



Simulating structure formation in modified gravity

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Modified gravity models (see also the review from Llinares, 2021)

Type of equation	Model	Method
$L(\phi, \dot{\phi}, \ddot{\phi}) = f(\rho, \phi)$ L, f linear	Dynamical DE	Modified expansion
	Coupled DE	Modified expansion + G_{eff}
	Non-local	Modified expansion + G_{eff}
	Vector DE	Modified expansion, power spectrum, and growth history
$L(\phi, \partial_x \phi, \partial_x^2 \phi) = f(\rho, \phi)$ L linear in $\partial_x^2 \phi$	$f(R)$	Multigrid (Jordan / Einstein frame)
	Generalized chameleon	Multigrid
	Symmetron	Multigrid
	Dilaton	Multigrid
$L(\phi, \partial_x \phi, \partial_x^2 \phi) = f(\rho, \phi)$ L non-linear in $\partial_x^2 \phi$	DGP	Multigrid or FFT-relaxation with smoothing / splitting
	Cubic Galileon	Multigrid with operator splitting
	Quartic Galileon	Multigrid with operator splitting
$L(\ddot{\phi}, \dot{\phi}, \partial_x \phi, \partial_x^2 \phi) = f(\rho, \dot{\phi}, \phi)$	Symmetron	Leap-frog
	$f(R)$	Implicit
	Disformal	Leap-frog

"GR" simulations

GR simulations ?

Cosmological action (GR + Λ CDM)

$$S = \int \left[\frac{1}{2\kappa} (R - 2\Lambda) + \mathcal{L}_m \right] \sqrt{-g} d^4x \quad (1)$$

Weakly-perturbed FLRW metric in Newtonian gauge

$$ds^2 = a^2(\eta) \left[- (1 + 2\psi) d\eta^2 + (1 - 2\phi) d\mathbf{x}^2 \right] \quad (2)$$

In GR, $\phi = \psi$

$$\nabla^2 \phi - 3\mathcal{H} (\phi' + \mathcal{H}\phi) = 4\pi G a^2 \delta\bar{\rho} \quad (3)$$

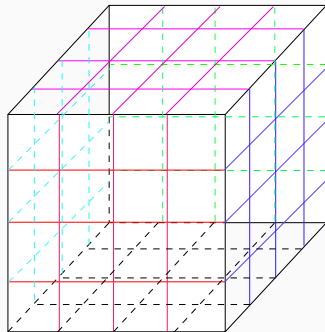
GR simulations ?

$$ds^2 = a^2(\eta) \left[- (1 + 2\psi) d\eta^2 + \underbrace{(1 - 2\phi)}_{\text{Problem!}} d\mathbf{x}^2 \right] \quad (4)$$

GR simulations ?

$$ds^2 = a^2(\eta) \left[-(1 + 2\psi) d\eta^2 + \underbrace{(1 - 2\phi)}_{\text{Problem!}} d\mathbf{x}^2 \right] \quad (4)$$

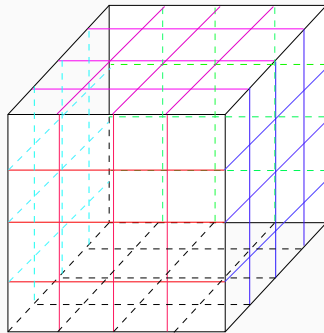
Not consistent with N-body simulations
(fixed grid geometry) !



GR simulations ?

$$ds^2 = a^2(\eta) \left[- (1 + 2\psi) d\eta^2 + \underbrace{(1 - 2\phi)}_{\text{Problem!}} d\mathbf{x}^2 \right] \quad (4)$$

Not consistent with N-body simulations
(fixed grid geometry) !



Solution

- Use the so-called N-body gauge
- Newtonian gauge + corrections in post-processing

From GR...

$$\nabla^2 \phi - 3\mathcal{H}(\phi' + \mathcal{H}\phi) = 4\pi G a^2 \delta \bar{\rho} \quad (5)$$

Newtonian simulations

From GR...

$$\nabla^2 \phi - 3\mathcal{H} \left(\overbrace{\cancel{\phi}}^{\text{Quasi-static}} + \mathcal{H}\phi \right) = 4\pi G a^2 \delta \bar{\rho} \quad (5)$$

Newtonian simulations

From GR...

$$\nabla^2\phi - 3\mathcal{H} \left(\overbrace{\cancel{\phi}}^{\text{Quasi-static}} + \underbrace{\mathcal{H}\cancel{\phi}}_{\text{Sub-horizon}} \right) = 4\pi G a^2 \delta\bar{\rho} \quad (5)$$

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... to Newton !

$$\boxed{\nabla^2\phi = 4\pi G a^2 \delta\bar{\rho}}$$

Can be solved using particle-tree or particle-mesh methods.

Solving PDEs in N-body simulations

Particle methods

- Particle-Particle, Tree codes (Barnes & Hut), Fast Multipole Methods (FMM)
- Solves $1/r^2$ interactions
- **Pros** : very fast, accurate on small scales
- **Con** : only works for Newtonian simulations

Particle-Mesh (PM) methods

- FFT, **Multigrid**, Conjugate gradient
- Solves $\nabla^2 \phi = 4\pi G a^2 \delta \bar{\rho}$
- **Pros** : extremely fast, **solves non-linear elliptic/parabolic PDEs**
- **Con** : inaccurate on small scales

PM with Adaptive-Mesh Refinement (PM-AMR) methods

- **Multigrid**, Conjugate gradient
- **Pros** : accurate on small scales, **solves non-linear elliptic/parabolic PDEs**
- **Cons** : Not as fast as previous methods, anisotropic

PM (+ AMR) solver

- Project the particle mass (density) onto a grid
- Compute the gravitational potential (solver!)
- Compute the force
- Interpolate back to particles
- Update velocities, then positions

First possibility : FFT

$$\phi(k) = -k^{-2}W(k) \times 4\pi G a^2 \delta\bar{\rho} \quad (6)$$

- "Exact" case $W(k) = 1$ is inaccurate on small scales
- Several possibilities for $W(k)$
- Also possible to directly compute the force !

Multigrid

Discretize the Laplacian operator : $\nabla^2 \phi_i = \frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{h^2}$

Iterative method (smoothing) : $\phi_i = \frac{1}{2} (\phi_{i+1} + \phi_{i-1} - h^2 4\pi G a^2 \delta \bar{\rho})$

Principles of Multigrid

- Smooth on fine grid (pre-smoothing)
- Compute residual $r_h = f_h(\delta) - \nabla_h^2 \phi_h \rightarrow$ Interpolate to coarser grid
- Compute the correction term $\nabla_H^2 \tilde{\phi}_H = -\mathcal{R}_H(r_h)$
- Apply to correction on the fine grid $\phi_h = \phi_h + \mathcal{P}_h(\tilde{\phi}_H)$
- Smooth on fine grid (post-smoothing)

Can solve non-linear equations using the Newton-Raphson method

$$\phi^{\text{new}} = \phi^{\text{old}} - \frac{\mathcal{L}(\phi^{\text{old}})}{\partial \mathcal{L} / \partial \phi^{\text{old}}} \quad (7)$$

Modified gravity simulations

Minimally coupled Quintessence

Action

$$S = \int \sqrt{-g} d^4x \left[\frac{1}{2\kappa} R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) + \mathcal{L}_m \right] \quad (8)$$

Dark-energy equation of state :

$$w = \frac{\dot{\phi}^2/2 - V(\phi)}{\dot{\phi}^2/2 + V(\phi)} \quad (9)$$

Popular models : Ratra-Peebles, SUGRA...

How to simulate ?

- Newtonian dynamical solver (no change)
- Modification of cosmological tables (background expansion)

Parametrized gravity

Set of equations

$$\nabla^2 \psi = 4\pi G \mu(a) a^2 \delta \bar{\rho} \quad (10)$$

$$\nabla^2 (\psi + \phi) = 8\pi G \Sigma(a) a^2 \delta \bar{\rho} \quad (11)$$

How to simulate ?

- Trivial change to the dynamical solver (one multiplication)
- Depending on applications, the second equation is not mandatory
- Background expansion is the same, but D_+ is not (initial rescaling)

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$$\ddot{D} + \mathcal{H}\dot{D} - \beta D = 0 \quad \text{first-order growth factor} \quad (12)$$

$$\ddot{E} + \mathcal{H}\dot{E} - \beta [E - D^2] = 0 \quad \text{second-order growth factor} \quad (13)$$

with $\beta = \frac{3}{2} \mu(a) \Omega_m(a)$

Modified Newtonian Dynamics (MOND)

Modified Poisson equation (AQUAL)

$$\nabla \left[\mu \left(\frac{|\nabla \phi|}{g_0} \right) \nabla \phi \right] = 4\pi G \delta \bar{\rho} \quad (14)$$

Modified Newtonian Dynamics (MOND)

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Quasi-linear MOND (QUMOND)

$$\nabla^2 \phi^N = 4\pi G \rho \quad (15)$$

$$\nabla^2 \phi = \nabla \left[\gamma \left(|\nabla \phi^N| / g_0 \right) \nabla \phi^N \right] \quad (16)$$

Modified Newtonian Dynamics (MOND)

- Simple :

$$\mathfrak{v}(y) = \frac{1}{2} + \frac{\sqrt{1 + 4/y}}{2} \quad (17)$$

Modified Poisson equation (AQUAL)

$$\nabla \left[\mu \left(\frac{|\nabla \phi|}{g_0} \right) \nabla \phi \right] = 4\pi G \delta \bar{\rho} \quad (14)$$

- The n -family :

$$\mathfrak{v}(y) = \left[\frac{1}{2} + \frac{\sqrt{1 + 4/y^n}}{2} \right]^{1/n} \quad (18)$$

- The β -family :

$$\mathfrak{v}(y) = \left(1 - e^{-y} \right)^{-1/2} + \beta e^{-y} \quad (19)$$

Quasi-linear MOND (QUMOND)

$$\nabla^2 \phi^N = 4\pi G \rho \quad (15)$$

$$\nabla^2 \phi = \nabla \left[\mathfrak{v} \left(\left| \nabla \phi^N \right| / g_0 \right) \nabla \phi^N \right] \quad (16)$$

- The γ -family :

$$\mathfrak{v}(y) = \left(1 - e^{-y^{\gamma/2}} \right)^{-1/\gamma} \quad (20)$$

$$+ (1 - \gamma^{-1}) e^{-y^{\gamma/2}} \quad (21)$$

- The δ -family :

$$\mathfrak{v}(y) = \left(1 - e^{-y^{\delta/2}} \right)^{-1/\delta} \quad (22)$$

Modified Newtonian Dynamics (MOND)

QUMOND equations

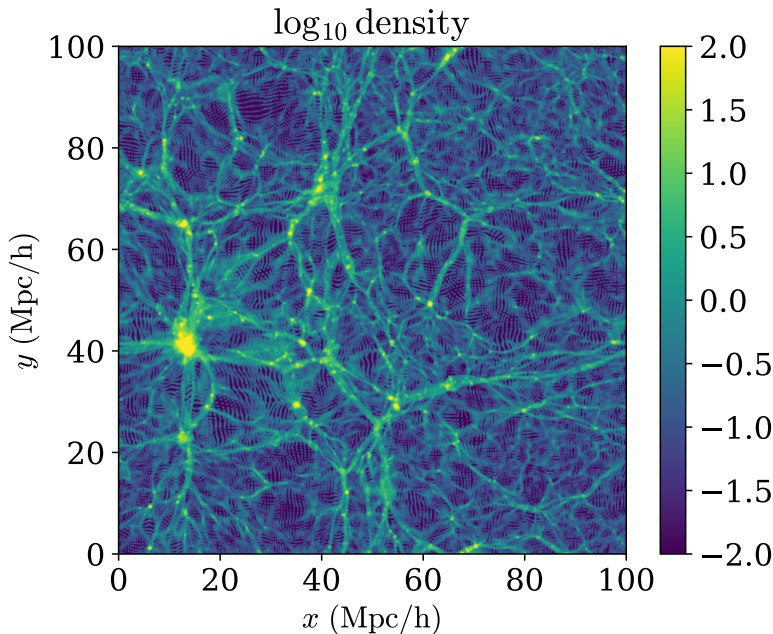
$$\nabla^2 \phi^N = 4\pi G \rho \quad (23)$$

$$\nabla^2 \phi = \nabla \left[\gamma \left(|\nabla \phi^N| / g_0 \right) \nabla \phi^N \right] \quad (24)$$

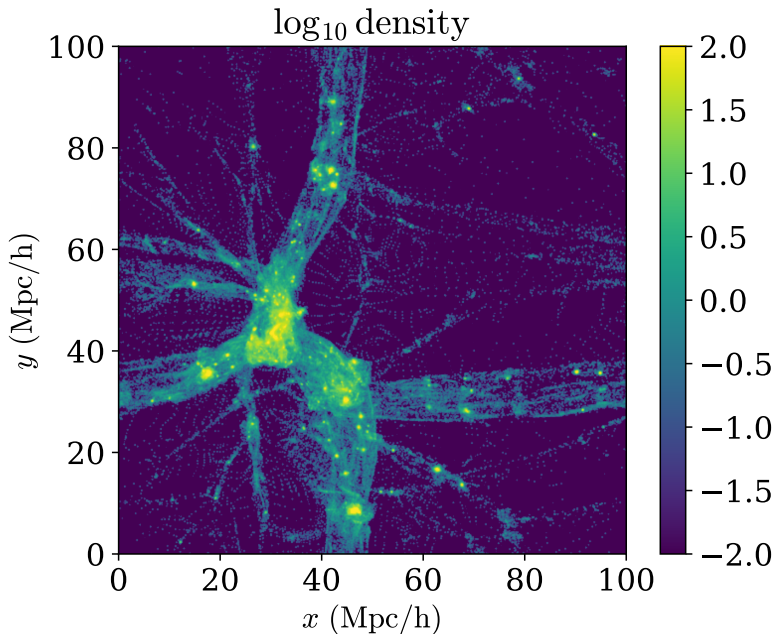
How to simulate ?

- Solve for a Newtonian potential (no change)
- Solve for a MOND potential (compute new source term)
- No modification of cosmological tables
- Arbitrary initial conditions rescaling

Newtonian simulation at $z = 0$



MOND simulation at $z = 0$



$f(R)$ gravity

$f(R)$ action

$$S = \int \left[\frac{1}{2\kappa} (R + f(R)) + \mathcal{L}_m \right] \sqrt{-g} d^4x \quad (25)$$

Hu & Sawicki (2007) model

$$f(R) = -m^2 \frac{c_1 (R/m^2)^n}{c_2 (R/m^2)^n + 1} \quad (26)$$

Only **two free parameters** (one is constrained by observations)

$$\frac{c_1}{c_2} = 6 \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \qquad \frac{c_1}{c_2^2} = -\frac{1}{n} \left[3 + 12 \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \right] f_{R0} \quad (27)$$

Modified set of equations

$$\nabla^2 \phi = \underbrace{4\pi G a^2 \delta \bar{\rho}}_{\text{Newton}} - \frac{c^2}{2} \nabla^2 f_R \quad (28)$$

$$\nabla^2 f_R = A \delta \bar{\rho} + B + C f_R^{-\frac{1}{n+1}} \quad (29)$$

How to simulate ?

- Solve the non-linear equation for the f_R field with Multigrid
- Then, solve a linear equation with modified source term
- In practice, we add the contribution directly in the Force

Modified set of equations

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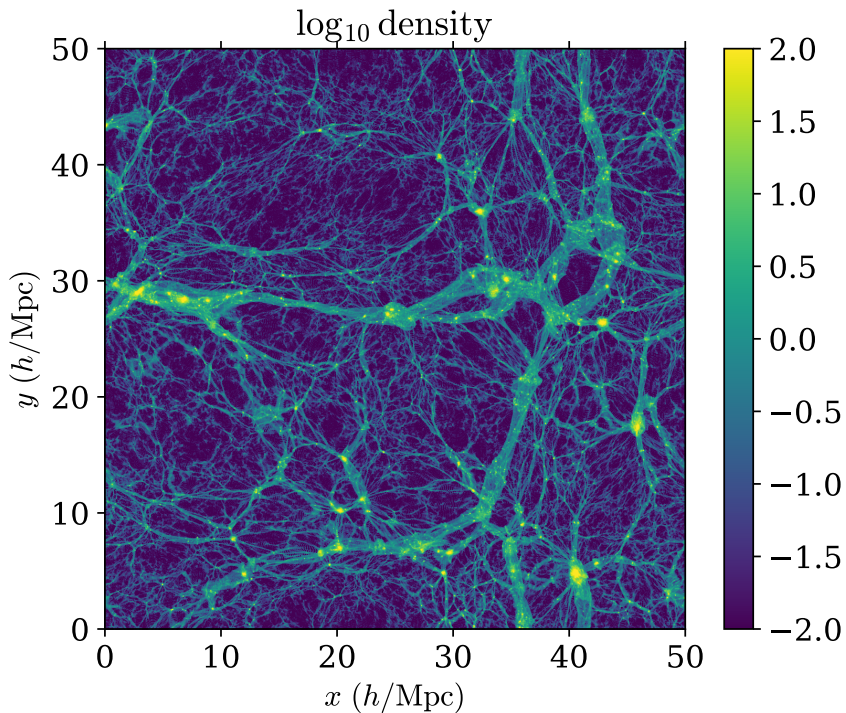
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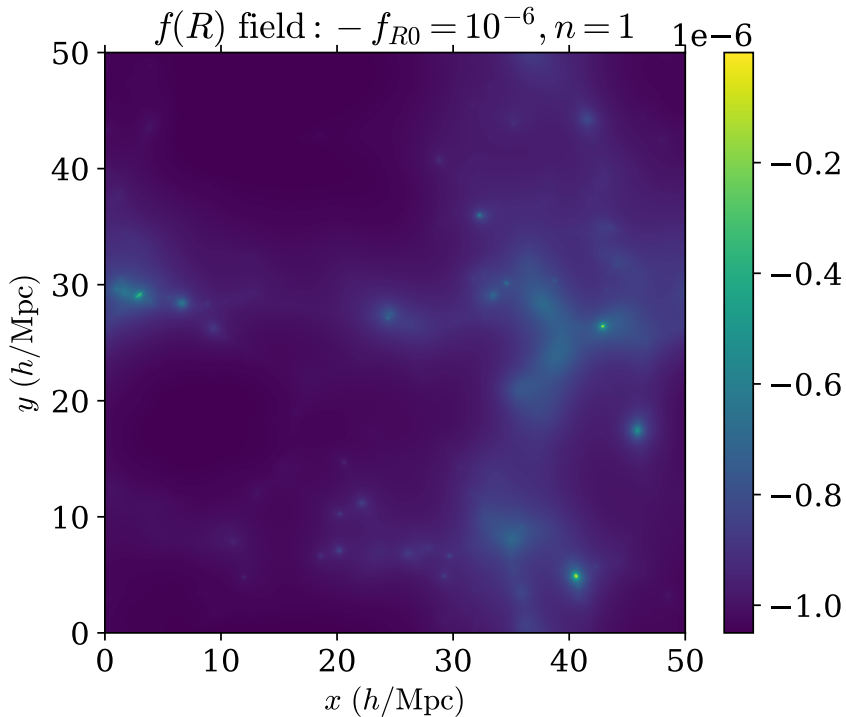
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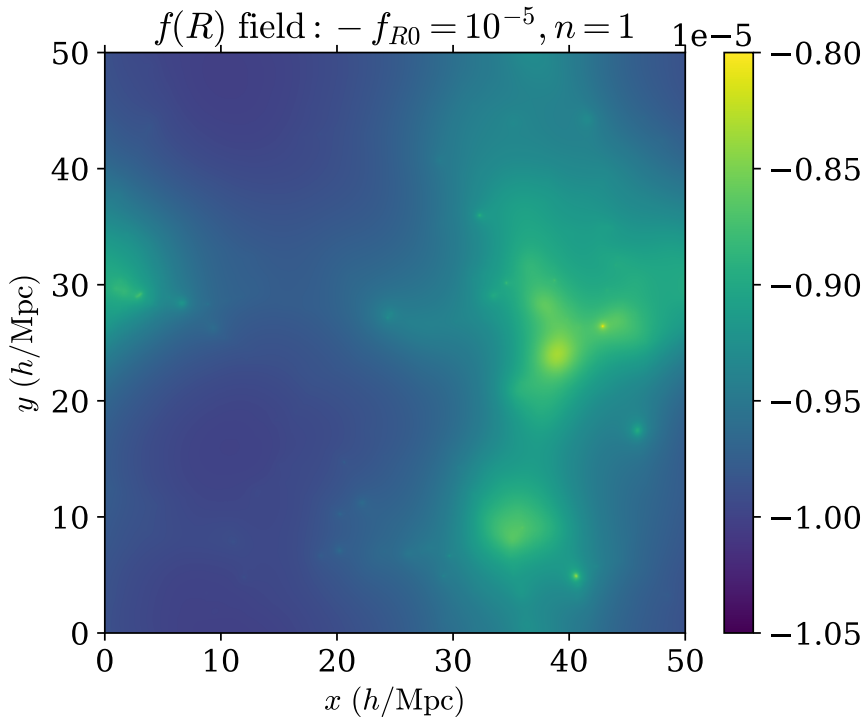
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Approximation

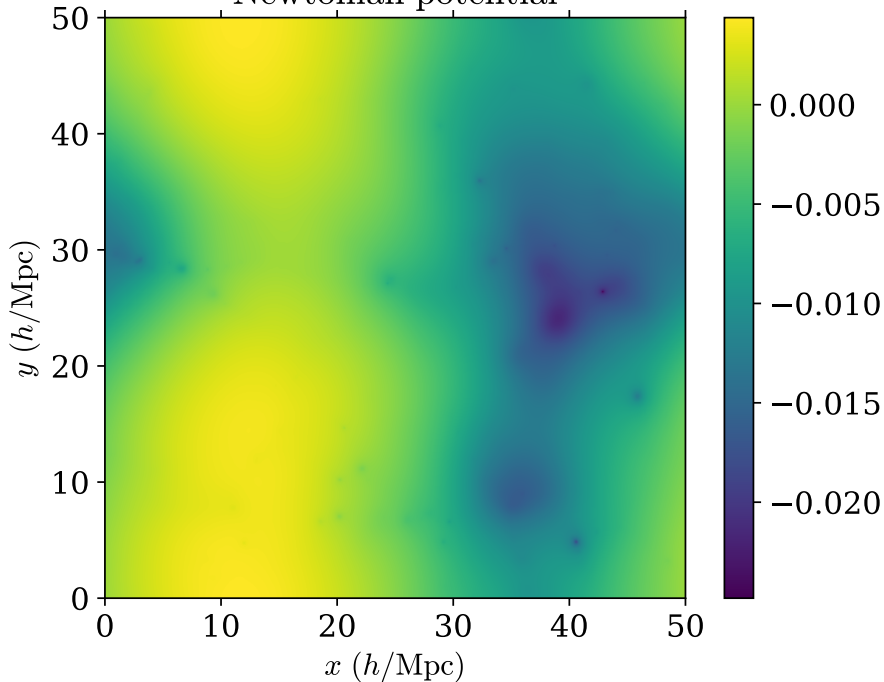
$$-k^2 \phi = 4\pi G_{\text{eff}}(a, k) a^2 \delta \bar{\rho} \quad (30)$$





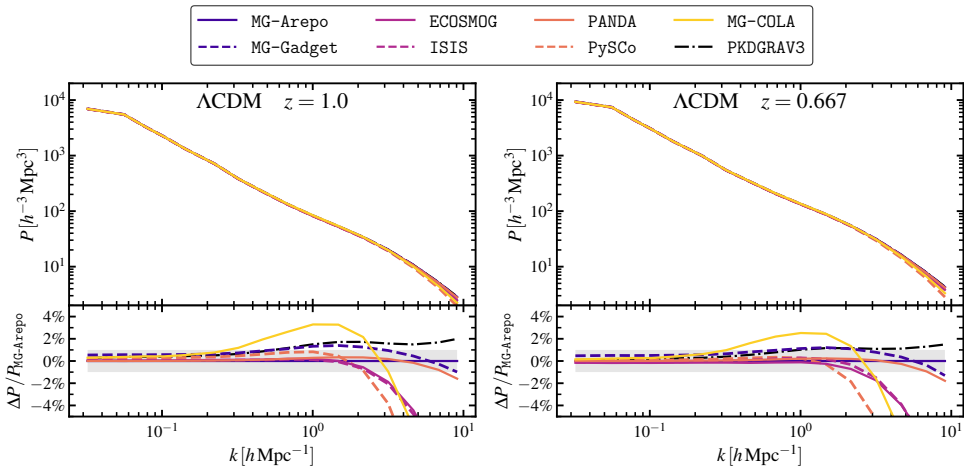


Newtonian potential

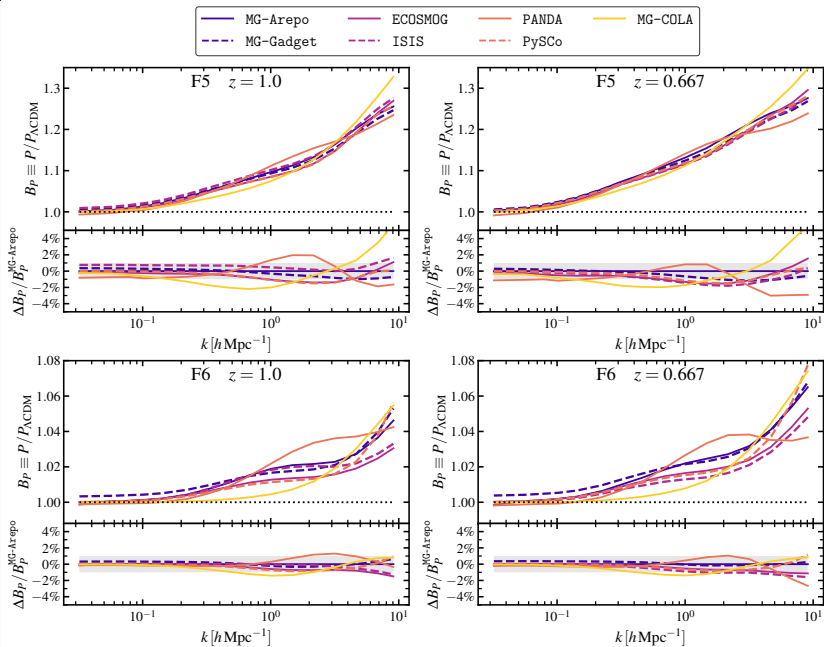


Some code-comparison challenge

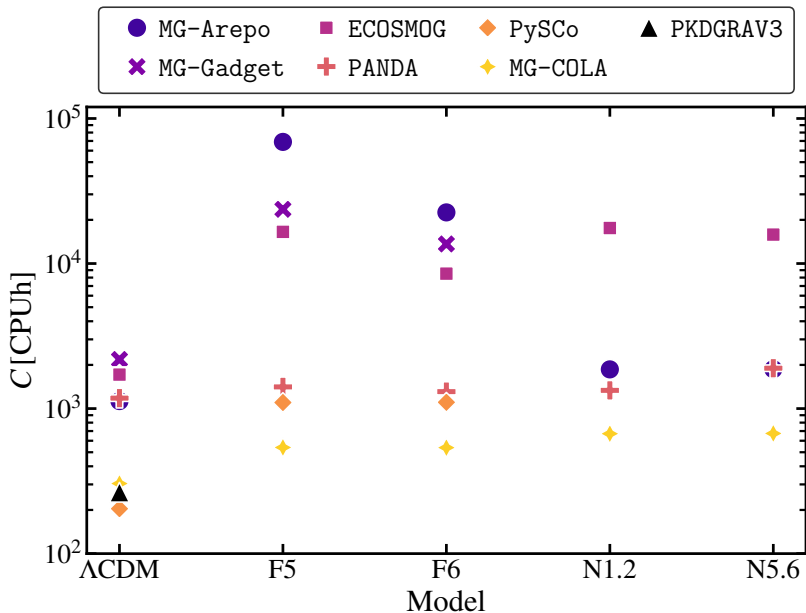
Λ CDM Power spectrum (Euclid Collaboration : Adamek et al., 2025)



$f(R)$ Power spectrum (Euclid Collaboration : Adamek et al., 2025)



Runtime comparison (Euclid Collaboration : Adamek et al., 2025)



Summary

Current status

- Different theories need appropriate numerical tools
- Fair consensus among different numerical implementation
- Increase the range of models we can simulate
- Ongoing work to implement EFTofDE in RAMSES and PySCo by Himanish Ganjoo at LUX

Limitations

- MG simulations are generally more expensive than Newtonians
- MG simulations usually restricted to snapshots
- Need to produce lightcones !
- Current MG N-body codes do not run on GPUs

Run MG simulations with PySCo github.com/mianbreton/pysco

Installation as package (also installs dependencies for you !)

```
python -m pip install pysco-nbody
```

Run from Python script

```
import pysco

param = {
    "nthreads": 1,
    "theory": "newton", # mond, fr...
    "H0": 72,
    "z_out": "[10, 5, 2, 1, 0.5, 0]"
    # Add other parameters
}

pysco.run(param)
```

Appendix

Preliminary work : EFTofDE (by Himanish Ganjoo)

The Effective-field theory of dark energy framework

General action (second order)

$$S = \int d^4x \sqrt{-g} \left[\frac{M_*^2}{2} f(t) R - \Lambda(t) - c(t) g^{00} + \frac{M_2^4(t)}{2} (\delta g^{00})^2 - \frac{\bar{M}_1^3(t)}{2} \delta g^{00} \delta K - \frac{\bar{M}_2^2(t)}{2} (\delta K)^2 - \frac{\bar{M}_3^2(t)}{2} \delta K_{ij} \delta K^{ij} + \frac{\hat{M}^2(t)}{2} \delta g^{00} \delta R^{(3)} + \dots \right] \quad (31)$$

Rewritten with 4 time-varying α -functions, encoding deviations from GR/ Λ CDM

Cubic screening (simplest non-linear)

$$\nabla^2 \psi = 4\pi G a^2 \delta \bar{\rho} + (\alpha_B - \alpha_M) \nabla^2 \chi \quad (32)$$

$$p \nabla^2 \chi + q 4\pi G a^2 \delta \bar{\rho} + r \left[(\nabla^2 \chi)^2 - \partial_i \partial_j \chi \partial^i \partial^j \chi \right] = 0 \quad (33)$$

