



# Probing a class of scotogenic models via Z and Higgs boson decays

Based in work to appear soon

[Darricau, Lee, Orloff, Teixeira, 2505.xxxx]



[A. Alvarez & al., 2301.08485]

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# **Standard Model: an incomplete picture**

### Despite its many successes, experimental evidence of the **incompleteness** of the SM:

- $\rightarrow$  Neutrino oscillations
- $\rightarrow$  Dark Matter (DM)
- $\rightarrow$  Baryon Asymmetry of the Universe (BAU)









### Also some theoretical issues...



# **Standard Model: an incomplete picture**

### Also some **theoretical issues** such as:

- $\rightarrow$  Fundamental scales gap
- $\rightarrow$  Hierarchy problem
- $\rightarrow$  Flavour puzzle

If NP at Planck scale, 30 order cancellation with bare mass!





# $\Delta M_H^2 |_{1-loop}^{SM} = \frac{3\Lambda_{UV}^2}{8\pi^2 v^2} [M_H^2 + 2M_W^2 + M_Z^2 - 4m_t^2 + ...]$

 $\rightarrow$  How to overcome these experimental caveats?  $\rightarrow$  Can we also **improve these theoretical issues**?





# **Scotogenic models: motivation**

Extension at the **Tera** eV **Scale** 

A lot of mechanisms to explain  $\nu$  mass generation!

### $\rightarrow$ Seesaw realisation, radiative mechanisms...

Can both be related?

A plethora of models putting forward DM candidates!

 $\rightarrow$  WIMPs, FIMPs...





### Yes ! A "Dark generation" of $\nu$ mass: "Scotogenic"









### A simple radiative seesaw $\nu$ mass generation !

 $\rightarrow$  A global  $\mathbb{Z}_2$  symmetry to stabilise DM candidates  $\rightarrow$  Neutrino mass generation mechanism constrained by the global  $\mathbb{Z}_2$  symmetry, exclusion of tree level neutrino masses









Many interesting realisations of the scotogenic paradigm T1-1-A, **T1-2-A**...

[D. Restrepo, 1308.3655]





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# **Description of the T1-2-A setup**

 $\rightarrow$  T1-2-A setup: adds 2  $SU(2)_L$  doublets and (1 fermionic, 1 scalar) charged under  $\mathbb{Z}_2$ 

 $\rightarrow$  Extended: addition of fermion singlet and c

$$\begin{split} \mathcal{V}_{\text{scalar}} &= \frac{1}{2} M_S^2 \, \mathbf{S}^2 + \frac{1}{2} \lambda_{4S} \, \mathbf{S}^4 + M_\eta^2 \, |\boldsymbol{\eta}|^2 + \lambda_{4\eta} \, |\boldsymbol{\eta}|^4 + \frac{1}{2} \lambda_S \, \mathbf{S} \\ &+ \lambda_\eta \, |\boldsymbol{\eta}|^2 \, |\Phi|^2 + \lambda_\eta' \, |\boldsymbol{\eta} \Phi^\dagger|^2 + \frac{1}{2} \lambda_\eta'' \left[ \left( \Phi \boldsymbol{\eta}^\dagger \right)^2 + \text{H.c.} \right] + \lambda_\eta'' \, \mathbf{M} \, \mathbf{M}$$



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d 2 singlets	Field	η	$\boldsymbol{S}$	$F_1$	$F_2$	$\Psi_1$	
	$\mathrm{SU}(2)_L$	2	1	1	1	2	
doublet	$\mathrm{U}(1)_Y$	1	0	0	0	-1	



### Which states after EWSB?









# **Description of the T1-2-A setup**

- $\rightarrow$  T1-2-A setup: adds 2  $SU(2)_L$  doublets and (1 fermionic, 1 scalar) charged under  $\mathbb{Z}_2$
- $\rightarrow$  Extended: addition of fermion singlet and c
- After EWSB:
  - $\rightarrow$  Scalars: 2 CP even ( $\phi_{1/2}$ ), 1 CP odd ( $A^0$ ) neutral states and charged  $\eta^{\pm}$  state  $\rightarrow$  Fermions: 4 Majorana states ( $\chi_i$ ) and 1 charged vector-like  $\psi^{\pm}$

  - In short:  $\rightarrow$  3 potential **DM candidate** :  $\phi_1$ ,  $\chi_1$  and  $A^0$  $\rightarrow$  No mixing between CP even and odd state in scalar sector





d 2 singlets	Field	$\eta$	S	$F_1$	$F_2$	$\Psi_1$	
	$\mathrm{SU}(2)_L$	2	1	1	1	2	
doublet	$\mathrm{U}(1)_Y$	1	0	0	0	-1	

 $\rightarrow$  4 Majorana states ( $\chi_i$ ) needed for 3 massive neutrinos (and leptogenesis)







# **Description of the T1-2-A setup**

### $\checkmark \nu$ mass mechanism, lepton decays and muon magnetic dipole moment ( $\Delta a_{\mu} \sim 4\sigma$ ) ✓ Found **2 viable DM** candidates ( $\phi_1, \chi_1$ ) Extensive study of BAU via leptogenesis

### Viable DM candidates ( $\phi_1, \chi_1$ )



### [A. Alvarez & al., 2301.08485]



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# The scotogenic T1-2-A setup: revisited

- A flavour physics approach  $\rightarrow$  keep points with unviable DM candidate **Relax** certain (driving) assumptions  $\rightarrow$  generic  $\Delta a_{\mu}$  - from SM-like to NP (at ~ 4.2 $\sigma$ ); no BAU
- **Thorough** exploration of **flavoured** and **ElectroWeak Precision Observables** (EWPO)!
  - $\rightarrow$  cLFV decays: leptonic (radiative, 3 body, conversion in Nuclei),  $Z/H \rightarrow \ell_{\alpha} \ell_{\beta}$
  - $\rightarrow$  **EWPO**: sensitive probes of new interactions (scalar, vector, fermion...) Oblique parameters (S,T,U),  $Z/H \rightarrow inv$ ,  $Z/H \rightarrow \ell_{\alpha}\ell_{\alpha}$
  - $\rightarrow$  Lepton Flavour Universality Violation (LFUV)
  - $\rightarrow$  In short: revisit a well studied model from a flavour perspective with updated  $\Delta a_{\mu}$





Our goals ?



### $\nu$ mass generation (at 1 loop)







 $\nu$  mass generation

- $\rightarrow$  A "seesaw" mechanism but at 1 loop !
- $\rightarrow$  Flavour violating coupling  $\mathscr{G}$  giving neutrino masses

$$\mathscr{L}_{\text{fermion}} \supset -g_{\psi}^{\alpha} \widetilde{\overline{\psi_{2}}} L_{L}^{\alpha} S - g_{F_{i}}^{\alpha} \widetilde{\overline{L_{L}^{\alpha}}} \eta F_{i}$$

$$M_{\nu} = \mathscr{G}^T M_L \mathscr{G}$$

**Naturally suppressed** by a loop prefactor !









# Viable Dark Matter candidate



 $\rightarrow$  Most points excluded by DM-related constraints...







 $\rightarrow$  MicrOmegas to generate relic density and direct detection observables  $\rightarrow$  CP even  $\phi_1$ , Majorana fermion  $\chi_1$  and **new viable candidate** CP odd  $A^0$ !



# Muon dipole moment



[A. Boccaletti & al. 2407.10913]

 $\rightarrow$  If large  $\Delta a_{\mu}$  (4.2 $\sigma$ )  $\rightarrow$  boosts dipole operator  $\rightarrow$  If small  $\Delta a_{\mu}$  (SM-like)  $\rightarrow$  constraints on NP What is the impact of a more suppressed  $\Delta a_{\mu}$ ?



Lattice QCD results appeases tension with experiments





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### **cLFV** lepton processes









### $\mu \rightarrow e\gamma, \mu \rightarrow 3e$ ; effect of NP vs SM-like $\Delta a_{\mu}$





# **cLFV** lepton processes

### Neutrino less Mu-e conversion in the Nuclei









## **cLFV** lepton processes

### Leptonic tau decays











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### Z cLFV decays







# $\rightarrow$ **Representative** of the other Z cLFV decays $\rightarrow$ Driven by the NP coupling to right-handed leptons $g_R^{\alpha}$

$$\begin{aligned} \mathcal{V}_{\text{fermion}} &= M_{\nu}^{\alpha\beta} \, \overline{\nu_{\alpha}^{c}} \, \nu_{\beta} - M_{\psi} \, \overline{\psi_{1}} \, \overline{\psi_{2}} + \frac{1}{2} \, M_{F_{ii}} \, \overline{F_{i}^{c}} \, F_{i} - y_{1i}^{*} \, \overline{F_{i}} \, \Phi^{\dagger} \, \overline{\psi_{1}} - y_{2i}^{*} \\ &+ g_{\psi}^{\alpha} \, \overline{\widetilde{\psi_{2}}} \, L_{L}^{\alpha} S + g_{F_{i}}^{\alpha} \, \overline{\widetilde{L_{L}^{\alpha}}} \, \eta F_{i} + g_{R}^{\alpha} \overline{e_{R}^{\alpha}} \, \eta^{\dagger} \, \psi_{1} \right) + \text{H.c.} \end{aligned}$$

Beyond future sensitivity ( $\tau \rightarrow 3\mu$ )

Not as important as expected...









# Higgs cLFV decays





- → **Representative** of the other H cLFV
- $\rightarrow$  Driven by the NP trilinear coupling  $\alpha$ when large

$$\begin{split} \mathcal{V}_{\text{scalar}} &= \frac{1}{2} M_S^2 \, \mathbf{S}^2 + \frac{1}{2} \lambda_{4S} \, \mathbf{S}^4 + M_\eta^2 \, |\boldsymbol{\eta}|^2 + \lambda_{4\eta} \, |\boldsymbol{\eta}|^4 + \frac{1}{2} \lambda_S \, \mathbf{S}^2 |\Phi|^2 + \frac{1}{2} \lambda_S \, \mathbf{S}$$

Again, beyond future sensitivity...

### Maybe EWPO ? LFUV ?









# $\rightarrow$ **EWPO**: sensitive probes of new interactions (scalar, vector, fermion...)

### Why important? $\rightarrow$ Towards high precision tests of the SM (FCC-ee...)

- $\rightarrow Z/H \rightarrow inv$
- $\rightarrow$  Oblique parameters
- $\rightarrow Z/H \rightarrow \ell_{\alpha}\ell_{\alpha} + LFUV$





- Oblique parameters (S,T,U),  $Z/H \rightarrow inv$ ,  $Z/H \rightarrow \ell_{\alpha}\ell_{\alpha}$  and their sensitive ratios



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# ElectroWeak Precisient in the set of the se $= \sum_{i_{1}, i_{2}, i_{3}} = \sum_{i_{1}, i_{2}, i_{3}} \sum_{i_{1}, i_{2}, i_{3}} \sum_{i_{2}, i_$ $E_{TX^{(4)}} = \sum_{i=1}^{N} \frac{32^{i}}{32^{i}} M_{i}^{(2)} M_{i}^{$ $\left(\left(\left(\frac{1}{2}\right)^{p\times W}\right)^{p\times W}$

 $\sum_{k=1}^{N} \sum_{k=1}^{N} \sum_{$ 

 $\begin{pmatrix} (W_{2}^{2}, W_{1}^{2}, W_{2}^{2}, W_{3}^{2}, W_{3}$ 

 $+ \operatorname{Le}^{(k)} \left( \operatorname{Le}^{(k)} \right)^{*}$ 

(1 Ety.

) # 1000 (( +

 $\begin{array}{c} + \sum_{k} \left( \left( 1, \frac{1}{k} \right)^{k} \right)^{k} \left( \left( 1, \frac{1}{k} \right)^{k} \right)^{k} \left( 1, \frac{1}{k} \right$ 

 $\sum_{k=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{k=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{i=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{i$ 

 $\frac{zW + \left( \stackrel{p_{\chi}}{zW}, \stackrel{q_{\phi}}{zW}, \stackrel{i_{\eta}}{zW}, \stackrel{i_{\eta}}{zW} \right)^{-12} \left( \stackrel{M_{Z}}{zW}, \stackrel{q_{\phi}}{zW}, \stackrel{i_{\eta}}{zW}, \stackrel{i_{\eta}}{zW} \right)^{-114} \left( \stackrel{M_{Z}}{zW}, \stackrel{m_{\phi}}{zW}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}} \right)^{-12} \left( \stackrel{M_{Z}}{zW}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{M_{Z}}{M_{\phi}} \right)^{-12} \left( \stackrel{M_{Z}}{zW}, \stackrel{M_{Z}}{M_{\phi}}, \stackrel{$ 

 $+ \sum_{L}^{12} \left( \frac{M_{Z_{3}}}{M_{Z_{3}}} \right)^{n_{L}} \frac{M_{Z_{3}}}{M_{U}} C_{1} \left( \frac{m_{Z_{3}}}{m_{Z_{3}}} \right)^{n_{L}} \frac{M_{U}}{M_{U}} \left( \frac{m_{Z_{3}}}{m_{Z_{3}}} \right)^{n_{L}} \frac{M_{U}}{M_{U}} \frac{M_{U}}{M_{U}} \frac{M_{U}}{M_{U}} \frac{M_{U}}{M_{U}} \right)$ 

+ T<sup>Re</sup> (T<sup>Re</sup>)\* <sup>nne</sup><sup>g</sup> C<sup>12</sup> (M<sup>2</sup>; <sup>nne</sup><sup>g</sup>, <sup>nne</sup><sup>g</sup>, <sup>NP</sup>, <sup></sup>

 $F_{L}^{1S,00} = \frac{F_{U}}{16} \begin{pmatrix} T_{T}^{0} \\ T_{L}^{0} \\ T_{L}^{0}$ 

(THE CAN IN THE CONTRACT OF THE STATE OF THE

 $F_{SL}^{2S,\alpha\beta} = -\frac{eU_{\phi}^{2k}}{16\pi^2 c_w s_w} \Gamma_L^{\beta3} \left(\Gamma_L^{\alpha k}\right)^* \left(C_{00} \left(M_Z^2, m_{\ell_{\alpha}}^2, m_{\ell_{\beta}}^2, M_{\phi_k}^2, M_{A^0}^2, M_{\psi}^2\right)$ 

 $F_{SL}^{1\chi,\alpha\beta} = \frac{ie\left(c_{w}^{2} - s_{w}^{2}\right)}{16\pi^{2}c_{w}s_{w}} \Gamma_{L}^{\beta c}\left(\Gamma_{L}^{\alpha c}\right)^{*}C_{00}\left(M_{Z}^{2}, m_{\ell_{\alpha}}^{2}, m_{\ell_{\beta}}^{2}, M_{\eta^{\pm}}^{2}, M_{\eta^{\pm}}^{2}, M_{\chi_{c}}^{2}\right)$ 

 $-\Gamma_{L}^{\beta k} \left(\Gamma_{R}^{\alpha k}\right)^{*} M_{\psi} m_{\ell_{\alpha}} C_{0} \left(M_{Z}^{2}, m_{\ell_{\alpha}}^{2}, m_{\ell_{\beta}}^{2}, M_{\psi}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2}\right)$ 

 $-\Gamma_{R}^{\beta k} \left(\Gamma_{L}^{\alpha k}\right)^{*} M_{\psi} m_{\ell_{\beta}} C_{0} \left(M_{Z}^{2}, m_{\ell_{\alpha}}^{2}, m_{\ell_{\beta}}^{2}, M_{\psi}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2}\right)$ 

 $+ \Gamma_R^{\beta k} \left( \Gamma_R^{\alpha k} \right)^* m_{\ell_\alpha} m_{\ell_\beta} C_1 \left( m_{\ell_\alpha}^2, m_{\ell_\beta}^2, M_Z^2, M_{\psi}^2, M_{\phi_k}^2, M_{\psi}^2 \right) \right)$ 

 $+ \sum_{R} \sum_{k} \left( \sum_{i=1}^{n} \sum_{i=1}^{n}$ 

 $E^{\frac{1}{2} \cdot \alpha_{0}g} = \frac{ie(2^{2} - \frac{3}{2^{2}})}{16\pi^{2}c_{0}s_{0}} (\Gamma_{E}^{3}(\Gamma_{0}^{2})^{*}m_{e}^{*}C_{12}(M_{E}^{2},m_{e}^{2})^{*}m_{e}^{*}C_{12}(M_{E}^{2},m_{e}^{2},M_{e}^{2},M_{e}^{2},M_{e}^{2})}{2^{*}-2} N_{F2} N_{F2}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*}N_{e}^{*})$ 

\* 50 (Cb + 4) + C1 (M20, M2, M2, M2) (1 (M, (10, 1)) + C1 (M20, M2, M2, M2) (1 (M, (10, 1)) - M2, M2) (1 (M, (10, 1)) - M2, (10, 1)) - M2, (10, 1) - M

# Not the sam

 $\begin{array}{c} {}^{t} \mathcal{W} : \mathcal{P}_{\mathcal{Y}} \\ \mathcal{W} : \mathcal{P}_{\mathcal{Y}} \\ \mathcal{H} : \mathcal{P}_{\mathcal{$  $F_{SL}^{1S,\alpha\beta} = \frac{16\pi^{2}c_{w}s_{w}}{32\pi^{2}c_{w}s_{w}} \left( \Gamma_{L}^{\beta k} \left( \Gamma_{L}^{\alpha k} \right)^{*} \left( B_{0} \left( m_{\ell_{\beta}}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2} \right) - 2 C_{00} \left( M_{Z}^{2}, m_{\ell_{\alpha}}^{2}, m_{\ell_{\beta}}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2} \right) \right)$   $F_{SL}^{1S,\alpha\beta} = \frac{ie \left( c_{w}^{2} - s_{w}^{2} \right)}{32\pi^{2}c_{w}s_{w}} \left( \Gamma_{L}^{\beta k} \left( \Gamma_{L}^{\alpha k} \right)^{*} \left( B_{0} \left( m_{\ell_{\beta}}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2} \right) - 2 C_{0} \left( m^{2} - m^{2} - M^{2} - M^{2} - M^{2} - M^{2} - M^{2} \right) \right)$  $- m_{\nu_{i}} \Gamma_{ck}^{j} \left(\Gamma_{dk}^{i}\right)^{*} C_{12} \left(M_{Z}^{2}, m_{\nu_{i}}^{2}, m_{\nu_{j}}^{2}, M_{\chi_{c}}^{2}, M_{\chi_{d}}^{2}, M_{\phi_{k}}^{2}\right)\right),$  $32\pi^{2}c_{w}s_{w} + M_{Z}^{2}C_{1}\left(M_{Z}^{2}, m_{\ell_{\beta}}^{2}, m_{\ell_{\alpha}}^{2}, M_{\psi}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2}\right) + m_{\ell_{\alpha}}^{2}C_{1}\left(m_{\ell_{\alpha}}^{2}, m_{\ell_{\beta}}^{2}, M_{Z}^{2}, M_{\psi}^{2}, M_{\phi_{k}}^{2}, M_{\psi}^{2}\right)$  $F_{L}^{1\chi,ij} = \sum_{ck} \frac{-ieU_{\Phi}^{2k}}{16\pi^{2}c_{w}s_{w}} \left( \Gamma_{c3}^{j} \left( M_{\chi_{c}}\Gamma_{c3}^{i} \left( C_{1} \left( m_{\nu_{j}}^{2}, m_{\nu_{i}}^{2}, M_{Z}^{2}, M_{\phi_{k}}^{2}, M_{\chi_{c}}^{2}, M_{A^{0}}^{2} \right) \right) \right)$  $- C_1 \left( m_{\nu_j}^2, m_{\nu_i}^2, M_Z^2, M_{A^0}^2, M_{\chi_c}^2, M_{\phi_k}^2 \right) \right) + m_{\nu_i} \left( \Gamma_{ck}^i \right)^* \left( C_{12} \left( M_Z^2, m_{\nu_i}^2, m_{\nu_i}^2 \right) \right)$ +  $C_{12}\left(M_Z^2, m_{\nu_i}^2, m_{\nu_j}^2, M_{\phi_k}^2, M_{A^0}^2, M_{\chi_c}^2\right)\right)$  +  $m_{\nu_j}\Gamma_{ck}^i\left(\Gamma_{c3}^j\right)^*\left(C_{12}\left(M_Z^2, m_{\chi_c}^2\right)\right)$ +  $C_{12}\left(M_Z^2, m_{\nu_i}^2, m_{\nu_j}^2, M_{\phi_k}^2, M_{A^0}^2, M_{\chi_c}^2\right)\right)$ , 32 00 So TS

A.











### **EWPO:** $Z \rightarrow inv$





### $\rightarrow$ Strong constraint to NP and strongly disfavours light invisible states $M_{inv} < M_Z/2$

 $\rightarrow$  Small tension with the SM, points reaching toward the experimental measurements? **Sensible by future experiments** ?

Again, beyond future sensitivity...

 $\mathcal{V}_{\text{fermion}} = M_{\nu}^{\alpha\beta} \,\overline{\nu_{\alpha}^c} \,\nu_{\beta} - M_{\psi} \,\overline{\psi_1} \,\overline{\psi_2} + \frac{1}{2} \,M_{F\,ii} \,\overline{F_i^c} \,F_i - y_{1i}^* \,\overline{F_i} \,\Phi^\dagger \,\overline{\psi_1} - y_{2i}^* \,\overline{F_i} \,\Phi \,\psi_2{}^c$  $g_{\psi}^{\alpha} \widetilde{\overline{\psi_2}} L_L^{\alpha} S - g_{F_i}^{\alpha} \widetilde{\overline{L_L^{\alpha}}} \eta F_i + g_R^{\alpha} \overline{\overline{e_R^{\alpha}}} \eta^{\dagger} \psi_1 + \text{H.c.}$ 









# **Invisible Higgs decays**



![](_page_21_Picture_3.jpeg)

![](_page_21_Picture_4.jpeg)

![](_page_21_Picture_5.jpeg)

 $\rightarrow$  Strong constraint to NP and strongly disfavours light invisible states  $M_{inv} < M_H/2$ 

 $\rightarrow$  A lower floor driven by  $m_{\nu}$ 

Again, beyond future sensitivity...

$$\begin{split} \mathcal{V}_{\text{fermion}} &= M_{\nu}^{\alpha\beta} \, \overline{\nu_{\alpha}^{c}} \, \nu_{\beta} - M_{\psi} \, \overline{\psi_{1}} \, \widetilde{\psi_{2}} + \frac{1}{2} \, M_{F\,ii} \, \overline{F_{i}}^{c} \, F_{i} - y_{1i}^{*} \, \overline{F_{i}} \, \Phi^{\dagger} \, \widetilde{\psi_{1}} - y \\ &+ g_{\psi}^{\alpha} \, \widetilde{\overline{\psi_{2}}} \, L_{L}^{\alpha} S + g_{F_{i}}^{\alpha} \, \widetilde{\overline{L_{L}^{\alpha}}} \, \eta F_{i} + g_{R}^{\alpha} \overline{e_{R}^{\alpha}} \, \eta^{\dagger} \, \psi_{1} + \text{H.c.} \end{split}$$

![](_page_21_Picture_10.jpeg)

![](_page_21_Picture_11.jpeg)

![](_page_22_Figure_1.jpeg)

 $\rightarrow$  Still not within reach...

![](_page_22_Picture_6.jpeg)

![](_page_22_Picture_7.jpeg)

### $\rightarrow$ A reminder on the **importance of analytical derivation**

![](_page_22_Picture_9.jpeg)

# Z flavour conserving decays

![](_page_23_Figure_1.jpeg)

![](_page_23_Picture_3.jpeg)

![](_page_23_Picture_4.jpeg)

 $\rightarrow$  Representative of the other Z conserving decays

 $\rightarrow$  Large trilinear coupling  $\alpha$  yields sizeable contributions!

 $\rightarrow$  **Points excluded** by other bounds... Actually from  $H \to \ell_{\alpha} \ell_{\alpha}!$ 

$$\begin{split} \mathcal{V}_{\text{scalar}} &= \frac{1}{2} M_S^2 \, \mathbf{S}^2 + \frac{1}{2} \lambda_{4S} \, \mathbf{S}^4 + M_\eta^2 \, |\eta|^2 + \lambda_{4\eta} \, |\eta|^4 + \frac{1}{2} \lambda_S \, \mathbf{S}^2 |\Phi|^2 + \frac{1}{2} \\ &+ \lambda_\eta \, |\eta|^2 \, |\Phi|^2 + \lambda_\eta' \, |\eta \Phi^\dagger|^2 + \frac{1}{2} \lambda_\eta'' \left[ \left( \Phi \eta^\dagger \right)^2 + \text{H.c.} \right] + \alpha \, \mathbf{S} \left[ \Phi \eta^\dagger \right] \end{split}$$

![](_page_23_Picture_9.jpeg)

![](_page_23_Picture_10.jpeg)

![](_page_23_Picture_11.jpeg)

![](_page_23_Picture_12.jpeg)

![](_page_23_Picture_13.jpeg)

![](_page_23_Picture_14.jpeg)

![](_page_24_Figure_1.jpeg)

 $\rightarrow$  SM-Like  $\Delta a_{\mu}$  relaxes constraints

![](_page_24_Picture_4.jpeg)

![](_page_24_Picture_5.jpeg)

![](_page_24_Picture_6.jpeg)

![](_page_24_Picture_10.jpeg)

# **Lepton Flavour Universality Violation**

![](_page_25_Figure_1.jpeg)

### A. Darricau, LPCA Clermont

![](_page_25_Picture_3.jpeg)

![](_page_25_Picture_4.jpeg)

 $R_{\alpha\beta}^{Z/H} \equiv \frac{\Gamma(Z/H \to \ell_{\alpha}\ell_{\alpha})}{\Gamma(Z/H \to \ell_{\beta}\ell_{\beta})}$ 

- $\rightarrow$  Representative of the other ratios
- $\rightarrow$  SM Predicts LFUC (other than mass effects)
- → Z LFUV corrections fully under control! → H LFUV corrections favours small trilinear coupling  $\alpha$

$$\mathcal{P}_{\text{scalar}} = \frac{1}{2} M_S^2 \, \mathbf{S}^2 + \frac{1}{2} \lambda_{4S} \, \mathbf{S}^4 + M_\eta^2 \, |\boldsymbol{\eta}|^2 + \lambda_{4\eta} \, |\boldsymbol{\eta}|^4 + \frac{1}{2} \lambda_S \, \mathbf{S}^2 |\Phi|^2 + \frac{1}{2} \lambda_{\eta} \, |\boldsymbol{\eta}|^2 \, |\Phi|^2 + \lambda_\eta' \, |\boldsymbol{\eta}\Phi^\dagger|^2 + \frac{1}{2} \lambda_\eta'' \left[ \left( \Phi \boldsymbol{\eta}^\dagger \right)^2 + \text{H.c.} \right] + \alpha \, \mathbf{S} \left[ \Phi \boldsymbol{\eta}^\dagger - \mathbf{M} \, \mathbf{S}^\dagger \right] \, \mathbf{M} \, \mathbf{S}^{\dagger} \, \mathbf{M} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{M} \, \mathbf{M}^{\dagger} \, \mathbf{M} \, \mathbf{$$

![](_page_25_Picture_10.jpeg)

# ffects) **ol**!

![](_page_25_Picture_12.jpeg)

![](_page_25_Picture_13.jpeg)

![](_page_26_Picture_0.jpeg)

### $\rightarrow$ **Thorough analysis** on the flavour phenomenology

- $\rightarrow$  Found a new viable DM candidate:  $A^0$
- $\rightarrow$  Put forward the **consequences** of relaxing  $\Delta a_{\mu}$
- $\rightarrow$  Parameter space favoured by leptogenesis **disfavoured by EWPO**

![](_page_26_Picture_6.jpeg)

![](_page_26_Picture_7.jpeg)

![](_page_26_Picture_9.jpeg)

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### Backup

![](_page_27_Figure_1.jpeg)

### A. Darricau, LPCA Clermont

![](_page_27_Picture_3.jpeg)

![](_page_27_Picture_4.jpeg)

1. to a to

 $10^{-9}$ 

 $|g_{\psi}^{e}|$ 

 $10^{-7}$ 

10-5

10-11

10-13

### Parameter space $g_{F_2}$ 100 $10^{-1}$ 10-2 10<sup>-3</sup> |**g**⊭<sub>₂</sub>| $10^{-4}$ 10-5 Parameter Range $10^{-6}$ $M_S^2, M_\eta^2$ $[5 \times 10^5, 5 \times 10^6]$ 10<sup>-7</sup> $M_1, M_2$ [100, 20000]10-7 10<sup>-4</sup> 10<sup>-3</sup> 10<sup>-2</sup> 10<sup>-1</sup> 10<sup>0</sup> 10<sup>-5</sup> $10^{-8}$ $10^{-6}$ $|g_{F_2}^e|$ [700, 2000] $M_\psi$ $g_{\psi}^{\alpha}$ $\pm \left[10^{-10}, 10^{-4}\right]$ $y_{11,12,21,22}$ 100 $\left[10^{-19}, 10^{-10}\right]$ $m_{ u_1}$ $10^{-}$ ا *g*<sup>µ</sup><sub>\[</sub>] 10<sup>-2</sup> ا

10-3

 $10^{-4}$ 

![](_page_27_Picture_6.jpeg)

![](_page_27_Figure_7.jpeg)

 $|g_{F_2}^{\tau}|$ 

![](_page_27_Figure_8.jpeg)

### Backup

![](_page_28_Figure_1.jpeg)

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![](_page_28_Picture_3.jpeg)

![](_page_28_Picture_4.jpeg)

![](_page_28_Picture_5.jpeg)

![](_page_28_Figure_6.jpeg)

![](_page_28_Figure_7.jpeg)

### Backup

- $\rightarrow$  Use extra degrees of freedom of *R* to set  $\Delta a_{\mu}$  and keep  $\ell_{\alpha} \rightarrow \ell_{\beta} \gamma$  under control
- $\rightarrow$  This **generates** a fine tuned  $\mathscr{G}$  matrix and  $g_R^{\alpha}$  vector as a byproduct

![](_page_29_Figure_5.jpeg)

![](_page_29_Picture_7.jpeg)

### Casas-Ibara

 $\rightarrow$  Links  $M_{\nu}^{\alpha\beta}$  to its experimental values through  $M_{L}$ ,  $\mathcal{G}$  and R a complex orthogonal matrix

![](_page_29_Picture_10.jpeg)