

Neutrino oscillations and decoherence: insights from microscopic models

Renata Ferrero

based on work in collaboration with

Joao Coelho, Alba Domi, Max Joseph Fahn, Kristina Giesel and Roman Kemper

Bridging high and low energies in search of quantum gravity - BridgeQG COST Action

Paris, 8. July 2025



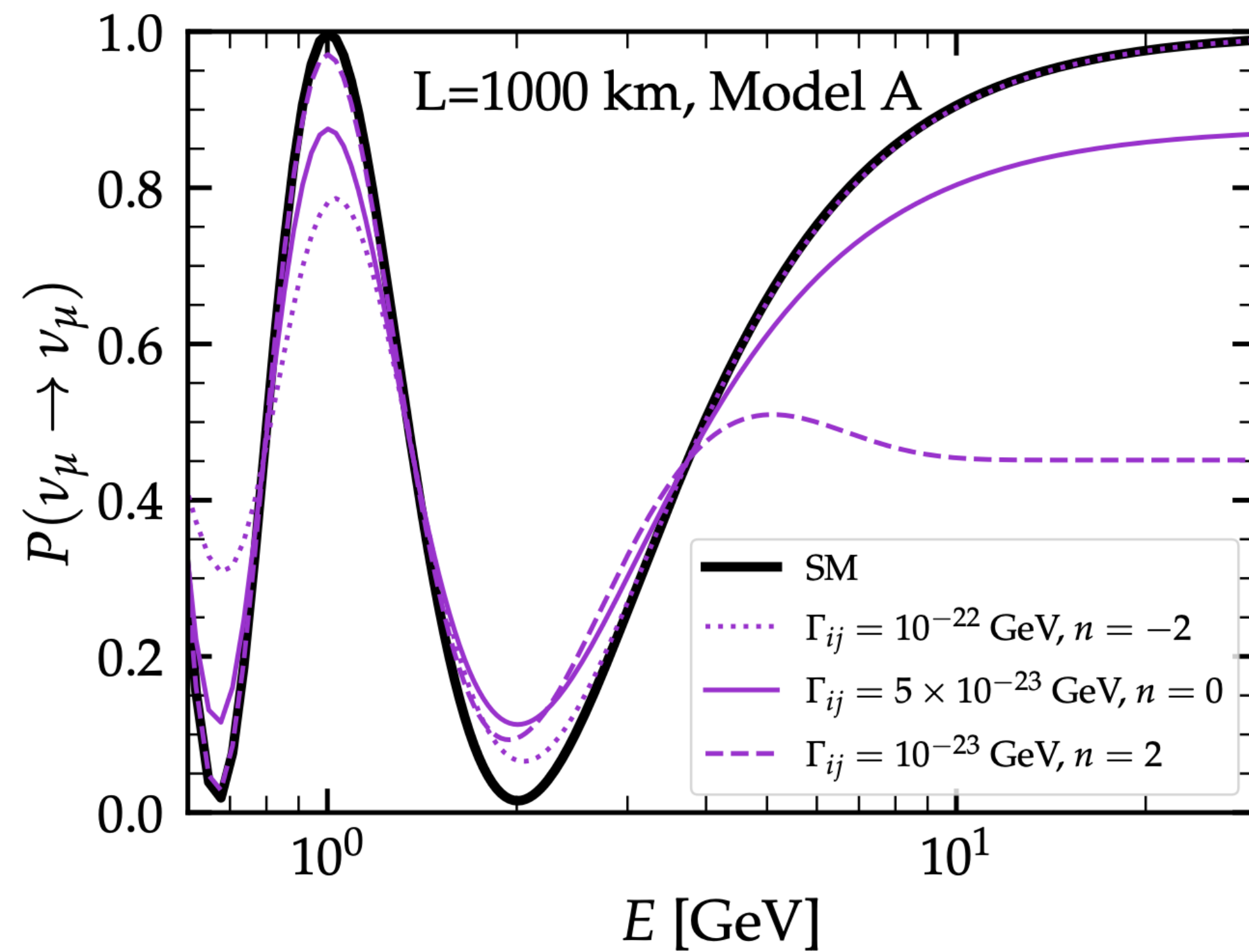
Friedrich-Alexander-Universität
Erlangen-Nürnberg



Erlangen Centre
for Astroparticle
Physics

Motivation

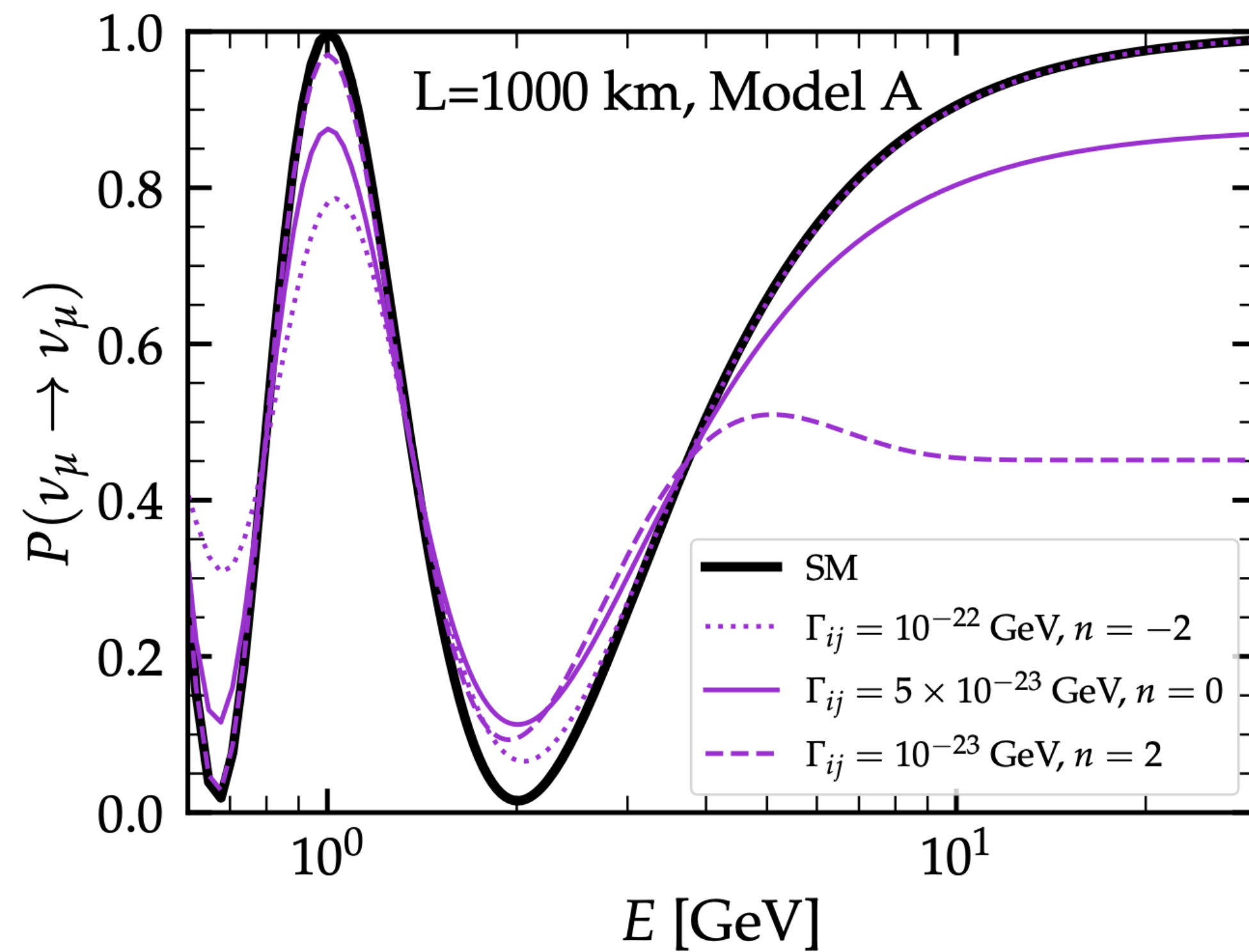
Precision era for decoherence in neutrino oscillations



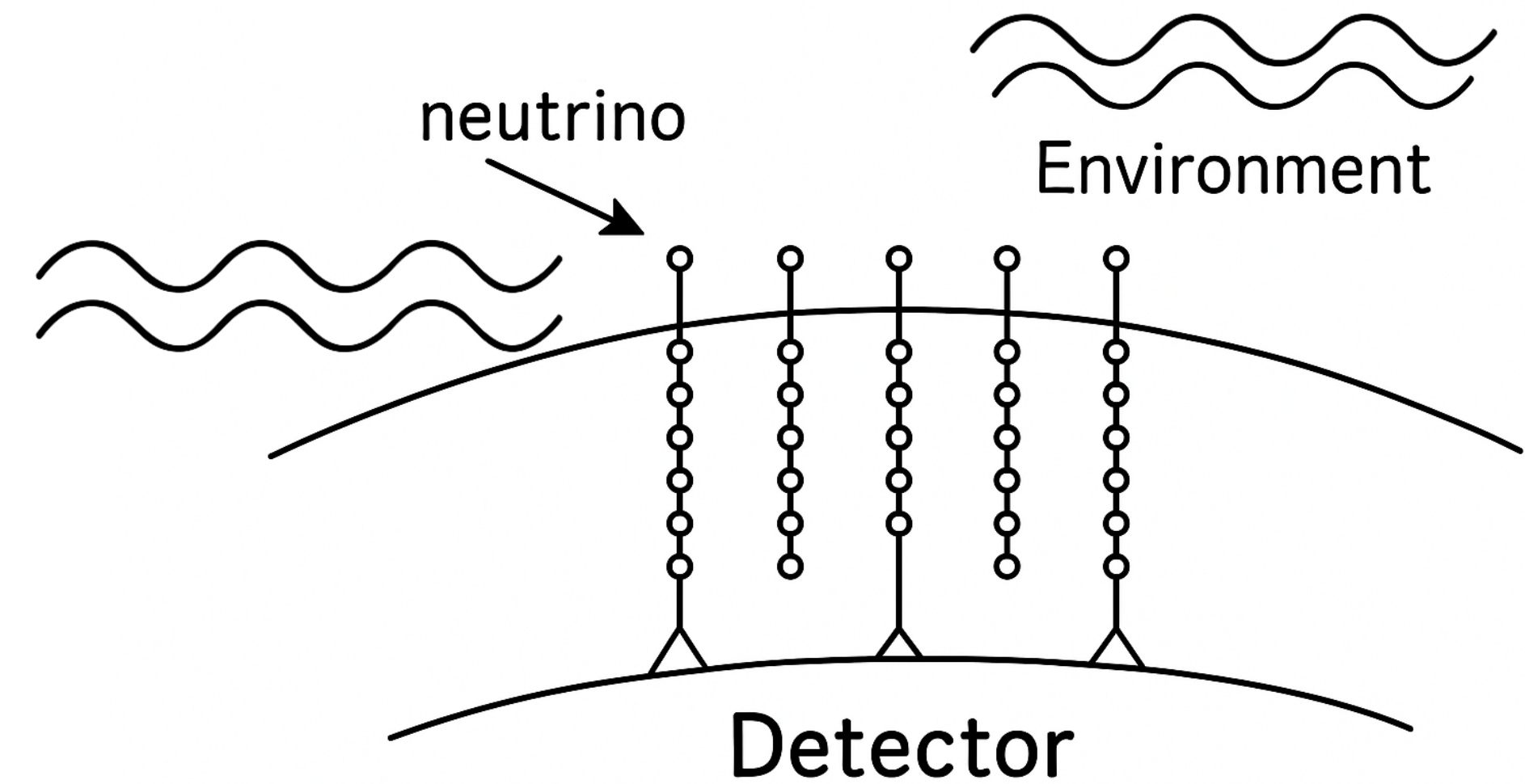
[De Romeri, Giunti, Stuttard, Ternes
2306.14699 [hep-ph]]

Motivation

Precision era for decoherence in neutrino oscillations

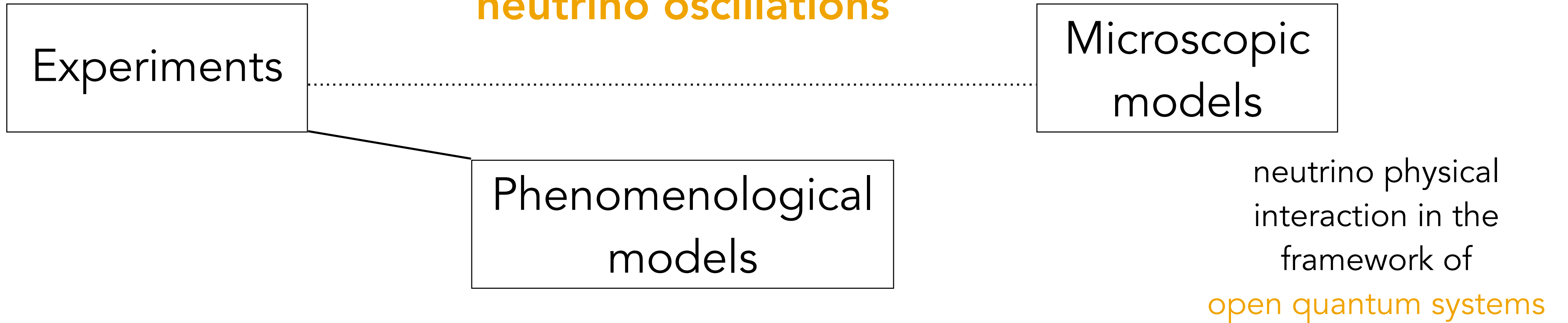


[De Romeri, Giunti, Stuttard, Ternes
2306.14699 [hep-ph]]



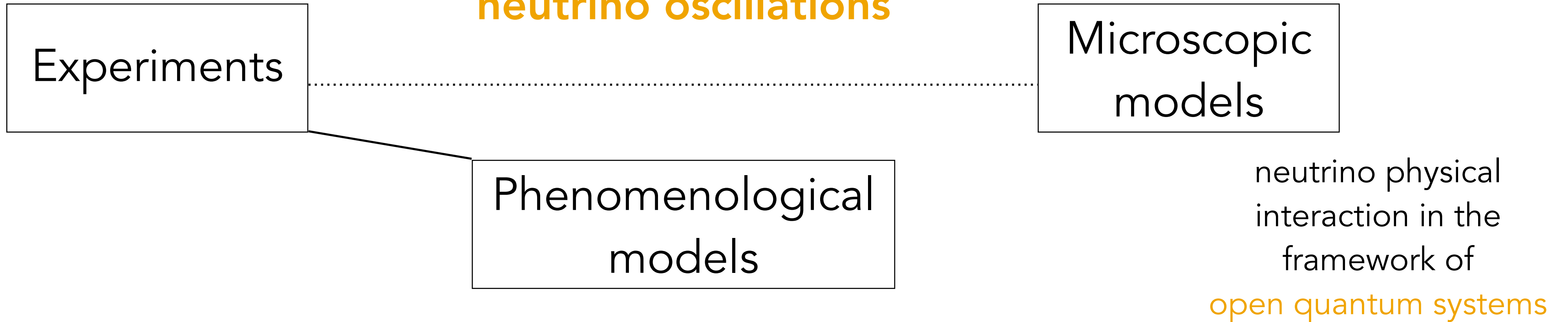
Motivation

Precision era for
decoherence in
neutrino oscillations



Motivation

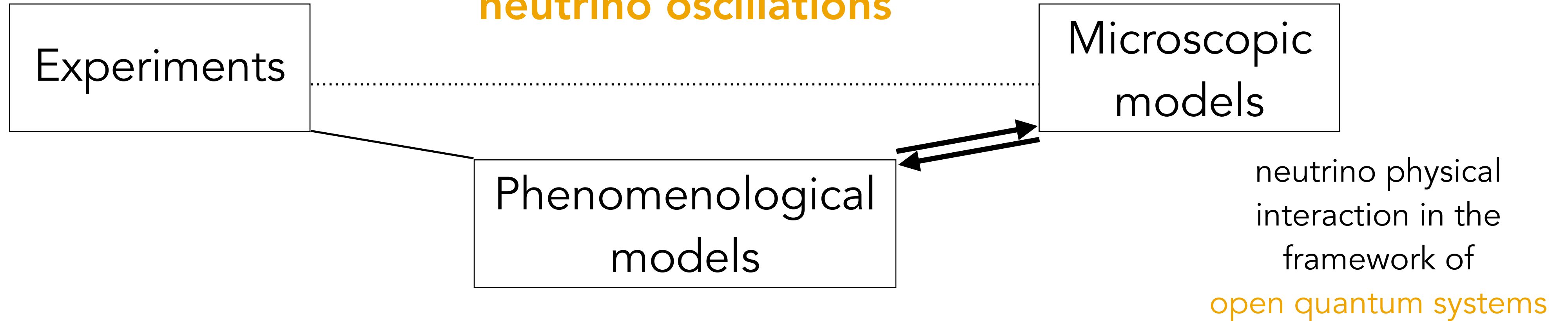
Precision era for
decoherence in
neutrino oscillations



Which is the underlying physics?
Maybe coupling to gravity $\rightarrow$ signature for quantum gravity?

Motivation

Precision era for
decoherence in
neutrino oscillations



Which is the underlying physics?
Maybe coupling to gravity $\rightarrow$ signature for quantum gravity?

Questions

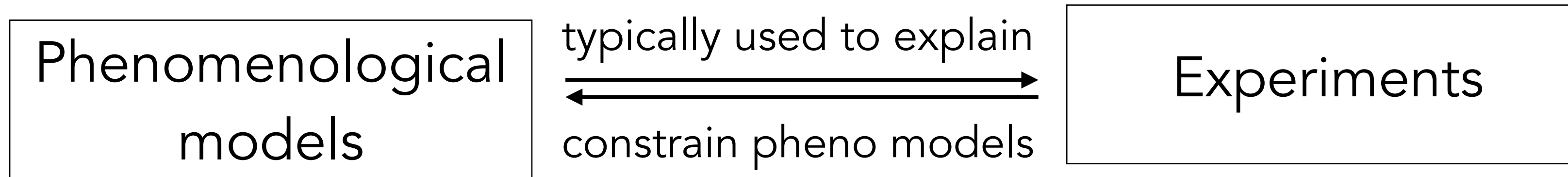
Can one justify via theoretical models a given phenomenological ansatz?

Can a phenomenological model select a preferred microscopic model?

Phenomenological
models

typically used to explain
←→
constrain pheno models

Experiments

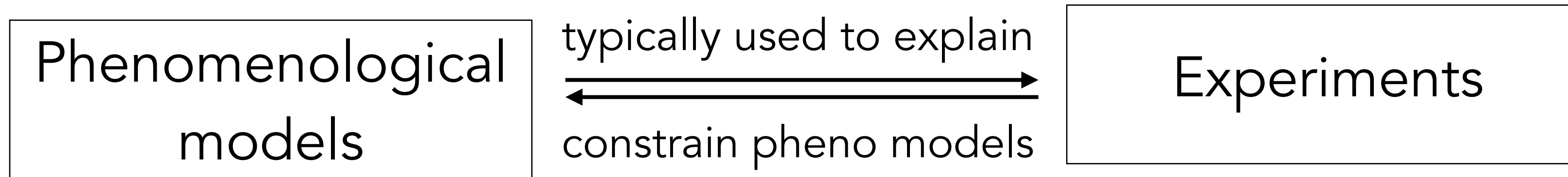


However, there are still some open points:

energy dependence

shape of the dissipator

→ these choices can affect the oscillation probabilities



However, there are still some open points:

energy dependence

shape of the dissipator

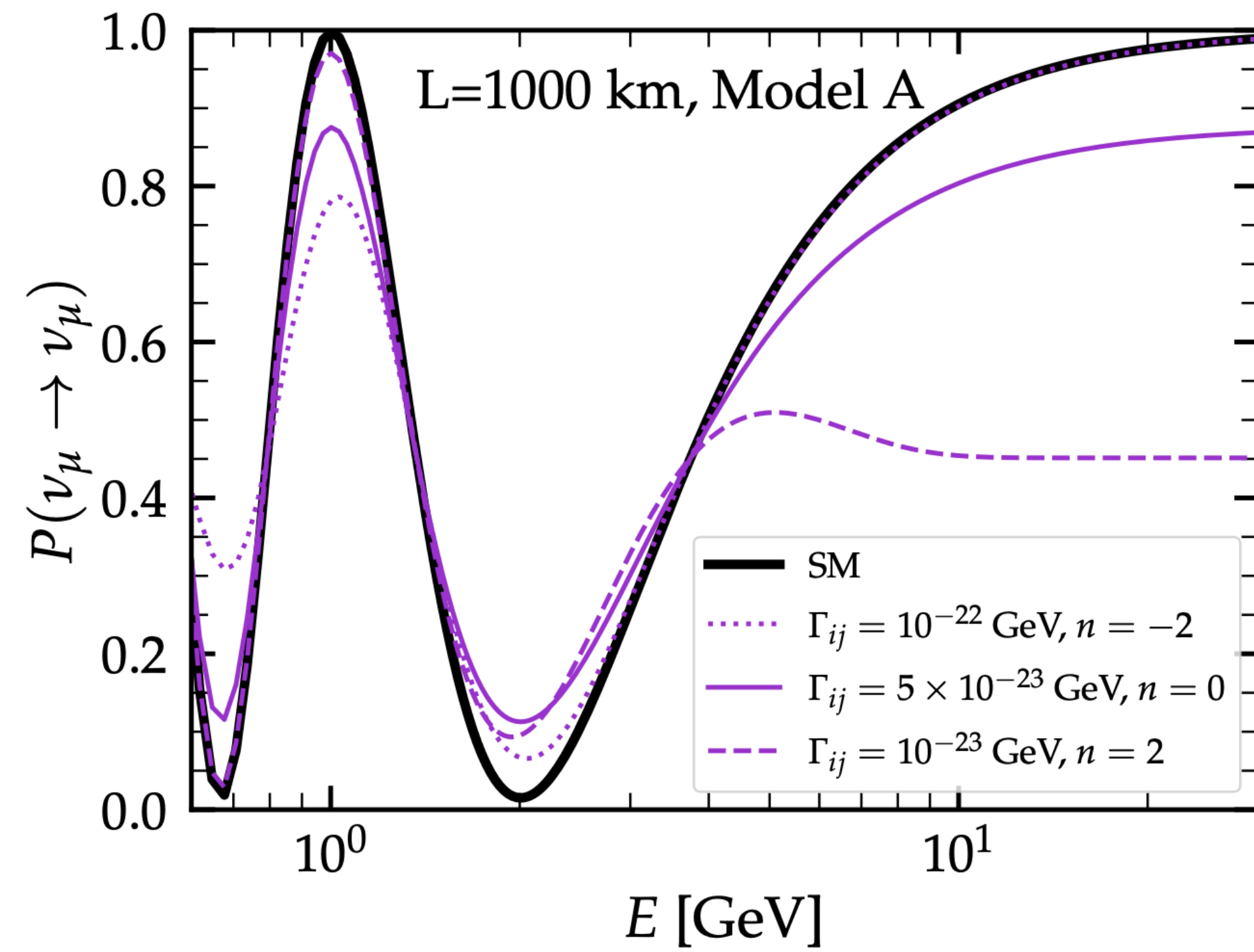
→ these choices can affect the oscillation probabilities

Microscopic
models

Start from **open QFT** models: **underlying microscopic QM** model arises in the limit

→ the coupling between environment and system is given to us by the **action**

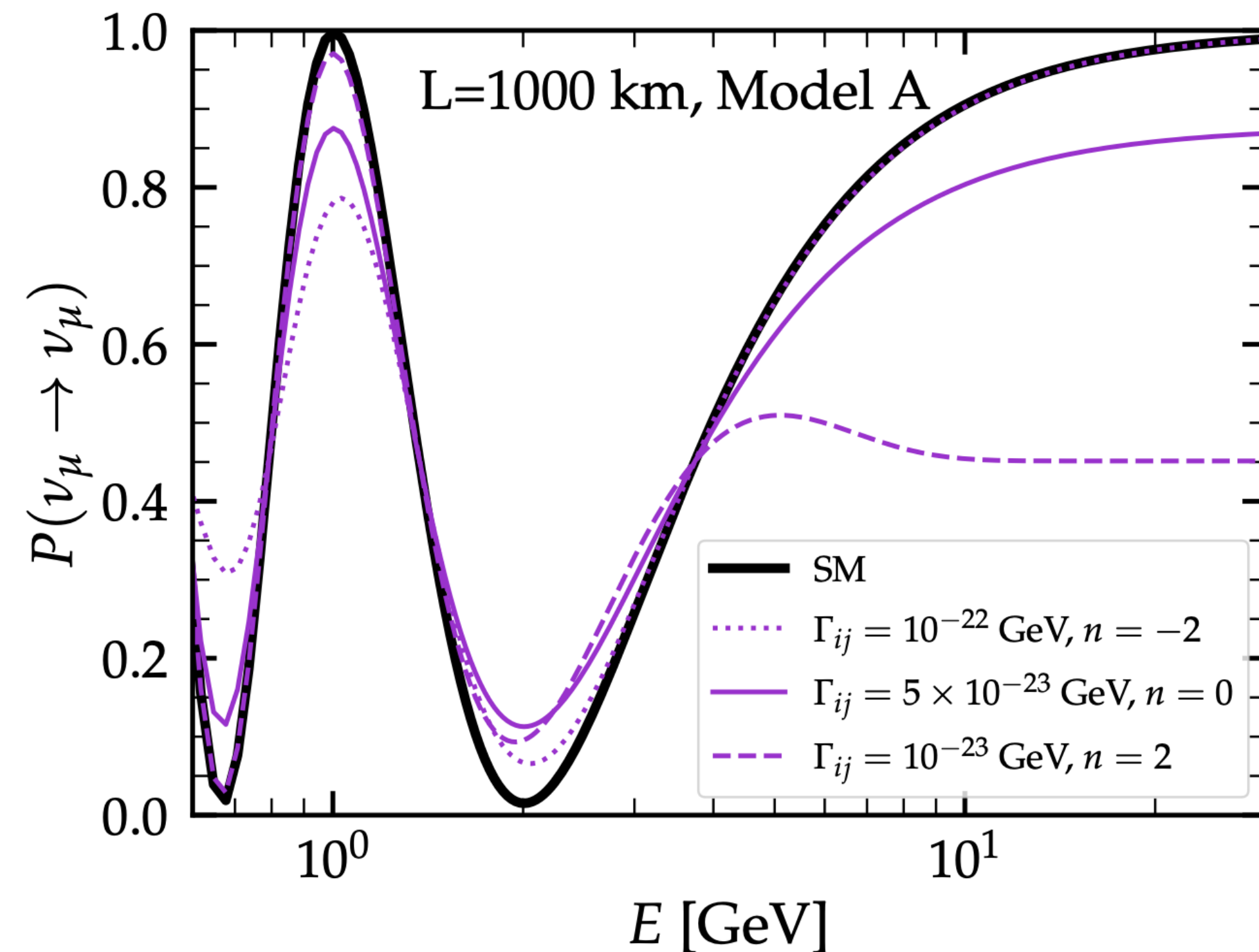
Phenomenological models



[De Romeri, Giunti, Stuttard, Ternes
2306.14699 [hep-ph]]

Phenomenological models

from the **dissipator**:
contains both
decoherence &
dissipation



[De Romeri, Giunti, Stuttard, Ternes
2306.14699 [hep-ph]]

$$\rho_{ij}(t) = \rho_{ij}(0) \cdot e^{-\frac{i}{\hbar}(H_i - H_j)t - \Gamma_{ij}t}$$

$$\Gamma_{ij} = \gamma_{ij} E^n, \quad n \in [-2, 3]$$

E^n : energy dependence

γ_{ij} : shape of the dissipator

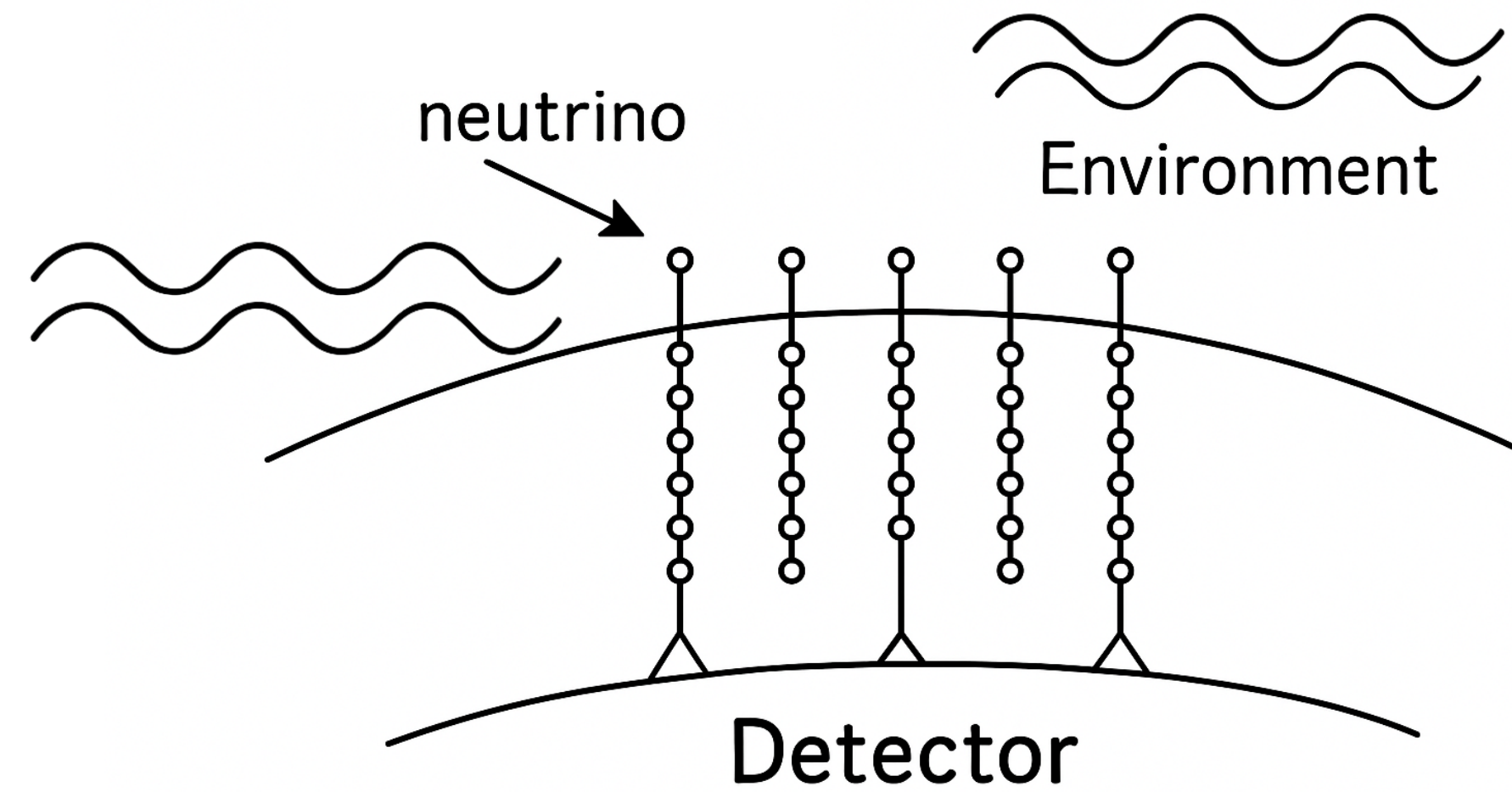
Microscopic
models

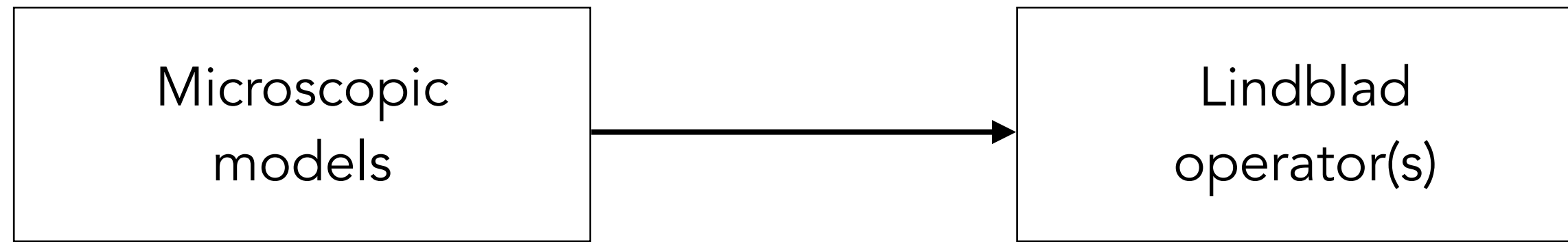
Microscopic
models

QFT model $\rightarrow$ QM model:

Hamiltonian (environment
+ system + interaction)

$$\hat{H} = \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}} , \quad \hat{H}_{\text{int}} = \hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$





QFT model → QM model:

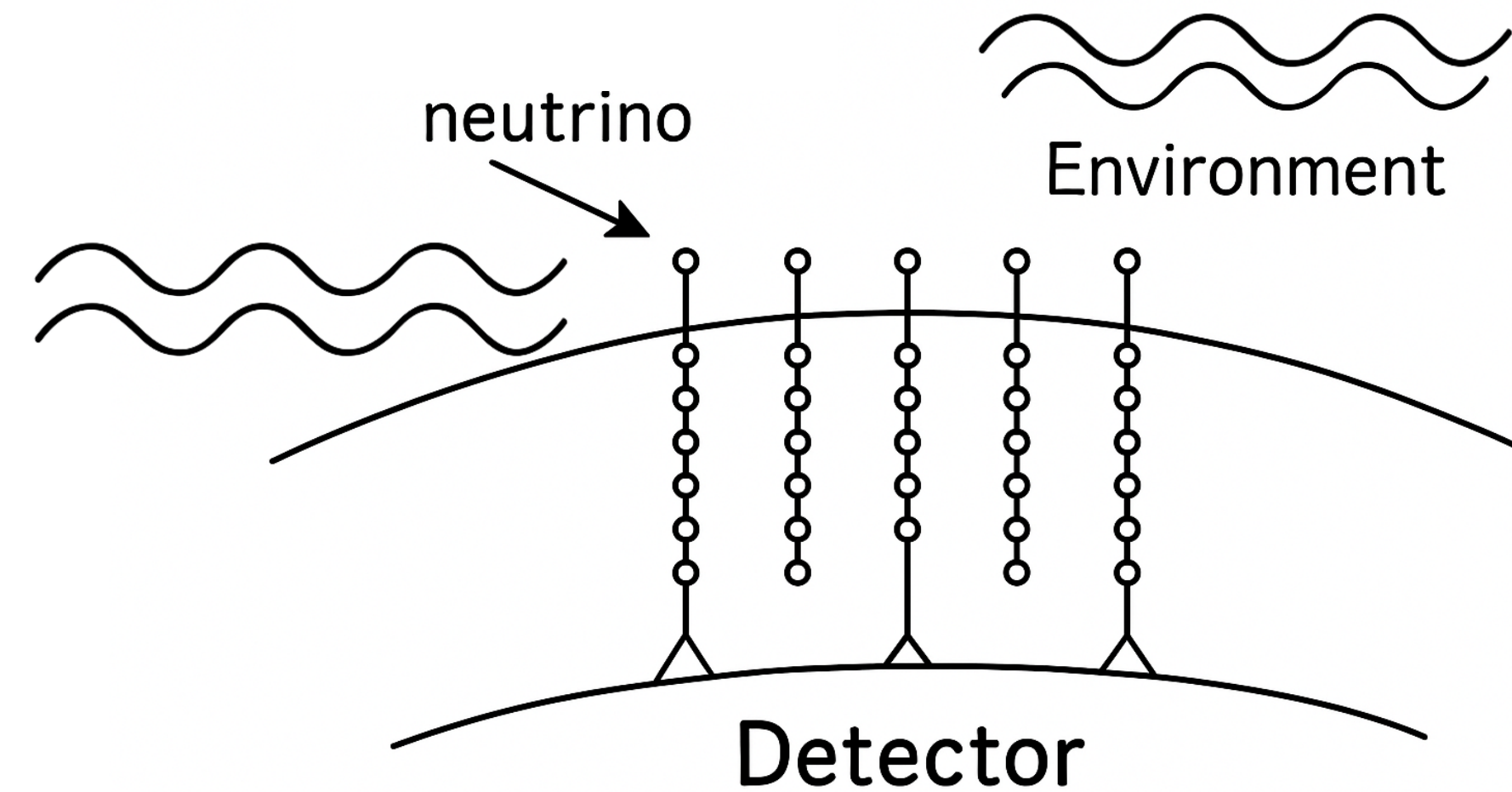
Hamiltonian (environment
+ system + interaction)

$$\hat{H} = \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}} , \quad \hat{H}_{\text{int}} = \hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$

Trace out environment dofs:

Lindblad master equation

(Lindblad operators)



$$\frac{d\hat{\rho}(t)}{dt} = -i[\hat{H}_S, \hat{\rho}(t)] + \underbrace{\frac{1}{2} \sum_j \left([\hat{A}_j, \hat{\rho}(t) \hat{A}_j^\dagger] + [\hat{A}_j \hat{\rho}(t), \hat{A}_j^\dagger] \right)}_{\text{dissipator}}$$

dissipator

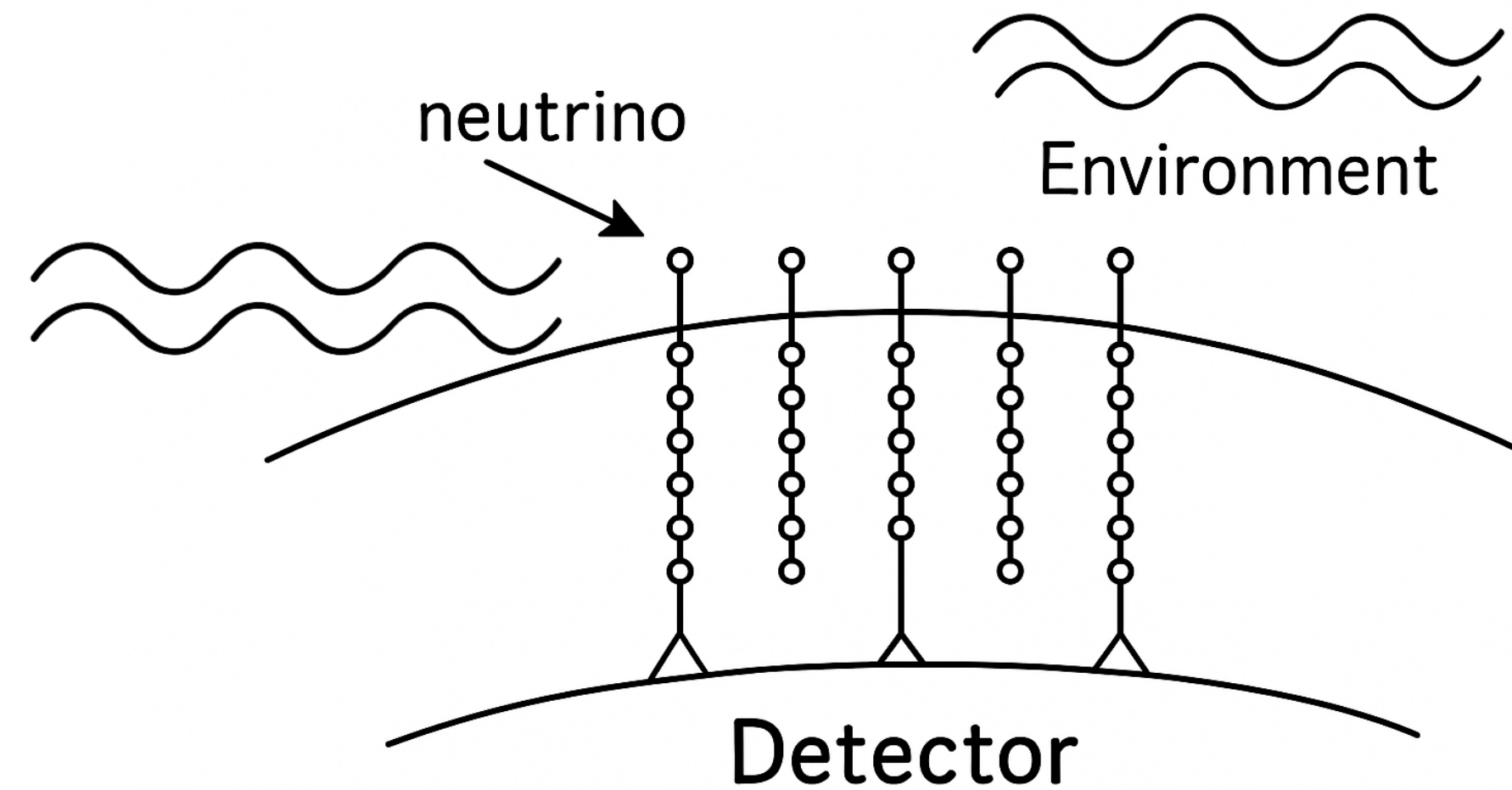
* under the assumption
that the master equation
is of Lindblad form



QFT model → QM model:

Hamiltonian (environment
+ system + interaction)

$$\hat{H} = \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}} , \quad \hat{H}_{\text{int}} = \hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$



Trace out environment dofs:
Lindblad master equation
(Lindblad operators)

$$\frac{d\hat{\rho}(t)}{dt} = -i[\hat{H}_S, \hat{\rho}(t)] + \underbrace{\frac{1}{2} \sum_j \left([\hat{A}_j, \hat{\rho}(t) \hat{A}_j^\dagger] + [\hat{A}_j \hat{\rho}(t), \hat{A}_j^\dagger] \right)}_{\text{dissipator}}$$

Vectorize the master
equation to write down
dissipator

$$\frac{d}{dt} |\rho\rangle\rangle = \mathbb{L} |\rho\rangle\rangle$$

$$\frac{d}{dt} \rho_\alpha(t) = (h^\mu f_{\mu\nu\alpha} + D_{\nu\alpha}) \rho^\nu(t)$$

dissipator * under the assumption
that the master equation
is of Lindblad form

Energy dependence

Rationale: positive powers, decoherence stronger at high E
negative powers, decoherence stronger at low E

high energy atmospheric
solar neutrinos, low energy
atmospheric

Energy dependence

Rationale: positive powers, decoherence stronger at high E
negative powers, decoherence stronger at low E

high energy atmospheric
solar neutrinos, low energy
atmospheric

Consider for instance the
QM model

$$\begin{aligned}\hat{H} &= \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}} \\ &= \hat{H}_S + \frac{1}{2} \sum_{i=1}^N [\hat{p}_i^2 + \omega_i^2 \hat{q}_i^2] - \hat{H}_S \otimes \sum_{i=1}^N g_i \hat{q}_i\end{aligned}$$

inspired by coupling
via linearized gravity

[Blencowe, Xu
2005.02554 [quant-
[Fahn, Giesel
2409.12790 [hep-th]]

Energy dependence

Rationale: positive powers, decoherence stronger at high E
negative powers, decoherence stronger at low E

high energy atmospheric
solar neutrinos, low energy
atmospheric

Consider for instance the
QM model

$$\rightarrow \left\{ \begin{array}{l} \gamma_{ij} = \frac{\eta^2 c^8 k_B T}{\hbar^3} (\Delta m_{ij}^2)^2 \\ n = -2 \end{array} \right.$$

$$\begin{aligned} \hat{H} &= \hat{H}_S + \hat{H}_{\mathcal{E}} + \hat{H}_{\text{int}} \\ &= \hat{H}_S + \frac{1}{2} \sum_{i=1}^N [\hat{p}_i^2 + \omega_i^2 \hat{q}_i^2] - \hat{H}_S \otimes \sum_{i=1}^N g_i \hat{q}_i \end{aligned}$$

inspired by coupling
via linearized gravity

[Blencowe, Xu
2005.02554 [quant-
[Fahn, Giesel
2409.12790 [hep-th]]

[Roman Kemper's talk]

[Domi, Eberl, Fahn, Giesel, Hennig, Katz, Kemper, Kobler 2403.03106 [gr-qc]]

Energy dependence

Rationale: positive powers, decoherence stronger at high E
negative powers, decoherence stronger at low E

high energy atmospheric
solar neutrinos, low energy
atmospheric

Consider for instance the
QM model

$$\rightarrow \begin{cases} \gamma_{ij} = \frac{\eta^2 c^8 k_B T}{\hbar^3} (\Delta m_{ij}^2)^2 \\ n = -2 \end{cases}$$

$$\begin{aligned} \hat{H} &= \hat{H}_S + \hat{H}_{\mathcal{E}} + \hat{H}_{\text{int}} \\ &= \hat{H}_S + \frac{1}{2} \sum_{i=1}^N [\hat{p}_i^2 + \omega_i^2 \hat{q}_i^2] - \hat{H}_S \otimes \sum_{i=1}^N g_i \hat{q}_i \end{aligned}$$

inspired by coupling
via linearized gravity

[Blencowe, Xu
2005.02554 [quant-
[Fahn, Giesel
2409.12790 [hep-th]]

[Roman Kemper's talk]

[Domi, Eberl, Fahn, Giesel, Hennig, Katz, Kemper, Kobler 2403.03106 [gr-qc]]

Generalized model

$$\hat{H} = \hat{H}_S + \frac{1}{2} \sum_{i=1}^N [\hat{p}_i^2 + \omega_i^2 \hat{q}_i^2] - f(\hat{H}_S) \otimes \sum_{i=1}^N g_i \hat{q}_i$$

Ansatz

$$f(\hat{H}_S) = \hat{H}_S^{k/2}, \quad k \geq 1 \text{ and } k \in \mathbb{N} \quad \rightarrow \quad \Gamma_{ij} \approx E^{k-4} \frac{k^2 (\Delta m_{ij}^2)^2}{16} \propto E^{k-4}$$

Energy dependence

If $f(\hat{H}_S) \sim \hat{H}_S^0 \rightarrow$ no decoherence

If $f(\hat{H}_S) \sim \hat{H}_S^{1/2} \rightarrow n = -3$

If $f(\hat{H}_S) \sim \hat{H}_S \rightarrow n = -2$

If $f(\hat{H}_S) \sim \hat{H}_S^{3/2} \rightarrow n = -1$

If $f(\hat{H}_S) \sim \hat{H}_S^2 \rightarrow n = 0$

If $f(\hat{H}_S) \sim \hat{H}_S^{5/2} \rightarrow n = 1$

If $f(\hat{H}_S) \sim \hat{H}_S^3 \rightarrow n = 2$

$\vdots$
higher order

Energy dependence

If $f(\hat{H}_S) \sim \hat{H}_S^0 \rightarrow$ no decoherence

If $f(\hat{H}_S) \sim \hat{H}_S^{1/2} \rightarrow n = -3$

If $f(\hat{H}_S) \sim \hat{H}_S \rightarrow n = -2$

If $f(\hat{H}_S) \sim \hat{H}_S^{3/2} \rightarrow n = -1$

If $f(\hat{H}_S) \sim \hat{H}_S^2 \rightarrow n = 0$

If $f(\hat{H}_S) \sim \hat{H}_S^{5/2} \rightarrow n = 1$

If $f(\hat{H}_S) \sim \hat{H}_S^3 \rightarrow n = 2$

$\vdots$
higher order

Linearized gravity model,
lightcone fluctuations

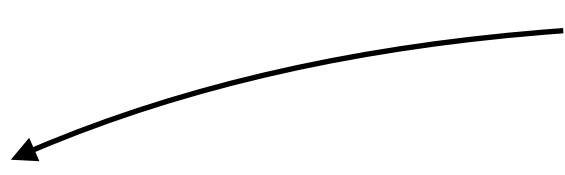
Lorentz invariant model

Energy independent

Cross-section-like,
quantum mechanics violation

String motivated model

different environments
(if at all specified),
different physics



[Studdard 2103.15313 [hep-ph]]

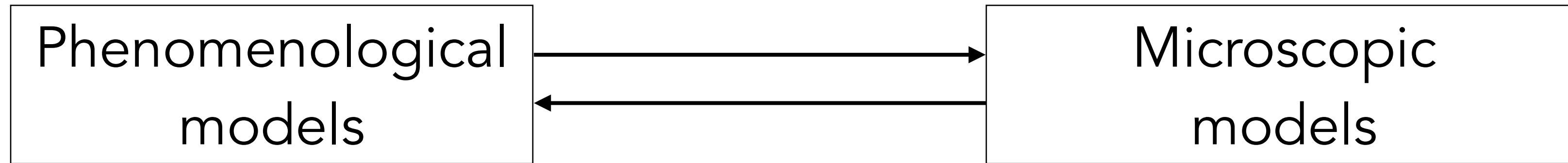
[Lisi, Marrone and Montanino hep-ph/0002053]

[Liu, Hu, and Ge, 1997]

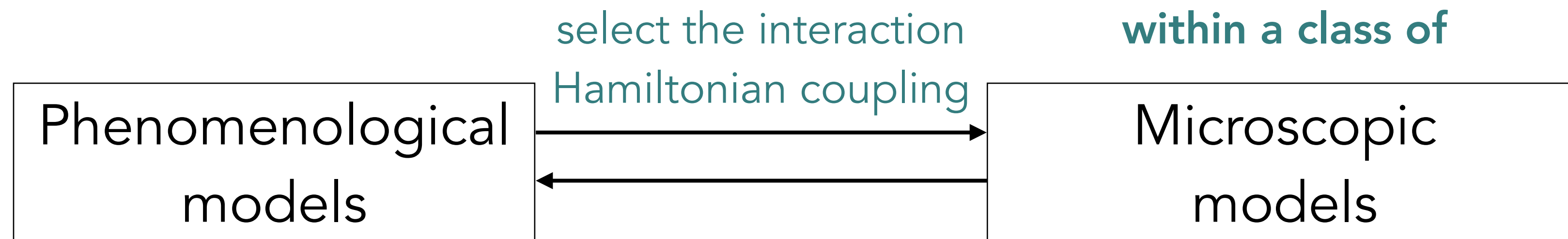
[Ellis, Lopez, Mavromatos, Nanopoulos,
Winstanley, hep-ph/9505340, gr-qc/9602011]

[D'Esposito, Gubitosi 2306.14778 [hep-ph]]

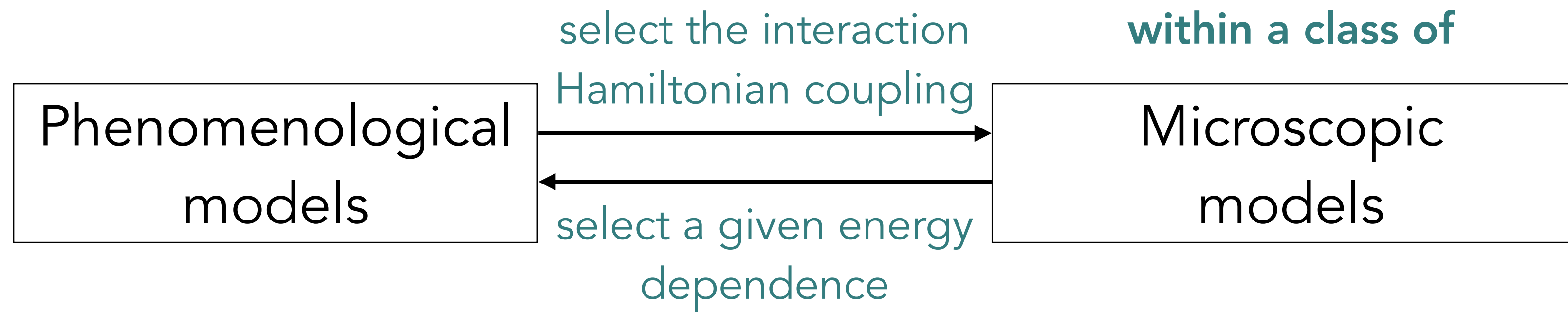
Energy dependence



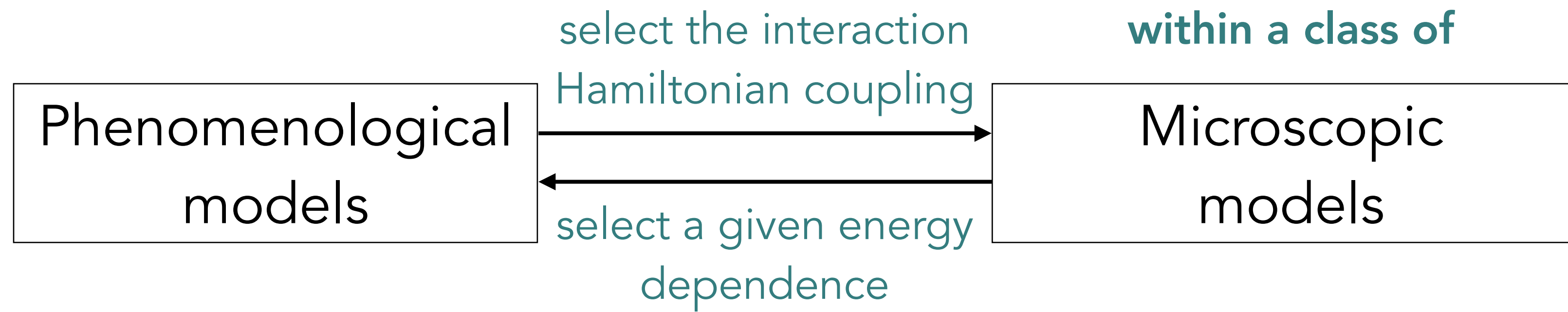
Energy dependence



Energy dependence

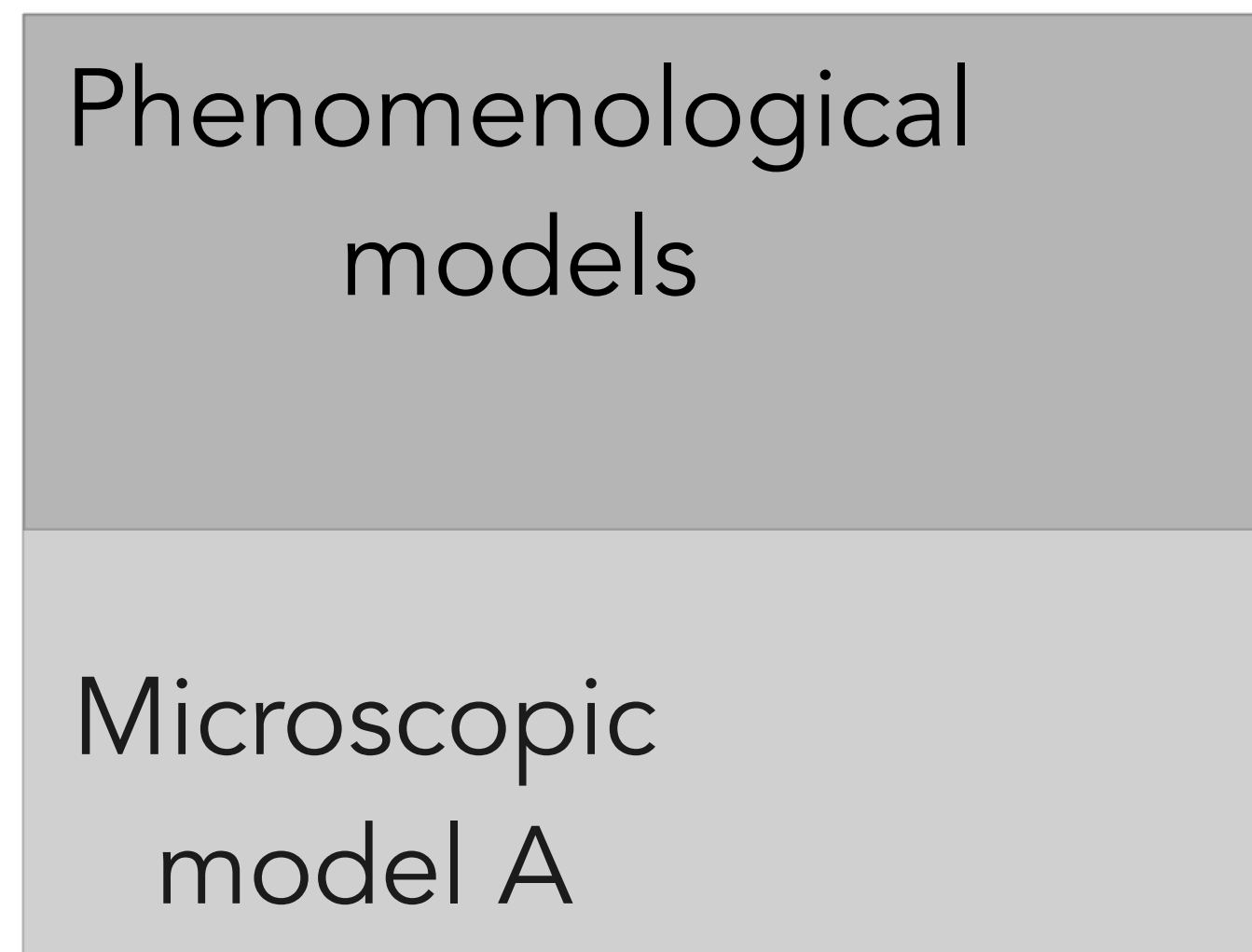
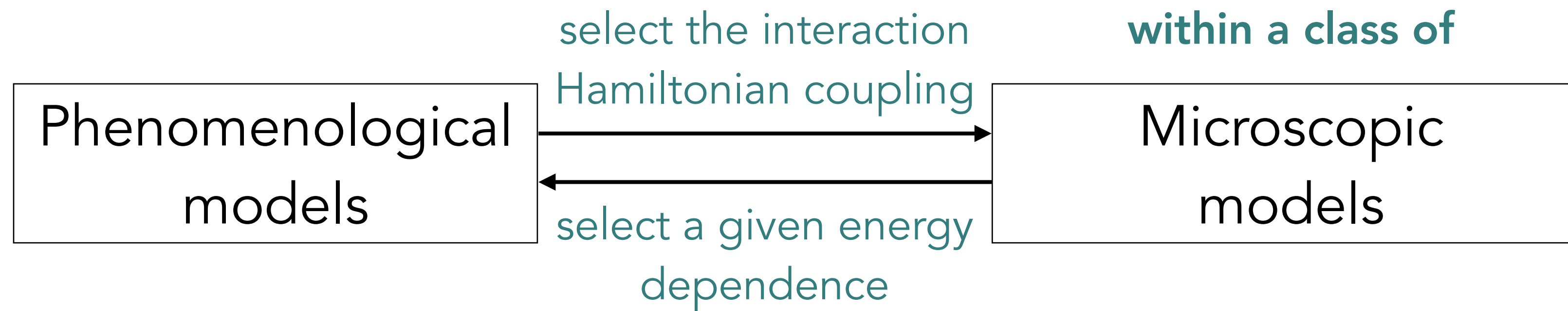


Energy dependence

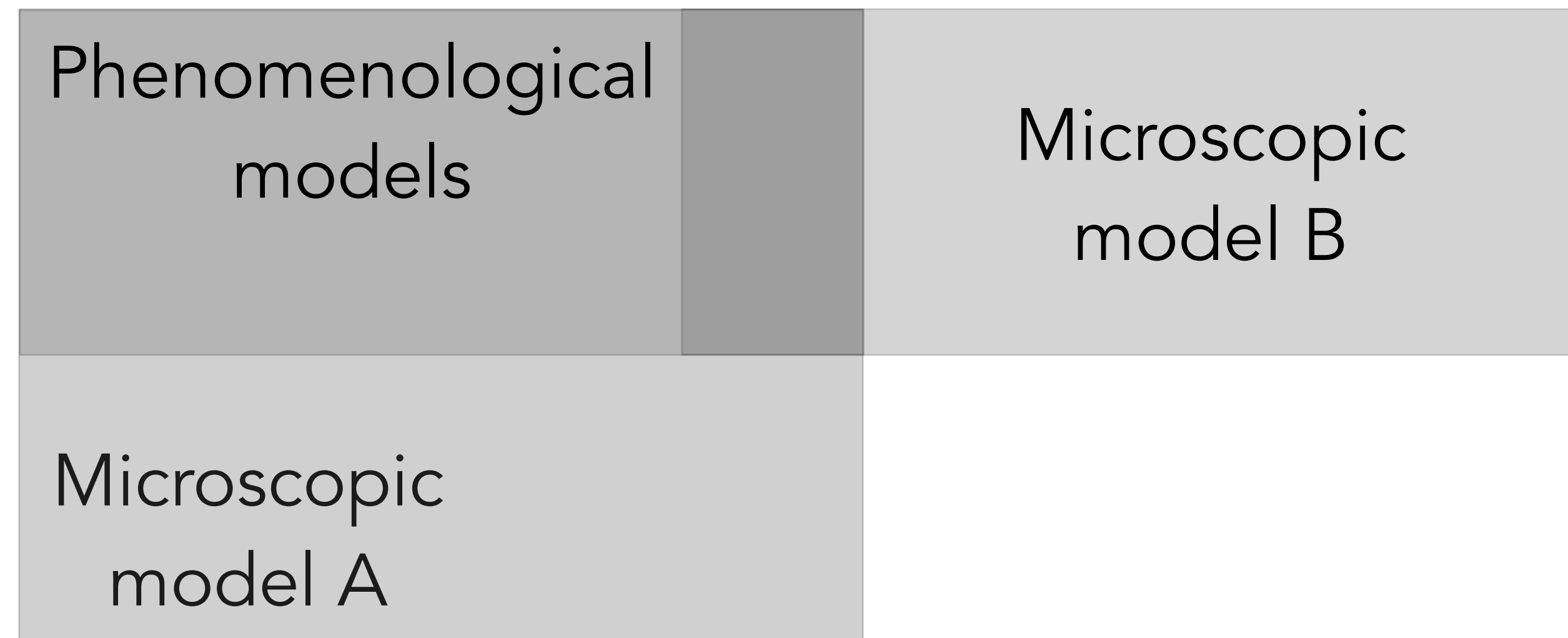
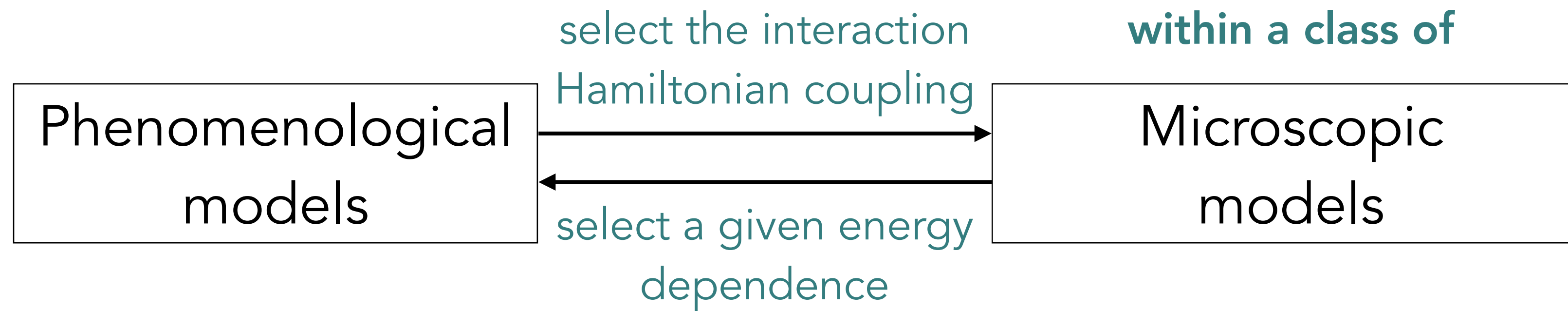


Phenomenological
models

Energy dependence



Energy dependence



Shape of the dissipator

Rationale: Relaxation time (energy loss) $\gg$ Decoherence time

Expectation: solar and high energy atmospheric neutrinos' oscillations will "see" also dissipation

[Gomes, Gomes, Peres 2001.09250 [hep-ph]]

Shape of the dissipator

Rationale: Relaxation time (energy loss) $\gg$ Decoherence time

Expectation: solar and high energy atmospheric neutrinos' oscillations will "see" also dissipation

[Gomes, Gomes, Peres 2001.09250 [hep-ph]]

In **vectorized** form

$$\frac{d}{dt}\rho_{\alpha}(t) = (h^{\mu}f_{\mu\nu\alpha} + D_{\nu\alpha})\rho^{\nu}(t)$$

ansatz for the
dissipator

motivated e.g. by
whether one expects
only astrophysical
neutrinos to have
dissipation

Shape of the dissipator

Rationale: Relaxation time (energy loss) $\gg$ Decoherence time

Expectation: solar and high energy atmospheric neutrinos' oscillations will "see" also dissipation

[Gomes, Gomes, Peres 2001.09250 [hep-ph]]

In **vectorized** form

$$\frac{d}{dt}\rho_\alpha(t) = (h^\mu f_{\mu\nu\alpha} + D_{\nu\alpha}) \rho^\nu(t)$$

ansatz for the
dissipator

motivated e.g. by
whether one expects
only astrophysical
neutrinos to have
dissipation

From the open quantum system perspective:



SU(2) $\hat{A}_S = \sum_{i=1}^3 a_i \hat{\sigma}_i \longrightarrow D_{\mu\nu} = -\frac{1}{2} \begin{pmatrix} (a_2)^2 + (a_3)^2 & -a_1 a_2 & -a_1 a_3 \\ -a_1 a_2 & (a_1)^2 + (a_3)^2 & -a_2 a_3 \\ -a_1 a_3 & -a_2 a_3 & (a_1)^2 + (a_2)^2 \end{pmatrix}$

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**



depend on the Lindblad
operator(s) $\hat{A}_S$ and on Γ

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**

Dissipation: from the microscopic model

Lindblad operators



specific form of the dissipator with
related non-zero off-diagonal entries

depend on the Lindblad
operator(s) $\hat{A}_S$ and on Γ

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**

Dissipation: from the microscopic model

depend on the Lindblad
operator(s) $\hat{A}_S$ and on Γ

Lindblad operators



specific form of the dissipator with
related non-zero off-diagonal entries

Consider the ansatz

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, -\chi)$$

[Oliveira 1603.08065 [hep-ph]]

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**

depend on the Lindblad
operator(s) $\hat{A}_S$ and on Γ

Dissipation: from the microscopic model

Lindblad operators



specific form of the dissipator with
related non-zero off-diagonal entries

Consider the ansatz

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, -\chi)$$

[Oliveira 1603.08065 [hep-ph]]

Comparing: it cannot be obtained with one Lindblad operator, **we need at least 2**

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**

depend on the Lindblad operator(s) $\hat{A}_S$ and on Γ

Dissipation: from the microscopic model

Lindblad operators



specific form of the dissipator with
related non-zero off-diagonal entries

Consider the ansatz

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, -\chi)$$

[Oliveira 1603.08065 [hep-ph]]

Comparing: it cannot be obtained with one Lindblad operator, **we need at least 2**

E.g. 2 Lindblad operators



$$D_{\mu\nu} = -\frac{1}{2} \text{diag}(a^2, a^2, 2a^2)$$

the entries are related

Shape of the dissipator

Energy conservation: The dissipator is diagonal.

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, 0)$$

two flavour case **SU(2)**

$$D_{\mu\nu} = \text{diag}(0, -\Gamma_{21}, -\Gamma_{21}, 0, -\Gamma_{31}, -\Gamma_{31}, -\Gamma_{32}, -\Gamma_{32}, 0)$$

three flavour case **SU(3)**

depend on the Lindblad operator(s) $\hat{A}_S$ and on Γ

Dissipation: from the microscopic model

Lindblad operators



specific form of the dissipator with
related non-zero off-diagonal entries

Consider the ansatz

$$D_{\mu\nu} = \text{diag}(-\gamma, -\gamma, -\chi)$$

[Oliveira 1603.08065 [hep-ph]]

Comparing: it cannot be obtained with one Lindblad operator, **we need at least 2**

E.g. 2 Lindblad operators



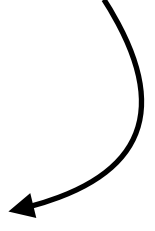
$$D_{\mu\nu} = -\frac{1}{2} \text{diag}(a^2, a^2, 2a^2)$$

the entries are related

Which is then the underlying system-environment physics?

Take home message

The interaction environment-system determines the dissipator,
energy dependence and shape of dissipator

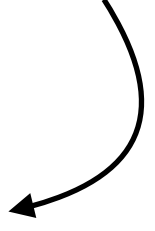
$$\hat{H} = \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}}$$
$$\hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$


* under the assumption that the master equation is of Lindblad form

Take home message

The interaction environment-system determines the dissipator,
energy dependence and shape of dissipator

$$\hat{H} = \hat{H}_S + \hat{H}_E + \hat{H}_{\text{int}}$$

$$\hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$


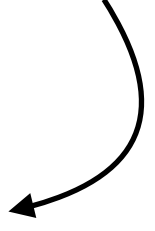
both affect the Lindblad
operators, hence the
shape of the dissipator

* under the assumption that the master equation is of Lindblad form

Take home message

The interaction environment-system determines the dissipator,
energy dependence and shape of dissipator

$$\hat{H} = \hat{H}_S + \hat{H}_\mathcal{E} + \hat{H}_{\text{int}}$$

$$\hat{S} \otimes \sum_{i=1}^N g_i \hat{E}_i$$


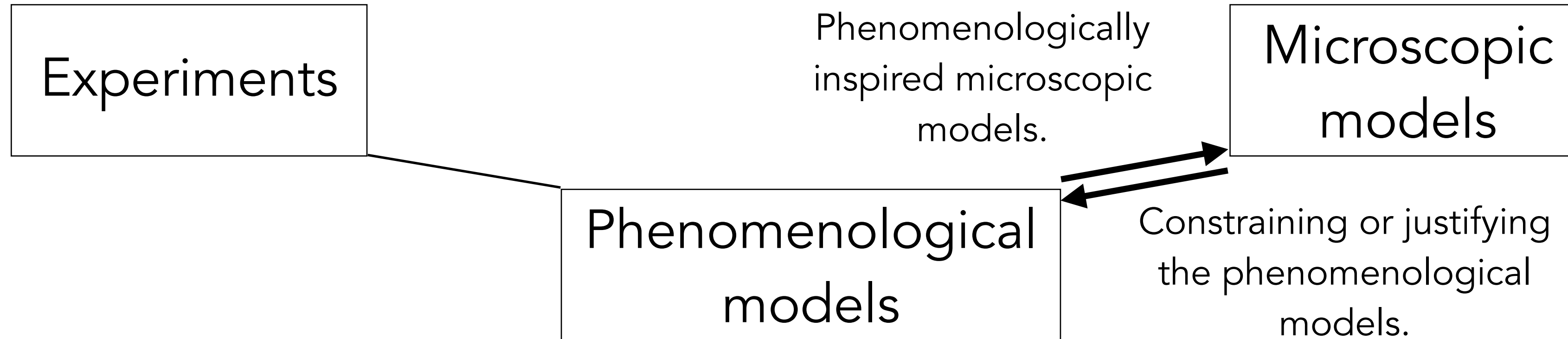
both affect the Lindblad
operators, hence the
shape of the dissipator

both affect the
decoherence rate, hence
the energy dependence

* under the assumption that the master equation is of Lindblad form

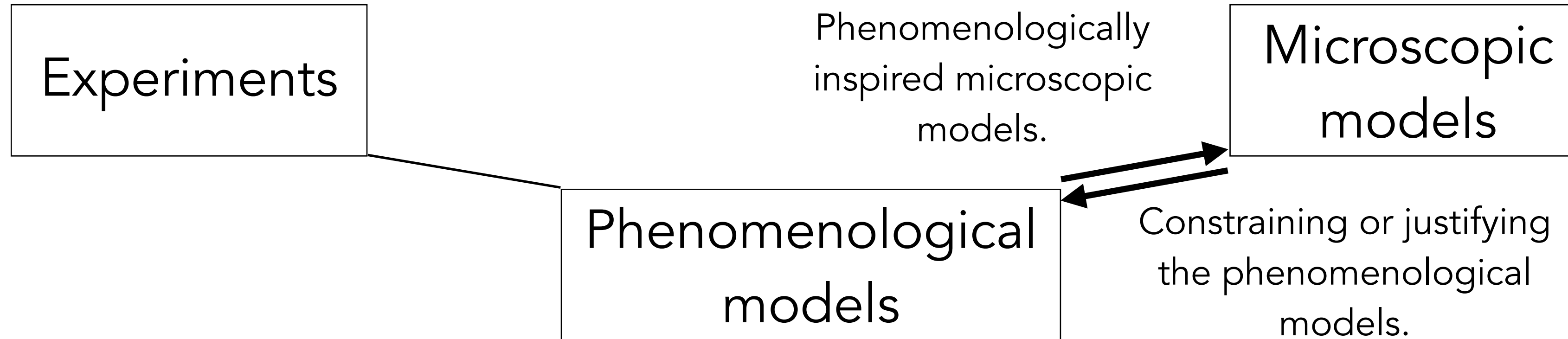
Conclusions

Precision era for decoherence in neutrino oscillations



Conclusions

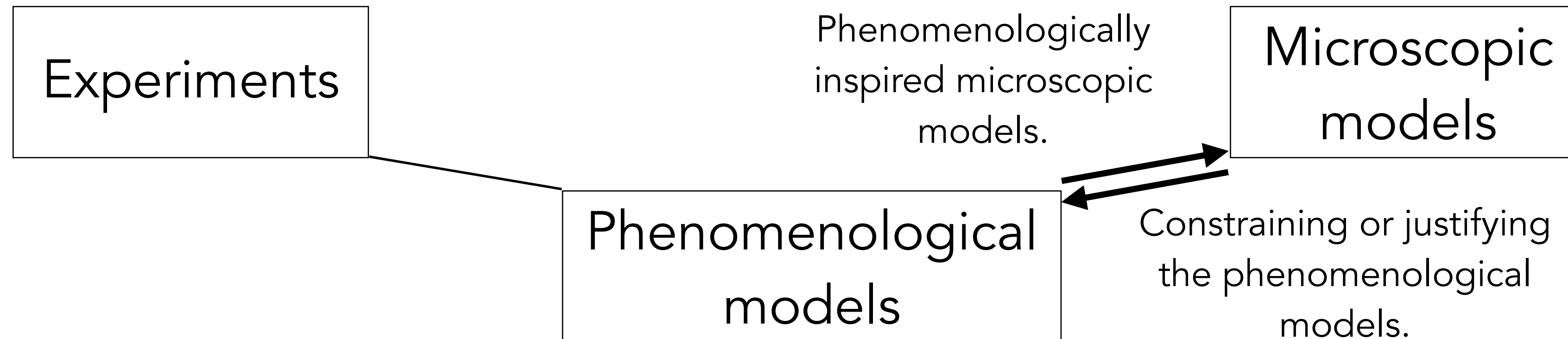
Precision era for decoherence in neutrino oscillations



Understanding the underlying physics is important to interpret the data.

Conclusions

Precision era for decoherence in neutrino oscillations

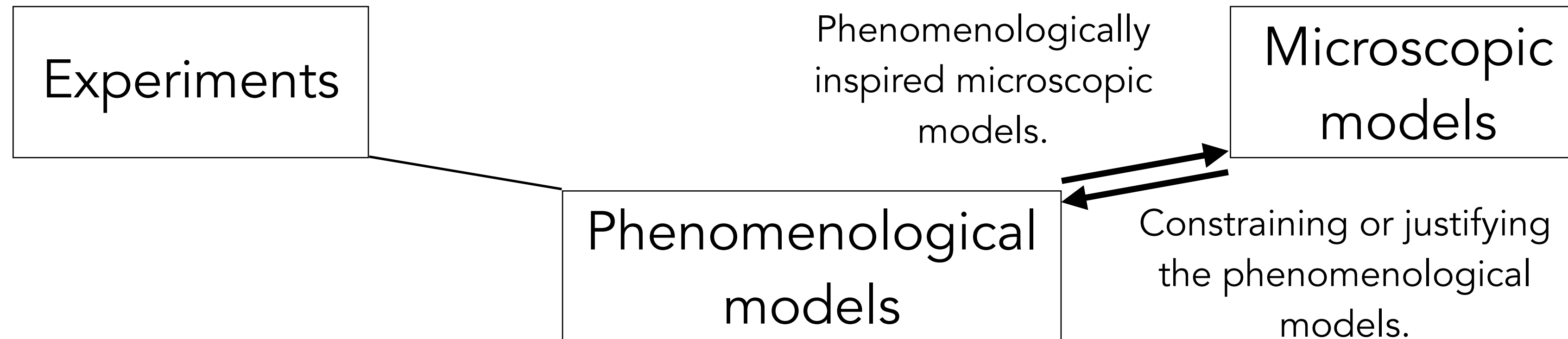


Understanding the underlying physics is important to interpret the data.

Study of the **environment-system interaction** and mapping to pheno models $\longrightarrow$
energy dependence
shape of the dissipator

Conclusions

Precision era for decoherence in neutrino oscillations



Understanding the underlying physics is important to interpret the data.

Study of the **environment-system interaction** and mapping to pheno models $\longrightarrow$
energy dependence
shape of the dissipator

Thank you!