Probing the general axion-nucleon-nucleon interaction in water Cherenkov experiments

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In collaboration with Diego Guadagnoli (LAPTh), Axel Iohner (LAPTh), Pablo Fernández-Menéndez (DIPC), Ludovico Vittorio (LAPTh) Among low-mass particles, the QCD axion is especially well-motivated

- solves, by a symmetry, the question why the strong force does not violate CP
- provides an excellent Dark Matter candidate



$$\mathcal{L}_{axion-quark} = rac{\partial_{\mu} a}{f_a} ar{q} \left( k_R \gamma_R^{\mu} + k_L \gamma_L^{\mu} 
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The interaction is derivative, so observables will be proportional to

$$\left(\frac{\text{external momenta}}{f_a}\right)^2$$

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Dense and hot enough astro-objects may radiate a large quantity of axions

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Isolated Neutron Star Cooling First generation: [Buschmann et al, 2022] Adding a new cooling mode accelerates cooling  $\Rightarrow$  Constraints on axion-matter couplings

#### Isolated Neutron Star Cooling

First generation: [Buschmann et al, 2022]

#### Supernova Axion Emission

First generation: [Ericson, Mathiot, 1989 ; Carenza et al, 2019 ; Carenza et al, 2020 ; Lella et al, 2022 ; Caputo, Raffelt, 2024] With strange matter: [Cavan-Piton et al, PRL, 2024]















Axions reaching Earth could leave traces in Cherenkov facilities (like neutrinos)



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Steps :Axion emissionCalculation of detectablespectrum calculationparticle spectrum $\frac{dN_a}{dE_a}$  $\frac{dN_i}{dE_i} = \frac{N_{target}}{4\pi d_{SN-Earth}^2} \int \frac{d\sigma}{dE_i} \frac{dN_a}{dE_a} dE_a$ 

#### Processes in the litterature

#### Axion emission :

 $\begin{array}{c} \underbrace{\text{pion-nucleon scattering:}}_{\pi + N \rightarrow N' + a} \quad (E_a \sim 240 \pm 80 \text{ MeV}) \\ \underbrace{\text{nucleon-nucleon scattering:}}_{N_1 + N_2 \rightarrow N_3 + N_4 + a} \quad (E_a \sim 60 \pm 50 \text{ MeV}) \end{array}$ 

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Our conclusions must be proved robust w.r.t. the modeling of:

- <u>axion-matter interactions:</u> we address this question within a consistent EFT approach
- SN-core equation of state (EoS) and thermodynamics: we consider different EoS and also vary the thermodynamic parameters

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*low*  $\downarrow$  *energy* 

parametrized in terms of field field field

axion-hadron interactions and their couplings fixed by global symmetries (i.e. as Noether currents)

In the first generation, all axion-matter couplings can be expressed in terms of the axion-nucleon-nucleon couplings  $C_{aNN}$  ( $N \in \{p; n\}$ ) In the first generation, all axion-matter couplings can be expressed in terms of the axion-nucleon-nucleon couplings  $C_{aNN}$  ( $N \in \{p; n\}$ )

KSVZ : 
$$C_{app} = -0.452(28)$$
 and  $C_{ann} = 0.012(28)$ 

$$\begin{array}{l} \mathsf{DFSZ}:\\ \mathcal{C}_{app} = -0.169(30) - 0.430(15)\sin^2(\beta),\\ \mathcal{C}_{ann} = -0.123(30) + 0.406(15)\sin^2(\beta)\\ \text{ with } \tan(\beta) \in [0.25;170] \end{array}$$

"Agnostic":  $C_{aNN}$  are only constrained by data (NSs or SNe)



Hyper-Kamiokande ( $M_{det} = 374$  kton), Betelgeuse (d = 0.2 kpc)



Hyper-Kamiokande ( $M_{det} = 374$  kton)

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- SNe are excellent probes of fundamental physics, in particular of well-motivated SM extensions.
- The process  $ap 
  ightarrow p\pi^0$  is a promising candidate.



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