

② Massless vector boson polarisation can be written in terms of spinor products/strings as follows

$$\epsilon_+^{*\mu}(p_i, r) = \frac{1}{\sqrt{2}} \frac{\langle r \gamma^\mu i \rangle}{\langle ri \rangle}$$

$$\epsilon_-^{*\mu}(p_i, r) = -\frac{1}{\sqrt{2}} \frac{\text{tr } \gamma^\mu i \rangle}{\text{tr } i \rangle}$$

where r is an arbitrary massless reference vector.

$$\begin{aligned} \text{a.} \quad \epsilon_+^{*}(p_i, r) \cdot p_i &= \frac{1}{\sqrt{2}} \frac{\langle r \gamma^\mu i \rangle}{\langle ri \rangle} p_{i\mu} \\ &= \frac{1}{\sqrt{2}} \frac{\langle r \not{p}_i i \rangle}{\langle ri \rangle} \\ &= \frac{1}{\sqrt{2} \langle ri \rangle} \frac{\langle r (i \rangle [i + i] \langle i \rangle) i \rangle}{\langle ri \rangle} \\ &= \frac{1}{\sqrt{2} \langle ri \rangle} \frac{\langle ri \rangle [ii] + \langle ri \rangle \langle ii \rangle}{\langle ri \rangle} \\ &= 0 \end{aligned}$$

$$\begin{aligned} \epsilon_-^{*}(p_i, r) \cdot p_i &= -\frac{1}{\sqrt{2}} \frac{\text{tr } \gamma^\mu i \rangle}{\text{tr } i \rangle} p_{i\mu} \\ &= -\frac{1}{\sqrt{2}} \frac{[\text{tr } i \rangle \langle ii \rangle + [\text{tr } i \rangle [ii]]}{\text{tr } i \rangle} \end{aligned}$$



$$= 0$$



$$\begin{aligned}\epsilon_+^{*\mu}(p_i, r) r_\mu &= \frac{1}{\sqrt{2}} \frac{\langle r \sigma^\mu i \rangle}{\langle r i \rangle} r_\mu \\ &= \frac{1}{\sqrt{2}} \frac{\langle r r i \rangle}{\langle r i \rangle} = 0\end{aligned}$$

$$\begin{aligned}\epsilon_-^{*\mu}(p_i, r) r_\mu &= -\frac{1}{\sqrt{2}} \frac{[r \sigma^\mu i]}{[r i]} r_\mu \\ &= -\frac{1}{\sqrt{2}} \frac{[r r i]}{[r i]} = 0\end{aligned}$$

b.) $\epsilon_+^{*\mu}(r_1, r_1) \cdot \epsilon_{+\mu}^*(p_2, r_2)$

$$= \frac{1}{\sqrt{2}} \frac{\langle r_1 \sigma^\mu 1 \rangle}{\langle r_1 1 \rangle} \frac{1}{\sqrt{2}} \frac{\langle r_2 \sigma_\mu 2 \rangle}{\langle r_2 2 \rangle}$$

$$= \frac{1}{2} \frac{2 \langle r_1 r_2 \rangle [21]}{\langle r_1 1 \rangle \langle r_2 2 \rangle} \quad \left(\text{using Fierz identity} \right)$$

$$\epsilon_-^{*\mu}(p_1, r_1) \cdot \epsilon_{-\mu}^*(p_2, r_2)$$

$$= \left(-\frac{1}{\sqrt{2}} \right) \frac{[r_1 \sigma^\mu 1]}{[r_1 1]} \left(-\frac{1}{\sqrt{2}} \right) \frac{[r_2 \sigma_\mu 2]}{[r_2 2]}$$

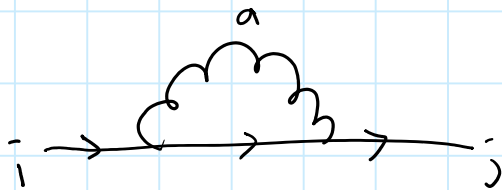
$$= \frac{[r_1 r_2] \langle 21 \rangle}{[r_1 1] [r_2 2]}$$

$$\epsilon_{+}^{\mu} (p_1, r_1) \epsilon_{-\mu} (p_2, r_2) = \frac{1}{\sqrt{2}} \frac{\langle r_1 \sigma^{\mu 1} \rangle}{\langle r_1 1 \rangle} \left(-\frac{1}{\sqrt{2}} \right) \frac{[r_2 \sigma_{\mu 2}]}{[r_2 2]} \\ = - \frac{\langle r_1 2 \rangle [r_2]}{\langle r_1 1 \rangle [r_2 2]}$$

$$\epsilon_{-}^{\mu} (p_1, r_1) \epsilon_{+\mu} (p_2, r_2) = \left(-\frac{1}{\sqrt{2}} \right) \frac{[r_1 \sigma_{\mu 1}]}{[r_1 1]} \frac{1}{\sqrt{2}} \frac{\langle r_2 \sigma^{\mu 2} \rangle}{\langle r_2 2 \rangle} \\ = - \frac{[r_2] \langle r_2 1 \rangle}{[r_1 1] \langle r_2 2 \rangle}$$

③  $t_{ii}^a = \text{tr}(t^a) = 0$

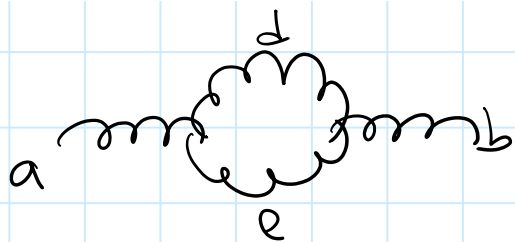
 $t_{ij}^a t_{ji}^b = \text{tr}(t^a t^b) \\ = \frac{1}{2} \delta^{ab}$



$$C_F = \frac{N_c^2 - 1}{2 N_c}$$

"quadratic Casimir operator
for $SU(N_c)$ in fundamental
reps."

$$\Rightarrow t_{ji}^a t_{mi}^a \\ = \frac{1}{2} \left[\delta_{ji} \delta_{mm} - \frac{1}{N_c} \delta_{jm} \delta_{mi} \right] \\ = \frac{1}{2} \left[N_c - \frac{1}{N_c} \right] \delta_{ji} \\ = C_F \delta_{ji}$$



$$\rightarrow f^{ade} f^{bde}$$

$$t^d [t^a, t^b] = i f^{abc} t^c t^d$$

$$\text{tr} (t^d [t^a, t^b]) = i f^{abc} \text{tr} (t^c t^d)$$

$$2 \text{tr} (t^d [t^a, t^b]) = i f^{abc}$$

$$f^{abc} = -2i \text{tr} (t^c [t^a, t^b])$$

$$f^{ade} f^{bde} = -4 \text{tr} (t^e [t^a, t^d]) \cdot \text{tr} (t^e [t^b, t^d])$$

$$= -4 t_{ij}^e [t^a, t^d]_{ji}$$

$$t_{kl}^e [t^b, t^d]_{lk}$$

$$= -4 \left(t_{ij}^e t_{jm}^a t_{mi}^d - t_{ij}^e t_{jm}^d t_{mi}^a \right)$$

$$\left(t_{kl}^e t_{ln}^b t_{nk}^d - t_{kl}^e t_{ln}^d t_{nk}^b \right)$$

```
In[26]:= tmp = -4 * (t[e, i, j] * t[a, j, m] * t[d, m, i] - t[e, i, j] * t[d, j, m] * t[a, m, i]) *
            (t[e, k, l] * t[b, l, n] * t[d, n, k] - t[e, k, l] * t[d, l, n] * t[b, n, k])
```

```
Out[26]= -4 (-t[a, m, i] * t[d, j, m] * t[e, i, j] + t[a, j, m] * t[d, m, i] * t[e, i, j])
            (-t[b, n, k] * t[d, l, n] * t[e, k, l] + t[b, l, n] * t[d, n, k] * t[e, k, l])
```

```
In[34]:= FierzId = {
            t[a_, i_, j_] * t[a_, k_, l_] => (1/2) * (d[i, l] * d[k, j] - 1/Nc * d[i, j] * d[k, l])
        };
```

```
In[35]:= Contract = {
            d[i1_, i2_] * d[i2_, i3_] => d[i1, i3],
            t[a_, i1_, i2_] * d[i2_, i3_] => t[a, i1, i3],
            t[a_, m_, m_] => 0,
            d[i1_, i1_] => Nc,
            t[a_, i1_, i2_] * t[b_, i2_, i1_] => trT[a, b]
        };
```

```
In[36]:= (Expand[tmp] //. FierzId // Expand) //. Contract /. trT[a_, b_] => (1/2) * d[a, b]
```

```
Out[36]= Nc d[a, b]
```

→ repeated use of Fierz identity

$$t_{ij}^a t_{kl}^a = \frac{1}{2} \left(\delta_{il} \delta_{kj} - \frac{1}{N_c} \delta_{ij} \delta_{kl} \right)$$

and $\text{tr}(t^a t^b) = \frac{1}{2} \delta^{ab}$

$$\begin{aligned}
 \text{cloud diagram} &= f^{ade} f^{bde} \\
 &= N_c \delta^{ab} \\
 &\quad \text{" } C_A
 \end{aligned}$$

