The regularization of spacetime singularities

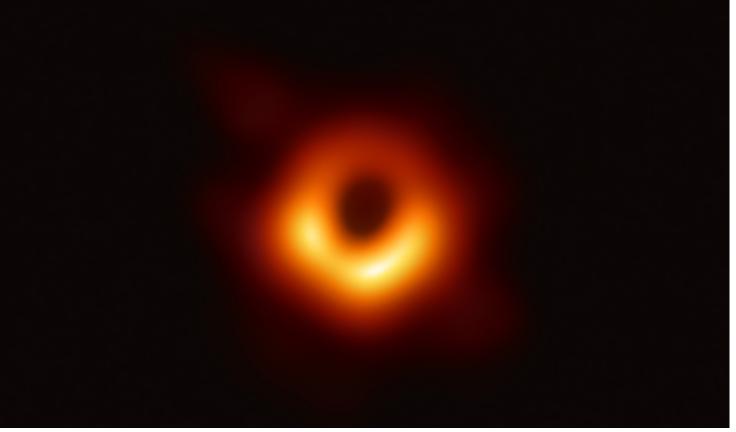
Vania Vellucci BridgeQG 2025, Paris

SDU - hotc

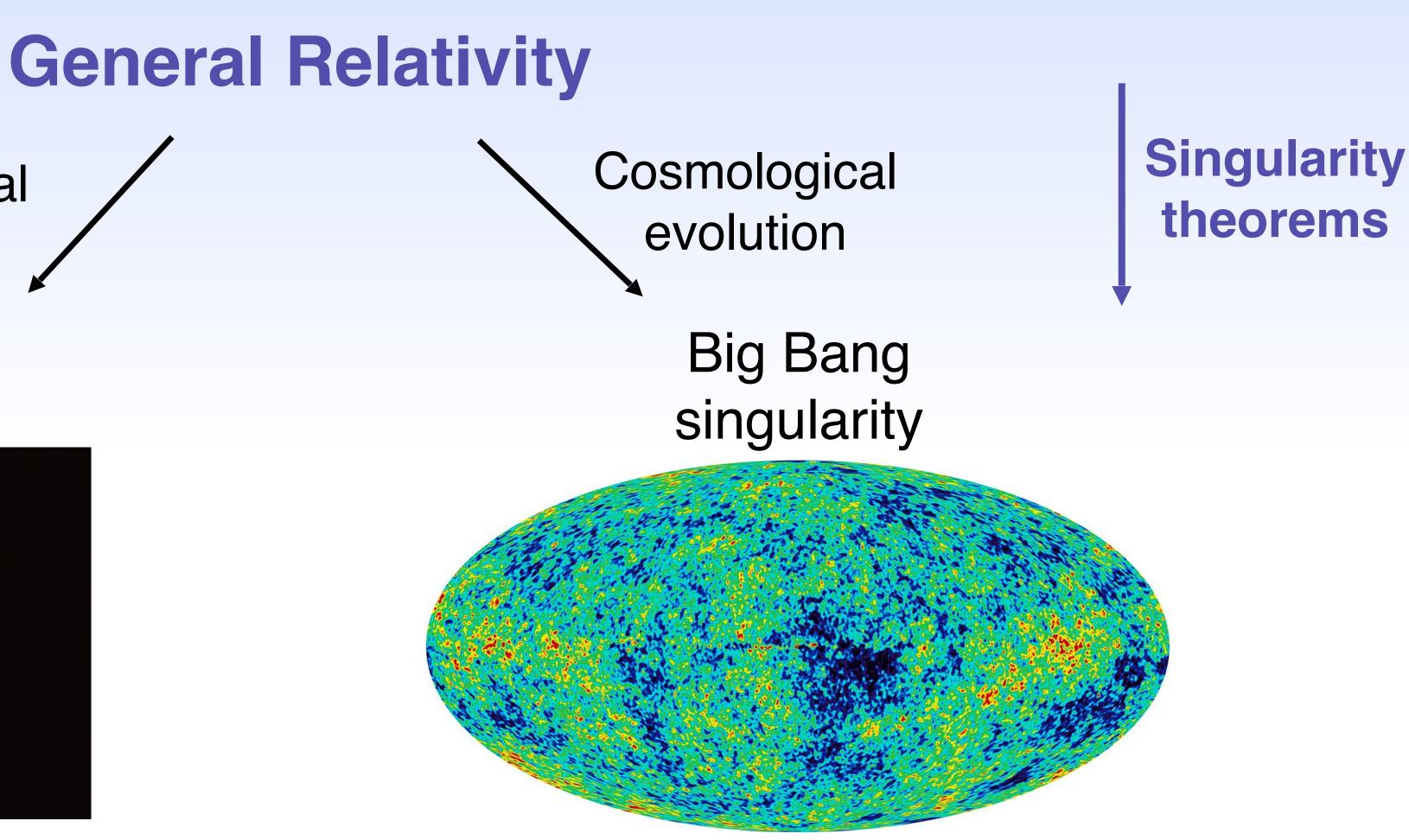




Black Hole singularities



A singular space-time is **geodesic incomplete**: there exist at least one geodesic that cannot be extended beyond a finite proper time or affine parameter (particles seem to disappear from existence!)









In a complete theory of quantum gravity we expect the formation of spacetime singularities to be prevented



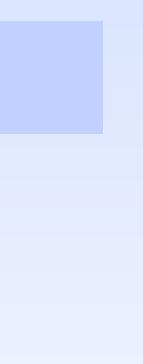
Horizon X Singularity 🗙

There are few possible alternatives to singular black holes to describe the ultra-compact objects that we see in the sky

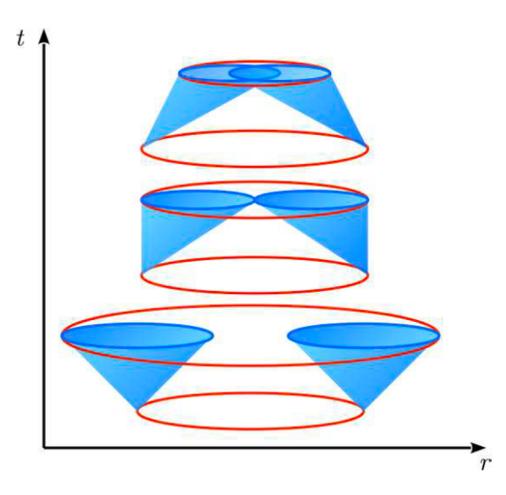
Regular Black holes

Horizon \checkmark Singularity 🗙 BHs with integrable singularites

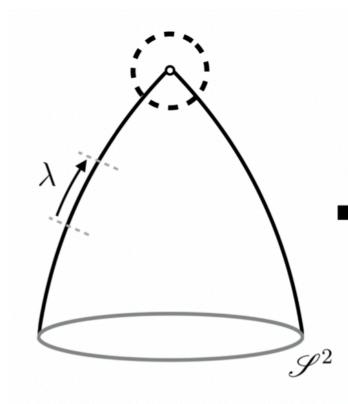
> Horizon \checkmark Singularity ~



Regular Black holes

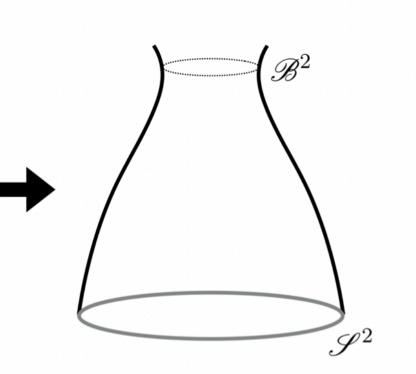


$\theta_+ < 0$ in some points GR + NEC +non-compact Cauchy surface



Singular focusing point! $\theta_+ \rightarrow - \infty$

When both the expansion of null ingoing θ_{-} and outgoing θ_{+} geodesics become negative, a trapped surface (horizon) is formed

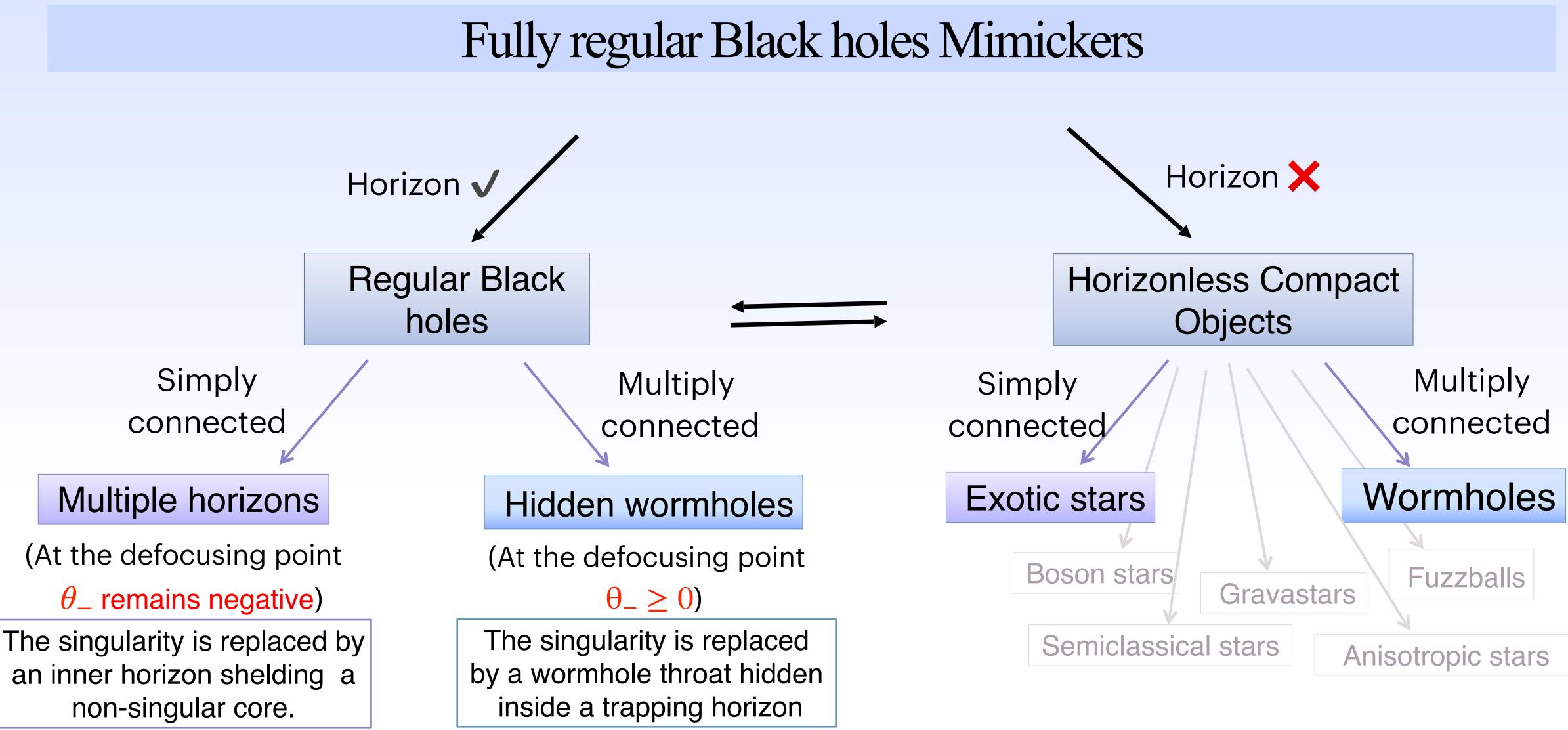


 $\theta_{+} < 0$ in some points

GR + NEC + non-compact Cauchy surface

Defocusing point at which the expansion changes sign again $\theta_{+} = 0$

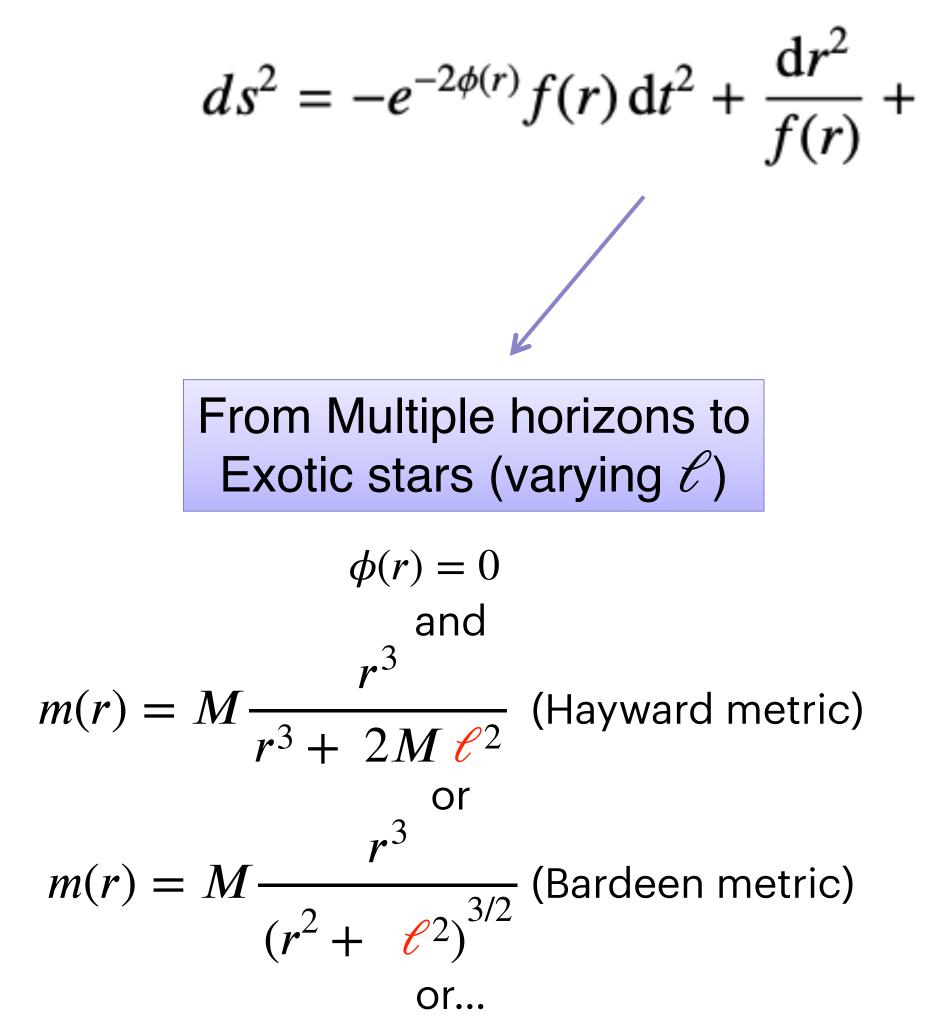




Carballo Rubio et al.,2019

*we are considering spherical symmetric objects





Carballo-Rubio et al. 2023

$$r^2 \left(\mathrm{d}\theta^2 + \sin^2\theta \,\mathrm{d}\varphi^2 \right), \quad f(r) = 1 - \frac{2m(r)}{r}.$$

From Hidden wormholes to traversable Wormholes (varying \mathscr{C})

$$\phi(r) = \frac{1}{2} \log \left(1 - \frac{\ell^2}{r^2} \right)$$

and
$$m(r) = M \left(1 - \frac{\ell^2}{r^2} \right) + \frac{\ell^2}{2r} \text{ (Simpson-Visser)}$$



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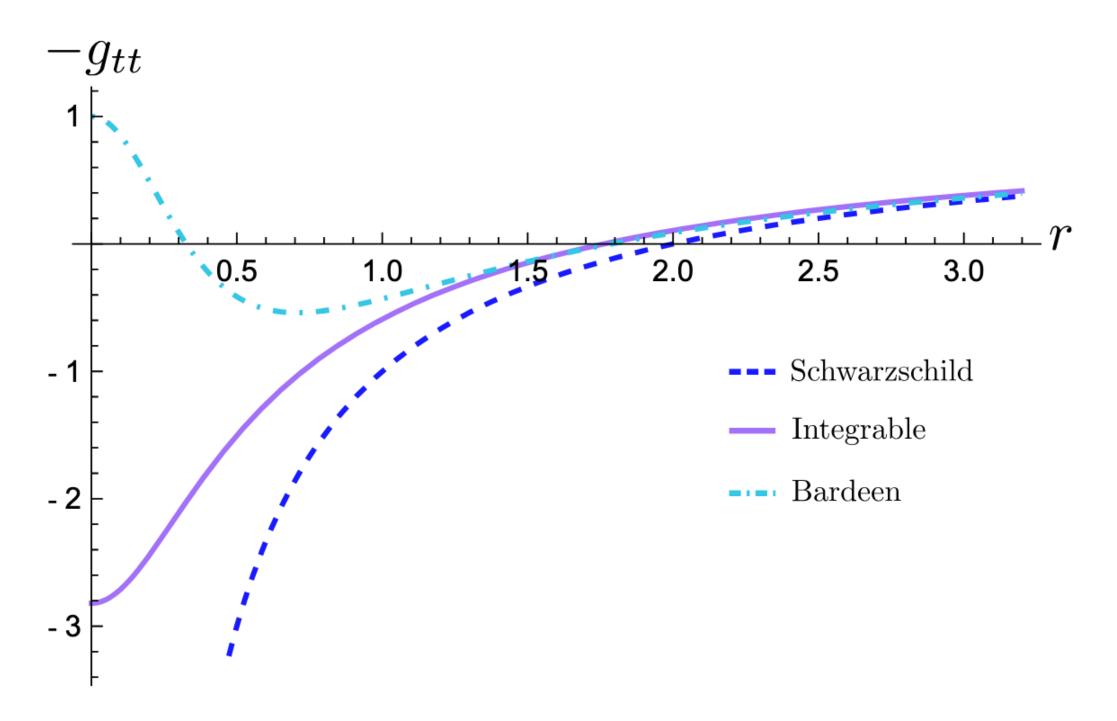
Black holes with integrable singularities

Lukash+Strokov 2011, Casadio 2024, Bonanno+Koch+Platania 2017...

Near the center

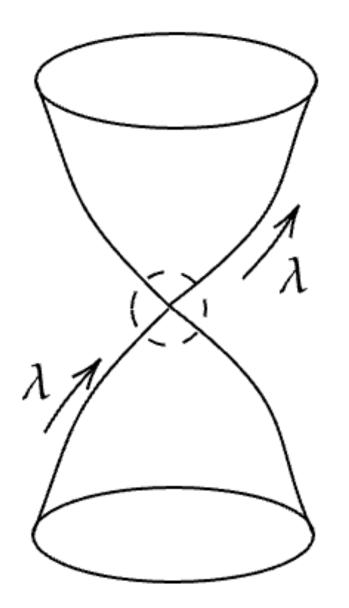
 $m(r) \sim r$

the metric is regular but the curvature scalars diverges at r = 0



No need of an inner Cauchy horizon!

The focusing point is still there but point particles feel finite tidal forces, so you can "continue" their path after the focusing point



However when you extend the spacetime beyond the focusing point, global hyperbolicity is still lost: it is a sort of "Cauchy point"

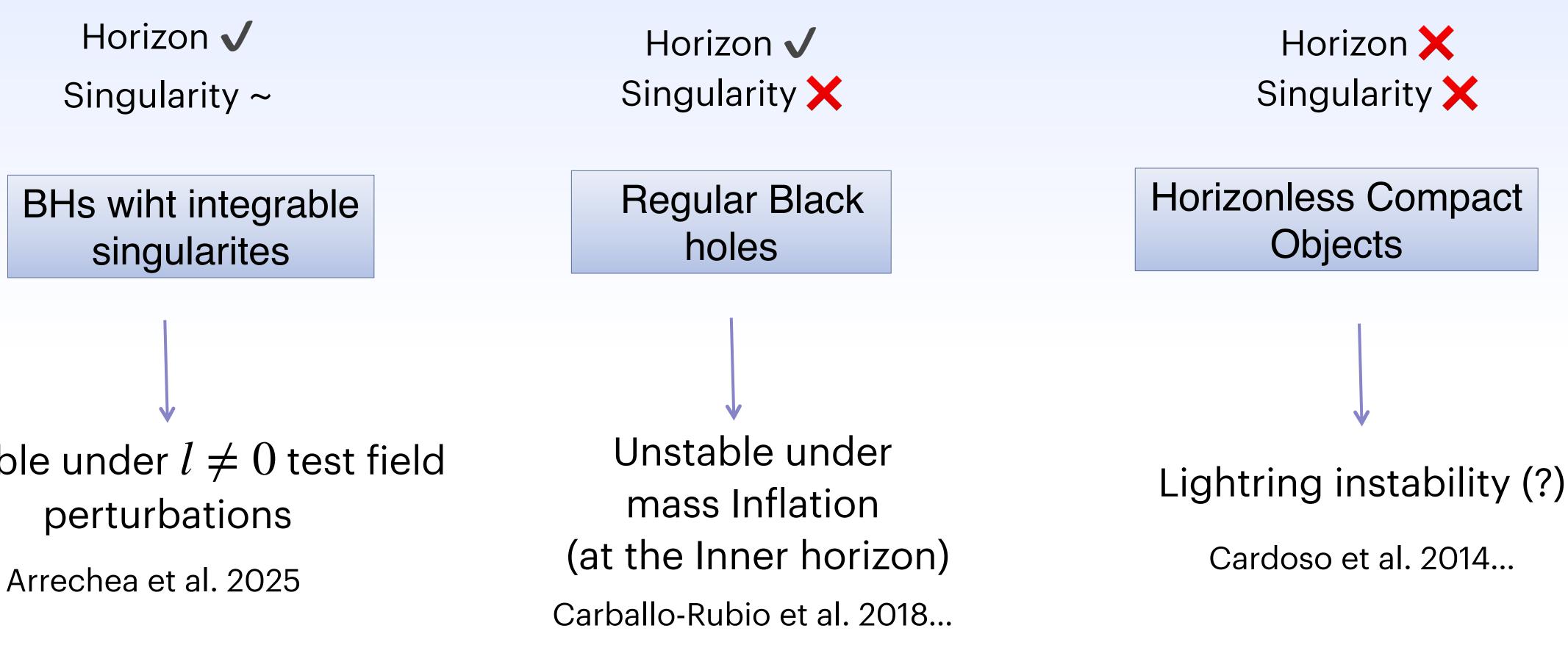
Arrechea et al. 2025





Disclaimer: instabilities

Within classical dynamics (GR) all these alternatives are unstable



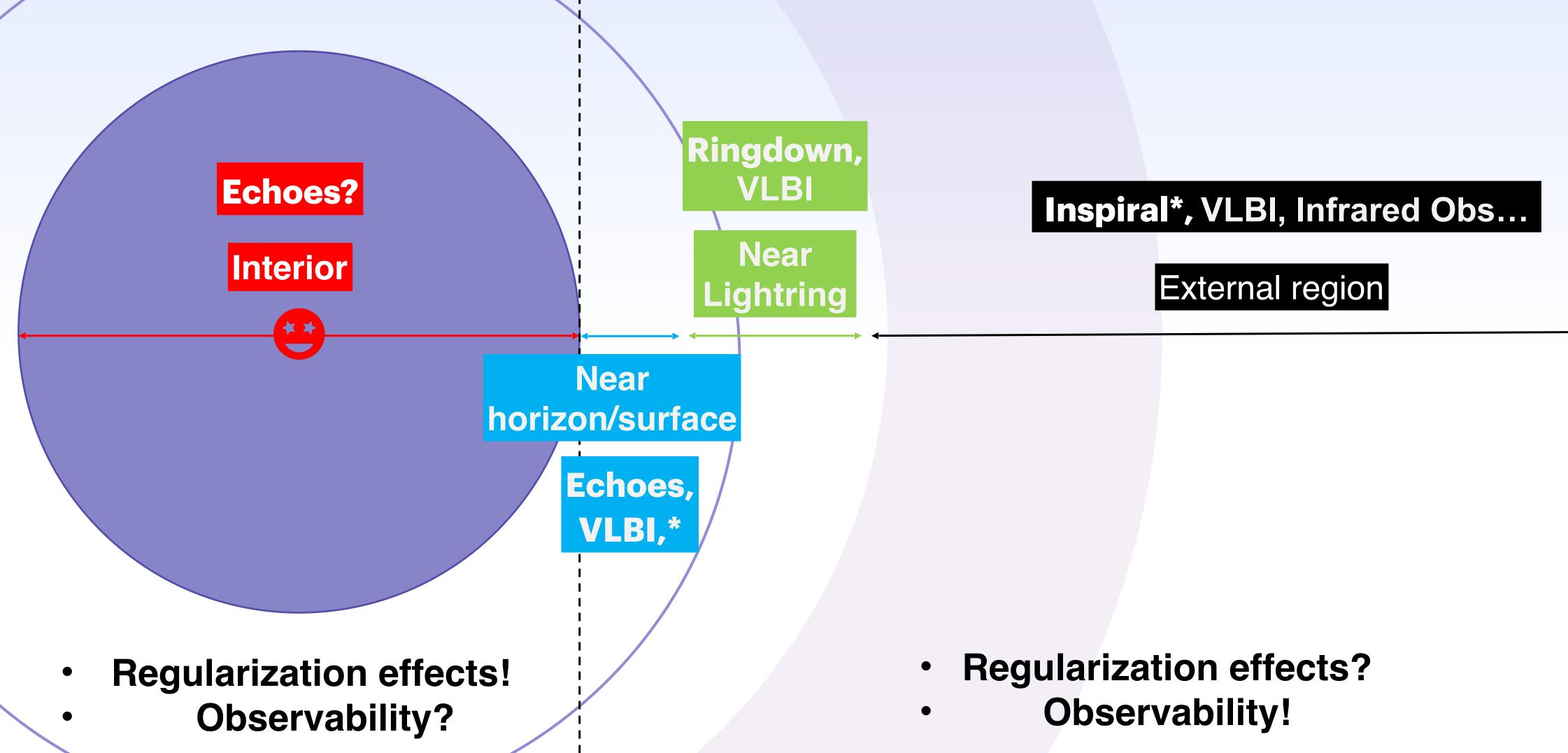
Unstable under $l \neq 0$ test field

(Quantum) corrected dynamics should prevent the collapse of these objects in singular BHs, taming these instabilities



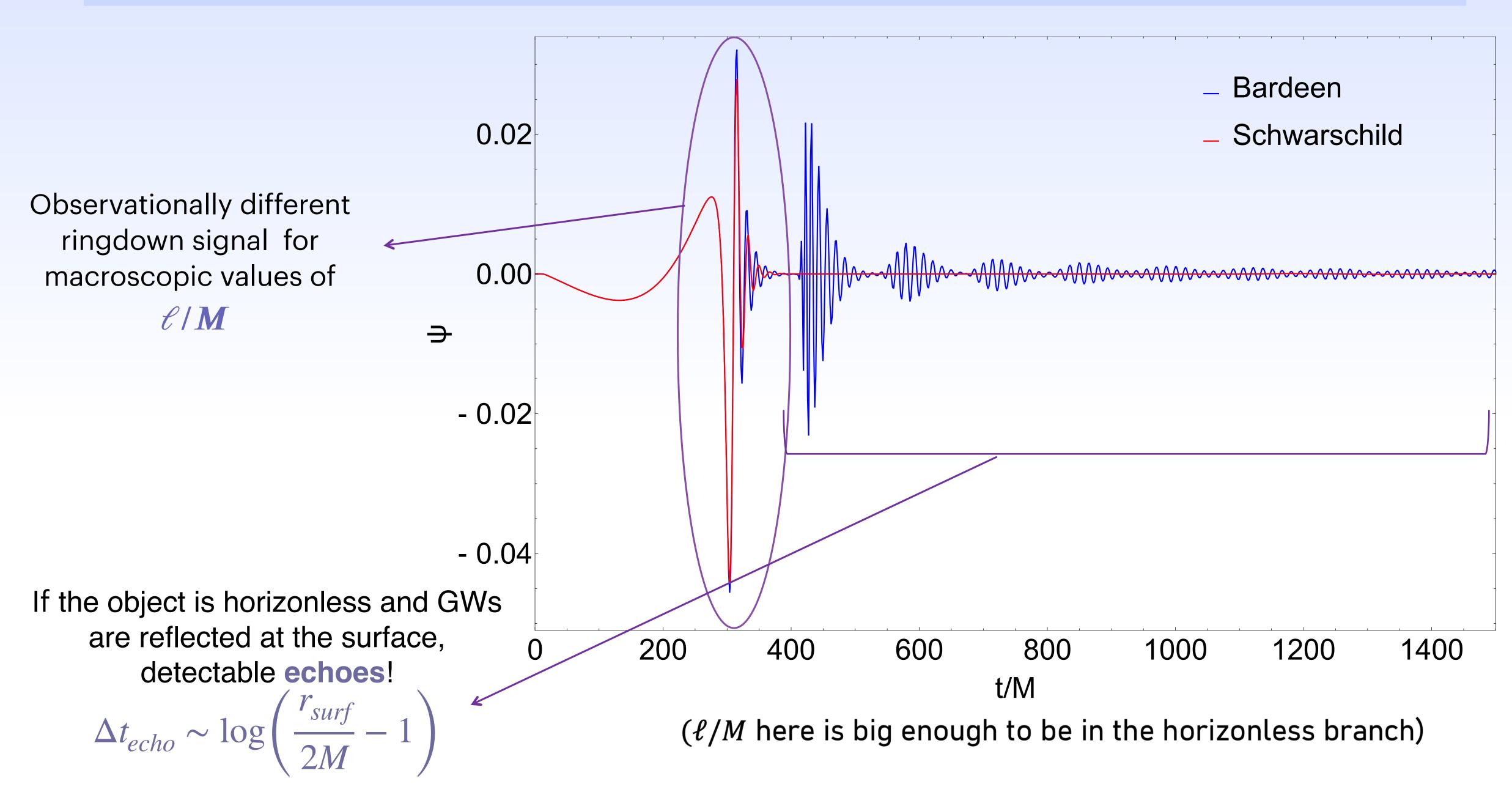


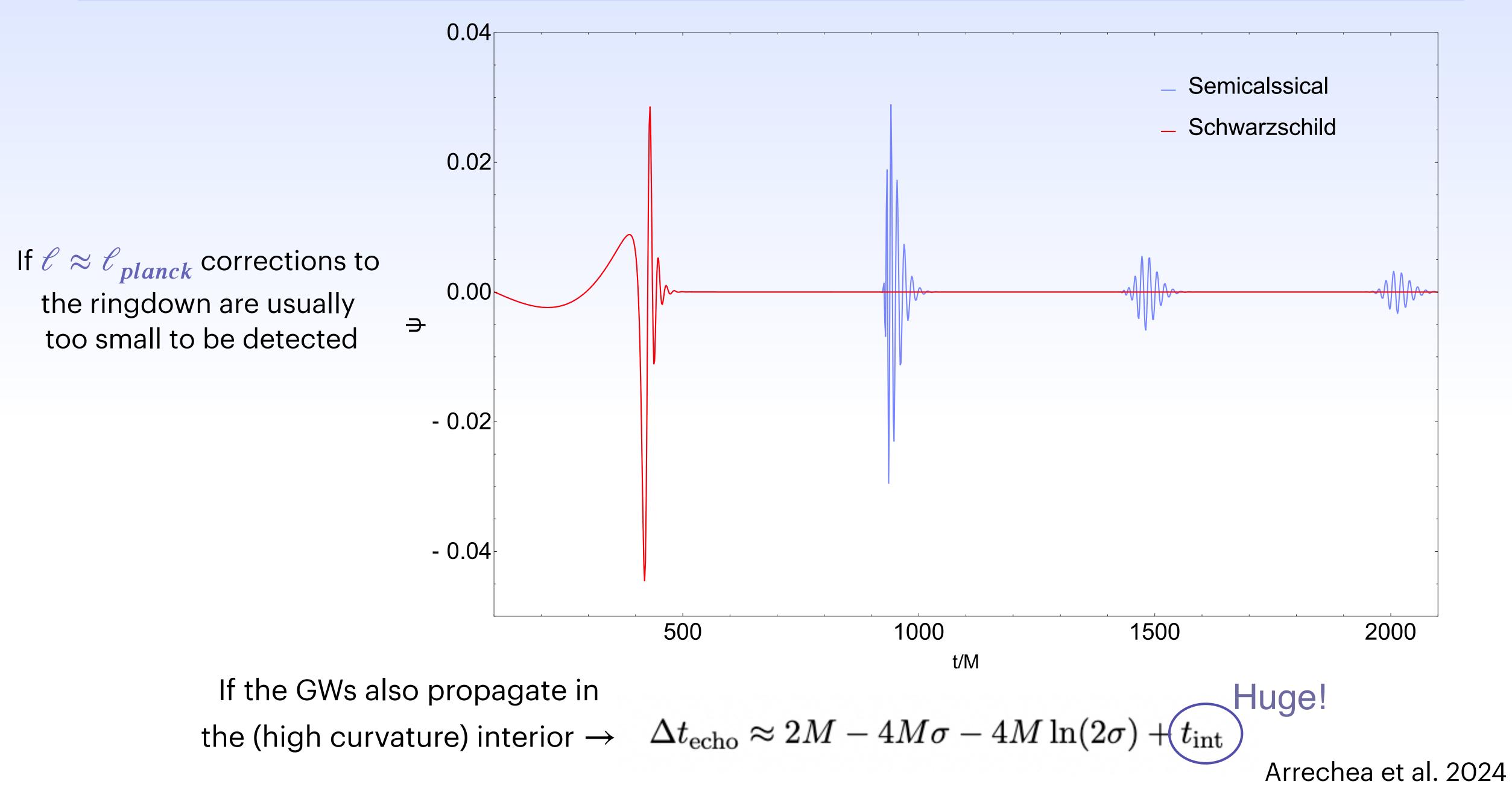
Observational signatures of the regularization?



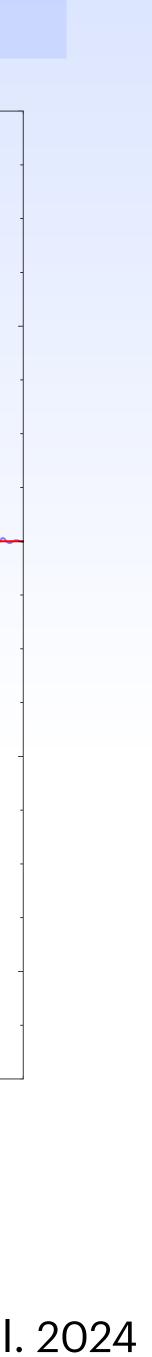


The post merger signal of a regular compact object





But...





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They all present phenomenological signatures with respect to singular BHs even if their observability is not guaranteed

Conclusions

In a complete theory of quantum gravity we expect the formation of spacetime singularities to be prevented

* There are few regular alternatives to describe ultra-compact objects

They all presents issues/instabilities when analysed in GR QG dynamics needed to really asses viability

Thank You

An example: the post-merger GWs signal

0.4

0.3

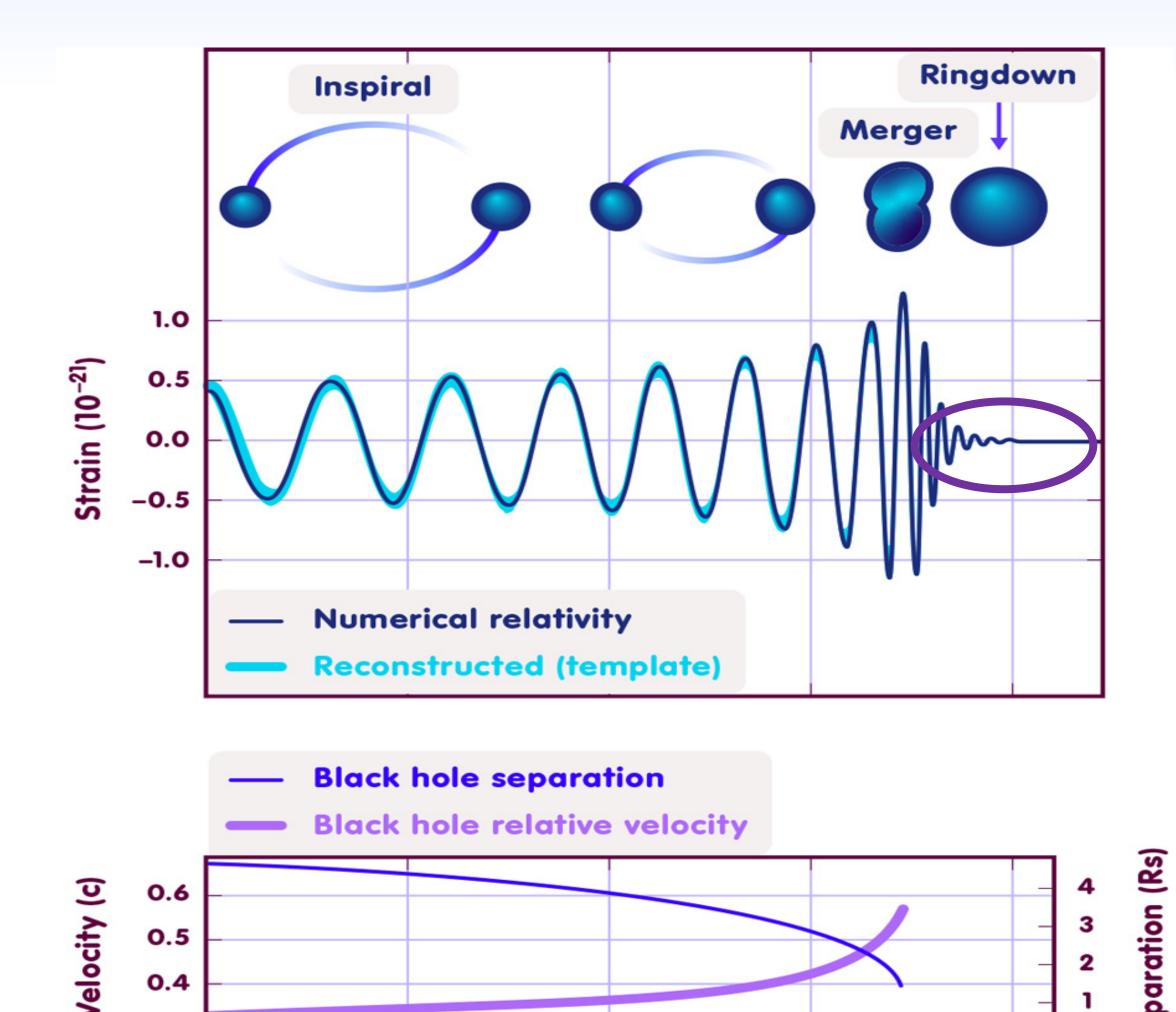
0.30

>

Ringdown signal:

- caused by the relaxation of the remnant object to its final, stationary state
- can be studied in perturbation theory:

$$g_{\mu\nu} = g^0_{\mu\nu} + h_{\mu\nu}$$



LIGO / Redesign: Daniela Leitner

0.40

0.35

Time (s)



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0.45

Perturbations around a spherically symmetric BH or an Horizonless Compact Object

$$\left(\frac{d^2}{dr_*^2} - \frac{d^2}{dt^2} - V_l(r)\right)h_{lm}(r) = 0$$

Feel very different potentials (different boundary conditions)

- We see echoes of the original ringdown signal
- We have a complete different spectrum of (trapped) QNMs

Echoes

Kokkotas (1996); Ferrari & Kokkotas (2000); Cardoso, Franzin, Pani (2016); Cardoso+ (2016)

