
Precision Tests of the MSSM at LHC

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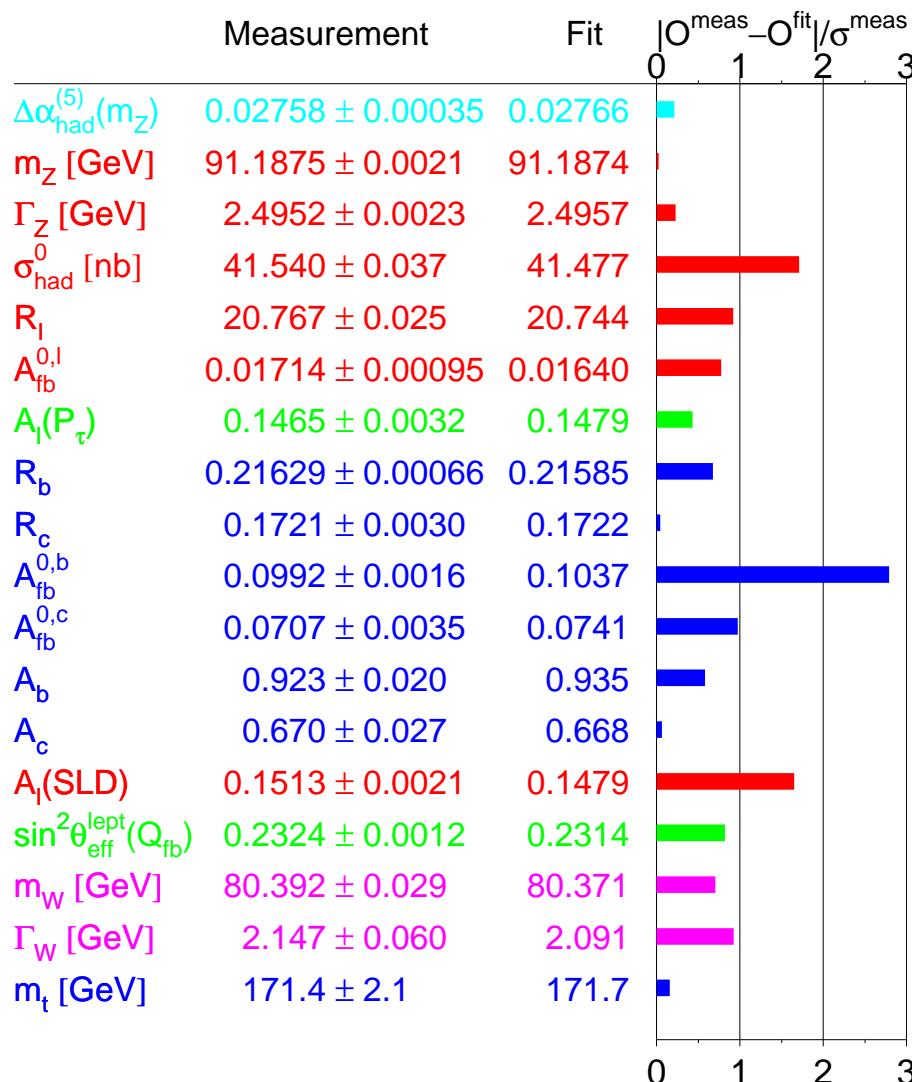
Outline:

- Higgs boson mass in the MSSM
- Coupling constant unification in the MSSM

Standard Model

The SM is amazingly successful:

[LEP and TEVATRON data]



Motivation

The SM has deficiencies. Many open questions:

- ...
- What is the origin of the mass spectrum?
- What is really the Higgs?
- Are the fundamental forces unified?
- What is the dark matter?
- Why is there matter-antimatter asymmetry?
- ...

Possible answers in physics *Beyond the Standard Model*

- Supersymmetry, GUT, Extra Dimensions, String Theory, ...

Supersymmetry

Unification of all forces of Nature \rightsquigarrow include Gravity around $M_{Plank} \simeq 10^{19}$ GeV .

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 - SUSY generators:

$$Q |boson\rangle = |fermion\rangle \quad Q |fermion\rangle = |boson\rangle$$

\Rightarrow a partial unification of matter(fermions) with forces (bosons) naturally achieved.

Supersymmetry

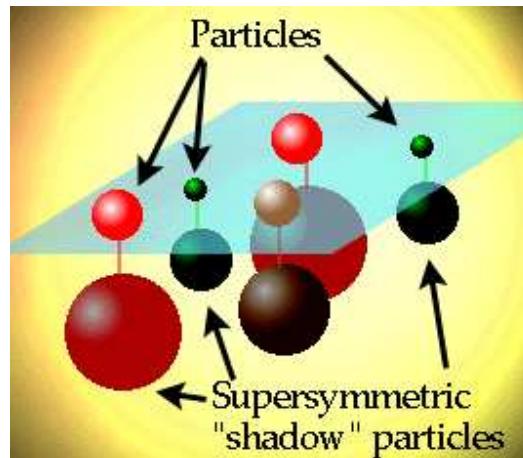
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 - Number of bosonic and fermionic degrees of freedom are equal.
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- Minimal Supersymmetric extension of SM (MSSM)

	Bosons (spin=0)	Bosons (spin=1)	Fermions
Gauge		gluon weak photon	gluino wino, zino photino
Matter	sleptons squarks		leptons quarks
Higgs	Higgses		higgsinos

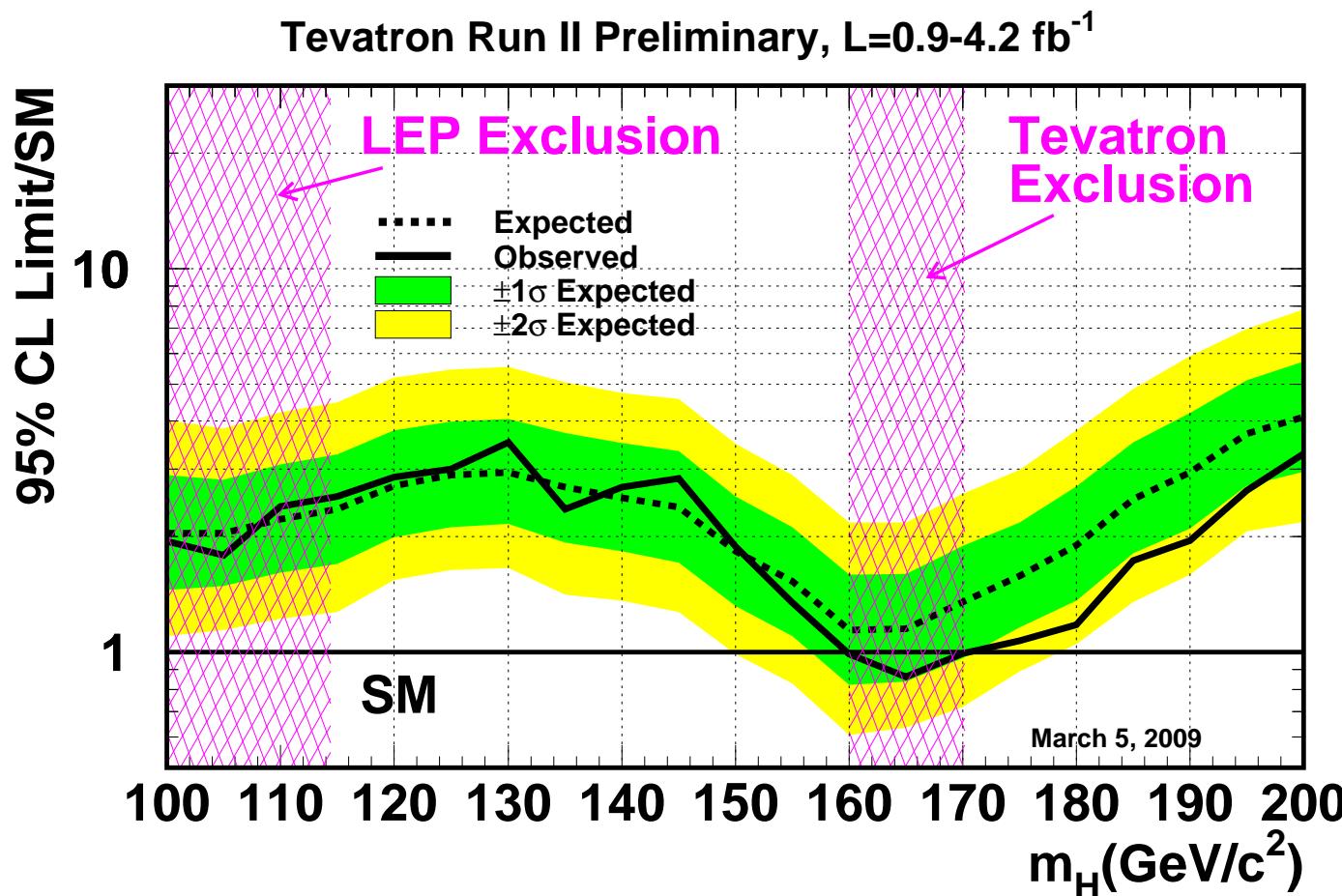
Supersymmetry

- SUSY can naturally address some of the basic questions:
 - the hierarchy problem ($M_W \ll M_{GUT}$) become natural
 - dark matter candidate
 - predicts a light Higgs boson !!
 - predicts gauge coupling unification !!
 - predicts SUSY particles at about 1 TeV
 - ...

Higgs boson mass in the MSSM

SM Higgs Search

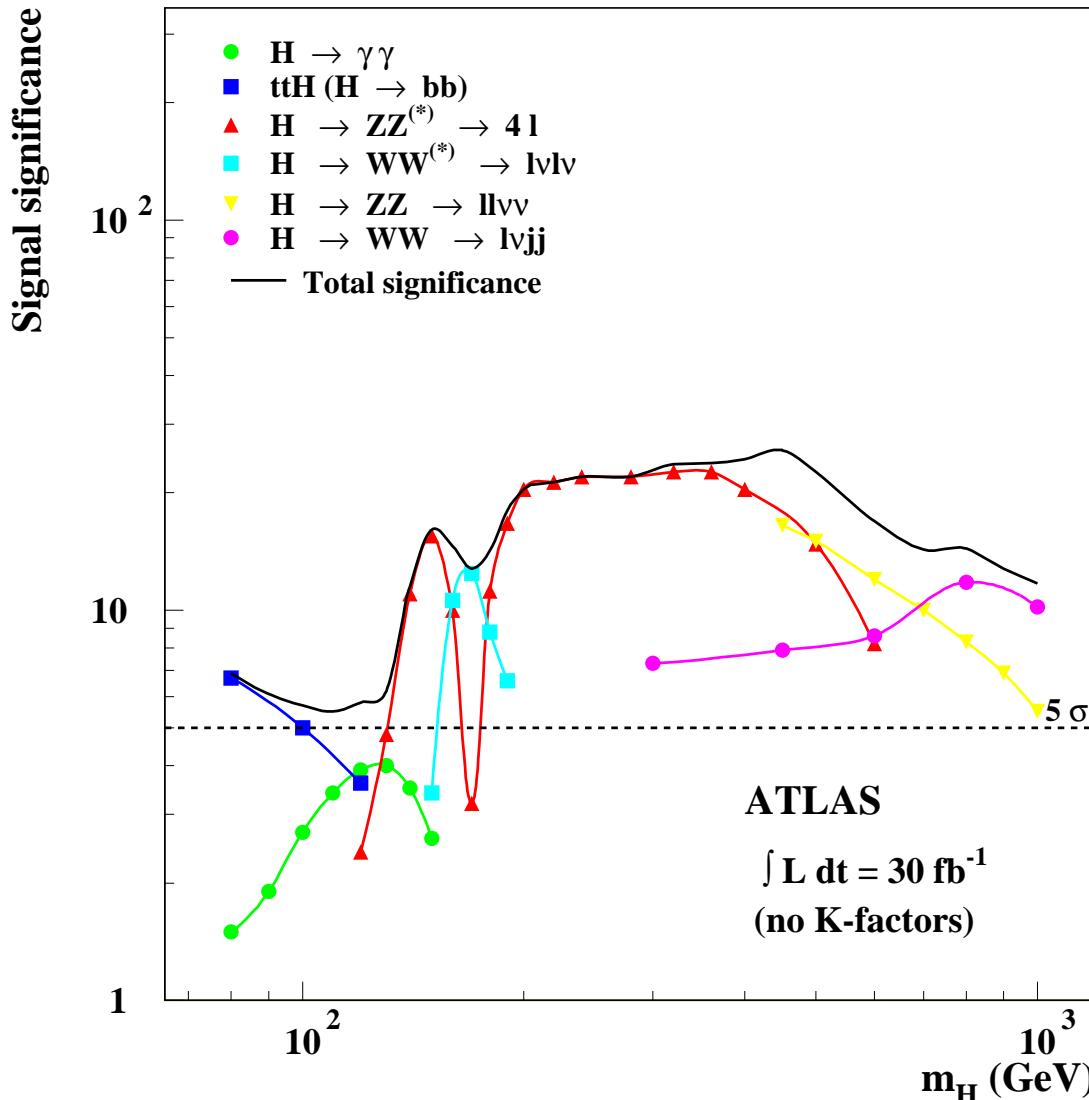
at present: direct searches



electroweak precision data: $m_H = 87^{+36}_{-27} \text{ GeV}$

SM Higgs Search

expected soon at the LHC



MSSM Higgs

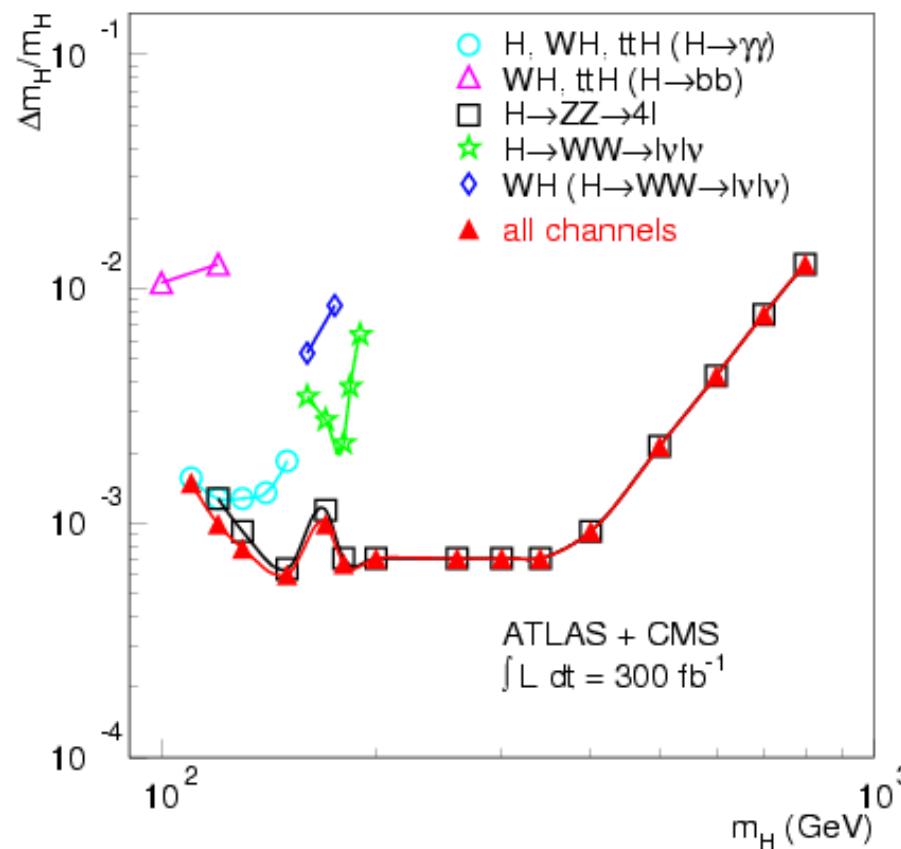
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- Higgs mass is very sensitive to top/stop sector through radiative corrections
 \Rightarrow the mass of the light Higgs used as a consistency check of the MSSM
- MSSM electroweak precision data + heavy flavour + dark matter:

$$m_h = 110^{+10}_{-8} \pm 3(\text{th}) \text{ GeV}$$

Theory vs. experiment

Experiment LHC: $\delta m_h^{\text{exp}} = 100 - 200 \text{ MeV}$ [CERN-LHCC-2006-21](#)
 ILC: $\delta m_h^{\text{exp}} = 50 \text{ MeV}$ [AguilarSaavedra et al '06](#)

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- 2-loop $\mathcal{O}(\alpha_t \alpha_s, \alpha_t^2, \alpha_b \alpha_s, \alpha_b \alpha_t)$ in effective potential approximation ($p^2 = 0$)
[Haber, Hempfling, Hoang '96], [Heinemeyer, Hollik and Weiglein '98], [Degrassi, Slavich, Zwirner '01], [Espinosa and Zang '00], [Brignole, Degrassi, Slavich, Zwirner '02] , [Carena et al '00], [Heinemeyer et al '05] , [S. Martin '03]
- Momentum-dependent corrections ($p^2 = m_h^2$): 2-loop SUSY-QCD [S.Martin '05]
- 3-loop LL and NLL $\mathcal{O}(\alpha_t \alpha_s^2, \alpha_t^2 \alpha_s, \alpha_t^3)$ [S. Martin '07]

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 - full 2-loop corrections
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Framework

MSSM: SUSY \Rightarrow two free parameters: $\tan \beta = v_2/v_1$, $M_A^2 = -m_{12}^2(\tan \beta + \cot \beta)$

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$$\mathcal{M}_{H,\text{tree}}^2 = \frac{\sin 2\beta}{2} \times \begin{pmatrix} M_Z^2 \cot \beta + M_A^2 \tan \beta & -M_Z^2 - M_A^2 \\ -M_Z^2 - M_A^2 & M_Z^2 \tan \beta + M_A^2 \cot \beta \end{pmatrix}$$

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Higher order corrections

$$\mathcal{M}_H^2 = \mathcal{M}_{H,\text{tree}}^2 - \begin{pmatrix} \hat{\Sigma}_{\phi_1} & \hat{\Sigma}_{\phi_1 \phi_2} \\ \hat{\Sigma}_{\phi_1 \phi_2} & \hat{\Sigma}_{\phi_2} \end{pmatrix}$$

$\hat{\Sigma}_{\phi_i}$ = renormalized self-energies

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$\Sigma_i(0)$ = bare self-energies

δV_i = Higgs potential counterterms

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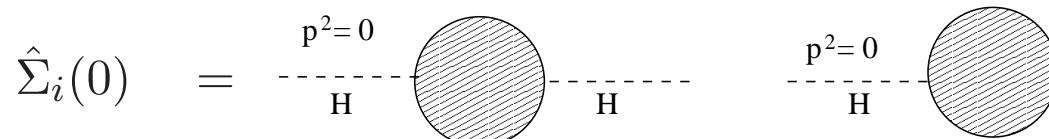
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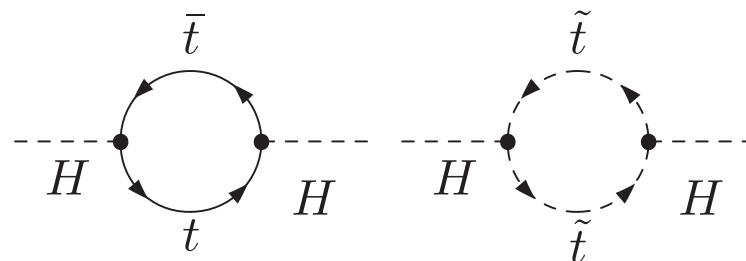
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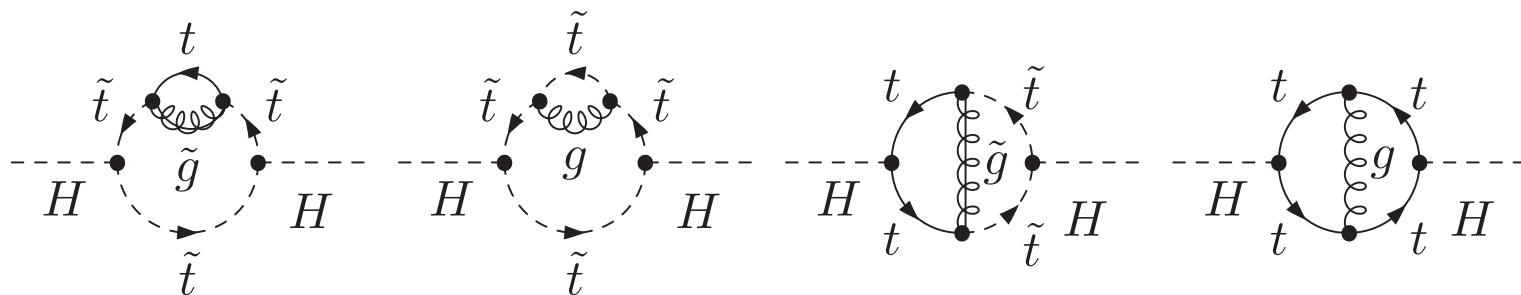
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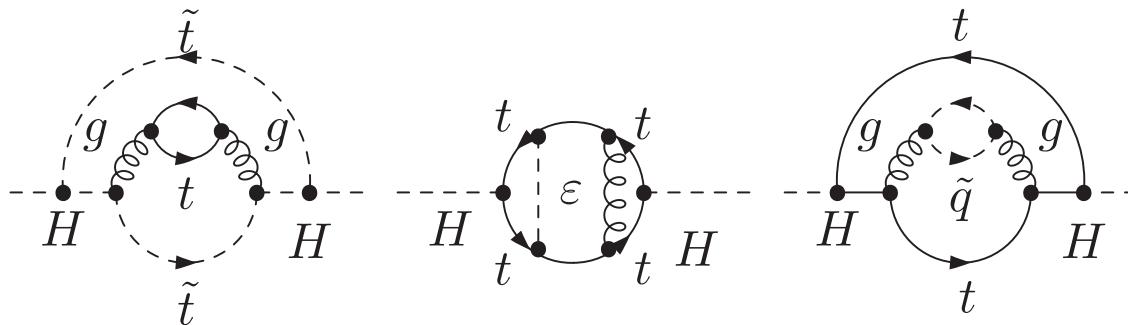
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Computation of $\hat{\Sigma}_{\phi_{ij}}(0)$ at 3-loops:

- $\simeq 28.000$ diagrams
- Computer programs: QGRAF, PERL, FORM, MINCER, MATAD, EXP, ...
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- Asymptotic expansion \rightsquigarrow 3-loop tadpole integrals \rightsquigarrow MATAD

Numerical Results (no stop-mixing)

Input SM parameters: $\mu = M_t = 172.4 \text{ GeV}$ $G_F = 1.16637 \times 10^{-5} \text{ GeV}^{-2}$
 $M_Z = 91.1876 \text{ GeV}$ $\alpha_s^{(5)}(M_Z) = 0.1189 \Rightarrow \alpha_s(M_t) = 0.0926$

MSSM parameters: $M_A = 1 \text{ TeV}$ $\tan \beta = 40$; $A_t = 0$ $M_{\tilde{q}} = 2 \text{ TeV}$
 $M_{\tilde{t}_2} = M_{\tilde{t}_1} = M_{\tilde{g}} = M_{\text{SUSY}}$

OS-scheme $\Delta M_h^{(n)} \equiv M_h^{(n-\text{loop})} - M_h^{\text{tree}}$

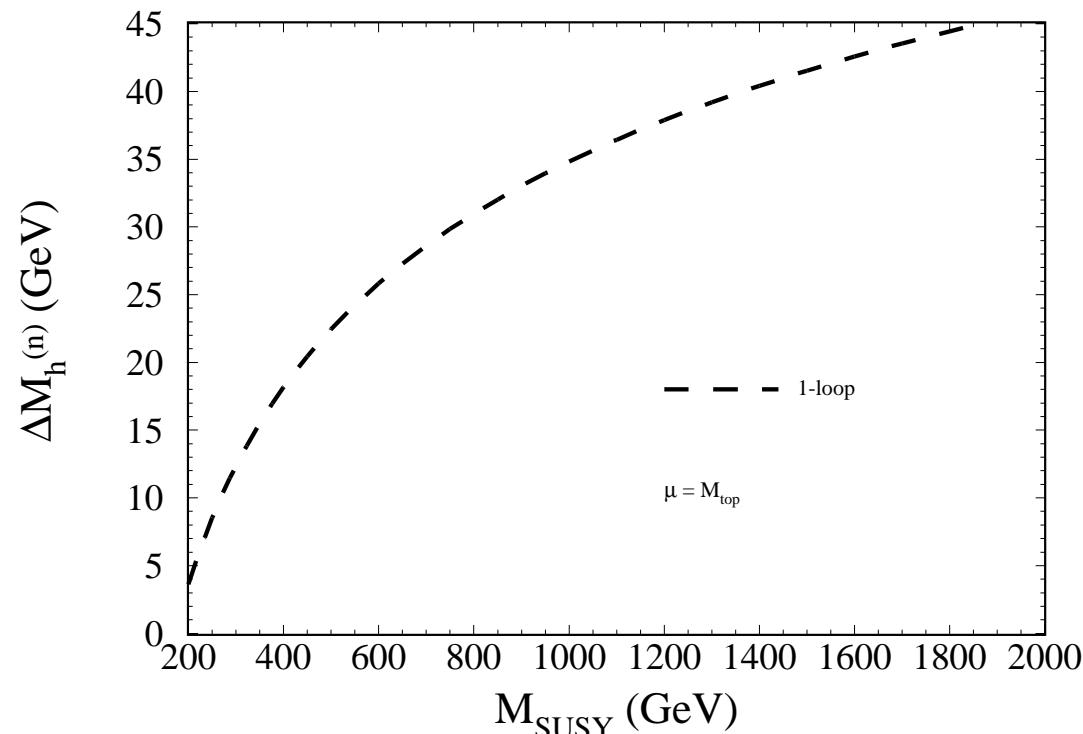
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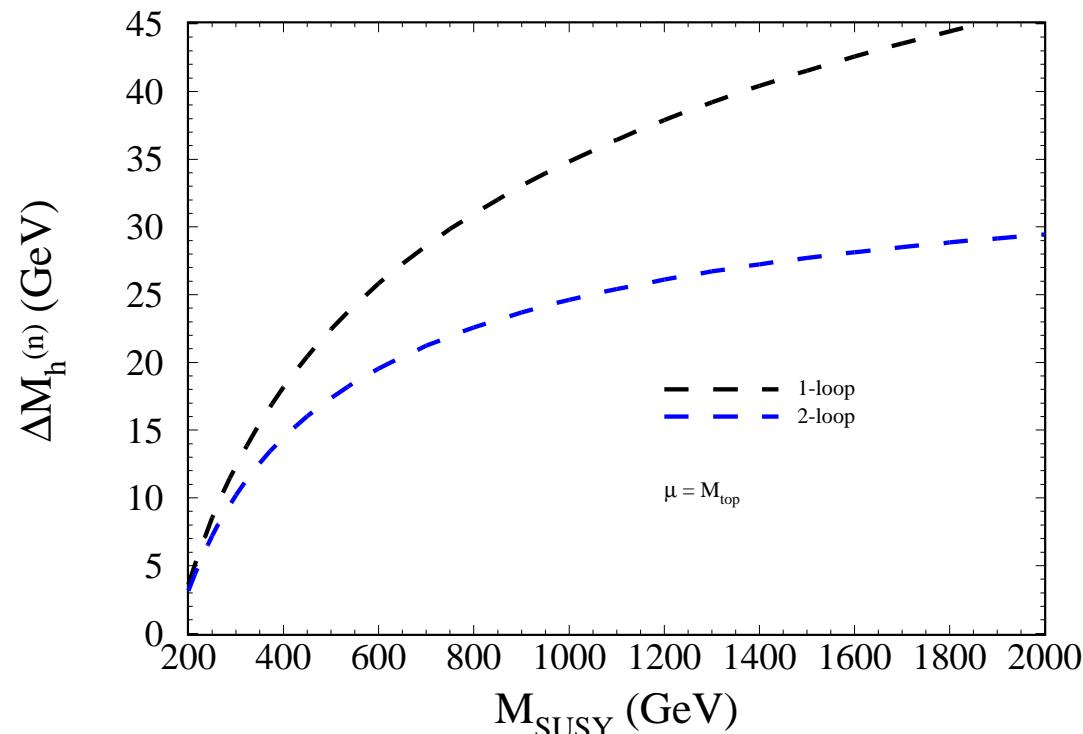
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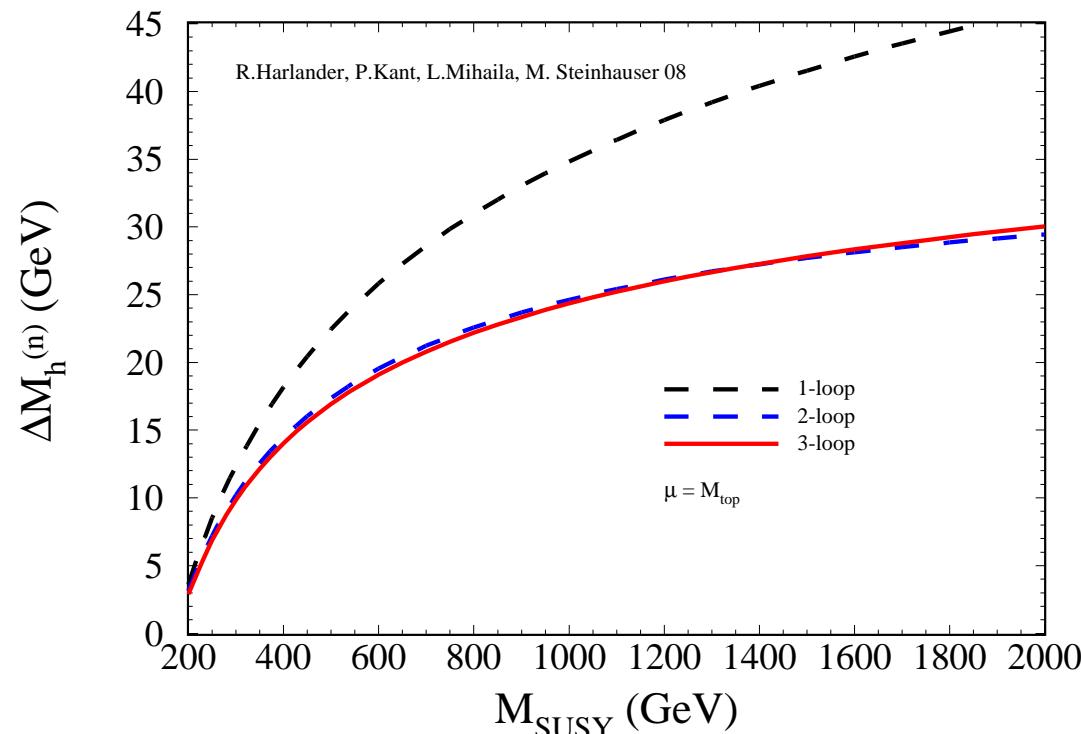
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$M_{\text{SUSY}} = 0.3 - 1 \text{ TeV} : \Delta M_h^{(3)} \simeq 500 \text{ MeV}$

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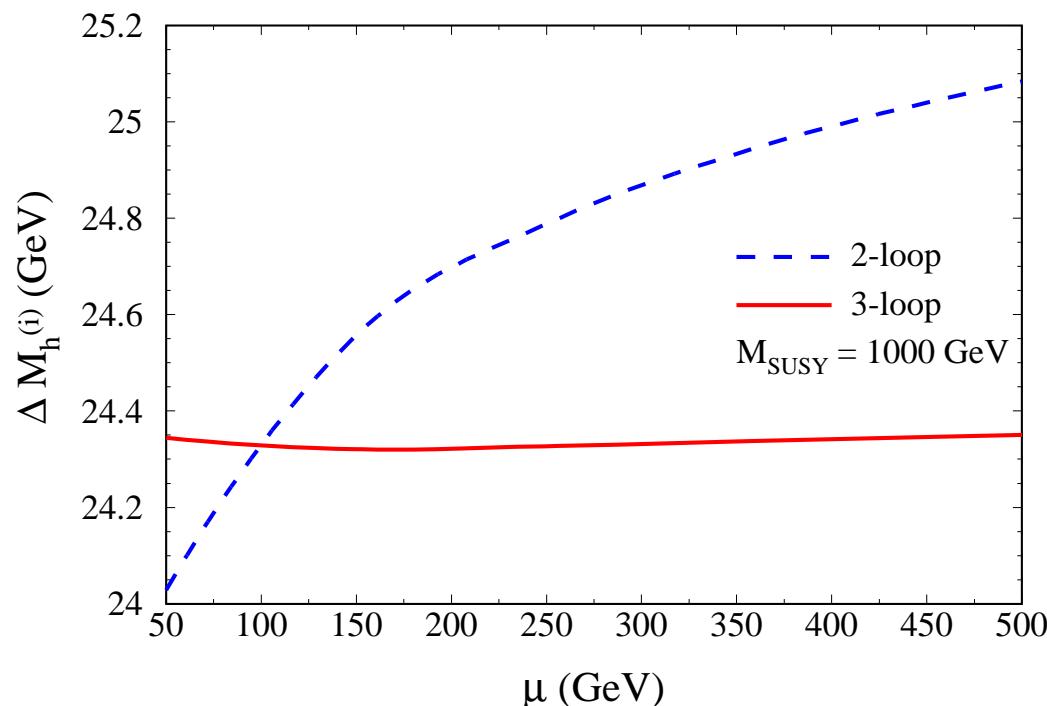
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Theoretical uncertainties : $\delta \Delta M_h^{(3)} \simeq 35 \text{ MeV}$



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Parametric uncertainties: dominated by M_t

$$\delta \Delta M_t^{\text{LHC}} \simeq 1 - 2 \text{ GeV} \Rightarrow \delta \Delta M_h \simeq 1 - 2 \text{ GeV}$$

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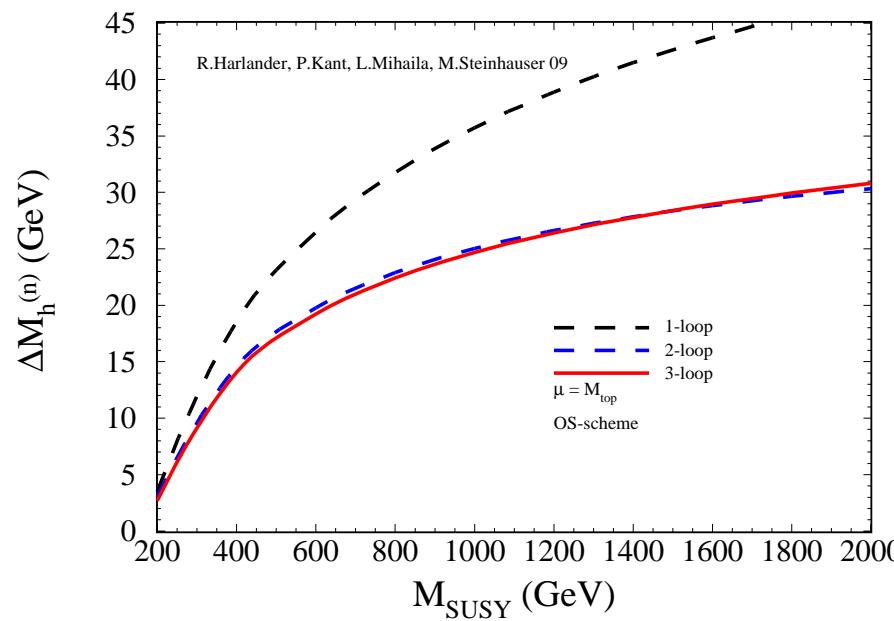
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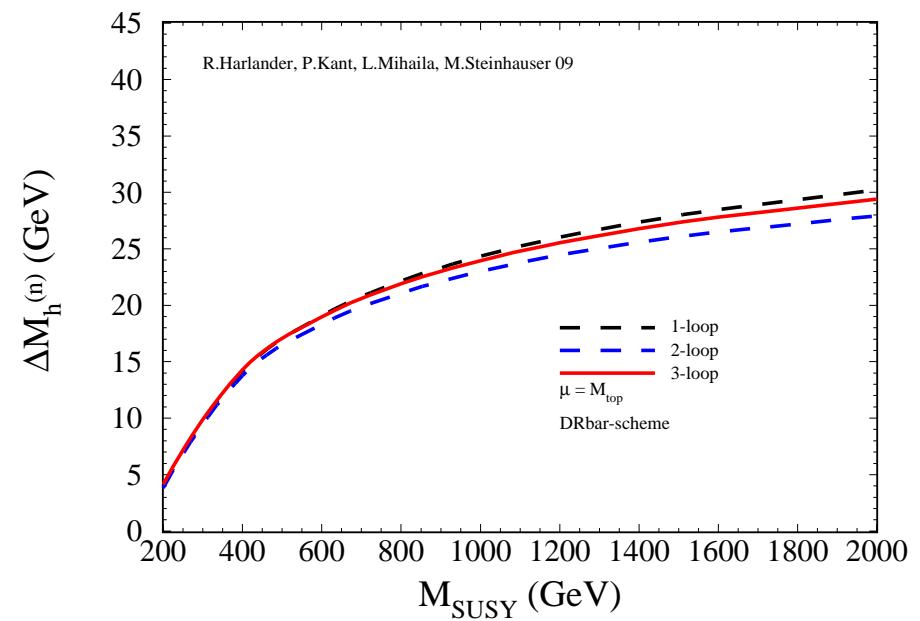
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Renormalization scheme dependence:

OS



$\overline{\text{DR}}$



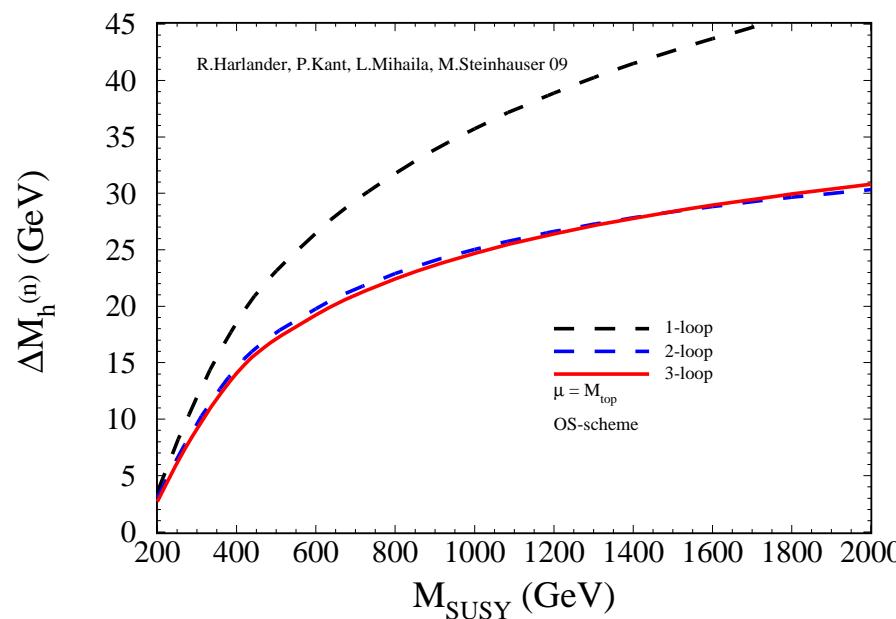
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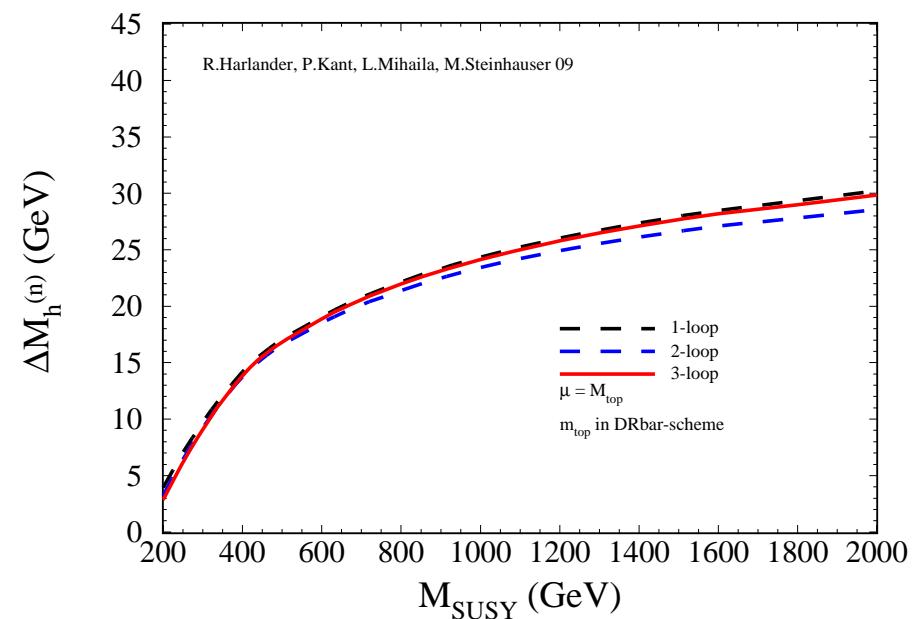
MSSM parameters: $M_A = 1 \text{ TeV}$ $\tan \beta = 40$; $A_t = 0$ $M_{\tilde{q}} = 2 \text{ TeV}$

Sensitivity to M_t renormalization scheme :

OS



m_t in $\overline{\text{DR}}$ -scheme: TSIL [S. Martin '05]



Numerical Results (stop-mixing)

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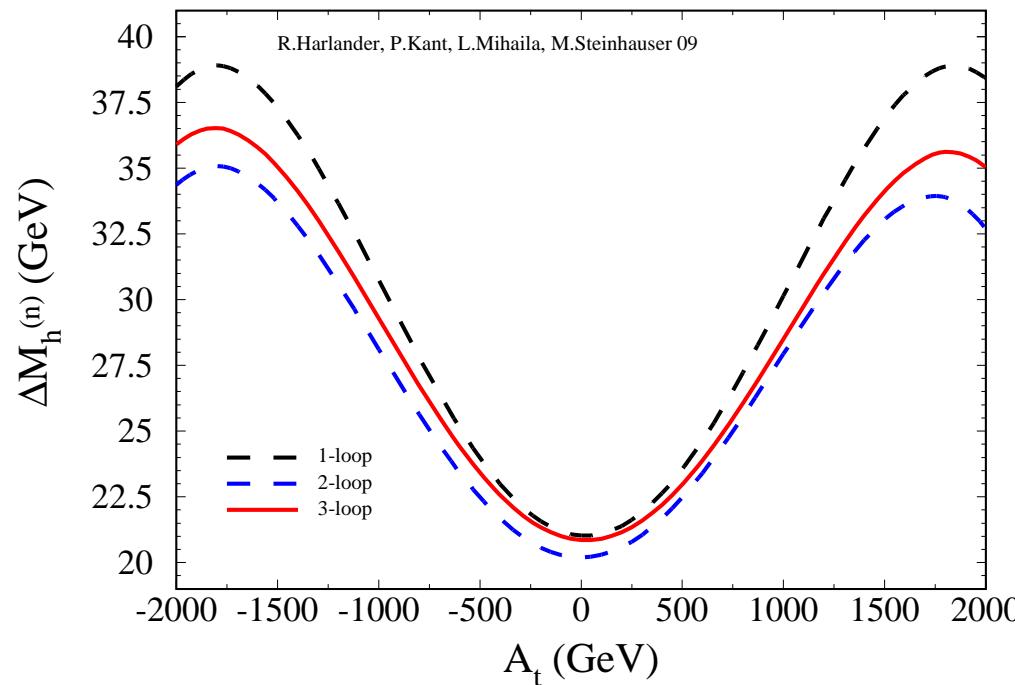
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$\overline{\text{DR}}$ scheme $A_t = 0 : \pm 2 \text{ TeV} : \Delta M_h^{(3)} = 0.5 - 1.5 \text{ GeV}$

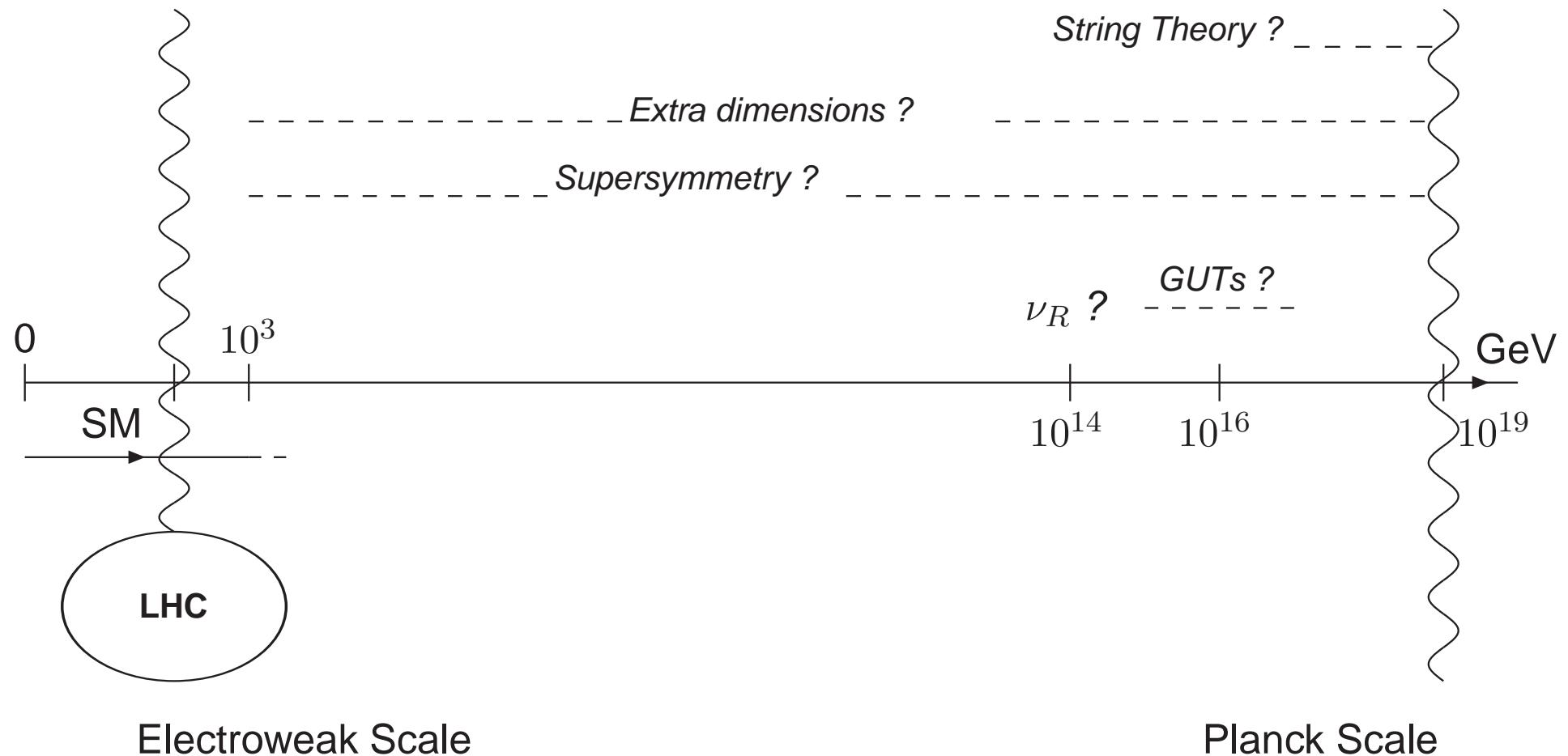


Conclusions

- m_h to 3-loop accuracy
 - 3-loop effects **larger** than experimental accuracy expected at LHC & ILC
 - 3-loop corrections **stabilize** the perturbative series
 - Stop-mixing & large $|A_t|$: $\Delta M_h^{(3)} \simeq 1.5 \text{ GeV}$
- ToDo:
 - Computer code to compute m_h for realistic SUSY mass spectrum
 - 3-loop effects due to CP-violation in the MSSM

Coupling constant unification

Open Questions



Grand Unification

- Gauge symmetry increases with energy [Georgi, Quinn, Weinberg '74]

Low energy			⇒	High energy
$SU_c(3) \otimes$	$SU_L(2) \otimes$	$U_Y(1)$	⇒	G_{GUT}
gluons	W, Z	photon	⇒	gauge bosons
quarks	leptons		⇒	fermions
g_3	g_2	g_1	⇒	g_{GUT}

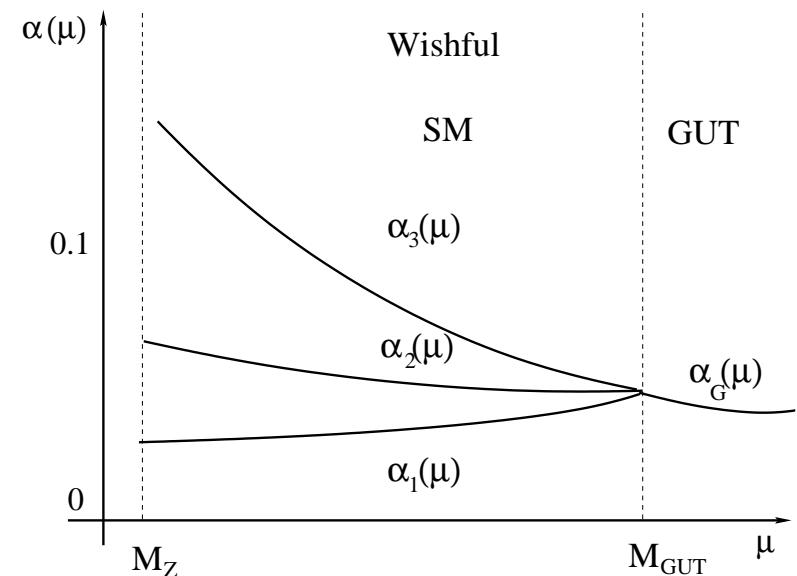
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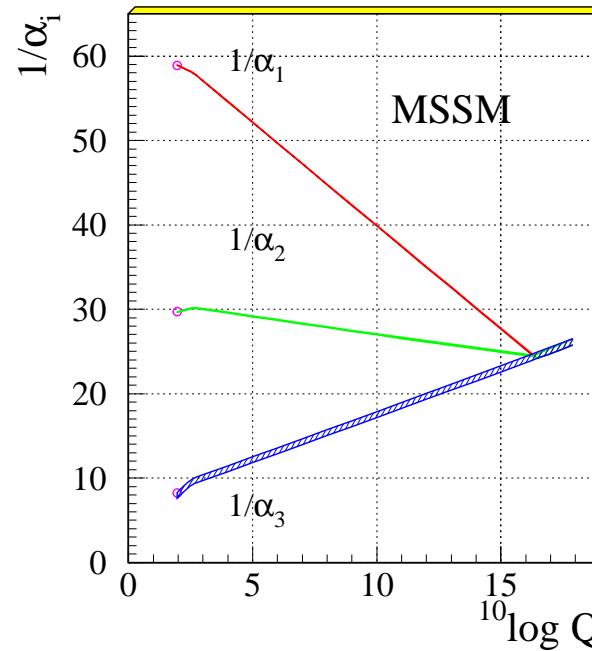
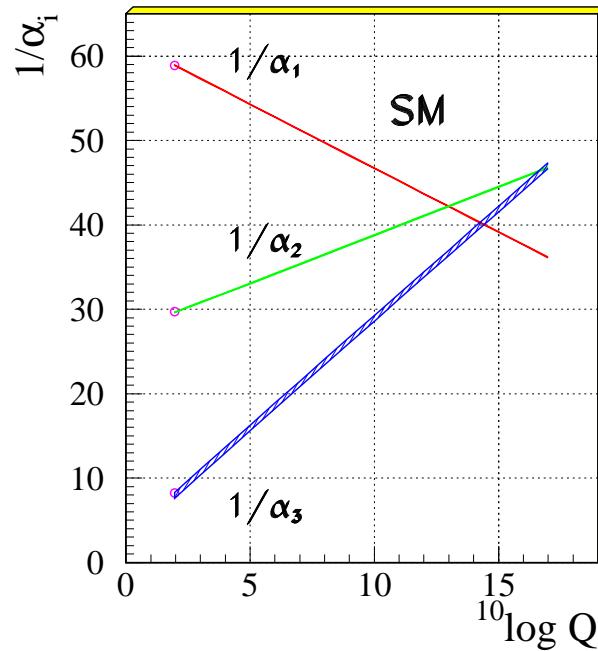
- Low energy interactions = branches of the unique interaction of a **simple** gauge group.

$$\begin{aligned} \alpha_i \equiv g_i^2 / 4\pi : \quad & \alpha_1(M_Z) = 0.017, \\ & \alpha_2(M_Z) = 0.034, \\ & \alpha_3(M_Z) = 0.118, \\ M_Z &= 91.1876 \text{ GeV}. \end{aligned}$$



MSSM and LEP data

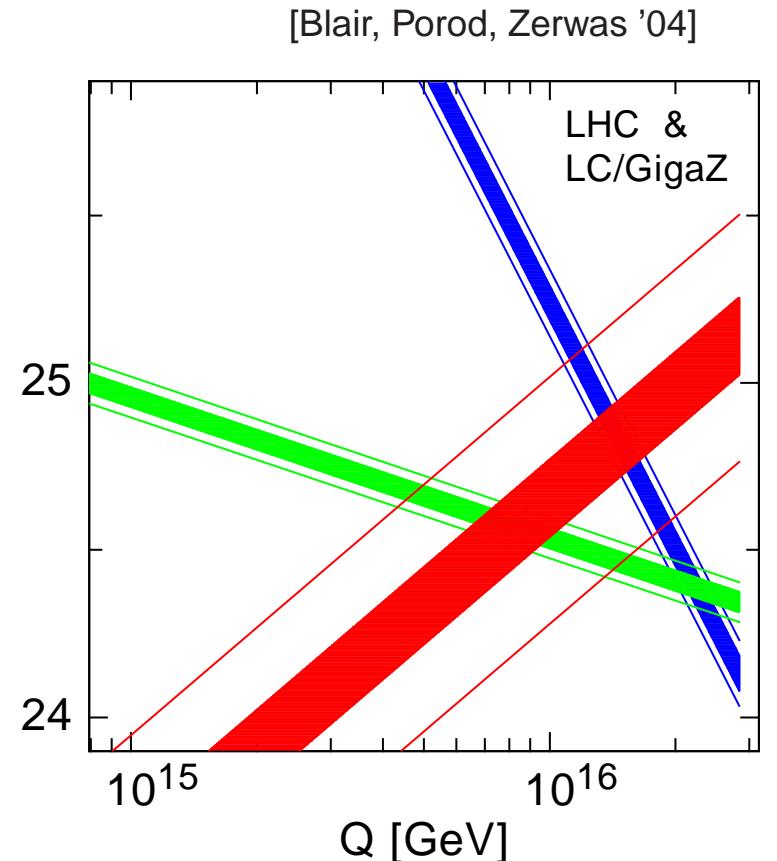
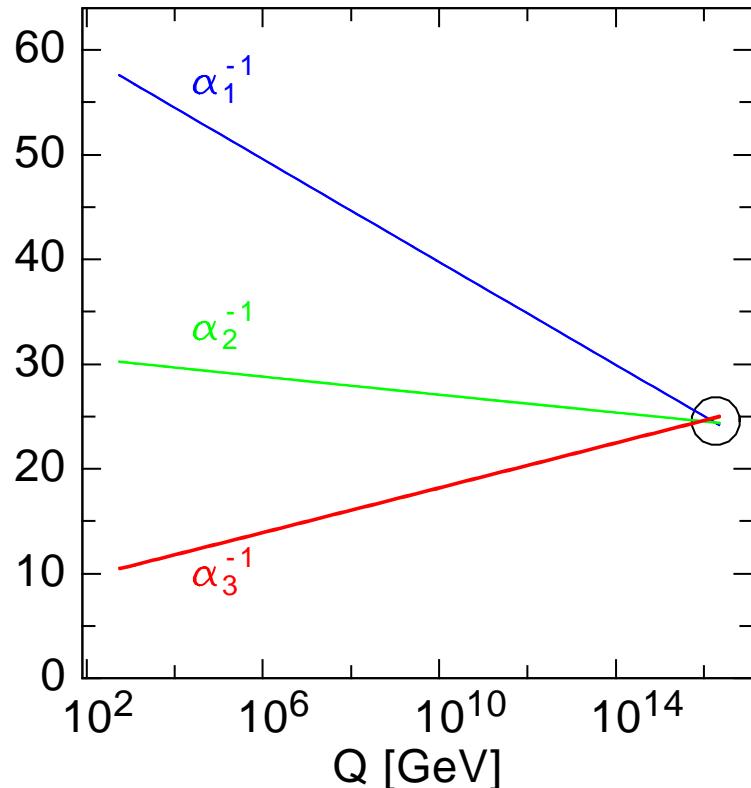
Unification of the Coupling Constants in the SM and the minimal MSSM



[Amaldi, Furstenau, de Boer]
[Langacker, Luo]
[Ellis, Kelley, Nanopoulos]

- Gauge Coupling unification within SM excluded by about 12σ .
- Gauge coupling Unification within SUSY GUTs works extremely well:
it fits within 3σ the present low energy data.

High precision data



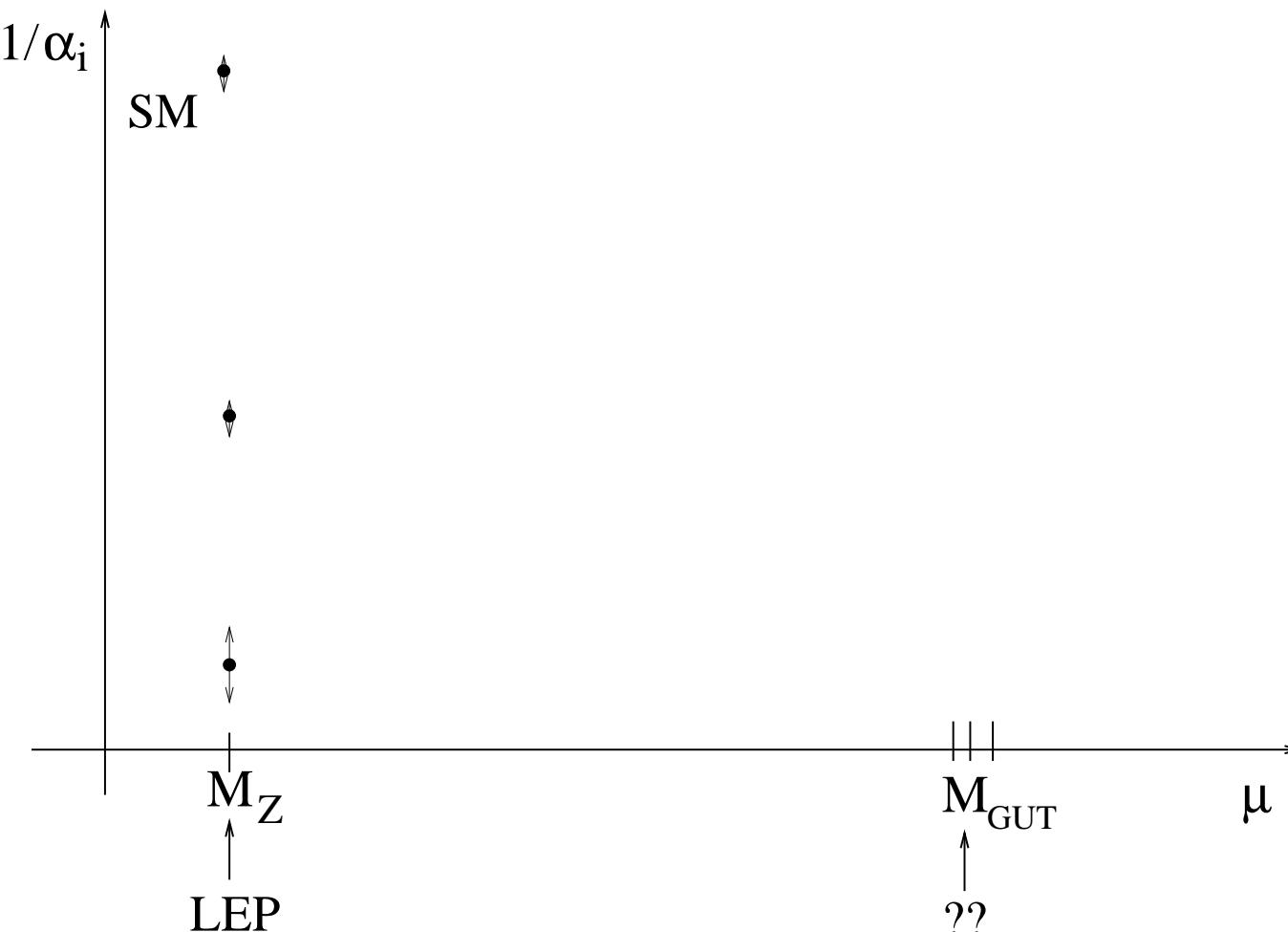
- Computation: common SUSY mass scale $\simeq 1$ TeV
2-loop Renormalization Group Running 1-loop threshold corrections at the weak scale (M_Z)
- Our aim: improve theoretical accuracy on $\alpha_s(M_{\text{GUT}})$ calculated from $\alpha_s(M_Z)$

GUTs

- GUT threshold corrections (effects of heavy particles $M \geq M_{\text{GUT}}$) \Rightarrow gauge coupling no longer meet at a point M_{GUT}

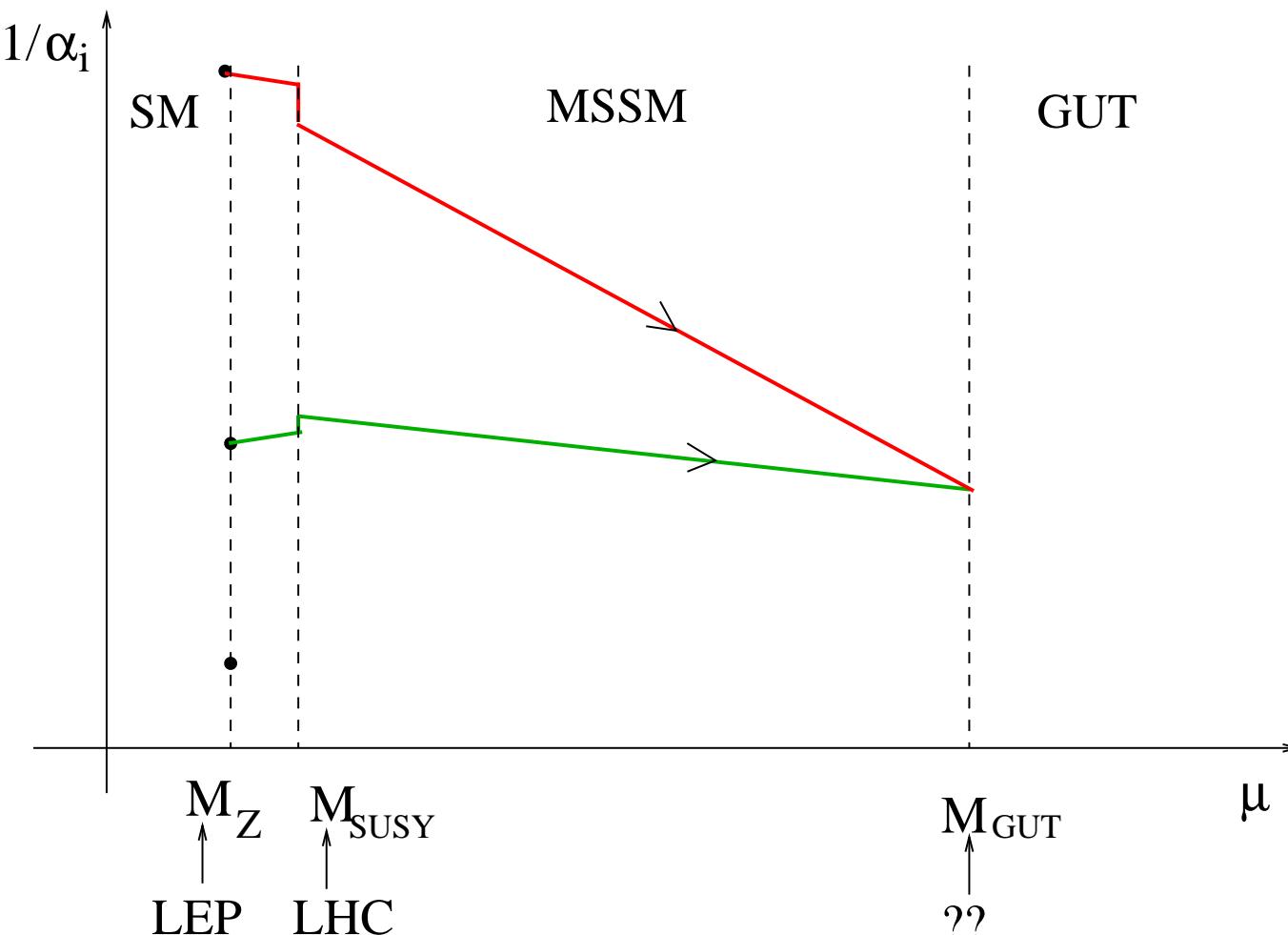
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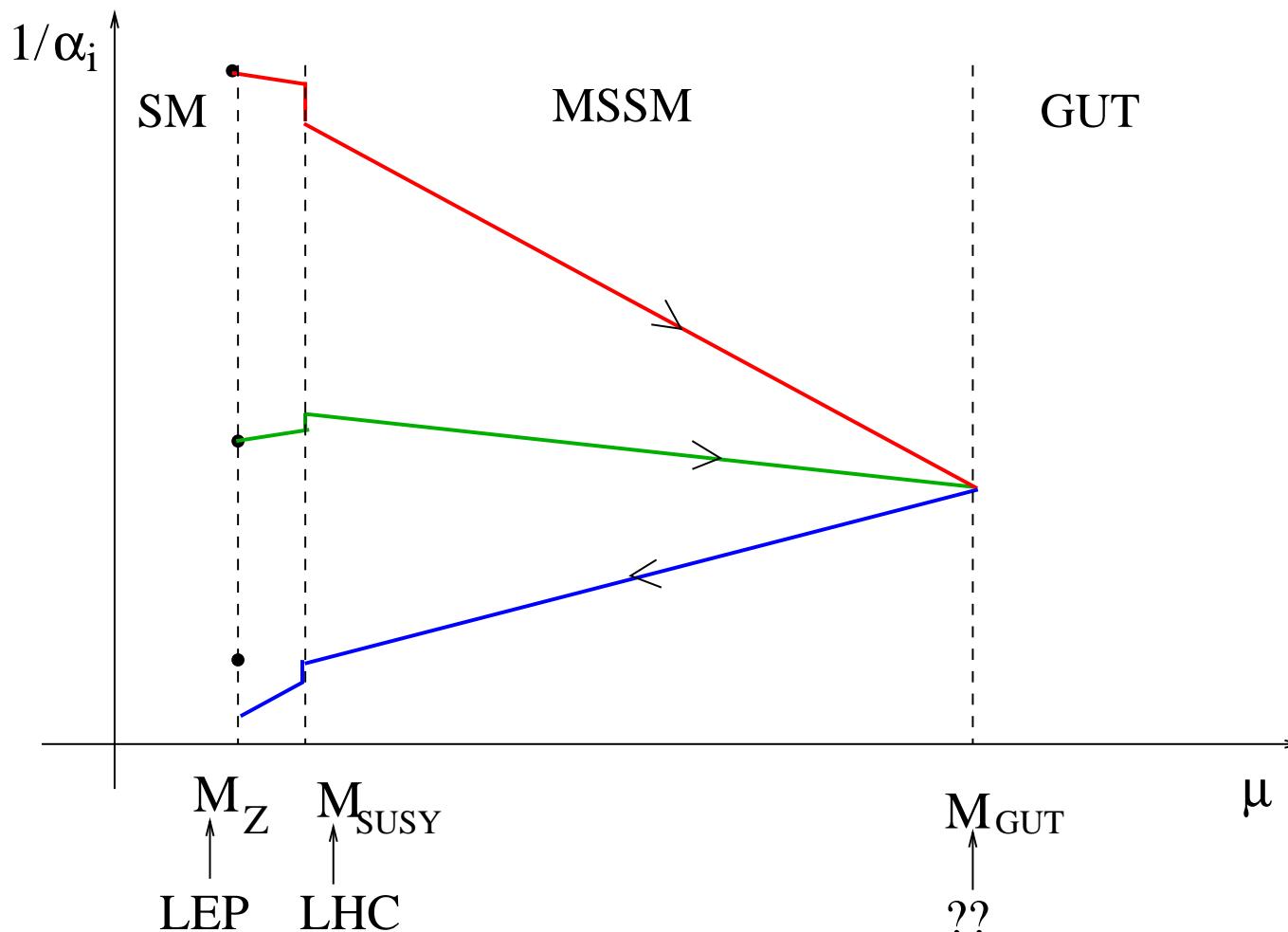
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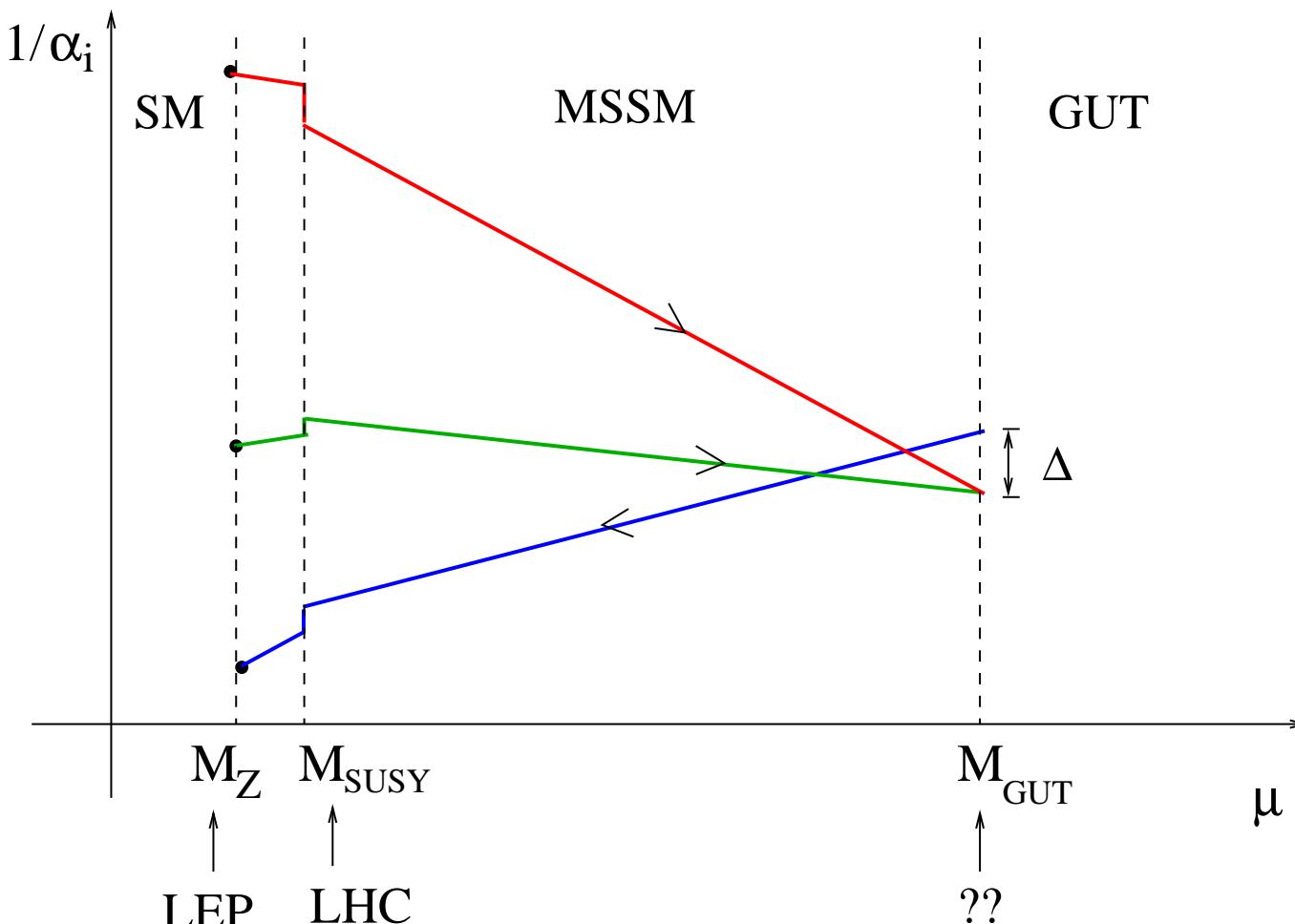
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- M_{GUT} : $\simeq 2 \times 10^{16}$ GeV in SUSY GUTs
 $\simeq 10^{15}$ GeV in nonSUSY GUTs
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 - 1-loop [K.Hagiwara, Y. Yamada '93],[J.Hisano, M. Murayama '94]:

$$\Delta \sim \alpha_G \ln \left(\frac{M_X}{M_{\text{GUT}}} \right) \quad X = \text{gauge bosons, Higgs ,...}$$

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- Proton decay in GUT:

$$\tau_p = f(\alpha_G, M_{\text{GUT}}, M_X, \dots)$$

- Experiment (Super-Kamiokande): $\tau_{p \rightarrow e^+ \pi^0} > 5. \times 10^{33}$ yrs (at 90% CL)

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⇒ severe constraints on possible local gauge symmetries in GUTs

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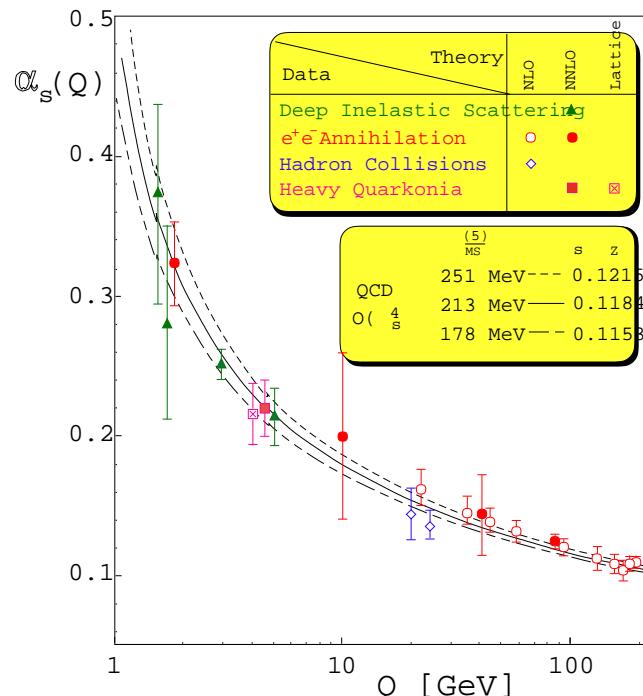
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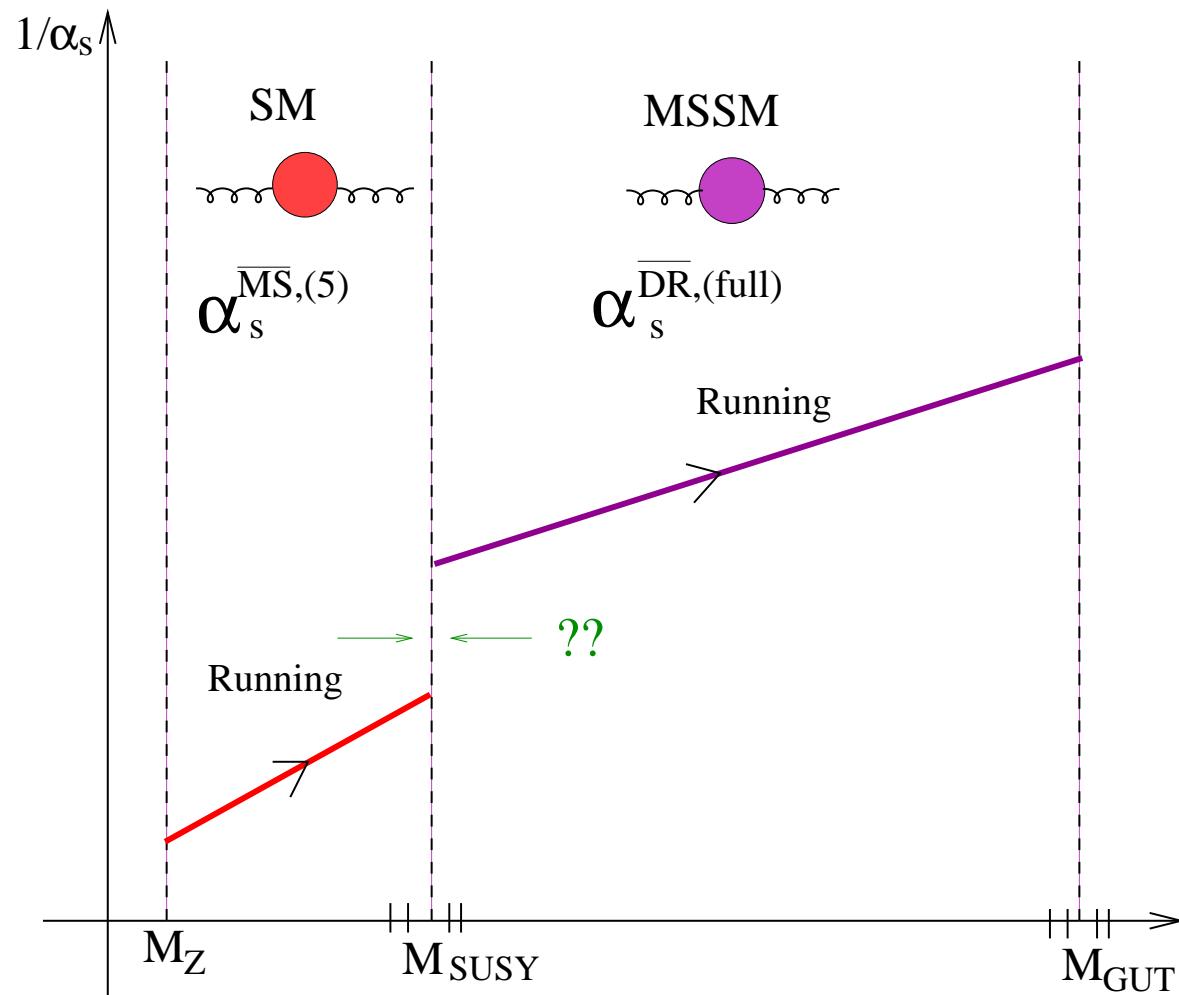
- Quantum Field Theory :

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- Vacuum can screen or anti-screen the gauge charges.
- Anti-screening gives rise to the asymptotic freedom of strong interactions.



Evolution of the strong coupling

● Input parameter: $\alpha_s^{\overline{\text{MS}},(5)}(M_Z)$ \Rightarrow Output parameter: $\alpha_s^{\overline{\text{DR}},(\text{full})}(M_{\text{GUT}})$



Running of α_s

$$\mu^2 \frac{d}{d\mu^2} \alpha_s(\mu) = \beta(\alpha_s)$$

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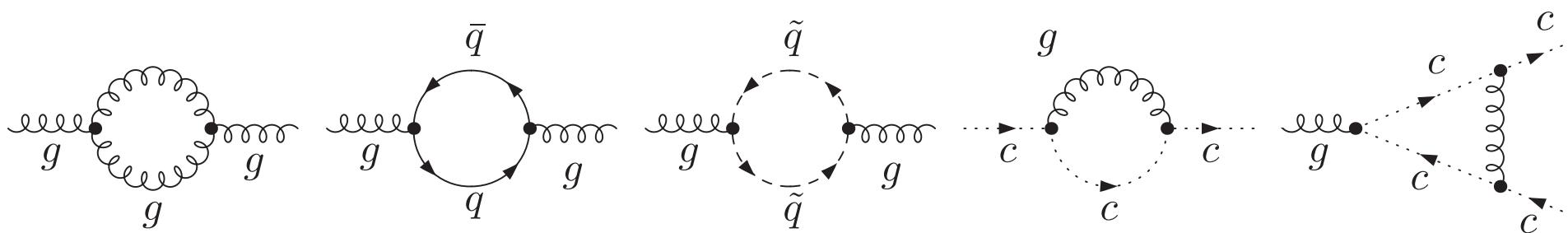
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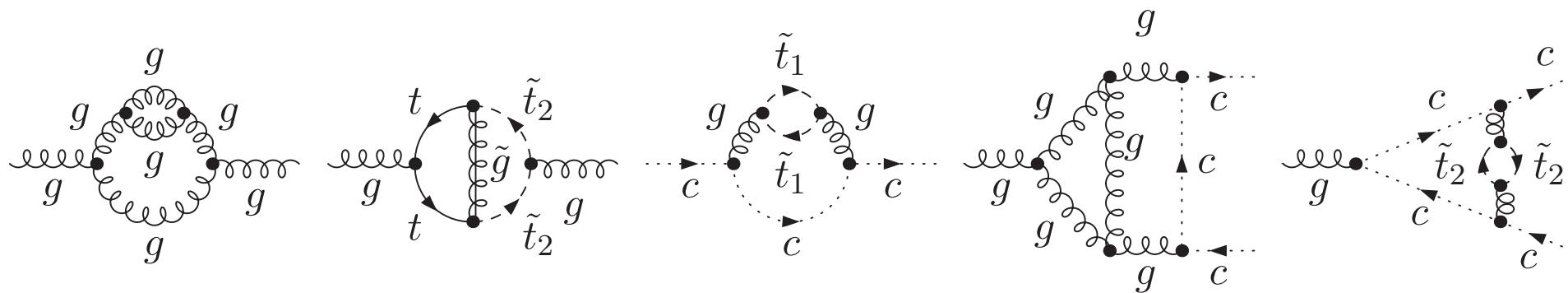
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2-loops:



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- 3-loop β_s in the MSSM

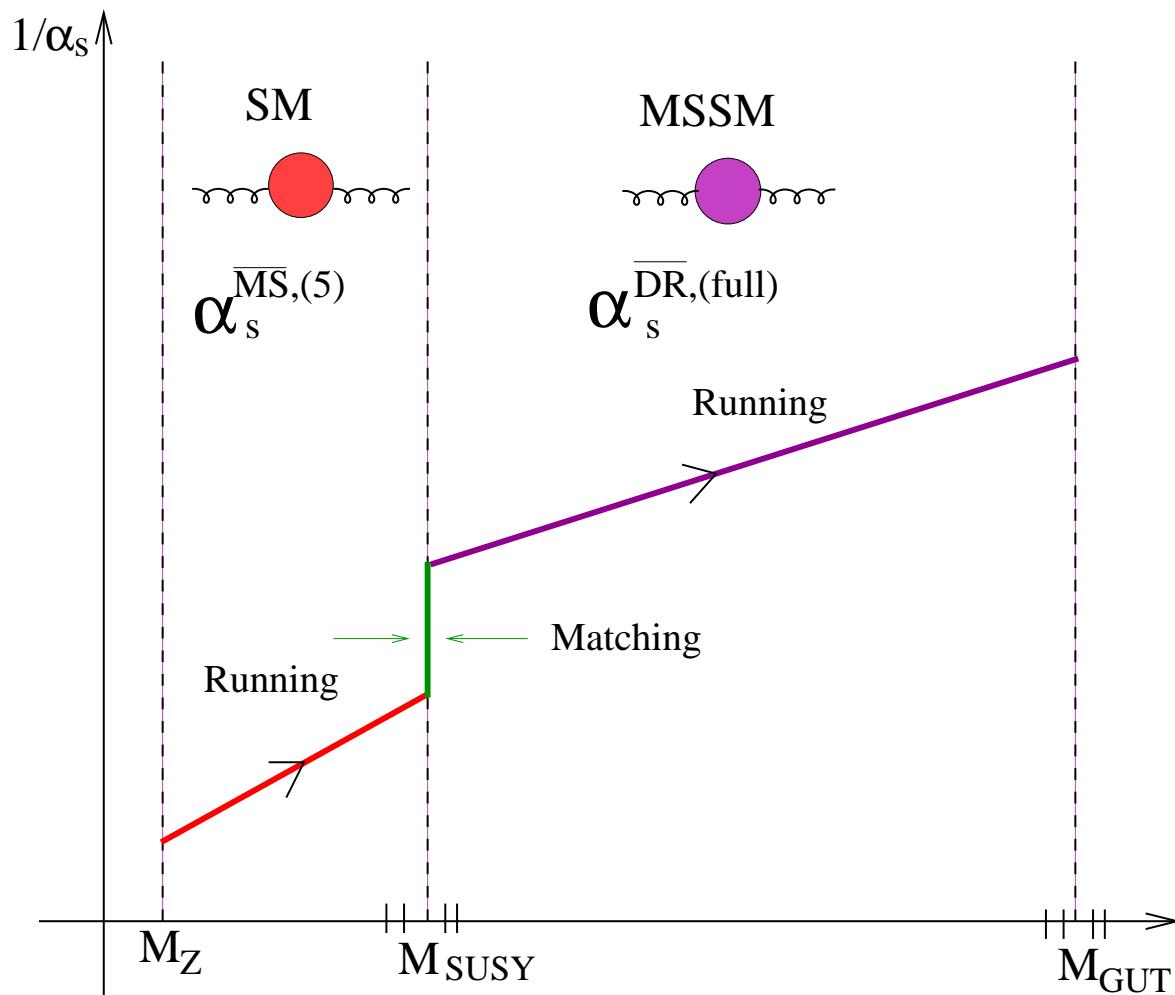
● $\simeq 100.000$ diagrams

● Computer programs: **QGRAF**, **FORM**, **MINCER**, **MATAD**, **EXP**, ...

[Noguiera; Vermaseren; Larin, Tkachov; Steinhauser; Seidensticker, Harlander; ...]

Matching

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$$\mathcal{L}_{\text{MSSM}}(\alpha_s^{(\text{full})}, \dots) \rightarrow \mathcal{L}(\alpha_s^{(5)}, \dots) \text{ at energy } \mu$$

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$$\begin{aligned}\alpha_s^{(5)} &= \zeta_s \alpha_s^{(\text{full})} \\ &\vdots \\ \zeta_s &= \zeta_s(\alpha_s, M_{\text{SUSY}}, m_t, \mu) \\ \zeta_s &= \text{matching coefficient}\end{aligned}$$

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- μ not predicted by theory
- Physical quantities must be independent of μ
- Quantum corrections improve stability

Matching coefficient for α_s

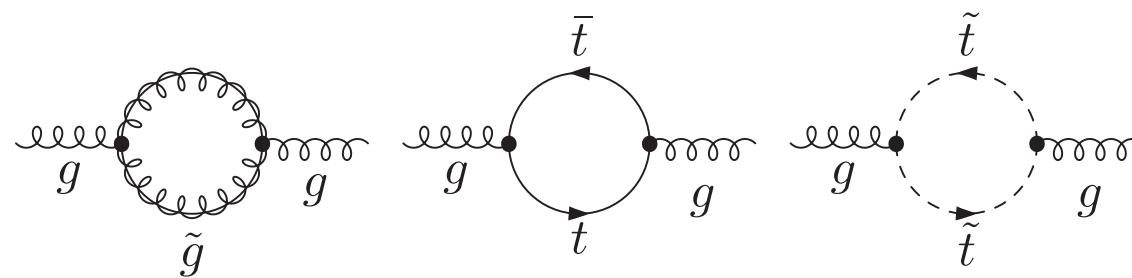
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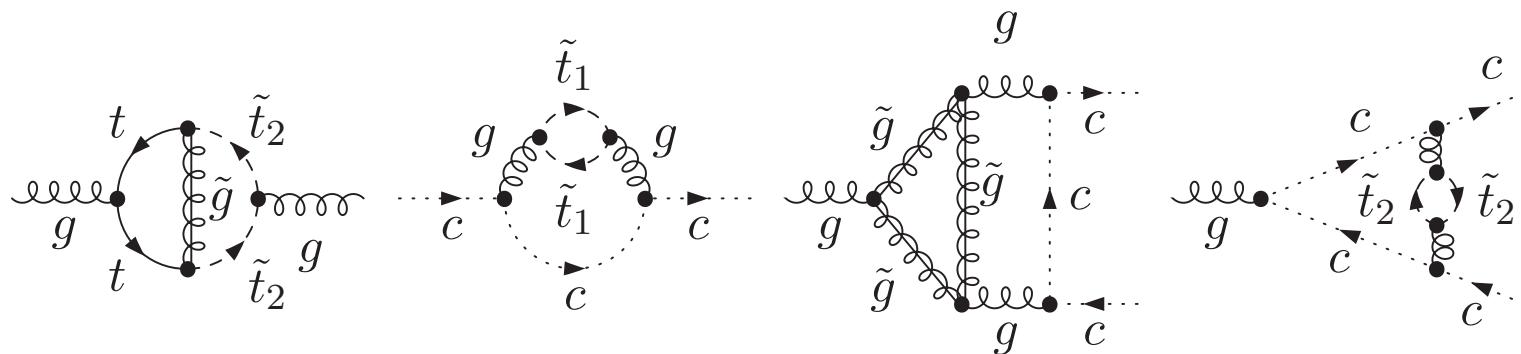
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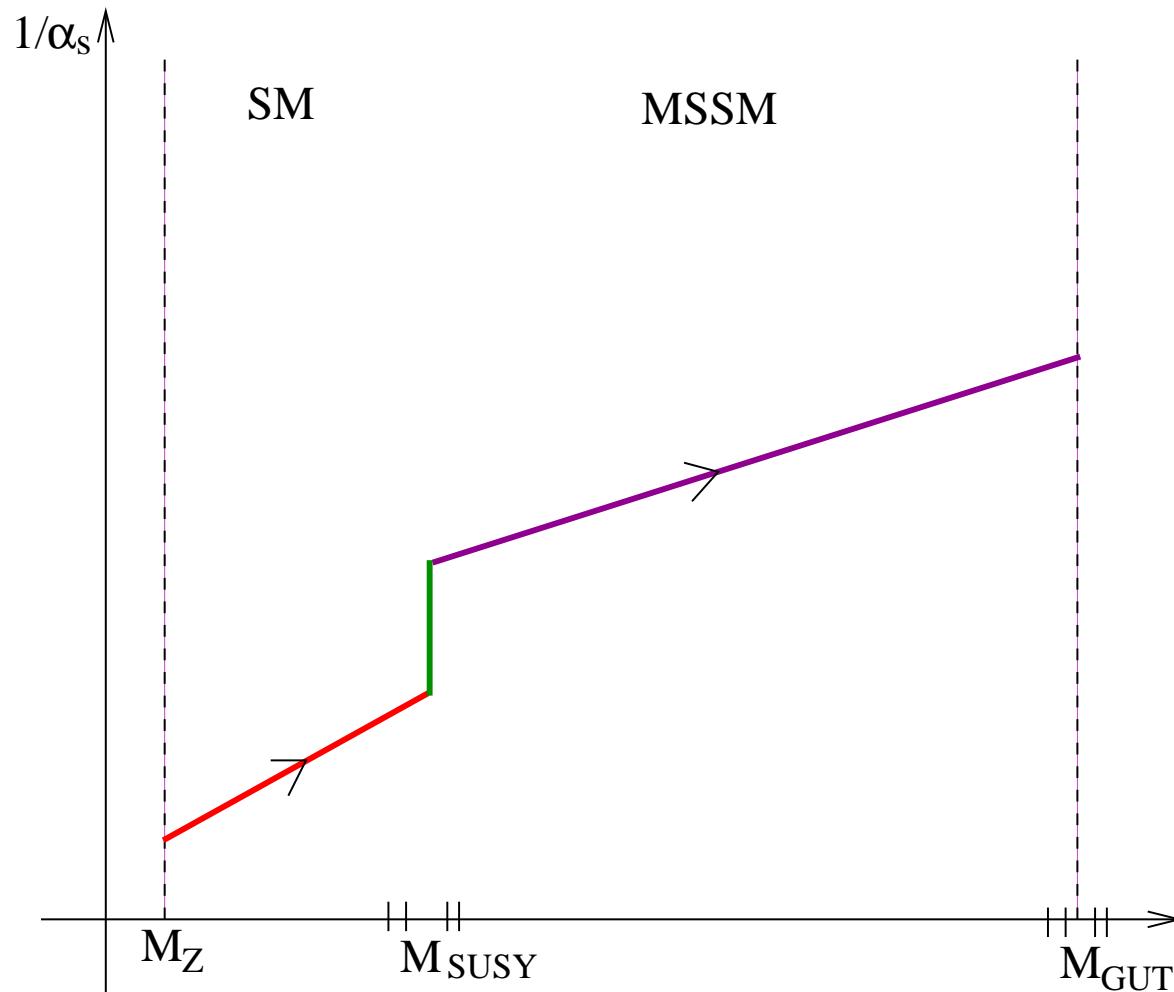
$$\begin{aligned} \zeta_{s2}^{(5)} = & -\frac{215}{96} - \frac{19}{24} \ln \frac{\mu^2}{m_t^2} - \frac{5}{2} \ln \frac{\mu^2}{\tilde{M}^2} + \left[\frac{1}{6} \ln \frac{\mu^2}{m_t^2} + \ln \frac{\mu^2}{\tilde{M}^2} \right]^2 \\ & + \left(\frac{m_t}{\tilde{M}} \right)^2 \left(\frac{5}{48} + \frac{3}{8} \ln \frac{m_t^2}{\tilde{M}^2} \right) - \frac{7\pi}{36} \left(\frac{m_t}{\tilde{M}} \right)^3 \\ & + \left(\frac{m_t}{\tilde{M}} \right)^4 \left(\frac{881}{7200} - \frac{1}{80} \ln \frac{m_t^2}{\tilde{M}^2} \right) + \frac{7\pi}{288} \left(\frac{m_t}{\tilde{M}} \right)^5 \end{aligned}$$

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- Wish: $\alpha_s(M_{\text{GUT}})$ independent of the matching scale

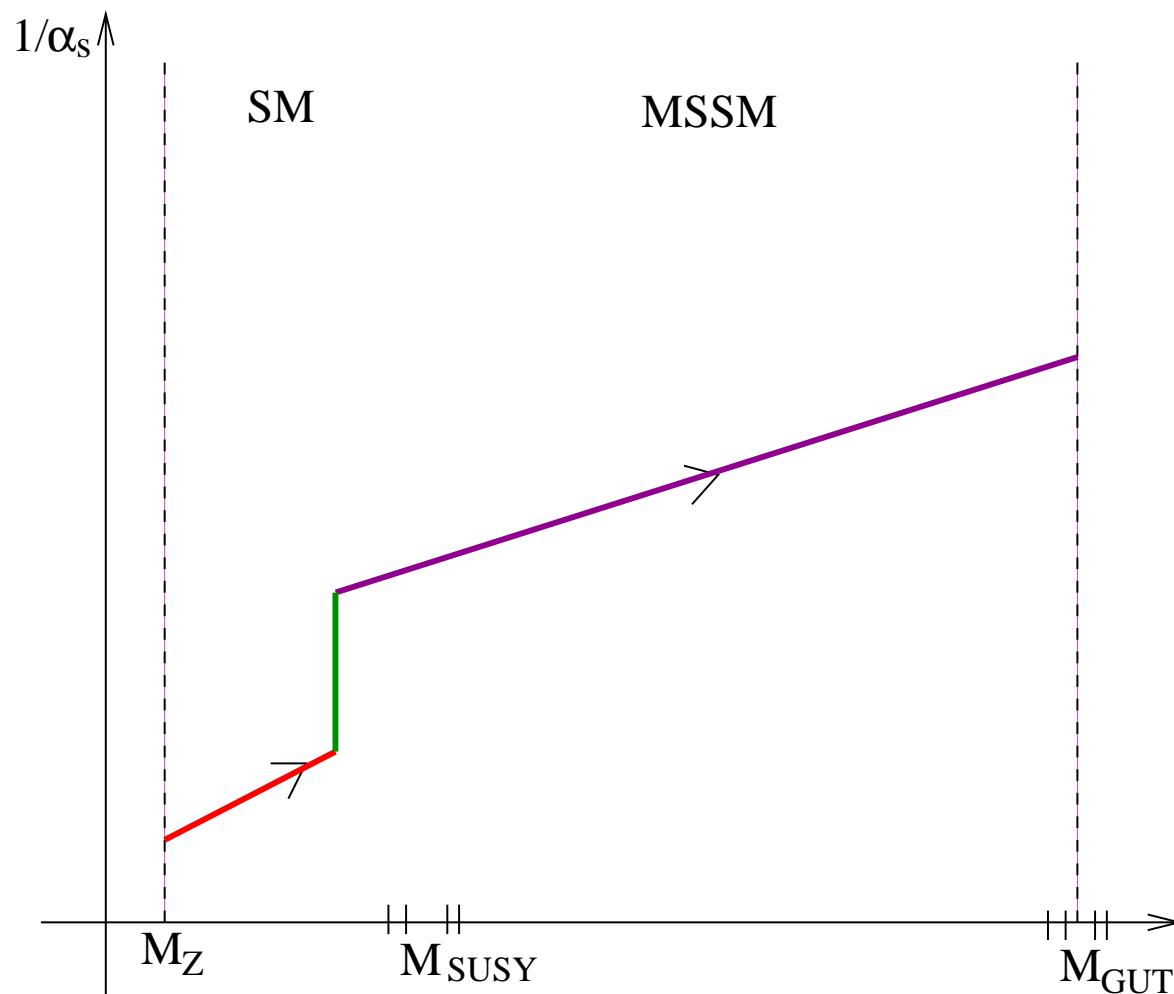
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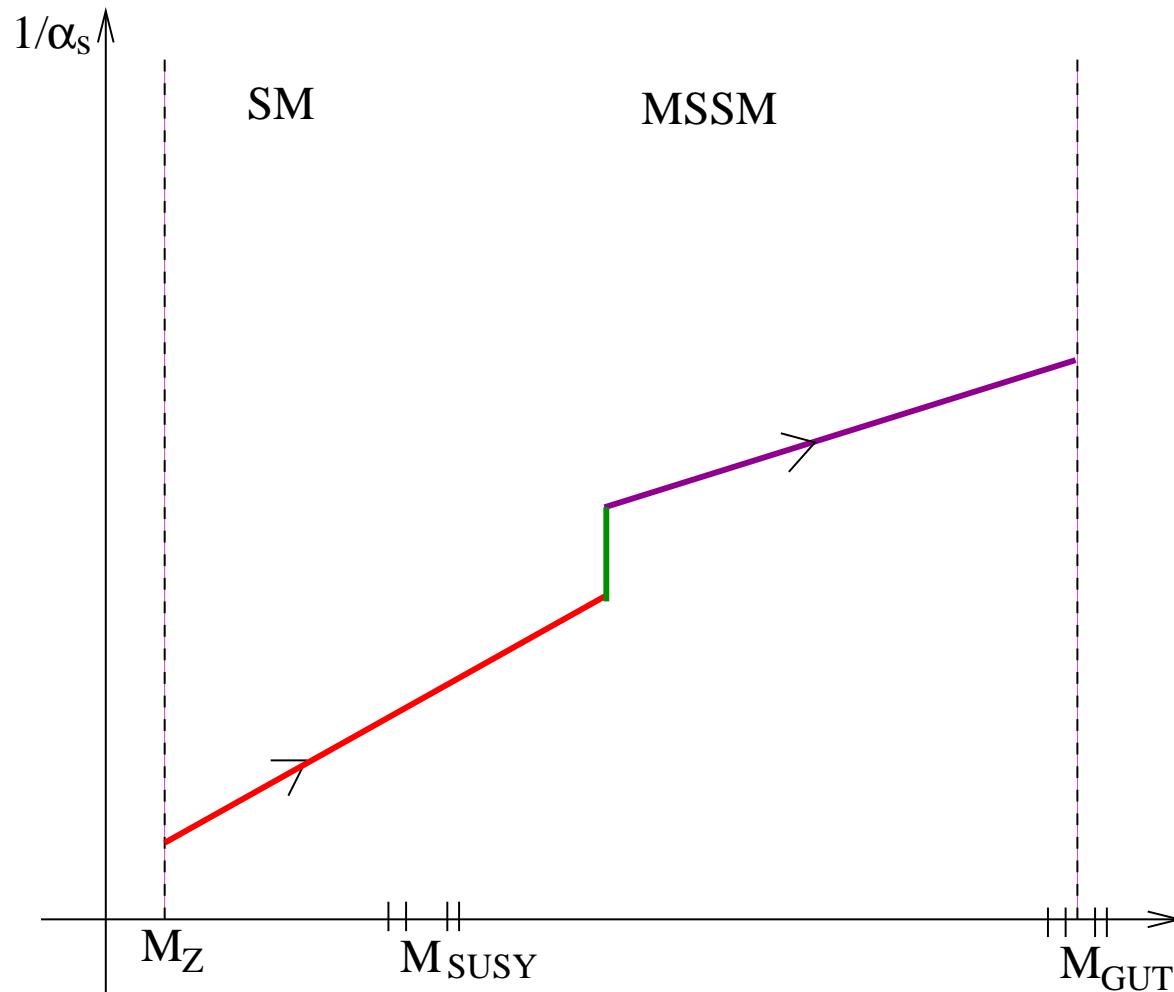
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$$\alpha_s(M_{\text{GUT}})$$

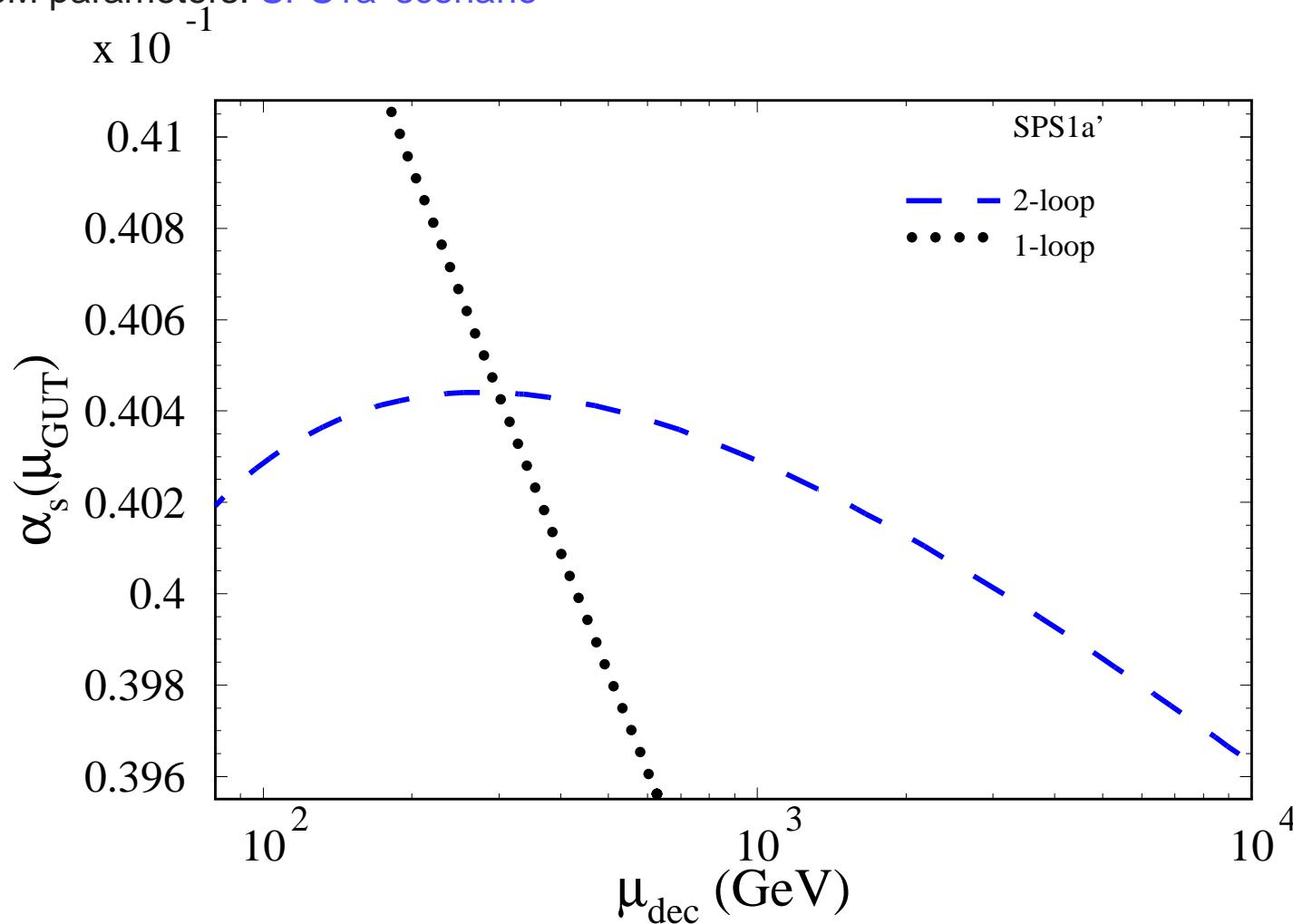
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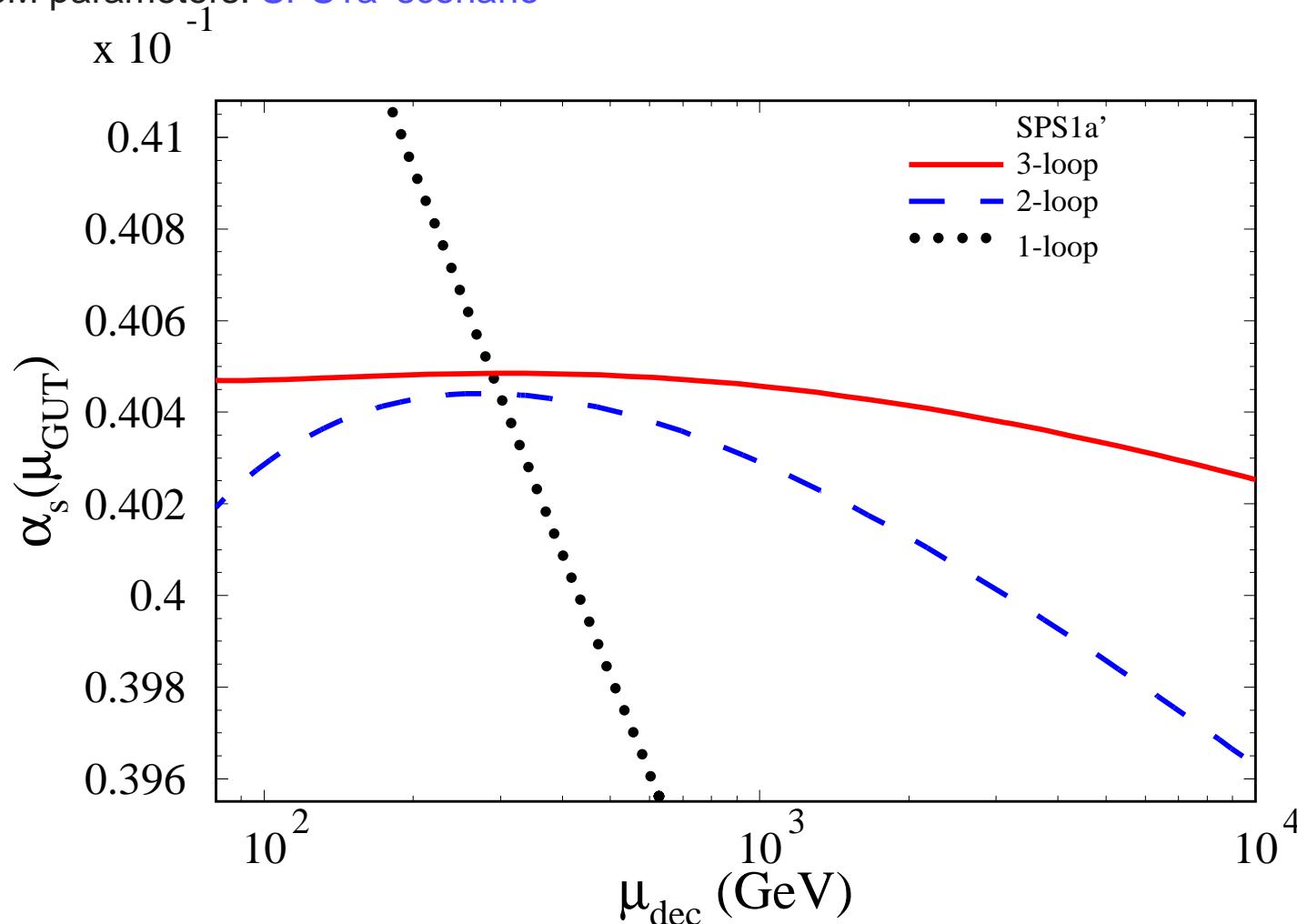
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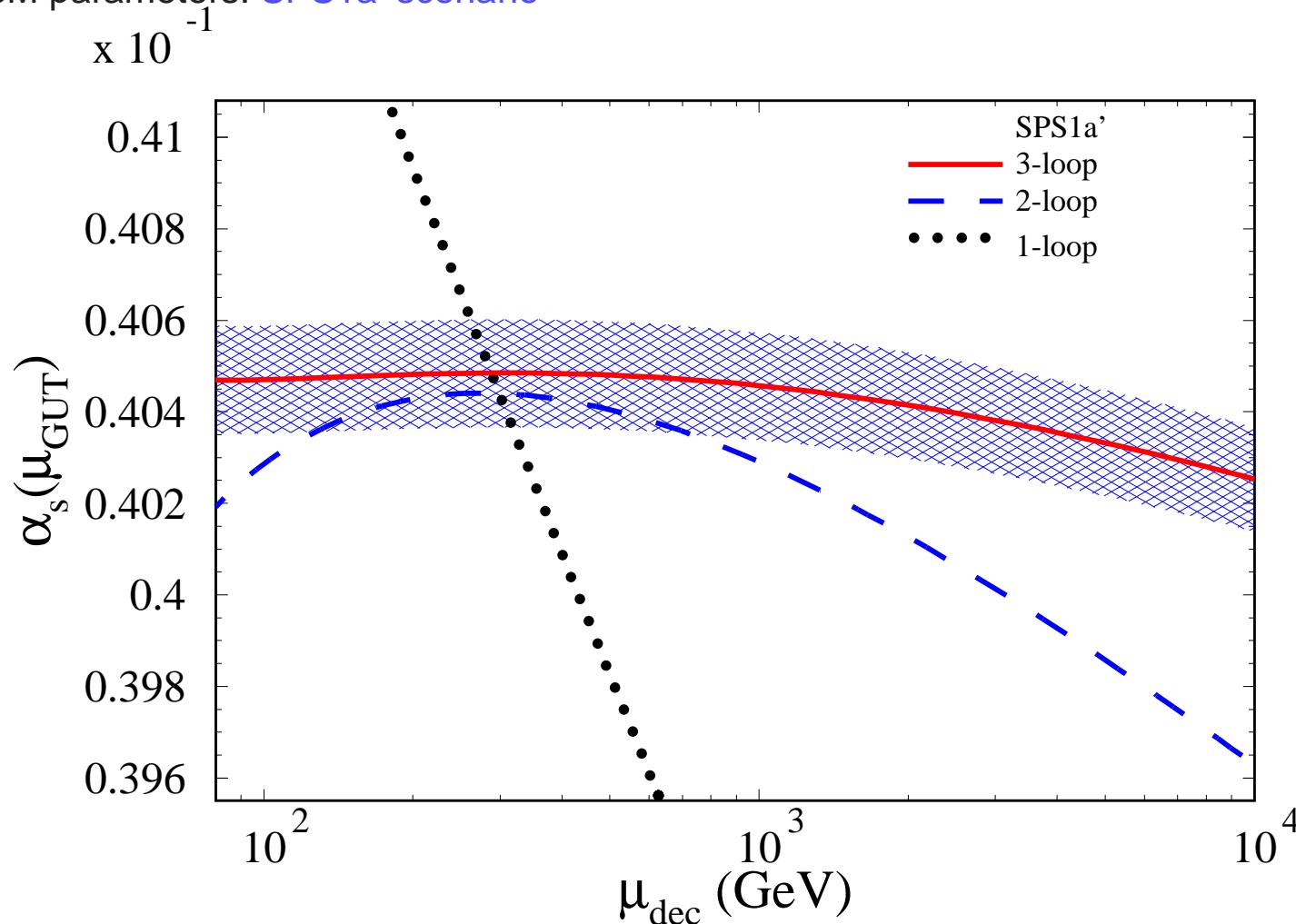
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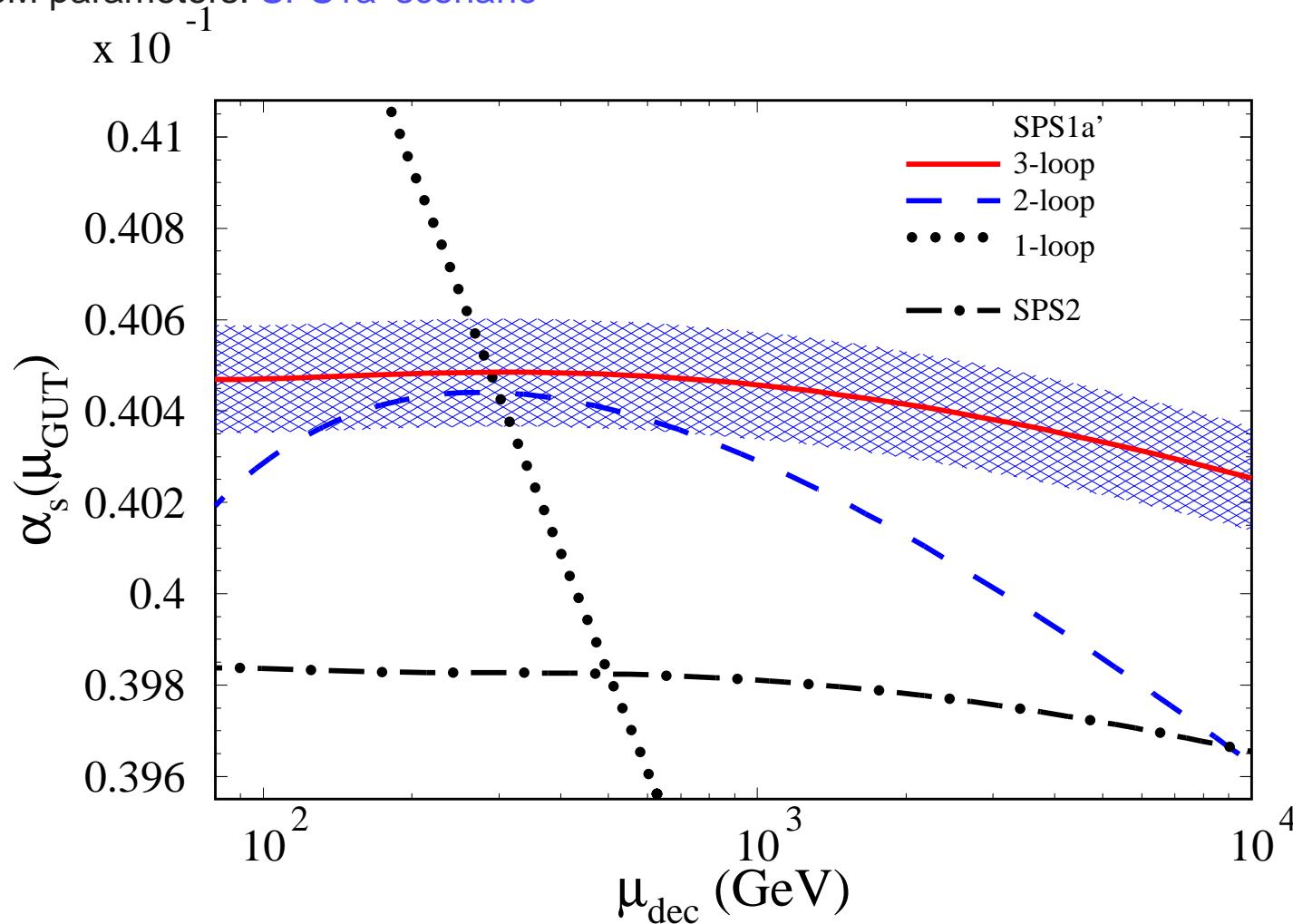
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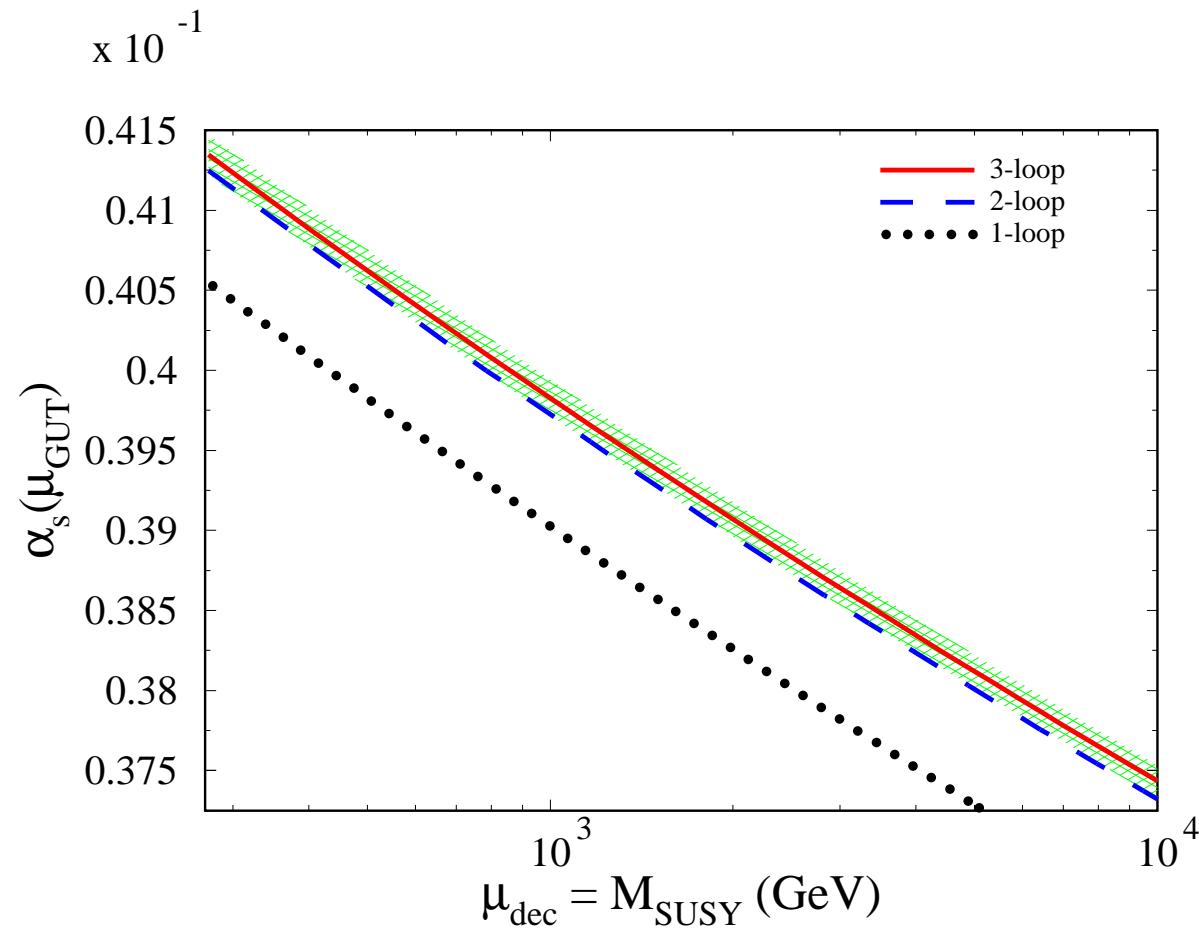
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M_{SUSY}

- Sensitivity of $\alpha_s(M_{\text{GUT}})$ to SUSY-mass scale:



Bottom quark mass

- (SUSY)GUT models \Rightarrow predictions for $m_t, m_b/m_\tau$
- Bottom quark mass in SM: known with 4-loop accuracy

$$\delta m_b^{\overline{\text{MS}}}(m_b) = 25 \text{ MeV} \quad [\text{J. H. K\"uhn, M. Steinhauser, C. Sturm '07}]$$

- Bottom quark in MSSM (models with large $\tan \beta$)
 - SUSY mass spectrum sensitive to bottom Yukawa coupling
 - $Y_b(\mu) \leftrightarrow m^{\overline{\text{DR}}}(\mu)$ affected by large SUSY radiative corrections

$m_b(M_{\text{GUT}})$

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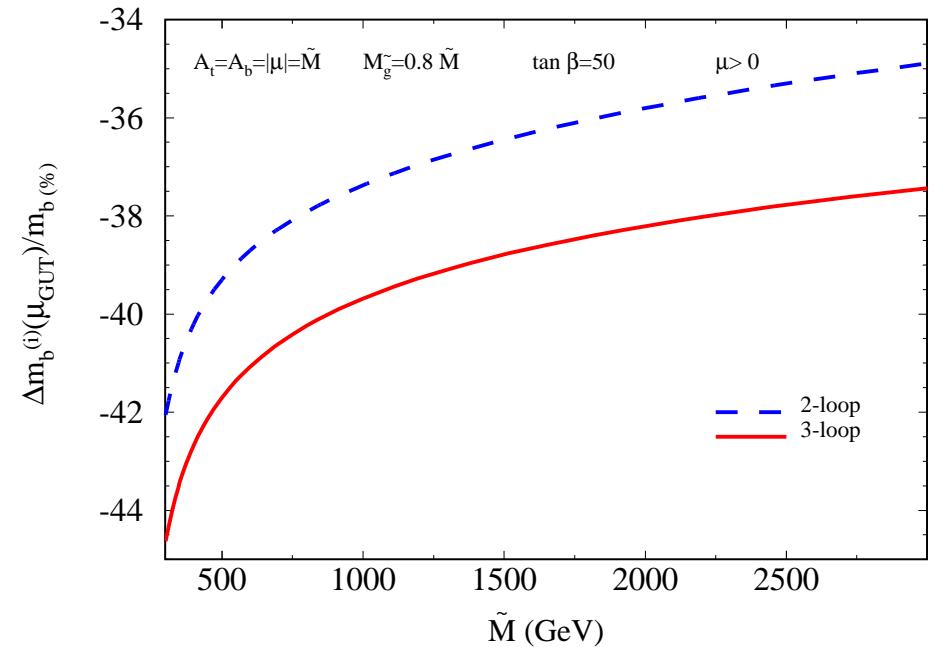
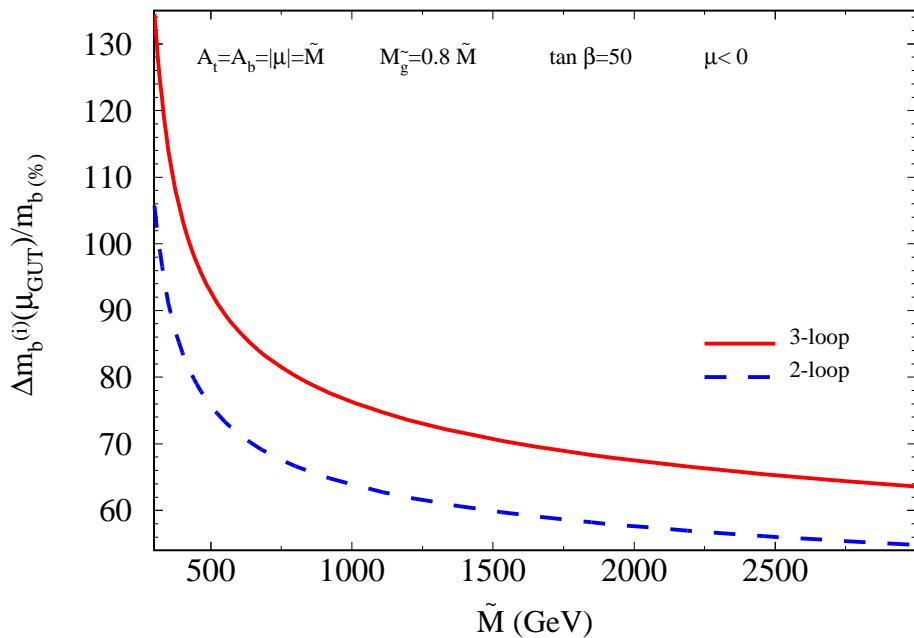
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simplified assumptions for the MSSM parameters

[A. Bauer, L. M., J. Salomon '08]



Conclusions

- $\alpha_s^{\overline{\text{DR}}}(M_{\text{GUT}})$
 - The **3-loop** effects comparable with the experimental accuracy for α_s
 - $\alpha_s(M_{\text{GUT}})$ very sensitive to SUSY-mass scale
- $m_b^{\overline{\text{DR}}}(M_{\text{GUT}})$
 - MSSM with large $\tan \beta$: **3-loop** effects reach up to 30%
- ToDo:
 - combine: 3-loop running analysis for the strong sector
known 2-loop running for the electroweak sector
 - extend analysis to SUSY-GUT models

“With the LHC, we will expand the frontiers of fundamental physics.

... We will learn whether existing indications for unification and supersymmetry have been Nature teaching us or Nature teasing us.”

[Franck Wilczek, SUSY'07]