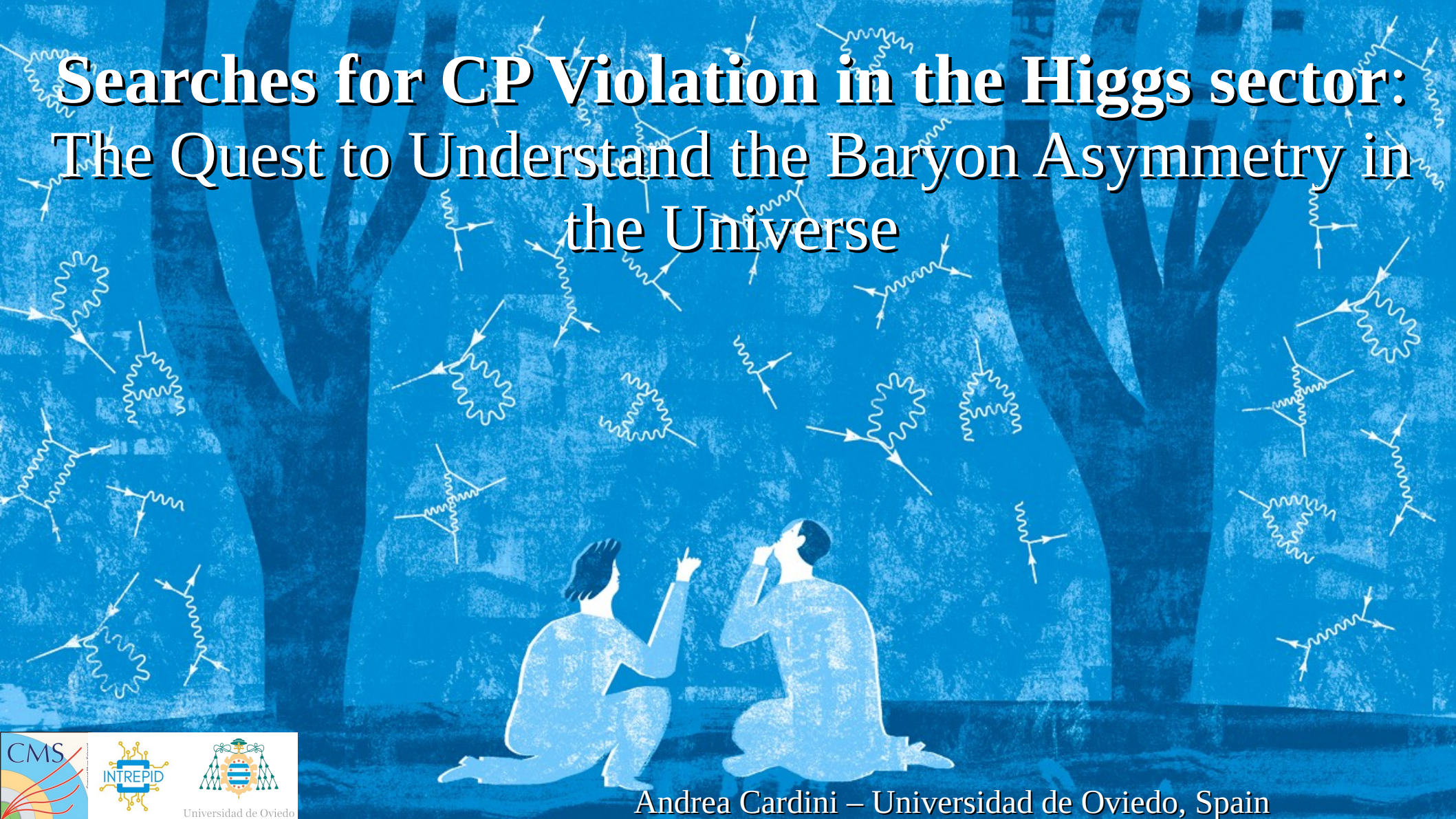


Searches for CP Violation in the Higgs sector: The Quest to Understand the Baryon Asymmetry in the Universe

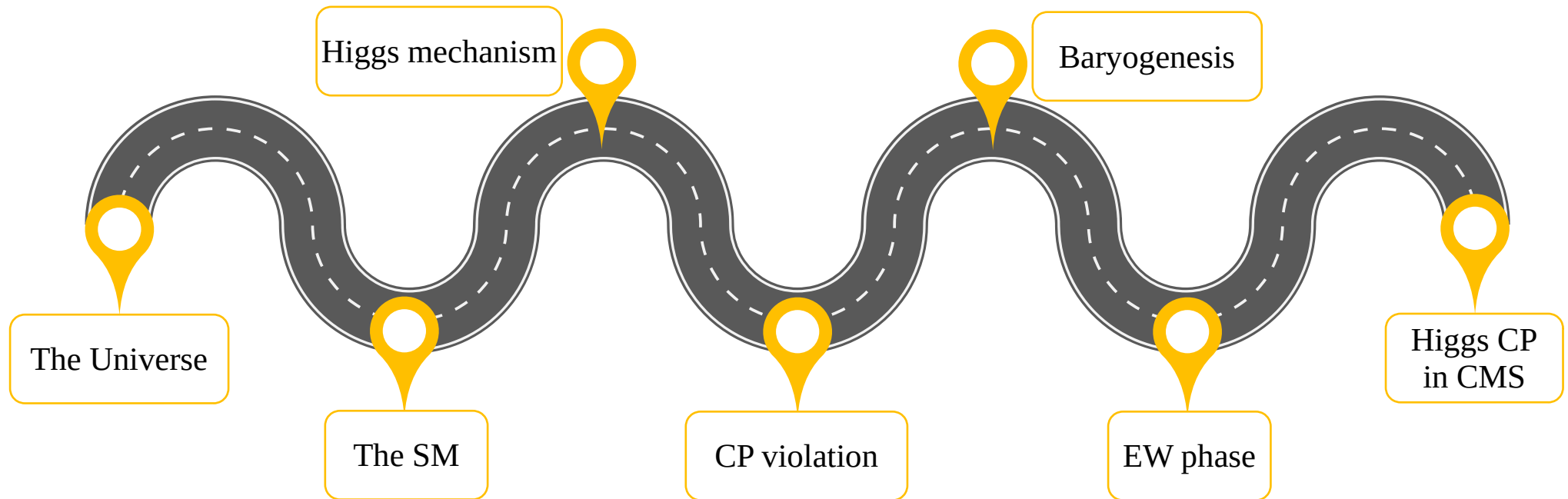




Road of the talk



- The aim:
 - > general understanding of the mystery of absence of antimatter in the Universe
 - > review of CP violation results in CMS that could help explain it

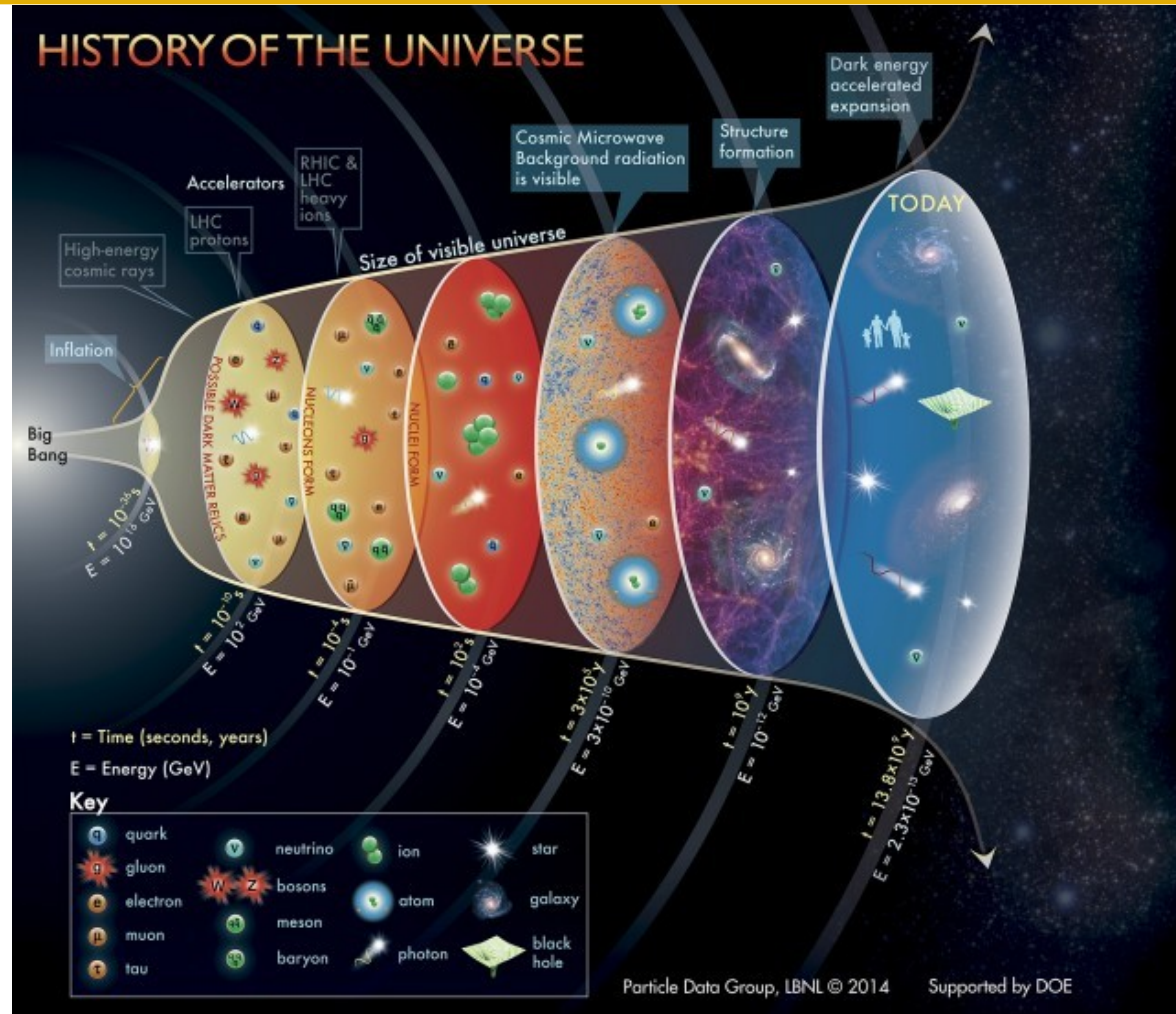




Once upon a time...



- the Universe was not like we see it today:
 - > No people
 - > No Earth
 - > No Galaxies
 - > and it was much much *smaller*
- The primordial Universe contained both **matter** and **anti-matter**

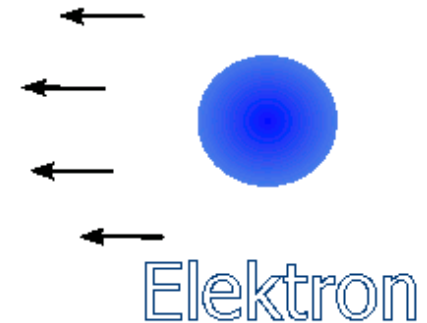
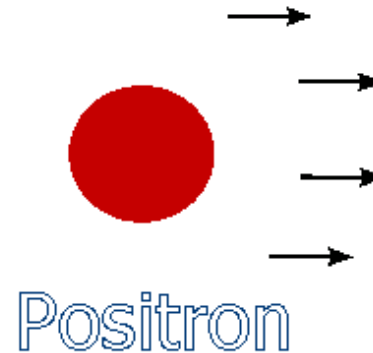




Once upon a time...



- the Universe was not like we see it today:
 - > No people
 - > No Earth
 - > No Galaxies
 - > and it was much much *smaller*
- The primordial Universe contained both **matter** and **anti-matter**
- Energy was converting into matter and anti-matter which then annihilated to create again energy
- This cycle was broken, and the Universe now is made just of matter
- Where did the anti-matter go?



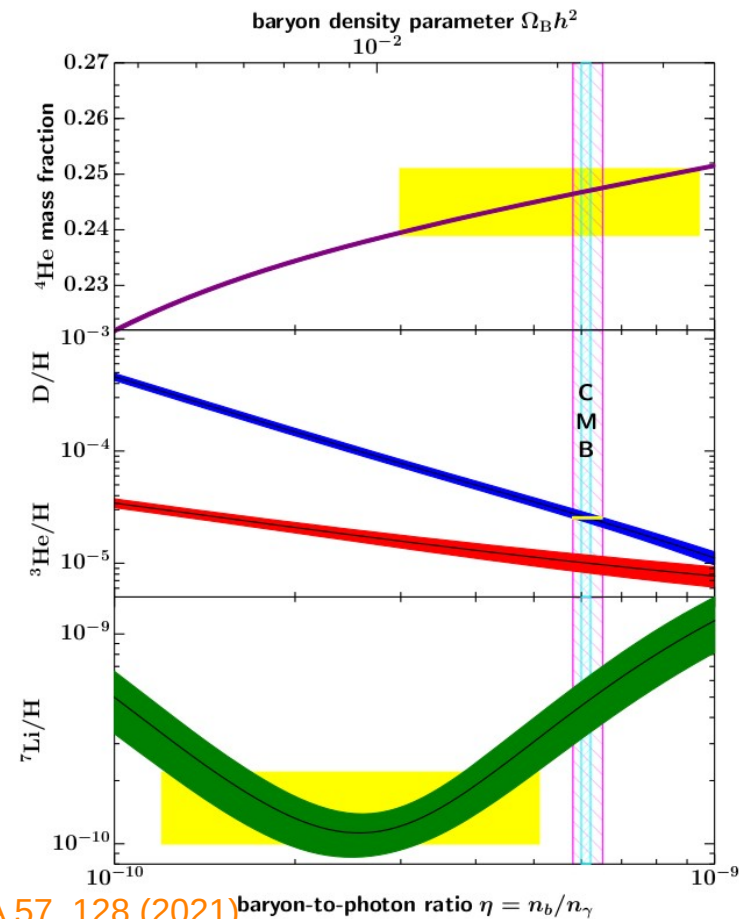


The Universe



- Our Universe is now made of matter, forming the baryons and leptons which make up the atoms and molecules we are made of
- This required an **asymmetry** between matter and anti-matter in the primordial Universe
 - > **Baryon Asymmetry in the Universe (BAU)**
 - > Matter : Antimatter
1'000'000'001 : 1'000'000'000
- More specifically, the asymmetry is encoded in the ratio of baryons over photons in the Universe
- From the cosmic microwave background this is

$$\eta = \frac{(n_B - n_{\bar{B}})}{n_\gamma} = \frac{n_{B,relic}}{n_\gamma} \approx 6 \times 10^{-10}$$



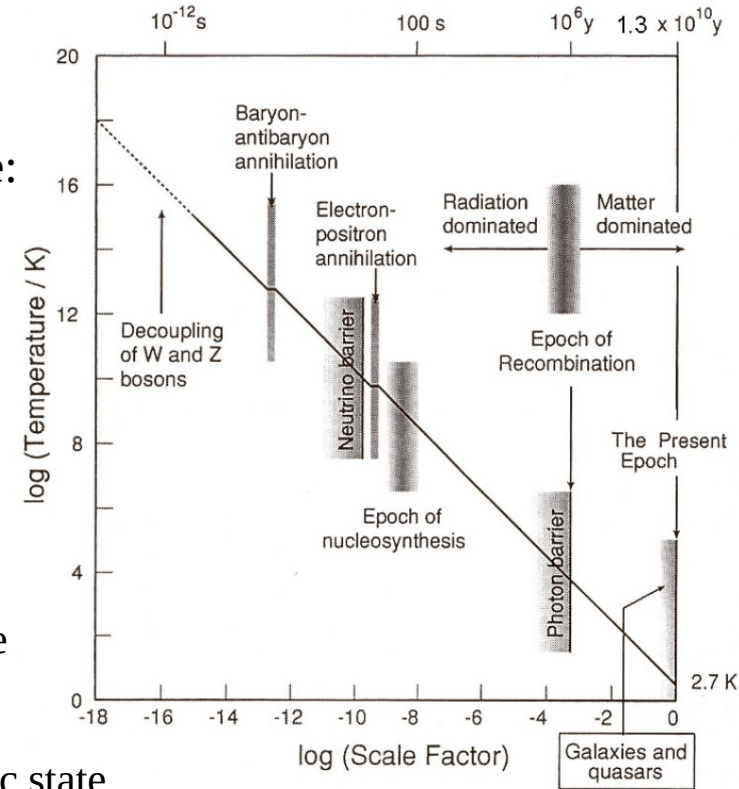
Eur. Phys. J. A 57, 128 (2021) baryon-to-photon ratio $\eta = n_b/n_\gamma$



Evolution of the Universe



- The cosmic microwave background originated when the Universe cooled off enough that the photons were free to propagate
 - > Epoch or recombination: formation of atoms
- Going backward in time particle density and temperature increase:
 - > formation of nuclei
 - > cosmic neutrino background
 - > **the baryon-antibaryon annihilation** → the matter asymmetry should have been present before this moment
- Sakharov theorized 3 conditions **necessary** for a CPT conserving theory to explain the BAU:
 - > Baryon number violation: generate asymmetry in a symmetric Universe
 - > C and CP violation: asymmetric interaction to remove antimatter
 - > Deviation from thermal equilibrium: to prevent reverting to a symmetric state
- Is there a theory that satisfies all 3 conditions?





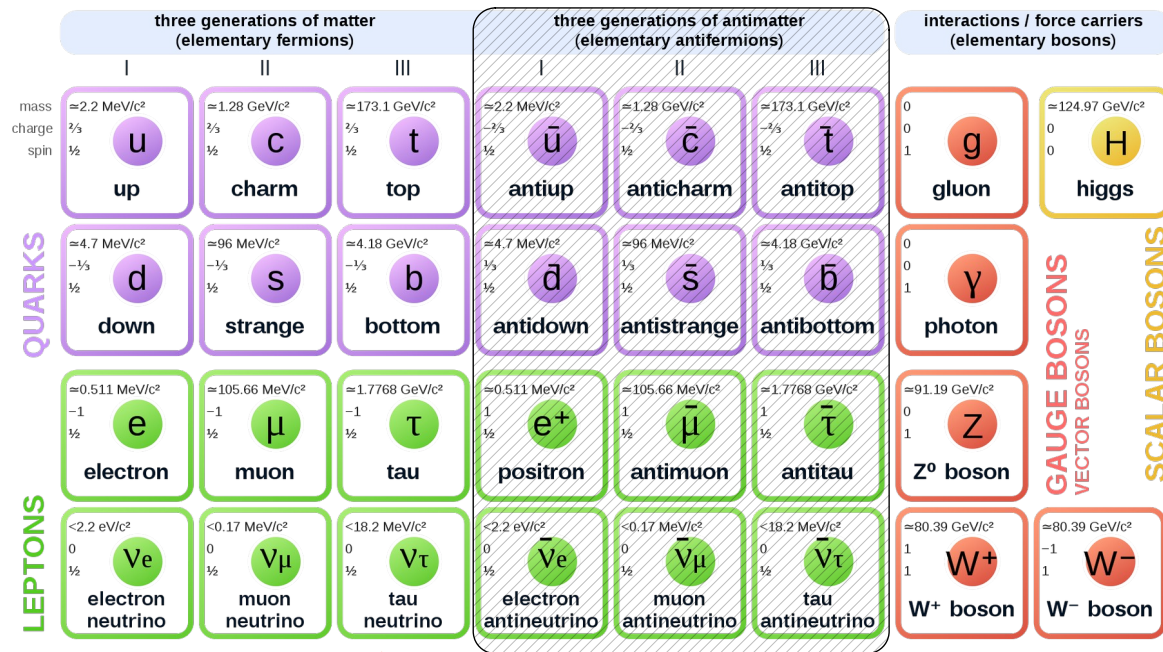
The Standard Model of particle physics



- Successful theory predicting the behavior of matter constituents (**fermions**)
- Interactions mediated by **gauge bosons**, gravity not included

- Relevant symmetries:

Standard Model of Elementary Particles



> Gauge: $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$

> Parity (P)

> Charge conjugation (C)



CP

> Time reversal (T)

> Chirality: χ_L, χ_R

- The decoupling of the weak and electromagnetic forces is a key moment in the history of the Universe

C

Inversion of additive quantum numbers



Electroweak Symmetry Breaking



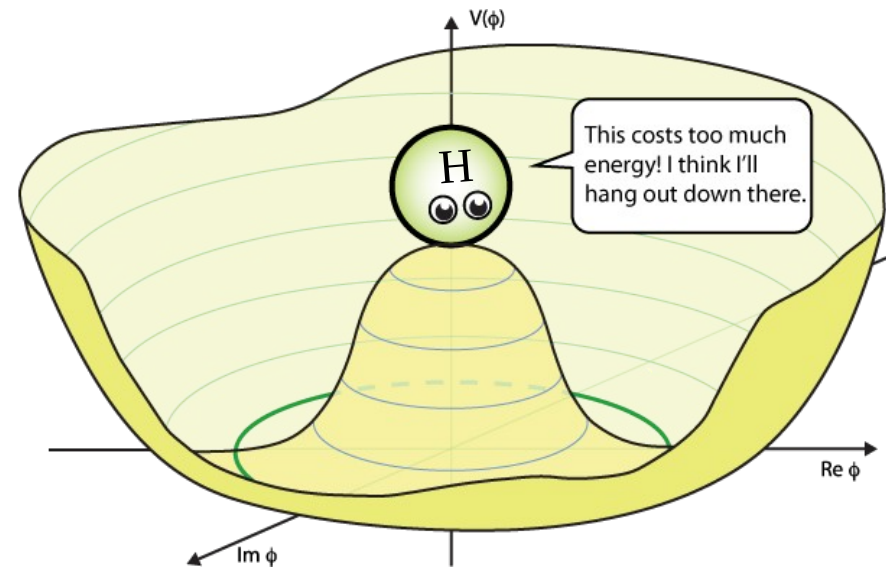
- The gauge invariance of the SM would prevent mass terms for the vector bosons and the fermions
- The Higgs boson was introduced to softly break the electroweak symmetry and give mass to the particles
- EWSB: $SU(2)_L \otimes U(1)_Y \rightarrow U(1)_{em}$

- > Gauge bosons gain mass directly through Higgs kinetic term

$$L_{mWZ} = \left(\frac{v^2 g^2}{4} W_\mu^\dagger W^\mu + \frac{v^2 (g^2 + g'^2)}{8} Z_\mu Z^\mu \right)$$

- > Fermions gain mass via **Yukawa coupling**

$$L_Y^f = - \underbrace{m_f (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)}_{\text{fermionic mass term}} - \underbrace{\frac{m_v}{v} (\bar{\psi}_L H \psi_R + \bar{\psi}_R H \psi_L)}_{\text{coupling of the Higgs to fermions}}$$



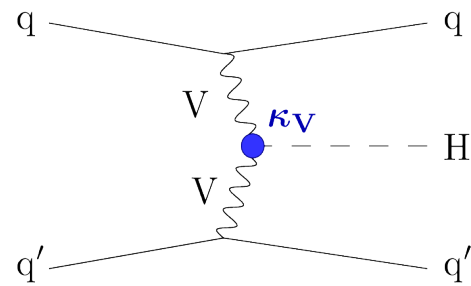
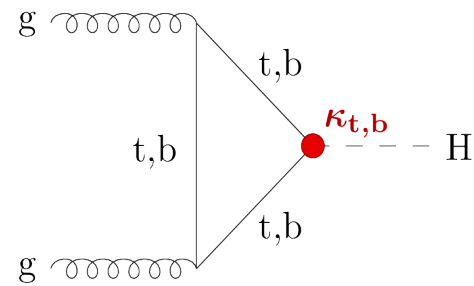
- Fermions form the matter our Universe is made of → Look for CP violation in the Higgs sector



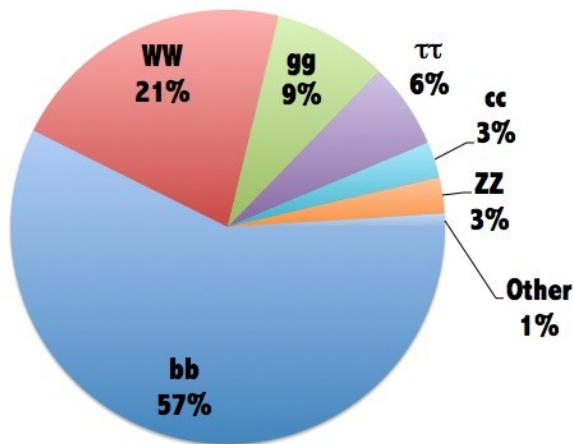
The Higgs boson – short recap



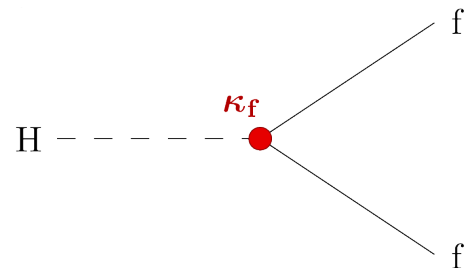
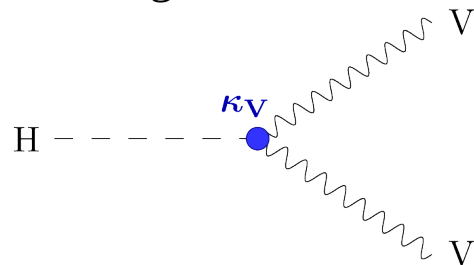
- At the LHC the main way to produce the Higgs boson is via gluon-gluon fusion
- Production via vector boson fusion, Higgs-strahlung and in association with top or bottom quarks are rarer



Higgs decays at $m_H=125\text{GeV}$



- The Higgs boson is not stable and decays immediately
- It can decay directly to W and Z bosons, and to all fermions via Yukawa coupling
- Via intermediate fermion or boson loop it can also decay to photons and gluons

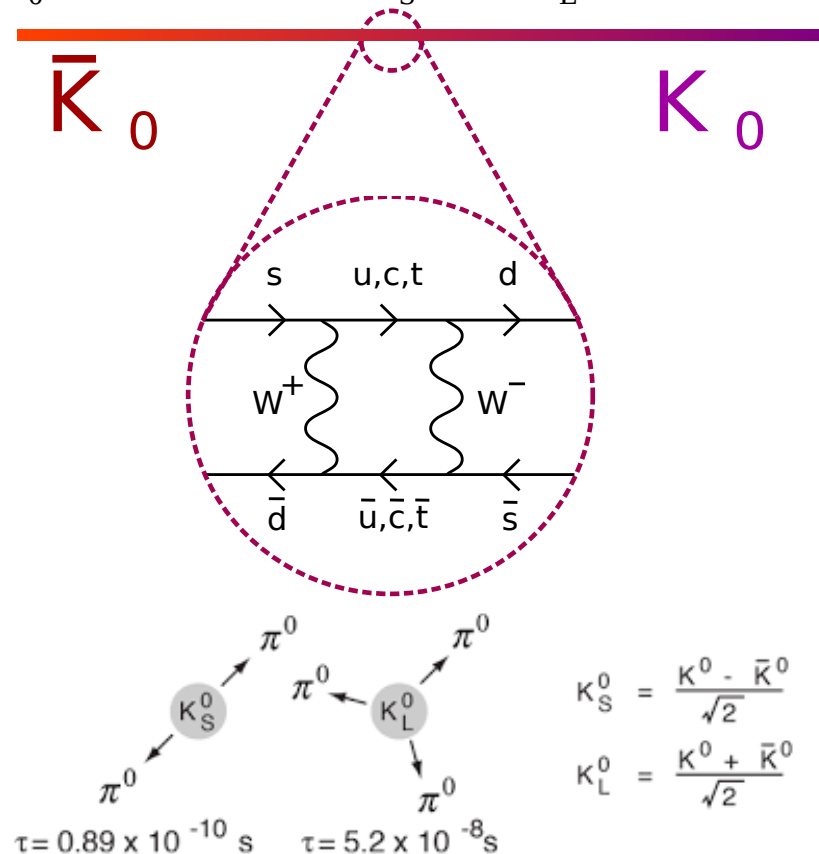




CP violation



- Weak interactions massively violates C and P symmetries, it also violates CP:
 K_0 oscillation $\rightarrow K_S$ and K_L decays – Cronin and Fitch in the 60s



- The most notable violation of C, P and CP in the SM comes from flavor changing currents in the weak interaction
- CP violation located in the complex phase in the Cabibbo-Kobayashi-Maskawa matrix
- The matrix was introduced to describe the quark mixing in the weak interaction



The CKM matrix and Yukawa interaction

- The quark mixing and the quark masses have a common origin: the Yukawa interaction
- The Higgs boson gives mass to the particles and defines the fermion generations

$$\begin{aligned} \mathcal{L}_{\phi q} = & - \sum_{ij} G_{ij} \begin{pmatrix} \bar{U}_{iL} & \bar{D}_{iL} \end{pmatrix} \cdot \begin{pmatrix} \phi^{0\dagger} \\ -\phi^- \end{pmatrix} U_{jR} & \longrightarrow & \text{Yukawa for up-type quarks} \\ & - \sum_{ij} H_{ij} \begin{pmatrix} \bar{U}_{iL} & \bar{D}_{iL} \end{pmatrix} \cdot \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} D_{jR} + H.c. & \longrightarrow & \text{Yukawa for down-type quarks} \end{aligned}$$

with the indices i, j running over the quark generations

- After EWSB the quarks acquire mass via the Higgs vacuum expectation value

$$\mathcal{L}_m = - \sum_{ij} \bar{U}_{iL} m_{ij}^U U_{jR} - \sum_{ij} \bar{D}_{iL} m_{ij}^D D_{jR} + H.c.$$

- No theoretical requirement for the m_{ij} matrices to be diagonal **but** the SM Lagrangian is invariant under unitarity transformations



The CKM matrix



- We can multiply all fields by unitary matrices and redefine the quark fields to diagonalize the mass matrices

$$\begin{array}{ll} m_{ij}^U = \frac{1}{\sqrt{2}} G_{ij} v & m^{U'} = A_L^U m^U A_R^{U\dagger} \\ m_{ij}^D = \frac{1}{\sqrt{2}} H_{ij} v & m^{D'} = A_L^D m^D A_R^{D\dagger} \end{array} \quad \longrightarrow$$

- Cannot diagonalize both mass matrices at once $\rightarrow$ diagonalize one mass matrix and redefines the fields for either up or down type quarks
- The lightest quarks are **u**, **d** and **s** $\rightarrow$ the quark mixing was observed for kaon decays between the **d** and **s** quarks $\rightarrow$ the mass term diagonalized is the one for up-type quarks

$$V_{CKM} = A_L^U (A_L^D)^{-1} \quad \longrightarrow \quad q_{uL} = \left(\frac{1 - \gamma_5}{2} \right) \left(V_{ud} d + V_{us} s + V_{ub} b \right)$$



CP violation in the SM



- The CKM matrix is unitary **but** is not required to be Real $\rightarrow$ it can have a complex phase leading to direct CP violation in the electroweak sector

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

- The complex phase $\delta = 1.196 \pm 0.045$ [rad] is the main source of CP violation allowed in the SM
- The strong interaction does not seem to violate CP experimentally... however in QCD CP-violating terms can arise leading to the question of why no signs of CP violation have been observed
- Another source of CP violation *could** be found in the neutrino oscillation if the PMNS matrix presents a complex phase

**there is no experimental consensus on a CP-violating phase in the PMNS matrix*

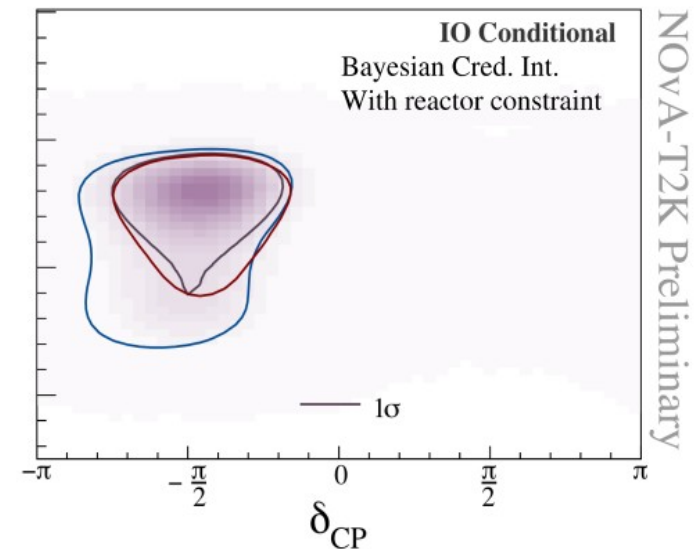
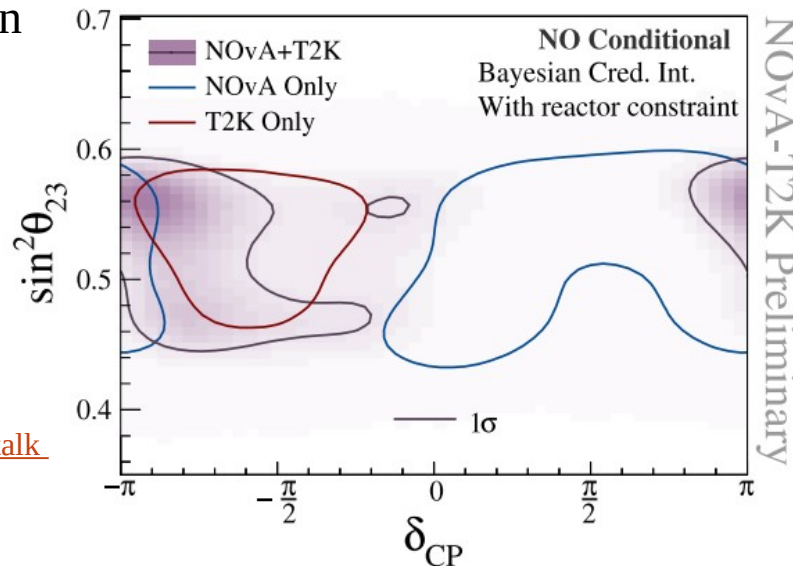


NovA and T2K results on the PMNS CP-phase



- The CP-phase in the PMNS matrix is investigated studying the neutrino oscillation process
- Most precise results at hand from the NOvA and the T2K experiments
- The measurement constrains the CP-phase together with the mass differences between the neutrino mass states
- At present the two experiments prefer different CP-phases in the normal mass ordering (NO), and maximum CP violation in the inverted mass ordering (IO)

Plots from
[J.P. Ochoa Ricoux's talk at ICHEP2024](#)

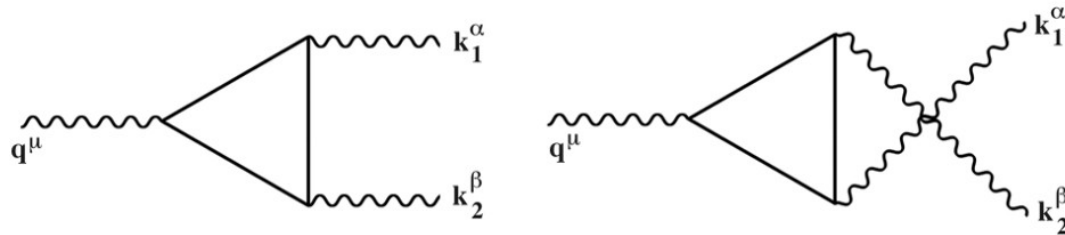




EW baryogenesis – B violation



- The main source of CP violation in the SM comes from the EW sector and is linked to the Higgs boson
- What about the other 2 Sakharov's conditions?
 - > Baryon number violation → is there baryon number violation in the SM?
 - In the SM the baryon (B) and lepton (L) number are conserved only via accidental symmetries – there are no gauge invariant operators of dim 4 that violate them
 - **Anomalies** appear at triangular loop level due to the chiral structure of the SM – Adler, Bell and Jackiw (ABJ) anomalies



$$\partial_\mu J_B^\mu = -\frac{\alpha_2}{8\pi} N_g W_i^{\mu\nu} \tilde{W}_{i\mu\nu} + \frac{\alpha_1}{8\pi} N_g \left(\frac{4}{9} + \frac{1}{9} - \frac{1}{18} \right) Y^{\alpha\beta} \tilde{Y}_{\alpha\beta}$$

$$\partial_\mu J_L^\mu = -\frac{\alpha_2 N_g}{8\pi} W_i^{\mu\nu} \tilde{W}_{i\mu\nu} + \frac{\alpha_1}{8\pi} N_g \left(1 - \frac{1}{2} \right) Y^{\alpha\beta} \tilde{Y}_{\alpha\beta} .$$



Violation of baryon and lepton number violation at loop level



Sphalerons



- $U(1)_B$ and $U(1)_L$ are broken at quantum level, as is $B+L$, however **$B-L$ is preserved**

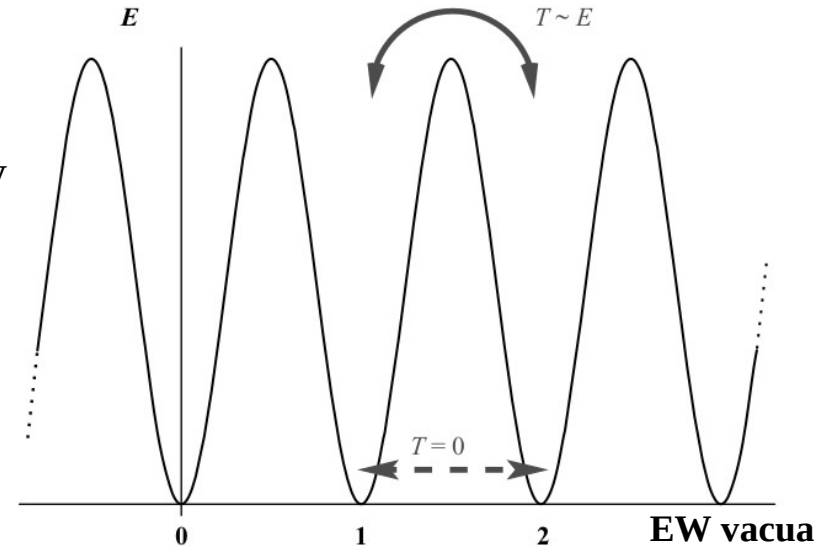
$$> \partial_\mu J_{B+L}^\mu = \frac{\alpha_1^2}{8\pi} N_g Y^{\alpha\beta} \tilde{Y}_{\alpha\beta} - \frac{\alpha_2^2}{4\pi} N_g W_i^{\alpha\beta} \tilde{W}_{i\alpha\beta}$$

$$> \partial_\mu J_{B-L}^\mu = 0 .$$

- At low temperatures the B violation is low, but at high temperature it can proceed via **sphaleron** production

[F. R. Klinkhamer and N. S. Manton Phys. Rev. D 30, 2212](#)

- > A Sphaleron – from Greek Σφαλερός ready to fall – is a static, but unstable, solution for the electroweak theory
- > It allows the transition between distinct EW vacuum states (different baryon or lepton number)
- > $E_{\text{sph}} \sim 8\pi v(T)/g_W$ – the energy is proportional to the vacuum expectation value of the Higgs field



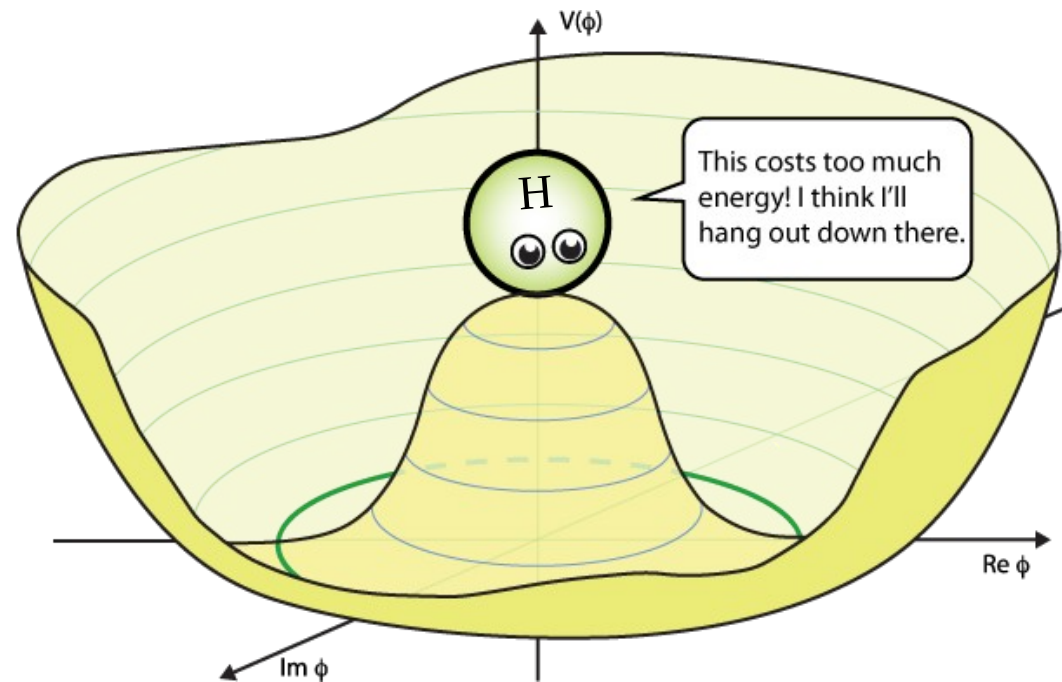
[M. Dine \(2016\). Supersymmetry and String Theory](#)



Electroweak phase transition



- The last condition necessary for matter-antimatter asymmetry to occur is that the CP violating interactions take place out of thermal equilibrium
- We need a **phase transition**: an interaction outside of thermal equilibrium
- The electroweak symmetry breaking provides the appropriate conditions
- The Higgs potential in the present has the usual *sombrero* shape you are familiar with

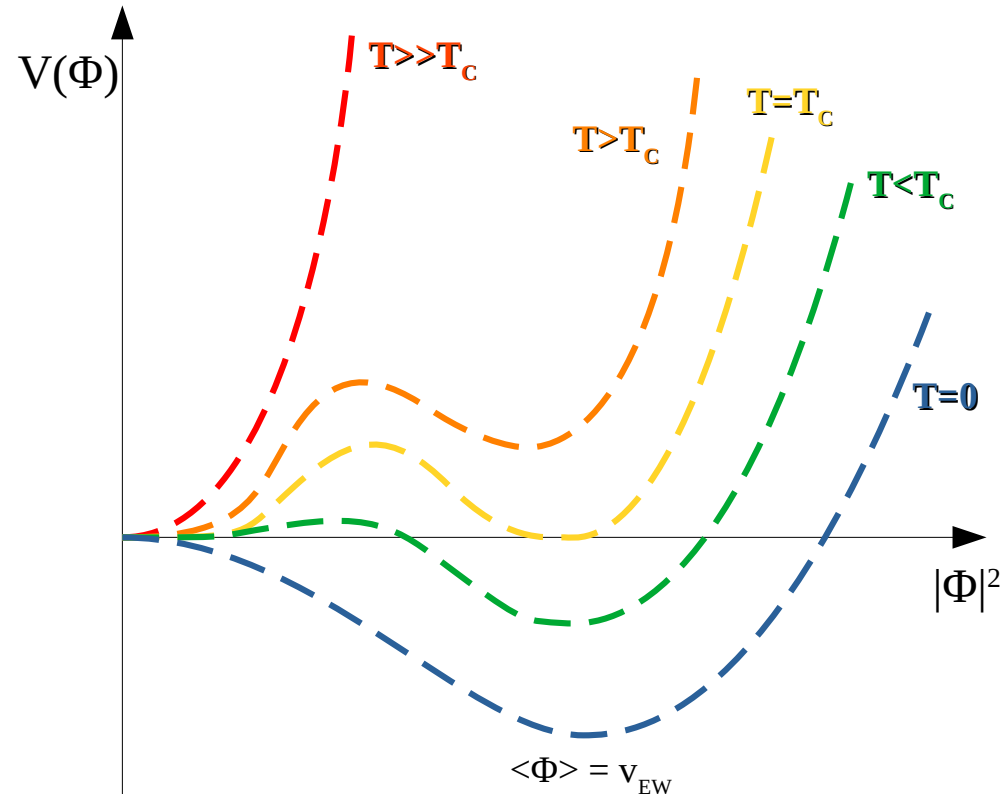




Electroweak phase transition



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- We need a **phase transition**: an interaction outside of thermal equilibrium
- The electroweak symmetry breaking provides the appropriate conditions
- The Higgs potential in the present has the usual *sombrero* shape you are familiar with
- The electroweak phase transition started at $T < T_c$ – when the the absolute minima of the system became $|\Phi|^2 > 0$
 - > 1st order PT: the system goes to the absolute minima via tunnelling

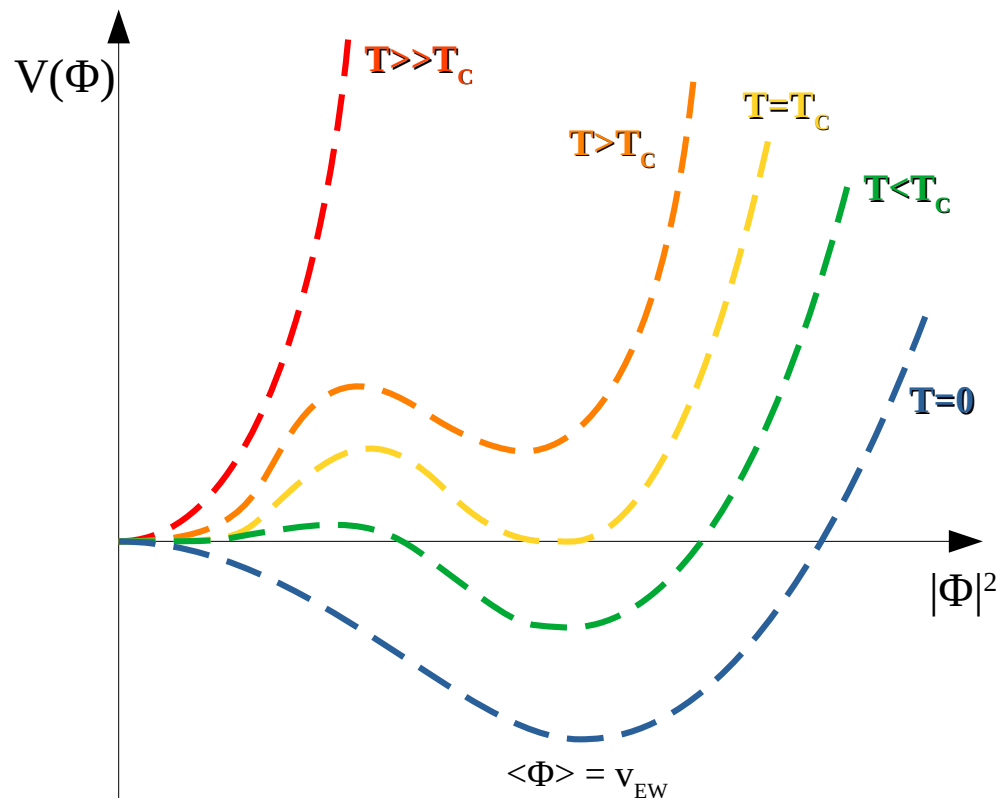




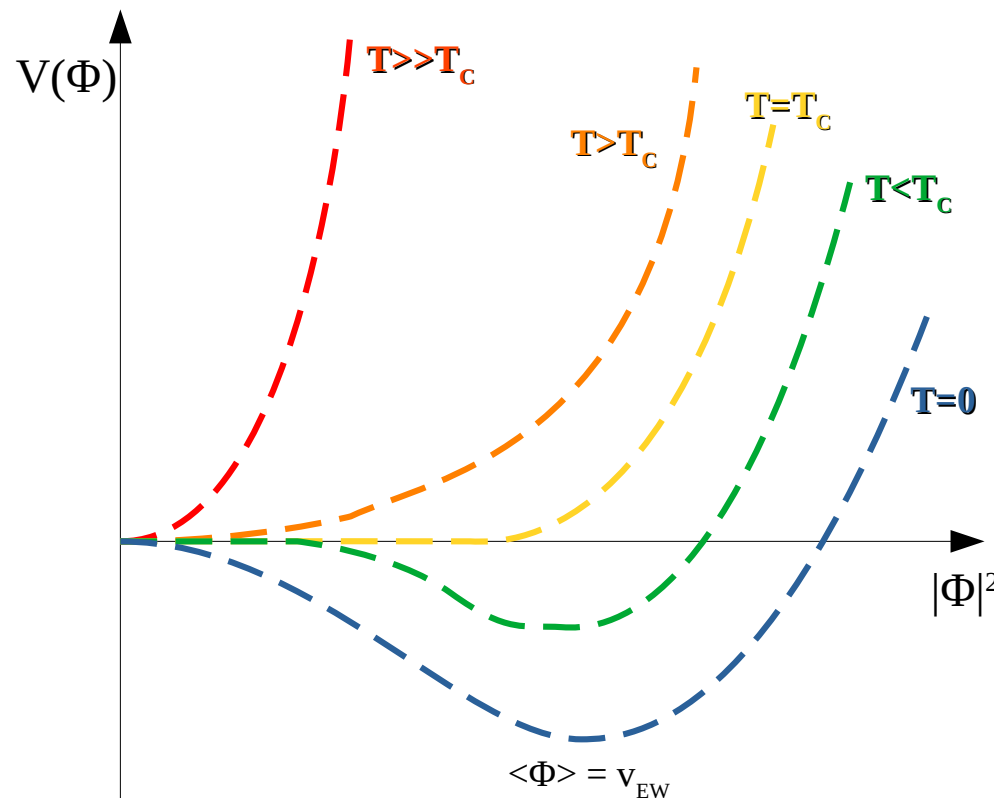
Electroweak phase transition



1st order PT: the system goes to the absolute minima via tunnelling



2nd order PT: there is no energy barrier and the system goes to the minimum *smoothly*

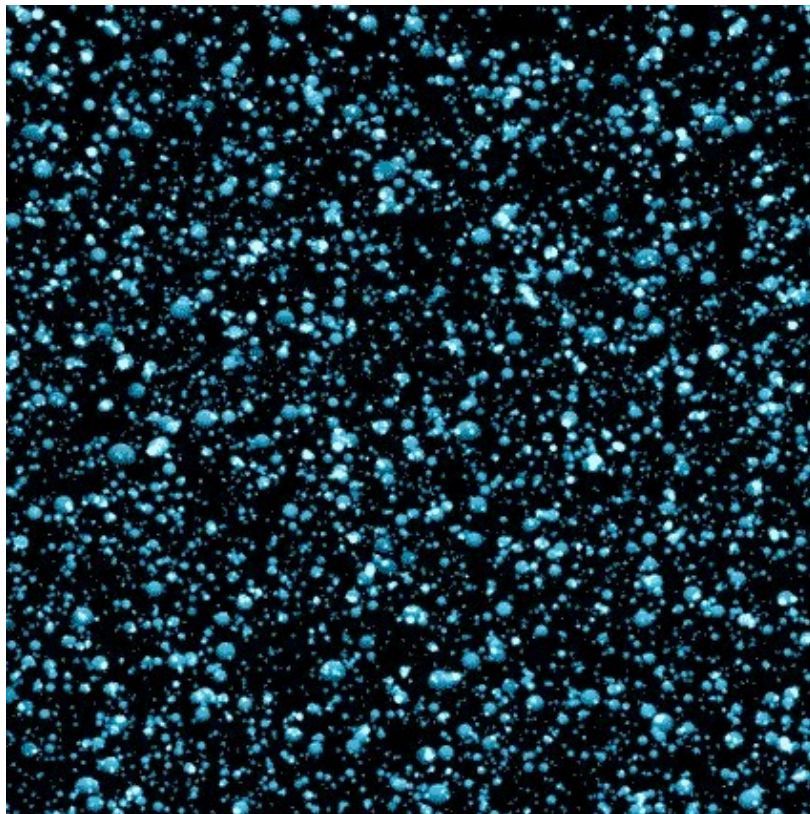




Higgs bubbles



- Like when boiling water, the entire Universe did not reach EWPT instantaneously everywhere



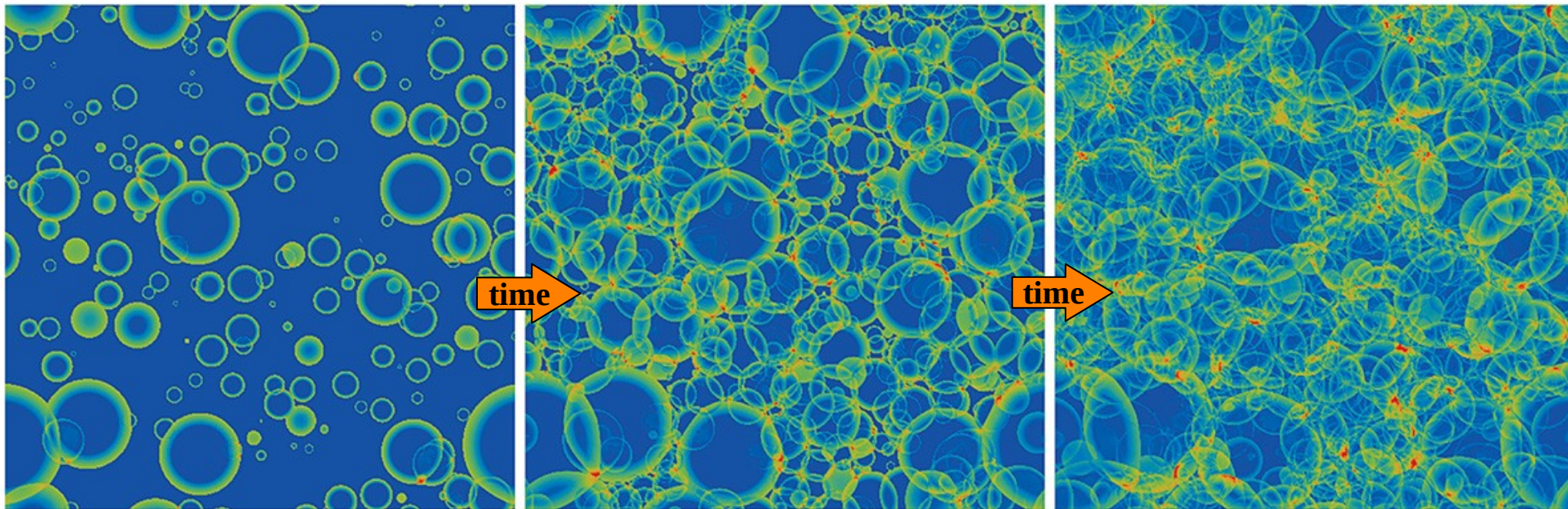
[giphy.com](https://www.giphy.com)



Higgs bubbles



- Like when boiling water, the entire Universe did not reach EWPT instantaneously everywhere
- The transition happened along expanding *bubbles* where the EW symmetry breaks

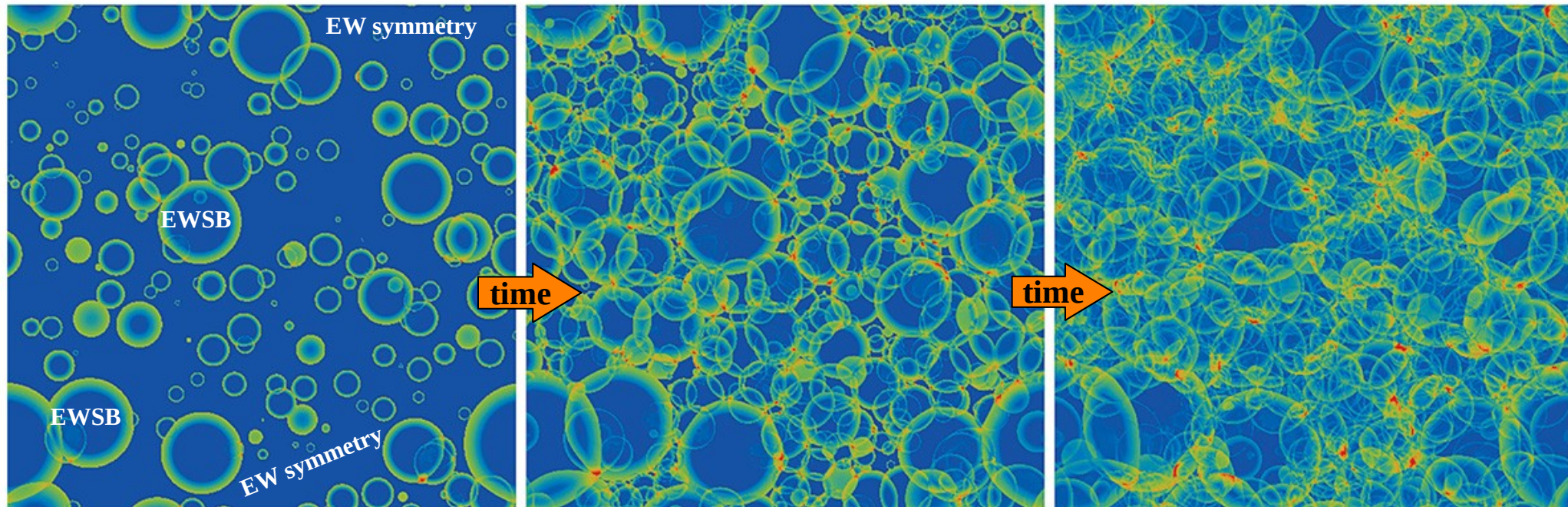




Higgs bubbles



- Like when boiling water, the entire Universe did not reach EWPT instantaneously everywhere
- The transition happened along expanding *bubbles* where the EW symmetry breaks
- Along the bubble walls – transition from the symmetric state to the broken EW one
- Fields on the bubble walls are highly out of equilibrium $\rightarrow$ satisfy the 3rd Sakharov's condition



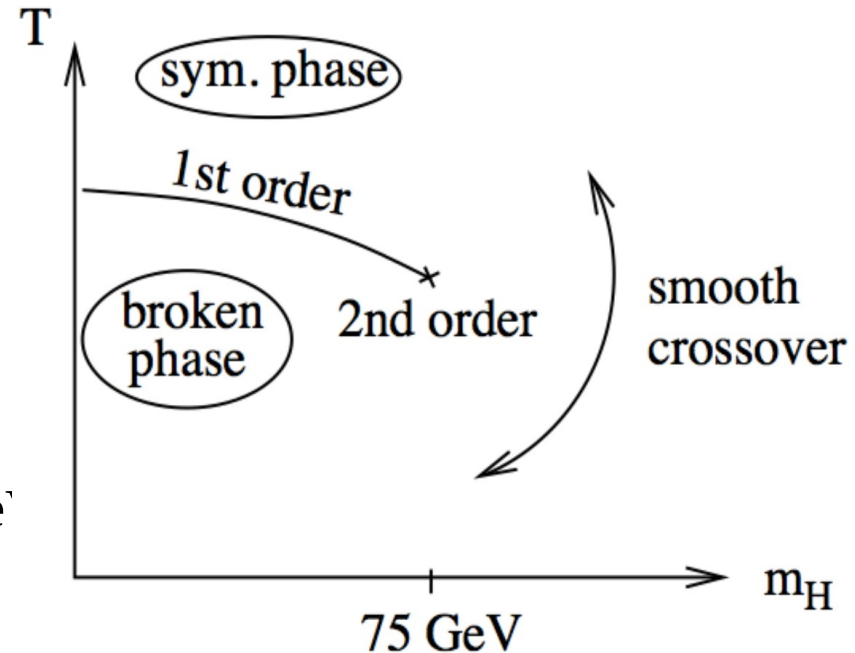


The SM Higgs and the EWPT



- B violating processes have different rates in the EW symmetric or broken phases
 - > In the unbroken phase T is high enough to have a continuous sphaleron production $\rightarrow$ B-L is preserved but the rate of B violating processes is high
 - > At the wall they are Boltzmann suppressed
 - > In the broken phase they are exponentially suppressed
- To have baryon asymmetry inside the bubble and freeze it, the sphaleron processes need to shut down quickly
 - $\rightarrow$ **require a strong 1st order phase transition**
- Problem: this would require a Higgs of mass < 80 GeV
 - > The SM provides the conditions for generating some baryon asymmetry but not enough for the observed BAU

Nucl.Phys. B443 (1995) 47
Phys. Rev. Lett. 77, 2887



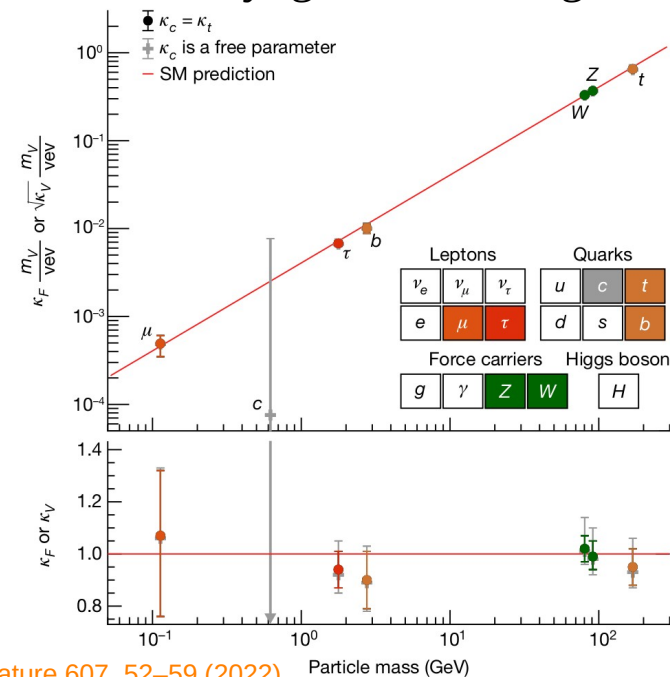


Explaining the BAU with CPV in the Higgs sector



- To explain the BAU we need either
 - > More Higgs bosons $\longrightarrow$ extra Higgs doublets/singlets, SUSY, etc.
 - > More sources of **CP violation** $\longrightarrow$ today's topic
- As the Higgs boson is intrinsically linked to the SM contributions to the baryogenesis it is a good candidate to look for additional CPV
- What do we know about the Higgs boson now?
 - > Discovery announced in 2012 by the ATLAS and CMS experiments – over 10 years ago
 - > Established couplings to 3rd generation fermions
 - > Coupling to W and Z well studied since Run 1
 - > More recently studied are decays to muons and charm quarks, and evidence for $Z\gamma$ has been established

} Not mutually exclusive



Nature 607, 52–59 (2022)

Particle mass (GeV)



CP violation in the Higgs sector



- Higgs boson predicted to have spin-parity 0^+ → direct coupling to Z and W bosons
- CP violation in the Higgs sector:

- > HVV couplings (V= Z,W bosons)

- studied in 4 lepton final state / VBF production

[Phys. Rev. D 104 \(2021\) 052004](#)

[Phys. Rev. D 108 \(2023\) 032013](#)

Pure CP odd excluded
already with Run 1 data →
search for anomalous couplings

- > Yukawa coupling of Higgs to fermions

- $gg \rightarrow H$ +jets: effective Higgs-gluon coupling via top quark loop

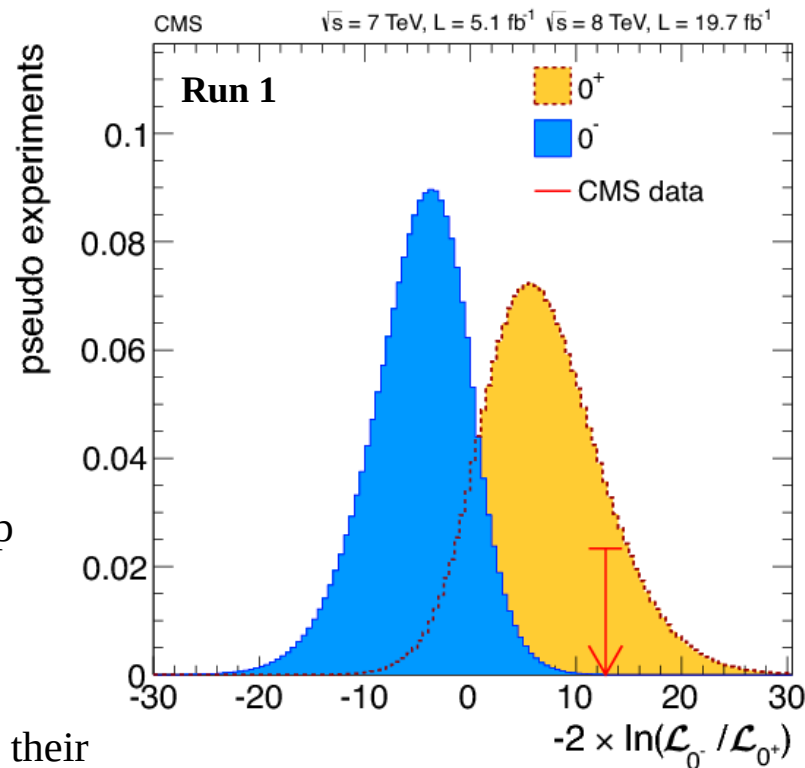
- [Phys. Rev. D 104 \(2021\) 052004](#)

- $gg \rightarrow t\bar{t}H$: direct Yukawa coupling of Higgs to top quarks

- [JHEP 07 \(2023\) 092](#)

- $H \rightarrow \tau\tau$ decays: exploit spin correlation between tau leptons and their decay products

- [JHEP 06 \(2022\) 012](#)



[Phys. Rev. D 89 \(2014\), 092007](#)



CP violation in HVV (and ggH)



- Investigating CP violation in HVV couplings with an EFT approach

- > Amplitude for Higgs coupling to two spin 1 particles (VV= WW, ZZ, Zγ, γγ, gg etc.) with operators up to dimension 6

Tree level CP-even coupling (=0 if absent in SM)

CP-even anomalous higher order couplings

$$\mathcal{A}(HVV) \simeq \left[a_1^{VV} + \frac{k_1^{VV} q_1^2 + k_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} + \frac{k_3^{VV} (q_1 + q_2)^2}{(\Lambda_Q^{VV})^2} \right] m_V^2 \varepsilon_{V1}^* \varepsilon_{V2}^* \\ + \left[a_2^{VV} \right] f_{\mu\nu}^{*(1)} f^{*(2)\mu\nu} + \left[a_3^{VV} \right] f_{\mu\nu}^{*(1)} \bar{f}^{*(2)\mu\nu} ,$$

CP-even anomalous coupling

CP-odd anomalous coupling

$$f_{a3} = \frac{\sigma(a_3=1, a_{i \neq 3}=0)}{\sum_i \sigma_i}$$

- > Effect on cross-section parametrized as the fractional contribution of the anomalous coupling to the total cross-section

- Theoretical approach includes also Hγγ and Hgg with tree level coupling being set to 0

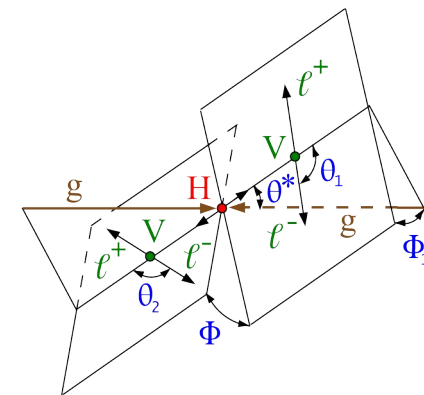
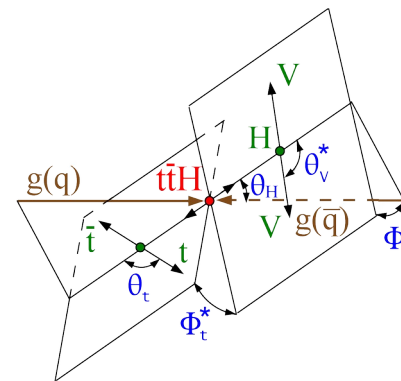
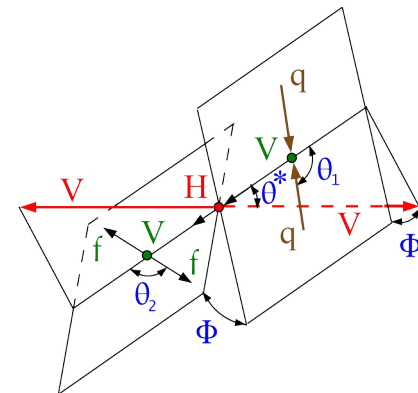
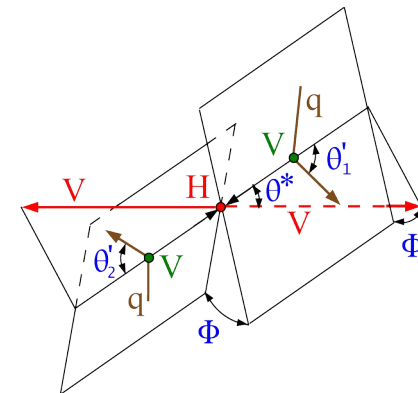


CP violation in HVV (and ggH)



- Higgs direct and effective couplings to vector bosons studied both in production and decay
- Decay:
 - > Target process: $H \rightarrow VV \rightarrow 4 \text{ leptons}$
 - > Uses kinematics of leptons in the final state
- Production:
 - > Target process: **Vector Boson Fusion (VBF)**
 - > Higgs CP nature affects kinematics of jets from the initial state
 - > Studied in Higgs decays to $\tau\tau$, 4 leptons and $\gamma\gamma$

[Phys. Rev. D 104 \(2021\) 052004](#)

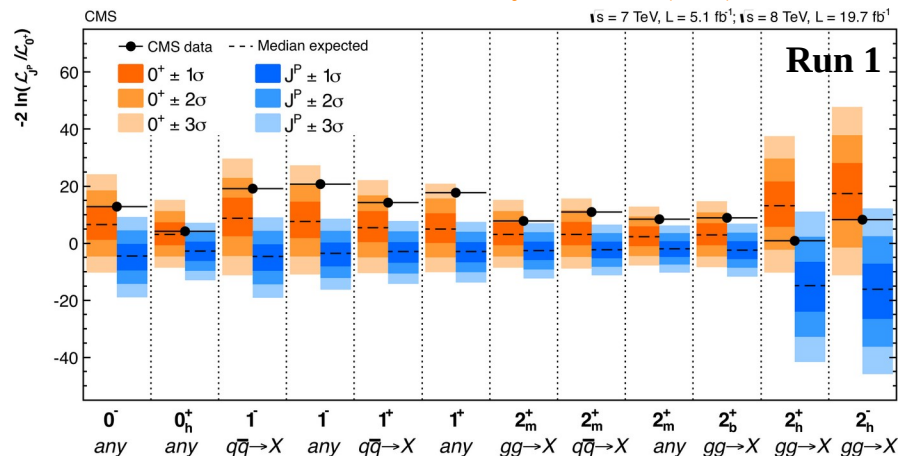




CP violation in $H \rightarrow 4$ leptons

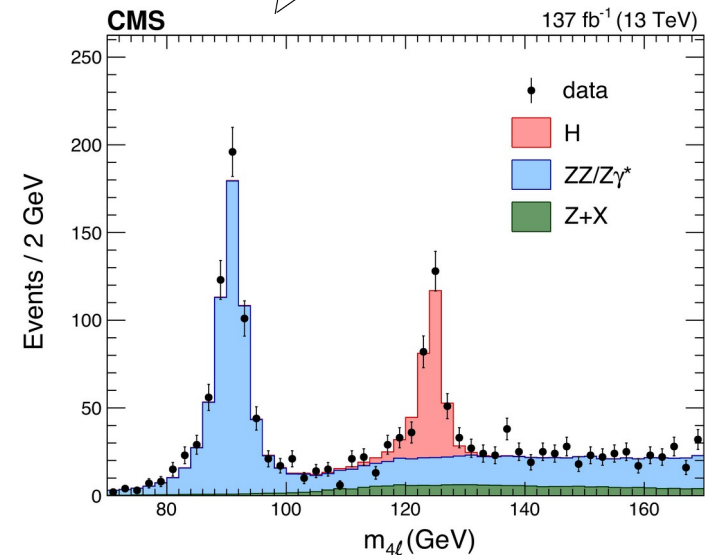
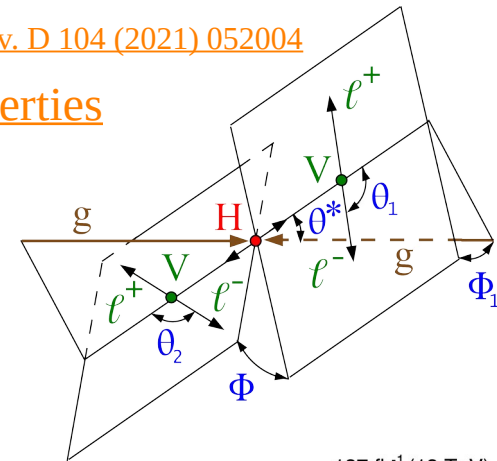


- Higgs decays to 4 leptons were used already in the Higgs discovery [Phys. Rev. D 104 \(2021\) 052004](#)
- Excellent handle to measure Higgs properties: [mass](#), [spin](#), [width](#), and [CP properties](#)
 - Used in Run 1 to exclude pure pseudoscalar HVV coupling and investigate other spin hypotheses [Phys. Rev. D 89 \(2014\) 092007](#)



- Analysis with full Run 2 also targets interference terms of anomalous couplings with other hypothesis using MVA classifiers

$$\mathcal{D}_{\text{int}}(\Omega) = \frac{\mathcal{P}_{\text{int}}(\Omega)}{2 \sqrt{\mathcal{P}_{\text{sig}}(\Omega) \mathcal{P}_{\text{alt}}(\Omega)}}$$





CP violation in $H \rightarrow 4$ leptons

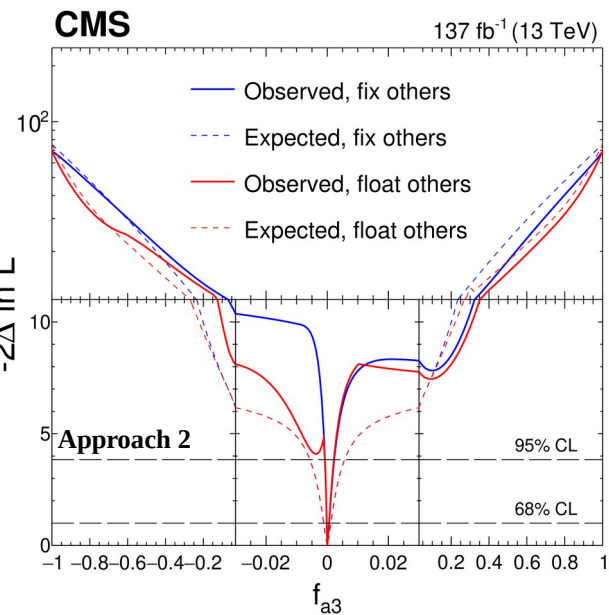
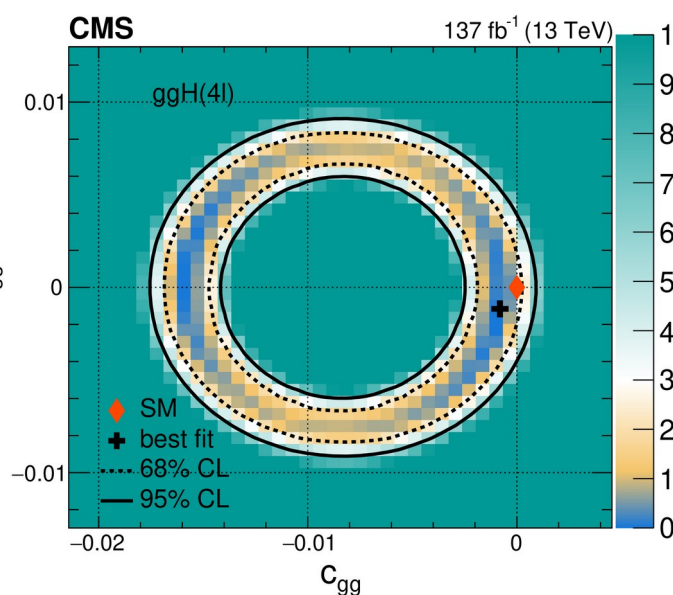
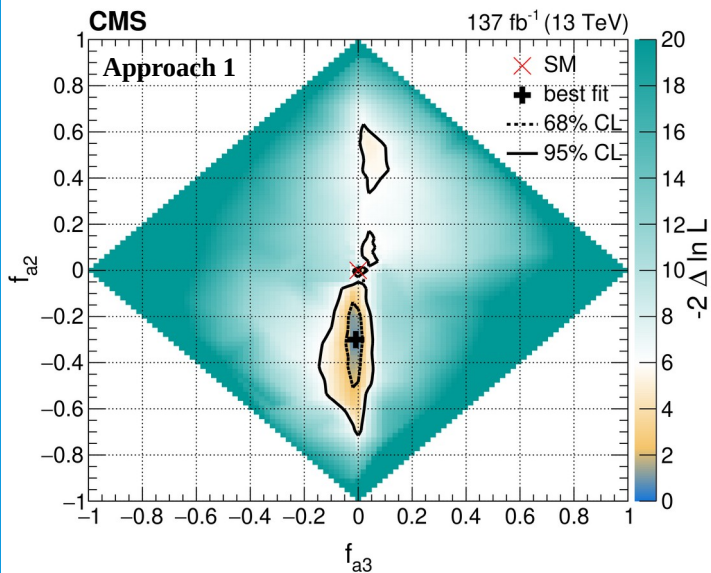


- Constraints expressed in several scenarios based on the possible symmetries imposed on the couplings:

- > Approach 1: $a_i^{WW} = a_i^{ZZ}$
- > Approach 2: $SU(2) \times U(1)$ symmetry

[Phys. Rev. D 104 \(2021\) 052004](#)

- The 4 lepton final state is also used to target H_{gg} and H_{tt} couplings

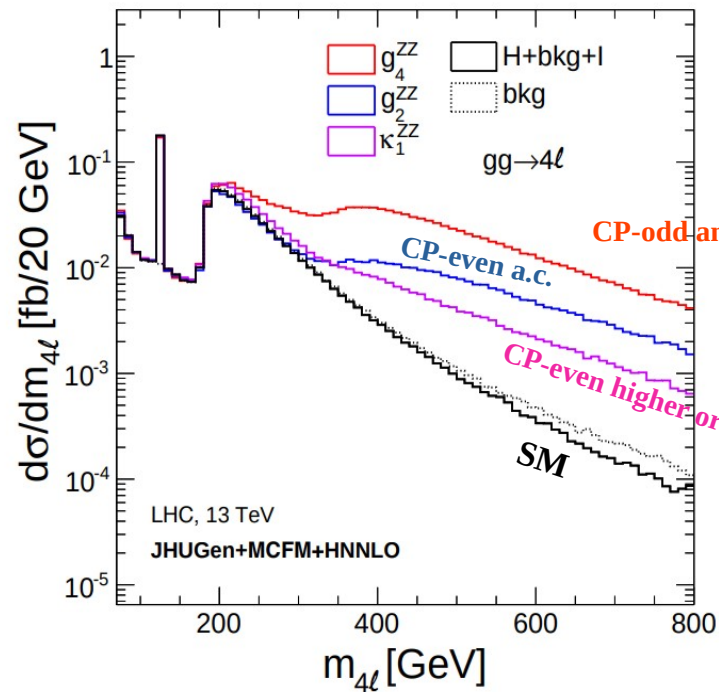




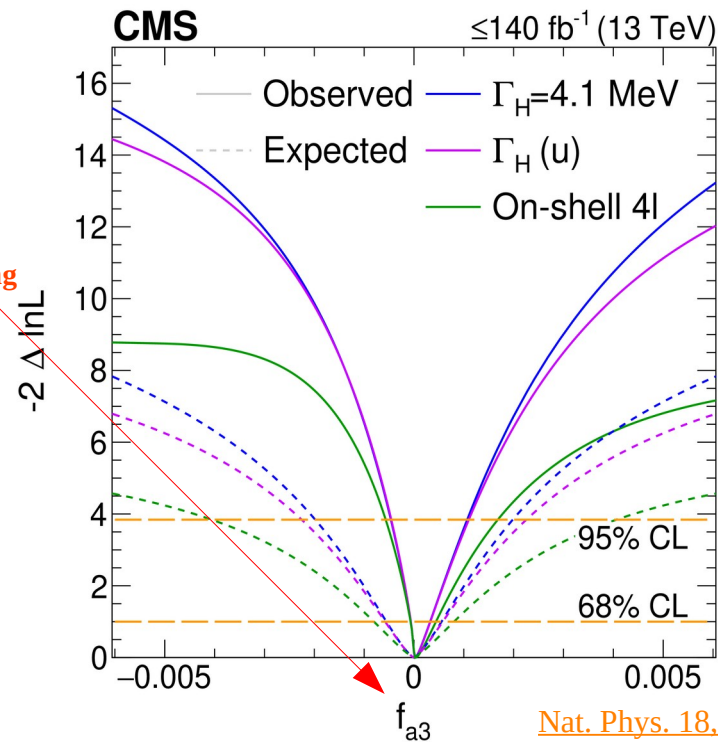
Combination with off-shell $H \rightarrow 2l2\nu/4\text{ lepton}$



- Off-shell Higgs production can provide additional sensitivity to anomalous couplings
 - > CP-odd anomalous coupling constrained between $[-4.6, 11] \times 10^{-4}$ at 95% CL



[Phys.Rev.D102\(2020\)056022](#)



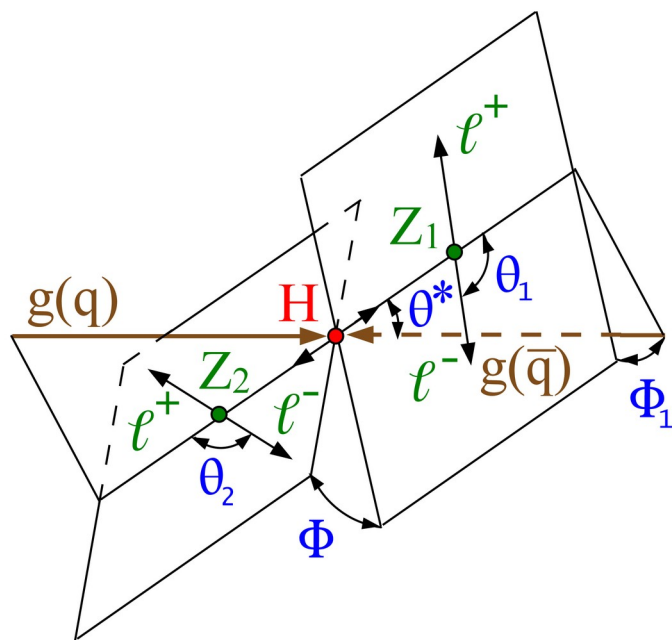
[Nat. Phys. 18, 1329–1334 \(2022\)](#)



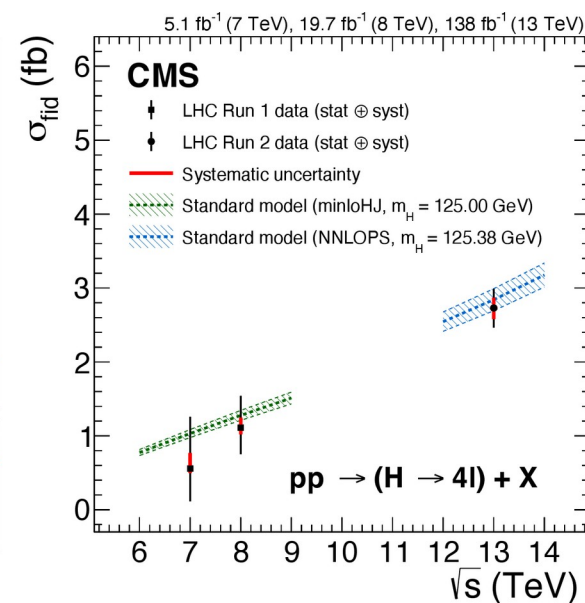
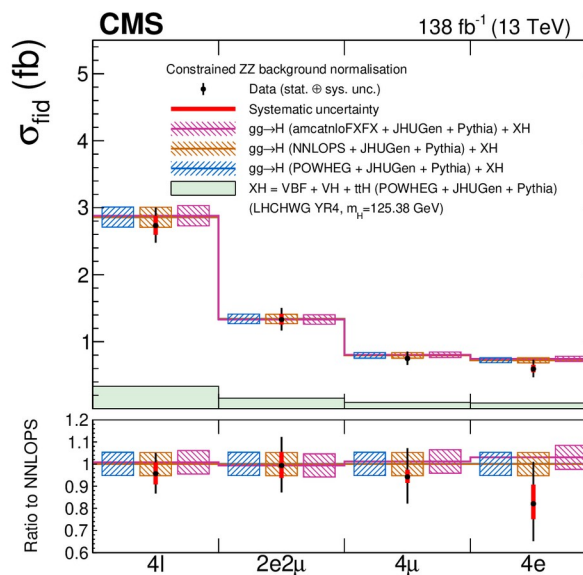
$H \rightarrow ZZ^* \rightarrow 4 \text{ leptons}$ differential cross-section



- $H \rightarrow 4 \text{ leptons}$ also studied differentially
- Differential measurement performed with respect to angular variables, Higgs momentum, and accompanying jet $p_T \rightarrow$ sensitivity to loop correction and BSM operators



Inclusive fiducial cross-section



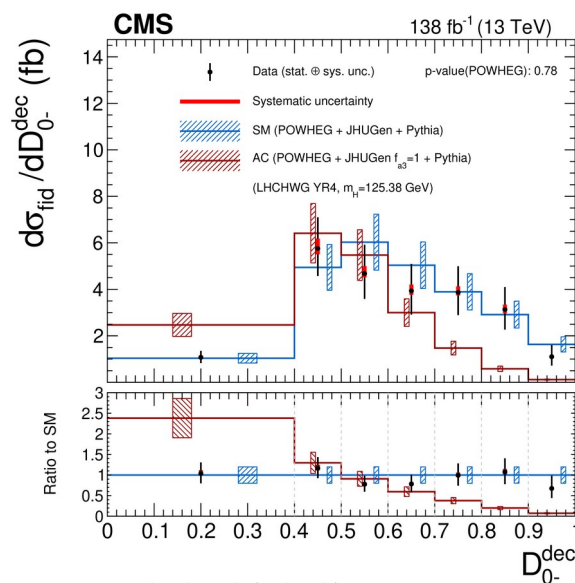
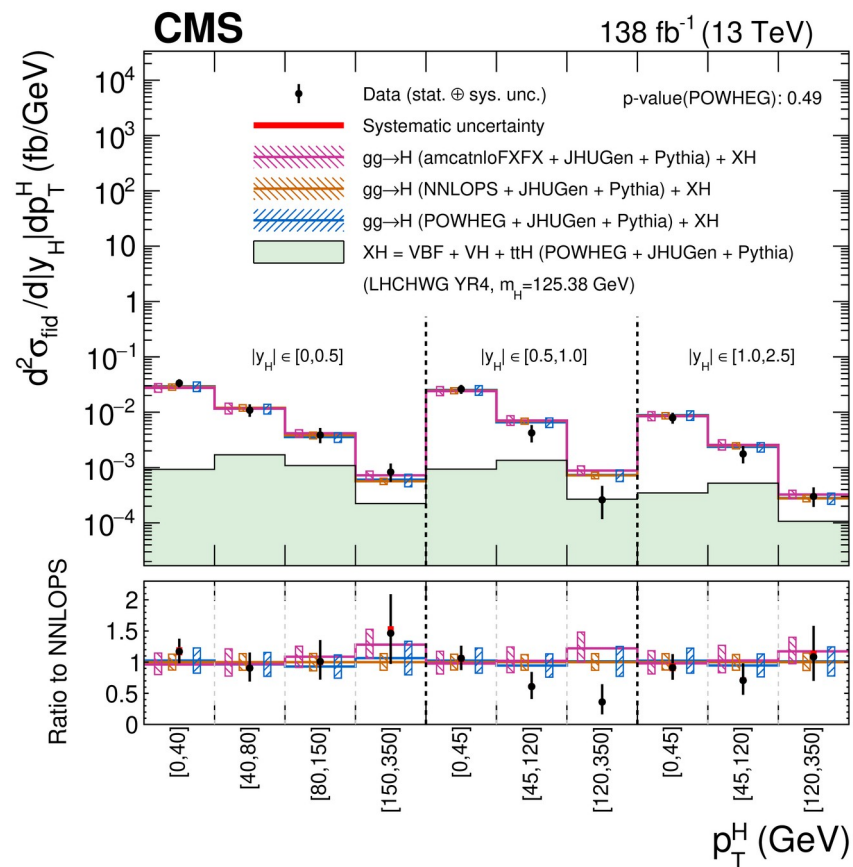
JHEP2023,40(2023)



$H \rightarrow ZZ^* \rightarrow 4 \text{ leptons}$ differential cross-section

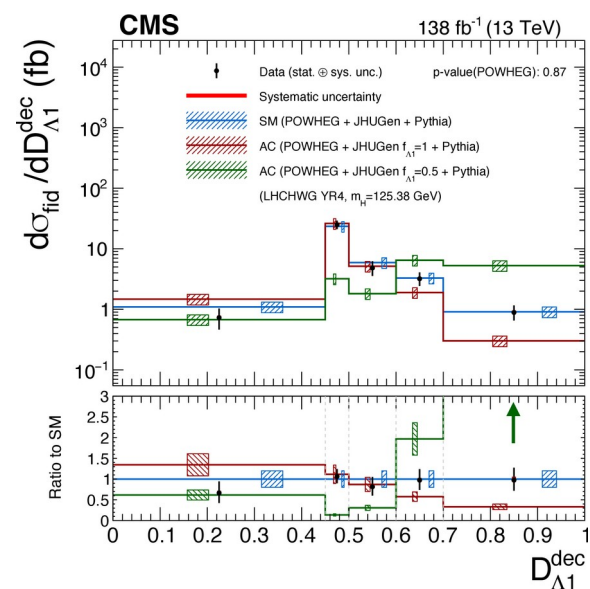


- The increased sensitivity allows to constrain better CP-violating terms and validate our current theoretical understanding



[JHEP2023,40\(2023\)](#)

- All differential distributions agree with the latest SM predictions

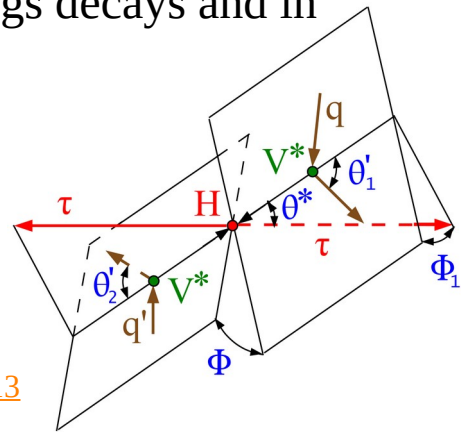




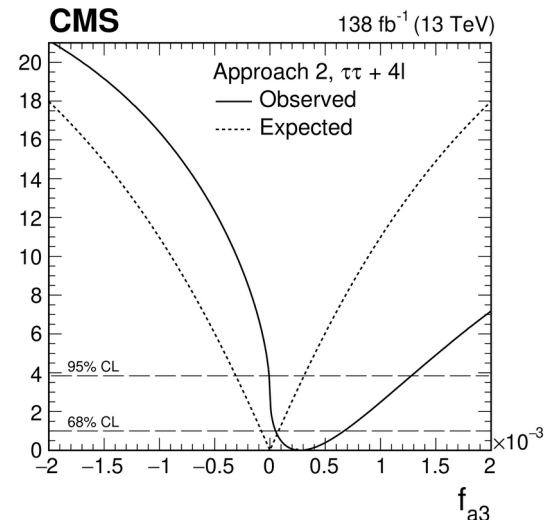
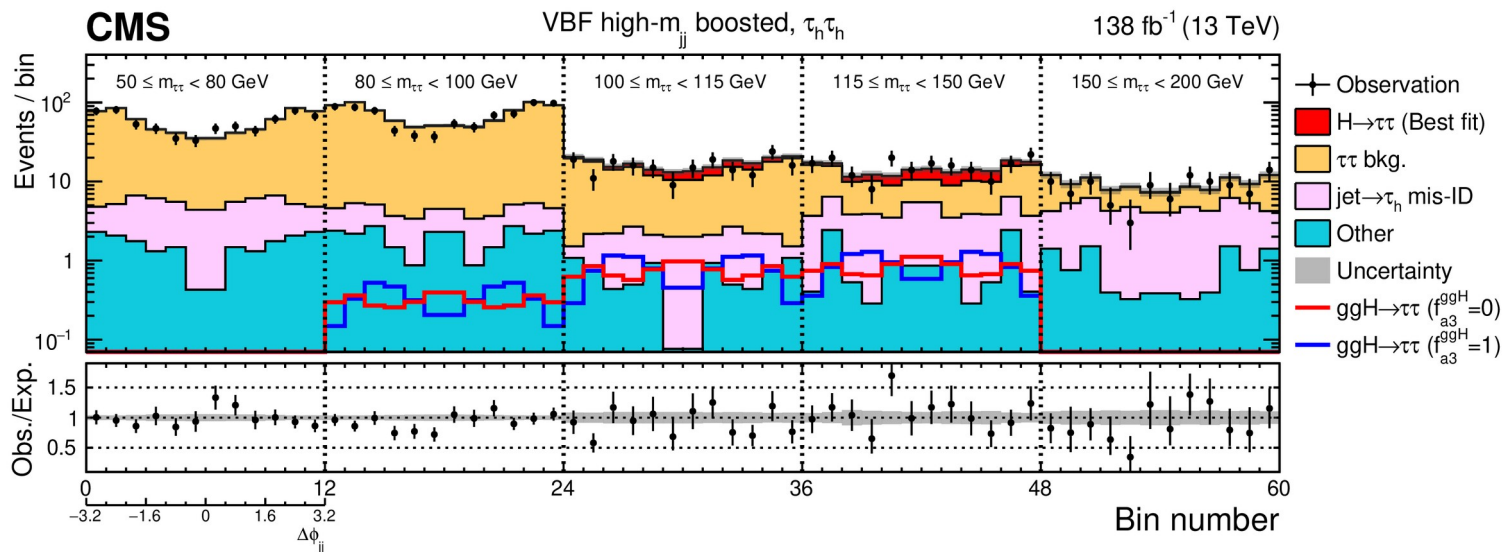
VBF production with $H \rightarrow \tau\tau$ decays



- Higgs decays to tau leptons can be used as a probe both for CP violation in Higgs decays and in combination with $H \rightarrow 4l$ for the search of CP violation in production
 - When studying Higgs CP in production the $H \rightarrow \tau\tau$ vertex is kept SM-like
 - The azimuthal angle between jets coming from the initial state offers sensitivity to the Higgs CP quantum number



[Phys. Rev. D 108 \(2023\) 032013](#)





CP properties of Yukawa couplings



- Higgs Yukawa interaction can be parametrized as follows:

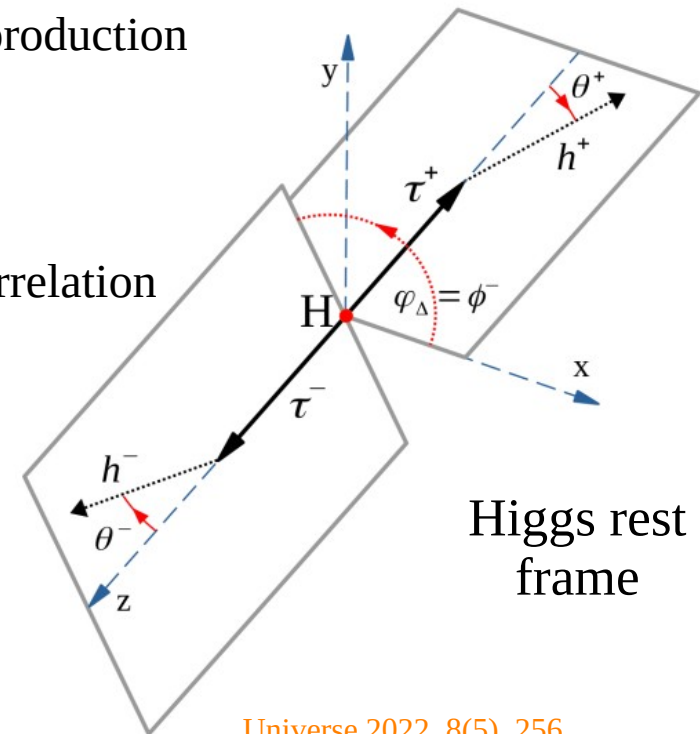
$$> \mathcal{L}_{Y,f} = -\frac{m_f}{v} \bar{\psi}_f (\kappa_f + i\gamma^5 \tilde{\kappa}_f) H \psi_f$$

- > Higgs direct coupling to top quarks investigated in ttH and tH production

- CP-odd anomalous coupling affects top quark(s) kinematics
- Investigated with MVA techniques

- > Higgs coupling to τ leptons investigated in decays via τ spin correlation

- CP mixing encoded in $\alpha^{H\tau\tau}$:
 - $\kappa_\tau = \sqrt{\mu^{\tau\tau}} \cos(\alpha^{H\tau\tau})$
 - $\tilde{\kappa}_\tau = \sqrt{\mu^{\tau\tau}} \sin(\alpha^{H\tau\tau})$
- Mixing angle can be accessed via the angle between τ decay planes



[Universe 2022, 8\(5\), 256](#)

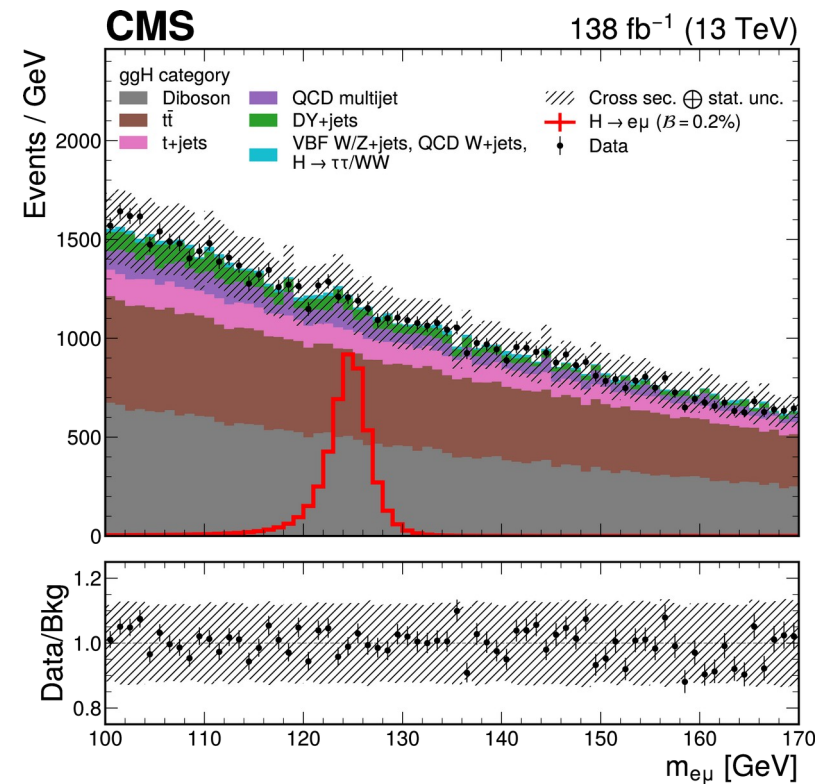


Lepton Flavor Violation



- Beyond the coupling to top quarks and tau leptons we also look for lepton flavor violation
 - > Lepton flavor is conserved $\rightarrow$ the Yukawa matrix is diagonal wrt lepton generations
- In analogy with the CKM matrix we can look for CP violation also by studying off-diagonal elements of the lepton Yukawa matrix
 - > A search for $H \rightarrow e\mu$ decays can hint towards sources of new physics
 - > Search performed in mass window $110 < m_{e\mu} < 160$ GeV
 - > Analysis exploits BDT to improve S/B ratio in two signal categories: ggH and VBF

[Phys. Rev. D 108, 072004](#)

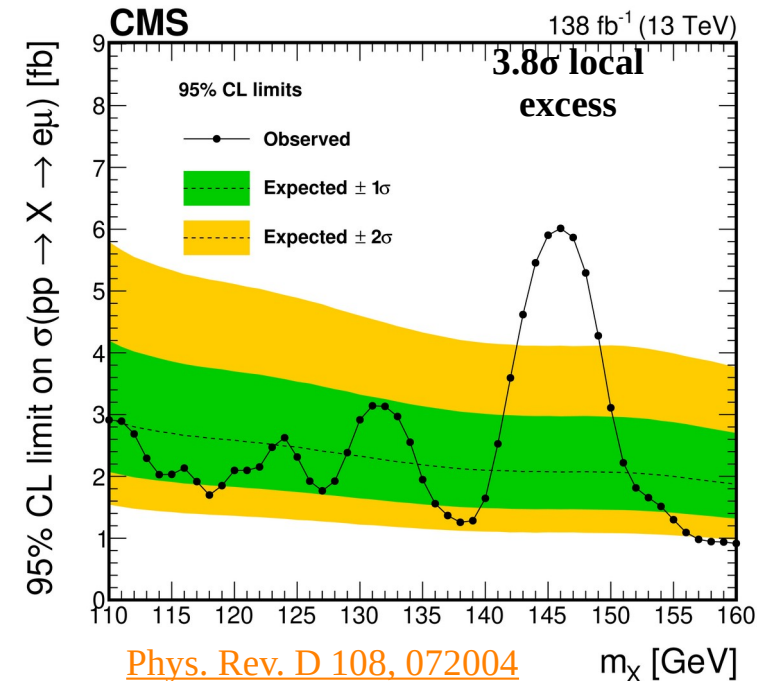
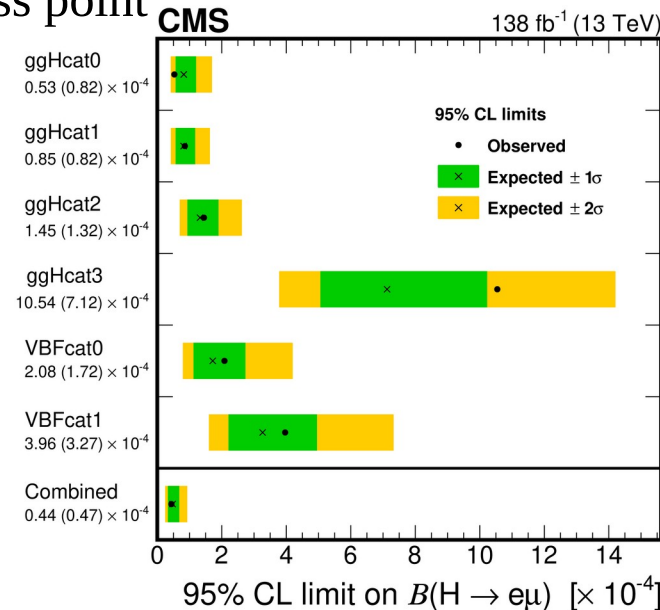




$H \rightarrow e\mu$ LFV decays

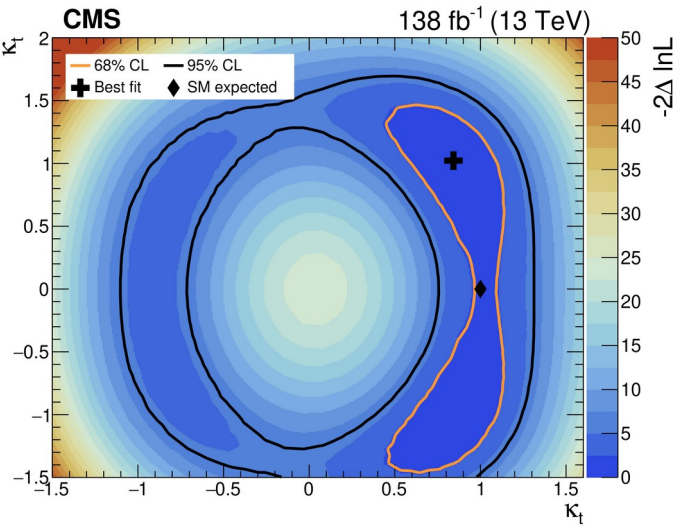
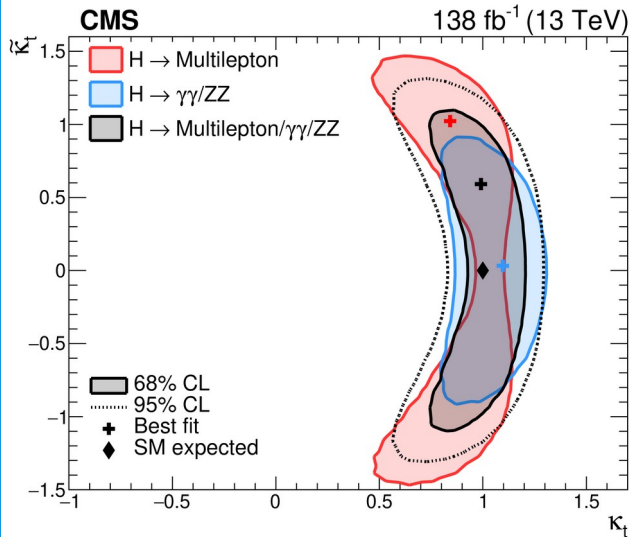


- Constraints on the $H \rightarrow e\mu$ branching fraction determined for SM-like $H(125)$ and more generally for generic scalar particle X of mass between 110 and 160 GeV
 - > **First direct search for $e\mu$ resonance in chosen mass range**
- Observed (expected) limit on $\mathcal{B}(H \rightarrow e\mu) < 4.4$ (4.7) $\times 10^{-5}$ at 95% for a SM-like Higgs
- The $X \rightarrow e\mu$ search presents a local (global) excess of 3.8σ (2.8σ) for the 146 GeV mass point
- The best fit signal for this mass point has a cross-section of 3.89 ± 1.25 fb
- More data is needed to investigate this excess





CP violation in Higgs-top couplings

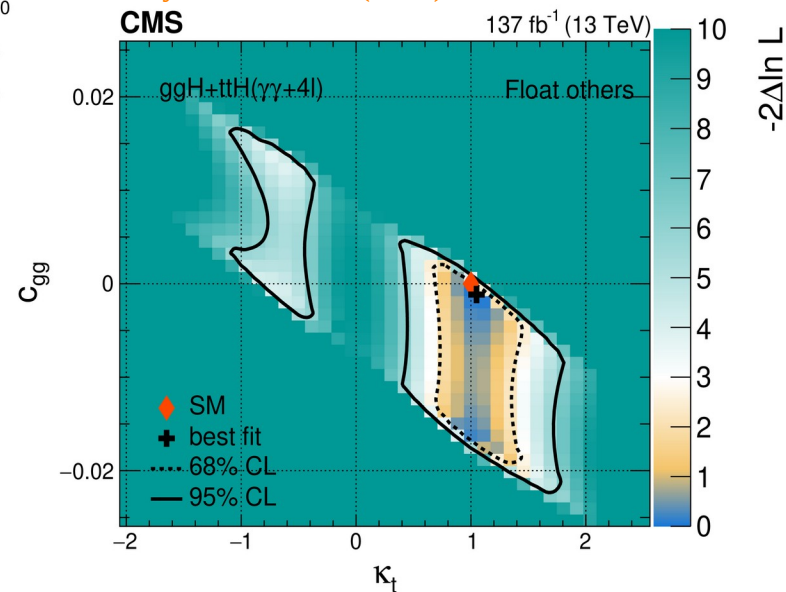


[JHEP 07 \(2023\) 092](#)

- Top quark loop has leading contribution to ggH effective coupling
 → simultaneous fit of CP-odd contribution ggH and ttH production
- Combination performed under hypothesis that ggH loop is dominated by top quark

- Higgs couplings are proportional to the mass → enhanced coupling with the top quark → target ttH and tH production
- Htt coupling CP structure was probed in the $H \rightarrow WW, \tau\tau$ (multilepton) $\gamma\gamma$ and ZZ decays

[Phys. Rev. D 104 \(2021\) 052004](#)





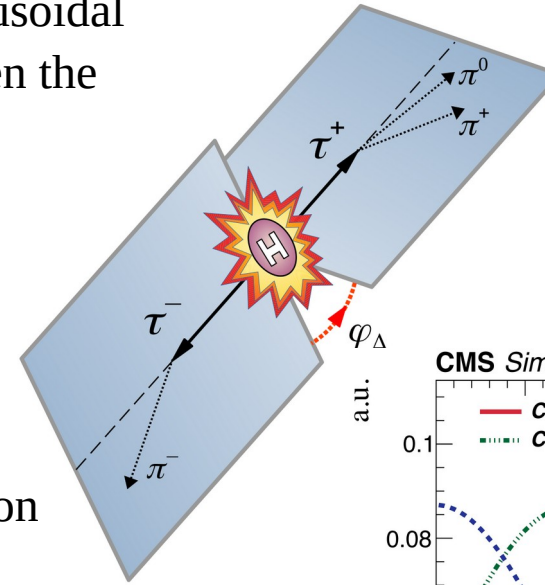
CP violation in $H \rightarrow \tau\tau$ decays



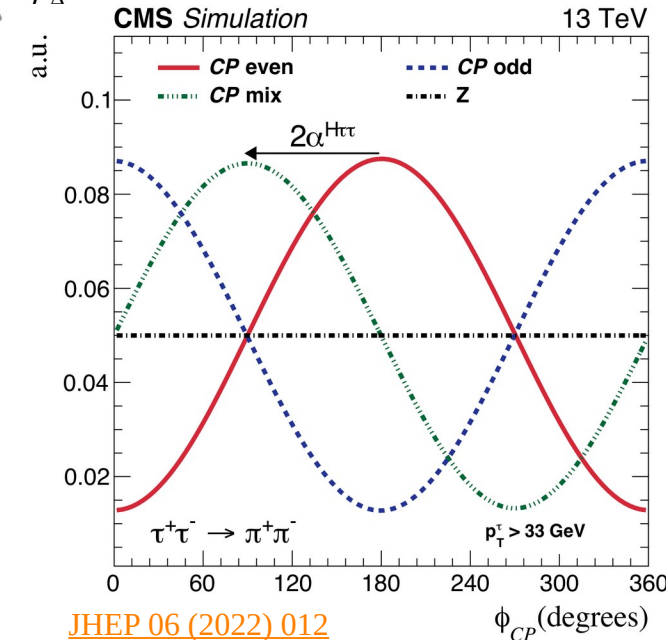
- The cross section of the $H \rightarrow \tau\tau$ decay acquires a sinusoidal dependence on the azimuthal angle spanned between the two τ polarimeters

$$\frac{d\sigma}{d\varphi_{CP}} \propto \text{const} - \cos(\varphi_{CP} - 2\varphi_{\tau\tau})$$

- CP mixing appears as a **phase-shift** in the distribution
- Measuring the Higgs CP mixing angle requires accessing the azimuthal difference between the polarimetric vectors
- This angle can be generalized to the angle between the τ decay planes, an **acoplanarity angle**



[Universe 2022, 8\(5\), 256](#)

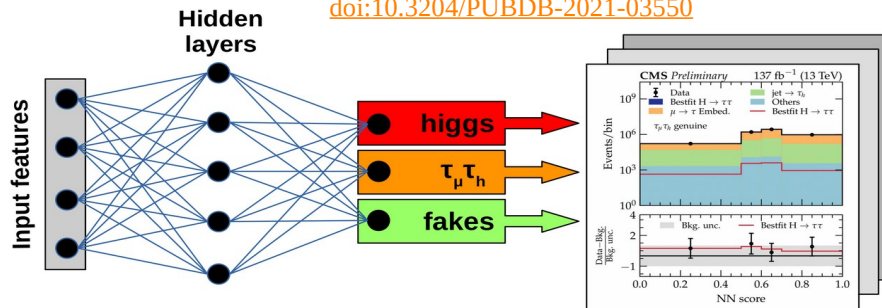




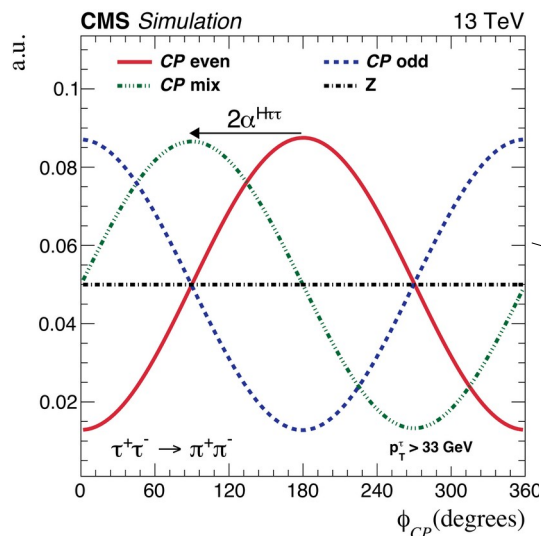
CP violation in $H \rightarrow \tau\tau$ decays



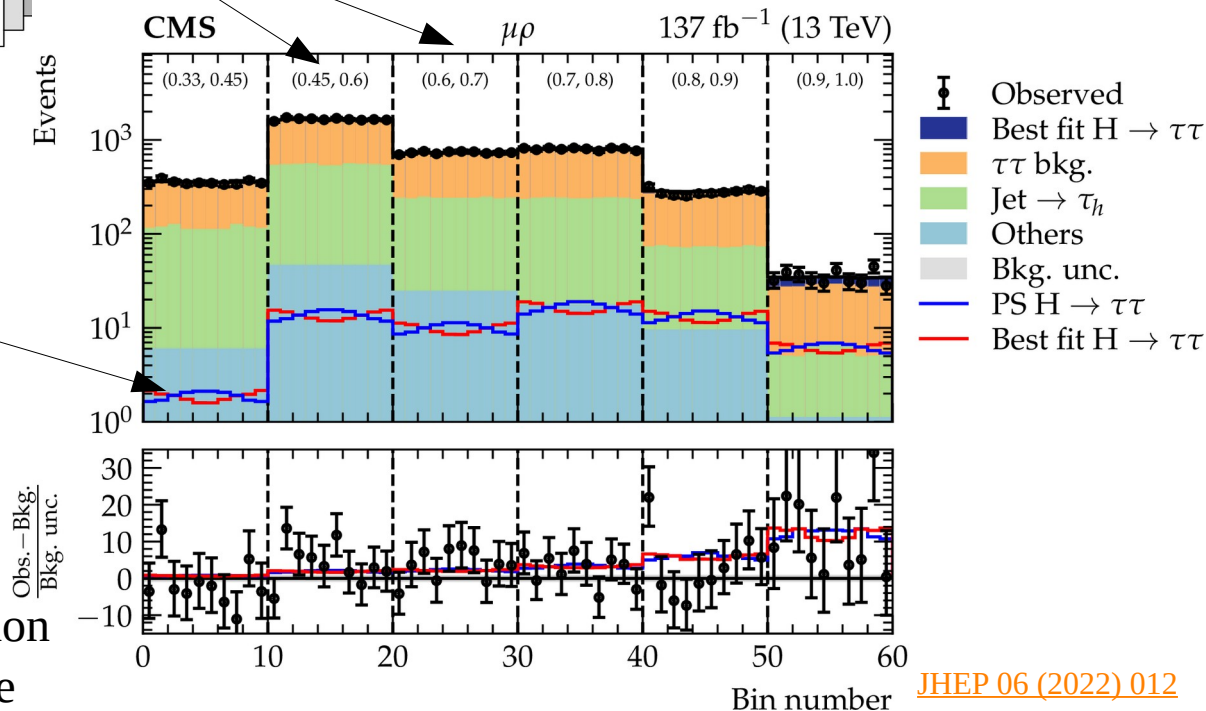
doi:10.3204/PUBDB-2021-03550



- Use of ML classifiers to identify the Higgs boson



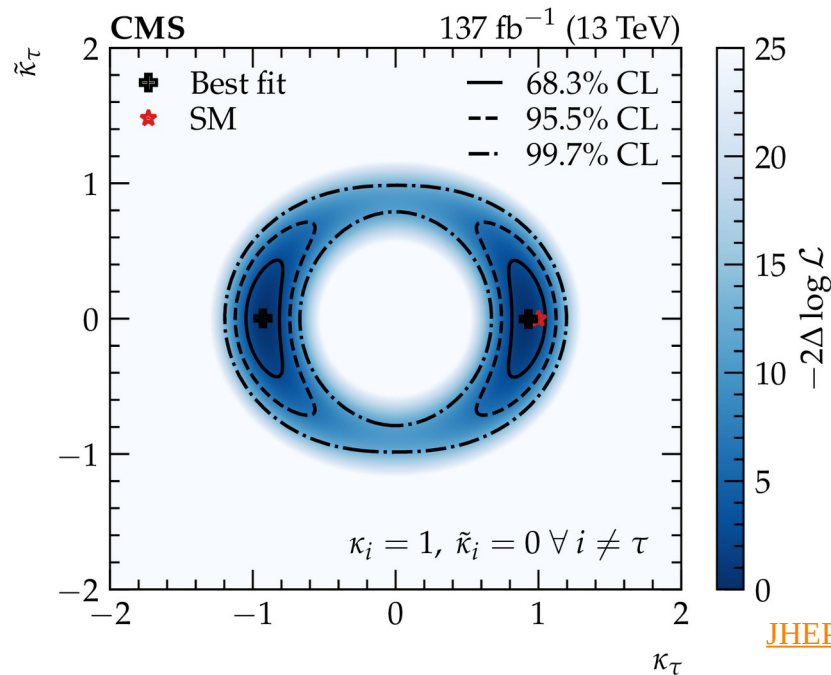
- Modulation of the differential cross-section can be used to probe the CP mixing angle



JHEP 06 (2022) 012

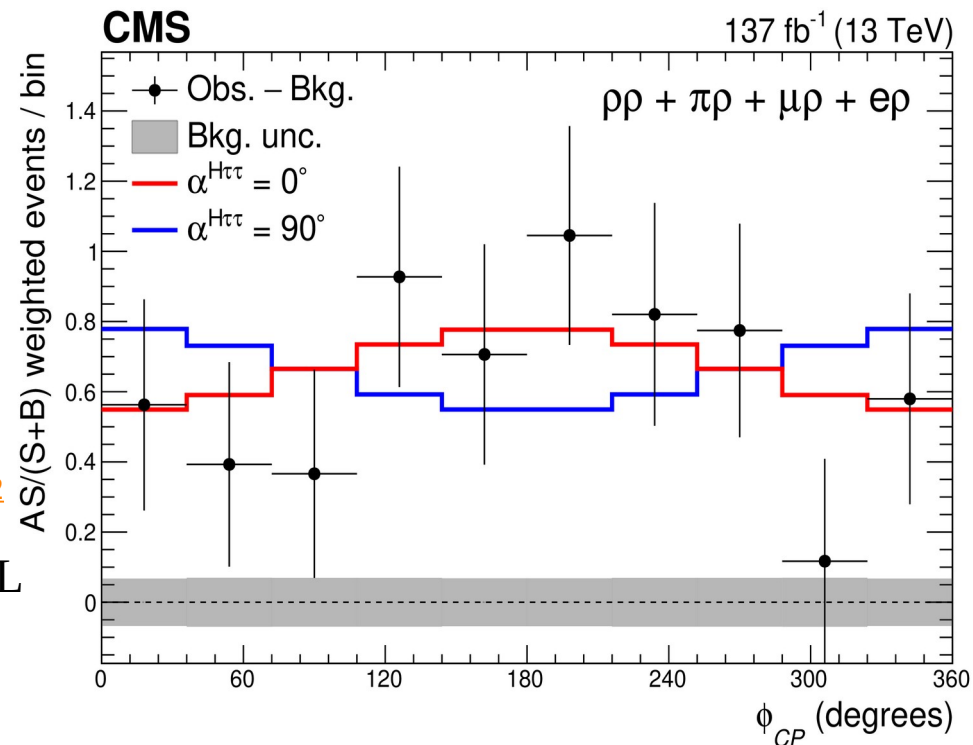


CP-even/odd couplings



- CP mixing angle:

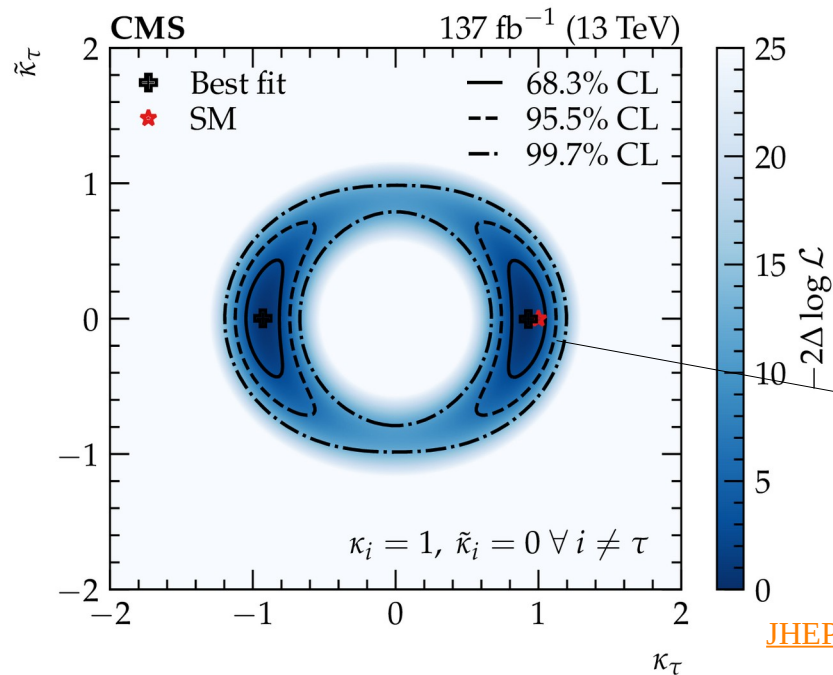
$$\alpha^{H\tau\tau} = -1 \pm 19 \text{ (stat)} \pm 1 \text{ (syst)} \pm 2 \text{ (bin-by-bin)} \pm 1 \text{ (theo)}^\circ$$



- Couplings consistent with SM hypothesis at 68% CL
 - Observed (exp.) exclusion of pure CP-odd hypothesis significance = **3.0σ** (2.6σ)



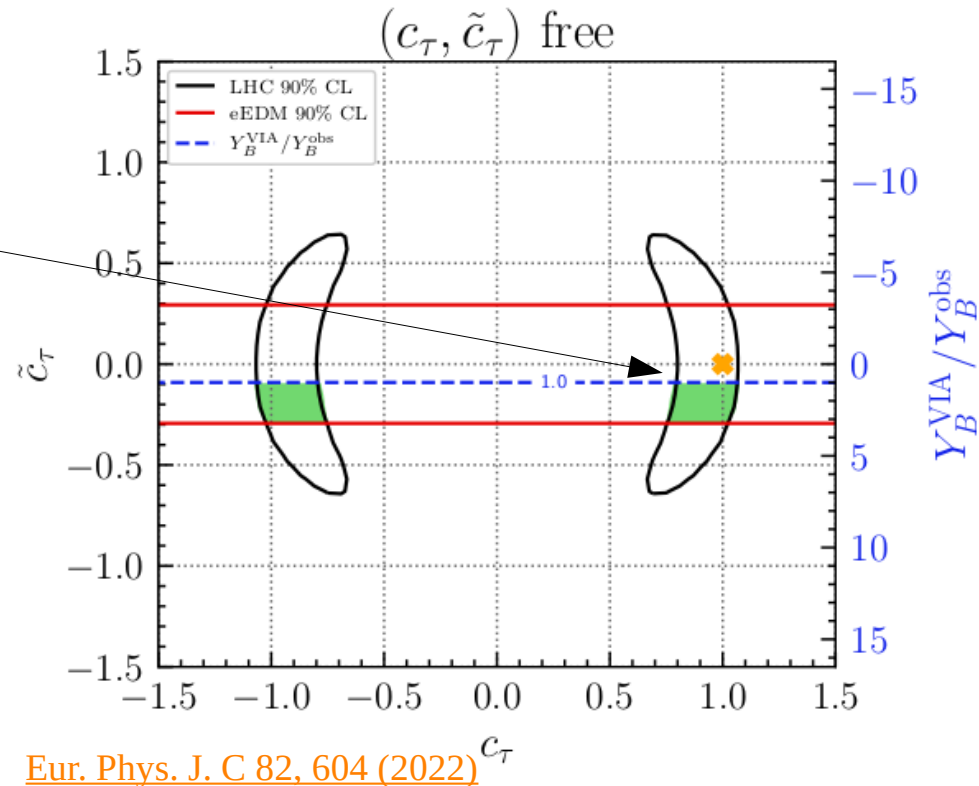
CP-even/odd couplings



[JHEP 06 \(2022\) 012](#)

- Couplings consistent with SM hypothesis at 68% CL
 - Observed (exp.) exclusion of pure CP-odd hypothesis significance = **3.0 σ** (2.6 σ)

- Interesting for MSSM interpretation to test CP mixing angle of 7° → good target for future measurement



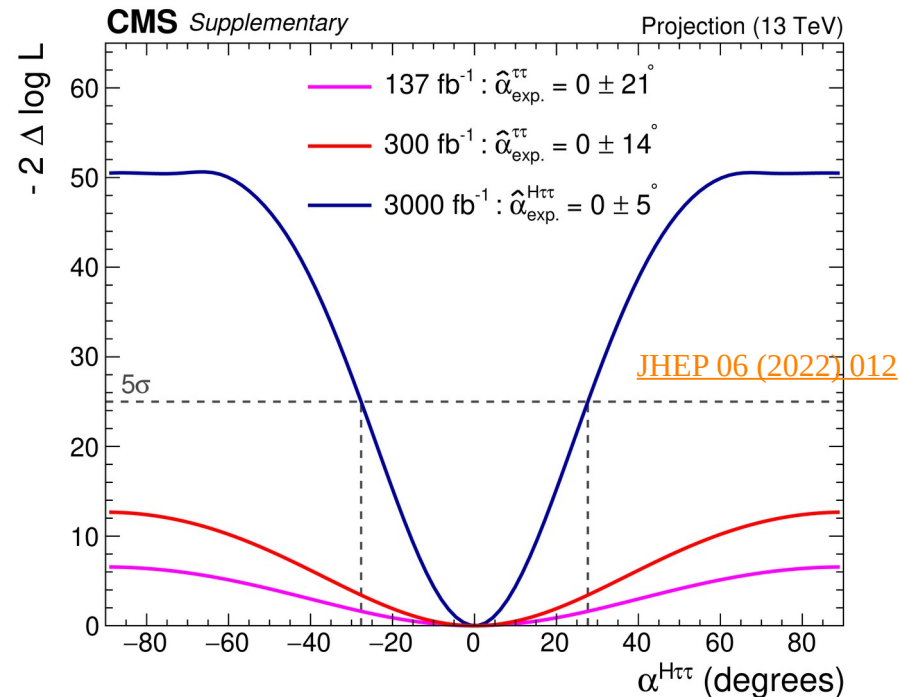
[Eur. Phys. J. C 82, 604 \(2022\)](#)



Summary and prospects



- Higgs boson couplings have been investigated looking for CP-odd contributions which could help explaining the baryon asymmetry in the Universe
- Exclusion limits for pure CP-odd Higgs couplings:
 - > HVV couplings: $|f_{a3}| < O(10^{-3})$ at 95% CL in combination with off-shell production
 - > Htt: $|f_{CP}^{Htt}| < 0.55$ at 68% CL
 - > H $\tau\tau$: $\tan(\alpha^{H\tau\tau}) < 0.36$ at 68% CL
- All analyses exclude pure CP-odd hypothesis and agree with SM prediction
- Studies are continuing in Run 3 and we will hopefully get new results about CP violation



Continuing

in Run 3

Searches for CP Violation in the Higgs sector: The Quest to Understand the Baryon Asymmetry in the Universe

Keep looking for more!