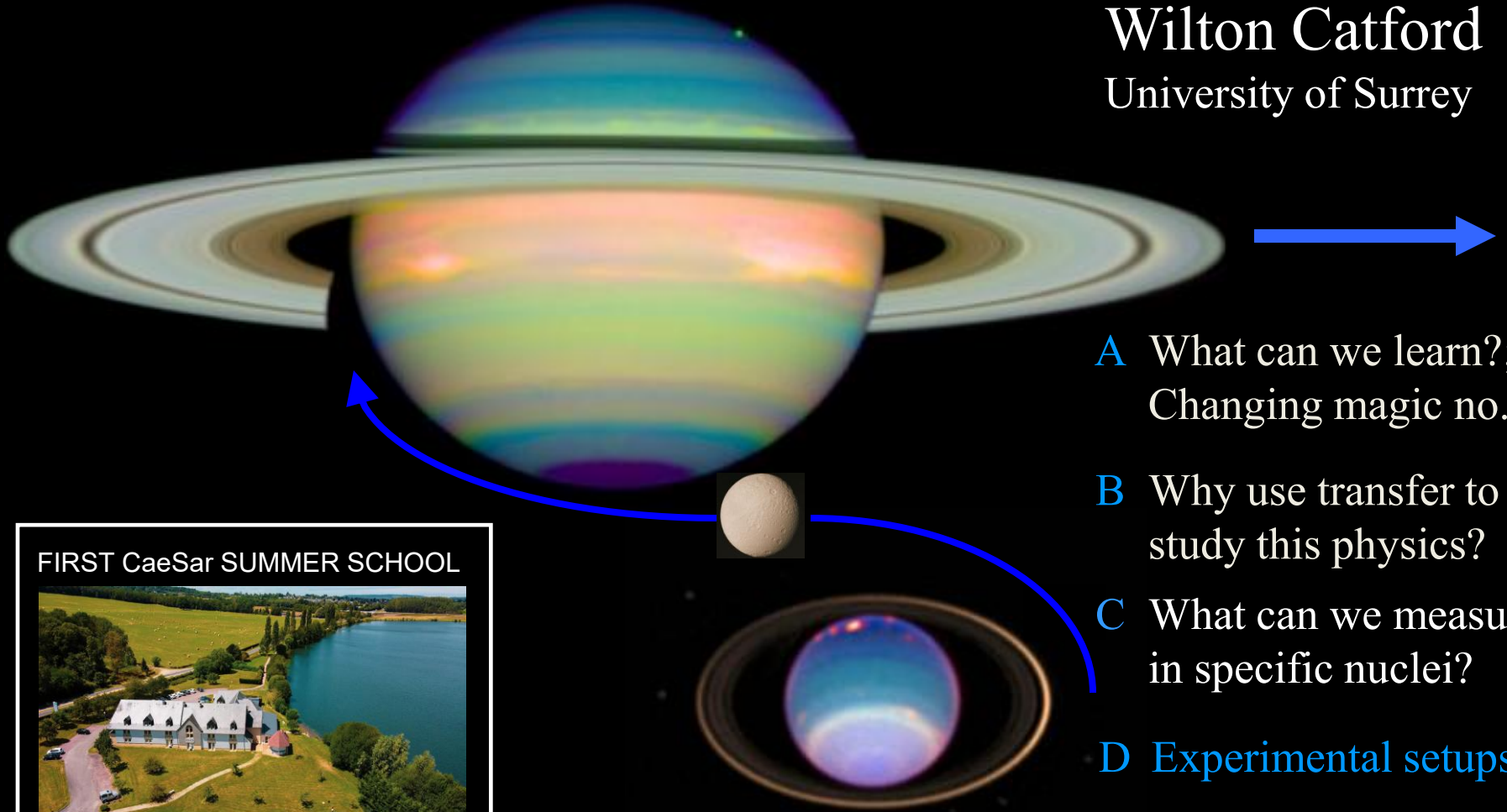


Exploring Unusual Nuclei with Nucleon Transfer Reactions

Wilton Catford
University of Surrey



- A What can we learn?;
Changing magic no.'s
- B Why use transfer to
study this physics?
- C What can we measure
in specific nuclei?
- D Experimental setups
- E Examples of data &
Interpretation

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SATURN HST/ IR 1998, TETHYS VOYAGER2 1981, URANUS HST/ IR 1986

A PLAN for how to study nuclear STRUCTURE :

- Use **transfer reactions** to identify strong single-particle states, measuring their spins and strengths
- Use the **energies** of these states to compare with theory
- **Refine** the structure (e.g. shell model, ab initio) theory
- Improve the extrapolation to very exotic nuclei
- Hence learn the structure of very exotic nuclei

N.B. The **shell model** is arguably the best theoretical approach for us to confront with our results, but it's **not the only one**. The experiments are needed, no matter which theory we use.

N.B. Transfer (as opposed to knockout) allows us to study orbitals that are empty, so **we don't need quite such exotic beams**.

- Motivation: nuclear structure reasons for transfer
- What quantities we actually measure
- What reactions/energies can we choose to use?
- Inverse Kinematics
- Implications for Experimental approaches
- Why do people make the choices that they do?
- Example experiments and results

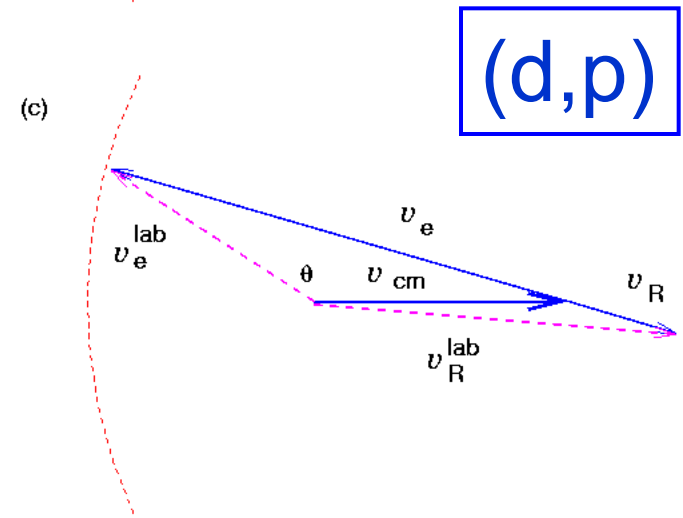
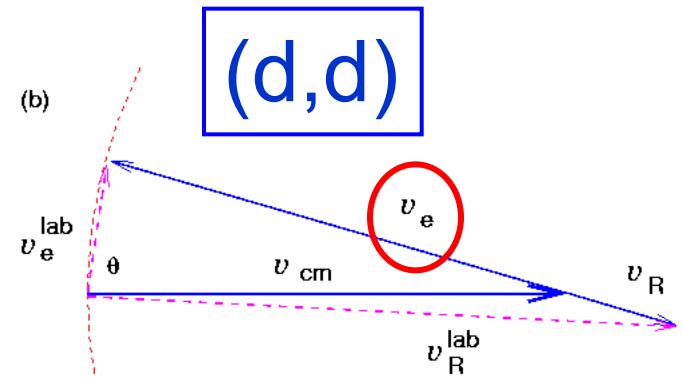
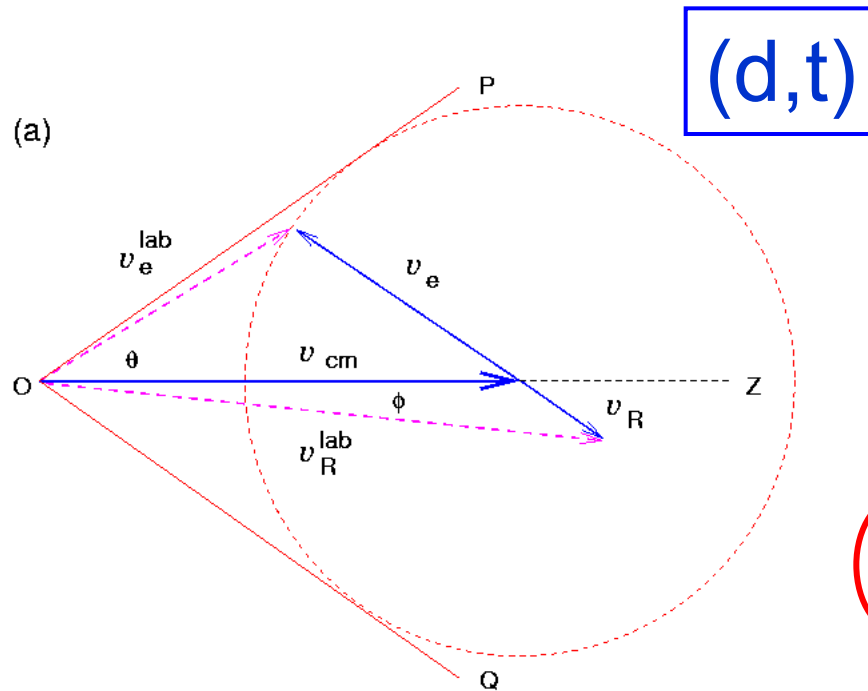
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LE LAC DE PONT-L'ÉVÊQUE

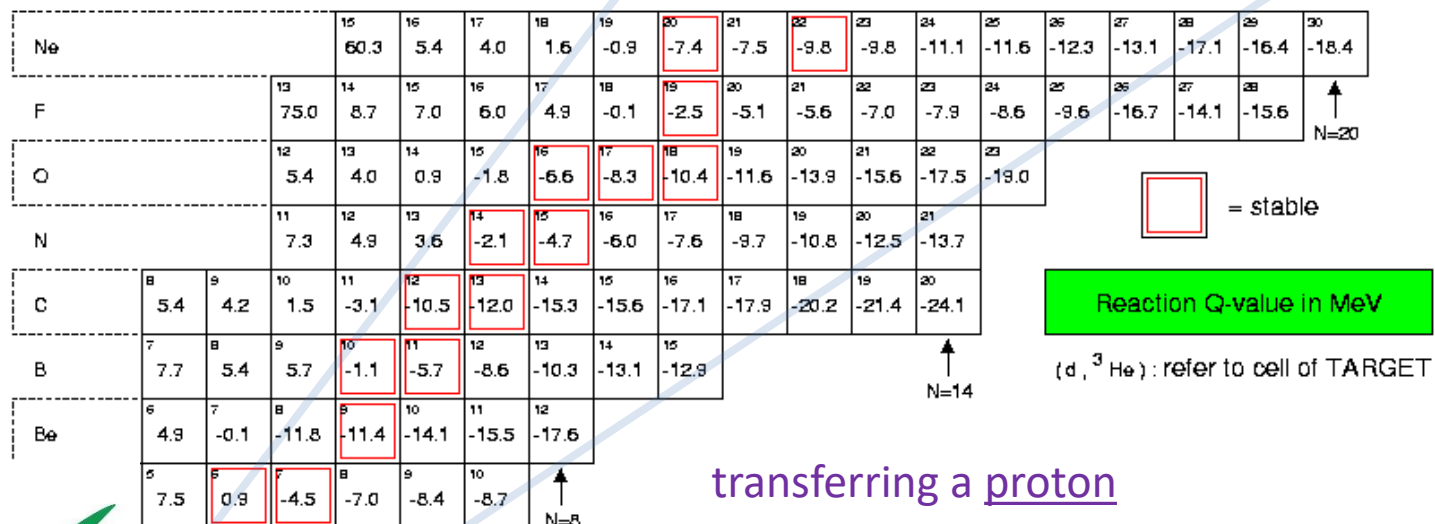
Using Radioactive Beams in Inverse Kinematics (heavy incident on light)



$$\frac{v_e}{v_{cm}} = \left(q f \frac{M_R}{M_P} \right)^{1/2} \cong \sqrt{q f}$$

$$\theta_{\max} = \sin^{-1} \sqrt{f}$$

$f = 1/2$ for (p,d), $2/3$ for (d,t)
 $q \cong 1 + Q_{\text{tot}} / (E/A)_{\text{beam}}$

$$\mathbf{Z}$$


close to zero

very negative

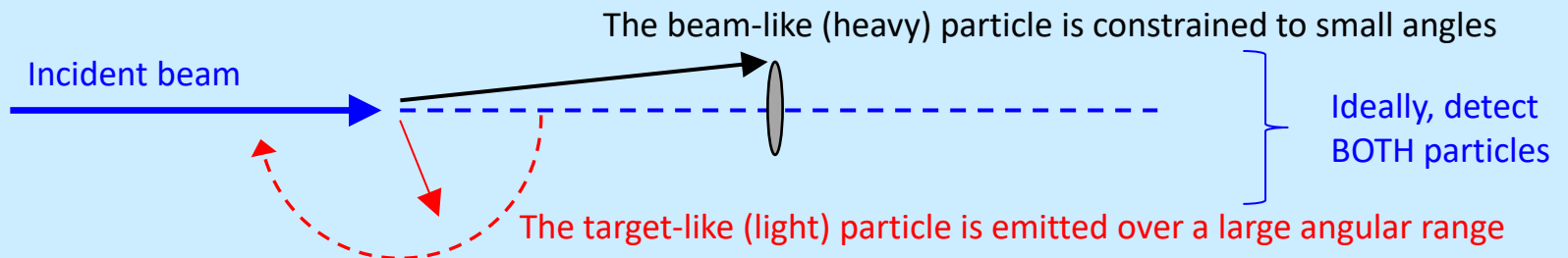
Possible experimental approaches to measuring transfer reactions...

1) Rely on detecting the **beam-like ejectile** in a spectrometer

- Kinematically favourable unless beam mass (and focussing) *too* great ✓
- Spread in beam energy (several MeV) translates to E_x measurement ✗
- Hence, need energy tagging, or a dispersion matching spectrometer ✗
- Spectrometer is subject to broadening from gamma-decay in flight ✗

2) Rely on detecting the **target-like ejectile** in a Si detector

- Kinematically less favourable for angular coverage ✗
- Spread in beam energy generally gives little effect on E_x measurement ✓
- Resolution limited by difference [$dE/dx(\text{beam}) - dE/dx(\text{ejectile})$] ✗
- Target thickness limited to 0.5-1.0 mg/cm² to maintain resolution ✗



Calculations of E_x resolution arising from particle detection...

152

J.S. Winfield et al. / Nucl. Instr. and Meth. in Phys. Res. A 396 (1997) 147–164

beamlike
particle
detected

Table 2

Major contributions in keV to the resolution of the excitation energy spectra of single neutron stripping and pickup reactions in inverse kinematics, where the heavy ion is detected in a spectrometer. The detection angle corresponds to 10°_{cm} . The last column is an approximate estimate as a sum in quadrature of the net effect of five non-Gaussian contributions. Other symbols are explained in the text

Reaction	E_i/A (MeV)	θ_{lab}	Origin of contribution					Σ_{quad}
			$\Delta\theta$	Δp	E_{stragg}	$\Theta_{1/2}$	dE/dx	
$p(^{12}\text{Be}, ^{11}\text{Be})d$	30	1.07°	172	147	101	74	23	259
$p(^{12}\text{Be}, ^{11}\text{Be})d$	15	1.06°	84	71	99	74	37	169
$p(^{77}\text{Kr}, ^{76}\text{Kr})d$	30	0.16°	1404	811	808	723	56	1952
$p(^{77}\text{Kr}, ^{76}\text{Kr})d$	10	0.10°	334	143	502	570	268	883
$d(^{76}\text{Kr}, ^{77}\text{Kr})p$	10	0.21°	1140	614	2177	1859	1321	3408

dominant contribution from ANGLE measurement

light
particle
detected

Table 3

Major contributions in keV to the resolution of the excitation energy spectra of single neutron pickup and stripping reactions in inverse kinematics, where the light particle is detected in a silicon detector. Symbols as described in text and Table 2

Reaction	E_i/A (MeV)	θ_{lab}	Origin of contribution					Σ_{quad}
			$\Delta\theta$	ΔE_f	ΔE_i	$\Theta_{1/2}$	dE/dx	
$p(^{12}\text{Be}, d)^{11}\text{Be}$	30	19.0°	136	74	114	96	649	685
$p(^{12}\text{Be}, d)^{11}\text{Be}$	15	17.8°	66	72	55	89	984	995
$p(^{77}\text{Kr}, d)^{76}\text{Kr}$	30	15.0°	124	55	64	63	186	249
$p(^{77}\text{Kr}, d)^{76}\text{Kr}$	10	6.0°	26	24	23	19	775	777
$d(^{76}\text{Kr}, p)^{77}\text{Kr}$	10	155.3°	52	93	37	60	1309	1316

energy loss leaving target

Lighter projectiles

Heavier projectiles

Some minor advantages to detect
beam-like particle
(angular resolution difficult at higher energies)

Better to detect light particle
(target thickness limits E_x resolution)

Possible Experimental Approaches to Nucleon Transfer

1) Rely on detecting the beam-like ejectile in a spectrometer

- Kinematically favourable unless beam mass (and focussing) *too* great
- Spread in beam energy (several MeV) translates to E_x measurement
- Hence, need energy tagging, or a dispersion matching spectrometer
- Spectrometer is subject to broadening from gamma-decay in flight

2) Rely on detecting the target-like ejectile in a Si detector

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- Spread in beam energy generally gives little effect on E_x measurement
- Resolution limited by difference [$dE/dx(\text{beam}) - dE/dx(\text{ejectile})$]
- Target thickness limited to 0.5-1.0 mg/cm² to maintain resolution

3) Detect decay gamma-rays in addition to particles

- Need exceptionally high efficiency, of order > 25%
- Resolution limited by Doppler shift and/or broadening
- Target thickness increased up to factor 10 (detection cutoff, mult scatt'g)

An experiment where we measured the beam-like particle...

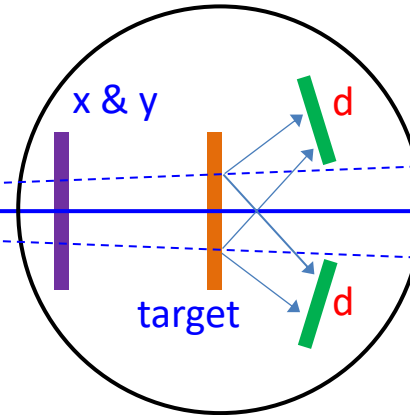
GANIL

$p(^{11}\text{Be}, d)^{10}\text{Be}$

Light projectile

^{11}Be beam

range of energies physically
dispersed across the target



dispersion-matched
spectrometer SPEG
(spectrometre à perte
d'énergie du GANIL)

angles:
 θ and φ

"in flight" beam
(range of energies)

spectrum of ^{10}Be states

beam stop

position
angles
 θ' and φ'
 ΔE
 E

An experiment where we measured the beam-like particle...

GANIL

$p(^{11}\text{Be}, d)^{10}\text{Be}$

Light projectile

^{11}Be beam

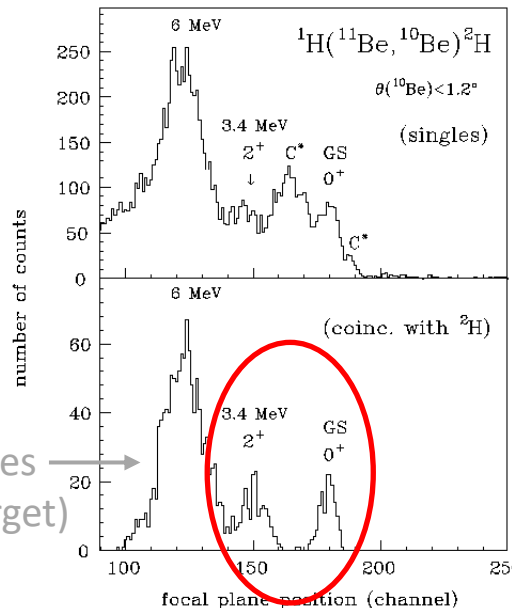
dispersion-matched spectrometer SPEG
(spectromètre à perte d'énergie du GANIL)

range of energies physically dispersed across the target

angles:
 θ and φ

"in flight" beam
(range of energies)

coincidence with d removes background (carbon in target)



spectrum of ^{10}Be states

beam stop

position
angles
 θ' and φ'
 ΔE
 E

An experiment where we measured the beam-like particle...

GANIL

$p(^{11}\text{Be}, d)^{10}\text{Be}$

Light projectile

^{11}Be beam

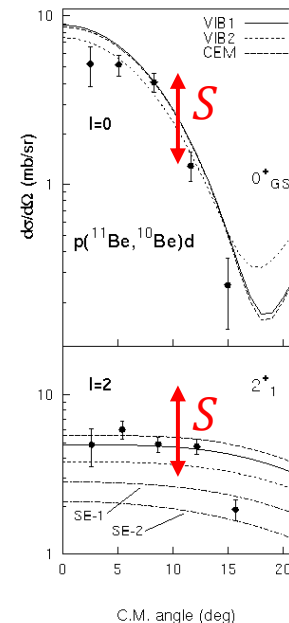
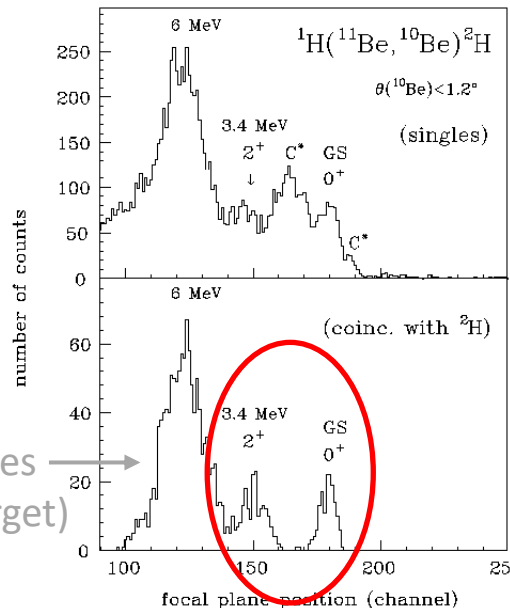
dispersion-matched spectrometer SPEG
(spectromètre à perte d'énergie du GANIL)

range of energies physically dispersed across the target

angles:
 θ and φ

"in flight" beam
(range of energies)

coincidence with d removes background (carbon in target)



$l=0$

$l=2$

beam stop

position

angles

θ' and φ'

ΔE

E

this is the angle that gets harder to measure for heavier projectiles

Possible experimental approaches to measuring transfer reactions...

1) Rely on detecting the **beam-like ejectile** in a spectrometer

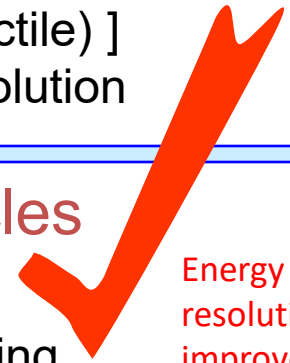
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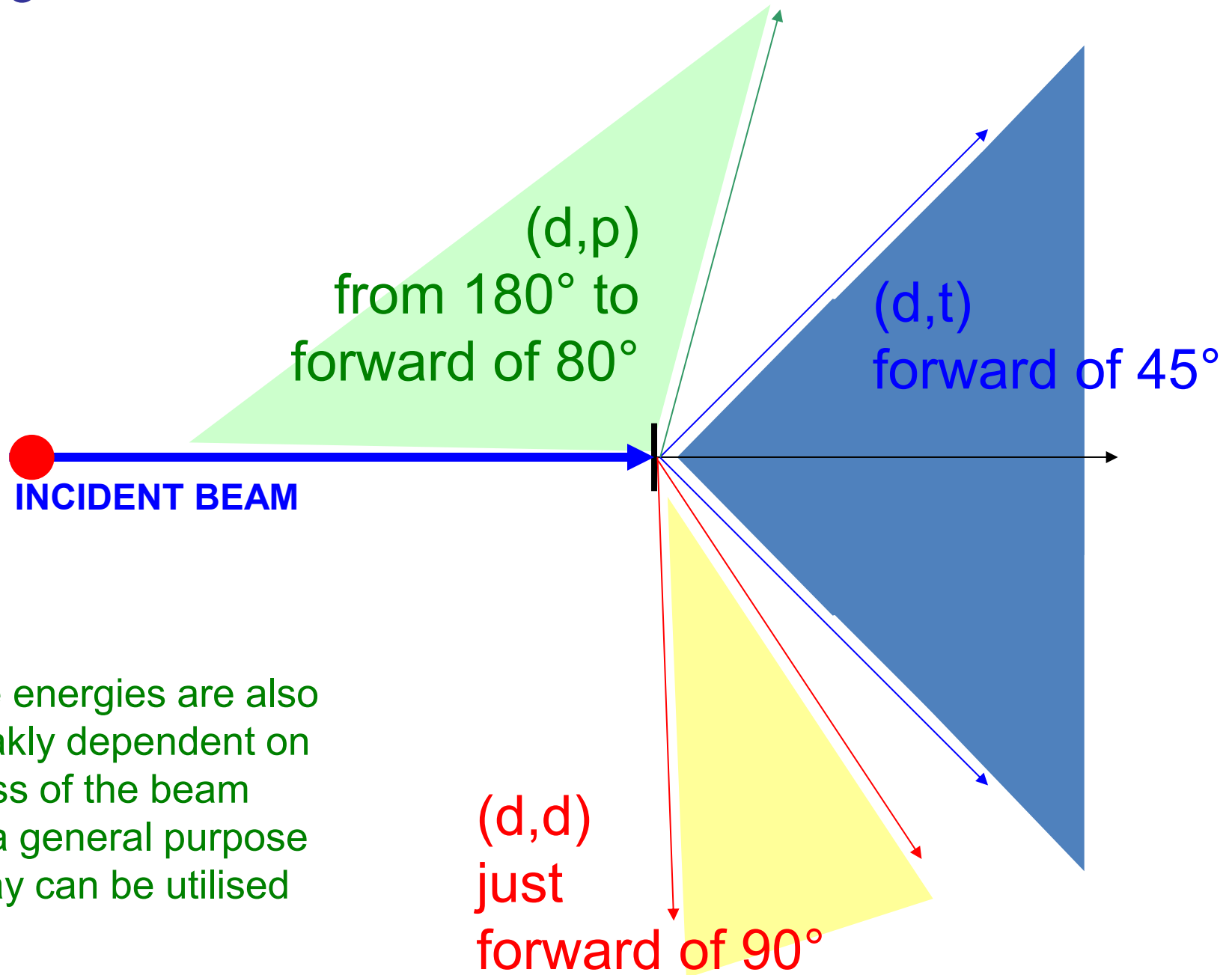
3) **Detect decay gamma-rays** in addition to **particles**

- Need exceptionally high efficiency, of order > 25%
- Resolution limited by only by Doppler shift and/or broadening
- Gamma-ray decay scheme gives valuable structure information



Energy
resolution
improved
by a factor
up to TEN

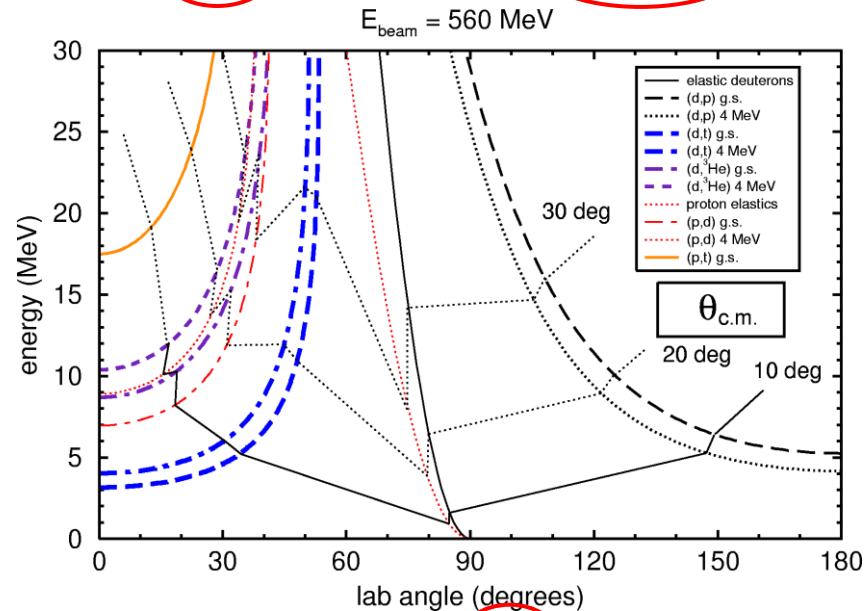
Using Radioactive Beams in Inverse Kinematics...



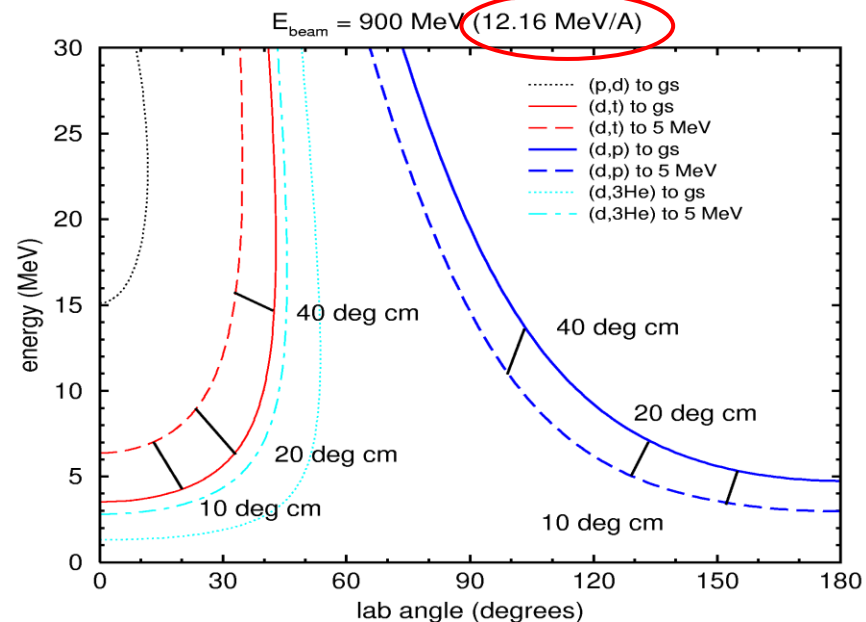
The energies are also weakly dependent on mass of the beam so a general purpose array can be utilised

The **general form** of the kinematic diagrams is determined by the **light particle masses**, and has **little dependence** on the beam mass or beam velocity

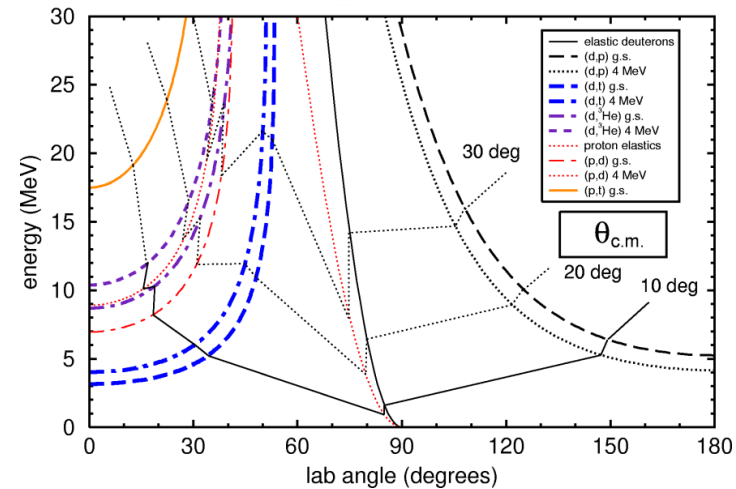
^{16}C incident on ^2H at 35 MeV/u



(p,d) and (d,t) and (d,p) on ^{74}Kr in inverse kinematics



The **general form** of the kinematic diagrams is determined by the **light particle masses**, and has **little dependence** on the beam mass or beam velocity



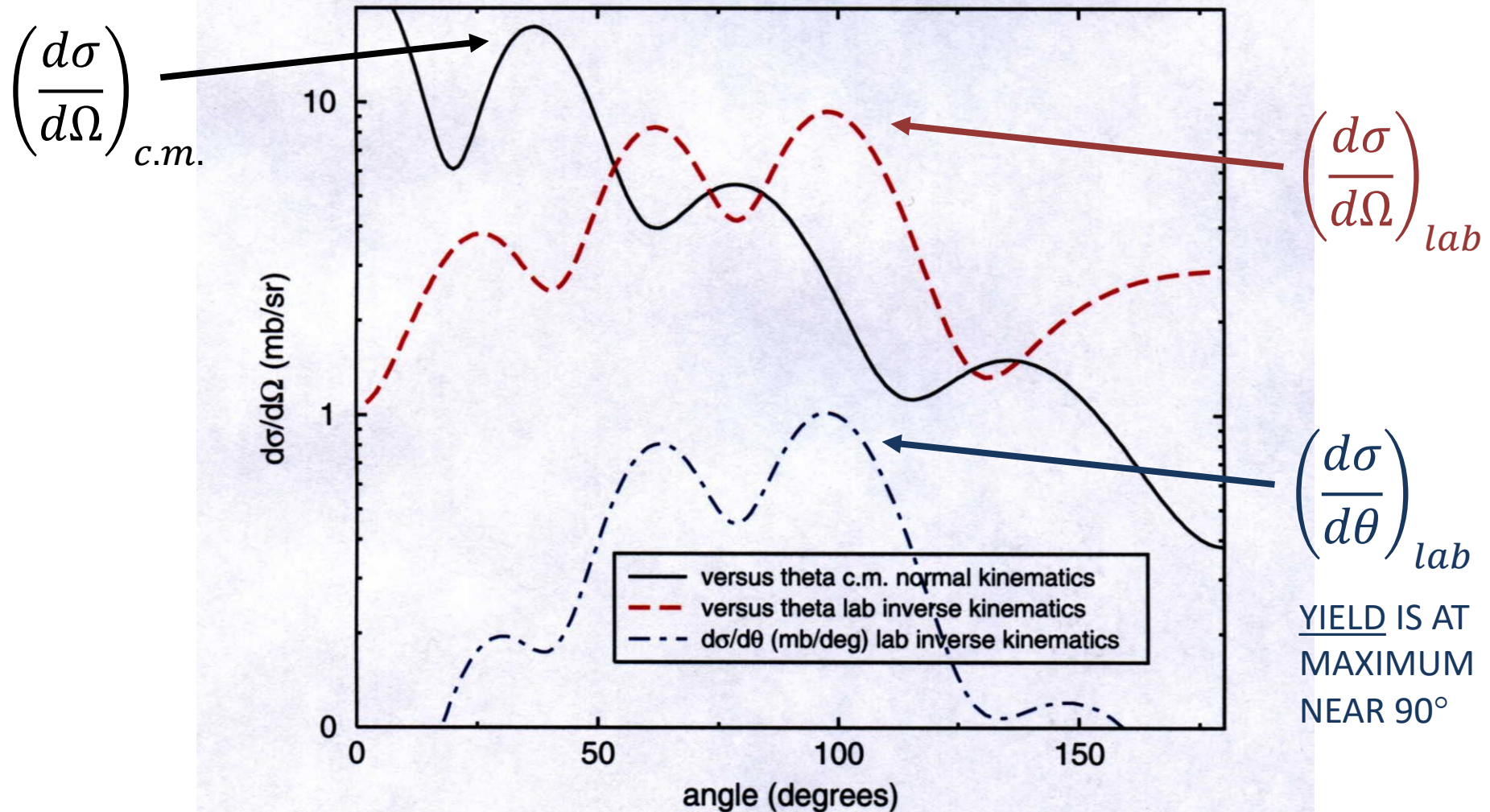
THIS IS IMPORTANT BECAUSE

it means that we can design a fixed experimental setup that will work for all experiments with all beams

By the way: remember the difference between $\frac{d\sigma}{d\Omega}$ and $\frac{d\sigma}{d\theta}$!!

DWBA ZR: $^{94}\text{Sr}(d,p)^{95}\text{Sr}^*(1.0\text{MeV};s_{1/2})$ at 4.894 MeV/u

Adiabatic deuteron potential (B-G) and Perey proton potential





- We need to measure:
 - Energies
 - Angles
 - Particle identification
- We know the kinematics
- We know the angular distributions
- Nobody will give us more than 10 days

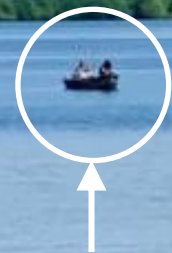
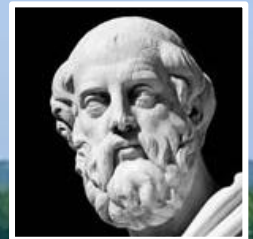


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let's explore some options...

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Which is the best way to implement the “light particle” option?

It turns out that the target thickness is a real limitation on the energy resolution...

Several hundred keV is implicit, when tens would be required,
So the targets should be as thin as possible...

But RIBs, as well as being heavy compared to the deuteron target, are:

(a) Radioactive

(b) Weak

Issues arising:

(a) Gamma detection useful for improving resolution

(b) Active target (TPC) to minimize loss of resolution

(c) Need MAXIMUM efficiency for detection

Experimental solutions can be classed roughly as:

(a) For beams $< 10^3$ pps ACTIVE TARGET

(b) $10^3 < \text{beam} < 10^6$ pps Si BOX in a γ -ARRAY

(c) For beams $> 10^6$ pps MANAGE RADIOACTIVITY

Solutions for beam intensities from 10^2 to 10^4 pps using TPCs*

* TIME PROJECTION CHAMBERS

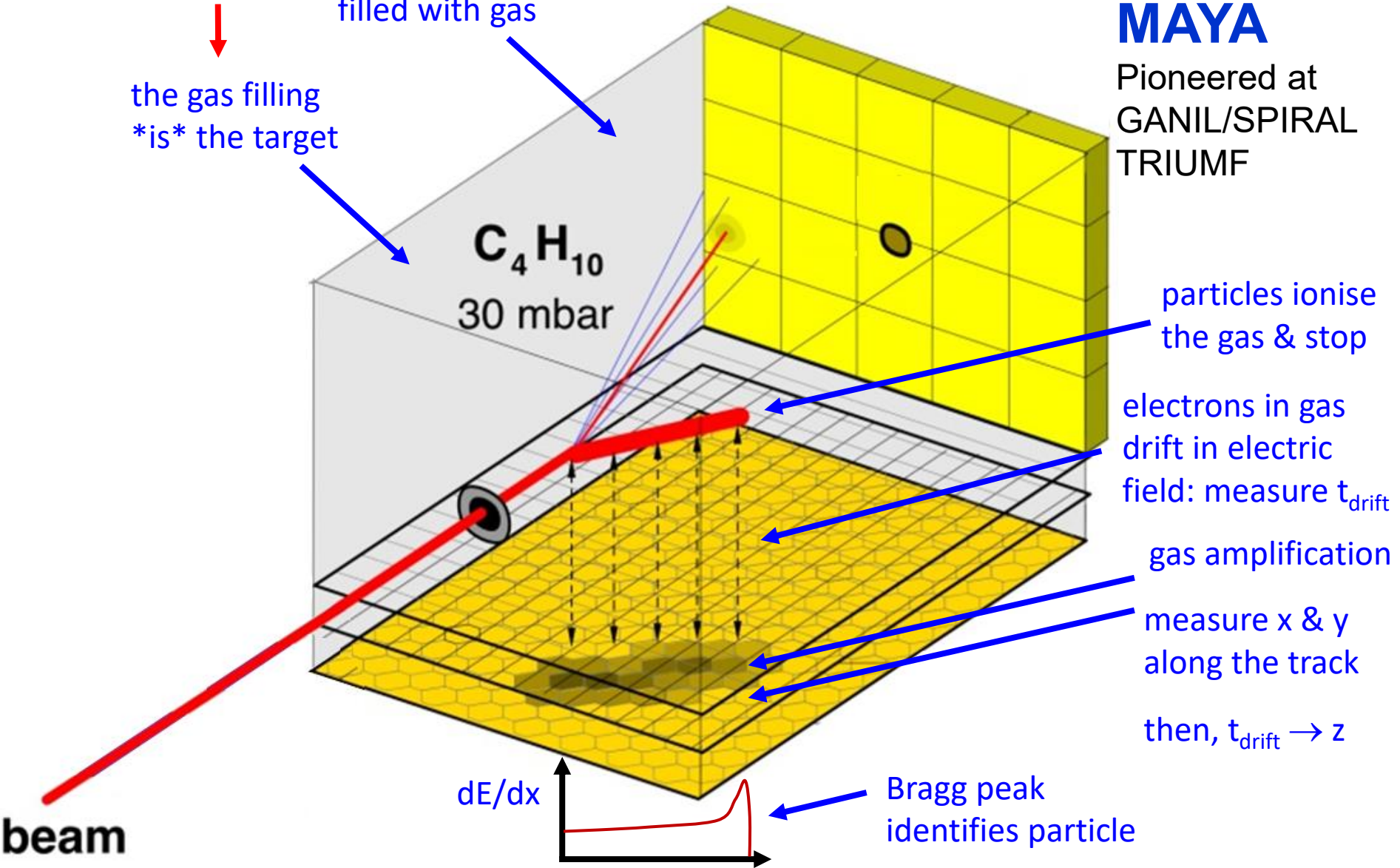
the TARGET is
the DETECTOR

the vessel is
filled with gas

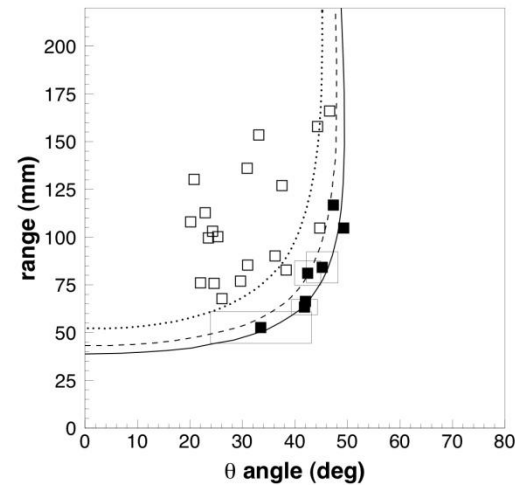
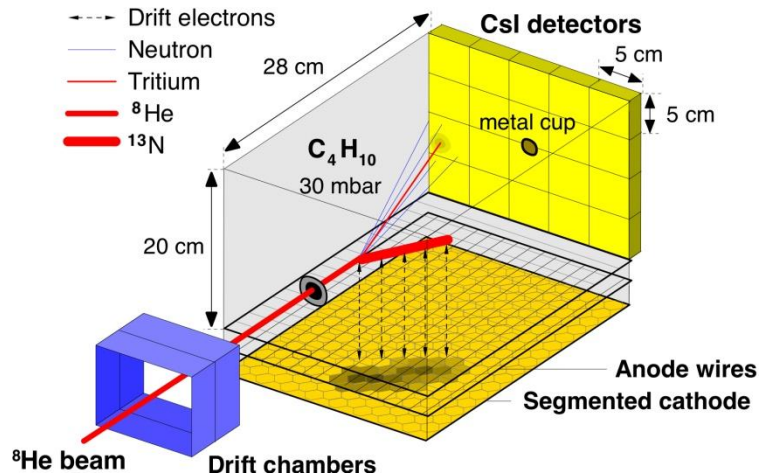
the gas filling
is the target

MAYA

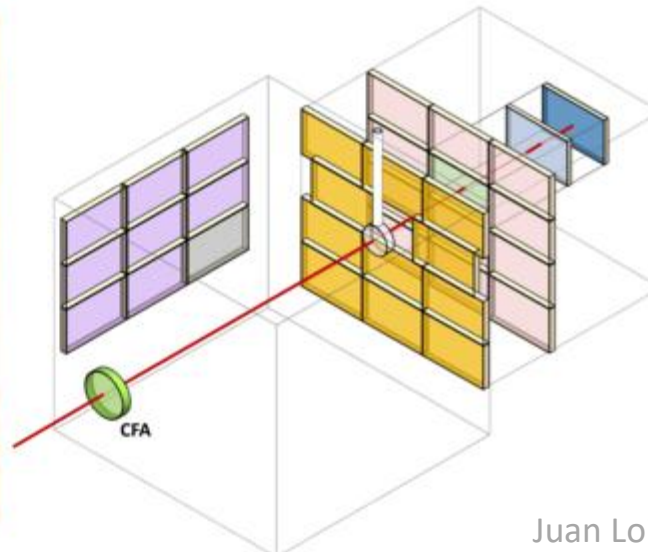
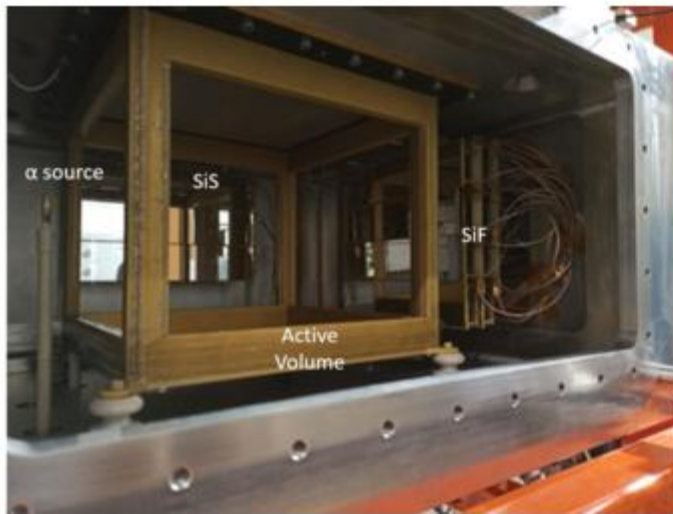
Pioneered at
GANIL/SPIRAL
TRIUMF



Solutions for beam intensities from 10^2 to 10^4 pps using TPCs*



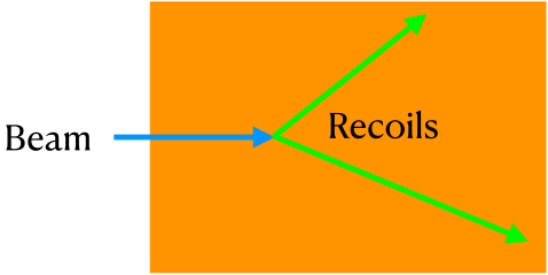
MAYA
Pioneered at
GANIL/SPIRAL
TRIUMF



ACTAR
Currently in
operation at
GANIL/SPIRAL
TRIUMF

Juan Lois Fuentes, PhD Thesis (2023)

* TIME PROJECTION CHAMBERS



Target = Detector

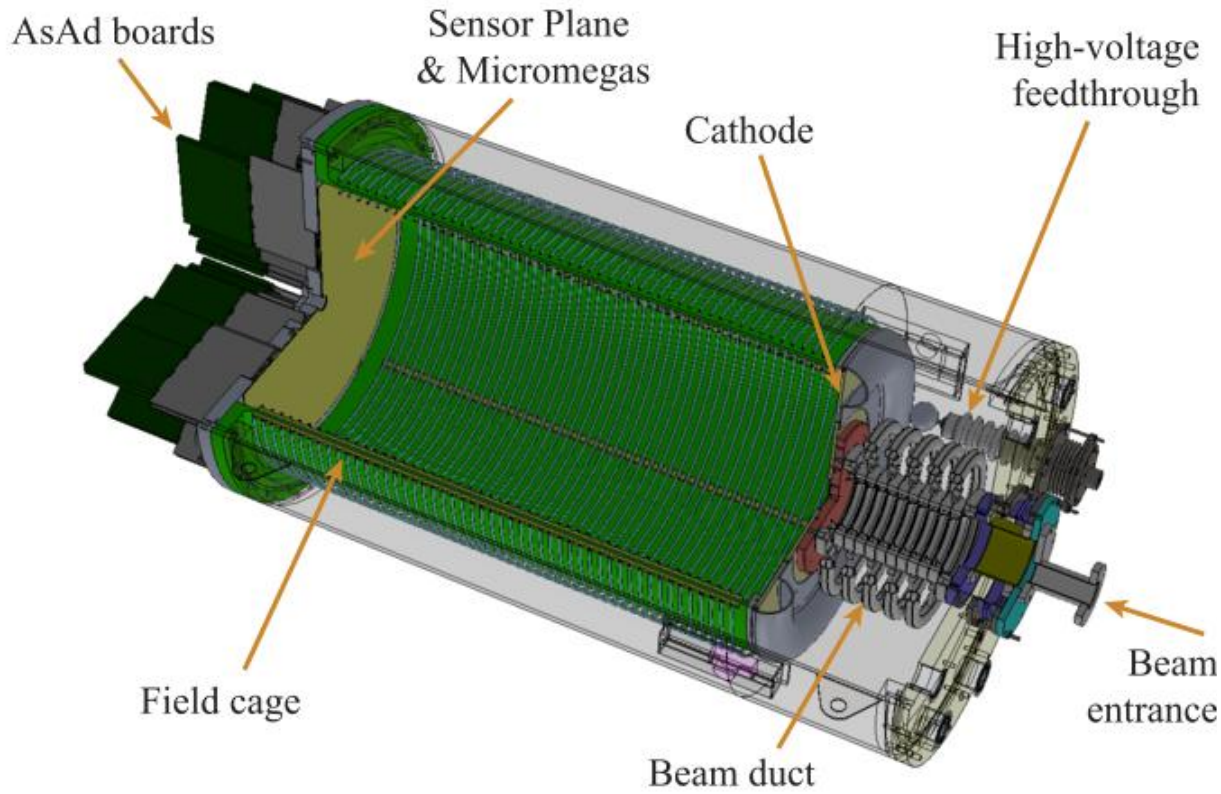
AT-TPC @ SOLARIS

D. Bazin, DREB 2024, June 24-28, 2024, Wiesbaden, Germany



FRIB + ReA6

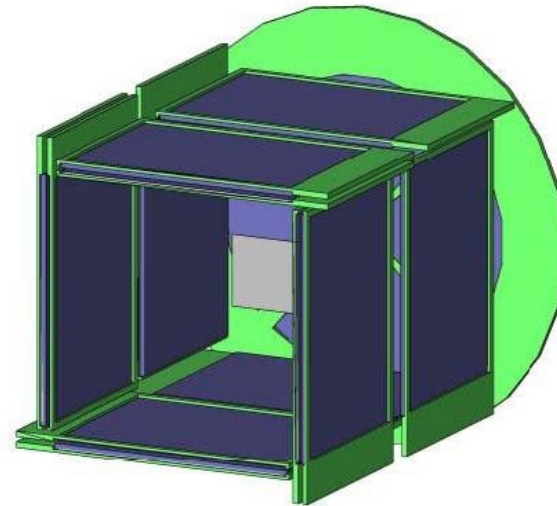
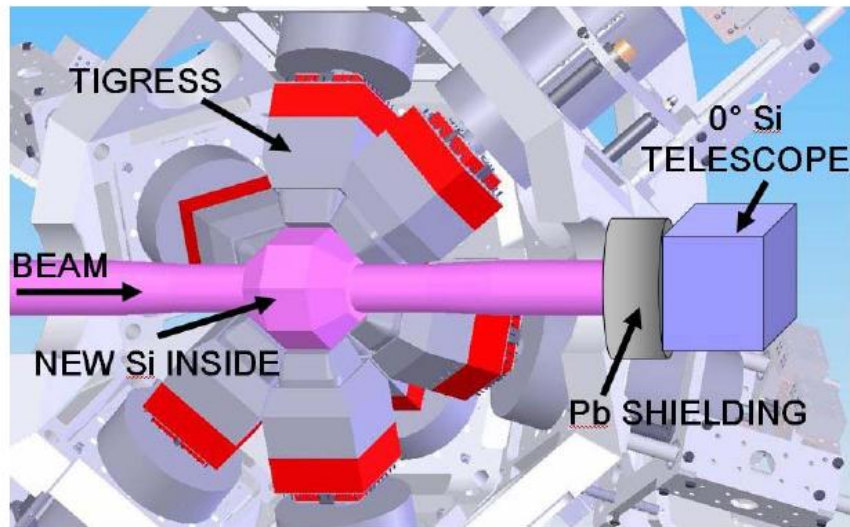
Active Target Time Projection Chamber



SOLARIS @ FRIB



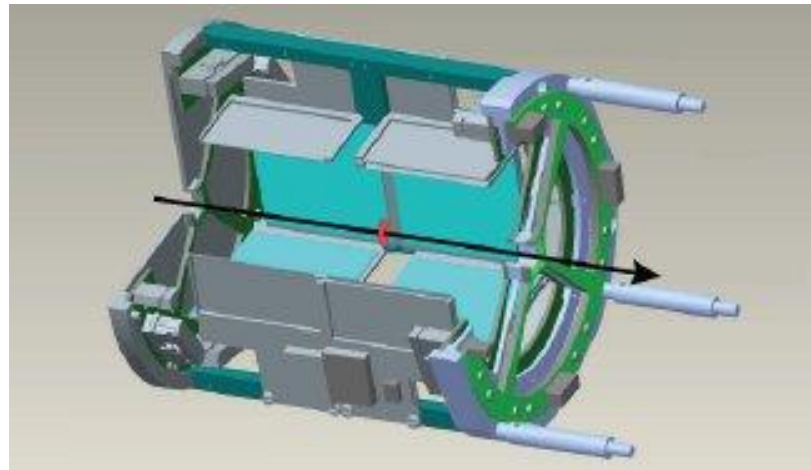
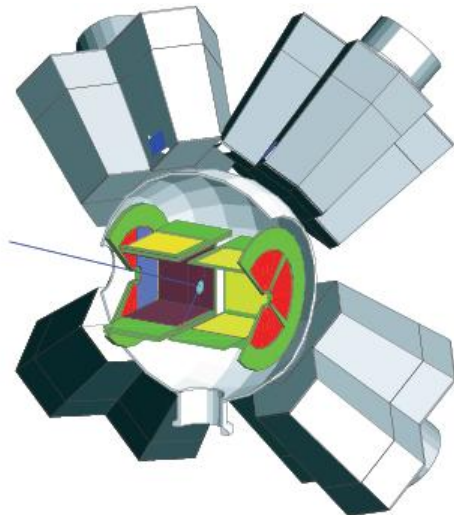
Solutions for more-intense beams of 10^4 to 10^6 pps, using gammas



SHARC

TIGRESS
TRIUMF

TIGRESS
COLLABORATION
York
Surrey



T-REX

MINIBALL
REX-ISOLDE

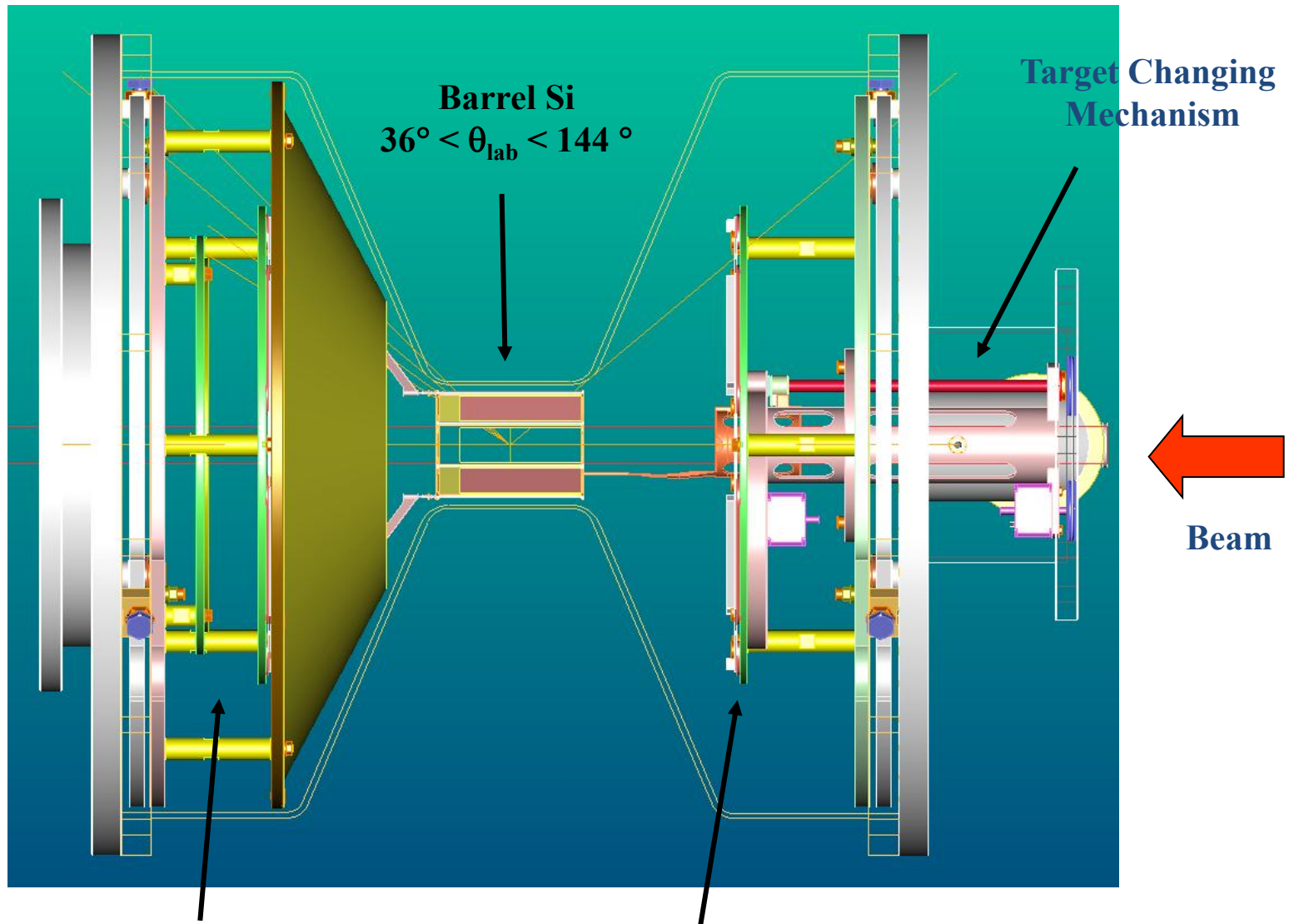
MINIBALL
COLLABORATION
Munich
Leuven

ORRUBA & GODDESS OAK RIDGE

Solutions for very intense beams of up to 10^9 pps, using gammas

TIARA^{☆☆☆}

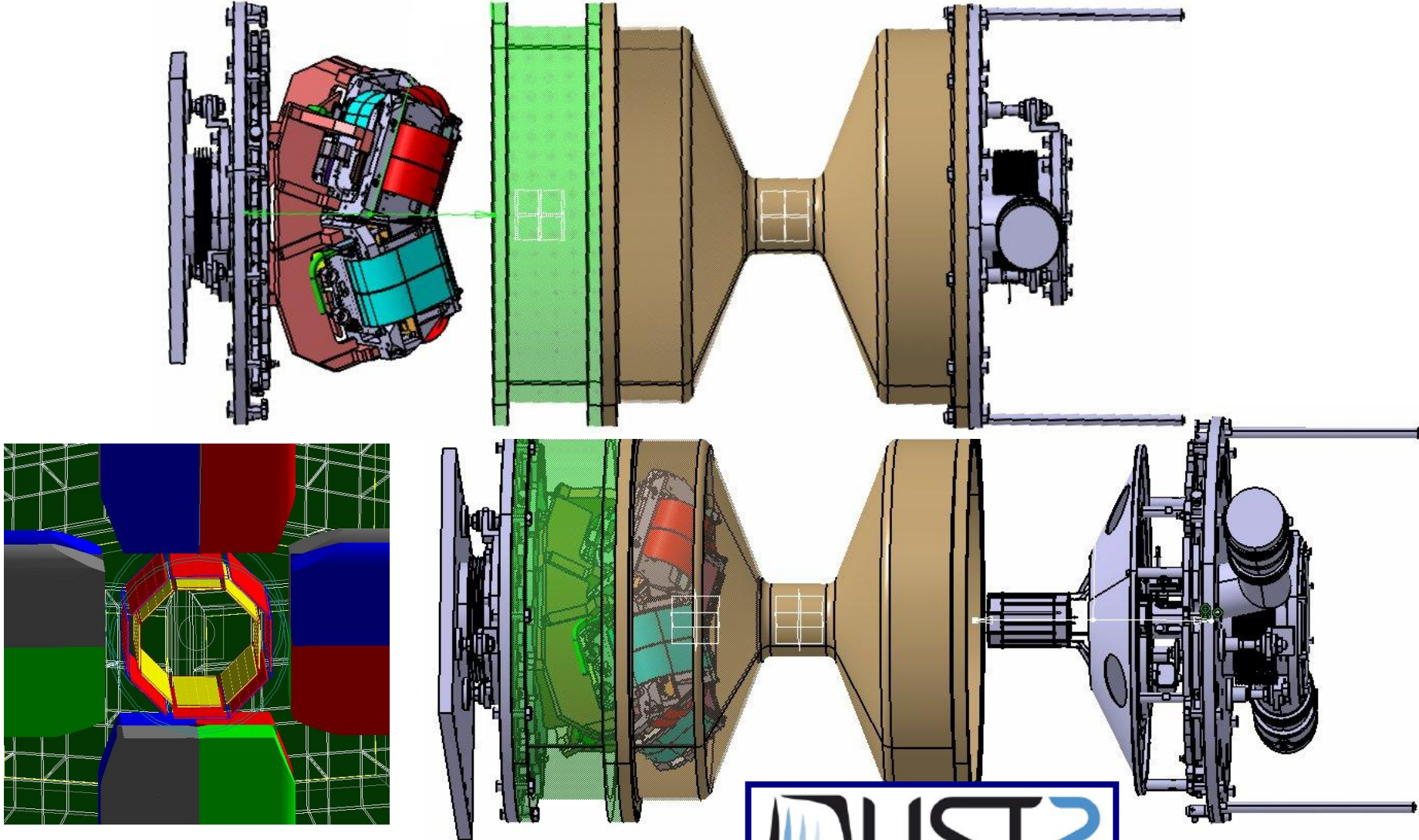
VAMOS



Forward Annular Si
 $5.6^\circ < \theta_{\text{lab}} < 36^\circ$

Backward Annular Si
 $144^\circ < \theta_{\text{lab}} < 168.5^\circ$

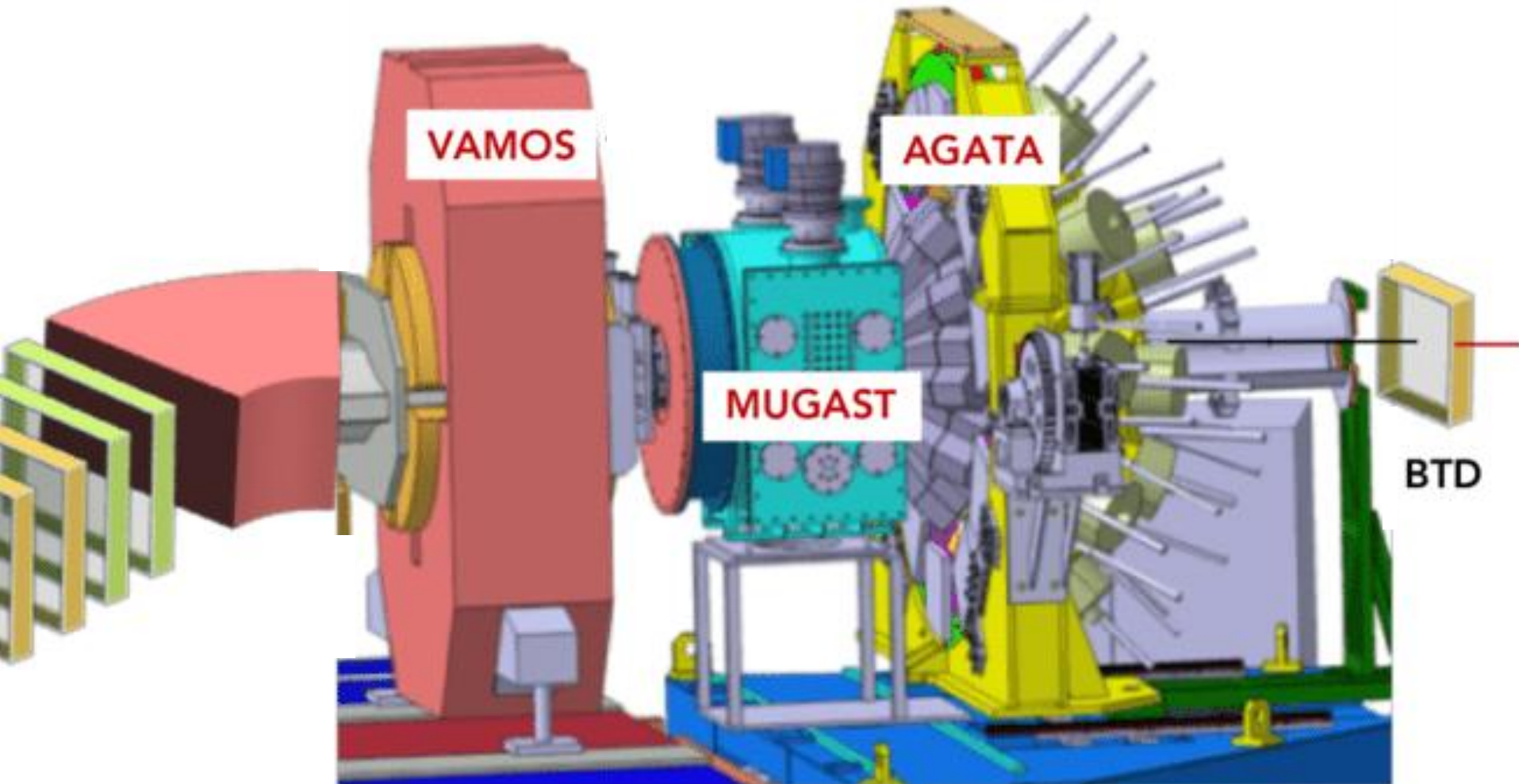
Solutions for very intense beams of up to 10^9 pps, using gammas



AUST2

TIARA 

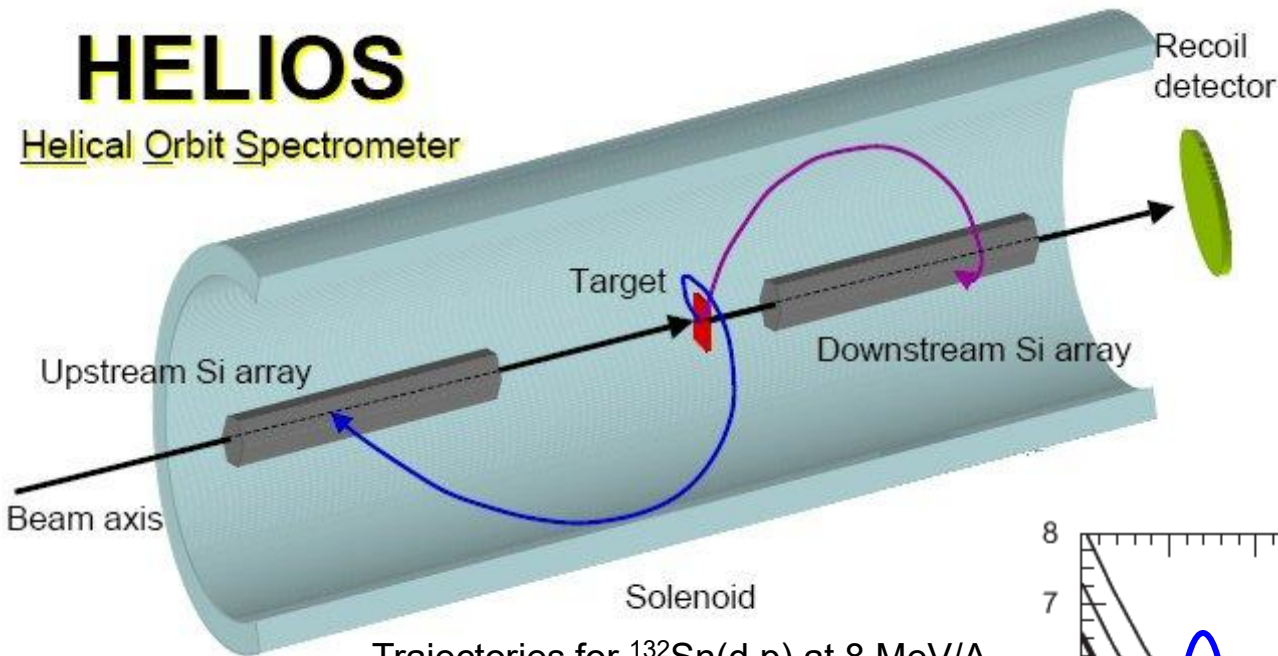
Beam: ^{47}K at 7.7 MeV/u (SPIRAL1)
Intensity: 5×10^5 pps
Purity: $\sim 100\%$



Solenoidal detector for 4π detection & to *de-compress* the kinematics

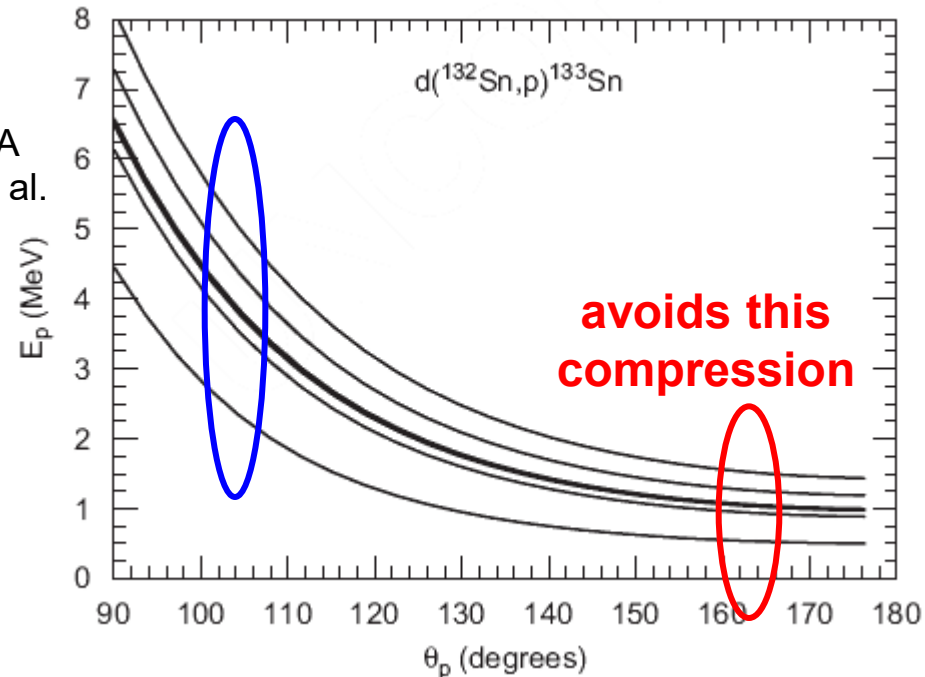
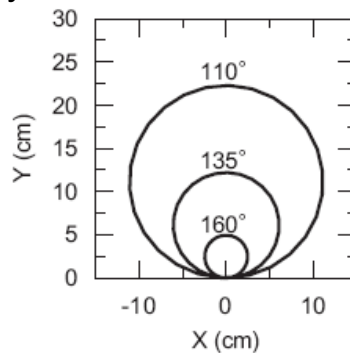
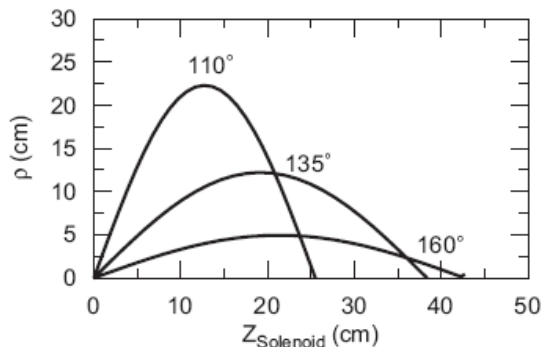
HELIOS

Helical Orbit Spectrometer



Actual solenoid – from MRI

Trajectories for $^{132}\text{Sn}(d,p)$ at 8 MeV/A
HELIOS: Kay, Wuosmaa, Schiffer et al.



Superconducting solenoid = second-hand MRI magnet (2.5 - 4.0 T)

"J'ai passé une IRM"



Superconducting solenoid = second-hand MRI magnet

"J'ai passé une IRM"

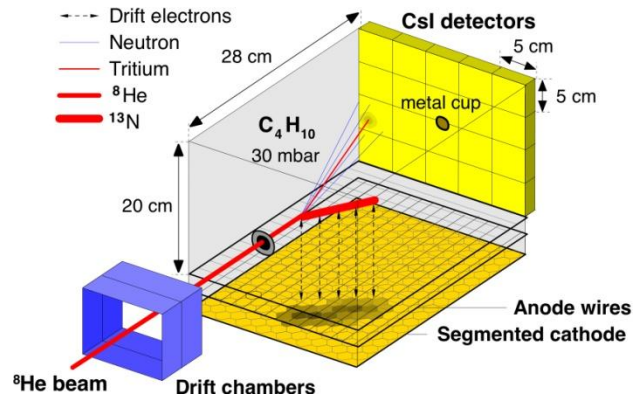


"C'est vrai"
(not as bad as I thought)

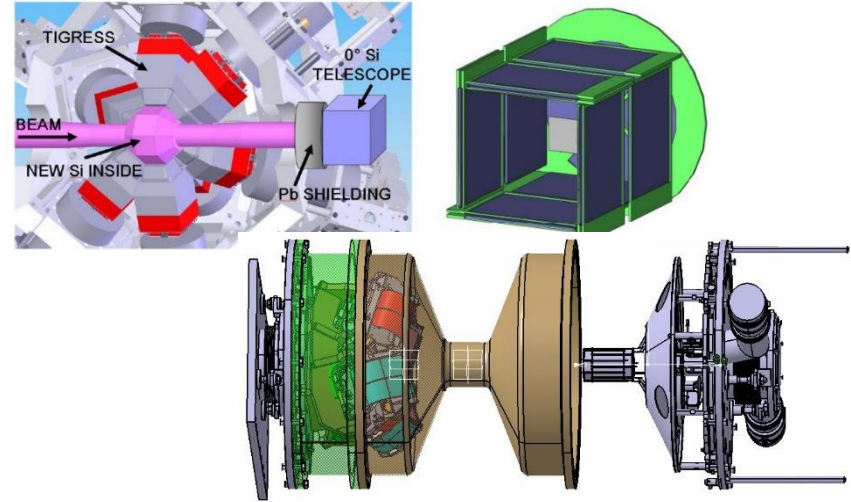


The experimental choice is strongly driven by the beam intensity:

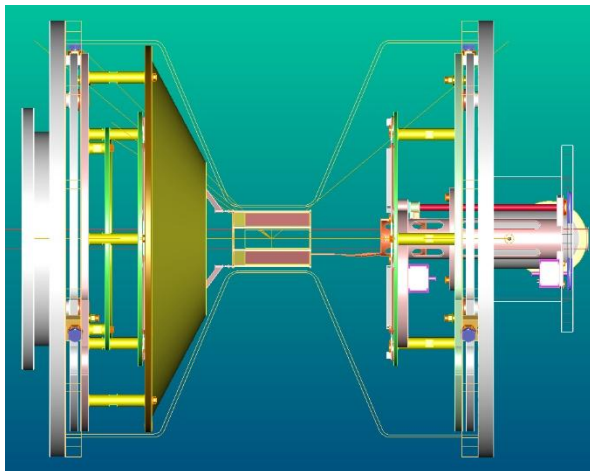
Below 10^4 pps MAYA, ACTAR...



Below 10^6 pps SHARC, MUGAST...

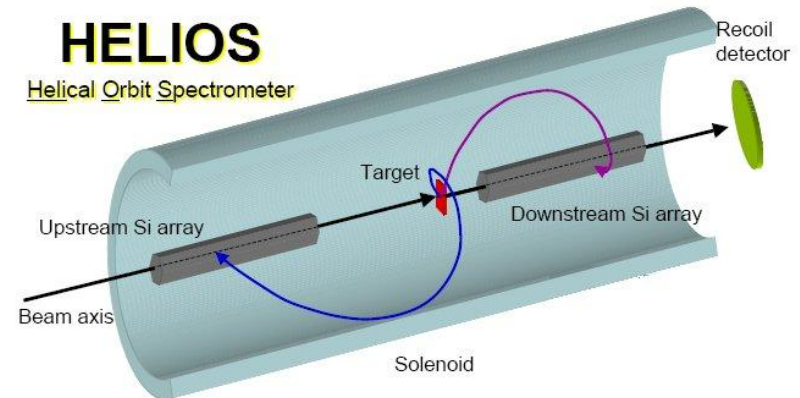


Up to 10^9 pps TIARA or MUGAST, or...



TIARA ★★★

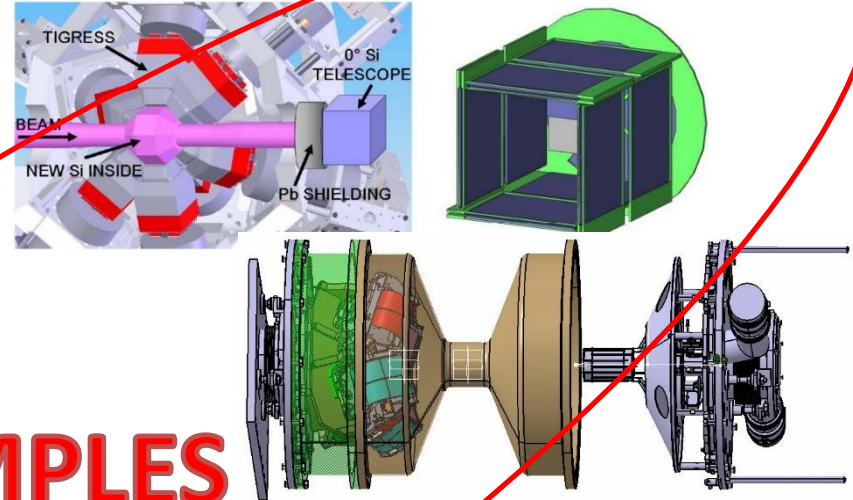
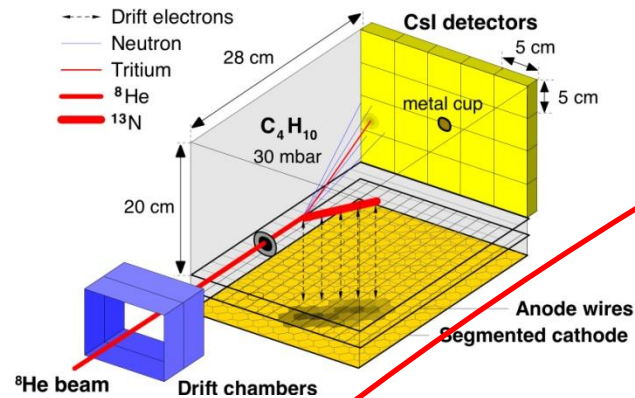
A solenoid device (no gammas)...



The experimental choice is strongly driven by the beam intensity:

Below 10^4 pps MAYA, ACTAR...

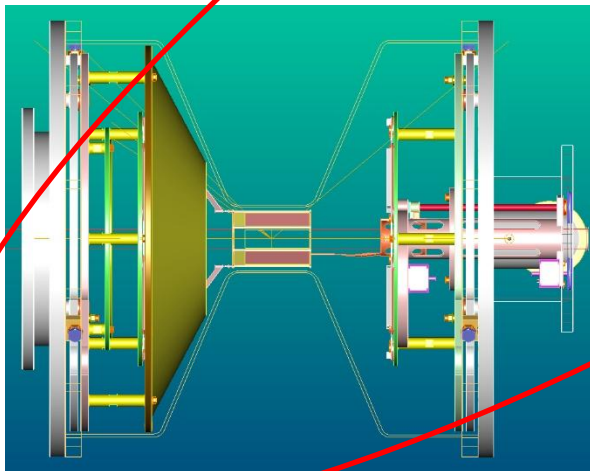
Below 10^6 pps SHARC, MUGAST...



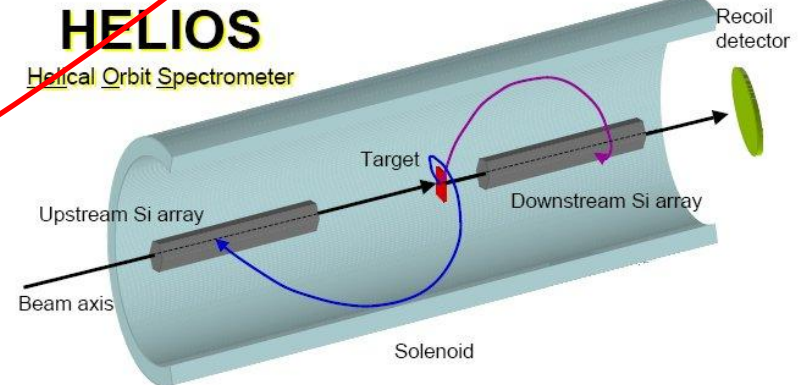
THREE EXAMPLES

Up to 10^9 pps TIARA or MUGAST, or...

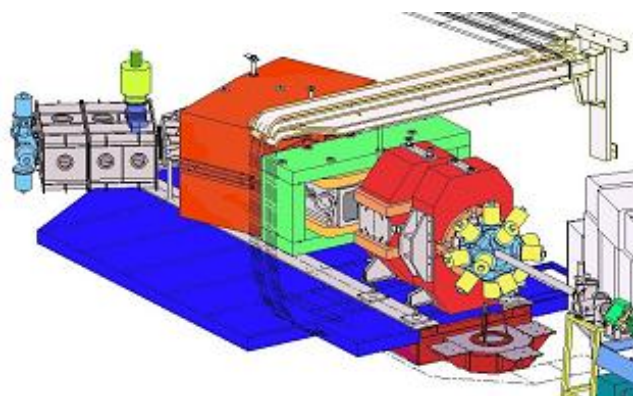
A solenoid device (no gammas)...



TIARA ★★★

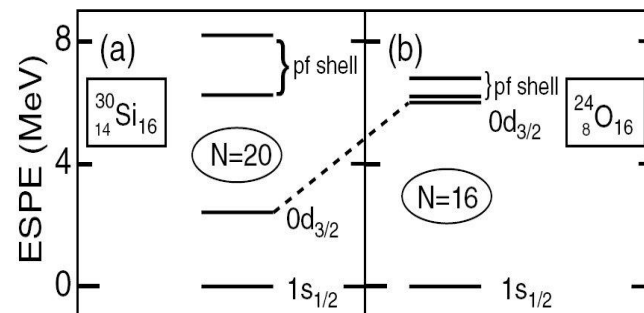
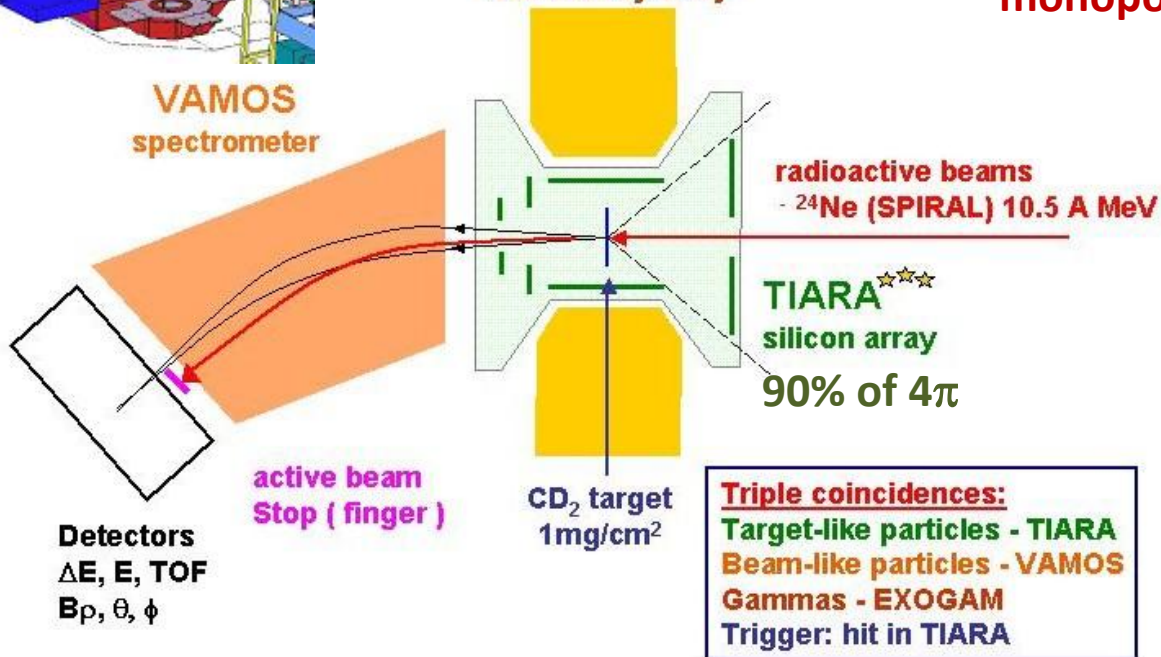


$^{24}\text{Ne}(d, p\gamma) ^{25}\text{Ne}$ – we discovered that N=16 replaces the broken N=20



VAMOS
spectrometer

EXOGAM
Gamma-ray array



monopole migration

10⁵ pps (pure)

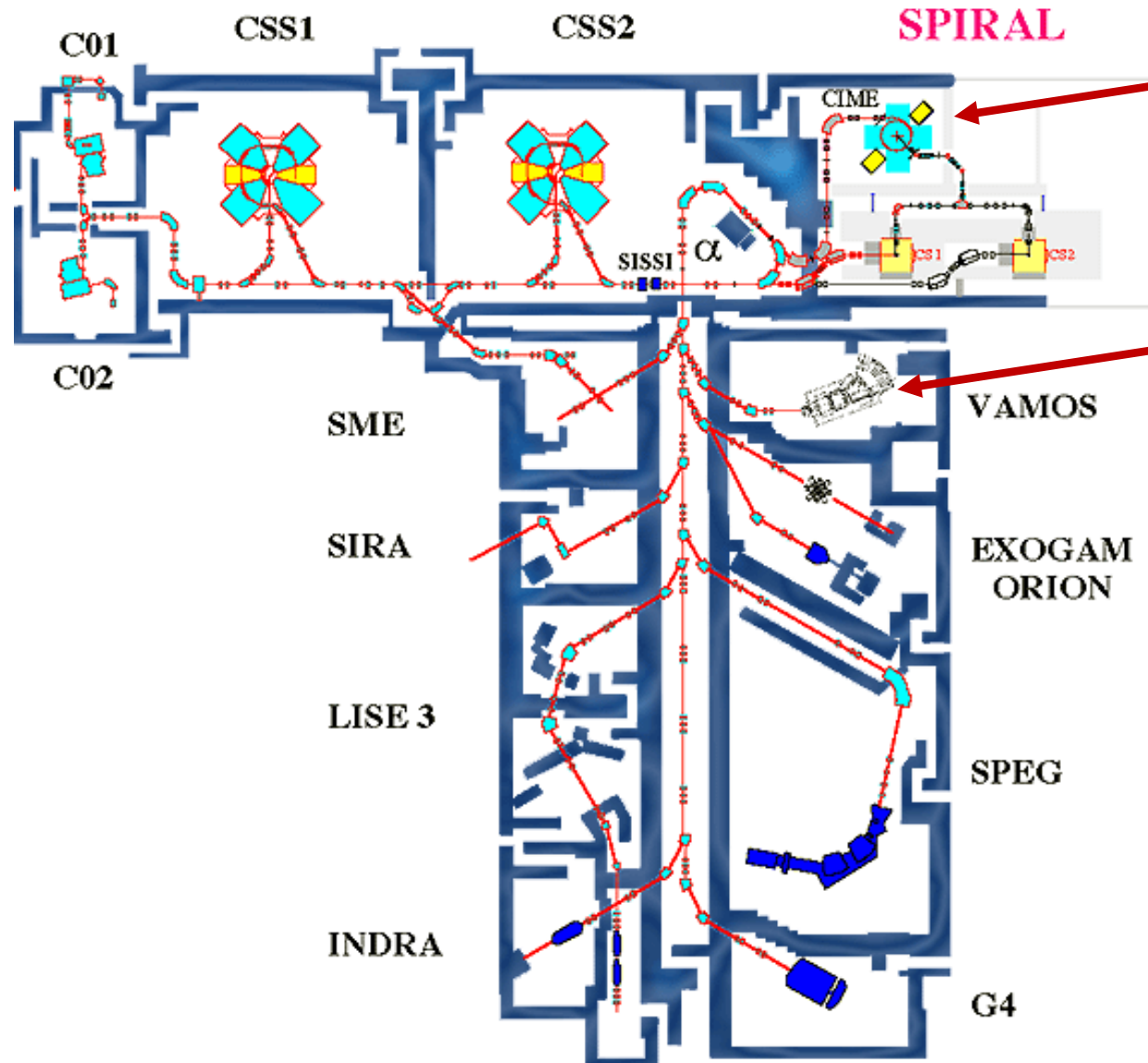
TIARA★★★

W.N. Catford *et al.*, Eur. Phys. J. **A25**, Suppl. 1, 245 (2005)

W.N. Catford *et al.*, Phys. Rev. Lett. **104**, 192501 (2010)

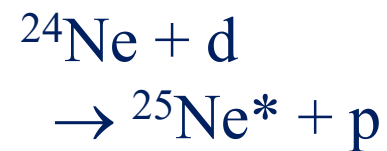
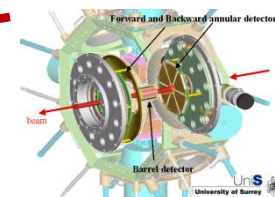
GANIL

GRAND ACCELERATEUR NATIONAL D'IONS LOURDS
LABORATOIRE COMMUN DSM/CEA-IN2P3/CNRS



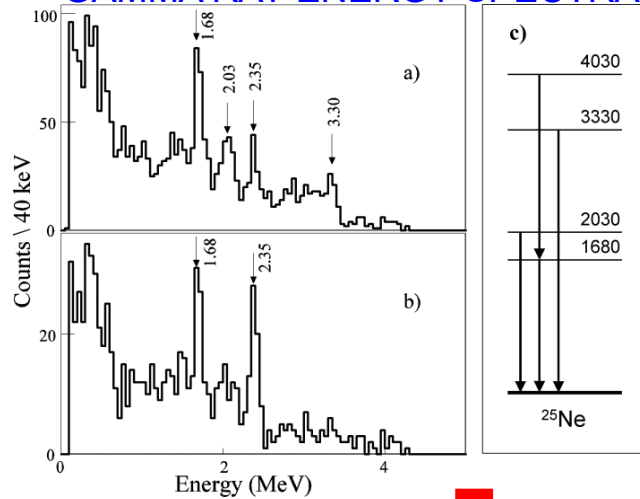
^{24}Ne ISOL
 $\tau = 3.38$ min
 100,000 pps
 SPIRAL 1

TIARA 



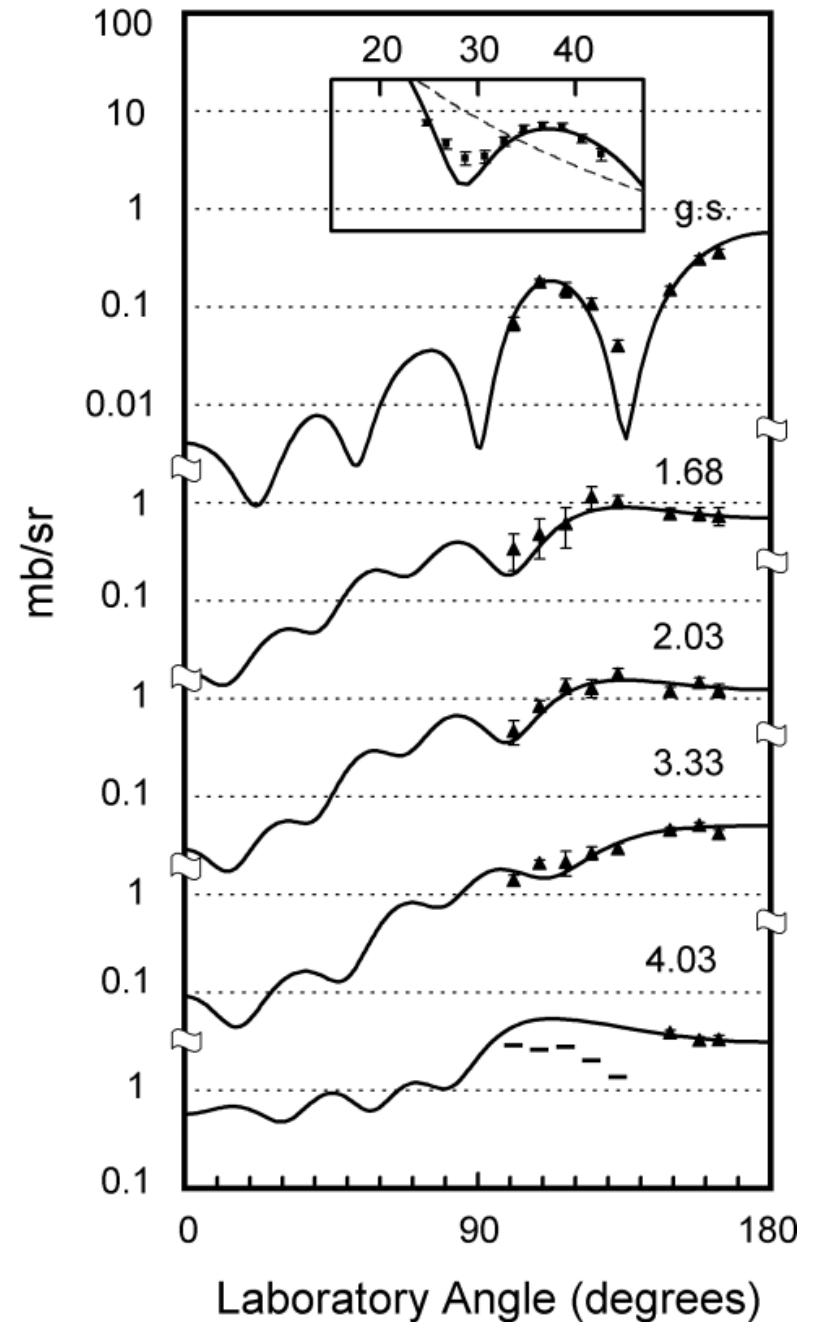
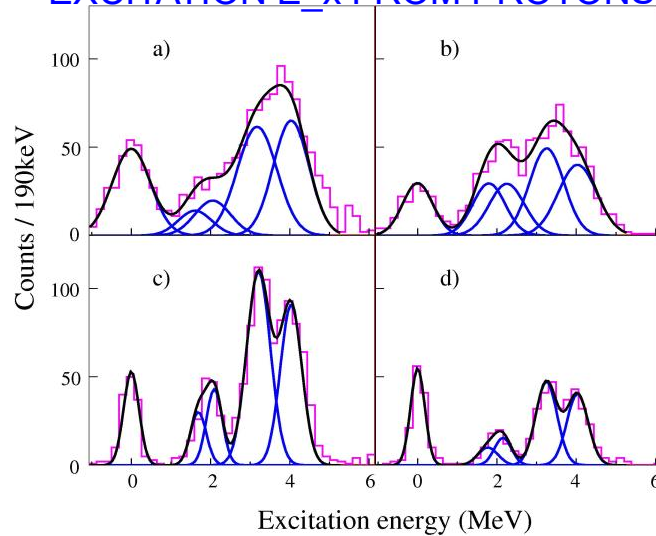
$^{24}\text{Ne}(d, p\gamma) ^{25}\text{Ne}$ – ANALYSIS

GAMMA RAY ENERGY SPECTRA

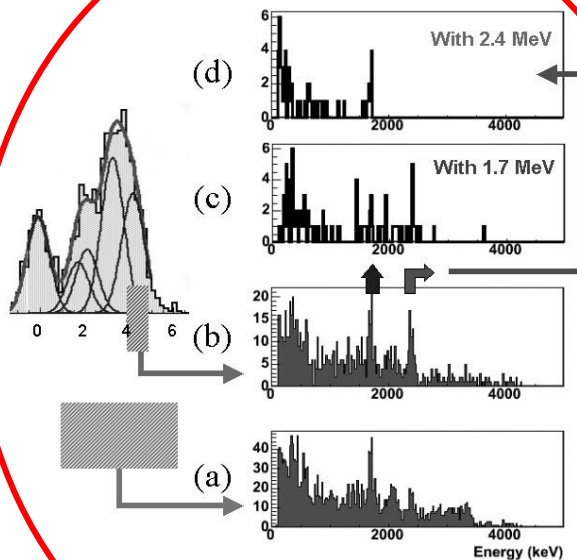


FIX E_x

EXCITATION E_x FROM PROTONS

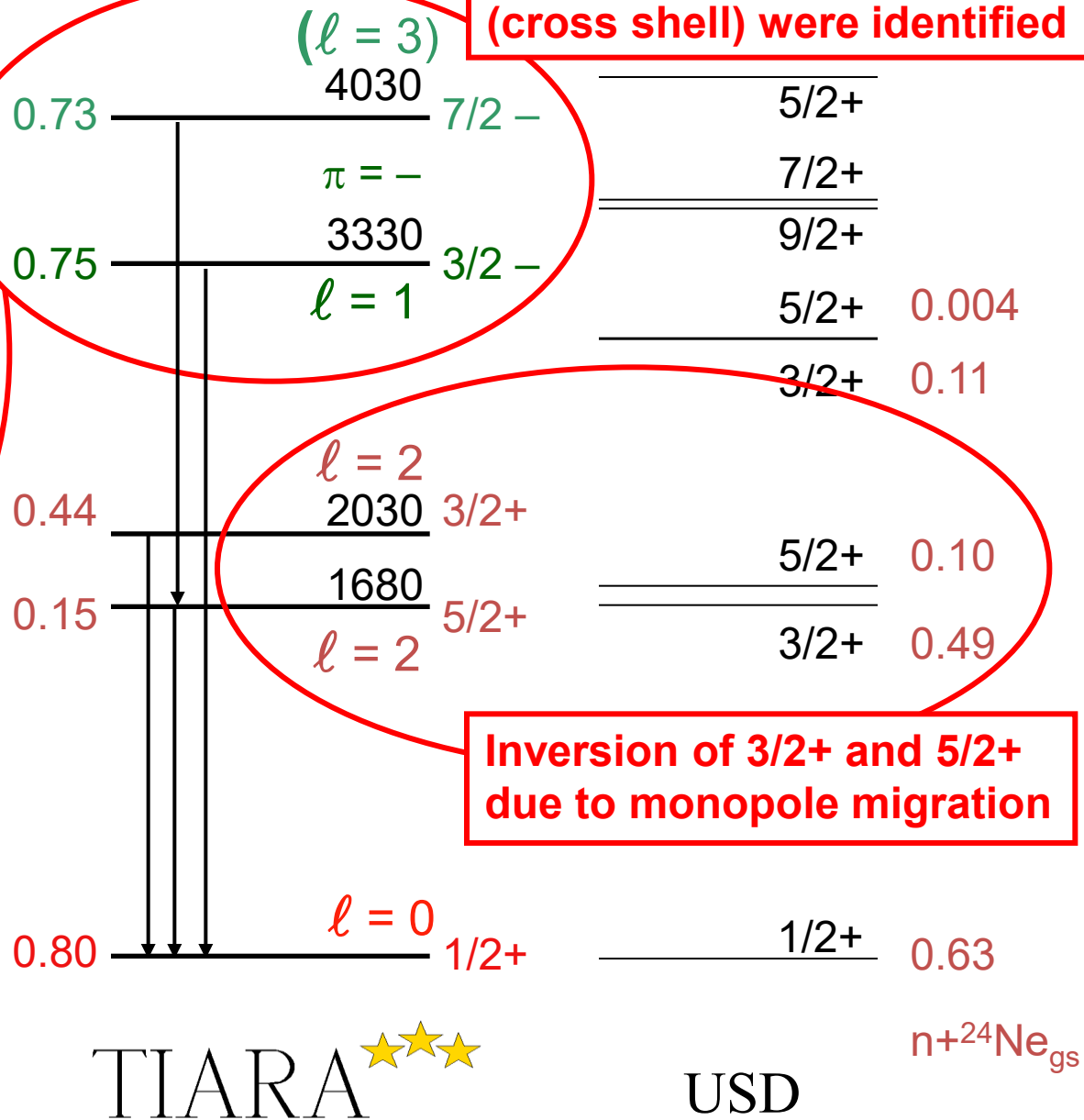


$^{24}\text{Ne}(d,p\gamma)^{25}\text{Ne}$ – ANALYSIS



In ^{25}Ne we used
gamma-gamma coincidences to **distinguish spins**
and go beyond orbital AM
FIRST QUADRUPLE
COINCIDENCE (p-HI- γ - γ)
RIB TRANSFER DATA

**Negative parity states
(cross shell) were identified**

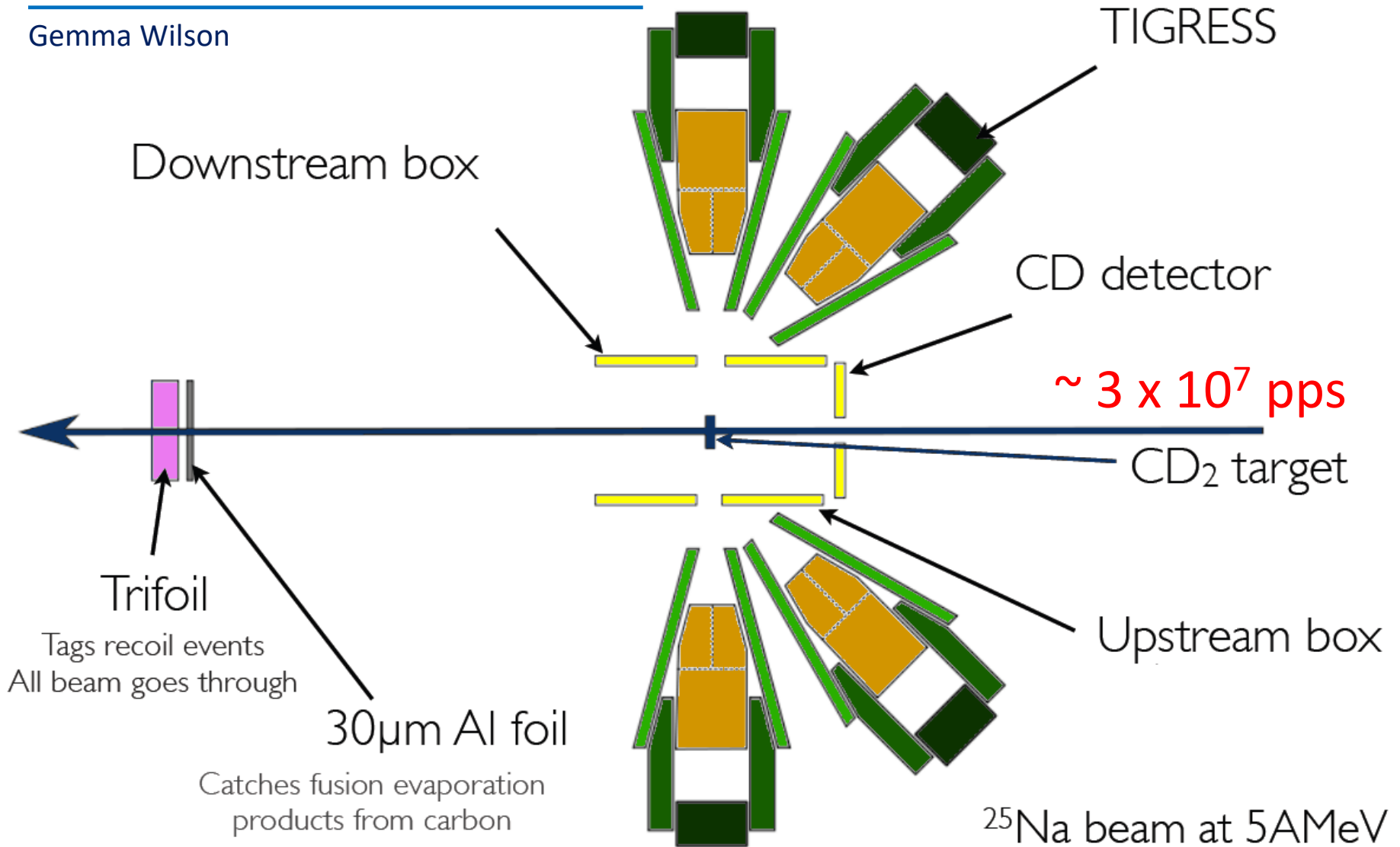


**Inversion of 3/2+ and 5/2+
due to monopole migration**

$^{25}\text{Na}(d,p\gamma)^{26}\text{Na}$ – we saw how one extra $d_{5/2}$ proton interacts near $N=16$

SHARC at ISAC2 at TRIUMF

Gemma Wilson



**SHARC chamber
(compact Si box)**

TIGRESS

TRIFOIL @ zero degrees

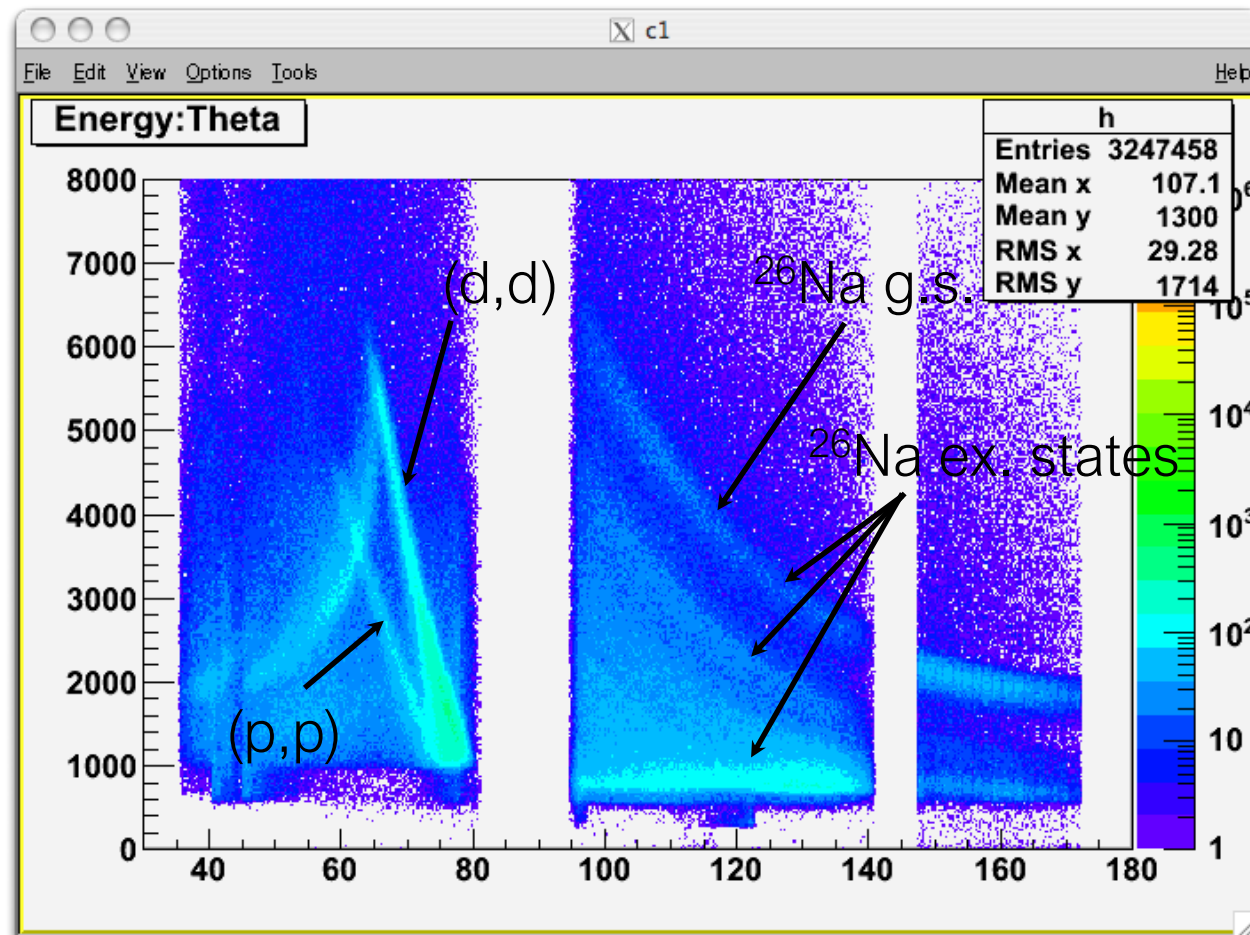
BEAM

TIGRESS

**Bank of 500 preamplifiers
cabled to TIG10 digitizers**

$^{25}\text{Na}(d,p\gamma)^{26}\text{Na}$

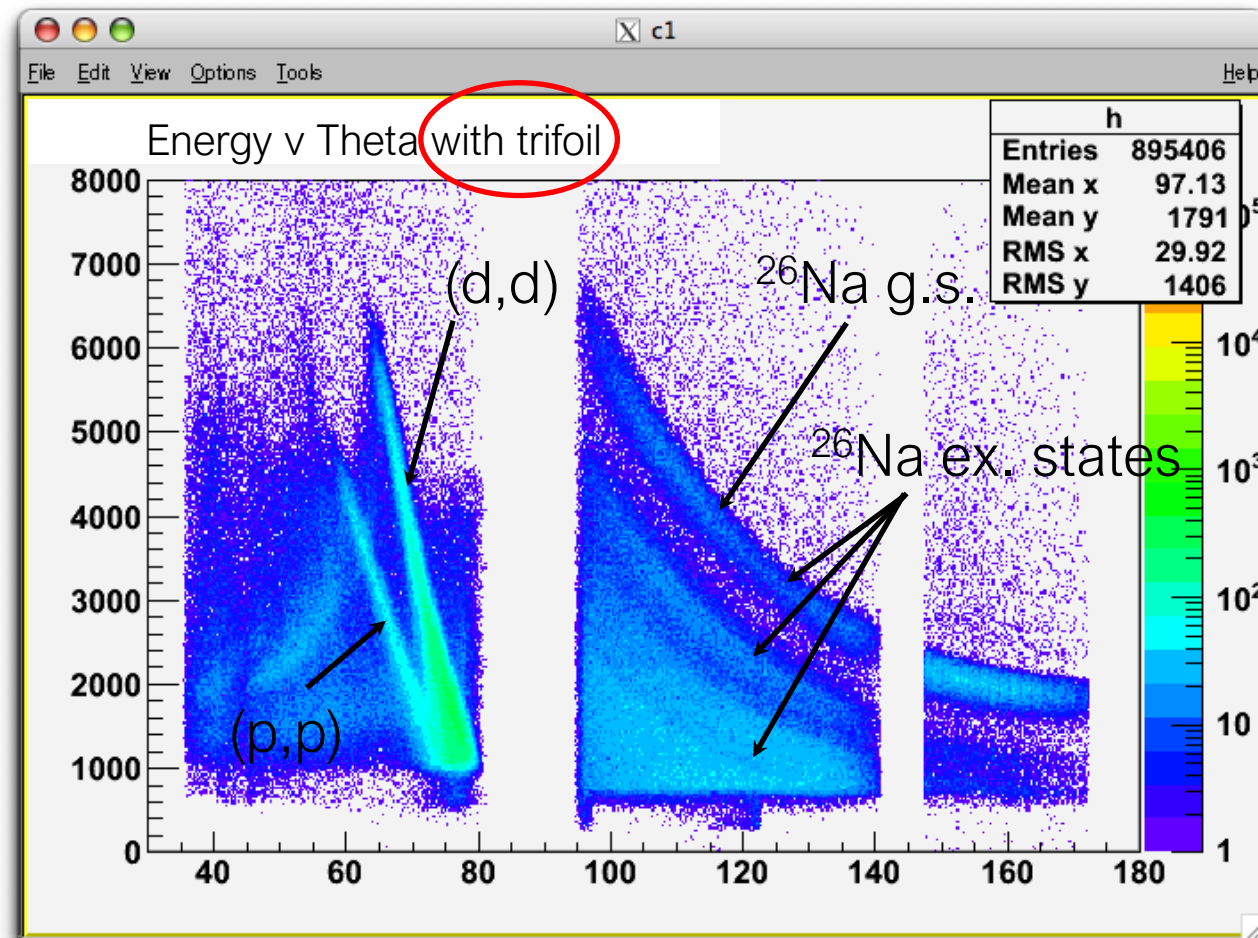
Kinematical Analysis: E vs θ



G.L. Wilson, W.N. Catford *et al.*, Phys. Lett. B759, 417 (2016)



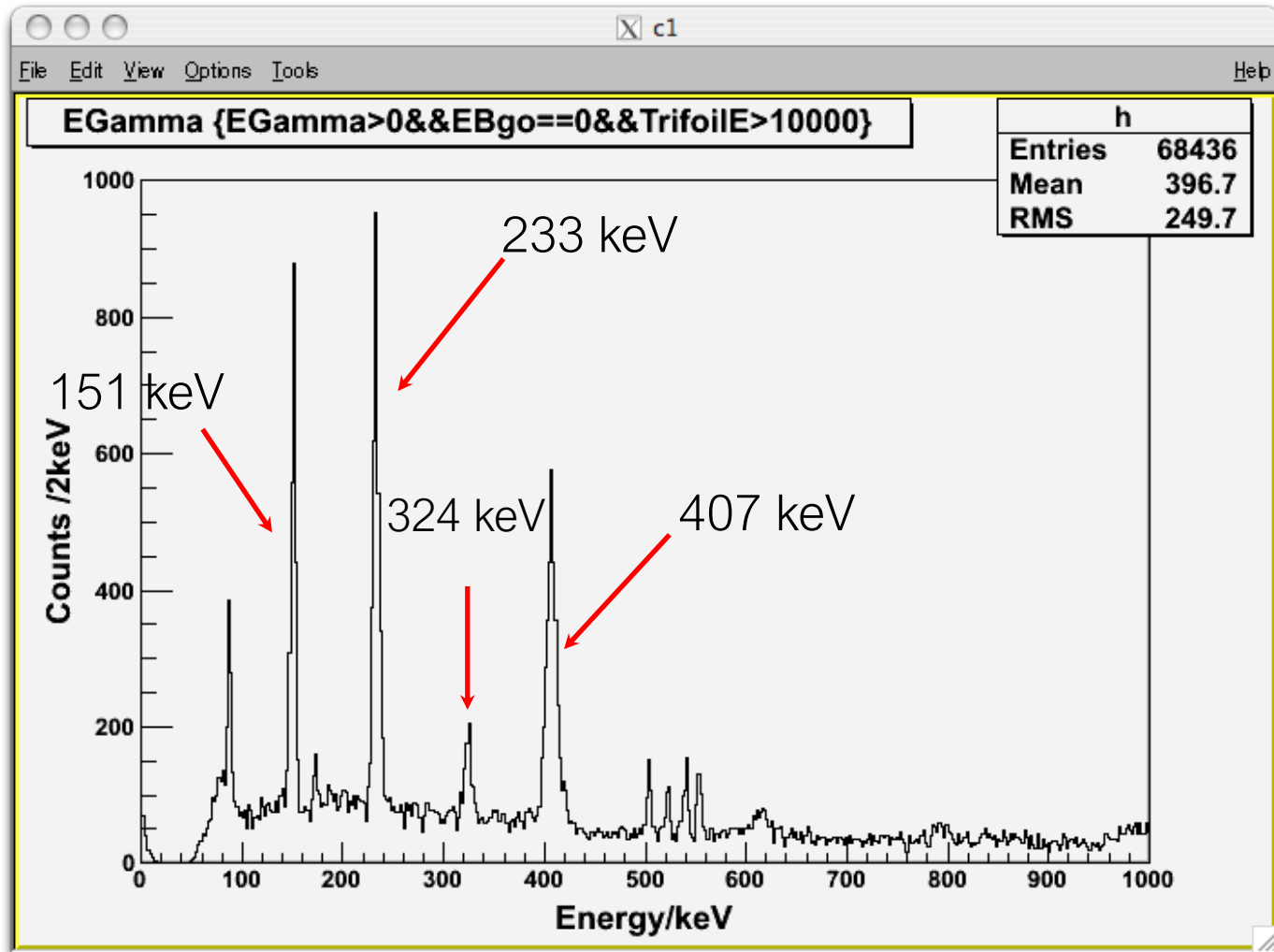
Kinematical Analysis: E vs θ



G.L. Wilson, W.N. Catford *et al.*, Phys. Lett. B759, 417 (2016)

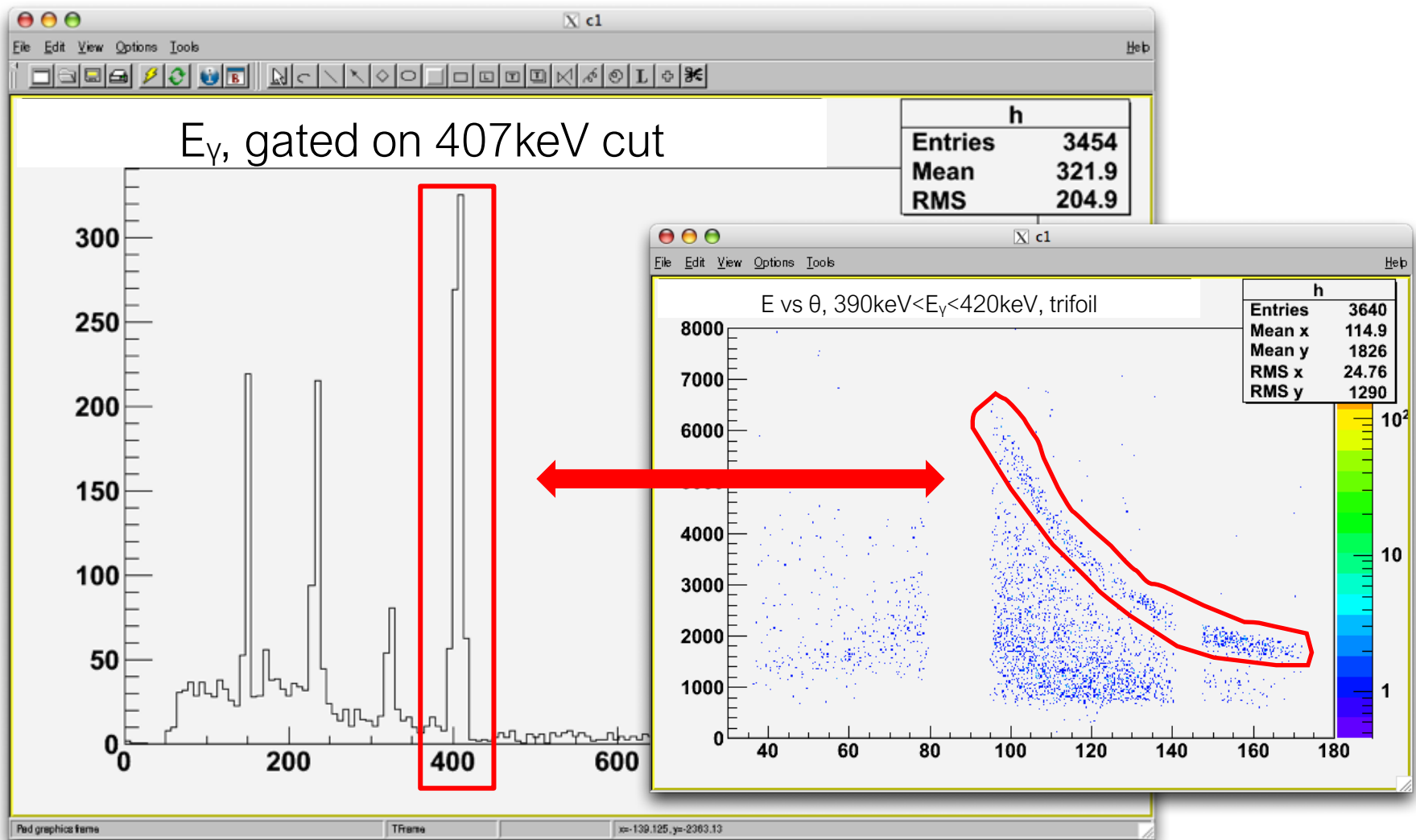
$^{25}\text{Na}(d,\gamma)^{26}\text{Na}$

Gamma-ray Analysis: precise energies



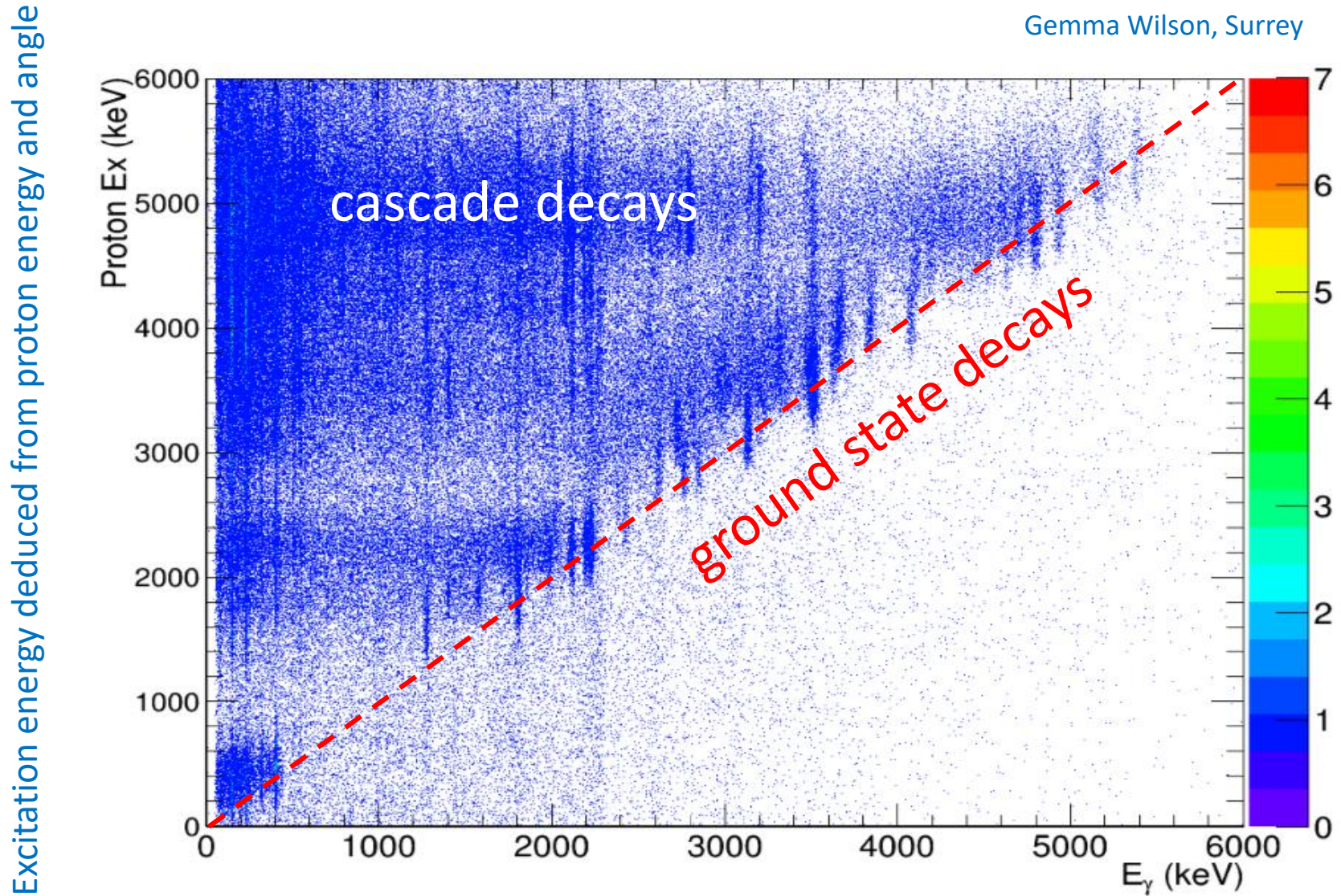
$^{25}\text{Na}(d,p\gamma)^{26}\text{Na}$

Combining the particle and γ -ray Analyses



Data from $d(^{25}\text{Na}, p)^{26}\text{Na}$ at 5 MeV/A using SHARC/ISAC2/TRIUMF

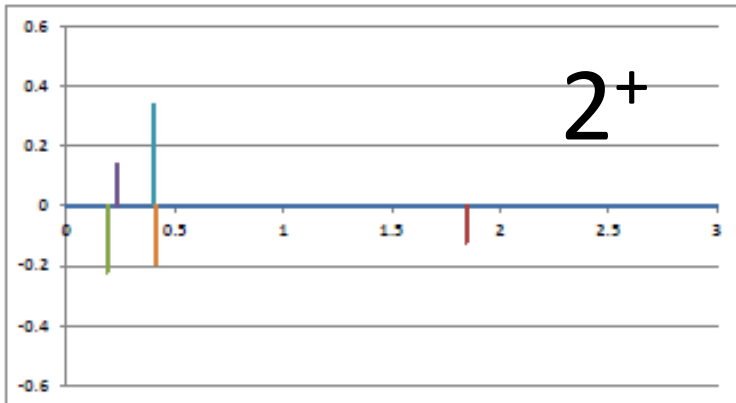
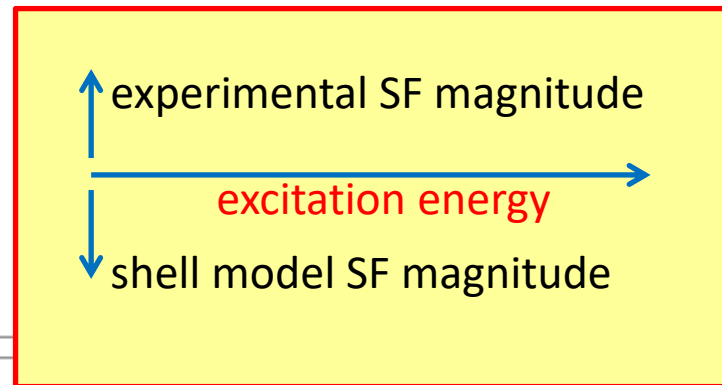
Gemma Wilson, Surrey



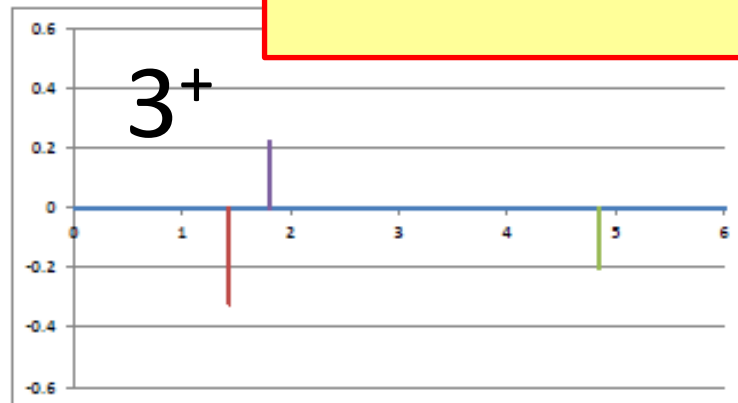
Doppler corrected ($\beta=0.10$) gamma ray energy measured in TIGRESS

Shell Model Predictions (and new candidates) for ^{26}Na states expected in (d,p)...

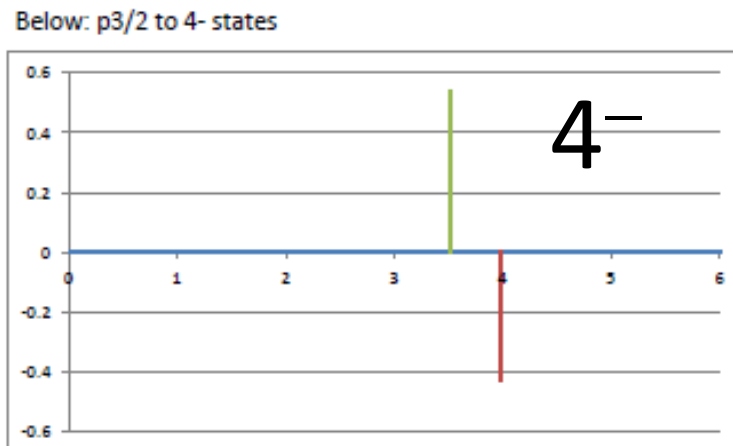
Comparison of spectroscopic strength in
theory and experiment



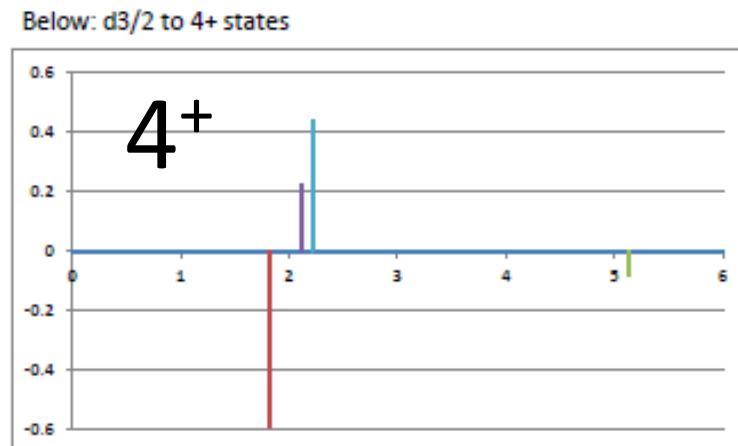
Above: s1/2 strength for 2^+ states



Above: d3/2 to 3^+ states



Below: p3/2 to 4^- states



Below: d3/2 to 4^+ states

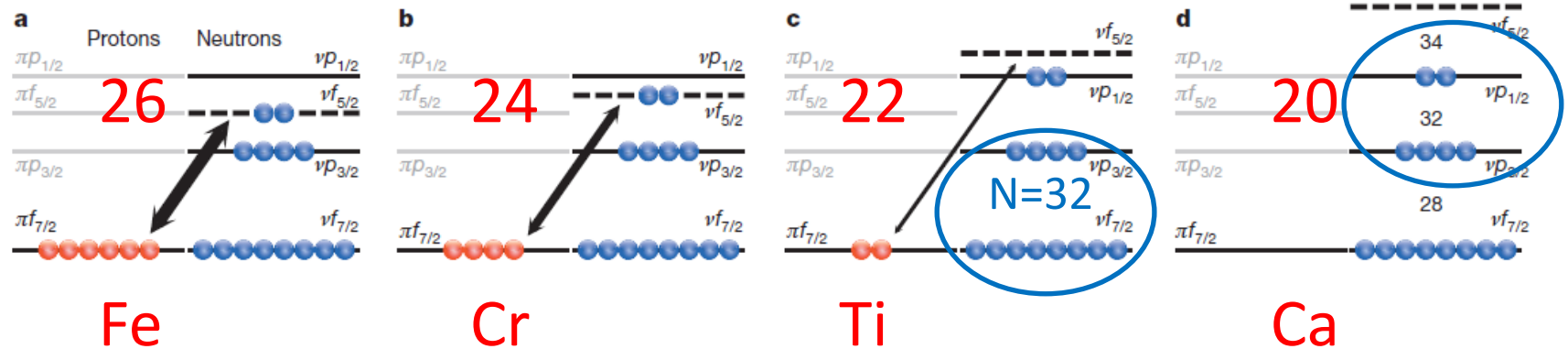
Study of $d(^{47}\text{K}, p)^{48}\text{K}$ at 7.7 MeV/A using MUGAST/AGATA/SPIRAL1

We measured how proton filling affects neutron energies; example –

Removing $\pi(f_{7/2})$

$\nu(f_{5/2})$ rises above $\nu(p_{1/2})$
N=32 and N=34 emerge as magic

RESEARCH LETTER



nature

doi:10.1038/nature12522

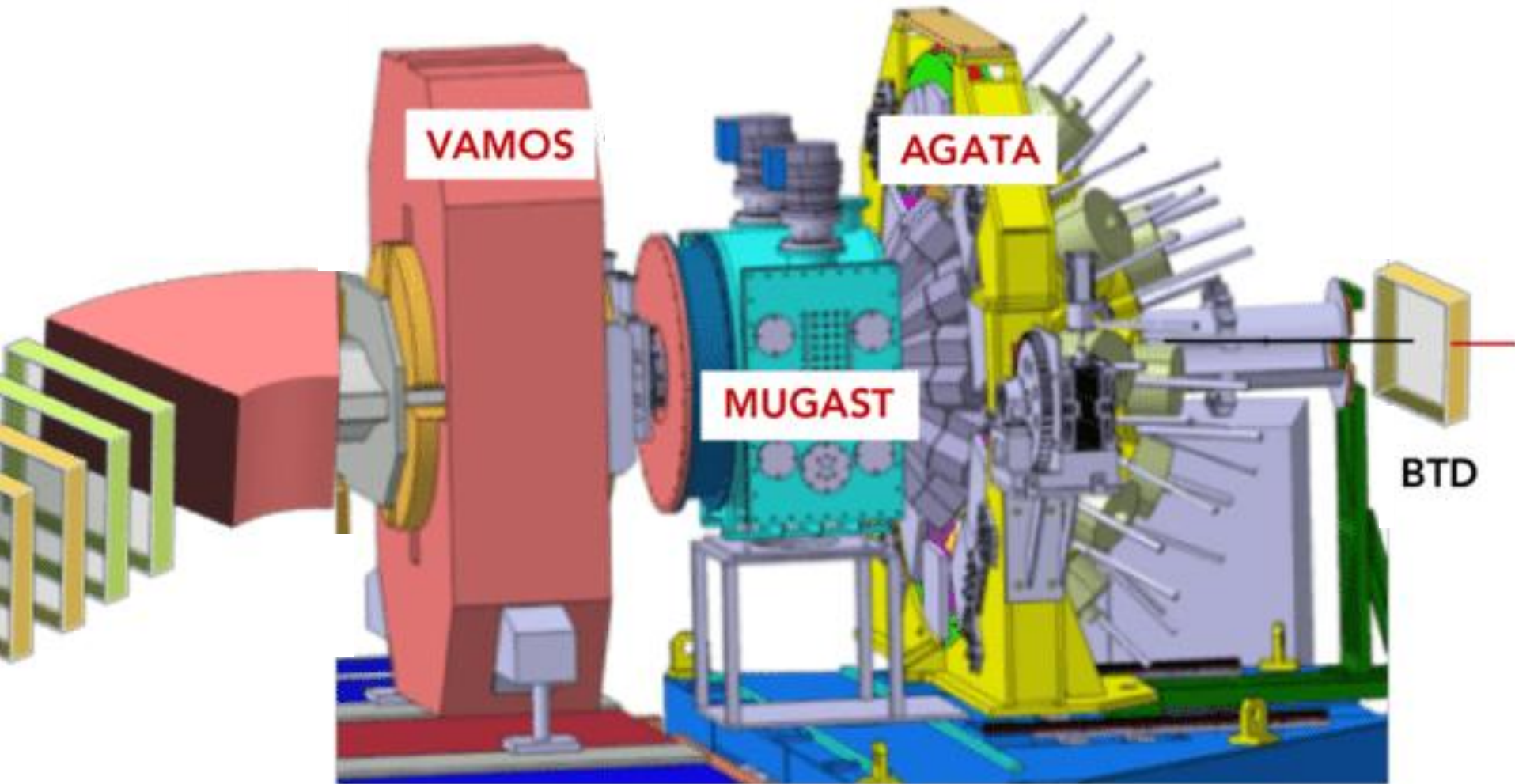
Evidence for a new nuclear ‘magic number’ from the level structure of ^{54}Ca

D. Steppenbeck¹, S. Takeuchi², N. Aoi³, P. Doornenbal², M. Matsushita¹, H. Wang², H. Baba², N. Fukuda², S. Go¹, M. Honma⁴, J. Lee², K. Matsui⁵, S. Michimasa¹, T. Motobayashi², D. Nishimura⁶, T. Otsuka^{1,5}, H. Sakurai^{2,5}, Y. Shiga⁷, P.-A. Söderström², T. Sumikama⁸, H. Suzuki², R. Taniuchi⁵, Y. Utsuno⁹, J. J. Valiente-Dobón¹⁰ & K. Yoneda²

Study of $d(^{47}\text{K}, p)^{48}\text{K}$

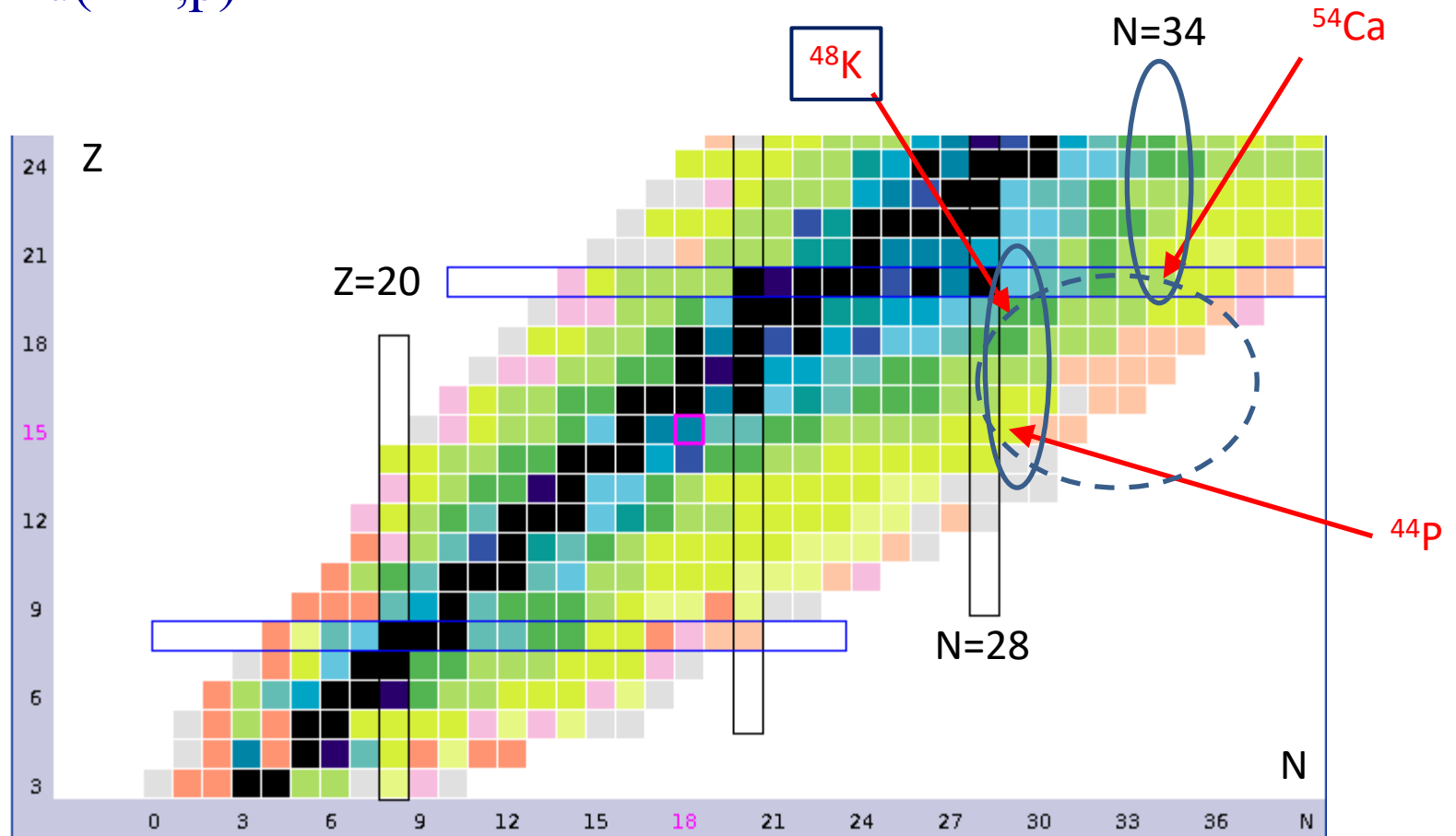


Beam: ^{47}K at 7.7 MeV/u (SPIRAL1)
Intensity: 5×10^5 pps
Purity: $\sim 100\%$



C.J. Paxman, A. Matta, W.N. Catford *et al.*, Phys. Rev. Lett. **134**, 162504 (2025)

Study of $d(^{47}\text{K}, p)^{48}\text{K}$

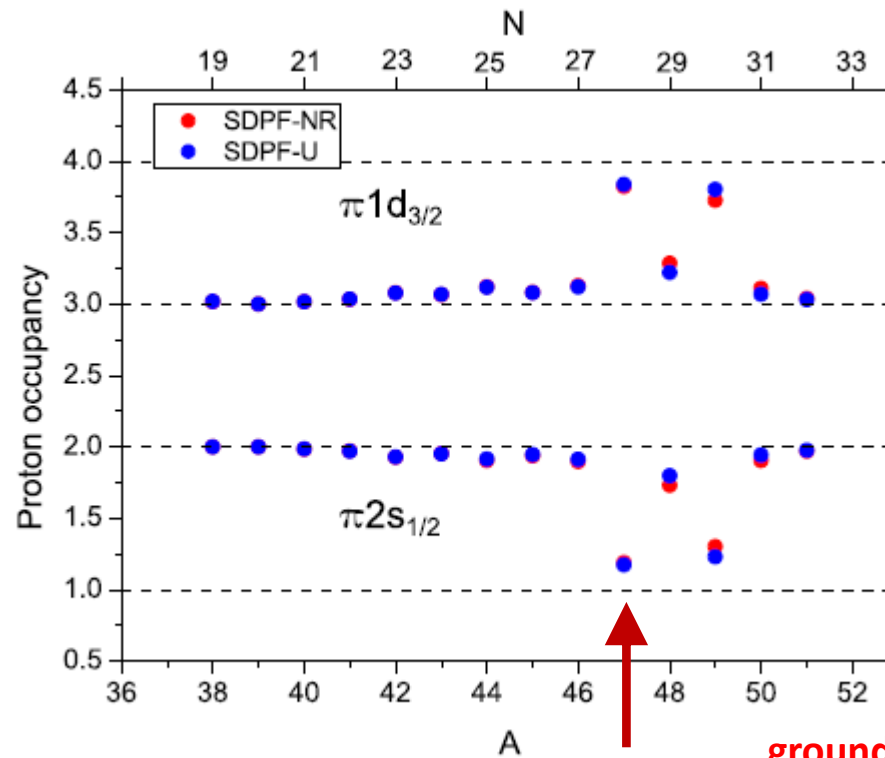


PLAN: use the reaction (d, p) to add a neutron to $p_{3/2}$, $p_{1/2}$ and $f_{5/2}$

... study the interaction with the odd proton way down in $s_{1/2}$

Shell structure of potassium isotopes deduced from their magnetic moments

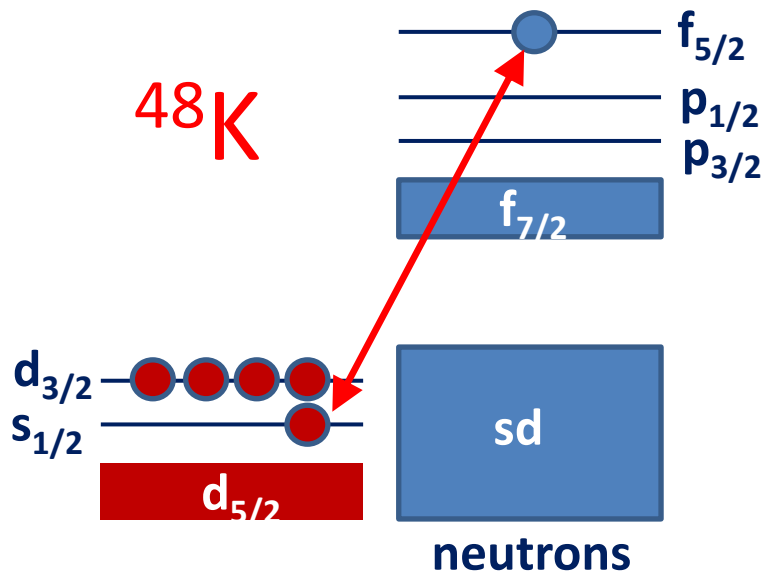
J. Papuga,^{1,*} M. L. Bissell,¹ K. Kreim,² C. Barbieri,³ K. Blaum,² M. De Rydt,¹ T. Duguet,^{4,5} R. F. Garcia Ruiz,¹ H. Heylen,¹
 M. Kowalska,⁶ R. Neugart,⁷ G. Neyens,¹ W. Nörtershäuser,^{7,8} M. M. Rajabali,¹ R. Sánchez,^{9,10}
 N. Smirnova,¹¹ V. Somà,⁴ and D. T. Yordanov^{2,12}



experimentally
confirmed
inversion of proton
orbital occupancies

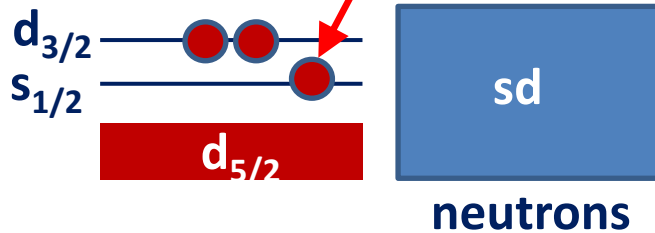
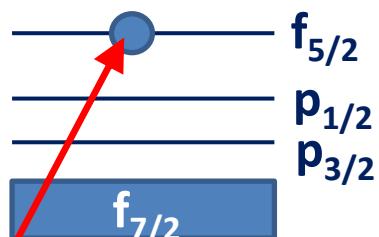
^{47}K
 ground state is $1/2^+$
 (odd proton in $s_{1/2}$)
 (same happens in ^{45}Cl)

^{48}K



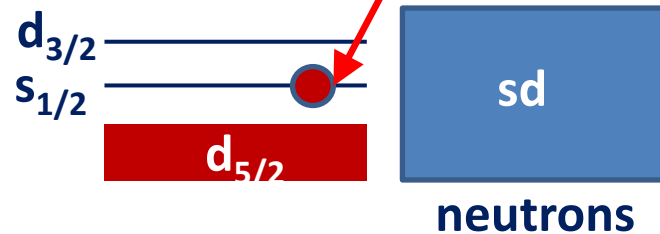
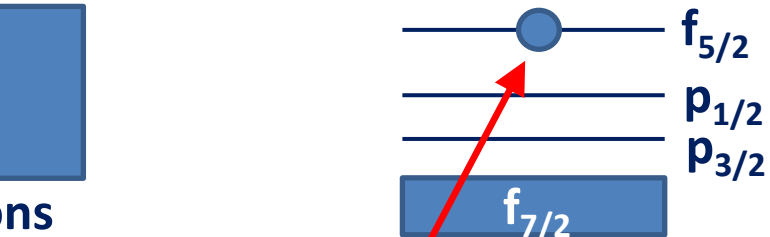
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^{46}Cl



neutrons

^{44}P

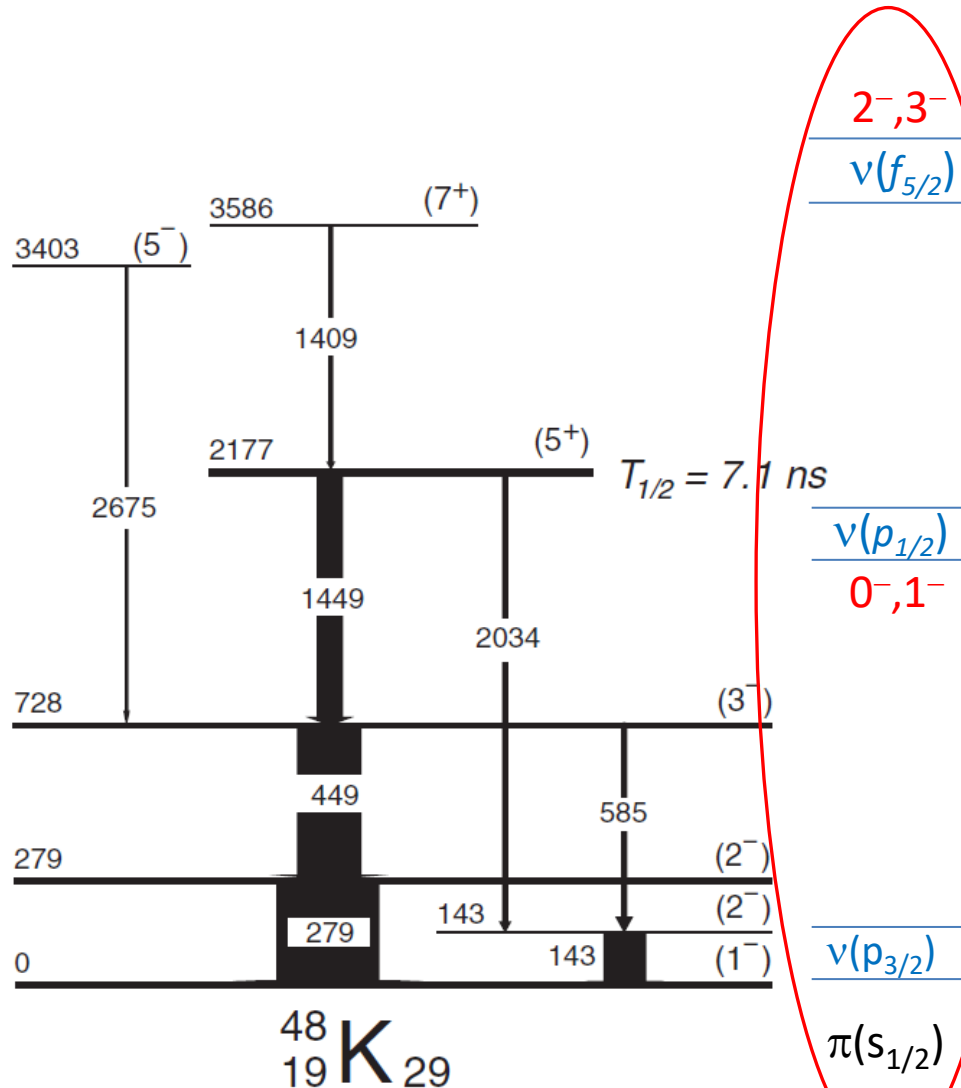


neutrons



Coupling of the proton-hole and neutron-particle states in the neutron-rich ^{48}K isotope

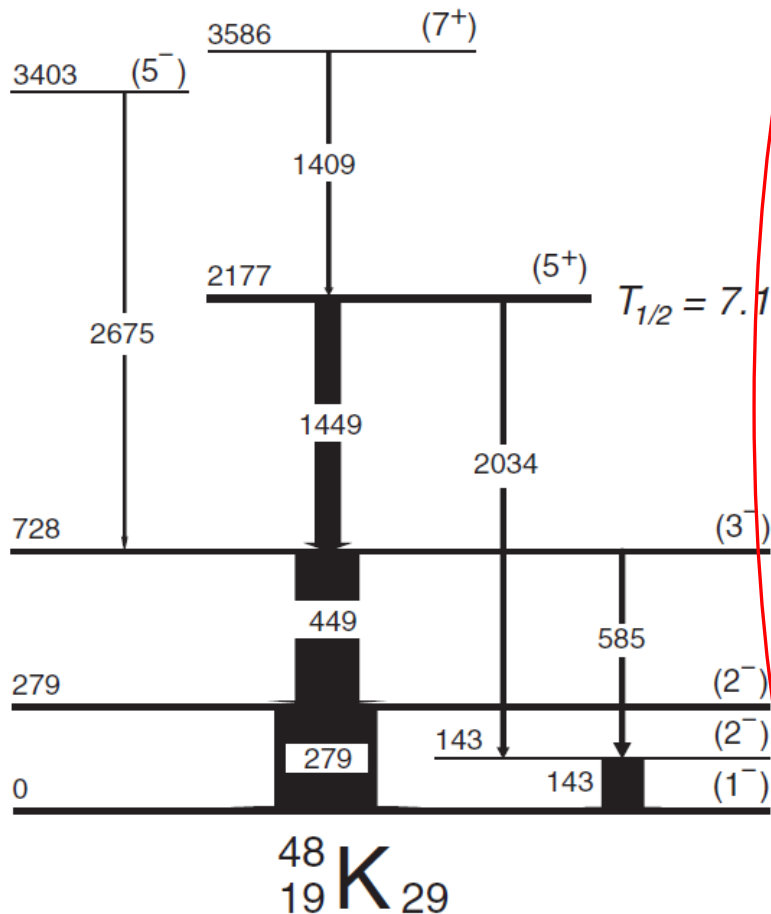
W. Królas,¹ R. Broda,¹ B. Fornal,¹ R. V. F. Janssens,² A. Gadea,^{3,4} S. Lunardi,⁵ J. J. Valiente-Dobon,³ D. Mengoni,⁵
 N. Márginean,^{3,6} L. Corradi,³ A. M. Stefanini,³ D. Bazzacco,⁵ M. P. Carpenter,² G. De Angelis,³ E. Farnea,⁵ E. Fioretto,³
 F. Galtarossa,⁵ T. Lauritsen,² G. Montagnoli,⁵ D. R. Napoli,³ R. Orlandi,³ T. Pawlat,¹ I. Pokrovskiy,³ G. Pollaro,⁷ E. Sahin,³
 F. Scarlassara,⁵ D. Seweryniak,² S. Szilner,⁸ B. Szpak,¹ C. A. Ur,⁵ J. Wrzesiński,¹ and S. Zhu²



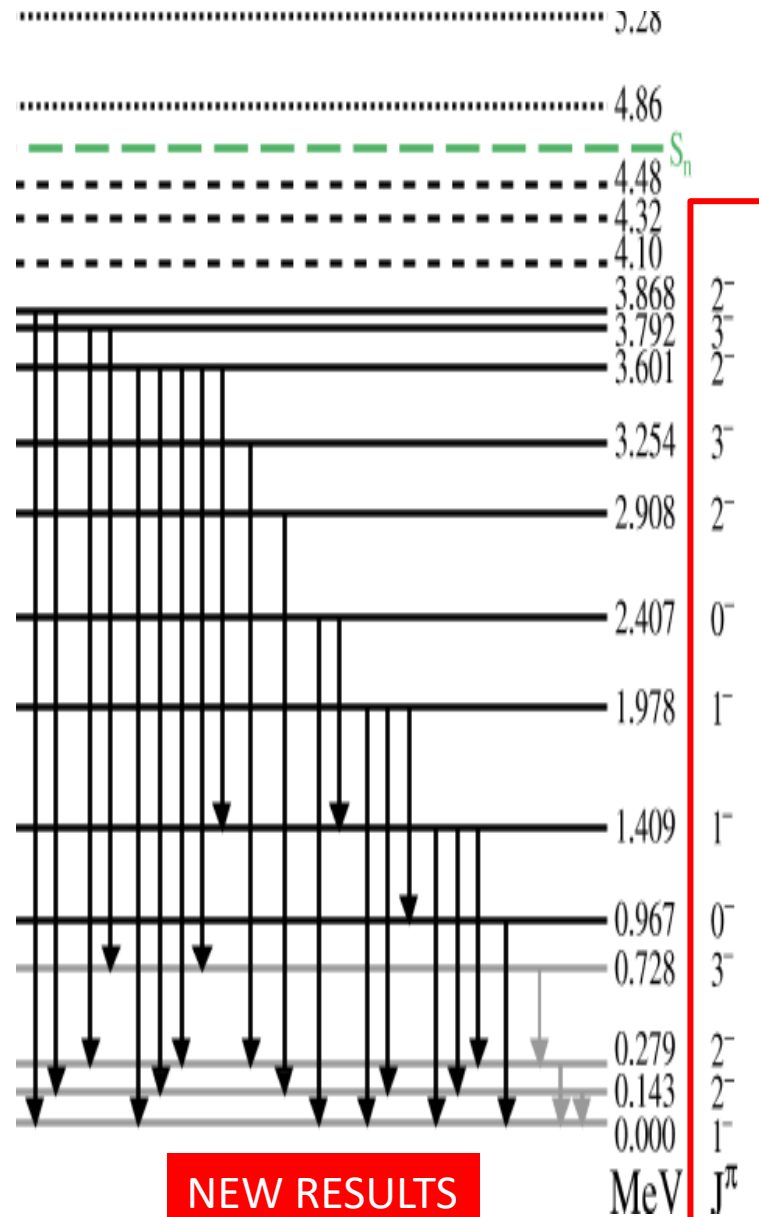
Doublet for each neutron orbital
 (coupling to odd proton in $s_{1/2}$)

Coupling of the proton-hole and neutron-particle states in the neutron-rich ^{48}K isotope

W. Królas,¹ R. Broda,¹ B. Fornal,¹ R. V. F. Janssens,² A. Gadea,^{3,4} S. Lunardi,⁵ J. J. Valiente-Dobon,³ D. Mengoni,⁵ N. Márginean,^{3,6} L. Corradi,³ A. M. Stefanini,³ D. Bazzacco,⁵ M. P. Carpenter,² G. De Angelis,³ E. Farnea,⁵ E. Fioretto,³ F. Galtarossa,⁵ T. Lauritsen,² G. Montagnoli,⁵ D. R. Napoli,³ R. Orlandi,³ T. Pawlat,¹ I. Pokrovskiy,³ G. Pollaro,⁷ E. Sahin,³ F. Scarlassara,⁵ D. Seweryniak,² S. Szilner,⁸ B. Szpak,¹ C. A. Ur,⁵ J. Wrzesiński,¹ and S. Zhu²

 $2^-, 3^-$ $\nu(f_{5/2})$ $\nu(p_{1/2})$ $0^-, 1^-$ $\nu(p_{3/2})$ $\pi(s_{1/2})$

C.J. Paxman, A. Matta, W.N. Catford *et al.*,
Phys. Rev. Lett. **134**, 162504 (2025)



sdpfmu

^{48}K

sdpfu

$$\pi(s_{1/2}) \otimes \nu(p_{3/2} p_{1/2} f_{5/2})$$

“weak coupling”

excitation energy (MeV)

6

5

4

3

2

1

0

$\nu(f_{5/2})$

$\nu(p_{1/2})$

$\nu(p_{3/2})$

$\pi(s_{1/2})$

3^-

2^-

1^-

0^-

1^-

2^-

$\pi(s_{1/2})$

$\nu(p_{1/2})$

$\nu(p_{3/2})$

sdpfmu

^{48}K

sdpfu

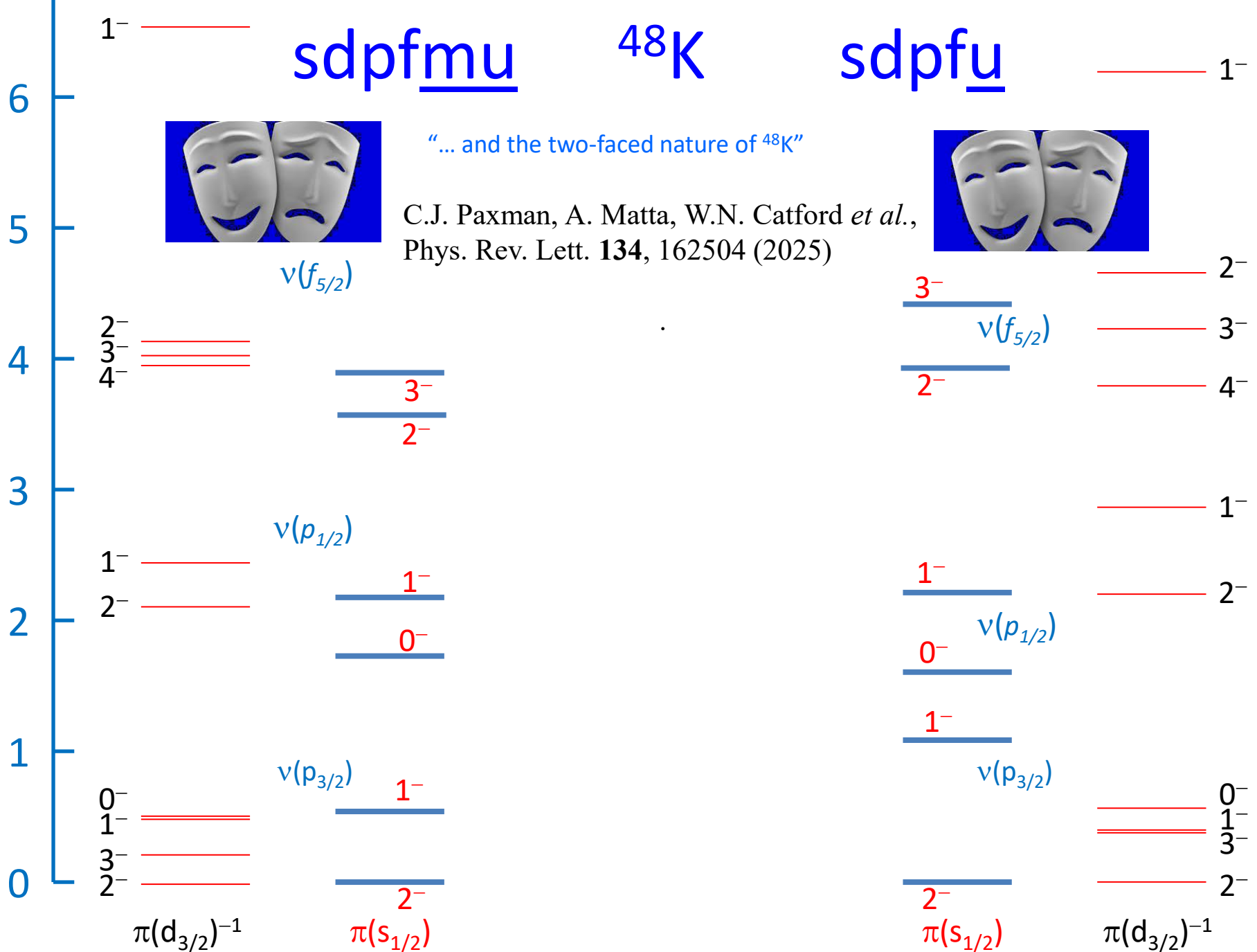


"... and the two-faced nature of ^{48}K "

C.J. Paxman, A. Matta, W.N. Catford *et al.*,
Phys. Rev. Lett. **134**, 162504 (2025)



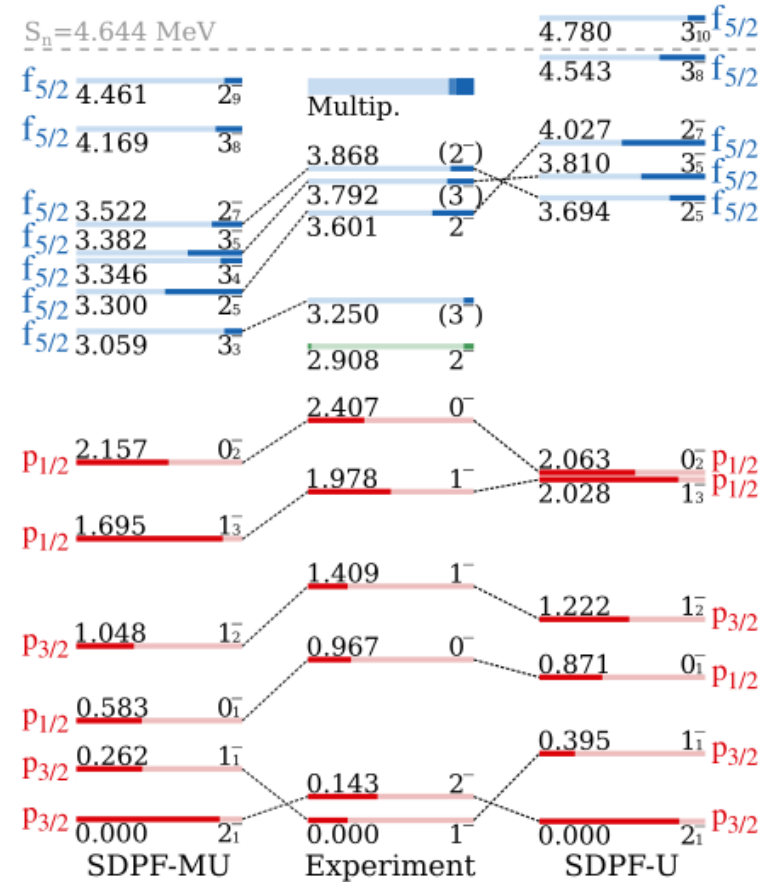
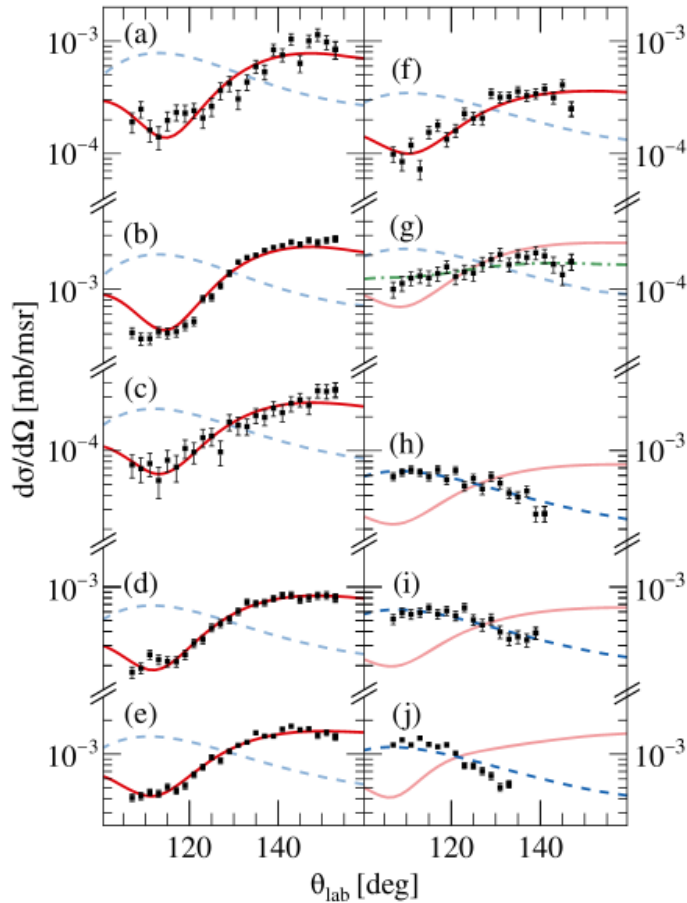
excitation energy (MeV)



Study of $d(^{47}\text{K},p)^{48}\text{K}$ at 7.7 MeV/A using MUGAST/AGATA/SPIRAL1

Angular momentum transfer from $d\sigma/d\Omega$,
Spin/parity using gamma-decay branching

Use measured spectroscopic factors to link
to shell model states and choose best model



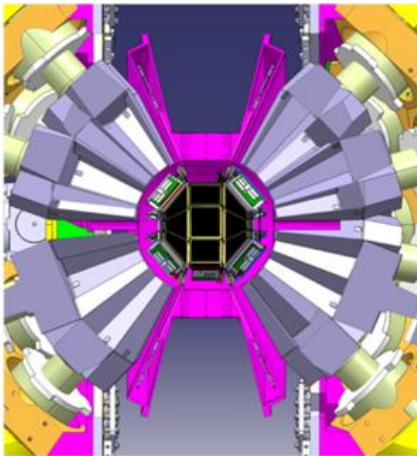
Some future perspectives...

Remember, for very low beam intensities we have ACTAR, AT-TPC
There are also solenoidal detectors (HELIOS, ISS, SOLARIS)

Whenever possible, we benefit from also measuring gamma-rays...

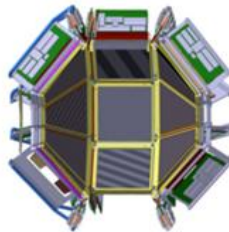
GANIL

AGATA



France, Italy, Spain, UK
2029

ISOL beams, 10 MeV/u



GRETA

FAUST

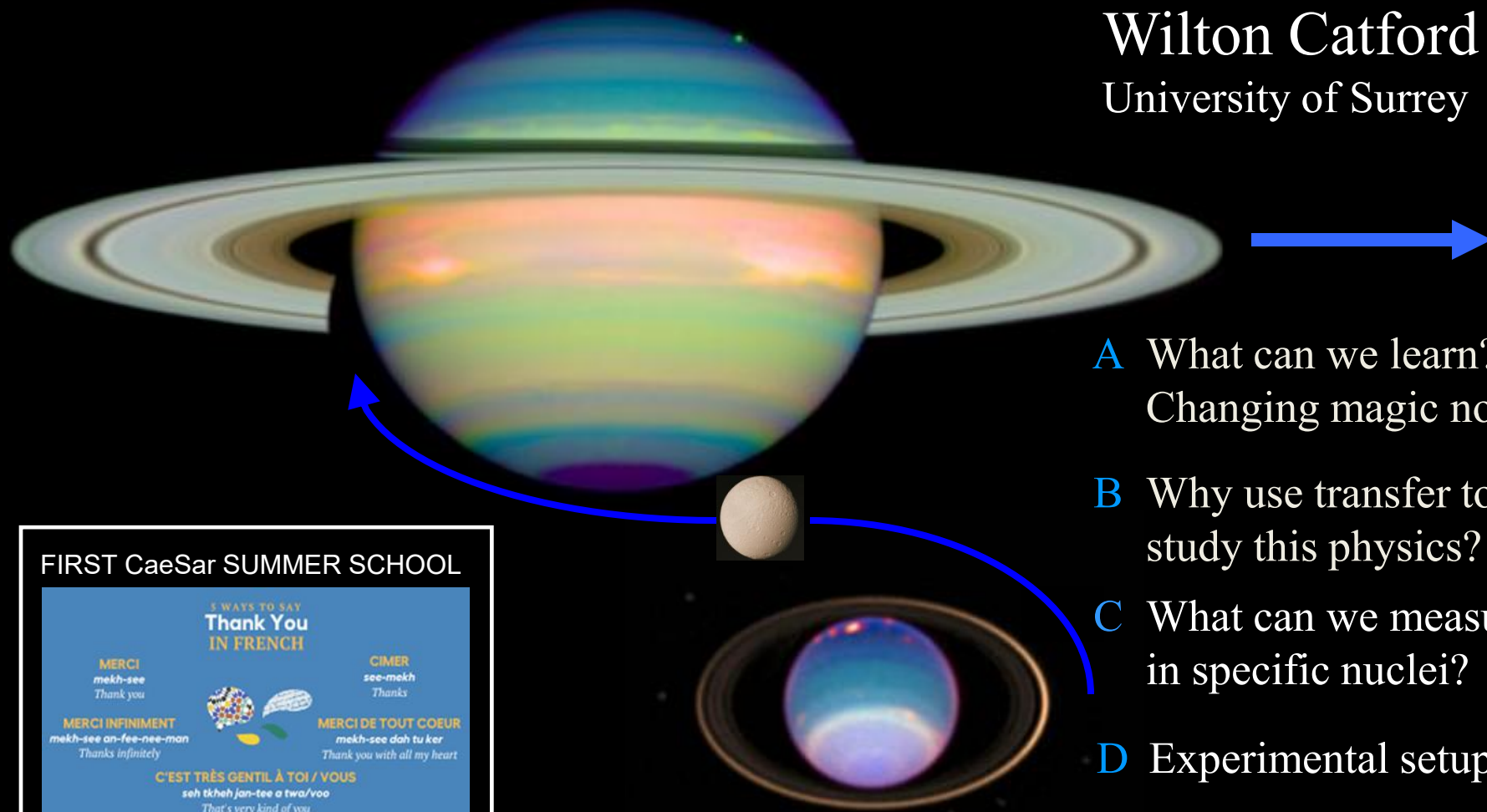


UK, USA
2028

In-flight beams, 40 MeV/u

Exploring Unusual Nuclei with Nucleon Transfer Reactions

Wilton Catford
University of Surrey



- A What can we learn?;
Changing magic no.'s
- B Why use transfer to
study this physics?
- C What can we measure
in specific nuclei?
- D Experimental setups
- E Examples of data &
Interpretation

FIRST CaeSar SUMMER SCHOOL



Pont L'Évêque**3-5 September 2025