

Introduction to the eikonal approximation and its applications

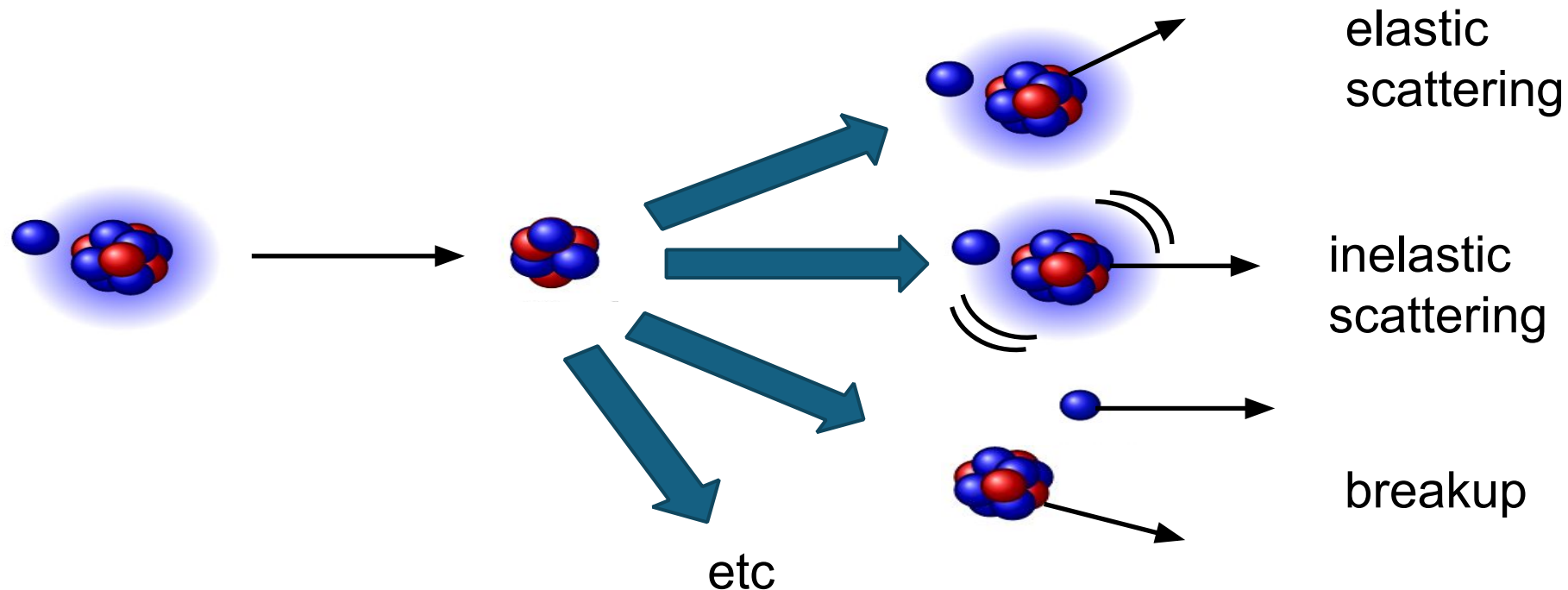
Ecole Joliot Curie

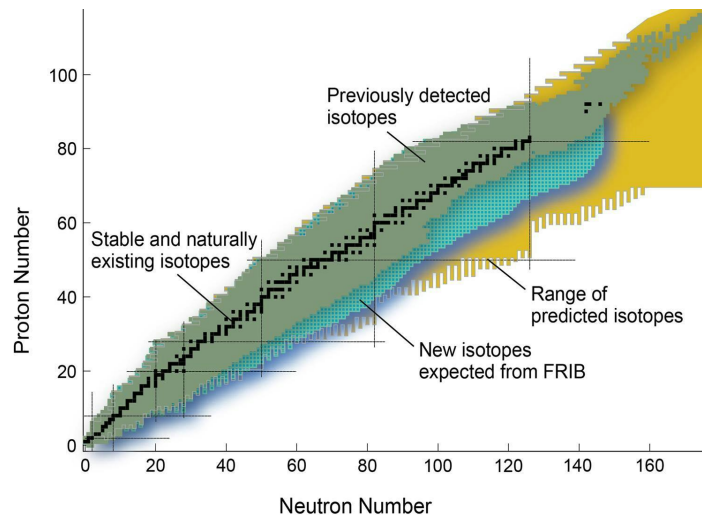
Chloë Hebborn – September 2026



Nuclear Reactions

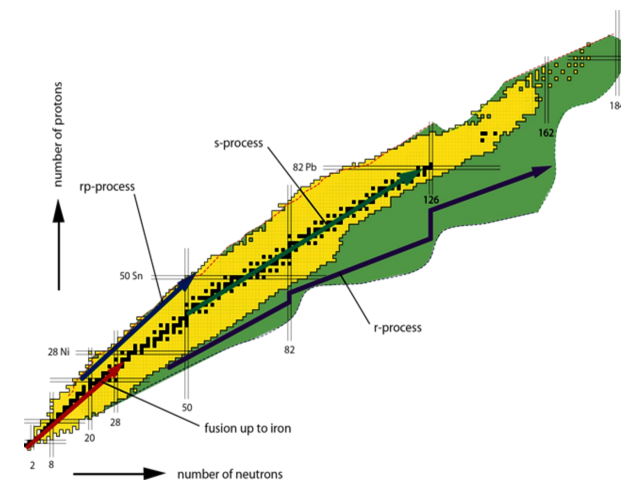
Nuclear reactions are processes in which two nuclear species collide, allowing them to interact, exchanging matter, energy or momentum





Studying properties of unstable nuclei
(Source: FRIB)

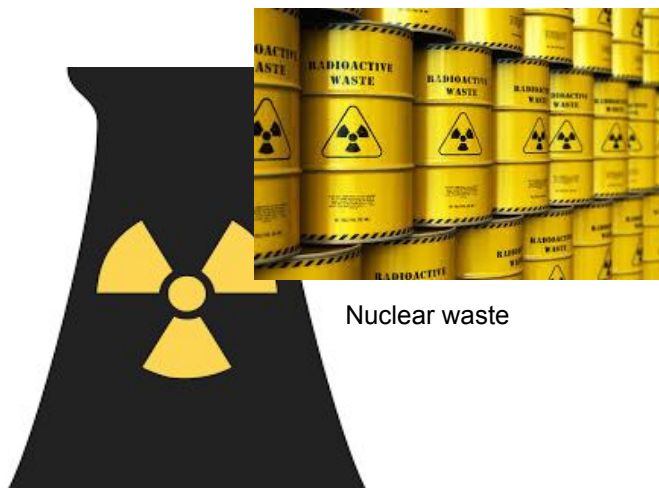
Why do we care?



Nuclear astrophysics
(source GSI)



Stockpile stewardship (Source: DOE)

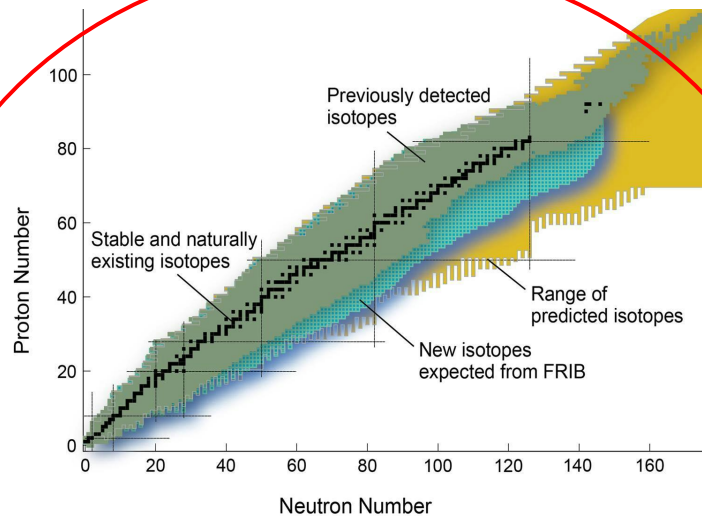


Nuclear waste

Nuclear energy

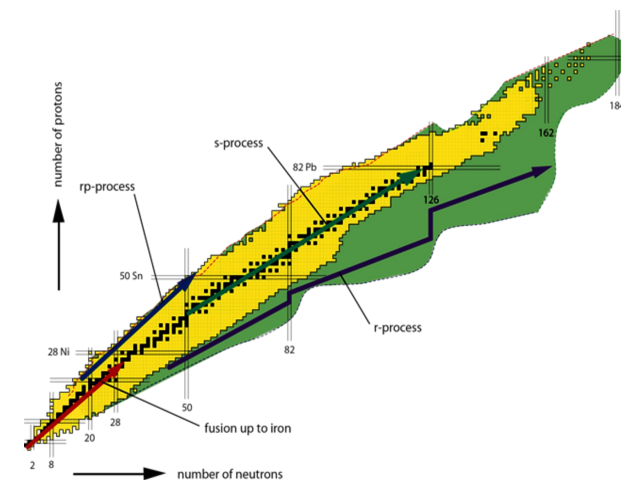


Nuclear medicine



Studying properties of unstable nuclei
(Source: FRIB)

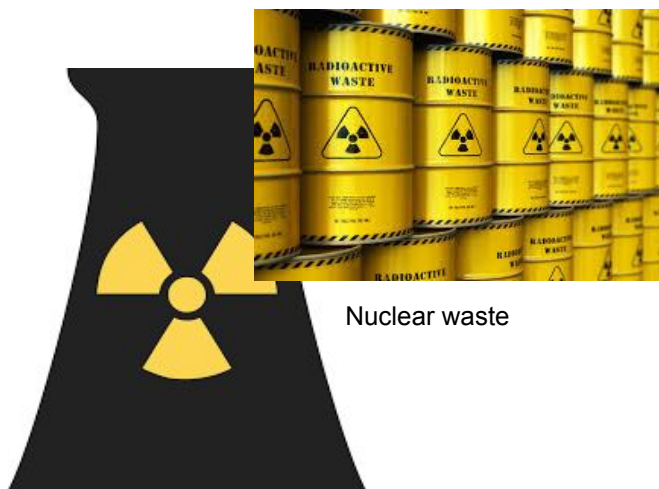
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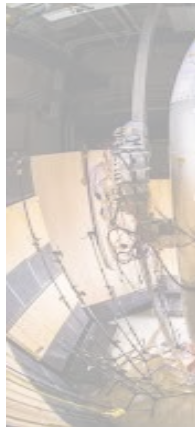
Nuclear waste



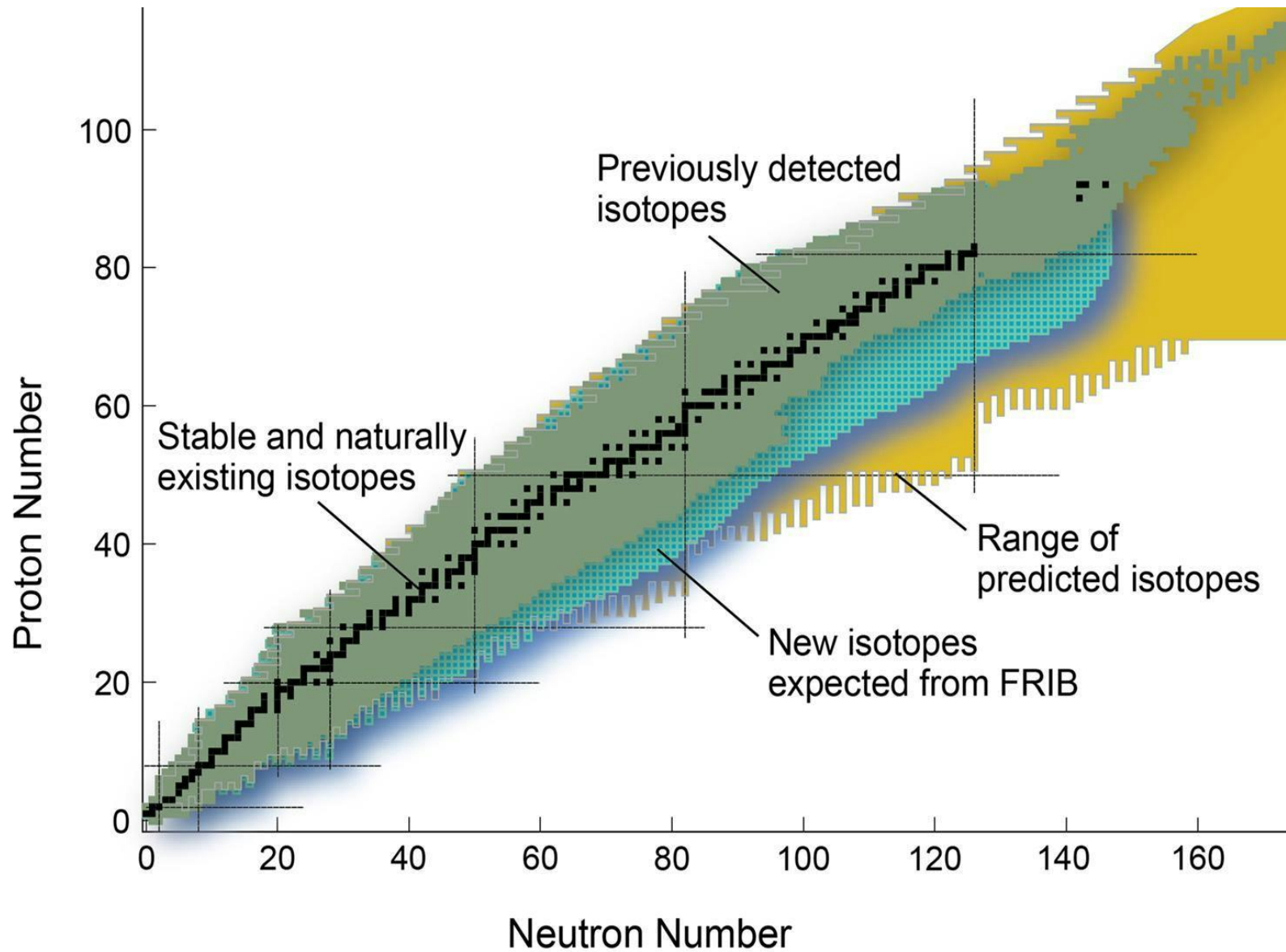
Nuclear energy



Nuclear medicine

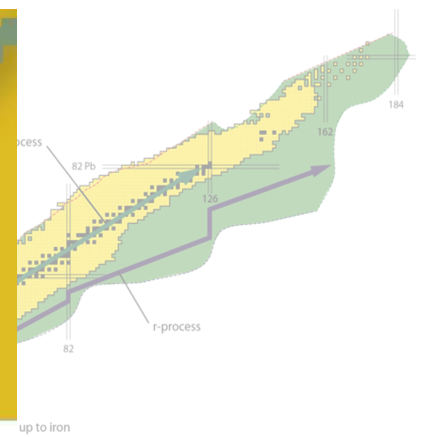


Stockpile stev
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Studying properties of unstable nuclei

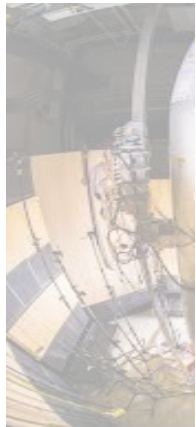
(Source: FRIB)



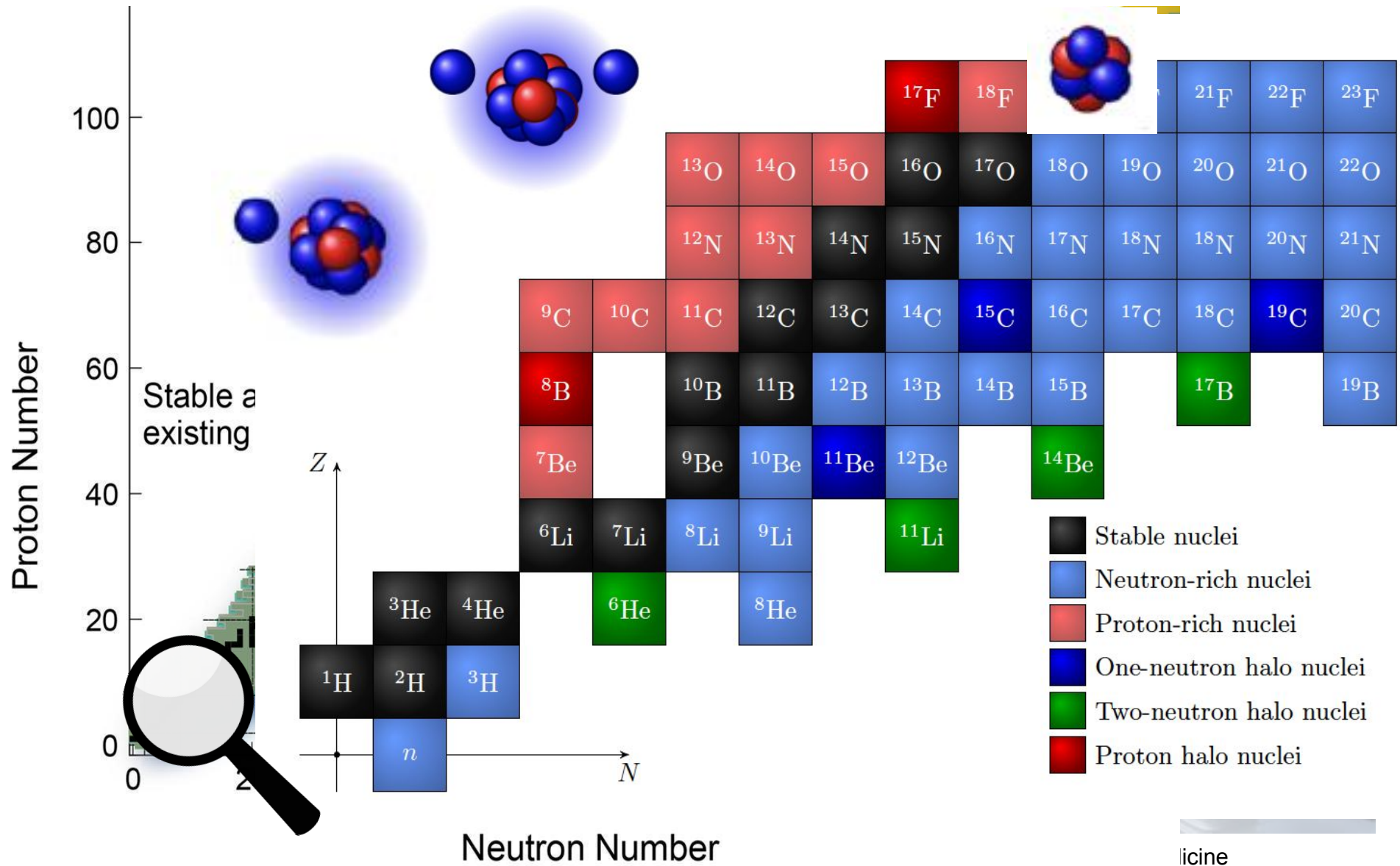
for astrophysics
source GSI)



licine

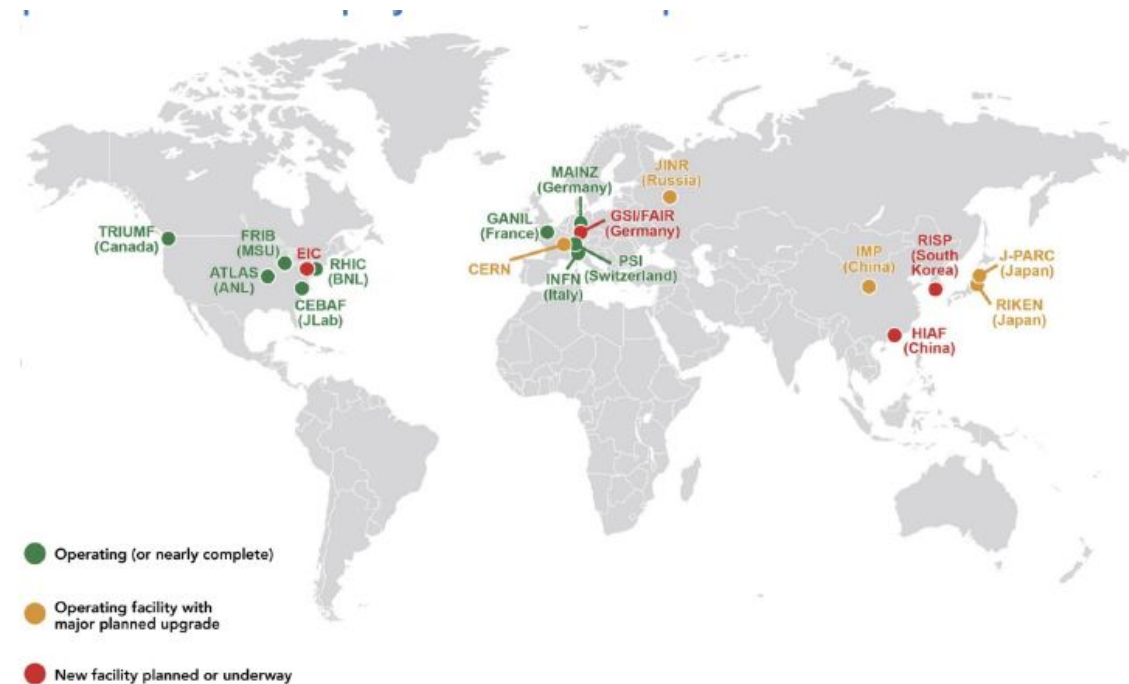
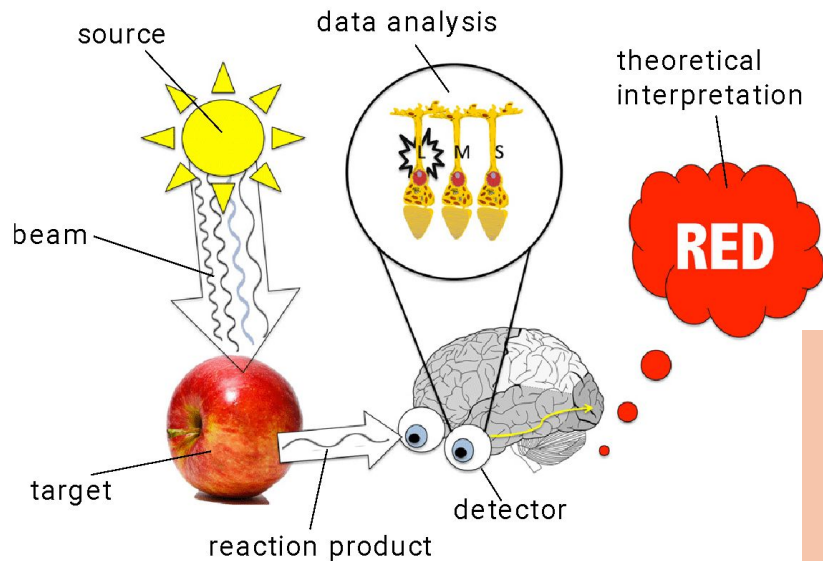
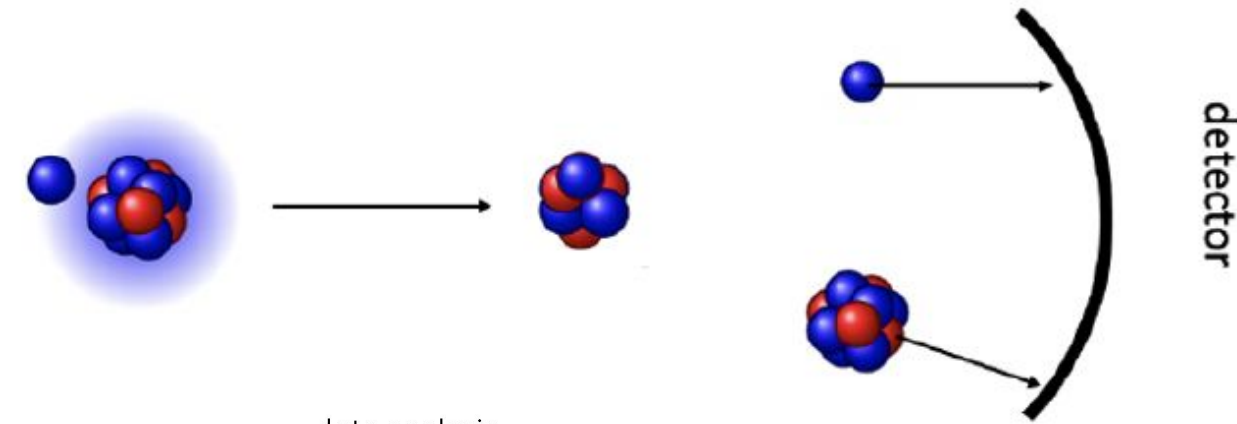


Stockpile stev
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Unstable nuclei are often produced and studied through reactions



**Accurate
measurements and
theory crucial!**

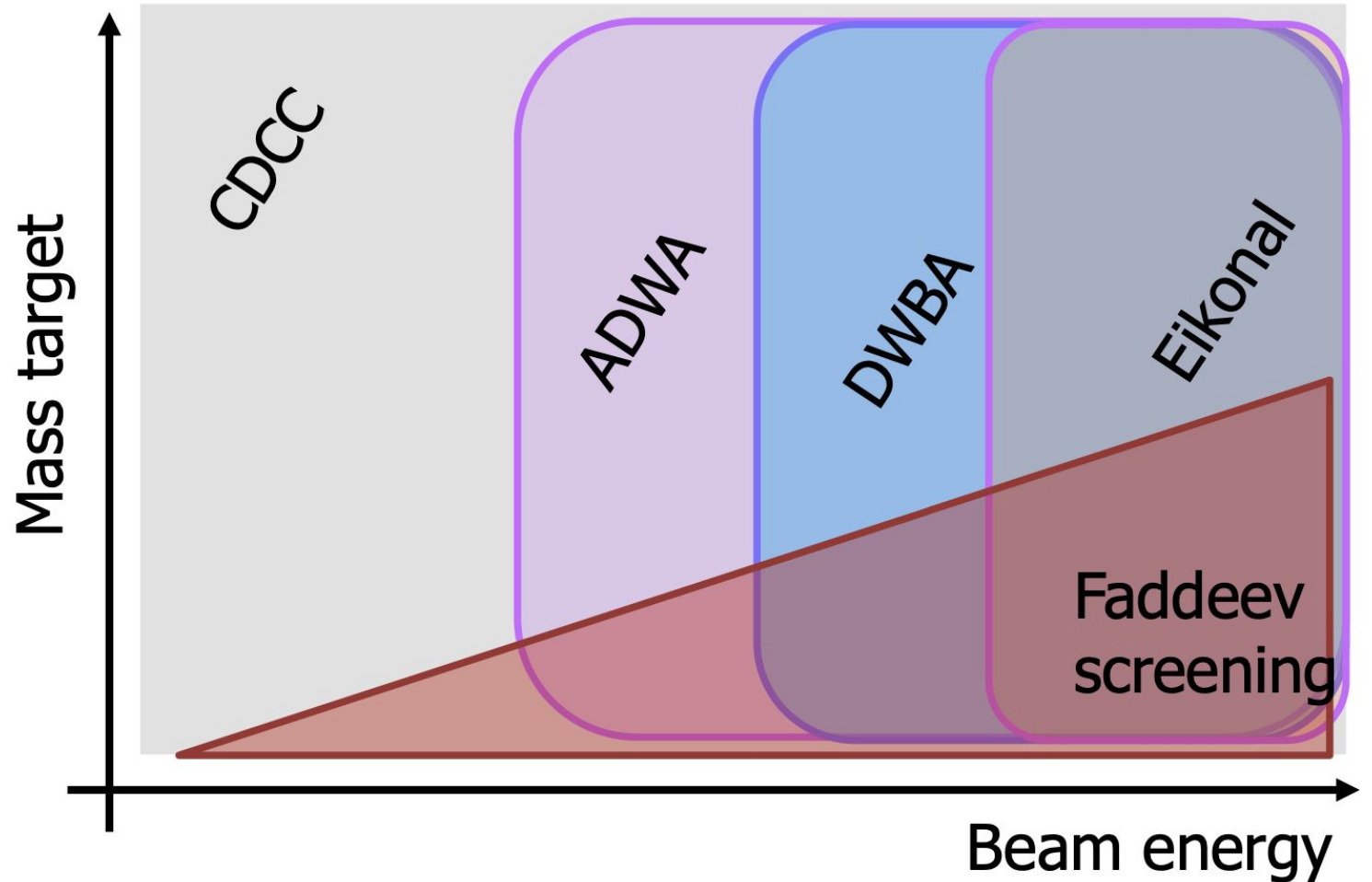
During this
hour and
1/2, we'll
try to reply
to few
questions..

- When is the eikonal approximation used?
For which reactions & at what energies?
- Formal derivation of the eikonal approximation for a two-body problem
- Extension to a few-body ones
- Hands on part with a python notebook



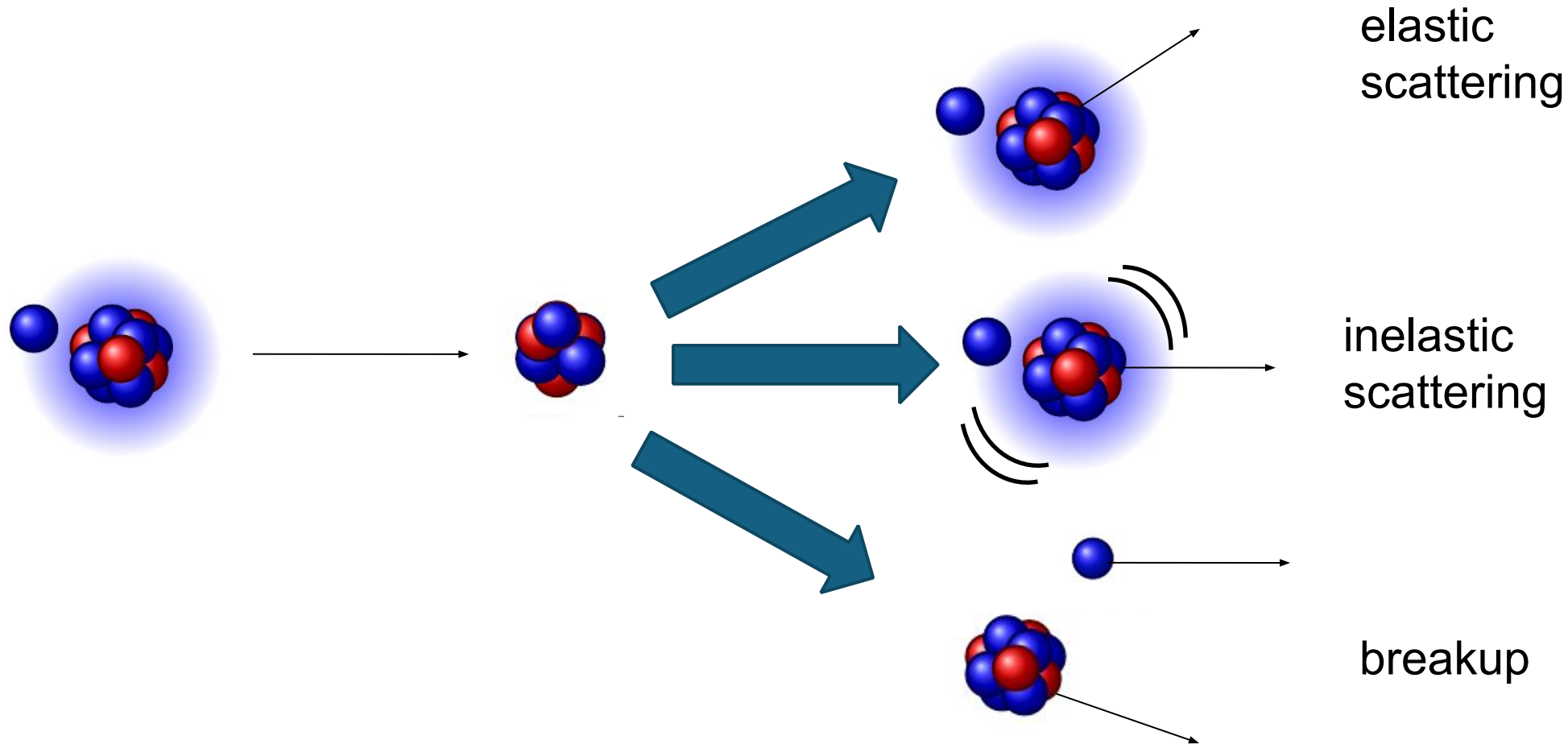
At what energy is the eikonal approximation used?

Image courtesy of
Filomena Nunes

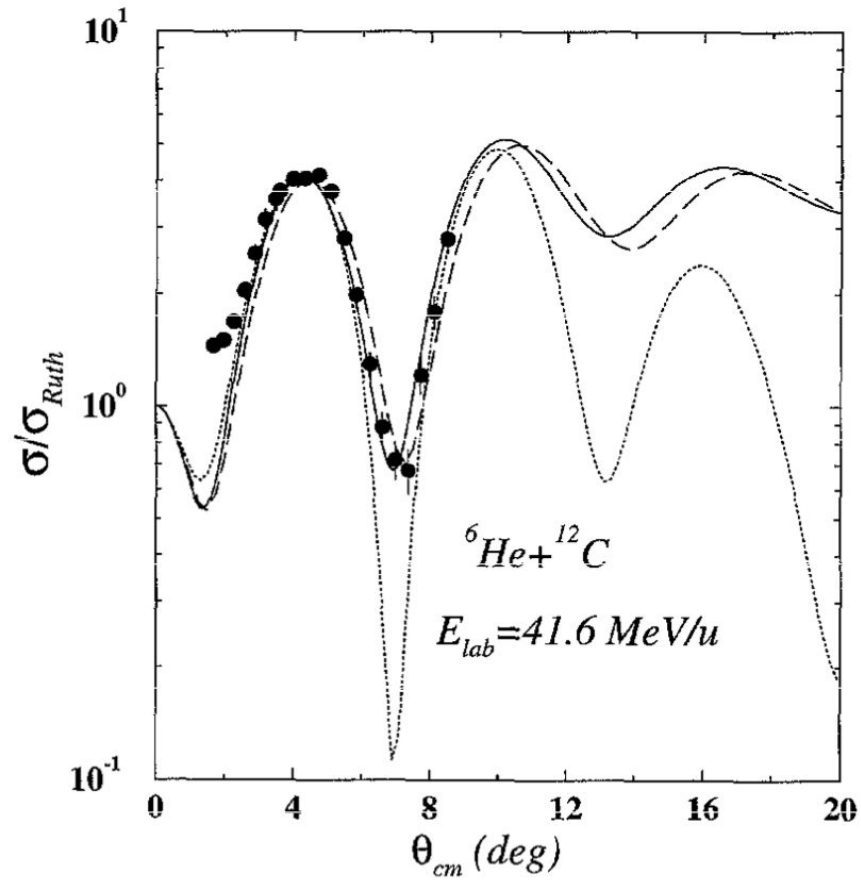


Eikonal model is cheap computationally, it can be generalized to complex reactions...

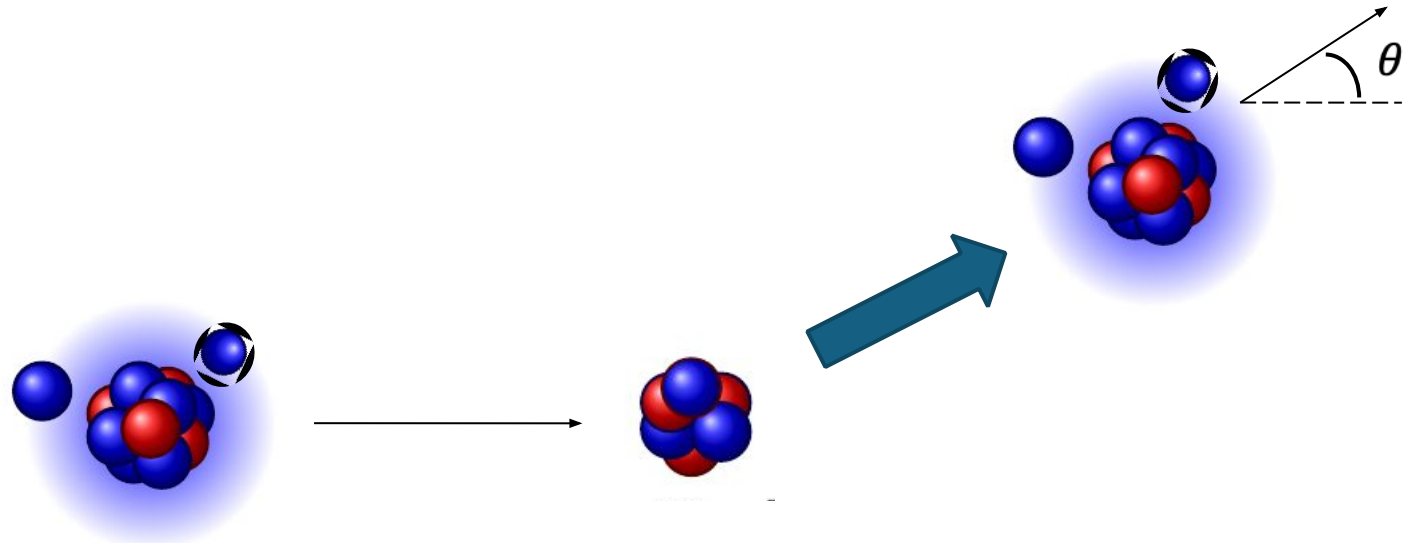
And for which reaction probes ?



Intermediate energy elastic scattering of ${}^6\text{He}$ in a four-body reaction model

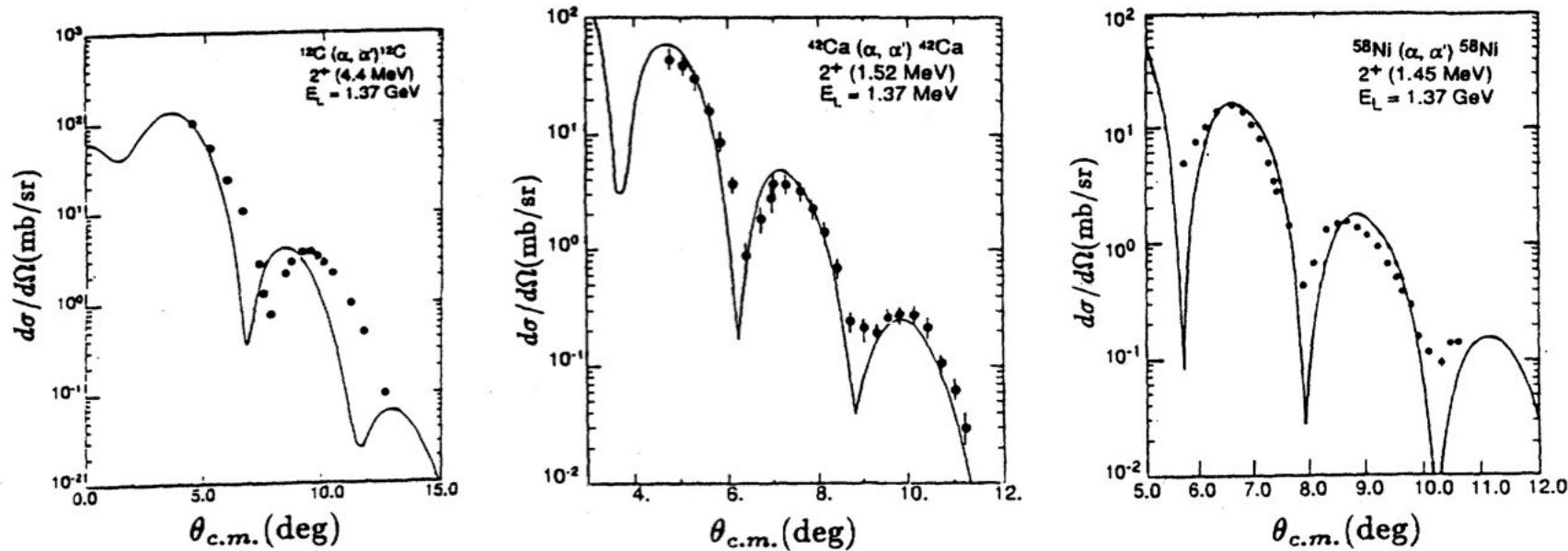


Al Khalili et al. PLB 378 (1996) 45-49



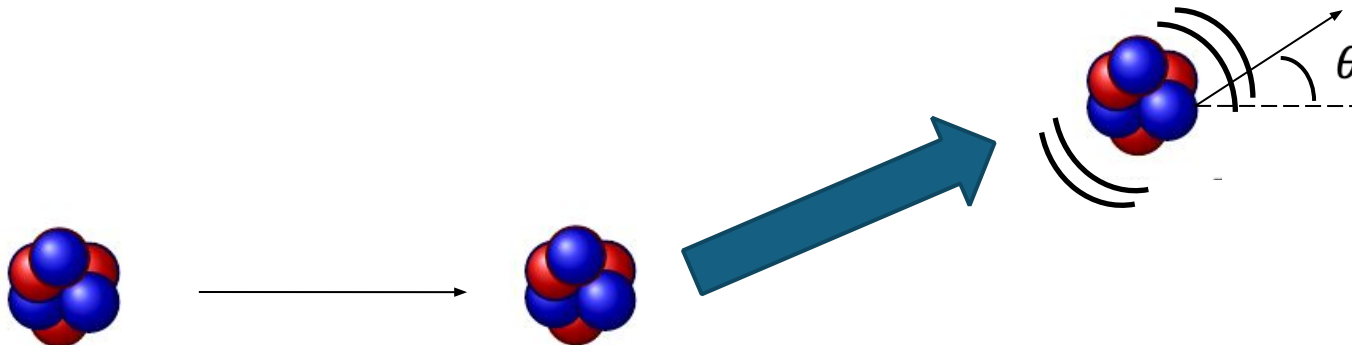
Elastic scattering observables can inform about the cluster structure in nuclei and their rms radii !

High energy inelastic scattering of nuclei

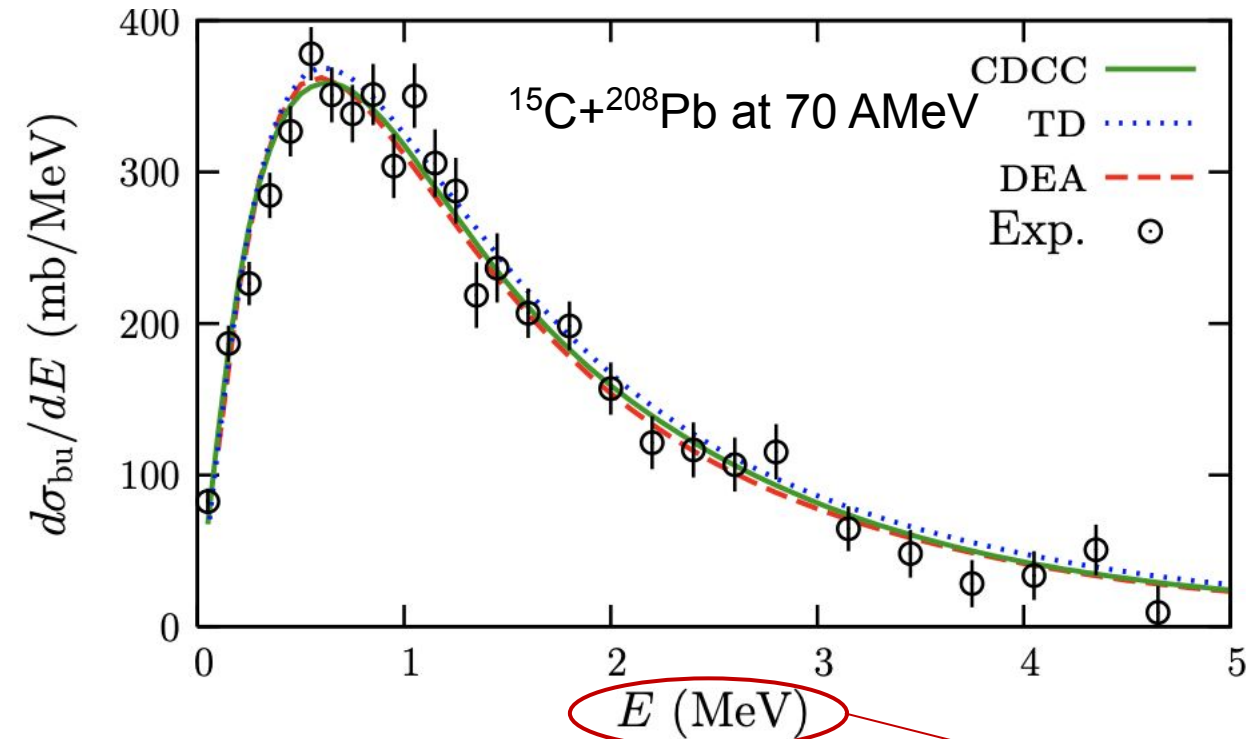


Inelastic scattering observables can inform about the electromagnetic transition (e.g. $B(E2)$) in nuclei !

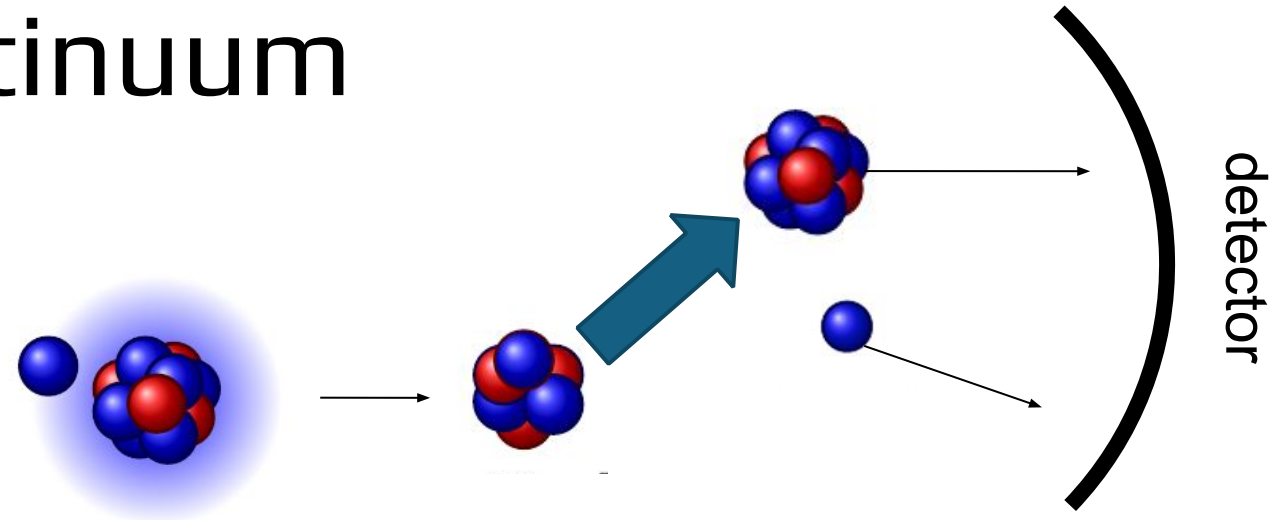
Lenzi et al. PRC 40 6 1989



Breakup to study electromagnetic response in the continuum



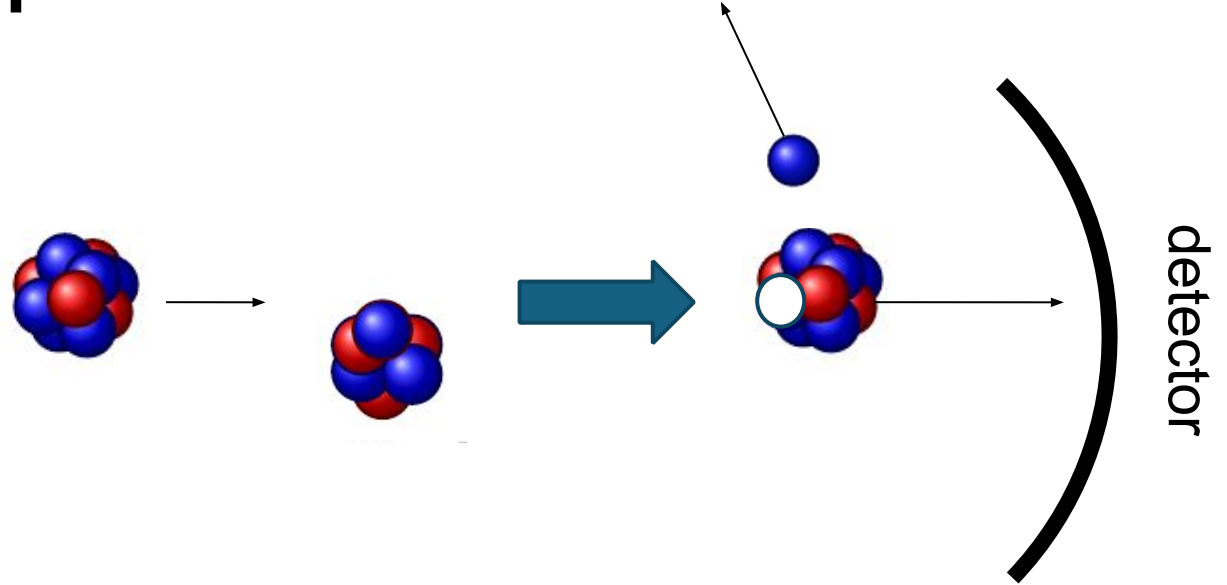
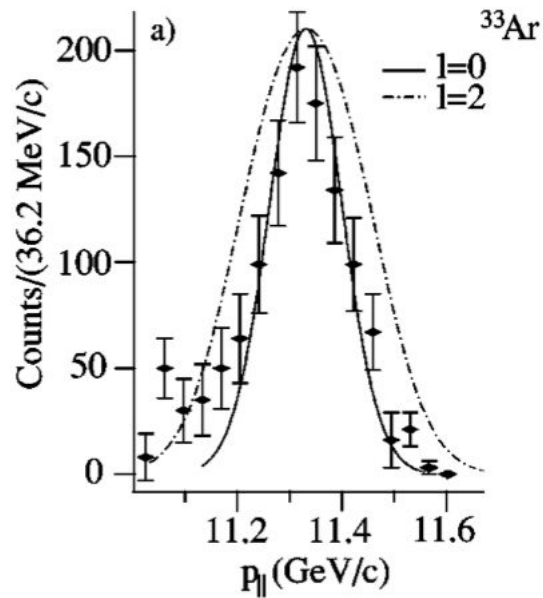
Capel et al., PRC 85, 044604 (2012)



Breakup reactions are useful to study the cluster structure and probe electromagnetic response of nuclei !

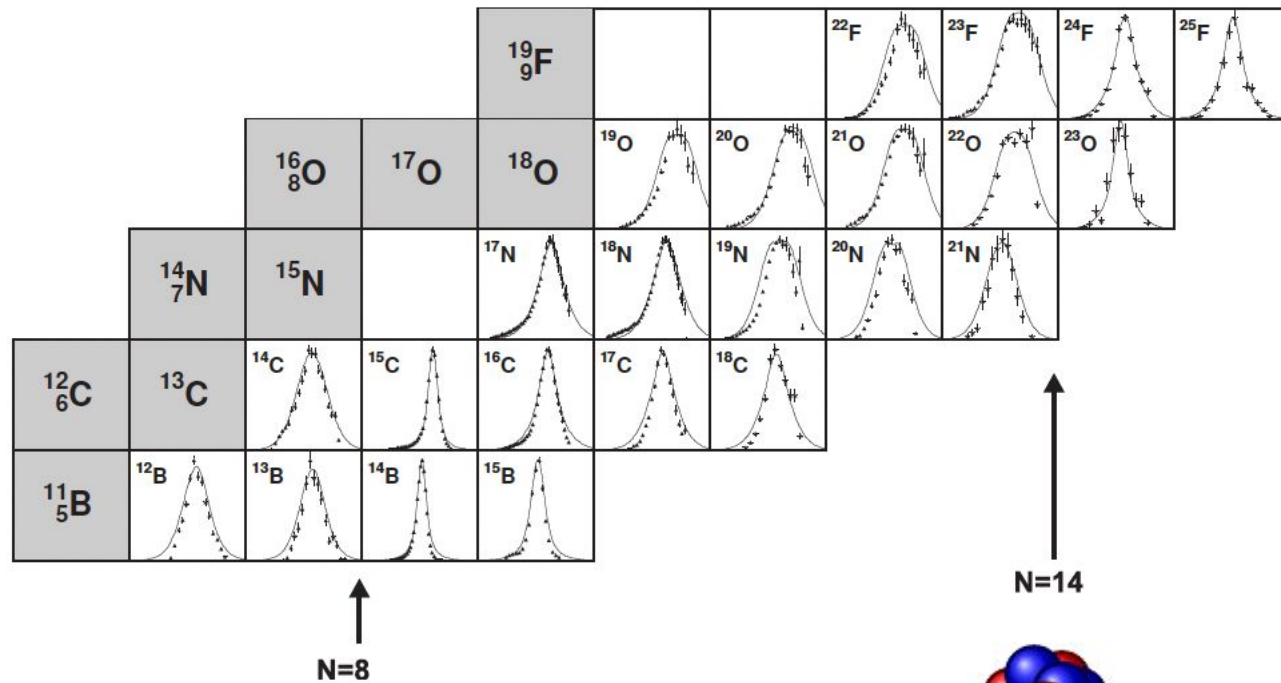
Relative ^{14}C -n energy after the breakup

Nucleon removal to study spectroscopic properties of nuclei



Gade et al., *PRC* 69 034311 (2004)

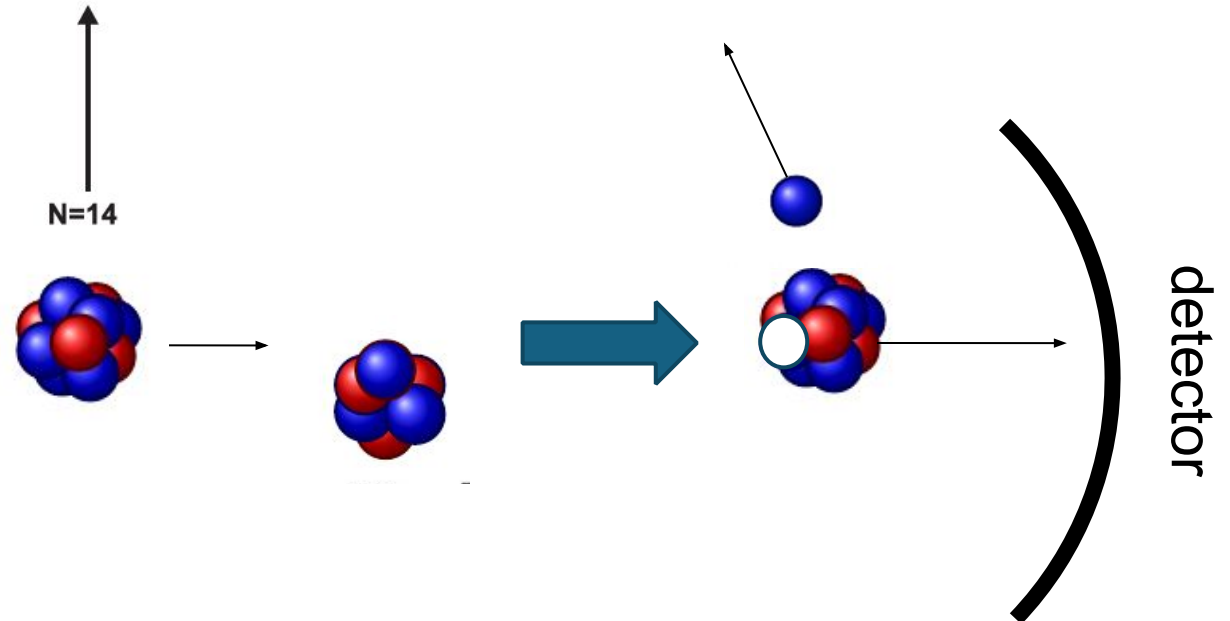
Nucleon removal also carries information about the nucleus size



Sauvan et al. PLB 491 1 (2000)

Heisenberg's principle:

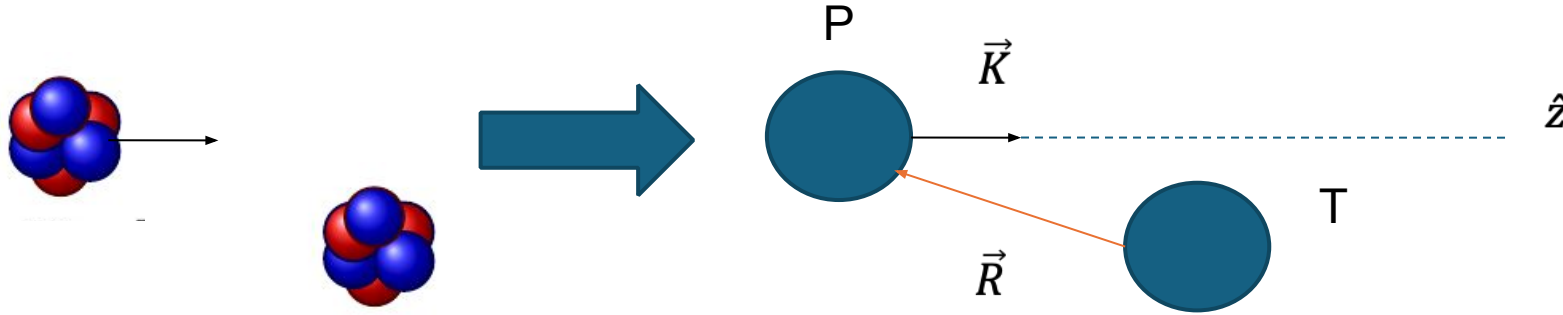
Narrow momentum distribution
=
Large spatial extension



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The reaction is often seen as a few-body problem



- Two-body Schrödinger equation

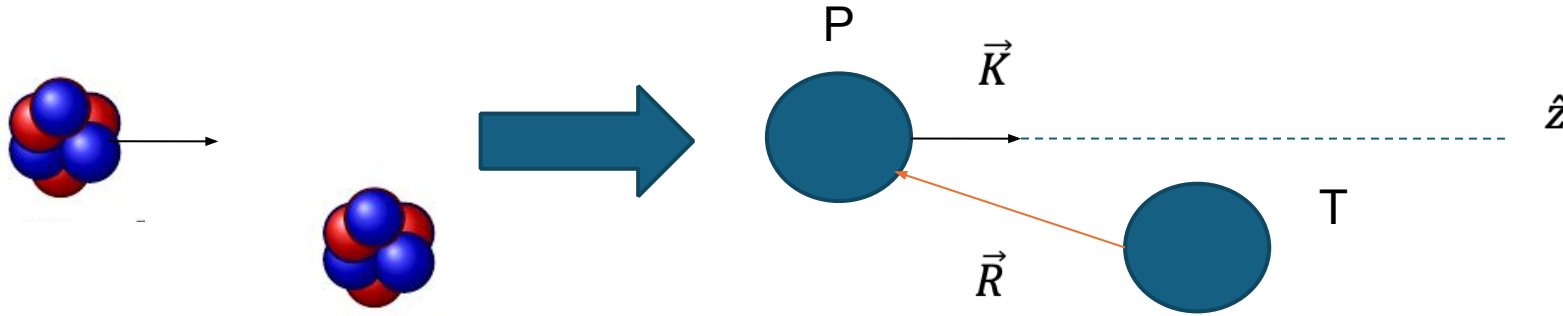
$$[T_{\vec{R}} + V_{PT}(R)]\Psi(\vec{R}) = E\Psi(\vec{R})$$

- Initial condition

$$\Psi(\vec{R}) \rightarrow_{Z \rightarrow -\infty} e^{iKZ}$$

With energy of the system $E = \frac{\hbar^2 K^2}{2\mu}$

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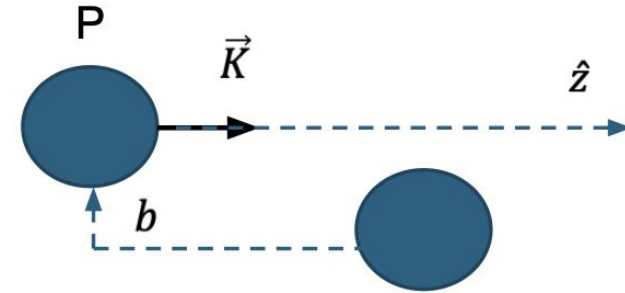
Optical potential:

complex effective interaction

with imaginary parts modeling
loss of flux in inelastic channels

At high energy, the projectile is not much deflected by the target

- Two-body Schrödinger equation $[T_{\vec{R}} + V_{PT}(R)]\Psi(\vec{R}) = E\Psi(\vec{R})$
- Initial condition $\Psi(\vec{R}) \xrightarrow{Z \rightarrow -\infty} e^{iKZ}$
- Eikonal factorization: $\Psi(\vec{R}) = e^{iKZ} \widehat{\Psi(\vec{R})}$



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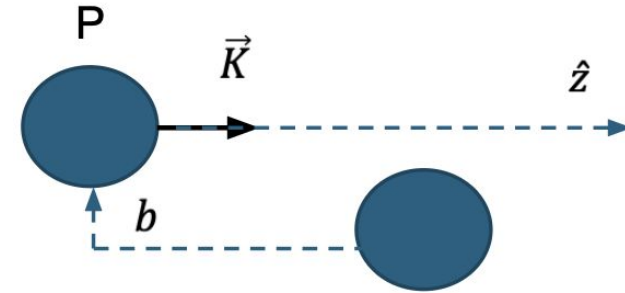
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$$T_{\vec{R}} e^{iKZ} \widehat{\Psi}(\vec{R}) = -\frac{\hbar^2}{2\mu} \Delta_{\vec{R}} e^{iKZ} \widehat{\Psi}(\vec{R}) = -\frac{\hbar^2}{2\mu} \left(\frac{\delta^2}{\delta X^2} + \frac{\delta^2}{\delta Y^2} + \frac{\delta^2}{\delta Z^2} \right) e^{iKZ} \widehat{\Psi}(\vec{R})$$

=?

$$= -\frac{\hbar^2}{2\mu} \left(-K^2 e^{iKZ} \widehat{\Psi}(\vec{R}) + 2iK e^{iKZ} \frac{\delta}{\delta Z} \widehat{\Psi}(\vec{R}) + e^{iKZ} \Delta_{\vec{R}} \widehat{\Psi}(\vec{R}) \right)$$



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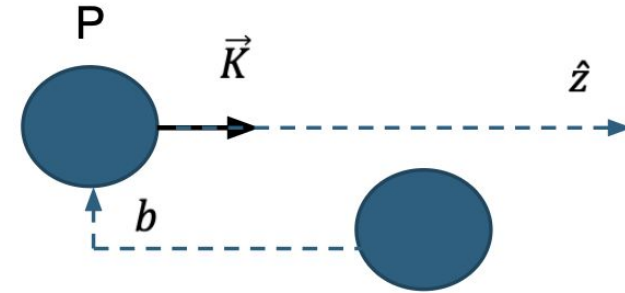
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$$= ?$$

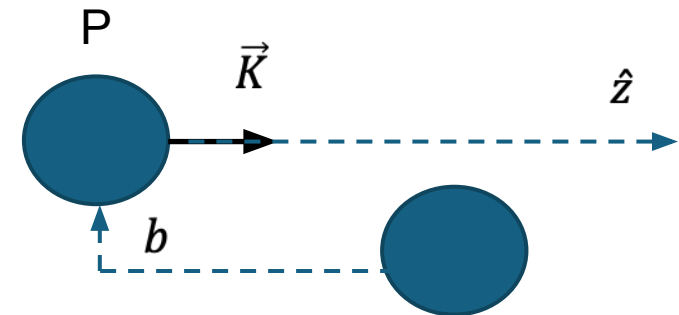
$$= -\frac{\hbar^2}{2\mu} \left(-K^2 e^{iKZ} \widehat{\Psi}(\vec{R}) + 2iK e^{iKZ} \frac{\delta}{\delta Z} \widehat{\Psi}(\vec{R}) + e^{iKZ} \Delta_{\vec{R}} \widehat{\Psi}(\vec{R}) \right)$$

- Eikonal approximation: $\Delta_{\vec{R}} \widehat{\Psi}(\vec{R}) \ll K \frac{\delta}{\delta Z} \widehat{\Psi}(\vec{R})$



Eikonal equations can be solved analytically!

... and one directly obtains $i \hbar v \frac{\delta}{\delta Z} \widehat{\Psi(\vec{R})} = V_{PT}(R) \widehat{\Psi(\vec{R})}$ with $v = \frac{\hbar K}{\mu}$

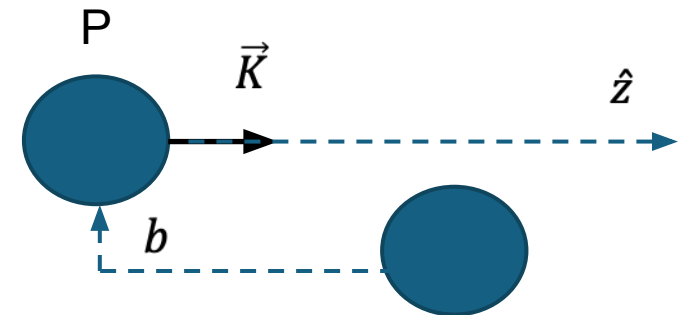


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Eikonal Solution: $\widehat{\Psi(\vec{R})} = e^{iKZ} \widehat{\Psi(\vec{R})} = e^{iKZ} e^{\frac{-i}{\hbar v} \int_{-\infty}^Z V_{PT}(R') dZ'}$

Analytical form → Fast



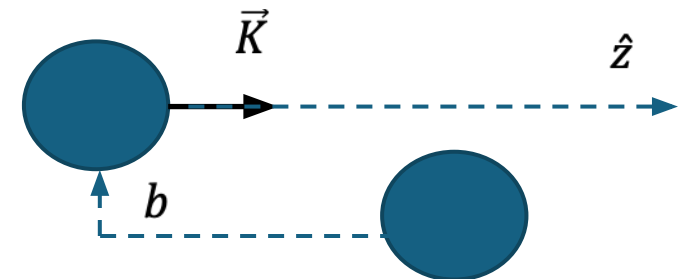
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Analytical form $\rightarrow$ Fast

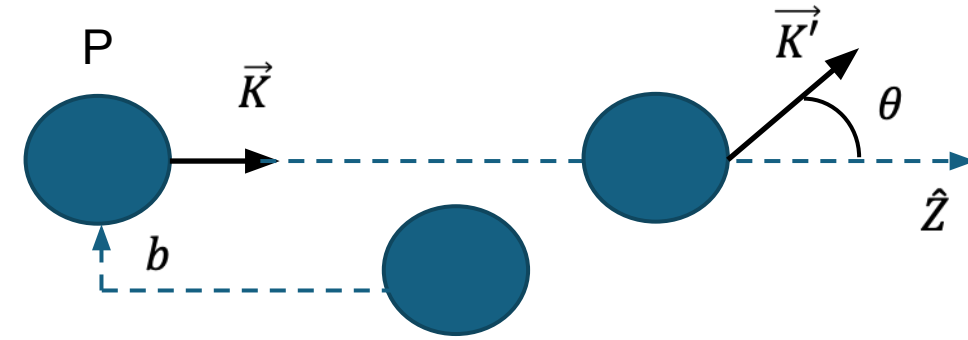
Semiclassical interpretation : P follows a straightline trajectory at constant impact parameter b along which it accumulates a phase due to its interaction with T



How to obtain elastic scattering cross sections?

Elastic scattering for a certain momentum:

$$\underline{\frac{d\sigma}{d\theta} = |f(\theta)|^2}$$



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Elastic scattering for a certain momentum:

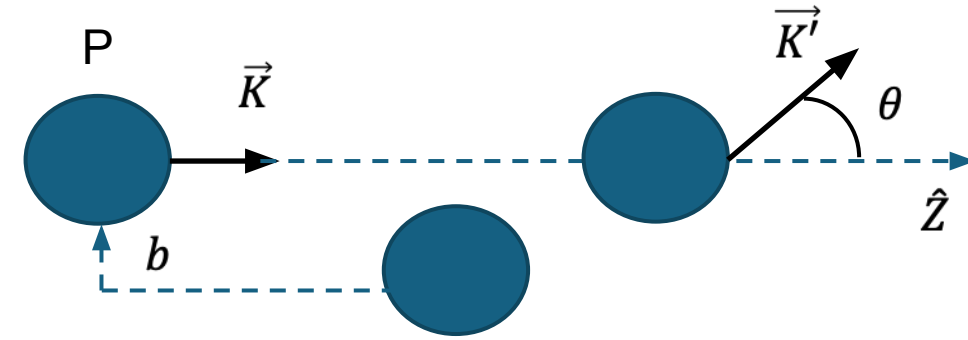
$$\frac{d\sigma}{d\theta} = |f(\theta)|^2$$

One can show

$$f^{eik}(\theta) = -iK \int_0^\infty db \, b J_0(qb) [e^{i\chi(b)} - 1]$$

$$\text{with } \vec{q} = \vec{K}' - \vec{K}$$

$$\text{Eikonal phase } \chi(b) = -\frac{1}{\hbar v} \int_{-\infty}^{\infty} V_{PT}(R) dZ$$



All the information is in the eikonal phase!

Integrated cross sections and questions

- Using the eikonal phase

$$\chi(b) = -\frac{1}{\hbar v} \int_{-\infty}^{\infty} V_{PT}(R) dZ$$

- Elastic scattering cross sections

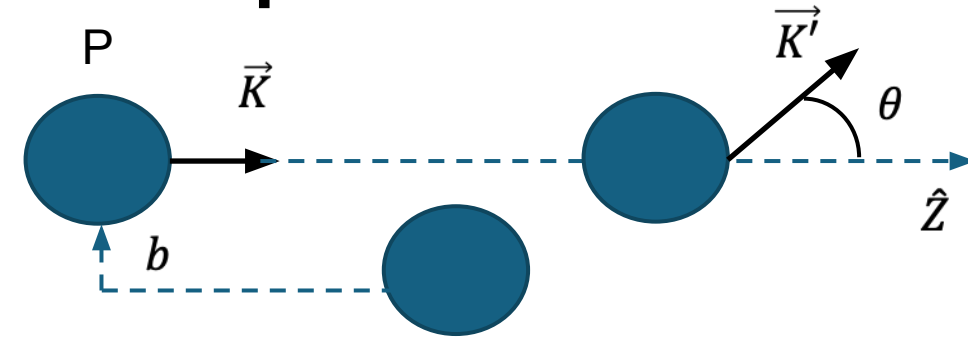
$$\sigma_{el} = 2\pi \int_0^{+\infty} b db P_{el}(b) \text{ with } P_{el}(b) = |1 - e^{i\chi(b)}|^2$$

- Absorption cross section

$$\sigma_{abs} = 2\pi \int_0^{+\infty} b db P_{abs}(b) \text{ with } P_{abs}(b) = 1 - |e^{i\chi(b)}|^2$$

- Total cross section

$$\sigma_{tot} = \sigma_{el} + \sigma_{abs} = 2\pi \int_0^{+\infty} b db 2[1 - \text{Re}(e^{i\chi(b)})]$$



1) What happens to P_{el} and P_{abs} if the potential is nil?

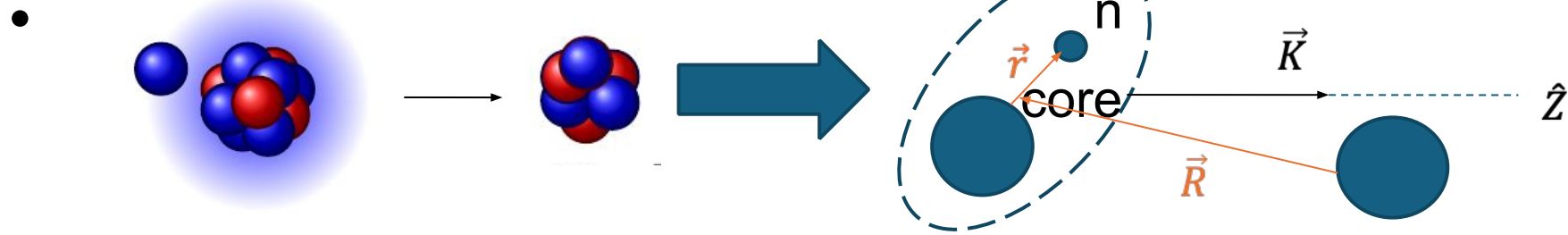
2) What happens to P_{el} and P_{abs} when b is very large (outside of the range of the potential)?

3) What happens to P_{abs} if the potential is real?

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- Three-body scattering problem

$$[T_{\vec{R}} + h_{cn} + V_{cT}(R_{cT}) + V_{nT}(R_{nT})]\Psi(\vec{R}, \vec{r}) = E\Psi(\vec{R}, \vec{r})$$

- Initial condition

$$\Psi(\vec{R}, \vec{r}) \rightarrow_{Z \rightarrow -\infty} \phi_{gs}(\vec{r})e^{iKZ}$$

With $h_{cn}\phi_{gs}(\vec{r}) = [T_{\vec{r}} + V_{cn}(r)]\phi_{gs}(\vec{r}) = \epsilon_{gs}\phi_{gs}(\vec{r})$

Similar process as before... except that now it depends on the **projectile** system

- Eikonal factorization: $\Psi(\vec{R}, \vec{r}) = e^{iKZ} \widehat{\Psi(\vec{R}, \vec{r})}$ and eikonal approximation $\Delta\Psi \ll K \frac{\partial\Psi}{\partial Z}$

- Using the eikonal factorization in the 3-body Schrödinger equation & energy conservation

$$i \hbar v \frac{\delta}{\delta Z} \widehat{\Psi(\vec{R}, \vec{r})} = (h_{cn} - \epsilon_{gs}) \widehat{\Psi(\vec{R}, \vec{r})} + (V_{cT}(R_{cT}) + V_{nT}(R_{nT})) \widehat{\Psi(\vec{R}, \vec{r})} \text{ with } v = \frac{\hbar K}{\mu}$$

(dynamical eikonal approximation)

- Usual Eikonal assumes the adiabatic approximation: $h_{cn} \sim \epsilon_{gs}$

$$i \hbar v \frac{\delta}{\delta Z} \widehat{\Psi(\vec{R}, \vec{r})} = (V_{cT}(R_{cT}) + V_{nT}(R_{nT})) \widehat{\Psi(\vec{R}, \vec{r})}$$

- Eikonal Solution: $\Psi(\vec{R}, \vec{r}) = e^{ikz} \widehat{\Psi(\vec{R}, \vec{r})} = e^{ikz} e^{\frac{-i}{\hbar v} \int_{-\infty}^Z [V_{cT}(R'_{cT}) + V_{nT}(R'_{nT})] dZ'} \phi_{gs}(\vec{r})$
 $\rightarrow$ 2 eikonal phases: one for each fragment of P !

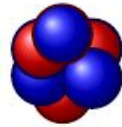
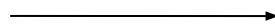
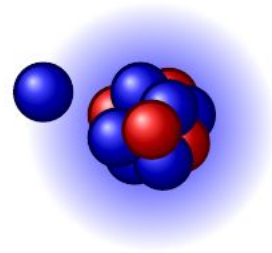
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Breakup and knockout cross sections

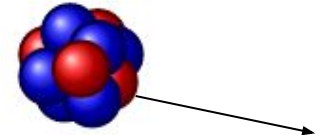
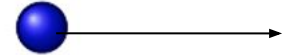
Diffractive breakup: neutron elastically scattered:

$$\sigma_{diff} = 2\pi \int d\epsilon \int b db \left| \langle \phi_\epsilon | 1 - S_n S_c | \phi_{gs} \rangle \right|^2$$

with ϕ_ϵ c-n scattering state

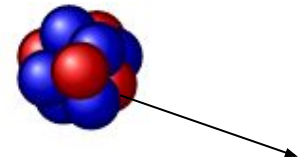


diffractive breakup



detector

knockout



detector

Stripping: neutron inelastically scatters with the target:

$$\sigma_{str} = 2\pi \int b db \left| \langle \phi_{gs} | (1 - |S_n|^2) | S_c |^2 | \phi_{gs} \rangle \right|^2$$

Knockout: $\sigma_{ko} = \sigma_{str} + \sigma_{diff}$

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Few questions to answer

1. Understanding:

- Do you understand what the python notebook `ComparisonEikonalExact.ipynb` is doing?
- Can you interpret the eikonal S-matrix?
- What happens to the eikonal S-matrix if you take a nil potential? Is it what you expect?
- What happens to the S-matrix & the eikonal phase when $b \rightarrow \infty$? Is it what you expect?
- What happens to the S-matrix if the potential is real? Comment on the value of the absorption cross section.

1. Develop a cell to compare the integrated and angular cross sections obtained with different beam energy? (make sure that you results are converged)

1. Develop a cell to compare the integrated and angular cross sections obtained with different rR ?

1. Develop a cell to compare the integrated and angular cross sections obtained with different imaginary depths?

Bonus :

1. Considering a knockout reaction, estimate from the S-matrix the knockout probability as a function of b . What can you deduce about the peripherality of knockout reactions
1. What happens if you change the range of the n-T potential?
1. What happens if you change the range of the c-T potential?

Some useful references about eikonal

- Suzuki, Y., Yabana, K., Lovas, R.G., & Varga, K. (2003). Structure and Reactions of Light Exotic Nuclei (1st ed.). CRC Press.
- P.G. Hansen and J. A. Tostevin, Annu. Rev. Nucl. Sci. 53, 219 (2003)
- D. Baye and P. Capel, in Clusters in Nuclei, Vol. 2, edited by C. Beck (Springer, Heidelberg, 2012), vol. 848.

Useful website for those who make reaction calculations

<https://sites.google.com/view/opticalpotentials/>



Optical potentials in nuclear phy...

Reaction codes

Optical potential parameterizatio...

Other tools



Optical potentials in nuclear physics