BH mergers and colliders: Two complementary windows to New Physics

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Do we still need New Physics?

While the SM is complete and extremely successful, observational and theoretical needs for NPh remain.

- Neutrino masses
- Dark Matter and Dark Energy
- Baryon asymmetry in the Universe
- Absence of strong CP violation



- Naturalness of the EW scale (hierarchy problem)
- How to incorporate Inflation?
- Quantum gravity

• ...



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I will discuss two complementary types of collisions:

- Black hole mergers: provide signals in multi-messenger astrophysics
- Particle colliders, present and futuresome digging required!





From the Heavens...

...to Earth.

New Physics from the Heavens?



Ingredients: Merging Black Holes, various telescopes, and ... a leap of faith.

Astrophysical Black Holes

Stellar mass BHs produced at the end of heavy star life

SMBH ~ 10^5- billions M_{\odot} and found at the centre of galaxies



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Stellar BH masses seen by GWs from the NS threshold to 200 solar masses.

Heavier masses may exist thanks to successive mergers.

Astrophysical Black Holes

Stellar mass BHs produced at the end of heavy star life

SMBH ~ 10^5 — billions M_{\odot} and found at the centre of galaxies





SMBHs exist at the centres of galaxies. They are responsible for AGNs and for the motion of central stars.

Recently photographed by EHT.

Note that all observations are indirect!

Hawking's Radiation

1974



Credit: Physics Feed

Hawking's Radiation



Hawking's Radiation

$$T_H = \frac{\hbar c^3}{8 \pi G k_B M}$$

Black holes evaporate

$$t_{evaporate} \approx 2.140 \times 10^{67} \text{ years} \left(\frac{M}{M_{\odot}}\right)^3$$

~~~~~~

1974

M

mm

my

M

 $t_{Universe-age} \approx 1.40 \times 10^{10}$  years

The smaller the BH the higher the temperature BHs heavier than 1/2 the moon are colder than the CMB!

No chance to observe HR for any ordinary astrophysical BH

### Can Hawking Radiation be measured?

# A Black Hole Analogy

#### Sonic black hole

#### The setup

A fluid of ultra-cooled atoms flows through a tube. The fluid undergoes quantum fluctuations that produce pairs of phonons, or units of sound, which quickly annihilate.

Pair creation VV and annihilation VV



Sonic Hawking radiation

A laser is used to accelerate the fluid to supersonic speeds partway along the tube. If a pair of phonons straddles the "sonic horizon," one phonon is swept into the supersonic side with no chance of annihilating with its partner, which propagates through the subsonic fluid. Fluid flow  $\rightarrow$ 









0.15

0.10

0.05

-0.05

-0.10

0

LUCY READING-IKKANDA FOR QUANTA MAGAZINE



b

## A Black Hole Analogy



## HR from smaller BHs

Smaller BHs can emit HR that could be potentially observed

Asteroid size BHs could be produced in the early universe (PBHs)

Potential candidates for dark matter, if sufficiently long lived

HR from PBHs constrained by diffuse gamma ray background

Kimura, Takahashi, Koma, 1607, 01964

Albert et al.(HAWC), JCAP 04 (2020), 026

## What if?

BH mergers could leave a trail of small BHs (<u>BH morsels</u>)

While not expected in general relativity, they may be related to the presence of strong non-linearities, or new physics effects...



Okunkova, Phys. Rev. D.96, 104054 (2017)

Colour shades measure non-linearity



### BH morsels

Is the Hawking radiation from BH morsels observable?

Cacciapaglia, Hohenegger, Sannino, 2405.12880



- The HR emission is isotropic, hence signal will not depend on the geometry of the emission.
- The particle flux only depends on the morsel masses (i.e. Hawking temperature).
- The energy of the emitted particles increases with time, hence giving a characteristic smoking-gun signature!
- Coincidence with gravitational wave observation (depending on the morsel masses...)

## Hawking emission

Emission rate for a given BH morsels distribution

$$J_p = \frac{1}{2} \int_{M_{\min}}^{M_{\max}} \frac{dM_{Bm}}{dM_{Bm}} \frac{dn_{Bm}}{dM_{Bm}} \frac{d^2 N_p}{dt \, dE_p} (\mu_t(M_{Bm}, t))$$

 $\mu_t(M_{\rm Bm},t)$  - Morsels mass at time t &  $M_{\rm Bm}$  is the initial mass

$$\frac{1}{2}$$
 - geometric factor

Differential primary flux for a given species "p"

## Secondary radiation

 $J_p$  computed via BlackHawk

Arbey, Aufflinger Eur Phys J. C 79, 693 (2019), 81, 910 (2021).

Emission, decays and hadronisation via Pythia.

Hadronisation tables reliable for primary between 5 GeV and a few TeVs

Assume population of BH morsels with equal masses and non rotating

Angular momentum dissipates faster than the evaporating mass

# Morsels & Mergers Energy Budget

LIGO/VIRGO/KAGRA pre-merger masses between a few and several  $M_{\odot}$  . Abbott et al. (LIGO), PRX 9, 031040 (2019), PRX 11, 021053 (2021),...

Distance 240 Mpc to 3 Gpc

Initial and final masses indirectly measured via GW spectrum

Example: GW170814 (first BH merger observed by all 3 detectors)

$$30.5_{-3.0}^{5.7} + 25.3_{-4.2}^{2.8} = 53.2_{-2.7}^{3.2}$$

GW energy emitted  $2.7^{0.4}_{-0.3}$ 

Several  $M_{\odot}$  can go into BH morsels, but conservatively we assume one.

### Particle emission

Neutral stable particles, reach Earth undeflected by galactic magnetic fields

Consider BH merger at D=300 Mpc (nearest detected BH mergers)

Photon flux on Earth 
$$F_{\gamma} = \frac{1}{4\pi D^2} \int dE_{\gamma} J_{\gamma}$$

Differential flux  $E_{\gamma}^2 dF_{\gamma}/dE_{\gamma}$ 

### Photon flux

Same mass Bm distribution normalised to  $M_{\odot}$ 

Solid lines  $2 \times 10^7 \, \text{kg}$  & 3400 sec evaporation time

Colours: different times from production 500 sec (blue) 3000 sec (purple) 3400 sec (red)



Emission constant up to 500 sec Explosive at end of BH lifetime Red curve exceeds 100 TeV! BlackHawk limit!

100TeV cutoff, photon optical transparency intergalactic medium

## On the mass dependence

Evaporation time  $\propto M_{Bm}^3$ 

Lighter Bm: more energetic GRBs and shorter evaporating times



### Evaporation time



# Multi-messenger approach

Fermi-GBM & Swift-BAT monitor photons within 30 sec from event alerts Coverage between keVs and MeVs (Neutron star merger range) In this range the signal is below exp sensitivity  $10^{-7}$  erg sec<sup>-1</sup> cm<sup>-2</sup>



# Preliminary bounds (HESS)

HESS followed four LIGO/VIRGO BH mergers (O2 and O3 runs)

1 - 10 TeVs, time: 10<sup>4</sup> - 10<sup>5</sup> sec after the BH merger

Energy flux from mergers below  $10^{-12} erg sec^{-1} cm^{-2}$ 



### PBH connection

BH morsels and PBHs both evaporate via TeV photons Ideal for HAWC, HESS, LHAASO (Atmospheric Cherenkov Telescopes) Search for point-like unaccounted GRBs



Strongest bound from HAWC

Albert et al.(HAWC), JCAP 04 (2020), 026

# Preliminary bounds (HAWC)

A morsel distribution at D is equivalent to a single PBH at  $D_{
m PHB}$ 

$$D_{\rm PBH} = \frac{D}{\sqrt{n_{\rm Bm}}}$$

Naive:

Rescaling LIGO/VIRGO BH mergers rate

$$\rho_{\rm LV} = 24^{+14}_{-9} \,\,{\rm Gpc}^{-3}\,\,{\rm yr}^{-1}$$

PBH rescaled densities for HAWC

$$ho_{\rm PBH} = 3400~{
m pc}^{-3}~{
m yr}^{-1}$$





### Next steps

- Golden opportunity for Cherenkov telescopes (work in progress with Fermi-LAT colleagues).
- New physics effects may appear in the spectrum at the end of the evaporation time (modelling of NPh effects in progress).
- Modelling of the BH morsel production necessary: we are exploring various ideas at the moment.

## New Physics from earth?



Ingredients: Particle colliders (the LHC and a future programme), a good theory of New Physics

#### Motivation for Higgs compositeness

- · Composite models 'solve' the Hierarchy problem...
- with new scale in the multi-TeV!





multi-TeV mountain

#### What are we looking for?

- -> Precision EW + Higgs observables
- -> light composite scalars
- -> multi-TeV resonances (top partners, pNGBs, spin-1)

### Composite spectra



# Low-hanging fruits: scalars!



How can light states emerge?



#### EW scalars: SU(5)/SO(5) benchmark





Dominantly EW pair-prod. Good targets for ee colliders?

#### Best exclusion from multi-photon searches



#### Typical ALP Lagrangian:

$$\mathcal{L}_{\text{eff}}^{D\leq 5} = \frac{1}{2} \left( \partial_{\mu} a \right) \left( \partial^{\mu} a \right) - \frac{m_{a,0}^2}{2} a^2 + \frac{\partial^{\mu} a}{\Lambda} \sum_{F} \bar{\psi}_F \, \mathcal{C}_F \, \gamma_{\mu} \, \psi_F + g_s^2 \, C_{GG} \, \frac{a}{\Lambda} \, G_{\mu\nu}^A \, \tilde{G}^{\mu\nu,A} + g^2 \, C_{WW} \, \frac{a}{\Lambda} \, W_{\mu\nu}^A \, \tilde{W}^{\mu\nu,A} + g'^2 \, C_{BB} \, \frac{a}{\Lambda} \, B_{\mu\nu} \, \tilde{B}^{\mu\nu} \,,$$

Composite Higgs scenario:

$$\frac{C_{WW}}{\Lambda} \sim \frac{C_{BB}}{\Lambda} \sim \frac{N_{\rm TC}}{64\sqrt{2} \pi^2 f}$$
$$(C_{\gamma\gamma} = C_{WW} + C_{BB})$$

 $\frac{C_{GG}}{\Lambda} = 0$ 

(Poor bounds at the LHC)

#### C<sub>F</sub> is loop-induced:

M.Bauer et al, 1708.00443



#### Tera-Z portal to compositeness (via ALPS) G.Cacciapagli

G.Cacciapaglia et al. 2104.11064

#### Photo-phobic

#### Photo-philic



No leading order coupling to Photons (WZW interaction is Zero!!)

#### eg. SU(4)/SP(4), $SU(4)\times SU(4)/SU(4)$

WZW interaction to photons (like the pion) eg. SU(5)/SO(5), SU(6)/SO(6)



#### ALPs at FCC-ee

Production via Z decays:
  $\Gamma(Z \to \gamma a) \propto C_{Z\gamma}^2$ 

o Lifetime  $\Gamma_a^{tot} \propto C_{\gamma\gamma}^2$ 





4 experimental regions depending on decay length L of ALP

- •100 events for L<10 mm (prompt)</pre>
- •4 events for 10<L<2000 mm (Long lived) Decay in ID
- •4 events for 2000<L<4500 mm (Calo) Decay in calorimeter

•100 events for L>4500 mm: ALP decays outside the detector, only accompanying photon detected (monophoton)

#### Combined plot FCC-ee



10/10/2024

#### Photophobic case



Hadronic decays interesting for FCC-ee/CepC

Tau and muon could be accessible at the LHC (LHCb)

#### Estimated reach of LHCb from 2106.12615





#### Roadmap to Higgs compositeness

- The HL-LHC will leave an important legacy, but NOT covering the whole interesting parameter space! (i.e. 10TeV is the target)
- A Tera-Z run will fully test the presence of a light composite ALP -> well beyond the 10 TeV mark
- Case 1 : discovery +
   EWPTs can fix the scale

Case 2 : non-discovery
+ EWPTs

 In both cases, the results will strongly constraint the model building, providing testable predictions for a high energy pp collider.



# Morsels emission dynamics

Each morsel emits particles with rates depending on its mass

Emission rate per species "p" follows

$$\frac{d^2 N_p}{dt \, dE_p} = \frac{1}{h} \frac{\Gamma_p(E_p, M_{\rm Bm})}{\exp \frac{E_p}{k_B T_{\rm Bm}} \pm 1}$$

 $T_{Bm}$  - BH morsel temperature

Denominator: Boltzmann statistic factor

 $\pm$  1 for fermions/bosons

 $\Gamma_p(E_p, M_{
m Bm})$  - grey body factor

### Neutrinos

Neutrino observatories ANTARES and IceCube monitor BH mergers

ANTARES and IceCube ~ 500 sec window

No excess found

However, flux limits are orders of magnitude above photon ones

Require a luminosity of  $10^{51} \, erg \, sec^{-1}$ 

BH morsels predict neutrino fluxes similar to photons.

Experiments not competitive

### Super Massive BH Mergers

