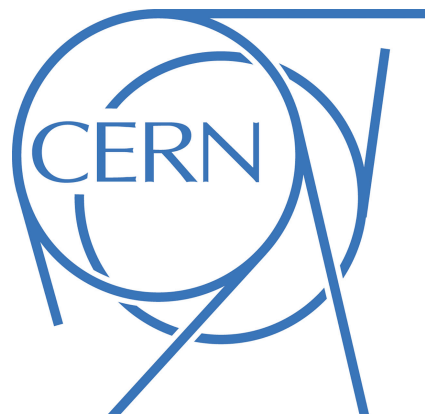


Swampland and Naturalness



Irene Valenzuela

CERN

IFT UAM-CSIC



“New approaches to naturalness”
Lyon, May 23, 2025



European Research Council
Established by the European Commission

MOTIVATION



Ferdi Rizkiyanto

MOTIVATION

Our universe contains gravity.

MOTIVATION

Our universe contains gravity.

Naturalness issues are based on a Wilsonian approach in QFT
(without gravity)

Our universe contains gravity.

Naturalness issues are based on a Wilsonian approach in QFT
(without gravity)

But gravity behaves differently than other interactions.

MOTIVATION

Our universe contains gravity.

Naturalness issues are based on a Wilsonian approach in QFT
(without gravity)

But gravity behaves differently than other interactions.

Can this help to explain naturalness problems of our universe?

Our universe contains gravity.

Naturalness issues are based on a Wilsonian approach in QFT
(without gravity)

But gravity behaves differently than other interactions.

Can this help to explain naturalness problems of our universe?

In this talk:

Requiring consistency with quantum gravity
principles can have implications at low energies.

(Swampland program)

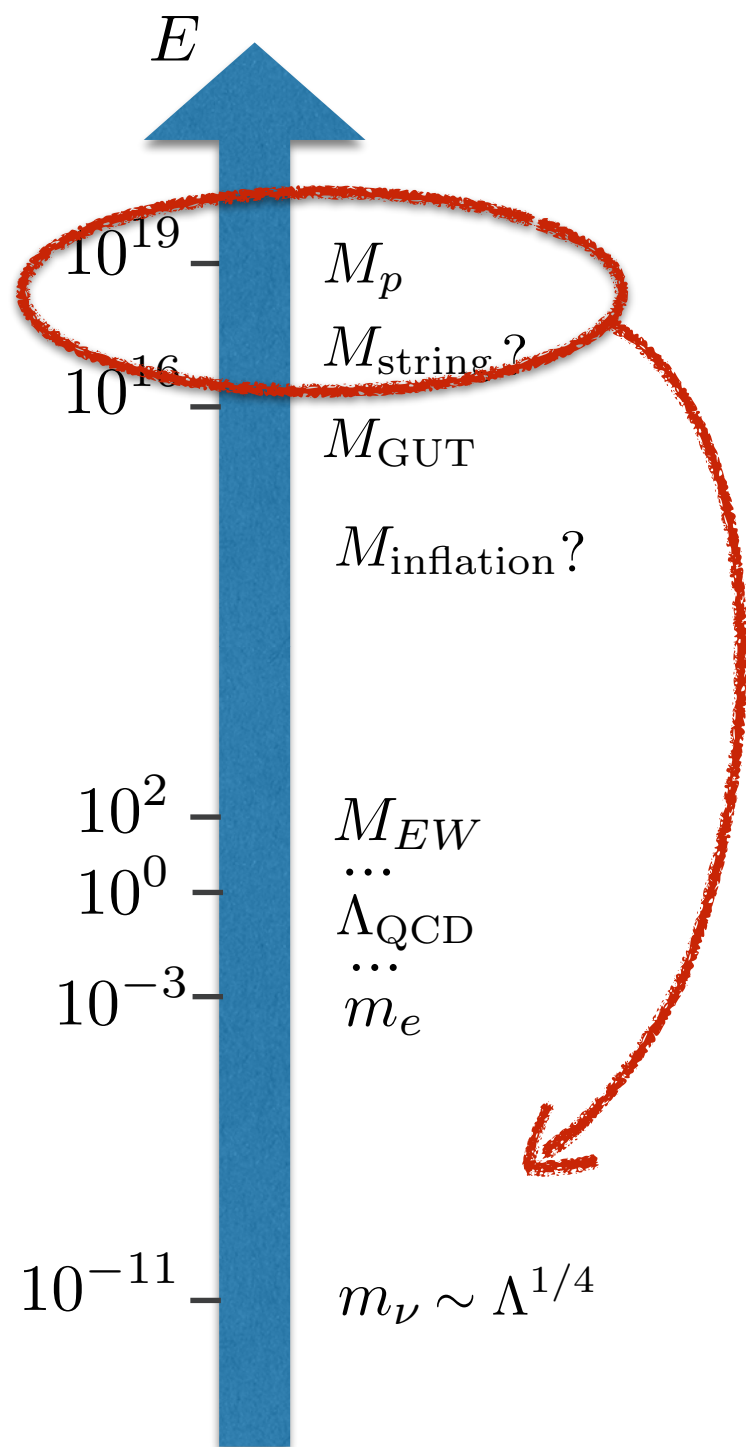
❖ Generic ideas/expectations

- UV/IR mixing of gravity
- Not everything goes (not every EFT can be consistently coupled to quantum gravity)

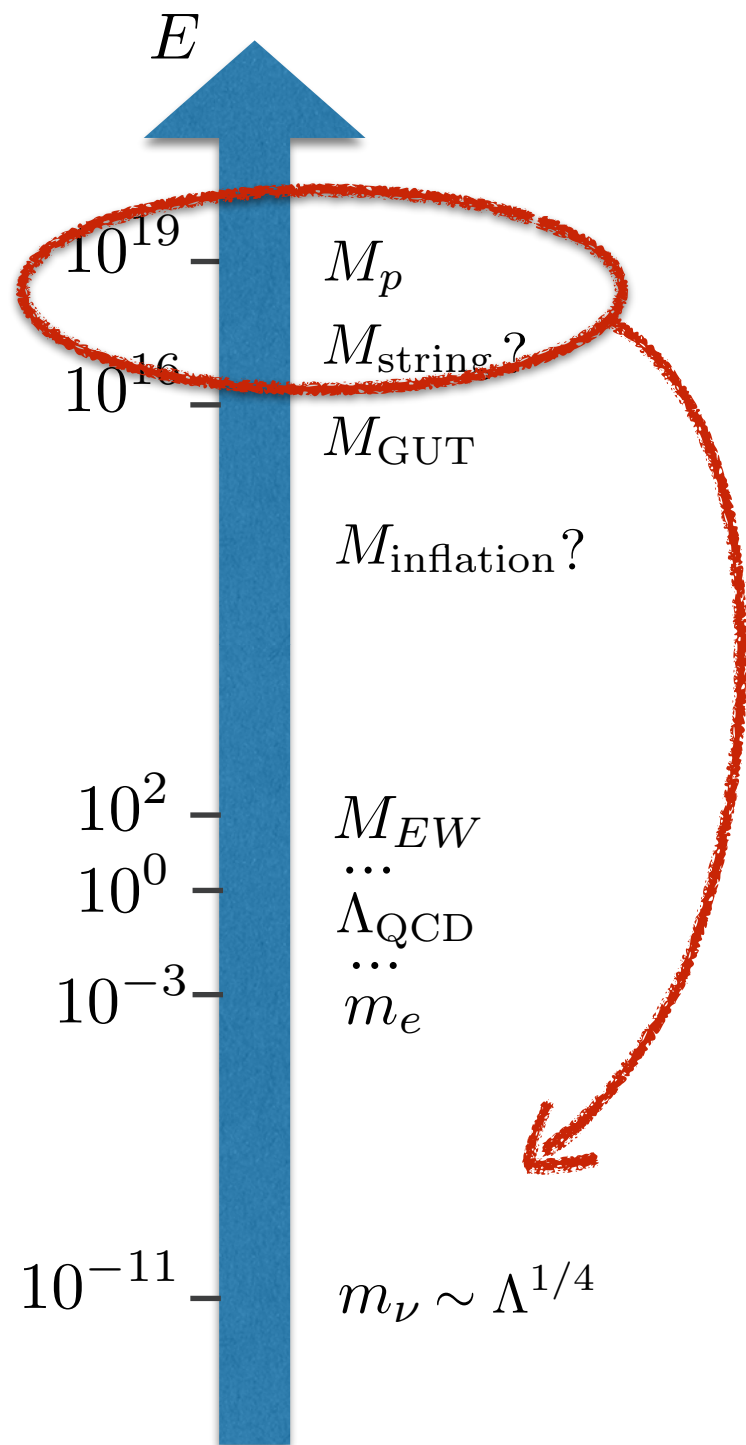
❖ Explicit examples of ‘unnatural parameters’ from IR perspective

- Breaking of global symmetries
- Light scalars and Weak Gravity Conjecture
- Dynamical parameters and Distance Conjecture
- Small vacuum energy and dark dimension
- AdS vacua from SM and neutrino masses

UV/IR mixing by gravity

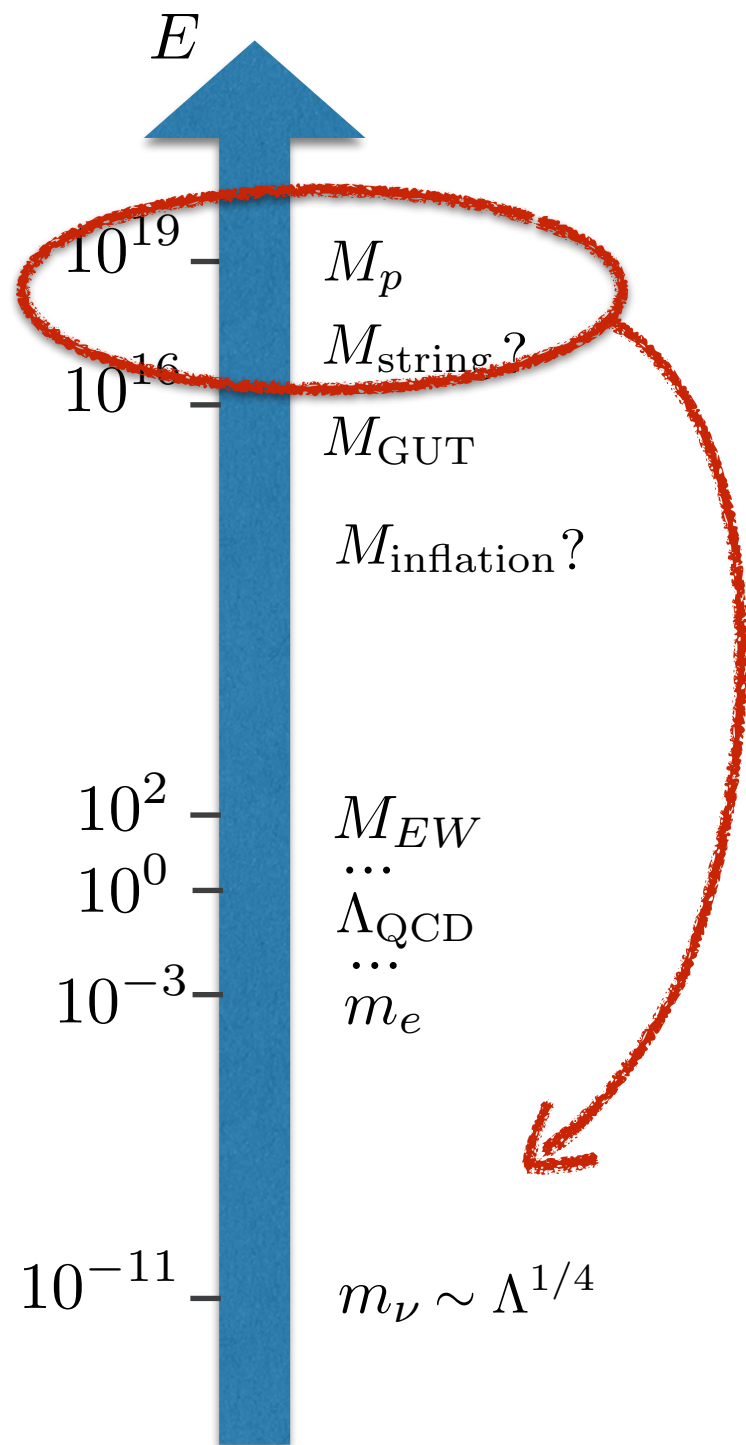


UV/IR mixing by gravity

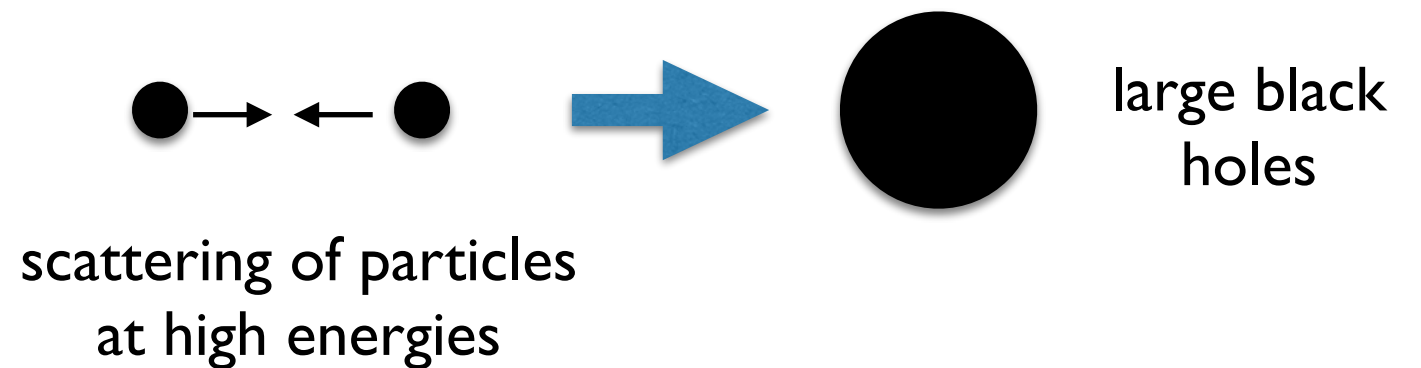


- ❖ **Black Hole Physics** makes manifest a correlation between high and low energy physics (UV/IR mixing)

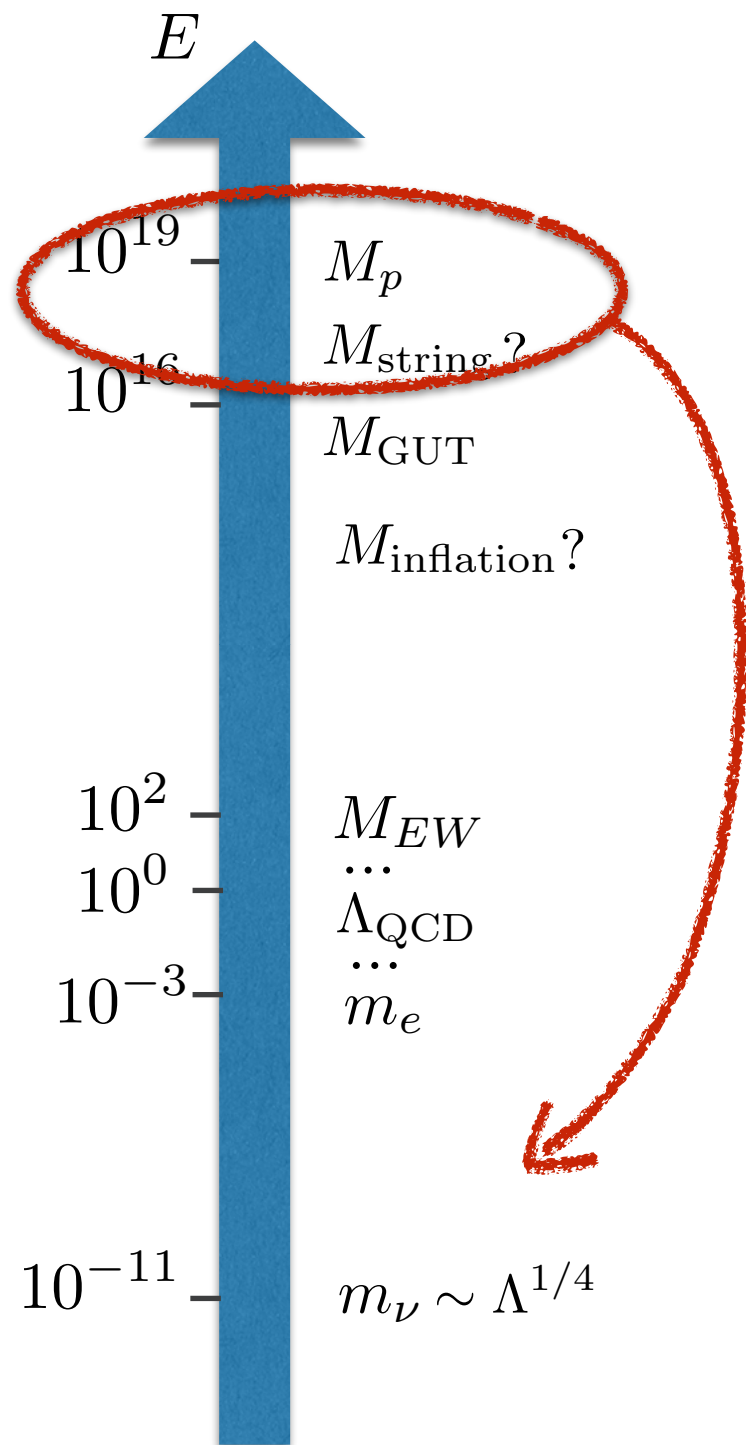
UV/IR mixing by gravity



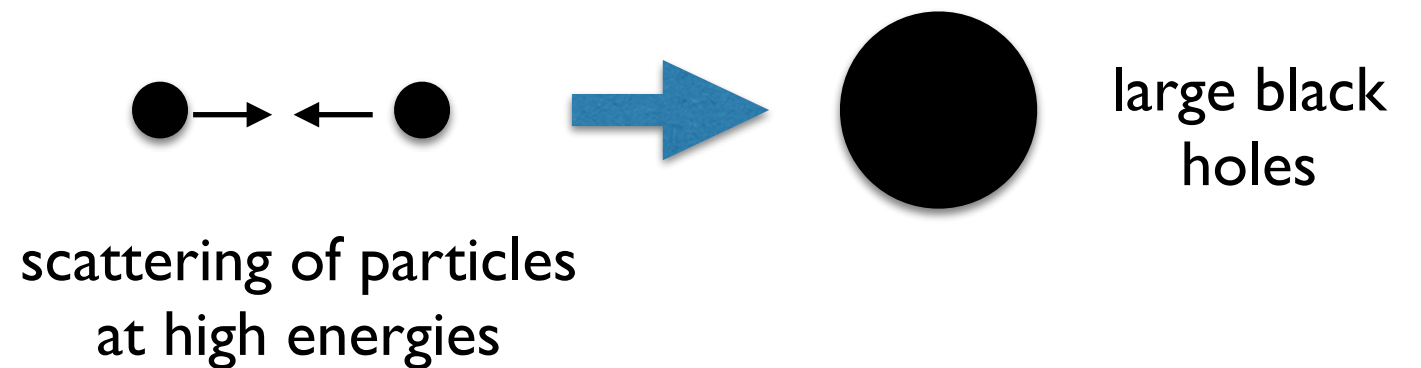
❖ **Black Hole Physics** makes manifest a correlation between high and low energy physics (UV/IR mixing)



UV/IR mixing by gravity

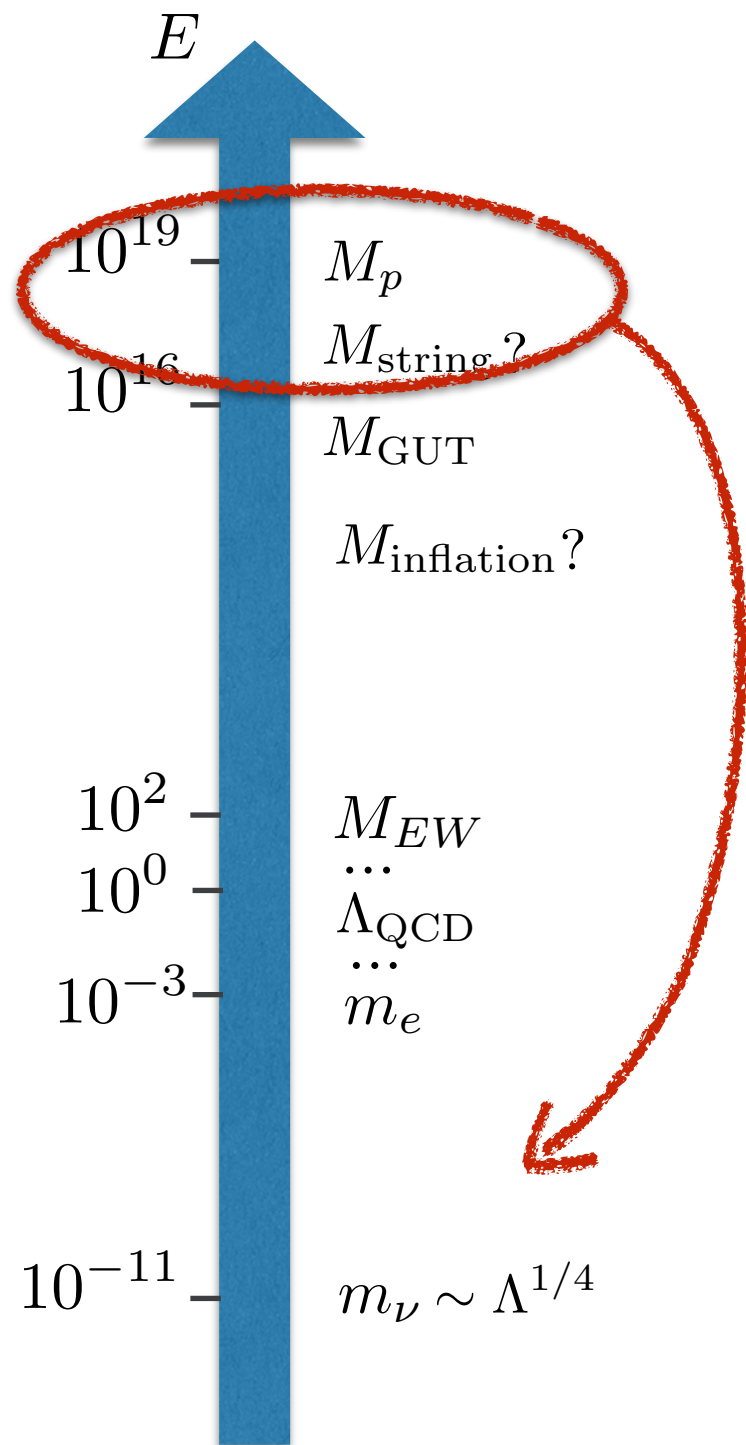


- ❖ **Black Hole Physics** makes manifest a correlation between high and low energy physics (UV/IR mixing)

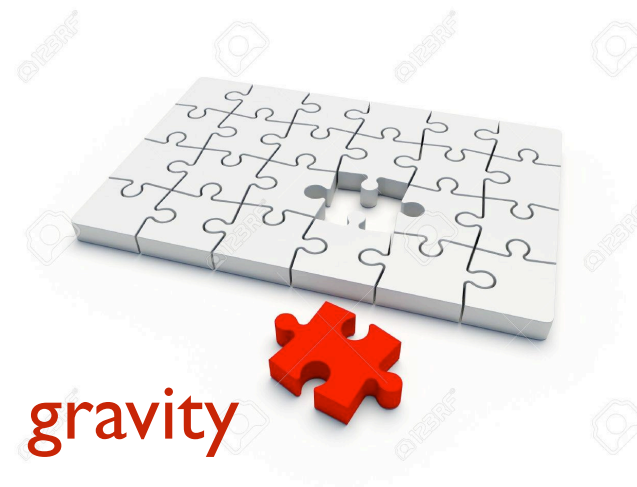


- ❖ Evidence from **String Theory** (dualities, modular invariance...)

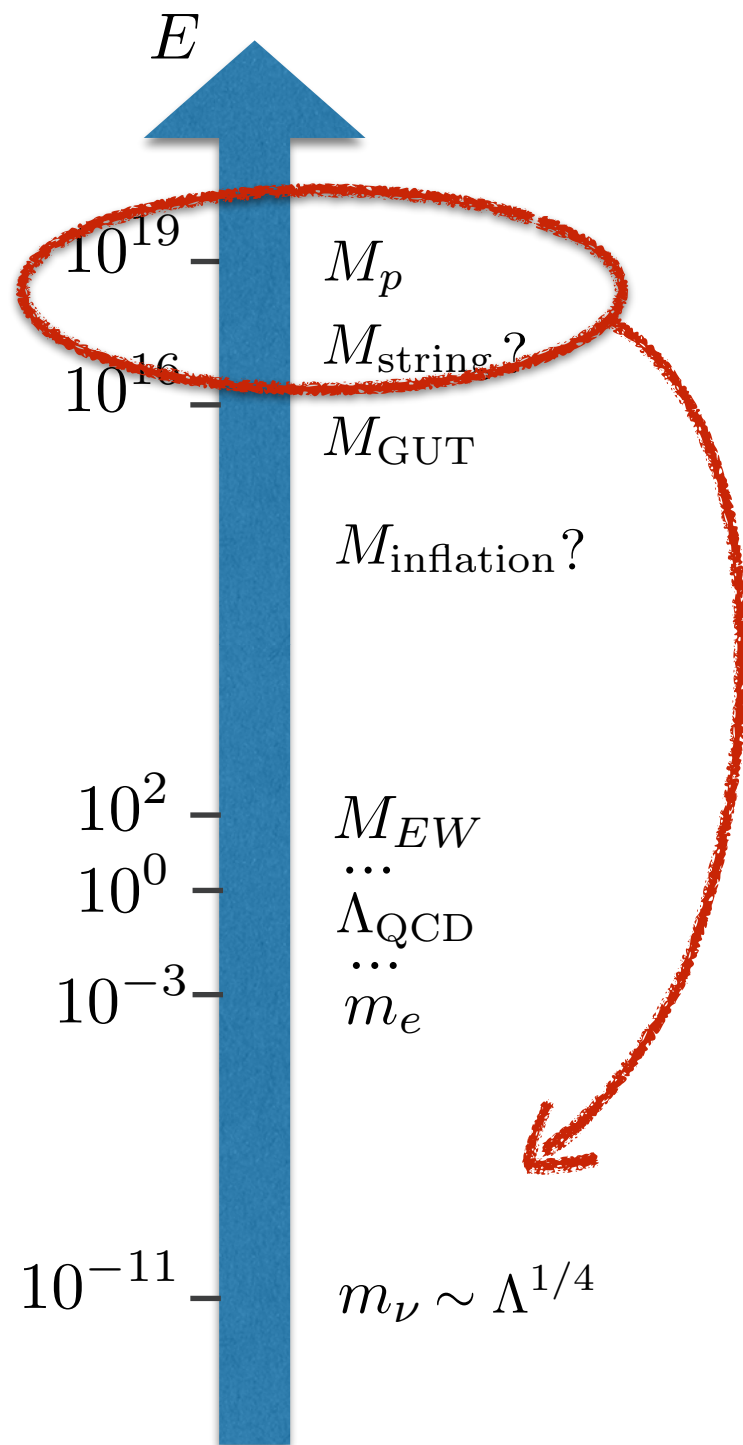
UV/IR mixing by gravity



Naturalness issues are an opportunity.



UV/IR mixing by gravity

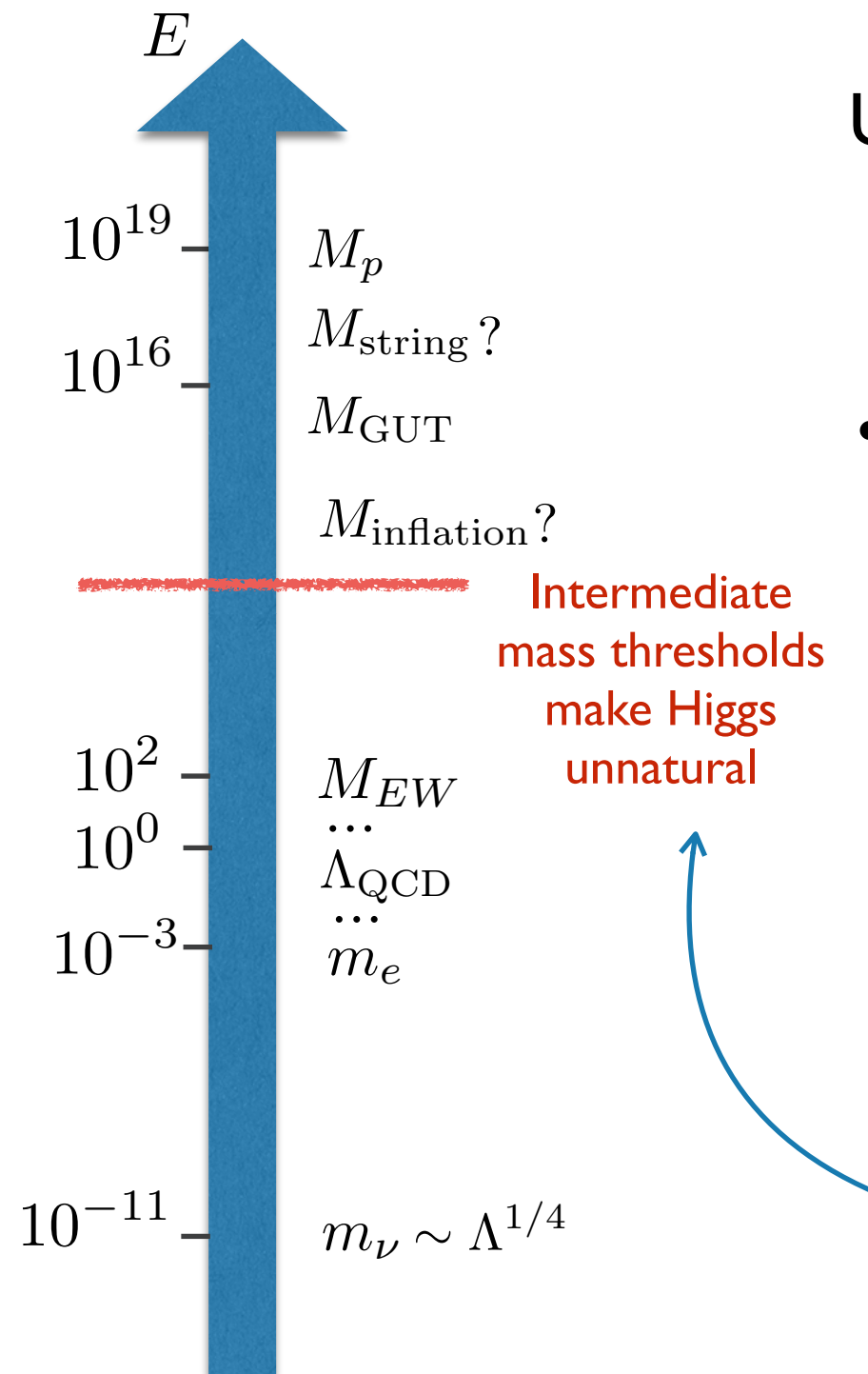


Using terminology of James Wells' talk:
(as I understood it)

- **Intrinsic hierarchy problem** gets solved

Once gravity kicks in, rules are different
(in fact, string theory implies infinity many states,
but divergences are cured by UV/IR mixing)

UV/IR mixing by gravity



Using terminology of James Wells' talk:
(as I understood it)

- **Intrinsic hierarchy problem** gets solved

Once gravity kicks in, rules are different
(in fact, string theory implies infinity many states,
but divergences are cured by UV/IR mixing)

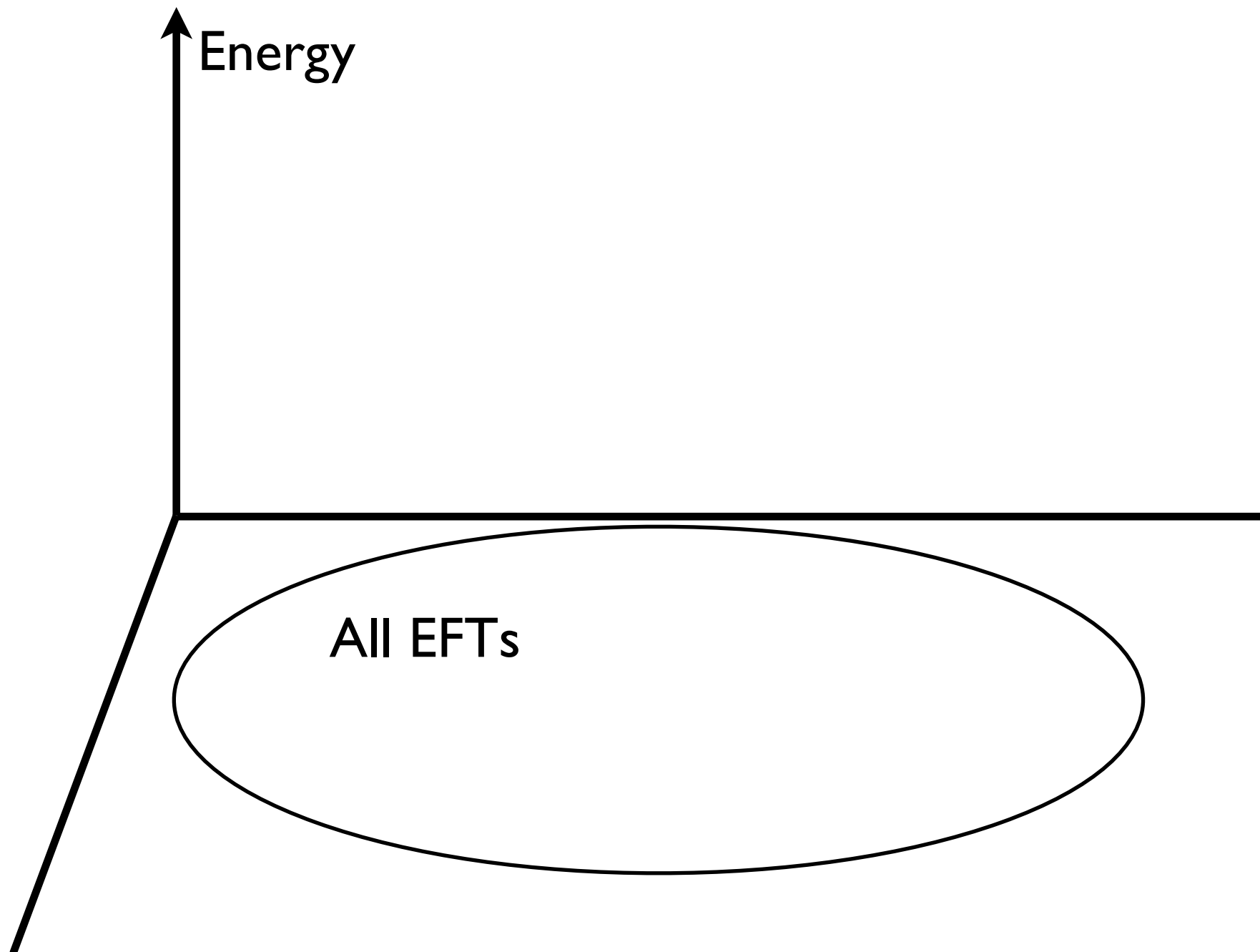
- However, **extrinsic hierarchy problem**
can still be there

At low energies, we get QFT coupled to Einstein
gravity and Wilsonian paradigm seems to hold

GENERIC ARGUMENTS

Not every low energy EFT coupled
to gravity can be UV completed
(in a quantum gravity theory)

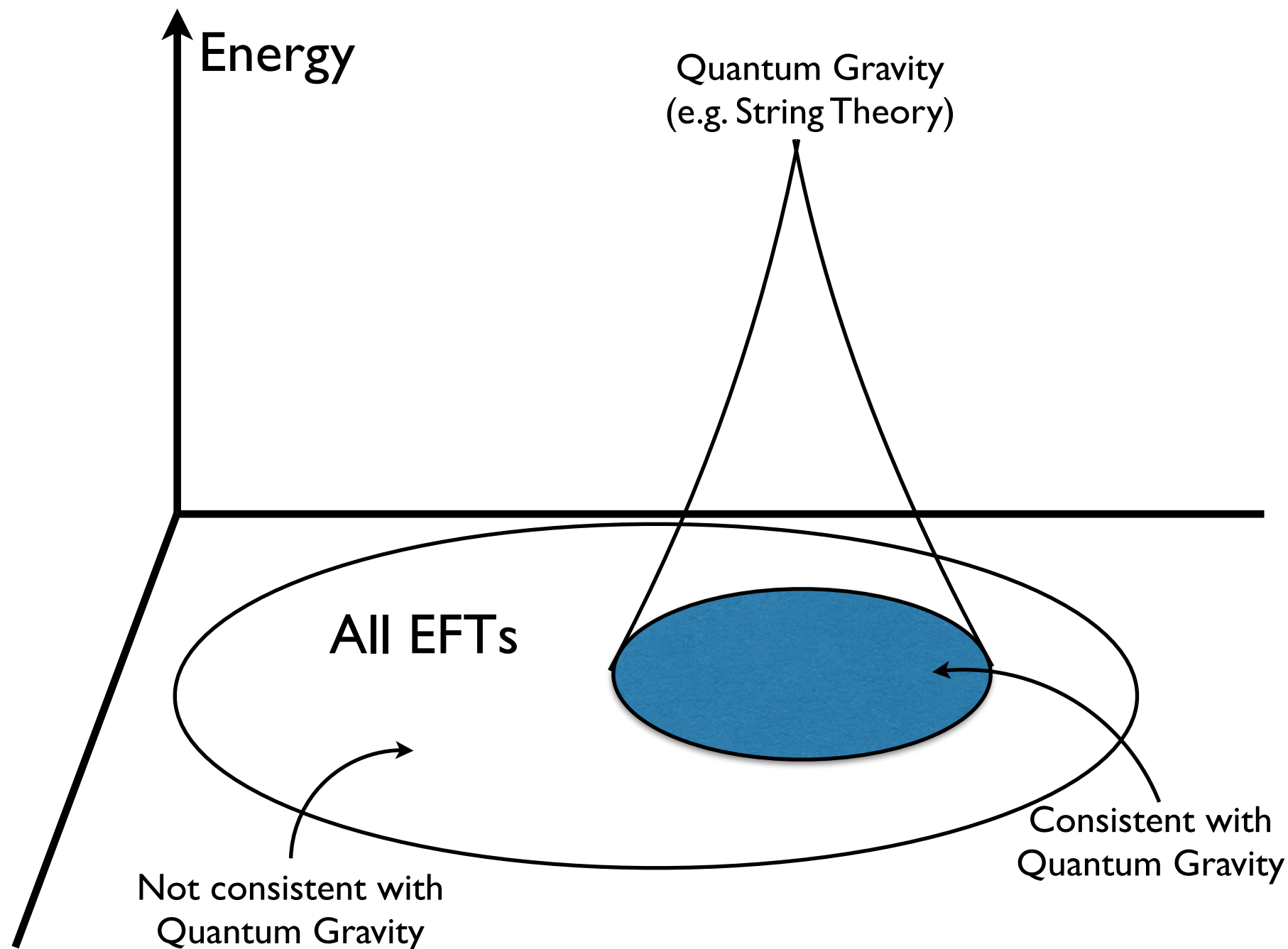
Not everything goes



Swampland Lectures/Reviews:
[Brennan,Vafa'17] [Palti'19]
[Van Beest,Calderon-
Infante,Mirfendereski,IV'21]

Not every low energy EFT coupled to gravity can be UV completed (in a quantum gravity theory)

Not everything goes

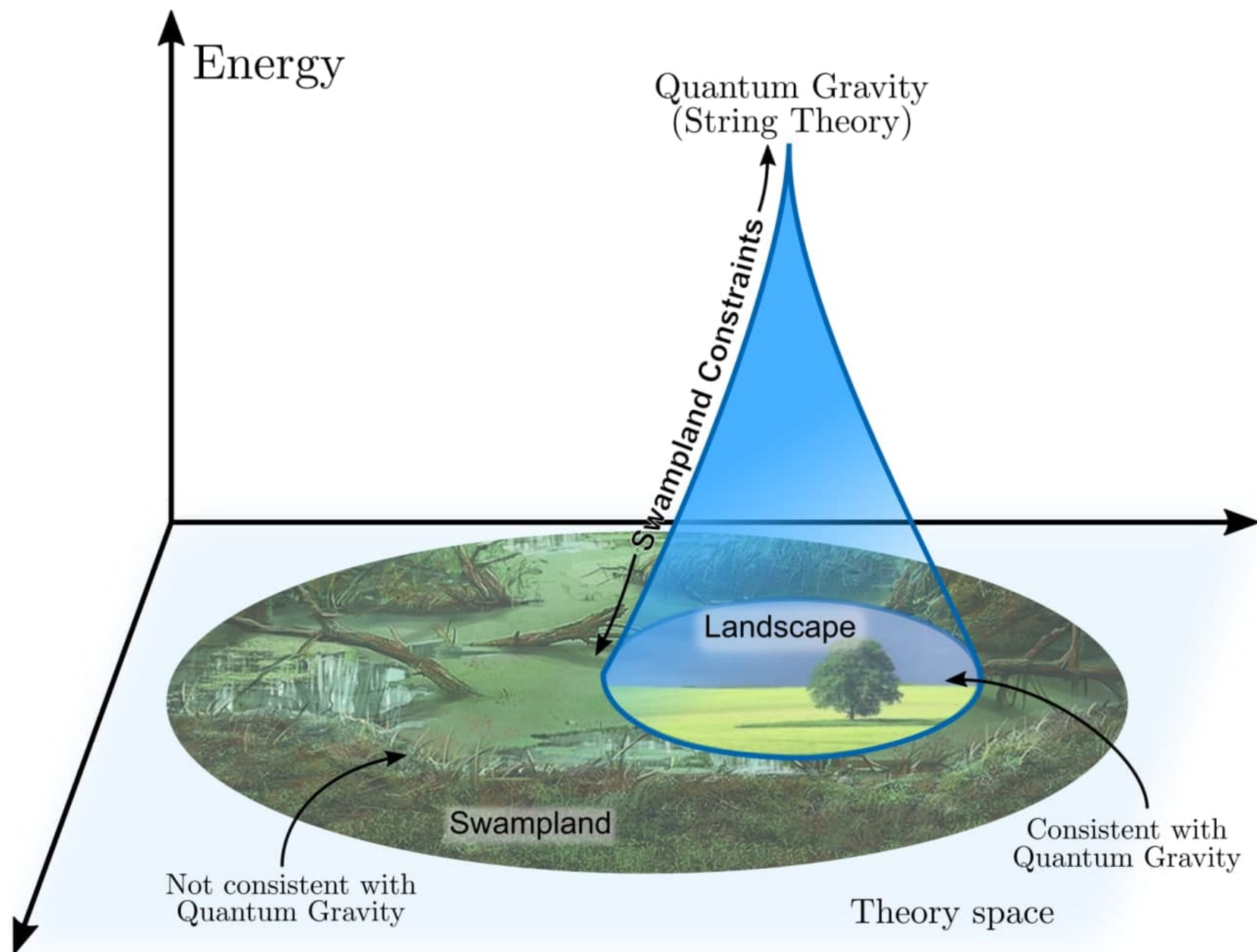


Swampland Lectures/Reviews:
[Brennan,Vafa'17] [Palti'19]
[Van Beest,Calderon-Infante,Mirfendereski,IV'21]

GENERIC ARGUMENTS

Not every low energy EFT coupled to gravity can be UV completed (in a quantum gravity theory)

Not everything goes



Swampland Lectures/Reviews:
[Brennan,Vafa'17] [Palti'19]
[Van Beest,Calderon-Infante,Mirfendereski,IV'21]

[Vafa'05]

Not everything goes

Goal of Swampland program:

Determine the constraints that an EFT must satisfy to be consistent with a UV completion in quantum gravity (QG)

Not everything goes

Goal of Swampland program:

Determine the constraints that an EFT must satisfy to be consistent with a UV completion in quantum gravity (QG)

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

What can it go here?

What is the quantum gravity cut-off?

Not everything goes

Goal of Swampland program:

Determine the constraints that an EFT must satisfy to be consistent with a UV completion in quantum gravity (QG)

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

What can it go here?

What is the quantum gravity cut-off?

Using:

- **String Theory** as a theoretical laboratory of QG to extract universal features
- Bottom-up arguments based on **black hole physics**, holography, scattering amplitudes/bootstrap,...

Not everything goes

Following James Wells' talk again:

Premise 2: Aleatory Parameters. Coefficients of the operators of the Ur-Theory are aleatorily assigned to each of the symmetry-allowed operators.

Not everything goes

Following James Wells' talk again:

Premise 2: Aleatory Parameters. Coefficients of the operators of the Ur-Theory are aleatorily assigned to each of the symmetry-allowed operators.

It is **false** if imposing quantum gravity UV consistency



not the entire space of parameters is allowed

Not everything goes

Following James Wells' talk again:

Premise 2: Aleatory Parameters. Coefficients of the operators of the Ur-Theory are aleatorily assigned to each of the symmetry-allowed operators.

It is **false** if imposing quantum gravity UV consistency



not the entire space of parameters is allowed

Let's understand the space of parameters/theories that are UV consistent before concluding that something is 'unnatural'

Maybe this can solve the extrinsic hierarchy problem!

- 1) Breaking of global symmetries
- 2) Light scalars and Weak Gravity Conjecture
- 3) Dynamical parameters and Distance Conjecture
- 4) Small vacuum energy and dark dimension
- 5) AdS vacua from SM and neutrino masses

I) **Breaking global symmetries**

I) **Breaking global symmetries**

Exact symmetries of the UV theory cannot be global (only “gauged”)

[Banks-Dixon'88] [Horowitz,Strominger,...] [Susskind] [Banks,Seiberg'11]

Exact symmetries of the UV theory cannot be global (only “gauged”)

[Banks-Dixon'88] [Horowitz,Strominger,...] [Susskind] [Banks,Seiberg'11]

Global symmetries

B-L in the SM

Continuous scalar shift symmetry

Flavour symmetries

Gauged symmetries

Electric charge

Modular symmetries, dualities

Discrete shift symmetry of a
fundamental axion

Exact symmetries of the UV theory cannot be global (only “gauged”)

[Banks-Dixon'88] [Horowitz,Strominger,...] [Susskind] [Banks,Seiberg'11]

Global symmetries

B-L in the SM

Continuous scalar shift symmetry

Flavour symmetries

Gauged symmetries

Electric charge

Modular symmetries, dualities

Discrete shift symmetry of a
fundamental axion

Evidence: String Theory + black hole arguments + proof in AdS/CFT

[Polchinski] [Harlow,Ooguri '18] [Harlow,Shaghoulian '20] [Chen,Lin '20] [Hsin et al '20] [Yonekura '20]

[Bah,Chen,Maldacena'22]

Exact symmetries of the UV theory cannot be global (only “gauged”)

[Banks-Dixon'88] [Horowitz,Strominger,...] [Susskind] [Banks,Seiberg'11]

Global symmetries

B-L in the SM

Continuous scalar shift symmetry

Flavour symmetries

Gauged symmetries

Electric charge

Modular symmetries, dualities

Discrete shift symmetry of a
fundamental axion

Evidence: String Theory + black hole arguments + proof in AdS/CFT

[Polchinski] [Harlow,Ooguri '18] [Harlow,Shaghoulian '20] [Chen,Lin '20] [Hsin et al '20] [Yonekura '20]

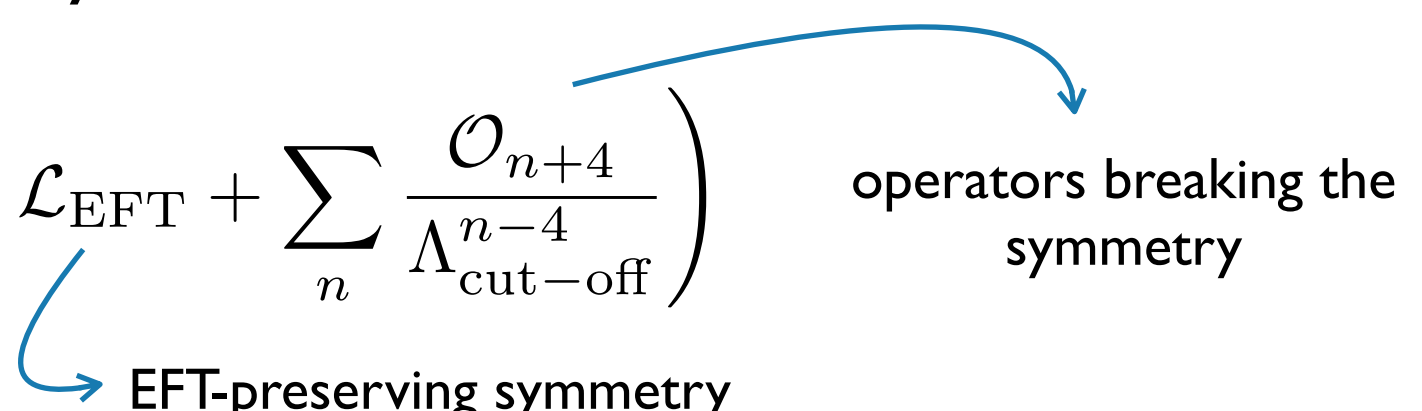
[Bah,Chen,Maldacena'22]

Immediate consequence:

We cannot invoke the existence of UV global symmetries to avoid
undesired higher dimensional operators

But we can get approximate symmetries in the IR

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

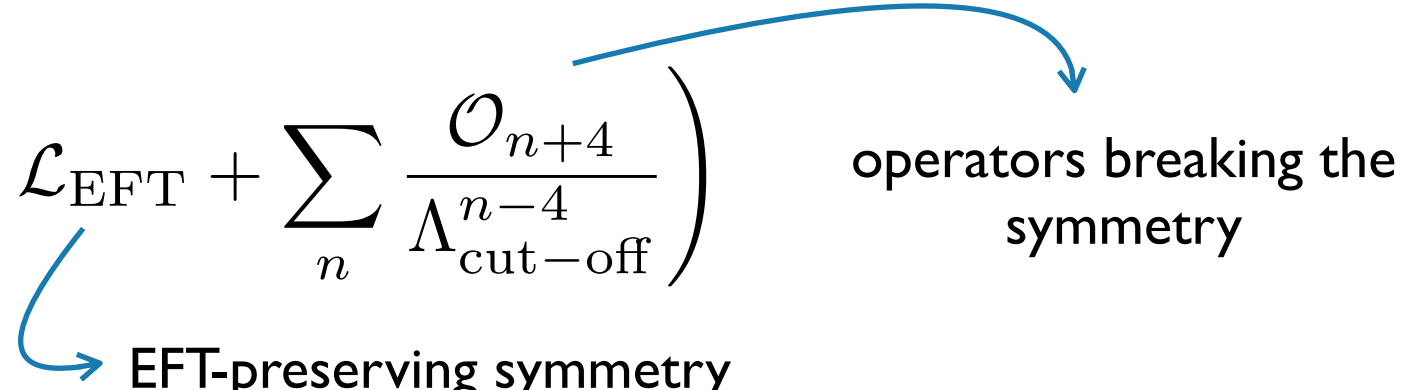


EFT-preserving symmetry

operators breaking the symmetry

But we can get approximate symmetries in the IR

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$



EFT-preserving symmetry

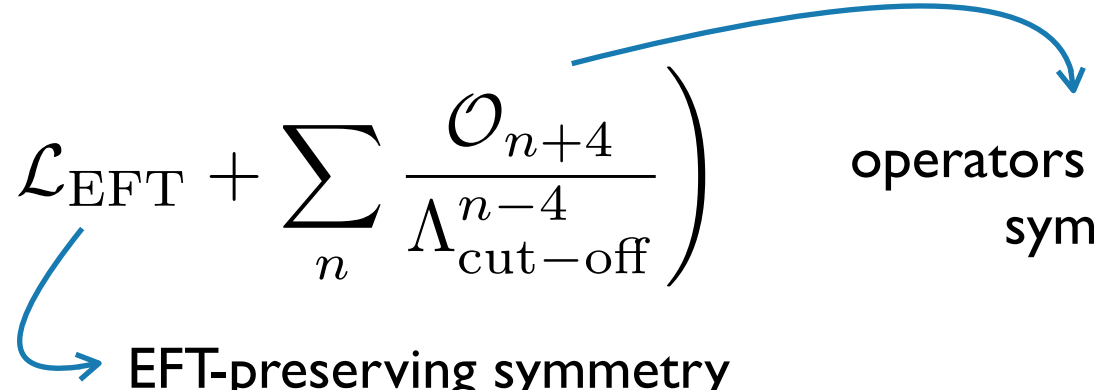
operators breaking the symmetry

We can learn a lot from understanding how symmetries are broken

I) How much are they broken?

But we can get approximate symmetries in the IR

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$


operators breaking the symmetry
EFT-preserving symmetry

We can learn a lot from understanding how symmetries are broken

I) How much are they broken?

[Daus et al '20] [ongoing ;)]

- Ordinary (0-form) symmetries : broken at least by $\sim e^{-\left(\frac{M_p}{\Lambda_{\text{cut-off}}}\right)^2}$

We can suppress the breaking by lowering the cut-off scale

But we can get approximate symmetries in the IR

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

We can learn a lot from understanding how symmetries are broken

I) How much are they broken?

[Daus et al '20] [ongoing ;)]

- Ordinary (0-form) symmetries : broken at least by $\sim e^{-\left(\frac{M_p}{\Lambda_{\text{cut-off}}}\right)^2}$

We can suppress the breaking by lowering the cut-off scale

- Non-invertible symmetries : broken by quantum effects in gravity

... but only at loop order $L \geq 1$!! [Kaidi, Tachikawa, Zhang '24]

[Heckman, McNamara, Montero, Sharon, Vafa, IV, '24]

2) There can be correlations to ensure all symmetries are broken

$$\mathcal{L} = \phi \operatorname{Tr}(F \wedge F) + \dots$$

d=dimension space-time

$j = \operatorname{Tr}(F \wedge F)$ Chern-Weyl current of a (d-5)-form global symmetry
which is gauged thanks to the coupling to the axion

[Heidenreich,McNamara,Montero,Reece,Rudelius,IV, '20]

2) There can be correlations to ensure all symmetries are broken

$$\mathcal{L} = \phi \operatorname{Tr}(F \wedge F) + \dots$$

d=dimension space-time

$j = \operatorname{Tr}(F \wedge F)$ Chern-Weyl current of a (d-5)-form global symmetry
which is gauged thanks to the coupling to the axion

[Heidenreich,McNamara,Montero,Reece,Rudelius,IV, '20]

- Existence of axion helps us to avoid a global symmetry

2) There can be correlations to ensure all symmetries are broken

$$\mathcal{L} = \phi \left(\text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots \right)$$

d=dimension space-time

$j = \text{Tr}(F \wedge F)$ Chern-Weyl current of a (d-5)-form global symmetry
which is gauged thanks to the coupling to the axion

[Heidenreich,McNamara,Montero,Reece,Rudelius,IV, '20]

- Existence of axion helps us to avoid a global symmetry
- If the axion couples to other fields, we get a new Chern-Weyl current for each term, and only $j = \text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots$ is gauged

2) There can be correlations to ensure all symmetries are broken

$$\mathcal{L} = \phi \left(\text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots \right)$$

d=dimension space-time

$j = \text{Tr}(F \wedge F)$ Chern-Weyl current of a (d-5)-form global symmetry
which is gauged thanks to the coupling to the axion

[Heidenreich,McNamara,Montero,Reece,Rudelius,IV,'20]

- Existence of axion helps us to avoid a global symmetry
- If the axion couples to other fields, we get a new Chern-Weyl current for each term, and only $j = \text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots$ is gauged

The other linear combinations remain as global symmetries and
must be broken by quantum gravity consistency

➡ there must be a way to transform one charge into another

$$\int \text{Tr}(F \wedge F) \longleftrightarrow \int \text{Tr}(\tilde{F} \wedge \tilde{F})$$

2) There can be correlations to ensure all symmetries are broken

$$\mathcal{L} = \phi \left(\text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots \right)$$

d=dimension space-time

$j = \text{Tr}(F \wedge F)$ Chern-Weyl current of a (d-5)-form global symmetry
which is gauged thanks to the coupling to the axion

[Heidenreich,McNamara,Montero,Reece,Rudelius,IV,'20]

- **Existence of axion** helps us to avoid a global symmetry
- If the axion couples to other fields, we get a new Chern-Weyl current for each term, and only $j = \text{Tr}(F \wedge F) + \text{Tr}(\tilde{F} \wedge \tilde{F}) + \text{Tr}(R \wedge R) + \dots$ is gauged

The other linear combinations remain as global symmetries and must be broken by quantum gravity consistency

➡ there must be a way to transform one charge into another

$$\int \text{Tr}(F \wedge F) \longleftrightarrow \int \text{Tr}(\tilde{F} \wedge \tilde{F})$$

Can this ameliorate the QCD axion quality problem?

EXPLICIT EXAMPLES

Weak Gravity Conjecture:

2) Light scalars, weak couplings and the Weak Gravity Conjecture

[Arkani-Hamed et al'06]

2) Light scalars, weak couplings and the Weak Gravity Conjecture

Weak Gravity Conjecture:

[Arkani-Hamed et al'06]

Given an EFT with a long-range gauge interaction and gravity, there must exist a particle satisfying:

mass in Planck units

$$\frac{Q}{m} \geq 1$$

electric charge

fixed by black hole extremality bound

so that the gravitational force acts weaker than the gauge force



$$F_g = \frac{m^2}{r^2} \leq F_e = \frac{Q^2}{r^2}$$

2) Light scalars, weak couplings and the Weak Gravity Conjecture

Weak Gravity Conjecture:

[Arkani-Hamed et al'06]

Given an EFT with a long-range gauge interaction and gravity, there must exist a particle satisfying:

mass in Planck units

$$\frac{Q}{m} \geq 1$$

electric charge

fixed by black hole extremality bound



so that the gravitational force acts weaker than the gauge force

$$F_g = \frac{m^2}{r^2} \leq F_e = \frac{Q^2}{r^2}$$

Satisfied by the electron in the SM, but puts constraints on BSM with weakly coupled dark photons

2) Light scalars, weak couplings and the Weak Gravity Conjecture

Weak Gravity Conjecture:

[Arkani-Hamed et al'06]

Given an EFT with a long-range gauge interaction and gravity, there must exist a particle satisfying:

mass in Planck units

$$\frac{Q}{m} \geq 1$$

electric charge

fixed by black hole extremality bound



so that the gravitational force acts weaker than the gauge force

$$F_g = \frac{m^2}{r^2} \leq F_e = \frac{Q^2}{r^2}$$

Satisfied by the electron in the SM, but puts constraints on BSM with weakly coupled dark photons


Moreover, the EFT breaks down at


$$\Lambda_{\text{cut-off}} \sim g M_p$$

(from applying WGC to magnetic gauge field)


2) Light scalars, weak couplings and the Weak Gravity Conjecture

Is this “natural” for a scalar? $Q \geq m$

 quadratic divergences


 logarithmic running (recall $Q = qg$)


How is this satisfied in explicit string theory examples?

- At weak coupling ($g \ll 1$) there is a tower of infinitely many particles satisfying WGC  changes scaling of quantum corrections [Grimm, Palti, IV'18]
[Heidenreich et al'18]


2) Light scalars, weak couplings and the Weak Gravity Conjecture

Is this “natural” for a scalar? $Q \geq m$

 quadratic divergences

 logarithmic running (recall $Q = qg$)

How is this satisfied in explicit string theory examples?

- At weak coupling ($g \ll 1$) there is a tower of infinitely many particles satisfying WGC  changes scaling of quantum corrections [Grimm, Palti, IV'18]
[Heidenreich et al'18]

$$\frac{1}{g_{IR}^2} = \frac{1}{g_{UV}^2} - \sum_k^N \left(\frac{8q_k^2}{3\pi^2} \log \frac{\Lambda_{UV}}{m_k} \right) \begin{cases} (\delta g)^{-2} \sim N \log m_0 & \text{for } N \text{ fields} \\ & \text{(logarithmic running)} \\ (\delta g)^2 \sim m^2 & \text{for infinite tower } m_k = k m \\ & \text{(power law!) } k = 0, 1, \dots, \infty \end{cases}$$

consistent with $Q \geq m$ across the RG flow

using $\Lambda_{UV} = \frac{M_p}{\sqrt{N}}$ (species scale) [Dvali'07]

2) Light scalars, weak couplings and the Weak Gravity Conjecture

Is this “natural” for a scalar? $Q \geq m$

quadratic divergences

logarithmic running (recall $Q = qg$)

How is this satisfied in explicit string theory examples?

- At weak coupling ($g \ll 1$) there is a tower of infinitely many particles satisfying WGC \rightarrow changes scaling of quantum corrections [Grimm, Palti, IV'18]
[Heidenreich et al'18]
- For $g \sim \mathcal{O}(1)$ we can have just a few WGC-states, but examples under control are typically supersymmetric

Open question whether they can be scalars with $m < \Lambda_{UV}$

3) Dynamical parameters and Distance Conjecture

Every continuous parameter is dynamical:
it is set by the vacuum expectation value of scalar fields

$m(\phi), g(\phi) \dots$



parametrize size/form extra
dimensions in string theory

Universal lesson from string theory, also
necessary to avoid (-1)-form global symmetries

3) Dynamical parameters and Distance Conjecture

Every continuous parameter is dynamical:
it is set by the vacuum expectation value of scalar fields

$m(\phi), g(\phi) \dots$



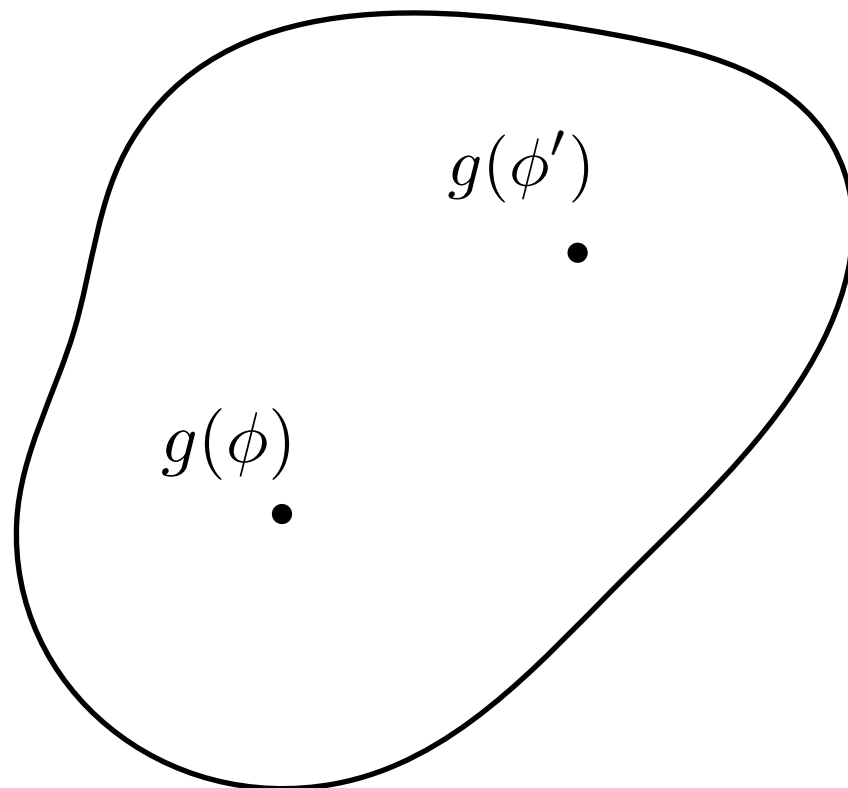
parametrize size/form extra
dimensions in string theory

Universal lesson from string theory, also
necessary to avoid (-1)-form global symmetries

parameter space of the EFT

=

scalar field space of string theory



3) Dynamical parameters and Distance Conjecture

Every continuous parameter is dynamical:
it is set by the vacuum expectation value of scalar fields

$m(\phi), g(\phi) \dots$

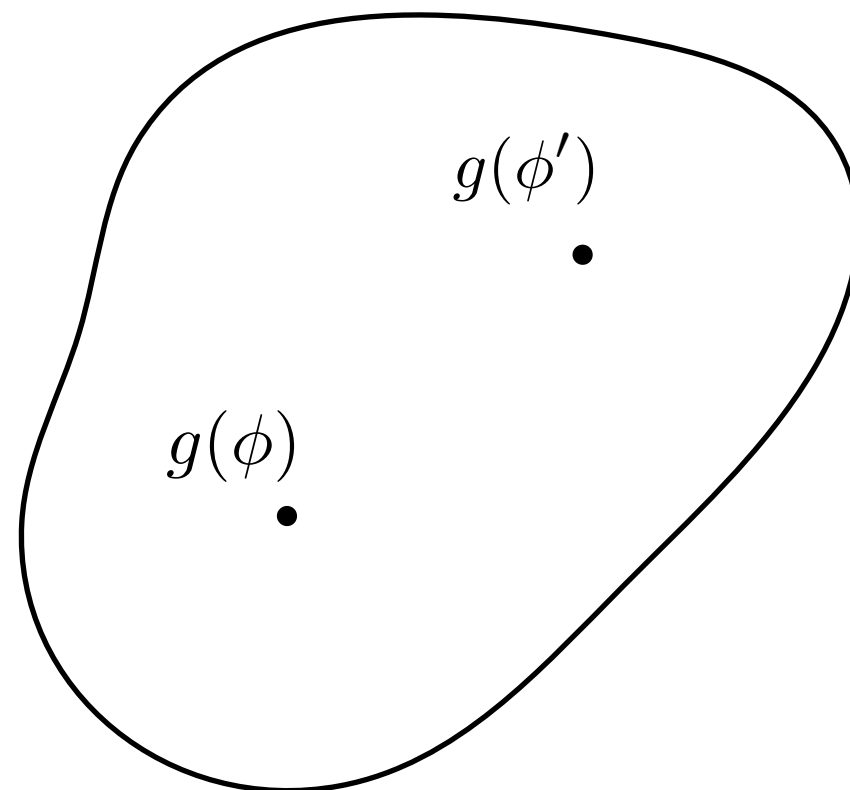
parametrize size/form extra dimensions in string theory

Universal lesson from string theory, also necessary to avoid (-1)-form global symmetries

parameter space of the EFT

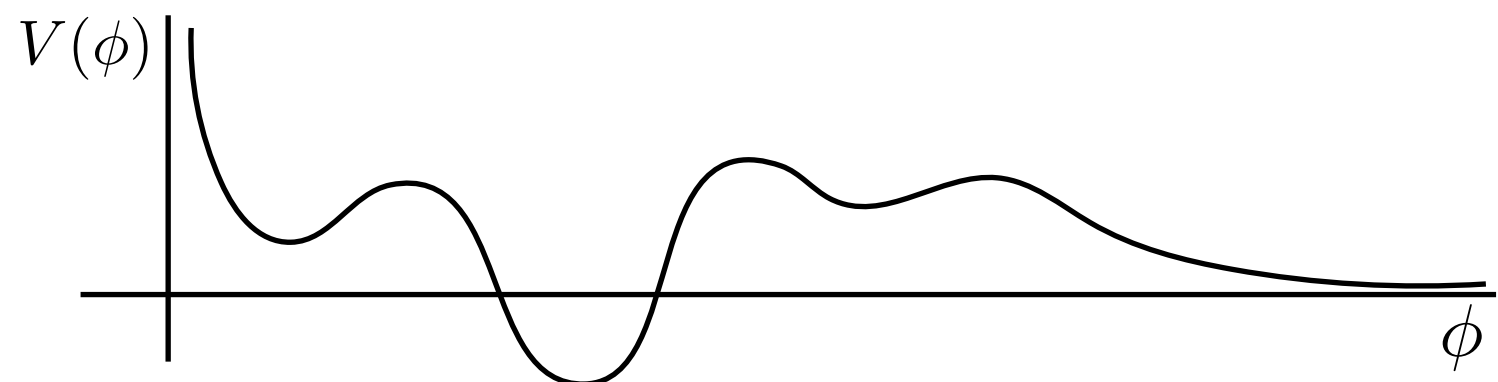
=

scalar field space of string theory



Big challenge!

Naturalness issues get translated to finding a dynamical mechanism that stabilizes the scalars precisely at the values required to reproduce observations



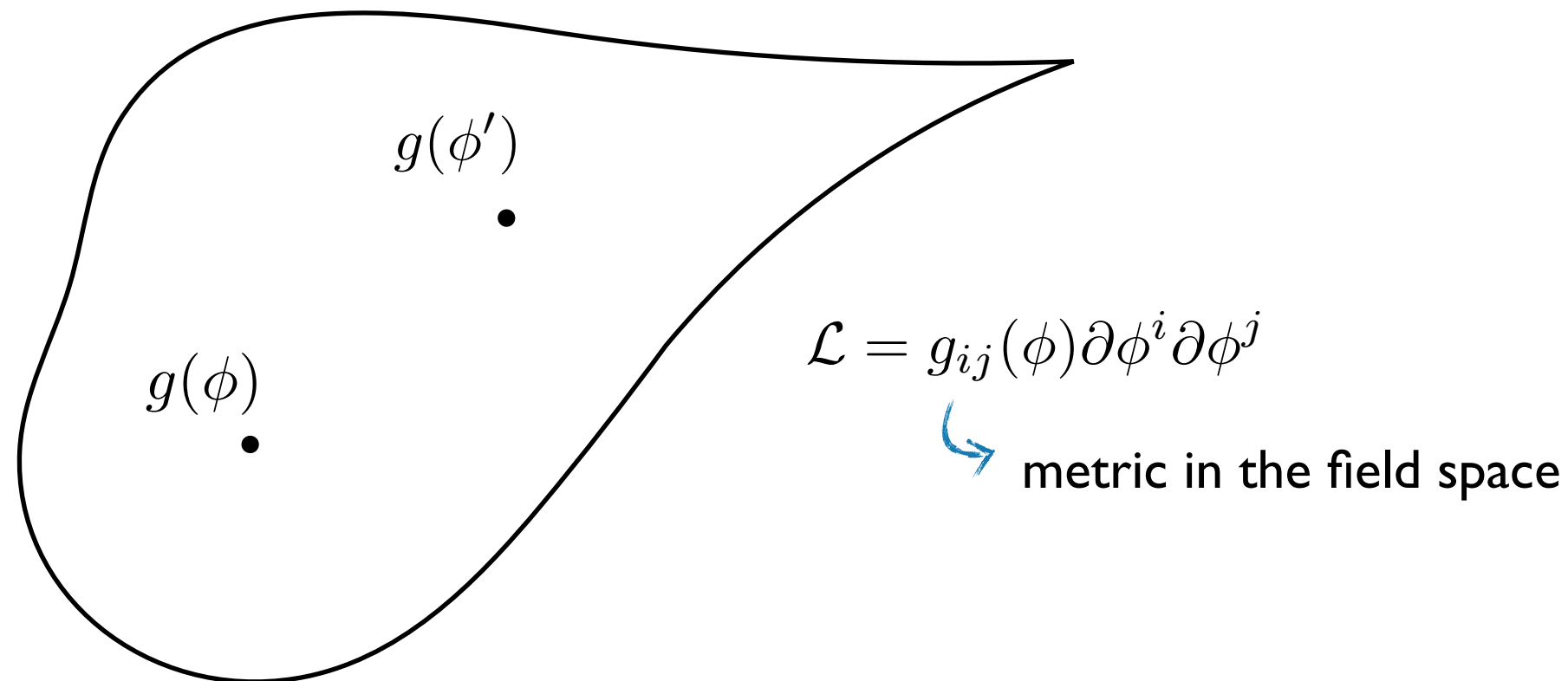
3) **Dynamical parameters and Distance Conjecture**

Universal features:

3) Dynamical parameters and Distance Conjecture

Universal features:

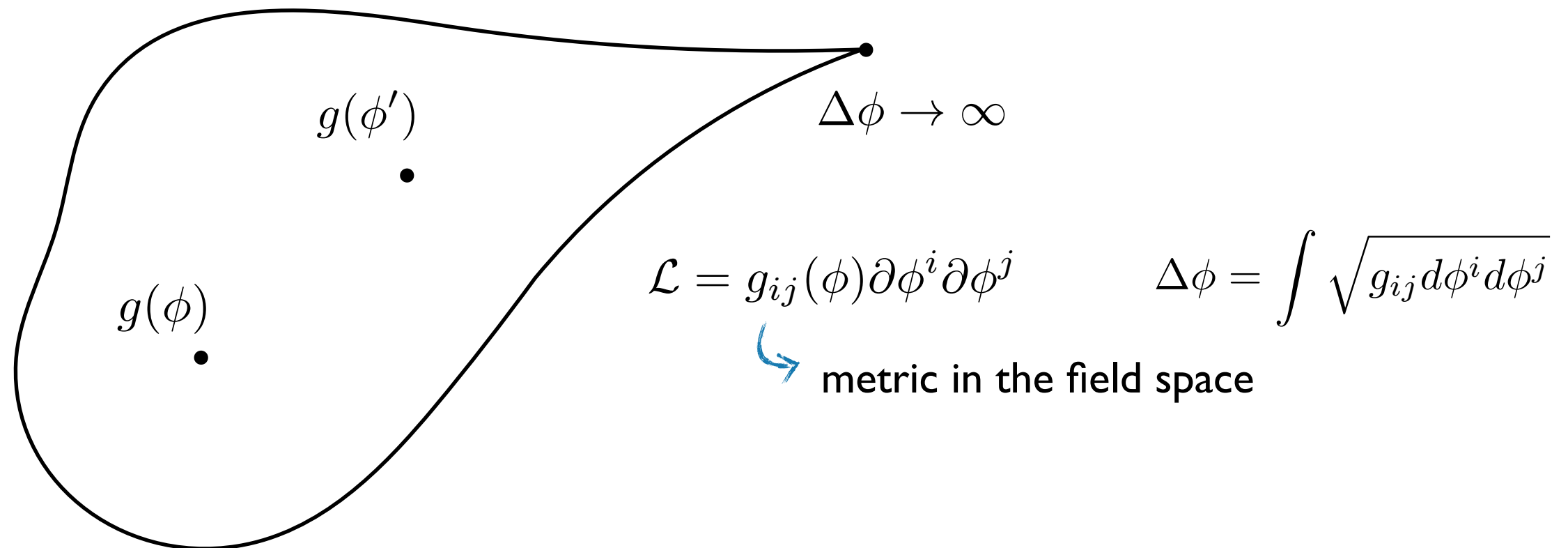
Global symmetries, vanishing couplings or large hierarchies occur at infinite distance in this field space



3) Dynamical parameters and Distance Conjecture

Universal features:

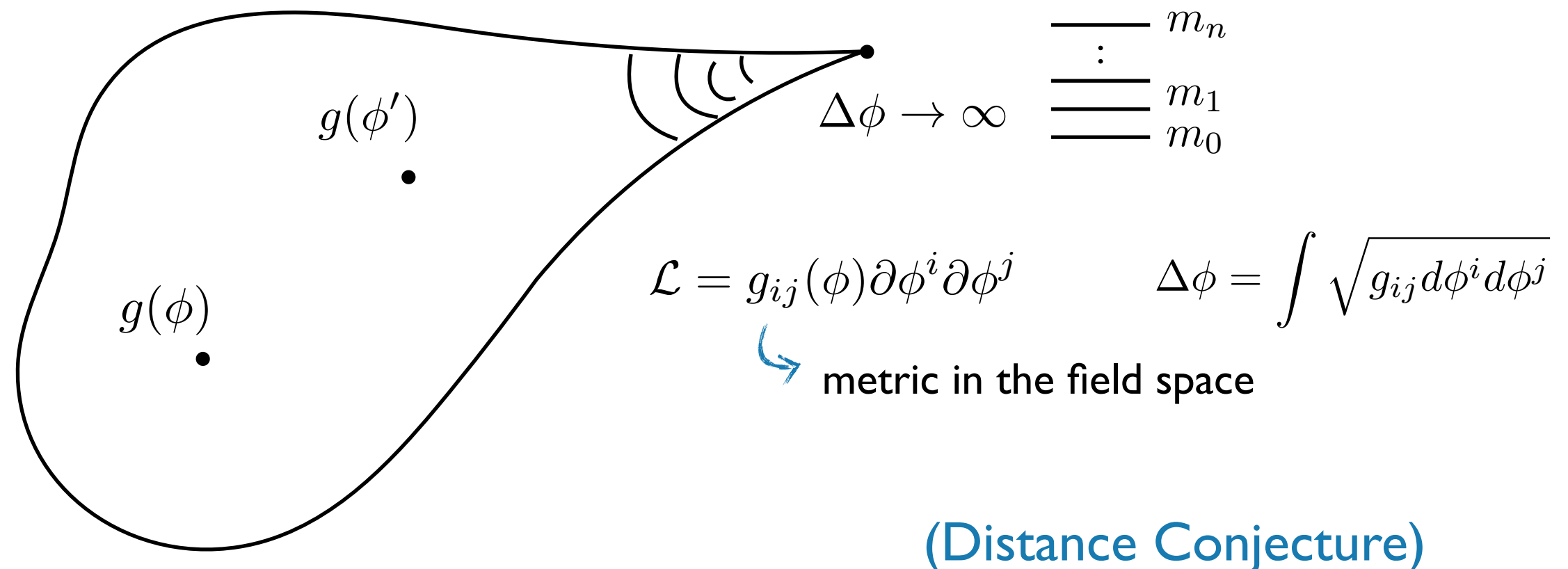
Global symmetries, vanishing couplings or large hierarchies occur at infinite distance in this field space



3) Dynamical parameters and Distance Conjecture

Universal features:

Global symmetries, vanishing couplings or large hierarchies occur at infinite distance in this field space



There is an **infinite tower of states** becoming **exponentially light** at every **infinite field distance** limit [Ooguri-Vafa'06]

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty$$

3) Dynamical parameters and Distance Conjecture

In all string theory models that are known, the tower is (in some duality frame) either a: [\[Lee,Lerche,Weigand'19\]](#)

- Kaluza-Klein tower → decompactification of extra dimensions
- Oscillator modes of a string → perturbative string limit

3) Dynamical parameters and Distance Conjecture

In all string theory models that are known, the tower is (in some duality frame) either a: [Lee,Lerche,Weigand'19]

- Kaluza-Klein tower → decompactification of extra dimensions
- Oscillator modes of a string → perturbative string limit

Example: One large extra dimension (from (d+1)- to d-dimensions)

$$\text{KK tower: } m_{\text{KK}} = \frac{M_{\text{pl},d+1}}{r} = M_{\text{pl},d} \exp\left(-\sqrt{\frac{d-1}{d-2}} \Delta\phi\right)$$

$$S = \frac{M_{\text{pl},d}^{d-2}}{2} \int d^d x \left(R - \frac{1}{2} \sqrt{\frac{d-1}{d-2}} \frac{(\partial r)^2}{r^2} + \dots \right) = \frac{M_{\text{pl},d}^{d-2}}{2} \int d^d x \left(R - \frac{1}{2} (\partial\phi)^2 + \dots \right)$$

As $r \rightarrow \infty$ KK tower becomes exponentially light in terms of canonically normalised radius

3) Dynamical parameters and Distance Conjecture

Take-home message:

Quantum gravity consistency requires this behaviour of Kaluza-Klein theory to be generic whenever any scalar field vev gets large

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty$$

depends on nature of the tower

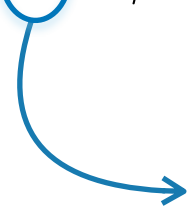
$$\frac{1}{\sqrt{2}} \leq \alpha \leq \sqrt{\frac{3}{2}} \quad \text{in 4d} \quad [\text{Etheredge et al'22}]$$

3) Dynamical parameters and Distance Conjecture

Take-home message:

Quantum gravity consistency requires this behaviour of Kaluza-Klein theory to be generic whenever any scalar field vev gets large

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty$$


 depends on nature of the tower

$$\frac{1}{\sqrt{2}} \leq \alpha \leq \sqrt{\frac{3}{2}} \quad \text{in 4d} \quad [\text{Etheredge et al'22}]$$

QFT breaks down at a cut-off scale known as the species scale Λ [Dvali'07]

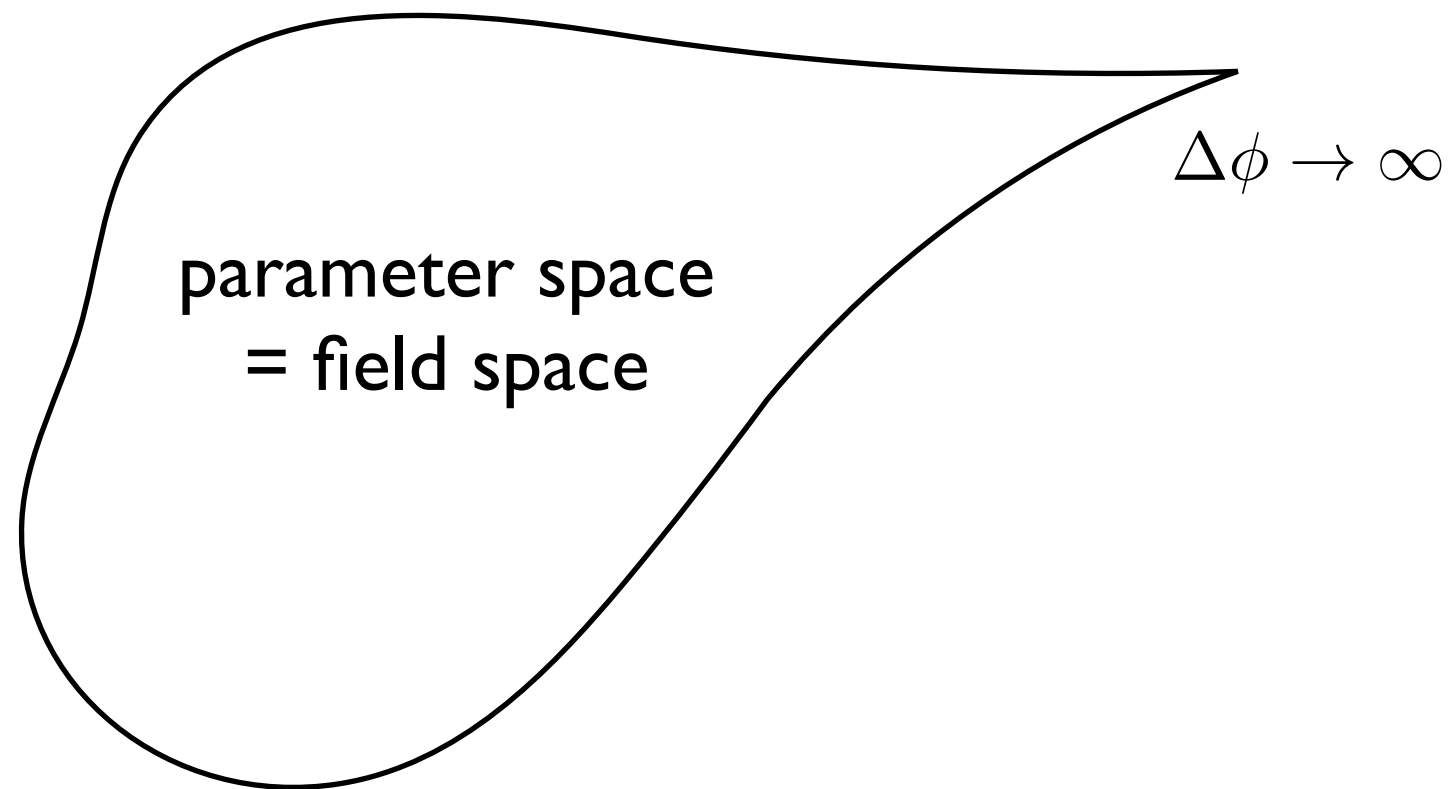
If KK tower: $\Lambda = M_{\text{pl},D} \ll M_{\text{pl},d}$

If string tower: $\Lambda \simeq M_{\text{str}} \ll M_{\text{pl},d}$

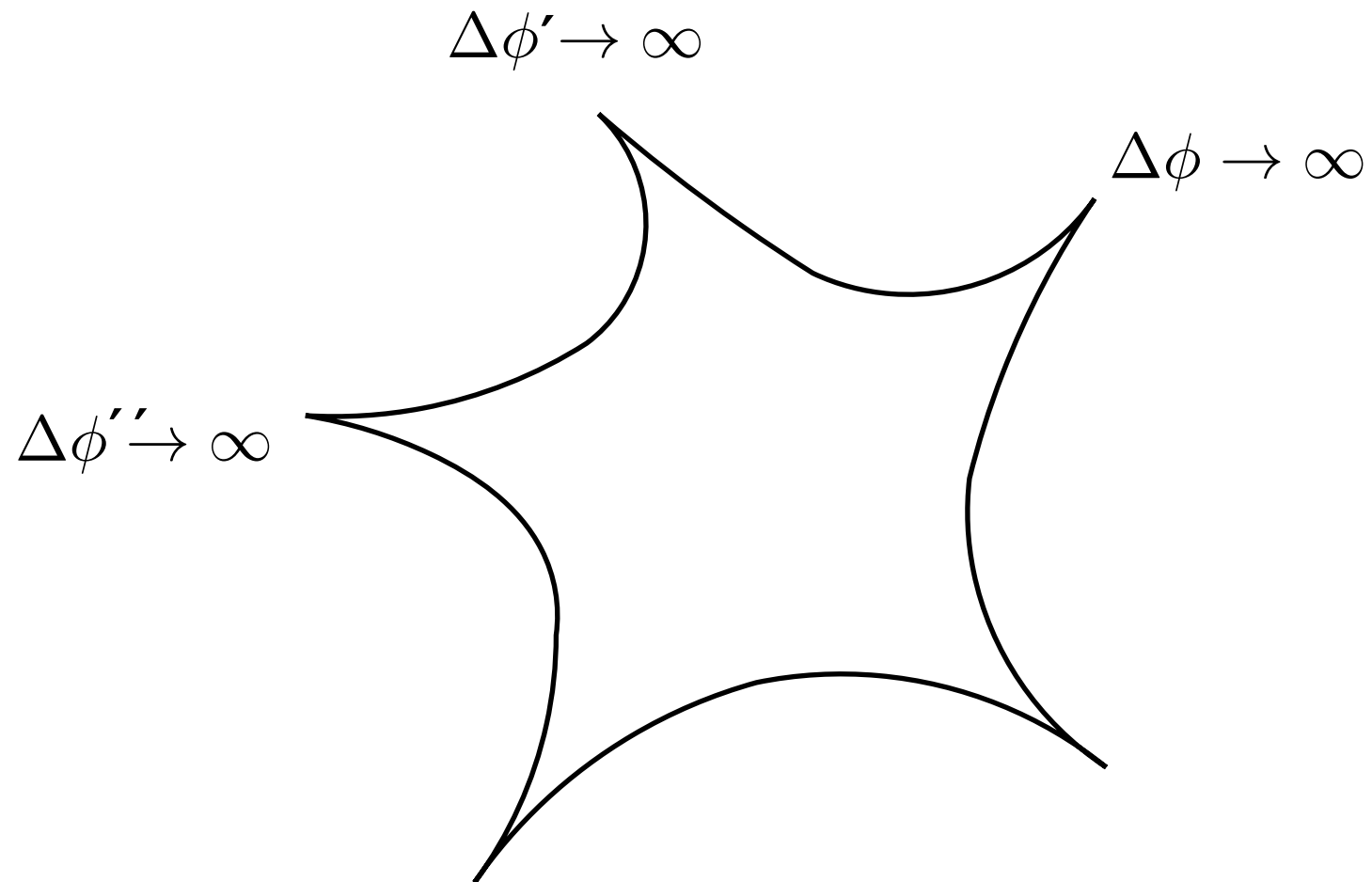
$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{2}$$

[Castellano,Ruiz,IV'23]

3) Dynamical parameters and Distance Conjecture

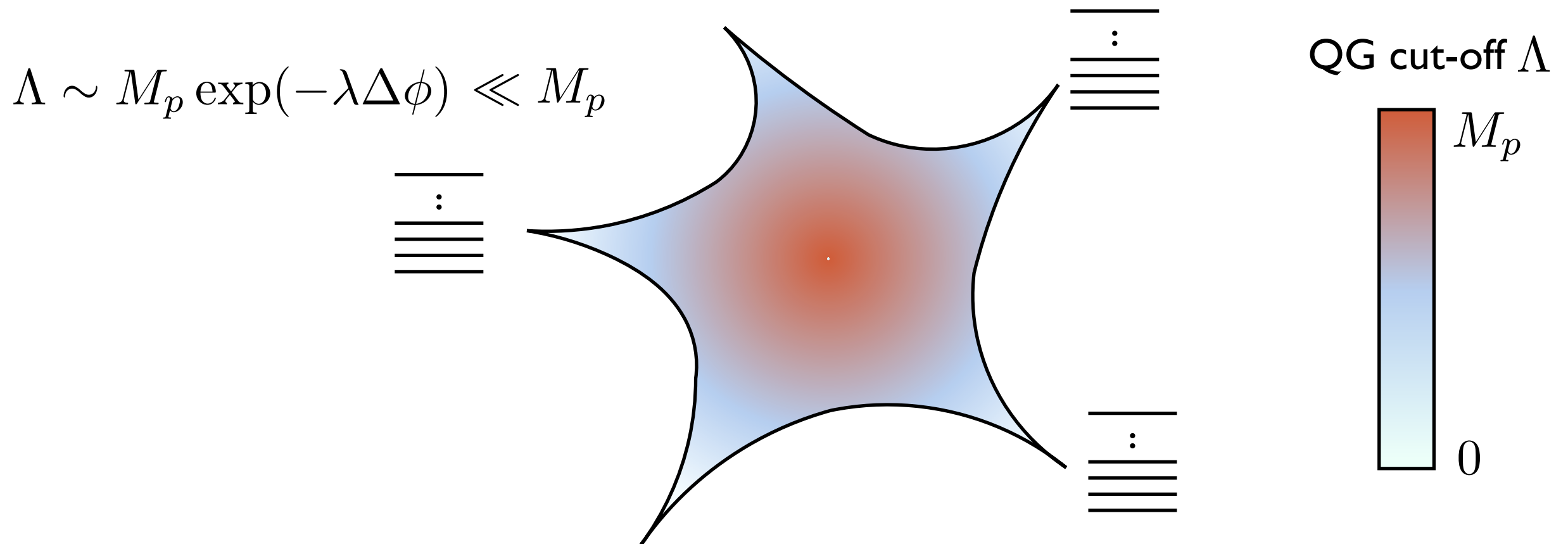


3) Dynamical parameters and Distance Conjecture



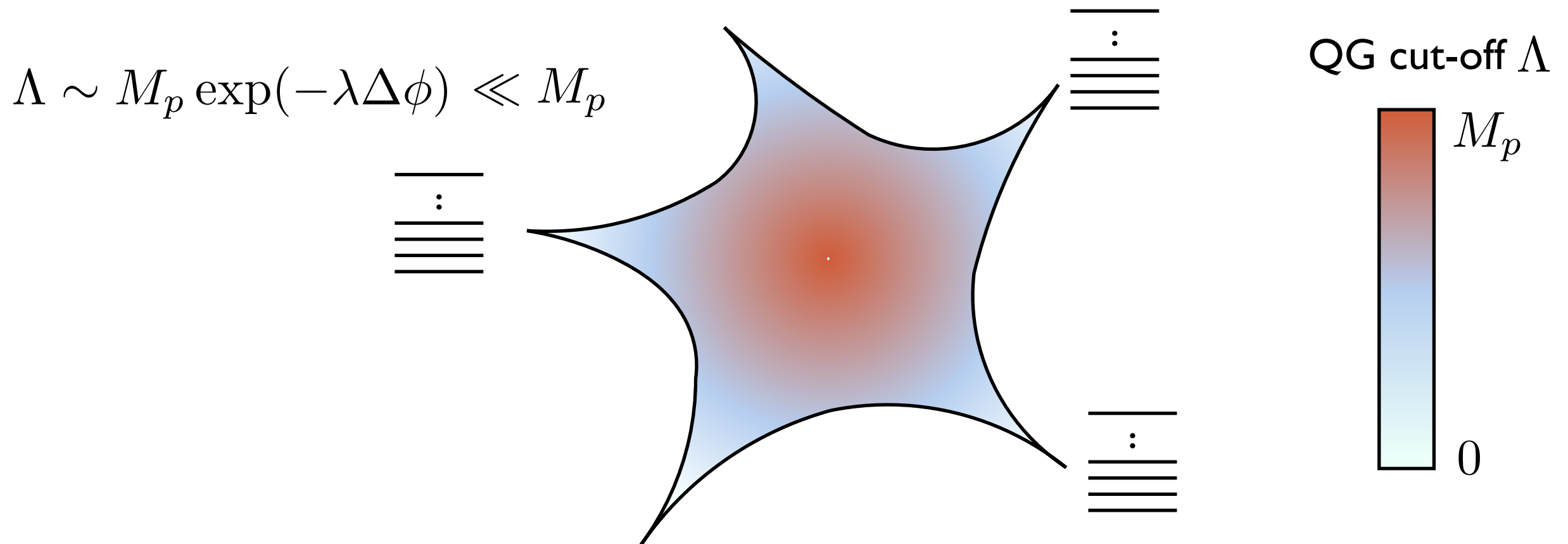
3) Dynamical parameters and Distance Conjecture

The cut-off Λ is field dependent and becomes parametrically smaller than M_p at the infinite field distance boundaries



3) Dynamical parameters and Distance Conjecture

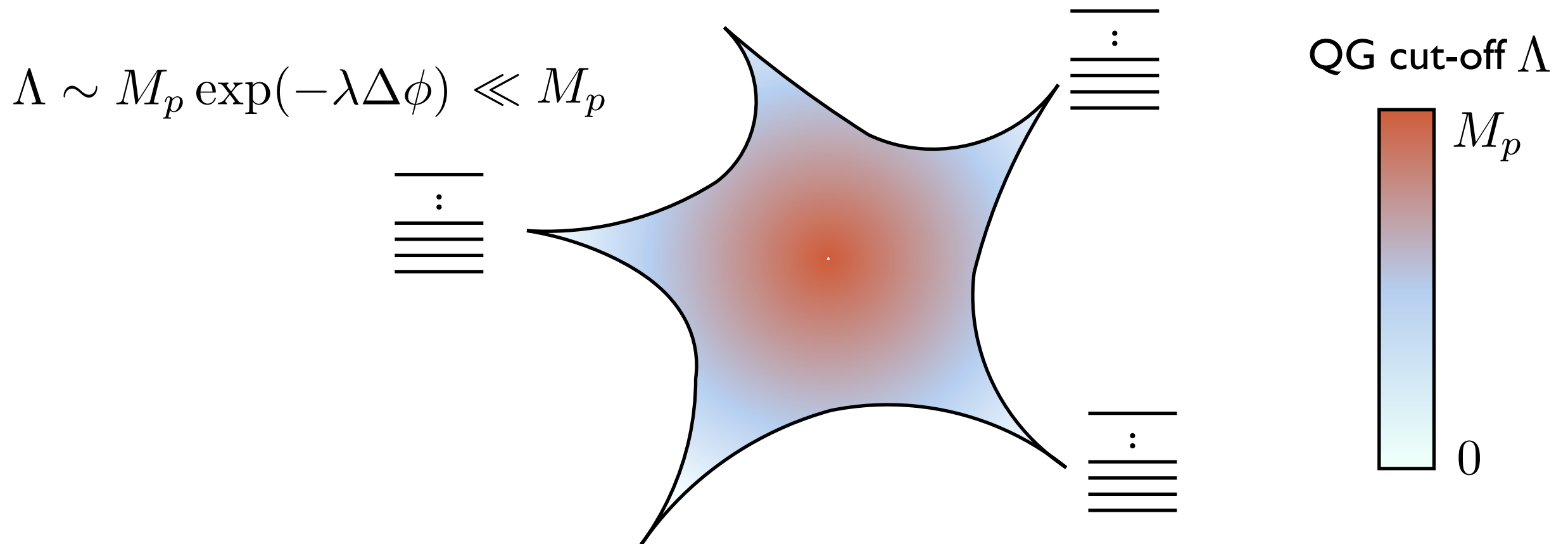
The cut-off Λ is field dependent and becomes parametrically smaller than M_p at the infinite field distance boundaries



$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}}[g(\phi), m(\phi), \dots] + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda^{n-4}(\phi)} \right)$$

3) Dynamical parameters and Distance Conjecture

The cut-off Λ is field dependent and becomes parametrically smaller than M_p at the infinite field distance boundaries

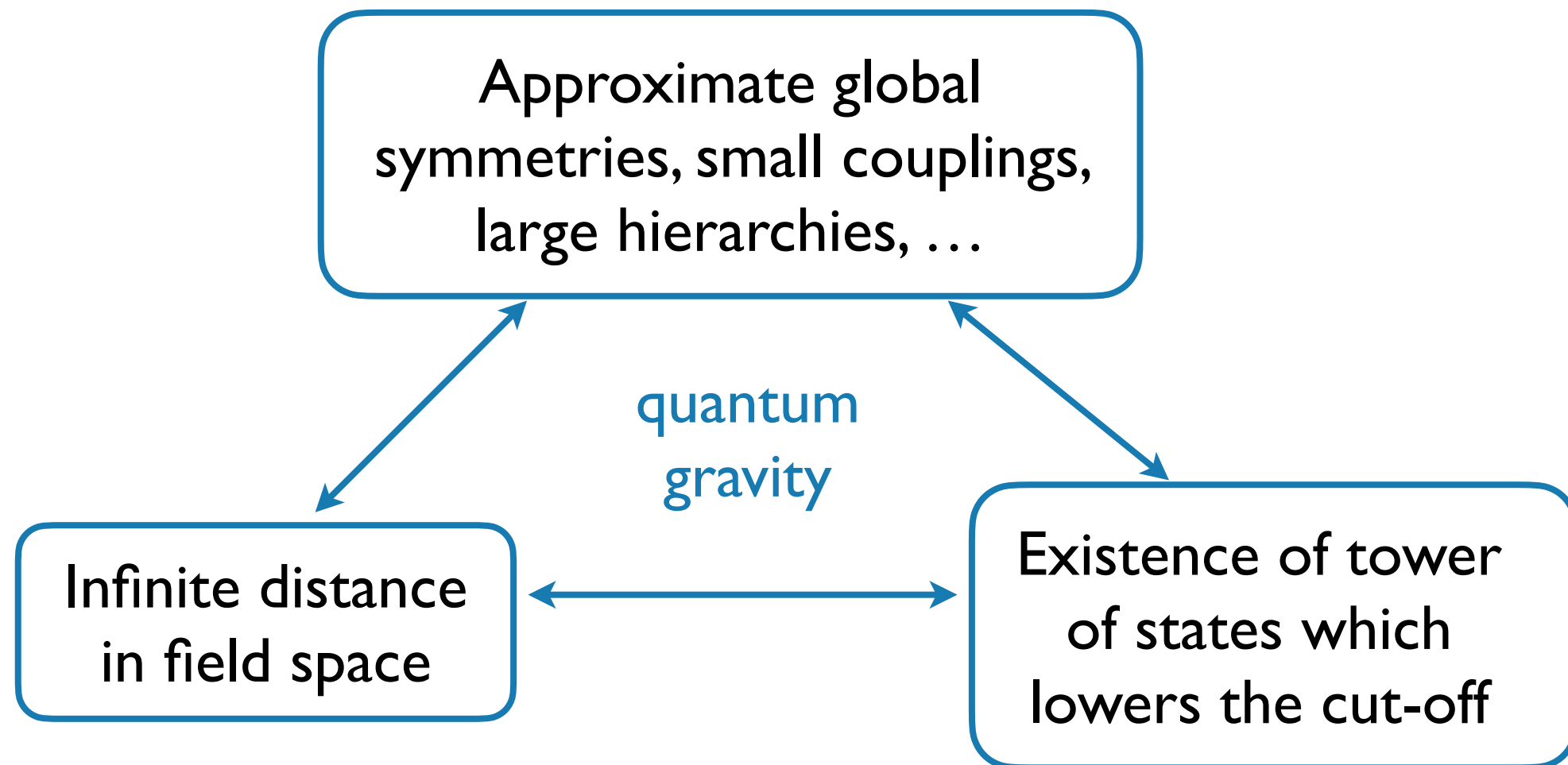


$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}}[g(\phi), m(\phi), \dots] + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda^{n-4}(\phi)} \right)$$

given this \nearrow

\nwarrow we can determine this (sometimes)

3) Dynamical parameters and Distance Conjecture



3) **Dynamical parameters and Distance Conjecture**

Implications:

3) Dynamical parameters and Distance Conjecture

Implications:

I) **Bound on the maximum scalar field range** that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi) \quad \rightarrow \quad \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

Constraints on large field inflation, cosmological relaxation, ...

[Ibanez et al'15] [Scalisi,IV18] , ...

3) Dynamical parameters and Distance Conjecture

Implications:

1) **Bound on the maximum scalar field range** that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi) \quad \rightarrow \quad \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

Constraints on large field inflation, cosmological relaxation, ...

[Ibanez et al'15] [Scalisi,IV18] , ...

2) **Bound on how small gauge couplings** can be $\Lambda \lesssim g^k M_p \quad 1/3 \leq k \leq 1$

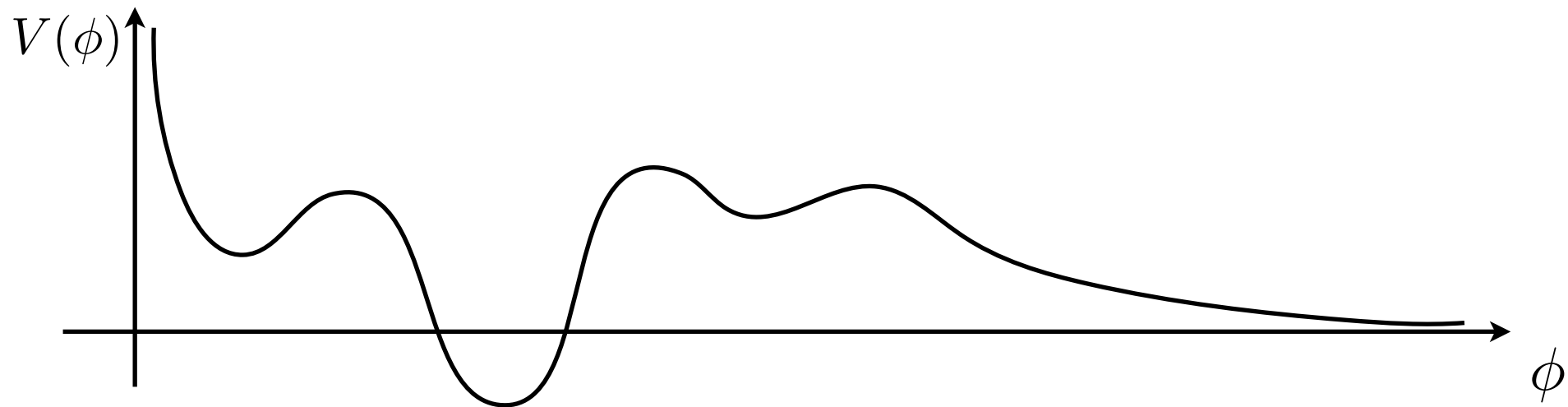
Constraints on weakly coupled dark photons, Stuckelberg couplings, ...

[Reece'18] [Montero,Muñoz,Obied'22] , ...

4) Small vacuum energy and dark dimension

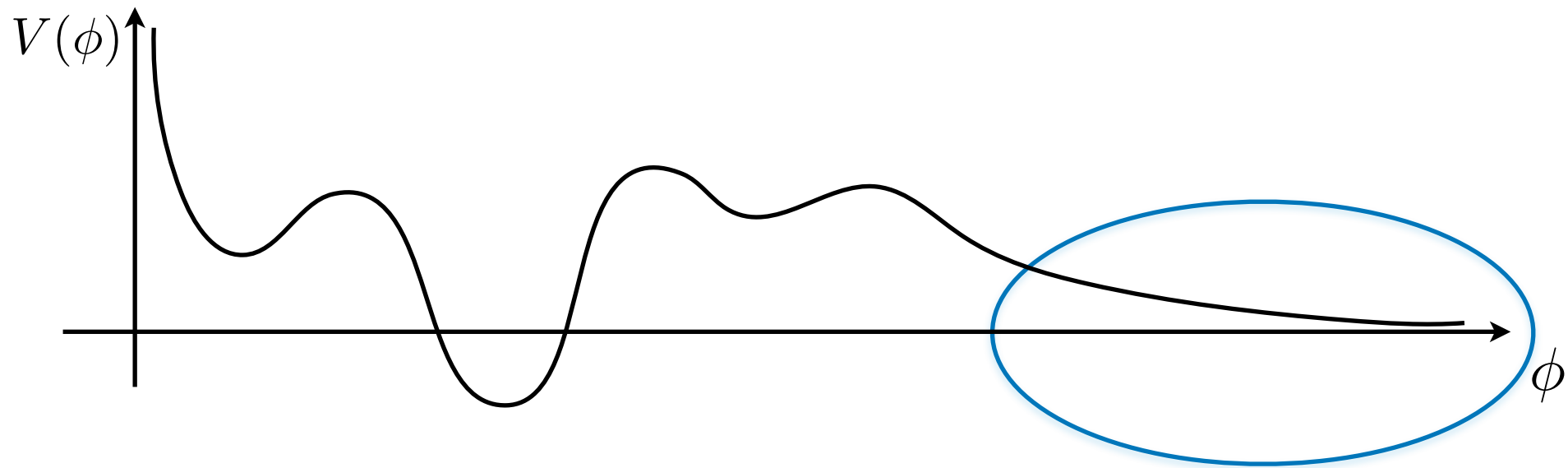
4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically (at infinite field distance)



4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically
(at infinite field distance)

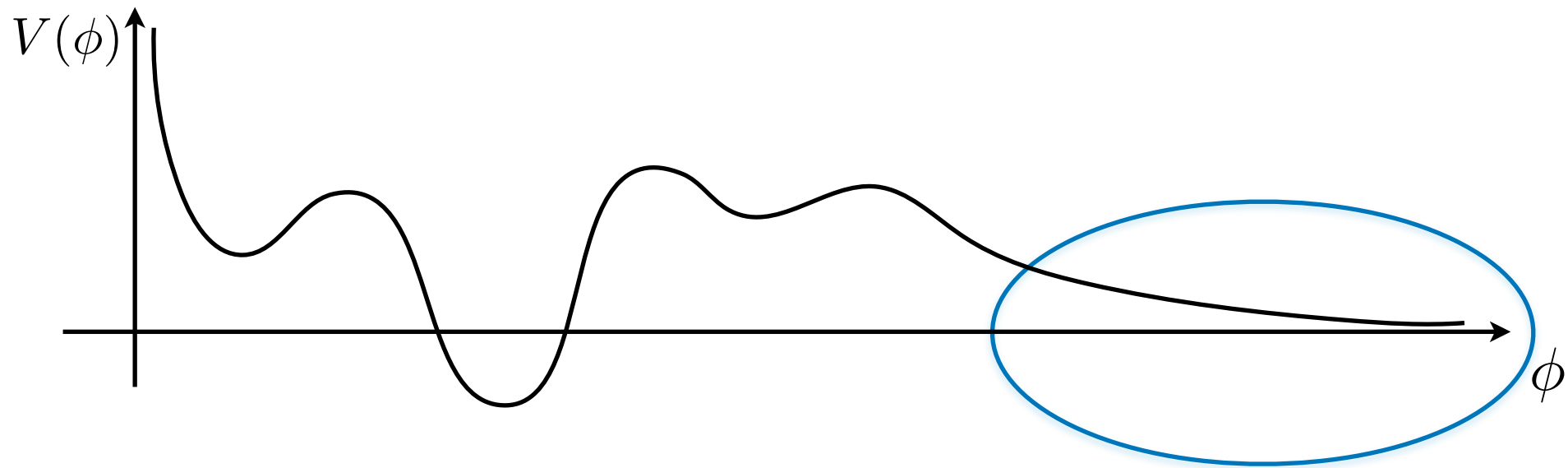


how come? it looks “unnatural”

string theory contains infinitely many massive states!

4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically
(at infinite field distance)



regardless of supersymmetry

how come? it looks “unnatural”

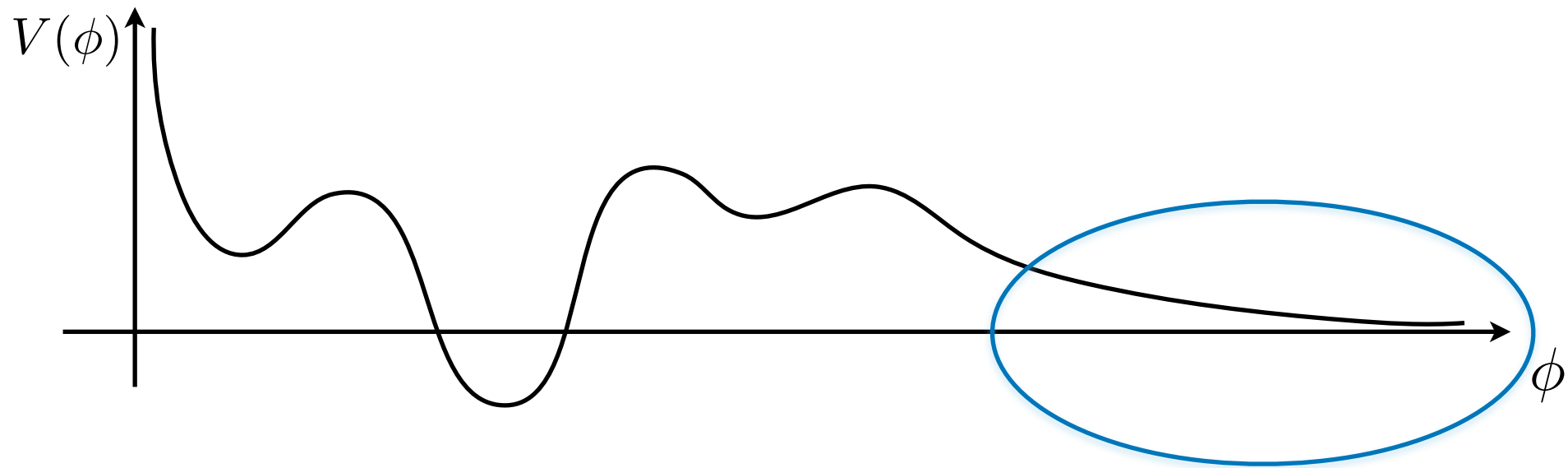
string theory contains infinitely many massive states!

Asymptotically (at parametrically late times) the universe
becomes nearly flat (Minkowski)

[Obied et al'18] [Dvali's work]

4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically
(at infinite field distance)



how come? it looks “unnatural”

regardless of supersymmetry

string theory contains infinitely many massive states!

e.g. Non-SUSY heterotic string theory

$$V = \cancel{V_{\text{tree}}} + V_{\text{one loop}} + \dots \sim \exp(-\gamma\phi)$$

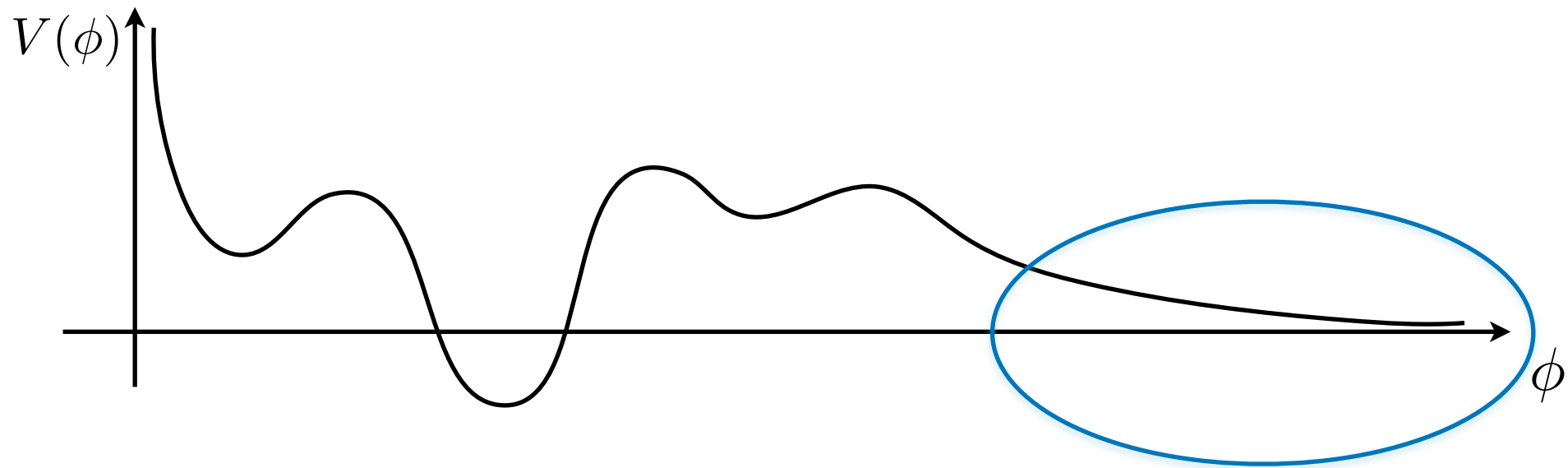
0

(conformal invariance of
the string worldsheet)

(“Casimir energy”
of string states)

4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically
(at infinite field distance)



regardless of supersymmetry

how come? it looks “unnatural”

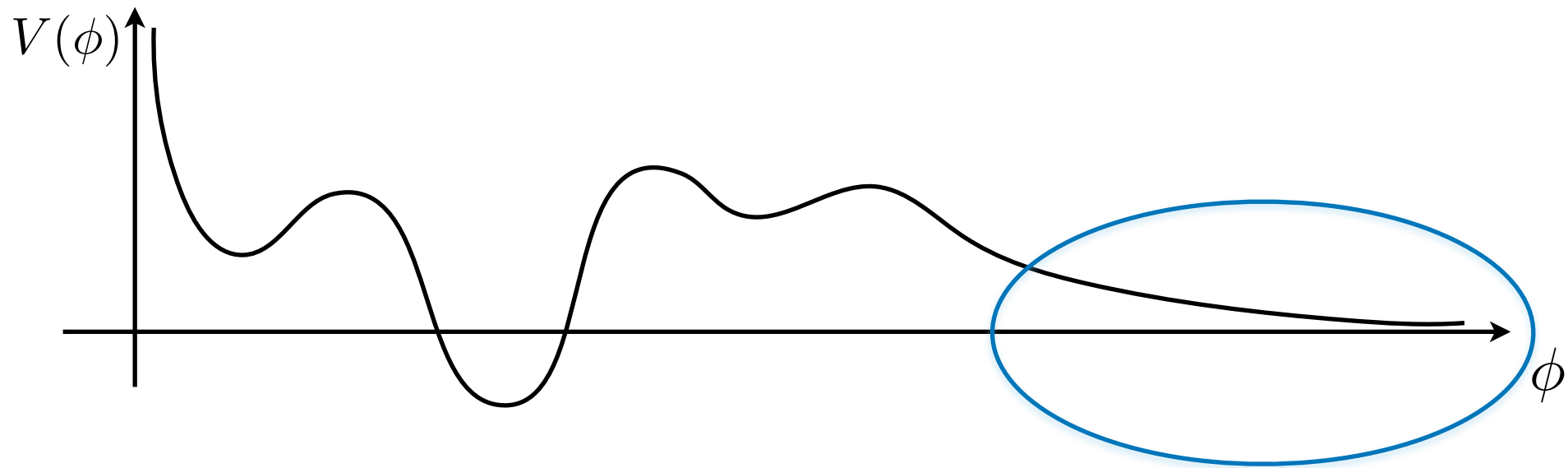
string theory contains infinitely many massive states!

It is crucial to integrate out the infinite towers of states to obtain a finite result that it is smaller than expected if you only considered quantum corrections from a finite number of fields

[see Steven's talk]

4) Small vacuum energy and dark dimension

Vacuum energy (if positive and below Λ) vanishes asymptotically
(at infinite field distance)

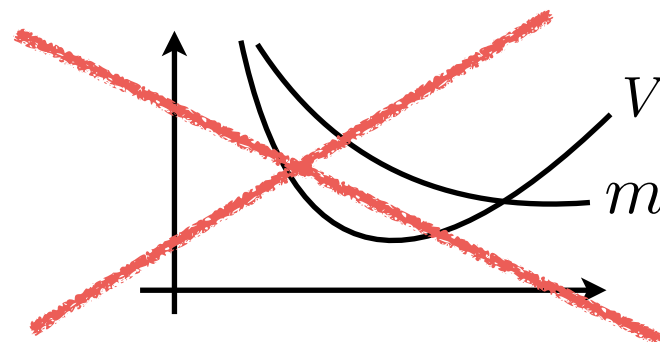
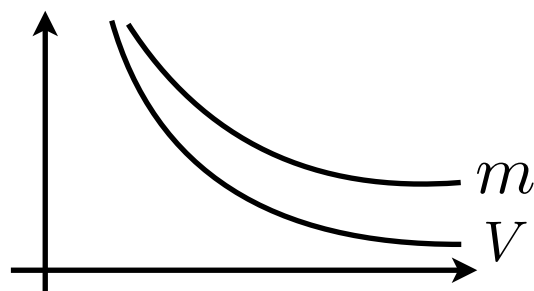


how come? it looks “unnatural”

regardless of supersymmetry

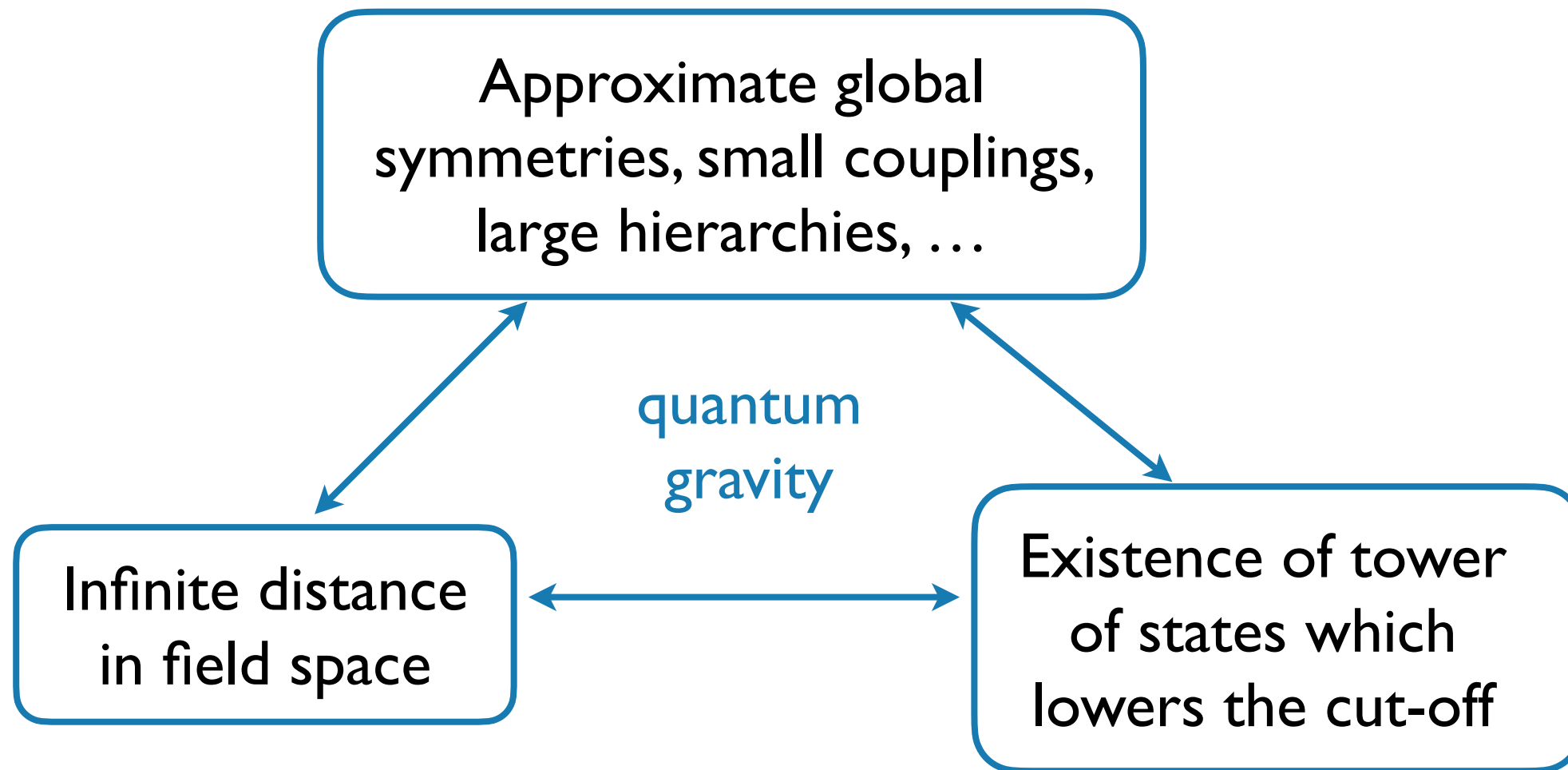
string theory contains infinitely many massive states!

This is a universal result in string theory and can also be motivated by the Swampland Distance conjecture (i.e. the existence of the tower of states)

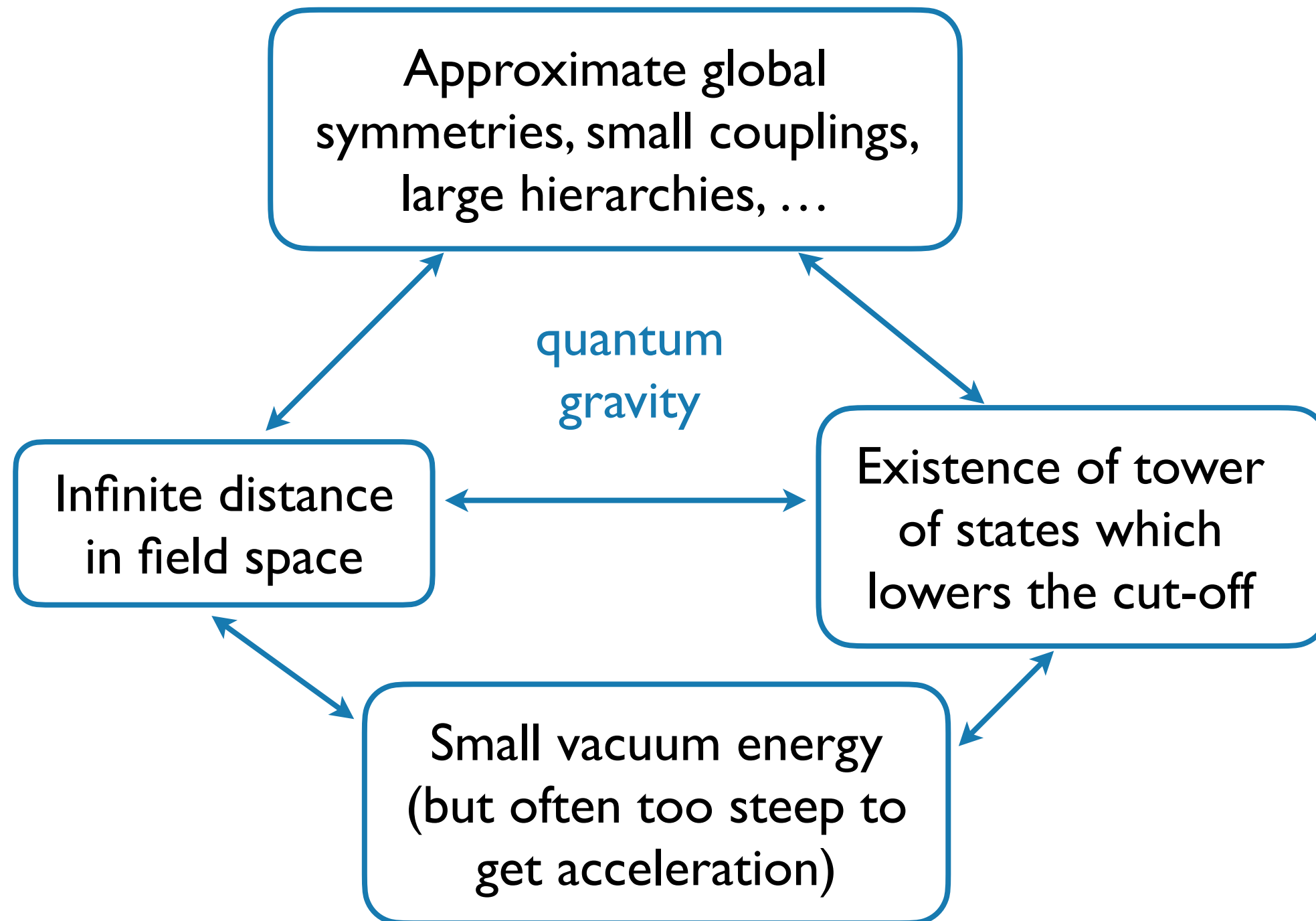


[Ooguri et al'18]
[Montero, Vafa, IV'22]

4) Small vacuum energy and dark dimension

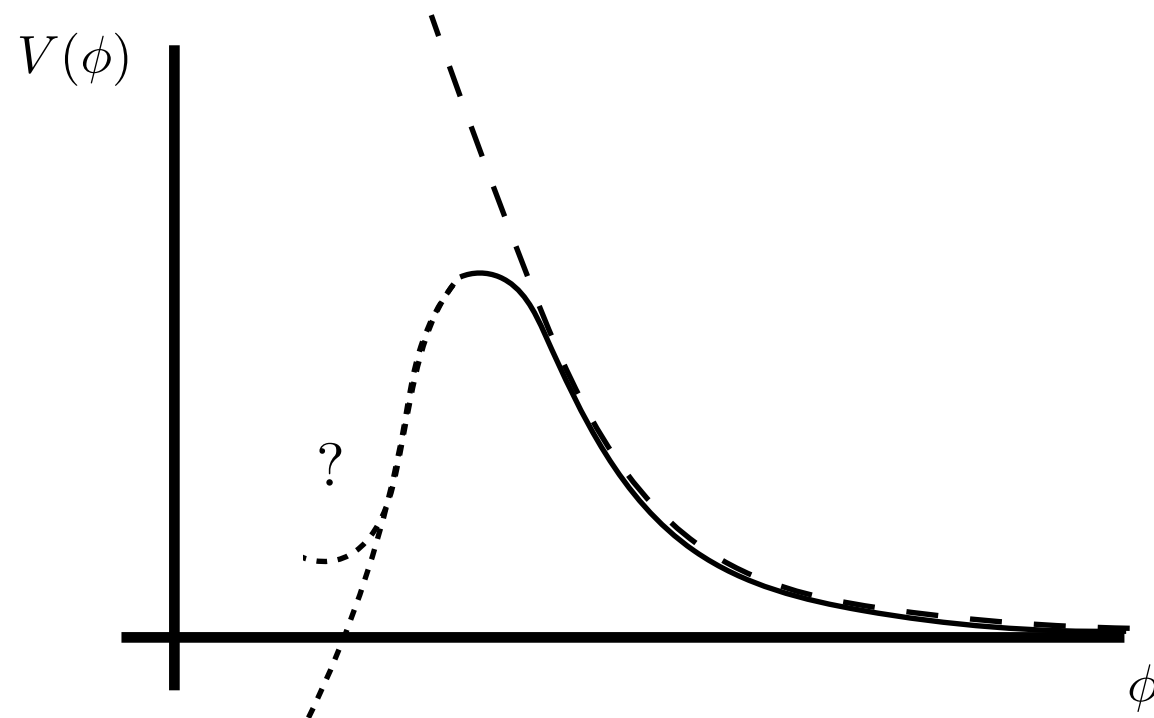


4) Small vacuum energy and dark dimension



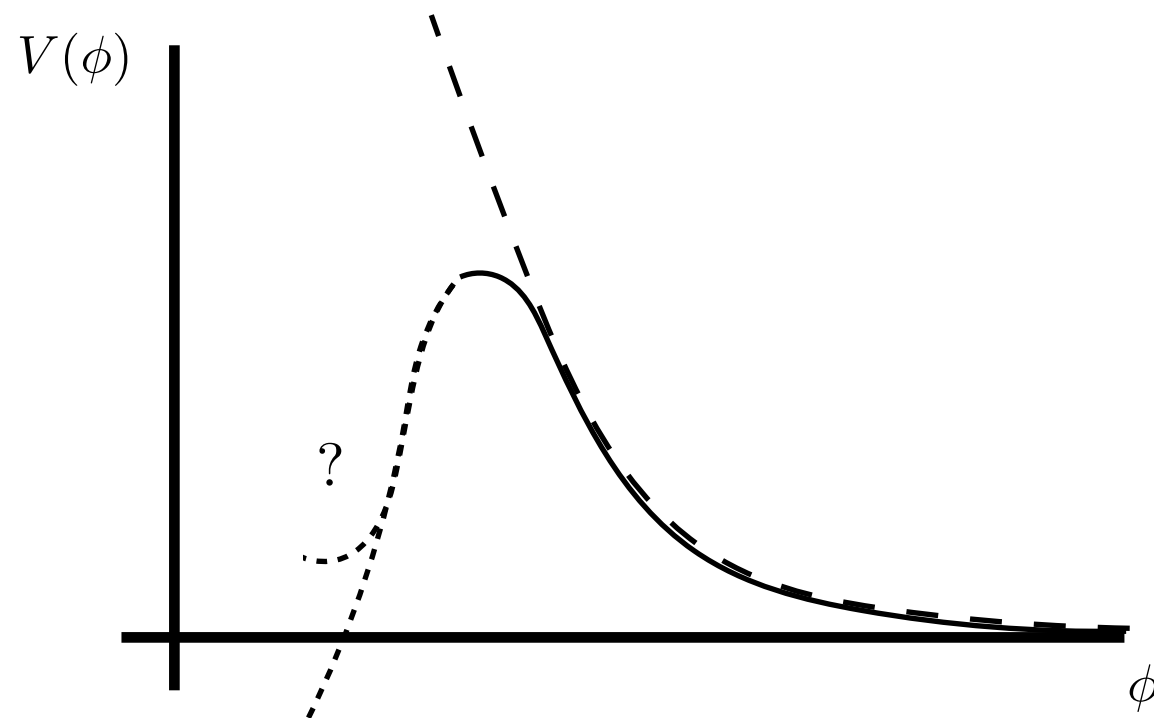
4) Small vacuum energy and dark dimension

Could it be that the **dark energy** (vacuum energy) in our universe is **small**



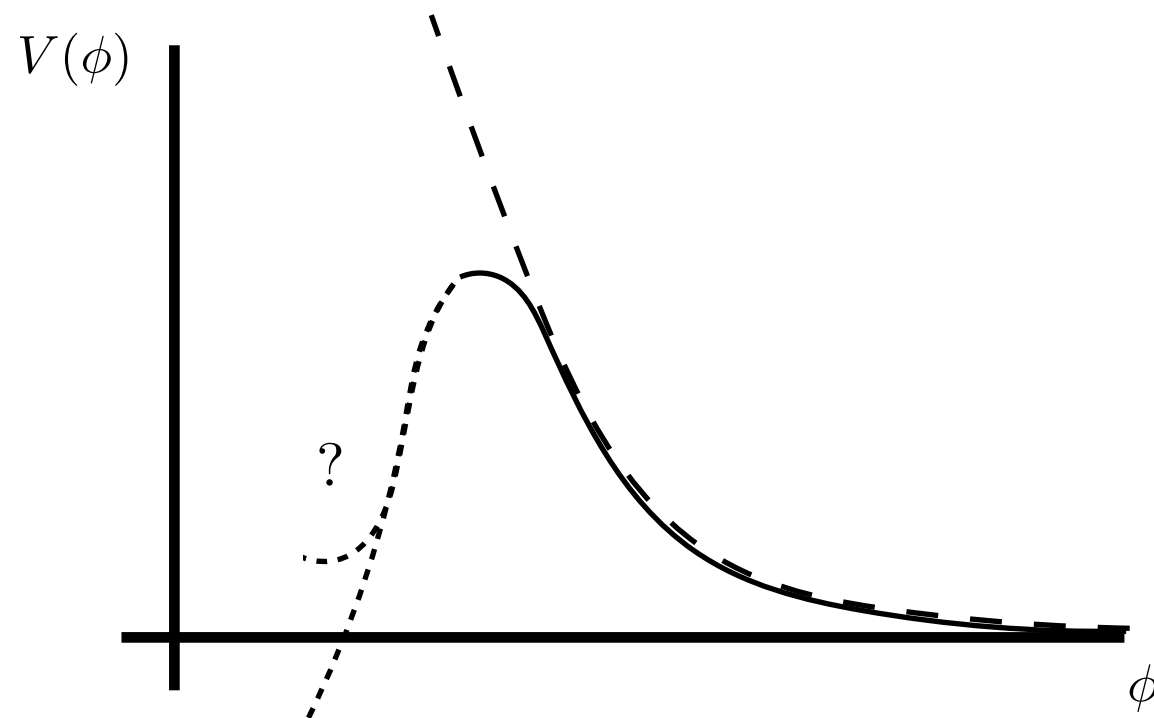
4) Small vacuum energy and dark dimension

Could it be that the **dark energy** (vacuum energy) in our universe is **small** not because of a huge fine-tuning of many contributions



4) Small vacuum energy and dark dimension

Could it be that the **dark energy** (vacuum energy) in our universe is **small** not because of a huge fine-tuning of many contributions but **because we live near and asymptotic limit** where it naturally goes to zero?

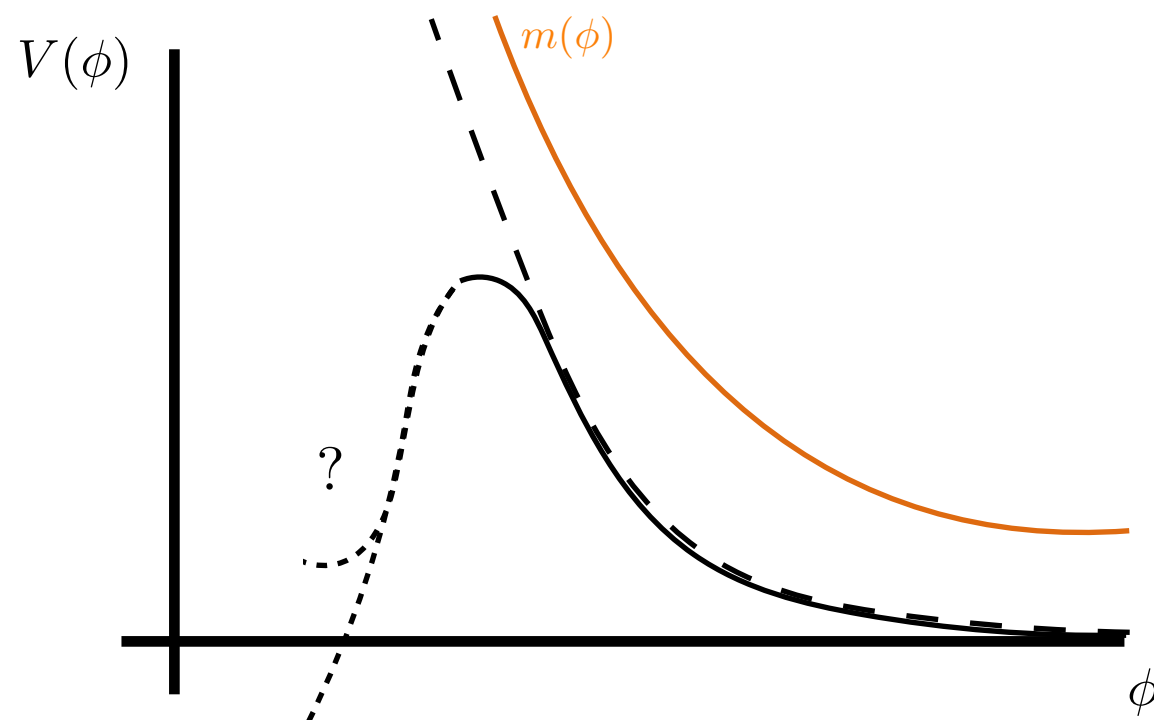


4) Small vacuum energy and dark dimension

Could it be that the **dark energy** (vacuum energy) in our universe is **small** not because of a huge fine-tuning of many contributions but **because we live near and asymptotic limit** where it naturally goes to zero?

If so, one thing is clear:

There should be a tower of states becoming light



4) Small vacuum energy and dark dimension

Combining:

Theoretical bounds

$$V^{1/2} \lesssim m \lesssim V^{1/4} \quad +$$

Higuchi bound   no fine-tuning of
quantum corrections

Experimental bounds on
deviations of Newton's law
and astrophysical bounds

4) Small vacuum energy and dark dimension

Combining:

Theoretical bounds

$$V^{1/2} \lesssim m \lesssim V^{1/4} \quad +$$

Higuchi bound  no fine-tuning of quantum corrections 

Experimental bounds on deviations of Newton's law and astrophysical bounds

The first state of the tower should have a mass:

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$



neutrino scale!

Tower of right handed neutrinos?

4) Small vacuum energy and dark dimension

Combining:

Theoretical bounds

$$V^{1/2} \lesssim m \lesssim V^{1/4} \quad +$$

Higuchi bound  no fine-tuning of quantum corrections 

Experimental bounds on deviations of Newton's law and astrophysical bounds

The first state of the tower should have a mass:

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$

 neutrino scale!

Tower of right handed neutrinos?

implying one large extra dimension $l \sim 0.1 - 10\mu m$

The Dark Dimension [Montero,Vafa,IV'22]

currently testing it in
New ISLE at the
Conrad Observatory

[Aspelmeyer,Adelberger,
Shayeghi,Zito...]

There is another instance in which a tower with $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$ seems necessary to comply with Swampland constraints

There is another instance in which a tower with $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$ seems necessary to comply with Swampland constraints

Take Standard Model + gravity and compactify to 3d on a circle:

$$V(R) = \frac{2\pi\Lambda_4}{R^2} - \frac{4}{720\pi R^6} + \sum_i \frac{(2\pi R)}{R^3} (-1)^{s_i} n_i \rho_i(R)$$

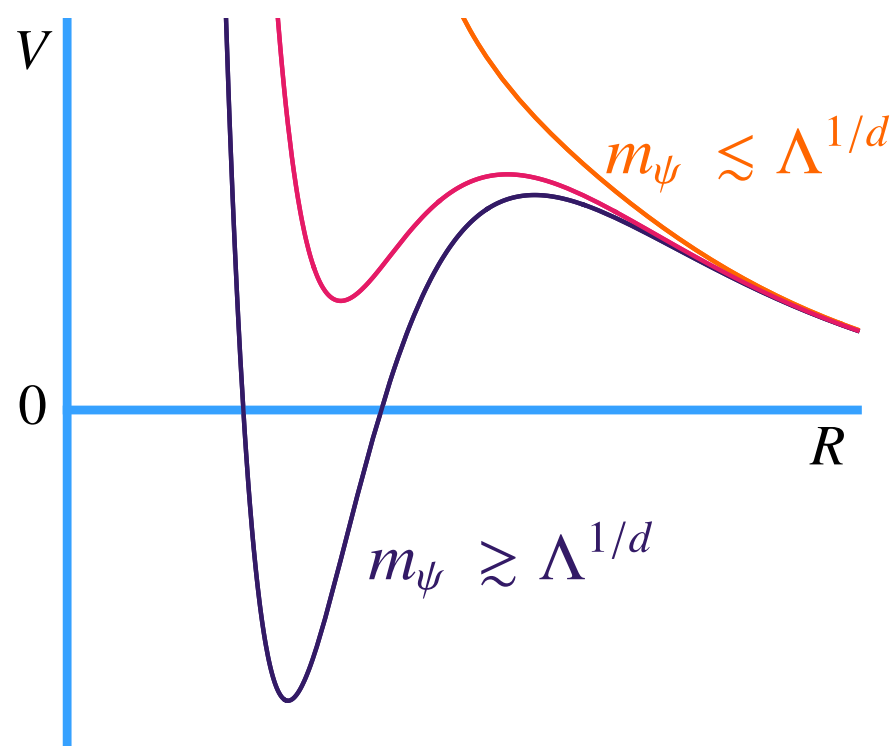
↑ graviton, photon
↑ massive particles:
neutrinos,...

There is another instance in which a tower with $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$ seems necessary to comply with Swampland constraints

Take Standard Model + gravity and compactify to 3d on a circle:

$$V(R) = \frac{2\pi\Lambda_4}{R^2} - \frac{4}{720\pi R^6} + \sum_i \frac{(2\pi R)}{R^3} (-1)^{s_i} n_i \rho_i(R)$$

↑ graviton, photon
↑ massive particles: neutrinos, ...



Depending on the value of Dirac neutrino masses, we get 3-dimensional AdS, Minkowski or dS vacua.

[Arkani-Hamed et al'07]

5) AdS SM vacua and neutrino masses

[Martin-Lozano,Ibanez,IV'17] [Gonzalo et al'18-21]

These AdS vacua are in tension with certain Swampland constraints
(like the AdS Distance conjecture) [Luest,Palti,Vafa'19]

[Martin-Lozano,Ibanez,IV'17] [Gonzalo et al'18-21]

These AdS vacua are in tension with certain Swampland constraints
(like the AdS Distance conjecture) [Luest,Palti,Vafa'19]

- 1) There are more than two Dirac fermions with $m \lesssim \Lambda_{cc}^{1/4}$ Neutrinos!!!
- 2) There is an infinite tower of states in 4d scaling as $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$

[Martin-Lozano,Ibanez,IV'17] [Gonzalo et al'18-21]

These AdS vacua are in tension with certain Swampland constraints
(like the AdS Distance conjecture) [Luest,Palti,Vafa'19]

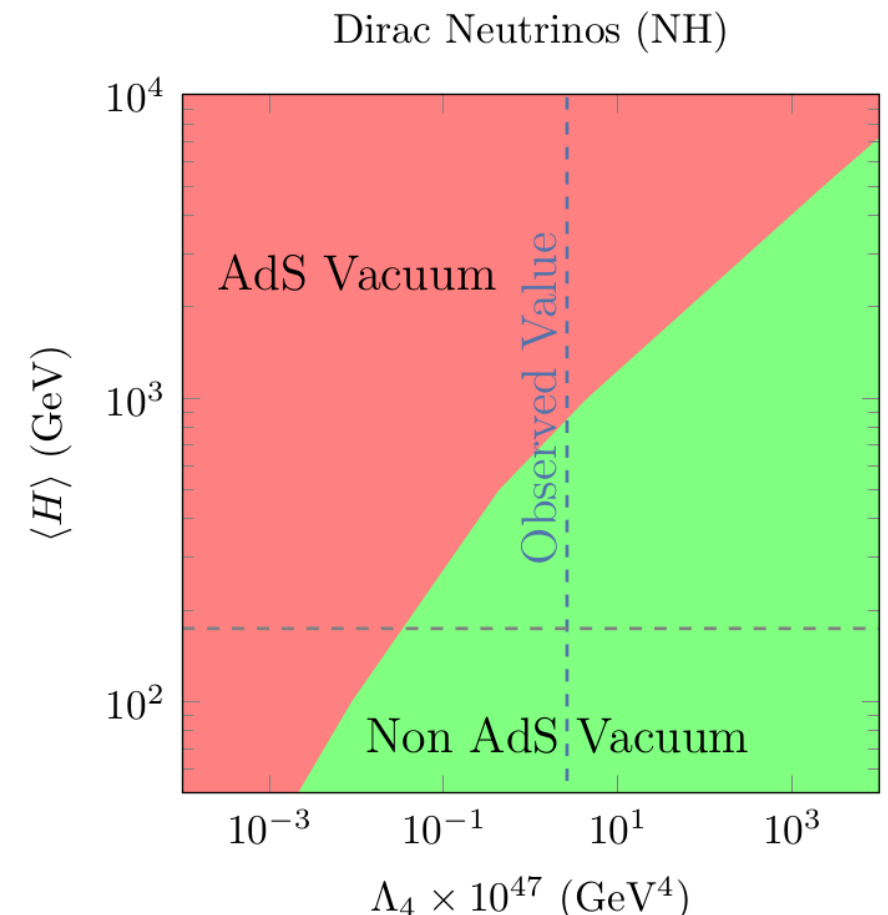
1) There are more than two Dirac fermions with $m \lesssim \Lambda_{cc}^{1/4}$ **Neutrinos!!!**

2) There is an infinite tower of states in 4d scaling as $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$

We can translate the bound on neutrino masses to an upper bound on the EW scale in terms of the cosmological constant:

$$\langle H \rangle \lesssim 1.6 \frac{\Lambda^{1/4}}{Y_{\nu_1}}$$

Parameters leading to a higher EW scale would not yield theories consistent with quantum gravity



[Martin-Lozano,Ibanez,IV'17] [Gonzalo et al'18-21]

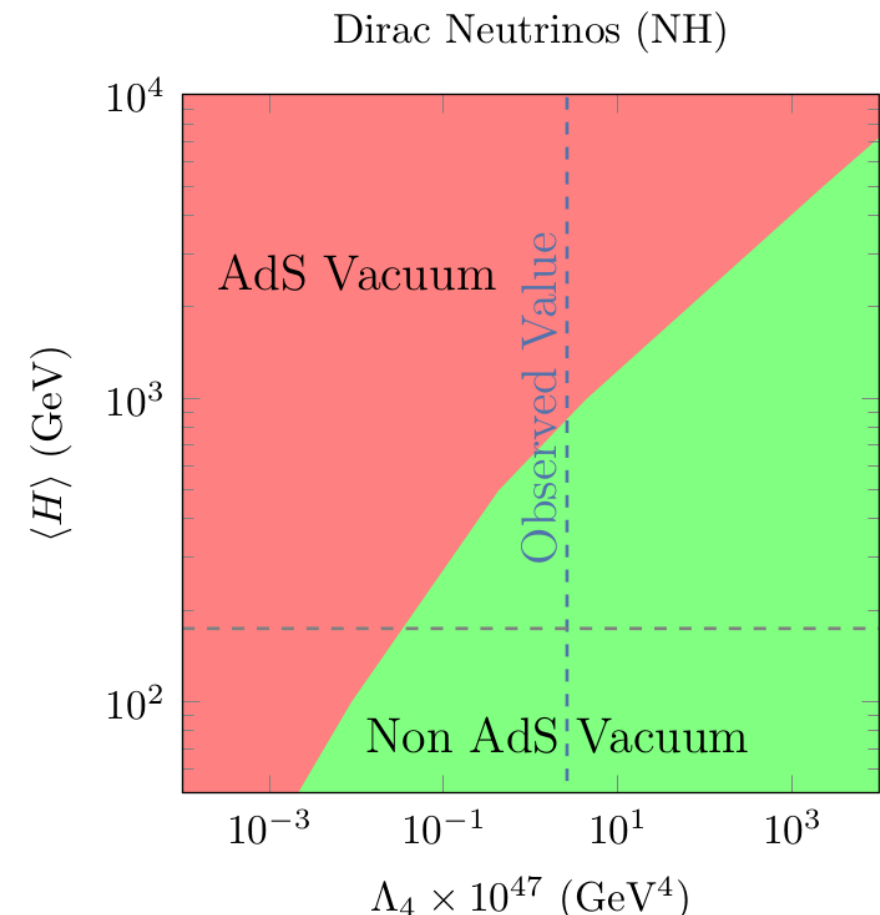
These AdS vacua are in tension with certain Swampland constraints
(like the AdS Distance conjecture) [Luest,Palti,Vafa'19]

- 1) There are more than two Dirac fermions with $m \lesssim \Lambda_{cc}^{1/4}$ **Neutrinos!!!**
- 2) There is an infinite tower of states in 4d scaling as $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$

We can translate the bound on neutrino masses to an upper bound on the EW scale in terms of the cosmological constant:

$$\langle H \rangle \lesssim 1.6 \frac{\Lambda^{1/4}}{Y_{\nu_1}} \quad \text{Parameters leading to a higher EW scale would not yield theories consistent with quantum gravity}$$

Potential loopholes. But proof of principle that space of parameters consistent with quantum gravity can be smaller than expected, not everything goes!



CONCLUSIONS

Requiring quantum gravity UV consistency can have implications at low energies and potentially shed new light into naturalness problems.

- We saw examples of Swampland constraints that imply ‘surprising’ or ‘unnatural’ results from an IR perspective, but they are natural from the UV complete theory
- None of them clearly solves the concrete naturalness issues of our universe (EW hierarchy problem, cosmological constant problem, etc.)
but they provide interesting avenues to pursue.
- This is only the beginning. What other guiding principles for BSM can we learn from quantum gravity?

Can we make the UV/IR mixing more manifest?

Thank you!

Thank you!

Online series of Swampland seminars / open mic discussions
on Tuesdays at 11 am ET (5 pm CET)

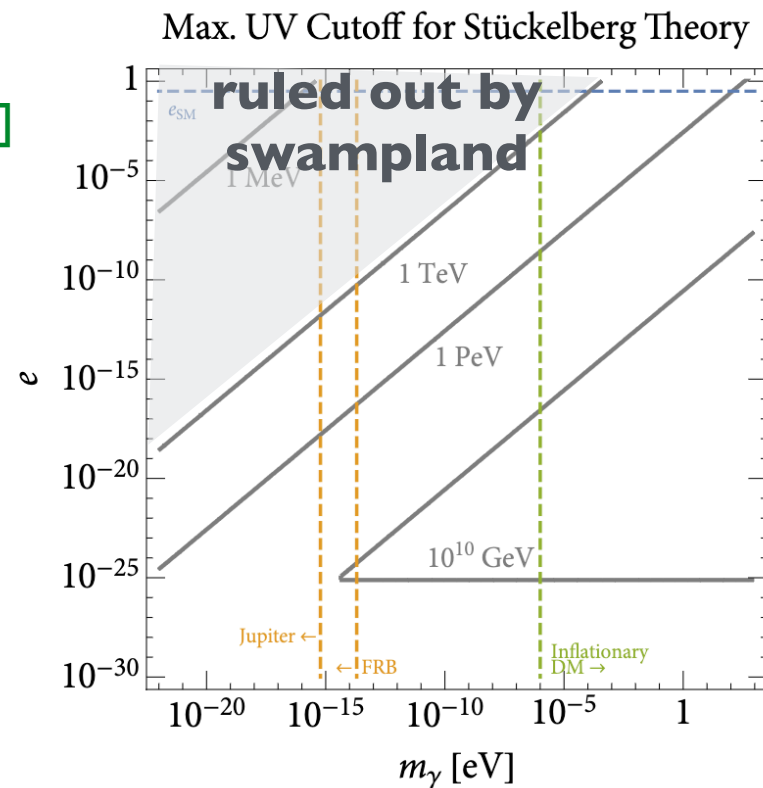
You can subscribe here: <https://web.lists.fas.harvard.edu/mailman/lists/hetg-swampland.lists.fas.harvard.edu/>

Everybody is welcome! :)

back-up slides

How does the tower/cut-off behaves in terms of EFT data?

[Reece'18]



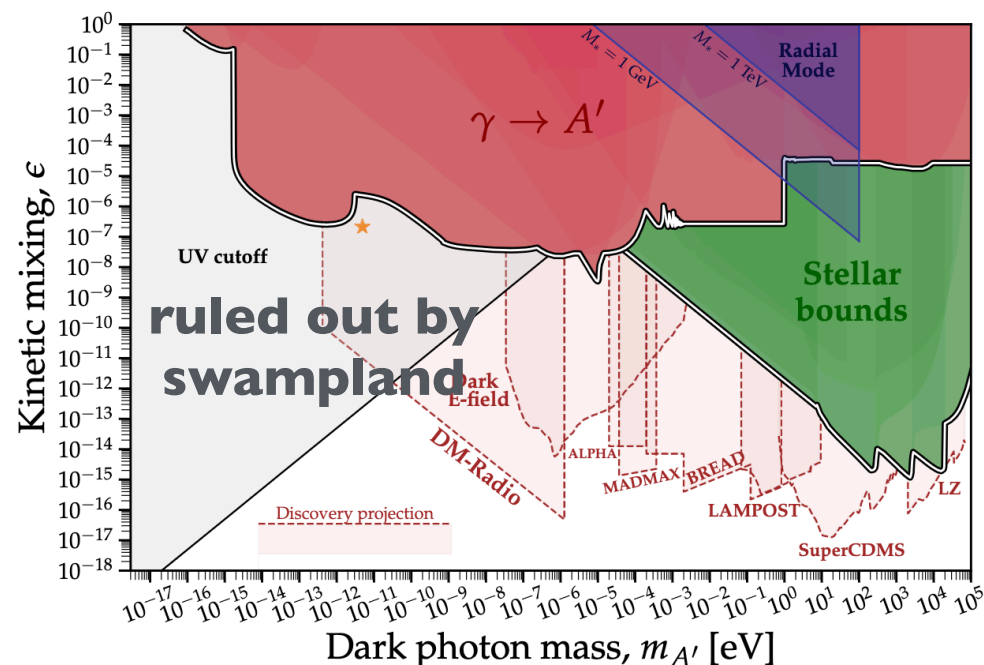
For weak coupling limits:

$$\exists \text{ tower with } m \leq \gamma_{\text{BH}} q g_{\text{YM}} M_p$$

$\xrightarrow{\mathcal{O}(1) \text{ factor}}$
 $\swarrow \text{quantized charge} \quad \searrow \text{gauge coupling}$

(Weak Gravity Conjecture)

[Arkani-Hamed et al'06]



[Montero,Muñoz,Obied'22]

→ EFT breaks down at $\Lambda \lesssim g_{\text{YM}} M_p$

Lower bound on gauge coupling!

It can rule out some BSM proposals
and been tested experimentally

Consistency under dimensional reduction

Background independence of quantum gravity implies that:

If we start with a theory
consistent with QG



Compactifications of that theory
should also be consistent

Consistency under dimensional reduction

Background independence of quantum gravity implies that:

If we start with a theory
consistent with QG



Compactifications of that theory
should also be consistent

Take a 4-dimensional theory, and compactify one dimension on a circle of radius R . The 3-dimensional EFT contains a new scalar with potential:

$$V(R) = \frac{2\pi\Lambda_4}{R^2} + \text{Casimir energy}(n_b, n_f, m_b, m_f)$$

\downarrow \downarrow

tree-level one-loop corrections

n_b : bosonic d.o.f
 n_f : fermionic d.o.f

Consistency under dimensional reduction

Background independence of quantum gravity implies that:

If we start with a theory consistent with QG  Compactifications of that theory should also be consistent

Take a 4-dimensional theory, and compactify one dimension on a circle of radius R . The 3-dimensional EFT contains a new scalar with potential:

$$V(R) = \frac{2\pi\Lambda_4}{R^2} + \text{Casimir energy}(n_b, n_f, m_b, m_f)$$

\downarrow \downarrow

tree-level one-loop corrections

n_b : bosonic d.o.f
 n_f : fermionic d.o.f

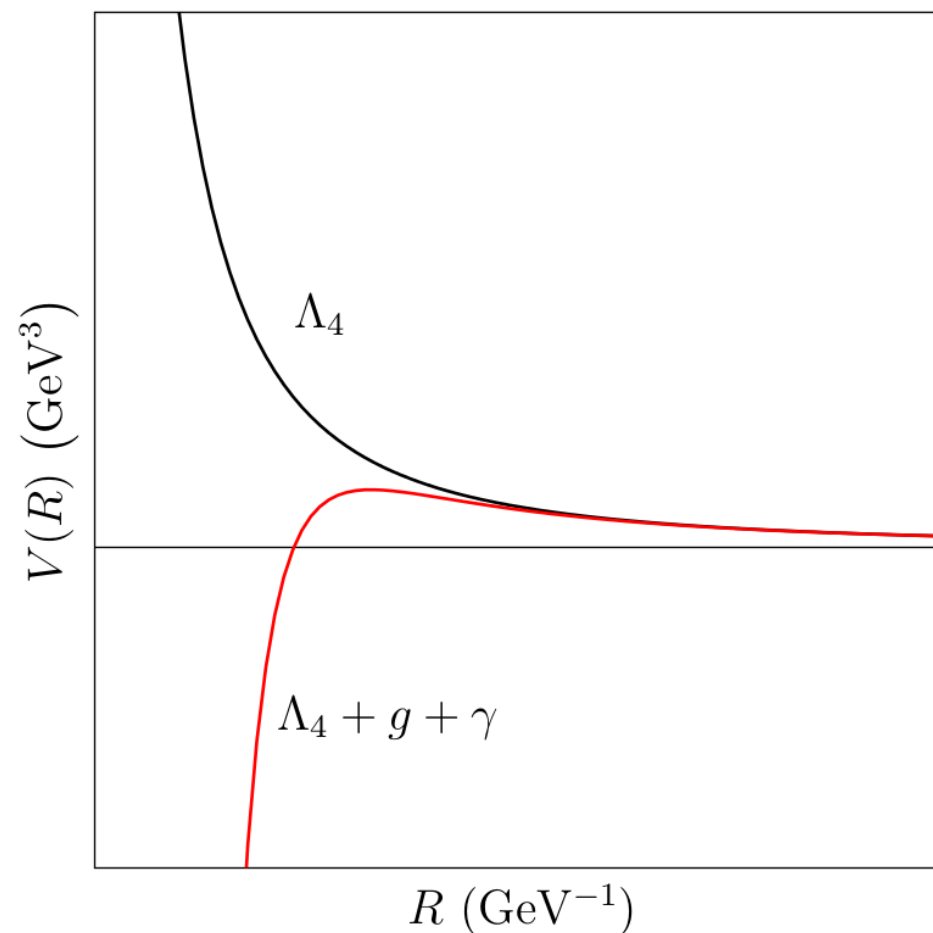
Goal: Determine constraints on the 4d field spectra that guarantee that $V(R)$ is consistent with Swampland conjectures

Compactification of the SM to 3d

Standard Model + Gravity on S^1 :

$$V(R) = \frac{2\pi\Lambda_4}{R^2} - \frac{4}{720\pi R^6}$$

massless particles:
graviton, photon



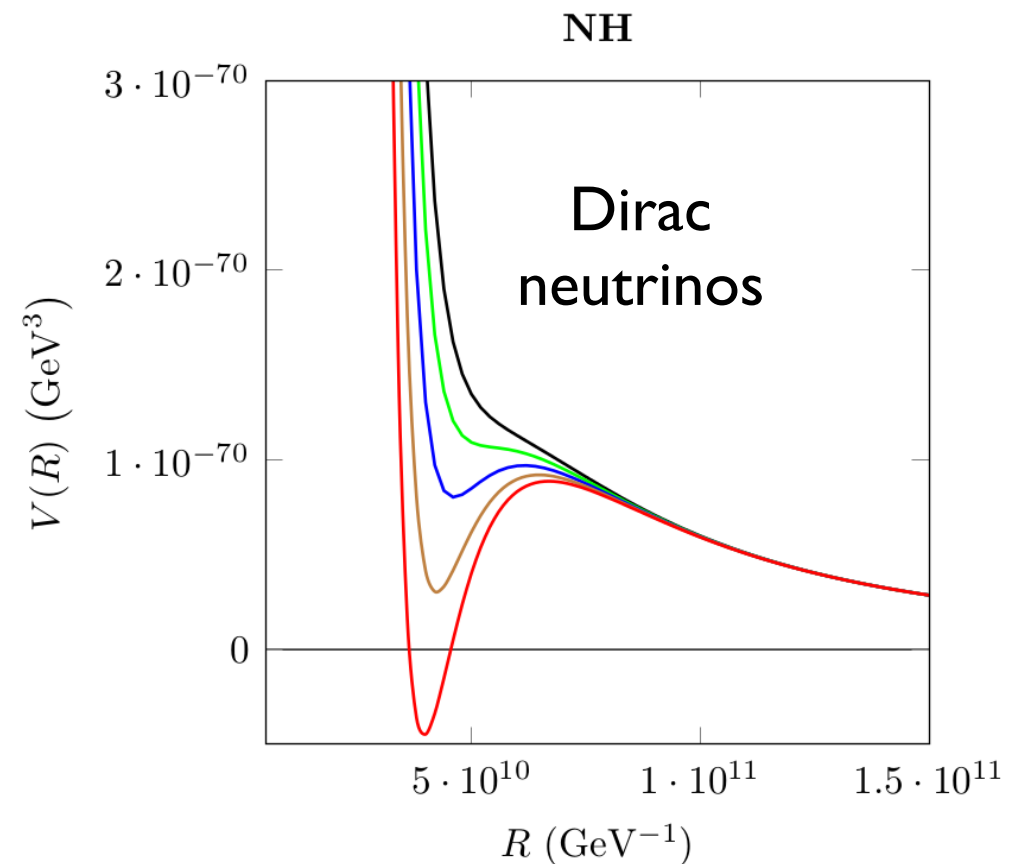
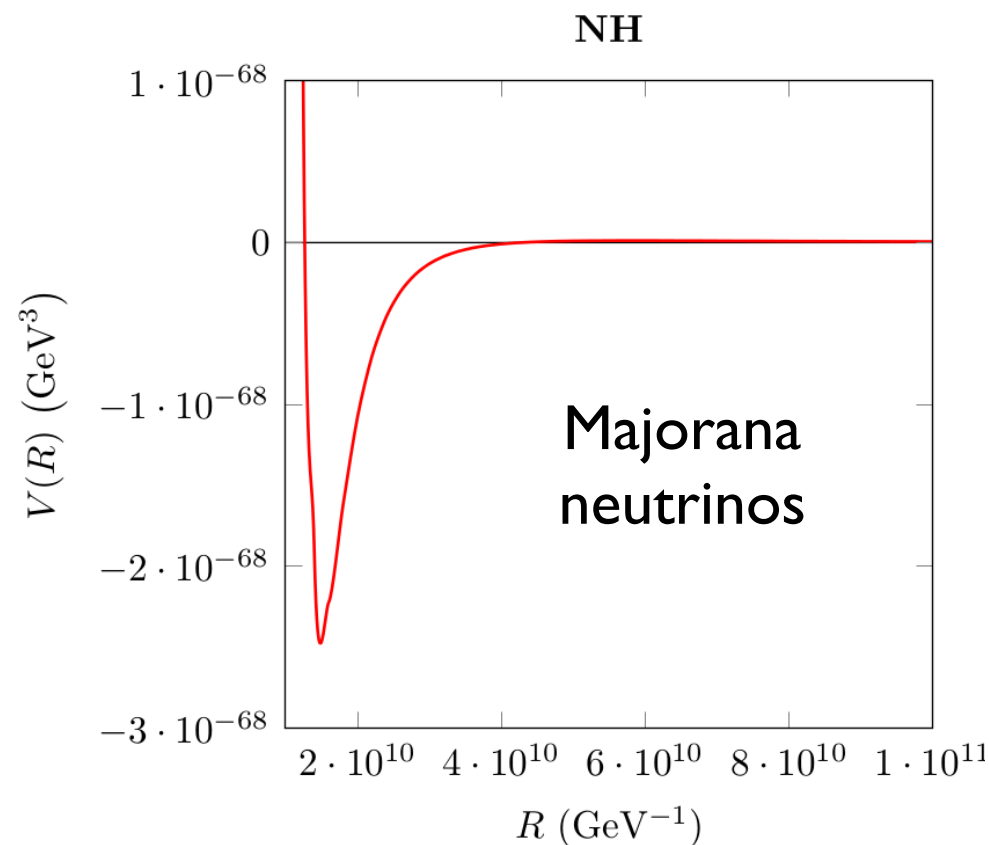
Compactification of the SM to 3d

Standard Model + Gravity on S^1 :

$$V(R) = \frac{2\pi\Lambda_4}{R^2} - \frac{4}{720\pi R^6} + \sum_i \frac{(2\pi R)}{R^3} (-1)^{s_i} n_i \rho_i(R)$$

massless particles:
graviton, photon

massive particles:
neutrinos,...

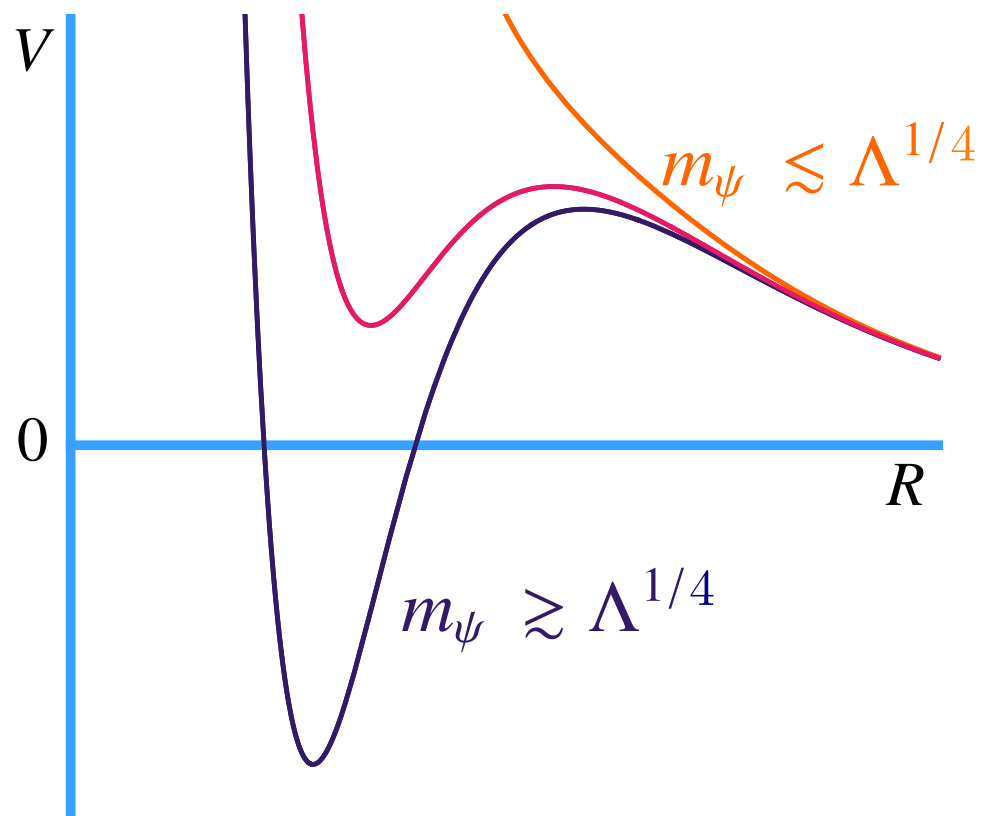


The more massive the neutrinos, the deeper the AdS vacuum

[Ibanez, Martin-Lozano, IV'17] (see also [Hamada-Shiu'17])

Compactification of the SM to 3d

- Results:**
- Depending on the value of Dirac mass for the neutrinos, we get 3-dimensional AdS, Minkowski or dS vacua.
 - By scanning a family of 4d theories with different values of the neutrino masses (e.g. by scanning the Higgs vev), we can cross the flat space limit.



We cross the flat space limit when neutrino masses:

$$m_\nu \sim \Lambda_{\text{cc}}^{1/4}$$

Compactification of the SM to 3d

According to the Generalized (AdS) Distance Conjecture, there should be a tower of states becoming light in the flat space limit

Two possible resolutions:

- 1) **There is a surplus of light fermions** with $m \lesssim \Lambda_{cc}^{1/4}$
(larger masses are obstructed by quantum gravity consistency)
- 2) **There is an infinite tower of states** in 4d scaling as $m_{\text{tower}} \sim \Lambda_{cc}^{1/4}$



The KK tower of the Dark Dimension ensures that compactifications of the SM are consistent with the Swampland conjectures

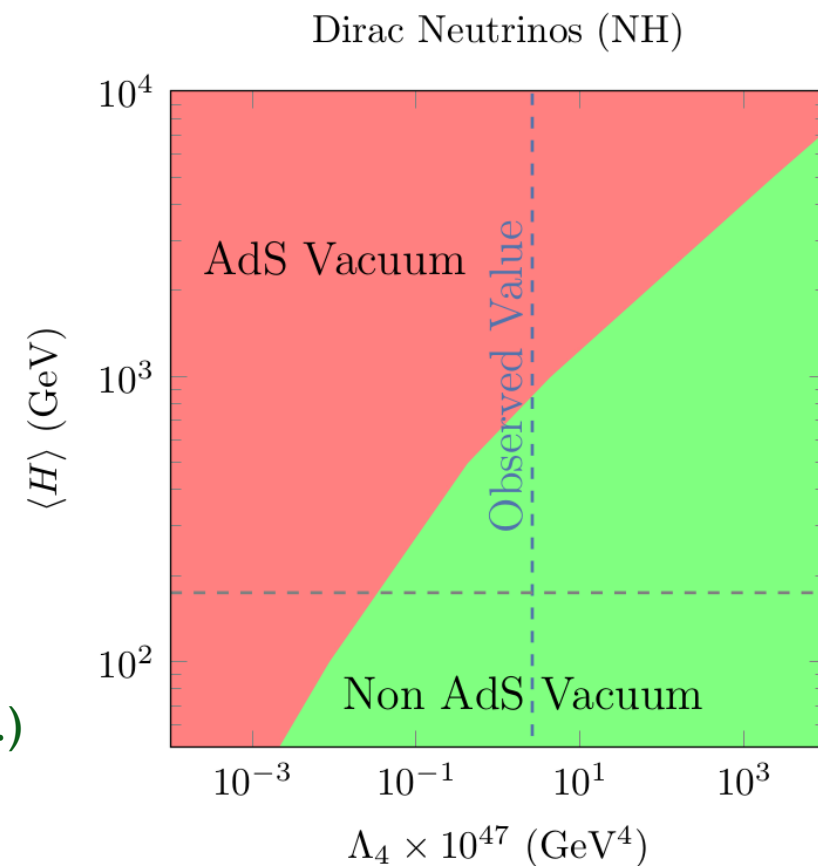
Naturalness issues

We can translate the bound on neutrino masses to an upper bound on the EW scale in terms of the cosmological constant:

$$\langle H \rangle \lesssim 1.6 \frac{\Lambda^{1/4}}{Y_{\nu_1}}$$

[Martin-Lozano, Ibanez, IV'17]

(see also [Gonzalo et al'18-21] [Rudelius'21]...)



Parameters leading to a higher EW scale do not yield theories consistent with quantum gravity

Solution to EW hierarchy problem?

Space of parameters consistent with quantum gravity is smaller than expected, not everything goes!

We need to revisit the logic of naturalness

String Theory Evidence

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

 All known holographic AdS vacua [Luest,Palti,Vafa'19]

String Theory Evidence

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

☑ All known holographic AdS vacua [Luest,Palti,Vafa'19]

☑ 4d N=1 flux string compactifications with positive runaways

[Lanza, Marchesano, Martucci, IV'20] [Andriot et al'20]

String Theory Evidence

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

☑ All known holographic AdS vacua [Luest,Palti,Vafa'19]

☑ 4d N=1 flux string compactifications with positive runaways

[Lanza, Marchesano, Martucci, IV'20] [Andriot et al'20]

☑ Non-SUSY string theories with positive runaways

String Theory Evidence

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

☑ All known holographic AdS vacua [Luest,Palti,Vafa'19]

☑ 4d N=1 flux string compactifications with positive runaways

[Lanza, Marchesano, Martucci, IV'20] [Andriot et al'20]

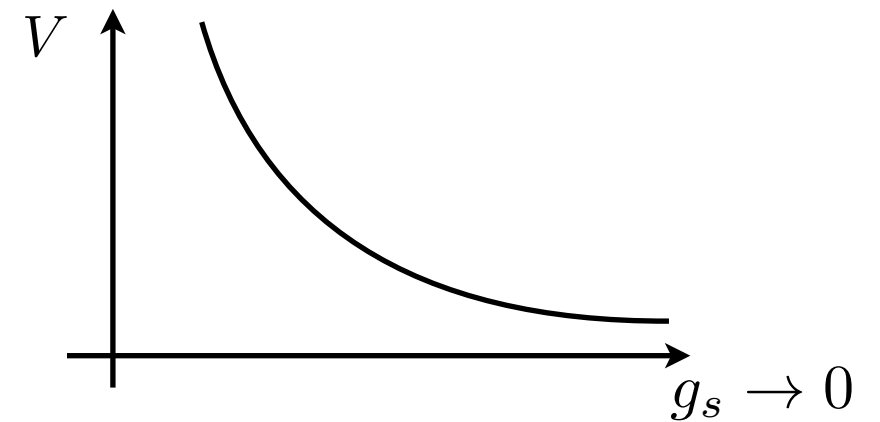
☑ Non-SUSY string theories with positive runaways

☑ KKLT-like proposals for dS in string theory [Blumenhagen et al'22]

Failure of EFT expectation

Non-SUSY example

Recall: $SO(16) \times SO(16)$ non-SUSY (tachyon-free) heterotic string theory:



Positive runaway on the dilaton

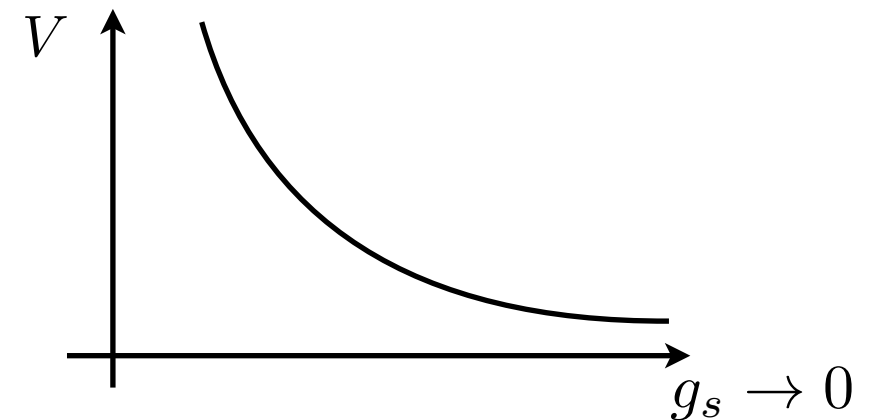
Failure of EFT expectation

Non-SUSY example

Recall: $SO(16) \times SO(16)$ non-SUSY (tachyon-free) heterotic string theory:

Tower of string modes becoming light in the weak coupling limit, starting at

$$m \sim M_s$$



Positive runaway on the dilaton

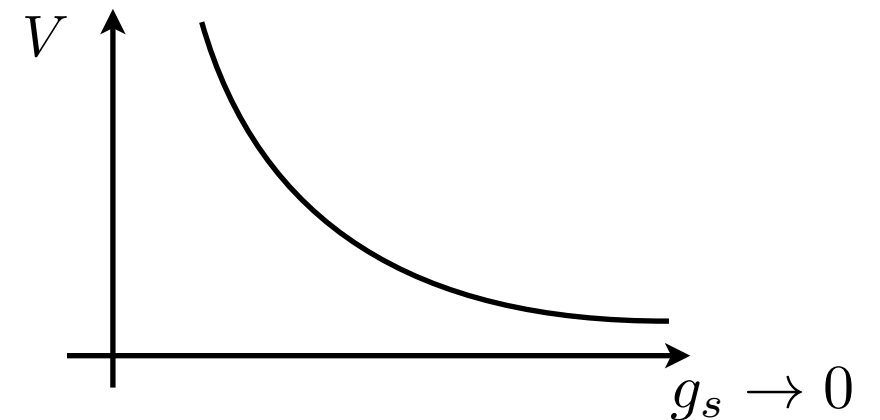
Failure of EFT expectation

Non-SUSY example

Recall: $SO(16) \times SO(16)$ non-SUSY (tachyon-free) heterotic string theory:

Tower of string modes becoming light in the weak coupling limit, starting at

$$m \sim M_s$$



Positive runaway on the dilaton

$$V_{1\text{-loop}} \sim - \sum_i (-1)^{F_i} \int_{\Lambda_{UV}^{-2}}^{\infty} \frac{ds}{s^6} \exp\left(-\frac{m_i^2 s}{2}\right)$$

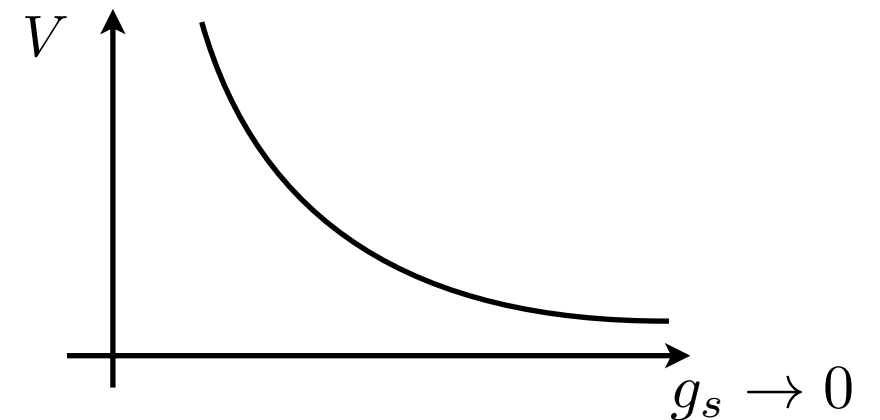
Failure of EFT expectation

Non-SUSY example

Recall: $SO(16) \times SO(16)$ non-SUSY (tachyon-free) heterotic string theory:

Tower of string modes becoming light in the weak coupling limit, starting at

$$m \sim M_s$$



Positive runaway on the dilaton

$$V_{1\text{-loop}} \sim - \sum_i (-1)^{F_i} \int_{\Lambda_{UV}^{-2}}^{\infty} \frac{ds}{s^6} \exp\left(-\frac{m_i^2 s}{2}\right) \quad \rightarrow \quad V \sim m^{10}$$

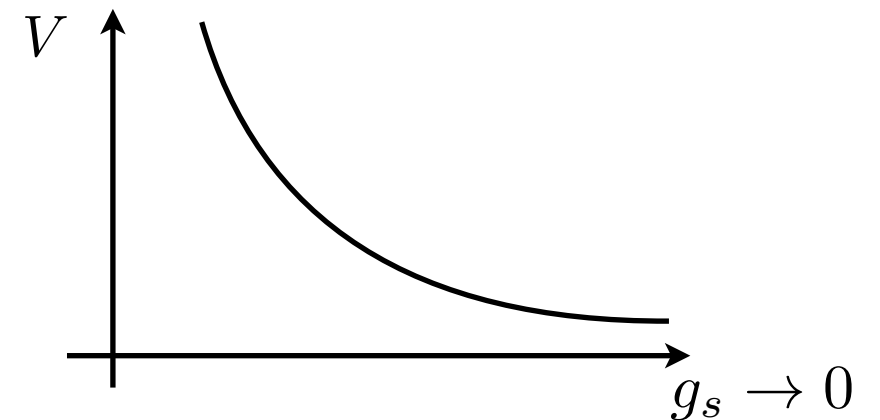
Failure of EFT expectation

Non-SUSY example

Recall: $SO(16) \times SO(16)$ non-SUSY (tachyon-free) heterotic string theory:

Tower of string modes becoming light in the weak coupling limit, starting at

$$m \sim M_s$$



Positive runaway on the dilaton

$$V_{1\text{-loop}} \sim - \sum_i (-1)^{F_i} \int_{\Lambda_{UV}^{-2}}^{\infty} \frac{ds}{s^6} \exp\left(-\frac{m_i^2 s}{2}\right) \quad \rightarrow \quad V \sim m^{10}$$

Contribution of massive string excitations is cut-off at M_s due to modular invariance

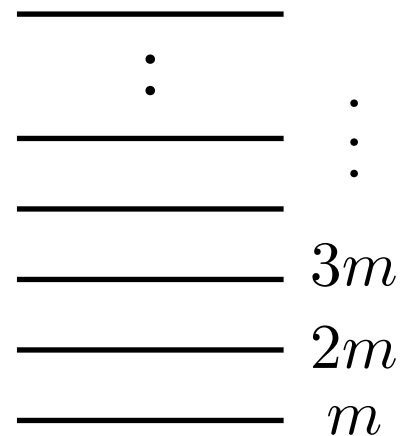
Failure of EFT expectation

It is very important to integrate out the entire infinite tower of states

Failure of EFT expectation

It is very important to integrate out the entire infinite tower of states

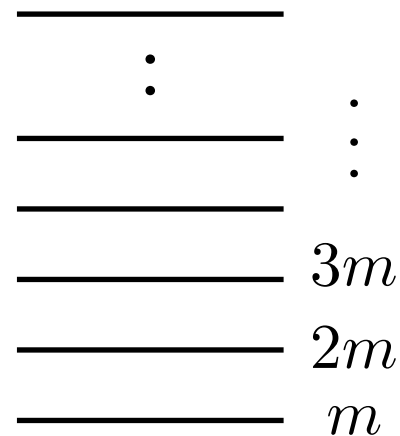
The result is drastically different than if integrating out only a finite number of them



Failure of EFT expectation

It is very important to integrate out the entire infinite tower of states

The result is drastically different than if integrating out only a finite number of them



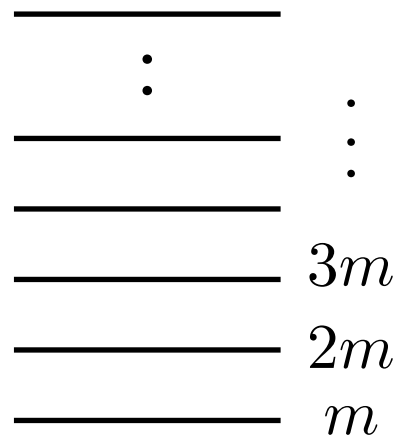
- If integrating out the infinite tower:

$$V_0 \sim m_1^d \quad \text{first light state of the tower!}$$

Failure of EFT expectation

It is very important to integrate out the entire infinite tower of states

The result is drastically different than if integrating out only a finite number of them



- If integrating out the infinite tower:

$$V_0 \sim m_1^d \quad \text{first light state of the tower!}$$

- If integrating a finite number of fields below a cut-off:

$$V_0 \sim m_{\text{heavy}}^d \quad \text{the heavy states dominate}$$

Experimental constraints

Is a tower with $V^{1/2} \lesssim m \lesssim V^{1/4}$

compatible with experimental constraints?

In our universe: $V^{1/4} \sim 2.31 \text{ meV} \rightarrow m^{-1} \gtrsim O(10) \mu m$

Experimental constraints:

❖ Astrophysical bounds: $m^{-1} \leq 10^{-4} \mu m \quad (n = 2)$

[Hannestad and Raffelt '03] $m^{-1} \leq 44 \mu m \quad (n = 1)$

❖ Dev. from Newton's laws ($n=1$): $m^{-1} \leq 30 \mu m$

[Lee et al '21]

Experimental constraints

Is a tower with $V^{1/2} \lesssim m \lesssim V^{1/4}$

compatible with experimental constraints?

In our universe: $V^{1/4} \sim 2.31 \text{ meV} \rightarrow m^{-1} \gtrsim O(10) \mu m$

Experimental constraints:

❖ Astrophysical bounds: $m^{-1} \leq 10^{-4} \mu m$ ~~$(n=2)$~~ ruled out
[Hannestad and Raffelt '03] $m^{-1} \leq 44 \mu m$ $(n=1)$

❖ Dev. from Newton's laws ($n=1$): $m^{-1} \leq 30 \mu m$
[Lee et al '21]

Experimental constraints

Is a tower with $V^{1/2} \lesssim m \lesssim V^{1/4}$

compatible with experimental constraints?

In our universe: $V^{1/4} \sim 2.31 \text{ meV} \rightarrow m^{-1} \gtrsim O(10) \mu m$

Experimental constraints:

❖ Astrophysical bounds: $m^{-1} \leq 10^{-4} \mu m$ ~~$(n=2)$~~ ruled out
[Hannestad and Raffelt '03] $m^{-1} \leq 44 \mu m$ $(n=1)$

❖ Dev. from Newton's laws ($n=1$): $m^{-1} \leq 30 \mu m$
[Lee et al '21]

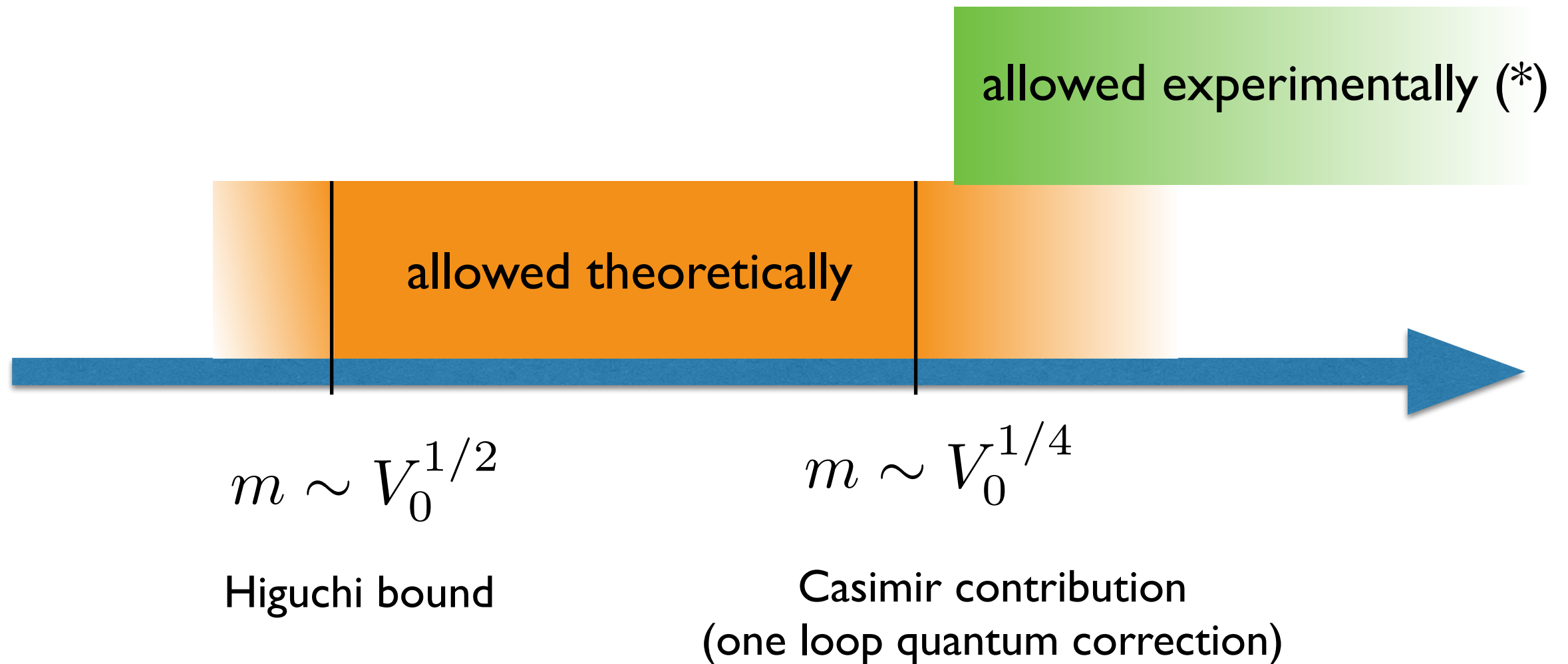
Only $n=1$ (one large extra dimension) is marginally compatible!

Dark Dimension

Mass scale of the tower of states:

Dark Dimension

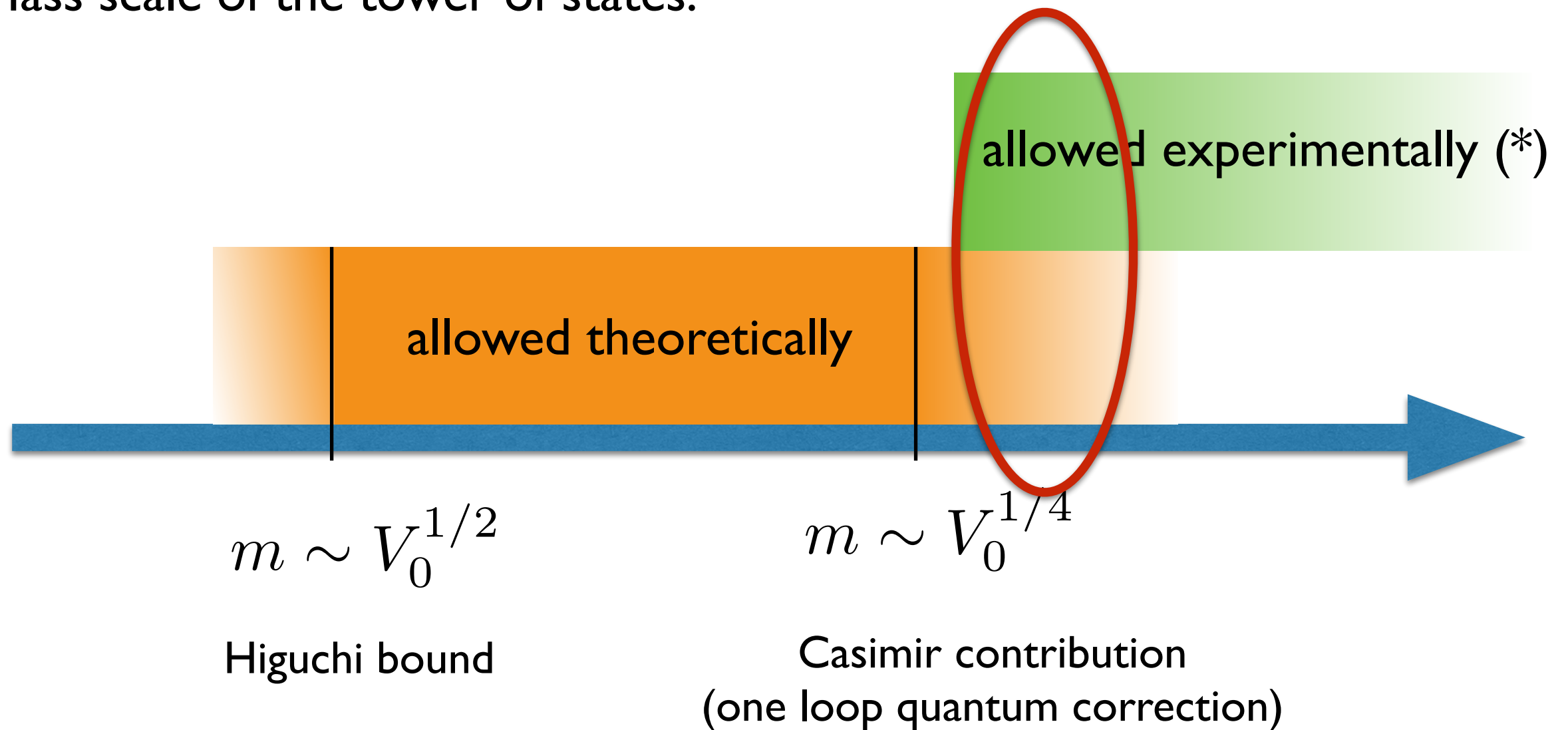
Mass scale of the tower of states:



(*) astrophysical bounds and deviations from Newton's law

Dark Dimension

Mass scale of the tower of states:



(*) astrophysical bounds and deviations from Newton's law

Evidence for WGC and SDC

❖ **String theory compactifications:** Plethora of quantitative tests!

- Systematic approach according to the level of supersymmetry
- Interesting connections to mathematics

[Grimm, Palti, IV'18]

[Grimm, Palti, Li'18]

[Lee, Lerche, Weigand'18-19]

...

Evidence for WGC and SDC

❖ String theory compactifications: Plethora of quantitative tests!

- Systematic approach according to the level of supersymmetry
- Interesting connections to mathematics

[Grimm, Palti, IV'18]

[Grimm, Palti, Li'18]

[Lee, Lerche, Weigand'18-19]

...

❖ AdS/CFT:

- WGC proven for AdS3 using modular invariance of the CFT
- WGC from QI theorems and entanglement entropy
- SDC formulated in terms of a CFT Distance conjecture

[Heidenreich et al'16]

[Montero et al'16]

[Montero'18]

[Perlmutter et al'20]

Evidence for WGC and SDC

❖ String theory compactifications: Plethora of quantitative tests!

- Systematic approach according to the level of supersymmetry
- Interesting connections to mathematics

[Grimm, Palti, IV'18]

[Grimm, Palti, Li'18]

[Lee, Lerche, Weigand'18-19]

...

❖ AdS/CFT:

- WGC proven for AdS3 using modular invariance of the CFT
- WGC from QI theorems and entanglement entropy
- SDC formulated in terms of a CFT Distance conjecture

[Heidenreich et al'16]

[Montero et al'16]

[Montero'18]

[Perlmutter et al'20]

❖ Black hole arguments:

- WGC follows from requiring black holes to decay
- WGC/SDC follows from entropy bounds associated to small BHs
- Connection between WGC and weak cosmic censorship

[Arkani-Hamed et al'06]

[Hamada et al'21]

[Crisford et al'17]

Evidence for WGC and SDC

❖ String theory compactifications: Plethora of quantitative tests!

- Systematic approach according to the level of supersymmetry
- Interesting connections to mathematics

[Grimm, Palti, IV'18]

[Grimm, Palti, Li'18]

[Lee, Lerche, Weigand'18-19]

...

❖ AdS/CFT:

- WGC proven for AdS3 using modular invariance of the CFT
- WGC from QI theorems and entanglement entropy
- SDC formulated in terms of a CFT Distance conjecture

[Heidenreich et al'16]

[Montero et al'16]

[Montero'18]

[Perlmutter et al'20]

❖ Black hole arguments:

- WGC follows from requiring black holes to decay
- WGC/SDC follows from entropy bounds associated to small BHs
- Connection between WGC and weak cosmic censorship

[Arkani-Hamed et al'06]

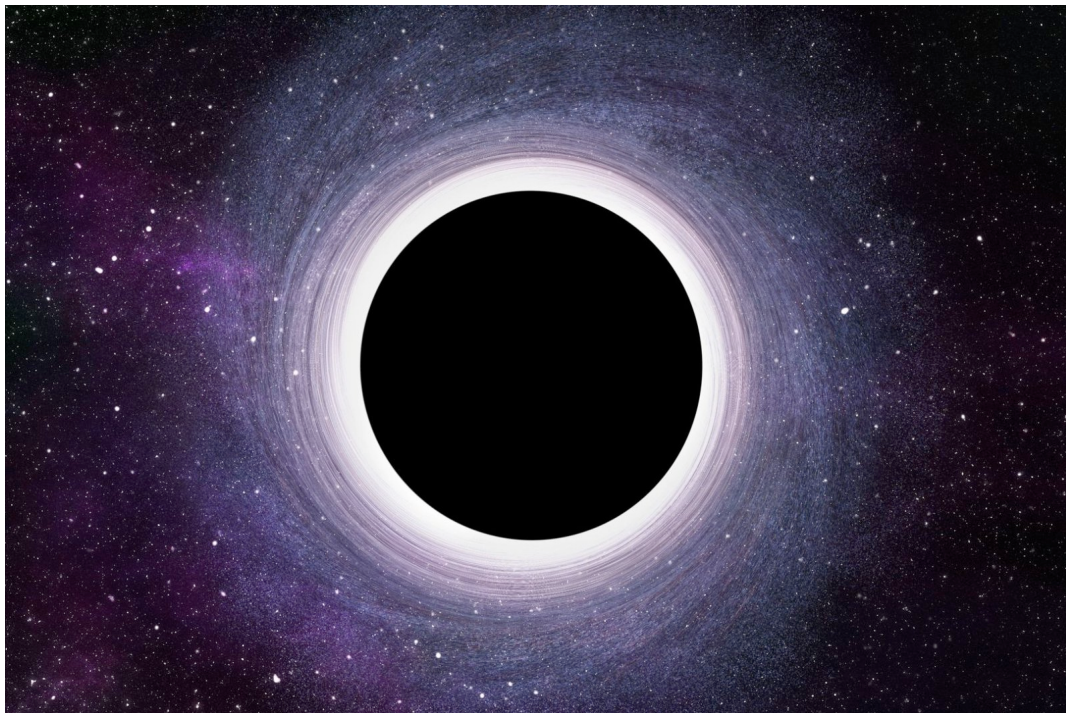
[Hamada et al'21]

[Crisford et al'17]

❖ Using positivity/unitarity bounds: lead to mild versions of the WGC

[Cheung et al'18][Hamada et al'18]...

Motivation from black holes

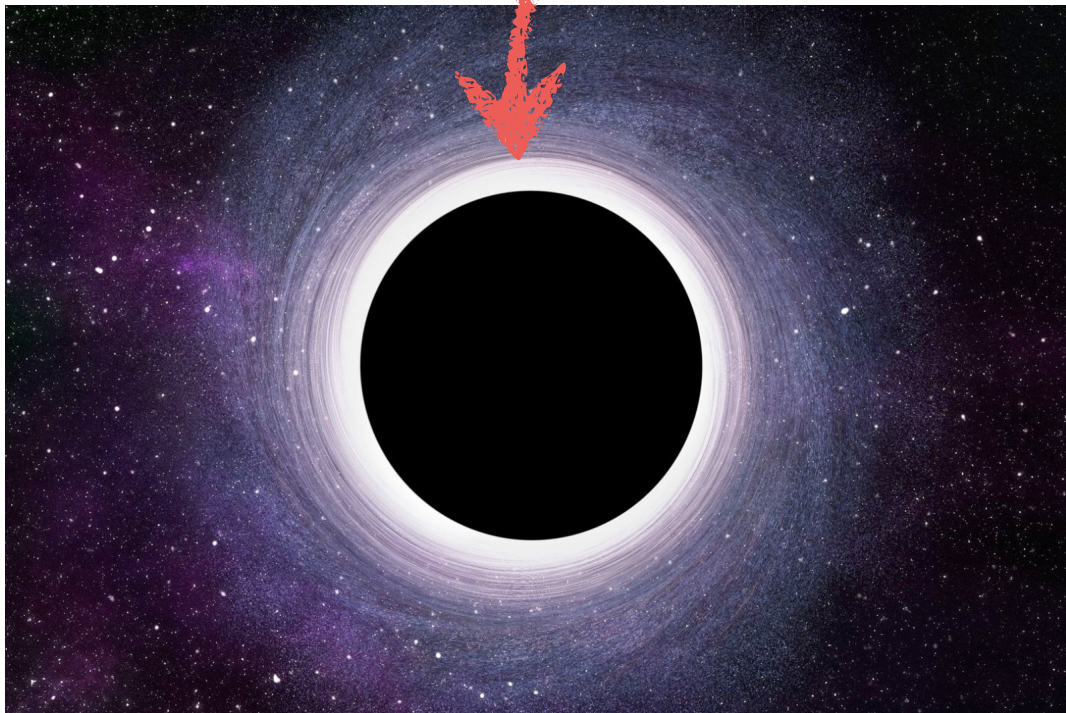


Black holes have
finite entropy



(“how many different ways lead to
the same black hole exterior”)

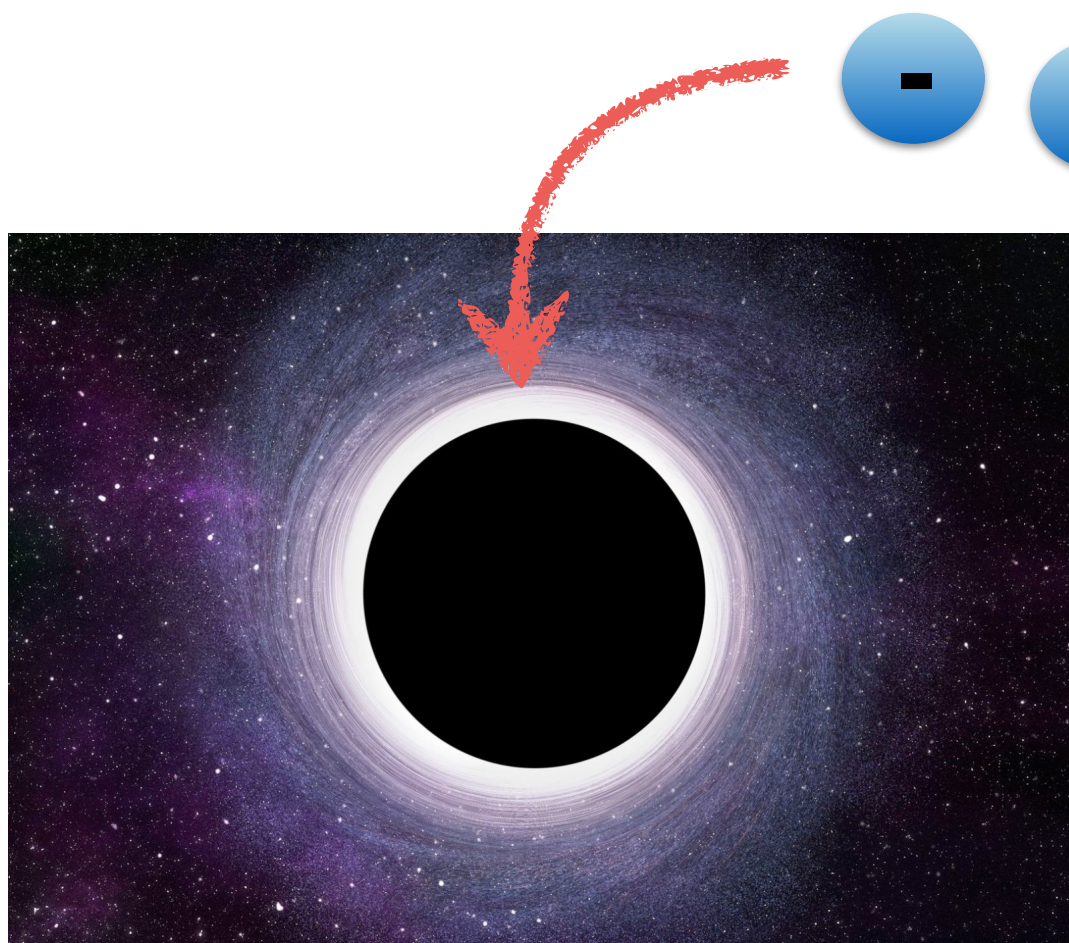
Motivation from black holes



Black holes have
finite entropy

(“how many different ways lead to
the same black hole exterior”)

Motivation from black holes

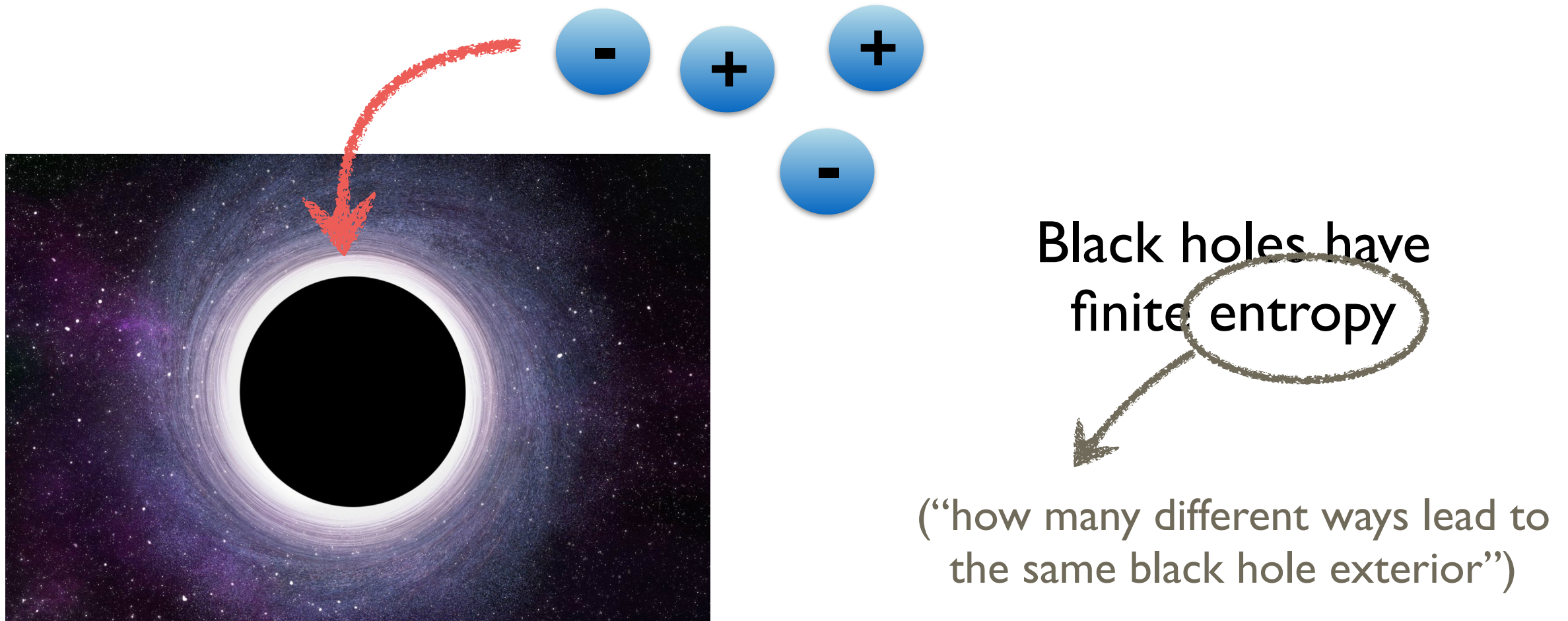


Black holes have
finite entropy

(“how many different ways lead to
the same black hole exterior”)

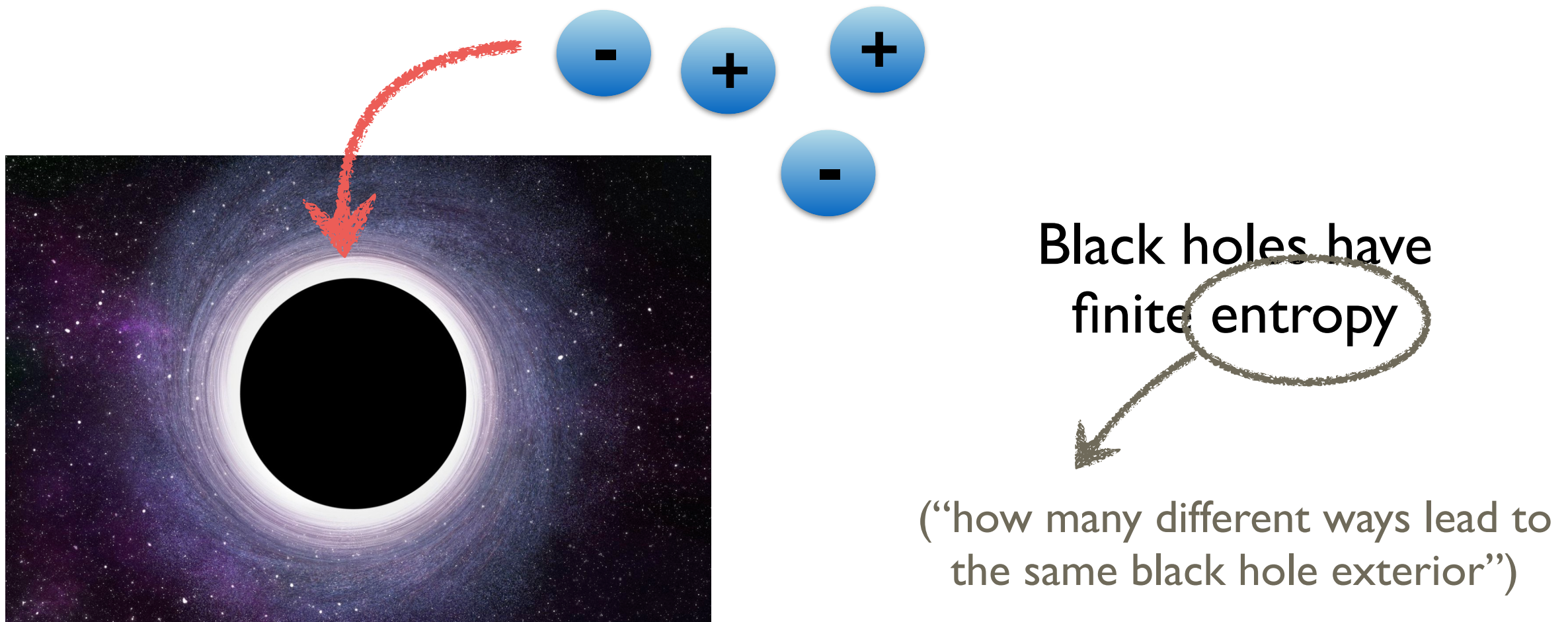
Since global symmetries cannot be detected from far away,

Motivation from black holes



Since global symmetries cannot be detected from far away, we could have infinitely many black holes with different values of the global charge, but that look the same from far way

Motivation from black holes



Since global symmetries cannot be detected from far away,
we could have infinitely many black holes with different values of
the global charge, but that look the same from far way



Infinite entropy

contradiction!

Evidence for “No global symmetries”

- Proof in perturbative string theory [Polchinski's book]
- Proof in AdS/CFT [Harlow,Ooguri '18]
- Correlation to unitary black hole evaporation (and topology changing processes)
[Harlow,Shaghoulian '20] [Chen,Lin '20] [Hsin et al '20] [Yonekura '20] ...


Evidence for “No global symmetries”

- Proof in perturbative string theory [Polchinski's book]
- Proof in AdS/CFT [Harlow,Ooguri '18]
- Correlation to unitary black hole evaporation (and topology changing processes)
[Harlow,Shaghoulian '20] [Chen,Lin '20] [Hsin et al '20] [Yonekura '20] ...

Global symmetries are not well defined in quantum gravity as the topology itself fluctuates

Implications for Cosmology

Maximum scalar field range that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off


$$\Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

Implications for Cosmology

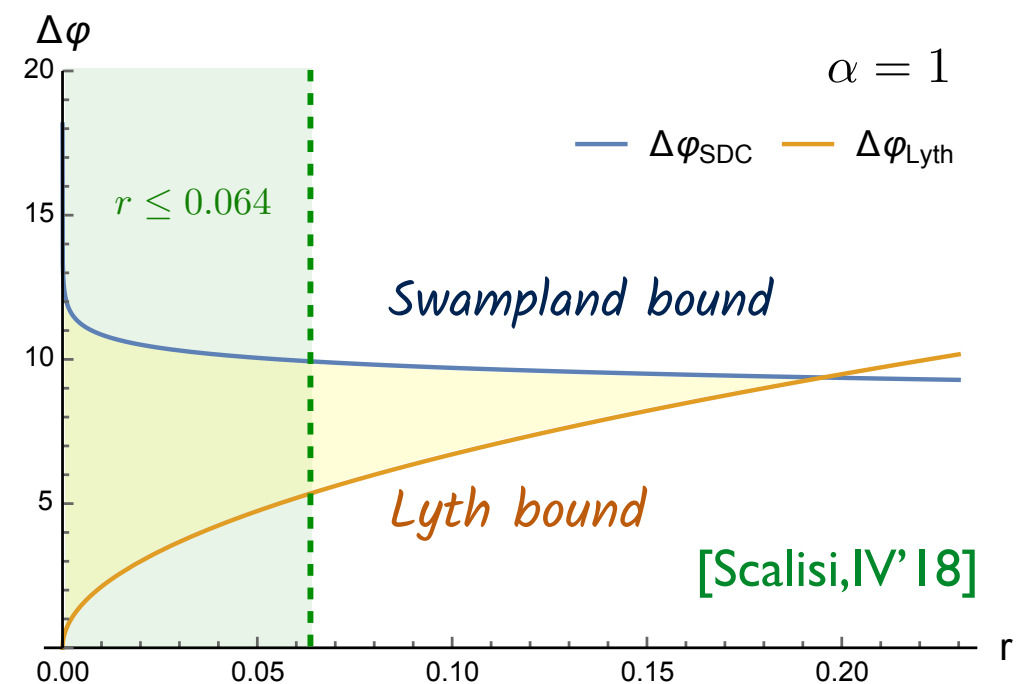
Maximum scalar field range that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

Example: Constraints on single field inflation

$$\Delta\phi \leq \frac{1}{\lambda} \log \frac{M_p}{H} = \frac{1}{\lambda} \log \sqrt{\frac{2}{\pi^2 A_s r}}$$

\downarrow
 $H \leq \Lambda$



Implications for Cosmology

Maximum scalar field range that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

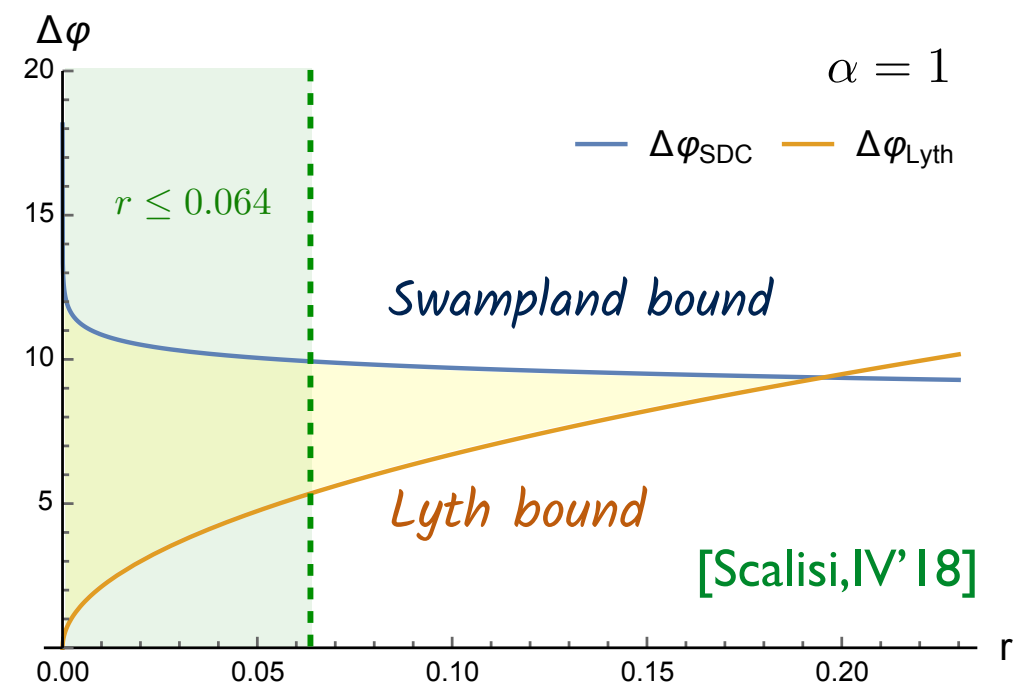
$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

Example: Constraints on single field inflation

$$\Delta\phi \leq \frac{1}{\lambda} \log \frac{M_p}{H} = \frac{1}{\lambda} \log \sqrt{\frac{2}{\pi^2 A_s r}}$$

\downarrow
 $H \leq \Lambda$

Large field inflation is not ruled out, but can be highly constrained



(constraints on multi-field inflation are more involved)

Implications for Cosmology

Maximum scalar field range that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

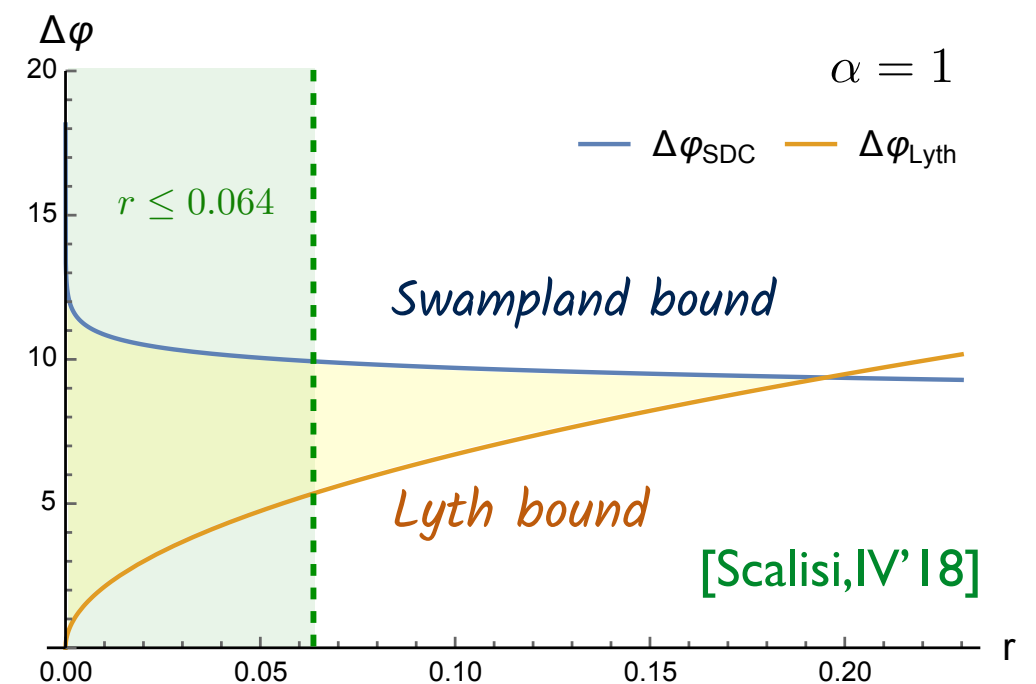
Example: Constraints on single field inflation

$$\Delta\phi \leq \frac{1}{\lambda} \log \frac{M_p}{H} = \frac{1}{\lambda} \log \sqrt{\frac{2}{\pi^2 A_s r}}$$

\downarrow
 $H \leq \Lambda$

Large field inflation is not ruled out,
but can be highly constrained

This triggered the revolution of the Swampland program in 2015



(constraints on multi-field inflation
are more involved)

Implications for Particle Physics

Constraints on new forces:



Implications for Particle Physics

Constraints on new forces:

- ❖ If (B-L) symmetry in the SM is gauged (there is a new gauge boson associated to it):

$$g < 10^{-24} \rightarrow \Lambda_{\text{QG}} < 10^{11} \text{ GeV}$$



experimental bound

Implications for Particle Physics

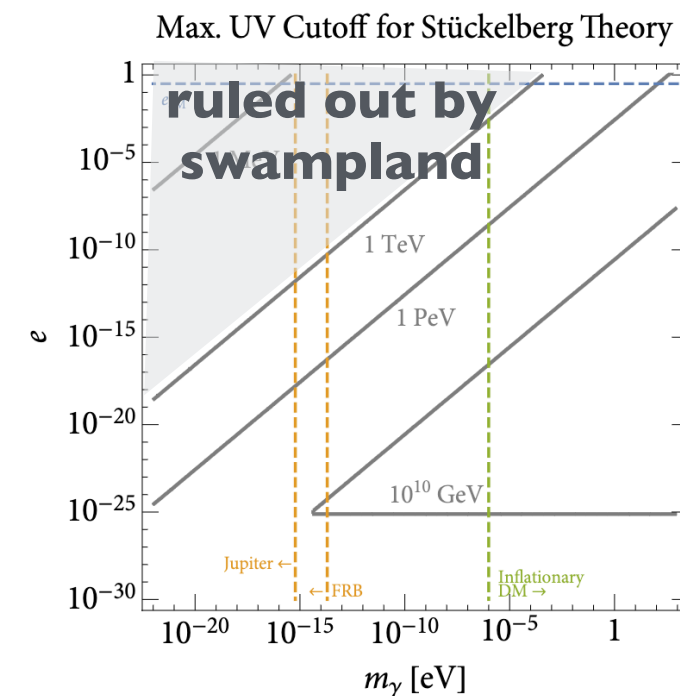
Constraints on new forces:

- ❖ If (B-L) symmetry in the SM is gauged (there is a new gauge boson associated to it):

$$g < 10^{-24} \rightarrow \Lambda_{\text{QG}} < 10^{11} \text{ GeV}$$

experimental bound

- ❖ Bounds on dark photons:



[Reece'18]

Implications for Particle Physics

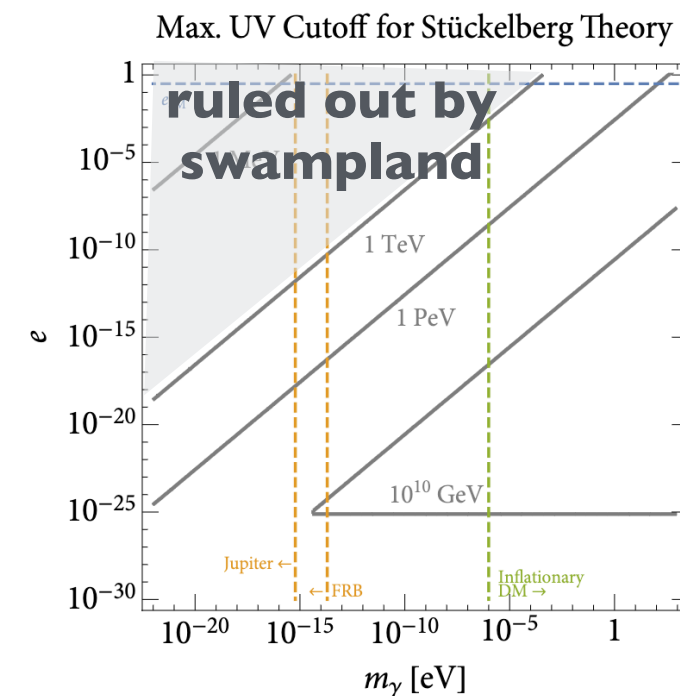
Constraints on new forces:

- ❖ If (B-L) symmetry in the SM is gauged (there is a new gauge boson associated to it):

$$g < 10^{-24} \rightarrow \Lambda_{QG} < 10^{11} \text{ GeV}$$

experimental bound

- ❖ Bounds on dark photons:



[Reece'18]

$\Lambda_{QG} \ll M_p$ for weakly coupled or very light dark photons