

Cosmological constraints obtained with the fourth gravitational wave transients catalog (GWTC-4.0) of the LVK collaboration

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For the LIGO-Virgo-KAGRA collaboration
GDR, Paris, October 2025



1. The data: GWTC-4.0
2. Analysis method
3. Cosmological constraints with GWTC-4.0

LVK publications:
<https://pnp.ligo.org/ppcomm/Papers.html>

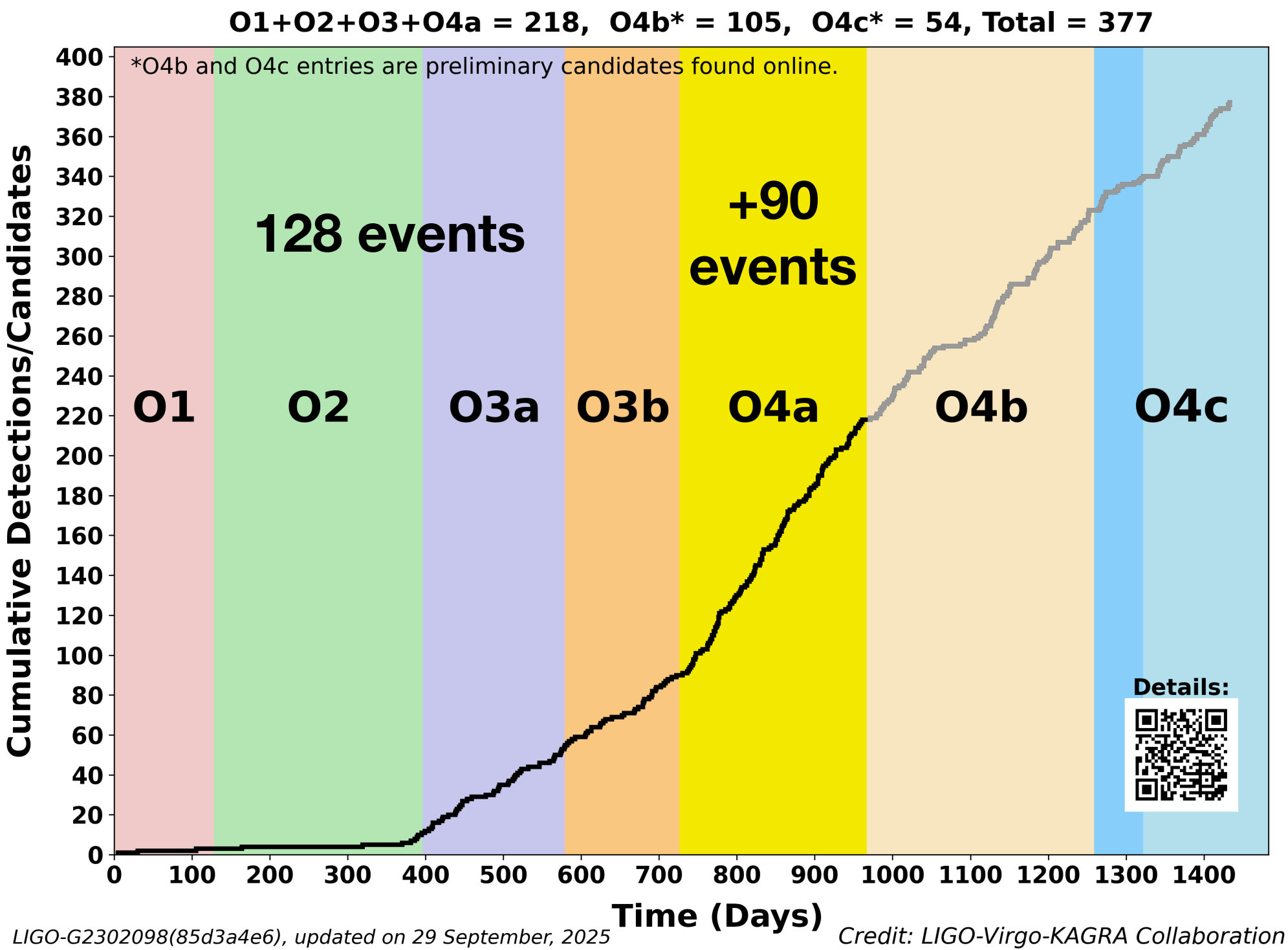
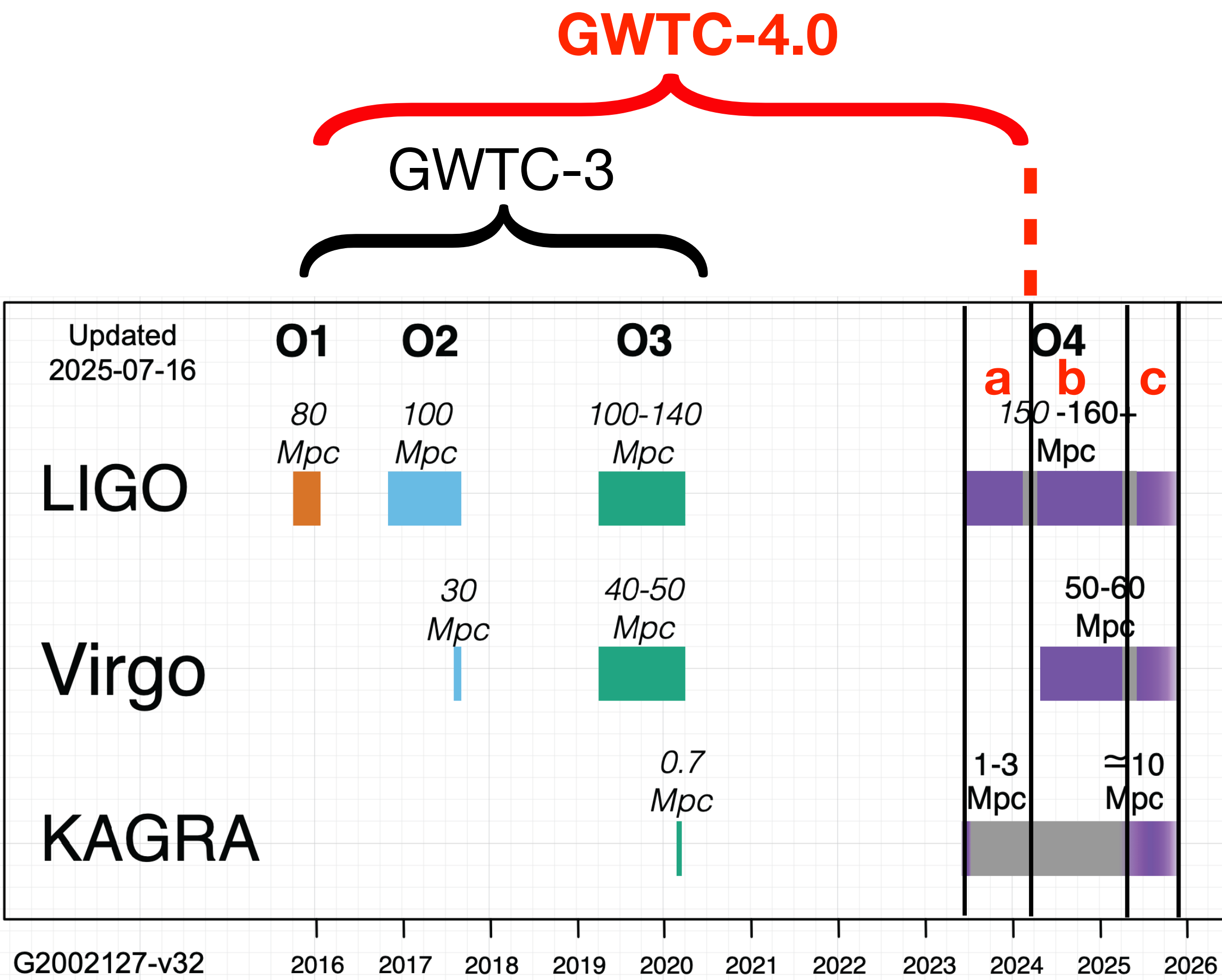
1. The data: GWTC-4.0

arXiv:2508.18080 (intro to GWTC-4.0)
arXiv:2508.18082 (GWTC-4.0 catalog paper)
arXiv:2508.18083 (astrophysical population)

The data: GWTC-4.0

Sept 2015-January 2024

O4a: May 24th, 2023 - January 16th, 2024



Total of 218 candidates in GWTC-4.0

Cosmology with GWs

From the GW signal alone, we can estimate the distance to the source:

$$d_L \propto \langle h_{\text{GW}} \rangle^{-1} \tau_{\text{coal}}^{-1} f_{\text{GW}}^{-2}$$

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In the Big-Bang model, d_L is related to cosmology.

In particular, assuming the flat Λ CDM model:

$$d_L = \frac{(1+z)c}{H_0} \int_0^z \frac{dz'}{\sqrt{\Omega_m(1+z')^3 + \Omega_\Lambda}} \quad z \ll 1 \quad \approx \quad \frac{cz}{H_0}$$

We measure d_L : any information on z provides constraints on H_0

$$H_0 = c \frac{z}{d_L}$$

← redshift of the host galaxy

← given by the GW signal

2. Analysis method

Hierarchical bayesian inference

$$p(m_{1,s}, m_{2,s}) p(\vec{s}_1, \vec{s}_2) R(z) \dots$$

second inference to estimate
the hyper-parameters Λ describing the full population

Λ : H_0 , parameters describing the mass spectra...

15-20 parameters, jointly inferred

$$\theta_i = \begin{pmatrix} m_{1,\text{det}} \\ m_{2,\text{det}} \\ d_L \\ \vec{s}_1 \\ \vec{s}_2 \\ \dots \end{pmatrix} \quad \begin{pmatrix} m_{1,\text{det}} \\ m_{2,\text{det}} \\ d_L \\ \vec{s}_1 \\ \vec{s}_2 \\ \dots \end{pmatrix} \quad \bullet \quad \bullet \quad \bullet \quad \begin{pmatrix} m_{1,\text{det}} \\ m_{2,\text{det}} \\ d_L \\ \vec{s}_1 \\ \vec{s}_2 \\ \dots \end{pmatrix}$$

GW #1 GW #2 GW #n

first inference

(parameter estimation, PE)
to estimate θ , 11 parameters, the
individual properties of each event

Hierarchical inference

For a single CBC event:

$$\mathcal{L}(x_{\text{GW}}|\Lambda) \propto \frac{\int \mathcal{L}(x_{\text{GW}}|z, \theta, \Lambda) p(z|\Lambda) p(\theta|\Lambda) dz d\theta}{\int_{\rho \geq \rho_{\text{thr}}} p(\text{det}|z, \theta, \Lambda) p(z|\Lambda) p(\theta|\Lambda) dz d\theta}$$

PE

choice of models
($z, m_s, \text{spins, merger rate...}$)

marginalise over the
true (unkown) parameters
of the CBC

the denominator corrects for
the selection effect

Bright siren

associated EM counterpart,
 $z^{\text{GW}}=z_{\text{host}}^{\text{EM}}$, only 1 event since
2015

$$p(z|\Lambda) = \delta(z - z_{\text{host}}^{\text{EM}})$$

Spectral siren

structures in the detector-frame
mass spectra, agnostic z prior

$$m_{\text{det}} = (1 + z)m_s$$

uniform
comoving

$$p(z|\Lambda) \propto \frac{dV_c}{dz}$$

Dark siren

structures in the detector-frame
mass spectra, z from a catalog

$$m_{\text{det}} = (1 + z)m_s$$

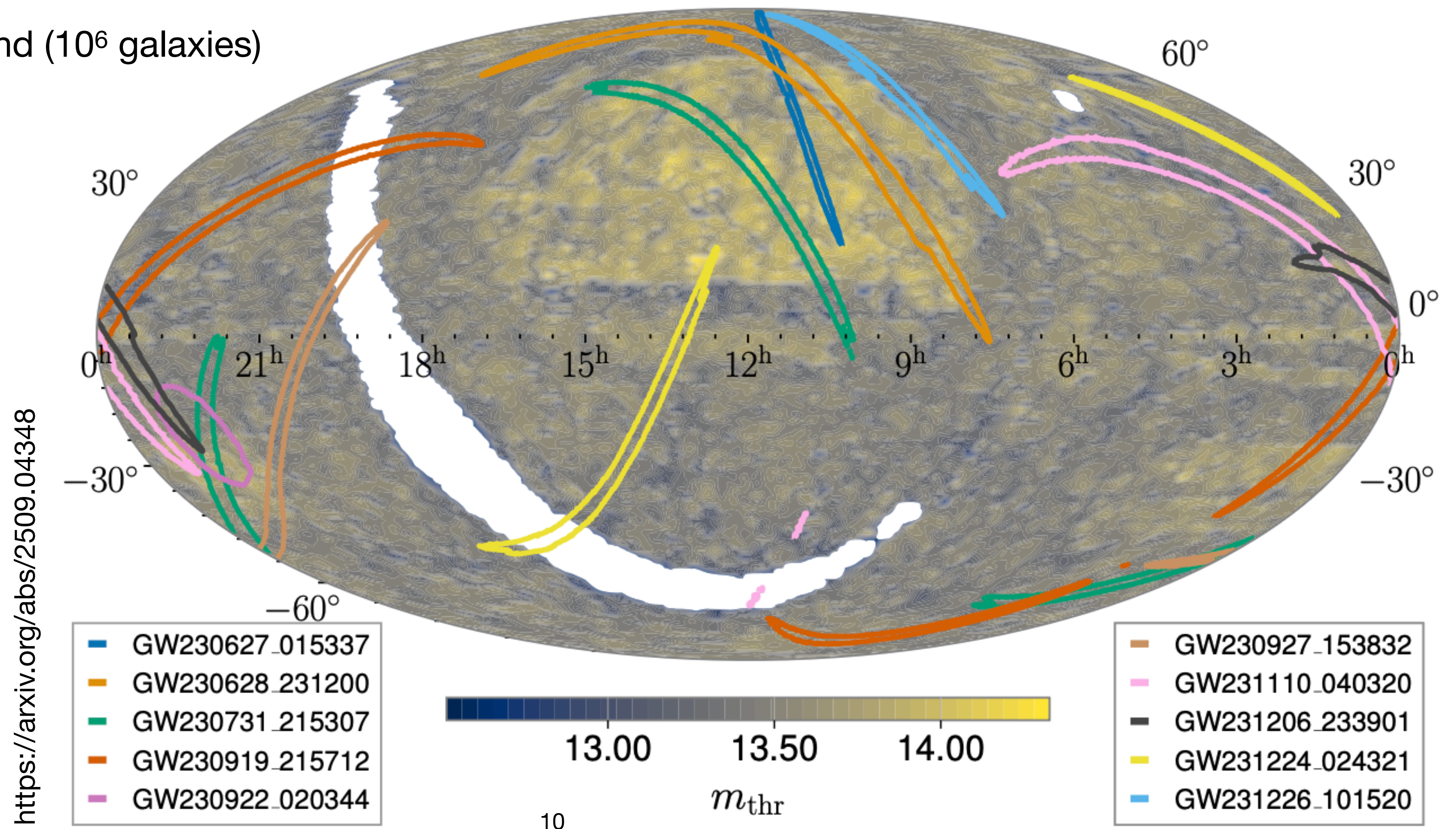
$$p(z|\Lambda) \propto p_{\text{EM catalog}}(z)$$

Galaxy catalog for the dark sirens analysis

Use GLADE+ Dályá et al., MNRAS 479, 2374 (2018) and 514, 1403 (2022):

GLADE+, full sky, 22 millions of galaxies: GWGC, 2MPZ, 2MASS XSC, HyperLEDA, and WISExSCOSPZ galaxy catalogues, and the SDSS-DR16Q quasar catalogue

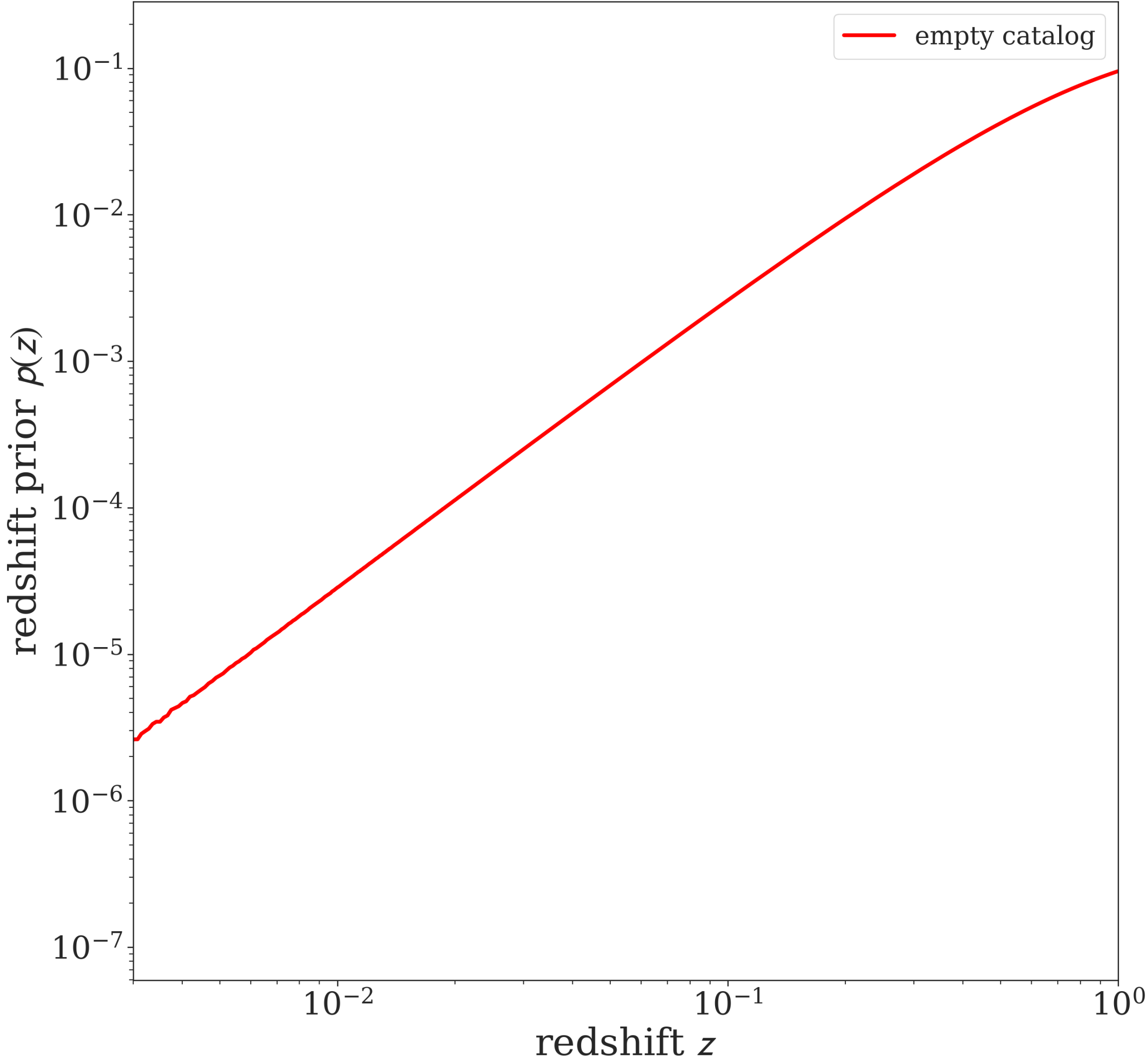
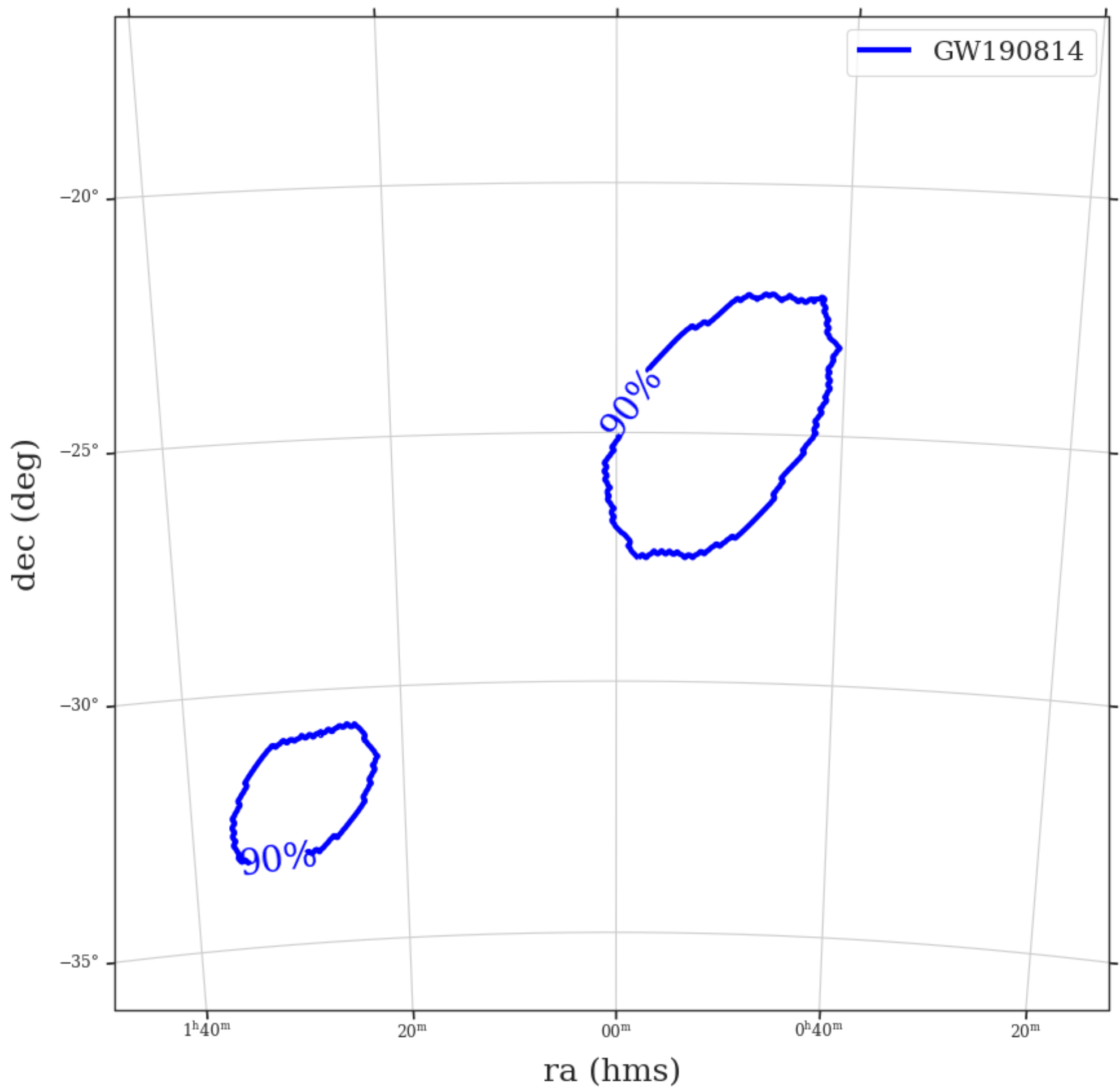
Choose the K-band (10^6 galaxies)



Constructing the redshift prior using the GLADE+ catalog

Dályá et al., MNRAS 479, 2374 (2018) and 514, 1403 (2022)

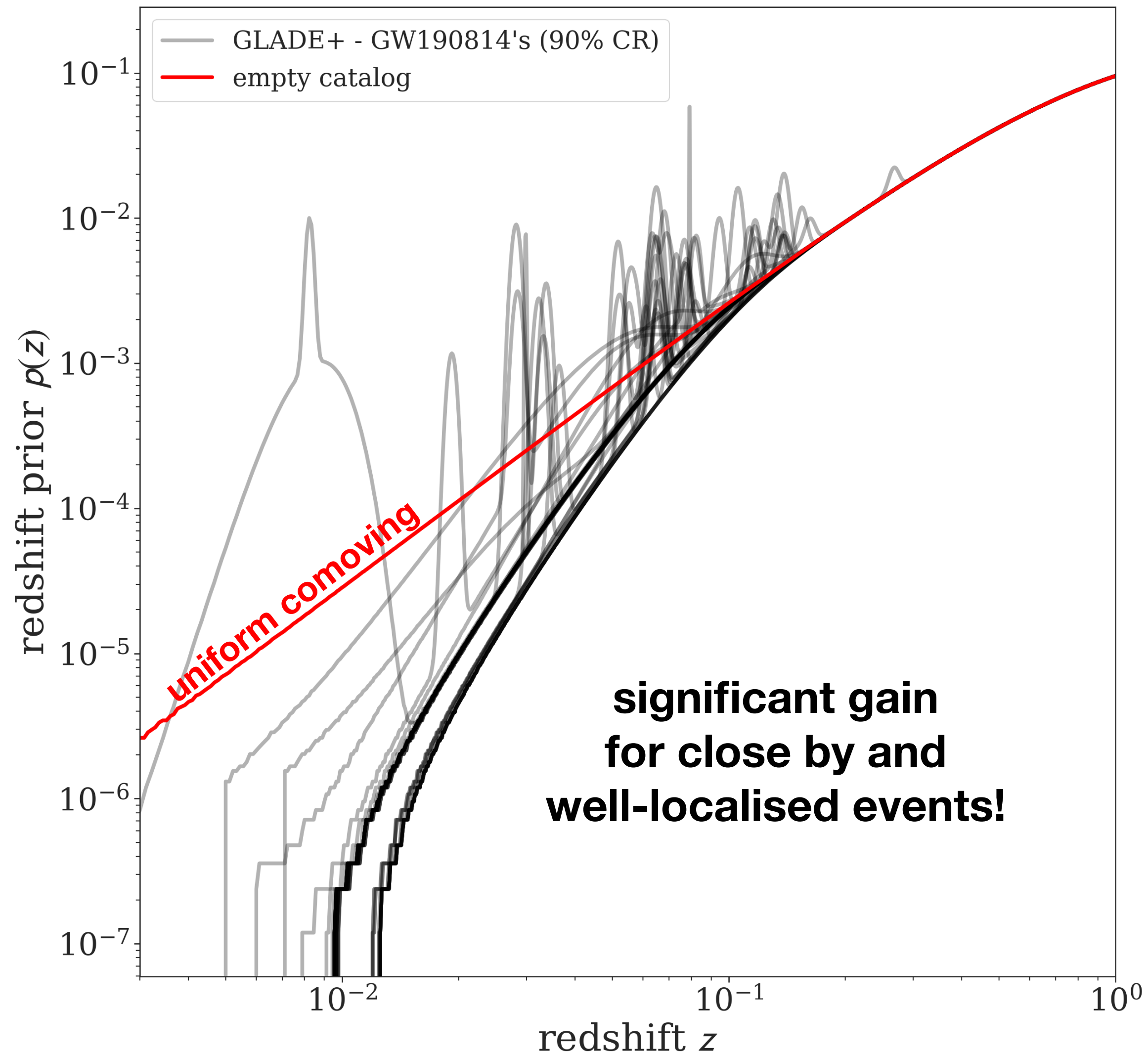
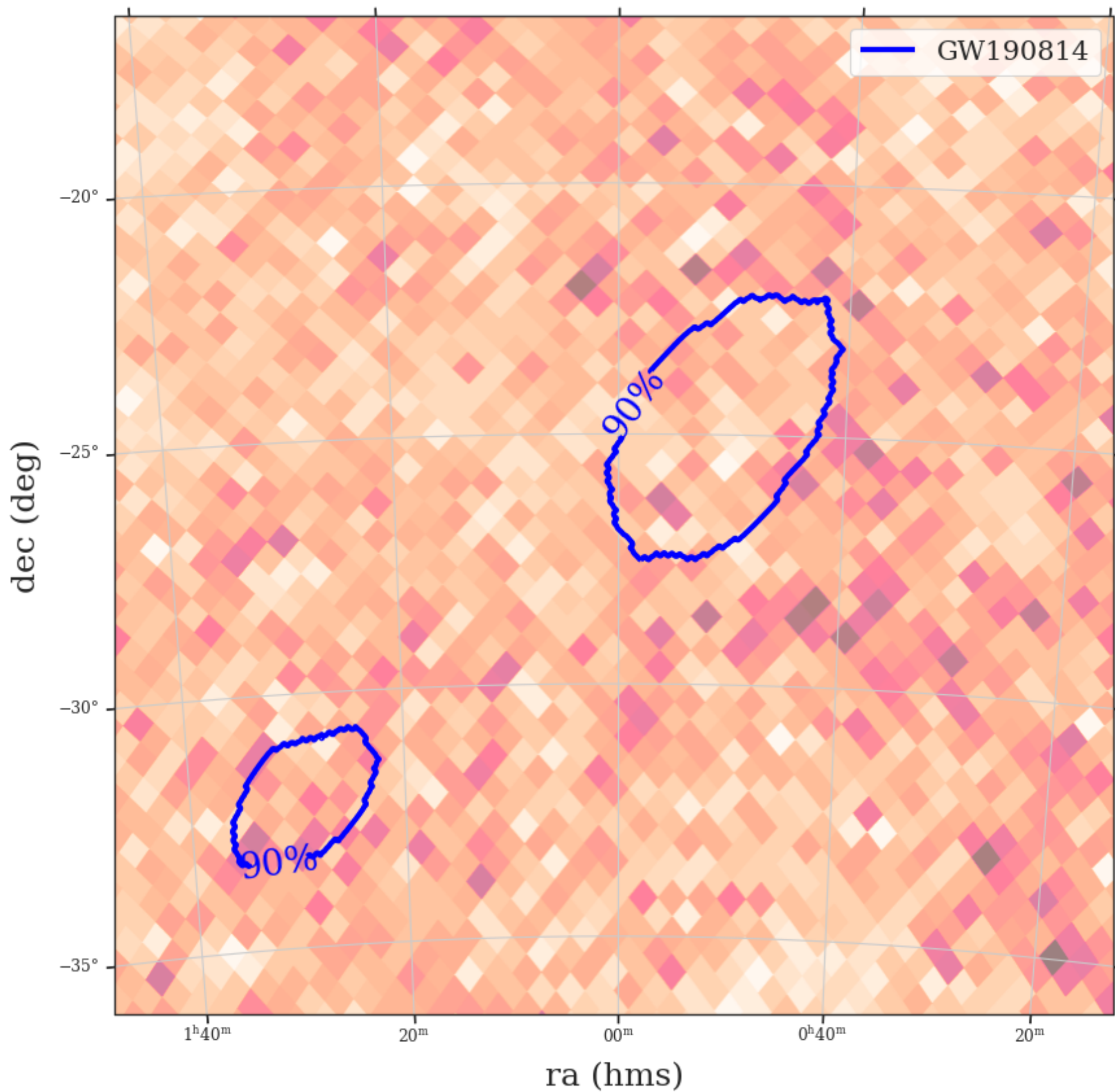
GW190814: $d_L \sim 230$ Mpc



Constructing the redshift prior using the GLADE+ catalog

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GW190814: $d_L \sim 230$ Mpc

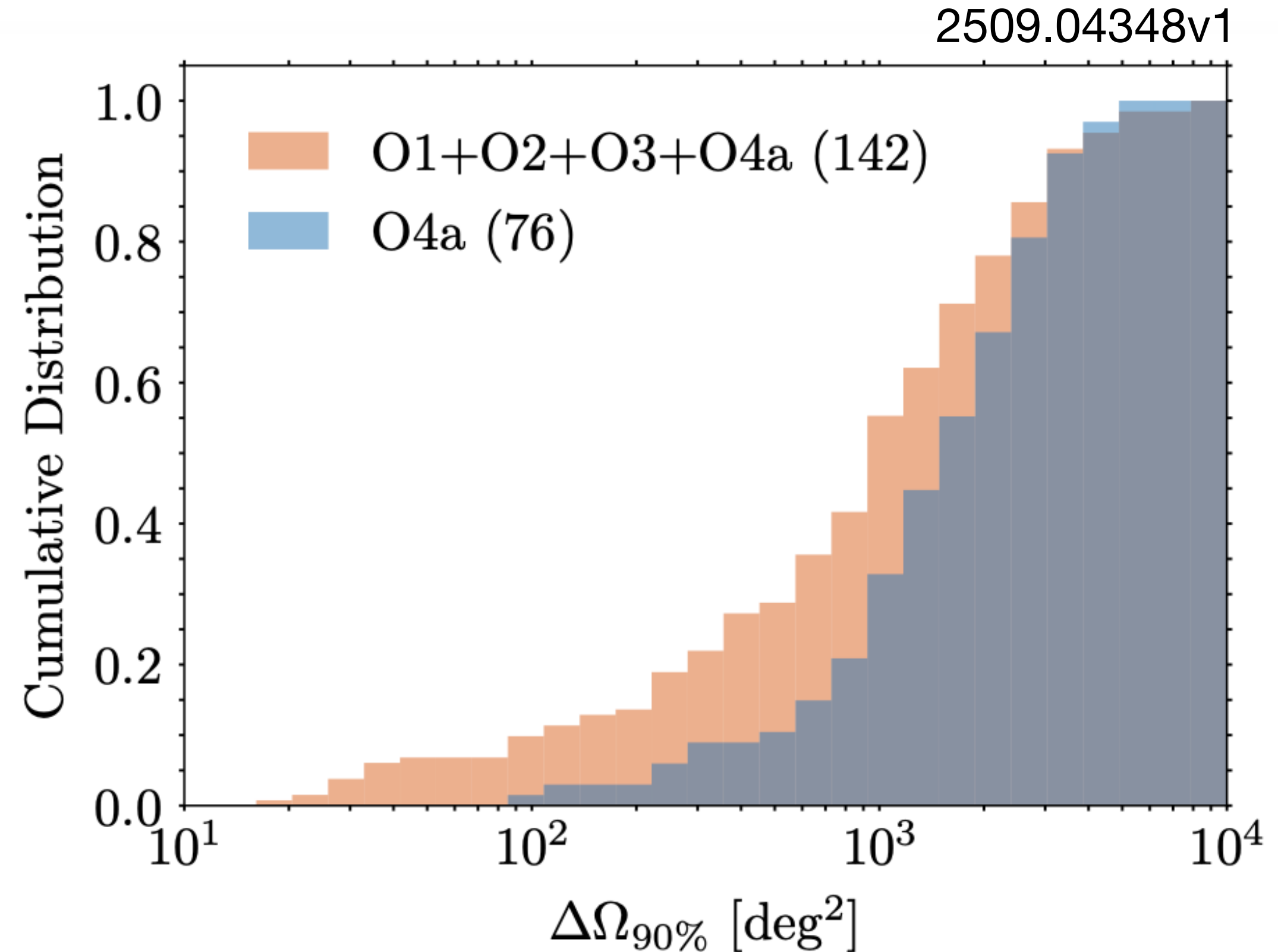


3. Cosmological constraints with GWTC-4.0

with GWTC-3: R. Abbott et al 2023 ApJ 949 76
with GWTC-4.0: arXiv:2509.04348

The data set: **142 CBCs, 141 dark sirens + the bright siren GW170817**
using a False Alarm Rate $< 1/ (4 \text{ years})$

Sky localisation is worse during O4a
(Virgo not running):
don't expect a large difference
between the spectral/dark analyses



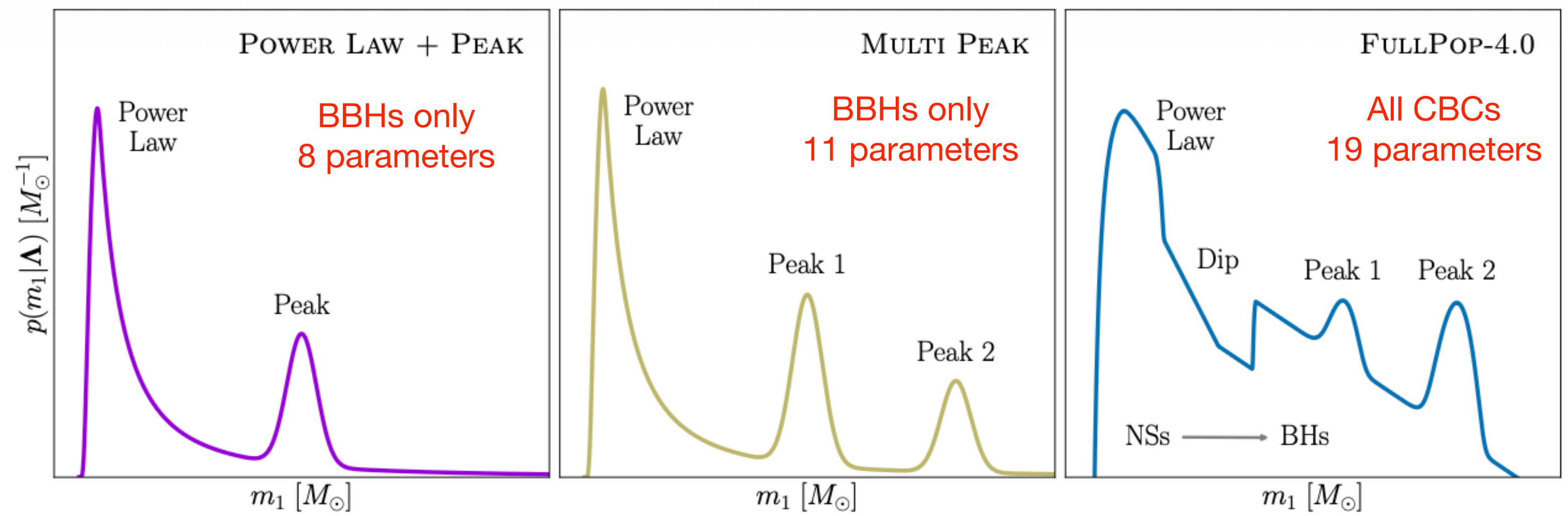
Two independent analysis pipelines:

- **gwcsmo** Gray et al, JCAP 12 (2023) 023, Gray et al, MNRAS 512 (2022) 1127
- **icarogw** Mastrogiovanni et al, A&A 682 (2024) A167 Mastrogiovanni et al, PRD, 104, 062009



We combine their results

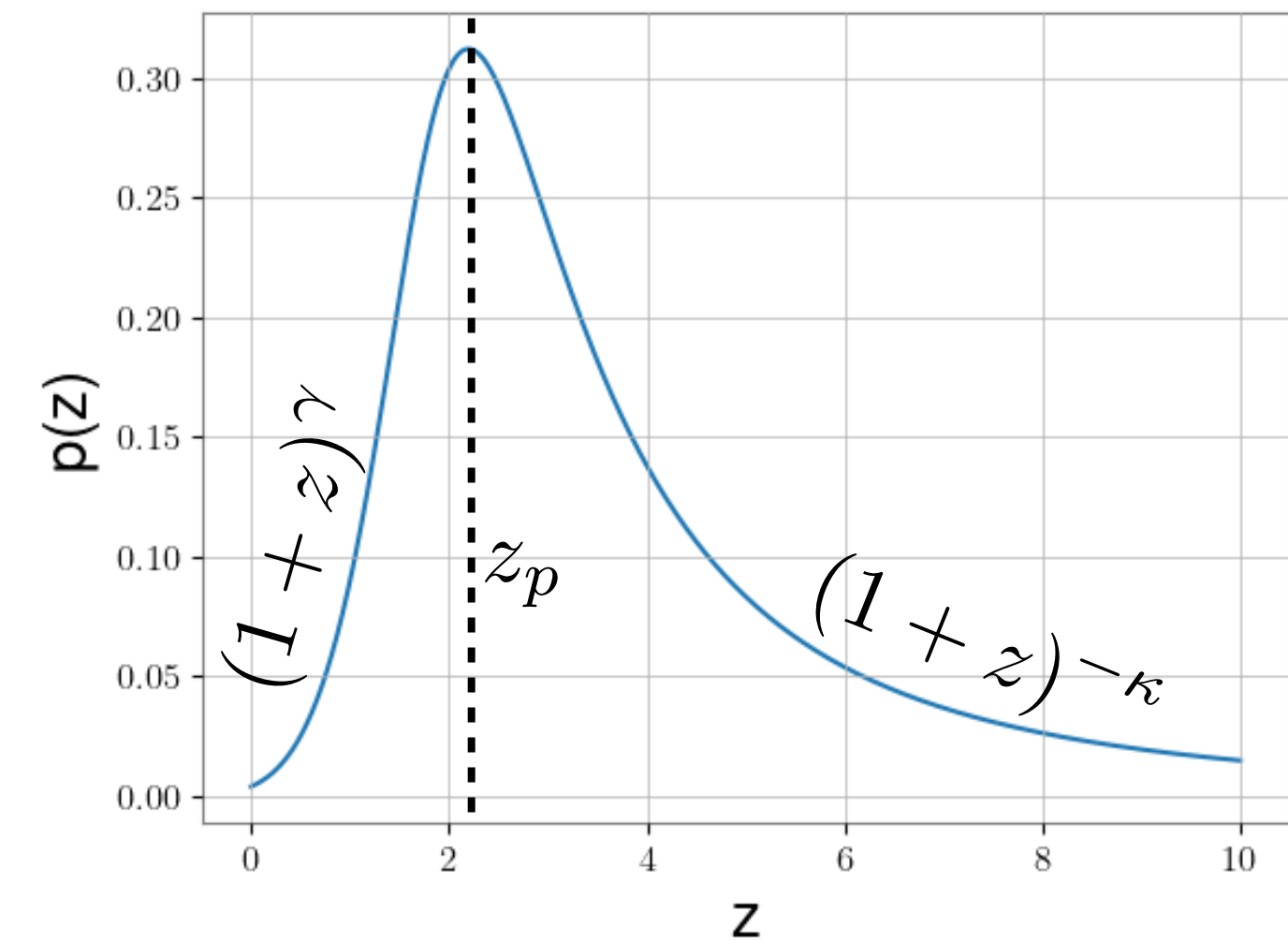
Choice of mass models:



Choice of a generic merger-rate model, Madau-Dickinson (SFR):

$$p(s|z) = R_0 \psi(z|\gamma, \kappa, z_p) = R_0 \frac{\left(1 + \frac{1}{(1+z_p)^{\gamma+\kappa}}\right) (1+z)^\gamma}{1 + \left(\frac{1+z}{1+z_p}\right)^{\gamma+\kappa}} \text{Gpc}^{-3} \text{yr}^{-1}$$

3 parameters



Choice of the redshift prior: spectral or dark analysis

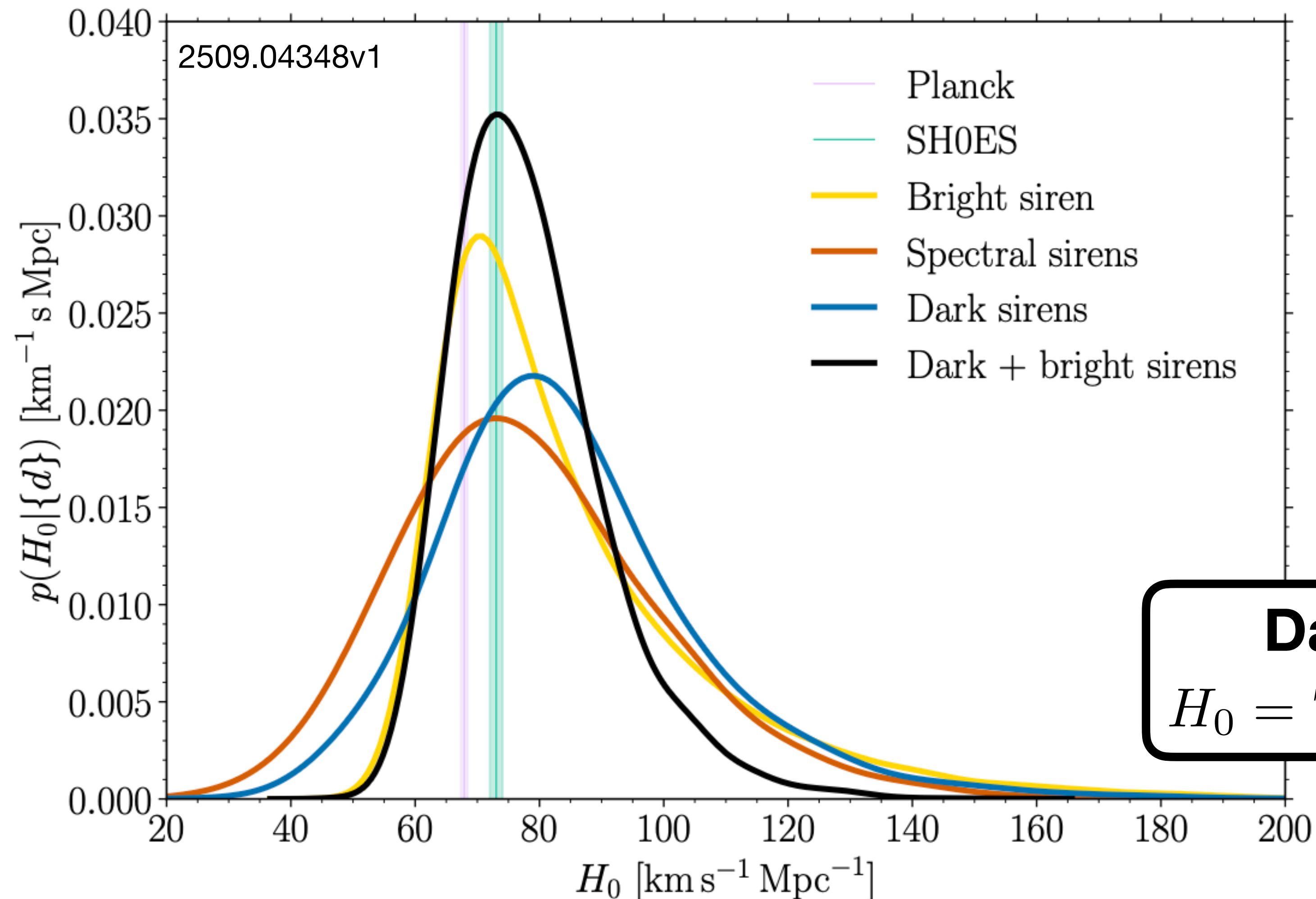
Choice of the host galaxies/luminosity relation:

$$p(\text{hosting CBC}|L) \propto 1$$

$$p(\text{hosting CBC}|L) \propto L$$

H_0 from GWTC-4.0

H_0 posterior after marginalisation over 22 hyper-parameters (mass and rate models)



GW170817 alone:

$$H_0 = 78.4^{+25.7}_{-12.0} \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Dark sirens only:

$$H_0 = 81.6^{+21.5}_{-15.9} \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Spectral sirens only:

$$H_0 = 76.4^{+23.0}_{-18.1} \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Dark+bright sirens:

$$H_0 = 76.6^{+13.0}_{-9.5} \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Modified gravity

Many modified-gravity models:

- modify the GW propagation equation (friction)
- change the GW amplitude wrt GR
- change the attributed distance d_L^{GW}

The ratio $d_L^{\text{GW}} / d_L^{\text{EM}}$ is a convenient probe of departures from GR.

We consider 2 parameterisations:

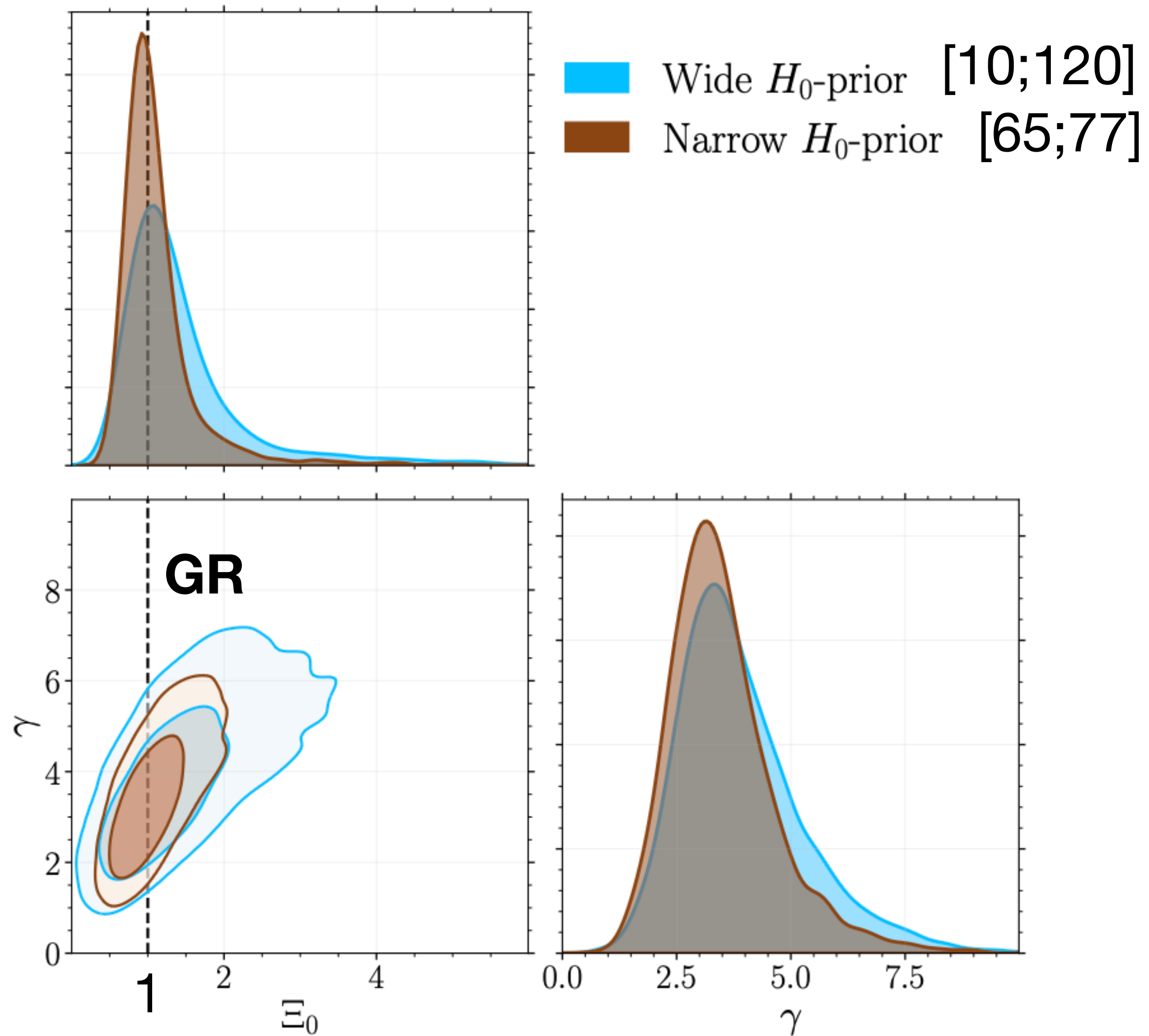
$$\frac{d_L^{\text{GW}}}{d_L^{\text{EM}}} = \begin{cases} \Xi_0 + \frac{1-\Xi_0}{(1+z)^n} & \text{Phenomenological approach} \\ \exp \left(\int_0^z \frac{c_M dz'}{E^2(z')(1+z')} \right) & \text{Horndeski gravity, running Planck mass} \end{cases}$$

(Belgacem, E., et al. 2019b, JCAP, 07, 024)

(Horndeski, G. W. 1974, Int. J. Theor. Phys., 10, 363)

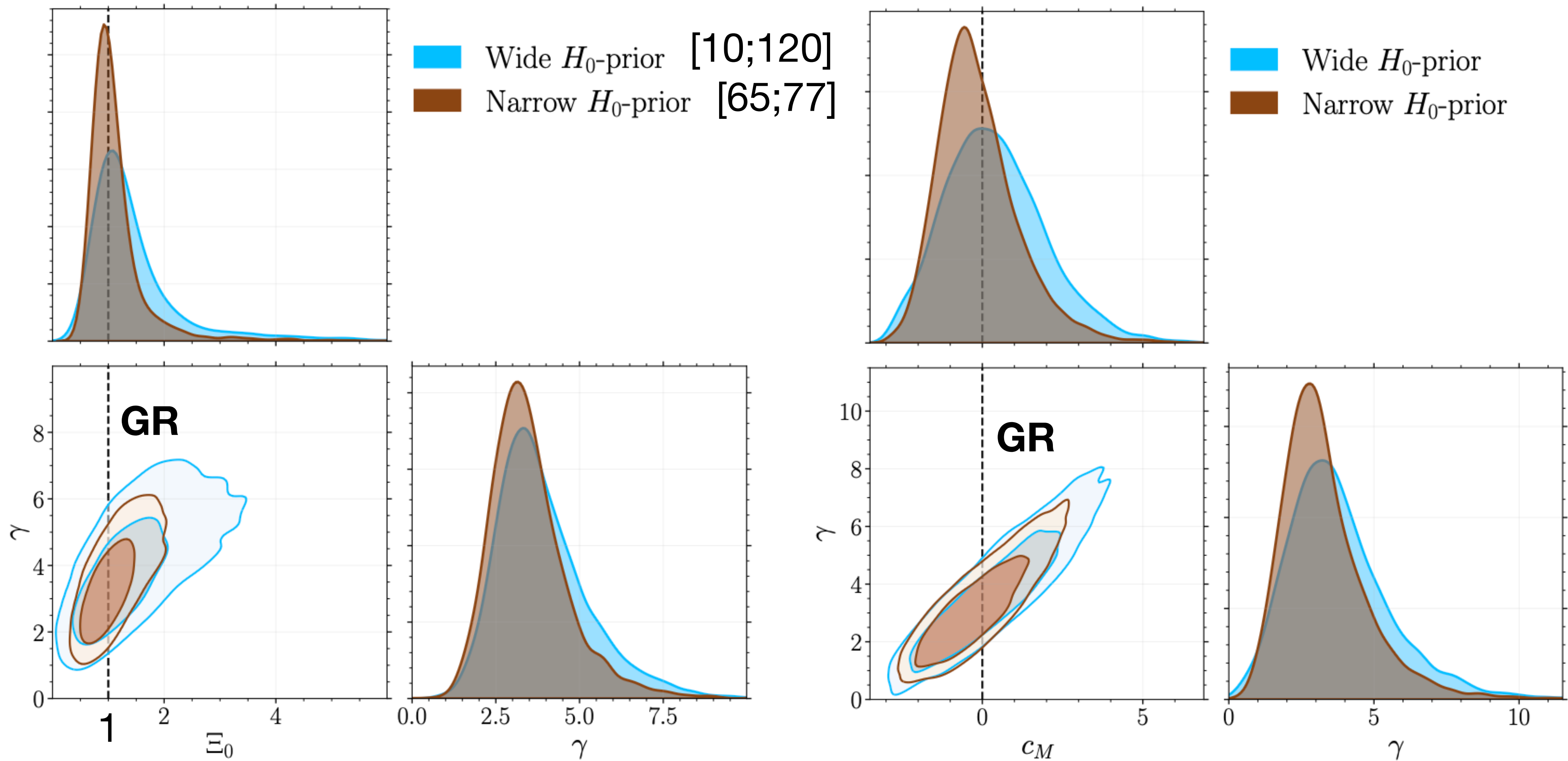
$$\text{GR: } \Xi_0 = 1 \text{ and } c_M = 0$$

Modified gravity



Strong correlation with γ as the detectability of GWs can be due to a change of d_L^{GW} or to a change in the merger rate (γ at low z): $p(s|z) \propto (1+z)^\gamma$

Modified gravity



Strong correlation with γ as the detectability of GWs can be due to a change of d_L^{GW}
or to a change in the merger rate (γ at low z): $p(s|z) \propto (1+z)^\gamma$

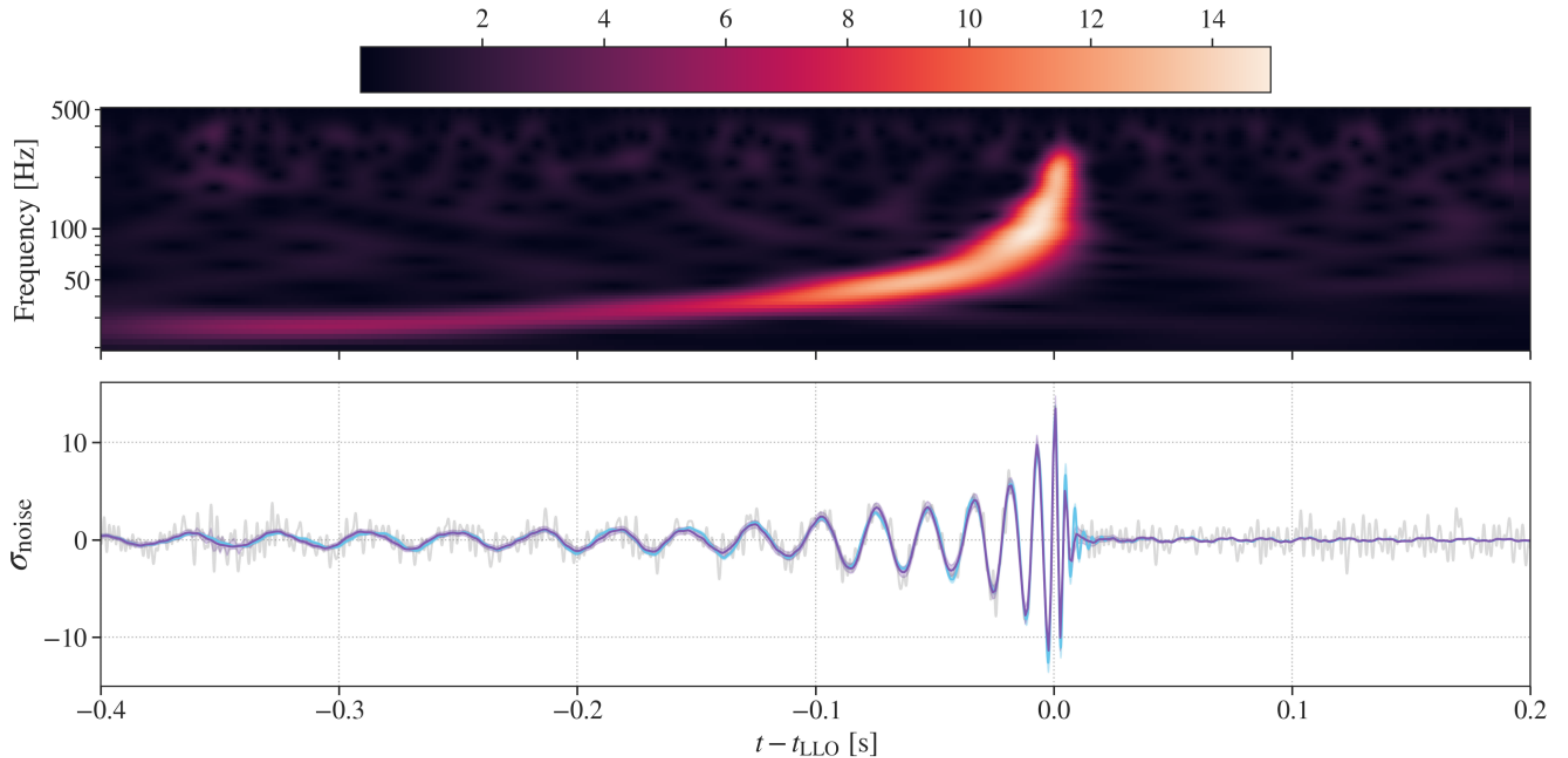
Conclusions

- Our best constraint (68%) with GWTC-4.0 on H_0 is $H_0 = 76.6^{+13.0}_{-9.5} \text{ km s}^{-1} \text{ Mpc}^{-1}$
- Improvement vs GWTC-3: 60% better for the spectral analysis (~3.3x more events + better mass model)
- Not yet competitive with the CMB/Standard candles measurements
- No departure from GR after tests on some modified gravity models
- Virgo was NOT participating in O4a: expect better constraints with O4b
- UpGLADE: add Pan-STARRS DR2, PS-STRM, CatWISE, Legacy Survey, Duncan22, much more complete catalog

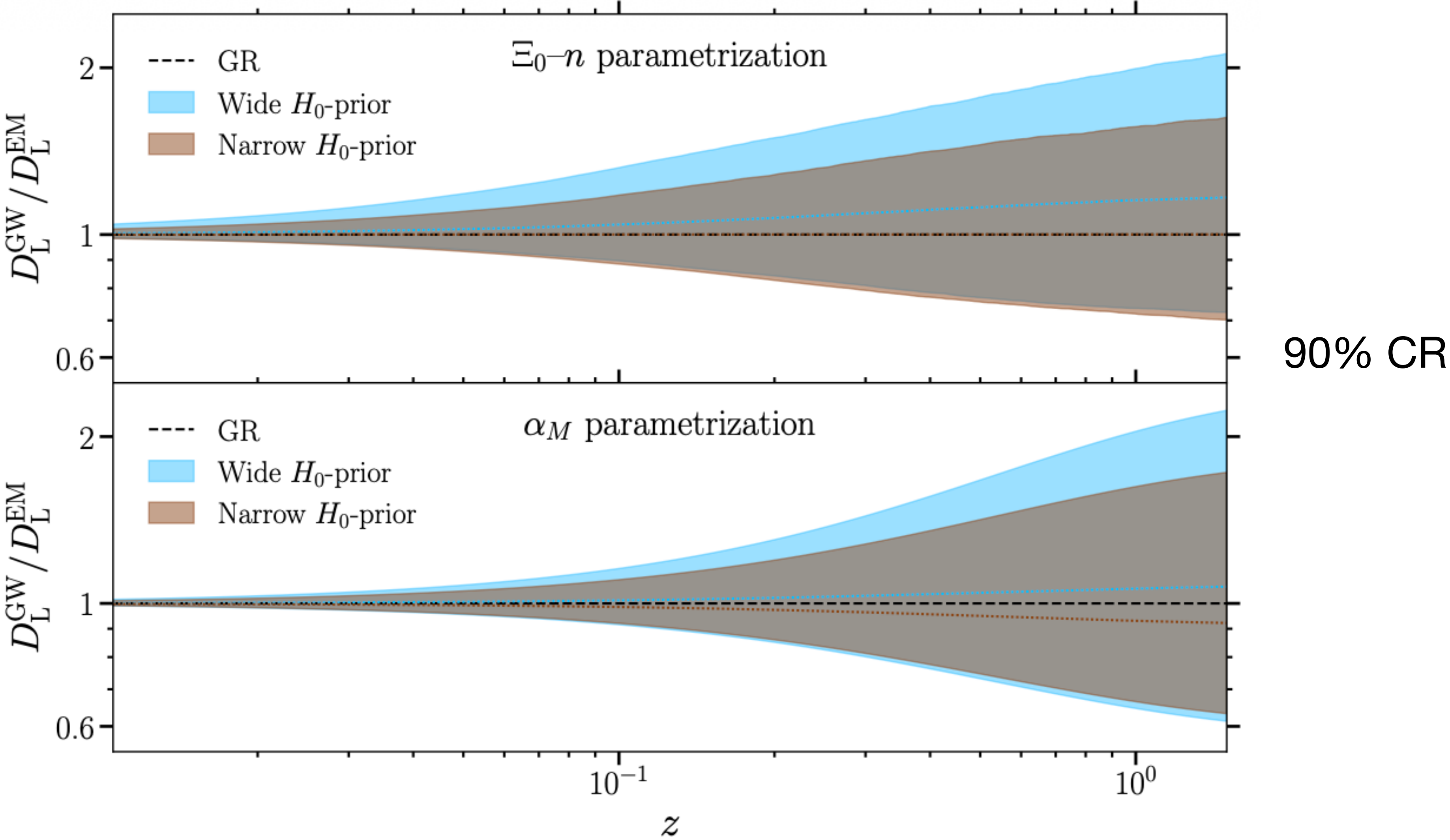
Backup

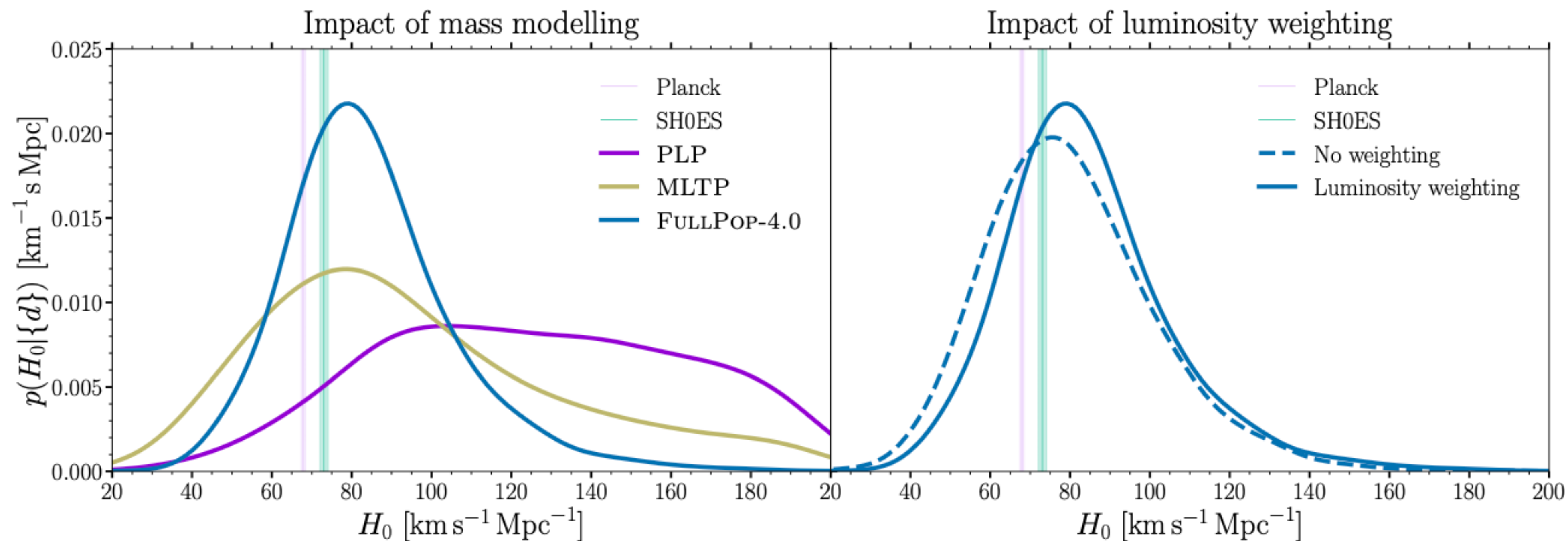
The data: GWTC-4.0

Largest SNR so far: GW230814_230901
SNR = 42.4



Modified gravity





Huge impact of mass modelling
the PLP model is strongly
disfavored by the Bayes factor

$$\log_{10} \mathcal{B} = 2.3$$

Small impact of luminosity
weighting

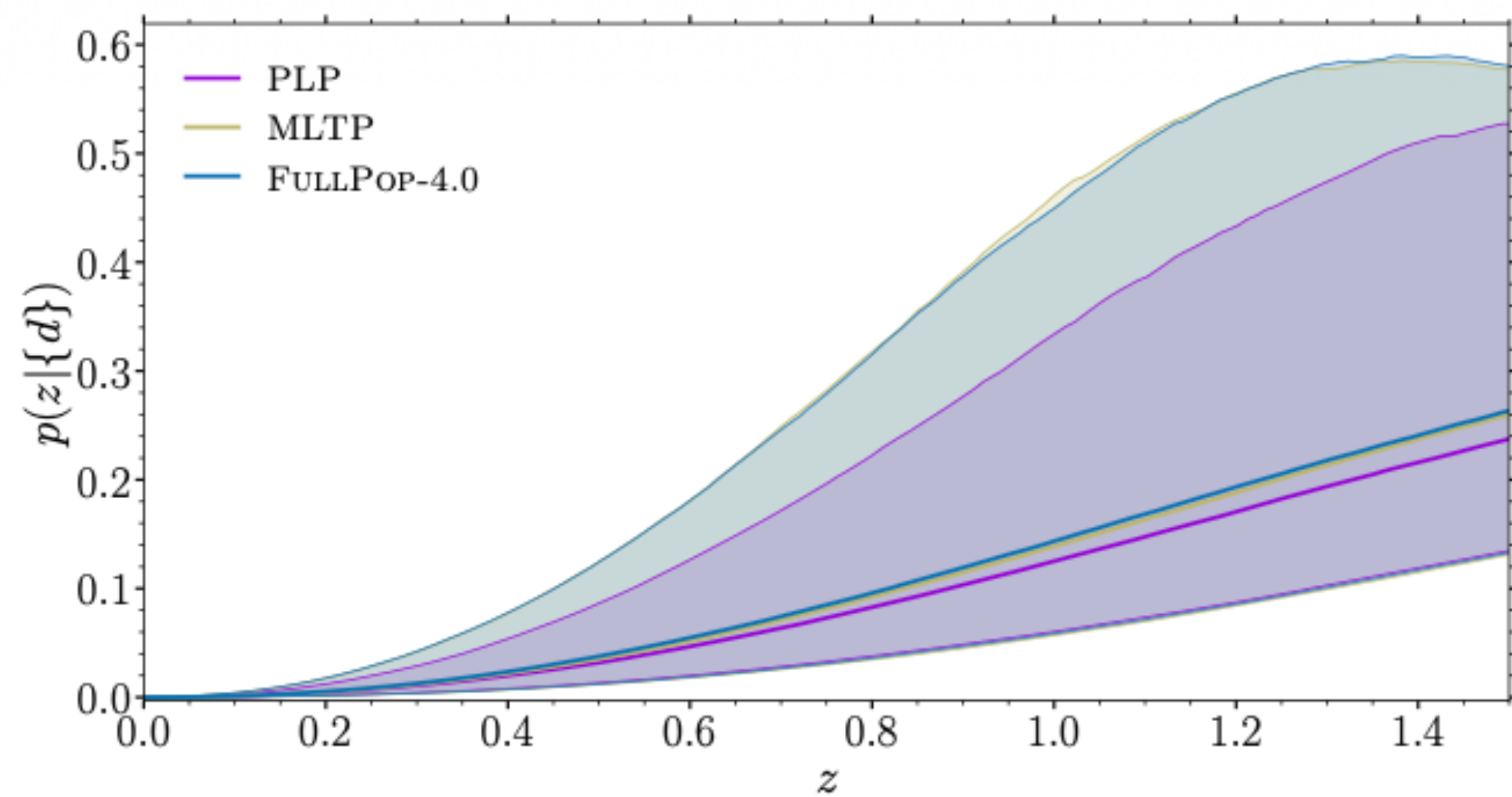
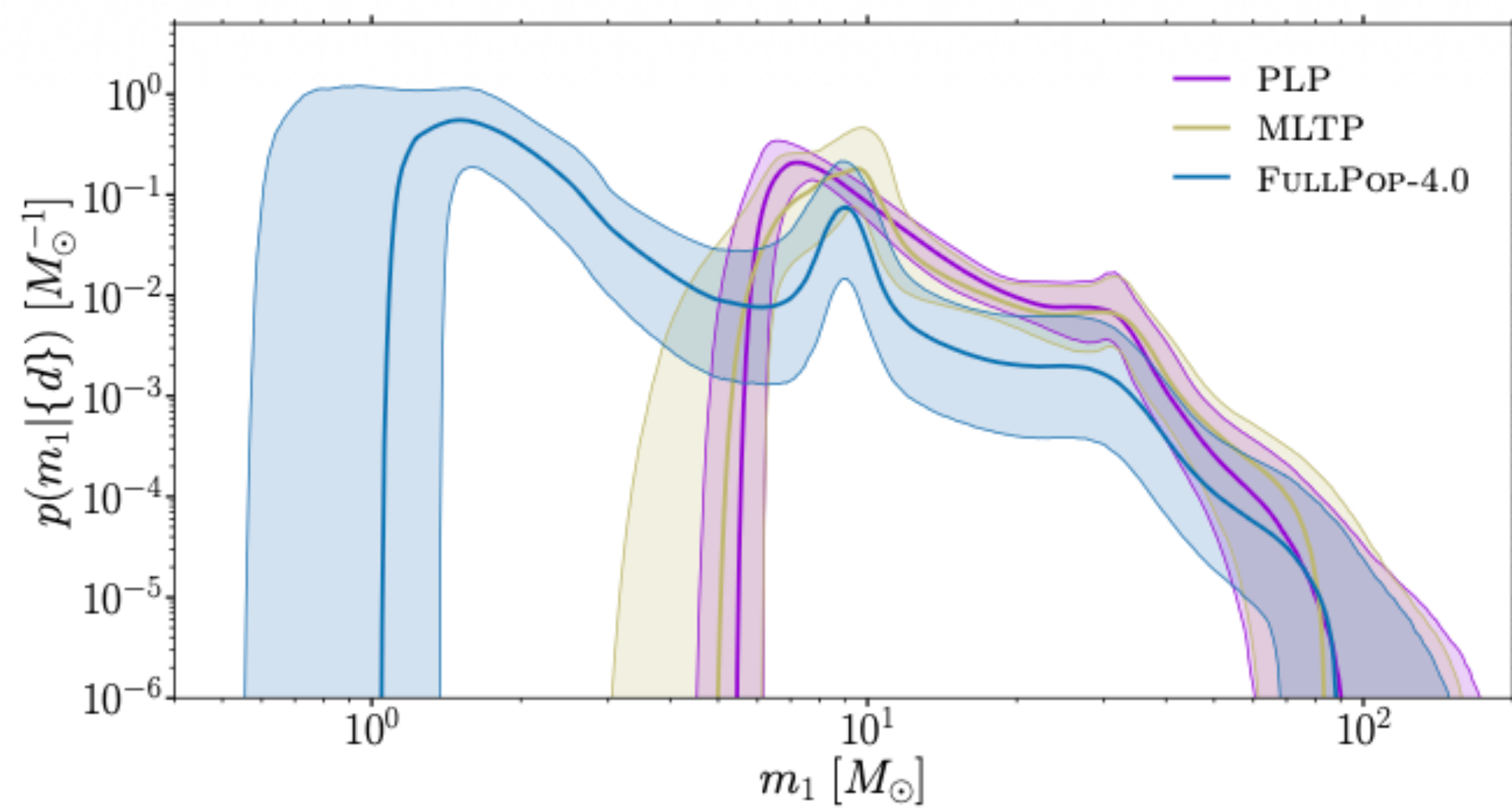
$$p(s|M) \propto L(M)^\epsilon$$

$$\text{no weighting: } \epsilon = 0$$

$$\text{luminosity weighting: } \epsilon = 1$$

no significant evidence for CBCs to occur in
more luminous galaxies in the present data

Dark siren, GR, eps=1



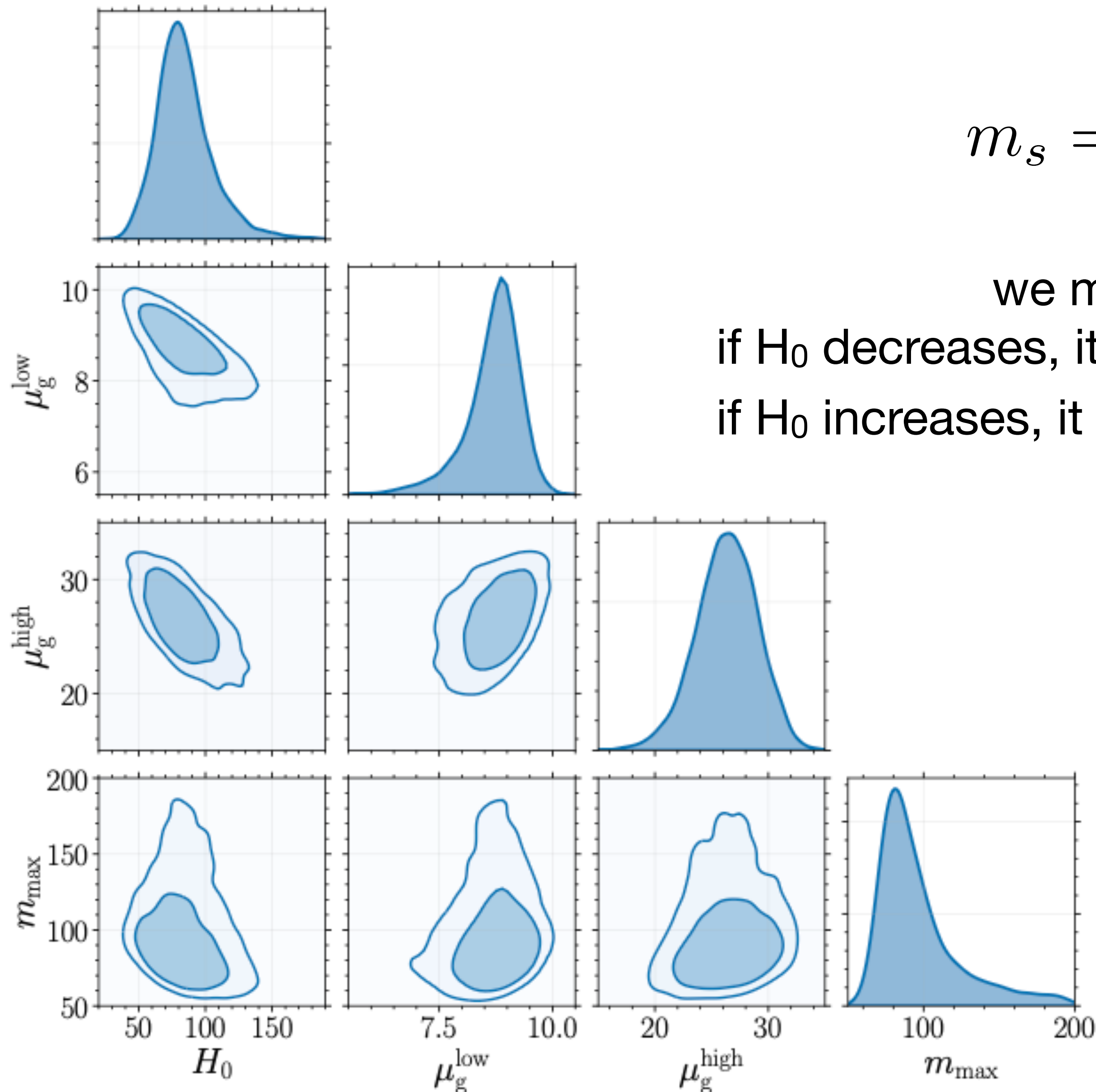
Dark siren, GR, eps=1

$$m_s = \frac{m_d}{1+z} = \frac{m_d}{1 + \frac{H_0 d_L}{c}}$$

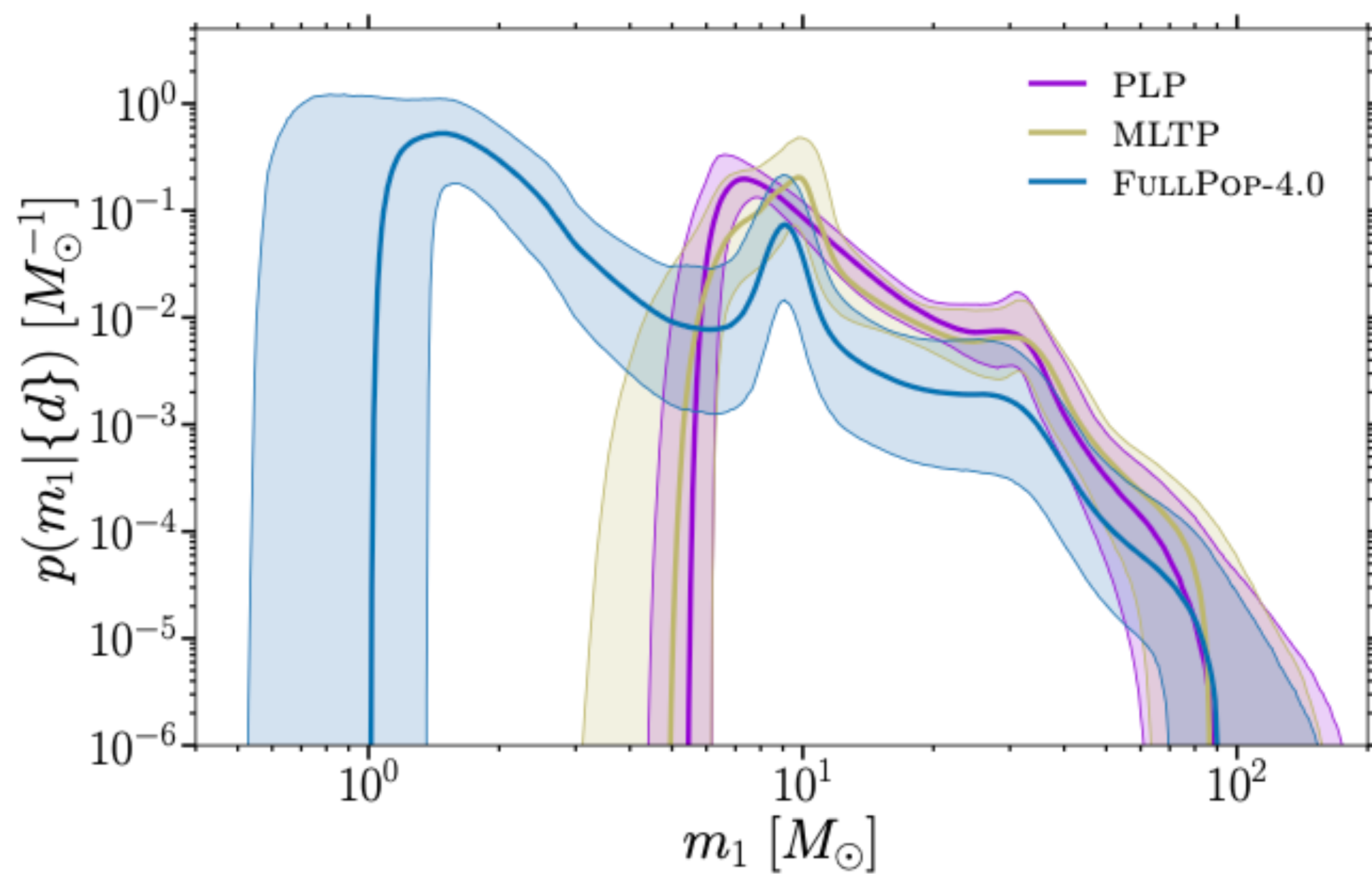
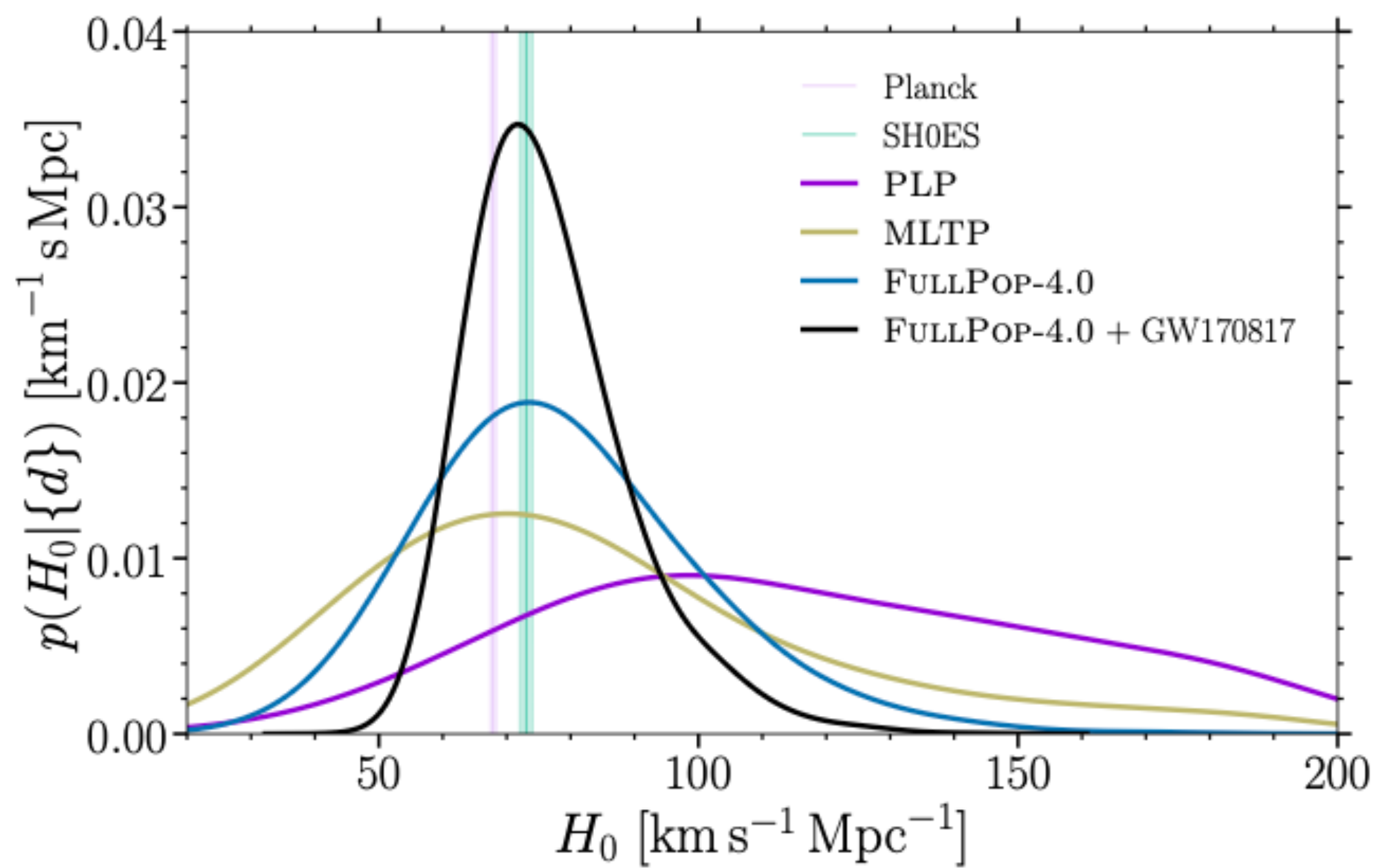
we measure m_d and d_L

if H_0 decreases, it gives a larger m_s (μ increases)

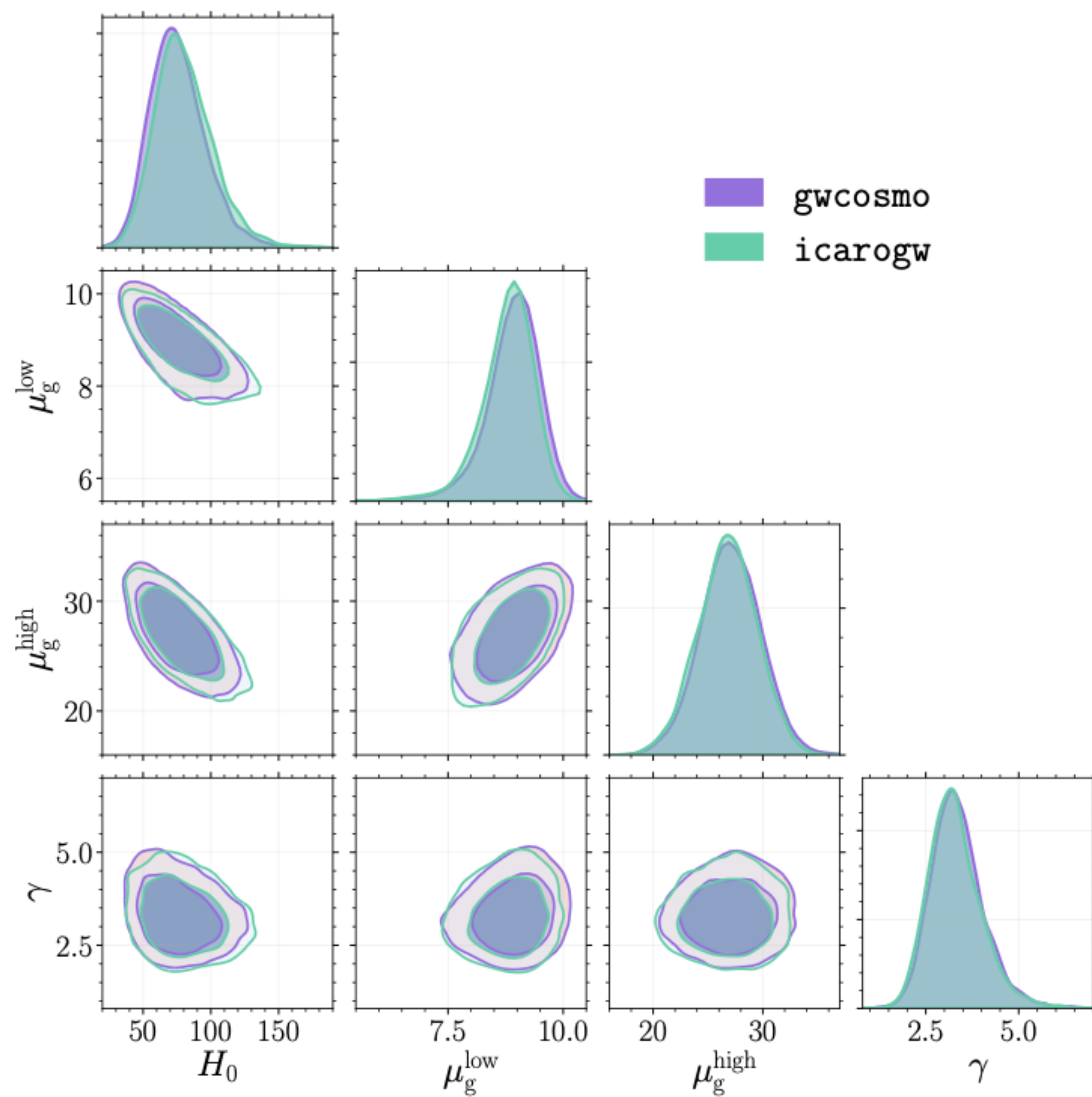
if H_0 increases, it gives a larger m_s (μ decreases)



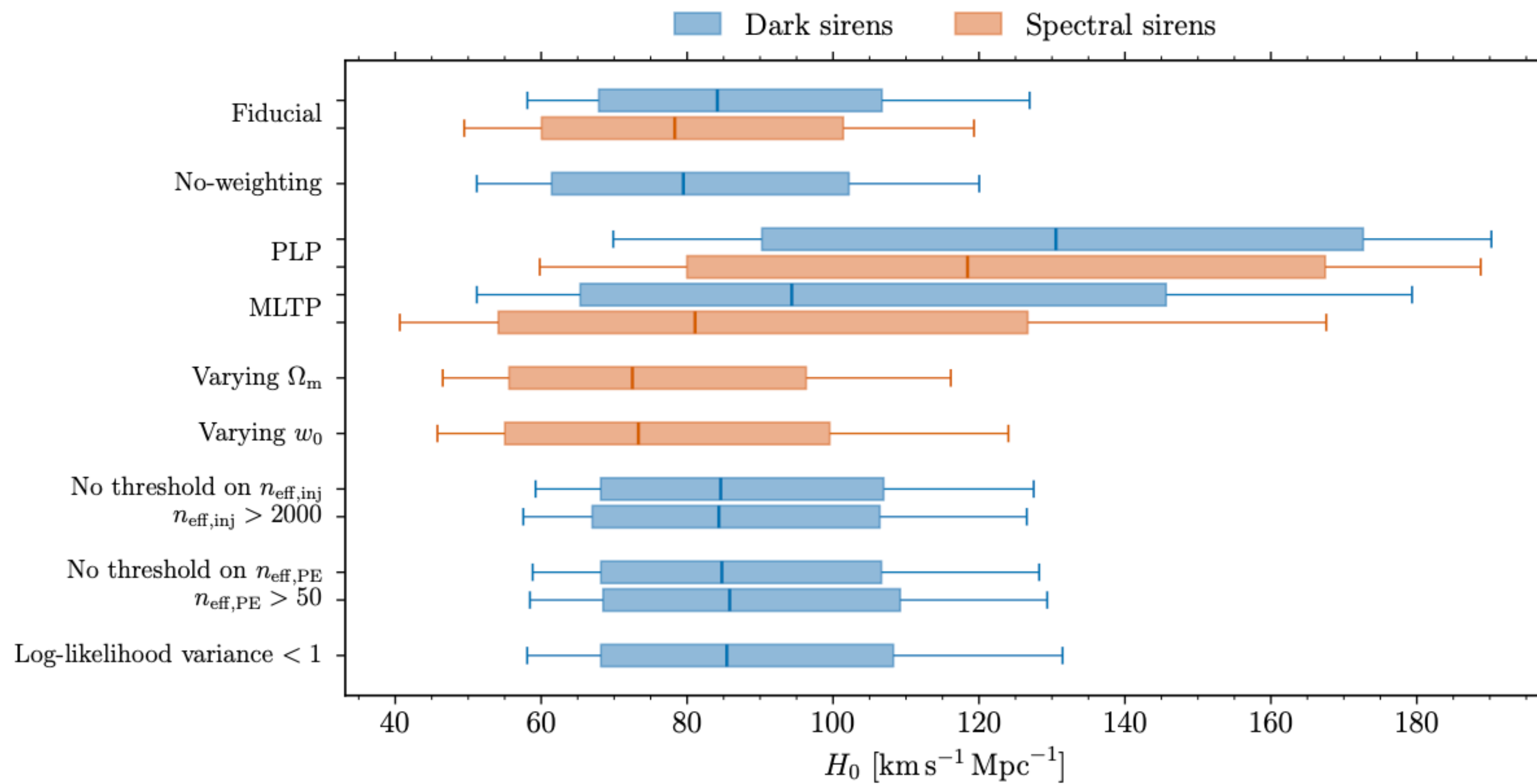
Spectral sirens - GR

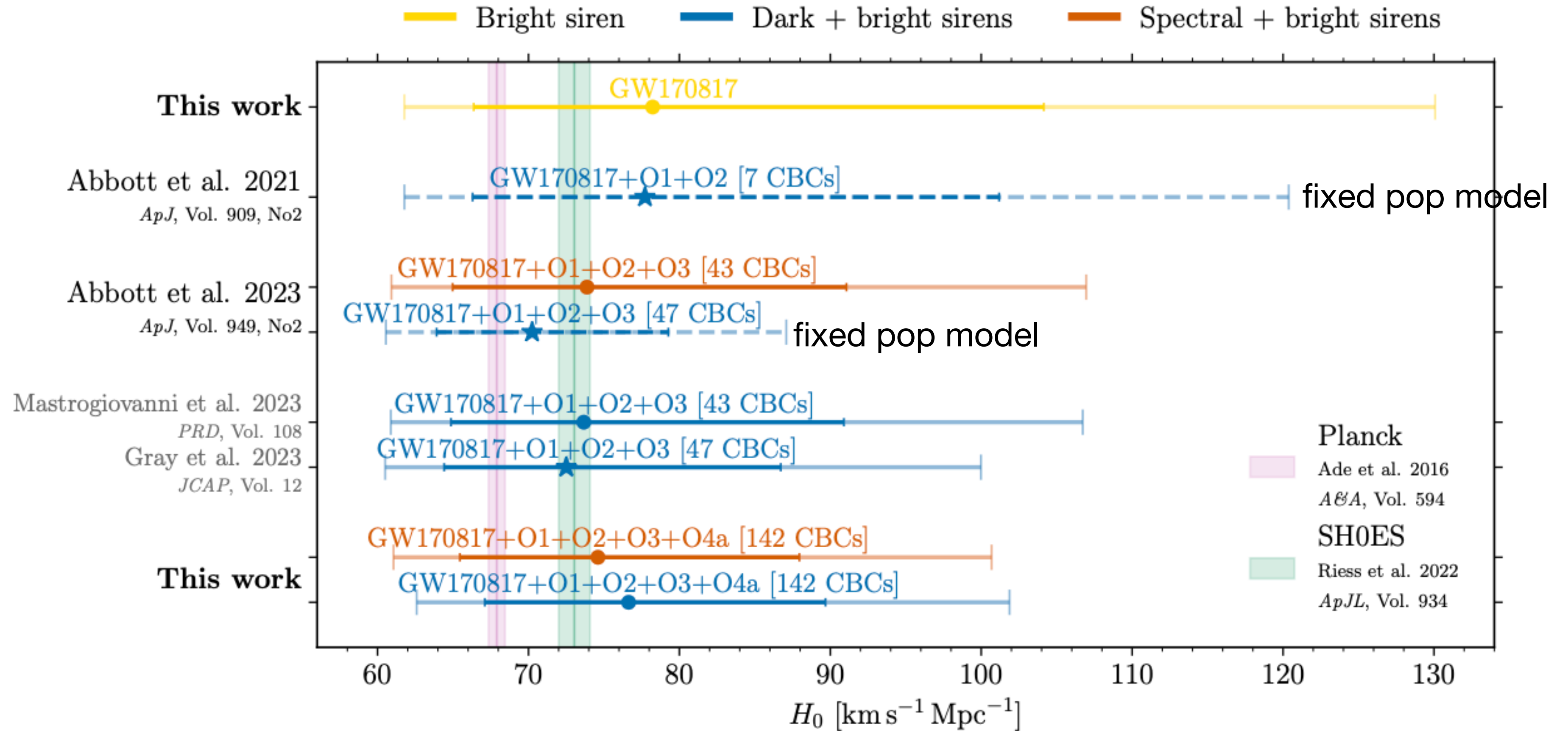


Spectral sirens - GR



Robustness tests





The current result corresponds to a marginalisation over all population and rate parameters, contrary to Abbott et al 2023 (47 CBCs only but using a fixed population and rate model)

The Line of Sight redshift prior: the full expression

$$p(z|\Omega_i, \Lambda, s, I) = \frac{1}{p(s|\Omega_i, \Lambda, I)} \left[p(G|\Omega_i, \Lambda, I) \frac{1}{N_{\text{gal}}(\Omega_i)} \sum_k^{N_{\text{gal}}(\Omega_i)} p(z|\hat{z}_k) p(s|z, M(z, \hat{m}_k, \Lambda), \Lambda, I) \right. \\ \left. + \left(\Theta[z_{\text{cut}} - z] \int_{M(z, m_{\text{th}}(\Omega_i), \Lambda)}^{M_{\text{max}}(H_0)} p(z, M|\Lambda, I) p(s|z, M, \Lambda, I) dM \right. \right. \\ \left. \left. + \Theta[z - z_{\text{cut}}] \int_{M_{\text{min}}(H_0)}^{M_{\text{max}}(H_0)} p(z, M|\Lambda, I) p(s|z, M, \Lambda, I) dM \right) \right].$$

<https://doi.org/10.48550/arXiv.2308.02281>

$$\Lambda = \{\Lambda_{\text{cosmo}}, \Lambda_{\text{mass}}, \Lambda_{\text{rate}}\}$$

if no evolution of the mass models with z ,
discard the Lambda-mass
Lambda-rate appears in $p(s|z, M\dots)$
Lambda-cosmo appears in the m, M, z
conversion

