

# — 59<sup>th</sup> Rencontres de Moriond 2025 —

*La Thuile, Italy - 29 / 03 / 25*

## CHARMING



### MAURO VALLI

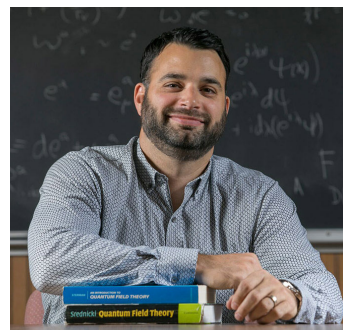
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INFN Rome



In collaboration w/: **Artemis Giannakopoulou & Patrick Meade**

arXiv:**2410.05236** — **JHEP 02 (2025) 067**



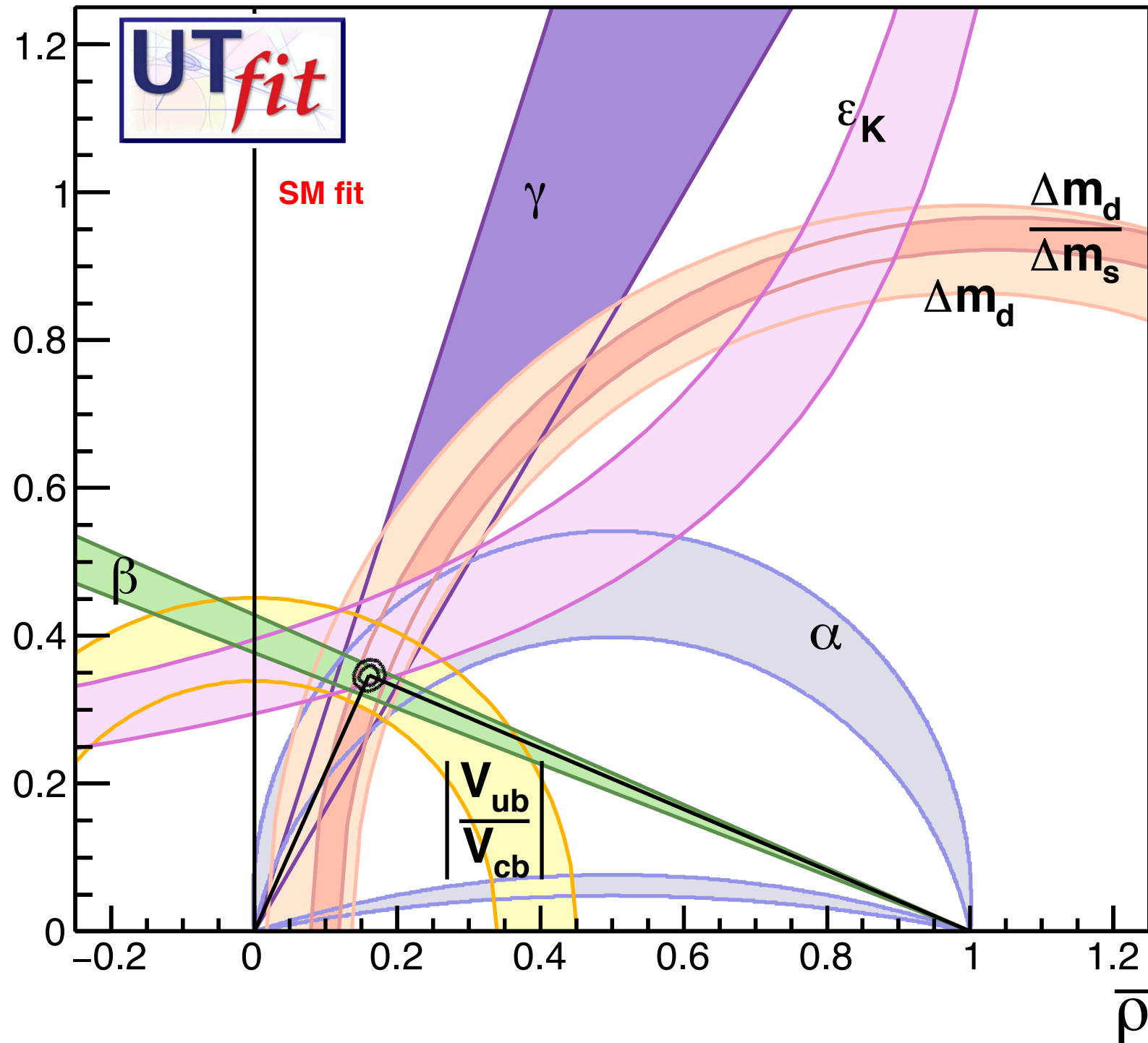
# UTA: Unitarity Triangle Analysis

*Rend.Lincei Sci.Fis.Nat.* 34 (2023)  
PoS WIFAI2023 (2024) 007

@ 95% prob

**UT**  
*fit*

SM fit



$$\bar{\rho} = 0.160 \pm 0.009 \sim 6\%$$

$$\bar{\eta} = 0.346 \pm 0.009 \sim 3\%$$

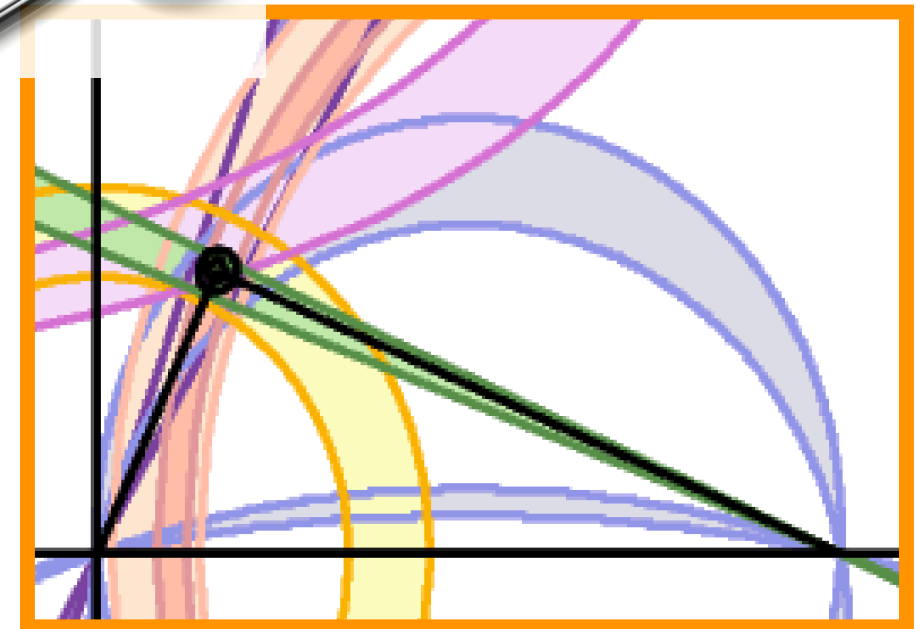
UT determination towards % precision

$$\lambda = 0.2251 \pm 0.0008$$

$$A = 0.827 \pm 0.010$$



Zoomed-in Unitarity Triangle



Analogous results from the CKMfitter collaboration.

# Flavor & BSM Physics

PoS WIFAI2023 (2024) 007

Most general  $|\Delta F| = 2$   $H_{\text{eff}}$ :

$$O_1^{q_i q_j} = \bar{q}_{jL}^\alpha \gamma_\mu q_{iL}^\alpha \bar{q}_{jL}^\beta \gamma^\mu q_{iL}^\beta$$

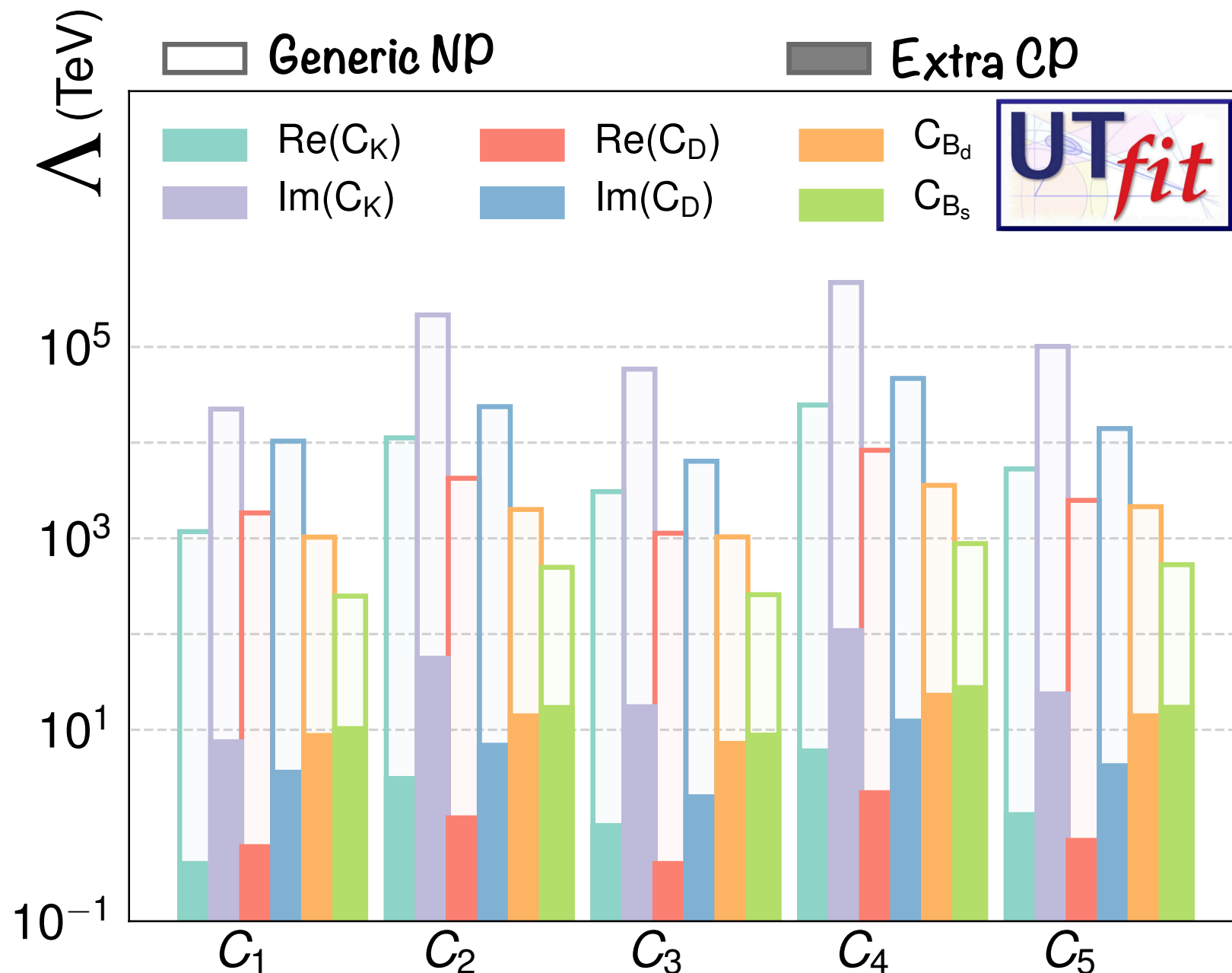
$$O_2^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jR}^\beta q_{iL}^\beta \quad \text{SM}$$

$$O_3^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jR}^\beta q_{iL}^\alpha$$

$$O_4^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jL}^\beta q_{iR}^\beta$$

$$O_5^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jL}^\beta q_{iR}^\alpha$$

+ chirally flipped  $\tilde{O}_{1,2,3}^{q_i q_j}$



○ **Generic NP** = no SM protection, i.e.:  $C(\Lambda) \sim 1/\Lambda^2$   $\Rightarrow$   $\Lambda > 4.7 \times 10^5$  TeV

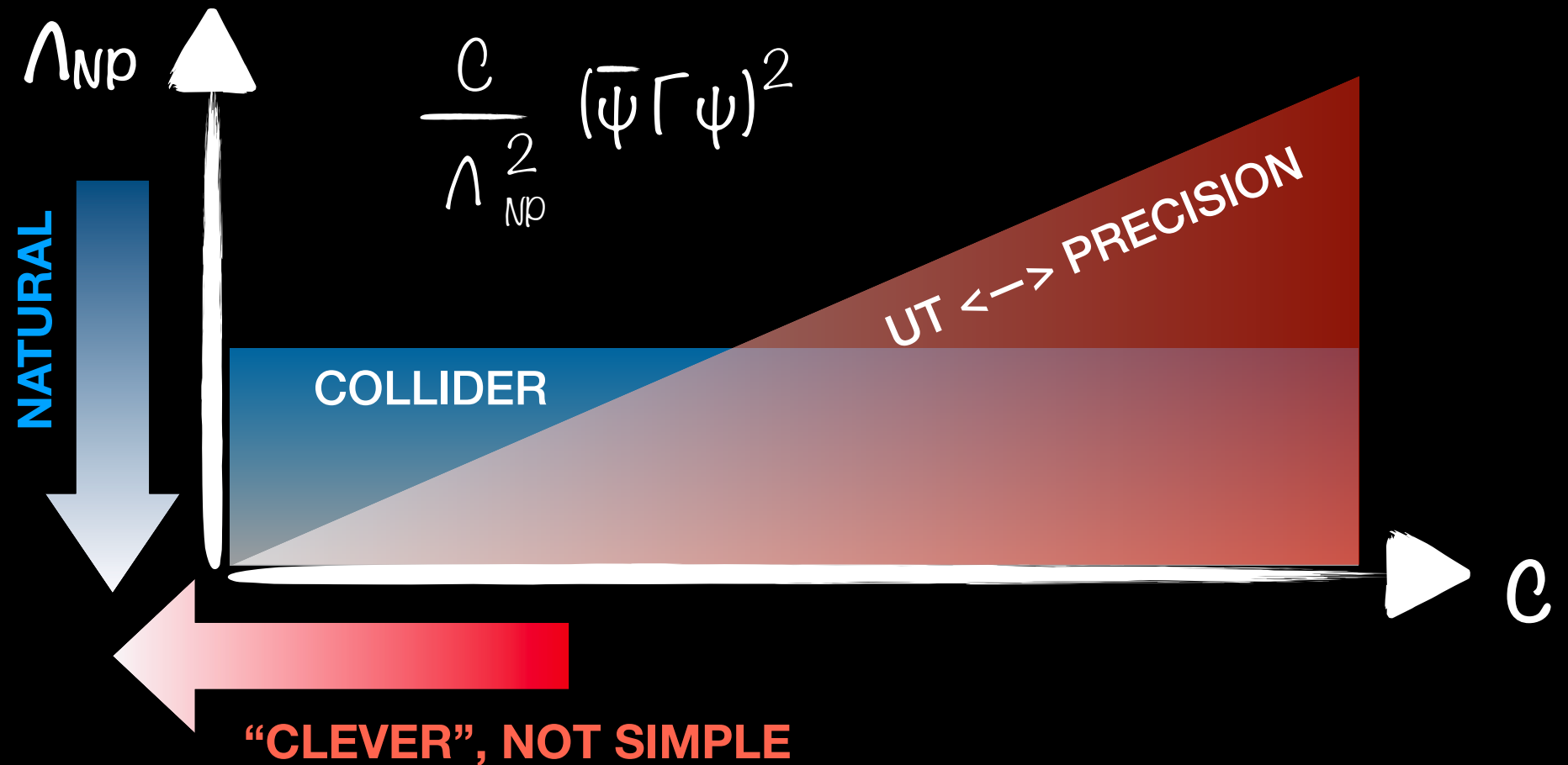
● **Extra CP** = SM-like protection but new  $O(1)$  phases  $\Rightarrow$   $\Lambda > 108$  TeV



# Lessons from Precision

- **SM UT:** Towards % precision ... overall remarkable consistency!  
—> *in the HL-LHC era we might aim at a permil test*

- **NP UT:**



**BOTTOM  
LINE**

A Theory of Flavor is either **VERY CLEVER** or **"JUST" UNNATURAL**

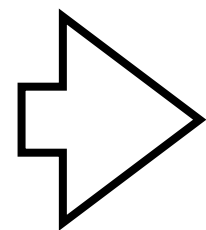
# How To Be Clever ...

**SEIBERGOLOGY:** Yukawas are spurions breaking the **flavor group**

$$U(3)^5 = SU(3)_Q \times SU(3)_d \times SU(3)_u \times SU(3)_L \times SU(3)_e \times U(1)^5$$

$$U(1)^5 = U(1)_B \times U(1)_L \times U(1)_Y \times U(1)_{PQ} \times U(1)_e$$

E.g.:  $-\mathcal{L}_{\text{SM}} \supset Y_{ij}^d \bar{Q}_i H d_j + Y_{ij}^u \bar{Q}_i \tilde{H} u_j + h.c.$



$$Y^d \sim (3, \bar{3}, 1) \quad Y^u \sim (3, 1, \bar{3})$$

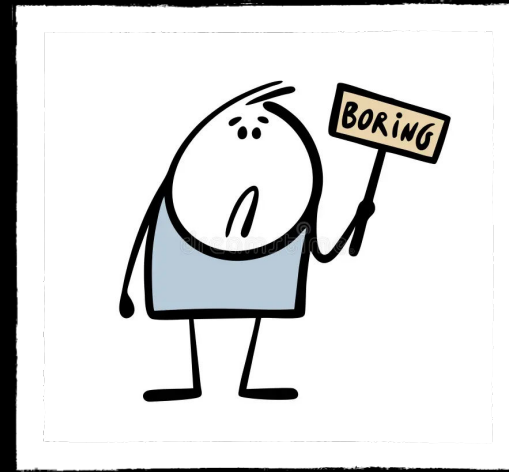
under  $SU(3)_Q \times SU(3)_d \times SU(3)_u$

**ASSUMPTION:** only spurions that break **flavor** are **SM Yukawas**

—> **MINIMAL FLAVOR VIOLATION** (*Nucl.Phys.B* 645 (2002) 155)

aka **MAXIMAL FLAVOR CONSERVATION** (see *arXiv:2402.09503*)

MFV  $\langle \text{---} \rangle$



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Can we avoid large FCNCs  
and have BSM with  
radically different flavor



# Another Way to Be Clever: **Alignment!**

See also  
V. Miralles  
talk!

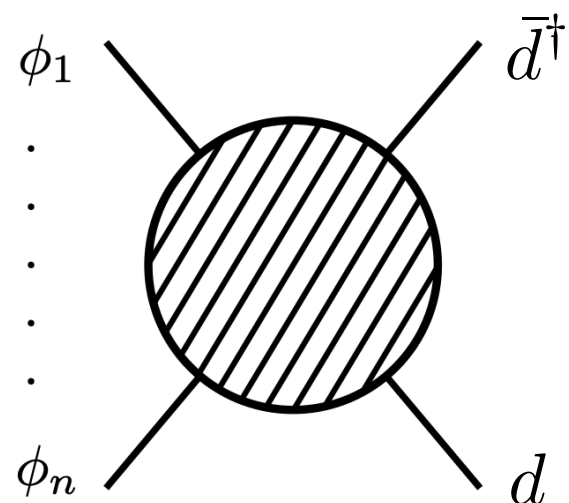
## — ANSATZ FOR ALIGNMENT IN THE UP-QUARK SECTOR —

U  
V

• **NO** breaking of the family number  $U(1)_f^3$  and of  $CP$  other than the wave-function renormalization  $Z^d$  of  $d$

• **NO** fields / spurions transforming under  $U(3)_d$  but  $d$  &  $Z^d$

Mixing with  
Spontaneously  
Broken Flavor  
Vacuum



$$\mathcal{L} \supset i Z_{ij}^d \bar{d}_i^\dagger \bar{\sigma}^\mu d_j + i \bar{u}_i^\dagger \bar{\sigma}^\mu u_i + i \bar{Q}_i^\dagger \bar{\sigma}^\mu Q_i + \text{Yukawa terms all flavor diagonal}$$

Canonical kinetic term implies:  $(Z^d)^{-\frac{1}{2}} \sim V y^d$

**The up-quark sector remains aligned.**

**New spurions** allowed:  $\lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$





# : Vector-like quark

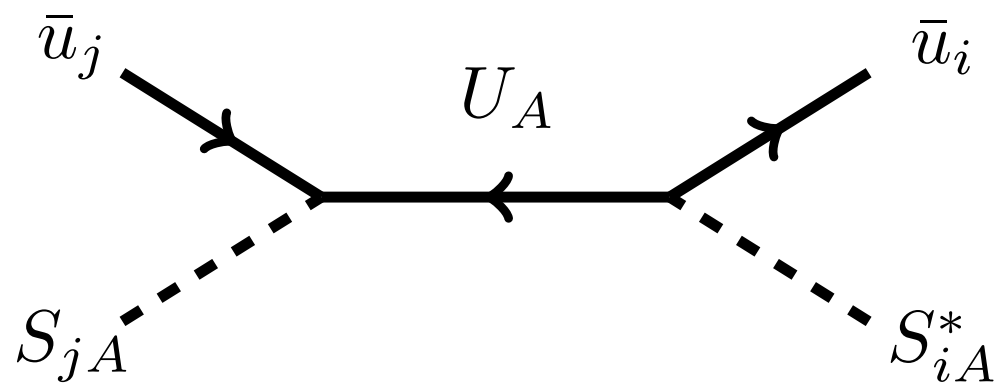
$$\mathcal{L} \supset M_{AB} U_A \bar{U}_B + \xi S_{iA} \bar{u}_i U_A$$

$$- [\eta_{ij}^u Q_i H \bar{u}_j - \eta_{ij}^d Q_i H^c \bar{d}_j + \text{h.c.}] + \mathcal{L}_{\text{BSM}}$$

No additional spurions/fields transforming under  $U(3)_{\bar{u}}$

Introduce mixing between up-quark and heavy VLQs in a flavor breaking vacuum

	$U(3)_U$	$U(3)_{\bar{U}}$	$U(3)_{\bar{u}}$	$U(1)_B$	$\mathbb{Z}_2$
$U$	3			1/3	-1
$\bar{U}$		3		-1/3	-1
$S$	$\bar{3}$		$\bar{3}$		-1



Integrating out heavy quarks leads to wave-function renormalization of the SM up-quarks

$$Z_{ij}^u = \delta_{ij} + \frac{\xi^* \xi}{M_A^* M_A} S_{iA}^* S_{jA}$$

The source of all flavor-breaking!  
CKM matrix arises from returning to canonical basis

S.Homiller@PPC2021



# SFV 2HDM

SFV physics may have generic family non-universal couplings to SM.

A simple framework: **SFV Two Higgs Doublet Models (2HDM).**

$$|D_\mu H_a|^2 - V(H_1, H_2) - \left( \mathcal{Y}_{aij}^u \bar{Q}_{Li} H_a U_{Rj} + \mathcal{Y}_{aij}^d \bar{Q}_{Li} H_a^c D_{Rj} + \mathcal{Y}_{aij}^\ell \bar{L}_{Li} H_a^c \ell_{Rj} + h.c. \right)$$

$$(a = 1, 2; i, j = 1, 2, 3)$$

2HDM features neutral 2 CP-even H:  
 $\cos(\beta - \alpha) \neq 0$  **interesting!**

$$h = H_1^0 \sin(\beta - \alpha) + H_2^0 \cos(\beta - \alpha)$$
$$H = H_1^0 \cos(\beta - \alpha) - H_2^0 \sin(\beta - \alpha)$$

+ a CP-odd and charged ones ... for simplicity:

$$m_H = m_{H^\pm} = m_A$$

## DOWN-TYPE SFV ANSATZ ON 2HDM

$$\mathcal{Y}_1^u = y^u = \text{diag}(y_u, y_c, y_t)$$

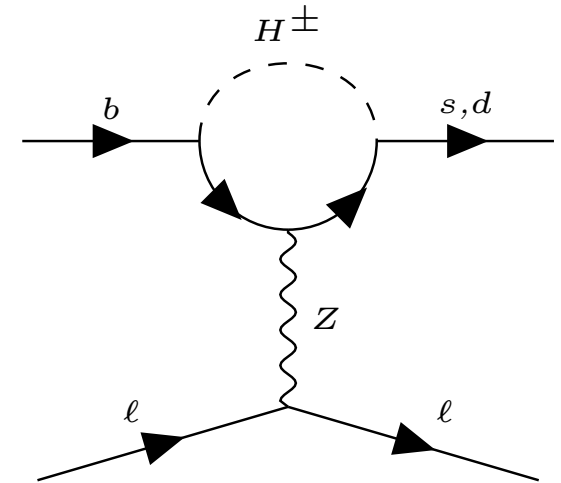
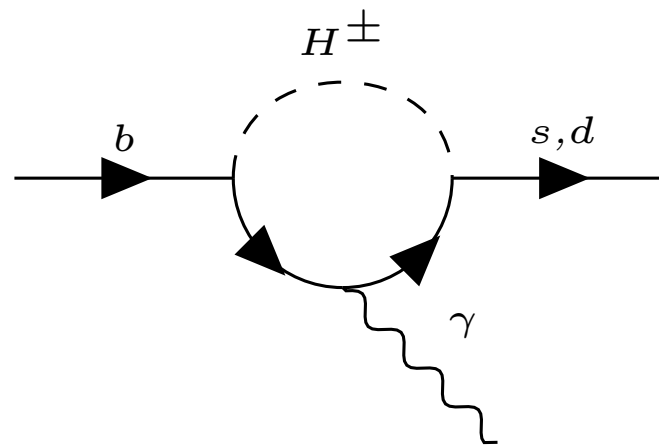
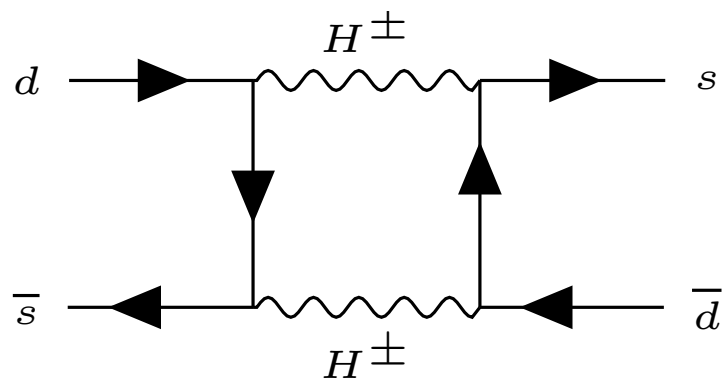
$$\mathcal{Y}_1^d = V y^d$$

$$\mathcal{Y}_1^\ell = y^\ell$$

$$\mathcal{Y}_2^u = \lambda^u = \text{diag}(\lambda_u, \lambda_c, \lambda_t)$$

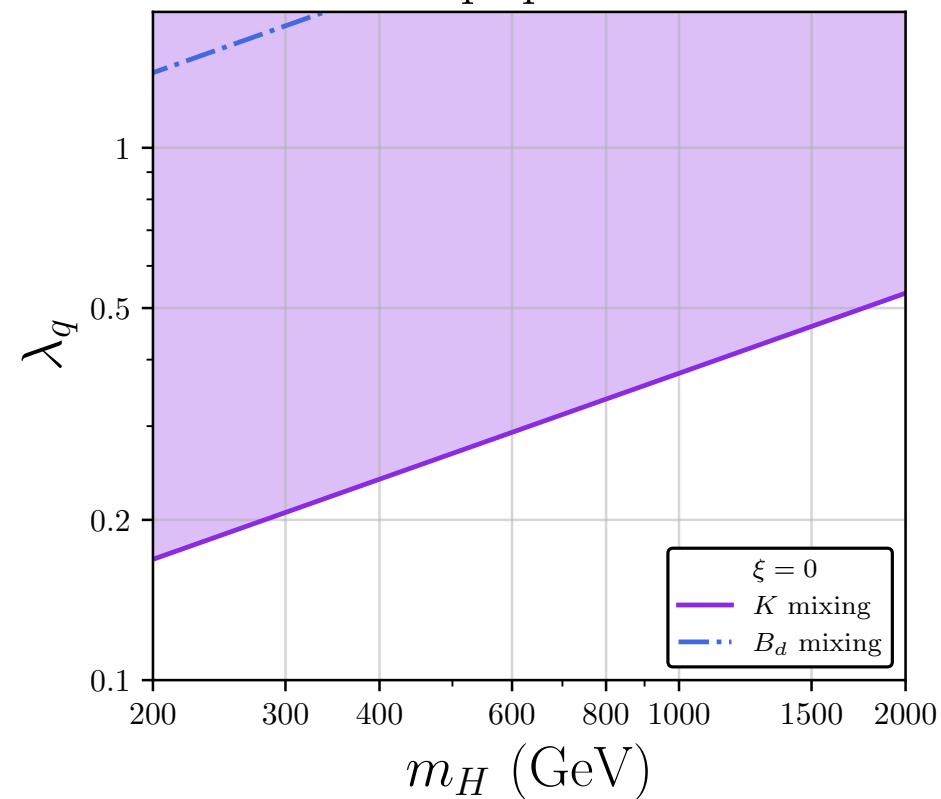
$$\mathcal{Y}_2^d = \xi V y^d$$

$$\mathcal{Y}_2^\ell = \xi^\ell y^\ell$$

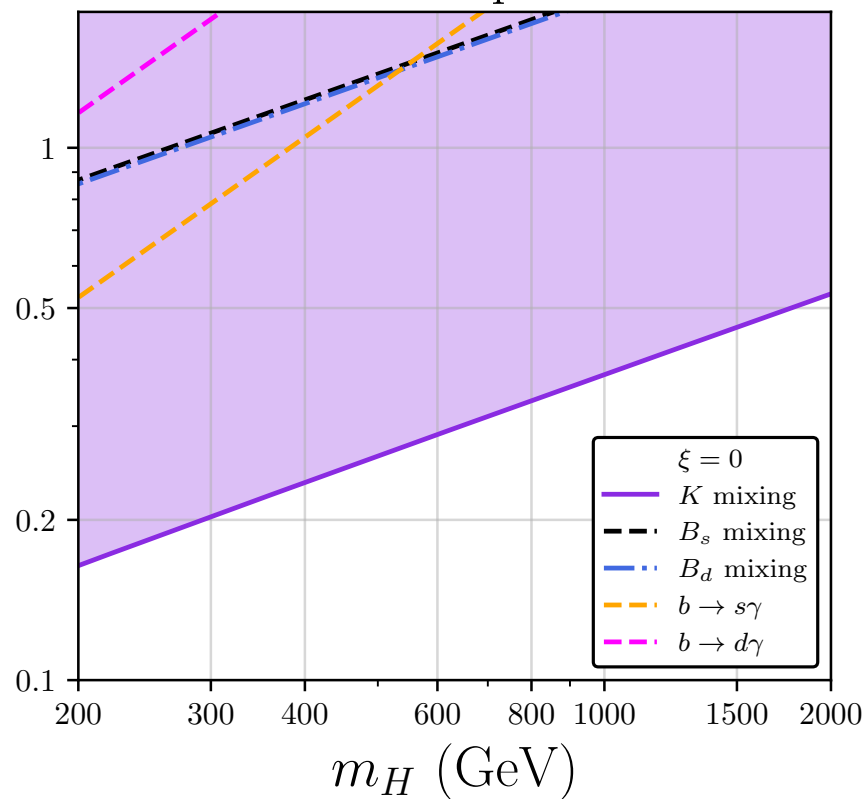


FLAVOR PROTECTION IN DOWN-TYPE SFV ALLOWS NEW HIGGSES  
 @ TEV TO COUPLE TO UP-TYPE QUARKS WITH  $\sim 0.1$  STRENGTH.

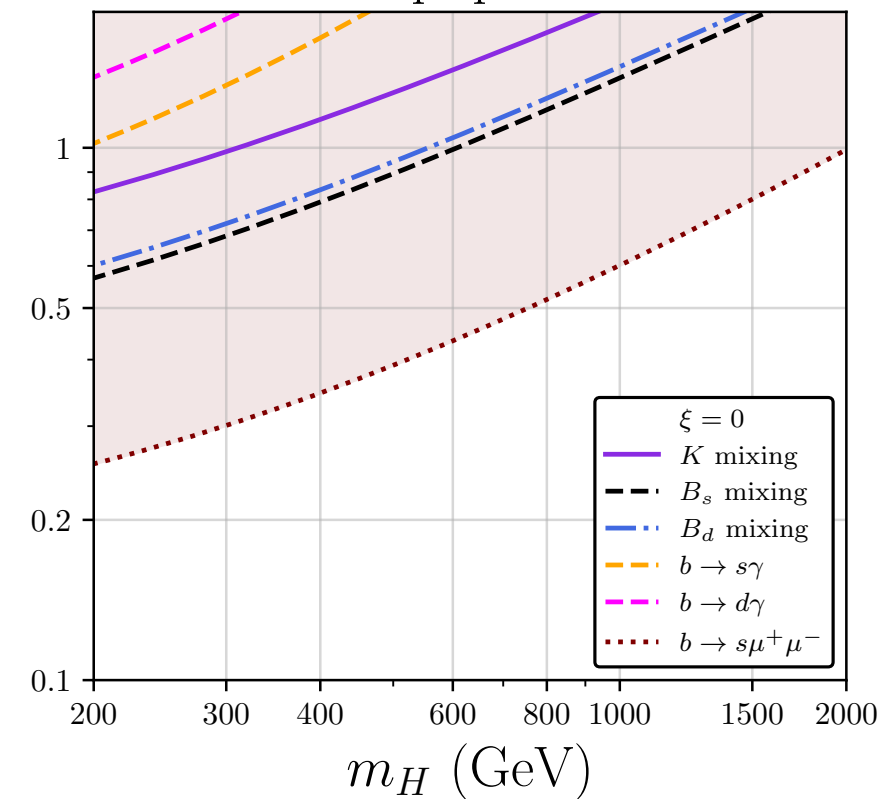
Up quark



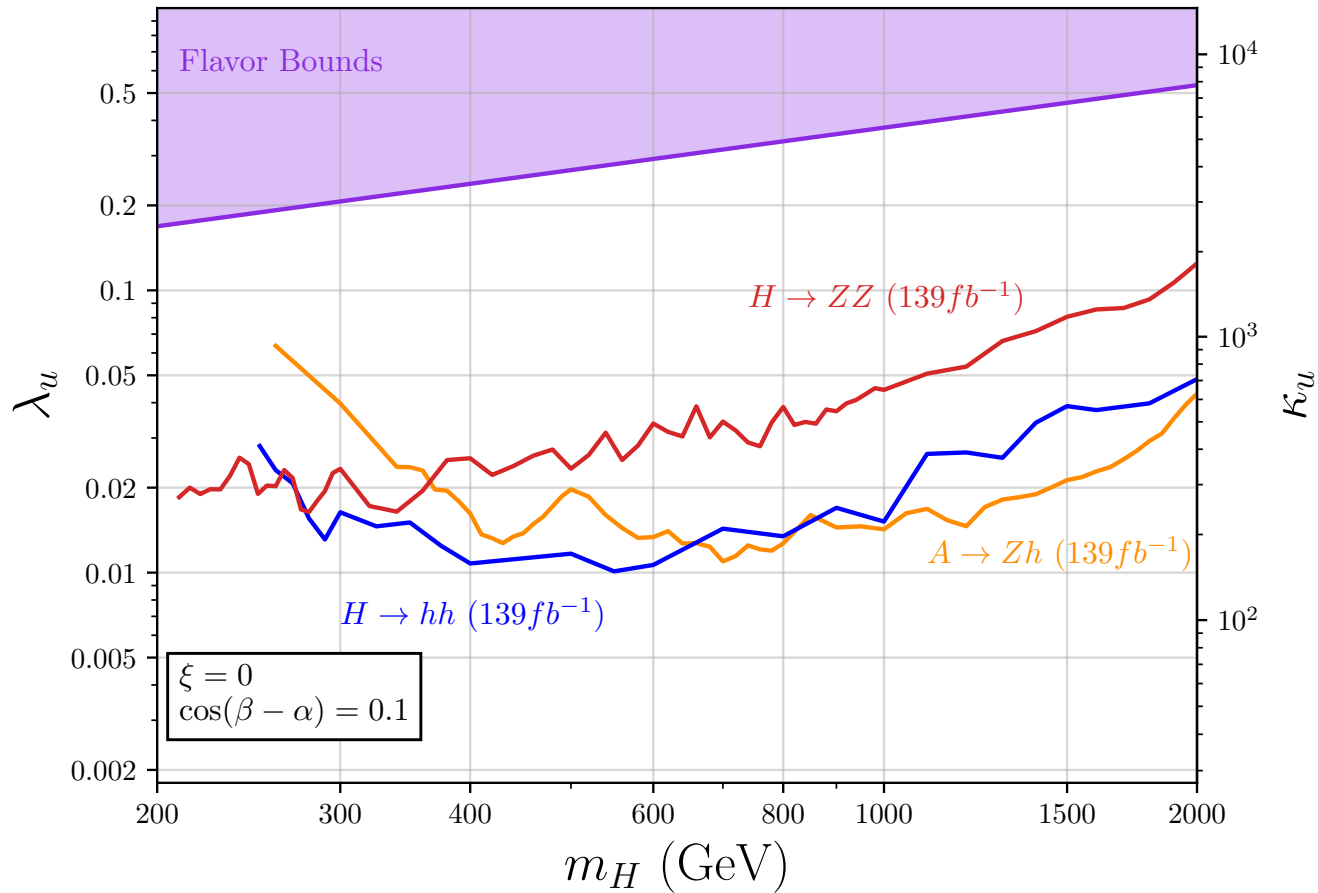
Charm quark



Top quark



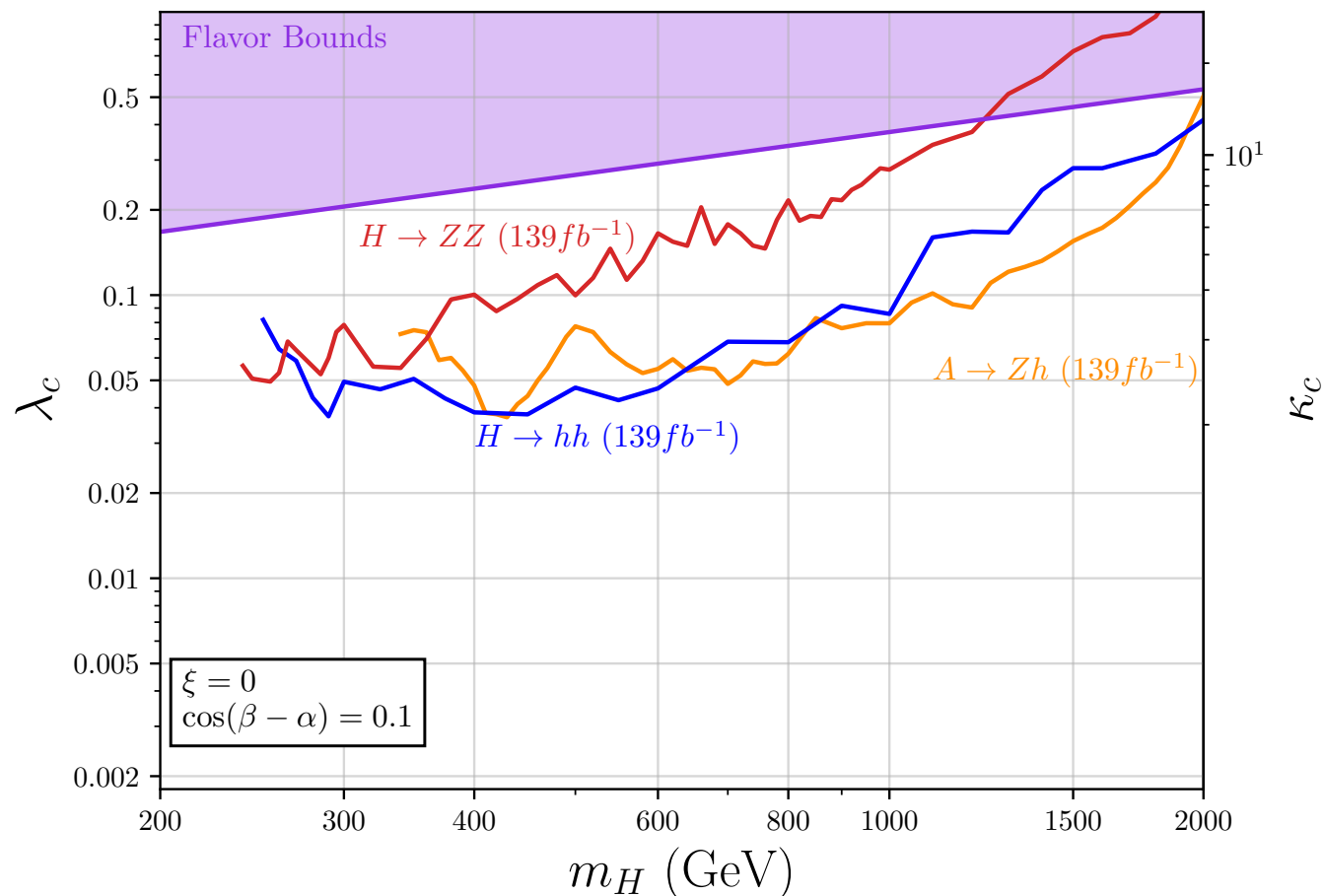
**Same outcome also for the up-type SFV 2HDM.**



Di-Higgs production is a very sensitive probe of enhanced Higgs couplings to light quarks

arXiv:2101.04119

**DOMINANT COLLIDER BOUNDS FROM SEARCH FOR NEW HIGGSES**



- Mainly from quark fusion, due to possibly large Yukawa
- Mostly into di-boson & di-Higgs due to non-zero mixing angle

OBS.

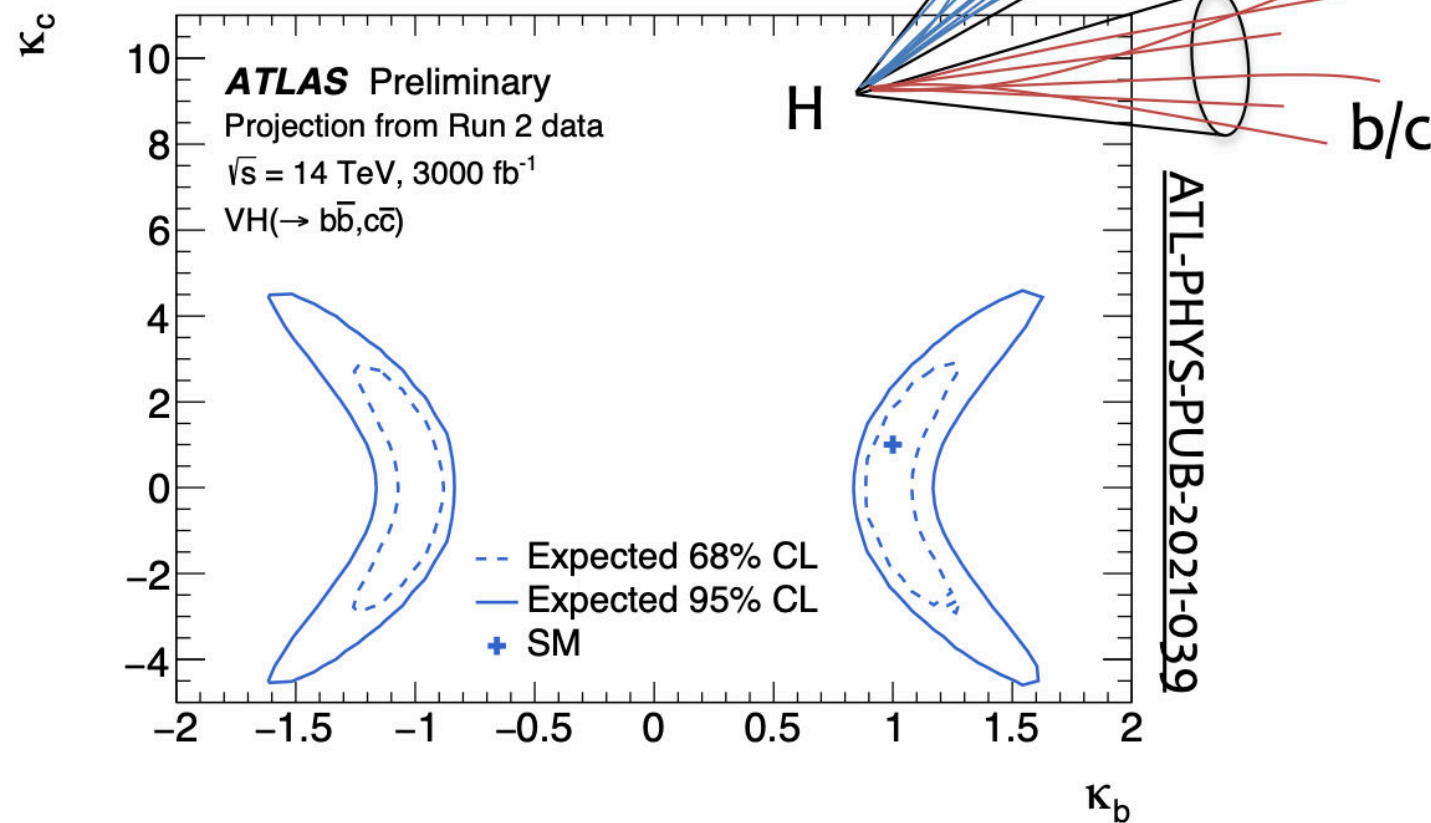
$\kappa_{u,c} \equiv y_{u,c}/y_{u,c}^{\text{SM}}$  is the Yukawa modifier

# LHC probing Charm Yukawa!

Boosted topologies allow to disentangle rare signals from very large background ( $H \rightarrow c\bar{c}$  20x smaller than  $H \rightarrow b\bar{b}$ )

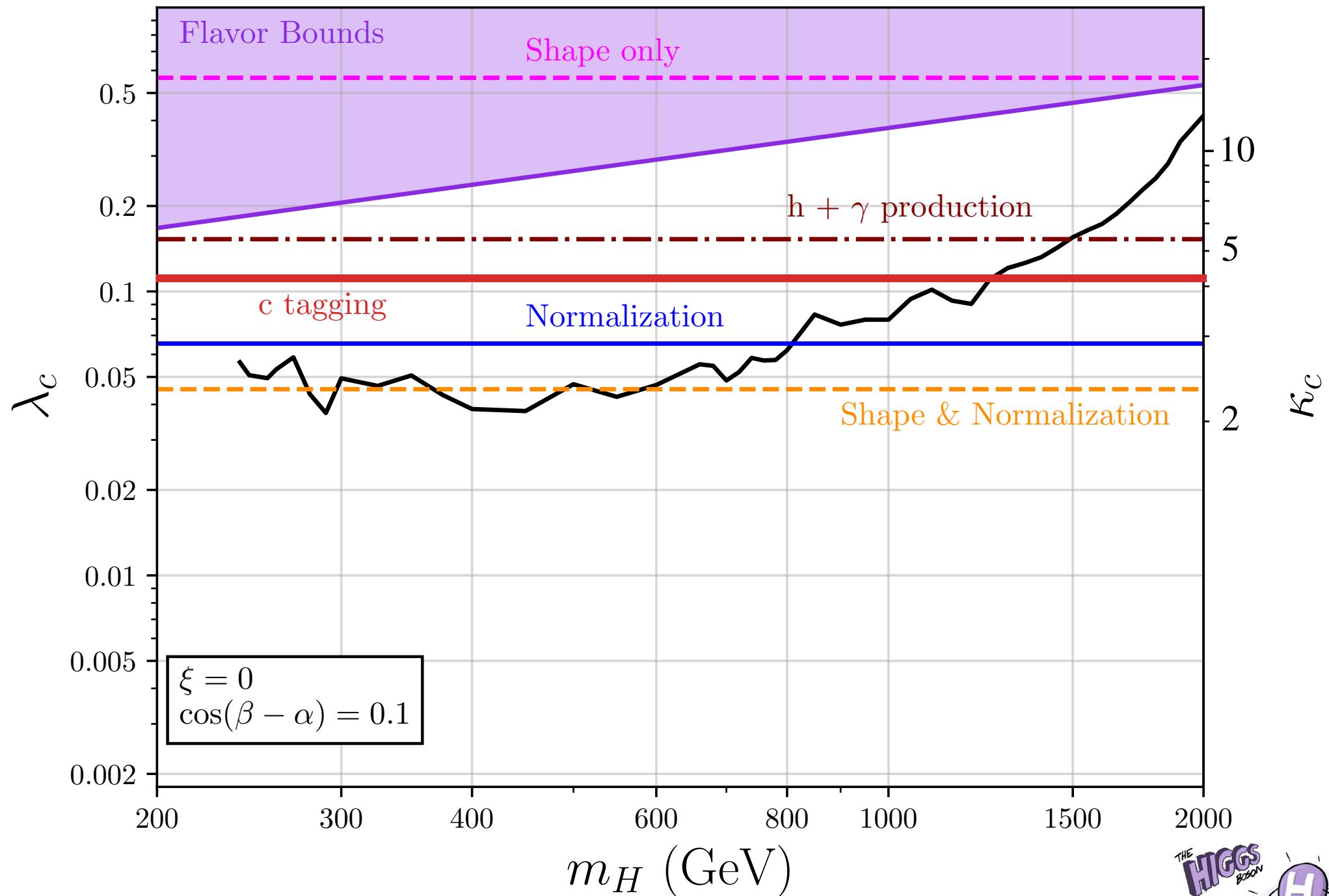
See also search in **CMS-HIG-21-008**

## H-b/c Yukawa



The advent of Deep Learning has been a game changer for ATLAS & CMS sensitivity to Yukawa coupling measurements.

# Charming Higgs @ LHC



arXiv:2410.05236



LHC **c tagging** turns out to be a remarkable probe of SFV scenarios!



- **Program of precision with flavor is still very rich, but naively points to very high scales for NP.**  
No clear tension emerging from flavor data
- **In absence of evidence for NP, the main direction we pursued for BSM may need a paradigm shift.**  
E.g.: **MFV** hypothesis  $\rightarrow$  **SFV** ansatz
- **Theory + EXP progress can lead to new targets.**  
Higgs Yukawa measurements insightful!

[ See 2410.05236 and also 2410.08272 ]