

IRN Terascale Montpellier 2025

Freeze-in with low reheating temperature

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Work in progress with A. Goudelis, A. Lessa



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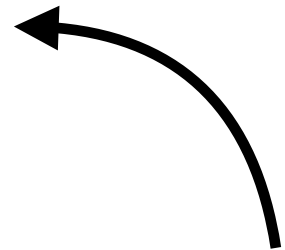
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Matrix elements, contains the particle physics information

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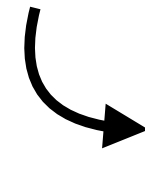
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Thermally averaged cross section, contains
 $|M|^2$



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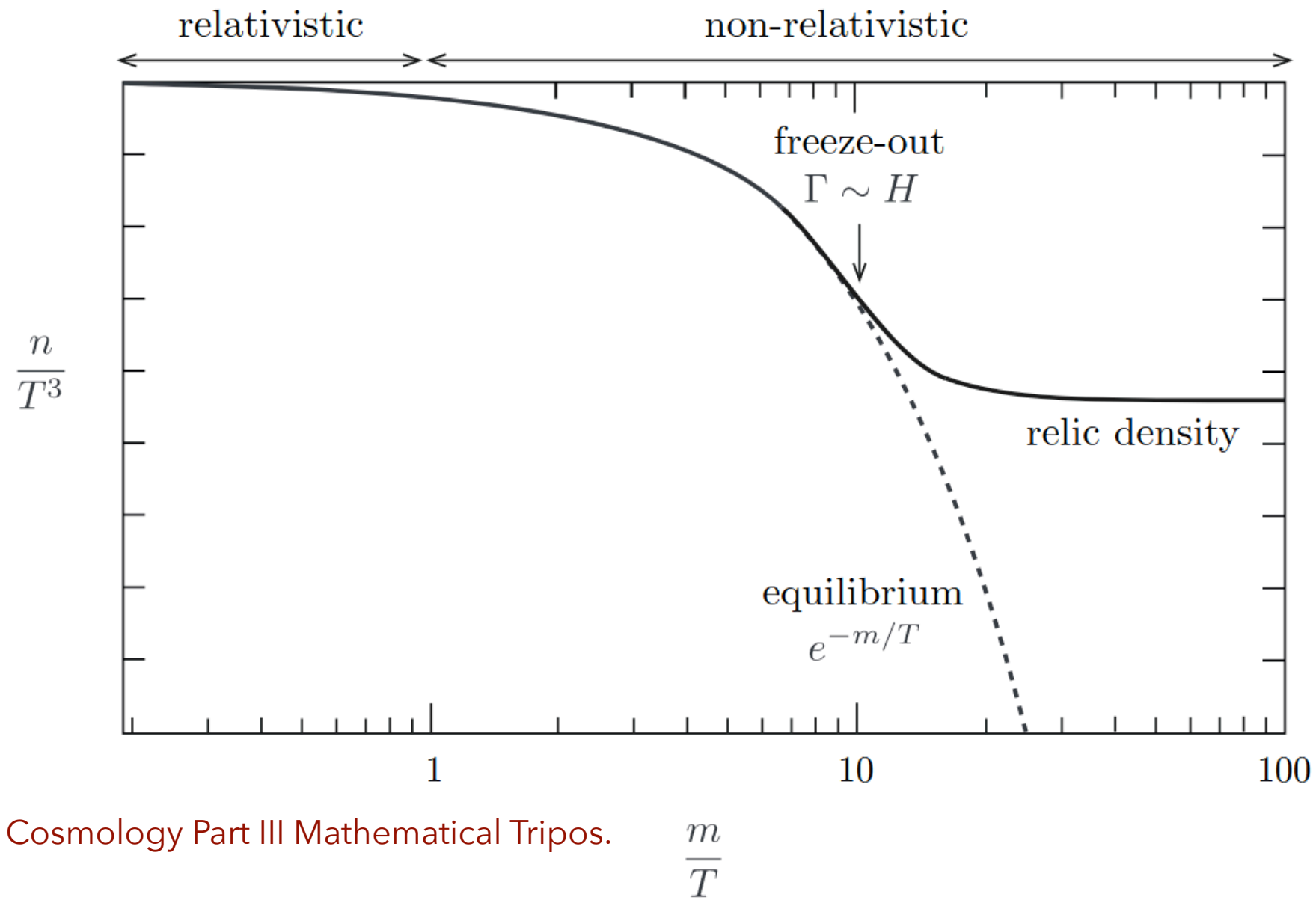
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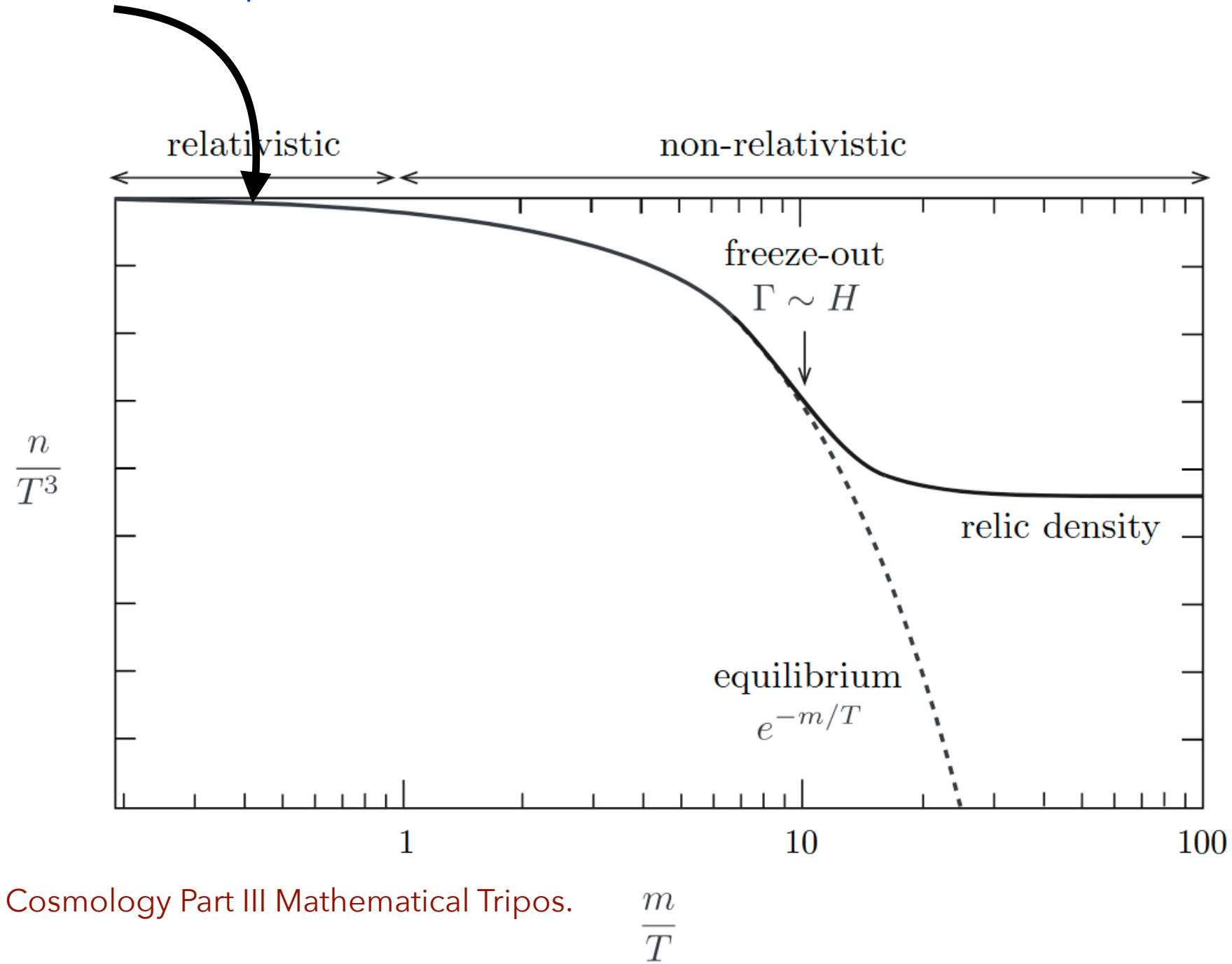


Daniel Baumann. Cosmology Part III Mathematical Tripos.

$\frac{m}{T}$

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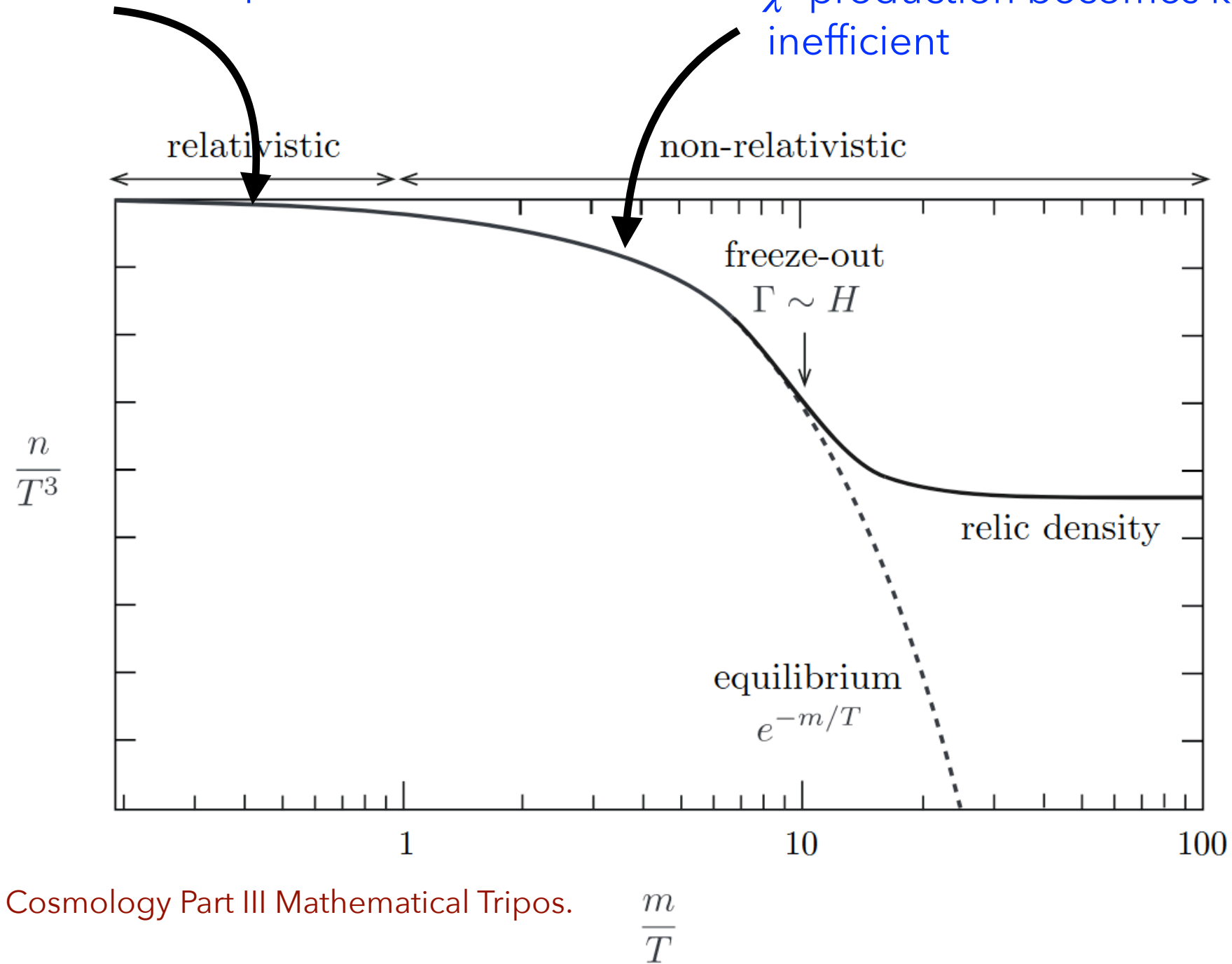
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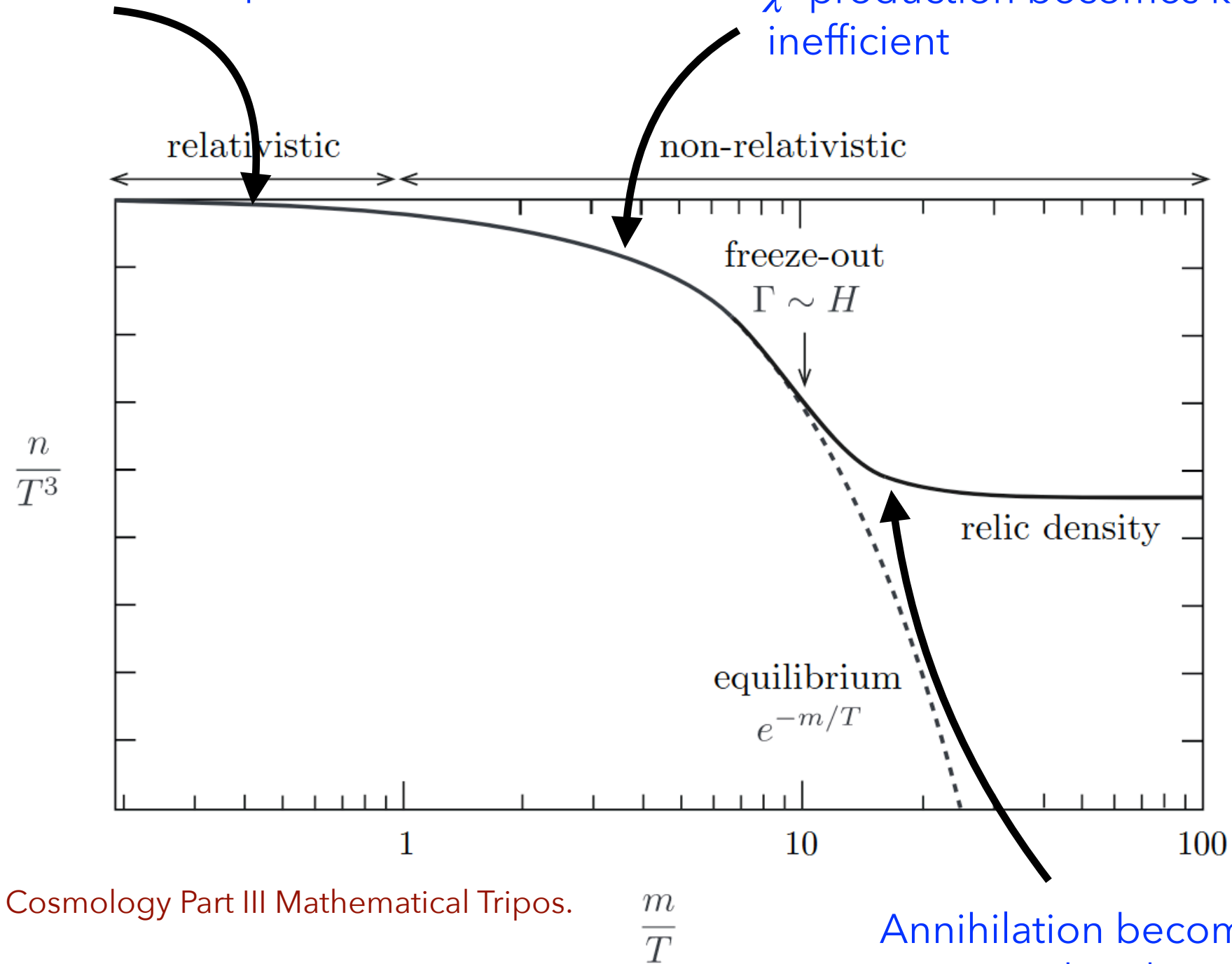


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Annihilation becomes inefficient compared to the expansion rate of the universe, χ freeze-out

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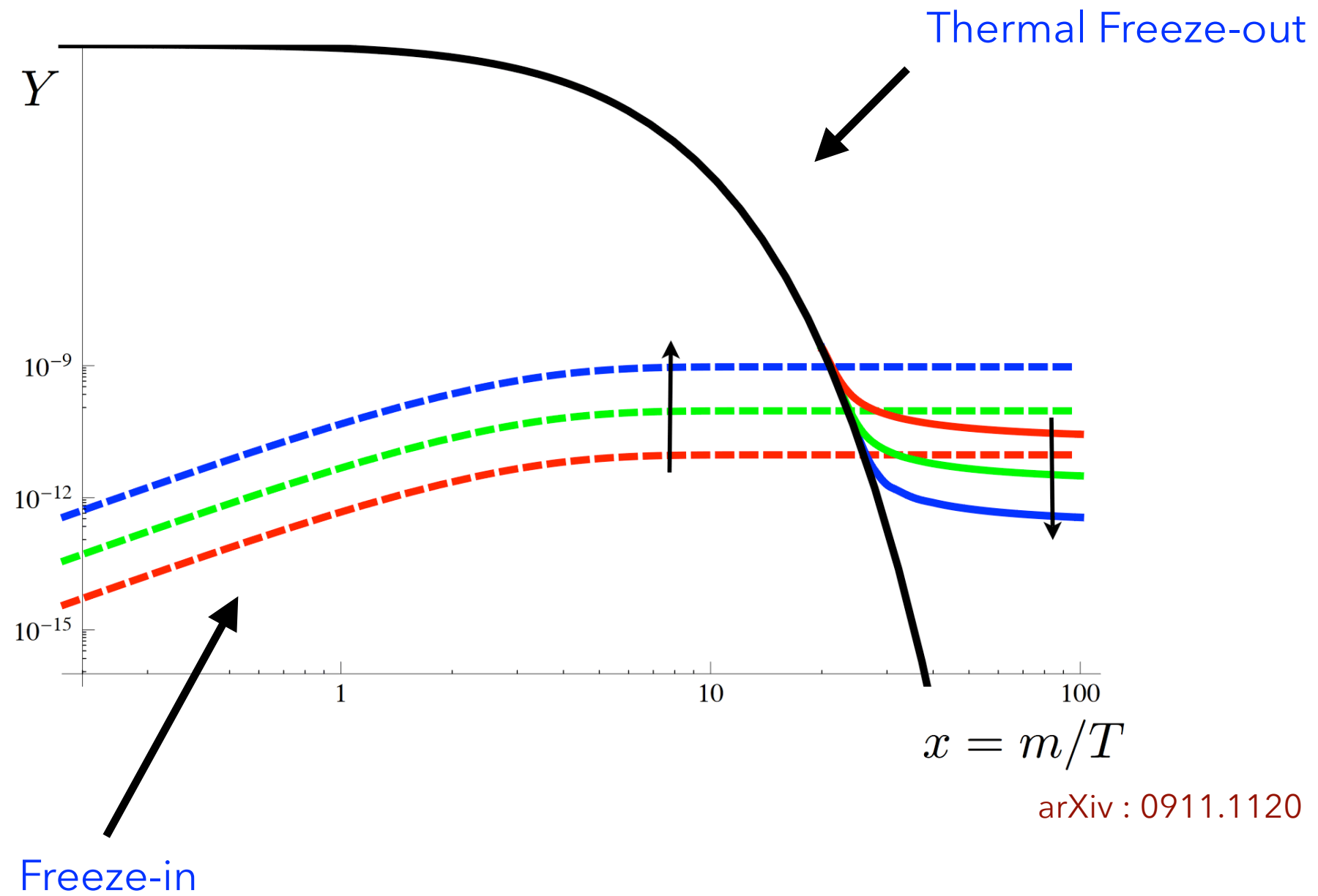
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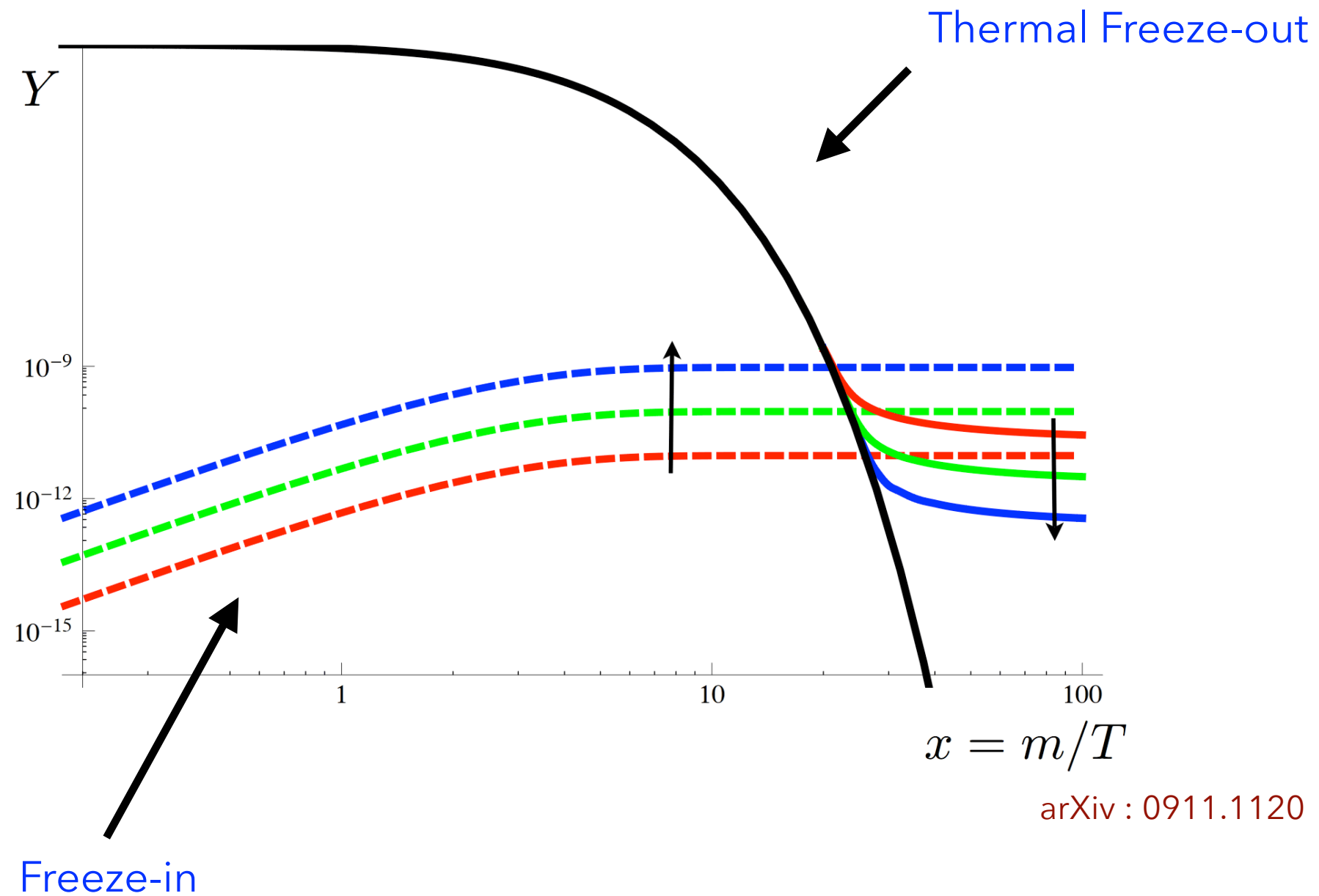
We can integrate this equation to obtain a solution as

$$Y_\chi = \int_{T_0}^{T_{rh}} \frac{dT}{T\bar{H}(T)s(T)} < \sigma v > Y_\chi^{eq} Y_X^{eq}$$

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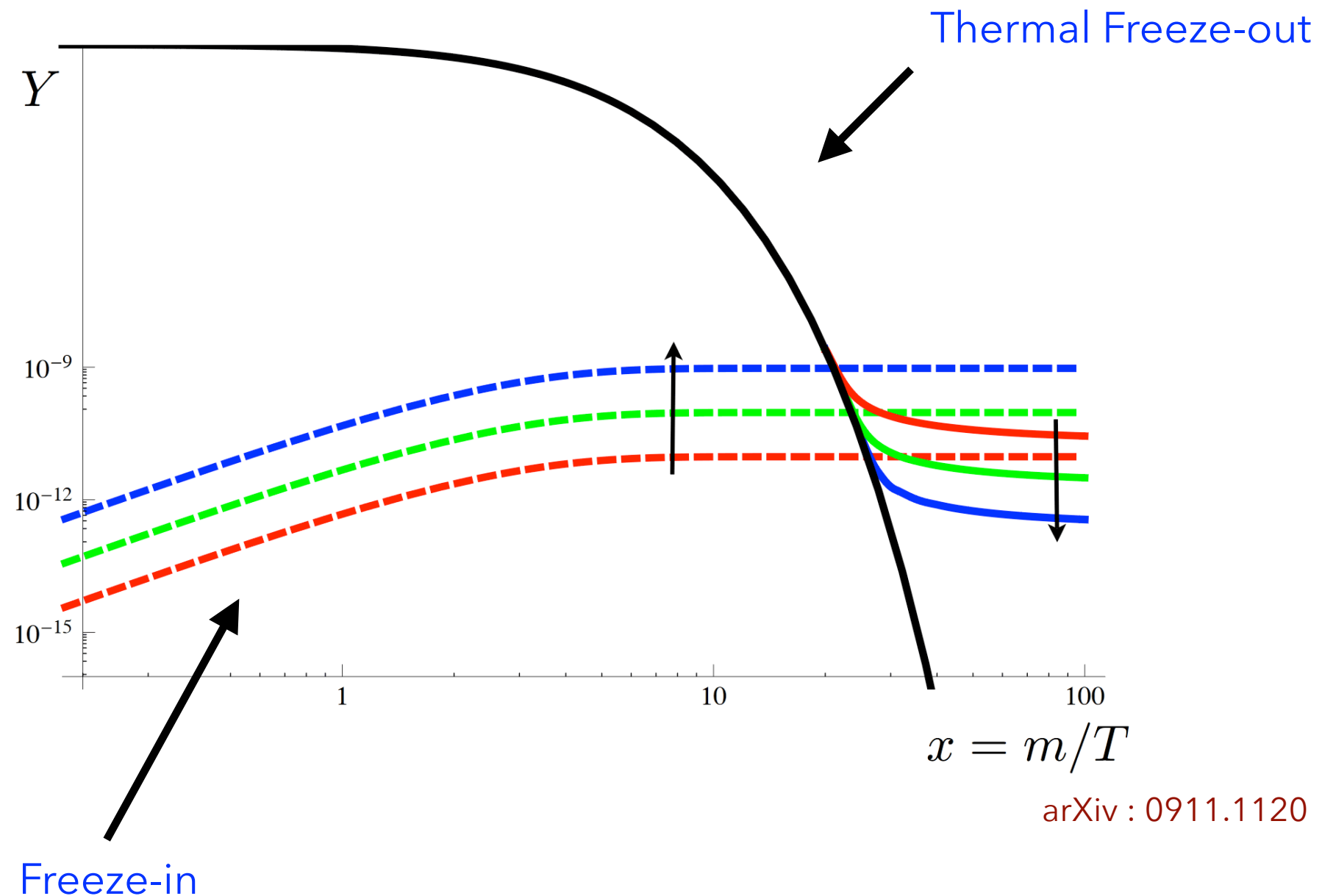
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What happens if one considers low T_{rh} ?

Case study : a charged parent model

Consider a real scalar singlet s , not charged under the SM gauge groups, and a vector-like fermion F , singlet under $SU(2)$. Both are odd under a \mathbf{Z}_2 symmetry to ensure the lightest state is stable.

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We will focus on the case in which F is color-neutral and carries hypercharge (Vector-like « lepton »)

Contributions to the relic density

There are usually two kinds of process contributing to the relic density

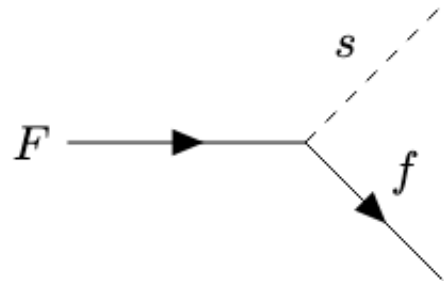
$$Y_s = \int_{T_0}^{T_{rh}} \frac{dT}{T \bar{H}(T) s(T)} \left(\mathcal{N}(F \rightarrow s_0, f) + 2\mathcal{N}(f, f \rightarrow s_0, s_0) \right)$$

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Production from the decay of
 F into s and SM fermion



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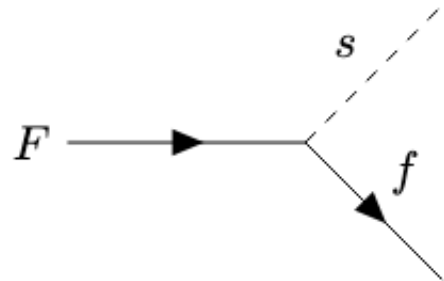
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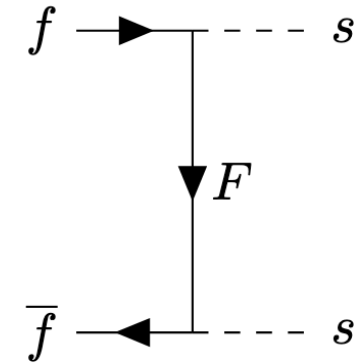


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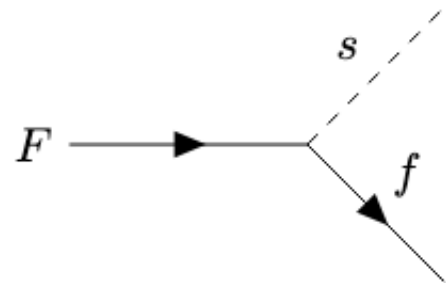
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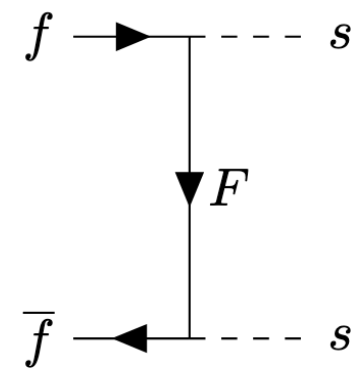


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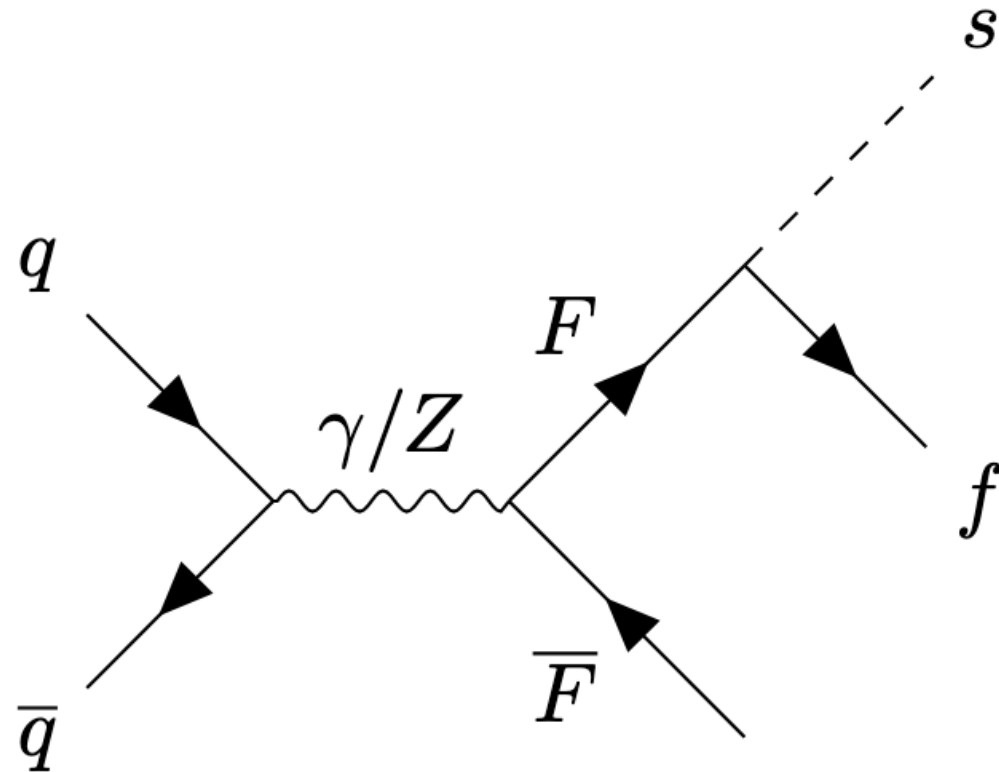
There are also additional contributions to the relic density such as the « super-WIMP » contribution (ie. mediator decay after its freeze-out) if $T_{rh} > T_{FO}$.

In the case $T_{rh} < T_{FO}$, the decay of frozen-in F 's

Constraints from LHC searches

For LHC searches, we focus on the detection of the mediator particle : F

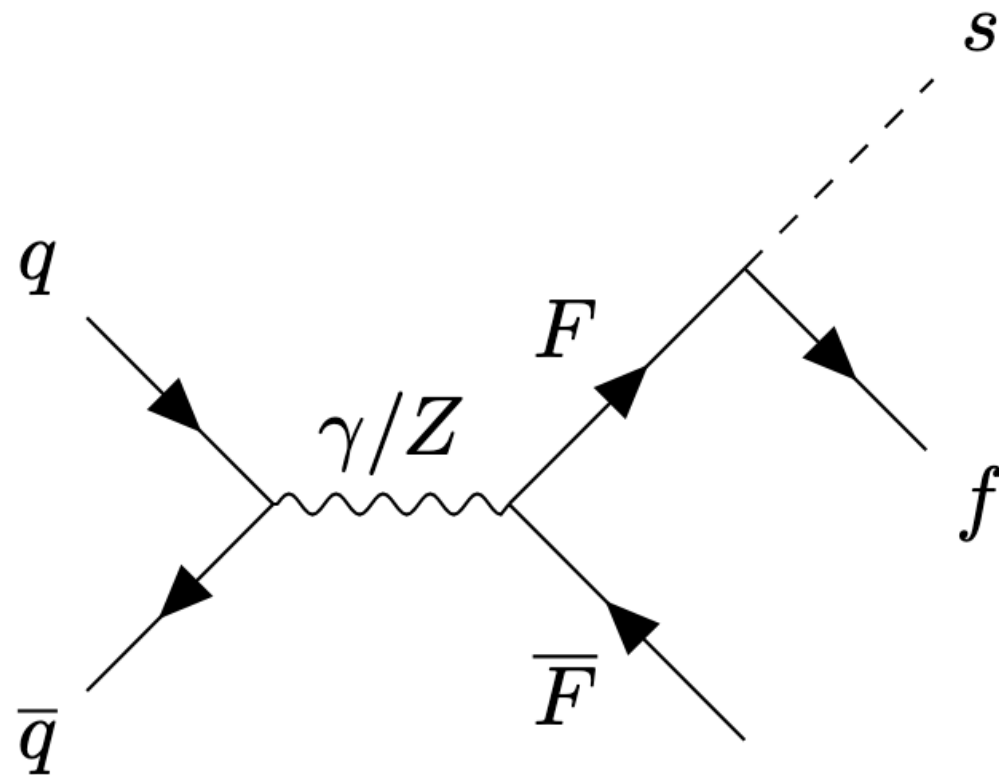
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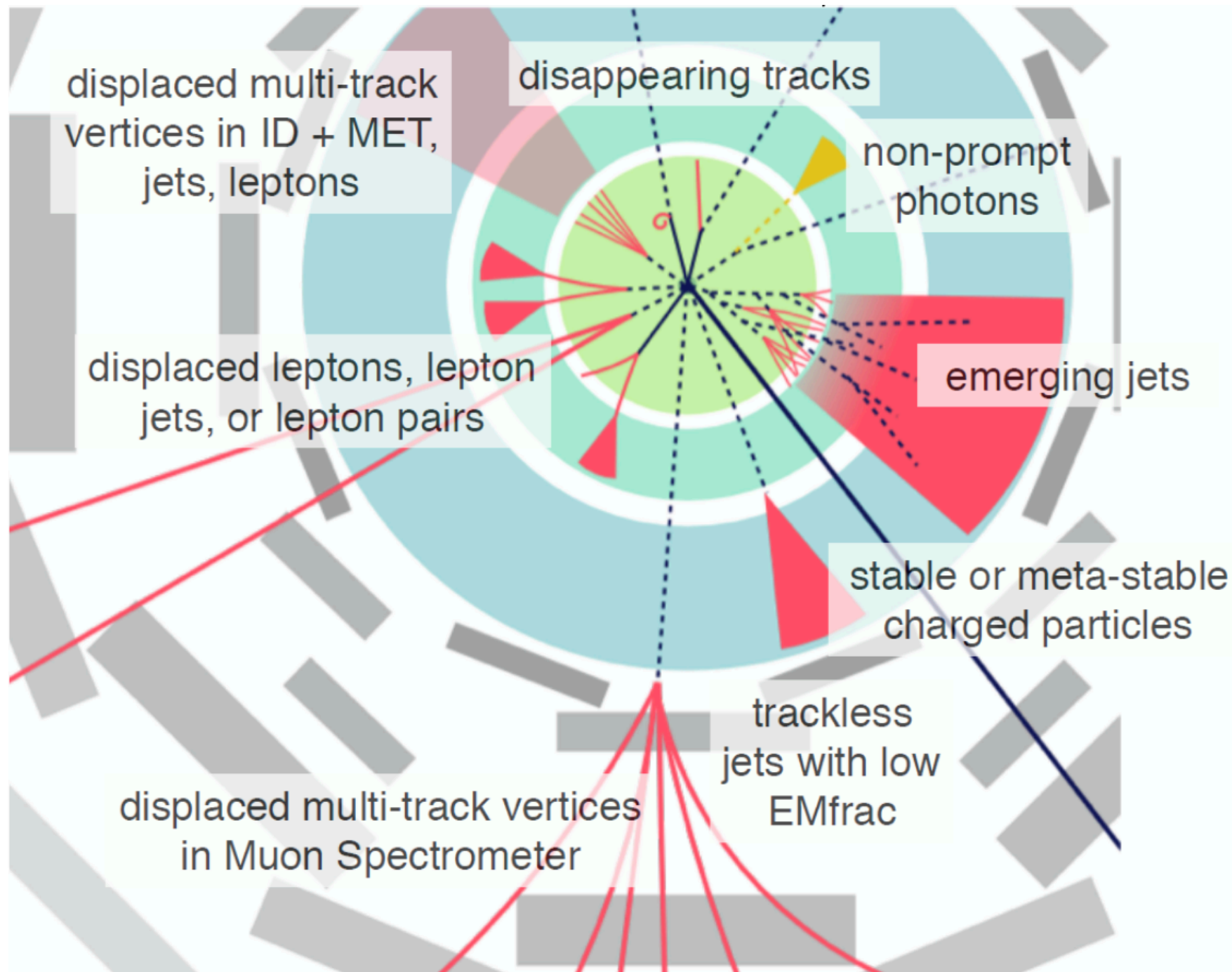
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In high T_{rh} freeze-in, this model predicts a very longed-lived F in order to saturate the relic density constraint

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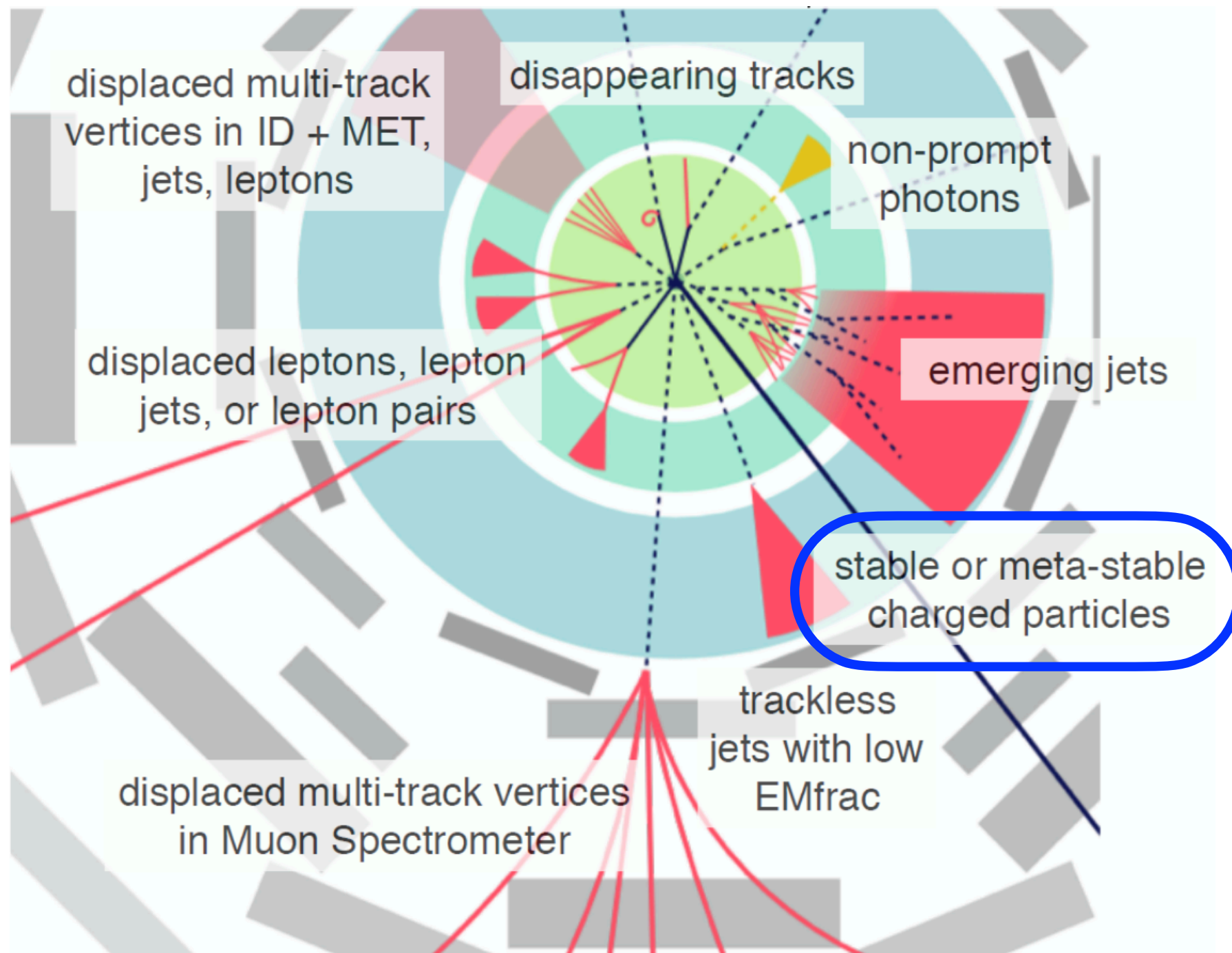
Depending on the lifetime of F , we can have different types of signatures in colliders



Graphic credits : Heather Russel

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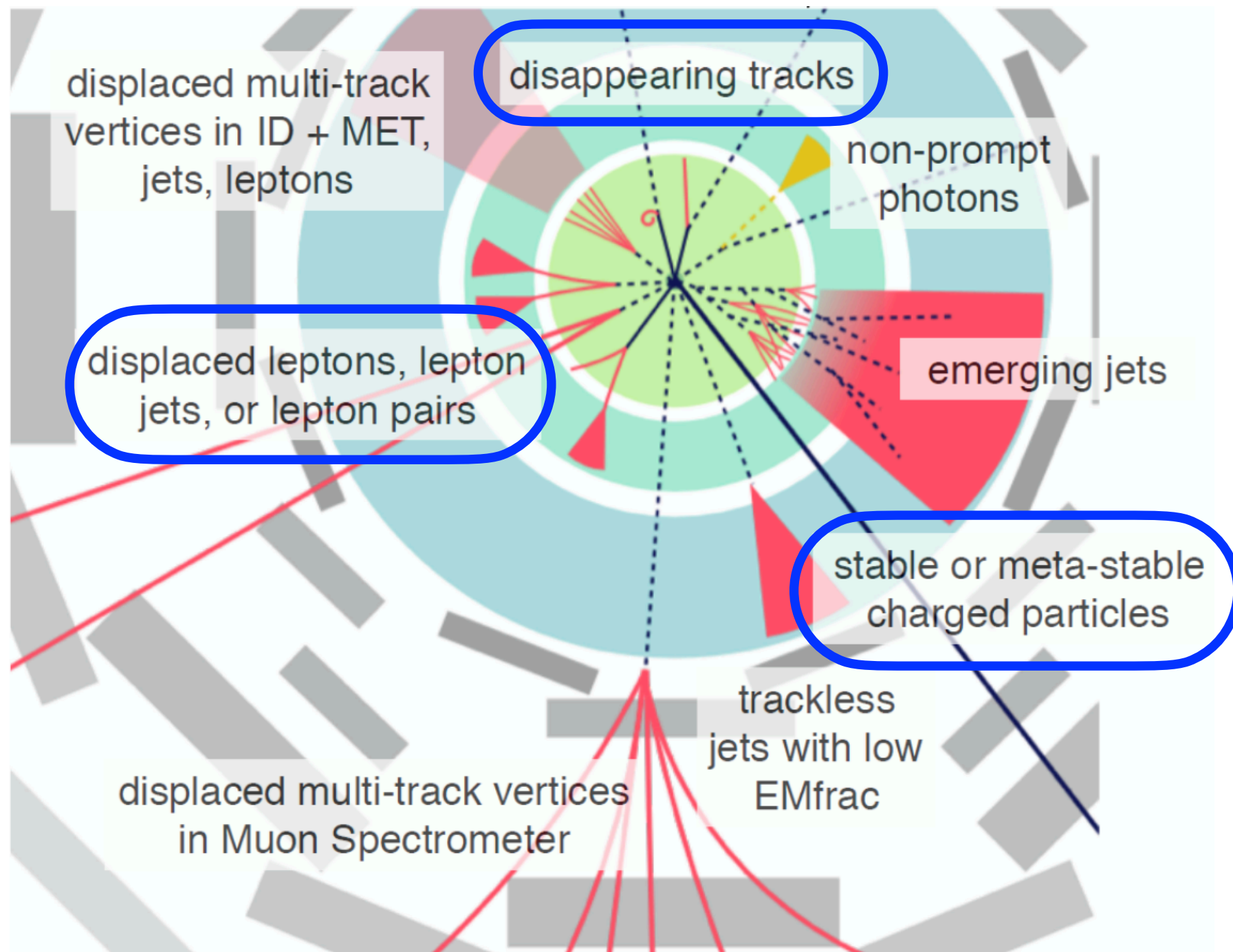
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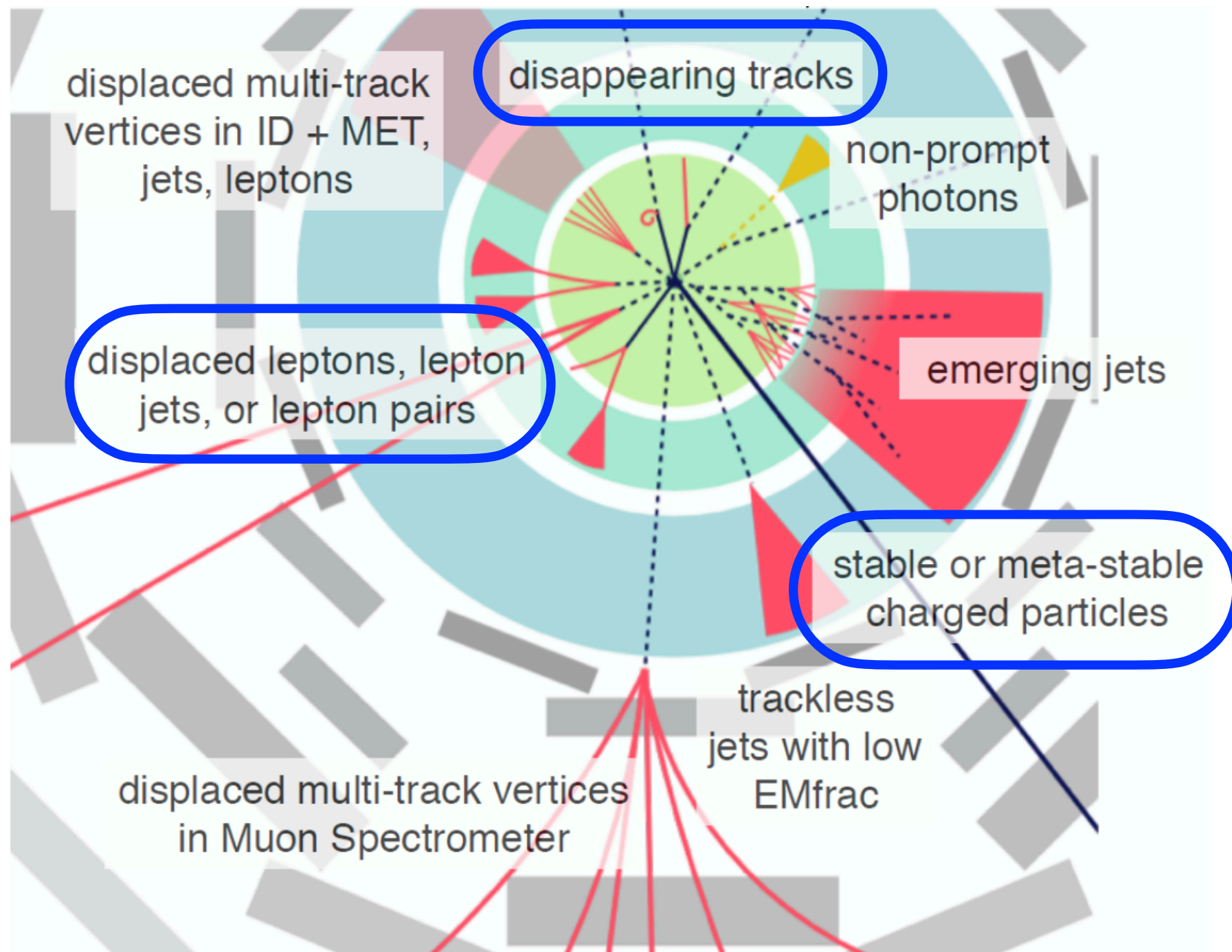
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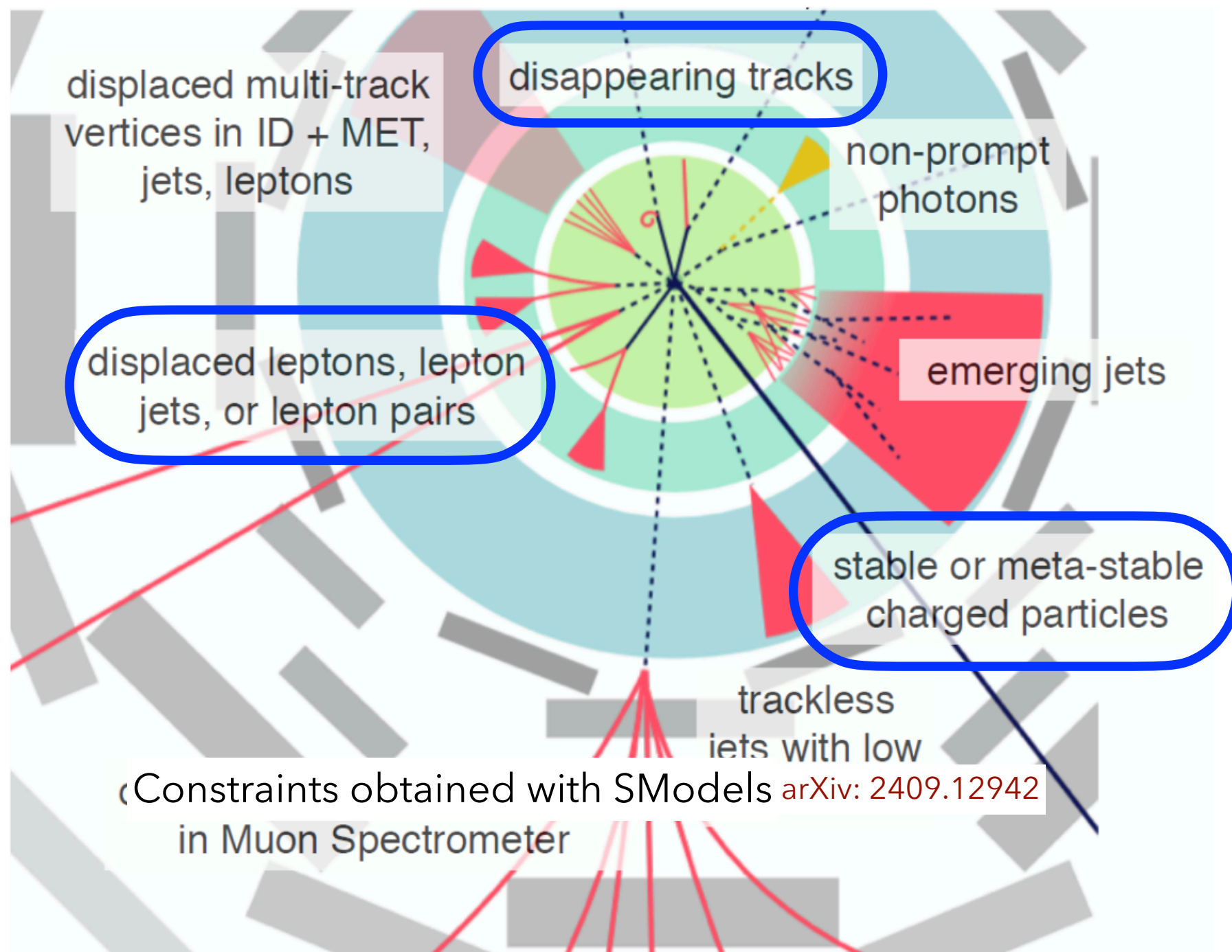


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+ eventually prompt searches for short lifetimes

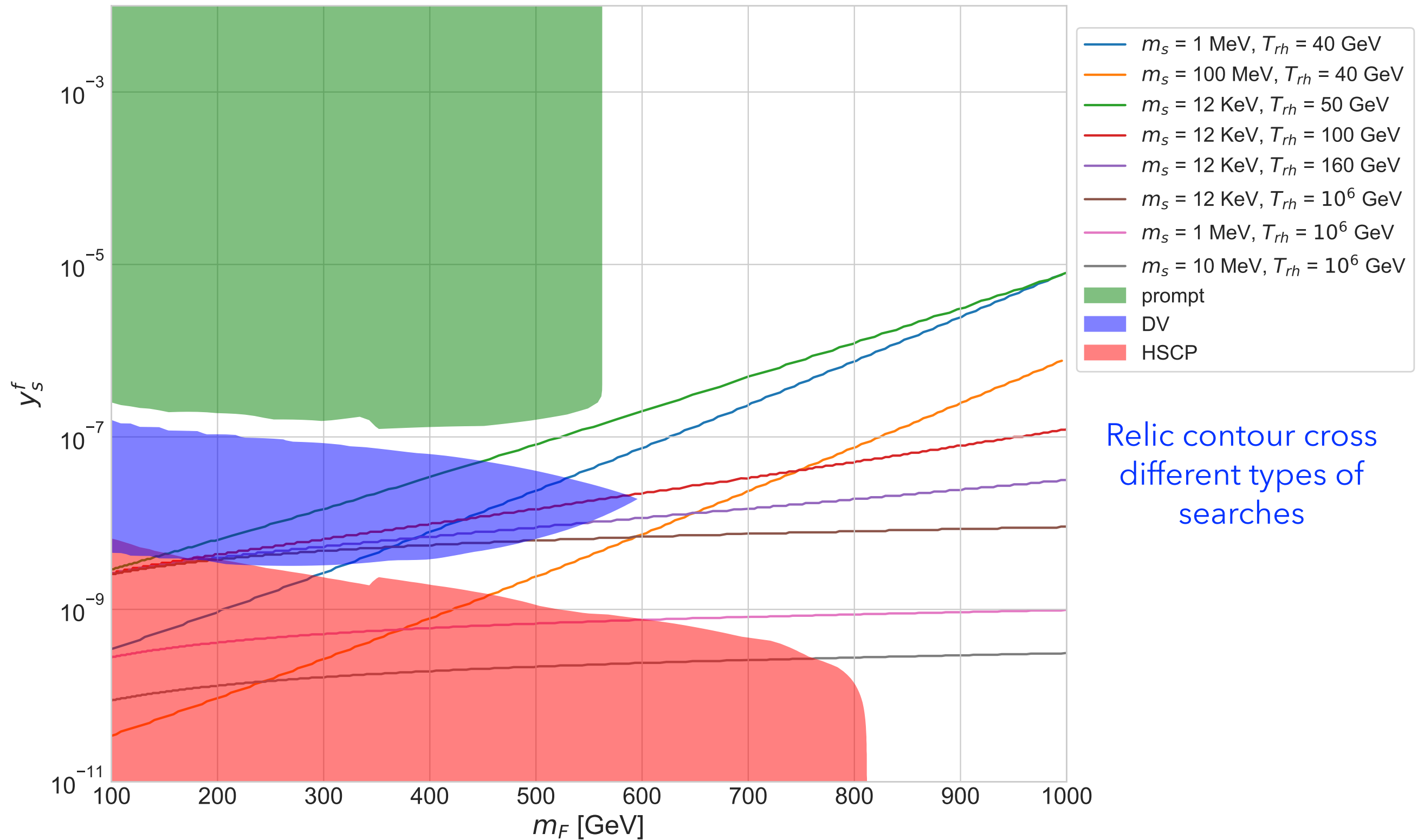
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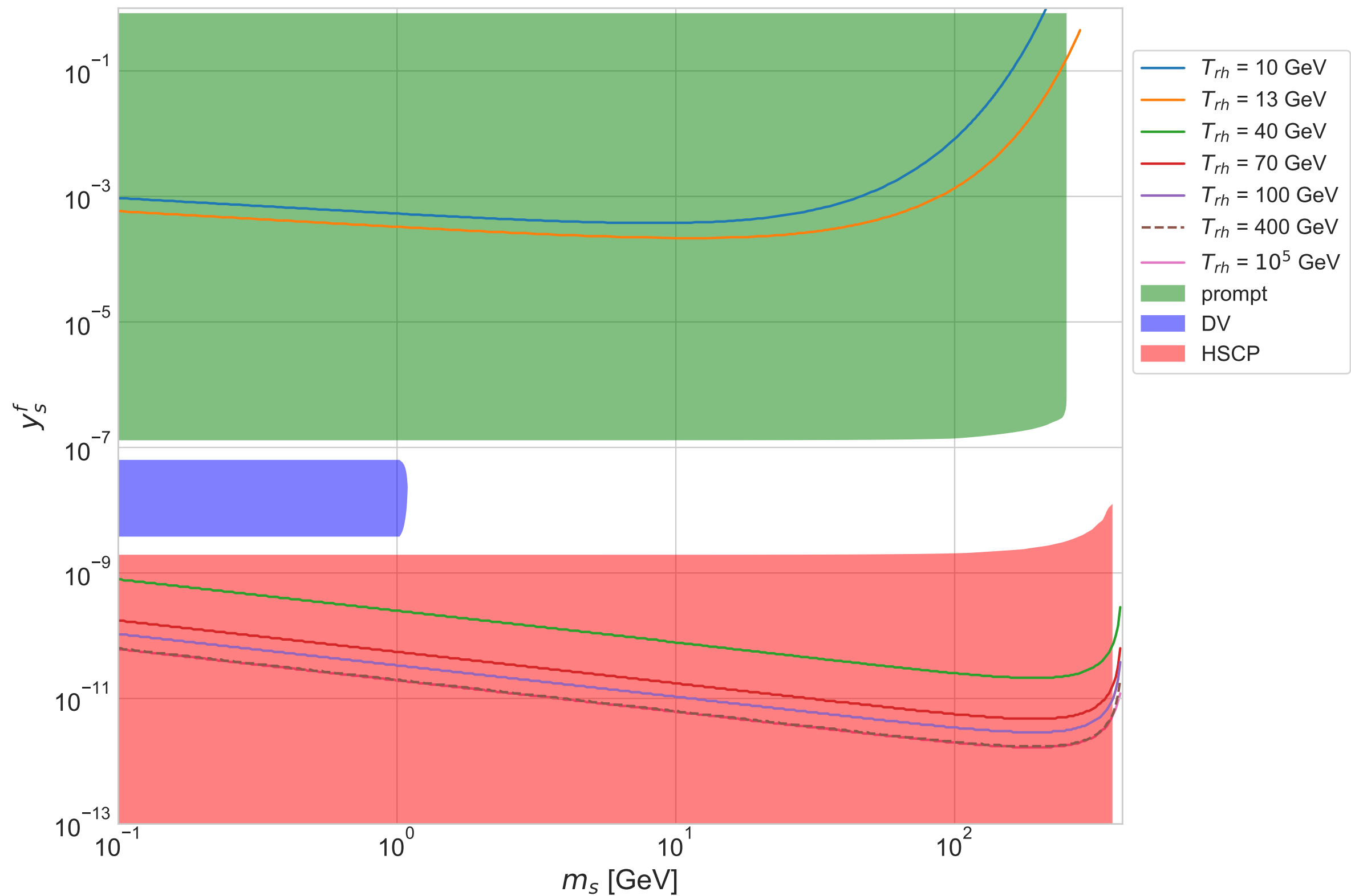


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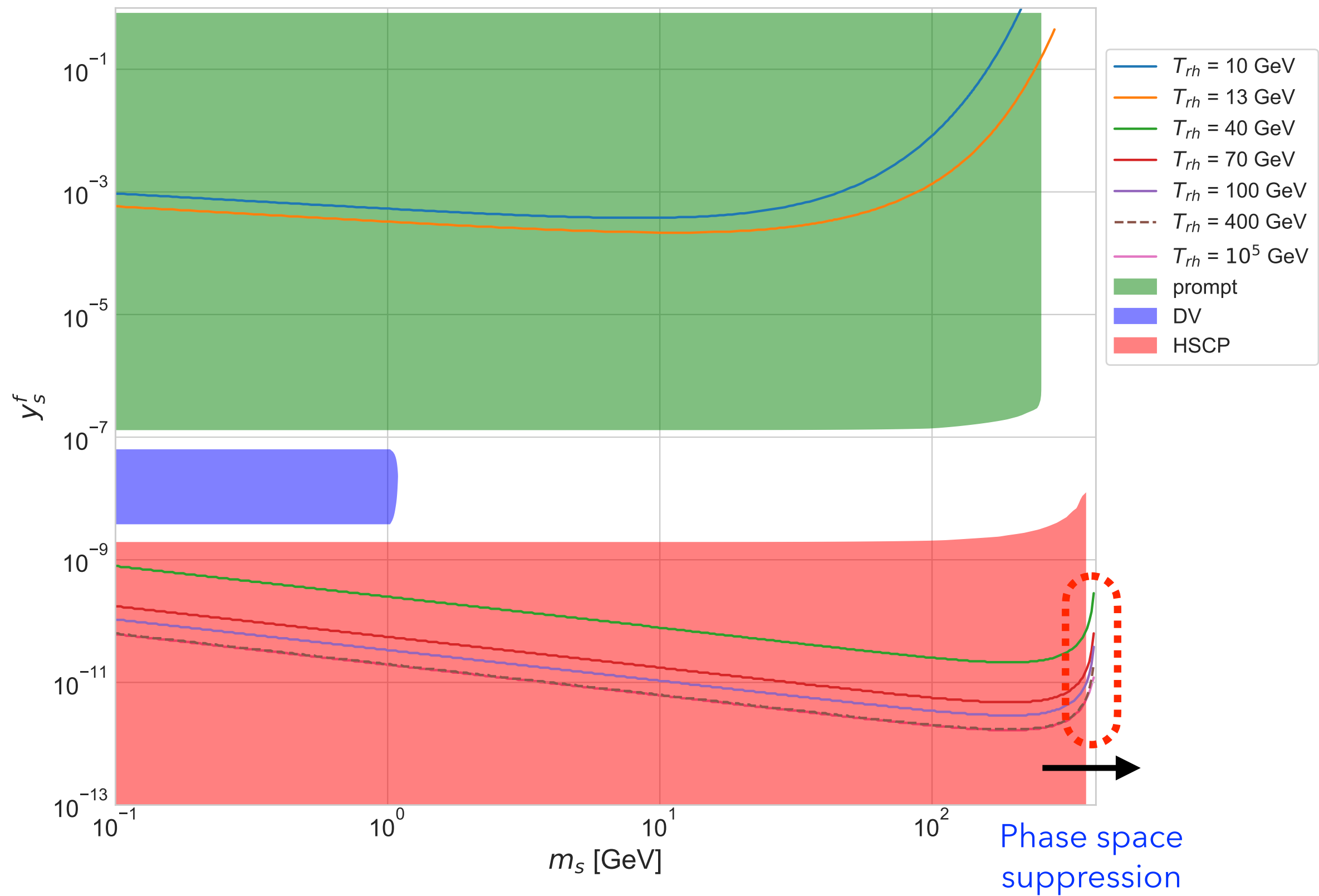
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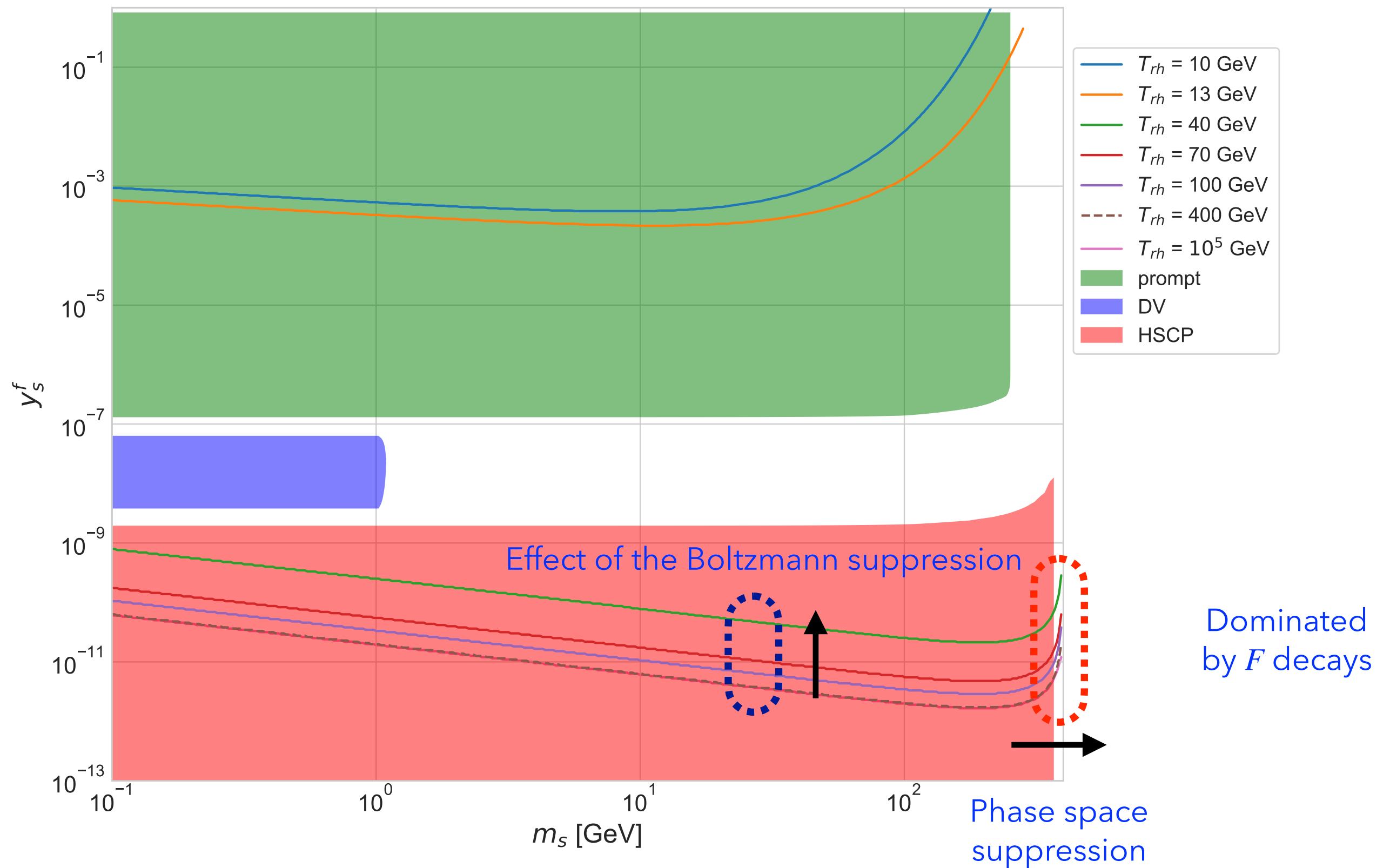
y_s^f vs m_s with $m_F = 400$ GeV



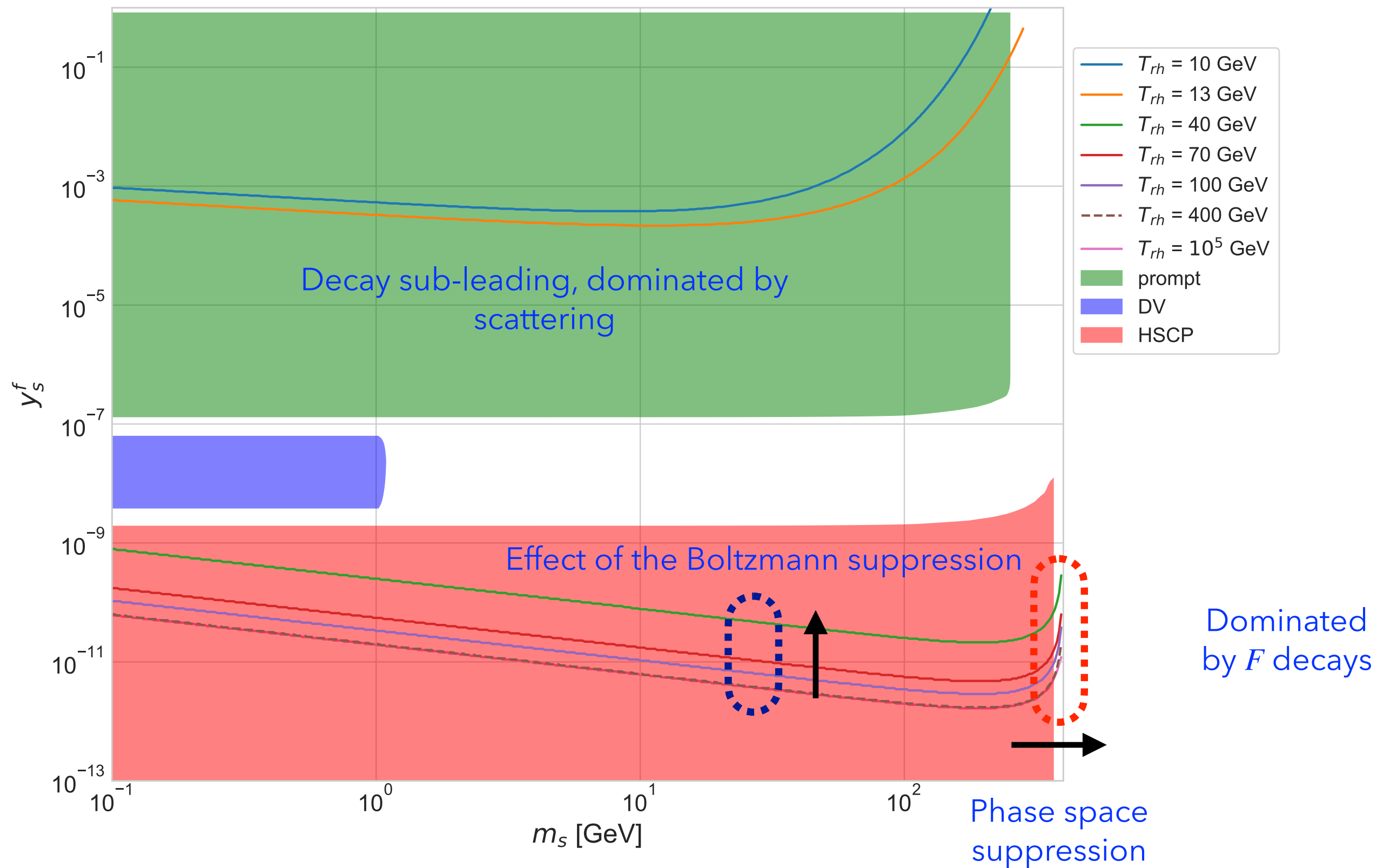
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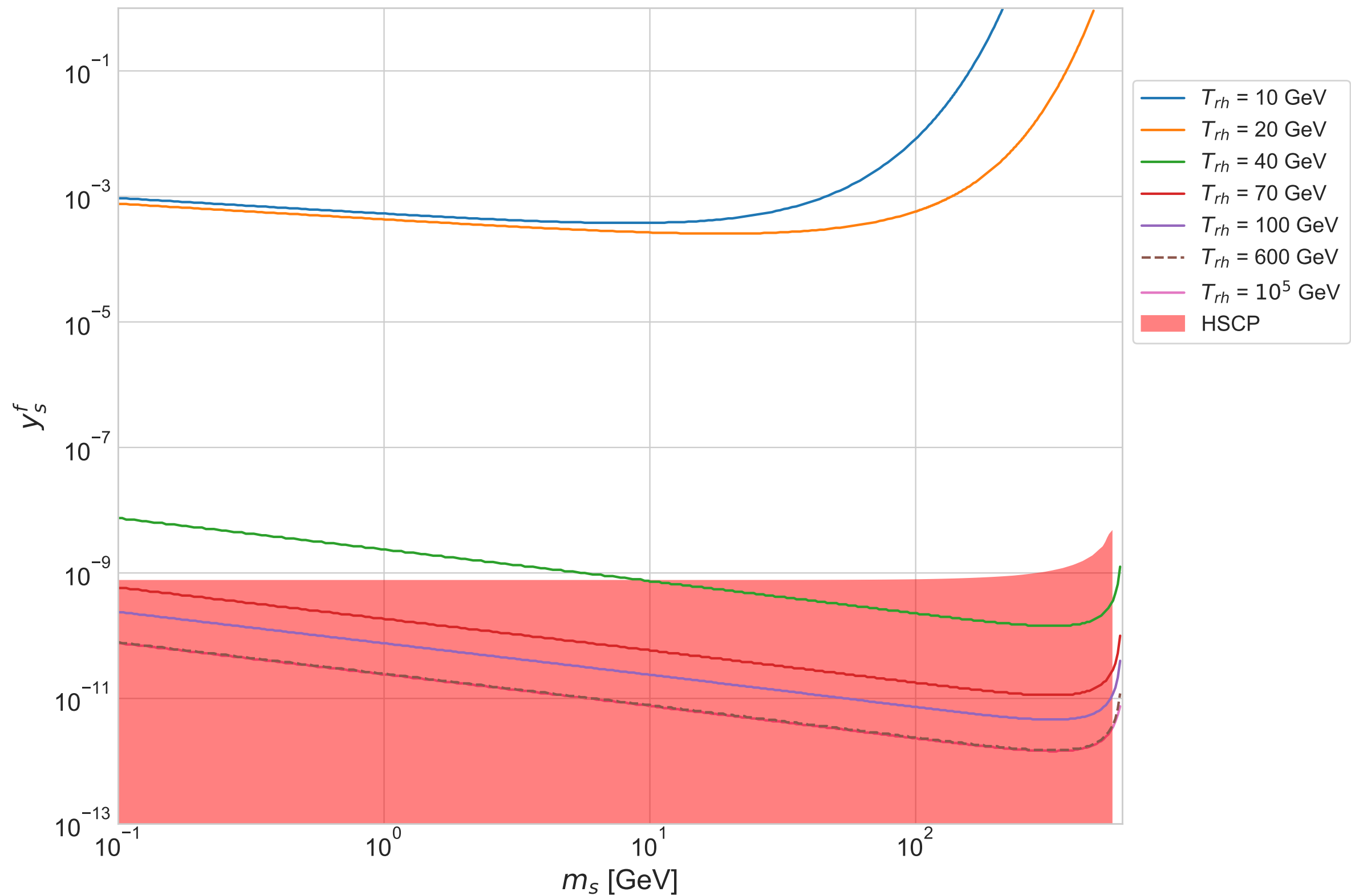
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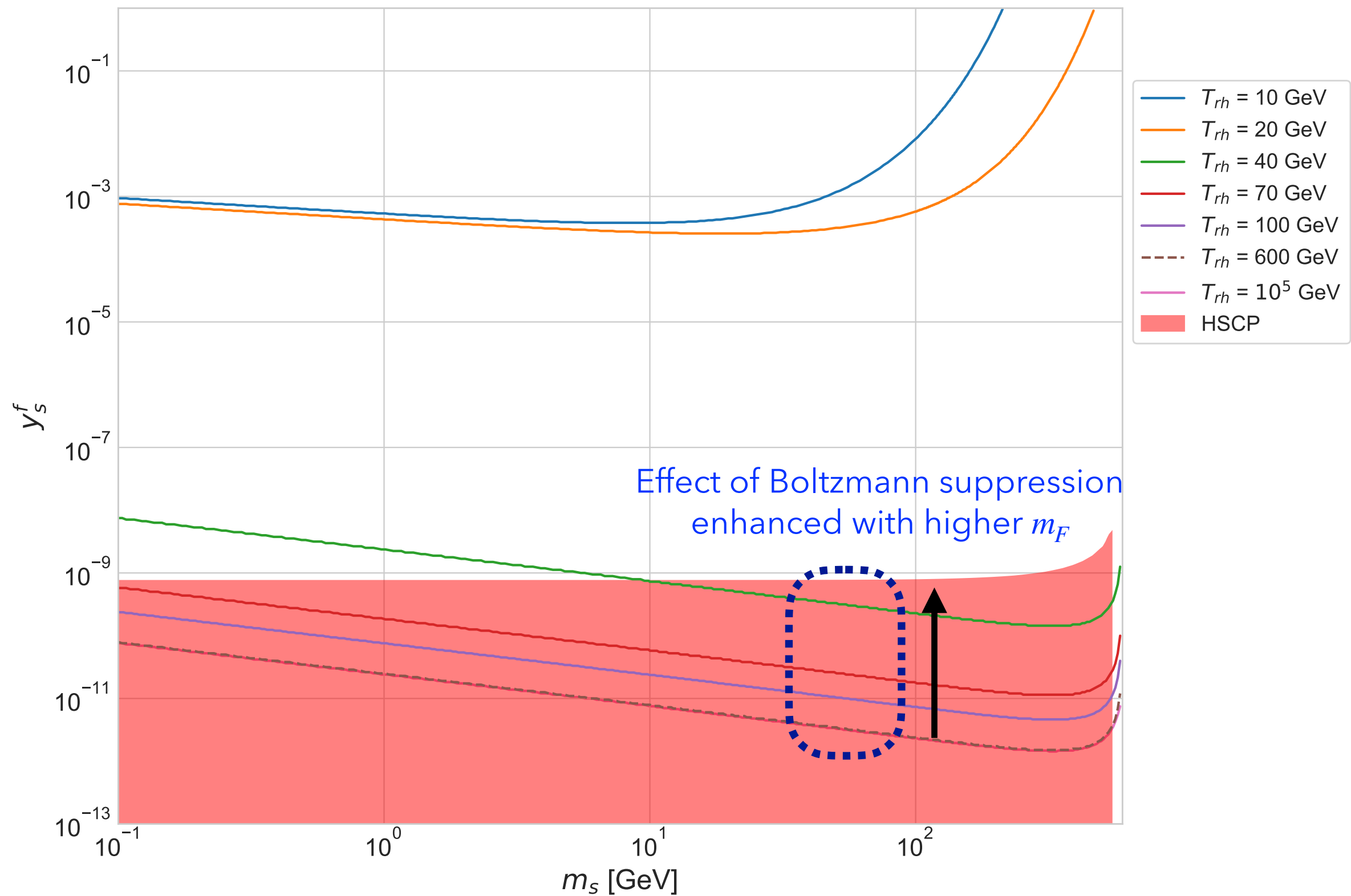
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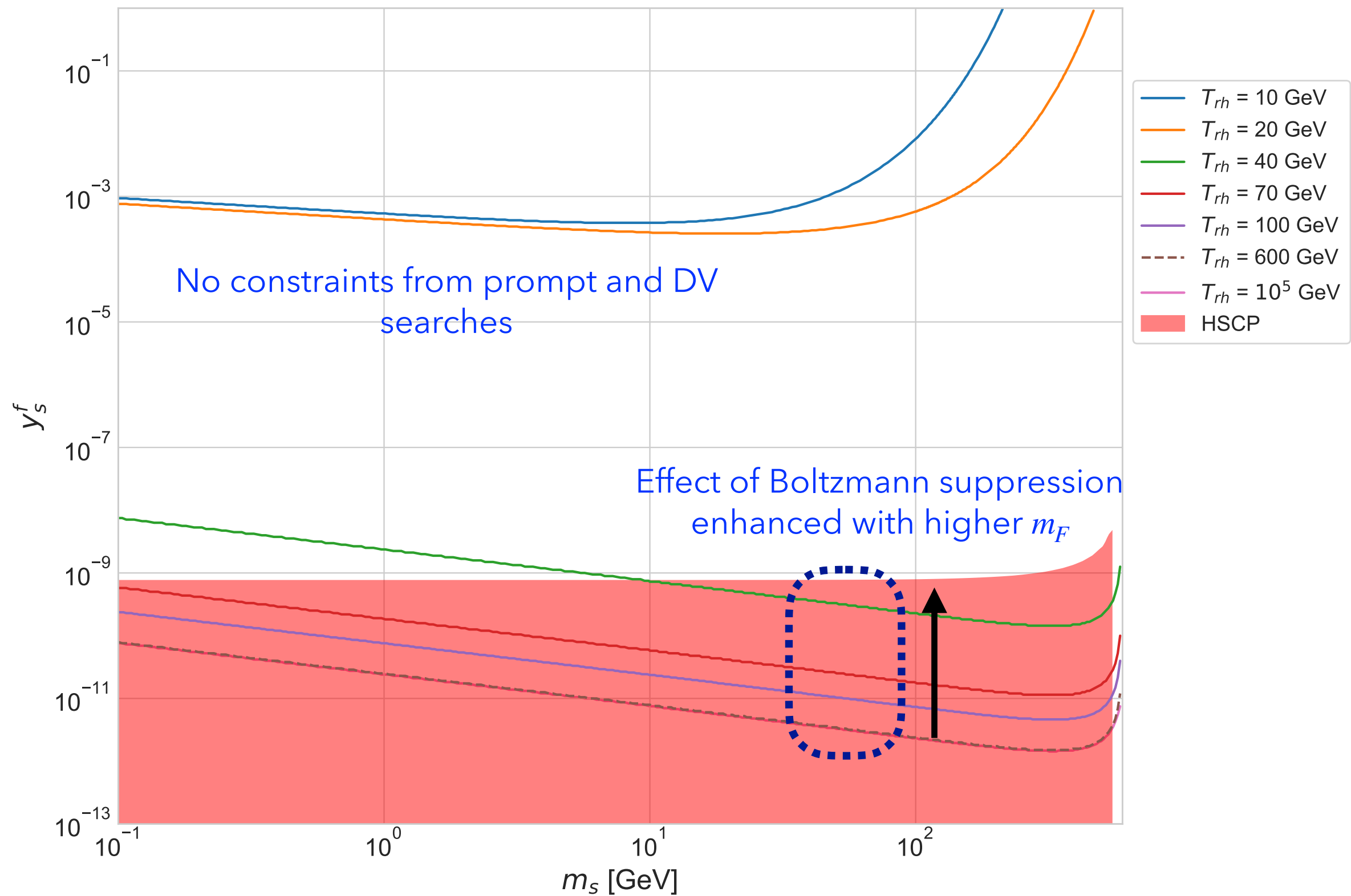
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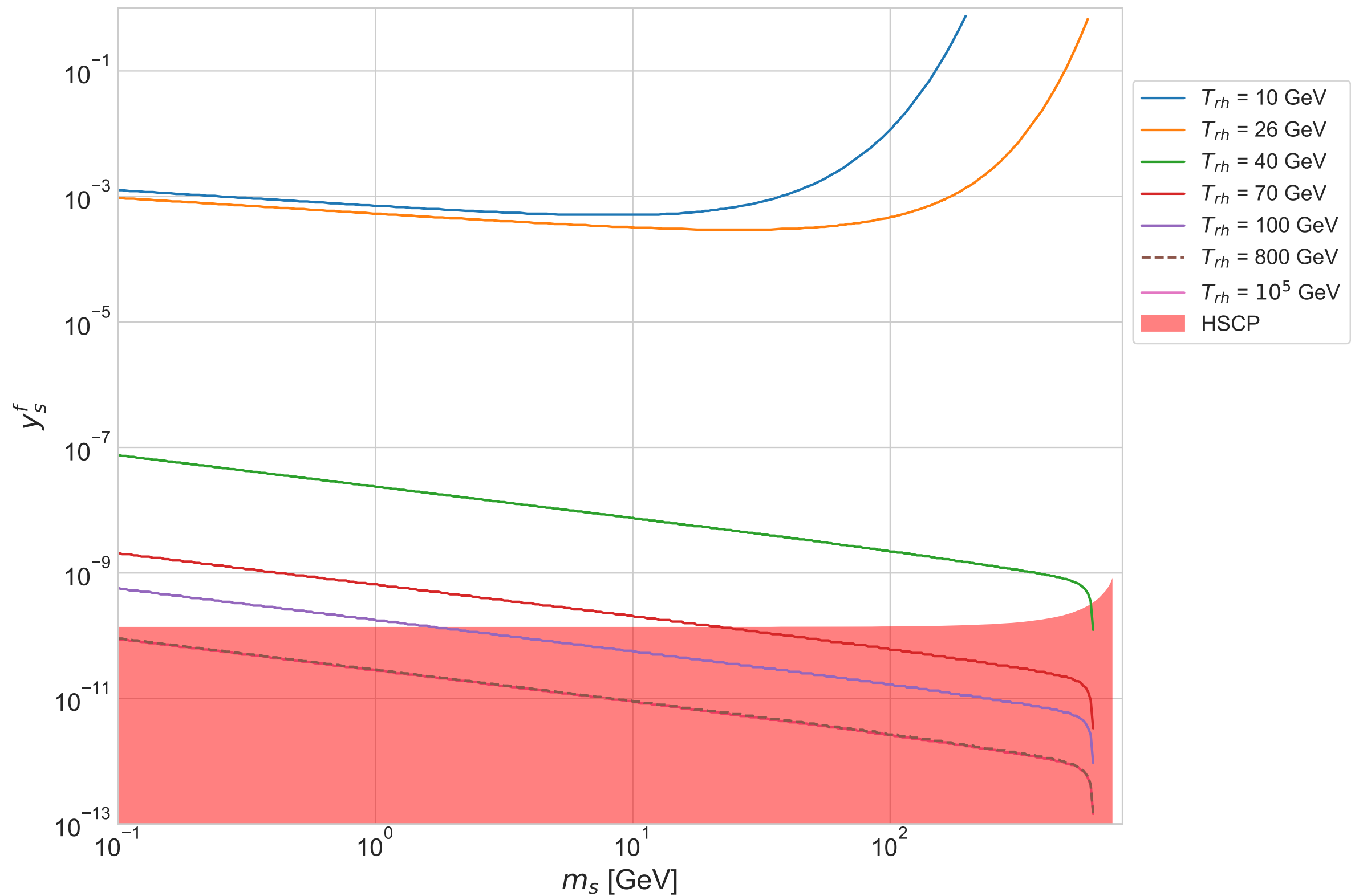
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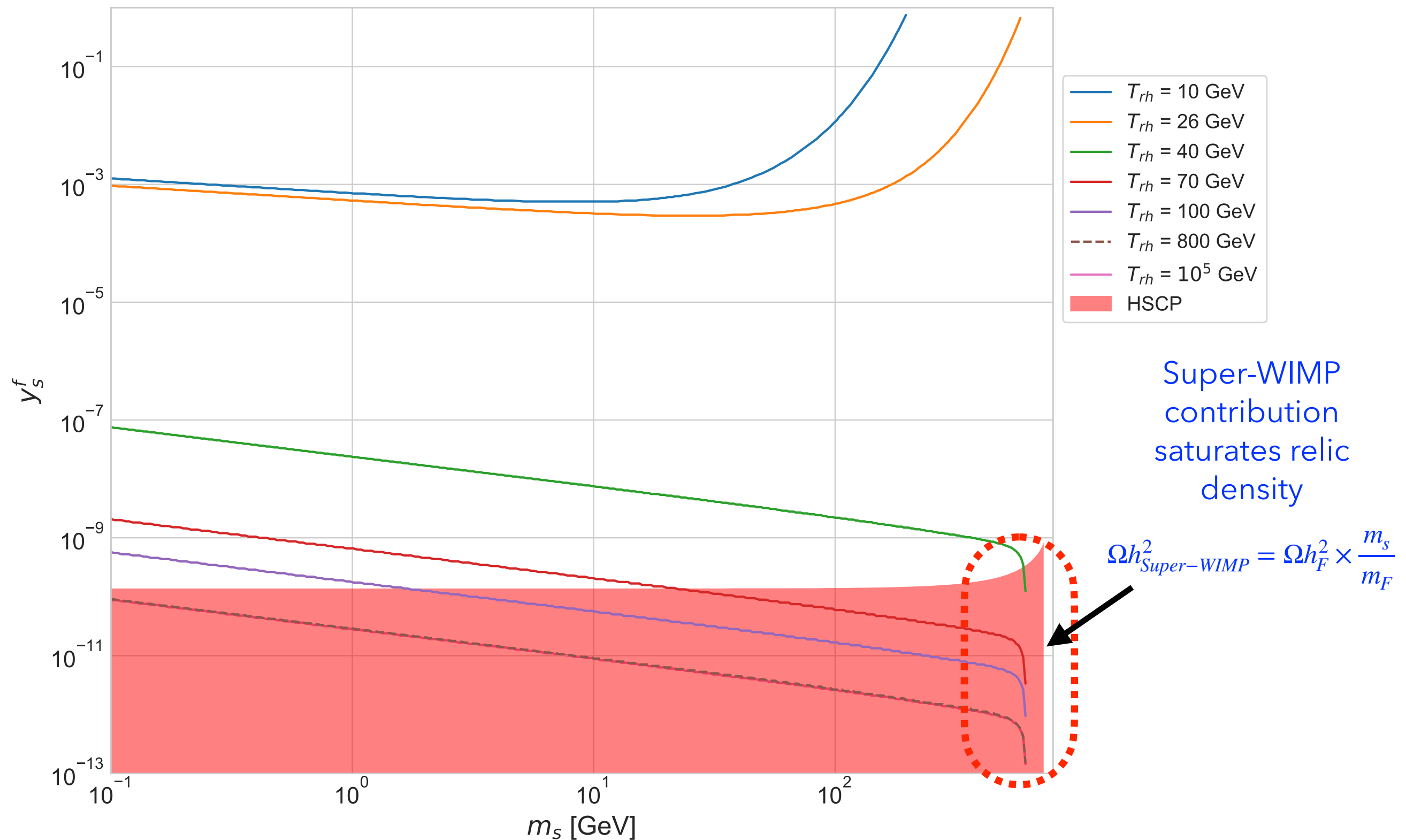
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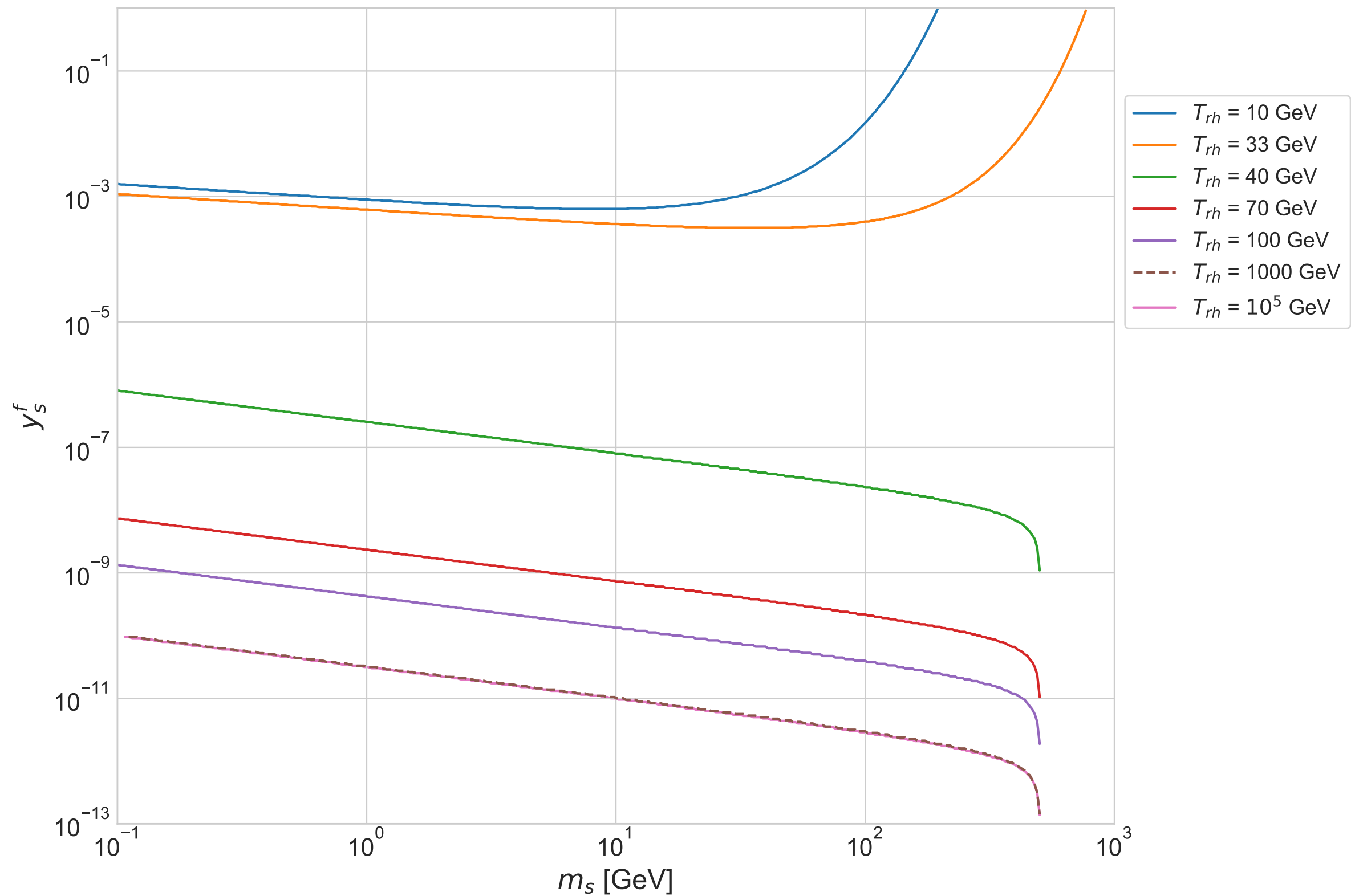
y_s^f vs m_s with $m_F = 800$ GeV



y_s^f vs m_s with $m_F = 800$ GeV



y_s^f vs m_s with $m_F = 1000$ GeV



Conclusion and outlook

- In the freeze-in mechanism, dark matter production depends on the reheating temperature which is poorly constrained.
- Usually taken to be effectively infinite, reducing its value leads to stronger coupling in order to reproduce the observed dark matter abundance in the universe
- Different contributions enter the predicted dark matter density (F decays, SM particle annihilation, Super-WIMP contribution) and become dominant in different regions of parameter space and for different reheating temperatures
 - ➔ The phenomenology of freeze-in models can be drastically modified
- Modified cosmological assumptions can lead to different phenomenological signatures for freeze-in models
- To appear : $\mu \rightarrow e\gamma$ + direct detection constraints, consider alternative coupling patterns