

INTERACTING DARK SECTOR WITH INTRINSIC ENTROPY COUPLINGS

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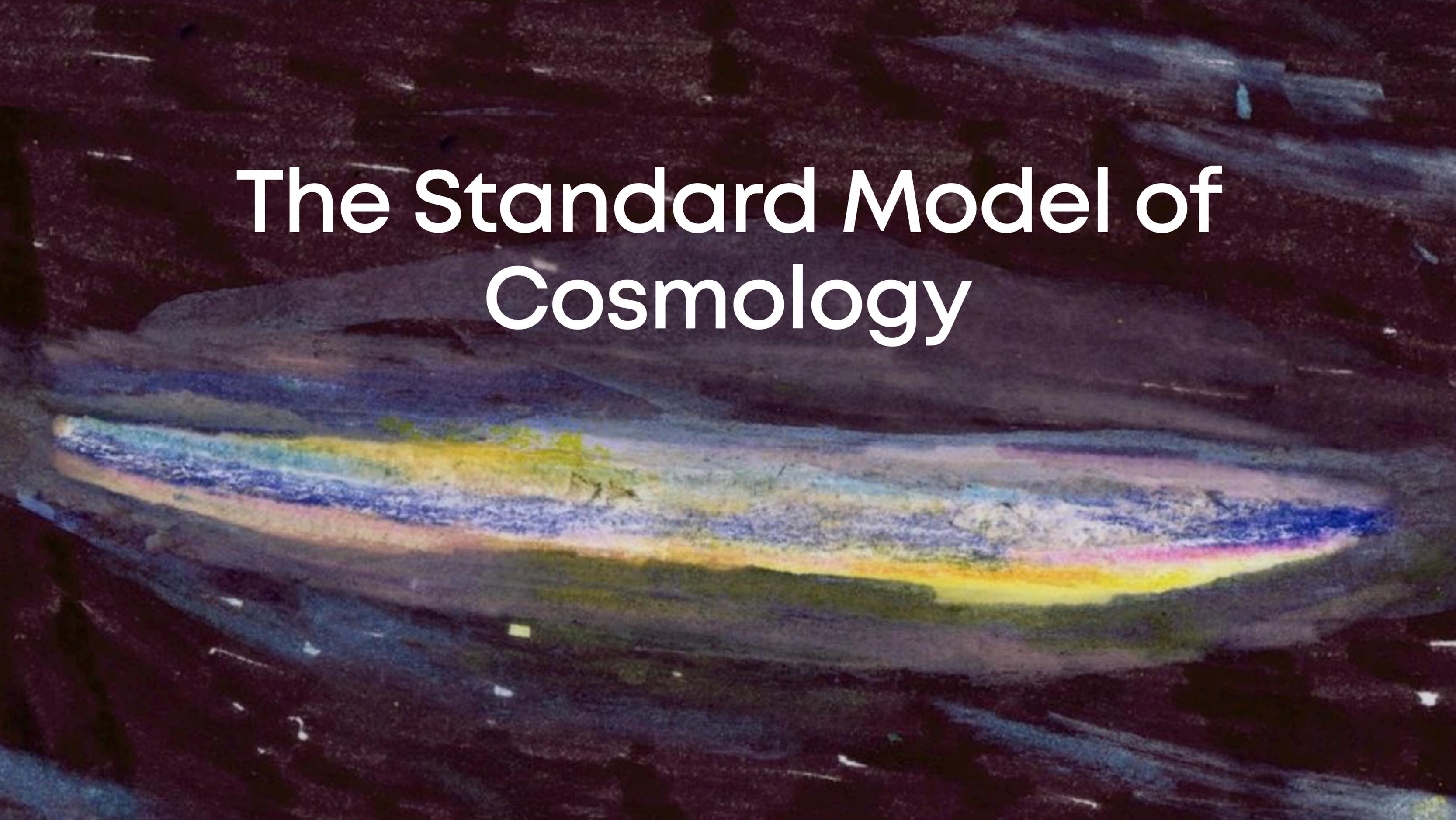
Illustrations: Inês Viegas Oliveira (ivoliveira.com)

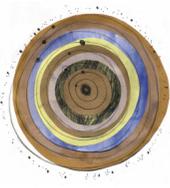
Based on ongoing work with:

- Erik Jensko, UCL
- Vivian Poulin, U. Montpellier

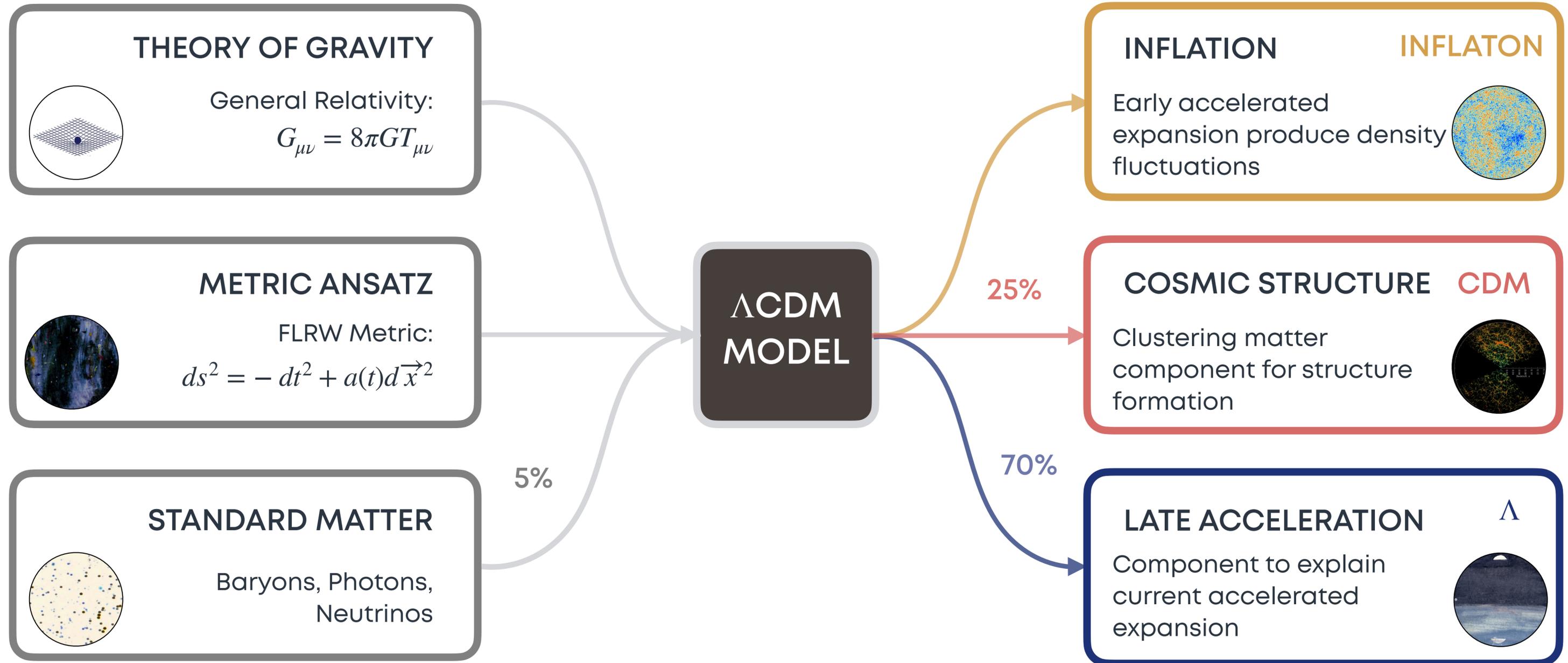
On arXiv tomorrow :)

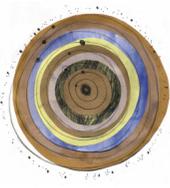
The Standard Model of Cosmology

The background of the slide is a dark, textured, and layered landscape. It features a prominent horizontal band of vibrant colors, including yellow, green, blue, and purple, which resembles a rainbow or a cross-section of geological strata. The overall appearance is that of a stylized, abstract natural scene.



The Lambda Cold Dark Matter Model





Challenges to the Λ CDM Model

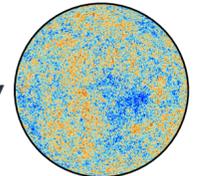
The Λ CDM model relies on:

- Inflation but needs firm theoretical grounds: primordial power spectrum of quantum fluctuations (simplest parameterisation in terms of spectral index and amplitude)
- Dark matter being a pressureless fluid of unknown nature/origin and no detection success (new particle(s) in the SM)
- Dark energy being a cosmological constant (Λ) with unknown nature/origin (vacuum energy, properties of empty space, etc)

INFLATION

INFLATON

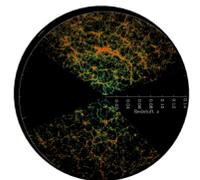
Early accelerated expansion produce density fluctuations



COSMIC STRUCTURE

CDM

Clustering matter component for structure formation

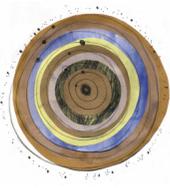


LATE ACCELERATION

Λ

Component to explain current accelerated expansion





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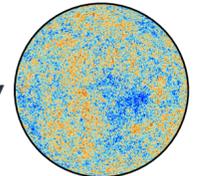
Cosmic tensions may signal that Λ CDM is incomplete:

- Anomalies in the CMB: lensing, curvature, etc
- The matter clustering S_8 tension
- The Hubble/ H_0 expansion rate tension

INFLATION

INFLATON

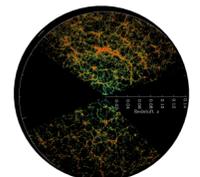
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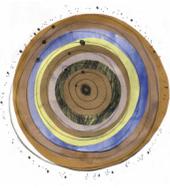


LATE ACCELERATION

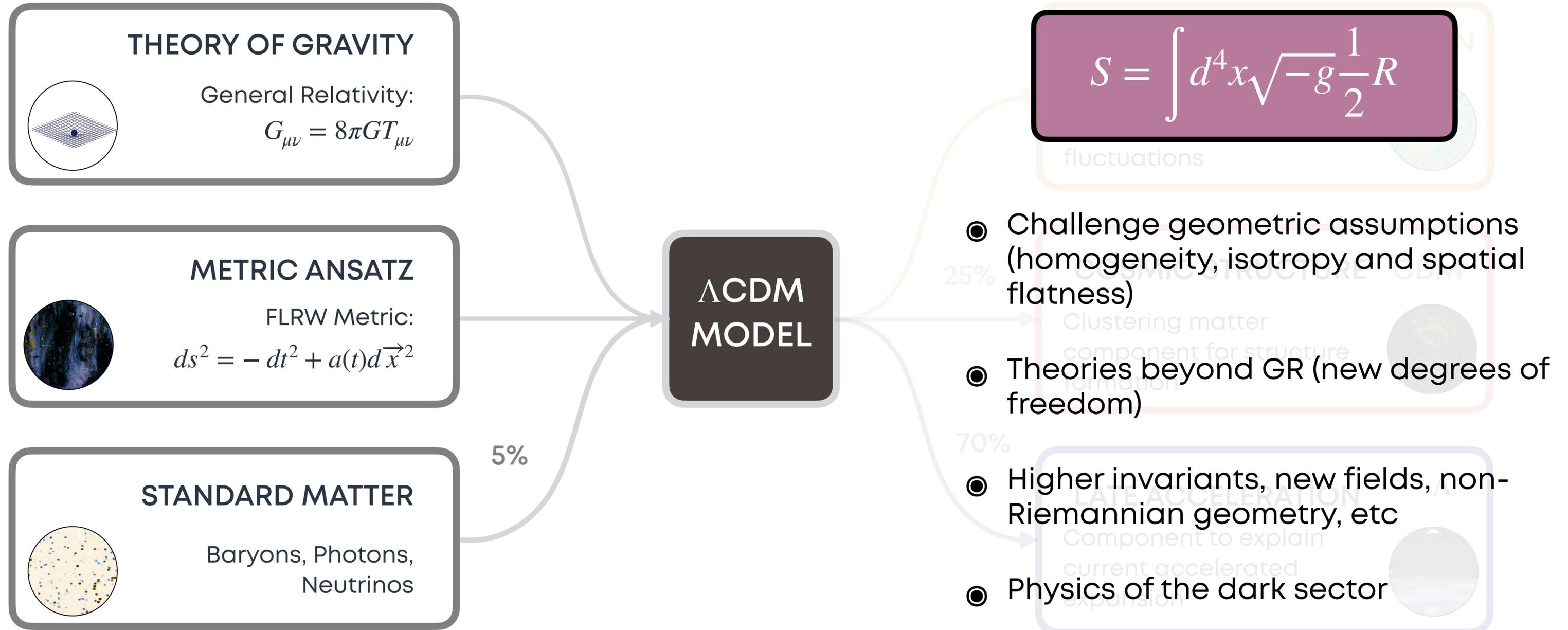
Λ

Component to explain current accelerated expansion



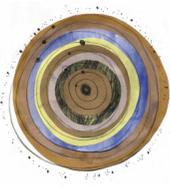


Going Beyond the Standard Model



Interacting dark sector with intrinsic entropy couplings

Based on: [E. Jensko, E. M. Teixeira, V. Poulin: [arxiv:2603.tomorrow](https://arxiv.org/abs/2603.tomorrow)]



Interacting dark sector

- Background cosmological couplings between dark matter (DM) and dark energy (DE) can be imposed at the level of the field equations:
- However, these are just phenomenological modifications:
- Prone to instabilities [Valiviita et. al. arxiv:0804.0232]
- Ambiguous definition of cosmological perturbations (e.g., added by hand)
- Couplings often heavily constrained by background observations
- Focus on **Lagrangian formulations**

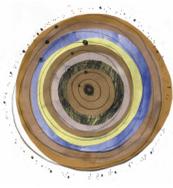


$$\dot{\rho}_{\text{DM}} + 3H(\rho_{\text{DM}} + p_{\text{DM}}) = Q$$

$$\dot{\rho}_{\text{DE}} + 3H(\rho_{\text{DE}} + p_{\text{DE}}) = -Q$$

$$\delta'_{\text{DM}} + 3\mathcal{H} \left(\frac{\delta p_{\text{DM}}}{\delta \rho_{\text{DM}}} + w_{\text{DM}} \right) \delta_{\text{DM}} = - (1 + w_{\text{DM}}) (\theta_{\text{DM}} - 3\Phi') + F(Q, \delta Q)$$

$$\theta'_{\text{DM}} + \left[\mathcal{H}(1 - 3w_{\text{DM}}) + \frac{w'_{\text{DM}}}{1 + w_{\text{DM}}} \right] \theta_{\text{DM}} = k^2 \left[\Psi + \frac{\delta p_{\text{DM}}}{\delta \rho_{\text{DM}}} \frac{\delta_{\text{DM}}}{1 + w_{\text{DM}}} + G(Q, \delta Q) \right]$$



Cosmological fluids

Perfect fluid energy-momentum tensor:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu + pg_{\mu\nu}, \quad \nabla^\mu T_{\mu\nu} = 0$$

Key thermodynamic quantities:

- ⦿ particle number density n
- ⦿ intrinsic entropy/particle s
- ⦿ energy density $\rho = \rho(n, s)$
- ⦿ pressure p
- ⦿ temperature T
- ⦿ chemical potential μ
- + four components of the 4-velocity u^μ with $u^\mu u_\mu = -1$

⦿ EoM (continuity & Euler equations) from projections along u_μ and $h_{\mu\nu}$

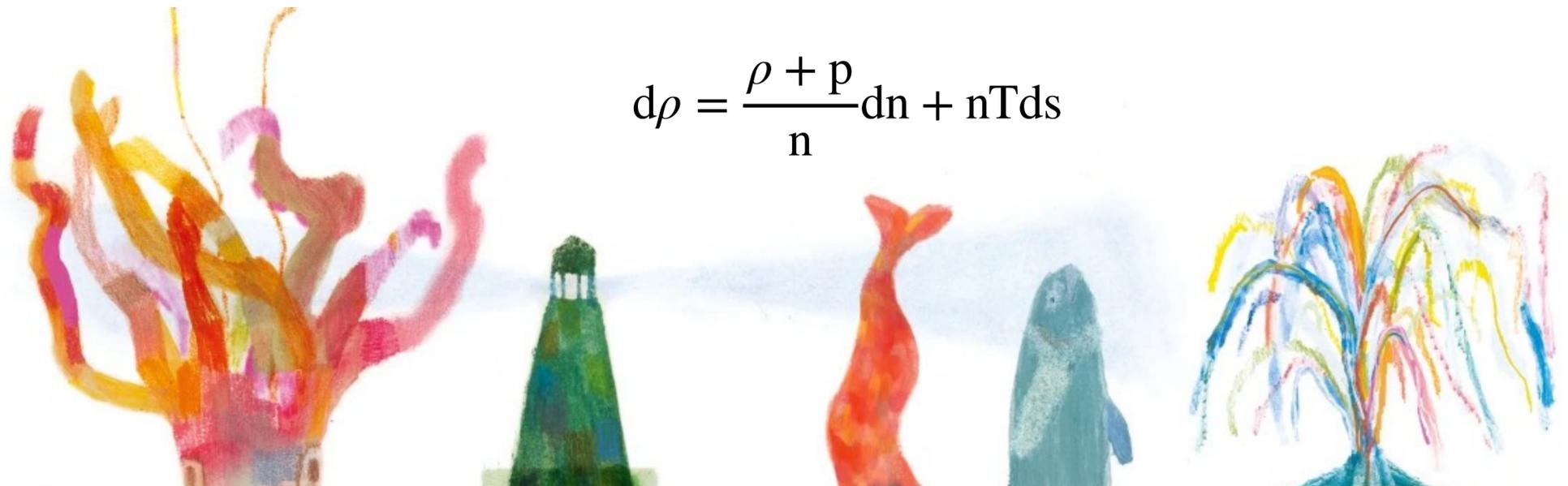
⦿ Thermodynamic identities

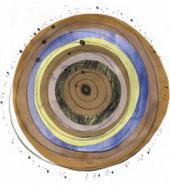
$$p(n, s) = n \left(\frac{\partial \rho}{\partial n} \right)_s - \rho$$

$$T(n, s) = \frac{1}{n} \left(\frac{\partial \rho}{\partial s} \right)_n$$

⦿ Satisfying 1st law of thermodynamics

$$d\rho = \frac{\rho + p}{n} dn + nTds$$





Perfect fluid action

- Model DM as a **perfect fluid** through action with variables $g_{\mu\nu}$, n^μ , s with Lagrange multipliers φ , θ - define the fluid's internal structure and and conservation laws covariantly

$$S_{\text{Brown}} = \int \sqrt{-g} \left[-\rho(n, s) + n^\mu (\varphi_{,\mu} + s\theta_{,\mu}) \right] d^4x$$

[J. D. Brown, Class. Quant. Grav. 10, 1579 (1993), arXiv:gr-qc/9304026]

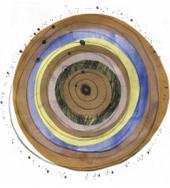
where we introduced the number flux $n^\mu := nu^\mu$ with $n = \sqrt{n^\mu n_\mu}$

constraints

- Variations give perfect fluid $T_{\mu\nu}$ with conservation equation $\nabla^\mu T_{\mu\nu} = 0$
- Additional constraints** preserving particle number and entropy fluxes

$$\nabla_\mu (nu^\mu) = 0, \quad \nabla_\mu (snu^\mu) = 0 \quad \implies \quad u^\mu \nabla_\mu s = 0$$





Cosmological fluids

$$\nabla_{\mu}(nu^{\mu}) = 0$$

$$\nabla_{\mu}(nsu^{\mu}) = 0$$

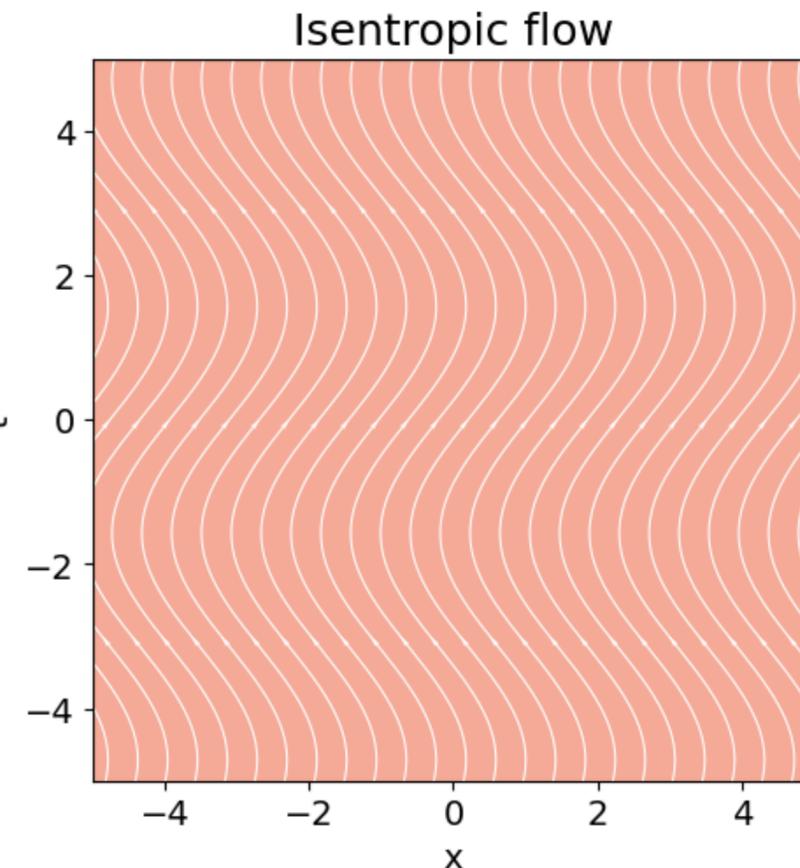
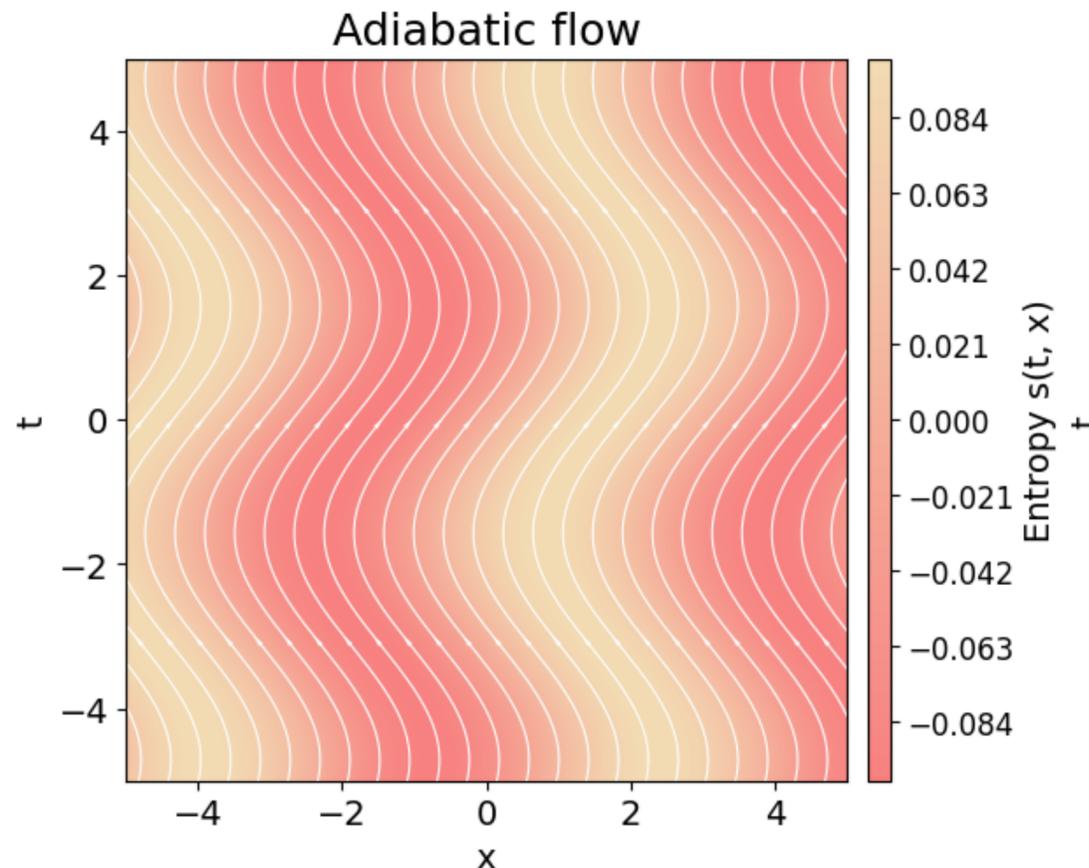
- Conservation of particle number and entropy along fluid flow lines
- Taken together these imply the entropy per particle is constant along fluid flow

adiabatic: $u^{\mu} \nabla_{\mu} s = 0 \not\Rightarrow$ isentropic: $\nabla_{\mu} s = 0$

General fluid condition:

uniform entropy per particle s along the fluid flow (no shocks or heat conduction):

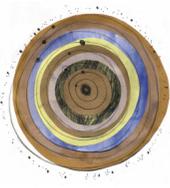
$$\rho = \rho(n, s)$$



Usual assumption in cosmology:

the entropy per particle s is itself time and spatially uniform (irrotational flow):

$$\rho = \rho(n)$$



Perfect fluid dark matter in FLRW

| Traditional CDM | General perfect fluid DM | “Entropic” CDM |
|--|---|---|
| Barotropic and isentropic ($\nabla_\mu s = 0$) | Non-barotropic and adiabatic ($u^\mu \nabla_\mu s = 0$) | Barotropic and adiabatic ($u^\mu \nabla_\mu s = 0$) |
| $\rho(n) = nm$ | $\rho(n, s)$ arbitrary | $\rho(n, s) = nm\gamma(s)$ |
| $p = 0$ | $p \neq 0$ | $p = 0$ |
| $c_s^2 = 0$ | $c_s^2 \neq 0$ | $c_s^2 = 0$ |
| $\Gamma = 0$ | $\Gamma \neq 0$ | $\Gamma = 0$ |
| $\bar{s} = 0$ | $\bar{s} = \text{constant}$ | $\bar{s} = \text{constant}$ |
| $\delta s = 0$ | $\delta s = \delta s(x, y, z)$ | $\delta s = \delta s(x, y, z)$ |
| $T = 0$ | $T \neq 0$ | $T = m\gamma'(s)$ |

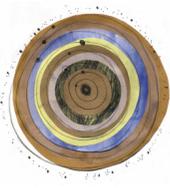
$$\delta p = c_a^2 \delta \rho + \Gamma$$

with non-adiabatic pressure perturbation

$$\Gamma = \left. \frac{\partial p}{\partial s} \right|_\rho \delta s$$

- intrinsic ~ fluctuations in fluid's internal entropy ($\Gamma = 0 \not\Rightarrow \delta s = 0$)

- relative ~ relative density isocurvature perturbations



Dark sector interactions

- Interactions in the dark sector accomplished through non-minimally coupled scalar fields ϕ

$$S_{\text{tot}} = \int \sqrt{-g} f(n, s, \phi, \dots) d^4x + S_{\text{Brown}} + S_{\phi} + S_{\text{SM}}$$

Many different choices of interacting terms in the literature:

- Algebraic (“coupled quintessence”): $f(n, \phi)$
- Derivative (“momentum coupling”): $f(n)n^\mu \partial_\mu \phi$
- General with $Y = \partial_\mu \phi \partial^\mu \phi$ and $Z = u^\mu \partial_\mu \phi$: $f(n, \phi, Y, Z)$

[Boehmer et al., arxiv:1501.06540]

[Boehmer et al., arxiv:1502.04030]

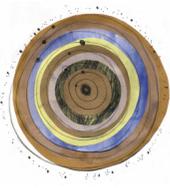
[Pourtsidou et al., arxiv:1307.0458]

Type-1: $Y + V(\phi) + e^{\alpha(\phi)}$ Type-2: $Y + V(\phi) + nh(Z)$ Type-3: $Y + V(\phi) + \gamma(Z) + f(n)$

E.g. Type-3 models is pure momentum-transfer up to linear order - background unchanged

Interactions





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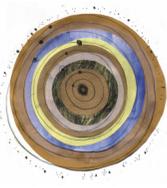
All past studies neglect the role of entropy in couplings $f(n, s, \phi, \partial\phi)$ by assuming $\partial_{\mu}s = 0$

Type-1: $Y + V(\phi) + e^{\alpha(\phi)}$ Type-2: $Y + V(\phi) + nh(Z)$ Type-3: $Y + V(\phi) + \gamma(Z) + f(n)$

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Interactions





Entropy couplings in the dark sector

- Couple the intrinsic entropy of dark matter s in $\rho_c(n_c, s) = m_c n_c \gamma(s)$ to the dark energy scalar ϕ - excite the otherwise hidden degrees of freedom associated with the entropy perturbations in barotropic fluids

- For this we consider pure entropy-scalar interactions:

$$\mathcal{L}_{\text{int}}(\phi, s, \nabla_\mu \phi, \nabla_\mu s, \dots)$$

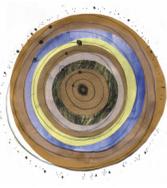
$$\mathcal{L}_{\text{alg}} = g(\phi, s)$$

$$\mathcal{L}_{\text{der}} = h(\nabla_\mu \phi, \nabla_\mu s)$$

- Later distinguish into algebraic and first-order derivative model
- Generally derived from an interaction action

$$S_{\text{int}} = - \int \sqrt{-g} f(s, \phi, \mathcal{S}) d^4x \quad \text{with} \quad \mathcal{S} := \nabla^\mu \phi \nabla_\mu s$$





Equations of motion

- Modified Einstein field equations

$$G_{\mu\nu} = \kappa \left(T_{\mu\nu} + T_{\mu\nu}^{\phi} + T_{\mu\nu}^{(\text{int})} \right)$$

with interaction EM tensor (with heat-fluxes $q_{\mu} = u^{\nu} h_{\mu}^{\lambda} T_{\nu\lambda}$):

$$T_{\mu\nu}^{(\text{int})} = -g_{\mu\nu} f + 2 \frac{\partial f}{\partial \mathcal{S}} \nabla_{(\mu} \phi \nabla_{\nu)} \mathcal{S}$$

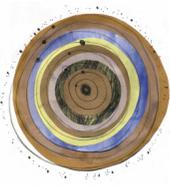
- Total effective dark energy tensor $\tilde{T}_{\mu\nu}^{(\phi)} := T_{\mu\nu}^{(\phi)} + T_{\mu\nu}^{(\text{int})}$ with coupling current

$$J_{\nu} := \nabla^{\mu} \tilde{T}_{\mu\nu}^{(\phi)} = \nabla_{\mu} (f_{, \mathcal{S}} \nabla^{\mu} \phi) \nabla_{\nu} \mathcal{S} - f_{, \mathcal{S}} \nabla_{\nu} \mathcal{S}$$

which acts only in the perpendicular direction $u^{\nu} J_{\nu} = 0$ (from $u^{\nu} \nabla_{\nu} \mathcal{S} = 0$)

Pure “momentum” transfer (up to linear order) - preserve background dynamics





Cosmology

● Conformal Newtonian gauge at linear order: $ds^2 = -a(\tau)^2(1 + 2\Phi)d\tau^2 + a^2(1 - 2\Psi)(dx^2 + dy^2 + dz^2)$

● Entropy conservation $u^\mu \nabla_\mu s = 0$ to linear order

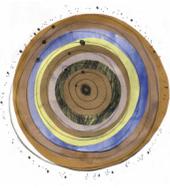


● Trivial background modification with effective potential $\hat{V}(\bar{\phi}) = V(\bar{\phi}) + f(\bar{\phi}, s_0, 0)$

$$\bar{\rho}_{\text{DE}} = \frac{1}{2a^2} \bar{\phi}'^2 + \hat{V}(\bar{\phi}), \quad \bar{p}_{\text{DE}} = \frac{1}{2a^2} \bar{\phi}'^2 - \hat{V}(\bar{\phi})$$

● Uncoupled dark sector in the background - perturbative effects only

$$\bar{\rho}'_I + 3\mathcal{H}(\bar{p}_I + \bar{\rho}_I) = 0 \quad \implies \quad \text{indistinguishable from uncoupled quintessence or } \Lambda\text{CDM}$$



Cosmological perturbations

- Split interaction term and take $p_c = \delta p_c = 0$ for simplicity

$$f(\phi, s, \nabla_\mu \phi \nabla^\mu s) = g(\phi, s) + h(\nabla_\mu \phi \nabla^\mu s)$$

Algebraic Derivative

- Unchanged continuity equation but modified Euler equation with scale-dependent source

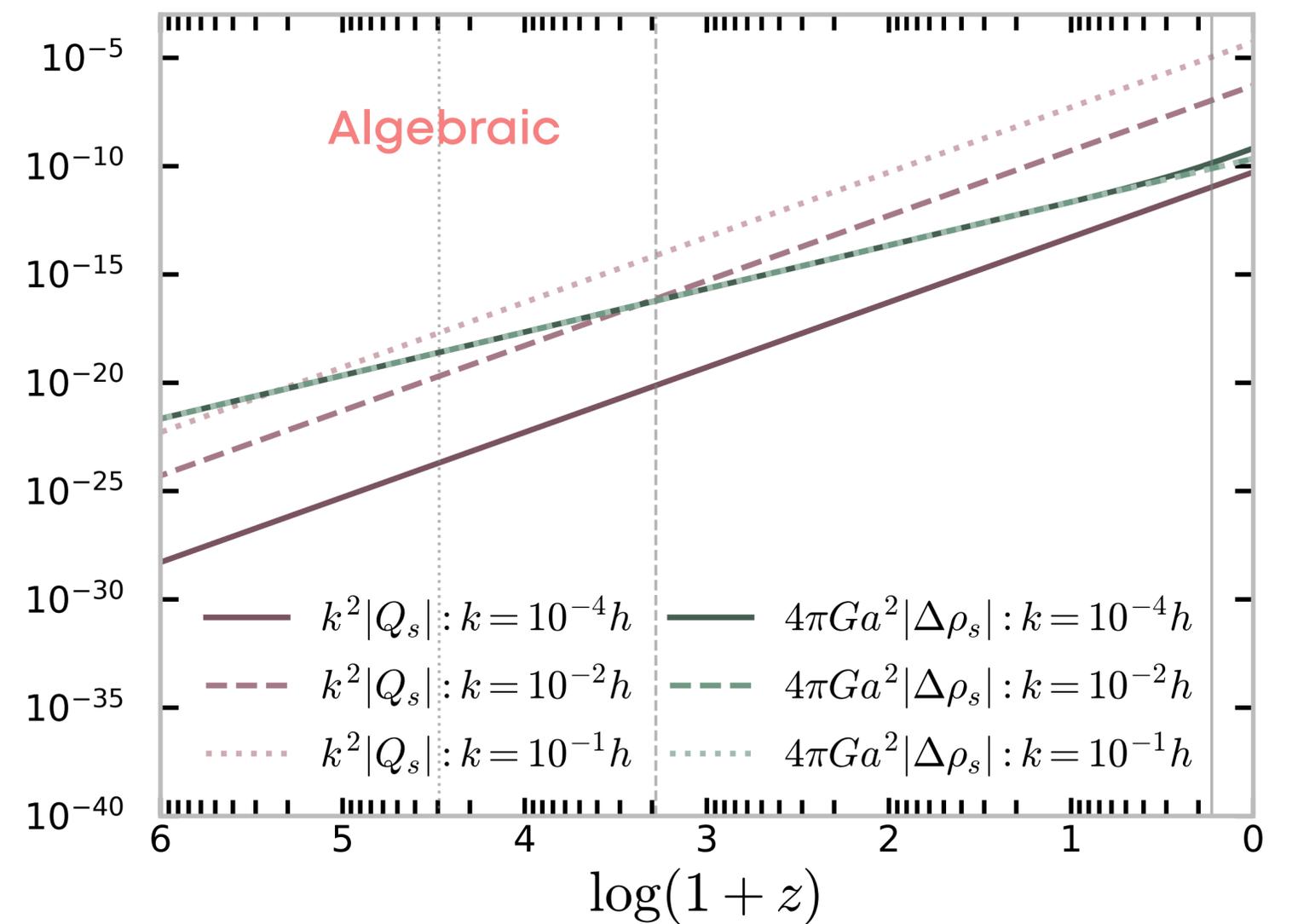
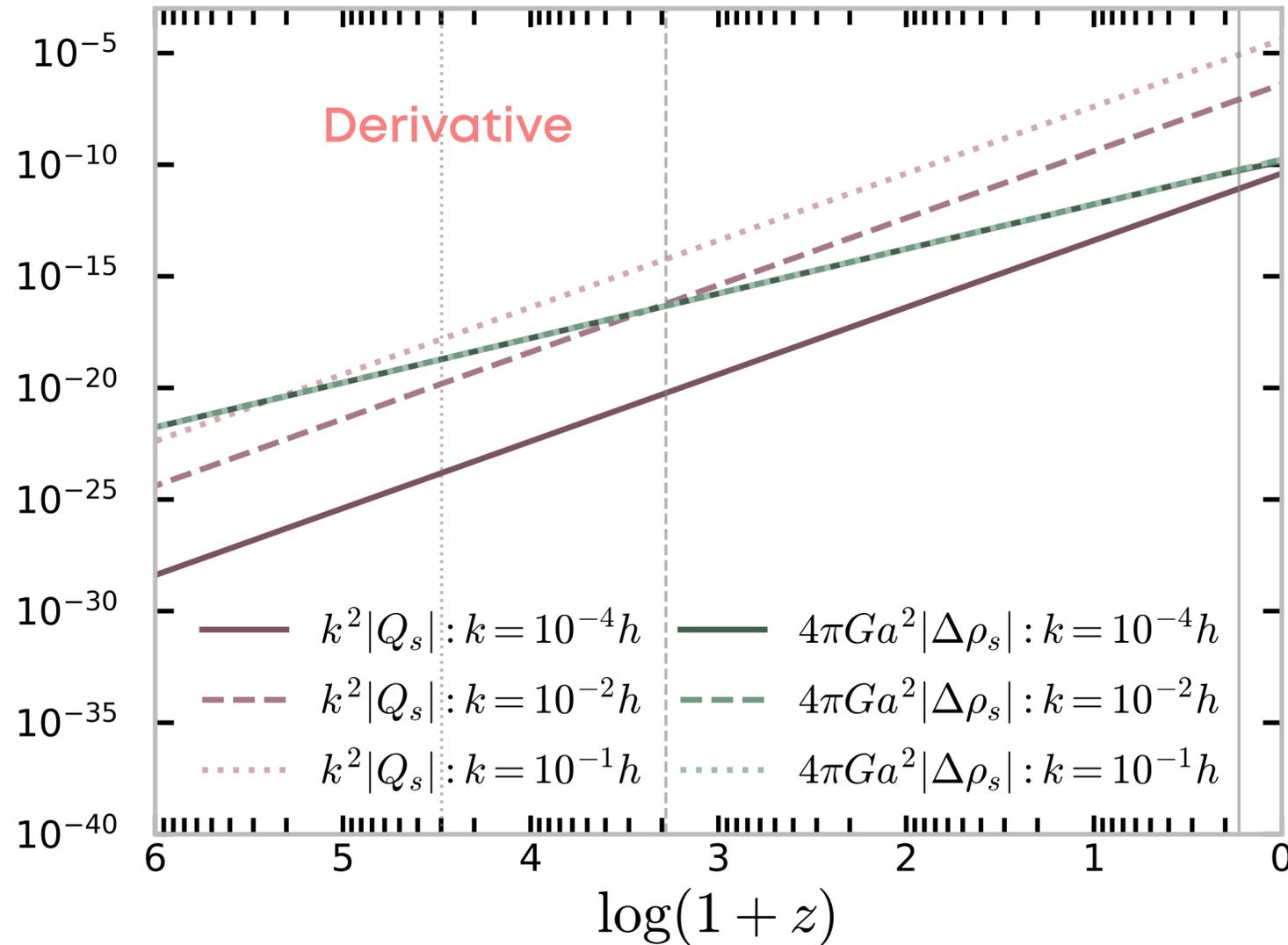
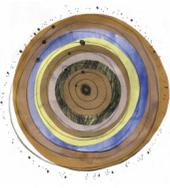
$$\delta'_c + \theta_c - 3\Psi' = 0$$

$$\theta'_c + \theta_c \mathcal{H} - k^2 \Psi = -k^2 Q_s \delta s(k), \quad \text{with} \quad Q_s := \frac{\bar{g}_{,s} - h_0 \hat{V}_{,\bar{\phi}}}{\bar{\rho}_c}$$

- In the quasi-static limit ($k \gg \mathcal{H}$) this becomes:

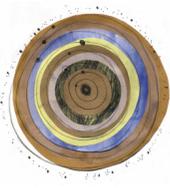
$$\delta''_c + \mathcal{H} \delta'_c - 4\pi G a^2 [\bar{\rho}_c \delta_c + \Delta \rho_s(k, a) \delta s] \simeq k^2 Q_s(a) \delta s$$

- The coupling appears as “fifth-force” $Q_s(a)$ and “effective density source” contribution $\Delta \rho_s(k, a)$



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© The coupling appears as “fifth-force” $Q_s(a)$ and “effective density source” contribution $\Delta\rho_s(k, a)$



Observational signatures

Pure derivative interaction with scale invariant entropy perturbations with standard DE exponential potential

- In the QS limit algebraic and derivative couplings reduce to similar effective source

$$k^2\Psi \propto -a^2(\delta\rho_m + \delta\rho_{\text{DE}})$$

$$k^2(\Psi' + \mathcal{H}\Psi) \propto a^2(\theta_m + \theta_{\text{DE}})$$

- Modification in θ_{DM} and $\theta_{\text{DE}} = -k^2(h_0\delta s + \delta\phi)/\phi'$

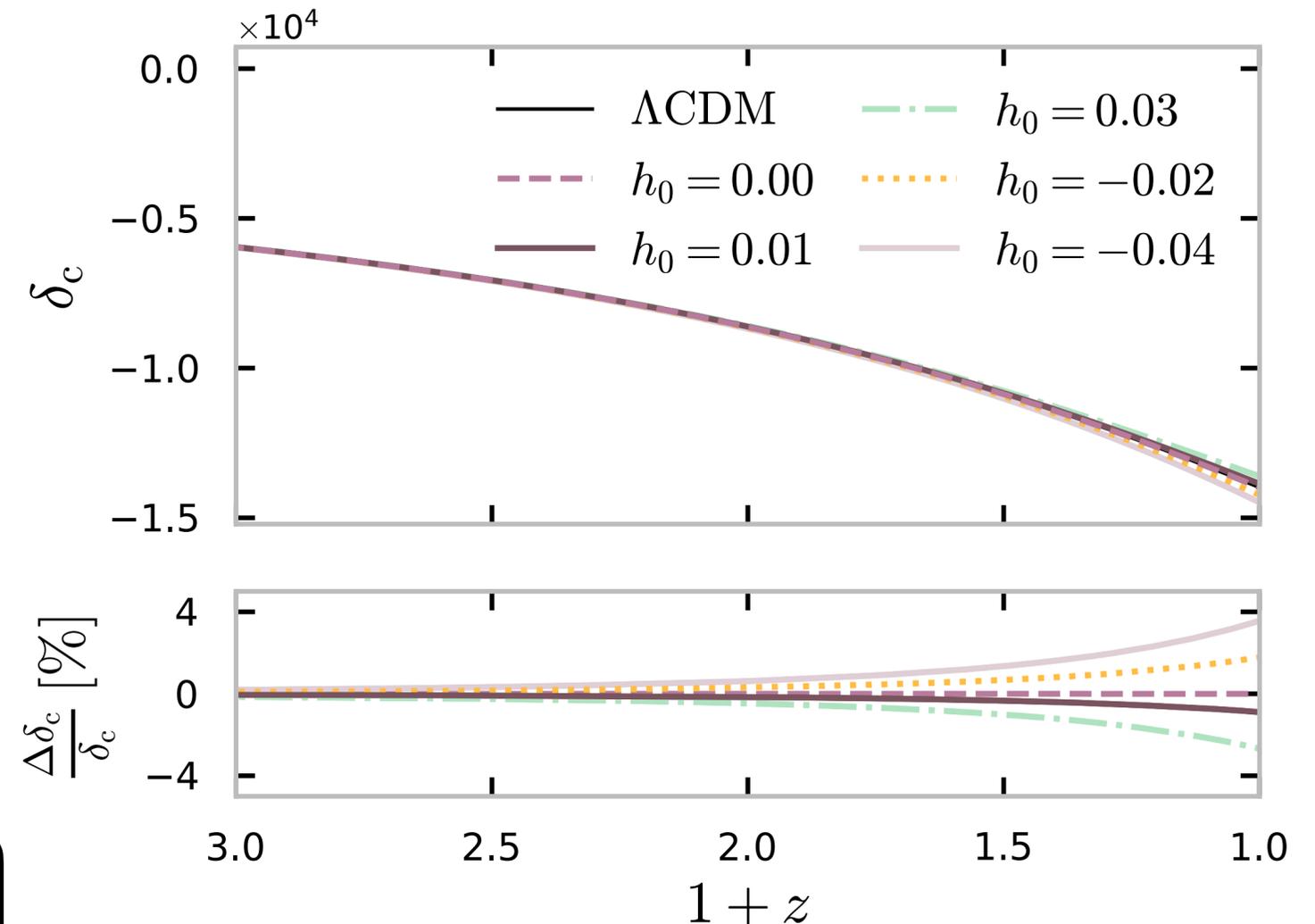
- Cut-off scale to tame large effects of coupling at small scales ($\propto k^2 Q_s$ in Euler equation)

$$Q_s^{\text{der}}(a) = -\frac{h_0}{\bar{\rho}_c} \hat{V}_{,\bar{\phi}}$$

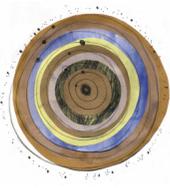
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$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

$$\delta s(k) = \mathcal{A}_e \left(k/k_p \right)^n \exp \left(- (k/k_c)^p \right)$$



$$[k_8 = 0.125h \text{ Mpc}^{-1}, \mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0]$$



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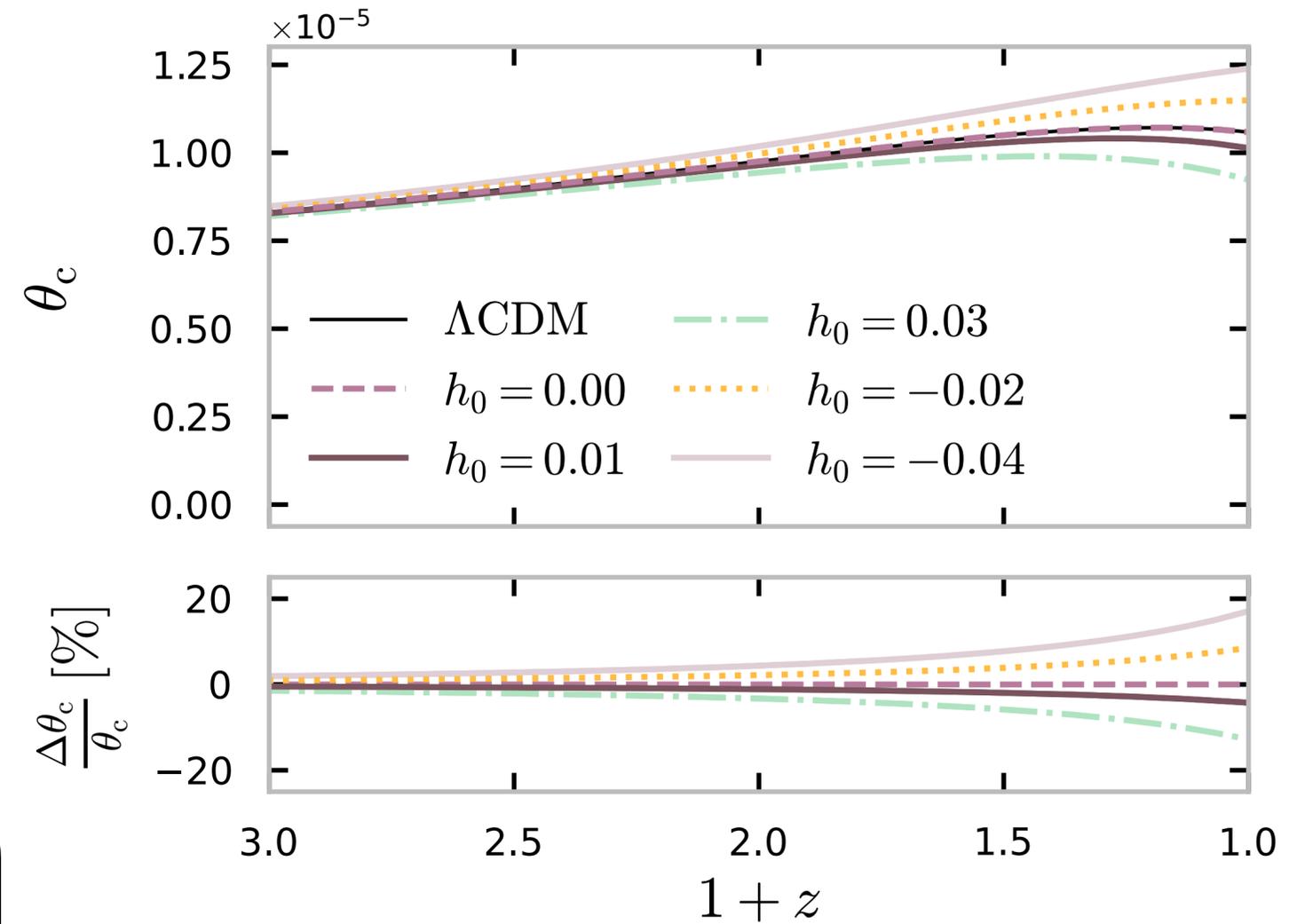
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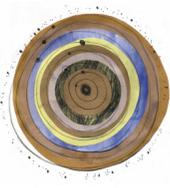
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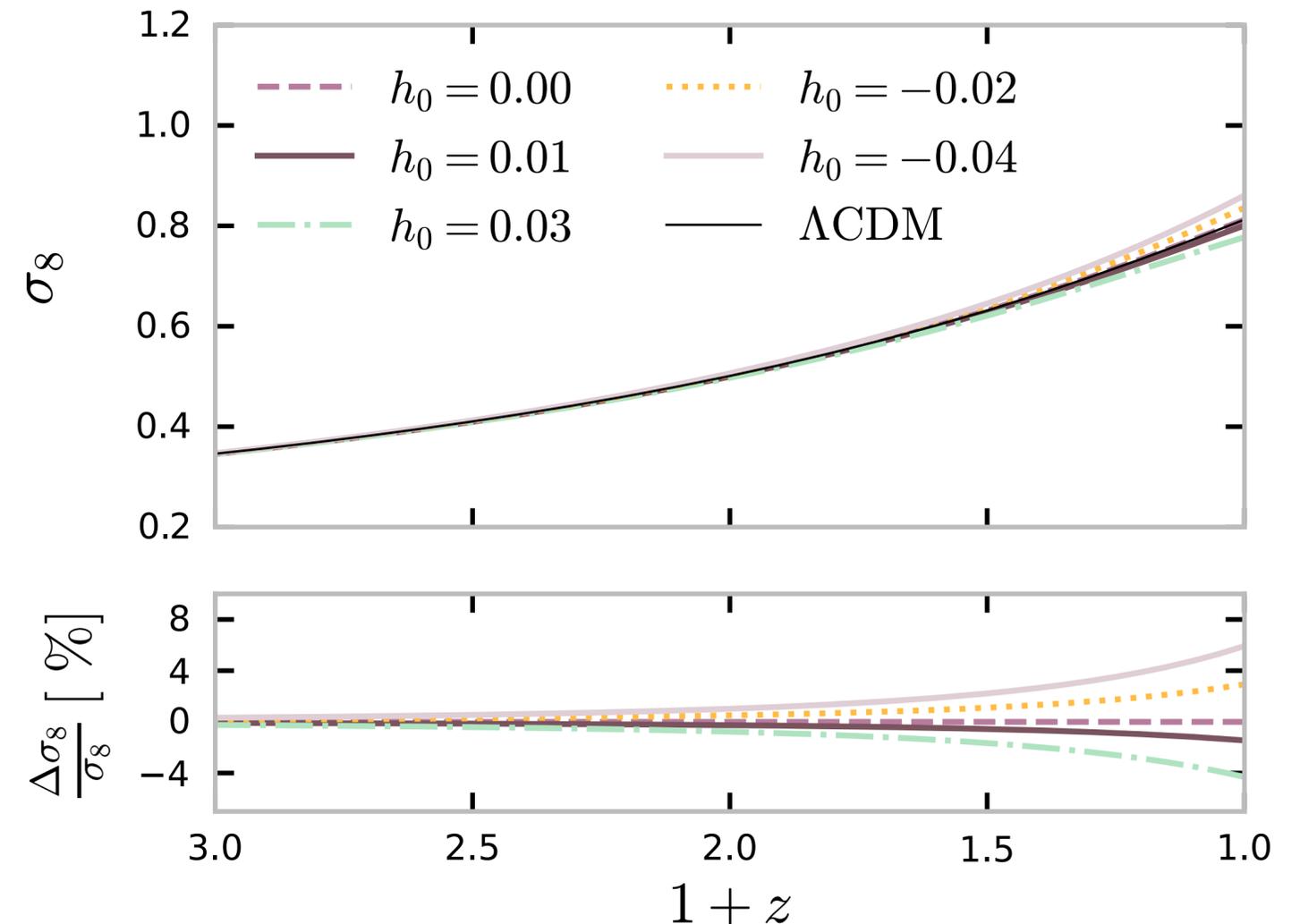
- Cut-off scale to tame large effects of coupling at small scales ($\propto k^2 Q_s$ in Euler equation)

$$Q_s^{\text{der}}(a) = -\frac{h_0}{\bar{\rho}_c} \hat{V}_{,\bar{\phi}}$$

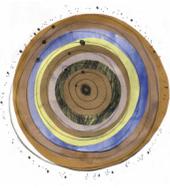
$$\Delta\rho_s^{\text{der}}(k, a) = -\frac{k^2 h_0}{k^2 + a^2 m_{\text{eff}}^2} \hat{V}_{,\bar{\phi}}$$

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

$$\delta s(k) = \mathcal{A}_e \left(k/k_p \right)^n \exp \left(- (k/k_c)^p \right)$$



$[k_8 = 0.125h \text{ Mpc}^{-1}, \mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0]$

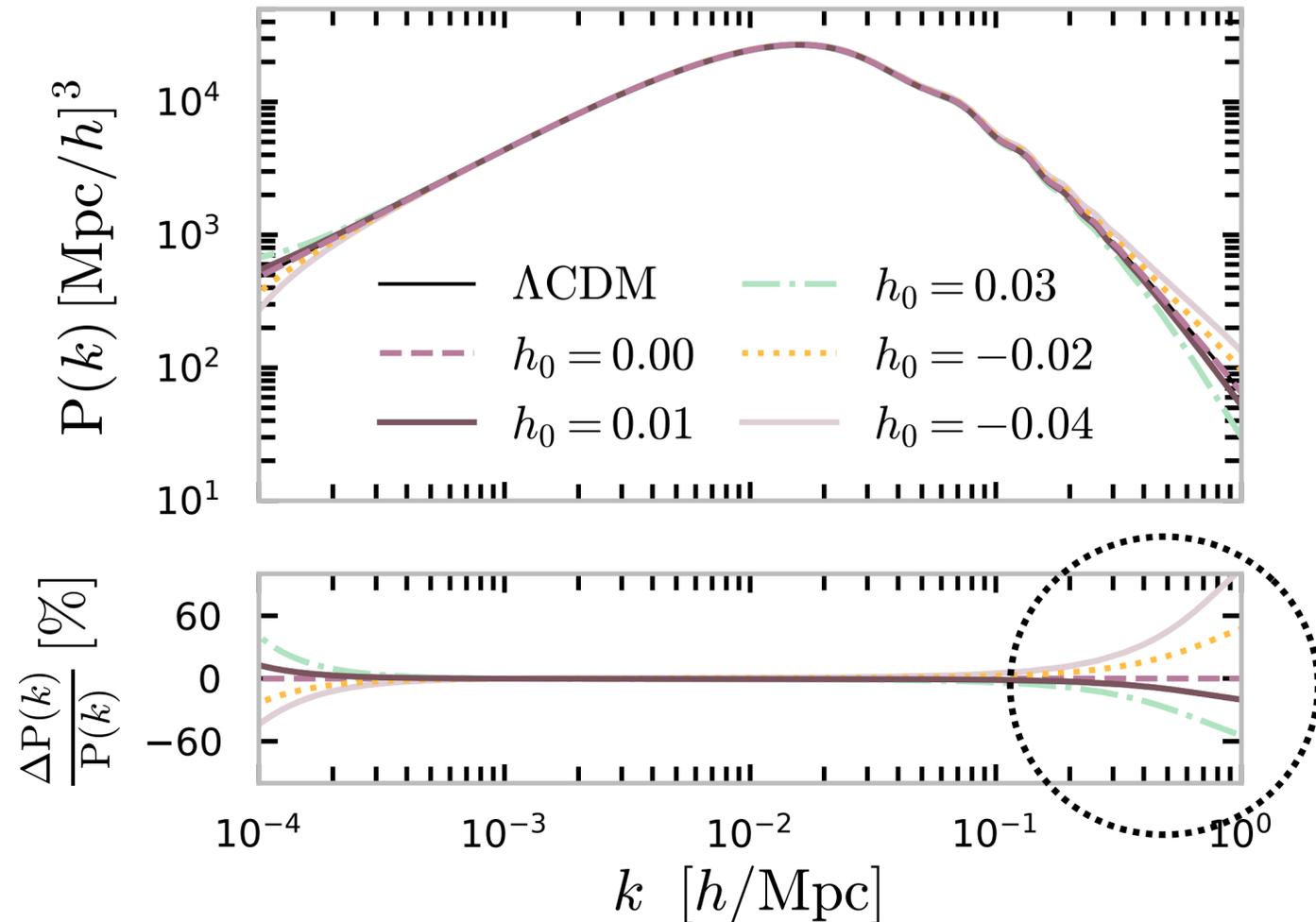


Observational signatures

Suppression/enhancement of matter power spectrum on small scales and large scales for $n=0$

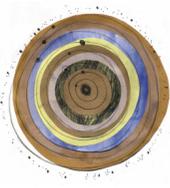
$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

$$\delta s(k) = \mathcal{A}_e \left(k/k_p \right)^n \exp \left(- (k/k_c)^p \right)$$



$$\theta'_c + \theta_c \mathcal{H} - k^2 \Psi = -k^2 Q_s \delta s(k), \quad Q_s := \frac{\bar{g}_{,s} - h_0 \hat{V}_{,\phi}}{\bar{\rho}_c}$$

$$[k_8 = 0.125h \text{ Mpc}^{-1}, \mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0]$$

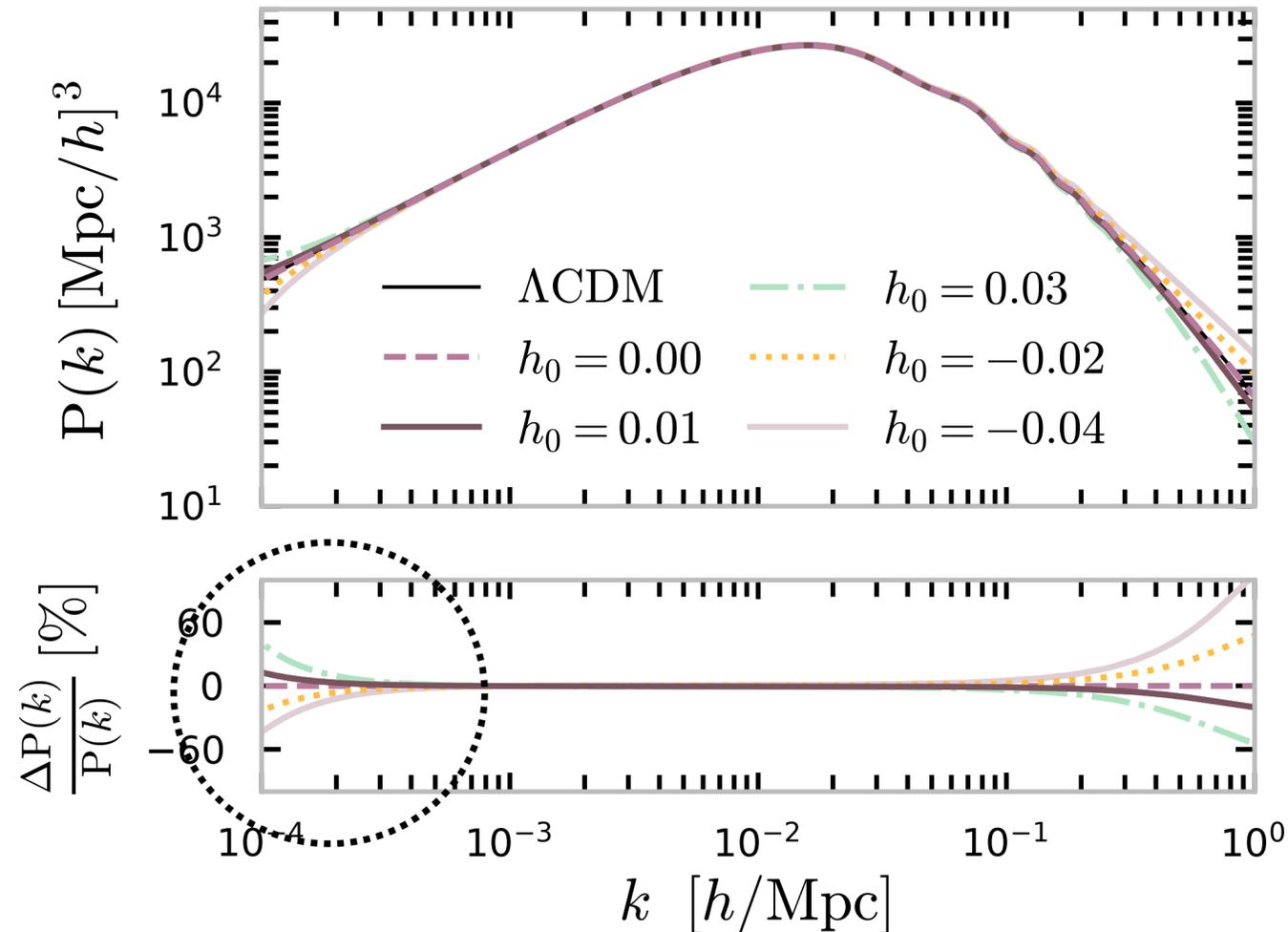


Observational signatures

Suppression/enhancement of matter power spectrum on small scales and large scales for $n=0$

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

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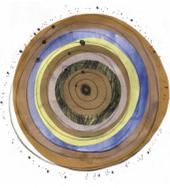


$$\theta'_c + \theta_c \mathcal{H} - k^2 \Psi = -k^2 Q_s \delta s(k), \quad Q_s := \frac{\bar{g}_{,s} - h_0 \hat{V}_{,\bar{\phi}}}{\bar{\rho}_c}$$

$$\Psi \propto -\frac{a^2}{k^2} (\delta \rho_m + \delta \rho_{DE})$$

$$\Psi' + \mathcal{H} \Psi \propto \frac{a^2}{k^2} (\theta_m + \theta_{DE})$$

$$\mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0$$

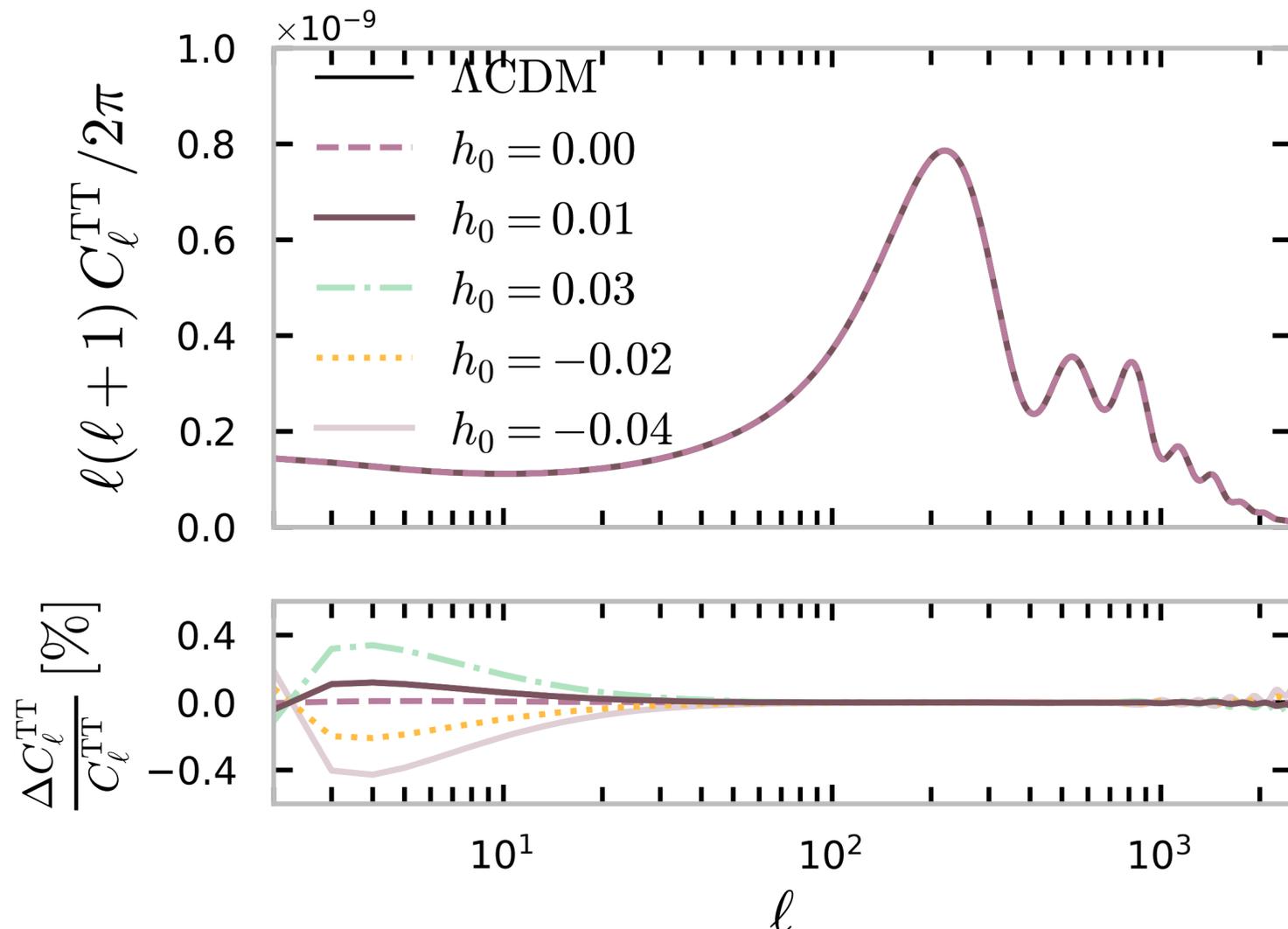


Observational signatures

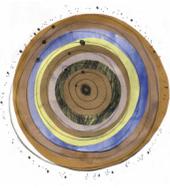
Tiny ISW changes the CMB temperature spectrum \mathcal{C}_ℓ^{TT} at low- ℓ (background unchanged and late time effects only)

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

$$\delta s(k) = \mathcal{A}_e \left(k/k_p \right)^n \exp \left(- (k/k_c)^p \right)$$



$$\mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0$$

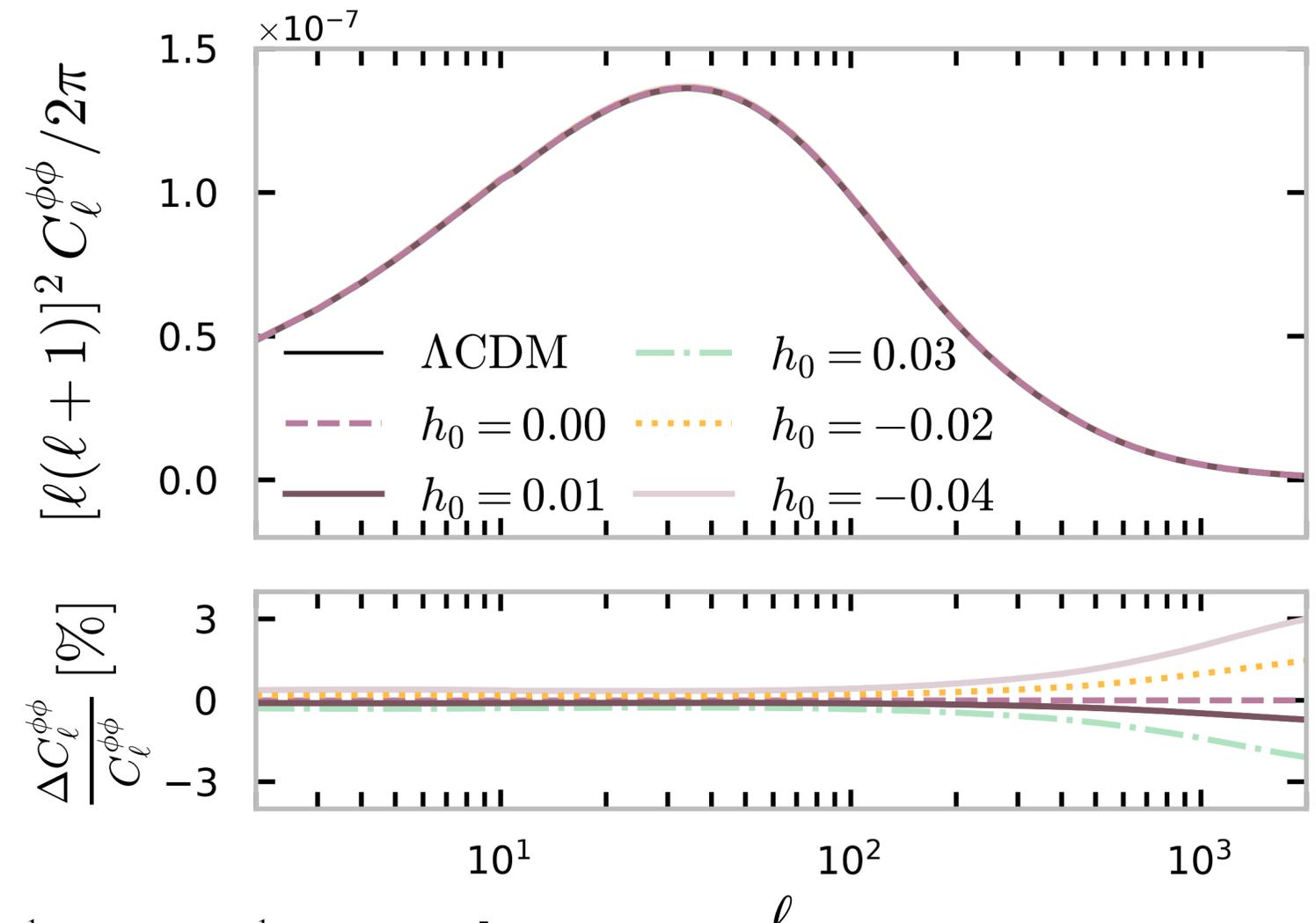
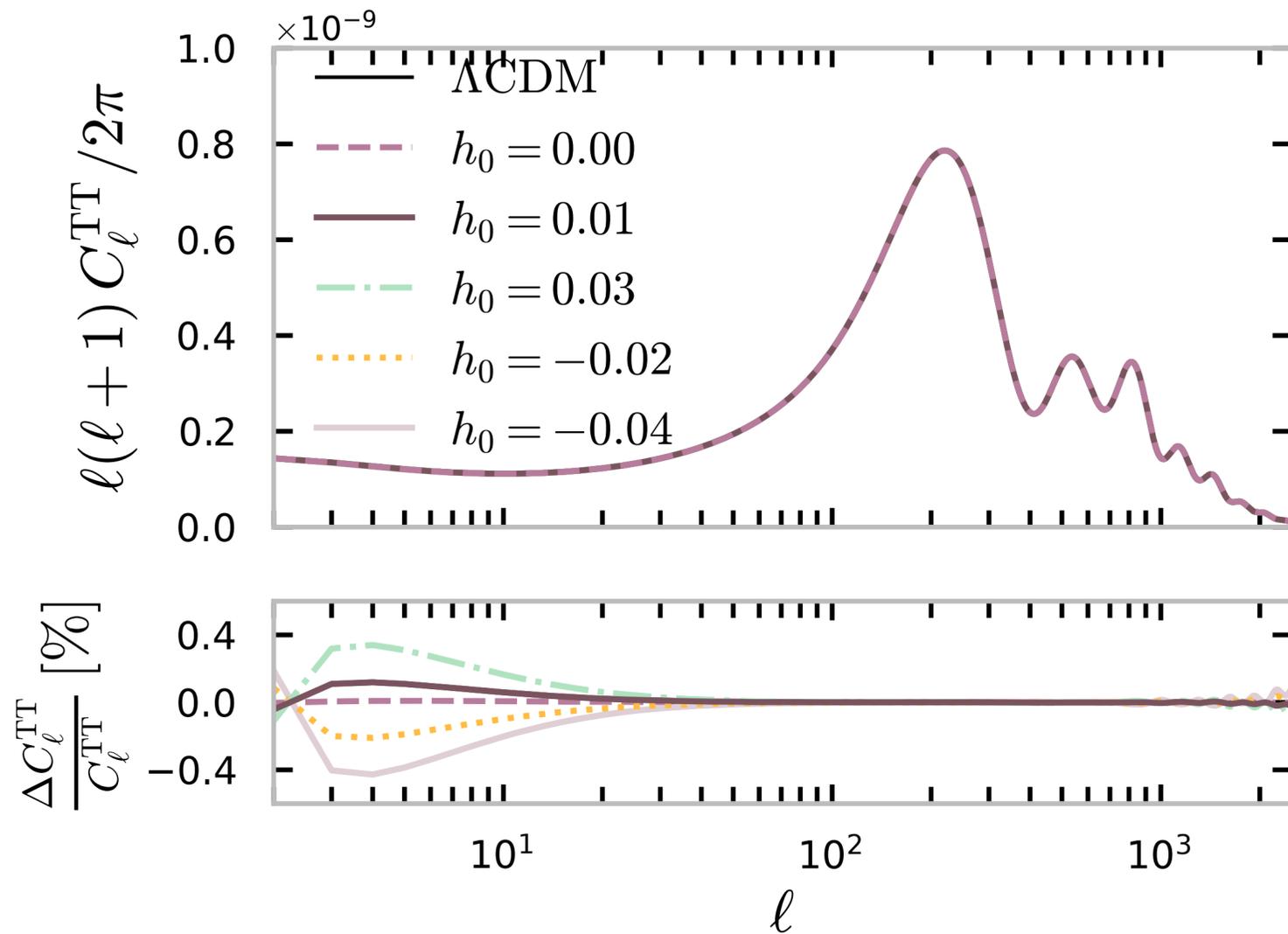


Observational signatures

Changes in the CMB lensing spectrum $\mathcal{C}_\ell^{\phi\phi}$ due to the cumulative changes in the growth and the metric perturbations at broad z

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

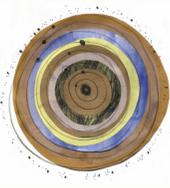
$$\delta s(k) = \mathcal{A}_e \left(k/k_p \right)^n \exp \left(- (k/k_c)^p \right)$$



$$\mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0$$

Summary and Conclusions

A photograph of a geological rock outcrop showing distinct horizontal sedimentary layers. The layers are colored in various shades, including blue, yellow, and pink, set against a dark background. The text "Summary and Conclusions" is overlaid in white on the upper part of the image.



Conclusions

- New interacting scenarios from coupling the dark matter intrinsic entropy to a dark energy scalar field
 - Unchanged background expansion
 - No energy transfer (pure-momentum only)
 - Scale-dependent suppression/enhancement of matter clustering at late-times due to k -dependent contributions of $\delta_s(k)$
-
- Key observational signatures: scale-dependent clustering
 - Observational constraints on allowed couplings including mixing with background energy/momentum exchange - cosmic tensions
 - Future work: explore different primordial power spectrum for $\delta_s(k)$ and early dark energy/neutrino scenarios

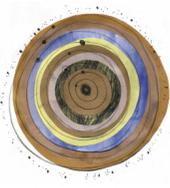


Thank you for your attention!

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CNRS & Université de Montpellier
elsa.teixeira@umontpellier.fr

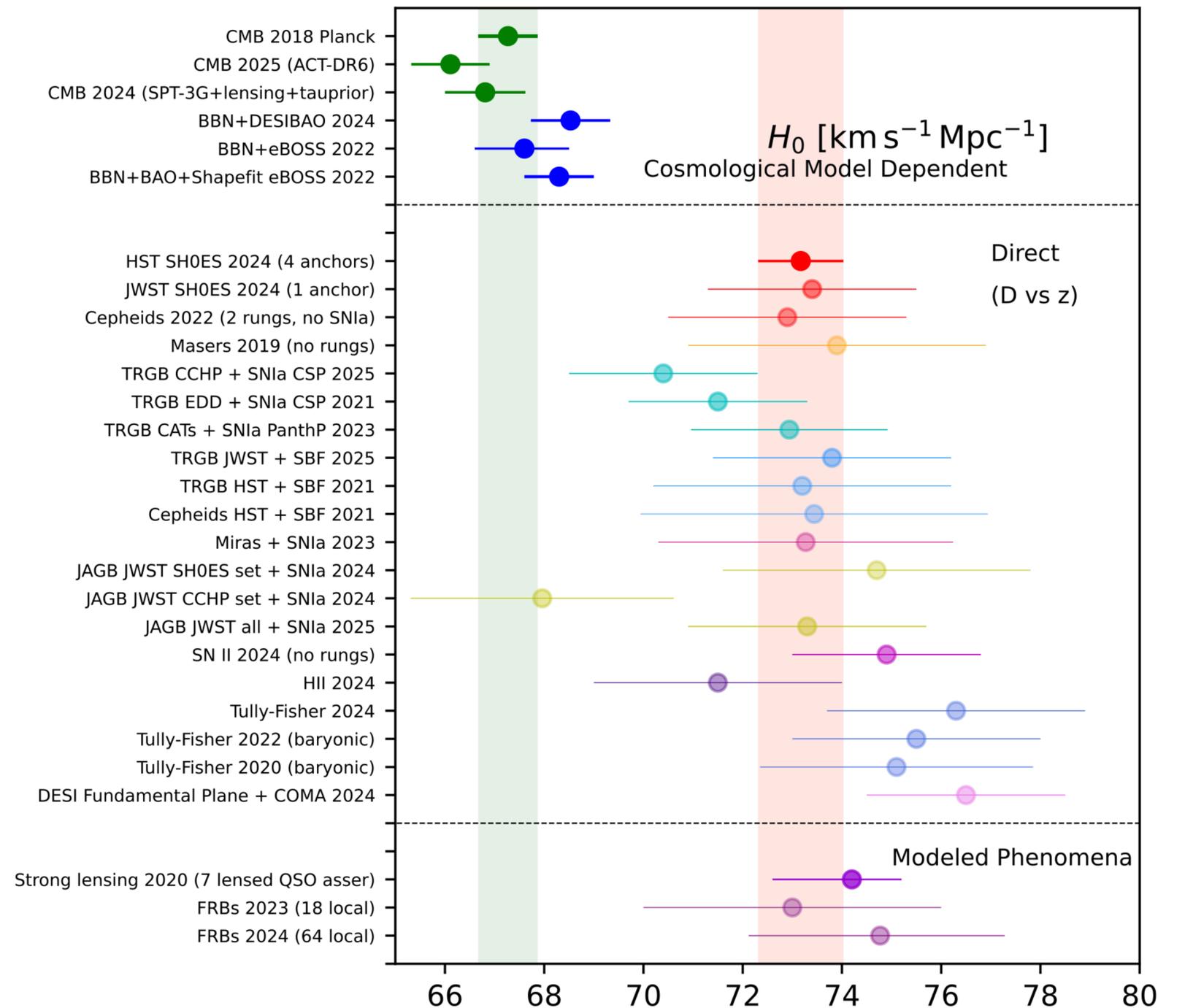
Illustrations: Inês Viegas Oliveira
(ivoliveira.com)

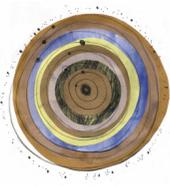


The Hubble Tension

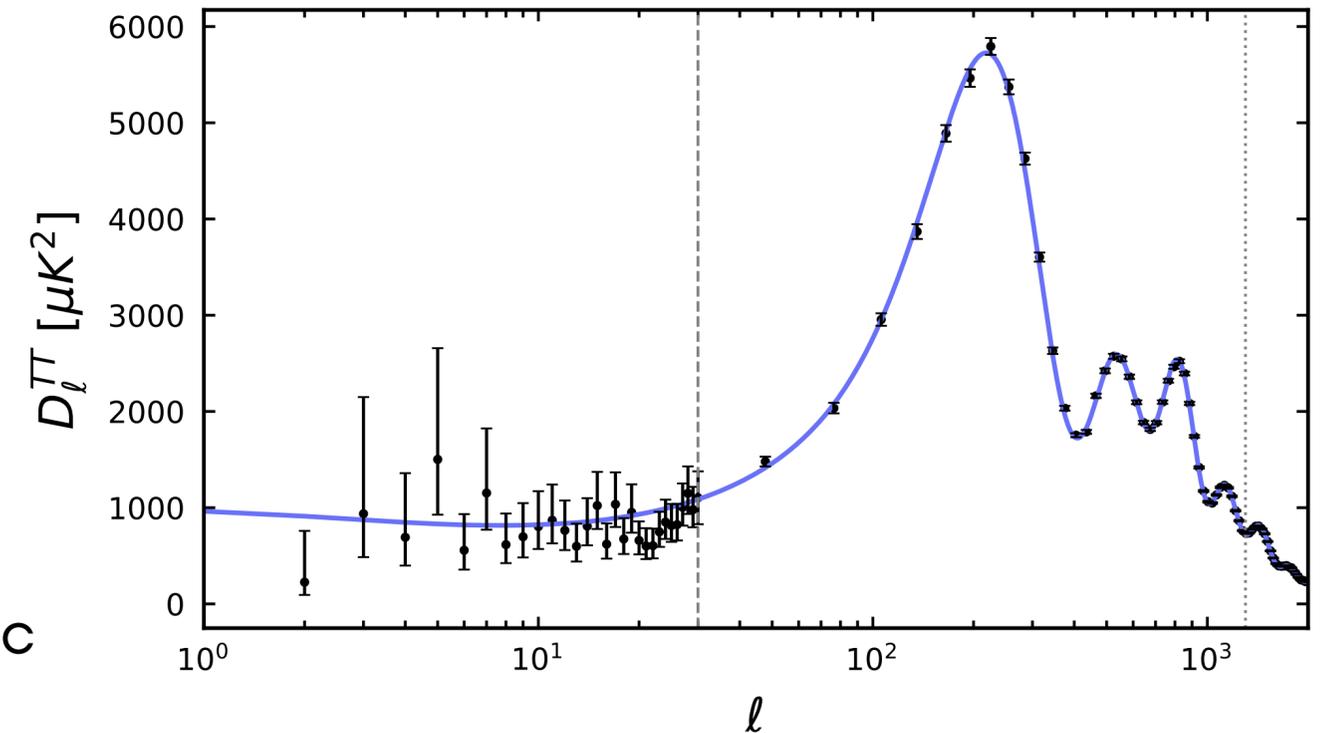
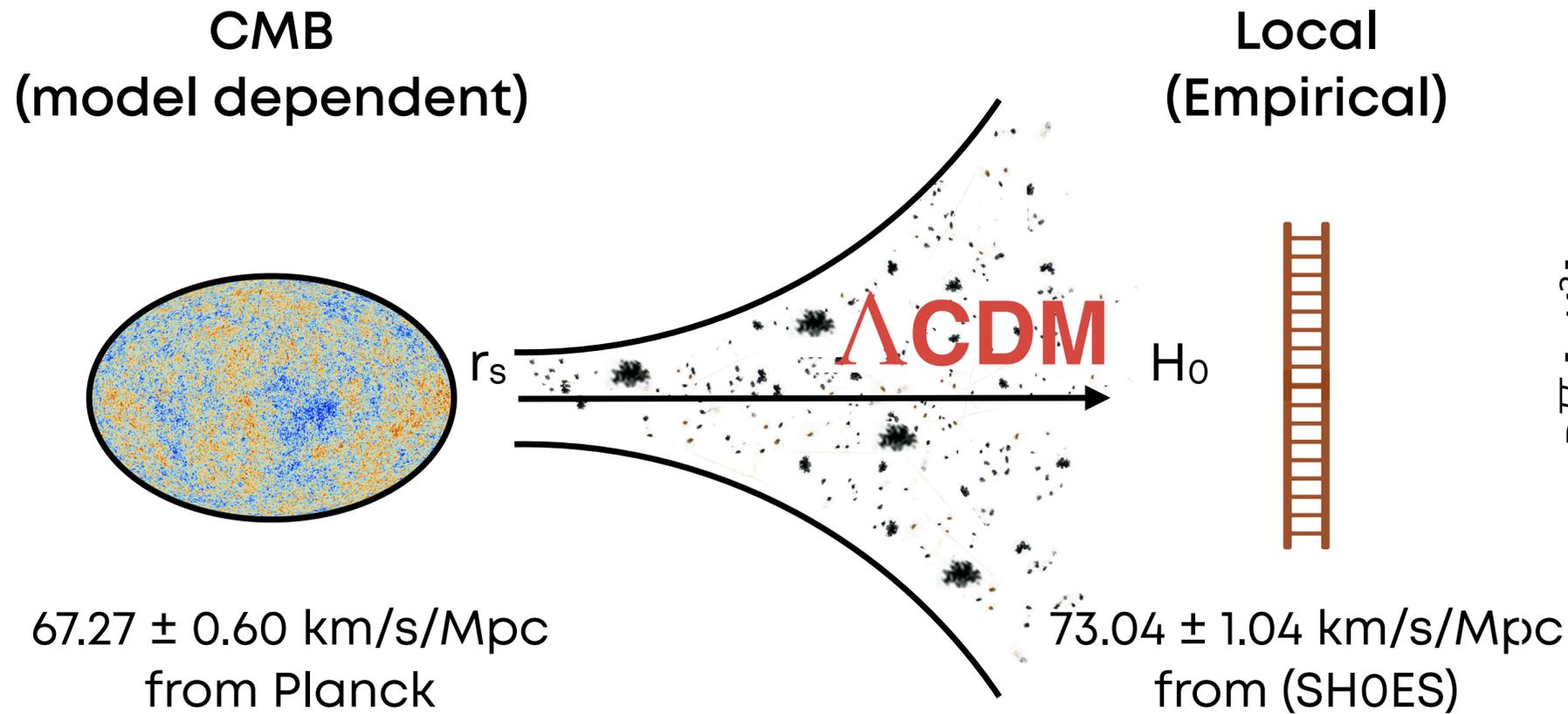
Unreconcilable values for H_0 from the CMB and from direct local distance ladder measurements

- $\sim 5\sigma$ tension between Planck 2018 and SH₀ES:
 - ▶ **CMB (Planck):** $H_0 = 67.27 \pm 0.60$ km/s/Mpc
 - ▶ **SNe (R22):** $H_0 = 73.04 \pm 1.04$ km/s/Mpc
- The CMB data assumes the Λ CDM model
- **DESI BAO (+BBN+CMB):** $H_0 = 68.45 \pm 0.47$ km/s/Mpc [DESI Collaboration DR2 2025: arXiv:2503.14738]
- Compilation of early vs late time data that disagree
- Could signal differences in the expansion history (nature of the dark sector)



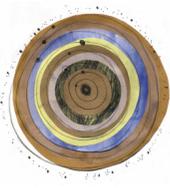


Cosmological Tensions

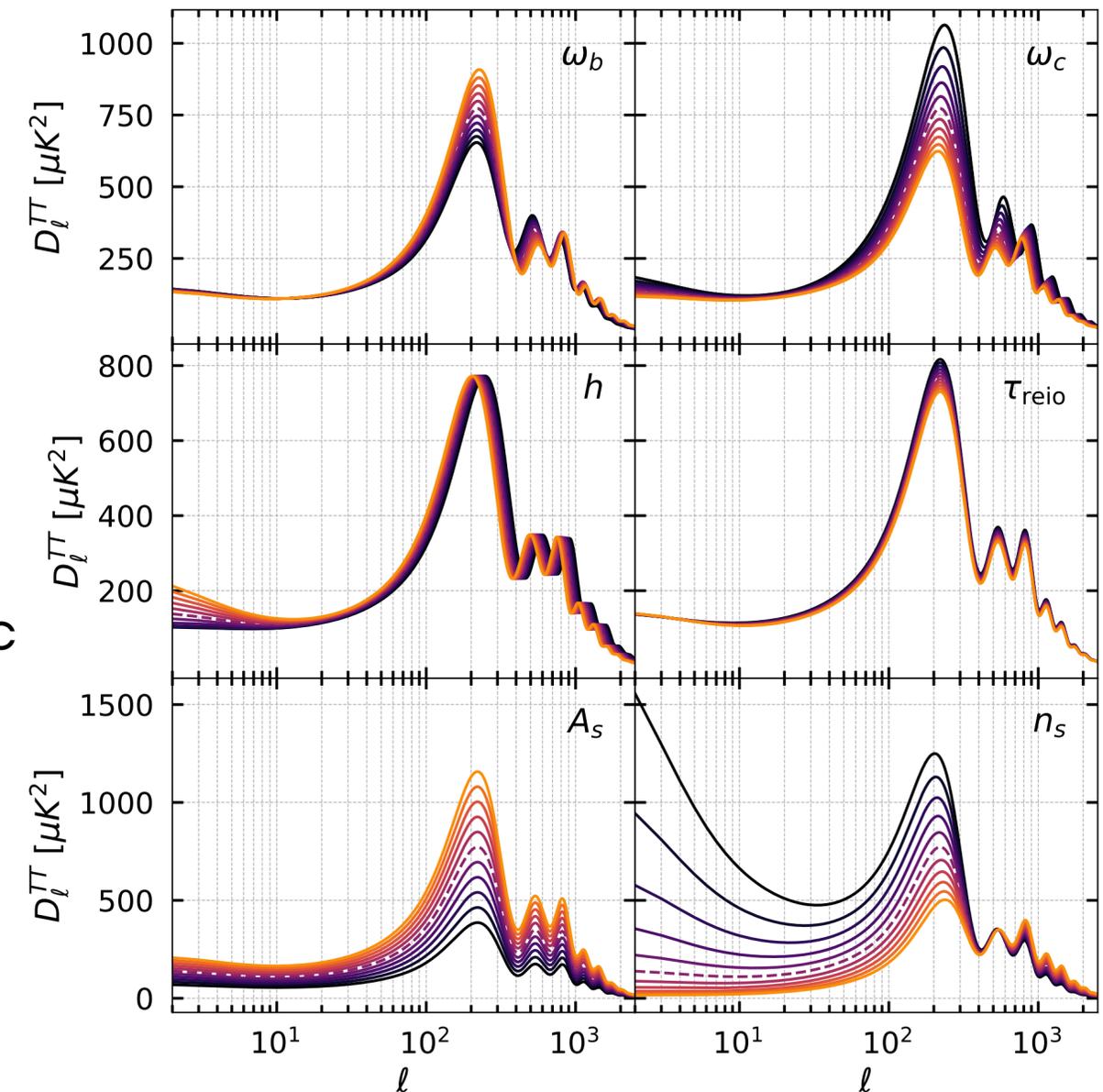
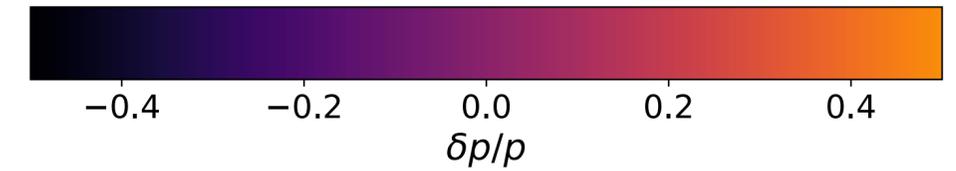
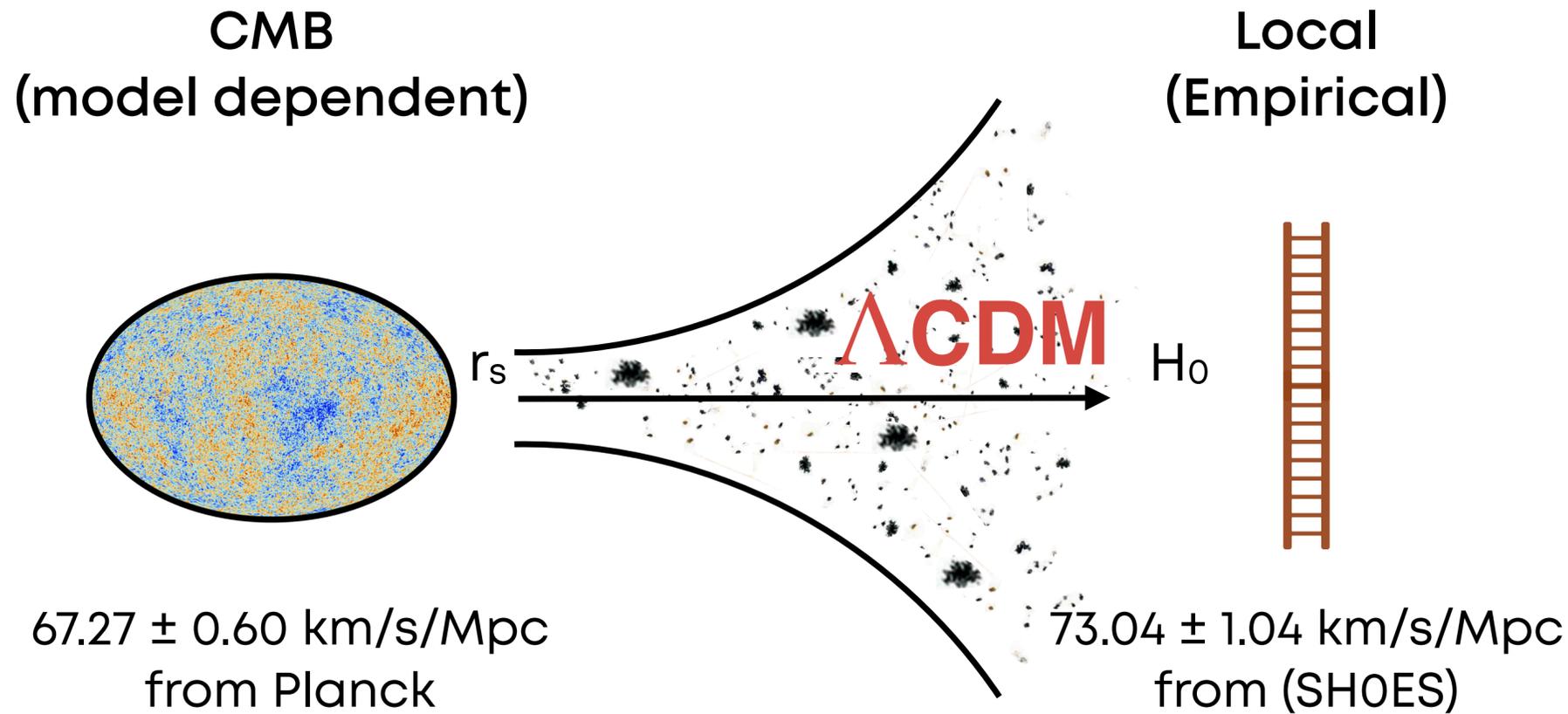


Missing Ingredients or New Physics?

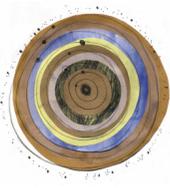
[Aghanim et al.: Astron.Astrophys. 641 (2020) A6]



Cosmological Tensions



Missing Ingredients or New Physics?



The S_8 Tension

Discrepancy between CMB data and lensing

surveys on combined quantity $S_8 = \sigma_8 \sqrt{\Omega_m / 0.3}$

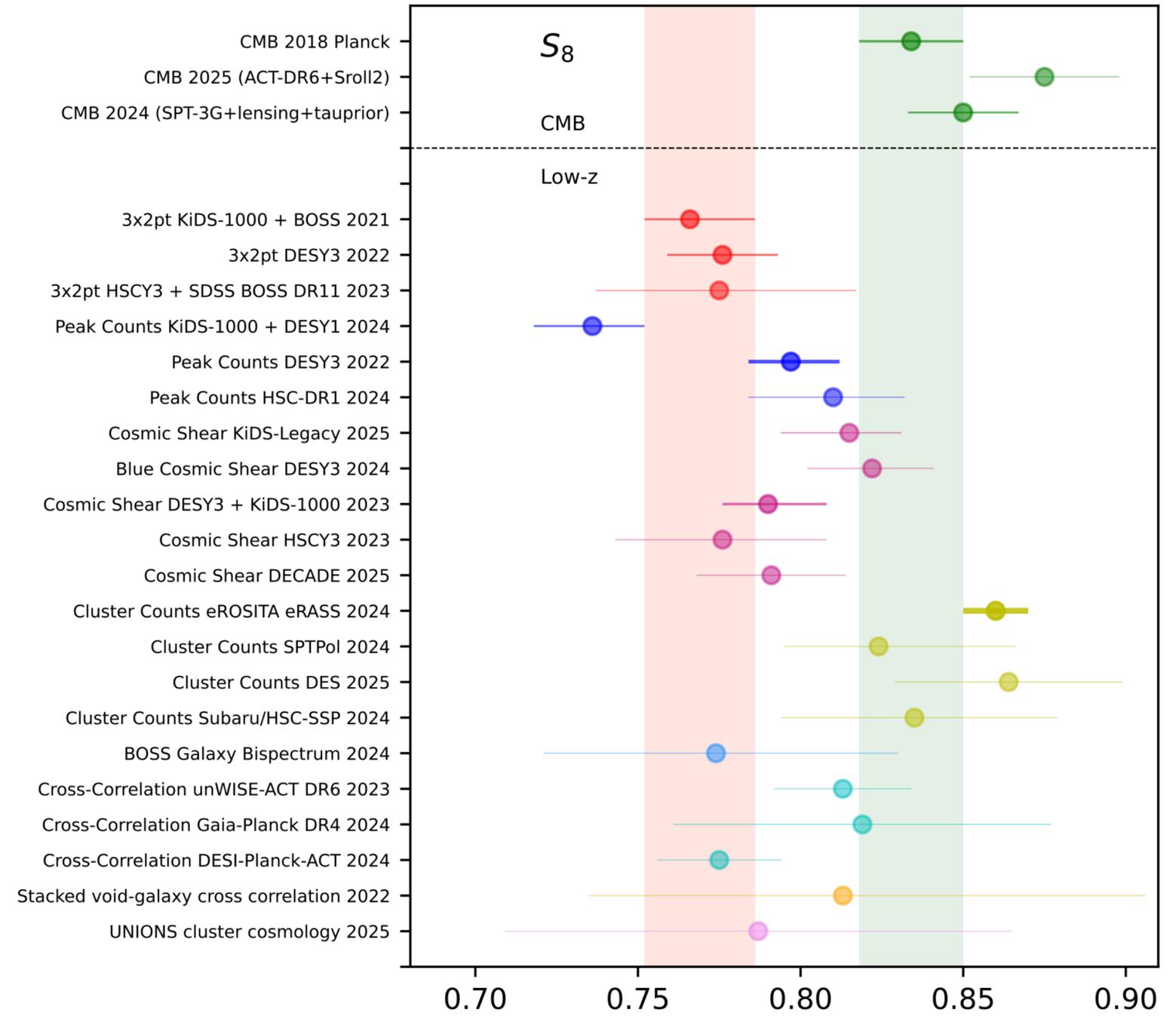
- ~ 3σ tension between Planck 2018 CMB data and KiDS-1000 combination of Cosmic Shear and Galaxy Clustering:

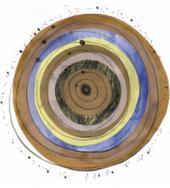
- ▶ **CMB (Planck 2018):** $S_8 = 0.832 \pm 0.013$
- ▶ **Cosmic Shear (DES-Y3):** $S_8 = 0.759^{+0.025}_{-0.023}$

- eRosita (eRASS1):** $S_8 = 0.86 \pm 0.01$ [Ghirardini et al. 2024]

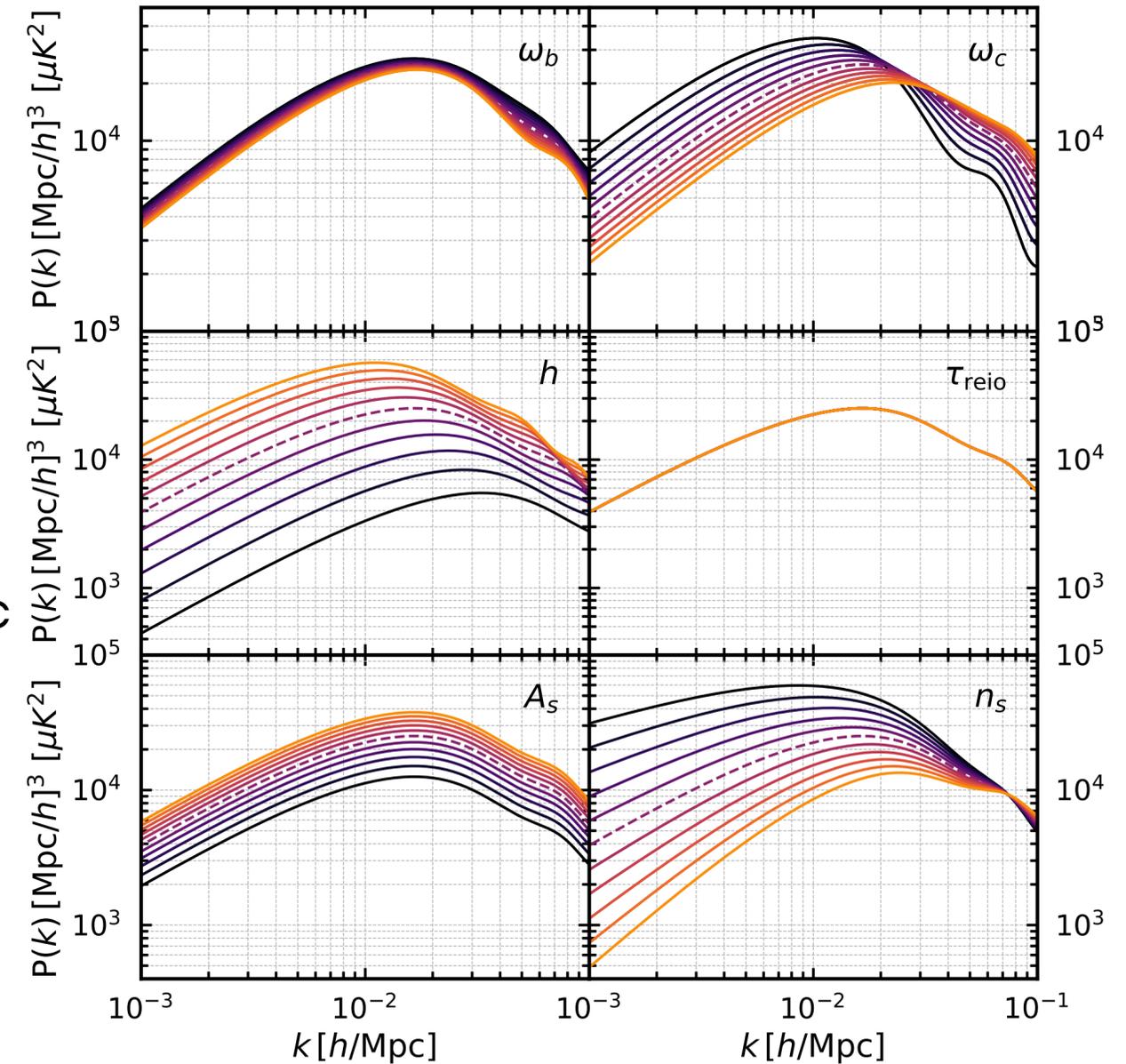
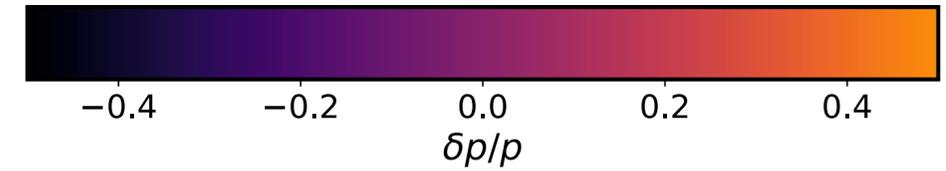
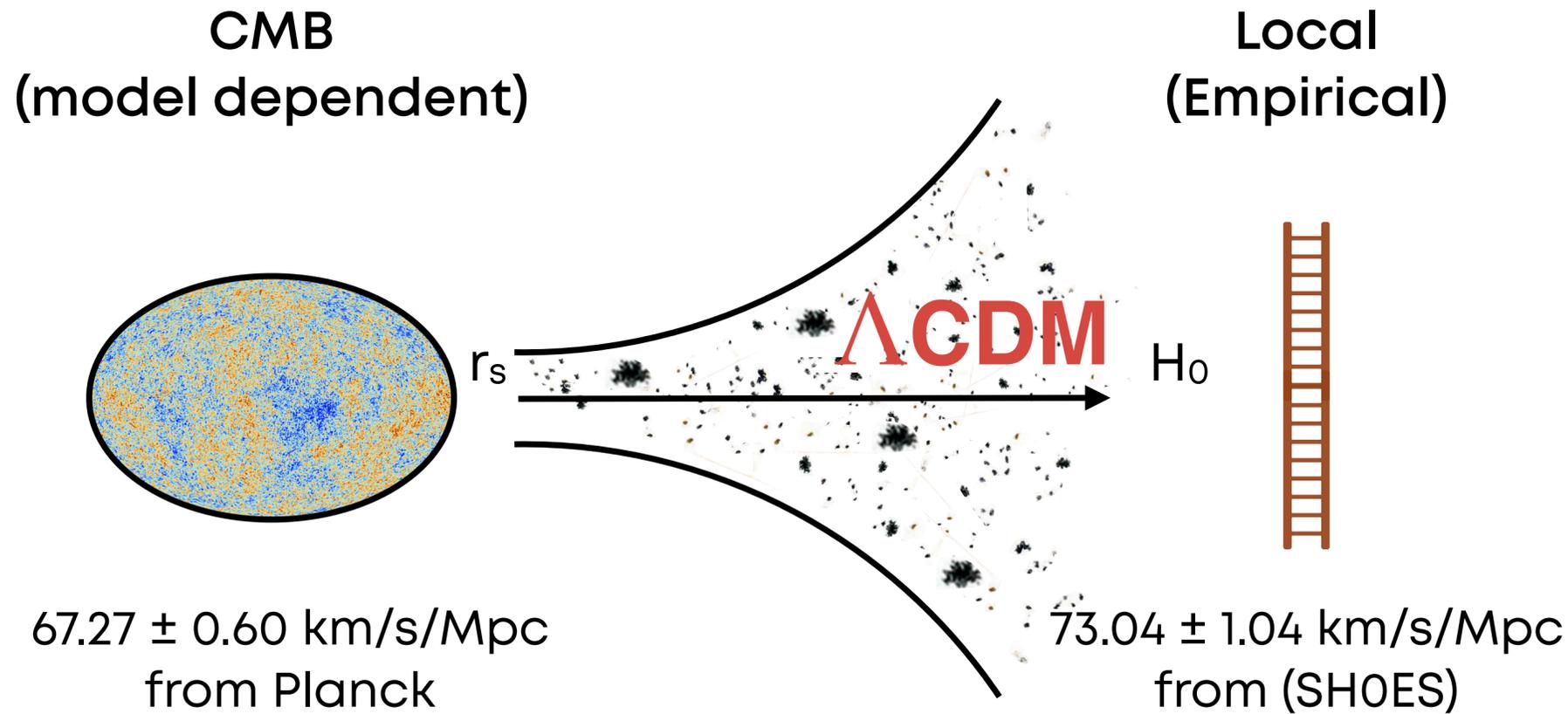
- Kids-Legacy** ($S_8 = 0.815^{+0.016}_{-0.021}$) find possible resolution with Planck but not for the other measurements (improved redshift distribution estimation and calibration, as well as new survey area and improved image reduction)

- Could signal changes in clustering of matter (nature of CDM)

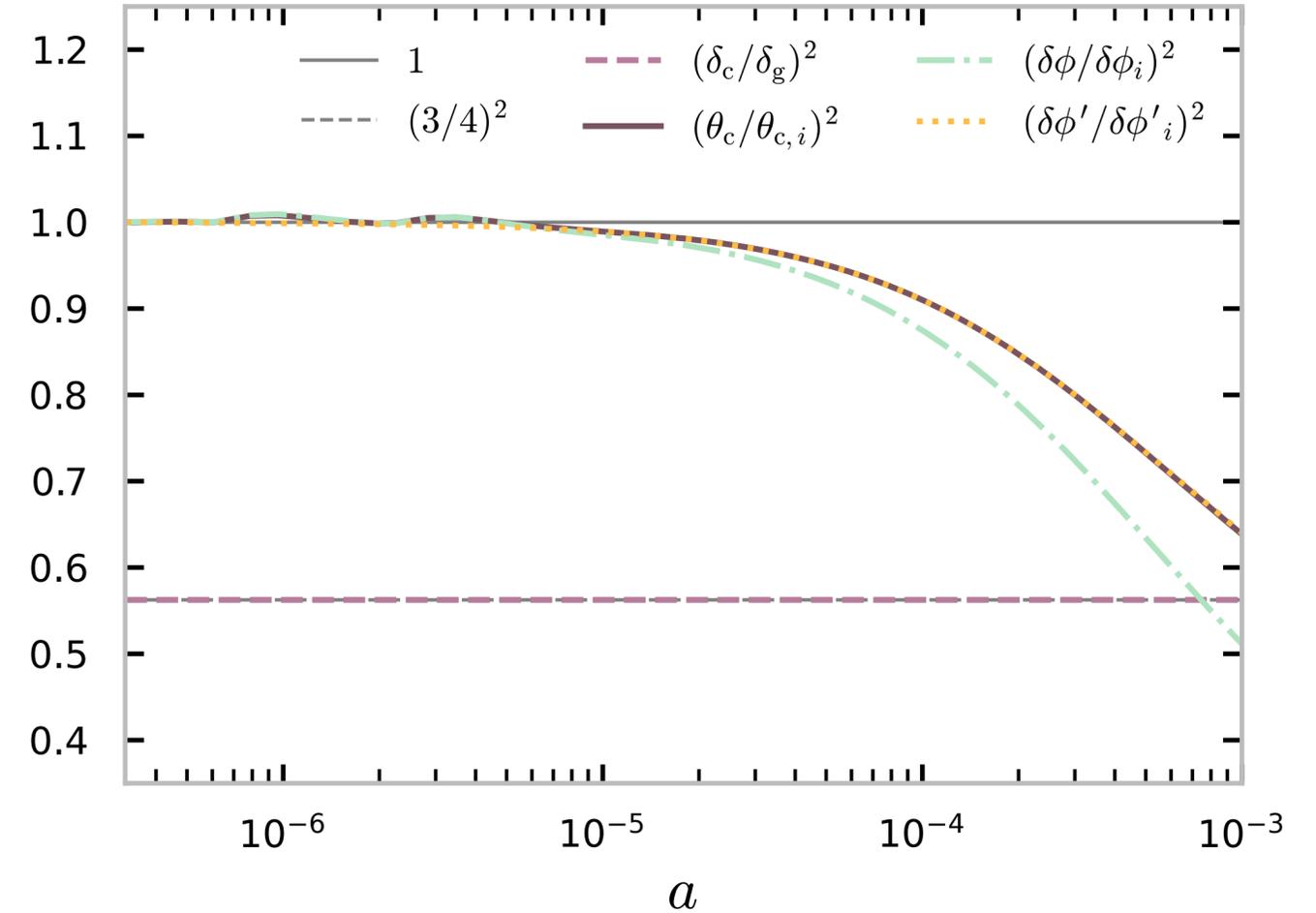
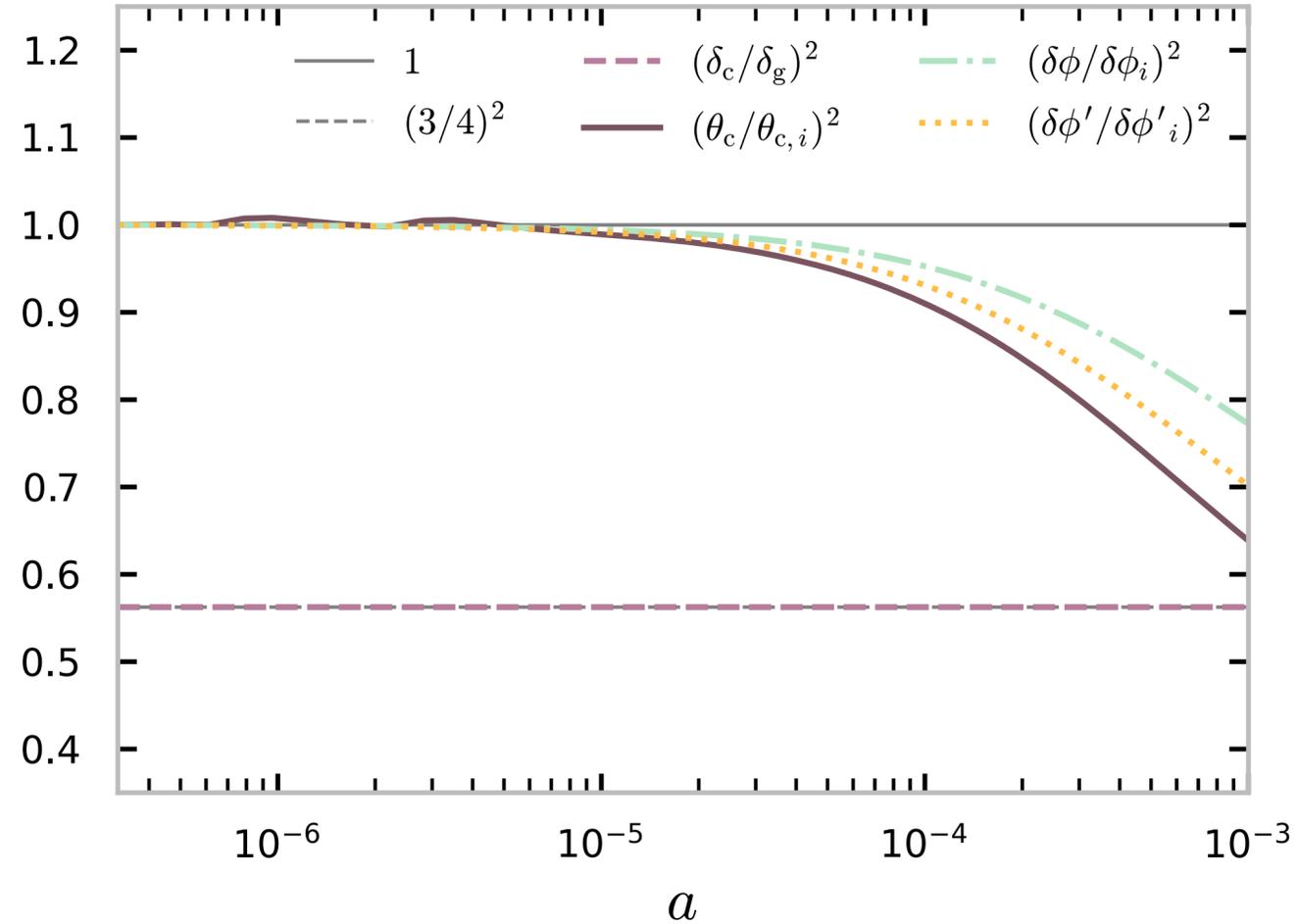
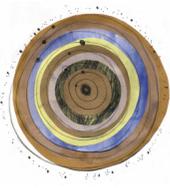




Cosmological Tensions



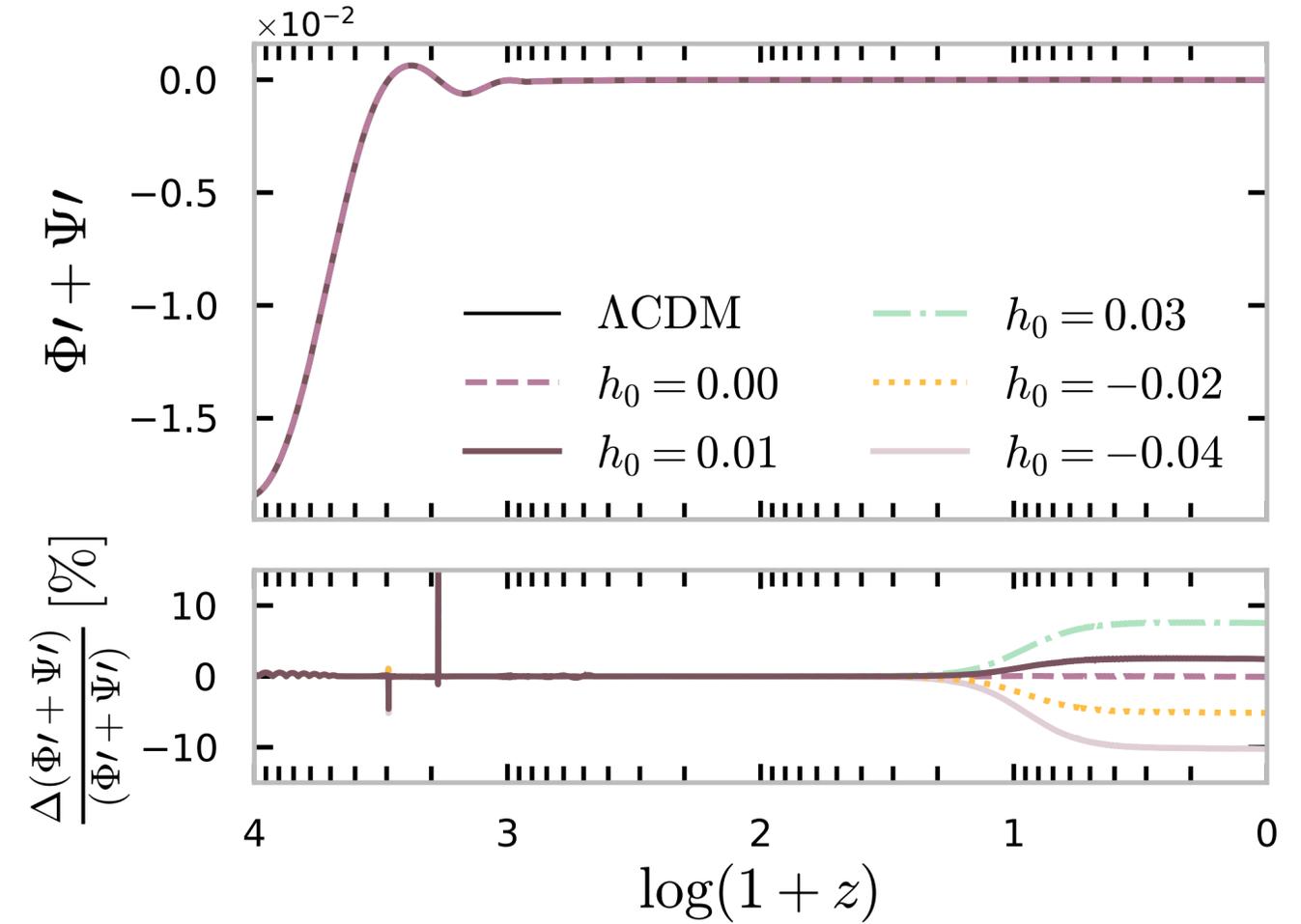
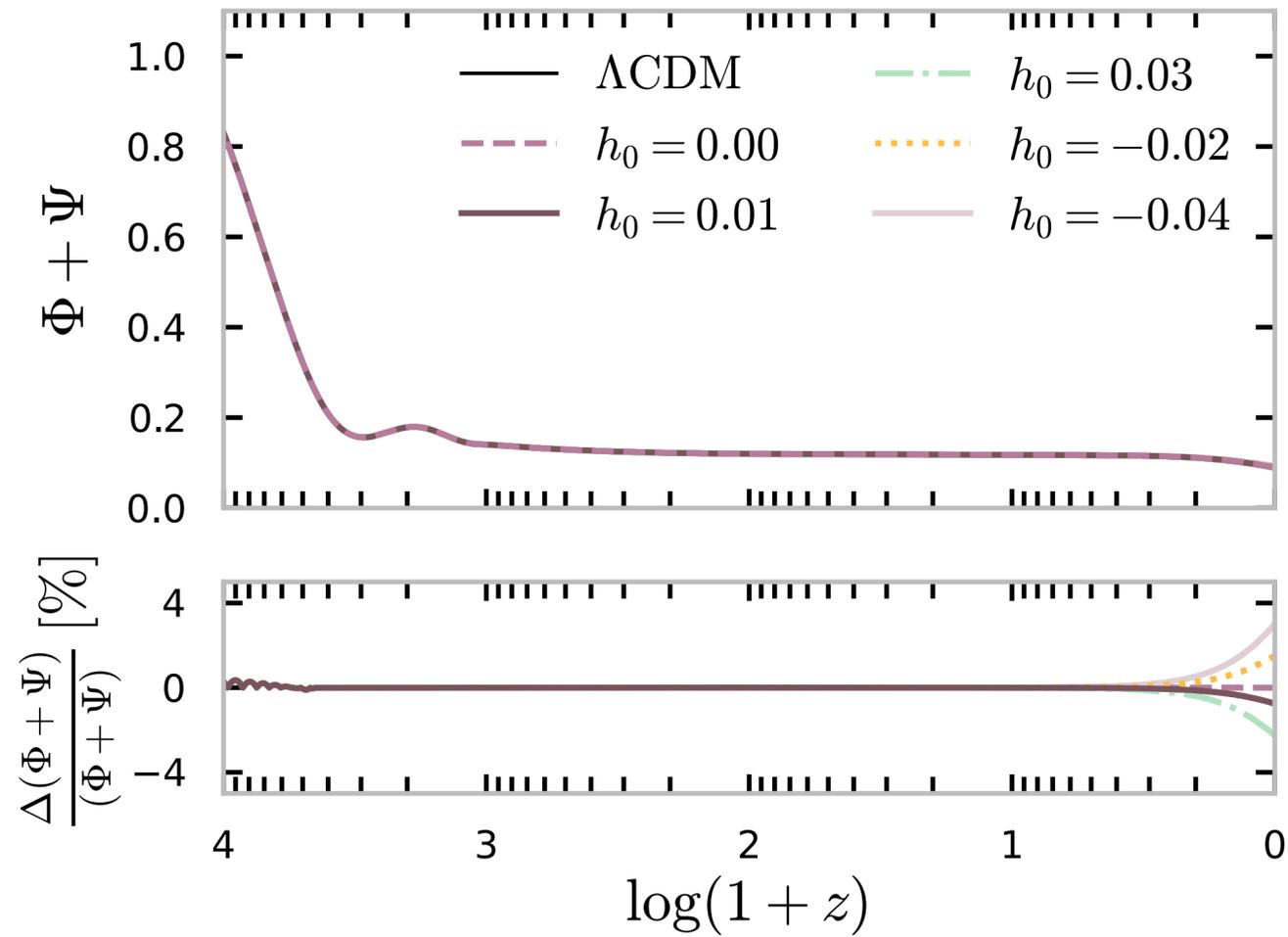
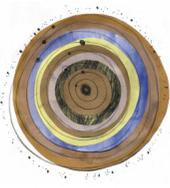
Missing Ingredients or New Physics?



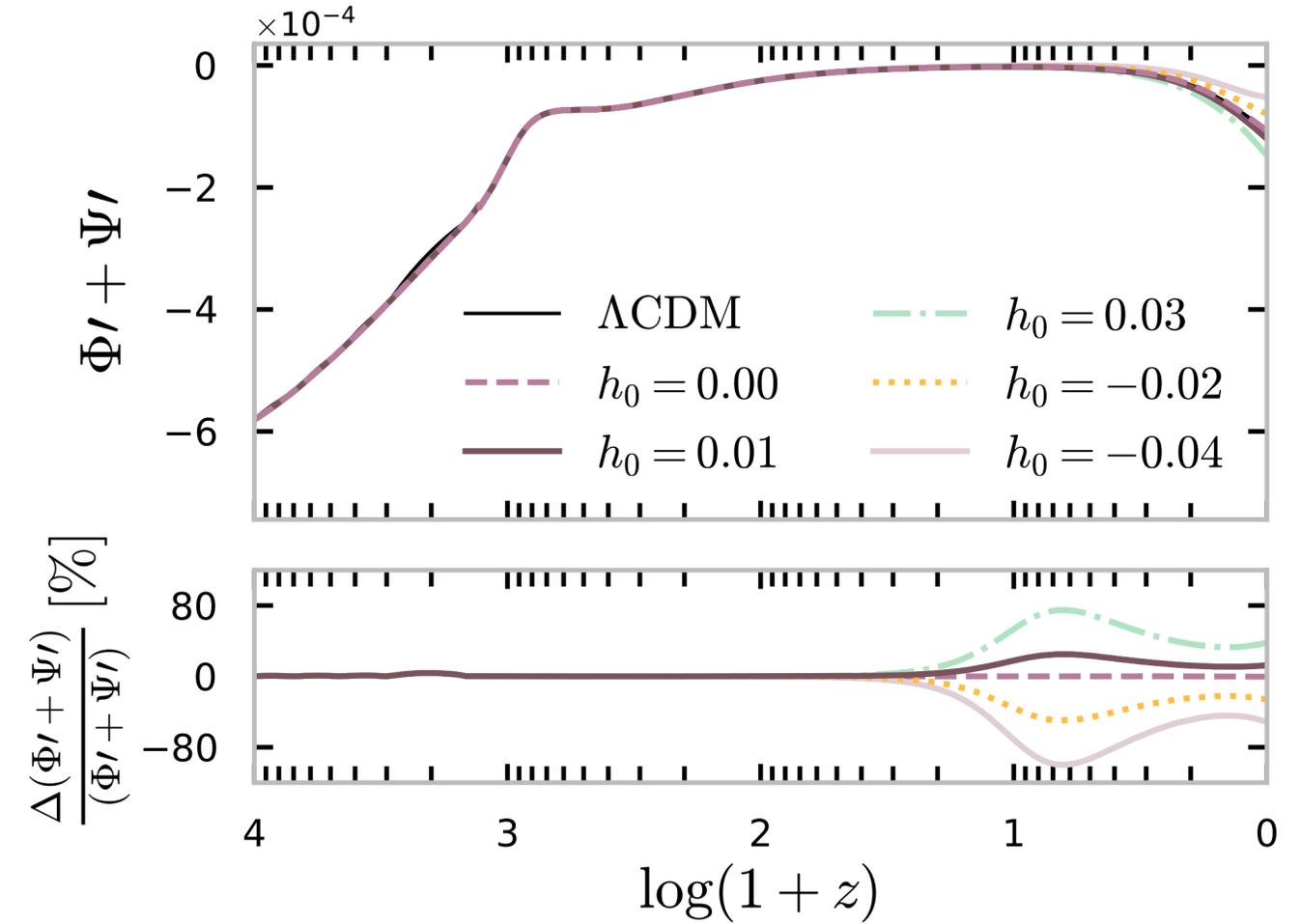
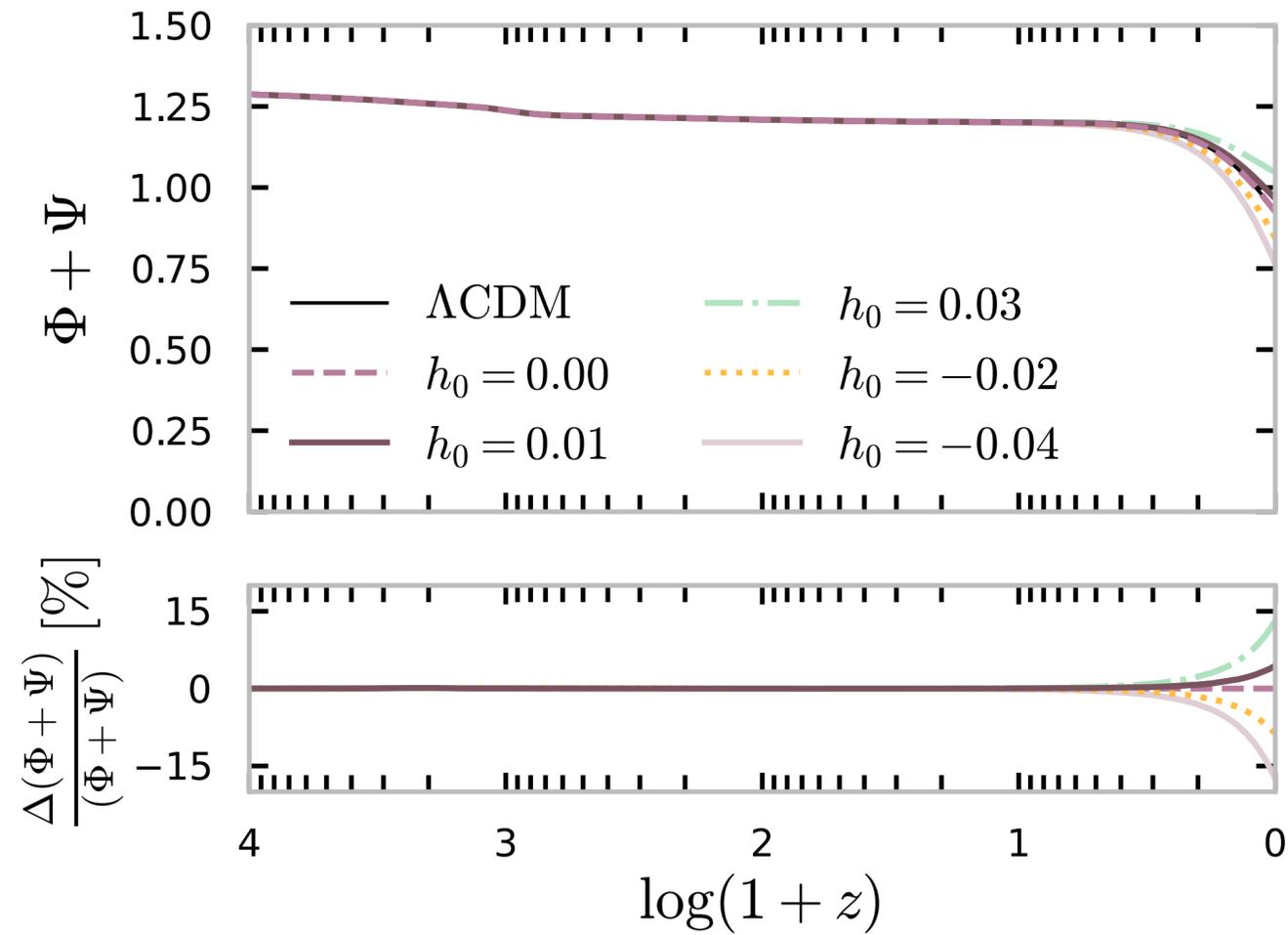
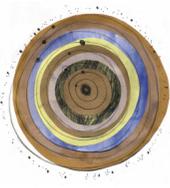
$$\delta_c(\tau) = -\frac{C}{2}(k\tau)^2, \quad \theta_c(\tau) = -\frac{D\delta s}{5k^2}(k\tau)^4$$

$$\delta\phi(\tau) = -\frac{h_0}{6}\delta s (k\tau)^2 - \frac{\alpha^2 g_{s\phi}\delta s}{20k^4}(k\tau)^4 + \mathcal{O}((k\tau)^5),$$

$$\delta\phi'(\tau) = -\frac{h_0}{3}\delta s k (k\tau) - \frac{\alpha^2 g_{s\phi}\delta s}{5k^3}(k\tau)^3 + \mathcal{O}((k\tau)^4)$$



$[k_8 = 0.125h \text{ Mpc}^{-1}, \mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0]$



$[k = 10^{-4} h \text{ Mpc}^{-1}, \mathcal{A}_e = 1, k_p = 0.05 \text{ Mpc}^{-1}, k_c = 1 \text{ Mpc}^{-1}, p = 2, n = 0]$