

MIXED QCD-EW CORRECTIONS TO THE DRELL-YAN PROCESS

GDR QCD workshop on W mass
Orsay, 01/07/2025



**GHENT
UNIVERSITY**

 **FACULTY
OF SCIENCES**

Simone Devoto



European Research Council
Established by the European Commission

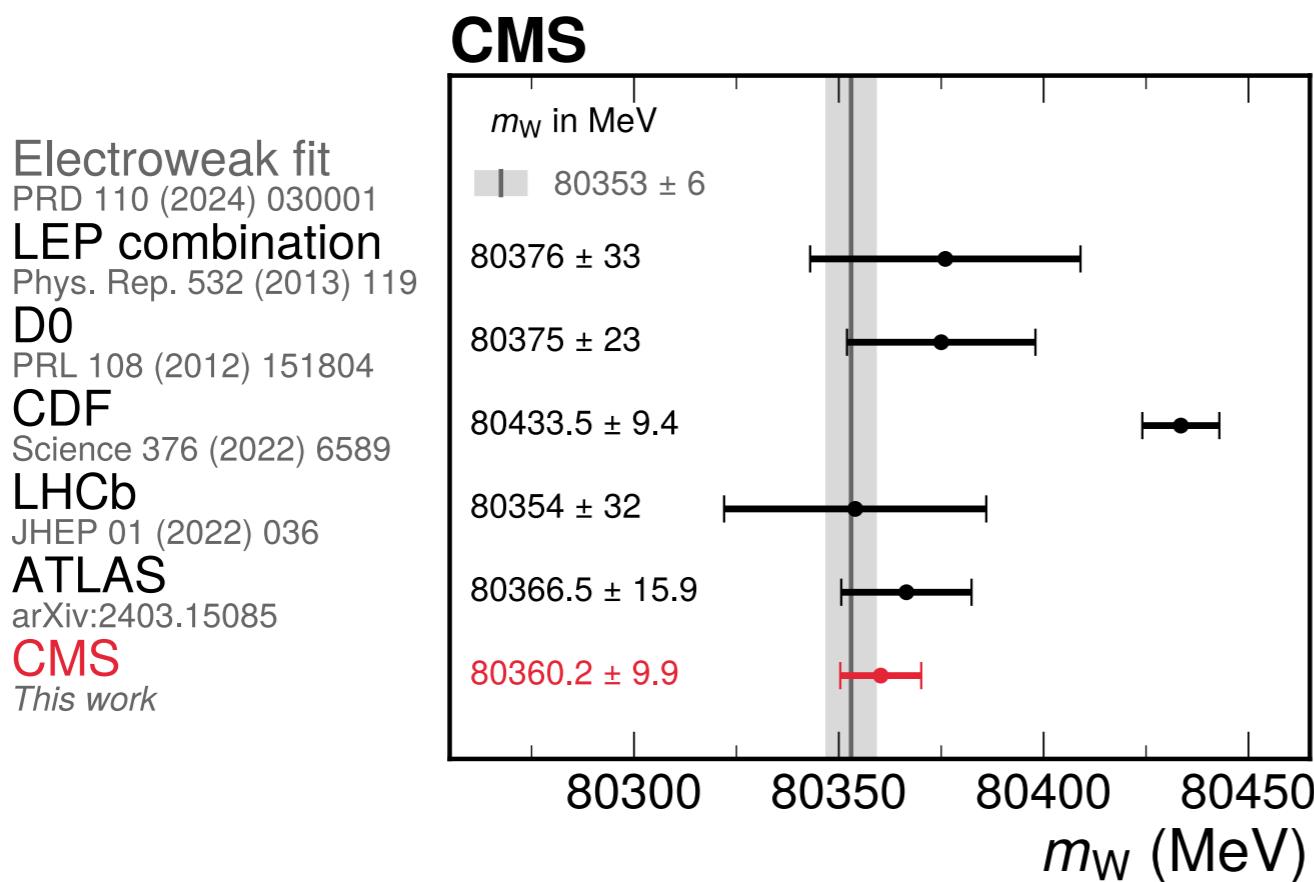
In collaboration with:
T. Armadillo, R. Bonciani, L. Buonocore,
M. Grazzini, S. Kallweit, N. Rana, A. Vicini

MOTIVATIONS

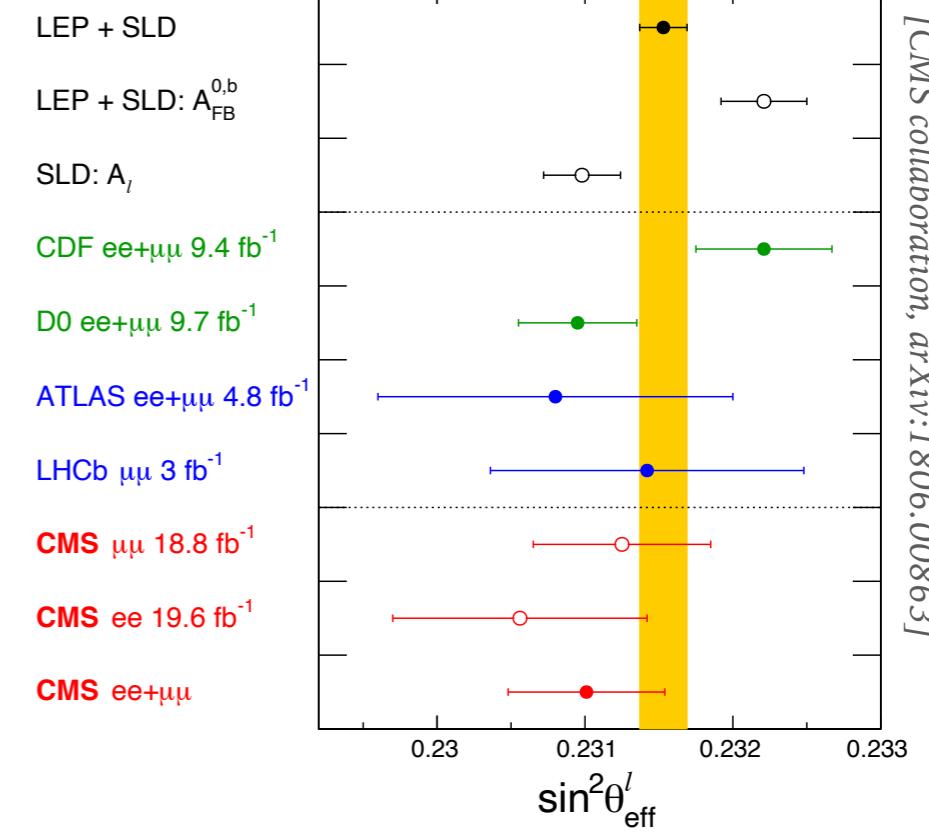
Standard Model Precision Studies

- Extremely precise measurement of W boson mass (10 MeV uncertainties!);
- Determination of $\sin^2 \theta_{\text{eff}}^{\text{lep}}$ is becoming competitive with LEP result: 0.23152(16);
- Crucial process for PDF determination.

Can we provide SM theoretical predictions matching experimental accuracy?



[CMS collaboration, arXiv:2412.13872]



MIXED QCDxEW CORRECTIONS

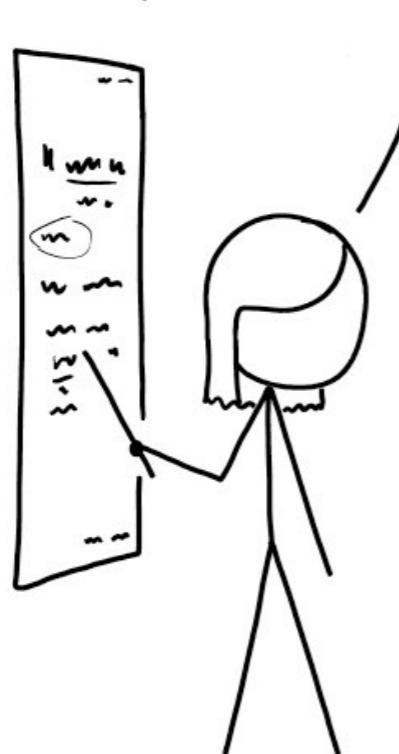
$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4) = \sigma^{(0,0)} \quad \text{Drell-Yan (1970)}$$

MIXED QCDxEW CORRECTIONS

$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4) = \sigma^{(0,0)} \quad \text{Drell-Yan (1970)}$$

AT THIS POINT, YOU'RE PROBABLY
THINKING, "I LOVE THIS EQUATION
AND WISH IT WOULD NEVER END!"

WELL, GOOD NEWS!



credits: xkcd (2605)

TAYLOR SERIES EXPANSION IS THE ~~WORST.~~
BEST!

$$\begin{aligned} &+ \alpha_S \sigma^{(1,0)} &+ \alpha \sigma^{(0,1)} \\ &+ \alpha_S^2 \sigma^{(2,0)} &+ \alpha \alpha_S \sigma^{(1,1)} &+ \alpha^2 \sigma^{(0,2)} \\ &+ \alpha_S^3 \sigma^{(3,0)} &+ \dots \end{aligned}$$

QCD

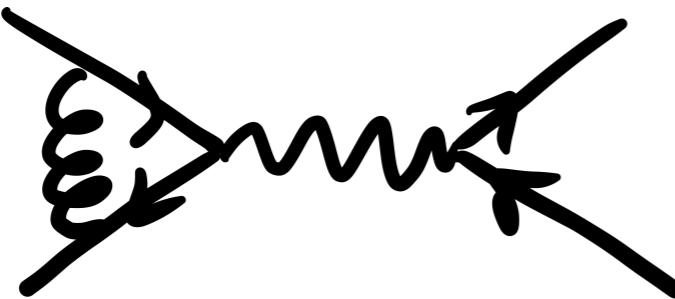
MIXED

EW

MIXED QCDxEW CORRECTIONS

$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4) = \sigma^{(0,0)}$$

QCD CORRECTIONS



$$\begin{aligned} & + \alpha_S \sigma^{(1,0)} + \alpha \sigma^{(0,1)} \\ & + \alpha_S^2 \sigma^{(2,0)} + \alpha \alpha_S \sigma^{(1,1)} + \alpha^2 \sigma^{(0,2)} \\ & + \alpha_S^3 \sigma^{(3,0)} + \dots \end{aligned}$$

► Dominant effect;

NLO:

[G.Altarelli, R.Ellis, G.Martinelli *Nucl.Phys.B* 157 (1979)];

NNLO:

[R.Hamberg, T.Matsuura, W.van Nerveen, *Nucl. Phys. B* 359 (1991)];

[C.Anastasiou, L.J.Dixon, K.Melnikov, F.Petriello, *hep-ph:0306192*];

[S.Catani, L.Cieri, G.Ferrera, D.de Florian, M.Grazzini
arXiv:0903.2120];

► Known up to N3LO.

N3LO:

[C.Duhr, F.Dulat, B.Mistlberger *arXiv:2007.13313*];

[X.Chen, T.Gehrmann, N.Glover, A.Huss, T.Yang, and H.Zhu
arXiv:2107.09085];

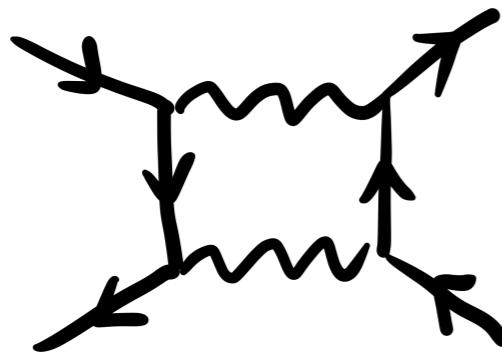
[S.Camarda, L.Cieri, G.Ferrera *arXiv:2103.04974*];

[X.Chen, T.Gehrmann, N.Glover, A.Huss, P.Monni, E.Re, L.Rottoli,
P.Torrielli *arXiv:2203.01565*].

MIXED QCDxEW CORRECTIONS

$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4) = \sigma^{(0,0)}$$

EW CORRECTIONS



$$+ \alpha_S \sigma^{(1,0)} + \alpha \sigma^{(0,1)}$$

$$+ \alpha_S^2 \sigma^{(2,0)} + \alpha \alpha_S \sigma^{(1,1)} + \alpha^2 \sigma^{(0,2)}$$

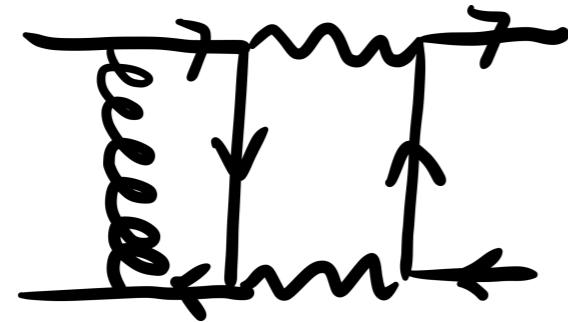
$$+ \alpha_S^3 \sigma^{(3,0)} + \dots$$

- Subdominant with respect to QCD: **physical counting** $\alpha_S \simeq \alpha^2$;
- NLO corrections known;
[U.Baur, O.Brein, W.Hollik, C.Schappacher, D.Wackerlo, *hep-ph:0108274*];
[S.Dittmaier, M.Kramer, *hep-ph:0109062*];
[U.Baur, D.Wackerlo, *hep-ph:0405191*];
- NNLO corrections still missing (available Sudakov high energy approximation).
[B.Jantzen, J.H.Kühn, A.A.Penin, V.A.Smirnov, *hep-ph:0509157*];

MIXED QCDxEW CORRECTIONS

$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4) = \sigma^{(0,0)}$$

MIXED CORRECTIONS



$$\begin{aligned} &+ \alpha_S \sigma^{(1,0)} + \alpha \sigma^{(0,1)} \\ &+ \alpha_S^2 \sigma^{(2,0)} + \alpha \alpha_S \sigma^{(1,1)} + \alpha^2 \sigma^{(0,2)} \\ &+ \alpha_S^3 \sigma^{(3,0)} + \dots \end{aligned}$$

NLO QCD and NLO EW separately large: what about mixed?

By physical counting, expected size comparable with N3LO QCD!

[R. Bonciani, L. Buonocore, M. Grazzini, S. Kallweit, N. Rana, F. Tramontano, A. Vicini, arXiv:2106.11953]

[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini, S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]

[F. Buccioni, F. Caola, H. Chawdhry, F. Devoto, M. Heller, A. von Manteuffel, K. Melnikov, R. Röntsch, C. Signorile-Signorile, arXiv:2203.11237]

RECENT PROGRESSES IN MIXED CORRECTIONS

► Theoretical Developments

- **2-loop virtual Master Integrals with internal masses** [*U. Aglietti, R. Bonciani, arXiv:0304028, arXiv:0401193*], [*R. Bonciani, S. Di Vita, P. Mastrolia, U. Schubert, arXiv:1604.08581*], [*M.Heller, A.von Manteuffel, R.Schabinger arXiv:1907.00491*], [*M.Long,R,Zhang,W.Ma,Y.Jiang,L.Han,,Z.Li,S.Wang, arXiv:2111.14130*], [*X.Liu, Y.Ma, arXiv:2201.11669*]
- **Altarelli-Parisi splitting functions including QCD-QED effects** [*D. de Florian, G. Sborlini, G. Rodrigo, arXiv:1512.00612*]
- **Renormalisation** [*G.Degrassi, A.Vicini, hep-ph/0307122*], [*S.Dittmaier,T.Schmidt,J.Schwarz, arXiv:2009.02229*], [*S.Dittmaier, arXiv:2101.05154*]

► On-shell Z and W production

- **pole approximation of the NNLO QCD-EW corrections** [*S.Dittmaier, A.Huss, C.Schwinn, arXiv:1403.3216, 1511.08016*]
- **analytical total Z production cross section including NNLO QCD-QED corrections** [*D. de Florian, M.Der, I.Fabre, arXiv:1805.12214*]
- **fully differential on-shell Z production including exact NNLO QCD-QED corrections** [*M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:1909.08428*] [*S.Hasan, U.Schubert, arXiv:2004.14908*]
- **analytical total Z production cross section including NNLO QCD-EW corrections** [*R. Bonciani, F. Buccioni, R.Mondini, A.Vicini, arXiv:1611.00645*], [*R. Bonciani, F. Buccioni, N.Rana, I.Triscari, A.Vicini, arXiv:1911.06200*], [*R. Bonciani, F. Buccioni, N.Rana, A.Vicini, arXiv:2007.06518, arXiv:2111.12694*]
- **fully differential Z and W production including NNLO QCD-EW corrections** [*F. Buccioni, F. Caola, M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:2005.10221*], [*A. Behring, F. Buccioni, F. Caola, M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:2009.10386, 2103.02671*]

► Complete Drell-Yan

- **neutrino-pair production including NNLO QCD-QED corrections** [*L. Cieri, D. de Florian, M.Der, J.Mazzitelli, arXiv:2005.01315*]
- **2-loop amplitudes** [*M.Heller, A.von Manteuffel, R.Schabinger, arXiv:2012.05918*], [*T.Armadillo, R.Bonciani, SD, N.Rana, A.Vicini, arXiv:2201.01754; arXiv:2412.16095*],
- **NNLO QCD-EW corrections to neutral-current DY including leptonic decay** [*R.Bonciani, L.Buonocore, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano, A.Vicini, arXiv:2106.11953*], [*F.Buccioni, F.Caola, H.Chawdhry, F.Devoto, M.Heller,A.von Manteuffel, K.Melnikov, R.Röntsch, C.Signorile-Signorile, arXiv:2203.11237*] [*T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini, S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095*]
- **NNLO QCD-EW corrections to charged-current DY including leptonic decay (2-loop contributions in pole approximation).** [*L.Buonocore, M.Grazzini, S.Kallweit, C.Savoini, F.Tramontano, arXiv:2102.12539*]

CONTENTS



- **Phenomenological application:**
 - Neutral Current Drell-Yan;
 - Forward-Backward Asymmetry;
- **Two-loop Amplitudes:**
 - Theoretical challenges;
 - Our workflow;
 - Neutral Current Drell-Yan;
 - Charged Current Drell-Yan.

PHENOMENOLOGY

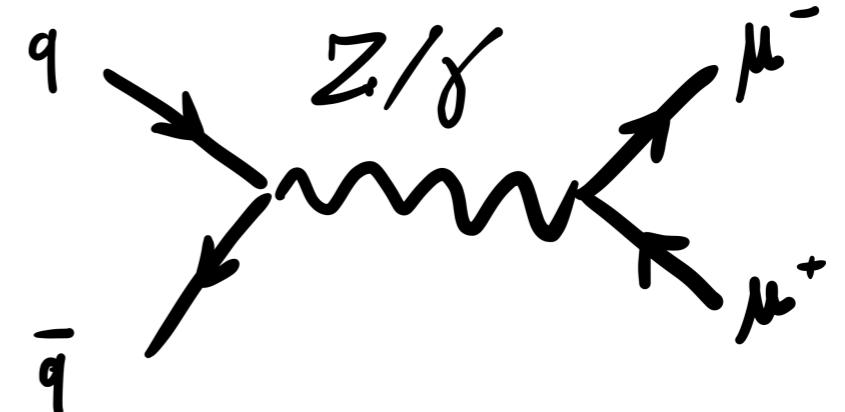
EXPLORING THE STANDARD MODEL



NEUTRAL CURRENT DRELL-YAN

[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]

$$q(p_1) + \bar{q}(p_2) \rightarrow \mu^-(p_3) + \mu^+(p_4)$$



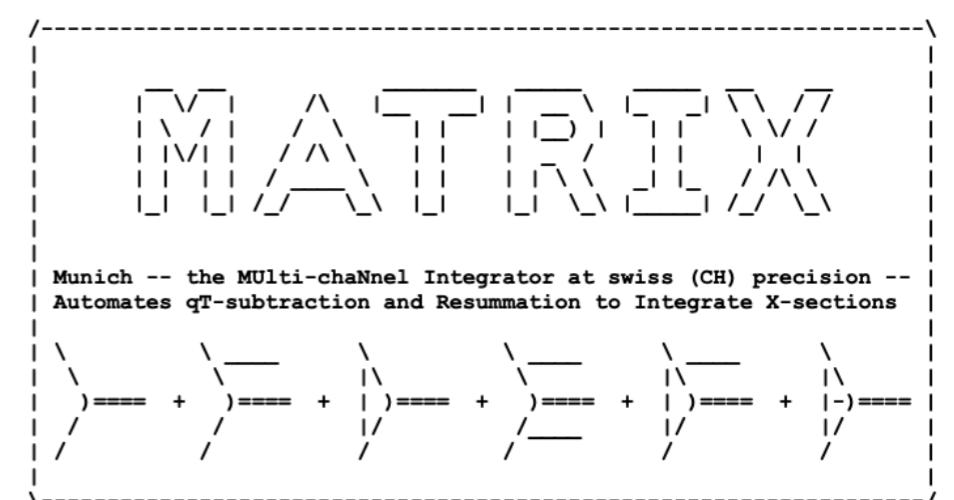
- IR singularities handled by **q_T-subtraction formalism**;

[S. Catani, M. Grazzini (2007)]

[L.Buonocore, M. Grazzini, F.Tramontano (2019)]

$$d\sigma_{(N)NLO}^F = \mathcal{H}_{(N)NLO}^F \otimes d\sigma_{LO}^F + \left[d\sigma_{(N)LO}^{F+jets} - d\sigma_{(N)LO}^{CT} \right]$$

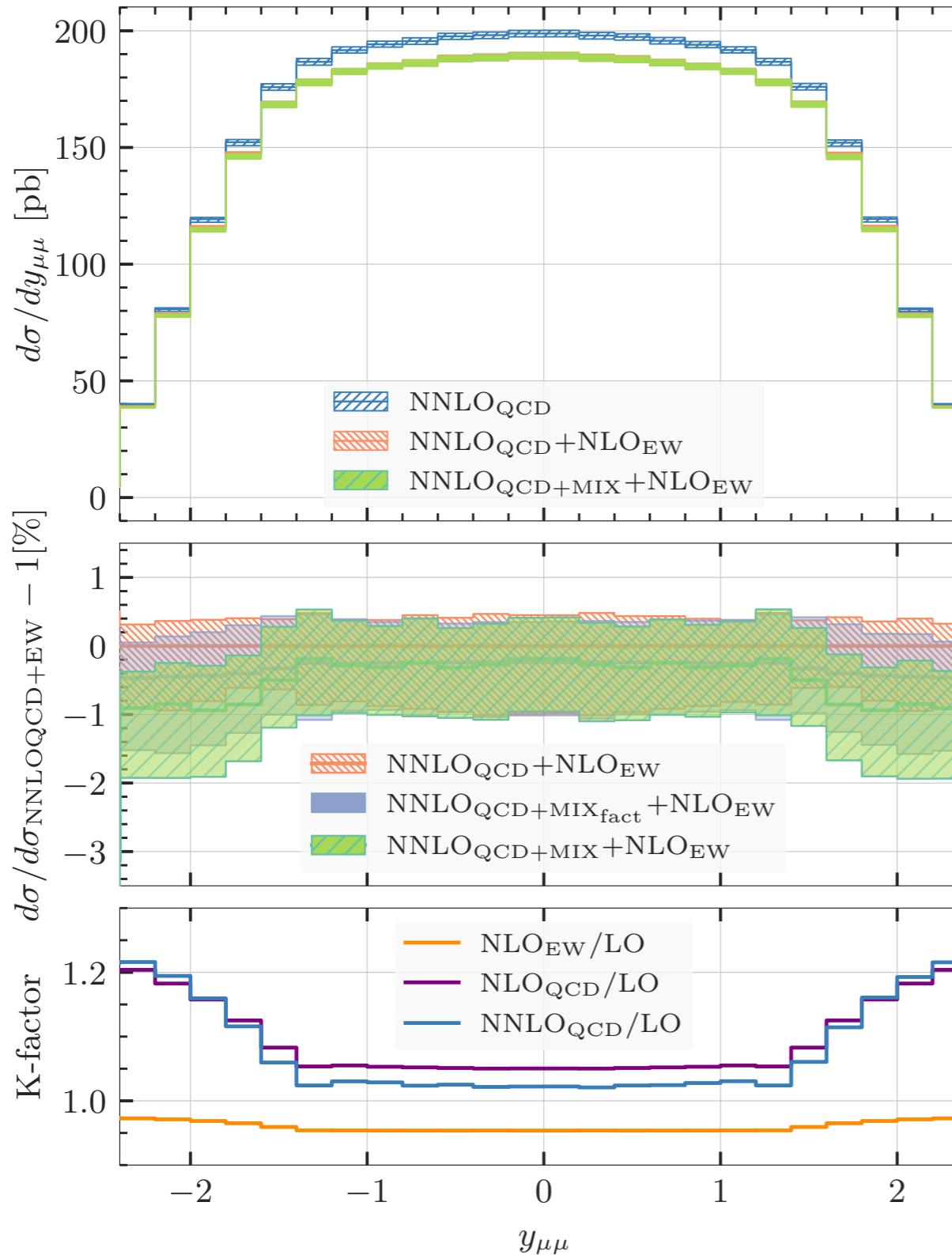
- Final-state soft singularities are regularised by the **muon mass** m_μ ;
- Process implemented in the **MATRIX** framework.



<https://matrix.hepforge.org>

DIFFERENTIAL RESULTS

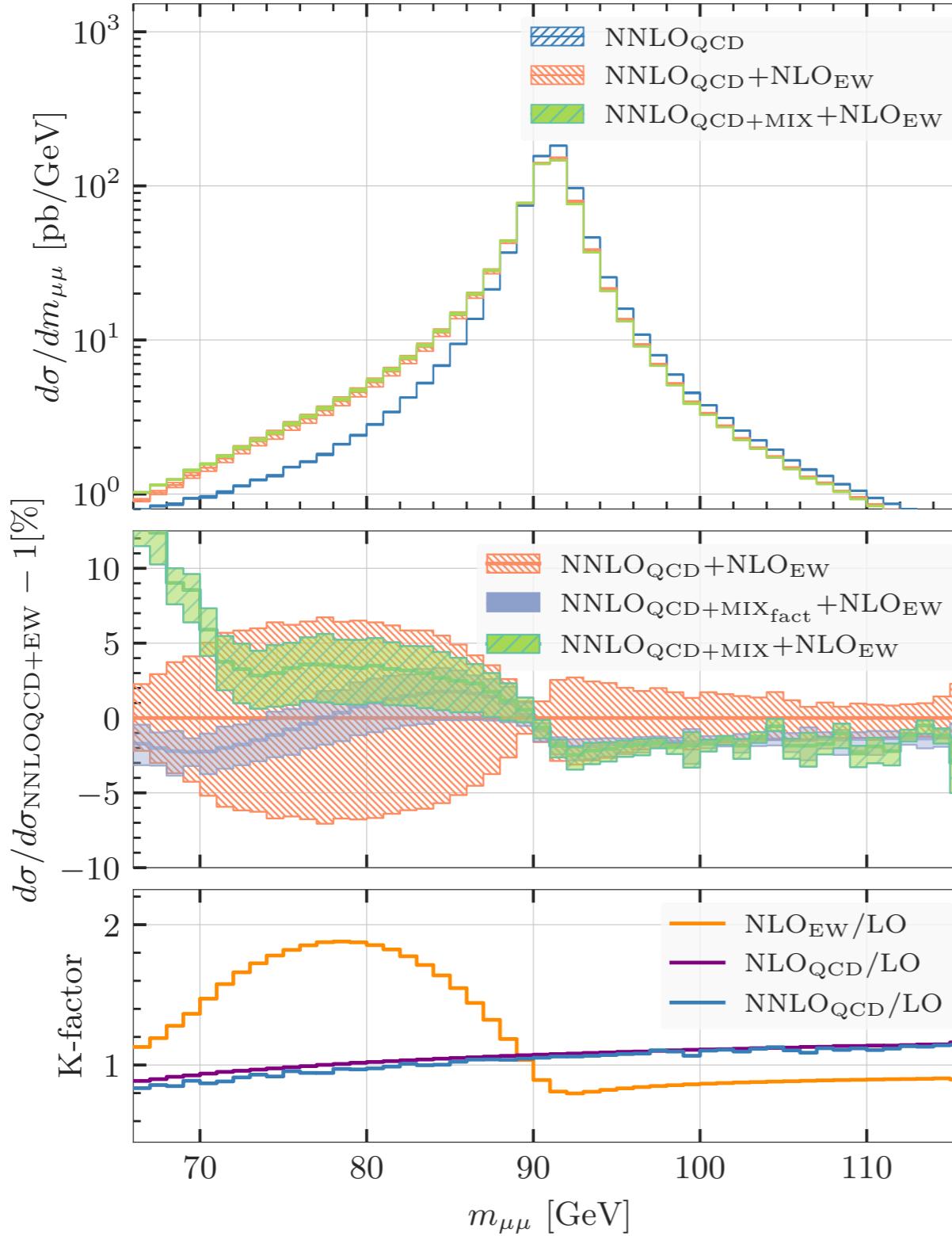
[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]



- We compare exact results MIX with the **factorised approximation** MIX_{fact}
- $$\frac{d\sigma_{\text{fact}}^{(1,1)}}{dX} = \left(\frac{d\sigma^{(1,0)}}{dX} \right) \times \left(\frac{d\sigma_{q\bar{q}}^{(0,1)}}{dX} \right) \times \left(\frac{d\sigma_{\text{LO}}}{dX} \right)^{-1};$$
- We observe a **shape distortion**: the strongest effect comes from NNLO QCD, but it is **emphasised** by the **mixed corrections**;
- Effect of the mixed corrections up to the **$\mathcal{O}(1\%)$ level**;
- Shape distortions are relevant for **PDF determinations**.

DIFFERENTIAL RESULTS

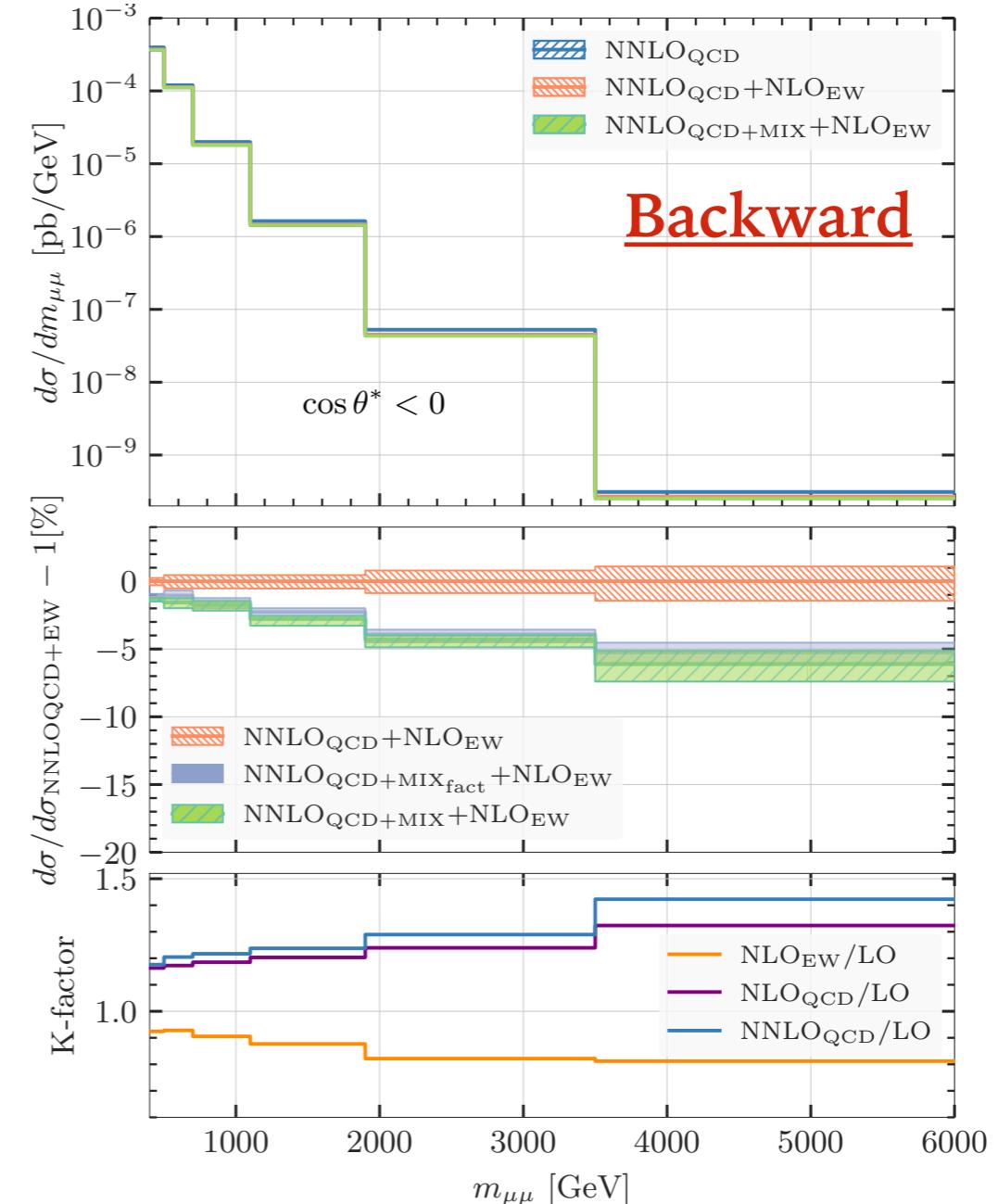
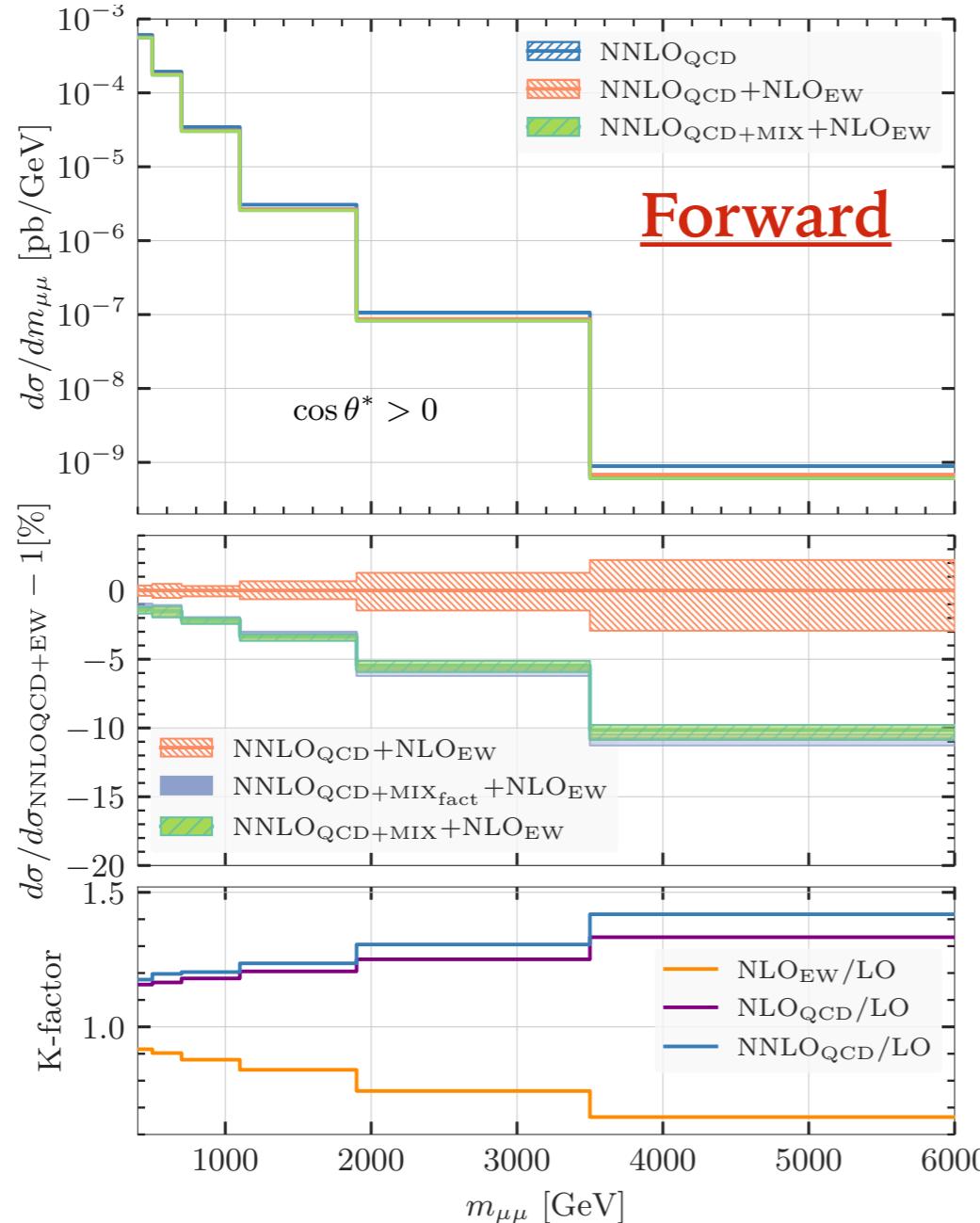
[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]



- **Mild** shape distortion from QCD corrections;
- **Strong** shape distortion from EW corrections, which generate a **radiative tail** before the peak;
- The **mixed corrections** provide a non-trivial **shape distortion**;
- The **factorised approximation** does **not** give a good description of the resonant region.

FORWARD-BACKWARD ASYMMETRY

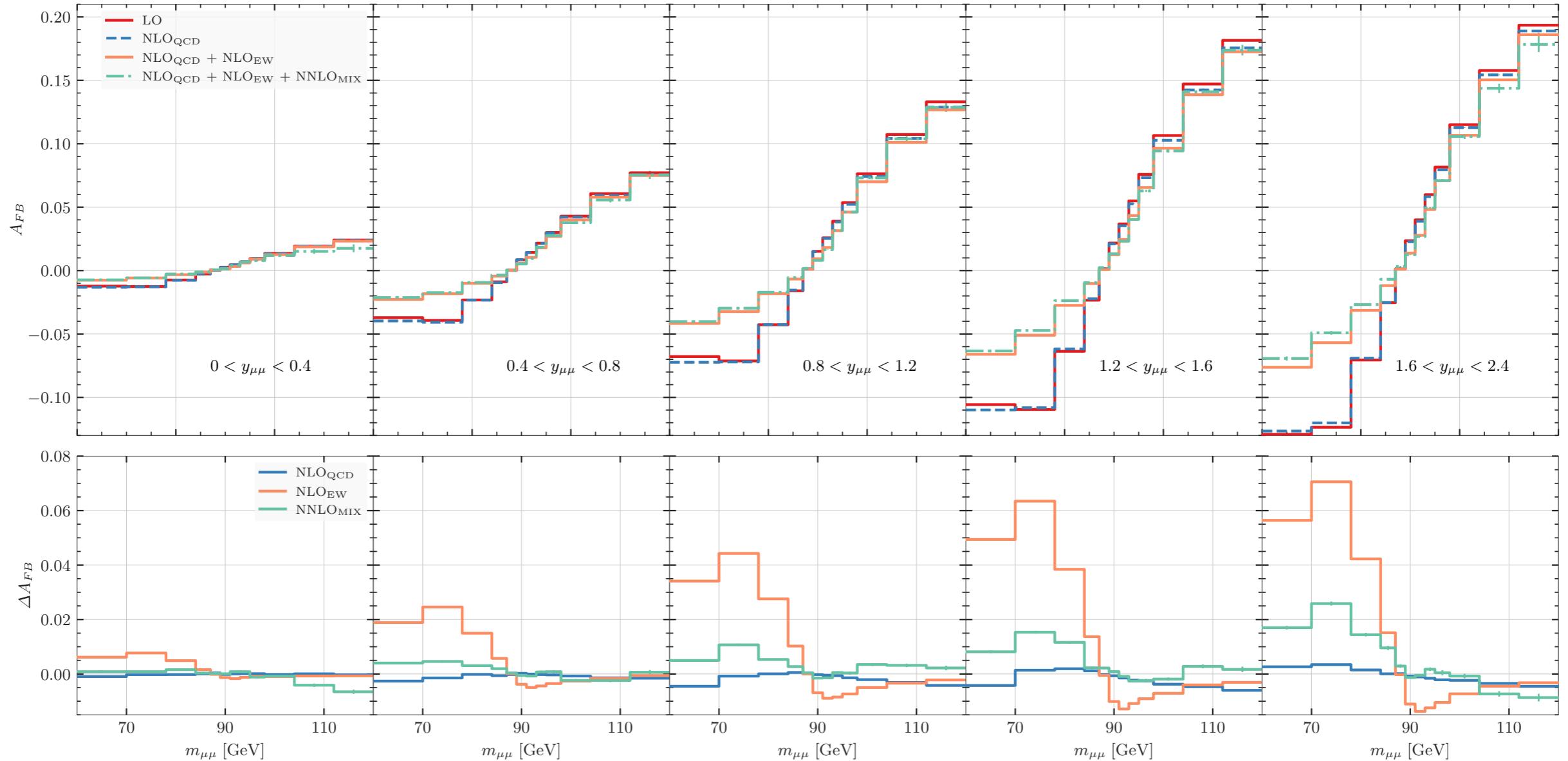
[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]



$$A_{FB}(m_{\mu\mu}) = \frac{F(m_{\mu\mu}) - B(m_{\mu\mu})}{F(m_{\mu\mu}) + B(m_{\mu\mu})}$$

FORWARD-BACKWARD ASYMMETRY

[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]



- Significant effect from **mixed corrections** to the **forward-backward asymmetry**;
- This measurement has a direct impact on $\sin^2 \theta_{\text{eff}}^\ell$.
- We apply the cuts used by the CMS collaboration for the extraction of $\sin^2 \theta_{\text{eff}}^\ell$:
 $p_{T,1} > 25 \text{ GeV}; p_{T,2} > 15 \text{ GeV}; |y_\mu| < 2.4; 60 < m_{\mu\mu} < 120 \text{ GeV}.$

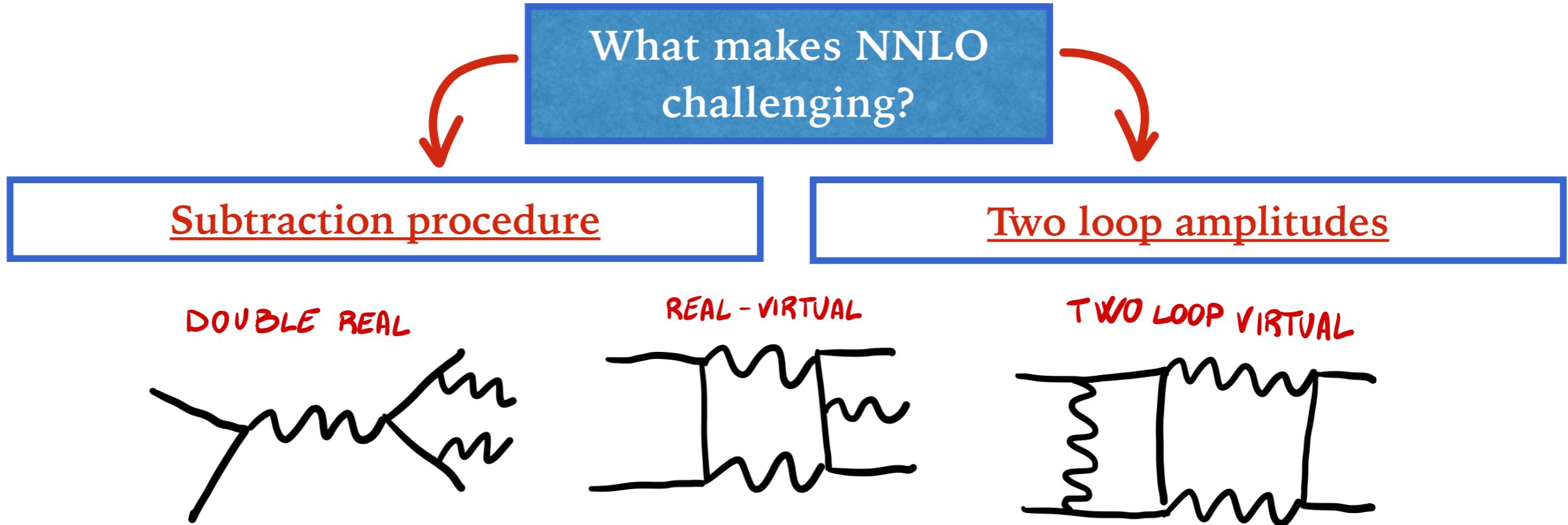
[CMS collaboration, arXiv:1806.00863]

TWO-LOOP AMPLITUDES

THEORETICAL CHALLENGES



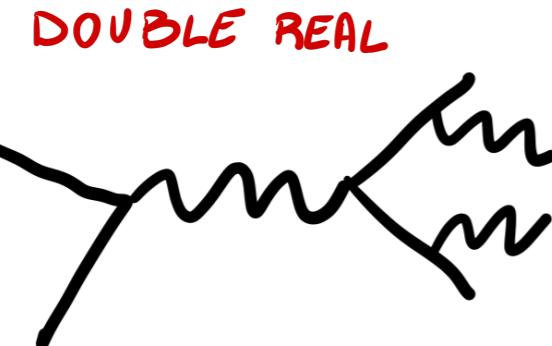
THEORY BOTTLENECKS



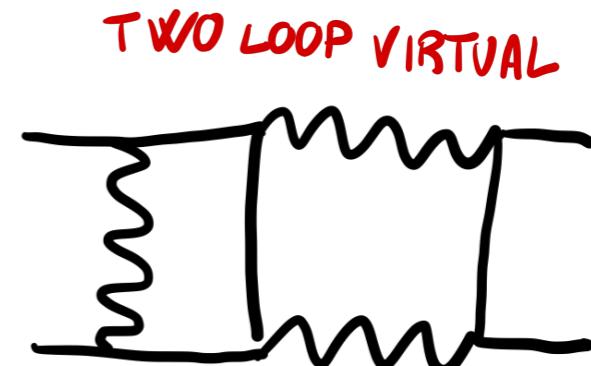
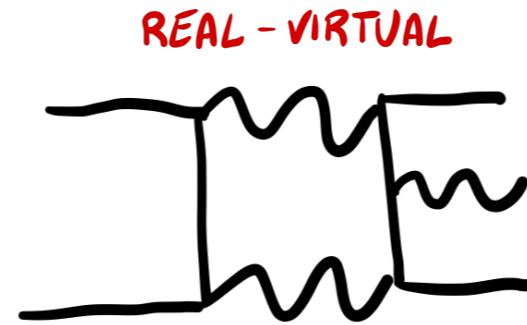
THEORY BOTTLENECKS

What makes NNLO challenging?

Subtraction procedure



Two loop amplitudes



What makes EW corrections challenging?

- additional internal **massive lines**;
additional scales in the problem (m_Z , m_W , $m_H\dots$) bring additional complications!
- treatment of γ_5 ;
how can γ_5 be consistently used in dimensional regularisation?
- need for the **complex mass scheme**;
requires to analytically continue the master integrals on the complex plane!

OUR WORKFLOW

THE BUILDING BLOCKS



OUR WORKFLOW

$$|\mathcal{M}|^2 = \sum_k c_k I_k$$

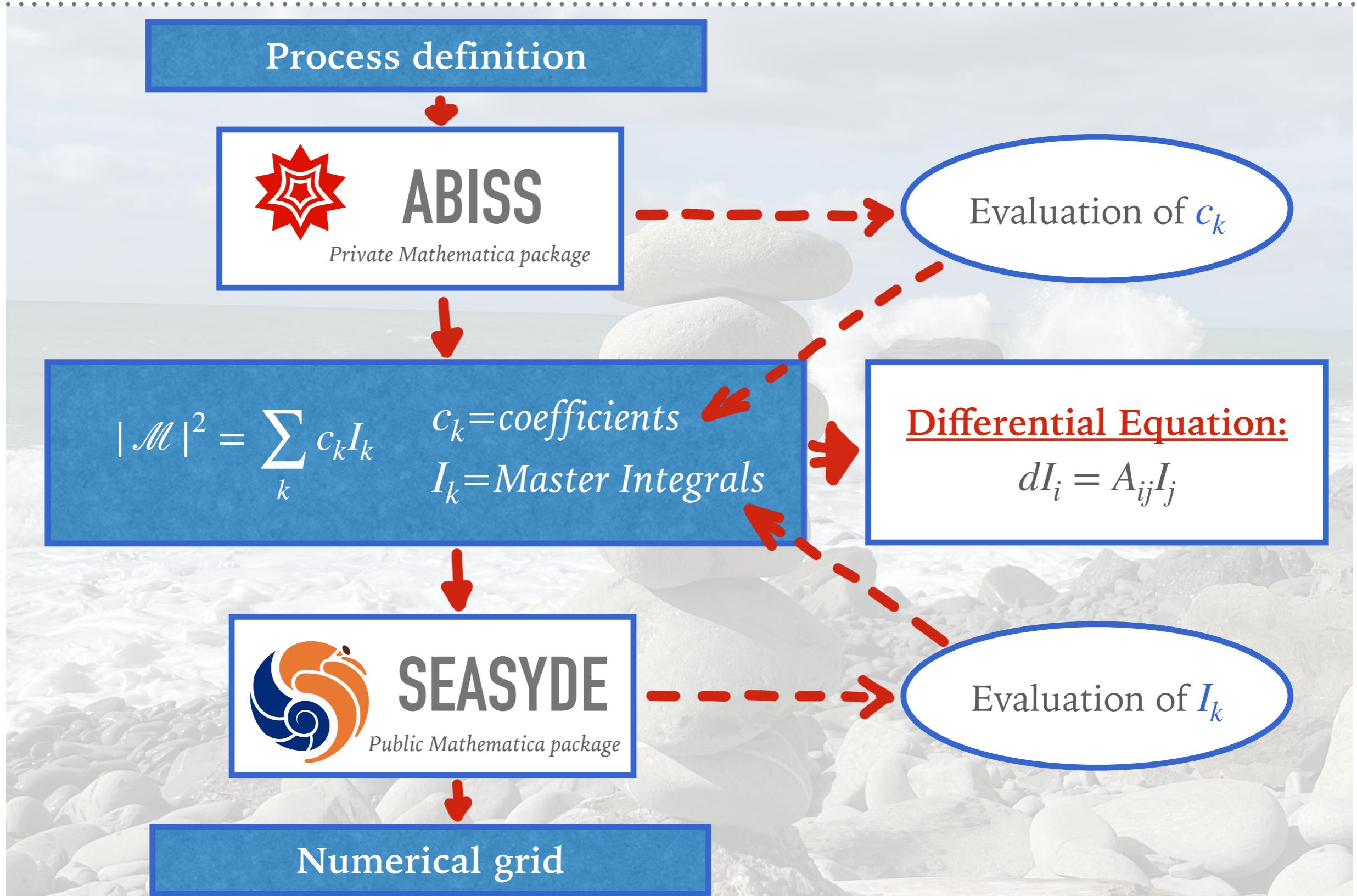
c_k =coefficients
 I_k =Master Integrals

Differential Equation:

$$dI_i = A_{ij} I_j$$



OUR WORKFLOW





SEASYDE - EVALUATING I_k

[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2205.03345]

Numerical Result

The result of the master integral is provided as a numerical **grid**.

Analytical Result

The result of the master integral can be expressed in closed form as a combination of elementary and special functions, whose **power expansion is known**.

Semi-Analytical Result

The result of the master integral can be expanded as a power series at every point of its domain, but without any additional functional relations.

We solve the master integral with **series expansion!**



SEASYDE - EVALUATING I_k

[*T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2205.03345*]

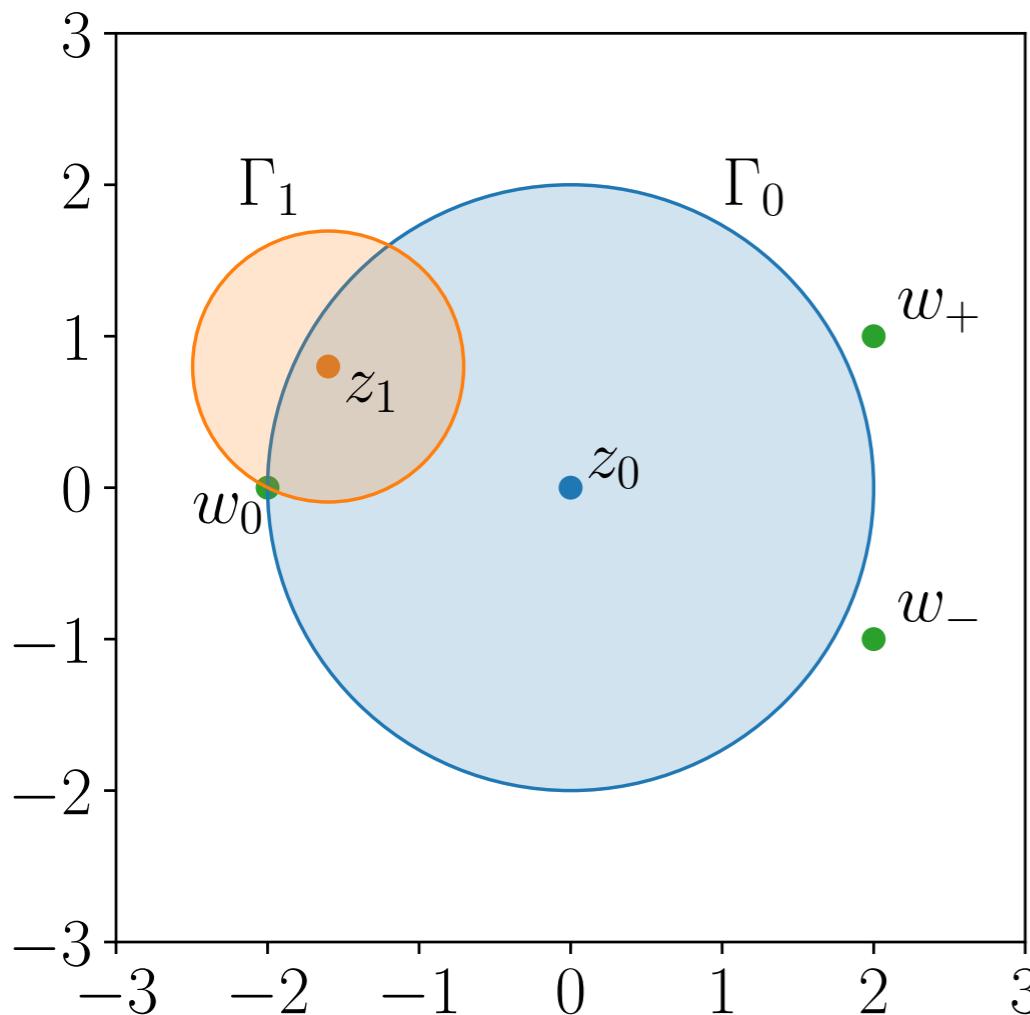


SEASYDE

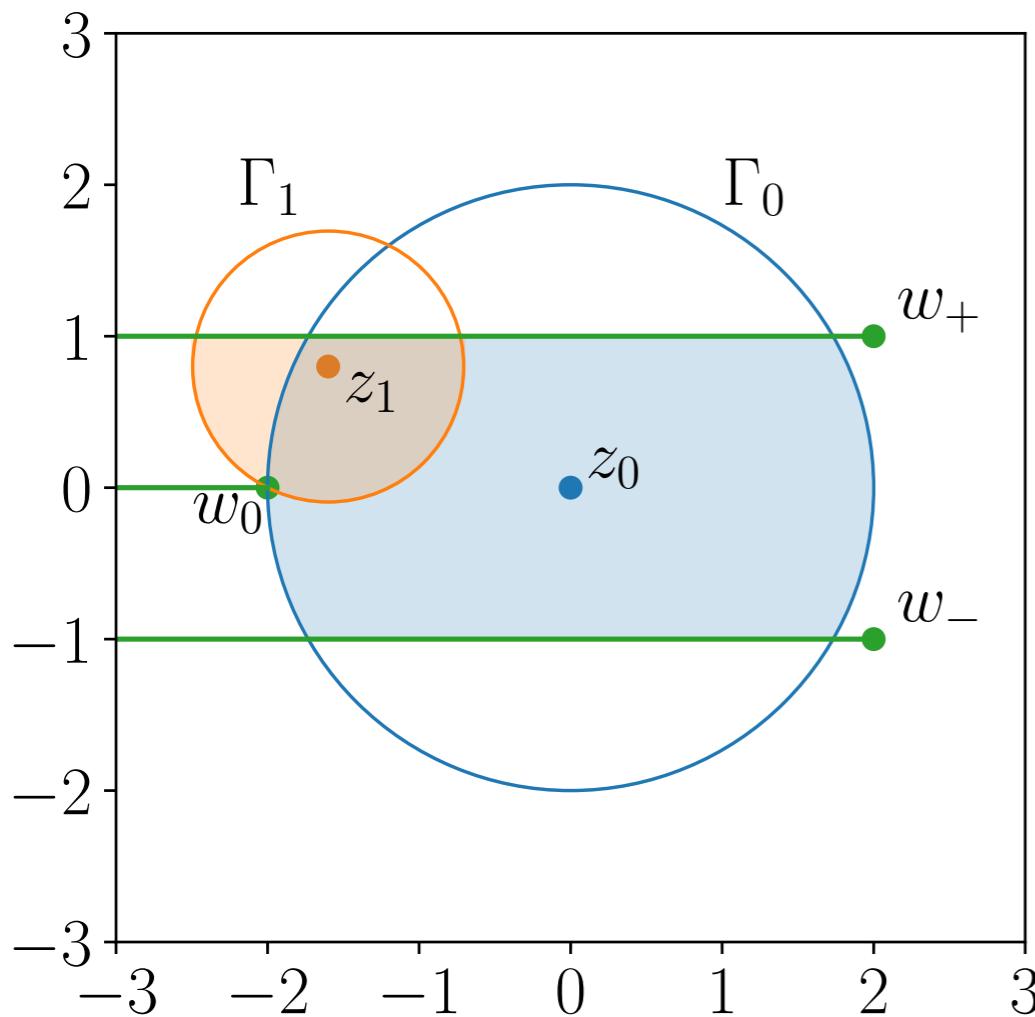
The latest version of **SEASYDE** can be
downloaded from:
github.com/TommasoArmadillo/SeaSyde

- **SEASYDE**(Series Expansion Approach for System of Differential Equations) is a **MATHEMATICA package** for solving the system of differential equation, associated to the Master Integrals of a given topology.
- **SEASYDE** can handle any **system of coupled differential equations**.
- Method implemented in the Mathematica package DiffExp for real kinematic variables [*F. Moriello, arXiv:1907.13234*], [*M. Hidding, arXiv:2006.05510*] (see also AMFLOW [*X. Liu and Y.-Q. Ma, arXiv: 2201.11669*]; LINE[*R.M. Prisco, J. Ronca, F. Tramontano, arXiv:2501.01943*])
- By using the auxiliary mass flow, AMFLOW can provide **boundary conditions**.

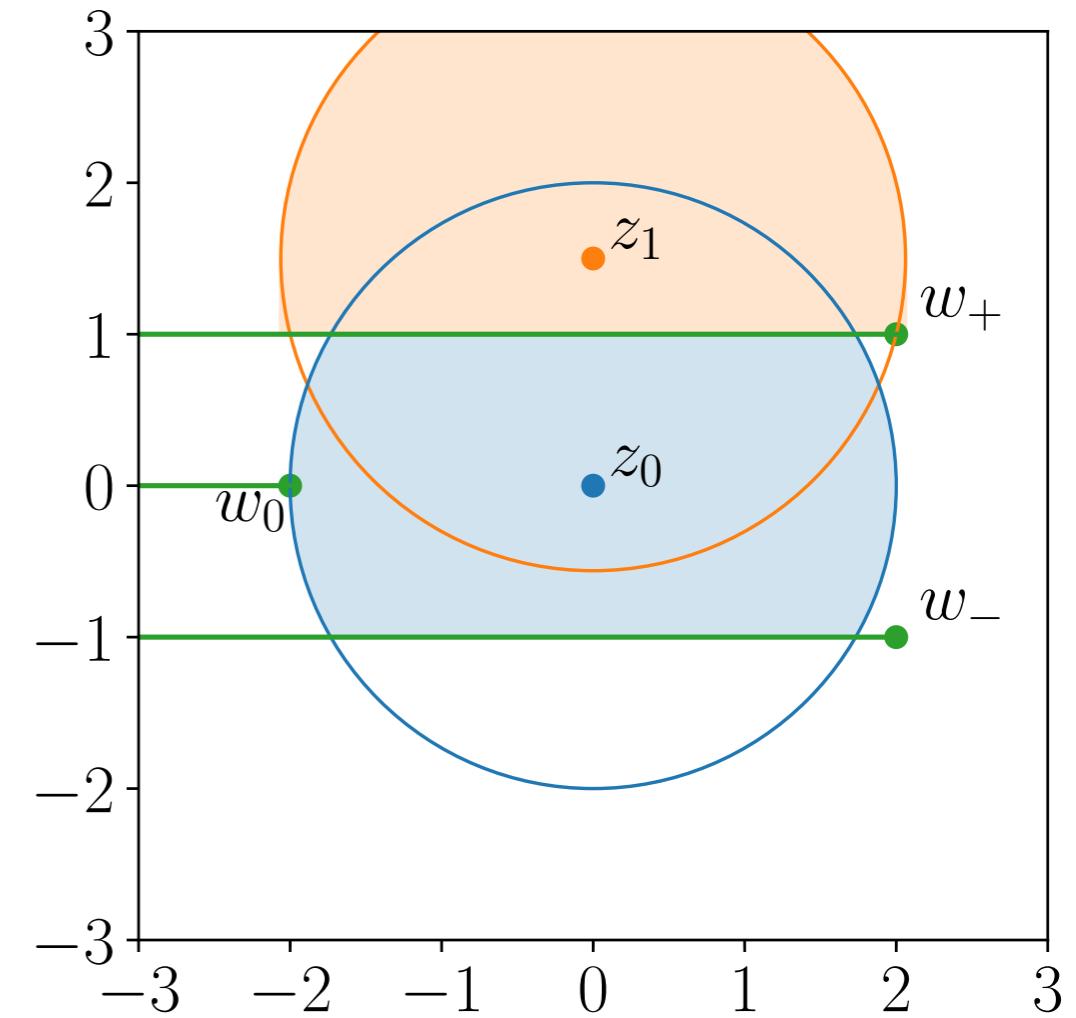
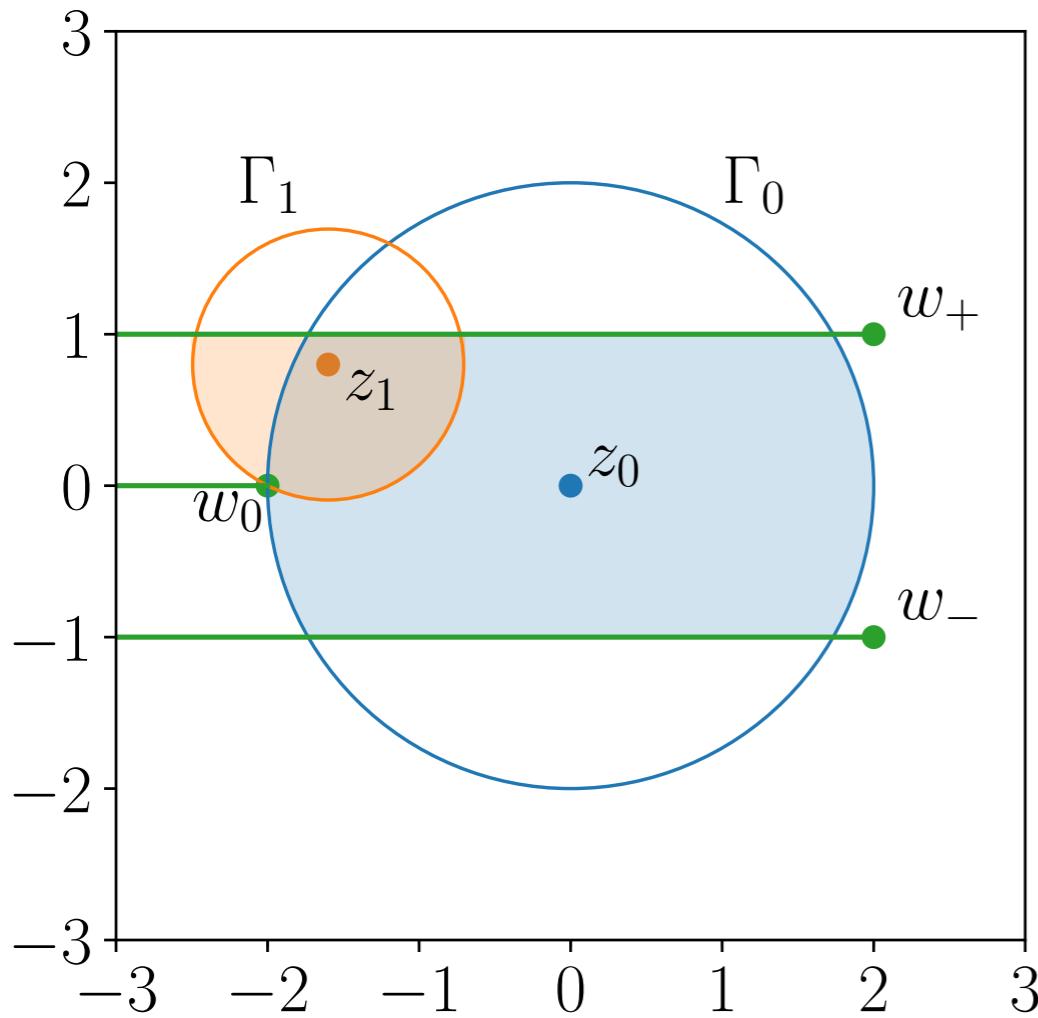
- We generalised the series expansion method to **arbitrary complex-valued masses** → **complex plane** of the kinematical invariants!
- The radius of convergence of the series is limited by the presence of **poles**;
- “Transport” of the boundary conditions need to consider **branch-cuts**.



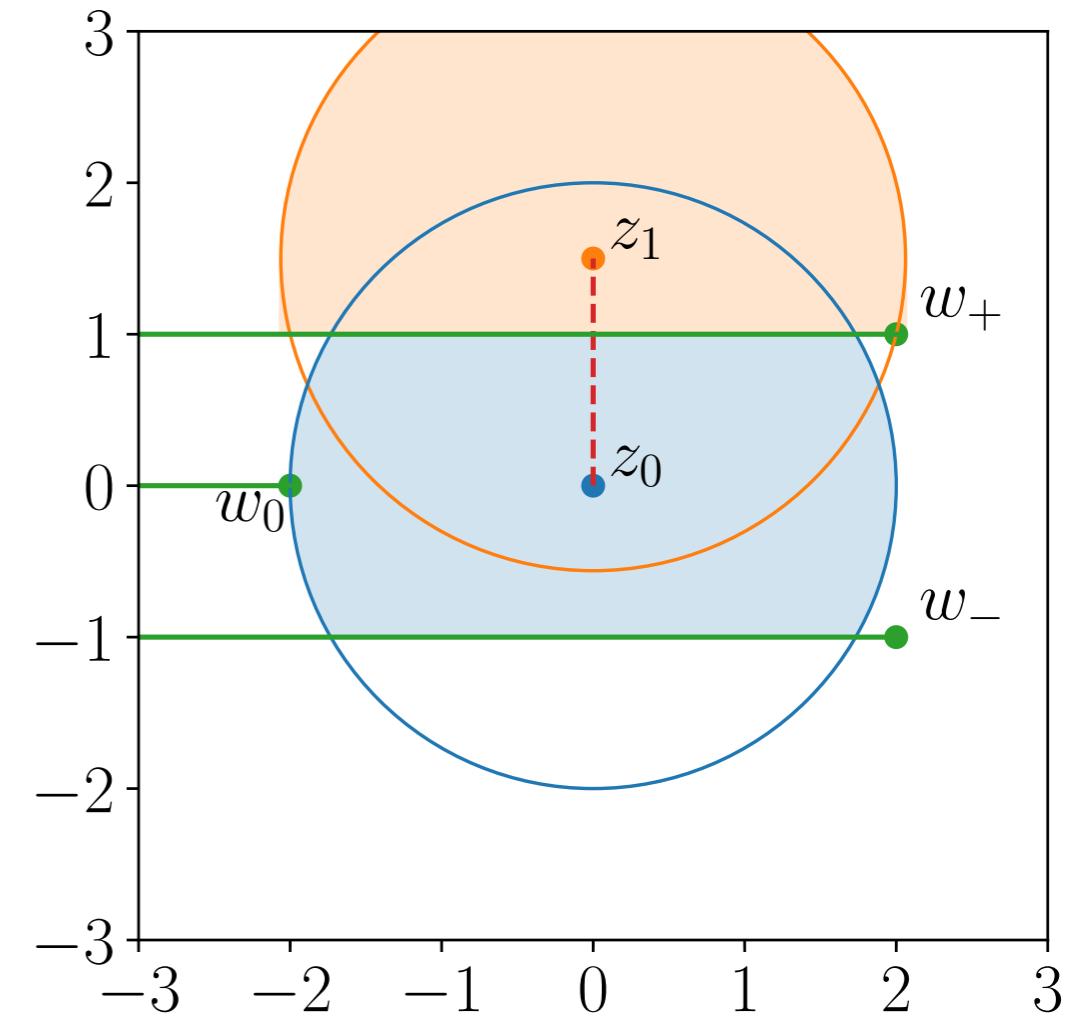
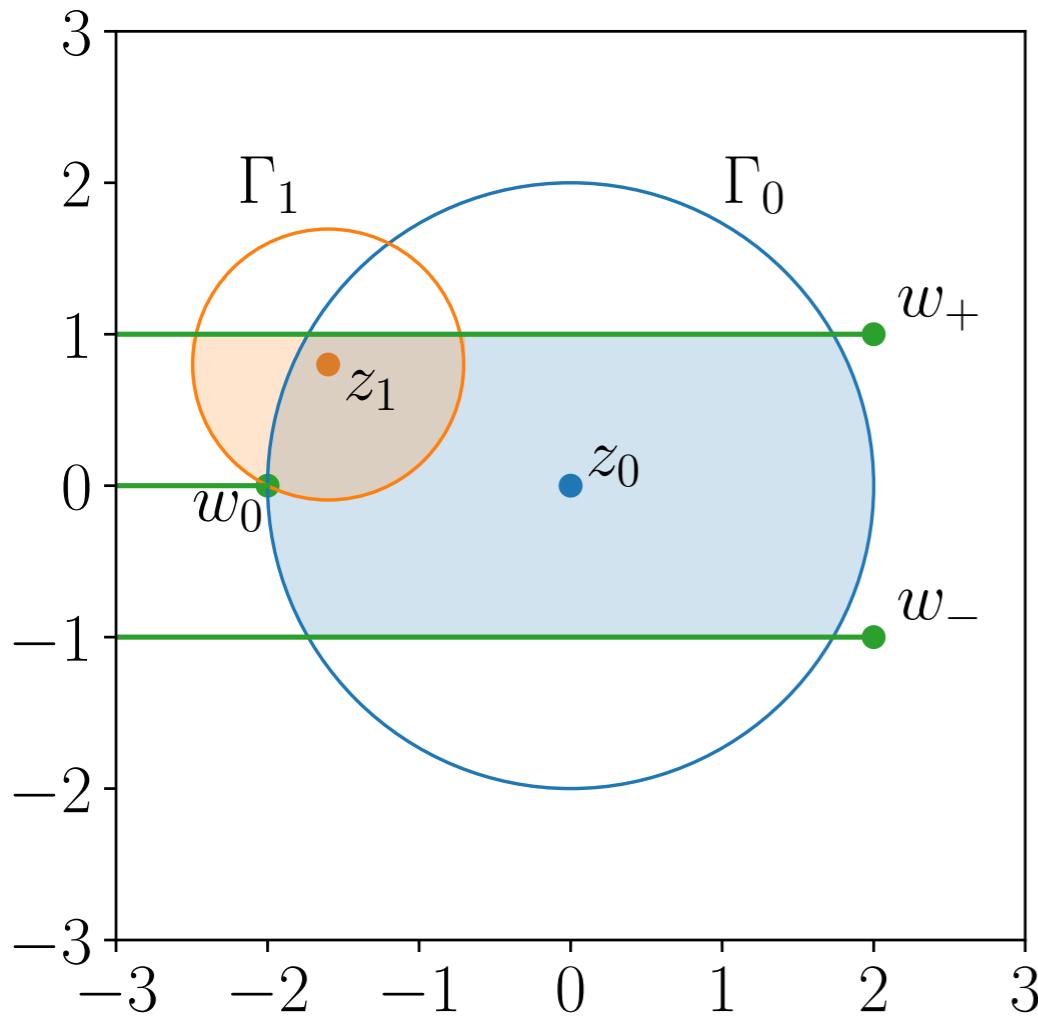
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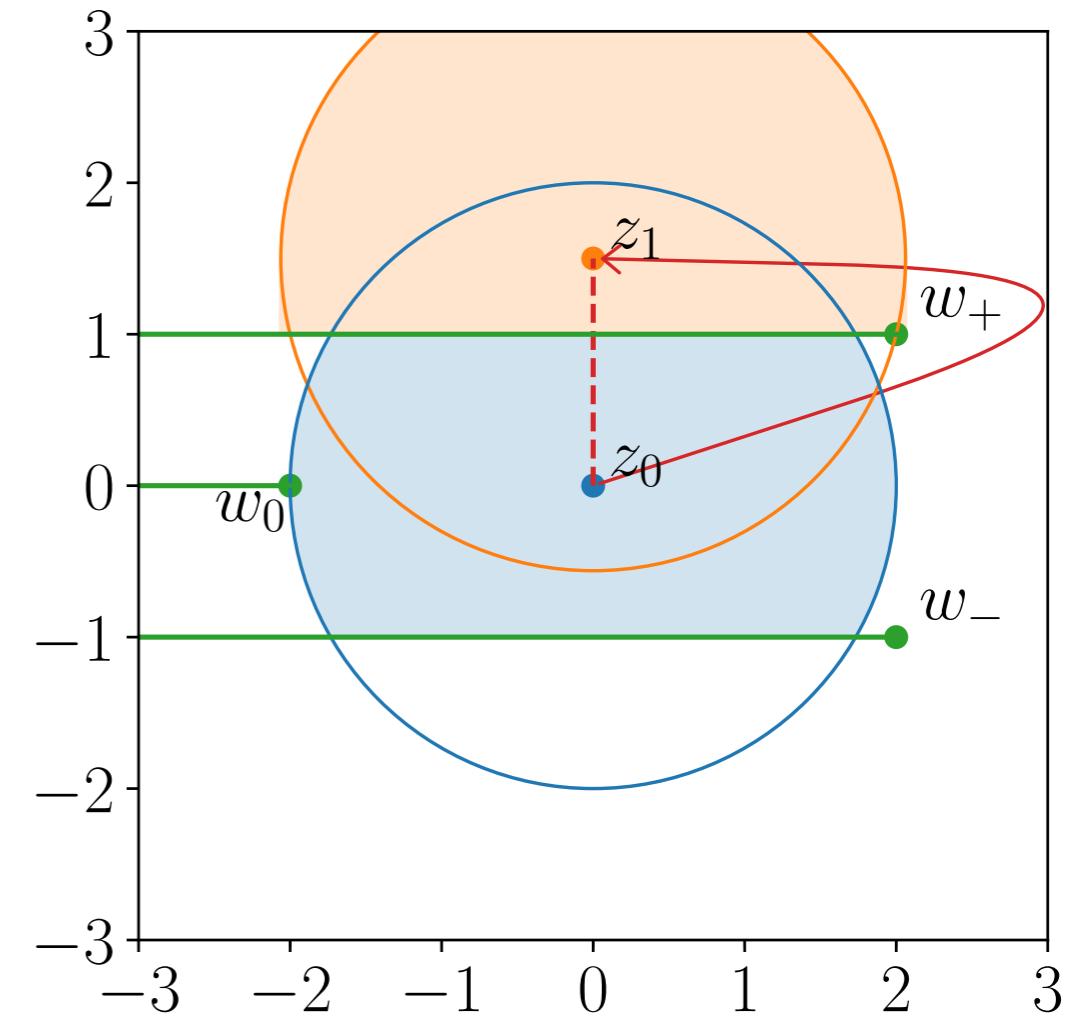
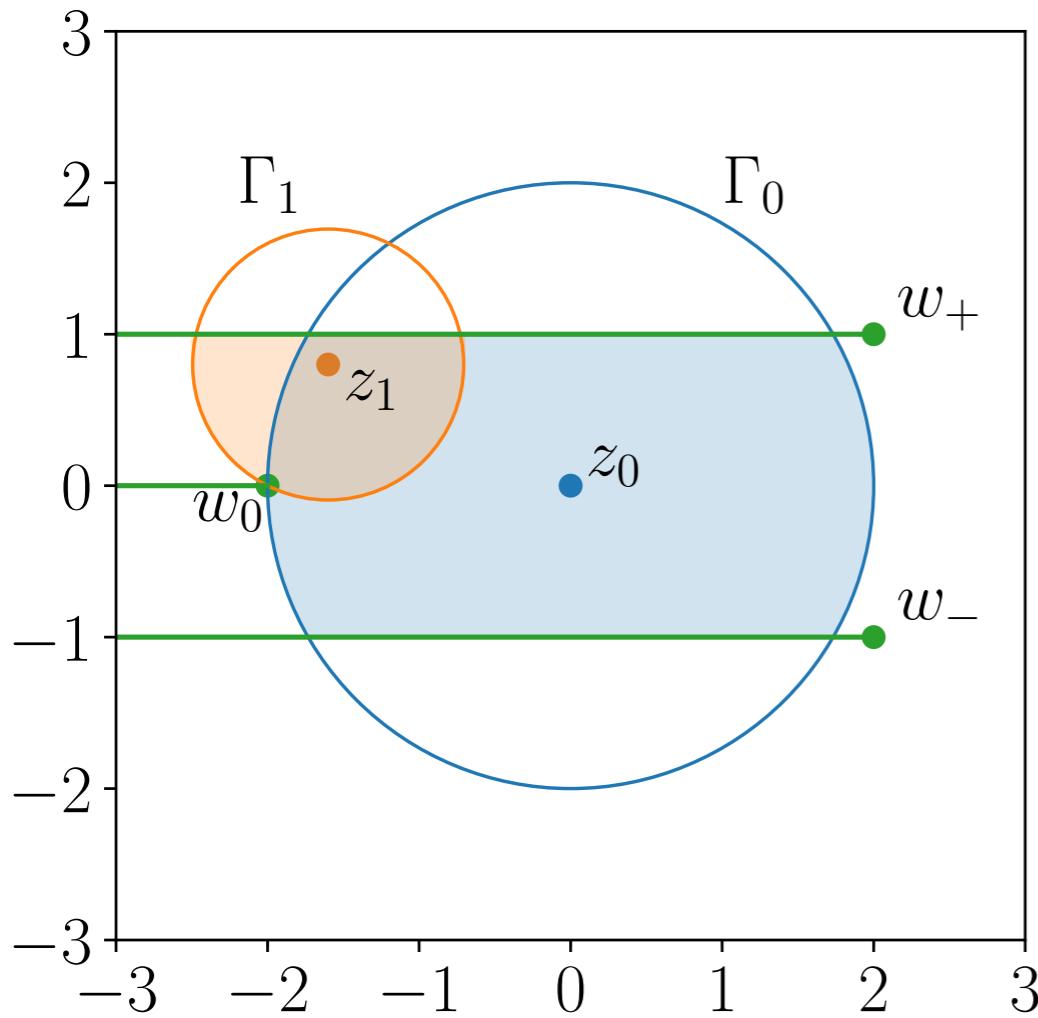
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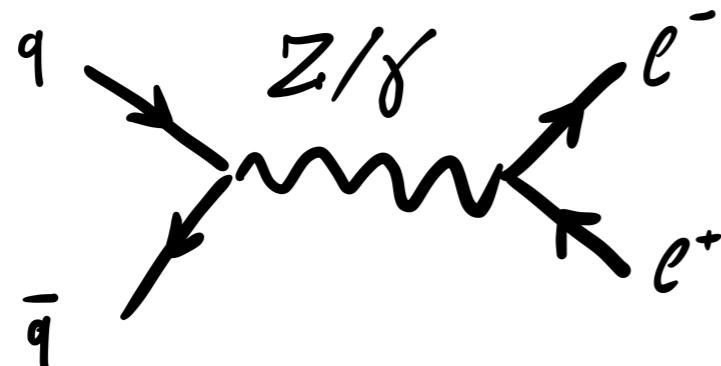


COMPUTATIONAL FRAMEWORK

[T. Armadillo, R. Bonciani, SD, N.Rana, A.Vicini, arXiv:2201.01754]

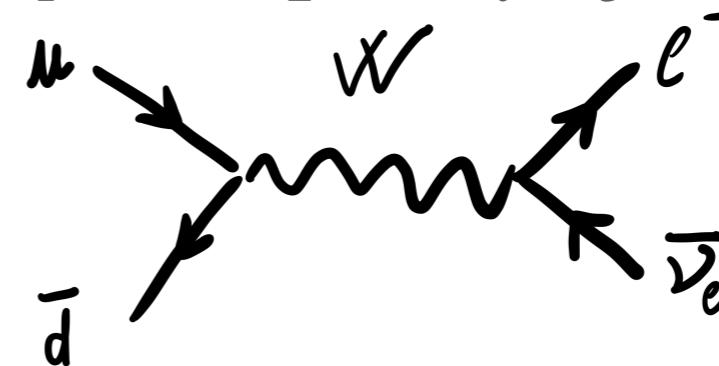
NCDY

$$q(p_1) + \bar{q}(p_2) \rightarrow l^-(p_3) + l^+(p_4)$$



CCDY

$$u(p_1) + \bar{d}(p_2) \rightarrow \nu_e(p_3) + l^+(p_4)$$



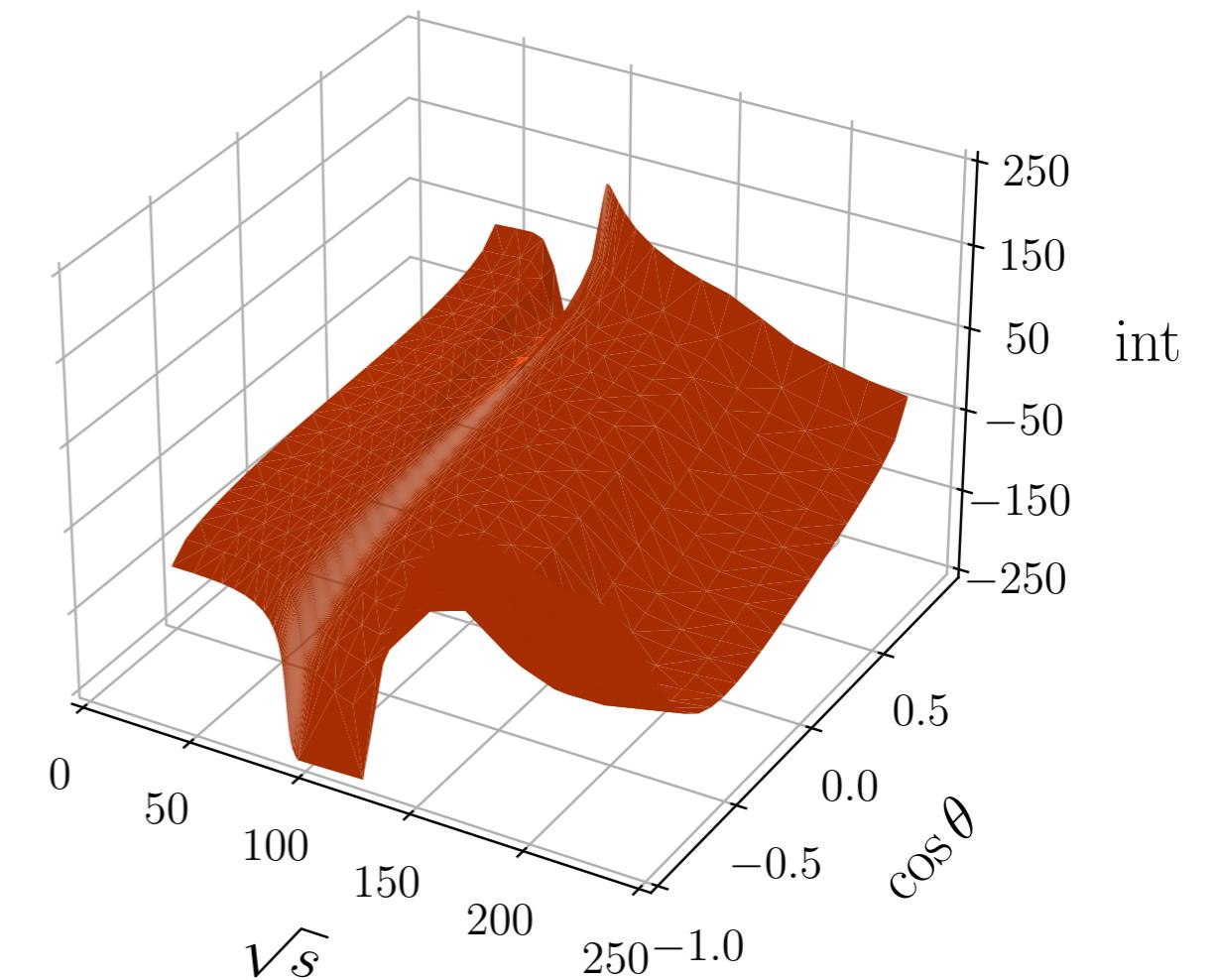
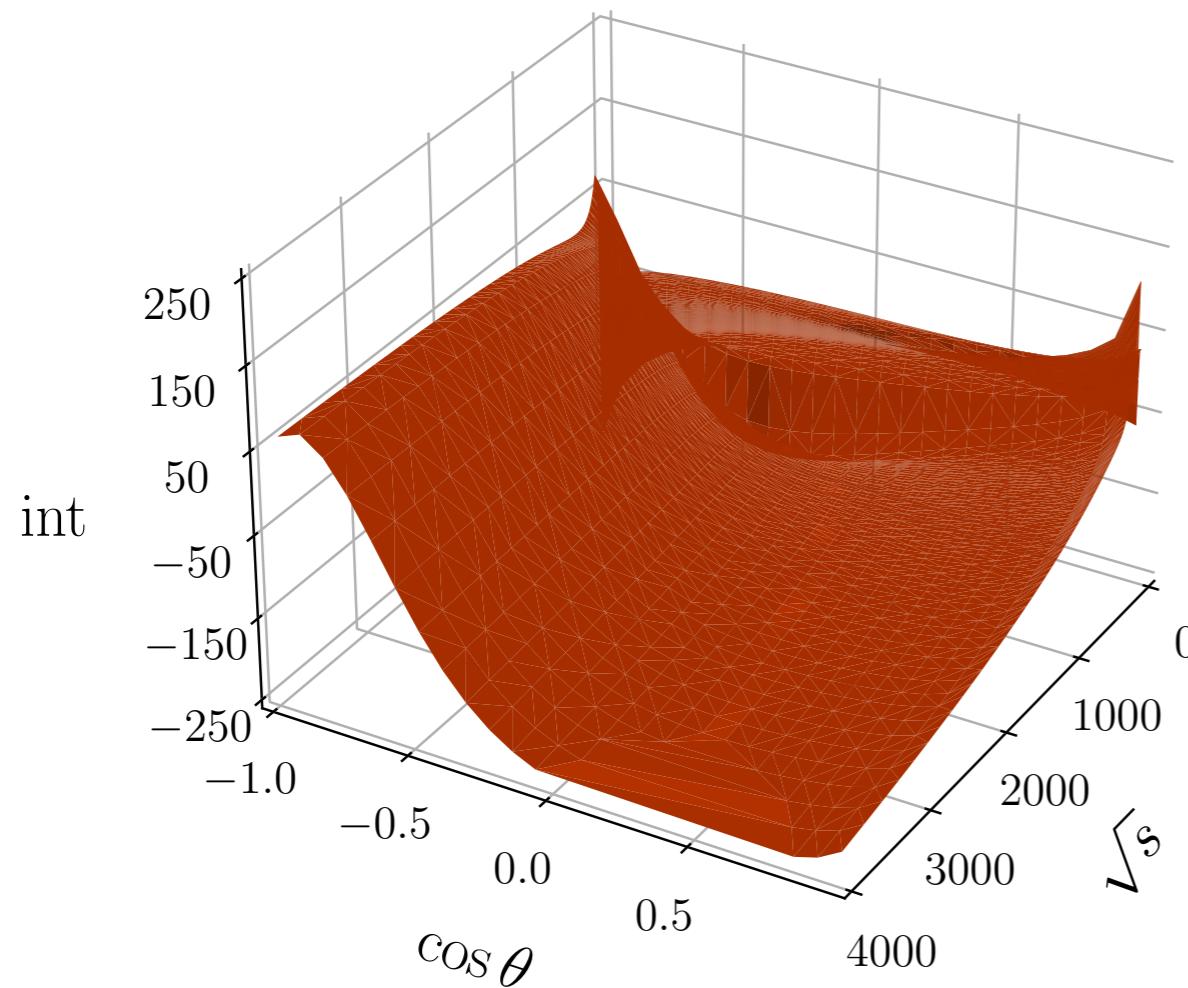
- IR singularities handled by **q_T-subtraction formalism** (straightforward implementation of **any other framework** by replacing the **subtraction operator**);
- Final-state collinear singularities regularised by the **lepton mass**;
- **small lepton mass limit**: consider the ratio m_l/\sqrt{s} and keep only **logarithmic terms** $\sim \log(m_l/\sqrt{s})$.

- When dealing with intermediate unstable particles, such as W and Z , it is useful to perform the calculations in the **complex-mass scheme**;
- We introduce the complex mass $\mu_V^2 = m_V^2 - i\Gamma_V m_V$ for both the Z and W bosons.

NCDY - NUMERICAL GRIDS

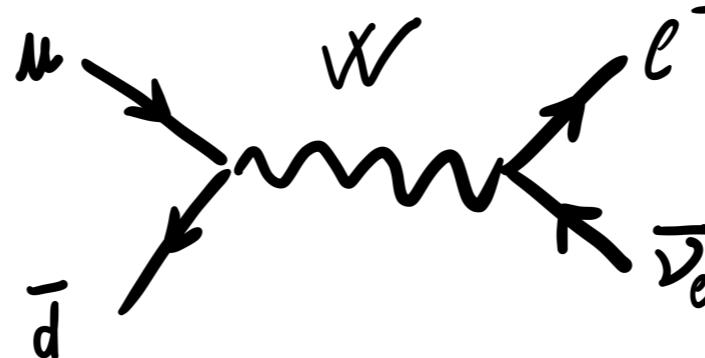
[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2201.01754]

- After subtracting IR and UV divergences, we obtain the **hard function**;
- **Publicly available** as a MATHEMATICA notebook;
- Subtraction of the IR poles done in the **qT-subtraction** formalism;
- Production of the grid (3250 points) required $\mathcal{O}(12h)$ on a 32-cores machine;
- Interpolation of the grid with excellent accuracy requires **negligible time**.

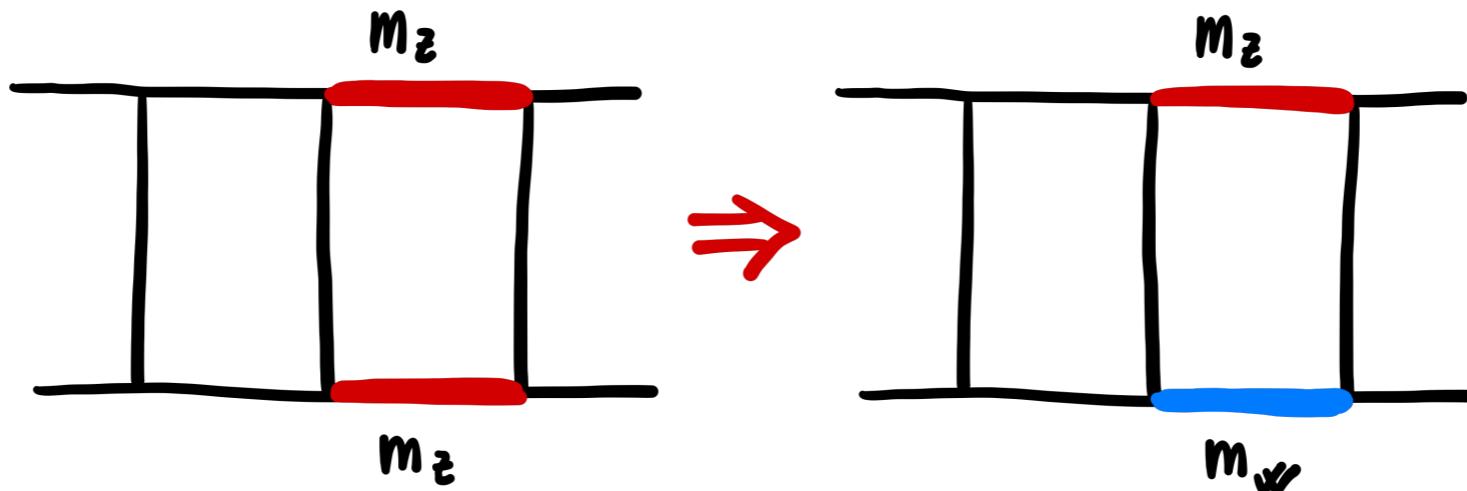


CCDY – BASIS OF MASTER INTEGRALS

[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2405.00612]



- Computationally **similar** to neutral current Drell-Yan;
- Extra complexity coming from new diagrams where **two different internal massive lines** appear:

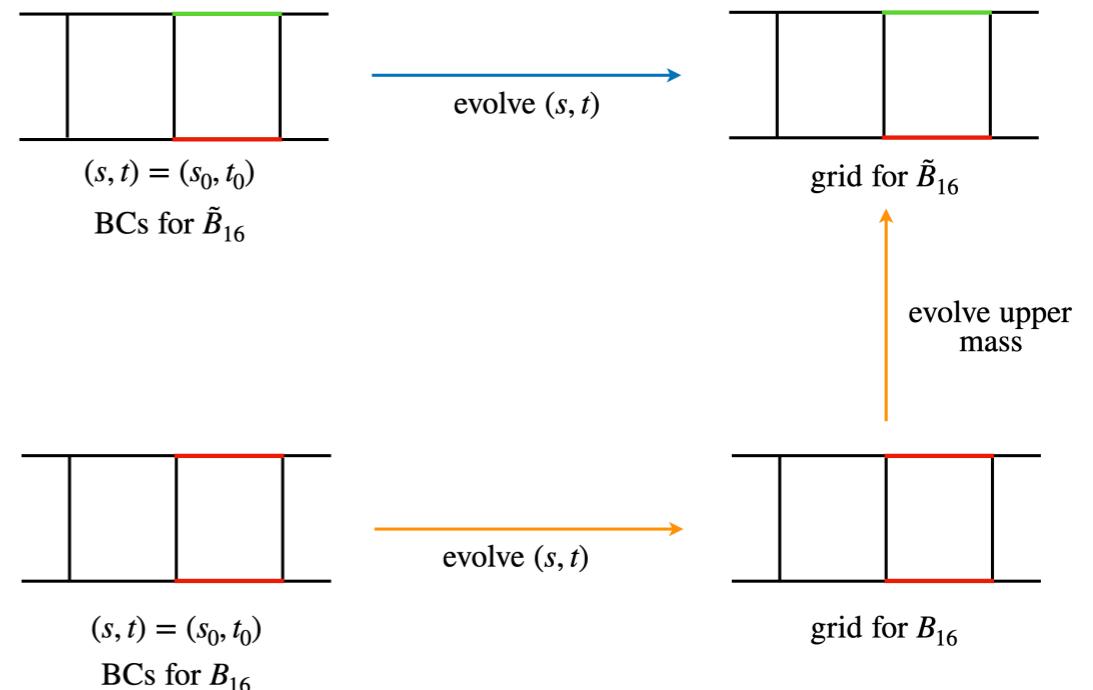


- **56** Master integrals with two internal masses:
 - 34 **identical** to the ones appearing in the neutral current case;
 - 22 **generalisation** of masters with two massive lines to the case of different masses;

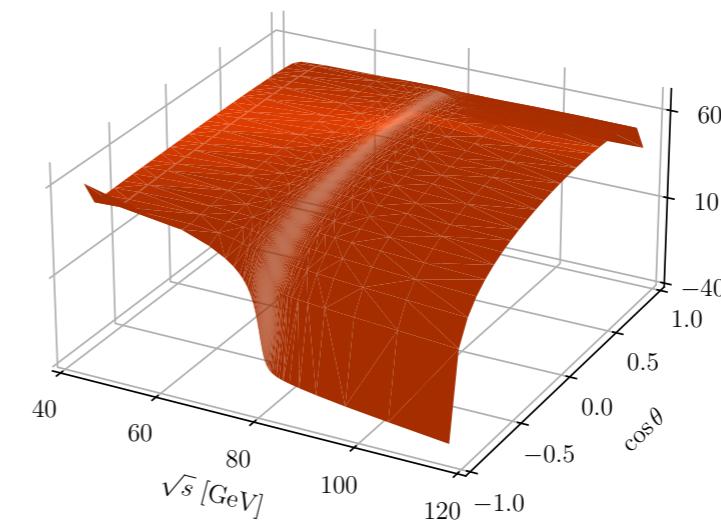
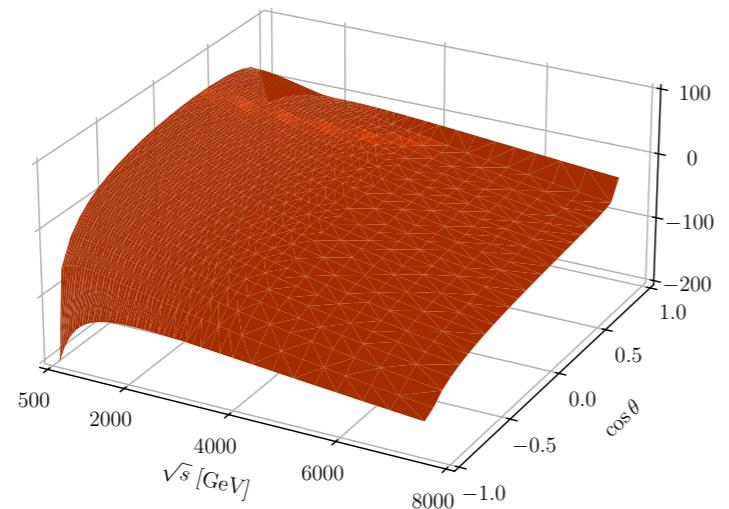
CCDY – NUMERICAL GRIDS

[T. Armadillo, R. Bonciani, SD, N.Rana, A.Vicini, arXiv:2405.00612]

- Two possible approaches:
 - Compute the boundary conditions with AMFLOW and use the **differential equations in s and t** to build the grid;
 - Use the grid of the neutral-current Drell-Yan as a boundary condition and use the **differential equation w.r.t. one of the masses**.



UV renormalisation & IR subtraction are analogous to the neutral current case.



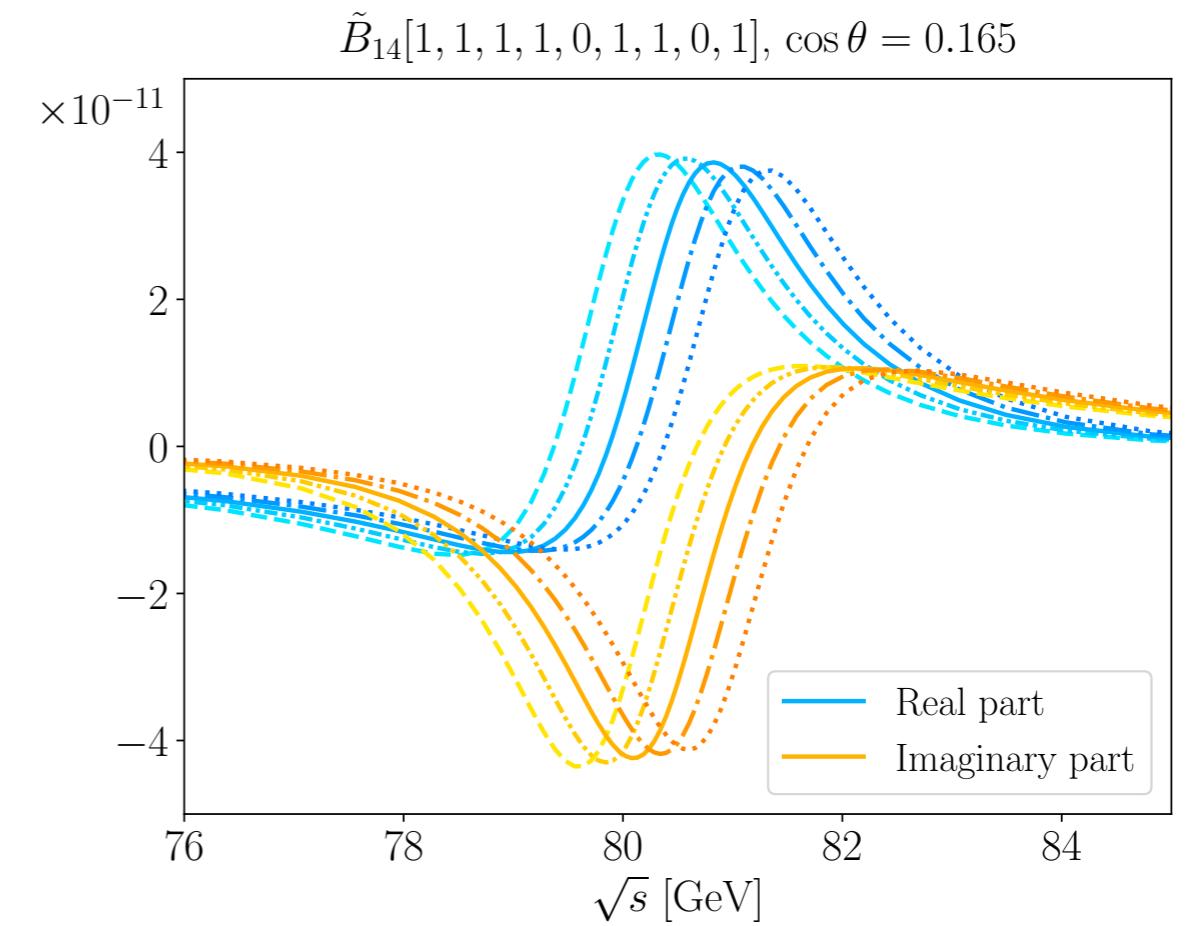
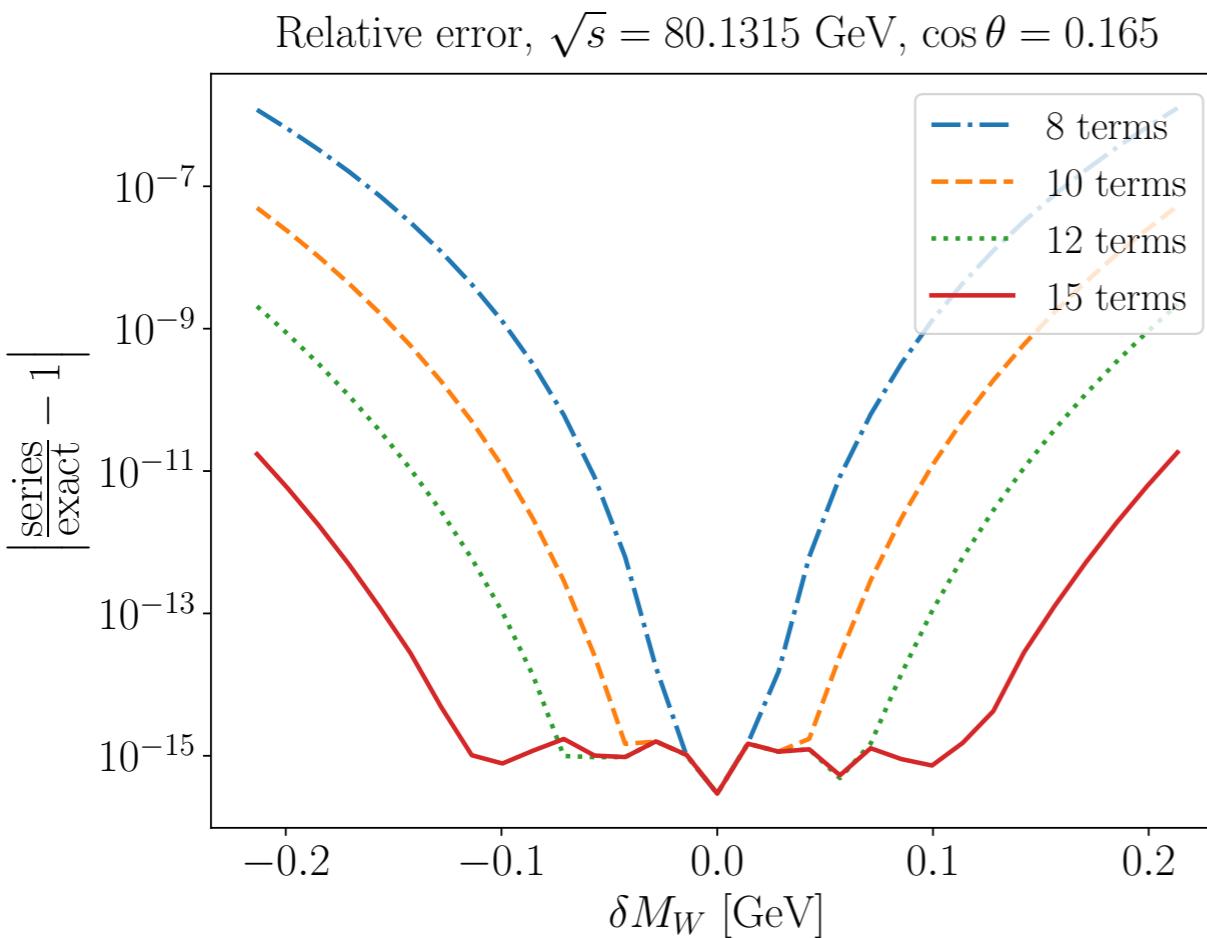
CCDY - NUMERICAL GRIDS

[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2405.00612]

- The mass evolution offers the possibility to also provide the **additional dependence on the W-boson mass**;
- Useful for frontier studies, such as M_W determination.

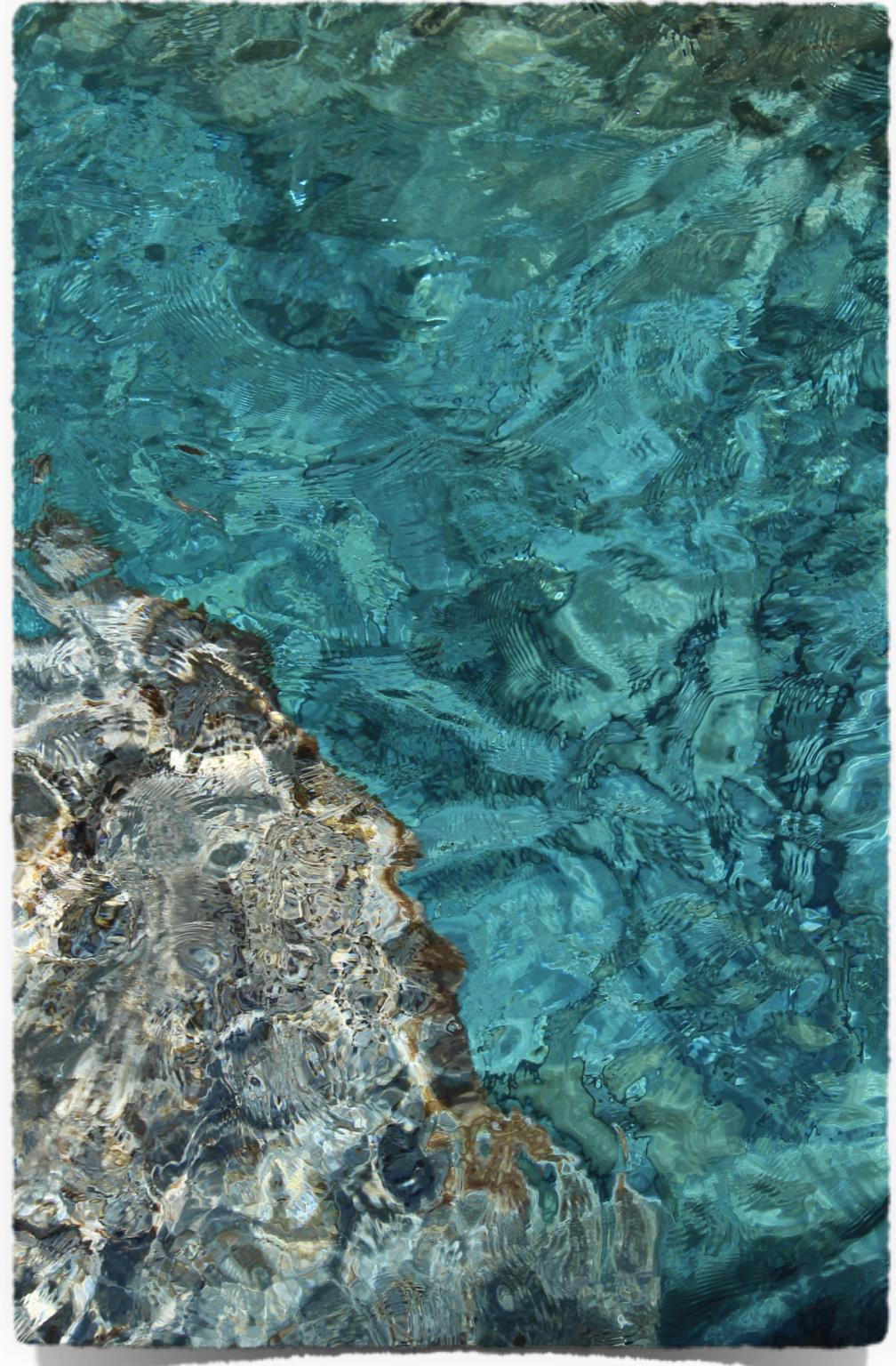
$$\delta M_W = M_W - \bar{M}_W$$

$\bar{M}_W \rightarrow$ reference value



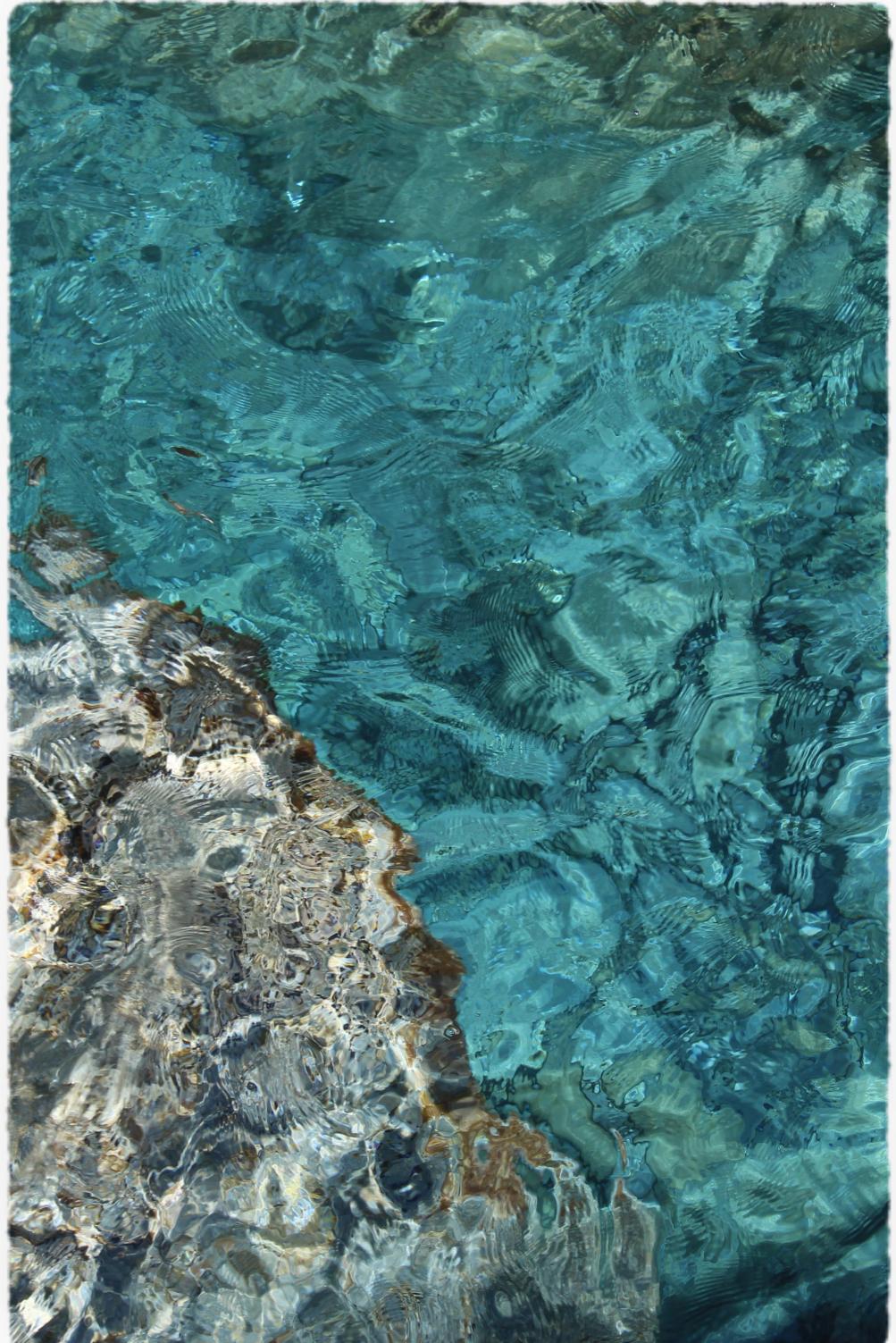
SUMMARY & OUTLOOK

- **EW corrections** will be of crucial importance for the future LHC scientific program;
- We developed a **framework** which relies on **ABISS** for the evaluation of rational coefficients and **SEASYDE** for the evaluation of the Master Integrals;
- We successfully applied our framework to the computation of **mixed QCD-EW corrections** to the Drell-Yan process.



SUMMARY & OUTLOOK

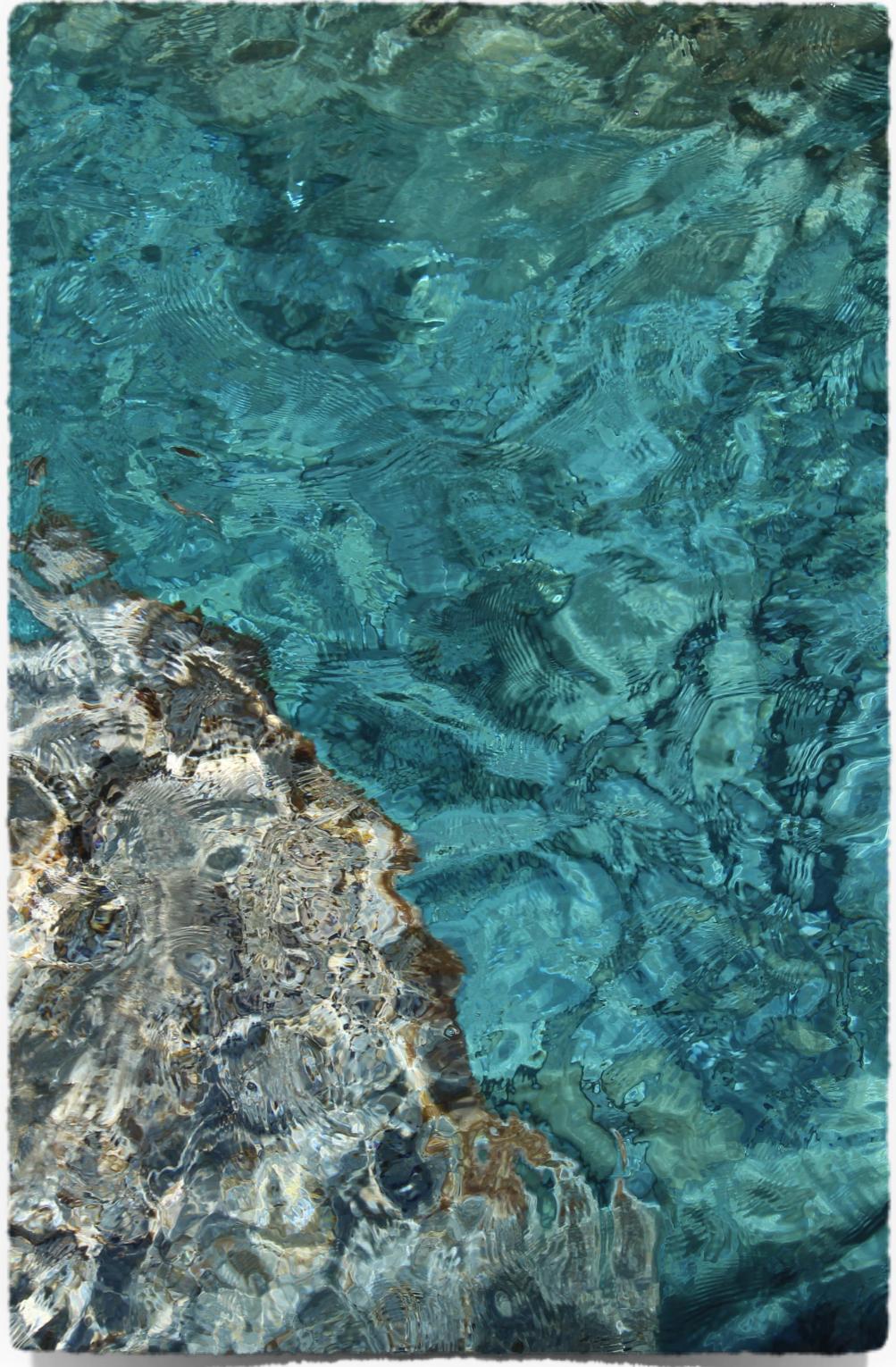
- Phenomenological studies for **charged-current Drell-Yan**;
- Towards **NNLO EW**?



SUMMARY & OUTLOOK

- Phenomenological studies for **charged-current Drell-Yan**;
- Towards **NNLO EW**?

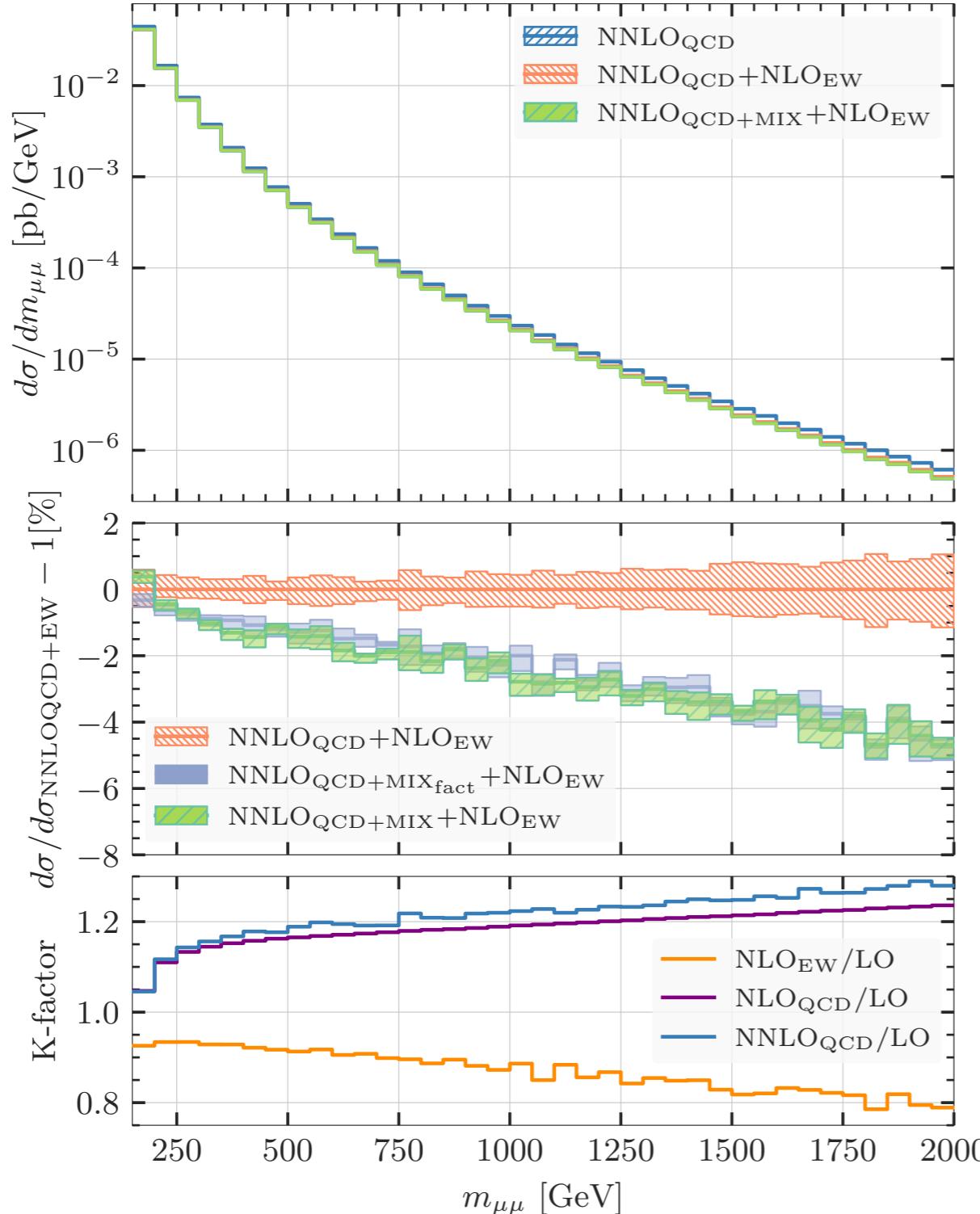
THANKS!



BACKUP SLIDES

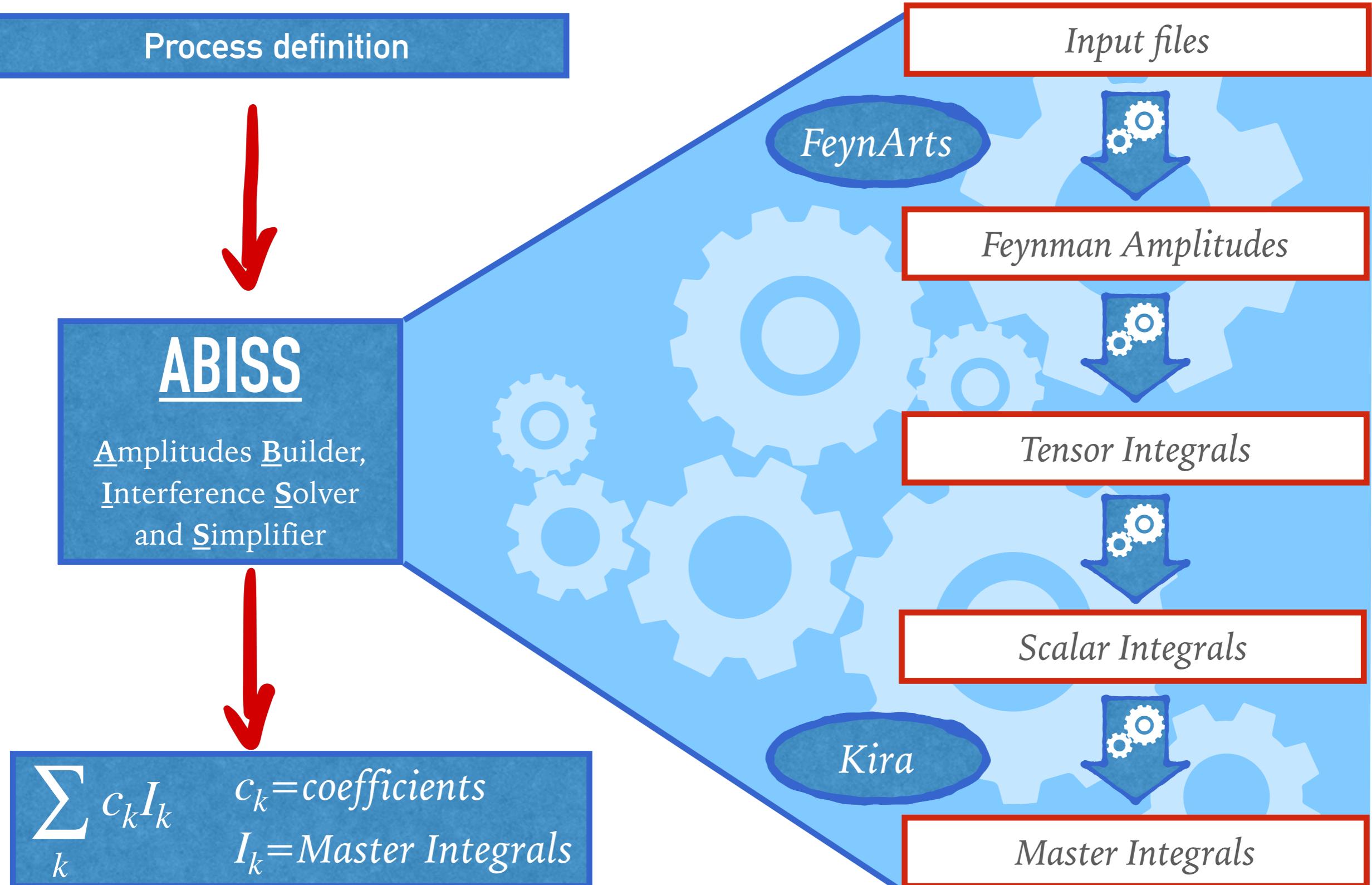
DIFFERENTIAL RESULTS

[T. Armadillo, R. Bonciani, L. Buonocore, SD, M. Grazzini,
S. Kallweit, N. Rana, A. Vicini, arXiv:2412.16095]



- We examine the **high invariant-mass region**, following cuts from [CMS collaboration, arXiv:2103.02708]:
 $p_{T,\mu^\pm} > 53 \text{ GeV}$, $|y_\mu| < 2.4$,
 $m_{\mu\mu} > 150 \text{ GeV}$
- The **mixed corrections** provide a relevant contribution, **increasing at higher invariant mass** and larger than the statistical error expected by the end of **HL-LHC**;
- The **factorised approximation** gives a **good description** of the high-invariant mass region.

ABISS - EVALUATING c_k



SOLVING D.E. B SERIES EXPANSION

Method implemented in the Mathematica package DiffExp for real kinematic variables [F.Moriello, arXiv:1907.13234], [M.Hidding, arXiv:2006.05510] (see also AMFLOW [X. Liu and Y.-Q. Ma, arXiv: 2201.11669])

A Simple Example

$$\begin{cases} f'(x) + \frac{1}{x^2 - 4x + 5} f(x) = \frac{1}{x+2} \\ f(0) = 1 \end{cases}$$

$$f_{hom}(x) = x^r \sum_{k=0}^{\infty} c_k x^k$$

$$f'_{hom}(x) = \sum_{k=0}^{\infty} (k+r) c_k x^{(k+r-1)}$$

$$\begin{cases} r c_0 = 0 \\ \frac{1}{5} c_0 + c_1(r+1) = 0 \\ \frac{4}{25} c_0 + \frac{1}{5} c_1 + c_2(2+r) = 0 \\ \dots \end{cases}$$

$$f_{hom}(x) = 5 - x - \frac{3}{10}x^2 + \frac{11}{150}x^3 + \dots$$

Expanded around $x' = 0$

$$f_{part}(x) = f_{hom}(x) \int_0^x dx' \frac{1}{(x'+2)} f_{hom}^{-1}(x')$$

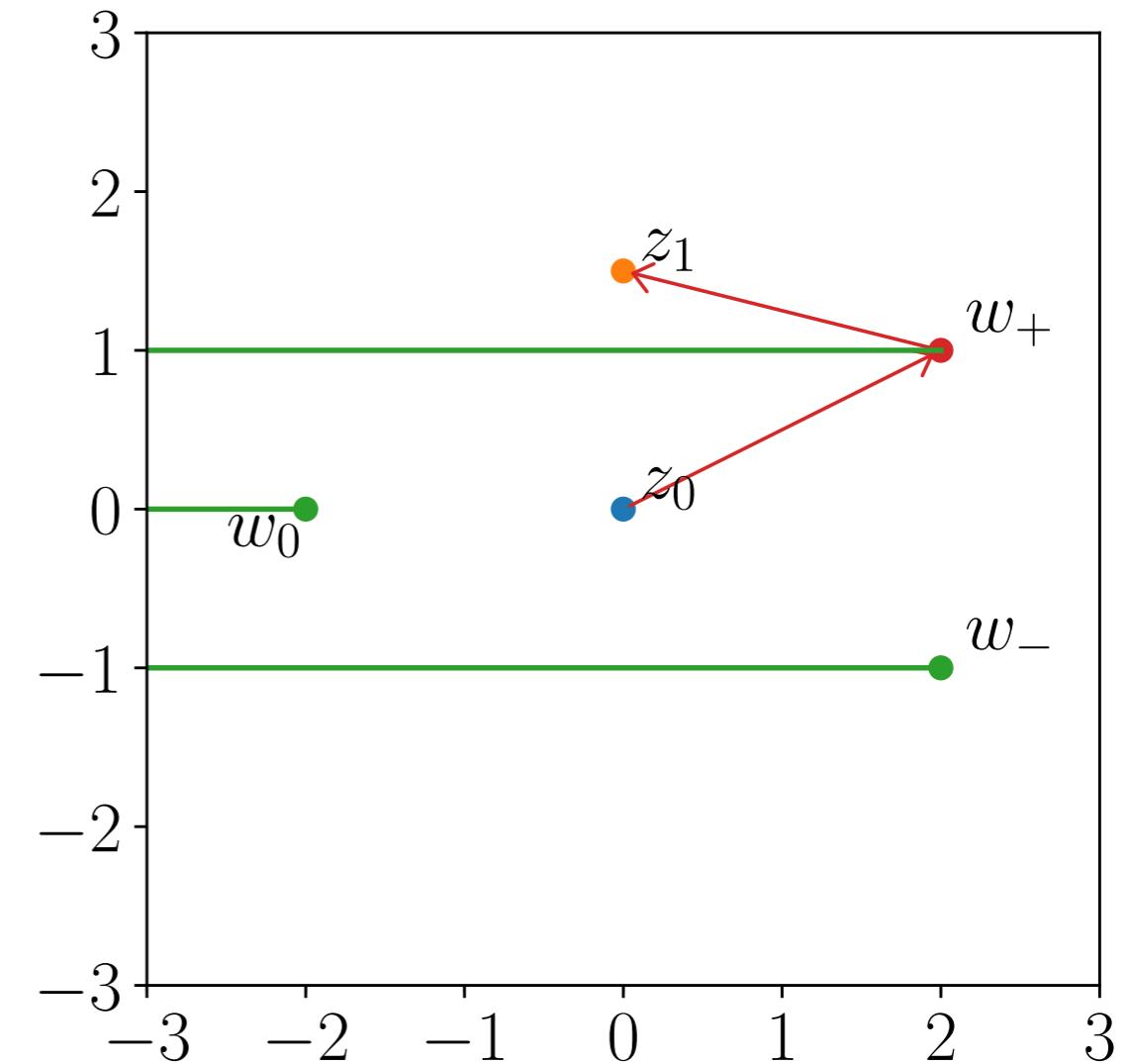
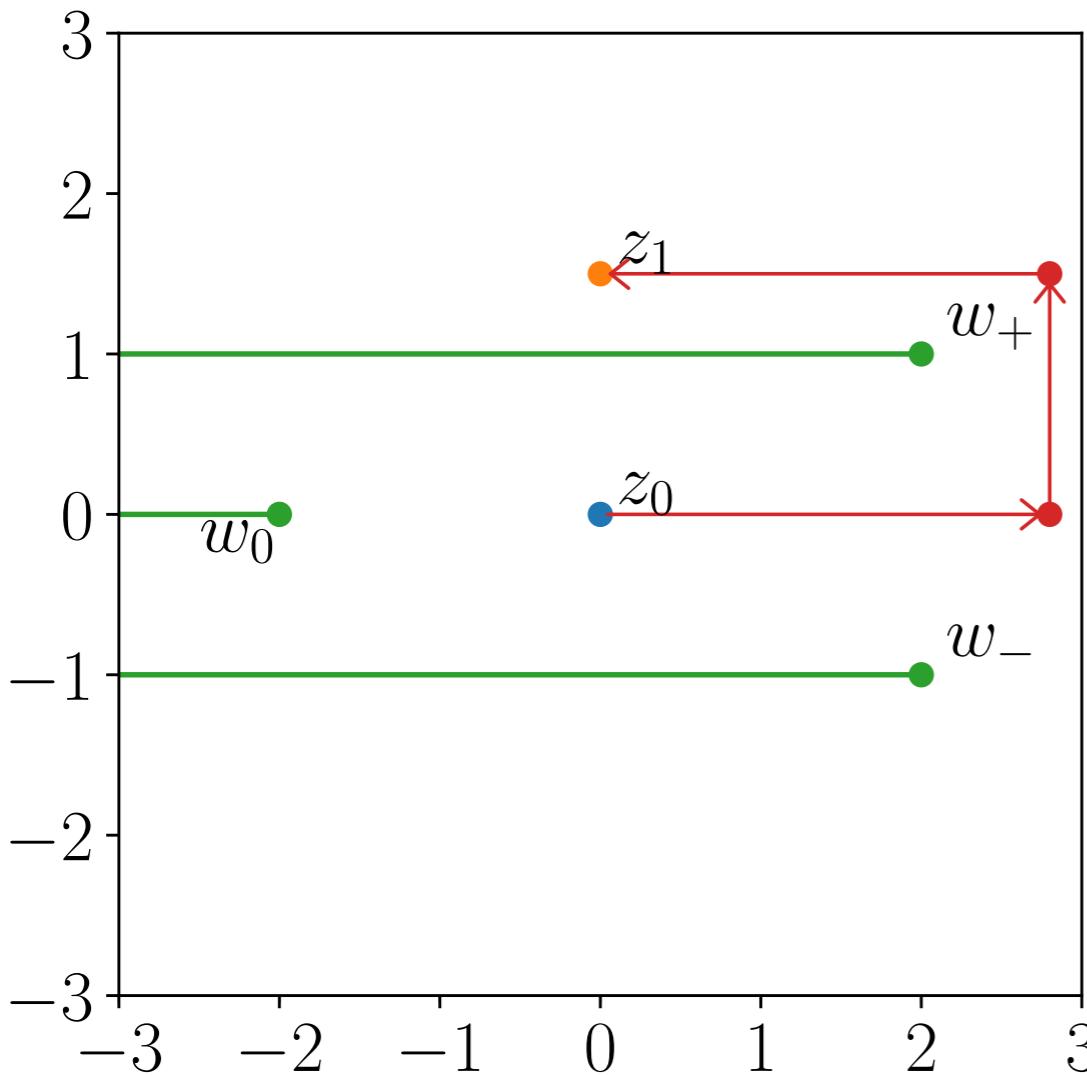
$$= \frac{1}{2}x - \frac{7}{40}x^2 + \frac{2}{75}x^3 + \dots$$

$$f(x) = f_{part}(x) + C f_{hom}(x)$$

$$f(0) = 1 \rightarrow C = \frac{1}{5}$$

TAYLOR VS LOGARITHMIC EXPANSION

- **Taylor expansion:** avoids the singularities;
- **Logarithmic expansion:** uses the singularities as **expansion points**.
- Logarithmic expansion has larger convergence radius but requires longer evaluation time. **We use Taylor expansion as default.**



TREATMENT OF γ_5

γ_5 is not well defined in a non integer number of dimensions!

	ANTICOMMUTATION $\{\gamma_\mu, \gamma_5\} = 0$	CYCLOCITY OF THE TRACE
't Hooft and Veltmann <i>Nucl. Phys. B</i> 44 (1972) 189–213	✗	✓
Kreimer et al. <i>Phys. Lett. B</i> 237 (1990) 59–62	✓	✗

For neutral-current Drell Yan proven that at 2loops the two prescriptions yield:

► **different** scattering amplitudes; ► **same** finite corrections after subtraction.

[M. Heller, A. von Manteuffel, R. M. Schabinger and H. Spiesberger, arXiv:hep-ph/2012.05918]

Our procedure:

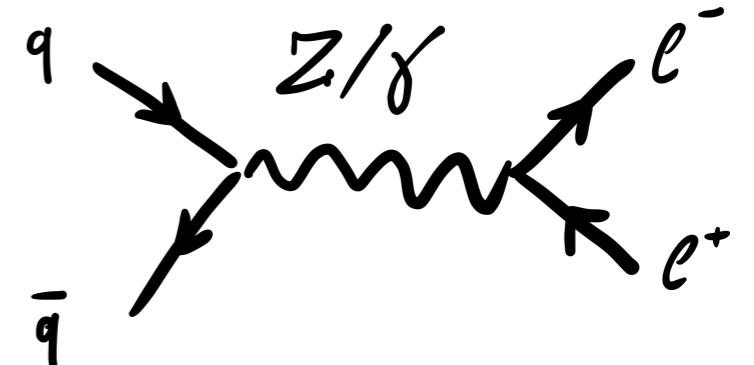
1. Use anticommutation relation, bring all γ_5 at the end of the Dirac trace;
2. Use $\gamma_5^2 = 1$, end up with zero or one γ_5 in each Dirac trace;
3. Replace the (single) leftover γ_5 with the relation: $\gamma_5 = \frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma$.

NCDY - BASIS OF MASTER INTEGRALS

[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2201.01754]

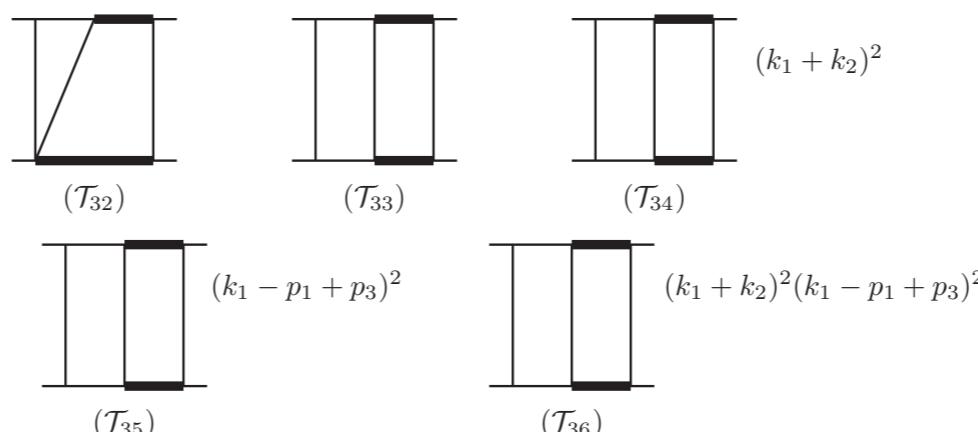
Basis of Master integrals composed by:

- MIs relevant for the **QCD-QED corrections**,
with massive final state; [R.Bonciani, A.Ferroglia, T.Gehrmann, D.Maitre,
C.Studerus, arXiv:0806.2301, 0906.3671]
- MIs with 1 or 2 internal mass relevant for the
EW form factor; [U.Aglietti, R.Bonciani, hep-ph/0304028, hep-ph/
0401193]
- 31 MIs with 1 mass and **36 MIs with 2 masses** including boxes, relevant for
the **QCD-EW corrections to the full Drell-Yan**.



5 box integrals are in **Chen-Goncharov** representation;

[R. Bonciani, S. Di Vita, P. Mastrolia,
U. Schubert, arXiv:1604.08581]



difficult numerical evaluation requires alternative strategy.

closed form in terms of GPLs found but not public

[M.Heller, A.von Manteuffel, R.M. Schabinger,
arXiv:1907.00491]

NCDY - NUMERICAL GRIDS

[T. Armadillo, R. Bonciani, SD, N.Rana,
A.Vicini, arXiv:2201.01754]

- 31 out of 36 masters known in terms of GPLs: **validation** of SEASYDE.
- 5 out of 36 masters are a genuine SEASYDE **prediction**;
- solution can be computed with **arbitrary number of significant digits**.

