



Heavy flavor correlations and Quarkonia production in high energy proton-proton collisions in the EPOS4 framework

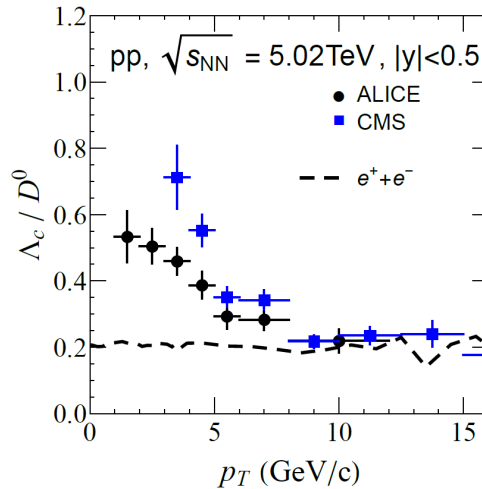
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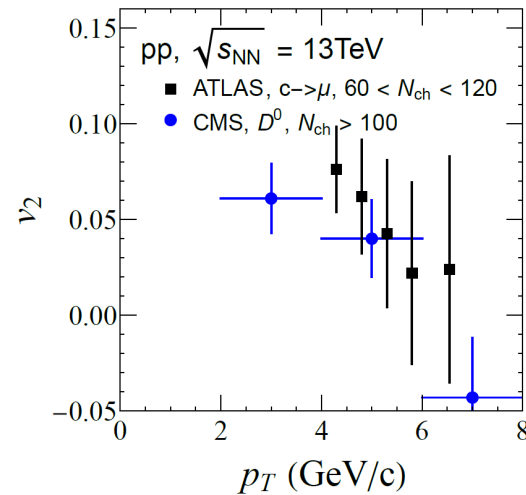
The study of heavy mesons in pp collisions offers big surprises

Λ_c / D^0 ratio



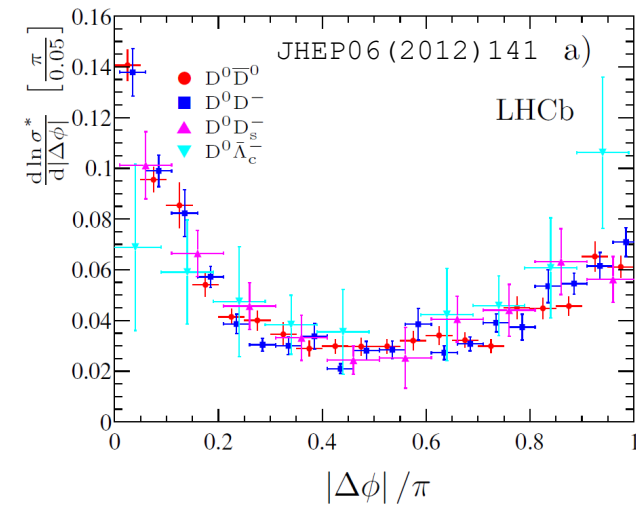
are fragmentation
functions not universal?

elliptic flow v_2



pQCD: $v_2 = 0$
Where does the finite
 v_2 come from?

azimuthal $\Delta\phi(p_D, p_{Dbar})$

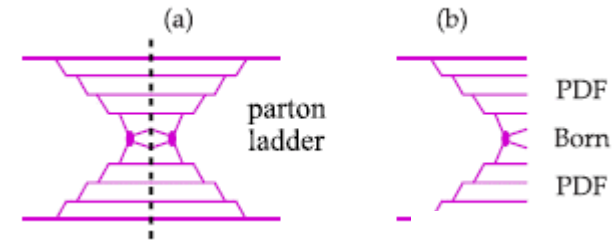
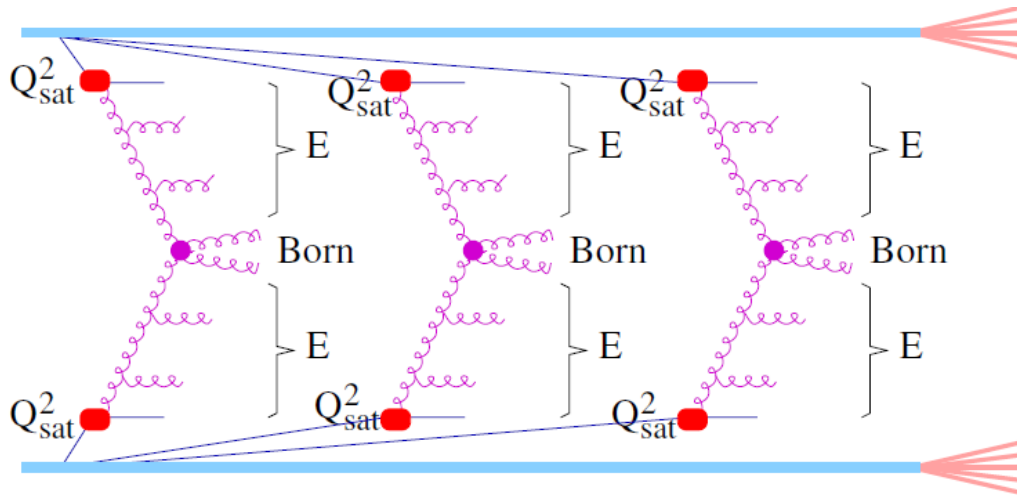


What causes this
structured correlation function?

EPOS4 (→ see literature below)

EPOS4: general purpose event generator for heavy ion collisions at RHIC and LHC

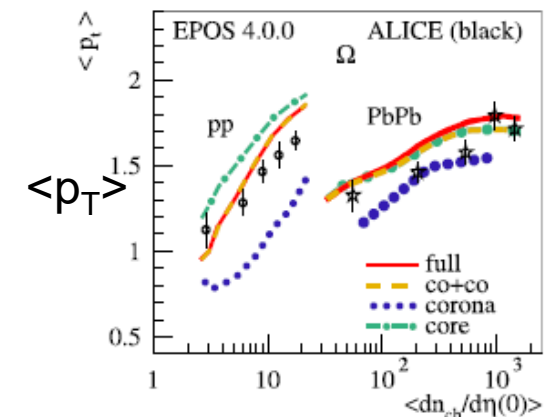
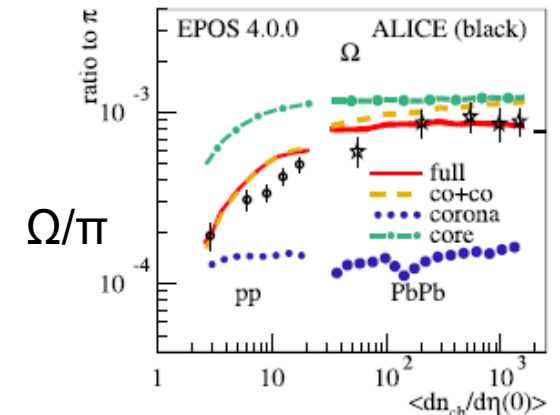
All scattering are rigorously treated in **parallel**
Overall **energy conservation** and factorisation
Binary scaling for hard processes
Saturation



EPOS4 describes well a multitude of observables:

Here those for Ω^-

quite difficult to reproduce



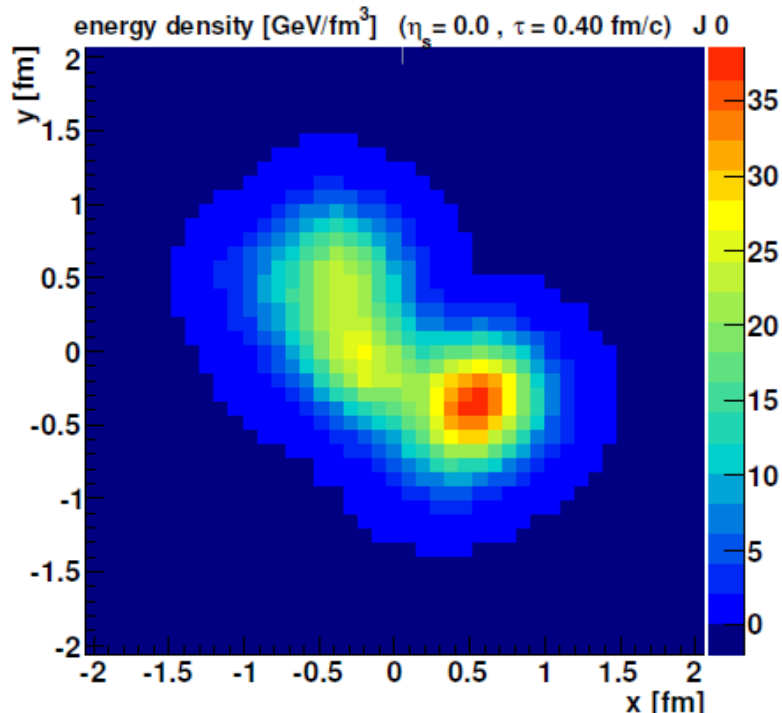
pp in EPOS4HQ

EPOS4(HQ) treats pp collisions like heavy-ion collisions: *

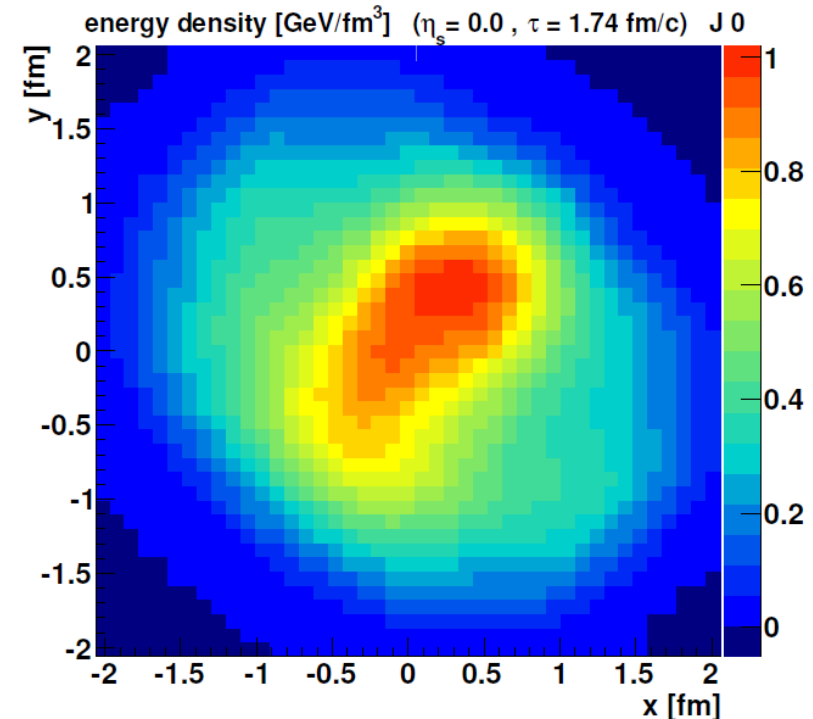
a Quark-Gluon-Plasma is created if energy density $> 0.57 \text{ GeV/fm}^2$

Energy density in the transverse plane of a typical pp event
(each event looks differently)

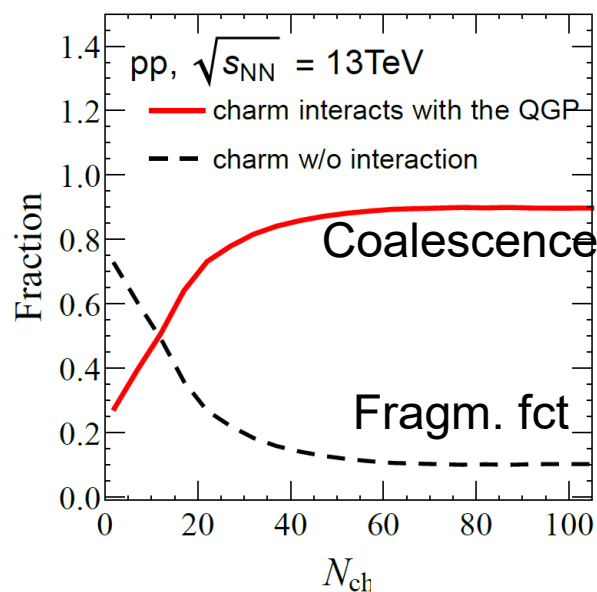
Initial



close to hadronization

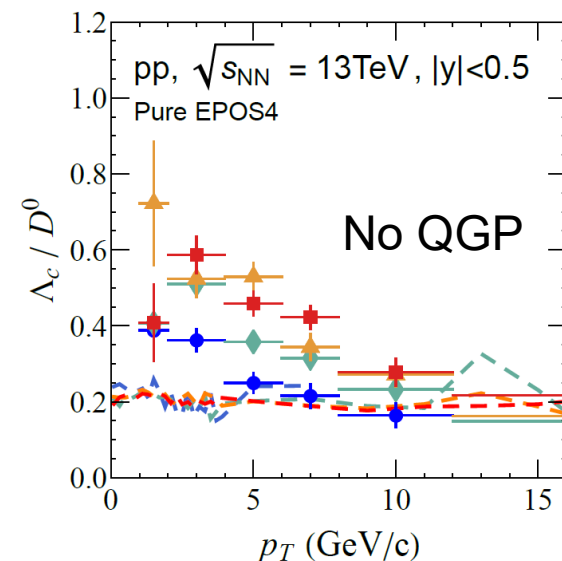
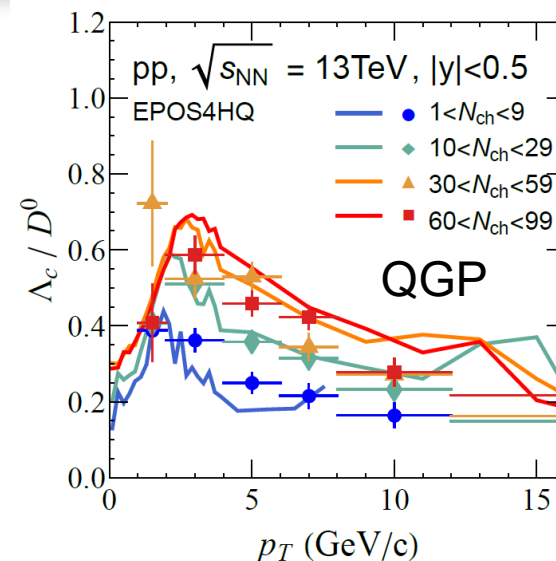
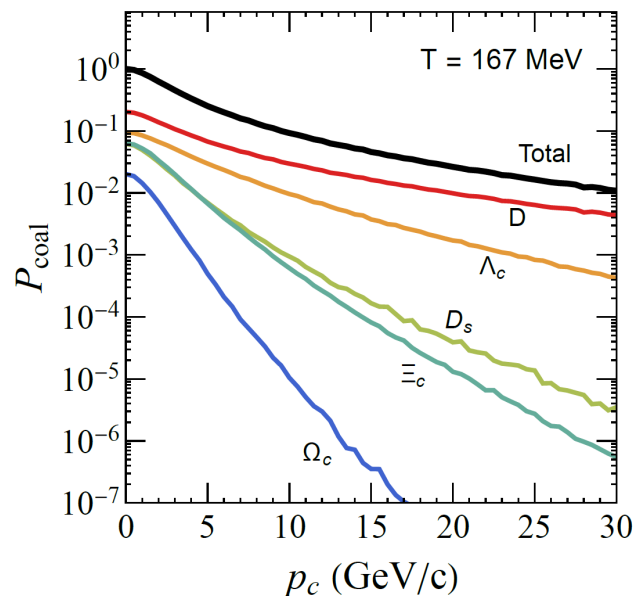


In high multiplicity pp collisions the critical energy density is reached what makes several experimental observations understandable



With increasing N_{ch}
more Q pass a QGP
saturates at $N_{ch} \approx 40$

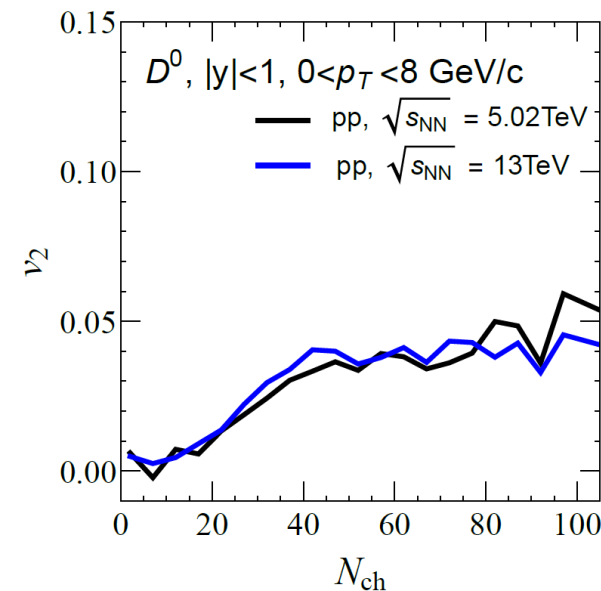
QGP hadronizes
by coalescence
and not by
fragmentation.



In EPOS4 without QGP: ratio as in e^+e^-

v_2 depends on N_{ch}
saturates when
all heavy quarks
pass a QGP ($N_{ch} \approx 40$)

High multiplicity pp collisions
seem not to be elementary collisions



Correlations between Q and Qbar

Correlations between Q and Qbar are important

if one wants to study/understand D Dbar correlation

if one wants to study hidden heavy flavour mesons like J/ψ

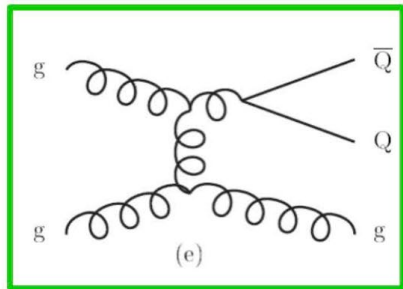
if one wants to understand the p_T distributions of heavy hadrons

FONLL only single particle p_T spectrum

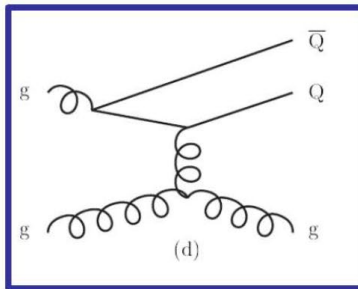
Pythia ISR and FSR can be added

EPOS4HQ separates the three different production mechanisms

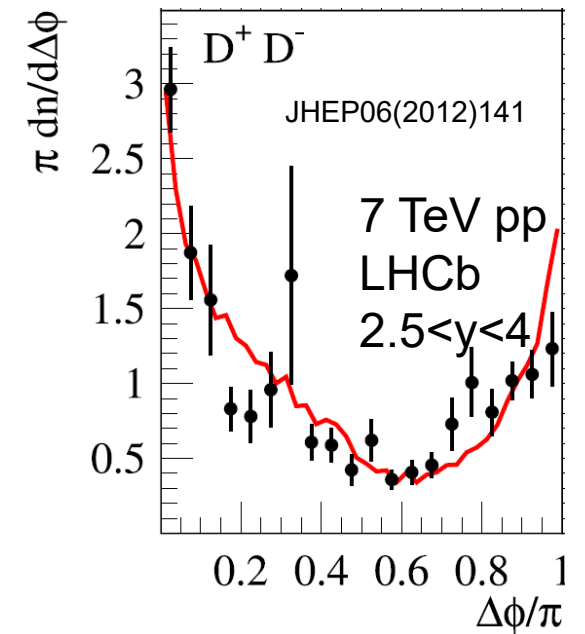
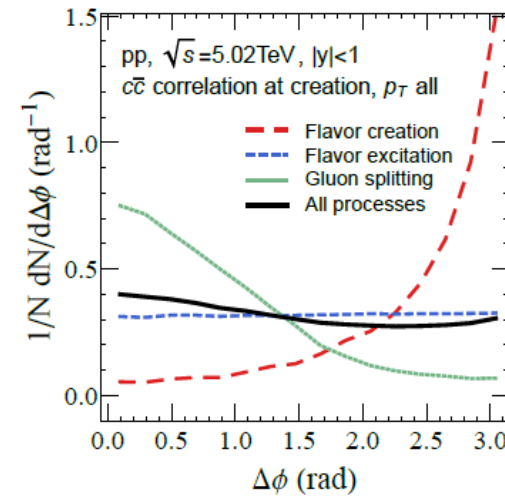
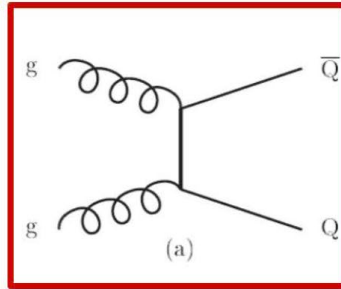
gluon splitting
time like



flavor excitation
space like



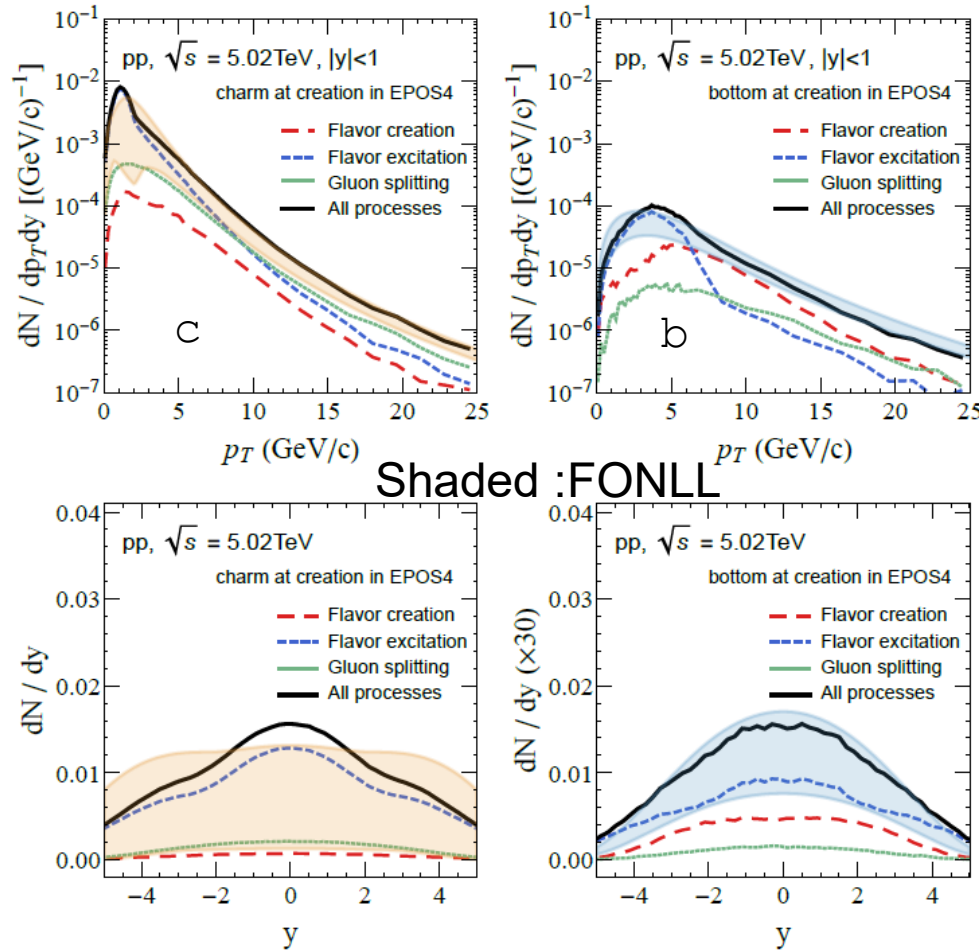
flavour creation
hard process (Born)



Different QCD processes → different correlations → seen in data

Correlations between Q and Qbar

p_T and y distributions depend on creation mechanism



For b and c quarks the importance of different contributions are different

High p_T c \rightarrow gluon splitting

High p_T b \rightarrow flavor creation
(more energy avail.)

low p_T flavor excitation dominating

Spectra (sum of all contributions)
agree with FONLL (shaded area)

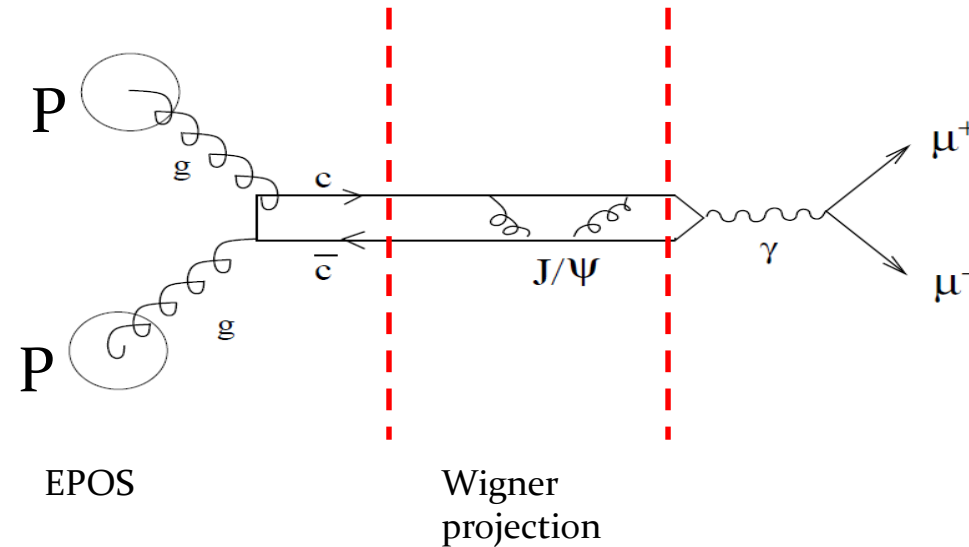
Formation of bound $Q\bar{Q}$ states

Correlations between Q and \bar{Q} show also up in quarkonium production

How to describe a **bound** state like a $c\bar{c}$ in QCD?

It involves low momenta and needs **non perturbative** input \rightarrow assumptions.

Our approach: **Wigner density** formalism (as successful at lower energies)



Formation of bound $Q\bar{Q}$ states

Wigner Density Formalism

c-cbar interaction depends on relative \mathbf{p} and \mathbf{r} only, \rightarrow plane wave of CM

Starting point: Wave function (w.f.) of the relative motion of state i : $|\Phi_i\rangle$

w.f. \rightarrow density matrix $\rho_i = |\Phi_i\rangle\langle\Phi_i|$

Wigner density of $\Phi_i^W(\mathbf{r}, \mathbf{p}) = \int d^3y e^{i\mathbf{p}\cdot\mathbf{y}} \langle \mathbf{r} - \frac{1}{2}\mathbf{y} | \Phi_i \rangle \langle \Phi_i | \mathbf{r} + \frac{1}{2}\mathbf{y} \rangle$.

(close to classical phase space density)

The probability to form a state $|\Phi_i\rangle$ if the system is described by $\rho^N = |\psi^N\rangle\langle\psi^N|$

$$P = \text{Tr}[\rho^N \rho_i] = \int \prod d^3p_i d^3r_i W^i(\mathbf{r}, \mathbf{p}) W^N(\mathbf{r}_1 \dots \mathbf{r}_N, \mathbf{p}_1 \dots \mathbf{p}_N)$$

Where \mathbf{r}, \mathbf{p} are relative coordinates.

Formation of bound $Q\bar{Q}$ states

Step 1: solution of $Q\bar{Q}$ Schrödinger eq.

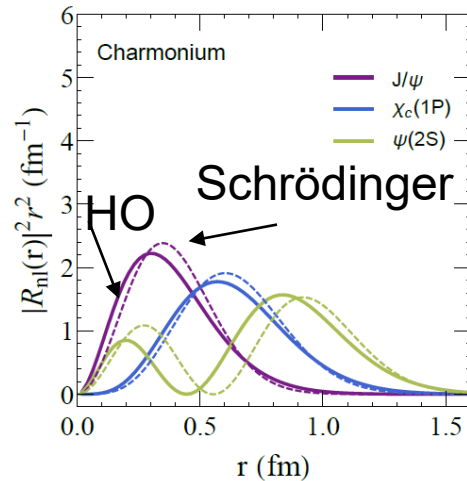
$$\left[-\frac{1}{2\mu} \left(\frac{d^2}{dr^2} + \frac{2}{r} \frac{d}{dr} \right) + \frac{l(l+1)}{2\mu r^2} + V(r) \right] R_{nl}(r) = E R_{nl}(r)$$

Step 2: Conversion into a 3d harmonic oscillator wave fct with same spin and rms radius $\psi^{HO}(\mathbf{r})$

$$V(r) = -\alpha/|\mathbf{r}| + \sigma|\mathbf{r}| \text{ with } \alpha = 0.513, \sigma = 0.17 \text{ GeV}^2, m_c = 1.5 \text{ GeV}, m_b = 5.2 \text{ GeV}.$$

$$\mathbf{r} = \mathbf{r}_i - \mathbf{r}_j$$

$$\psi(\mathbf{r}) = R_{nl}(r) Y_{l,m}(\theta, \phi).$$



and conversion into a 2 body Wigner density:

$$W_{nl}(\mathbf{r}, \mathbf{p}) = \int d^3y e^{i\mathbf{p}\mathbf{y}} \langle \mathbf{r} + \mathbf{y}/2 | \psi^{HO} \rangle \langle \psi^{HO} | \mathbf{r} - \mathbf{y}/2 \rangle$$

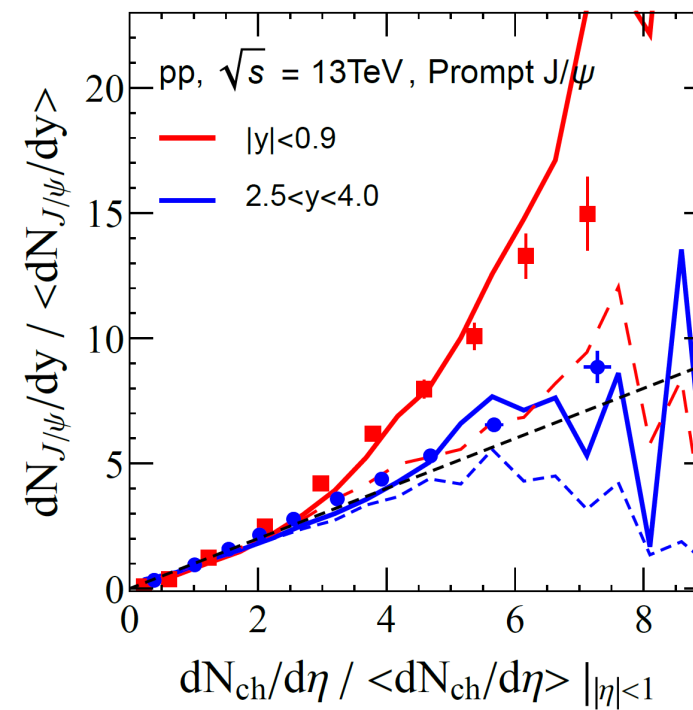
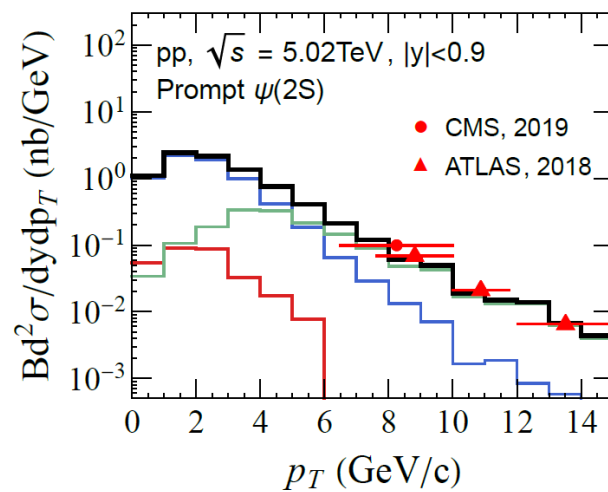
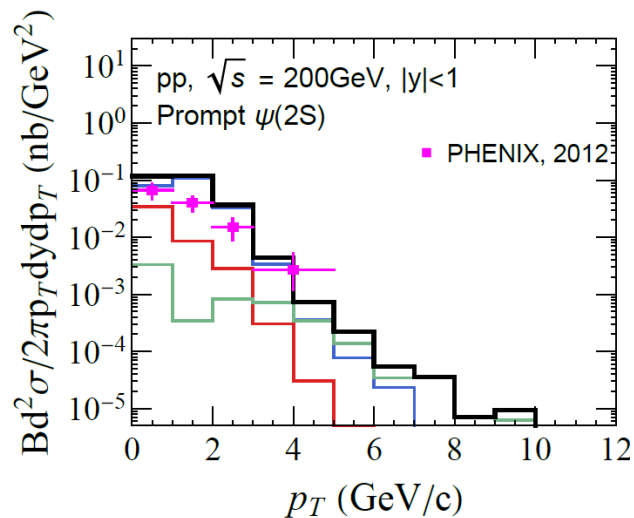
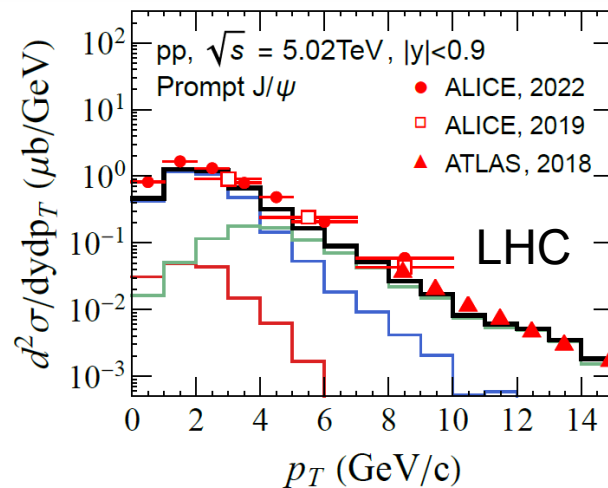
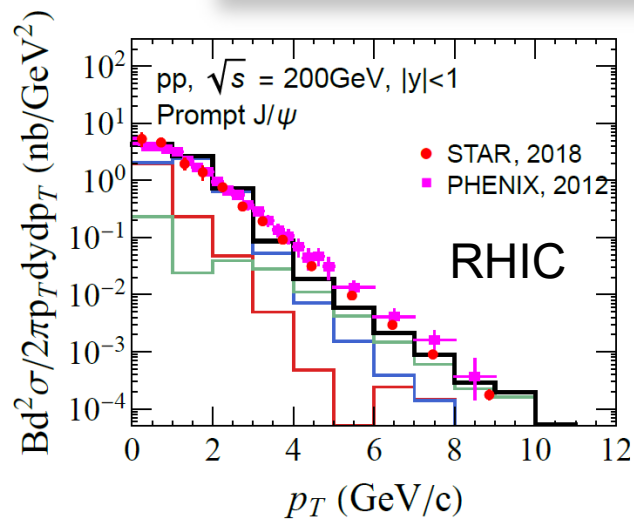
Step 3: Probability that a quarkonium with quantum numbers n, l is produced is given by the overlap integral

$$\frac{dP_{nl}^m}{d^3\mathbf{P}_{\text{cm}}} = \sum \int \frac{d^3r d^3p}{(2\pi)^6} W_{nl}^m(\mathbf{r}, \mathbf{p}) W^{(2)}(\mathbf{P}_{\text{cm}}, \mathbf{r}, \mathbf{p})$$

$$W^{(2)}(\mathbf{P}, \mathbf{r}, \mathbf{p}) \sim r^2 \exp \left(-\frac{r^2}{2\sigma_{Q\bar{Q}}^2} \right) f_{Q\bar{Q}}^{\text{EPOS4}}(\mathbf{P}, \mathbf{p})$$

= 2-body $Q\bar{Q}$ phase space distr. by EPOS4

J/ψ and ψ(2s) spectra at RHIC and LHC

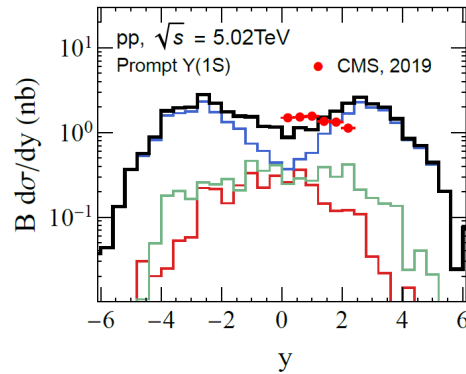
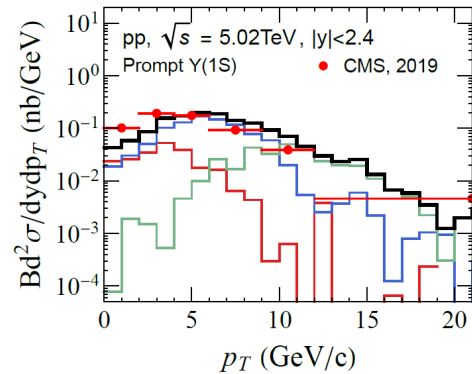


scaled multiplicities at
different rapidities

high p_T spectra dominated by gluon splitting
low p_T by flavor excitation

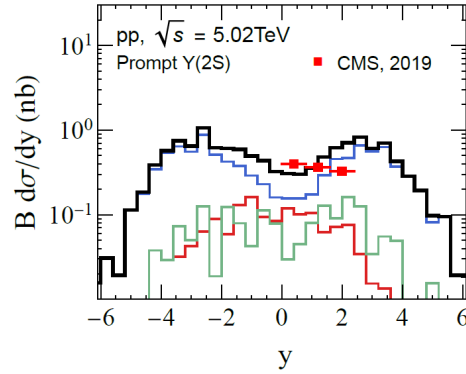
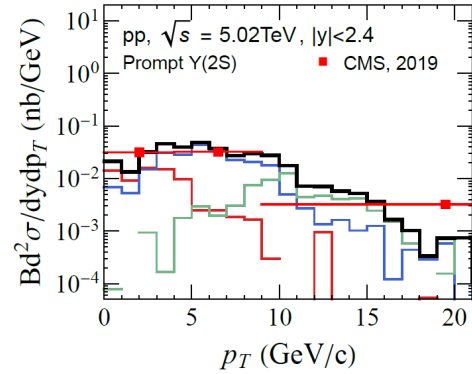
Without understanding of the QCD processes
we do not understand the spectra

Y(1), Y(2), Y(3) and B_c

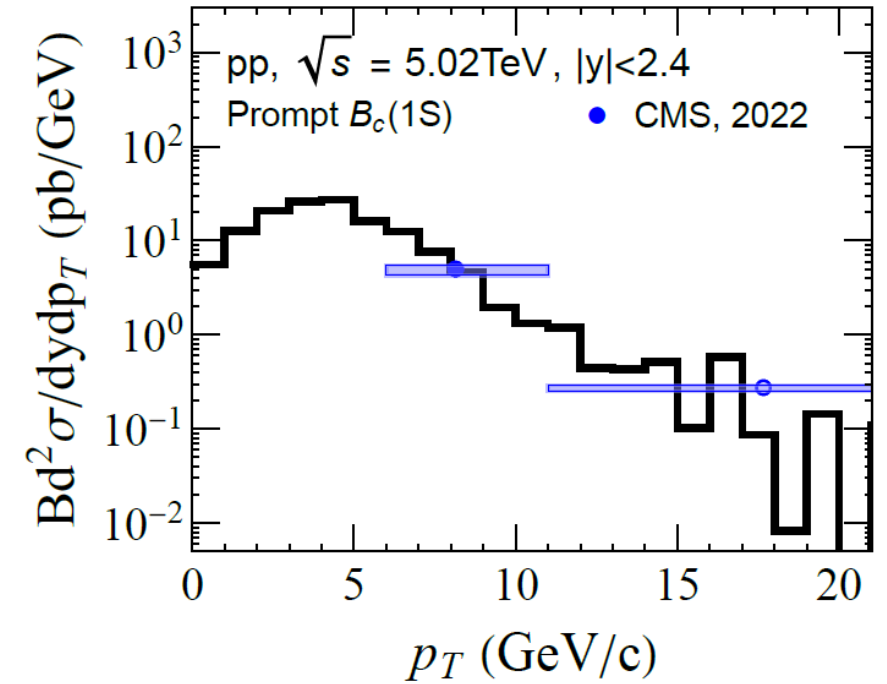
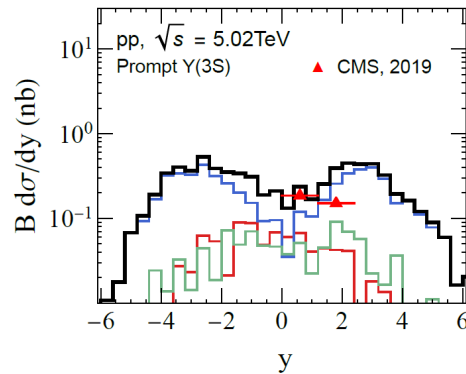
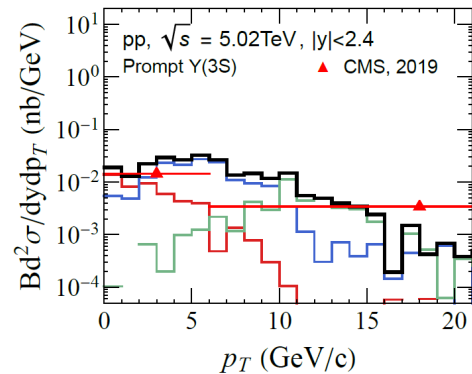


high p_T spectra
dominated by gluon
splitting

in difference to the



high p_T spectra
of B mesons,
dominated by flavor
production



Coming from different
production vertices
-> tests the spacial structure
of the interaction zone

Conclusion

astonishing heavy flavor results in pp collisions (finite v_2 , $\Lambda_c/D(p_T)$) can be understood if one assumes that in (high multiplicity) pp events, as in heavy ions, a quark gluon plasma is formed if

$$\varepsilon > \varepsilon_0 = 0.57 \text{ GeV/fm}^3$$

The azimuthal $D\bar{D}$ correlation points towards $D\bar{D}$ production by different QCD processes:

In EPOS4HQ The different QCD production processes contribute to the p_T spectra in different ways

- high p_T spectra dominated by gluon splitting
- low p_T spectra by flavor excitation
- the sum of all processes compatible with FONLL

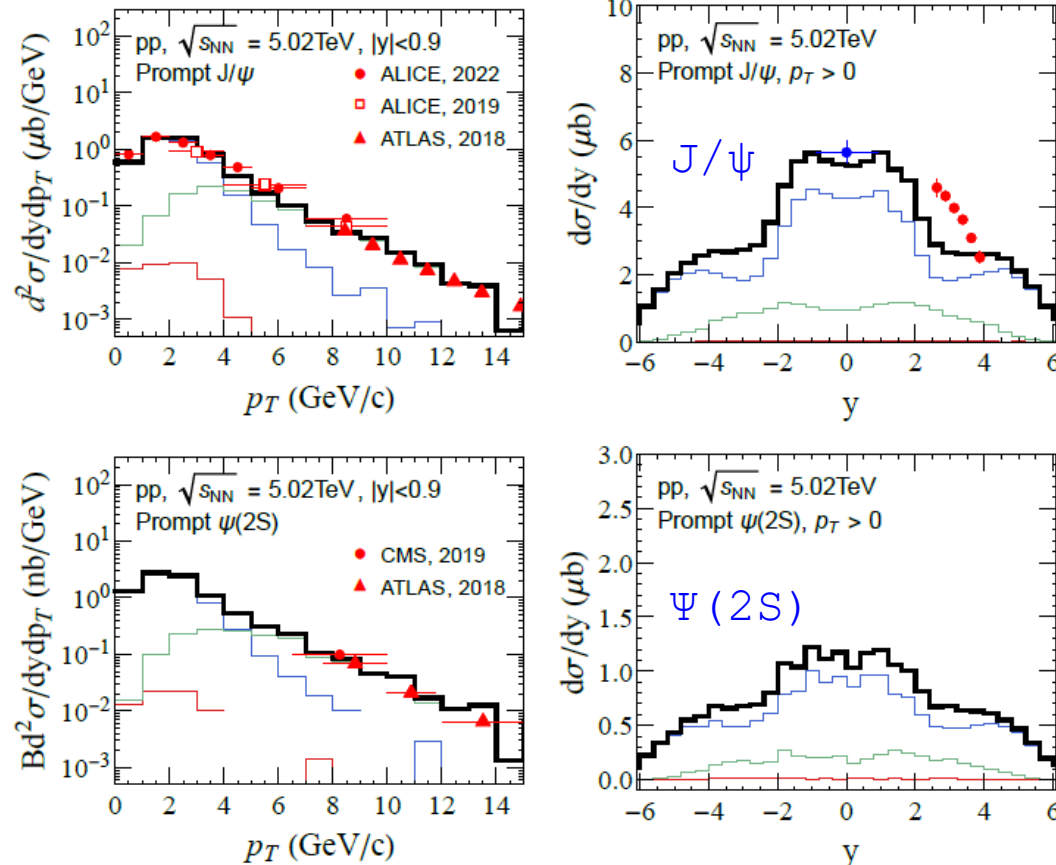
Describing of the production of hidden heavy flavor mesons by a density matrix approach we reproduce the experimental spectra of J/ψ , $\psi(2S)$, $Y(n)$ and B_c

Detailed analysis shows that the contribution of the QCD processes varies with y and p_T so

run 3 data: perspective to study different aspects of QGP/QCD in detail

Correlations between Q and Qbar

Prompt J/ψ spectrum and contribution of the different Q Qbar creation processes



high p_T :
dominated by gluon splitting

flavor creation does not
play a role

low p_T :
Dominated by flavor excitation

Without understanding the
correlations one cannot
understand J/ψ production