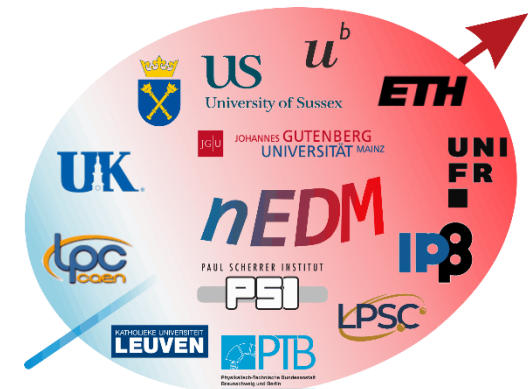


Ultra Cold Neutrons and mercury spins in magnetic fields. n2EDM experiment

Morgan Ferry, Katia Michielsen

LPSC (Grenoble)

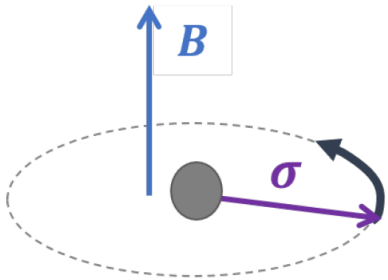
13th November 2025



I - Context

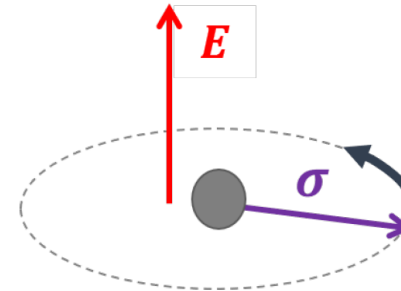
What is an EDM ?

Coupling of the spin to the magnetic field



Observable = magnetic moment μ

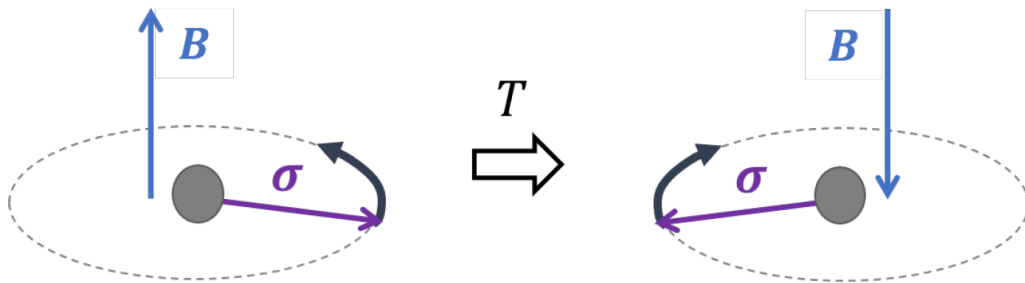
Coupling of the spin to the electric field



Observable = electric dipole moment (EDM) d

An **EDM** violates Charge Parity symmetry.

Coupling of the spin to the **magnetic field**



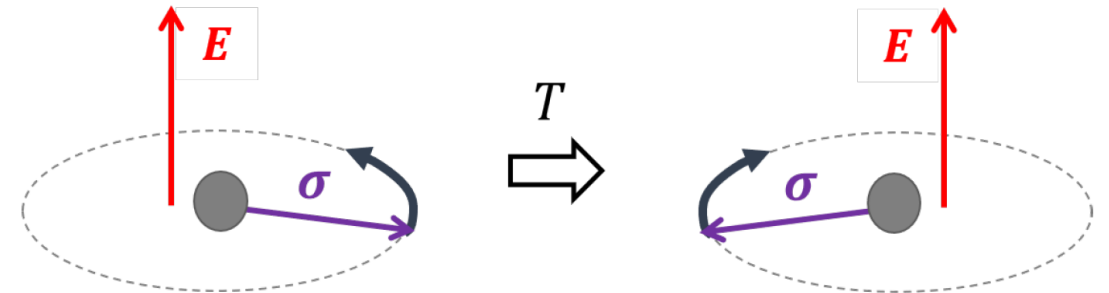
Observable = magnetic moment μ

$$\mu > 0$$

$$\mu > 0$$

T-invariant

Coupling of the spin to the **electric field**



Observable = electric dipole moment (**EDM**) d

$$d > 0$$

$$d < 0$$

not T-invariant

T violation \leftrightarrow CP violation
CP violation is one of the Sakharov criteria for matter anti-matter asymmetry during baryogenesis.

What do we know about the **nEDM**?

Theory:

$$d_n = \underbrace{10^{-32} e \text{ cm}}_{\text{weak interaction}} + \underbrace{\bar{\theta} \cdot 10^{-16} e \text{ cm}}_{\text{strong interaction}} + \underbrace{\hspace{1cm}}_{\text{new physics?}}$$

Current limit [1] : $d_n = (0.0 \pm 1.1_{\text{stat}} \pm 0.2_{\text{syst}}) 10^{-26} e \text{ cm}$

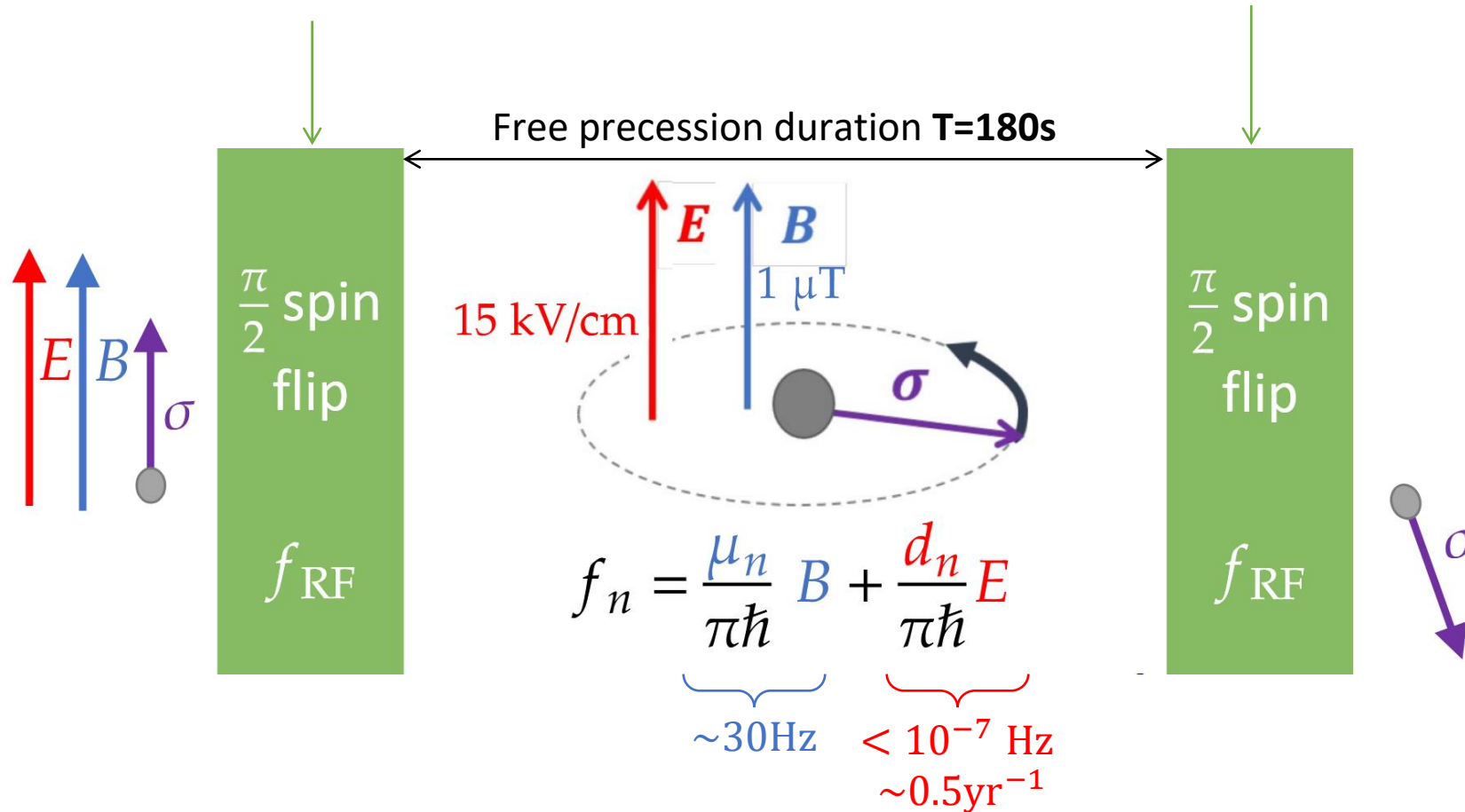
n2EDM sensitivity goal : $10^{-27} e \text{ cm}$

- CP violation probe in $\left\{ \begin{array}{l} \text{strong sector : } \bar{\theta} < 10^{-10} \text{ "strong CP problem"} \rightarrow \text{introduction of } \mathbf{axion} \\ \text{new physics with mass scale of } > 1 \text{ PeV (EFT)} \end{array} \right.$
- Generic probe
—> Complementary EDM measurements required : proton, electron, muon EDM

II – n2EDM overview

Larmor precession frequency f_n measurement

Ramsey method of oscillatory fields : 2 spin flips



Spin Asymmetry :

$$A(f_{\text{RF}}) = \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}} = \alpha \cos\left(\frac{f_{\text{RF}} - f_n}{\Delta\nu}\right)$$

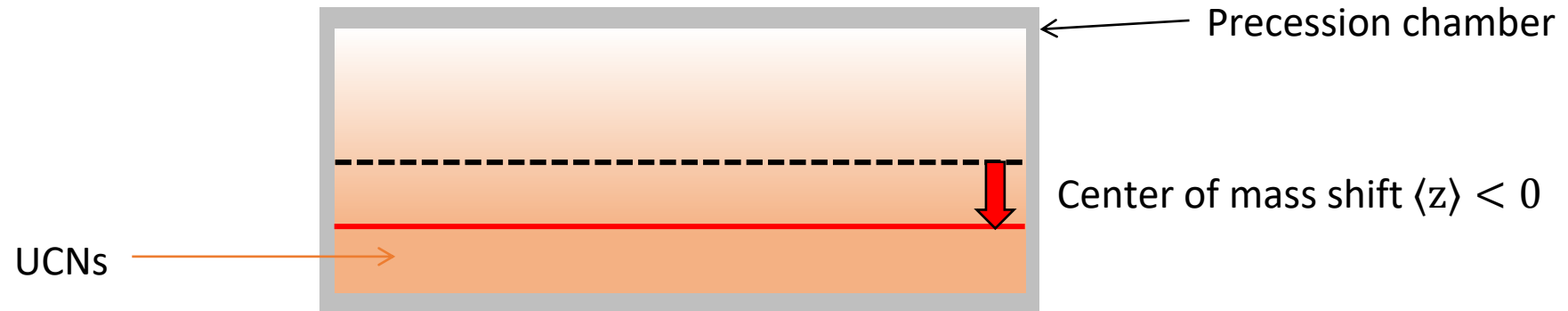
Number of neutrons spin aligned/anti-aligned with B

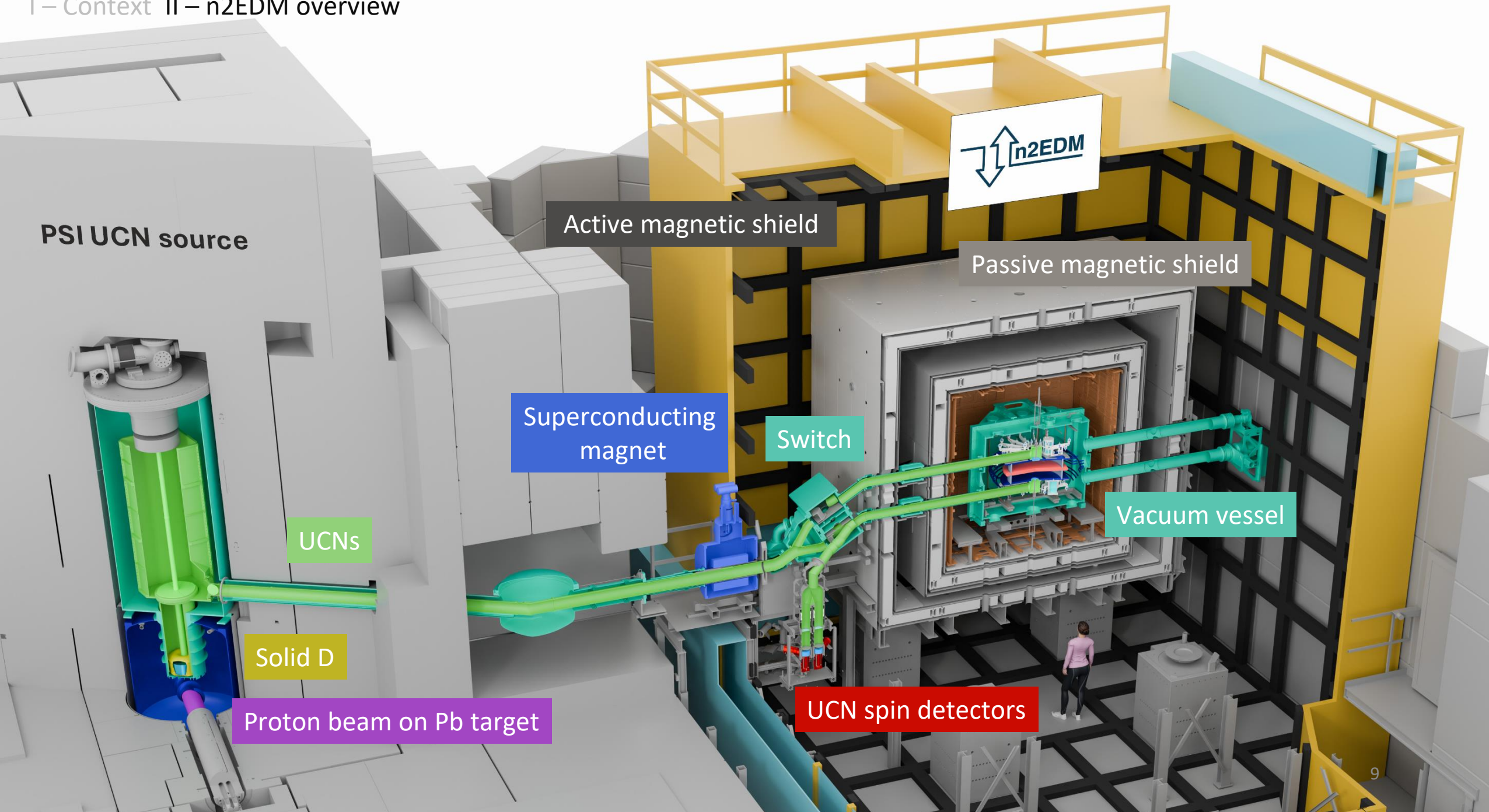
$$\sigma(f_n) = \frac{1}{2\pi\alpha T\sqrt{N}} \sim 10^{-6} \text{ Hz}$$

Measure f_n in a **known magnetic field** and **strong electric field**.

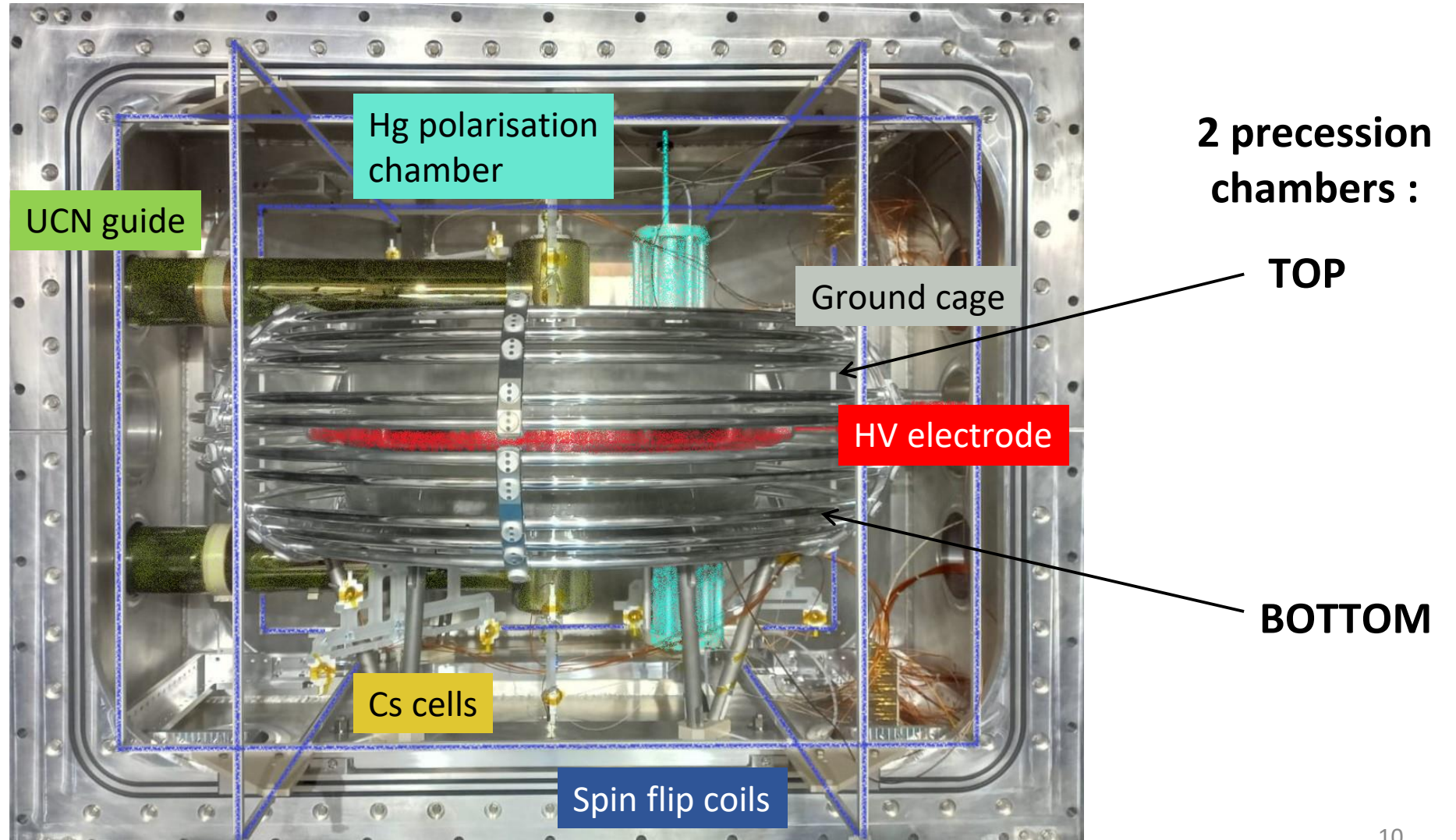
The use of Ultra Cold Neutrons (UCN)

- Challenge : **trap** the neutrons while measuring their precession. No charge...
 - **UCN** = Kinetic energy < 100 neV \rightarrow **bounce** off the walls
 - Gravitational potential energy ~ 100 neV !!
- \rightarrow **Sag in the gravitational potential** (unlike gases, their center of mass is below the geometric center)





Experiment overview



III- The Hg co-magnetometer

Concept, setup and performances.

Why do we need a co-magnetometer?

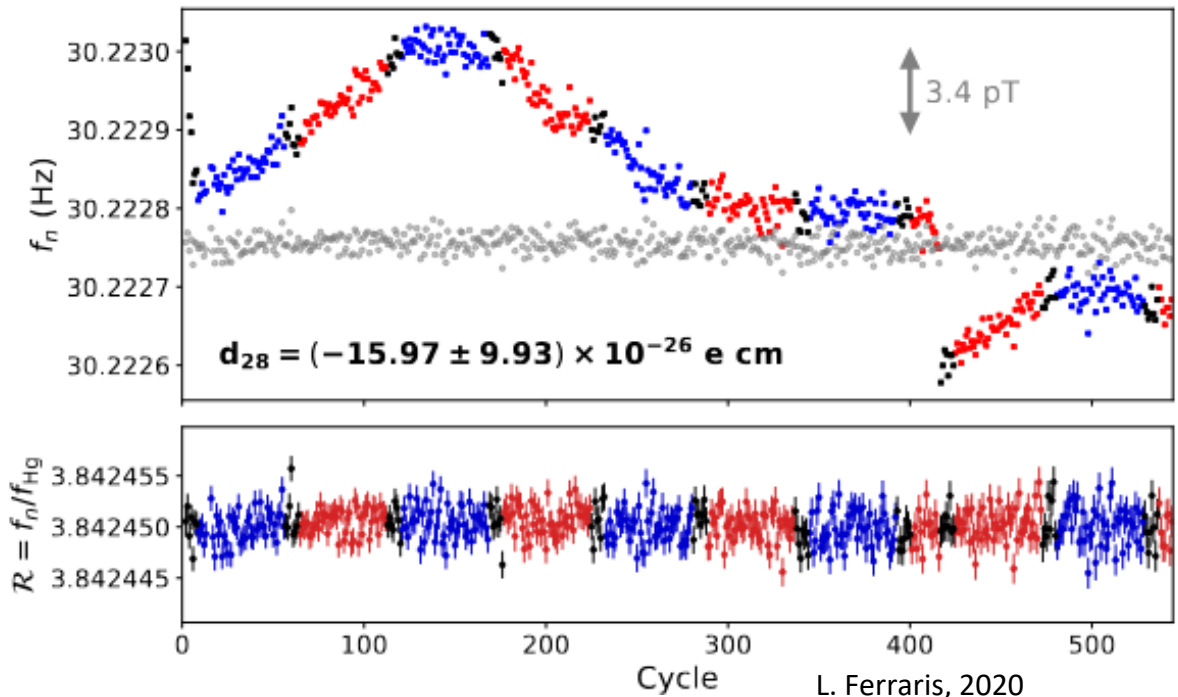
- Magnetic field drifts dominate f_n fluctuations
- Atom with nuclear spin and a negligible EDM [2]

$$f_n = \underbrace{\frac{\mu_n}{\pi\hbar} B}_{\sim 30\text{Hz}} + \underbrace{\frac{d_n}{\pi\hbar} E}_{< 10^{-7}\text{ Hz}}$$

-> Mercury vapour (same volume,

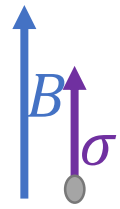
same time as n): $f_{\text{Hg}} = \frac{\mu_{\text{Hg}}}{\pi\hbar} B$

- We define $\mathcal{R} = \frac{f_n}{f_{\text{Hg}}}$



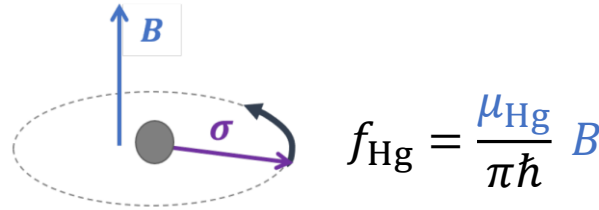
Hg magnetometry concept in n2EDM

Polarized Hg



$\frac{\pi}{2}$
spin
flip

Hg precession



$T = 180 \text{ s}$

time

Optical pumping

Precession

relevant ^{199}Hg quantum states

Excited states

-1/2 1/2

$h\nu$
spin 1

Ground states

-1/2 1/2



relevant ^{199}Hg quantum states

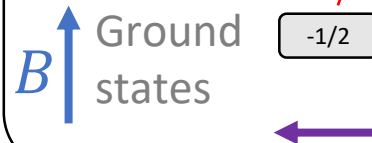
Excited states

-1/2 1/2

$h\nu$
spin 1

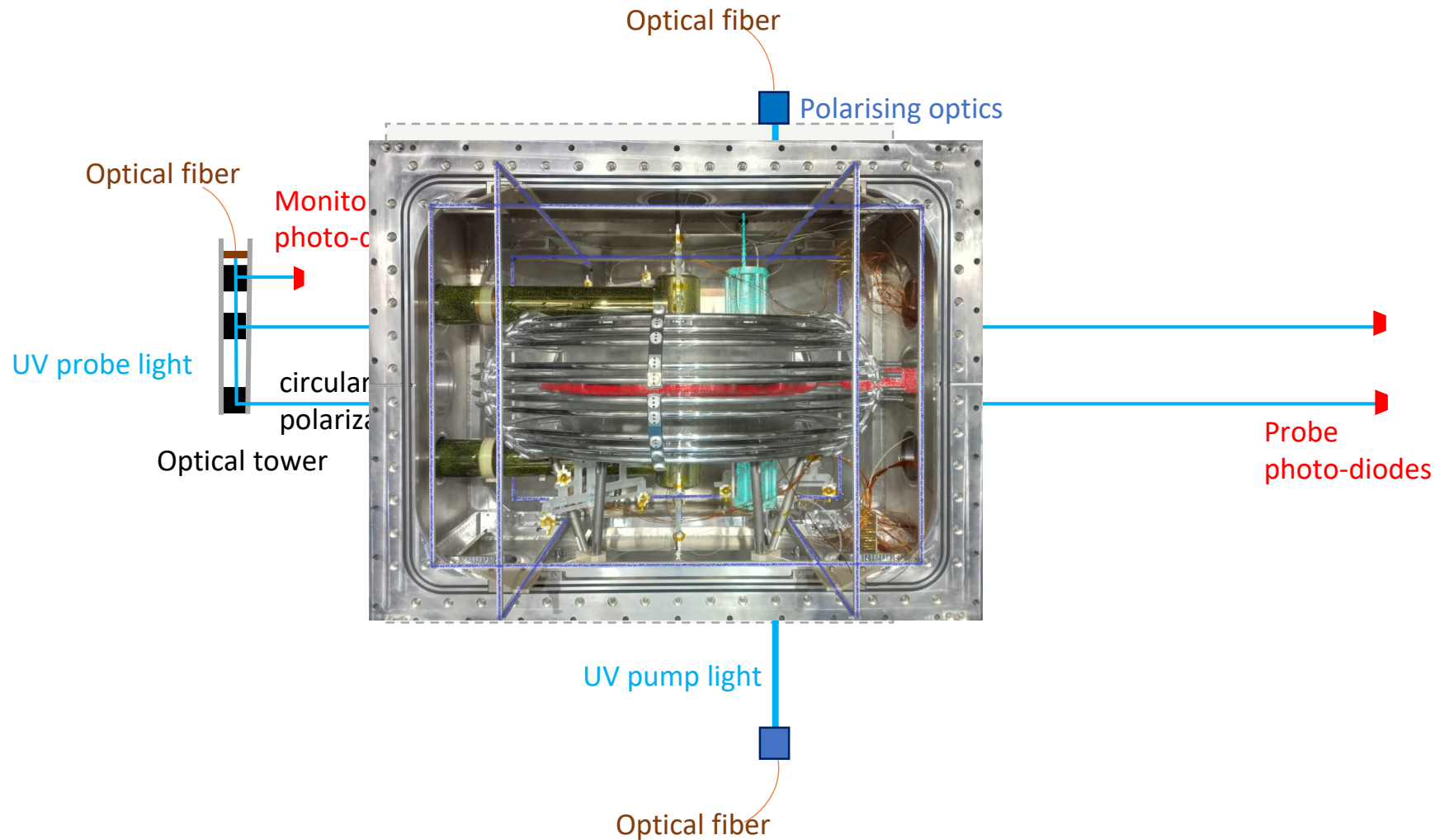
Ground states

-1/2 1/2

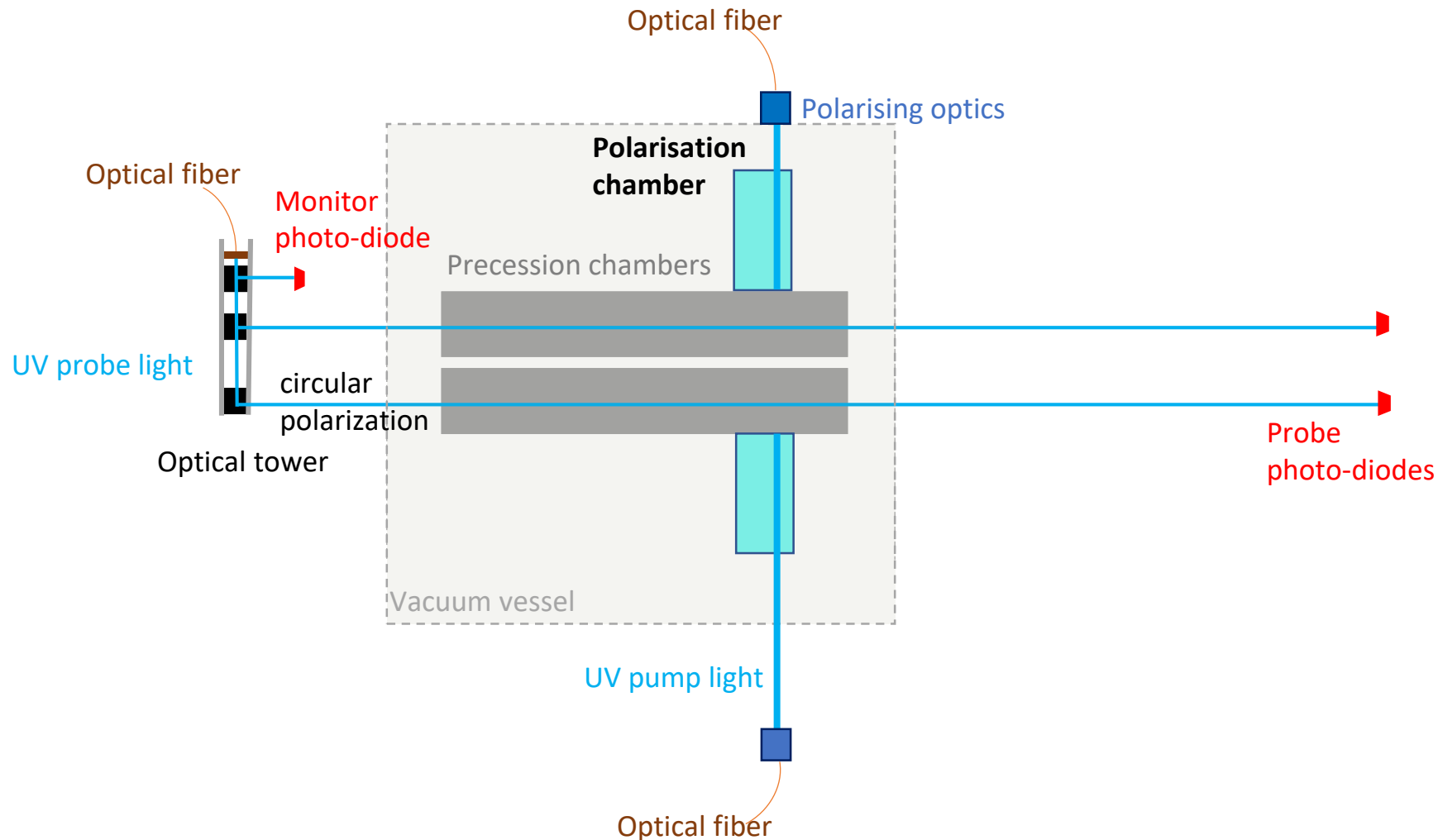


Absorption/transmission
modulated at the precession
frequency.

Hg magnetometry implementation in n2EDM

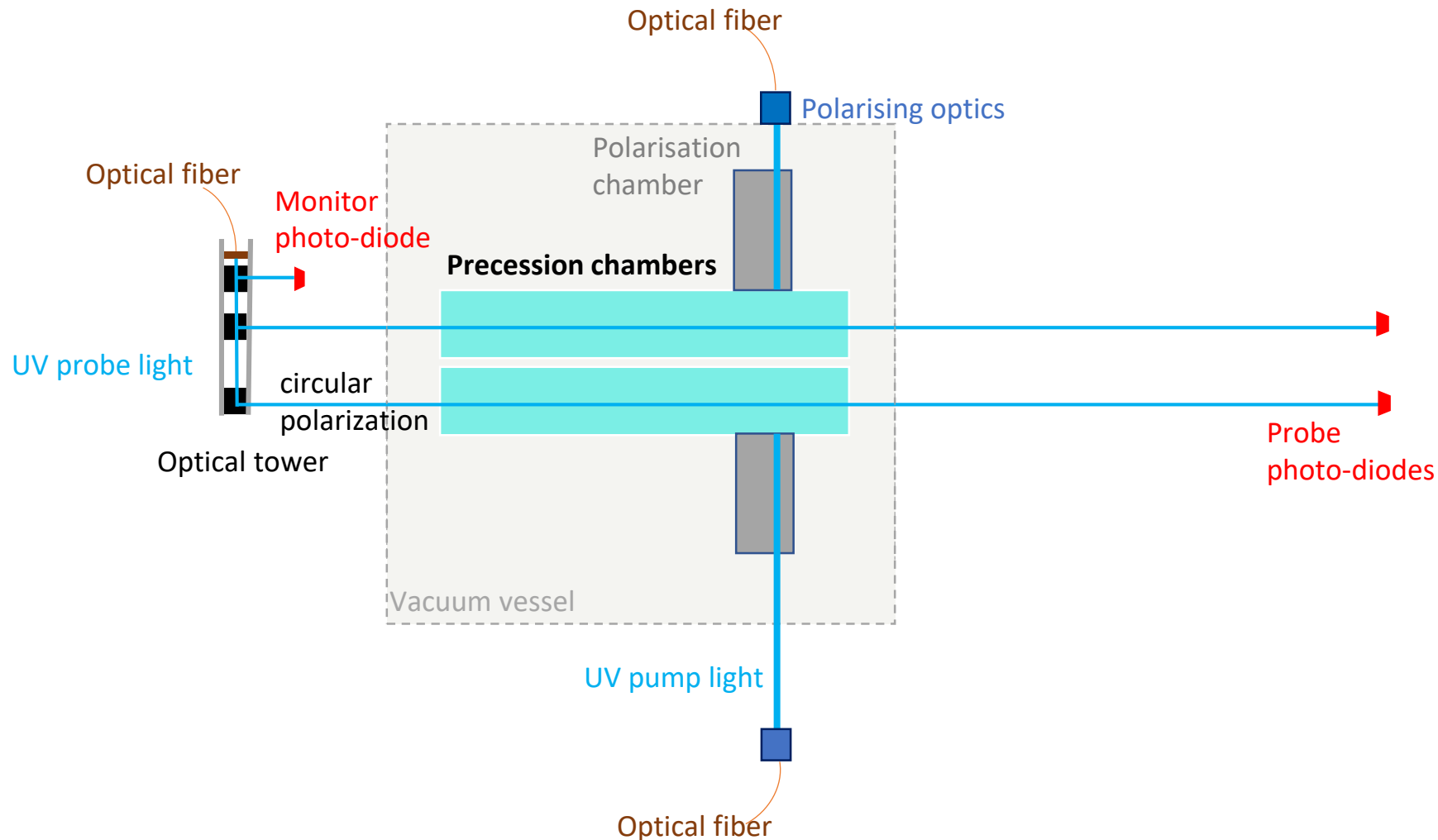


Hg magnetometry implementation in n2EDM



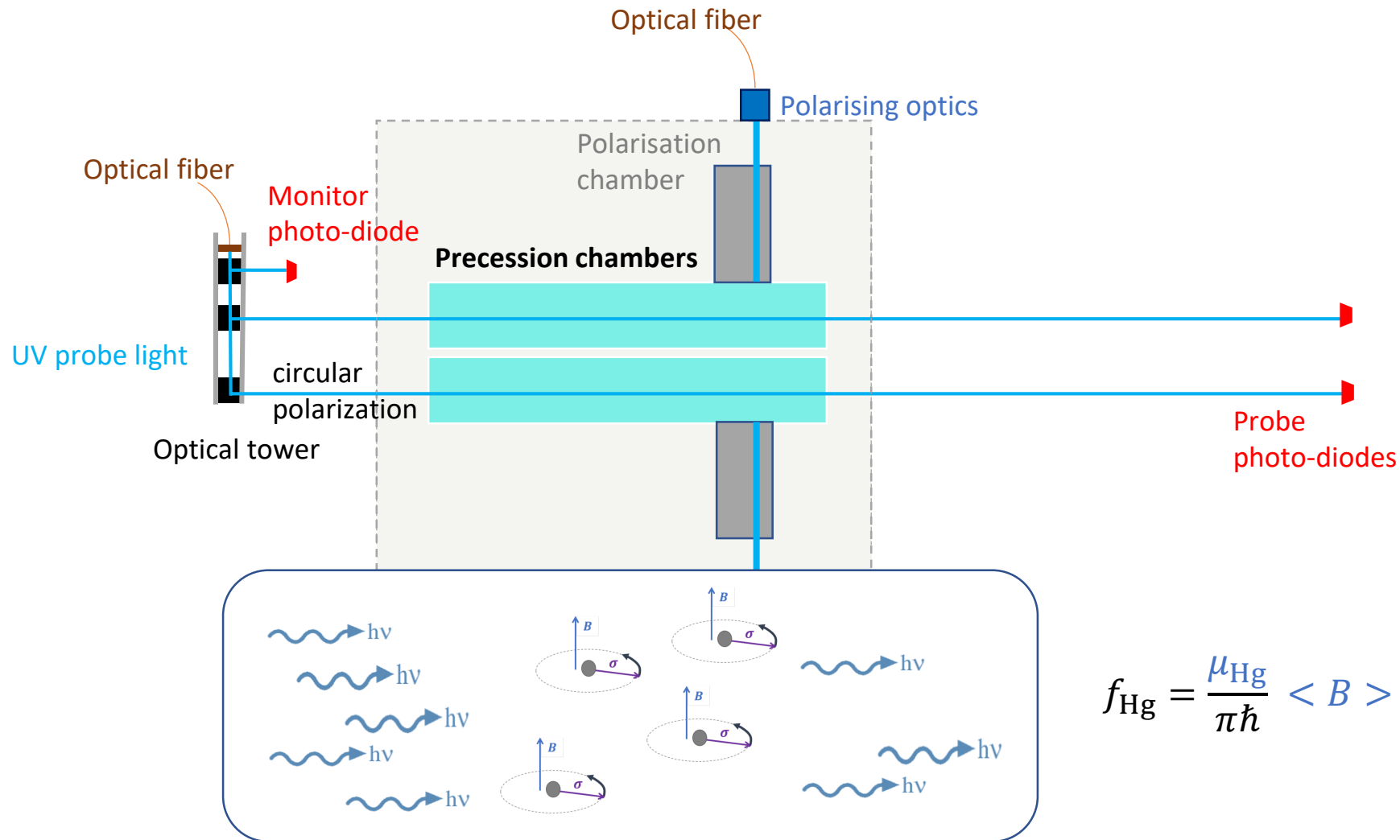
- Hg vapour is polarised prior to injection

Hg magnetometry implementation in n2EDM



- Hg vapour is polarised prior to injection
- **Hg and neutron precession occur simultaneously**

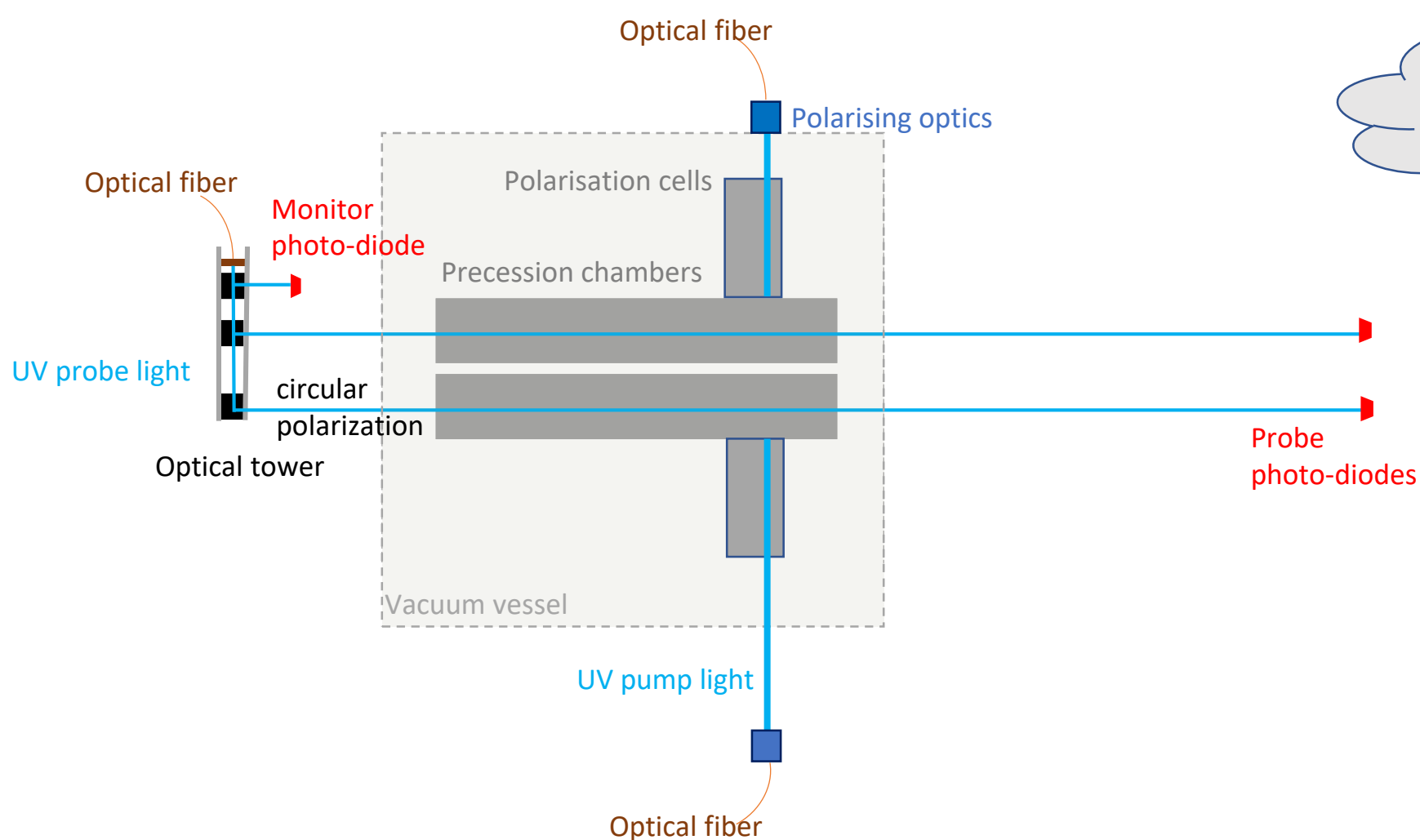
Hg magnetometry implementation in n2EDM



- Hg vapour is polarised prior to injection
- Hg and neutron precession occur simultaneously
- **Modulated transmitted light is continuously recorded.**

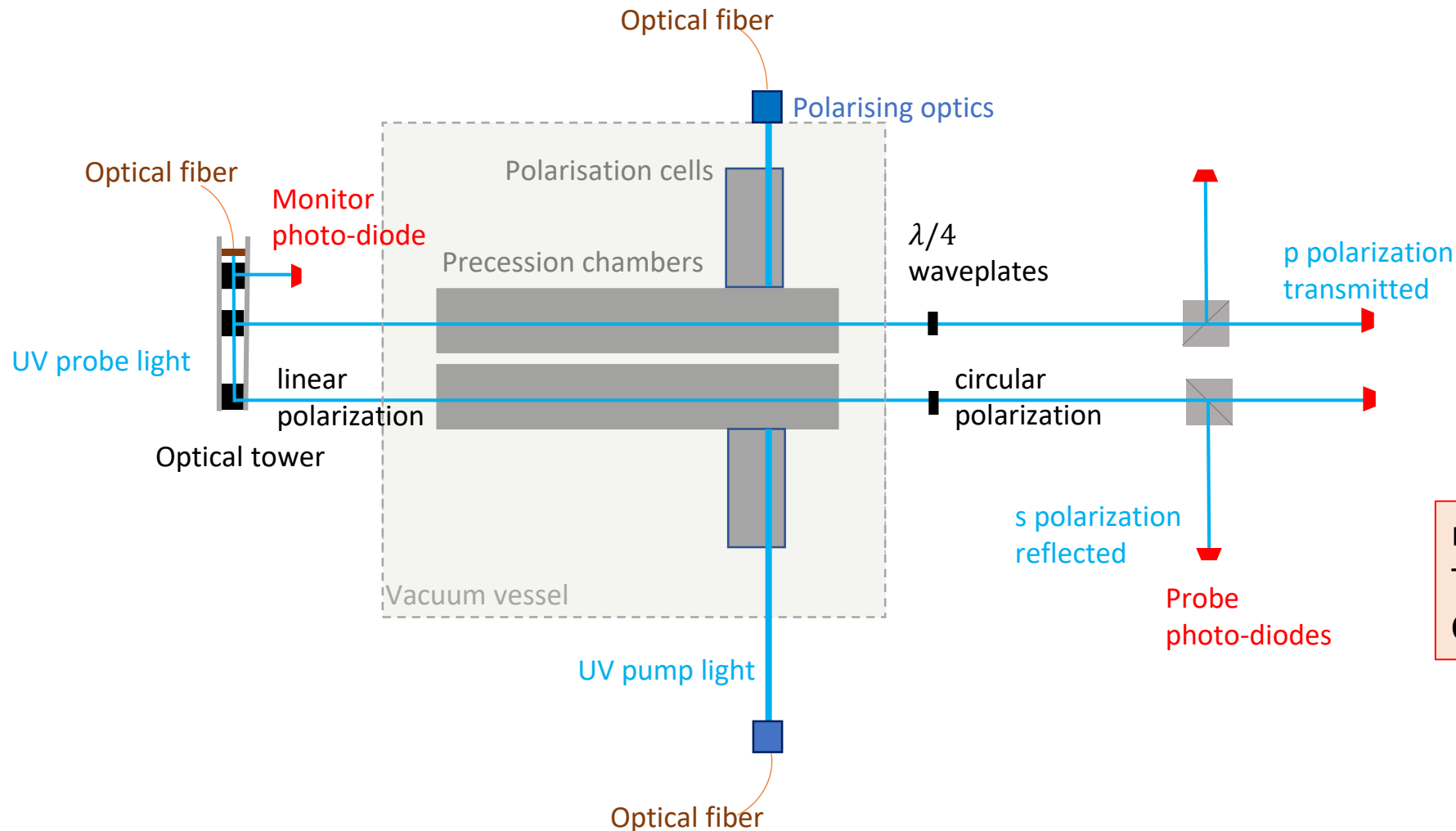
$$f_{\text{Hg}} = \frac{\mu_{\text{Hg}}}{\pi \hbar} \langle B \rangle$$

Optical reading of Hg precession



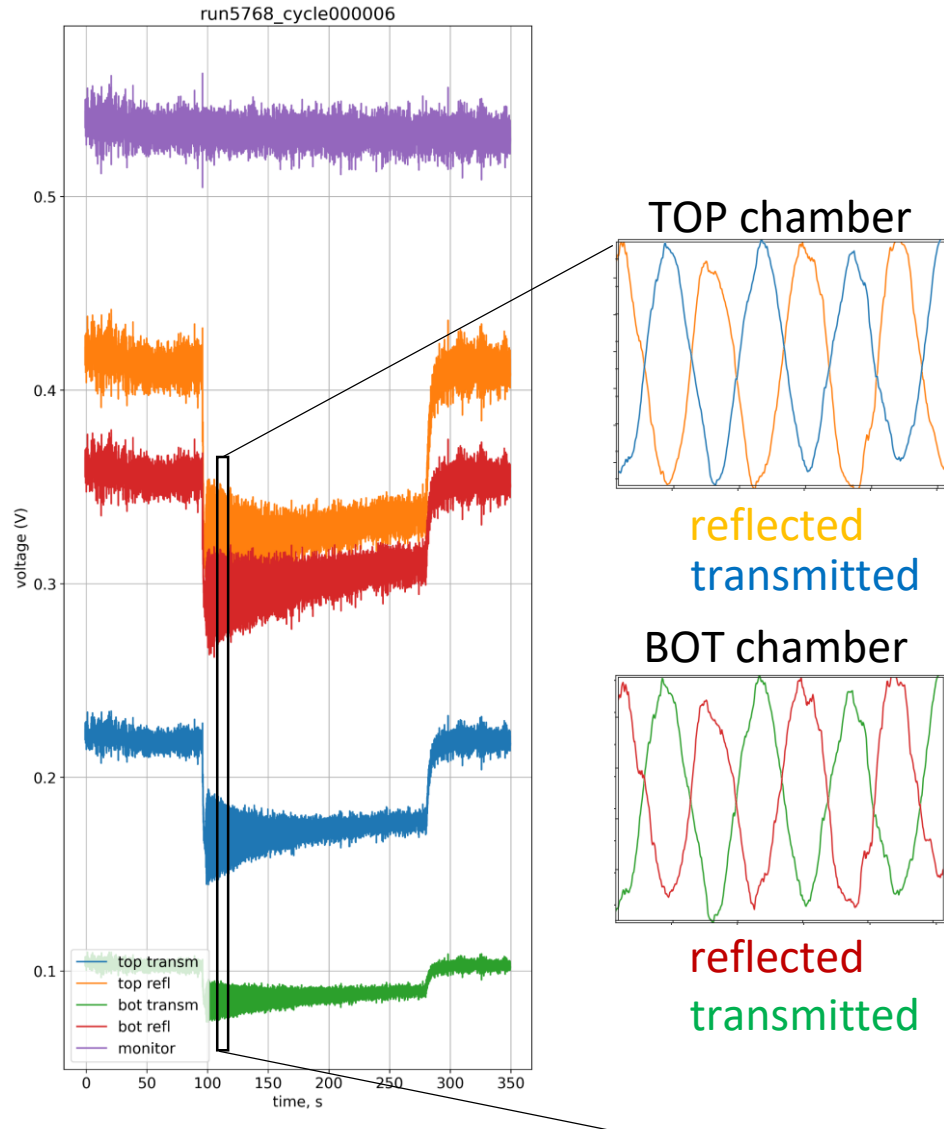
nEDM: Circularly polarised probe.
One signal per chamber.

Optical reading of Hg precession

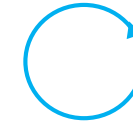


n2EDM: linearly polarised probe.
Two signals for each chamber.
Cancellation of common noise.

Optical reading of Hg precession



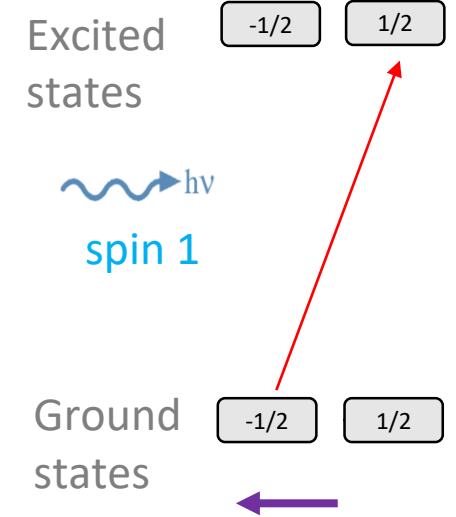
UV light



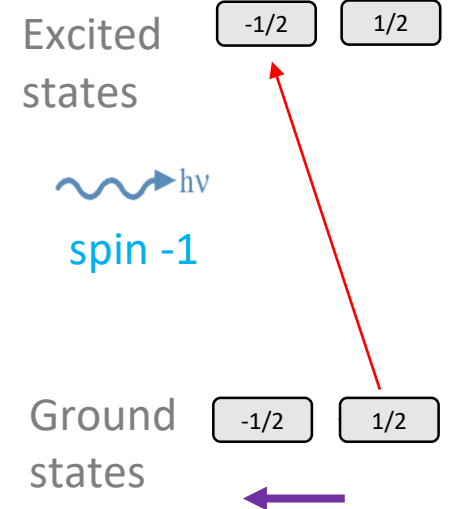
+



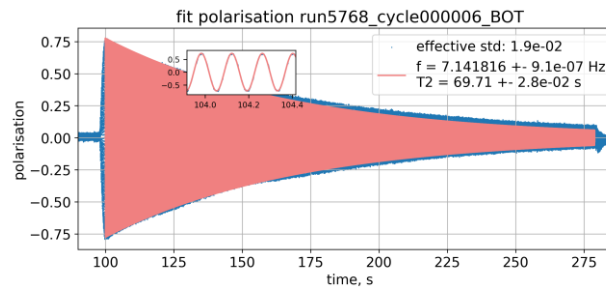
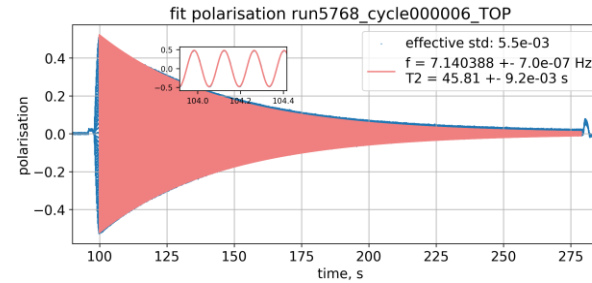
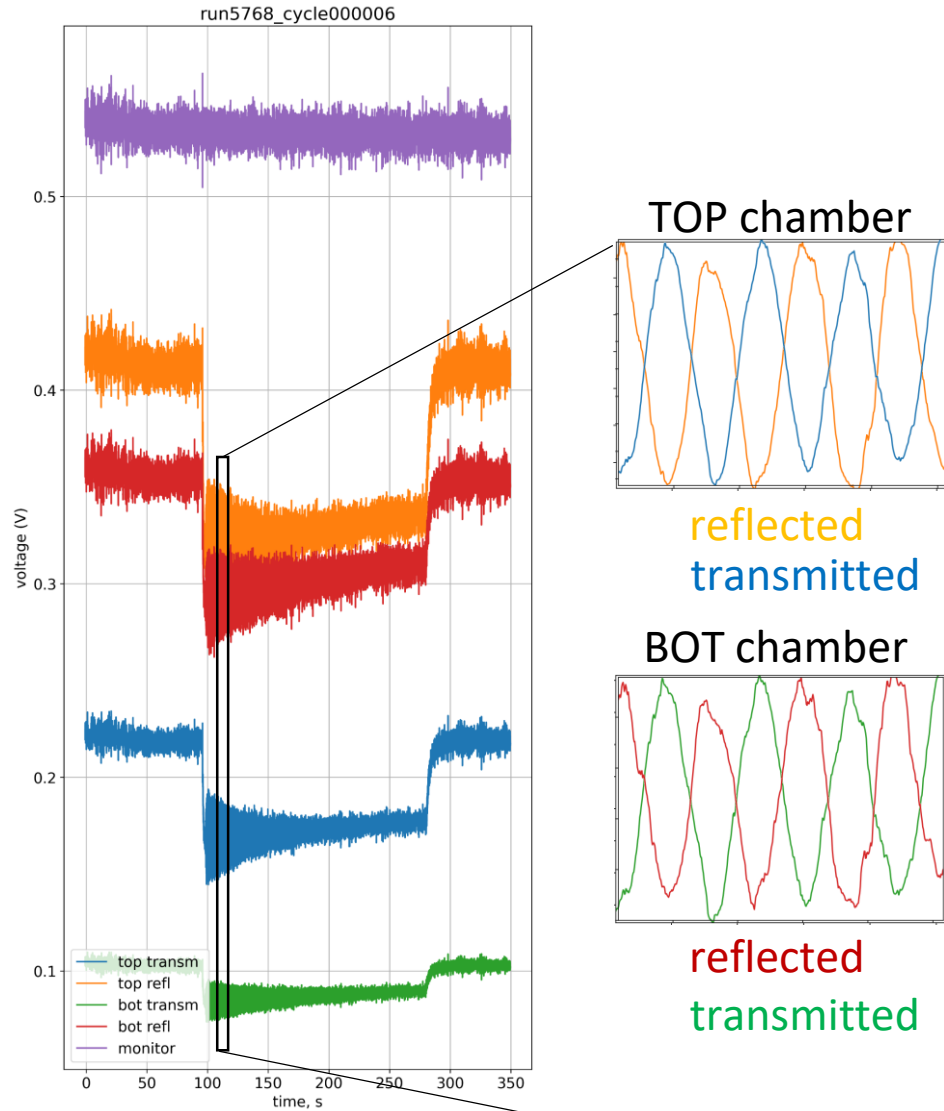
^{199}Hg spin during precession



^{199}Hg spin during precession

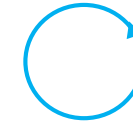


Optical reading of Hg precession

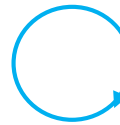


Linearly polarised probe lowered the uncertainty by a factor 3.

UV light



+



^{199}Hg spin during precession

Excited states

-1/2 1/2

$\hbar\nu$
spin 1

Ground states

-1/2 1/2

^{199}Hg spin during precession

Excited states

-1/2 1/2

$\hbar\nu$
spin -1

Ground states

-1/2 1/2

The great combination with linearly polarised light

Transmitted power

$$\Pi_t = a \Pi_0 e^{-\sigma n L (1+p)}$$

Reflected power

$$\Pi_r = (1-a) \Pi_0 e^{-\sigma n L (1-p)}$$

σ unpolarised cross section

n mercury density

L length of the precession volume

p polarisation of the mercury vapour

Transmission/reflection factors:

a and $1-a$

$$\left. \begin{array}{l} \Pi_t = a \Pi_0 e^{-\sigma n L (1+p)} \\ \Pi_r = (1-a) \Pi_0 e^{-\sigma n L (1-p)} \end{array} \right\} \begin{array}{l} \Pi_r \Pi_t = a(1-a) e^{-2\sigma n L} \quad \text{Derive } \sigma n_0 e^{-t/T_{leak}} L \\ \frac{\Pi_r}{\Pi_t} = e^{-2\sigma n L p} \quad \text{Derive } p \text{ knowing } \sigma n L \end{array}$$

$$-\frac{1}{2\sigma n L} \ln \left(\frac{\Pi_r}{\Pi_t} \right) = p_0 \cos(2\pi f_{\text{Hg}} t + \phi) e^{-\frac{t}{T_2}}$$

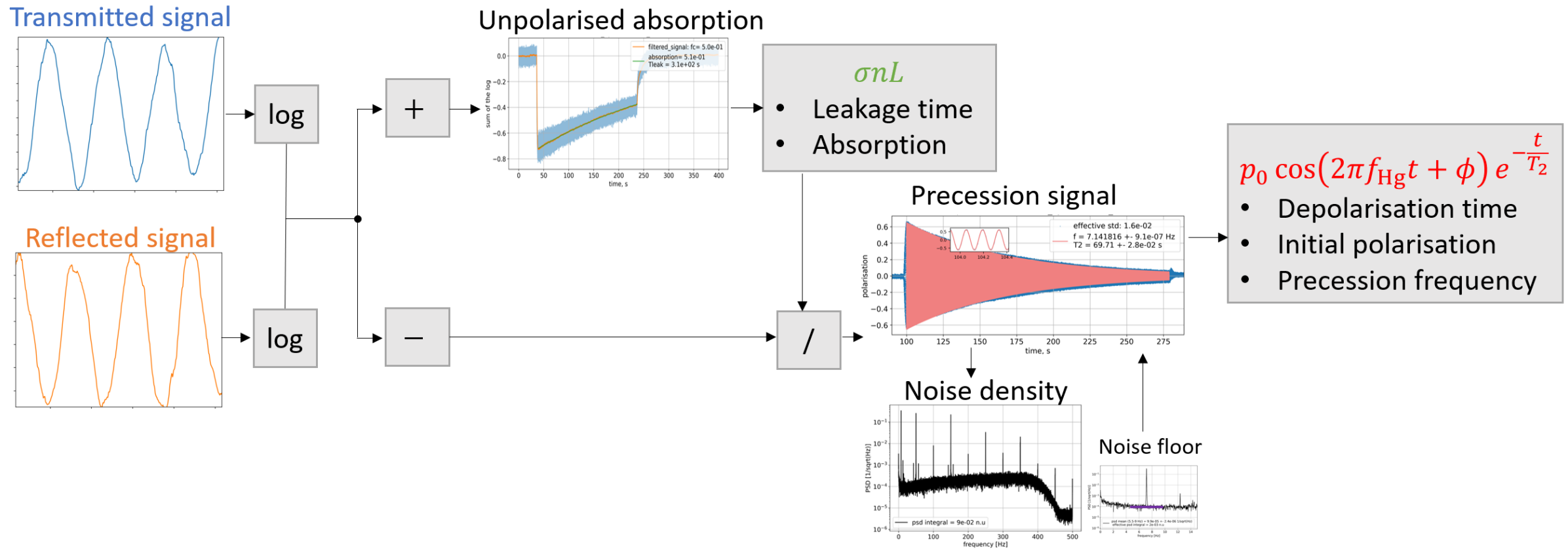
1st fit provides

- Absorption
- Leakage time of precession chamber

2nd fit provides

- Polarisation of mercury vapour p_0
- Depolarisation time T_2
- Noise spectrum
- Precession frequency f_{Hg}

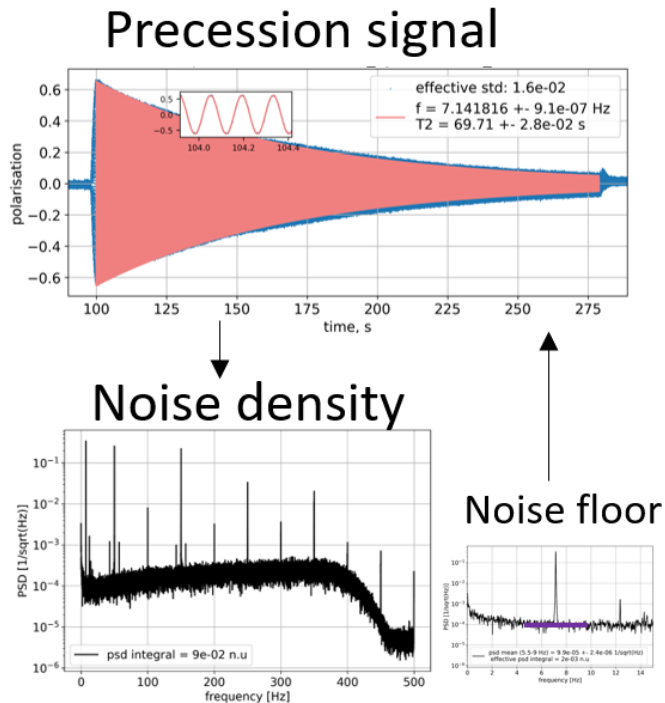
Hg precession analysis: summary



Uncertainty derivation of precession frequency

Philosophy: filtering without filtering

→ Only the noise at the frequency of the signal affects our estimator



Equivalent white noise as input of the χ^2 minimisation fit.


→ diagonal elements of the covariant matrix give the uncertainties.

Verified with simulated data. ✓

Performance of the Hg magnetometer

- Precision:

TOP chamber	$\frac{\sigma(f_{\text{Hg}})}{f_{\text{Hg}}} = 7 \times 10^{-8}$	$\frac{\sigma(f_{\text{n}})}{f_{\text{n}}} = 2 \times 10^{-7}$
BOT chamber	$\frac{\sigma(f_{\text{Hg}})}{f_{\text{Hg}}} = 4 \times 10^{-8}$	$\frac{\sigma(f_{\text{n}})}{f_{\text{n}}} = 2 \times 10^{-7}$
- Accuracy: unbiased estimator?
 - effect of non-white noise,
 - magnetic field drifts within one cycle

status

ongoing

Magnetic field measured with a precision of 70 fT !

Also vertical gradient G_{TB} is measured.

→ Study systematic effects affecting $\mathcal{R} = \frac{f_{\text{n}}}{f_{\text{Hg}}}$

IV – Physics cases

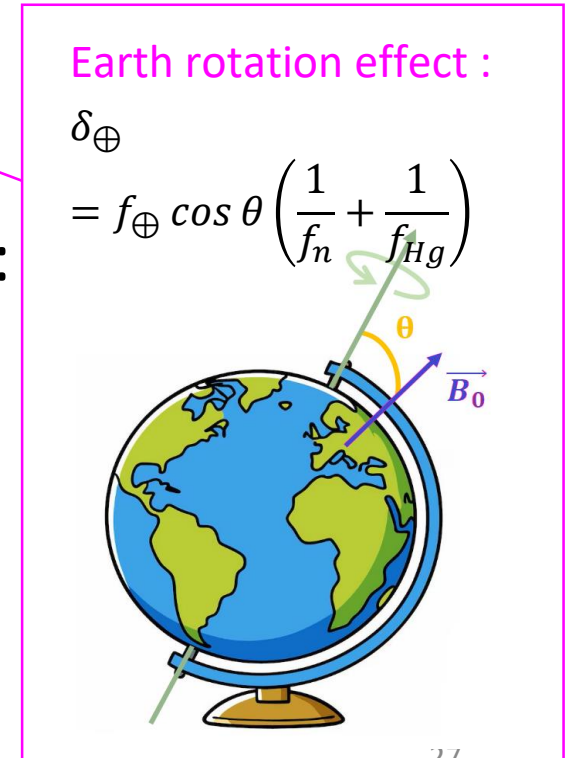
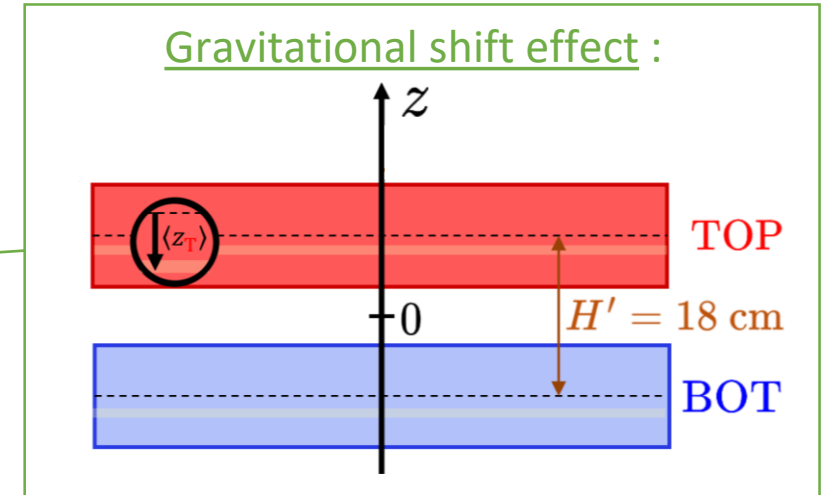
Expression of $\mathcal{R} = f_n/f_{Hg}$

$$\mathcal{R}_{T,B}^{\uparrow,\downarrow} = \frac{\gamma_n}{\gamma_{Hg}} \left[1 + \underbrace{\frac{\langle z_{T,B} \rangle}{B} G_{TB}}_{\sim 10^{-7}} \mp \underbrace{\delta_{\oplus}}_{\sim 10^{-6}} \right]$$

- At first order, $\mathcal{R} \sim \frac{\gamma_n}{\gamma_{Hg}}$ $\swarrow \searrow$ neutron & Hg gyromagnetic ratios

...but there are some minor corrections that depend on :

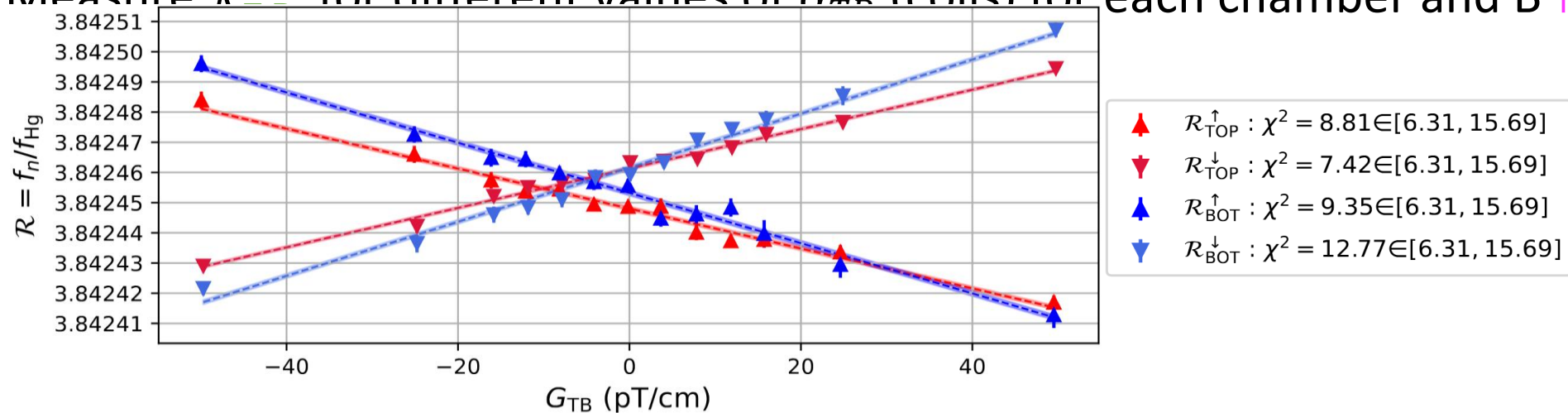
- G_{TB} : vertical magnetic gradient
- Chamber **TOP**, **BOT**
- B orientation \uparrow, \downarrow



\mathcal{R} – Curves for the probe of physics cases

- \mathcal{R} – Curves : $\mathcal{R}_{T,B}^{\uparrow,\downarrow} = \frac{\gamma_n}{\gamma_{Hg}} \frac{\langle z_{T,B} \rangle}{B} G_{TB} + \frac{\gamma_n}{\gamma_{Hg}} (1 \mp \delta_{\oplus})$
 $= a_{T,B} G_{TB} + b^{\uparrow,\downarrow} \rightarrow 4 \text{ affine functions } \mathcal{R}_{TOP}^{\uparrow}, \mathcal{R}_{TOP}^{\downarrow}, \mathcal{R}_{BOT}^{\uparrow}, \mathcal{R}_{BOT}^{\downarrow} \text{ VS } G_{TB}$

- Measure $\mathcal{R}_{\pm}^{\uparrow,\downarrow}$ for different values of G_{TB} (coils) for each chamber and $B^{\uparrow,\downarrow}$



γ_n / γ_{Hg}	$\langle z_{TOP} \rangle$ (cm)	$\langle z_{BOT} \rangle$ (cm)	T_{\oplus} (h)
3.8424557(3)	-0.161(4)	-0.214(6)	24(1)

Compatible with [1]

Compatible with [2]

[1] : *A measurement of the neutron to 199Hg magnetic moment ratio*, nEDM collaboration, 2014

[2] : *Ultracold neutron energy spectrum and storage properties from magnetically induced spin depolarization*, n2EDM collaboration, 2025

Conclusion

- First data-set not limited by Hg uncertainty
- Gravitational shift and Earth rotation measurements
- Newer of more precise data —> access to new systematic effects
- nEDM data taking at the end of 2025



Thank you for your attention!
Questions?

Additional slides

n2EDM status and outlook

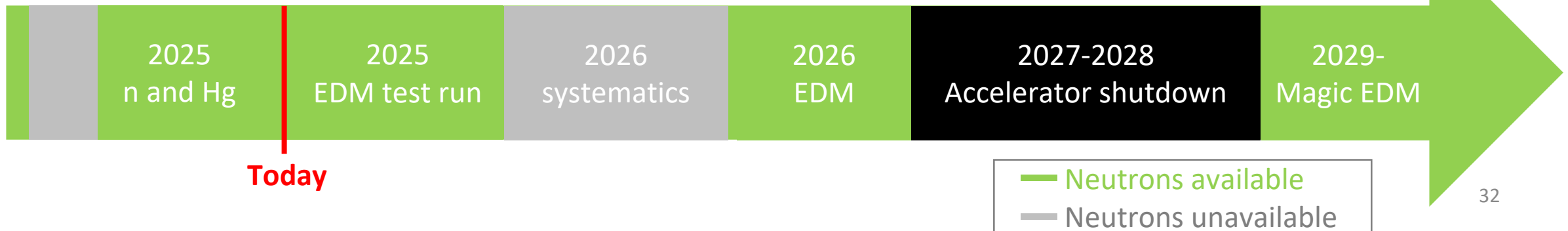
- UCN production, detection, polarization, spin transport
- Magnetic shielding and characterization
- Magnetic field generation and characterization
- Electric field
- Hg co-magnetometer injection, polarization, optical reading
- Cs magnetometers installation, polarisation & optical reading



Operational



Not yet operational



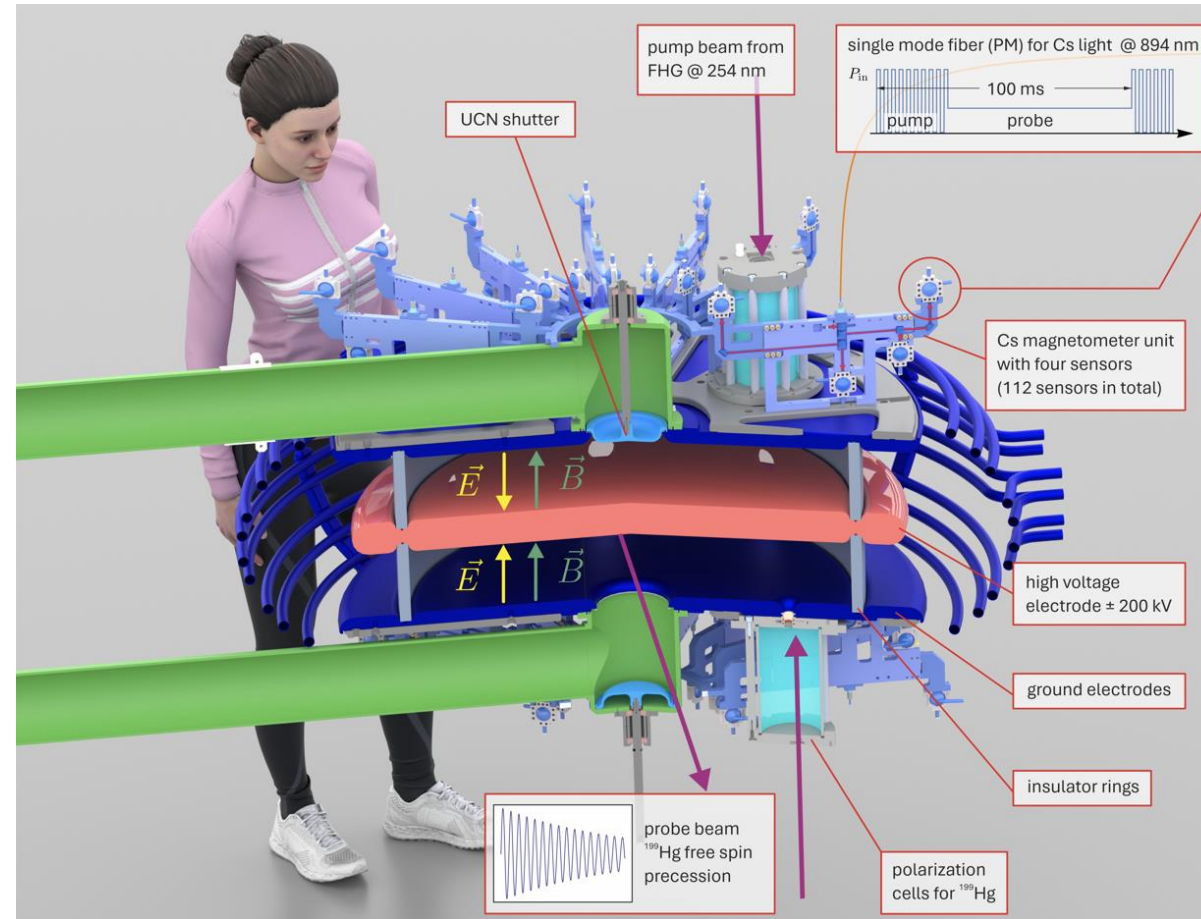
n2EDM Experiment overview

UCN source

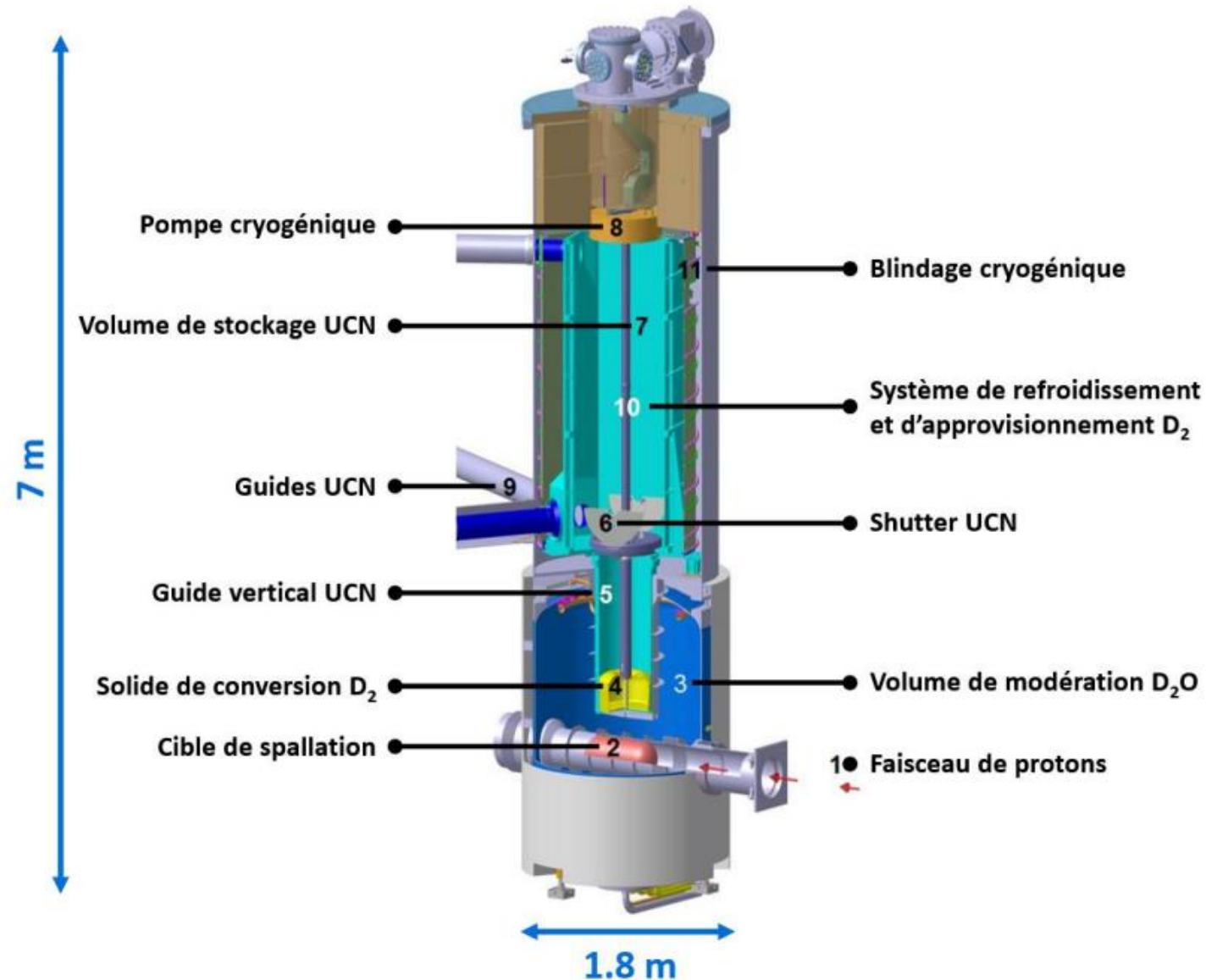
Superconducting
magnet

Vacuum vessel with
Precession chambers

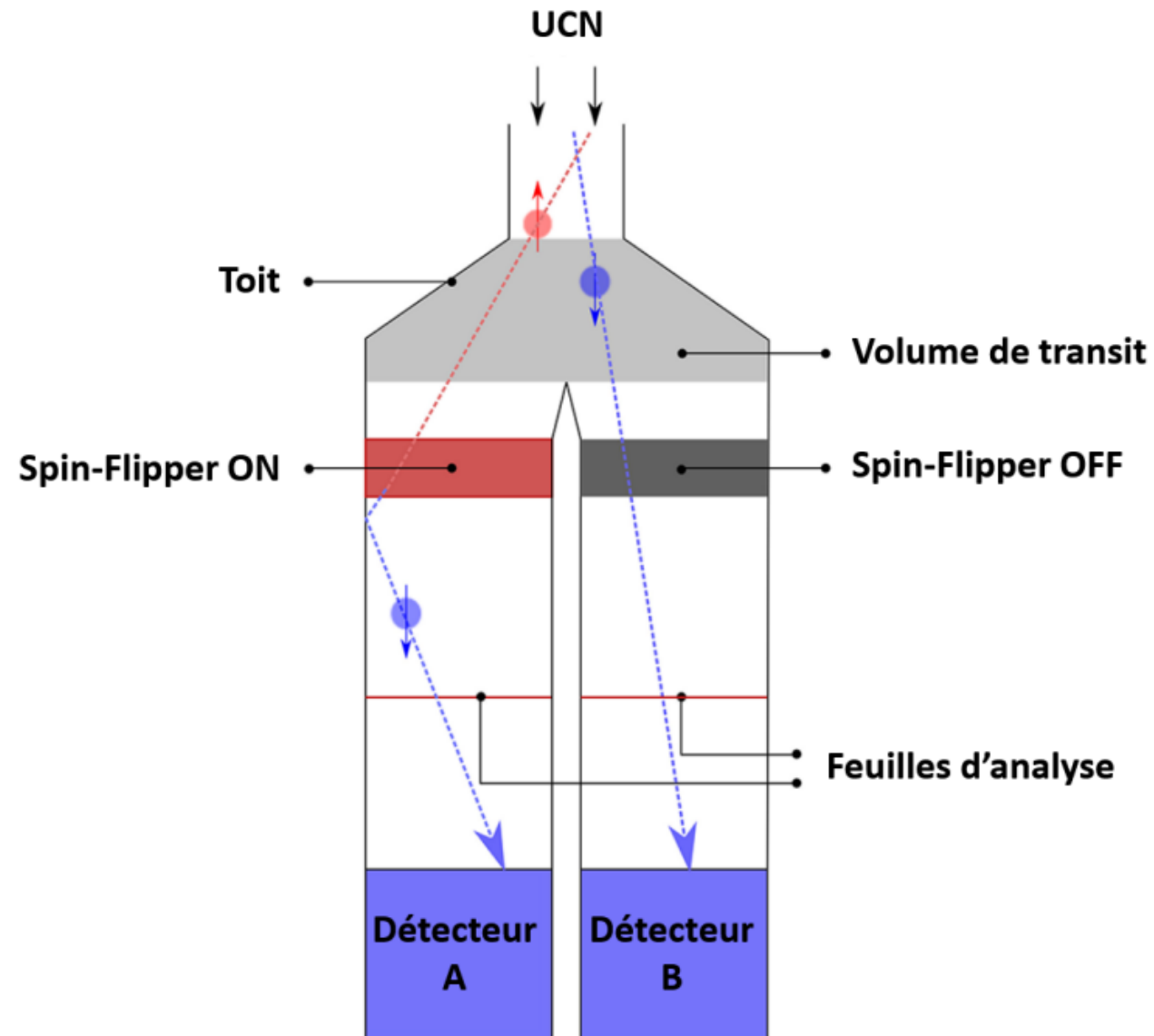
n2EDM Experiment overview



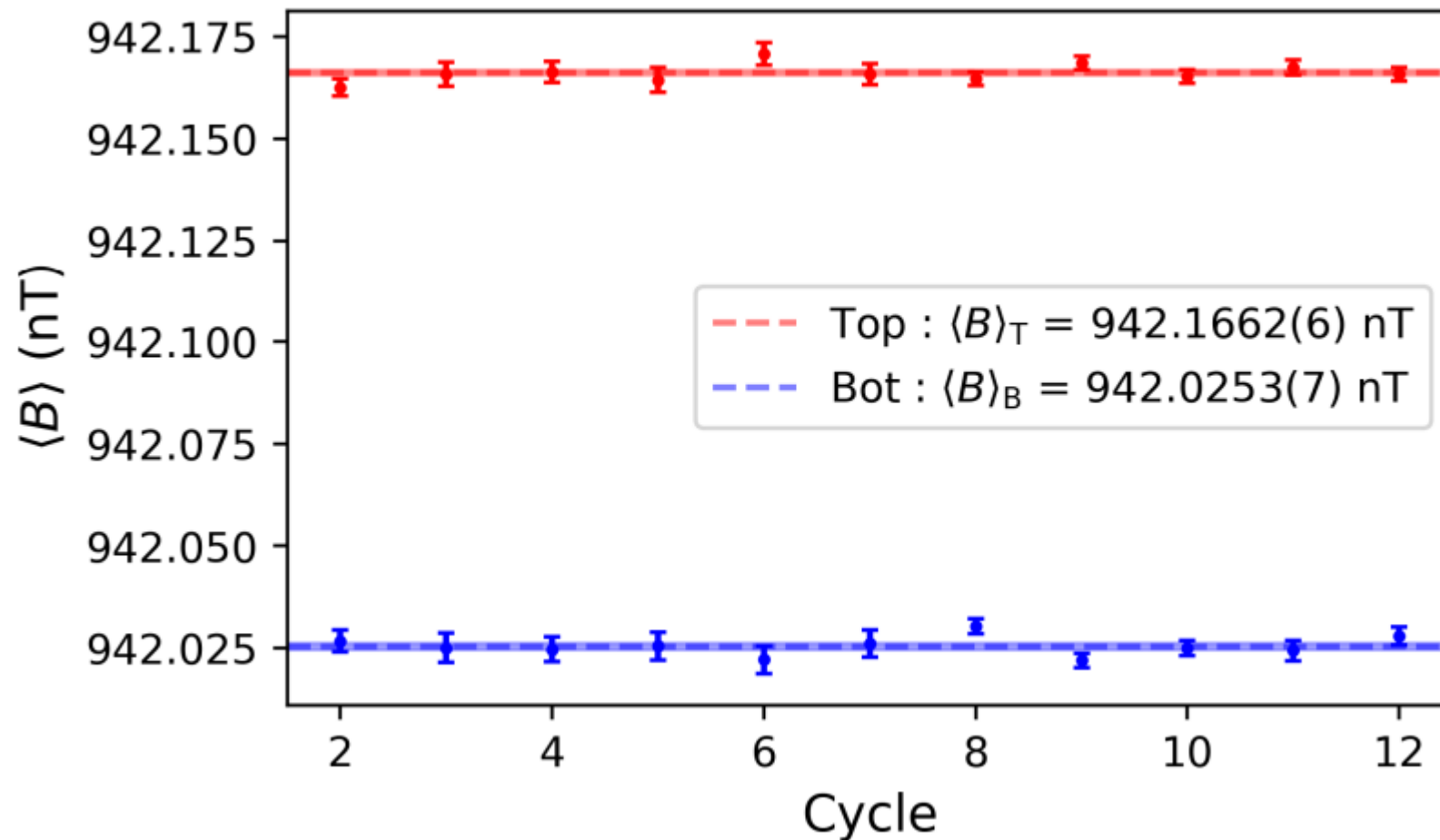
UCN production



UCN spin detectors



Magnetic field measurement in the presence of gradient



Spin asymmetry fit

$$A = \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}} = \alpha \cos\left[\frac{(f_{RF} - f_n)}{\Delta\nu}\right]$$

$$f_n \sim \frac{\gamma_n}{\gamma_{\text{Hg}}} f_{\text{Hg}}$$

