

Modified dispersion relations and relativistic gas dynamics

Class. Quant. Grav. **41** (2023) 015025 [arXiv:2310.01487 [gr-qc]]

Manuel Hohmann

Laboratory of Theoretical Physics, Institute of Physics, University of Tartu
Center of Excellence “Fundamental Universe”



BridgeQG first annual conference - 10. July 2025

Idea and motivation

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard ⚡.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard ⚡.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems ⚡.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\nexists$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\nexists$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\nexists$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\nexists$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\nexists$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\nexists$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard ⚡.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems ⚡.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible (✓).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard ⚡.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems ⚡.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible (✓).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work ⚡.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard ⚡.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems ⚡.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible (✓).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of **quantum corrections** to GR.
 - Study the phenomenology of quantum corrections.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\nexists$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\nexists$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\nexists$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.
 - Study the phenomenology of quantum corrections.
- How can we study quantum gravity phenomenology?
 - Find physical system which could amplify deviations from general relativity.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\nexists$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\nexists$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\nexists$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.
 - Study the phenomenology of quantum corrections.
- How can we study quantum gravity phenomenology?
 - Find physical system which could amplify deviations from general relativity.
 - Example: study compact system with very strong gravity.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\nexists$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\nexists$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\nexists$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.
 - Study the phenomenology of quantum corrections.
- How can we study quantum gravity phenomenology?
 - Find physical system which could amplify deviations from general relativity.
 - Example: study compact system with very strong gravity.
 - Think of possible observables in the chosen system.

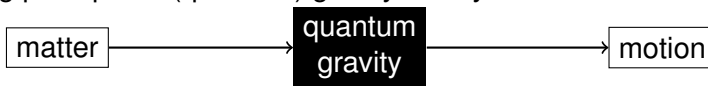
- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\downarrow$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\downarrow$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\downarrow$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.
 - Study the phenomenology of quantum corrections.
- How can we study quantum gravity phenomenology?
 - Find physical system which could amplify deviations from general relativity.
 - Example: study compact system with very strong gravity.
 - Think of possible observables in the chosen system.
 - Calculate how effective quantum gravity influences observables.

- How can we quantize gravity?
 - Use same methods as in QFT $\rightsquigarrow$ doesn't work $\downarrow$.
 - Guess a complete theory of quantum gravity $\rightsquigarrow$ hard $\downarrow$.
 - Assume gravity is classical $\rightsquigarrow$ leaves unsolved problems $\downarrow$.
 - Study effective quantum gravity model phenomenology $\rightsquigarrow$ maybe feasible ($\checkmark$).
- How can we study effective gravity models?
 - Don't care (too much) about fundamental laws of gravity.
 - Assume that general relativity (GR) is almost correct.
 - Think of possible sources of quantum corrections to GR.
 - Study the phenomenology of quantum corrections.
- How can we study quantum gravity phenomenology?
 - Find physical system which could amplify deviations from general relativity.
 - Example: study compact system with very strong gravity.
 - Think of possible observables in the chosen system.
 - Calculate how effective quantum gravity influences observables.

$\Rightarrow$ Here: effective quantum gravity phenomenology with gas dynamics near black holes.

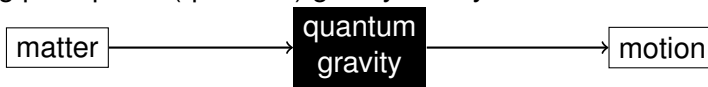
How does quantum gravity phenomenology work?

- Basic operating principle of (quantum) gravity theory:



How does quantum gravity phenomenology work?

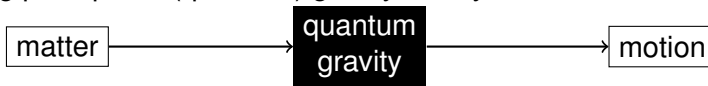
- Basic operating principle of (quantum) gravity theory:



⚡ Quantum gravity is a black box!

How does quantum gravity phenomenology work?

- Basic operating principle of (quantum) gravity theory:

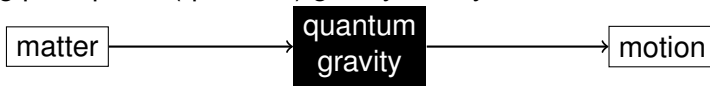


- ⚡ Quantum gravity is a black box!
- Quantum gravity must approximate general relativity:

$$\text{quantum gravity} = \text{general relativity} + \epsilon \cdot \text{quantum correction}$$

How does quantum gravity phenomenology work?

- Basic operating principle of (quantum) gravity theory:



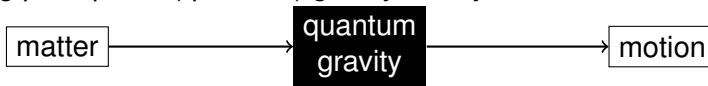
- ⚡ Quantum gravity is a black box!
- Quantum gravity must approximate general relativity:

$$\text{quantum gravity} = \text{general relativity} + \epsilon \cdot \text{quantum correction}$$

↪ We still have a black box, but it is multiplied by $\epsilon \ll 1 \rightsquigarrow$ perturbation.

How does quantum gravity phenomenology work?

- Basic operating principle of (quantum) gravity theory:



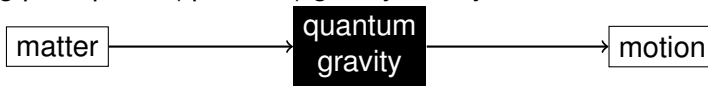
- ⚡ Quantum gravity is a black box!
- Quantum gravity must approximate general relativity:

$$\text{quantum gravity} = \text{general relativity} + \epsilon \cdot \text{quantum correction}$$

- ↪ We still have a black box, but it is multiplied by $\epsilon \ll 1 \rightsquigarrow$ perturbation.
- ✓ General relativity is a very simple theory!

How does quantum gravity phenomenology work?

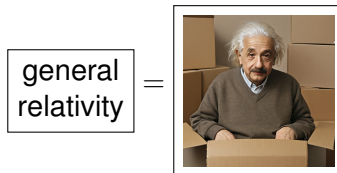
- Basic operating principle of (quantum) gravity theory:



- ⚡ Quantum gravity is a black box!
- Quantum gravity must approximate general relativity:

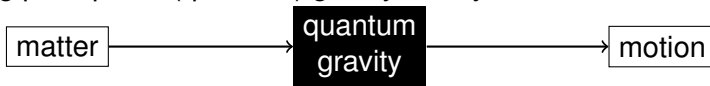
$$\text{quantum gravity} = \text{general relativity} + \epsilon \cdot \text{quantum correction}$$

- ⇒ We still have a black box, but it is multiplied by $\epsilon \ll 1 \rightsquigarrow$ perturbation.
- ✓ General relativity is a very simple theory!
- ⇒ We know what is in the white box:



How does quantum gravity phenomenology work?

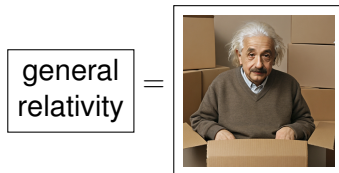
- Basic operating principle of (quantum) gravity theory:



- ⚡ Quantum gravity is a black box!
- Quantum gravity must approximate general relativity:

$$\text{quantum gravity} = \text{general relativity} + \epsilon \cdot \text{quantum correction}$$

- ↪ We still have a black box, but it is multiplied by $\epsilon \ll 1 \rightsquigarrow$ perturbation.
- ✓ General relativity is a very simple theory!
- ⇒ We know what is in the white box:



- ↪ Only need to study (all) possible quantum corrections!

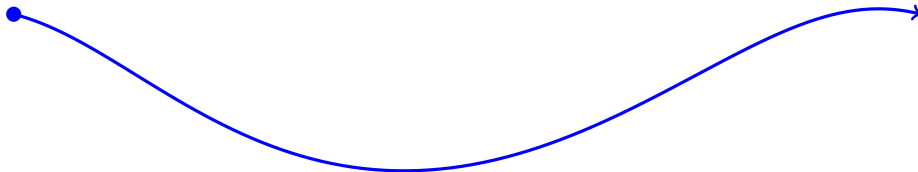
Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.



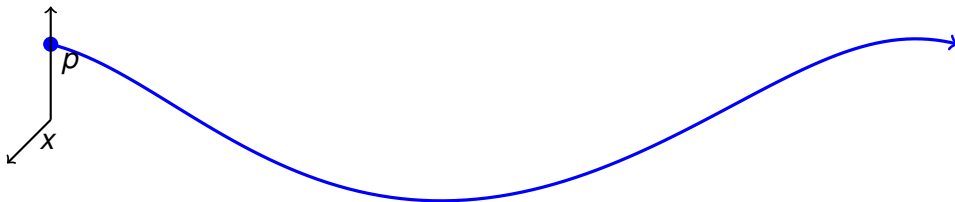
Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.
- Particle trajectories follow (relativistic) Hamiltonian dynamics.



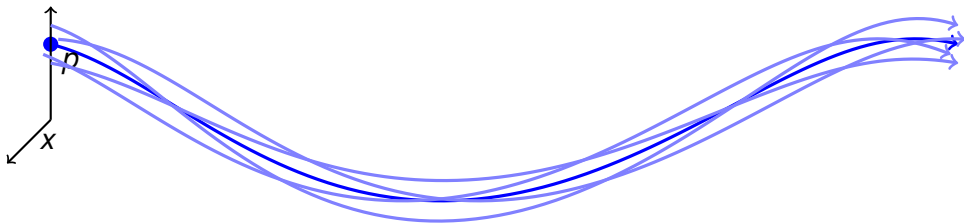
Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.
- Particle trajectories follow (relativistic) Hamiltonian dynamics.
- Each particle is described by (spacetime) position and four-momentum.



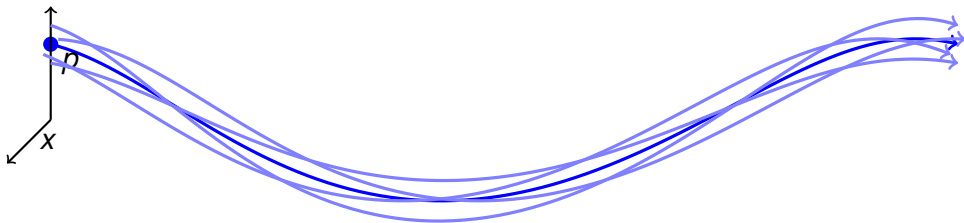
Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.
 - Particle trajectories follow (relativistic) Hamiltonian dynamics.
 - Each particle is described by (spacetime) position and four-momentum.
- ⇒ Kinetic gas is density distribution in 8-dimensional position-momentum phase space.



Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.
 - Particle trajectories follow (relativistic) Hamiltonian dynamics.
 - Each particle is described by (spacetime) position and four-momentum.
- ⇒ Kinetic gas is density distribution in 8-dimensional position-momentum phase space.
- ⇒ Gas dynamics follows from Hamiltonian particle dynamics.

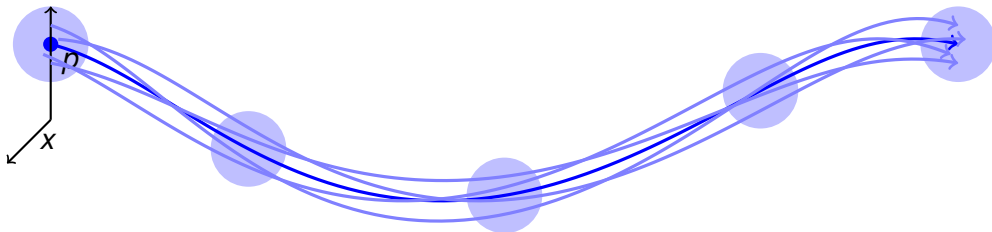


Kinetic gas to probe quantum gravity

- Gas is constituted by particles of equal mass.
 - Particle trajectories follow (relativistic) Hamiltonian dynamics.
 - Each particle is described by (spacetime) position and four-momentum.
- ⇒ Kinetic gas is density distribution in 8-dimensional position-momentum phase space.
- ⇒ Gas dynamics follows from Hamiltonian particle dynamics.

Collisionless gas

Particle density function is constant along particle trajectories in phase space.



Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.

Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.
- (Modified) dispersion relation: mass shell condition for Hamiltonian:

$$-\frac{m^2}{2} = H(x^\mu, \bar{x}_\mu). \quad (1)$$

Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.
- (Modified) dispersion relation: mass shell condition for Hamiltonian:

$$-\frac{m^2}{2} = H(x^\mu, \bar{x}_\mu). \quad (1)$$

- Particle trajectories derived from Hamilton's equations of motion:

$$\dot{x}^\mu = \bar{\partial}^\mu H, \quad \dot{\bar{x}}^\mu = -\partial_\mu H. \quad (2)$$

Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.
- (Modified) dispersion relation: mass shell condition for Hamiltonian:

$$-\frac{m^2}{2} = H(x^\mu, \bar{x}_\mu). \quad (1)$$

- Particle trajectories derived from Hamilton's equations of motion:

$$\dot{x}^\mu = \bar{\partial}^\mu H, \quad \dot{\bar{x}}^\mu = -\partial_\mu H. \quad (2)$$

- Canonical cotangent bundle geometry: symplectic form $\omega \in \Omega^2(T^*M)$ as

$$\theta = \bar{x}_\mu dx^\mu, \quad \omega = d\theta = d\bar{x}_\mu \wedge dx^\mu. \quad (3)$$

Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.
- (Modified) dispersion relation: mass shell condition for Hamiltonian:

$$-\frac{m^2}{2} = H(x^\mu, \bar{x}_\mu). \quad (1)$$

- Particle trajectories derived from Hamilton's equations of motion:

$$\dot{x}^\mu = \bar{\partial}^\mu H, \quad \dot{\bar{x}}^\mu = -\partial_\mu H. \quad (2)$$

- Canonical cotangent bundle geometry: symplectic form $\omega \in \Omega^2(T^*M)$ as

$$\theta = \bar{x}_\mu dx^\mu, \quad \omega = d\theta = d\bar{x}_\mu \wedge dx^\mu. \quad (3)$$

- Hamiltonian vector field X_H on T^*M : unique solution of

$$\iota_{X_H}\omega = -dH. \quad (4)$$

Hamiltonian formulation of relativistic particle dynamics

- Model particle dynamics on cotangent bundle T^*M with coordinates $(x^\mu, \bar{x}_\mu)$.
- (Modified) dispersion relation: mass shell condition for Hamiltonian:

$$-\frac{m^2}{2} = H(x^\mu, \bar{x}_\mu). \quad (1)$$

- Particle trajectories derived from Hamilton's equations of motion:

$$\dot{x}^\mu = \bar{\partial}^\mu H, \quad \dot{\bar{x}}^\mu = -\partial_\mu H. \quad (2)$$

- Canonical cotangent bundle geometry: symplectic form $\omega \in \Omega^2(T^*M)$ as

$$\theta = \bar{x}_\mu dx^\mu, \quad \omega = d\theta = d\bar{x}_\mu \wedge dx^\mu. \quad (3)$$

- Hamiltonian vector field X_H on T^*M : unique solution of

$$\iota_{X_H} \omega = -dH. \quad (4)$$

⇒ Particle trajectories are integral curves of X_H .

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3 . \quad (5)$$

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

- σ : hypersurface in T^*M which is transverse to X_H .

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

- σ : hypersurface in T^*M which is transverse to X_H .
- $N[\sigma]$: number of particle trajectories through σ .

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

- σ : hypersurface in T^*M which is transverse to X_H .
- $N[\sigma]$: number of particle trajectories through σ .
- $\Omega = \iota_{X_H} \Sigma$: particle measure.

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

- σ : hypersurface in T^*M which is transverse to X_H .
 - $N[\sigma]$: number of particle trajectories through σ .
 - $\Omega = \iota_{X_H} \Sigma$: particle measure.
- Collisionless gas: particles follow Hamilton's equations of motion, no interactions.

From particle dynamics to gas dynamics

- Introduce symplectic volume form:

$$\Sigma = \frac{1}{4!} \omega \wedge \omega \wedge \omega \wedge \omega = dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3 \wedge d\bar{x}_0 \wedge d\bar{x}_1 \wedge d\bar{x}_2 \wedge d\bar{x}_3. \quad (5)$$

- Kinetic gas: particles of equal mass described by phase space trajectories on T^*M .
- One particle distribution function $\phi : T^*M \rightarrow \mathbb{R}^+$:

$$N[\sigma] = \int_{\sigma} \phi \Omega. \quad (6)$$

- σ : hypersurface in T^*M which is transverse to X_H .
 - $N[\sigma]$: number of particle trajectories through σ .
 - $\Omega = \iota_{X_H} \Sigma$: particle measure.
 - Collisionless gas: particles follow Hamilton's equations of motion, no interactions.
- $\Rightarrow$ 1-PDF follows Liouville equation: $\mathcal{L}_{X_H} \phi = 0$.

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.
- ⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu . \quad (7)$$

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:

- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
- Hamiltonian and 1-PDF depend only on r, P, E, L .

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

- ⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:
- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
 - Hamiltonian and 1-PDF depend only on r, P, E, L .
 - Also Hamiltonian is constant of motion: $X_H H = 0$.

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

- ⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:
- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
 - Hamiltonian and 1-PDF depend only on r, P, E, L .
- Also Hamiltonian is constant of motion: $X_H H = 0$.
- ⇒ Replace P by H in new coordinates $(t, r, \Theta, \Phi, \Psi, E, H, L)$ such that:
- Θ, Φ, E, L, H are constant along trajectories.
 - 1-PDF depends only on r, E, L, H .

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:

- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
- Hamiltonian and 1-PDF depend only on r, P, E, L .

- Also Hamiltonian is constant of motion: $X_H H = 0$.

⇒ Replace P by H in new coordinates $(t, r, \Theta, \Phi, \Psi, E, H, L)$ such that:

- Θ, Φ, E, L, H are constant along trajectories.
- 1-PDF depends only on r, E, L, H .

⇒ Liouville equation becomes $\partial_r \phi = 0$.

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:

- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
- Hamiltonian and 1-PDF depend only on r, P, E, L .

- Also Hamiltonian is constant of motion: $X_H H = 0$.

⇒ Replace P by H in new coordinates $(t, r, \Theta, \Phi, \Psi, E, H, L)$ such that:

- Θ, Φ, E, L, H are constant along trajectories.
- 1-PDF depends only on r, E, L, H .

⇒ Liouville equation becomes $\partial_r \phi = 0$.

⇒ Most general solution to static spherically symmetric gas: $\phi = \phi(E, L, H)$.

Static spherically symmetric system

- Symmetry generated by vector fields $(X_I) = (X_0, X_1, X_2, X_3)$.

⇒ Hamiltonian and 1-PDF invariant under complete lift: $\mathcal{L}_{\hat{X}_I} H = \mathcal{L}_{\hat{X}_I} \phi = 0$ with

$$\hat{X} = X^\mu \partial_\mu - \bar{x}_\nu \partial_\mu X^\nu \bar{\partial}^\mu. \quad (7)$$

⇒ Introduce new coordinates $(t, r, \Theta, \Phi, \Psi, E, P, L)$ such that:

- Θ, Φ, E, L are constant along trajectories (Noether symmetries).
- Hamiltonian and 1-PDF depend only on r, P, E, L .

- Also Hamiltonian is constant of motion: $X_H H = 0$.

⇒ Replace P by H in new coordinates $(t, r, \Theta, \Phi, \Psi, E, H, L)$ such that:

- Θ, Φ, E, L, H are constant along trajectories.
- 1-PDF depends only on r, E, L, H .

⇒ Liouville equation becomes $\partial_r \phi = 0$.

⇒ Most general solution to static spherically symmetric gas: $\phi = \phi(E, L, H)$.

⇒ Consider gas $\phi \sim \delta(E)\delta(L)\delta(H)$ of identical energy, angular momentum, mass.

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} Z^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell Z^\mu \bar{x}_\mu} (g^{\mu\nu} + Z^\mu Z^\nu) \bar{x}_\mu \bar{x}_\nu . \quad (8)$$

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} Z^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell Z^\mu \bar{x}_\mu} (\textcolor{red}{g}^{\mu\nu} + Z^\mu Z^\nu) \bar{x}_\mu \bar{x}_\nu . \quad (8)$$

- Spacetime metric $\textcolor{red}{g}_{\mu\nu}$.

κ -Poincaré correction of Schwarzschild spacetime

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} \mathbf{Z}^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell \mathbf{Z}^\mu \bar{x}_\mu} (g^{\mu\nu} + \mathbf{Z}^\mu \mathbf{Z}^\nu) \bar{x}_\mu \bar{x}_\nu . \quad (8)$$

- Spacetime metric $g_{\mu\nu}$.
- Unit timelike vector field $\mathbf{Z}^\mu$ satisfying $\mathbf{Z}^\mu \mathbf{Z}^\nu g_{\mu\nu} = -1$.

κ -Poincaré correction of Schwarzschild spacetime

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} Z^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell Z^\mu \bar{x}_\mu} (g^{\mu\nu} + Z^\mu Z^\nu) \bar{x}_\mu \bar{x}_\nu. \quad (8)$$

- Spacetime metric $g_{\mu\nu}$.
- Unit timelike vector field Z^μ satisfying $Z^\mu Z^\nu g_{\mu\nu} = -1$.
- Planck length ℓ .

κ -Poincaré correction of Schwarzschild spacetime

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} Z^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell Z^\mu \bar{x}_\mu} (g^{\mu\nu} + Z^\mu Z^\nu) \bar{x}_\mu \bar{x}_\nu. \quad (8)$$

- Spacetime metric $g_{\mu\nu}$.
- Unit timelike vector field Z^μ satisfying $Z^\mu Z^\nu g_{\mu\nu} = -1$.
- Planck length ℓ .

⇒ Static spherically symmetric case defined by functions a, b, c, d of r :

$$H = -\frac{2}{\ell^2} \sinh^2 \left[\frac{\ell}{2} (-cE + dP) \right] + \frac{1}{2} e^{\ell(-cE+dP)} \left[(-a + c^2)E^2 - 2cdEP + (b + d^2)P^2 + \frac{L^2}{r^2} \right]. \quad (9)$$

κ -Poincaré correction of Schwarzschild spacetime

- General κ -Poincaré modification of metric dispersion relation:

$$H = -\frac{2}{\ell^2} \sinh^2 \left(\frac{\ell}{2} Z^\mu \bar{x}_\mu \right) + \frac{1}{2} e^{\ell Z^\mu \bar{x}_\mu} (g^{\mu\nu} + Z^\mu Z^\nu) \bar{x}_\mu \bar{x}_\nu. \quad (8)$$

- Spacetime metric $g_{\mu\nu}$.
- Unit timelike vector field Z^μ satisfying $Z^\mu Z^\nu g_{\mu\nu} = -1$.
- Planck length ℓ .

⇒ Static spherically symmetric case defined by functions a, b, c, d of r :

$$H = -\frac{2}{\ell^2} \sinh^2 \left[\frac{\ell}{2} (-cE + dP) \right] + \frac{1}{2} e^{\ell(-cE+dP)} \left[(-a + c^2)E^2 - 2cdEP + (b + d^2)P^2 + \frac{L^2}{r^2} \right]. \quad (9)$$

⇒ Minimal modification of Schwarzschild spacetime of mass M :

$$a^{-1} = b = c^{-2} = 1 - \frac{2M}{r}, \quad d = 0. \quad (10)$$

Example: swarm of orbiting particles

- Properties of particle ensemble:
 - Identical angular momentum $L > 0$ (motion has angular component).
 - Energy E such that particles are gravitationally bound.

Example: swarm of orbiting particles

- Properties of particle ensemble:
 - Identical angular momentum $L > 0$ (motion has angular component).
 - Energy E such that particles are gravitationally bound.
- ⇒ Orbit oscillates between two radii $R_{1,2}$.

Example: swarm of orbiting particles

- Properties of particle ensemble:

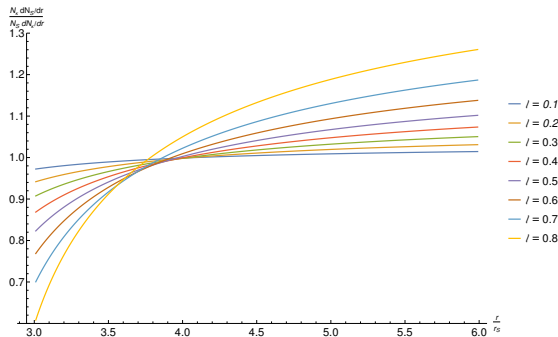
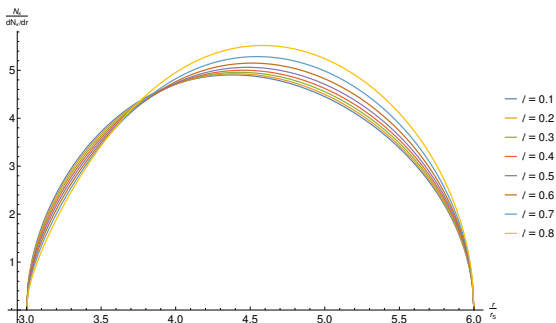
- Identical angular momentum $L > 0$ (motion has angular component).
- Energy E such that particles are gravitationally bound.

⇒ Orbit oscillates between two radii $R_{1,2}$.

- Calculate number of trajectories through time slice σ with $R < r < R + dr$.

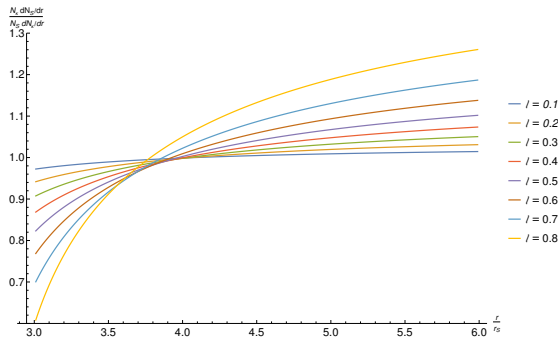
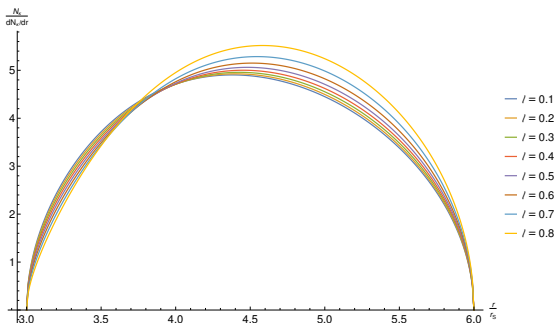
Example: swarm of orbiting particles

- Properties of particle ensemble:
 - Identical angular momentum $L > 0$ (motion has angular component).
 - Energy E such that particles are gravitationally bound.
- ⇒ Orbit oscillates between two radii $R_{1,2}$.
- Calculate number of trajectories through time slice σ with $R < r < R + dr$.
- Plot (inverse of) relative particle density $N/(dN/dr)$:



Example: swarm of orbiting particles

- Properties of particle ensemble:
 - Identical angular momentum $L > 0$ (motion has angular component).
 - Energy E such that particles are gravitationally bound.
- ⇒ Orbit oscillates between two radii $R_{1,2}$.
- Calculate number of trajectories through time slice σ with $R < r < R + dr$.
- Plot (inverse of) relative particle density $N/(dN/dr)$:



⇒ κ -Poincaré modification shifts particles inward.

Example: radial free fall from infinity

- Properties of particle ensemble:
 - Identical angular momentum $L = 0$ (purely radial motion).
 - Energy E such that particles are marginally bound (drop from rest at $r = \infty$).

Example: radial free fall from infinity

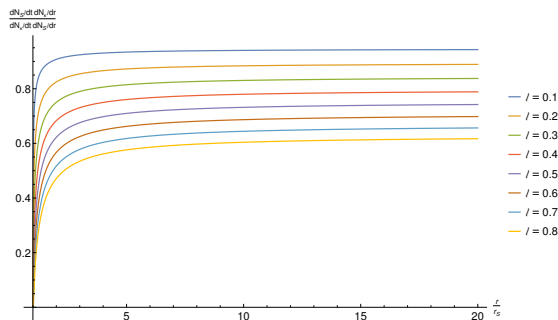
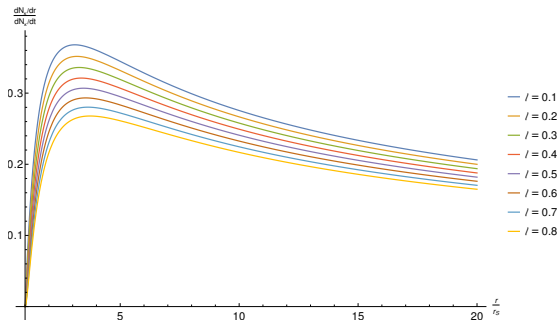
- Properties of particle ensemble:
 - Identical angular momentum $L = 0$ (purely radial motion).
 - Energy E such that particles are marginally bound (drop from rest at $r = \infty$).
- Assume constant flow rate through radial slice.

Example: radial free fall from infinity

- Properties of particle ensemble:
 - Identical angular momentum $L = 0$ (purely radial motion).
 - Energy E such that particles are marginally bound (drop from rest at $r = \infty$).
- Assume constant flow rate through radial slice.
- Calculate number of trajectories through time slice σ with $R < r < R + dr$.

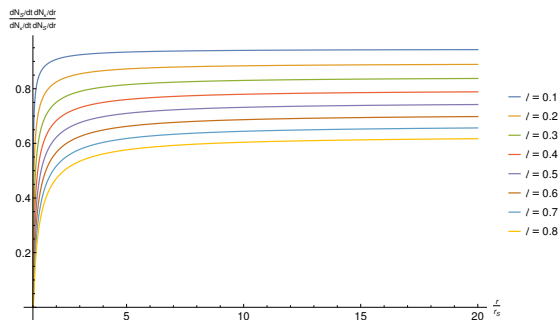
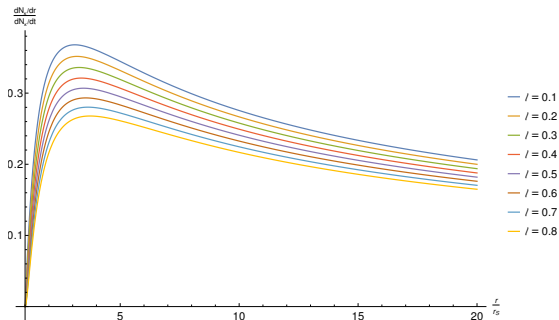
Example: radial free fall from infinity

- Properties of particle ensemble:
 - Identical angular momentum $L = 0$ (purely radial motion).
 - Energy E such that particles are marginally bound (drop from rest at $r = \infty$).
- Assume constant flow rate through radial slice.
- Calculate number of trajectories through time slice σ with $R < r < R + dr$.
- Plot particle density dN/dr per flow rate dN/dt :



Example: radial free fall from infinity

- Properties of particle ensemble:
 - Identical angular momentum $L = 0$ (purely radial motion).
 - Energy E such that particles are marginally bound (drop from rest at $r = \infty$).
- Assume constant flow rate through radial slice.
- Calculate number of trajectories through time slice σ with $R < r < R + dr$.
- Plot particle density dN/dr per flow rate dN/dt :



⇒ κ -Poincaré modification decreases particle density.

- Summary:

- ⚡ Fundamental theory of quantum gravity is unknown.
- ⇒ Consider effective quantum gravity models instead.
 - Effective model is small correction to general relativity.
- ⇒ Study observable effects of possible quantum corrections.
 - κ -Poincaré modification changes matter density near black hole.

- Summary:

- ⚡ Fundamental theory of quantum gravity is unknown.
- ⇒ Consider effective quantum gravity models instead.
 - Effective model is small correction to general relativity.
- ⇒ Study observable effects of possible quantum corrections.
 - κ -Poincaré modification changes matter density near black hole.

- Outlook:

- Consider more general quantum corrections.
- Consider spinning black holes.
- Consider more general gases or matter distributions with less symmetry:
 - Accretion disks and jets $\rightsquigarrow$ blazars.
 - Tidal disruption events.
 - Stellar wake of passing black hole and dynamical friction.
- Derive observable properties of black holes, quasars, AGN. . .

- Summary:

- ⚡ Fundamental theory of quantum gravity is unknown.
- ⇒ Consider effective quantum gravity models instead.
 - Effective model is small correction to general relativity.
- ⇒ Study observable effects of possible quantum corrections.
 - κ -Poincaré modification changes matter density near black hole.

- Outlook:

- Consider more general quantum corrections.
- Consider spinning black holes.
- Consider more general gases or matter distributions with less symmetry:
 - Accretion disks and jets $\rightsquigarrow$ blazars.
 - Tidal disruption events.
 - Stellar wake of passing black hole and dynamical friction.
- Derive observable properties of black holes, quasars, AGN. . .
- MH, “Kinetic gases in static spherically symmetric modified dispersion relations,” *Class. Quant. Grav.* **41** (2024) no.1, 015025 [arXiv:2310.01487 [gr-qc]].