

# Bootstrapping Quantum Field Theories

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LPENS (Axe Interactions fondamentales)



Introduction to the Bootstrap

Extending the Bootstrap: Some Contributions

Uncharted 2d CFTs

Conclusions and ongoing work

# Introduction to the Bootstrap

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## General Motivation

- The physics of many interacting degrees of freedom (QFT) is hard.
- The renormalization group organizes the structure: CFTs are endpoints in the space of QFTs

$$S = S_{\text{CFT}} + \lambda \int d^d x O(x) \quad (1)$$

- Microscopic formulations, in terms of a few fields  $\phi$  sometimes exist

$$\langle O(\phi) \rangle = \int D\phi O(\phi) e^{-S(\phi)} \quad (2)$$

- But are usually only useful in perturbation theory. The IR degrees of freedom are often non-trivially related to the UV fields.

# Why the Bootstrap?

- There are microscopically-minded non-perturbative approaches: Lattice Monte Carlo, Hamiltonian Truncation, Tensor Networks, etc.
- We follow the opposite philosophical viewpoint:  
**Bootstrap** principles/"Nuclear Democracy" [Chew and Frautschi, 1961].  
Focus on observables. All excitations are treated on the same footing.  
The system should interact self-consistently.
- The more constraints/symmetry, the better. Success in 2d CFTs [Belavin et al., 1984], and  $d > 2$  CFTs [Rattazzi et al., 2008]

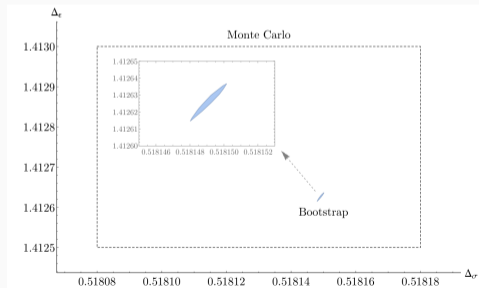
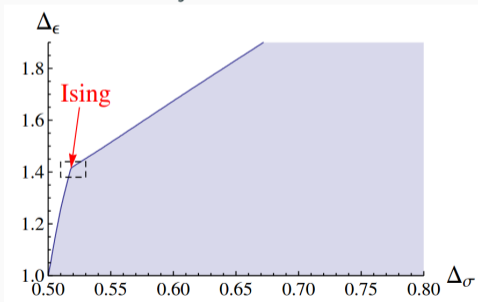
$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle = \langle \phi(x_1)\phi(x_4)\phi(x_3)\phi(x_2) \rangle \quad (3)$$

# Paradigmatic example: Four-point functions in CFT

- Crossing ( $x_2 \leftrightarrow x_4$ ), along with the convergent OPE and unitarity gives

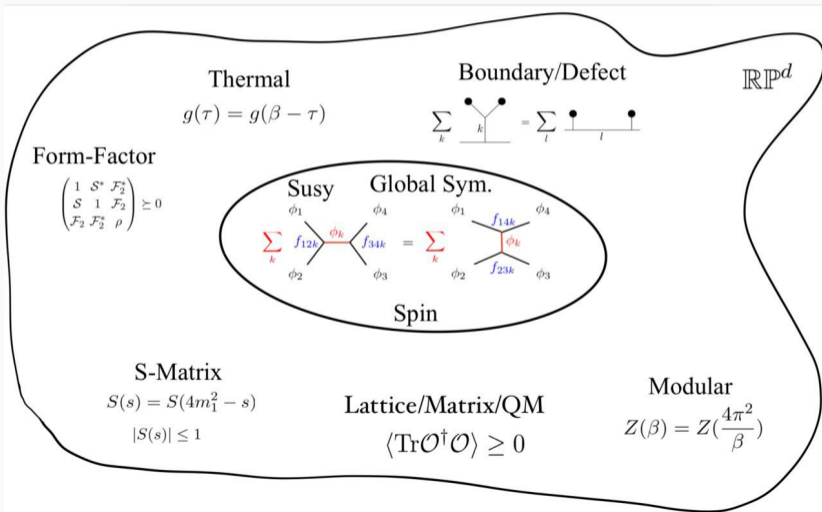
$$\sum_O c_{\phi\phi O}^2 \left( v^{\Delta_\phi} G_{\Delta_O, J_O}(u, v) - u^{\Delta_\phi} G_{\Delta_O, J_O}(v, u) \right) \equiv \sum_O c_{\phi\phi O}^2 F_O(u, v) = 0 \quad (4)$$

- Dual strategy: Narrow down the space of CFTs by excluding inconsistent CFT data  $\Delta_i, c_{ijk}$ . Find linear functional  $\alpha$  such that  $\alpha(F_O) > 0$  for  $O \in \text{TrialSpec}$ .



# Map of the Bootstrap

Other observables/conditions lead to the space of Bootstrappable physics:



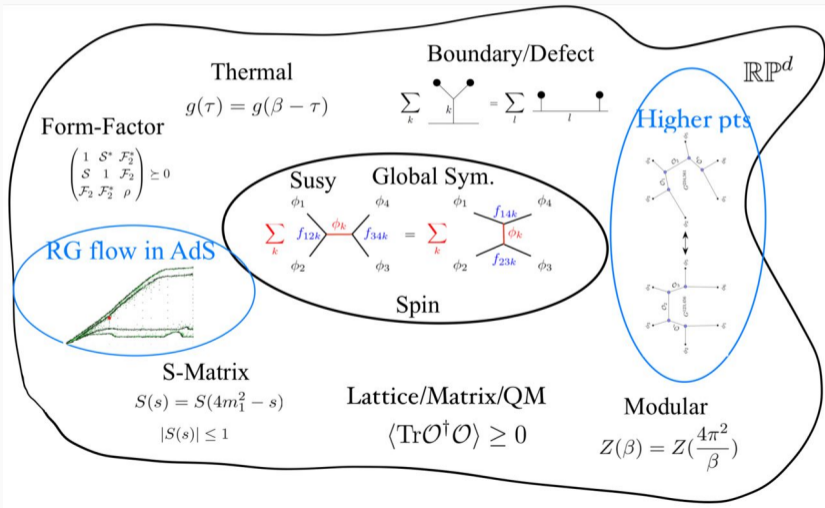
## **Extending the Bootstrap: Some Contributions**

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# Extended Map of the Bootstrap

Primal effort: extend the space of Bootstrap applications

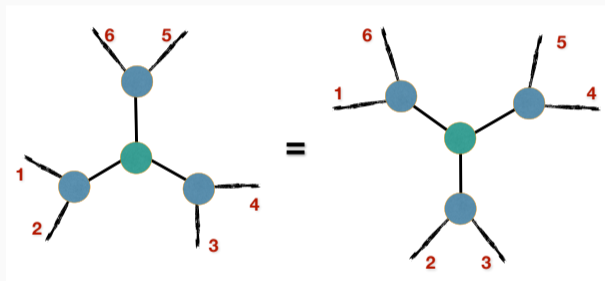


# Higher-point Bootstrap

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## Lightcone Bootstrap at higher points

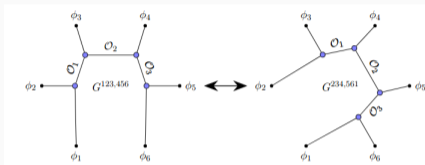
- Observables:  $\langle O(x_1)O(x_2)O(x_3)O(x_4)O(x_5)O(x_6) \rangle$  in Euclidean  $CFT_d$
- Consistency condition: Crossing ( $x_i \rightarrow x_{i+1}$ ) in snowflake OPE



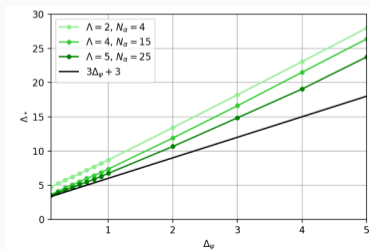
- Relate dominant terms in direct channel to large-spin data  $C_{J_1 J_2 O}$ ,  $C_{J_1 J_2 J_3}$  in the crossed channel, obtaining large spin behaviour.

# Lining up a positive semi-definite six-point bootstrap

- Positive  $\langle O(x_1)O(x_2)O(x_3)|O(-x_3)O(-x_2)O(-x_1)\rangle$  in  $\text{CFT}_1$  (comb-channel)



- Can bound dimensions  $\Delta_2$  in  $O \times O \times O$  OPE as well as  $\langle OOOO_2\rangle$ .

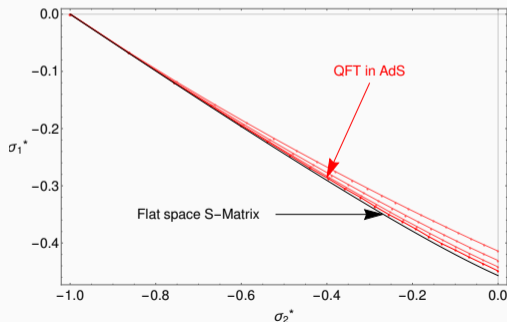


## RG flows in AdS

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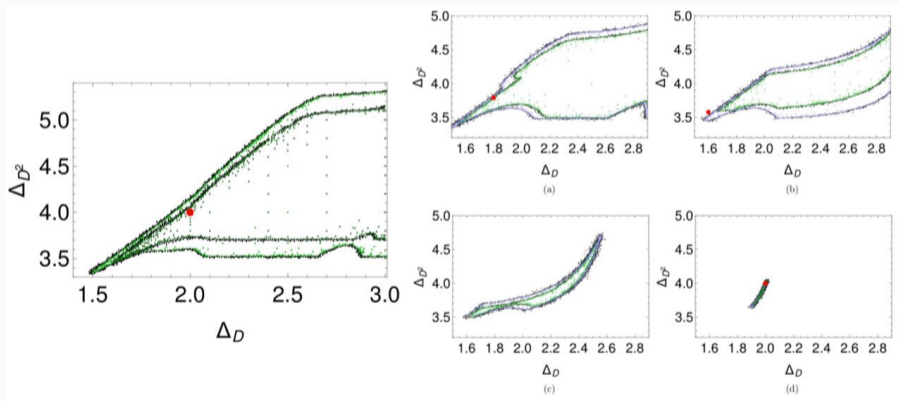
## Towards Bootstrapping RG flows: sine-Gordon in AdS

- Setup: QFT in fixed  $\text{AdS}_2$  background, where  $L_{\text{AdS}}$  sets the scale.
- Observables: "masses"  $\Delta_i/\Delta_j(L_{\text{AdS}})$ , "couplings"  $c_{ijk}$ , S-matrices  $\langle OOOO \rangle$ .
- Consistency condition: CFT axioms of boundary four-point functions.
- Non-perturbative bounds from UV to IR.



# Towards Bootstrapping RG flows: Minimal models in AdS

- Can also study setup where IR is gapless. Example: RG flow between Minimal Models.



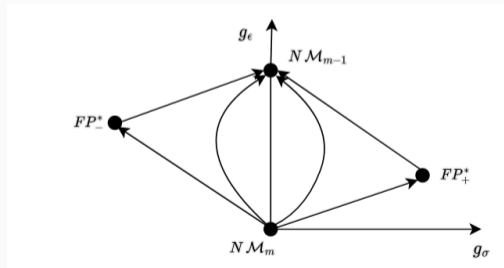
## Uncharted 2d CFTs

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## RG and irrational 2d CFTs

- The Bootstrap is an amazing tool to chart out the space of QFTs but older tools, like perturbative RG are also useful!
- A particularly unexplored region is the space of irrational 2d CFTs, with only conformal symmetry; they have  $\infty$ -many primaries and are presumably chaotic.
- We proposed a perturbative construction of infinite families of such CFTs.



## Coupled minimal models

- We find the theories as IR fixed points of  $N$  coupled rank  $m$  minimal models

$$S_{\text{CMM}} = \sum_{i=1}^N S_m^i + g_\epsilon \int d^2x N^{-\frac{1}{2}} \sum_{i=1}^N \phi_{(1,3)}^i + g_\sigma \int d^2x \binom{N}{4}^{-\frac{1}{2}} \sum_{i < j < k < l}^N \phi_{(1,2)}^i \phi_{(1,2)}^j \phi_{(1,2)}^k \phi_{(1,2)}^l. \quad (5)$$

- We argue that there is no enhanced symmetry by checking that all additional singlet currents up to spin 10 acquire anomalous dimensions ( $\bar{\partial} T_\ell = K_\ell$ )

$\ell \backslash N$	4	5	6	7
4	(1, 1)	(1, 2)	(1, 2)	(1, 2)
6	(2, 2)	(2, 5)	(2, 6)	(2, 6)
8	(4, 7)	(4, 17)	(4, 22)	(4, 23)
10	(5, 18)	(7, 50)	(7, 69)	(7, 75)

## **Conclusions and ongoing work**

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- Bootstrap methods have broad applications in physics.
- We developed extensions for higher-point correlators and RG flows in AdS.
- We proposed new CFTs that deserve to be bootstrapped.

## What I am working on right now

- (No) Integrable RG flows in AdS
- Improved six-point numerics, five-point Polyakov Bootstrap
- Bootstrapping  $\lambda\phi^4$  in AdS

Thank you!

Thank you for your attention!





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