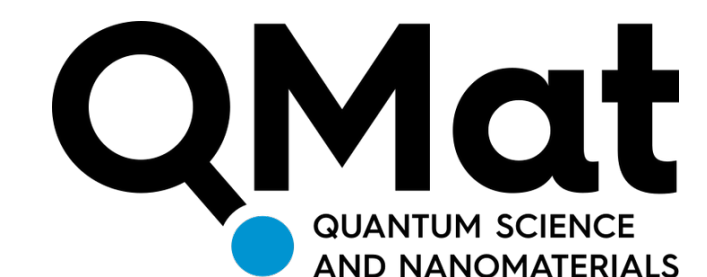


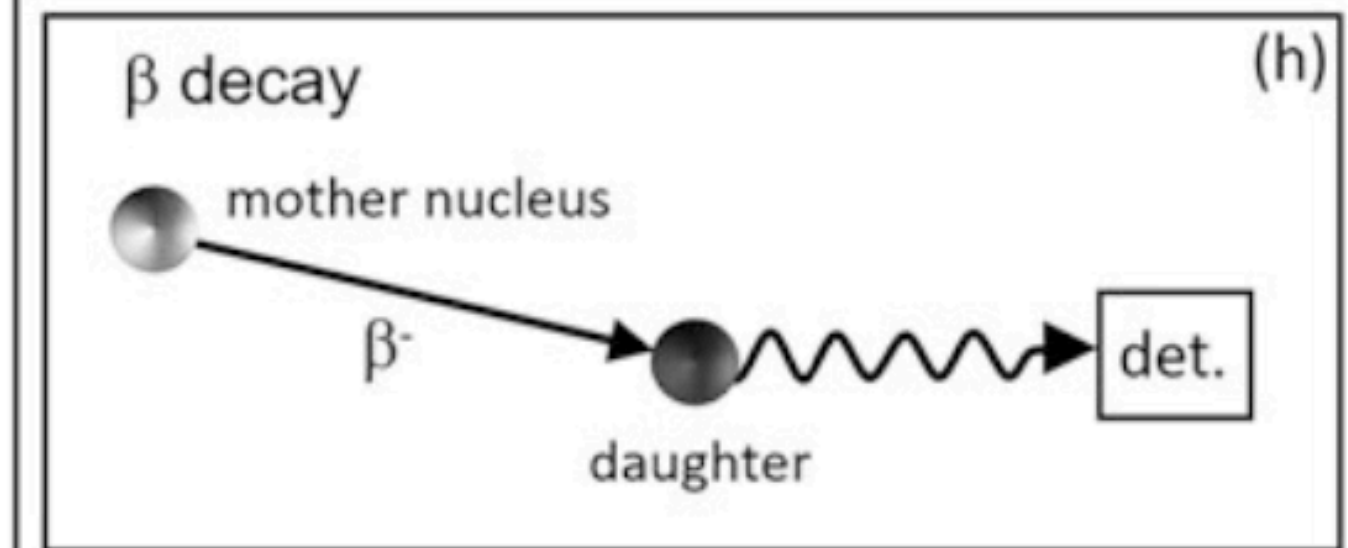
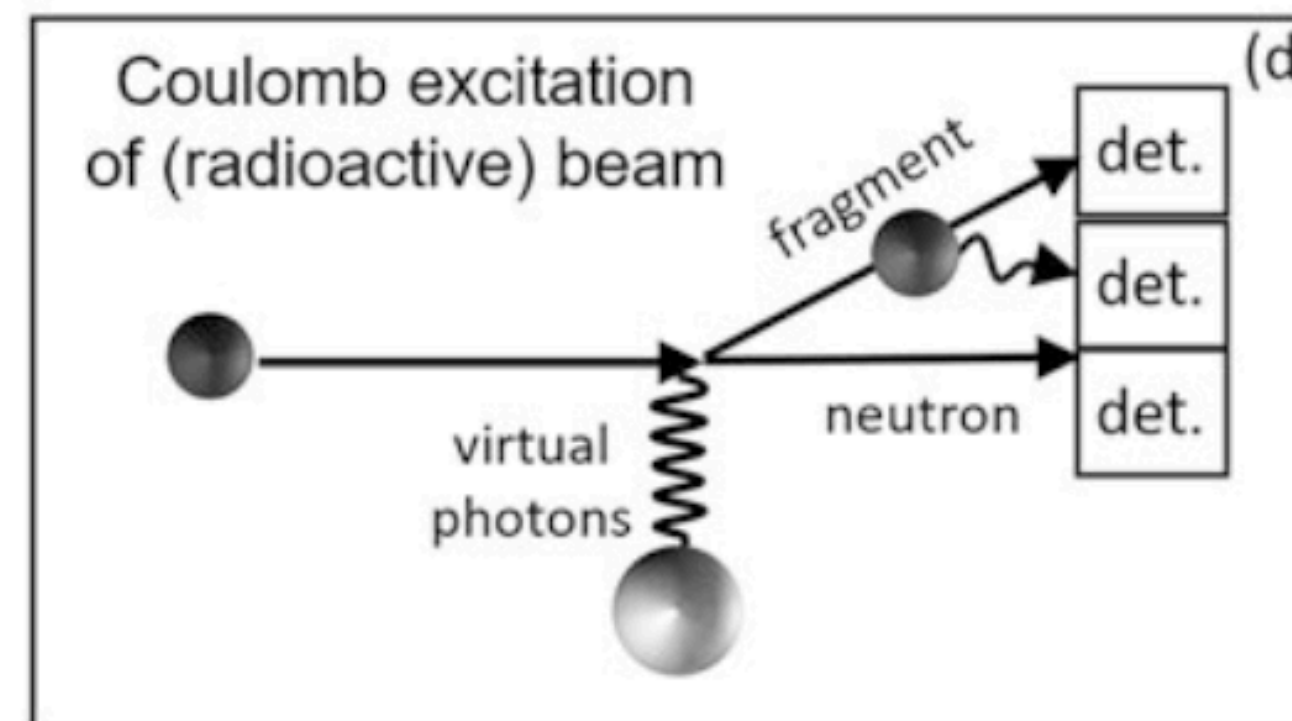
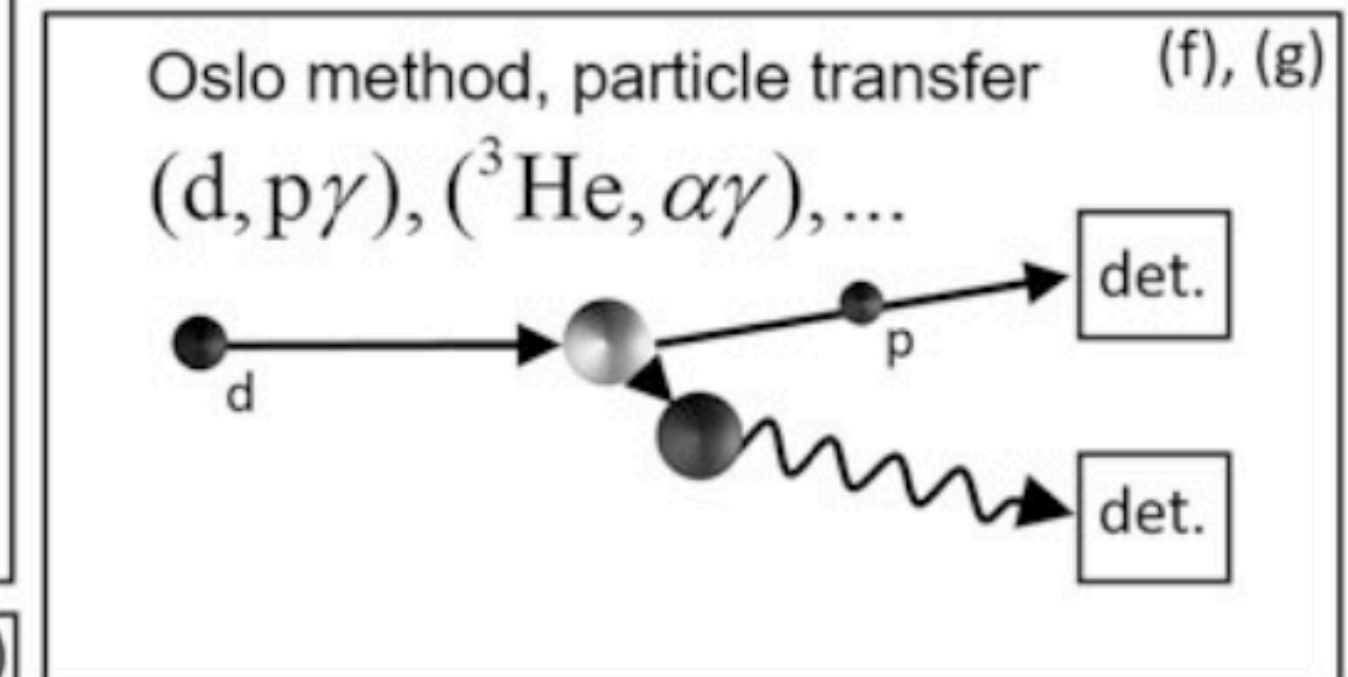
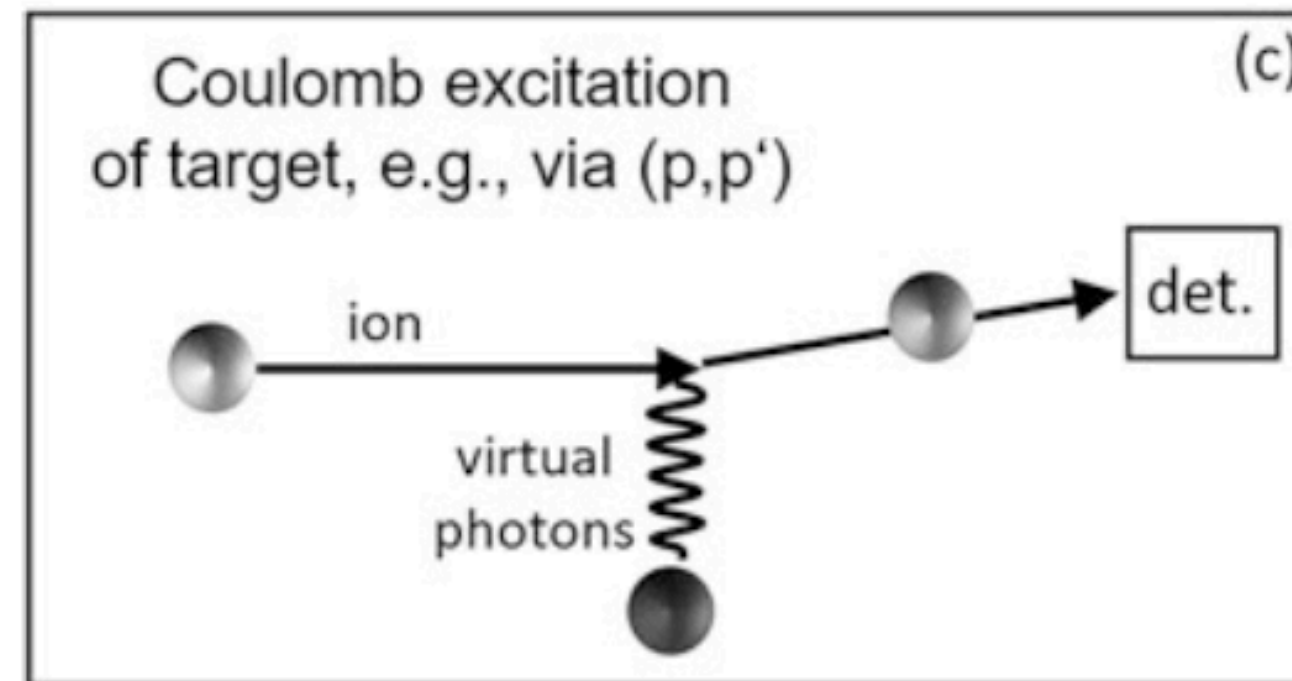
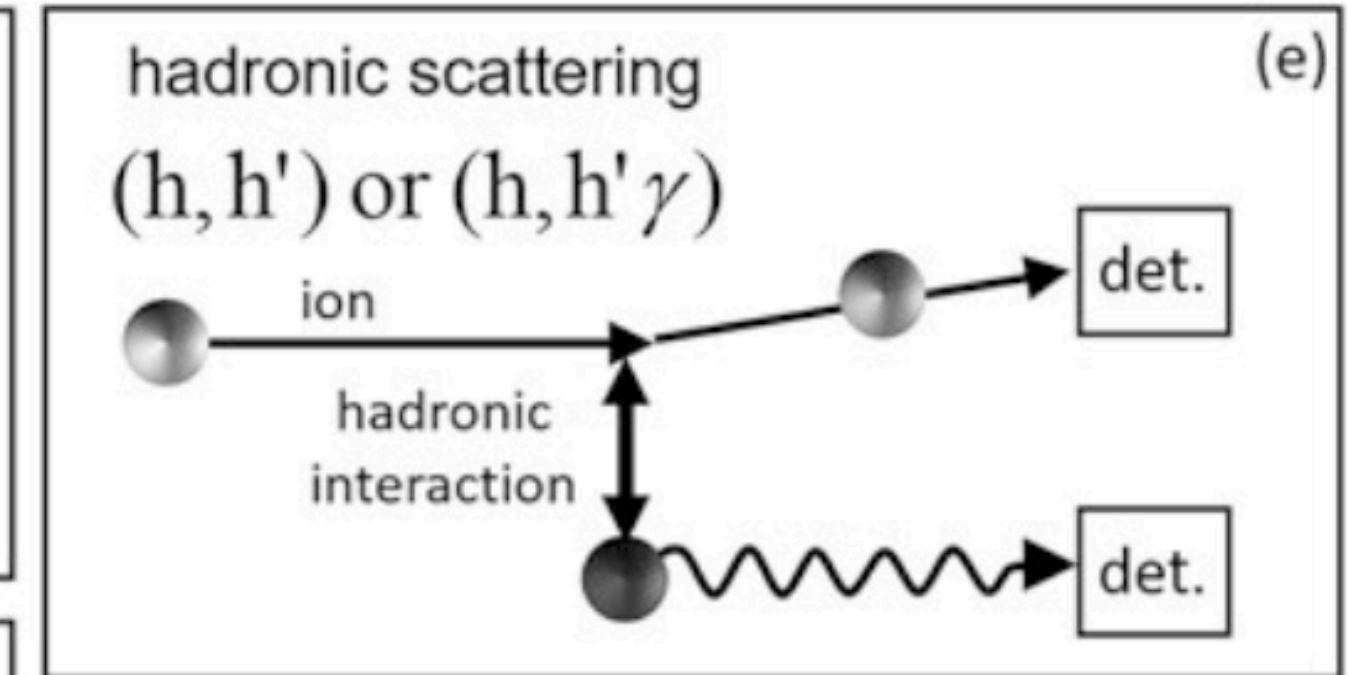
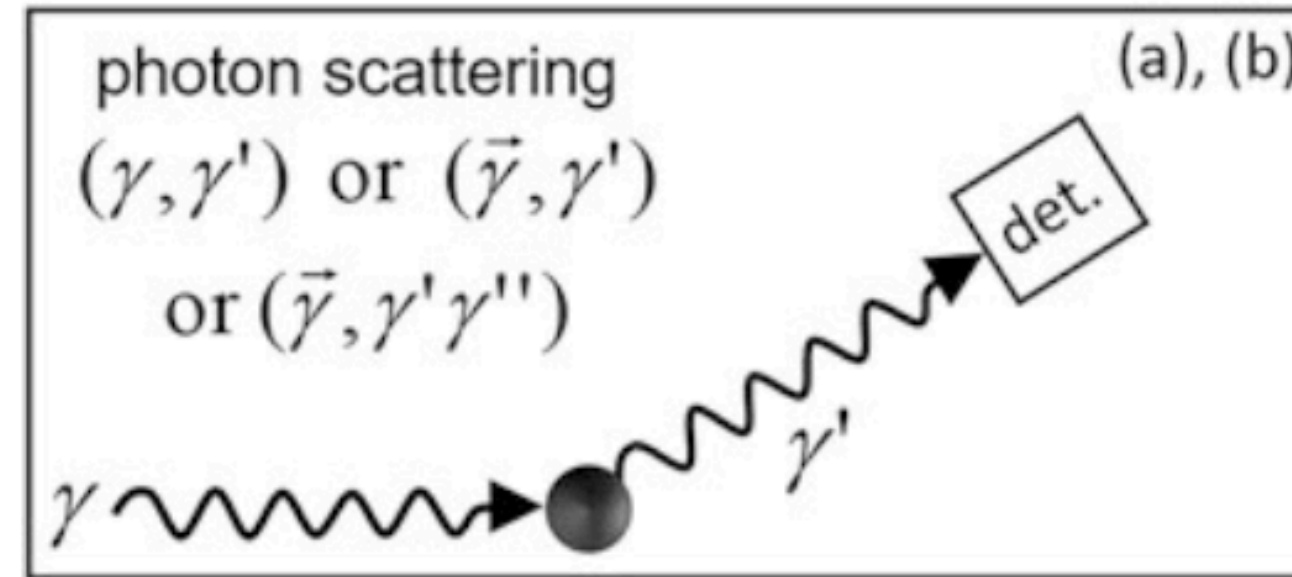
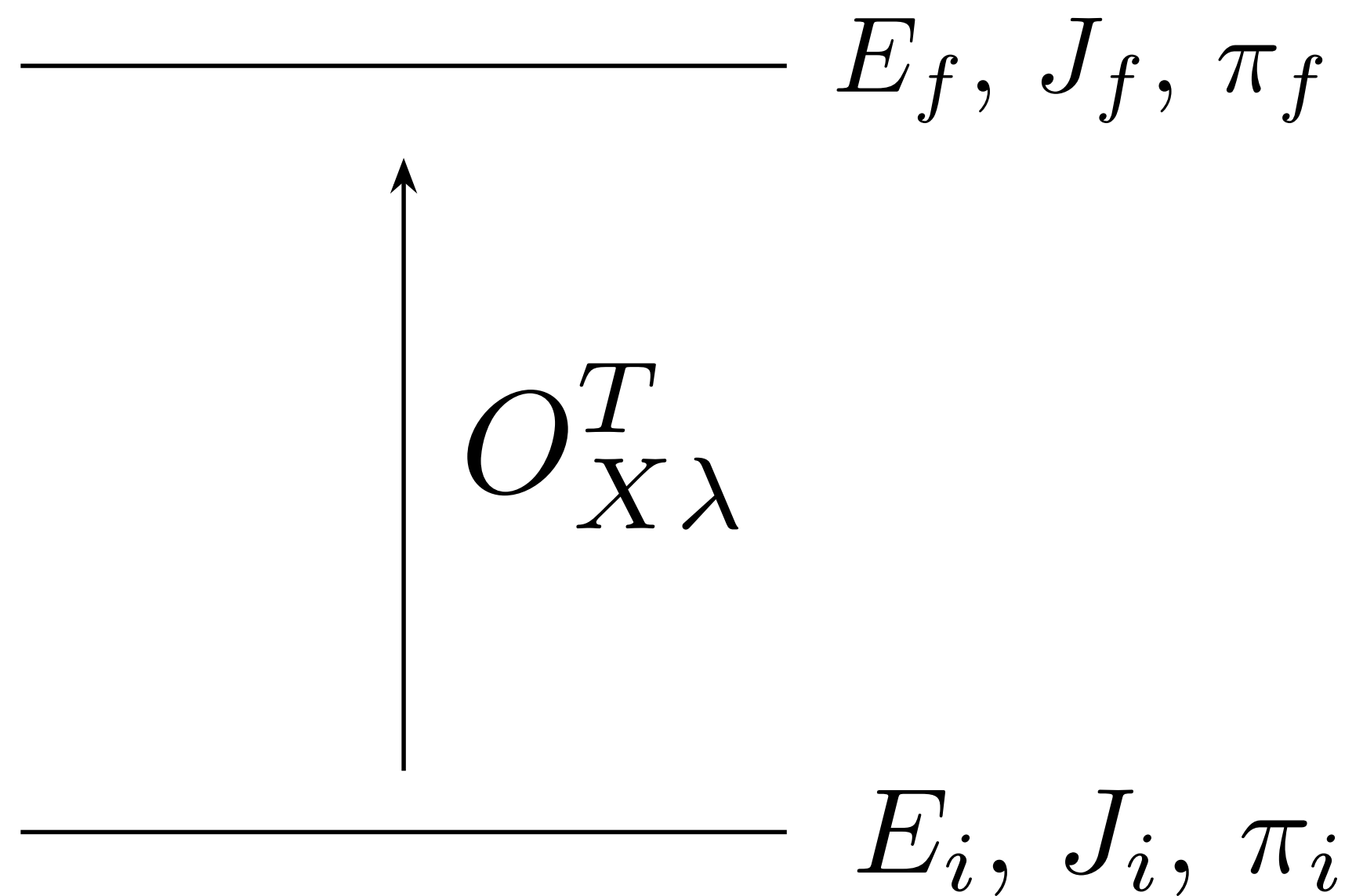
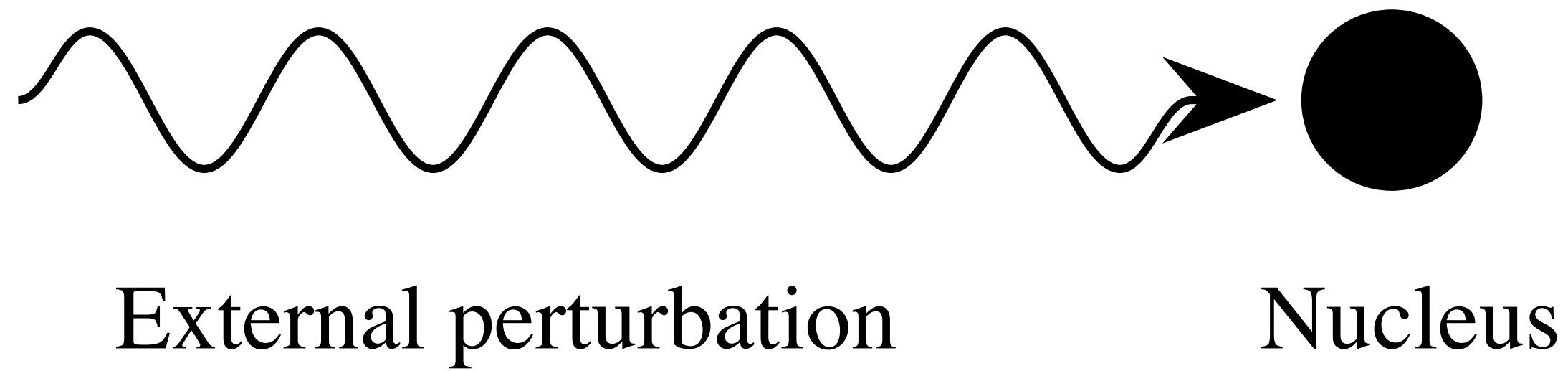
# A Configuration-Interaction Shell Model approach to the Pygmy Dipole Resonance

**Oscar Le Noan**

In collaboration with  
**Kamila Sieja**

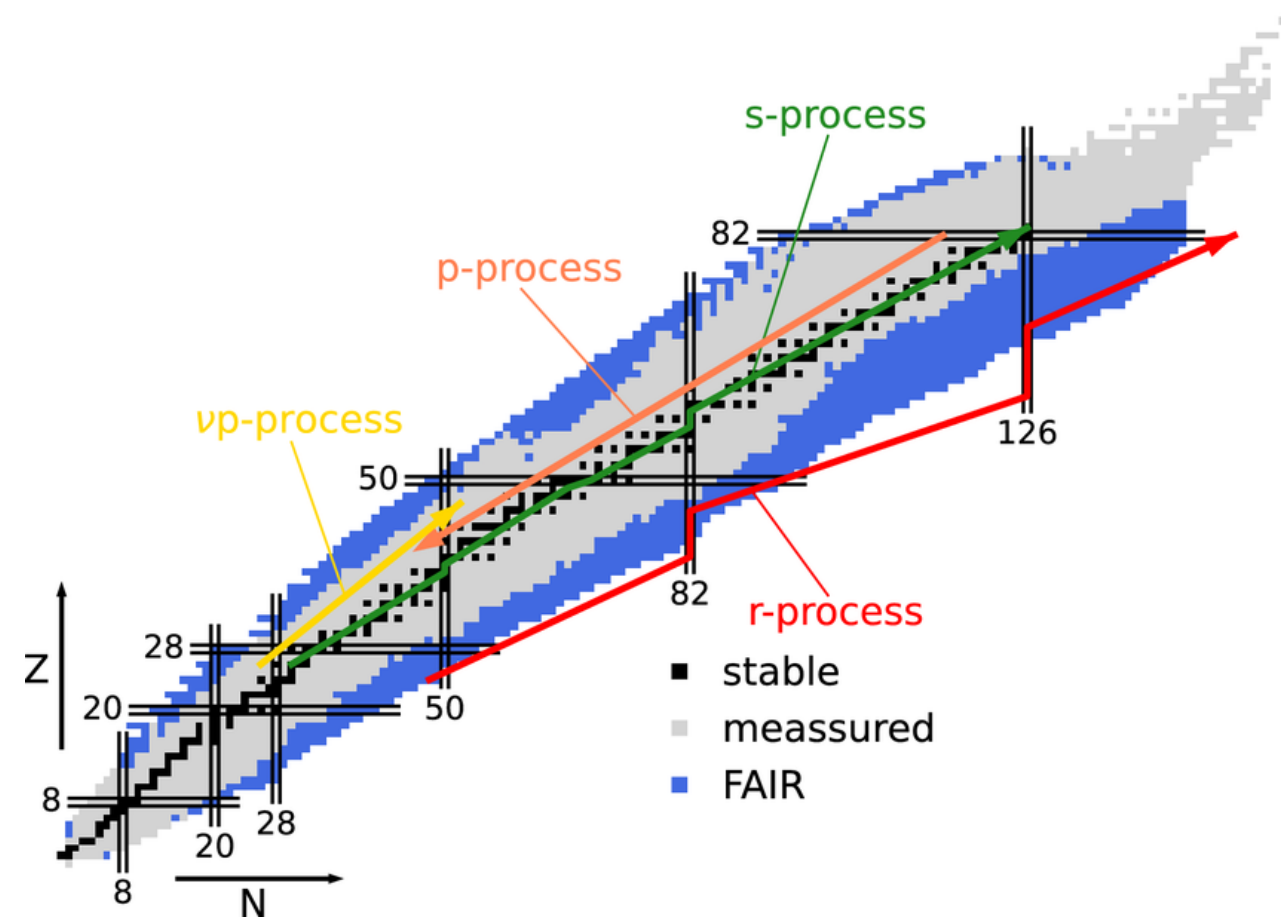


# Nuclear electric dipole response (E1): definition

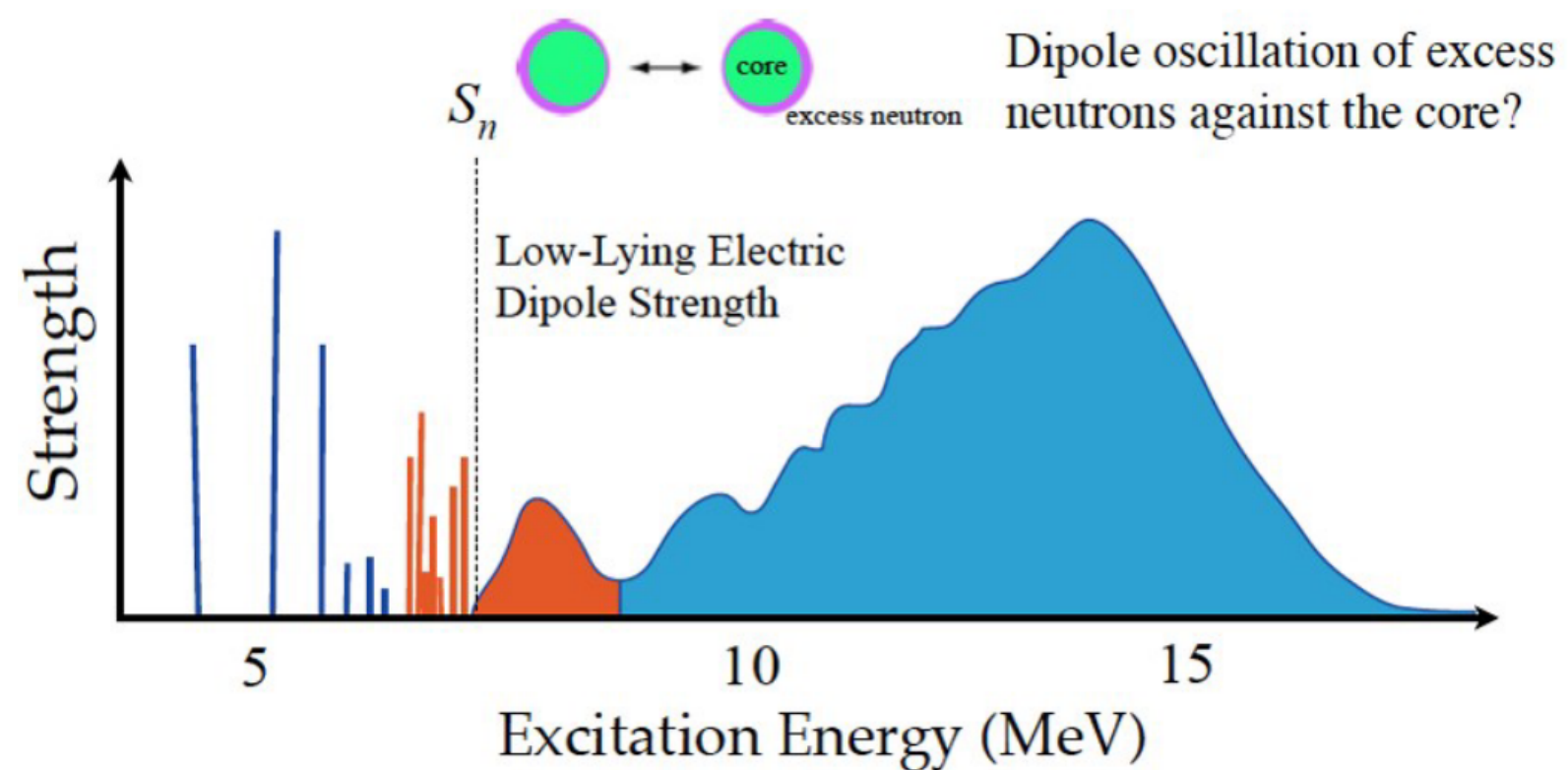


# Nuclear dipole response: motivation

## R-process nucleosynthesis



## Nuclear structure



## Neutron matter equation of state

$$\mathcal{E}(\rho, \alpha) = \mathcal{E}_{\text{SNM}}(\rho) + \alpha^2 S(\rho) + \mathcal{O}(\alpha^4)$$

$$\rho = (\rho_n + \rho_p) \quad \alpha = (\rho_n - \rho_p) / \rho$$

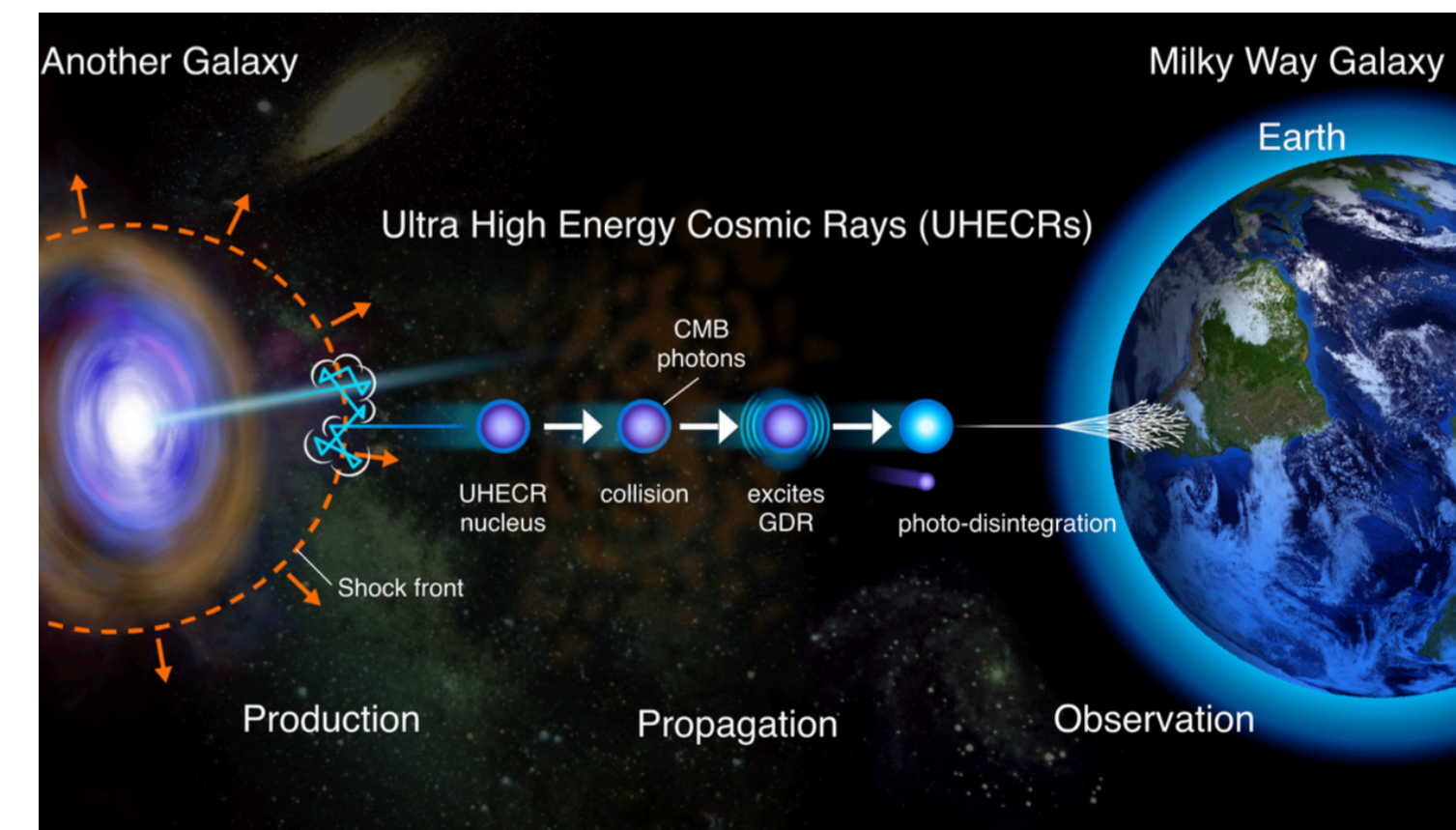
where the **density-dependent symmetry energy** is:

$$S(\rho) = J + L \frac{(\rho - \rho_0)}{3\rho_0} + \dots$$

**symmetry energy**  
at saturation  
density

**slope parameter**,  
related to **pressure of pure neutron matter** at saturation  
density

## Ultra High Energetic Cosmic Rays



# Theoretical framework: CI-SM

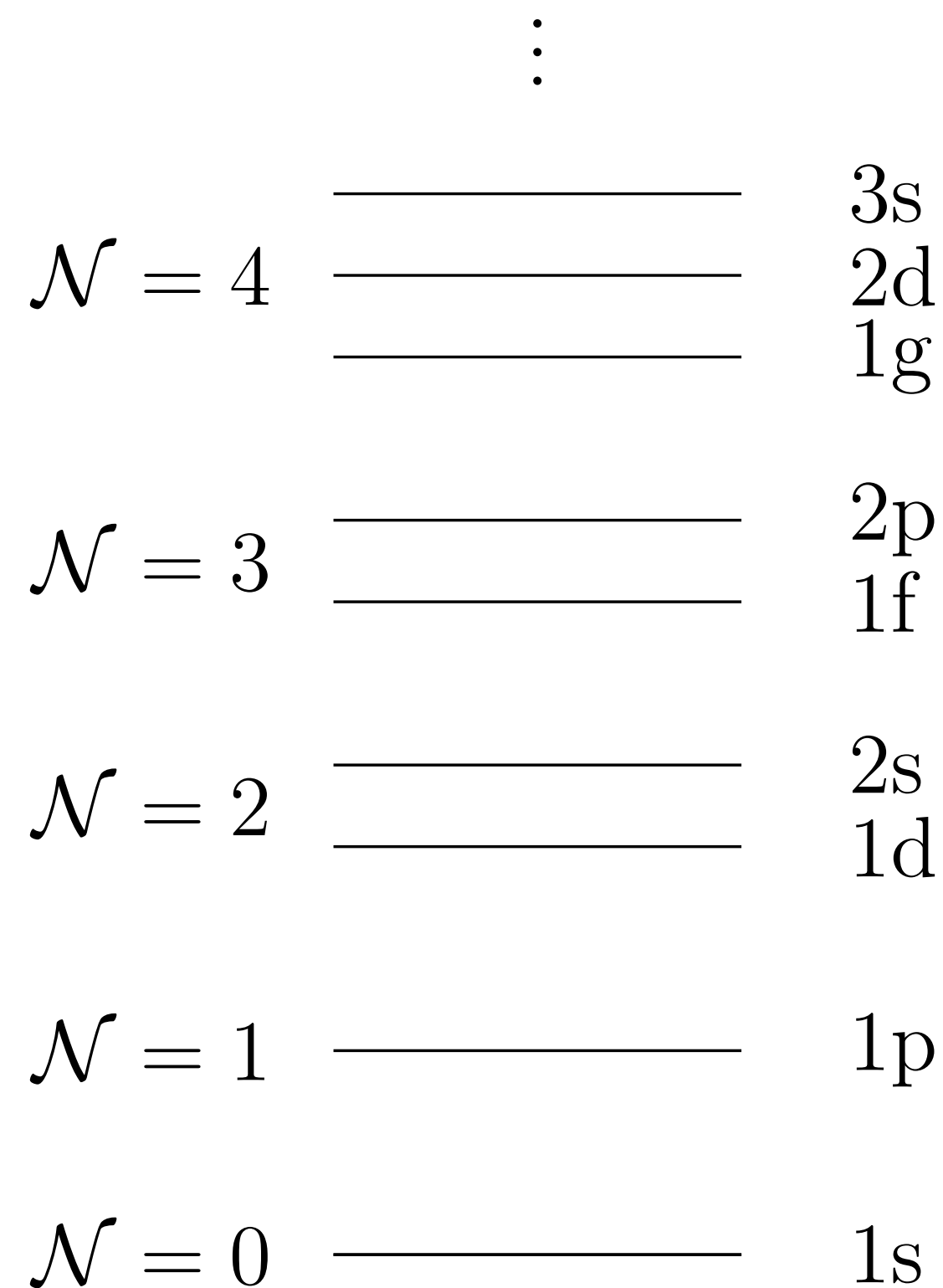
Quantity of interest: transition  
probability amplitudes

$$\langle \psi_f | |O_\lambda| | \text{GS} \rangle$$



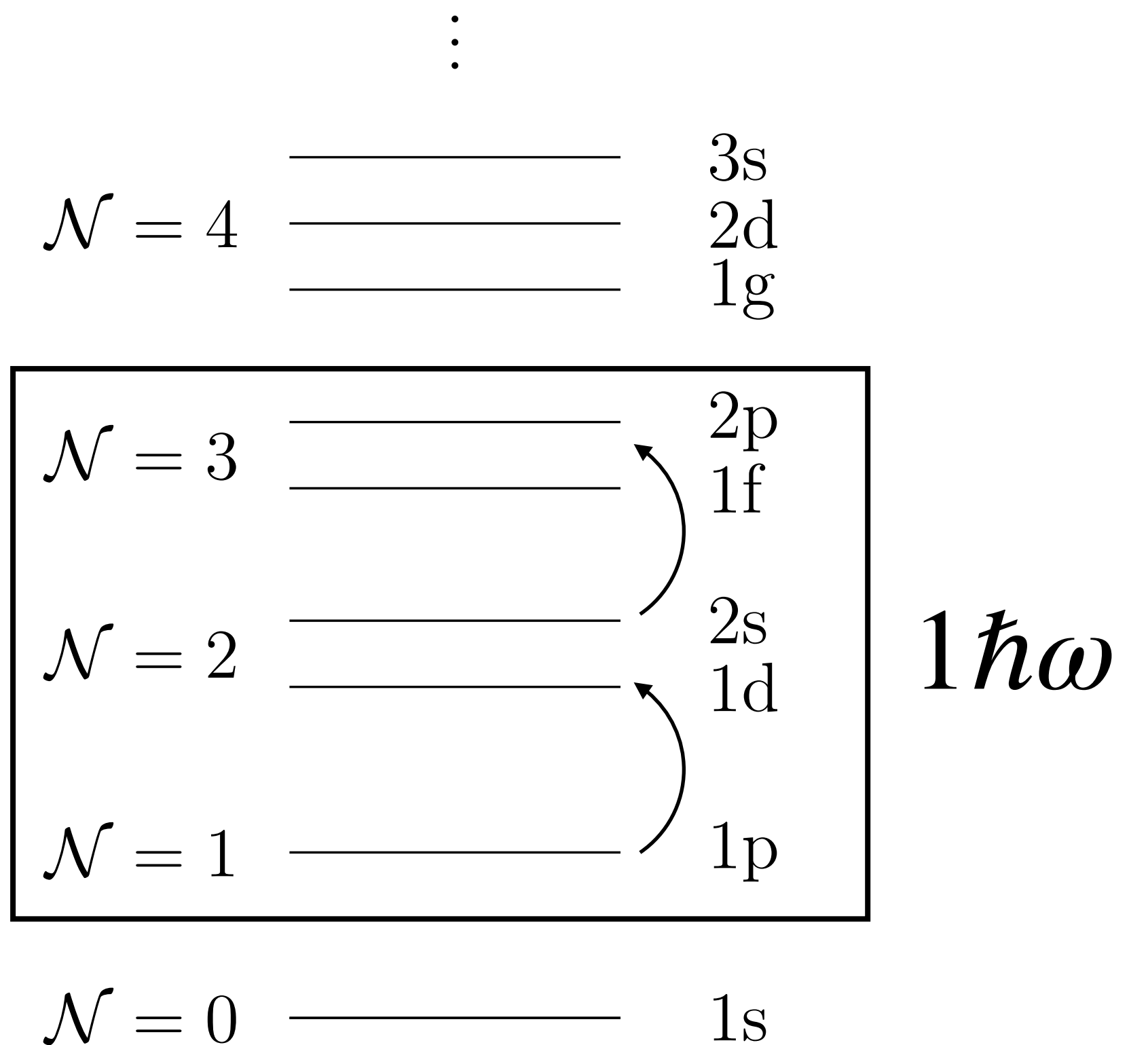
Discrete strength function

$$S(E) = \sum_f \frac{|\langle \psi_f | |O_\lambda| | \text{GS} \rangle|^2}{2J_{\text{GS}} + 1} \delta(E_f - E)$$



Valence space mapping:  
E1 sd-shell nuclei

$$\begin{aligned} \mathcal{H} &\longrightarrow \mathcal{H}_{\text{valence}} \\ H &\longmapsto H_{\text{eff}} \end{aligned}$$



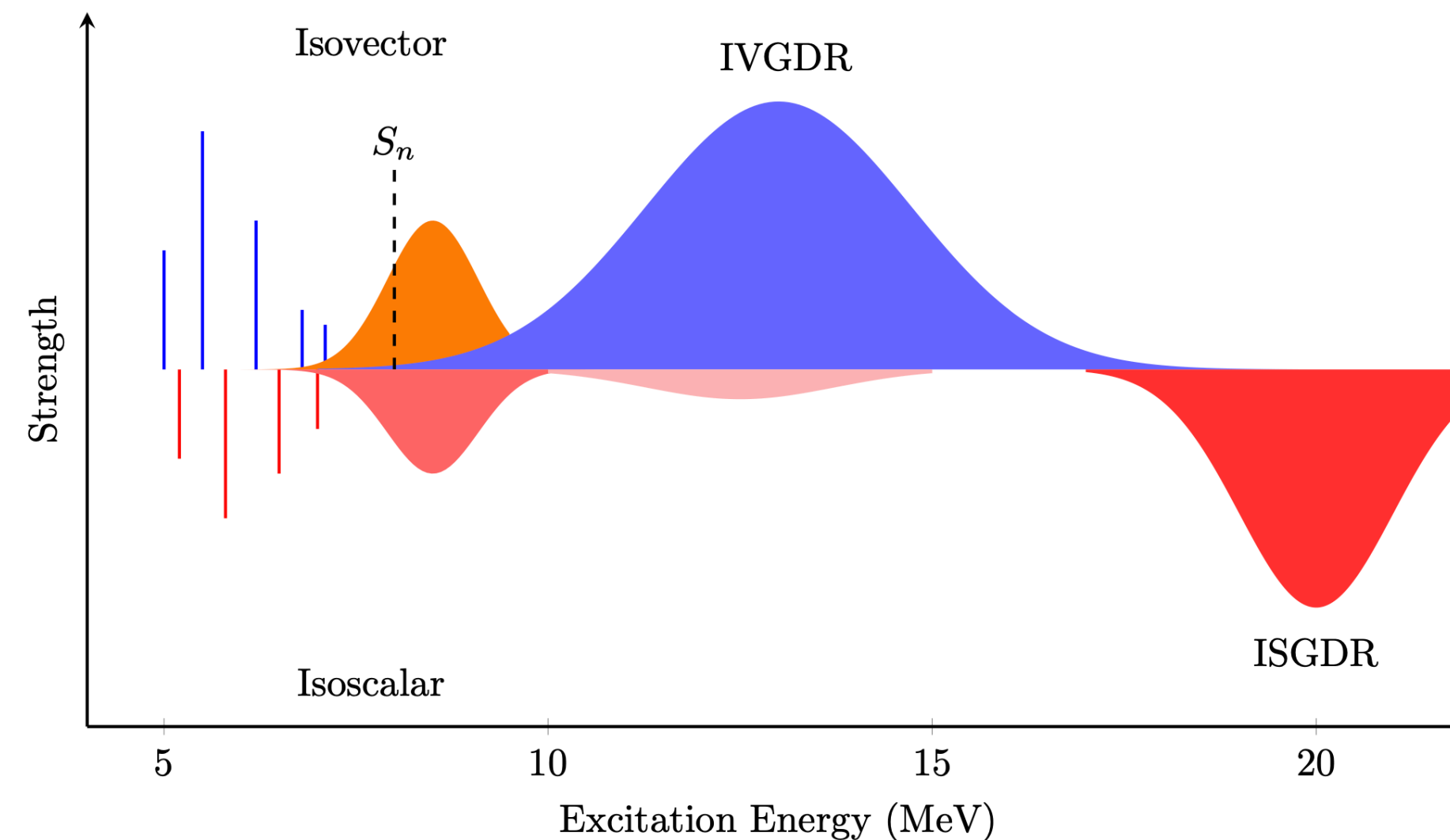
# Study of the Pygmy dipole resonance

## Definition

Low-lying E1 strength appearing in neutron-rich nuclei

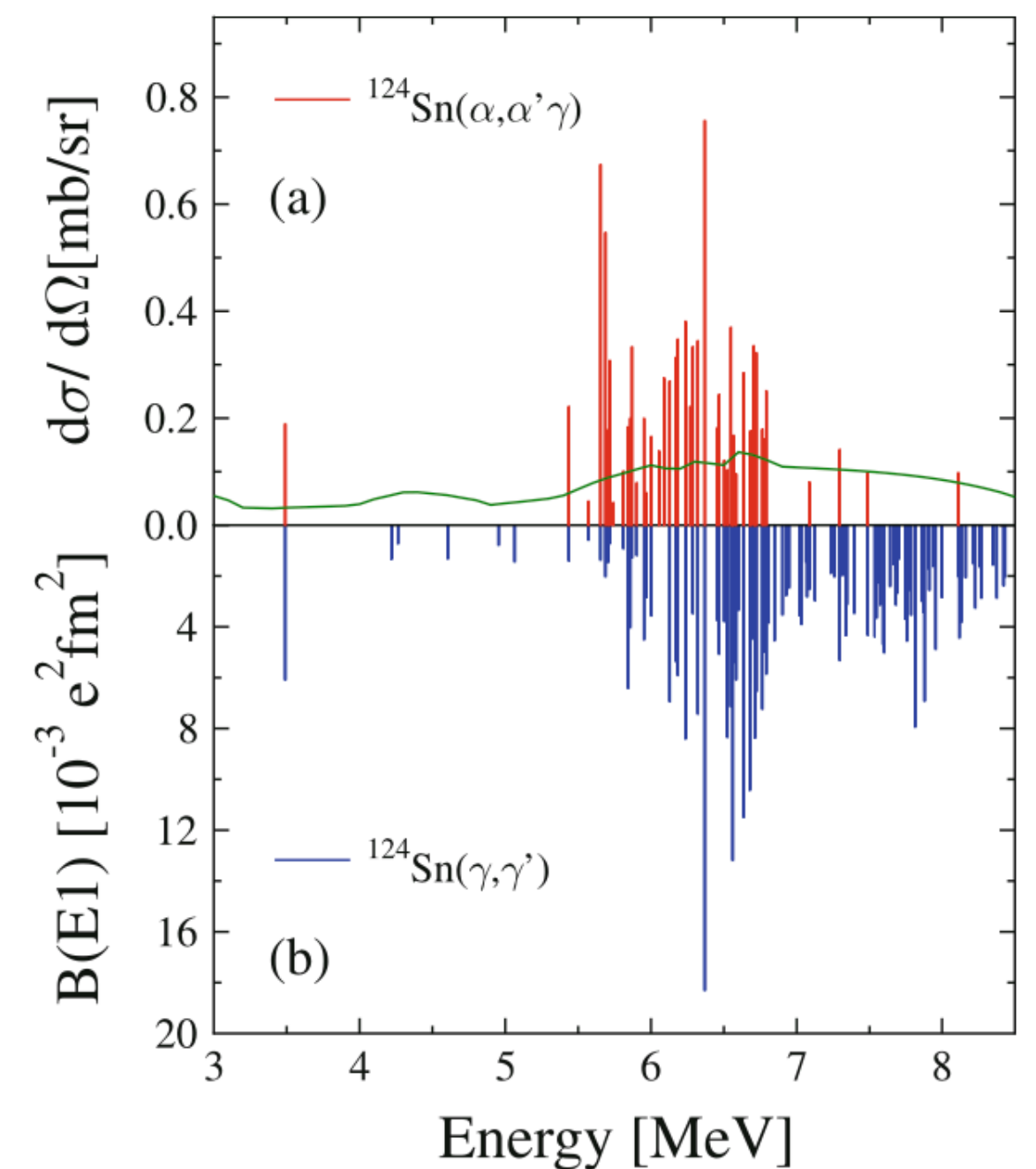
## State of the art

- Observed only in neutron-rich nuclei, both stable and unstable.
- Found both below and above the particle emission threshold and exhibits strong isospin mixing.
- In medium-mass and heavy nuclei, there is an isospin splitting of the PDR.



$$O_{E1} = O_{E1}^{(0)} + O_{E1}^{(1)}$$

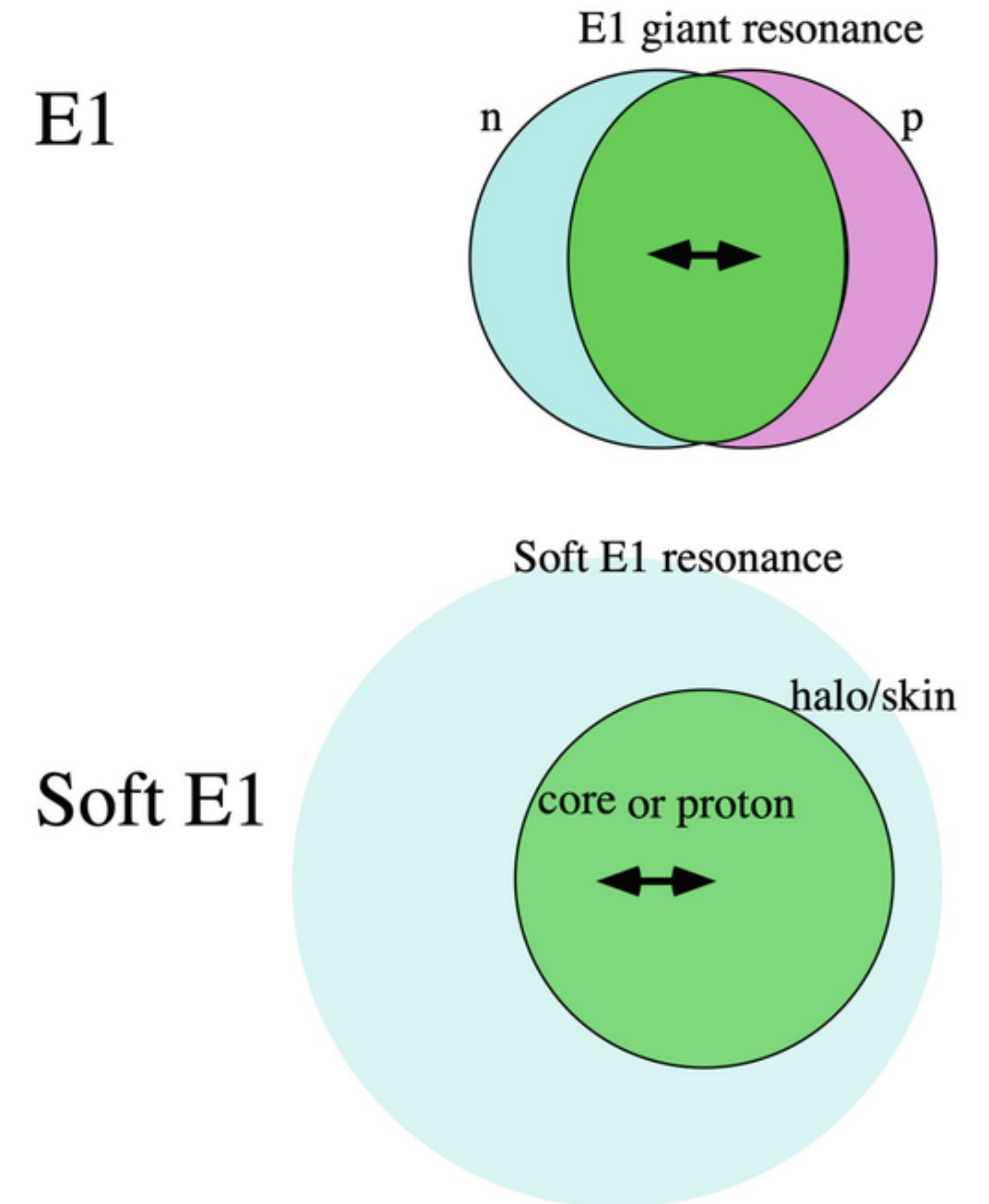
$$\begin{array}{cc} \text{Isoscalar} & \text{Isovector} \\ \sim r^3 Y_1 & \sim r Y_1 \tau_z \end{array}$$



# Study of the Pygmy dipole resonance

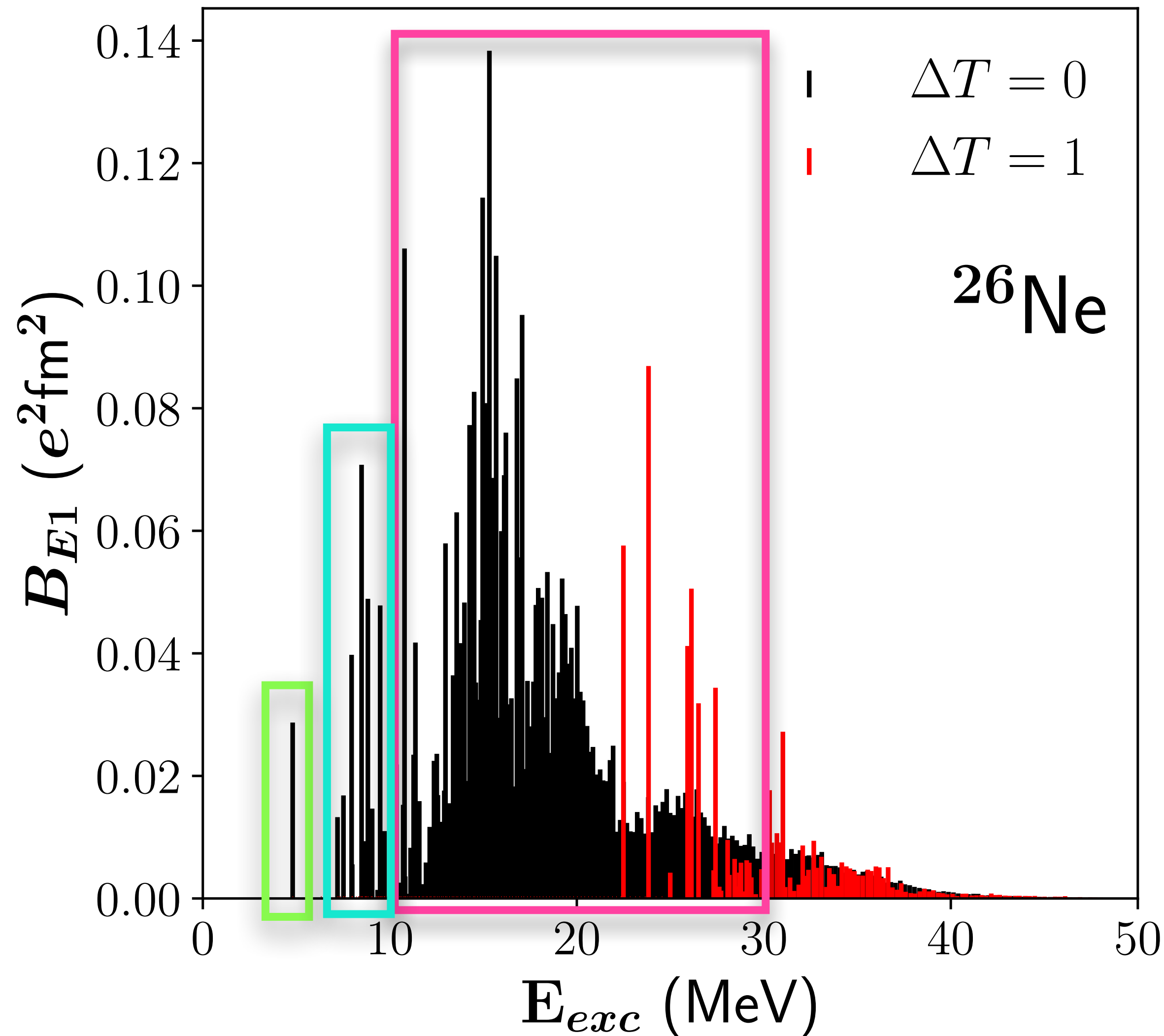
## Open questions:

- Degree of collectivity ?
- Is the isospin splitting of the PDR an ubiquitous property ?
- Classical representation : neutron excess oscillation ?
- In the PDR energy region, are there states that cannot be interpreted as neutron-skin oscillation ?
- How is the transition between PDR and GDR states ?
- ...



# Study of the Pygmy dipole resonance

Comparison to experiment



CI-SM

$$\bar{S}_{\text{PDR}} = 8.63 \text{ MeV}$$

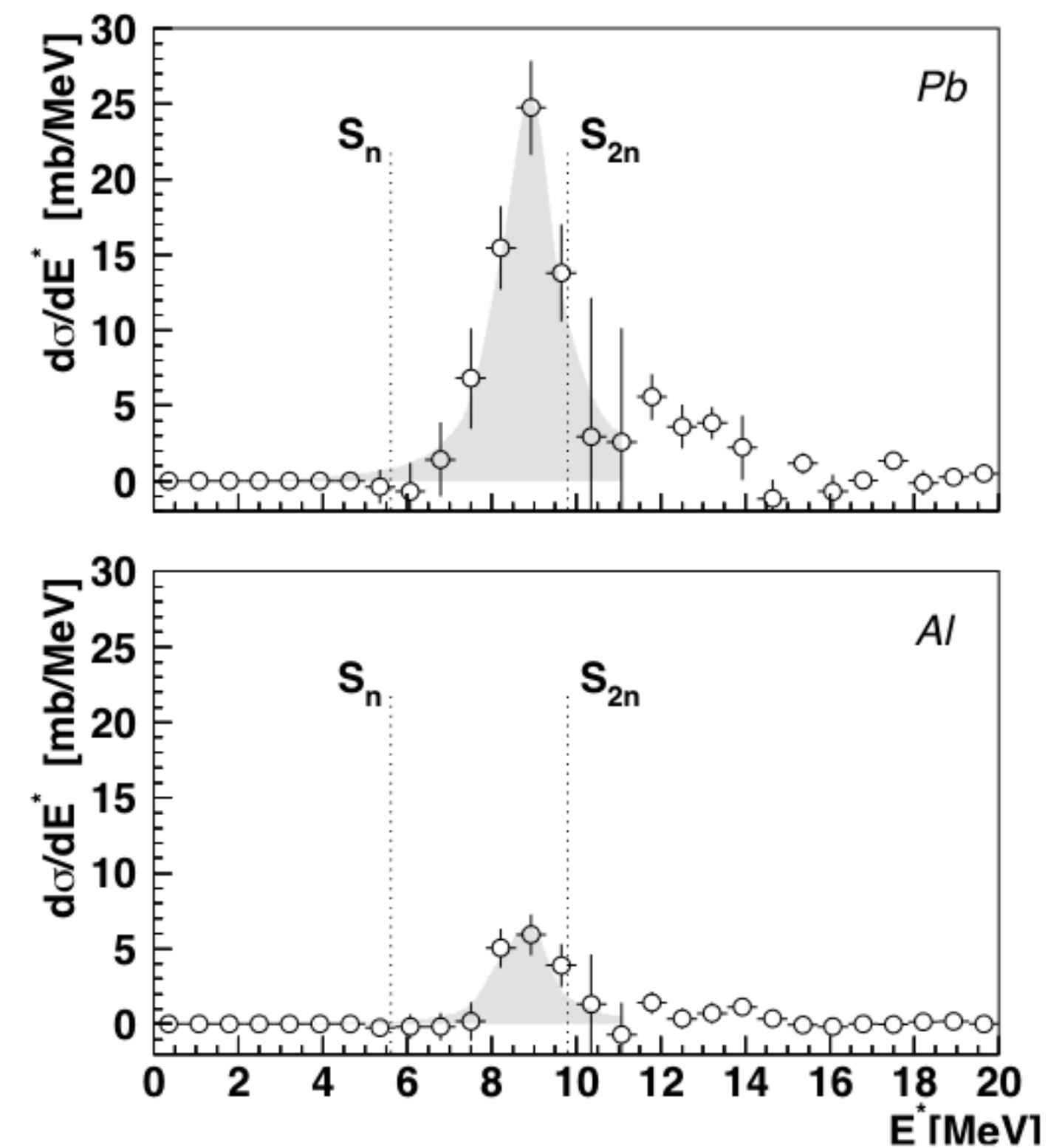
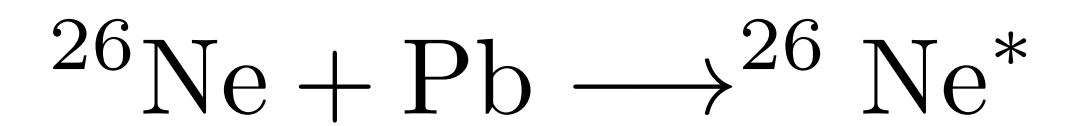
$$S_0^{\text{PDR}} = 0.37 e^2\text{fm}^2$$

$\sim 4\%$  TRK

EXP

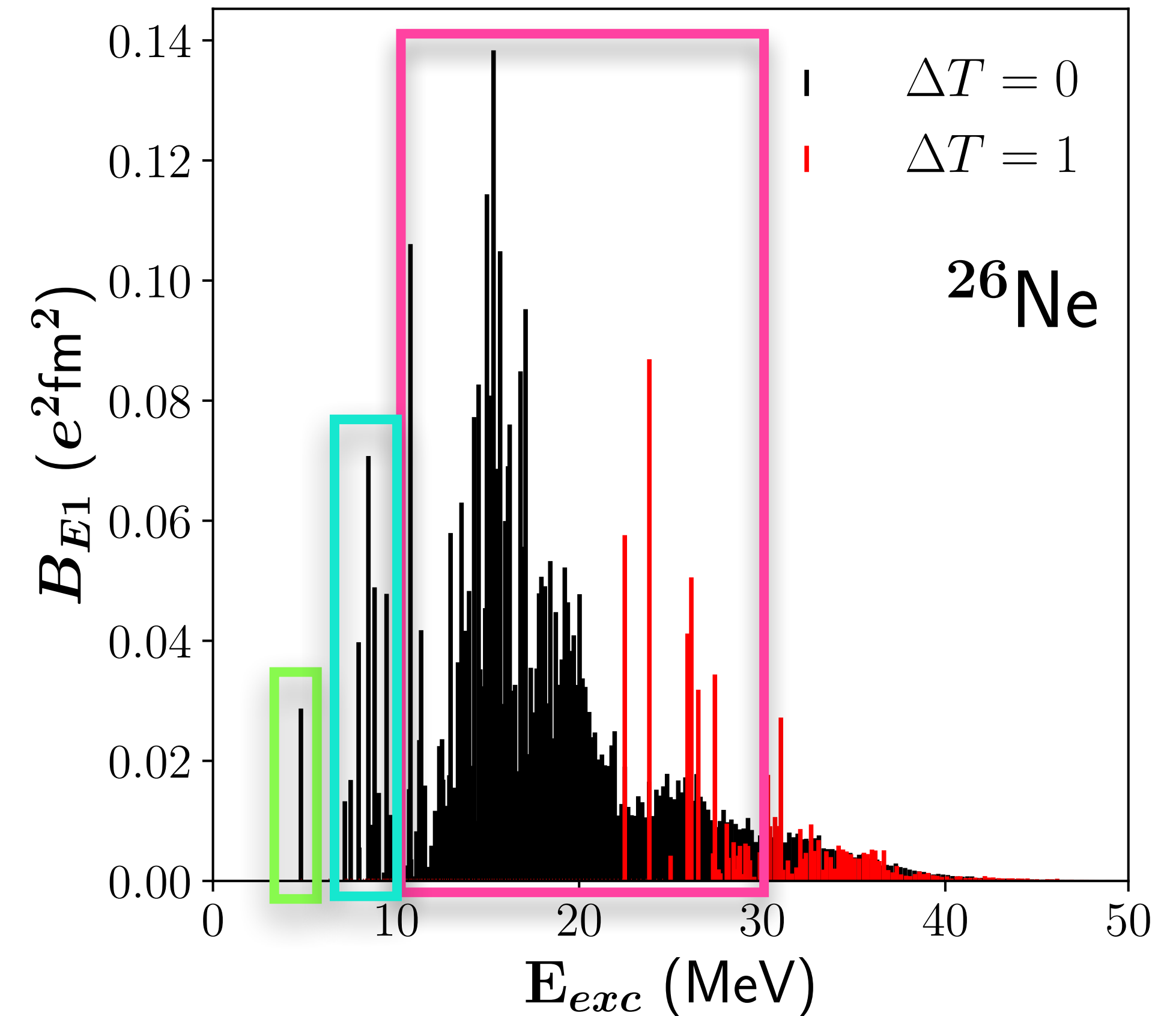
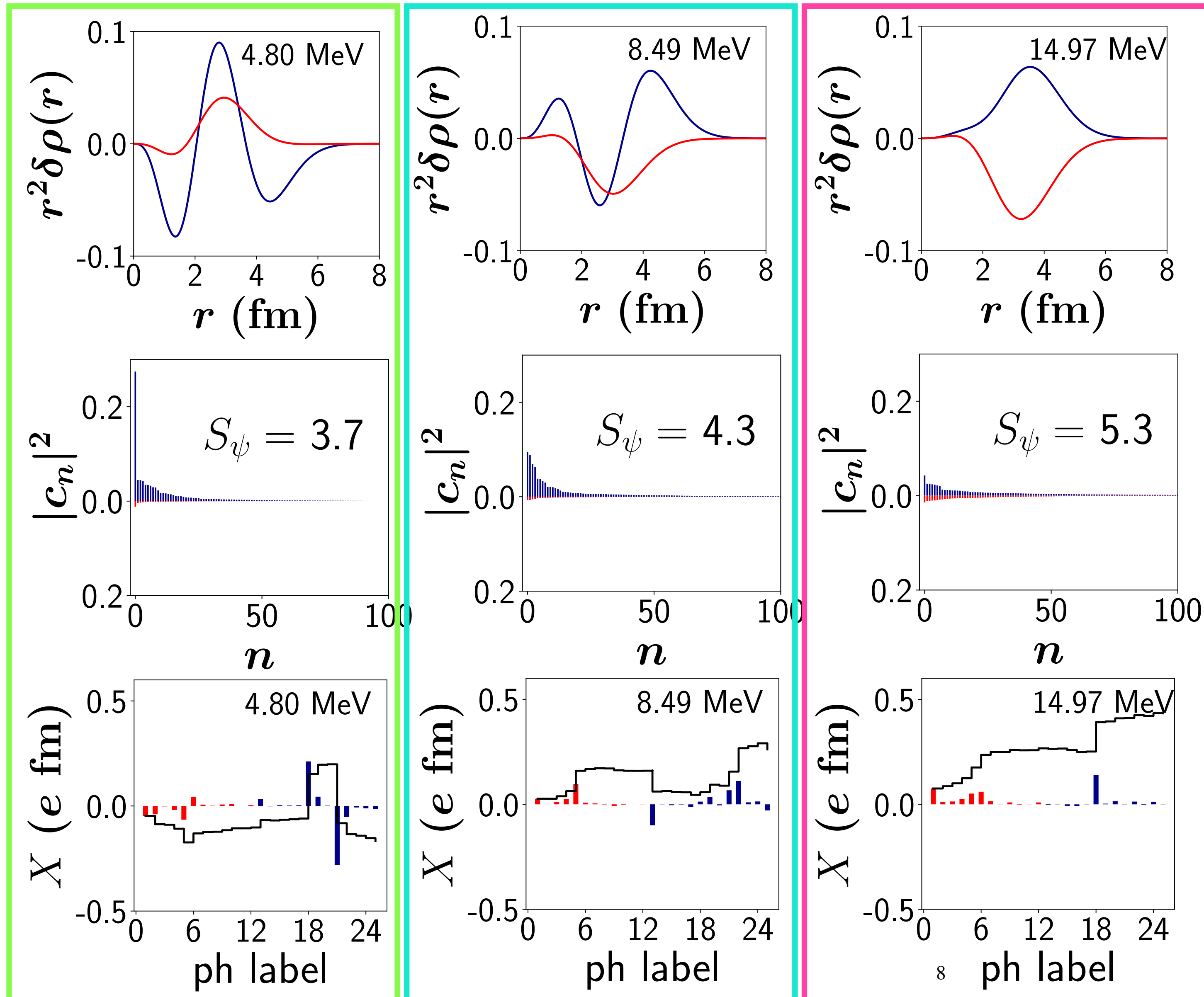
$$\bar{S}_{\text{PDR}} \approx 9 \text{ MeV}$$

$$S_0^{\text{PDR}} = 0.49 \pm 0.16 e^2\text{fm}^2$$



Gibelin, J. et al. Nuclear Physics A 2007, 788, 153–158.

# Study of the Pygmy dipole resonance $^{26}\text{Ne}$



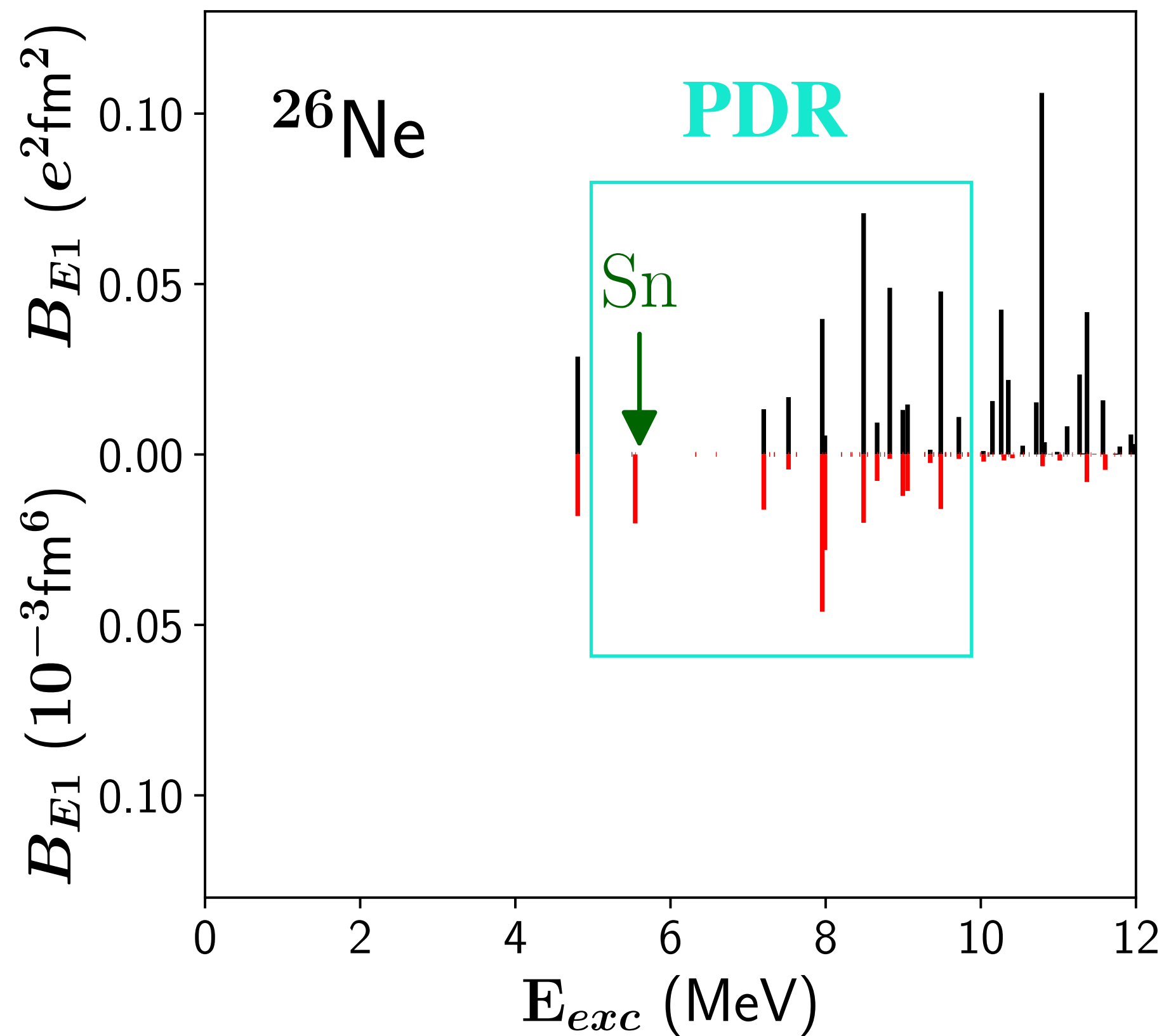
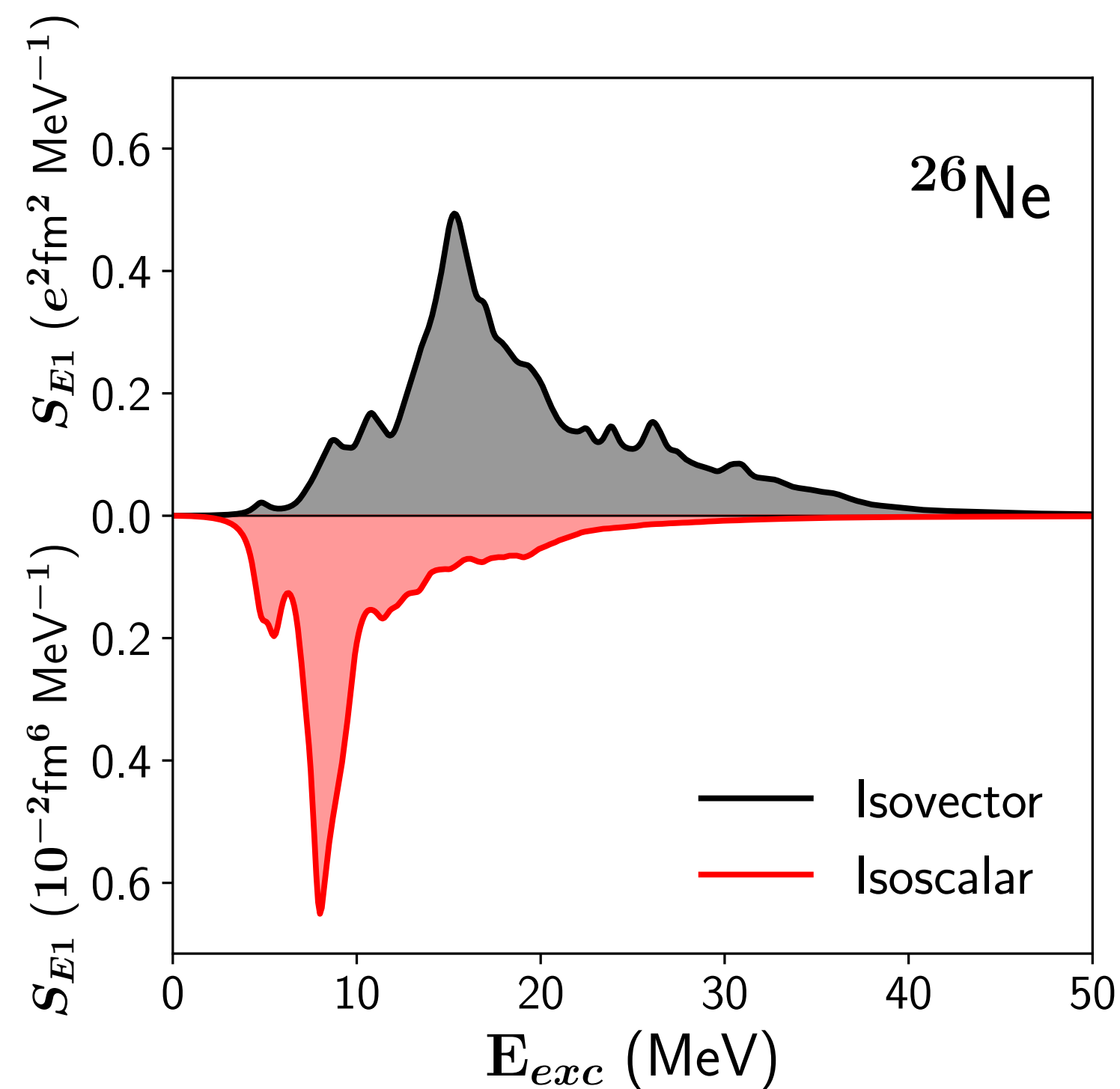
O. Le Noan and K. Sieja,  
Physical Review C 111,  
064308 (2025)

# Study of the Pygmy dipole resonance $^{26}\text{Ne}$

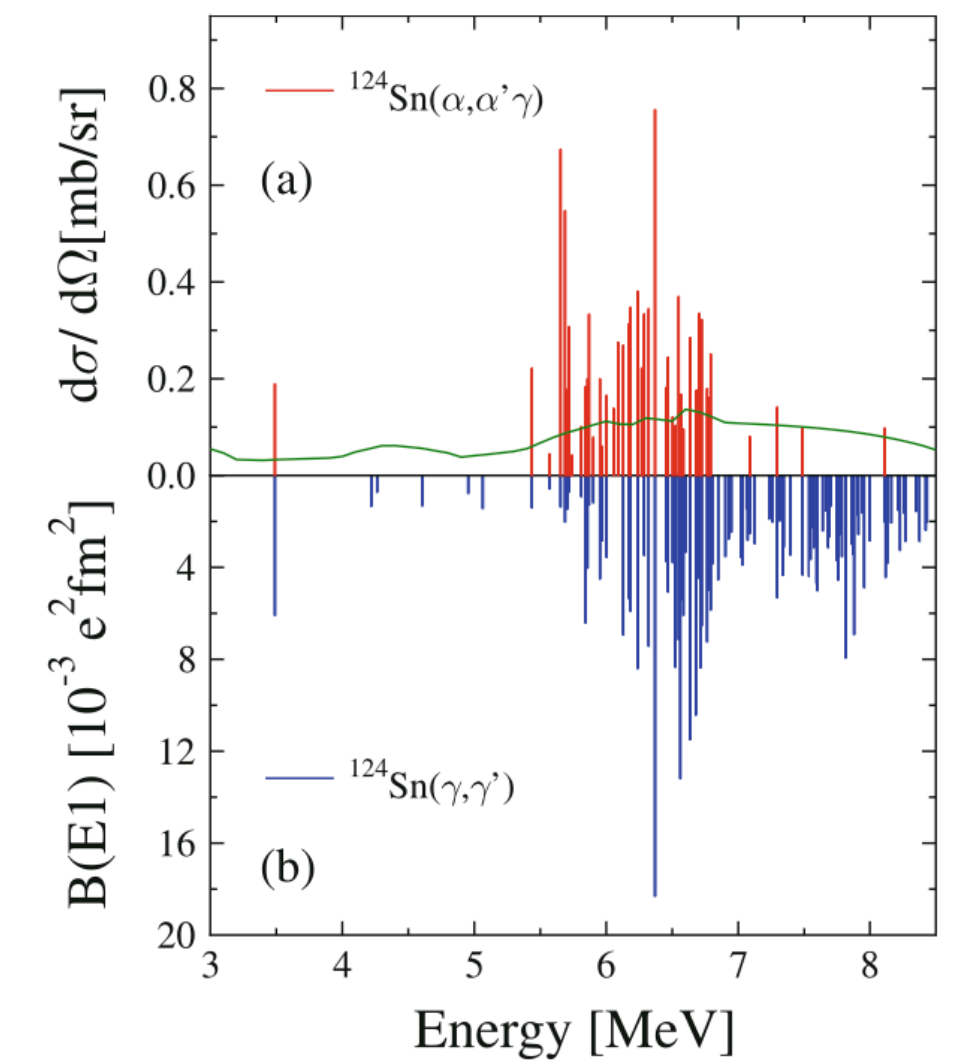
Isospin splitting of the PDR

$$O_{E1} = O_{E1}^{(0)} + O_{E1}^{(1)}$$

Isoscalar      Isovector  
 $\sim r^3 Y_1$      $\sim r Y_1 \tau_z$



Experimentally:



# Conclusion

## Systematics & UHECR

- CI-SM provides good prediction of isovector E1 response for  $A < 40$ .
- UHECRs modelisation is essentially sensitive to the centroid position in nuclei close to stability.
- First CI-SM description of low-lying isoscalar E1 strength, good agreement with exp.
- Future work: CI-SM for pf-shell nuclei ( $A > 40$ ).

## PDR

- We predict three kinds of low-lying dipole state: quadrupole-octupole, IS-LED and Pygmy.
- The PDR is collective in a shell-model picture ( $< \text{GDR}$ ).
- The classical picture of a neutron skin oscillating against an  $N=Z$  core holds.
- We predict an isospin splitting of the PDR above particle emission threshold (! depends on the definition of the PDR...)
- Future work: toroidal nature of the PDR, E1 from PGCM

**Back up**

# Theoretical landscape

Models commonly used to predict PSFs

## Purely phenomenological approaches

- Simple Modified Lorentzian (SMLO)
- Artificial Neural Network (ANN)

## Microscopic approaches

- Quasi-particle Random Phase approximation (QRPA)
- Quasi-particle Finite Amplitude Method (QFAM)
- Second QRPA (SQRPA)
- QRPA + phonon coupling
- Projected Generator Coordinates Method (PGCM)
- Configuration Interaction Shell Model (CI-SM)
- ...

Linear response

Beyond linear response

# Theoretical framework: CI-SM

Dressing of the E1 operator  $O(E1) = e \sum_{j=1}^Z r Y_{10}$

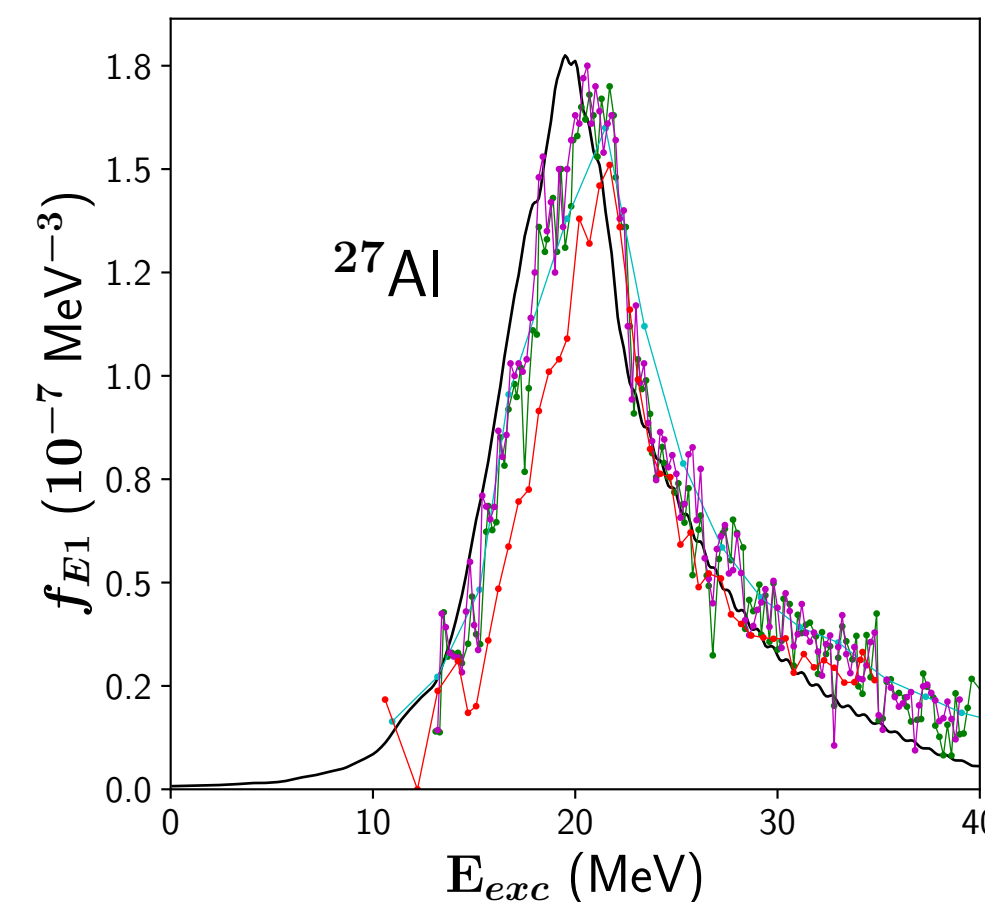
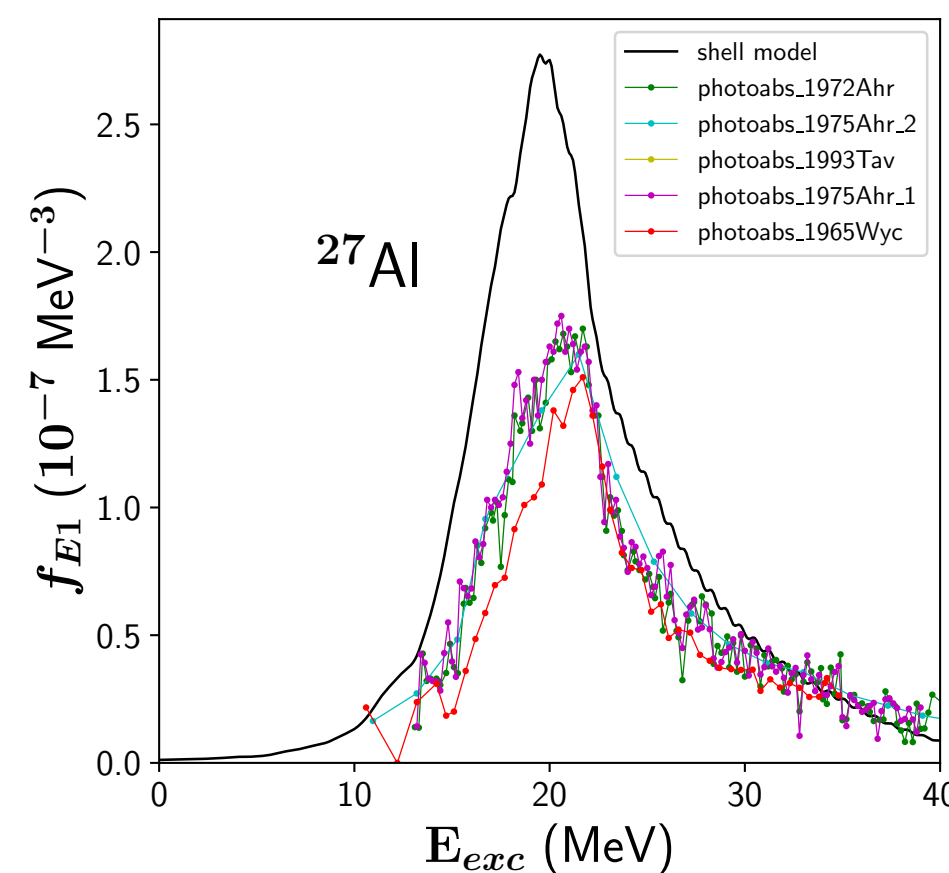
Valence space mapping:  $\mathcal{H} \longrightarrow \mathcal{H}_{\text{valence}}$  Likewise for all observables  $\mathcal{H} \longrightarrow \mathcal{H}_{\text{valence}}$   
 $H \longmapsto H_{\text{eff}}$   $O(E1) \longmapsto O_{\text{eff}}(E1)$

How to address this issue ?

Microscopic approaches (many body perturbation): Lee-Suzuki, IMSRG

Phenomenological approaches: scaling to data (effective charges),

scaling to TRK  $S_1^{TRK} = \frac{9\hbar^2 e^2}{8\pi m} \frac{NZ}{A}$



Details of numerical calculations

300 Lanczos iterations

Lawson's method for COM

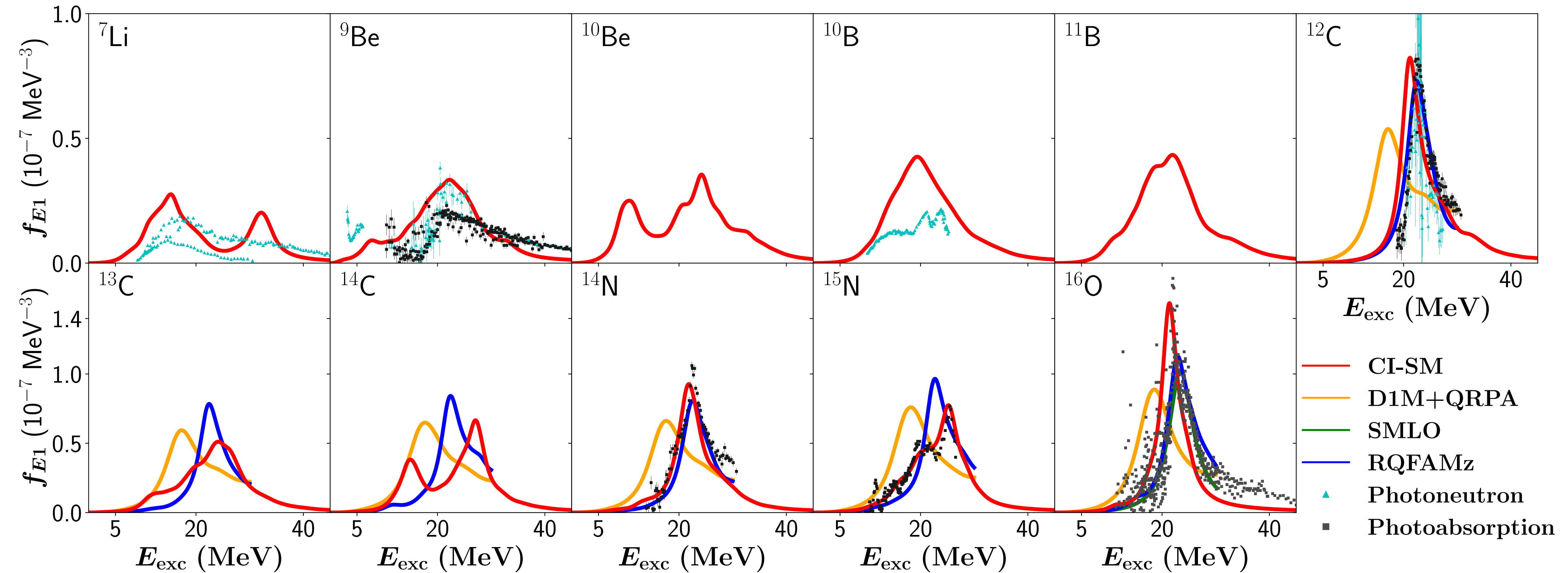
Effective interactions:

- E1 sd-shell: PSDPF
- E1 p-shell: WBP

# E1 response p-shell nuclei

O. Le Noan and K. Sieja, Physical Review C, vol.111, Jun.2025

O. Le Noan S. Goriely, E. Khan, K. Sieja, To be published in PRC.



# Theoretical framework: CI-SM

## Lanczos Strength function

Step 1: compute the GS  $[H_{\text{eff}}, J^z] = 0 \implies H_{\text{eff}}$  bloc diagonal in  $M$

For a  $0^+$  GS we diagonalize  $H_{\text{eff}}$  in the  $M = 0$  subspace using Lanczos algorithm  $\implies |\text{GS}\rangle$

Step 2: compute the sum rule state  $|\text{SR}\rangle = O |\text{GS}\rangle$

Step 3: Lanczos calculation using  $|\text{SR}\rangle$  Unitary matrix  $U_{ij} = \langle \mathcal{L}_i | \psi_j \rangle$  diagonalizing  $H_{\text{eff}}$  such that  $U_{1j} = \langle \text{GS} | O | \psi_j \rangle$

### Details of numerical calculations

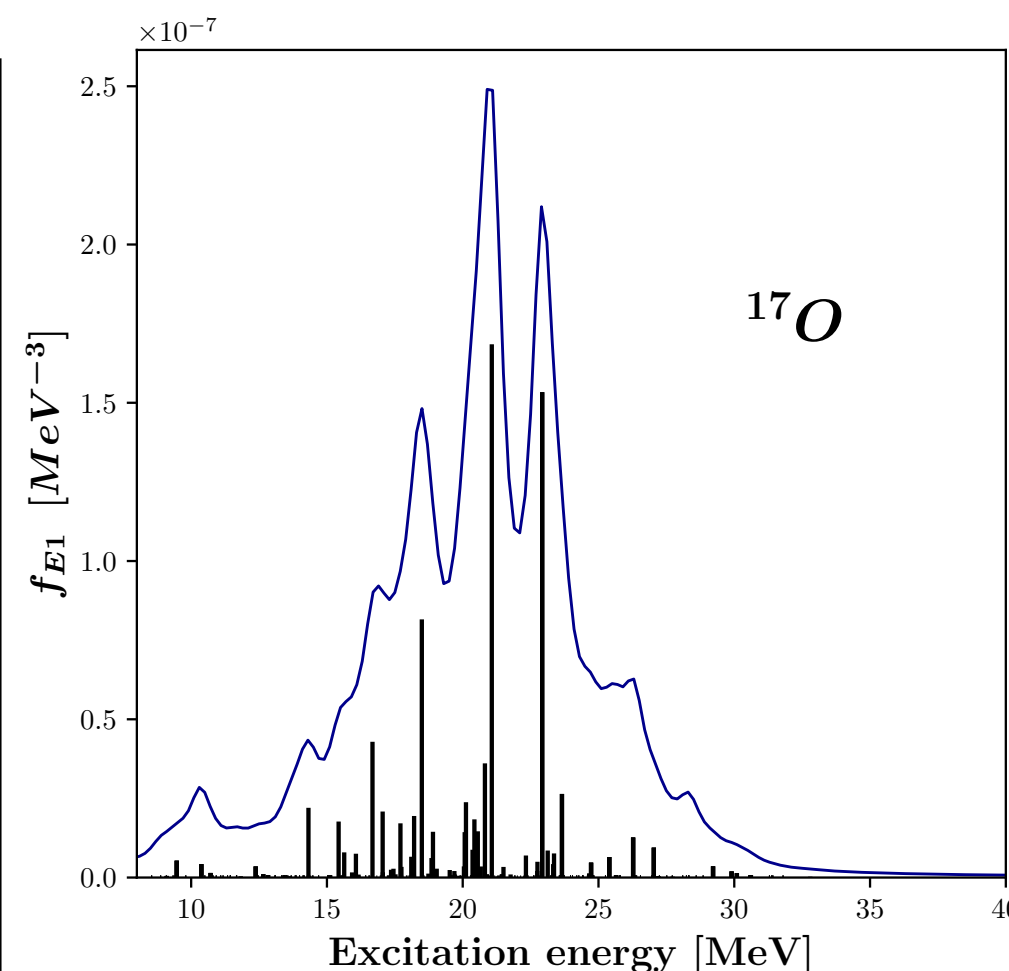
300 Lanczos iterations

Lawson's method for COM

Effective interactions:

- E1 sd-shell: PSDPF

- E1 p-shell: WBP



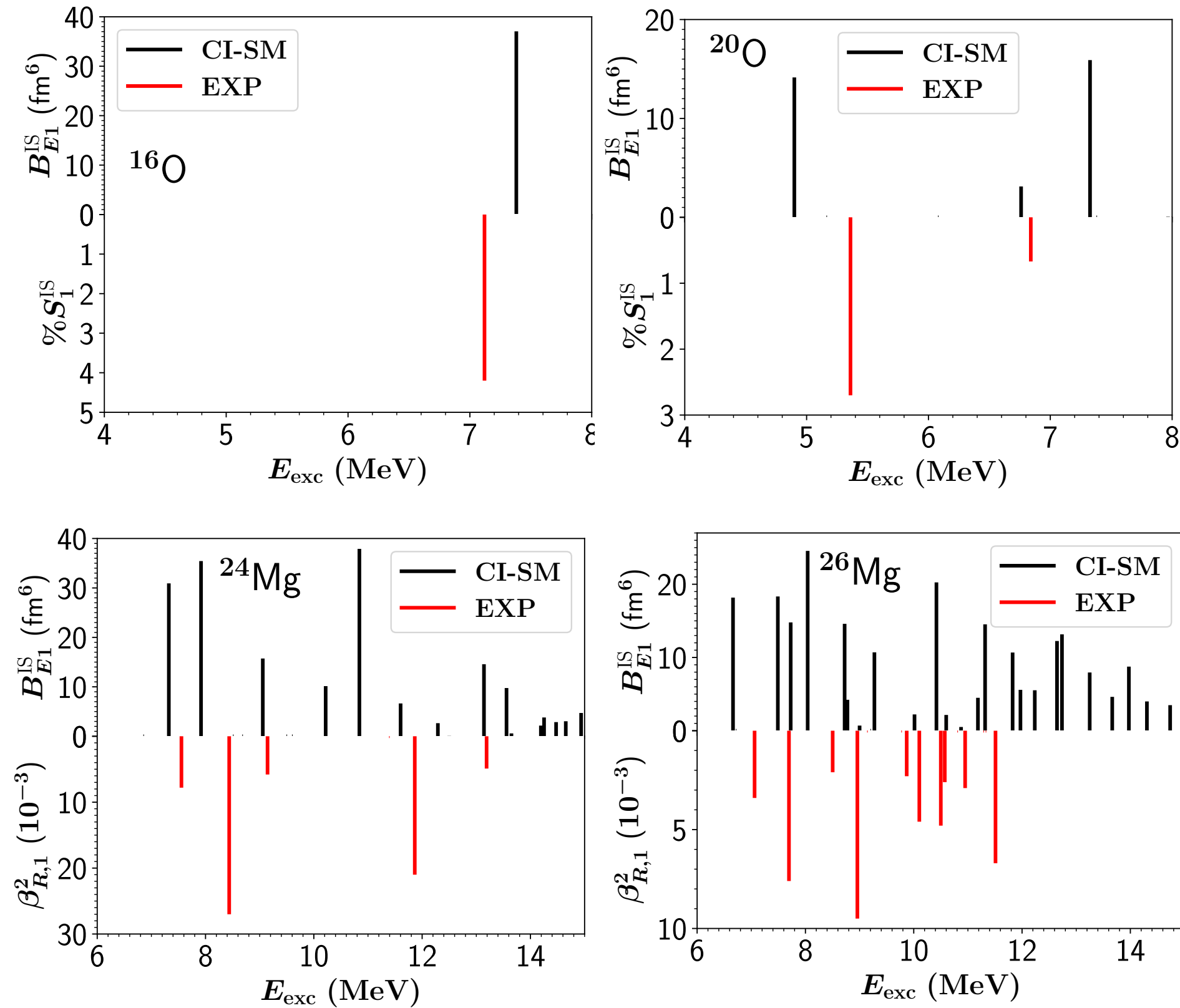
Theorem: for  $N_{\text{it}}$  Lanczos iterations the first  $2N_{\text{it}}$  moments  $S_k$  of the distribution are exact

$$S_k = \sum_j E_j^k |\langle \psi_j | O | \text{GS} \rangle|^2 = \langle \text{SR} | H^k | \text{SR} \rangle$$

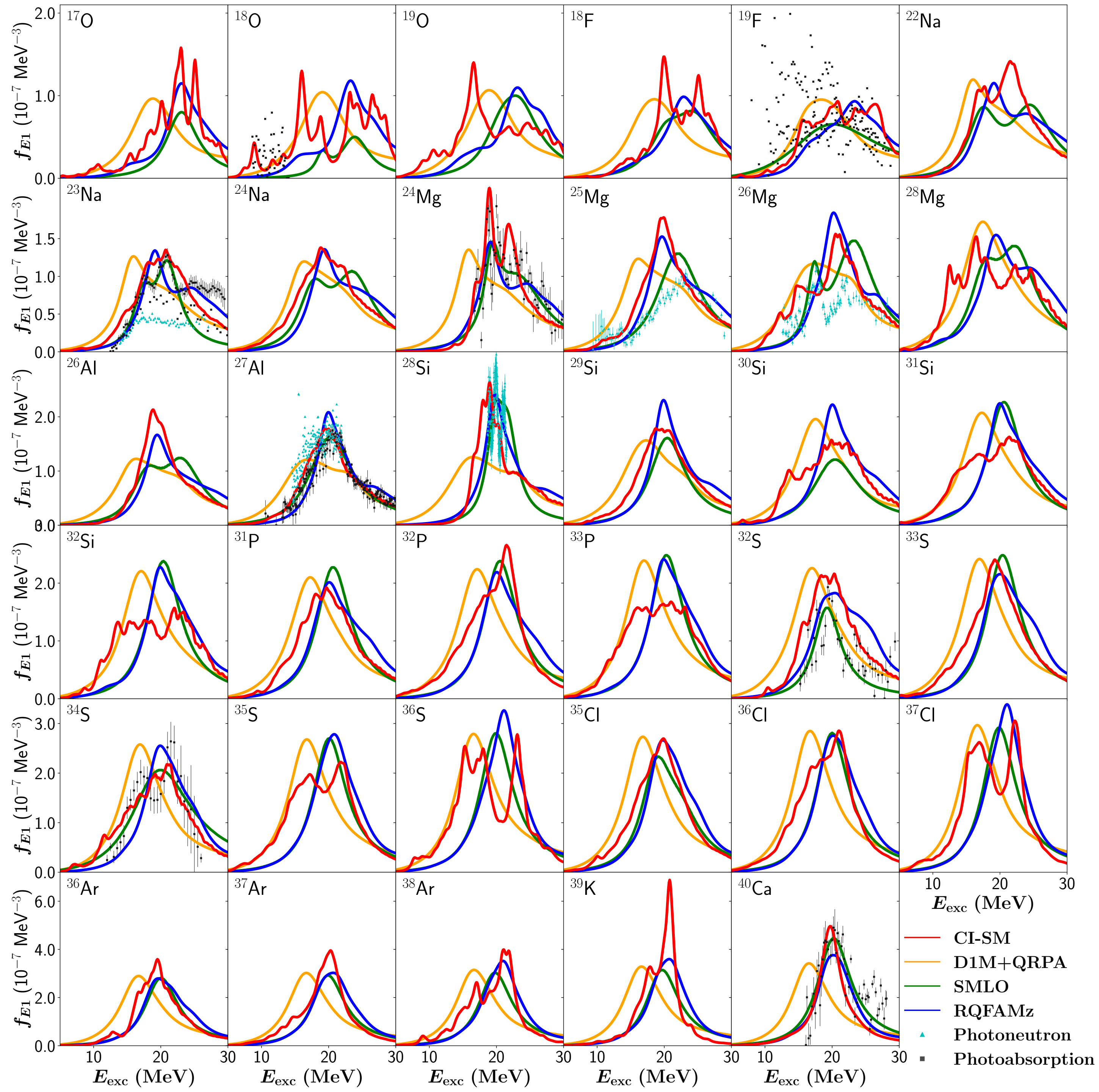
# E1 response sd-shell nuclei

132 nuclei (all isotopic chains within the valence space)

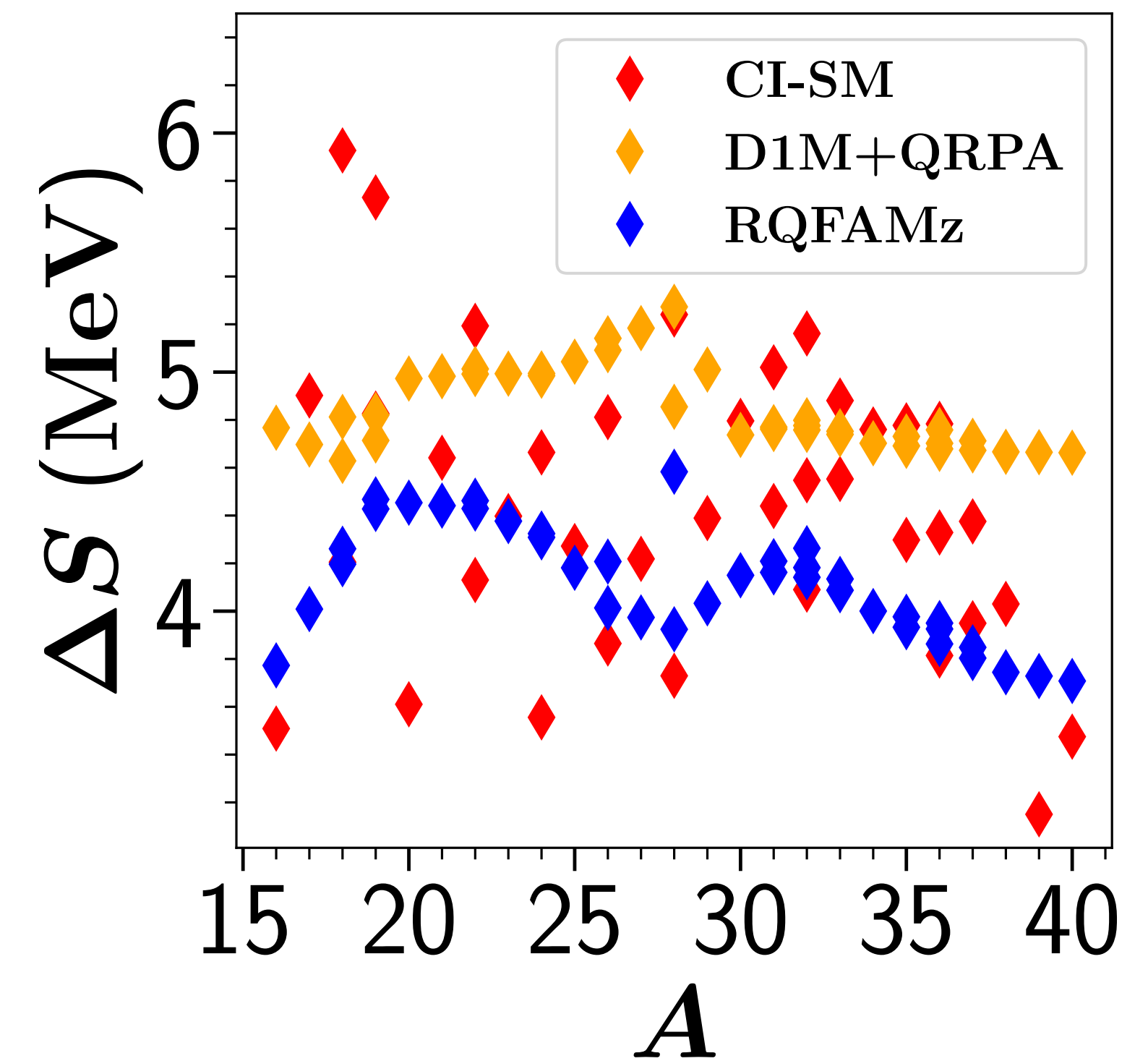
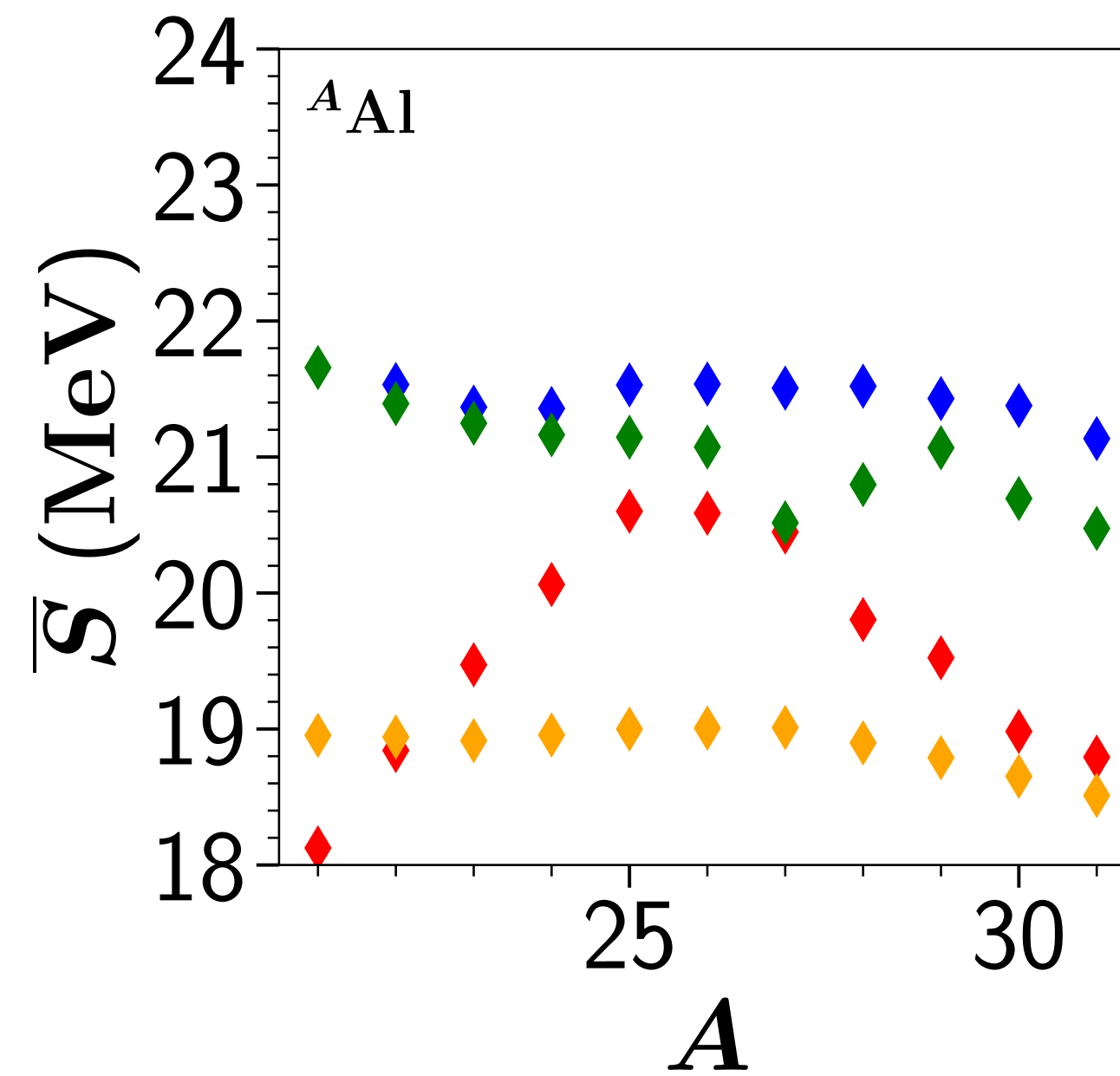
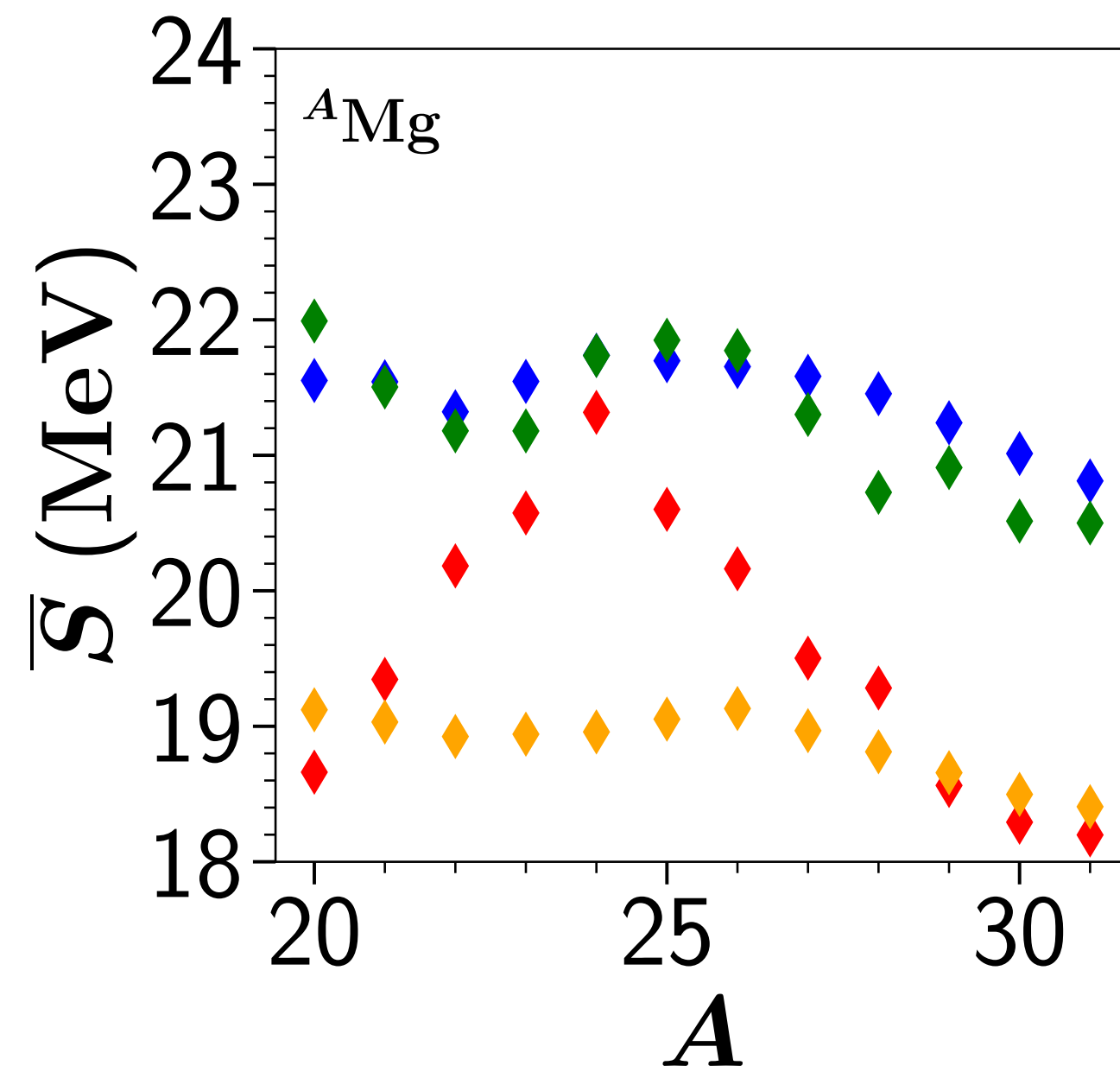
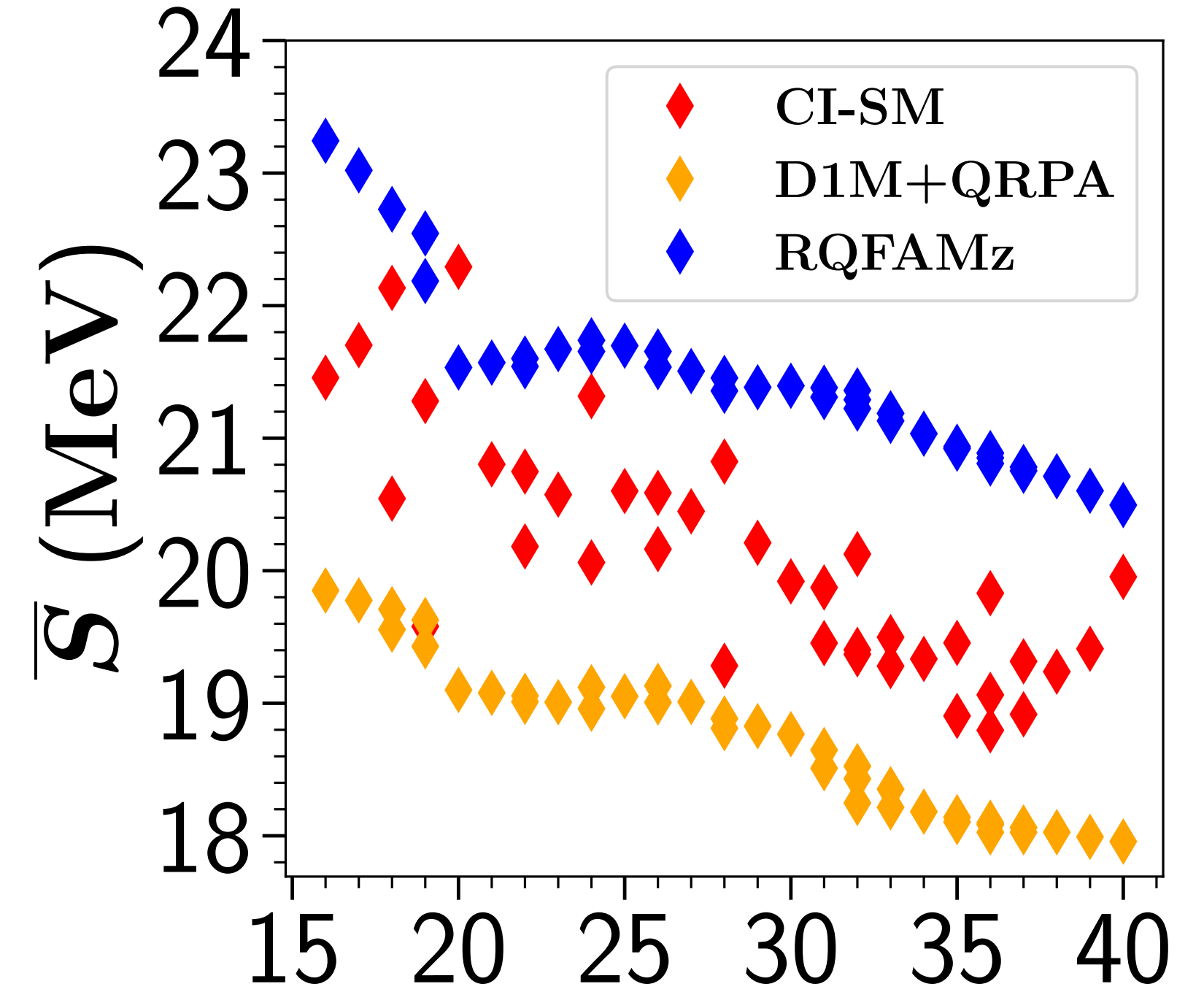
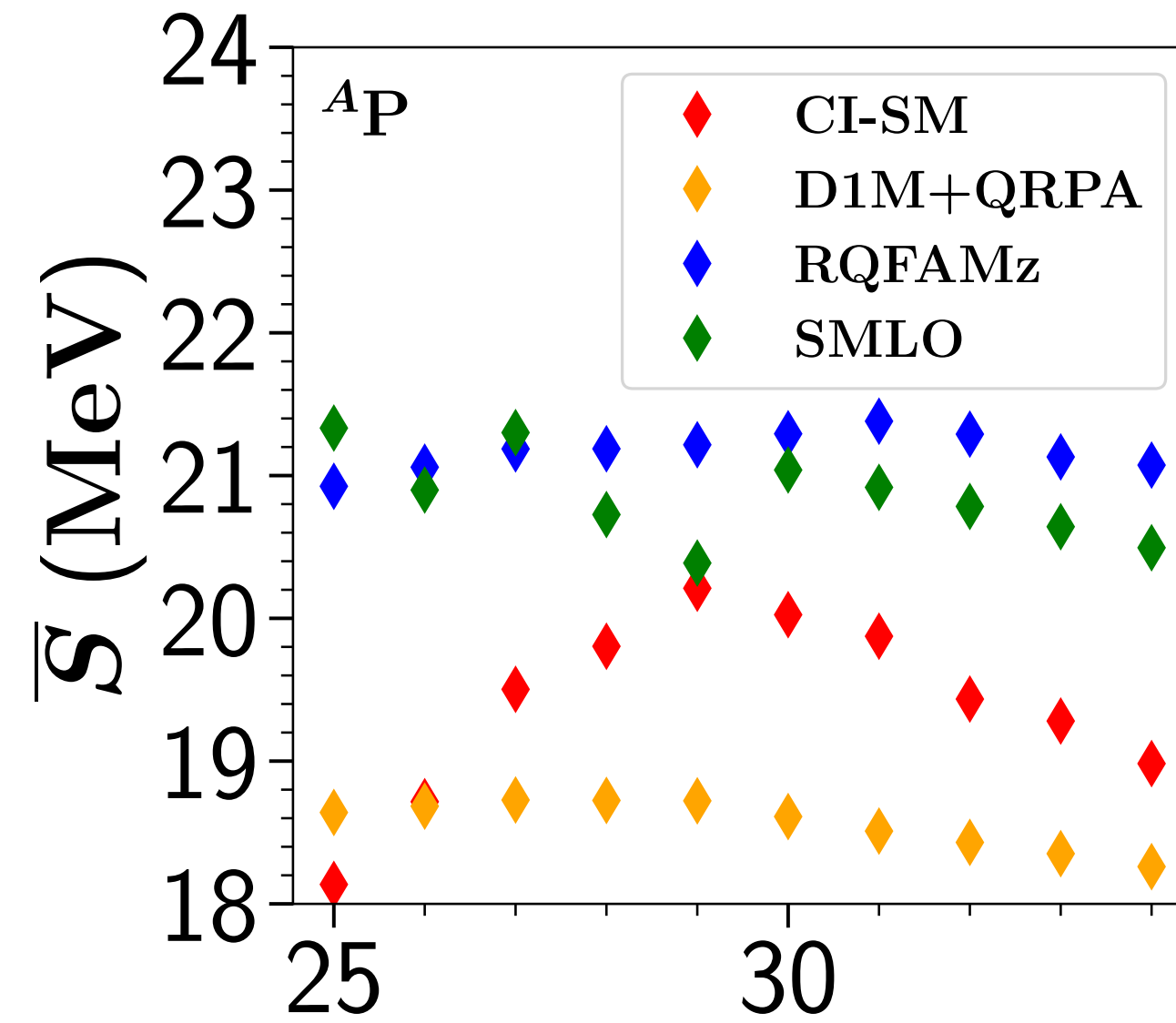
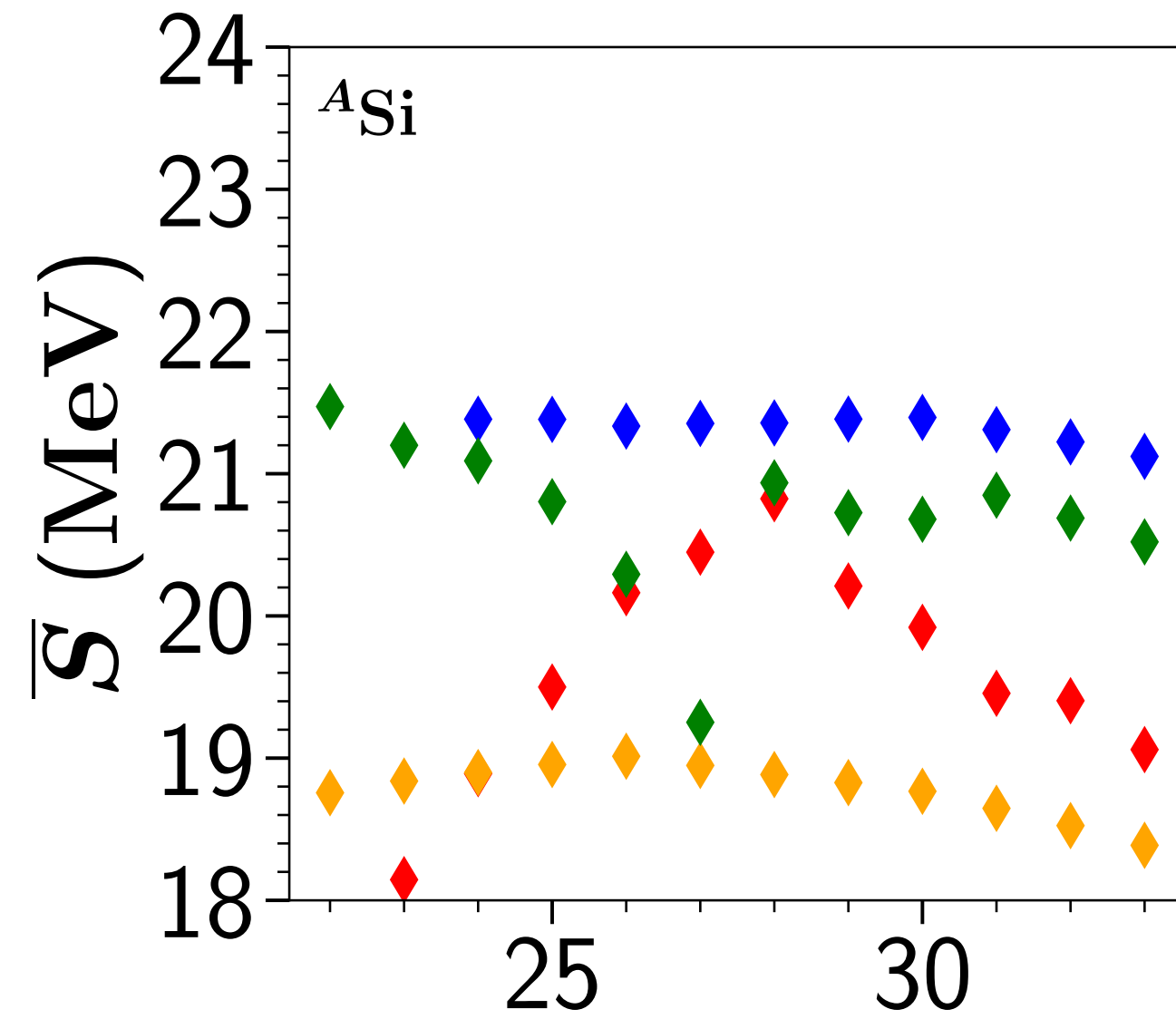
E1 PSF available in the Talys database, M1 to come.



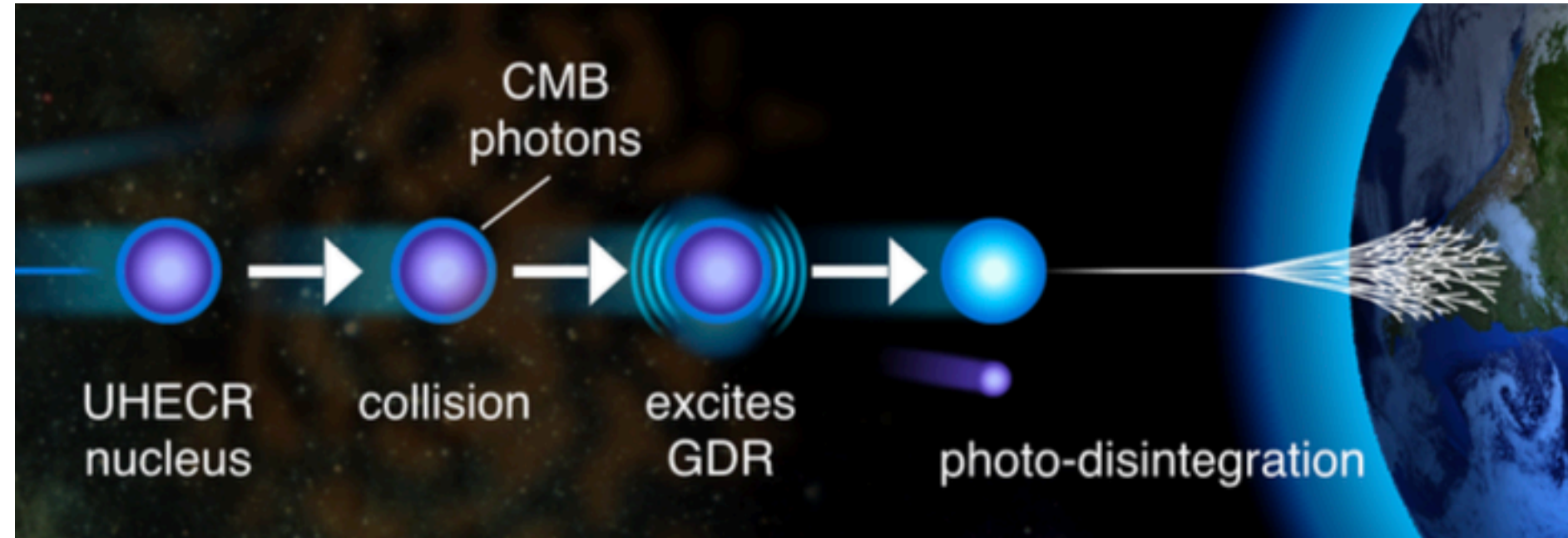
RMS  $\bar{S}$  0.84 MeV  
 $\Delta S$  0.56 MeV



# E1 sd-shell nuclei systematics



# Application to UHECR propagation

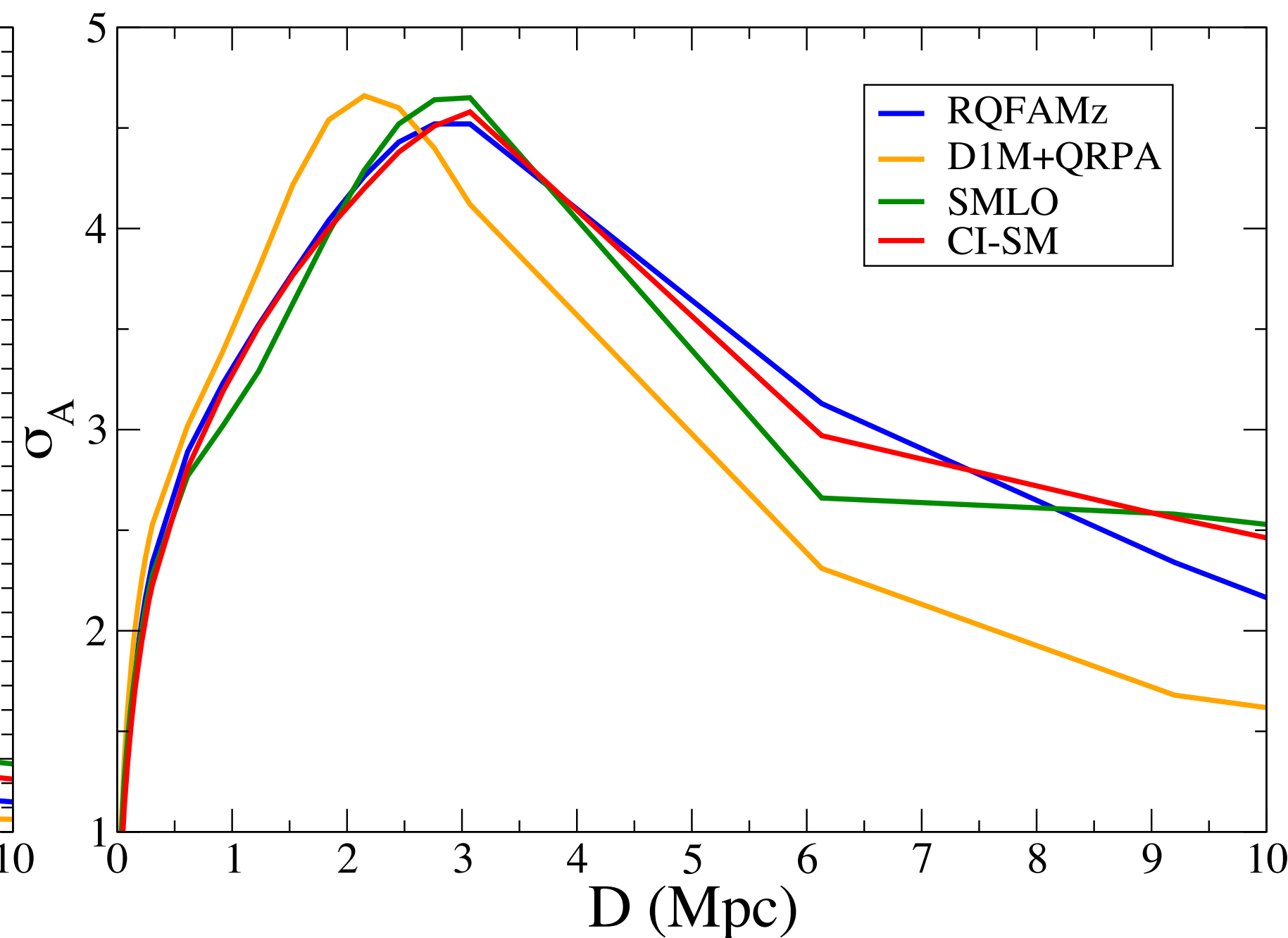
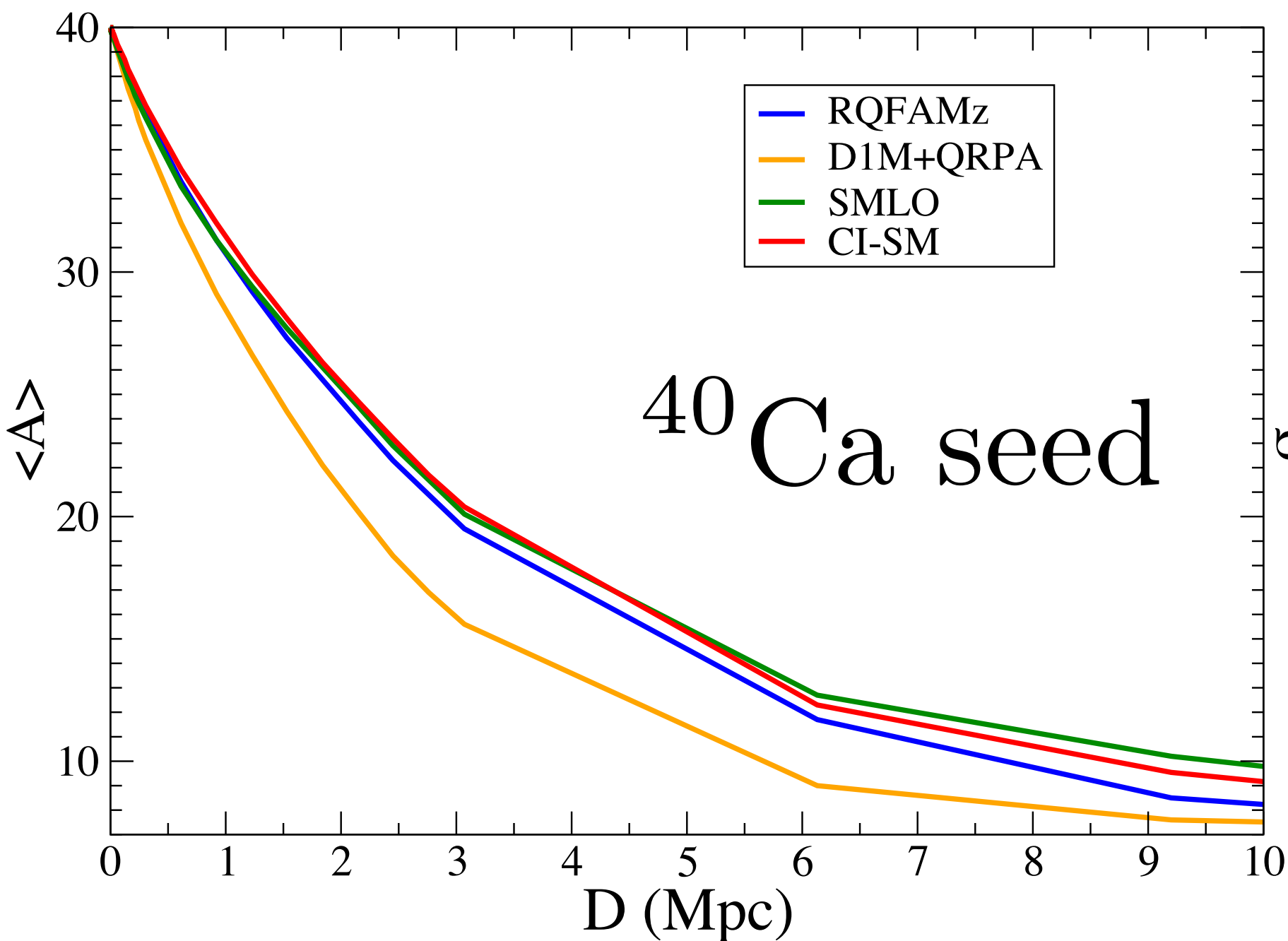


$$\begin{aligned} \frac{dN_{Z,A}}{dt} = & N_{Z+1,A} \lambda_{\beta}^{Z+1,A} + N_{Z-1,A} \lambda_{\beta}^{Z-1,A} \\ & + N_{Z,A+1} \lambda_{\gamma,n}^{Z,A+1} + N_{Z+1,A+1} \lambda_{\gamma,p}^{Z+1,A+1} \\ & + N_{Z+2,A+4} \lambda_{\gamma,\alpha}^{Z+2,A+4} + N_{Z,A+2} \lambda_{\gamma,2n}^{Z,A+2} \\ & + N_{Z+2,A+2} \lambda_{\gamma,2p}^{Z+2,A+2} + N_{Z+4,A+8} \lambda_{\gamma,2\alpha}^{Z+4,A+8} \\ & + N_{Z+1,A+2} \lambda_{\gamma,np}^{Z+1,A+2} + N_{Z+2,A+5} \lambda_{\gamma,n\alpha}^{Z+2,A+5} \\ & + N_{Z+3,A+5} \lambda_{\gamma,p\alpha}^{Z+3,A+5} \\ & - N_{Z,A} \left[ \lambda_{\beta}^{Z,A} + \sum_x \lambda_{\gamma,x}^{Z,A} \right], \end{aligned}$$

$$\lambda_{\gamma,x} = \int n(\epsilon) \sigma_{\gamma,x}(\epsilon) c d\epsilon$$

TALYS

CI-SM  
PSFs

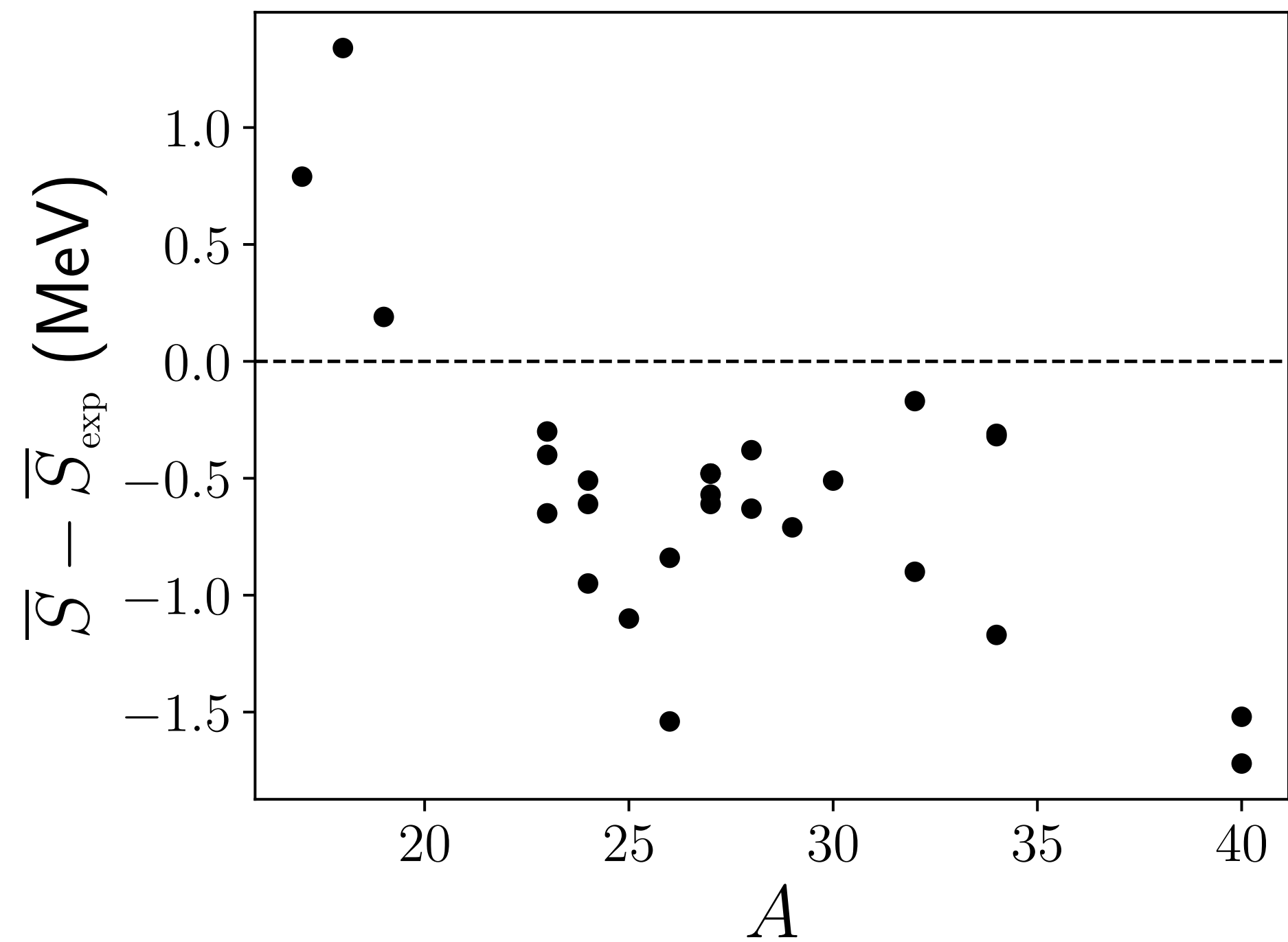


O. Le Noan, S. Goriely, E. Khan, and  
K. Sieja, Physical Review C 113,

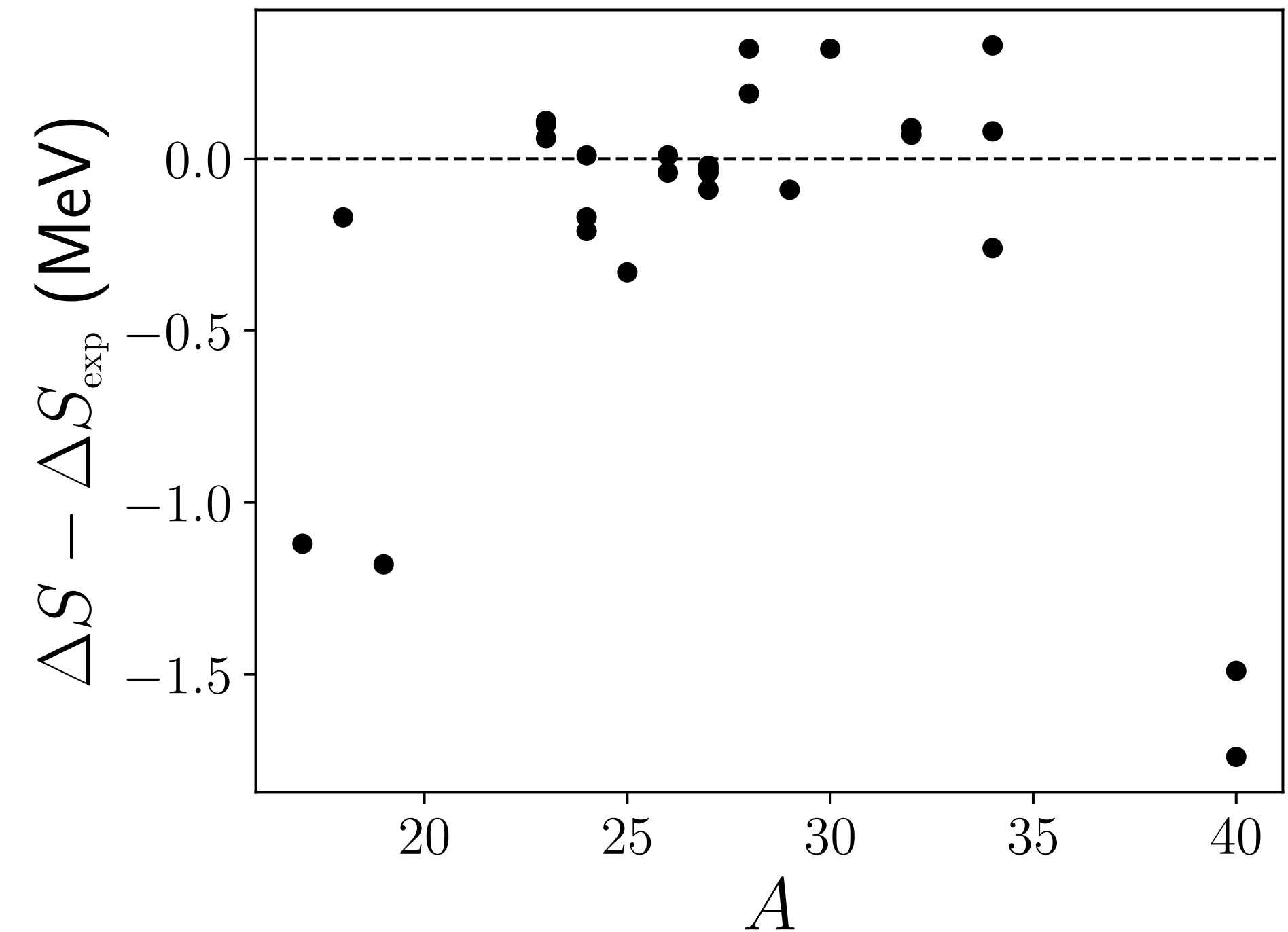
# E1 sd-shell nuclei systematics

Models	$\varepsilon_S$	$\sigma_S$	$\varepsilon_{\Delta S}$	$\sigma_{\Delta S}$
CI-SM	-0.36	0.85	0.27	0.51
D1M + QRPA	-1.47	1.71	0.13	0.49
RQFAMz	0.54	0.89	-0.33	0.48

Centroids  $\bar{S} = \frac{S_1}{S_0}$



Width  $\Delta S = \sqrt{\frac{S_2}{S_0} - \bar{S}^2}$



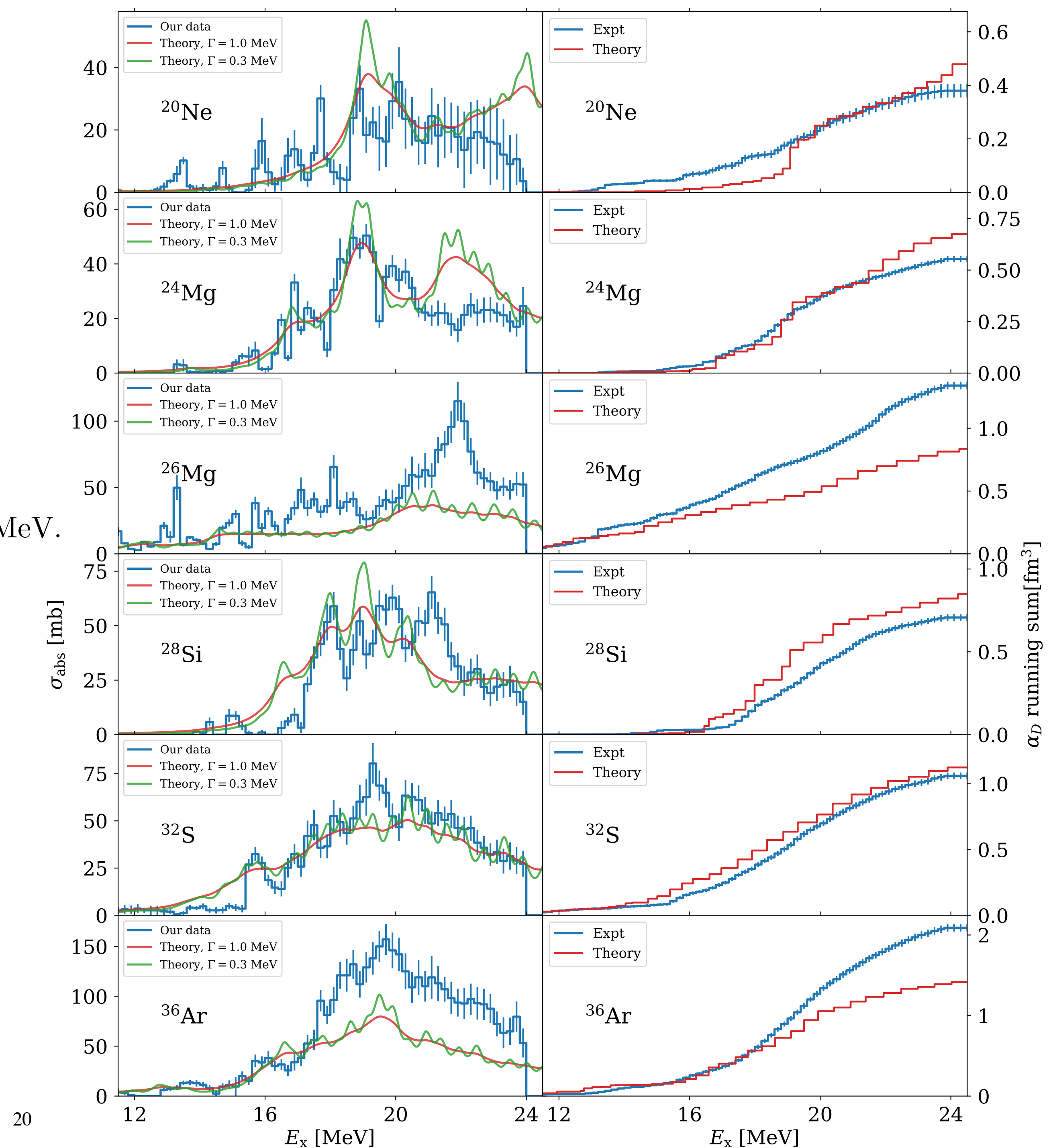
# E1 response sd-shell nuclei

## Electric dipole strength in sd-shell nuclei from small-angle proton scattering

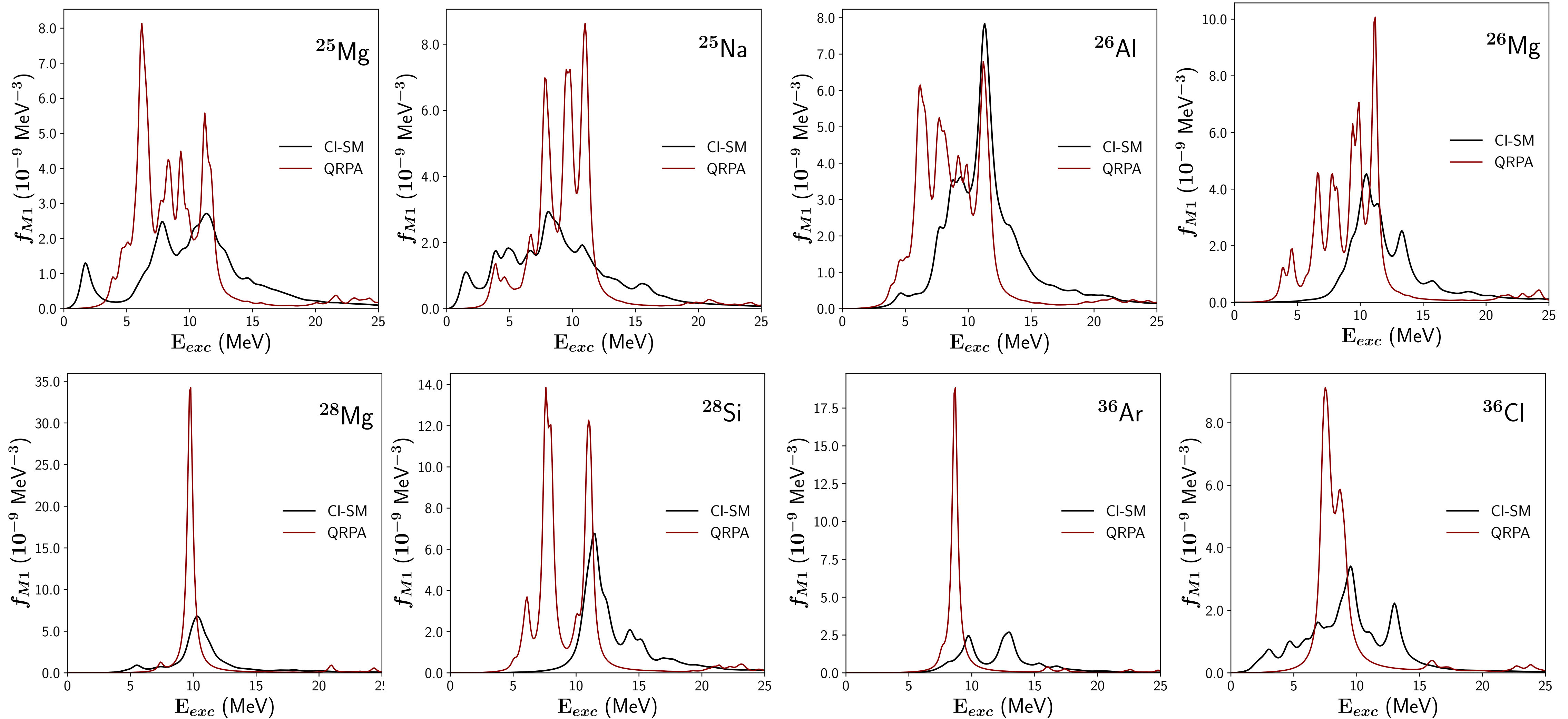
R. W. Fearick, O. Le Noan, H. Matsubara, P. von Neumann-Cosel, K. Sieja, and A. Tamii. To be published soon.

Table 1: Centroids and widths (in MeV), calculated between 12 and 20 MeV.

	$\bar{S}$	$\bar{S}_{\text{exp}}$	$\bar{S} - \bar{S}_{\text{exp}}$	$\Delta S$	$\Delta S_{\text{exp}}$	$\Delta S - \Delta S_{\text{exp}}$
$^{20}\text{Ne}$	18.40	17.50	0.90	1.40	1.99	-0.59
$^{24}\text{Mg}$	18.06	18.15	-0.09	1.43	1.34	0.09
$^{26}\text{Mg}$	16.49	16.65	-0.16	2.22	2.15	0.06
$^{28}\text{Si}$	17.95	18.41	-0.45	1.38	1.22	0.23
$^{32}\text{S}$	17.23	17.82	-0.59	1.86	1.65	0.21
$^{36}\text{Ar}$	17.49	18.01	-0.52	1.87	1.64	0.23



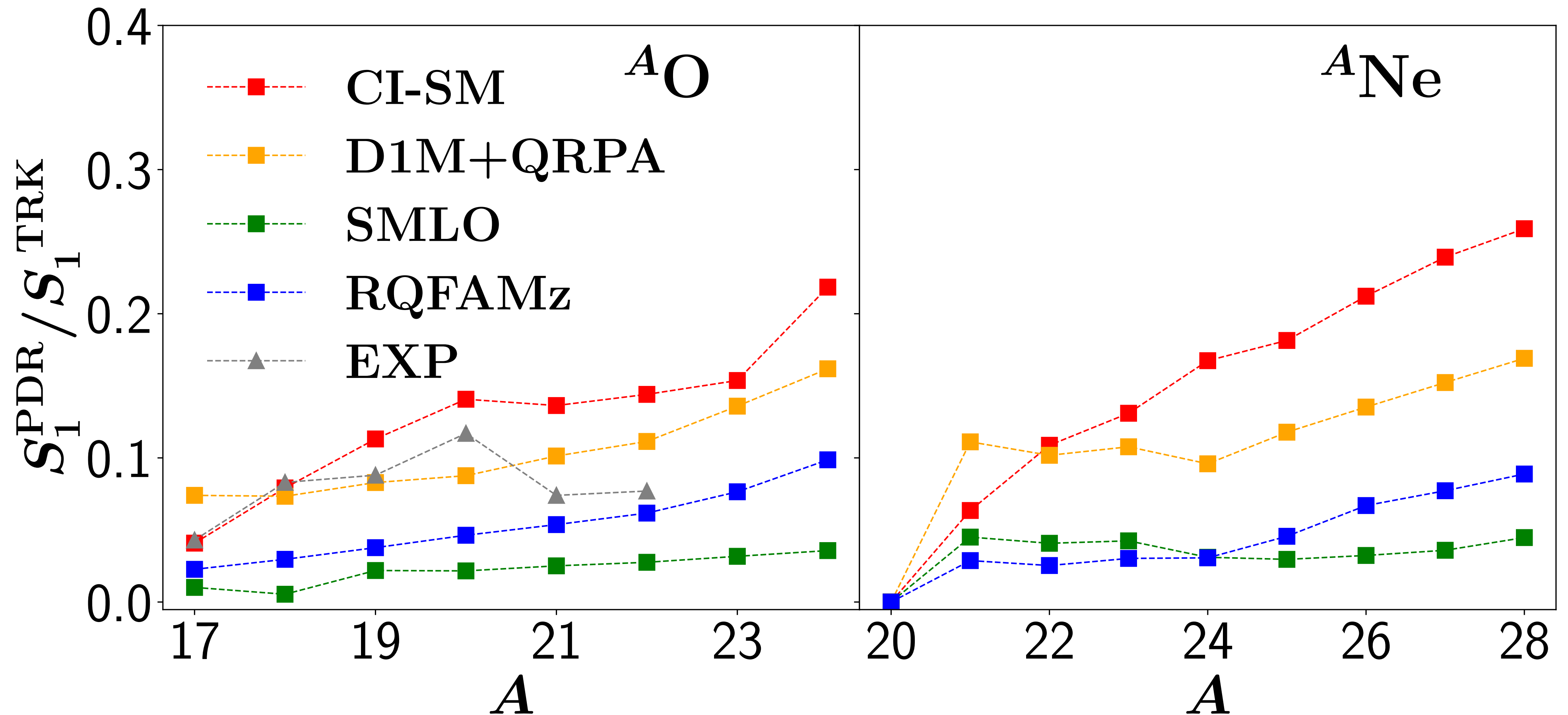
# M1 p, and sd-shell nuclei



# E1 response sd-shell nuclei

Neutron rich nuclei: model comparison

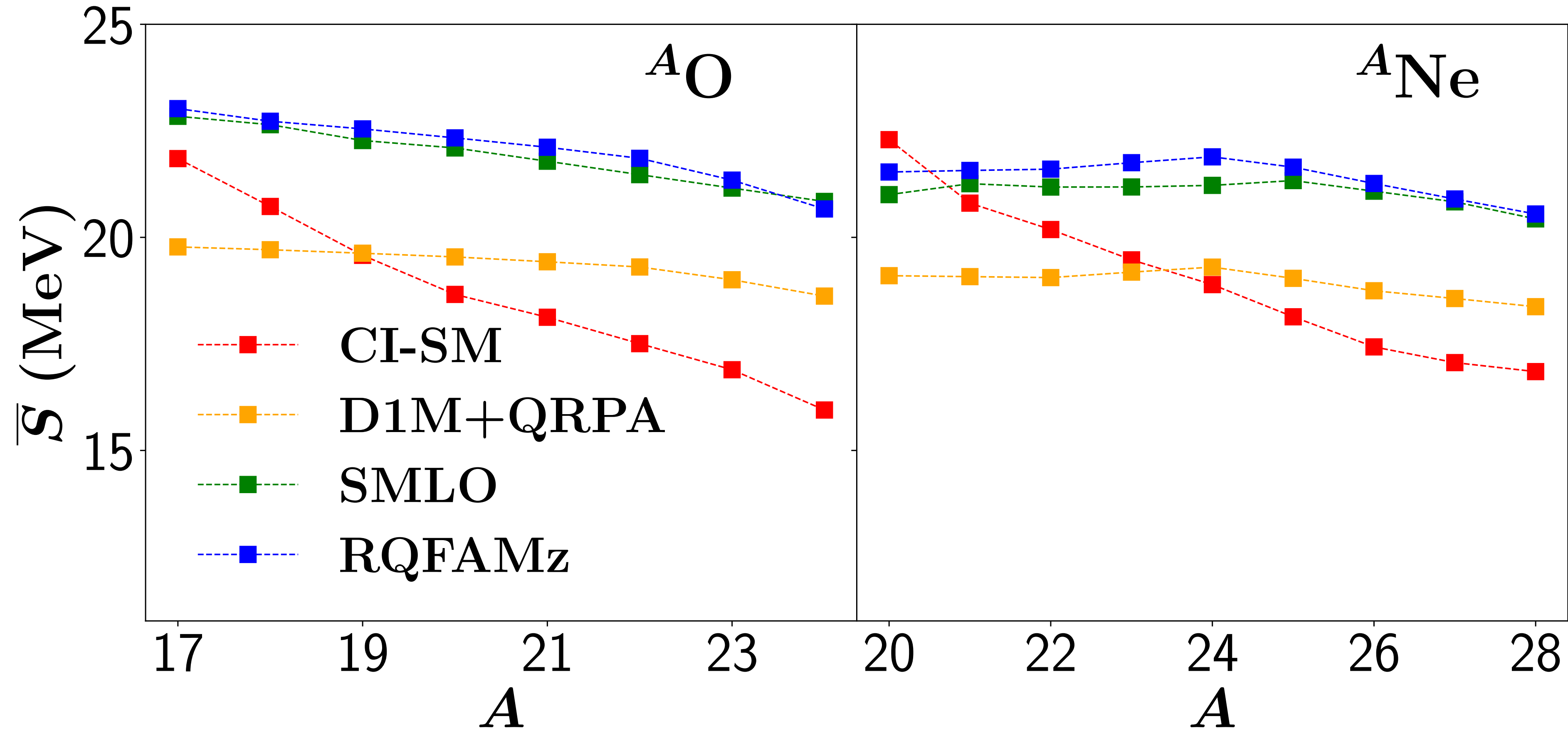
up to 30 MeV



Leistenschneider  
, A. et al.  
Physical Review  
Letters 2001, 86.

22 from  $S_n$  to 15 MeV

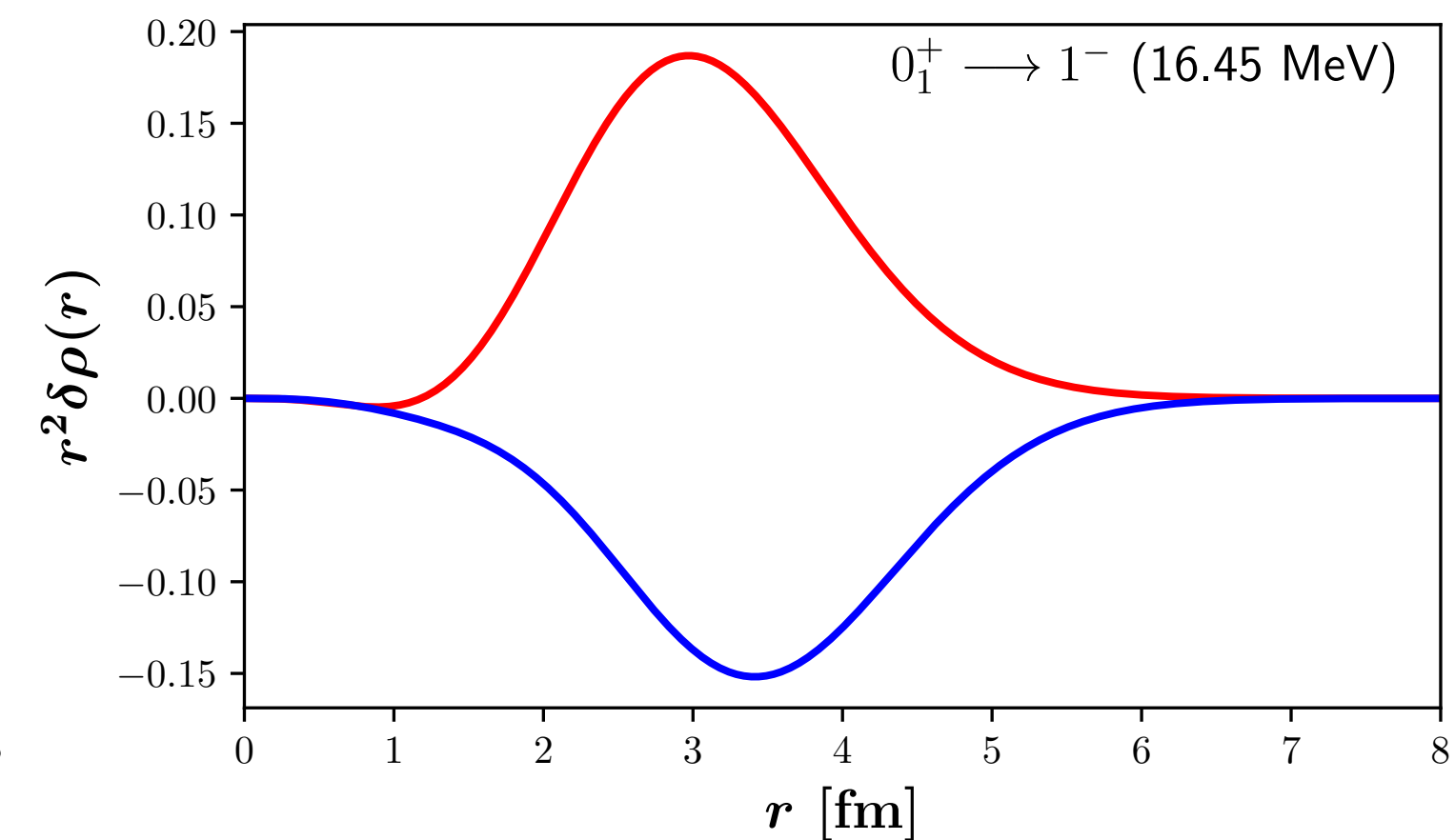
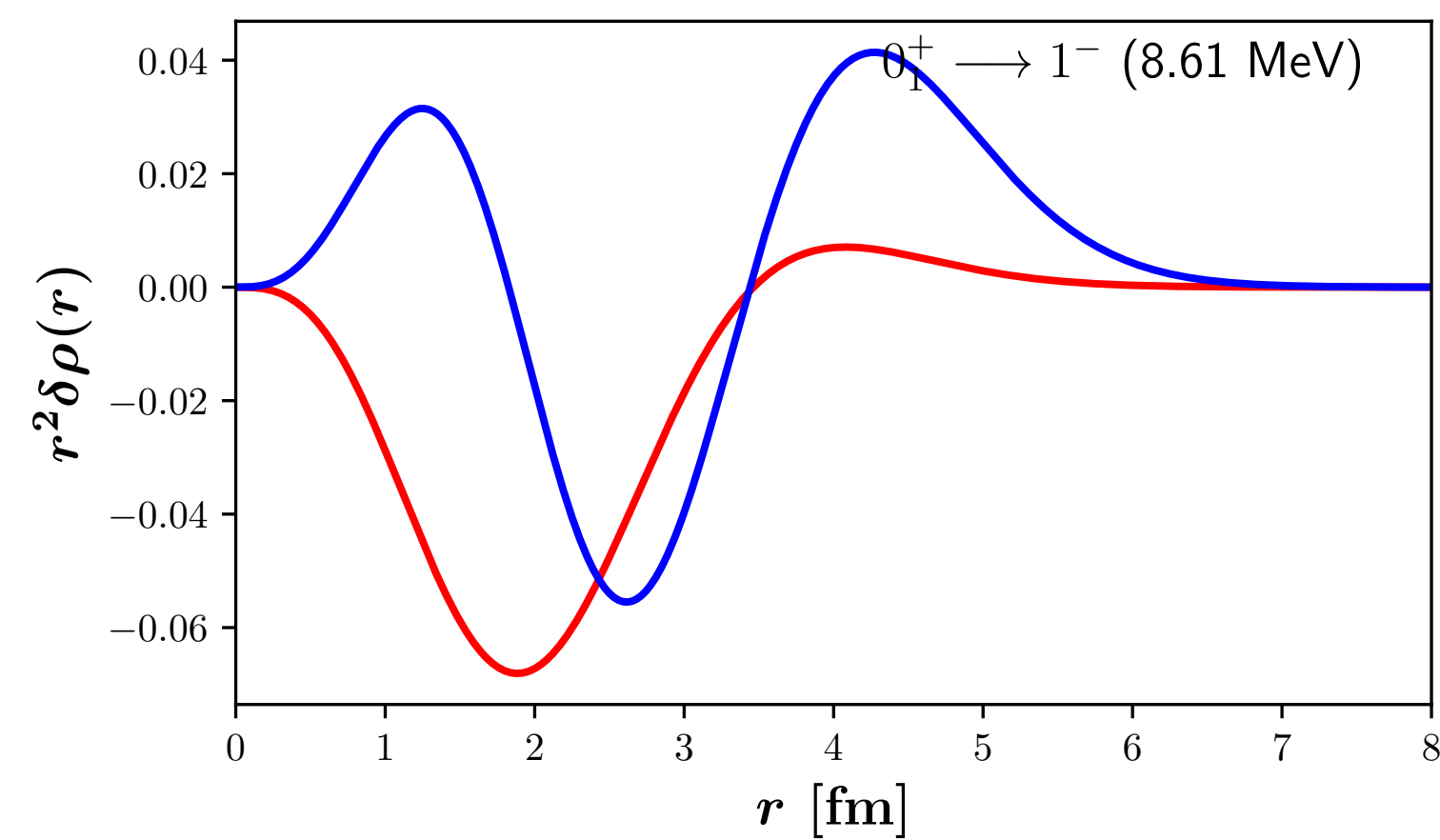
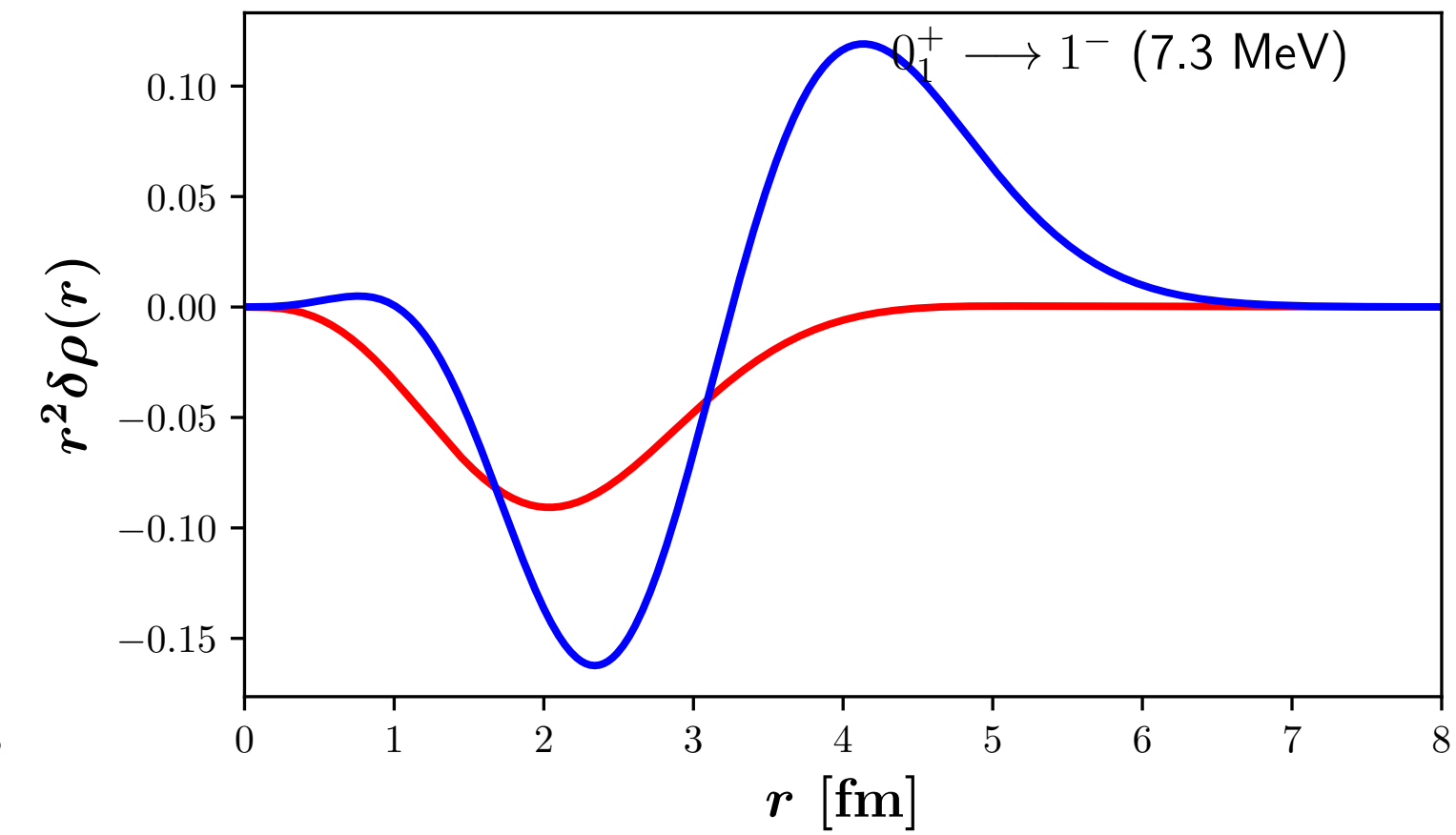
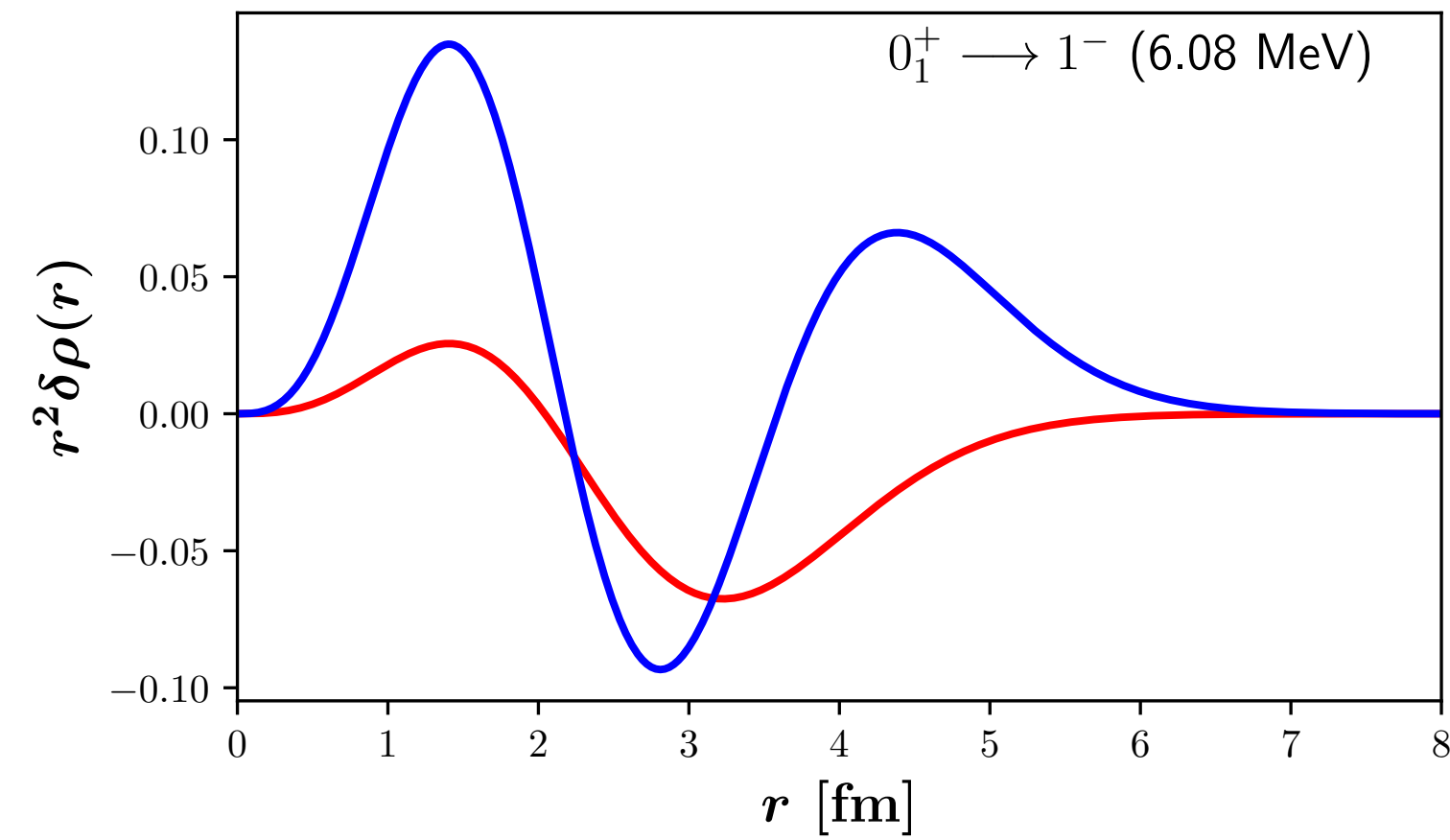
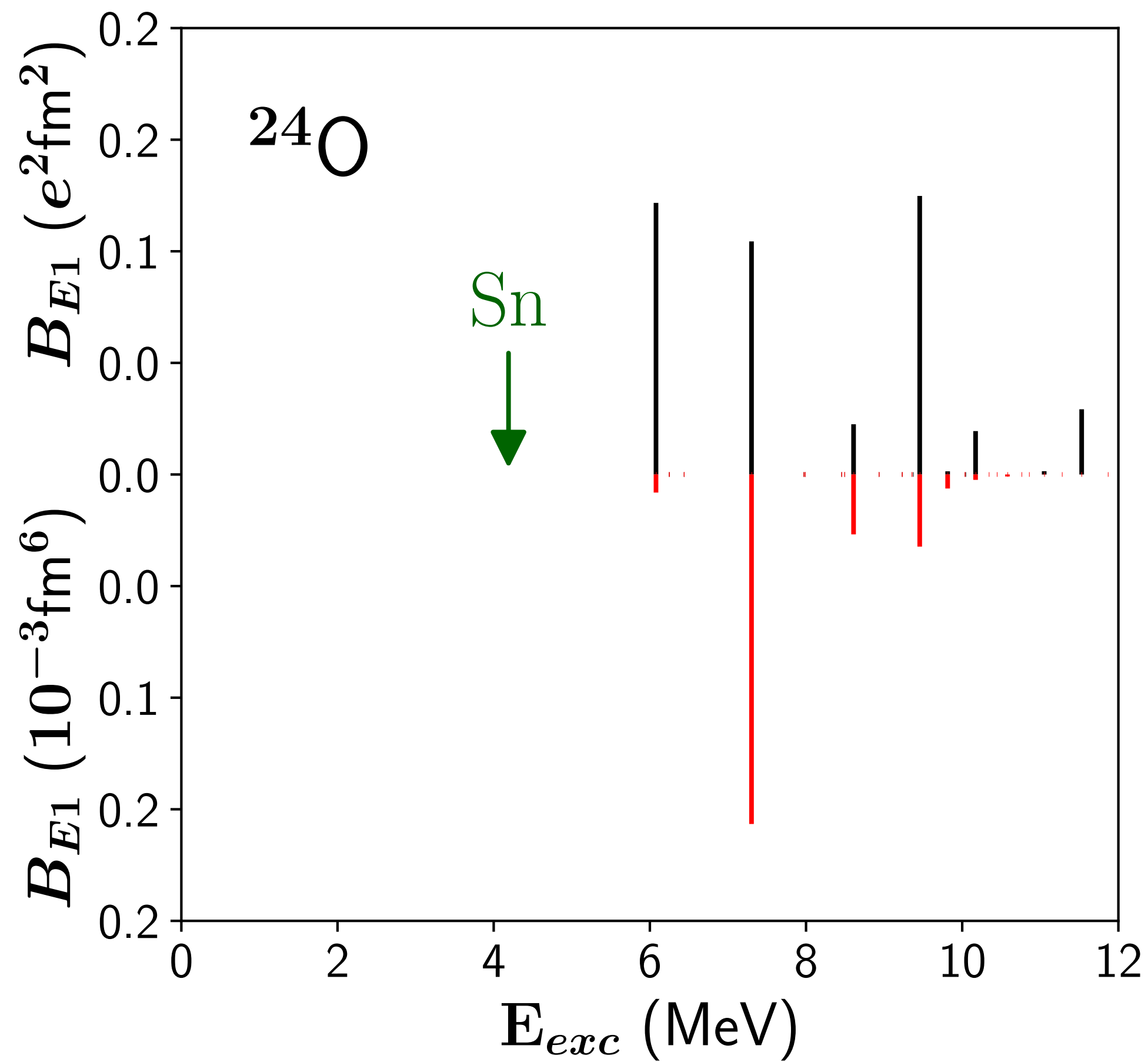
# Study of the Pygmy dipole resonance



up to 30 MeV

# Study of the Pygmy dipole resonance $^{24}\text{O}$

Oxygen chain: model comparison



# Study of the Pygmy dipole resonance

## Isospin mixing in the neon chain

