

# Solar disk gamma-rays emission via synthetic magnetic field from photosphere to low corona



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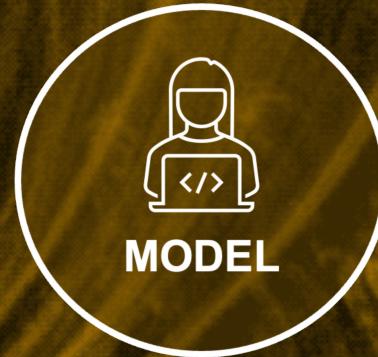
*Atelier "Physique cinétique des plasmas astrophysiques"*



Project supported by  
Grant 80NSSC22K0040



INTRO



MODEL



RESULTS



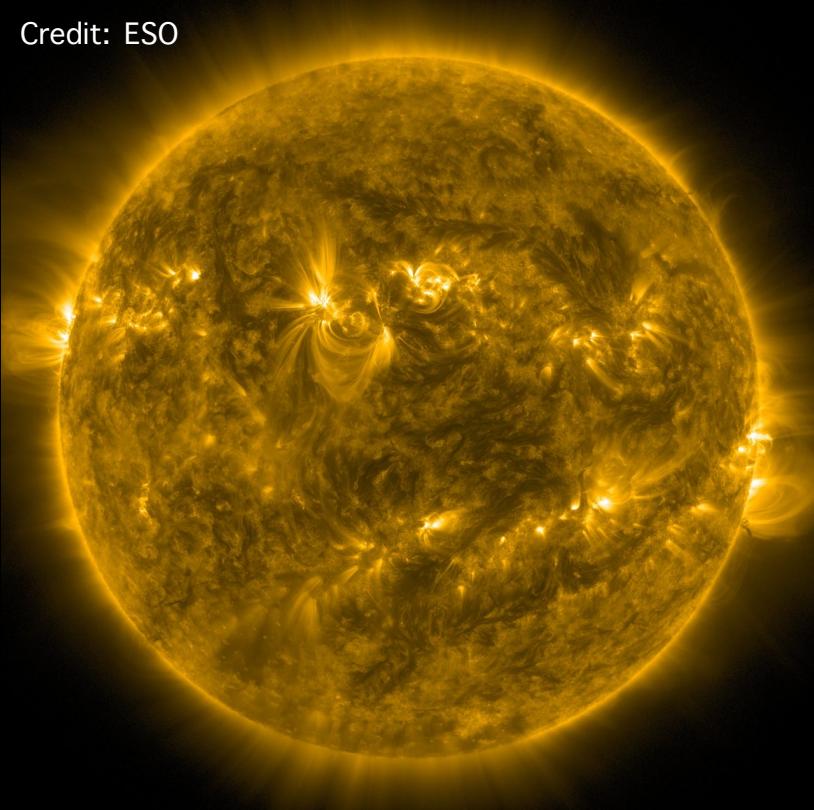
INTRO



# INTRO

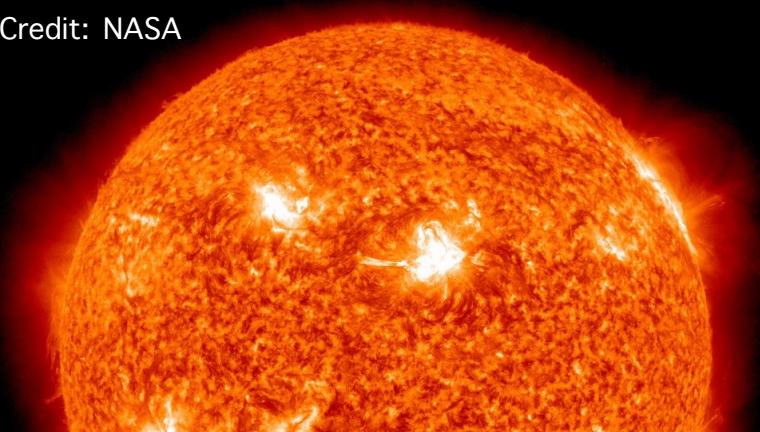
# SOLAR EMISSION MECHANISMS

Credit: ESO



- **Inverse-Compton scattering:** cosmic-ray electrons interacting with solar photons (solar **halo**)
- **Bombardment by hadronic cosmic rays** (mostly protons) interacting with solar gas (solar **disk**)
- **Solar flares and coronal mass ejections**

Credit: NASA





# GAMMA-RAYS PRODUCTION



Credit: Scott Wiessinger

 INTRO

# THE EXISTING THEORETICAL MODEL

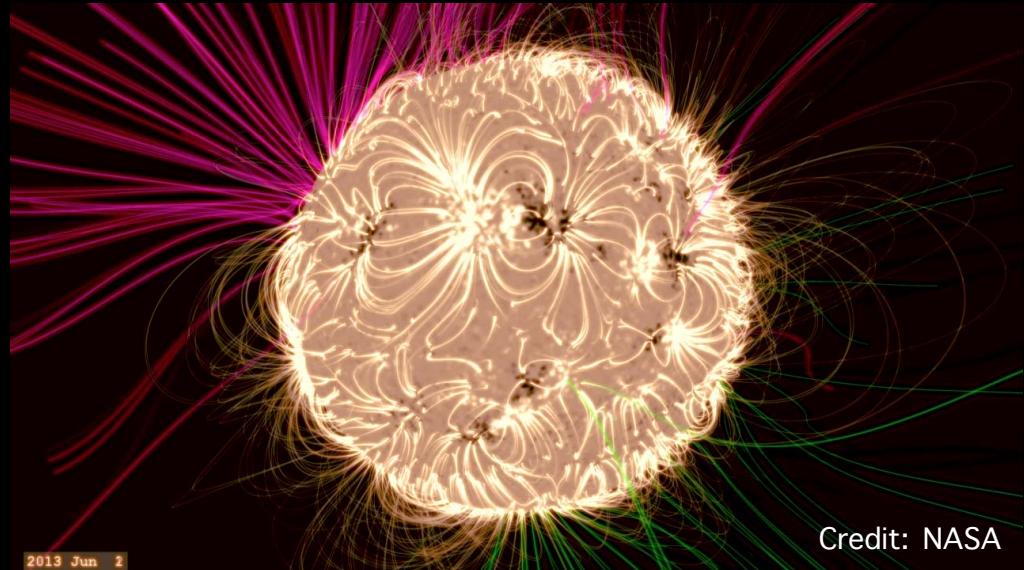
**Seckel, Stanev, and Gaisser model (1991)**

Magnetic flux tubes can reverse incoming protons deep within the solar atmosphere, where they have an appreciable probability of producing outgoing  $\gamma$ -rays.



$\gamma$ -ray flux is greatly enhanced!

Magnetic fields are crucial!

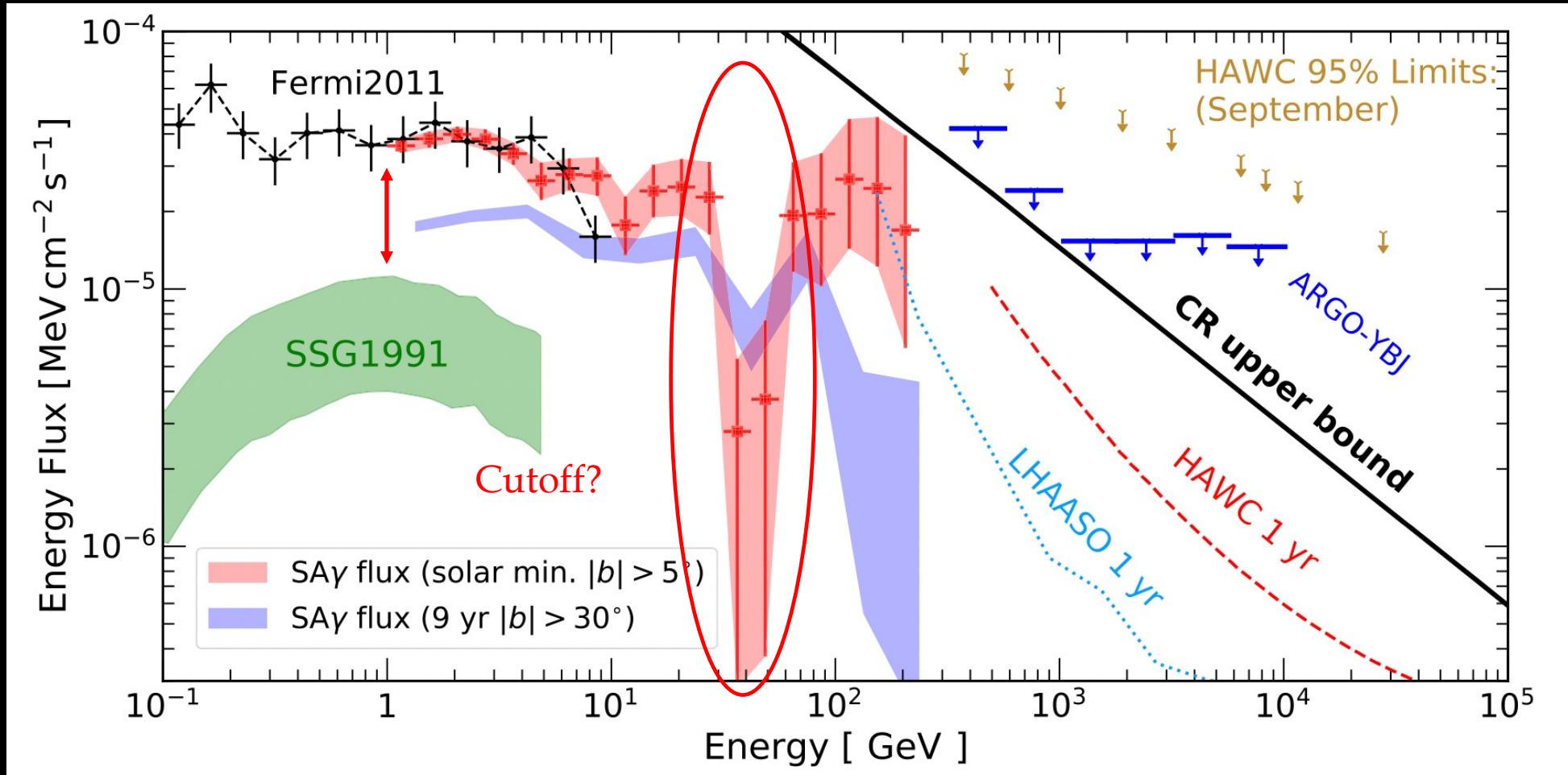


*“Current technology is improving to the point that such a flux of GeV  $\gamma$ -rays should be detectable by the EGRET instrument of the Gamma Ray Observatory (GRO)”*

- Fermi Gamma-ray Space Telescope
- High-Altitude Water Cherenkov Gamma-ray Observatory (HAWC)

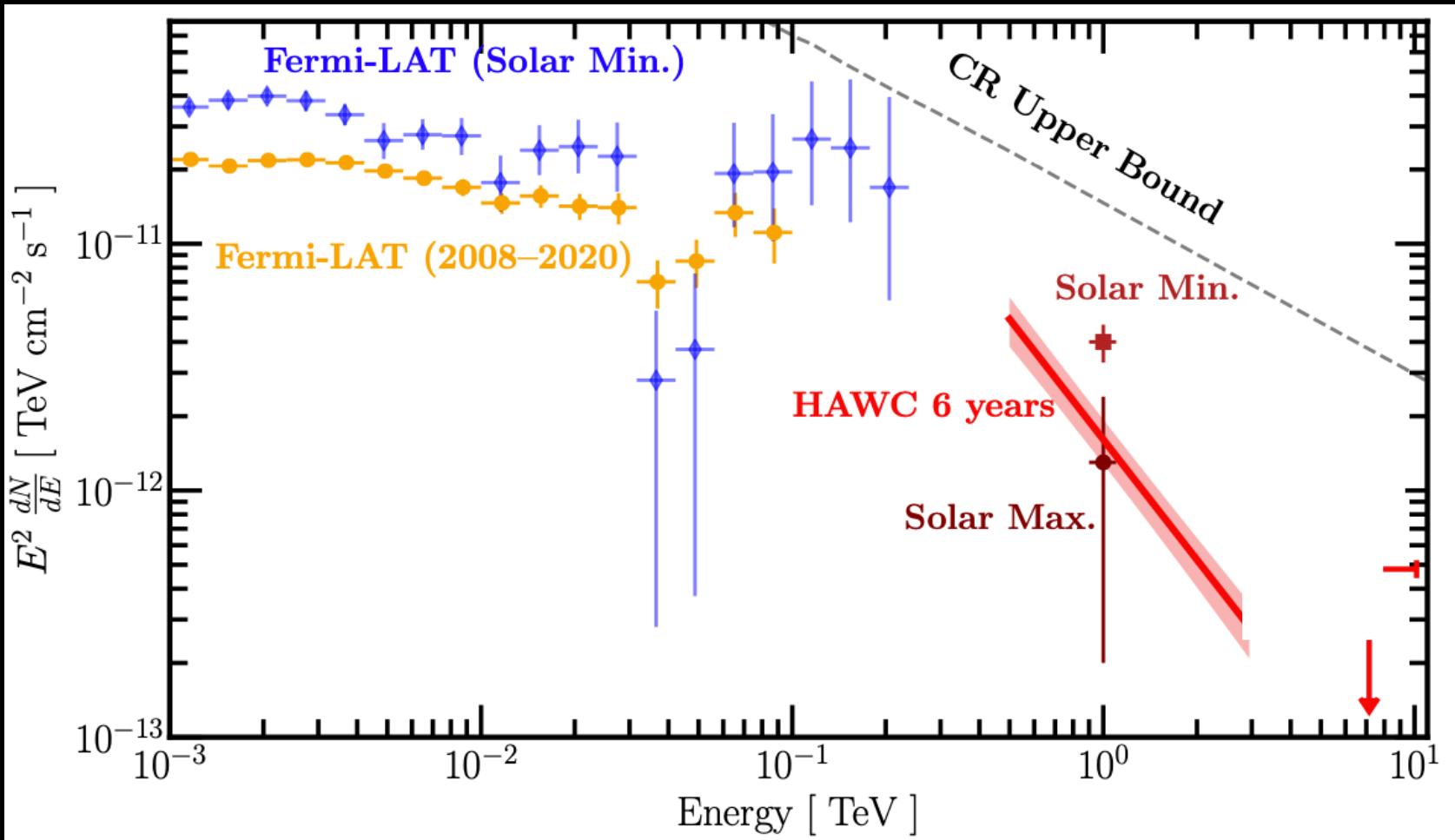
 INTRO

## MODEL VS OBSERVATIONS



 INTRO

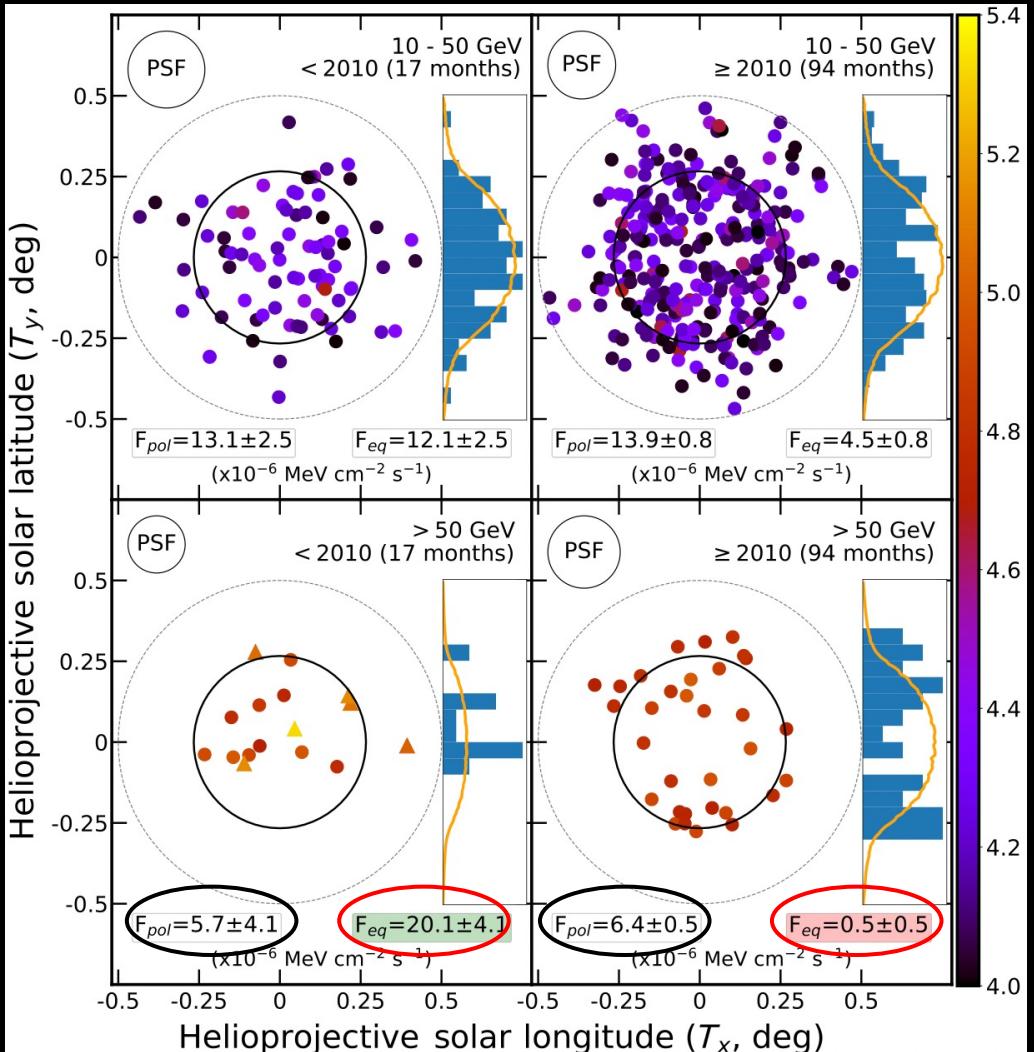
## MODEL VS OBSERVATIONS



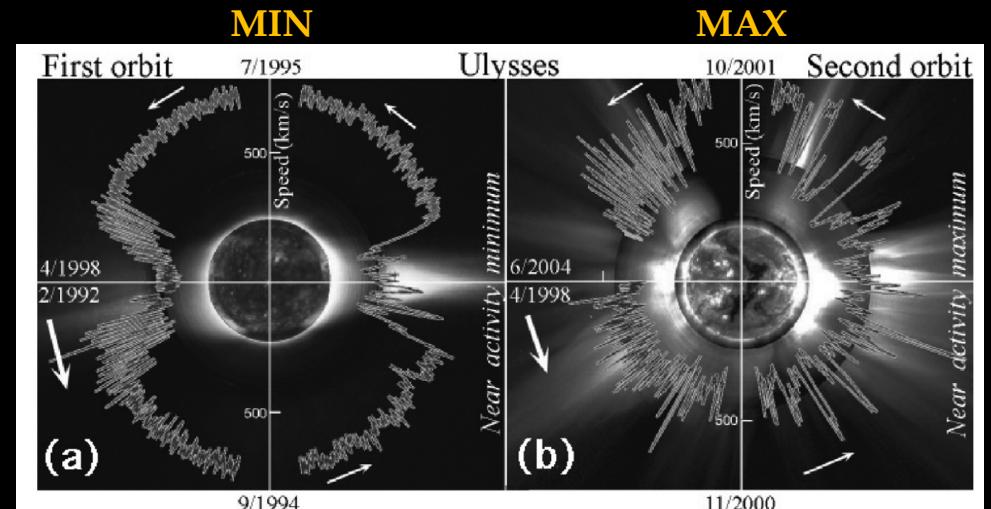


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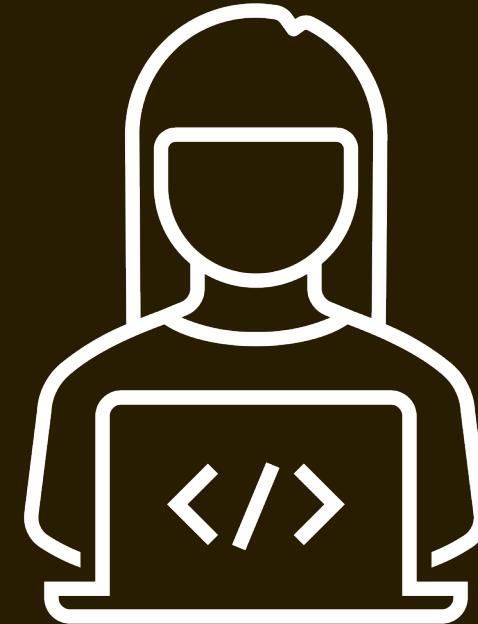
# SOLAR CYCLE DEPENDENCE



Credits: Linden et al. 2018



Credits: Priest 2012

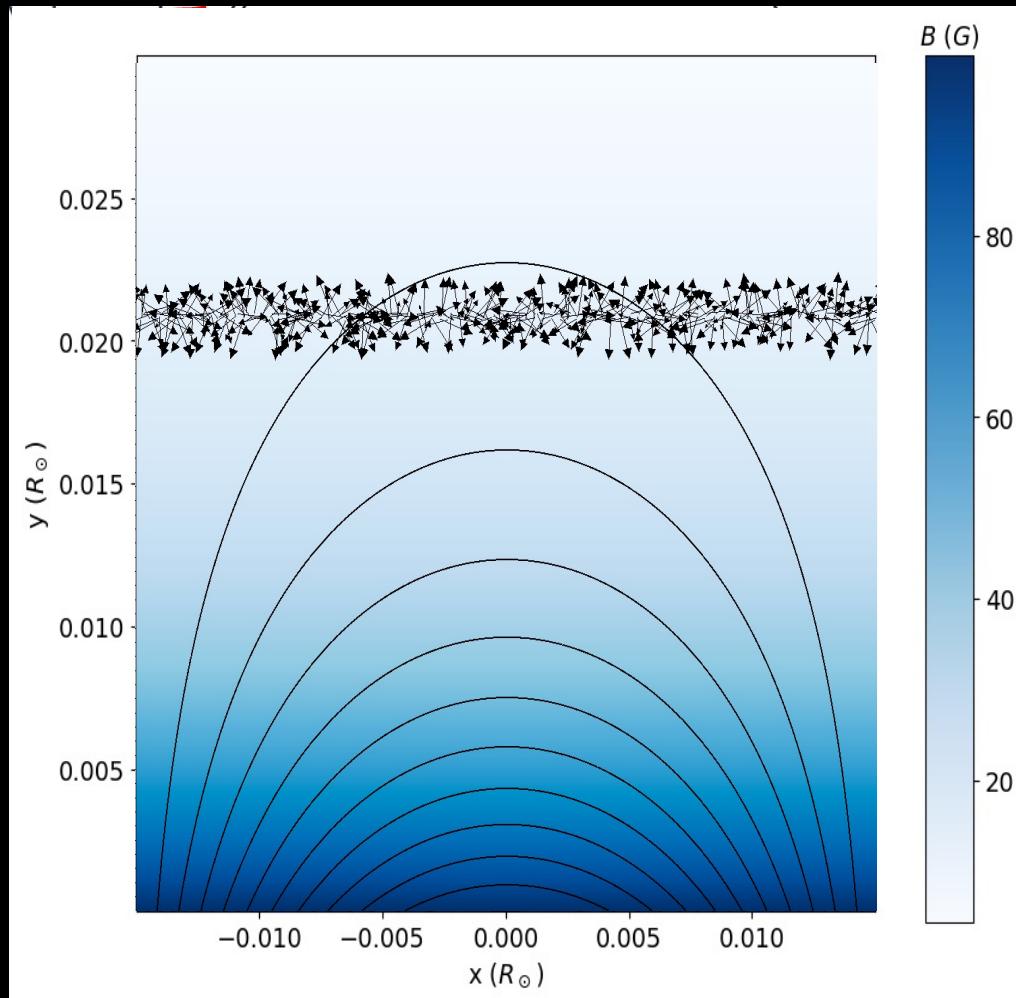


MODEL



# MODEL

## PREVIOUS WORK



*“Role of Magnetic Arcades in Explaining the Puzzle of the Gamma-Ray Emission from the Solar Disk”*

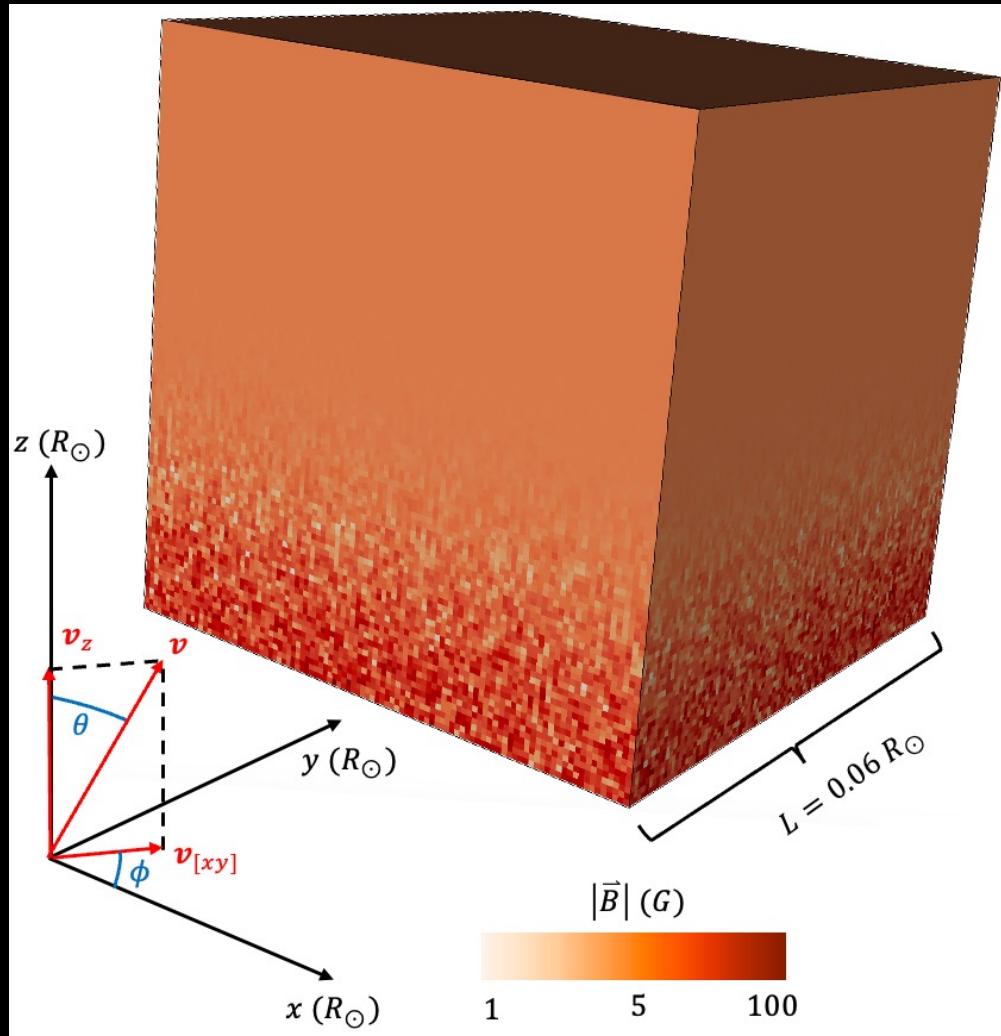


SCAN ME



# MODEL

## SYNTHETIC MAGNETIC FIELD



$$\mathbf{B}(x, y, z) = B_0 \hat{\mathbf{k}} + \nabla \times [f(z) \mathbf{S}(x, y, z)]$$

$$B_0 = 5 \text{ Gauss} \quad f(z) = e^{-\left(\frac{z}{\Lambda}\right)^2} \quad \Lambda \sim 10^{-2} R_\odot$$

Credits: Giacalone 2021

$$\mathbf{S}(x, y, z) = \sum_{n=1}^{N_m} \frac{A_n}{k_n} \left( \sin \alpha_n \widehat{x'_n} + i \cos \alpha_n \widehat{y'_n} \right) e^{ik_n z'_n + i\beta_n}$$

$$\delta \mathbf{B}(x, y, z) = \nabla \times \mathbf{S}(x, y, z)$$

Credits: Giacalone & Jokipii 1999

$$\delta \mathbf{B}(x, y, z) = \sum_{n=1}^{N_m} A_n \left( \cos \alpha_n \widehat{x'_n} + i \sin \alpha_n \widehat{y'_n} \right) e^{ik_n z'_n + i\beta_n}$$

Credits: Puzzoni et al. 2025 (submitted to ApJL)



# MODEL

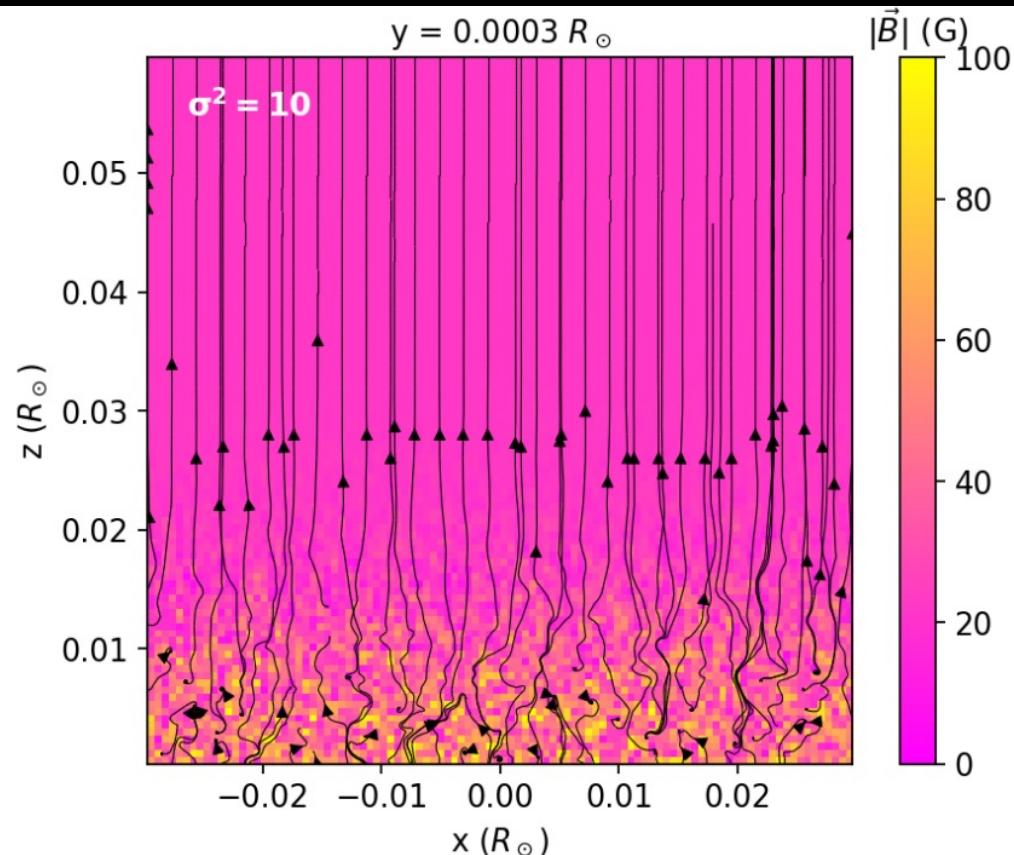
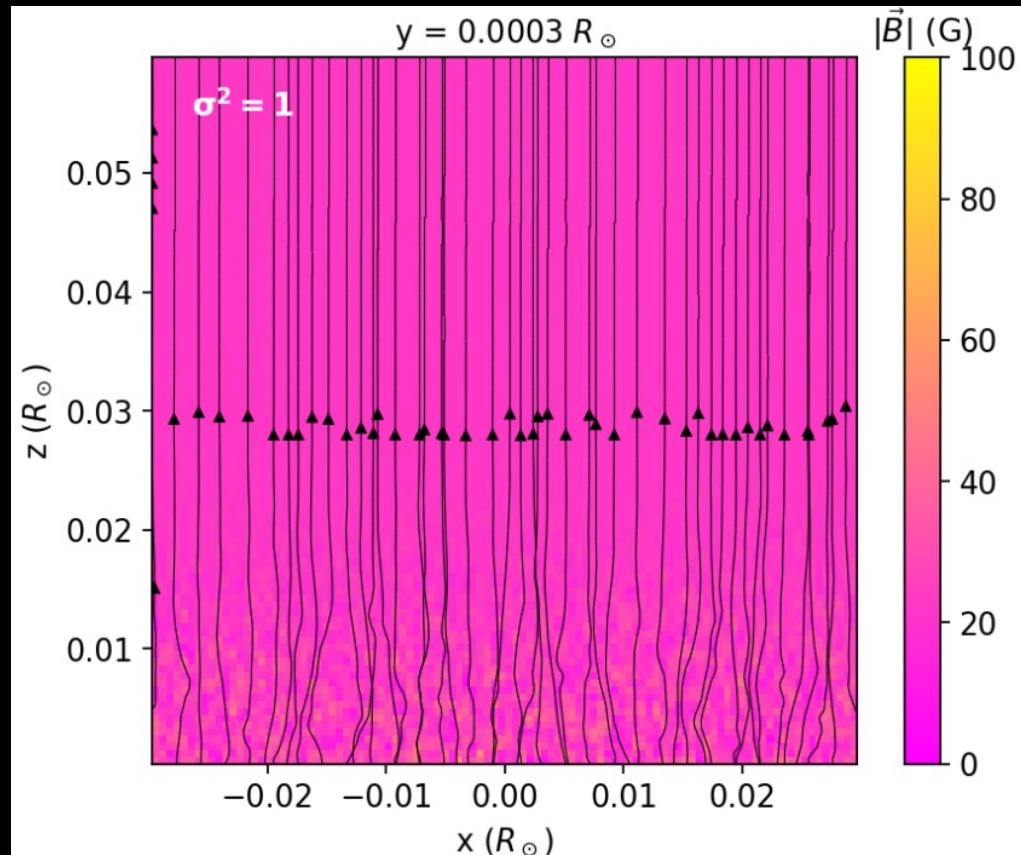
$$L_s = 0.003 R_\odot \sim 2 \text{ Mm}$$

$$L_{\min} = 10^{-3} R_\odot$$

$$L_{\max} = 0.06 R_\odot$$

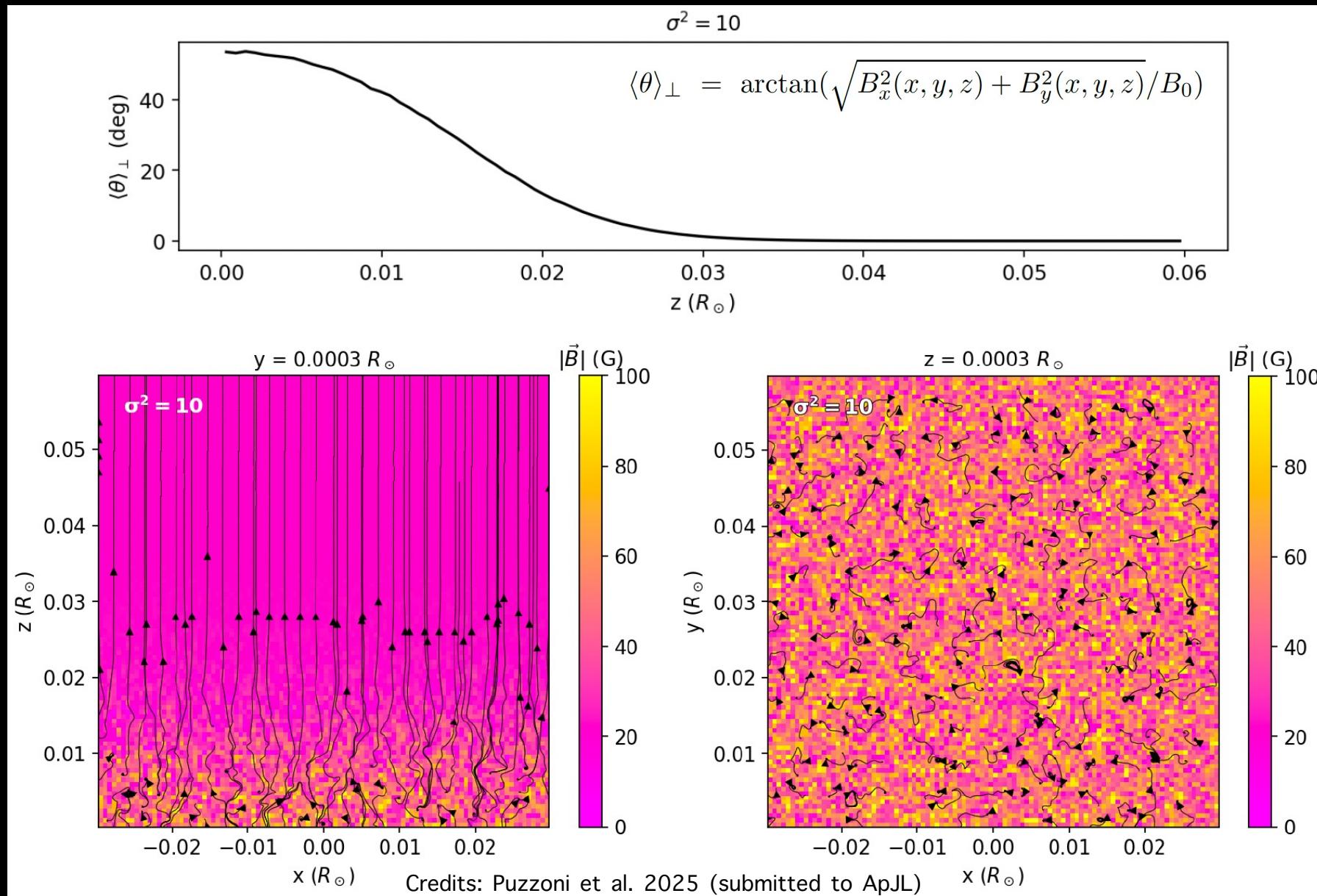
$$\sigma^2 = \frac{\sigma'^2}{B_0^2} = 0.1, 1.0, 10$$

Credits: Puzzoni et al. 2025 (submitted to ApJL)





# MODEL

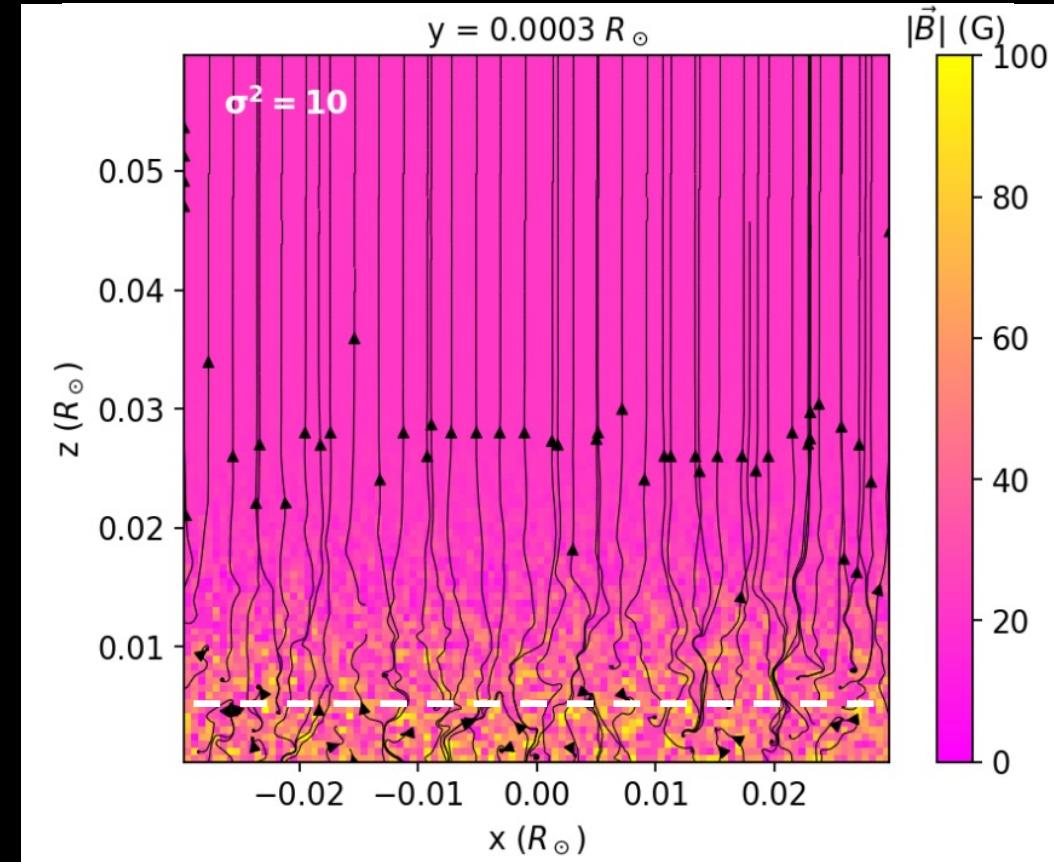
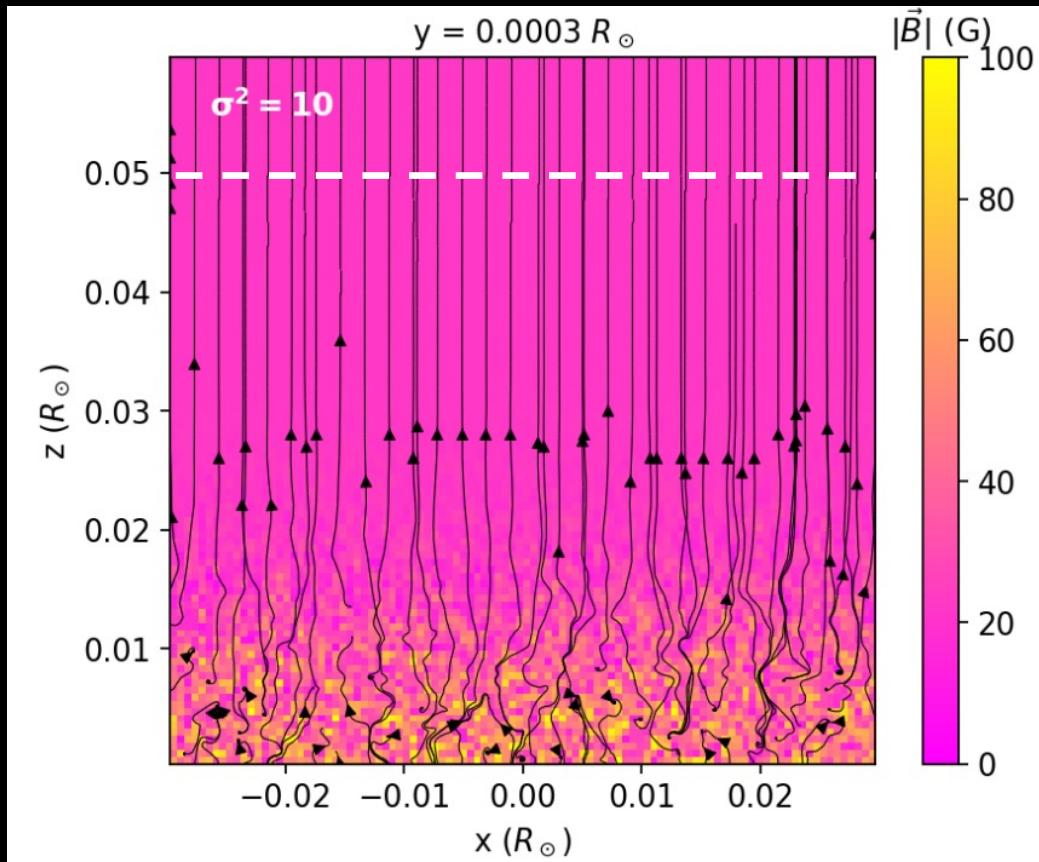


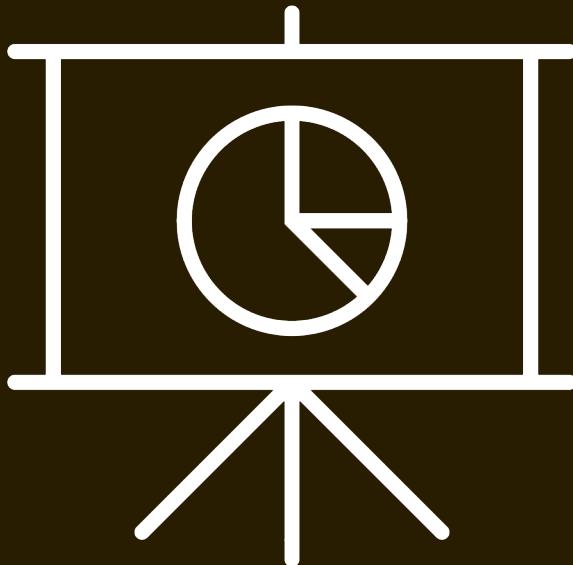


# MODEL

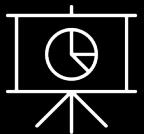
## TEST-PARTICLES INITIALIZATION

- Initialized with an **isotropic velocity distribution**
- Magnetic field configuration **reshuffled** every **100** particles





**RESULTS**

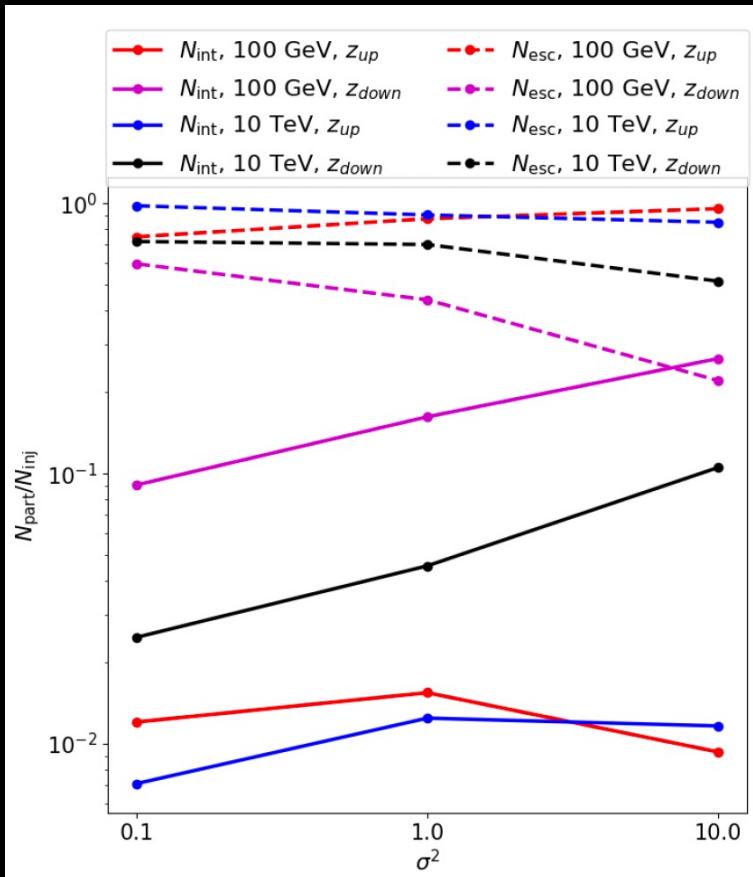


# RESULTS

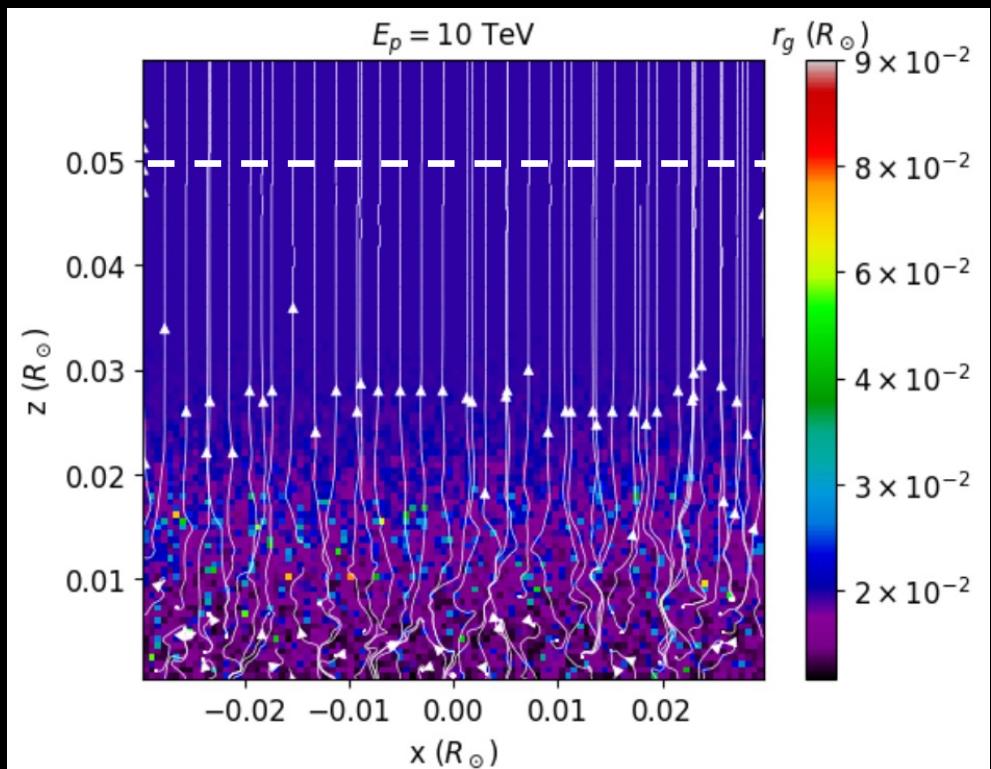
## INTERACTING GCR PROTONS

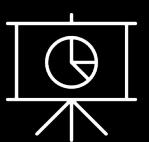
$$\frac{\Delta t}{t_{int}} = \int_{t_i}^{t_f} \frac{dt}{t_{int}(z)} = \sigma_{pp} v \int_{t_i}^{t_f} n[z(t)] dt > 1$$

$v_z > 0$  Focus: protons escaping from the Sun as they can produce gamma-rays observed at Earth



Credits: Puzzoni et al. 2025 (submitted to ApJL)

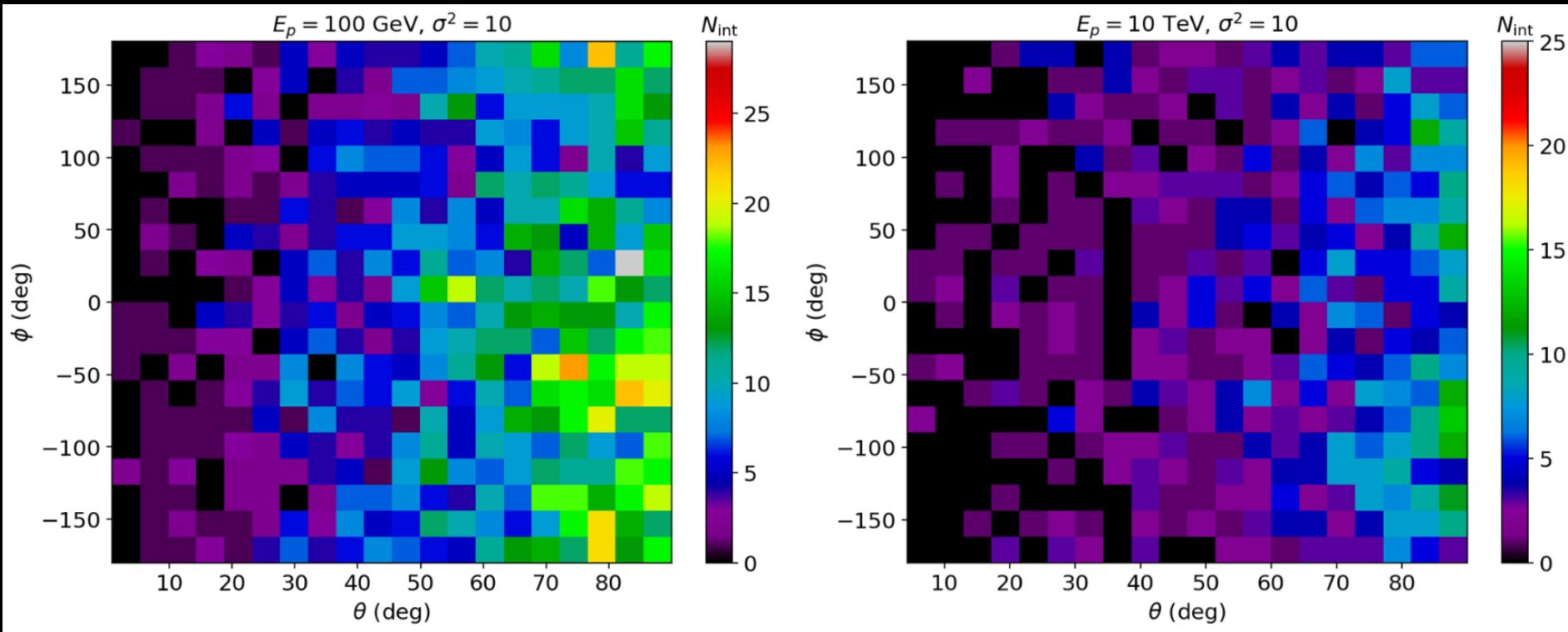
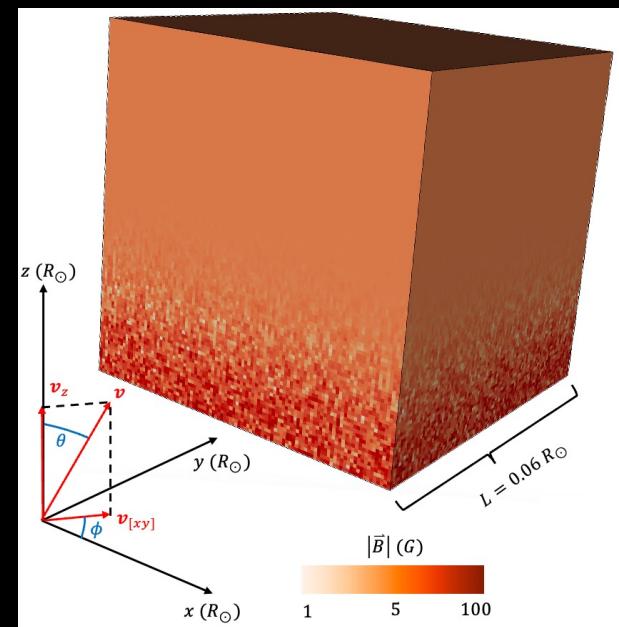




# RESULTS

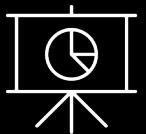
## $\gamma$ -RAYS EMISSION DIRECTION

Credits: Puzzoni et al. 2025 (submitted to ApJL)

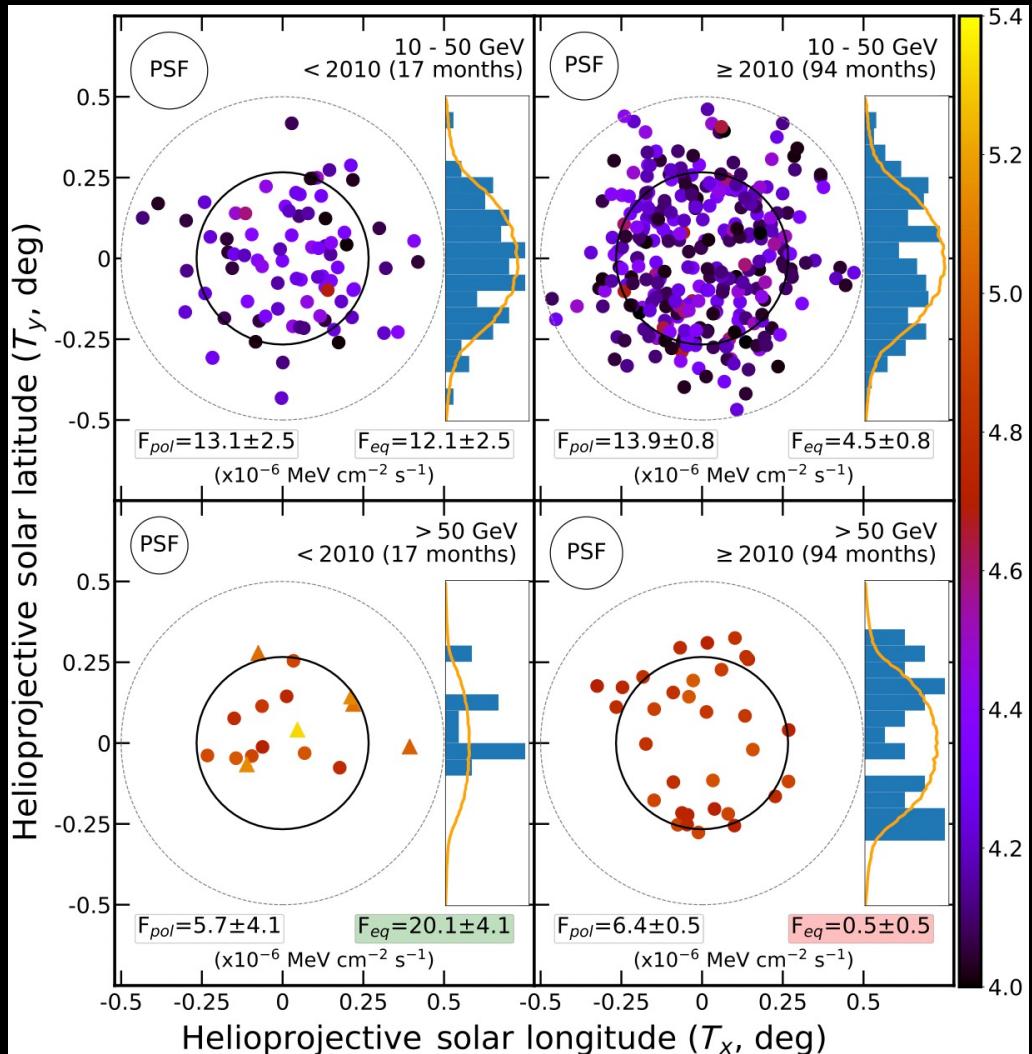


- $\theta$  range broadens to  $\approx 40^\circ$
- $70^\circ \lesssim \theta \lesssim 90^\circ$

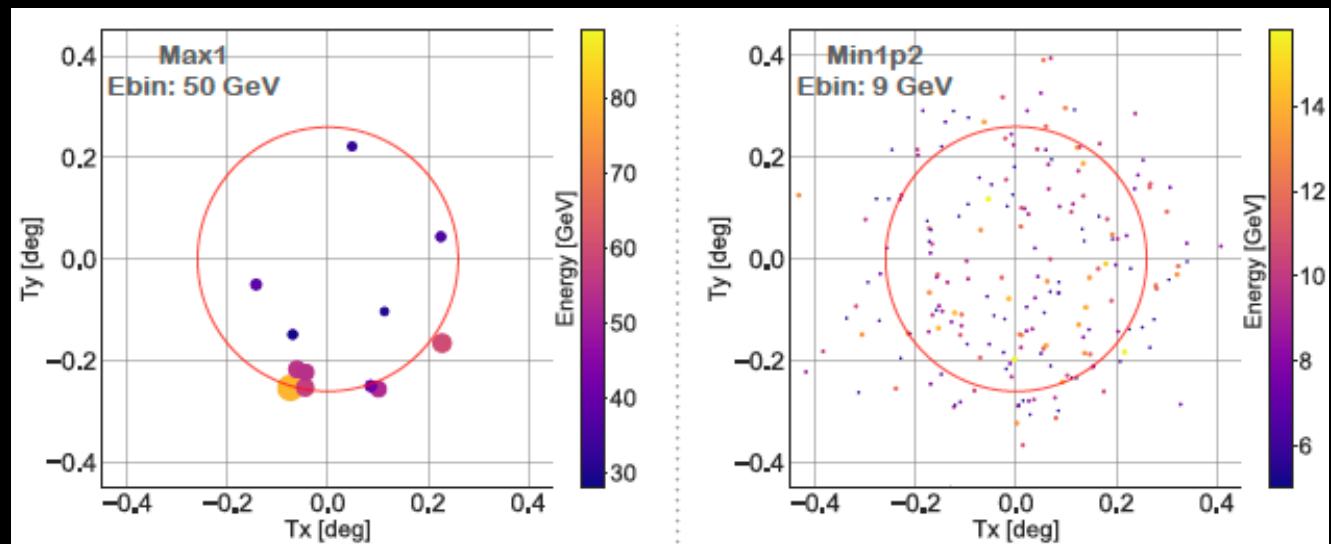
Anisotropic gamma-ray emission pattern



# RESULTS

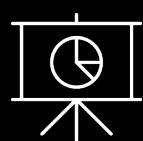


Credits: Linden et al. 2018



Anisotropic gamma-ray emission pattern

Credits: Arsioli & Orlando 2024

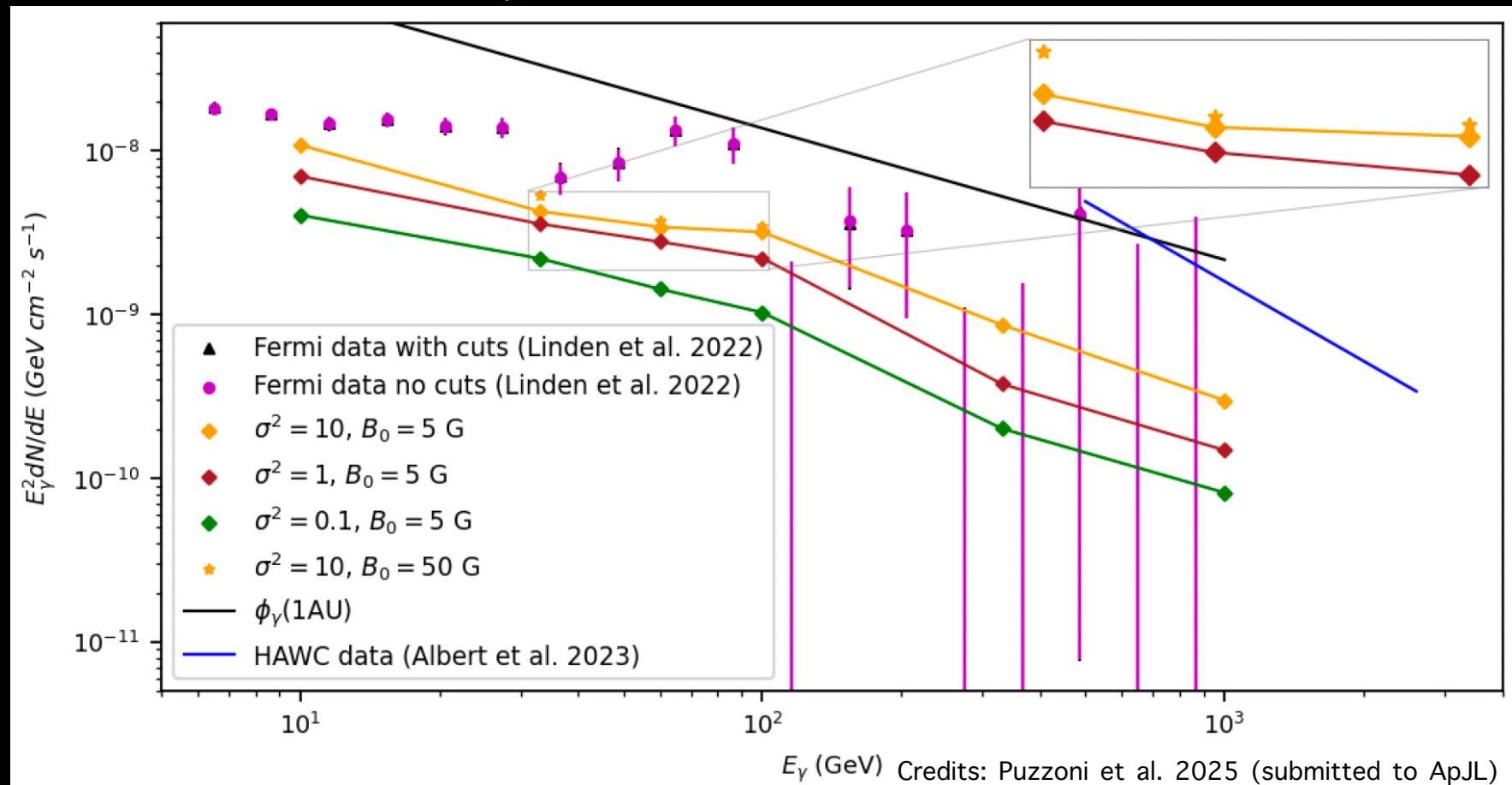


# RESULTS

## $\gamma$ -RAYS FLUX

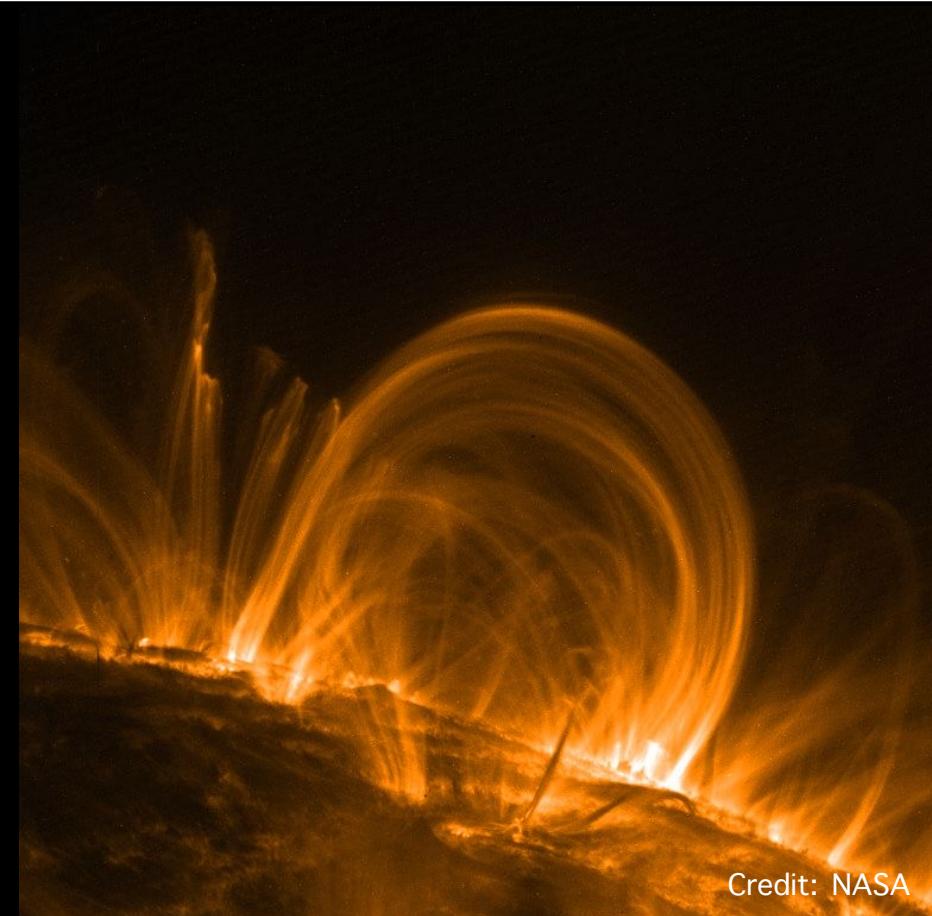
Credits: Kelner et al. 2006

$$\Phi_\gamma(E_\gamma) = \frac{dN_\gamma}{dE_\gamma} = cn_p \int_{E_\gamma}^{\infty} \sigma_{pp}(E_p) J_p(E_p) F_\gamma\left(\frac{E_\gamma}{E_p}, E_p\right) \frac{dE_p}{E_p} \times \frac{2\pi R_\odot^2}{(L \times L) R_\odot^2} \left(\frac{R_\odot^2}{R_{1AU}}\right)^2$$



## *future works*

- Synthetic model reproduces **overall trend** and **order of magnitude** of  $\gamma$ -ray flux both in **Fermi-LAT** and **HAWC** energy range
- Flatness of  $\gamma$ -ray flux in the spectral dip **region** increases with  $\sigma^2$
- Effect of **interchange reconnection** on  $\gamma$ -ray flux



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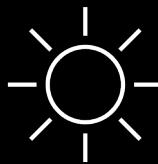
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*Atelier "Physique cinétique des plasmas astrophysiques"*



# MODEL EQUATIONS

Credits: Giacalone and Jokipii 1999

$$\delta \mathbf{B}(x, y, z) = \sum_{n=1}^{N_m} A(k_n) \hat{\xi}_n \exp(i k_n z'_n + i \beta_n), \quad (3)$$

where

$$\hat{\xi}_n = \cos \alpha_n \hat{x}'_n + i \sin \alpha_n \hat{y}'_n \quad (4)$$

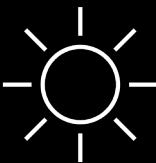
and

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos \theta_n \cos \phi_n & \cos \theta_n \sin \phi_n & -\sin \theta_n \\ -\sin \phi_n & \cos \phi_n & 0 \\ \sin \theta_n \cos \phi_n & \sin \theta_n \sin \phi_n & \cos \theta_n \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}. \quad (5)$$

$$A^2(k_n) = \sigma^2 G(k_n) \left[ \sum_{n=1}^{N_m} G(k_n) \right]^{-1}, \quad (6)$$

where

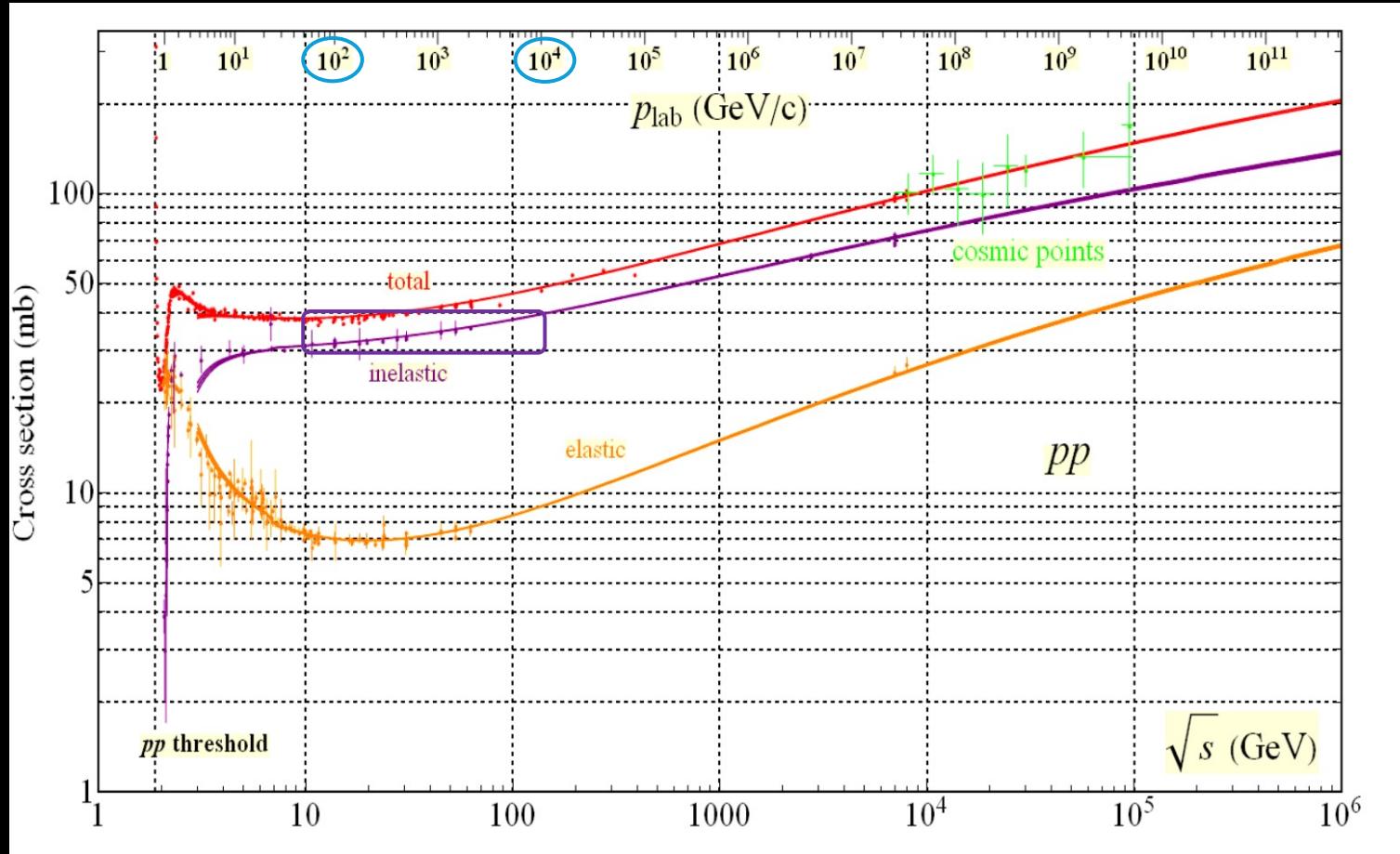
$$G(k_n) = \frac{\Delta V_n}{1 + (k_n L_c)^\gamma}. \quad (7)$$



# P – P INELASTIC CROSS SECTION

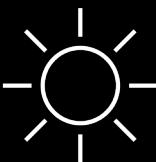
Credits: Kafexhiu et al. 2014

$$\sigma_{pp} = \left[ 30.7 - 0.96 \log\left(\frac{T}{T^{th}}\right) + 0.18 \log^2\left(\frac{T}{T^{th}}\right) \right] \times \left[ 1 - \left(\frac{T^{th}}{T}\right)^{1.9} \right]^3 \text{ mb}$$

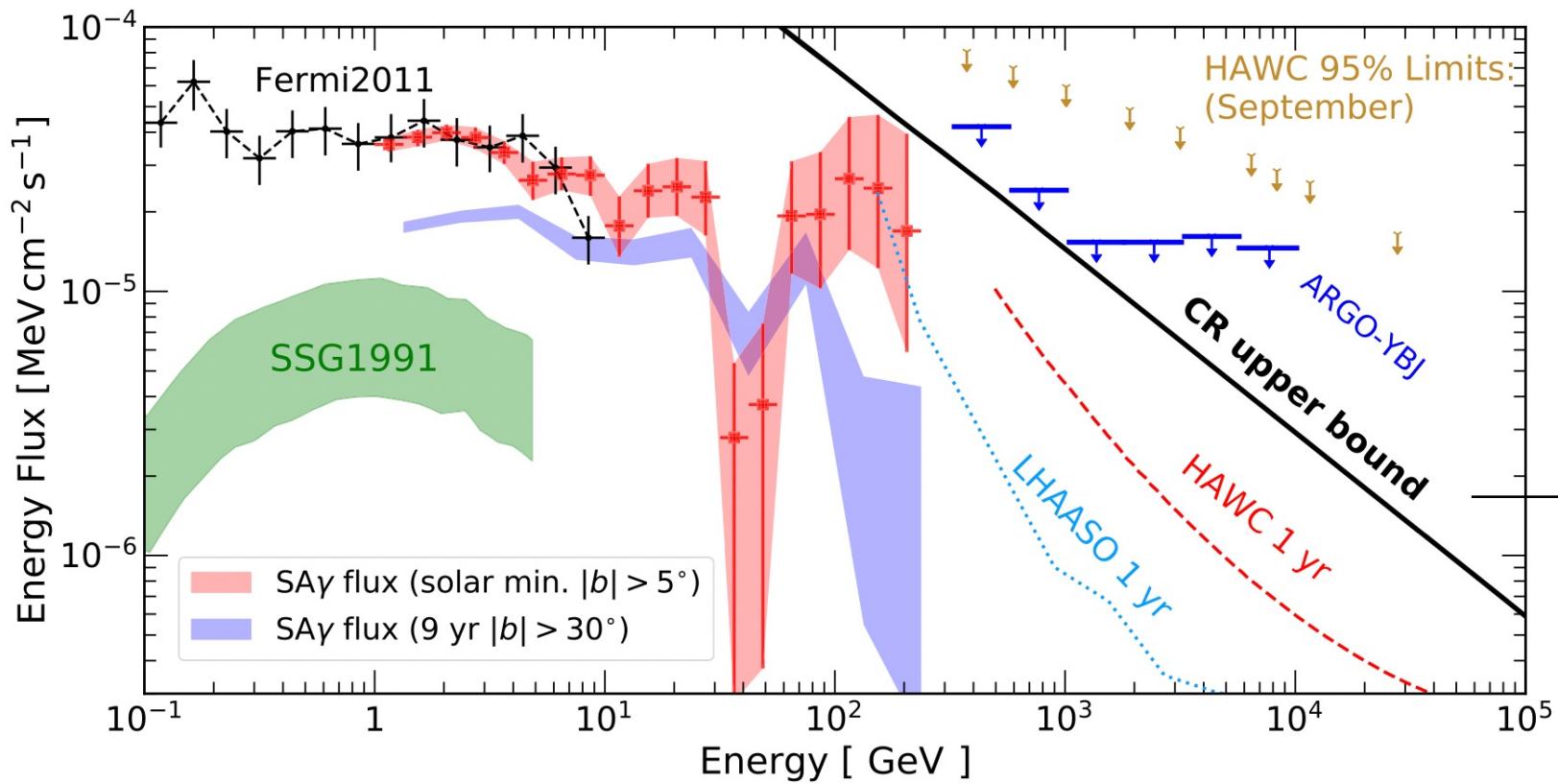


$T$ : proton kinetic energy in the laboratory frame

$T^{th} = 0.2797$  GeV: threshold kinetic energy for the production of neutral and charged pions

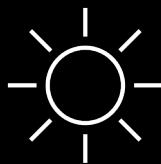


# FERMI DATA



- Fermi 2011
- Fermi solar minimum 2008-2010
- Fermi 9 years 2008-2017

Theoretical maximum gamma-ray flux the Sun can produce with cosmic rays



# DENSITY PROFILE

Credits: Morton et al. 2023

