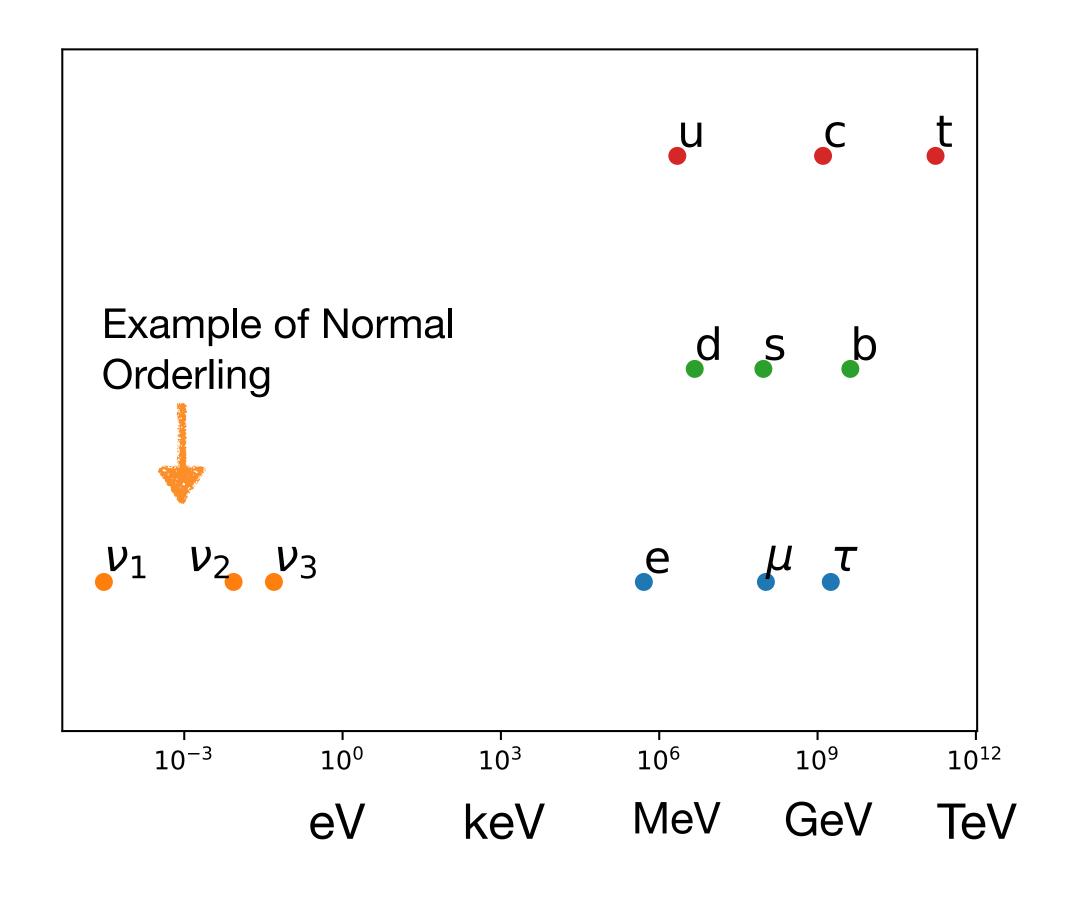
Leptogenesis via scalar doublet decay in the Scotogenic model

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In collaboration with Joe Sato, Kento Asai, Masato Yamanaka

work in progress

Why neutrino mass so small?



neutrinos are Dirac fermions? (SM+right handed neutrino)

Neutrino mass

$$y_{\nu}\bar{L}\tilde{\Phi}\nu_{\rm R} + \text{h.c.} \longrightarrow \mathcal{M}_{\nu} = y_{\nu}\frac{v}{\sqrt{2}}$$

 $\mathcal{O}(0.1)eV$ needs $y_{\nu} \sim 10^{-13}$

Seesaw mechanism? (right-handed neutrinos are Majorana)

$$y_{\nu}\bar{L}\tilde{\Phi}\nu_{R} + \frac{1}{2}M\bar{\nu}_{R}^{c}\nu_{R} + \text{h.c.}$$

$$\longrightarrow \mathcal{M}_{\nu} \simeq \frac{y_{\nu}^{2}v^{2}}{M}$$

Requires very heavy Majorana masses

Neutrinos are extremely light compared to charged leptons and quarks

Some mechanism?

Neutrino mass generated radiatively

Radiative Seesaw Model

- massless at tree level
- generated neutrino mass radiatively at one or more loops

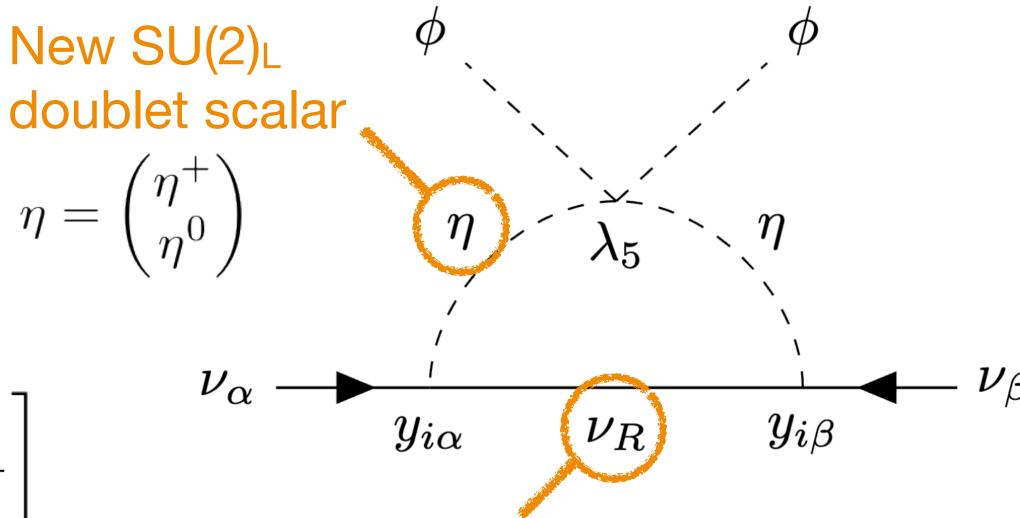
Suppression from loop factors and couplings

Scotogenic Model

♦ Generate neutrino masses via dark sector (η) radiatively

Neutrino mass

$$[\mathcal{M}_{\nu}]_{\alpha\beta} = \frac{\lambda_5 v^2}{16\pi^2} \sum_{i} \frac{y_{i\alpha} y_{i\beta} M_i}{m_{\eta}^2 - M_i^2} \left[1 - \frac{M_i^2}{m_{\eta}^2 - M_i^2} \log \frac{m_{\eta}^2}{M_i^2} \right]$$



Right-handed neutrino (i=1,2,3)

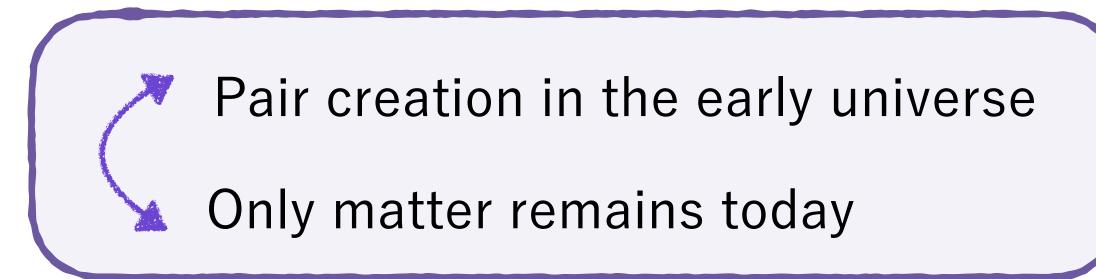
@ ICISE (Aug. 21. 2025)

E. Ma, Phys.Rev.D.73.077301 (2006) Z. Tao, Phys.Rev.D.54.5693 (1996)

Neutrinos as window into BSM physics

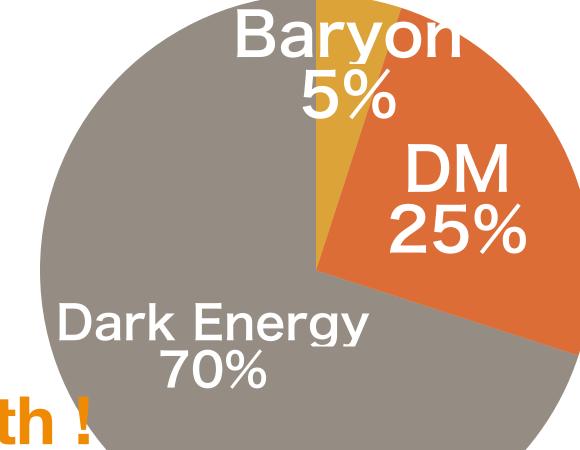
Cosmological issues suggesting BSM

Matter-antimatter asymmetry



why did particles remain, while antiparticles disappeared?

 Dark Matter Massive, long-lived, weakly interacting particle



Scotogenic model has the potential to explain them both !

Scenario

♦ DM candidate → N₁ or η

<u>n as DM candidate</u> $m_{\eta} < M_1 < M_2 < M_3$

- Neutrino masses
- ✓ DM abundance
- ✓ Baryon asymmetry

 N_1 as DM candidate

? Neutrino masses

? DM abundance

? Baryon asymmetry

$$M_1 < m_{\eta} < M_2 < M_3$$

$$M_1 < M_2 < m_{\eta} < M_3$$

$$M_1 < M_2 < M_3 < m_{\eta}$$

Scenario

♦ DM candidate → N₁ or η

η as DM candidate $m_{\eta} < M_1 < M_2 < M_3$

- Neutrino masses
- ✓ DM abundance
- ✓ Baryon asymmetry

N_1 as DM candidate

- ? Neutrino masses
- ? DM abundance
- ? Baryon asymmetry

$$M_1 < m_{\eta} < M_2 < M_3$$

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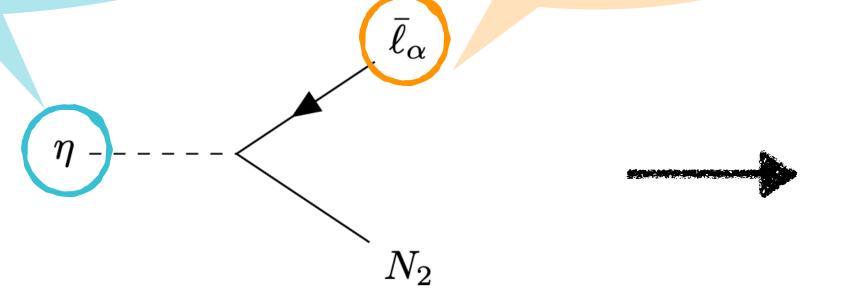
$$M_1 < m_{\eta} < M_2 < M_3$$

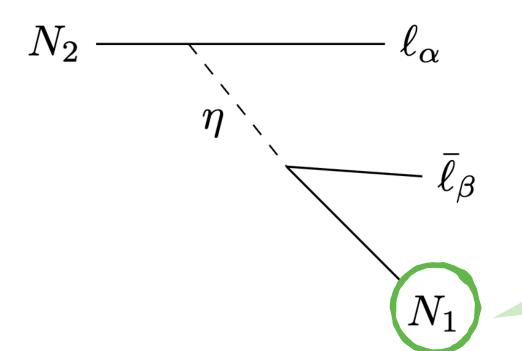
$$M_1 < M_2 < m_{\eta} < M_3$$

$$M_1 < M_2 < M_3 < m_{\eta}$$

η is thermaly produced in the early universe

Generation of lepton asym.





Production of Dark Matter

Baryon asymmetry from η decay

Baryon-asymmetry : the differences of the number of baryon and anti-baryon $B=n_b-n_{ar{b}}$

$$B = n_b - n_{\bar{b}}$$

Leptogenesis

Lepton asym. Baryon asym. sphaleron process

CP is conserved→no asymmetry

$$X \to A + B$$

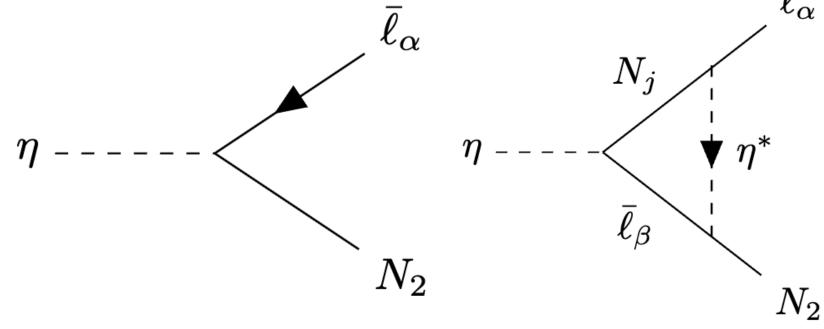
 $\bar{X} \to \bar{A} + \bar{B}$

- Need CP violation
- CP asym. from interference of tree and 1-loop

CP asym. in this scenario

$$\epsilon_{i\alpha} = \frac{\Gamma(\eta \to N_i \bar{\ell}_{\alpha}) - \Gamma(\bar{\eta} \to N_i \ell_{\alpha})}{\Gamma(\eta \to N_i \bar{\ell}_{\alpha}) + \Gamma(\bar{\eta} \to N_i \ell_{\alpha})}$$

$$= \frac{1}{8\pi} \sum_{\beta,j} \frac{Im[Y_{i\alpha}^* Y_{i\beta}^* Y_{j\alpha} Y_{j\beta}]}{|Y_{\alpha i}|^2} \frac{M_i M_j}{m_{\eta}^2 - M_i^2} \left\{ 1 - \left(1 + \frac{m_{\eta}^2 - M_j^2}{m_{\eta}^2 - M_i^2}\right) \log\left(\frac{m_{\eta}^2 - M_i^2}{m_{\eta}^2 - M_j^2} + 1\right) \right\}$$



A parameter for Neutrino mass

Neutrino masses from radiative seesaw

$$[\mathcal{M}_{\nu}]_{\alpha\beta} = \frac{\lambda_5 v^2}{16\pi^2} \sum_{i} \frac{y_{i\alpha} y_{i\beta} M_i}{m_{\eta}^2 - M_i^2} \left[1 - \frac{M_i^2}{m_{\eta}^2 - M_i^2} \log \frac{m_{\eta}^2}{M_i^2} \right]$$



$$M_1 = 1 \,\mathrm{MeV}$$

$$\lambda_5 = 1$$

$$M_2 = 9.5 \times 10^9 \,\mathrm{GeV}$$

 $M_3 = 9.8 \times 10^9 \, \text{GeV}$

$$u_{i\alpha} = \begin{cases} 6.03 \times 10^{-7} + 1.06 \\ -1.17 \times 10^{-2} + 5.95 \end{cases}$$

$$-1.71 \times 10^{-6} + 2.29 \times 10^{-6}$$
$$-2.34 \times 10^{-2} - 2.21 \times 10^{-2}$$

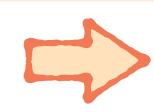
 $m_3 = 5.03 \times 10^{-2} eV$

$$-1.96 \times 10^{-6} - 5.23 \times 10^{-7} i$$
$$6.87 \times 10^{-3} - 2.22 \times 10^{-2} i$$

$$m_{\eta} = 10^{10} \, \mathrm{GeV}$$

$$y_{i\alpha} = \begin{cases} -1.17 \times 10^{-2} + 5.93 \times 10^{-3}i \\ -1.33 \times 10^{-2} - 5.85 \times 10^{-5}i \end{cases}$$

$$y_{i\alpha} = \begin{pmatrix} 6.03 \times 10^{-7} + 1.06 \times 10^{-6}i & -1.71 \times 10^{-6} + 2.29 \times 10^{-6}i & -1.96 \times 10^{-6} - 5.23 \times 10^{-7}i \\ -1.17 \times 10^{-2} + 5.93 \times 10^{-3}i & -2.34 \times 10^{-2} - 2.21 \times 10^{-2}i & 6.87 \times 10^{-3} - 2.22 \times 10^{-2}i \\ -1.33 \times 10^{-2} - 5.85 \times 10^{-5}i & 4.14 \times 10^{-2} - 1.23 \times 10^{-2}i & 4.17 \times 10^{-2} + 3.60 \times 10^{-3}i \end{pmatrix}$$



neutrino masses and mixing can be reproduced!

$$\sin \theta_{12}^2 = 0.307$$

$$m_1 = 1.00 \times 10^{-31} eV$$

$$\sin \theta_{23}^2 = 0.561$$

$$m_2 = 8.65 \times 10^{-3} eV$$

$$\sin \theta_{13}^2 = 0.02195$$

$$\delta_{CP}/^{\circ} = 177$$

(From NuFit6.0)

Calculation of Lepton asym.

L-asym. is calculated by solving Boltzmann equation

equation of time evolution of number density

$$\frac{dn_L}{dt} + 3Hn_L = -\left\{n_{\Delta\eta} + n_{\eta}^{eq} \frac{n_L}{n_{\ell}^{eq}} \frac{n_N}{n_N^{eq}}\right\} \langle \Gamma_{\eta} \rangle - \left\{n_+ - n_+^{eq} \frac{n_N}{n_N^{eq}}\right\} \langle \Gamma_{\eta} \rangle + \left\{n_{\Delta\eta} n_+ - n_{\eta}^{eq} n_+^{eq} \frac{n_L}{n_{\ell}^{eq}}\right\} C_{spec} \langle \sigma v \rangle_{\eta\eta \to \phi\phi}$$

$$n_{\Delta\eta} = n_{\bar{\eta}} - n_{\eta}$$
$$n_{+} = n_{\bar{\eta}} + n_{\eta}$$

 n^{eq} : number density in thermal equilibrium

 C_{spec} : Coefficient converting Higgs asym. to lepton asym.

solve a set of three coupled differential equations numerically

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$$+ \left\{n_{\Delta\eta} n_+ - n_{\eta}^{eq} n_+^{eq} \frac{n_L}{n_{\ell}^{eq}}\right\} C_{spec} \langle \sigma v \rangle_{\eta\eta \to \phi\phi}$$

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$$\frac{dn_{\Delta\eta}}{dt} + 3Hn_{\Delta\eta} = \left\{ n_{\Delta\eta} + n_{\eta}^{eq} \frac{n_L}{n_{\ell}^{eq}} \frac{n_N}{n_N^{eq}} \right\} \langle \Gamma_{\eta} \rangle + \left\{ n_+ - n_+^{eq} \frac{n_N}{n_N^{eq}} \right\} \epsilon \langle \Gamma_{\eta} \rangle
+ \left\{ n_{\Delta\eta} n_+ - n_{\eta}^{eq} n_+^{eq} \frac{n_L}{n_{\ell}^{eq}} \right\} C_{spec} \langle \sigma v \rangle_{\eta\eta \to \phi\phi}$$

$$\frac{dn_{+}}{dt} + 3Hn_{+} = -\left\{n_{+} + n_{+}^{eq} \frac{n_{N}}{n_{N}^{eq}}\right\} \langle \Gamma_{\eta} \rangle - \frac{1}{2} \left\{n_{+}^{2} - n_{+}^{eq2}\right\} \epsilon \langle \Gamma_{\eta} \rangle_{\eta\eta \to \phi\phi, \eta\eta \to AA}$$

$$\frac{dn_{N_2}}{dt} + 3Hn_{N_2} = \langle \Gamma_{\eta} \rangle \left\{ n_{\eta} - \frac{n_{N_2}}{n_{N_2}^{eq}} \frac{n_{\ell}}{n_{\ell}^{eq}} n_{\eta}^{eq} \right\}$$

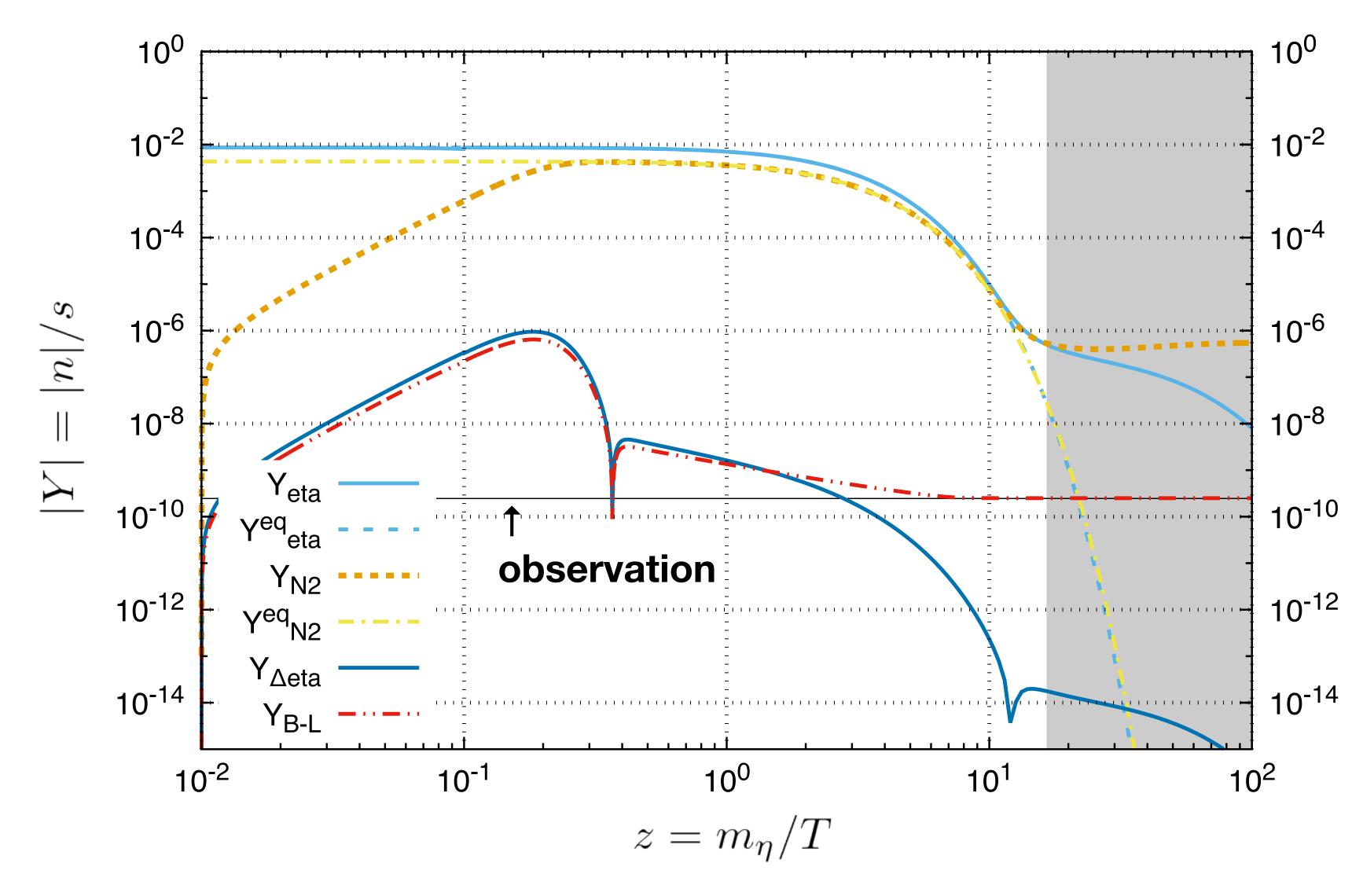
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L asym. time evolution



- lacktriangle Unified framework for $m_{
 u}$, mixing, DM, Baryon Asym.
- The favored parameter space is completely different from the scenario with η as the DM candidate.

Summary

- → The radiative seesaw model is an attractive scenario that explains the smallness of neutrino masses through one or more loops.
- Scotogenic model generate neutrino mass via dark sector.
- In this model, neutrino masses, the baryon asymmetry, and DM abundance can all be explained in unified framework when the lightest right-handed neutrino is the DM candidate.
- The favored parameter space is completely different from the scenario with η as the DM candidate.