









Search for high frequency gravitational waves in electromagnetic cavities

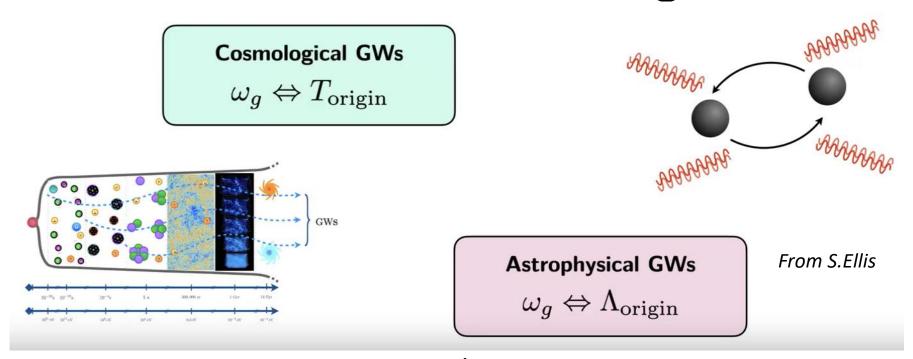
Jordan Gué

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In collab. with T. Krokotsch (Universität Hamburg)

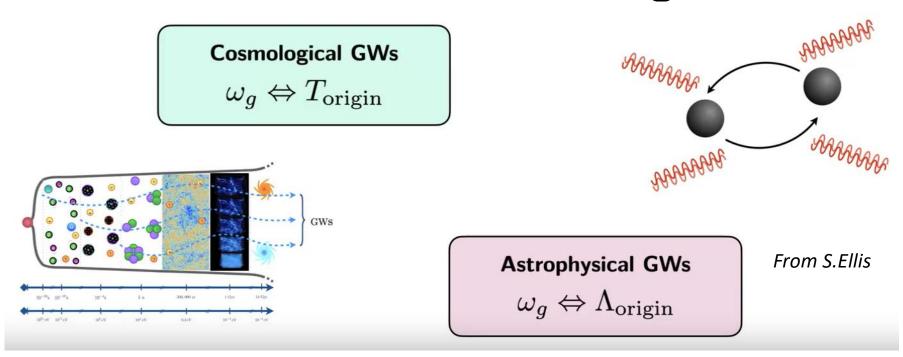


HFGW astro/cosmo signals



→ Higher GW frequency ⇔ Higher energy scale/Lower length scale we can probe

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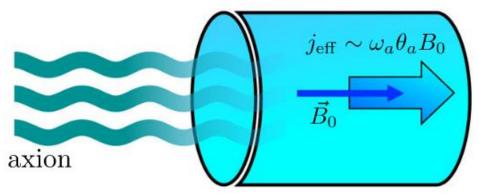
Some astrophysical sources:

- Mergers of PBH
- Mergers of ECO (boson stars,...)
- First order phase transition in neutron stars
- Superradiant boson clouds orbiting SMBH (monochromatic)

See N. Aggarwal et al, arXiv 2501.11723

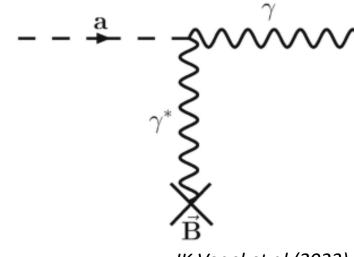
→ Many of those sources are BSM

Analogies with axion dark matter



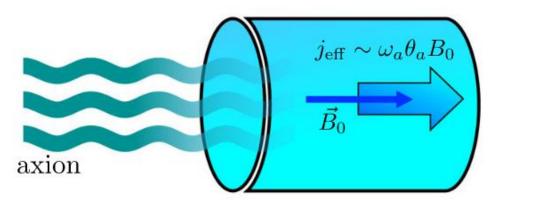
A. Berlin et al, PRD **105** (2022)

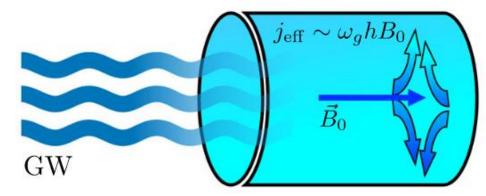
- Simple way of looking for axions coupled to EM is through inverse Primakoff effect
- → Use of microwave cavities to search for GHz axions



JK Vogel et al (2023)

Analogies with axion dark matter

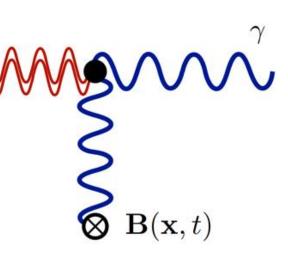


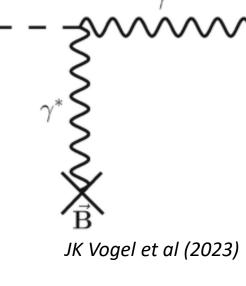


A. Berlin et al, PRD **105** (2022)

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- → Use of microwave cavities to search for GHz axions

- GW analog: inverse Gertsenshtein effect
- → Same apparatus is sensitive to HFGW





Credit: S. Ellis

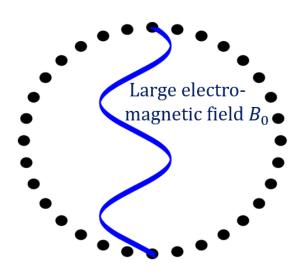
Expected GW signals

•
$$S = \int d^4x \sqrt{-g} \left(-\frac{1}{4} g^{\mu\alpha} g^{\nu\beta} F_{\mu\nu} F_{\alpha\beta} \right) \rightarrow \partial_{\nu} \delta F^{\mu\nu} \equiv j_{\rm eff}^{\mu} \propto \omega_g h B_0$$

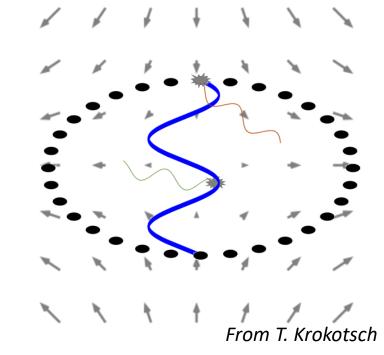
→ GW couples to EM energy

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$$\delta \ddot{x}_i - \partial_j \sigma_{ij} = F_i^h$$
 M. Hudelist et al, CQG **40** (2023)

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A. Berlin et al, PRD **105** (2022) V. Domcke et al, PRL **129** (2022)

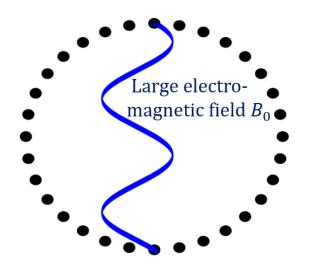


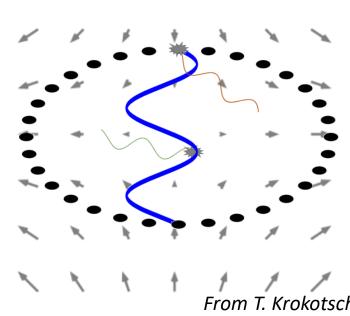
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How do these quantities enter the observable electric field?

$$E_{\underline{a}} = F_{\mu\nu}e^{\mu}_{\underline{a}}u^{\nu} \rightarrow \delta E_{\underline{a}} = \delta F_{\mu\nu}e^{\mu}_{\underline{a}}u^{\nu} + F_{\mu\nu}\delta e^{\mu}_{\underline{a}}u^{\nu} + F_{\mu\nu}e^{\mu}_{\underline{a}}\delta u^{\nu} + \delta x^{\rho}(\partial_{\rho}F_{\mu\nu})e^{\mu}_{\underline{a}}u^{\nu}$$

$$W. Ratzinger et al. JHEP 2024 (2024)$$

Observer infinitesimal coord. System = Tetrad $\frac{\mu}{2} = \frac{\nu}{2} = \frac{\nu}{2}$

$$g_{\mu\nu}e^{\mu}_{\underline{a}}e^{\nu}_{\underline{b}}=\eta_{\underline{a}\underline{b}}$$
 ; $e^{\nu}_{\underline{0}}=u^{\nu}$

 $\rightarrow \delta F^{0a}$ is **not** the observed field when dealing with GW

Frames

In general, to compute GW signals, choice between 2 frames: TT and PD

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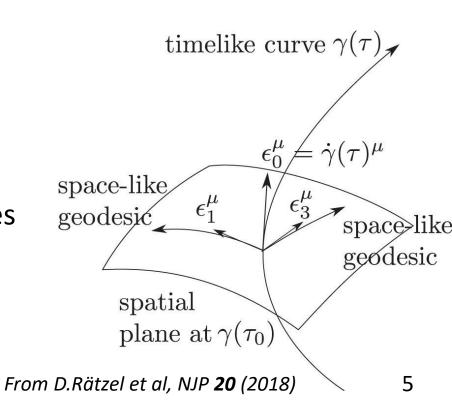
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- Traceless-Transverse (TT) gauge : global coord. system set by freely falling test masses $h_{0\mu}=\partial_i h^{ij}=h=0$
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- Proper Detector (PD) Frame: coordinate system built by extending observer's tetrads into geodesics
- → More intuitive : GW acts as a Newtonian force on rigid bodies
- → Metric perturbation more involved



Which frame should we use?

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What about haloscopes?

In conductor, $v_{\rm S}\sim 10^{-5}$, i.e for a cavity with $L\sim 0.1$ m, and $\omega_g\sim {\rm GHz},~\omega_g L/v_{\rm S}\gg 1$

→ TT more convenient

$$\delta E_{\underline{a}}^{TT,FF} = \delta F_{a0} + F_{\mu 0} \delta e_{\underline{a}}^{\mu} + F_{a\nu} \delta u^{\nu} + \delta x^{\rho} (\partial_{\rho} F_{a0}) = \delta F_{a0}$$

Signal power

In TT, at high frequency, one can solve $\partial_{\nu}\delta F^{\mu\nu}=j^{\mu}_{\rm eff}$ by expanding δF_{a0} in cavity eigenmodes.

On resonance, the signal power in a mode
$$\vec{E}_n$$
 is given by
$$P_{sig} = \frac{1}{2} Q \omega_g V \eta_g^2 h^2 \left| \vec{B} \right|^2_{A.~Berlin~et~al,~PRD~105~(2022)}$$

with the coupling

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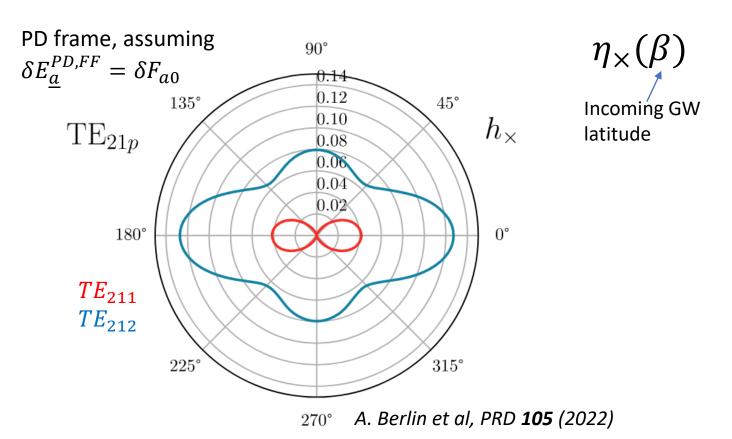
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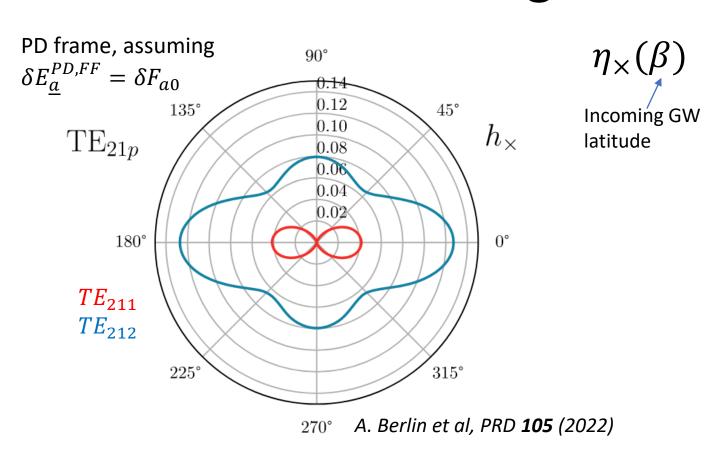
The signal power from axion DM is $P_{sig}^a = \frac{1}{2} g_{a\gamma}^2 Q \omega_a V \eta_a^2 a^2 \left| \vec{B} \right|^2$ with $\eta_a = \frac{|\text{J } av E_n.B|}{\sqrt{V \int dV \left| \vec{E}_n \right|^2}}$

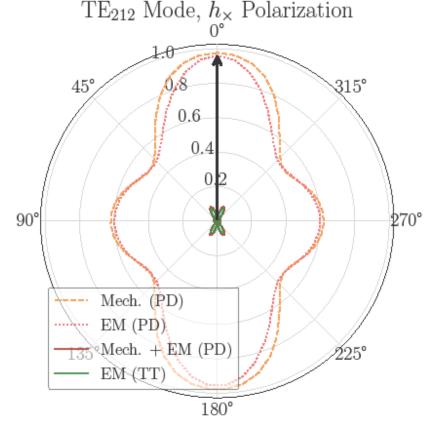
P. Sikivie, RMP 93 (2021)

Gauge invariance



Gauge invariance





T. Krokotsch, JG, in preparation

To get correct result in PD, need to take into account $F_{aj}\delta u^j$ as a boundary term

→ Signal is gauge invariant as expected.

- EM signal $\delta E_{\underline{a}} = \delta F_{a0} + F_{\mu 0} \delta e^{\mu}_{\underline{a}} + F_{a\nu} \delta u^{\nu} + \delta x^{\rho} (\partial_{\rho} F_{a0})$
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- Heterodyne/Homodyne setups

Conclusion

• Microwave cavities are powerful probes of monochromatic HFGW ($h \sim 10^{-22}$)

GW couples to all types of energy,
 care must be taken to model all effects

- With current quantum technology, this is not enough to probe cosmological GW
- → Cross correlate multiple cavities : GravNet ERC
- → Use of Earth modulation for persistent signals
- → Quantum enhancement techniques (e.g. squeezing)

