

Probing quantum gravity at all scales

Bridging high and low energies in search of quantum gravity
COST Action CA23130 Annual Conference, July 8, 2025

Astrid Eichhorn, Heidelberg University



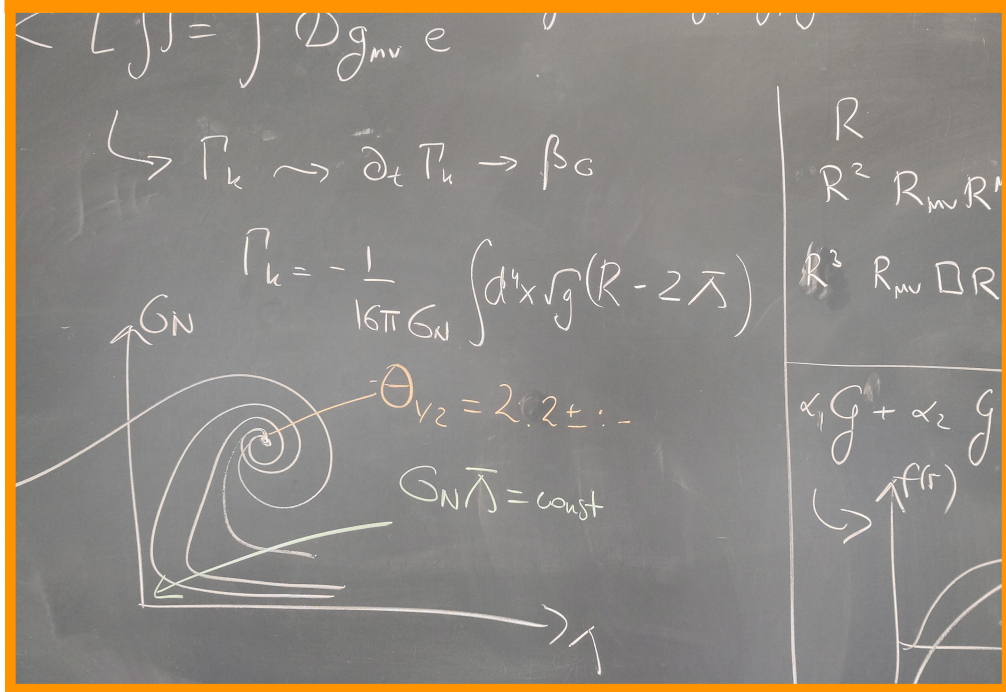
European Research Council



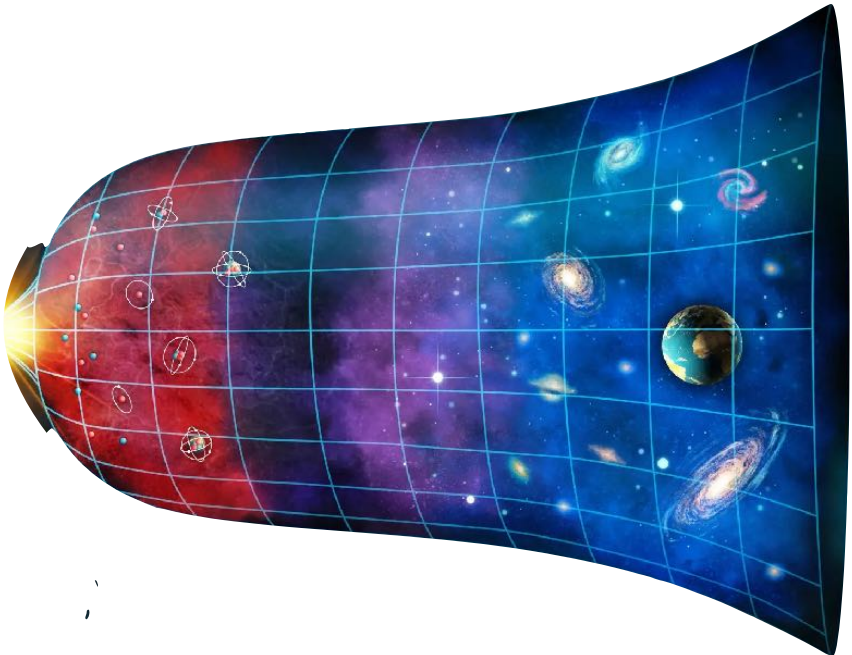
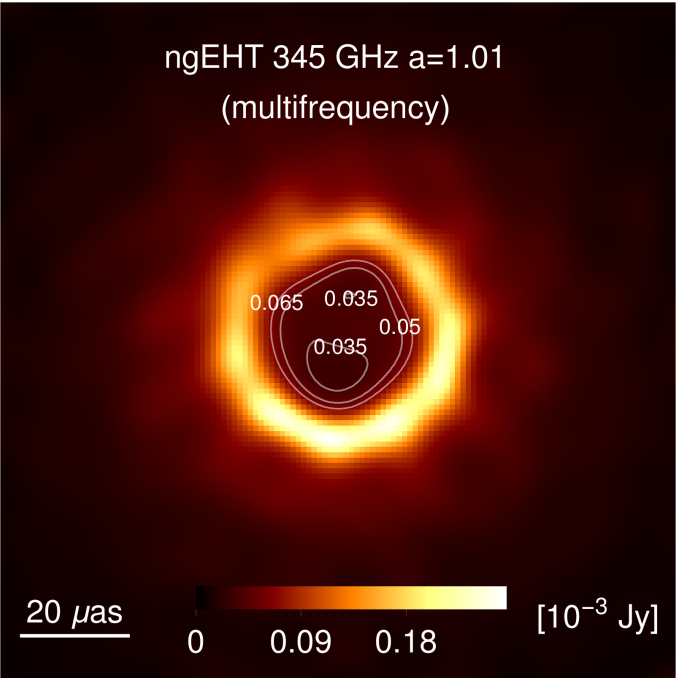
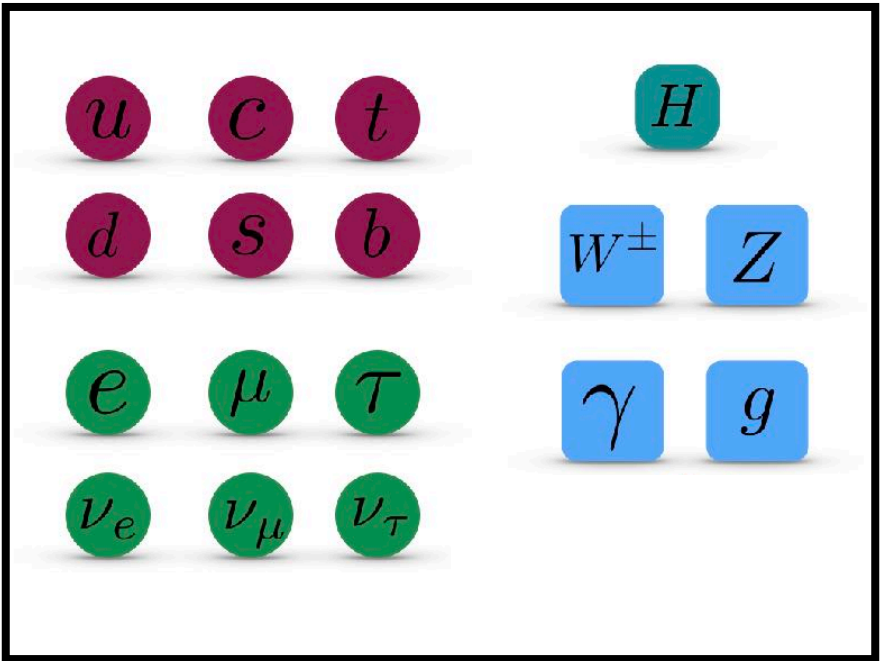
**UNIVERSITÄT
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ZUKUNFT
SEIT 1386

Motivation: How to test proposed theories of quantum gravity?

Theory of quantum gravity



Key challenge: gap in scales



Planckian scales

10^{-35} m

Particle physics scales

10^{-17} m

Black-hole scales

10^{11} m

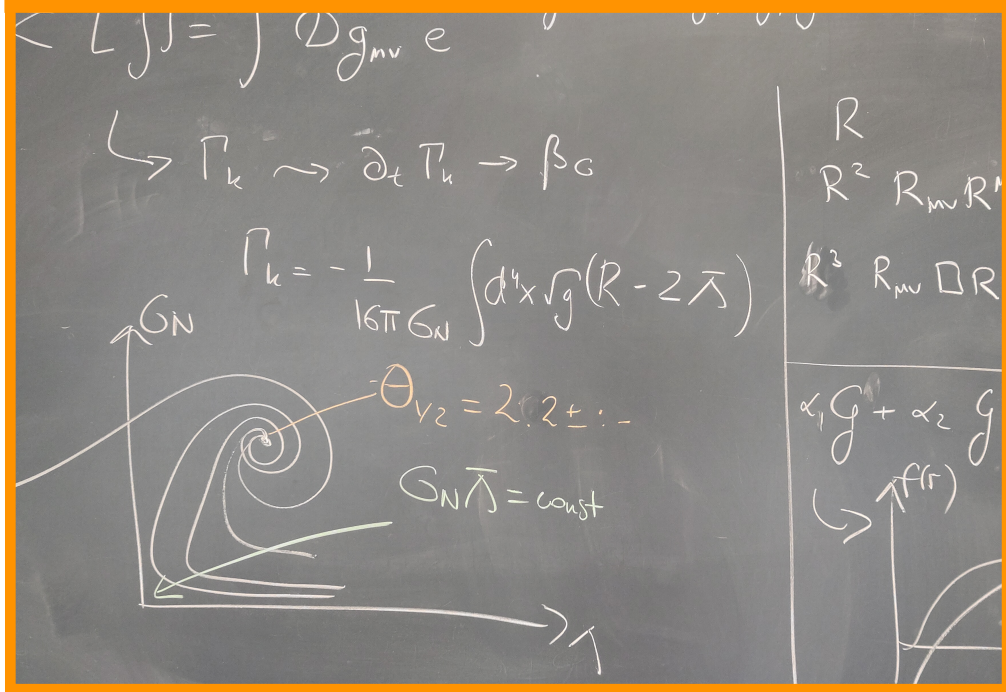
Cosmological scales

$> 10^{20} \text{ m}$

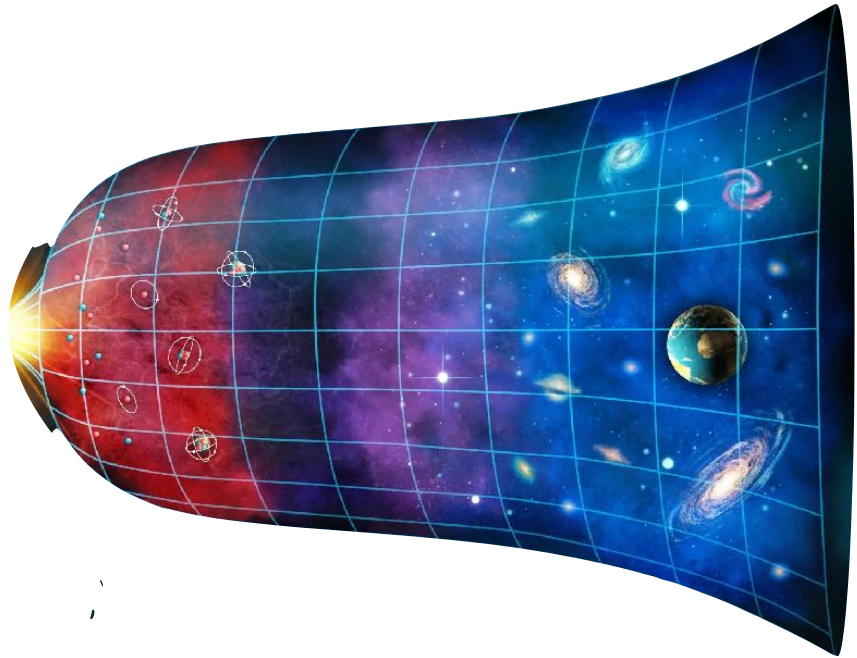
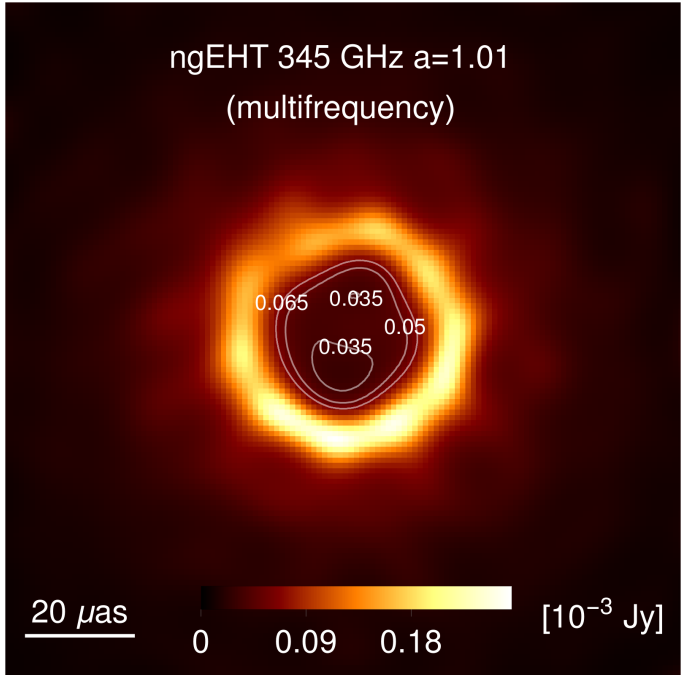
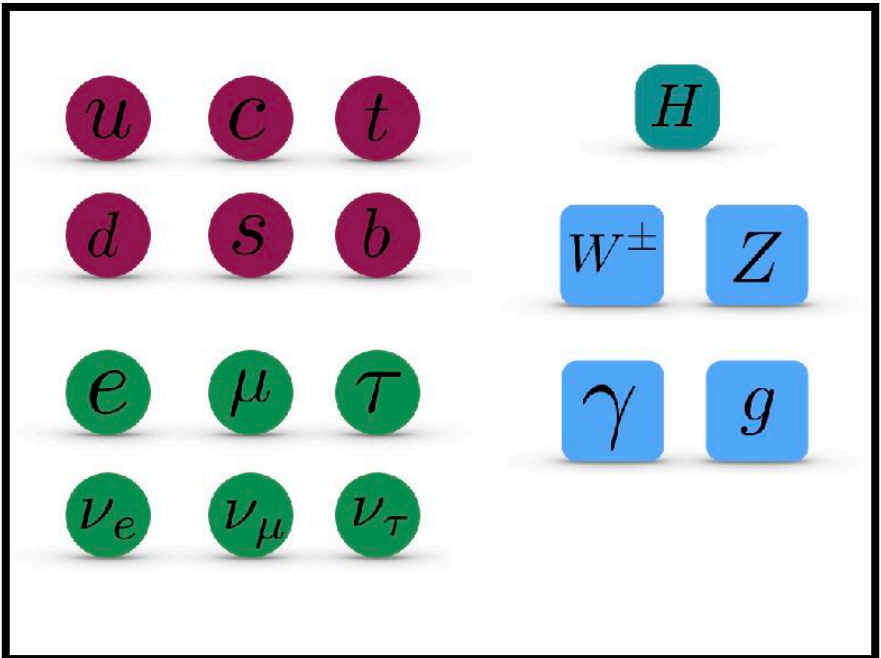
distance scale

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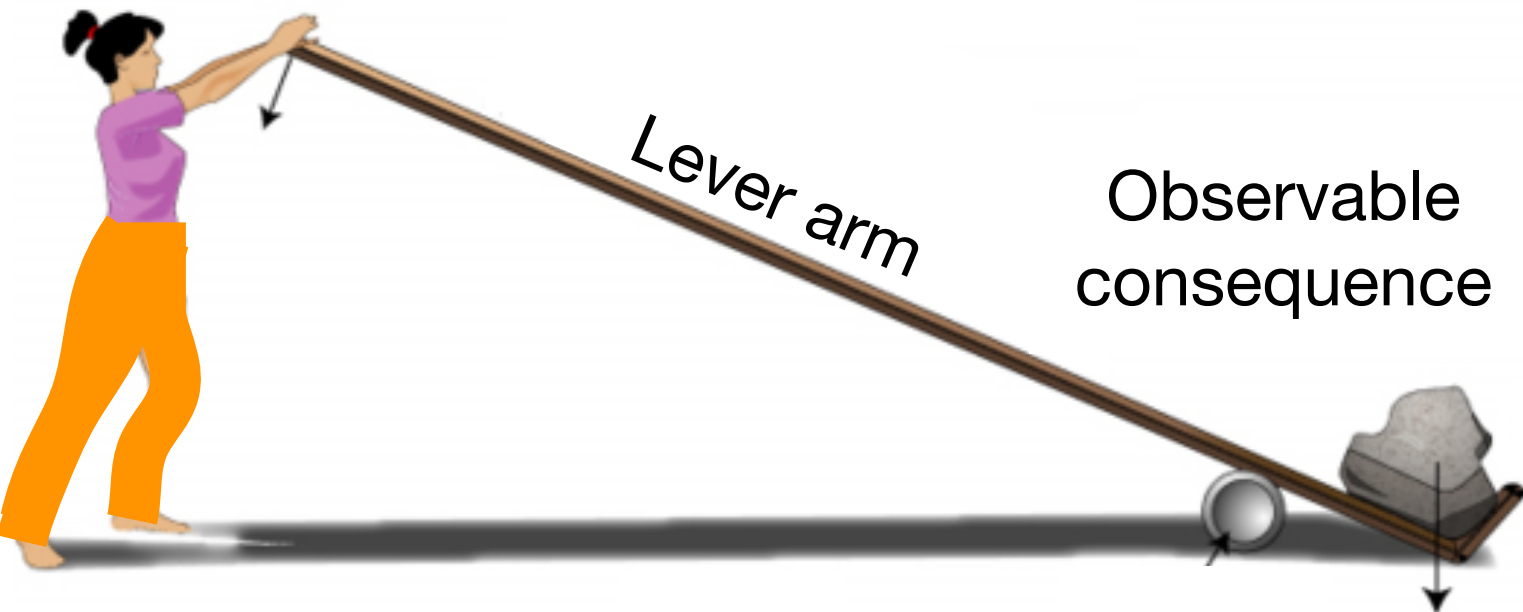
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Quantum-gravity effect

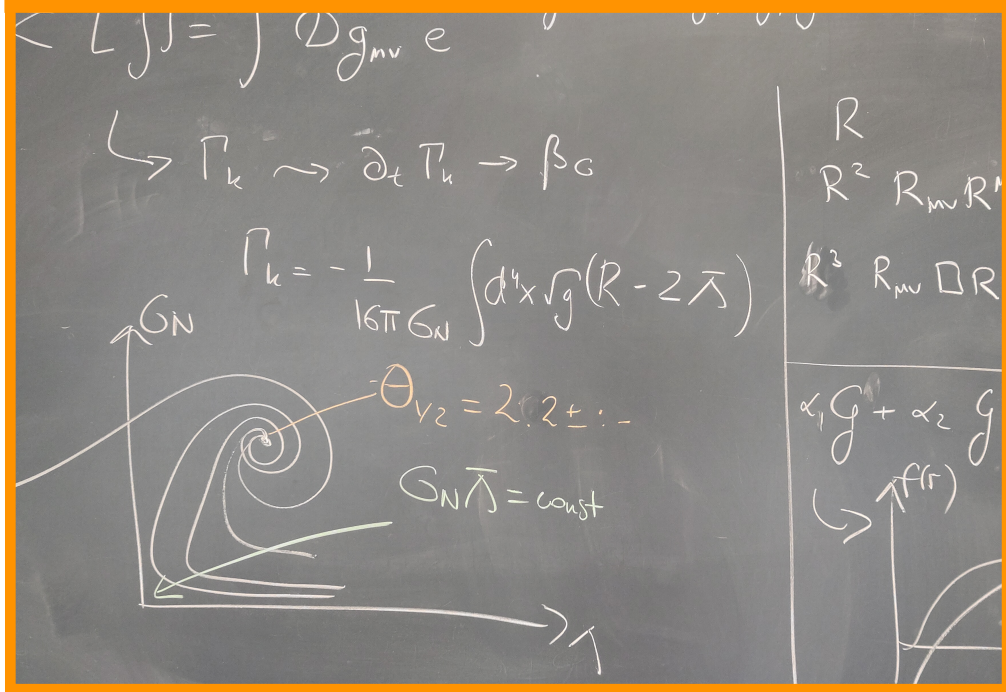


Examples:

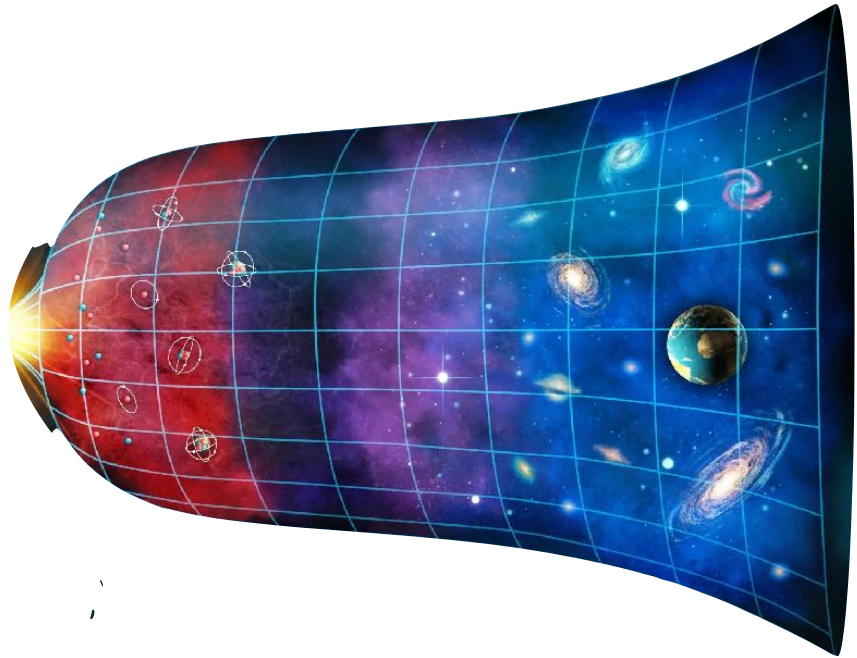
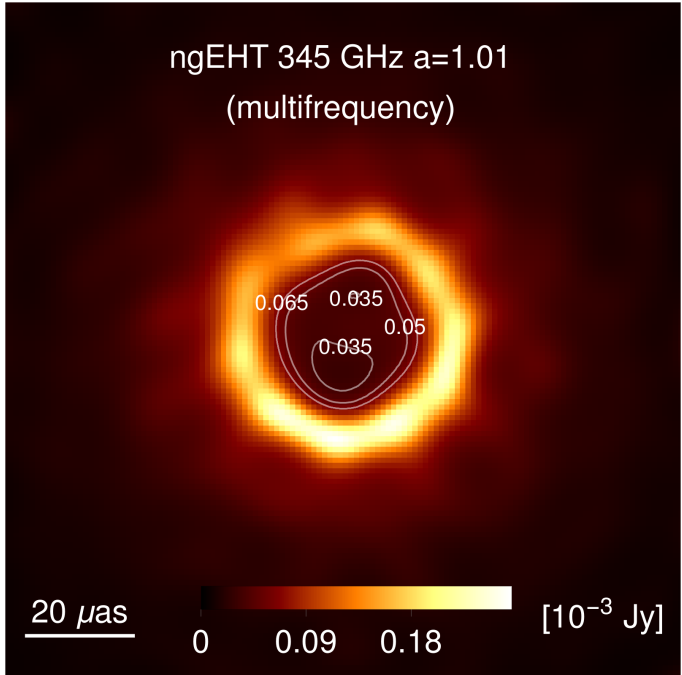
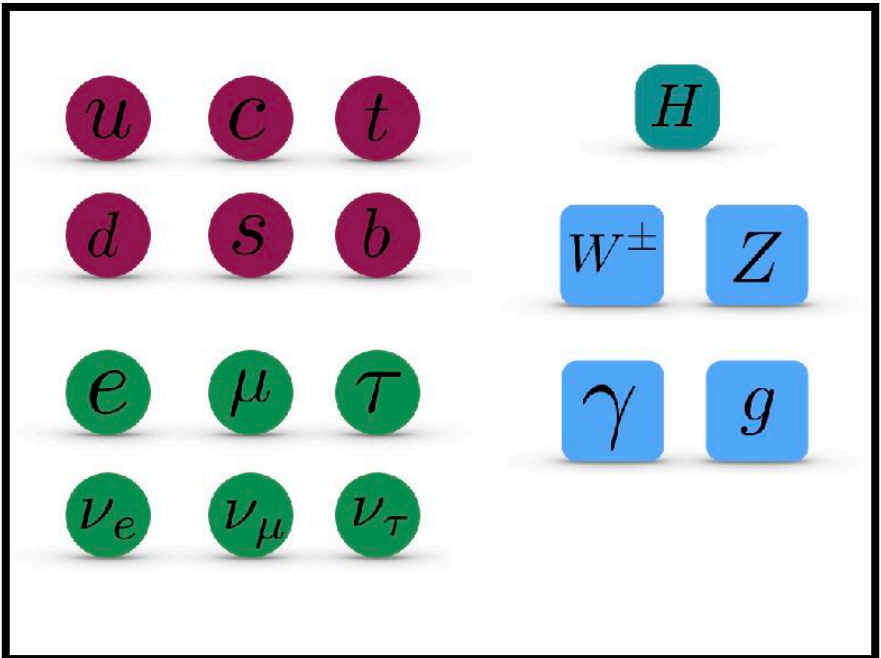
- Accumulation of Lorentz-Invariance-Violation over astrophysical distances of photons from short Gamma-Ray-Bursts
[Amelino-Camelia, Ellis, Mavromatos Nanopoulos, Sakar '97]
- High sensitivity in probes of CPT symmetry breaking in the Standard Model
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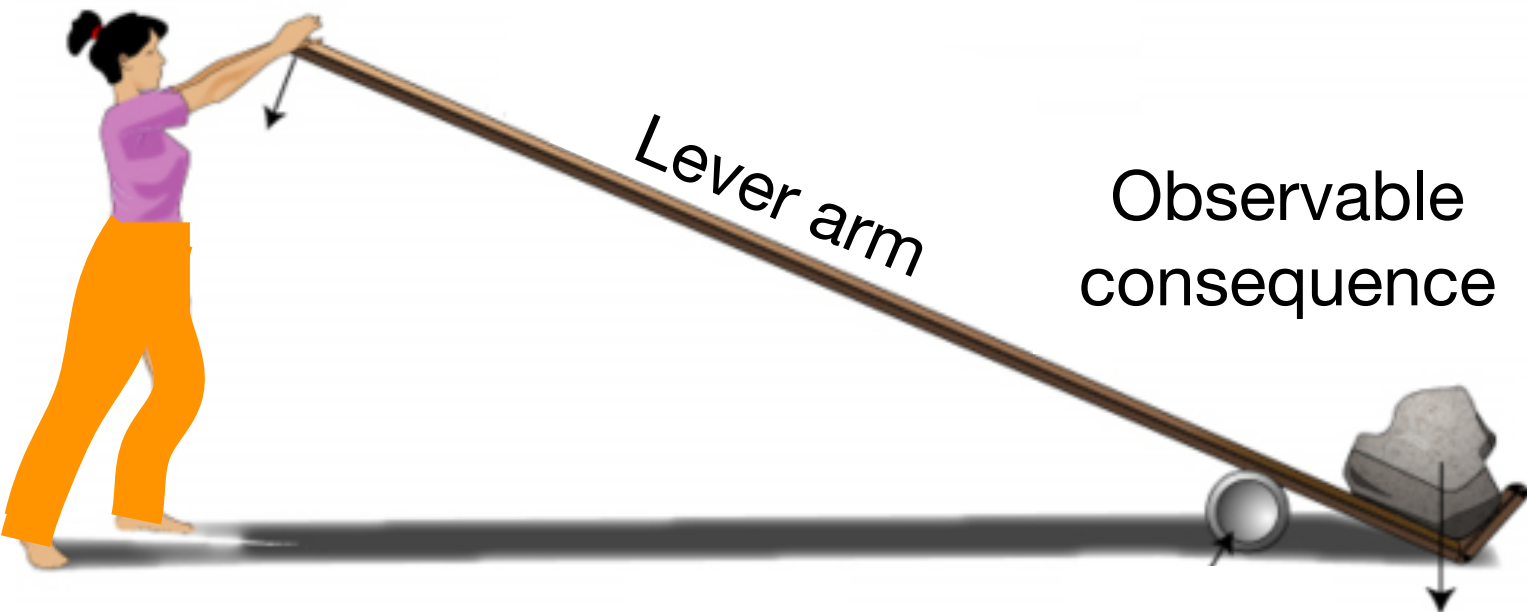
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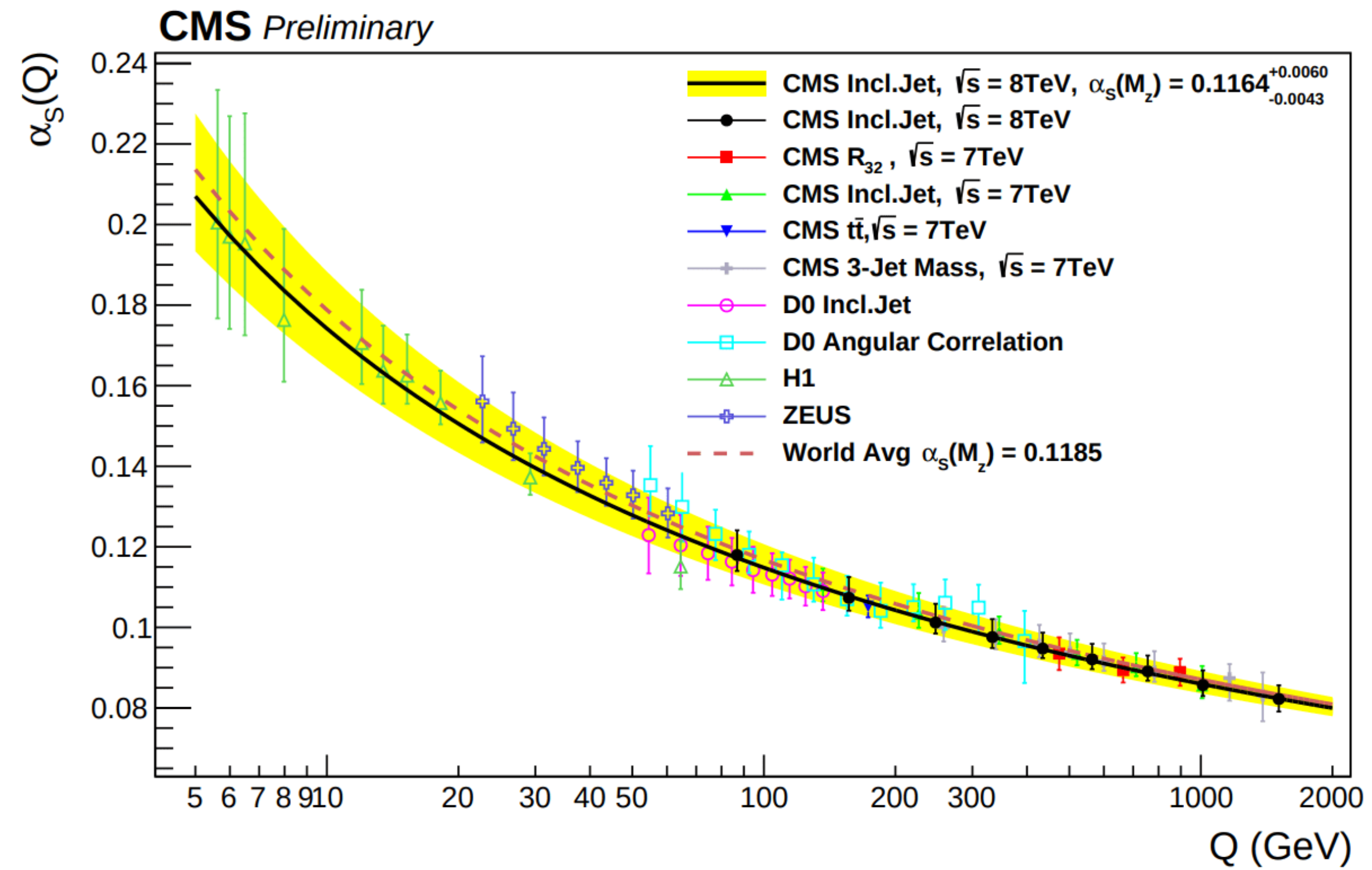


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- Renormalization Group flow of couplings

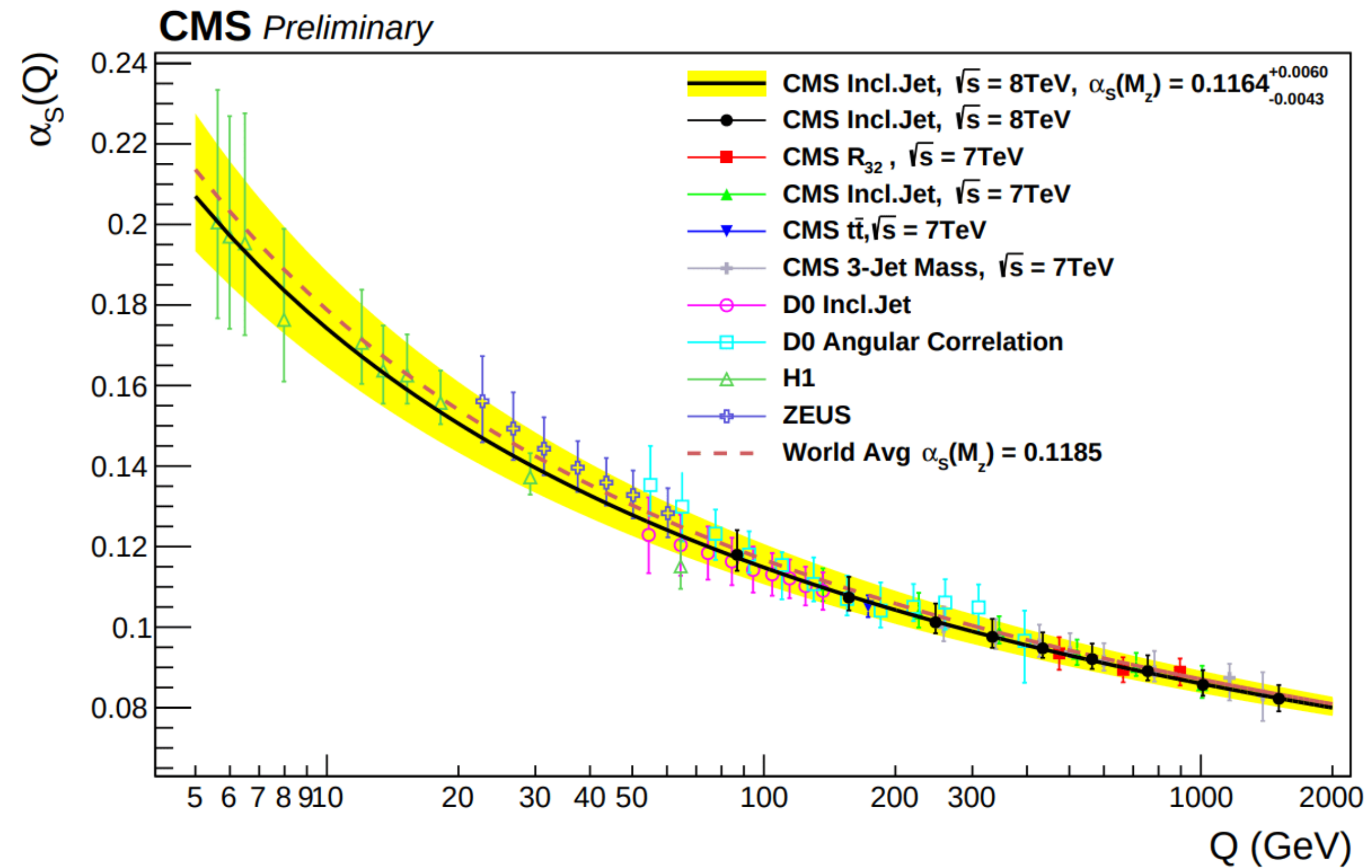
Renormalization Group flow

Coupling constants are not constant, instead they *run* as a function of scale due to quantum fluctuations

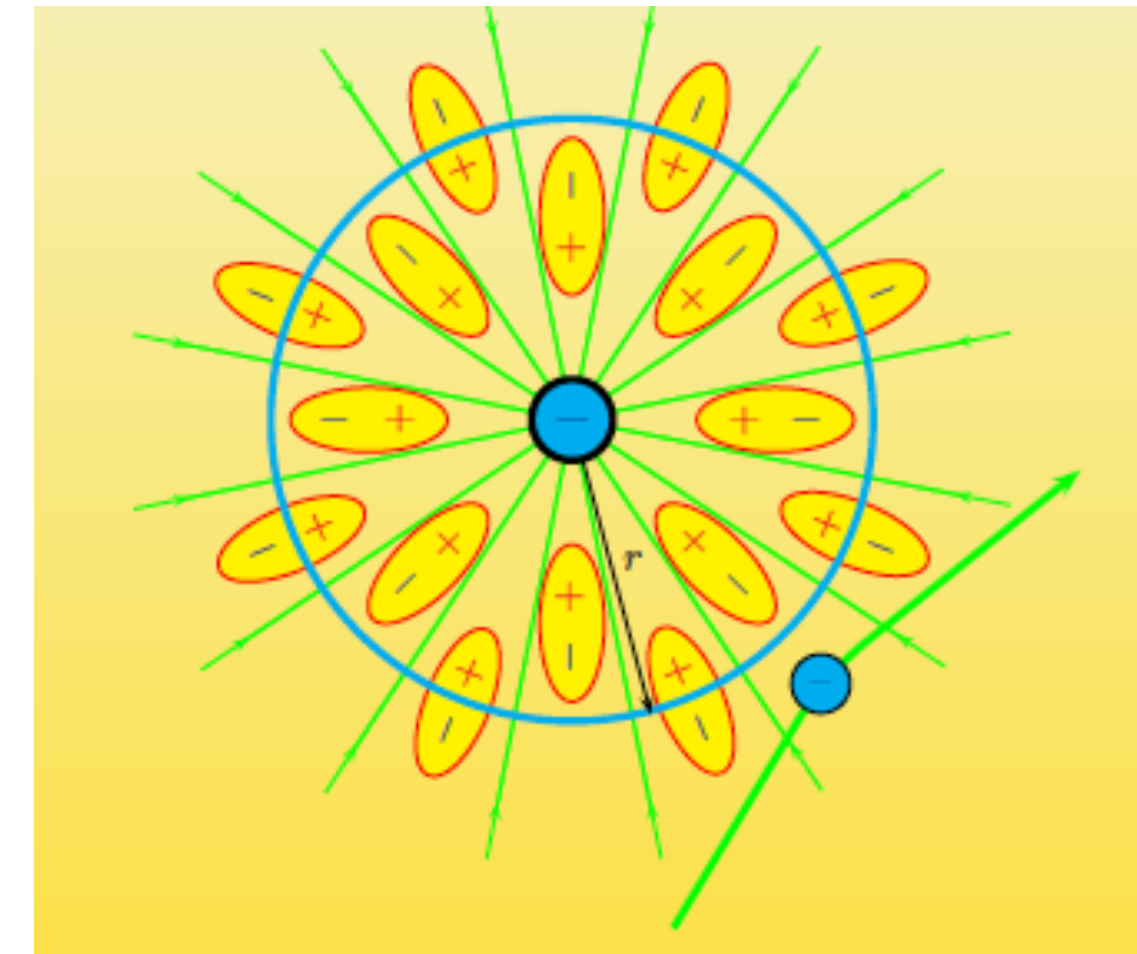


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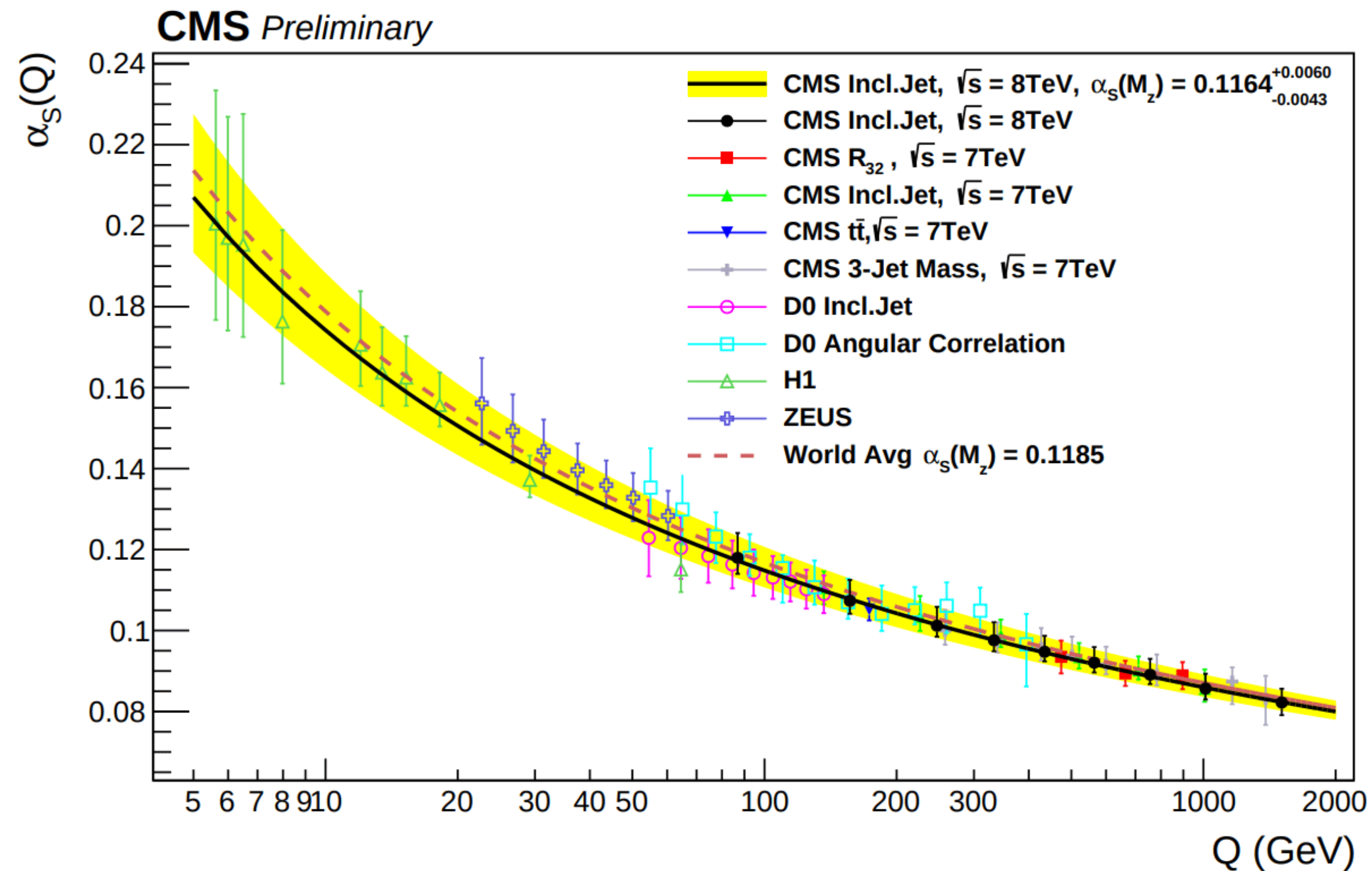


Quantum fluctuations screen or antiscreen interactions

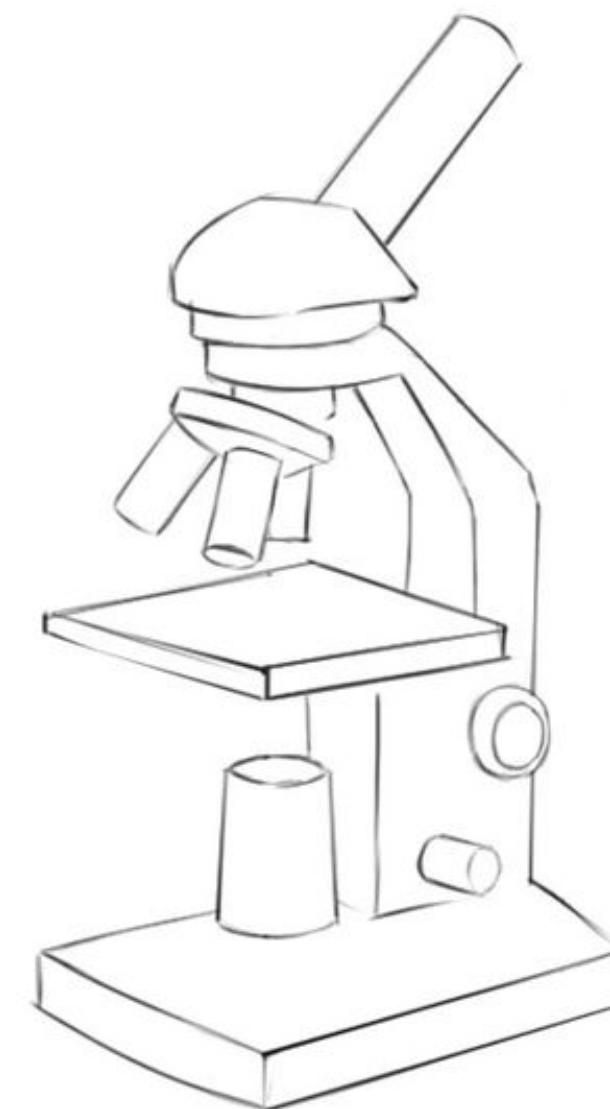
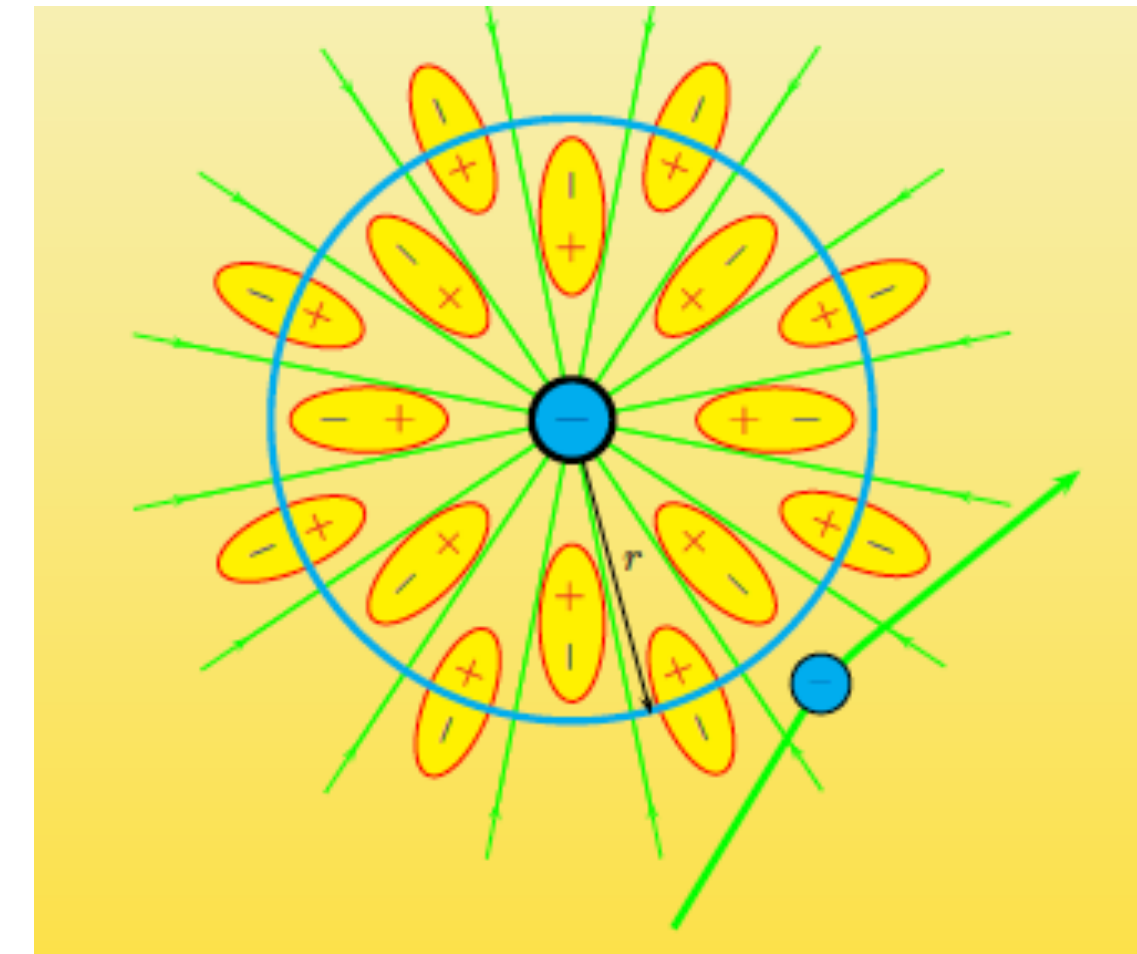


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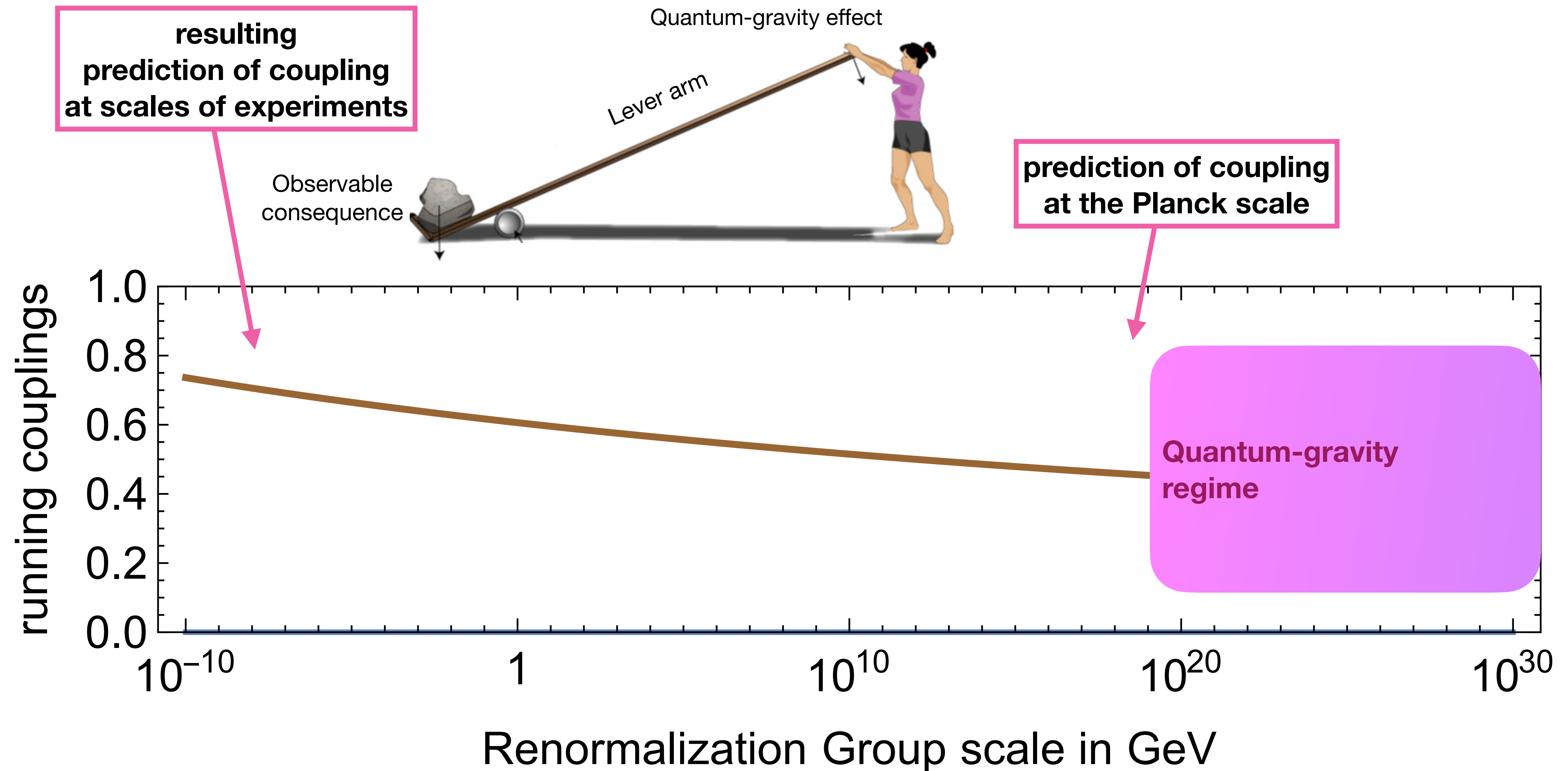
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Renormalization Group flow
= "microscope" for the theory

$\Rightarrow$ can calculate how couplings
change as function of scale

Renormalization Group flow as a lever arm

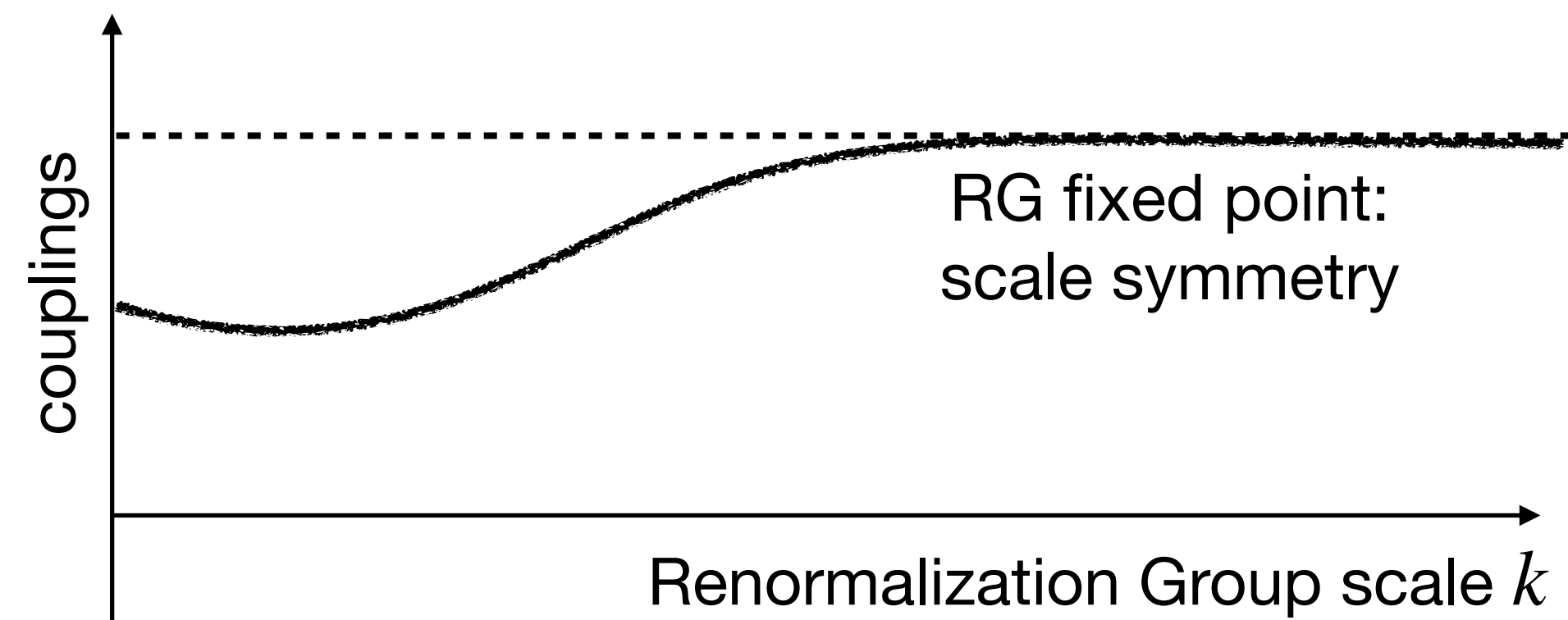


Key idea of asymptotic safety

Asymptotic safety = quantum scale symmetry

scale symmetry realized in transplanckian regime

signature: Renormalization Group fixed point



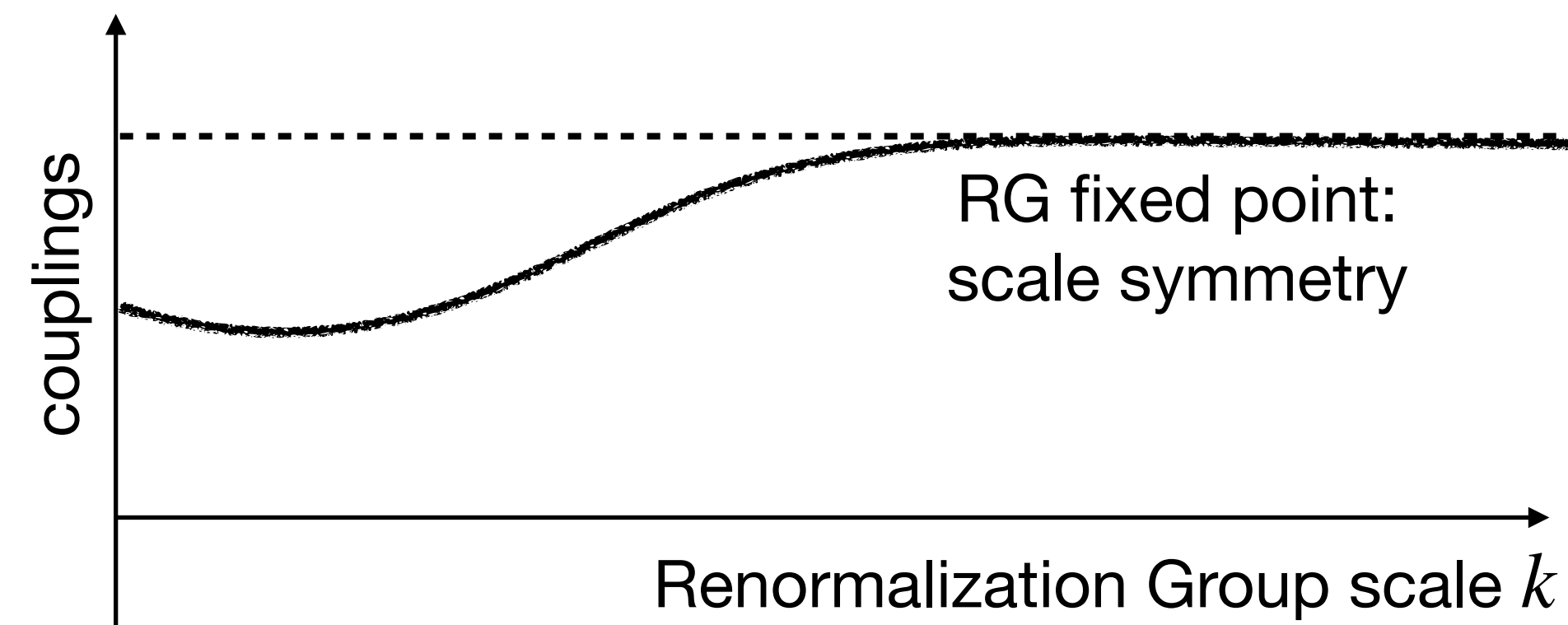
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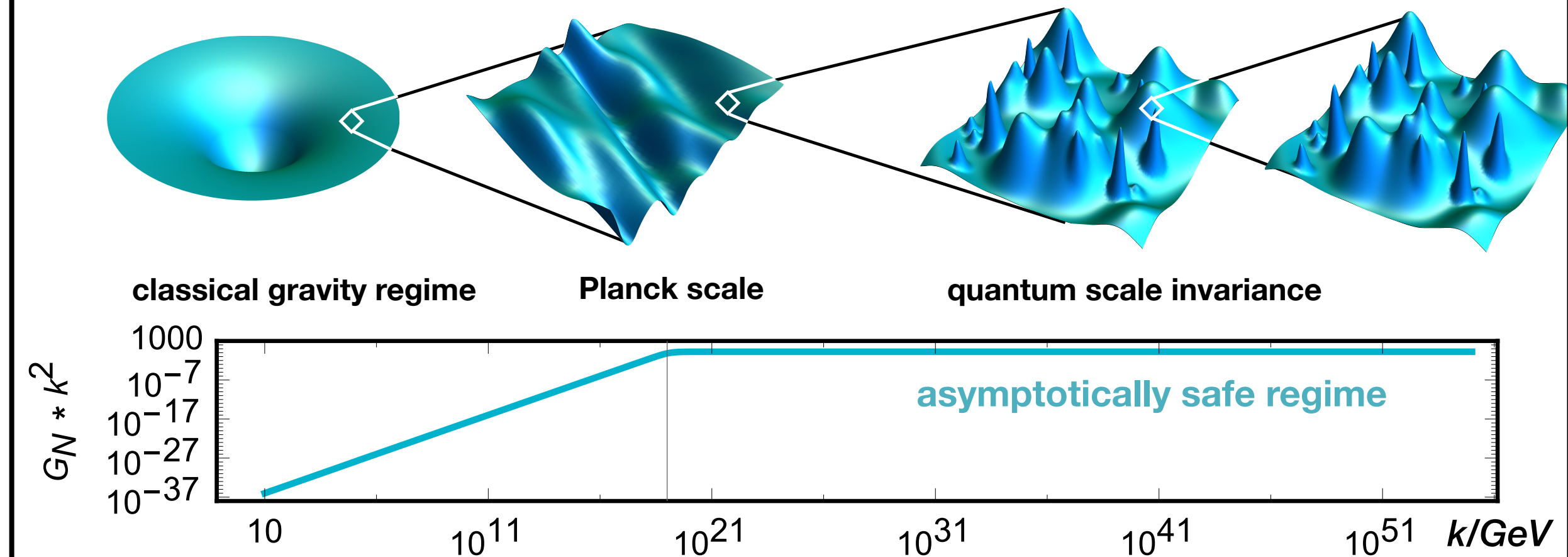
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scale symmetry = deep principle of nature

Asymptotic safety in gravity: [Weinberg '79; Reuter '96]

$$S = \int d^4x \sqrt{-g} \left(\Lambda(k) + \frac{1}{16\pi G_N(k)} R + a(k) R^2 + b(k) R_{\mu\nu} R^{\mu\nu} + \dots \right)$$

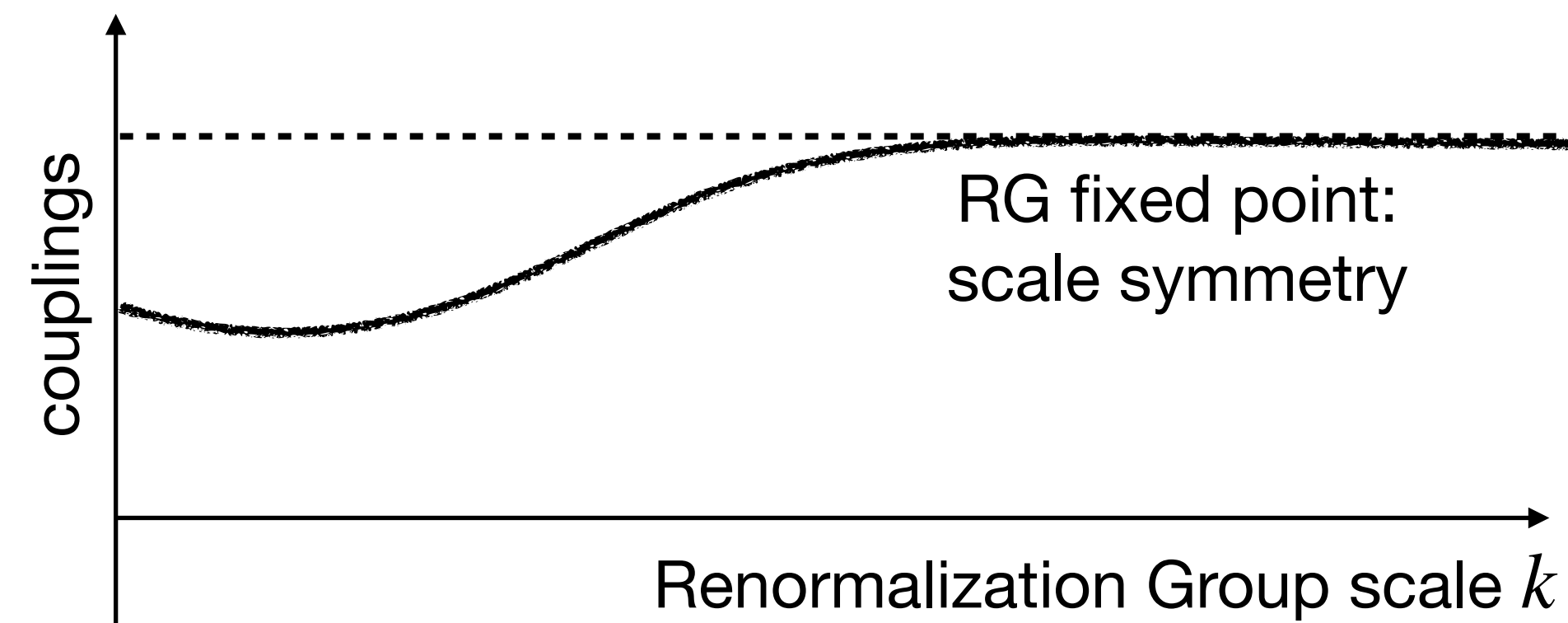


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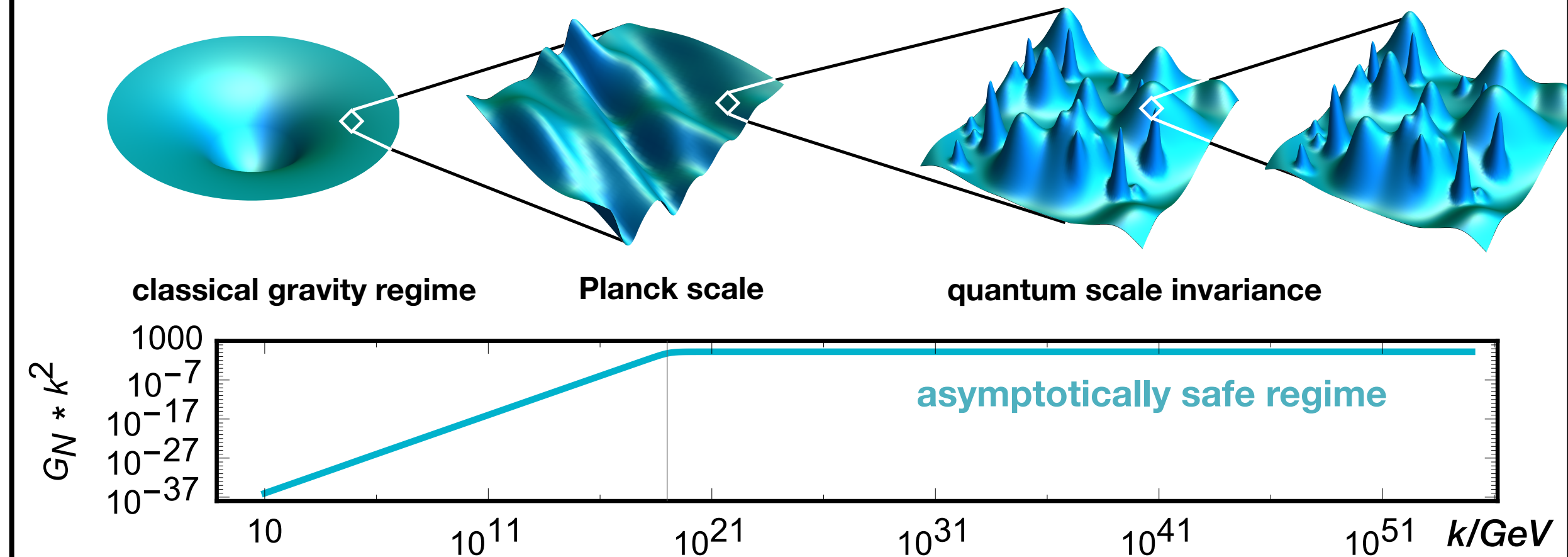
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Theoretical evidence for asymptotic safety in gravity:

Fixed point regarded as established (hundreds of papers)

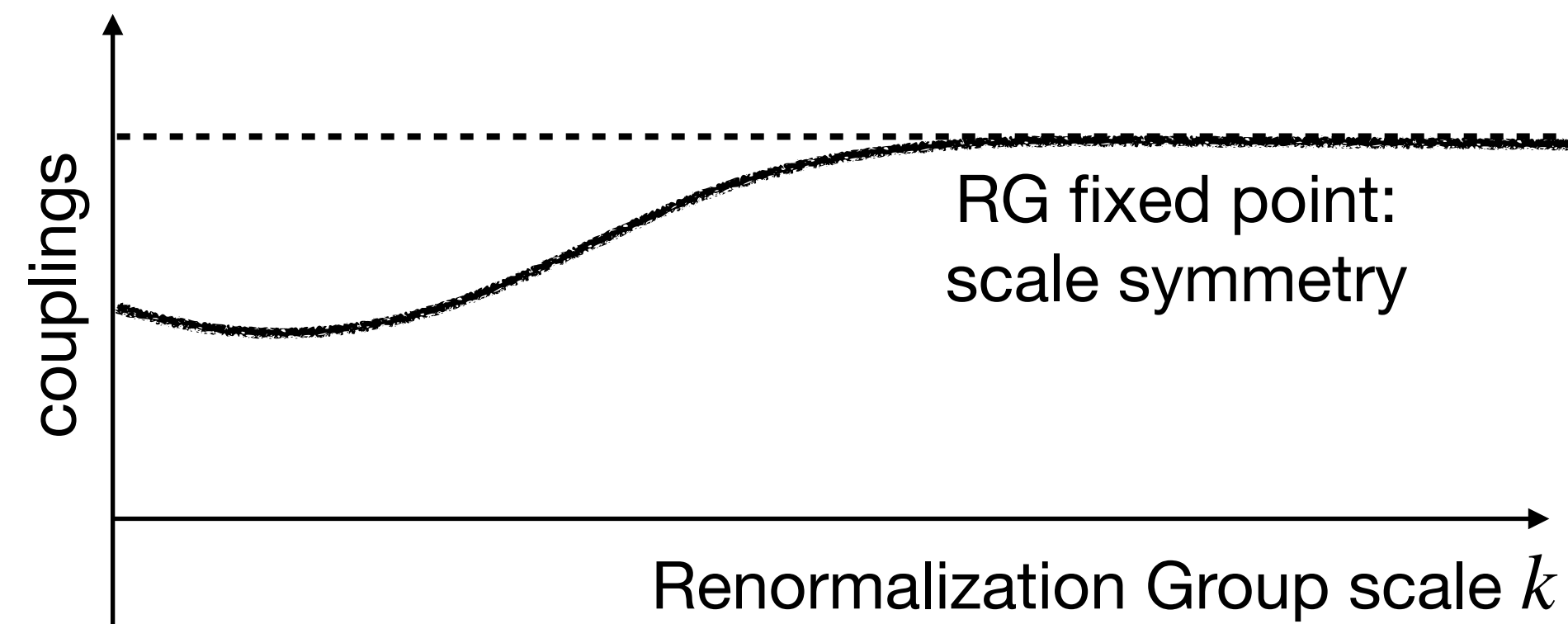
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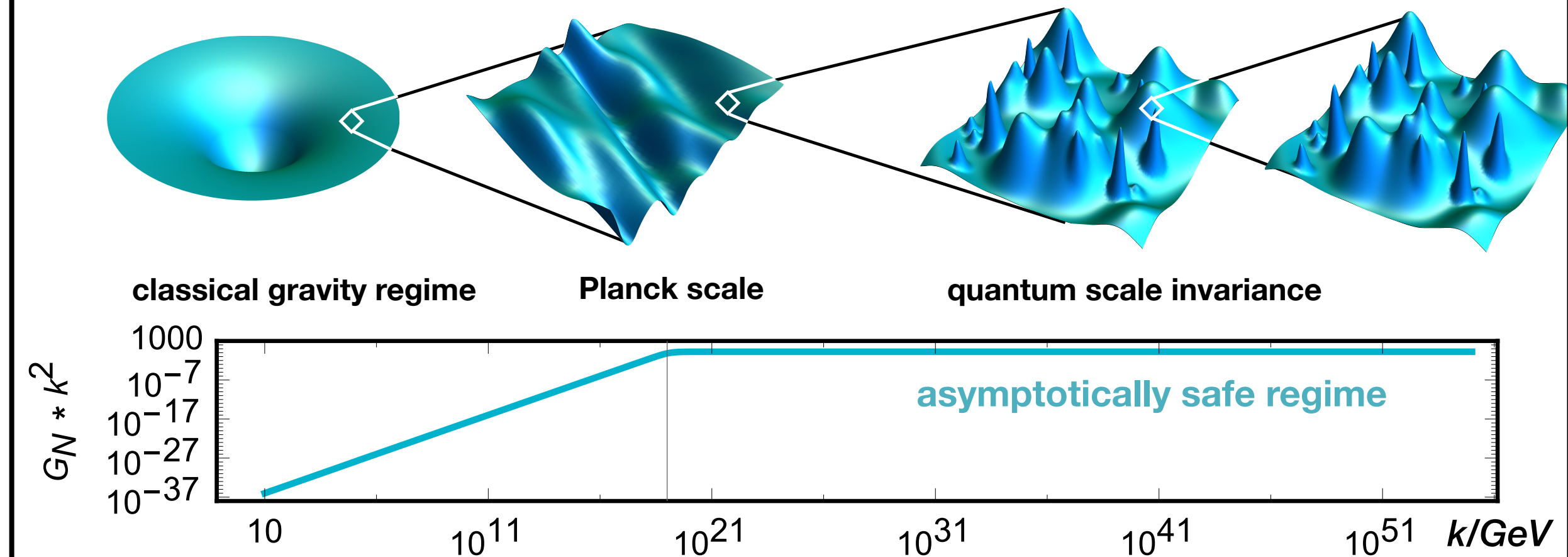
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Open questions: Lorentzian signature, unitarity under investigation

[e.g., Fehre, Litim, Pawłowski, Reichert '21; Platania '22; Saueressig, Wang '23; d'Angelo '23; Pastor-Guterres, Pawłowski, Reichert, Risi '24; d'Angelo, Ferrero, Fröb '25]

Predictive power in asymptotic safety

Origin of predictions at the Planck scale

Quantum fluctuations

screen or antiscreen interactions, e.g.,

$$\text{QED: } \beta_e = k \partial_k e(k) = \frac{1}{12\pi^2} e^3 + \dots$$

→ $e(k)$ decreases as k is lowered

$$\text{QCD: } \beta_g = k \partial_k g(k) = -\frac{7}{16\pi^2} g^3 + \dots$$

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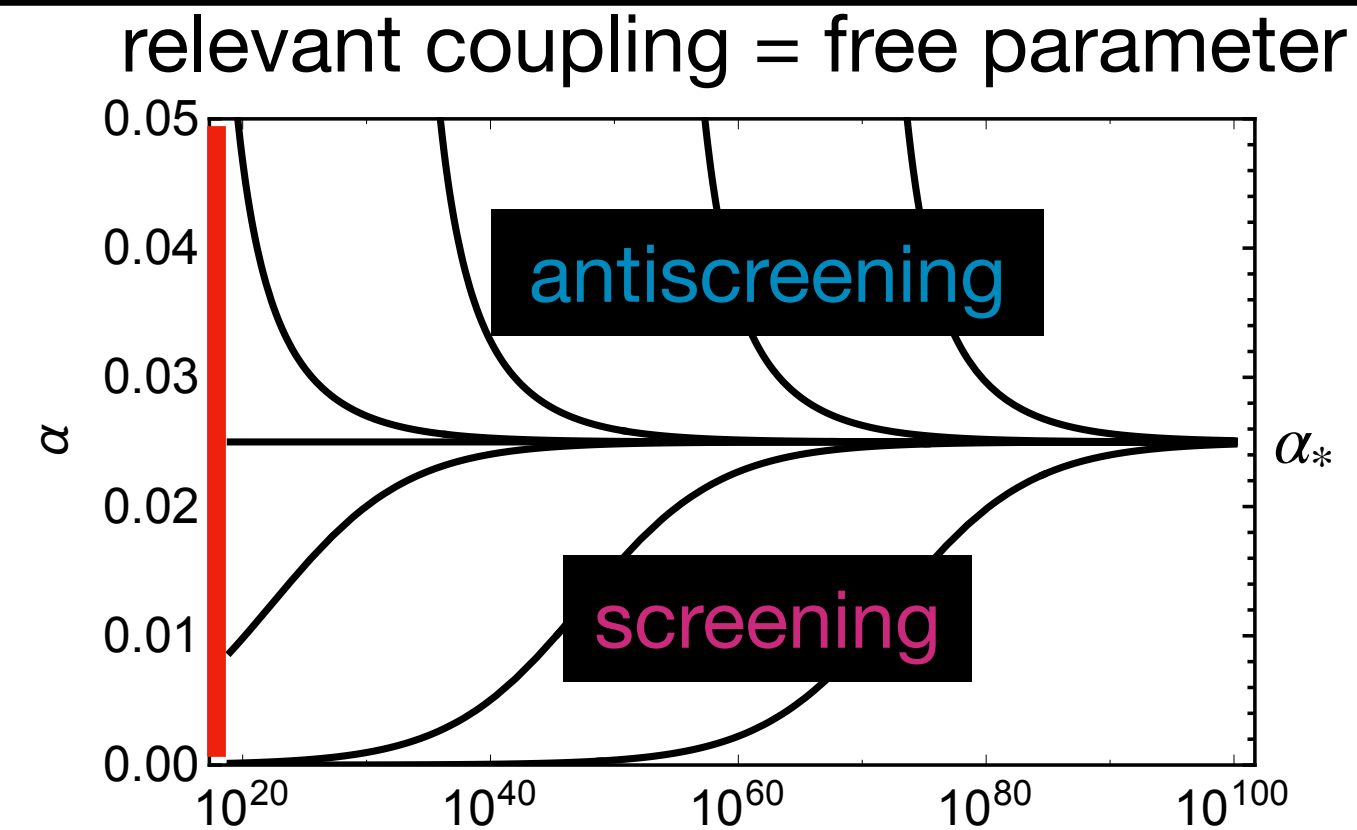
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quantum fluctuations drive coupling **away**
from scale symmetry

→ a range of coupling values
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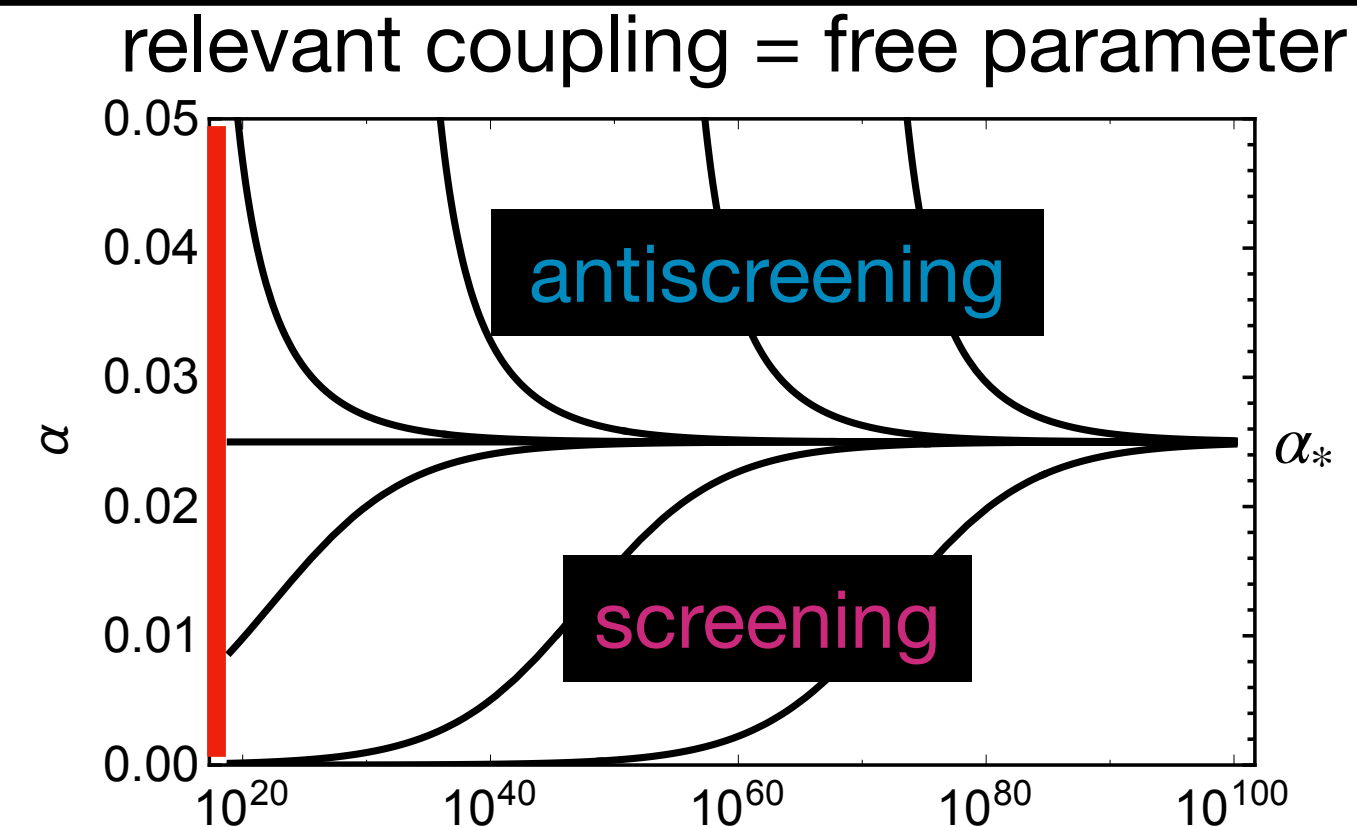
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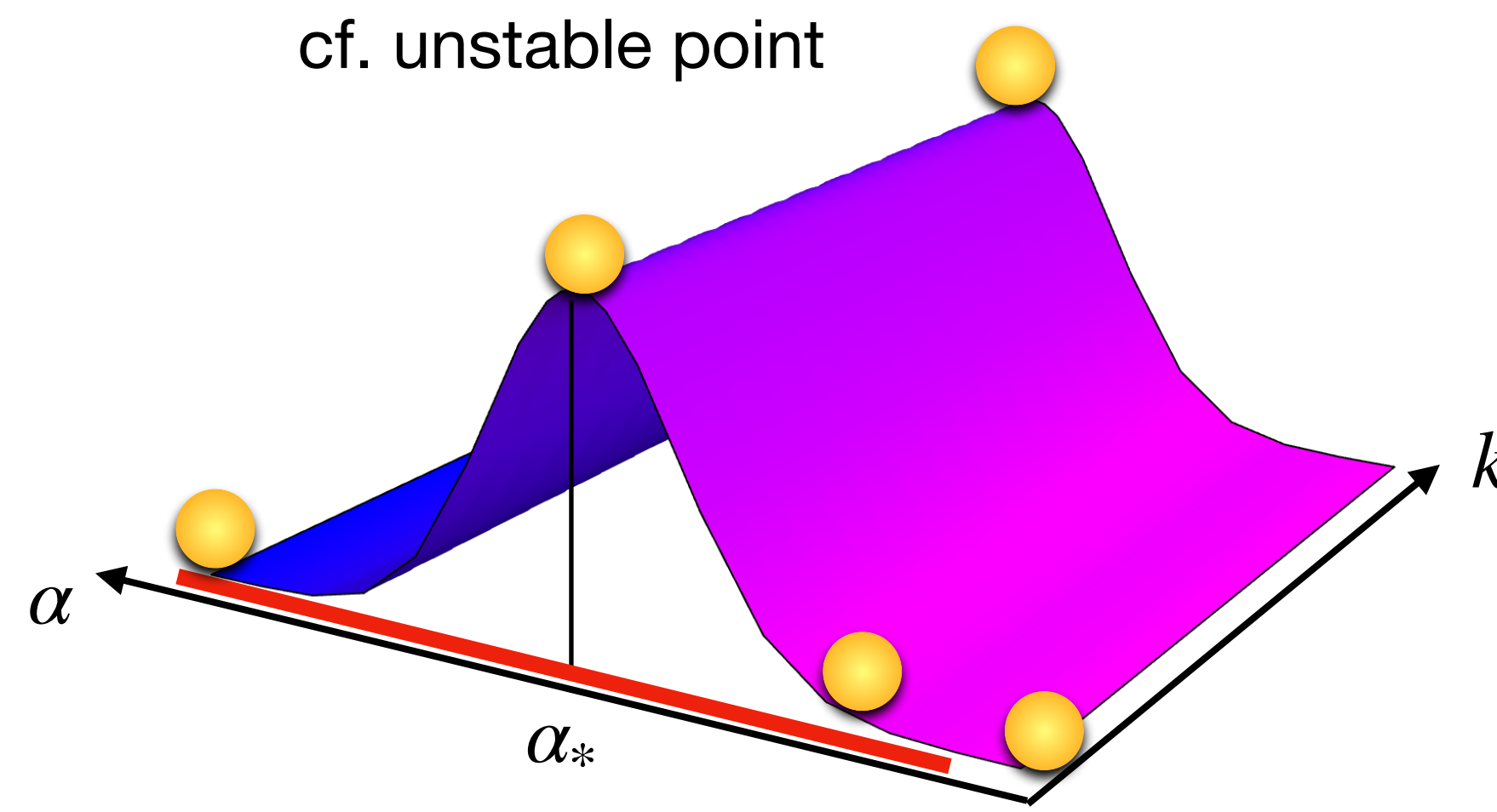
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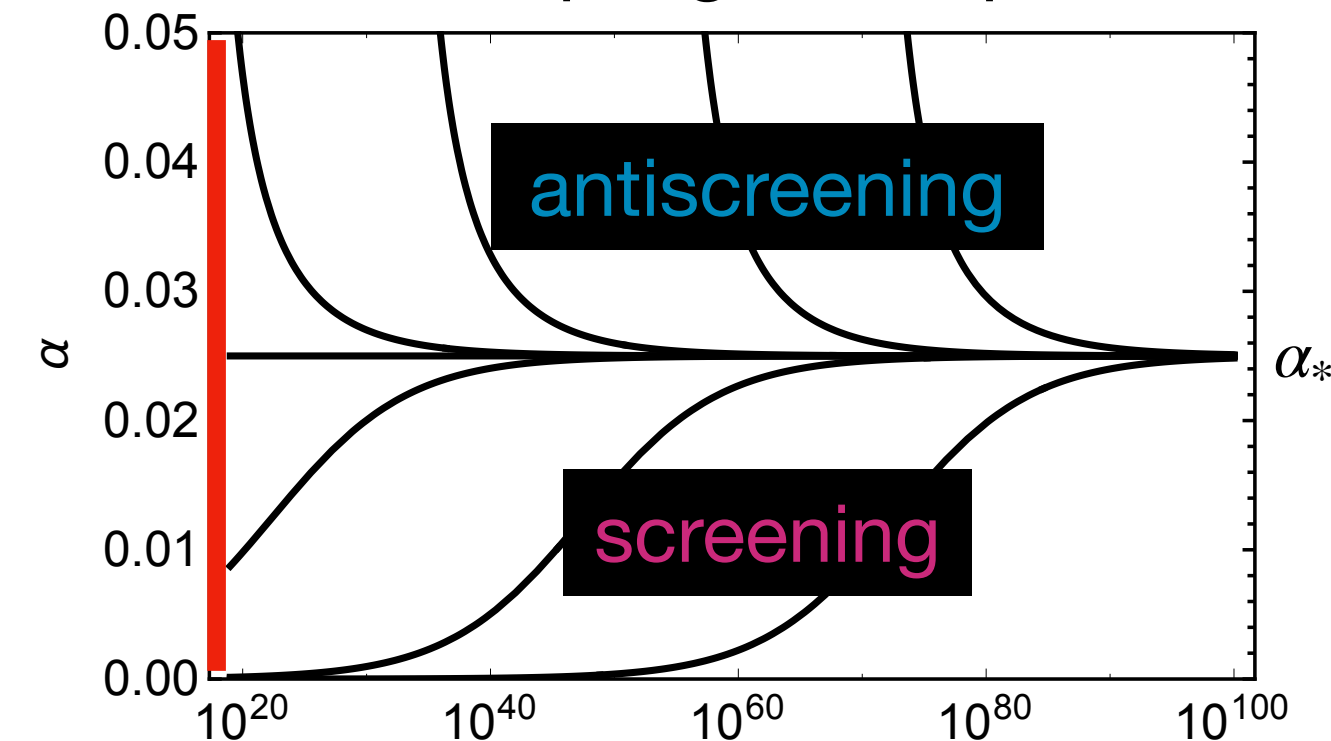
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relevant coupling = free parameter

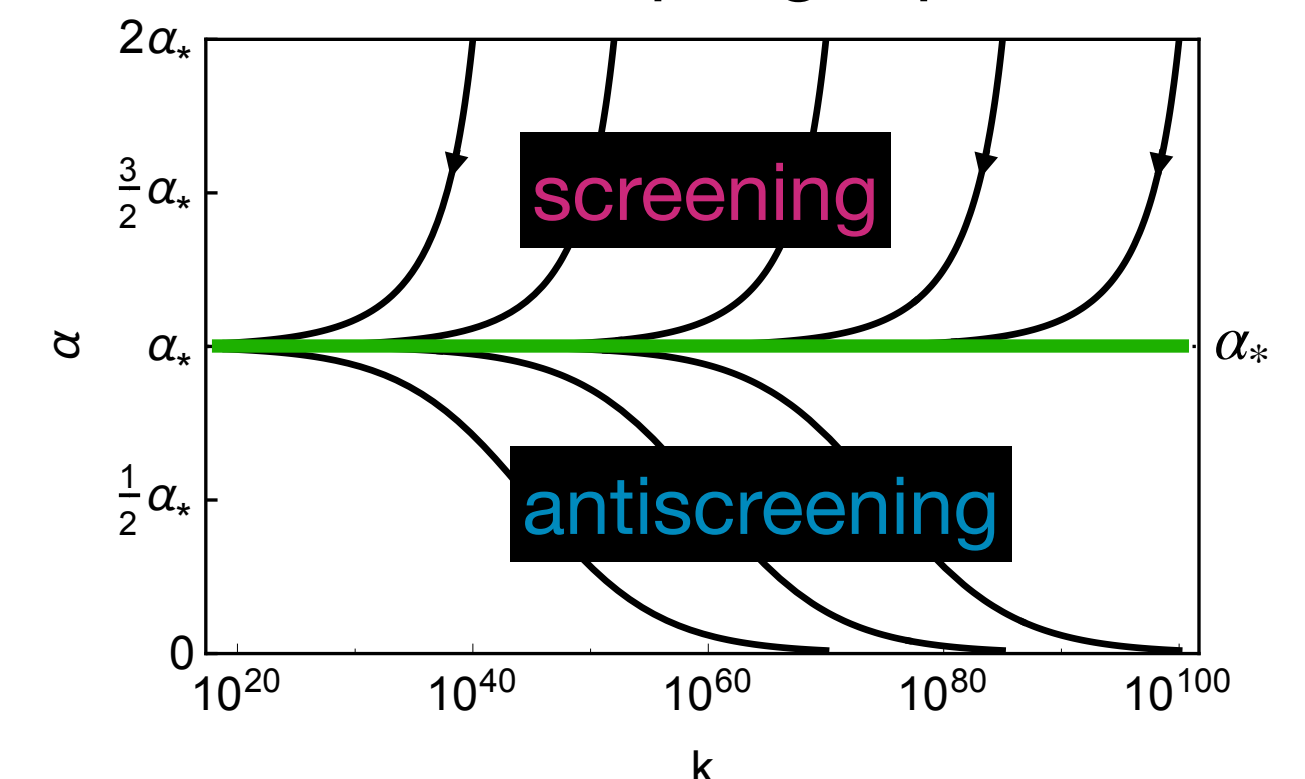


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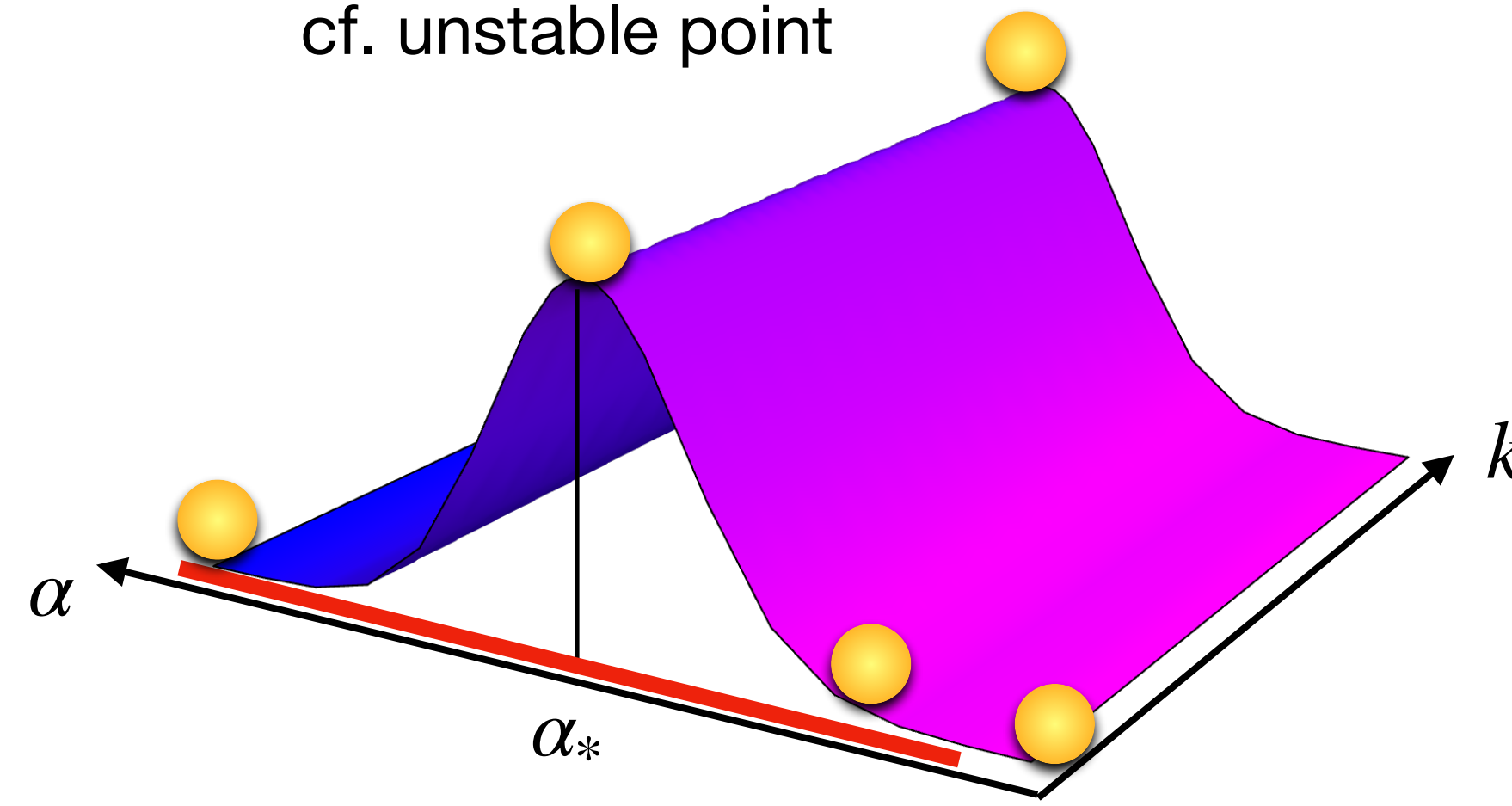


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quantum fluctuations drive coupling **towards** scale symmetry

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cf. unstable point



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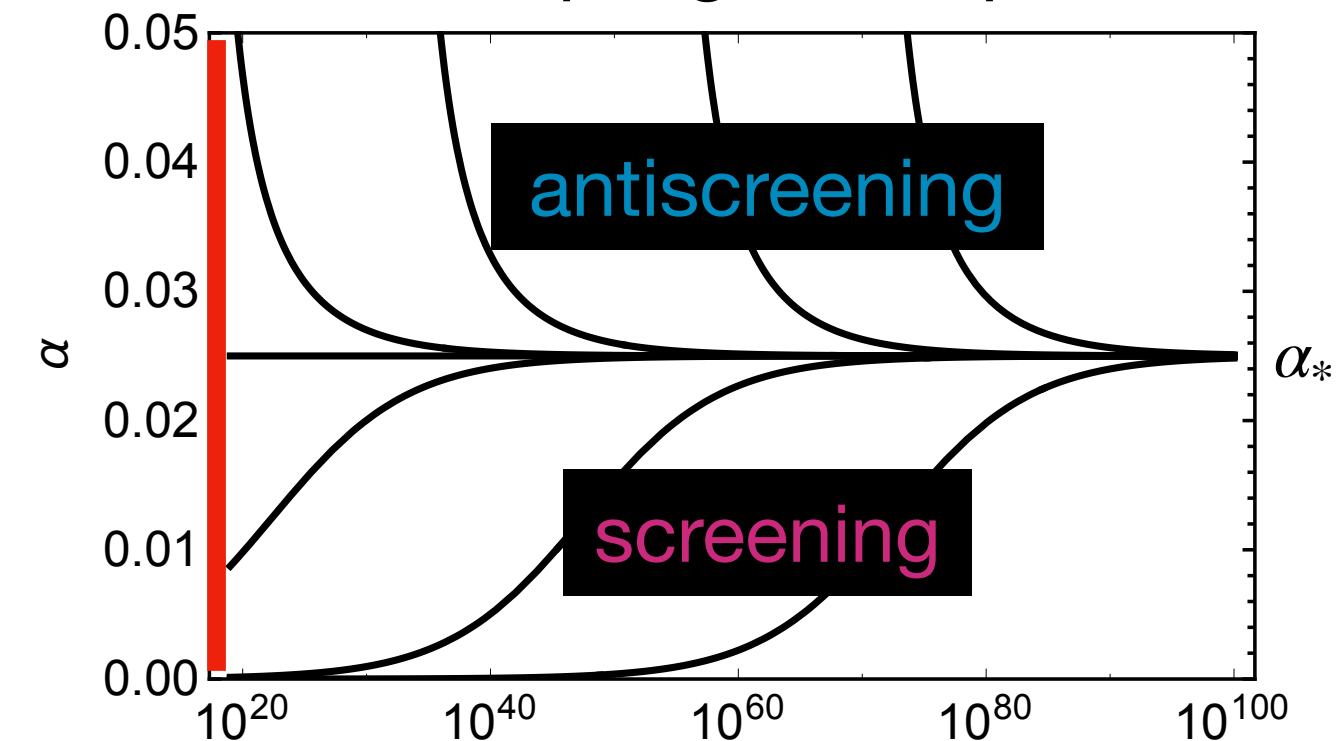
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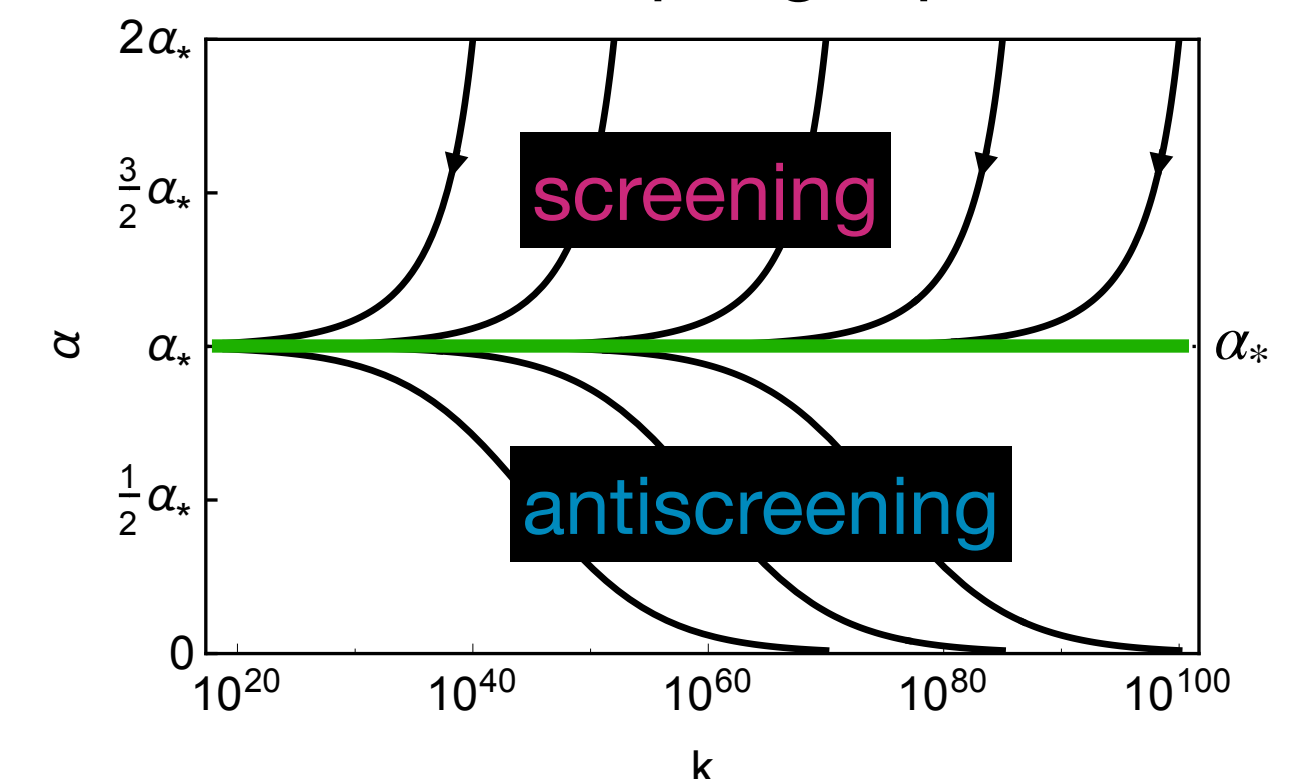


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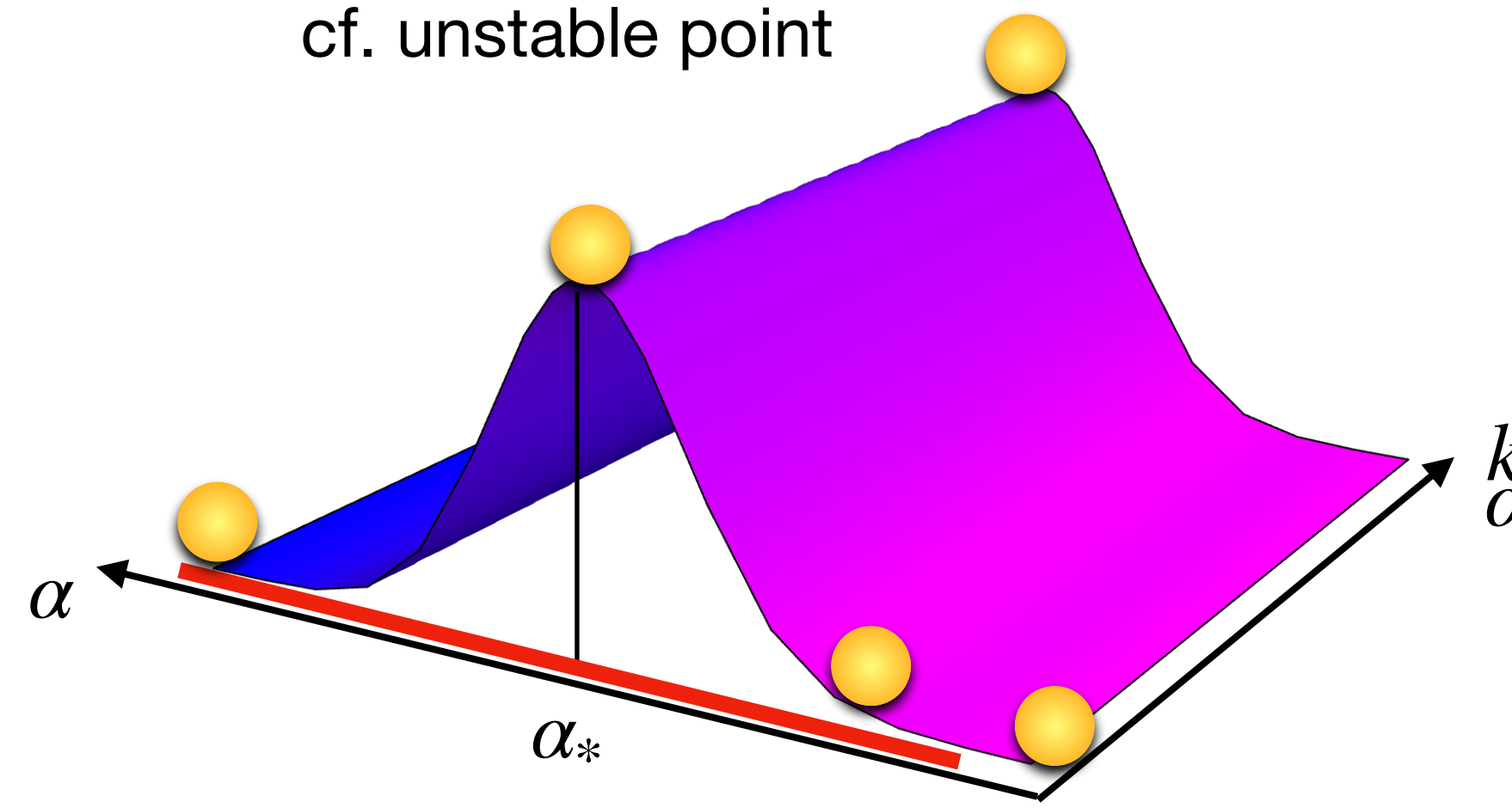


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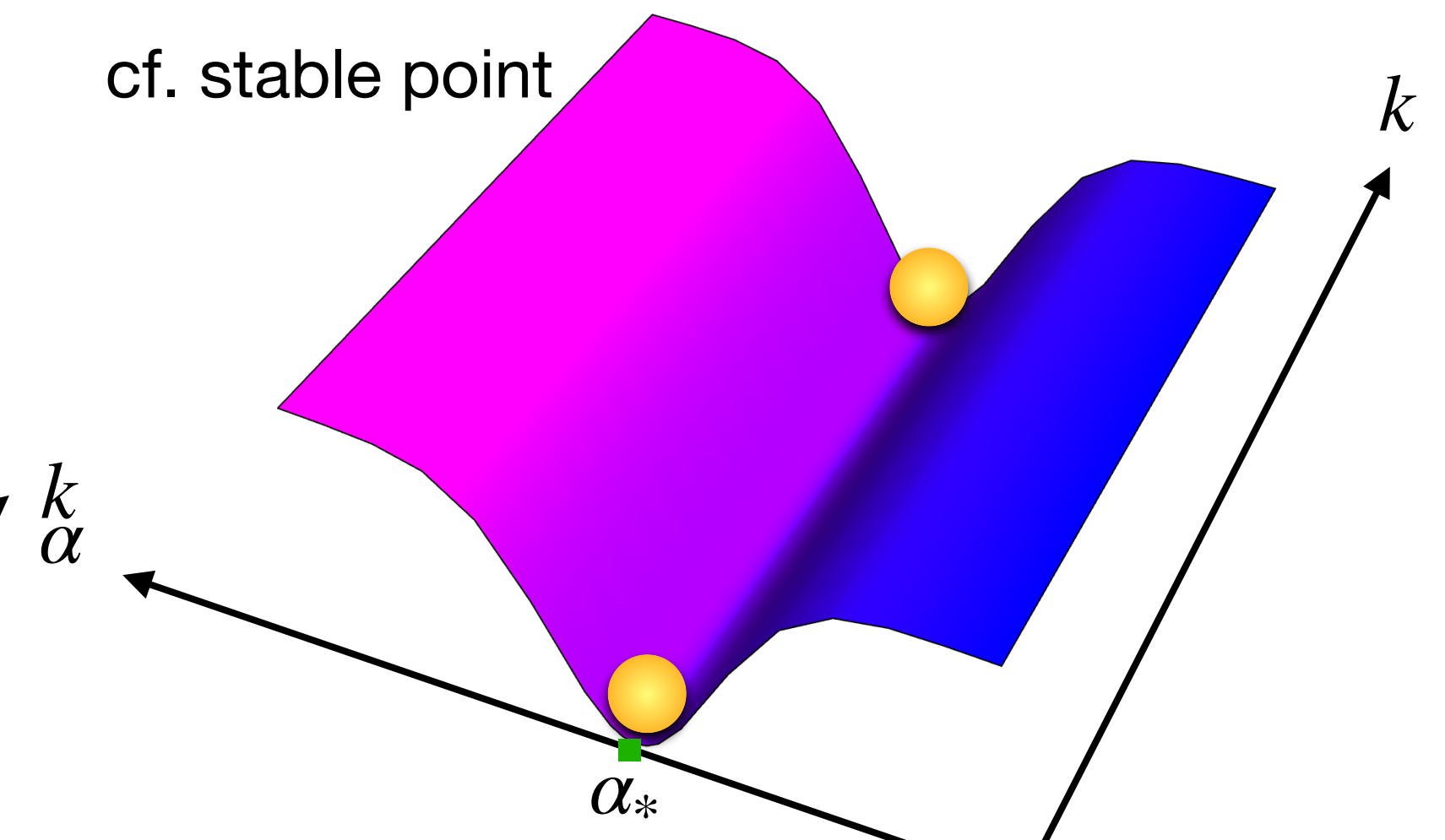
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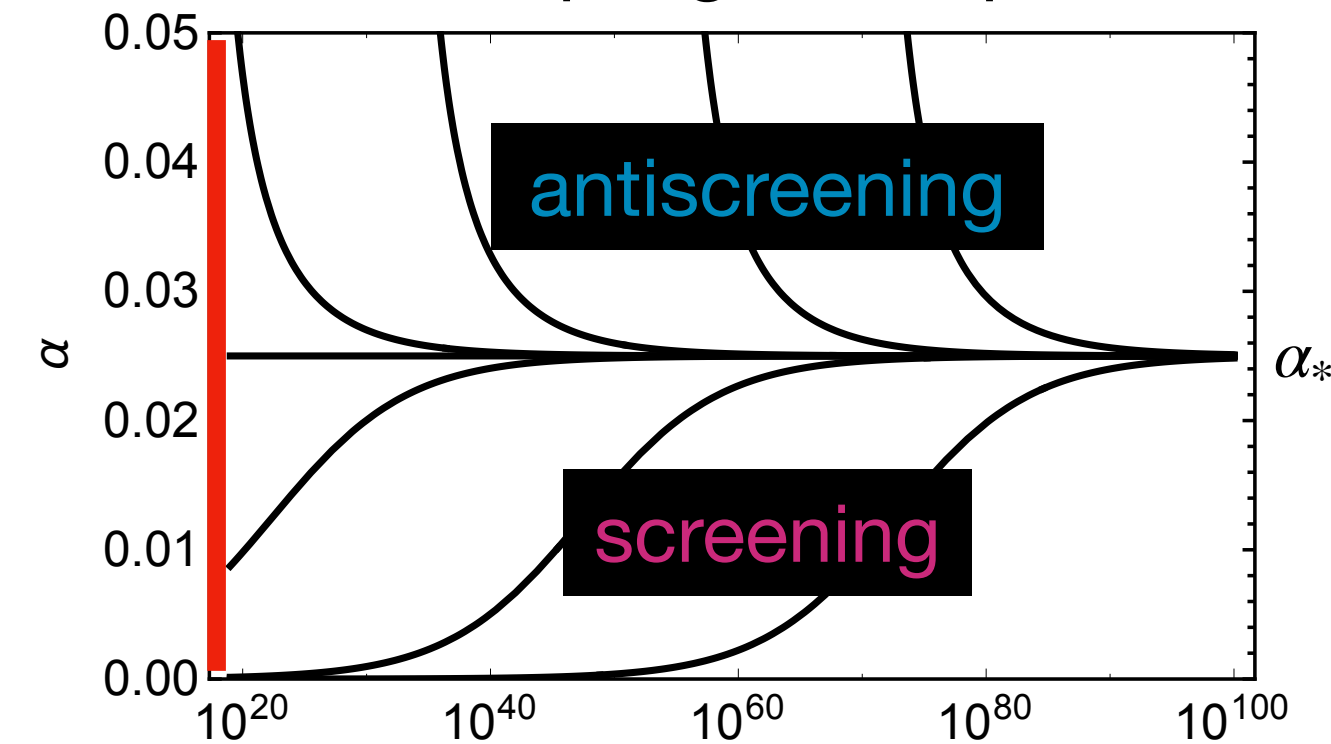
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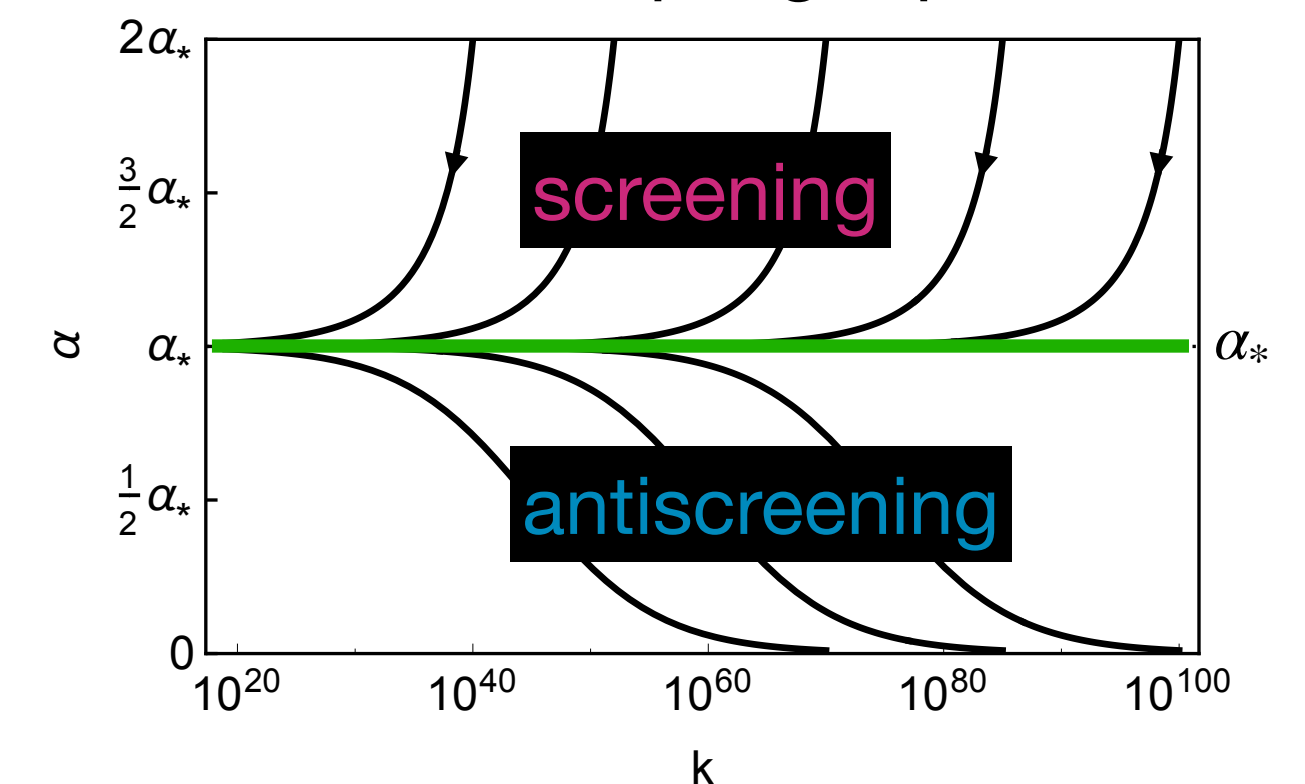


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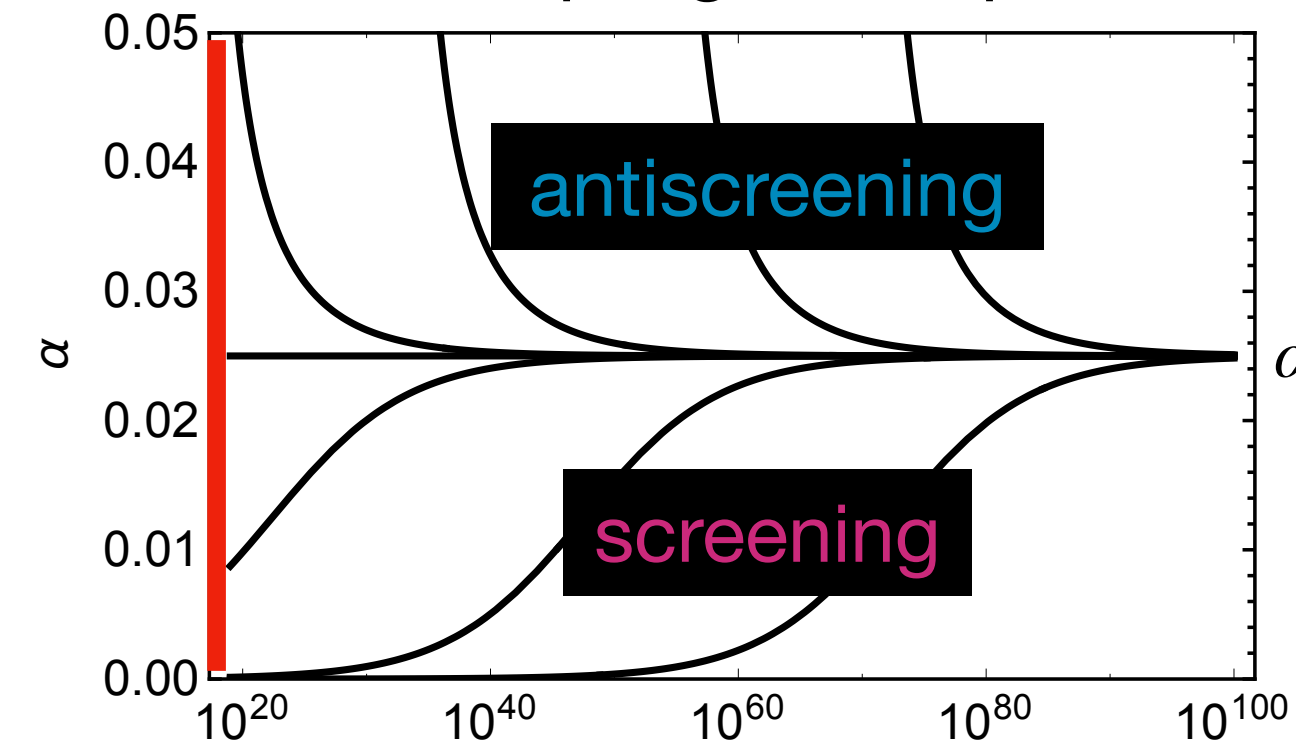
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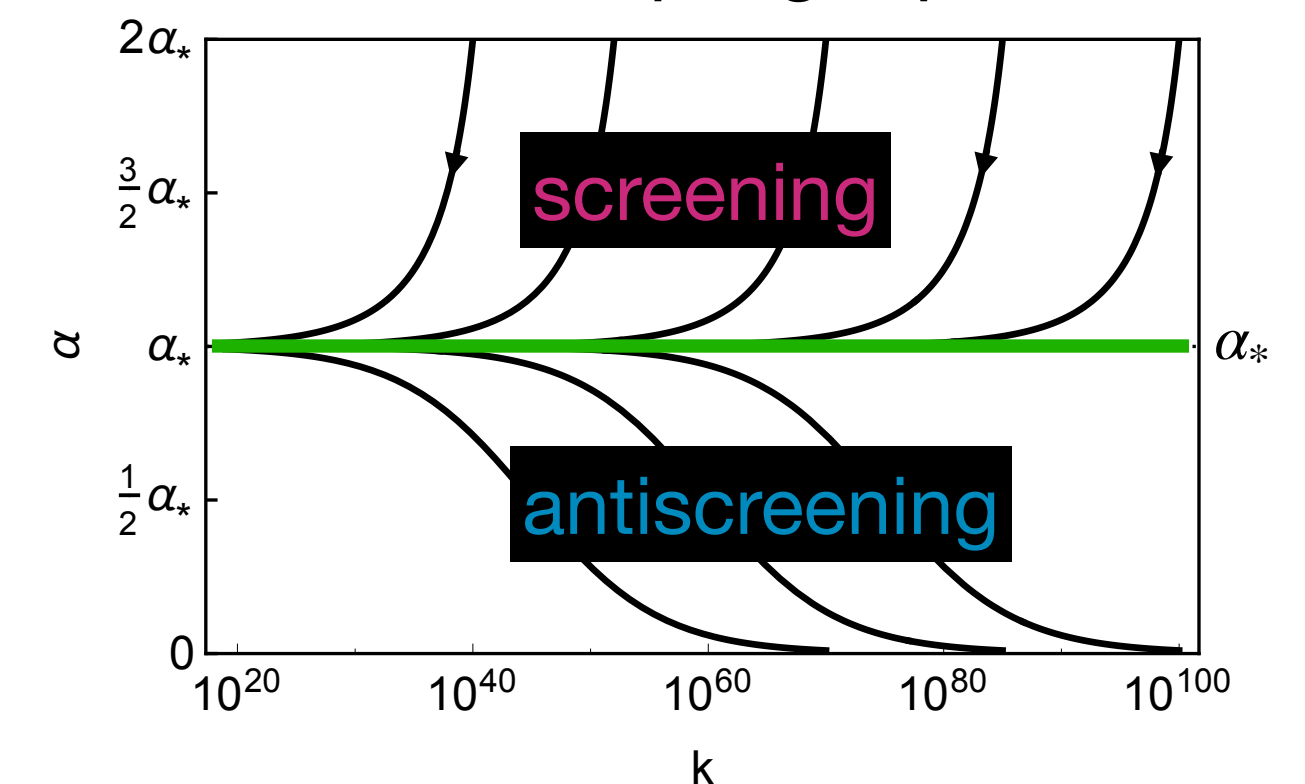


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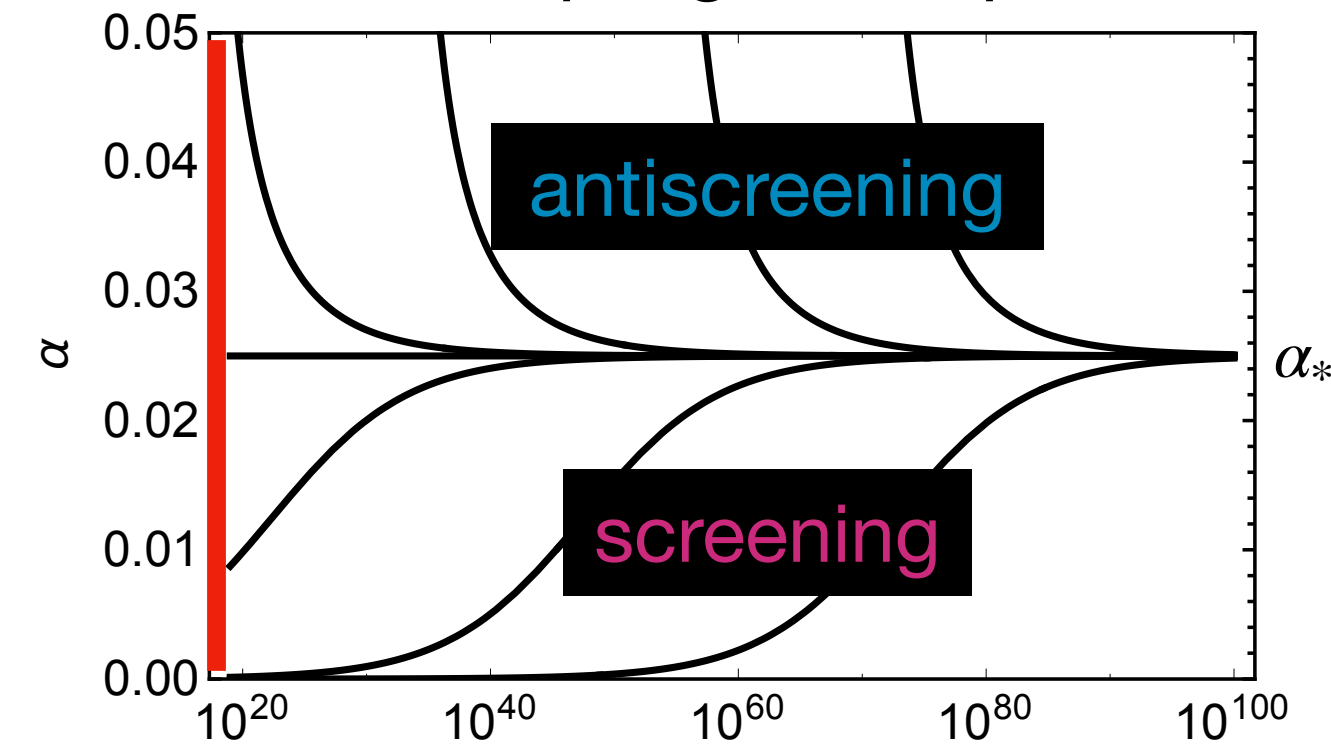
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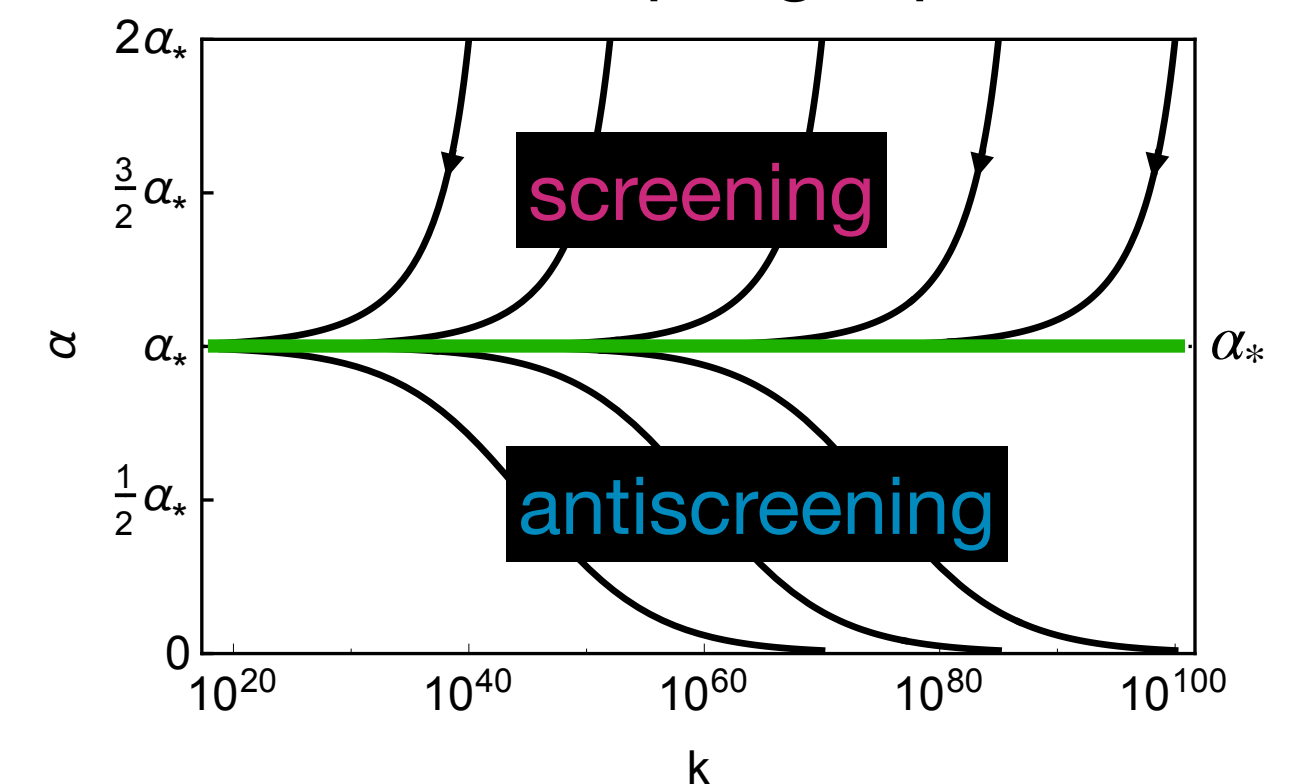


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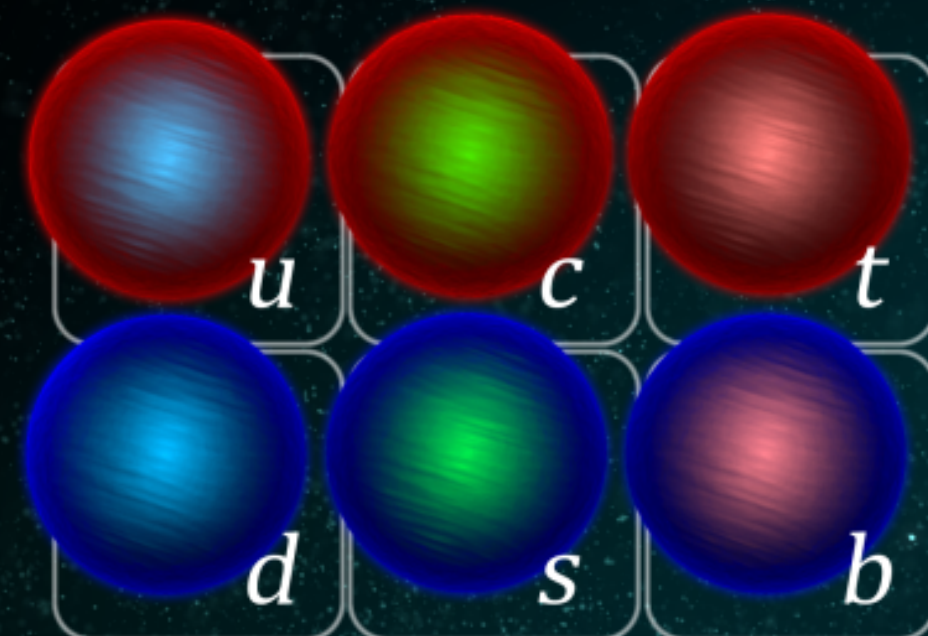
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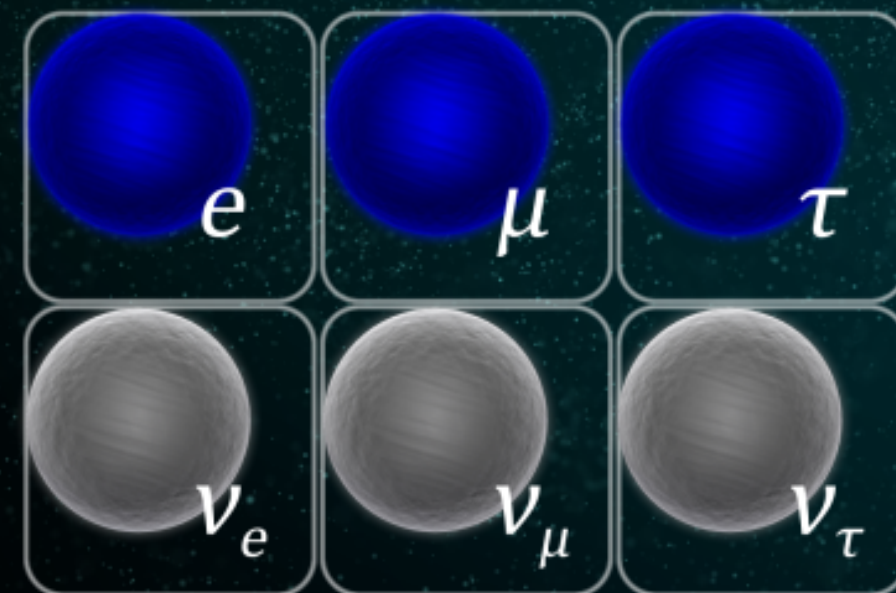
for matter:

rest of this talk!

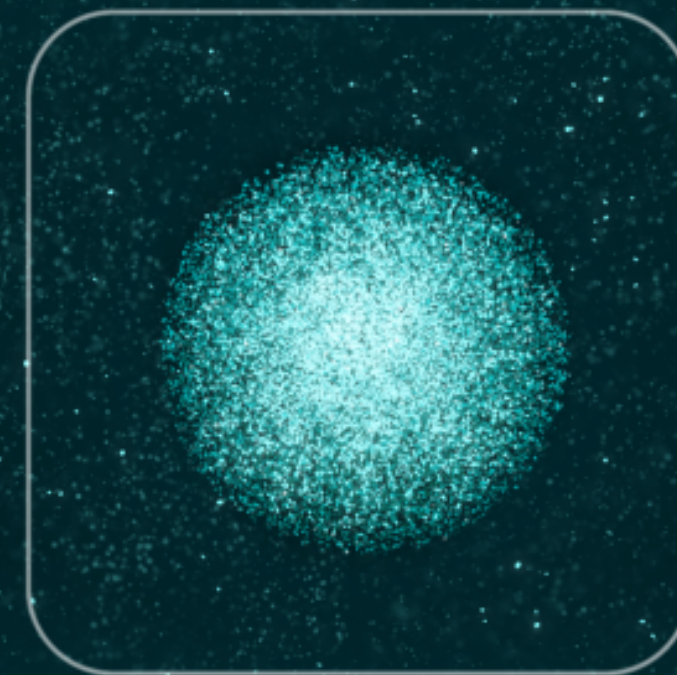
Standard Model couplings



Quarks



Leptons



Higgs boson



Forces

The microscopic structure of quantum space-time and matter from a renormalization group perspective

[Astrid Eichhorn](#)

[Nature Physics](#) **19**, 1527–1529 (2023) | [Cite this article](#)

[arXiv](#) > hep-th > arXiv:2212.07456

High Energy Physics – Theory

[Submitted on 14 Dec 2022]

Asymptotic safety of gravity with matter

[Astrid Eichhorn](#), [Marc Schiffer](#)

[arXiv](#) > hep-th > arXiv:1810.07615

High Energy Physics – Theory

[Submitted on 17 Oct 2018 ([v1](#)), last revised 21 Feb 2019 (this version, v2)]

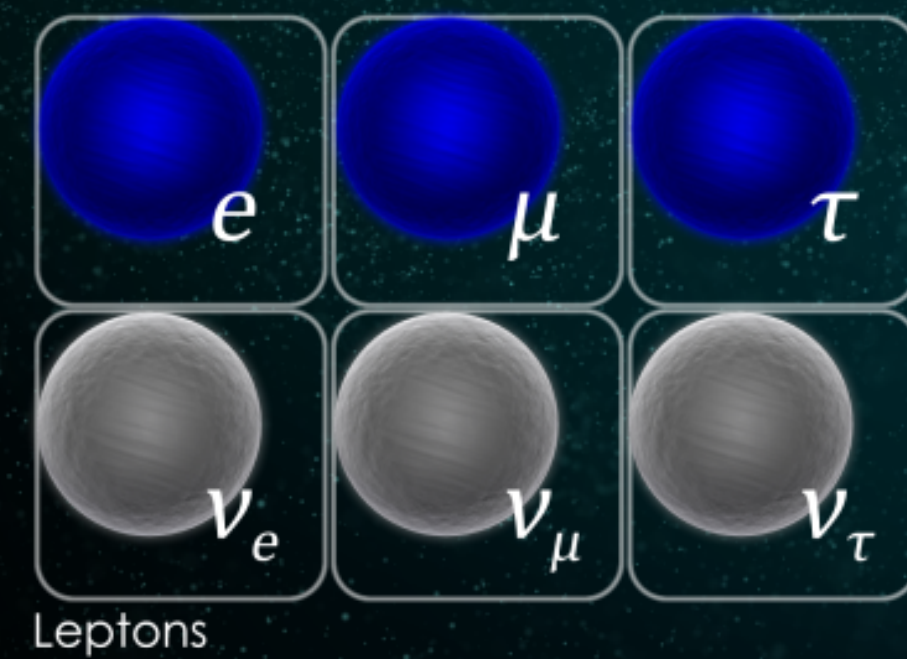
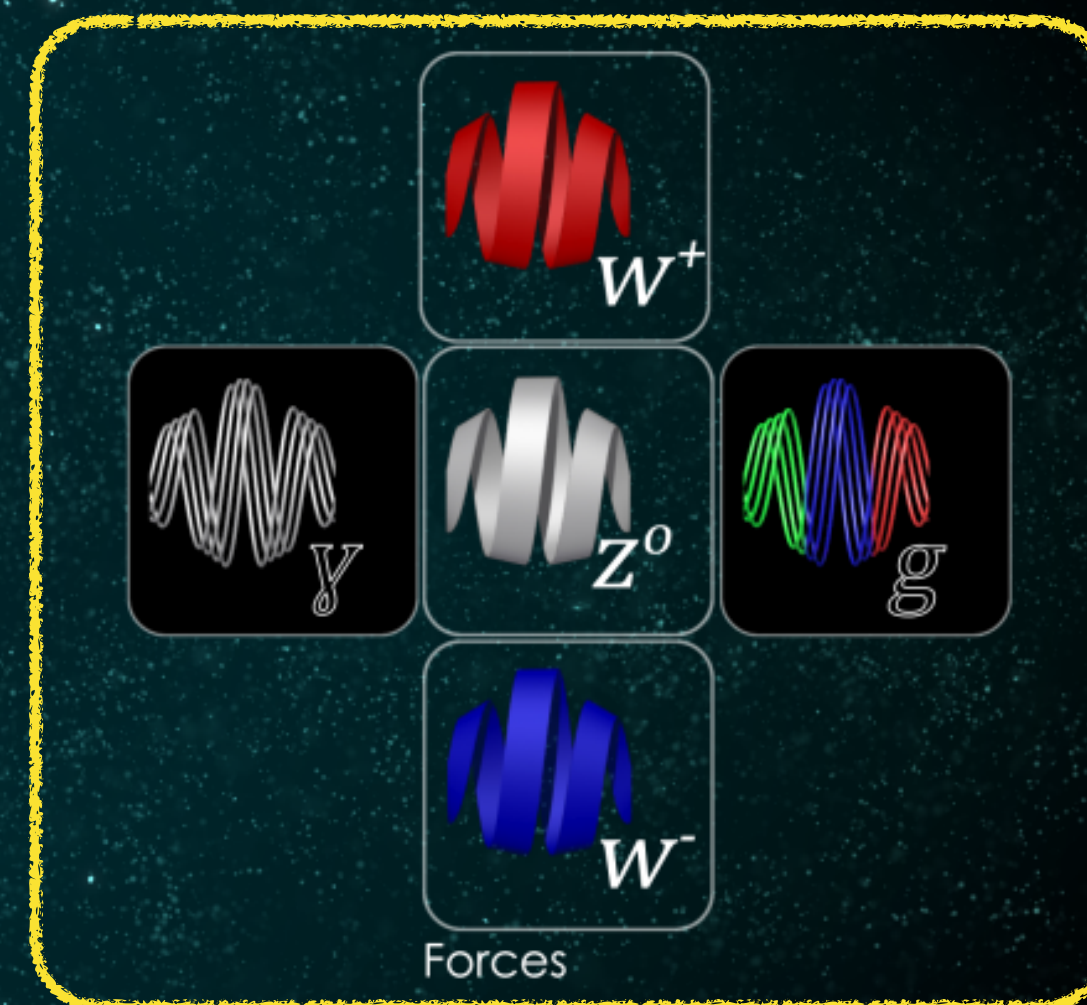
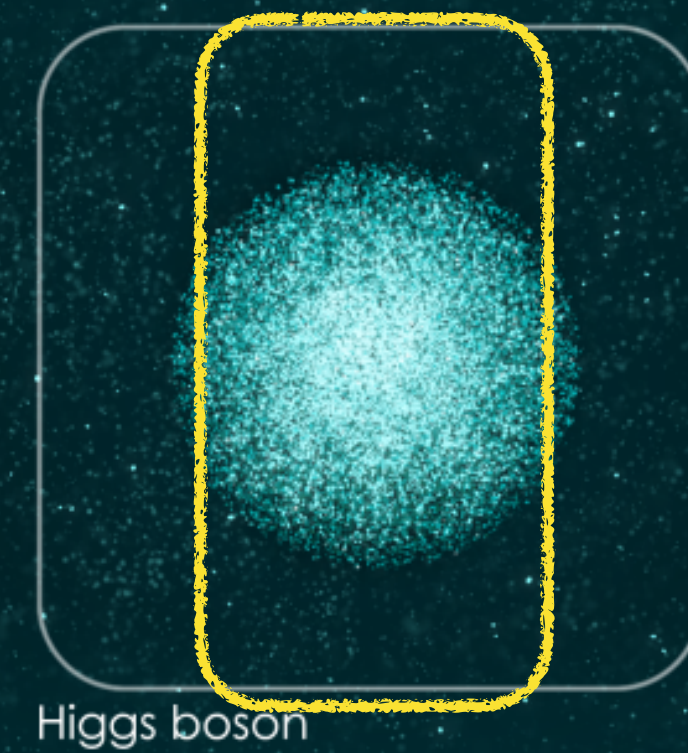
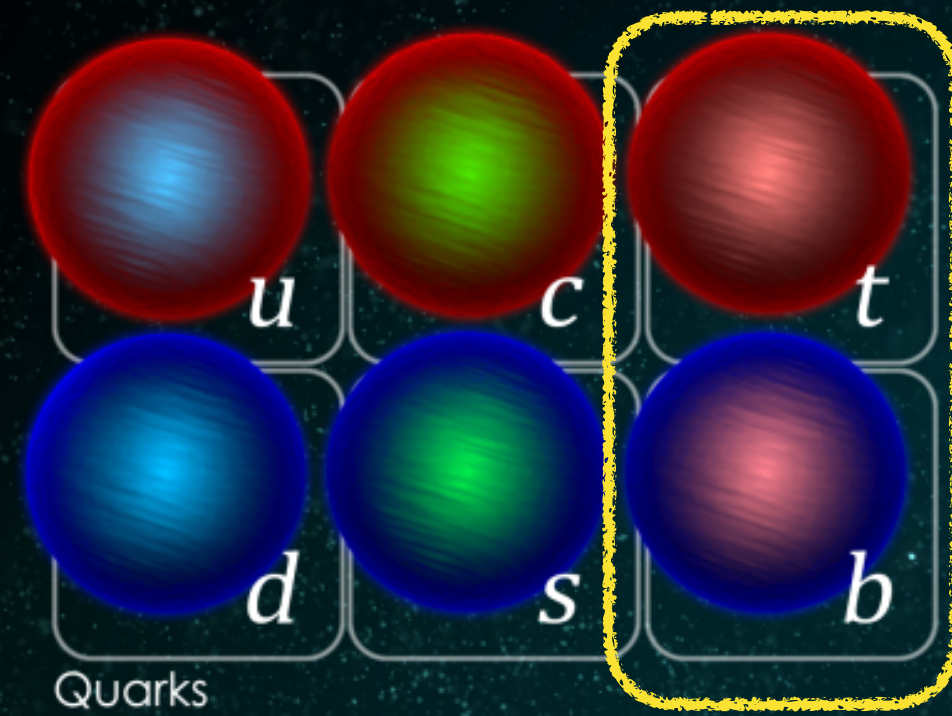
An asymptotically safe guide to quantum gravity and matter

[Astrid Eichhorn](#)

Standard Model couplings

Part 1: “heavy” Standard Model

[Harst, Reuter '11; Shaposhnikov, Wetterich '09, AE, Held '17, '18, AE, Versteegen '17]



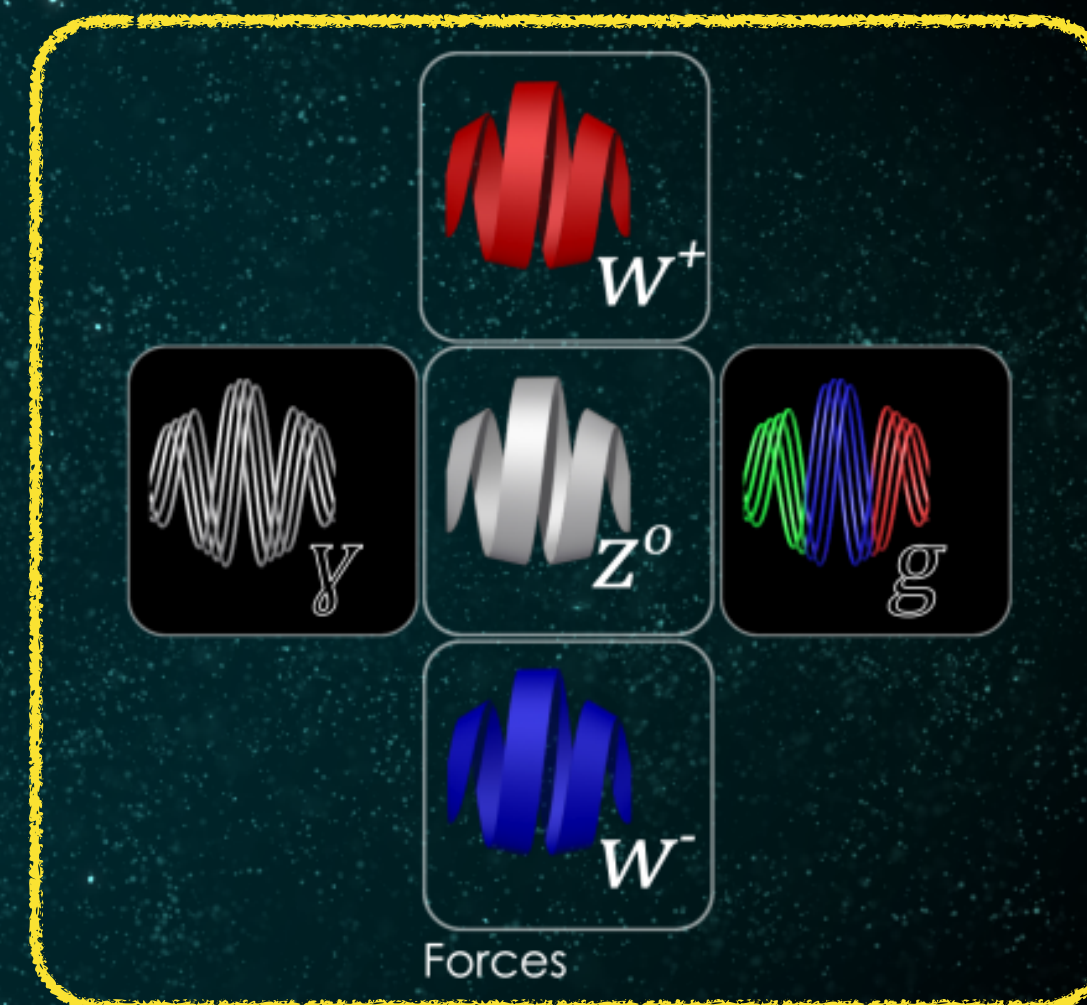
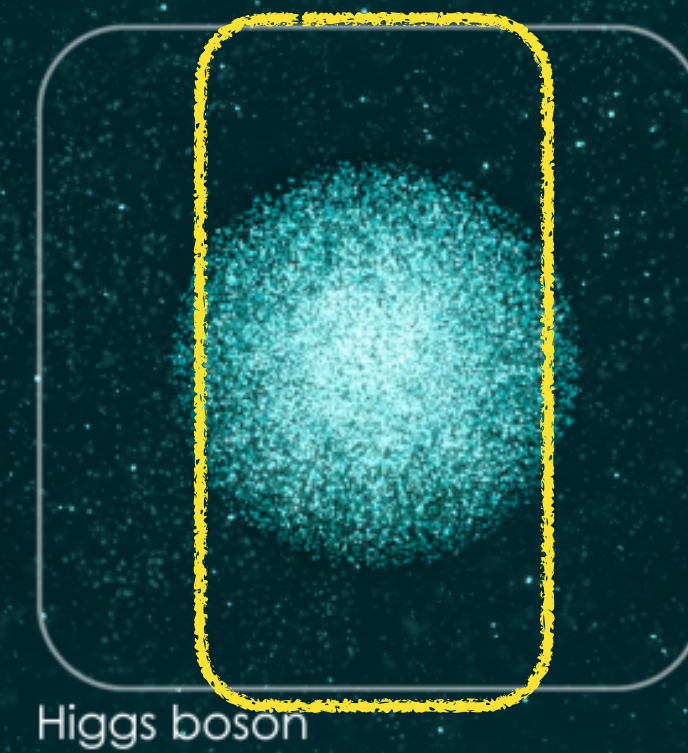
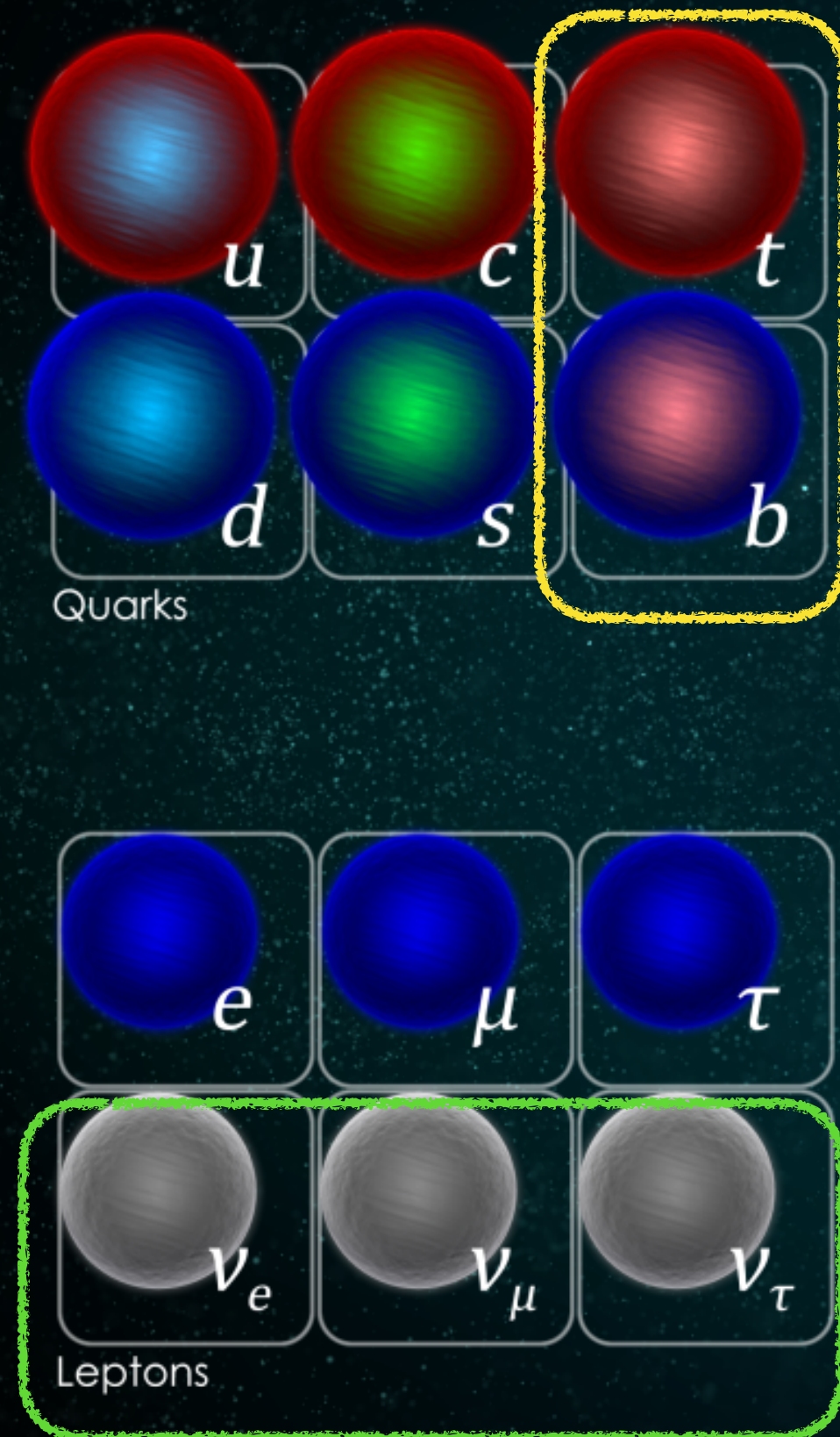
Standard Model couplings

Part 1: “heavy” Standard Model

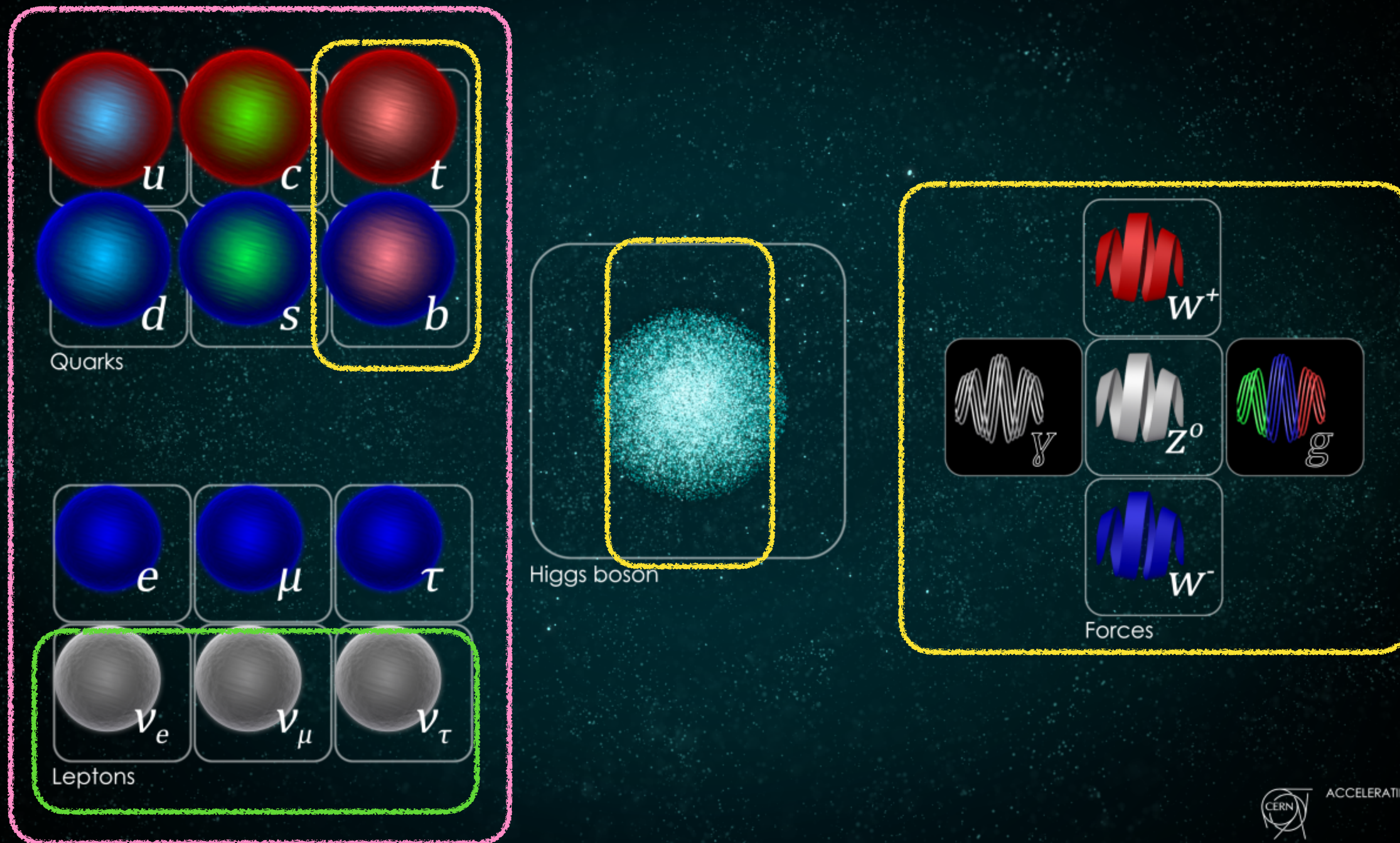
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Part 2: the lightest fermions: neutrinos

[Held, PhD thesis '19;
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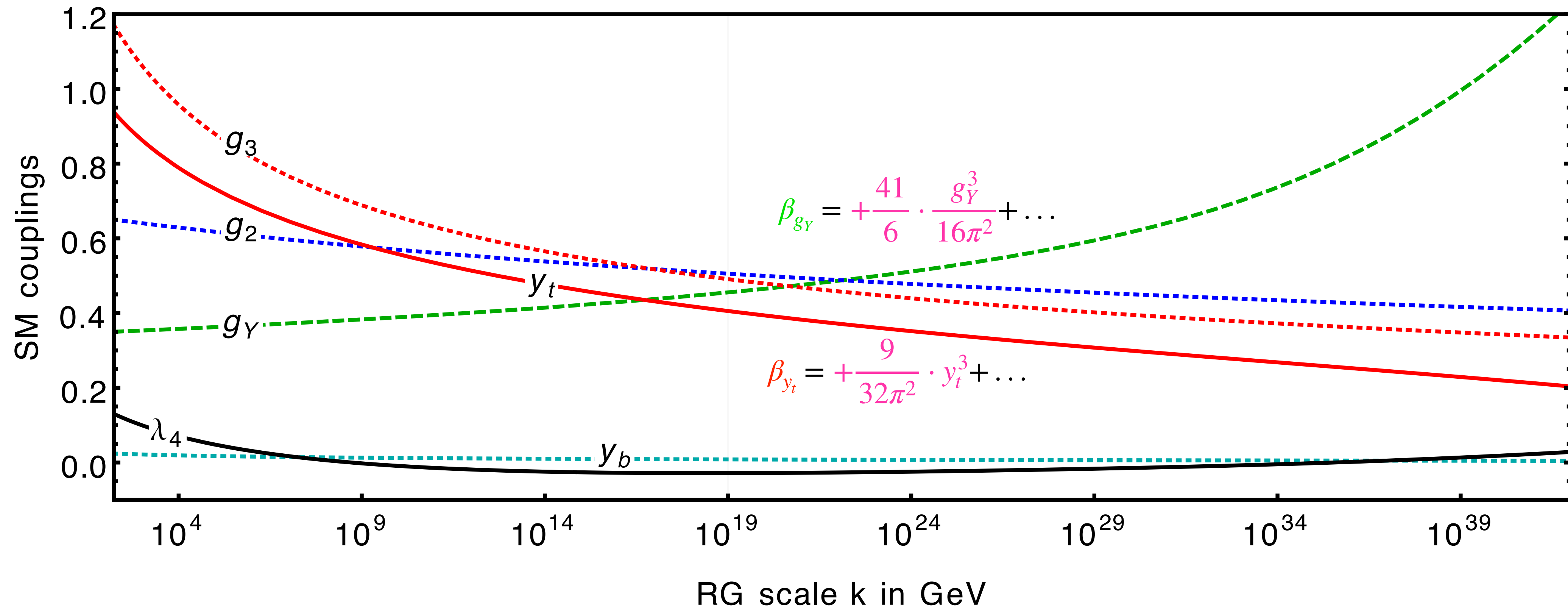
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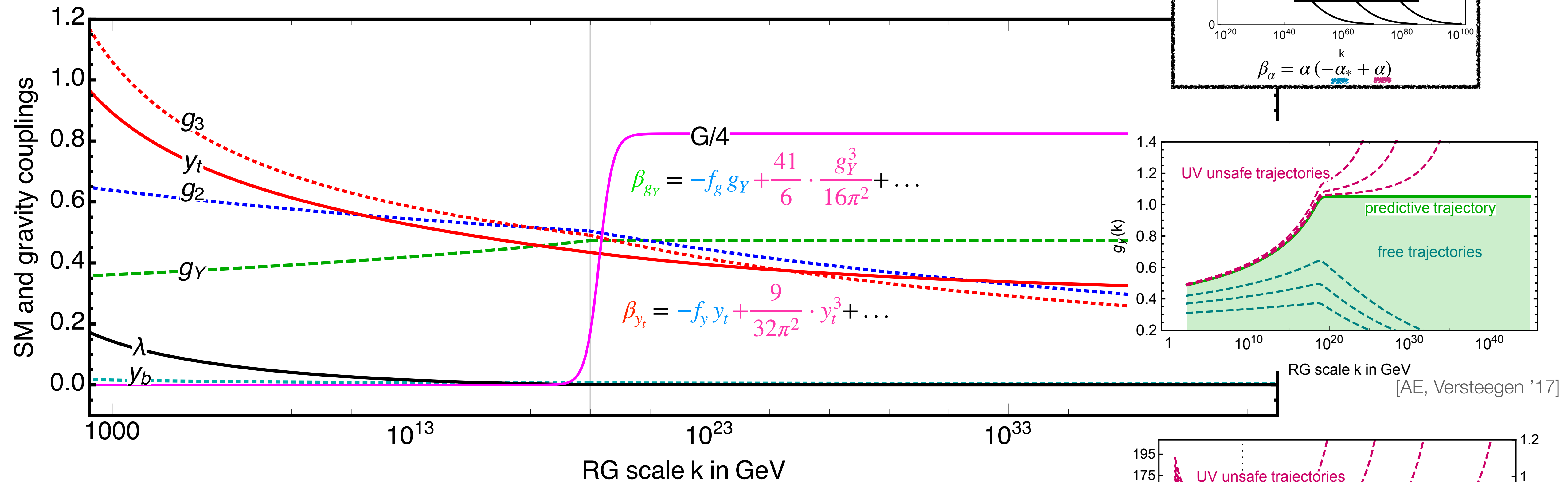
Part 1: Heavy Standard Model



without gravity:

- not ultraviolet complete (Landau pole/triviality problem)
- measured values are free parameters

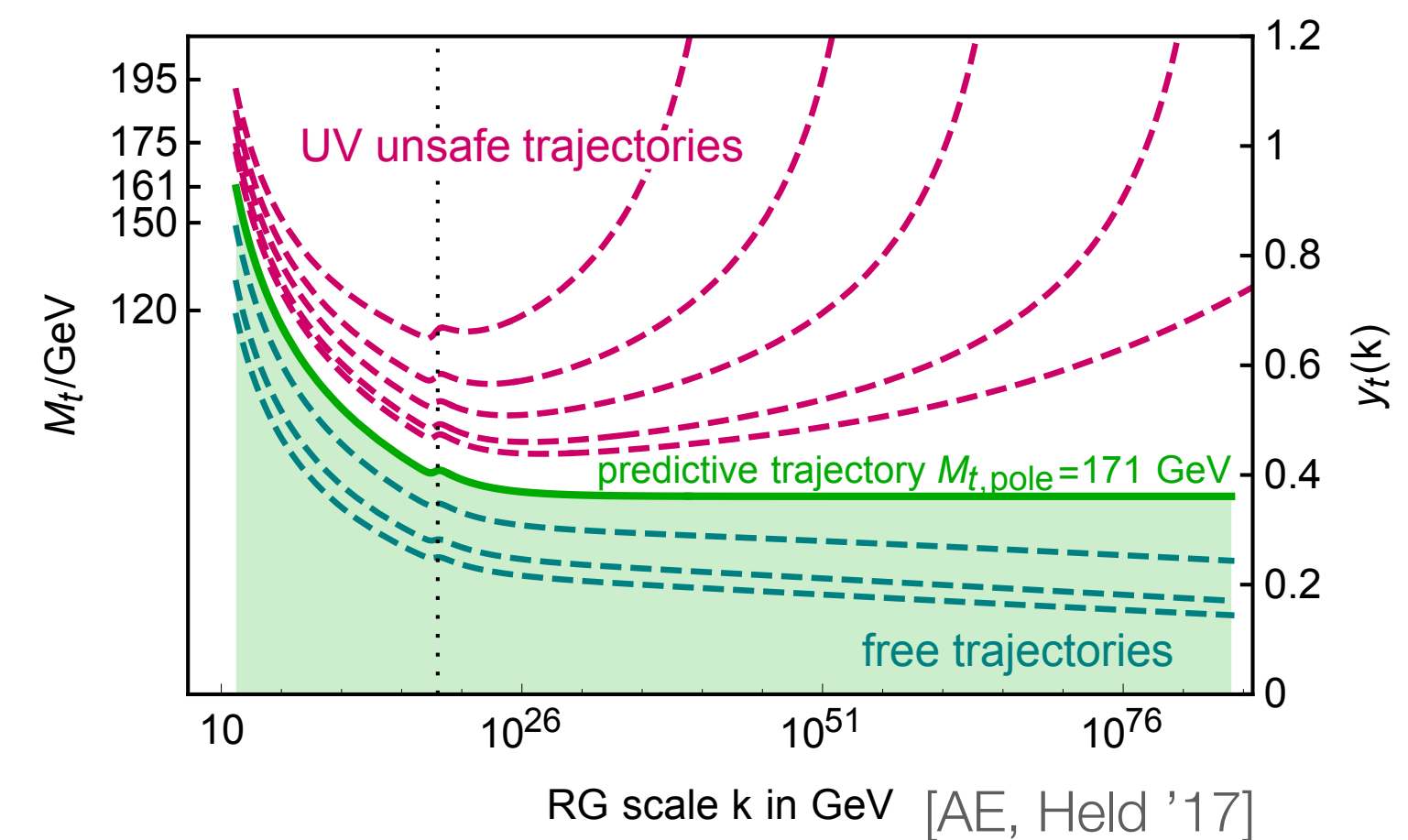
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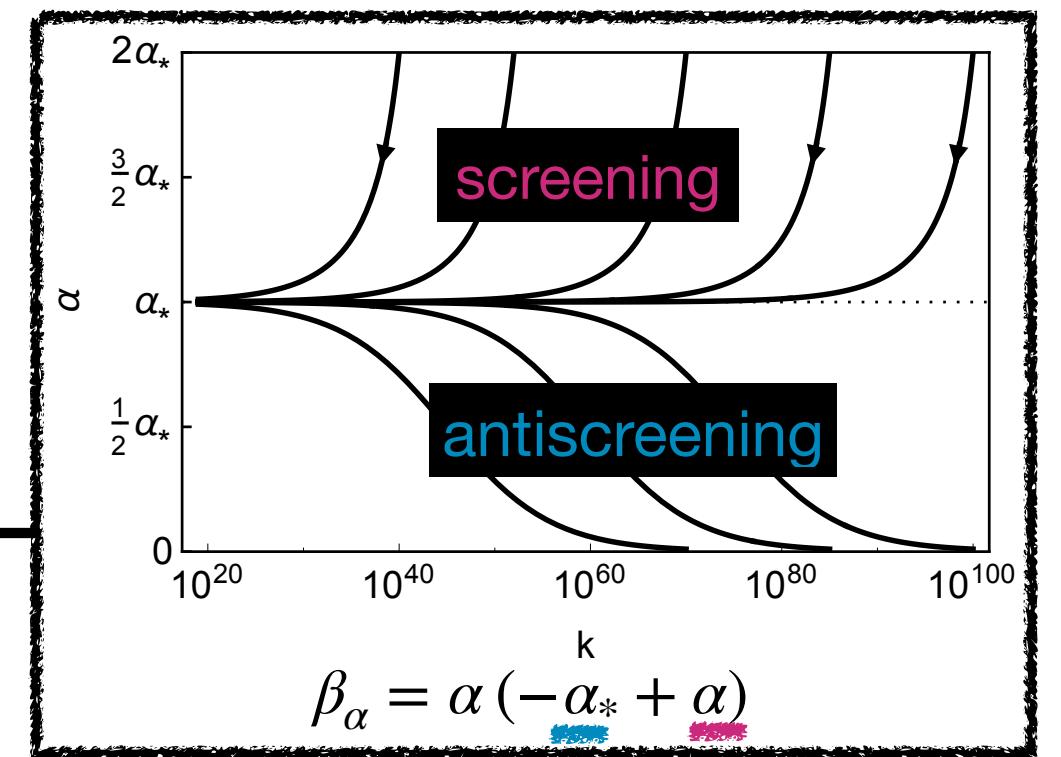
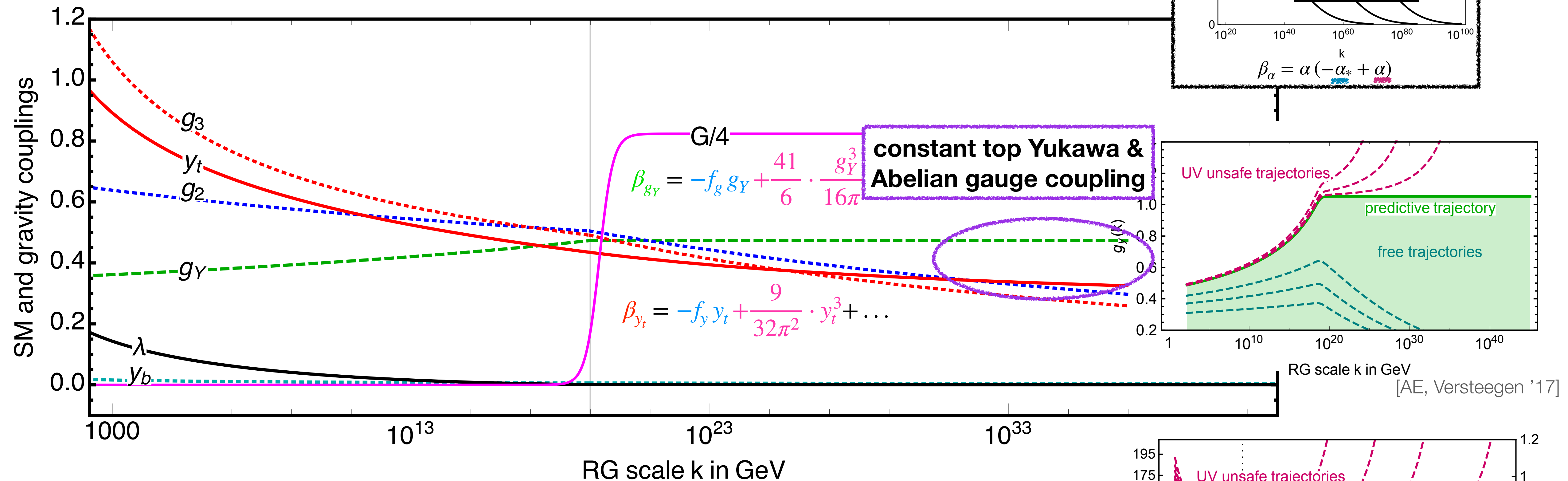
with gravity:

- ultraviolet complete (no Landau pole/triviality problem)
- y_t and g_Y are bounded from above

& can be constant above Planck scale for appropriate f_g, f_y



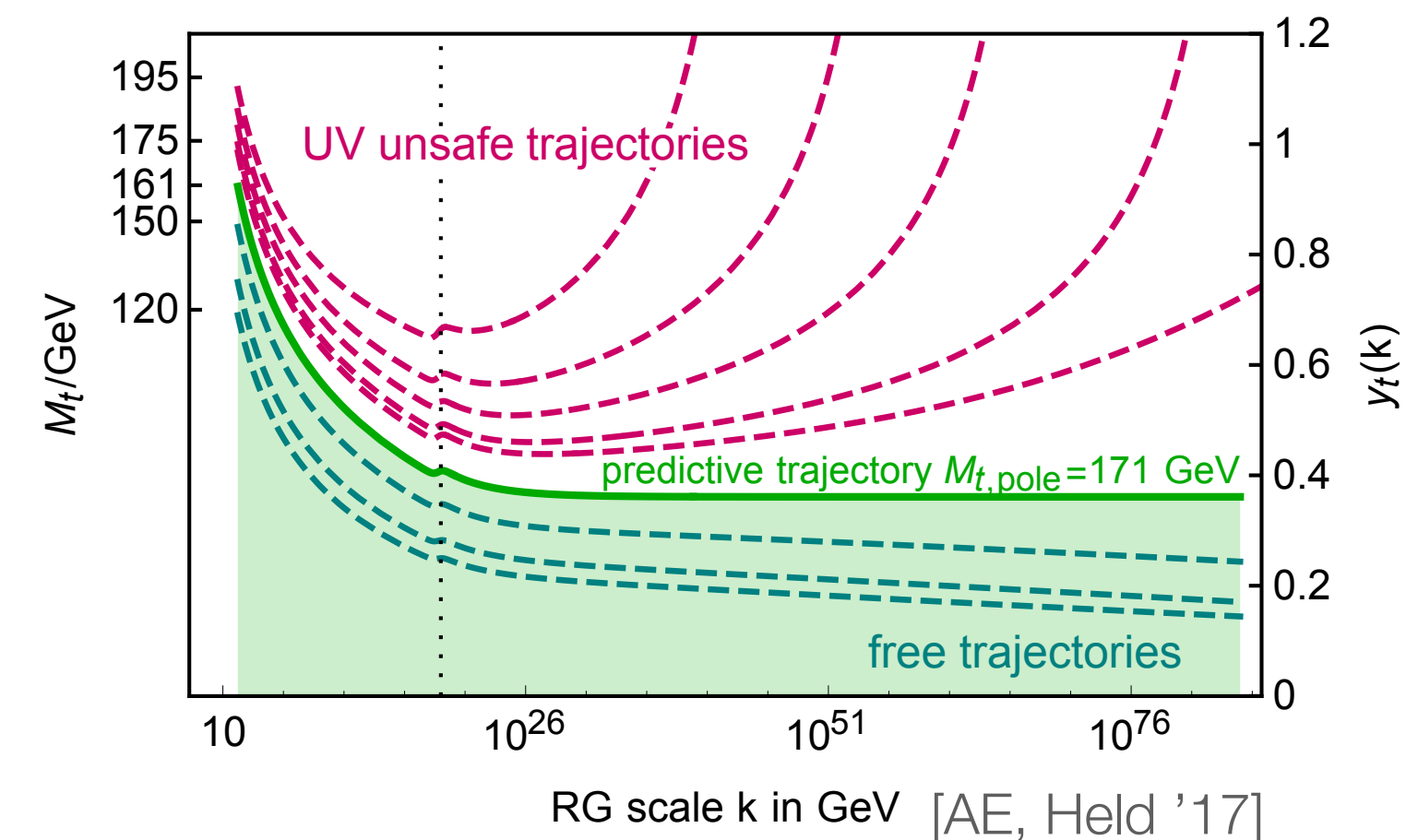
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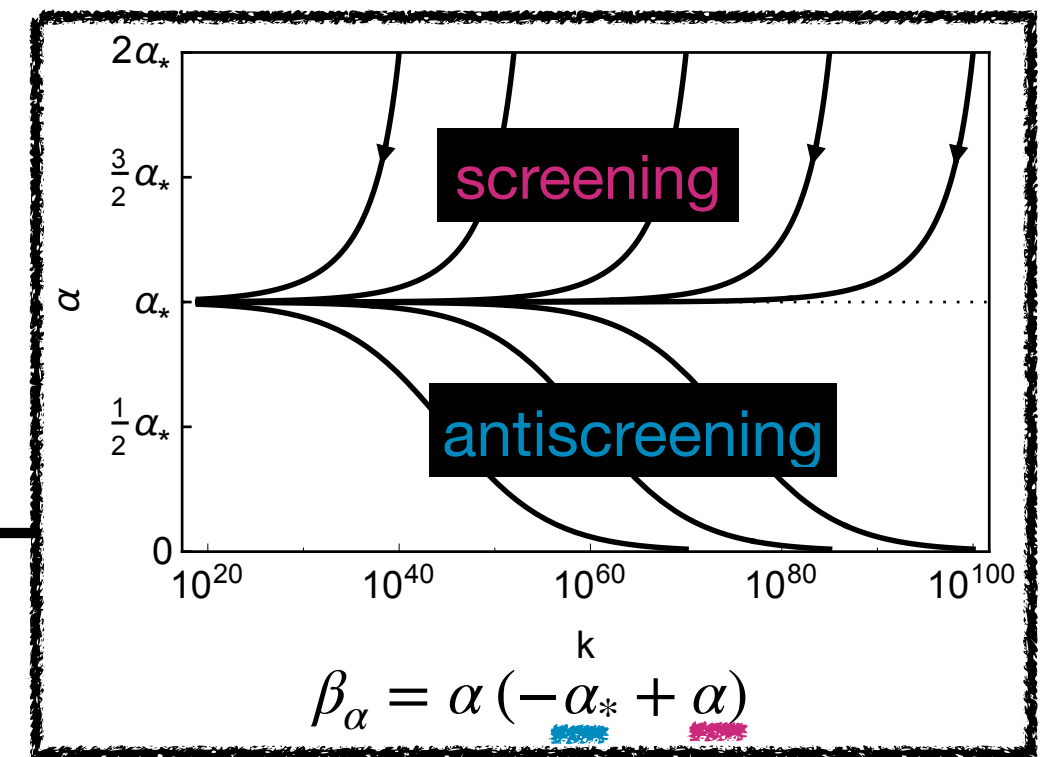
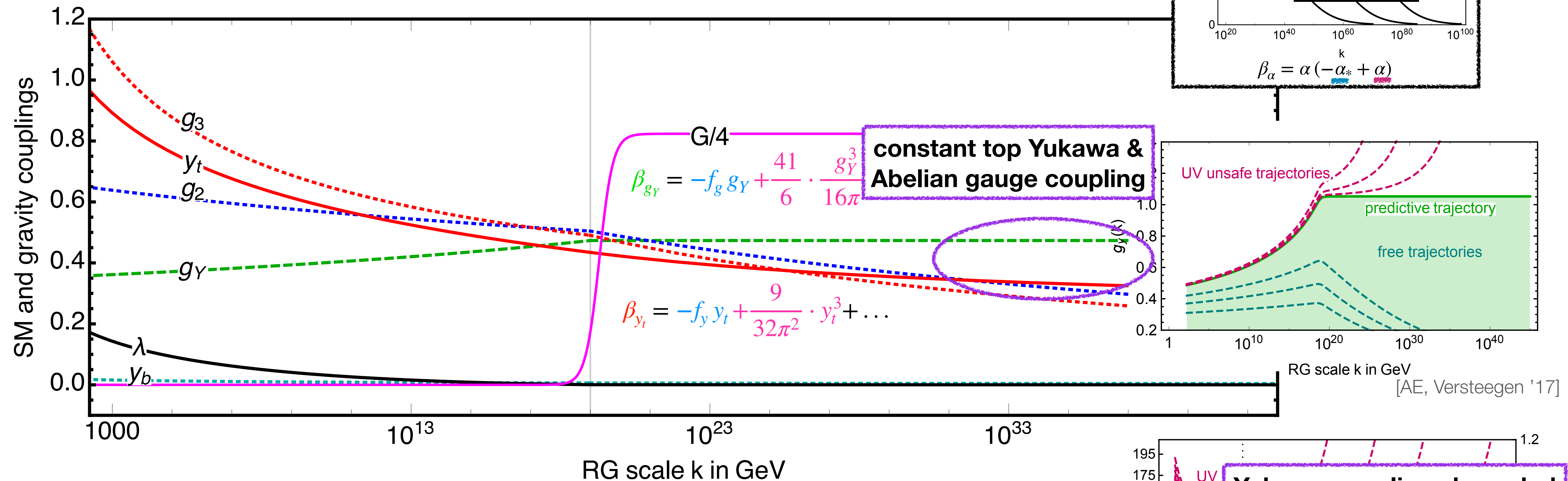
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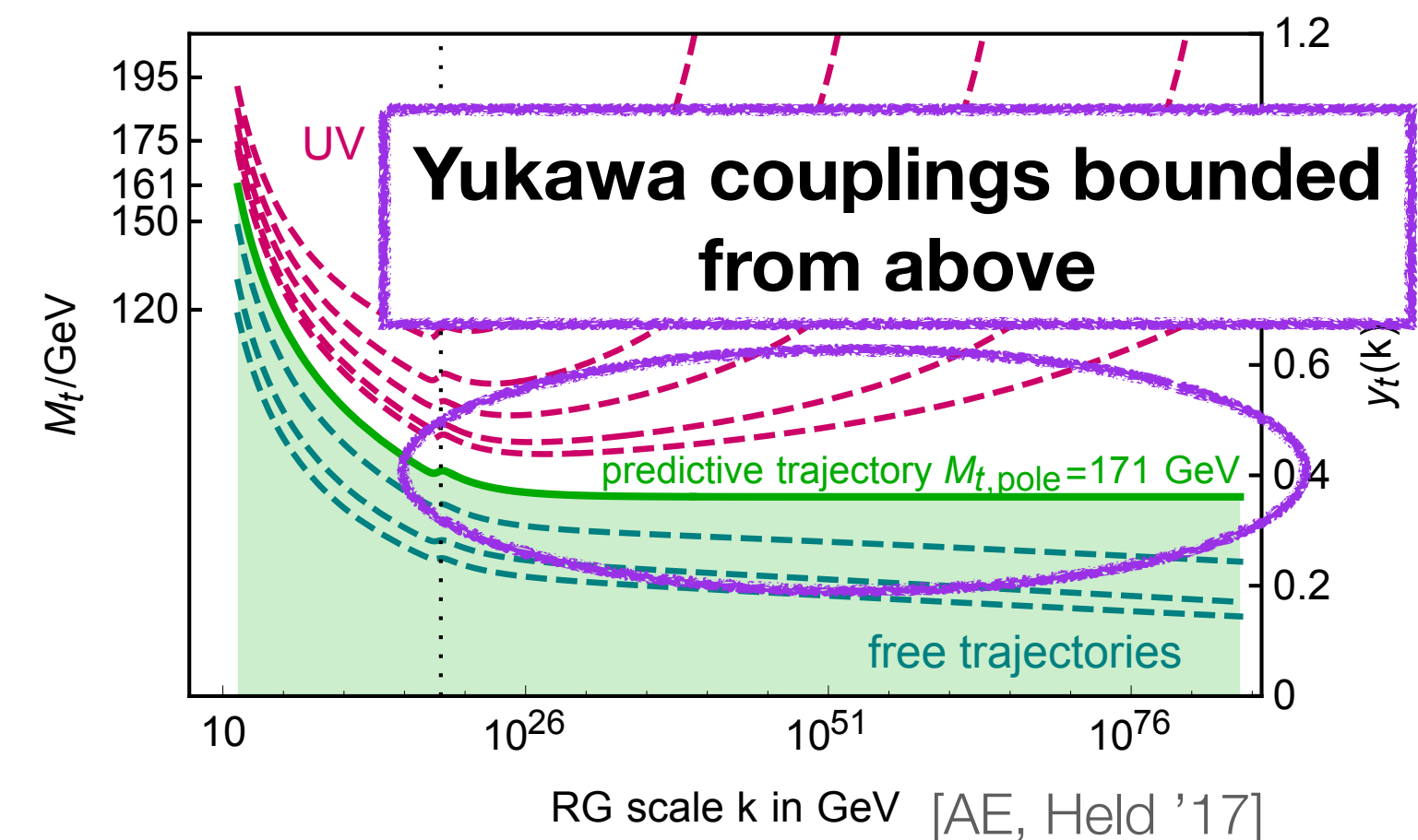
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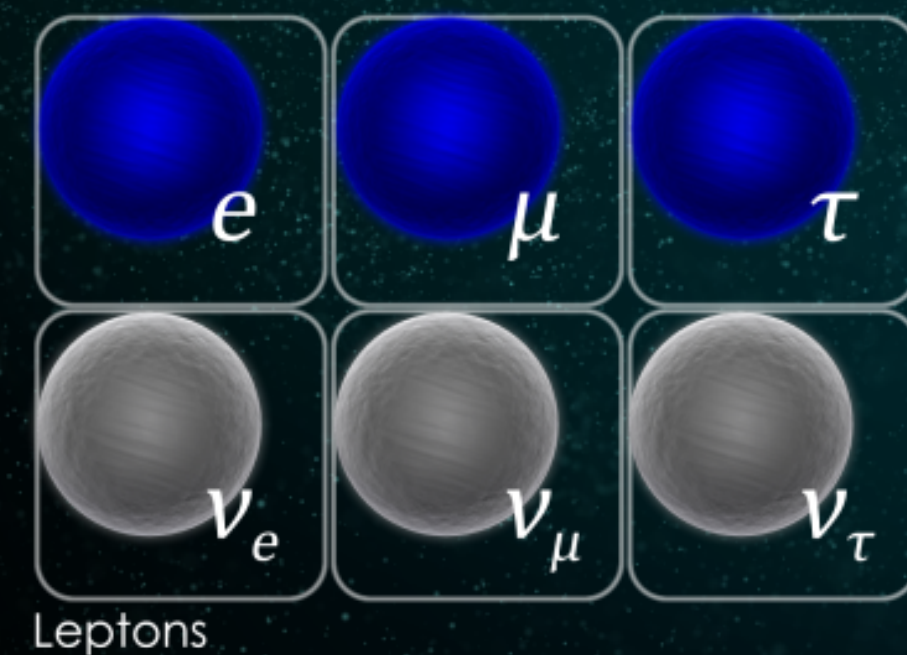
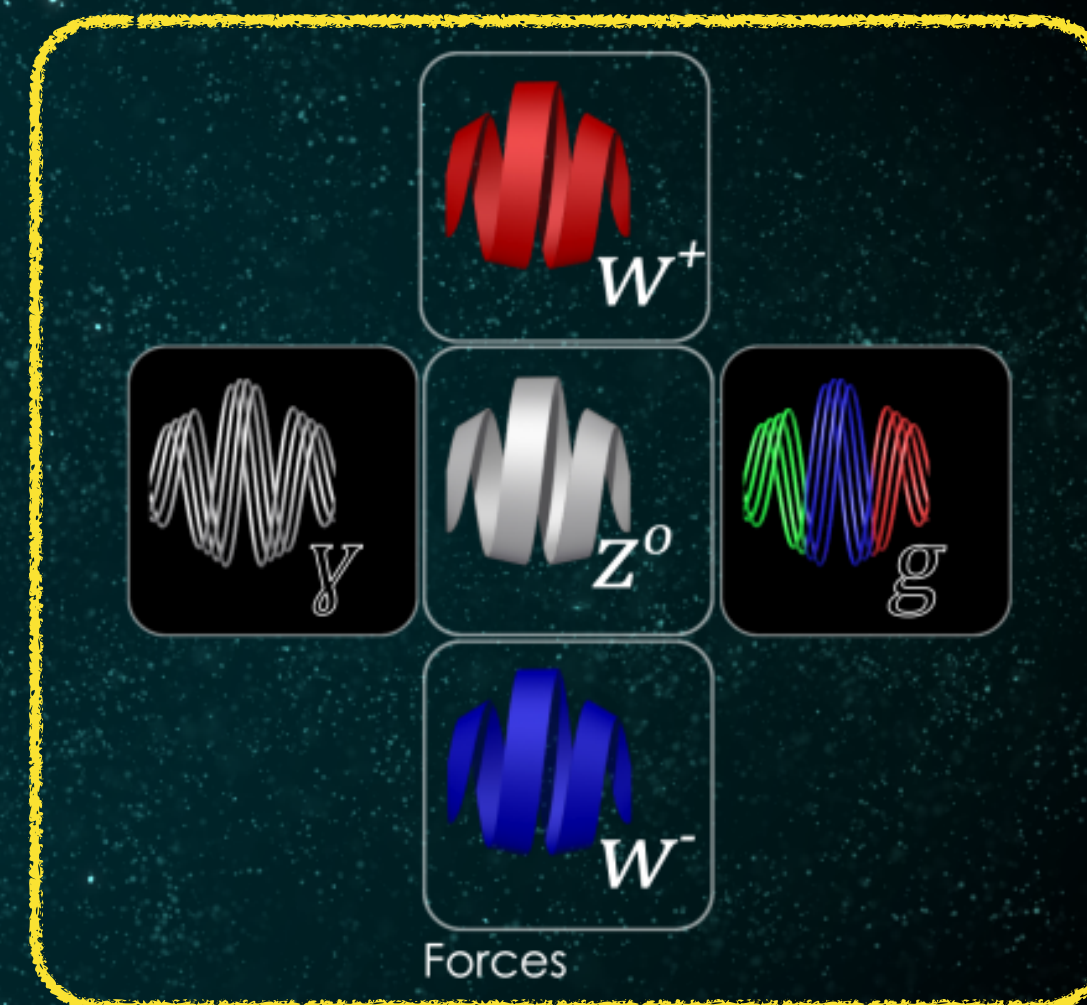
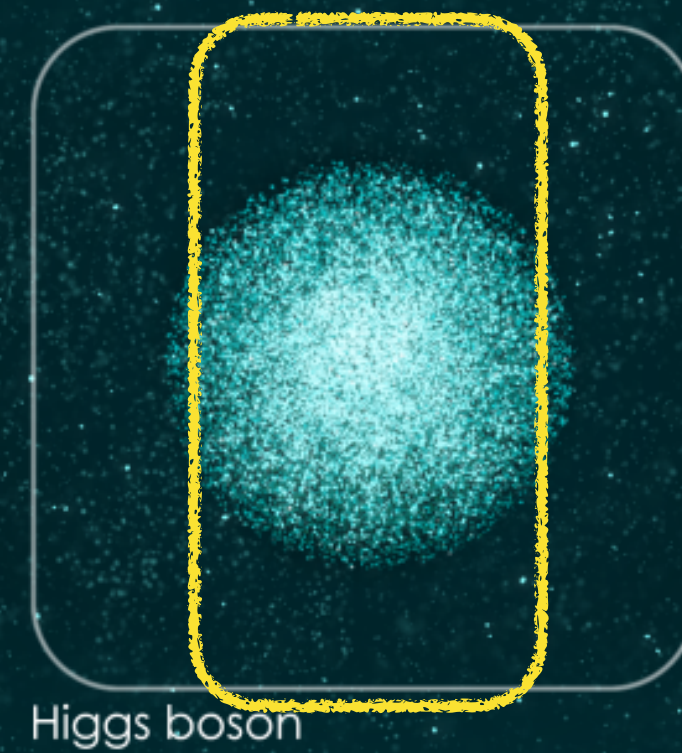
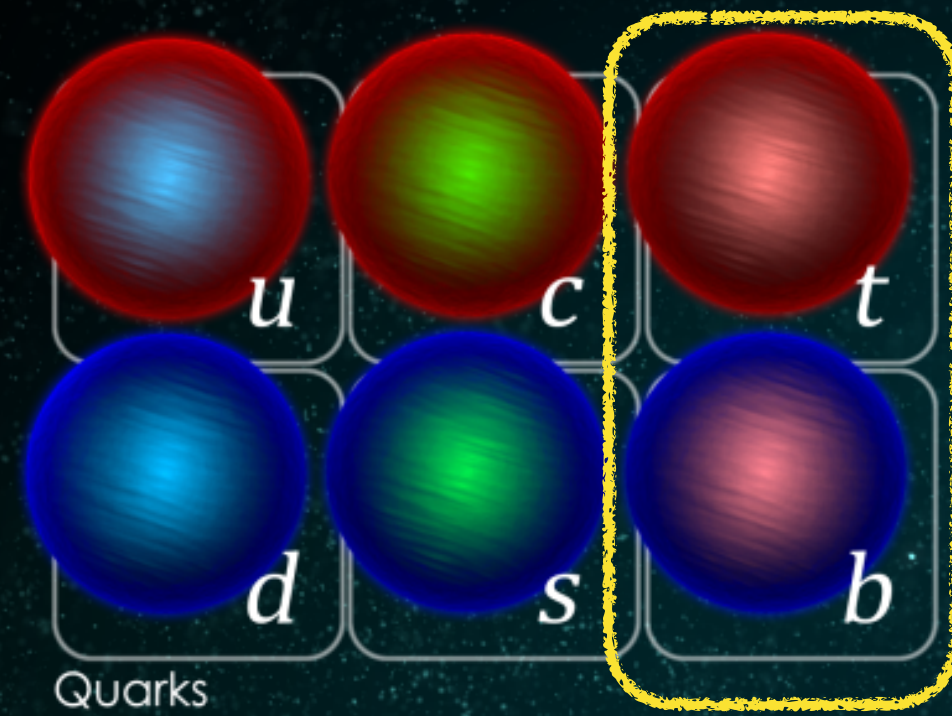
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Part 2: the lightest fermions: neutrinos

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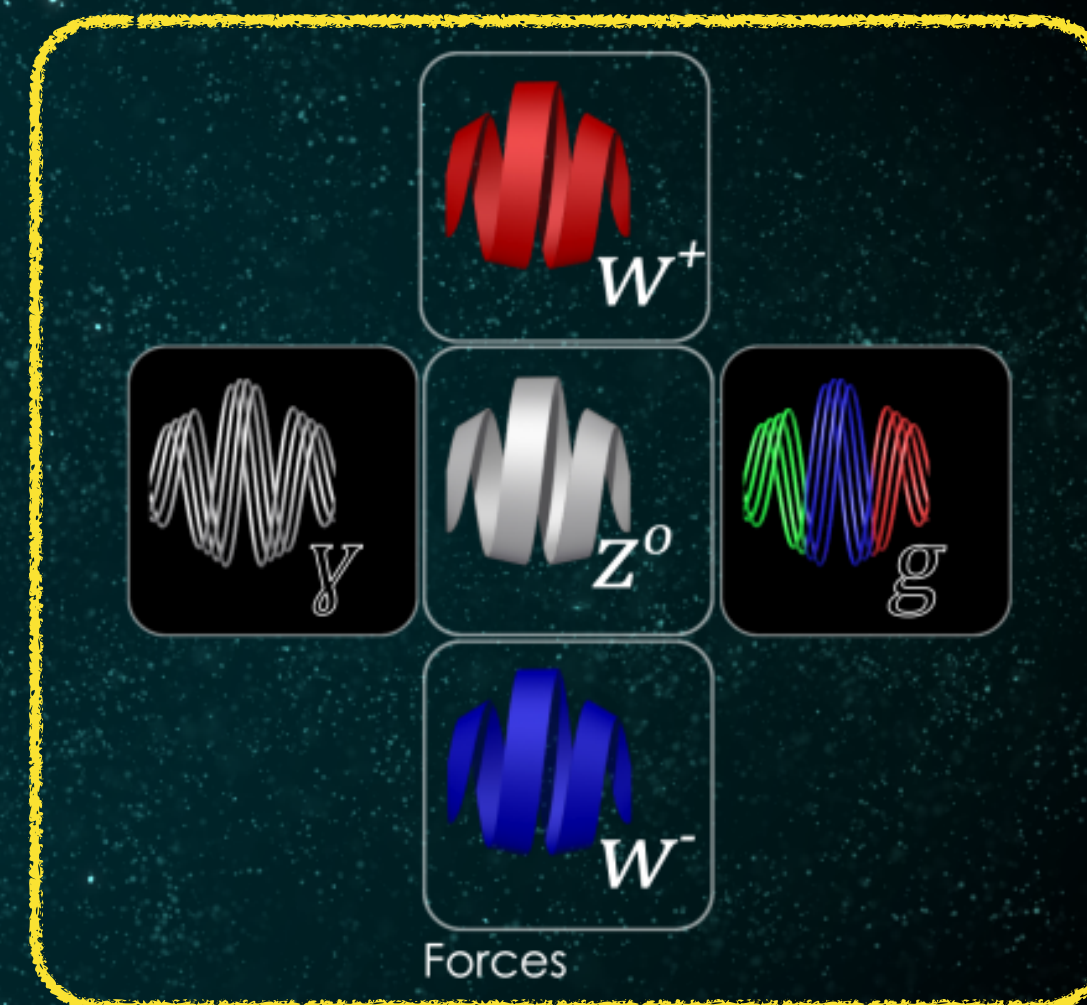
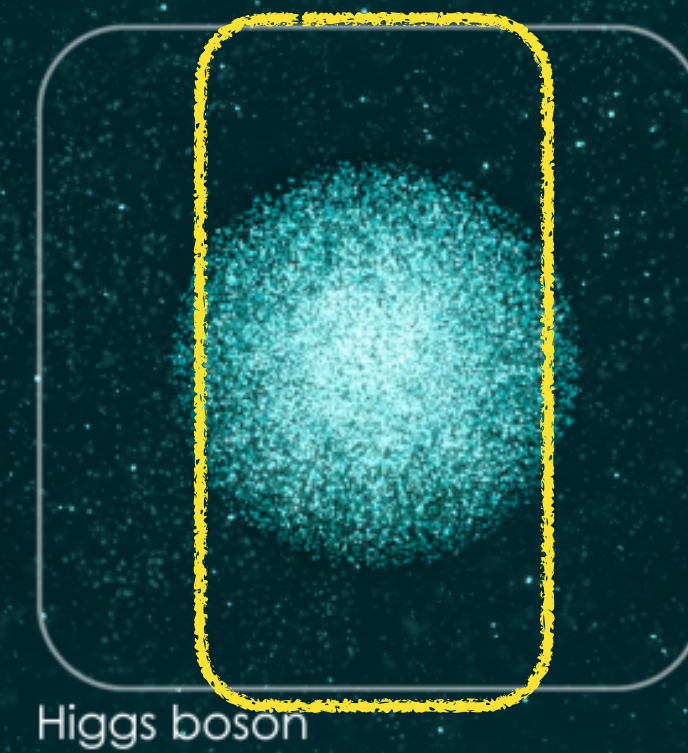
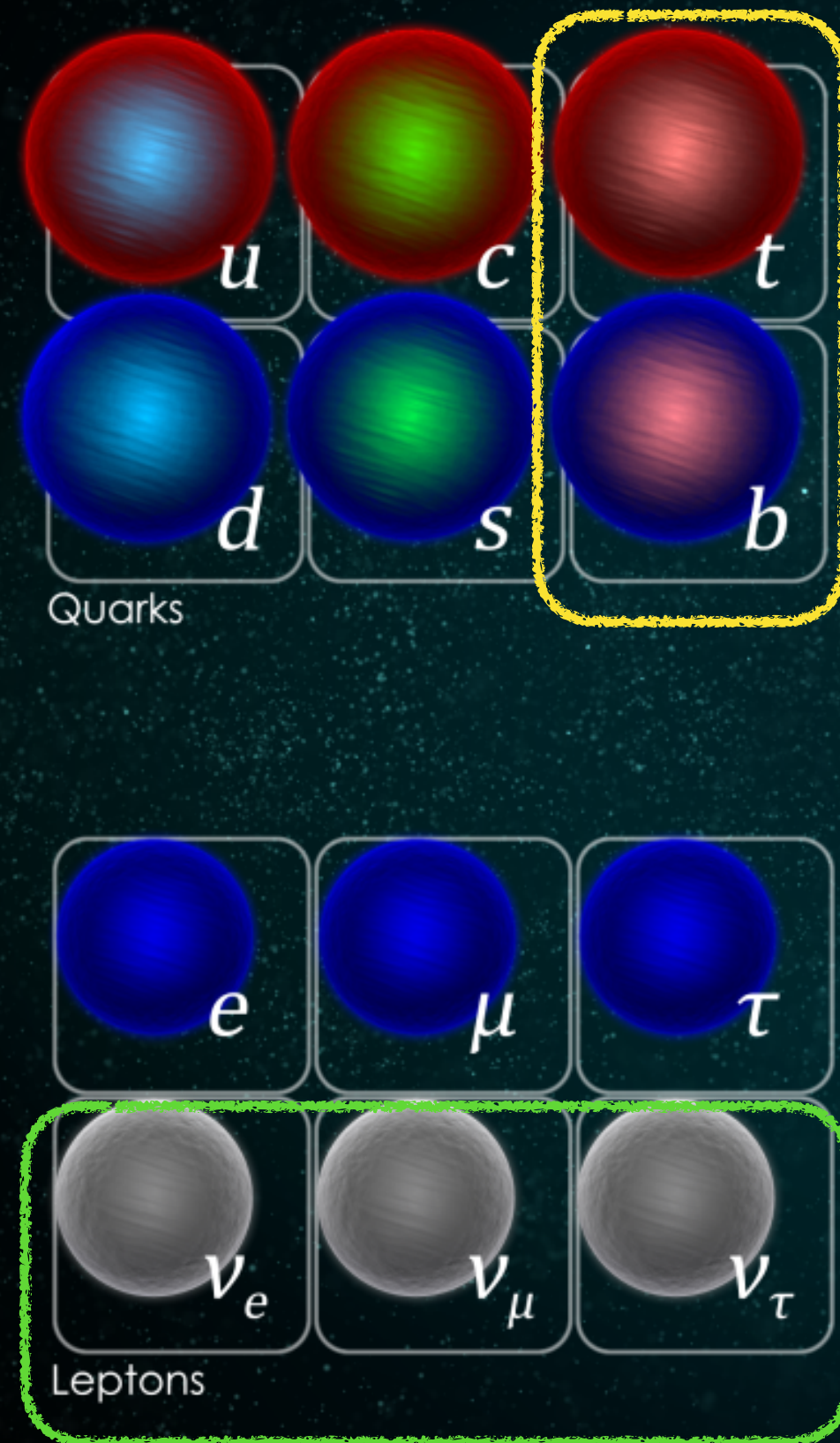
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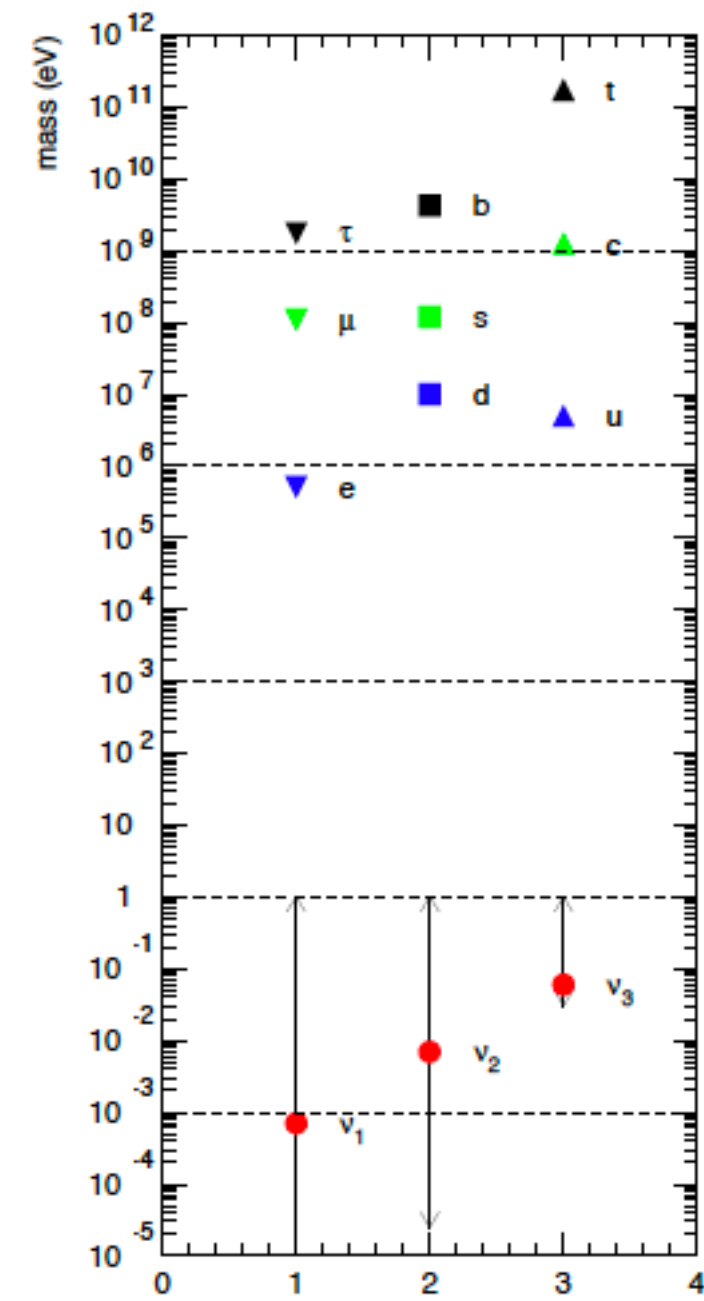
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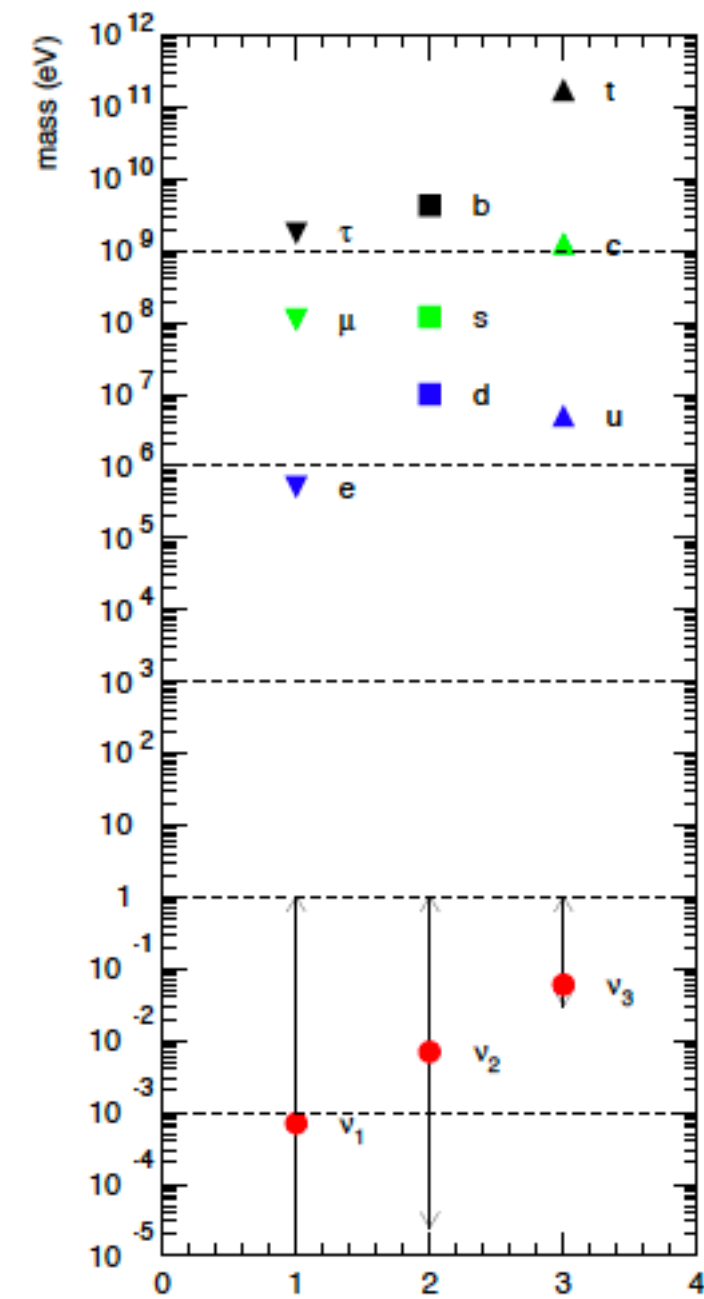
Part 2: neutrino masses

Standard Model fermion masses



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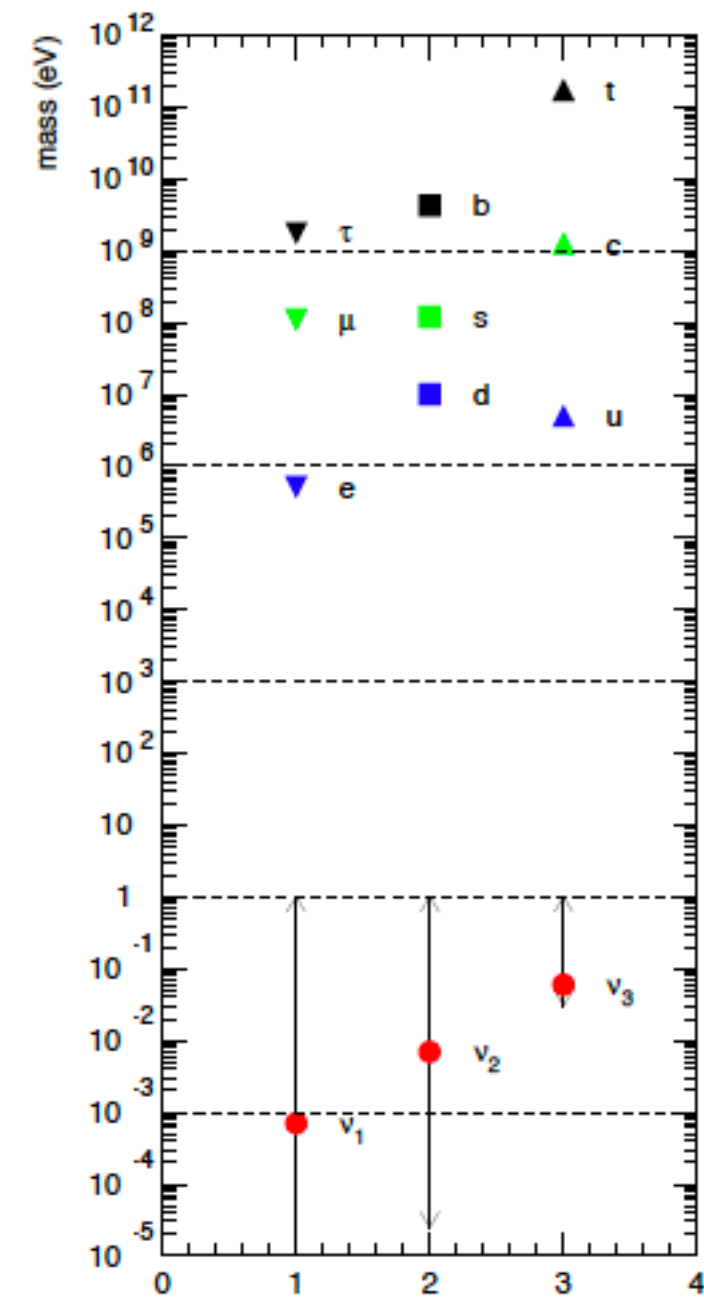


- Option 1: no new degrees of freedom, just Weinberg operator

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Part 2: neutrino masses

Standard Model fermion masses



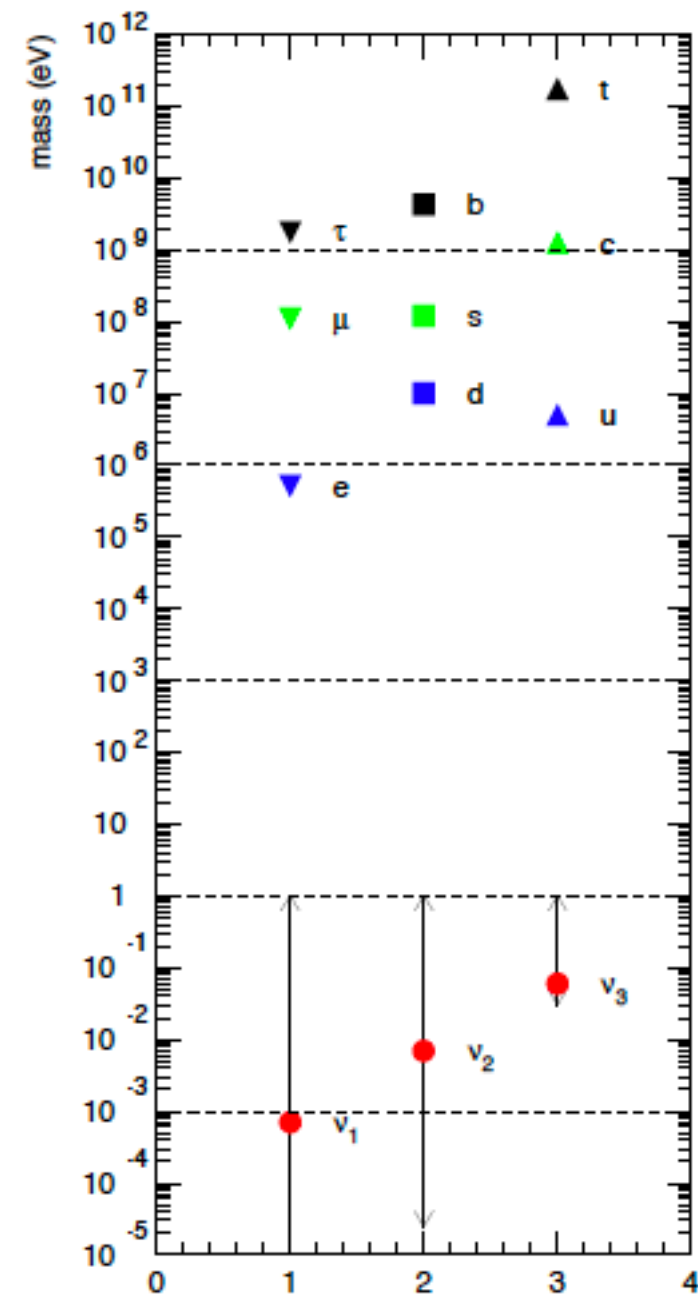
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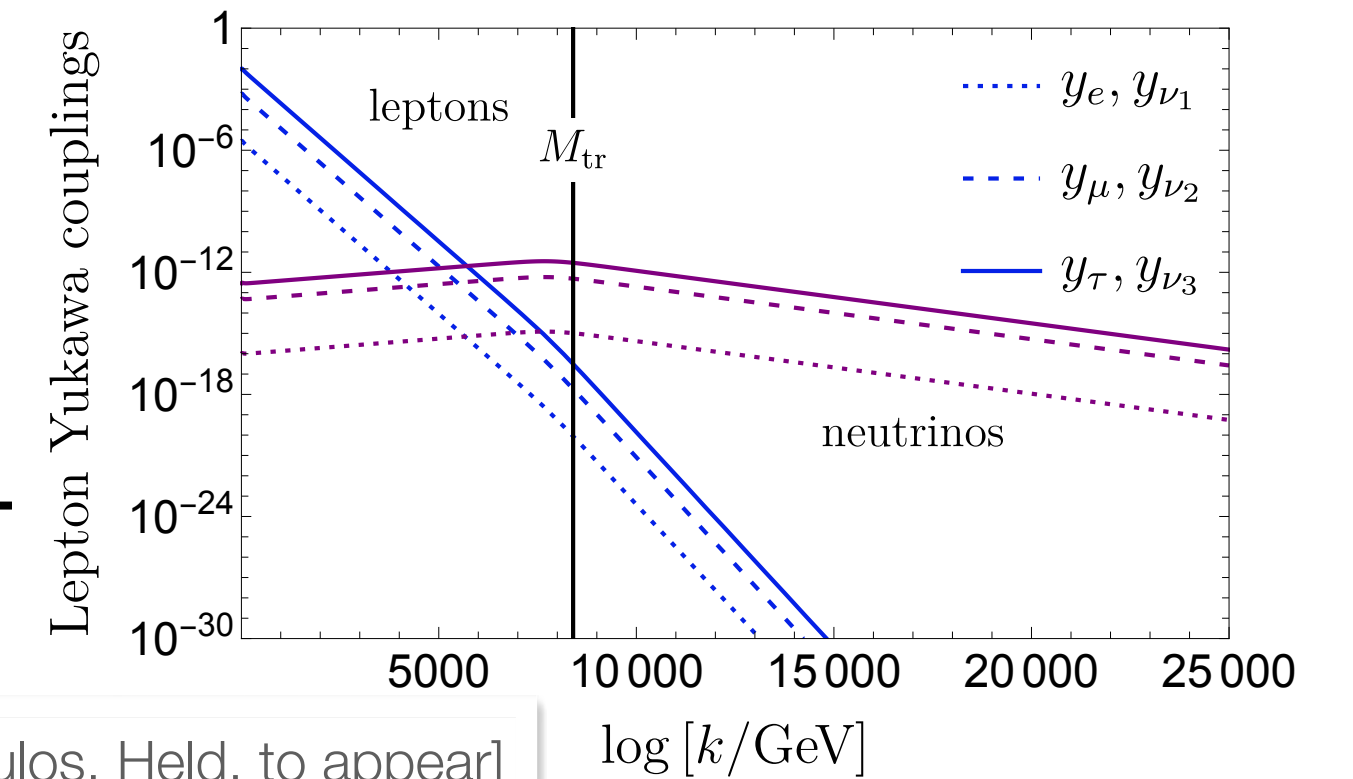
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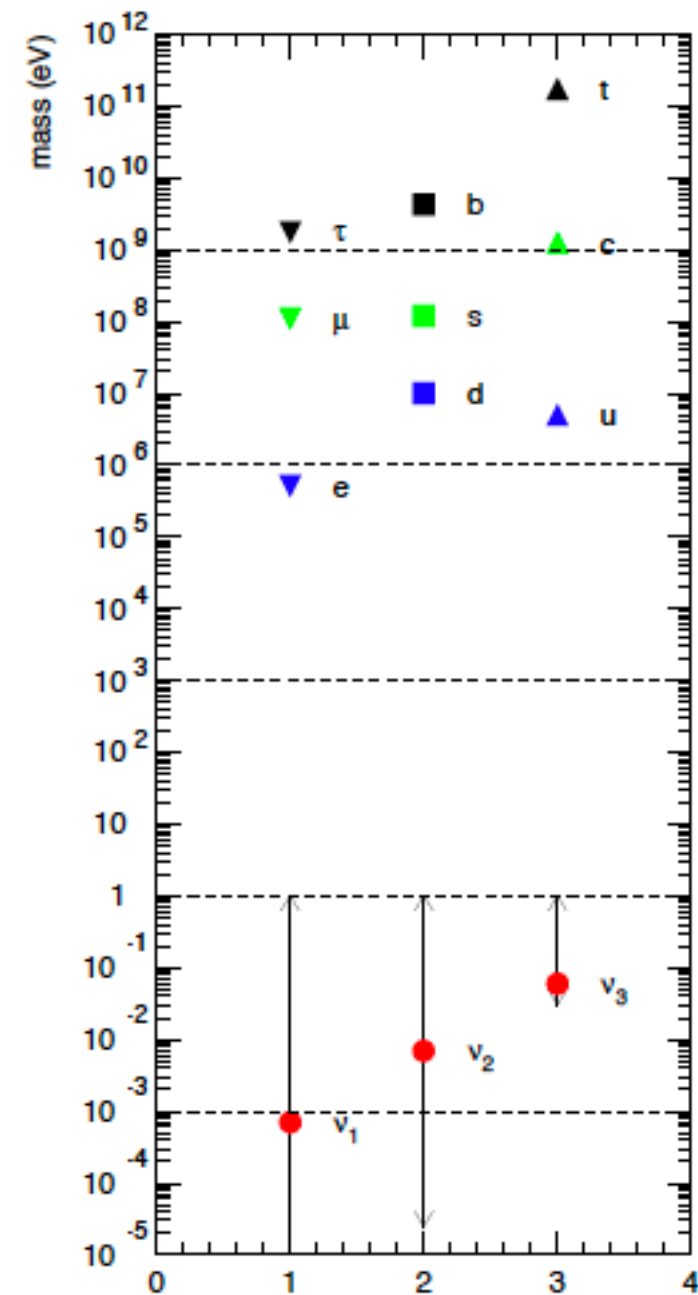
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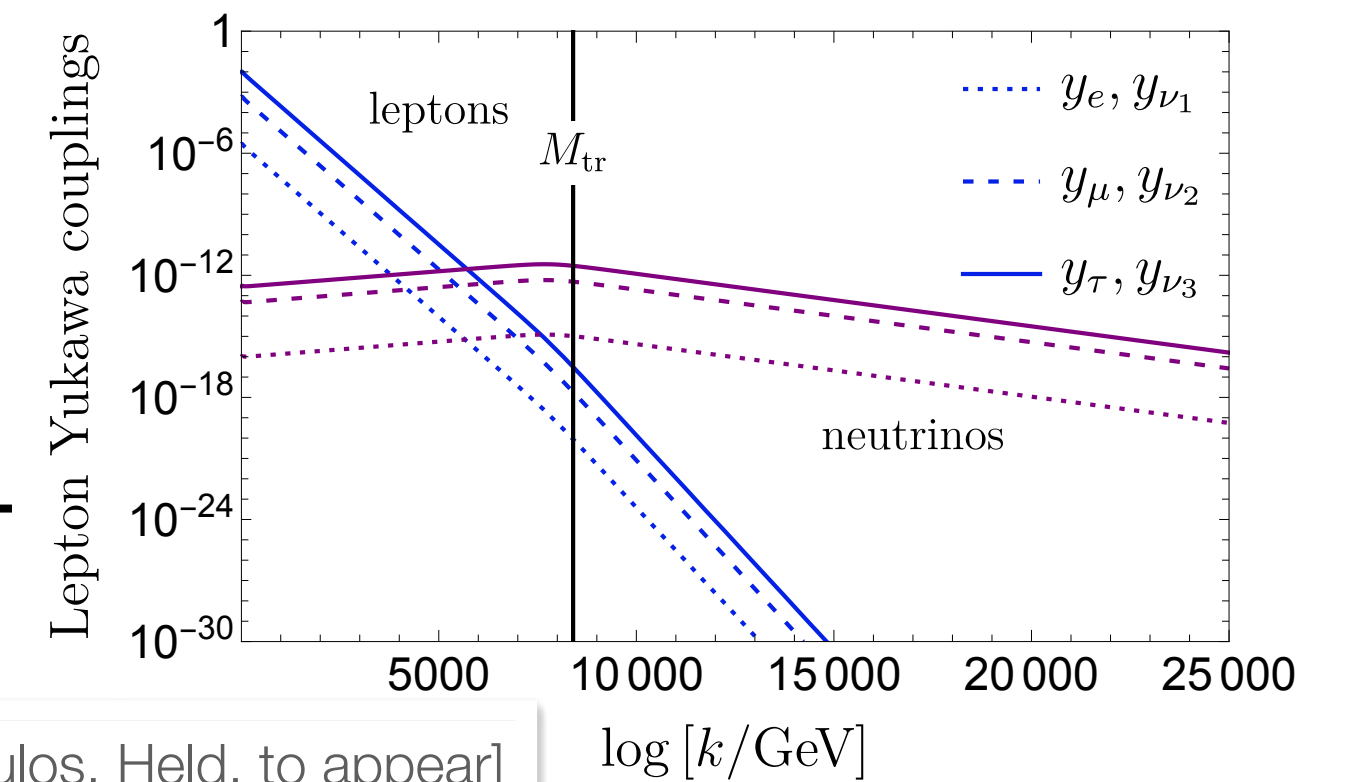
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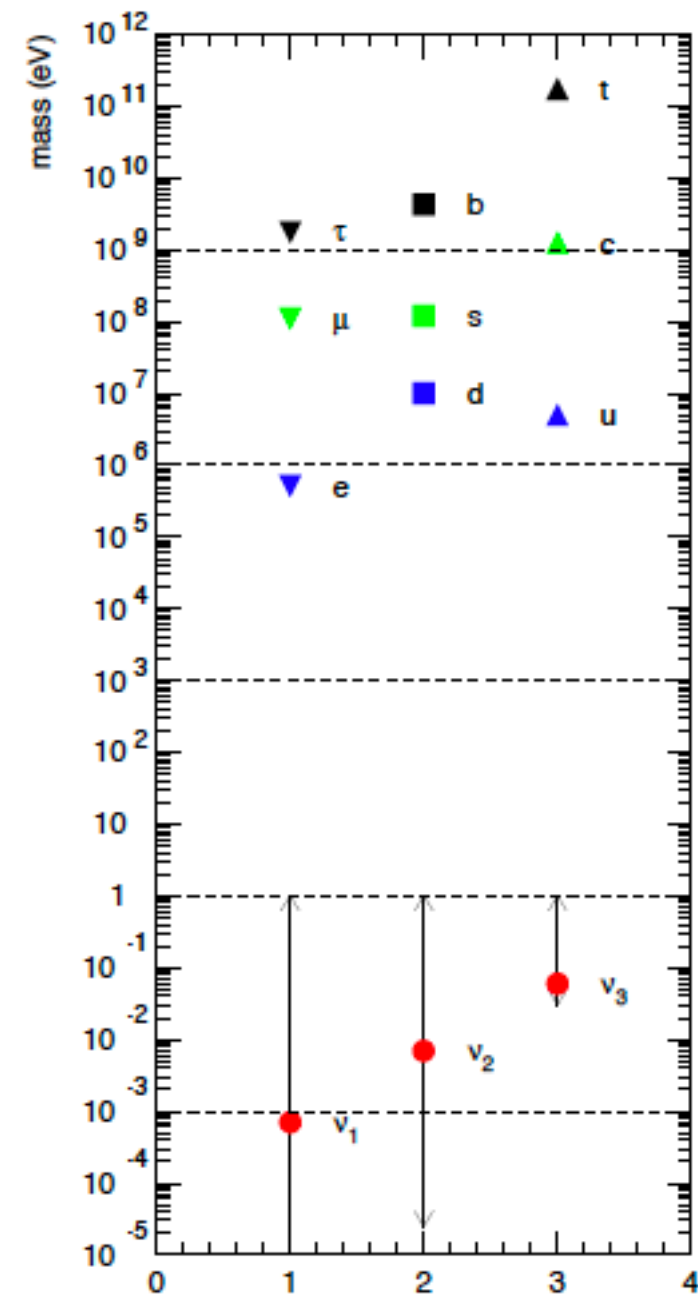
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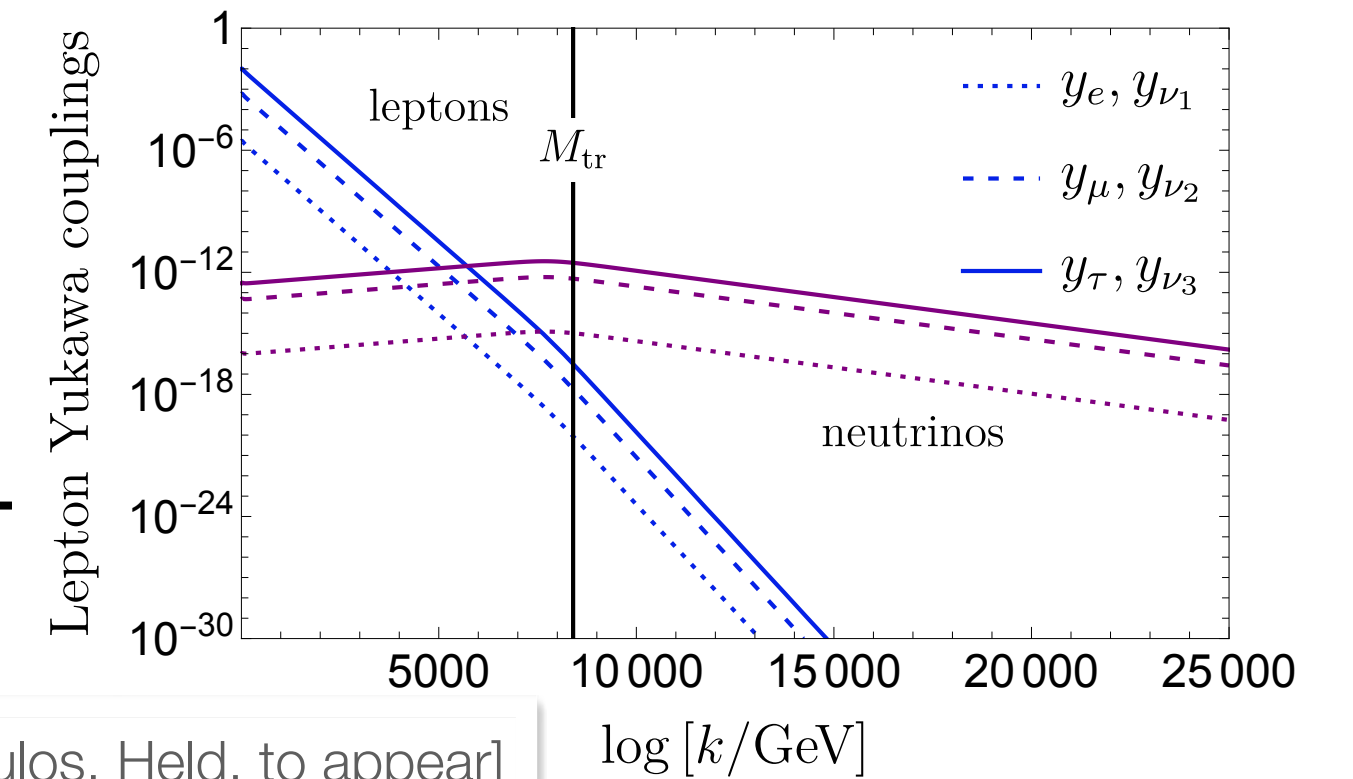
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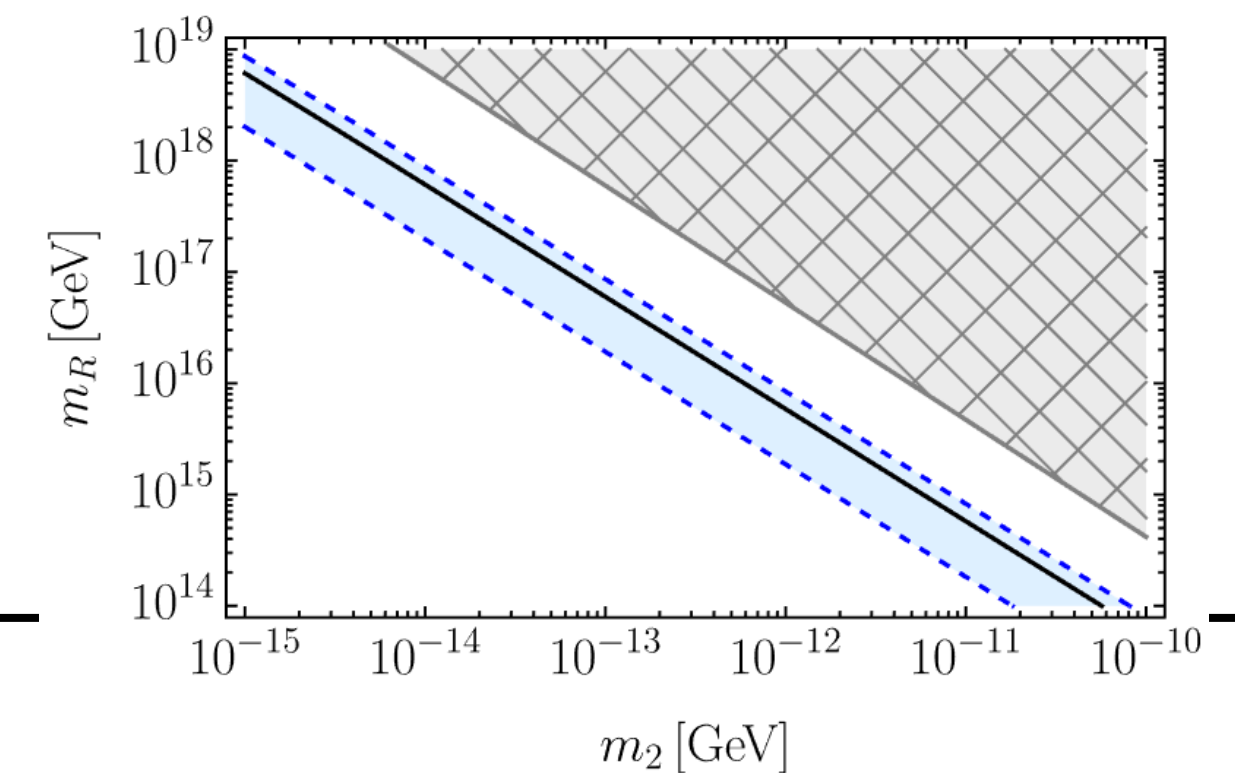
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Standard Model couplings

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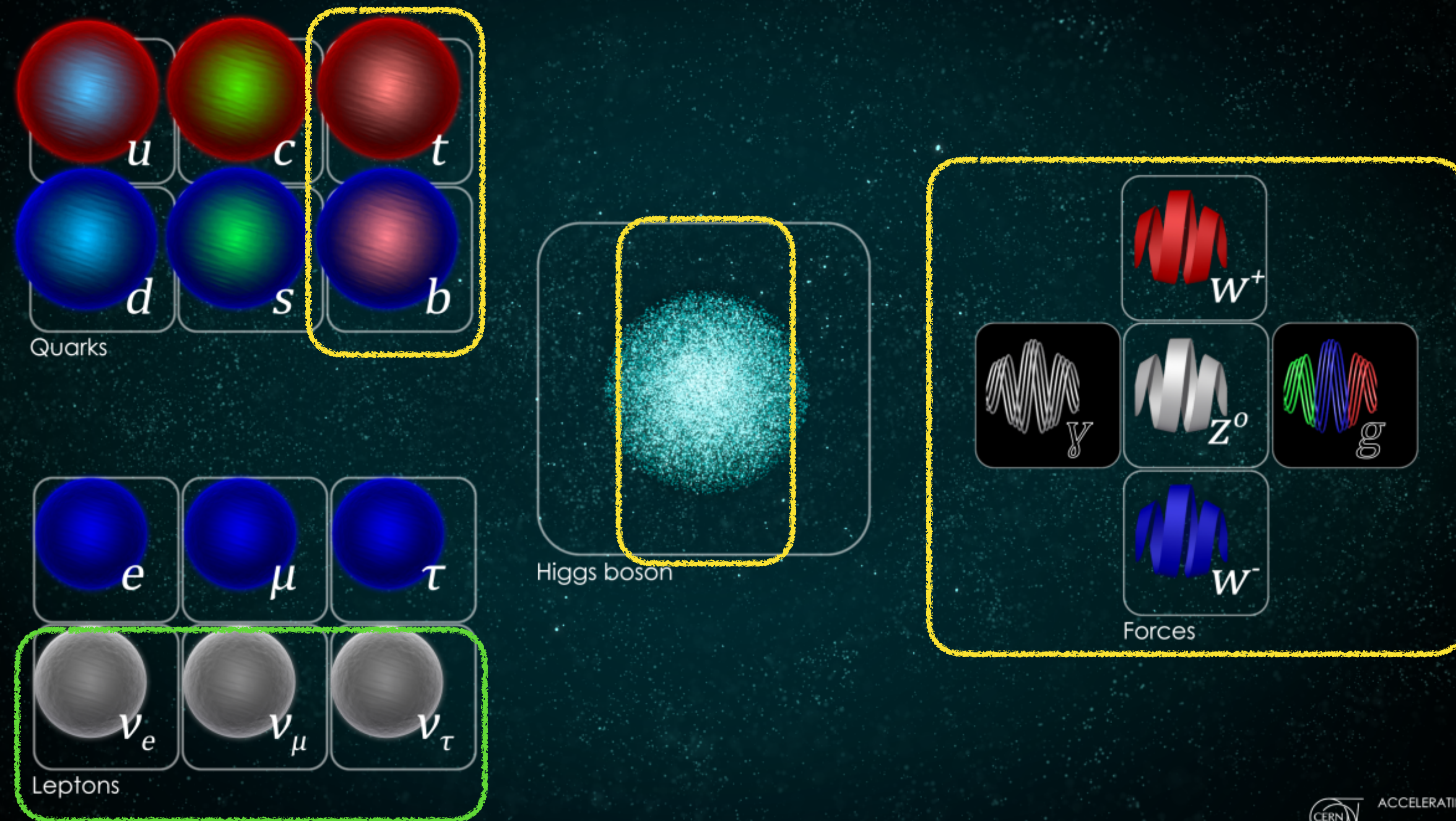
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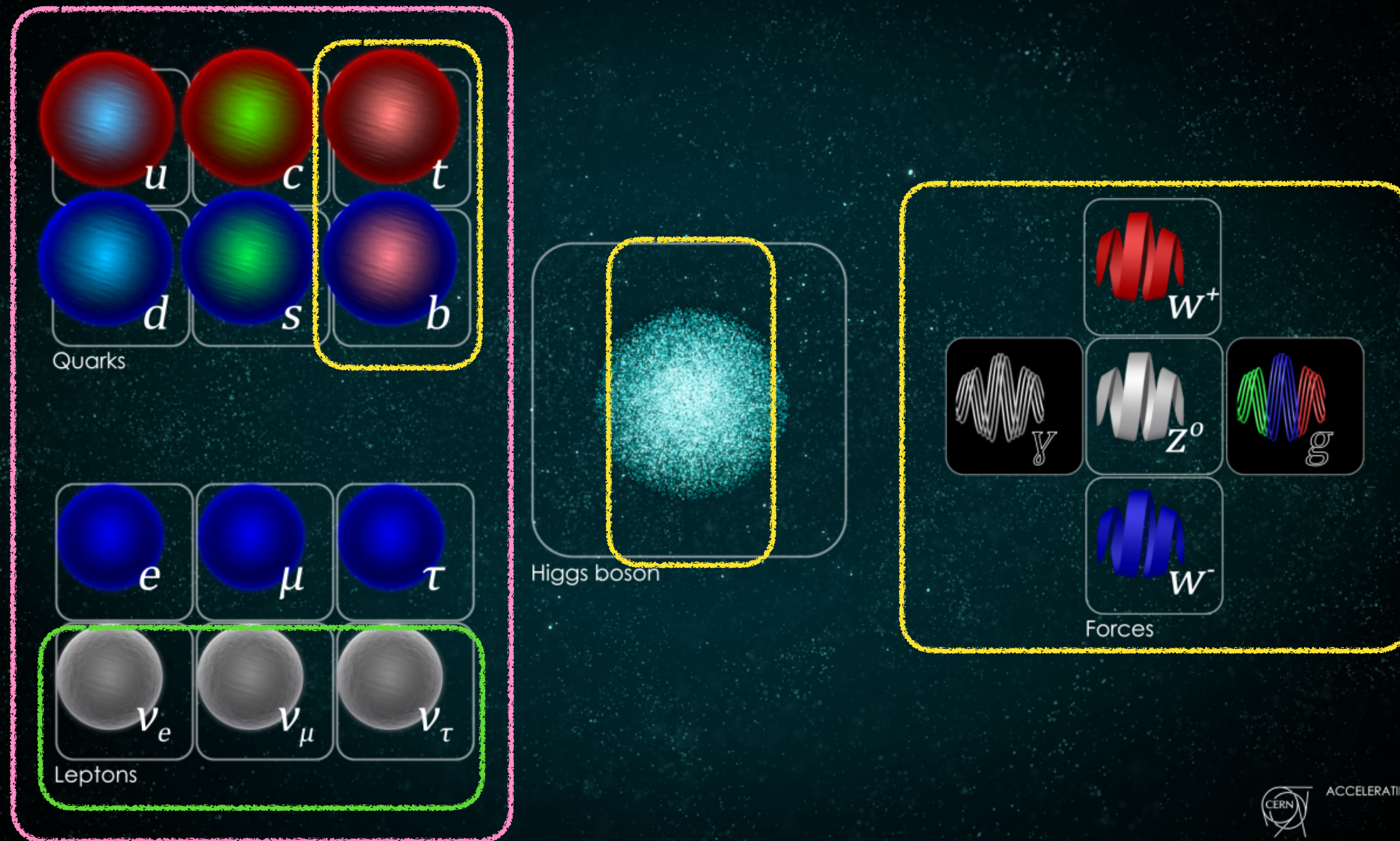
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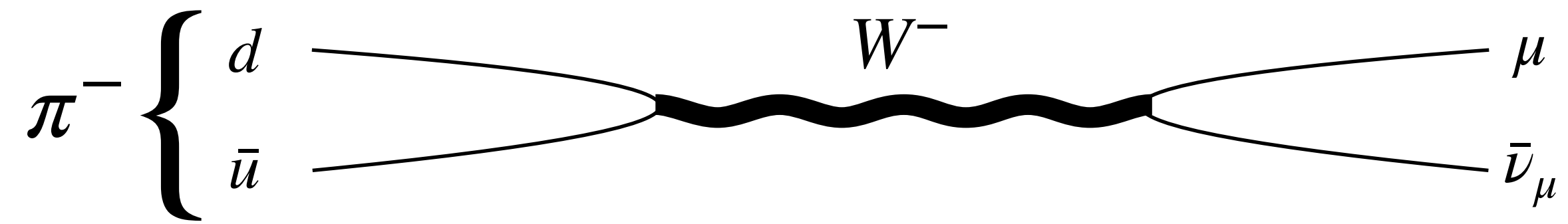
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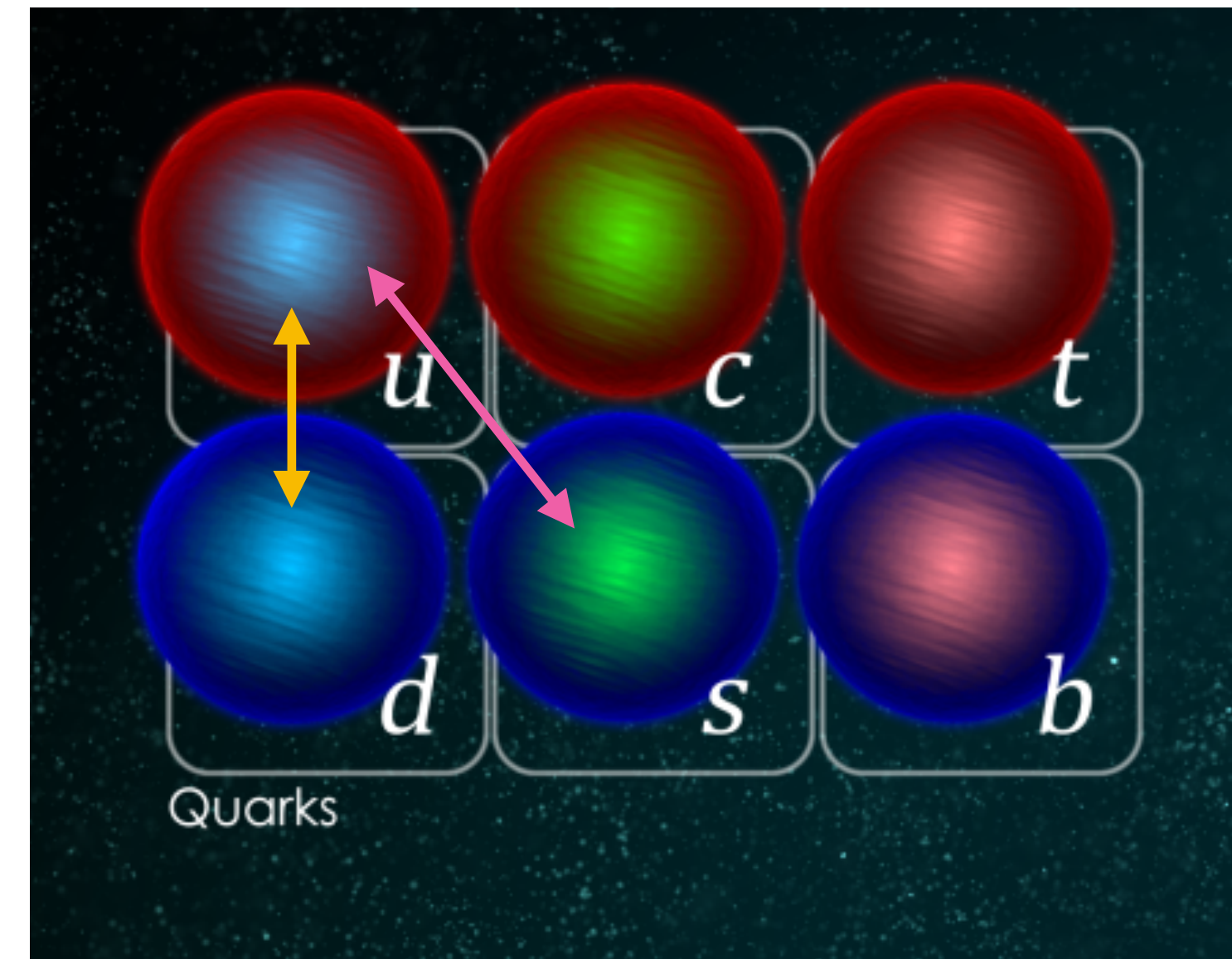
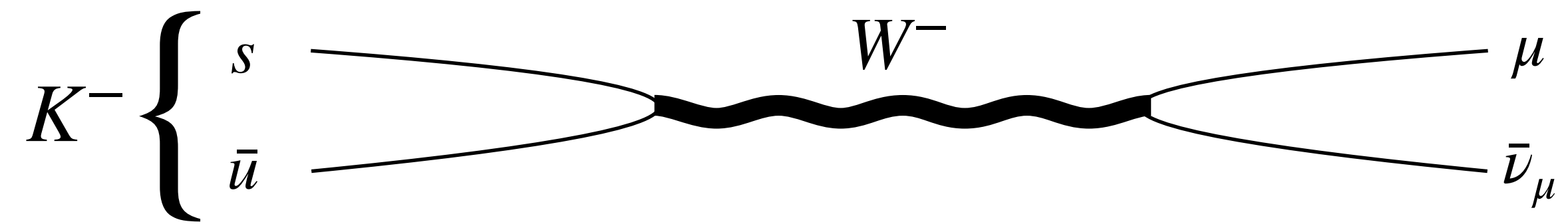
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Part 3: mixing matrices

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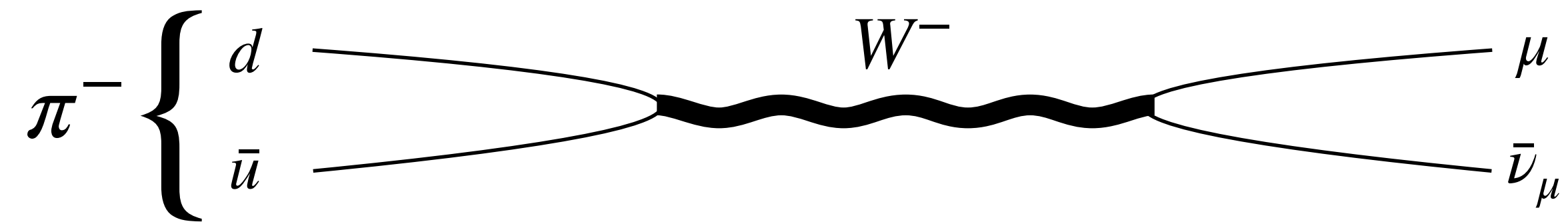


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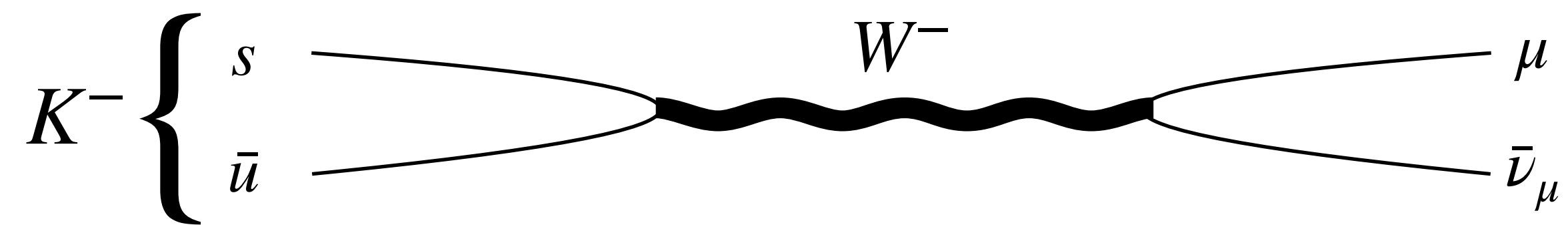


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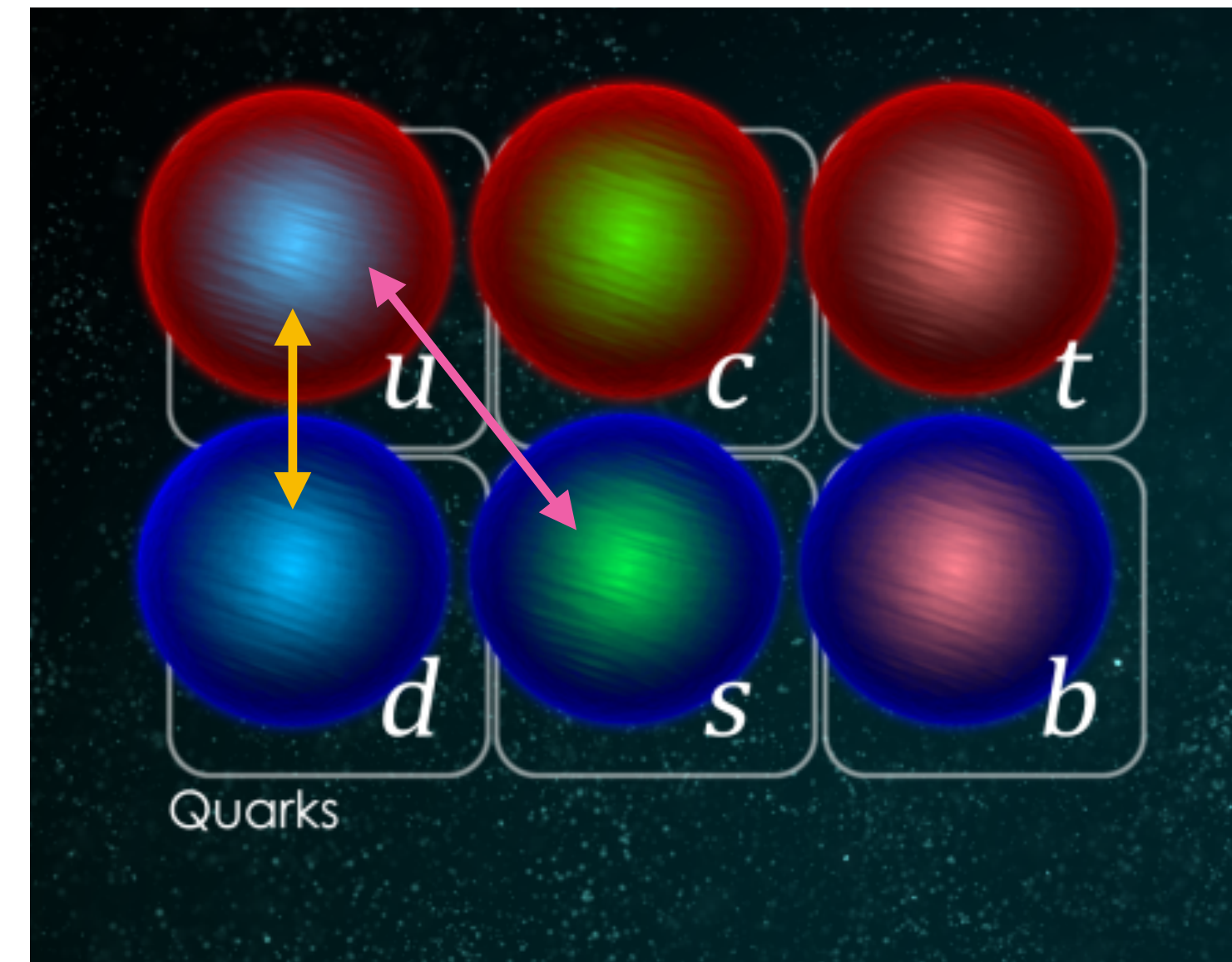
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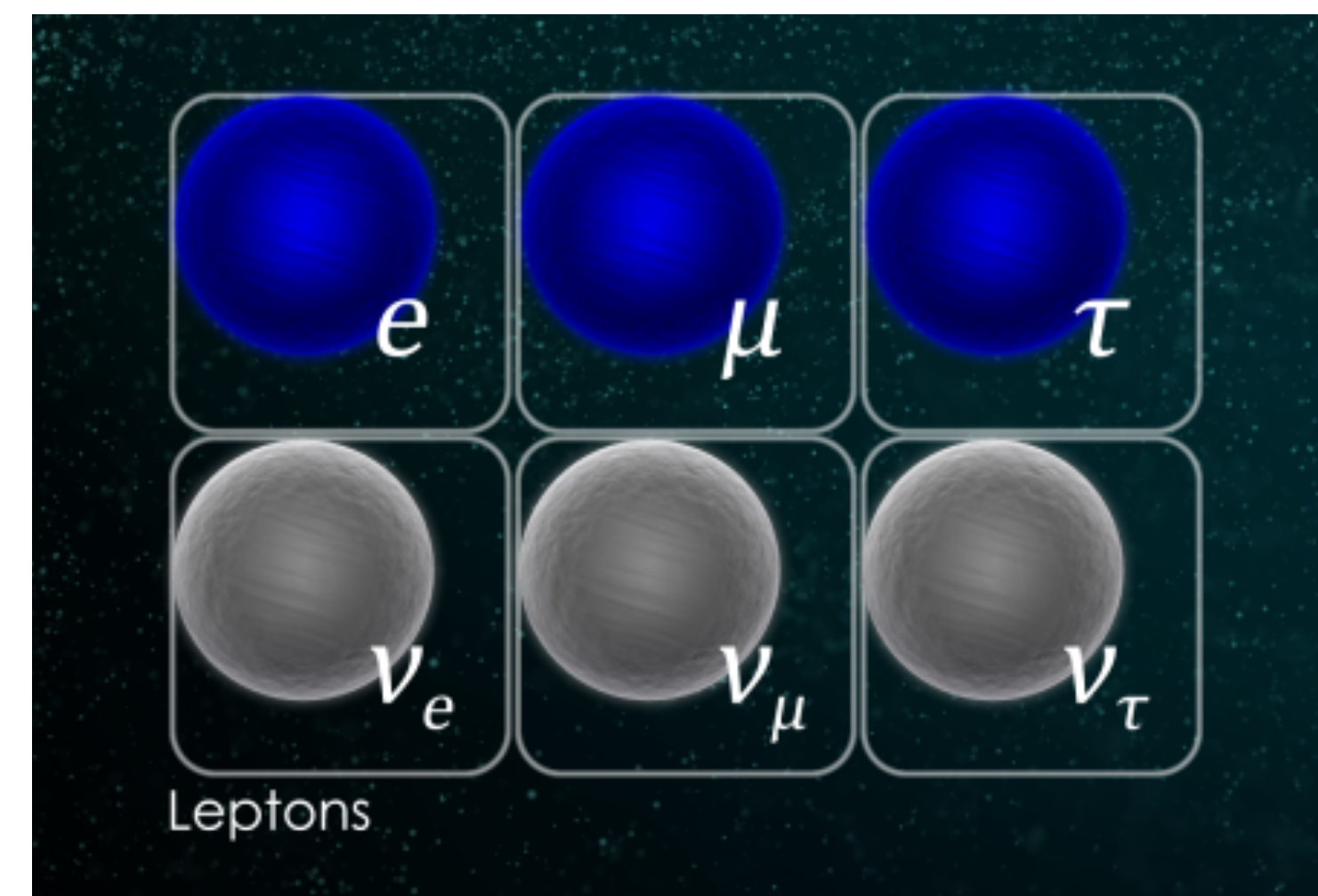
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CKM

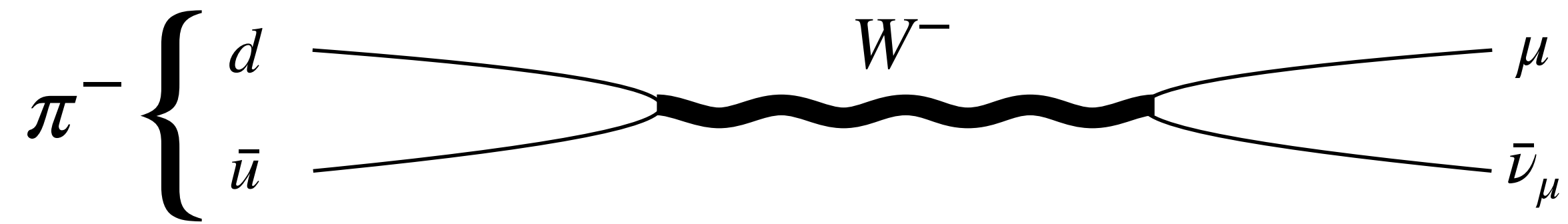


PMNS

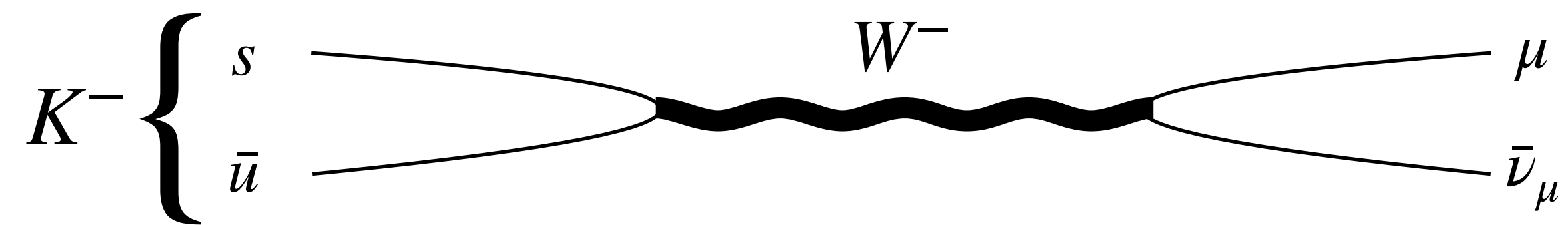


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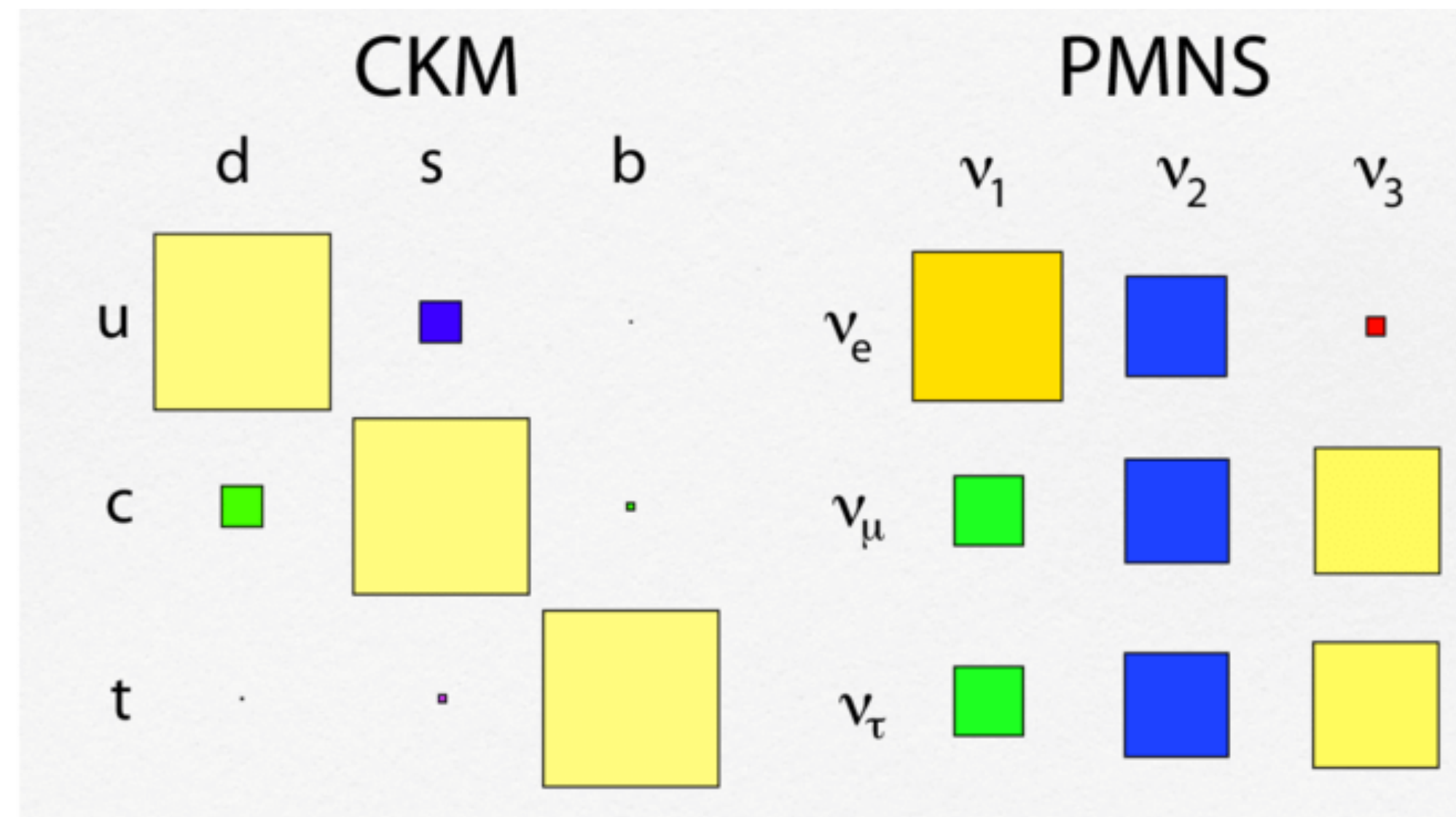
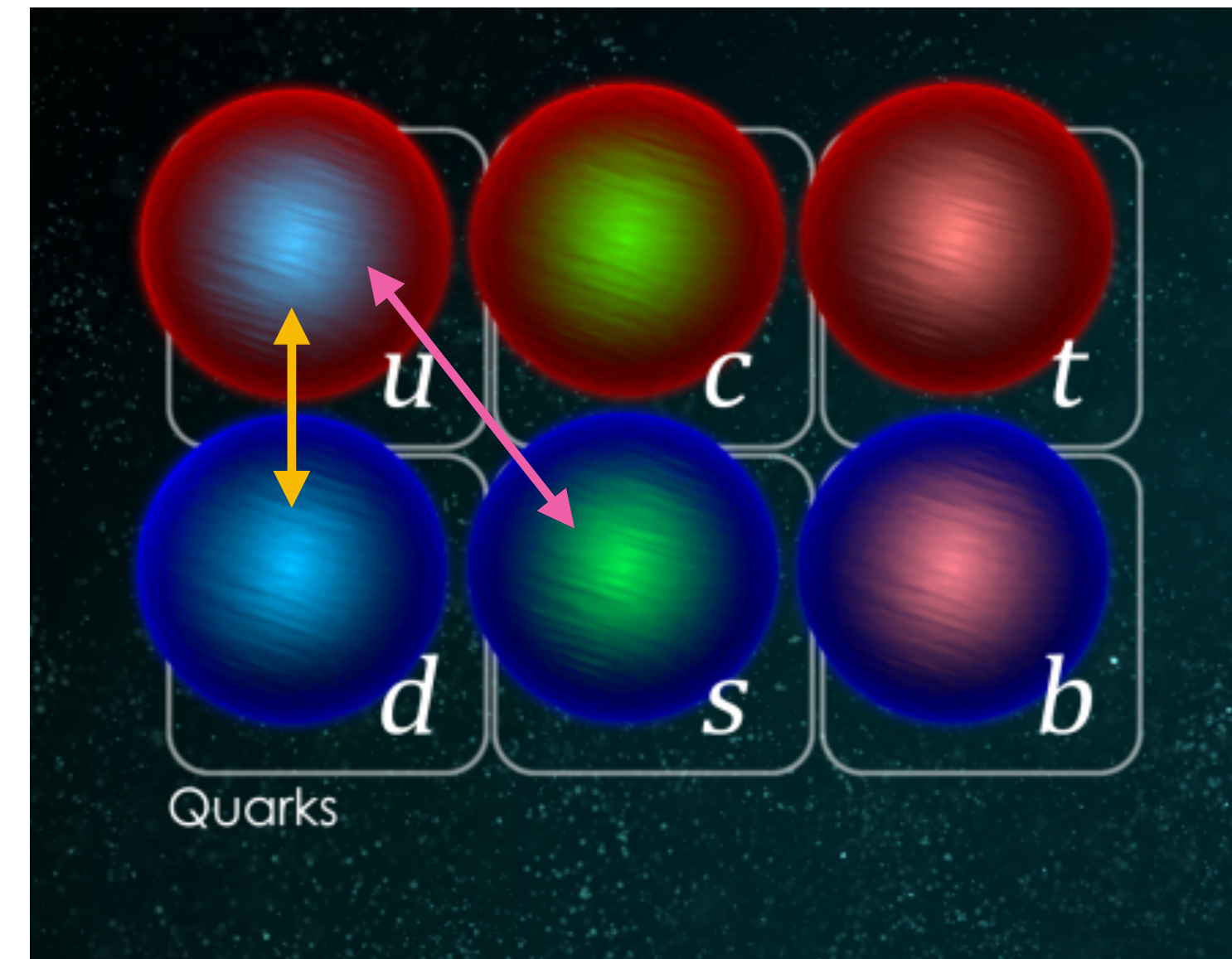
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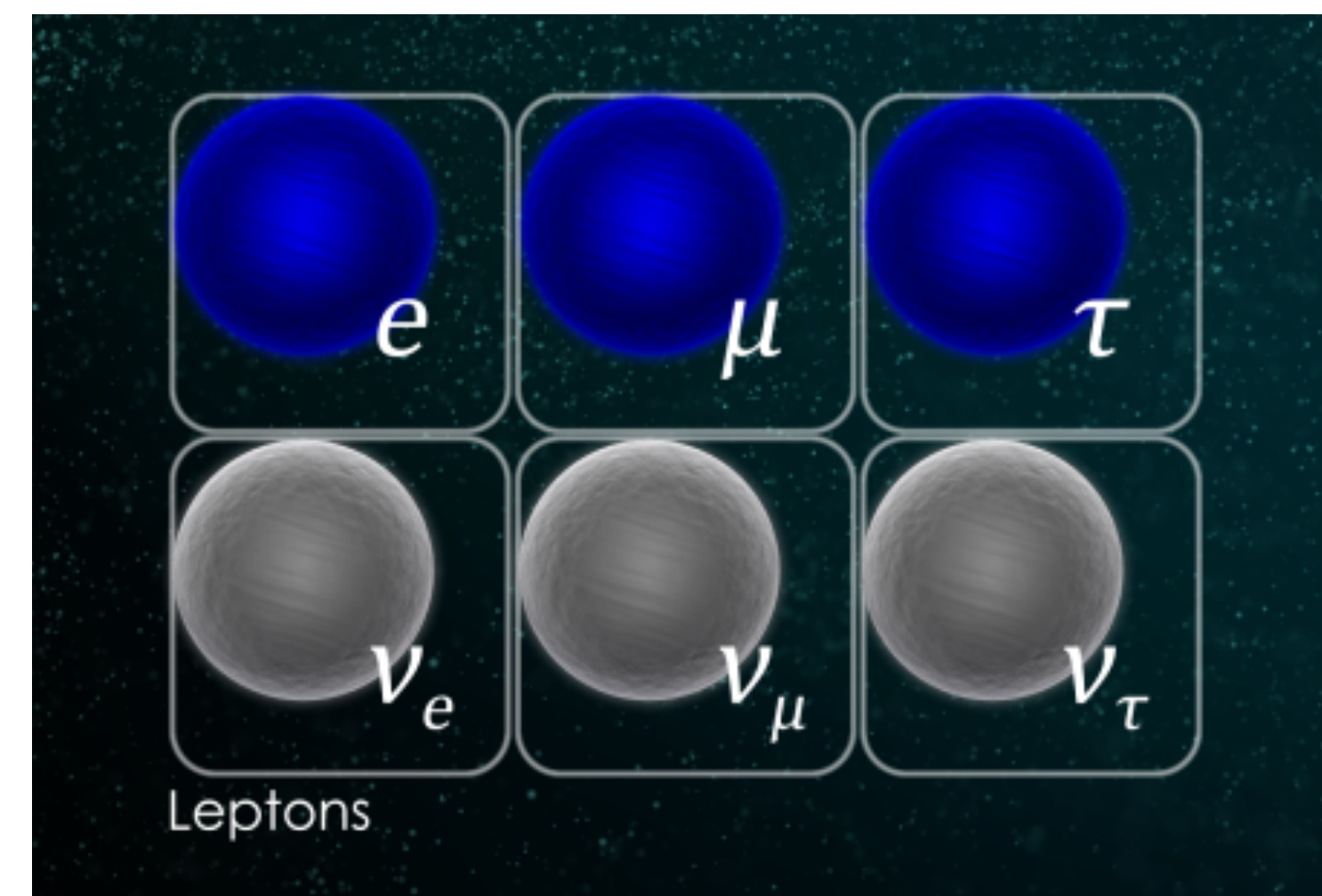
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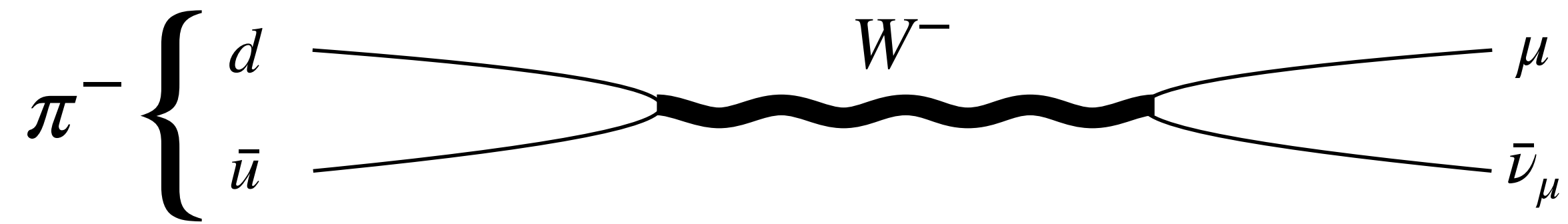


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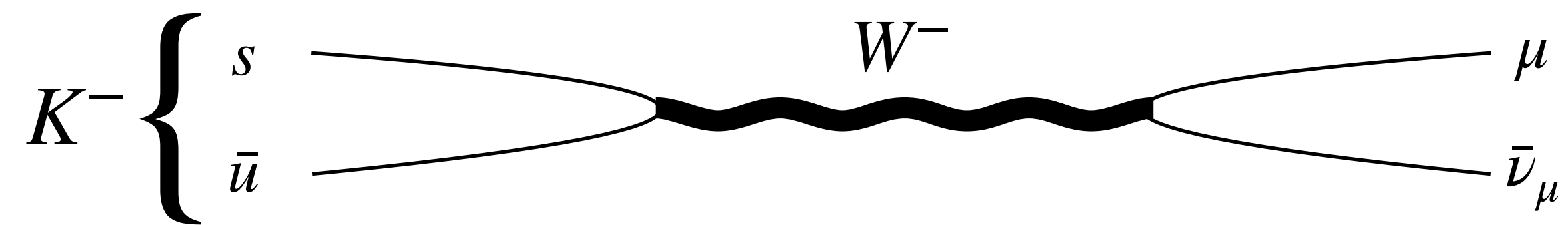


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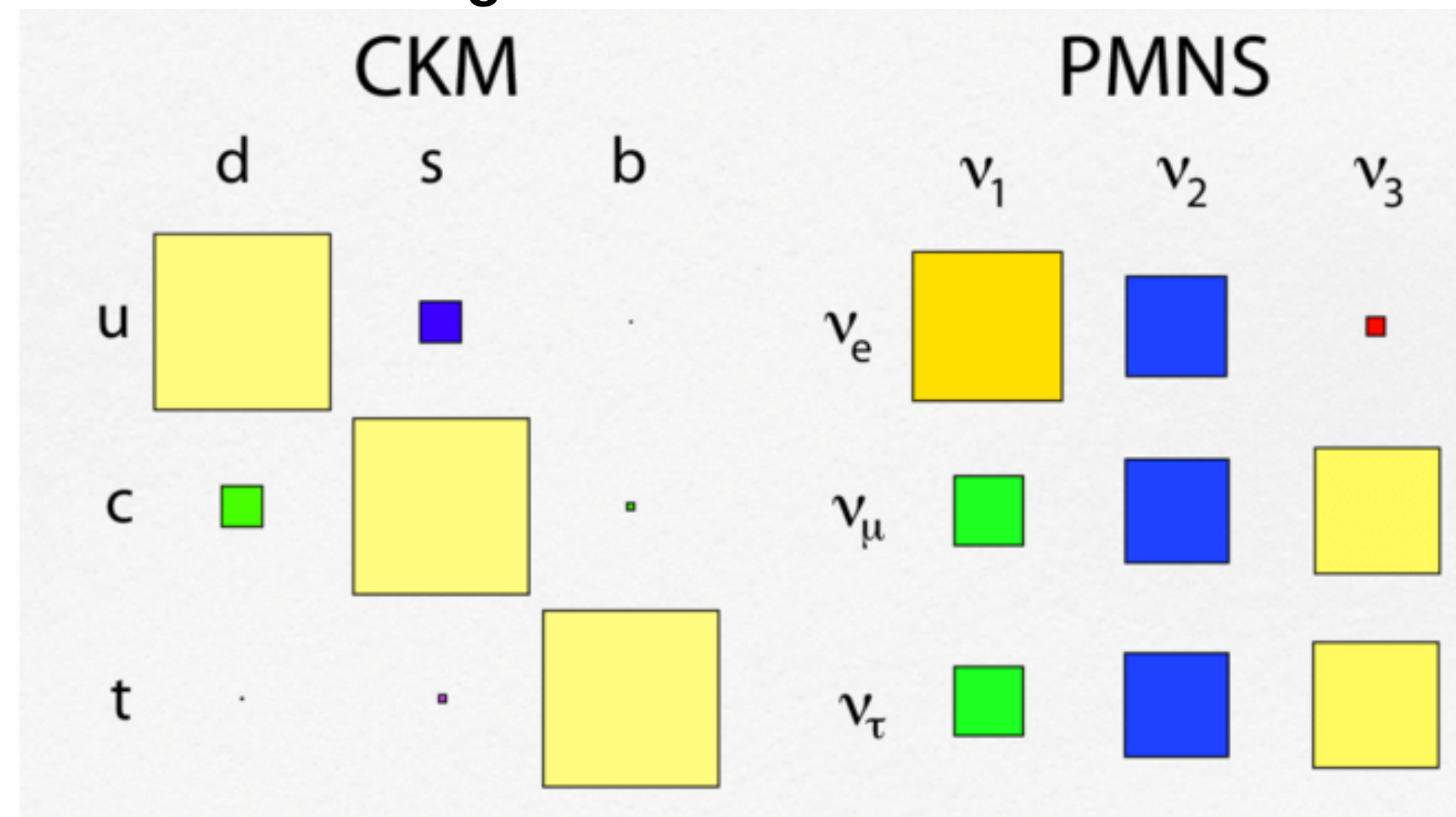
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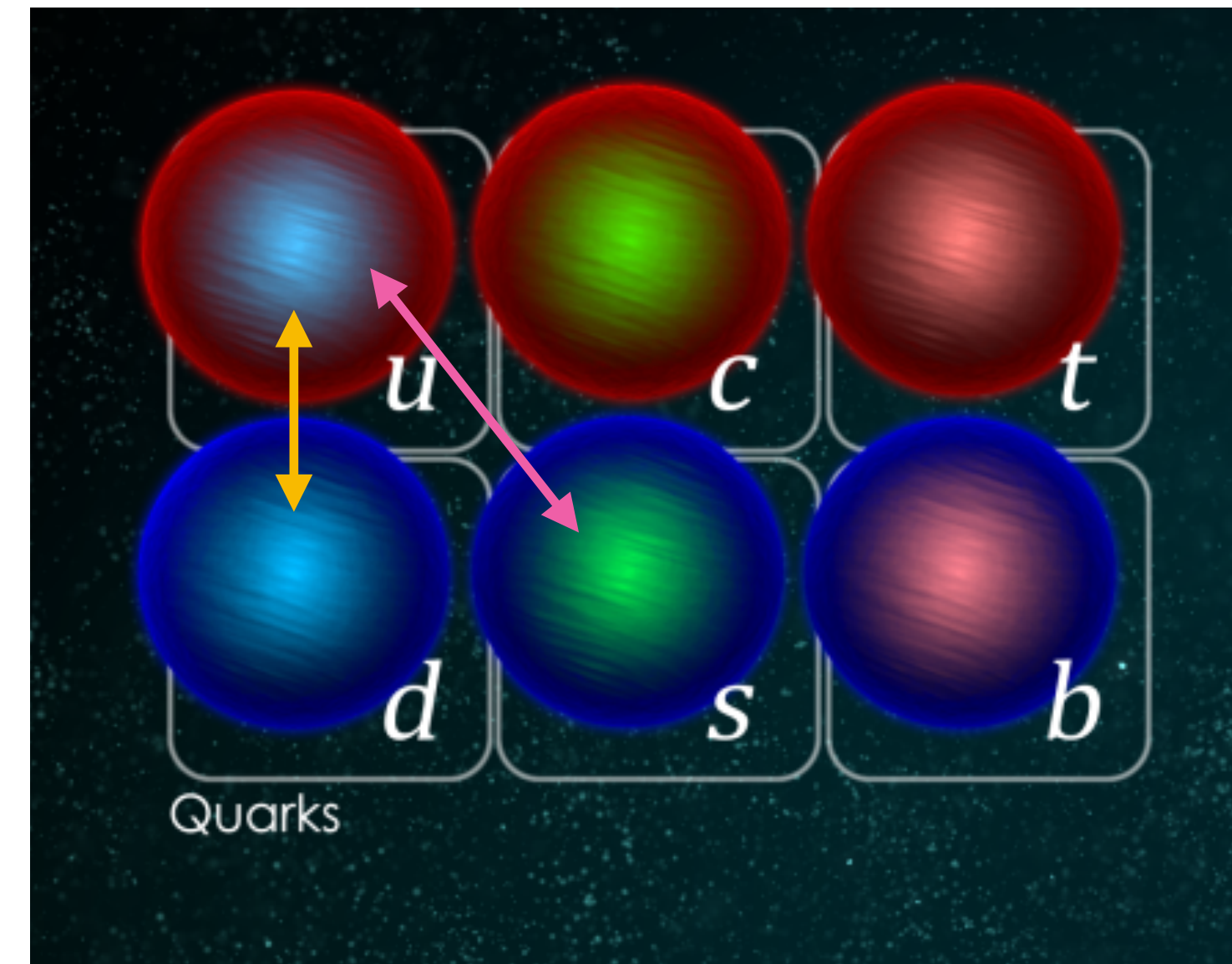
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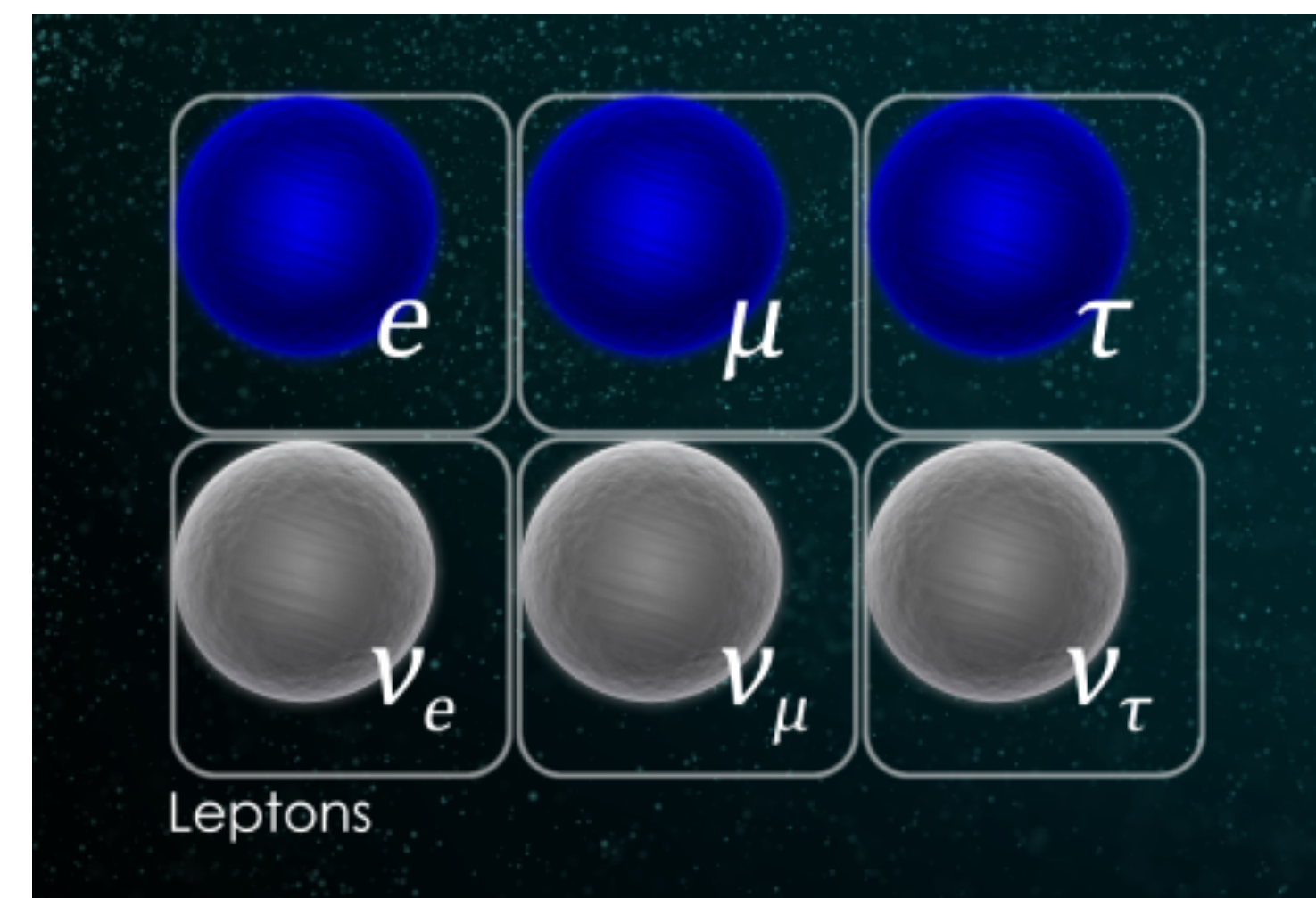
What generates these structures?



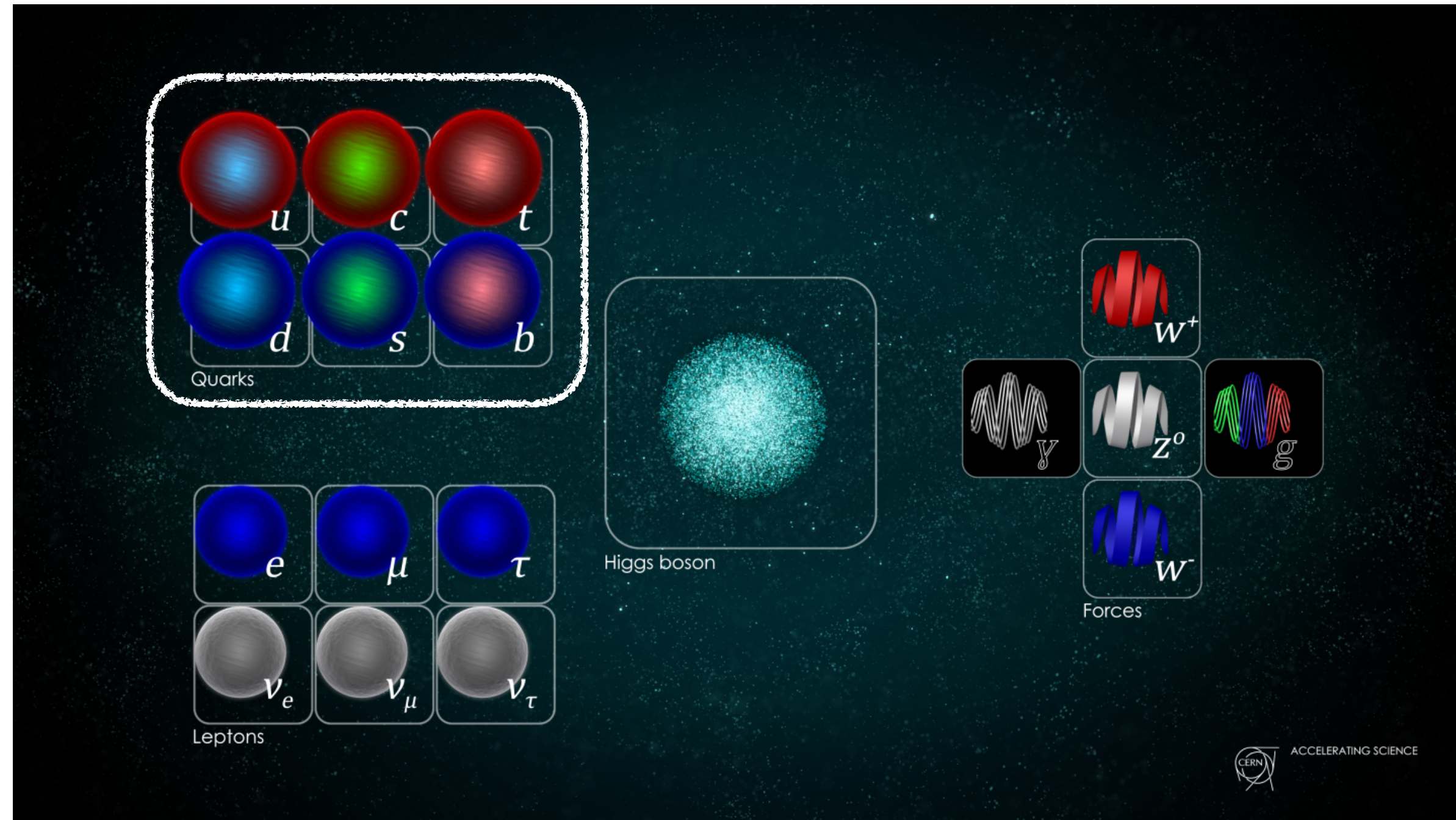
CKM



PMNS



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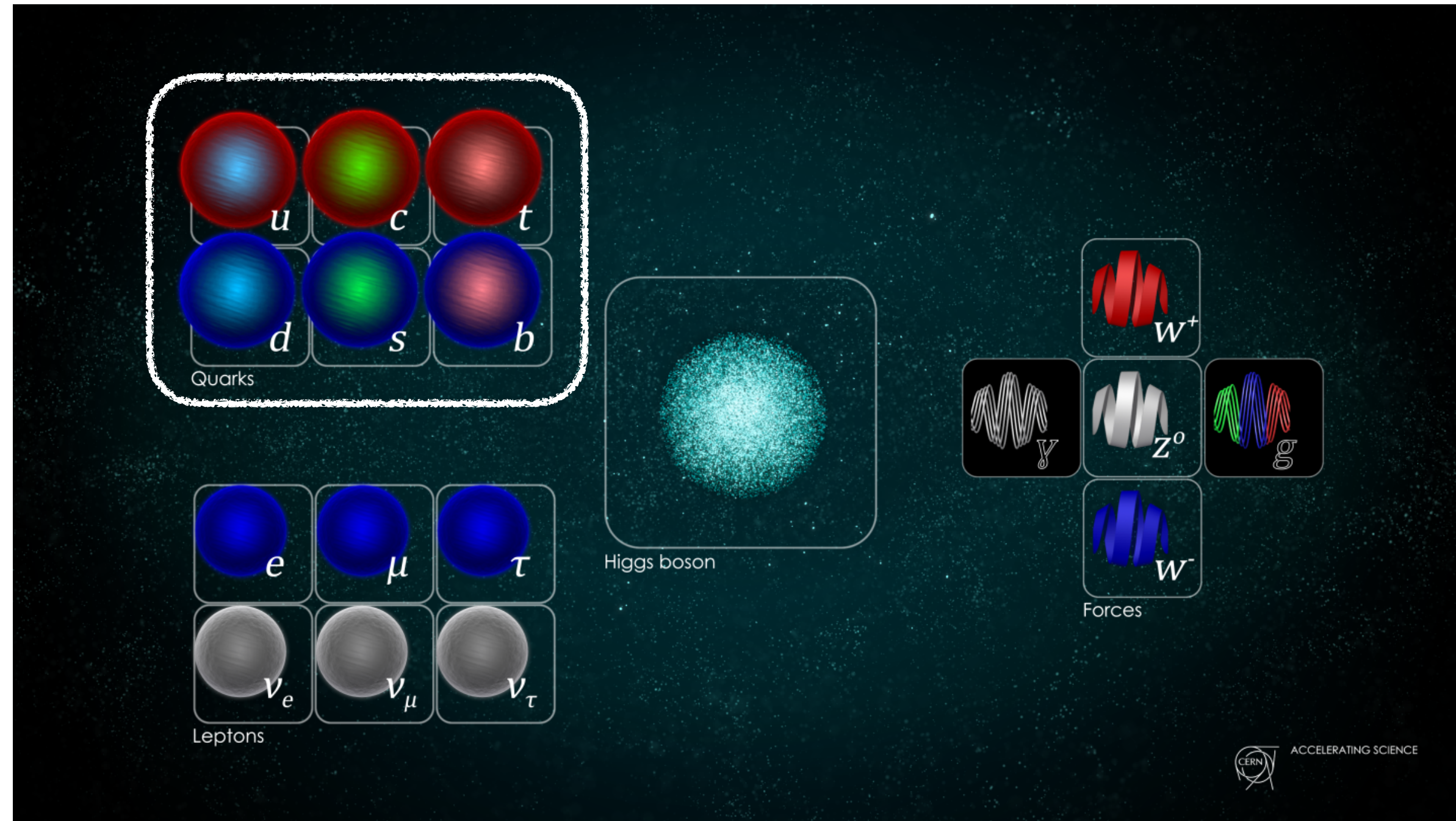
- Yukawa couplings (determine quark masses)
- CKM mixing matrix (determines mixing of flavors in weak interactions)

$$|V|^2 = \begin{pmatrix} X & Y & 1 - X - Y \\ Z & W & 1 - Z - W \\ 1 - X - Z & 1 - Y - W & X + Y + Z + W - 1 \end{pmatrix}$$

$$X = 0.94936 \pm 0.00031 \quad , \quad Y = 0.05063 \pm 0.00031$$

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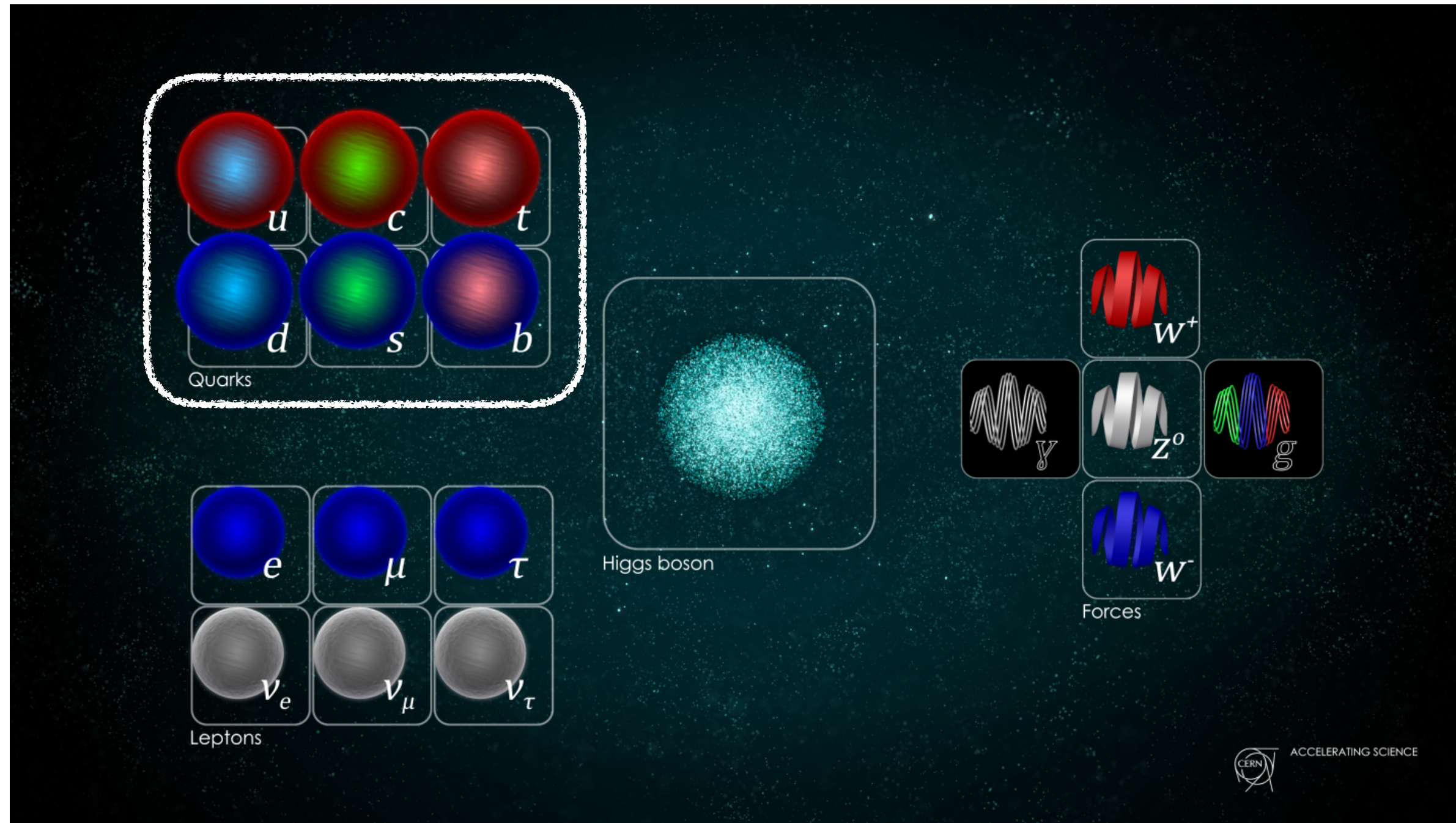
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observation: $X + Y = 0.99999 \pm 0.00044$, $W + Z = 0.99825 \pm 0.00044$.

satisfies: $X + Y = 1$, $W + Z = 1$ within 4σ .

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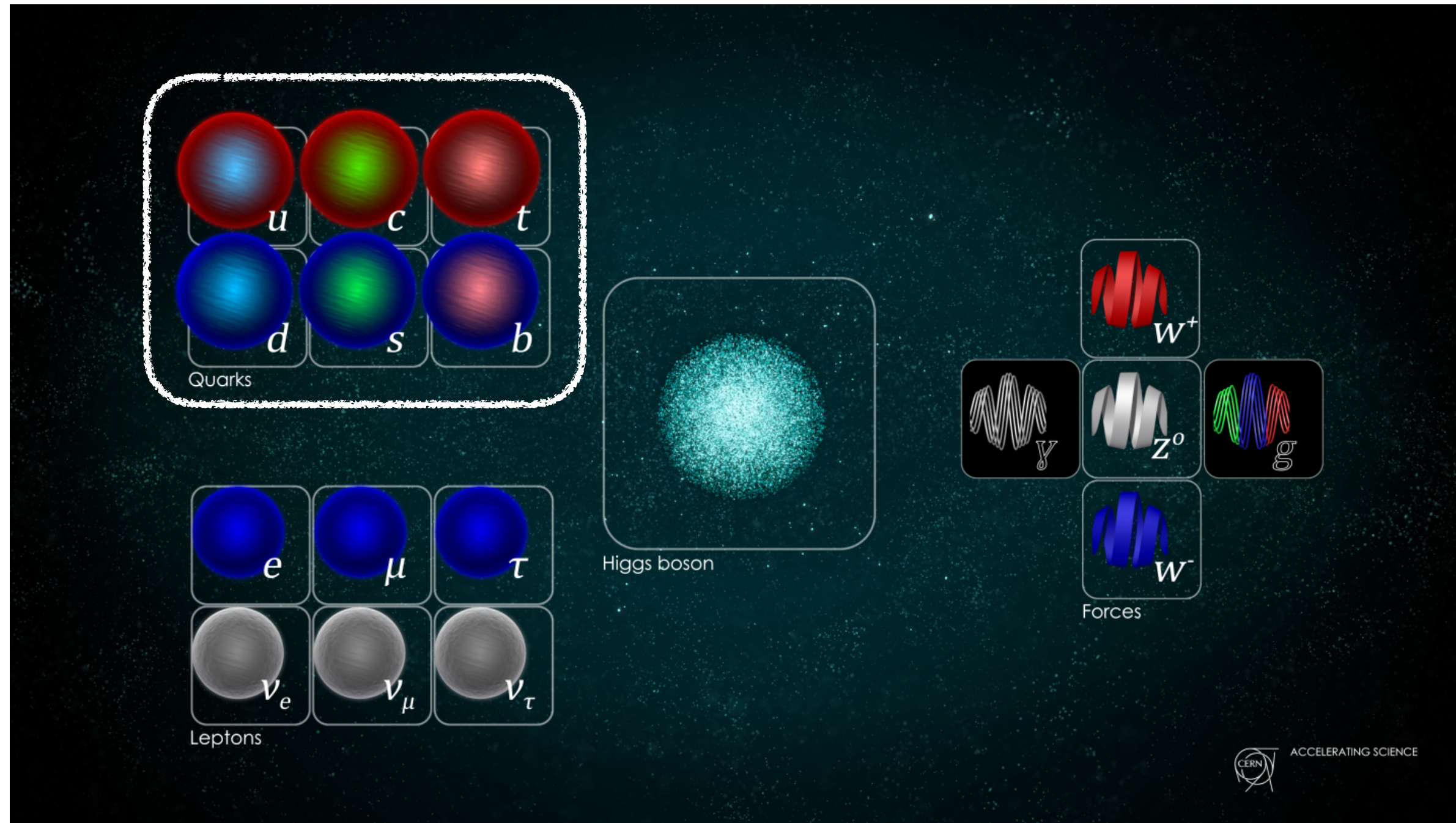
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$$\partial_t(X + Y - 1) = \frac{3}{16\pi^2} y_t^2 (X + Y - 1).$$

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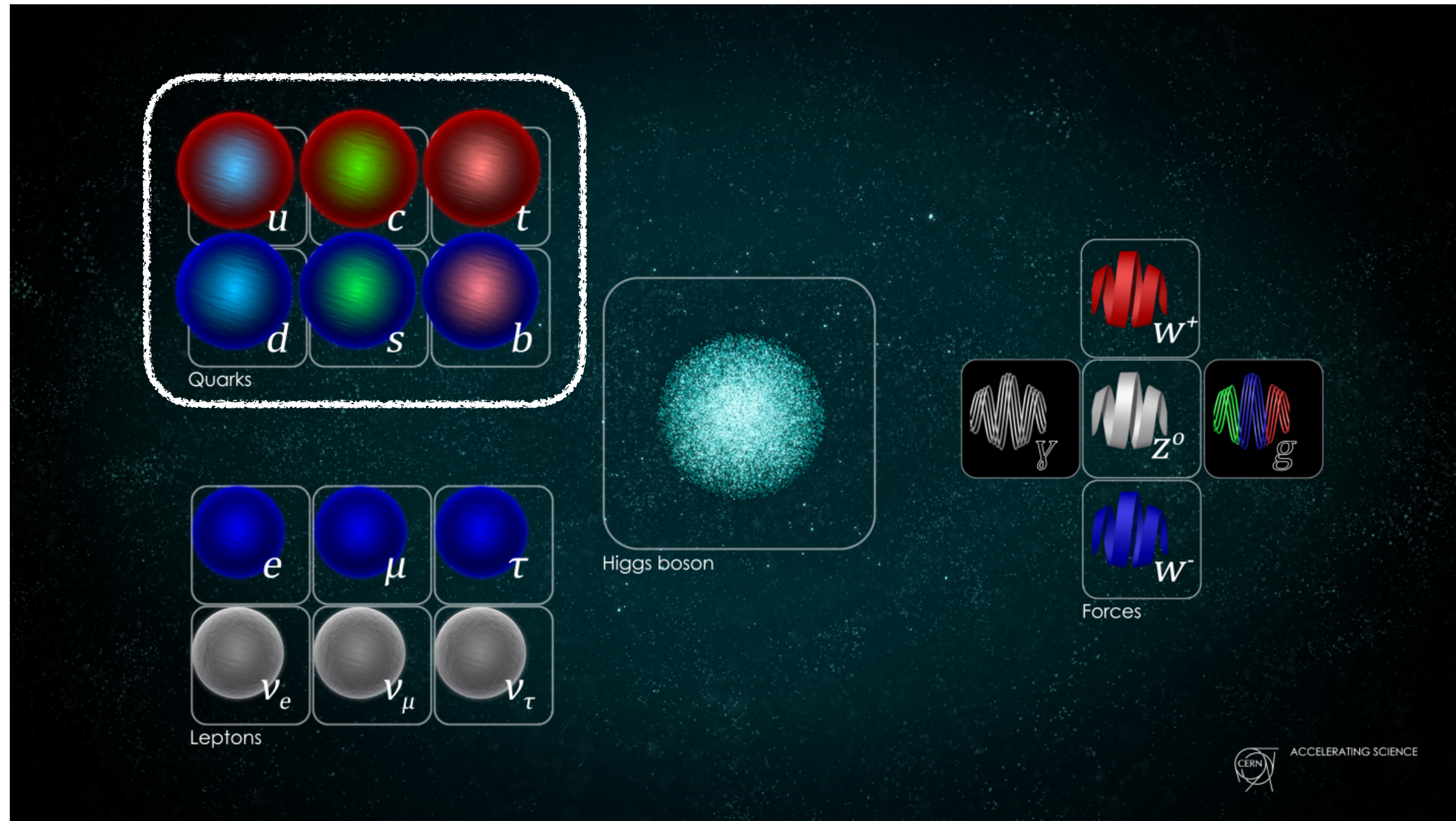
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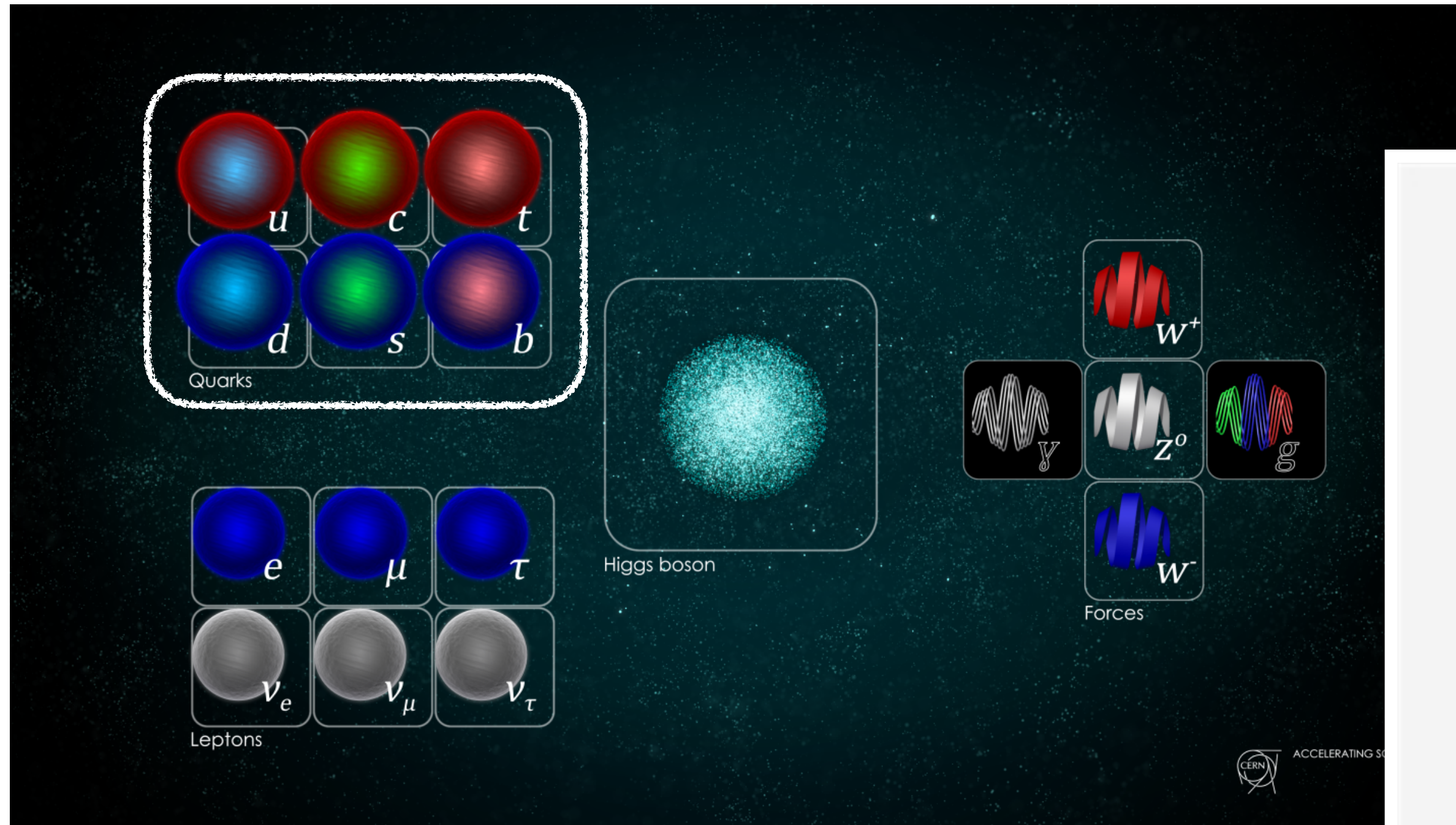
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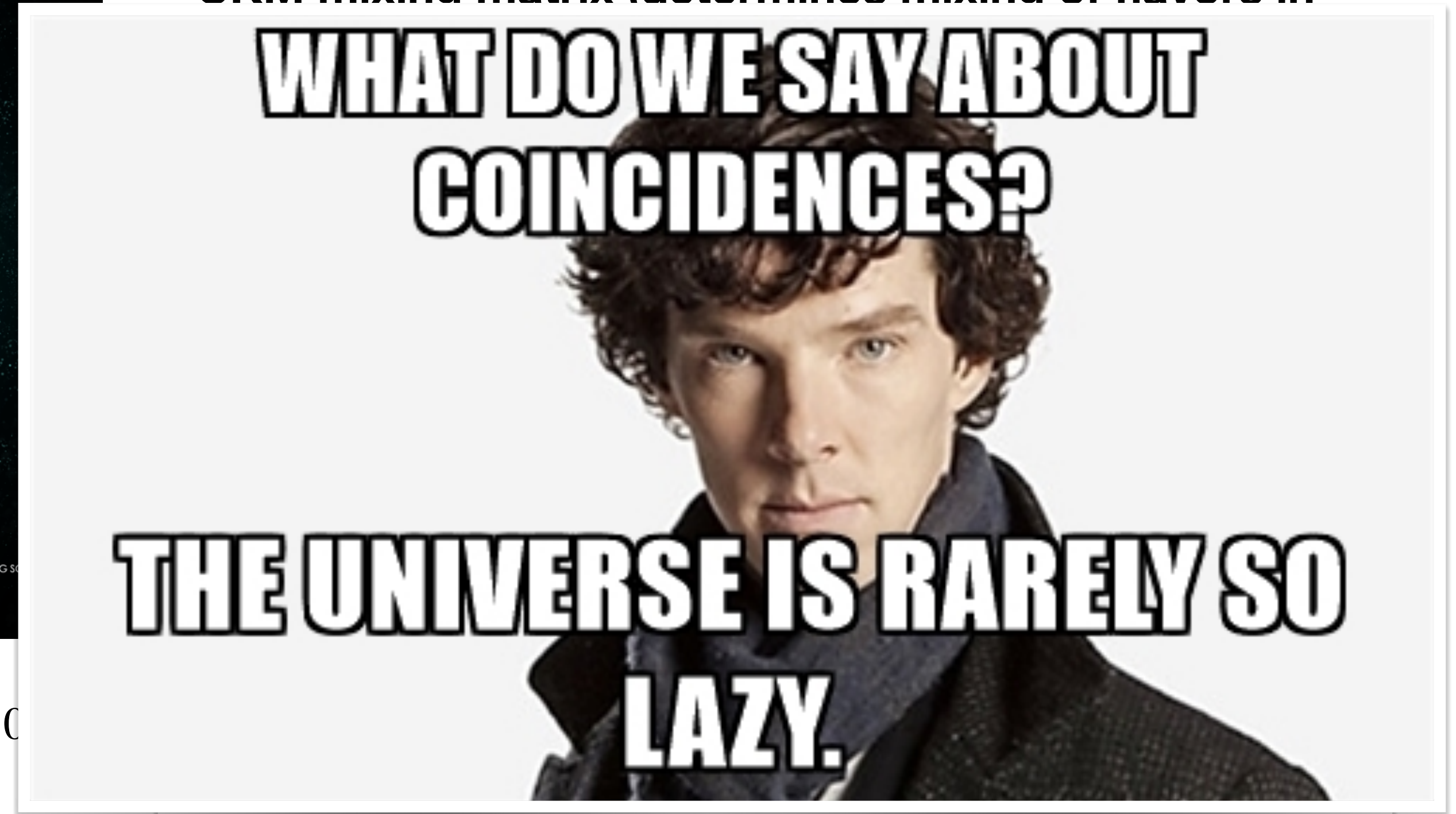
**Option 1): QFT has a cut-off at the Planck scale,
and the fact is just pure coincidence**

**Option 2): QFT holds far beyond the Planck scale
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Part 3: mixing matrices



- Yukawa couplings (determine quark masses)
- CKM mixing matrix (determines mixing of flavors in



observation: $X + Y = 0.99999 \pm 0.00044$, $W + Z = 0.99825 \pm 0.00044$

satisfies: $X + Y = 1$, $W + Z = 1$ within 4σ .

fact: [Alkofer, AE, Held, Nieto, Percacci, Schröfl '19]

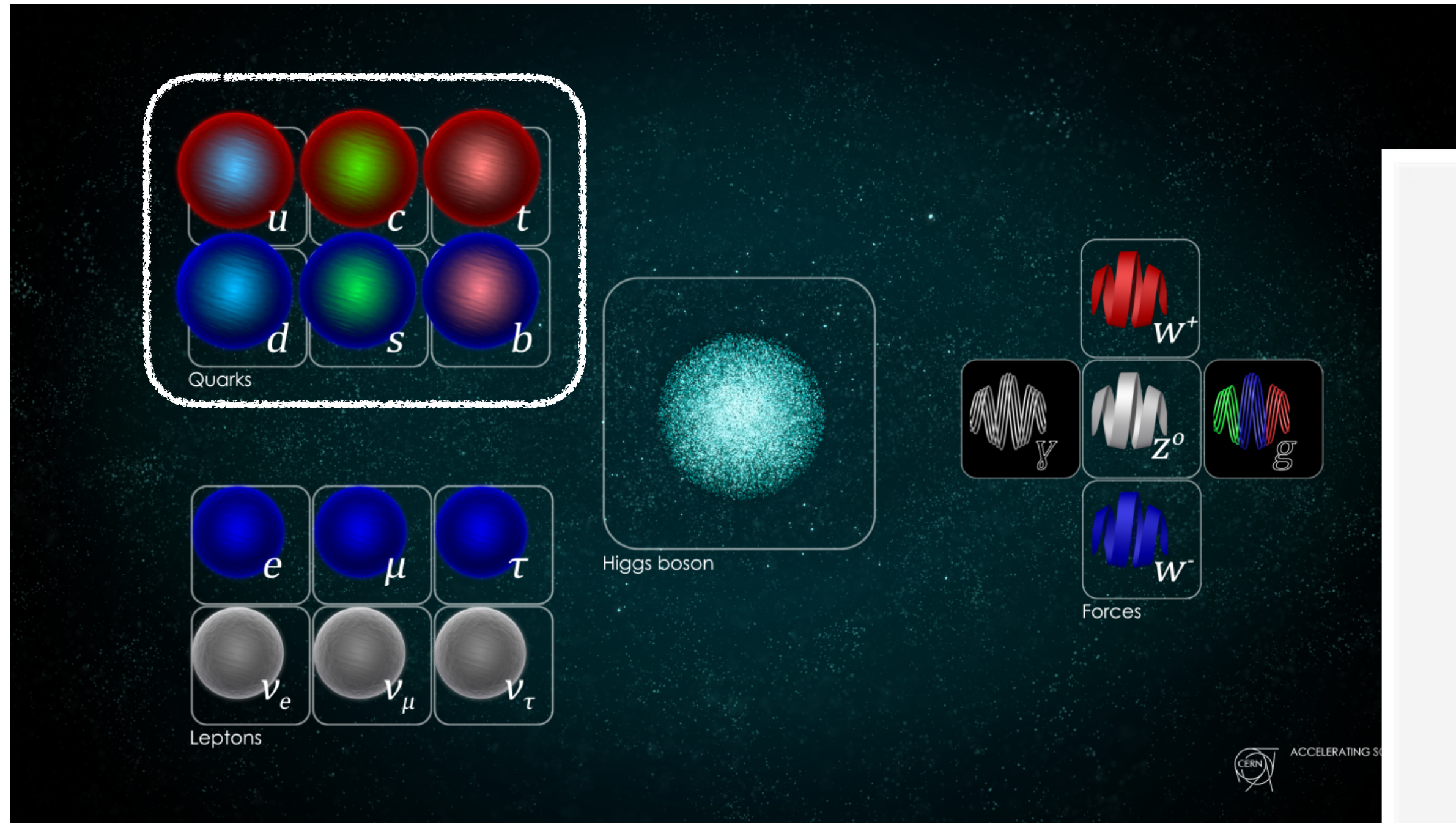
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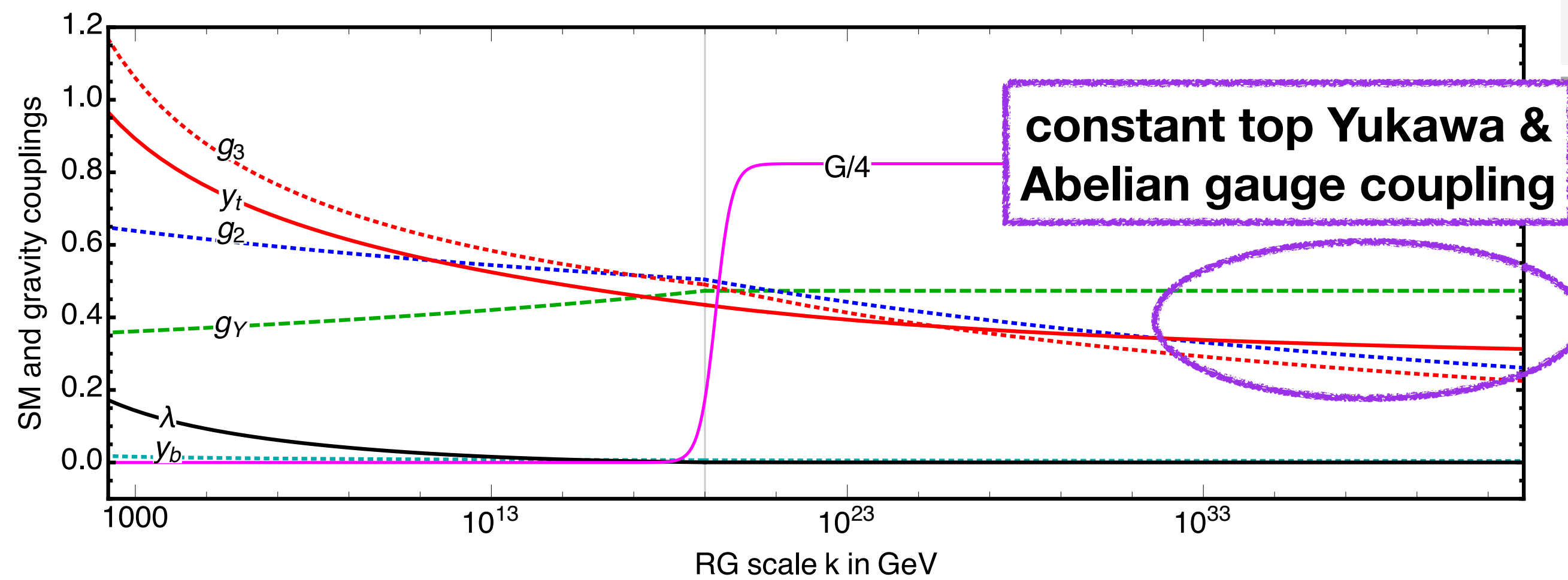
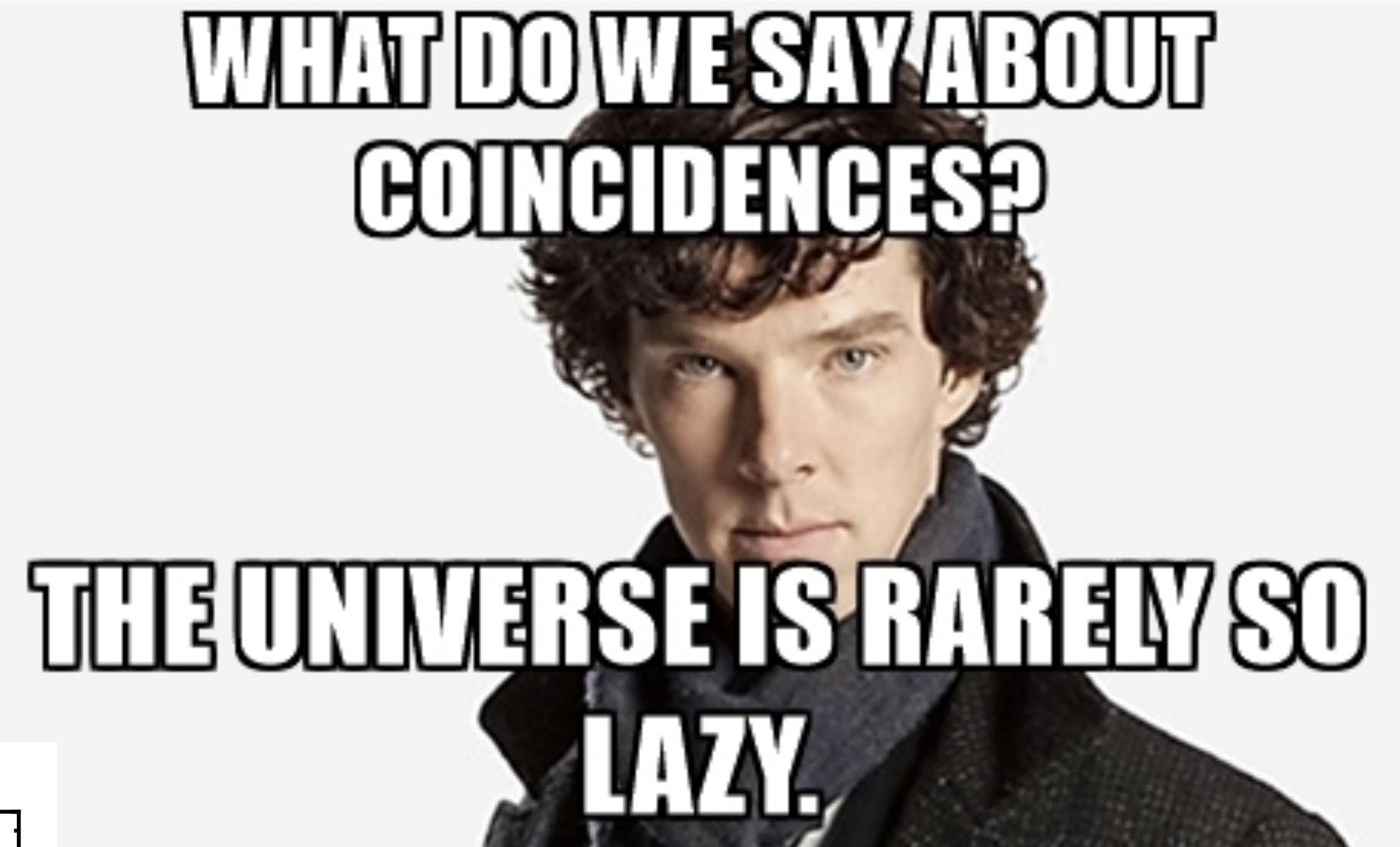
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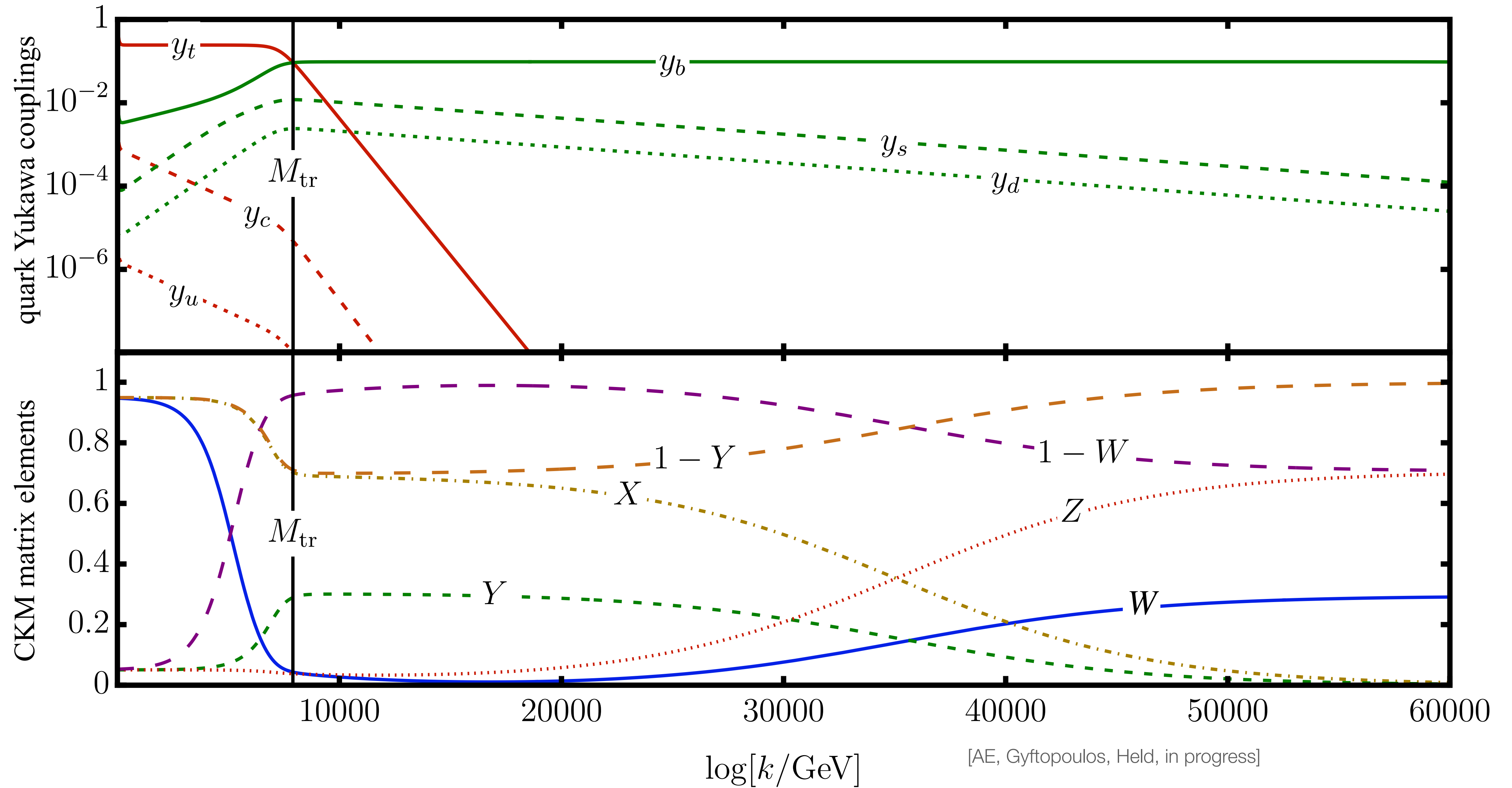
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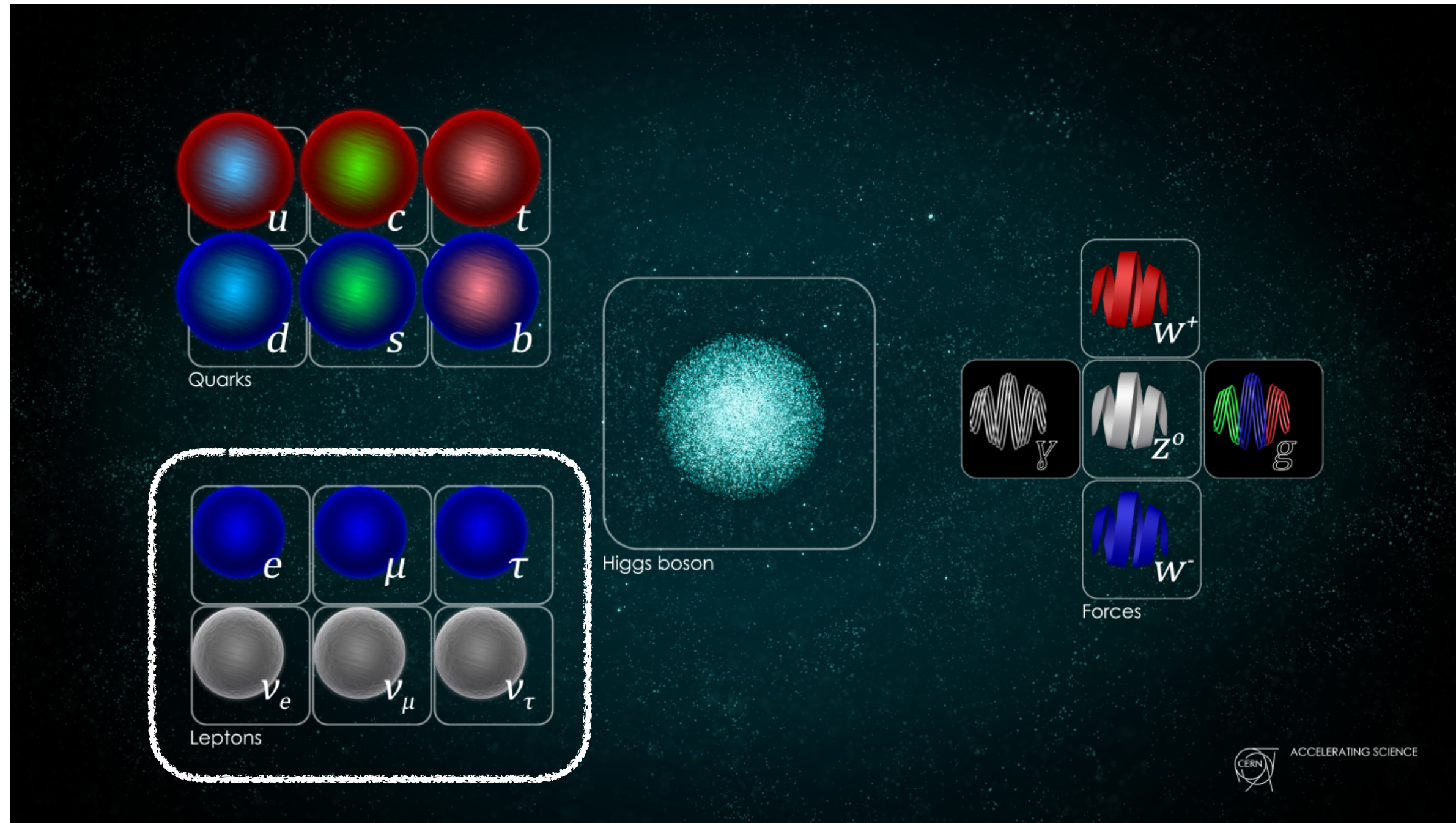
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Part 3: mixing matrices



- PMNS mixing matrix (determines mixing of leptons in weak interactions)

$$|V_{PMNS}| \approx \begin{pmatrix} 0.82343 & 0.54806 & 0.14697 \\ 0.47366 & 0.61638 & 0.62907 \\ 0.31243 & 0.56543 & 0.76333 \end{pmatrix}$$

observation: very far from fixed lines

But: $X + Y = 1$, $W + Z = 1$.

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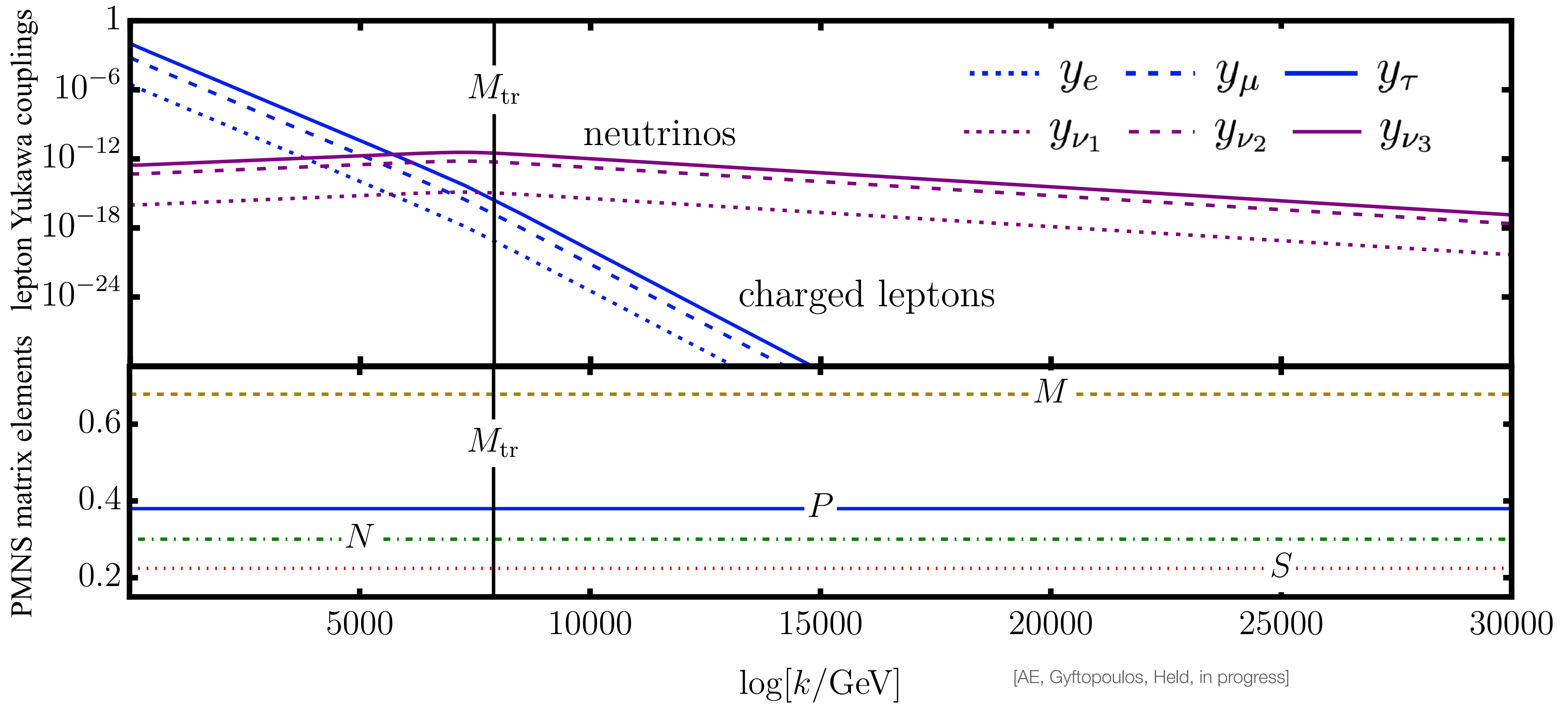
during which $y_\tau \gg y_{\mu, e, \nu_\tau, \nu_\mu, \nu_e}$ and $y_\tau = \text{const}$

Renormalization Group flow (simplified):

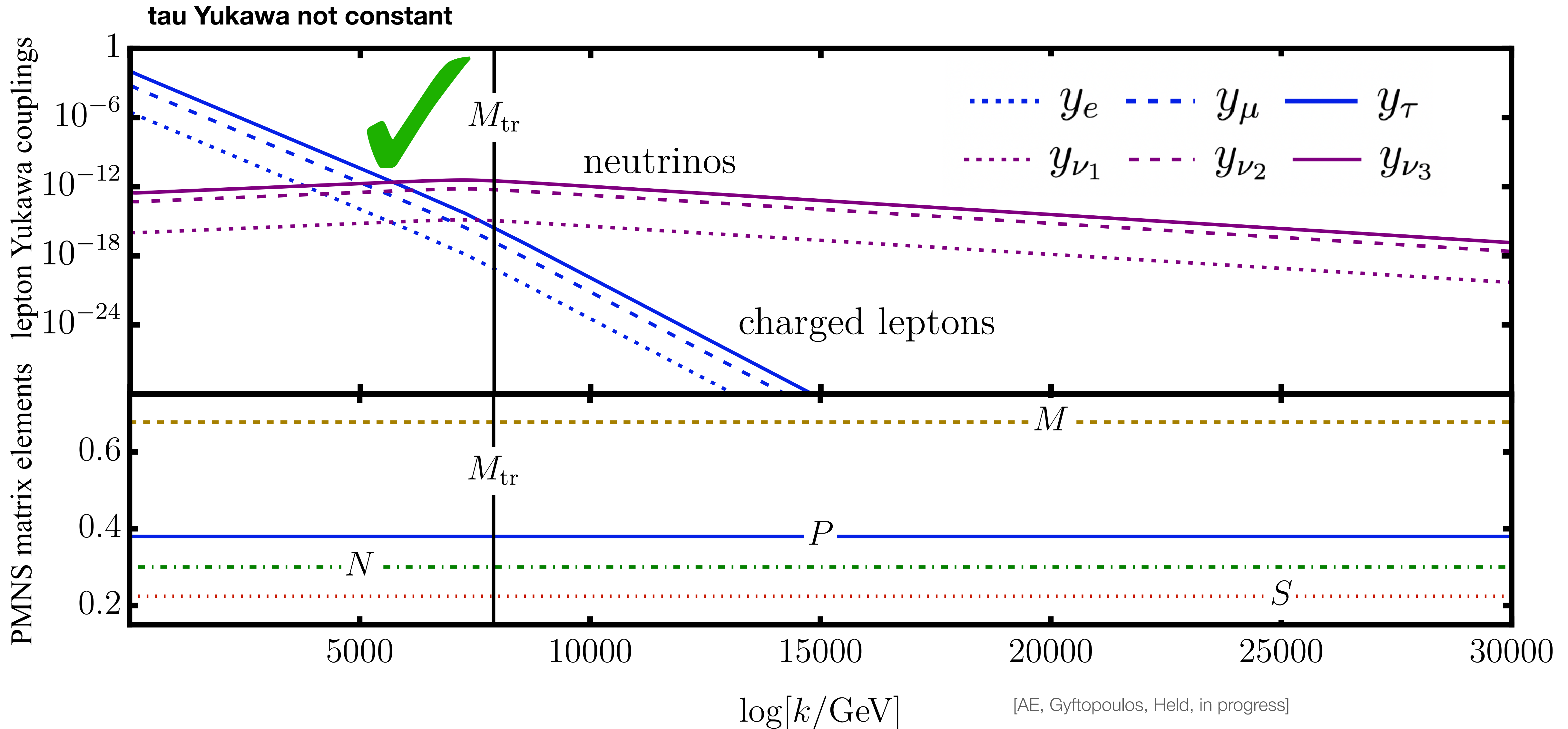
$$\partial_t(X + Y - 1) = \frac{3}{16\pi^2} y_\tau^2 (X + Y - 1).$$

⇒ if the UV completion has such a regime, it generically produces a wrong PMNS matrix

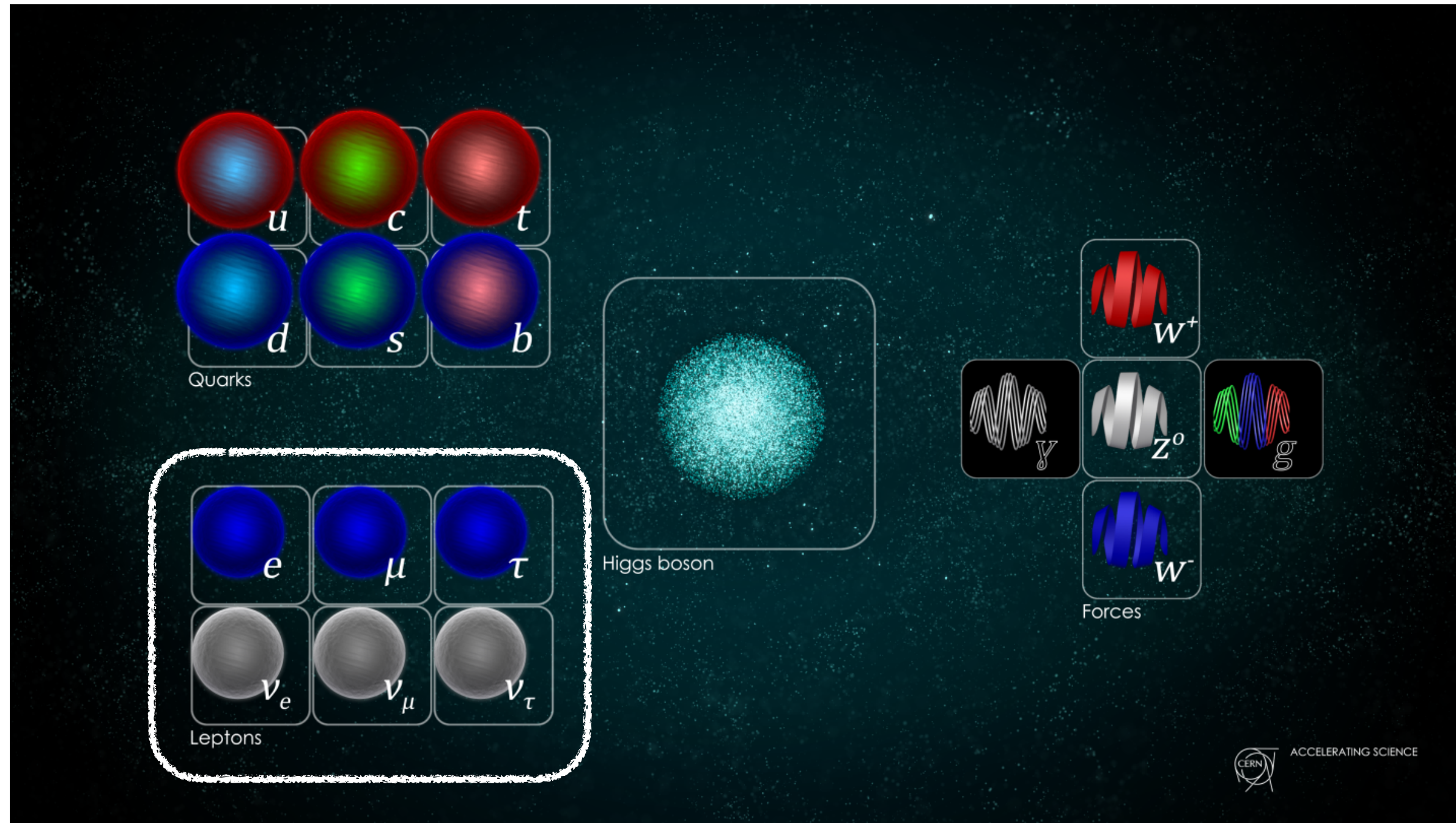
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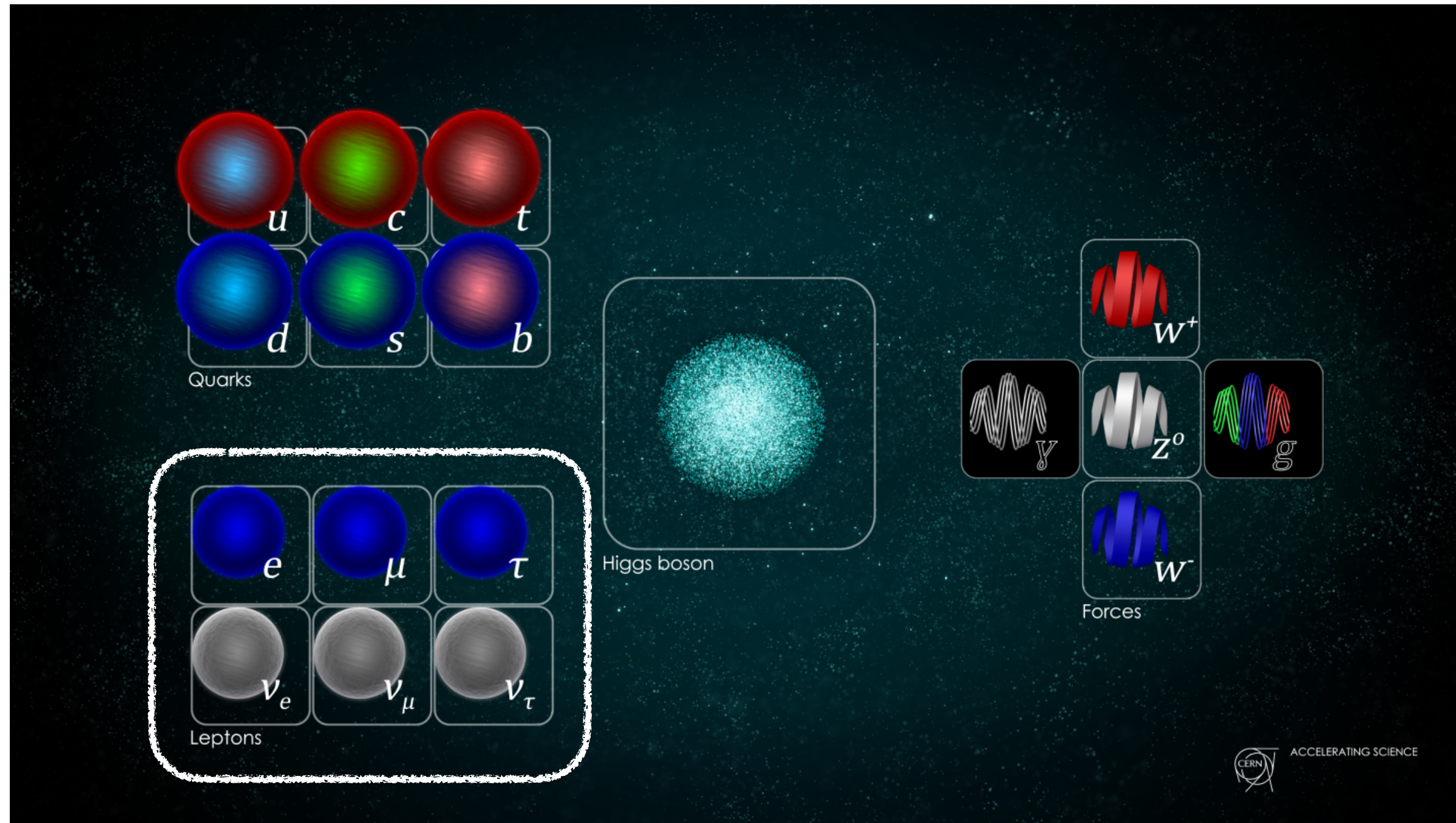
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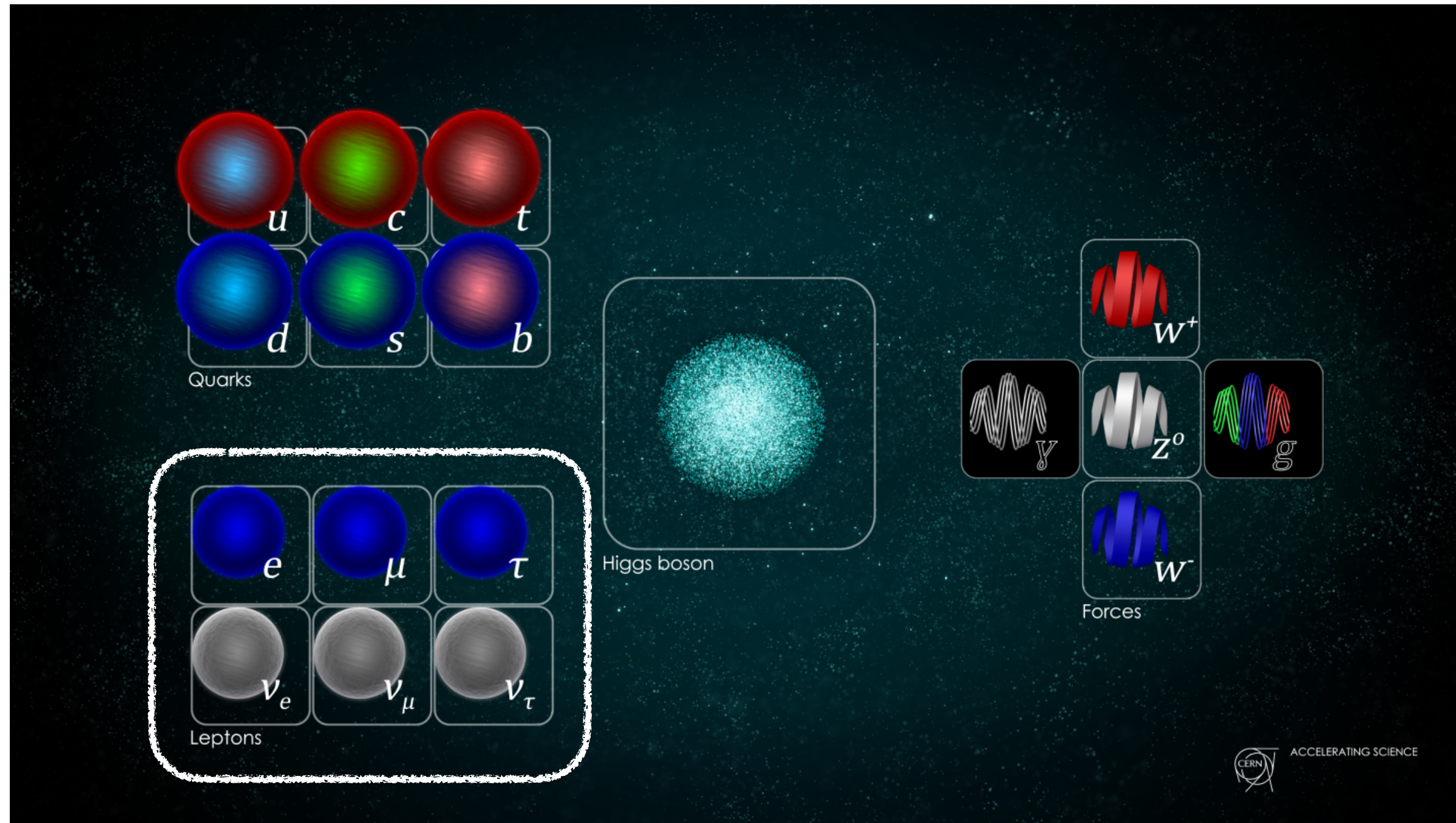
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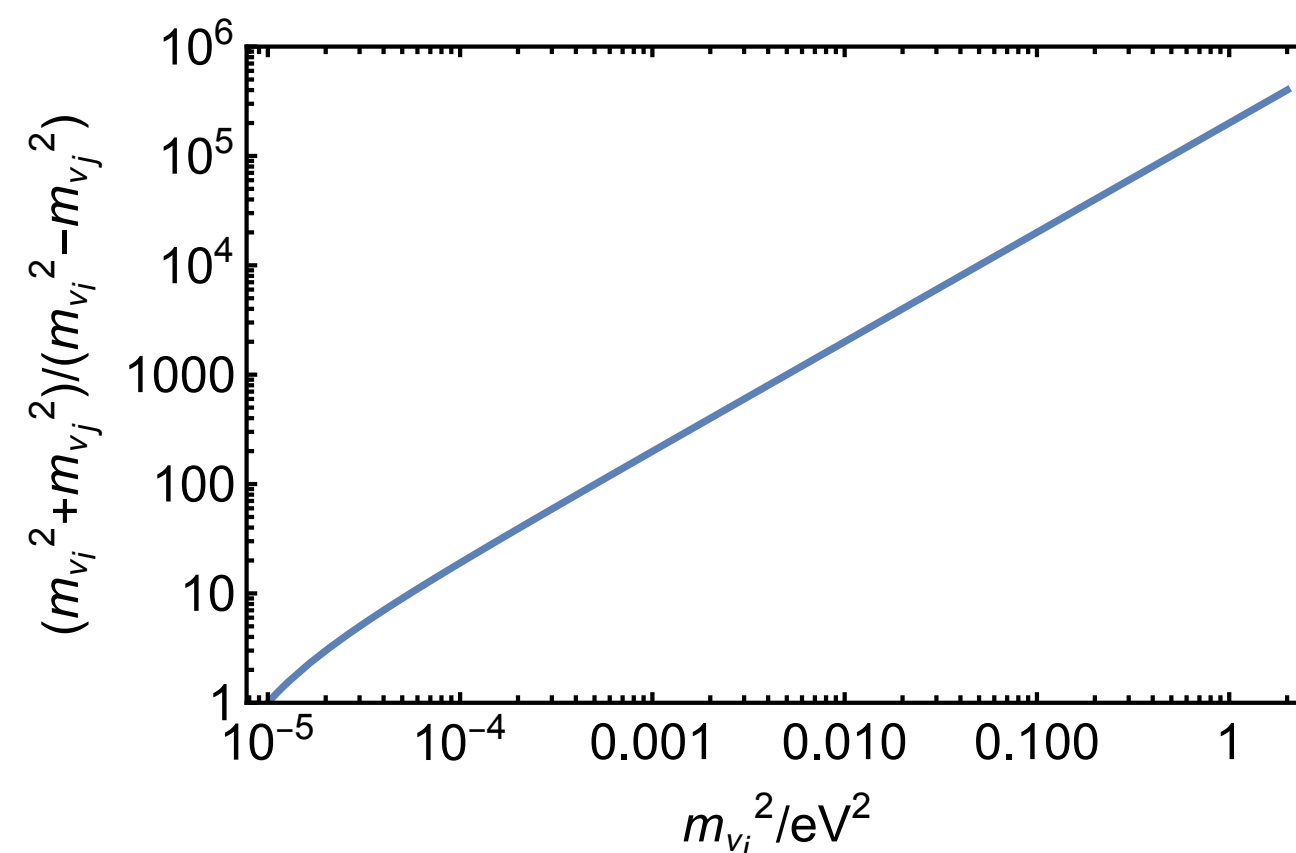


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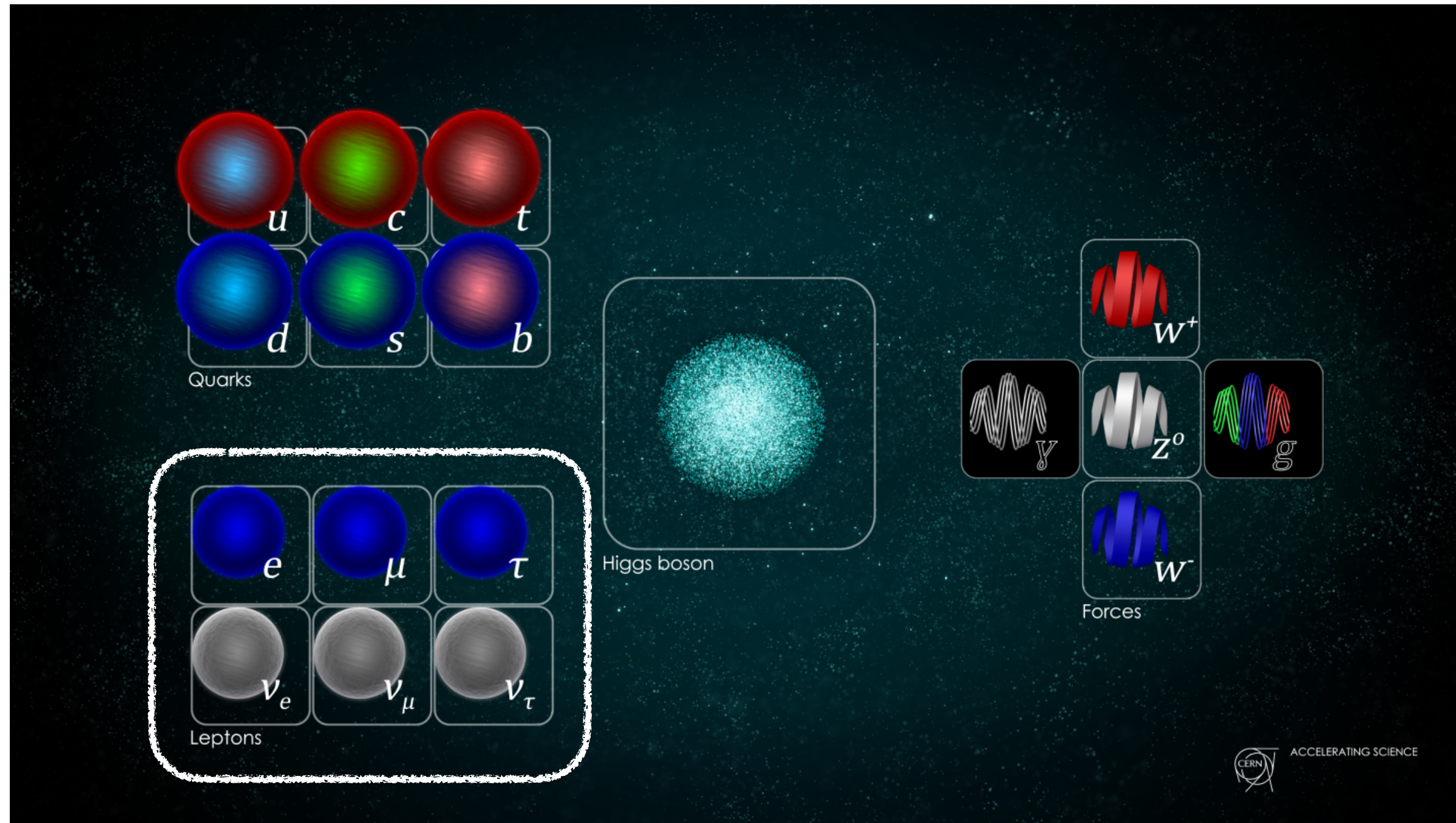
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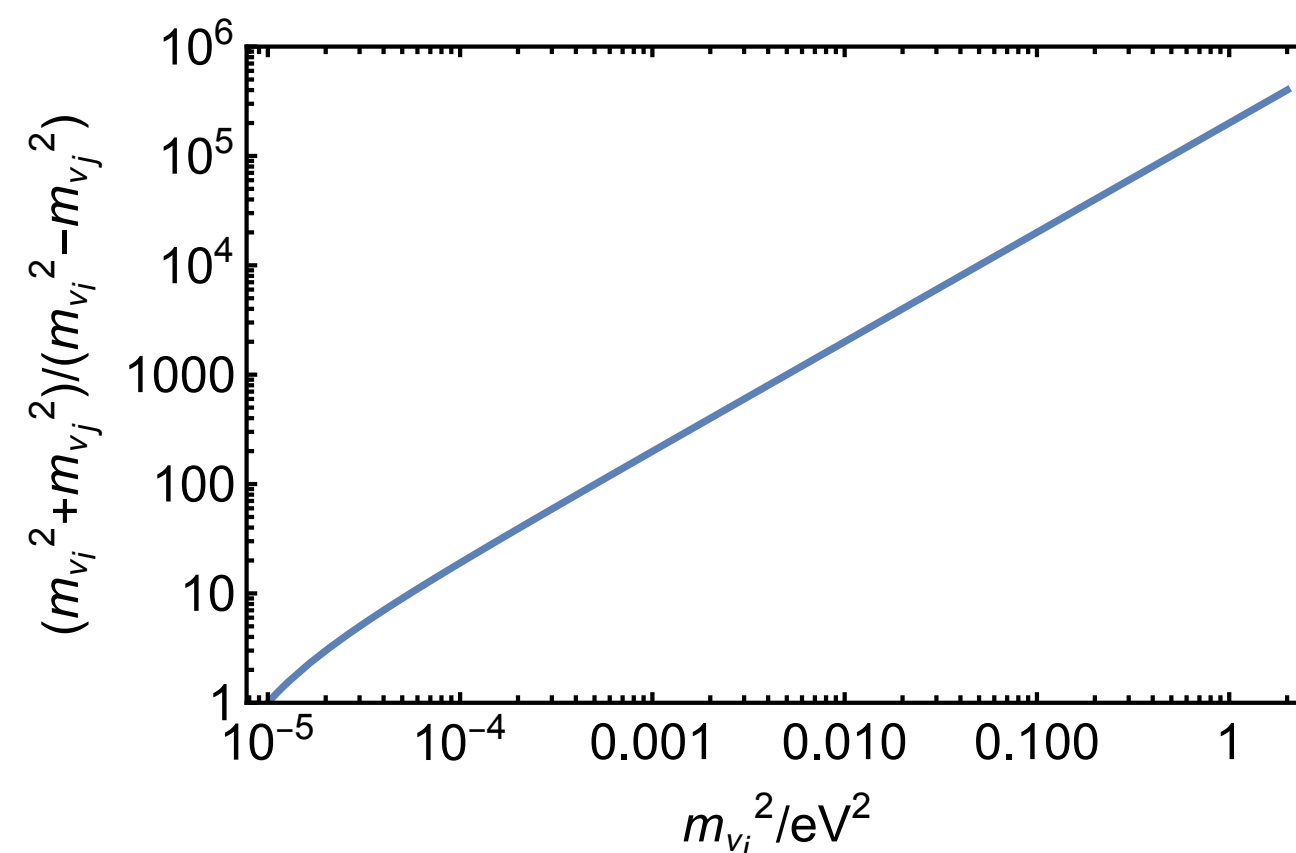


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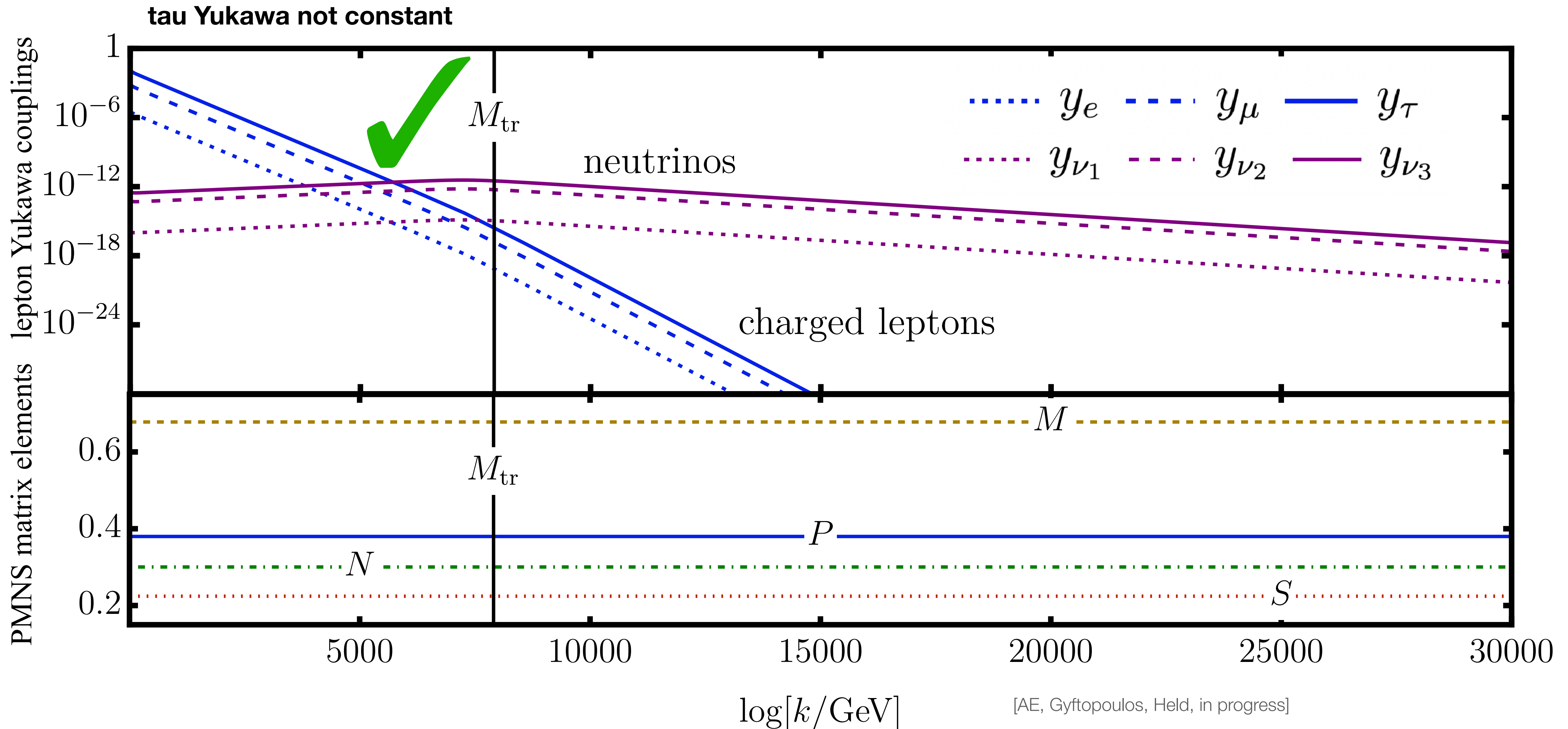


if the UV completion has
neutrino masses $\gtrsim 10$ eV,
it generically produces
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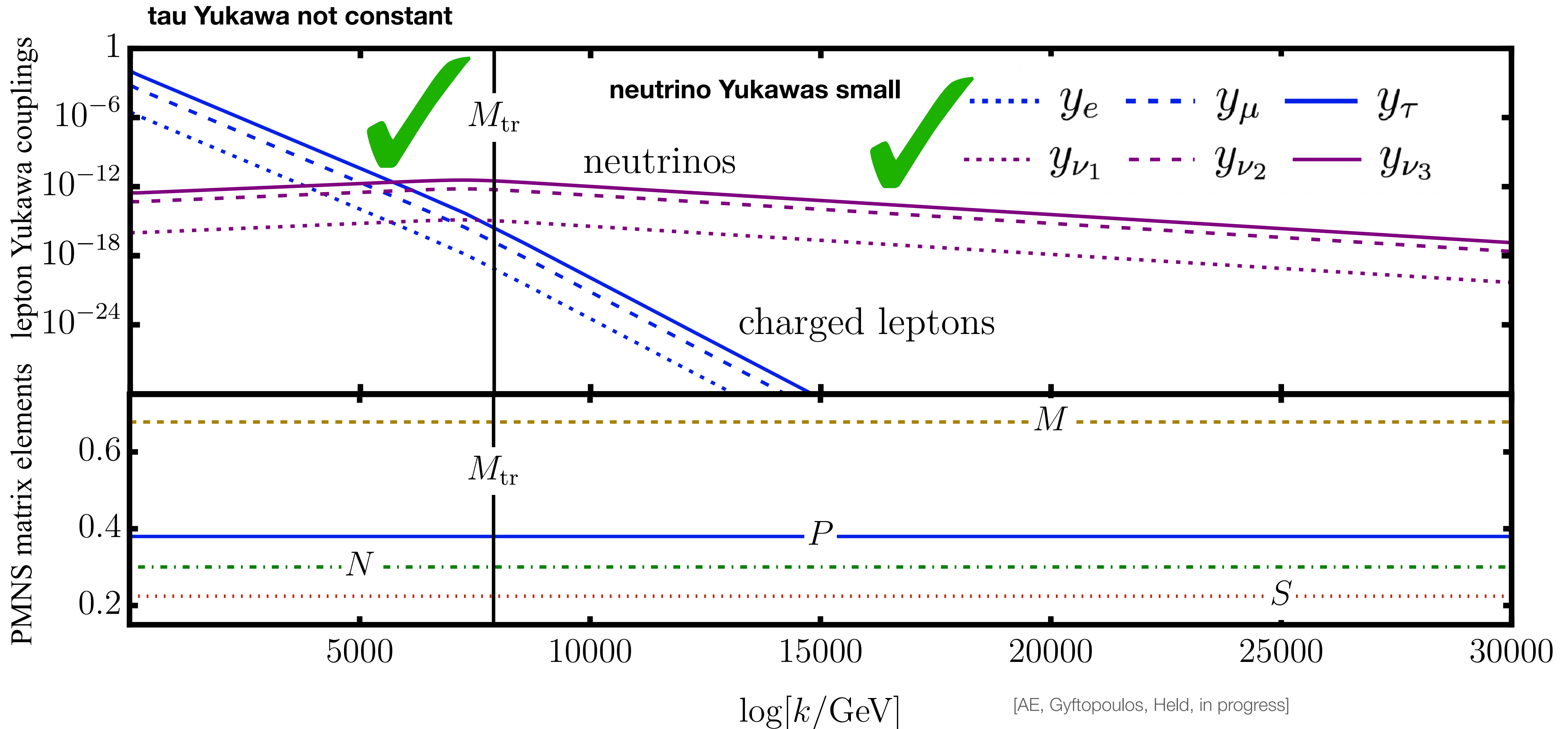
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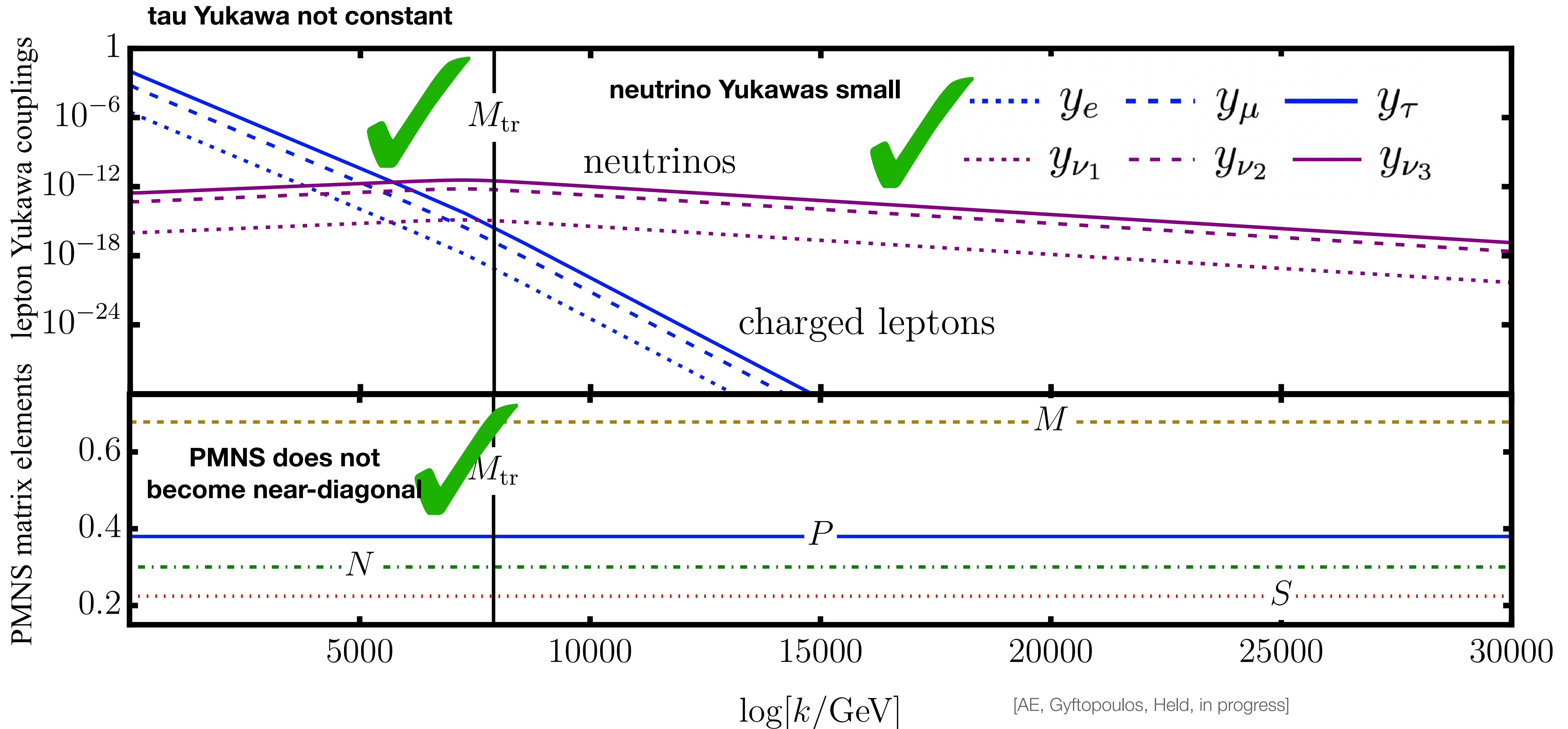
Part 3: mixing matrices



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Mixing matrix summary

Mechanism for mixing matrices to approach near-diagonal configuration, if:

**heaviest fermion dominates the RG flow over long range of scales,
or relative differences between fermion Yukawa couplings become tiny.**

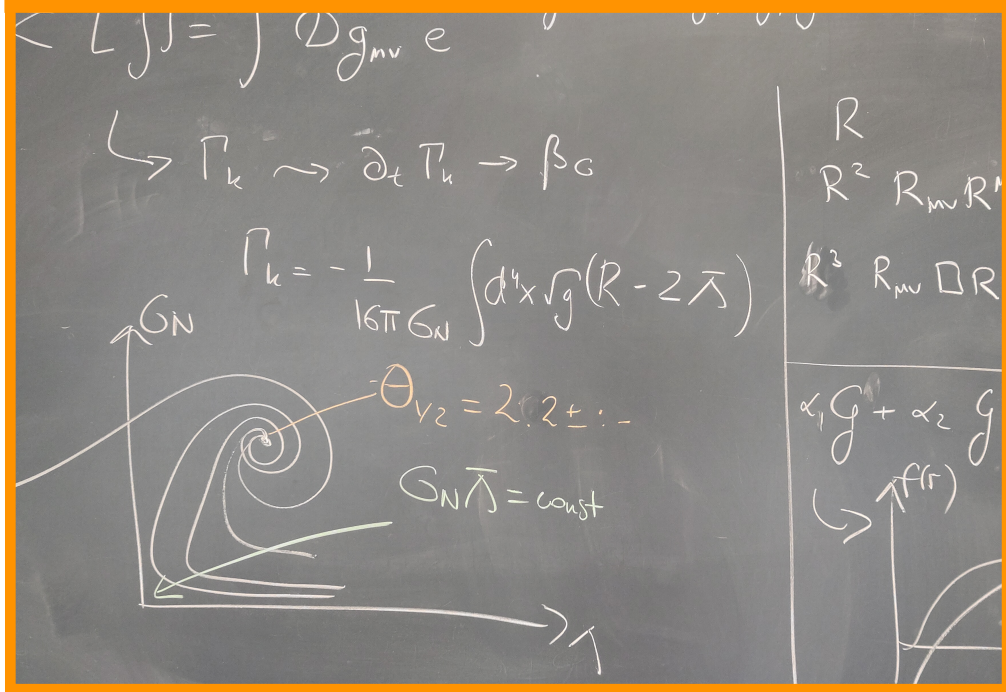
Asymptotically safe Standard Model:

- **realizes this mechanism for CKM, because top Yukawa coupling is constant and large over huge range in scales**
- **avoids this mechanism for PMNS, because tau Yukawa coupling is non-constant and overall neutrino mass scale is kept low, so relative differences do not become tiny**

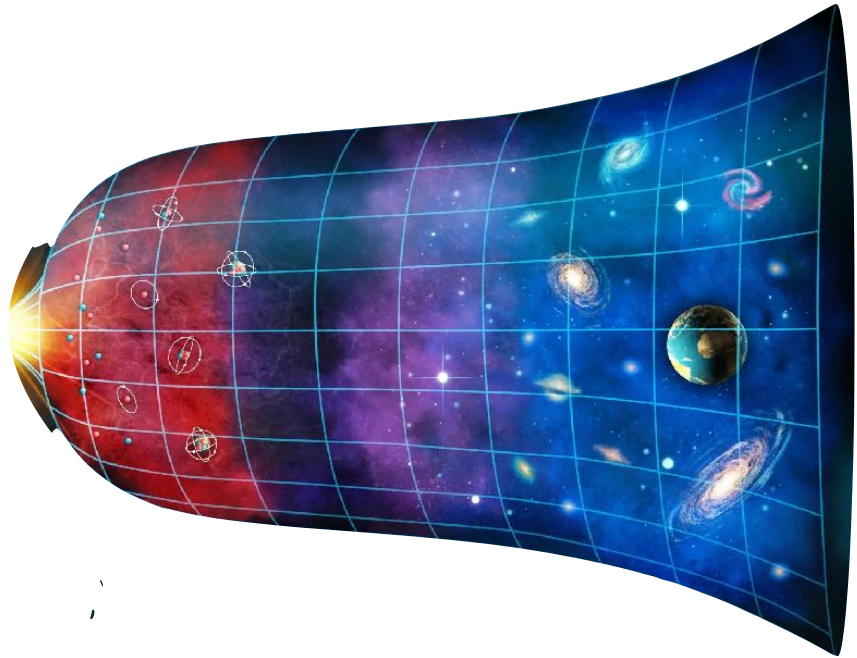
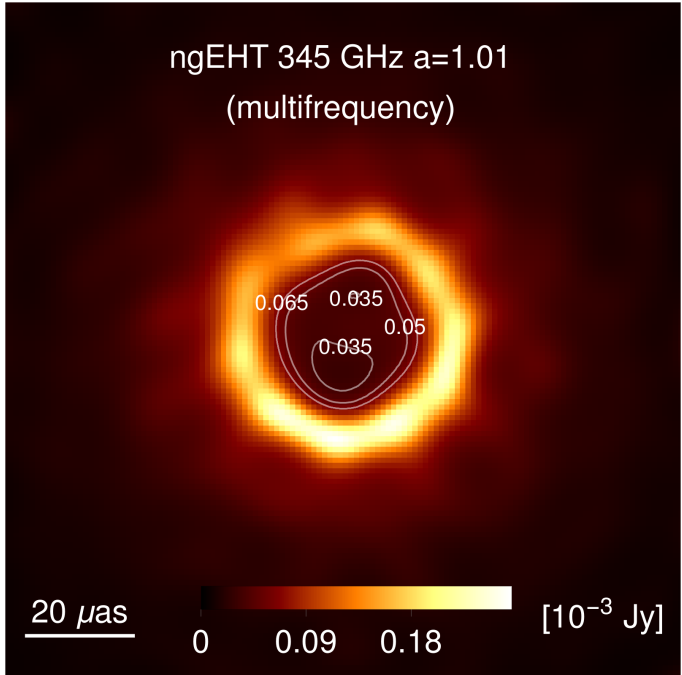
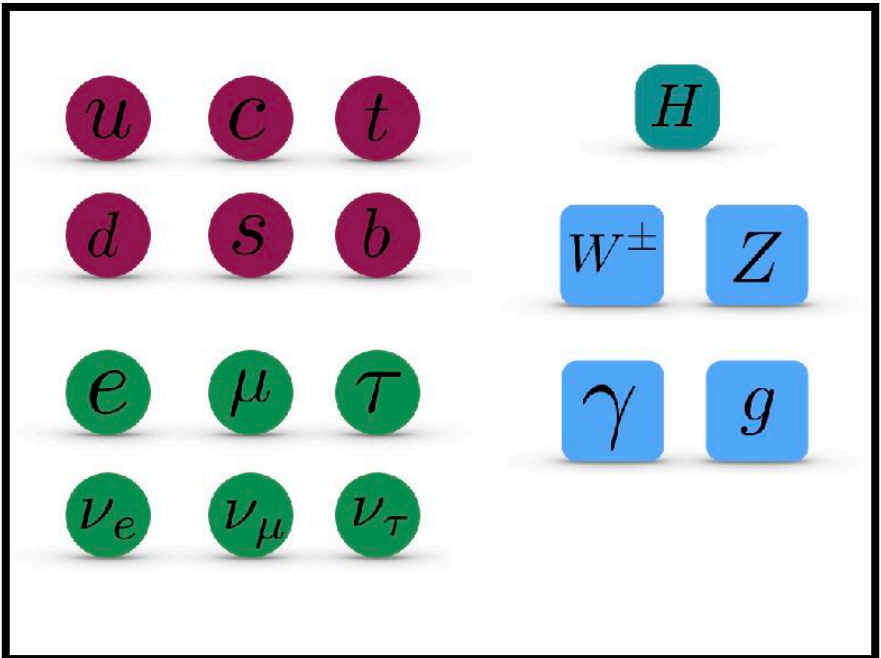
Consequence: RG flows over huge ranges of scales are suggested by observed patterns in mixing matrices

Motivation: How to test proposed theories of quantum gravity?

Theory of quantum gravity



Key challenge: gap in scales



Planckian scales
 10^{-35} m

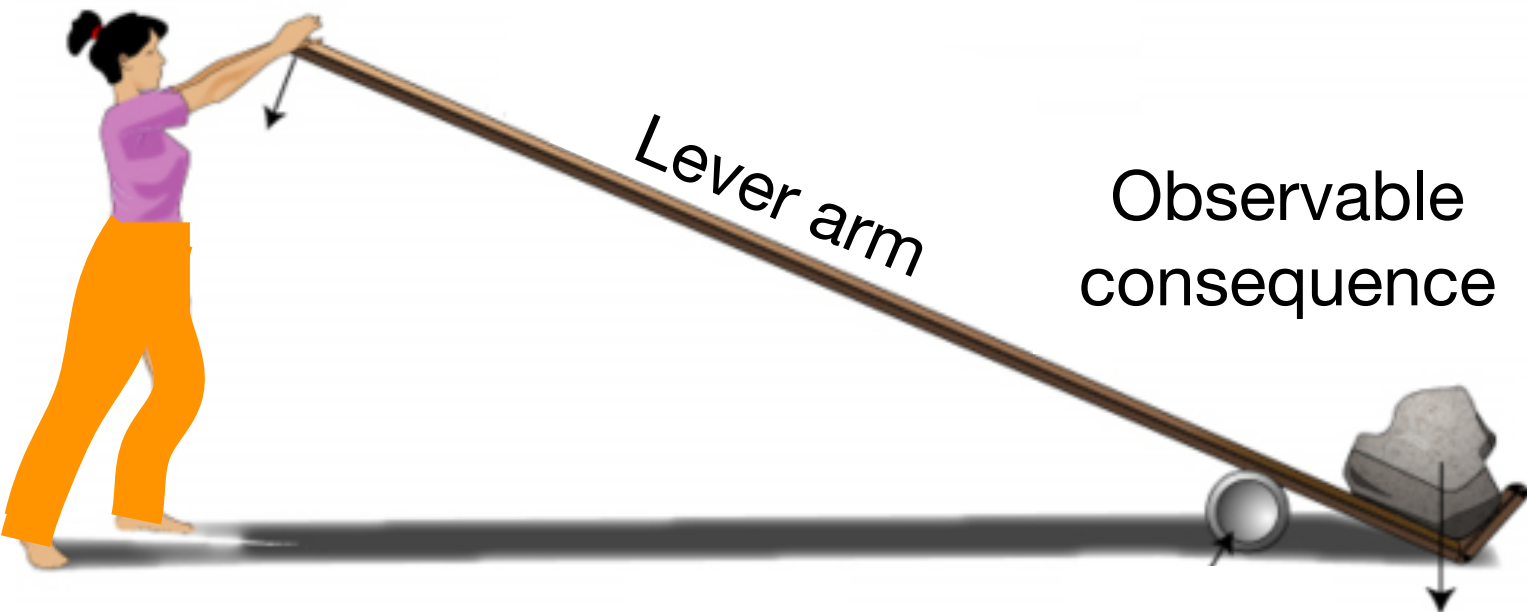
Particle physics scales
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Black-hole scales
 10^{11} m

Cosmological scales
 $> 10^{20}$ m

distance scale

Quantum-gravity effect

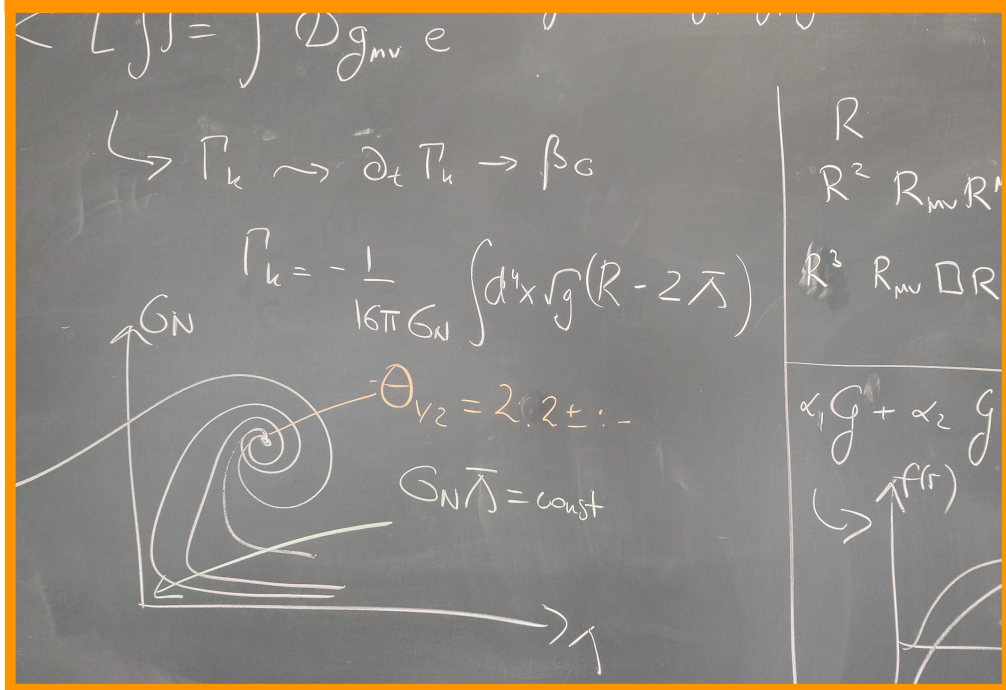


Examples:

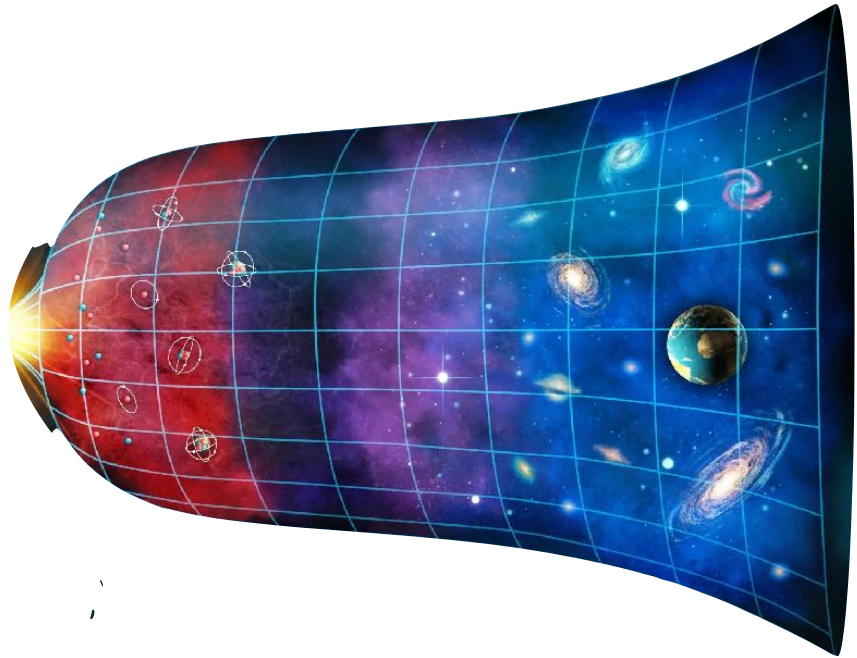
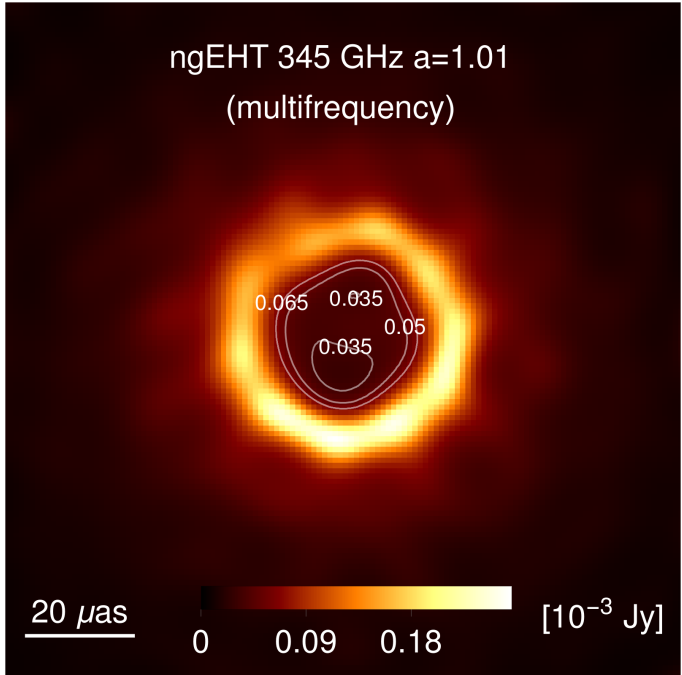
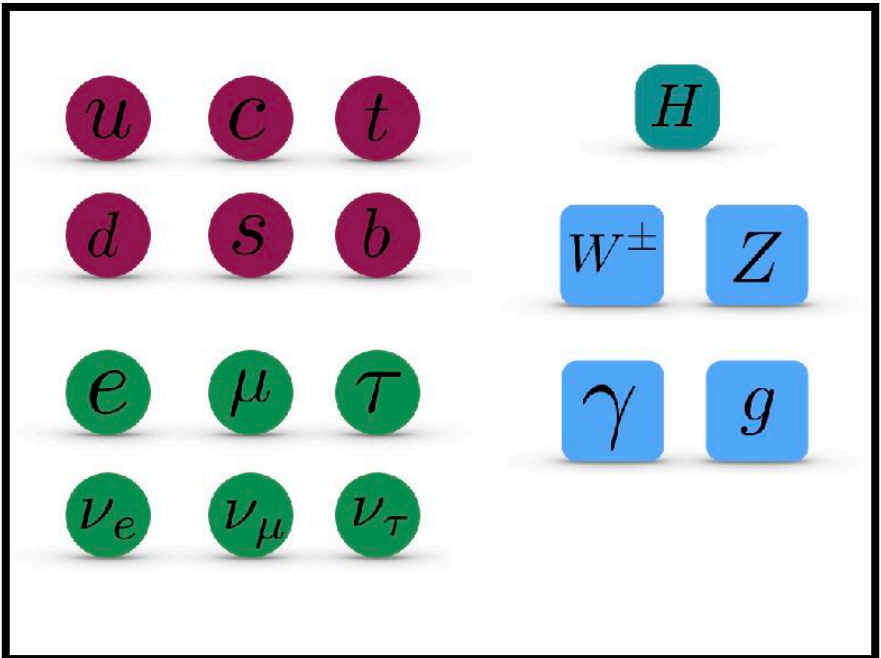
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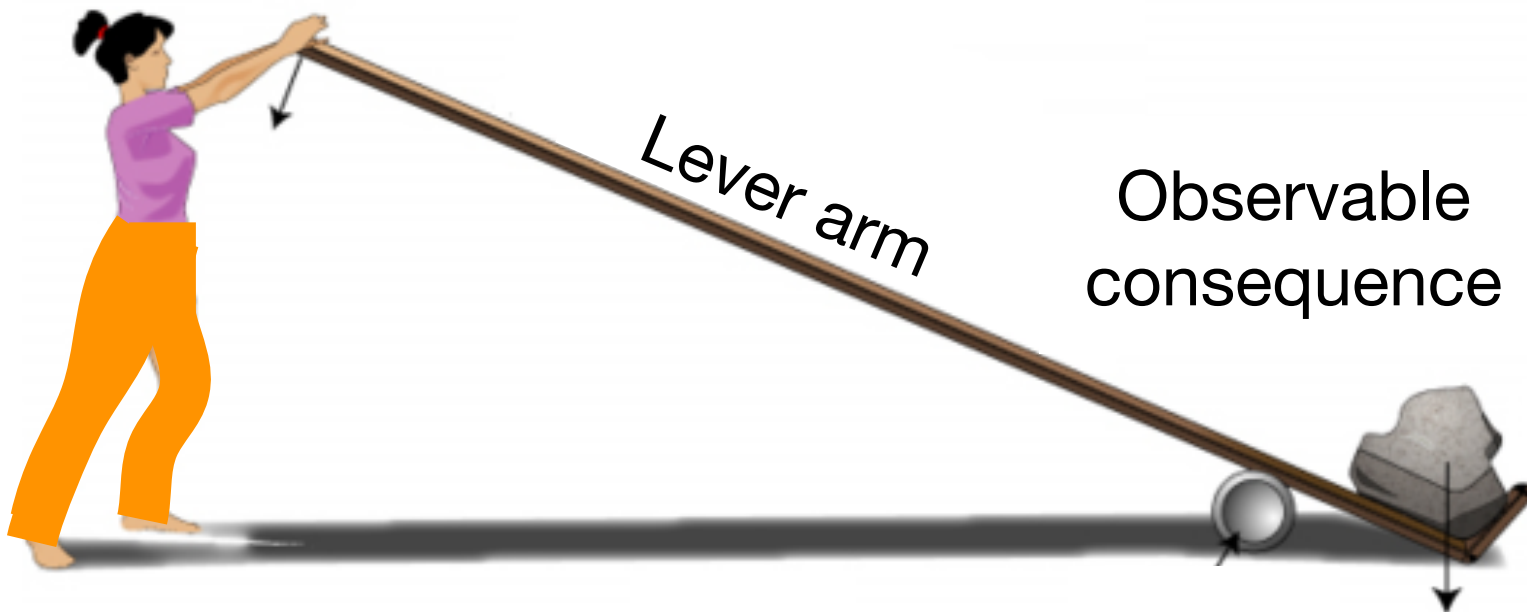
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What about the Renormalization Group perspective on Lorentz-invariance-violations and CPT violation?

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CPT & Lorentz violation and Renormalization Group flows

[AE, Schiffer '25]

$$S_{\text{scalar}}^{\text{CPT-odd}} = i\, g_\phi \int d^4x \sqrt{g} \, n_\mu \left(\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^* \right) .$$

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Experimental status: numerous constraints,

some examples: (see Kostelecky)

[Submitted on 1 Jan 2008 ([v1](#)), last revised 13 Jan 2025 (this version, v18)]

Data Tables for Lorentz and CPT Violation

[Alan Kostelecky](#), [Neil Russell](#)

This work tabulates measured and derived values of coefficients for Lorentz and CPT violation in the Standard-Model Extension. Summary tables are extracted listing maximal attained sensitivities in the matter, photon, neutrino, and gravity sectors. Tables presenting definitions and properties are also compiled.

Comments: 160 pages, 2025 edition

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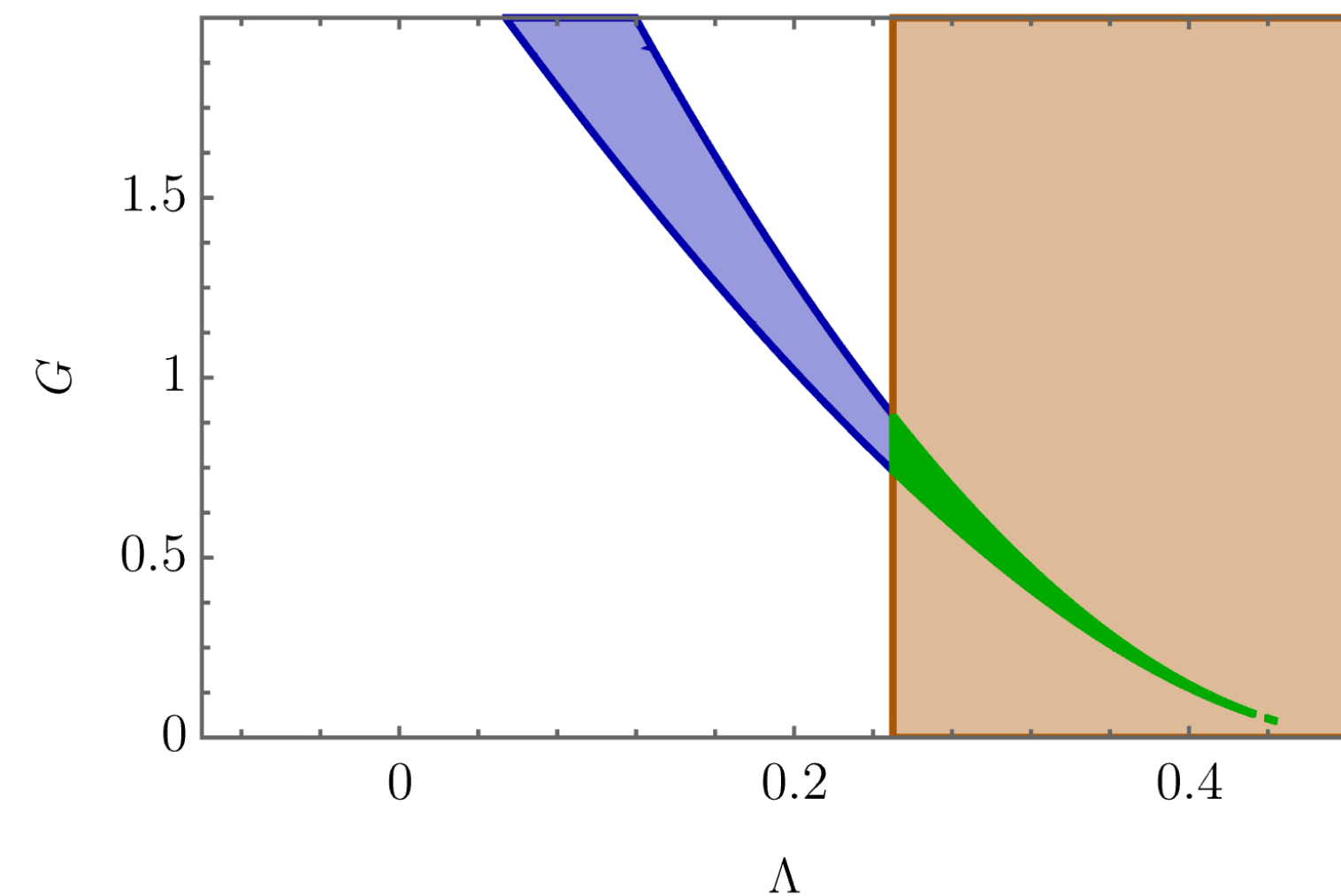
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$$f_\gamma > 0$$

$$f_\gamma > 0 \text{ \& } f_\phi > 0$$

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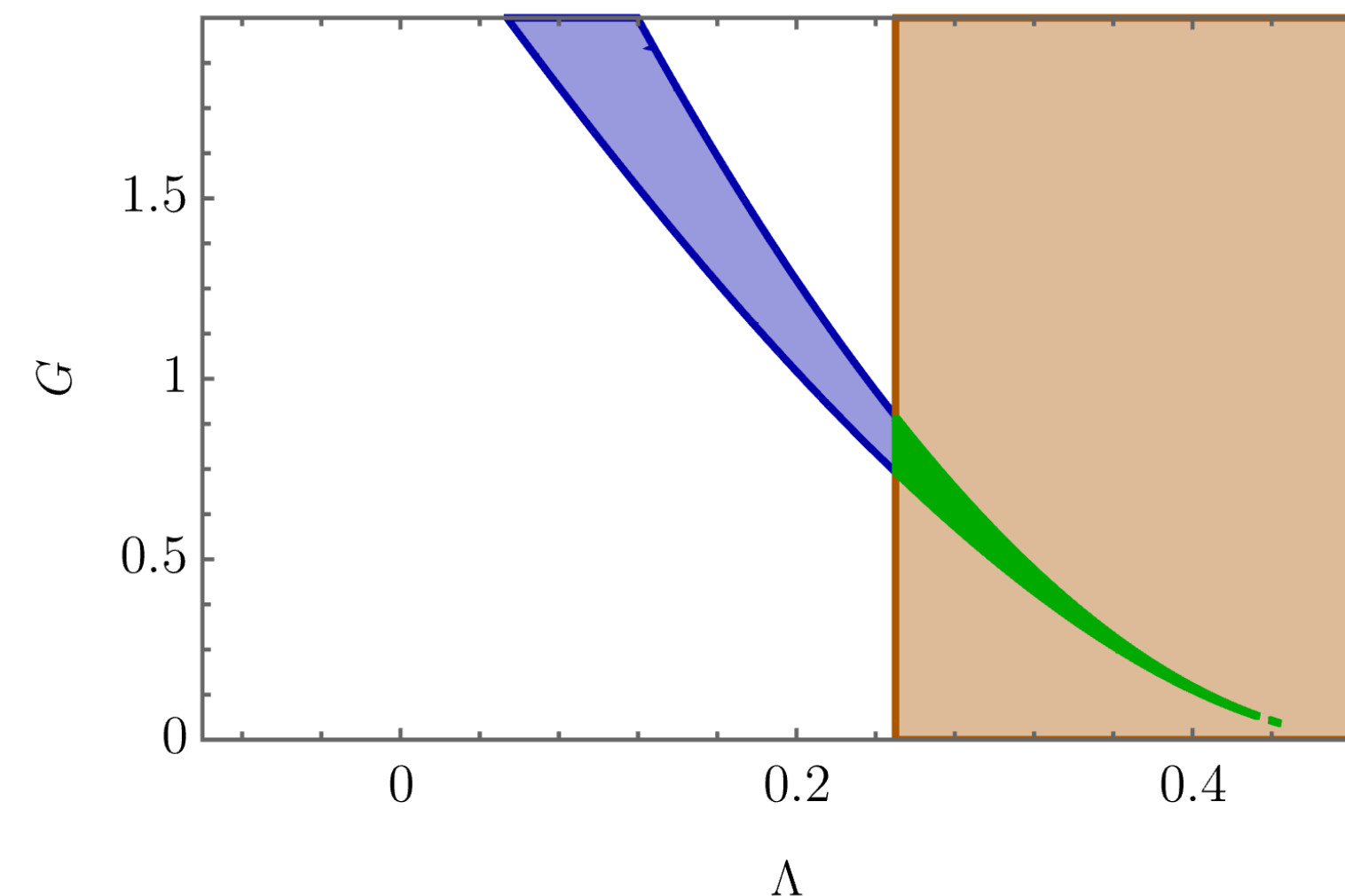
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$\Rightarrow$ quantum-gravity theories need a mechanism to ensure $g_i(\Lambda_{\text{UV}}) = 0$

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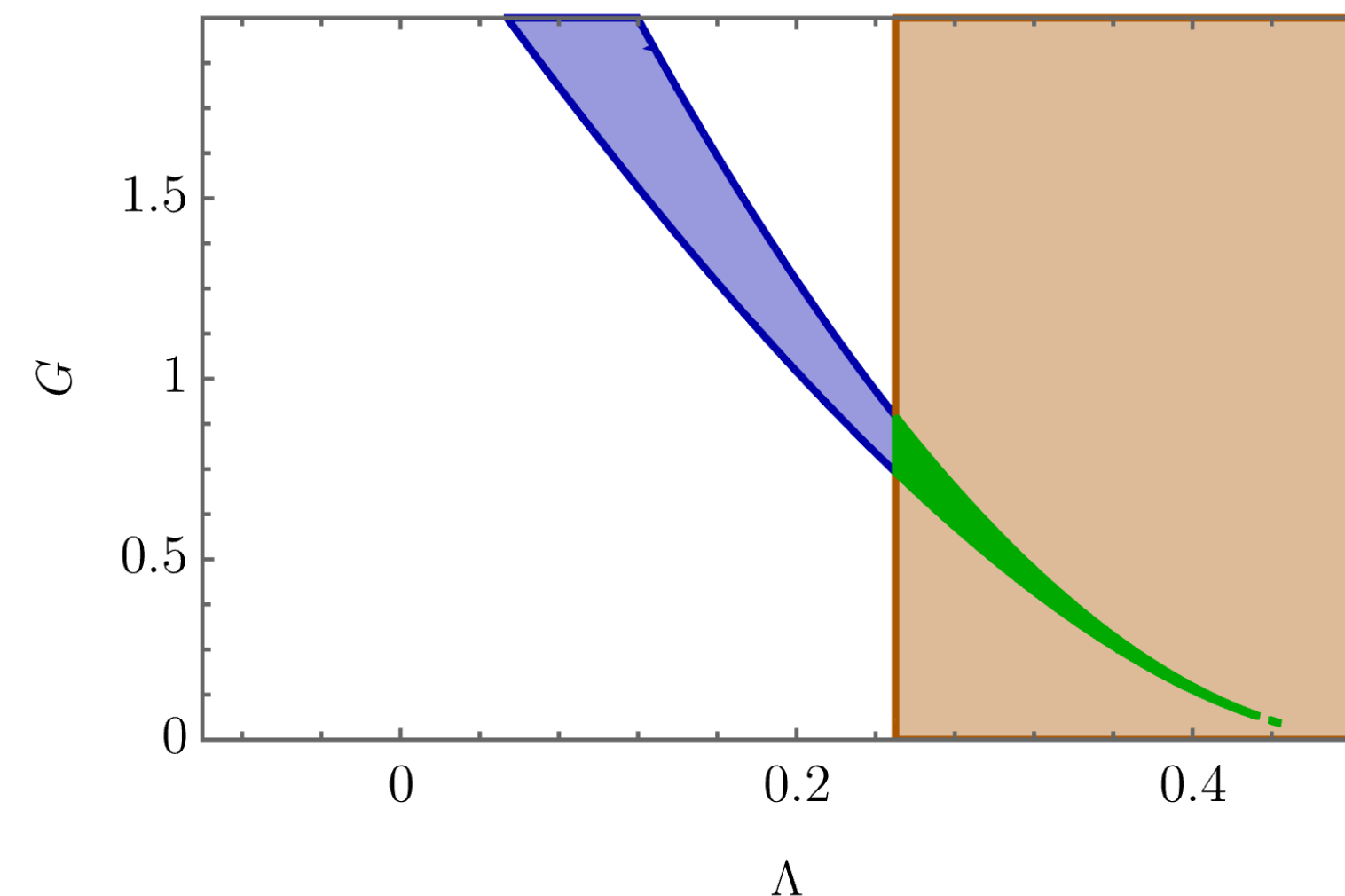
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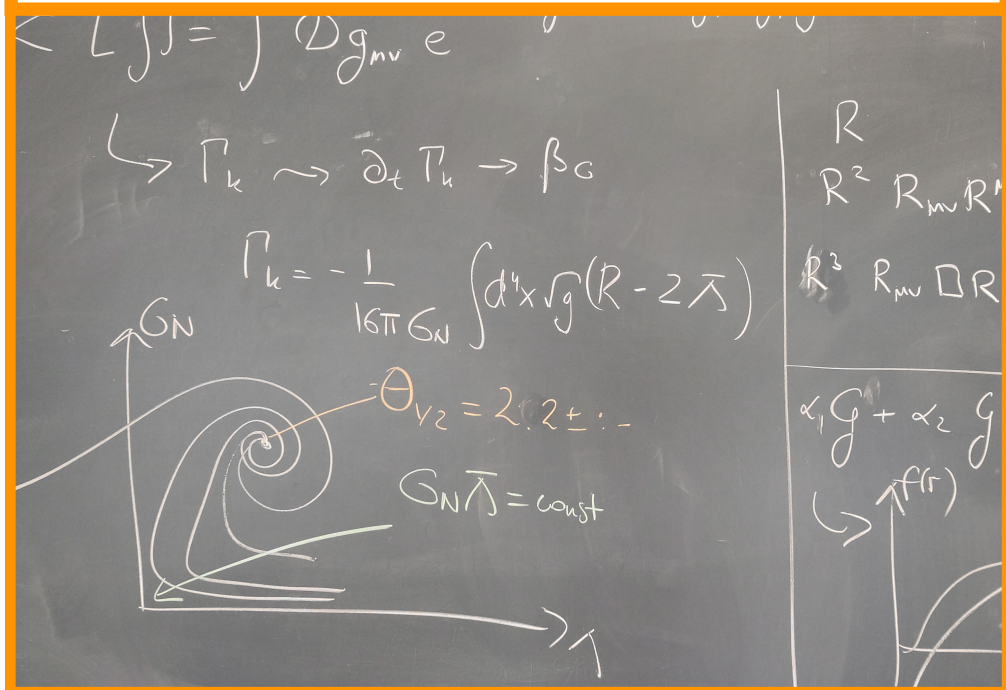
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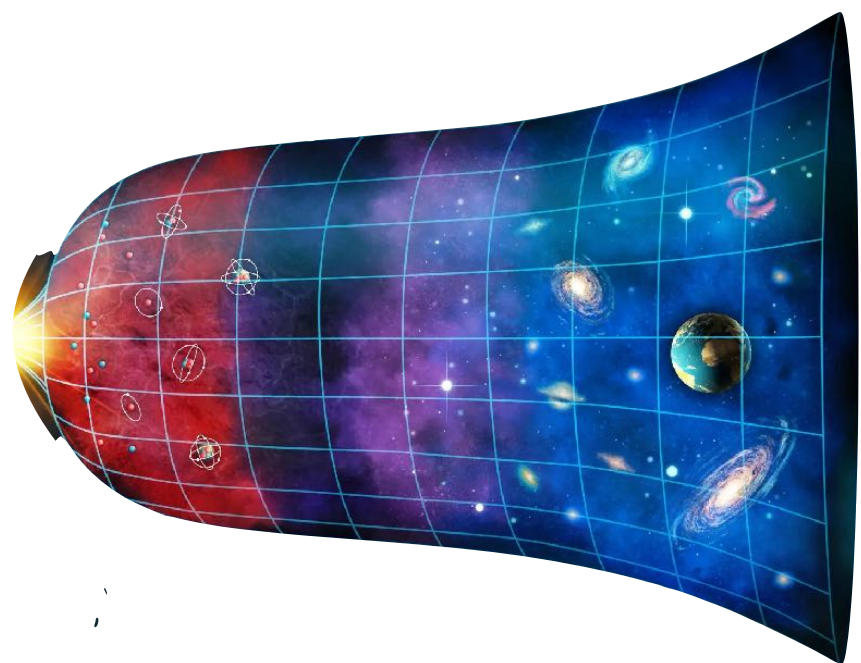
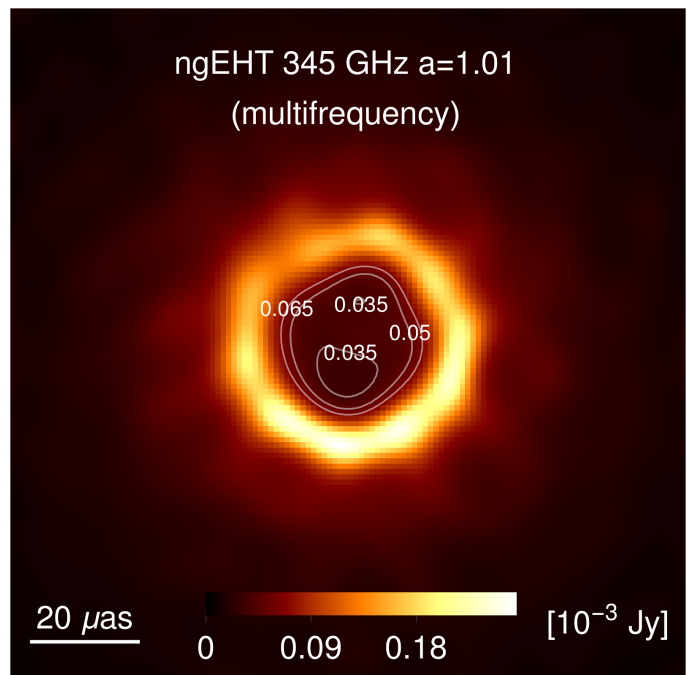
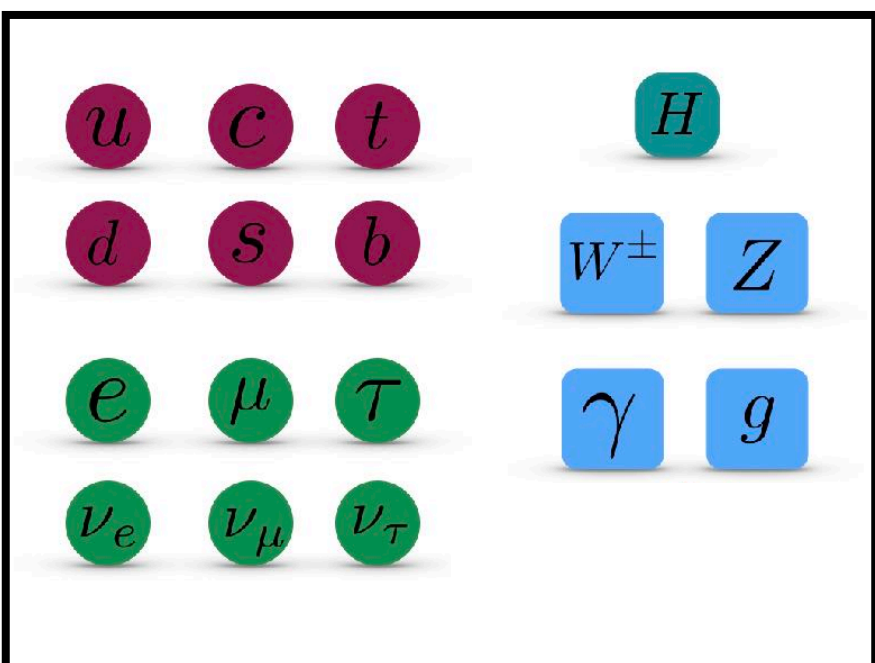
**Two examples: string theory (all symmetries are gauged),
asymptotic safety (global symmetries can be imposed)**

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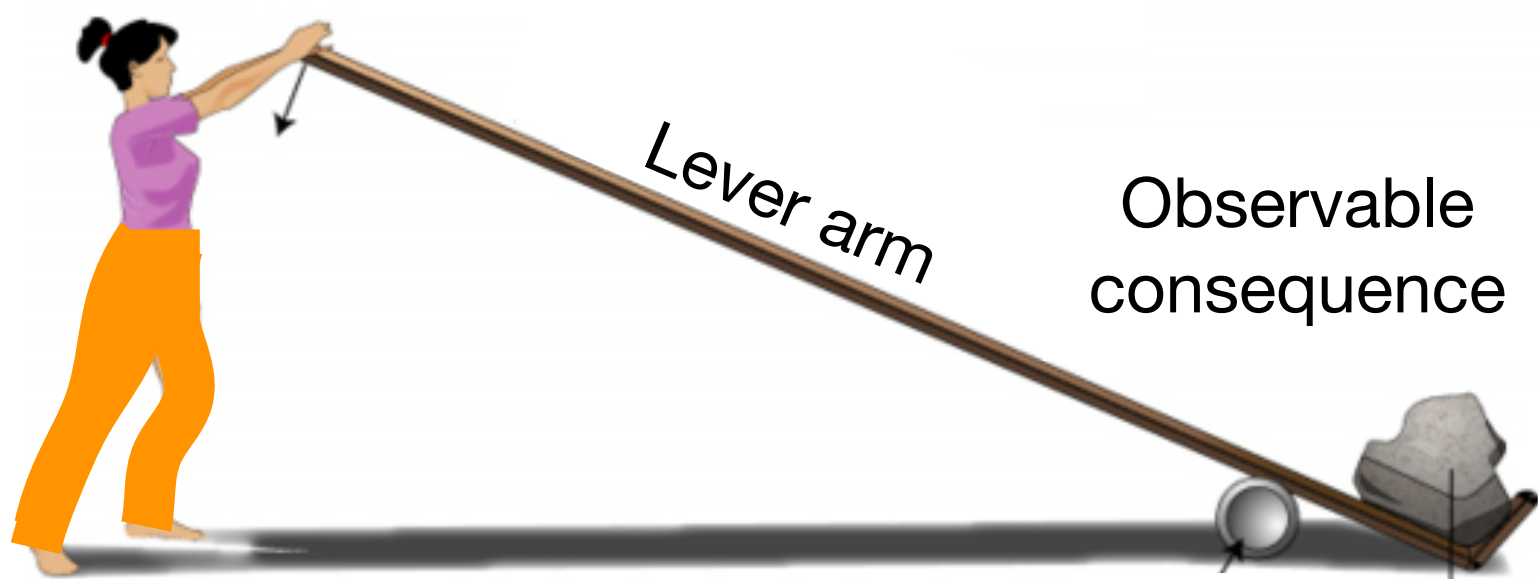
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Renormalization Group flow of couplings can be powerful lever arm for Bridging high & low energies



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[Amelino-Camelia, Ellis, Mavromatos Nanopoulos, Sakar '97]
- High sensitivity in probes of CPT symmetry breaking in the Standard Model
[Colladay, Kostelecky '96]
- Renormalization Group flow of couplings

Thanks to current and former group members!

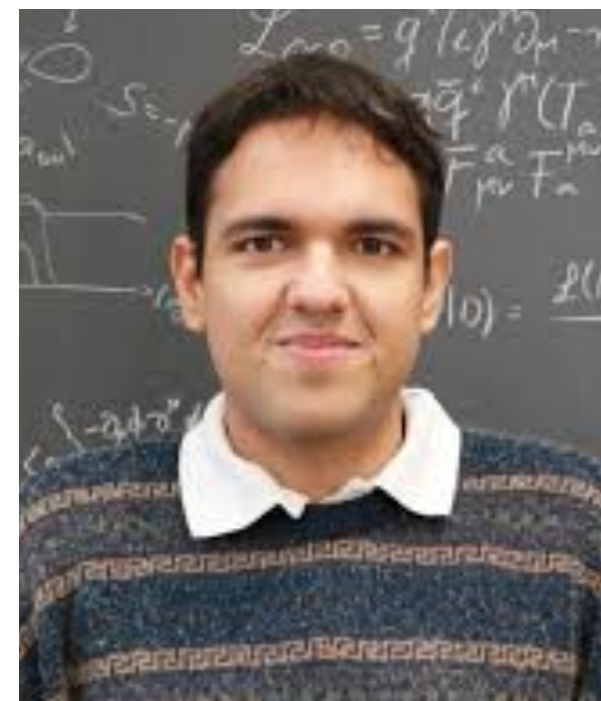
Former PhD students:



Aaron Held,
now ENS Paris



Marc Schiffer,
now Radboud U.



Rafael R. Lino dos Santos,
now Warsaw U.



Johannes Lumma,
now Oxford U.



Fleur Versteegen,
now ASML



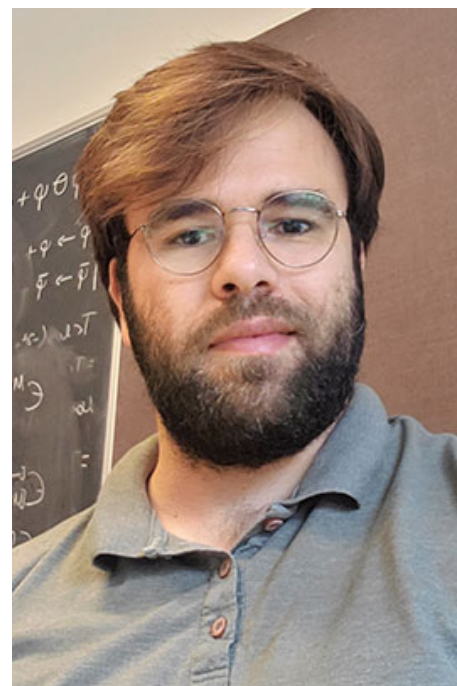
Martin Pauly,
now exnaton



Héloïse Delaporte,
now U. of Faroe Islands



Antonio Pereira,
assist. prof. at
Fluminense Federal U.,
Brazil



Gustavo P. de Brito,
assist. prof. at
São Paulo State U.



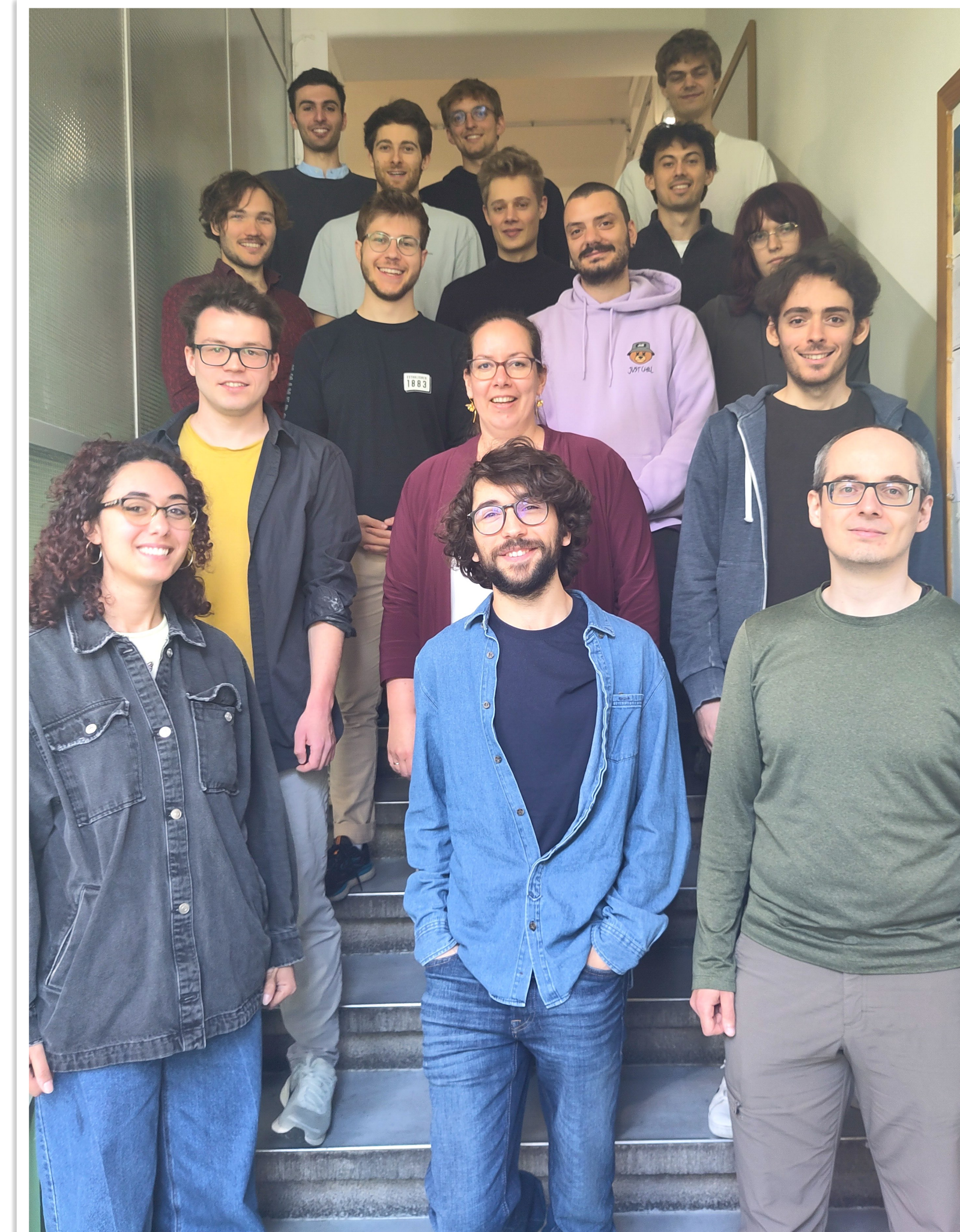
Alessia Platania,
assist. prof.
at Niels-Bohr-
Institute,
soon prof. at U. Graz



Raúl Carballo-Rubio
associate prof.
at U of Granada, Spain



Shouryya Ray
assist. prof.
at U. of Faroe Islands



Near-perturbative nature of asymptotic safety

How non-perturbative is the fixed point?

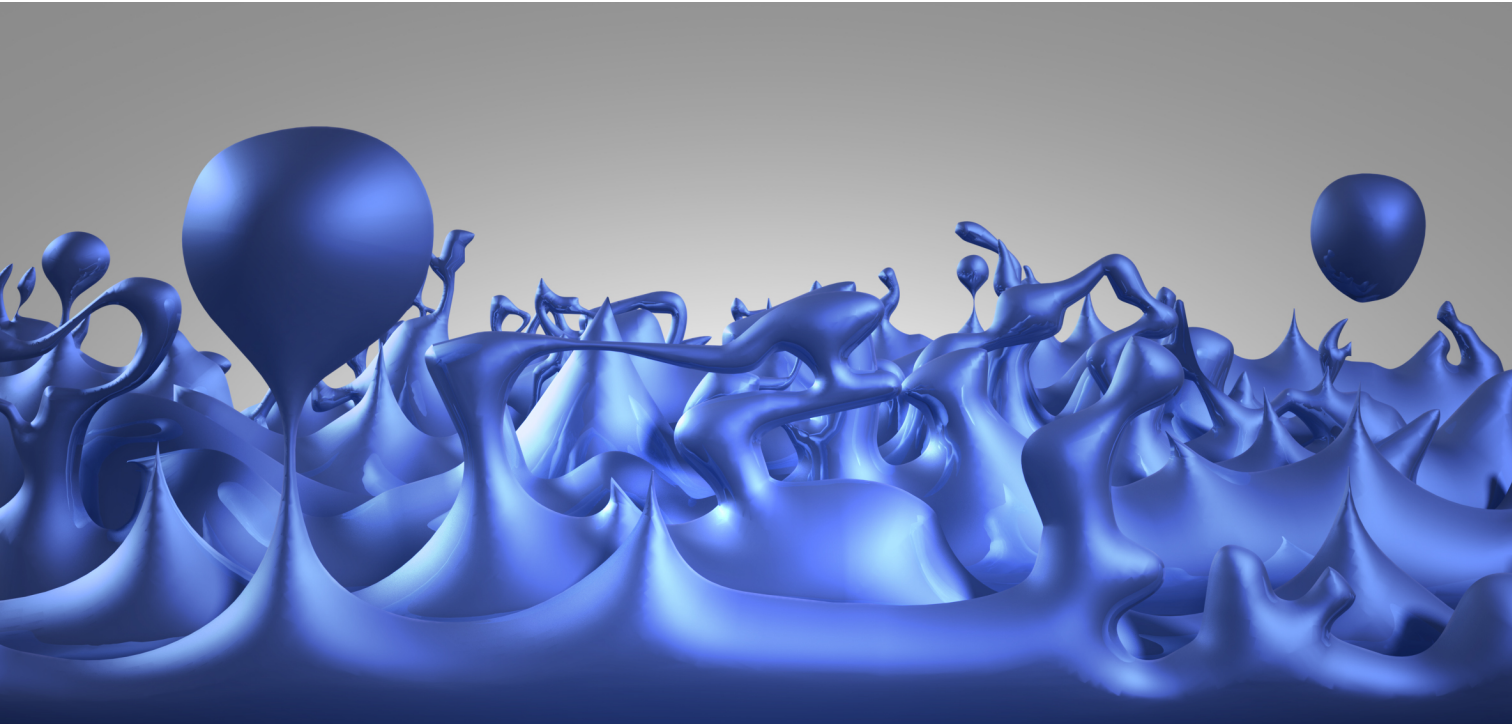
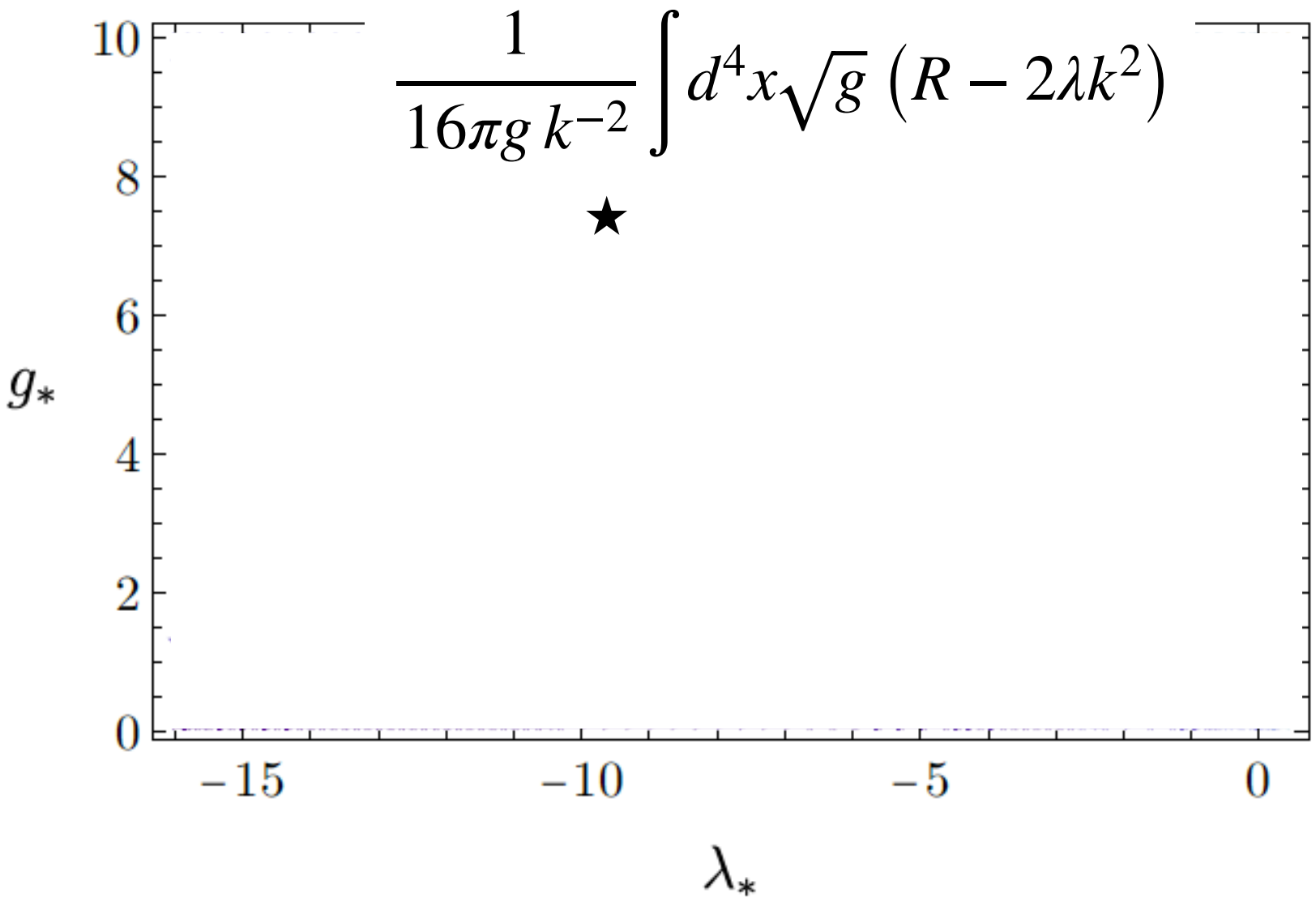


Image Credit: NASA/CXC/M. Weiss

or



?



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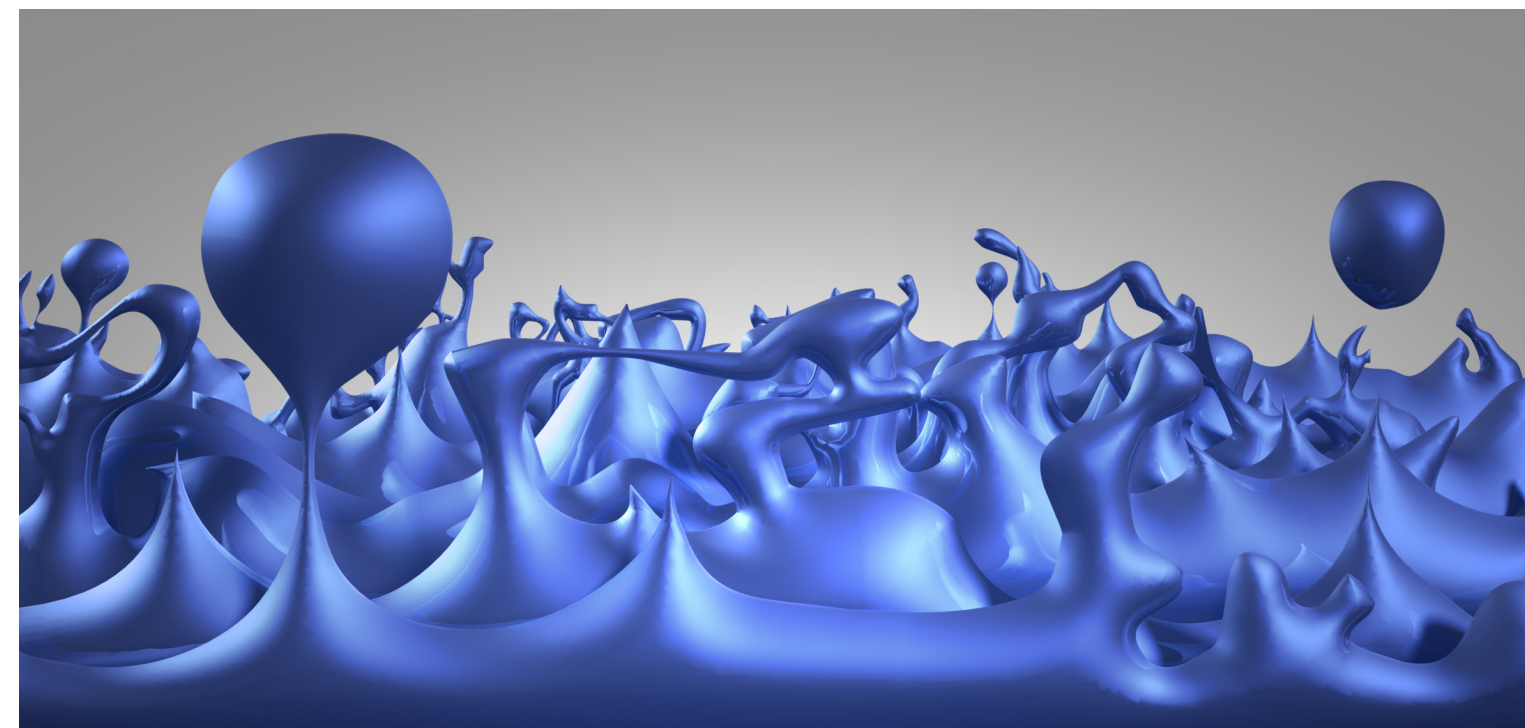
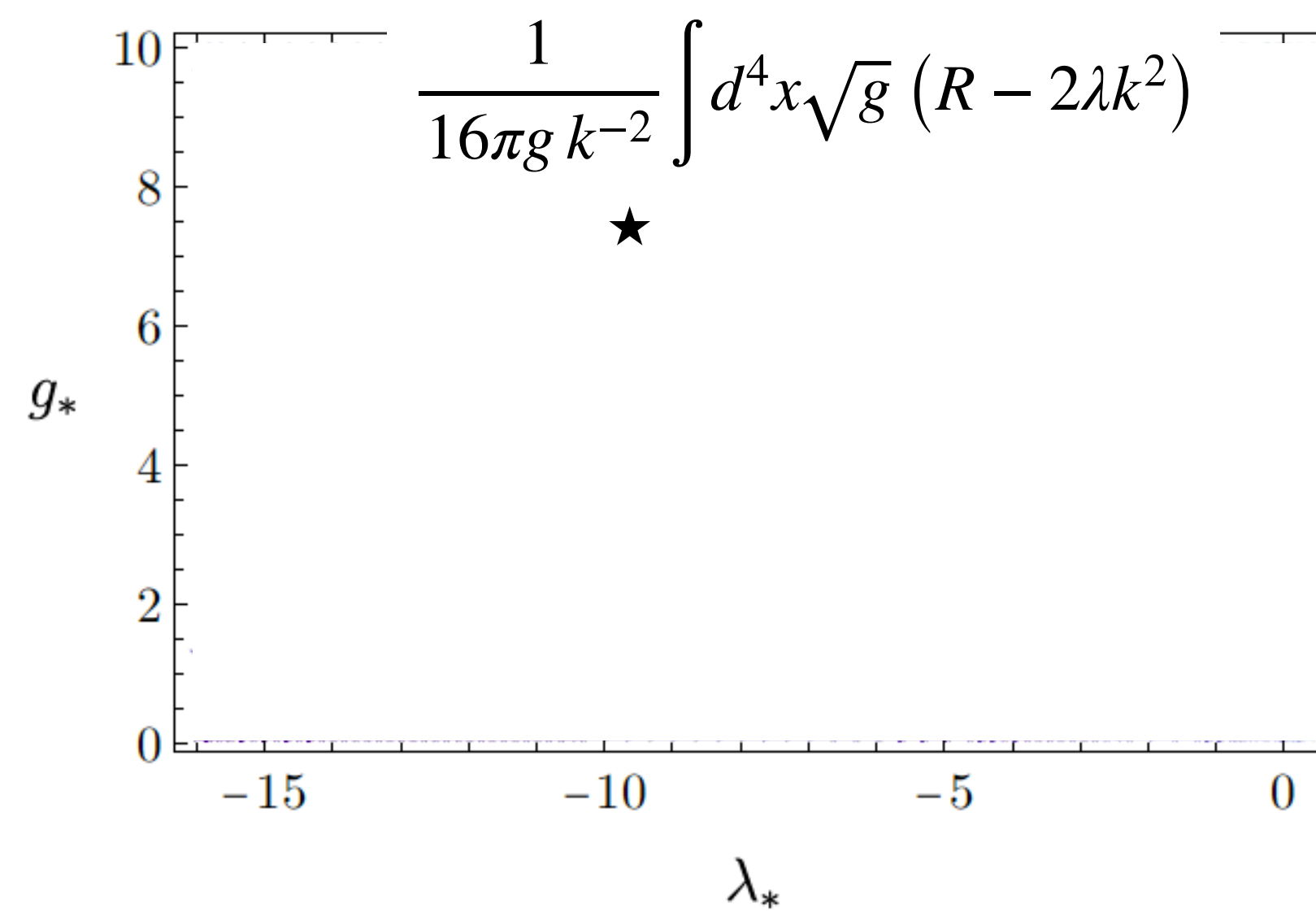


Image Credit: NASA/CXC/M. Weiss

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?



metric propagator:

$$\frac{g_*}{1 - 2\lambda_* + \dots}$$

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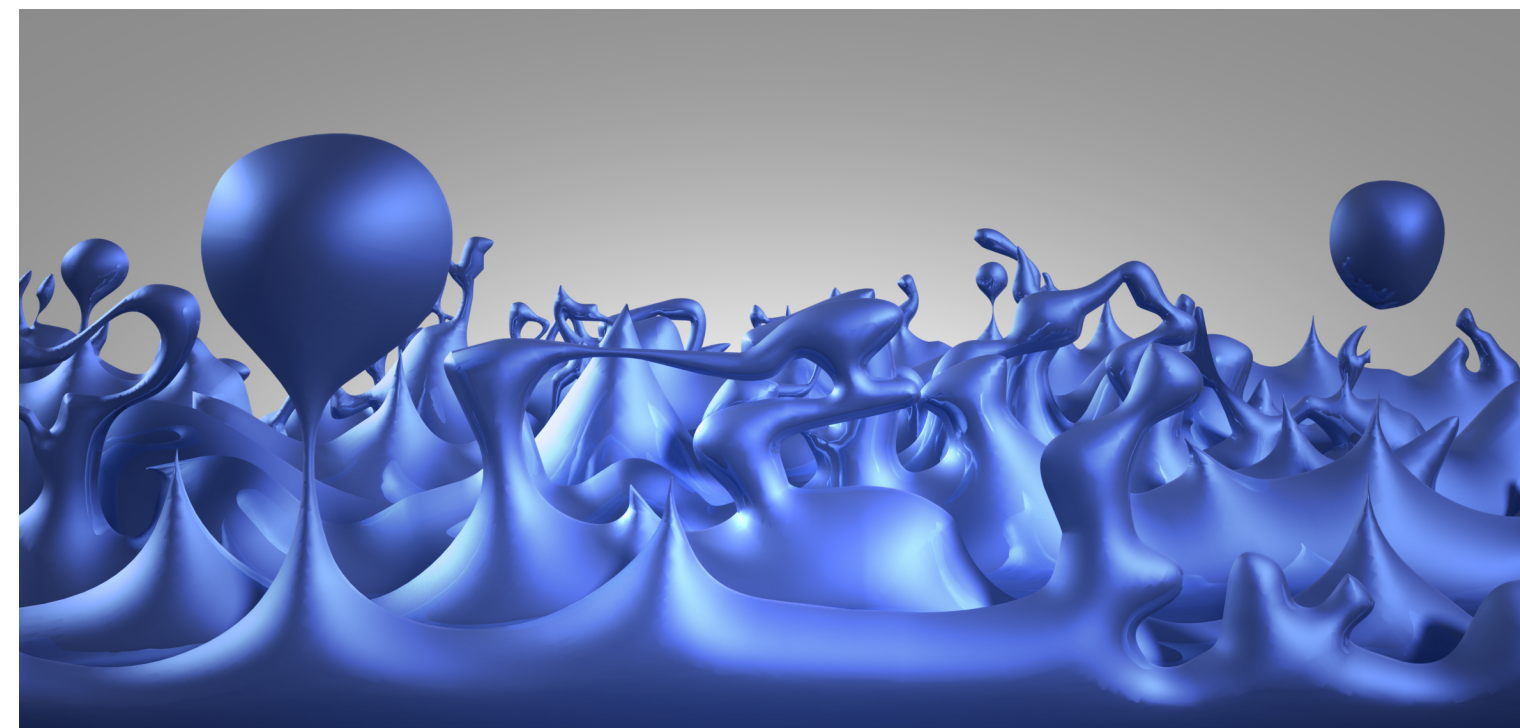


Image Credit: NASA/CXC/M. Weiss

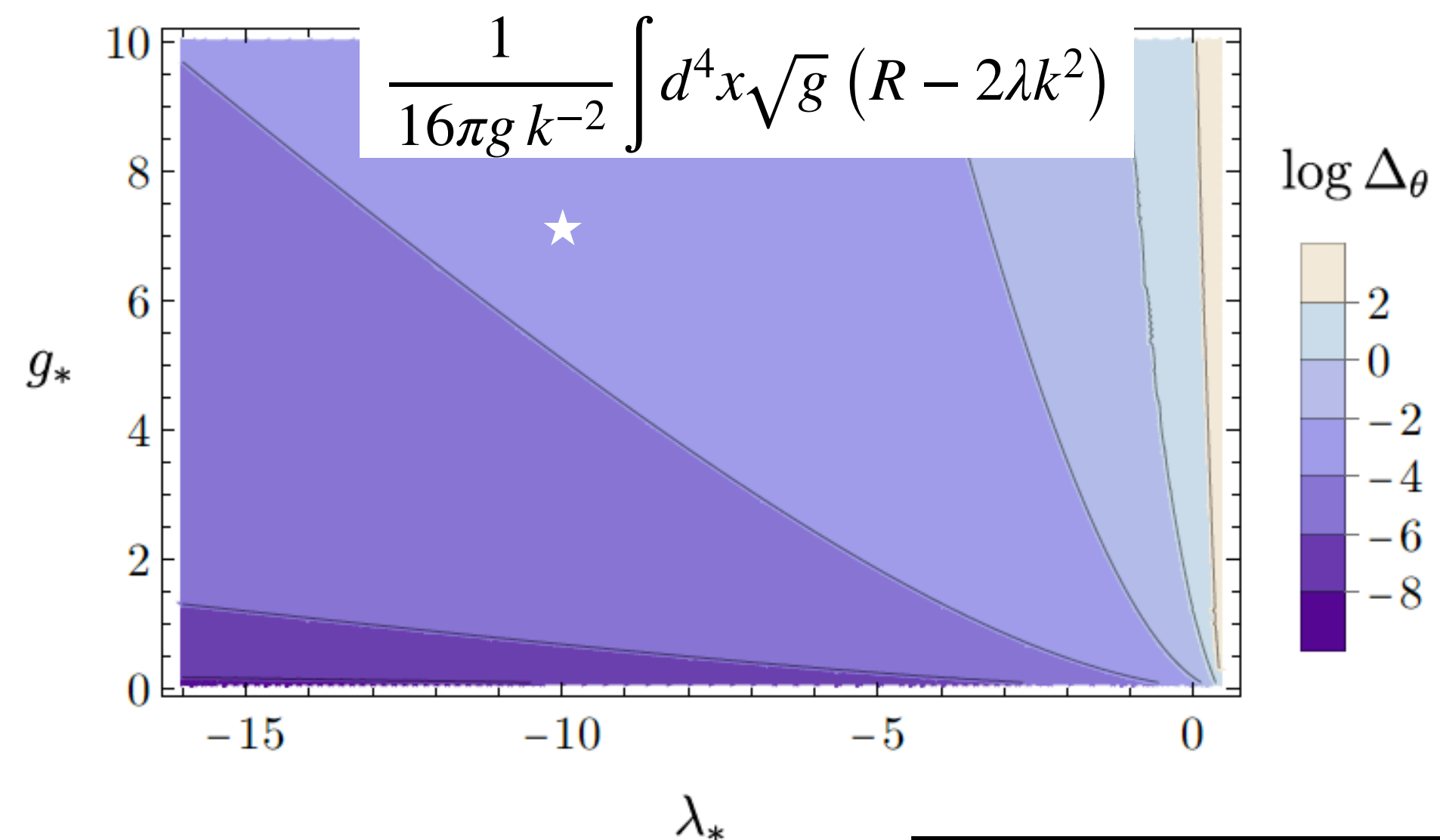
or



?

Key property: near-perturbative

- free parameters $\simeq$ dimension-4-interactions
- similar set as free parameters at perturbative (Gaussian) fixed point

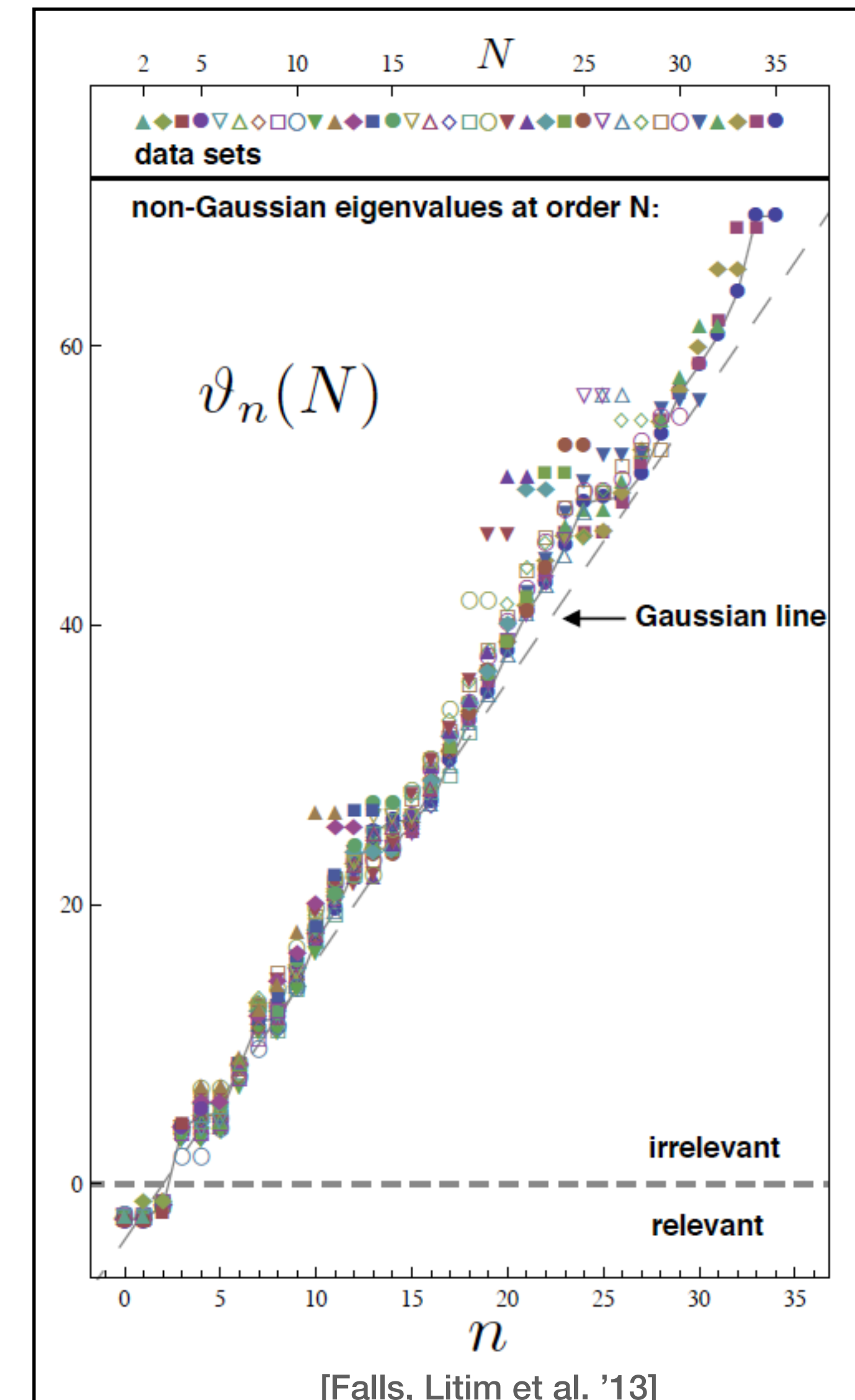


metric propagator:

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$$\Delta_\theta = \sqrt{\frac{\sum_i (\text{Re}(\theta^{(i)}) - \theta_{\text{Gauss}})^2}{\sum_i}}$$

[AE, Pauly '18]



Method & key assumptions

Functional Renormalization Group: based on Euclidean path integral

Γ_k : analog of classical action, but with quantum fluctuations above k included

$$k \partial_k \Gamma_k = \frac{1}{2} \text{Tr} \left[\left(\Gamma_k^{(2)} + R_k \right)^{-1} k \partial_k R_k \right] \quad \rightarrow \quad \beta_g = k \partial_k g(k)$$

[Wetterich '93; Reuter '96]

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Quantitative precision achievable

Example: Fixed point in the Ising model, derivative expansion

derivative expansion	ν	η
$s = 0$ (LPA)	0.651(1)	0
$s = 2$	0.6278(3)	0.0449 (6)
$s = 4$	0.63039(18)	0.0343(7)
$s = 6$	0.63012(5)	0.0361 (3)
$s \rightarrow \infty$	0.6300(2)	0.0358(6)
conformal bootstrap	0.629971(4)	0.0362978(20)

[Balog, Chaté, Delamotte, Marohnić, Wschebor '19]

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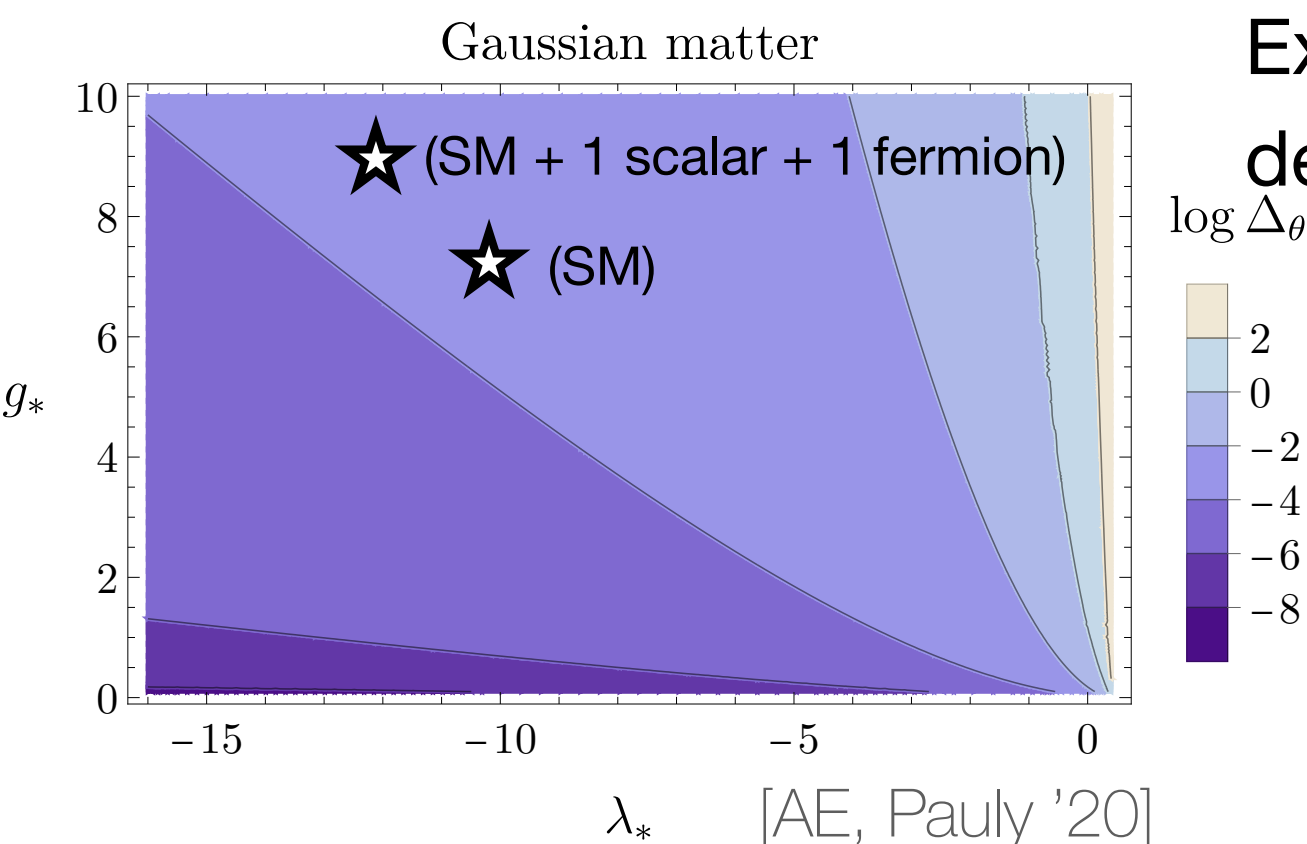
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Truncation scheme for matter-gravity: near-perturbativity as a bootstrap

- assume near-perturbativity:
quantum corrections are subleading compared to canonical scaling
- use canonical power-counting to set up truncations
- check that near-perturbativity holds at fixed point in truncation



Example (SM & BSM Yukawa sector):
deviation from perturbative scaling:

$$\Delta_\theta = \frac{1}{N} \sqrt{\sum_{i=1}^N \left(\theta_i - d_{\bar{g}_i} \right)^2}$$

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Gaussian matter

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[AE, Pauly '20]

Key assumption: Euclidean vs. Lorentzian signature

First hints of Lorentzian asymptotic safety

- impact of foliation on fixed-point structure small

[Biemans, Platania, Saueressig '16 '17; Saueressig, Wang '23]
- calculation in Einstein-Hilbert truncation in Lorentzian signature yields fixed point

[Fehre, Litim, Pawłowski, Reichert '21]