Precision predictions for hadroproduction processes beyond the Standard Model

### Benjamin Fuks (IPHC Strasbourg / Université de Strasbourg)

In collaboration with G. Bozzi, J. Debove, M. Klasen, F. Ledroit, Q. Li & J. Morel

Physics seminar @ Centre de Physique des Particules de Marseille November 2, 2009

# Outline

### Models and motivation

Motivation for resummation calculations

### 2

Parton showers and resummation

- Parton showers
- Transverse-momentum, threshold and joint resummation formalisms
- Matching to the fixed order

### 3

#### Numerical results, with uncertainties

- The Drell-Yan and the Tevatron
- Grand Unified Theories and Z' bosons
- The Minimal Supersymmetric Standard Model (MSSM)

### 4 Summary - conclusions

# Outline

### Models and motivation

Motivation for resummation calculations

### 2

#### Parton showers and resummation

- Parton showers
- Transverse-momentum, threshold and joint resummation formalisms
- Matching to the fixed order

#### <sup>3</sup> Numerical results, with uncertainties

- The Drell-Yan and the Tevatron
- Grand Unified Theories and Z' bosons
- The Minimal Supersymmetric Standard Model (MSSM)

#### Summary - conclusions

# Simple questions... and a proposal for answers

• One of the LHC purposes: which model of new physics is the correct one?

- \* We need data [which are hopefully coming this next year].
- \* We need theoretical predictions.
- \* Reliable predictions seem reasonnable. [that's the aim of this talk].

Confront data and theory.

# Simple questions... and a proposal for answers

• One of the LHC purposes: which model of new physics is the correct one?

- \* We need data [which are hopefully coming this next year].
- \* We need theoretical predictions.
- \* Reliable predictions seem reasonnable. [that's the aim of this talk].

Confront data and theory.

- How to make reliable predictions? toy case
  - \* Process: Drell-Yan lepton pair production at the Tevatron.
  - \* Considered observables:
    - $\diamond$  the lepton-pair invariant-mass distribution  $\frac{d\sigma}{dM}$ .
    - $\diamond$  the lepton-pair transverse-momentum distribution  $\frac{d\sigma}{d\sigma_{\tau}}$ .
  - \* No new physics [for the moment, the talk's topic is not changed... ].

### Showers/resummatio

## Fixed-order perturbative theory

• QCD factorization theorem.

$$\sigma = \sum_{ab} \int \mathrm{d}x_a \, \mathrm{d}x_b \, f_{a/\rho_1}(x_a;\mu_F) \, f_{b/\rho_2}(x_b;\mu_F) \, \hat{\sigma}_{ab}$$

- \* Long-distance and short-distance physics factorize.
- \* Long-distance physics: parton densities  $f_a$ ,  $f_b$ .
- \* Short-distance physics: hard scattering matrix-element  $\hat{\sigma}_{ab}$ .
- \* Introduction of the unphysical factorization scale  $\mu_F$ .

### Showers/resummatio

## Fixed-order perturbative theory

• QCD factorization theorem.

$$\sigma = \sum_{ab} \int \mathrm{d}x_a \, \mathrm{d}x_b \, f_{a/p_1}(x_a;\mu_F) \, f_{b/p_2}(x_b;\mu_F) \, \hat{\sigma}_{ab}$$

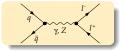
- \* Long-distance and short-distance physics factorize.
- \* Long-distance physics: parton densities  $f_a$ ,  $f_b$ .
- \* Short-distance physics: hard scattering matrix-element  $\hat{\sigma}_{ab}$ .
- \* Introduction of the unphysical factorization scale  $\mu_F$ .
- Partonic cross section: QCD parturbation theory.

$$\hat{\sigma} = \hat{\sigma}^{(0)} + \alpha_s \hat{\sigma}^{(1)} + \dots$$

Introduction 00●0000000		

# First guess: leading order predictions

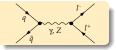
• Easy naive approach: matrix element calculation at leading order:



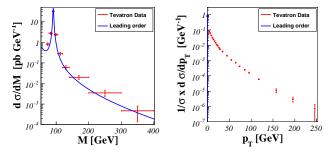
Introduction 00●00000000		

### First guess: leading order predictions

• Easy naive approach: matrix element calculation at leading order:



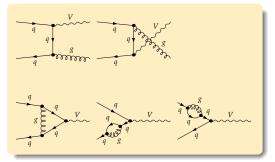
• Confrontation between theory and Tevatron data. [D\$\$\varphi\$ collaboration (2005, 2008)]



Disagreement between theory and experiment.

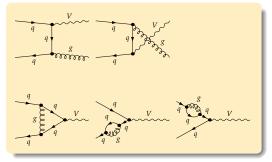
# Second try: next-to-leading order predictions (1)

• Improvement of the predictions: next-to-leading order calculation.



# Second try: next-to-leading order predictions (1)

• Improvement of the predictions: next-to-leading order calculation.



• Partonic invariant-mass and transverse-momentum distributions at  $\mathcal{O}(\alpha_s)$ ,

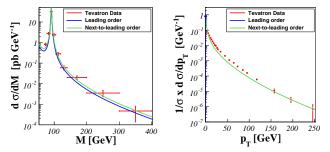
$$\begin{aligned} \frac{\mathrm{d}\hat{\sigma}}{\mathrm{d}M} &= \hat{\sigma}^{(0)}(M)\,\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z) + \mathcal{O}(\alpha_s^2),\\ \frac{\mathrm{d}^2\hat{\sigma}}{\mathrm{d}M\,\mathrm{d}p_T} &= \hat{\sigma}^{(0)}(M)\,\delta(p_T)\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z,p_T) + \mathcal{O}(\alpha_s^2), \end{aligned}$$
where  $z = M^2/s$ .

Precision predictions for BSM processes



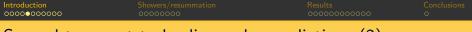
## Second try: next-to-leading order predictions (2)

● Confrontation between theory and Tevatron data. [DØ collaboration (2005, 2008)]



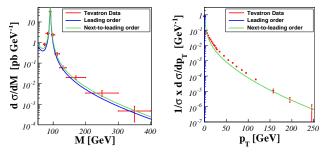
\* Invariant-mass distribution: good agreement.

- p\_-distribution:
  - ♦ good agreement in the large- $p_T$  region.
  - $\diamond$  diverges in the small- $p_T$  region.



# Second try: next-to-leading order predictions (2)

● Confrontation between theory and Tevatron data. [DØ collaboration (2005, 2008)]



\* Invariant-mass distribution: good agreement.

- $p_T$ -distribution:
  - ♦ good agreement in the large- $p_T$  region.
  - ♦ diverges in the small- $p_T$  region.

• How to improve NLO predictions? [in particular for the small-p<sub>T</sub> region.]

• Partonic invariant-mass and transverse-momentum distributions at  $\mathcal{O}(\alpha_s)$ ,

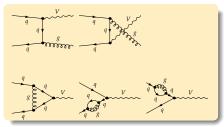
$$\begin{split} \frac{\mathrm{d}\hat{\sigma}}{\mathrm{d}M} &= \hat{\sigma}^{(0)}(M)\,\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z) + \mathcal{O}(\alpha_s^2),\\ \frac{\mathrm{d}^2\hat{\sigma}}{\mathrm{d}M\,\mathrm{d}p_T} &= \hat{\sigma}^{(0)}(M)\,\delta(p_T)\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z,p_T) + \mathcal{O}(\alpha_s^2),\\ \end{split}$$
 where  $z = M^2/s.$ 

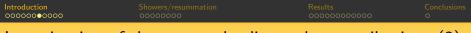
• Partonic invariant-mass and transverse-momentum distributions at  $\mathcal{O}(\alpha_s)$ ,

$$\frac{\mathrm{d}\hat{\sigma}}{\mathrm{d}M} = \hat{\sigma}^{(0)}(M)\,\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z) + \mathcal{O}(\alpha_s^2),$$
$$\frac{\mathrm{d}^2\hat{\sigma}}{\mathrm{d}M\,\mathrm{d}p_T} = \hat{\sigma}^{(0)}(M)\,\delta(p_T)\delta(1-z) + \alpha_s\,\hat{\sigma}^{(1)}(M,z,p_T) + \mathcal{O}(\alpha_s^2),$$

where  $z = M^2/s$ .

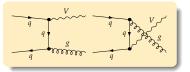
- $\hat{\sigma}^{(1)}$  contains two different pieces.
  - \* Real gluon emission diagrams.
  - \* Virtual loop contributions.



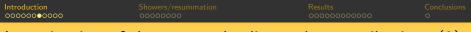


## Investigation of the next-to-leading order contributions (2)

• Amplitude for soft real gluon emission.

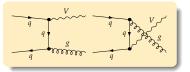


$$\begin{split} iM &= g_s T^a \ \bar{v}(k_2) \left[ \frac{\not(k_g) \ \left( k_g + k_2 \right) \ \Gamma^{\mu}_{qqV}}{2k_2 \cdot k_g} - \frac{\Gamma^{\mu}_{qqV} \ \left( k_g + k_1 \right) \ \not(k_g)}{2k_1 \cdot k_g} \right] u(k_1) \\ &\approx g_s T^a \left[ \frac{\epsilon \cdot k_2}{k_2 \cdot k_g} - \frac{k_1 \cdot \epsilon}{k_1 \cdot k_g} \right] \bar{v}(k_2) \ \Gamma^{\mu}_{qqV} u(k_1) \\ &= g_s T^a \left[ \frac{\epsilon \cdot k_2}{k_0^2 \mathbf{k}_0^2 (1 + \cos \theta)} - \frac{k_1 \cdot \epsilon}{k_0^2 \mathbf{k}_0^2 (1 - \cos \theta)} \right] \mathbf{iM}^{\mathrm{Born}} \end{split}$$



## Investigation of the next-to-leading order contributions (2)

• Amplitude for soft real gluon emission.



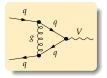
$$\begin{split} iM &= g_s T^a \ \bar{v}(k_2) \left[ \frac{\not(k_g) \ \left( \not k_g + \not k_2 \right) \ \Gamma^{\mu}_{qqV}}{2k_2 \cdot k_g} - \frac{\Gamma^{\mu}_{qqV} \ \left( \not k_g + \not k_1 \right) \ \not(k_g)}{2k_1 \cdot k_g} \right] u(k_1) \\ &\approx g_s T^a \left[ \frac{\epsilon \cdot k_2}{k_2 \cdot k_g} - \frac{k_1 \cdot \epsilon}{k_1 \cdot k_g} \right] \bar{v}(k_2) \ \Gamma^{\mu}_{qqV} u(k_1) \\ &= g_s T^a \left[ \frac{\epsilon \cdot k_2}{k_0^2 \mathbf{k}_g^0 (1 + \cos \theta)} - \frac{k_1 \cdot \epsilon}{k_0^2 \mathbf{k}_g^0 (1 - \cos \theta)} \right] \mathbf{iM}^{\mathrm{Born}} \end{split}$$

Soft and collinear radiation diverges and factorizes.



## Investigation of the next-to-leading order contributions (3)

• Amplitude for the virtual contribution (soft gluons in the loop).

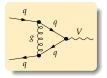


$$\begin{split} iM &= (i\,g_s^2)\bar{\nu}(k_2) \int \mathrm{d}k_g \frac{\gamma_{\nu}\left(k_2 + k_g\right)\Gamma^{\mu}_{qqV}\left(k_1 - k_g\right)\gamma^{\nu}}{k_g^2\left(2k_1 \cdot k_g\right)\left(2k_2 \cdot k_g\right)} u(k_1) \\ &\approx (i\,g_s^2) \int \mathrm{d}k_g \frac{k_1 \cdot k_2}{k_g^2\left(k_1 \cdot k_g\right)\left(k_2 \cdot k_g\right)} iM^{\mathrm{Born}} \\ &= (i\,g_s^2) \int \mathrm{d}k_g \frac{k_1 \cdot k_2}{k_g^2\left(k_1^0 \mathbf{k}_g^0(1 - \cos\theta)\right)\left(k_2^0 \mathbf{k}_g^0(1 + \cos\theta)\right)} \mathbf{i} \mathbf{M}^{\mathrm{Born}} \end{split}$$



## Investigation of the next-to-leading order contributions (3)

• Amplitude for the virtual contribution (soft gluons in the loop).



$$\begin{split} iM &= (i\,g_s^2)\bar{\nu}(k_2) \int \mathrm{d}k_g \frac{\gamma_\nu \left(k_2 + k_g\right)\Gamma^{\mu}_{qqV}\left(k_1 - k_g\right)\gamma^\nu}{k_g^2\left(2k_1 \cdot k_g\right)\left(2k_2 \cdot k_g\right)} u(k_1) \\ &\approx (i\,g_s^2) \int \mathrm{d}k_g \frac{k_1 \cdot k_2}{k_g^2\left(k_1 \cdot k_g\right)\left(k_2 \cdot k_g\right)} iM^{\mathrm{Born}} \\ &= (i\,g_s^2) \int \mathrm{d}k_g \frac{k_1 \cdot k_2}{k_g^2\left(k_1^0 \mathbf{k}_g^0(1 - \cos\theta)\right)\left(k_2^0 \mathbf{k}_g^0(1 + \cos\theta)\right)} \mathbf{i} \mathbf{M}^{\mathrm{Born}} \end{split}$$

The virtual contributions diverge and factorize.

• Sum of the two contributions.

$$\hat{\sigma}^{(1)} = \hat{\sigma}^{(1,\text{loop})} + \hat{\sigma}^{(1,\text{real})}$$

- \* Cancellation of the poles.
- \* Infrared behaviour: logarithmic terms in the distributions,

$$\alpha_s \left( \frac{\ln(1-z)}{1-z} \right)_+$$
 and  $\frac{\alpha_s}{p_T^2} \ln \frac{M^2}{p_T^2}$ 

\* Problems at  $z \leq 1$  or small  $p_T$ .

• Sum of the two contributions.

$$\hat{\sigma}^{(1)} = \hat{\sigma}^{(1,\text{loop})} + \hat{\sigma}^{(1,\text{real})}$$

- \* Cancellation of the poles.
- \* Infrared behaviour: logarithmic terms in the distributions,

$$\alpha_s \left( \frac{\ln(1-z)}{1-z} \right)_+$$
 and  $\frac{\alpha_s}{p_T^2} \ln \frac{M^2}{p_T^2}$ 

\* Problems at  $z \leq 1$  or small  $p_T$ .

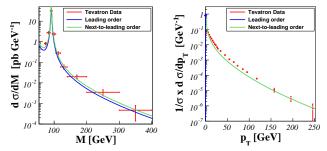
The fixed-order theory is unreliable in these kinematical regions.

 Introduction
 Showers/resummation
 Results
 Conclusions

 0000000000
 0000000000
 0
 0
 0

## The problem of the soft and collinear radiation (2)

• Confrontation between theory and Tevatron data. [D\$\varphi\$ collaboration (2005, 2008)]

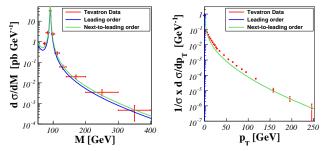


- \* Invariant-mass distribution:
  - $\diamond$  Convolution with the steeply falling parton densities at large z.
  - ♦ Next-to-leading order calculation reliable.

Introduction Showers/resummation Results Conclusions

# The problem of the soft and collinear radiation (2)

• Confrontation between theory and Tevatron data. [DØ collaboration (2005, 2008)]



- \* Invariant-mass distribution:
  - $\diamond$  Convolution with the steeply falling parton densities at large z.
  - ♦ Next-to-leading order calculation reliable.
- \* *p*<sub>T</sub>-distribution:
  - $\diamond~$  Next-to-leading order calculation reliable for the large  $p_T.$
  - ♦ Behaviour in the small- $p_T$  region:  $\propto \frac{\alpha_s}{p_T^2} \ln \frac{M^2}{p_T^2}$ .

### Improvements

Improvements of the next-to-leading order calculation.

- Matching with a resummation calculation.
  - \* Correct treatment of the soft and collinear radiation.
  - \* Perturbative method.
  - \* Soft and collinear radiation taken into account to all orders.
  - \* Parton-level calculation.
- Matching with a parton shower algorithm.
  - \* Approximation of the resummation calculation.
  - \* Suitable for a proper description of the collision.

# Outline

#### Models and motivation

Motivation for resummation calculations

### 2

### Parton showers and resummation

- Parton showers
- Transverse-momentum, threshold and joint resummation formalisms
- Matching to the fixed order

#### Numerical results, with uncertainties

- The Drell-Yan and the Tevatron
- Grand Unified Theories and Z' bosons
- The Minimal Supersymmetric Standard Model (MSSM)

#### Summary - conclusions

	Showers/resummation ●0000000	
Parton showers	(1)	

• The parton splitting factorizes  $\Rightarrow$  iterative splitting.

where t is the ordering variable and z the momentum fraction.

### Showers/resummation

Conclusions o

# Parton showers (1)

• The parton splitting factorizes  $\Rightarrow$  iterative splitting.

where t is the ordering variable and z the momentum fraction.

- No emission probability the Sudakov form factor.
  - \* Conservation of probability for the branching of a parton:

$$\begin{split} 1 &= P_{\rm no\ emis}(t+{\rm d}t,t) + P_{\rm emis}(t+{\rm d}t,t) \\ &= P_{\rm no\ emis}(t+{\rm d}t,t) + \frac{{\rm d}t}{t}\sum_b \int {\rm d}z \frac{\alpha_s(t)}{4\pi} P_{ab}(z) \end{split}$$

### Showers/resummation

Conclusions o

# Parton showers (1)

• The parton splitting factorizes  $\Rightarrow$  iterative splitting.

where t is the ordering variable and z the momentum fraction.

- No emission probability the Sudakov form factor.
  - \* Conservation of probability for the branching of a parton:

$$\begin{split} \mathbf{1} &= P_{\mathrm{no} \ \mathrm{emis}}(t + \mathrm{d}t, t) + P_{\mathrm{emis}}(t + \mathrm{d}t, t) \\ &= P_{\mathrm{no} \ \mathrm{emis}}(t + \mathrm{d}t, t) + \frac{\mathrm{d}t}{t} \sum_{b} \int \mathrm{d}z \frac{\alpha_{\mathsf{s}}(t)}{4\pi} P_{\mathsf{ab}}(z) \end{split}$$

\* Solving the equation defines the Sudakov form factor,

$$\mathbf{\Delta}(\mathbf{t}_1, \mathbf{t}_2) = P_{\text{no emis}}(t_1, t_2) = \exp\left[-\int_{t_1}^{t_2} \frac{\mathrm{d}t}{t} \sum_b \int \mathrm{d}z \frac{\alpha_s}{4\pi} P_{ba}(z)\right]$$

### Showers/resummation

Conclusions

# Parton showers (1)

• The parton splitting factorizes  $\Rightarrow$  iterative splitting.

where t is the ordering variable and z the momentum fraction.

- No emission probability the Sudakov form factor.
  - \* Conservation of probability for the branching of a parton:

$$\begin{split} 1 &= P_{\rm no\ emis}(t+{\rm d}t,t) + P_{\rm emis}(t+{\rm d}t,t) \\ &= P_{\rm no\ emis}(t+{\rm d}t,t) + \frac{{\rm d}t}{t}\sum_b \int {\rm d}z \frac{\alpha_s(t)}{4\pi} P_{ab}(z) \end{split}$$

\* Solving the equation defines the Sudakov form factor,

$$\Delta(\mathbf{t}_1, \mathbf{t}_2) = P_{\text{no emis}}(t_1, t_2) = \exp\left[-\int_{t_1}^{t_2} \frac{\mathrm{d}t}{t} \sum_b \int \mathrm{d}z \frac{\alpha_s}{4\pi} P_{ba}(z)\right]$$

\* Logarithmic dependence  $\Leftrightarrow$  leading-log approximation.

Showers/resummation ⊙●○○○○○○	
$\langle 0 \rangle$	

Evolution equation for the parton a to the cut-off scale  $t_0$ 

$$\phi_{a}(t,E) = \Delta_{a}(t,t_{0}) + \sum_{b} \int_{t_{0}}^{t} \frac{\alpha_{s}(t')}{4\pi} \frac{\mathrm{d}t'}{t'} \mathrm{d}z \,\Delta_{a}(t,t') \,P_{ab}(z) \,\phi_{b}(t',zE) \,\phi_{c}(t',(1-z)E)$$

Showers/resummation ⊙●○○○○○○	
$\langle 0 \rangle$	

Evolution equation for the parton a to the cut-off scale  $t_0$ 

$$\phi_{a}(t,E) = \Delta_{a}(t,t_{0}) + \sum_{b} \int_{t_{0}}^{t} \frac{\alpha_{s}(t')}{4\pi} \frac{\mathrm{d}t'}{t'} \mathrm{d}z \,\Delta_{a}(t,t') \,P_{ab}(z) \,\phi_{b}(t',zE) \,\phi_{c}(t',(1-z)E)$$

#### • Derivation of a parton shower algorithm

\* Ordered Markov chain (t-variable)

$$Q_0^2 \ll t_1 \ll t_2 \ll \ldots \ll t_N \ll Q^2$$

\* Choice of *t*: different shower algorithms.

Showers/resummation ⊙●○○○○○○	

Evolution equation for the parton a to the cut-off scale  $t_0$ 

$$\phi_{a}(t,E) = \Delta_{a}(t,t_{0}) + \sum_{b} \int_{t_{0}}^{t} \frac{\alpha_{s}(t')}{4\pi} \frac{\mathrm{d}t'}{t'} \mathrm{d}z \,\Delta_{a}(t,t') \,P_{ab}(z) \,\phi_{b}(t',zE) \,\phi_{c}(t',(1-z)E)$$

#### • Derivation of a parton shower algorithm

\* Ordered Markov chain (t-variable)

$$Q_0^2 \ll t_1 \ll t_2 \ll \ldots \ll t_N \ll Q^2$$

- \* Choice of t: different shower algorithms.
- Limitations.
  - \* Leading logarithms,
  - \* Large number of colors,
  - \* Collinear and/or soft-collinear radiation.

Showers/resummation ⊙●○○○○○○	

Evolution equation for the parton a to the cut-off scale  $t_0$ 

$$\phi_{a}(t,E) = \Delta_{a}(t,t_{0}) + \sum_{b} \int_{t_{0}}^{t} \frac{\alpha_{s}(t')}{4\pi} \frac{\mathrm{d}t'}{t'} \mathrm{d}z \,\Delta_{a}(t,t') \,P_{ab}(z) \,\phi_{b}(t',zE) \,\phi_{c}(t',(1-z)E)$$

#### • Derivation of a parton shower algorithm

\* Ordered Markov chain (t-variable)

$$Q_0^2 \ll t_1 \ll t_2 \ll \ldots \ll t_N \ll Q^2$$

- \* Choice of t: different shower algorithms.
- Limitations.
  - \* Leading logarithms,
  - \* Large number of colors,
  - \* Collinear and/or soft-collinear radiation.
- Improvements require matrix exponentiation  $\Rightarrow$  soft-gluon resummation.

# Main features of the resummation

#### • Reorganization of the cross section:

$$\mathrm{d}\sigma = \mathrm{d}\sigma^{(\mathrm{res})} + \mathrm{d}\sigma^{(\mathrm{fin})} \; .$$

### • $d\sigma^{(res)}$ :

- \* Contains all the logarithmic terms.
- \* Resummed to all orders in  $\alpha_s$ .
- \* Exponentiation (Sudakov form factor).
- $d\sigma^{(fin)}$ :
  - \* Remaining (regular) contributions.

# The resummed component (1)

#### • Based on factorization properties.

- \* Holds in non-physical conjugate spaces.
- \* Mellin N-space (N conjugate to  $M^2/S_h$ ).
- \* Impact parameter b (conjugate to  $p_T$ ).

$$\begin{split} \mathrm{d}\sigma^{(\mathrm{res})}(N,b) &= \sum_{a,b} f_{a/h_1}(N+1) f_{b/h_2}(N+1) \mathcal{W}_{ab}(N,b), \\ \mathcal{W}_{ab}(N,b) &= \mathcal{H}_{ab}(N) \exp\Big\{\mathcal{G}(N,b)\Big\}. \end{split}$$

# The resummed component (1)

#### • Based on factorization properties.

- \* Holds in non-physical conjugate spaces.
- \* Mellin N-space (N conjugate to  $M^2/S_h$ ).
- \* Impact parameter b (conjugate to  $p_T$ ).

$$\begin{split} \mathrm{d}\sigma^{(\mathrm{res})}(N,b) &= \sum_{a,b} f_{a/h_1}(N+1) f_{b/h_2}(N+1) \mathcal{W}_{ab}(N,b), \\ \mathcal{W}_{ab}(N,b) &= \mathcal{H}_{ab}(N) \exp\left\{\mathcal{G}(N,b)\right\}. \end{split}$$

- The *H*-coefficient:
  - \* Contains real and virtual collinear radiation, hard contributions.
- The Sudakov form factor  $\mathcal{G}$ :
  - \* Contains the soft-collinear radiation.

Showers/resummation

# The resummed component (2)

$$\mathcal{W}_{ab}(N,b) = \mathcal{H}_{ab}(N) \exp \left\{ \mathcal{G}(N,b) \right\}.$$

- The *H*-coefficient:
  - \* Contains real and virtual collinear radiation, hard contributions.
  - \* Can be computed perturbatively as series in  $\alpha_s$ , from fixed-order results.
  - \* Is process-dependent.

Showers/resummation

# The resummed component (2)

$$\mathcal{W}_{ab}(N,b) = \mathcal{H}_{ab}(N) \exp \left\{ \mathcal{G}(N,b) \right\}.$$

- The *H*-coefficient:
  - \* Contains real and virtual collinear radiation, hard contributions.
  - \* Can be computed perturbatively as series in  $\alpha_s$ , from fixed-order results.
  - \* Is process-dependent.
- The Sudakov form factor  $\mathcal{G}$ :
  - \* Contains the soft-collinear radiation.
  - \* Can be computed perturbatively as series in  $\alpha_s \log$ .
  - \* Is process-independent (universal).
  - \* Contains the full color and spin structure.

### References

- p<sub>T</sub>-resummation [Catani, de Florian, Grazzini (2001); Bozzi, Catani, de Florian, Grazzini (2006)]
  - \* Universal formalism  $\equiv$  process-independent Sudakov form factor.
  - \* Resums  $\frac{\alpha_s}{p_T^2} \ln \frac{M^2}{p_T^2}$ .
- Threshold resummation [Sterman (1987); Catani, Trentadue (1989,1991)]

\* Resums 
$$\left(\frac{\ln(1-z)}{1-z}\right)_+$$

- Joint resummation [Bozzi, BenjF, Klasen (2008)]
  - \* Universal formalism  $\equiv$  process-independent Sudakov form factor.
  - \* Resums both types of logarithms.

- Fixed-order calculations.
  - \* Reliable far from the critical kinematical regions.
  - \* Spoiled in the critical regions.

- Fixed-order calculations.
  - \* Reliable far from the critical kinematical regions.
  - \* Spoiled in the critical regions.
- Resummation (parton showers).
  - \* Needed in the critical regions.
  - \* Not justified far from the critical regions.

- Fixed-order calculations.
  - \* Reliable far from the critical kinematical regions.
  - \* Spoiled in the critical regions.
- Resummation (parton showers).
  - \* Needed in the critical regions.
  - \* Not justified far from the critical regions.
- Intermediate kinematical regions:
  - \* Both fixed order and resummation / parton showers contribute.

- Fixed-order calculations.
  - \* Reliable far from the critical kinematical regions.
  - \* Spoiled in the critical regions.
- Resummation (parton showers).
  - \* Needed in the critical regions.
  - \* Not justified far from the critical regions.
- Intermediate kinematical regions:
  - \* Both fixed order and resummation / parton showers contribute.

Information from both fixed order and resummation (parton showers) is required.  $\Rightarrow$  consistent matching procedure.

#### • Matching procedure:

- \* Addition of both resummation and fixed-order results.
- \* Subtracting the expansion in  $\alpha_s$  of the resummed result.
- \* No double-counting of the logarithms.

$$\mathrm{d}\sigma = \mathrm{d}\sigma^{(\mathrm{F.O.})} + \mathrm{d}\sigma^{(\mathrm{res})} - \mathrm{d}\sigma^{(\mathrm{exp})}.$$

#### • Matching procedure:

- \* Addition of both resummation and fixed-order results.
- \* Subtracting the expansion in  $\alpha_s$  of the resummed result.
- \* No double-counting of the logarithms.

$$d\sigma = d\sigma^{(F.O.)} + d\sigma^{(res)} - d\sigma^{(exp)}.$$

#### • Effects of the matching procedure:

- \* Far from the critical regions,  $d\sigma^{(res)} \approx d\sigma^{(exp)} \equiv$  perturbative theory.
- \* In the critical regions,  $d\sigma^{(F.O.)} \approx d\sigma^{(exp)} \equiv$  pure resummation.
- \* In the intermediate regions: **both contribute**.

### Outline

### Models and motivation

Motivation for resummation calculations

### 2

#### Parton showers and resummation

- Parton showers
- Transverse-momentum, threshold and joint resummation formalisms
- Matching to the fixed order

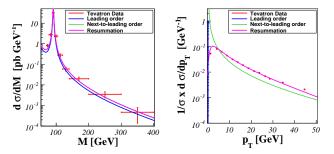
#### Numerical results, with uncertainties

- The Drell-Yan and the Tevatron
- Grand Unified Theories and Z' bosons
- The Minimal Supersymmetric Standard Model (MSSM)

#### 4 Summary - conclusions

### Resummation vs. Tevatron data

● Confrontation between theory and Tevatron data. [DØ collaboration (2005, 2008)]



- Invariant-mass distribution: good agreement. (no change with respect to next-to-leading order).
- *p<sub>T</sub>*-distribution: good agreement. (improvement with respect to next-to-leading order).

### Showers/resummation

# Grand Unified Theories and Z' bosons

- Generalities of the Grand Unified Theories.
  - \* Unification of the Standard Model gauge groups:

 $G \supset SU(3)_{\mathsf{C}} \times SU(2)_{\mathsf{L}} \times U(1)_{\mathsf{Y}}.$ 

# Grand Unified Theories and Z' bosons

- Generalities of the Grand Unified Theories.
  - \* Unification of the Standard Model gauge groups:

 $\label{eq:G_states} G \supset SU(3)_{C} \times SU(2)_{L} \times U(1)_{Y}.$ 

- \* Breaking to the SM at high energy scale:
  - $\diamond$  Appearance of additional U(1) symmetries.
  - ♦ Extra neutral gauge bosons Z'.

# Grand Unified Theories and Z' bosons

- Generalities of the Grand Unified Theories.
  - \* Unification of the Standard Model gauge groups:

 $\label{eq:G_states} G \supset SU(3)_{C} \times SU(2)_{L} \times U(1)_{Y}.$ 

- \* Breaking to the SM at high energy scale:
  - $\diamond$  Appearance of additional U(1) symmetries.
  - ♦ Extra neutral gauge bosons Z'.
- Considered theoretical model. [Green, Schwarz (1984); Hewett, Rizzo (1989)]
  - \* Ten-dimensional string theories  $E_8 \times E_8$ :
    - ♦ Anomaly-free and contains chiral fermions.
    - $\diamond$  Compactified to  $E_6$ .

# Grand Unified Theories and Z' bosons

- Generalities of the Grand Unified Theories.
  - \* Unification of the Standard Model gauge groups:

 $\label{eq:G_states} G \supset SU(3)_{C} \times SU(2)_{L} \times U(1)_{Y}.$ 

- \* Breaking to the SM at high energy scale:
  - $\diamond$  Appearance of additional U(1) symmetries.
  - ♦ Extra neutral gauge bosons Z'.

• Considered theoretical model. [Green, Schwarz (1984); Hewett, Rizzo (1989)]

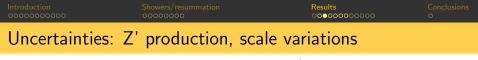
\* Ten-dimensional string theories  $E_8 \times E_8$ :

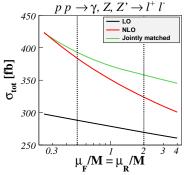
- ♦ Anomaly-free and contains chiral fermions.
- $\diamond$  Compactified to  $E_6$ .
- \* Breaking to the SM gauge groups

$$E_6 
ightarrow SO(10) imes {f U(1)}_\psi$$

- ightarrow *SU*(5) imes U(1) $_{\chi}$  imes U(1) $_{\psi}$
- ightarrow SU(3)<sub>C</sub> imes SU(2)<sub>L</sub> imes U(1)<sub>Y</sub> imes U(1)<sub> $\chi$ </sub> imes U(1)<sub> $\psi$ </sub>.

Additional bosons  $Z_{\psi}$  and  $Z' \equiv Z_{\chi}$ .



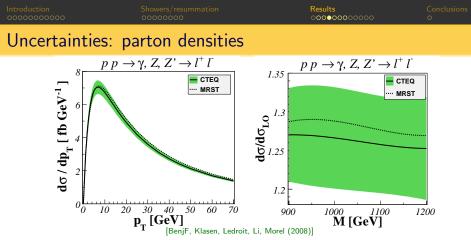


[BenjF, Klasen, Ledroit, Li, Morel (2008)]

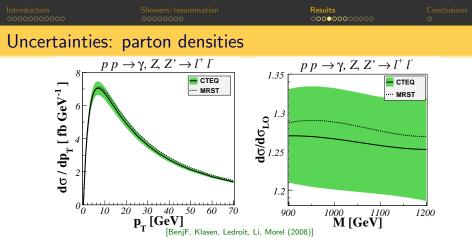
• Scenario.

' 1 TeV Z'.

- \* LHC collider @ 14 TeV.
- Total cross section (900 GeV  $\leq M \leq 1200$  GeV).
  - \* Leading order: full dependence related to  $\mu_F$  (~ 7%).
  - \* Next-to-leading order: introduction of  $\mu_R$  and the qg channel ( $\sim 17\%$ ).
  - \* Resummation: reduction of scale dependence ( $\sim$  9%).



- Scenario: 1 TeV Z'; LHC collider.
- CTEQ vs. MRST.
  - \* p<sub>T</sub>-spectrum: similar shapes but a bit harder for MRST.
  - \* Mass-spectrum: different shapes.



- Scenario: 1 TeV Z'; LHC collider.
- CTEQ vs. MRST.
  - \* p<sub>T</sub>-spectrum: similar shapes but a bit harder for MRST.
  - \* Mass-spectrum: different shapes.
- Variations along 20 directions for the CTEQ densities.
  - \* Variations along the PDF fits: modest uncertainties ( $\sim 10\%$ ).
  - \* Similar to scale dependence.

	Results ○000●00○○○○○	

### Non-perturbative effects

- Important non-perturbative effects in the  $p_T$ -distributions.
  - \* Intrinsic  $p_T$  of the partons inside the hadrons.
  - \* Modification of the Sudakov form factor,

$$\mathcal{G}(N,b) \rightarrow \mathcal{G}(N,b) + \mathcal{F}_{ab}^{\mathrm{NP}}.$$

	Results ○000●000000	

### Non-perturbative effects

- Important non-perturbative effects in the  $p_T$ -distributions.
  - \* Intrinsic  $p_T$  of the partons inside the hadrons.
  - \* Modification of the Sudakov form factor,

$$\mathcal{G}(N,b) \rightarrow \mathcal{G}(N,b) + \mathbf{F}_{ab}^{\mathrm{NP}}.$$

- Form factors [Ladinsky, Yuan (94); Landry, Brock, Nadolsky, Yuan (03); Konyshev, Nadolsky (06)].
  - \* Obtained from experimental data (fits) and assumed universal.

	<b>Results</b> ○ <b>○○○</b> ○○○○○○○	

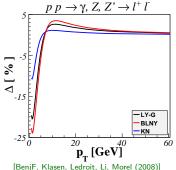
### Non-perturbative effects

- Important non-perturbative effects in the  $p_T$ -distributions.
  - \* Intrinsic p<sub>T</sub> of the partons inside the hadrons.
  - \* Modification of the Sudakov form factor,

$$\mathcal{G}(N,b) \rightarrow \mathcal{G}(N,b) + \mathbf{F}_{ab}^{\mathrm{NP}}.$$

• Form factors [Ladinsky, Yuan (94); Landry, Brock, Nadolsky, Yuan (03); Konyshev, Nadolsky (06)].

\* Obtained from experimental data (fits) and assumed universal.



• Non-perturbative effects under good control for p<sub>T</sub> > 5 GeV.

# Monte Carlo and resummation for BSM processes

### • Soft and collinear radiation $\equiv$ Sudakov form factor.

- \* Parton showers in general: leading logarithms, color,...
- \* Momentum conservation at each branching: (leading logs)+, e.g. PYTHIA.
- \* Resummation: next-to-leading logarithms.

# Monte Carlo and resummation for BSM processes

### • Soft and collinear radiation $\equiv$ Sudakov form factor.

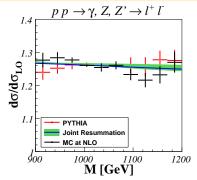
- \* Parton showers in general: leading logarithms, color,...
- \* Momentum conservation at each branching: (leading logs)+, e.g. PYTHIA.
- \* Resummation: next-to-leading logarithms.
- Matched with matrix elements.
  - \* Monte Carlo codes in general: leading order.
  - \* Sometimes next-to-leading order: e.g. MC@NLO and POWHEG.
  - \* Resummation: next-to-leading order.

# Monte Carlo and resummation for BSM processes

### • Soft and collinear radiation $\equiv$ Sudakov form factor.

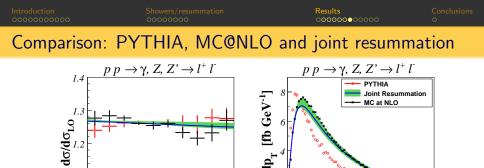
- \* Parton showers in general: leading logarithms, color,...
- \* Momentum conservation at each branching: (leading logs)+, e.g. PYTHIA.
- \* Resummation: next-to-leading logarithms.
- Matched with matrix elements.
  - \* Monte Carlo codes in general: leading order.
  - \* Sometimes next-to-leading order: e.g. MC@NLO and POWHEG.
  - \* Resummation: next-to-leading order.
- Comparison: resummation vs. PYTHIA vs. MC@NLO.
  - \* **PYTHIA**: virtuality-ordered showers; nice process library.
  - \* MC@NLO: angular-ordered showers; precision MC generator.
  - \* Resummation: best precision.

# Comparison: PYTHIA, MC@NLO and joint resummation



[BenjF, Klasen, Ledroit, Li, Morel (2008)]

- 1 TeV Z'; PYTHIA (LO/LL+), MC@NLO (NLO/LL), resummation (NLO/NLL).
- Mass-spectrum normalized to leading order:
  - \* PYTHIA (power shower): mass-spectrum multiplied by a K-factor of 1.26.
  - \* Good agreement between MC@NLO and resummation.



 $d \, \sigma / dp_{_{\rm T}}$ 

\* PYTHIA (*power shower*): mass-spectrum multiplied by a K-factor of 1.26.

10 20

 $p_{T}^{30}$  [GeV]

\* PYTHIA spectrum much too soft, peak not well predicted.

1100

M [GeV]

• Mass-spectrum normalized to leading order:

\* Good agreement between MC@NLO and resummation.

\* Good agreement between MC@NLO and resummation.

1200

[BenjF, Klasen, Ledroit, Li, Morel (2008)] ● 1 TeV Z'; PYTHIA (LO/LL+), MC@NLO (NLO/LL), resummation (NLO/NLL).

1.1

*φυυ* 

PYTHIA Joint Resummation MC at NLO

1000

Transverse-momentum distribution:

60 70

# The Minimal Supersymmetric Standard Model (MSSM)

- High energy extension to Standard Model.
- Symmetry between fermions and bosons.

$$\begin{split} & Q|\text{Boson}\rangle = |\text{Fermion}\rangle \\ & Q|\text{Fermion}\rangle = |\text{Boson}\rangle \quad \text{where } Q \text{ is a SUSY generator.} \end{split}$$

# The Minimal Supersymmetric Standard Model (MSSM)

- High energy extension to Standard Model.
- Symmetry between fermions and bosons.

 $Q|Boson\rangle = |Fermion\rangle$  $Q|Fermion\rangle = |Boson\rangle$  where Q is a SUSY generator.

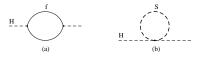
• The MSSM: one single supersymmetric (SUSY) generator Q.

#### The MSSM: one SUSY partner for each SM particle.

- \* Quarks  $\Leftrightarrow$  squarks.
- \* Leptons  $\Leftrightarrow$  sleptons.
- \* Gauge/Higgs bosons ⇔ gauginos/higgsinos ⇔ charginos/neutralinos.
- \* Gluon ⇔ gluino.

### Some features of the MSSM

- Introduction of the SUSY particles in the theory.
  - \* Solution to the hierarchy problem (stabilization of the Higgs mass).



◊ Fermionic loop (Fig. a):

$$\Delta M_{H}^{(a)} = -\frac{y_{f}^{2}}{16 \pi^{2}} \left[ 2 \Lambda^{2} + 6 m_{f}^{2} \ln \frac{\Lambda}{m_{f}} + \dots \right]$$

◊ Scalar loop (Fig. b):

$$\Delta M_H^{(b)} = \frac{\lambda_s}{16 \, \pi^2} \left[ \Lambda^2 - 2 \, m_s^2 \ln \frac{\Lambda}{m_s} + \dots \right]$$

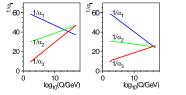
- ♦ MSSM: two scalars per fermion and  $\lambda_S = y_f^2$ .
- ♦ Sum: logarithmic divergences.

### Showers/resummatic

Results	
000000000000000000000000000000000000000	

### Some features of the MSSM

- Introduction of the SUSY particles in the theory.
  - \* Solution to the hierarchy problem (stabilization of the Higgs mass).
  - \* Gauge coupling unification at high energy.



[Fig. from de Boer, Sander (2004)].

- $\diamond~$  SUSY particles added in the RGE.
- $\diamond~$  Gauge couplings unify at  $Q\sim 10^{16}~{\rm GeV}.$

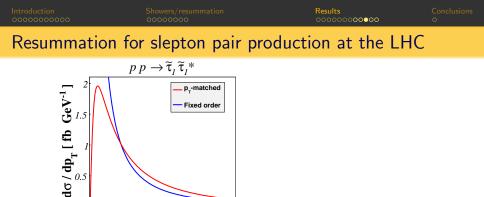
### Some features of the MSSM

- Introduction of the SUSY particles in the theory.
  - \* Solution to the hierarchy problem (stabilization of the Higgs mass).
  - \* Gauge coupling unification at high energy.
  - \* **Dark matter candidate**  $\Leftrightarrow$  lightest SUSY particle stable and neutral.

Results

### Some features of the MSSM

- Introduction of the SUSY particles in the theory.
  - \* Solution to the hierarchy problem (stabilization of the Higgs mass).
  - \* Gauge coupling unification at high energy.
  - \* **Dark matter candidate** ⇔ lightest SUSY particle stable and neutral.
- No SUSY discovery until now!
  - \* SUSY must be broken.
  - \* SUSY masses at a higher scale than Standard Model (SM) masses.
  - \* More than 100 new free parameters.
  - \* Simplified benchmark scenarios:
    - ♦ Minimal supergravity (mSUGRA).
    - ◊ Gauge-mediated SUSY-breaking (GMSB).
    - ۰...



[Bozzi, BenjF, Klasen (2006, 2007)]

100

• SUSY scenario: slepton masses  $\approx$  100-200 GeV.

60

• Resummation effects:

20

\* Finite results at small  $p_T$ .

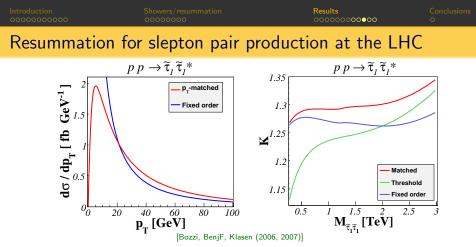
40

р<sub>т</sub> [GeV]

\* Matching: important effects at intermediate p<sub>T</sub>.

80

00

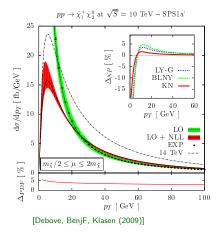


• SUSY scenario: slepton masses  $\approx$  100-200 GeV.

#### • Resummation effects:

- \* Finite results at small  $p_T$ .
- \* Matching: important effects at intermediate  $p_T$ .
- \* Small M:  $d\sigma^{(res)} \approx d\sigma^{(exp)} \equiv$  perturbative theory.
- \* Large *M*:  $d\sigma^{(F.O.)} \approx d\sigma^{(exp)} \equiv pure resummation.$

### Uncertainties: chargino-neutralino associated production



Scenario.

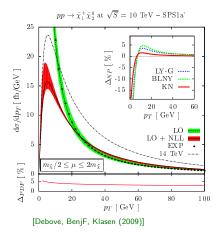
- \*  $\,pprox\,$  180 GeV gauginos.
- \* LHC collider (10 TeV & 14 TeV).

#### • *p*<sub>T</sub>-**spectrum**

- \* Next-to-leading logarithms.
- \*  $\mathcal{O}(\alpha_s)$  fixed-order.
- \* Small  $p_T$ : expansion  $\approx$  fixed-order.
- \* Large  $p_T$ : expansion  $\approx$  resummation.
- \* Intermediate  $p_T$ : enhancement.

#### Showers/resummatio 00000000

### Uncertainties: chargino-neutralino associated production



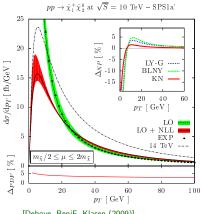
Scenario.

- \*  $\approx$  180 GeV gauginos.
- \* LHC collider (10 TeV & 14 TeV).

#### • *p*<sub>T</sub>-**spectrum**

- \* Next-to-leading logarithms.
- O(α<sub>s</sub>) fixed-order.
- \* Small  $p_T$ : expansion  $\approx$  fixed-order.
- \* Large  $p_T$ : expansion  $\approx$  resummation.
- \* Intermediate *p*<sub>T</sub>: enhancement.
- Scale dependence  $(M/2 \le \mu_R = \mu_F \le 2M)$ .
  - \* Reduction of the uncertainties.
  - \* Less than 5% for  $p_T > 5$  GeV.

# Uncertainties: chargino-neutralino associated production



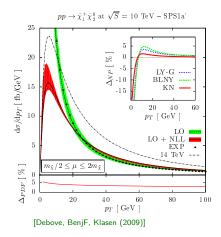
[Debove, BeniF, Klasen (2009)]

Scenario.

- \*  $\approx 180 \text{ GeV}$  gauginos.
- LHC collider (10 TeV & 14 TeV).

#### p<sub>T</sub>-spectrum

- \* Next-to-leading logarithms.
- \*  $\mathcal{O}(\alpha_s)$  fixed-order.
- \* Small  $p_T$ : expansion  $\approx$  fixed-order.
- **Large**  $p_T$ : expansion  $\approx$  resummation.
- \* Intermediate p<sub>T</sub>: enhancement.
- Scale dependence  $(M/2 \le \mu_R = \mu_F \le 2M)$ .
  - **Reduction** of the uncertainties
  - Less than 5% for  $p_T > 5$  GeV.
- Parton densities dependence (44 CTEQ sets).
  - 4-5% uncertainties for all  $p_T$ .
  - Similar to weak boson production.

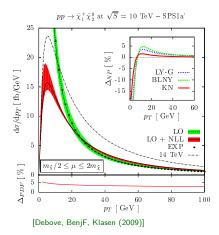


• Scenario.

- \*  $\,pprox\,$  180 GeV gauginos.
- \* LHC collider (10 TeV & 14 TeV).

### • *p*<sub>T</sub>-spectrum

- \* Next-to-leading logarithms.
- \*  $\mathcal{O}(\alpha_s)$  fixed-order.
- \* Small  $p_T$ : expansion  $\approx$  fixed-order.
- \* Large  $p_T$ : expansion  $\approx$  resummation.
- \* Intermediate *p*<sub>T</sub>: enhancement.
- Scale dependence  $(M/2 \le \mu_R = \mu_F \le 2M)$ .
  - \* **Reduction** of the uncertainties.
  - \* Less than 5% for  $p_T > 5$  GeV.
- Parton densities dependence (44 CTEQ sets).
  - \* 4-5% uncertainties for all  $p_T$ .
  - \* Similar to weak boson production.
- Non perturbative effects at low  $p_T$ .
  - <sup>\*</sup> **Under control** for  $p_T > 5$  GeV.



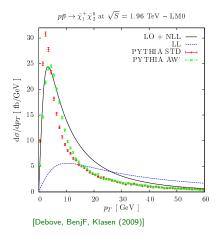
• Scenario.

- \*  $\,pprox\,$  180 GeV gauginos.
- \* LHC collider (10 TeV & 14 TeV).

### • $p_T$ -spectrum

- \* Next-to-leading logarithms.
- \*  $\mathcal{O}(\alpha_s)$  fixed-order.
- \* Small  $p_T$ : expansion  $\approx$  fixed-order.
- \* Large  $p_T$ : expansion  $\approx$  resummation.
- \* Intermediate *p*<sub>T</sub>: enhancement.
- Scale dependence  $(M/2 \le \mu_R = \mu_F \le 2M)$ .
  - \* Reduction of the uncertainties.
  - \* Less than 5% for  $p_T > 5$  GeV.
- Parton densities dependence (44 CTEQ sets).
  - \* 4-5% uncertainties for all  $p_T$ .
  - \* Similar to weak boson production.
- Non perturbative effects at low p<sub>T</sub>.
  - \* **Under control** for  $p_T > 5$  GeV.
- Uncertainties under control for  $p_T > 5$  GeV.

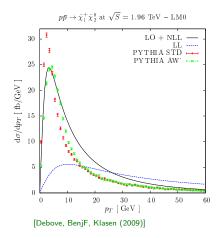
## Comparison: PYTHIA and $p_T$ -resummation



Scenario.

- pprox = 110 GeV gauginos.
- \* Tevatron collider.
- PYTHIA predictions.
  - \* Used for SUSY experimental analyses.
  - \* Leading log Sudakov form factor.
  - \* Two tunes.
    - ♦ CDF-AW.
    - ♦ Our tune AW'.
- Two set of resummed predictions.
  - \* Leading logaritmic approximation.
  - \* Next-to-leading logaritmic results.

# Comparison: PYTHIA and $p_T$ -resummation



Scenario.

- $^{\circ}~pprox$  110 GeV gauginos.
- \* Tevatron collider.
- PYTHIA predictions.
  - \* Used for SUSY experimental analyses.
  - \* Leading log Sudakov form factor.
  - Two tunes.
    - ◊ CDF-AW.
    - ♦ Our tune AW'.
- Two set of resummed predictions.
  - \* Leading logaritmic approximation.
  - \* Next-to-leading logaritmic results.
- Pythia results.
  - \* Improves the LL picture.
  - \* Intrinsic  $p_T$  helps to reproduce NLL.
  - \* **Underestimation** for intermediate  $p_T$ .
  - \* Direct impact for experimental analyses.

# Outline

### Models and motivation

Motivation for resummation calculations

### 2

#### Parton showers and resummation

- Parton showers
- Transverse-momentum, threshold and joint resummation formalisms
- Matching to the fixed order

#### <sup>8</sup> Numerical results, with uncertainties

- The Drell-Yan and the Tevatron
- Grand Unified Theories and Z' bosons
- The Minimal Supersymmetric Standard Model (MSSM)

### Summary - conclusions

### Summary - conclusions

- **Considered processes**: slepton-pair, gaugino-pair and Z' production.
- Soft and collinear radiation:
  - \* Large logarithmic corrections in  $p_T$  and invariant-mass spectra.
  - \* Need for resummation (or parton showers).
- p<sub>T</sub>, threshold and joint resummations have been implemented.
  - \* Reliable perturbative results.
  - \* Correct quantification of the soft-collinear radiation.
  - \* Important effects, even far from the critical regions.
  - \* Uncertainties from scales and parton densities under good control.
  - \* Reduced dependence on non-perturbative effects.
- Comparison with Monte Carlo generators
  - \* Significant shortcomings in normalization and shapes for PYTHIA.
  - \* MC@NLO reaches (almost) the same precision level as resummation. BUT: easier implentation in the analysis chains of any experiment.
- Implementation of other processes in the precision tools.