

INTERACTING DARK SECTOR WITH INTRINSIC ENTROPY COUPLINGS

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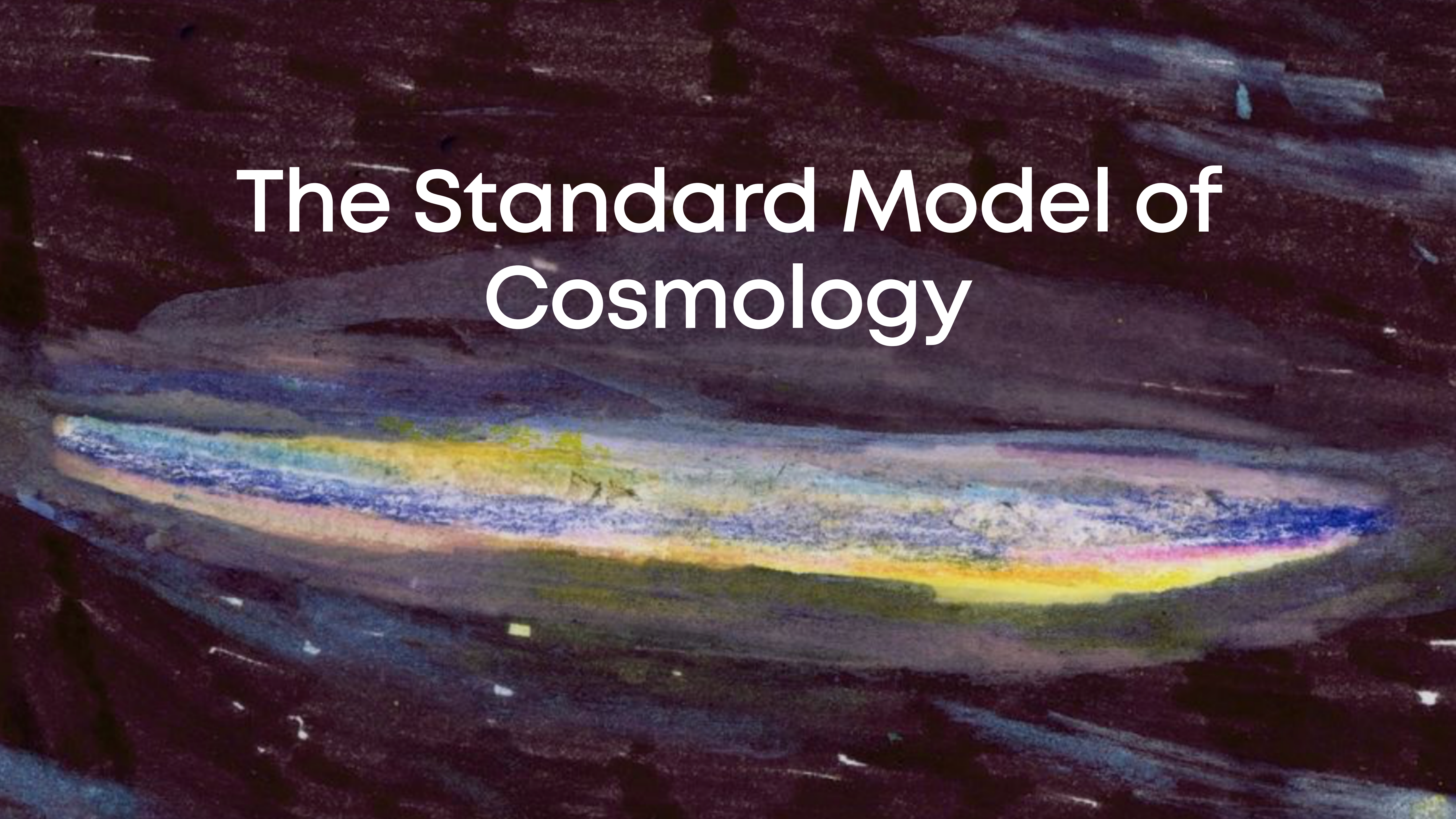
Based on ongoing work with:

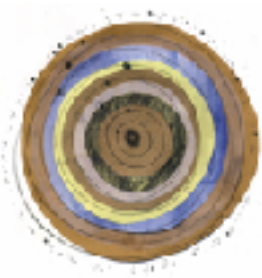
- Erik Jensko, UCL
- Vivian Poulin, U. Montpellier
- Tristan L. Smith, Swarthmore U.



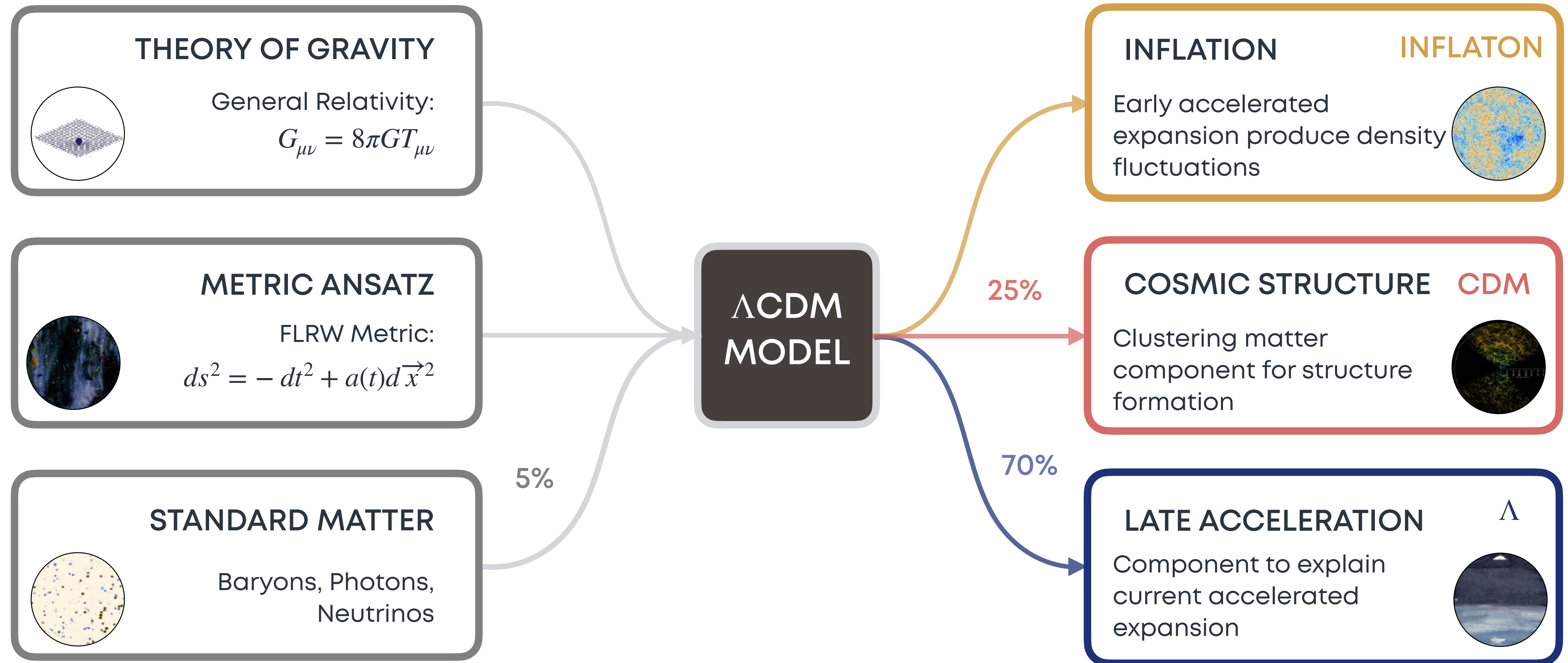
Illustrations: Inês Viegas Oliveira (ivoliveira.com)

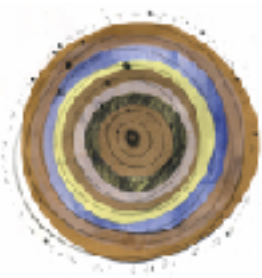
The Standard Model of Cosmology





The Lambda Cold Dark Matter Model





Challenges to the Λ CDM Model

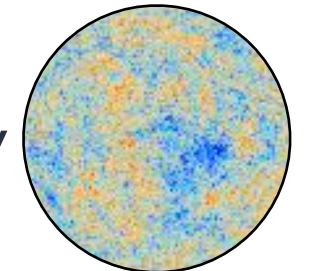
The Λ CDM model relies on:

- Inflation but needs firm theoretical grounds: primordial power spectrum of quantum fluctuations (simplest parameterisation in terms of spectral index and amplitude)
- Dark matter being a pressureless fluid of unknown nature/origin and no detection success (new particle(s) in the SM)
- Dark energy being a cosmological constant (Λ) with unknown nature/origin (vacuum energy, properties of empty space, etc)

INFLATION

INFLATON

Early accelerated expansion produce density fluctuations



COSMIC STRUCTURE

CDM

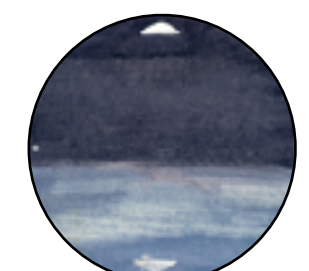
Clustering matter component for structure formation

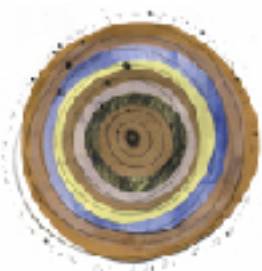


LATE ACCELERATION

Λ

Component to explain current accelerated expansion





Challenges to the Λ CDM Model

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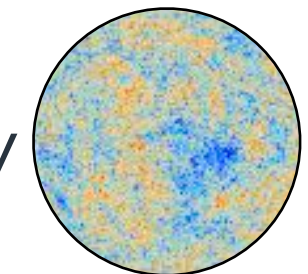
Cosmic tensions may signal that Λ CDM is incomplete:

- Anomalies in the CMB: lensing, curvature, etc
- The matter clustering S_8 tension
- The Hubble/ H_0 expansion rate tension

INFLATION

INFLATON

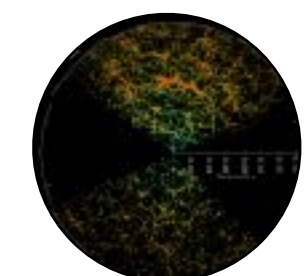
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COSMIC STRUCTURE

CDM

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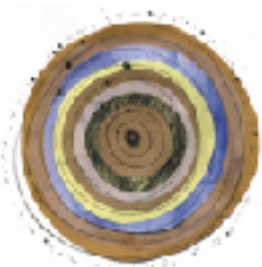


LATE ACCELERATION

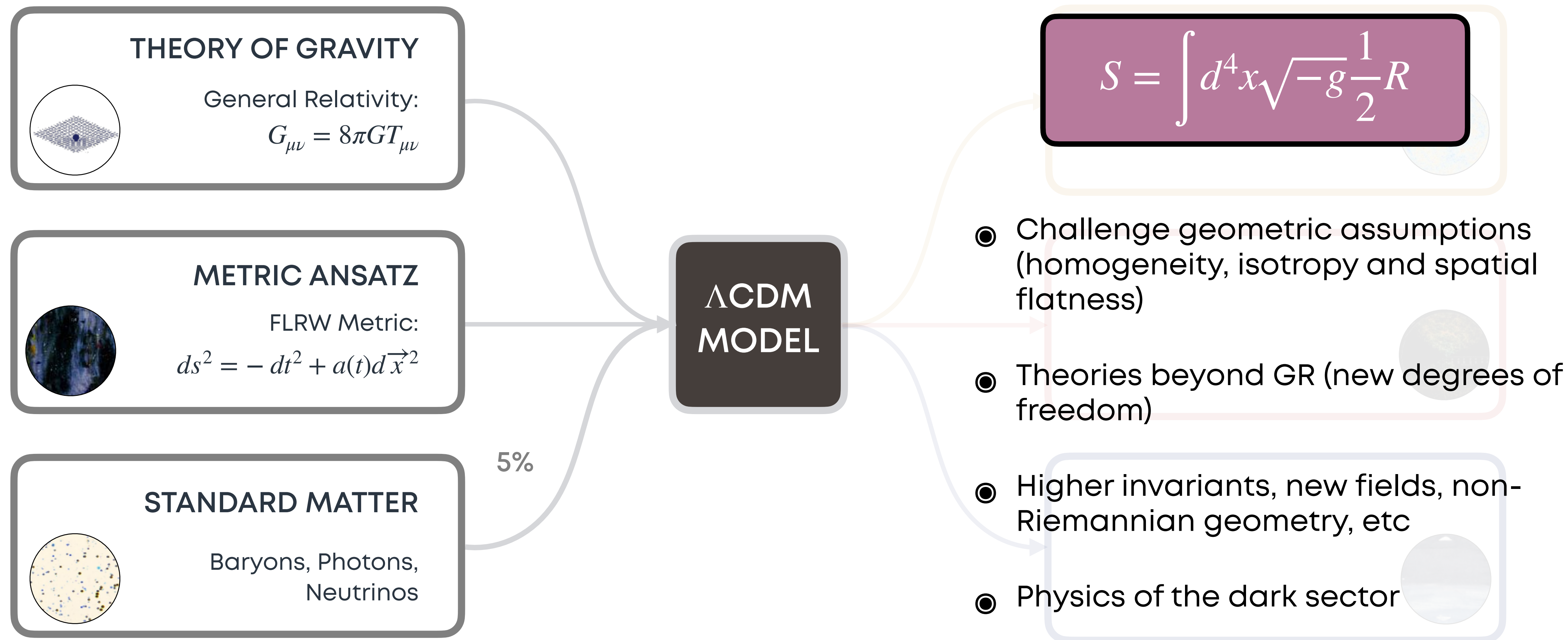
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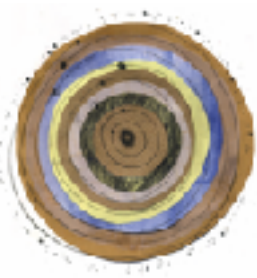


Going Beyond the Standard Model



Interacting dark sector with intrinsic entropy couplings

Based on: [E. Jensko, E. M. Teixeira, V. Poulin, T.L. Smith: [arxiv:25xx.xxxxx](#)]



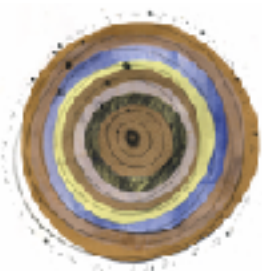
Interacting dark sector

- Background cosmological couplings between dark matter (DM) and dark energy (DE) can be imposed at the level of the field equations:
- However, these are just phenomenological modifications:
- Prone to instabilities [Valiviita et. al. arxiv:0804.0232]
- Ambiguous definition of cosmological perturbations (e.g., added by hand)
- Couplings often heavily constrained by background observations
- Focus on **Lagrangian formulations**



$$\begin{aligned}\dot{\rho}_{\text{DM}} + 3H(\rho_{\text{DM}} + p_{\text{DM}}) &= Q \\ \dot{\rho}_{\text{DE}} + 3H(\rho_{\text{DE}} + p_{\text{DE}}) &= -Q\end{aligned}$$

$$\begin{aligned}\delta'_{\text{DM}} + 3\mathcal{H} \left(\frac{\delta p_{\text{DM}}}{\delta \rho_{\text{DM}}} + w_{\text{DM}} \right) \delta_{\text{DM}} &= -(1 + w_{\text{DM}})(\theta_{\text{DM}} - 3\Phi') + F(Q, \delta Q) \\ \theta'_{\text{DM}} + \left[\mathcal{H}(1 - 3w_{\text{DM}}) + \frac{w'_{\text{DM}}}{1 + w_{\text{DM}}} \right] \theta_{\text{DM}} &= k^2 \left[\Psi + \frac{\delta p_{\text{DM}}}{\delta \rho_{\text{DM}}} \frac{\delta_{\text{DM}}}{1 + w_{\text{DM}}} + G(Q, \delta Q) \right]\end{aligned}$$



Cosmological fluids

Perfect fluid energy-momentum tensor:

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu + pg_{\mu\nu}, \quad \nabla^\mu T_{\mu\nu} = 0$$

Key thermodynamic quantities:

- particle number density n
- intrinsic entropy/particle s
- energy density $\rho = \rho(n, s)$
- pressure p
- temperature T
- chemical potential μ
- + four components of the 4-velocity u^μ with $u^\mu u_\mu = -1$

● EoM (continuity & Euler equations) from projections along u_μ and $h_{\mu\nu}$

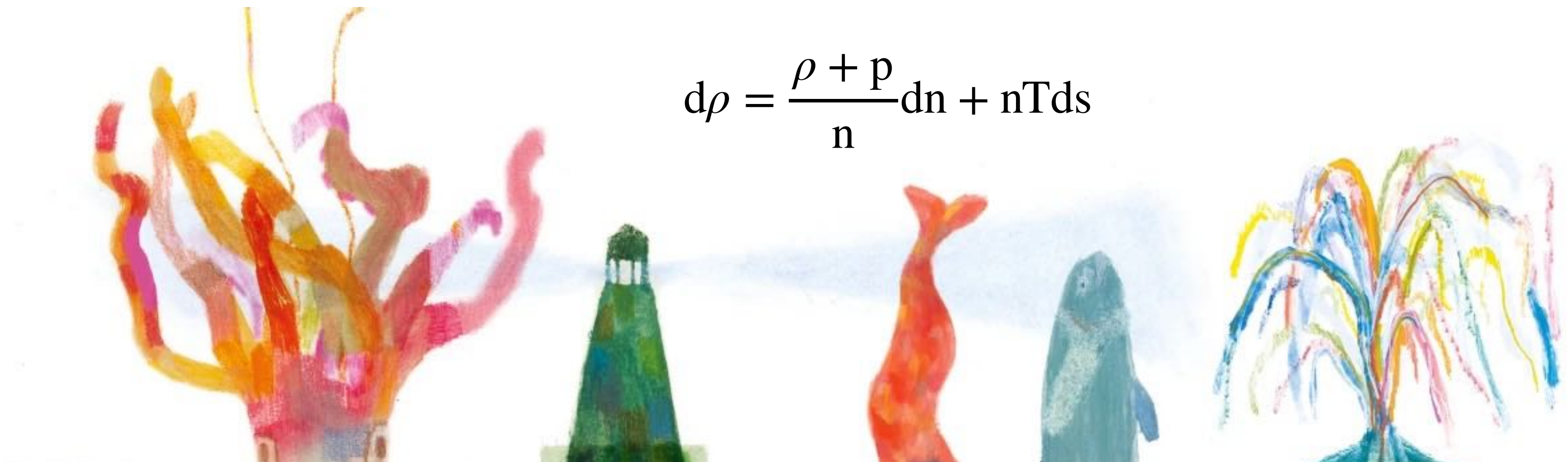
● Thermodynamic identities

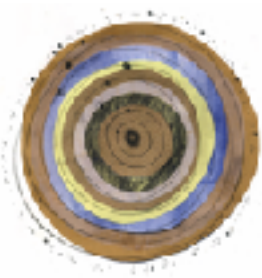
$$p(n, s) = n \left(\frac{\partial \rho}{\partial n} \right)_s - \rho$$

$$T(n, s) = \frac{1}{n} \left(\frac{\partial \rho}{\partial s} \right)_n$$

● Satisfying 1st law of thermodynamics

$$d\rho = \frac{\rho + p}{n} dn + nTds$$





Perfect fluid action

- Model DM as a **perfect fluid** through action with variables $g_{\mu\nu}$, n^μ , s with Lagrange multipliers φ , θ - define the fluid's internal structure and conservation laws covariantly

$$S_{\text{Brown}} = \int \sqrt{-g} \left[-\rho(n, s) + n^\mu (\varphi_{,\mu} + s\theta_{,\mu}) \right] d^4x$$

[J. D. Brown, Class. Quant. Grav. 10, 1579 (1993), arXiv:gr-qc/9304026]

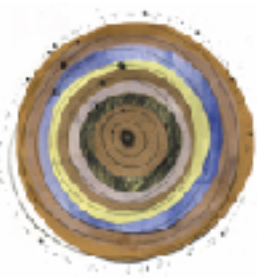
where we introduced the number flux $n^\mu := nu^\mu$ with $n = \sqrt{n^\mu n_\mu}$

constraints

- Variations give perfect fluid $T_{\mu\nu}$ with conservation equation $\nabla^\mu T_{\mu\nu} = 0$
- Additional constraints** preserving particle number and entropy fluxes

$$\nabla_\mu (nu^\mu) = 0, \quad \nabla_\mu (snu^\mu) = 0 \quad \implies \quad u^\mu \nabla_\mu s = 0$$





Cosmological fluids

$$\nabla_\mu(nu^\mu) = 0$$
$$\nabla_\mu(nsu^\mu) = 0$$

- Conservation of particle number and entropy along fluid flow lines
- Taken together these imply the entropy per particle is constant along fluid flow

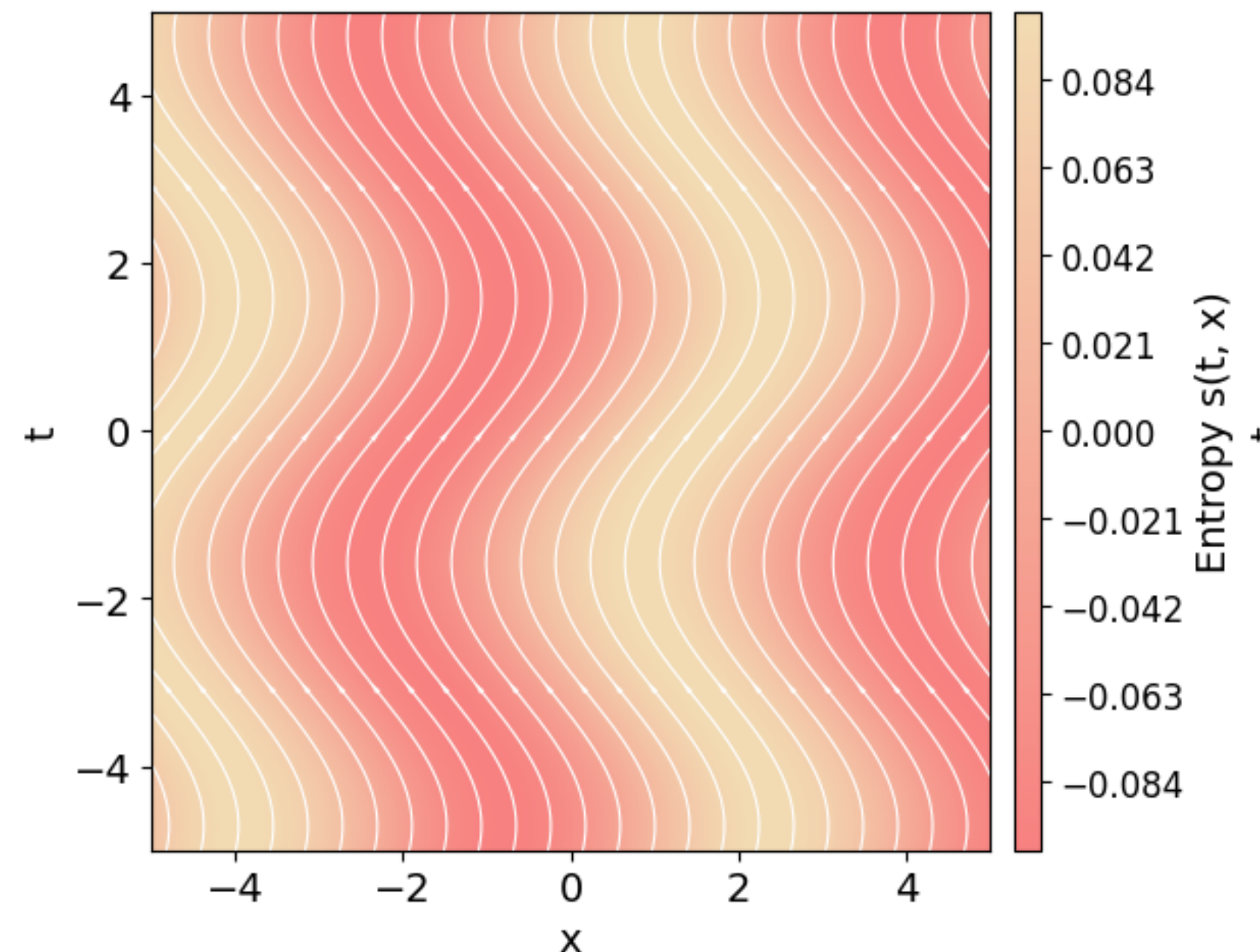
$$\text{adiabatic: } u^\mu \nabla_\mu s = 0 \quad \Rightarrow \quad \text{isentropic: } \nabla_\mu s = 0$$

General fluid condition:

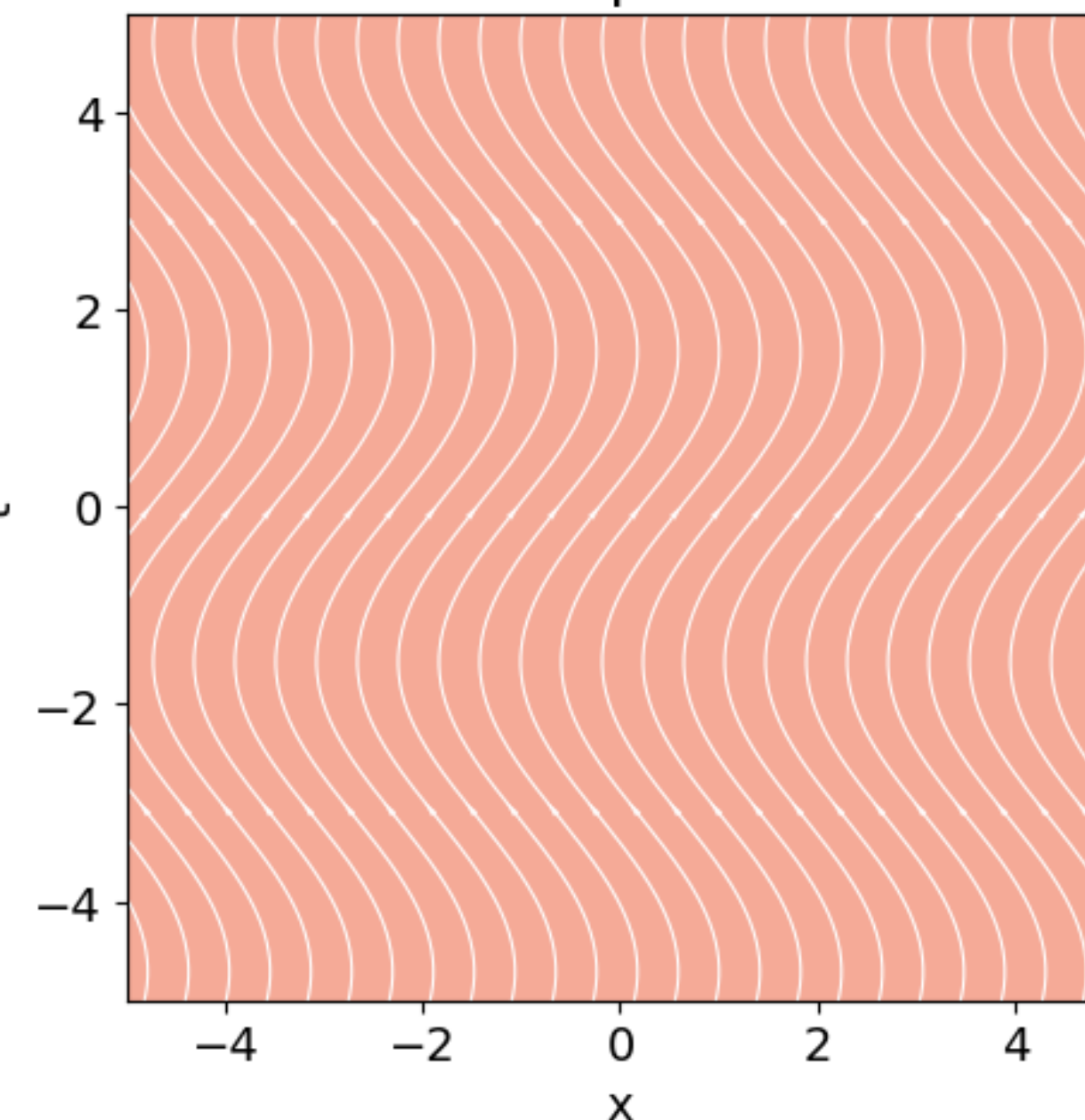
uniform entropy per particle s along the fluid flow (no shocks or heat conduction):

$$\rho = \rho(n, s)$$

Adiabatic flow



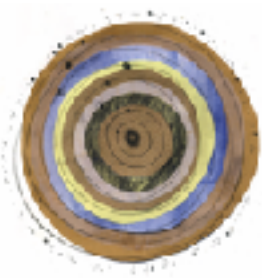
Isentropic flow



Usual assumption in cosmology:

the entropy per particle s is itself time and spatially uniform (irrotational flow):

$$\rho = \rho(n)$$



Perfect fluid dark matter in FLRW

Traditional CDM	General perfect fluid DM	“Entropic” CDM
Barotropic and isentropic ($\nabla_\mu s = 0$)	Non-barotropic and adiabatic ($u^\mu \nabla_\mu s = 0$)	Barotropic and adiabatic ($u^\mu \nabla_\mu s = 0$)
$\rho(n) = nm$	$\rho(n, s)$ arbitrary	$\rho(n, s) = nm\gamma(s)$
$p = 0$	$p \neq 0$	$p = 0$
$c_s^2 = 0$	$c_s^2 \neq 0$	$c_s^2 = 0$
$\Gamma = 0$	$\Gamma \neq 0$	$\Gamma = 0$
$\bar{s} = 0$	$\bar{s} = \text{constant}$	$\bar{s} = \text{constant}$
$\delta s = 0$	$\delta s = \delta s(x, y, z)$	$\delta s = \delta s(x, y, z)$
$T = 0$	$T \neq 0$	$T = m\gamma'(s)$

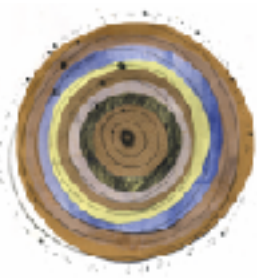
$$\delta p = c_a^2 \delta \rho + \Gamma$$

with non-adiabatic pressure perturbation

$$\Gamma = \left. \frac{\partial p}{\partial s} \right|_\rho \delta s$$

- intrinsic ~ fluctuations in fluid's internal entropy ($\Gamma = 0 \nRightarrow \delta s = 0$)

- relative ~ relative density isocurvature perturbations



Dark sector interactions

- Interactions in the dark sector accomplished through non-minimally coupled scalar fields ϕ

$$S_{\text{tot}} = \int \sqrt{-g} f(n, s, \phi, \dots) d^4x + S_{\text{Brown}} + S_{\phi} + S_{\text{SM}}$$

Many different choices of interacting terms in the literature:

Interactions

- Algebraic (“coupled quintessence”): $f(n, \phi)$
- Derivative (“momentum coupling”): $f(n) n^{\mu} \partial_{\mu} \phi$
- General with $Y = \partial_{\mu} \phi \partial^{\mu} \phi$ and $Z = u^{\mu} \partial_{\mu} \phi$: $f(n, \phi, Y, Z)$

[Boehmer et al., arxiv:1501.06540]

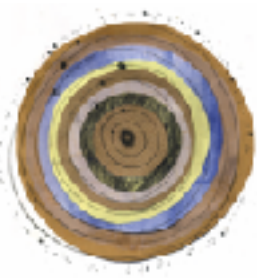
[Boehmer et al., arxiv:1502.04030]

[Pourtsidou et al., arxiv:1307.0458]

Type-1: $Y + V(\phi) + e^{\alpha(\phi)}$ Type-2: $Y + V(\phi) + nh(Z)$ Type-3: $Y + V(\phi) + \gamma(Z) + f(n)$

E.g. Type-3 models is pure momentum-transfer up to linear order - background unchanged





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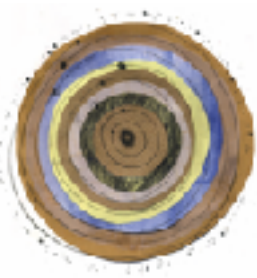
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All past studies neglect the role of entropy in couplings $f(n, s, \phi, \partial\phi)$ by assuming $\partial_{\mu} s = 0$

Type-1: $Y + V(\phi) + e^{\alpha(\phi)}$ Type-2: $Y + V(\phi) + nh(Z)$ Type-3: $Y + V(\phi) + \gamma(Z) + f(n)$

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Entropy couplings in the dark sector

- Couple the intrinsic entropy of dark matter s in $\rho_c(n_c, s) = m_c n_c \gamma(s)$ to the dark energy scalar ϕ - excite the otherwise hidden degrees of freedom associated with the entropy perturbations in barotropic fluids

- For this we consider pure entropy-scalar interactions:

$$\mathcal{L}_{\text{int}}(\phi, s, \nabla_\mu \phi, \nabla_\mu s, \dots)$$

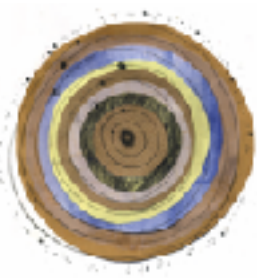
$$\mathcal{L}_{\text{alg}} = g(\phi, s)$$

$$\mathcal{L}_{\text{der}} = h(\nabla_\mu \phi, \nabla_\mu s)$$

- Later distinguish into algebraic and first-order derivative model
- Generally derived from an interaction action

$$S_{\text{int}} = - \int \sqrt{-g} f(s, \phi, \mathcal{S}) d^4x \quad \text{with} \quad \mathcal{S} := \nabla^\mu \phi \nabla_\mu s$$





Equations of motion

- Modified Einstein field equations

$$G_{\mu\nu} = \kappa \left(T_{\mu\nu} + T_{\mu\nu}^{\phi} + T_{\mu\nu}^{(\text{int})} \right)$$

with interaction EM tensor (with heat-fluxes $q_{\mu} = u^{\nu} h_{\mu}^{\lambda} T_{\nu\lambda}$):

$$T_{\mu\nu}^{(\text{int})} = -g_{\mu\nu} f + 2 \frac{\partial f}{\partial \mathcal{S}} \nabla_{(\mu} \phi \nabla_{\nu)} \mathcal{S}$$

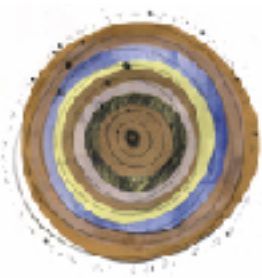
- Total effective dark energy tensor $\tilde{T}_{\mu\nu}^{(\phi)} := T_{\mu\nu}^{(\phi)} + T_{\mu\nu}^{(\text{int})}$ with coupling current

$$J_{\nu} := \nabla^{\mu} \tilde{T}_{\mu\nu}^{(\phi)} = \nabla_{\mu} (f_{,\mathcal{S}} \nabla^{\mu} \phi) \nabla_{\nu} \mathcal{S} - f_{,\mathcal{S}} \nabla_{\nu} \mathcal{S}$$

which acts only in the perpendicular direction $u^{\nu} J_{\nu} = 0$ (from $u^{\nu} \nabla_{\nu} \mathcal{S} = 0$)

Pure “momentum” transfer (up to linear order) - preserve background dynamics

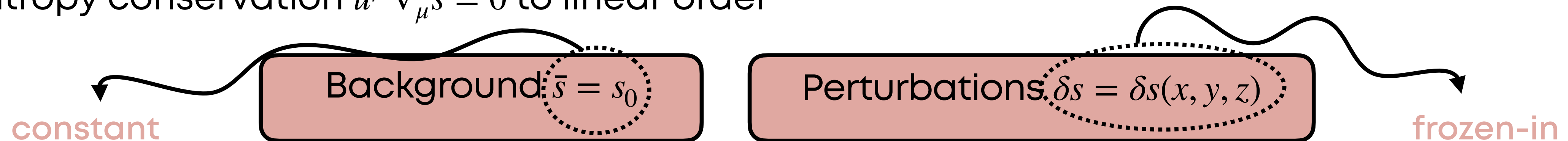




Cosmology

● Conformal Newtonian gauge at linear order: $ds^2 = -a(\tau)^2(1 + 2\Phi)d\tau^2 + a^2(1 - 2\Psi)(dx^2 + dy^2 + dz^2)$

● Entropy conservation $u^\mu \nabla_\mu s = 0$ to linear order

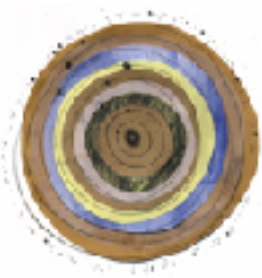


● Trivial background modification with effective potential $\hat{V}(\bar{\phi}) = V(\bar{\phi}) + f(\bar{\phi}, s_0, 0)$

$$\bar{\rho}_{\text{DE}} = \frac{1}{2a^2} \bar{\phi}'^2 + \hat{V}(\bar{\phi}), \quad \bar{p}_{\text{DE}} = \frac{1}{2a^2} \bar{\phi}'^2 - \hat{V}(\bar{\phi})$$

● Uncoupled dark sector in the background - perturbative effects only

$$\bar{\rho}_I' + 3\mathcal{H}(\bar{p}_I + \bar{\rho}_I) = 0 \quad \Rightarrow \quad \text{indistinguishable from uncoupled quintessence or } \Lambda\text{CDM}$$



Cosmological perturbations

- Split interaction term and take $p_c = \delta p_c = 0$ for simplicity

Algebraic

$$f(\phi, s, \nabla_\mu \phi \nabla^\mu s) = g(\phi, s) + h(\nabla_\mu \phi \nabla^\mu s)$$

Derivative

- Unchanged continuity equation but modified Euler equation with scale-dependent source

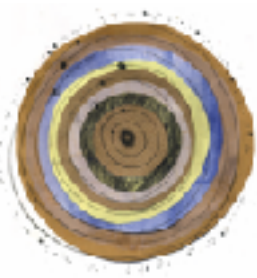
$$\delta'_c + \theta_c - 3\Psi' = 0$$

$$\theta'_c + \theta_c \mathcal{H} - k^2 \Psi = -k^2 Q_s \delta s(k), \quad \text{with} \quad Q_s := \frac{\bar{g}_{,s} - h_0 \hat{V}_{,\bar{\phi}}}{\bar{\rho}_c}$$

- In the quasi-static limit ($k \gg \mathcal{H}$) this becomes:

$$\delta''_c + \mathcal{H} \delta'_c - 4\pi G a^2 [\bar{\rho}_c \delta_c + \Delta \rho_s(k, a) \delta s] \simeq k^2 Q_s(a) \delta s$$

- The coupling appears as “fifth-force” $Q_s(a)$ and “effective density source” contribution $\Delta \rho_s(k, a)$



Observational signatures

Pure derivative interaction with scale invariant entropy perturbations with standard DE exponential potential

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

$$\delta s(k) = \mathcal{A} \left(k/k_s \right)^n$$

- In the QS limit algebraic and derivative couplings reduce to similar effective source

$$k^2 \Psi \propto -a^2 (\delta \rho_m + \delta \rho_{\text{DE}})$$

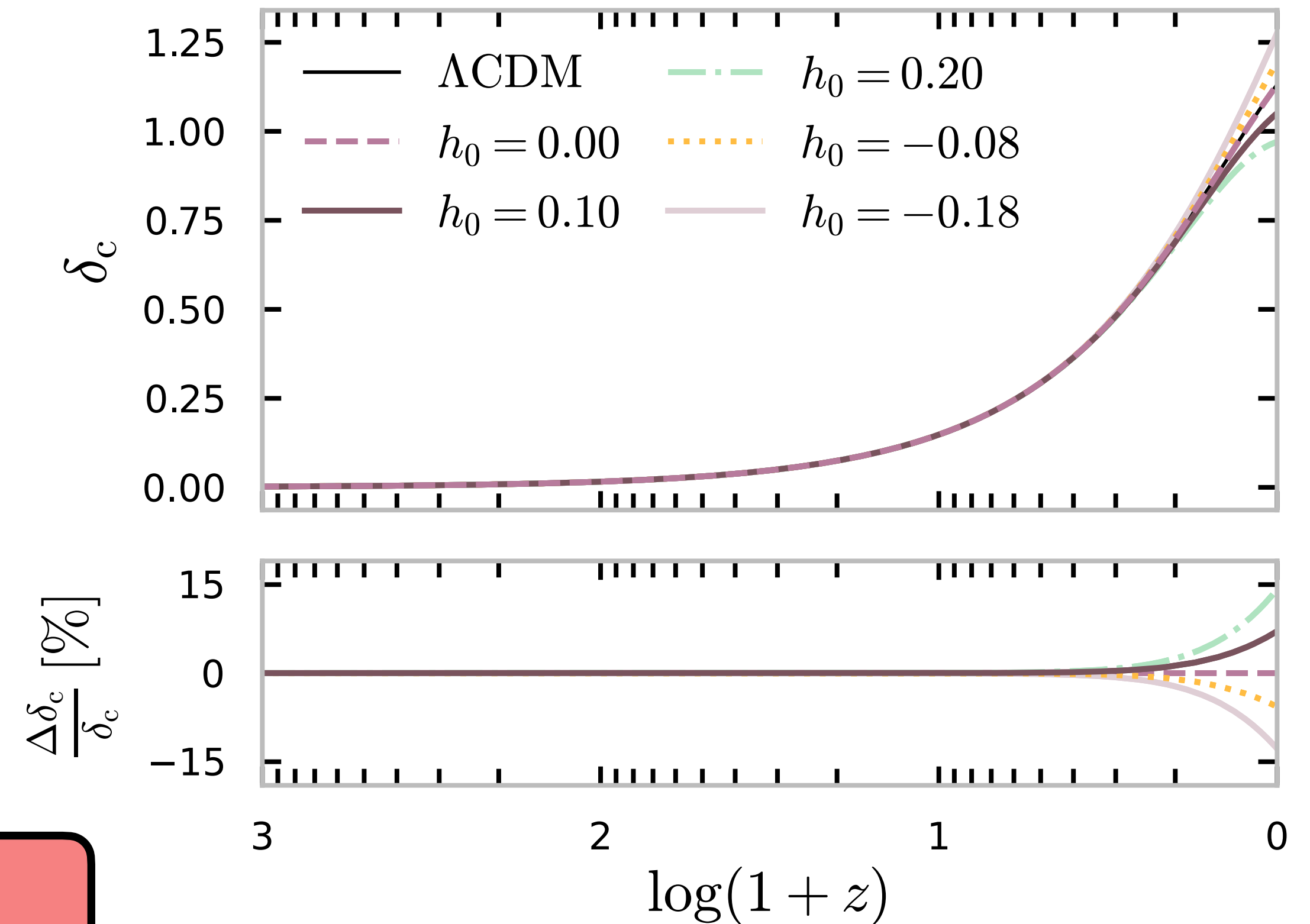
$$k^2 (\Psi' + \mathcal{H} \Psi) \propto a^2 (\theta_m + \theta_{\text{DE}})$$

- Modification in $\theta_{\text{DE}} = -k^2 (h_0 \delta s + \delta \phi) / \phi'$

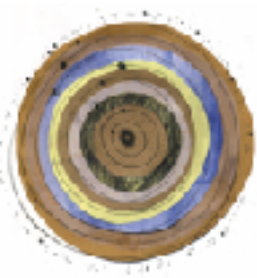
- If $n \neq 0$ need to tame large effects of the coupling at small scales

$$Q_s^{\text{der}}(a) = -\frac{h_0}{\bar{\rho}_c} \hat{V}_{,\bar{\phi}}$$

$$\Delta \rho_s^{\text{der}}(k, a) = -\frac{k^2 h_0}{k^2 + a^2 m_{\text{eff}}^2} \hat{V}_{,\bar{\phi}}$$



$$[k = 0.1h, \mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, n = 0]$$



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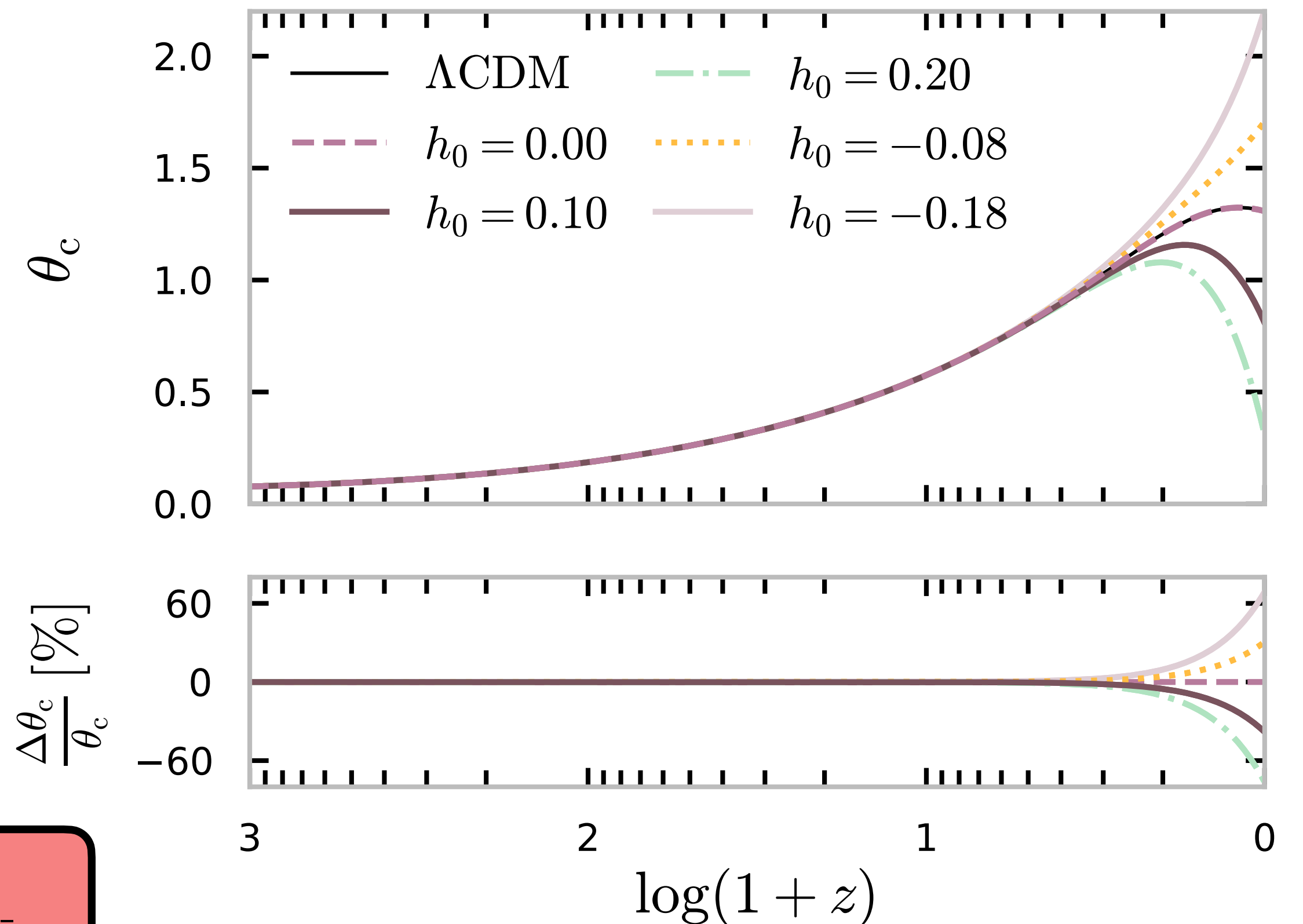
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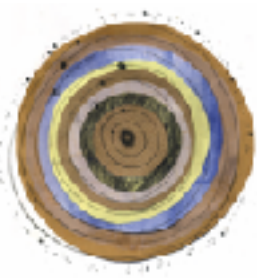


$$\log(1+z)$$

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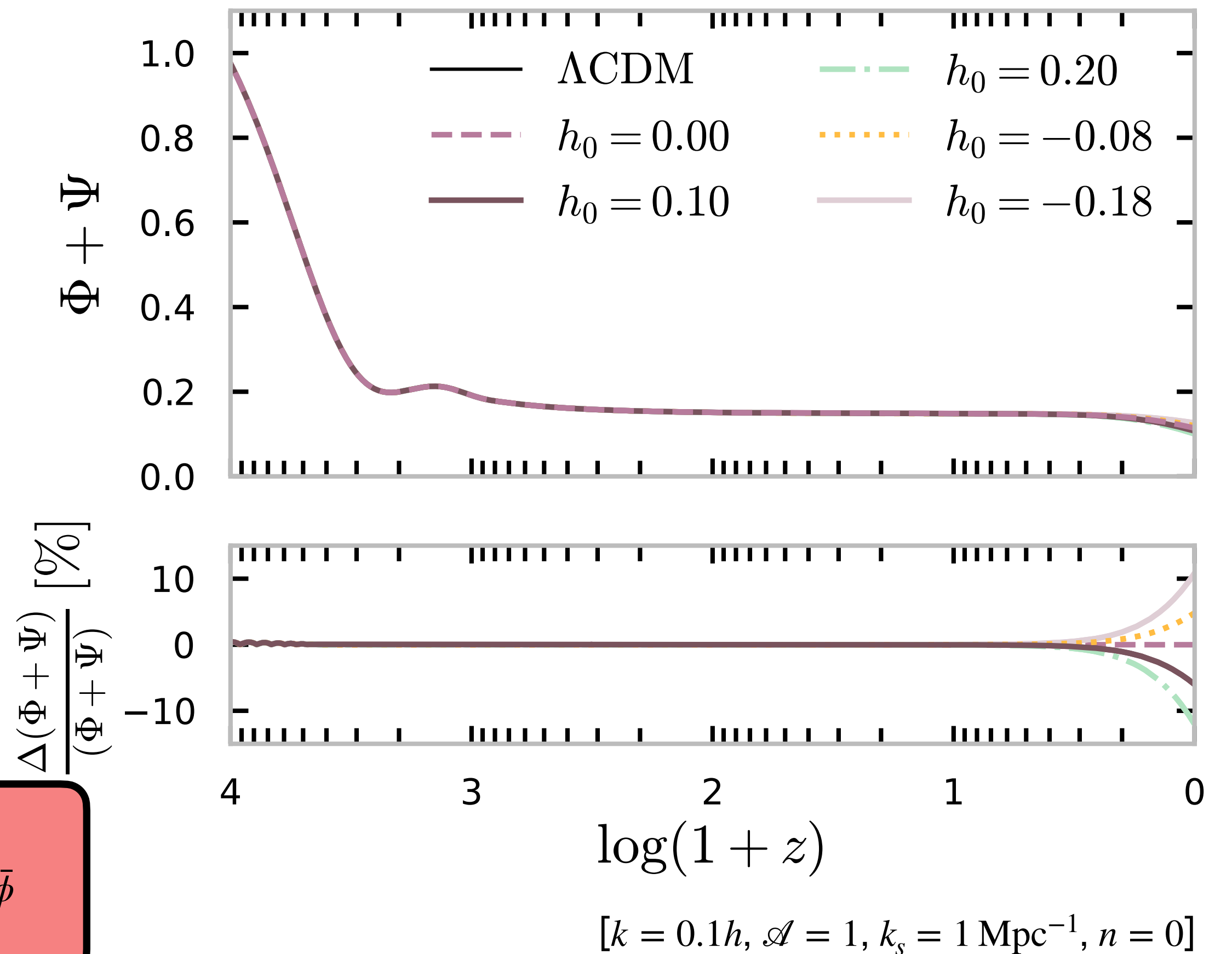
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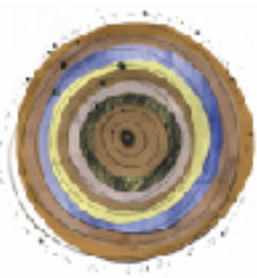
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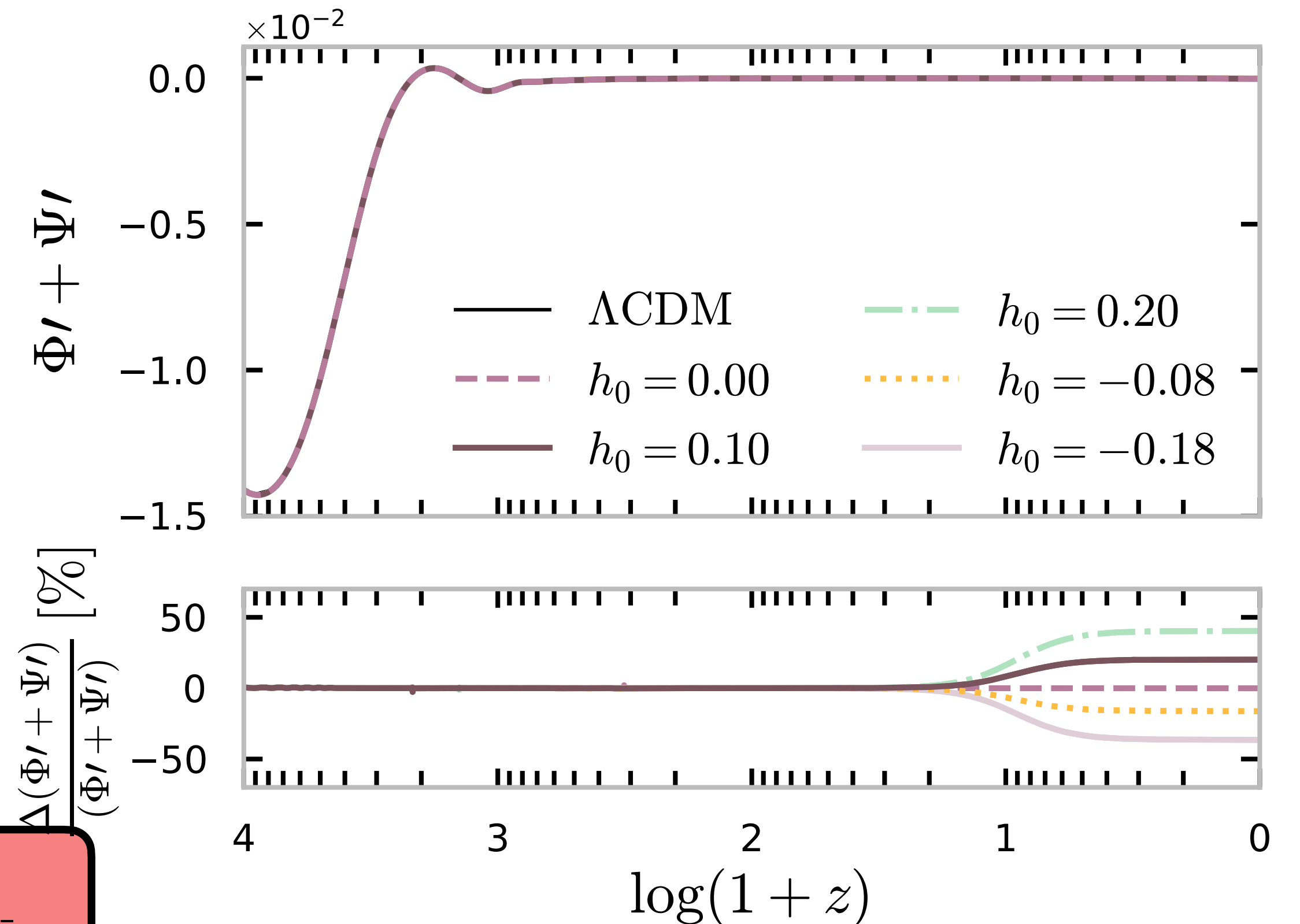
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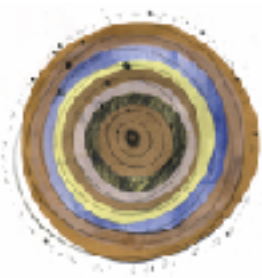
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$$[k = 0.1h, \mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, n = 0]$$

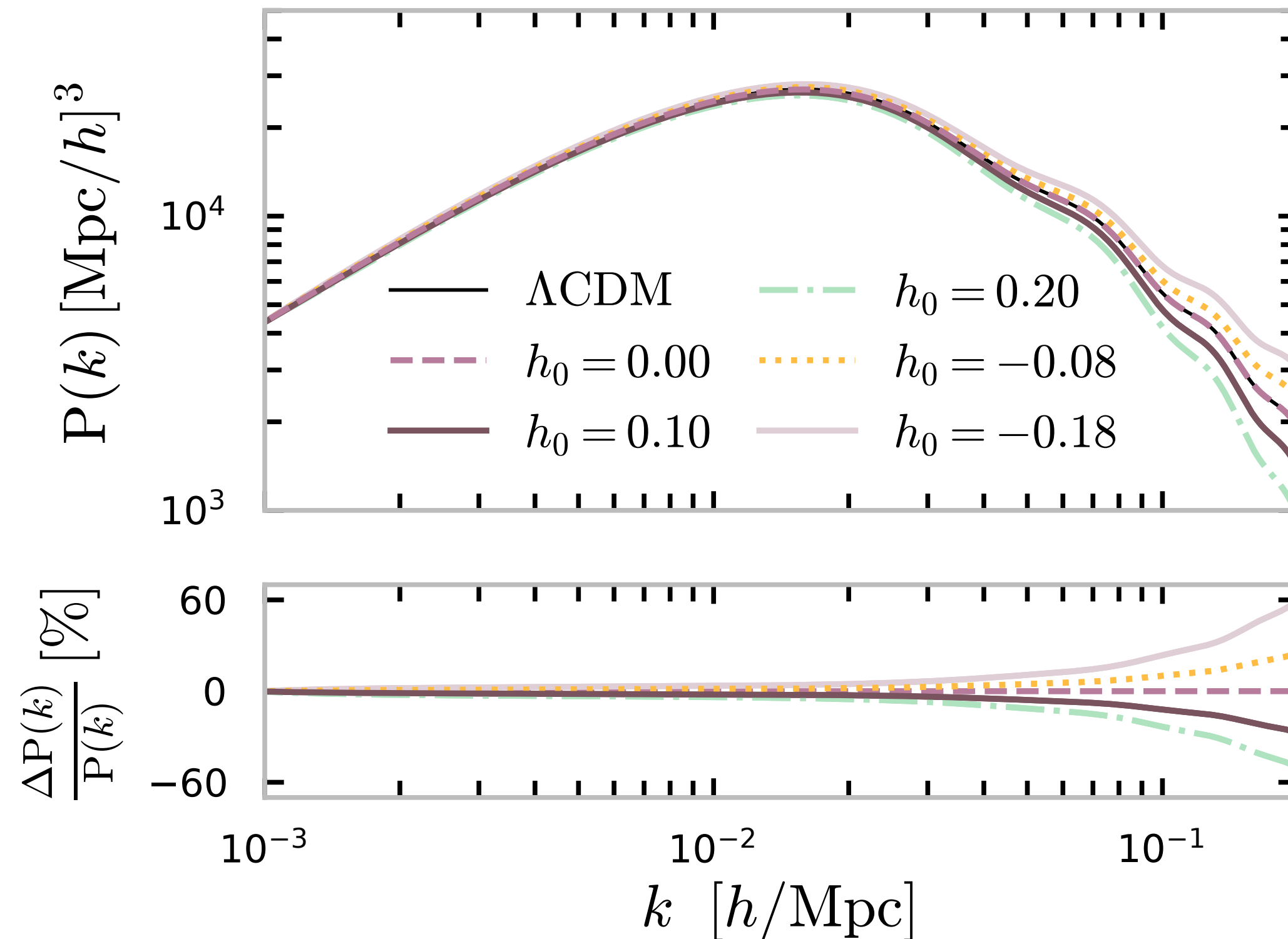


Observational signatures

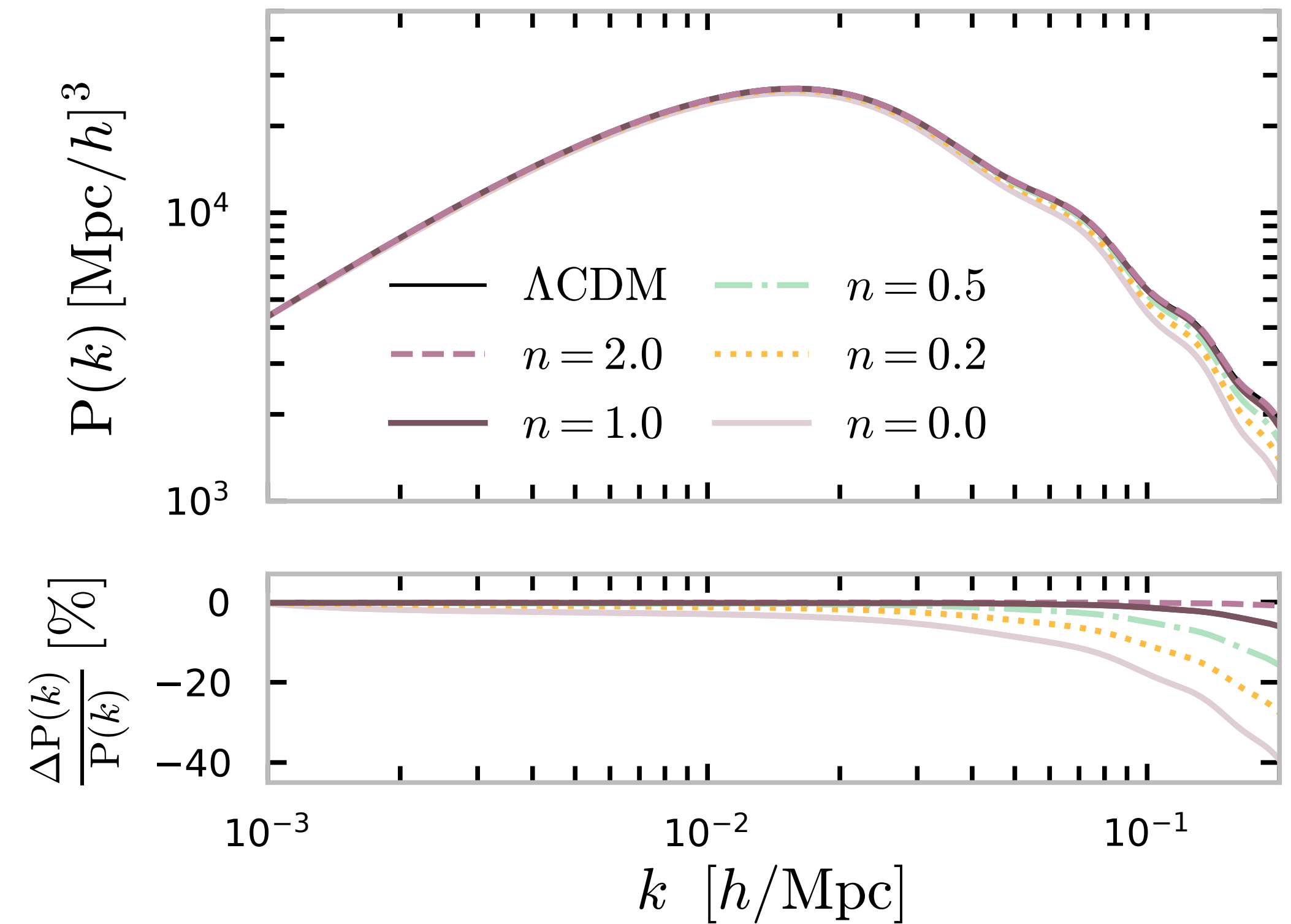
Suppression/enhancement of matter power spectrum on small scales

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

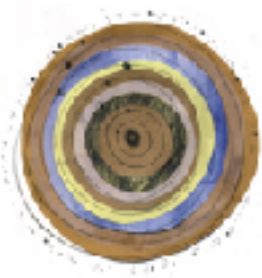
$$\delta s(k) = \mathcal{A} \left(k/k_s \right)^n$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, n = 0]$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, h_0 = 0.15]$$

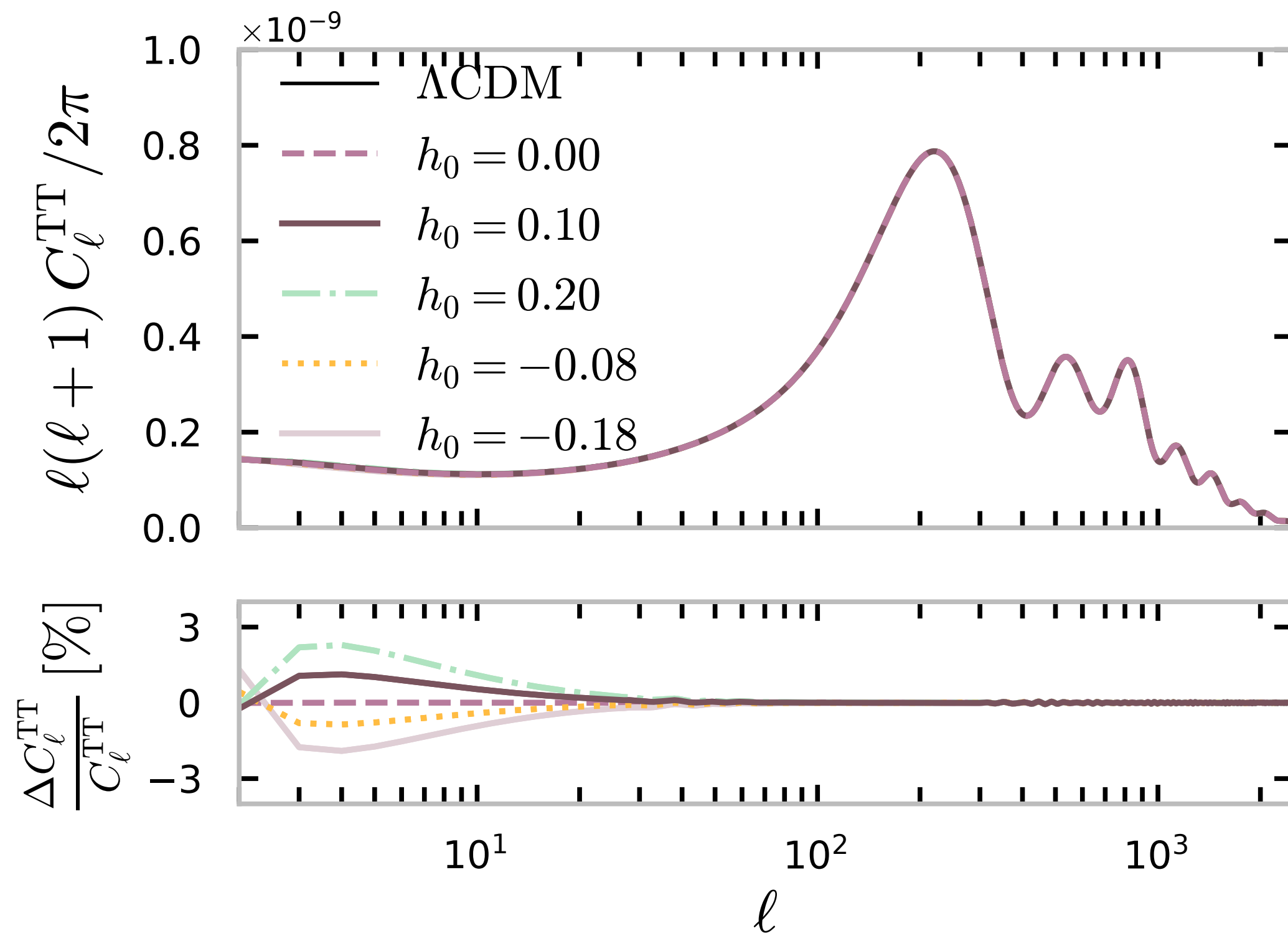


Observational signatures

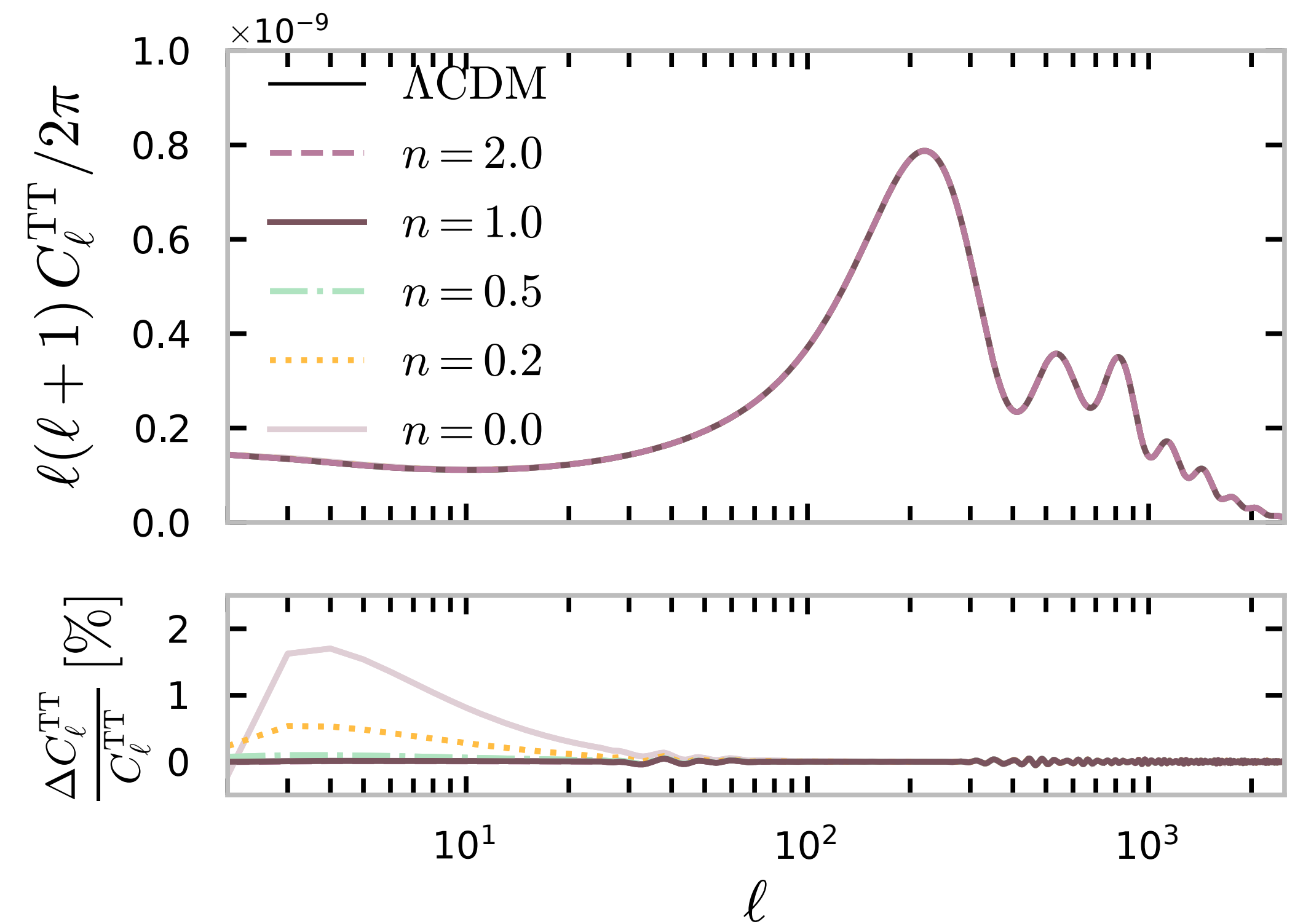
Tiny ISW changes to CMB temperature spectrum $\mathcal{C}_\ell^{TT}$ at low- ℓ

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

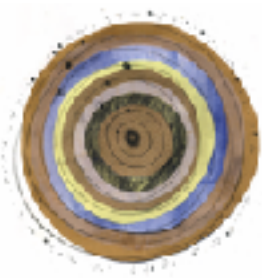
$$\delta s(k) = \mathcal{A} (k/k_s)^n$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, n = 0]$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, h_0 = 0.15]$$

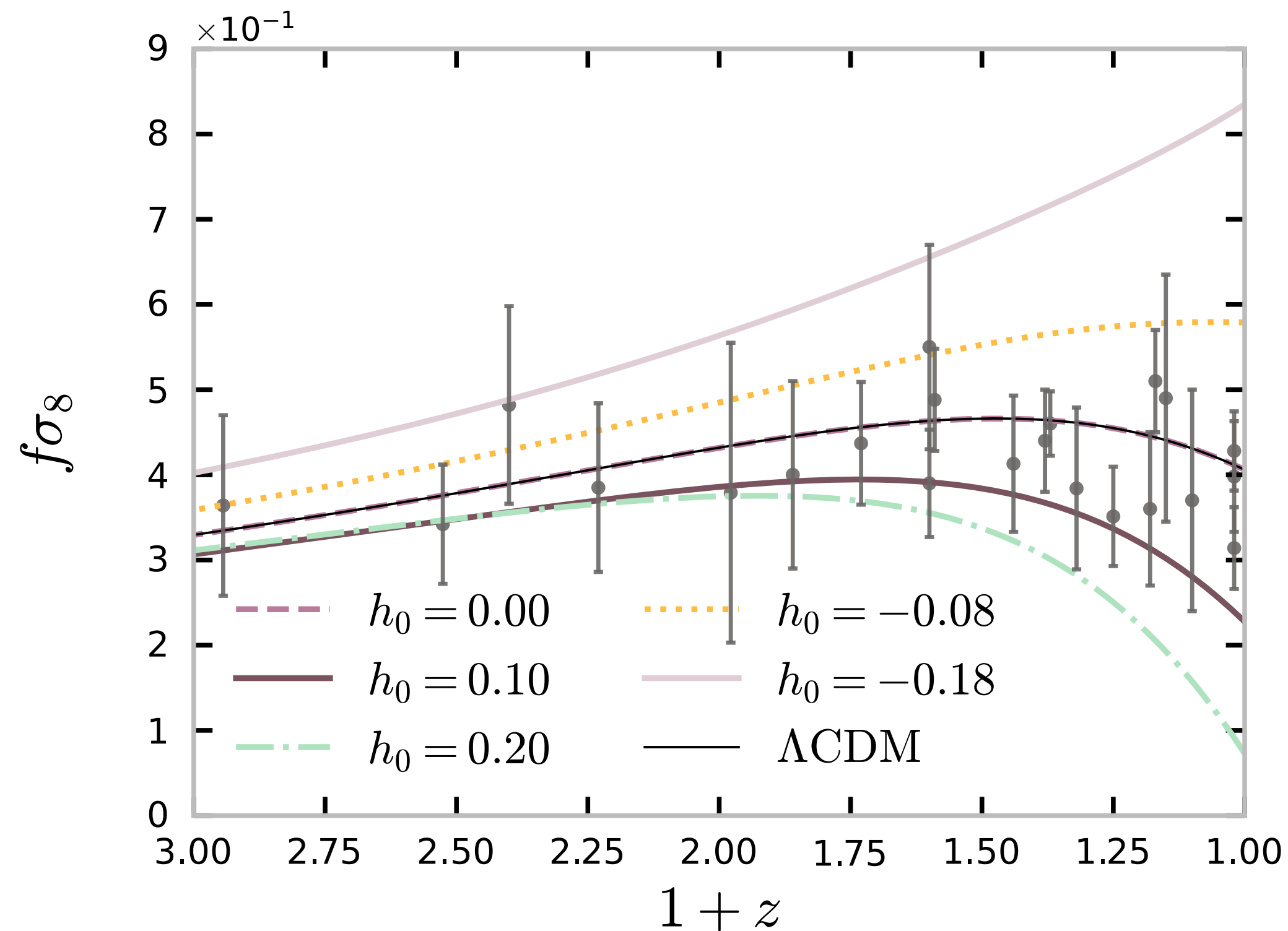


Observational signatures

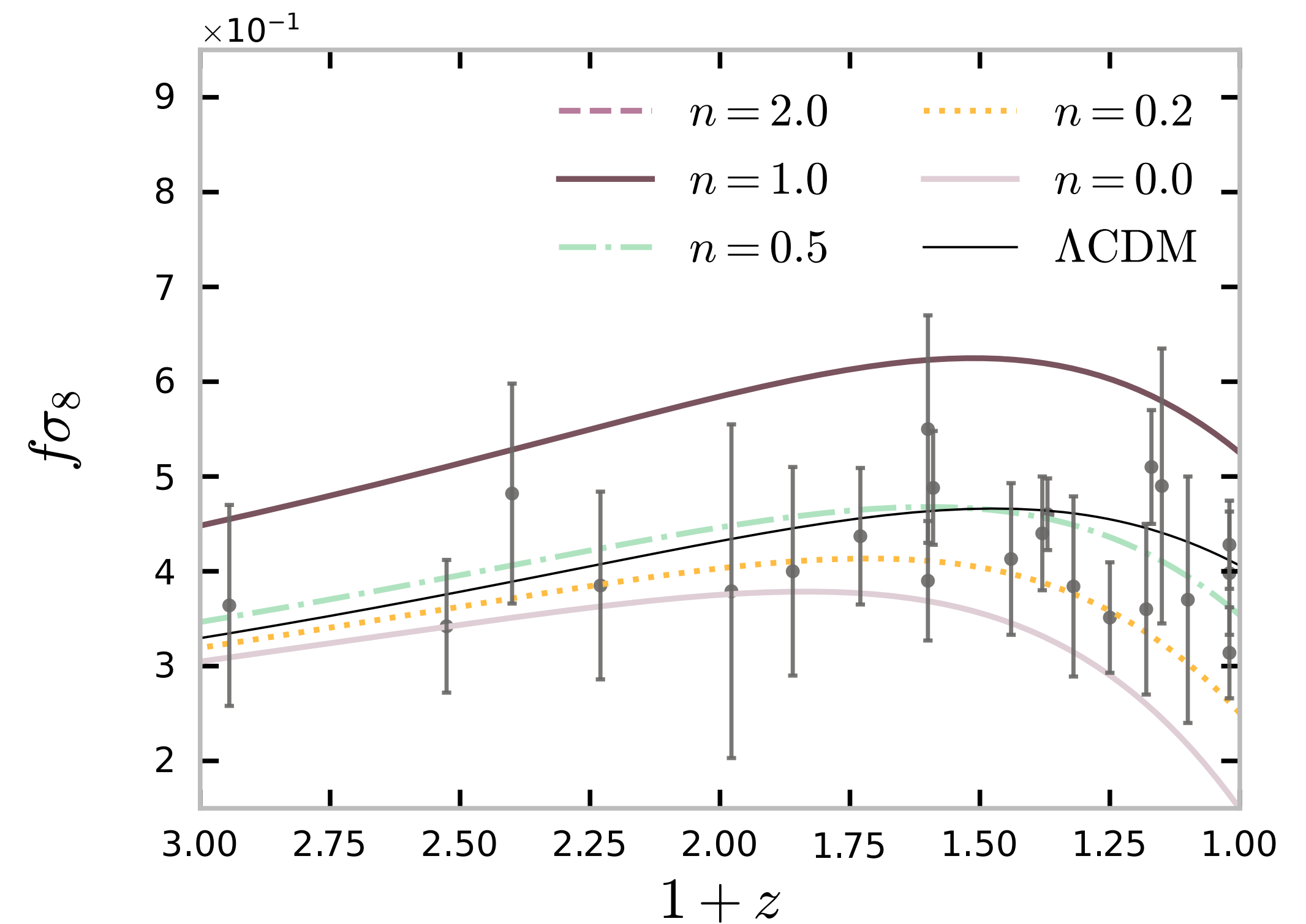
Significant changes to the growth - observational constraints

$$f(\phi, s, \mathcal{S}) = h_0 \mathcal{S} = h_0 \nabla_\mu \phi \nabla^\mu s$$

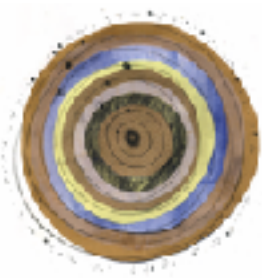
$$\delta s(k) = \mathcal{A} \left(k/k_s \right)^n$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, n = 0]$$



$$[\mathcal{A} = 1, k_s = 1 \text{ Mpc}^{-1}, h_0 = 0.15]$$



Observational signatures

Similar behaviour for algebraic models except for scale-dependent effects

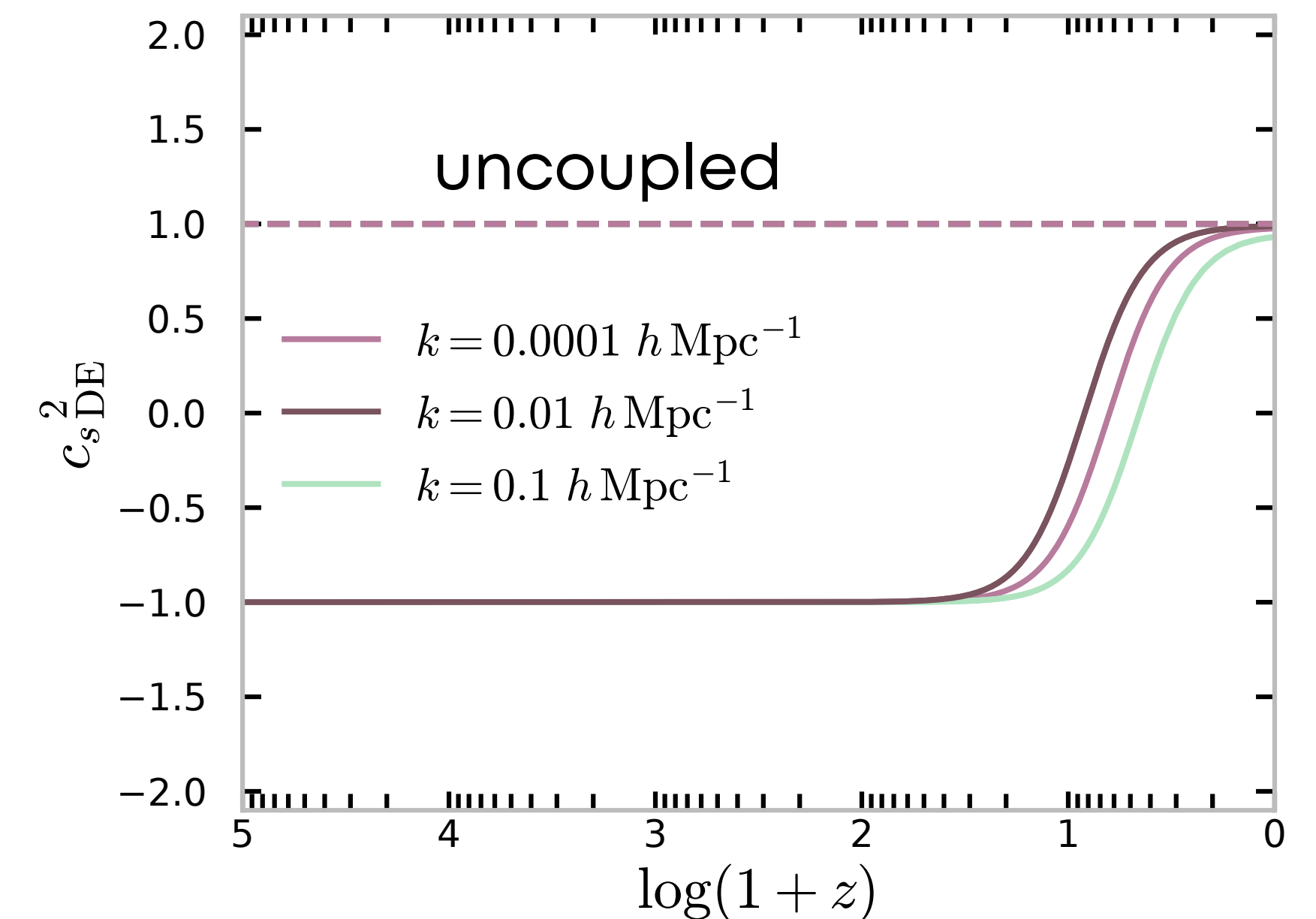
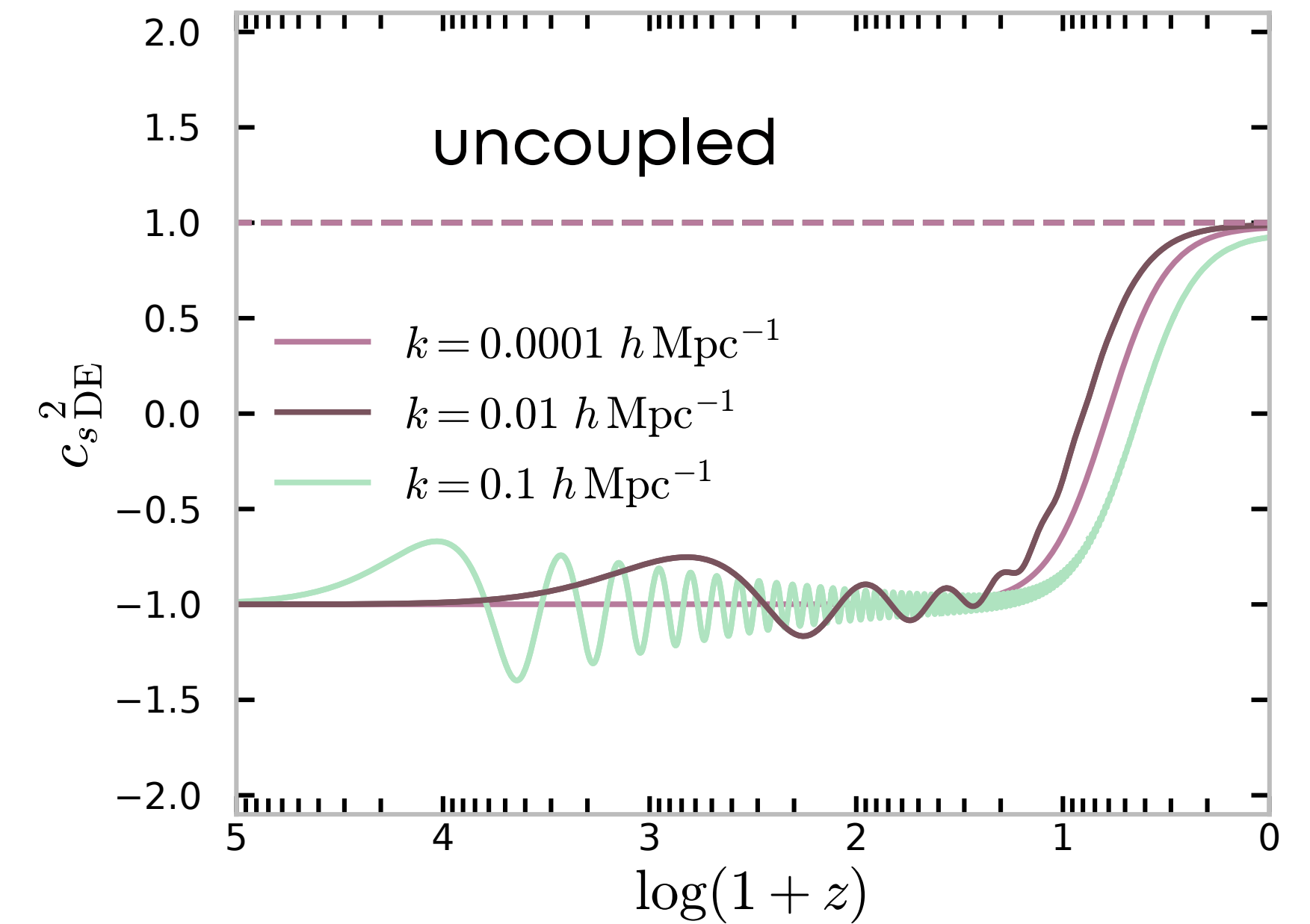
- Different scale-dependance of source term

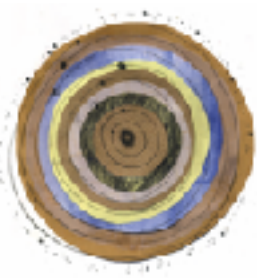
$$\Delta\rho_s^{\text{der}}(k, a) = -\frac{k^2 h_0}{k^2 + a^2 m_{\text{eff}}^2} \hat{V}_{,\bar{\phi}}$$

$$\Delta\rho_s^{\text{alg}}(k, a) = \bar{g}_{,s} - \frac{a^2 \bar{g}_{,s\phi}}{k^2 + a^2 m_{\text{eff}}^2} \hat{V}_{,\bar{\phi}}$$

- Different sound speed in the fluid rest frame ($c_s^2 \neq 1$):

$$c_s^2 := \left. \frac{\delta p}{\delta \rho} \right|_{\text{restframe}} = 1 - \frac{2a^2 \delta s (h_0 \hat{V}_{,\bar{\phi}} - \bar{g}_{,s})}{a^2 \delta s (h_0 \hat{V}_{,\bar{\phi}} - \bar{g}_{,s}) + \bar{\phi}' (\delta \phi' - \bar{\phi}' \Psi)}$$





Observational signatures

$$\delta\phi'' + 2\mathcal{H}\delta\phi' + (k^2 + a^2\hat{V}_{,\bar{\phi}\bar{\phi}})\delta\phi = -2a^2\Psi\hat{V}_{,\bar{\phi}} + 4\Psi'\bar{\phi}' - \left(a^2\bar{g}_{,s\phi} + \cancel{k^2 h_0}\right)\delta s$$

- Different scale-dependance of source term

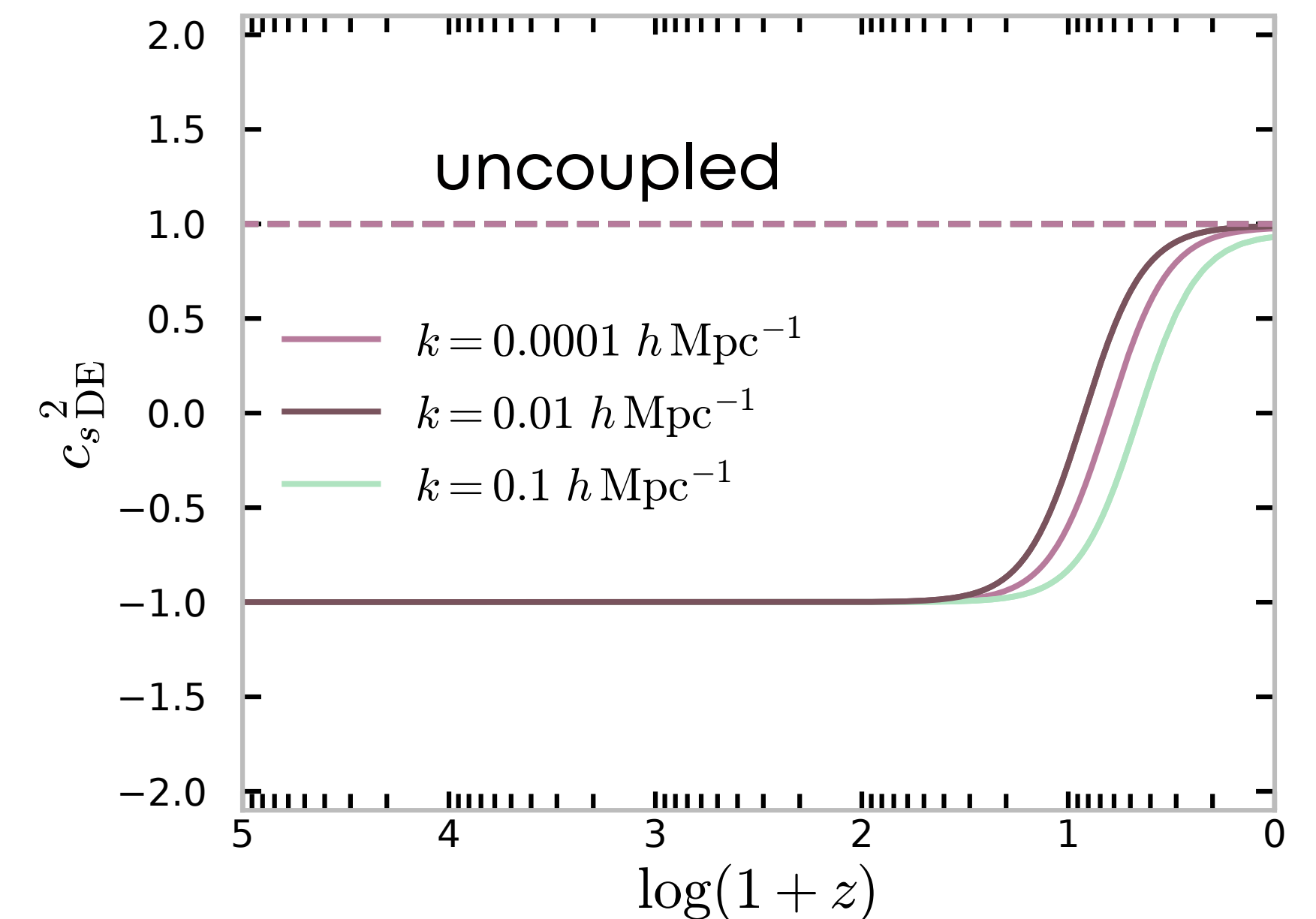
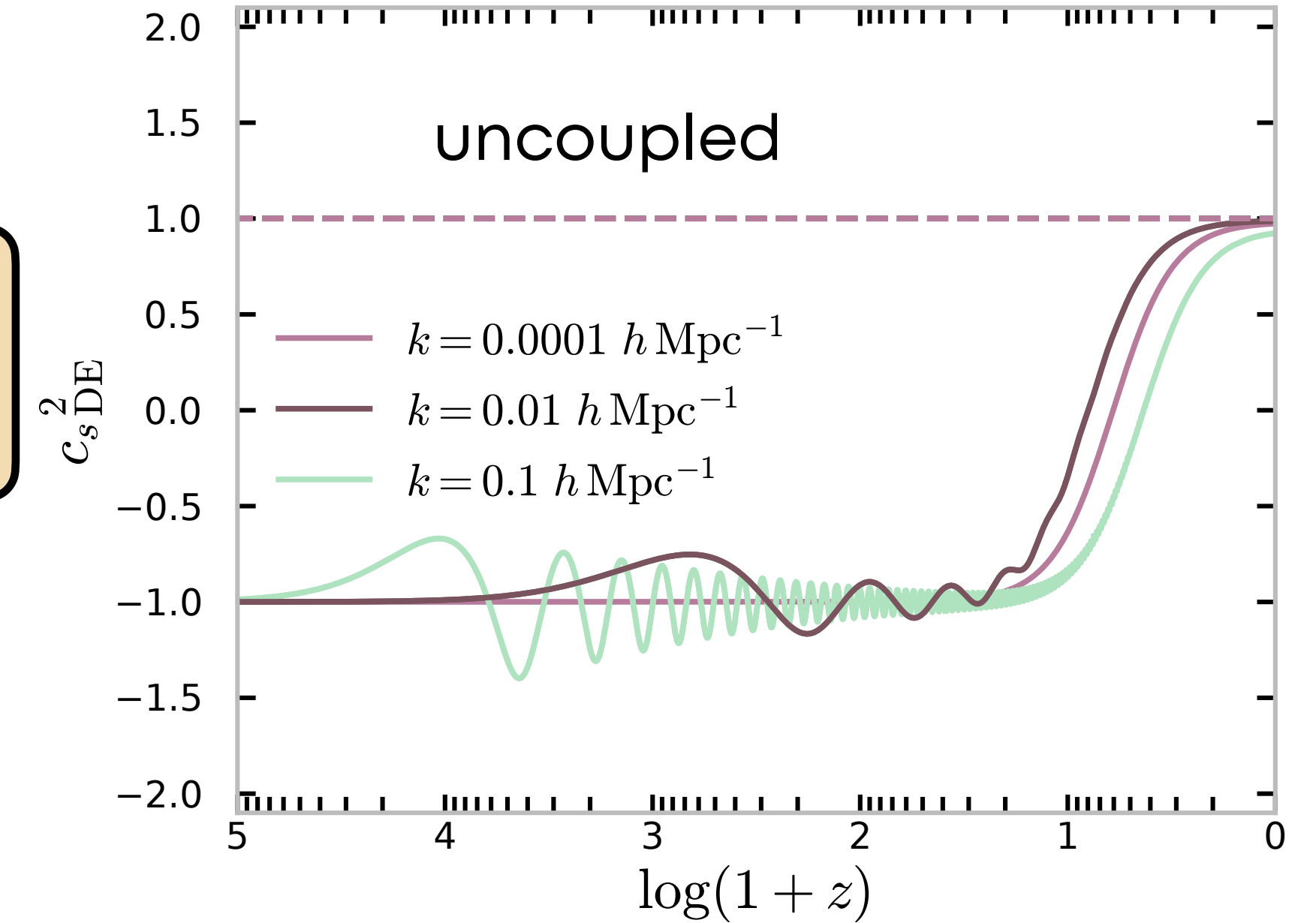
$$\Delta\rho_s^{\text{der}}(k, a) = -\frac{k^2 h_0}{k^2 + \cancel{a^2 m_{\text{eff}}^2}} \hat{V}_{,\bar{\phi}}$$

$$\Delta\rho_s^{\text{alg}}(k, a) = \cancel{\bar{g}_{,s}} \frac{a^2 \bar{g}_{,s\phi}}{k^2 + a^2 m_{\text{eff}}^2} \hat{V}_{,\bar{\phi}}$$

Oscillations in $\delta\phi$ in
the derivative
coupling for small
scales

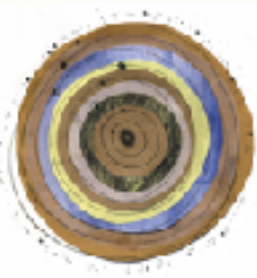
- Different sound speed in the fluid rest frame ($c_s^2 \neq 1$):

$$c_s^2 := \frac{\delta p}{\delta \rho} \Big|_{\text{restframe}} = 1 - \frac{2a^2 \delta s (\cancel{h_0 \hat{V}_{,\bar{\phi}} - \bar{g}_{,s}})}{a^2 \delta s (\cancel{h_0 \hat{V}_{,\bar{\phi}} - \bar{g}_{,s}}) + \bar{\phi}'(\delta\phi' - \bar{\phi}'\Psi)}$$



Summary and Conclusions





Conclusions

- New interacting scenarios from coupling DM intrinsic entropy to DE
 - Unchanged background expansion
 - No energy transfer (pure-momentum only)
 - Scale-dependent suppression/enhancement of matter clustering at late-times
-
- Observational constraints on allowed couplings - cosmic tensions
 - Explore different primordial power spectrum for $\delta_s(k)$ and initial conditions
 - Early dark energy/neutrino scenarios
 - Key observational signatures: scale-dependent clustering

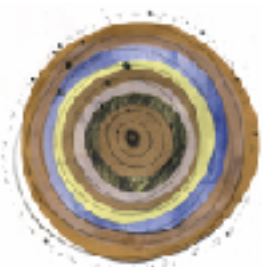


Thank you for your attention!

ELSA M. TEIXEIRA

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CNRS & Université de Montpellier
elsa.teixeira@umontpellier.fr

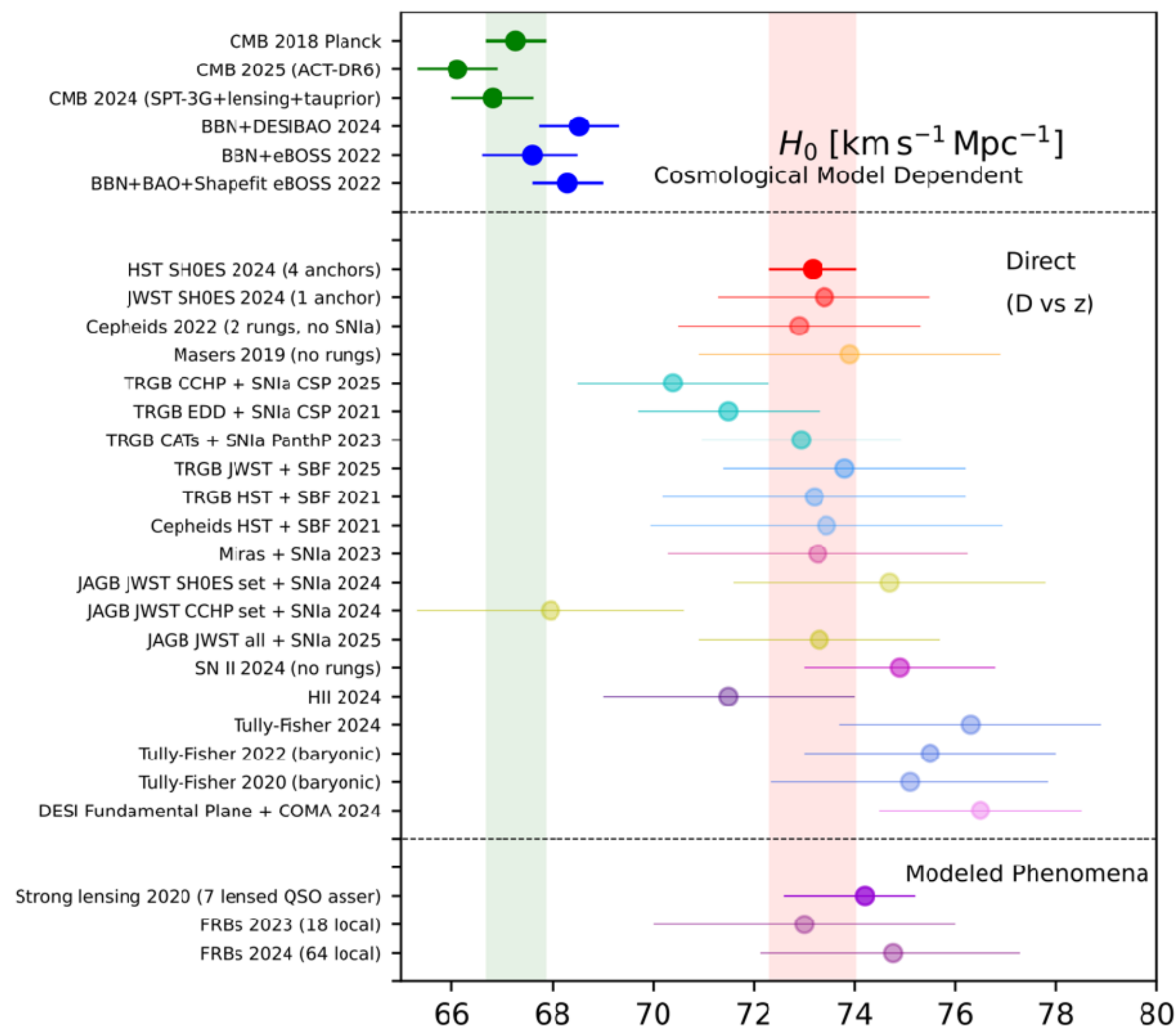
Illustrations: Inês Viegas Oliveira
(ivoliveira.com)

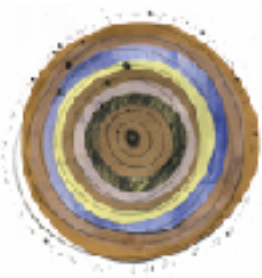


The Hubble Tension

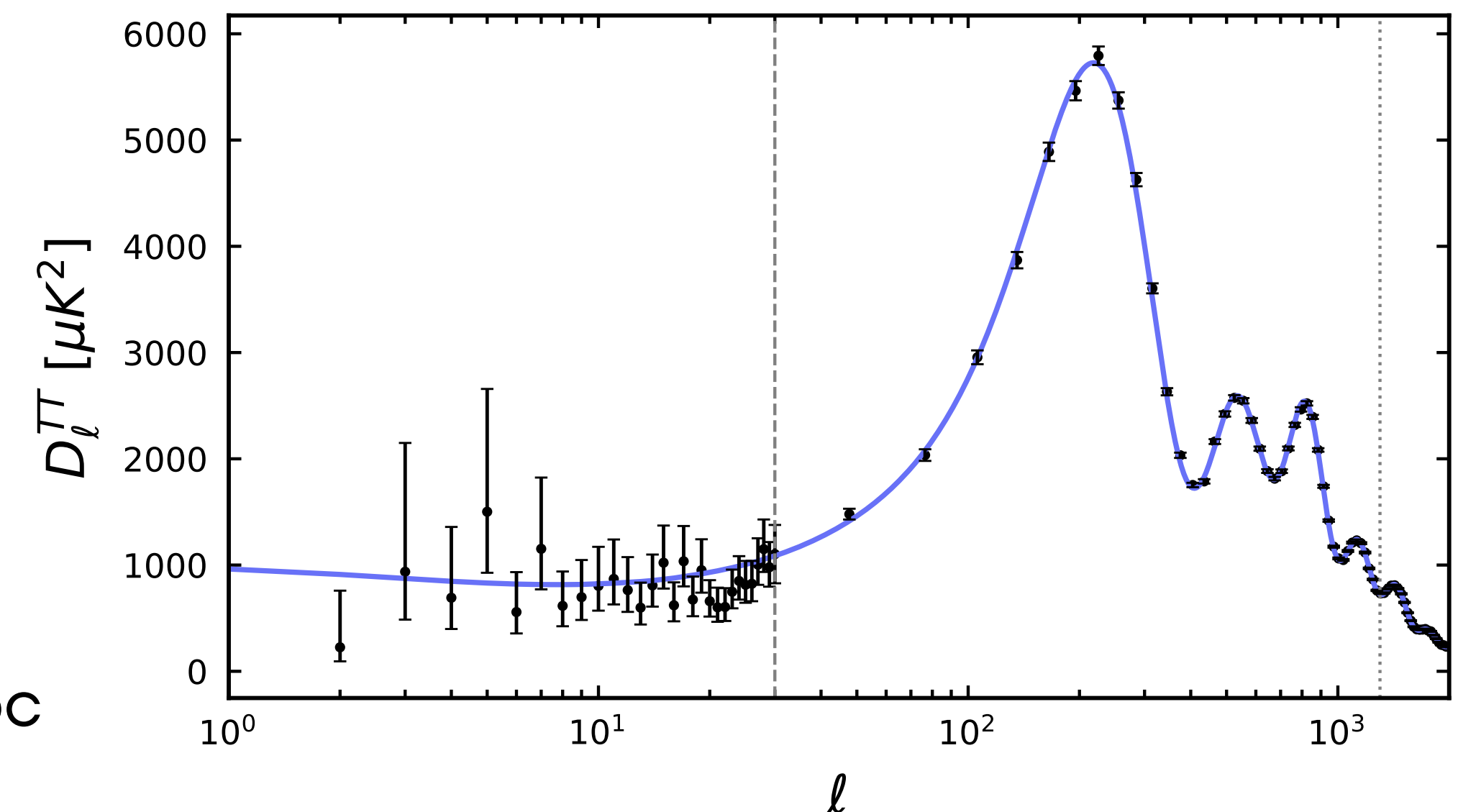
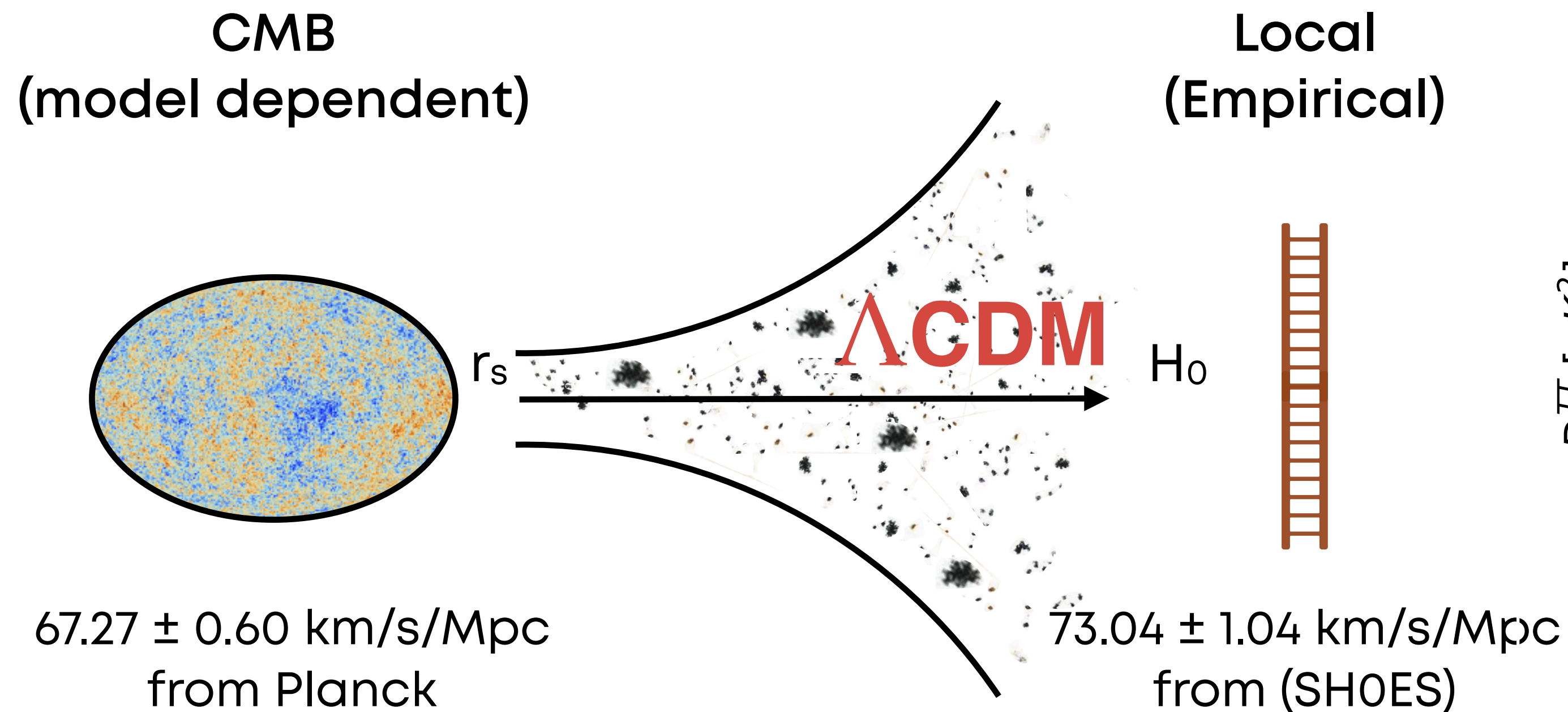
Unreconcilable values for H_0 from the CMB and from direct local distance ladder measurements

- $\sim 5\sigma$ tension between Planck 2018 and SH₀ES:
 - ▶ **CMB (Planck):** $H_0 = 67.27 \pm 0.60$ km/s/Mpc
 - ▶ **SNe (R22):** $H_0 = 73.04 \pm 1.04$ km/s/Mpc
- The CMB data assumes the Λ CDM model
- **DESI BAO (+BBN+CMB):** $H_0 = 68.45 \pm 0.47$ km/s/Mpc [DESI Collaboration DR2 2025: arXiv:2503.14738]
- Compilation of early vs late time data that disagree
- Could signal differences in the expansion history (nature of the dark sector)

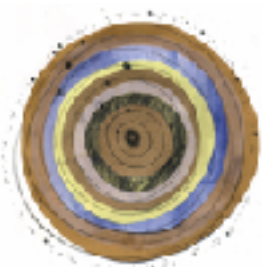




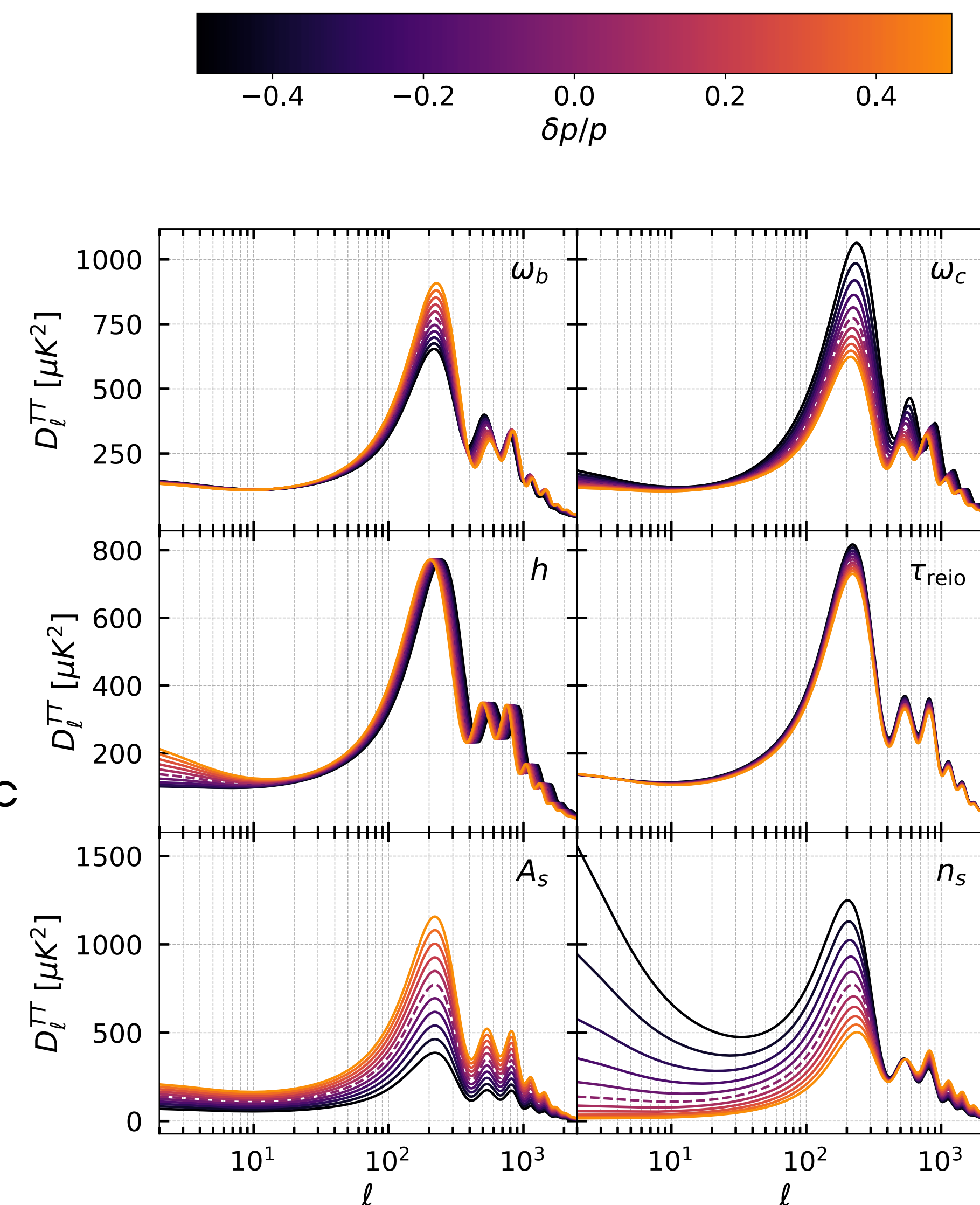
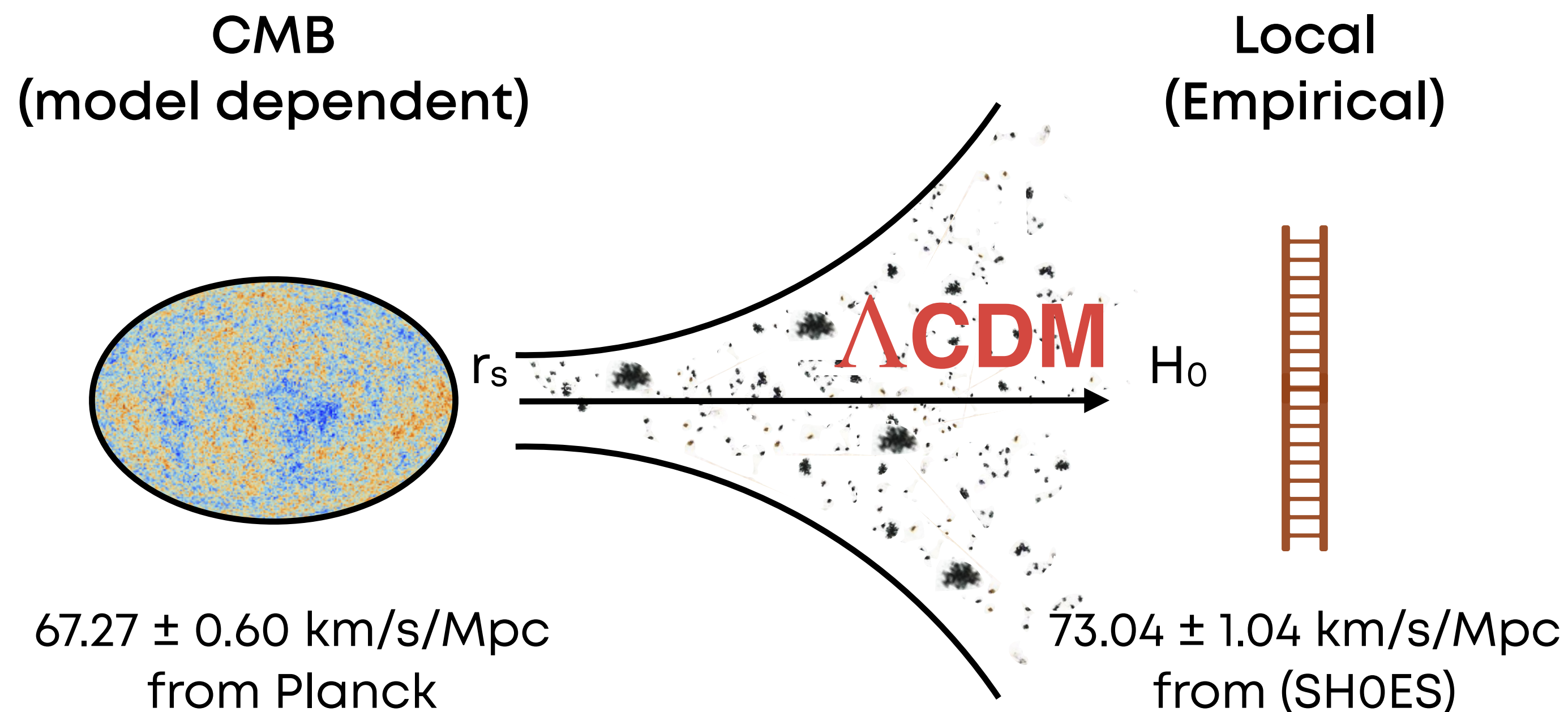
Cosmological Tensions

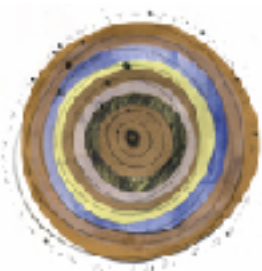


[Aghanim et al.: Astron.Astrophys. 641 (2020) A6]



Cosmological Tensions

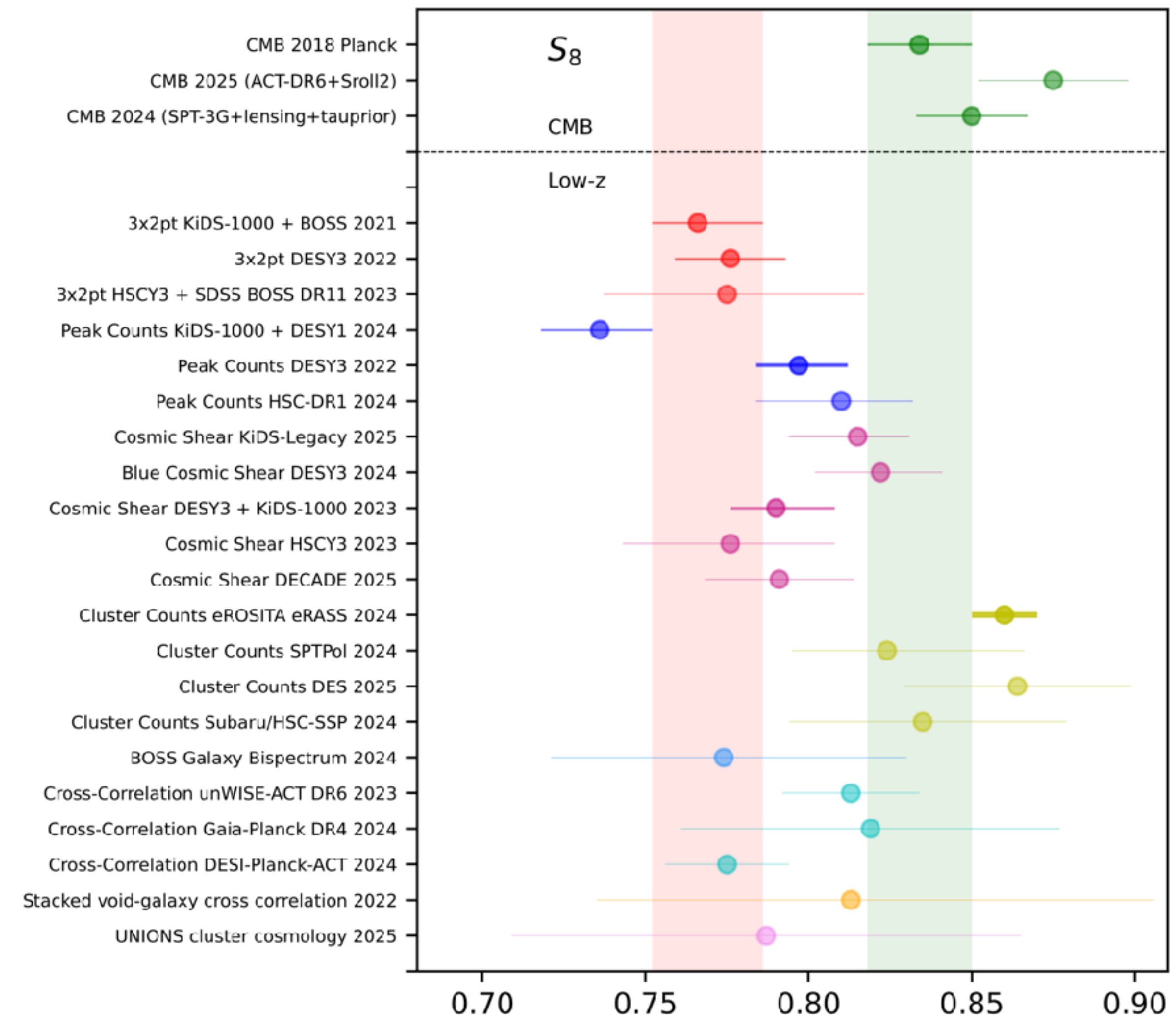


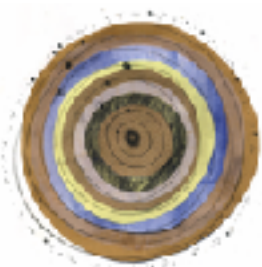


The S_8 Tension

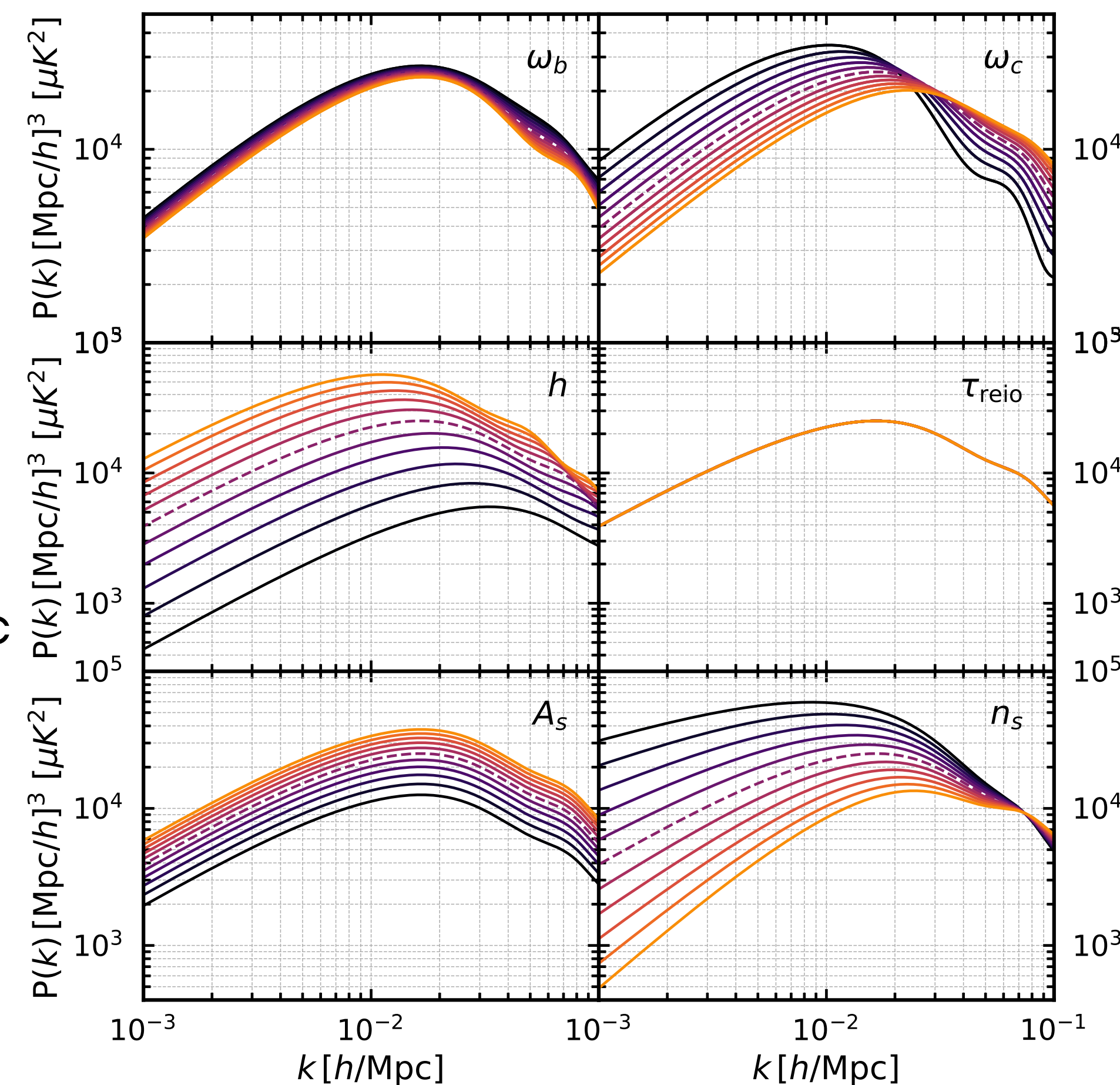
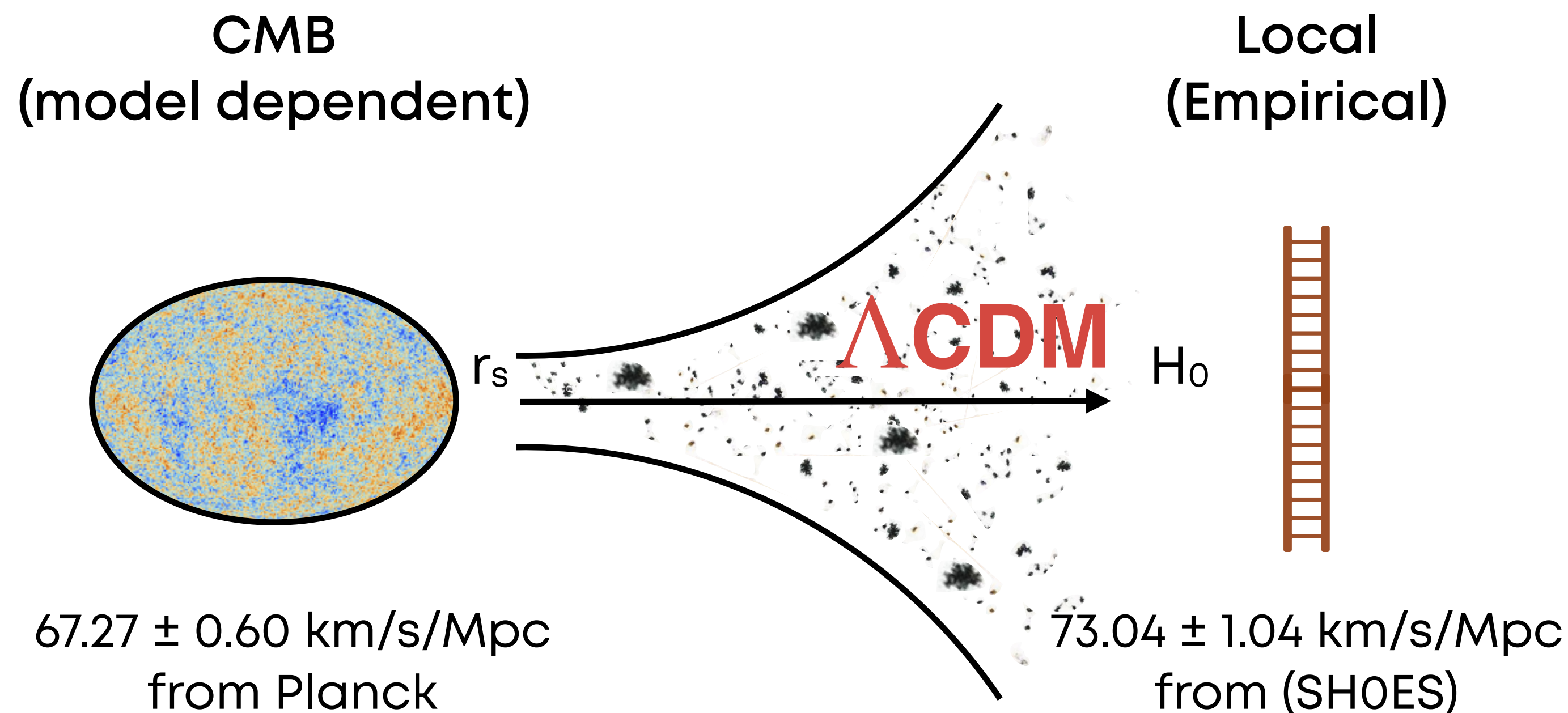
Discrepancy between CMB data and lensing surveys on combined quantity $S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$

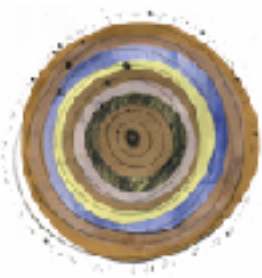
- ~ 3σ tension between Planck 2018 CMB data and KiDS-1000 combination of Cosmic Shear and Galaxy Clustering:
 - CMB (Planck 2018): $S_8 = 0.832 \pm 0.013$
 - Cosmic Shear (DES-Y3): $S_8 = 0.759^{+0.025}_{-0.023}$
- eRosita (eRASS1): $S_8 = 0.86 \pm 0.01$ [Ghirardini et al. 2024]
- Kids-Legacy ($S_8 = 0.815^{+0.016}_{-0.021}$) find possible resolution with Planck but not for the other measurements (improved redshift distribution estimation and calibration, as well as new survey area and improved image reduction)
- Could signal changes in clustering of matter (nature of CDM)





Cosmological Tensions





Conformal Transformation

- Simplest way to relate two geometries
- Rescaling of the metric that preserves angles
- Functions present in the metric and equations
- Map non-standard theories of gravity into GR plus a scalar field ϕ minimally coupled to the geometry
- Preserve the structure of Scalar-Tensor theories of the Jordan-Brans-Dicke form, such as $f(R)$

Non-Universal Coupling in the Dark Sector

$$\bar{g}_{\mu\nu} = C(\phi)g_{\mu\nu}$$

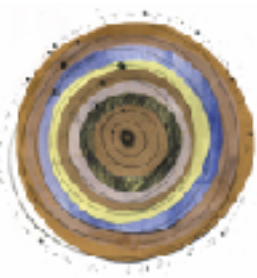
[Jordan: Z. Phys. 157 (1959), 112;
Brans and Dicke: Phys. Rev. 124 (1961), 925]

Disformal Transformation

- Distortion of both angles and lengths related with the gradient of ϕ
- The most general covariant effective metric
- The form of the non-derivative Lagrangian is preserved under disformal transformations
- Many cosmological applications

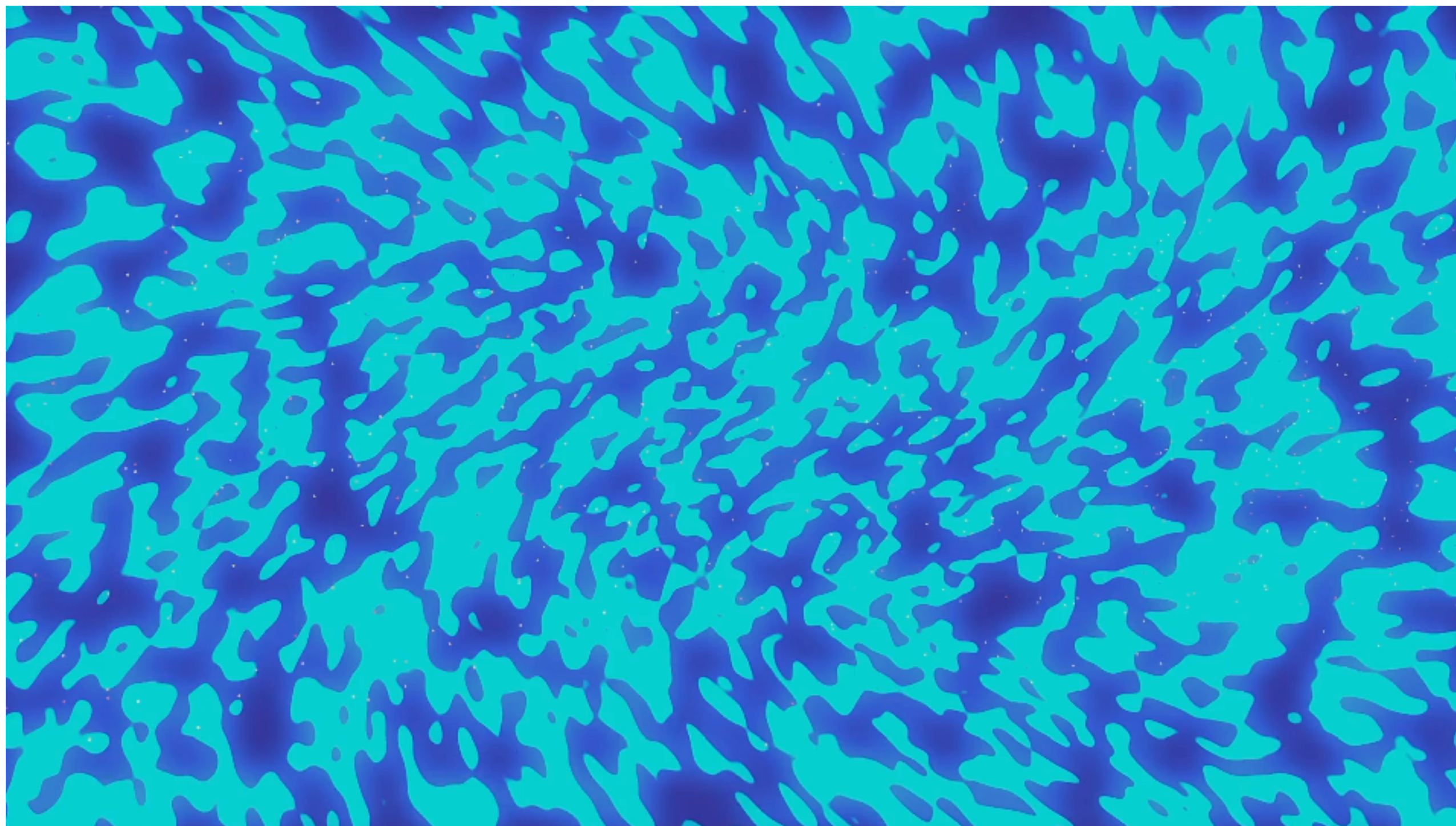
$$\bar{g}_{\mu\nu} = C(\phi)g_{\mu\nu} + D(\phi)\partial^\mu\phi\partial_\mu\phi$$

[Bettoni and Liberati: Phys. Rev. D88 (2013) 084020]

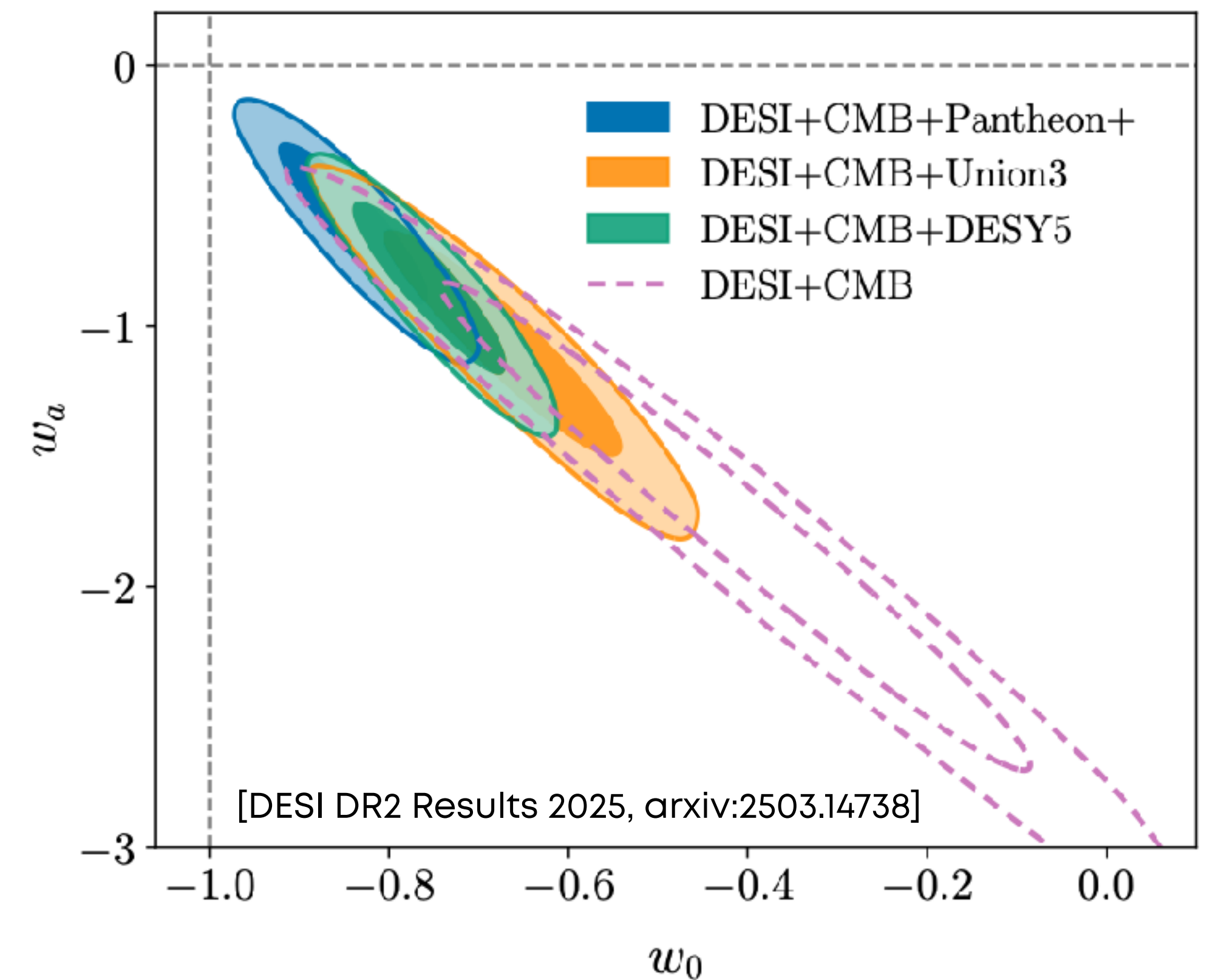


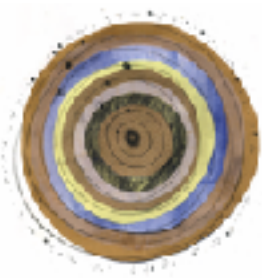
Extensions to Λ CDM

Hints of **dynamical DE** in **DESI BAO (baryonic acoustic oscillations) data** when combined with CMB and SN data



$$\begin{cases} \dot{\rho}_{\text{DE}} + 3H\rho_{\text{DE}}(1 + w_{\text{DE}}) = 0, \\ w_{\text{DE}} = w_0 + w_a(1 - a) \end{cases}$$





Extensions to Λ CDM

The observational tensions hint at **missing ingredients** or need for completely **new physics**

- “**Quintessence**” (ϕ) - dynamical scalar field that evolves in space and time, as opposed to Λ
- More physically motivated than a parametric fluid
- No fundamental principle/observational constraints which forbid **interactions between the dark species**
- Modified predictions for the evolution of the dark sector could naturally address the cosmic tensions

$$\begin{cases} \dot{\rho}_\phi + 3H\rho_\phi(1 + w_\phi) = 0, \\ \rho_\phi = \frac{1}{2}\dot{\phi}^2 + V(\phi), \quad p_\phi = \frac{1}{2}\dot{\phi}^2 - V(\phi) \end{cases}$$

