

# Will we ever see any GRBs from the very first stars?

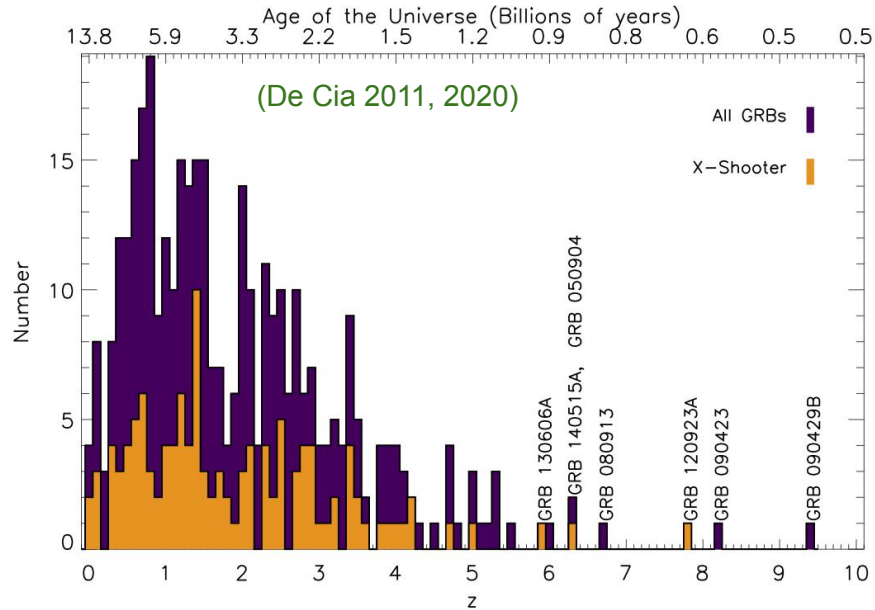
Ramandeep Gill  
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Paper

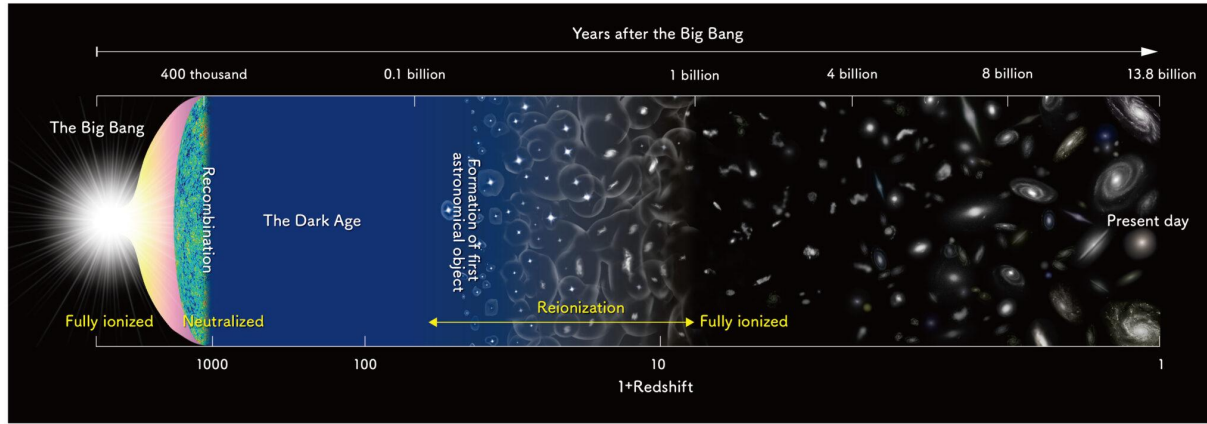
Morales-Rivera, G., R.G., Jane, Beniamini, Granot  
(2025, MNRAS, in press)

# GRBs as cosmological probes



- The intrinsic high luminosity of GRBs allows them to be observed to very high redshifts.
- The prompt gamma-ray emission can be detected up to  $z \sim 100$  (Lamb & Reichart 2000) and afterglows to  $z \sim 30$  (IR, X-rays, Radio) (Ciardi & Loeb 2000; Ioka 2003; Inoue 2004)
- This makes GRBs useful probes of the early Universe, that can provide constraints on:
  - Star formation history
  - Stellar evolution
  - Extragalactic background light
  - Lorentz invariance violation

# GRBs as probes of the first stars

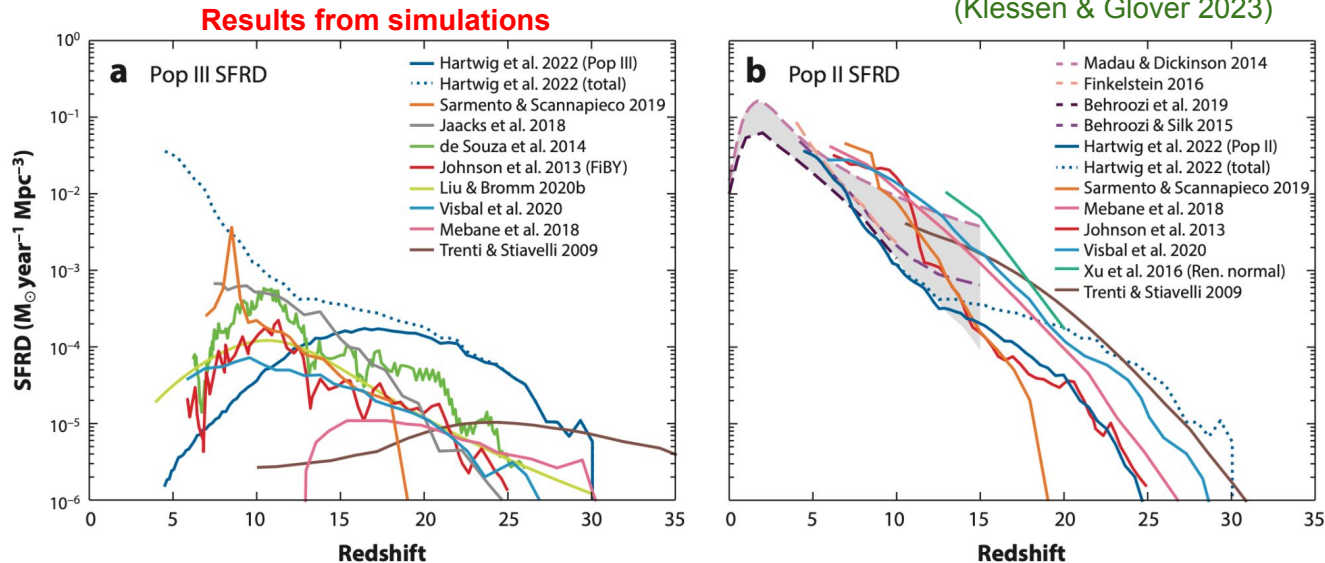


(NAOJ/NAO)

- The very first stars (Pop III) are very challenging to observe directly, even with JWST!
- In the collapsar channel, long GRBs originate in the core collapse of massive stars that only live for  $\sim 10^{6.5} - 10^7$  yr
- Since stars tend to be more massive in the past, some are expected to produce bright GRBs

# Pop III stars

- Pop III stars formed from the **metal-free** primordial gas comprising H + He.
- The  $\Lambda$ CDM model predicts their **formation at  $z \sim 30$** , where they **dominated the SFRD at  $z \sim 15 - 20$** , with peak around  $z \sim 10$ .



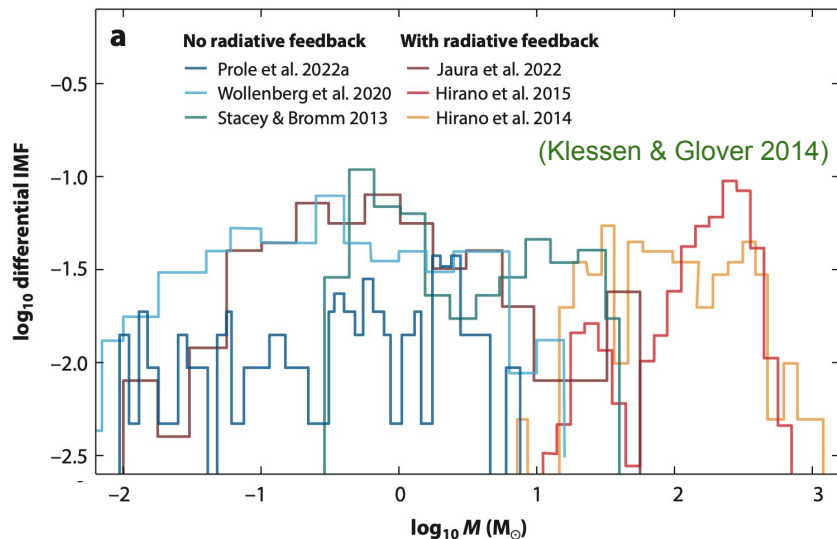
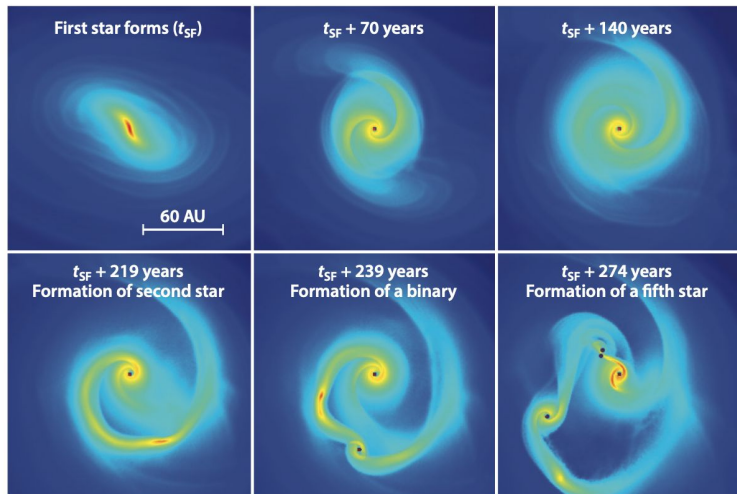
# Pop III stars: Initial mass function

- The IMF of Pop III is expected to be **top-heavy**, with  $\alpha \sim 1.13$  (Stacy+2016; Wollenberg+2020), and masses in the range  $1 \lesssim M/M_\odot \lesssim 10^3$  (Hirano+2014; Hosokawa+2016)

$$\frac{dN}{d \ln M} \propto M^{1-\alpha}$$

- The present-day IMF is **bottom-heavy** with  $\alpha = 2.35$  (Salpeter 1955) with average mass  $\langle M \rangle = 0.5 M_\odot$

(Clark+2011)

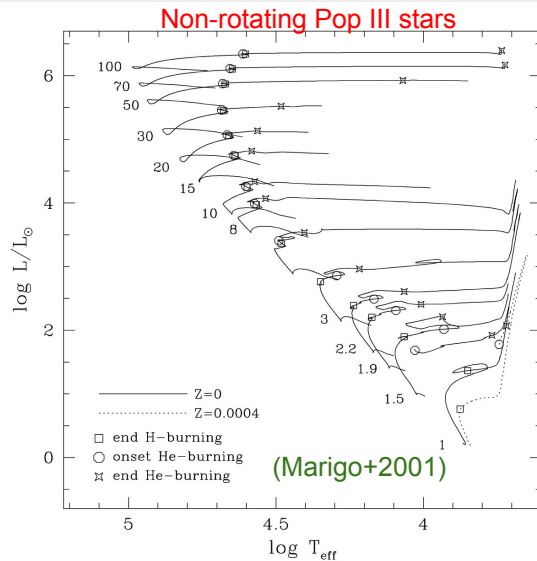
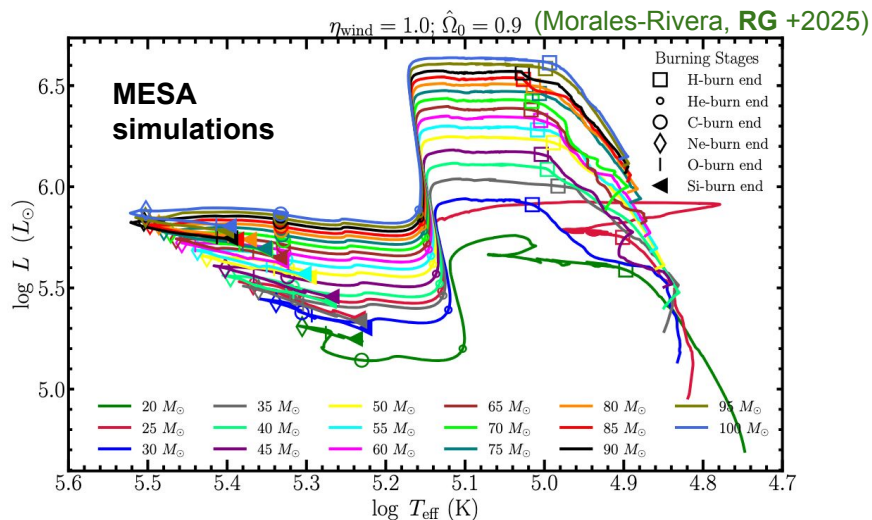


# Stellar evolution of rapidly rotating stars

- Pop III stars were also expected to be rapid rotators, with surface equatorial angular velocities approaching the critical limit:

$$\frac{\Omega_0}{\Omega_{\text{crit}}} \lesssim 1 \quad \Omega_{\text{crit}} = \left(1 - \frac{L}{L_{\text{Edd}}}\right)^{1/2} \Omega_k \quad \Omega_k = \sqrt{\frac{GM}{R_\star^3}}$$

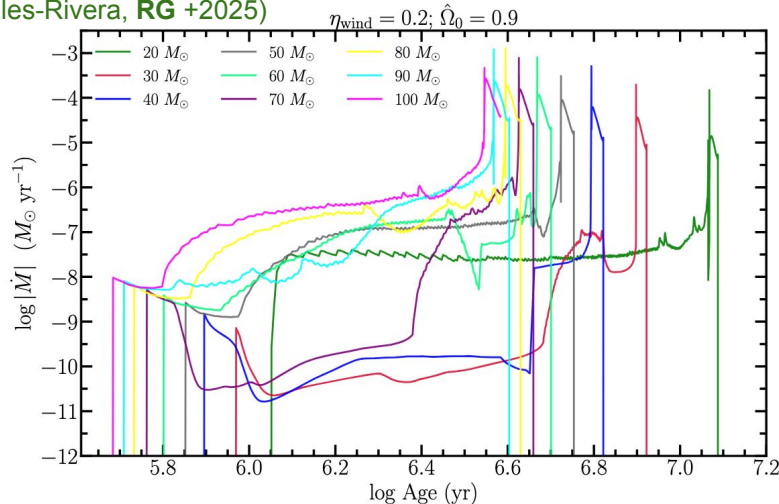
- Rapid rotators undergo a **chemically homogeneous evolution**, that leads to a bluward evolution on the HR diagram resulting in hotter, luminous stars with more massive cores.



# Mass loss

- Pop III stars can also **lose mass both mechanically and to line-driven winds** even though they initially don't have any metals in their atmospheres.
- Rapid rotation lowers the threshold to mechanical mass loss by providing additional centrifugal acceleration.
- Rotation also drives CHE that mixes metals produced in the core with outer stellar layers.

(Morales-Rivera, RG +2025)

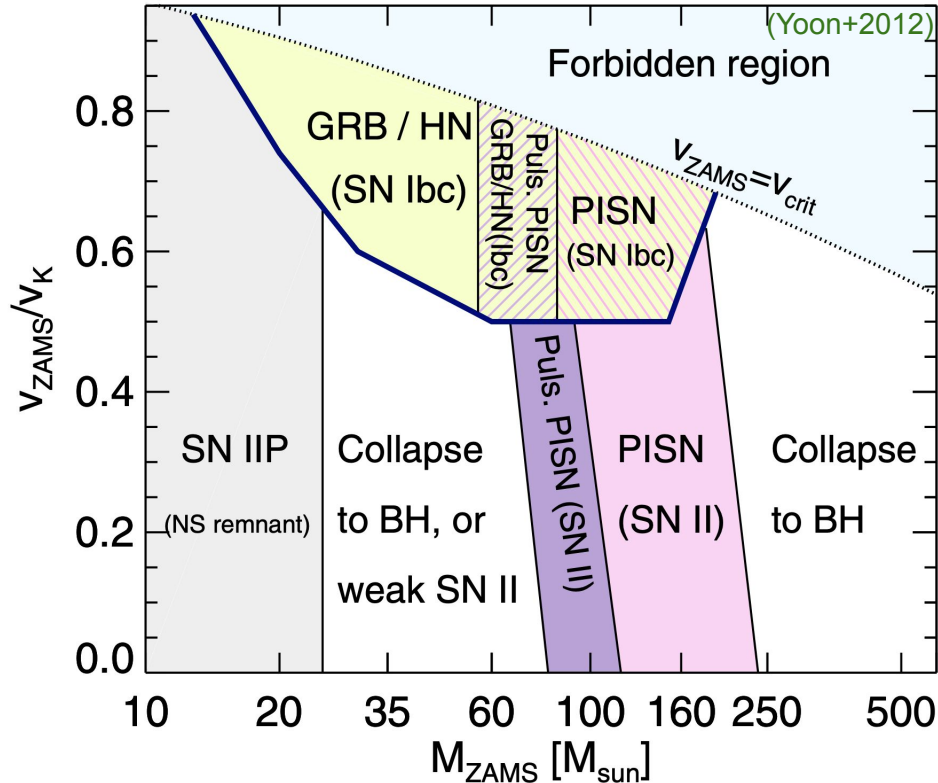


- **How much mass should be lost to line-driven winds** depends on the wind scheme and its efficiency.
- We use the Dutch wind scheme:

$$\dot{M}_{\text{wind}} = \eta_{\text{wind}} \dot{M}_{\text{Dutch}}$$

- The amount of mass loss is consistent with that expected in more metal-rich Wolf-Rayet stars

# Final outcome of rotating Pop III stars



## Requisites to produce a GRB

- The collapsar engine requires the formation of a **rapidly spinning BH** accreting stellar material at **high mass accretion rates**
- This requires that the specific angular momentum of the infalling stellar material exceeds

$$j \gtrsim 3 \times 10^{16} \text{ cm s}^{-1}$$

for a BH with mass  $M_{\text{BH}} = 5M_{\odot}$

- The jet power exceeds a critical value below which the jet can be choked.
- The engine remains active for long enough to break out of the star and power the GRB:

$$t_{\text{eng}} = t_{\text{bo}} + t_{\text{GRB}}$$

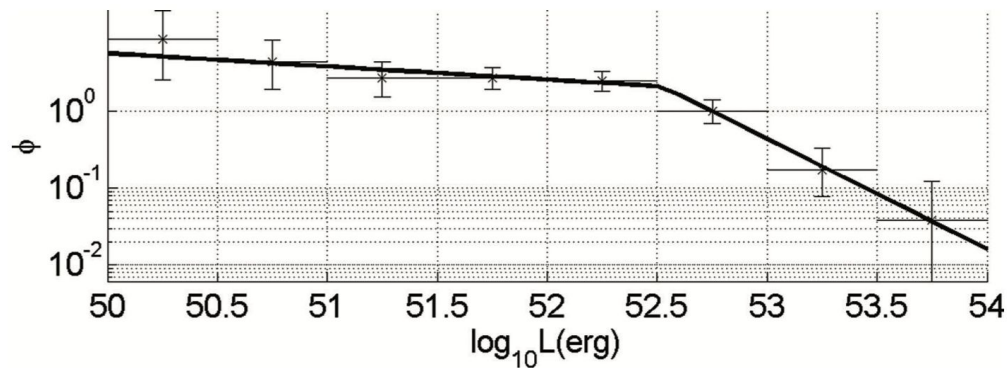
# Detection rate of Pop III GRBs

- The observed rate of GRBs is obtained from:  $\frac{dN_{\text{GRB}}}{dt_{\text{obs}}} = \eta_{\text{beam}} \int_0^z dz \Psi(z) \int_{L_{\text{lim}}(z)} \Phi(L) d\log L$

$$\eta_{\text{beam}} = \frac{\Delta\Omega}{4\pi} \approx \frac{\theta_j^2}{2} \sim \frac{1}{500} = \text{beaming fraction (Frail+2001)}$$

$\Phi(L)$  = luminosity function of GRBs above a detector-dependent limiting threshold luminosity  $L_{\text{lim}}(z)$

(Wanderman & Piran 2010)



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- The local comoving rate of GRBs at any redshift is given by:

$$\Psi(z) = \frac{d^2 N_{\text{GRB}}}{dt_{\text{obs}} dz} = \eta_{\text{GRB}}(z) \frac{\dot{\rho}_*(z)}{1+z} \frac{dV}{dz}$$

$\dot{\rho}(z)$  = local SFR,  $\frac{dV}{dz}$  = local volume element per unit redshift

- The efficiency of producing GRBs per unit stellar mass depends on the IMF and the rotation distribution of Pop III stars:**

$$\eta_{\text{GRB}} = \frac{\int_{m_{\text{low}}}^{m_{\text{high}}(\hat{\Omega}_0)} \xi(m) P_{\text{rot}}(\hat{\Omega}_0) dm d\hat{\Omega}_0}{\int_{m_{\text{min}}}^{m_{\text{max}}} m \xi(m) dm} \quad P_{\text{rot}}(\hat{\Omega}_0) \propto \hat{\Omega}_0^\delta \quad 0.6 \leq \hat{\Omega}_0 \leq 1$$

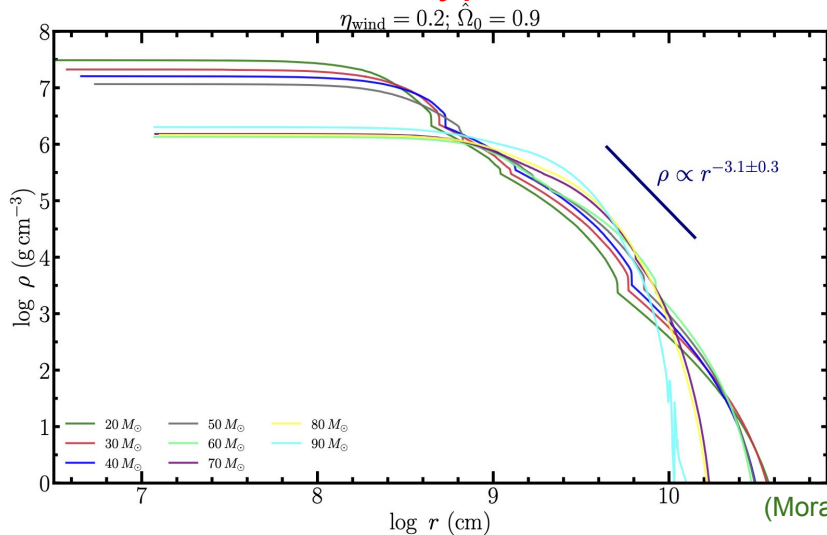
# MESA simulations: Stellar structure at core collapse

- We perform MESA simulations for a grid of Pop III stars with:

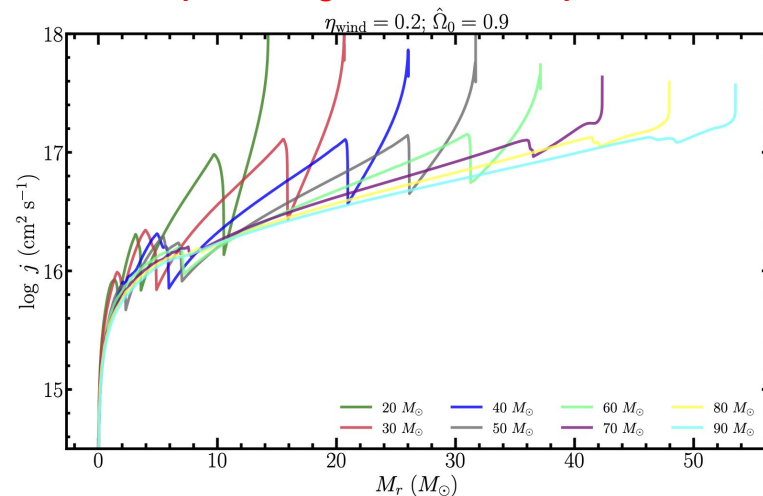
Initial stellar mass:  $20 \leq \frac{M_{\text{ZAMS}}}{M_{\odot}} \leq 100$  rotation rate:  $0.6 \leq \frac{\Omega_0}{\Omega_{\text{crit}}} \leq 0.9$  Wind efficiency:  $0.2 \leq \eta_{\text{wind}} \leq 1$

- Stars are evolved to just before core-collapse or depletion of core Silicon in the advanced burning stage.

**Stellar density profile**



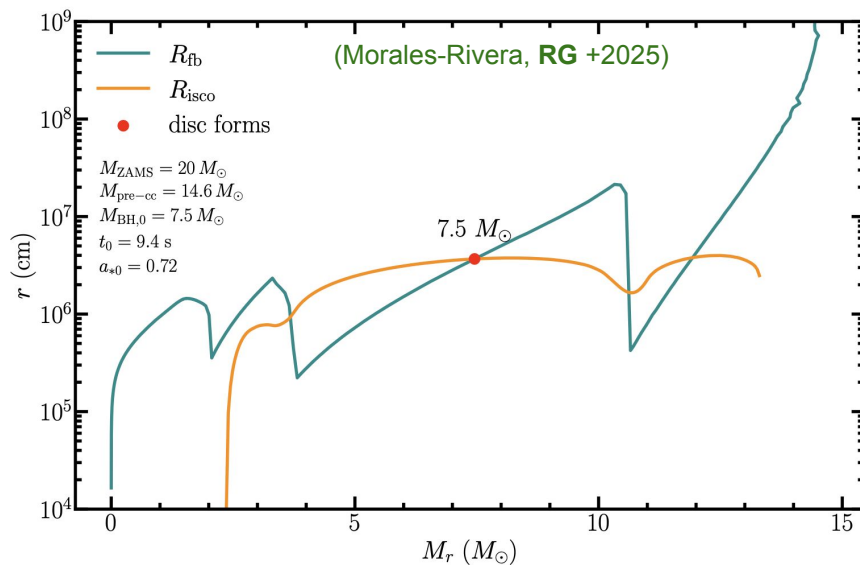
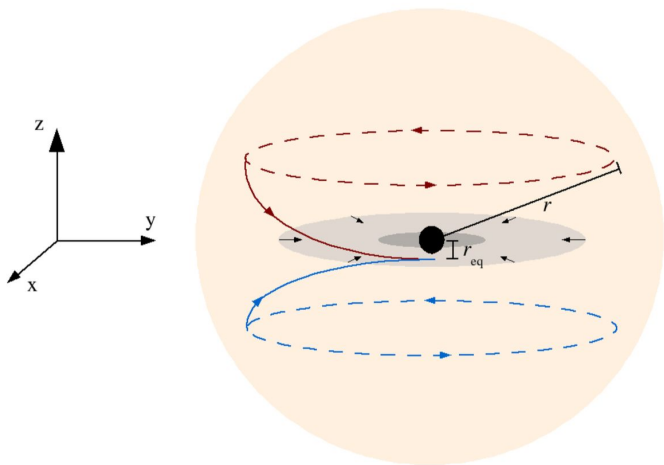
**Specific angular momentum profile**



(Morales-Rivera, RG +2025)

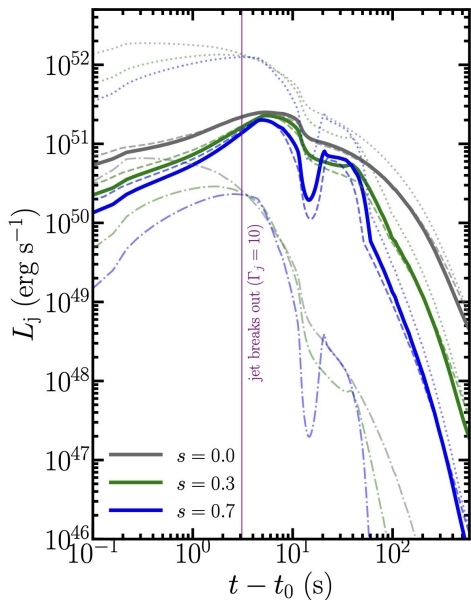
# Model of fallback and accretion

- We construct a semi-analytical model of stellar material fallback onto the equatorial plane and BH if lacking sufficient angular momentum (Kumar, Narayan, Johnson 2008).
- Initially, the core forms a proto-NS that promptly collapses into a BH due to fallback accretion
- Material continues to plunge directly into the BH until it circularizes outside of  $R_{\text{ISCO}}$  and forms a disc.



# BH properties & jet power

- From the mass fallback rate, we calculate the rate of accretion and growth of BH mass and its spin.
- We account for the mode of accretion: NDAF, ADAF, or combination, and disc-driven outflows.



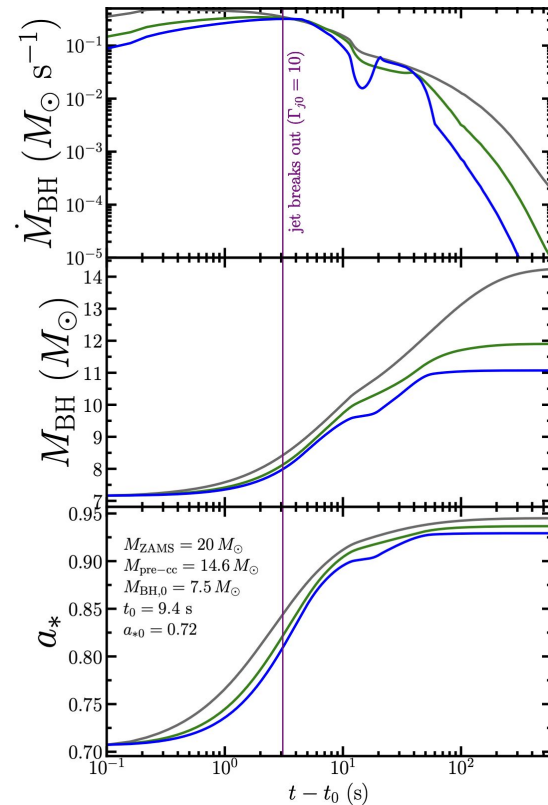
- Since we don't know how GRB jets are launched, we calculate the jet power for two launching mechanisms:

- Neutrino and anti-neutrino annihilation
- Jet efficiency obtained from GRMHD simulations

$$L_j = \eta_j (a_*) \dot{M} c^2$$

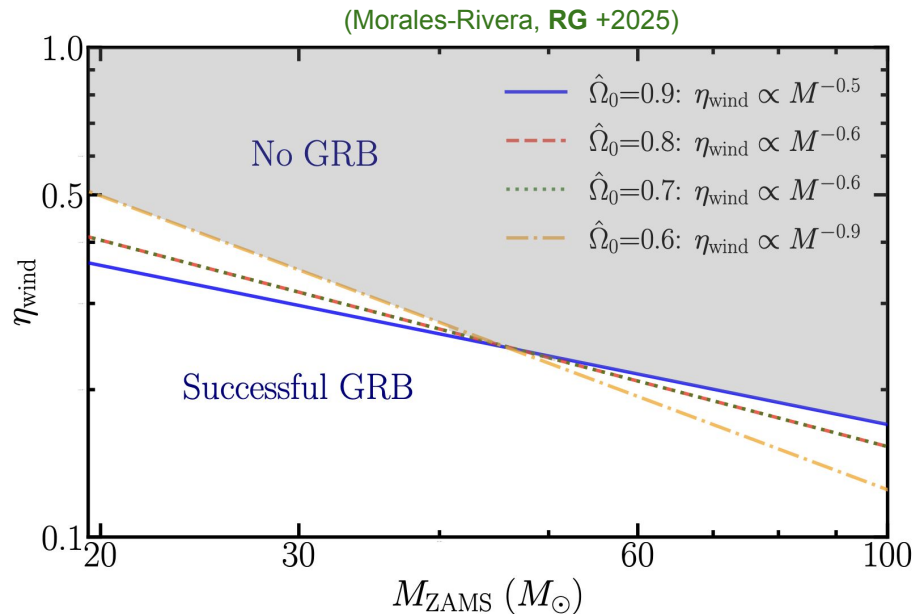
- Jets from magnetically arrested discs

(Morales-Rivera, RG +2025)



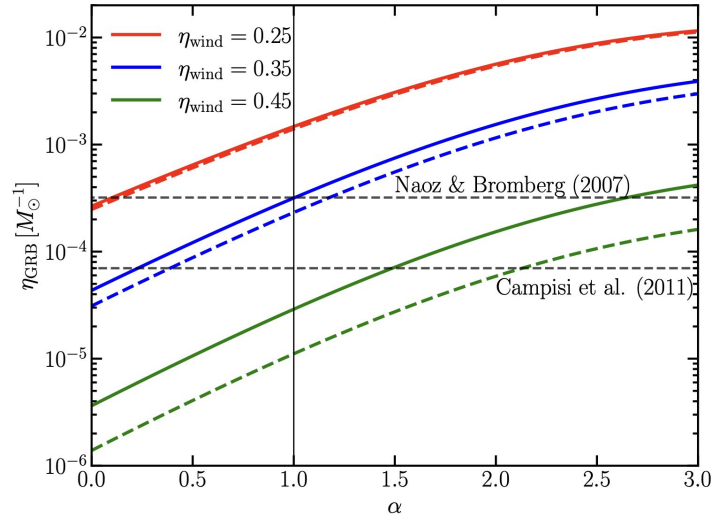
# Phase space for producing a GRB

- This new phase space factors into the effects of **rotation, wind mass loss, jet launching** and **breakout**.
- For any given mass, the wind efficiency must be smaller than some critical value. Otherwise, the star will lose too much mass and angular momentum.
- The maximum ZAMS mass up to which Pop III stars can produce a GRB depends on the initial rotation rate for a given wind efficiency.

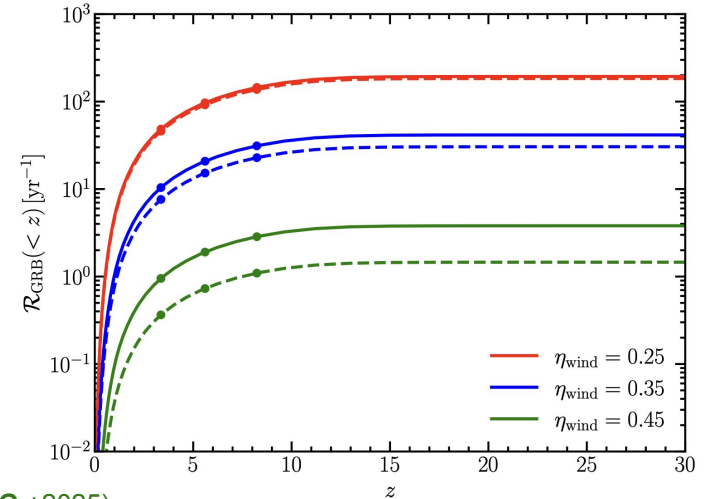


# GRB production efficiency & rate

- GRB production rate drops with increased wind mass-loss efficiency as more angular momentum is lost.
- Top-heavy IMF also leads to a low production efficiency due to the cap on the high-mass for a given wind mass-loss efficiency.
- 75% of the GRBs in the total integrated rate come from  $z < 8$ .



(Morales-Rivera, RG +2025)



# Conclusions

- The efficiency of GRB production per unit mass (or rate of GRBs) is very sensitive to the wind mass-loss efficiency and/or scheme.
  - Observed rate of high- $z$  GRBs can constrain this highly uncertain mass-loss efficiency.
  - For  $\eta_{\text{wind}} = 0.45$  and a flat IMF,  $\eta_{\text{GRB}} \sim 10^{-5} M_{\odot}^{-1}$  that yields a **detection rate of ~4 - 5 Pop III GRBs over the past 20 years of Swift/BAT**
- 
- We find that GRBs from Pop III progenitors are **similar in their intrinsic properties with standard GRBs**, which makes it hard to discern their high- $z$  origin.
  - This degeneracy may be broken with **high-resolution IR spectroscopy** and by looking for absence of absorption from iron-group elements.
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- If we don't detect an all-sky rate of 2 GRBs per year for a Swift/BAT-type detector, then the wind mass-loss must be high and/or these stars are born as slow rotators.
  - In that case, only **progenitors in a binary** may be able to produce the right conditions for GRB production.