Entanglement in solvable quantum many-body systems

GILLES PAREZ

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A. Einstein

B. Podolsky

N. Rosen



A. Einstein





N. Rosen

EPR paradox (1935)







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$EPR \overline{paradox} (\overline{1935})$

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$$|\psi\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$







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Two systems **A** and **B** are entangled if performing a measurement on **A** influences the state of **B** independently from their physical separation.







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Quantum system in the state $|\psi^{AB}\rangle$ with subsystems A and B

Entanglement

$$\Leftrightarrow$$

$$|\psi^{\mathbf{A}\mathbf{B}}\rangle \neq |\phi^{\mathbf{A}}\rangle \otimes |\varphi^{\mathbf{B}}\rangle$$

No entanglement

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What is the difference between these two states?

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Answer:

- $|\psi^1\rangle$: A sees $|\uparrow\rangle \Rightarrow$ B sees $|\downarrow\rangle$: Entanglement
- $|\psi^2\rangle$: measure in **A** does not affect **B**: **No entanglement**

• Quantum physics

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 - ▶ Resource for quantum computation

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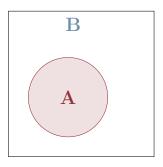
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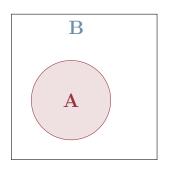
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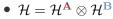
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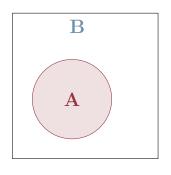
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- Relations to orthogonal polynomials, asymptotic analysis, ...

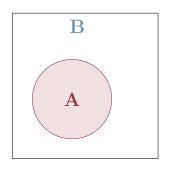




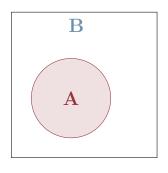




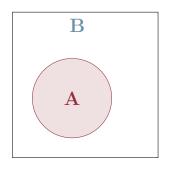
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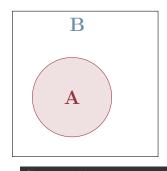
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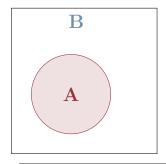
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Separability

Entanglement in many-body systems



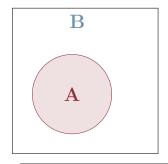
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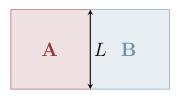
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A state is **entangled** if it is **not separable**: $S_A > 0$

Equilibrium and quantum phase transitions







Entanglement entropy and criticality



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• Logarithmic violation for 1 + 1d CFT



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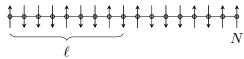
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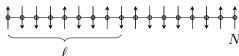
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$$S_{\mathbf{A}} = \frac{c}{3} \log \ell + \dots$$

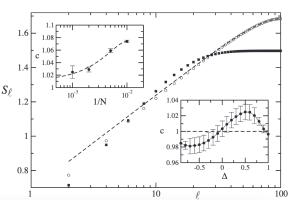
[De Chiara, Montangero, Calabrese, Fazio, 2006] $\underbrace{ \begin{pmatrix} \downarrow & \downarrow & \downarrow \\ \downarrow & \downarrow & \downarrow \\ \ell & \end{pmatrix} }_{\ell} \underbrace{ \begin{pmatrix} \downarrow & \downarrow \\ \downarrow & \downarrow \\ \end{pmatrix} }_{\ell} \underbrace{ \begin{pmatrix} \downarrow & \downarrow \\ \downarrow & \downarrow \\ \end{pmatrix} }_{N}$

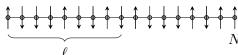


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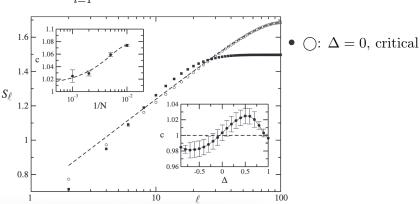


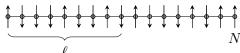
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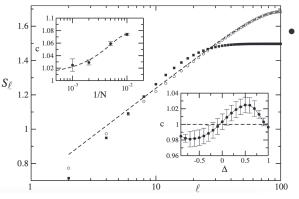


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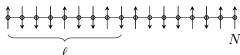
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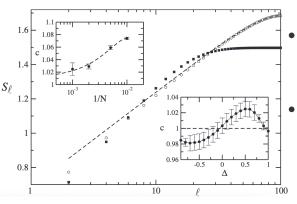
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: $\Delta = 0$, critical

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[De Chiara, Montangero, Calabrese, Fazio, 2006]



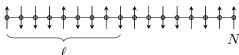
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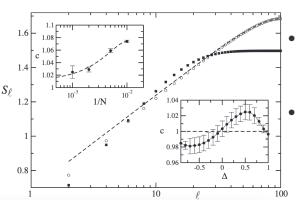
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$$S_{\ell} \sim \mathrm{cst}$$

Out of equilibrium and quantum quench

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Protocol to study an **isolated quantum system** out of equilibrium

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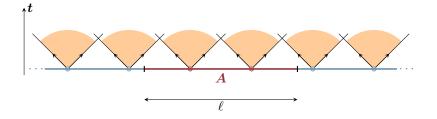
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Question: what happens between between t = 0 and $t = \infty$?

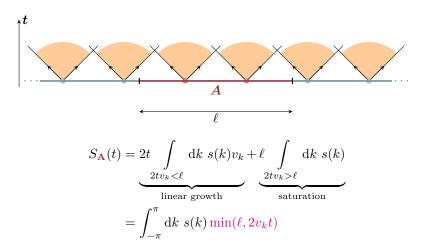
Quasiparticle picture: Single interval

[Calabrese, Cardy, 2005], [Alba, Calabrese, 2017-2018]



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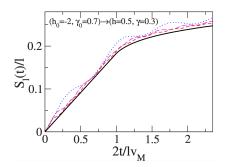
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Tests of the QPP

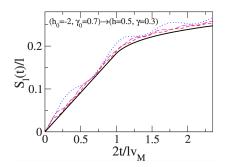
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Exact results for the XY spin chain [Fagotti, Calabrese, 2008]



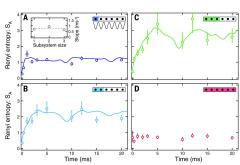
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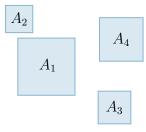


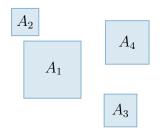
Experimental verification

[Kaufman et al., 2016]



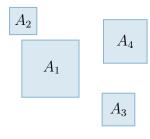
Multipartite setting





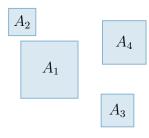
Physical states ρ satisfy

• $\operatorname{Tr}\rho = 1$



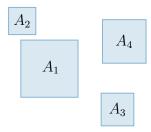
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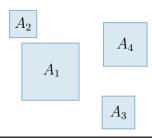
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Separability



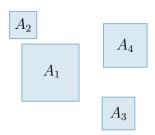
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Separability

The most classical state for $A_1 \cup A_2 \cup \cdots \cup A_m$ is a product,

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Separable states are convex combinations of product states,

$$\rho_{\text{sep}} = \sum_{k} p_k \ \rho_{A_1}^{(k)} \otimes \rho_{A_2}^{(k)} \otimes \cdots \otimes \rho_{A_m}^{(k)}$$

with $p_k \geqslant 0$ and $\sum_k p_k = 1$

Let us consider the 4-spin state

$$|GHZ\rangle = \frac{1}{\sqrt{2}}(|\uparrow_1\uparrow_2\uparrow_3\uparrow_4\rangle + |\downarrow_1\downarrow_2\downarrow_3\downarrow_4\rangle)$$

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The 3-spin density matrix is

$$\rho_{123} = \operatorname{Tr}_{4} |GHZ\rangle\langle GHZ| = \frac{1}{2} |\uparrow_{1}\uparrow_{2}\uparrow_{3}\rangle\langle\uparrow_{1}\uparrow_{2}\uparrow_{3}| + \frac{1}{2} |\downarrow_{1}\downarrow_{2}\downarrow_{3}\rangle\langle\downarrow_{1}\downarrow_{2}\downarrow_{3}|$$
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Multipartite entanglement **structure** in many-body systems is non-trivial. It is important to go **beyond the bipartite case!**

[GP, Witczak-Krempa, 2024]

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Let us consider

• A state $\rho(s)$ with an evolution parameter s

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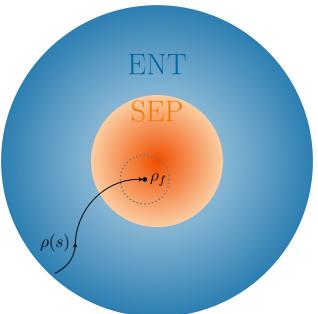
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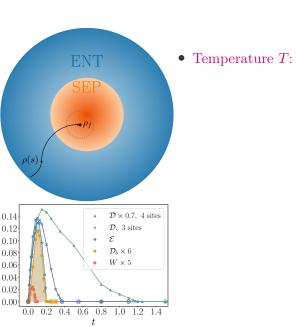
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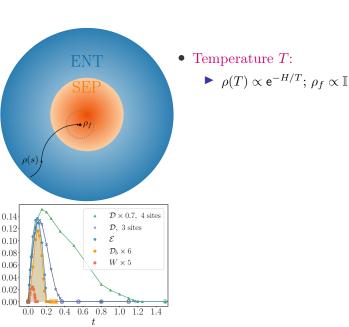
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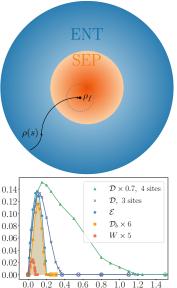
What is the **fate of entanglement** during such evolutions?

Separable continent

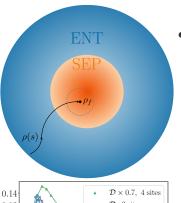




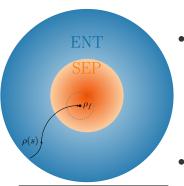


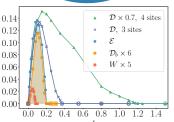


- Temperature T:
 - $ho(T) \propto \mathrm{e}^{-H/T}; \, \rho_f \propto \mathbb{I}$
 - ▶ Existence of a temperature sudden death

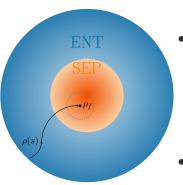


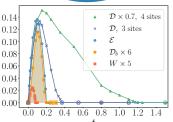
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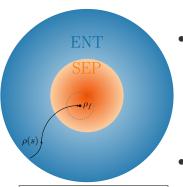


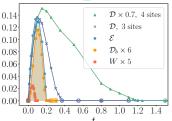
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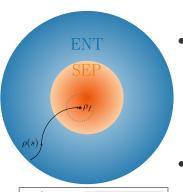


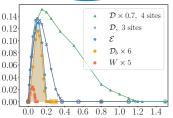
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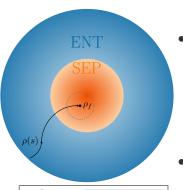


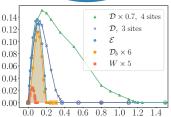
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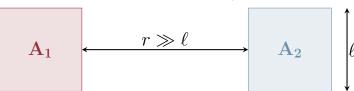


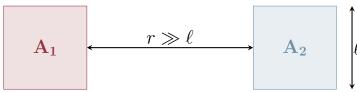
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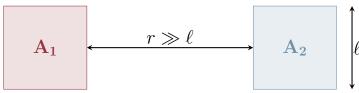
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- ► Rise and fall behavior



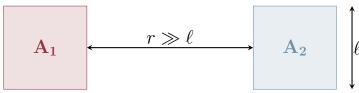


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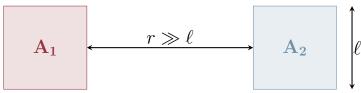


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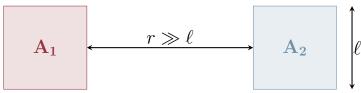
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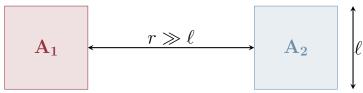
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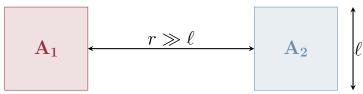
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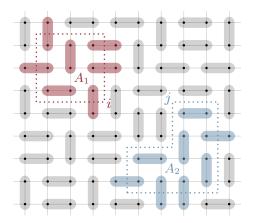
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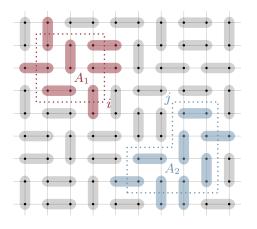
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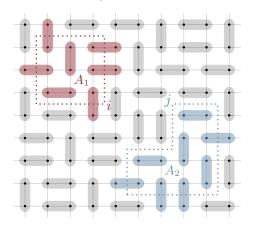
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Long-range correlations do not imply long-range entanglement!

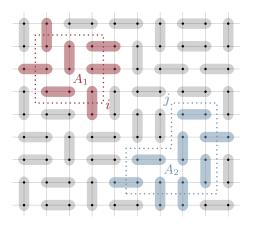




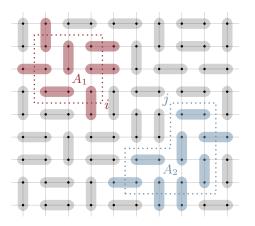
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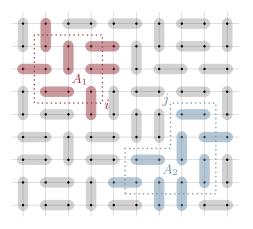
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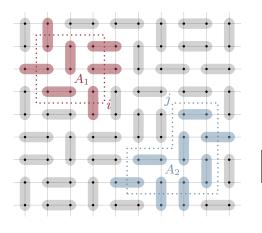
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States built from your favourite statistical model [Rokhsar, Kivelson, 1988]



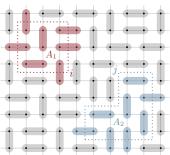
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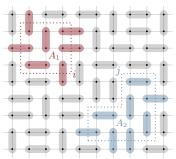
In the following, we focus on dimer RK states

$$E(c) = 0$$
, $Z = \#$ dimer configurations

[GP, Berthiere, Witczak-Krempa, 2023]

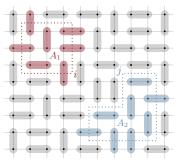


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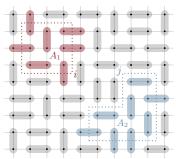
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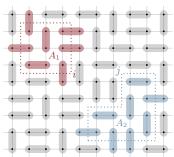
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The matrix $\rho_{A_1 \cup A_2}$ is separable! (for disjoints A_1 and A_2)

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