



Particle Energization in Magnetic Reconnection Jets

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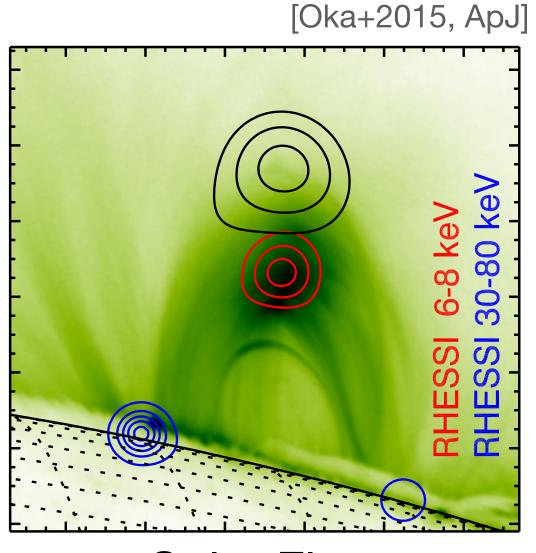
louis.richard@irf.se

Workshop on Kinetic Physics of Astrophysical Plasmas, June 18-20, 2025

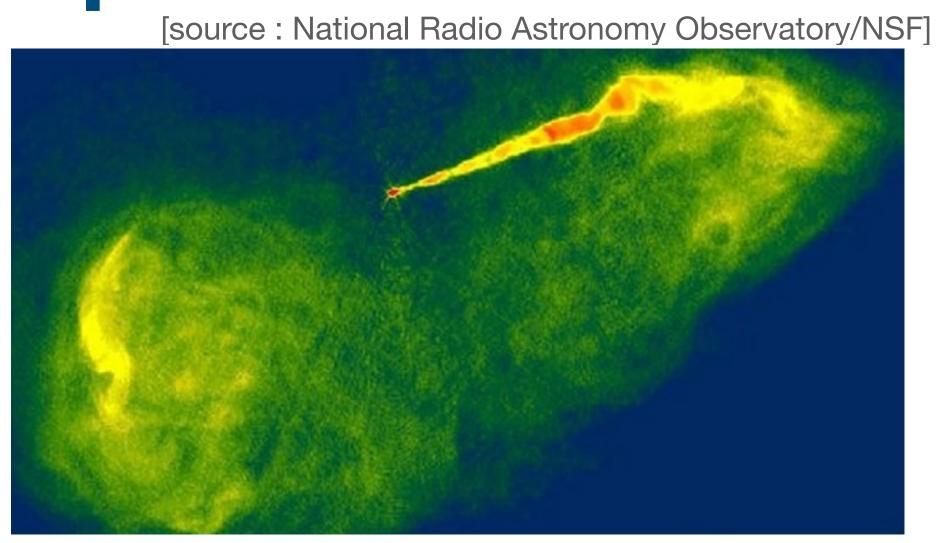
Introduction

Motivation

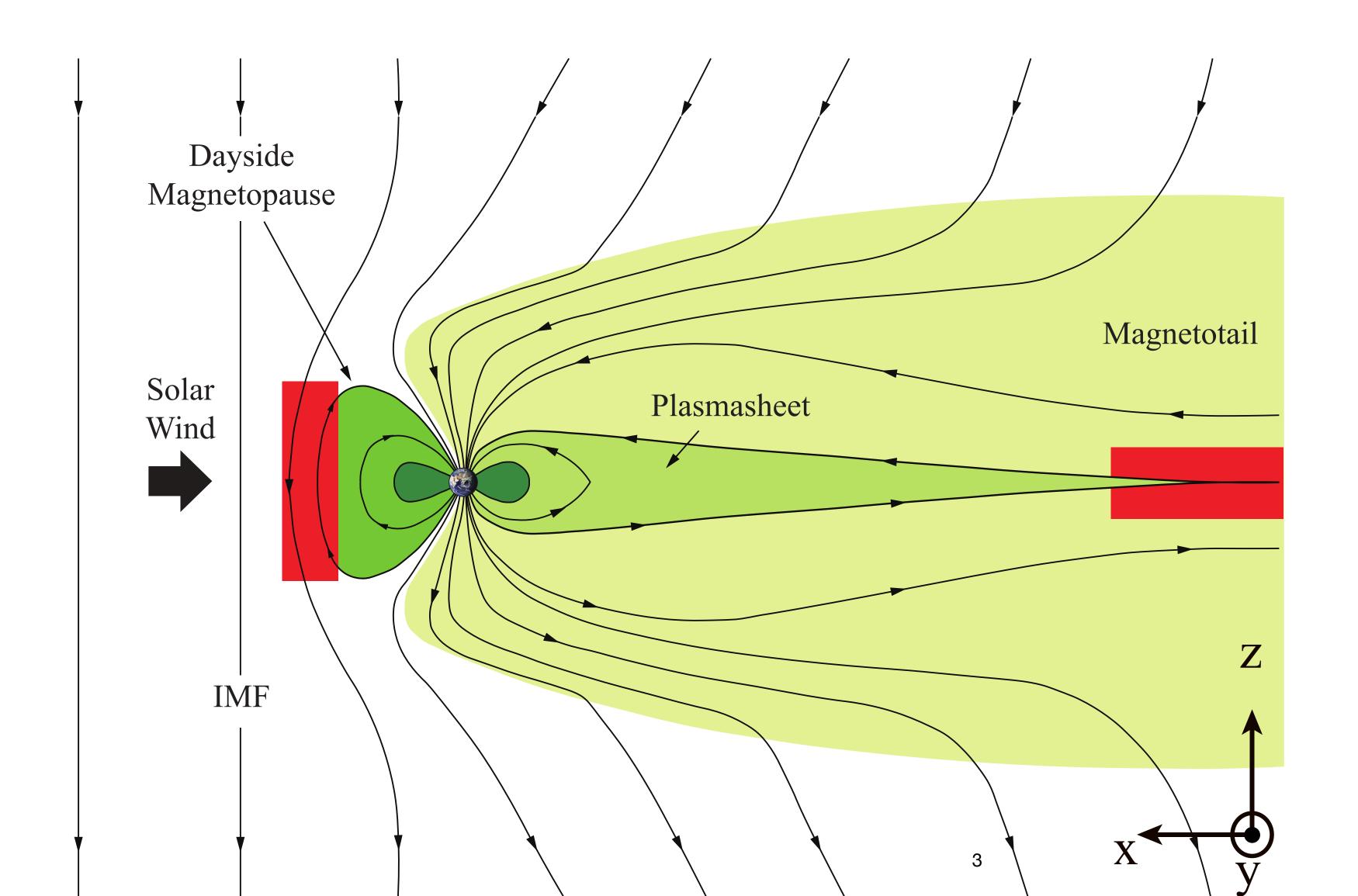
How do magnetic reconnection and turbulence heat the collisionless plasma and accelerate particles?

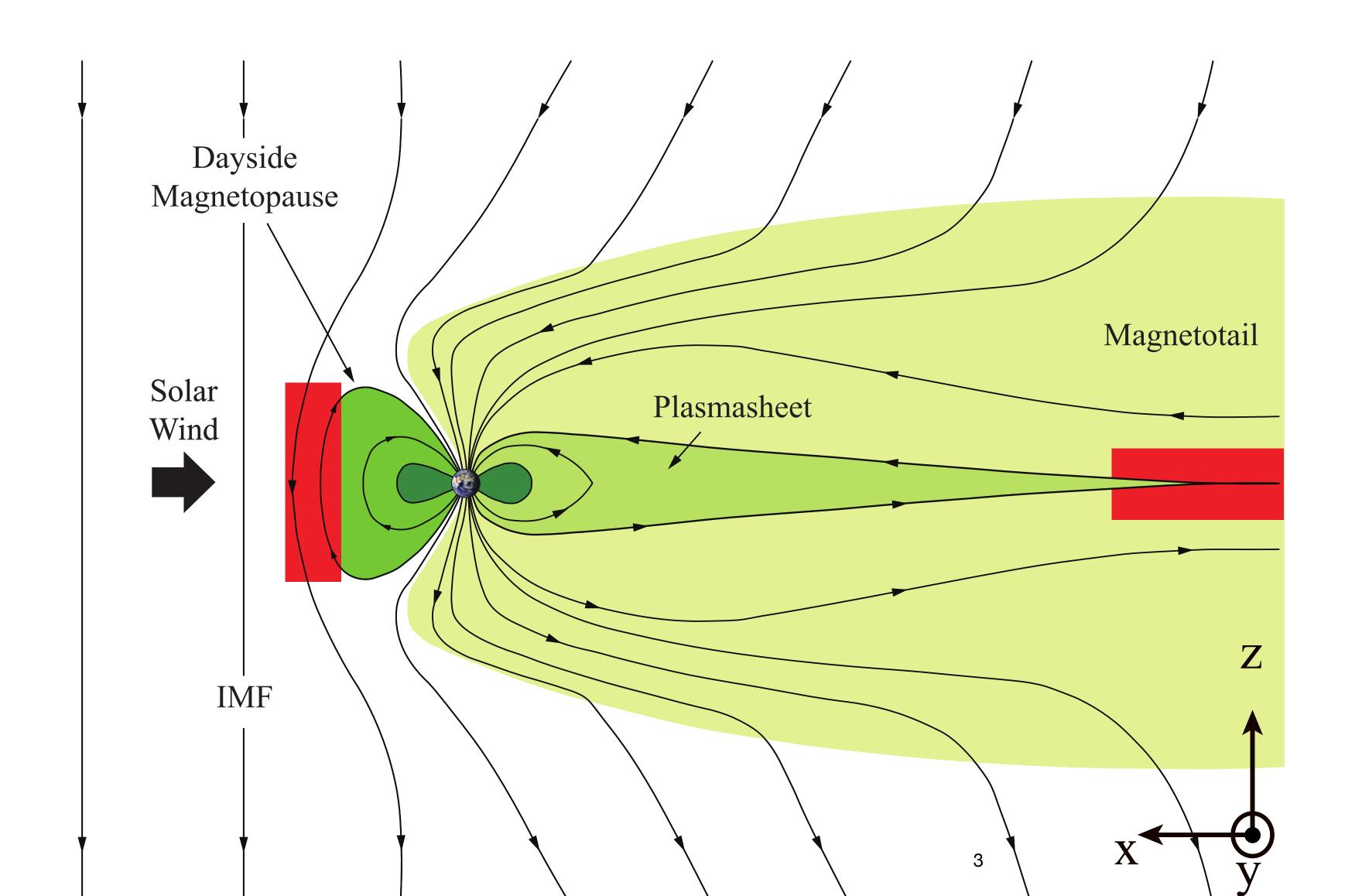


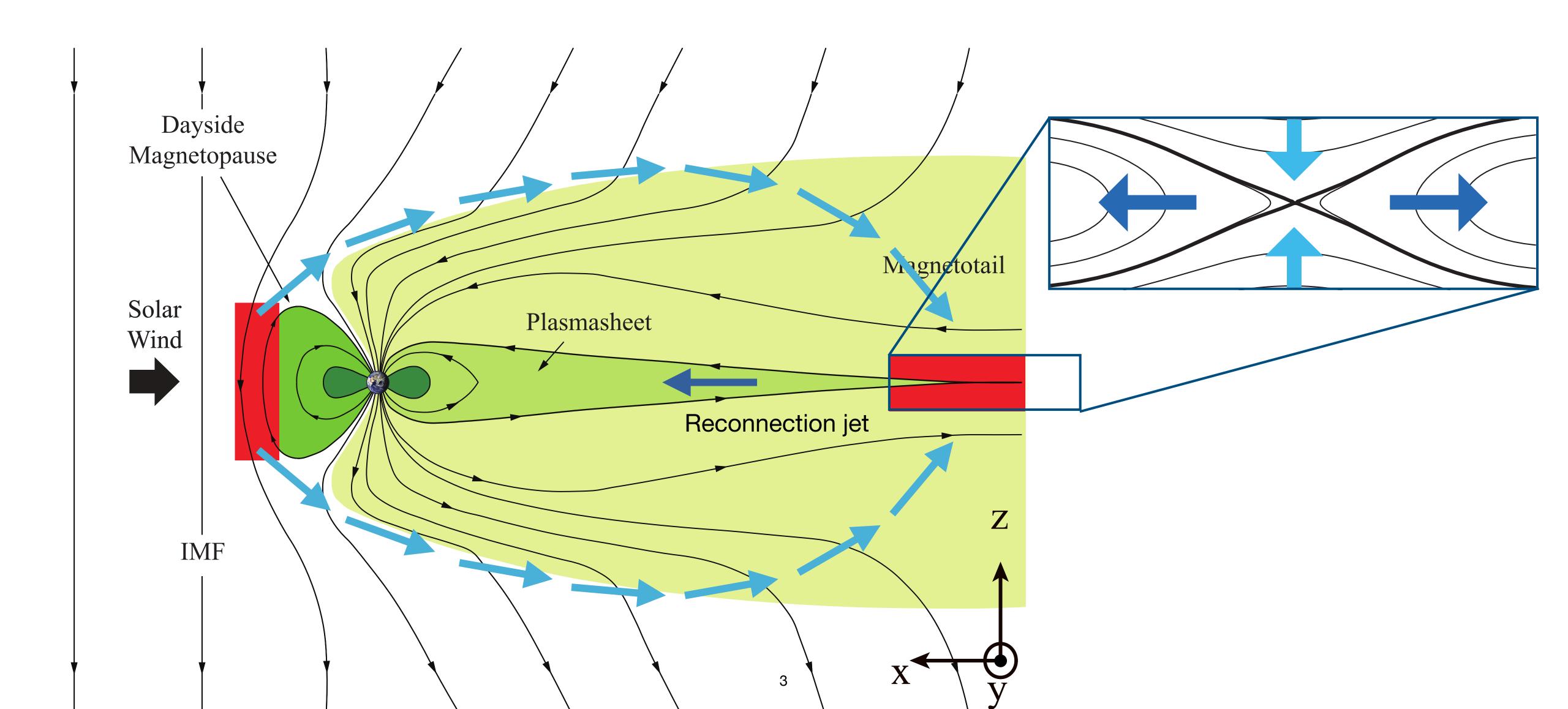
Solar Flares

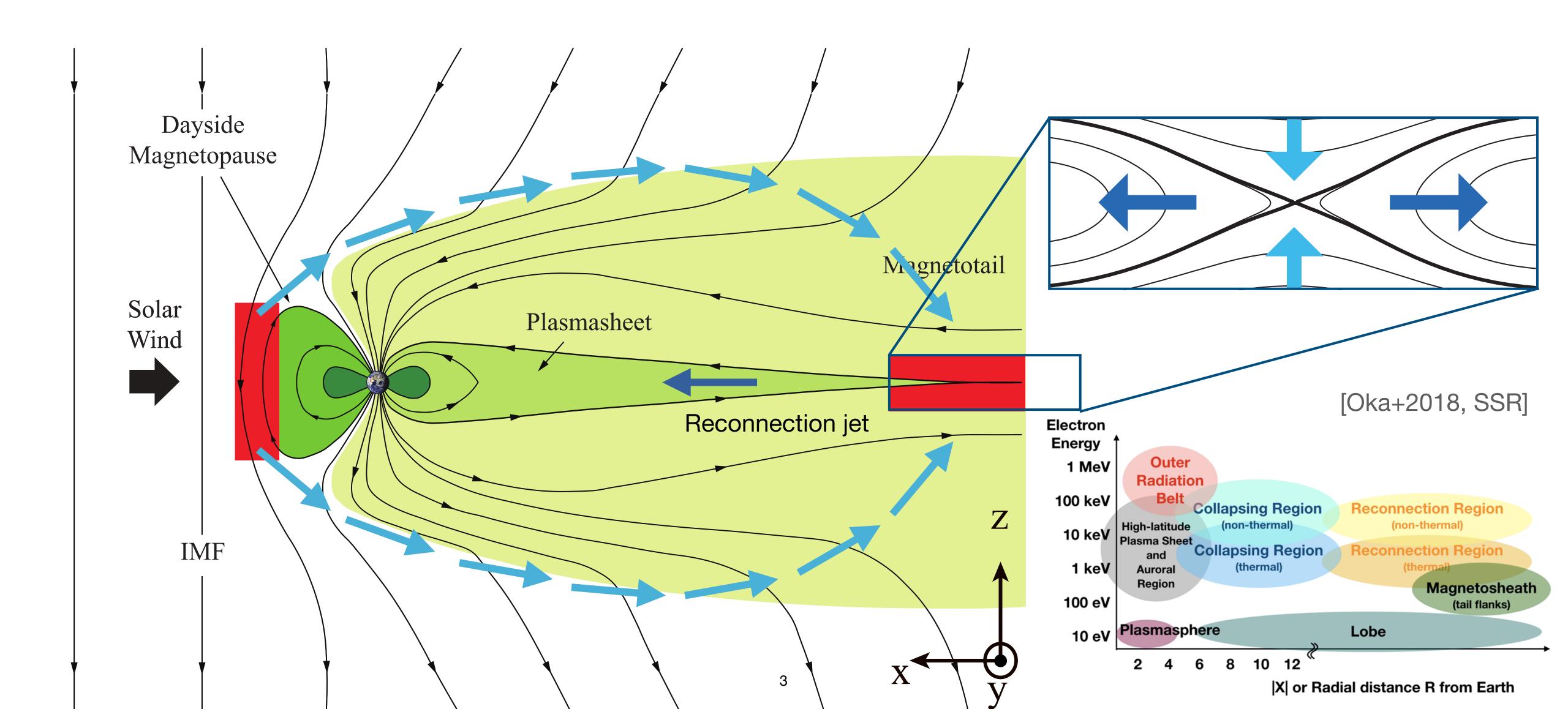


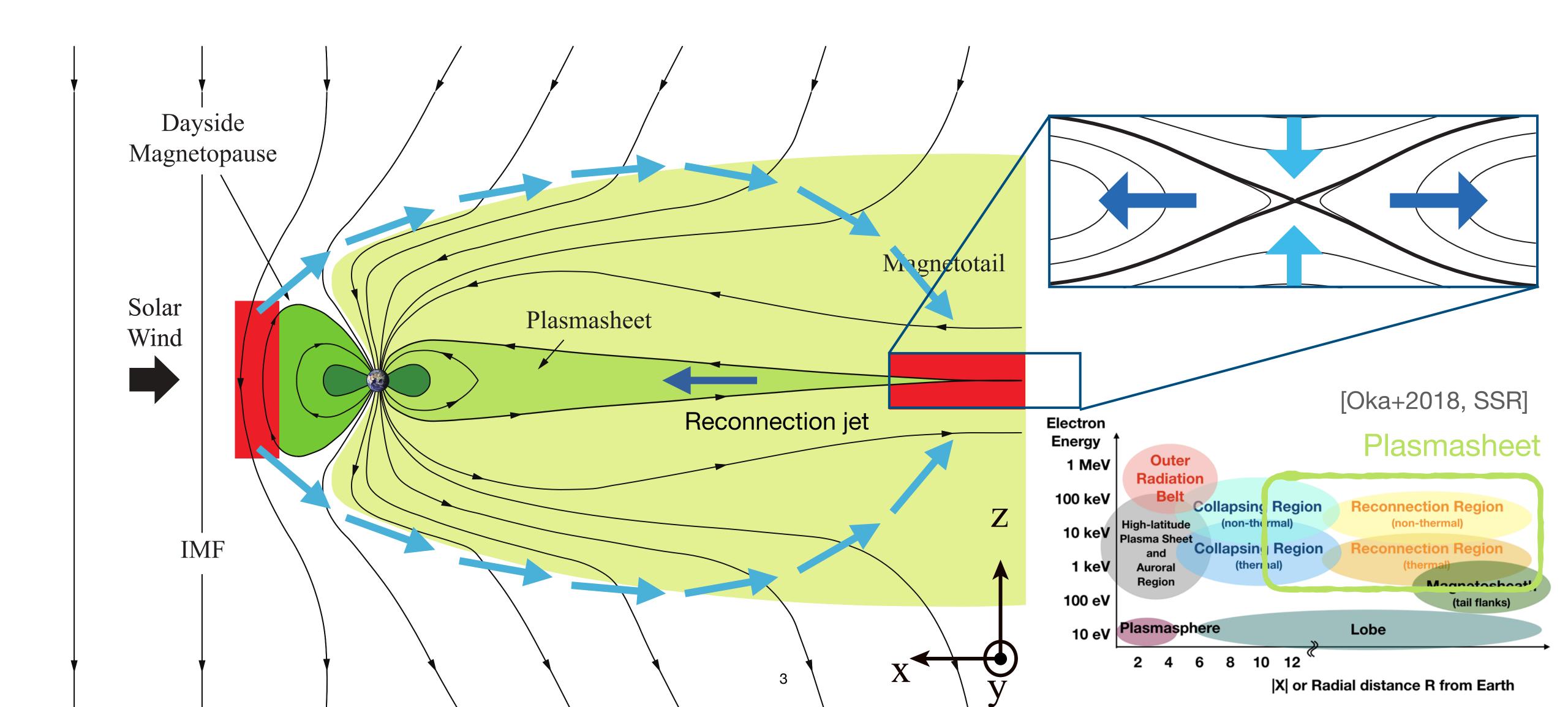
Active Galactic Nuclei (AGN) Jets







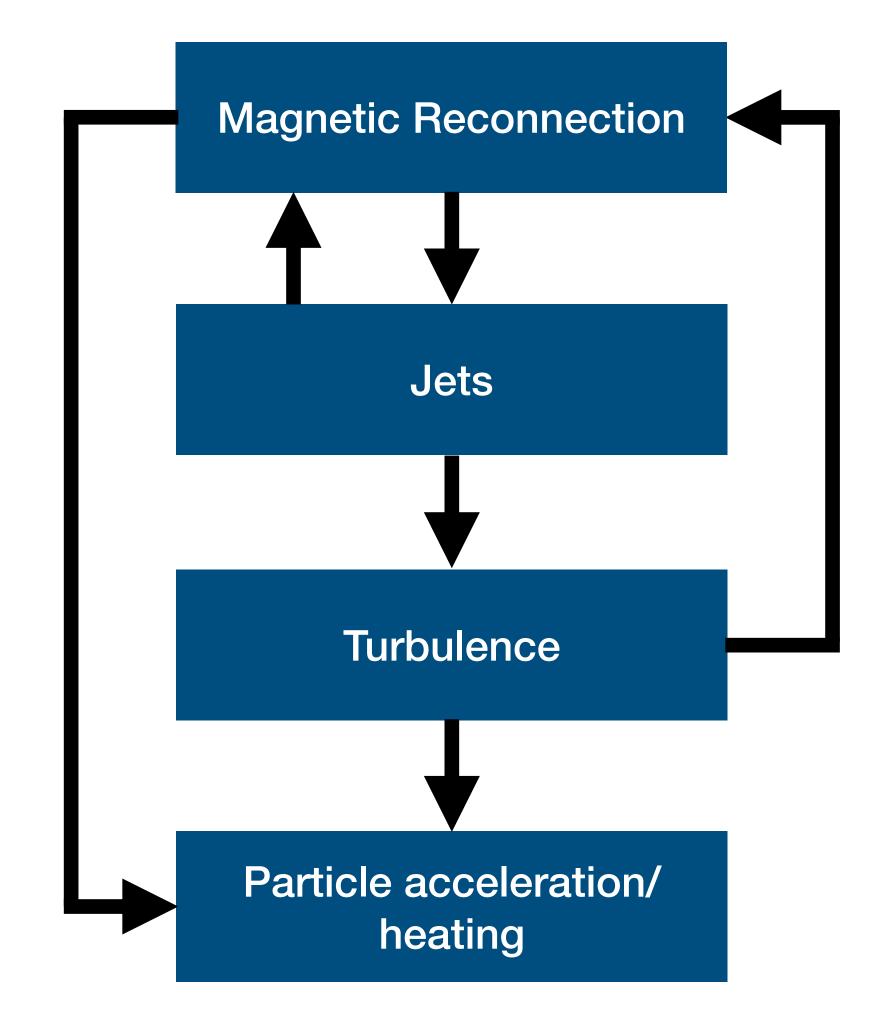




Introduction

Outline

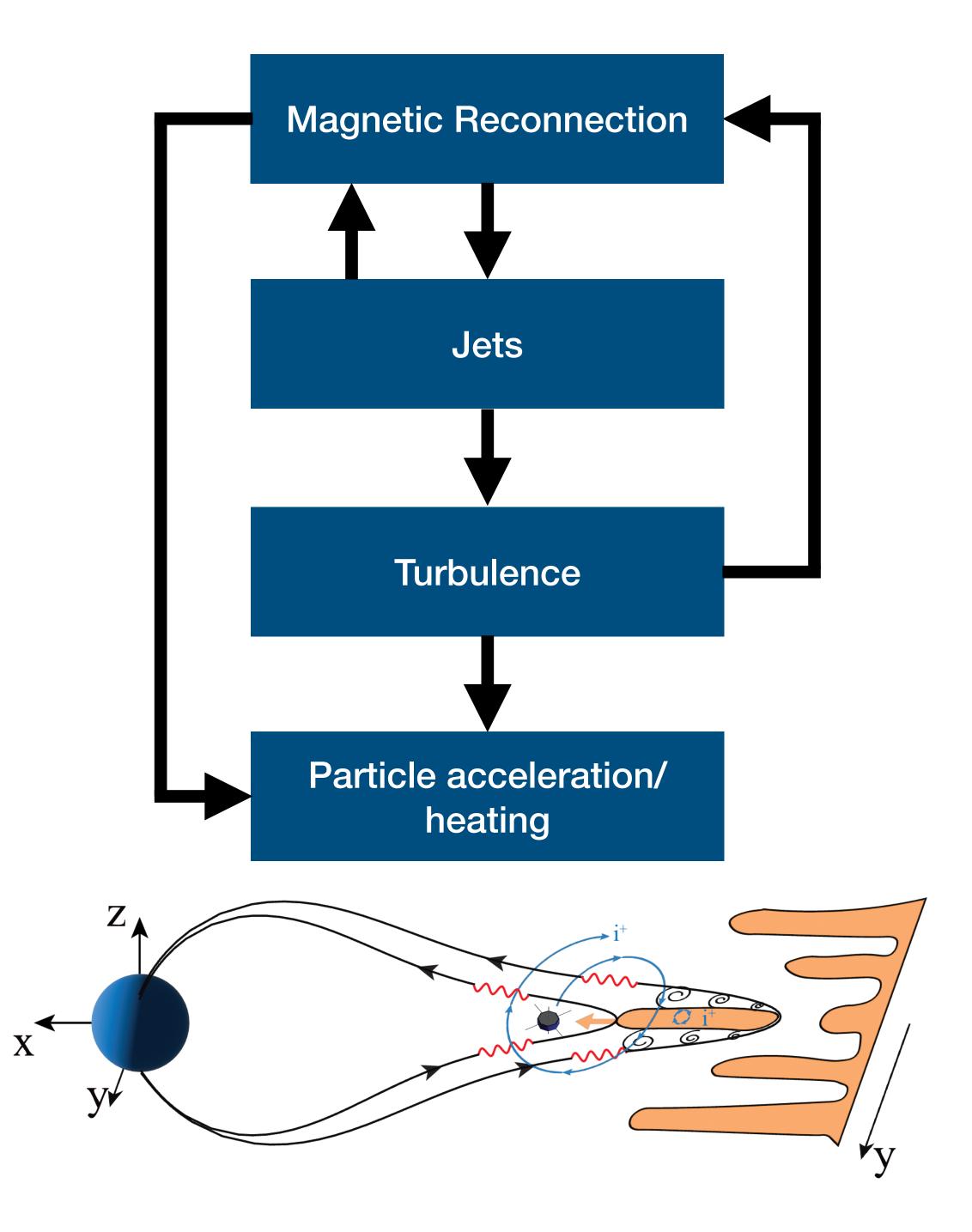
- The MMS mission & in situ observation of reconnection jets
- Turbulence
- Ion heating
- Electron heating
- Conclusions



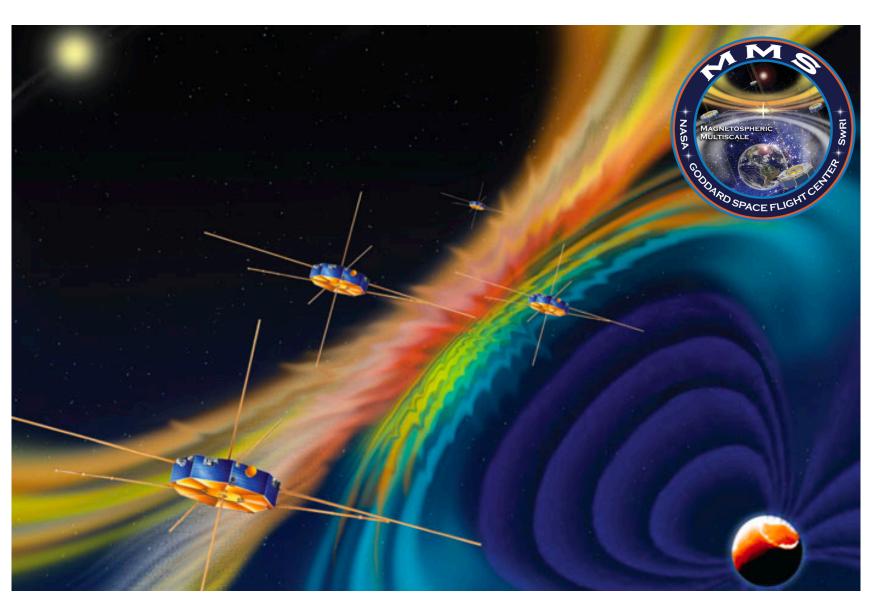
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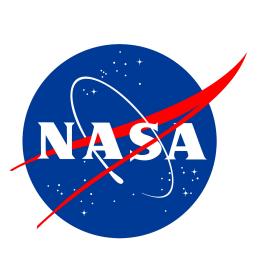


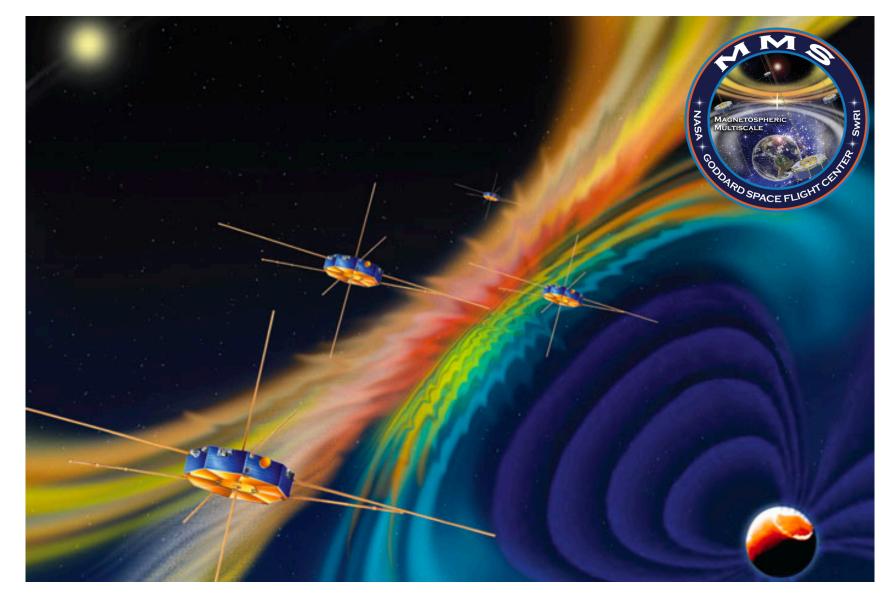
The Magnetospheric Multiscale (MMS) mission[†]

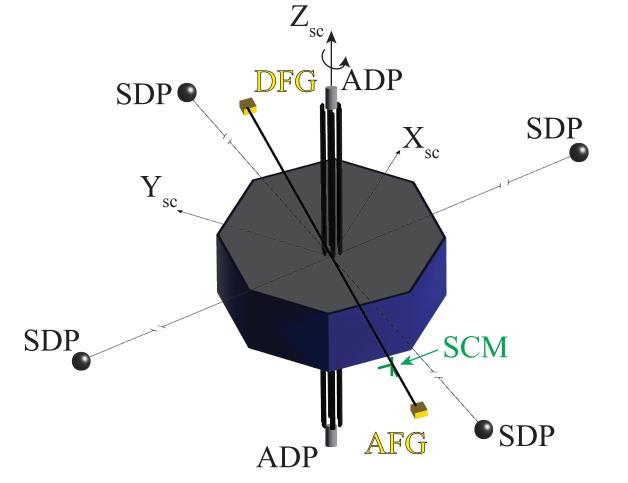


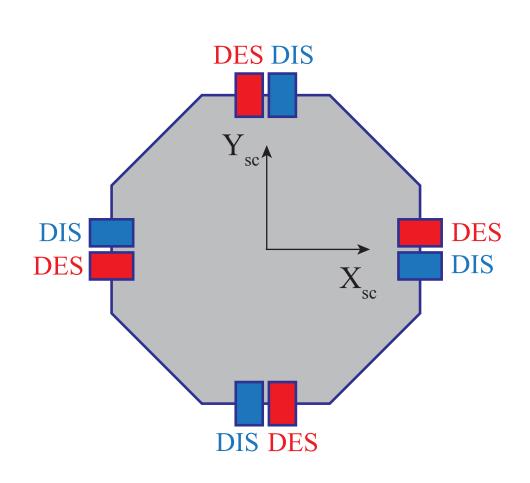
The Magnetospheric Multiscale (MMS) mission[†]

- NASA mission launched in 2015, designed to study magnetic reconnection in Earth's magnetosphere.
- Four spacecraft flying in a tetrahedral formation with ~10 km separation, enabling 3D measurements.
- Fields instruments provide high-cadence measurements of electromagnetic fields:
 - Electric and magnetic field data at up to 64 kHz (electric) and 8 kHz (magnetic).
- Particle instruments deliver high-resolution plasma measurements:
 - Thermal protons sampled every 150 ms and electrons every 30 ms.
 - Spin-resolution measurements (20 s cadence) of massresolved ions.
 - Suprathermal particles: ions at 10 s / 30 s and electrons at 30 s cadence.



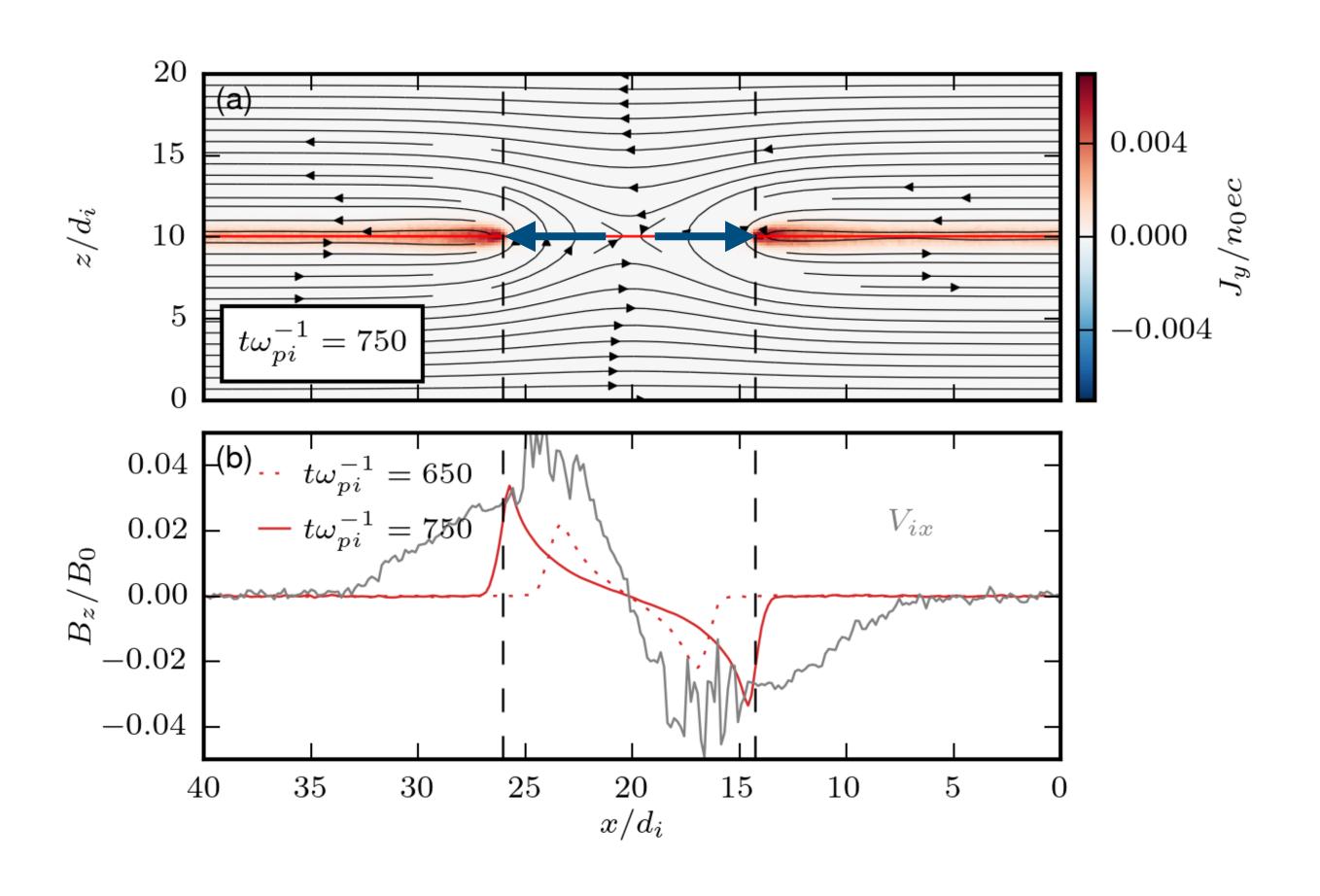






Magnetic Reconnection Jets

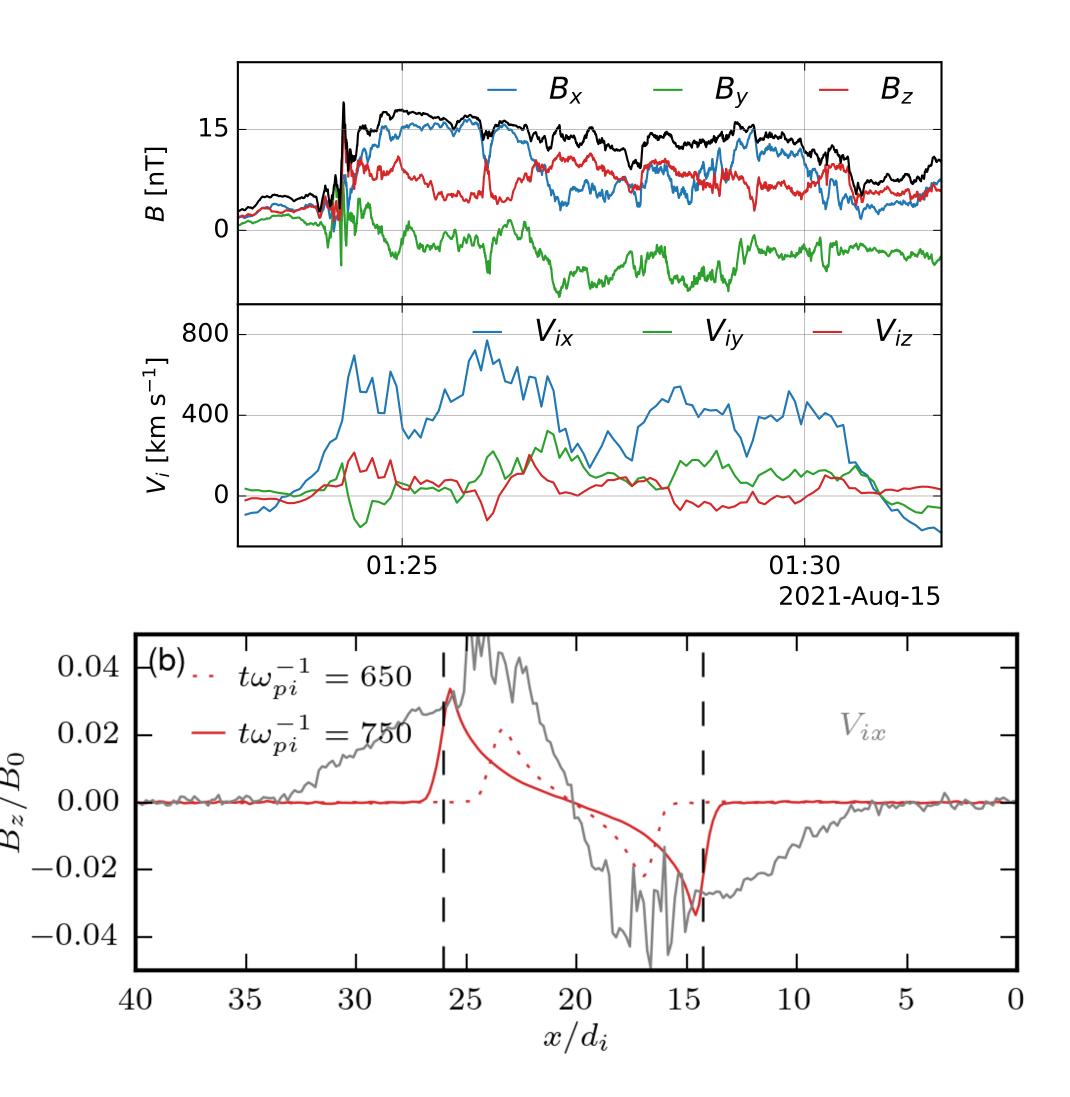
How do they look like?



Magnetic Reconnection Jets

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• Reconnection jets are short (~1 min) and fast ($|\mathbf{V}_i| \ge 300 \text{ km s}^{-1}$) ion flows.

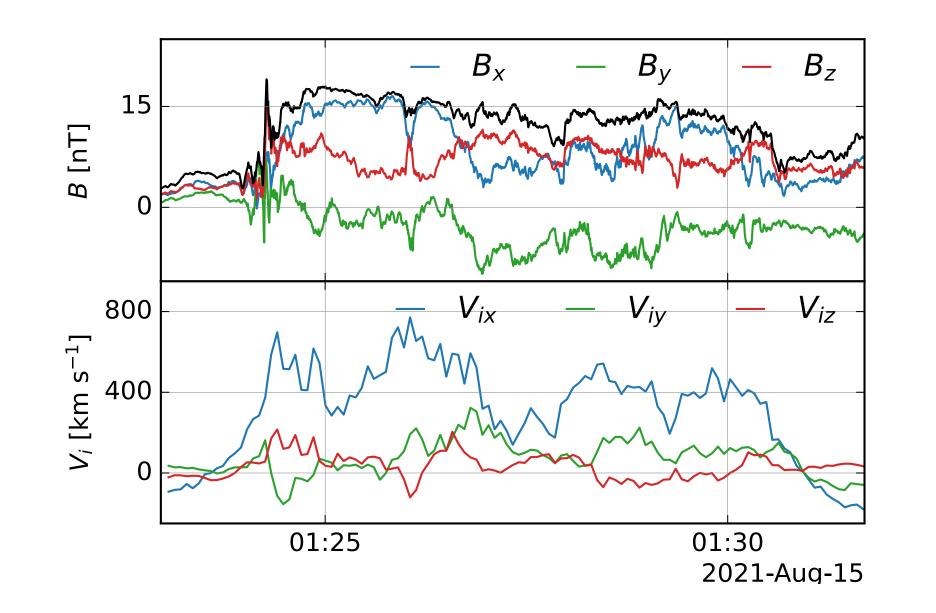


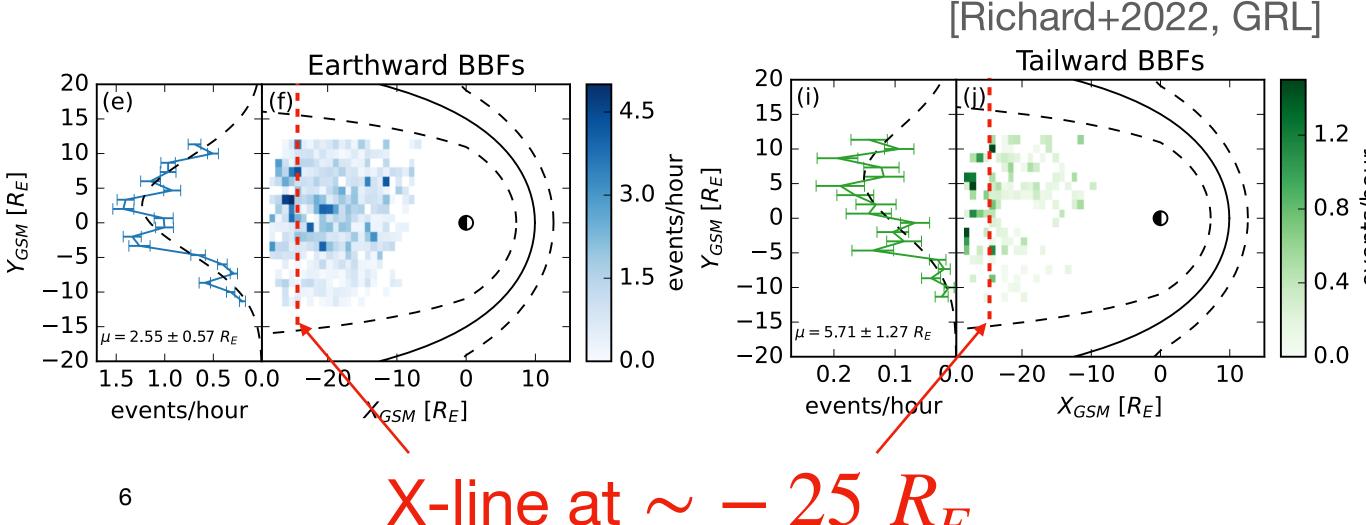
Magnetic Reconnection Jets

How do they look like?

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 Database of 2394 reconnection jets in the Earth's plasma sheet compiled from MMS orbits 2017-2021.

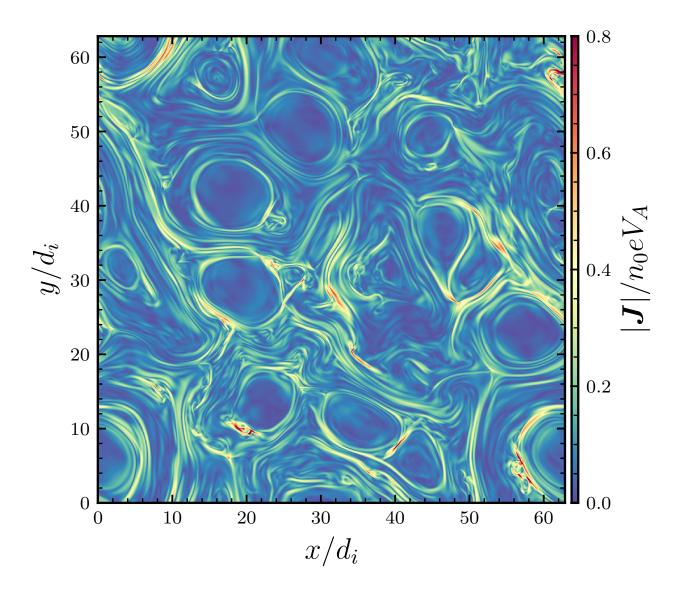




Why do we care?

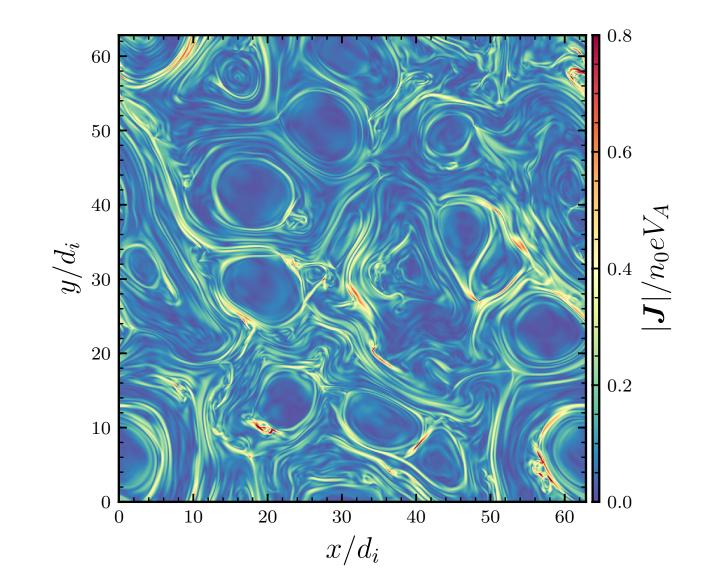
Why do we care?

• Turbulence-driven magnetic reconnection.



Why do we care?

- Turbulence-driven magnetic reconnection.
- Turbulent magnetic reconnection.

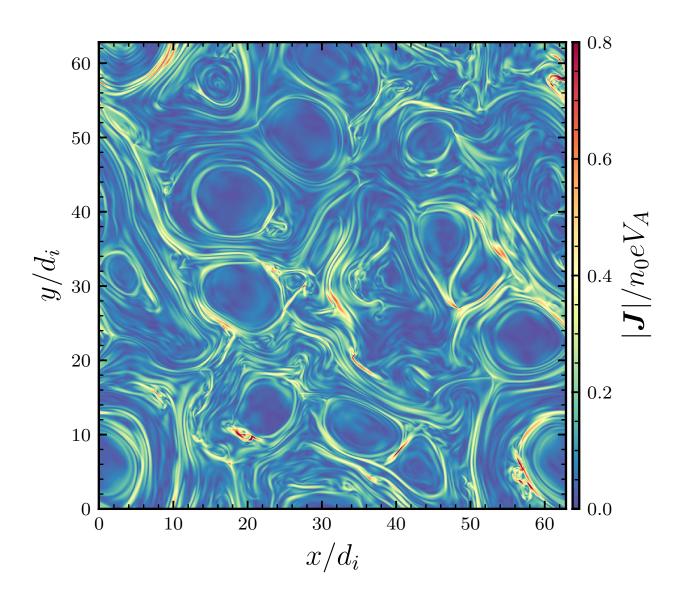


[Daughton+2011, Nat. Phys.]

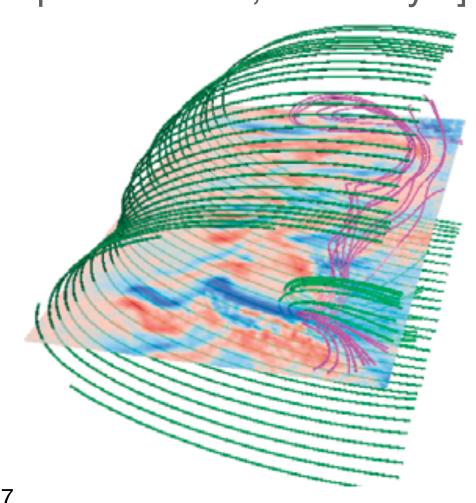


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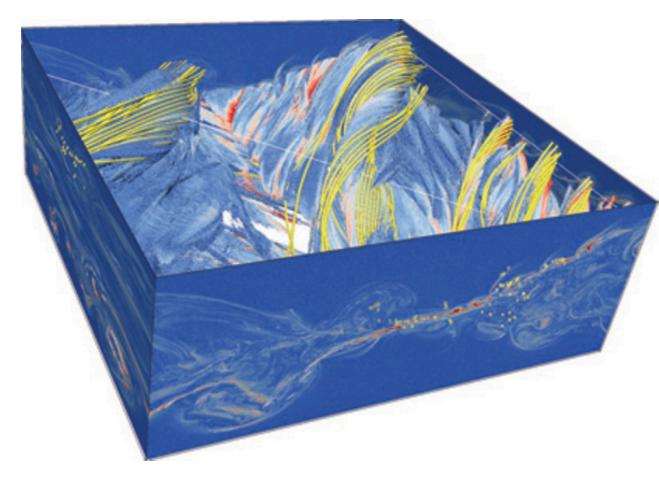
- Turbulence-driven magnetic reconnection.
- Turbulent magnetic reconnection.
- Secondary reconnection sites in turbulent outflows.



[Lapenta+2015, Nat. Phys.]

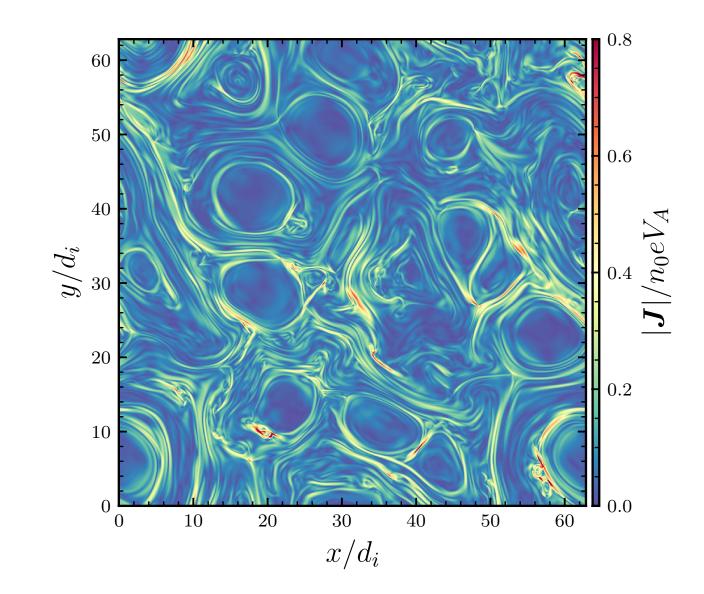


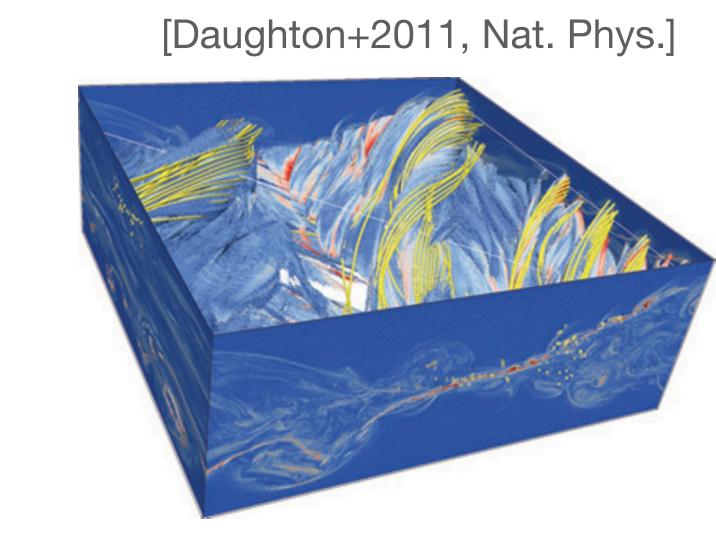
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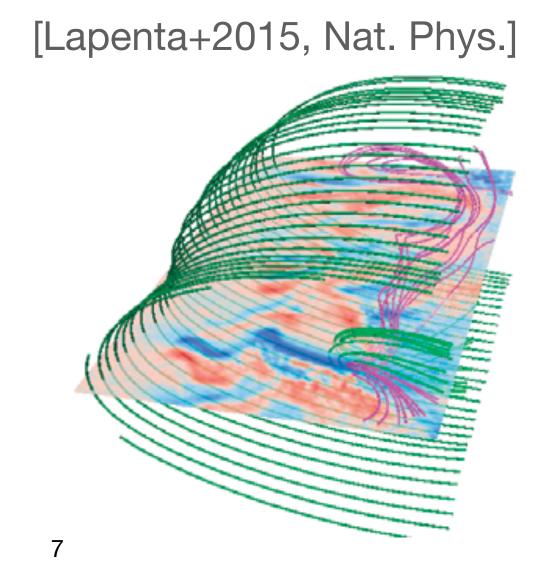


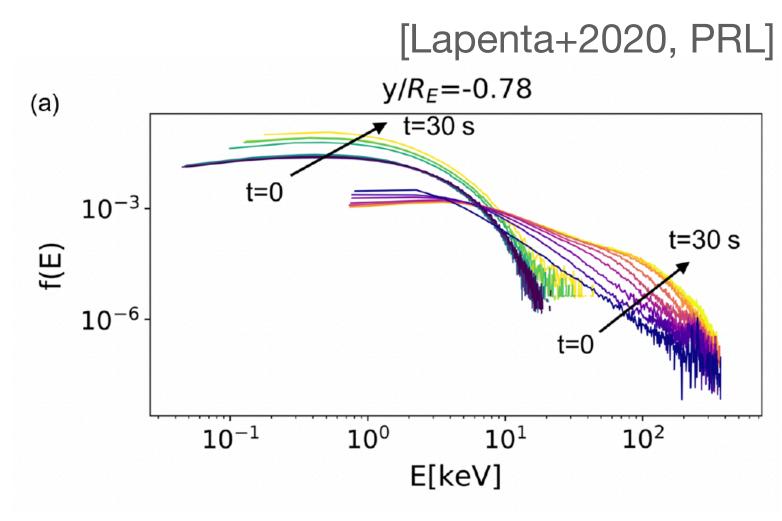
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- Turbulence-driven magnetic reconnection.
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- Secondary reconnection sites in turbulent outflows.
- Non-thermal particle acceleration in turbulent electromagnetic fields





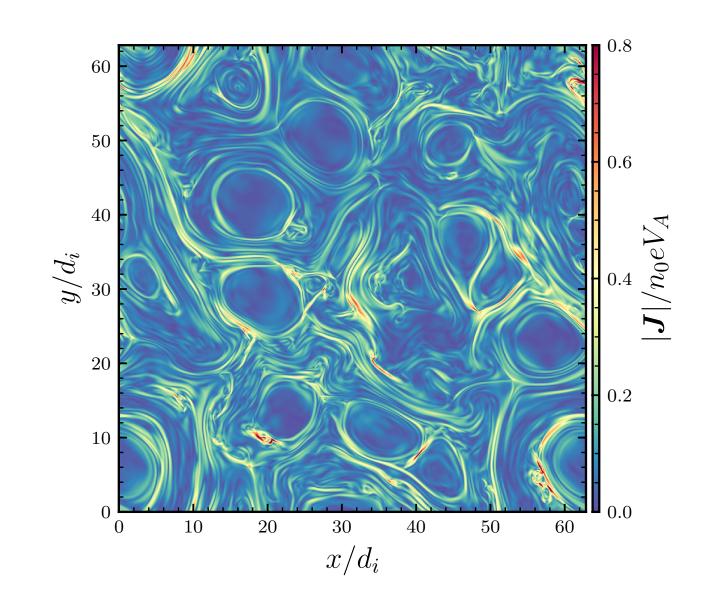


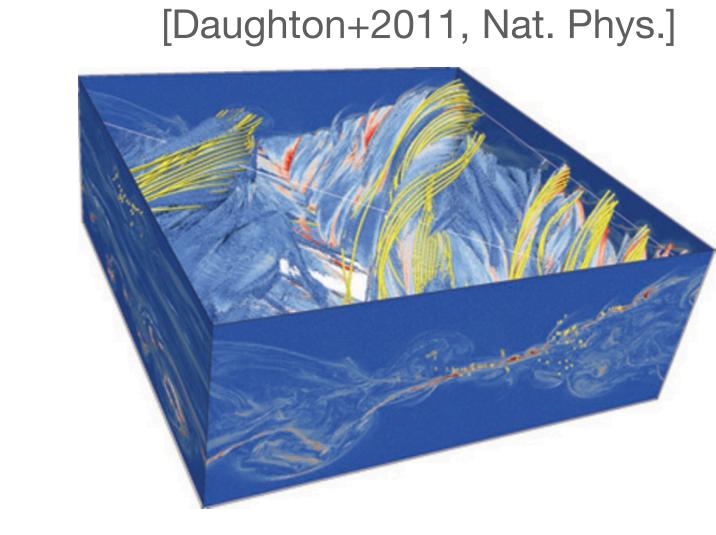


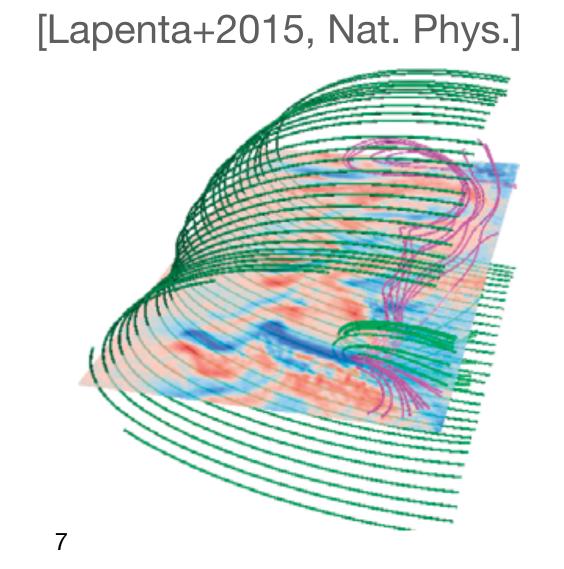
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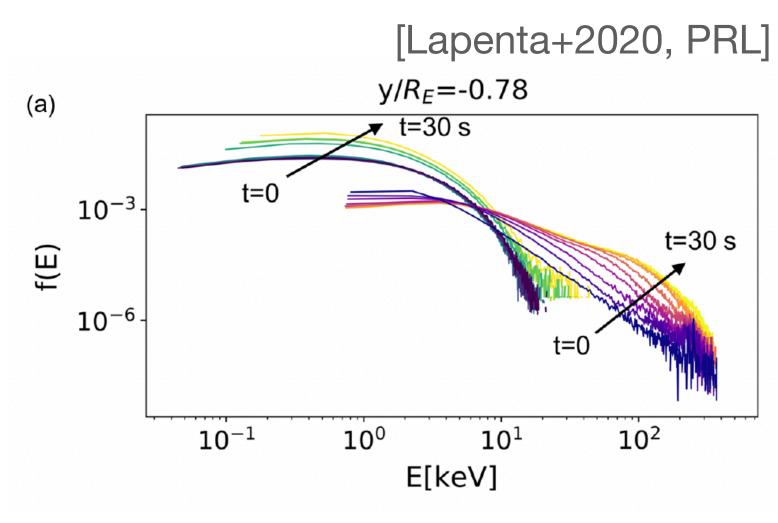
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How much does turbulence contribute to energy conversion in reconnection jets?



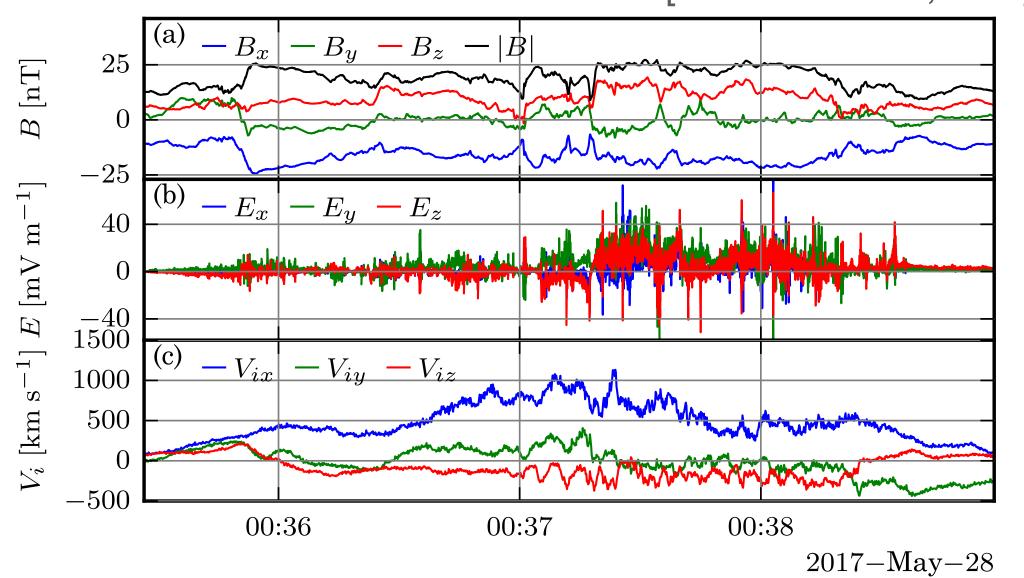






Example

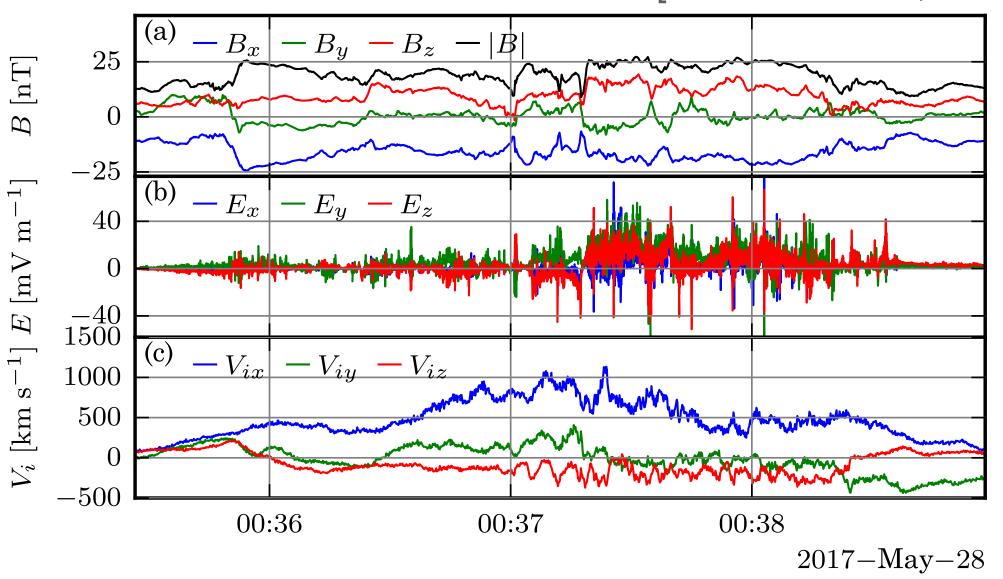
[Richard+2024, PRL]

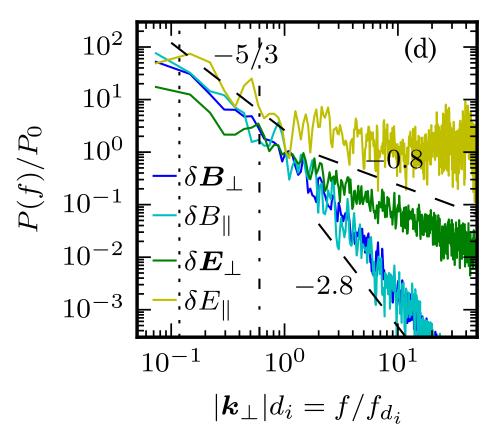


Example

Kolmogorov-like scaling at large (fluid) scales

[Richard+2024, PRL]



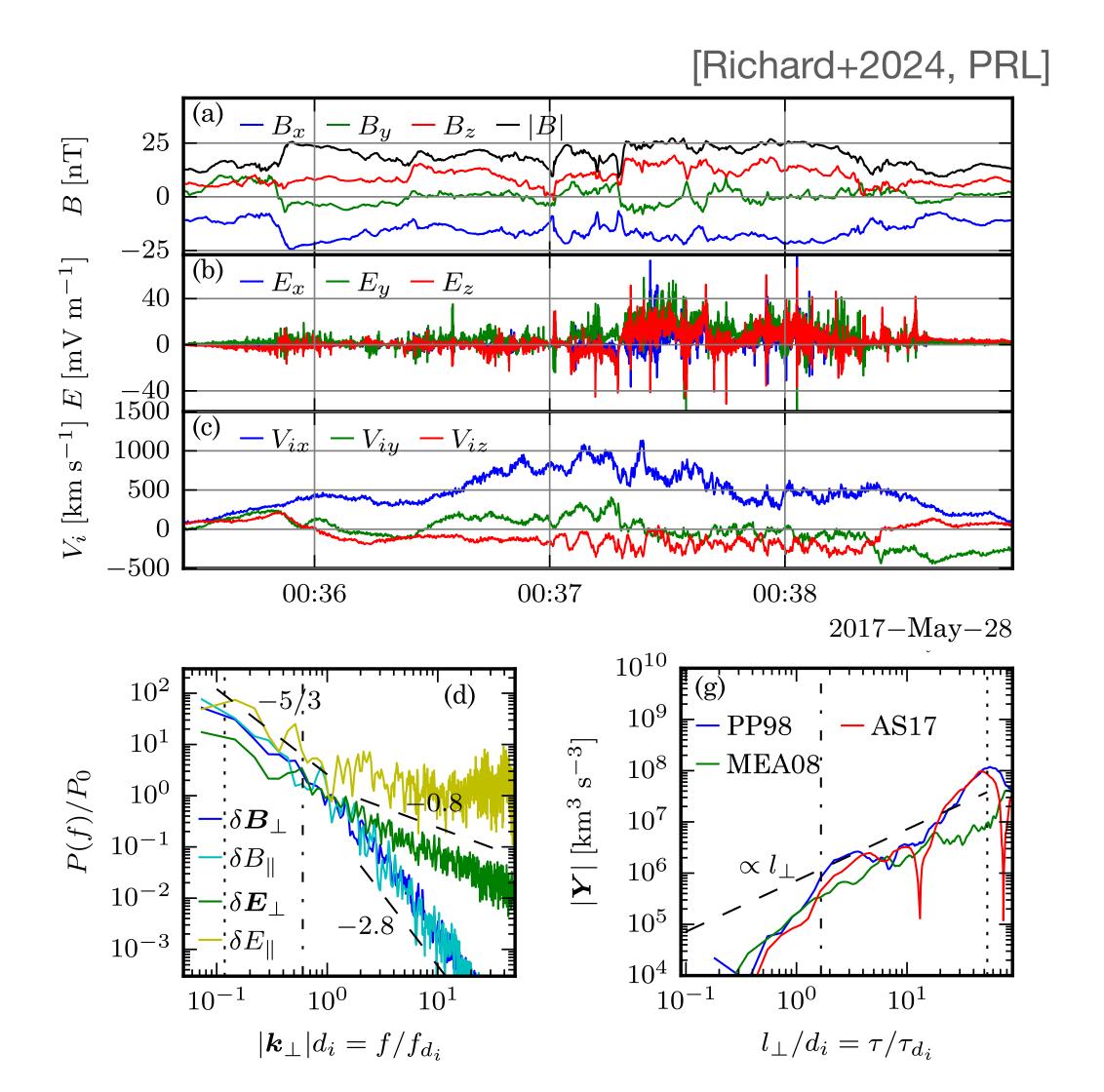


Example

- Kolmogorov-like scaling at large (fluid) scales
- Third-order energy transfer for MHD turbulence [e.g., Politano & Pouquet, 1993, PRE]

$$-2\varepsilon = \frac{1}{2}\nabla_l \cdot \mathbf{Y} + S,$$

with ε the energy "cascade" rate, Y the total energy flux, S, the total energy source



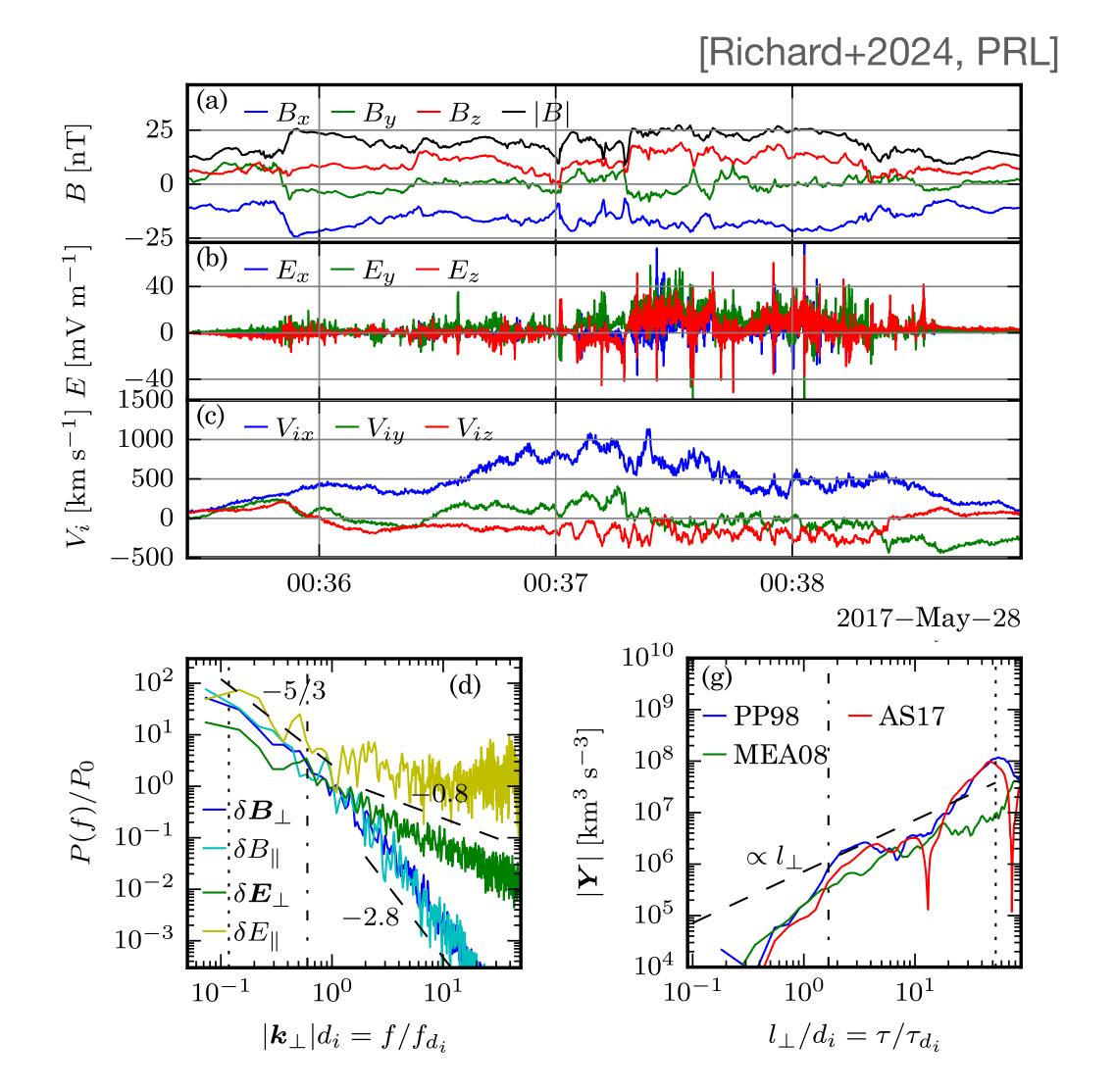
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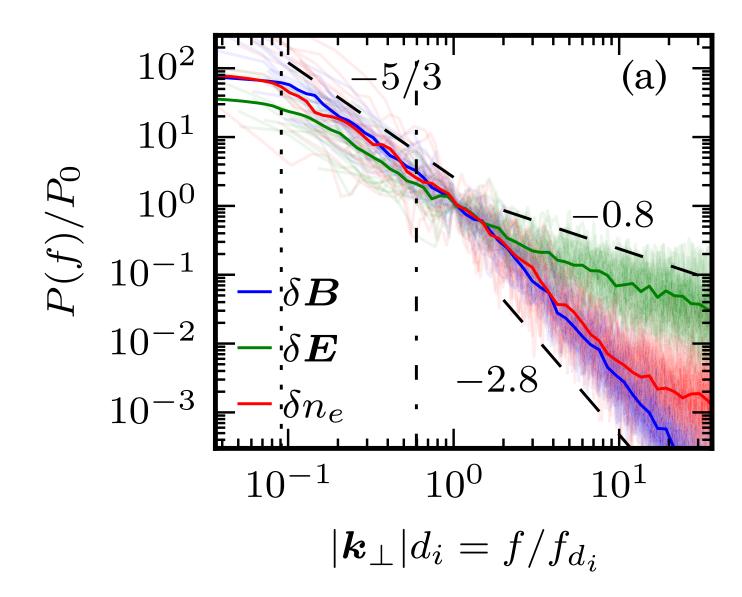
The reconnection outflow shows signs of fully developed turbulence



Statistics

Statistics

 Kolmogorov-like scaling at large (fluid) scales.



[Richard+2024, PRL]

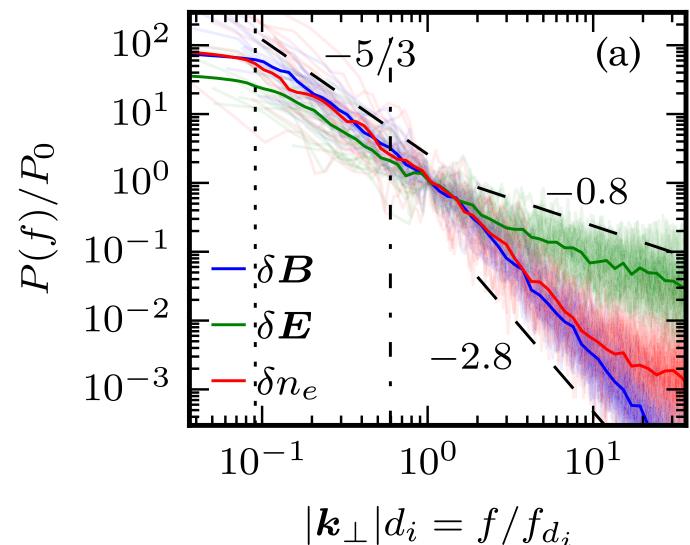
Statistics

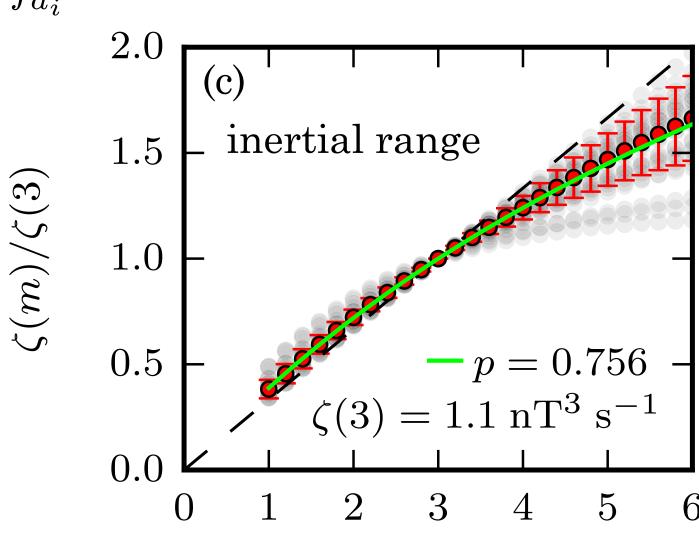
- Kolmogorov-like scaling at large (fluid) scales.
- Breakdown of Kolmogorov's selfsimilarity assumption, i.e., intermittency
- Degree of intermittency estimated using the p-model [Meneveau and Sreenivasan, 1987, PRL]

$$\frac{\zeta(m)}{\zeta(3)} = 1 - \log_2(p^{mh} + (1-p)^{mh})$$

The large-scale fluctuations indicate fully developed turbulence.





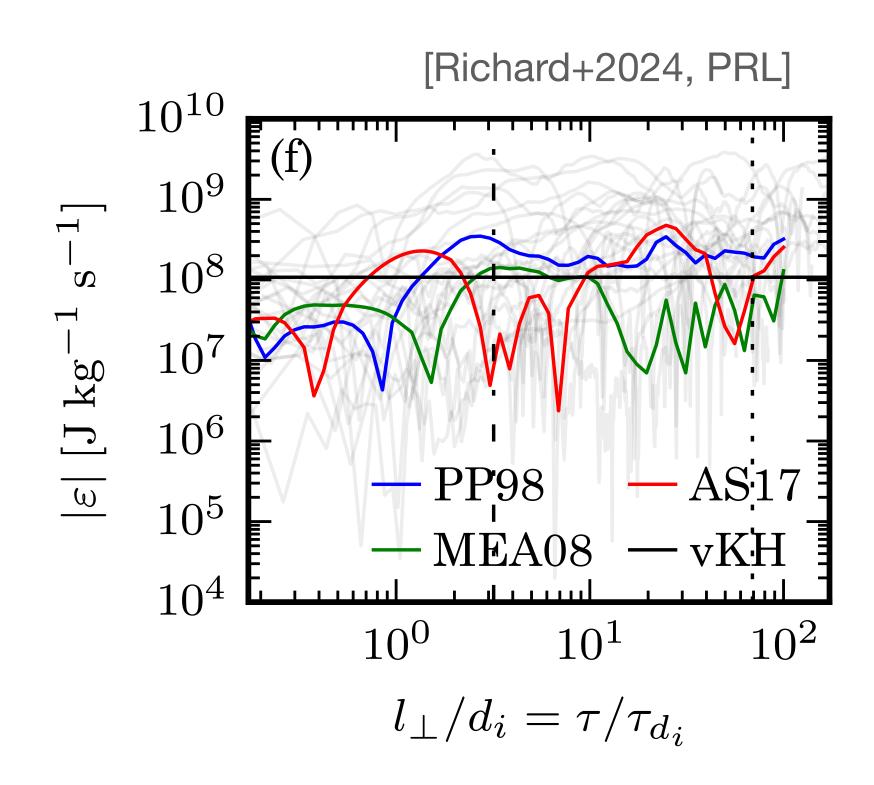


m

Energy cascade rate

Energy cascade rate

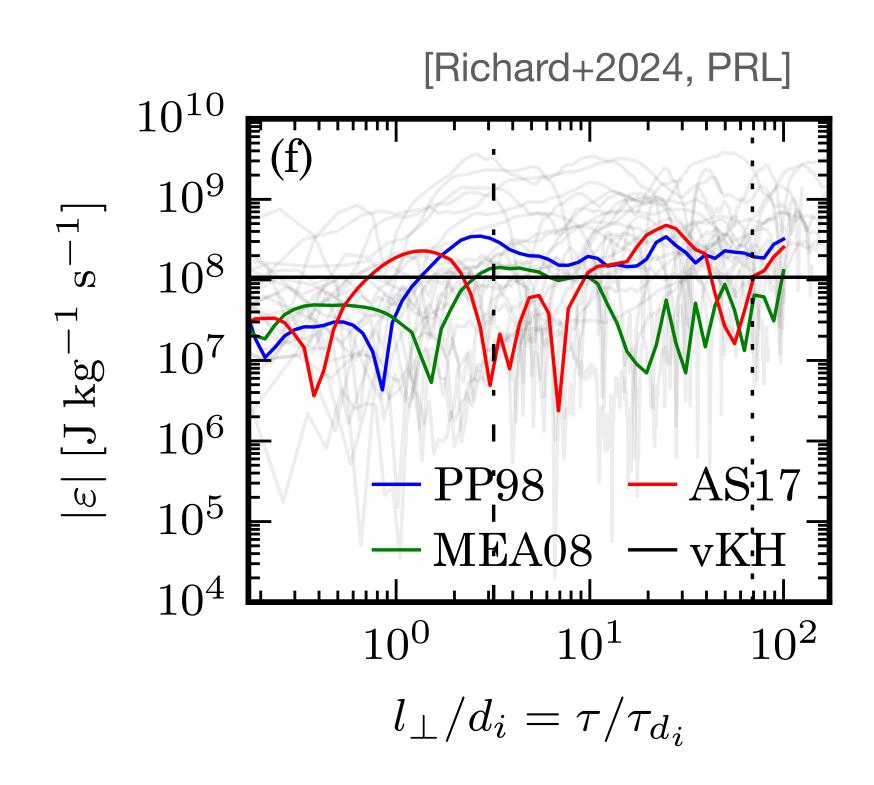
• Energy cascade rate $\varepsilon \sim 10^8~\mathrm{J~kg^{-1}~s^{-1}}$



Energy cascade rate

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	SW	CME sheath	MSH	PSBL	KH
$arepsilon_{jet}/arepsilon$	10 ⁵	10 ³	10 ²	10	4

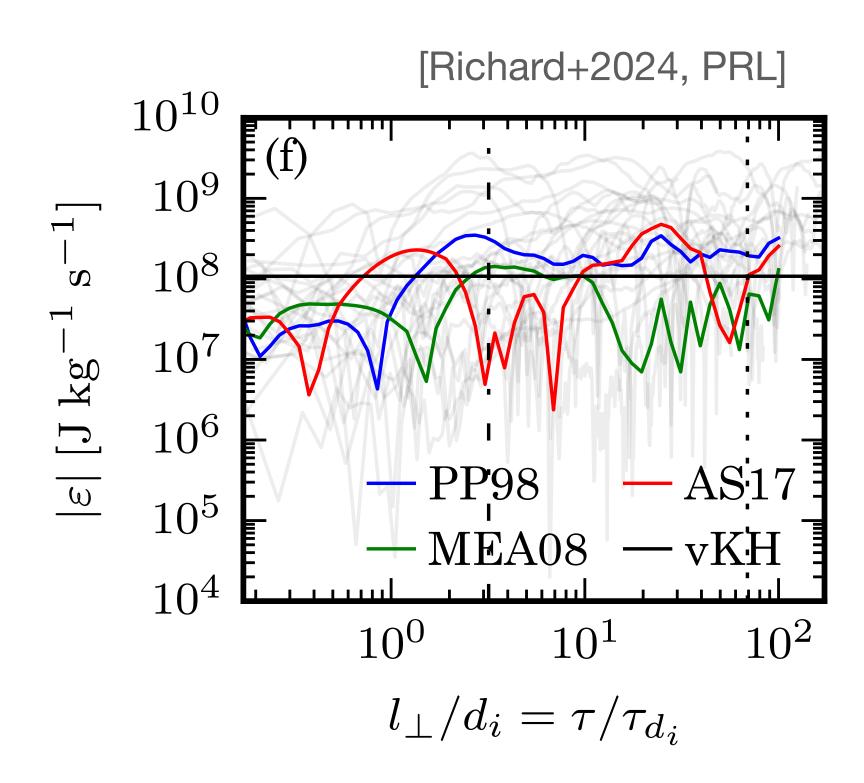


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 Energy injection (Von-Kármán-Howarth energy decay rate) is **balanced** by energy transfer across large scales.



Energy cascade rate

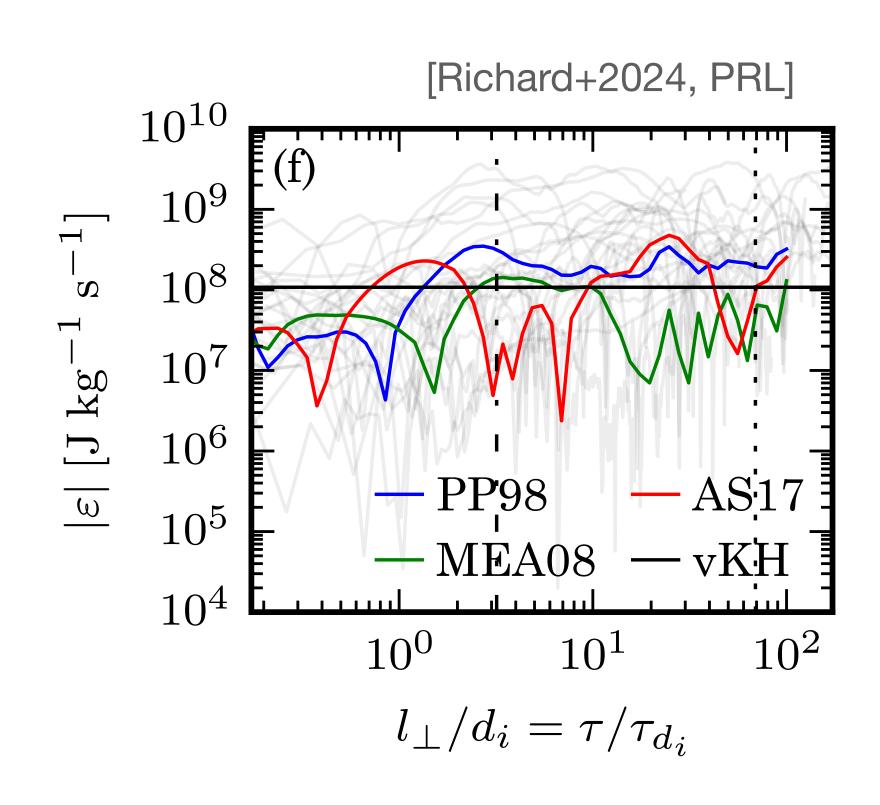
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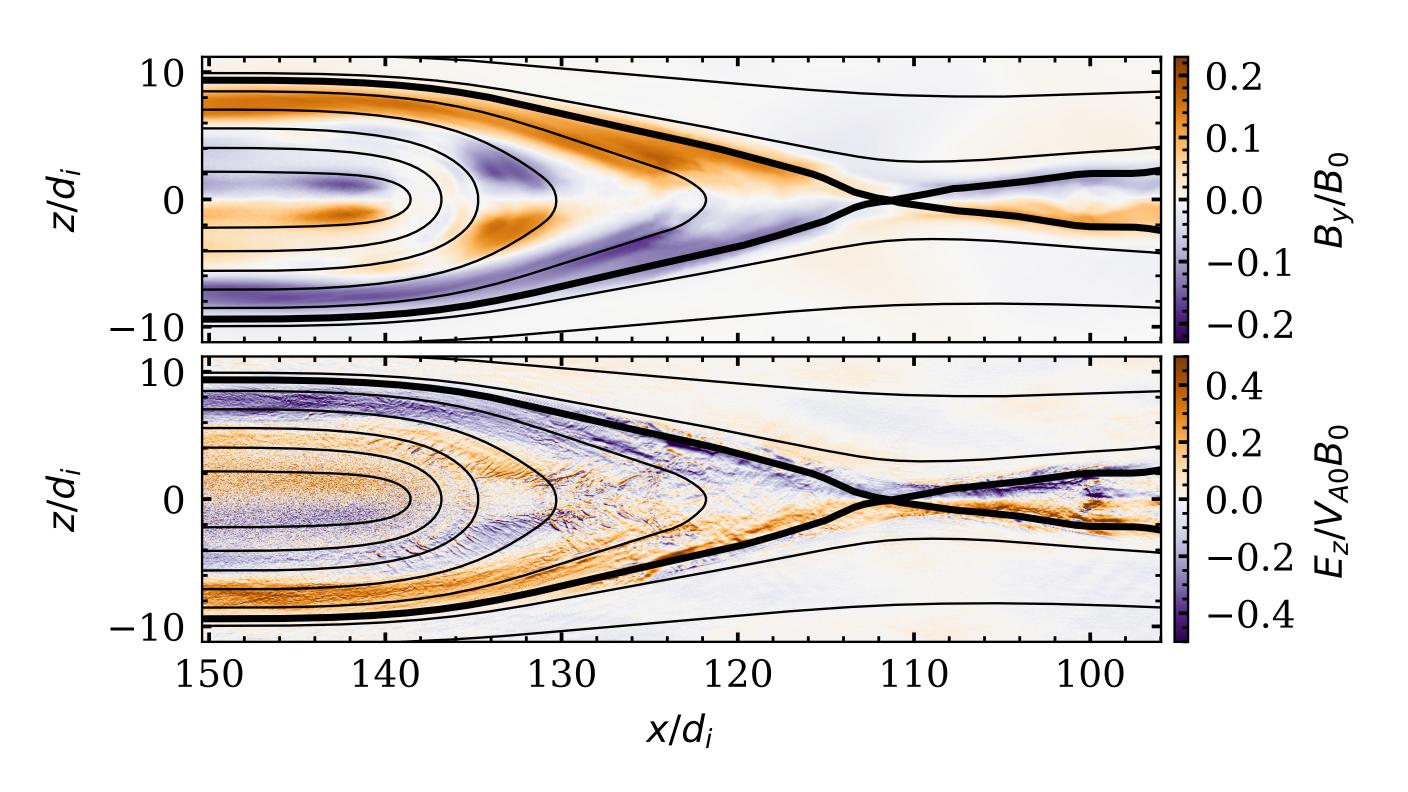
- Energy injection (Von-Kármán-Howarth energy decay rate) is **balanced** by energy transfer across large scales.
- Rate of decrease of magnetic energy in the reconnection inflow

$$\dot{\mathcal{E}}_b = (B_r^2/2\mu_0)/\Delta t \sim 10^9 \text{ J kg}^{-1} \text{ s}^{-1}$$

Turbulent energy transfer rate is ~10% of the reconnection energy input rate.

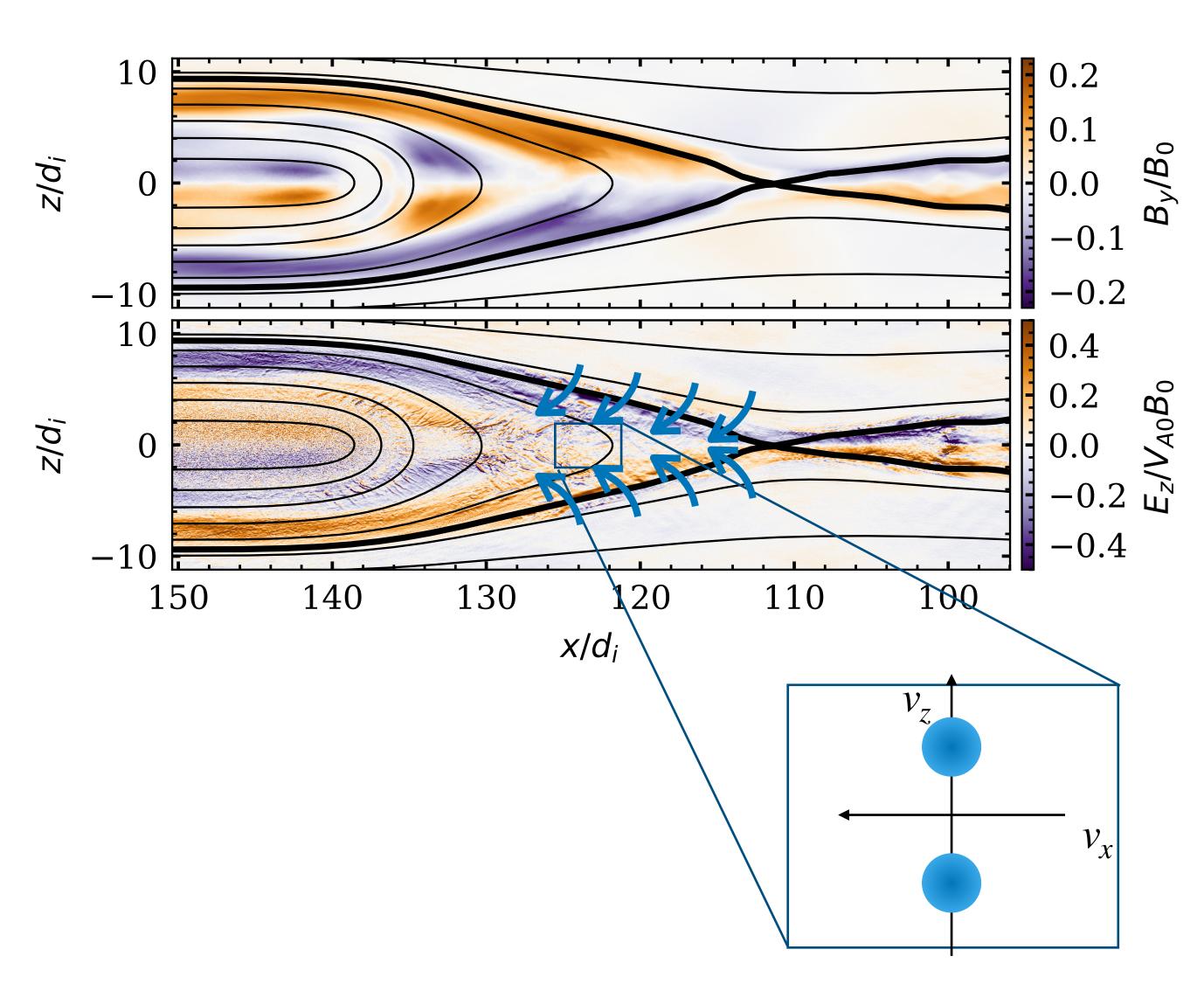


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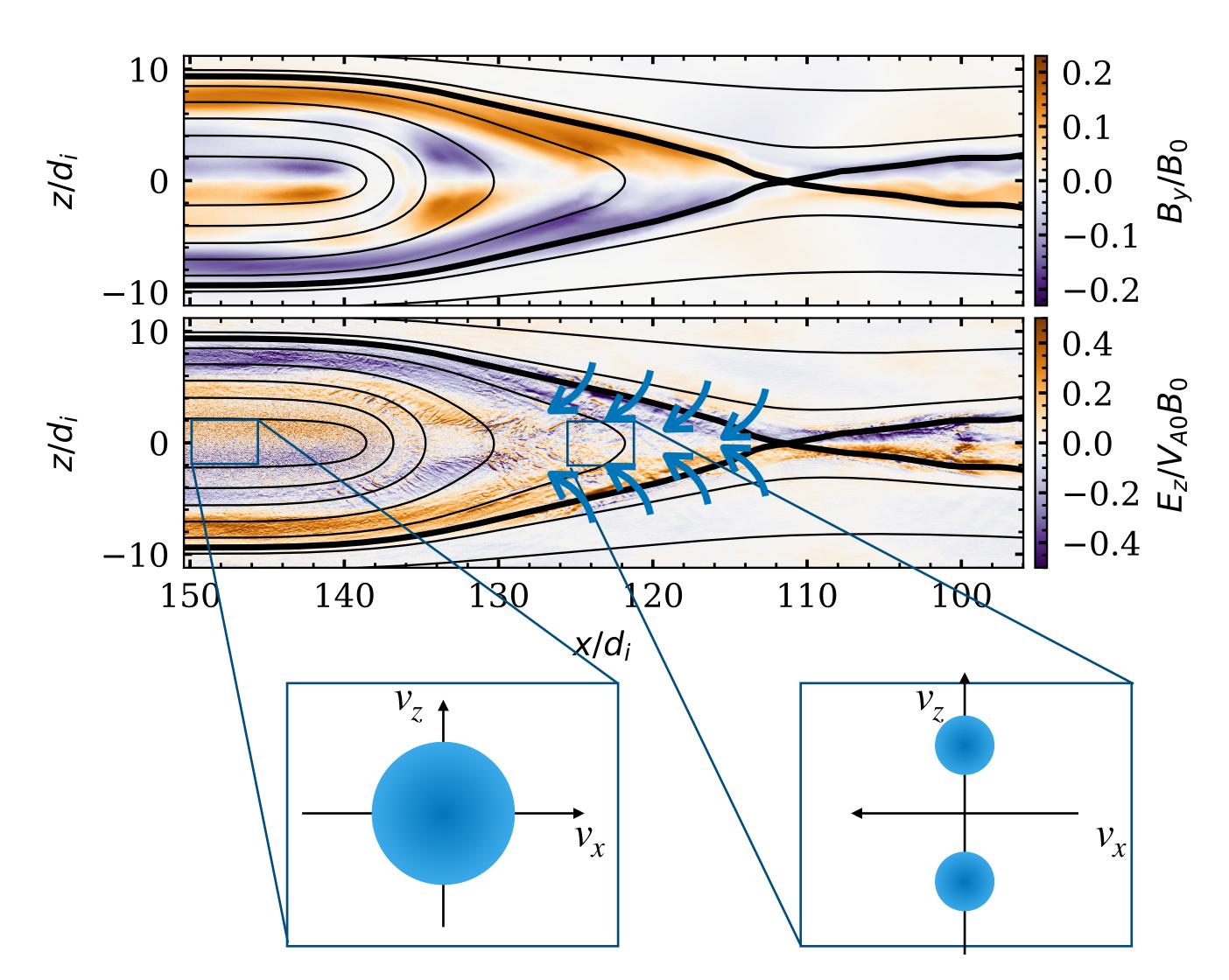
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 lons in the reconnection jet are accelerated by the Hall electric field, forming counter-streaming beams [e.g., Wygant+2005, JGR].



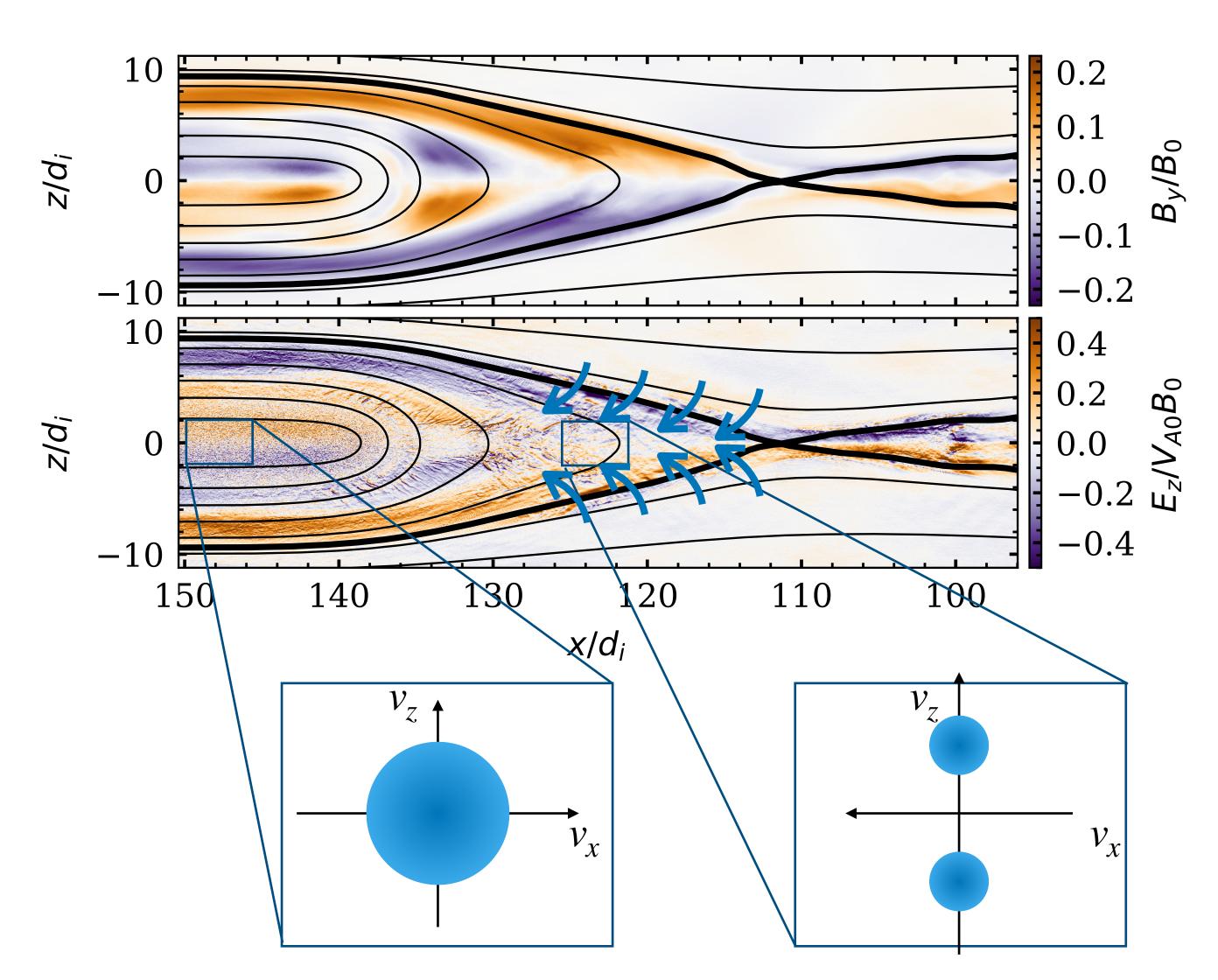
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- Temperature anisotropy decreases with distance from the reconnection region [Wu+2013, JGR].



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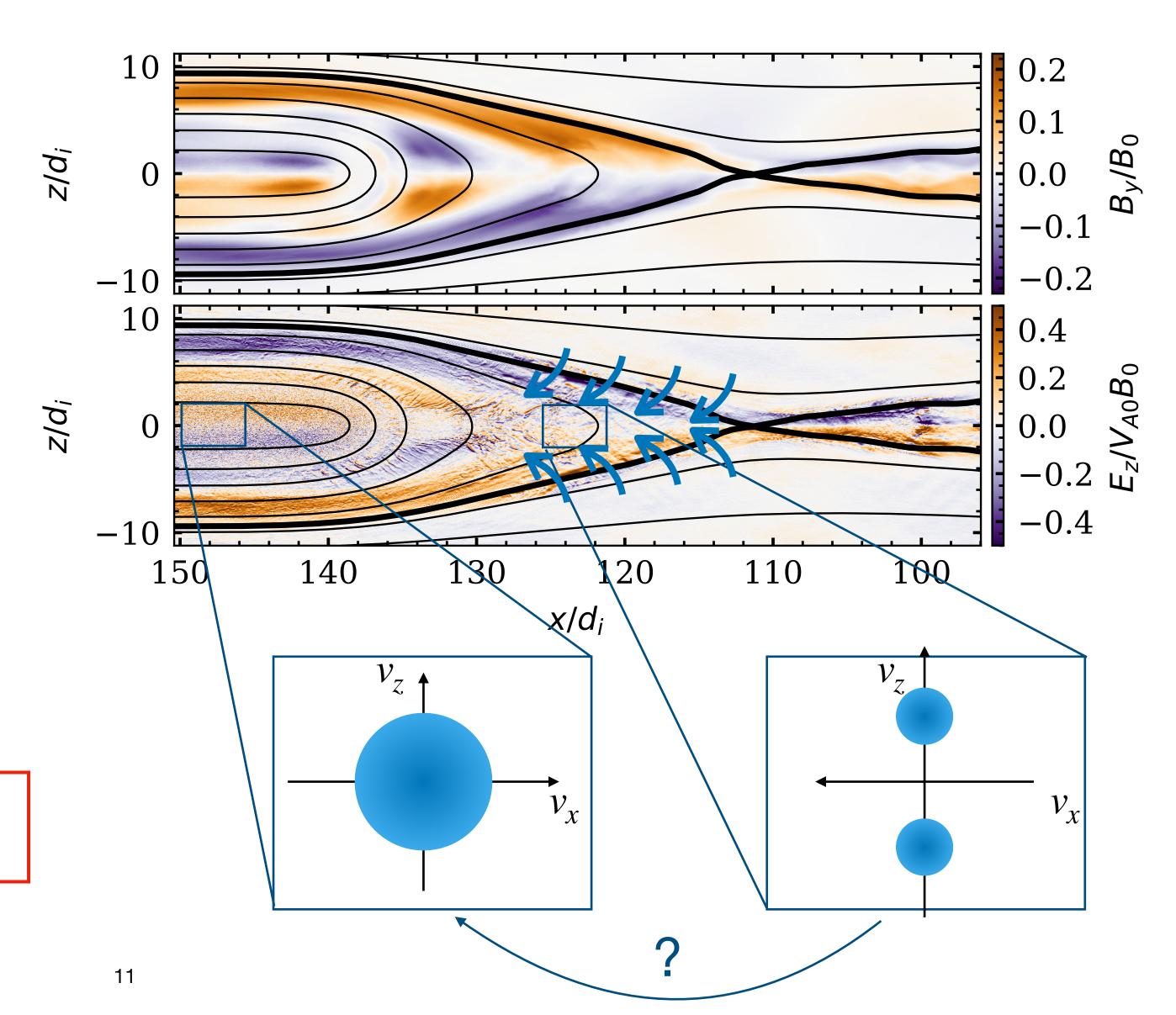
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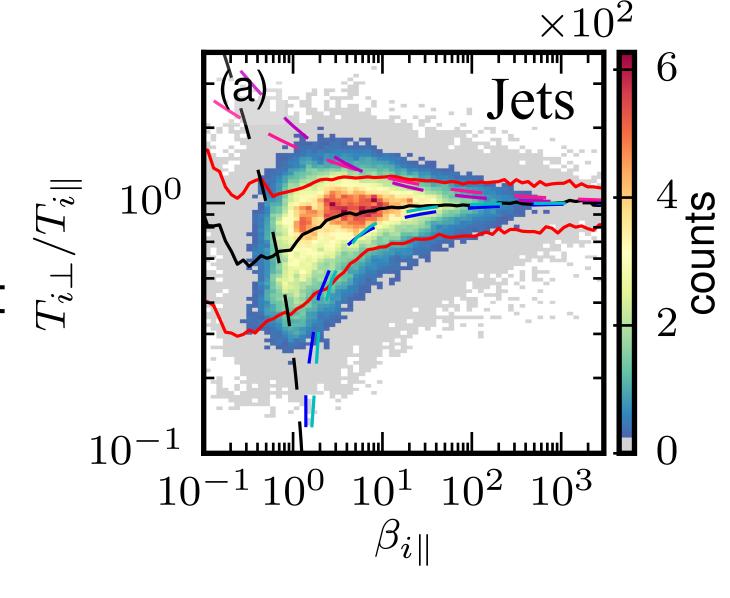
How are ion VDFs isotropized in the reconnection jet?



Data Distribution

Data Distribution

• >70% of the measurements are within the instability thresholds, indicating that the plasma conditions are efficiently regulated.

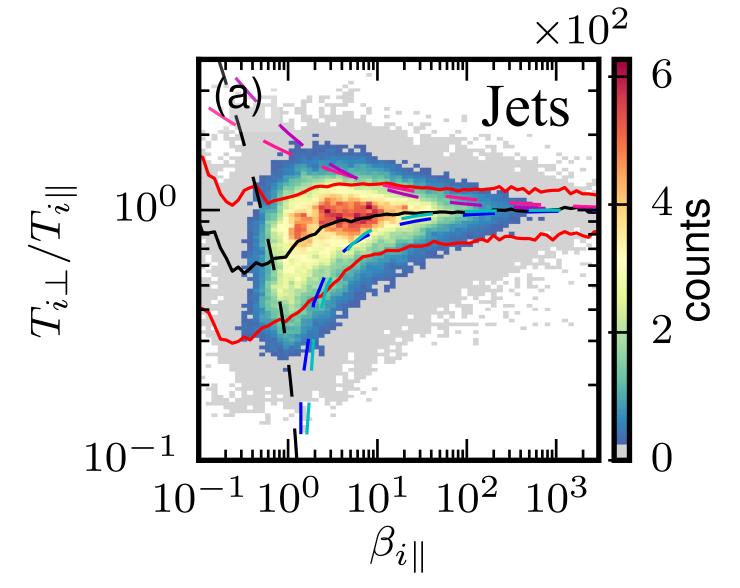


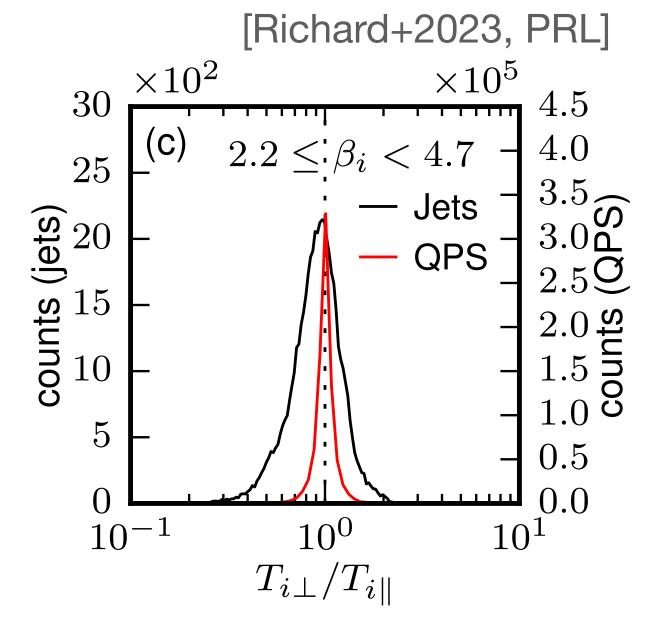
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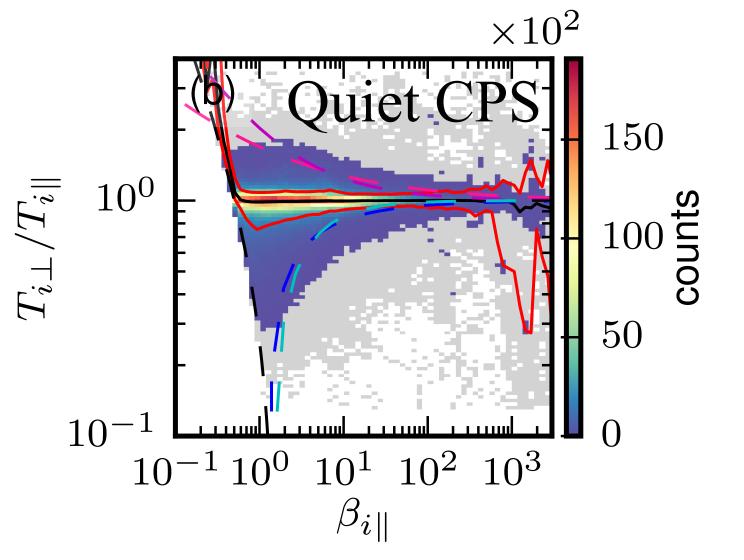
Data Distribution

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Magnetic reconnection drives ion VDFs out of LTE.

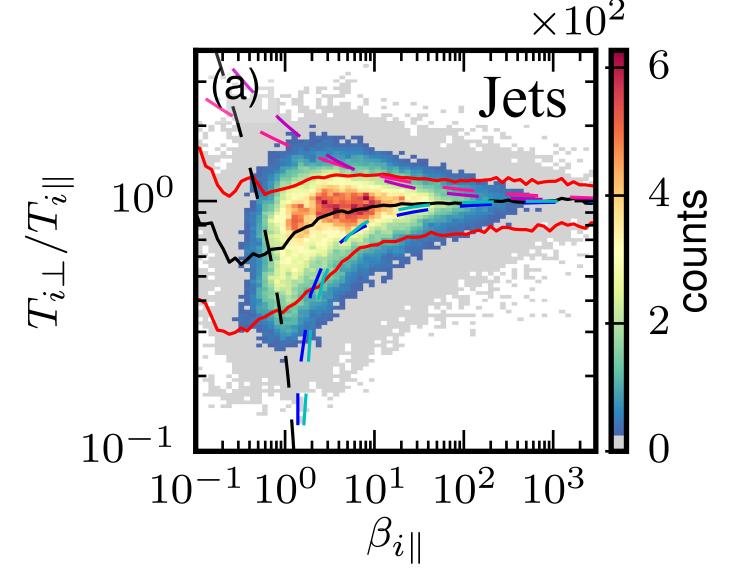


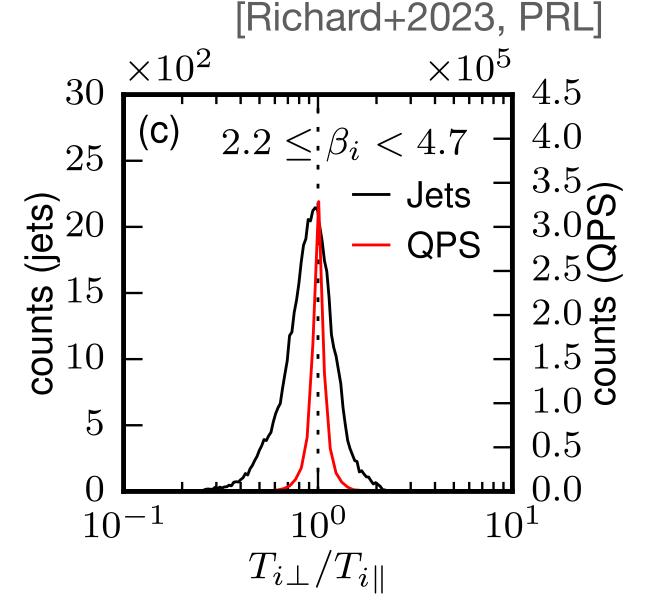




Data Distribution

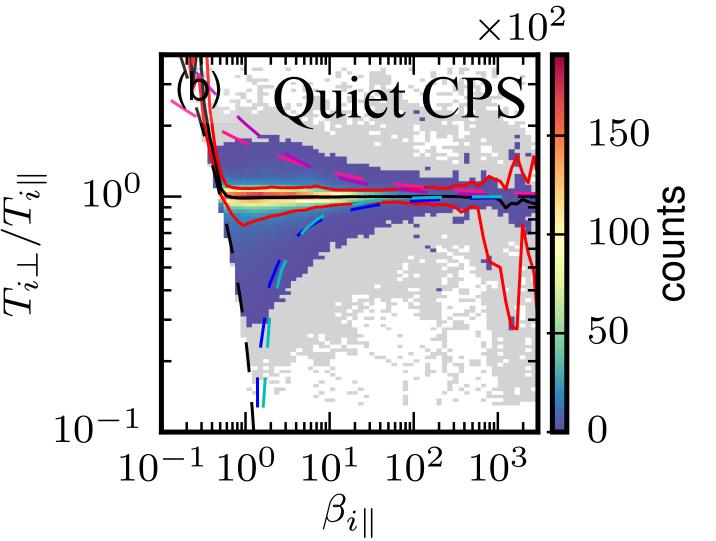
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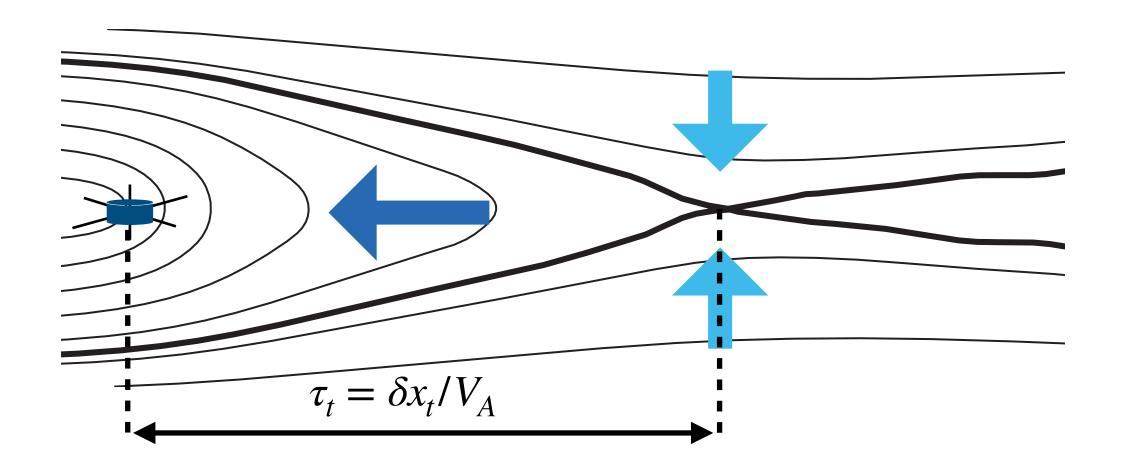
Ion VDFs are nearly isotropic in the quiet plasma sheet and in the reconnection jets.



Relaxation time

Relaxation time

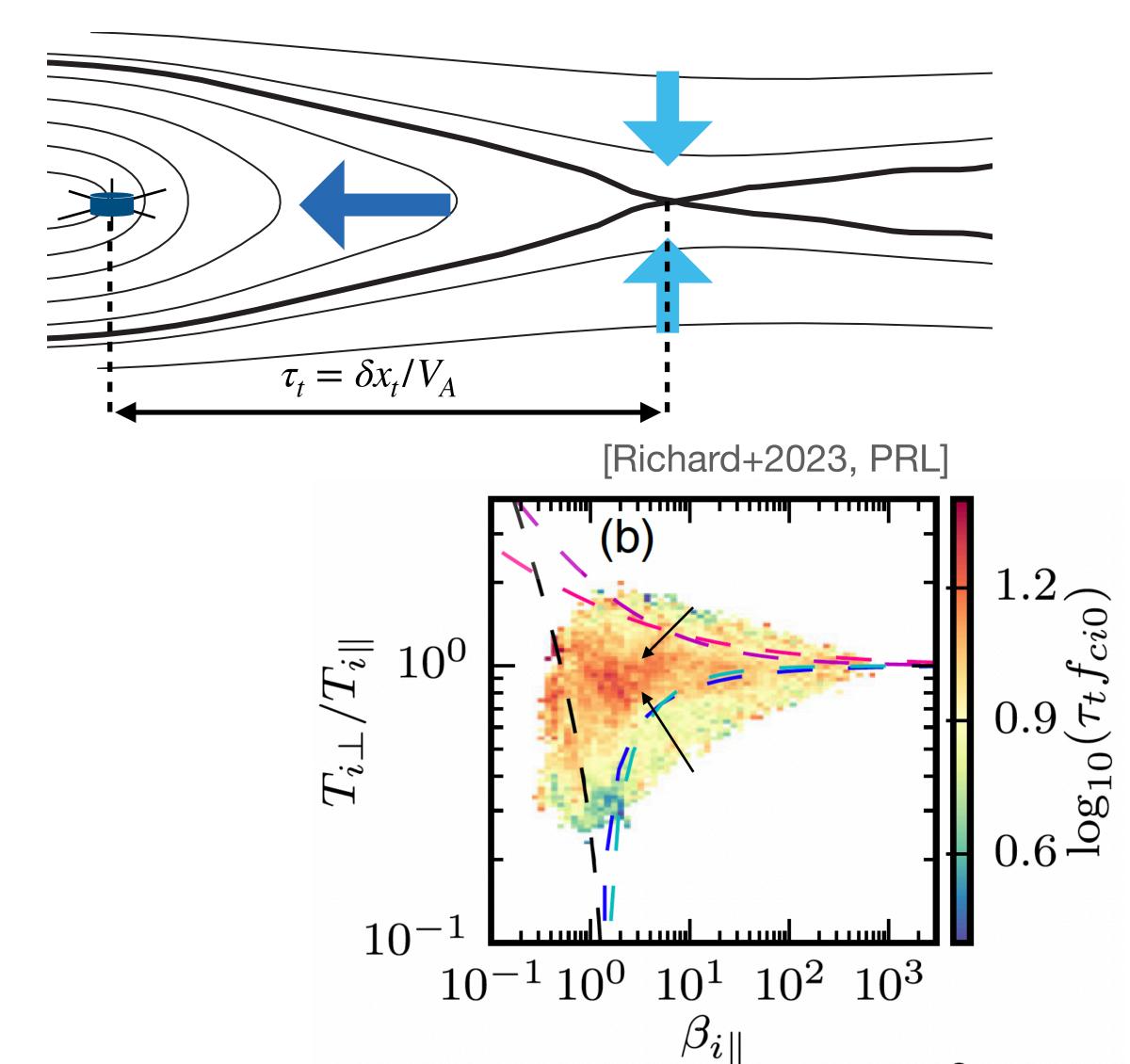
- How long does it take to isotropize the ion VDFs?
 - At most, the travel time from the X-line to the spacecraft.



Relaxation time

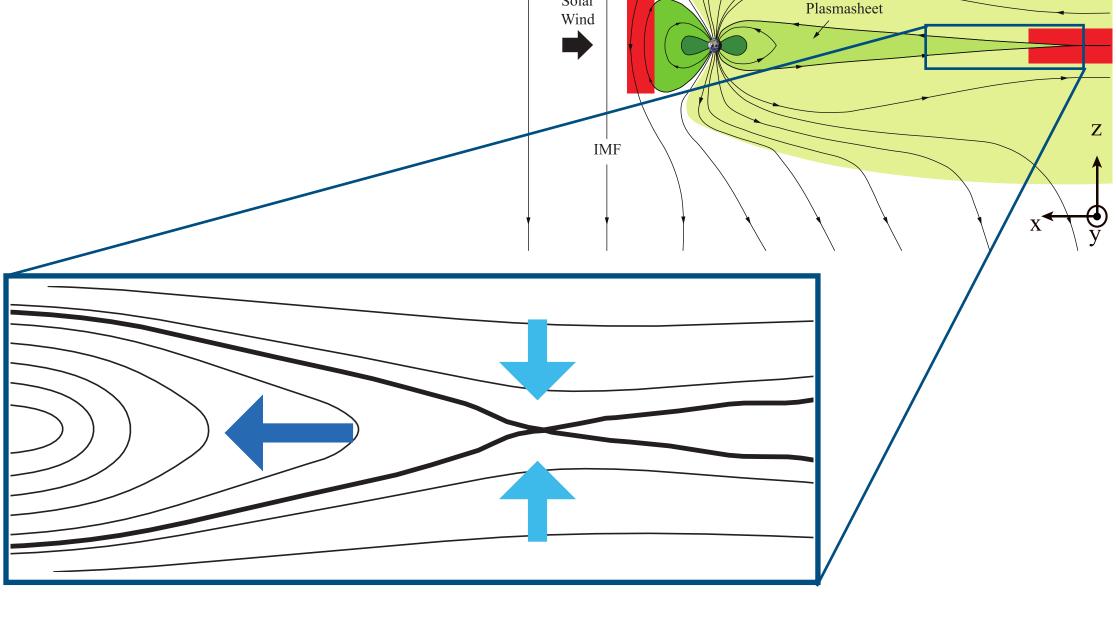
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 - At most, the travel time from the X-line to the spacecraft.

Non-LTE Ion VDFs are relaxed after $\tau_t f_{ci0} \sim 10$, i.e., $\delta x \sim 60 d_i$.



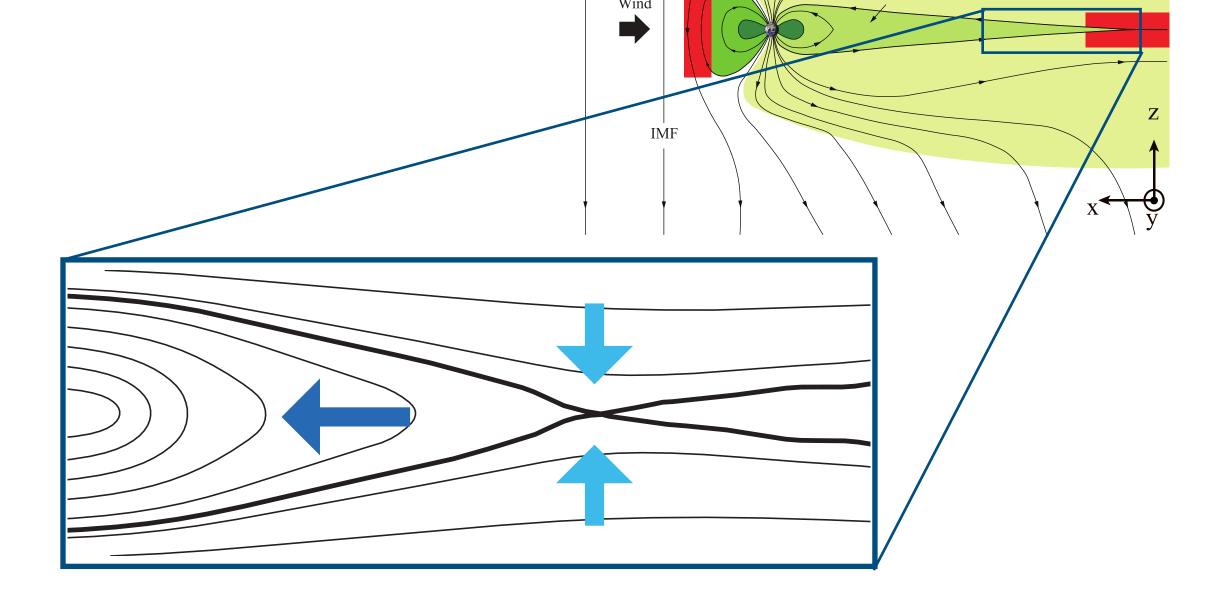
Current sheet scattering

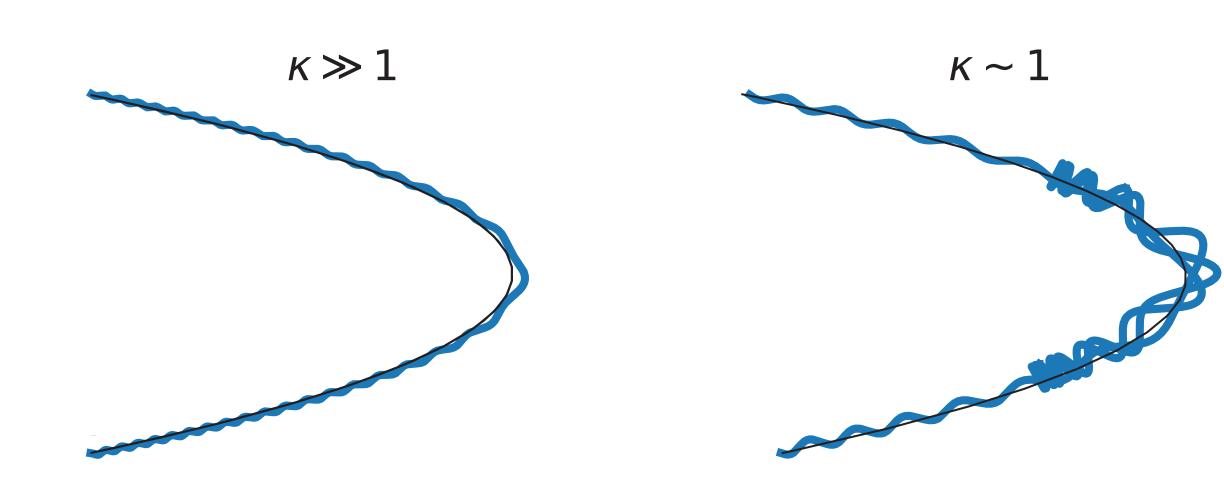
Current sheet scattering



Current sheet scattering

- Ion motion in current sheets are identified based on $\kappa = \sqrt{r_c/\rho_i}$ [Büchner and Zelenyi, 1989, JGR]:
 - Adiabatic motion $\kappa \gg 1$
 - Chaotic motion $\kappa \approx 1$ -> strong scattering
 - Quasi-adiabatic motion $\kappa \ll 1$ -> weak scattering

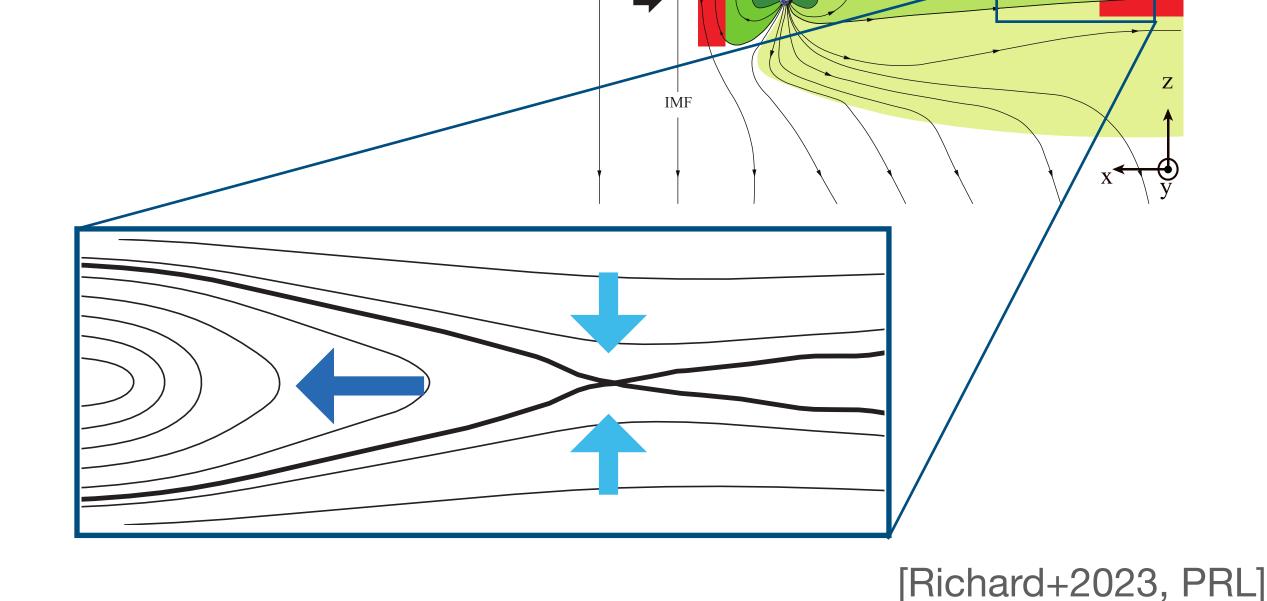


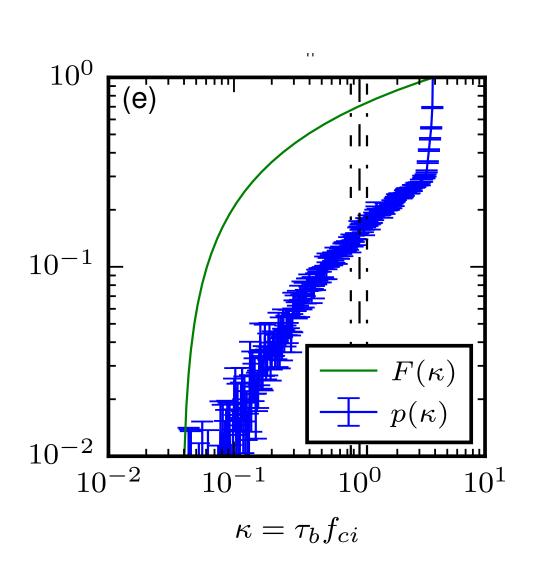


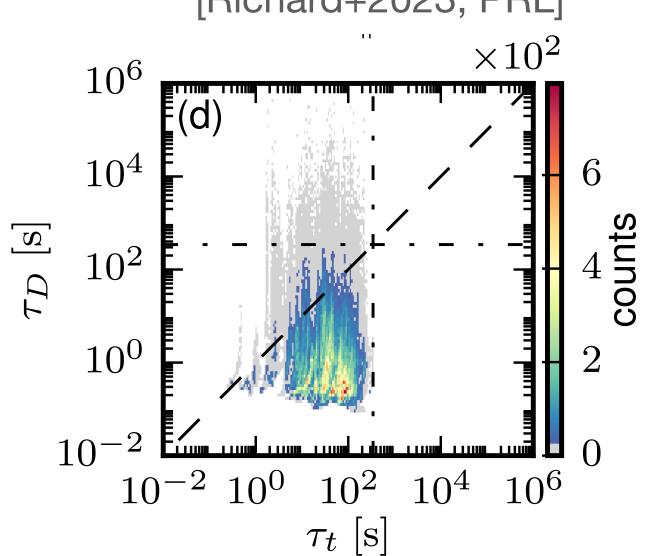
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- 73% of ion VDFs in reconnection jets show chaotic or quasi-adiabatic motion.

Chaotic ion motion in the CS results in fast pitch-angle scattering







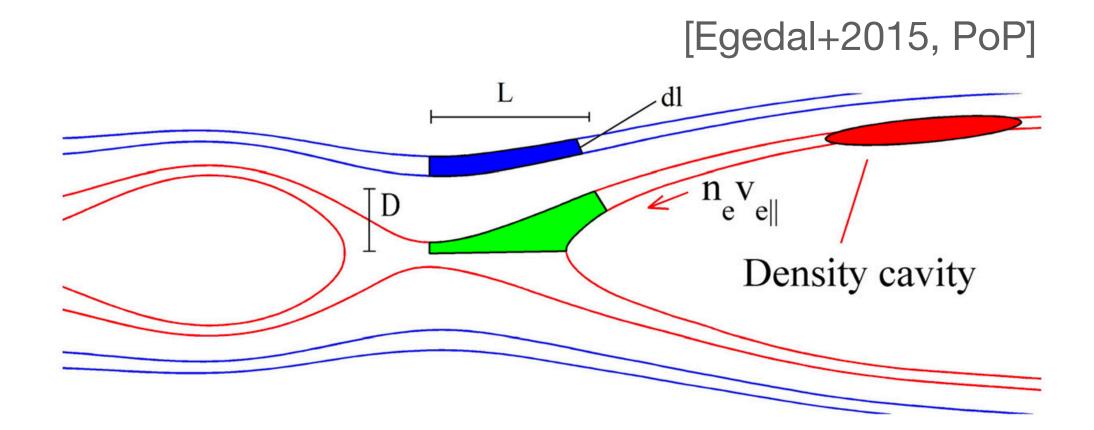
Parallel Electric Fields in Reconnection

Theory

Parallel Electric Fields in Reconnection

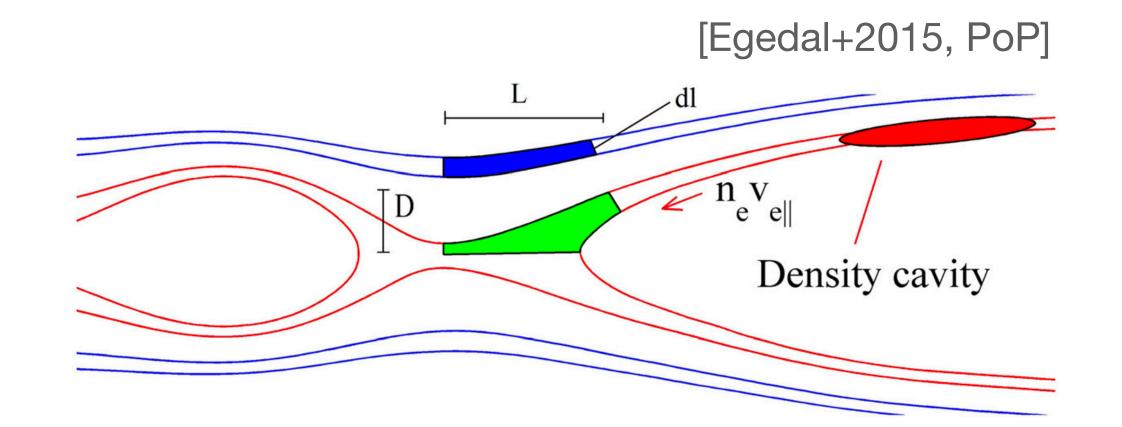
Theory

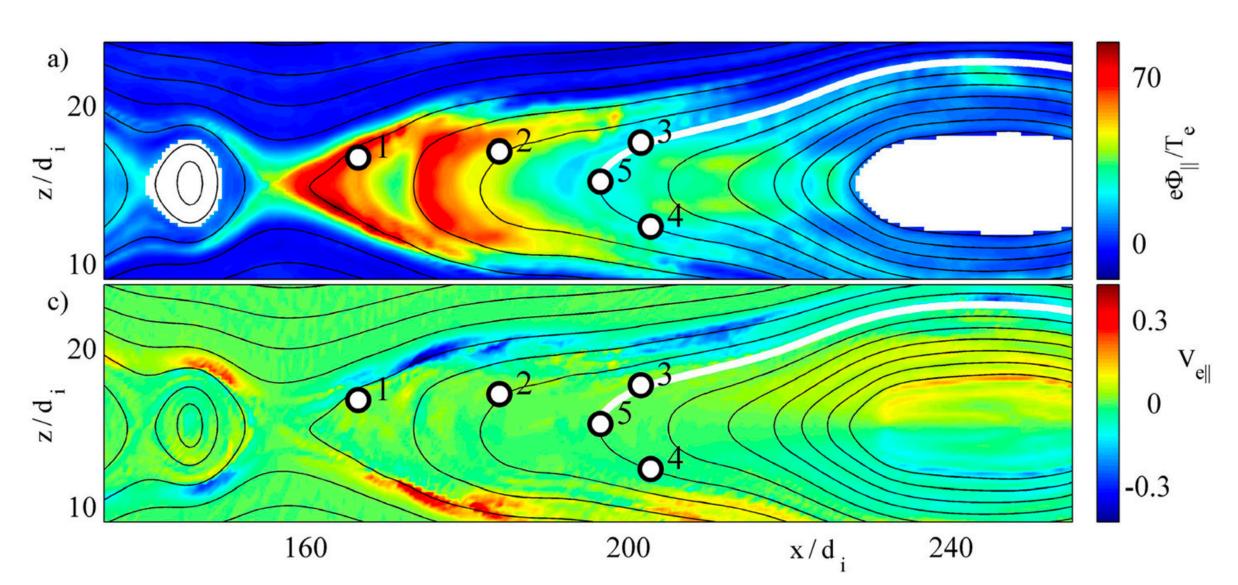
 As magnetic flux tubes expand, electron density should drops, while density of demagnetized ions stays constant resulting in positive charge density.



Parallel Electric Fields in Reconnection Theory

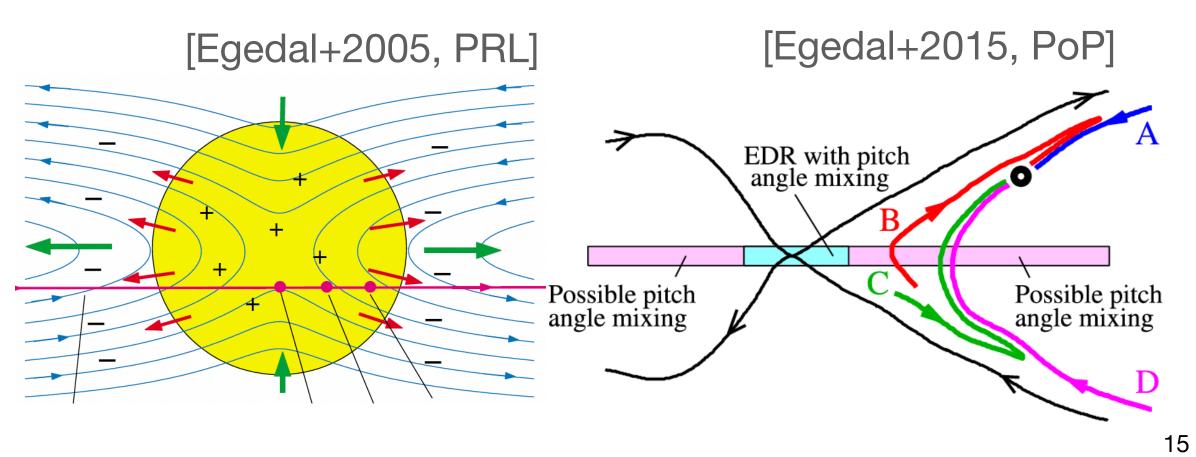
- As magnetic flux tubes expand, electron density should drops, while density of demagnetized ions stays constant resulting in positive charge density.
- Parallel electric fields E_{\parallel} form in the reconnection region to maintain quasi-neutrality.

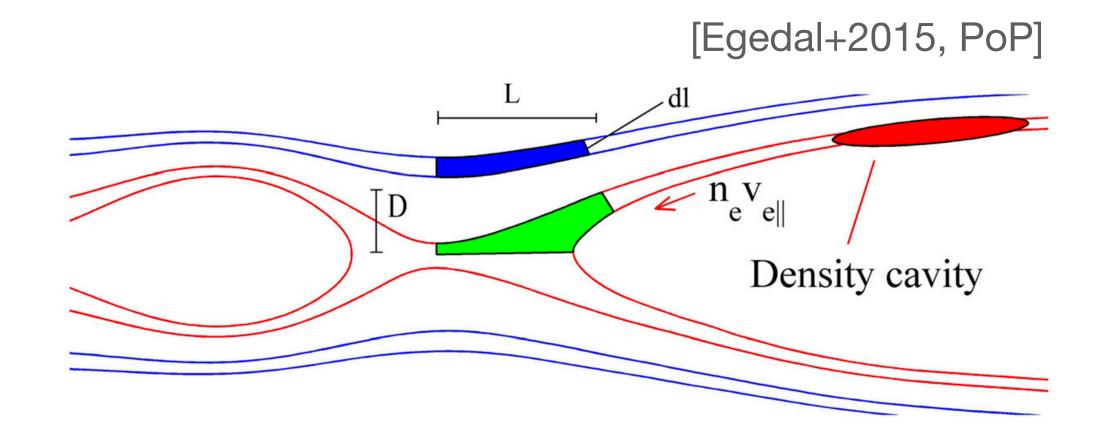


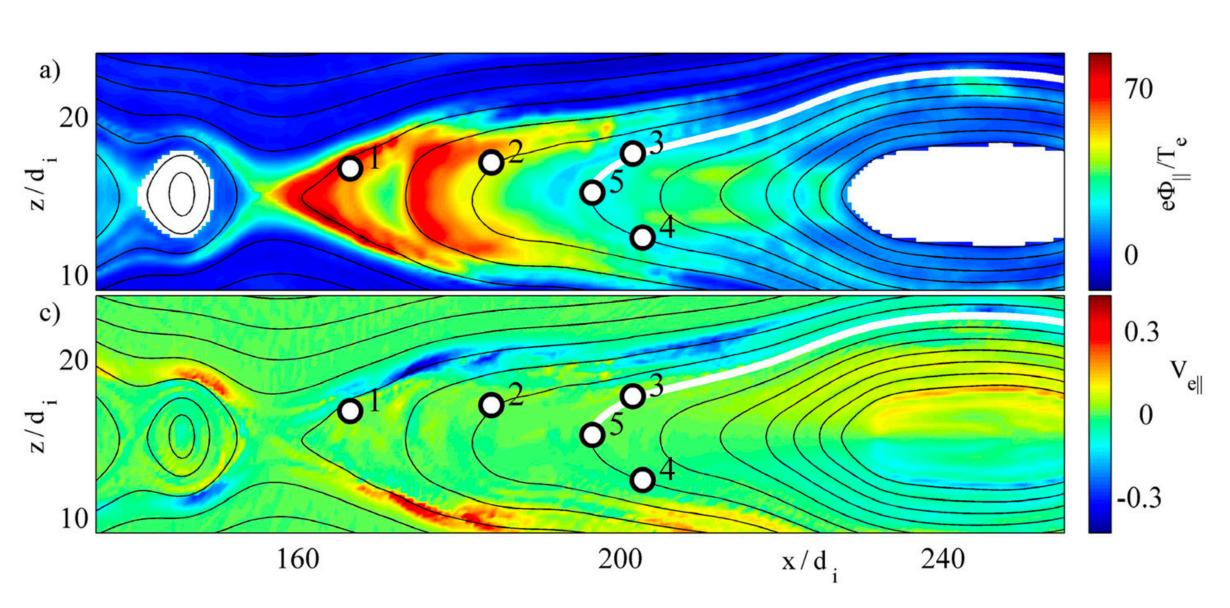


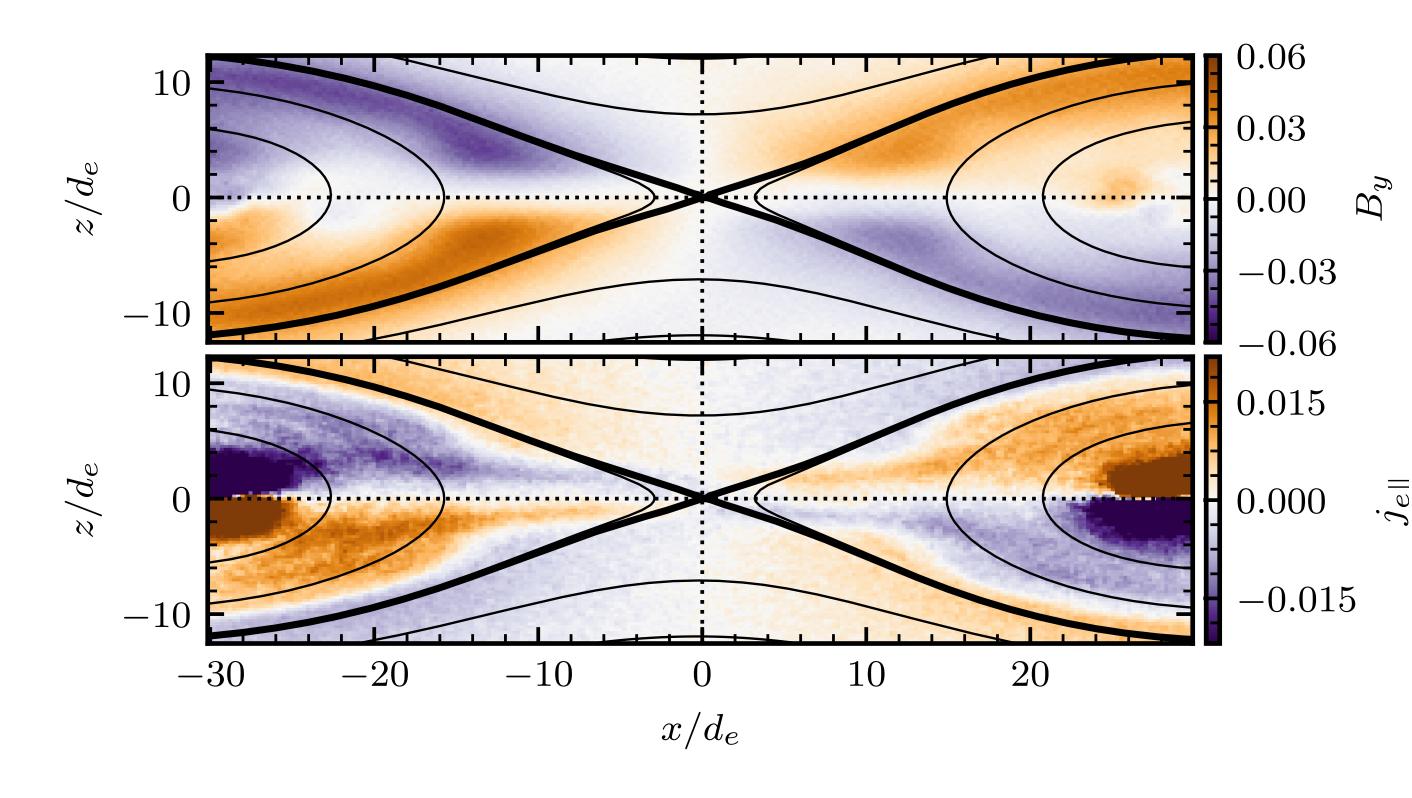
Parallel Electric Fields in Reconnection Theory

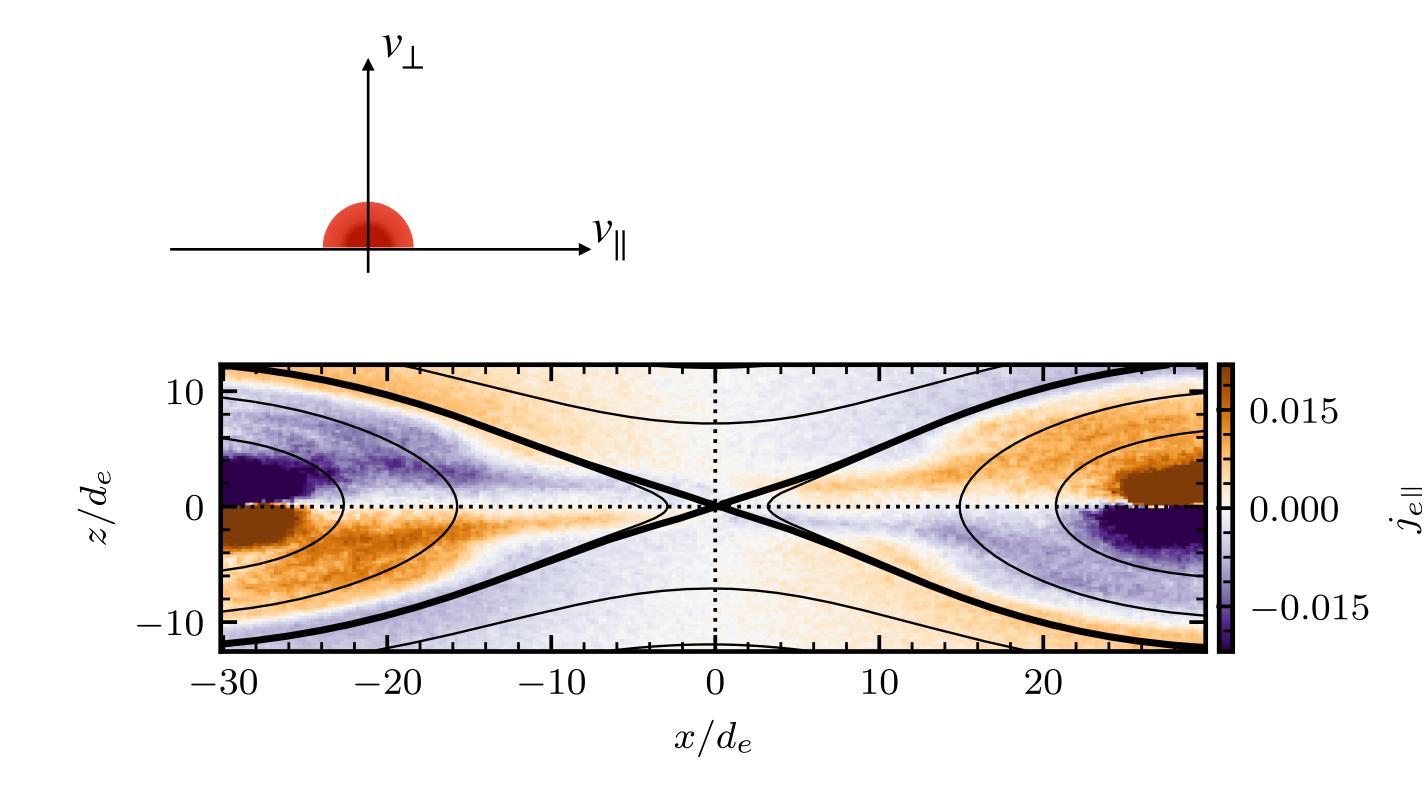
- As magnetic flux tubes expand, electron density should drops, while density of demagnetized ions stays constant resulting in positive charge density.
- Parallel electric fields E_{\parallel} form in the reconnection region to maintain quasi-neutrality.
- E_{\parallel} trap electrons locally to balance charge and maintain quasi-neutrality.



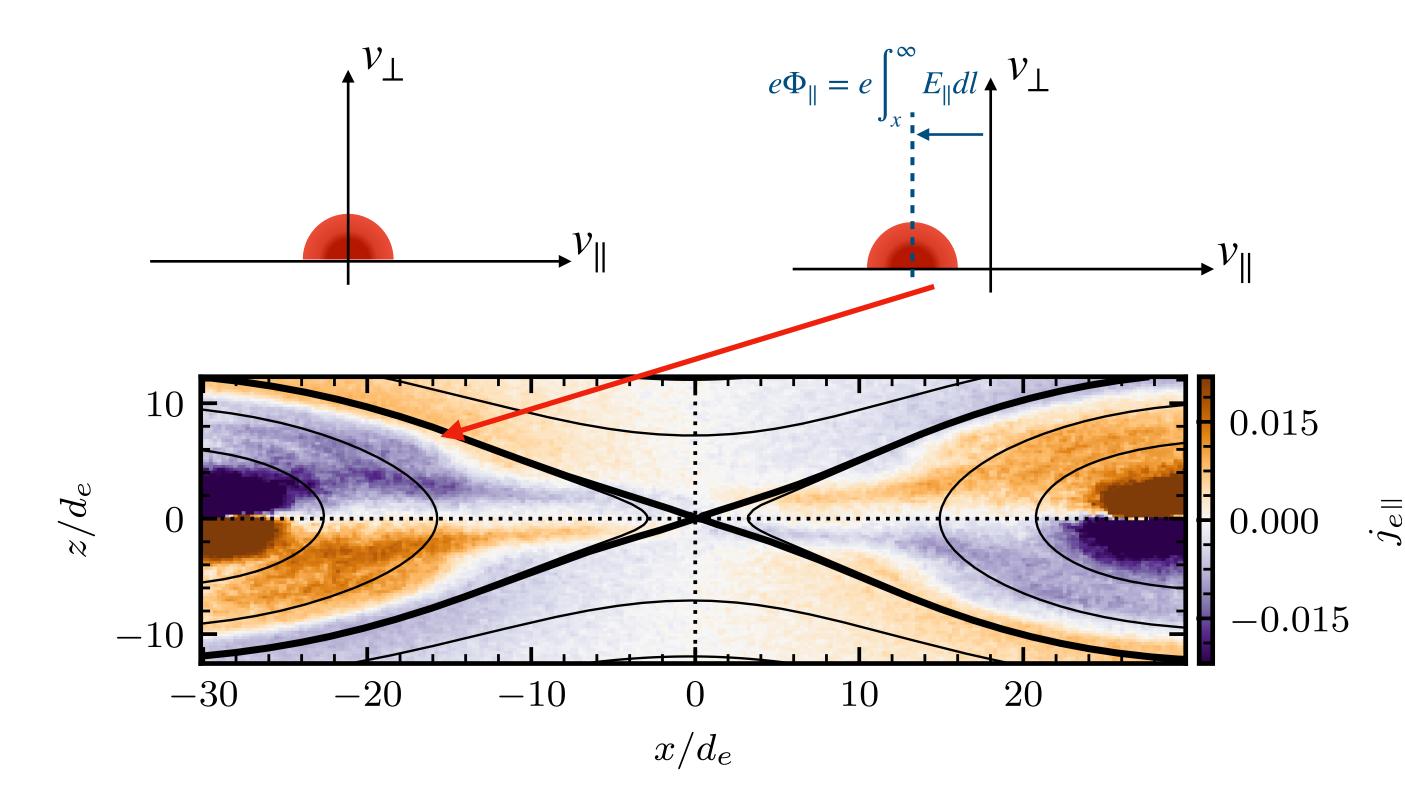




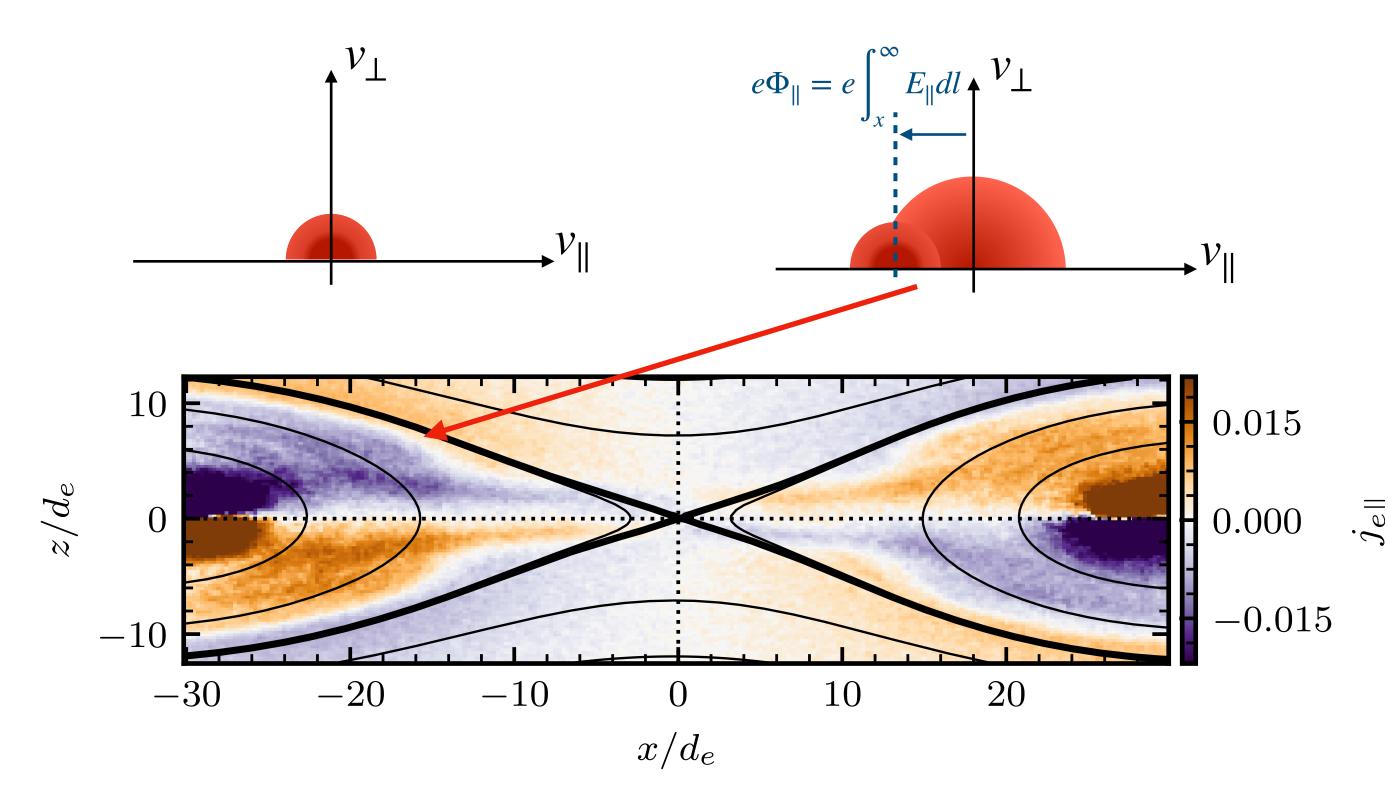




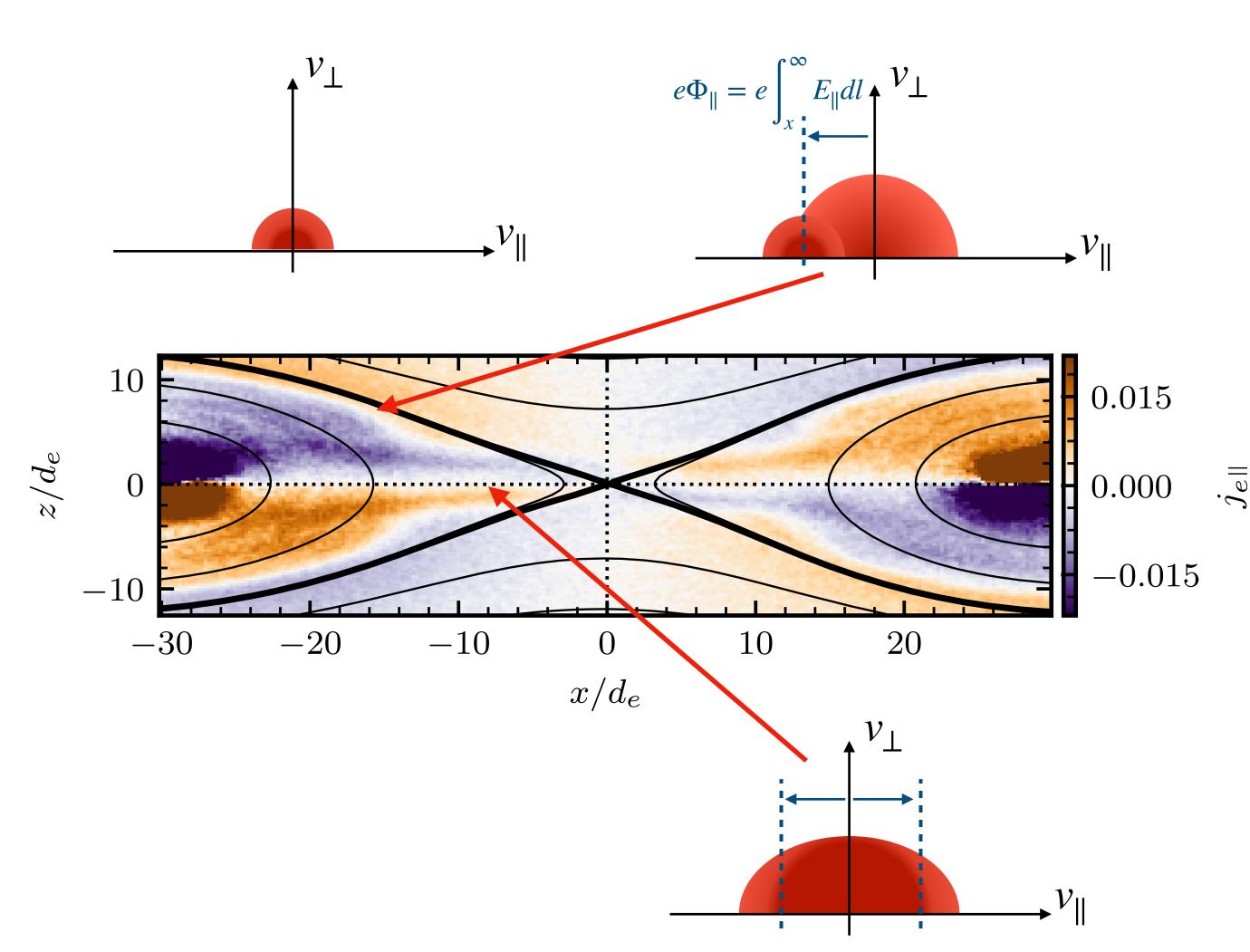
 Parallel electric fields accelerate electrons toward the X-line, forming beam-type electron VDFs.



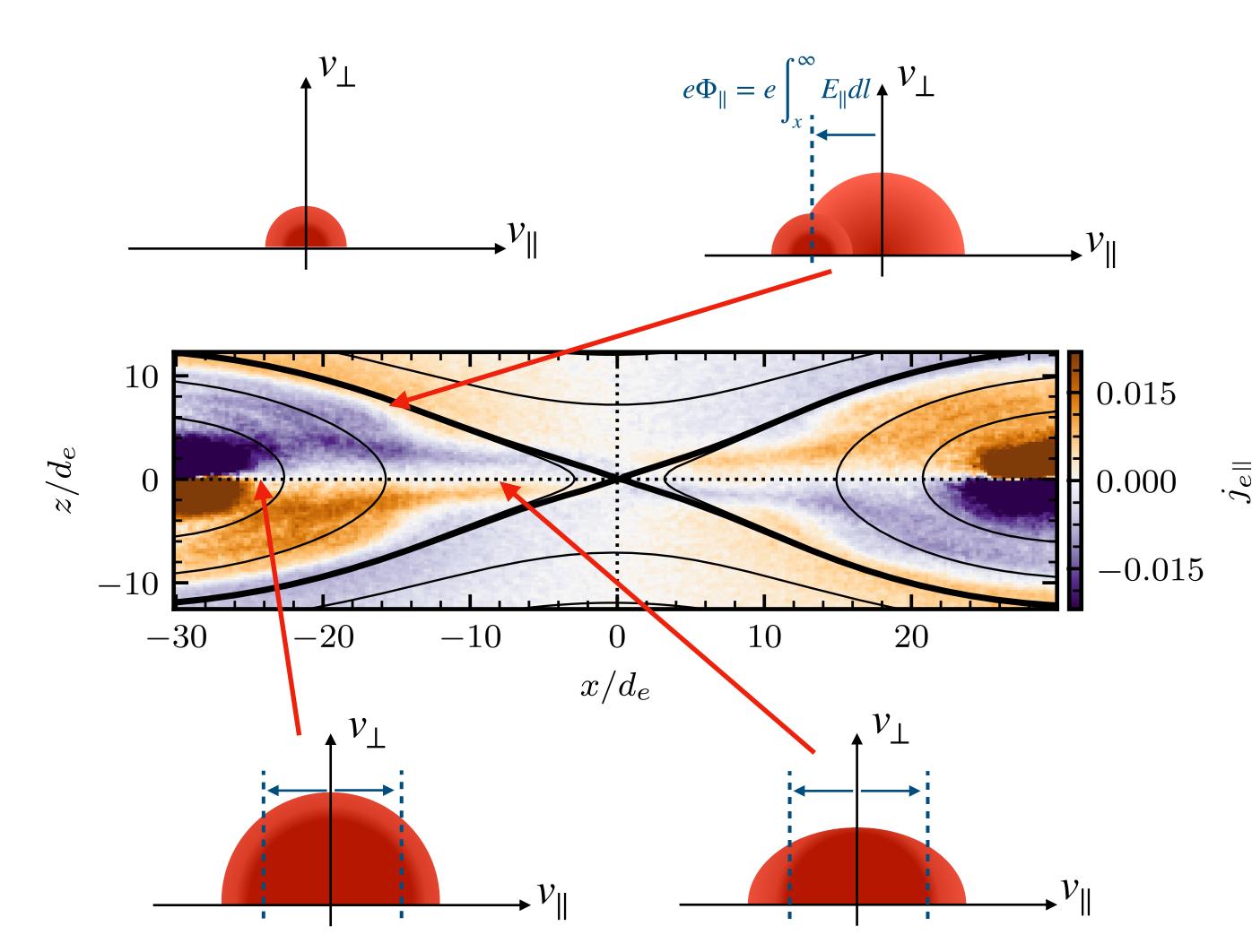
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- Unstable beam + thermalized hot electron excites electrostatic waves and forming flat-top VDFs [Fujimoto, 2014, GRL].

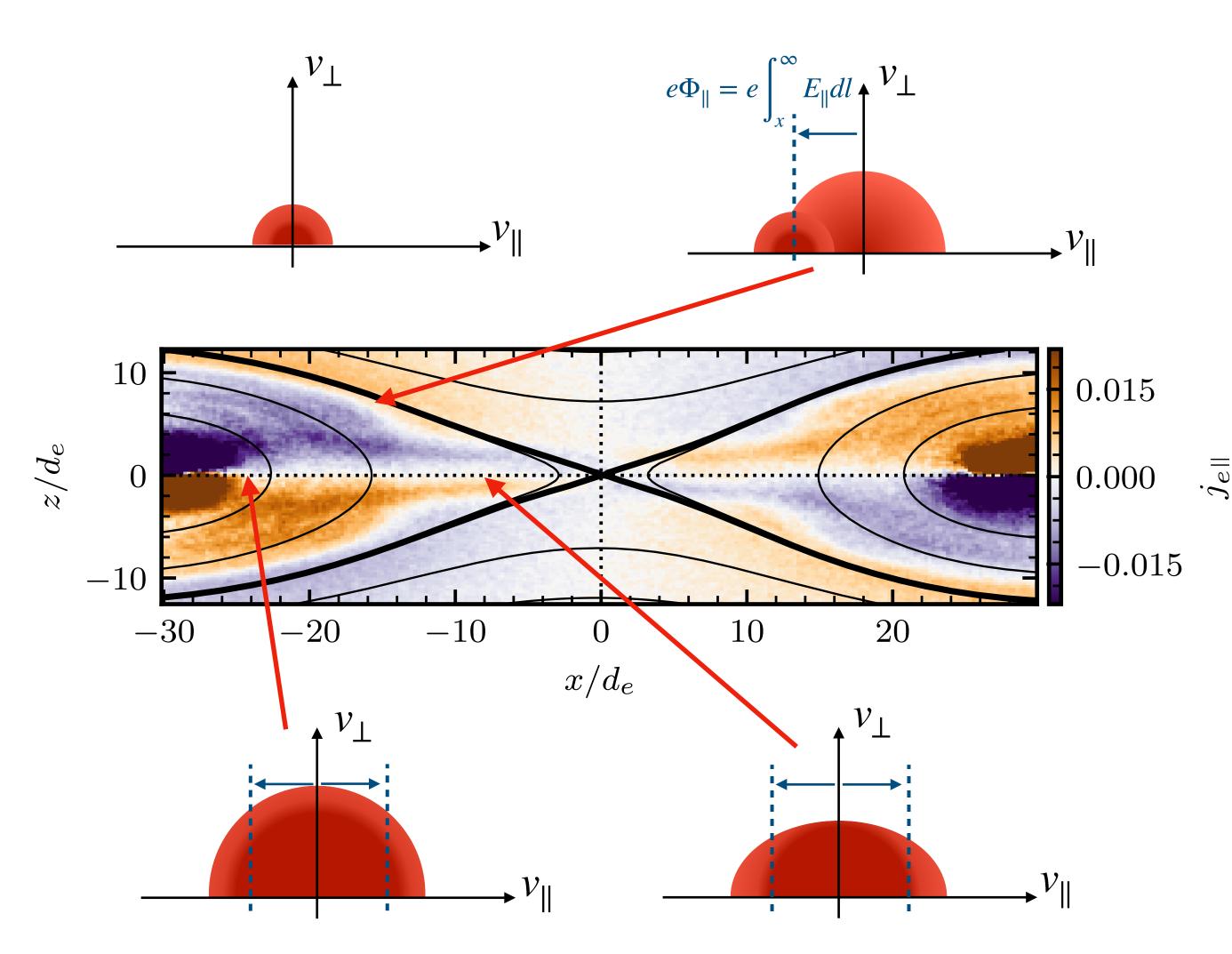


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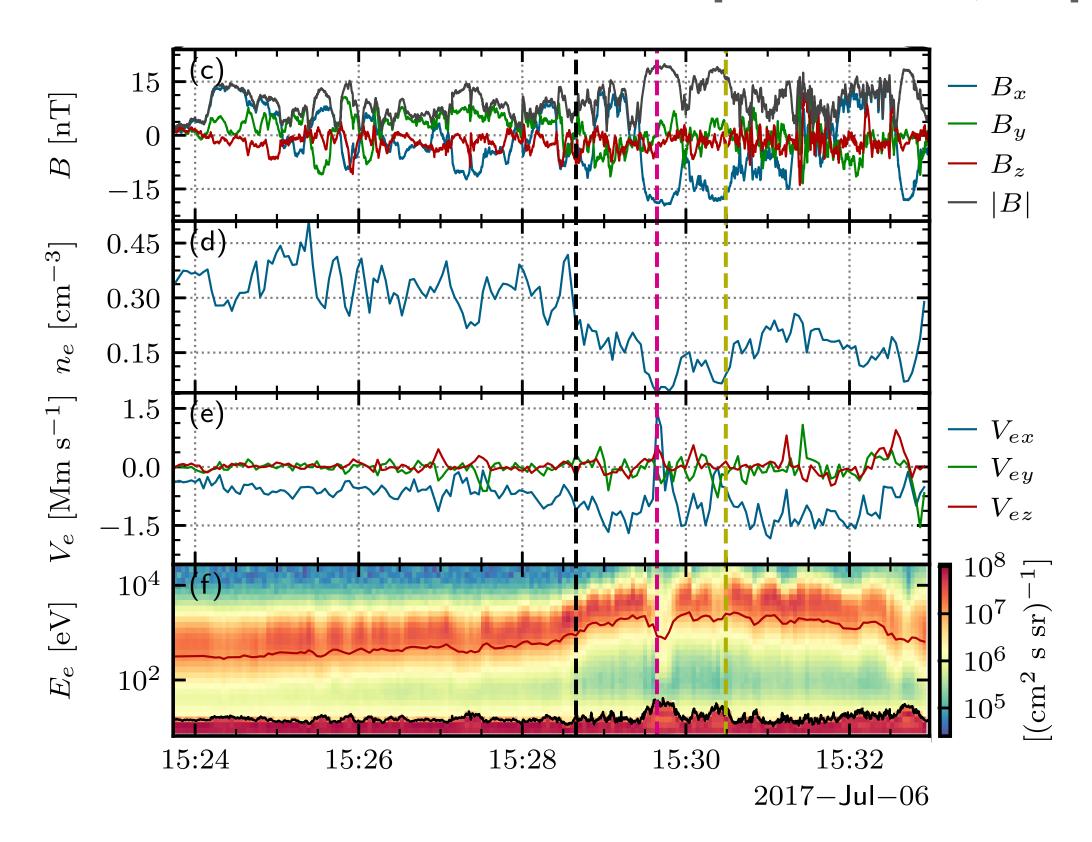
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What is the nature and the role of parallel electric fields in magnetic reconnection?



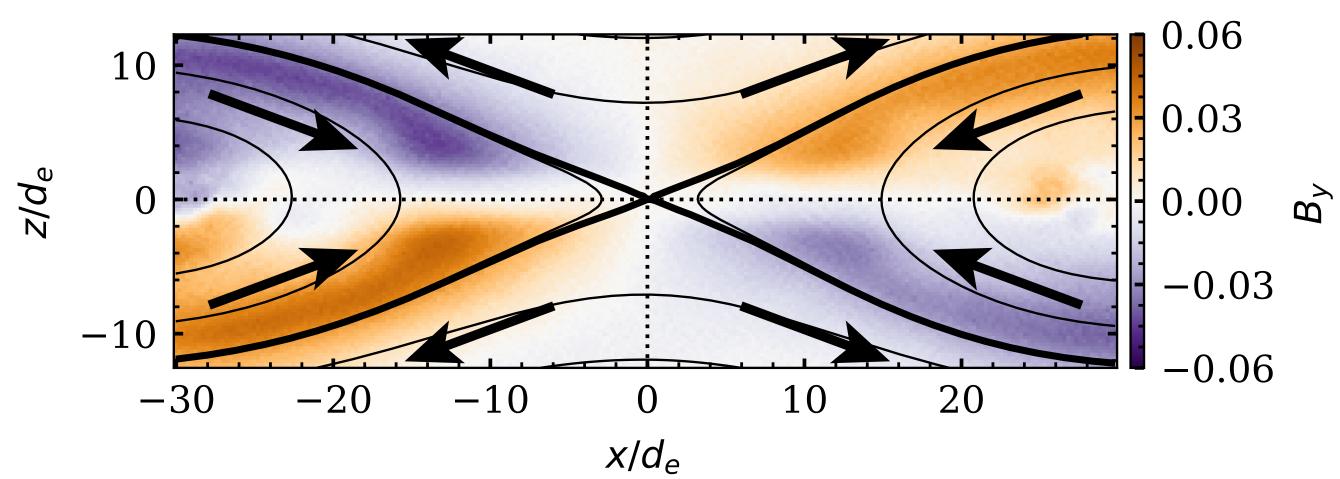
[Richard+2025, PRL]

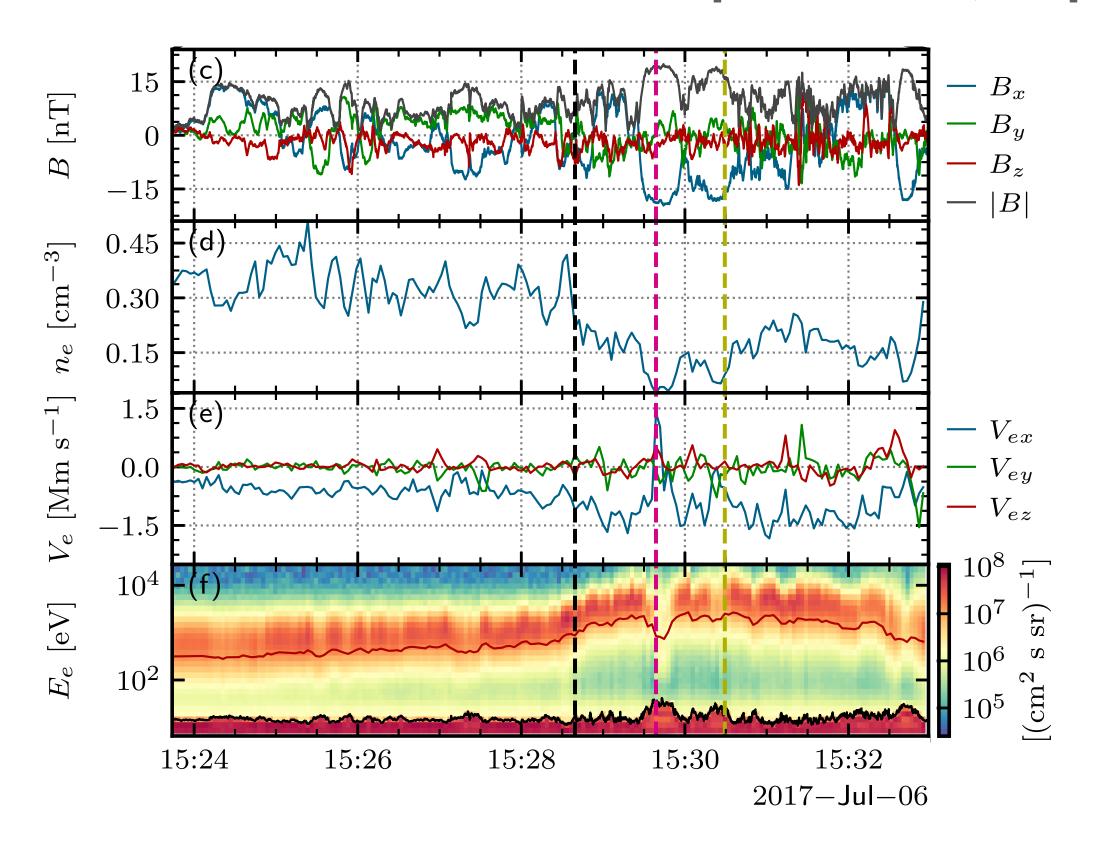
Example



[Richard+2025, PRL]

Example

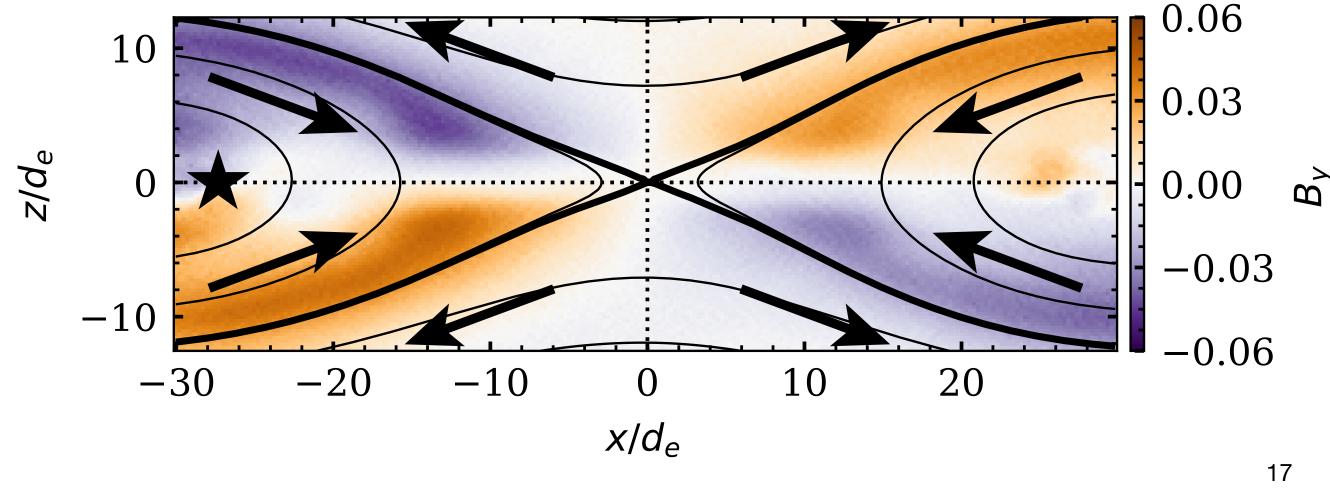


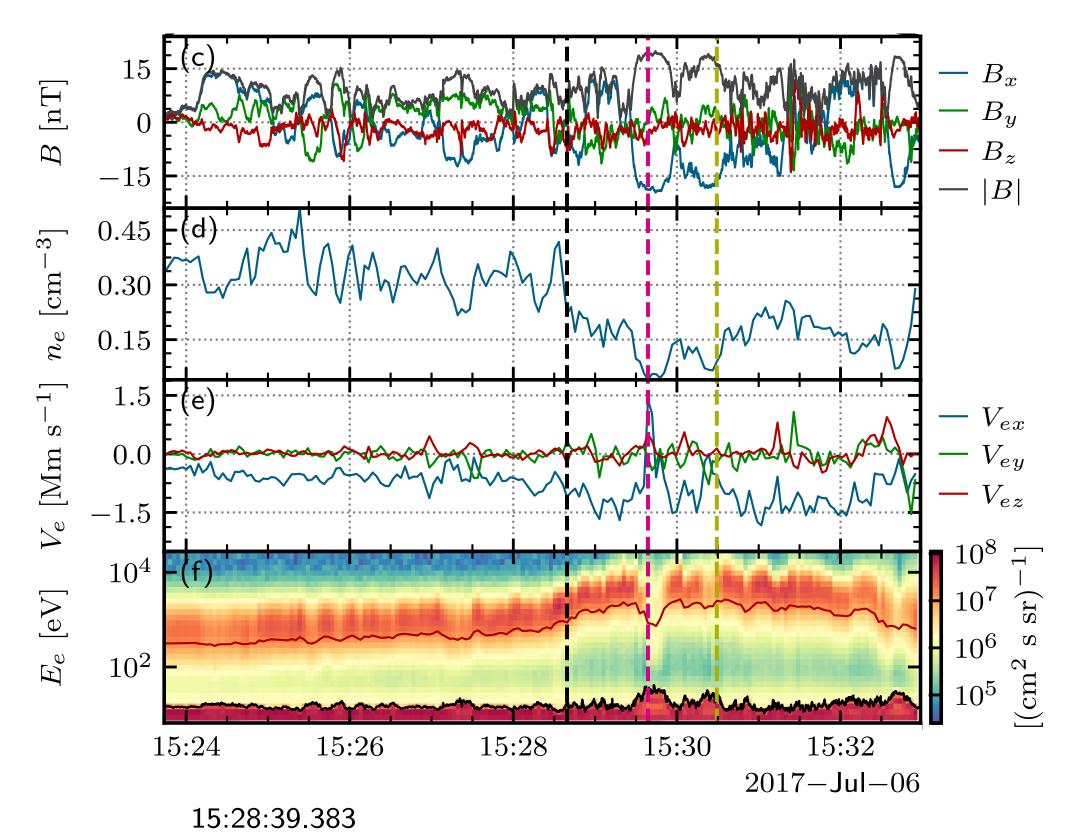


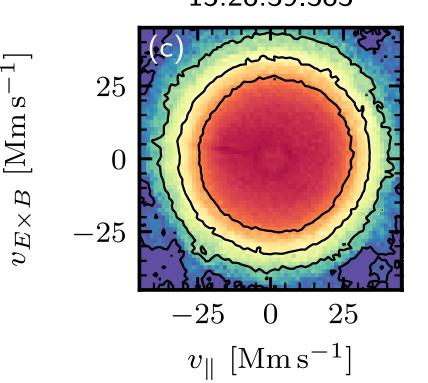
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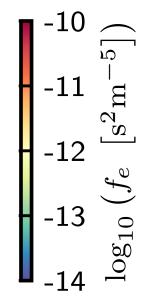
Example

Flat-top eVDF in the reconnection jet





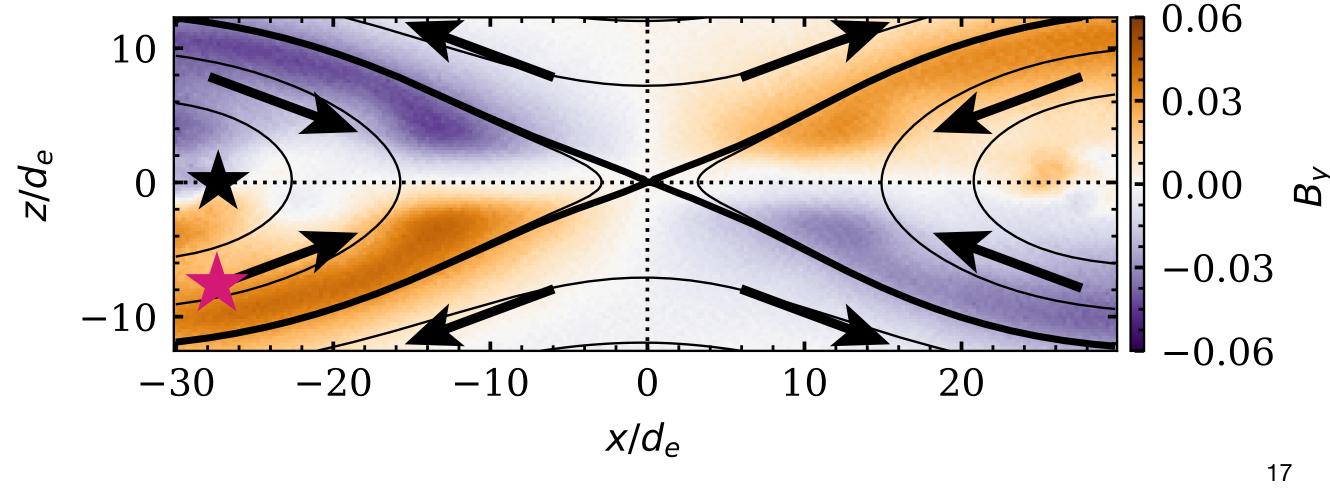


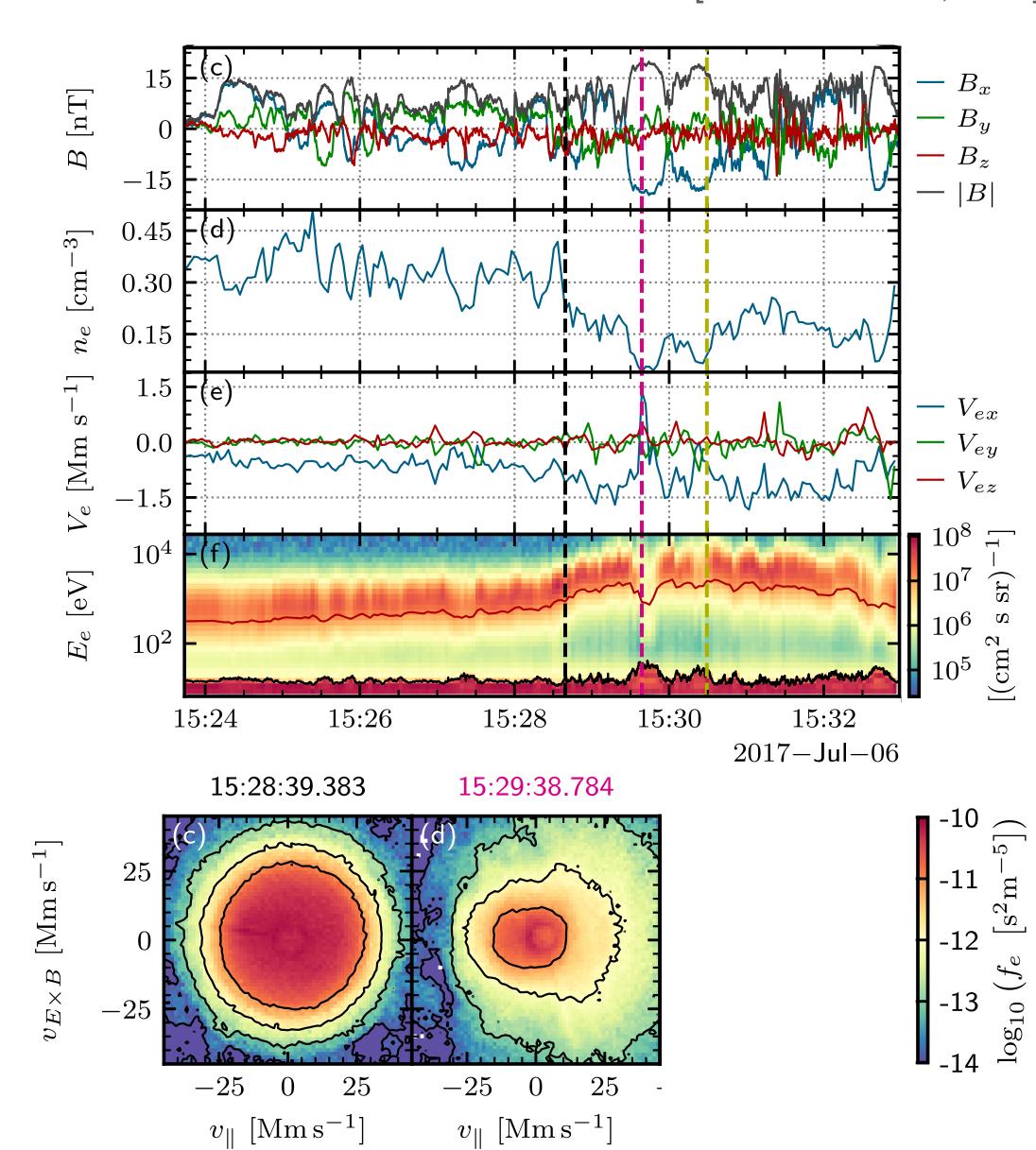


[Richard+2025, PRL]

Example

- Flat-top eVDF in the reconnection jet
- Beam + thermalized core at the separatrix flow (producing the Hall current)



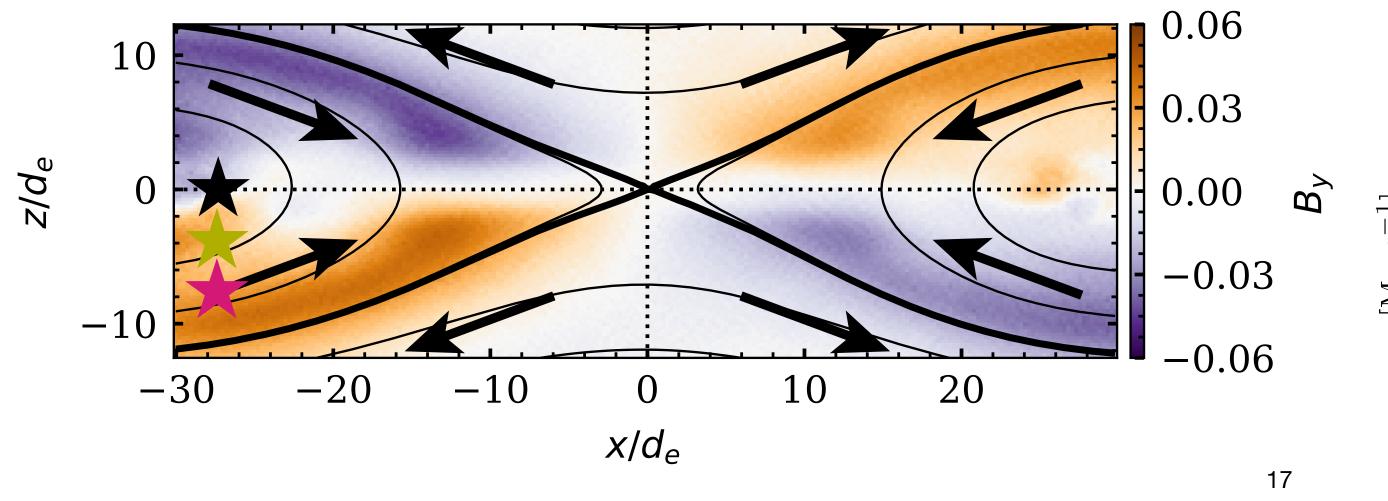


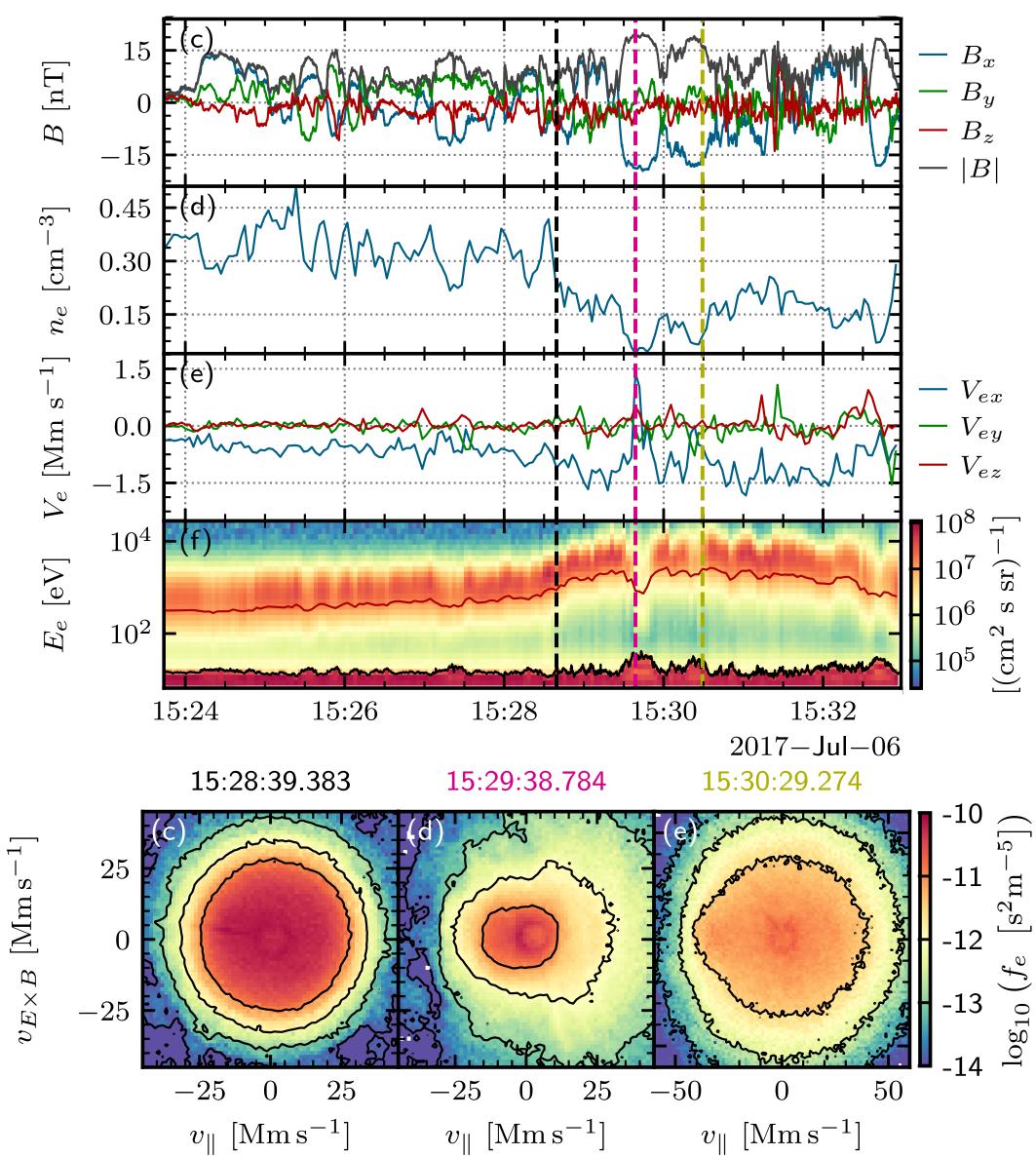
 $v_{\parallel} \, [\mathrm{Mm \, s^{-1}}]$

[Richard+2025, PRL]

Example

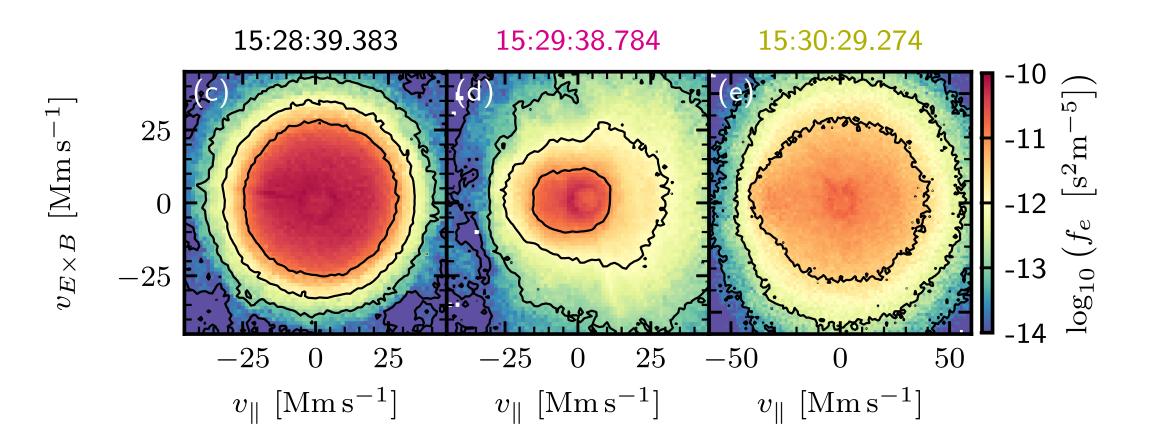
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[Richard+2025, PRL]

Modeling the electron VDFs

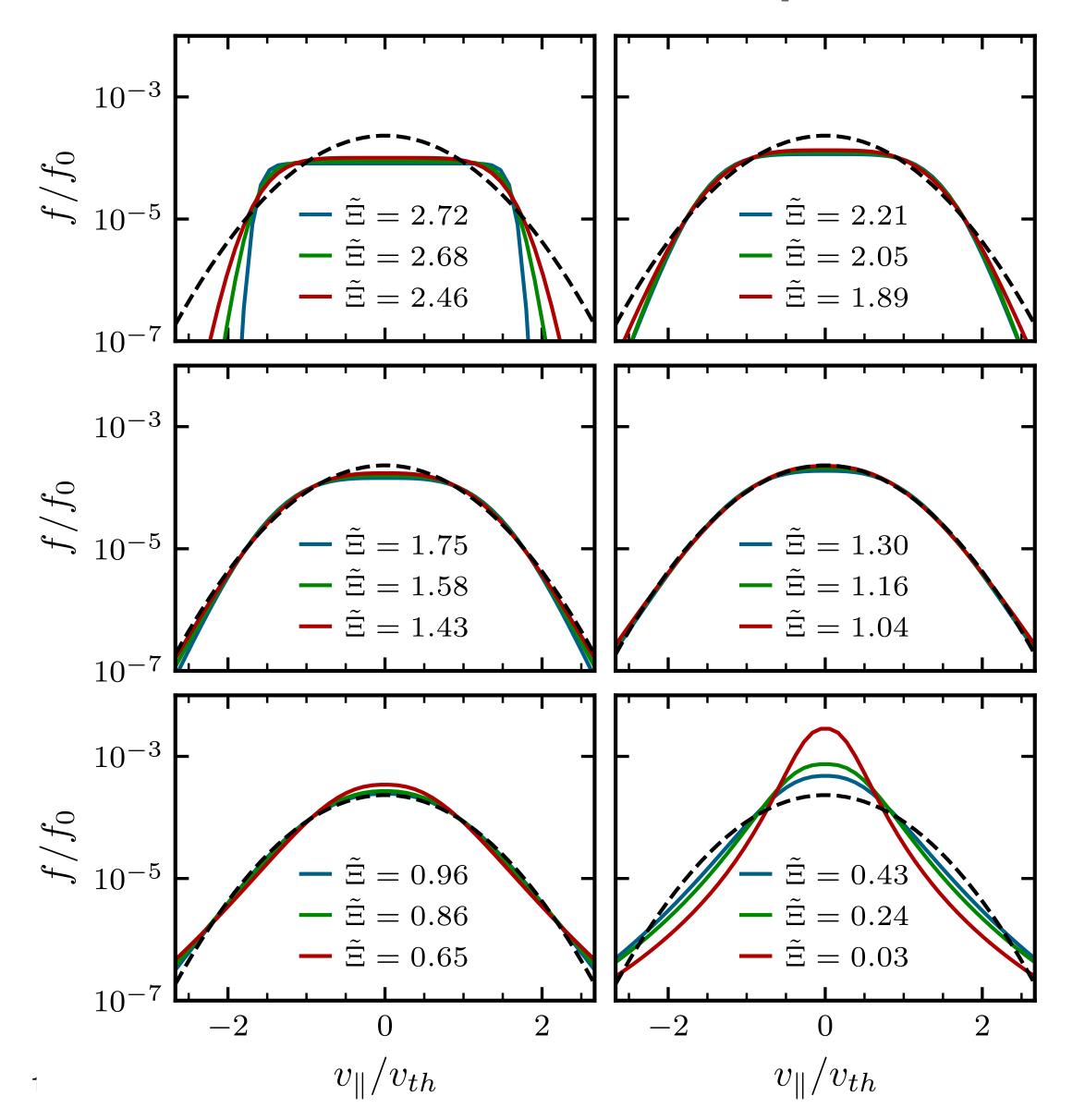


[Richard+2025, PRL]

Modeling the electron VDFs

- Model electron VDF (r,q) [Qureshi+2004, PoP] for $v_{\perp} \simeq 0$

$$f_{(r,q)} = f_0 \left[1 + \left(\frac{v_{\parallel}^2}{\xi(r,q)v_{te,\parallel}^2} \right)^{r+1} \right]^{-q}$$



[Richard+2025, PRL]

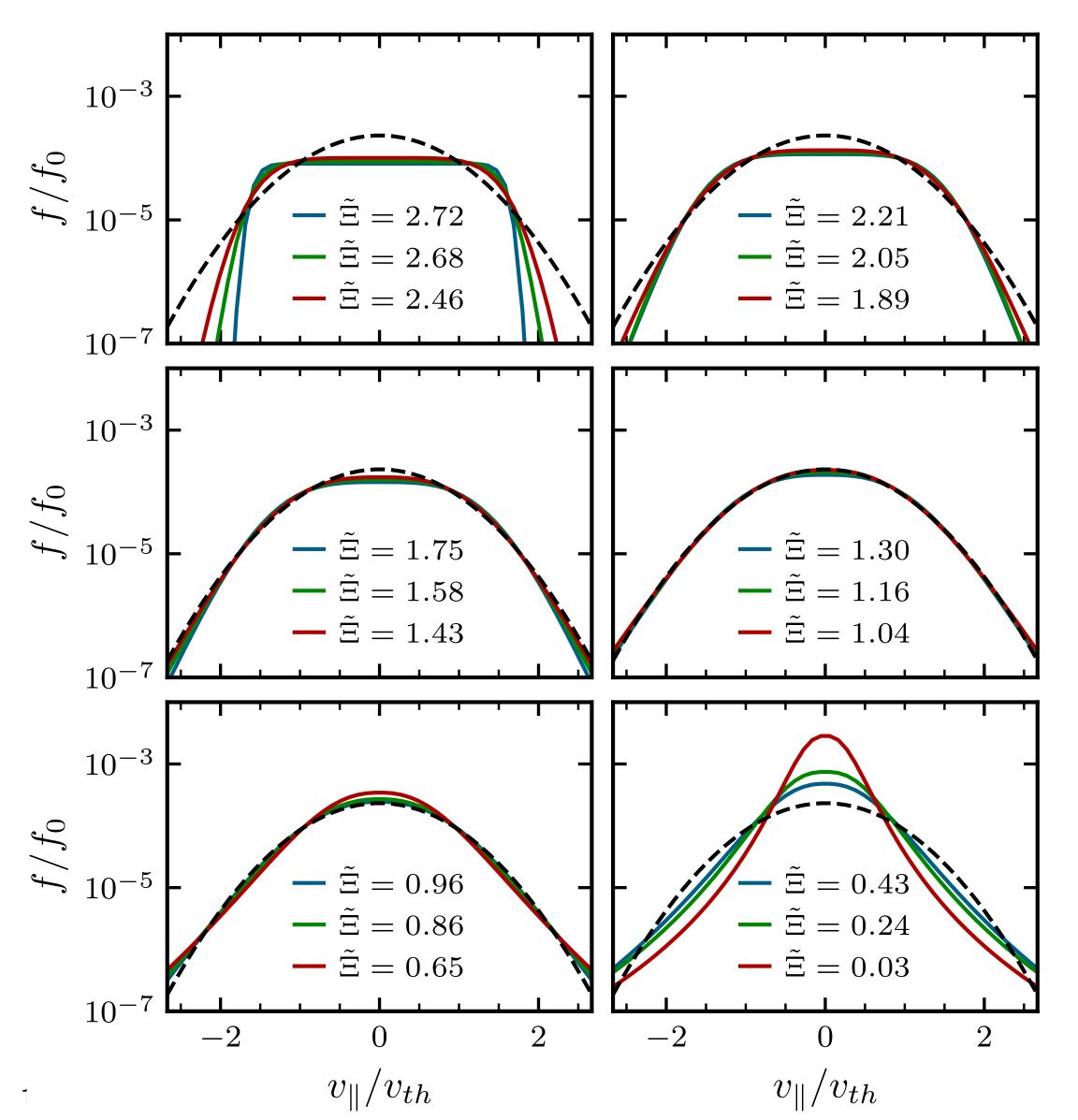
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• Flatness factor: $\Xi = f(v_{th})/f_0$:

$$\tilde{\Xi} \equiv \frac{\Xi_{(r,q)}}{\Xi_{bM}} = e \left(1 + \frac{1}{\xi^{r+1}} \right)^{-q}$$



[Richard+2025, PRL]

Modeling the electron VDFs

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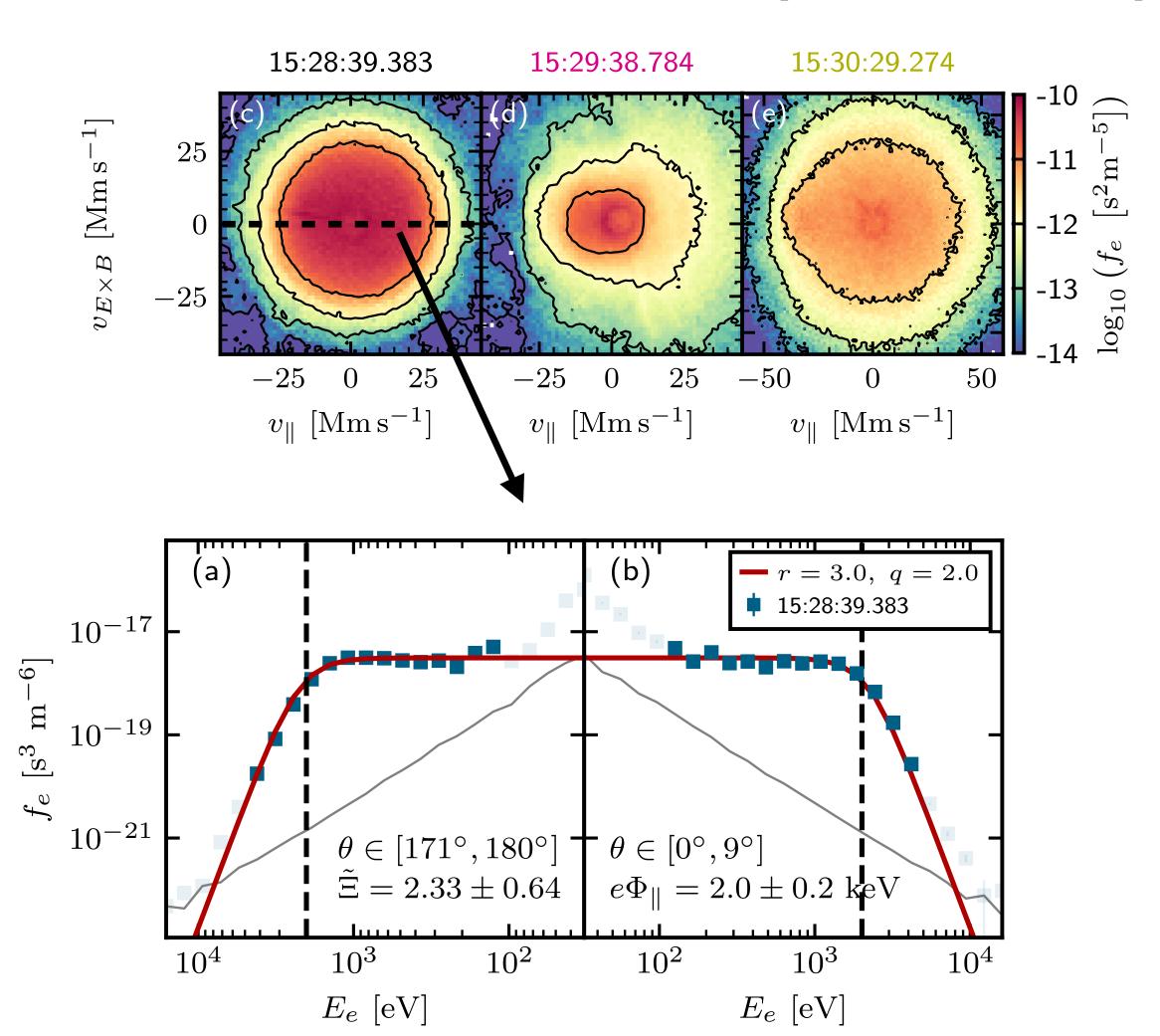
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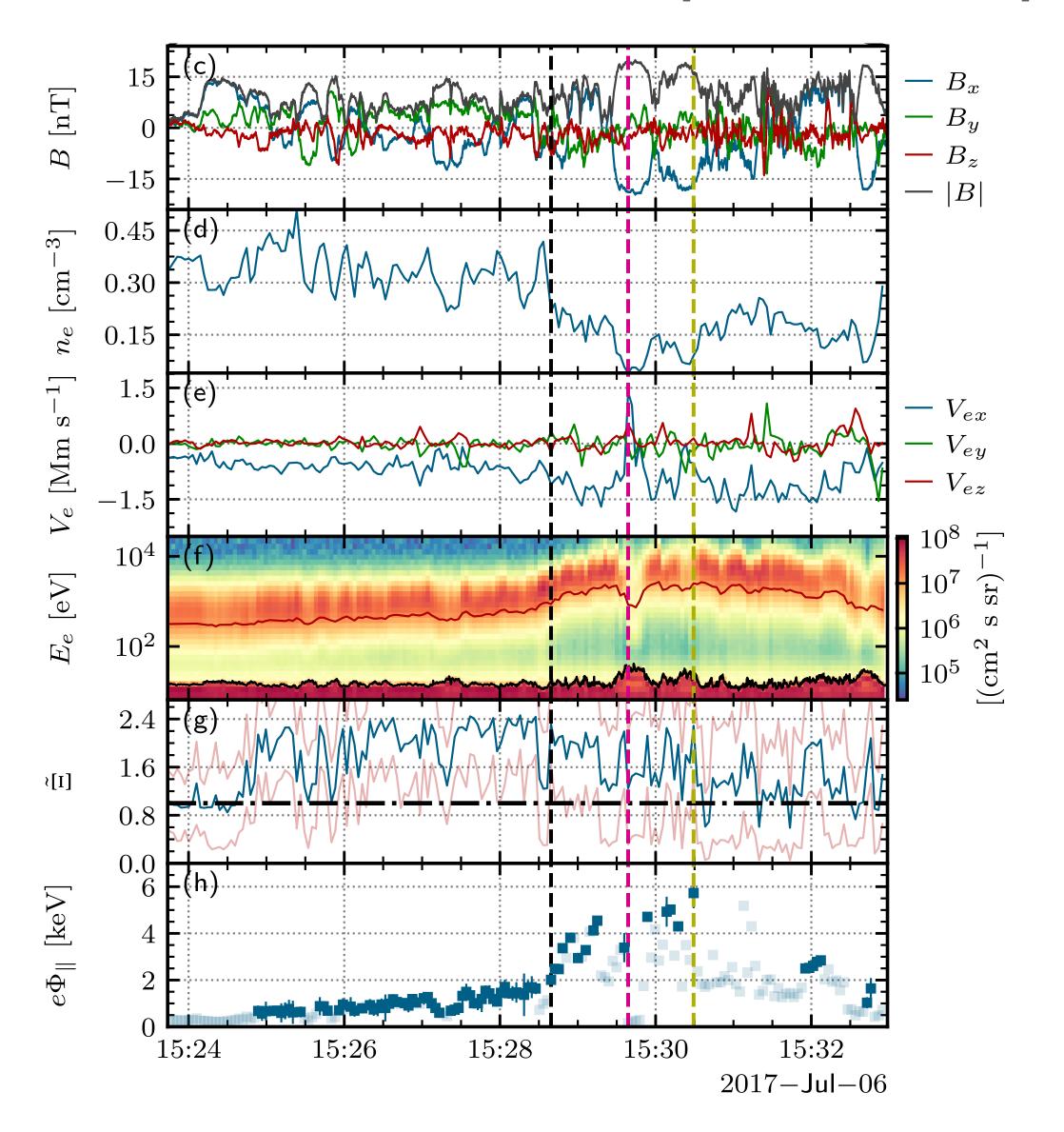
• Knee velocity v_Φ corresponding to the acceleration potential $f(v_\Phi)/f_0=\varepsilon$

$$v_{\Phi} = (\varepsilon^{-1/q} - 1)^{1/2(r+1)} \xi^{1/2} v_{te,\parallel}$$



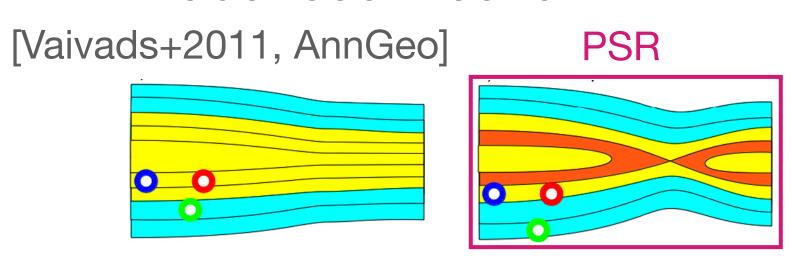
Acceleration potential

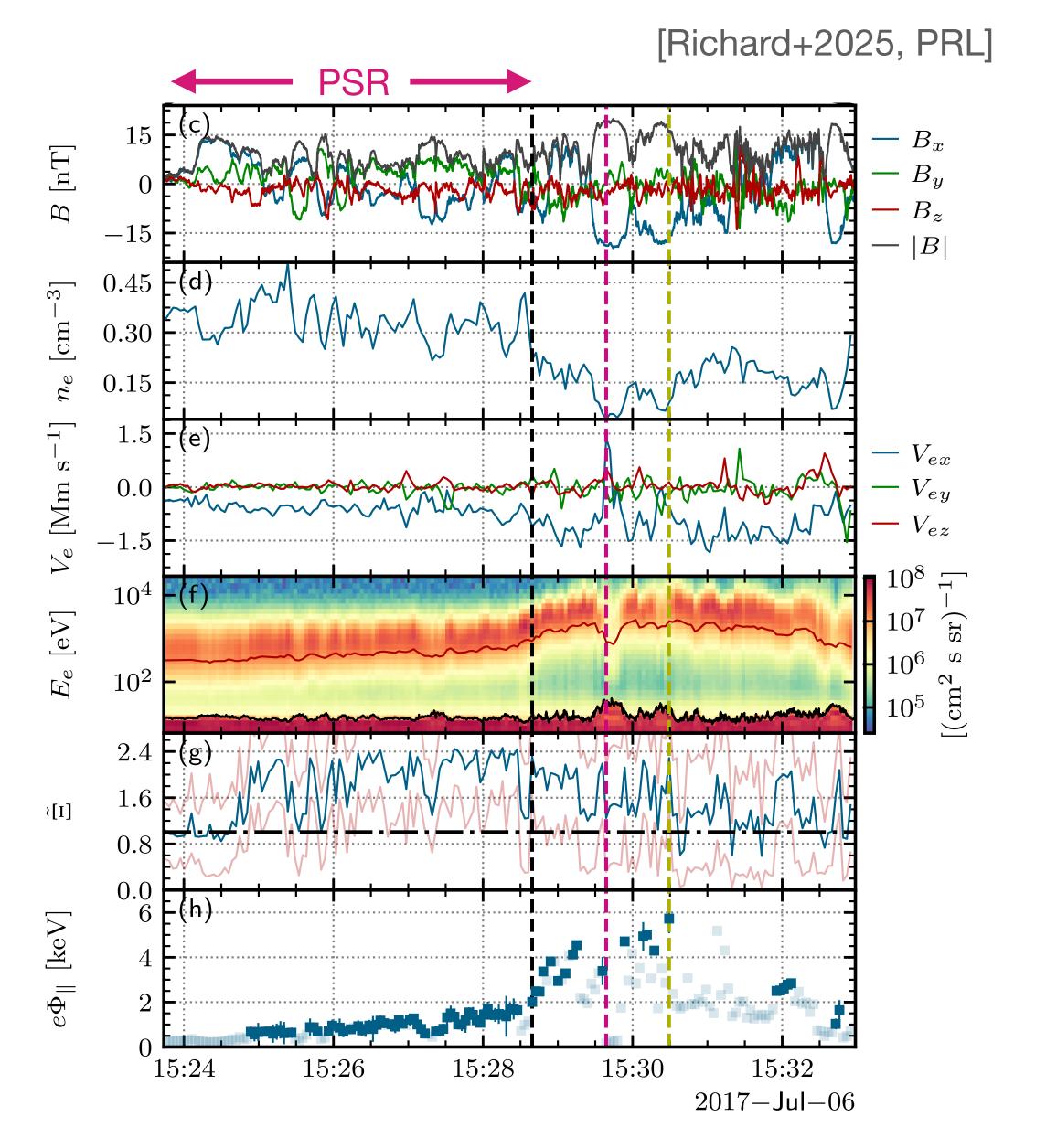
[Richard+2025, PRL]



Acceleration potential

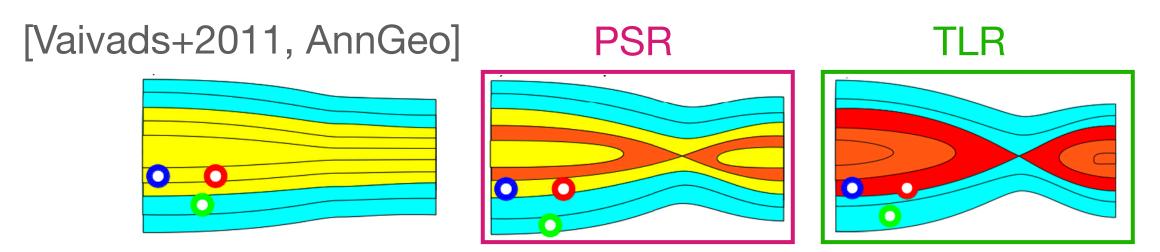
 Plasma sheet reconnection followed by lobe reconnection

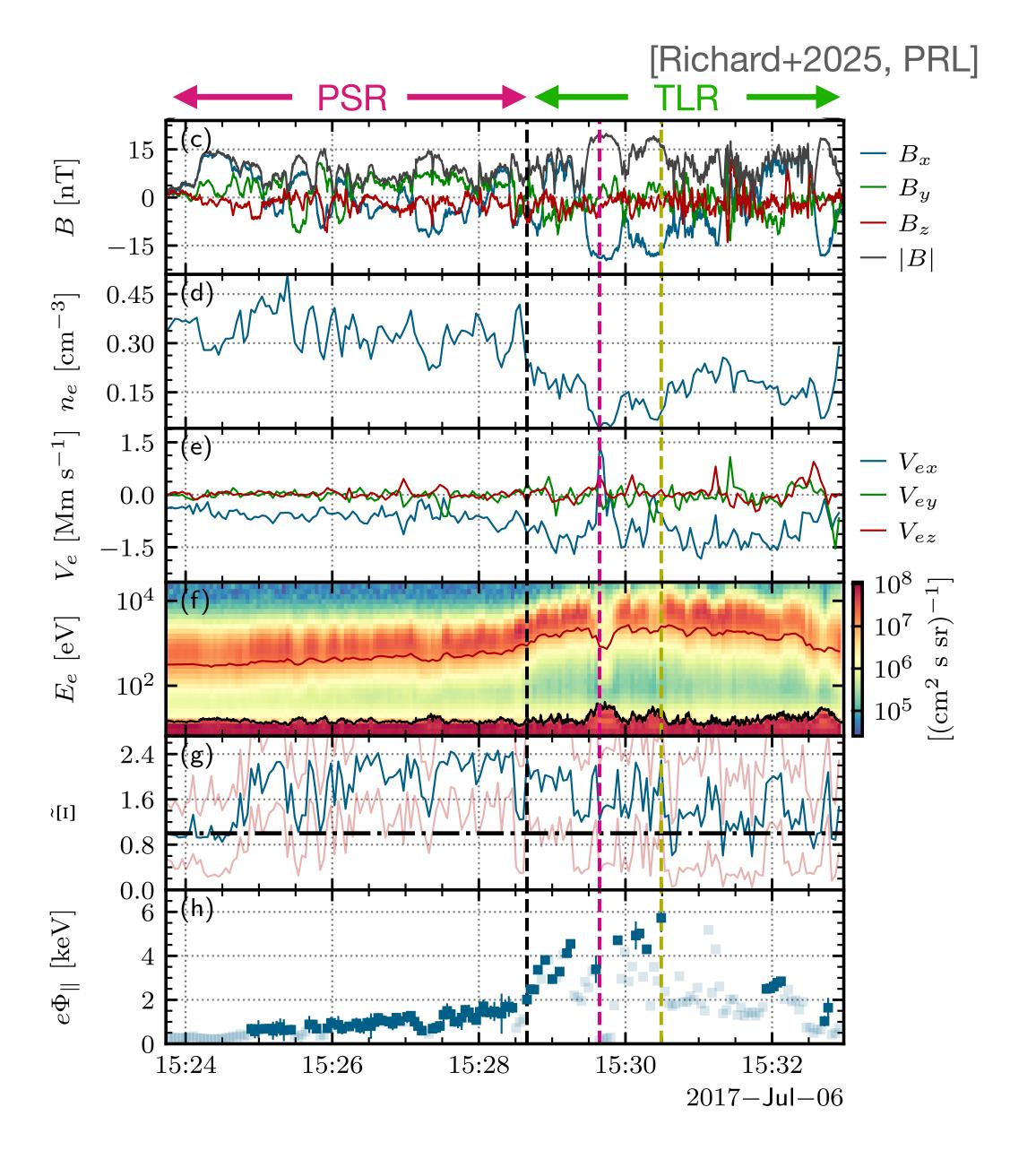




Acceleration potential

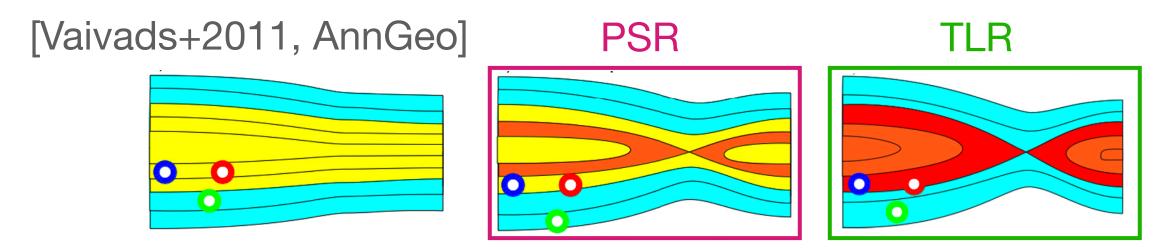
 Plasma sheet reconnection followed by lobe reconnection





Acceleration potential

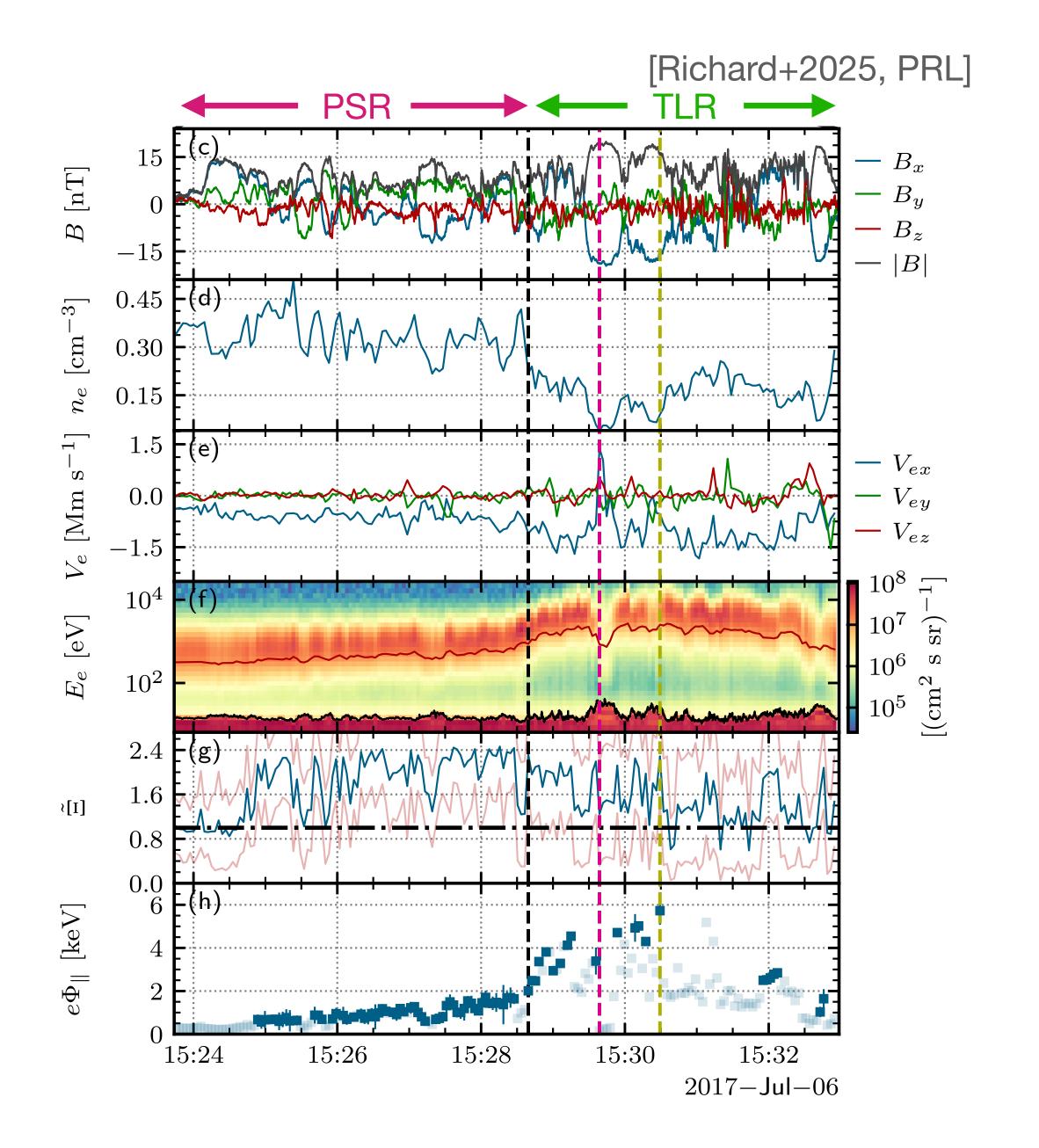
 Plasma sheet reconnection followed by lobe reconnection



 The increase in acceleration potential is proportional to the increase in inflow Alfvén speed

$$e\Phi_{\parallel}^{TLR}/e\Phi_{\parallel}^{PSR} \sim V_{Ae\infty}^{TLR}/V_{Ae\infty}^{PSR}$$

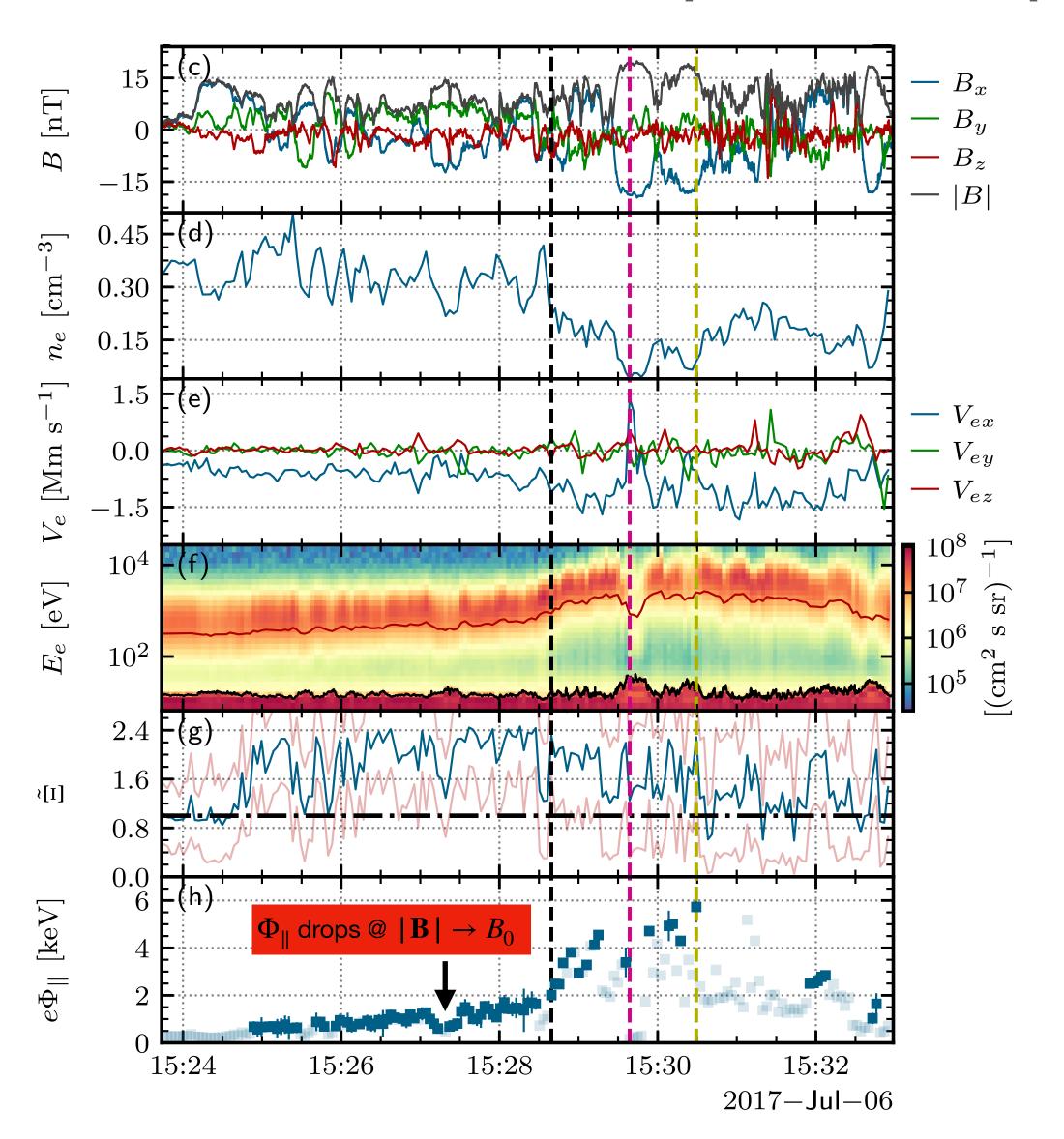
The acceleration potential increases with the inflow Alfvén speed



Results

The origin of E_{\parallel}

[Richard+2025, PRL]



Results

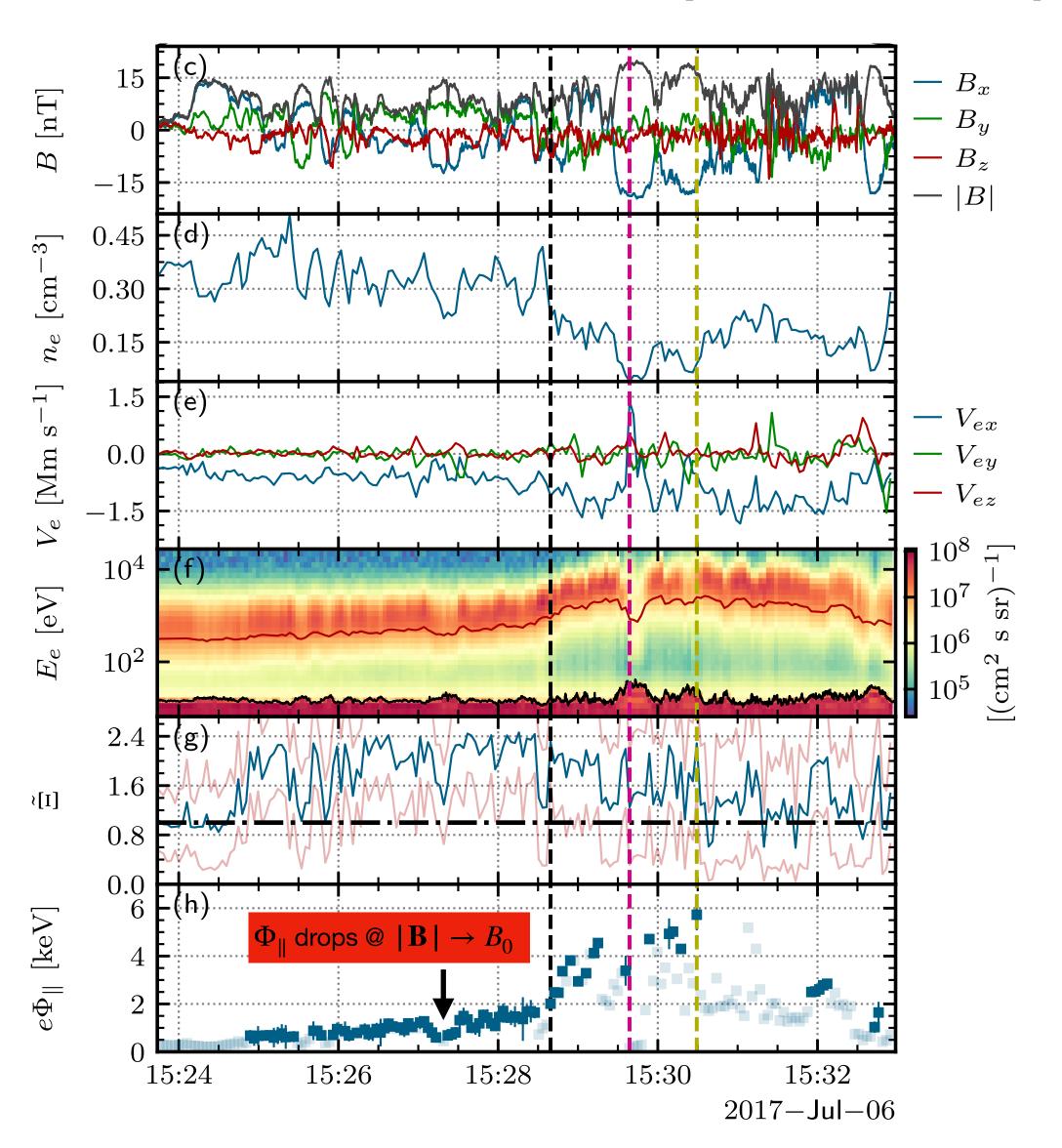
The origin of E_{\parallel}

 Steady-state local electron momentum balance along the field line

$$eE_{\parallel} = - \underbrace{\nabla_{\parallel} T_{e\parallel} - T_{e\parallel} \nabla_{\parallel} \ln n + (T_{e\parallel} - T_{e\perp}) \nabla_{\parallel} \ln B}_{\Phi_{\parallel T}}$$

$$\Phi_{\parallel T} \qquad \Phi_{\parallel n} \qquad \Phi_{\parallel B}$$

[Richard+2025, PRL]



Results

The origin of E_{\parallel}

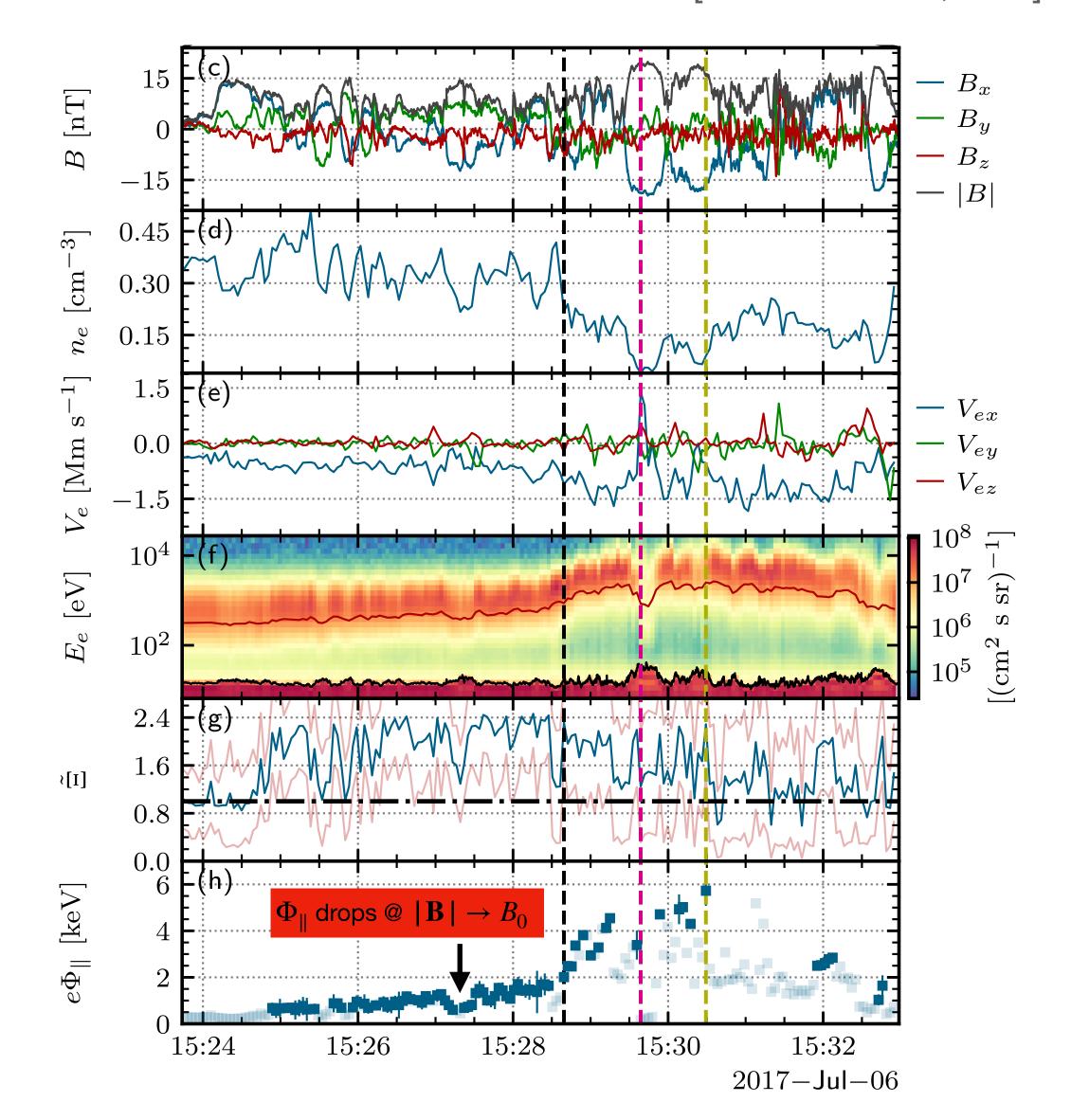
 Steady-state local electron momentum balance along the field line

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 Potential drop (incremental) across the outflow [Haggerty+2015, GRL]

$$e\Delta\Phi_{\parallel} = -\int_{b}^{a} E_{\parallel} dl \simeq \left(\frac{e\Delta\Phi_{\parallel T}}{10 \text{ eV}}\right) + \left(\frac{e\Delta\Phi_{\parallel n}}{160 \text{ eV}}\right) + \left(\frac{e\Delta\Phi_{\parallel B}}{50 \text{ eV}}\right)$$

Field-aligned ambipolar electric field is primarily due to electron density gradients

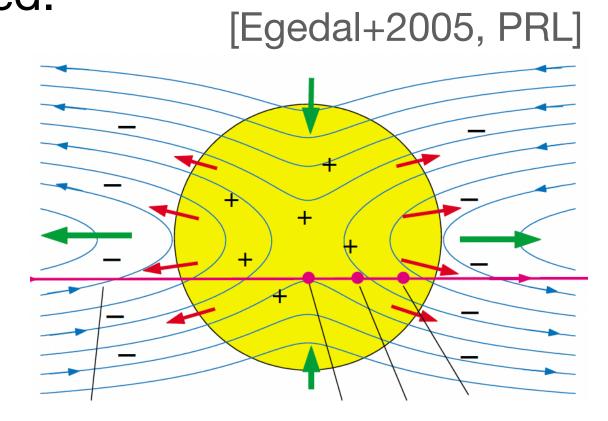


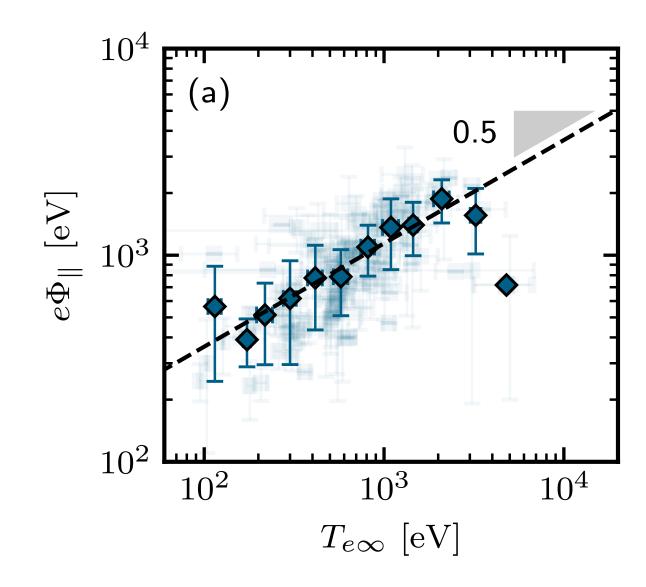
Results Statistics

Results

Statistics

• $e\Phi_{\parallel}$ increases with $T_{e\infty}$ to keep electrons trapped.

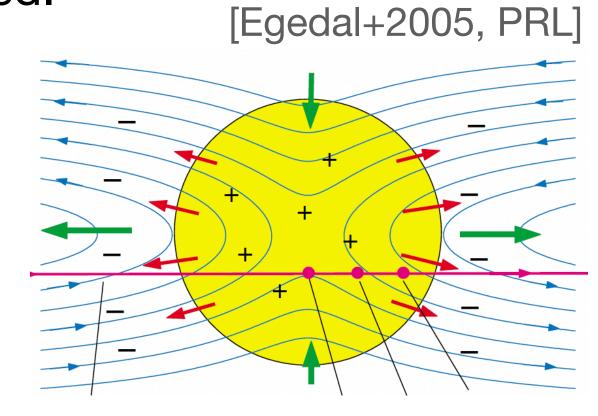




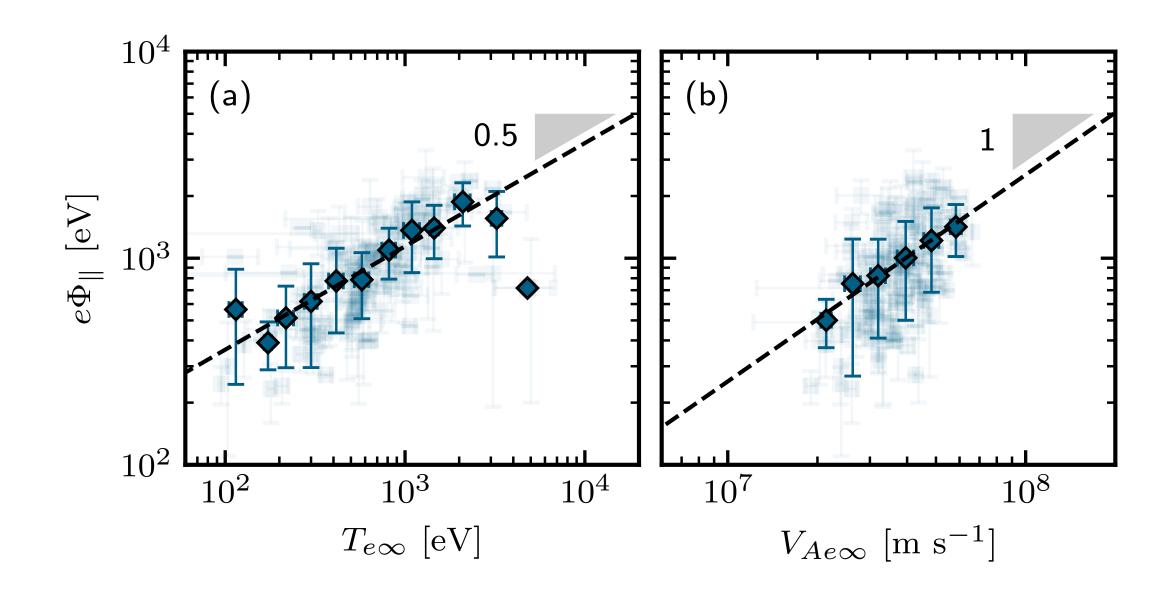
Results

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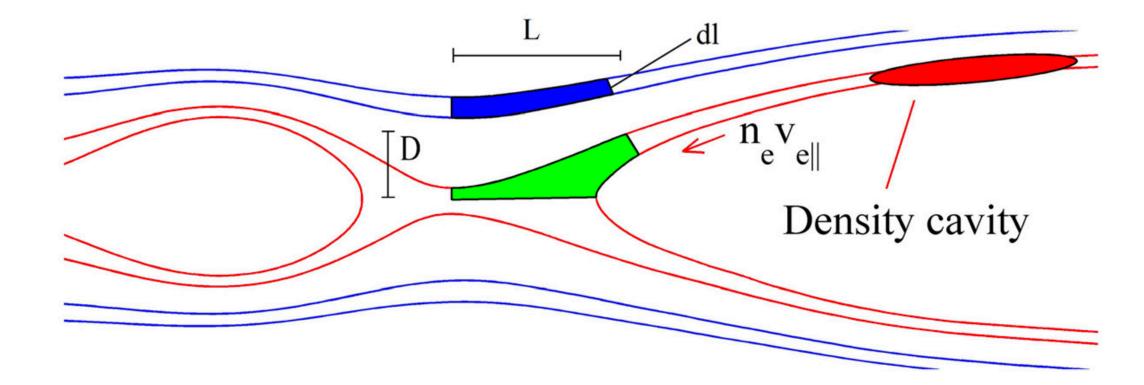
• $e\Phi_{\parallel}$ increases with $T_{e\infty}$ to keep electrons trapped.



- $e\Phi_{\parallel}$ increases with $V_{Ae\infty}$ to balance the flux tube expansion



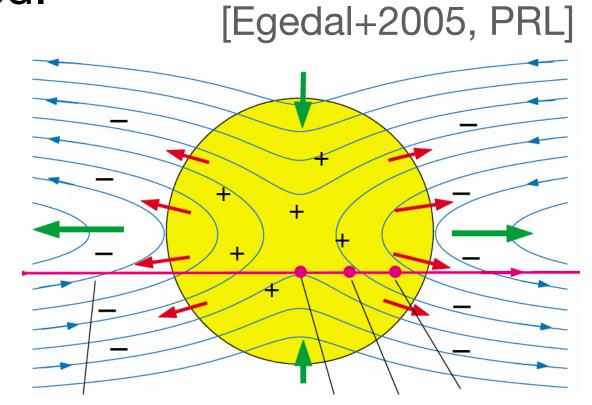
[Egedal+2015, PoP]



Results

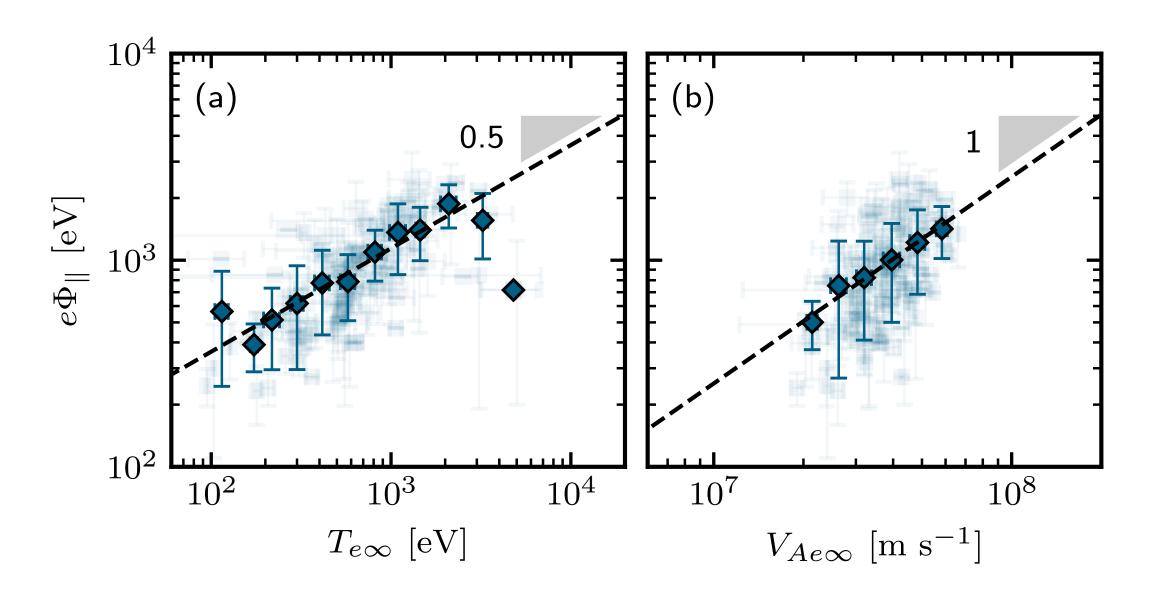
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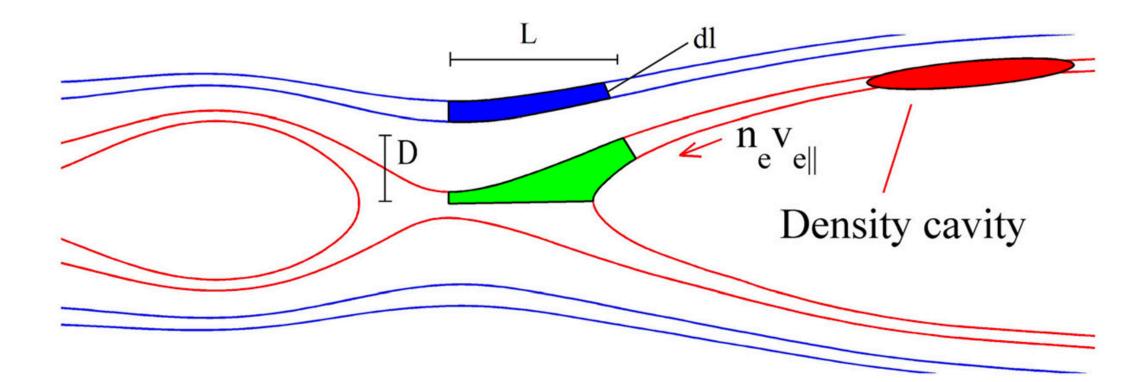


- $e\Phi_{||}$ increases with $V_{Ae\infty}$ to balance the flux tube expansion

 $e\Phi_{||}$ increases with $T_{e\infty}$ and $V_{Ae\infty}$ to maintain quasi-neutrality

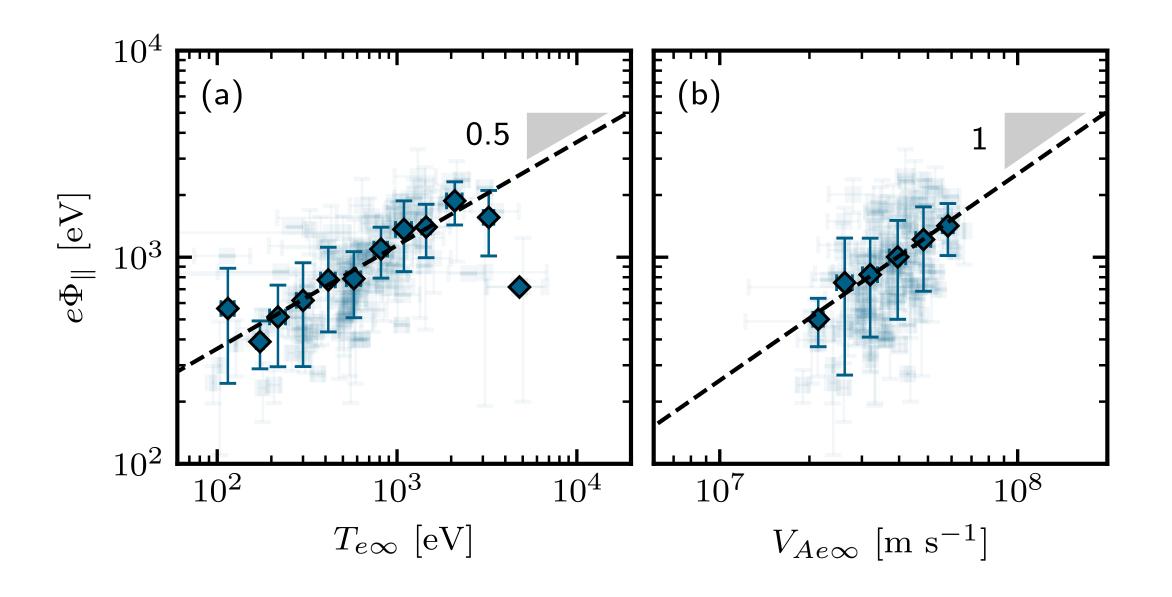


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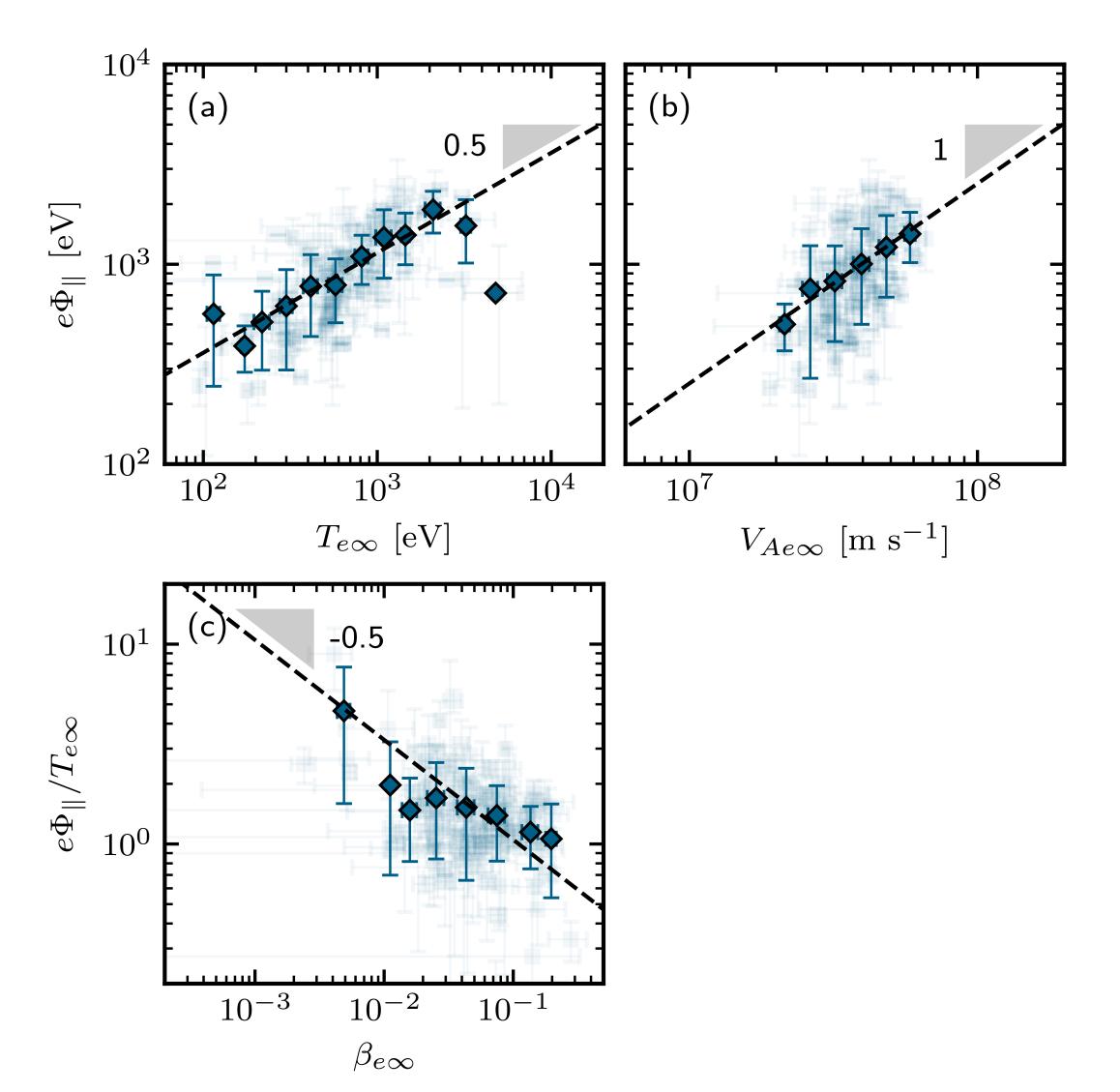
Role of E_{\parallel} in the energy partition



Results

Role of $E_{||}$ in the energy partition

• The acceleration potential scales as $e\Phi_{||}/T_{e\infty}=\alpha_{\Phi}\beta_{e\infty}^{-1/2}$

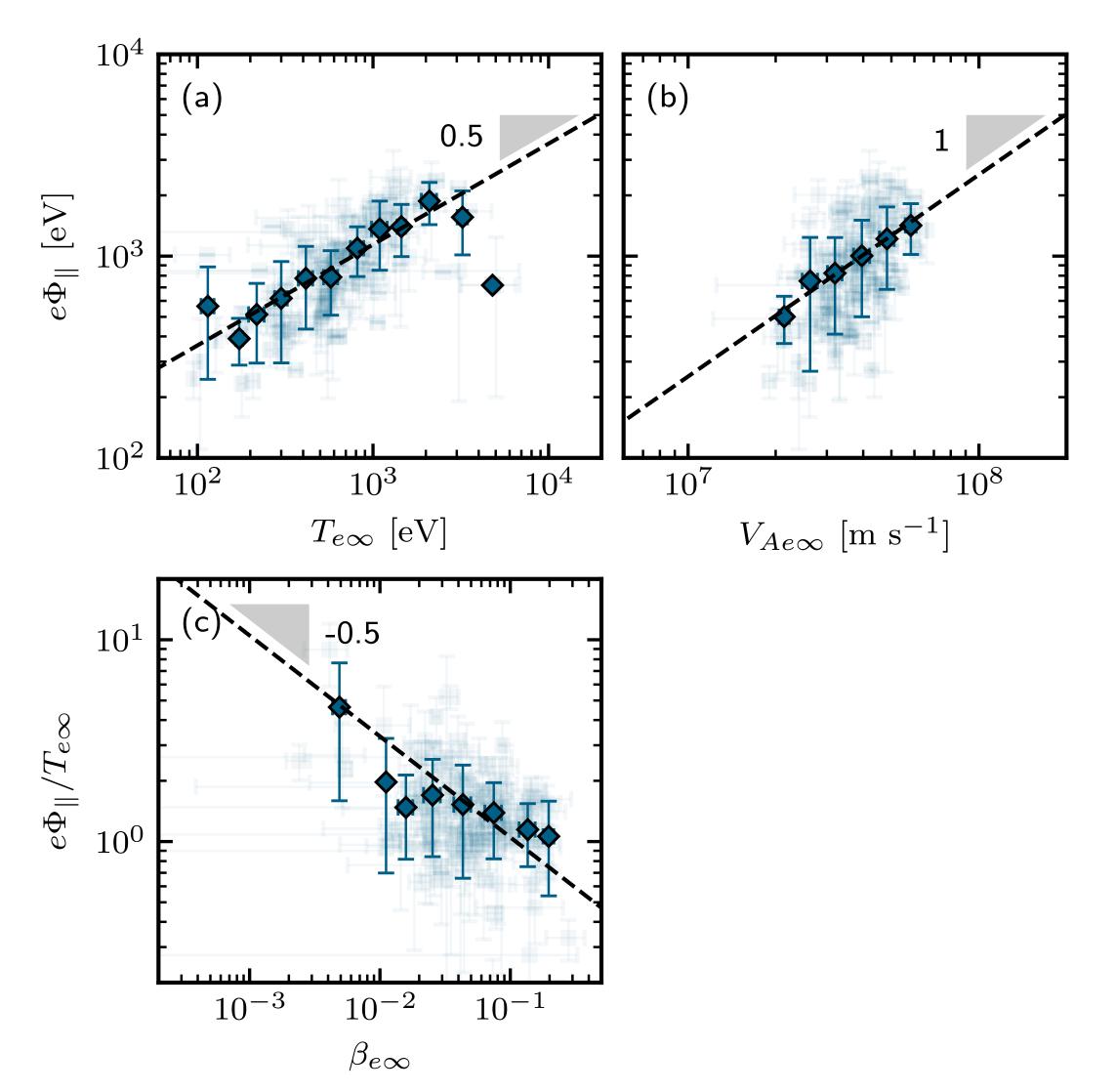


Results

Role of E_{\parallel} in the energy partition

- The acceleration potential scales as $e\Phi_{\parallel}/T_{e\infty}=\alpha_{\Phi}\beta_{e\infty}^{-1/2}$
- Electron heating: E_{rec} + E_{\parallel} [Le+2016, PoP, Øieroset+2020, ApJ]

$$\Delta T_e = \alpha_e m_i V_{Ai\infty}^2$$



Results

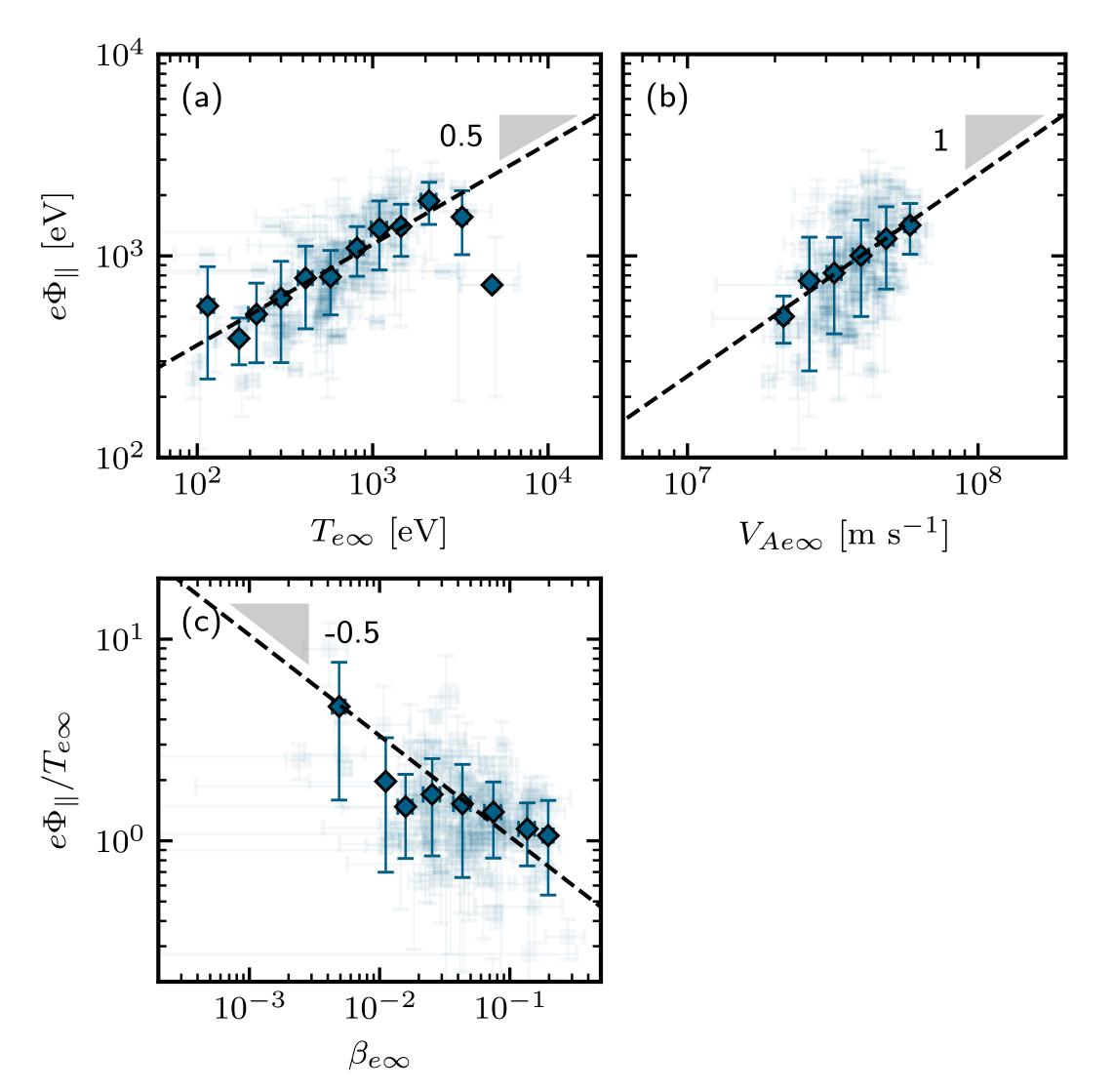
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• Ion heating: Pick-up + E_{\parallel} [Drake+2009, JGR, Haggerty+2015, GRL]

$$\Delta T_i = \alpha_{i1} m_i V_{Ai\infty}^2 - 2\alpha_{i2} e \Phi_{\parallel}$$



Results

Role of E_{\parallel} in the energy partition

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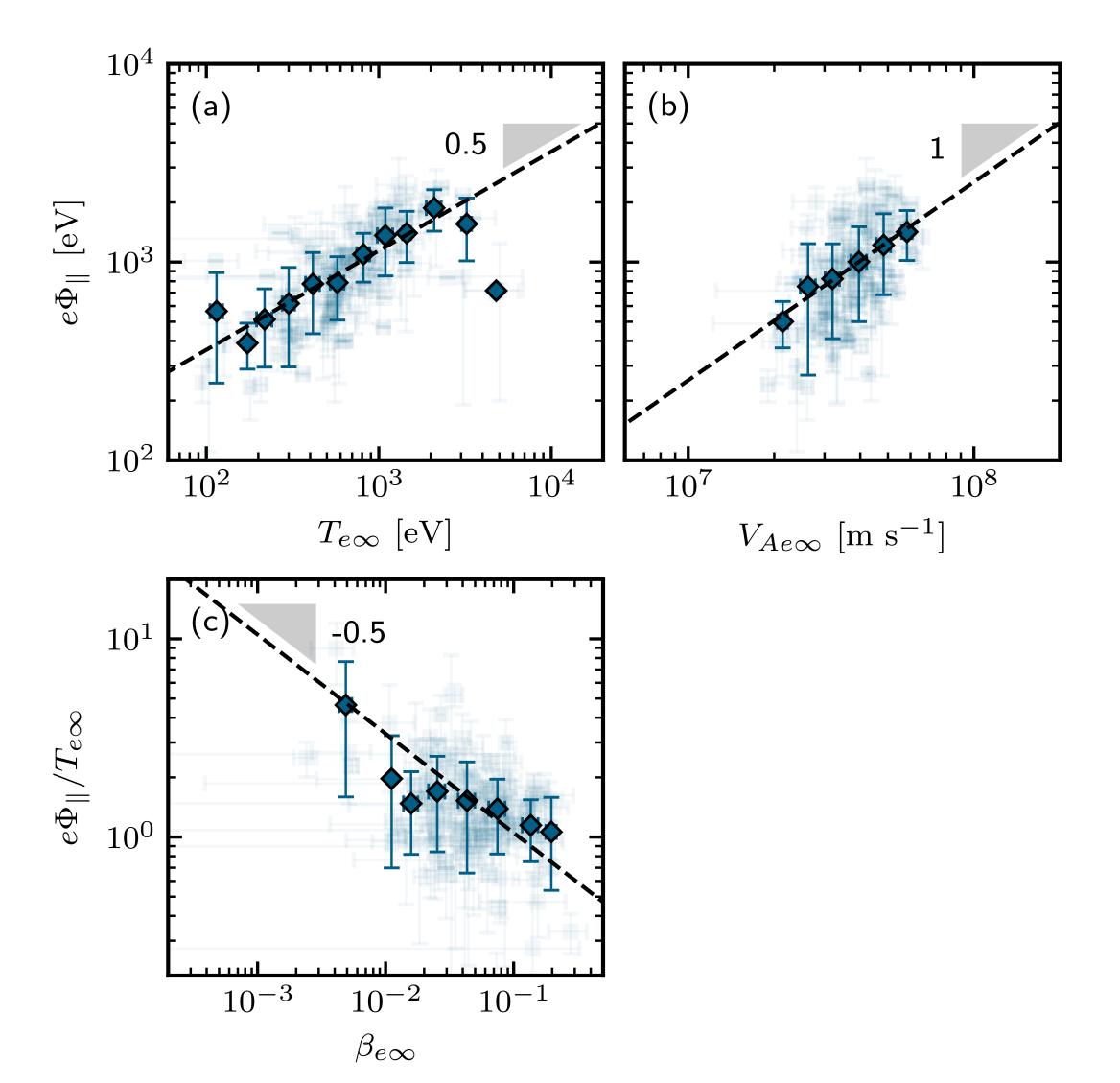
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$$\Delta T_i = \alpha_{i1} m_i V_{Ai\infty}^2 - 2\alpha_{i2} e \Phi_{\parallel}$$

Empirical model for the ion-to-electron energy partition

$$\frac{\Delta T_i}{\Delta T_e} = \frac{\alpha_{i1}}{\alpha_e} \left(1 - \frac{\alpha_{i2} \alpha_{\Phi}}{\alpha_{i1}} \sqrt{\beta_{e\infty}} \right)$$



Conclusions

 Turbulence develops rapidly in the reconnection jet, contributing approximately 10% to energy conversion during magnetic reconnection.

• Ion VDFs in the reconnection jet are efficiently heated and maintained near LTE due to chaotic interactions in sharply curved magnetic fields.

• Electron VDFs in the reconnection jet are heated up to 10 times by parallel electric fields that develop to maintain quasi-neutrality.