

Revisiting Quantum Criticality and Naturalness with the Standard Model

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IMPRS
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Fundamental Symmetries
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Standard notion of Naturalness

- No protection mechanism for the Higgs mass in the SM

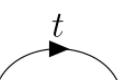
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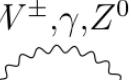
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 $(3y_t^2/8\pi^2)\Lambda^2$



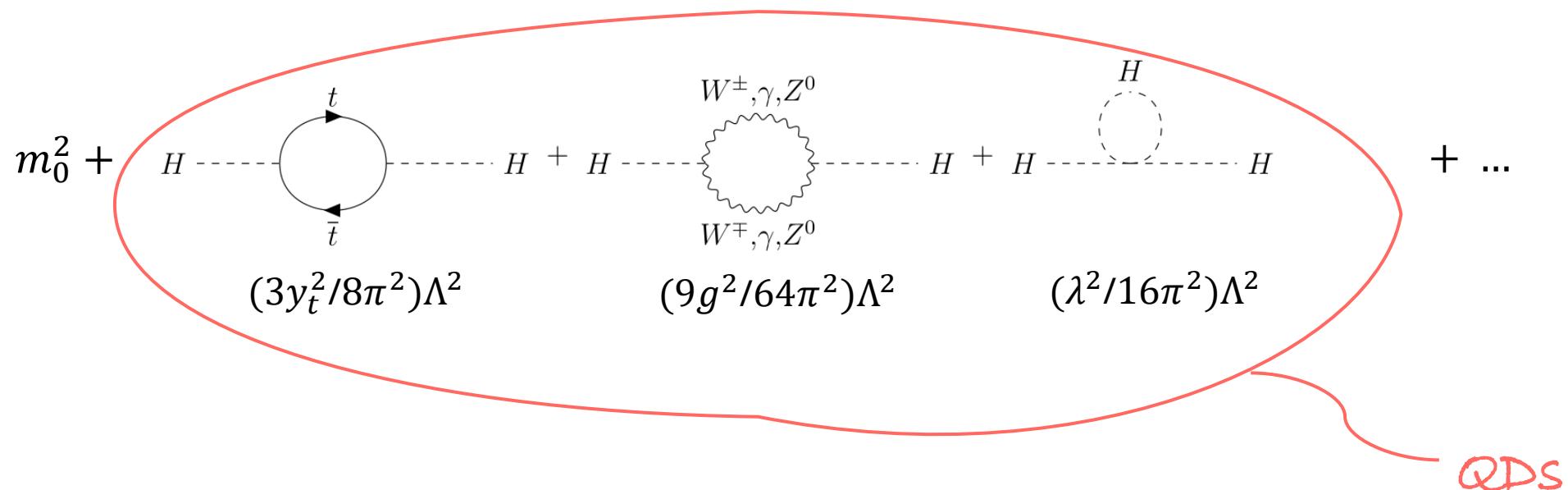
 $(9g^2/64\pi^2)\Lambda^2$



 $(\lambda^2/16\pi^2)\Lambda^2$

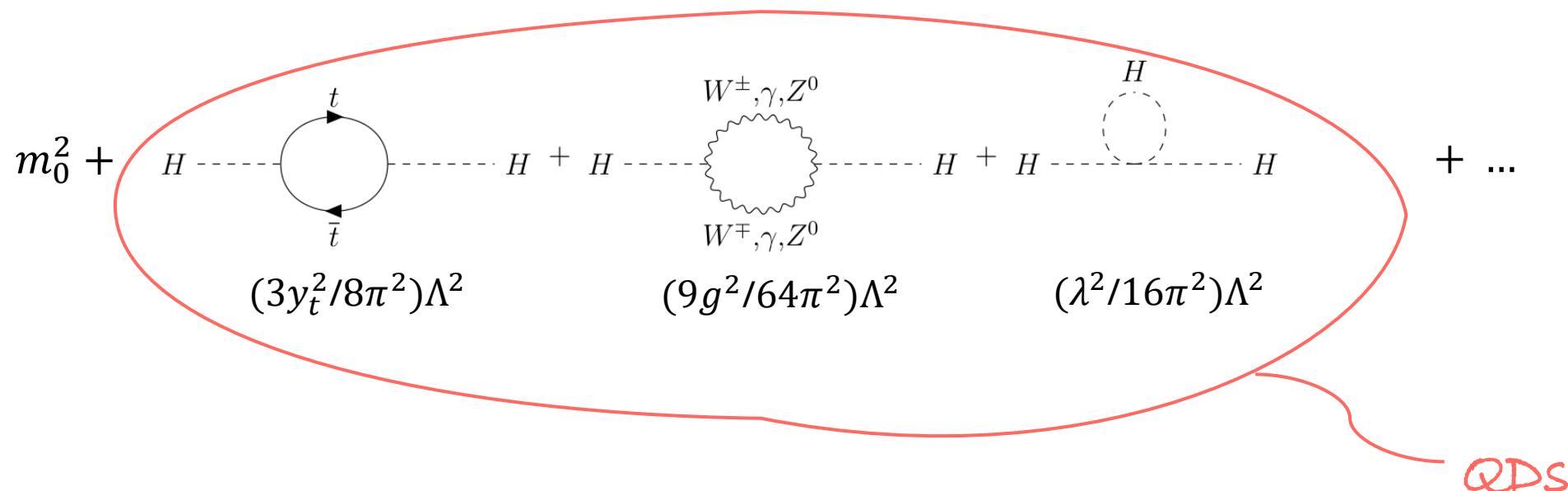
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- These QDs are the main source of **UV-sensitivity**

Naturalness in mass-independent schemes

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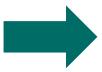
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 - QDs are manifested at the threshold scale
 - Manifesting the QDs requires an **Ansatz** of the UV embedding

Naturalness in mass-dependent schemes

- Physical effects of QDs contribute to the RGEs

Naturalness in mass-dependent schemes

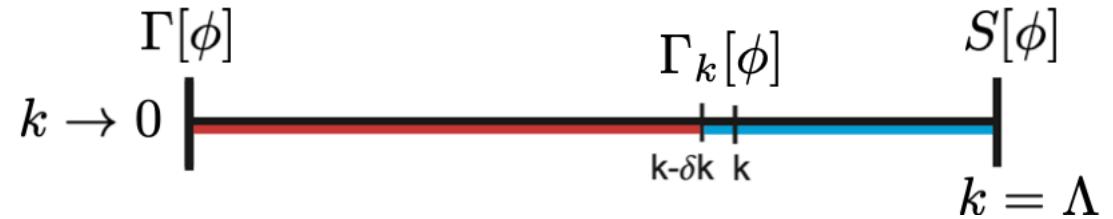
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 - Dynamical decoupling of heavy particles
 - UV sensitivity incorporated **along** the RG flow
 - And...

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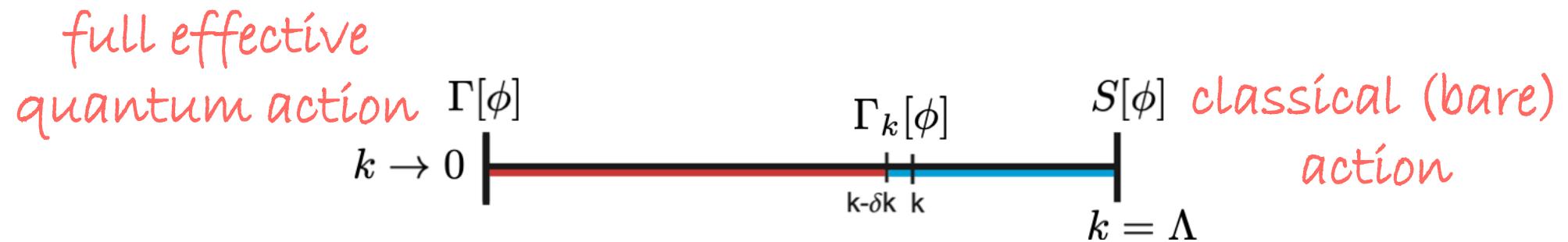
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Allows to study fine-tuning without an Ansatz of the UV embedding!

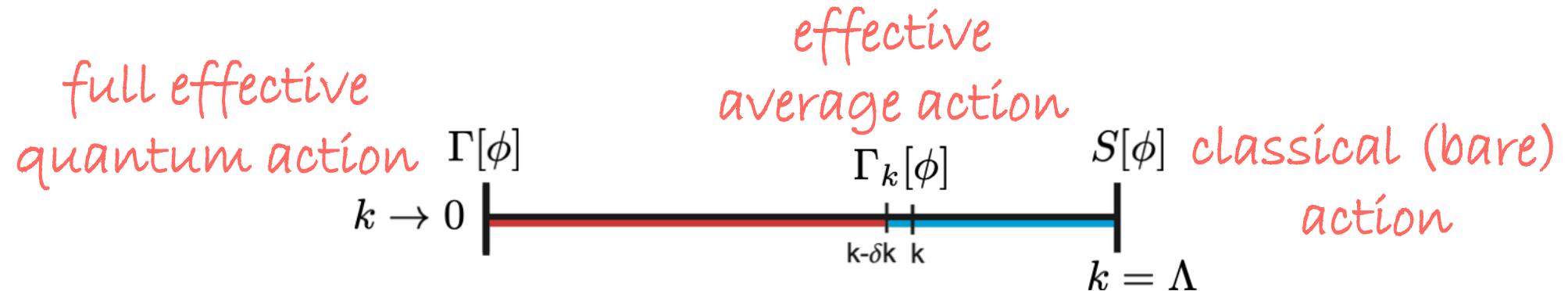
The functional RG



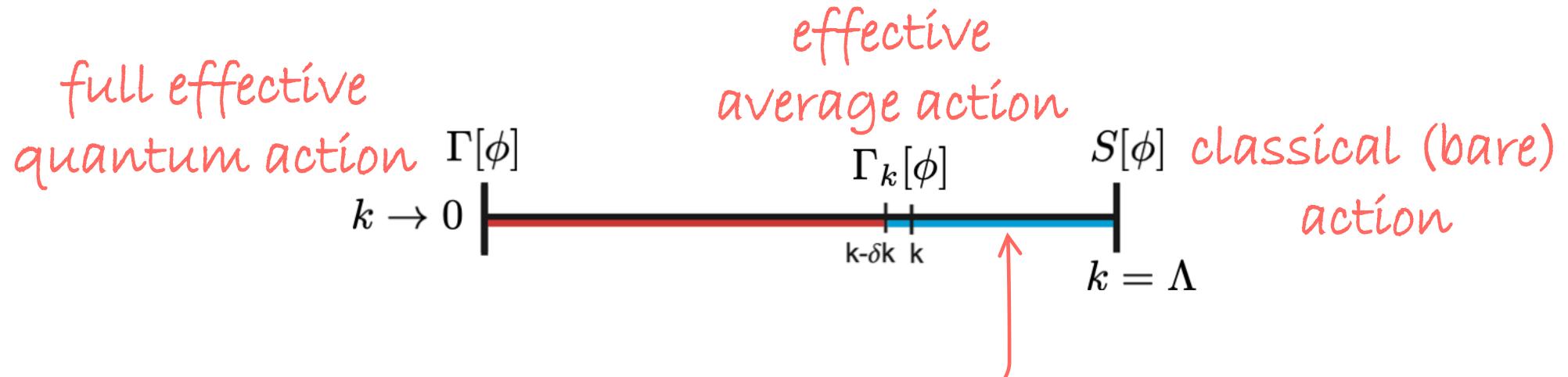
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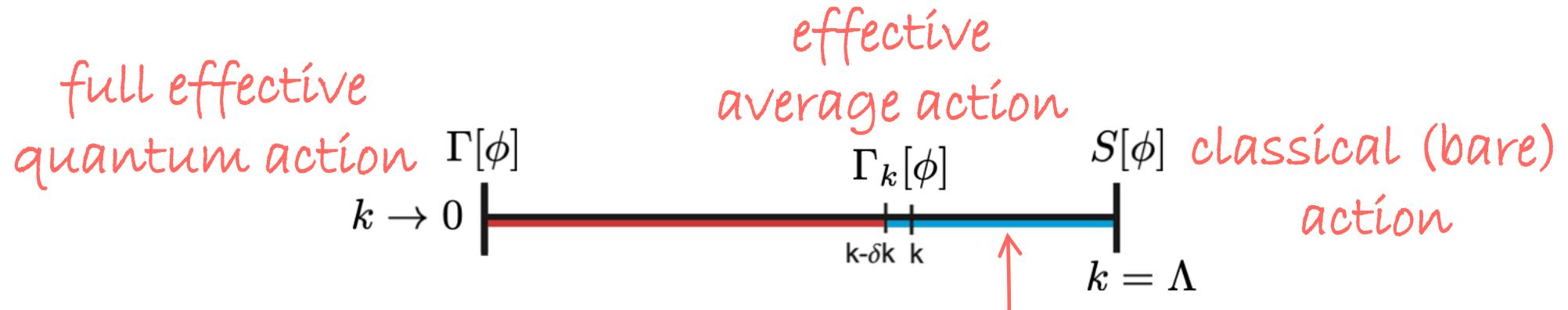
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(Wetterich '93) Flow equation

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left[\frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k \right]$$

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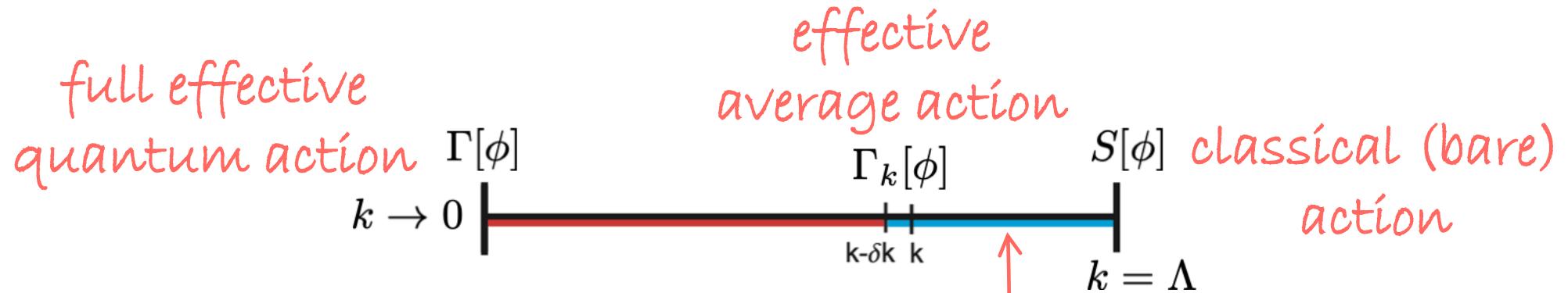


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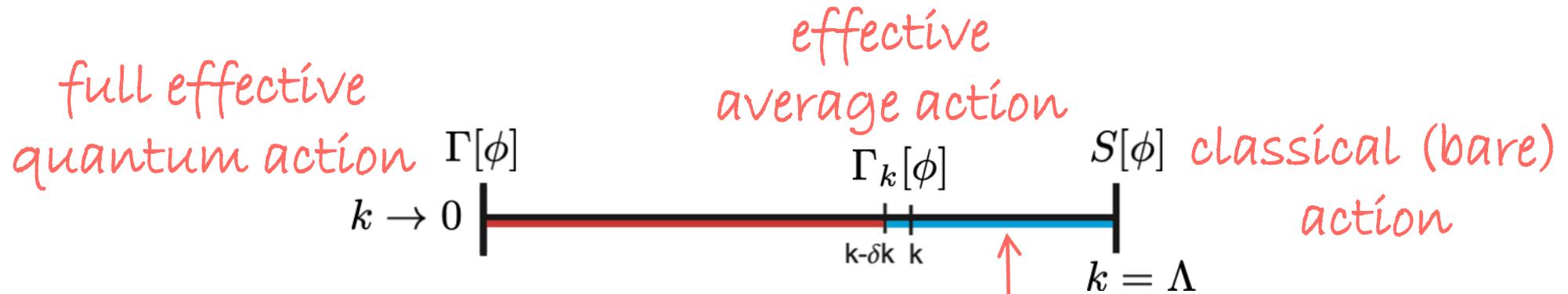
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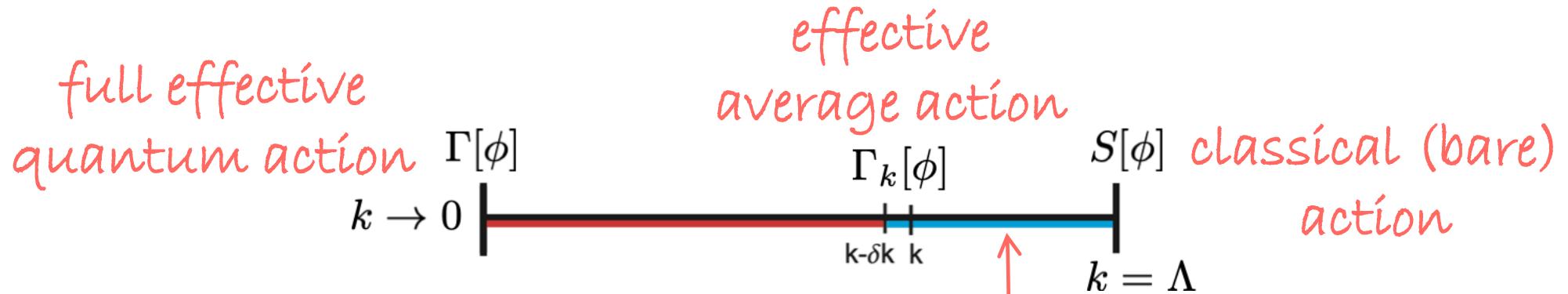
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mass-like regulator

$$\Delta S_{k[\phi]} = \int_p \phi(p) R_k \phi(-p)$$

Scalar couplings

$$\Gamma_k \supset V_{eff,k} = \mu_k^2 Z_{\Phi,k} \text{tr} \Phi^\dagger \Phi + \lambda_k (Z_{\Phi,k} \text{tr} \Phi^\dagger \Phi)^2$$

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \mathcal{G}_1 & i\mathcal{G}_2 \\ H & i\mathcal{G}_3 \end{pmatrix}$$

Scalar couplings

curvature

mass

quartic

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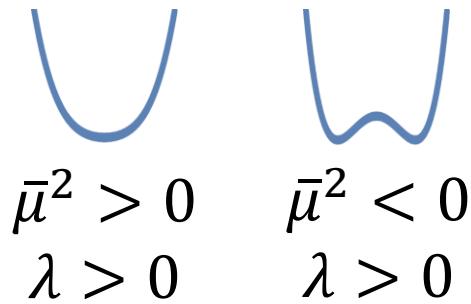
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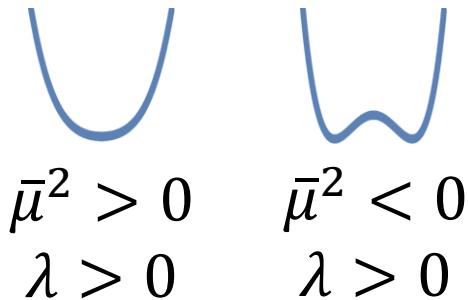
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$$\partial_t \bar{\mu}^2 \Big|_{\bar{\rho}_0=0} = (-2 + \eta_H) \bar{\mu}^2 + \partial_{\bar{\rho}} \overline{\text{Flow}}[V_{eff}] \Big|_{\bar{\rho}_0=0}$$

Scalar couplings

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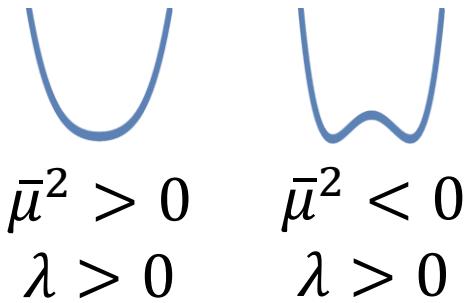
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Log running

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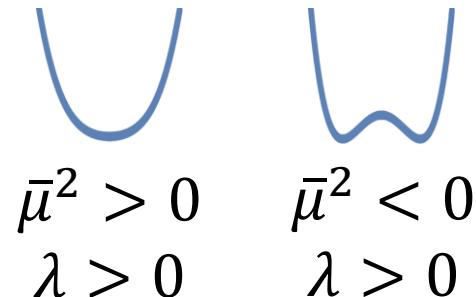
polynomial
running (QD)

Scalar couplings

$$\Gamma_k \supset V_{eff,k} = \text{curvature mass} + \lambda_k (\text{quartic coupling})$$

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polynomial running (QD)

$$\overline{\text{Flow}}[V_{\text{eff}}] \times k^4 = \frac{1}{2} \text{Diagram } H, \mathcal{G}^\pm, \mathcal{G}^0 + \frac{1}{2} \text{Diagram } W^\pm, Z^0, A^\gamma + \frac{1}{2} \text{Diagram } G^a - \text{Diagram } c_G^a, c_{W^\pm}, c_{Z^0} - \frac{1}{2} \sum_{q,l} \text{Diagram } q, l$$

Why the fRG?

- Mass-dependent renormalization scheme

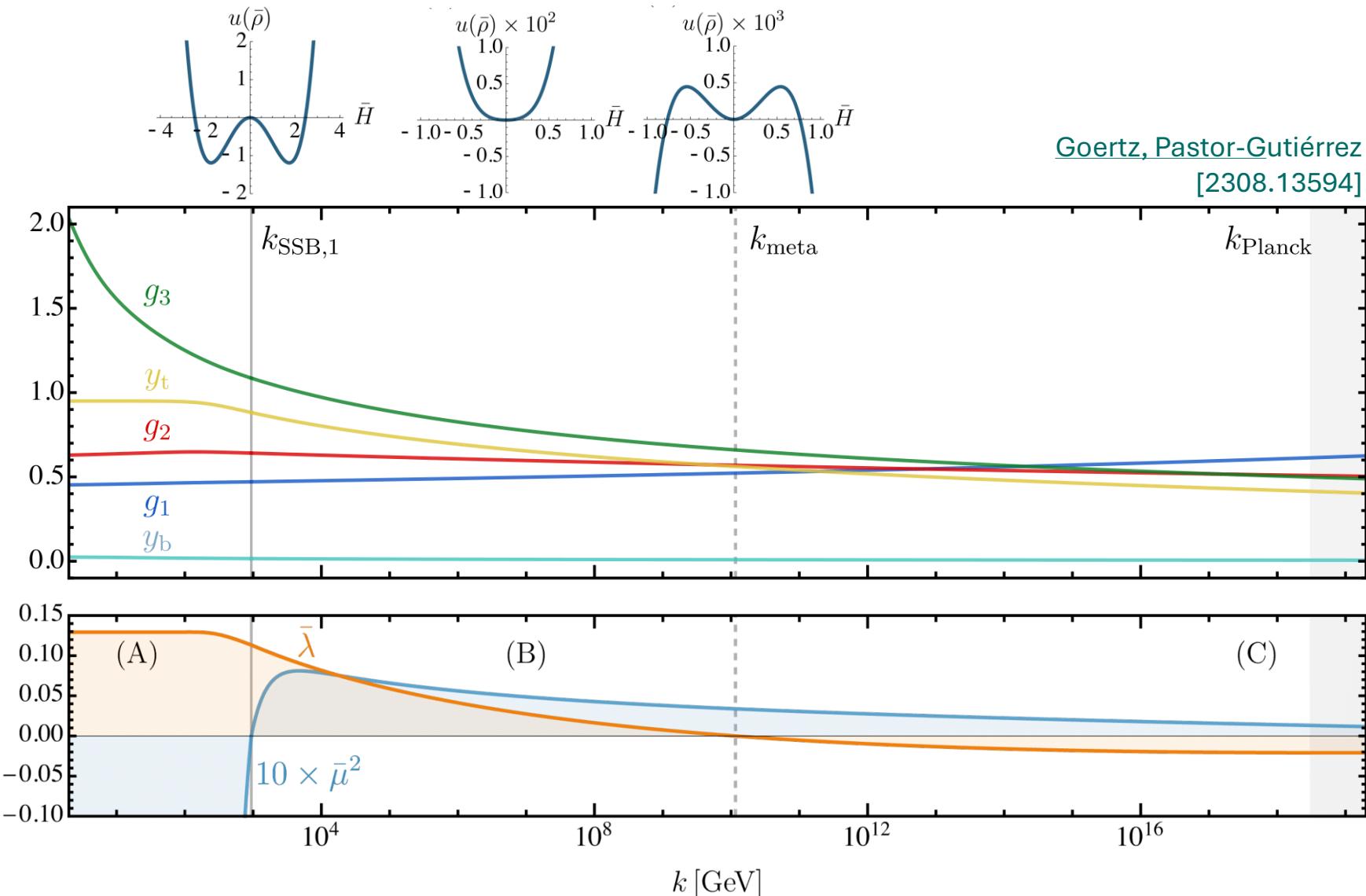
Why the fRG?

- Mass-dependent renormalization scheme
 - Incorporates QDs, the main source of fine-tuning, smoothly along the flow
 - Model-independent study of fine-tuning
 - Allows access to the phase structure and critical surface

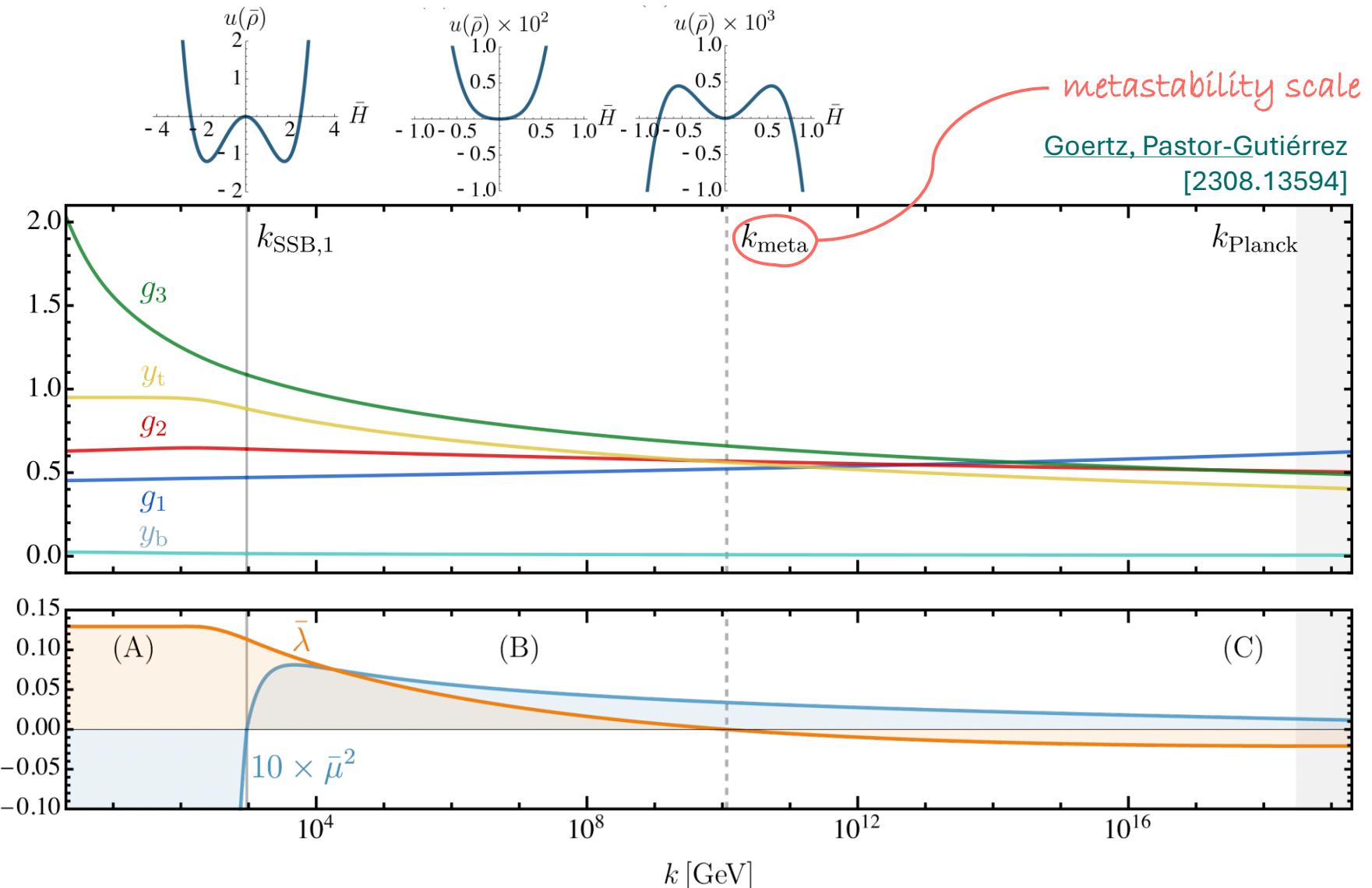
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- Possibility to explore non-perturbative physics
 - E.g. large anomalous dimensions

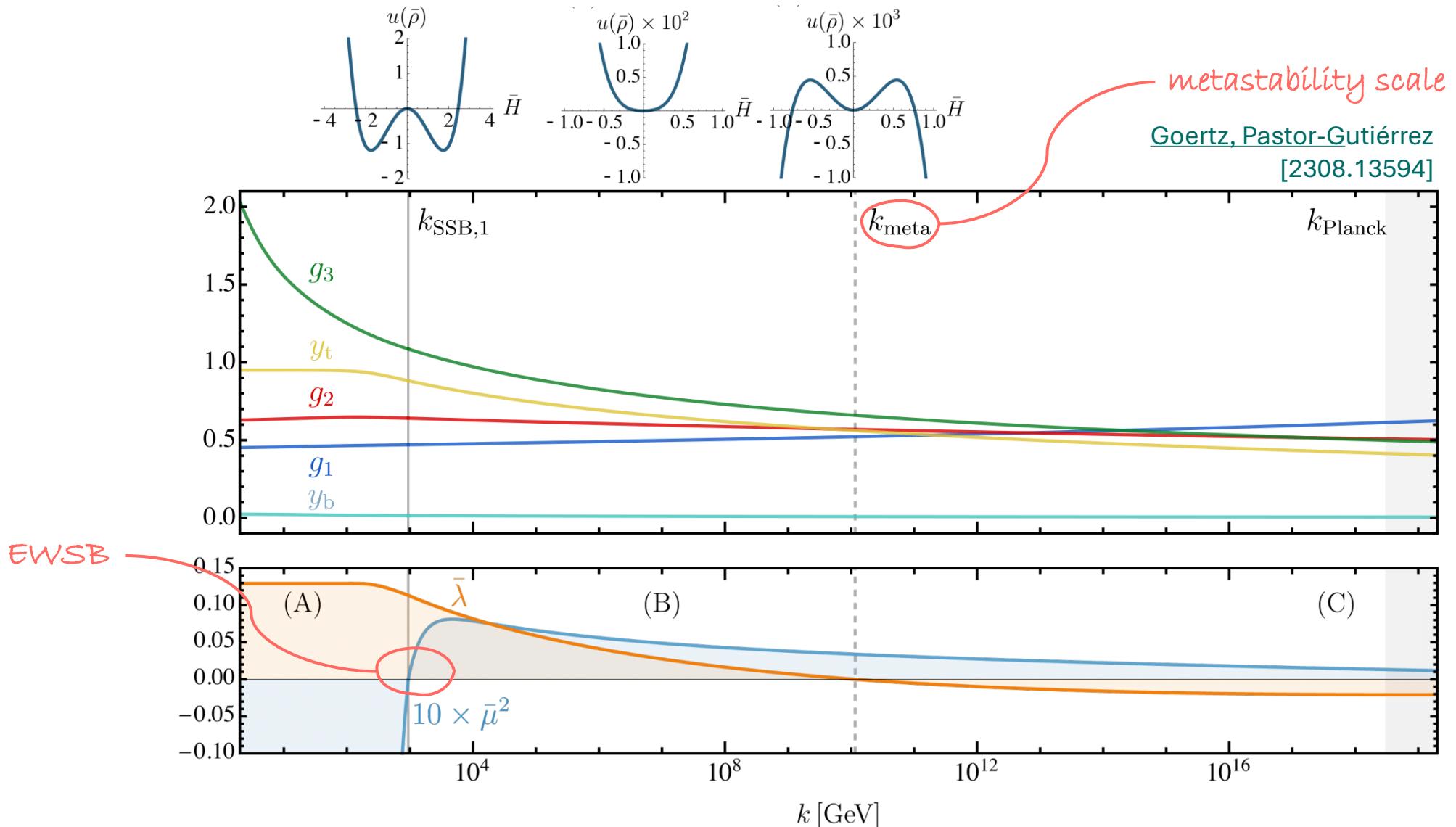
The Standard Model with fRG



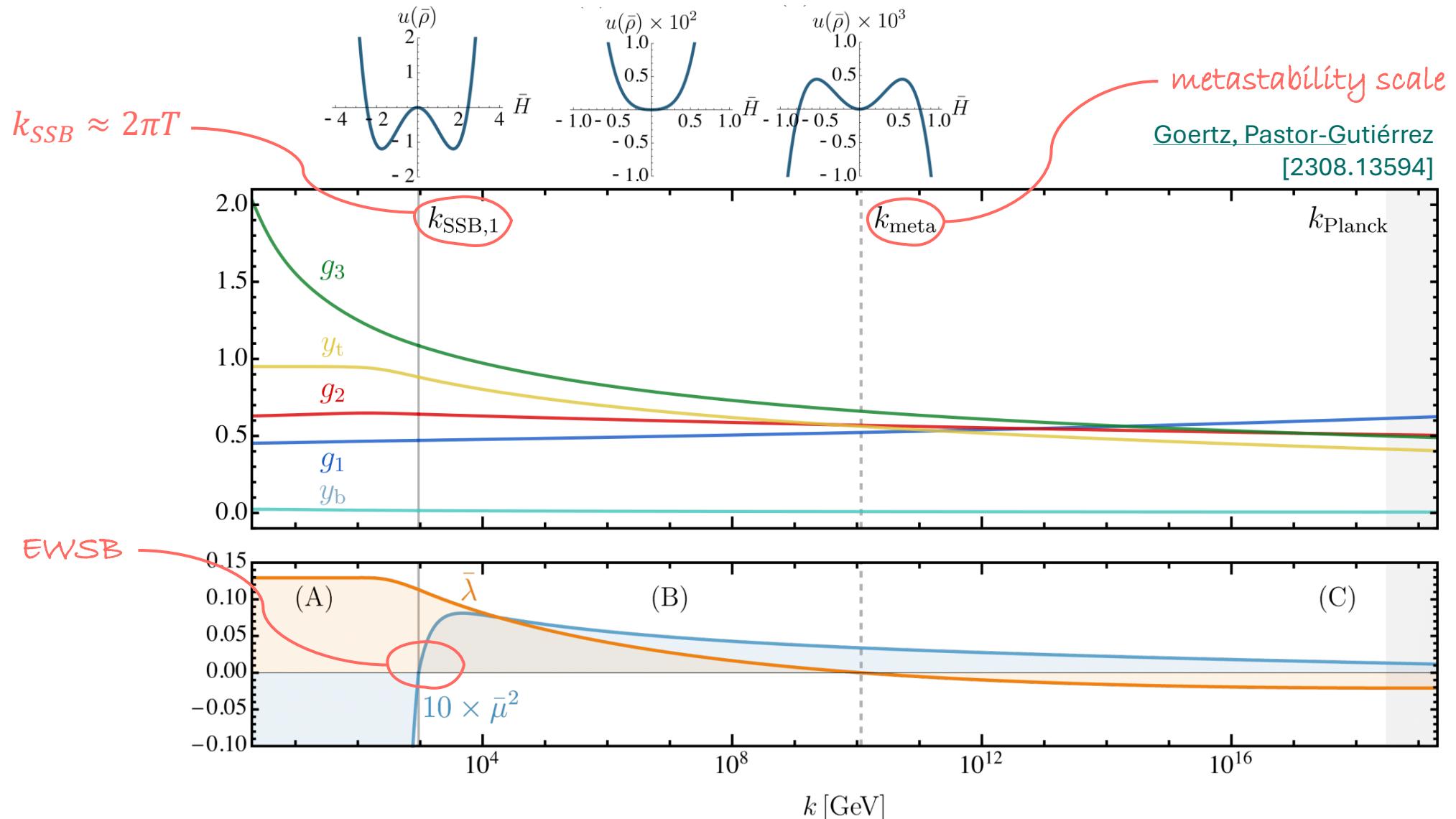
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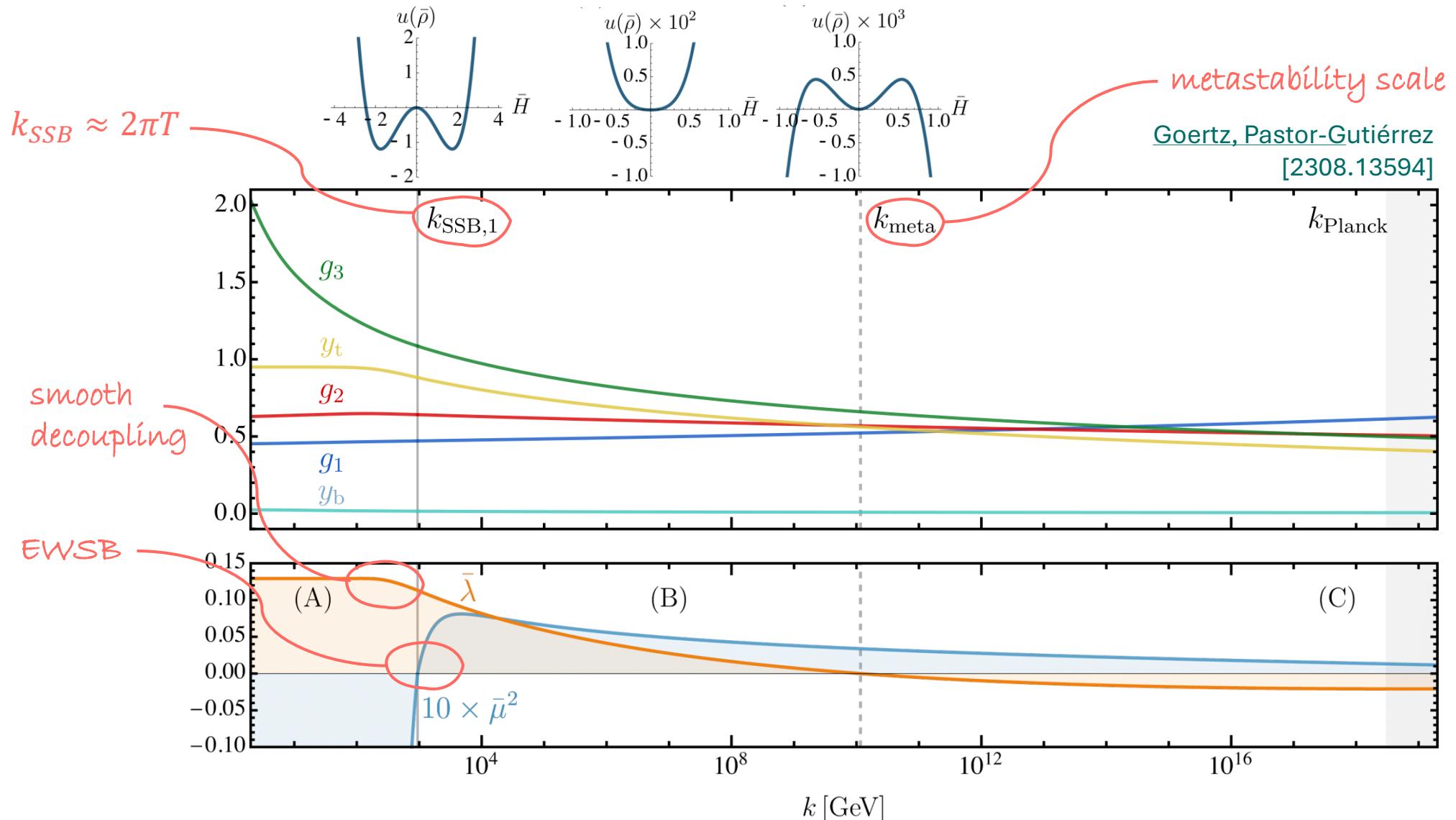
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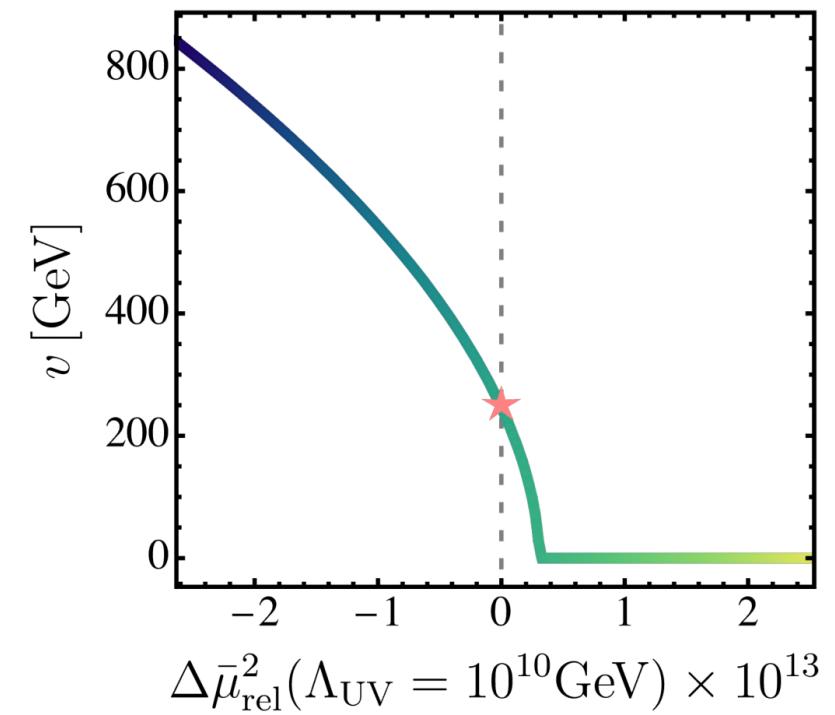
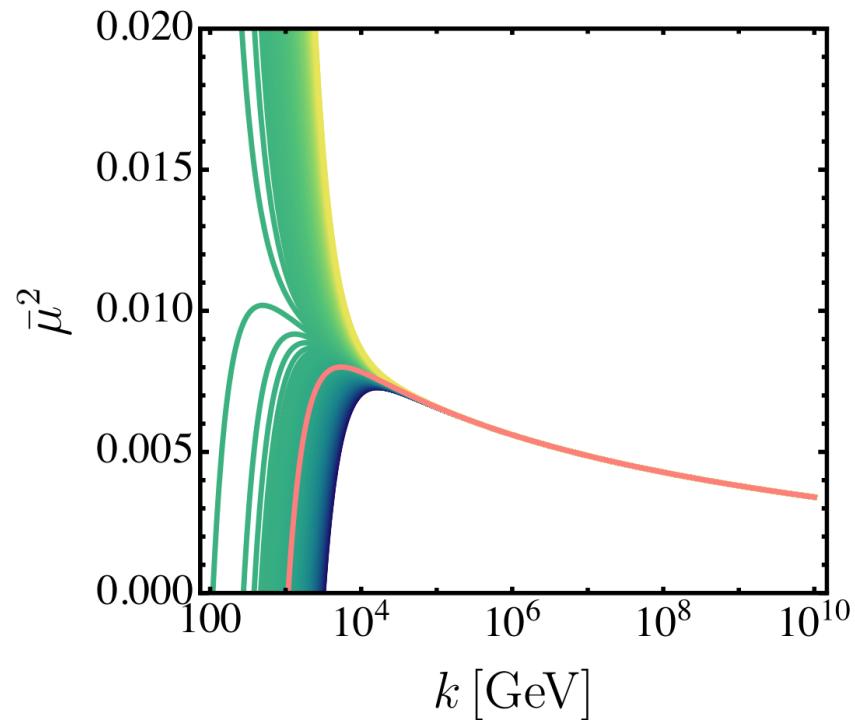


The fine-tuning problem with fRG

$$\bar{\mu}^2(\Lambda_{UV}) = \frac{\bar{\mu}_{\Lambda_{UV}}^2 - \bar{\mu}_{\Lambda_{UV,SM}}^2}{\bar{\mu}_{\Lambda_{UV}}^2}$$

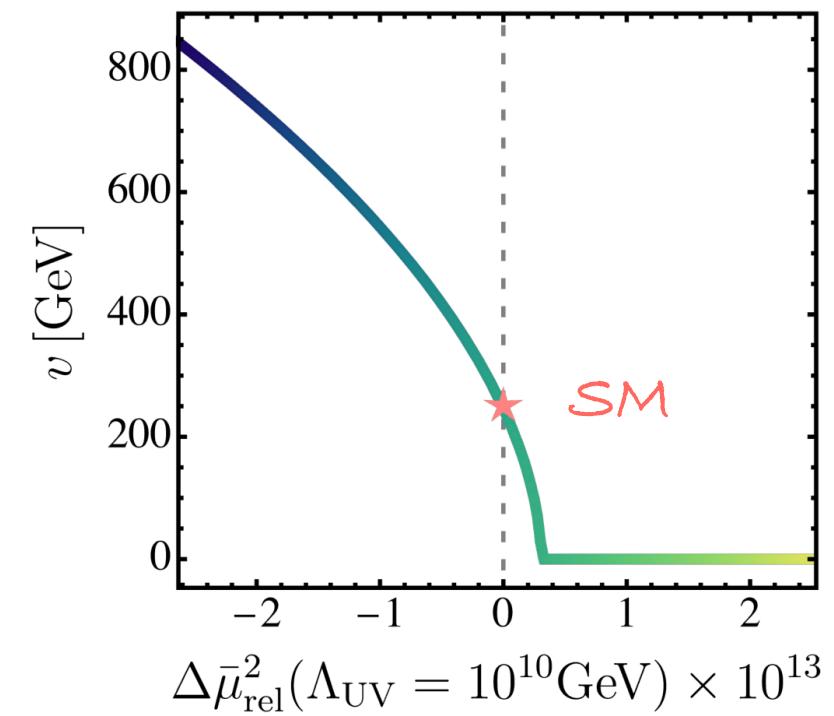
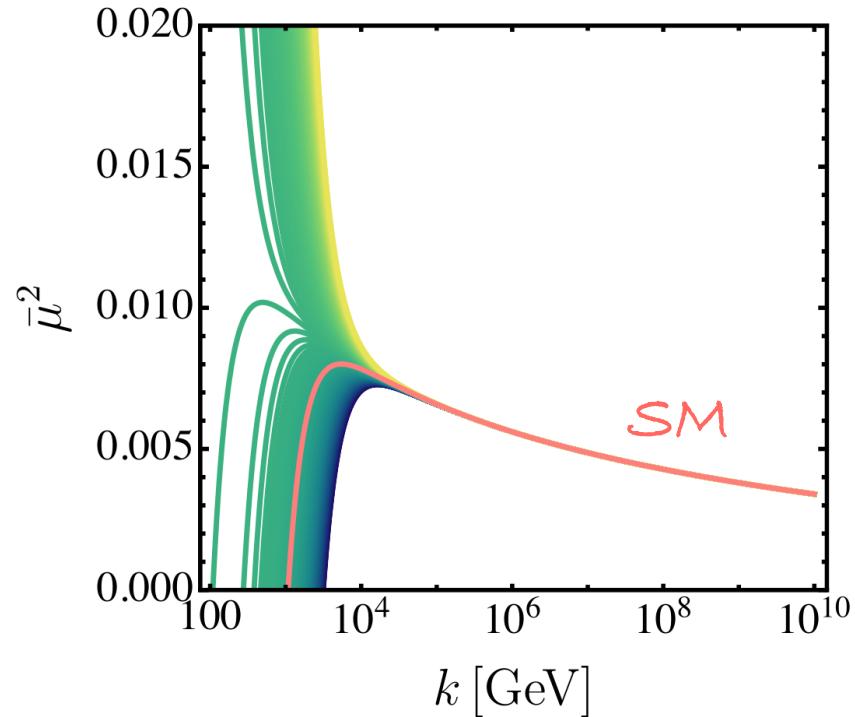
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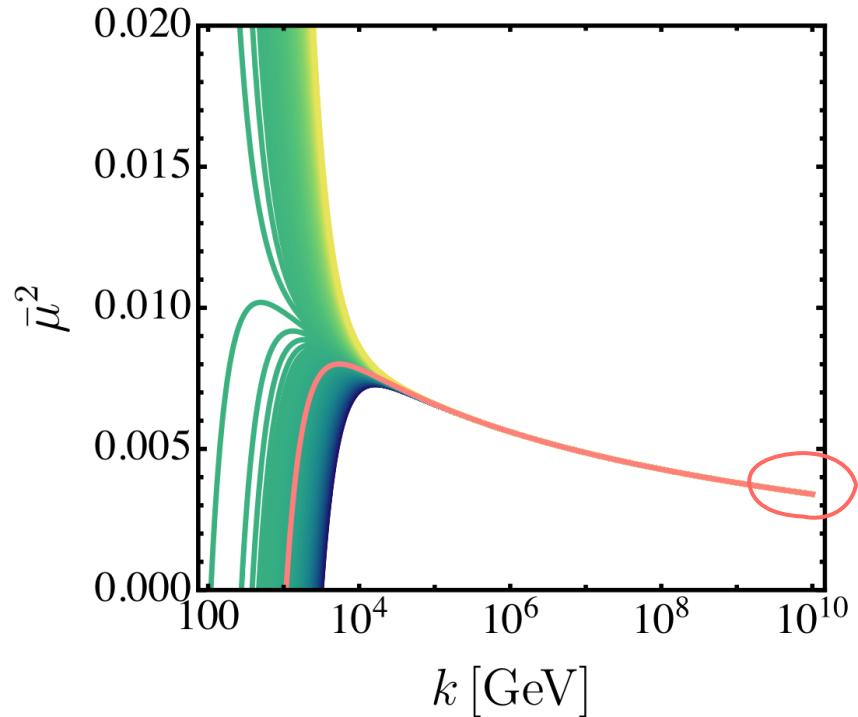
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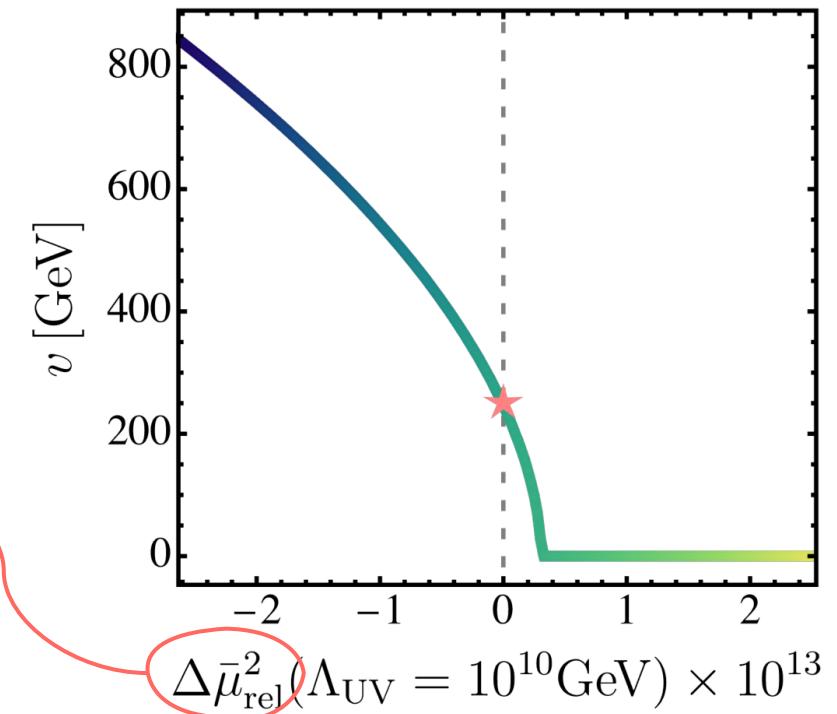


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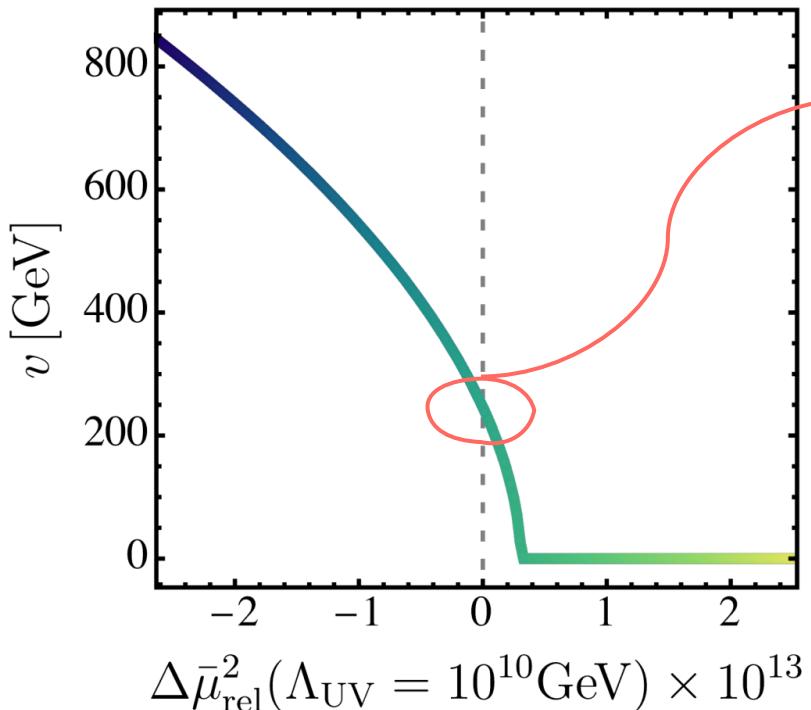
$$\bar{\mu}^2(\Lambda_{UV}) = \frac{\bar{\mu}_{\Lambda_{UV}}^2 - \bar{\mu}_{\Lambda_{UV,SM}}^2}{\bar{\mu}_{\Lambda_{UV}}^2}$$



$$\Delta \bar{\mu}^2(\Lambda_{UV}) = \mathcal{O}\left(\frac{v^2}{\Lambda_{UV}^2}\right)$$

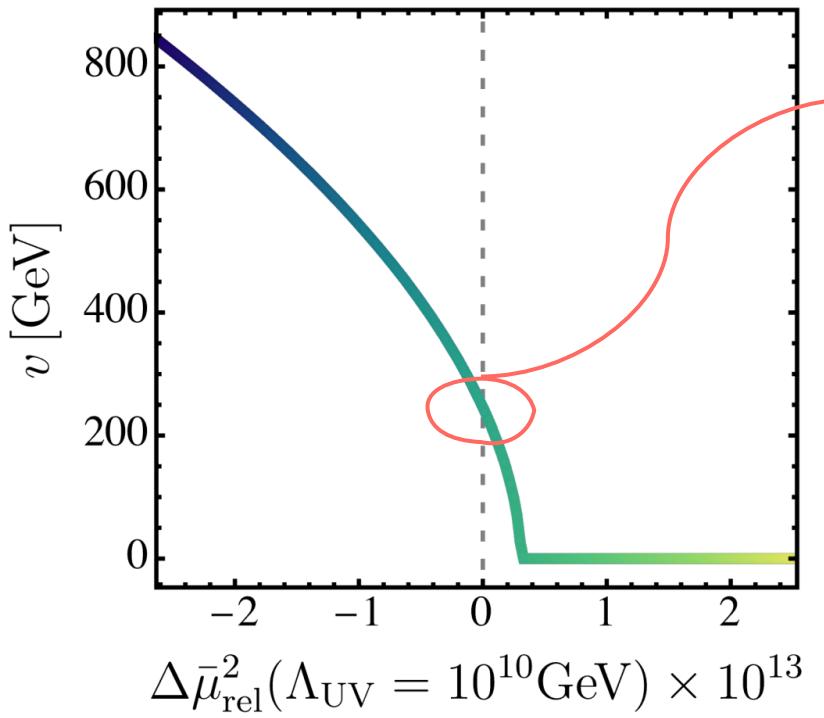


Quantifying the fine-tuning



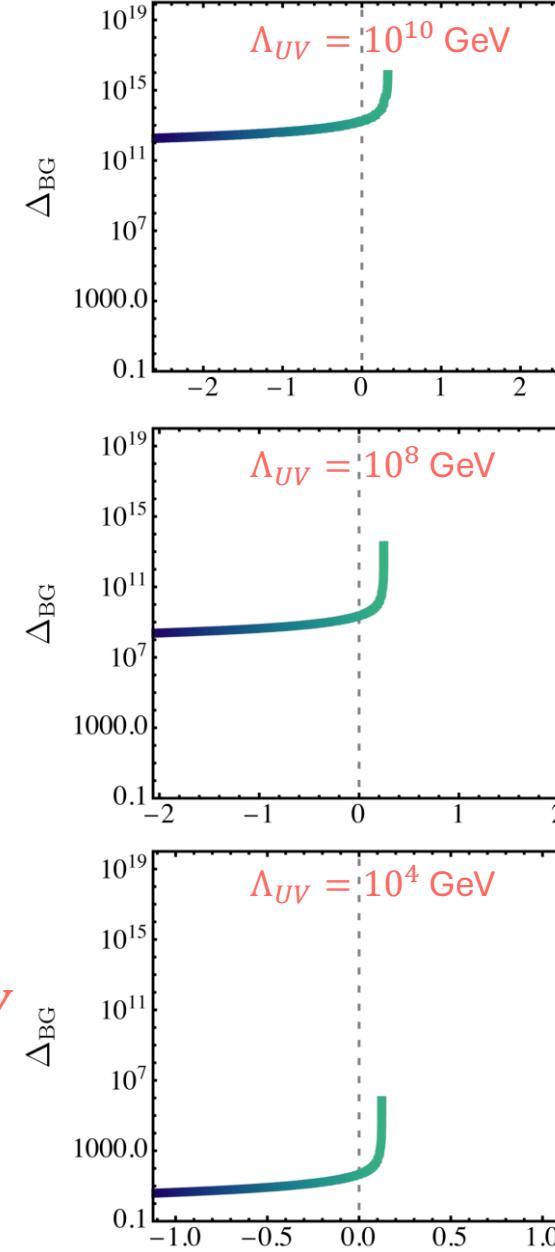
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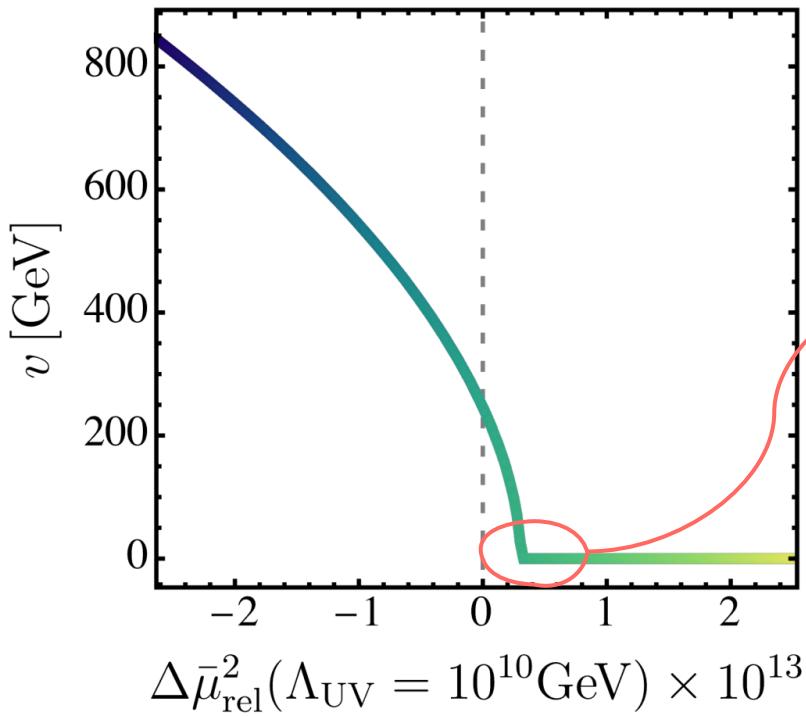
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Λ_{UV}

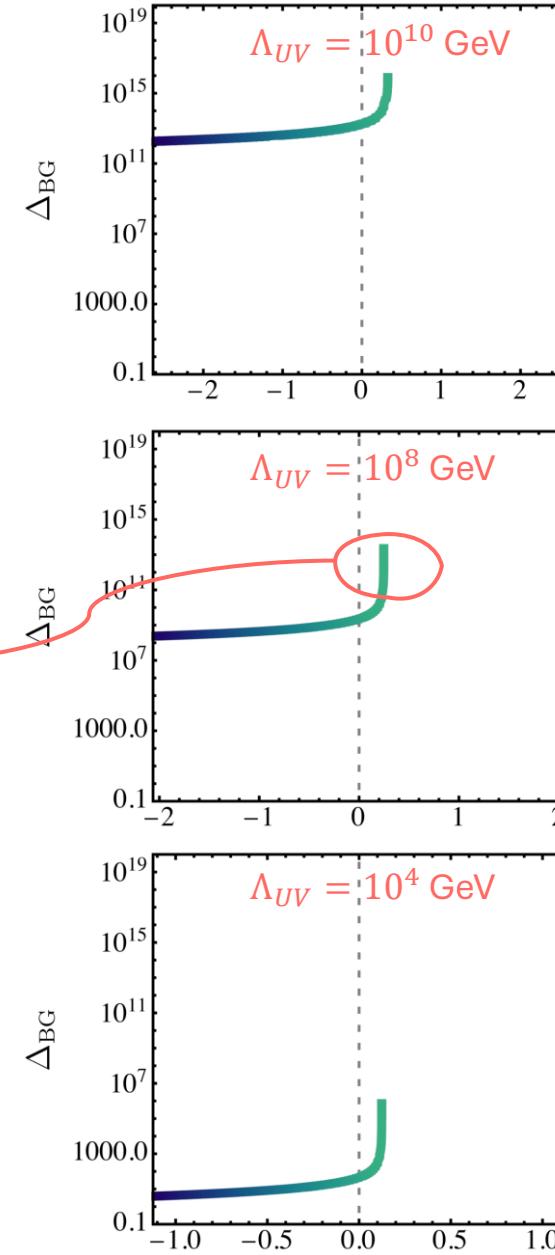


$$\Delta \bar{\mu}^2(\Lambda_{UV}) \times \frac{\Lambda_{UV}^2}{v^2}$$

Quantifying the fine-tuning



critical surface

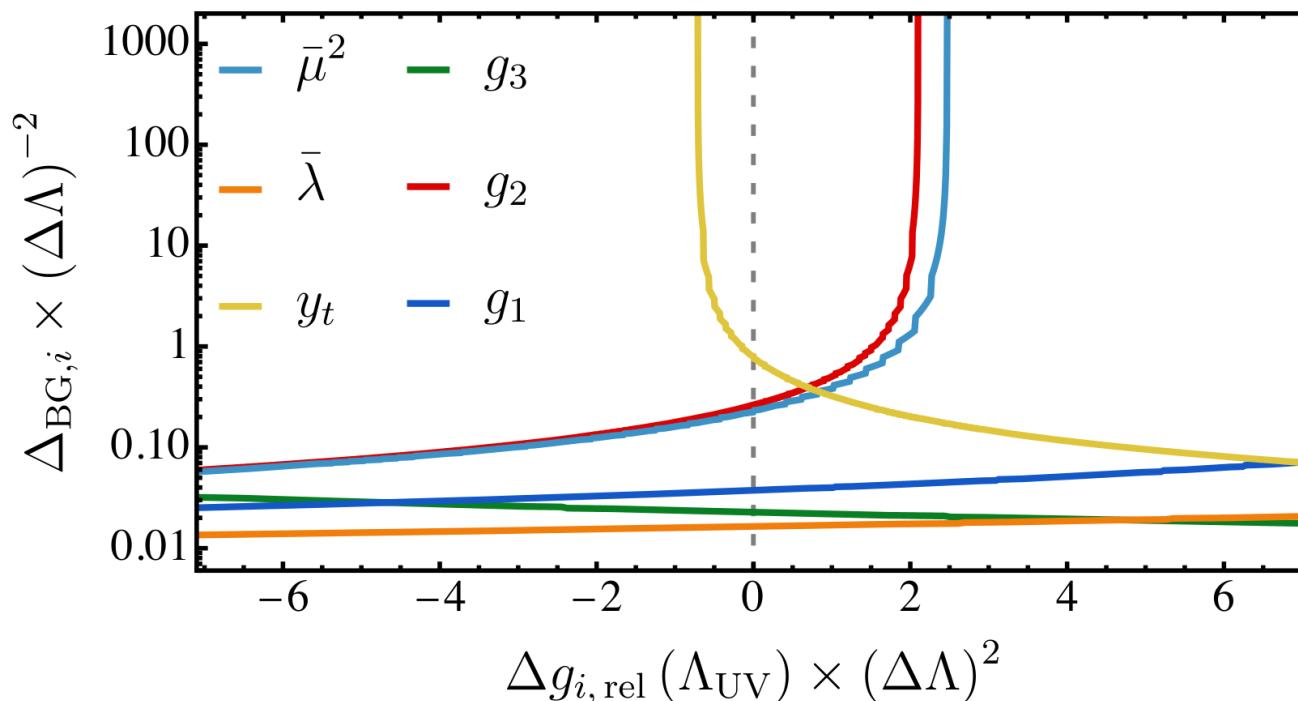


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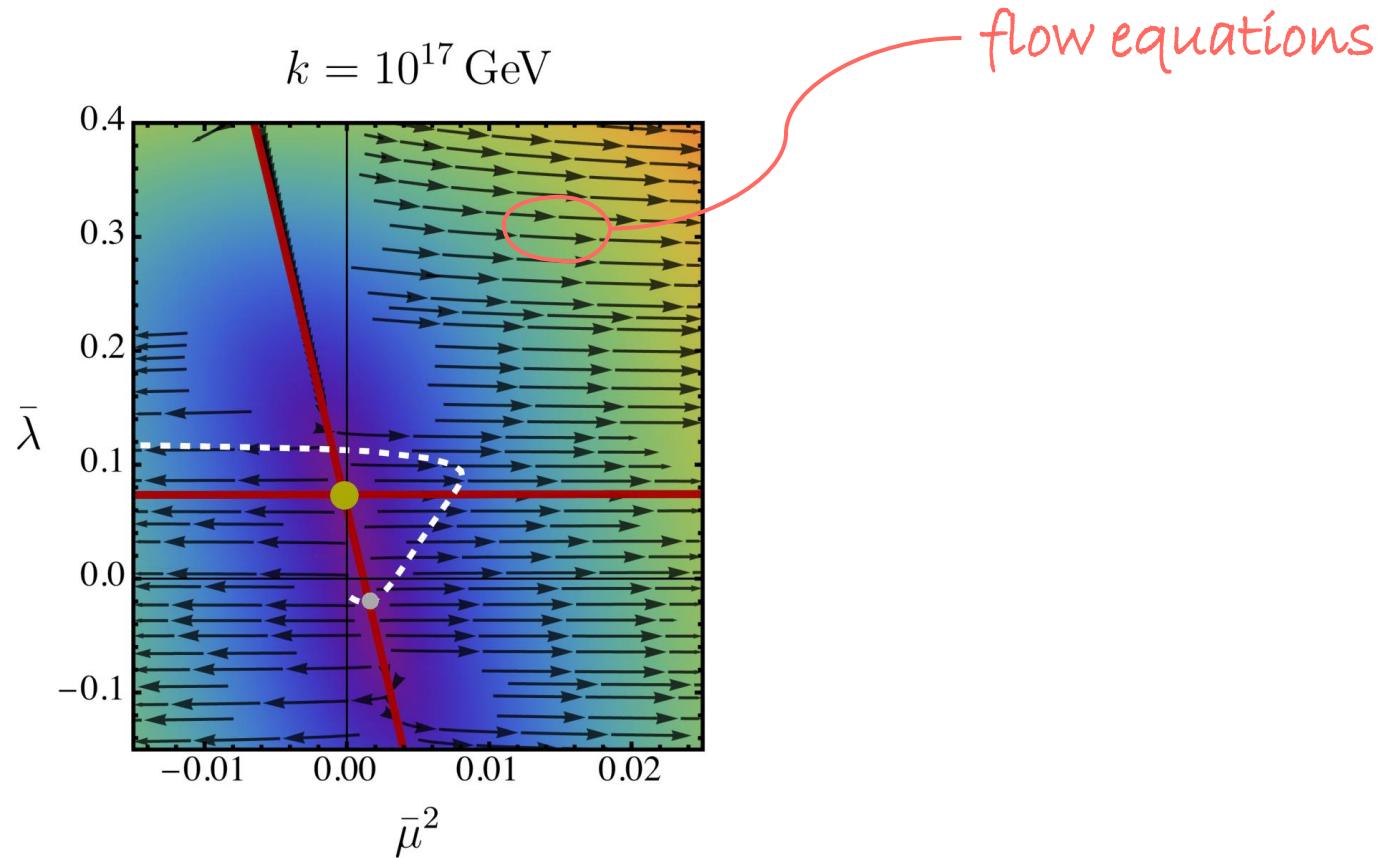
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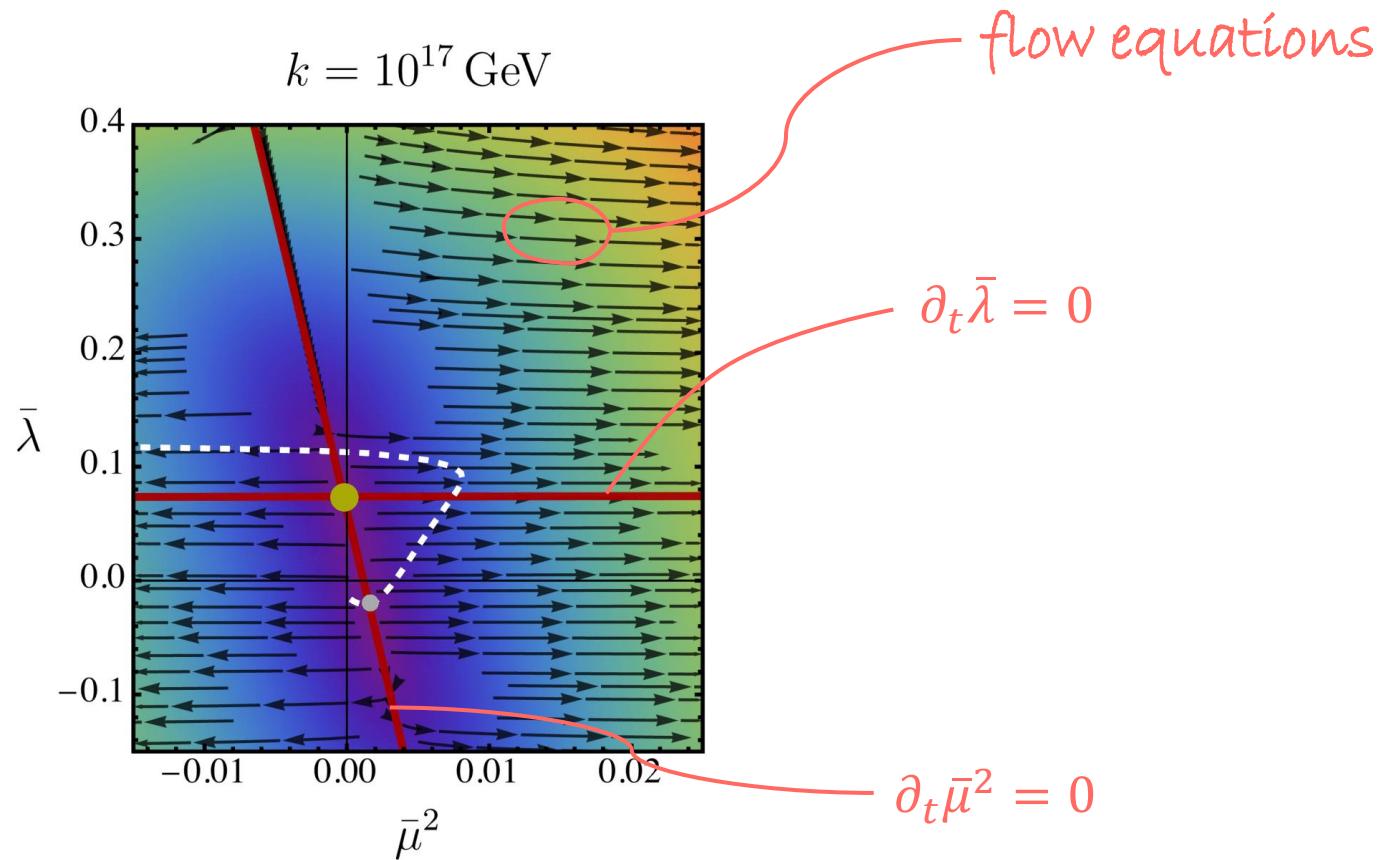
$$\Delta_{BG} = \max_{i,j} \left| \frac{\delta \log o_i}{\delta \log g_j} \right|$$



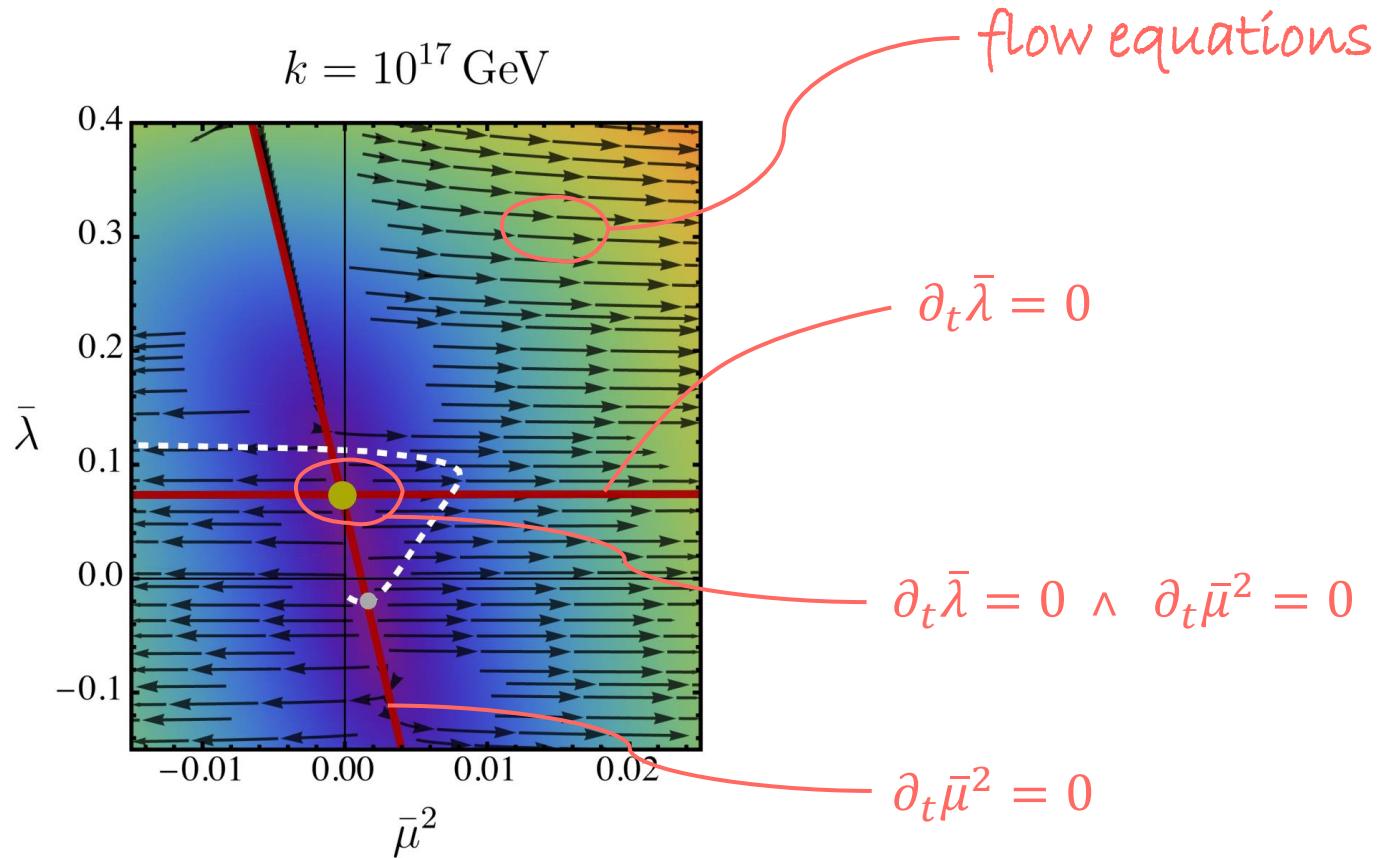
Phase diagram and critical surface



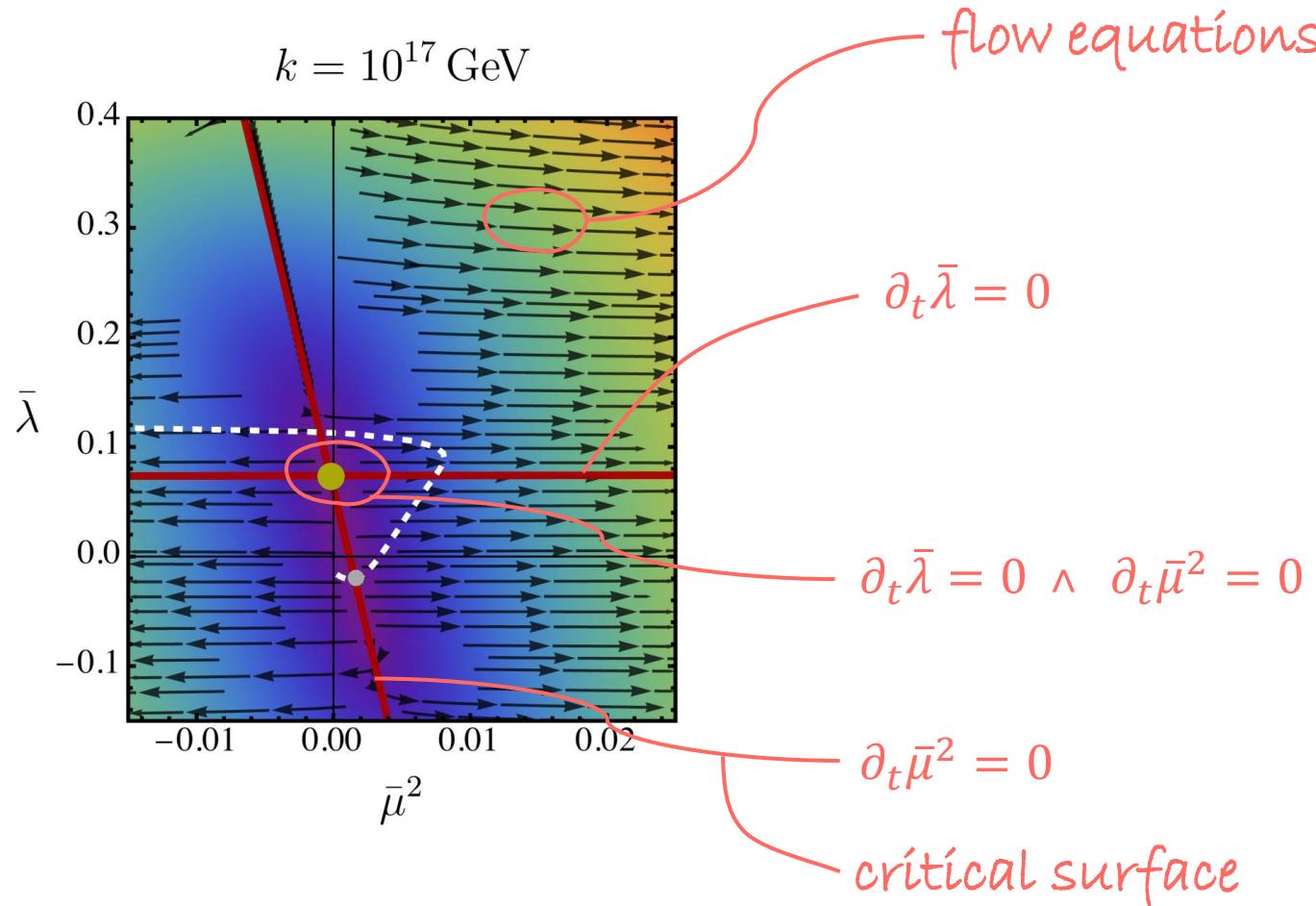
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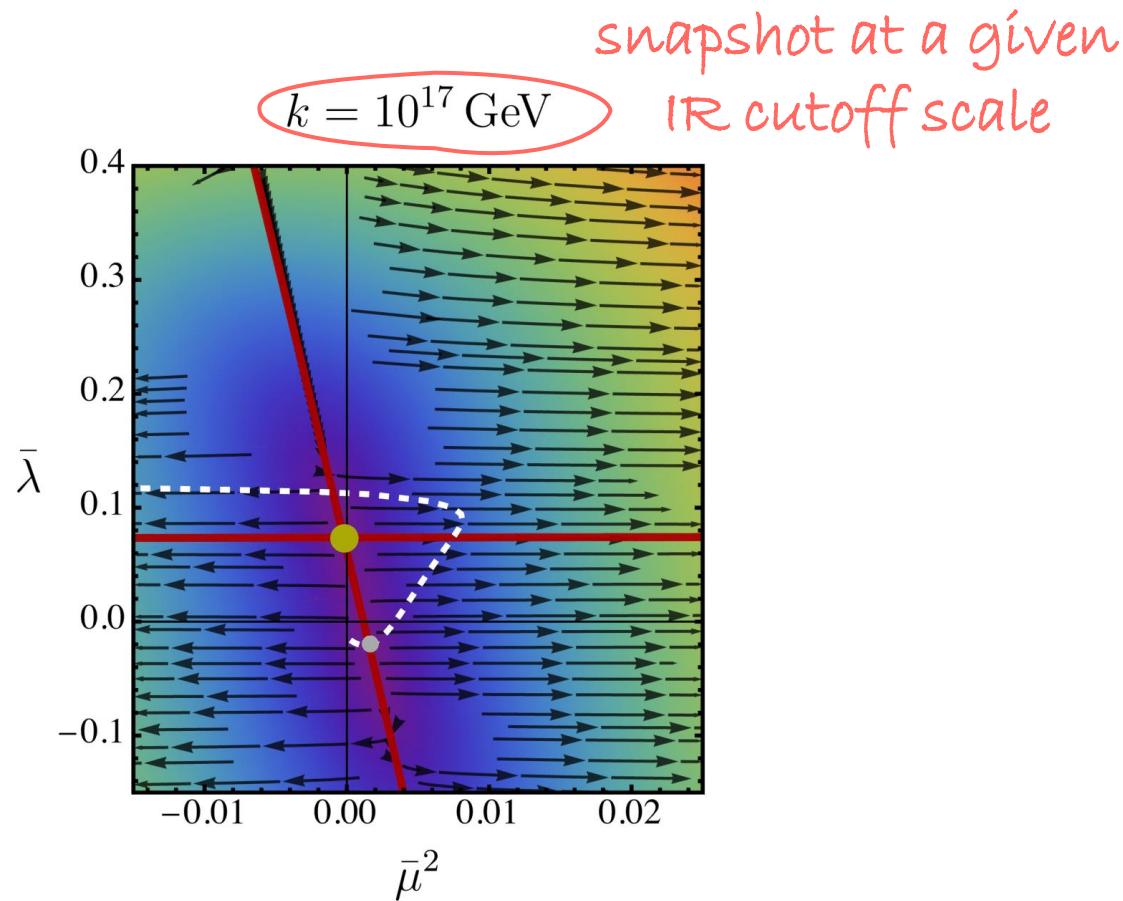
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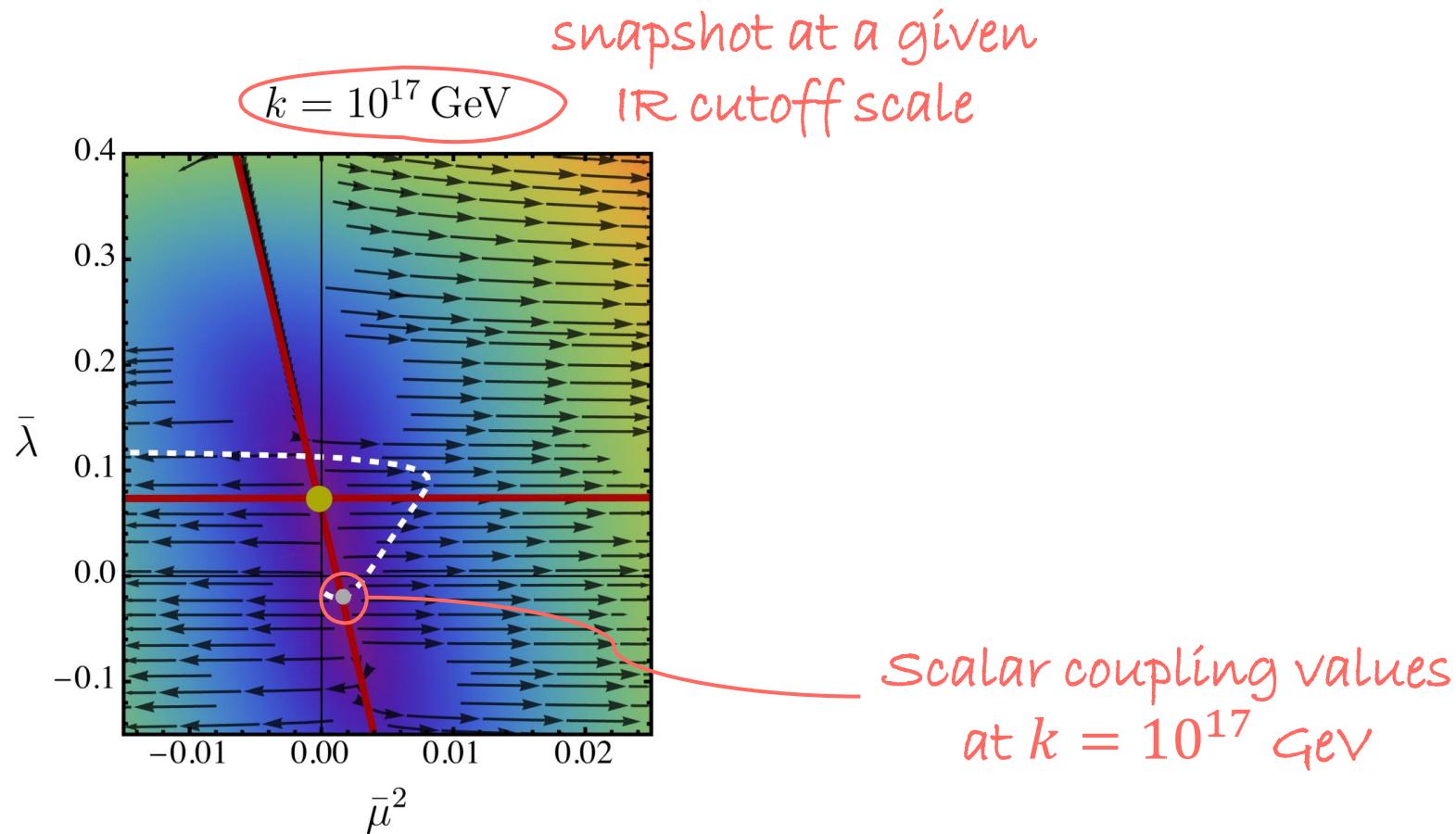
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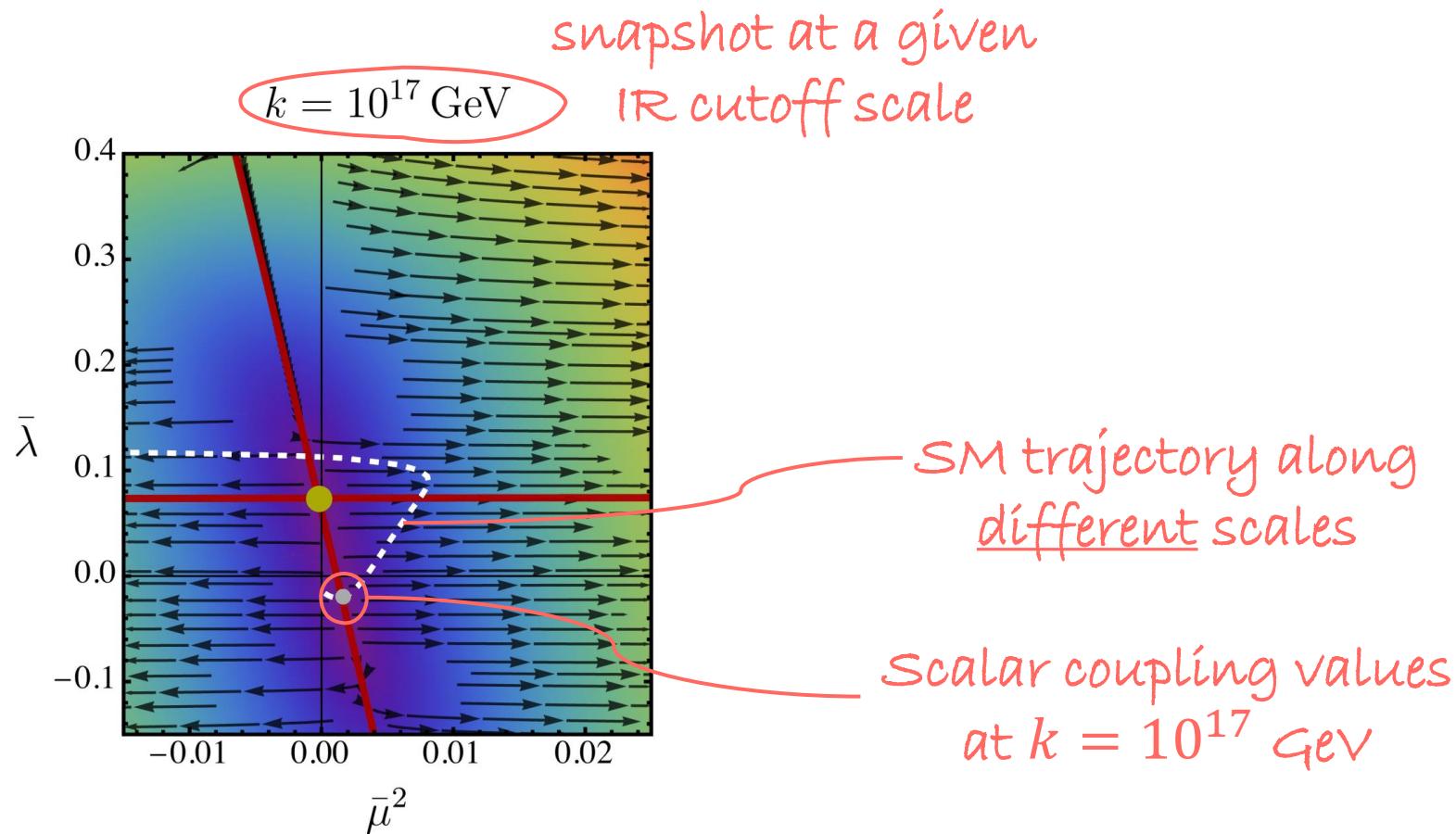
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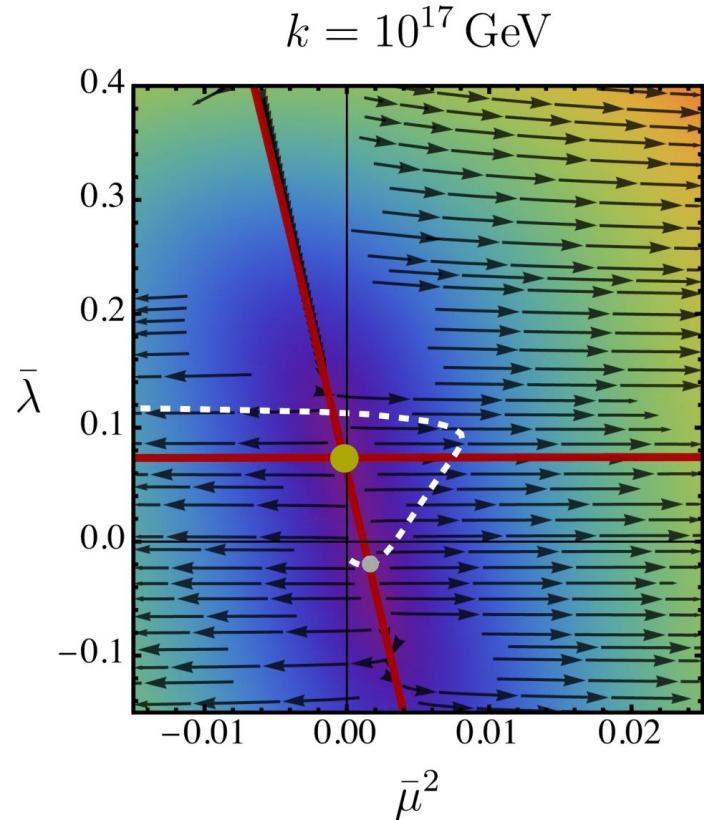
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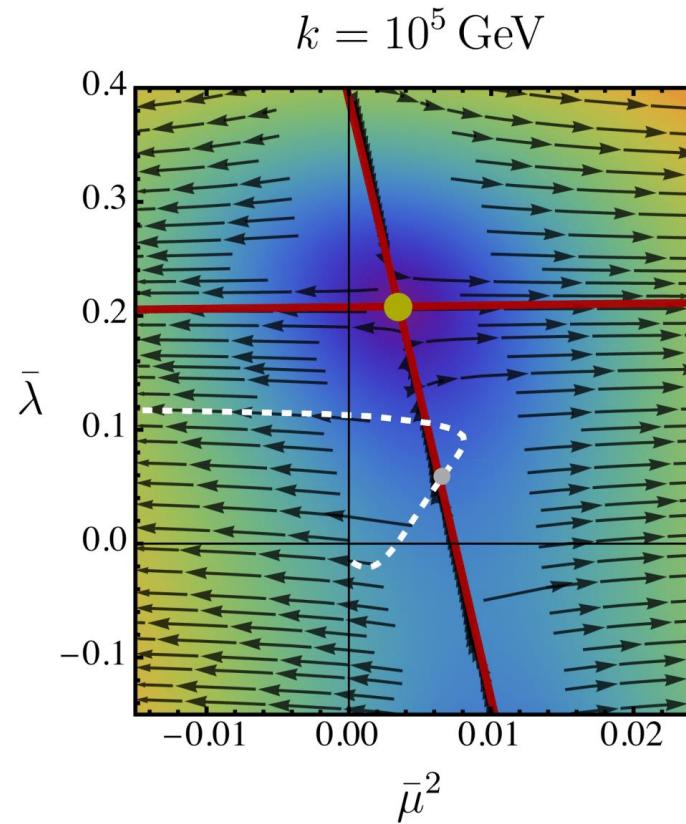
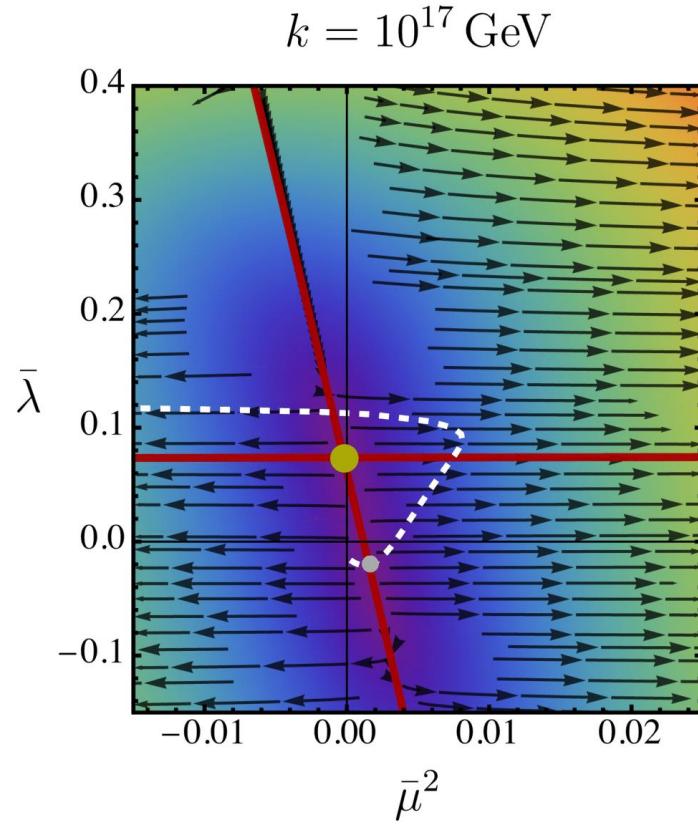
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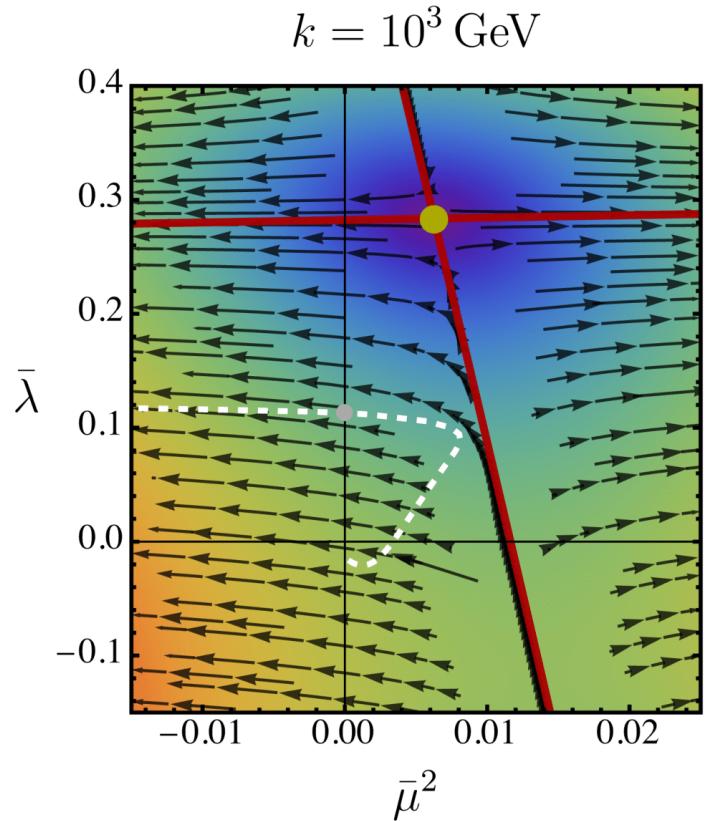
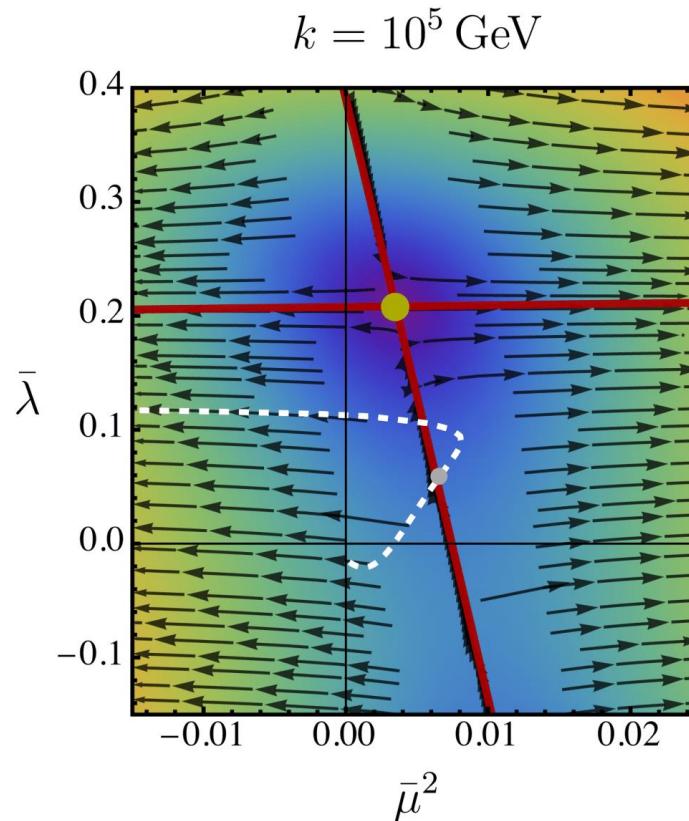
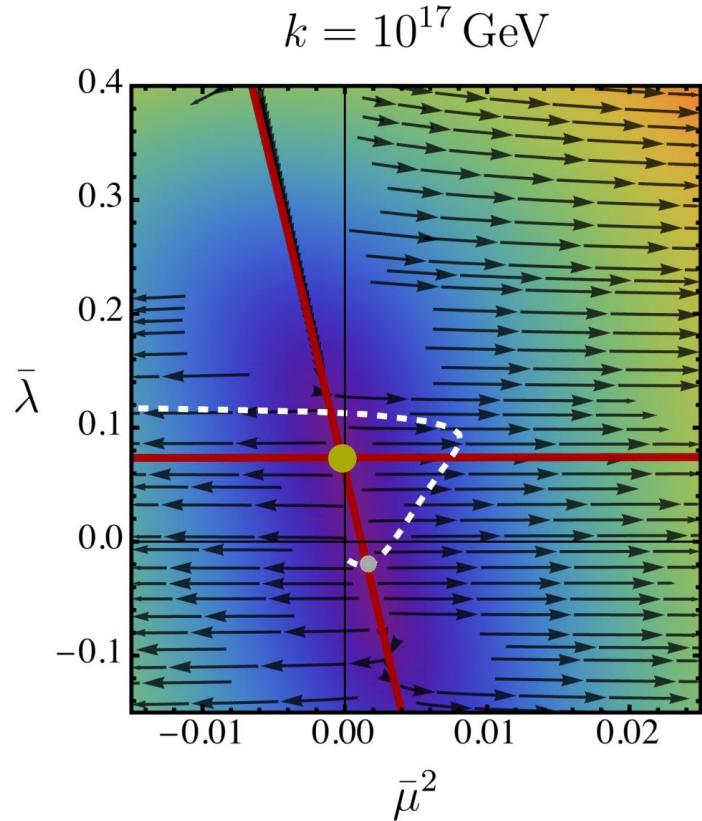
Phase diagram and critical surface



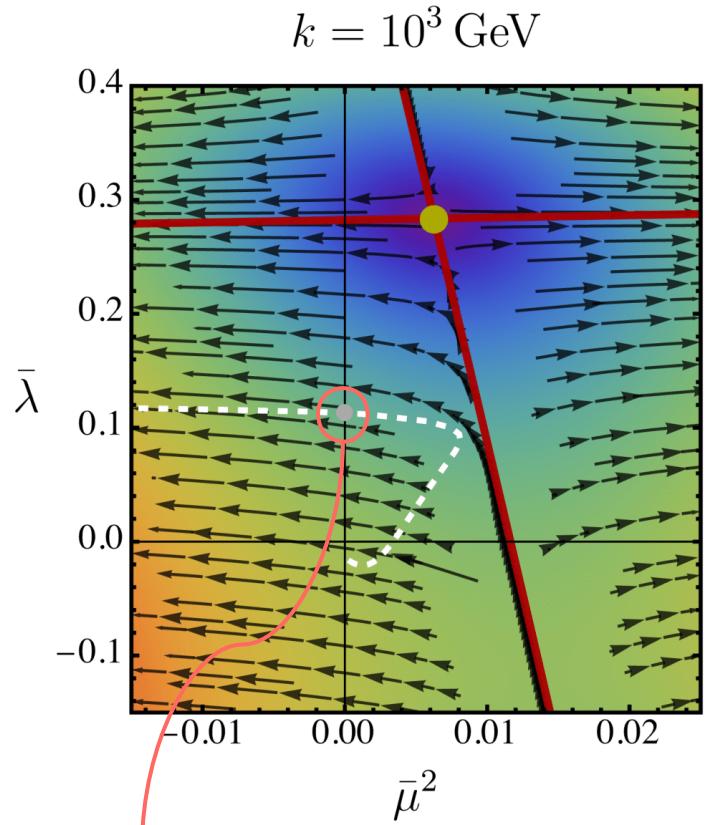
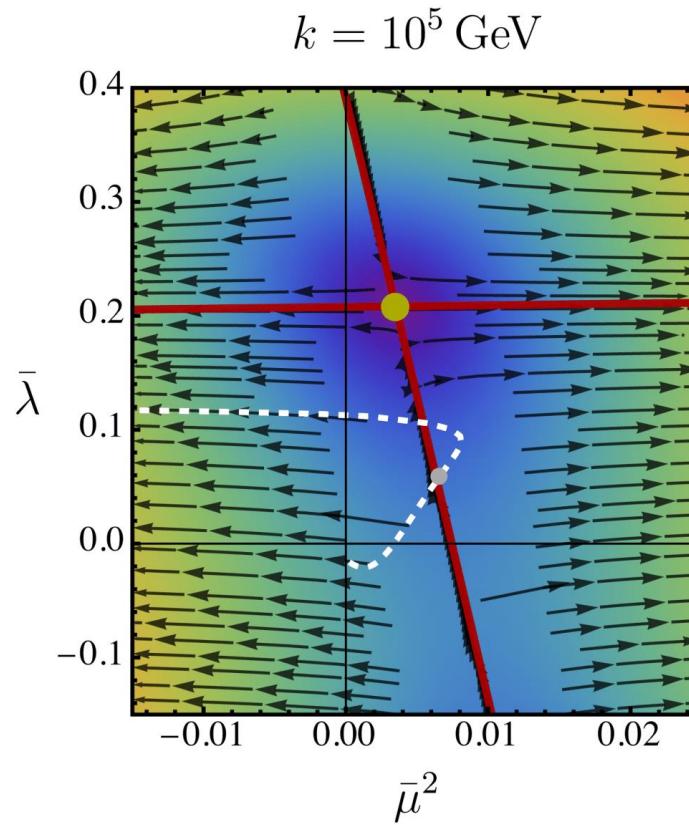
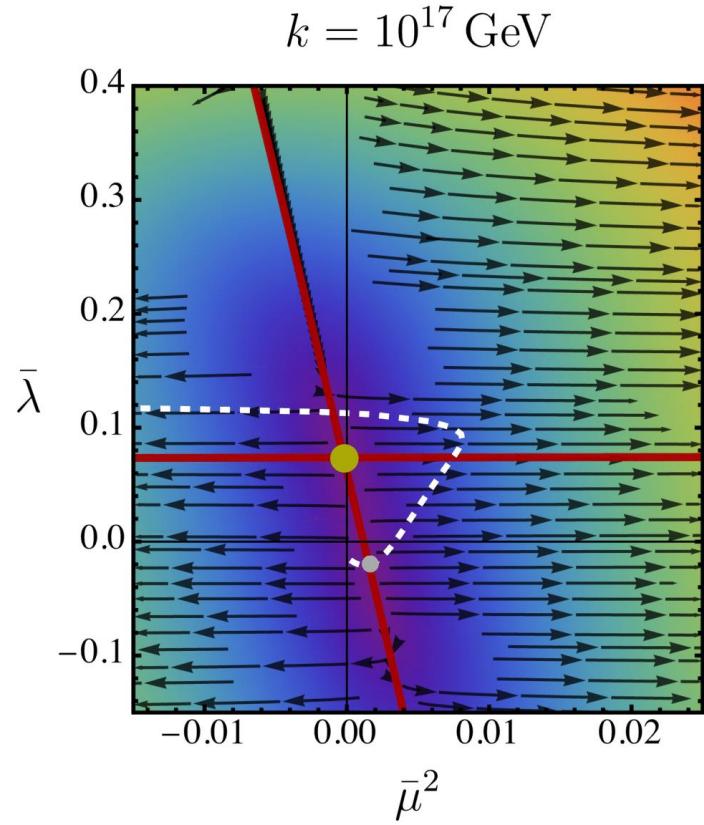
Phase diagram and critical surface



Phase diagram and critical surface

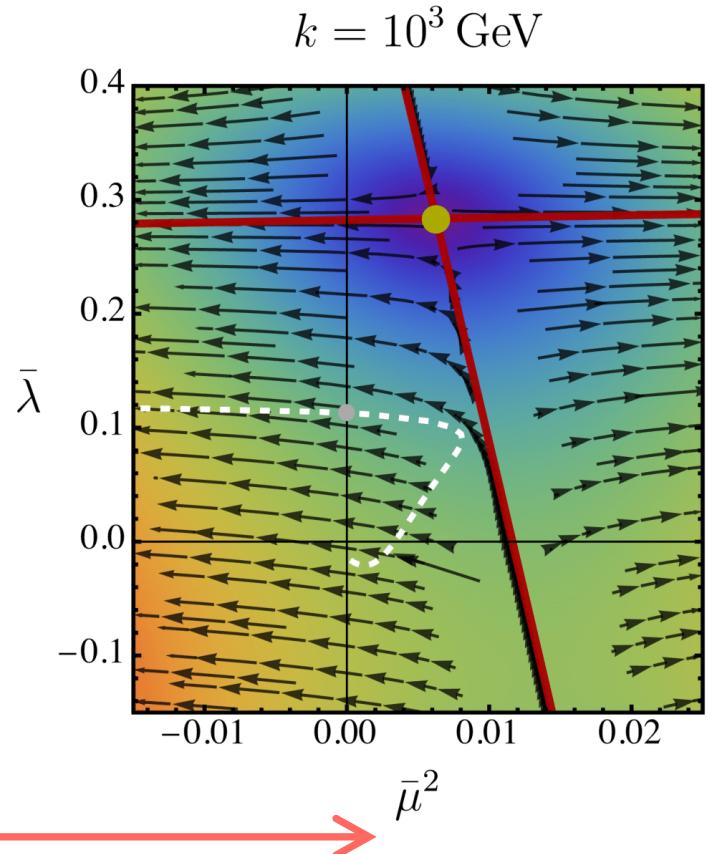
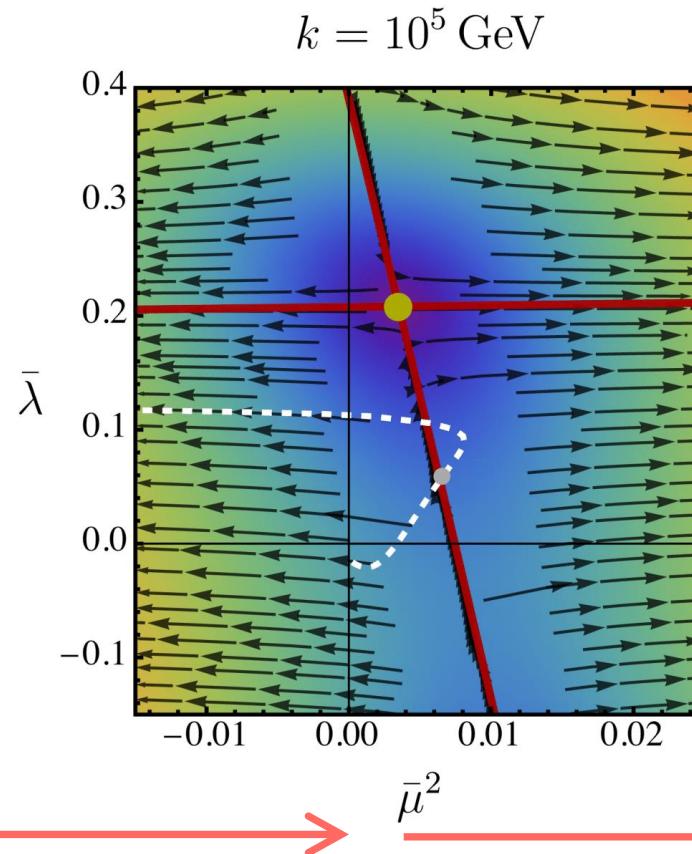
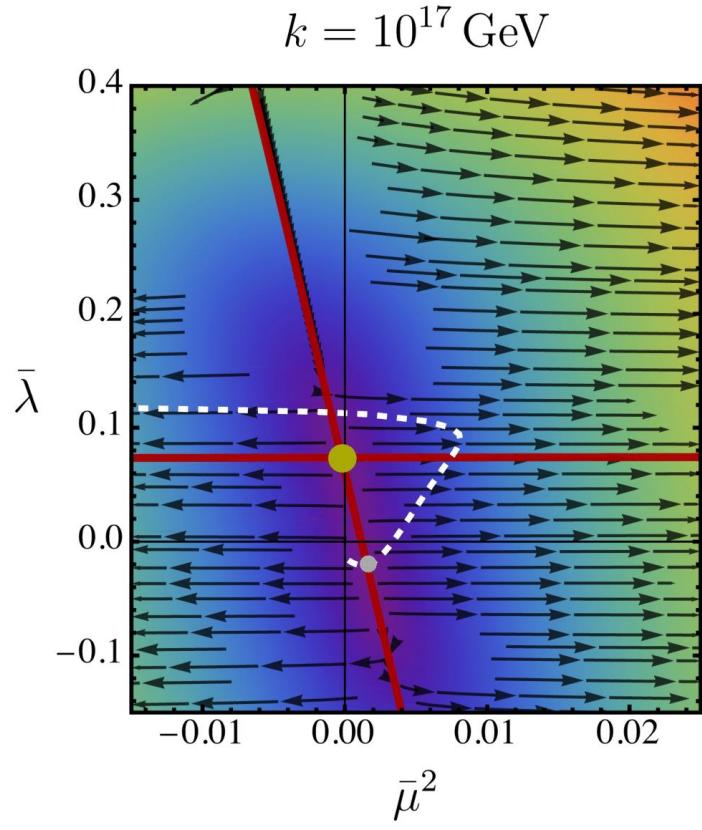


Phase diagram and critical surface



EWSB

Phase diagram and critical surface



$$\frac{k_1}{k_2} = 10^{12}$$

$$\frac{k_2}{k_3} = 10^2$$

Quantum criticality

1. Fine-tuning due to the magnitude of Λ_{UV}

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2. Fine-tuning due to the distance between the theory at Λ_{UV} w.r.t the critical surface

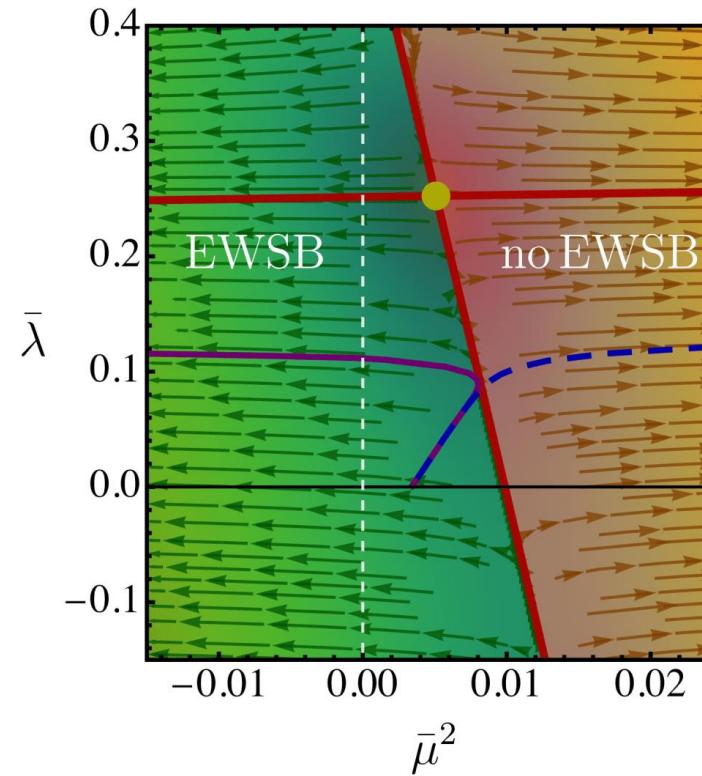
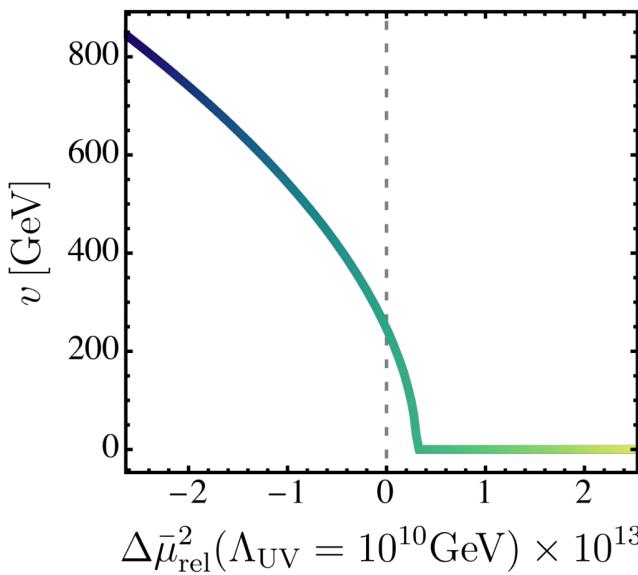
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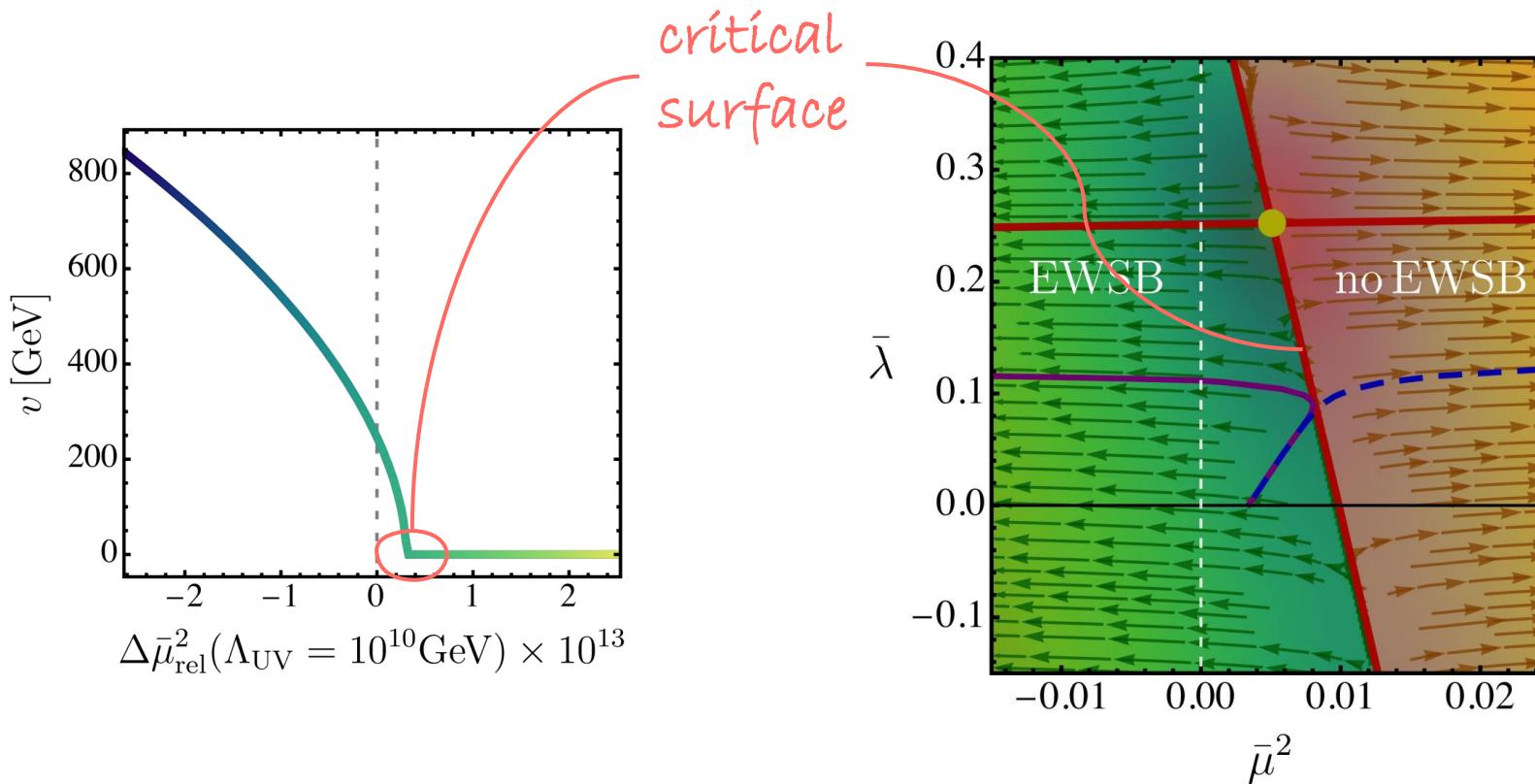
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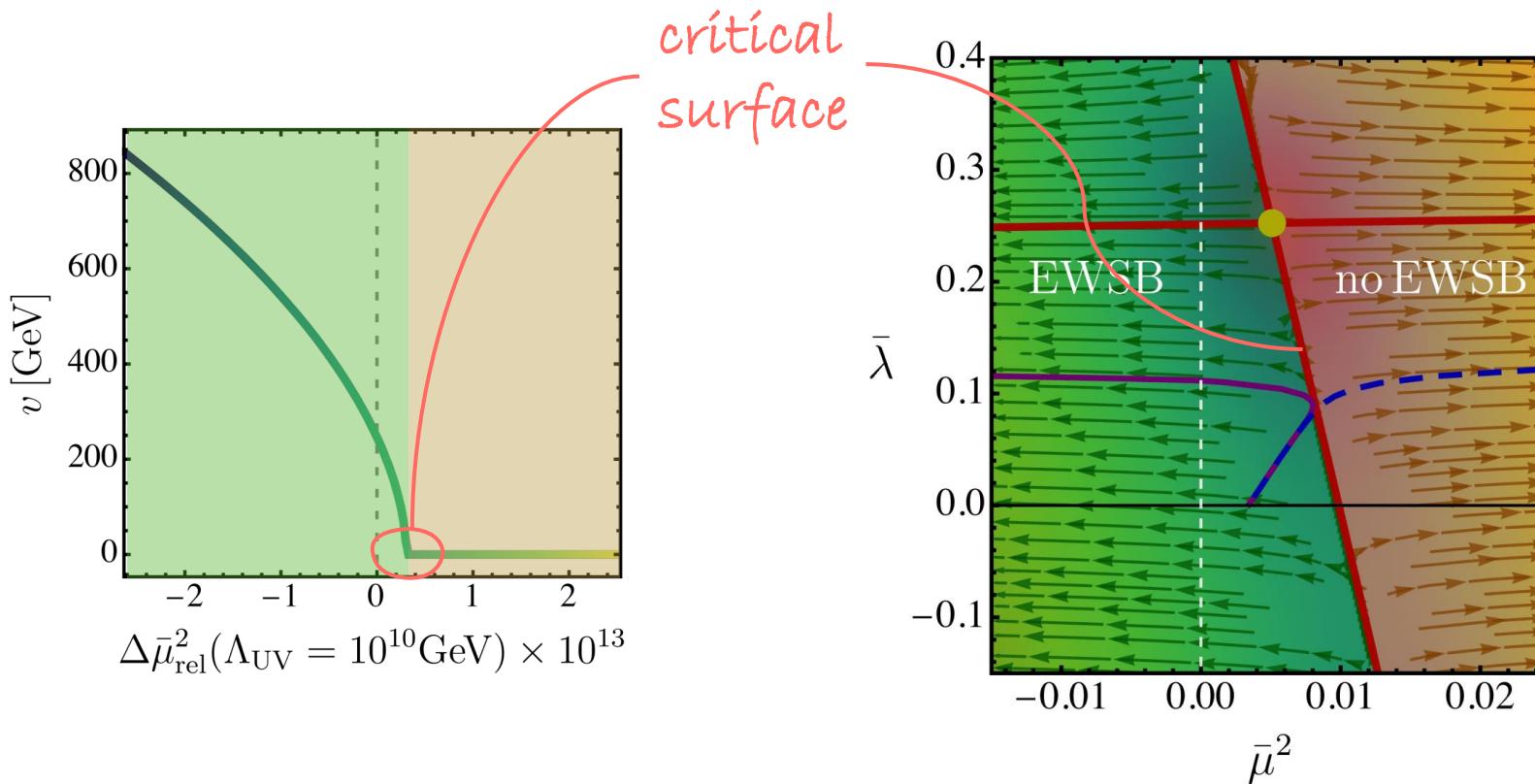
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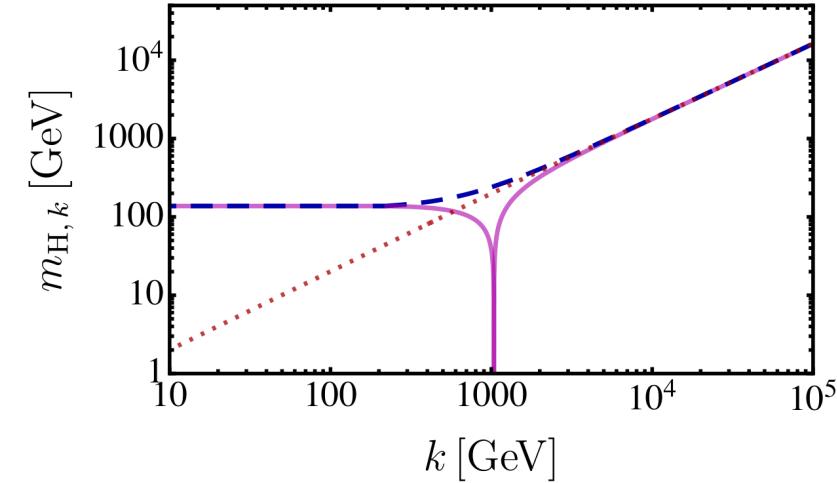
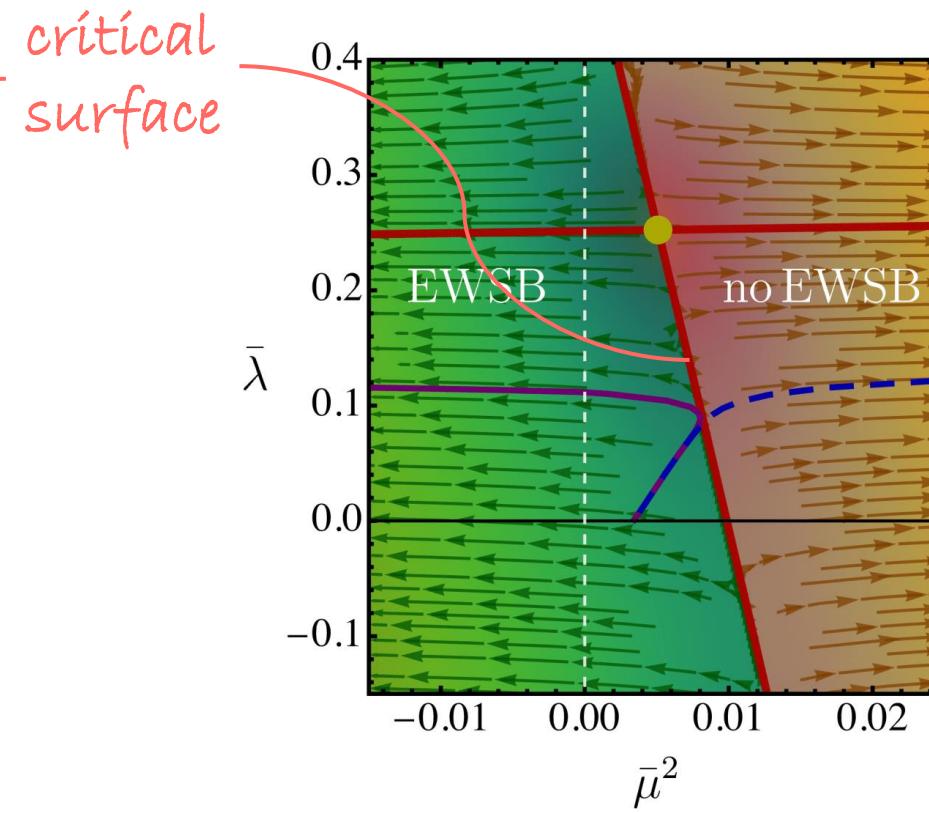
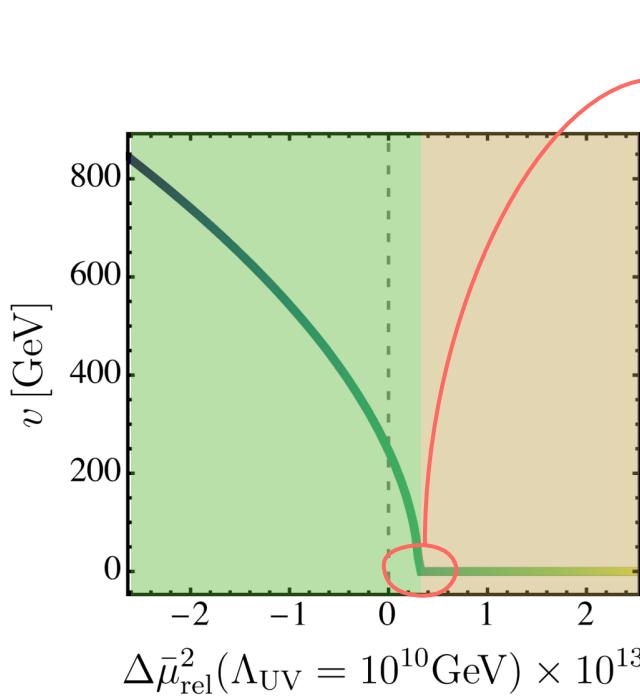
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Quantum criticality

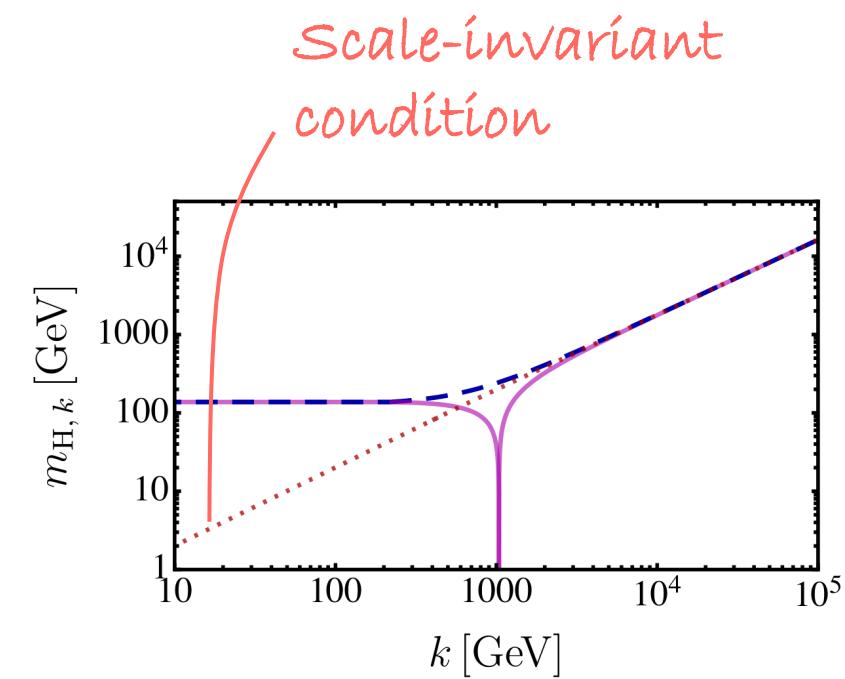
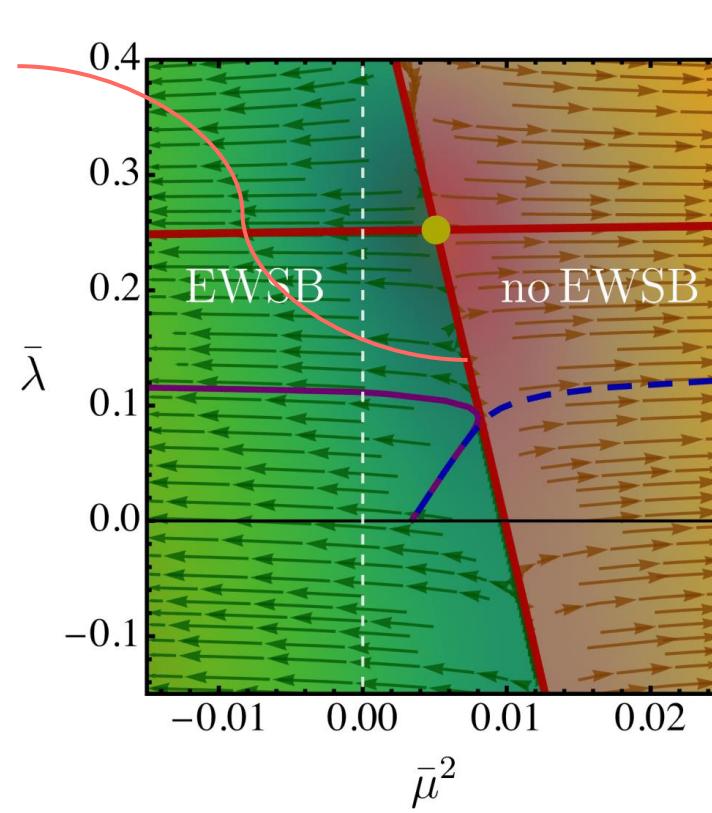
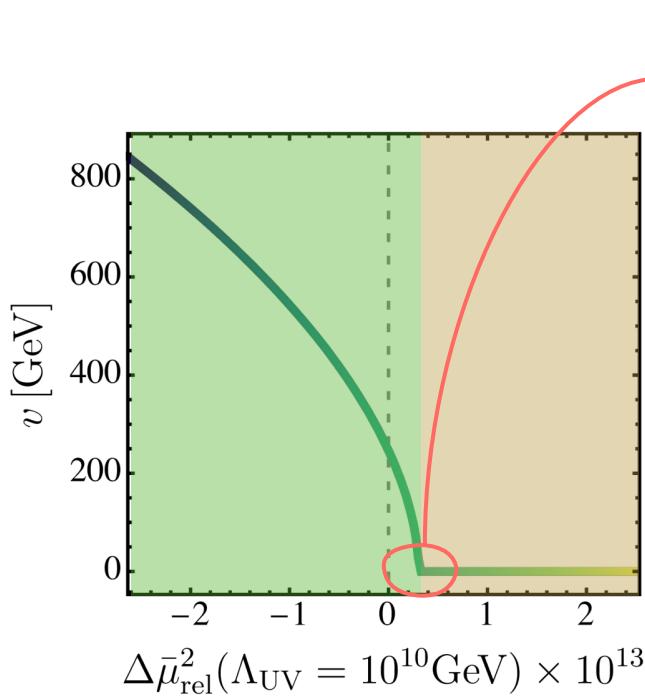
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Quantum criticality

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New-physics deformations



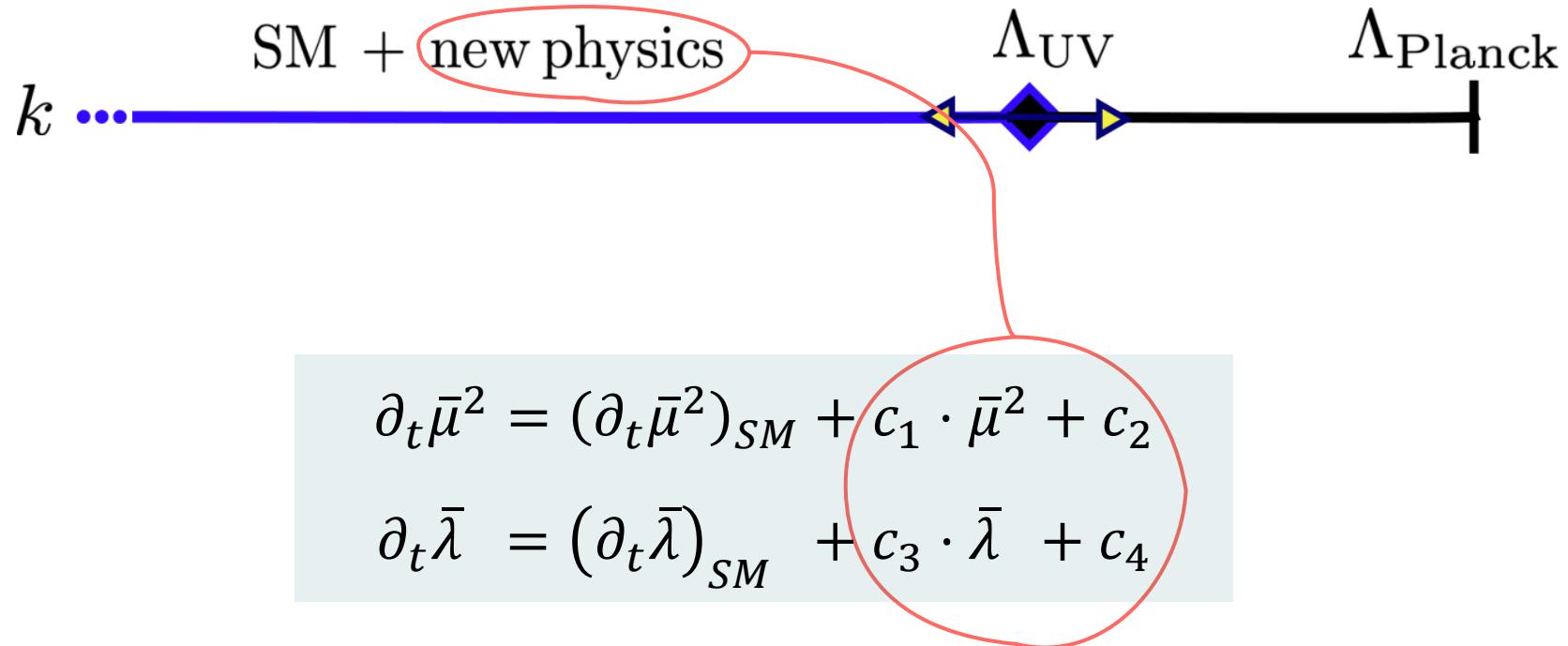
New-physics deformations



$$\partial_t \bar{\mu}^2 = (\partial_t \bar{\mu}^2)_{SM} + c_1 \cdot \bar{\mu}^2 + c_2$$

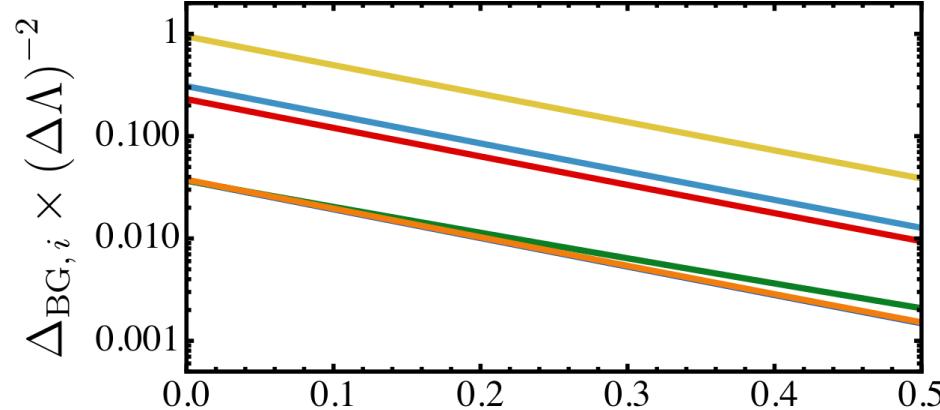
$$\partial_t \bar{\lambda} = (\partial_t \bar{\lambda})_{SM} + c_3 \cdot \bar{\lambda} + c_4$$

New-physics deformations

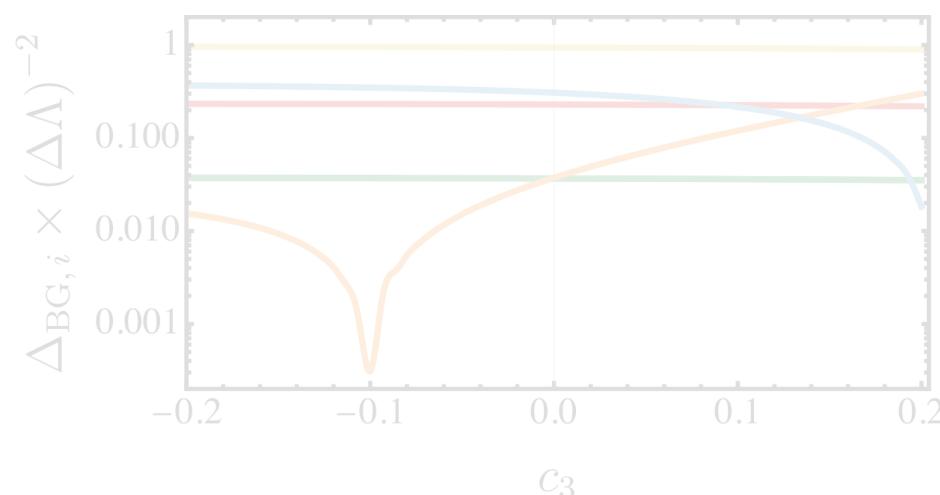


New-physics deformations

— $\bar{\mu}^2$ — $\bar{\lambda}$ — y_t — g_3 — g_2 — g_1

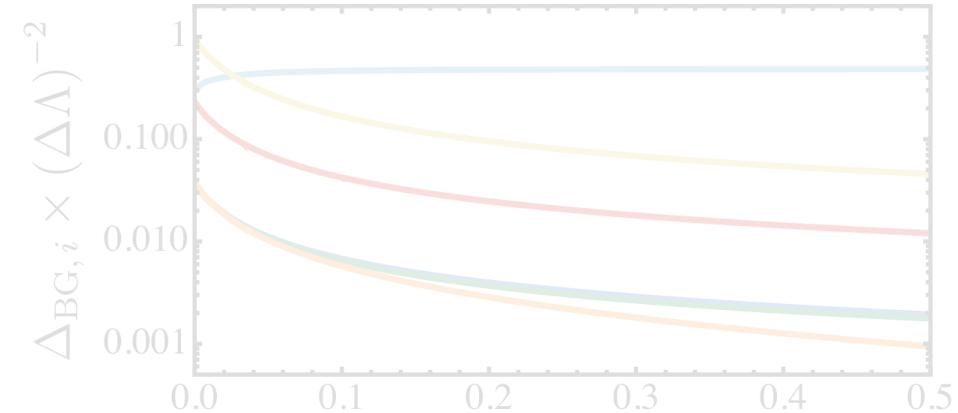


c_1

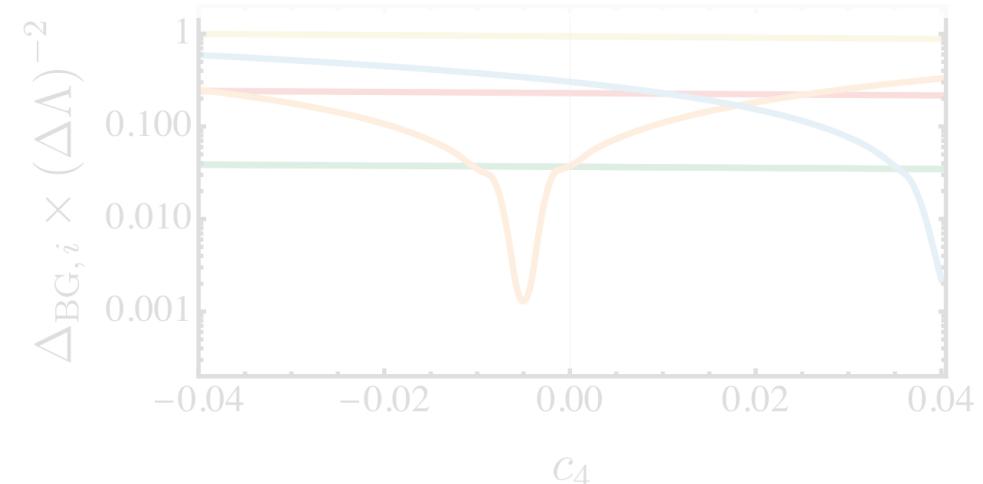


c_3

$$\partial_t \bar{\mu}^2 = (\partial_t \bar{\mu}^2)_{SM} + c_1 \cdot \bar{\mu}^2 + c_2$$
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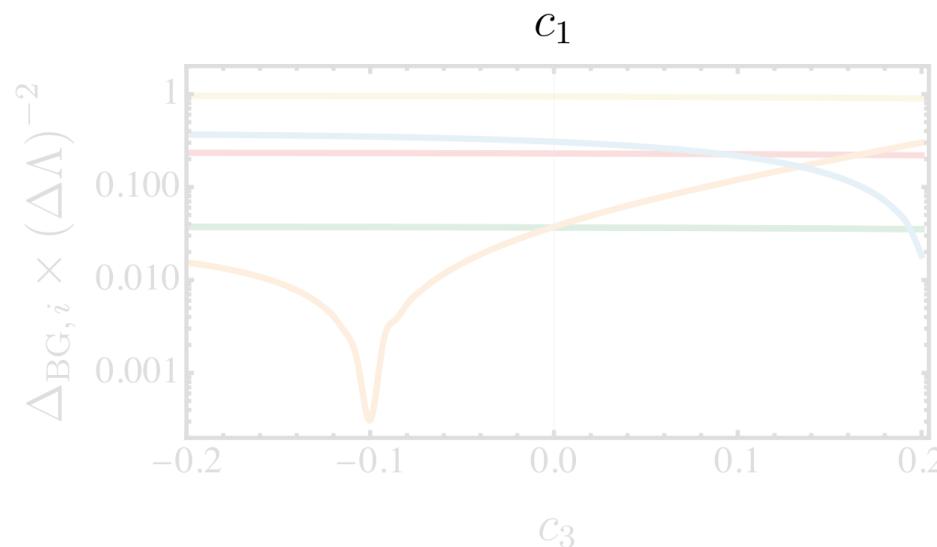
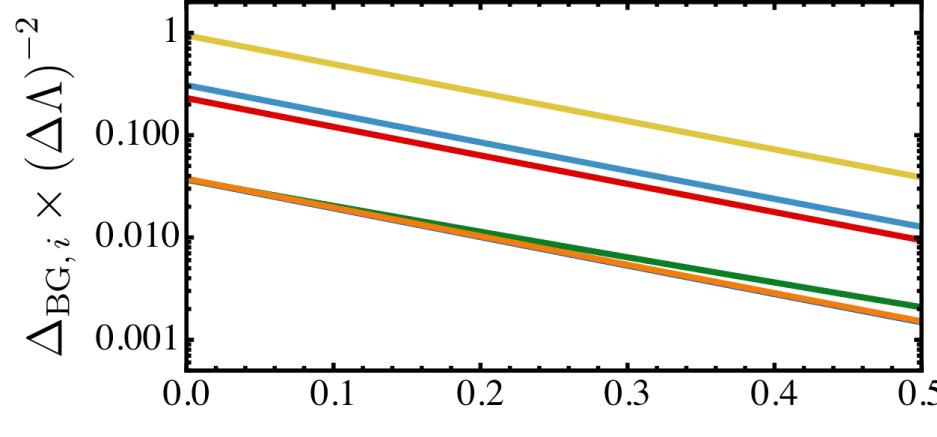
c_2



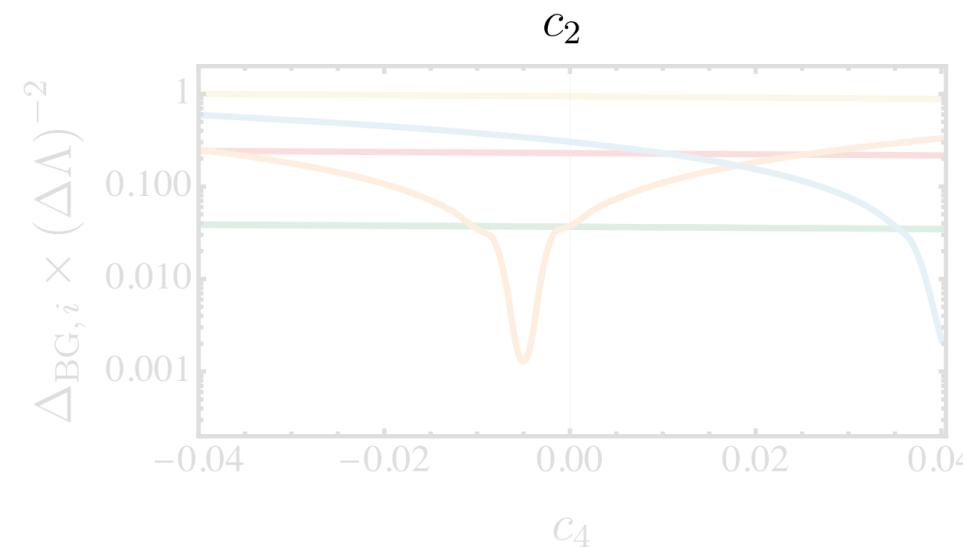
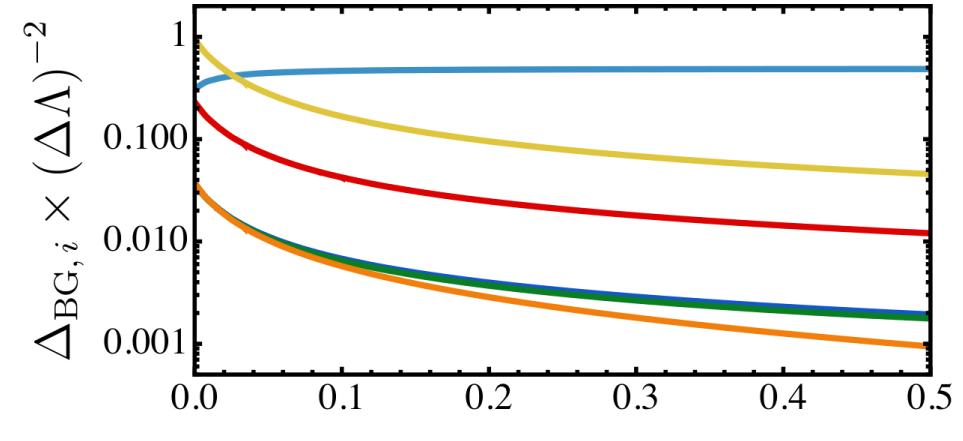
c_4

New-physics deformations

— $\bar{\mu}^2$ — $\bar{\lambda}$ — y_t — g_3 — g_2 — g_1

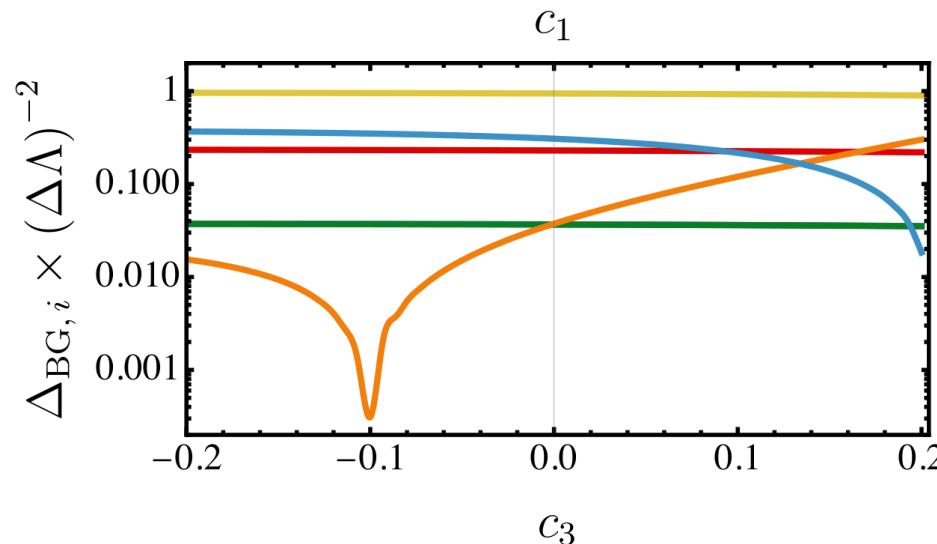
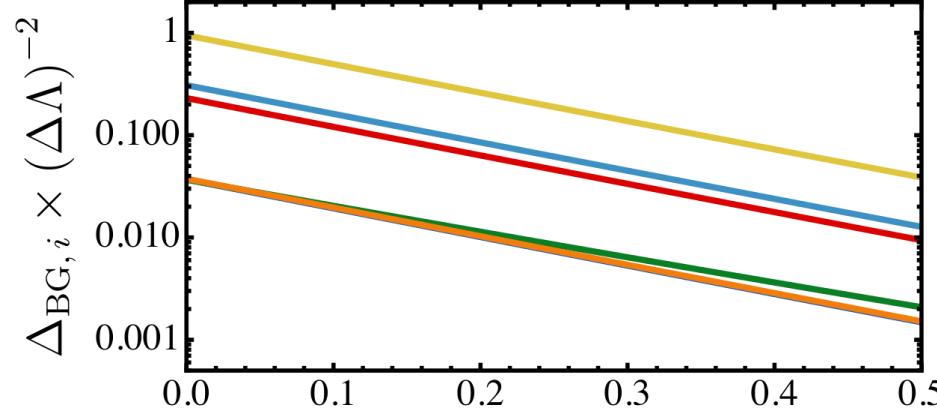


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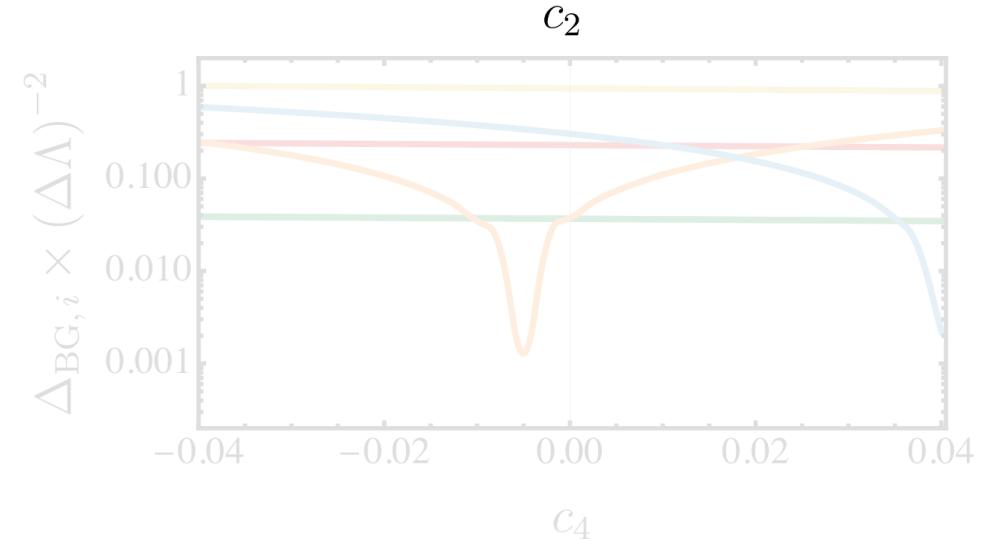
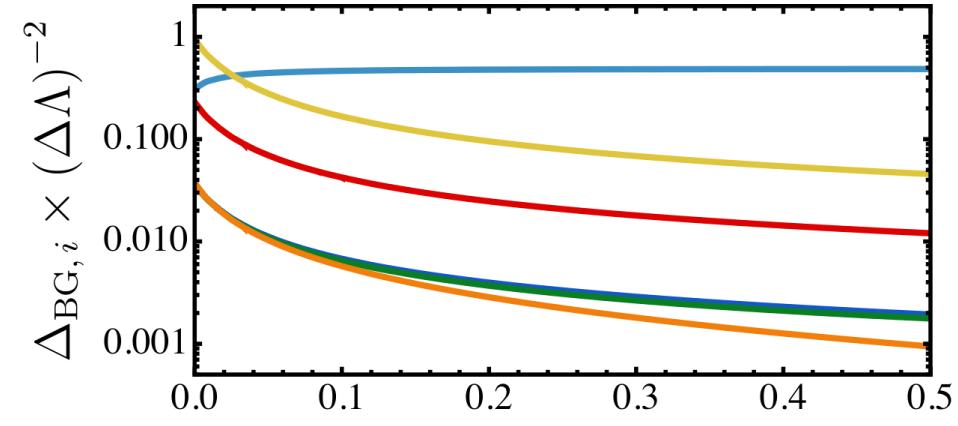


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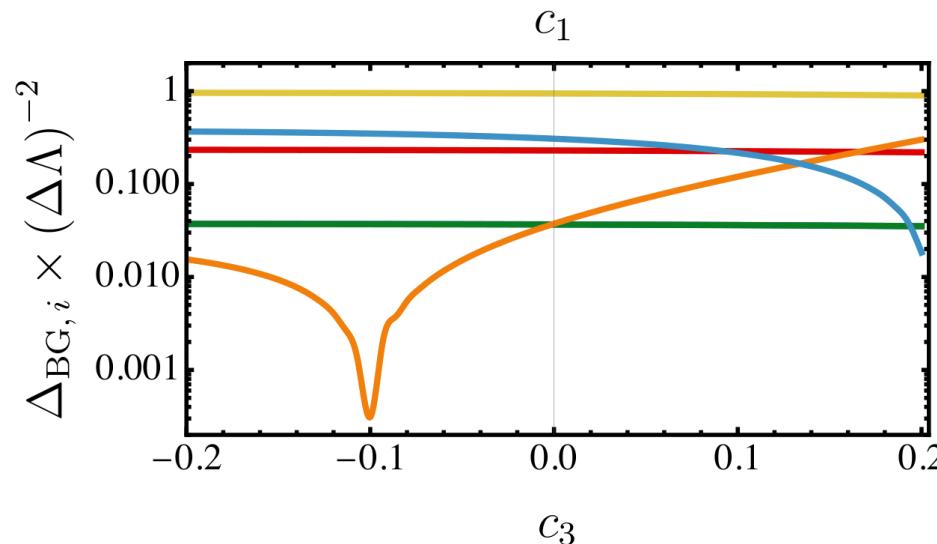
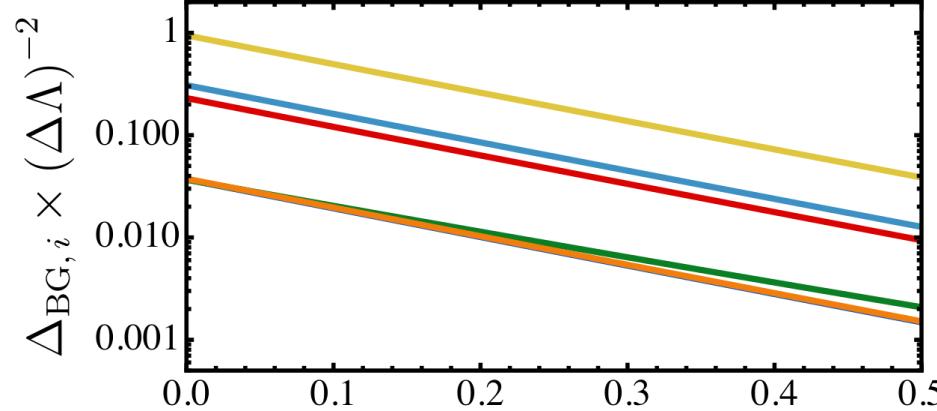


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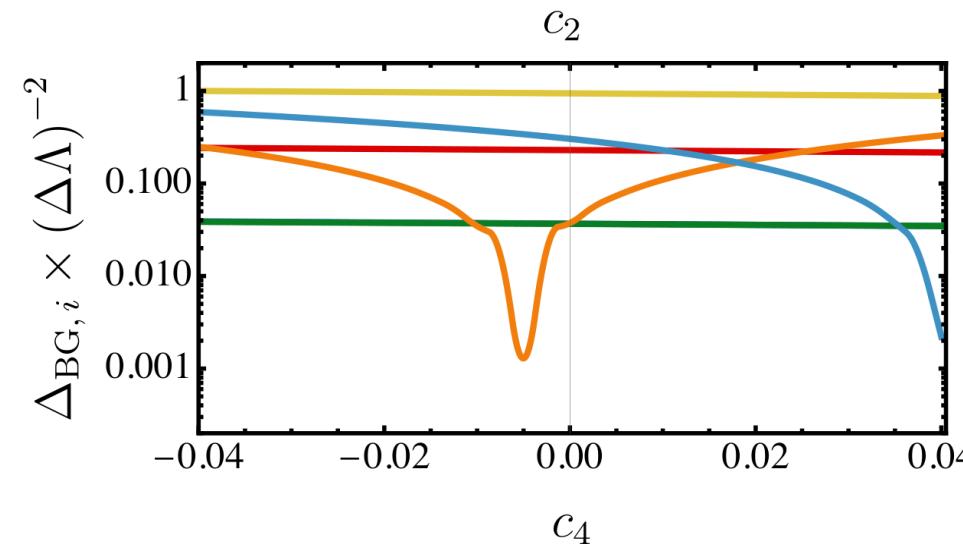
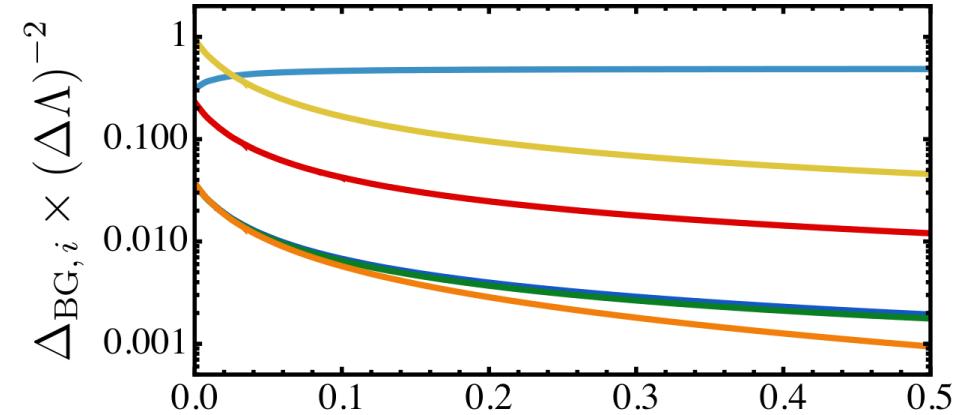


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Summary

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Summary

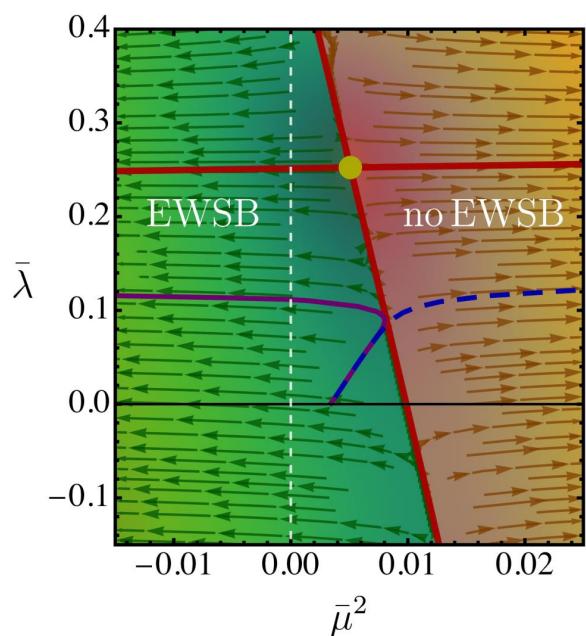
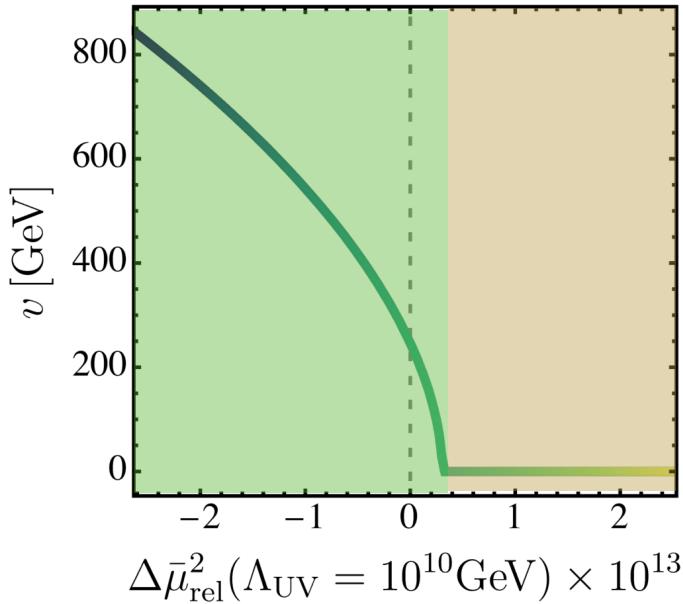
- The mass-like character of the fRG allowed us to study the fine-tuning problem in a model-independent way
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Summary

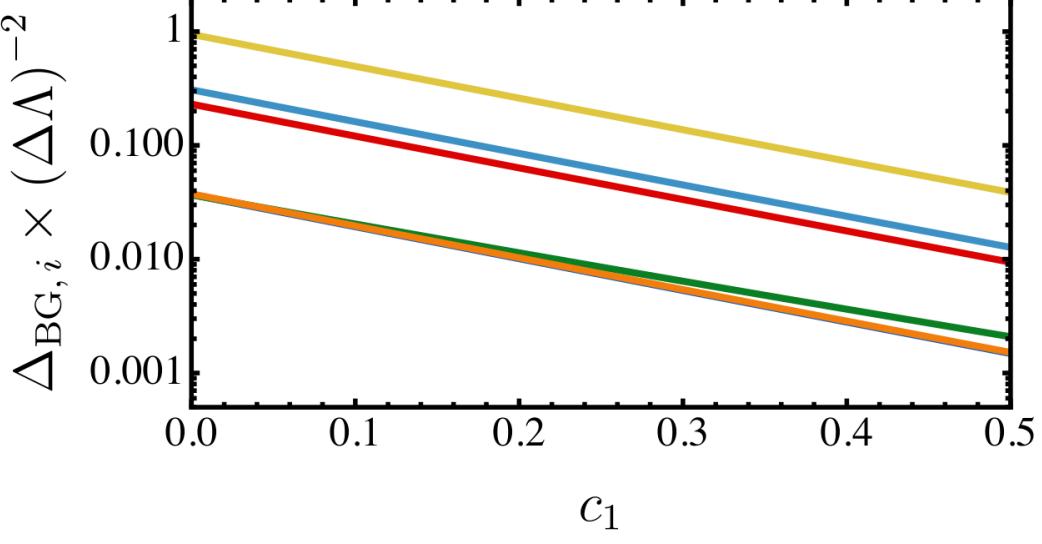
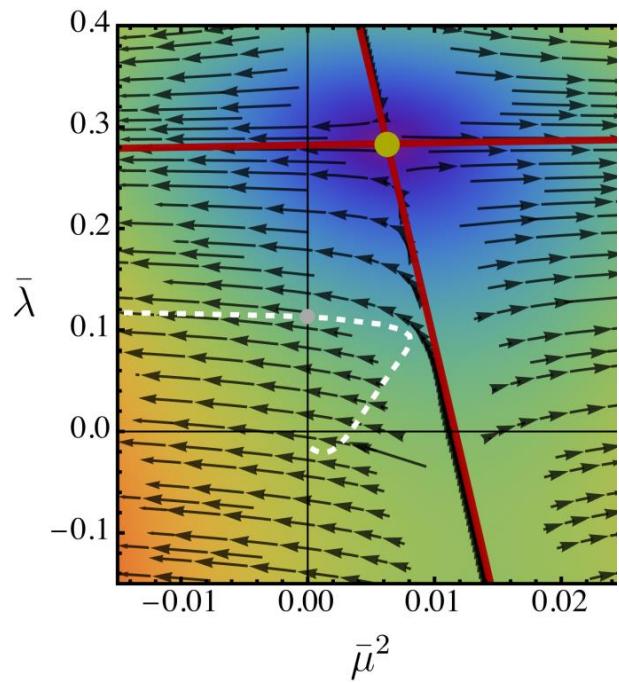
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Summary

- The mass-like character of the fRG allowed us to study the fine-tuning problem in a model-independent way
- Extended the standard fine-tuning discussion by connecting to the notion of quantum criticality
 - SM-like theories with large scale separations are not only fine-tuned but also near quantum critical
- A simplified approach to study the effect of new physics on naturalness reproduces the large anomalous dimension solution



Thank you!



Backup slides

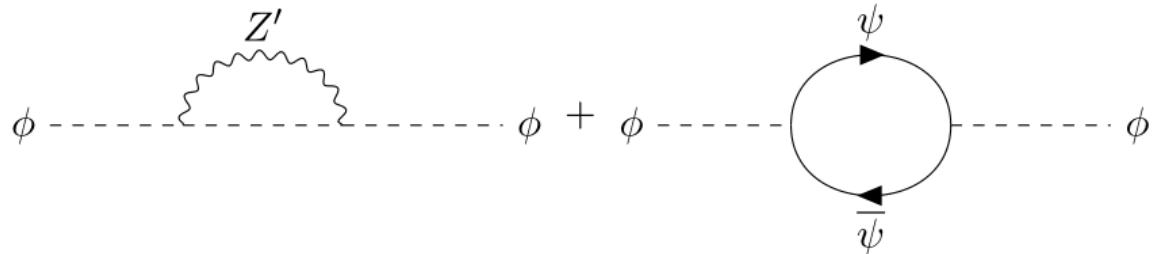
New-physics deformations

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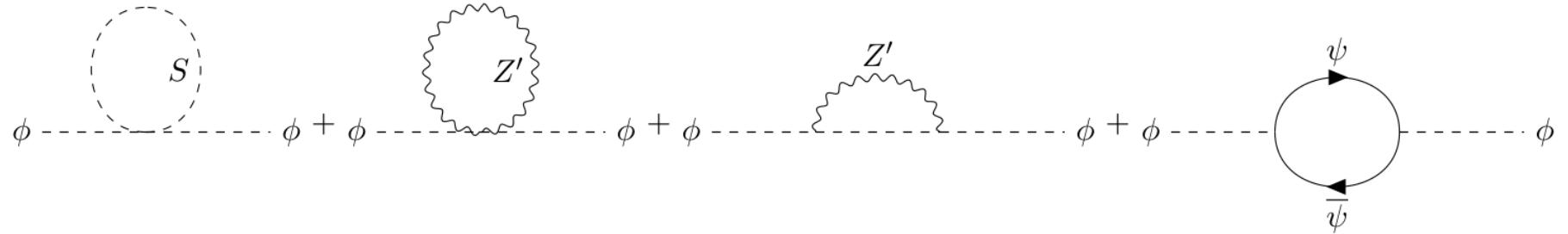
$c_{1,3}$

\supset



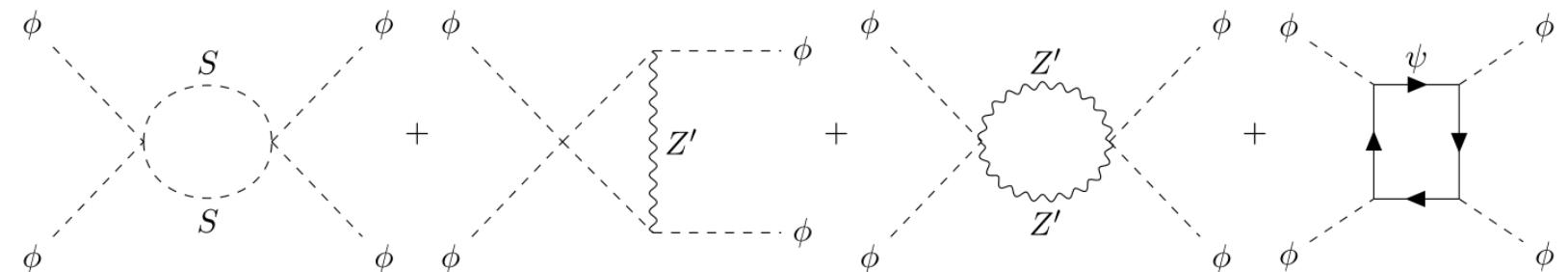
c_2

\supset

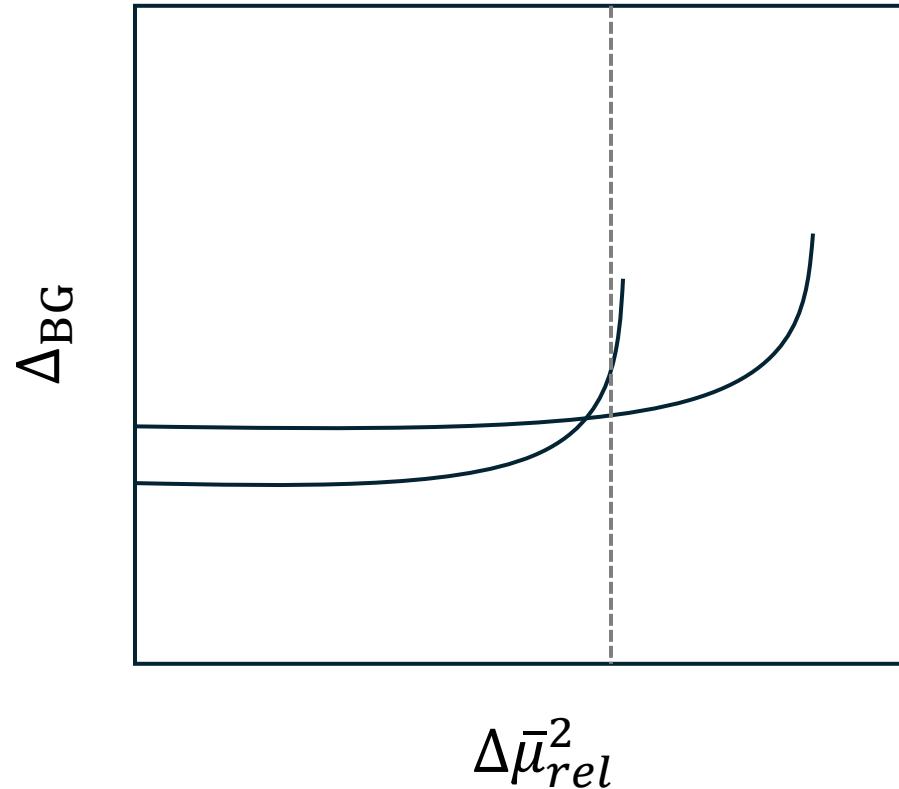


c_4

\supset



Consequences for the (little) hierarchy problem

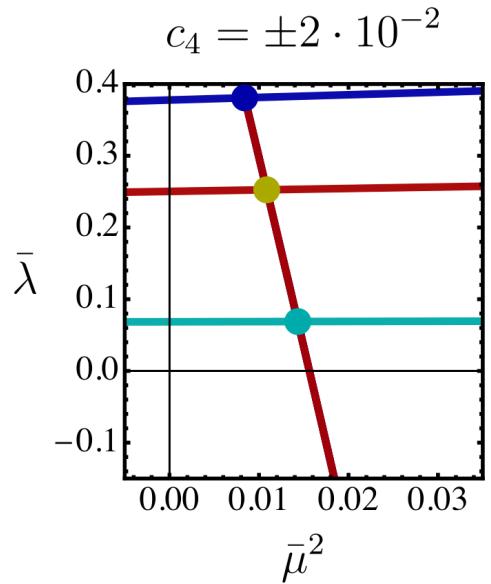
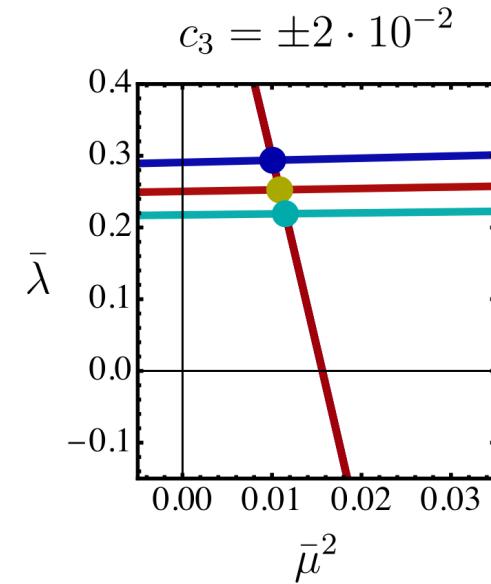
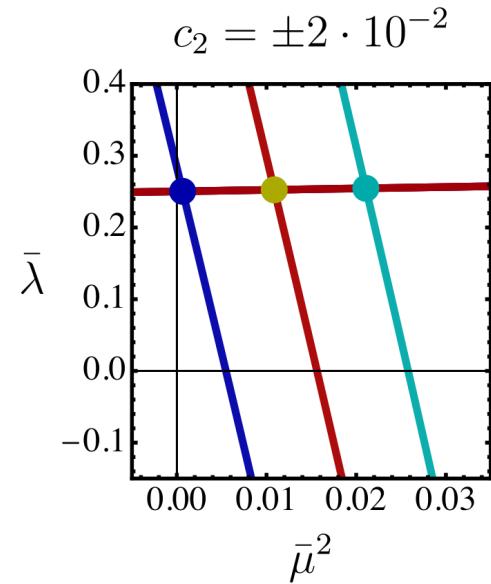
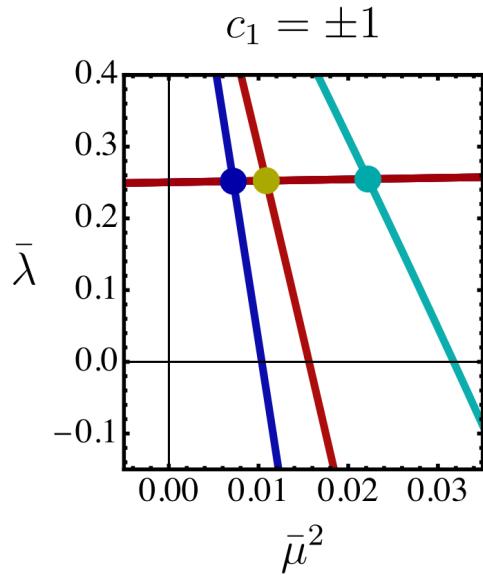


New-physics deformations

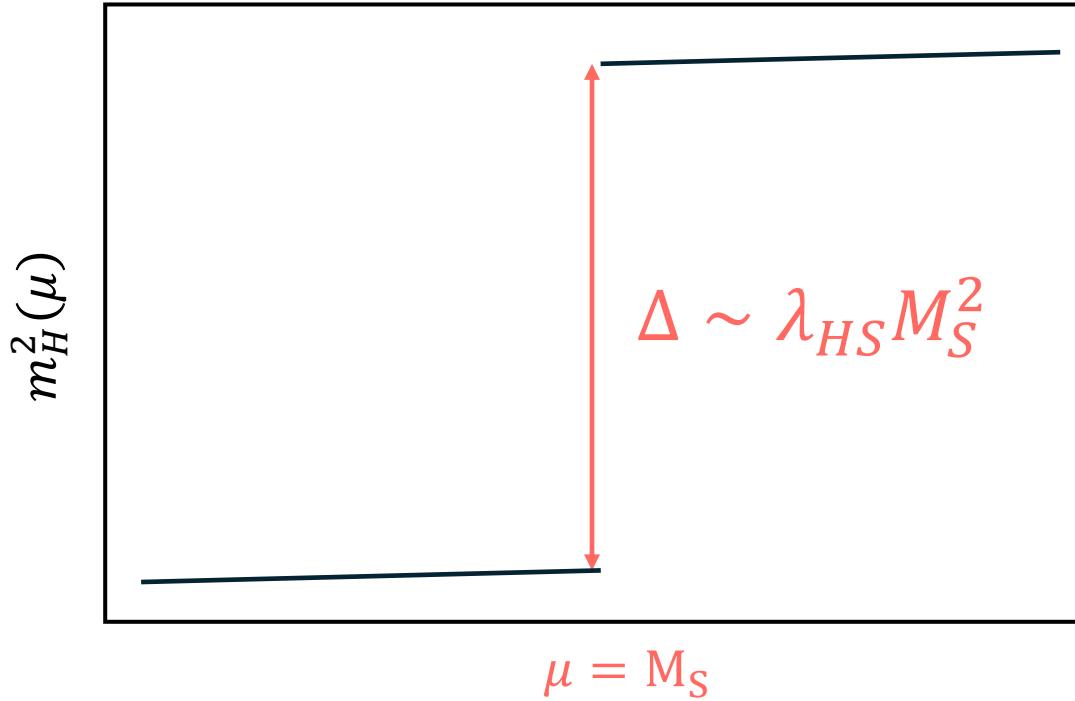


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Standard view of Naturalness



Fine-temperature analogy

$$G_{H,k}(T, p^2) = \frac{1}{Z_{H,k}} \frac{1}{\omega_n^2 + \vec{p}^2 (1 + r_k(\vec{p})) + m_{H,k}^2}$$

$$T_{\text{SSB},1}^{\text{crit}} \sim k_{\text{SSB},1}/(2\pi) \approx 150 \text{ GeV}$$

Previous Wilsonian studies of Naturalness

Aoki, Iso [1201.0857]

- Physical contributions of QD diagrams cannot be projected out

Krajewski, Lalak [1411.6435]

- Fine-tuning in a Higgs-Yukawa model using the fRG

Yamada [2004.00142]

- Gravity effects does not necessarily spoil the SM naturalness

Gies, Schmieden, Zambelli [2306.05943]

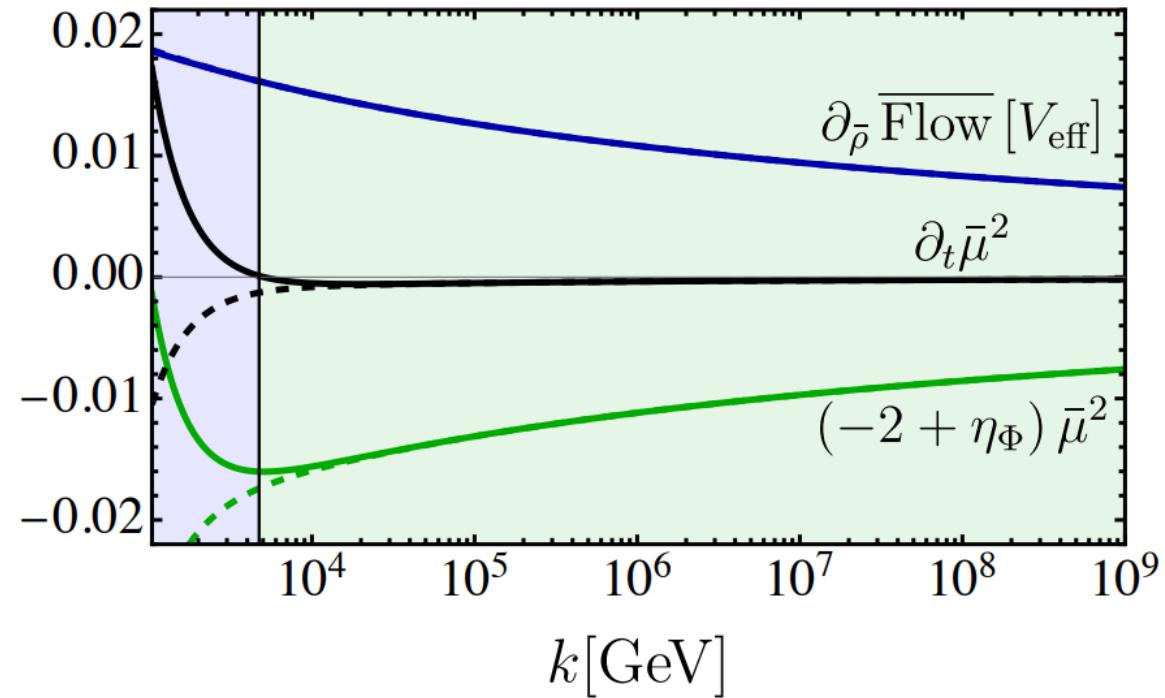
- Moving Λ_{QCD} closer to v_{EW} in the SM reduces fine-tuning

Mass and flow

$$m_H^2 = k^2 [\partial_{\bar{\rho}} u(\bar{\rho}) + 2\bar{\rho} \partial_{\bar{\rho}}^2 u(\bar{\rho})]_{\bar{\rho}=\bar{\rho}_0} = \begin{cases} \mu^2 & \bar{\rho}_0 = 0 \\ 2\bar{\lambda}v^2 & \bar{\rho}_0 > 0 \end{cases}$$

$$\begin{aligned} \partial_{\bar{\rho}} \overline{\text{Flow}} [V_{\text{eff}}] \Big|_{\bar{\rho}_0=0} = & \frac{1}{8\pi^2} \left[- \frac{\bar{\lambda}}{(1+\bar{\mu}^2)^2} \left(3 - \frac{\eta_G^\pm}{6} - \frac{\eta_G^0}{12} - \frac{\eta_H}{4} \right) + 3y_t^2 \left(1 - \frac{\eta_t}{5} \right) \right. \\ & \left. - \frac{18g_1^2}{80} \left(1 - \frac{\eta_{Z^0}}{6} \right) - \frac{18g_2^2}{16} \left(1 - \frac{\eta_{Z^0}}{18} - \frac{\eta_{W^\pm}}{9} \right) \right] \end{aligned}$$

Los vs. polynomial running



Critical exponents

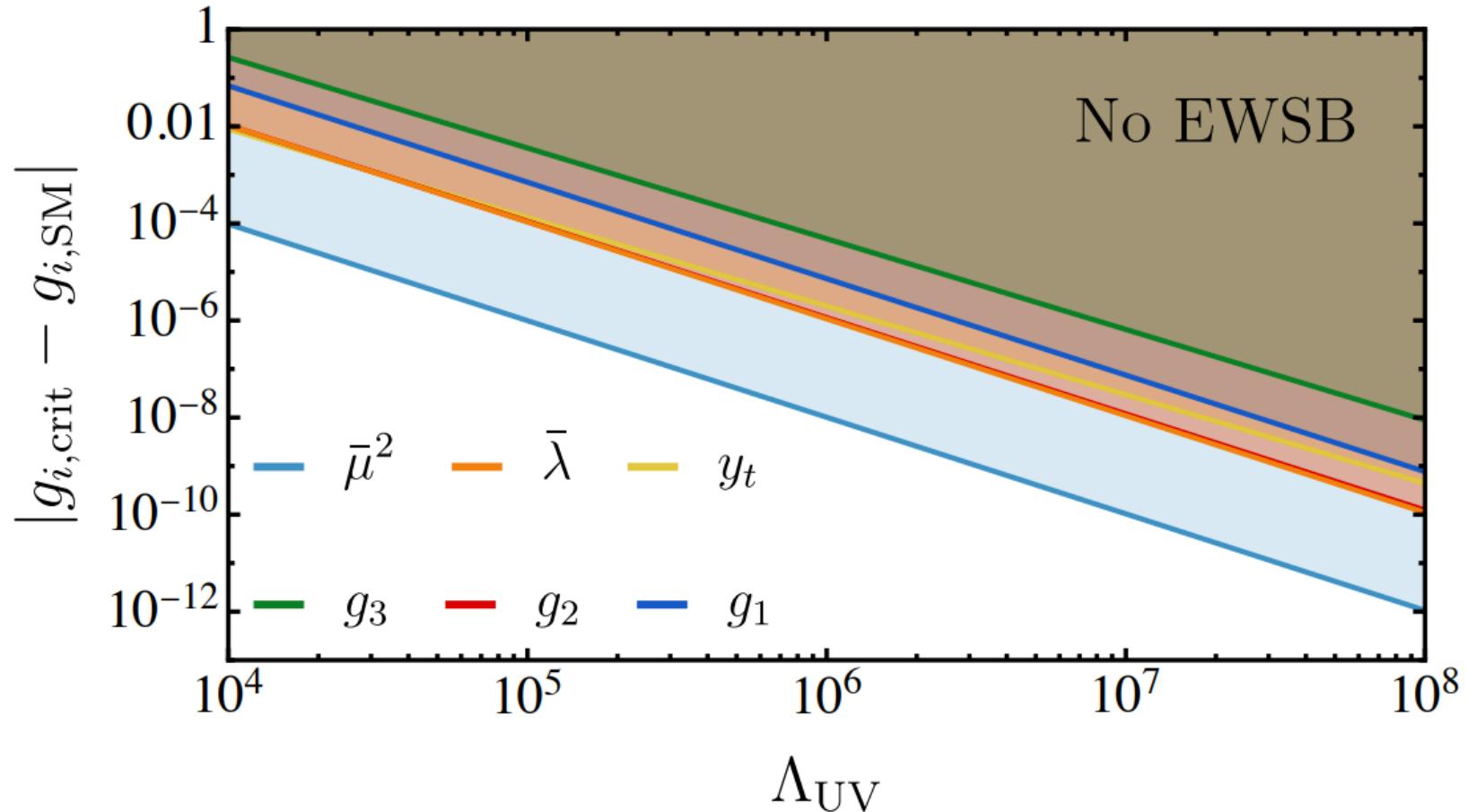
$$v \propto |\Delta \bar{\mu}_{\text{SP}}^2(\Lambda_{\text{UV}}) \times \Lambda_{\text{UV}}^2|^\beta$$

$$\Delta \bar{\mu}_{\text{SP}}^2 = \bar{\mu}_{\Lambda_{\text{UV}}}^2 - \bar{\mu}_{\Lambda_{\text{UV}}, \text{SP}}^2$$

$$\beta = \frac{\nu}{2} \left(d - 2 + \eta_H|_{k_{\text{SSB},1}} \right)$$

$$\nu = -\Theta_{\bar{\mu}^2}^{-1} \approx \left(2 - \eta_H|_{k_{\text{SSB},1}} \right)^{-1}$$

Distance to critical surface



IR observables

