



Status of V_{ub} and V_{cb} measurements at LHCb



GDR Intensity Frontier Annual Workshop
13/11/2025

Alois Caillet (LPNHE)



The CKM Matrix

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

- Describes the quark **flavour changing transition couplings**

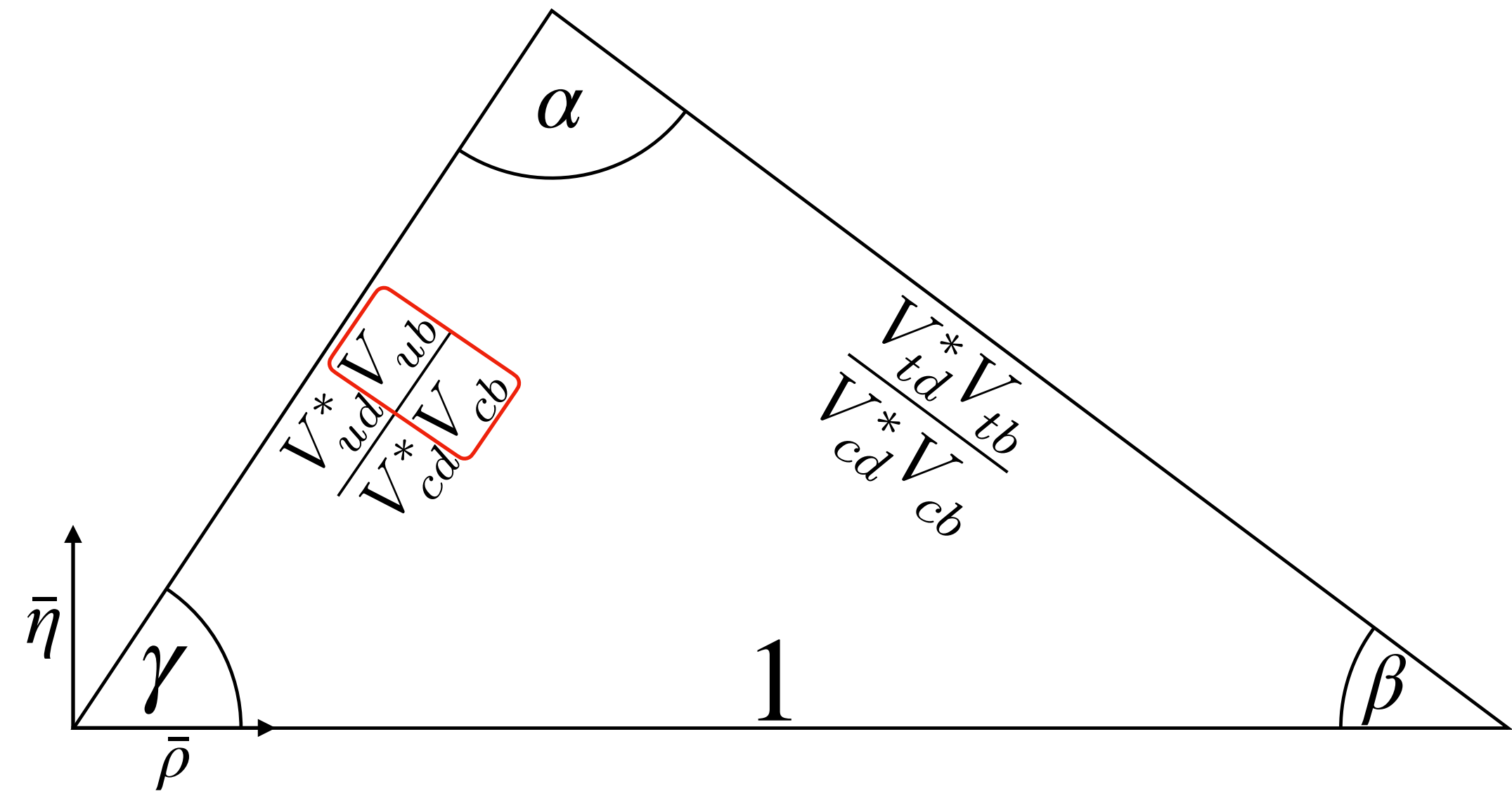
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- Describes the quark **flavour changing transition couplings**

$$\Rightarrow V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

- Its **unitarity** translates into the **unitarity triangle**



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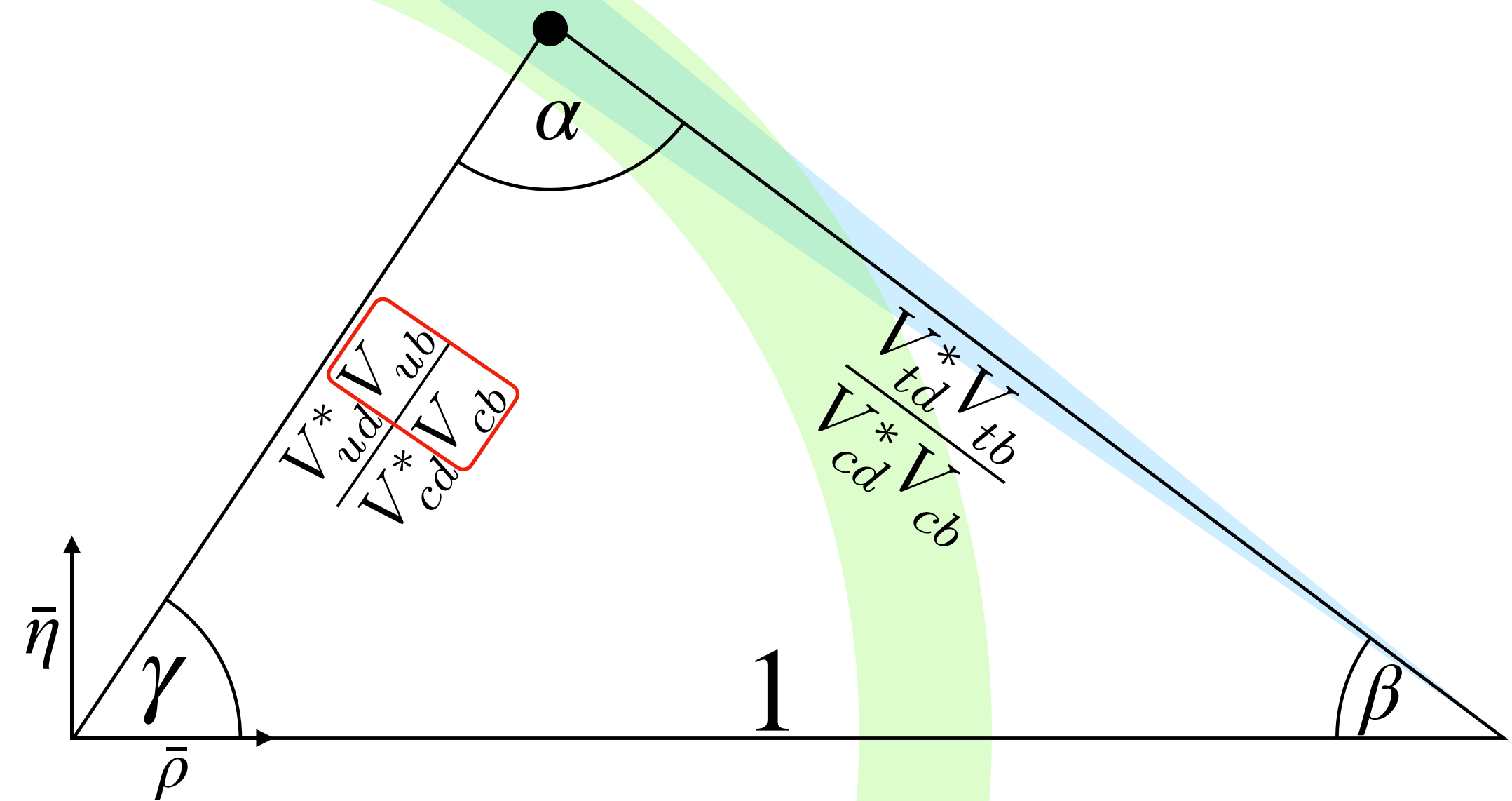
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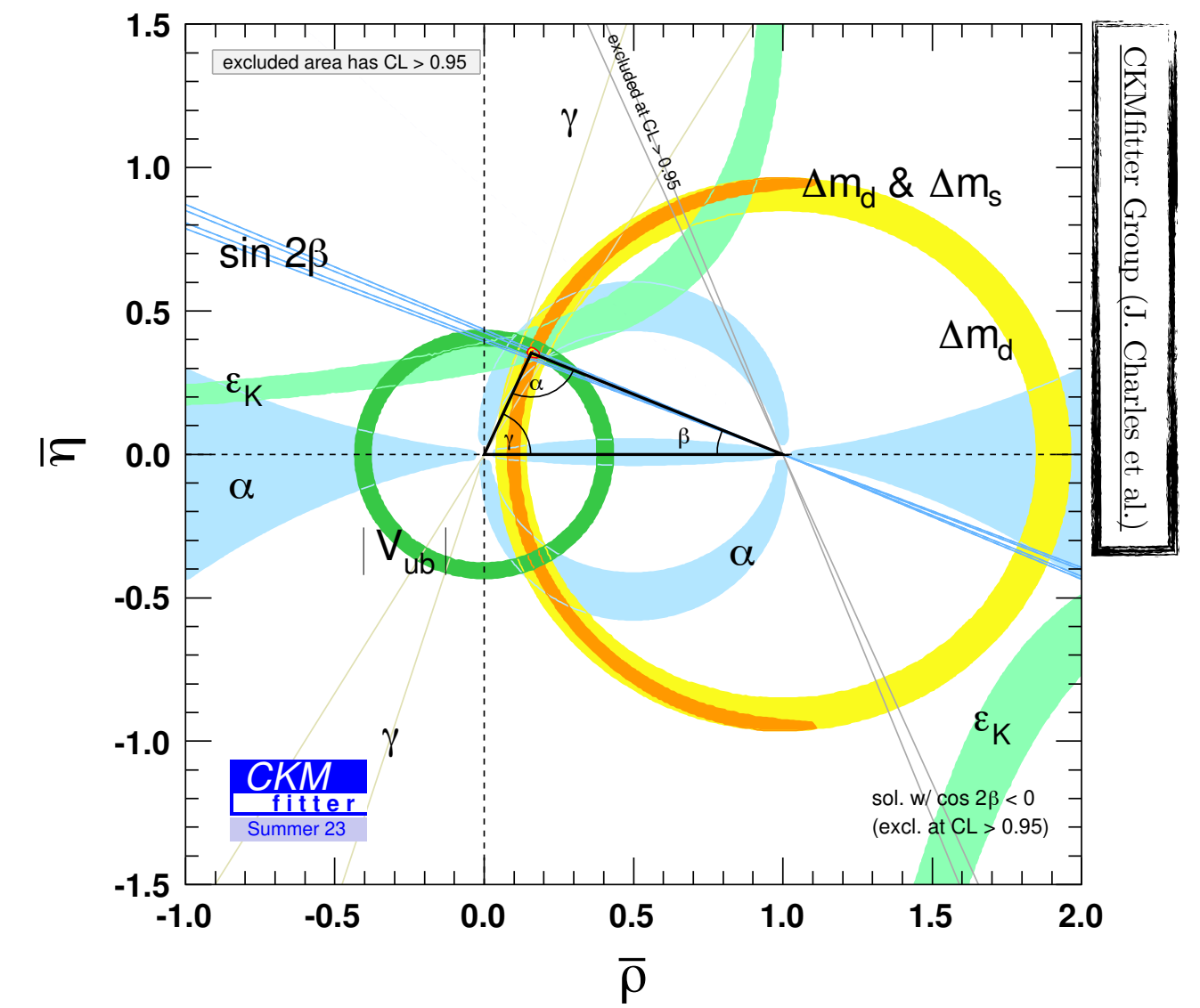
- V_{ub} and V_{cb} both connect β and its opposite side



Motivation

◆ These measurements can allow:

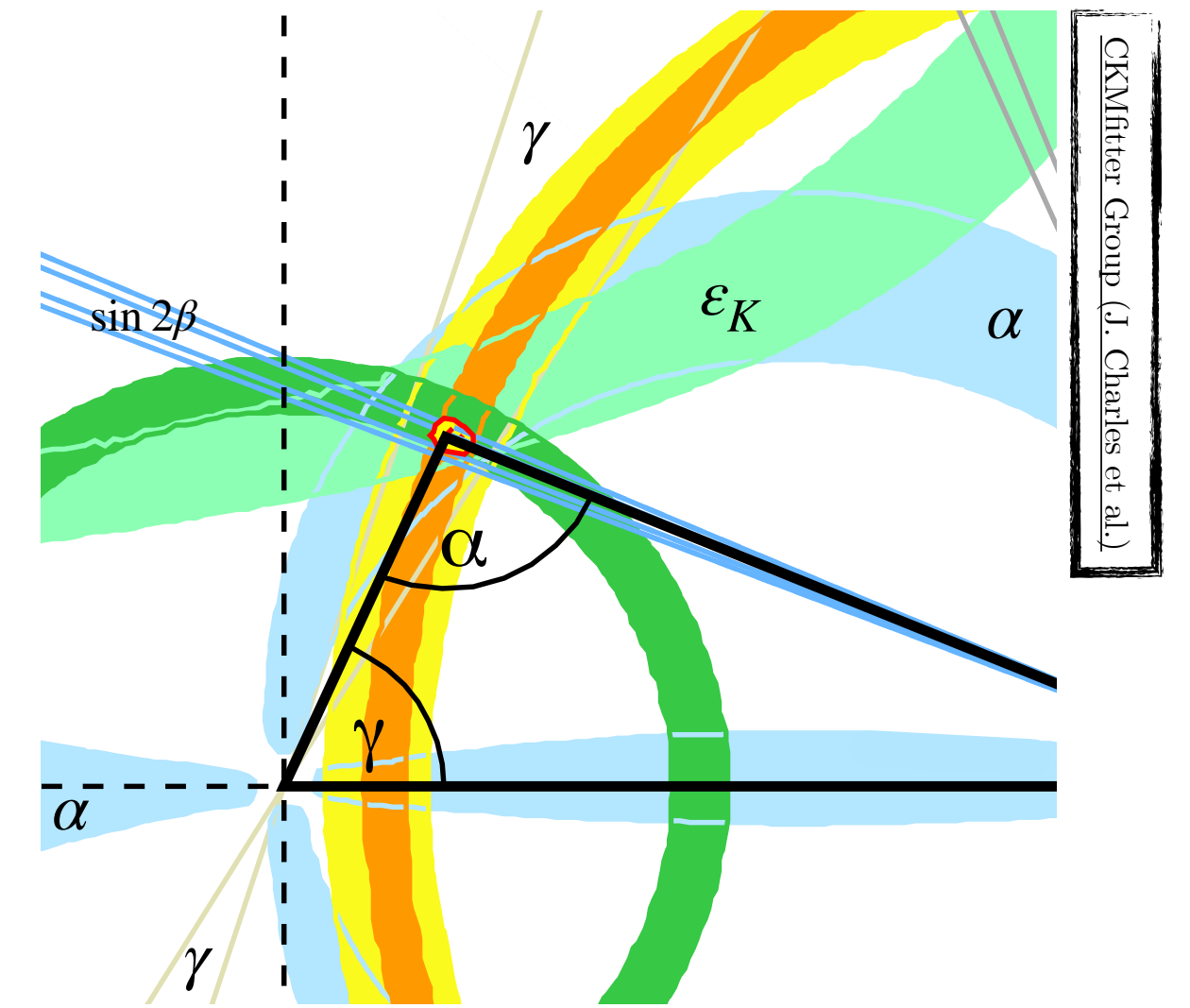
- Constraining the **Unitarity Triangle** of the CKM matrix



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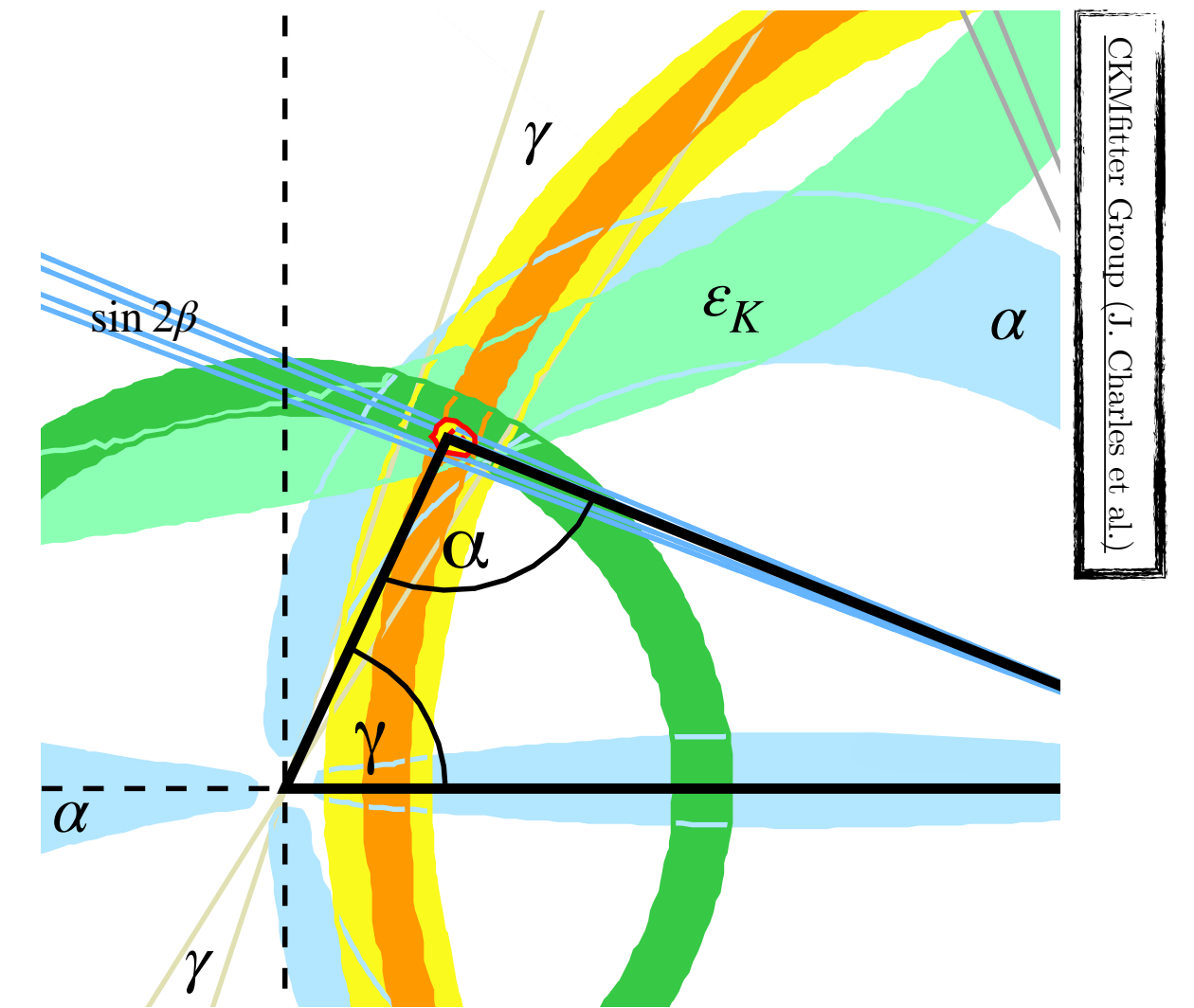
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Motivation

◆ These measurements can allow:

- Constraining the **Unitarity Triangle** of the CKM matrix
- Improving **precision** on **not well known** CKM elements

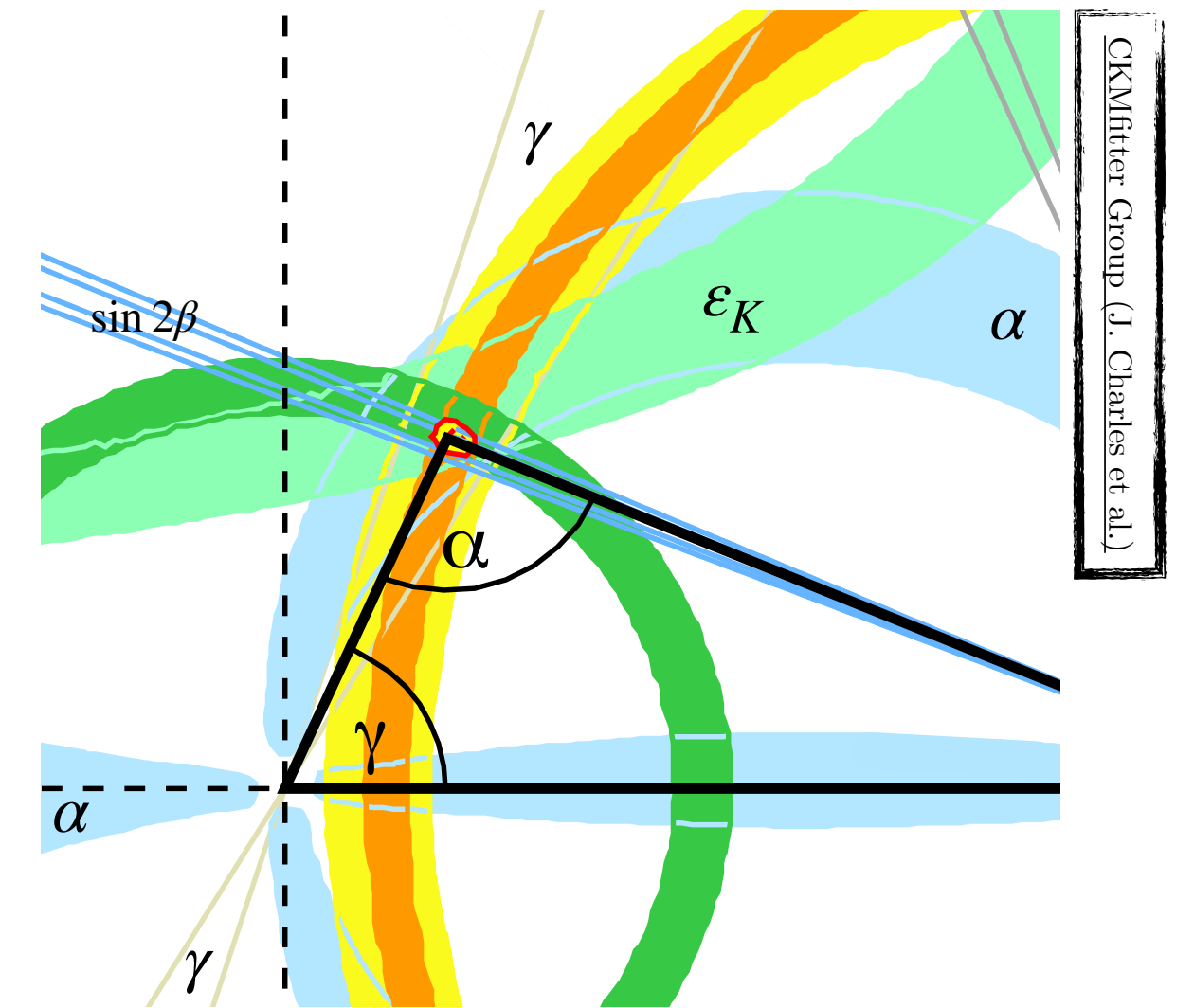


$$|V_{CKM}| = \begin{bmatrix} 0.97435 \pm 0.00016 & 0.22500 \pm 0.00067 & 0.00369 \pm 0.00011 \\ 0.22486 \pm 0.00067 & 0.97349 \pm 0.00016 & 0.04182^{+0.00085}_{-0.00074} \\ 0.00857^{+0.00020}_{-0.00018} & 0.04110^{+0.00083}_{-0.00072} & 0.999118^{+0.000031}_{-0.000036} \end{bmatrix}$$

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- Constraining the **Unitarity Triangle** of the CKM matrix
- Improving **precision** on **not well known** CKM elements
- Probing **New Physics**



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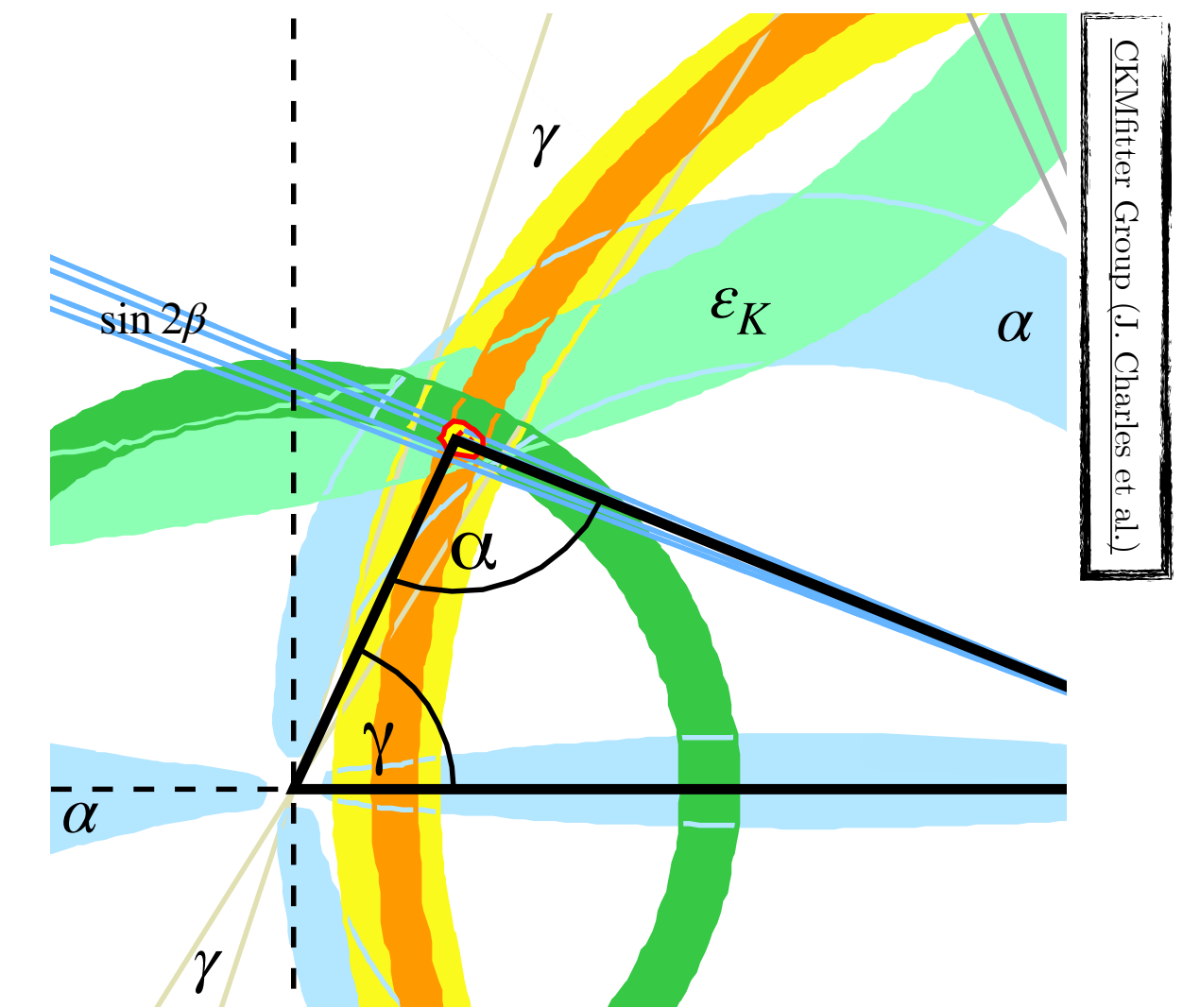
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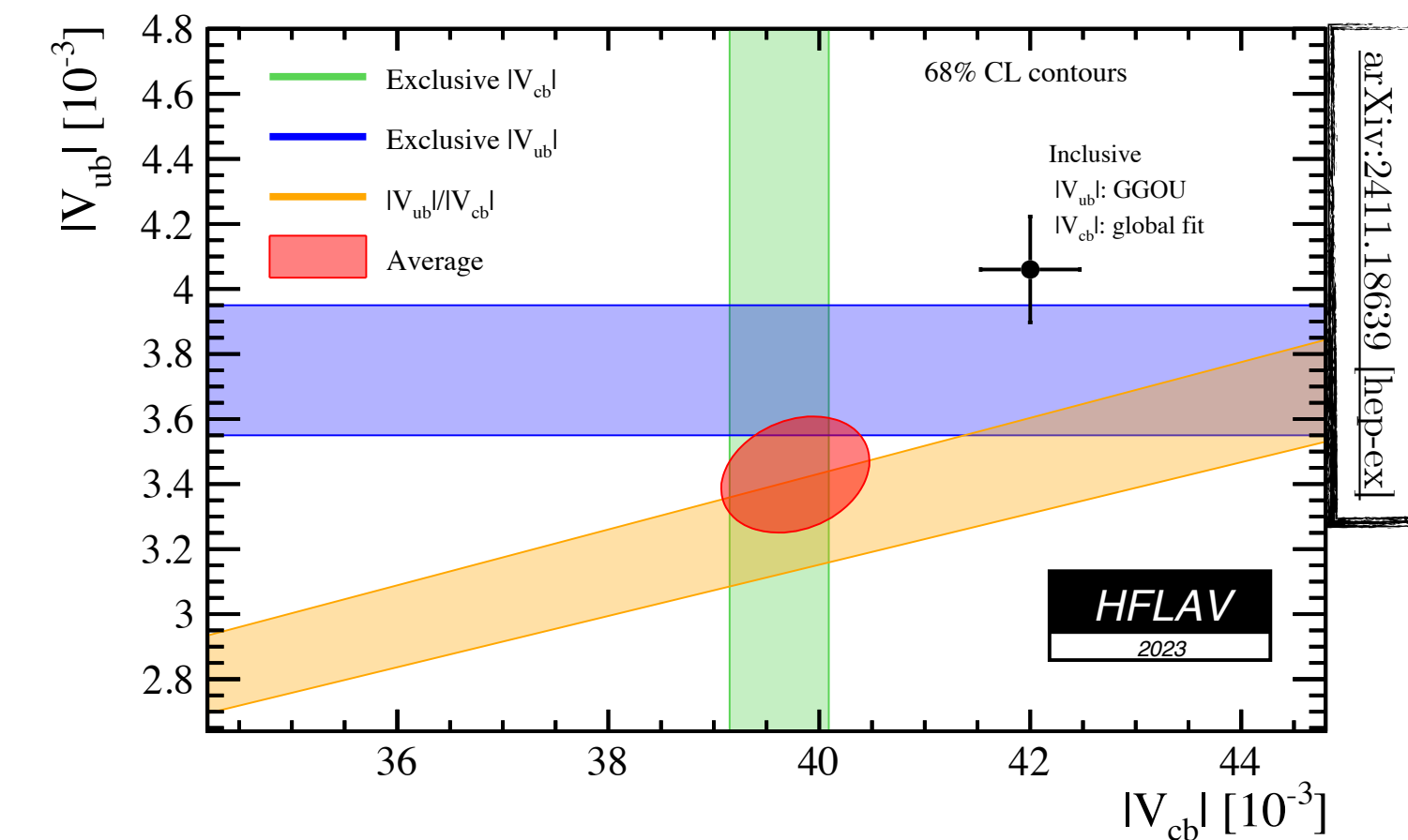
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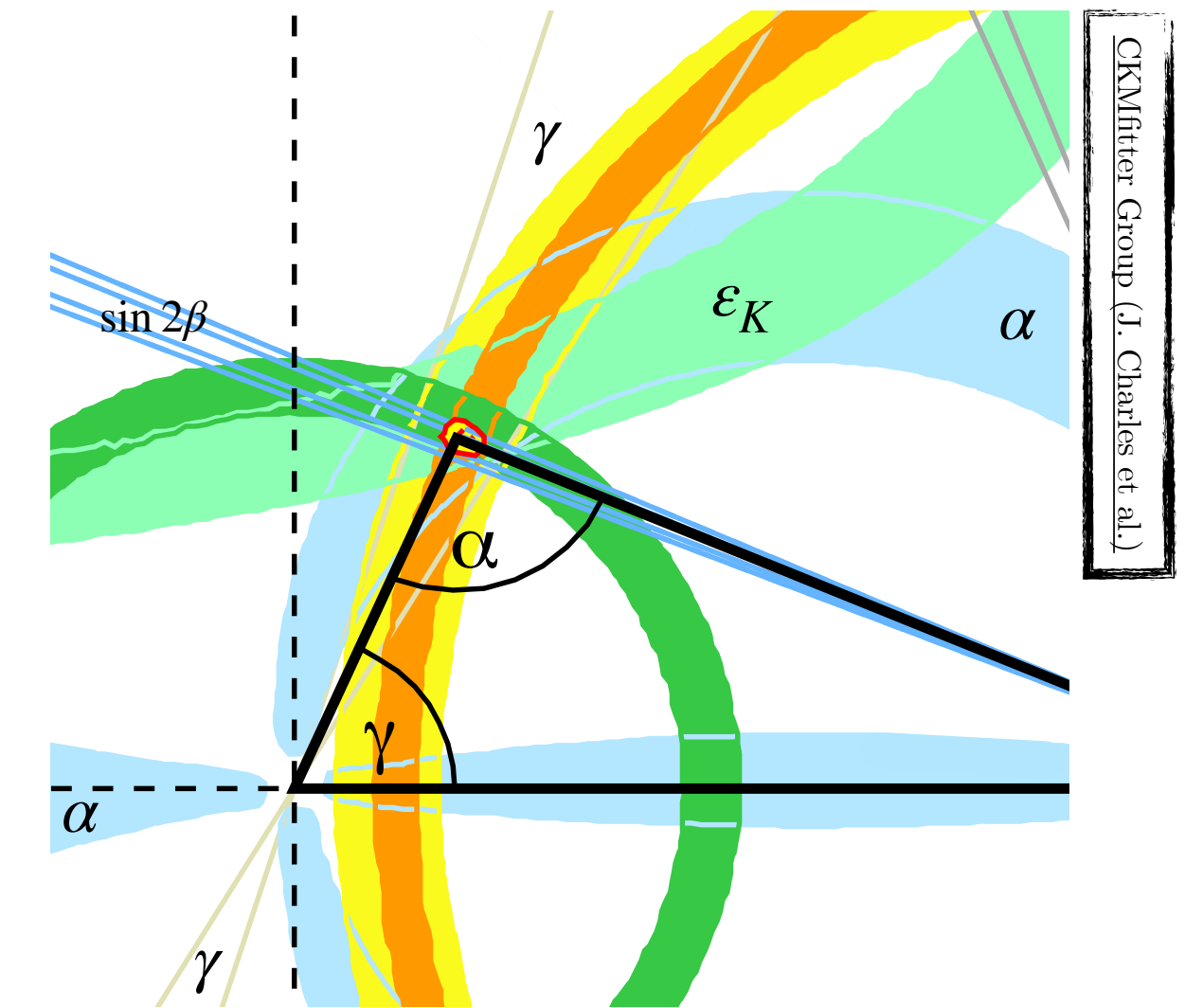
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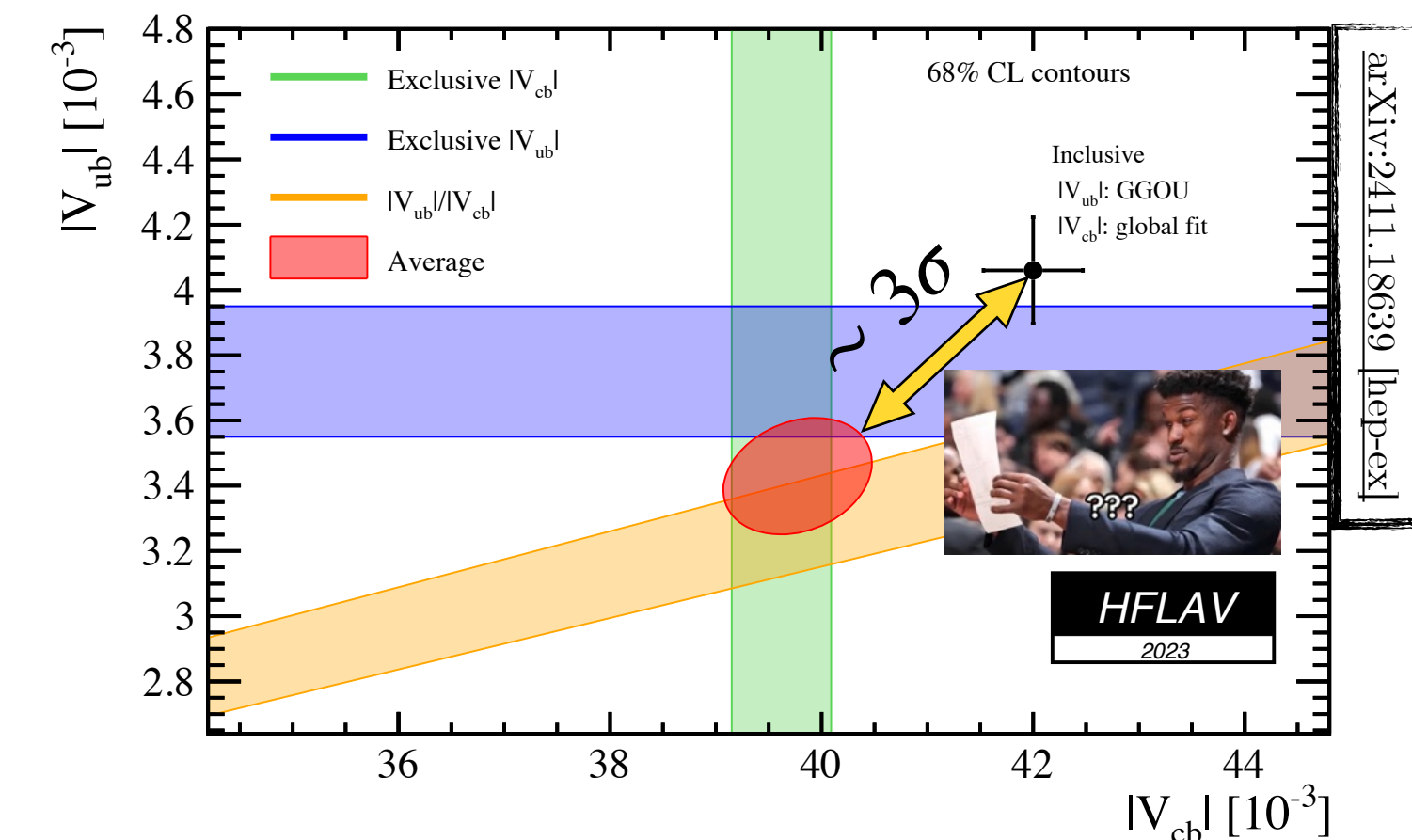
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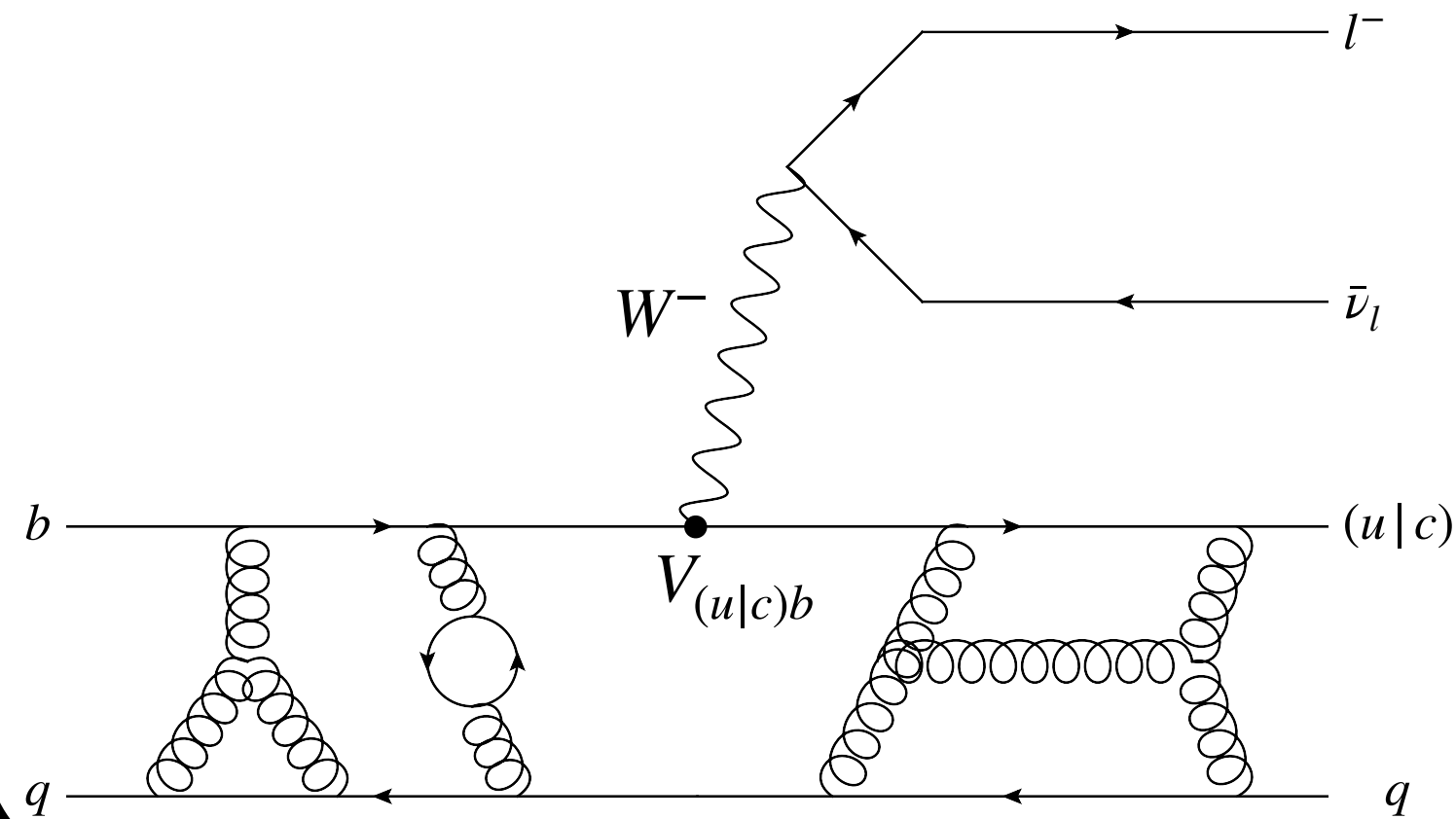


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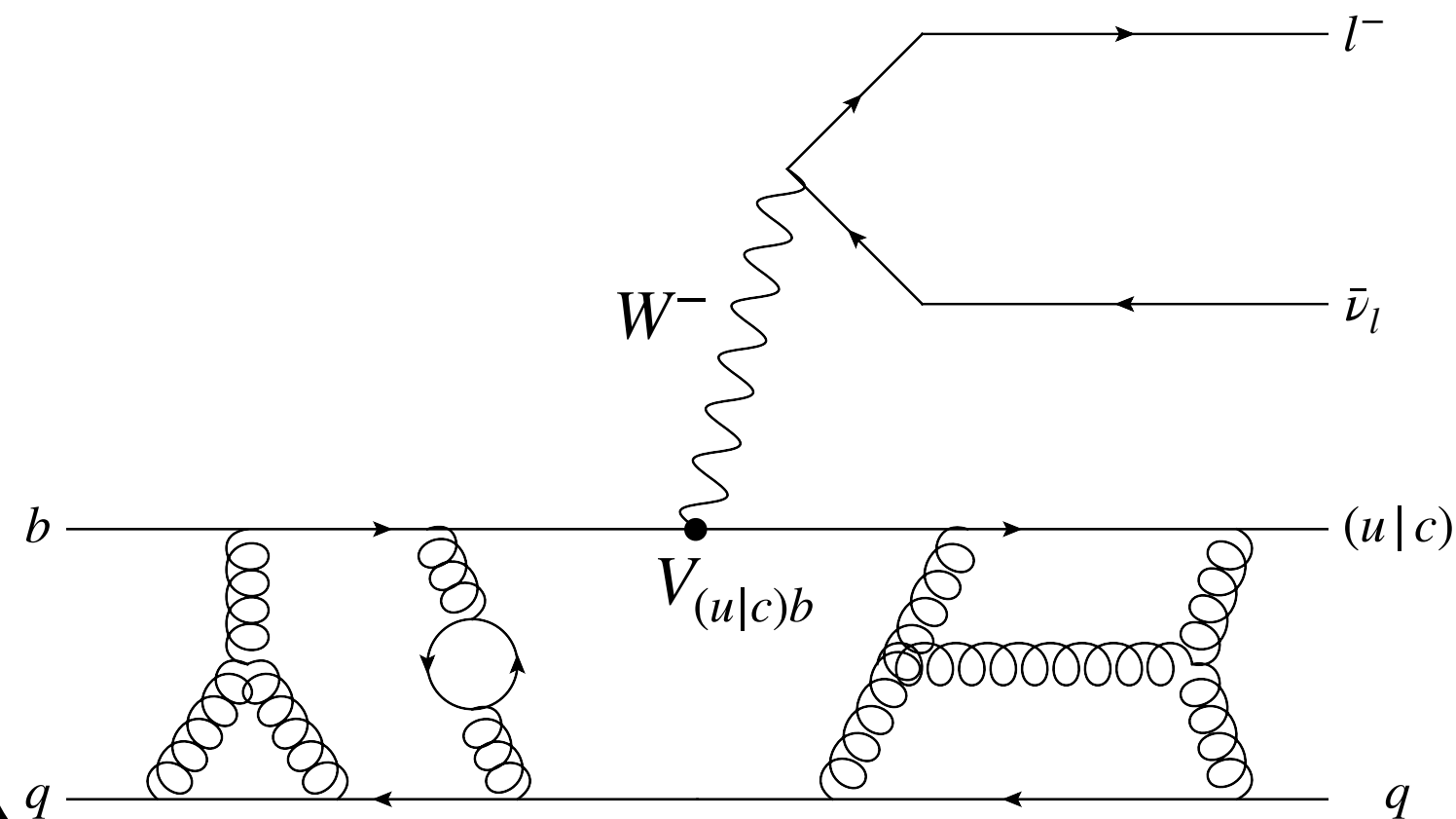
Measurement strategy for $|V_{ub}|$ and $|V_{cb}|$

Semi-Leptonic Decays



Measurement strategy for $|V_{ub}|$ and $|V_{cb}|$

Semi-Leptonic Decays

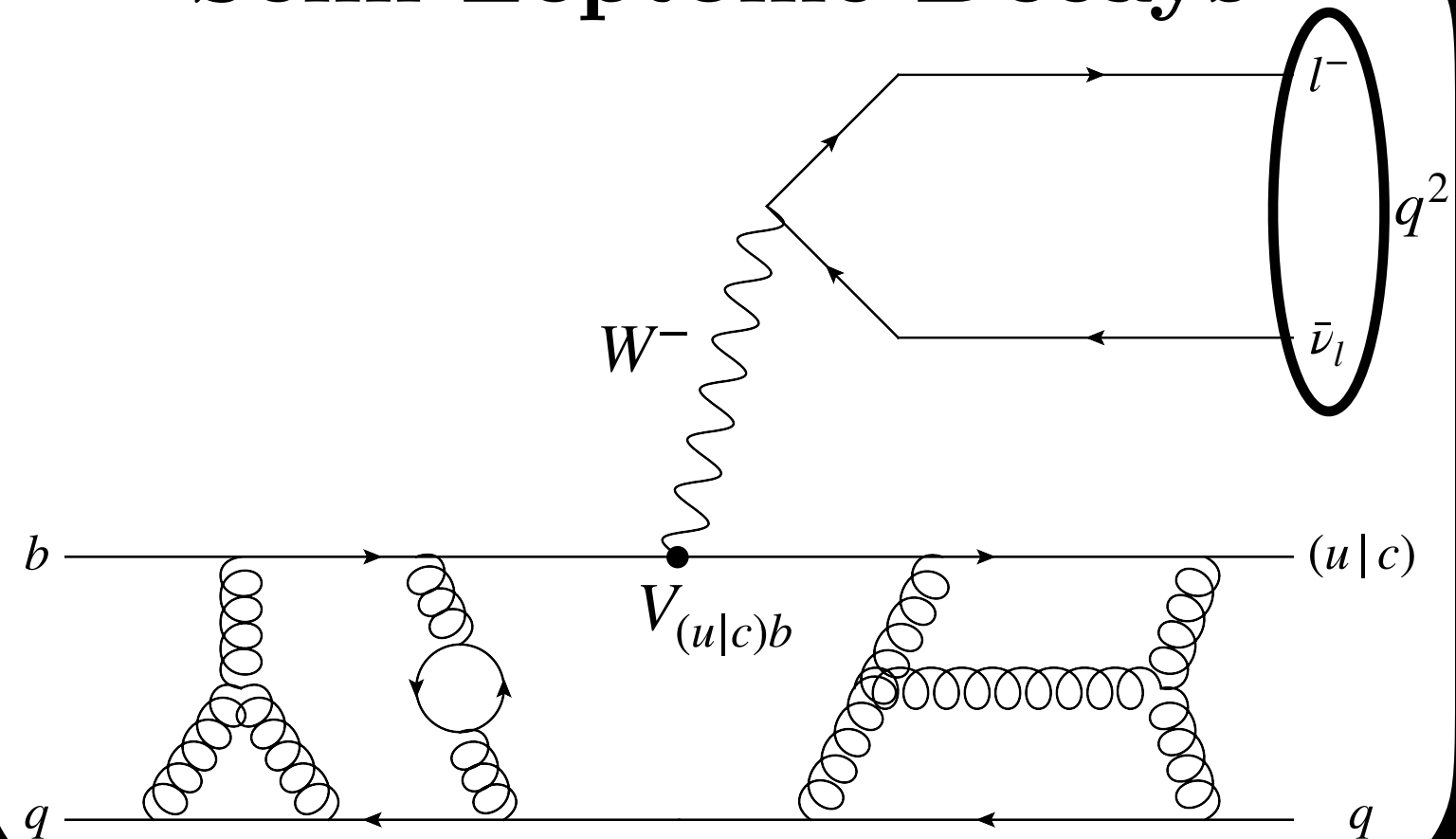


♦ Inclusive decays :

- ◉ Final hadronic states not reconstructed : $B \longrightarrow X_{(u|c)} l \bar{\nu}_l$
- ◉ Experimentally challenging and theoretically cleaner than exclusives
- ◉ $BR = |V_{(u|c)b}|^2 [\Gamma(b \rightarrow (u|c) l \bar{\nu}_l) + 1/m_{c,b} + \alpha_s + \dots]$
Heavy quark expansion

Measurement strategy for $|V_{ub}|$ and $|V_{cb}|$

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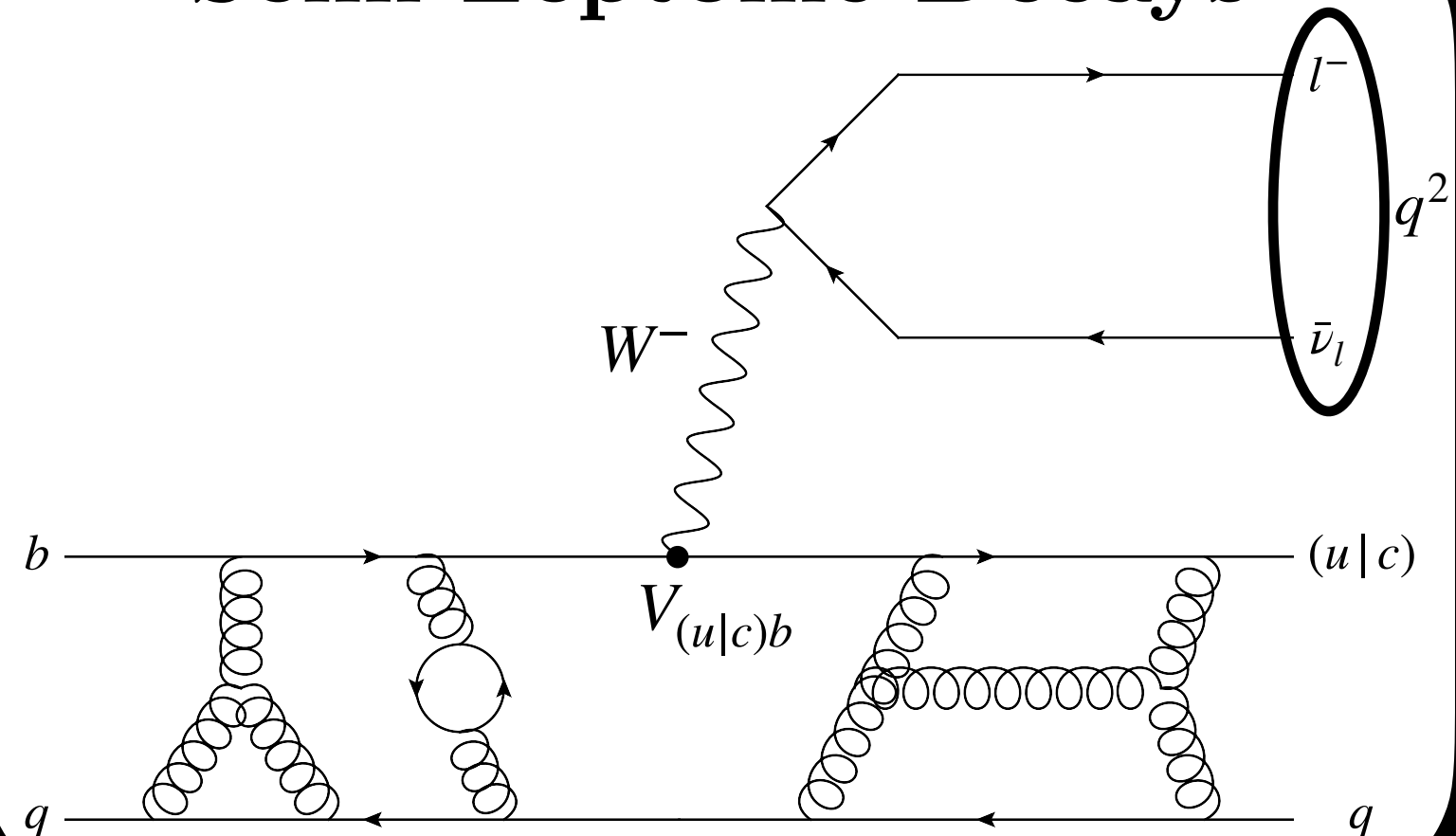
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♦ Exclusive decays :

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- ◉ $BR \propto |V_{(u|c)b}|^2 f(q^2)^2$ with f : form factor

Measurement strategy for $|V_{ub}|$ and $|V_{cb}|$

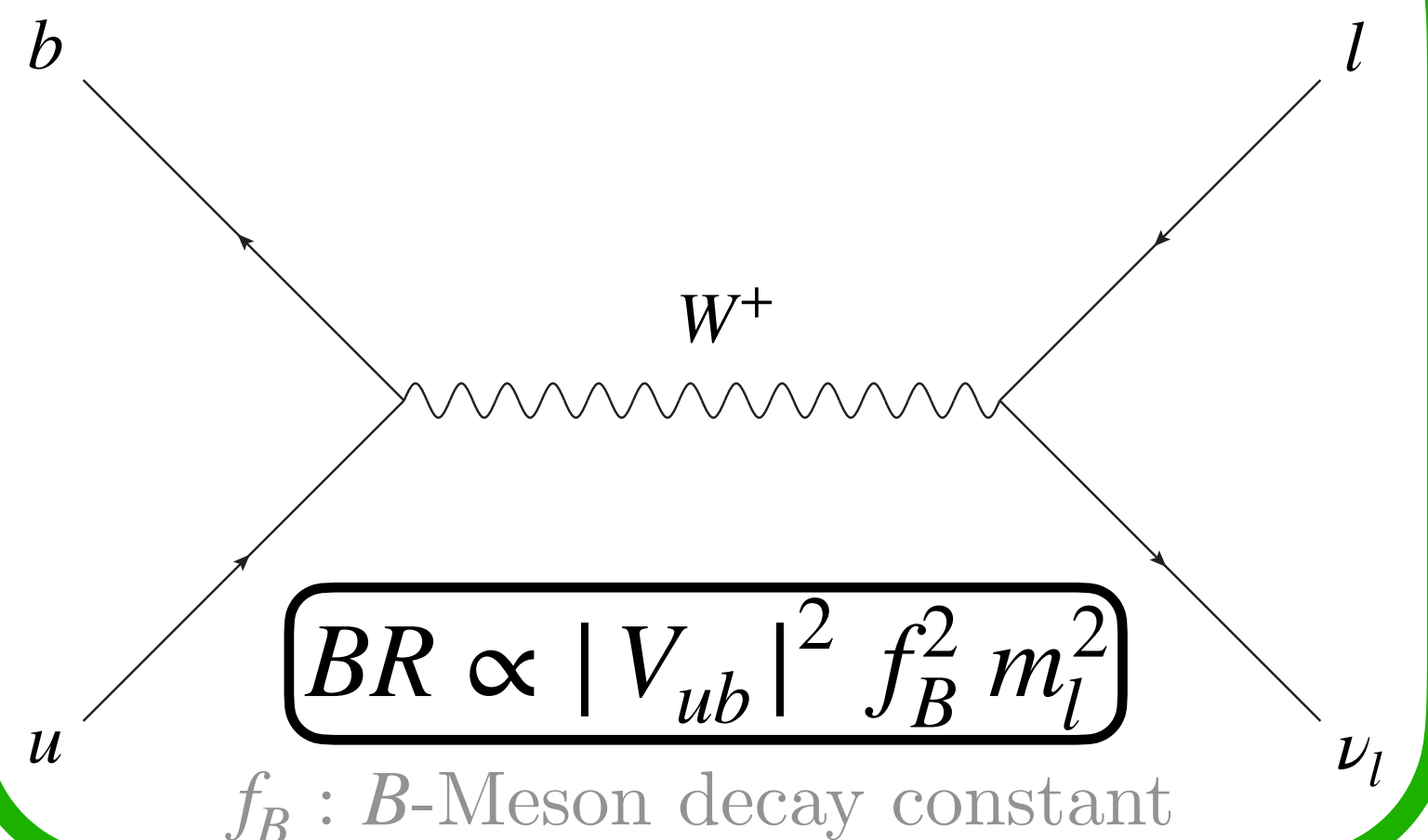
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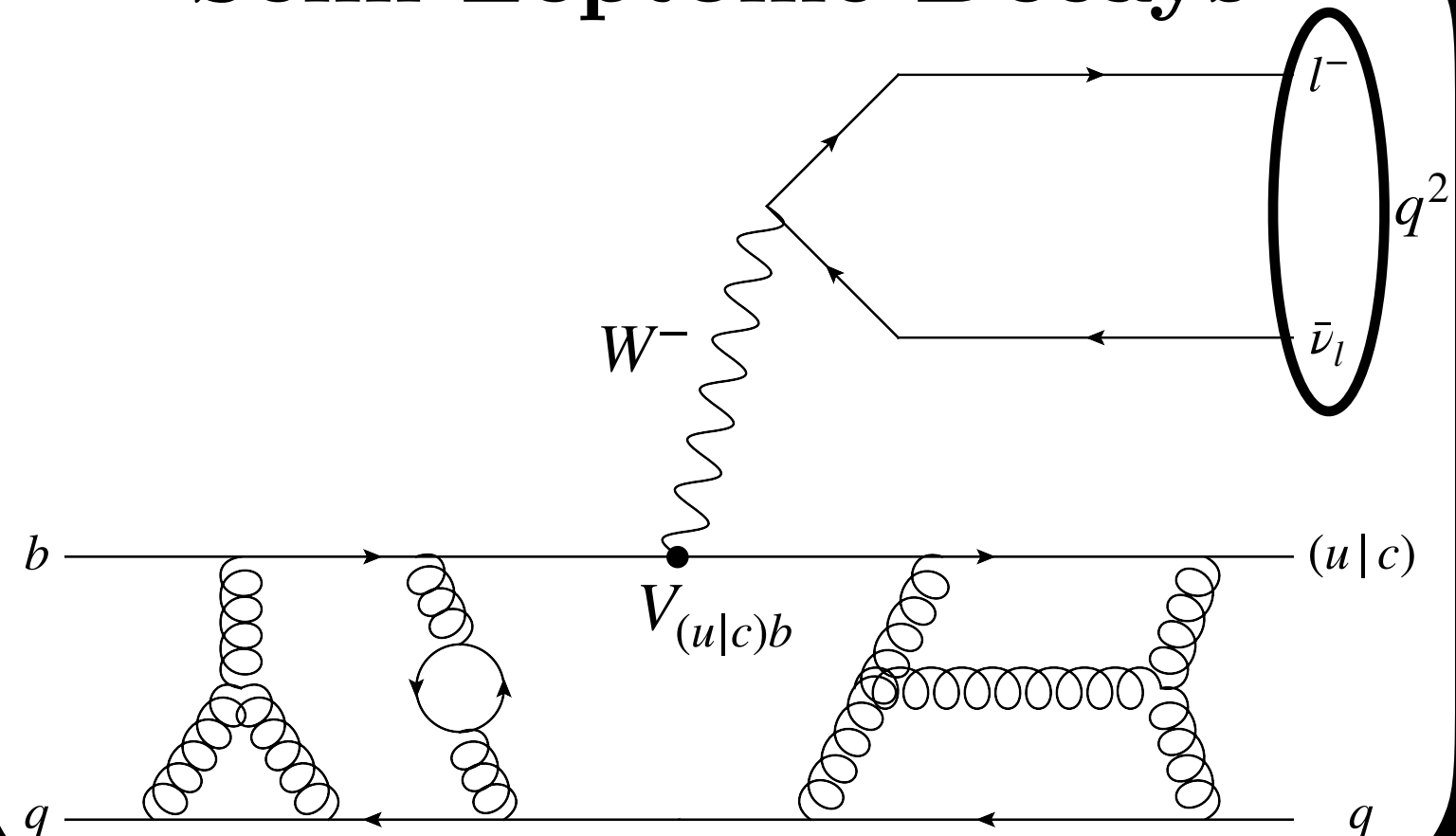


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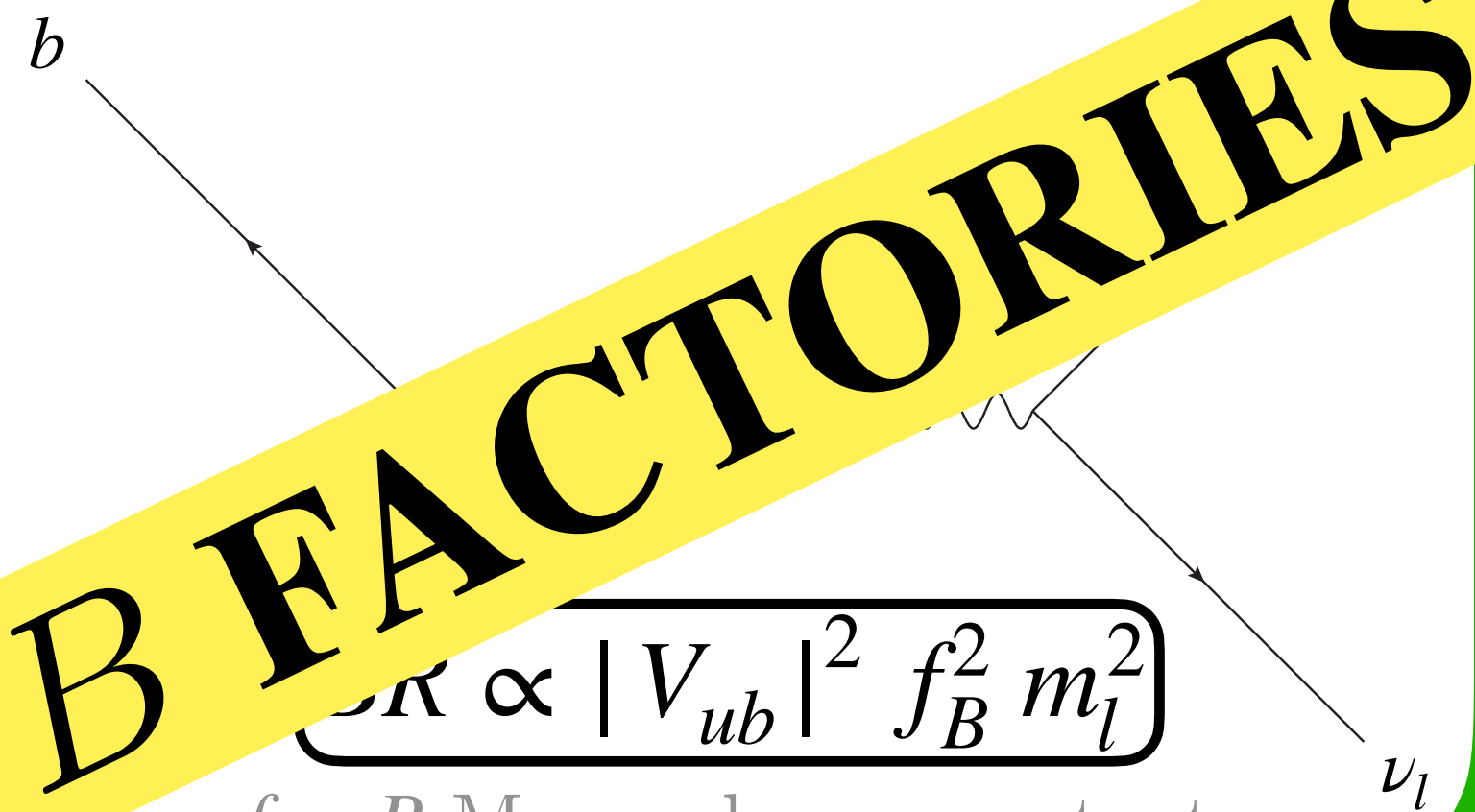
♦ Inclusive decays :

- ◉ Final hadron states

B FACTORIES

$[1/m_{c,b} + \alpha_s + \dots]$
Heavy quark expansion

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B FACTORIES

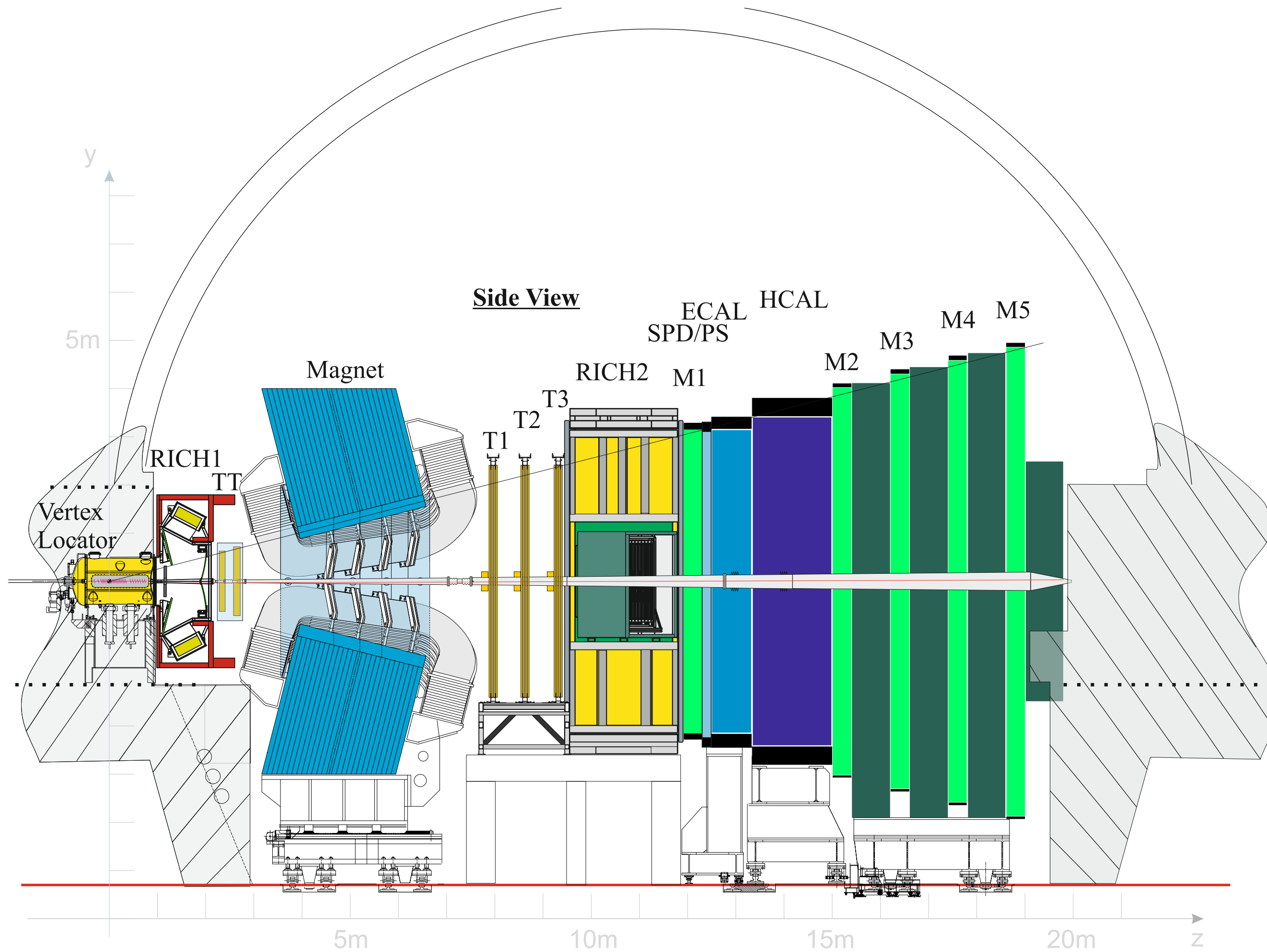
$BR \propto |V_{ub}|^2 f_B^2 m_l^2$

f_B : B -Meson decay constant

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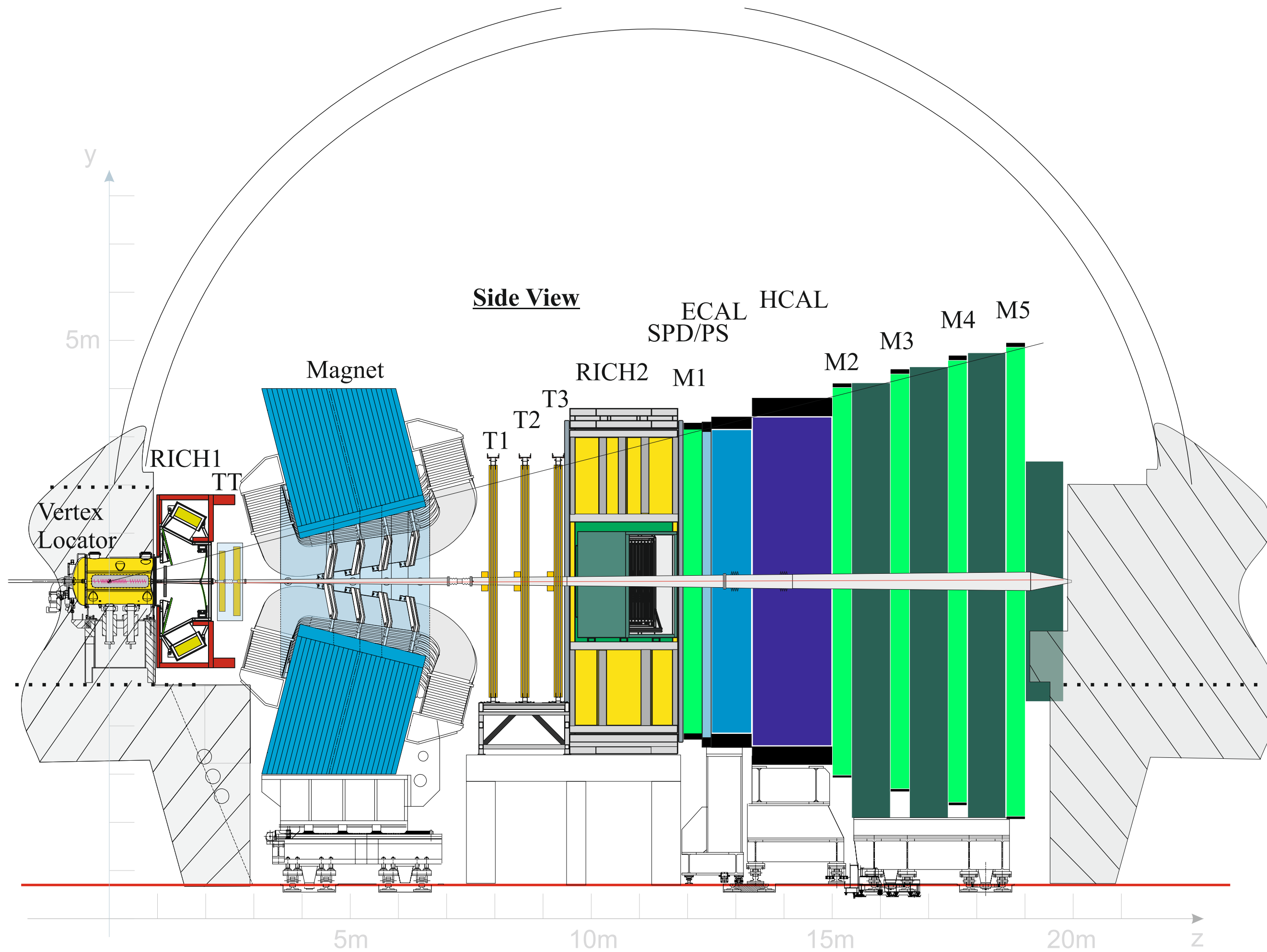
$|V_{ub}|$ and $|V_{cb}|$ measurement @ LHCb



Int. J. Mod. Phys. A 30 (2015), 1530022

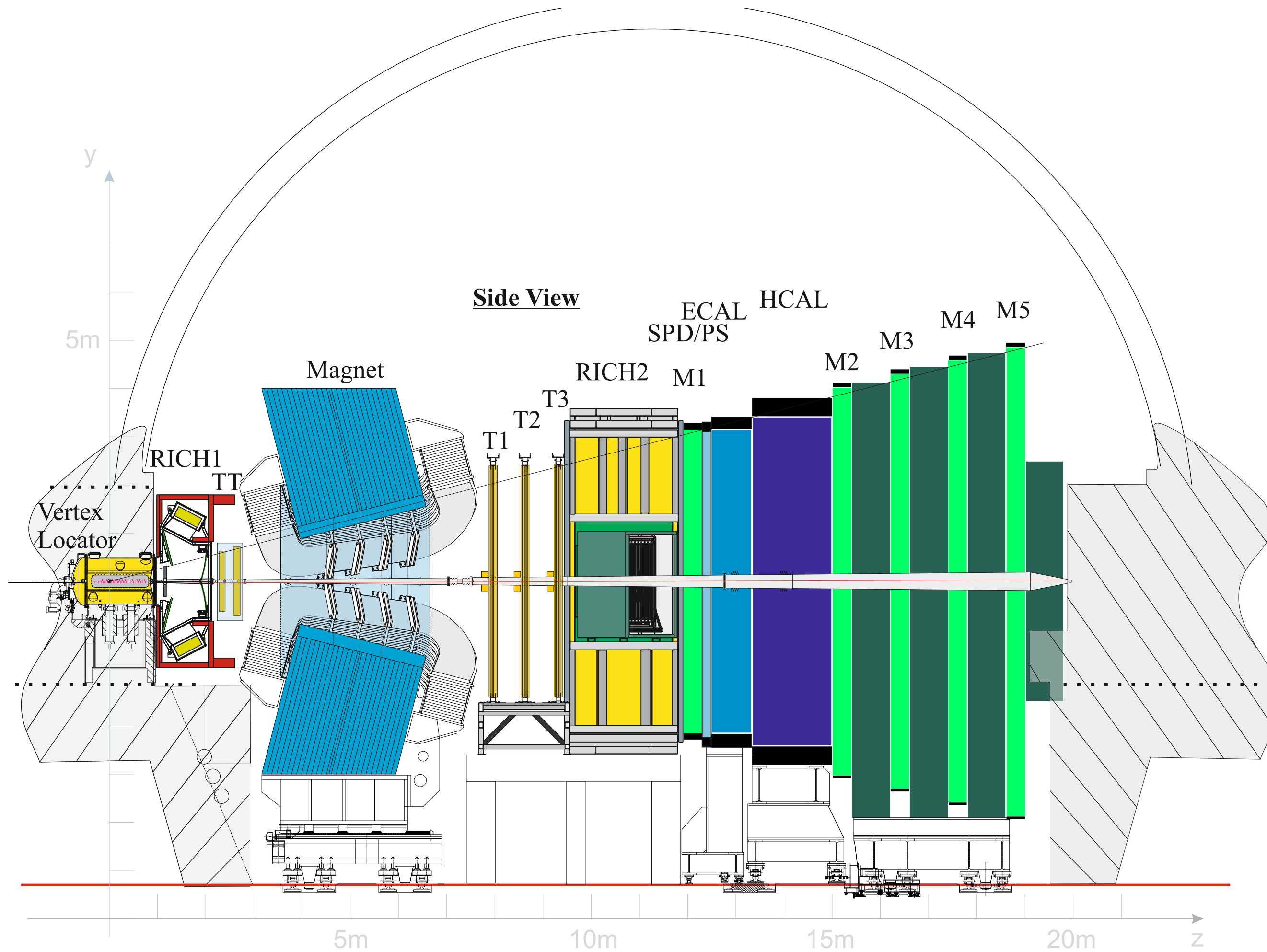
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- ✓ Large samples of $B^{\pm,0}$ mesons and heavier b -hadrons like B_s^0 , B_c and Λ_b^0



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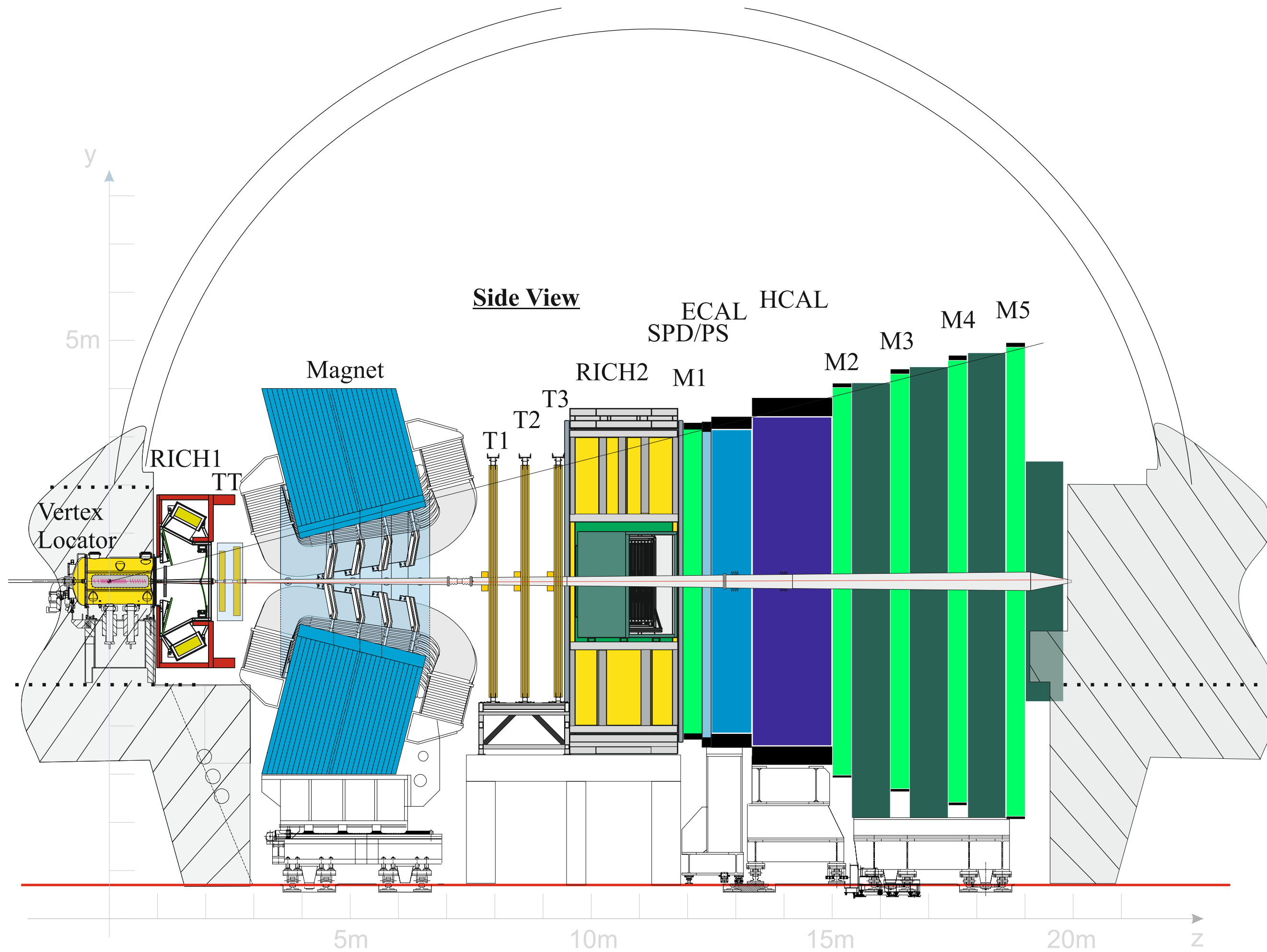
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- ✗ Dirty hadronic environment and unreconstructed $\nu \rightarrow$ **large background**

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Side View

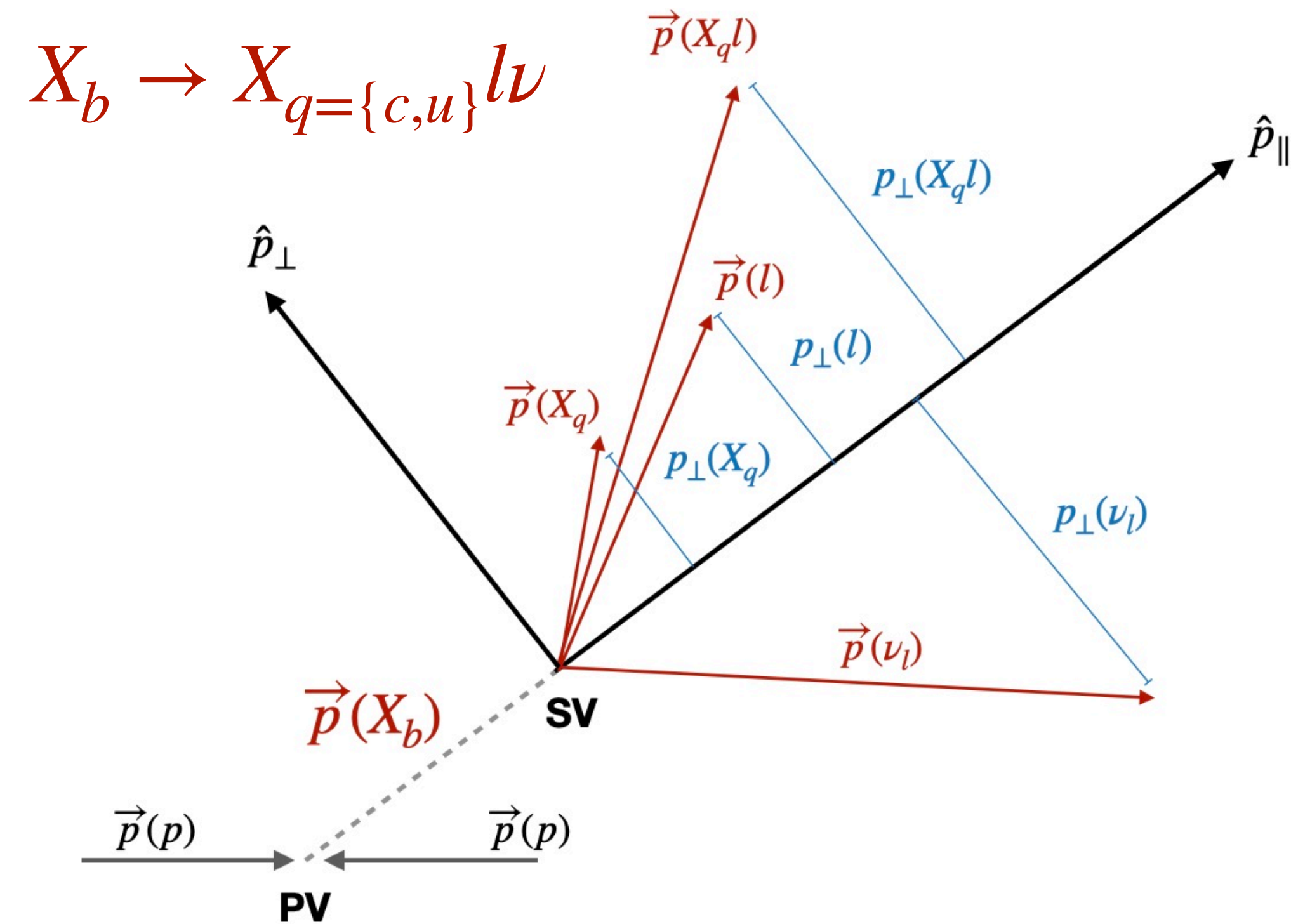
- ✓ Large samples of $B^{\pm,0}$ mesons and heavier b -hadrons like B_s^0 , B_c and Λ_b^0
- × Dirty hadronic environment and unreconstructed $\nu \rightarrow$ large background
- × Luminosity and $b\bar{b}$ production rate cannot be determined precisely $\Rightarrow$ large uncertainty of measured absolute BF's

Int. J. Mod. Phys. A 30 (2015), 1530022

Semileptonic analysis @ LHCb

X_b mass obtained with a reconstruction technique using the **visible decay products** $X_{(u|c)}l$ and **kinematic constraints**

$$\Rightarrow M_{corr}(X_b) = \sqrt{M(X_{(u|c)}l)^2 + p_{\perp}(X_{(u|c)}l)^2 + p_{\perp}(X_{(u|c)}l)}$$



Phys.Rev.Lett.80:660-665,(1998)

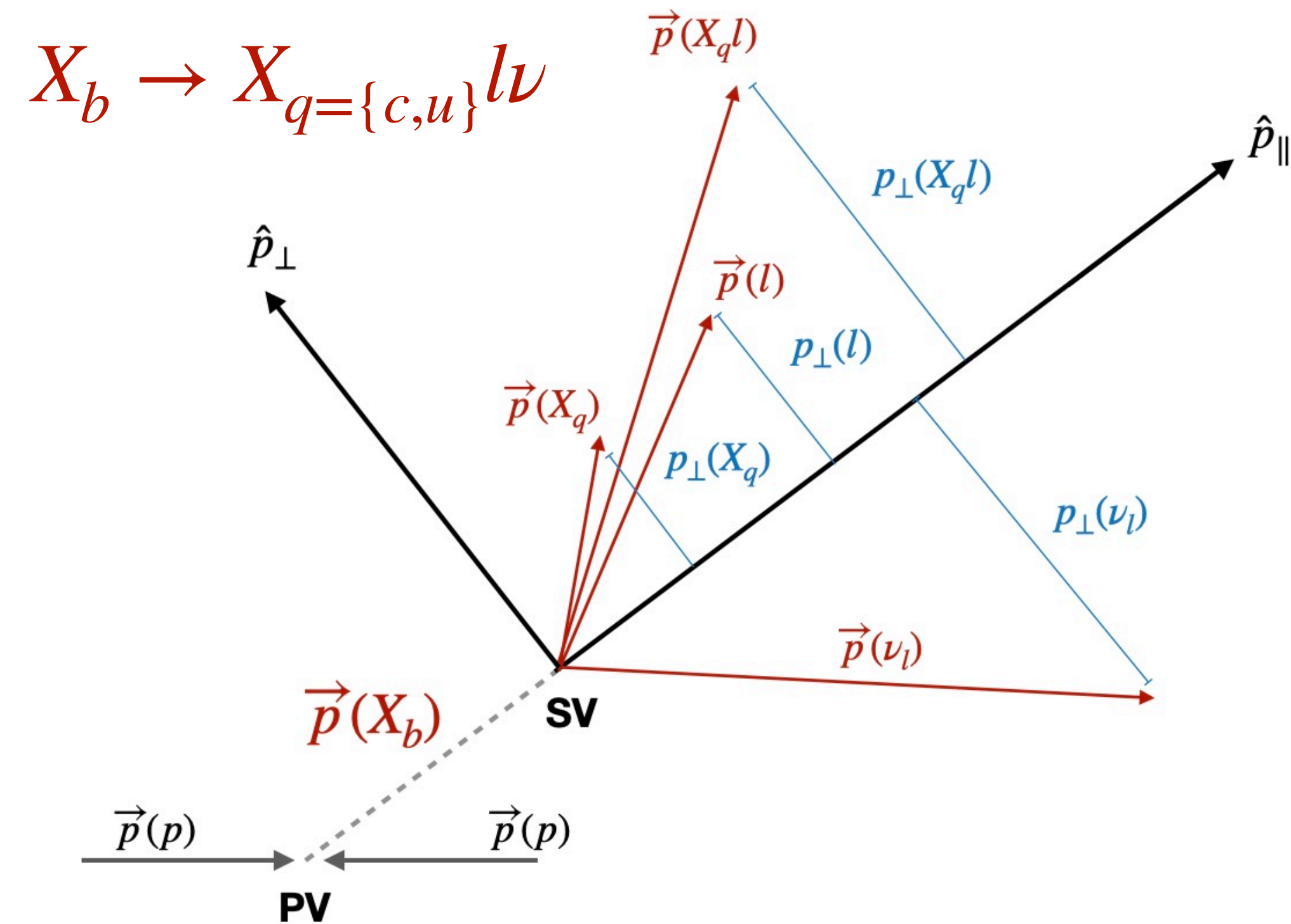
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Two fold ambiguity wrt q^2 leads to a **degraded resolution**

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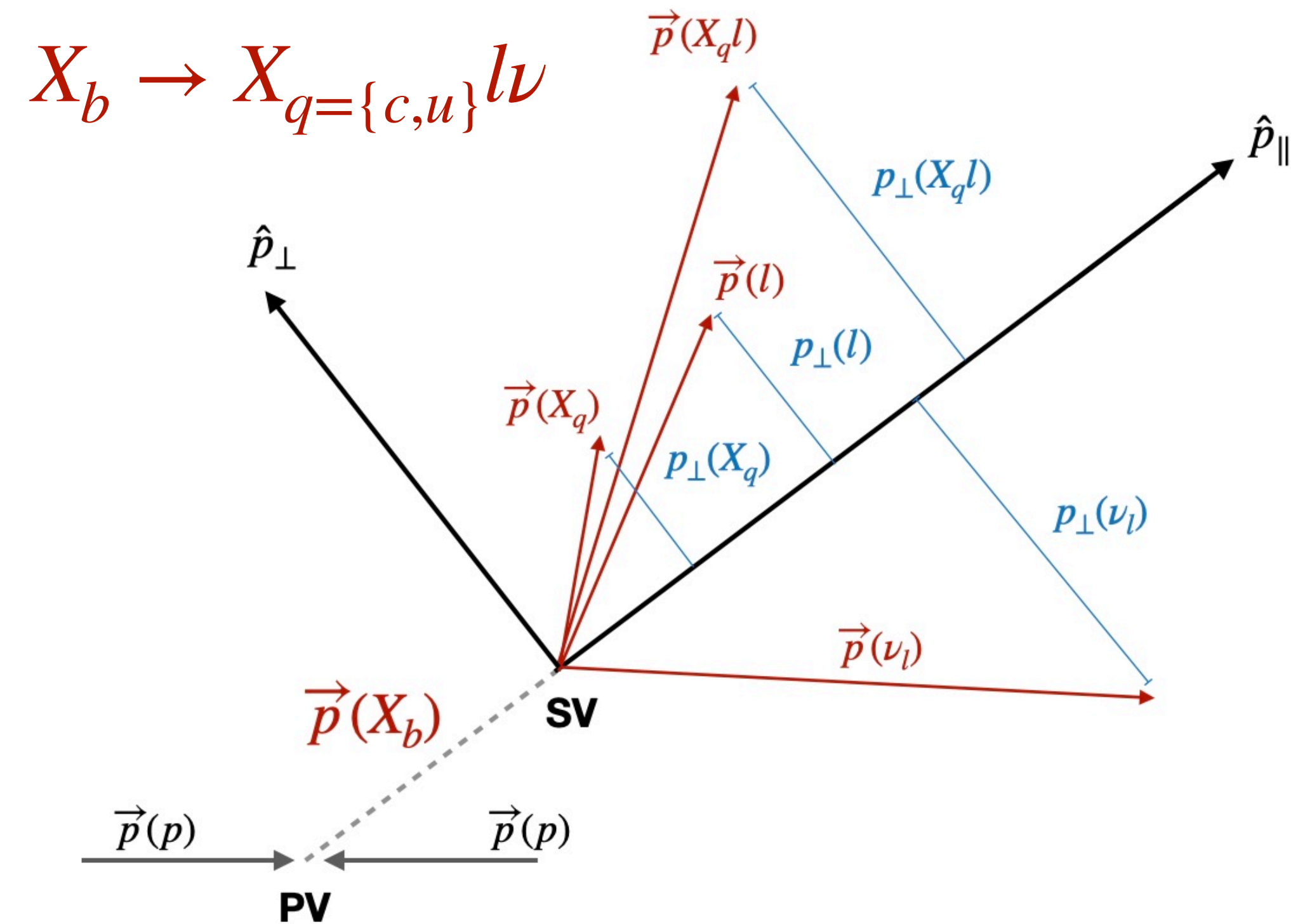
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Normalisation decays used to cancel the $b\bar{b}$ -production uncertainties and reduces the systematic ones $\Rightarrow$ Imported input : BFs Normalisation, fragmentation fraction, ...



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Published $|V_{(u|c)b}|$ measurements at LHCb

$|V_{cb}|$ from $B_s^0 \rightarrow D_s^{(*)-}(\rightarrow [KK]_\phi \pi^-) \mu \nu_\mu$ decays

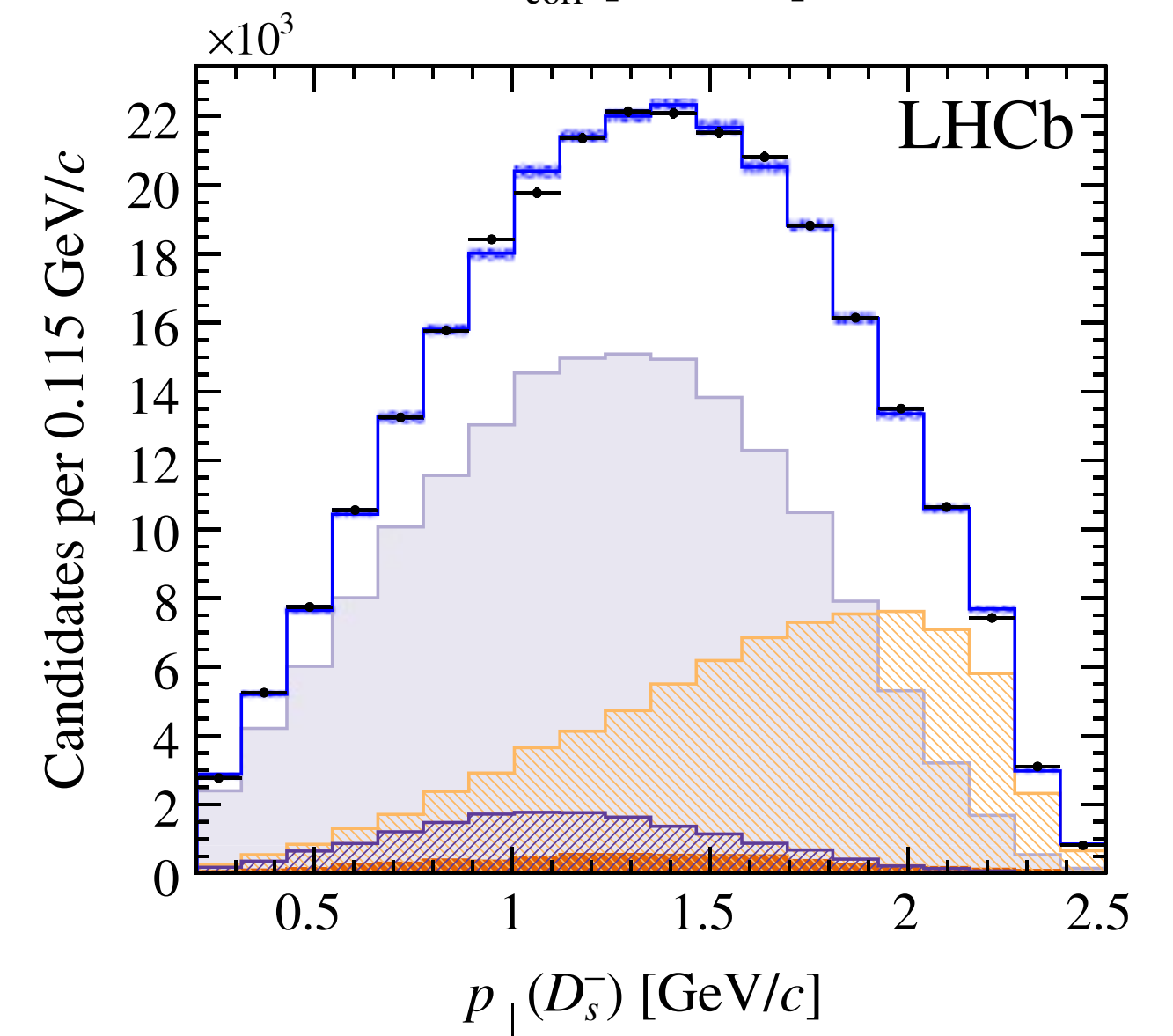
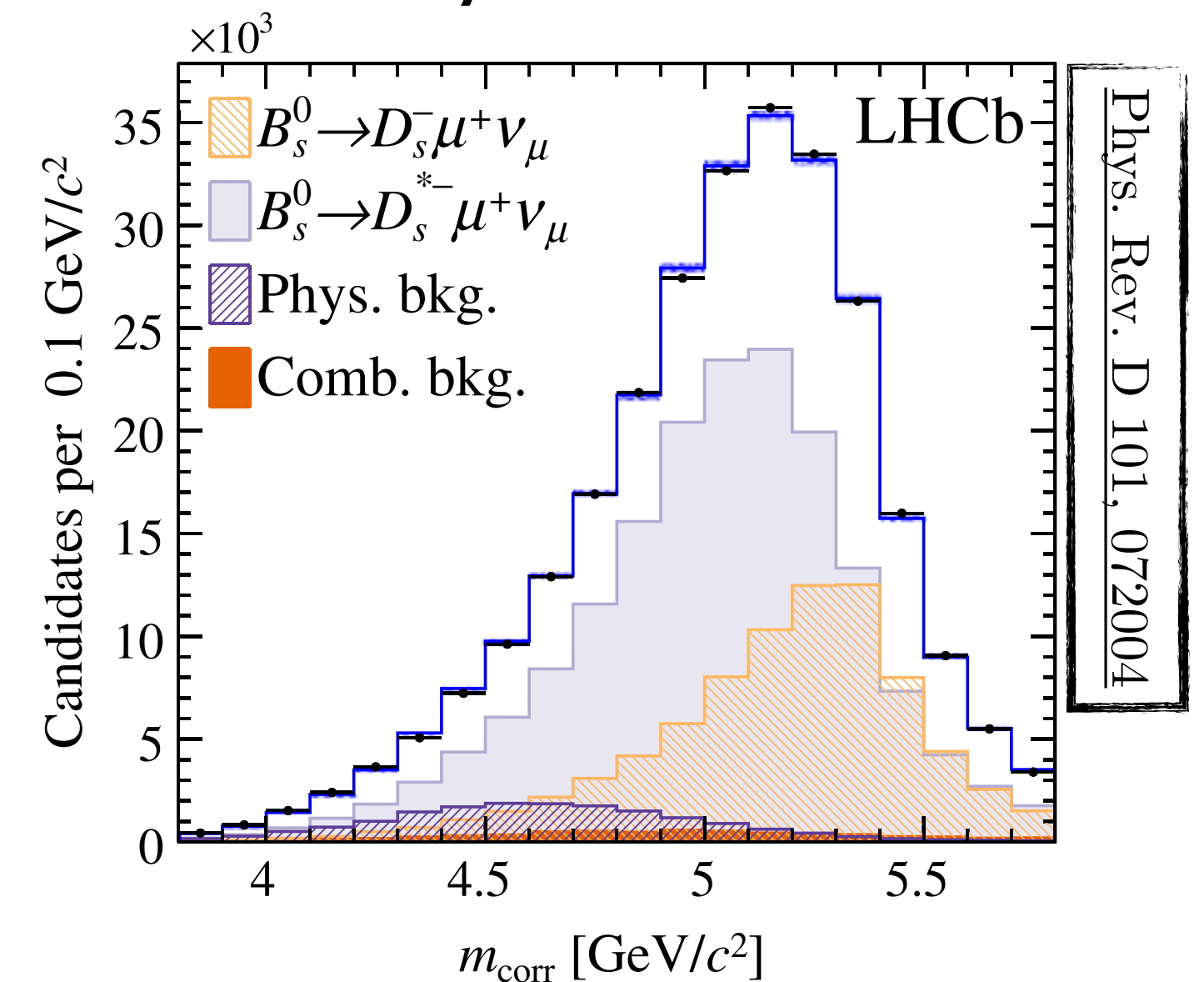
- First $|V_{cb}|$ measurement from B_s^0 decays
- Dataset : 1 fb^{-1} @ $\sqrt{s} = 7 \text{ TeV}$ + 2 fb^{-1} @ $\sqrt{s} = 8 \text{ TeV}$
- Normalisation mode : $B^0 \rightarrow D^{(*)-}(\rightarrow [KK]_\phi \pi^-) \mu \nu_\mu$
- Signal extraction : 2D fit wrt $M_{corr}(B_s^0)$ and $p_\perp(D_s^-)$

Results :

$$|V_{cb}|_{\text{CLN}} = (41.4 \pm 0.6_{\text{stat}} \pm 0.9_{\text{sys}} \pm 1.2_{\text{ext}}) \times 10^{-3}$$

$$|V_{cb}|_{\text{BGL}} = (42.3 \pm 0.8_{\text{stat}} \pm 0.9_{\text{sys}} \pm 1.2_{\text{ext}}) \times 10^{-3}$$

CLN and BGL in mutual agreement



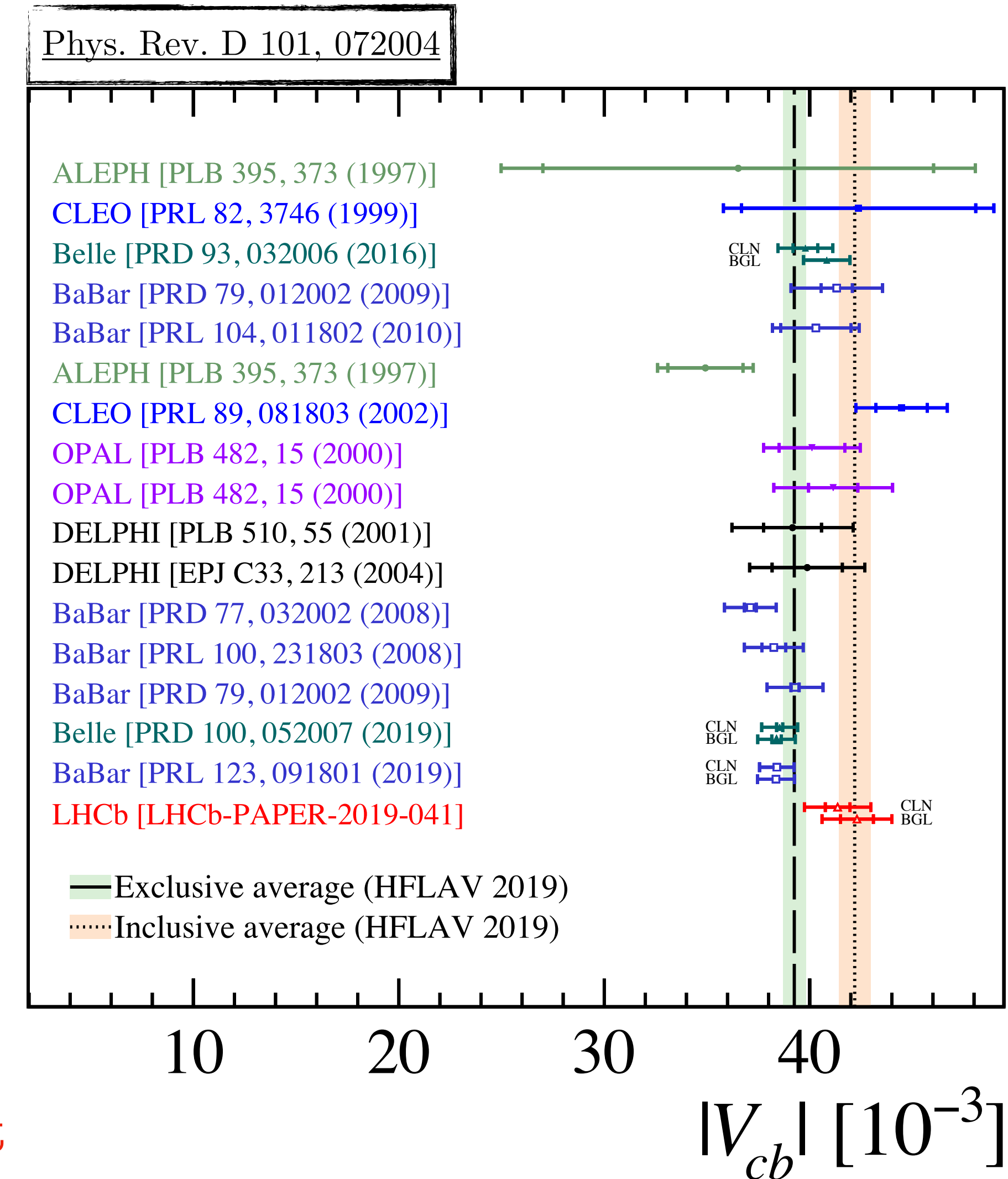
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Agreement with inclusives and exclusives measurement

$|V_{ub}|$ from $\Lambda_b^0 \rightarrow p \mu \nu_\mu$ decays

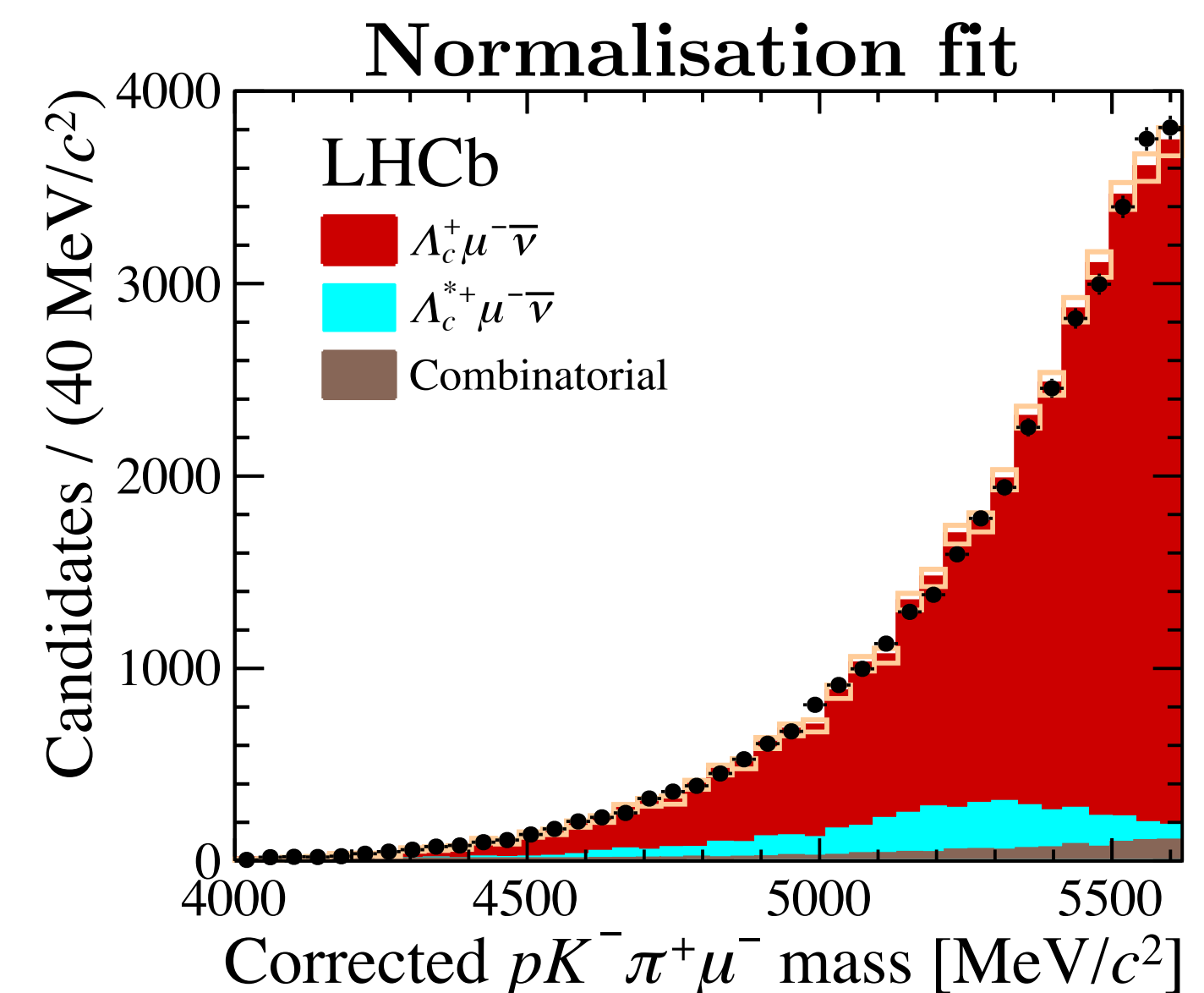
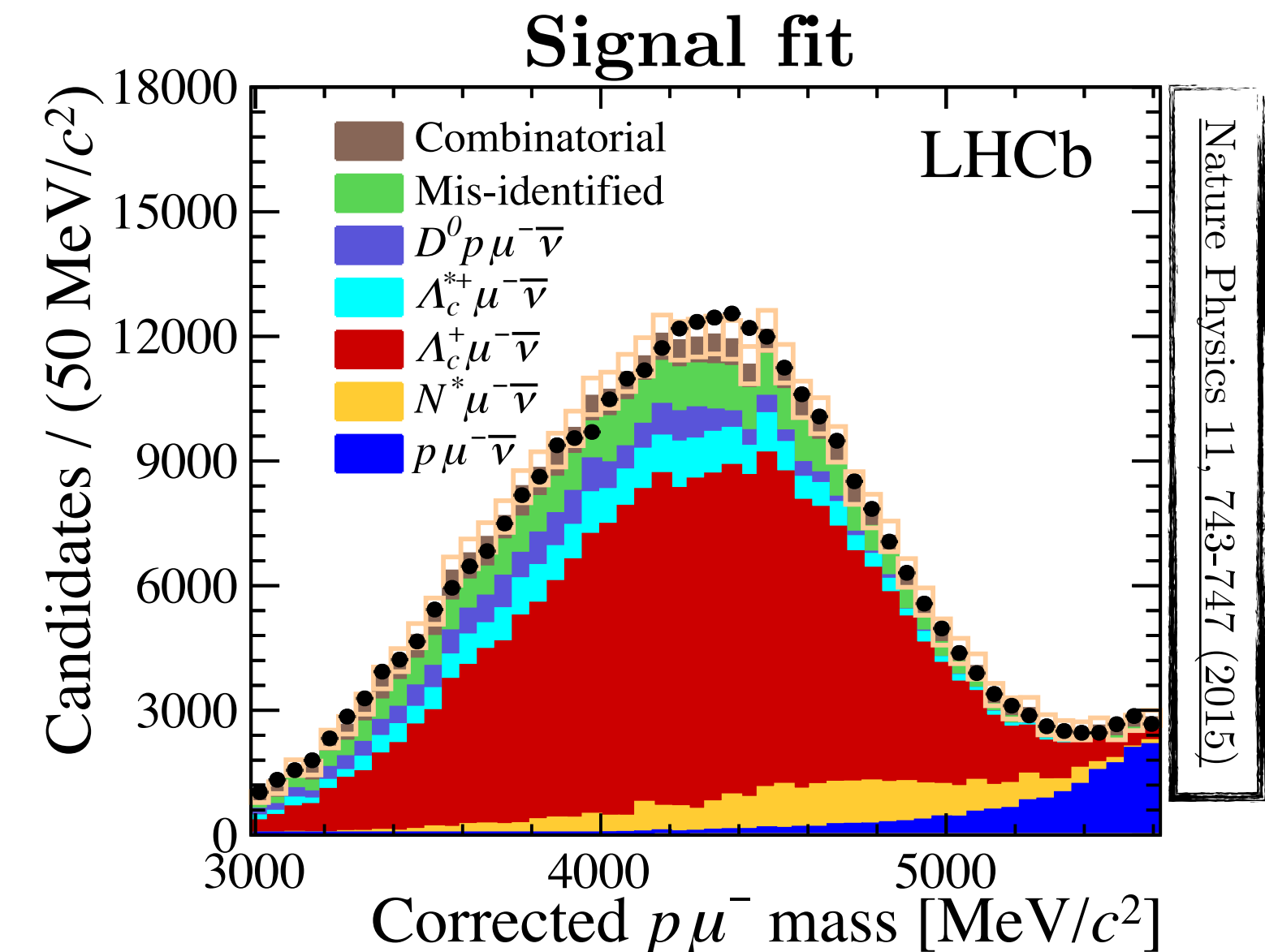
- Dataset : 2 fb^{-1} @ $\sqrt{s} = 8 \text{ TeV}$
- Normalisation mode : $\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow p K^- \pi^+) \mu \nu_\mu$
- Signal extraction : Fit on $M_{corr}(B_s^0)$ wrt the :
 - $p \mu$ system
 - $p K^- \pi^+ \mu^-$ system but only for the normalization mode

Results :

$$\frac{|V_{ub}|}{|V_{cb}|} = 0.083 \pm 0.004_{\text{exp}} \pm 0.004_{\text{theo}}$$

$\rightarrow V_{cb}$ coming from the normalization mode

Form Factors have been estimated by Lattice QCD



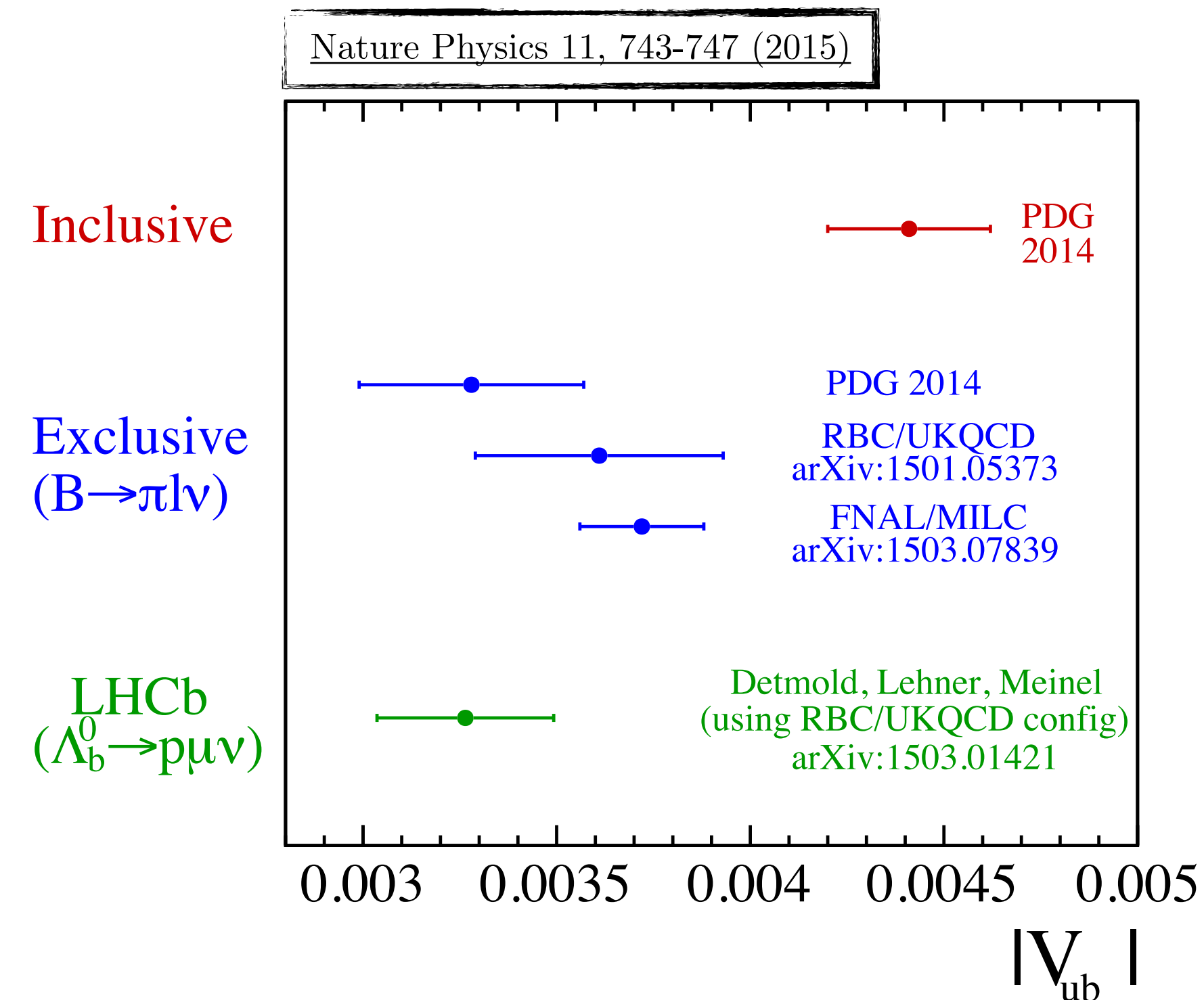
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Results :

$$|V_{ub}| = (3.27 \pm 0.15_{\text{exp}} \pm 0.16_{\text{theo}} \pm 0.06_{\text{norm}}) \times 10^{-3}$$

→ Using the world average $|V_{cb}| = (39.5 \pm 0.8) \times 10^{-3}$ using exclusive decays



Not in agreement with inclusives measurement

$|V_{ub}|/|V_{cb}|$ from $B_s^0 \rightarrow K \mu \nu_\mu$ decays

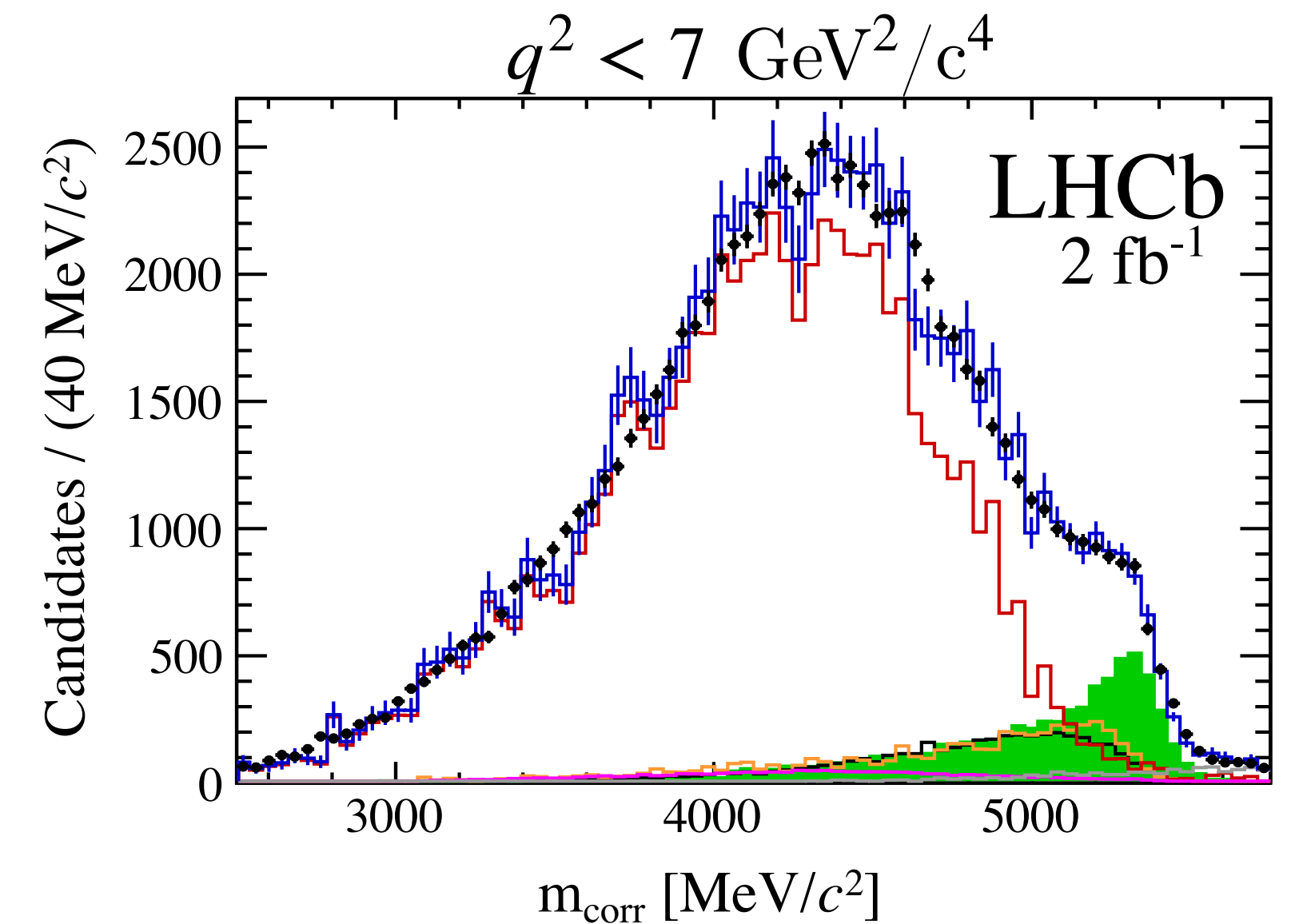
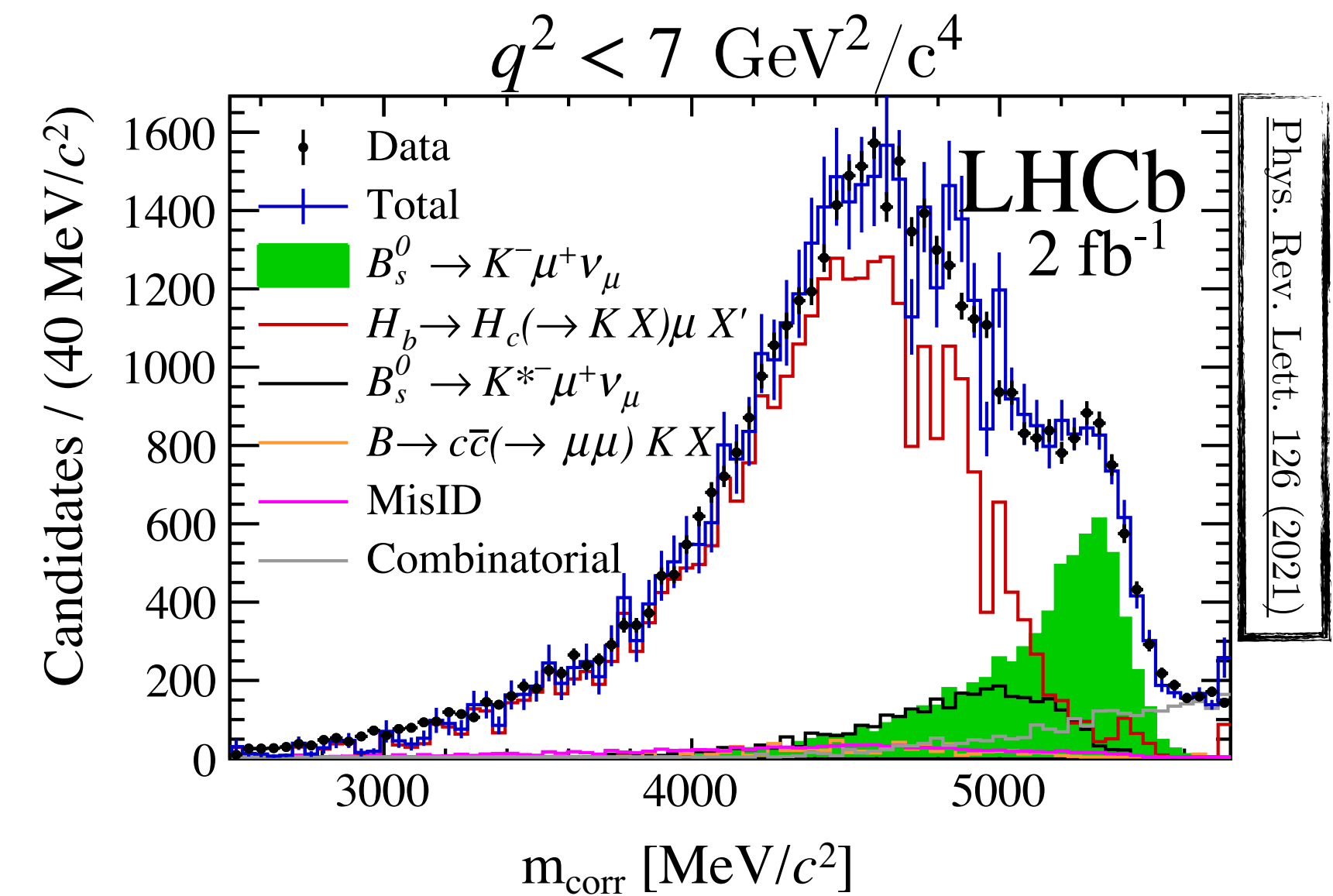
- First observation of the $B_s^0 \rightarrow K \mu \nu_\mu$ decay
- Dataset : 2 fb^{-1} @ $\sqrt{s} = 8 \text{ TeV}$
- Normalisation mode : $B_s^0 \rightarrow D_s^- (\rightarrow KK\pi^-) \mu^+ \nu_\mu$
- Signal extraction : $M_{\text{corr}}(B_s^0)$ in two q^2 regions $\lesssim 7 \text{ GeV}^2/c^4$

Results :

$$|V_{ub}|/|V_{cb}|_{\text{LCSR}, q^2 < 7 \text{ GeV}^2} = 0.0607 \pm 0.0015_{\text{stat}} \pm 0.0013_{\text{sys}} \pm 0.0008_{\text{norm}} \pm 0.0030_{\text{theo}}$$

$$|V_{ub}|/|V_{cb}|_{\text{LQCD}, q^2 > 7 \text{ GeV}^2} = 0.0946 \pm 0.0030_{\text{stat}} \pm {}^{+0.0024}_{-0.0025}_{\text{sys}} \pm 0.0013_{\text{norm}} \pm 0.0068_{\text{theo}}$$

Form Factors has been estimated by
LCSR at low q^2 and LQCD at high q^2



$|V_{ub}|/|V_{cb}|$ from $B_s^0 \rightarrow K \mu \nu_\mu$ decays

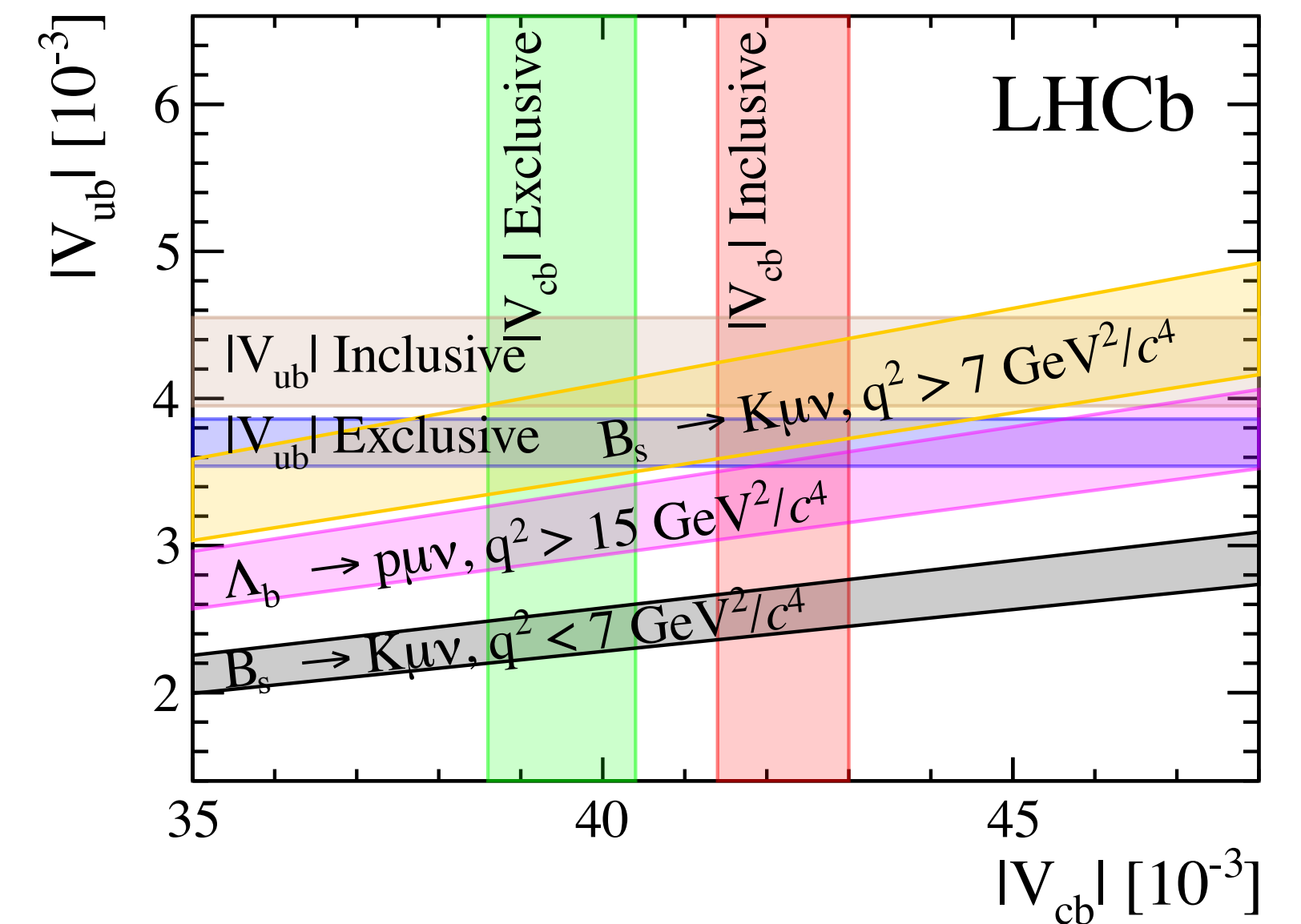
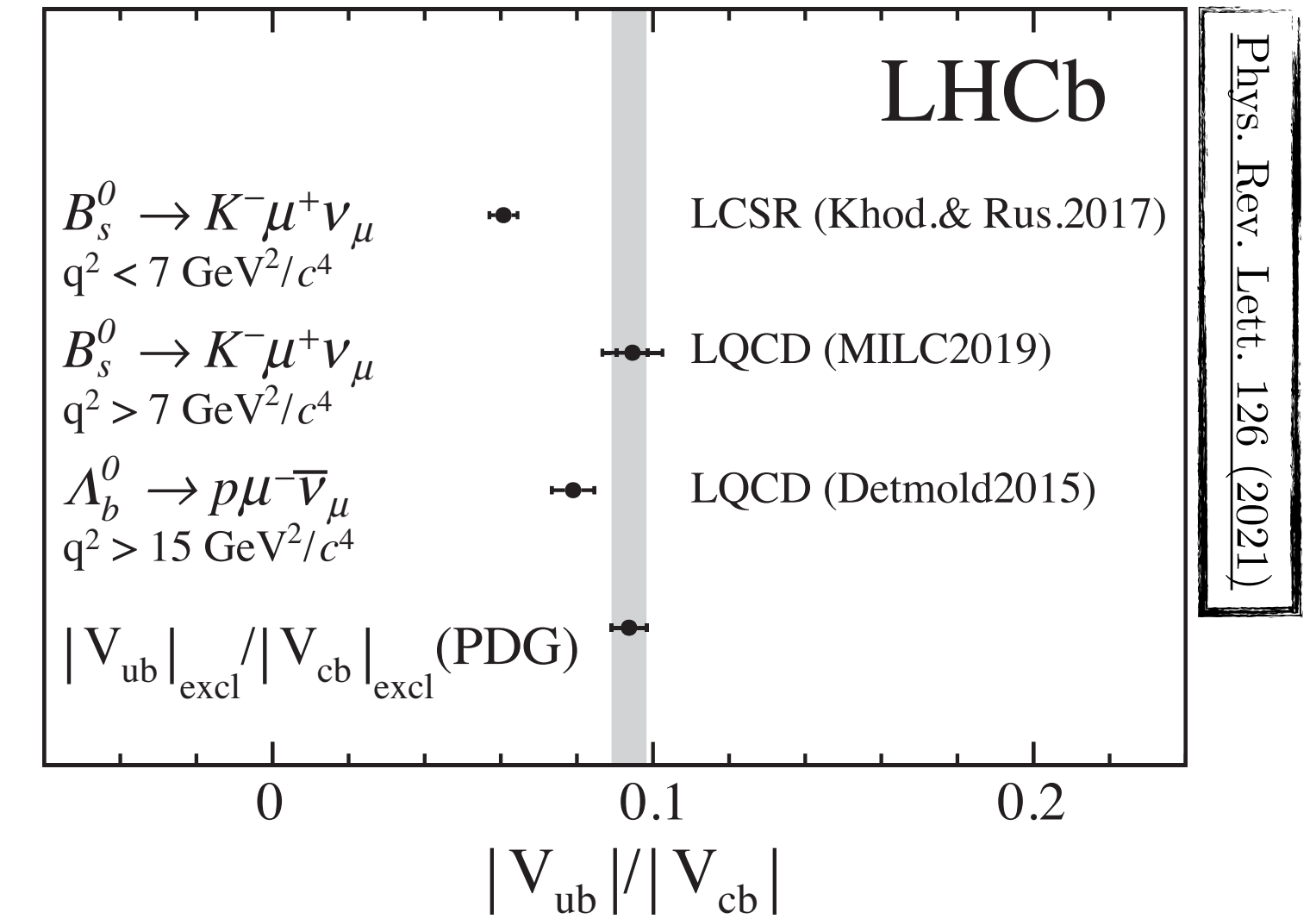
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$|V_{ub}|$ at low q^2 not in agreement with the incl/excl averages



Future $|V_{(u|c)b}|$
measurements at
LHCb

Future Run 2 | $V_{(u|c)b}$ | measurements at LHCb

- ◉ **Extracting $|V_{ub}|$** and a part of the $B^+ \rightarrow \rho^0$ BCL Form Factors from $B^+ \rightarrow \rho^0 \mu \nu_\mu$ in 10 q^2 bins
→ Expecting > 50 times higher signal yield wrt Belle [Phys. Rev. D 88, 032005 (2013)]

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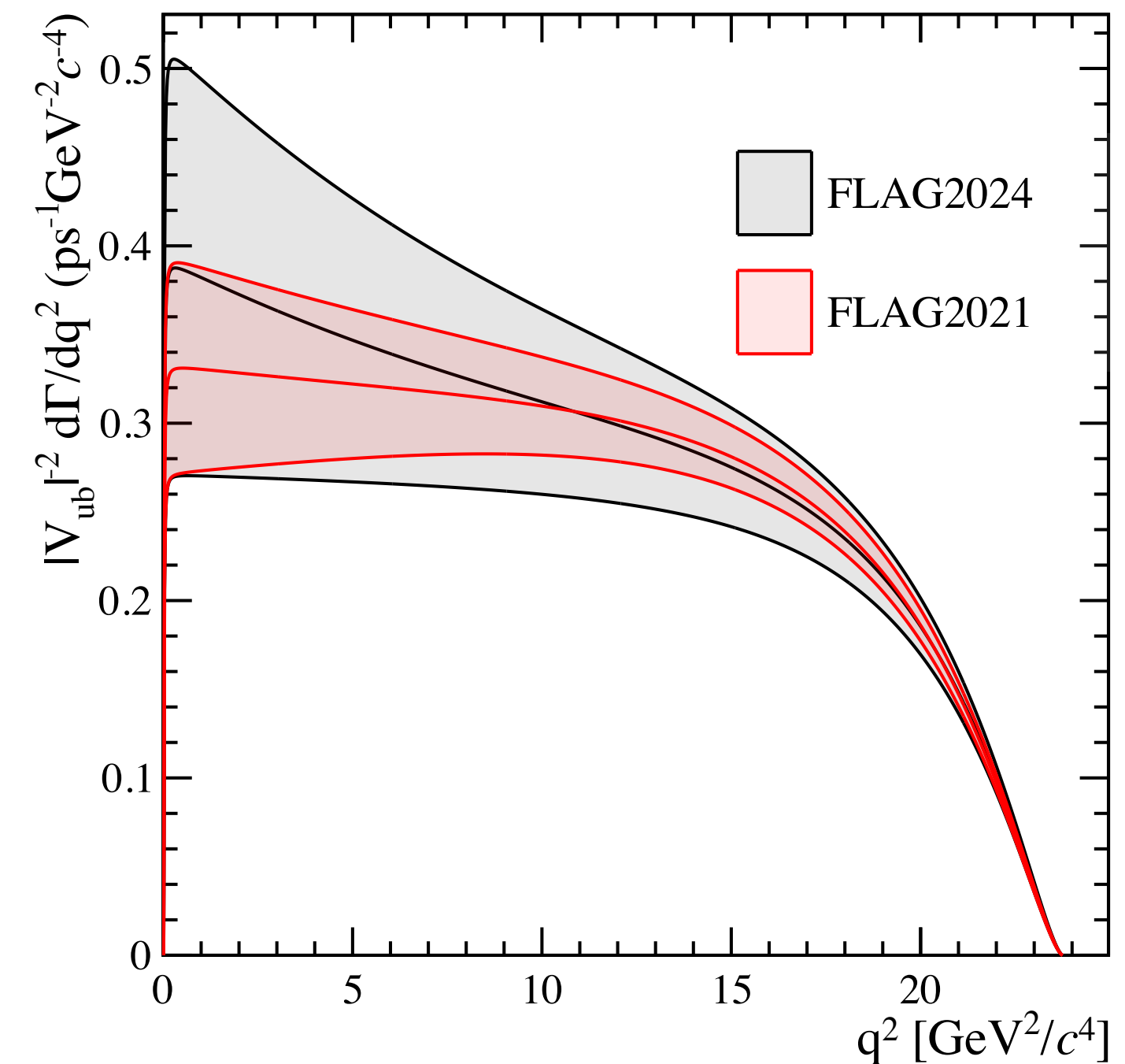
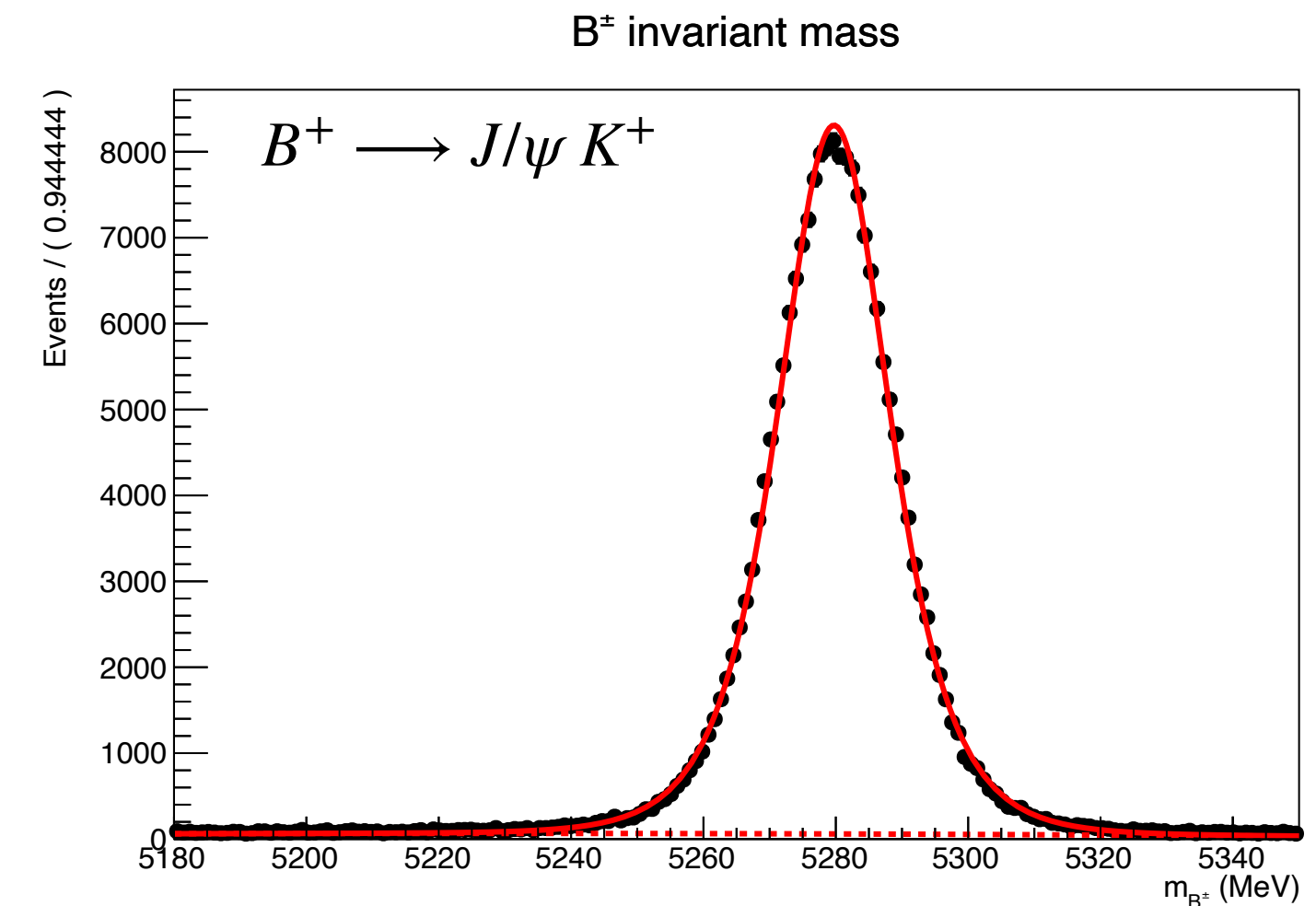
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- ◉ Extracting $|V_{ub}|/|V_{cb}|$ from $B_c^+ \rightarrow D^{*0} \mu \nu_\mu$ by normalizing to $B_c^+ \rightarrow J/\psi \mu \nu_\mu$
→ First CKM matrix element determined from B_c^+

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→ Expecting > 50 times higher signal yield wrt Belle [Phys. Rev. D 88, 032005 (2013)]
- ◉ Extracting $|V_{ub}|/|V_{cb}|$ from $B_c^+ \rightarrow D^{*0} \mu \nu_\mu$ by normalizing to $B_c^+ \rightarrow J/\psi \mu \nu_\mu$
→ First CKM matrix element determined from B_c^+
- ◉ Extracting $|V_{cb}|$ from $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu \nu_\mu$
→ First determination of $|V_{cb}|$ from a baryonic semileptonic decay

Status of $|V_{ub}|$
measurements from
 $B_s^0 \rightarrow K \mu \nu_\mu$ Run 2 dataset

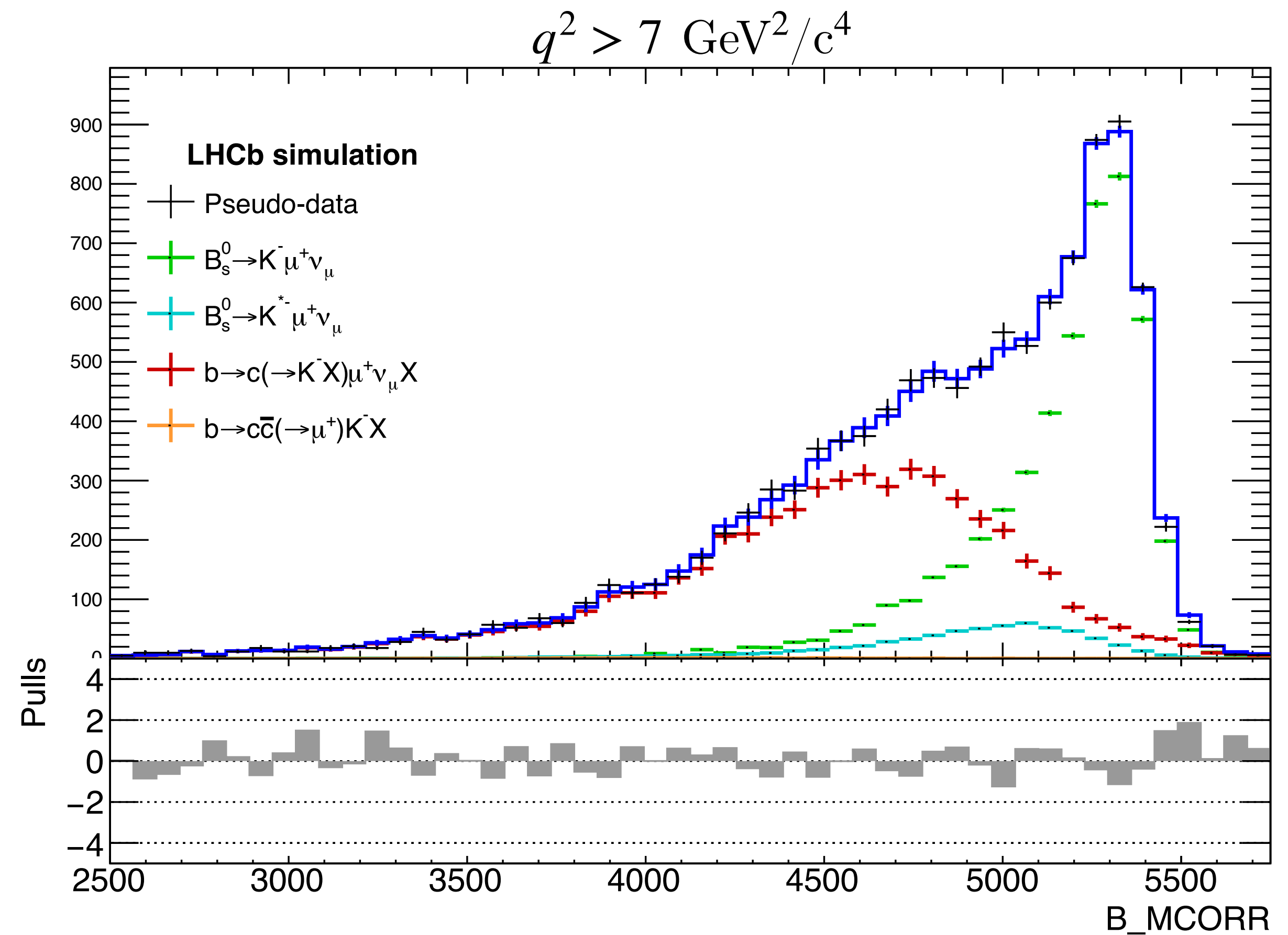
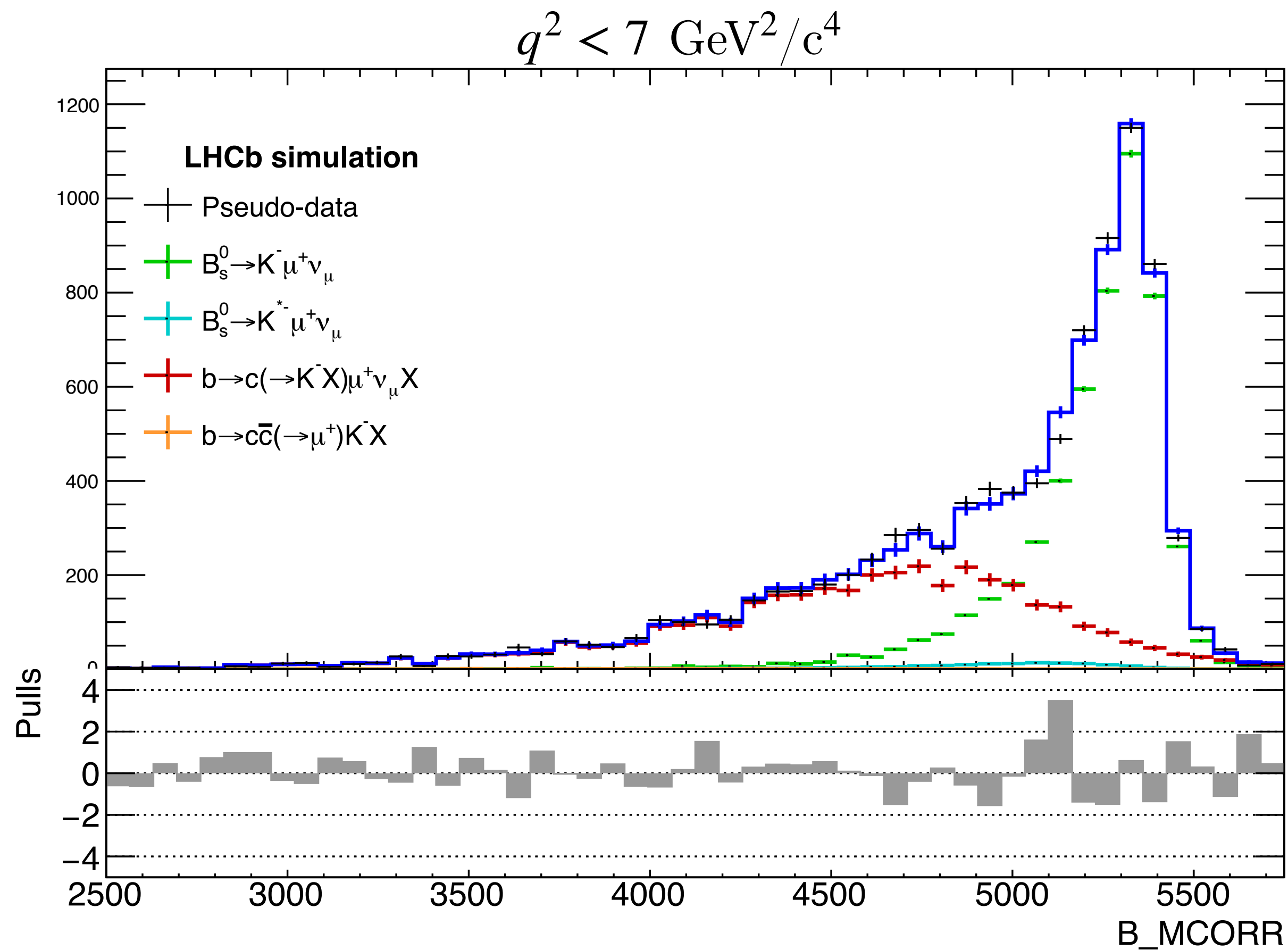
Analysis overview



- ◎ **Goal :** Determination of $\frac{d\Gamma}{dq^2}$ and $|V_{ub}|$ from $B_s^0 \rightarrow K^- \mu^+ \nu_\mu$
- ◎ **Normalisation channel :** $B^+ \rightarrow J/\psi K^+$
 - ⇒ Absolute measurement of $|V_{ub}|$ & high statistics
- ◎ **Dataset :** Full Run 2 $\sim 5.67 \text{ fb}^{-1}$ (factor 6 increase wrt Run1 analysis)
 - ⇒ Possibility to increase the number of q^2 bins from 2 (Run1) to $\sim 8 - 10$
- ◎ **FF parameterization :**
 - ⇒ FF schemes taken from FLAG 2024 and test others when available
- ◎ Based on Run1 strategy, working to optimise the various steps

Current M_{corr} plots ?

Toys



Conclusion

- $|V_{ub}|$ measurements @ LHCb are still in disagreement with the inclusive average, but not for the $|V_{cb}|$ results

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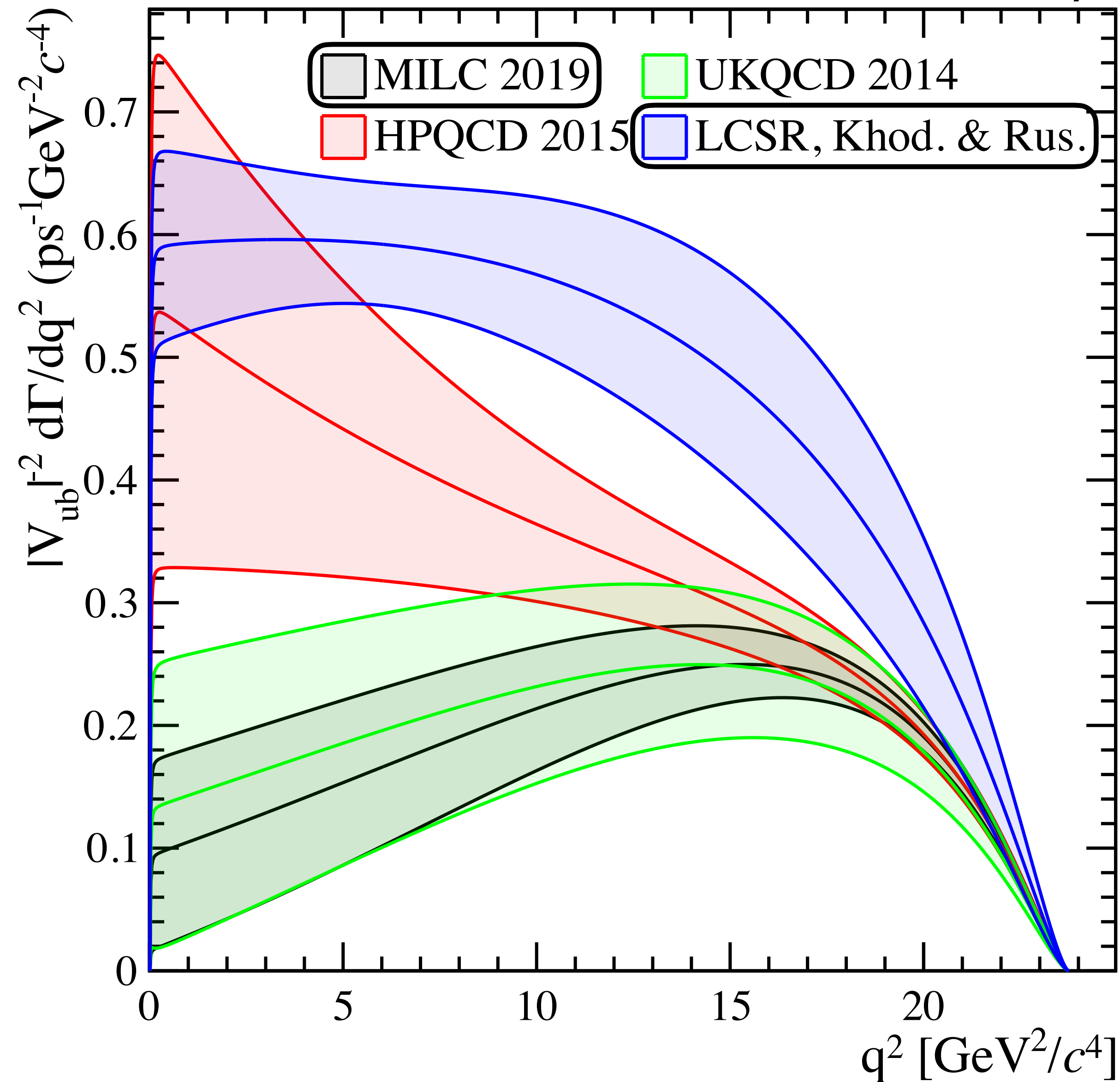
- ◉ $|V_{ub}|$ measurements @ LHCb are still in disagreement with the inclusive average, but not for the $|V_{cb}|$ results
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Conclusion

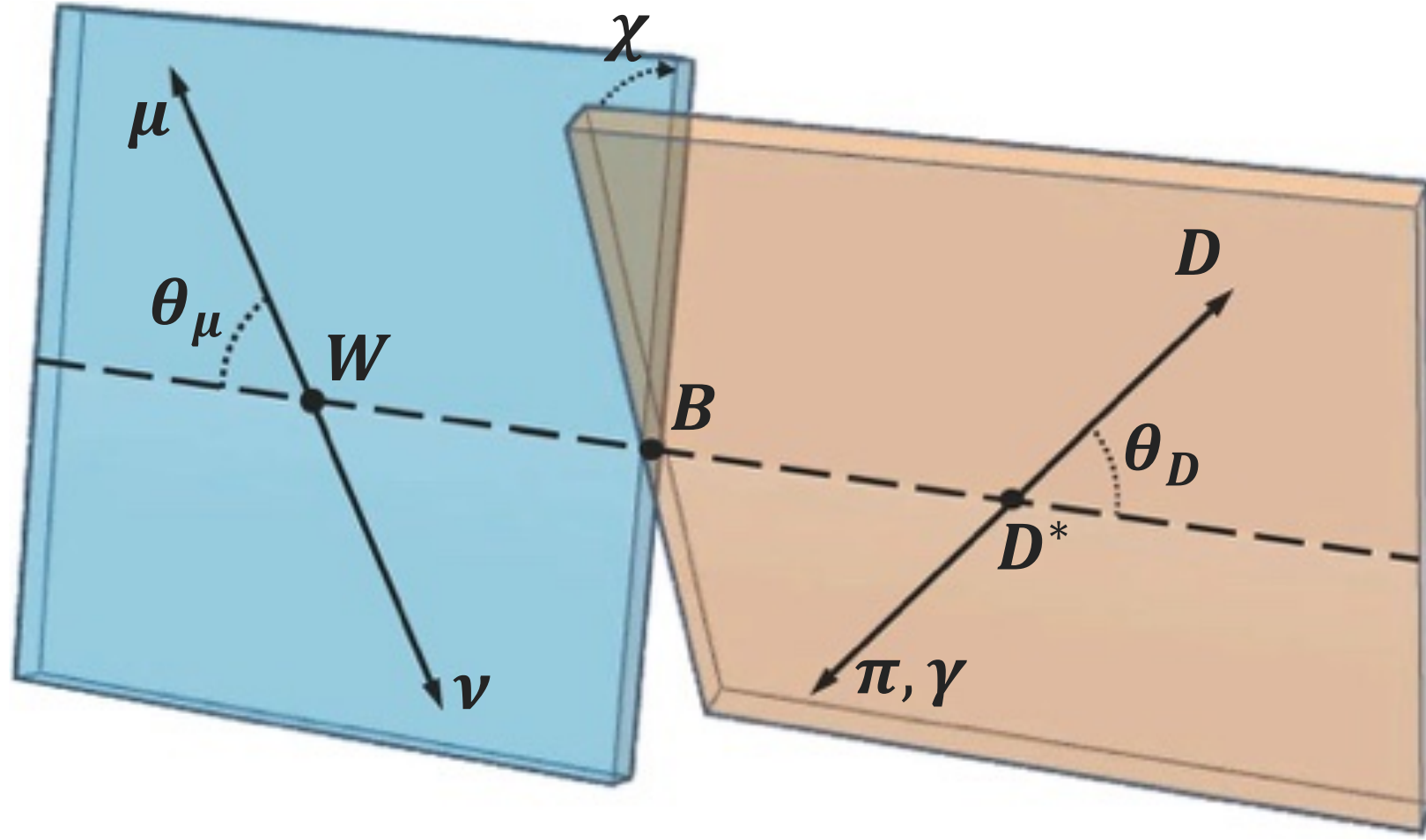
- ◉ $|V_{ub}|$ measurements @ LHCb are still in disagreement with the inclusive average, but not for the $|V_{cb}|$ results
- ◉ Measurement uncertainties mostly lead by the theory
- ◉ Need to be careful when using different FFs in a single analysis
- ◉ Waiting for Belle 2 new results !

BACK-UP

Form Factors used in $B_s^0 \rightarrow K \mu \nu_\mu$ from Run 1



$B \rightarrow D^{*-} \mu \nu_\mu$ analysis theory : CLN



$$\frac{d^4\Gamma(B \rightarrow D^* \mu \nu)}{dw d\cos\theta_\mu d\cos\theta_D d\chi} = \frac{3m_B^3 m_{D^*}^2 G_F^2}{16(4\pi)^4} \eta_{EW}^2 |V_{cb}|^2 |\mathcal{A}(w, \theta_\mu, \theta_D, \chi)|^2$$

$$|\mathcal{A}(w, \theta_\mu, \theta_D, \chi)|^2 = \sum_i^6 \mathcal{H}_i(w) k_i(\theta_\mu, \theta_D, \chi)$$

$$w = v_B \cdot v_{D^*} = (m_B^2 + m_{D^*}^2 - q^2)/(2m_B m_{D^*})$$

$$H_{\pm/0}(w) = 2 \frac{\sqrt{m_B m_{D^*}}}{m_B + m_{D^*}} (1 - r^2)(w + 1)(w^2 - 1)^{1/4} \times h_{A_1}(w) \tilde{H}_{\pm/0}(w),$$

with $r = m_{D^*}/m_B$ and

$$\tilde{H}_{\pm}(w) = \frac{\sqrt{1 - 2wr + r^2}}{1 - r} \left[1 \mp \sqrt{\frac{w - 1}{w + 1}} R_1(w) \right],$$

$$\tilde{H}_0(w) = 1 + \frac{(w - 1)(1 - R_2(w))}{1 - r}.$$

$$h_{A_1}(w) = h_{A_1}(1)[1 - 8\rho^2 z + (53\rho^2 - 15)z^2 - (231\rho^2 - 91)z^3],$$

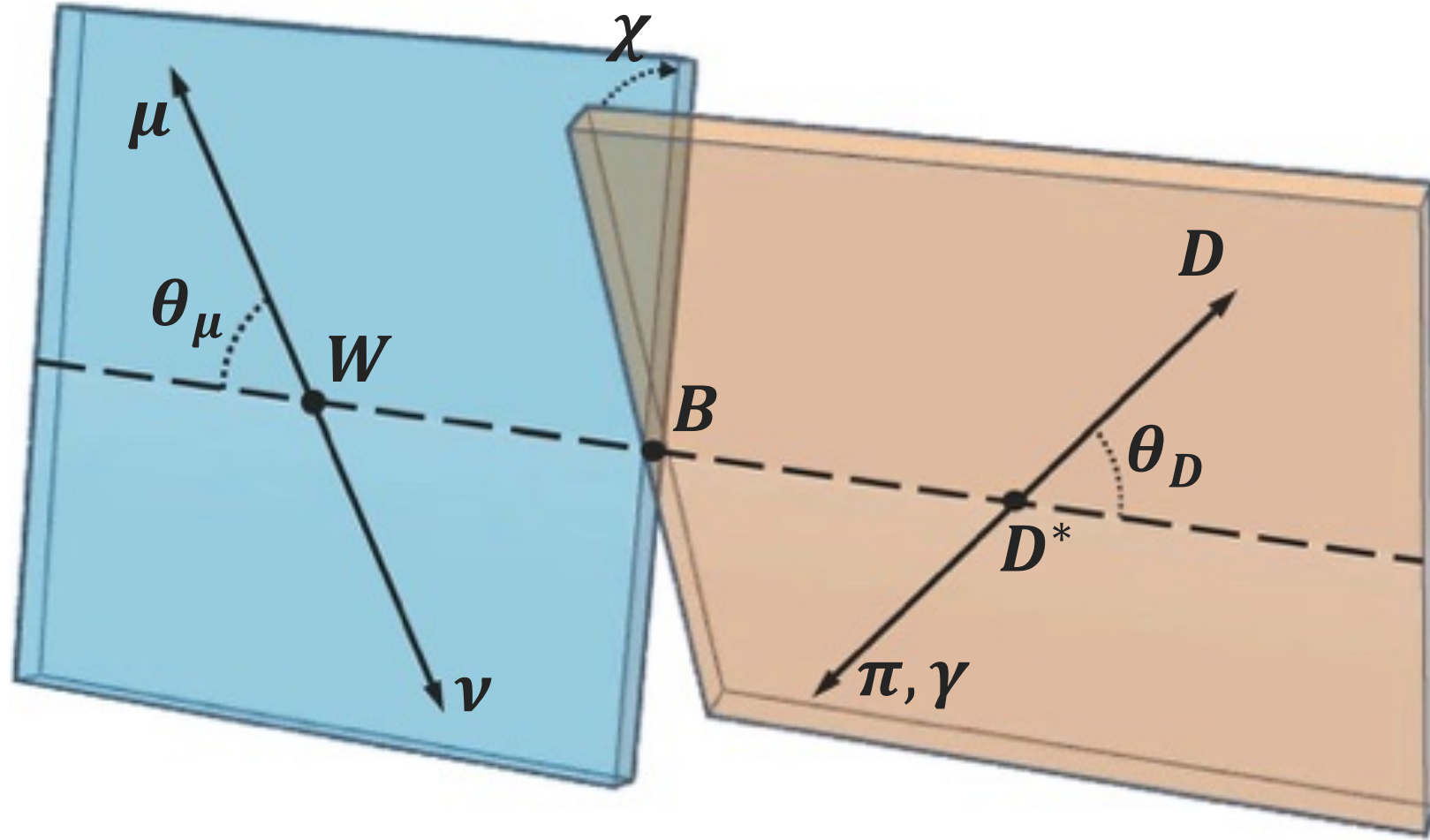
$$R_1(w) = R_1(1) - 0.12(w - 1) + 0.05(w - 1)^2$$

$$R_2(w) = R_2(1) - 0.11(w - 1) - 0.06(w - 1)^2$$

$$z \equiv \frac{\sqrt{w + 1} - \sqrt{2}}{\sqrt{w + 1} + \sqrt{2}}$$

i	$\mathcal{H}_i(w)$	$k_i(\theta_\mu, \theta_D, \chi)$	
		$D^* \rightarrow D\gamma$	$D^* \rightarrow D\pi^0$
1	H_+^2	$\frac{1}{2}(1 + \cos^2\theta_D)(1 - \cos\theta_\mu)^2$	$\sin^2\theta_D(1 - \cos\theta_\mu)^2$
2	H_-^2	$\frac{1}{2}(1 + \cos^2\theta_D)(1 + \cos\theta_\mu)^2$	$\sin^2\theta_D(1 + \cos\theta_\mu)^2$
3	H_0^2	$2\sin^2\theta_D \sin^2\theta_\mu$	$4\cos^2\theta_D \sin^2\theta_\mu$
4	$H_+ H_-$	$\sin^2\theta_D \sin^2\theta_\mu \cos 2\chi$	$-2\sin^2\theta_D \sin^2\theta_\mu \cos 2\chi$
5	$H_+ H_0$	$\sin 2\theta_D \sin\theta_\mu(1 - \cos\theta_\mu) \cos\chi$	$-2\sin 2\theta_D \sin\theta_\mu(1 - \cos\theta_\mu) \cos\chi$
6	$H_- H_0$	$-\sin 2\theta_D \sin\theta_\mu(1 + \cos\theta_\mu) \cos\chi$	$2\sin 2\theta_D \sin\theta_\mu(1 + \cos\theta_\mu) \cos\chi$

$B \rightarrow D^{*-} \mu \nu_\mu$ analysis theory : BGL



$$\frac{d^4\Gamma(B \rightarrow D^* \mu \nu)}{dw d\cos\theta_\mu d\cos\theta_D d\chi}$$

$$= \frac{3m_B^3 m_{D^*}^2 G_F^2}{16(4\pi)^4} \eta_{EW}^2 |V_{cb}|^2 |\mathcal{A}(w, \theta_\mu, \theta_D, \chi)|^2$$

$$|\mathcal{A}(w, \theta_\mu, \theta_D, \chi)|^2 = \sum_i^6 \mathcal{H}_i(w) k_i(\theta_\mu, \theta_D, \chi)$$

$$w = v_B \cdot v_{D^*} = (m_B^2 + m_{D^*}^2 - q^2)/(2m_B m_{D^*})$$

$$H_{\pm/0}(w) = 2 \frac{\sqrt{m_B m_{D^*}}}{m_B + m_{D^*}} (1 - r^2)(w + 1)(w^2 - 1)^{1/4}$$

$$\times h_{A_1}(w) \tilde{H}_{\pm/0}(w),$$

with $r = m_{D^*}/m_B$ and

$$\tilde{H}_\pm(w) = \frac{\sqrt{1 - 2wr + r^2}}{1 - r} \left[1 \mp \sqrt{\frac{w - 1}{w + 1}} R_1(w) \right],$$

$$\tilde{H}_0(w) = 1 + \frac{(w - 1)(1 - R_2(w))}{1 - r}.$$

$$h_{A_1}(w) = \frac{f(w)}{\sqrt{m_B m_{D^*}}(1 + w)},$$

$$R_1(w) = (w + 1)m_B m_{D^*} \frac{g(w)}{f(w)},$$

$$R_2(w) = \frac{w - r}{w - 1} - \frac{\mathcal{F}_1(w)}{m_B(w - 1)f(w)}$$

$$f(z) = \frac{1}{P_{1^+}(z)\phi_f(z)} \sum_{n=0}^{\infty} b_n z^n$$

$$g(z) = \frac{1}{P_{1^-}(z)\phi_g(z)} \sum_{n=0}^{\infty} a_n z^n$$

$$\mathcal{F}_1(z) = \frac{1}{P_{1^+}(z)\phi_{\mathcal{F}_1}(z)} \sum_{n=0}^{\infty} c_n z^n$$

$$P_{1^\pm}(z) = C_{1^\pm} \prod_{k=1}^{\text{poles}} \frac{z - z_k}{1 - z z_k}$$

$$t_\pm = (m_B \pm m_{D^*})^2 \quad n_I = 2.6$$

$$\sum_{n=0}^{\infty} a_n^2 \leq 1, \quad \sum_{n=0}^{\infty} (b_n^2 + c_n^2) \leq 1$$

$$b_0 = 2\sqrt{m_B m_{D^*}} P_{1^+}(0) \phi_f(0) h_{A_1}$$

$$c_0 = (m_B - m_{D^*}) \frac{\phi_{\mathcal{F}_1}(0)}{\phi_f(0)} b_0$$

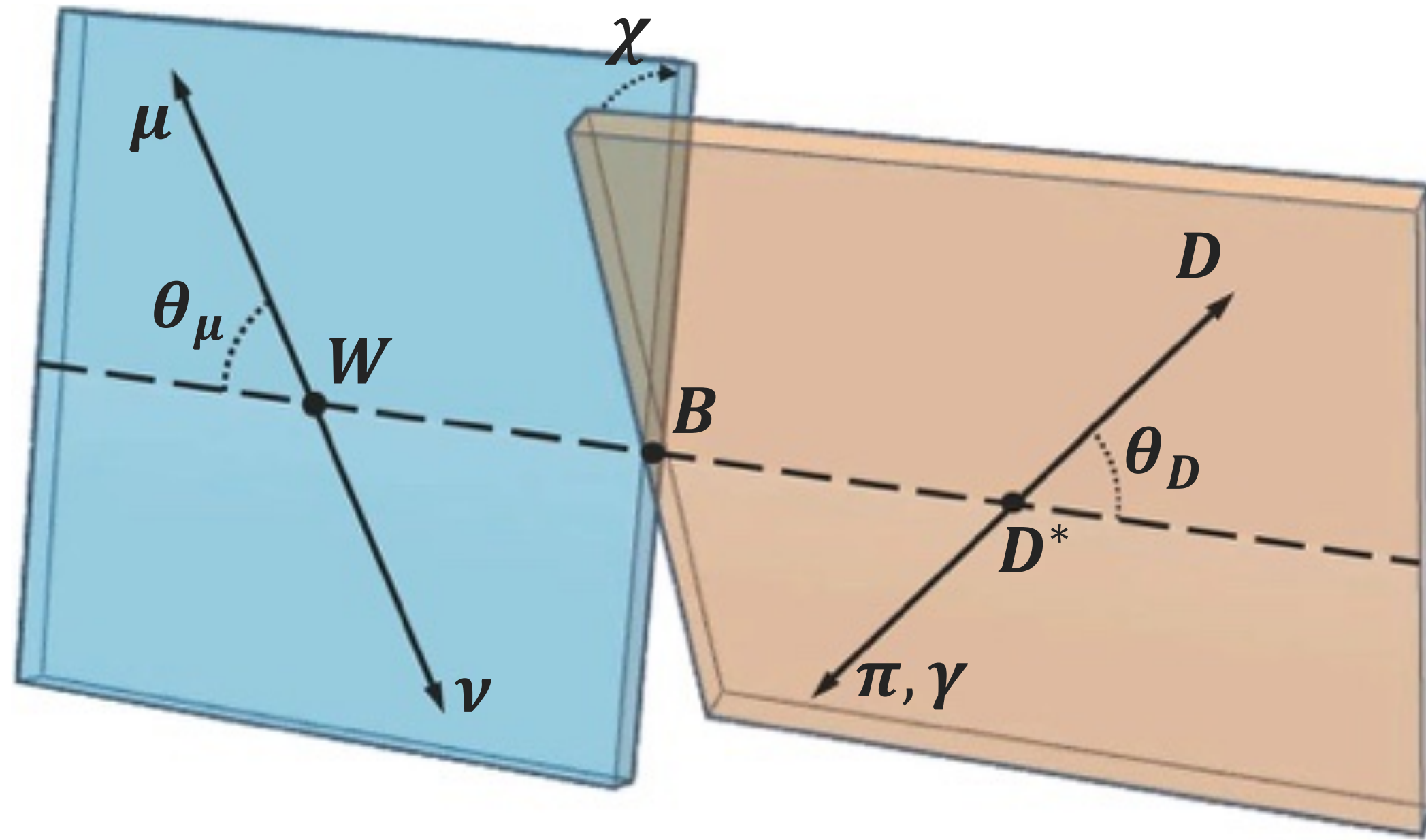
$$\phi_f(z) = \frac{4r}{m_B^2} \sqrt{\frac{n_I}{3\pi\tilde{\chi}_{1^+}(0)}} \frac{(1+z)\sqrt{(1-z)^3}}{[(1+r)(1-z) + 2\sqrt{r}(1+z)]^4},$$

$$\phi_g(z) = 16r^2 \sqrt{\frac{n_I}{3\pi\tilde{\chi}_{1^-}(0)}} \times \frac{(1+z)^2}{\sqrt{(1-z)[(1+r)(1-z) + 2\sqrt{r}(1+z)]^4}},$$

$$\phi_{\mathcal{F}_1}(z) = \frac{4r}{m_B^3} \sqrt{\frac{n_I}{6\pi\tilde{\chi}_{1^+}(0)}} \frac{(1+z)\sqrt{(1-z)^5}}{[(1+r)(1-z) + 2\sqrt{r}(1+z)]^5}$$

$$z_k = \frac{\sqrt{t_+ - m_k^2} - \sqrt{t_+ - t_-}}{\sqrt{t_+ - m_k^2} + \sqrt{t_+ - t_-}}$$

$B \rightarrow D^- \mu \nu_\mu$ analysis theory : CLN

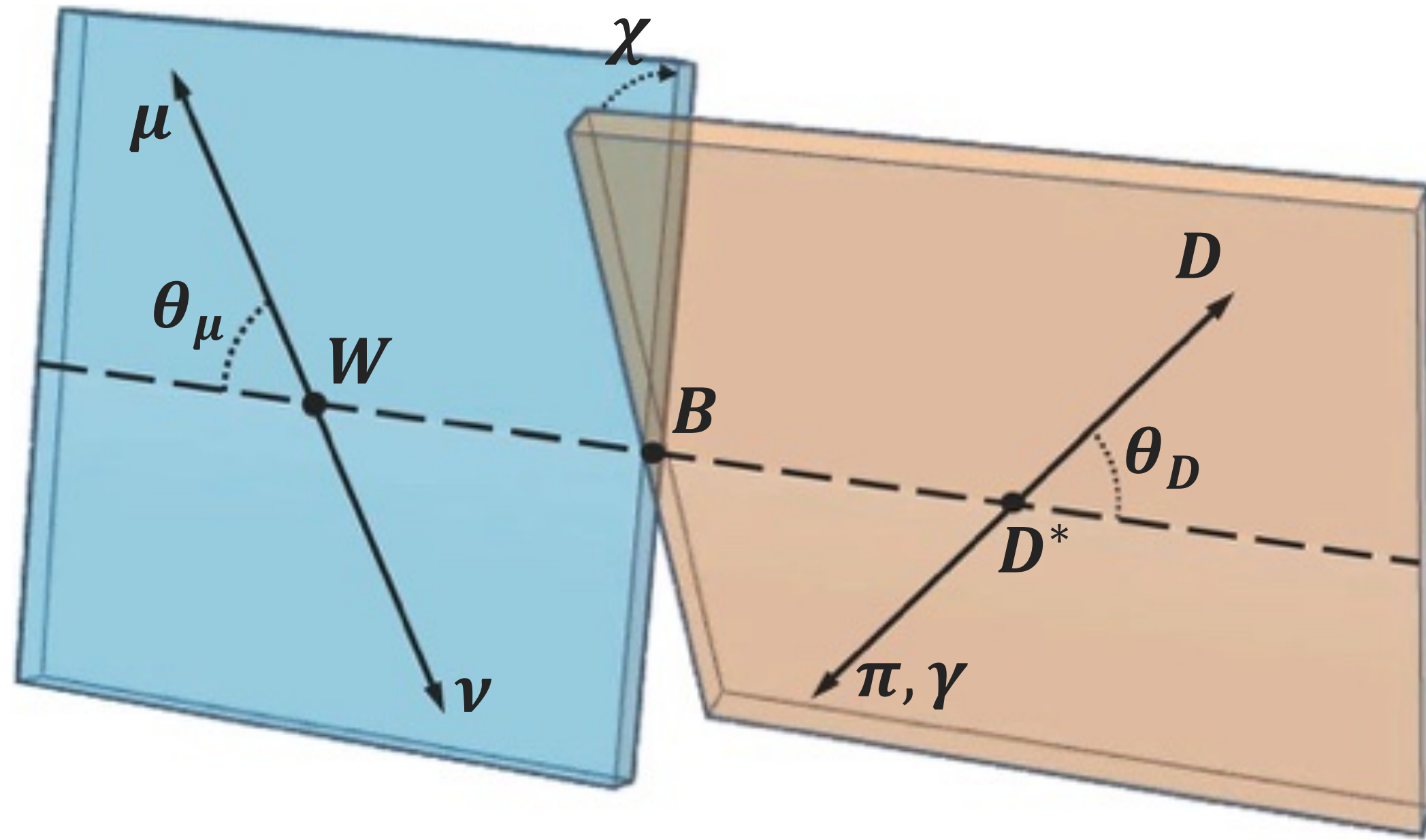


$$\frac{d\Gamma(B \rightarrow D\mu\nu)}{dw} = \frac{G_F^2 m_D^3}{48\pi^3} (m_B + m_D)^2 \eta_{EW}^2 \times |V_{cb}|^2 (w^2 - 1)^{3/2} |\mathcal{G}(w)|^2$$

$$\mathcal{G}(z) = \mathcal{G}(0)[1 - 8\rho^2 z + (51\rho^2 - 10)z^2 - (252\rho^2 - 84)z^3]$$

$$w = v_B \cdot v_{D^*} = (m_B^2 + m_{D^*}^2 - q^2)/(2m_B m_{D^*})$$

$B \rightarrow D^- \mu \nu_\mu$ analysis theory : BGL



$$\frac{d\Gamma(B \rightarrow D\mu\nu)}{dw} = \frac{G_F^2 m_D^3}{48\pi^3} (m_B + m_D)^2 \eta_{EW}^2 \times |V_{cb}|^2 (w^2 - 1)^{3/2} |\mathcal{G}(w)|^2$$

$$|\mathcal{G}(z)|^2 = \frac{4r}{(1+r)^2} |f_+(z)|^2 \quad r = m_D/m_B$$

$$f_+(z) = \frac{1}{P_{1-}(z)\phi(z)} \sum_{n=0}^{\infty} d_n z^n \quad \sum_{n=0}^{\infty} d_n^2 \leq 1$$

$$d_0 = \frac{1+r}{2\sqrt{r}} \mathcal{G}(0) P_{1-}(0) \phi(0)$$