

ATLAS Results on Higgs Physics and Electroweak Symmetry Breaking

Gustaaf Brooijmans
on behalf of the ATLAS Collaboration

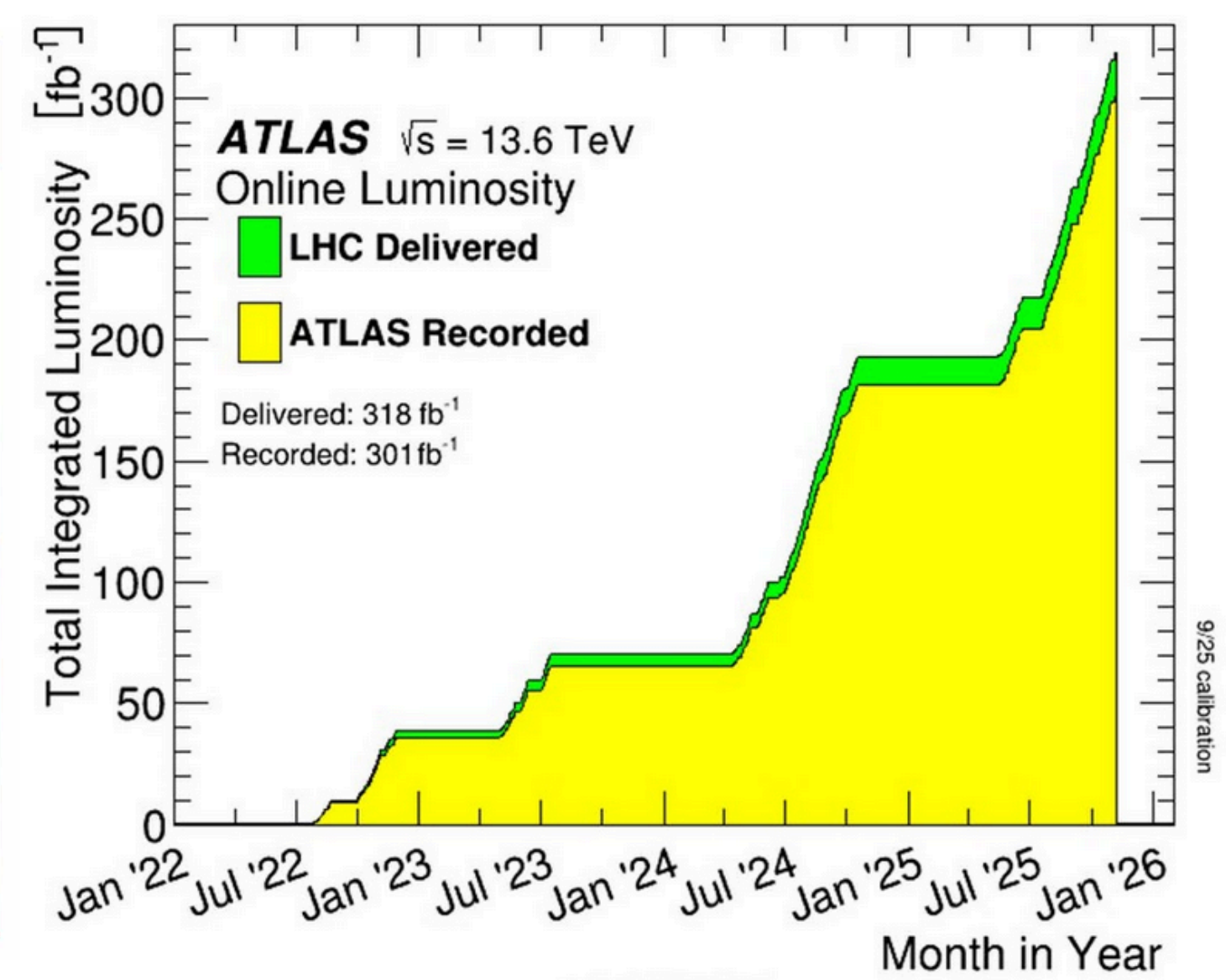
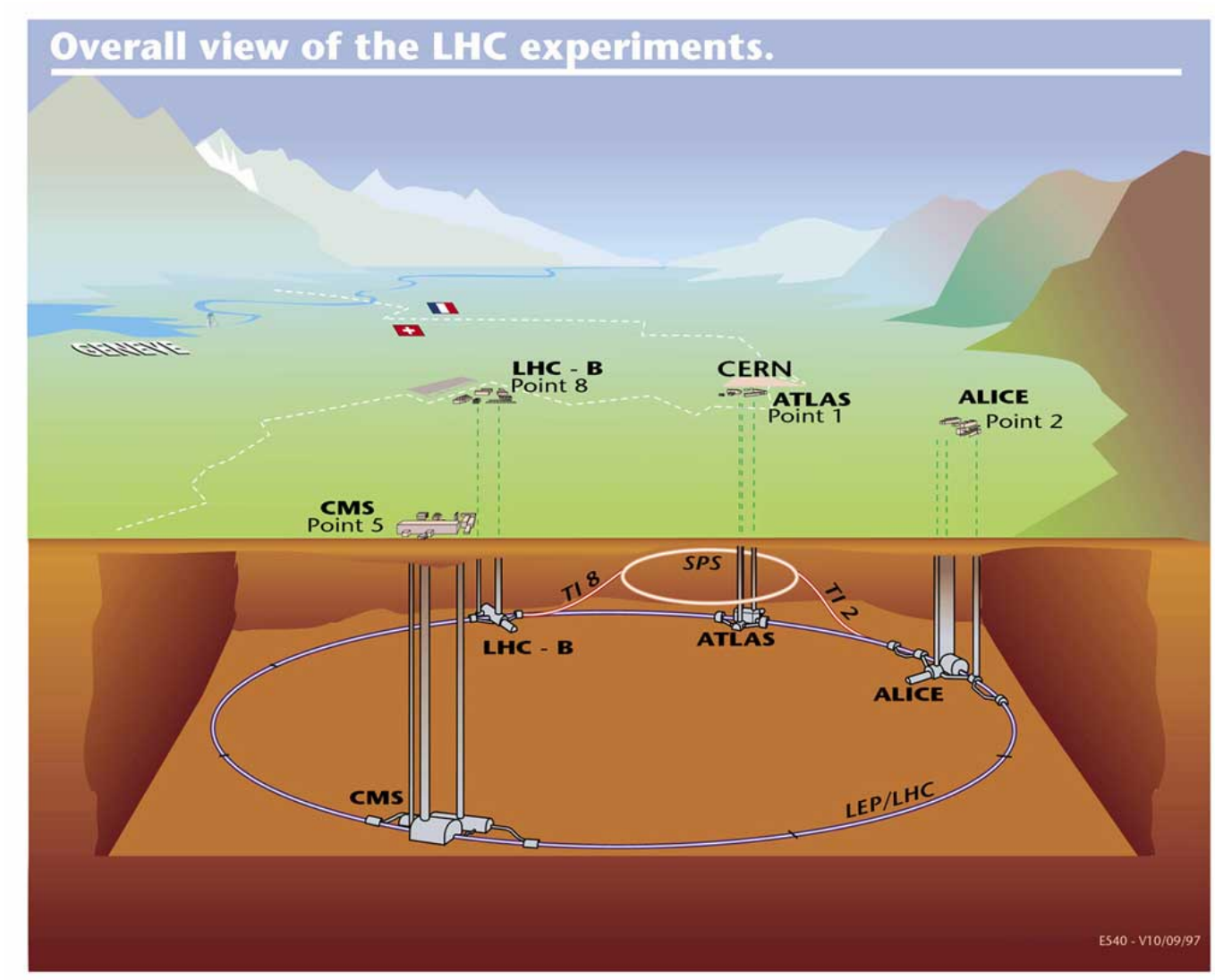
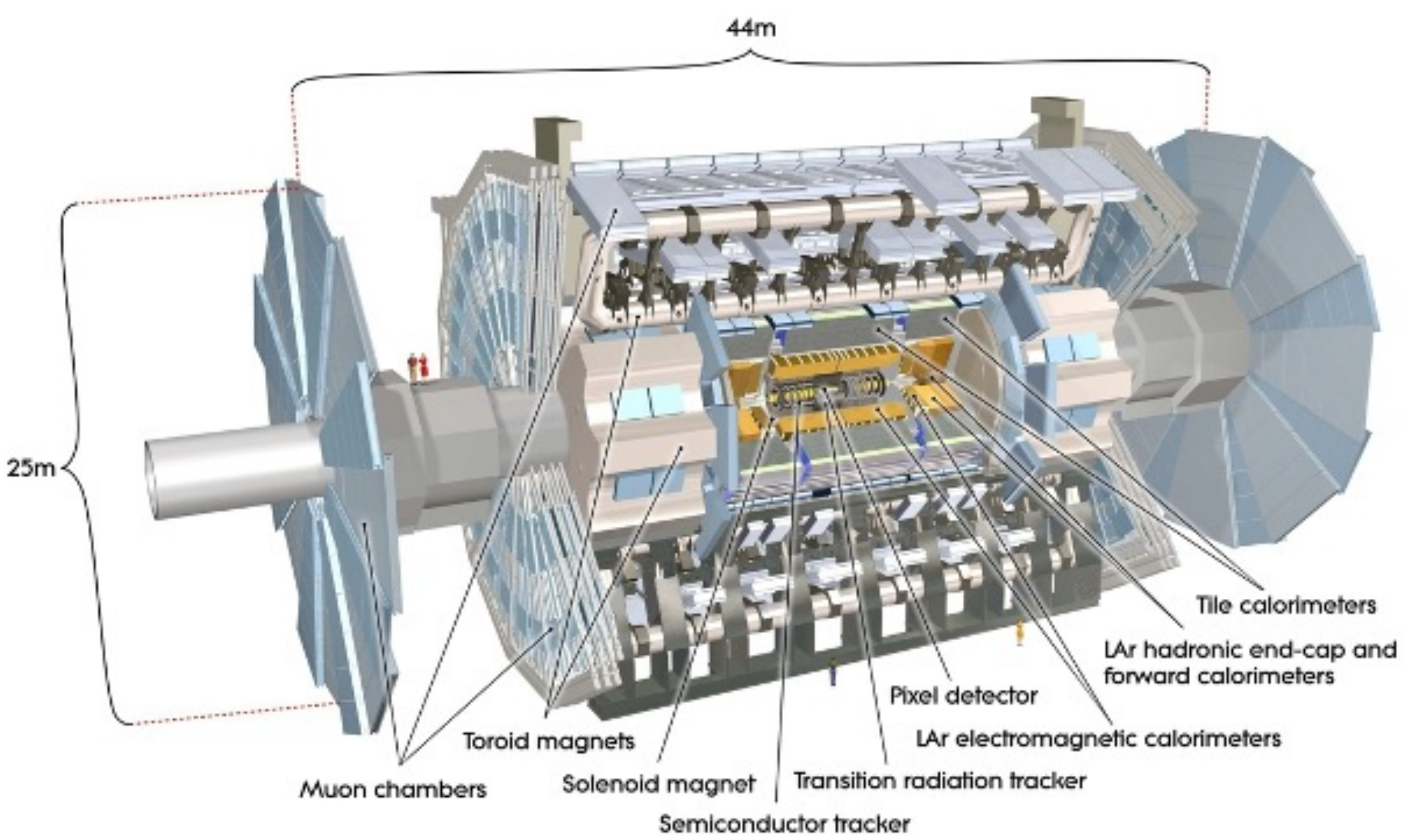
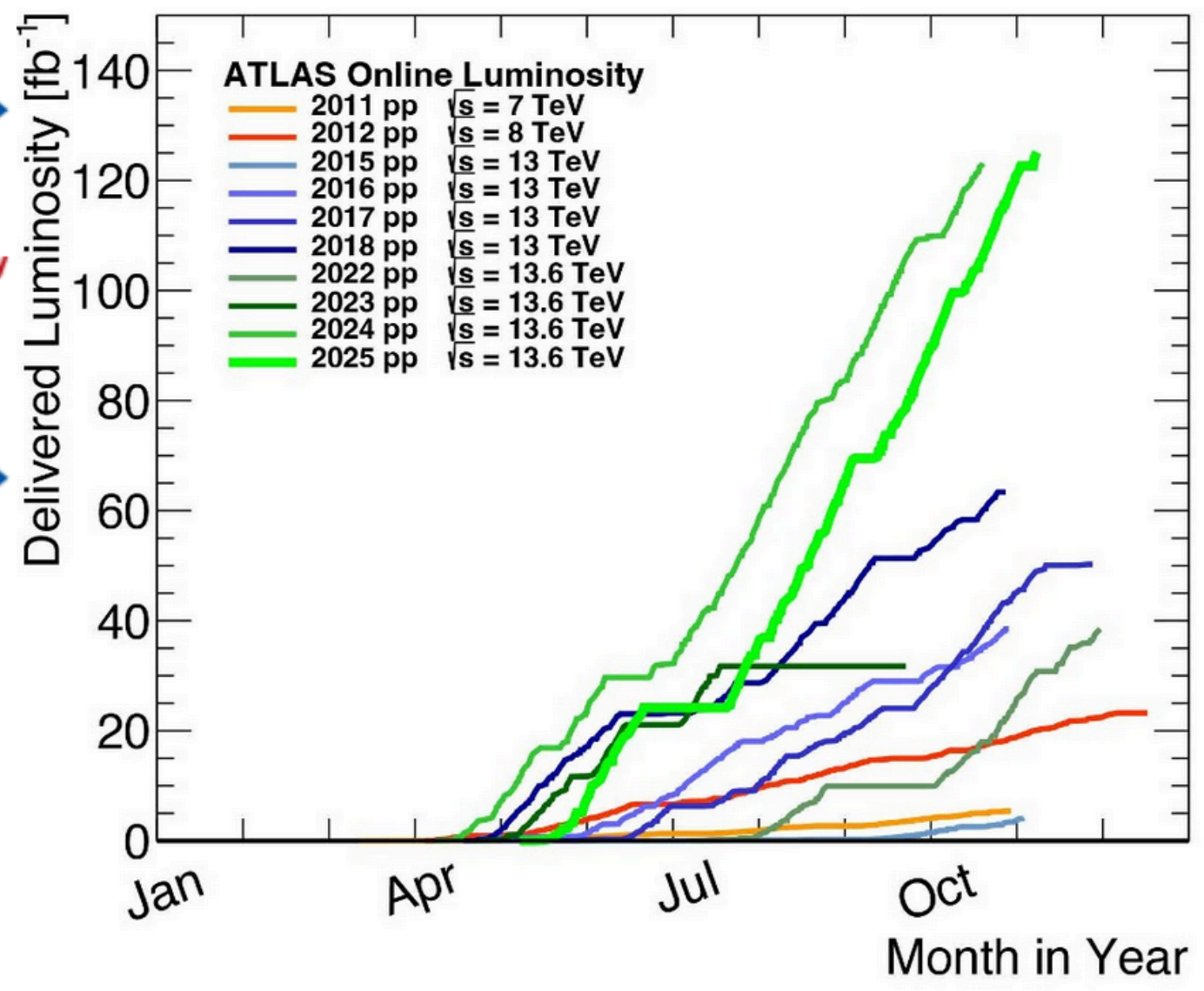
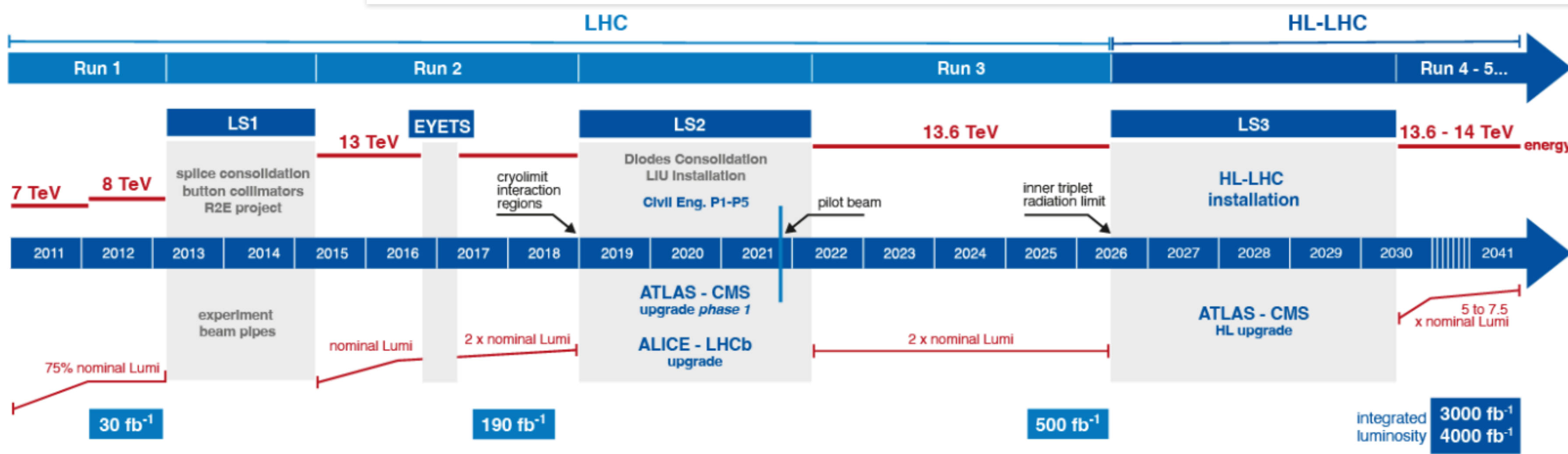


International Conference on the Physics of the Two Infinities
Tokyo, November 17-21, 2025

Outline

- Electroweak Symmetry Breaking
- Vector Boson Scattering
- Higgs Boson Properties
 - Higgs-Top Coupling *Sign?*
- Higgs Boson Rare Decays
 - $H \rightarrow \mu\mu$
 - $H \rightarrow Z\gamma$
- The Higgs Potential
 - Multi-Higgs Boson Production
- Concluding Remarks

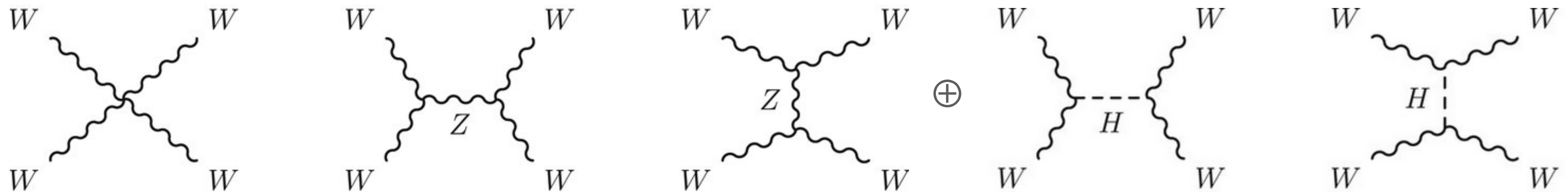
The LHC, ATLAS, Data



Electroweak Symmetry Breaking

- Without a Higgs boson, or other new physics, $\sigma(WW \rightarrow WW) \sim s$
 - This is similar to e.g. low energy neutrino scattering, where this behavior is “regulated” by the W boson propagator as the energy increases
 - To solve this for WW scattering, introduce an even-spin boson with $m \lesssim 1 \text{ TeV}$

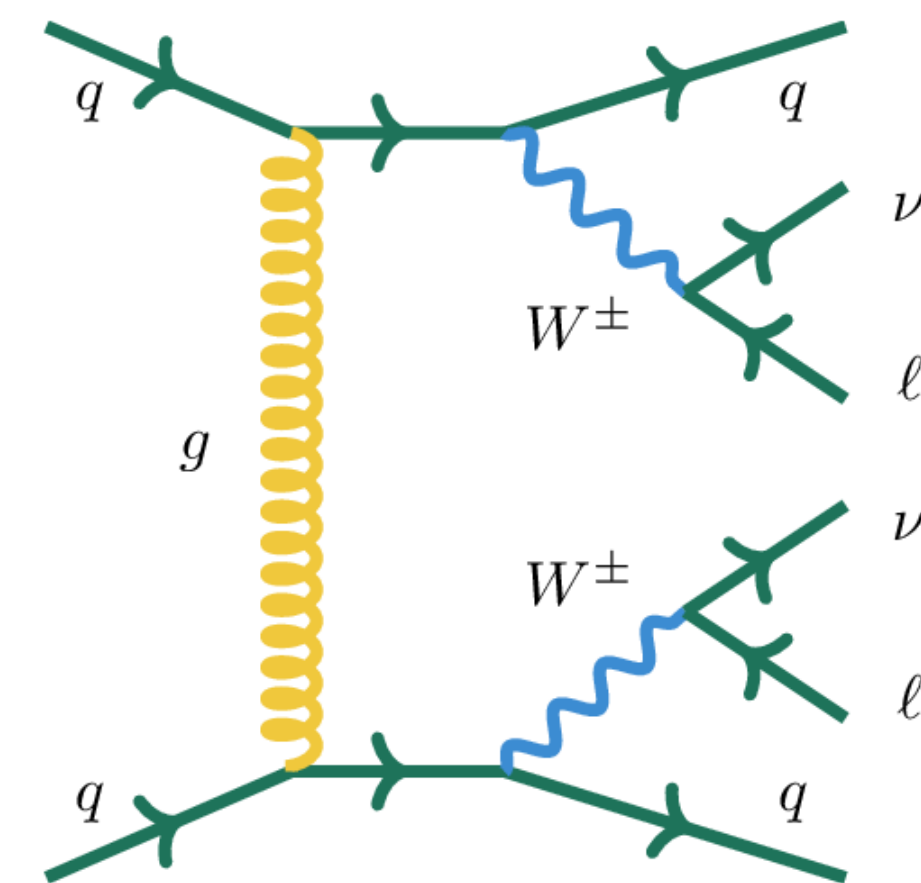
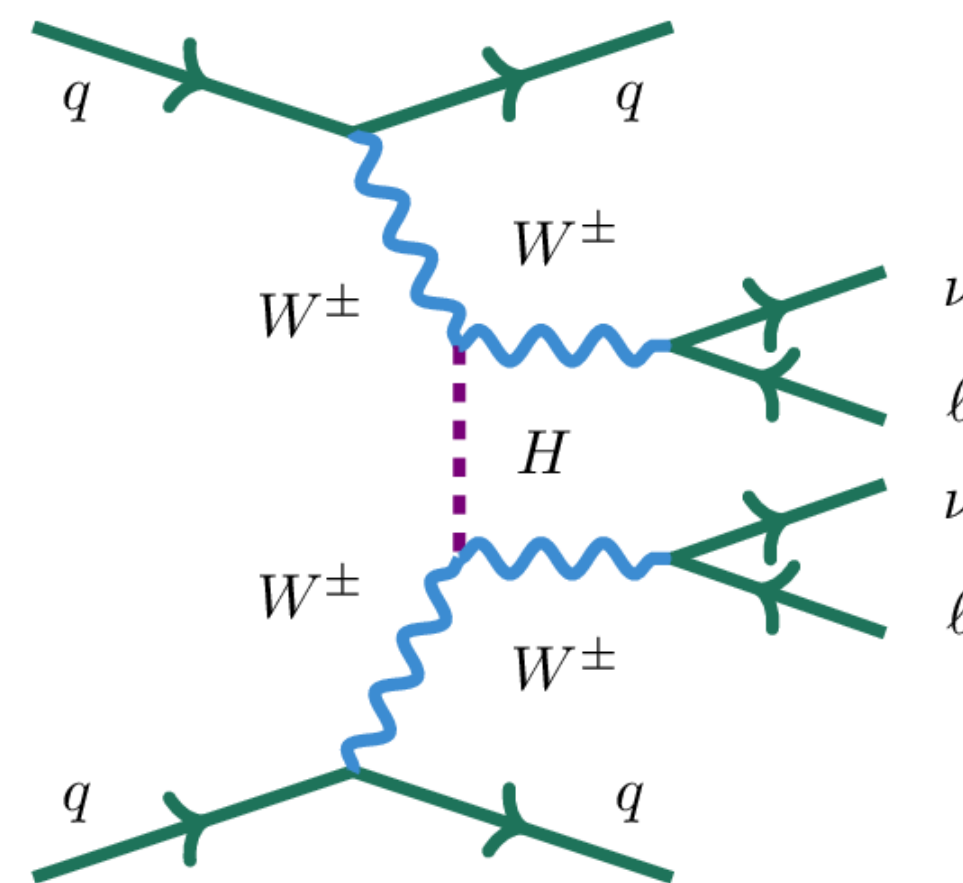
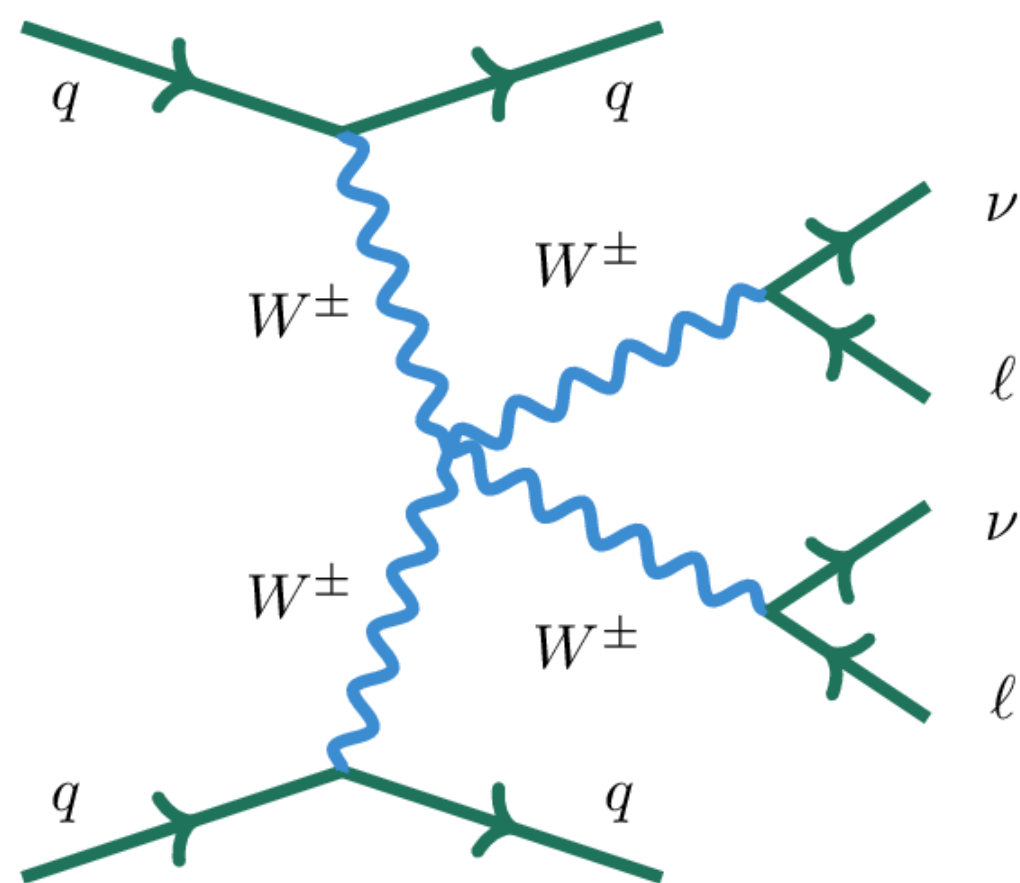
Lee, Quigg, Thacker, 1977



- A similar, but “softer”, issue arises in fermion scattering: $\sigma(f\bar{f} \rightarrow WW) \sim m_f \sqrt{s}$
 - To be economical, solve both problems at once: new boson needs to be a neutral, colorless scalar \Rightarrow quantum numbers of the vacuum!
- From a Standard Model lagrangian point of view, it needs to be embedded in (at least) an $SU(2)_L$ doublet...

Probe #1: Vector Boson Scattering

- Before the Higgs boson observation, vector boson scattering formed the basis of the no-lose theorem: something had to happen in VBS
 - In particular, scattering of longitudinally polarized vector bosons, since these are the degrees of freedom that only exist because of non-zero vector boson mass
 - At hadron colliders, W/Z bosons can be radiated from incoming quarks
 - Vector boson scattering (and fusion) is characterized by two “forward” jets
 - Jet transverse momentum (p_T) is $O(m_W/2)$

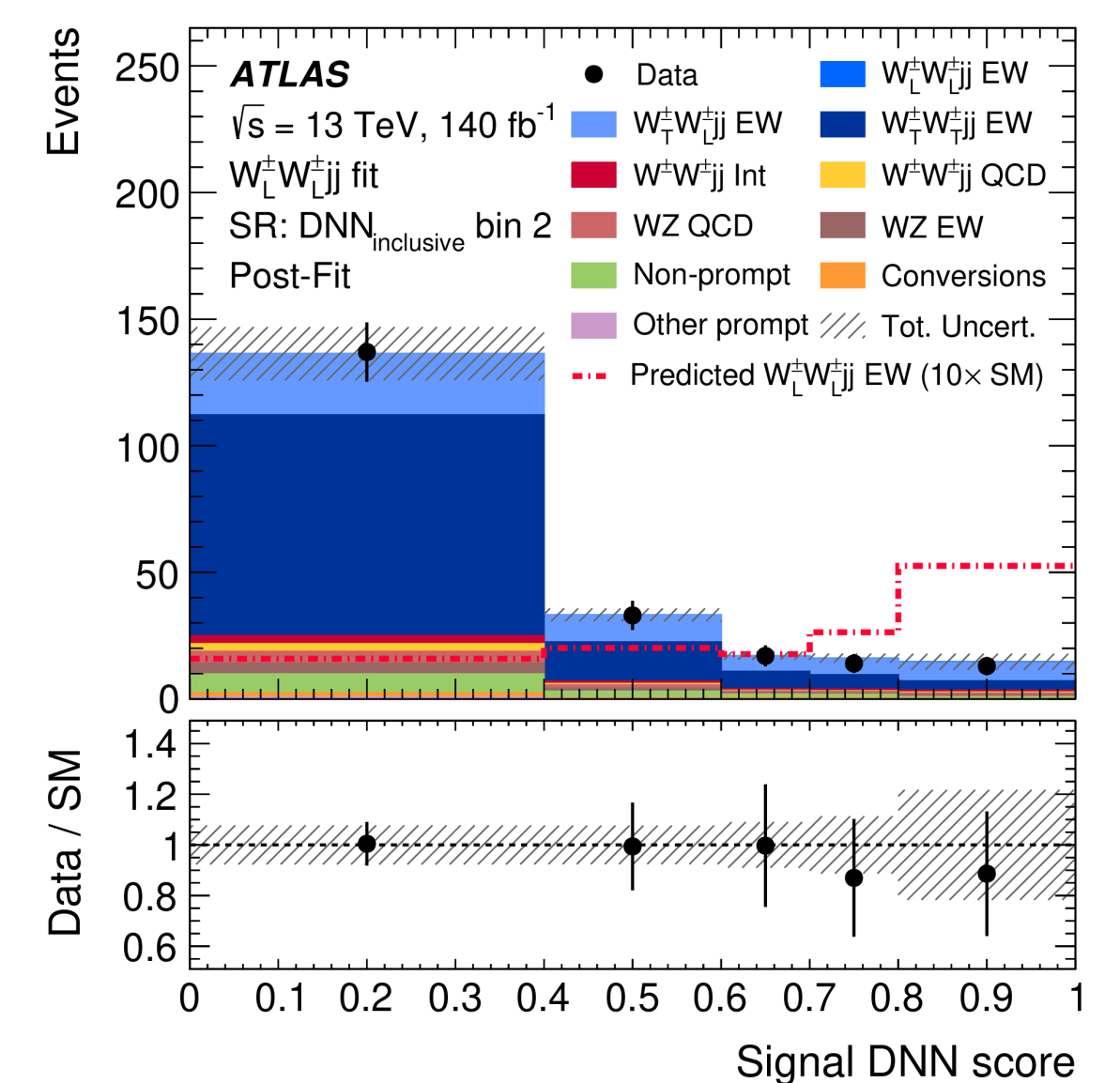
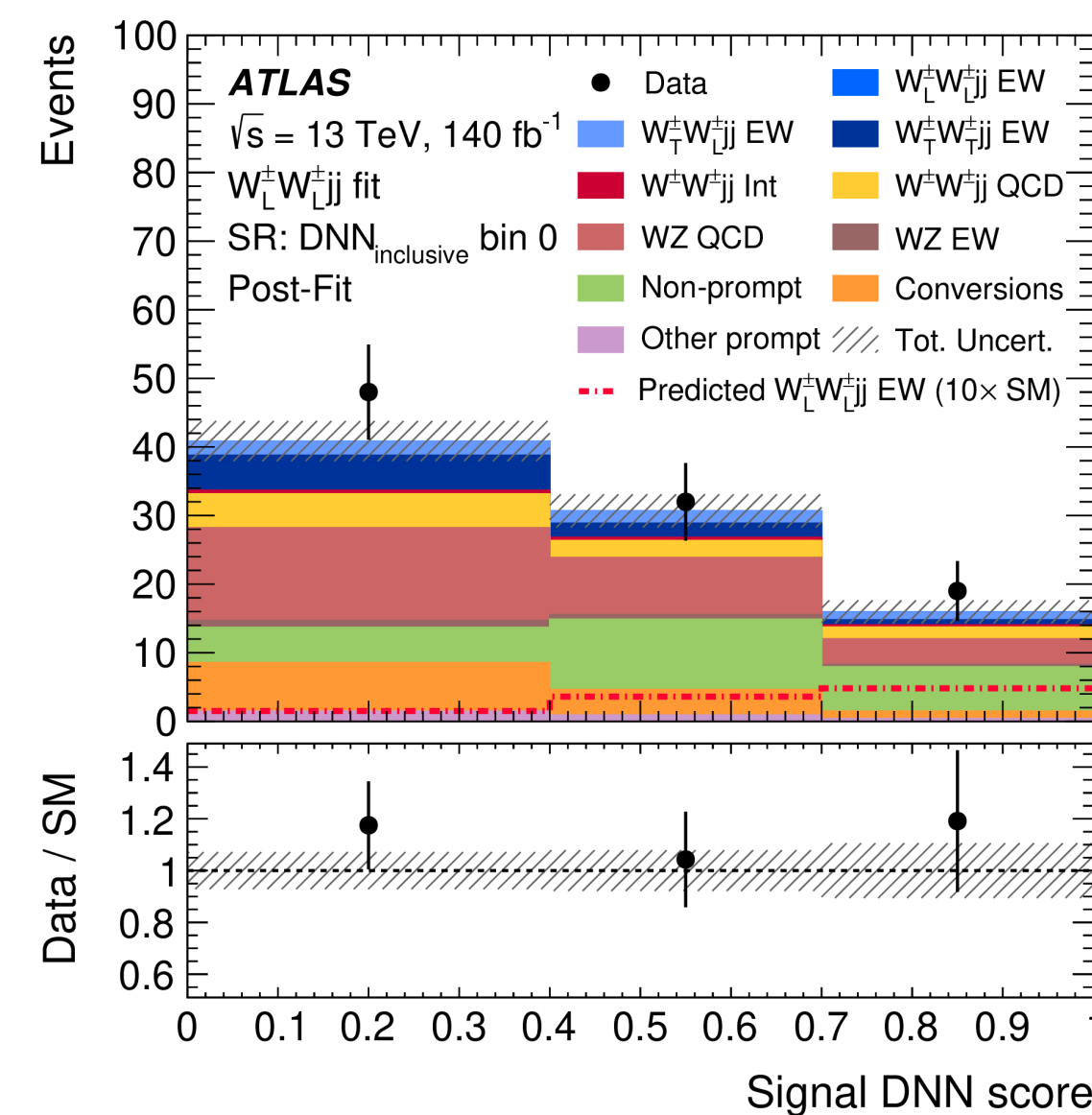
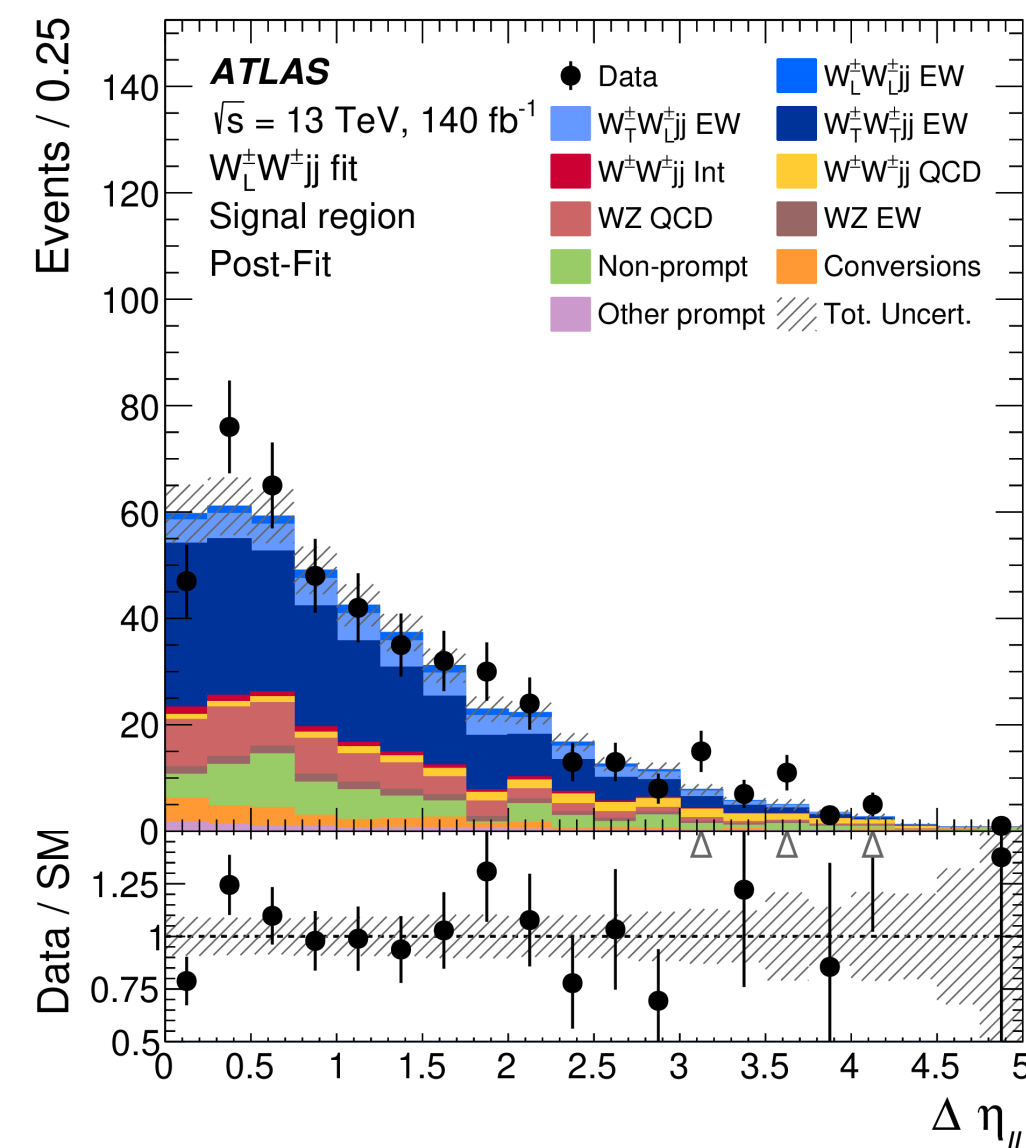
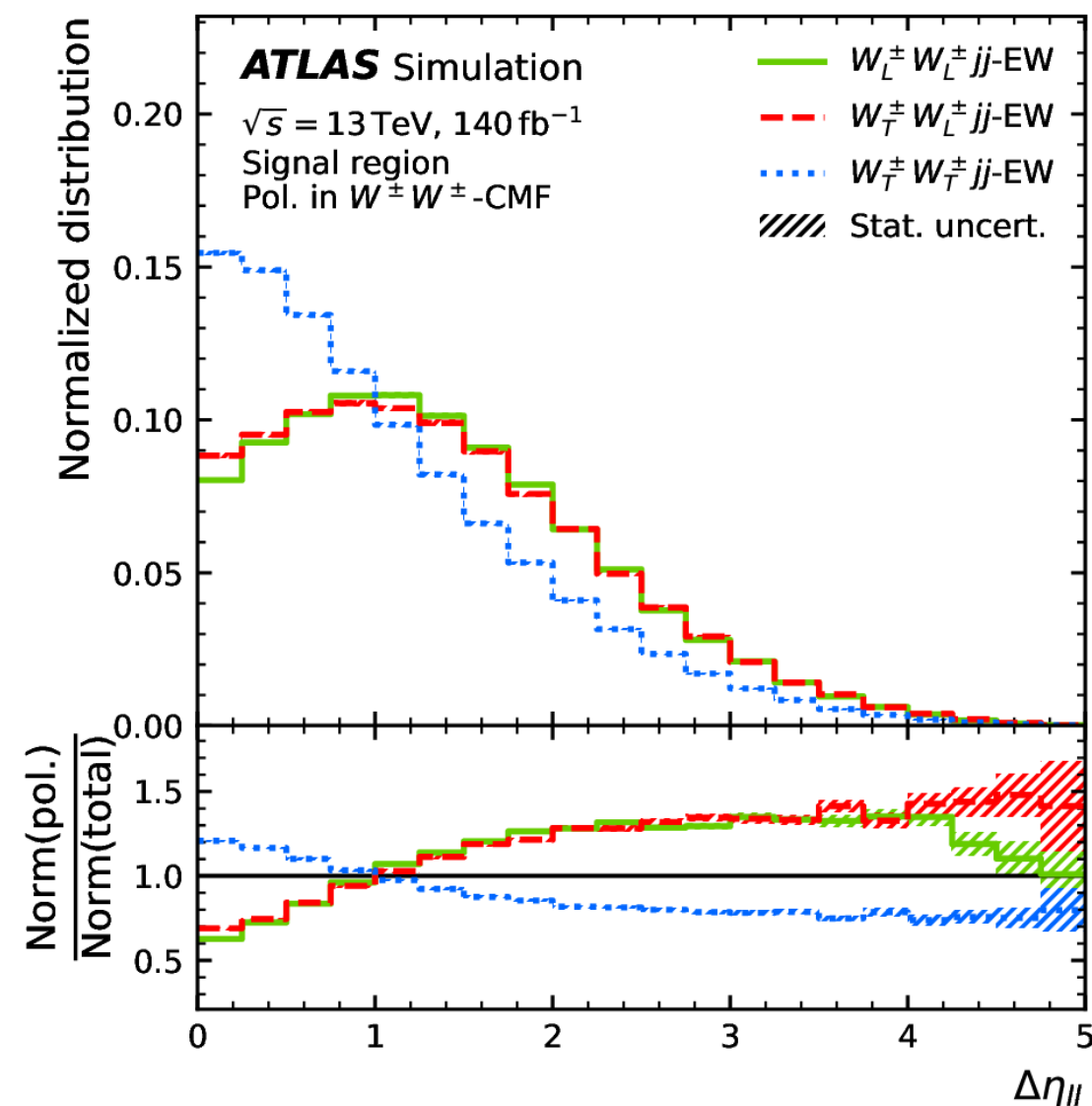


ATLAS Search for $W_L W_L$ Scattering

- To minimize backgrounds, look for same-sign $W^\pm W^\pm$ scattering, with each boson decaying to an electron or muon and a neutrino:
 - A pair of same-charge leptons, large MET, a pair of forward jets with large m_{jj}
- Then two-component analysis:
 - NN to separate WW scattering signal from non-VBS backgrounds, 3 bins
 - NNs to separate $W_L W_L$ from $W_T W$ and $W_L W$ from $W_T W_T$
 - Use all NN outputs in final fit

[Phys. Rev. Lett. 135 \(2025\) 111802](#)

[ATLAS webpage with plots](#)



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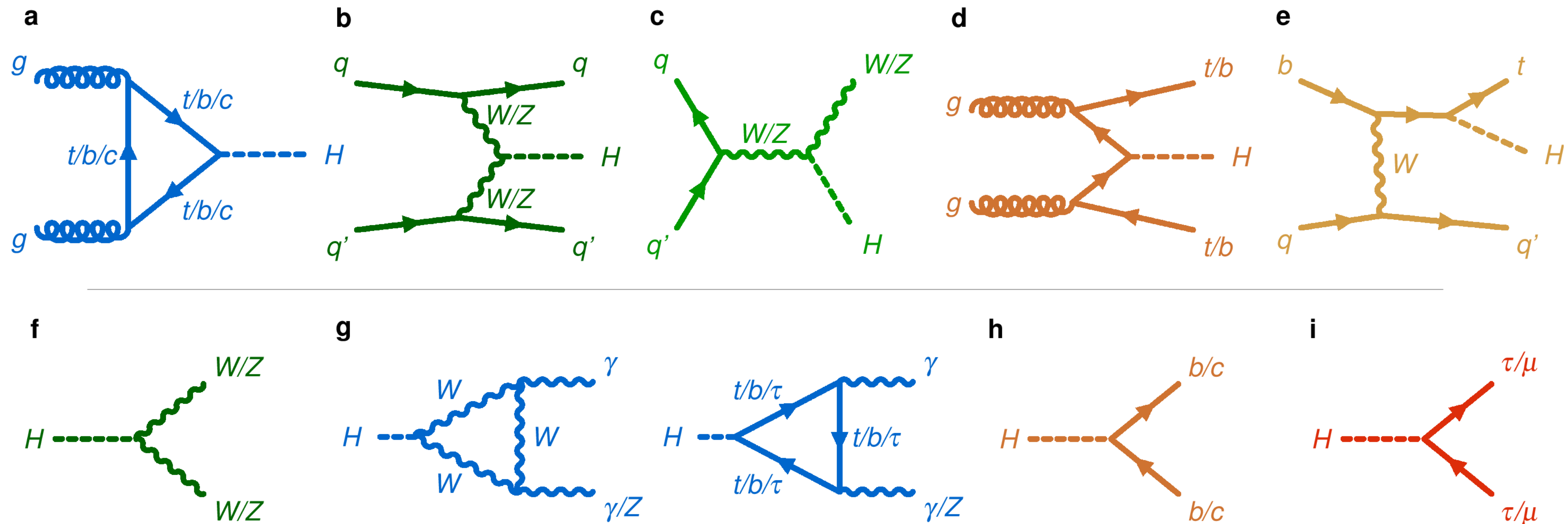
[ATLAS webpage with plots](#)

| Process | Predicted $\sigma\mathcal{B}$ (fb) | Measured $\sigma\mathcal{B}$ (fb) | Uncertainty breakdown (fb) |
|----------------------|------------------------------------|-----------------------------------|--|
| $W_L^\pm W_L^\pm jj$ | 0.29 ± 0.07 | 0.01 ± 0.21 (tot.) | ± 0.20 (stat.) ± 0.02 (mod. syst.) ± 0.05 (exp. syst.) |
| $W_T^\pm W^\pm jj$ | 2.56 ± 0.64 | 3.39 ± 0.35 (tot.) | ± 0.30 (stat.) ± 0.08 (mod. syst.) ± 0.16 (exp. syst.) |
| $W_L^\pm W^\pm jj$ | 1.18 ± 0.29 | 0.88 ± 0.30 (tot.) | ± 0.28 (stat.) ± 0.05 (mod. syst.) ± 0.08 (exp. syst.) |
| $W_T^\pm W_T^\pm jj$ | 1.67 ± 0.40 | 2.49 ± 0.32 (tot.) | ± 0.30 (stat.) ± 0.05 (mod. syst.) ± 0.12 (exp. syst.) |

3σ evidence for VBS with at least one longitudinally polarized W boson

“Legacy” Run 2 Higgs Measurements

- Combination of Run 2 (2015-2018) results to constrain the properties of the observed Higgs boson, and compare with SM expectations

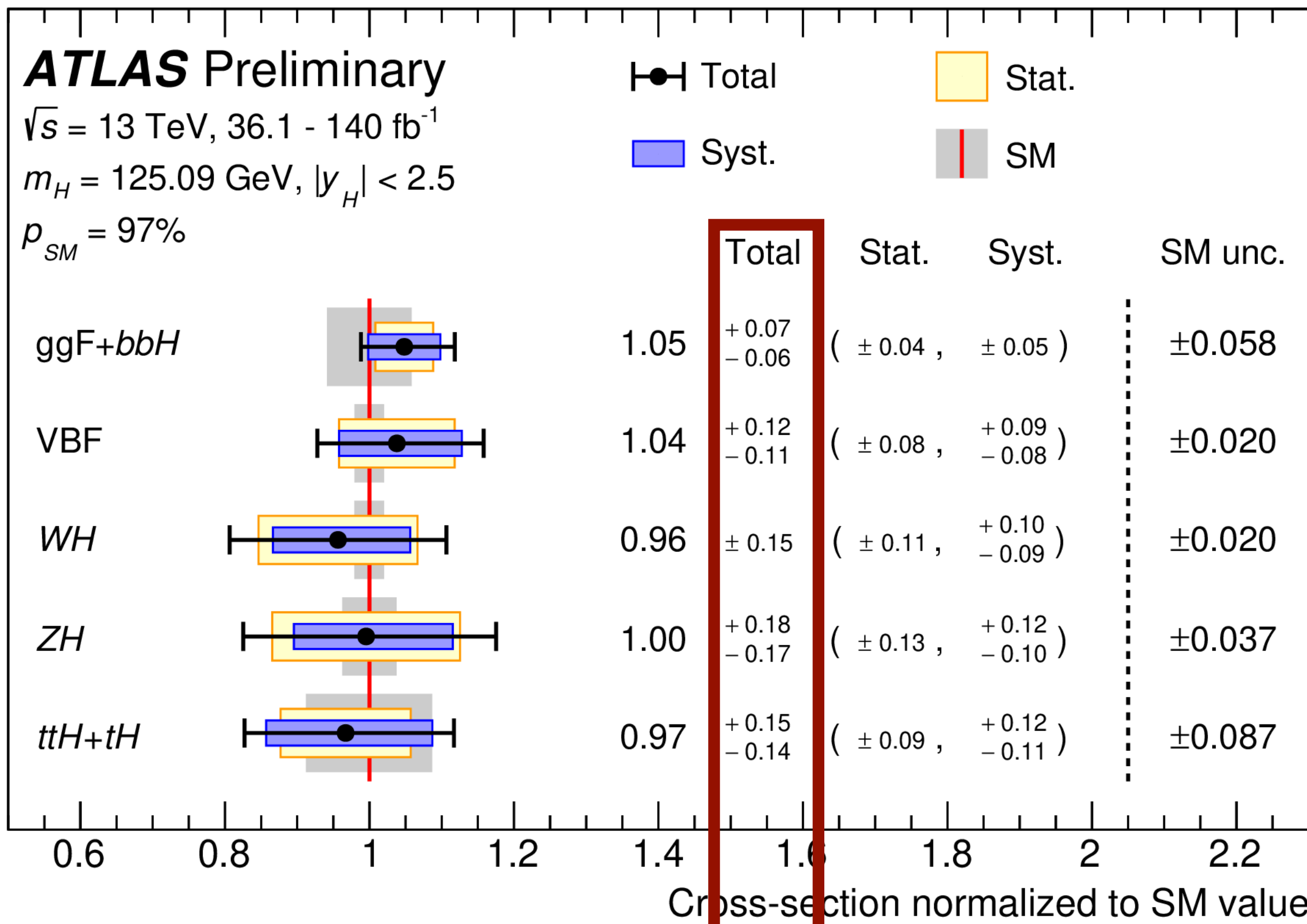


- Probes multiple production and decay modes

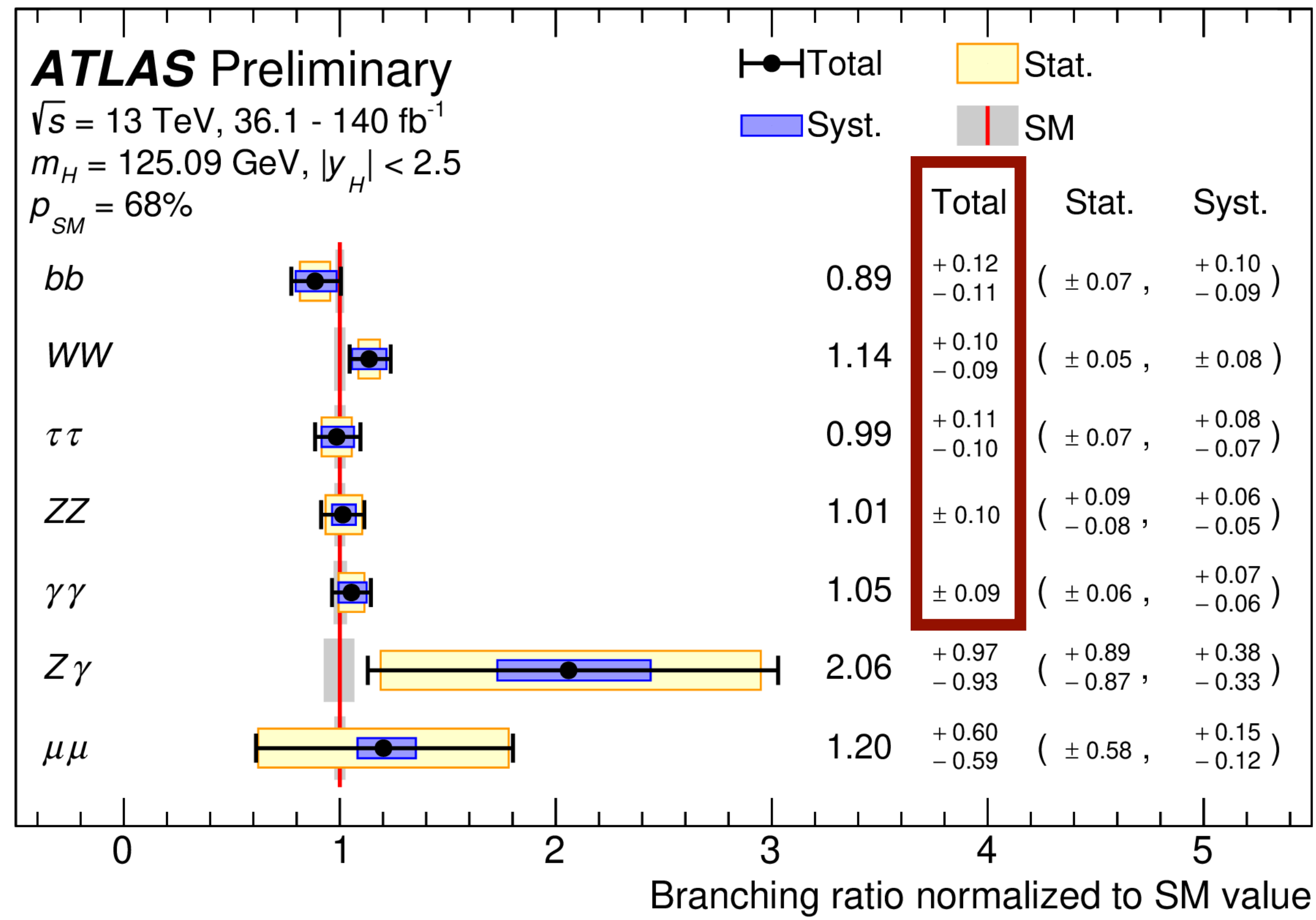
[ATLAS-CONF-2025-006](#)
[ATLAS webpage with Plots](#)

- Given the significant contributions of loops in both production and decay, combination adds significant constraining power

Production and Decay



10-20%



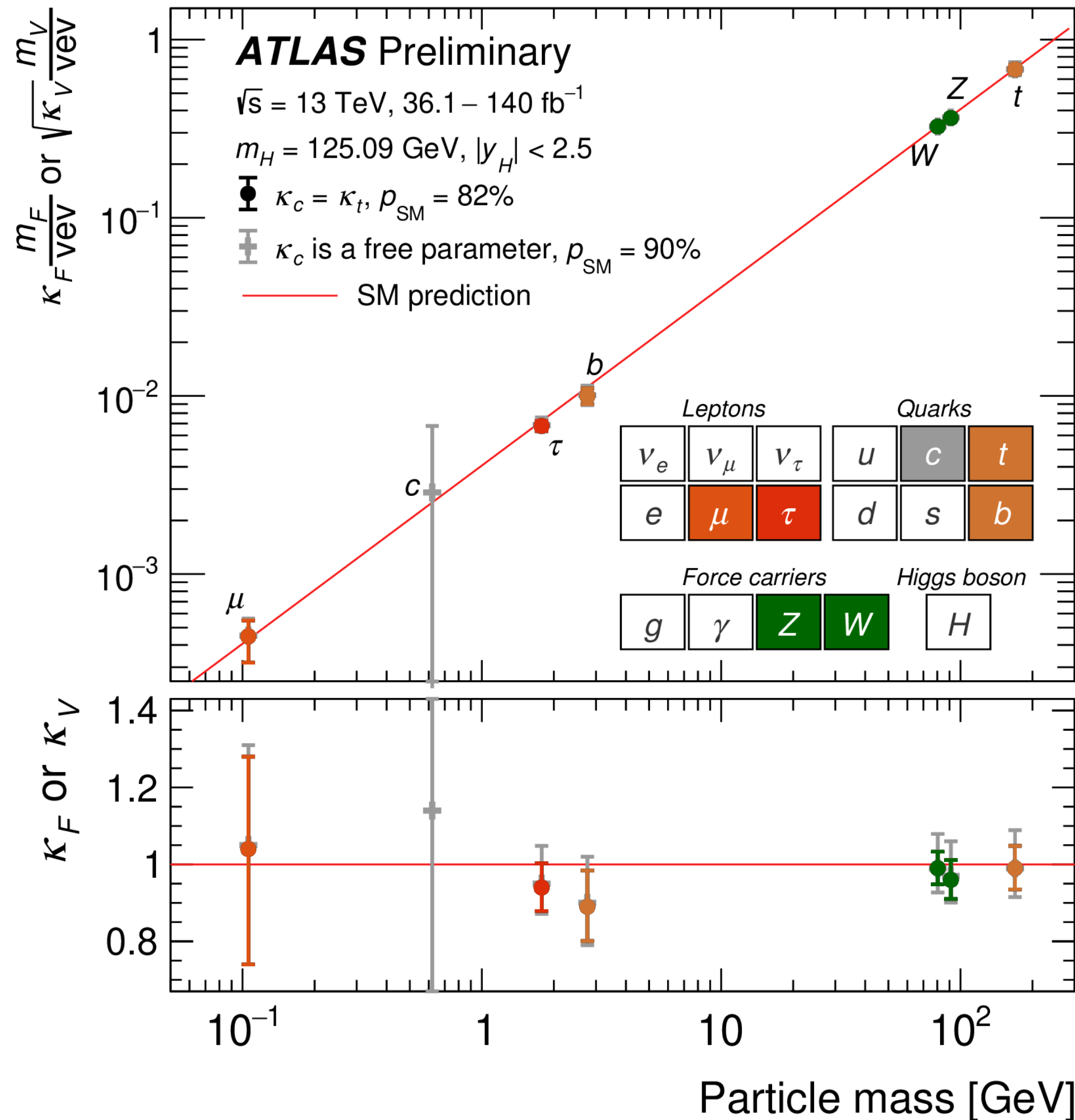
~10%

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[ATLAS webpage with Plots](#)

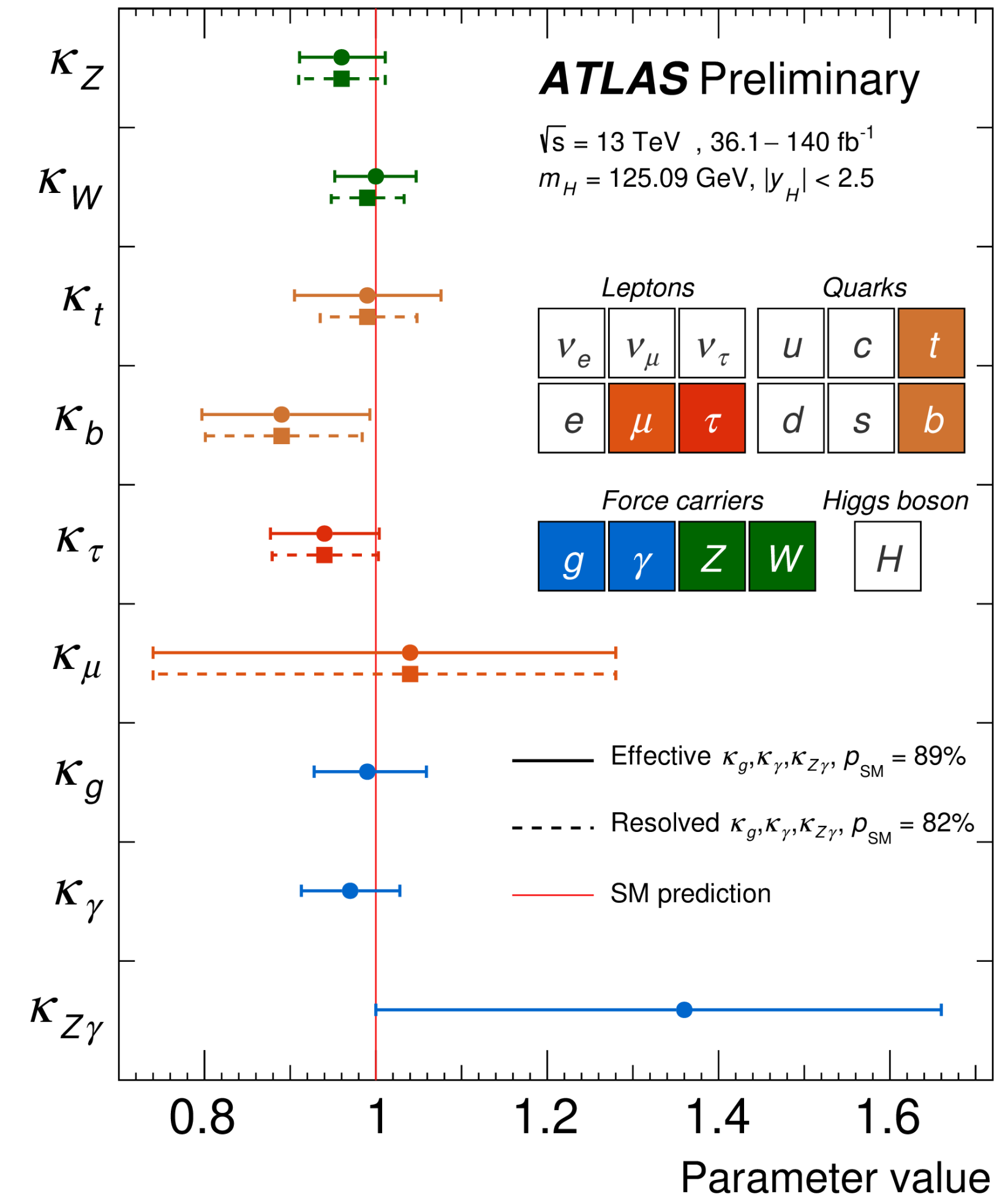
Couplings

[ATLAS-CONF-2025-006](#)
[ATLAS webpage with Plots](#)

Proportionality



Coupling “Modifiers”

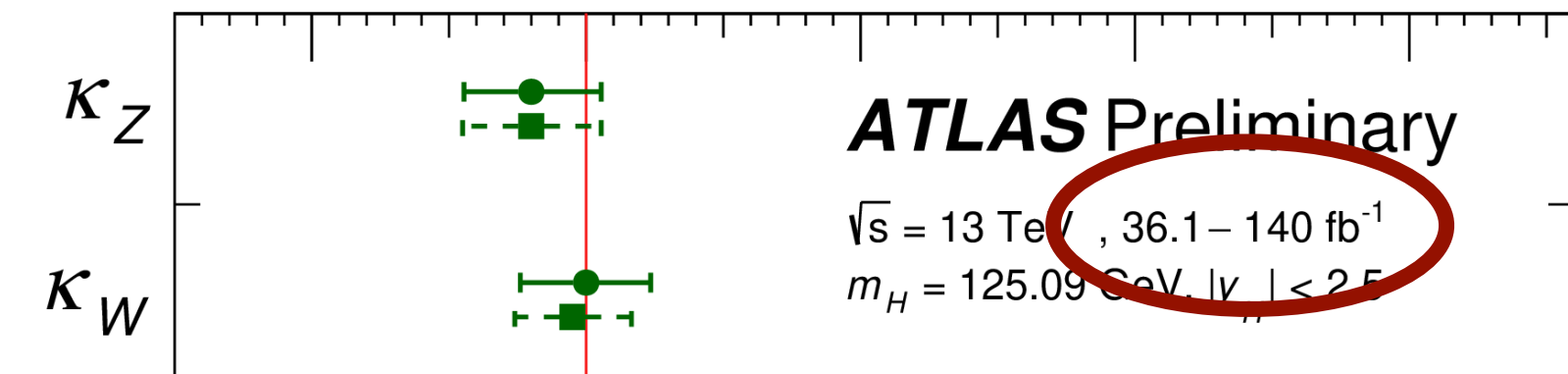
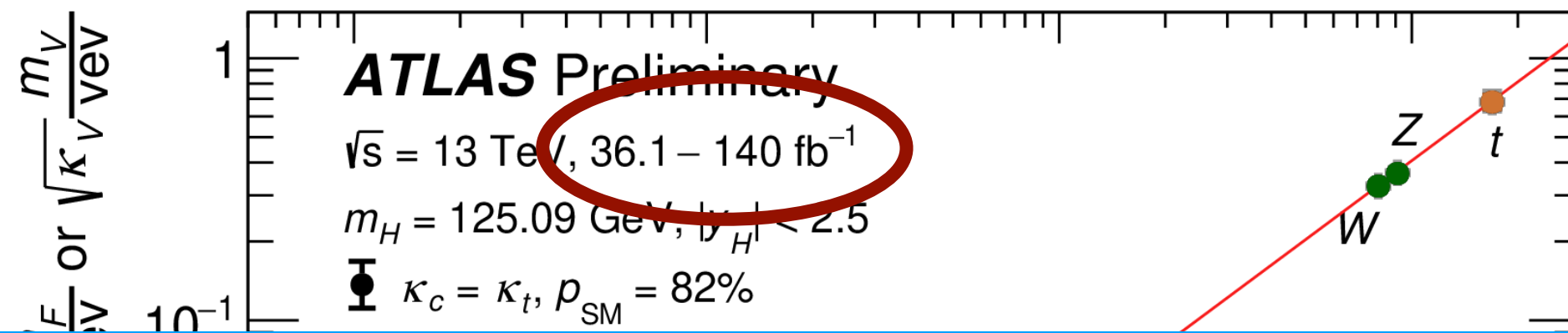


Couplings

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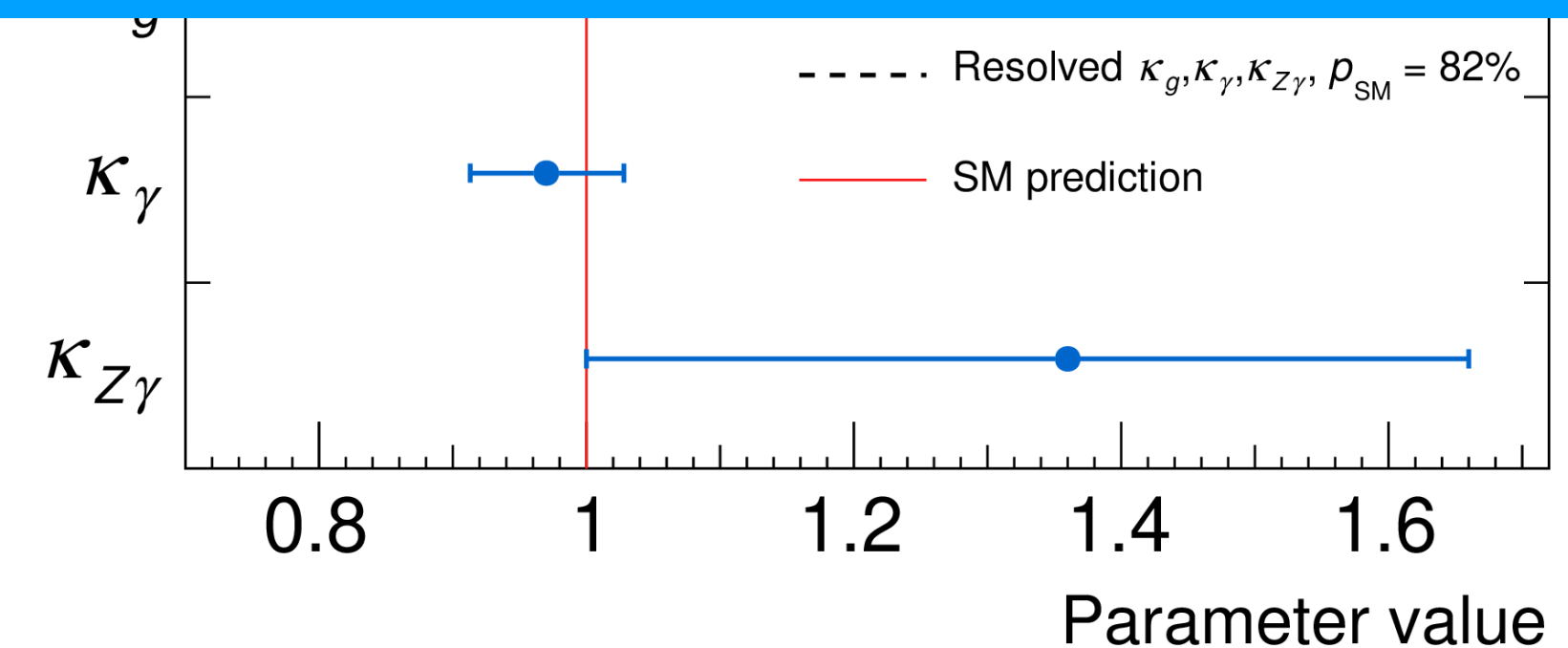
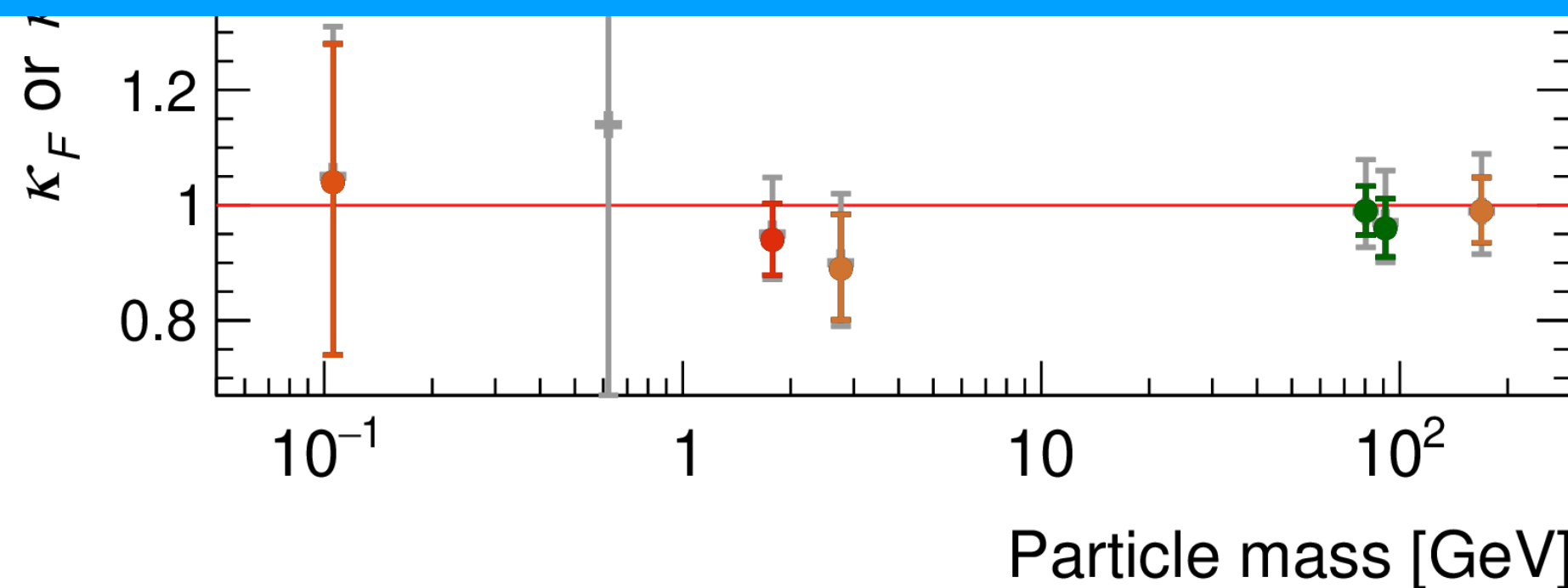
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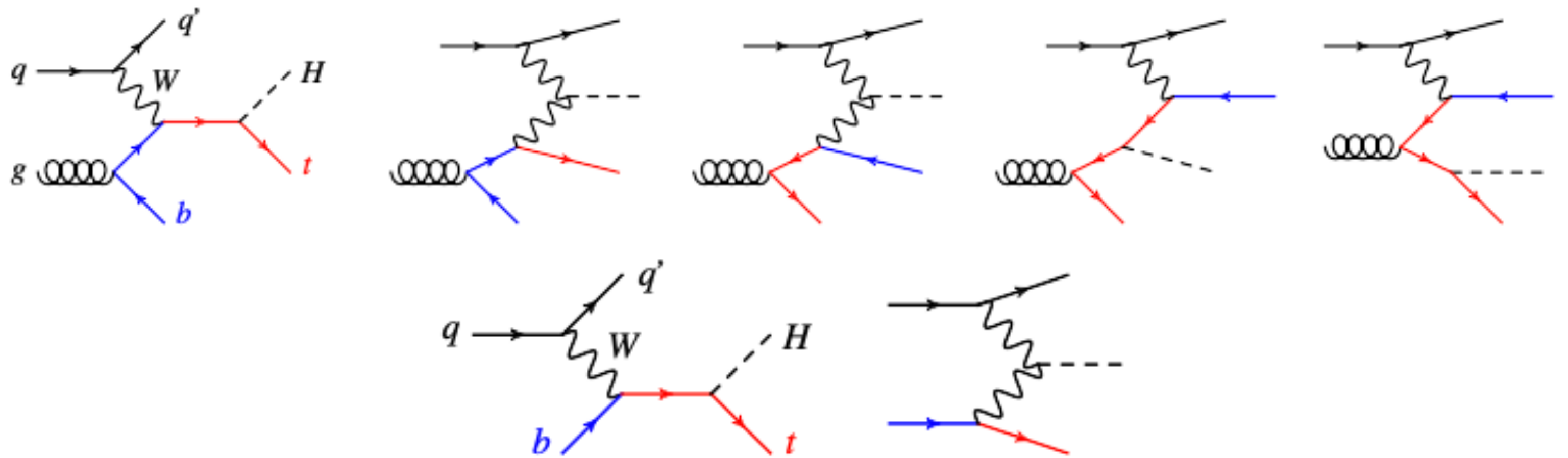
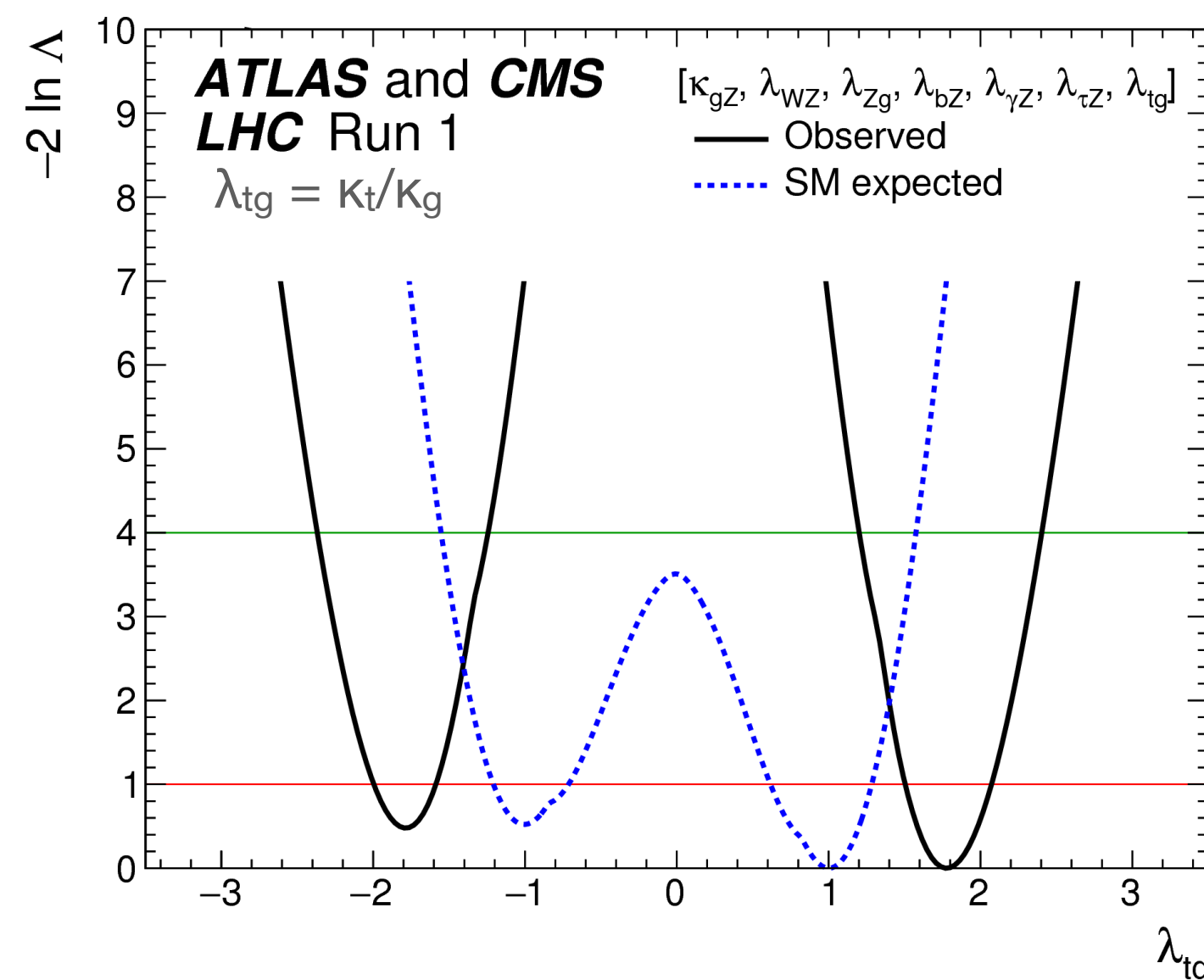
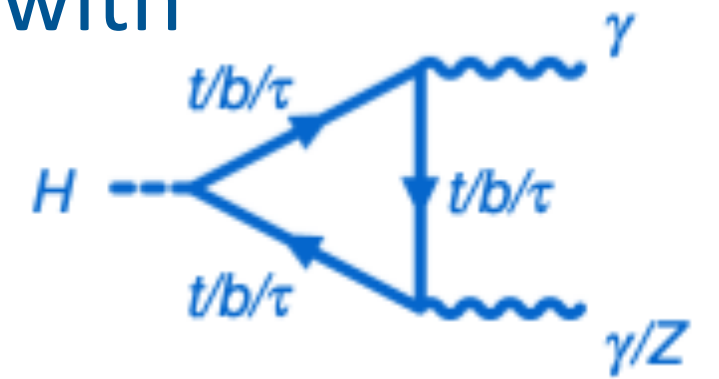
Couplings to 3rd generation fermions are within ~20% of SM prediction
 Couplings to SM interaction bosons are within ~10% of SM prediction

20x more data coming!



tH Production: y_t Sign

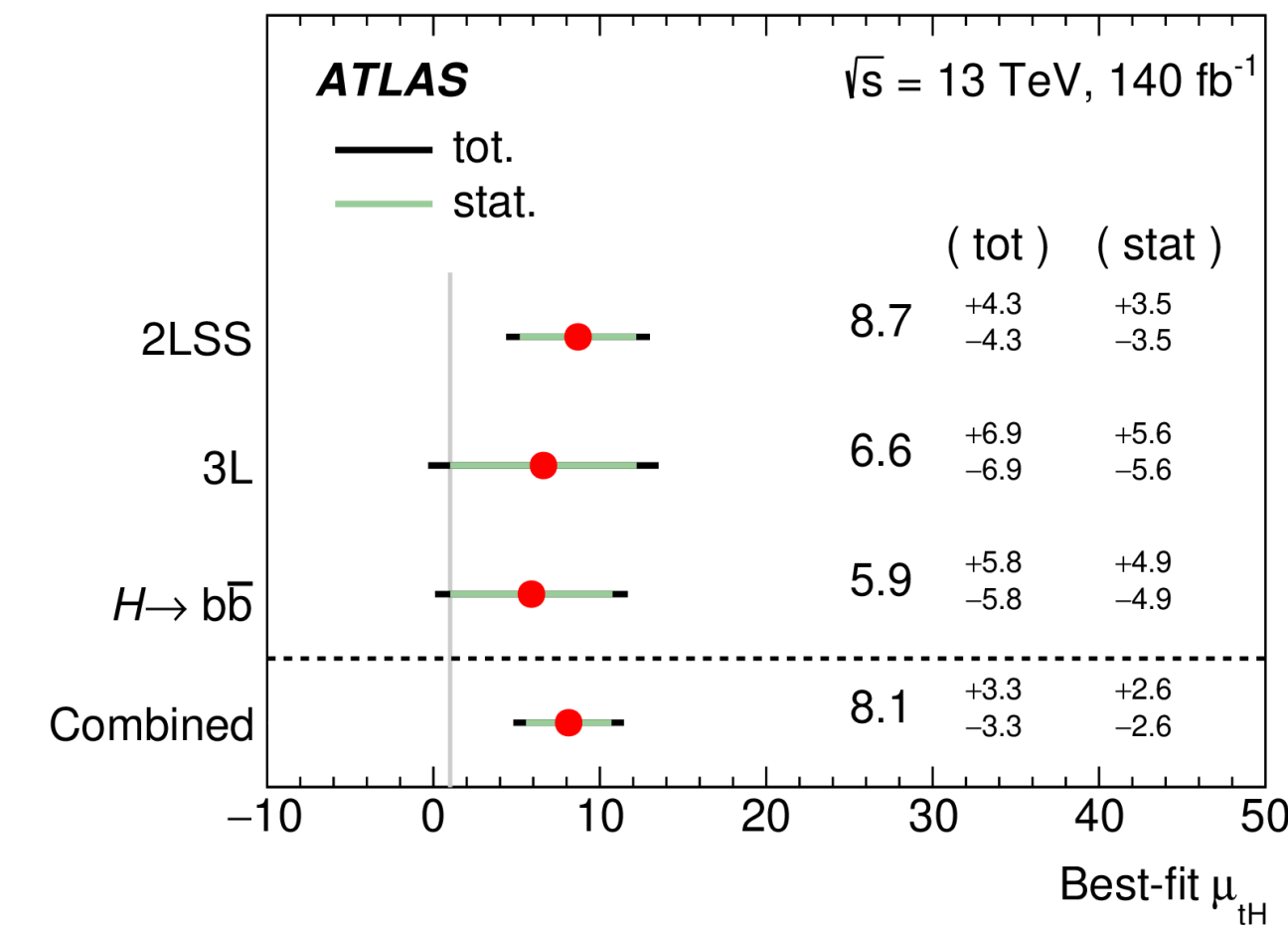
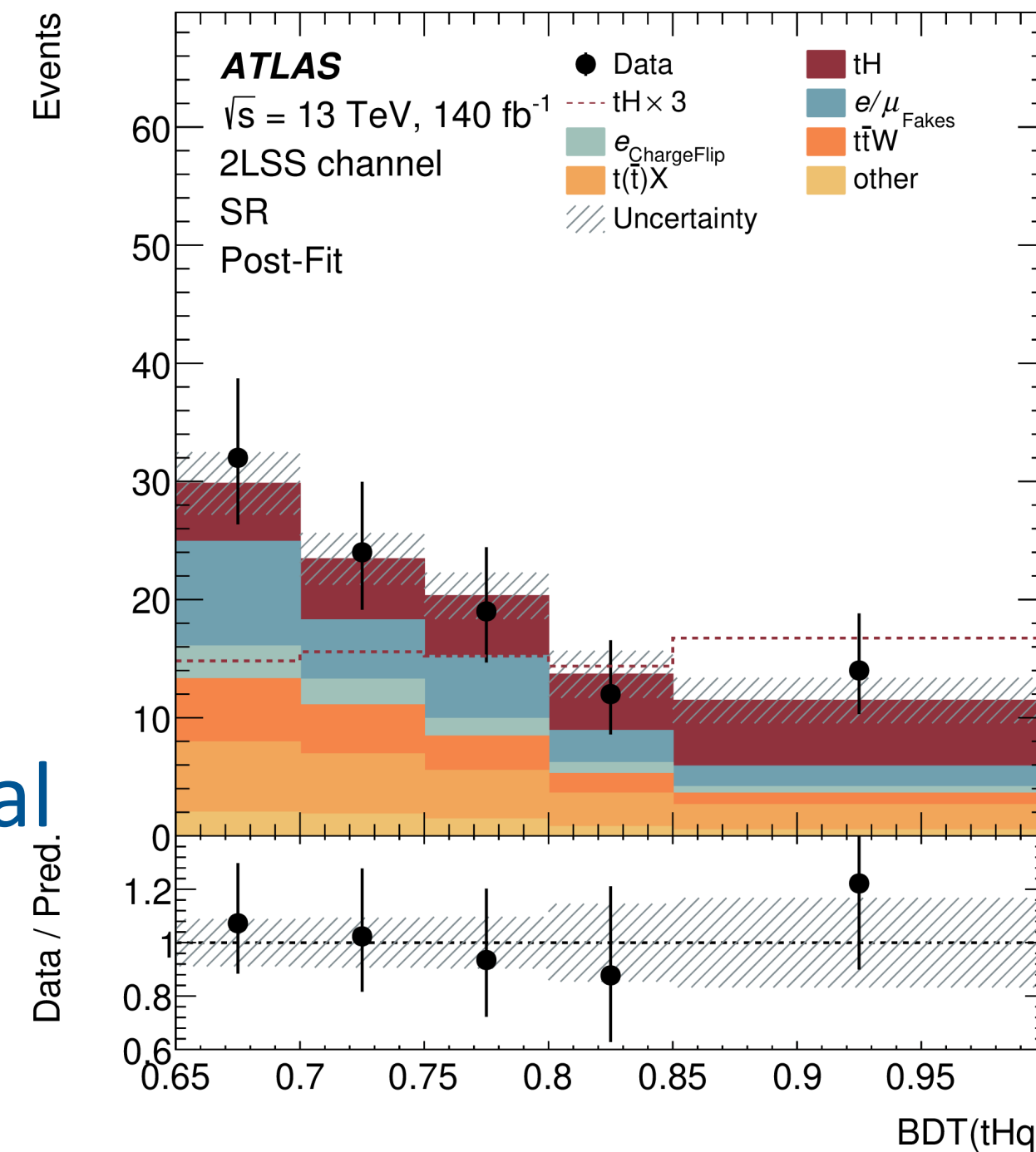
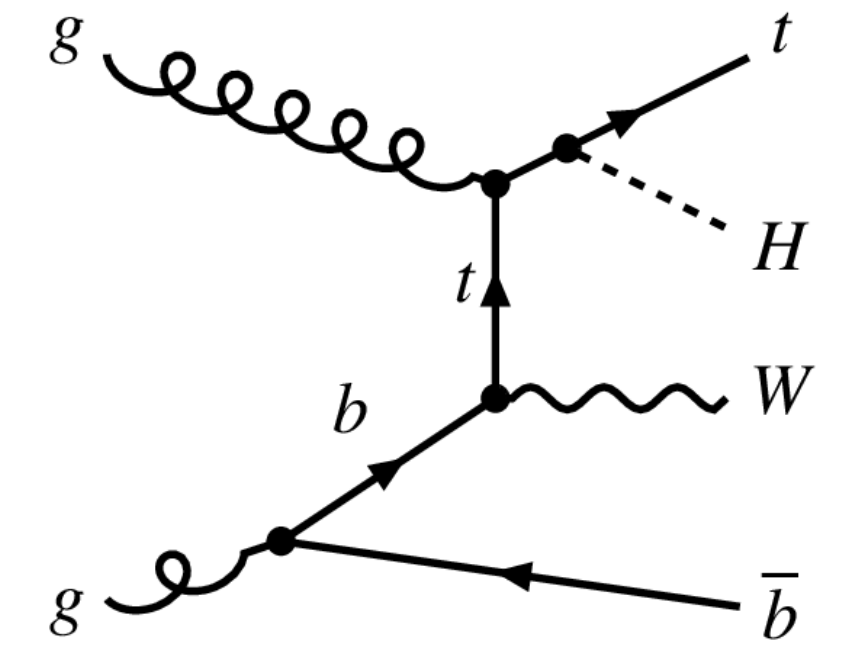
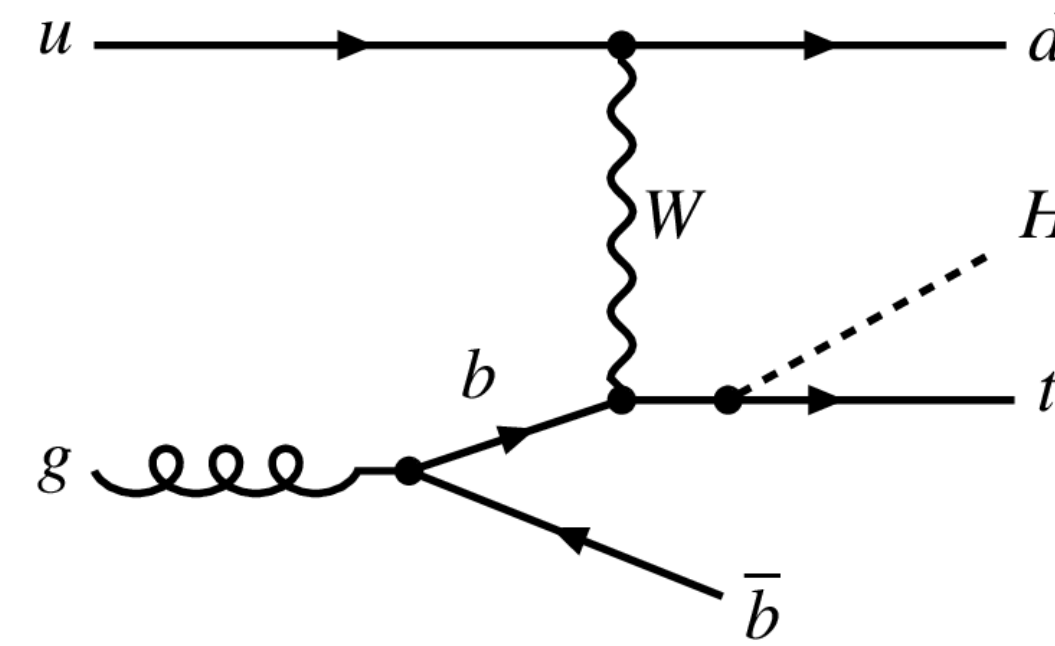
- Higgs coupling strength to the top quark (y_t) gets probed through multiple processes: not just $t\bar{t}H$, but also e.g. $H \rightarrow \gamma\gamma$ or gluon fusion H production (top loop)
 - These processes are mainly sensitive to $(y_t)^2$ - to probe a coupling's sign one needs a process with interference
 - $H \rightarrow \gamma\gamma$ (t-b interference) and ggF H production (t-W) have some sensitivity
- tH production however, is small in the SM *because* of significant destructive interference: with opposite sign the interference becomes constructive and $\sigma(tH)$ increases x 10



From Demartin et al, [Eur.Phys.J.C 75 \(2015\) 6, 267](#)

ATLAS Search for tH Production

- tHq and tHW yield distinct topologies
 - !tHW is same final state as ttH!
- Separate by number of leptons
 - 1 lepton ($H \rightarrow b\bar{b}$)
 - 2 same-charge leptons (mainly $H \rightarrow WW^*$)
 - 3 leptons (mainly $H \rightarrow WW^*$)
- Dedicated signal and control regions to constrain $t\bar{t}+b$, $t\bar{t}W$, “fake” leptons
- BDTs to separate signal from backgrounds
 - Only train to separate tHq from background: tHW is too close to $t\bar{t}H$ and including it in signal weakens separation!
 - But of course tHW contribution fully taken into account as signal

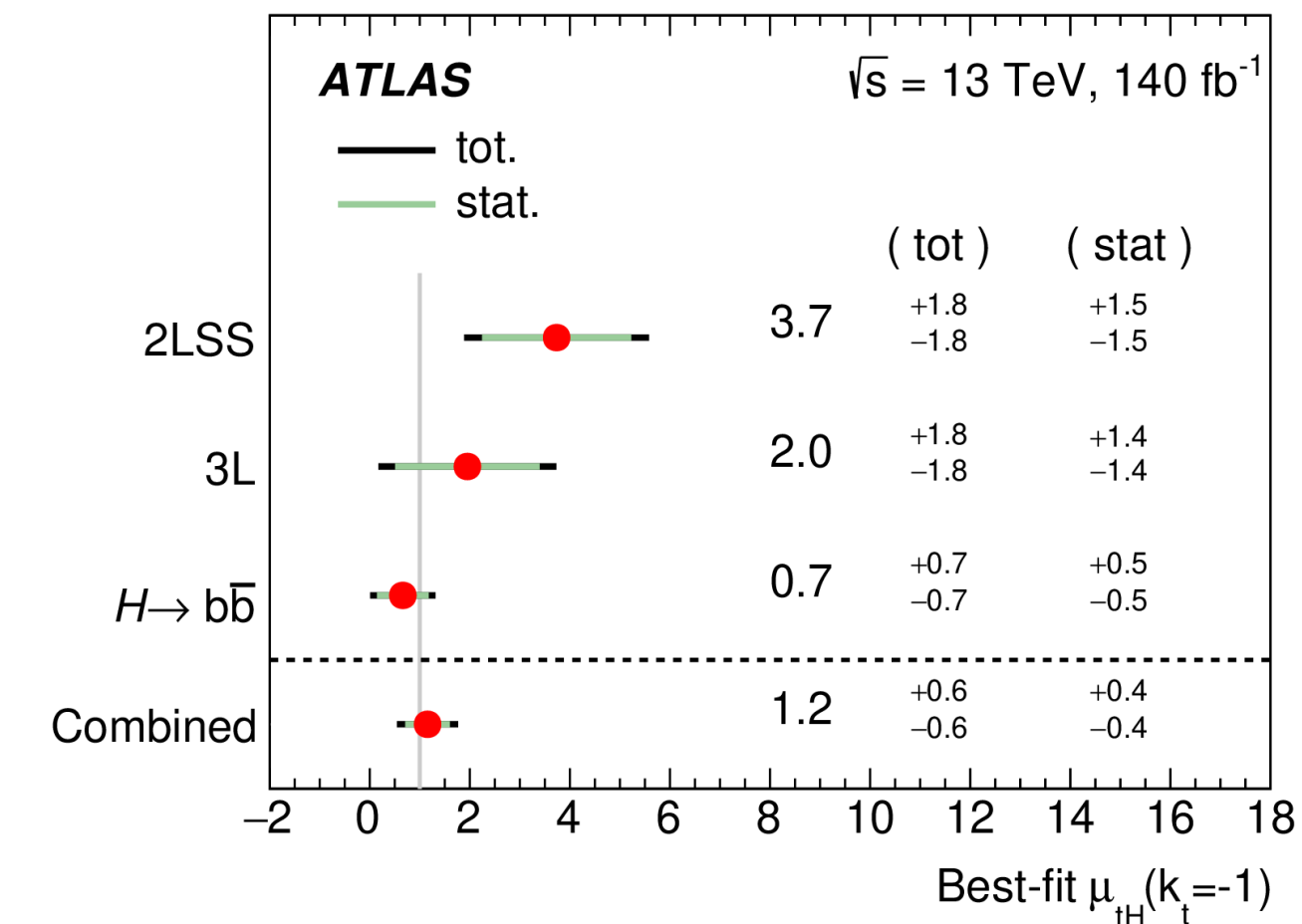
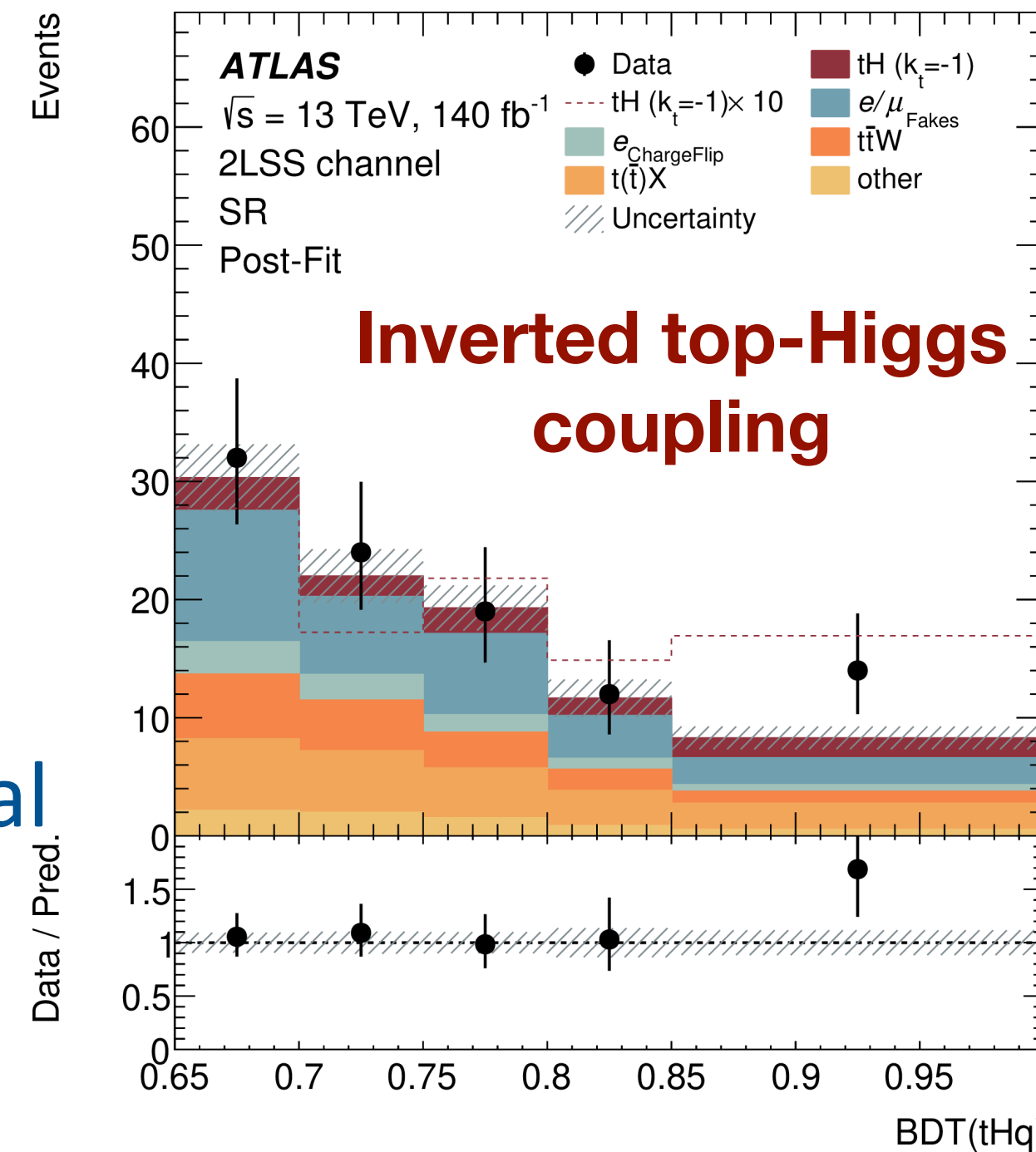
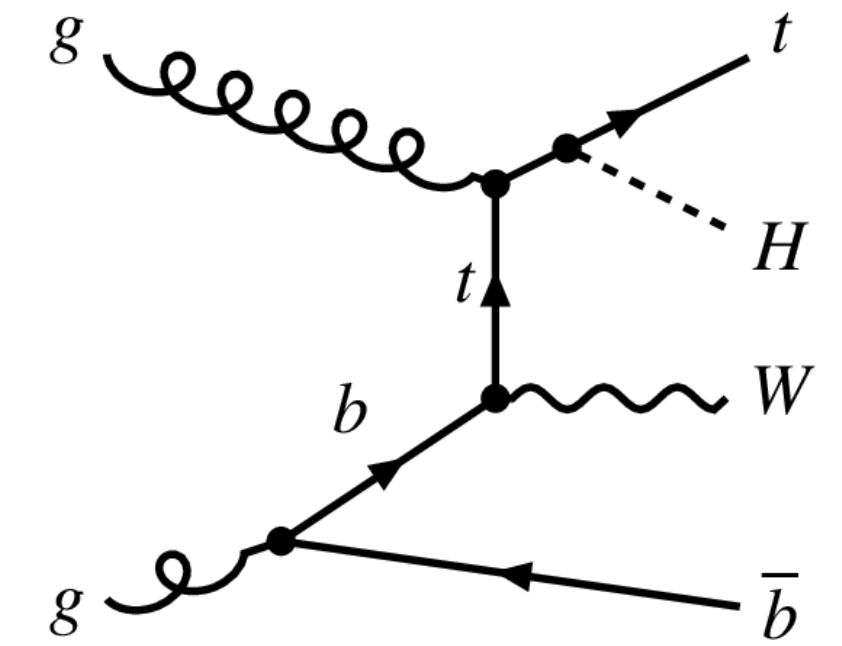
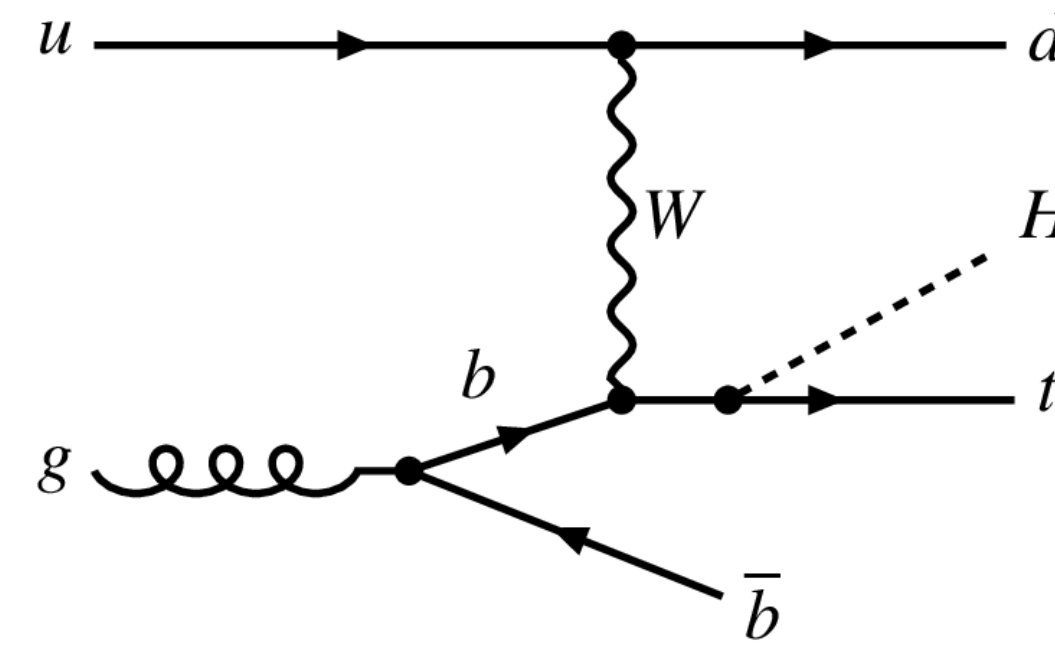


2.8 σ significance!
0.4 σ expected...

ATLAS, [arXiv:2508.14695](https://arxiv.org/abs/2508.14695)
[ATLAS webpage with Plots etc](#)

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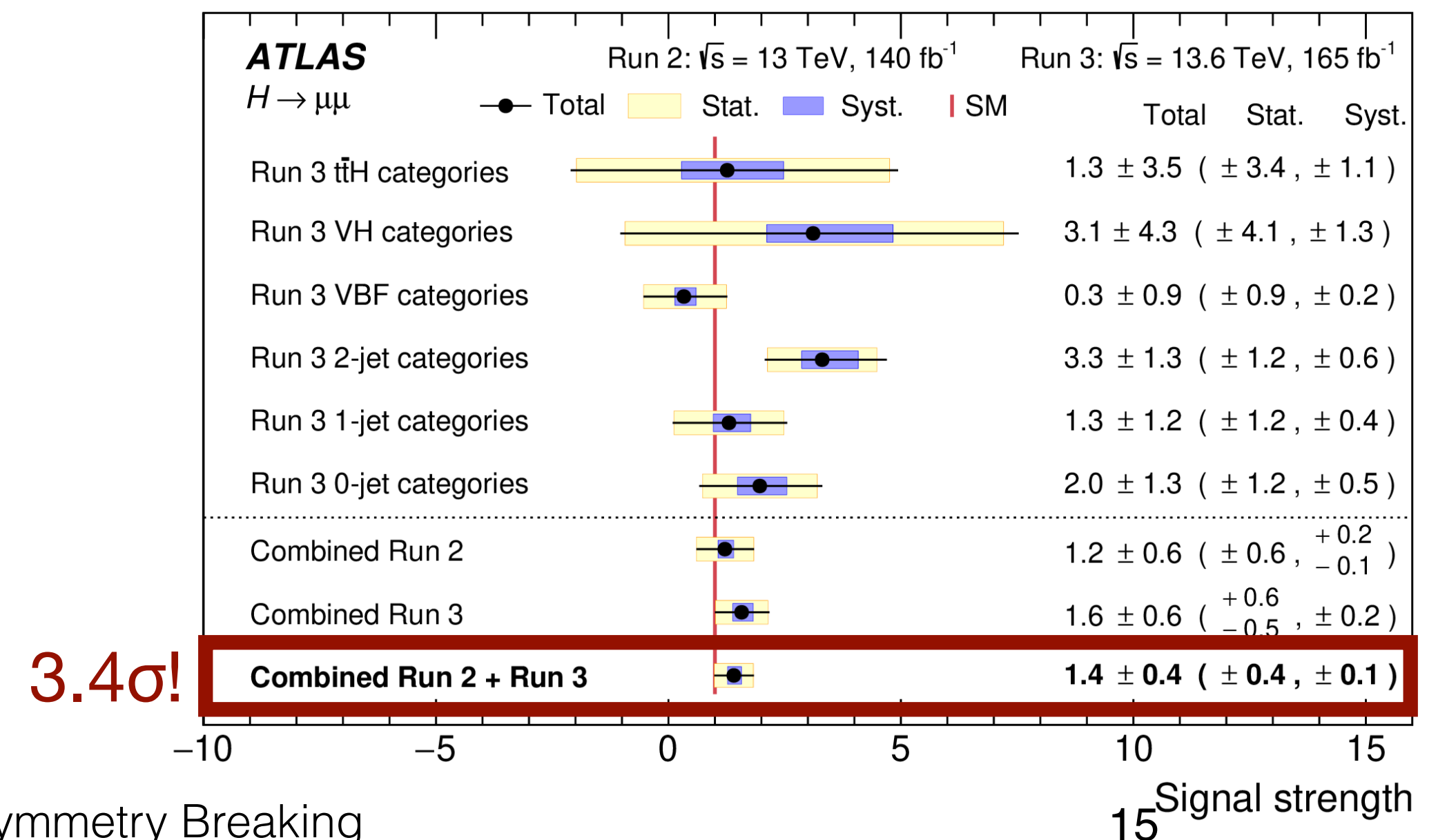
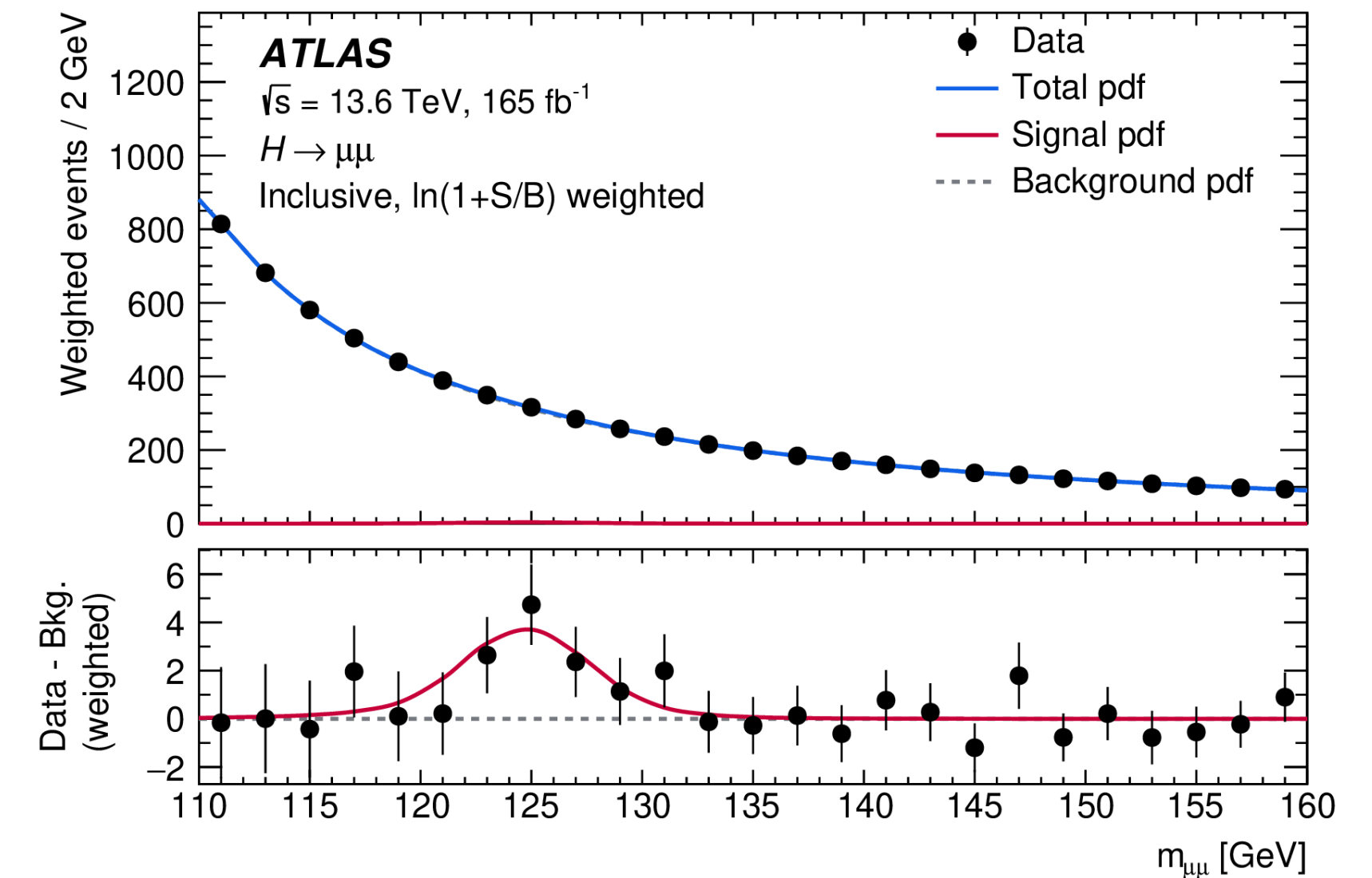
BDTs not retrained,
biggest changes in
multi-l channels

ATLAS, [arXiv:2508.14695](https://arxiv.org/abs/2508.14695)
 ATLAS webpage with Plots etc

Rare Decays: $H \rightarrow \mu\mu$

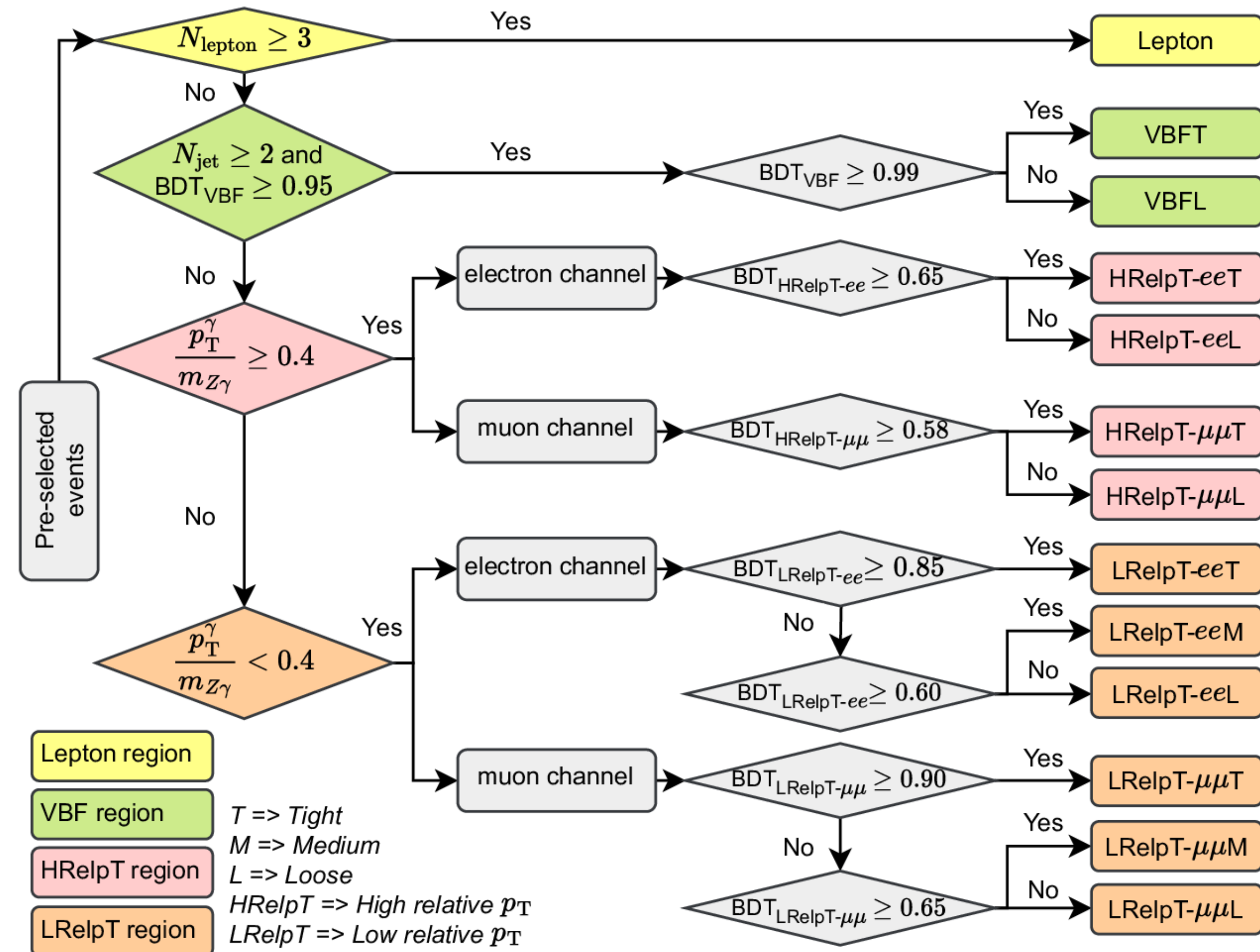
- Total Higgs production cross-section is > 50 pb, so 300 fb^{-1} of data $\Rightarrow 15 \times 10^6$ Higgs bosons produced
 - Now enough luminosity to probe rare decays: in the SM, $H \rightarrow \mu\mu$ branching ratio is $\sim 2 \times 10^{-4}$
 - That's 3000 potential $H \rightarrow \mu\mu$ events, reduced by detector acceptance and selection & reconstruction efficiency factors (47% for Run 3 data)
 - Primary background is Drell-Yan continuum $\mu\mu$ production
- New analysis improves background description (huge DY MC sample), mass resolution from improved $\mu\mu$ vertex fit; adds VH/ttH events with V/tt decaying hadronically; new NN for ttH event categorization
- 23 categories: separate events by production process and purity based on process-specific BDT or NN score
 - Fitted simultaneously

ATLAS, [arXiv:2507.03595](https://arxiv.org/abs/2507.03595)
 ATLAS webpage with plots



Rare Decays: $H \rightarrow Z\gamma$

- In the SM, $BR(H \rightarrow Z\gamma) = 1.5 \times 10^{-3}$, so 10x larger than $\mu\mu$
 - But this is loop-induced, so more sensitive to new heavy particles running in loops, and $BR(H \rightarrow Z\gamma)/BR(H \rightarrow \gamma\gamma)$ makes for a particularly sensitive probe
- Experimentally harder than $H \rightarrow \mu\mu$: to suppress very large γ +jets backgrounds, best option is $Z \rightarrow ee/\mu\mu$, but those BRs are just 3.4%
 - And decay products are “softer”, since effectively a 3-body decay
 - Classify events based on number of leptons and jets, and photon p_T
 - Background dominated by $Z\gamma$ and Z +jets



ATLAS, [arXiv:2507.12598](https://arxiv.org/abs/2507.12598)
 ATLAS webpage with plots

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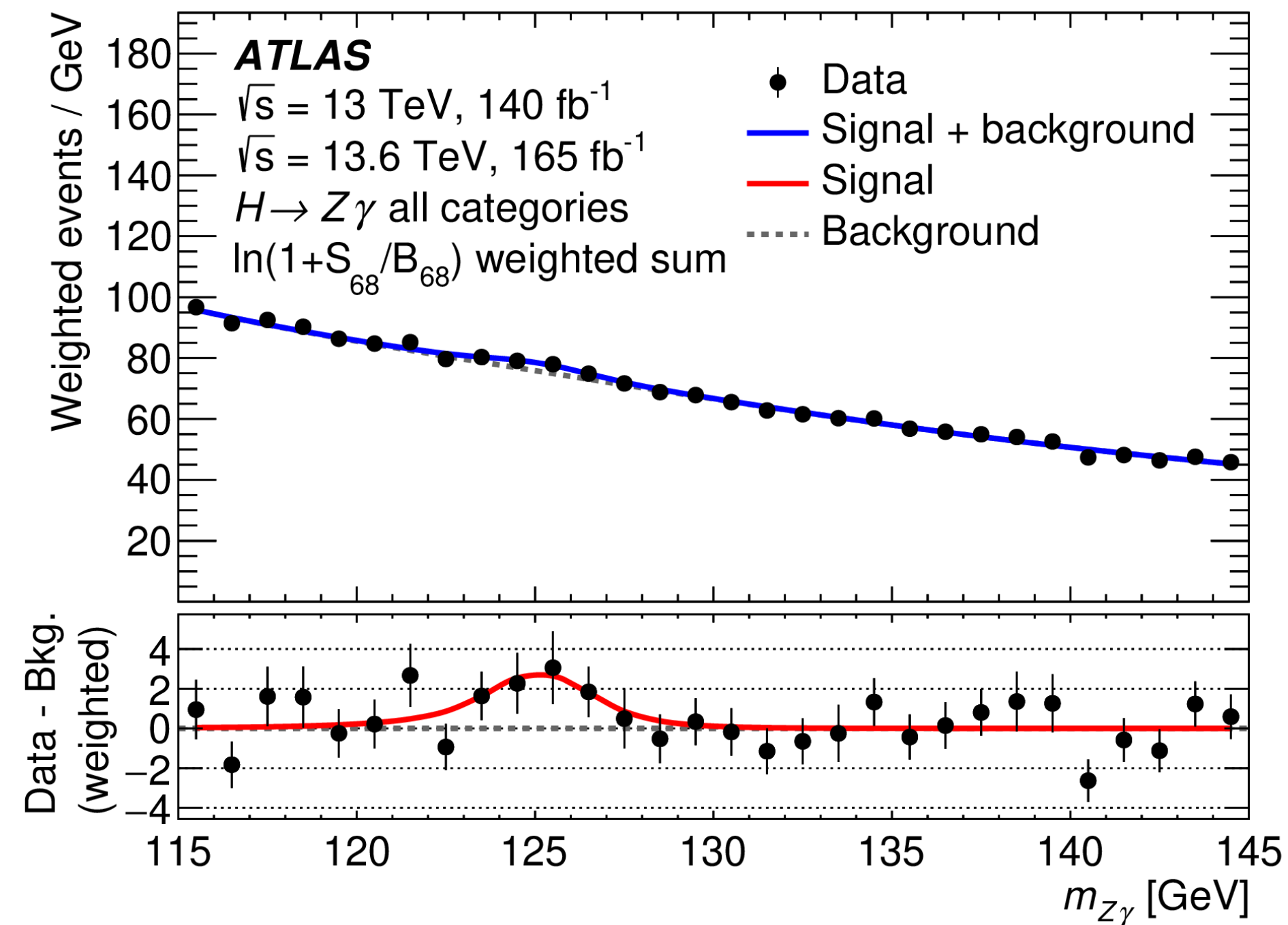
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| Category | S_{68}^{exp} | B_{68}^{exp} | N_{68} | w_{68} [GeV] | $S_{68}^{\text{exp}} / \sqrt{S_{68}^{\text{exp}} + B_{68}^{\text{exp}}}$ |
|--------------------|-----------------------|-----------------------|----------|----------------|--|
| Lepton | 1.5 ± 1.1 | 76.3 ± 2.8 | 78 | 4.4 | 0.17 |
| VBFT | 1.5 ± 1.1 | 1.2 ± 0.4 | 3 | 3.8 | 0.91 |
| VBFL | 2.8 ± 2.0 | 27.6 ± 1.8 | 23 | 4.0 | 0.51 |
| HRelpT- ee T | 1.2 ± 0.8 | 6.6 ± 0.9 | 11 | 3.1 | 0.43 |
| HRelpT- ee L | 3.0 ± 2.1 | 54.10 ± 1.8 | 77 | 4.0 | 0.40 |
| HRelpT- $\mu\mu$ T | 2.4 ± 1.7 | 20.3 ± 1.7 | 33 | 3.9 | 0.50 |
| HRelpT- $\mu\mu$ L | 2.4 ± 1.7 | 56.5 ± 1.7 | 72 | 4.1 | 0.31 |
| LRelpT- ee T | 9 ± 6 | 234 ± 6 | 251 | 3.8 | 0.57 |
| LRelpT- ee M | 29 ± 20 | $2\,591 \pm 19$ | 3\,806 | 4.1 | 0.56 |
| LRelpT- ee L | 24 ± 17 | $13\,260 \pm 50$ | 17\,435 | 4.5 | 0.21 |
| LRelpT- $\mu\mu$ T | 4.9 ± 3.4 | 96 ± 4 | 127 | 3.9 | 0.49 |
| LRelpT- $\mu\mu$ M | 34 ± 24 | $2\,545 \pm 19$ | 3\,133 | 4.1 | 0.67 |
| LRelpT- $\mu\mu$ L | 37 ± 26 | $16\,960 \pm 40$ | 19\,331 | 4.4 | 0.28 |
| Inclusive | 150 ± 110 | $35\,930 \pm 70$ | 44\,380 | 4.0 | 1.81 |

ATLAS, [arXiv:2507.12598](https://arxiv.org/abs/2507.12598)
 ATLAS webpage with plots

Rare Decays: $H \rightarrow Z\gamma$

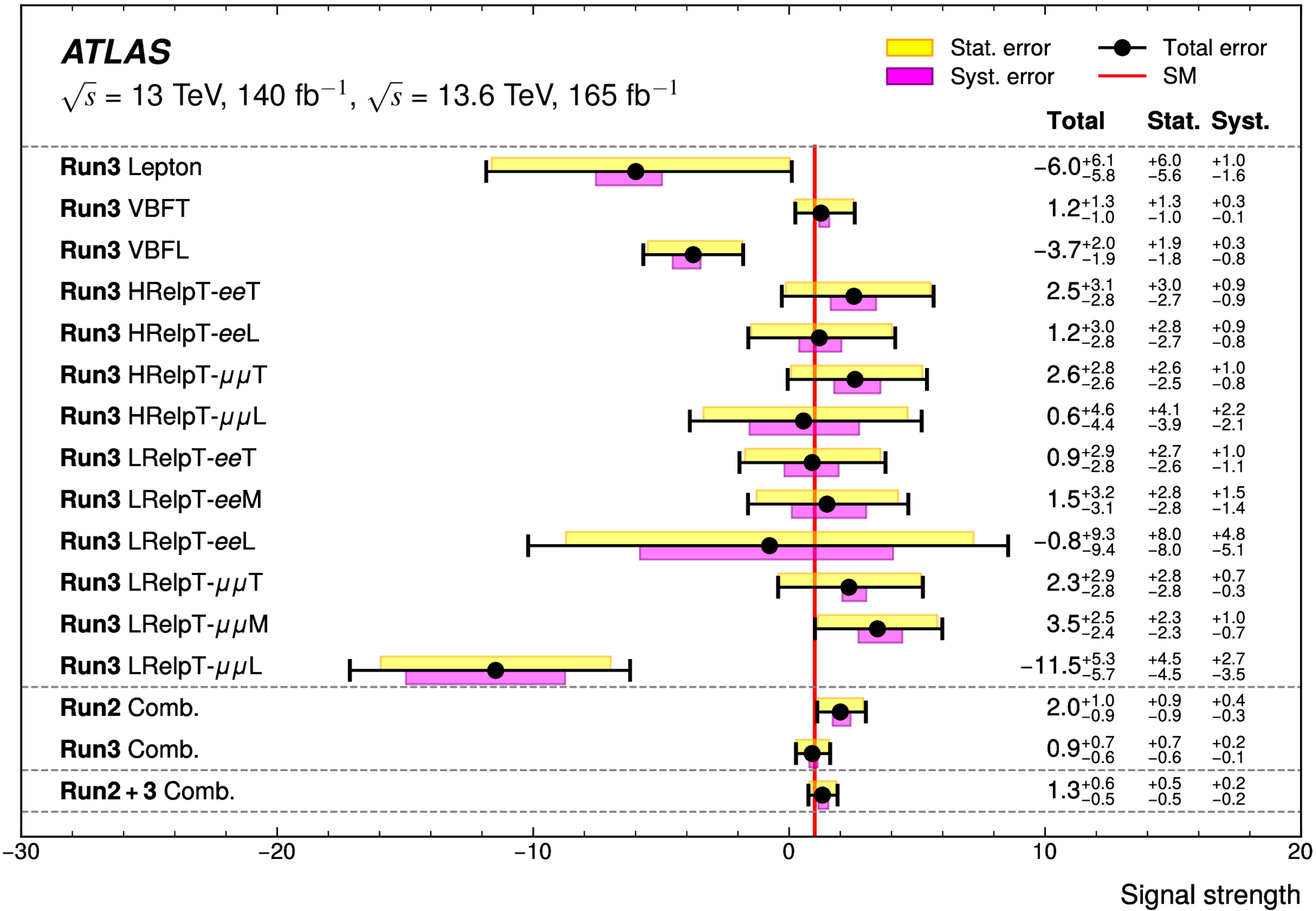
- Signal extraction done through S+B fit to $m_{Z\gamma}$ spectrum:



ATLAS, [arXiv:2507.12598](https://arxiv.org/abs/2507.12598)
[ATLAS webpage with plots](#)

- For the background, parametrize using an analytic function:
 - Require decent fit, spurious signal < 20% of expected signal statistical uncertainty
 - While minimizing the number of free parameters
- Then do unbinned ML fit to $m_{Z\gamma}$ across all categories simultaneously

Rare Decays: $H \rightarrow Z\gamma$



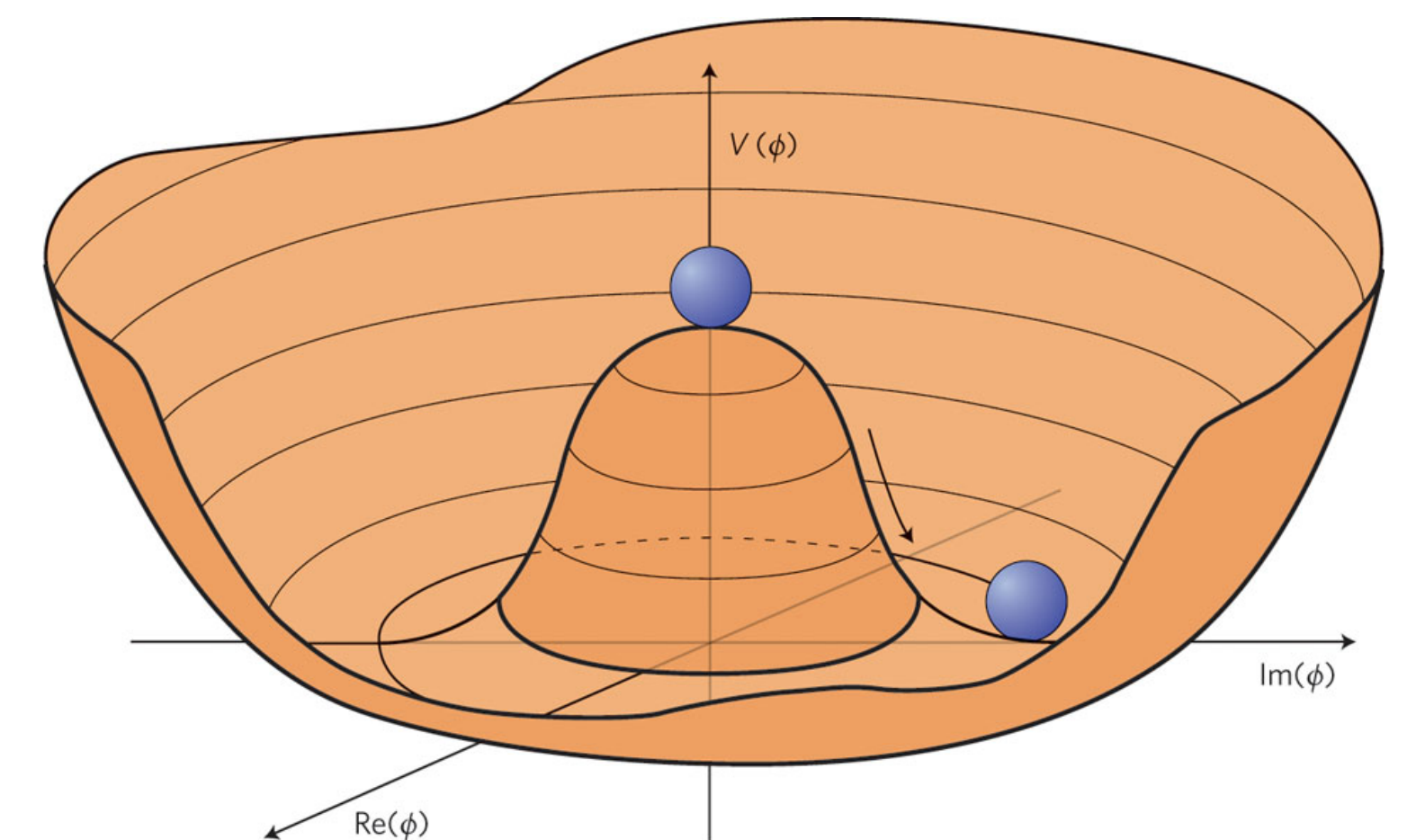
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2.5 σ

(Run 2 ATLAS+CMS
 combination was 3.4 σ)

The Higgs Potential

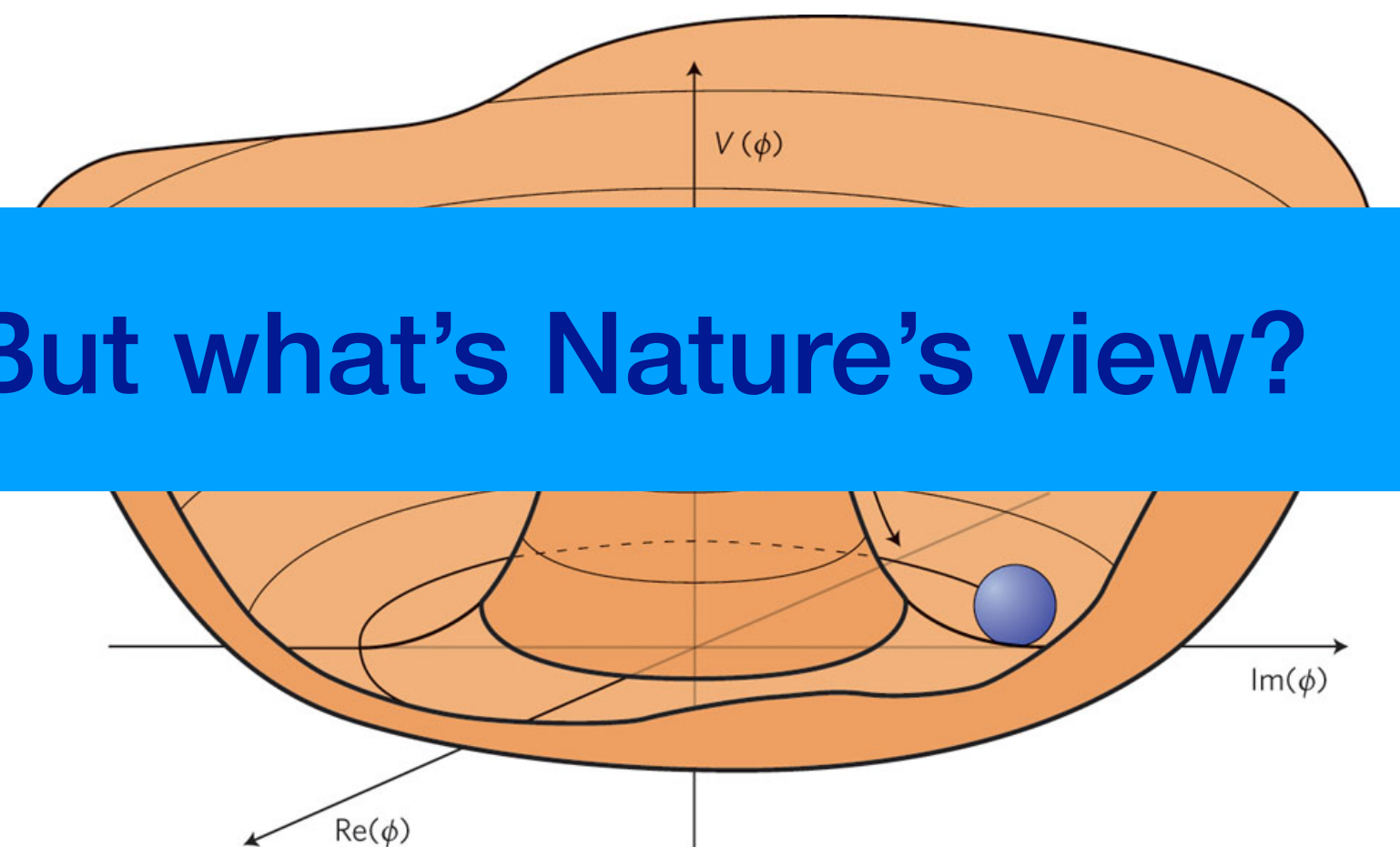
- WW scattering, and WW pair production from fermion annihilation, strongly pointed to the existence of a Higgs-like object
 - The most economical way to solve both problems
- From a more technical point of view, replacing $m_f \bar{\psi}_L \psi_R$ with $y_f \bar{\psi}_L \phi \psi_R$ requires ϕ to be (at least) a (complex) $SU(2)_L$ doublet
 - 3 degrees of freedom are the longitudinal W/Z components, the 4th the Higgs boson
- But for all this to work, and masses to be generated, the Higgs doublet needs to acquire a vacuum expectation value; need a potential term with multiple minima
 - From a human point of view, 4th order polynomial with only even terms is the most economical $\Rightarrow V = \mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2$ with $\mu^2 < 0$ and $\lambda > 0$ yields the desired result
 - m_H and v (246 GeV) determine μ and λ



The Higgs Potential

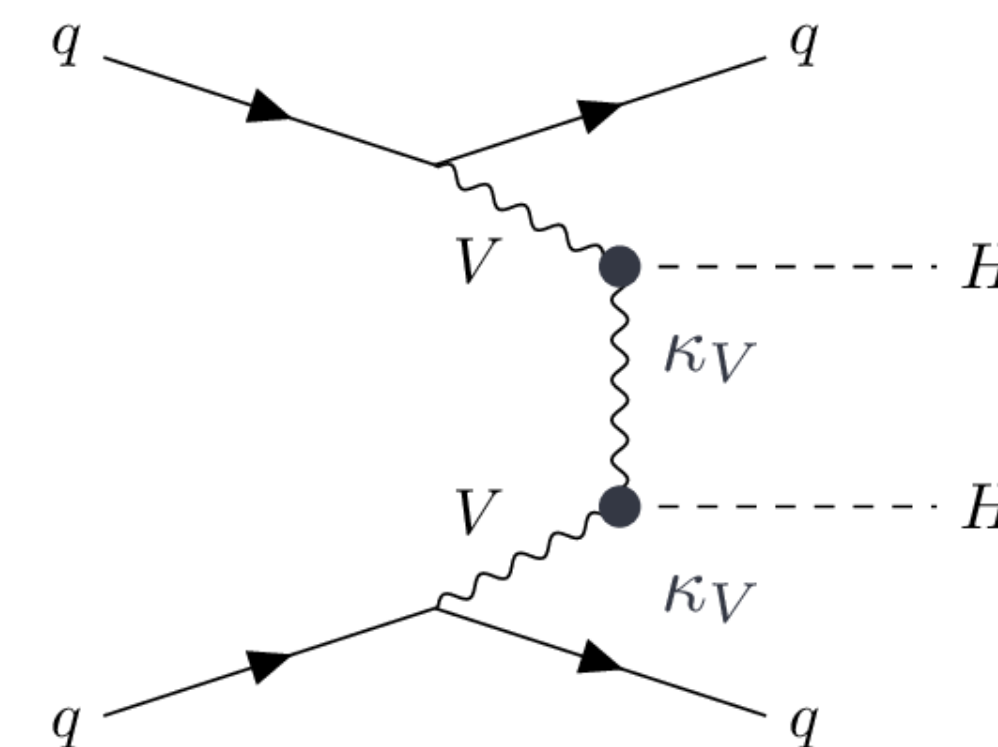
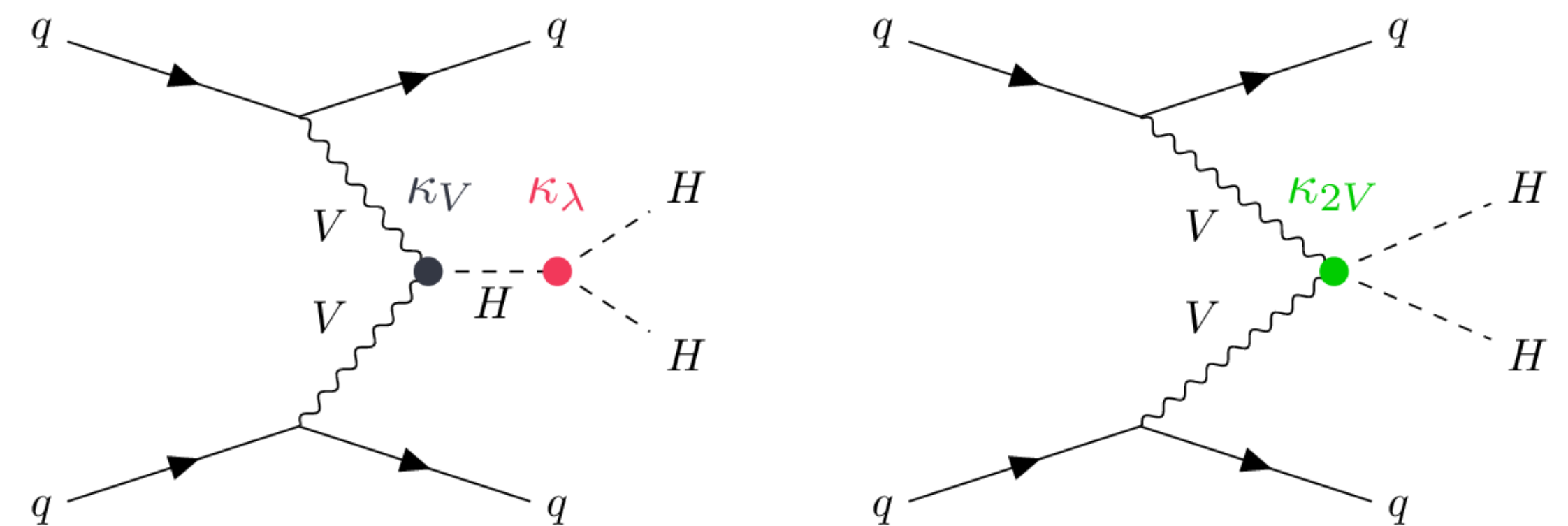
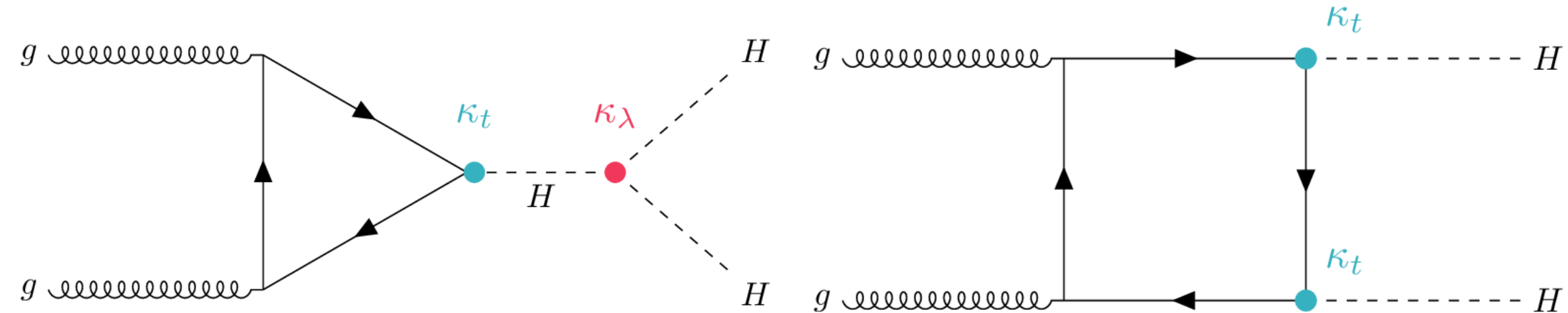
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But what's Nature's view?



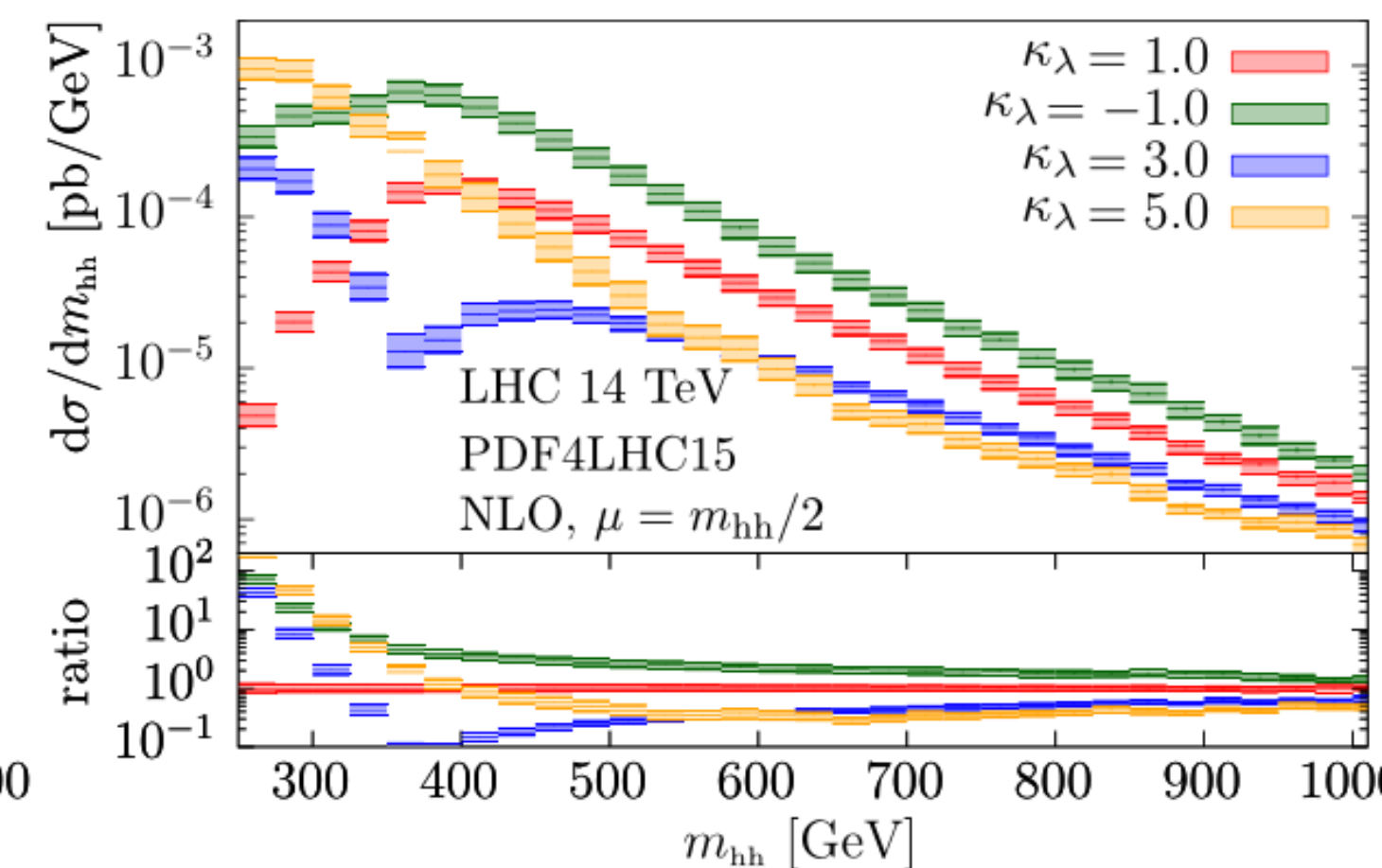
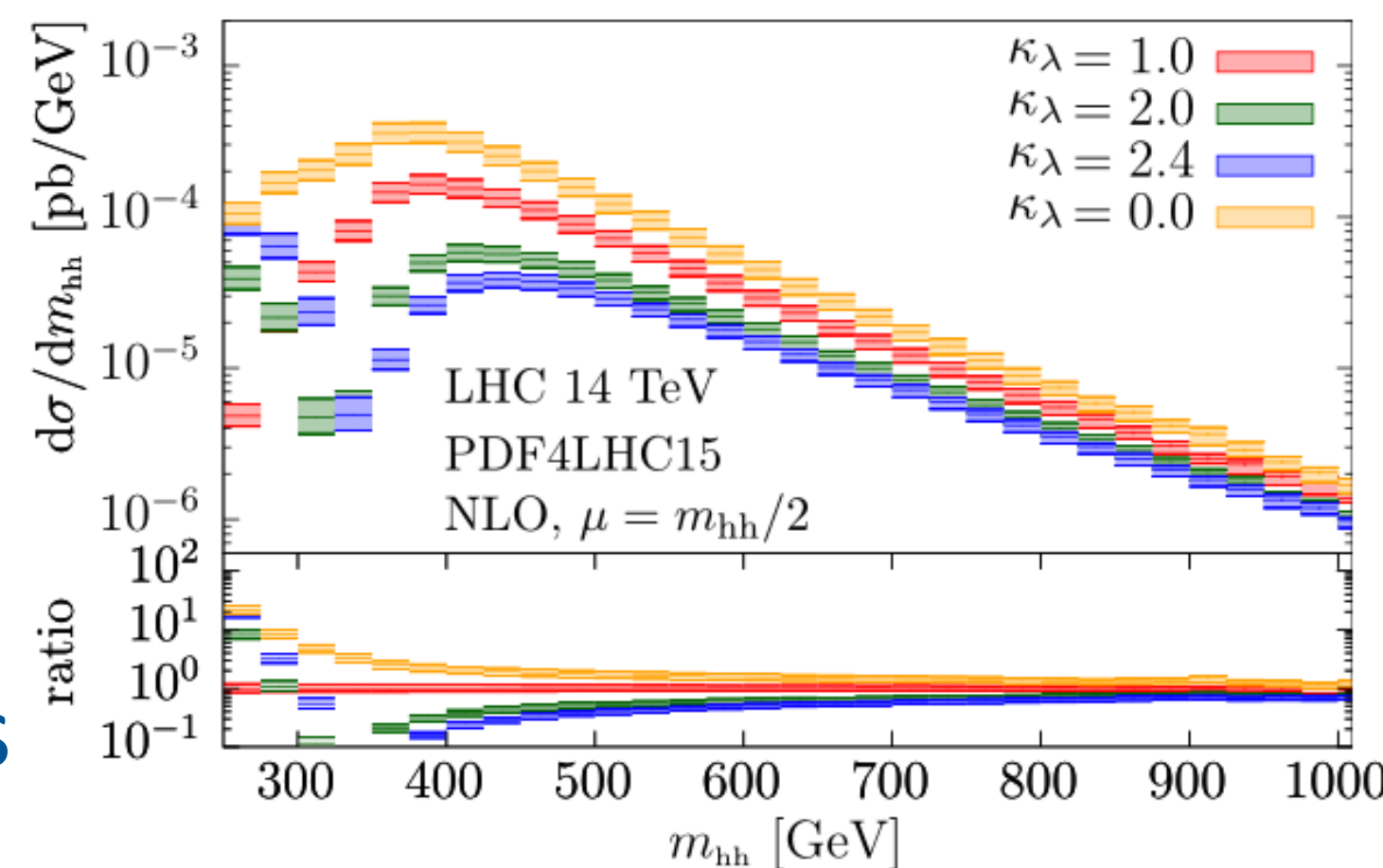
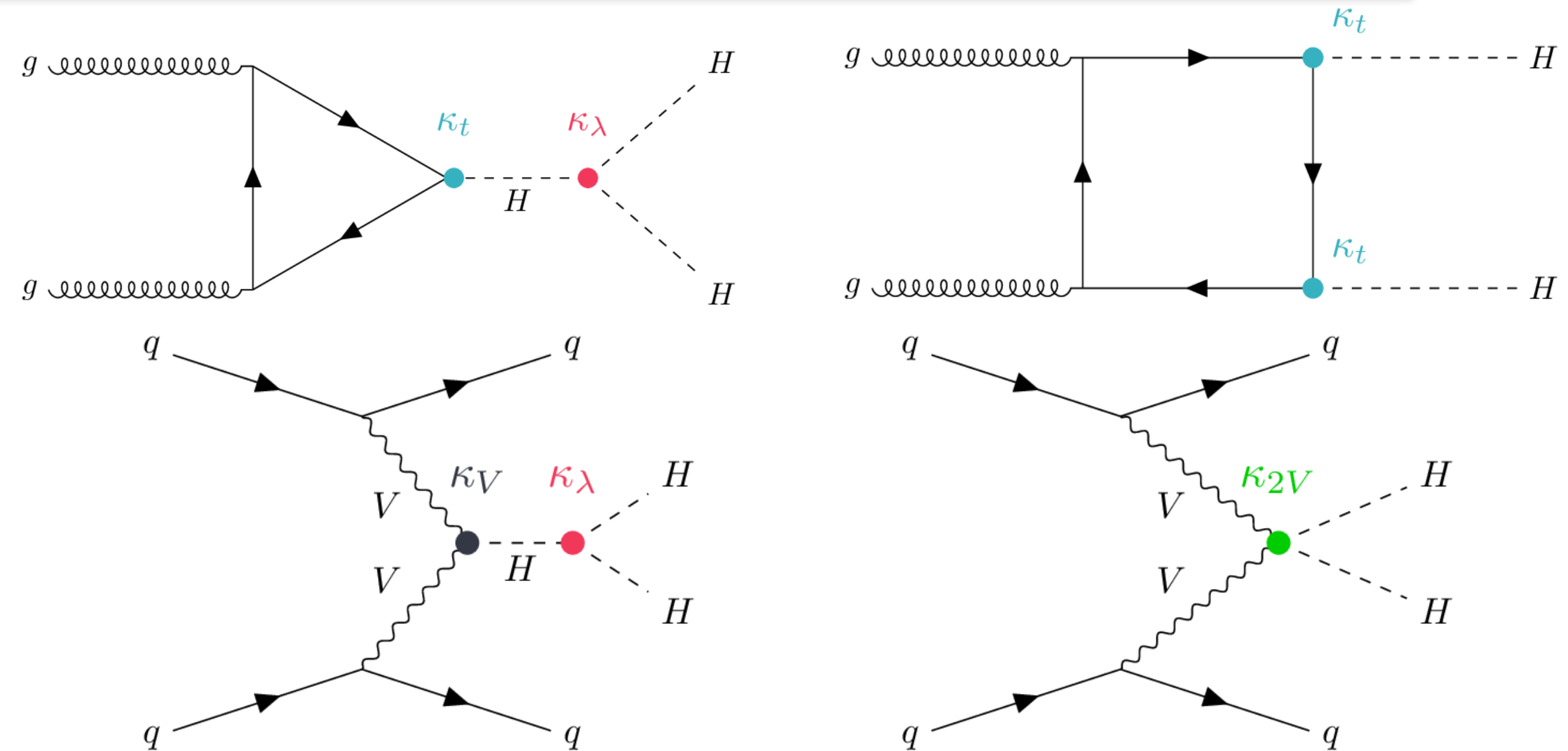
Probing $V(\phi)$: Di-Higgs Production

- Probing Higgs potential requires processes with multi-Higgs vertices
 - Most sensitive is di-Higgs production
- Whether through gluon fusion or vector boson fusion, there is destructive interference \Rightarrow tiny cross-sections (31 and 2 fb) (\Rightarrow HL-LHC!)
 - Also sensitivity to deviations from SM primarily close to threshold, hardest region to probe experimentally
 - $HH \rightarrow 4b$ yields 4 b-jets with $p = O(60 \text{ GeV})$
 - But deviation from SM typically means less destructive interference



Probing $V(\phi)$: Di-Higgs Production

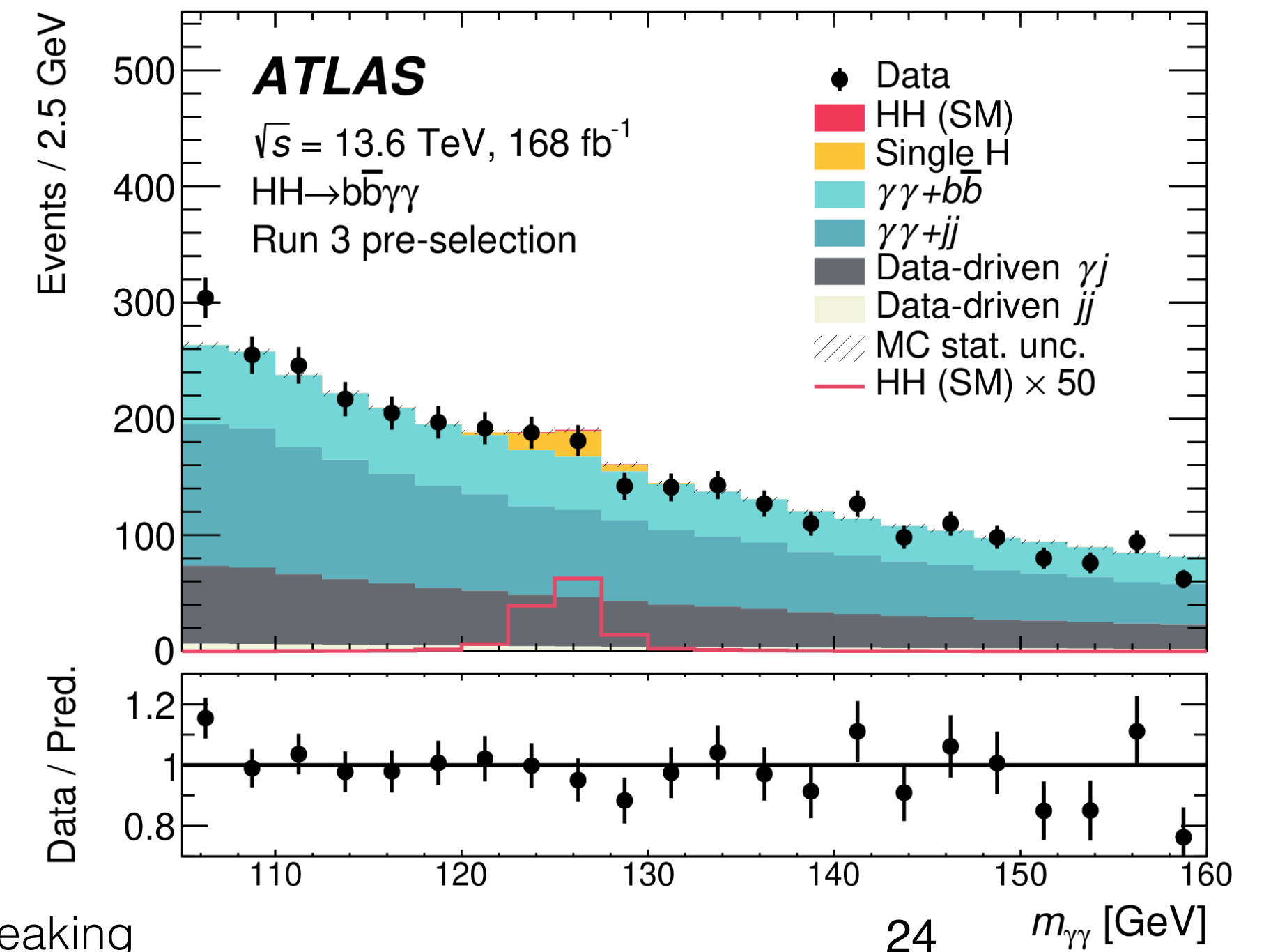
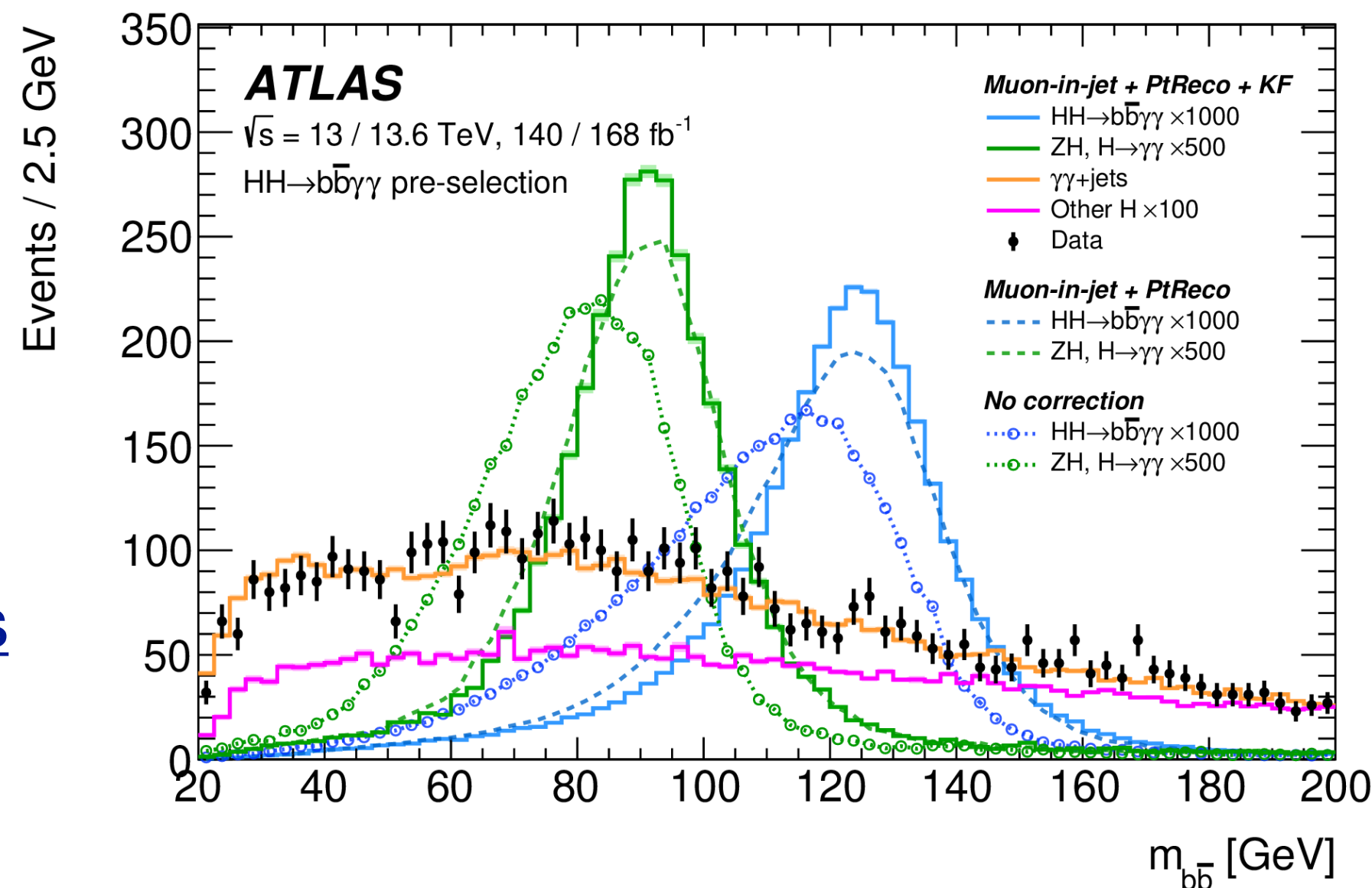
- Probing Higgs potential requires processes with multi-Higgs vertices
 - Most sensitive is di-Higgs production
- Whether through gluon fusion or vector boson fusion, there is destructive interference \Rightarrow tiny cross-sections (31 and 2 fb) \Rightarrow HL-LHC
 - Also sensitivity to deviations from SM primarily close to threshold, hardest region to probe experimentally
 - $HH \rightarrow 4b$ yields 4 b-jets with $p = O(60 \text{ GeV})$
 - But deviation from SM typically means less destructive interference



B. Di Micco, M. Gouzevitch, J. Mazzitelli, C. Vernieri (eds) et al, [arXiv:1910.00012](https://arxiv.org/abs/1910.00012)

Di-Higgs with 300 fb⁻¹

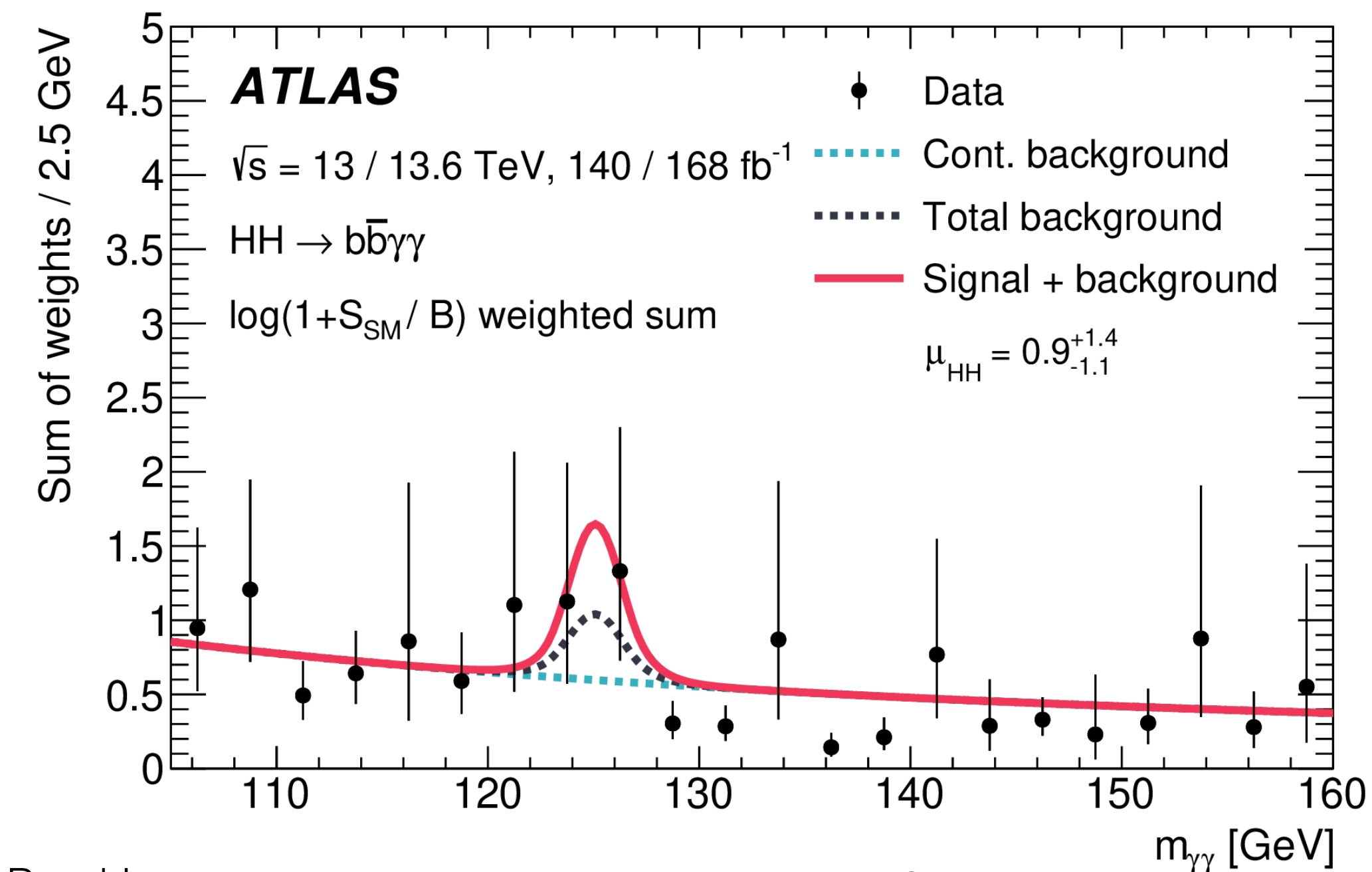
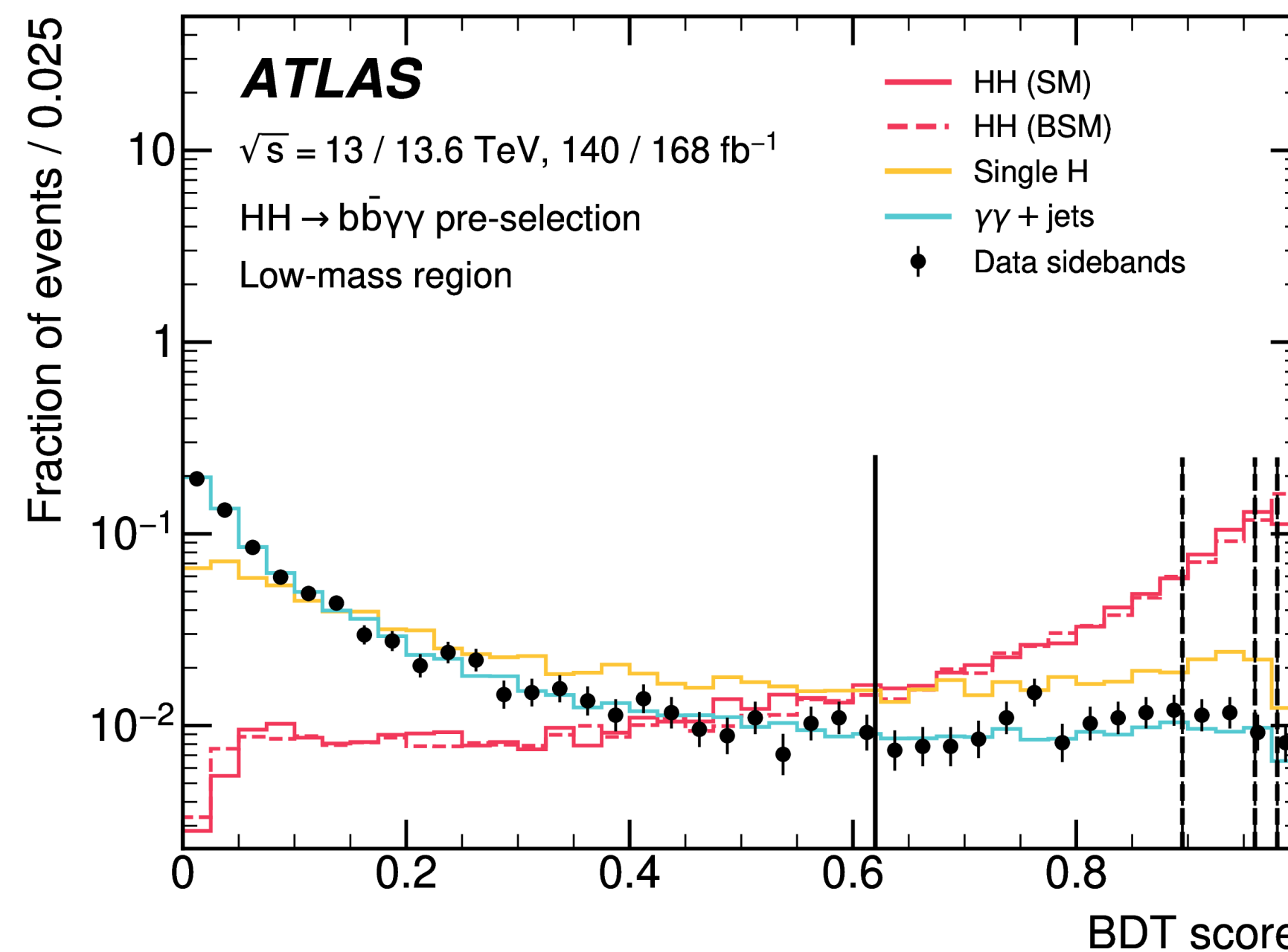
- New ATLAS bby γ result
 - bby γ channel has the best sensitivity near threshold: can trigger on photons and exploit excellent diphoton mass resolution
 - Dominant background is $\gamma\gamma$ + (b-)jets
 - Use kinematic fit to improve bb mass: aim for event balance in transverse plane, including (in a second step) a term $(m_{bb} - 125 \text{ GeV})^2$



ATLAS, [arXiv:2507.03495](https://arxiv.org/abs/2507.03495)
 ATLAS webpage with plots

Di-Higgs with 300 fb⁻¹

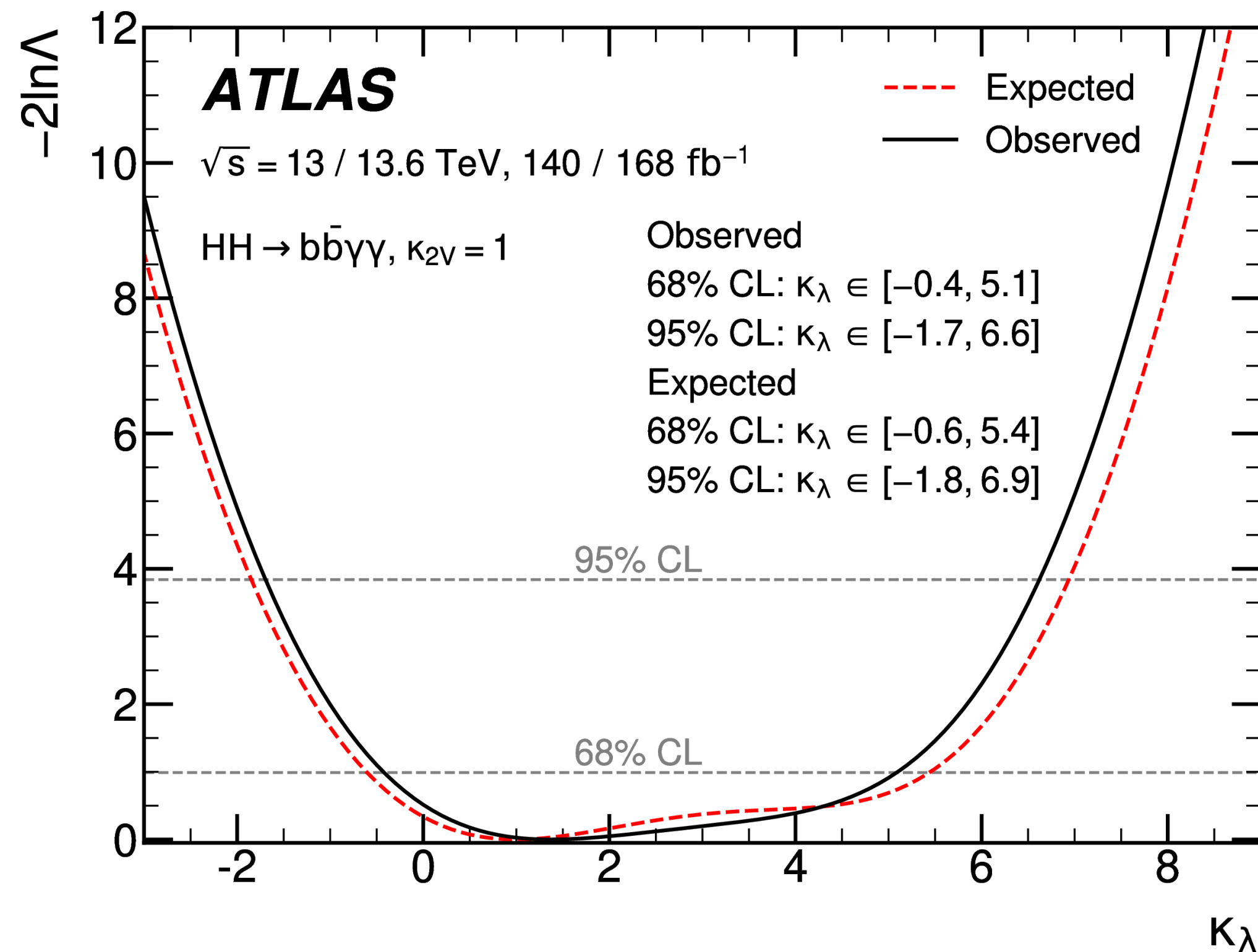
- New ATLAS bby $\gamma\gamma$ result
 - Split sample in two mass regions (for sensitivity to κ_λ and $\kappa_{2\nu}$), then use BDTs to separate signal from background (fitted from side bands/simulation), and define regions with different purities
 - Final fit to all categories simultaneously
 - Find $\mu_{HH} = 0.9^{+1.4}_{-1.1}$ (statistics dominated)



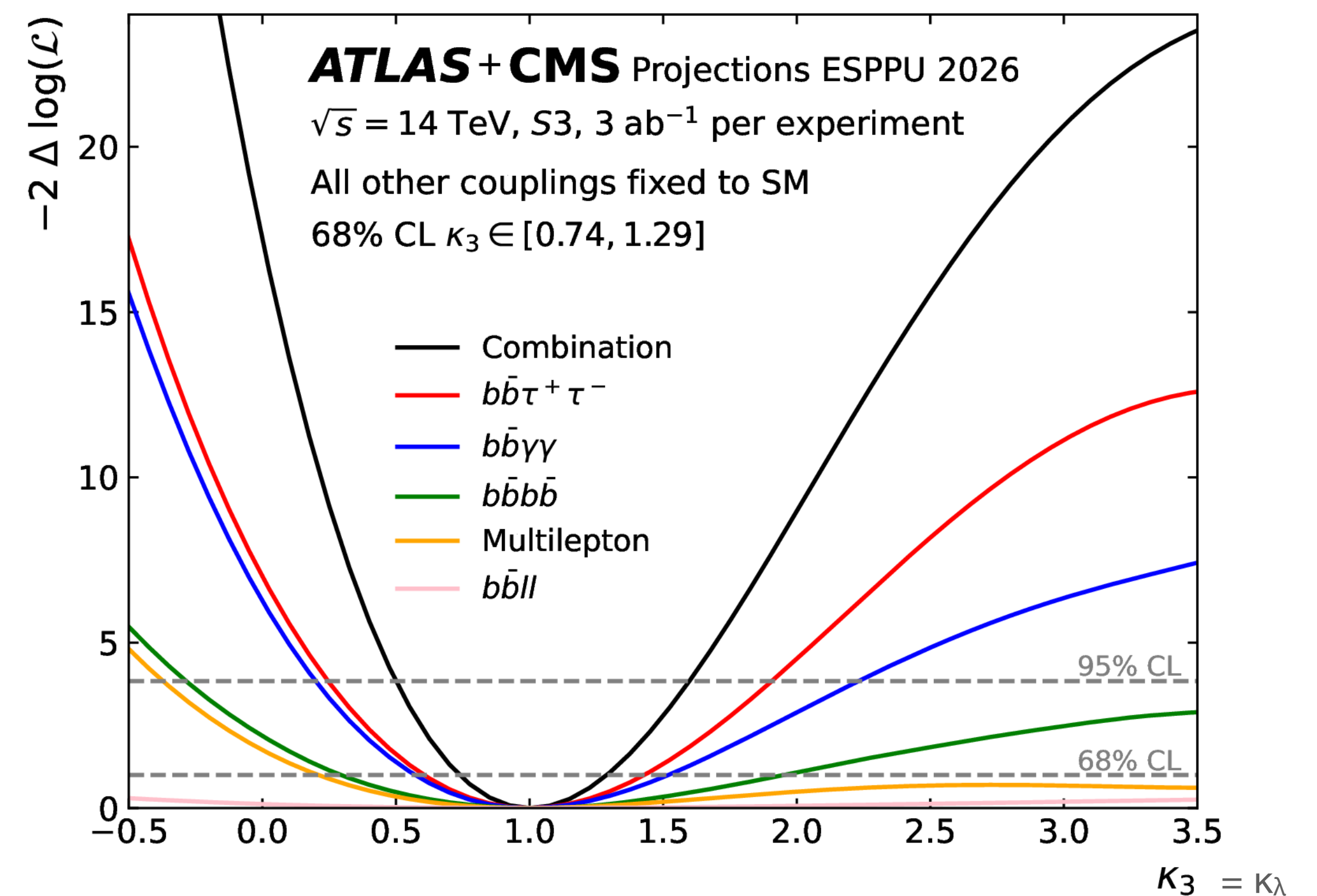
[ATLAS, arXiv:2507.03495](#)
[ATLAS webpage with plots](#)

Di-Higgs with 300 fb⁻¹, 3 ab⁻¹

- Still improving analyses, more channels, and 10x more data to come with HL-LHC



ATLAS, [arXiv:2507.03495](https://arxiv.org/abs/2507.03495)
[ATLAS webpage with plots](#)



ATLAS & CMS
[ATL-PHYS-PUB-2025-018](#)

Closing Remarks

- Higgs boson was first observed a little over a decade ago
 - Most economical way to realize electroweak symmetry breaking
- We are very lucky its mass is 125 GeV
 - Almost hardest point for discovery... but many observable production and decay channels
 - Now extensively probed: couplings to SM interaction bosons and third generation fermions are within $\sim 15\%$ of SM predictions
 - And we are starting to probe couplings to muon and charm quark
 - Also 10-120x more data coming!
- Next big endeavor is probing Higgs potential
 - SM prediction is just “simplest” form that meets requirements
 - Observing multi-Higgs production is key to probing $V(\phi)$
 - Improving analyses, so far consistent with SM, but uncertainties are large
 - BSM typically leads to larger total cross-section, but in kinematic range that’s hardest to probe

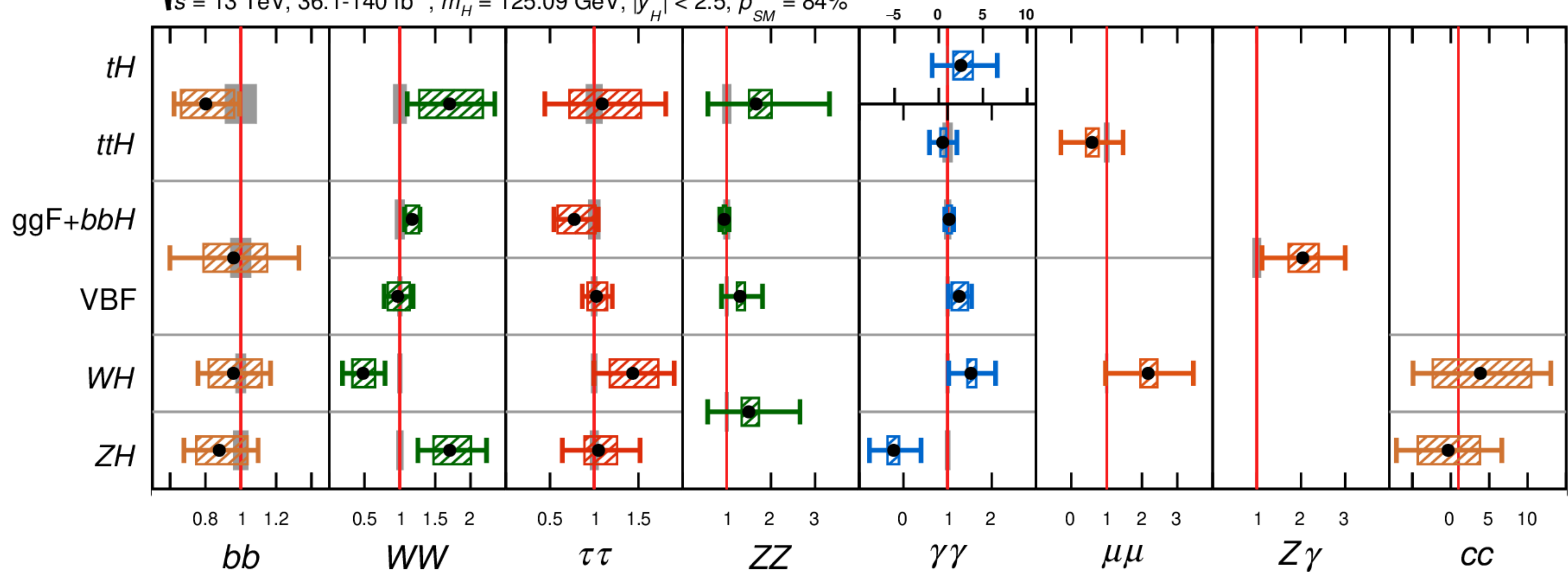
Just In Case

At Detailed Level

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$, $36.1\text{-}140 \text{ fb}^{-1}$, $m_H = 125.09 \text{ GeV}$, $|y_H| < 2.5$, $p_{SM} = 84\%$

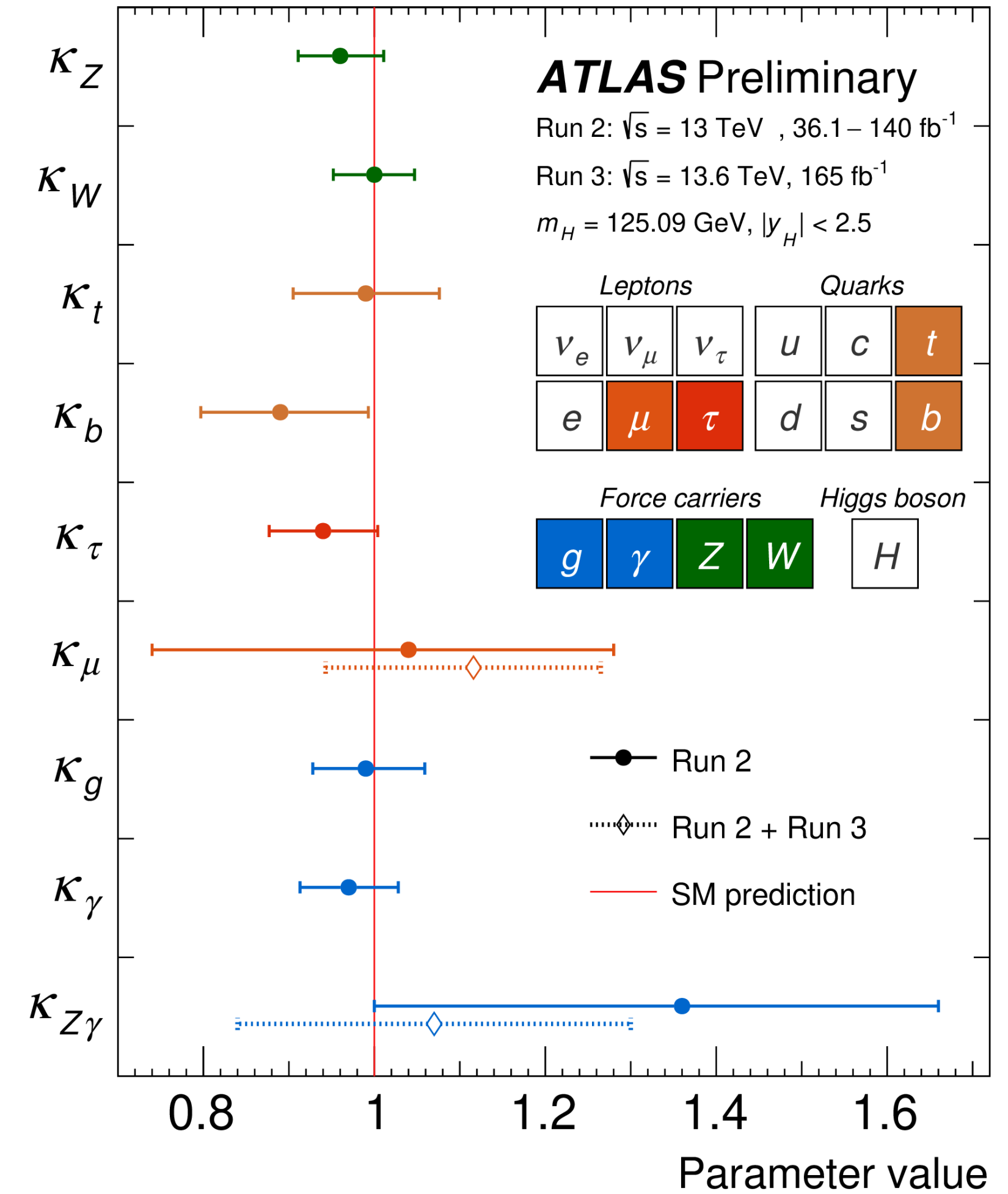
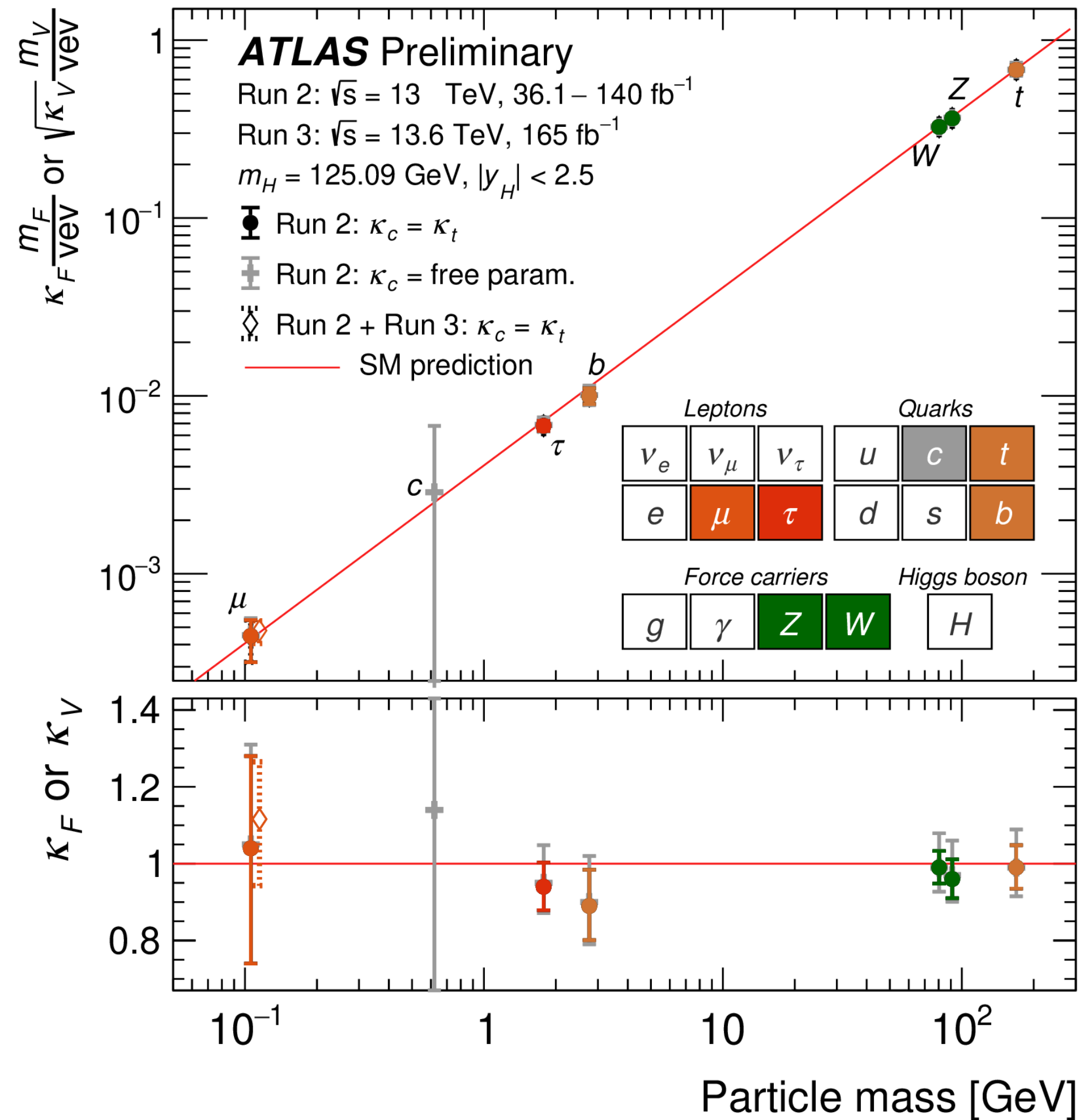
● Data (Total Unc.) ▨ Syst. Unc. ■ SM prediction



$\sigma \times B$ normalized to SM prediction

[ATLAS-CONF-2025-006](#)
[ATLAS Webpage with Plots](#)

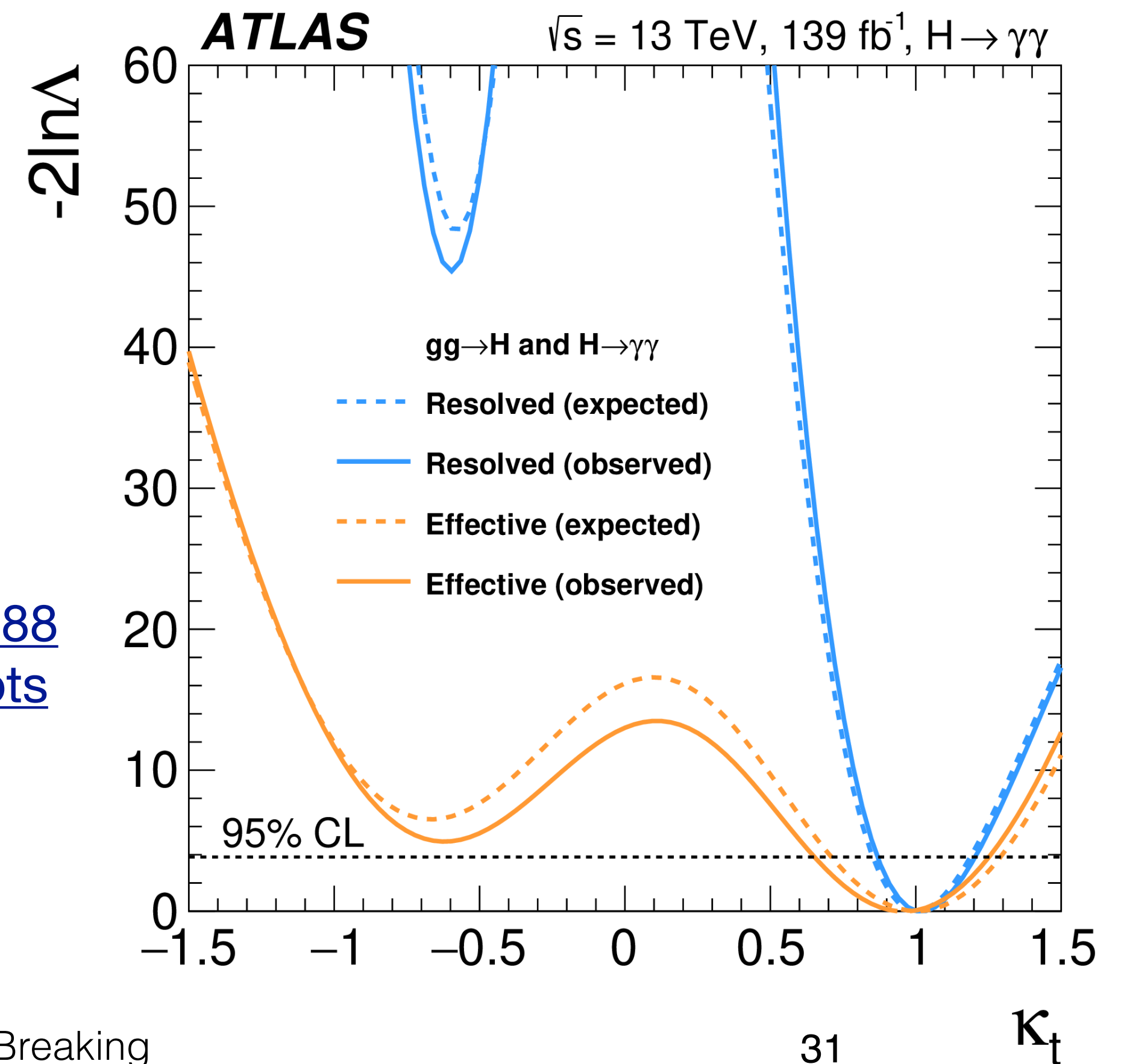
Couplings, Run 2 + Early Run 3



κ_t from $H \rightarrow \gamma\gamma$

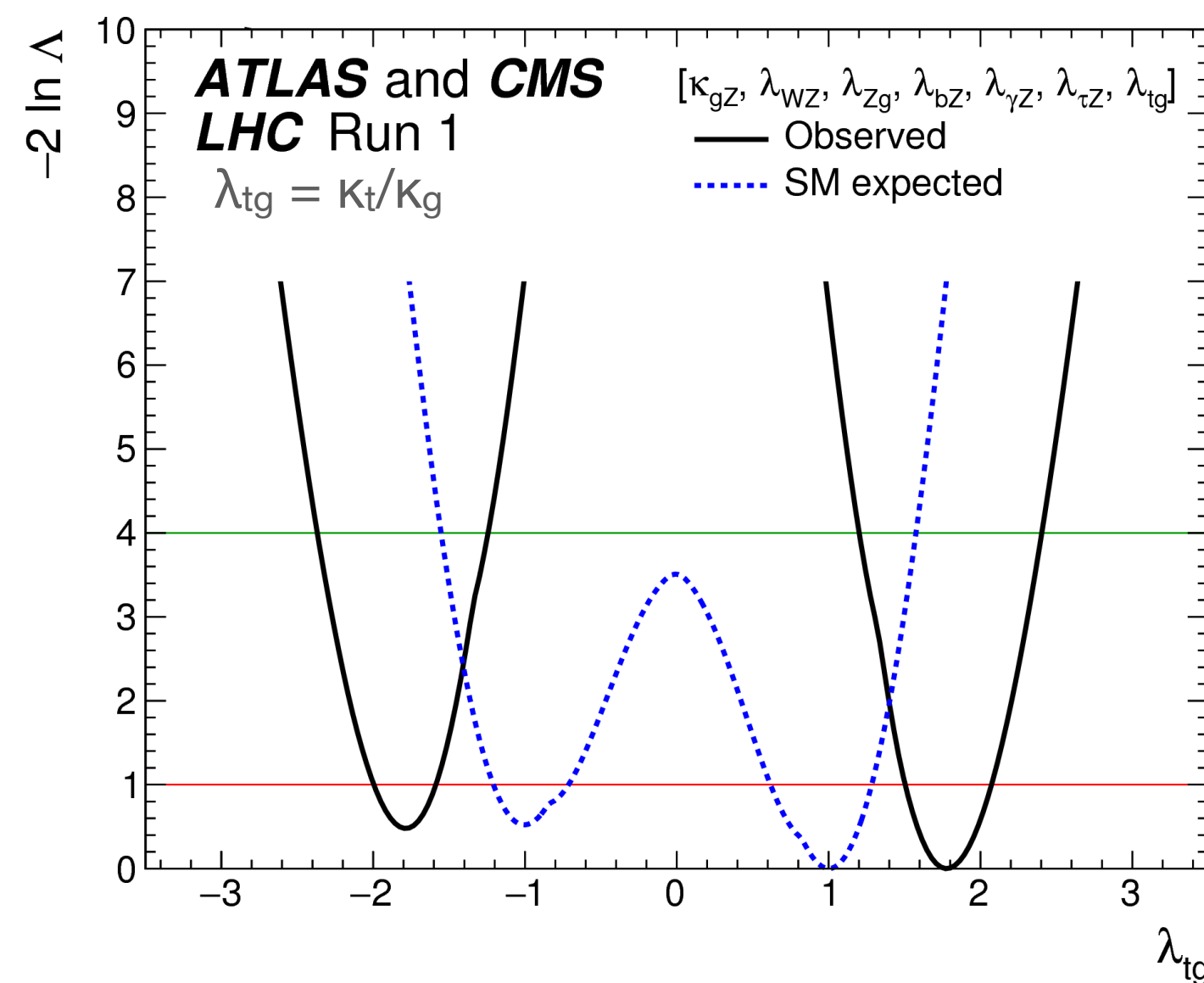
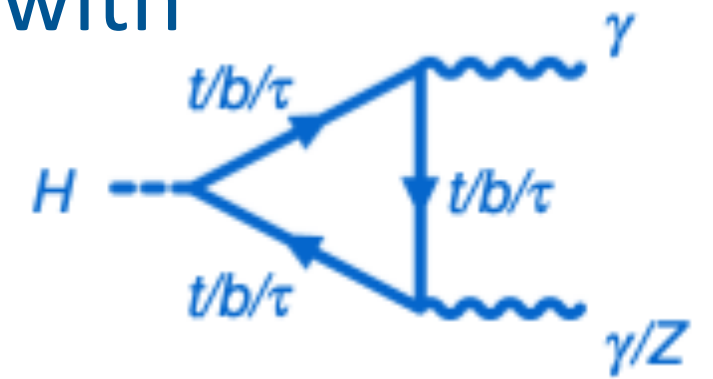
- Can parametrize $gg \rightarrow H$ and $H \rightarrow \gamma\gamma$ as a function of the couplings in the loop, and constrain κ_t , fixing all other couplings to their SM values
- Or can parametrize using effective κ_g and κ_γ and extract κ_t (with sensitivity mainly from tH production)
- $\kappa_t < 1$ is excluded by first approach

ATLAS, [JHEP 07 \(2023\) 088](#)
[ATLAS webpage with plots](#)

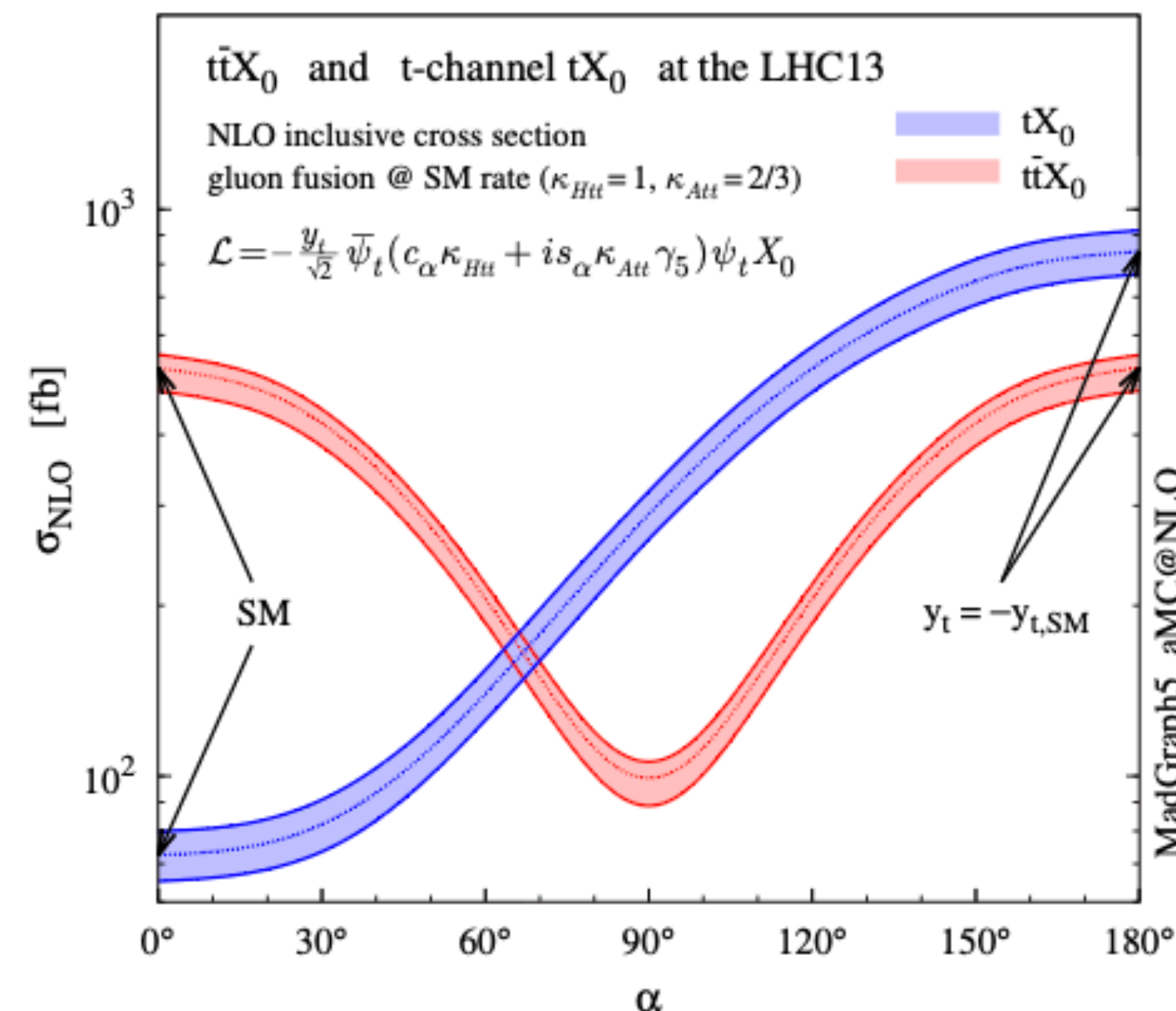


tH Production: y_t Sign

- Higgs coupling strength to the top quark (y_t) gets probed through multiple processes: not just $t\bar{t}H$, but also e.g. $H \rightarrow \gamma\gamma$ or gluon fusion H production (top loop)
 - These processes are mainly sensitive to $(y_t)^2$ - to probe a coupling's sign one needs a process with interference
 - $H \rightarrow \gamma\gamma$ (t-b interference) and ggF H production (t-W) have some sensitivity
- tH production however, is small in the SM *because* of significant destructive interference: with opposite sign the interference becomes constructive and $\sigma(tH)$ increases x 10



ATLAS+CMS Run 1 (7-8 TeV), no tH, JHEP 08 (2016) 045

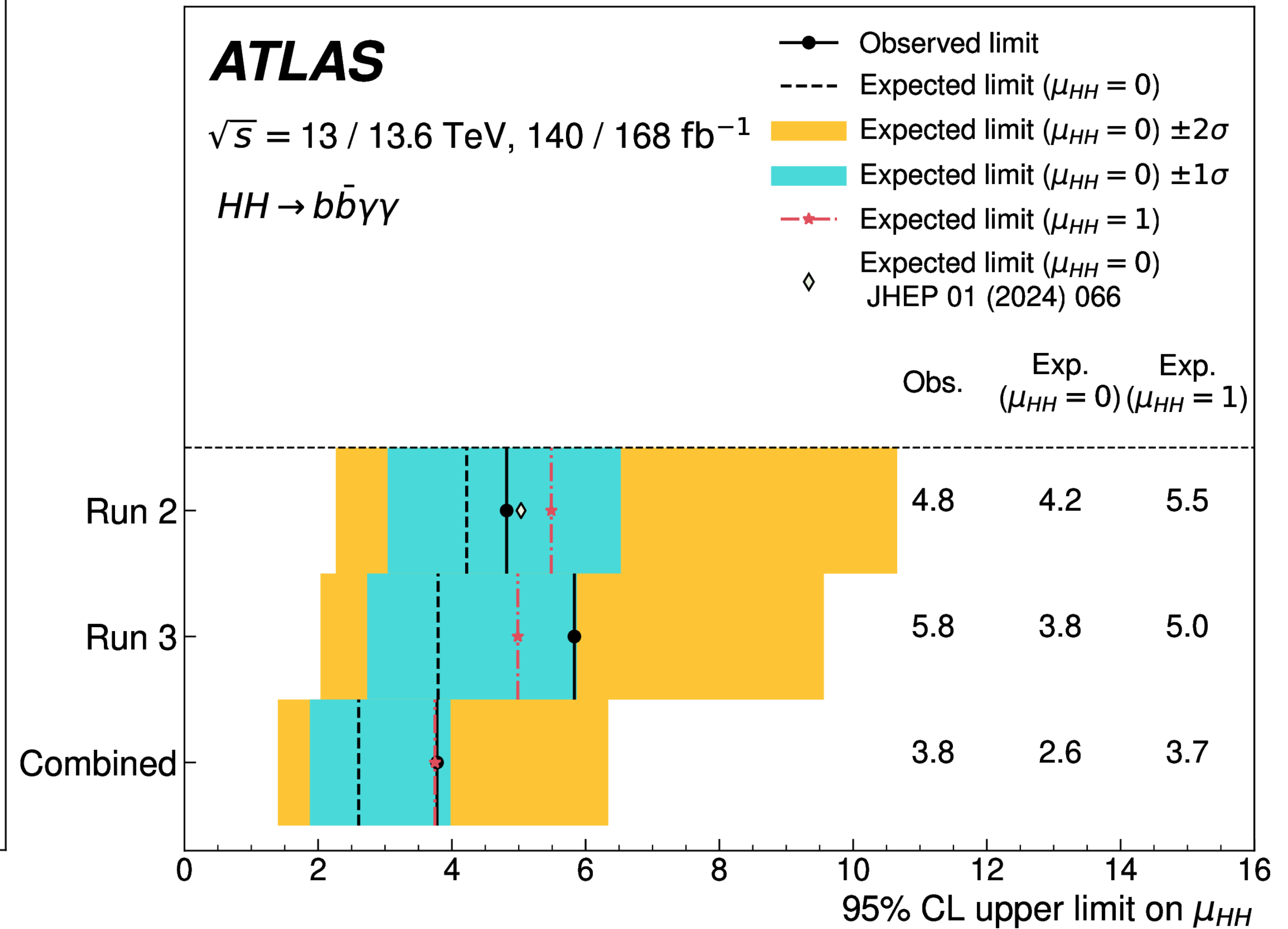
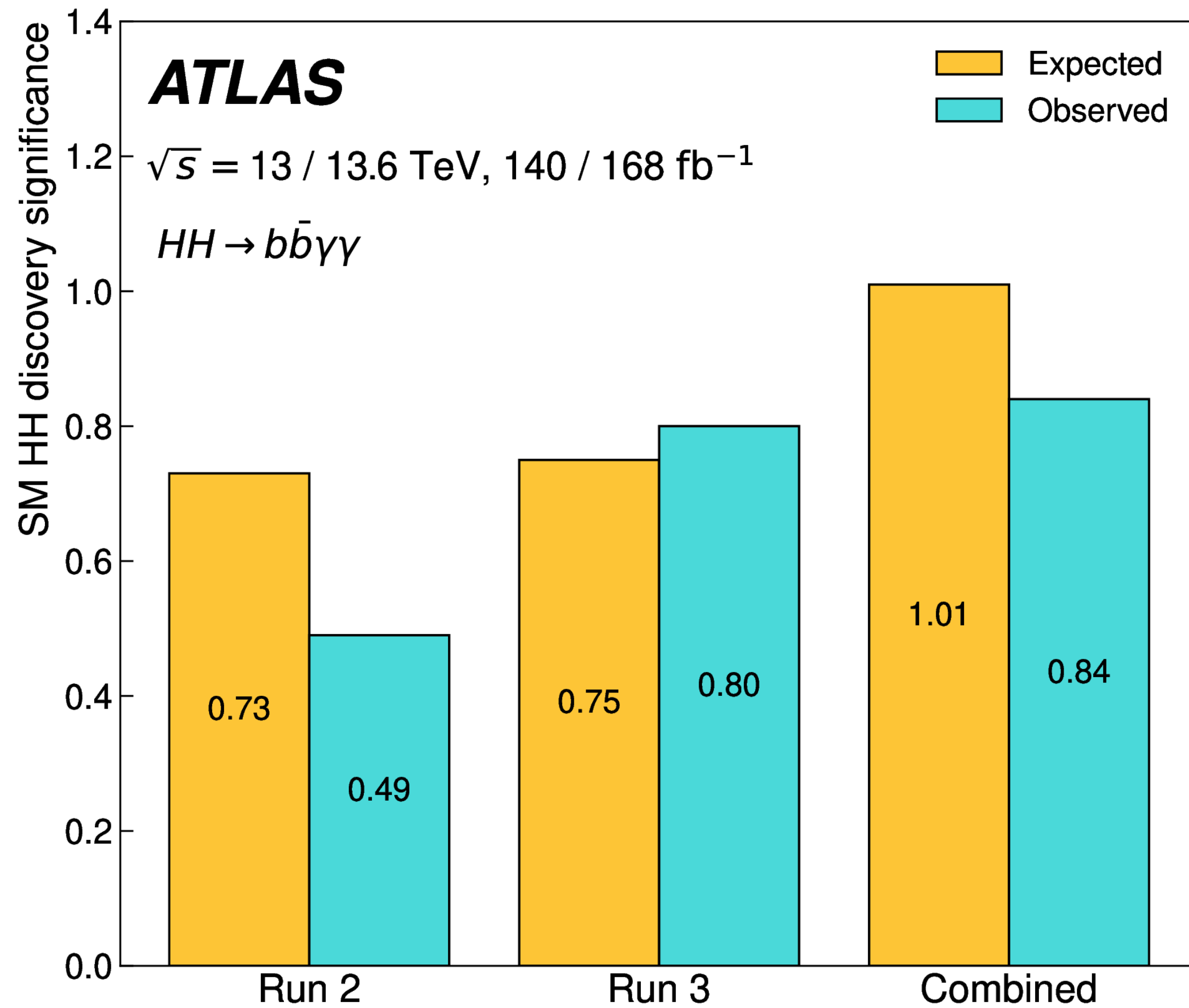


From Demartin et al, Eur.Phys.J.C 75 (2015) 6, 267

ATLAS Search for tH Production

| Channel | Region | BDT score | Conversion cut | Jets | Leptons flavour (p_T ordered) | Other |
|--------------------------|----------------------------|--|----------------|---|--|---------------------------------|
| $H \rightarrow b\bar{b}$ | SR | - | - | $N_{\text{not-}b}^{70} < 2$ $N_{\text{forward}}^{\text{jet}} \geq 1$ | - | - |
| | CR($t\bar{t} + \geq 1b$) | - | - | $N_{\text{not-}b}^{70} \geq 2$ $N_{\text{forward}}^{\text{jet}} = 0$ | - | - |
| 2ℓ SS | SR | BDT(tHq) > 0.65 BDT($t\bar{t}$) < 0.5 BDT($t\bar{t}W$) < 0.6 and BDT(VV) < 0.8 | yes | - | - | - |
| | CR(μ_{HF}) | BDT(tHq) < 0.65 BDT($t\bar{t}$) > 0.3 BDT($t\bar{t}W$) < 0.6 and BDT(VV) < 0.9 | yes | - | $\mu\mu$ | - |
| | CR(e_{HF}) | BDT(tHq) < 0.65 BDT($t\bar{t}$) > 0.3 BDT($t\bar{t}W$) < 0.6 and BDT(VV) < 0.9 | yes | - | $\mu e, ee$ | $H_T(\ell) < 225 \text{ GeV}$ |
| | CR(e_{conv}) | BDT($t\bar{t}$) > 0.3 | inverted | - | $\mu e, ee$ | $m(\ell\ell) < 150 \text{ GeV}$ |
| | CR($t\bar{t}W$) | BDT($t\bar{t}W$) > 0.6 BDT($t\bar{t}$) < 0.3 | yes | - | - | - |
| 3ℓ | SR | BDT(tHq) > 0.7 BDT($t\bar{t}$) < 0.9 BDT($t\bar{t}W$) < 0.8 | yes | - | - | - |
| | CR(μ_{HF}) | BDT(tHq) < 0.7 BDT($t\bar{t}$) > 0.5 BDT($t\bar{t}W$) < 0.8 | yes | - | $\mu\mu\mu, e\mu\mu,$ $\mu e\mu, ee\mu$ | - |
| | CR(e_{HF}) | BDT(tHq) < 0.7 BDT($t\bar{t}$) > 0.5 BDT($t\bar{t}W$) < 0.8 | yes | - | $\mu\mu e, e\mu e$ $\mu ee, eee$ | - |
| | CR(e_{conv}) | - | inverted | - | $\mu\mu e, e\mu e,$ $\mu ee, eee$ | - |
| | CR($t\bar{t}W$) | BDT($t\bar{t}W$) > 0.8 | yes | - | - | - |

Di-Higgs



Di-Higgs at HL-LHC

$$V(H) = \frac{1}{2}m_H^2 H^2 + \lambda_3 v H^3 + \frac{1}{4}\lambda_4 H^4$$

$$\lambda_3 = \lambda_4 = \lambda_{\text{SM}} = \frac{m_H^2}{2v^2}$$

