general!

Open questions

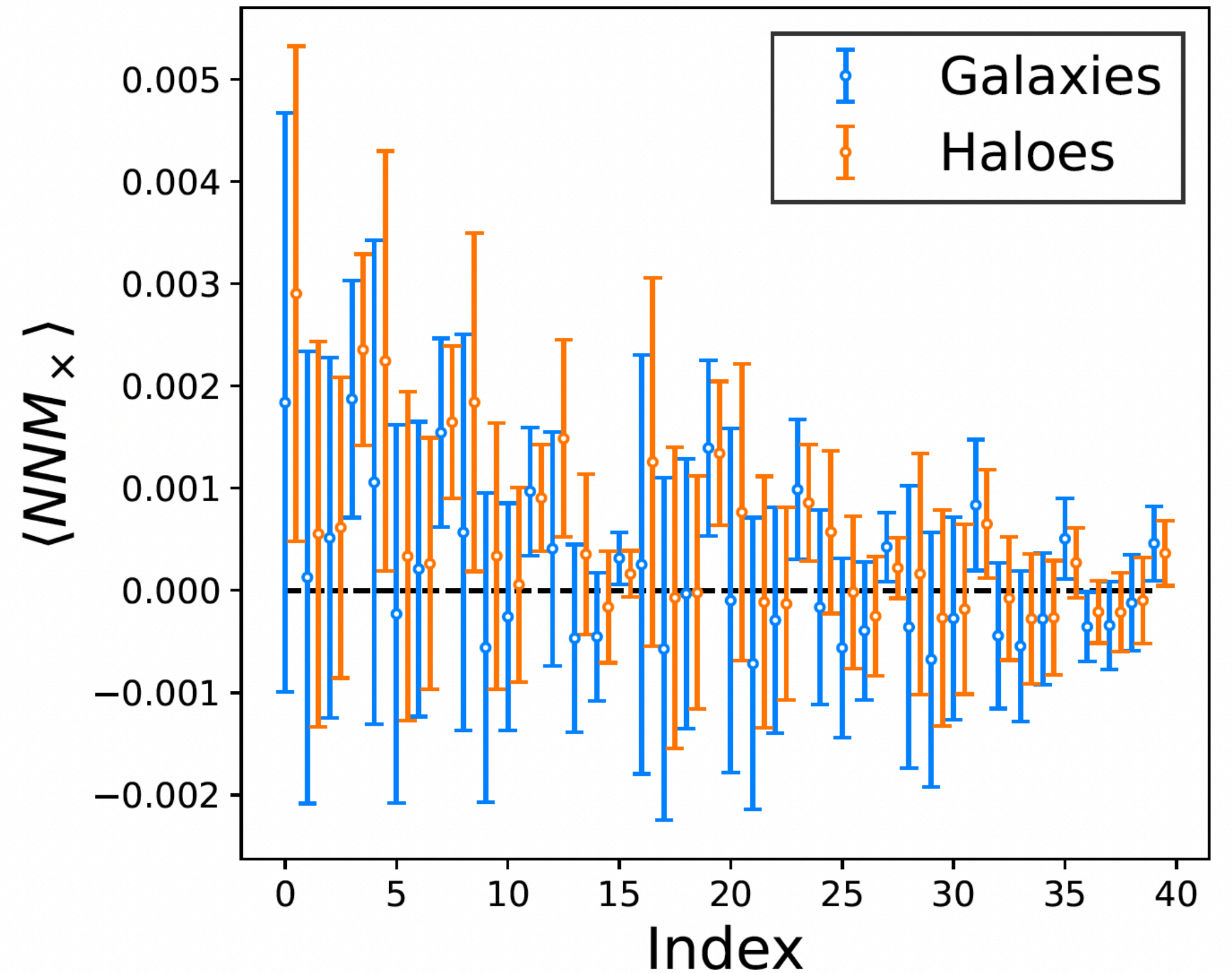
- Can we push to **smaller scales**?
- How does this affect **cosmological inference** at stage IV? (Some work already by Gomes et al, 2026)
- Plenty more! (Redshift dependence, interplay with other systematics, ...)

Extra slides

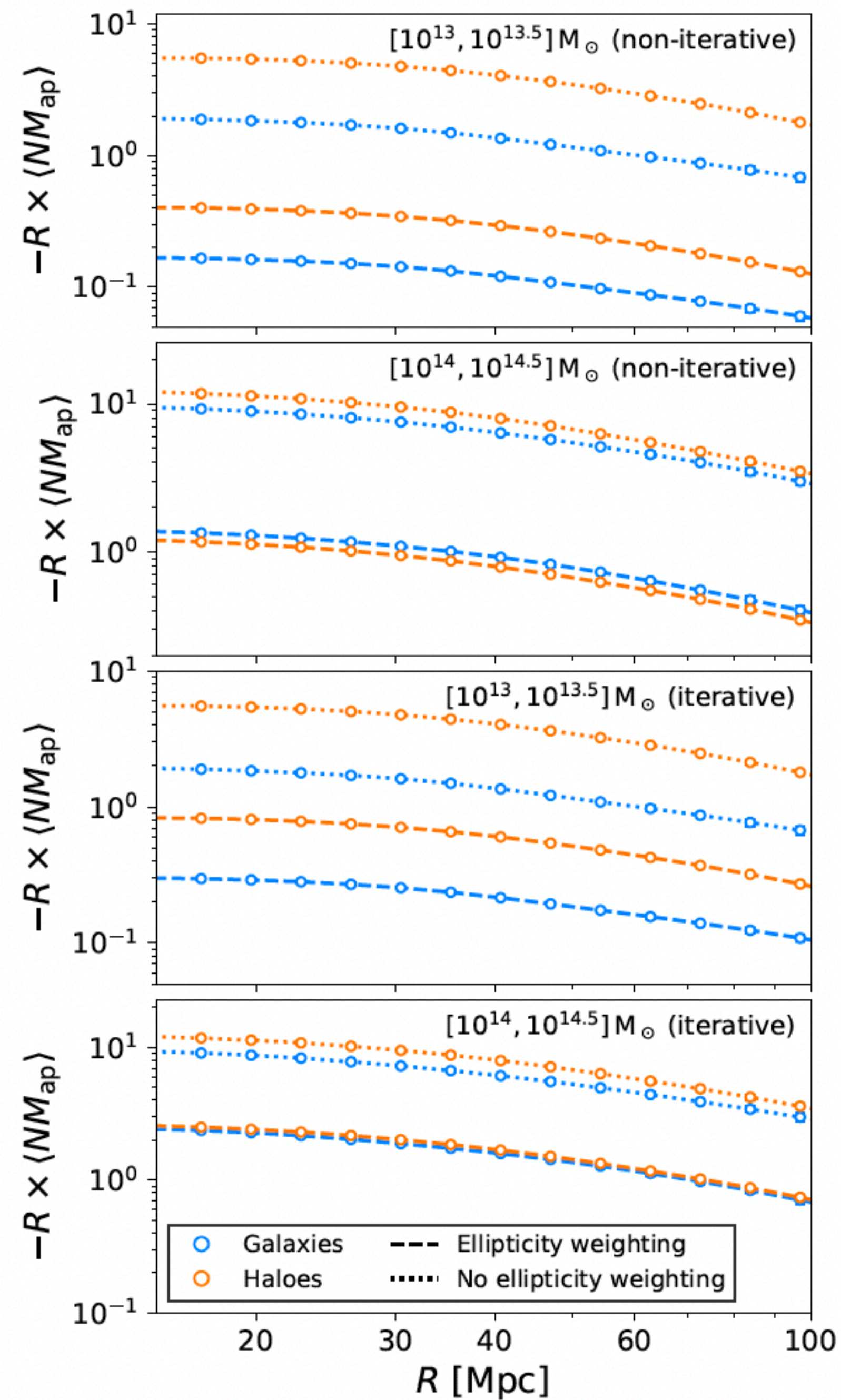
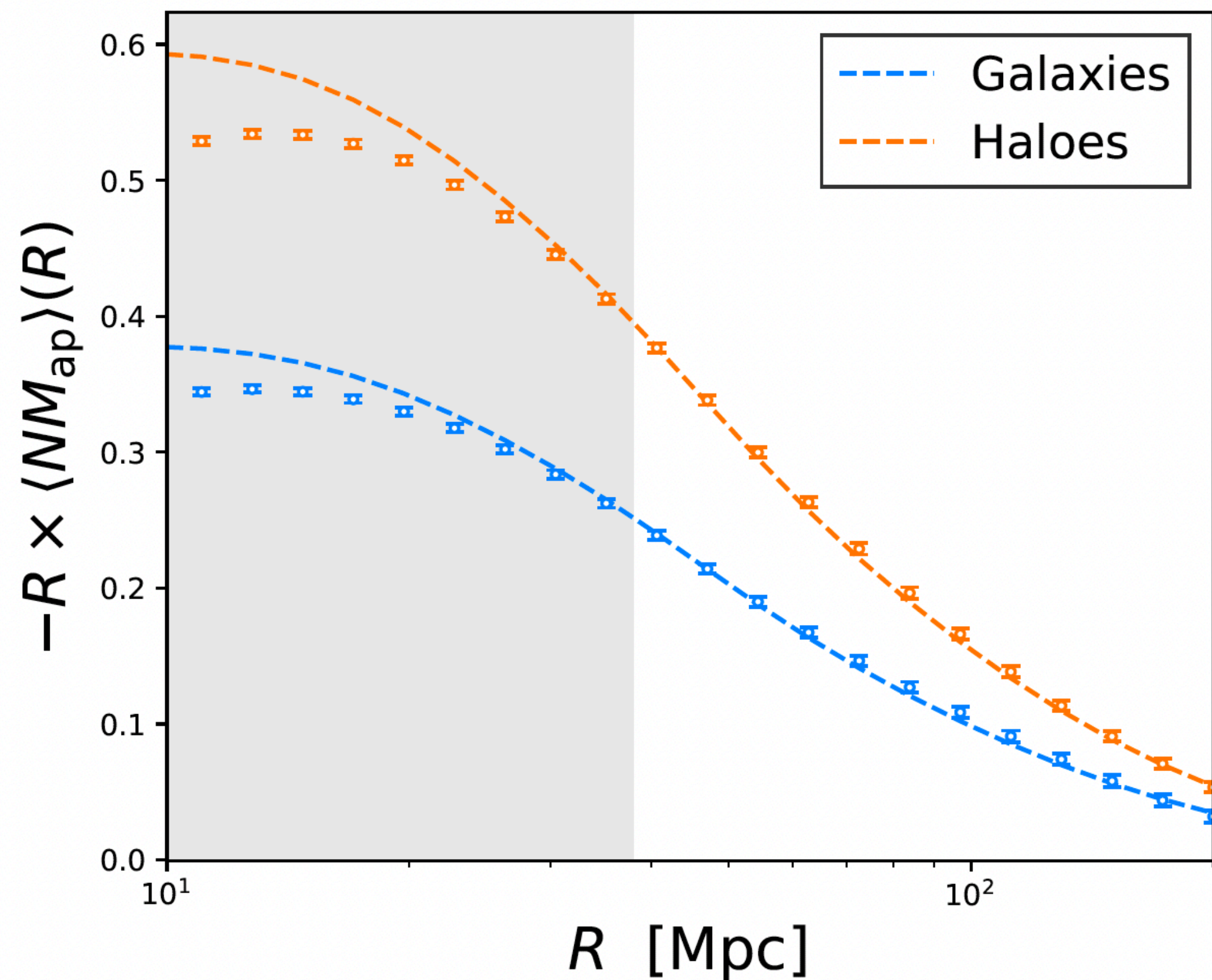
B-modes

- $\langle NNM_{\times} \rangle$ vanishes by construction due to angular averaging

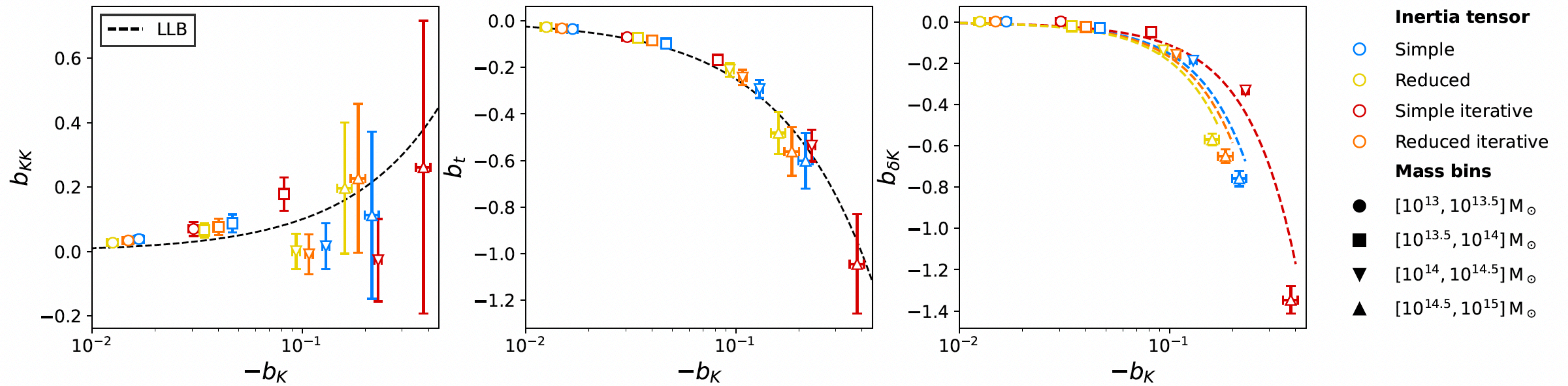
$$\begin{aligned}
 \langle NNM_{\times} \rangle &= \frac{1}{(2\pi)^5} \int_0^{\infty} k_{\perp}^{(1)} dk_{\perp}^{(1)} \int_0^{\infty} k_{\perp}^{(2)} dk_{\perp}^{(2)} \\
 &\times \int_{-\infty}^{\infty} dk_{\parallel}^{(1)} \int_{-\infty}^{\infty} dk_{\parallel}^{(2)} \int_0^{2\pi} d\phi \tilde{W}_{\Pi}(k_{\parallel}^{(1)}) \tilde{W}_{\Pi}(k_{\parallel}^{(2)}) \\
 &\times \tilde{U}_{R_1}(k_{\perp}^{(1)}) \tilde{U}_{R_2}(k_{\perp}^{(2)}) \tilde{U}_{R_3}(\|\mathbf{k}_{\perp}^{(1)} + \mathbf{k}_{\perp}^{(2)}\|) \\
 &\times B_{\delta\delta B}(k_{\perp}^{(1)}, k_{\perp}^{(2)}, k_{\parallel}^{(1)}, k_{\parallel}^{(2)}, \phi),
 \end{aligned}$$



2-pt. Measurements



Dependence on inertia tensor



Projections

