

Searches for Gravitational Wave Background and Future Prospects

on behalf of the LIGO-Virgo-KAGRA collaboration

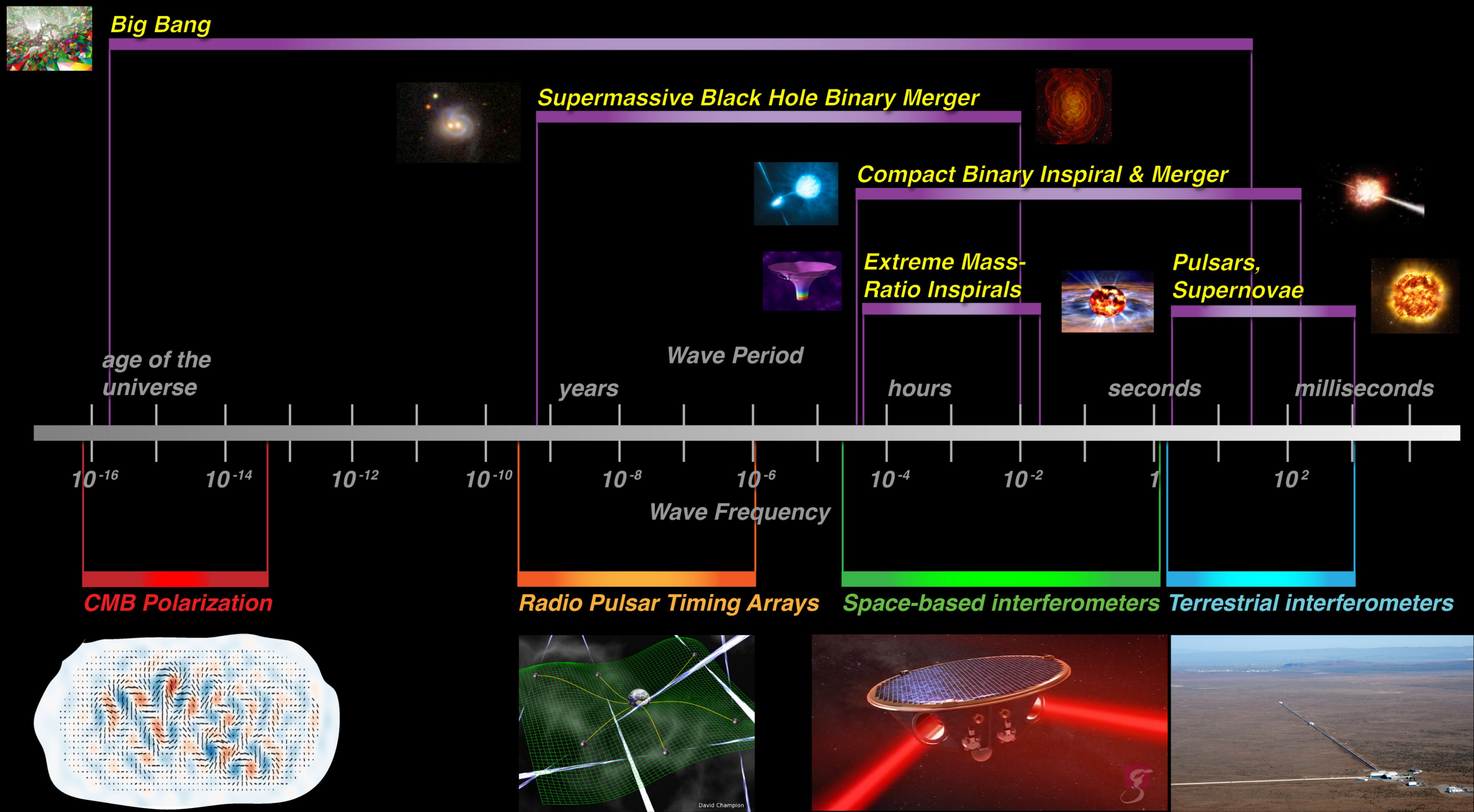
Jishnu Suresh

Henri Poincaré Fellow,

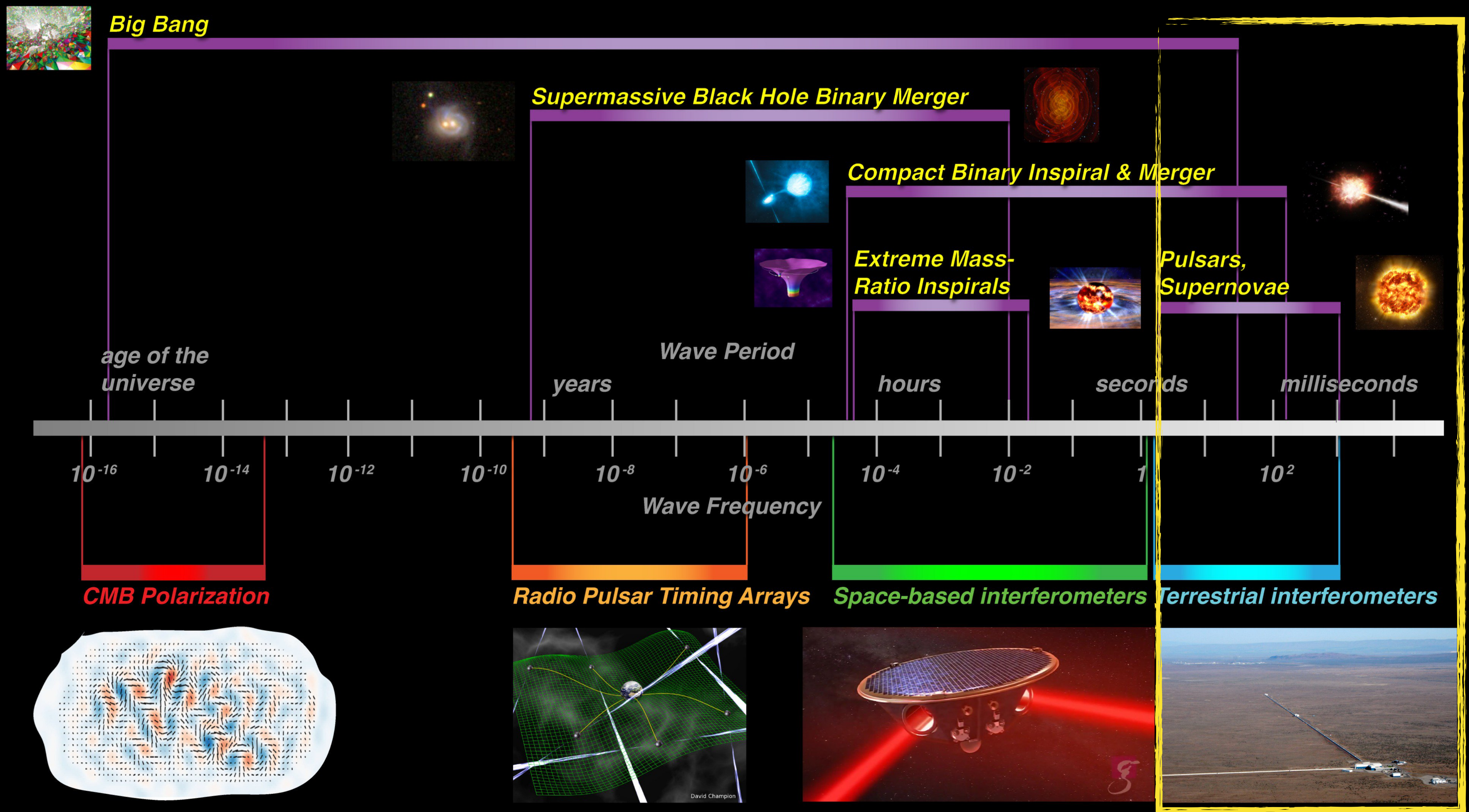
ARTEMIS, Observatoire de la Côte d'Azur, Nice



GRAVITATIONAL WAVE SPECTRA



GRAVITATIONAL WAVE SPECTRA



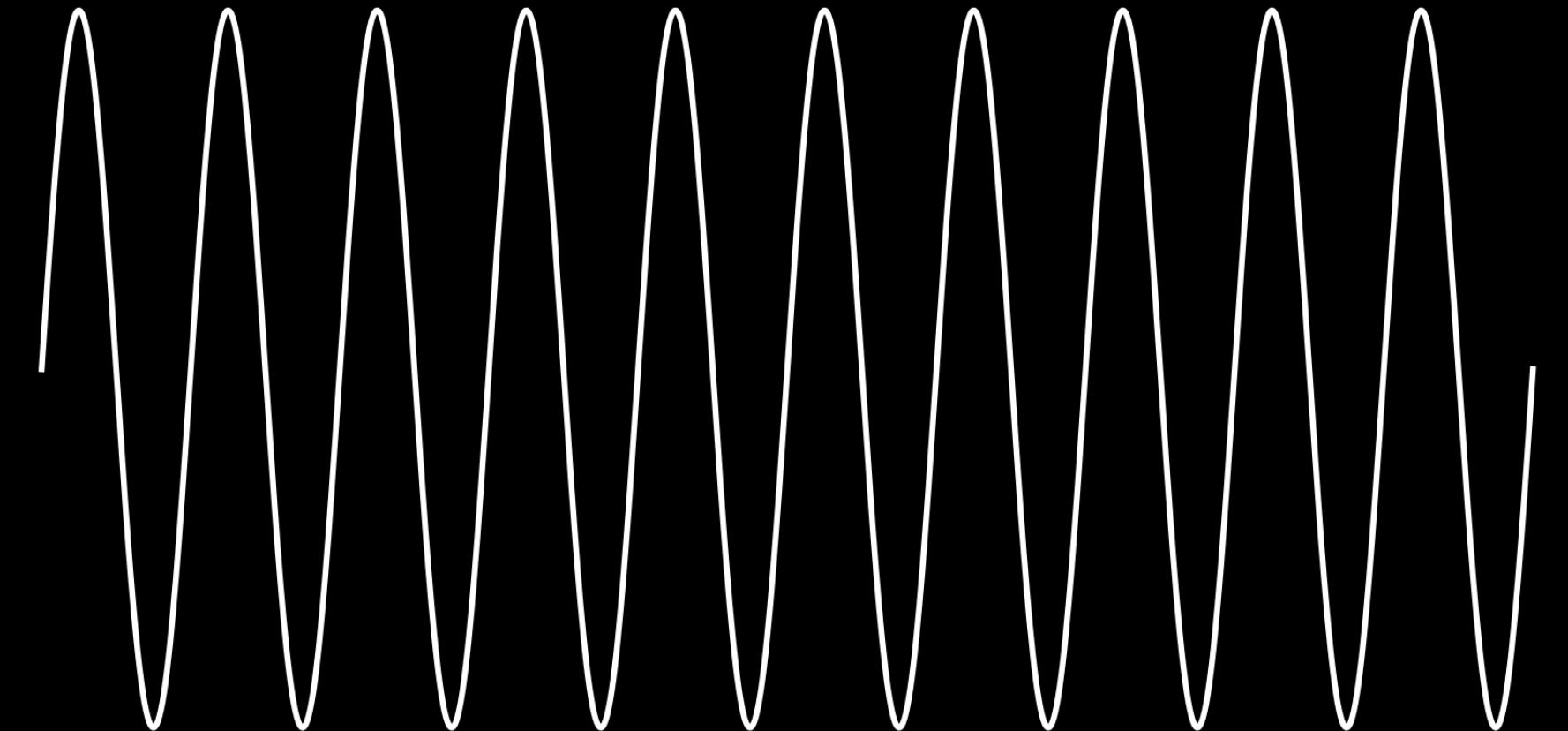
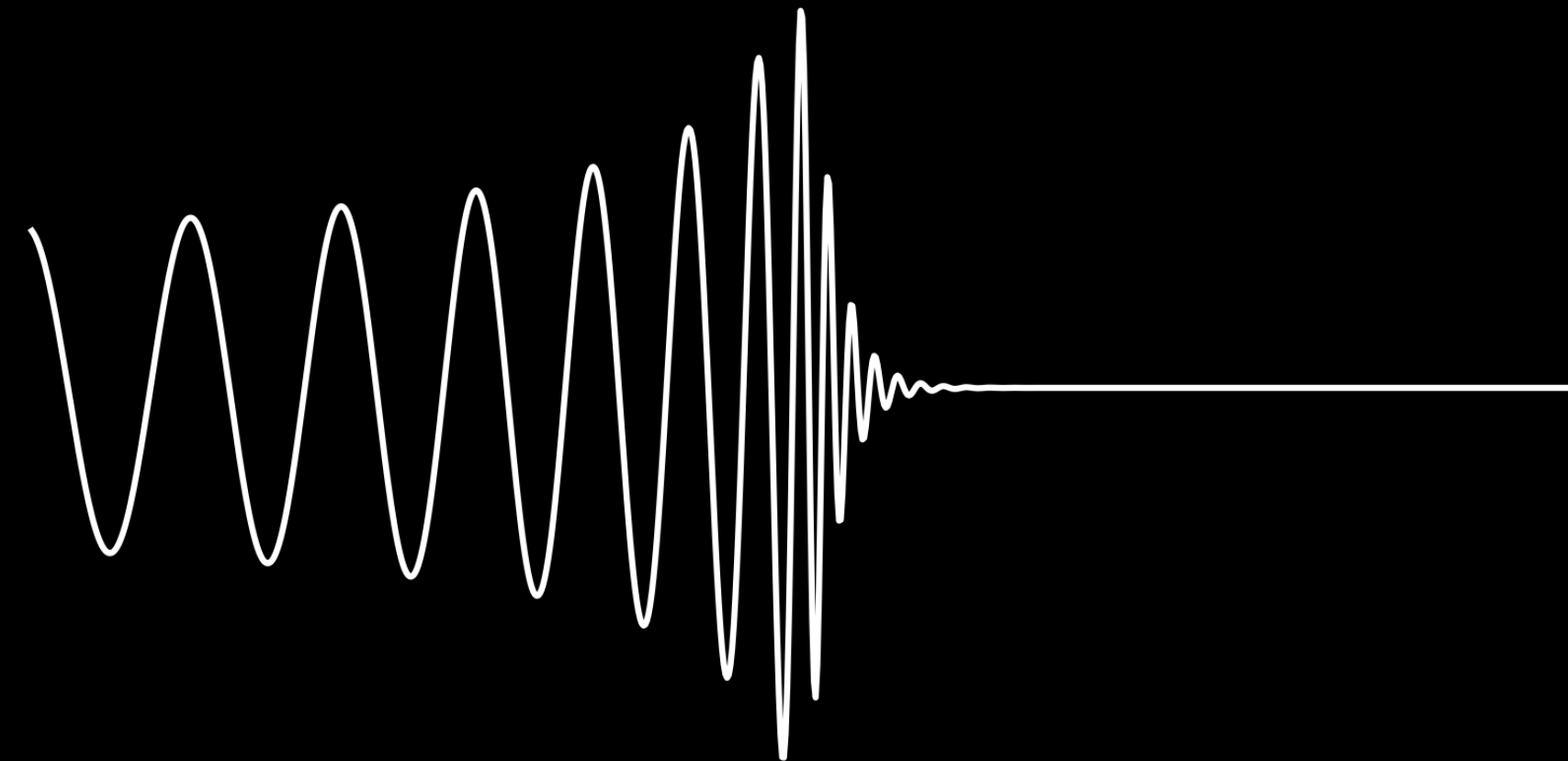
Transient

Persistent

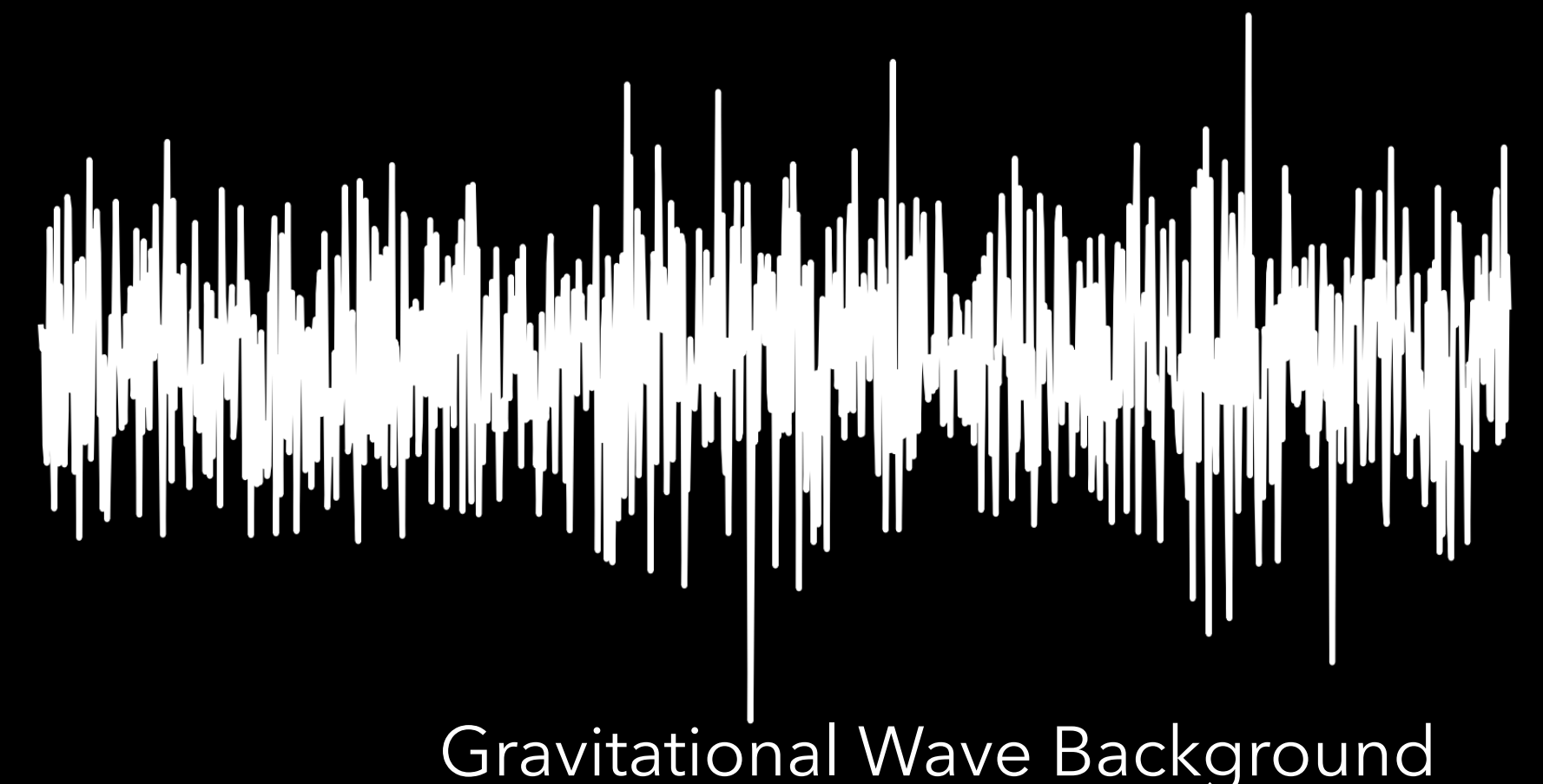
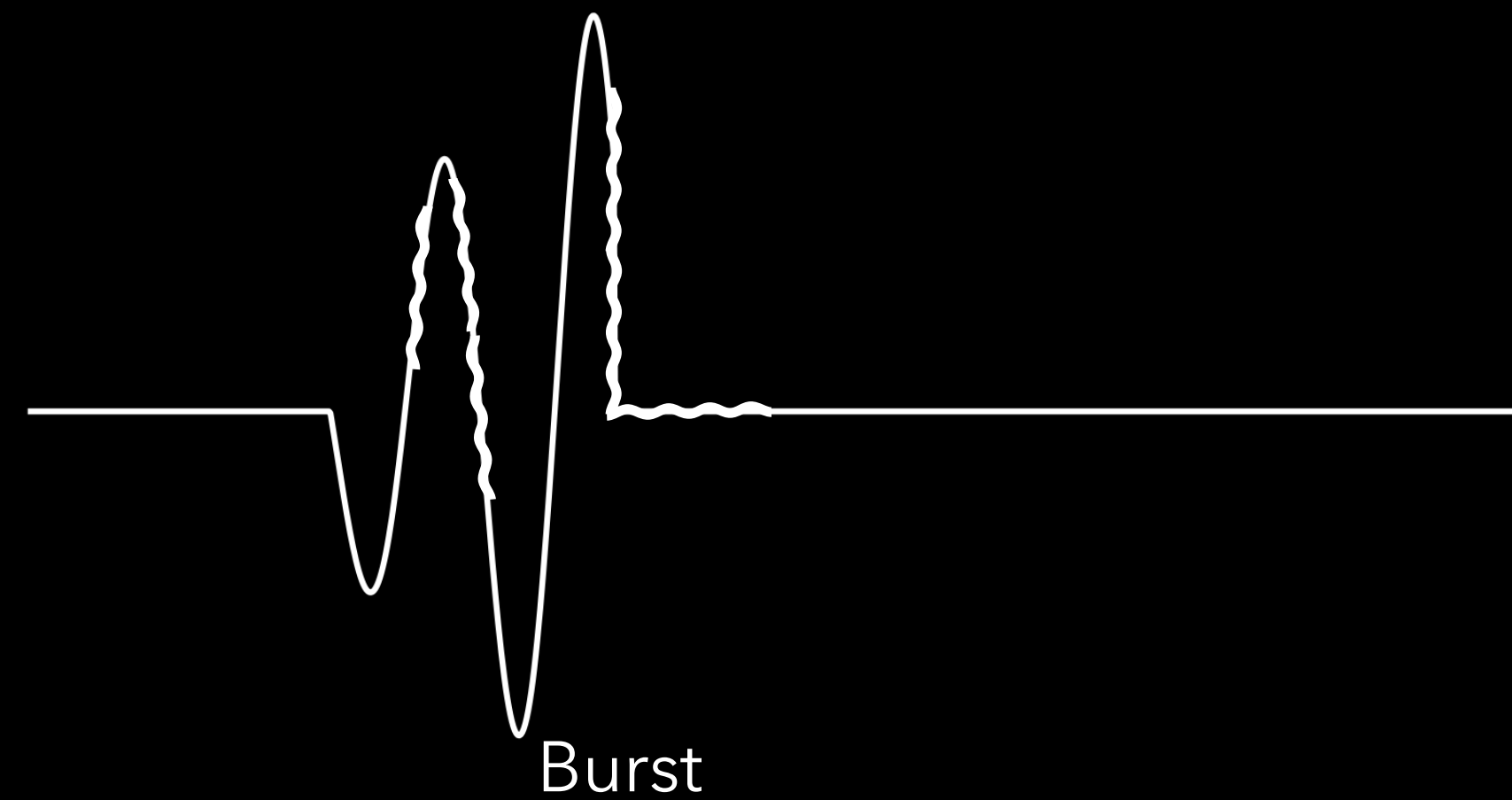
Binary Merger

Continuous Wave

Phase modelled



Phase unmodelled



Gravitational Wave Background

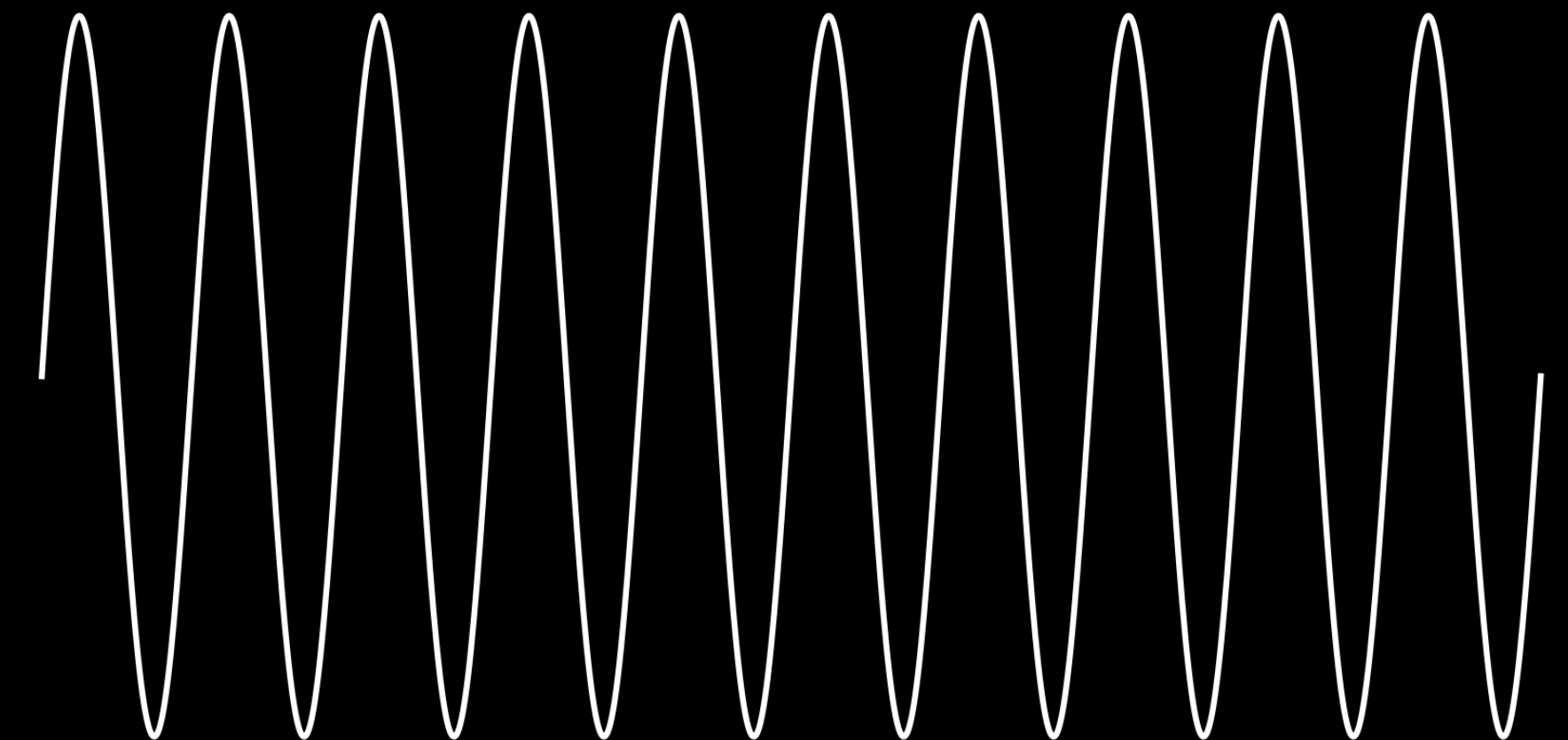
Transient

Persistent

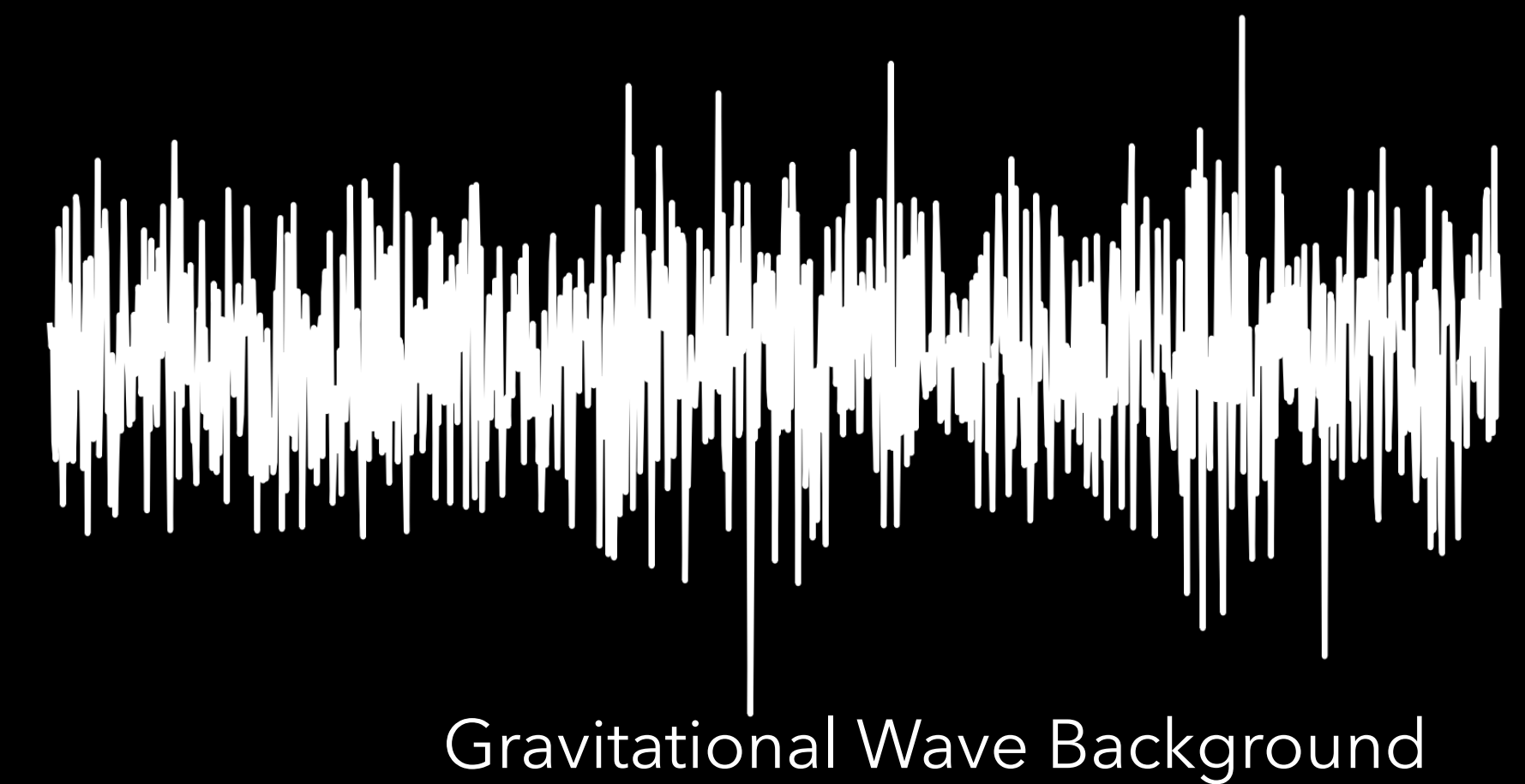
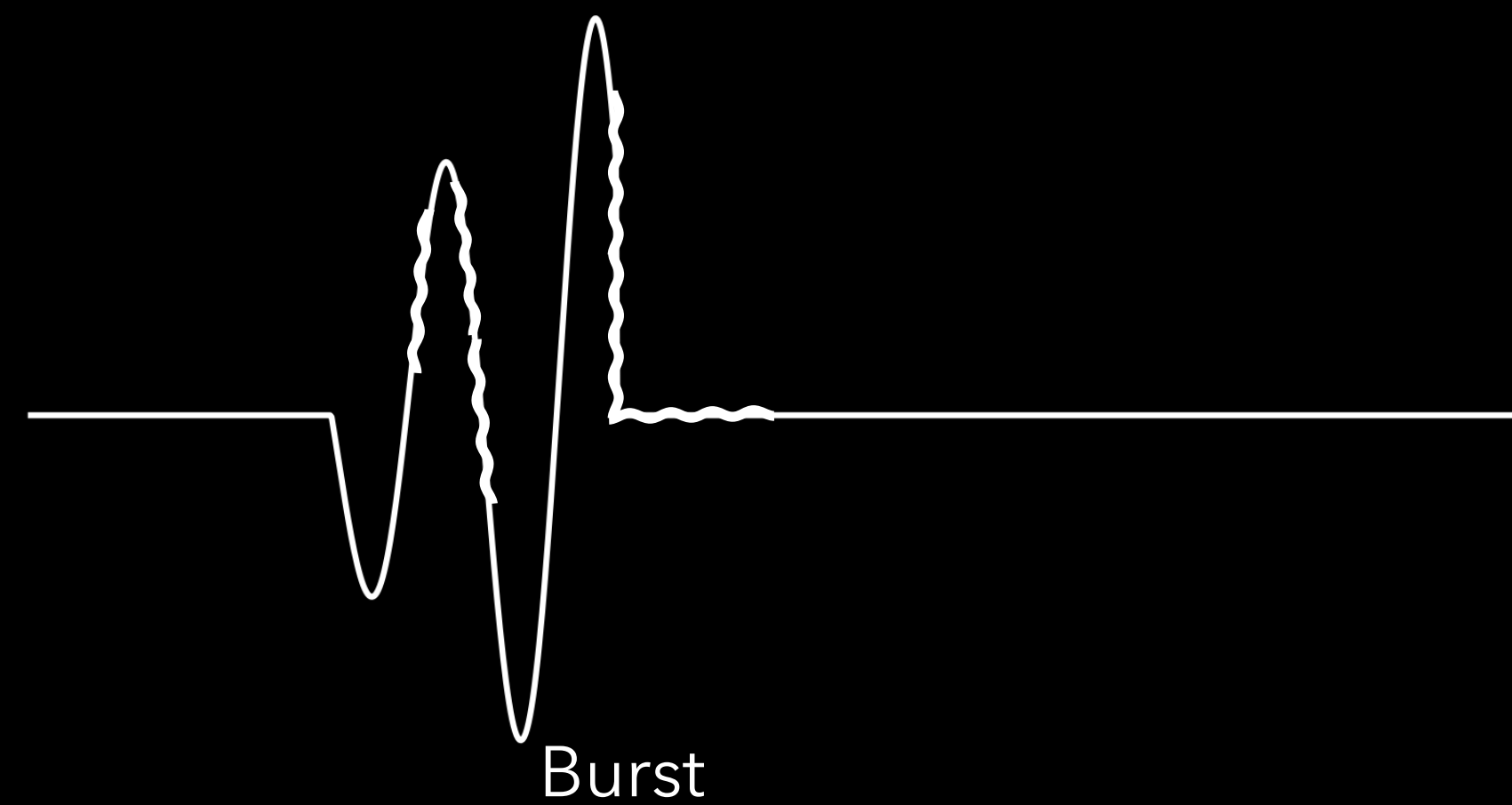
Binary Merger

Continuous Wave

Phase modelled



Phase unmodelled

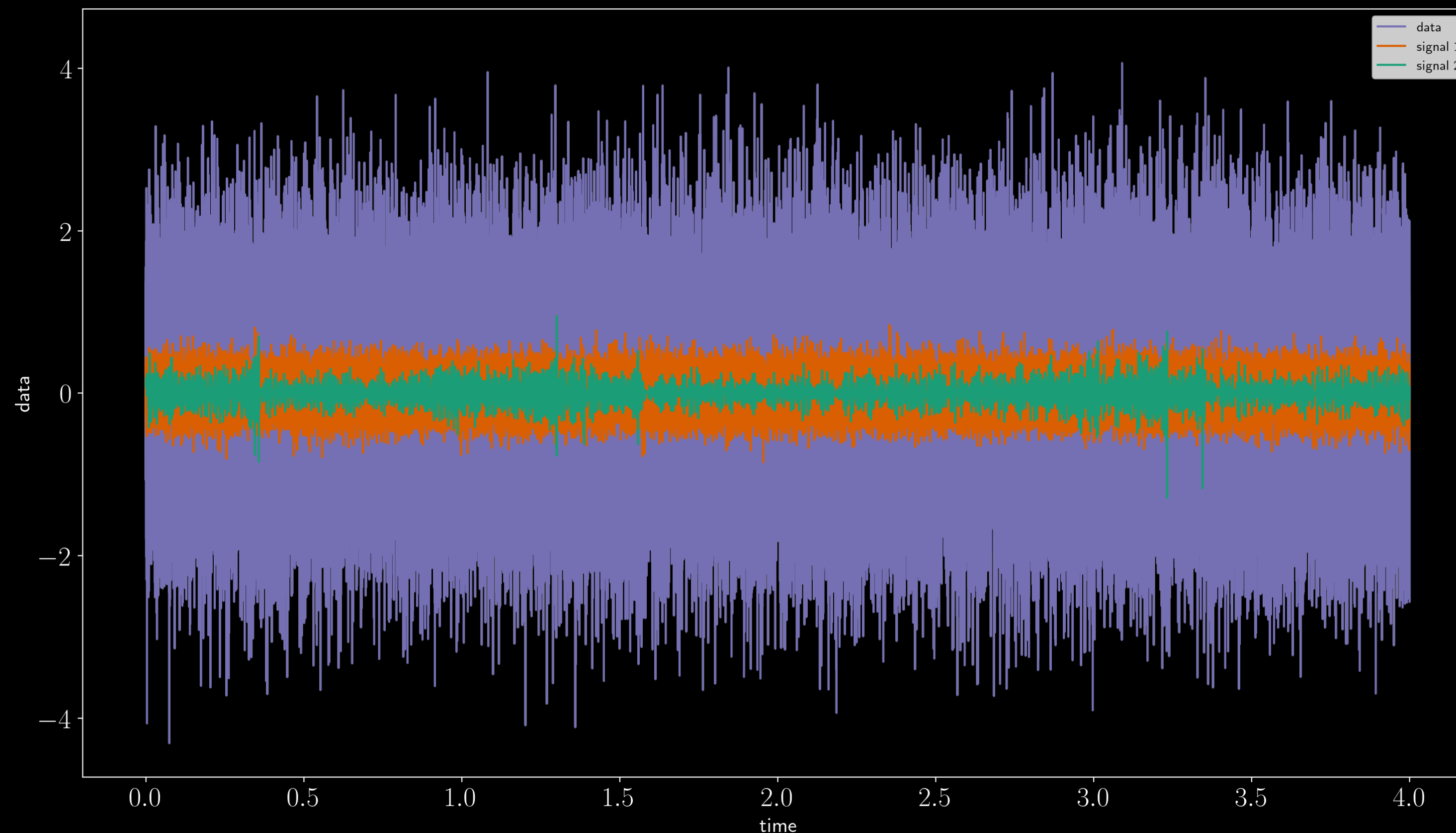


OPERATIONAL DEFINITION

Superposition of signals **too weak** or **too numerous** to individually detect

Looks **like noise** in a single detector

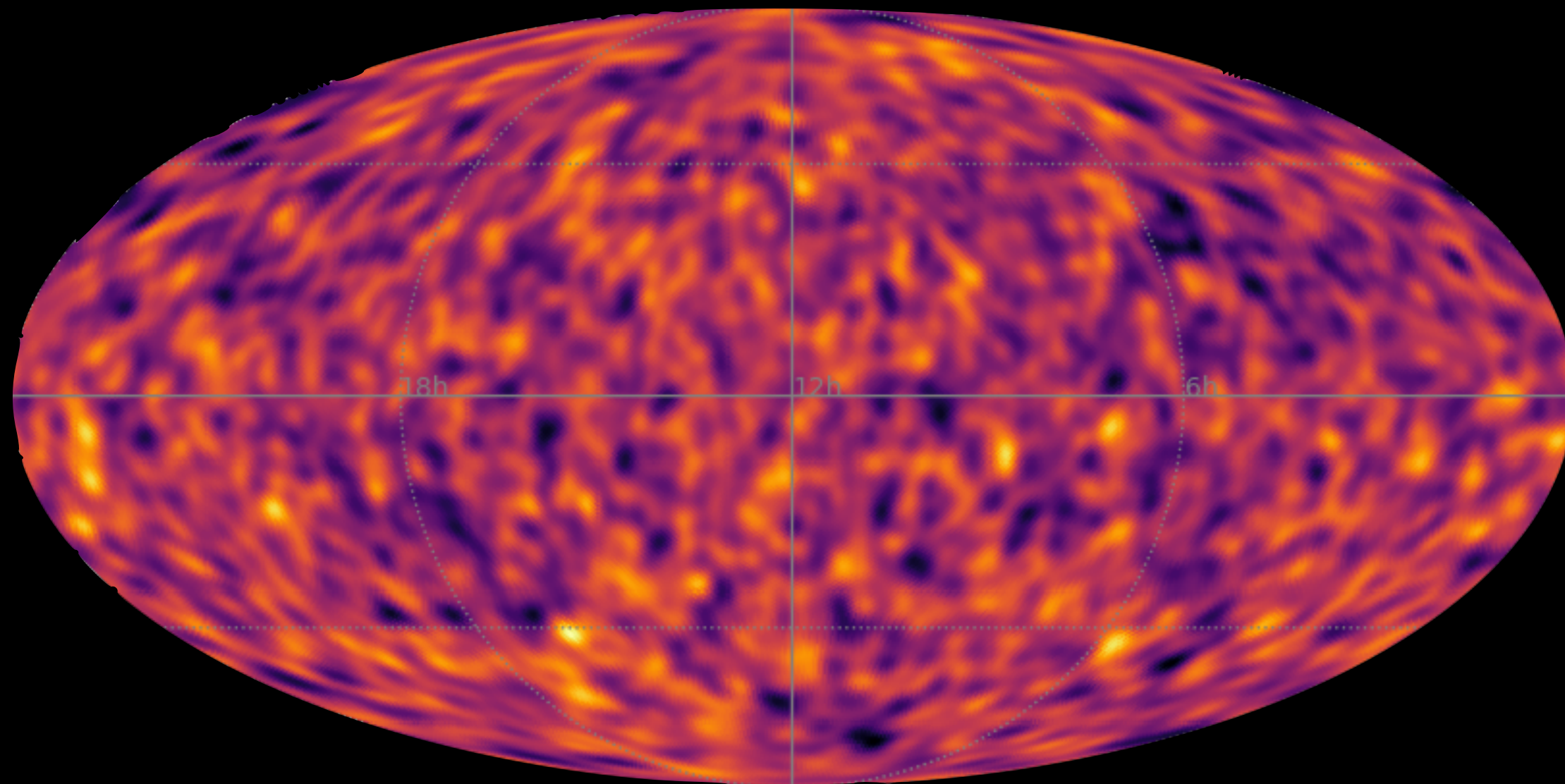
Characterized **statistically** in terms of ensemble averages of the metric perturbations



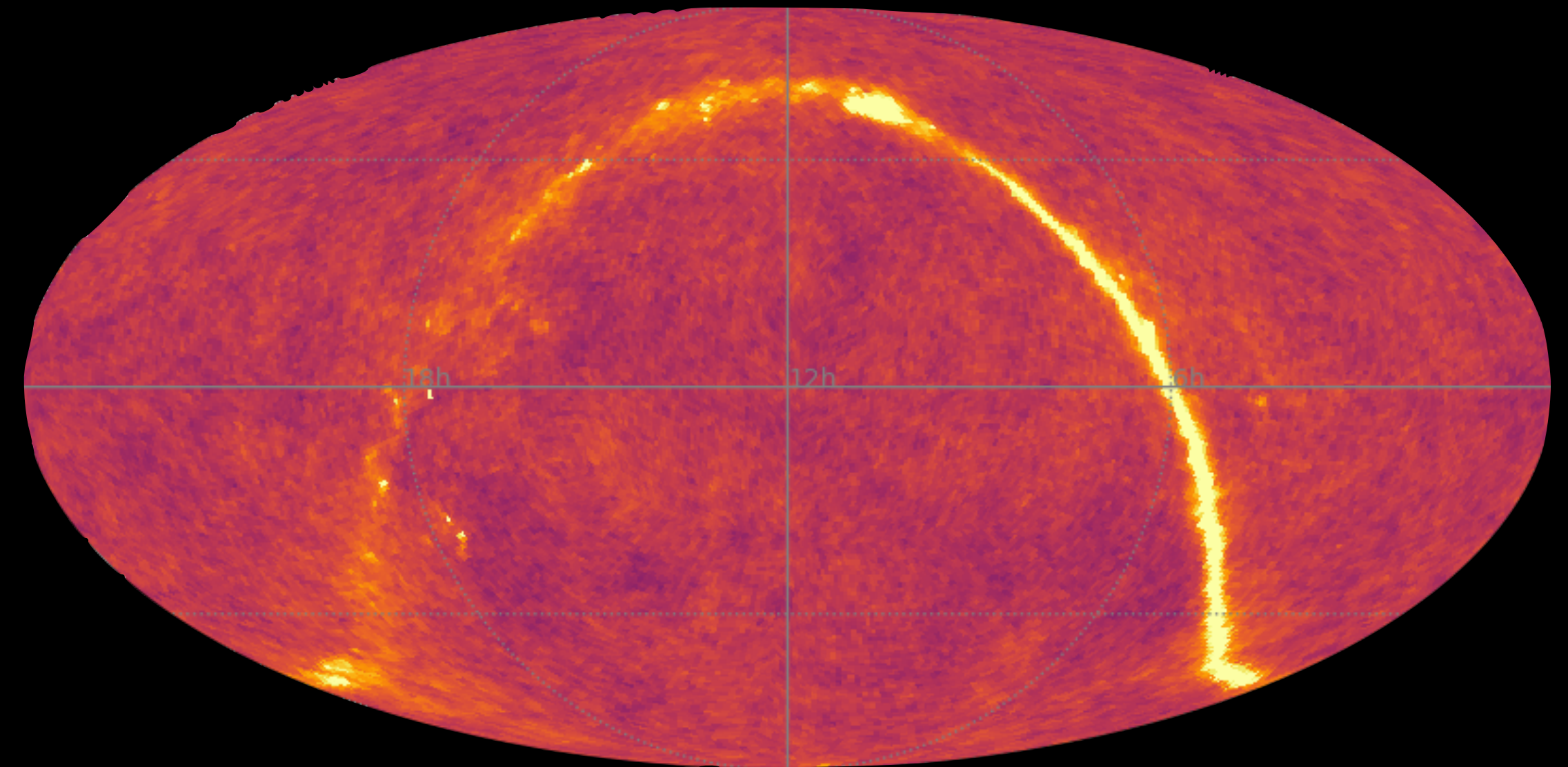
TYPES OF GRAVITATIONAL WAVE BACKGROUND

(1) Spatial distribution

isotropic



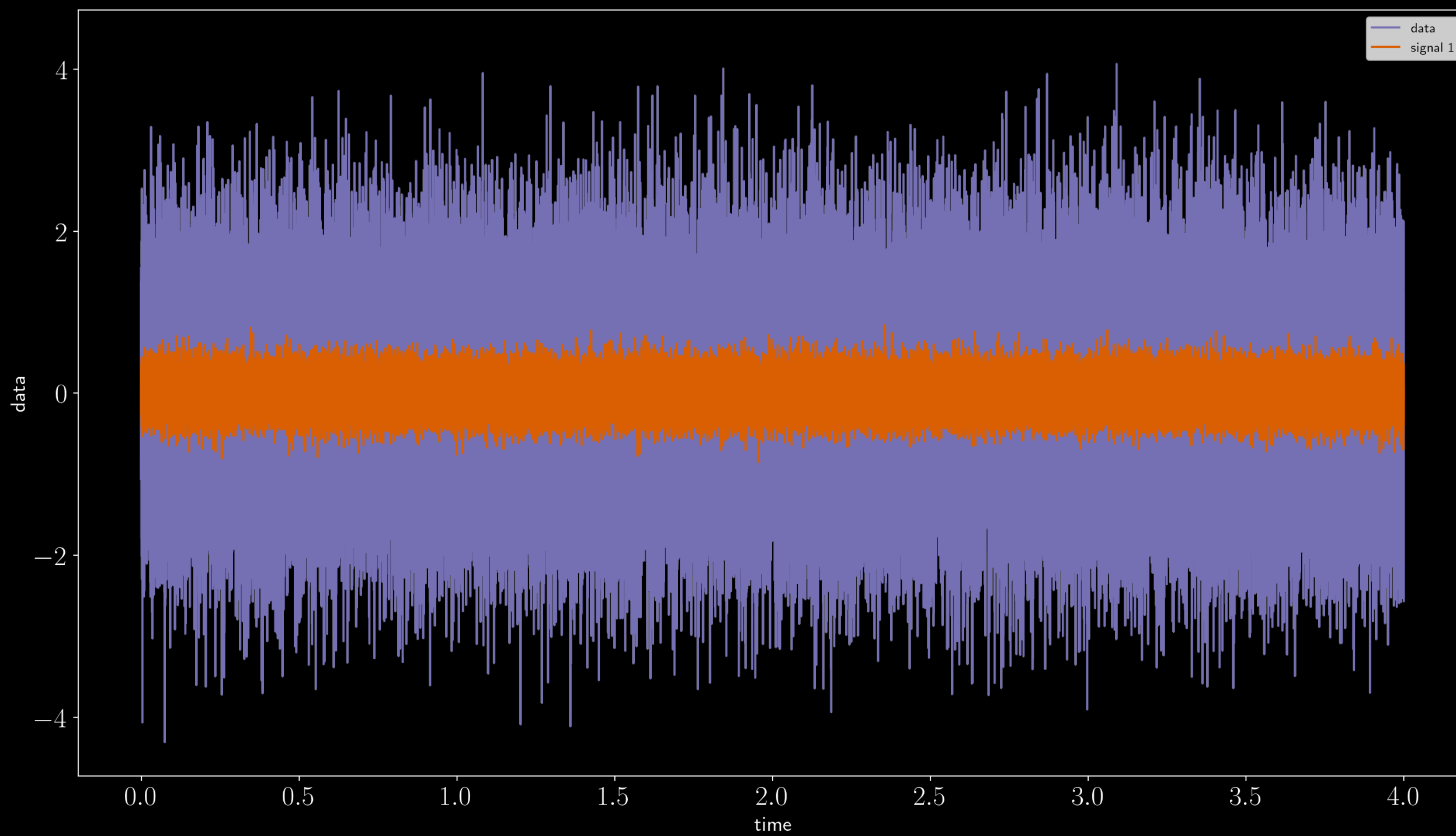
anisotropic



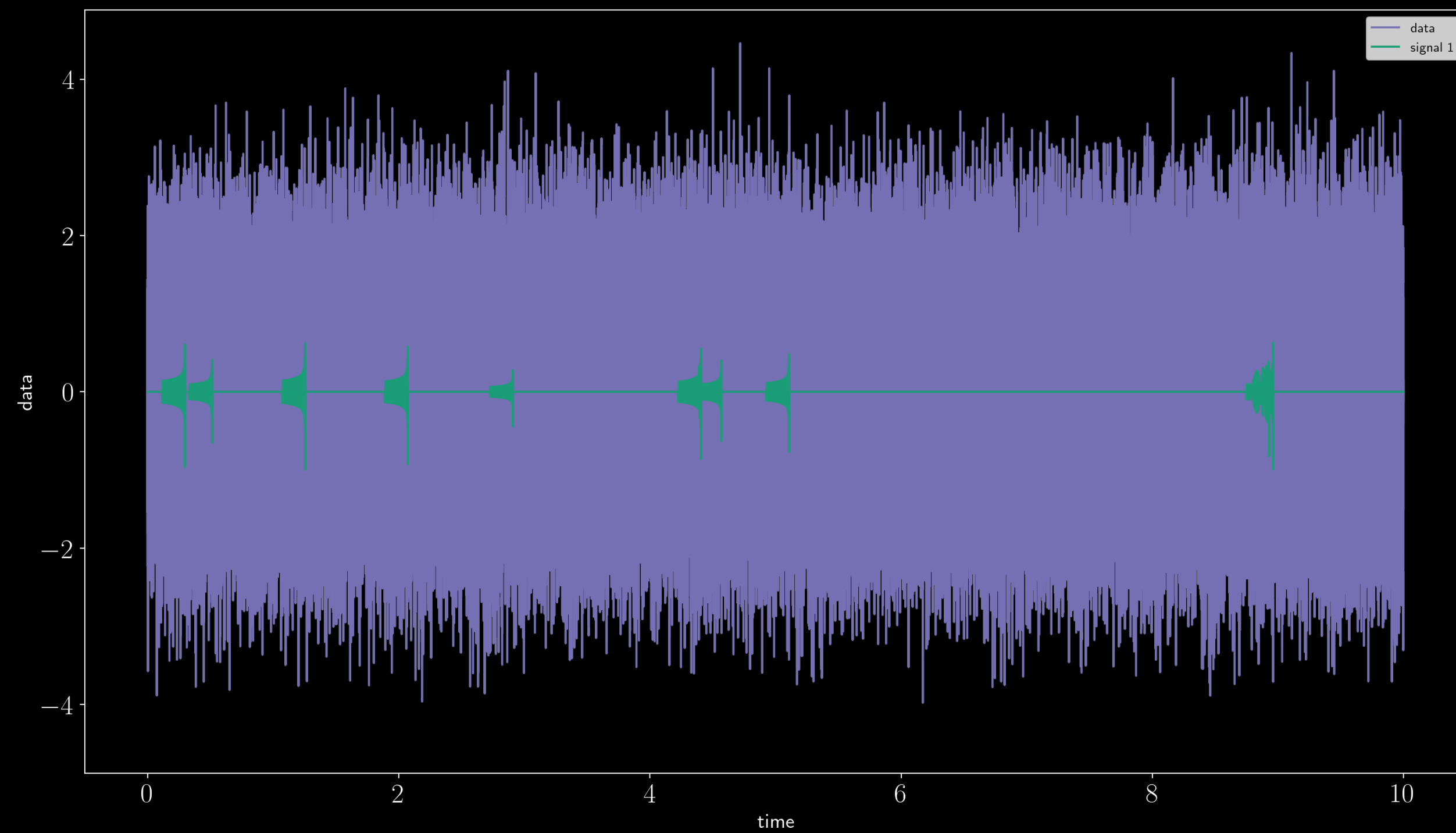
TYPES OF GRAVITATIONAL WAVE BACKGROUND

(2) Temporal distribution

Stationary Gaussian



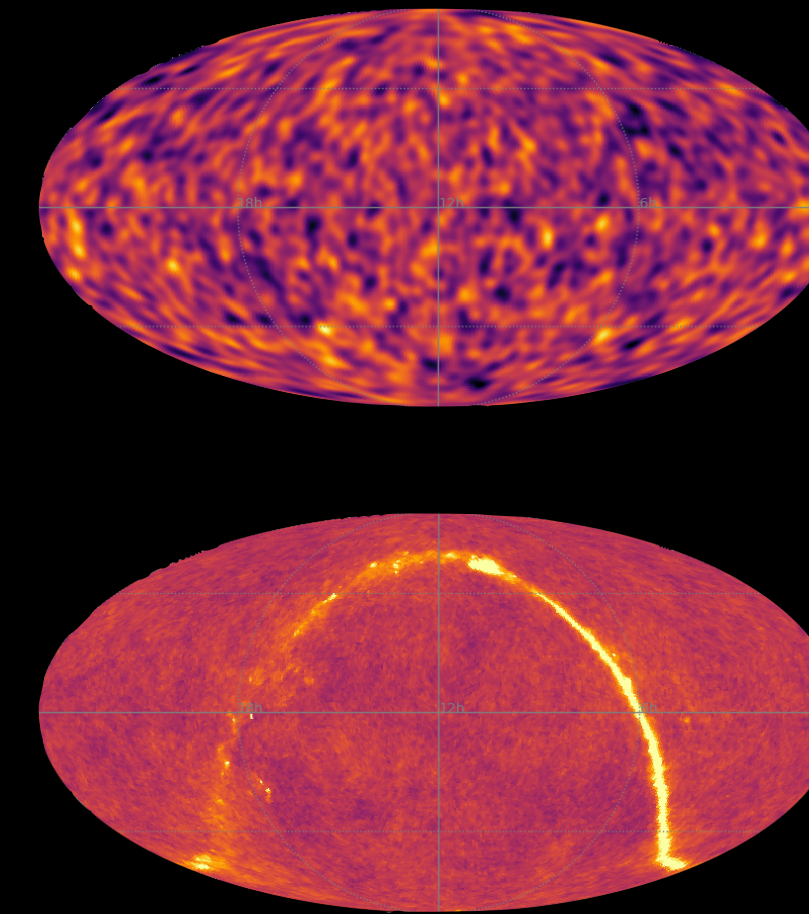
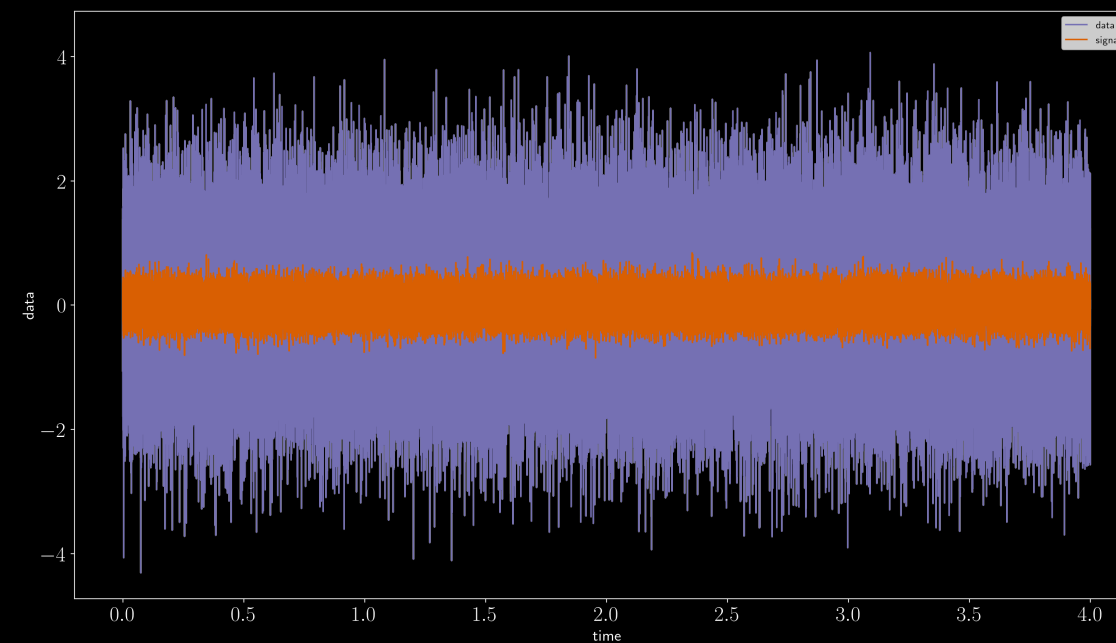
Non-stationary (non-gaussian)



GRAVITATIONAL WAVE BACKGROUND

In this talk, we will only consider the following cases

Unpolarized,



Unpolarized, stationary, **isotropic**:

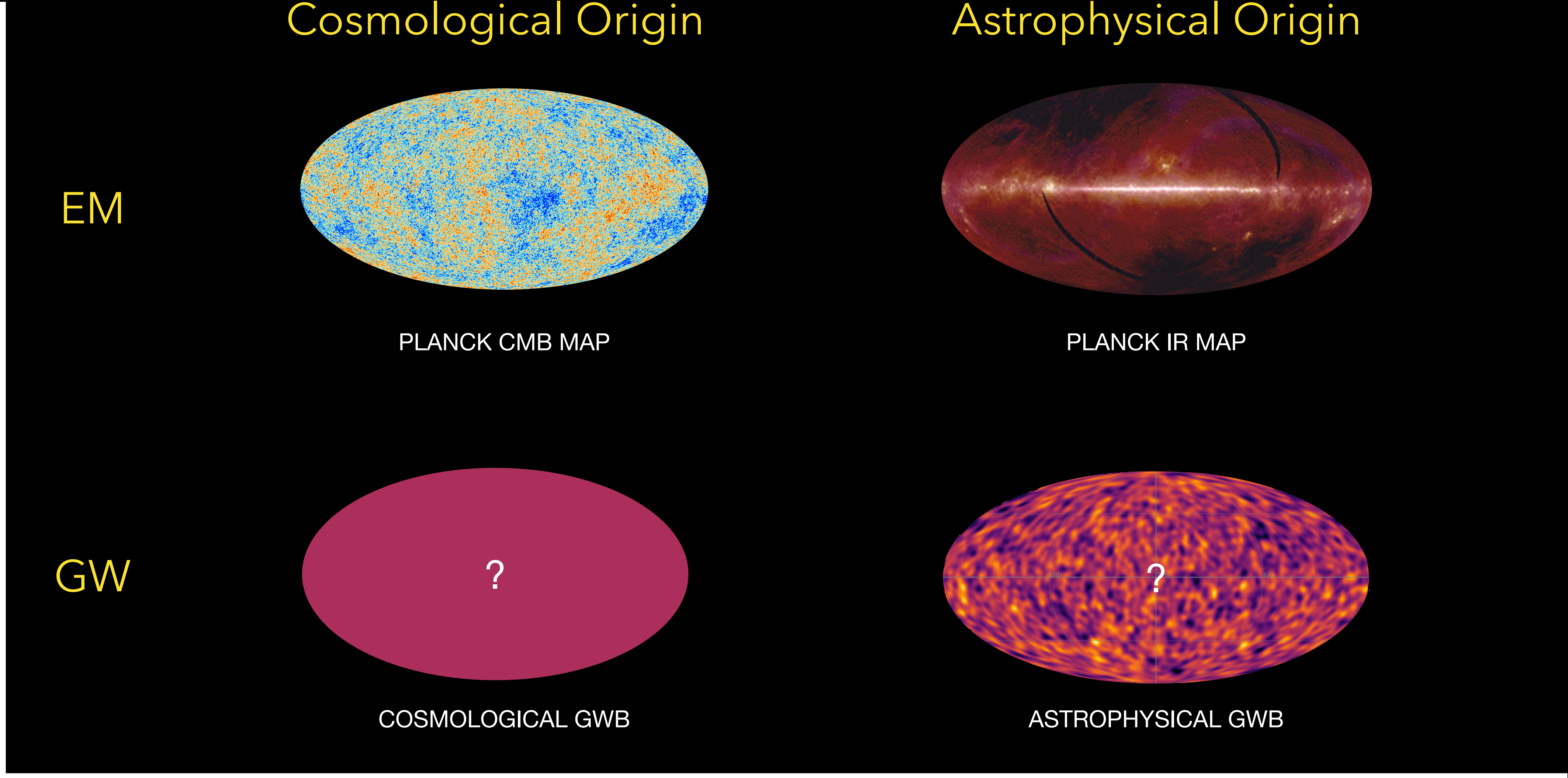
$$\langle h_A(f, \hat{n}) h_{A'}^*(f', \hat{n}') \rangle = \frac{1}{16\pi} S_h(f) \delta(f - f') \delta_{AA'} \delta^2(\hat{n}, \hat{n}')$$

Unpolarized, stationary, **anisotropic**:

$$\langle h_A(f, \hat{n}) h_{A'}^*(f', \hat{n}') \rangle = \frac{1}{4} \mathcal{P}(f, \hat{n}) \delta(f - f') \delta_{AA'} \delta^2(\hat{n}, \hat{n}')$$

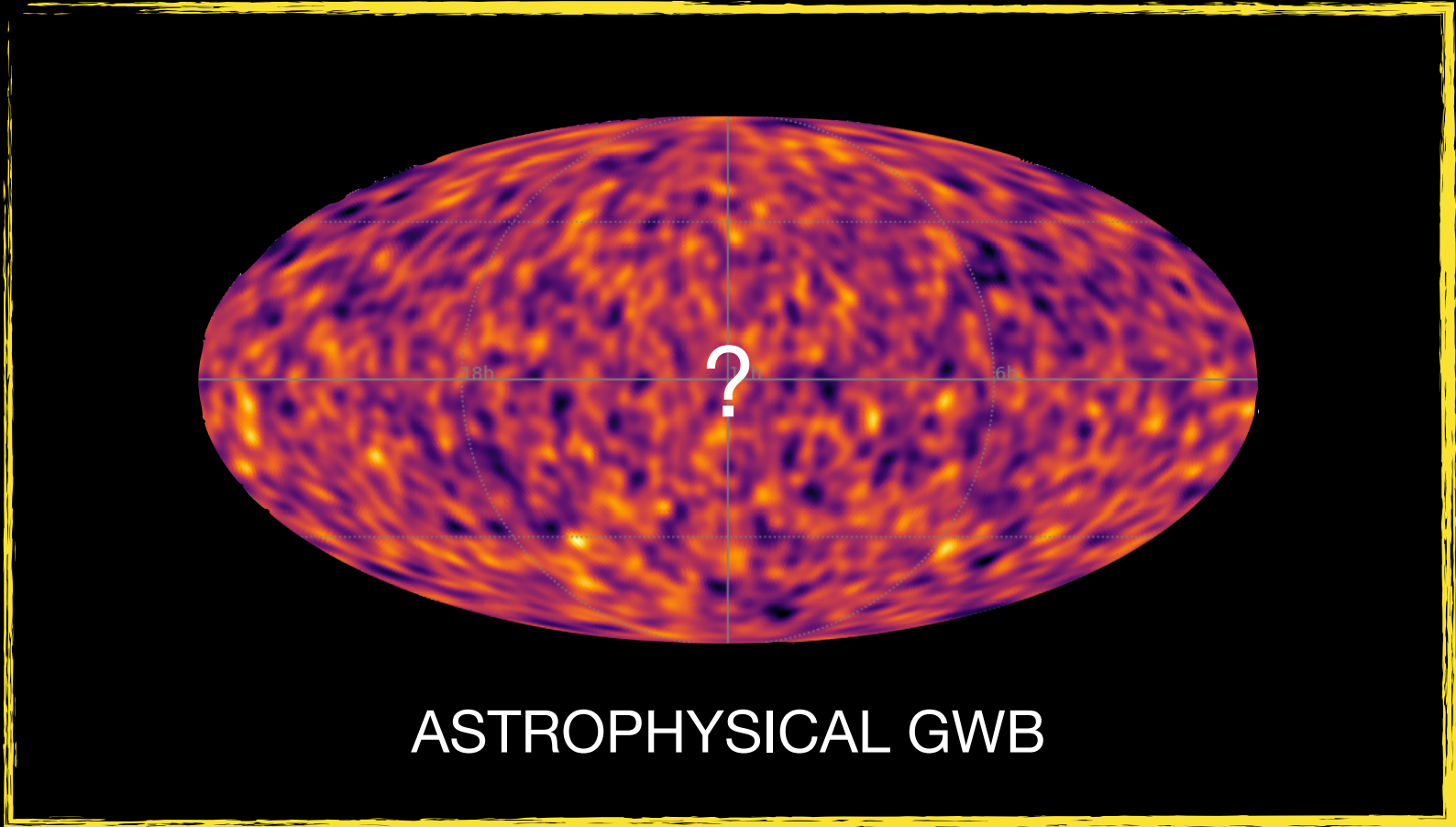
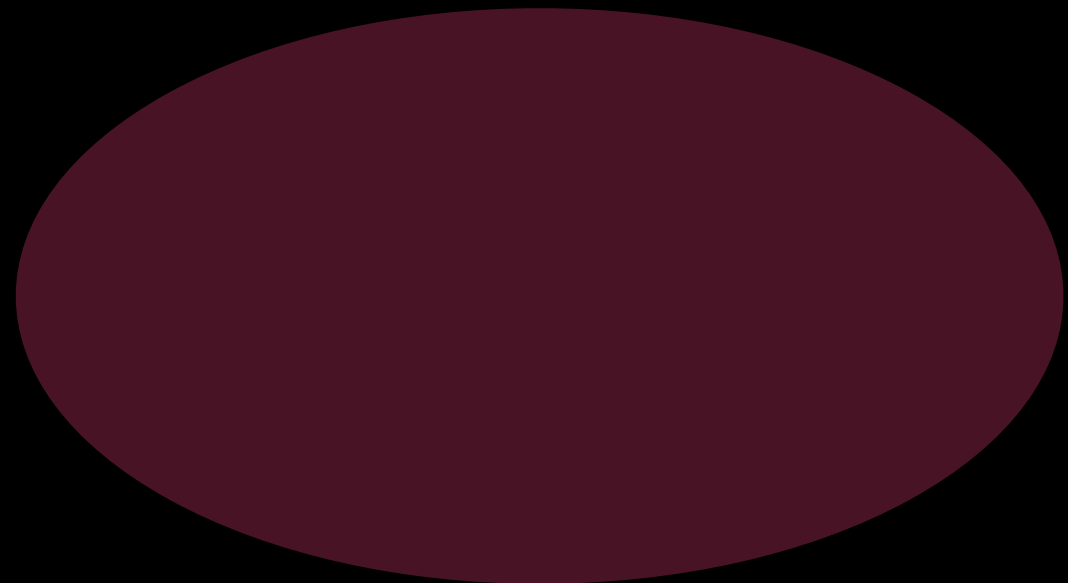
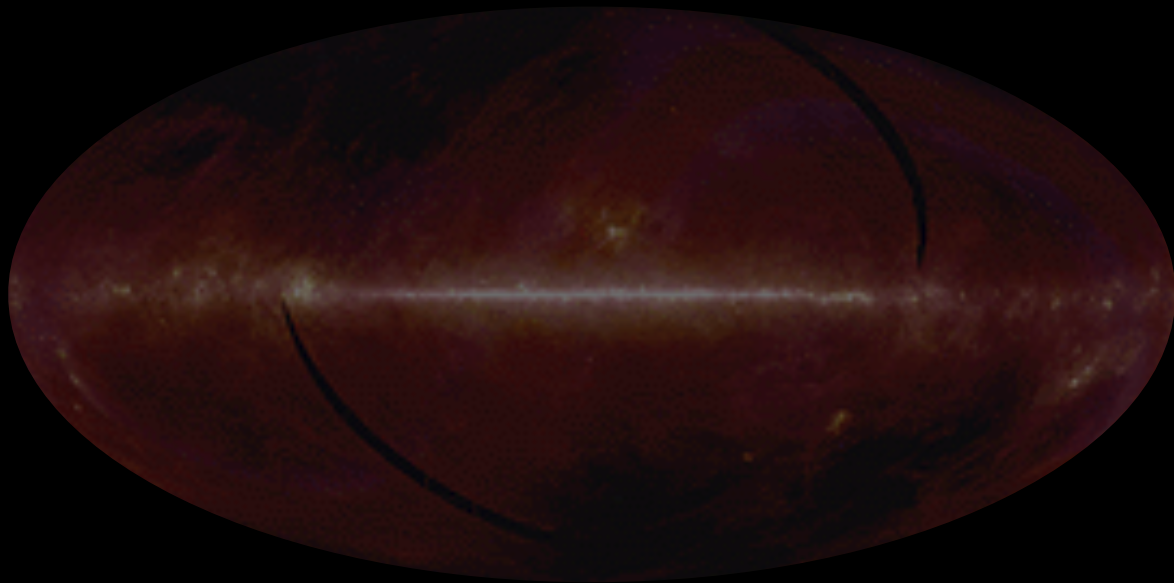
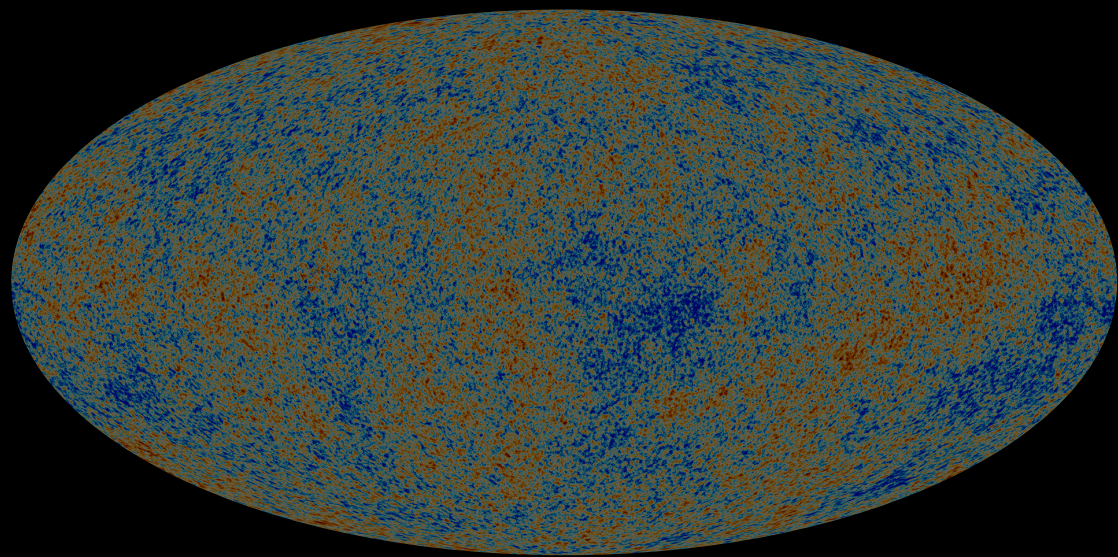
where $S_h(f) = \frac{3H_0^2}{2\pi^2} \frac{\Omega_{\text{gw}}(f)}{f^3}$

WHICH GWBs WE ARE SENSITIVE TO?

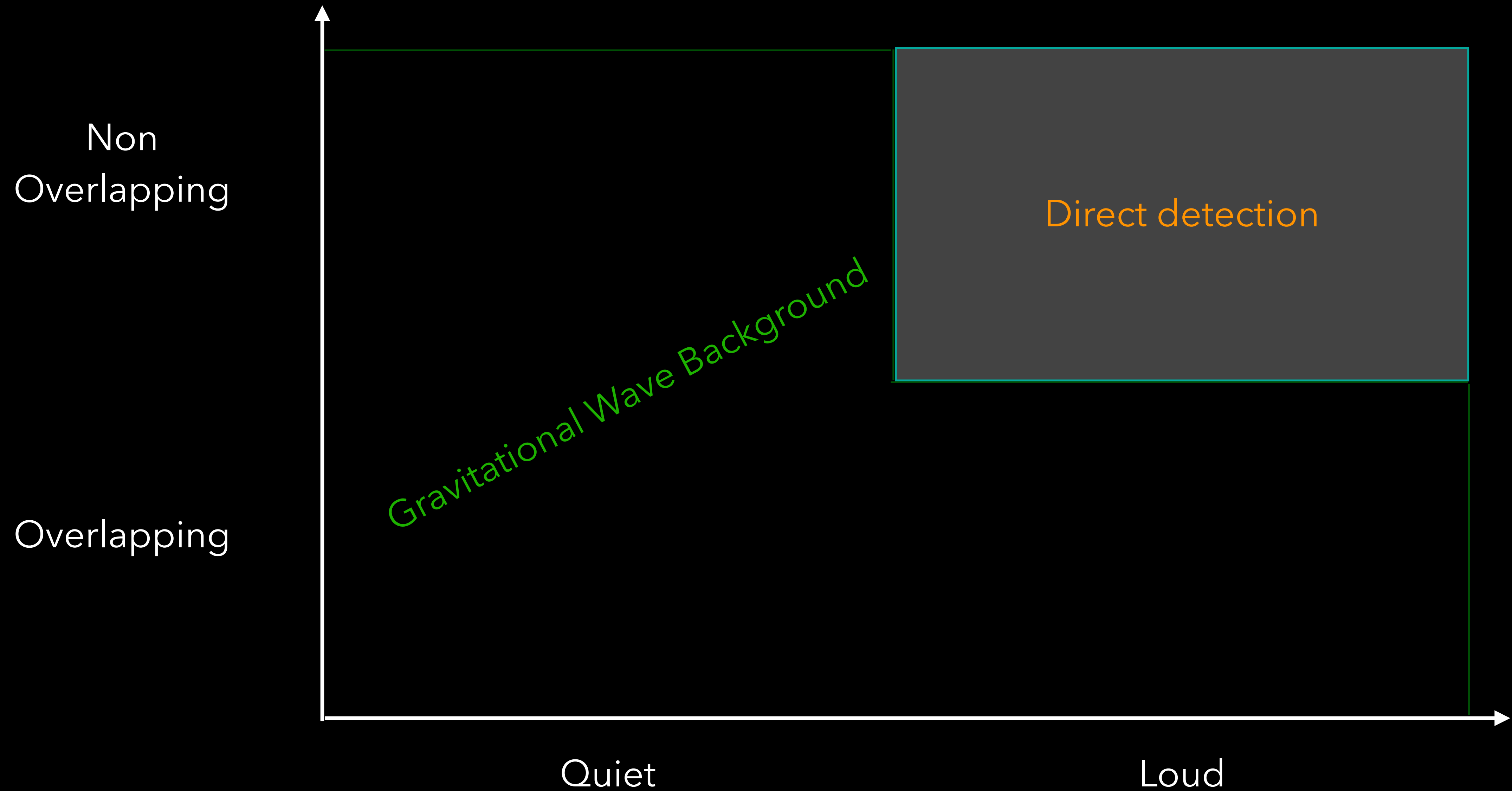


WHICH GWBs WE ARE SENSITIVE TO?

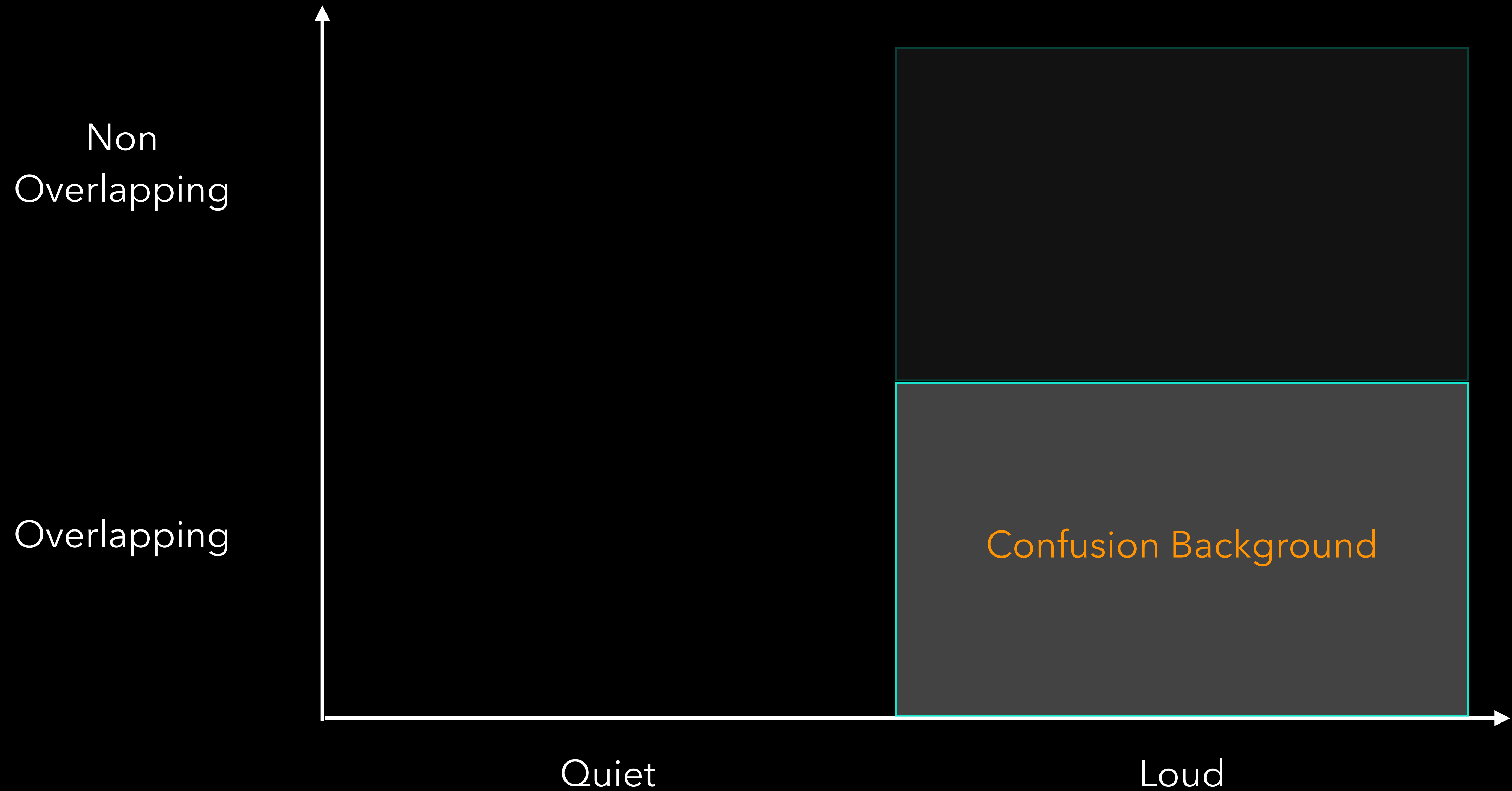
Astrophysical Origin



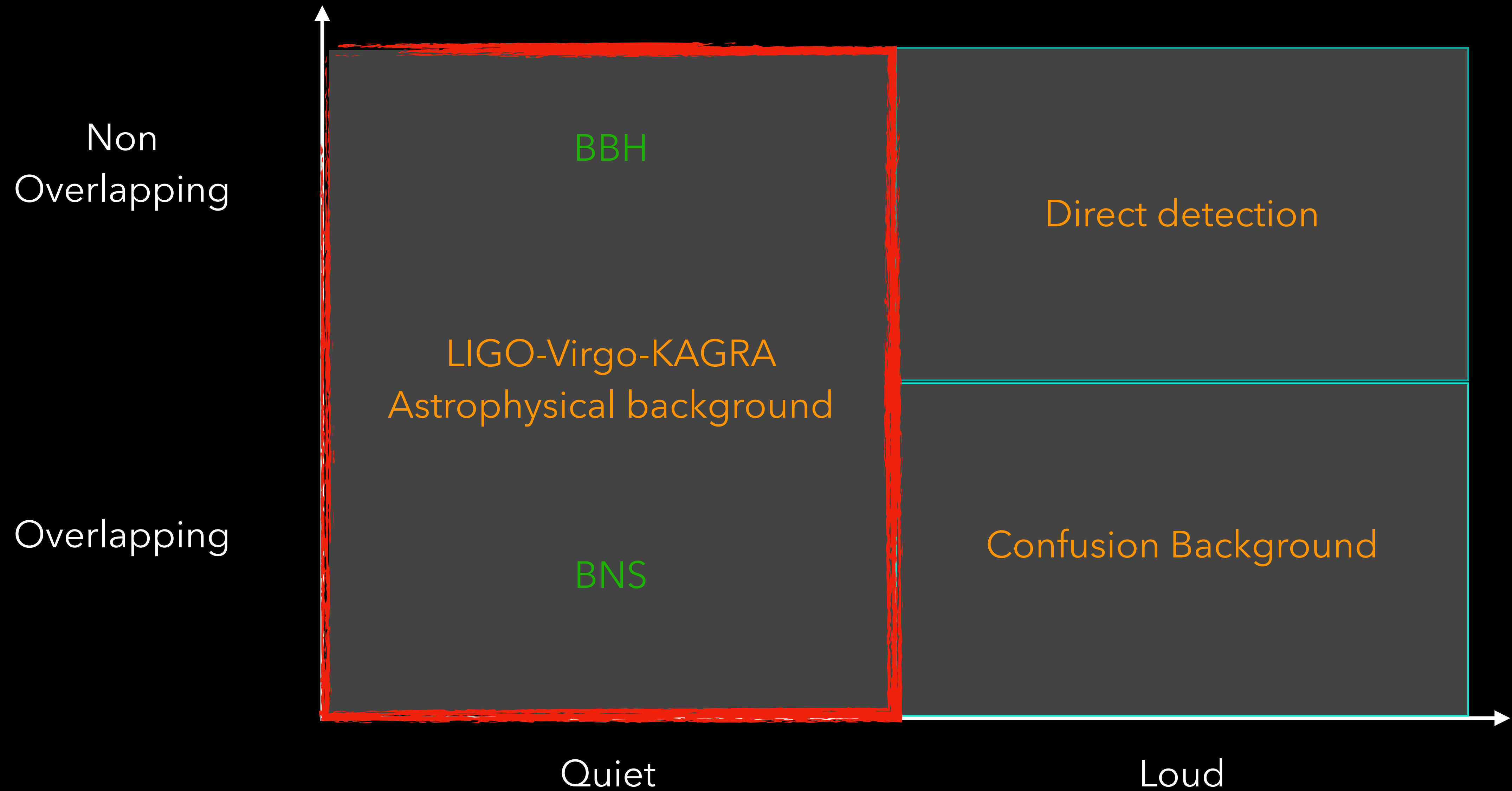
ASTROPHYSICAL GWB



ASTROPHYSICAL GWB



ASTROPHYSICAL GWB



WHAT DETECTION METHODS CAN WE USE?

The GWB signal looks more like noise in a single detector.

What can be done:

- Identify features that distinguish between the expected signal and noise.
- Measure our detector's noise sources well enough in amplitude and spectral shape.
- Detectors with uncorrelated noise: cross-correlation separates the signal from the noise.

WHAT DETECTION METHODS CAN WE USE?

The GWB signal looks more like noise in a single detector.

What can be done:

- Identify features that distinguish between the expected signal and noise.
- Measure our detector's noise sources well enough in amplitude and spectral shape.
- **Detectors with uncorrelated noise: cross-correlation separates the signal from the noise.**

Data from two detectors:

$$d_1 = h + n_1 \quad d_2 = h + n_2 \quad h \rightarrow \text{common GW signal component}$$

Expected value of cross-correlation:

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle + \langle n_1 n_2 \rangle + \cancel{\langle h n_2 \rangle} + \cancel{\langle n_1 h \rangle} = \langle h^2 \rangle + \langle n_1 n_2 \rangle$$

0 0

Assuming detector noise is uncorrelated*:

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle + \cancel{\langle n_1 n_2 \rangle}$$

0

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle \equiv S_h$$

Cross-correlation separates the signal from the noise

Intensity of the background

Cross correlation estimator:

$$\hat{C}^{IJ}(f) \propto \frac{2}{T} \frac{\Re[\tilde{s}_I^*(f) \tilde{s}_J(f)]}{\Gamma_{IJ}(f) S_0(f)}$$

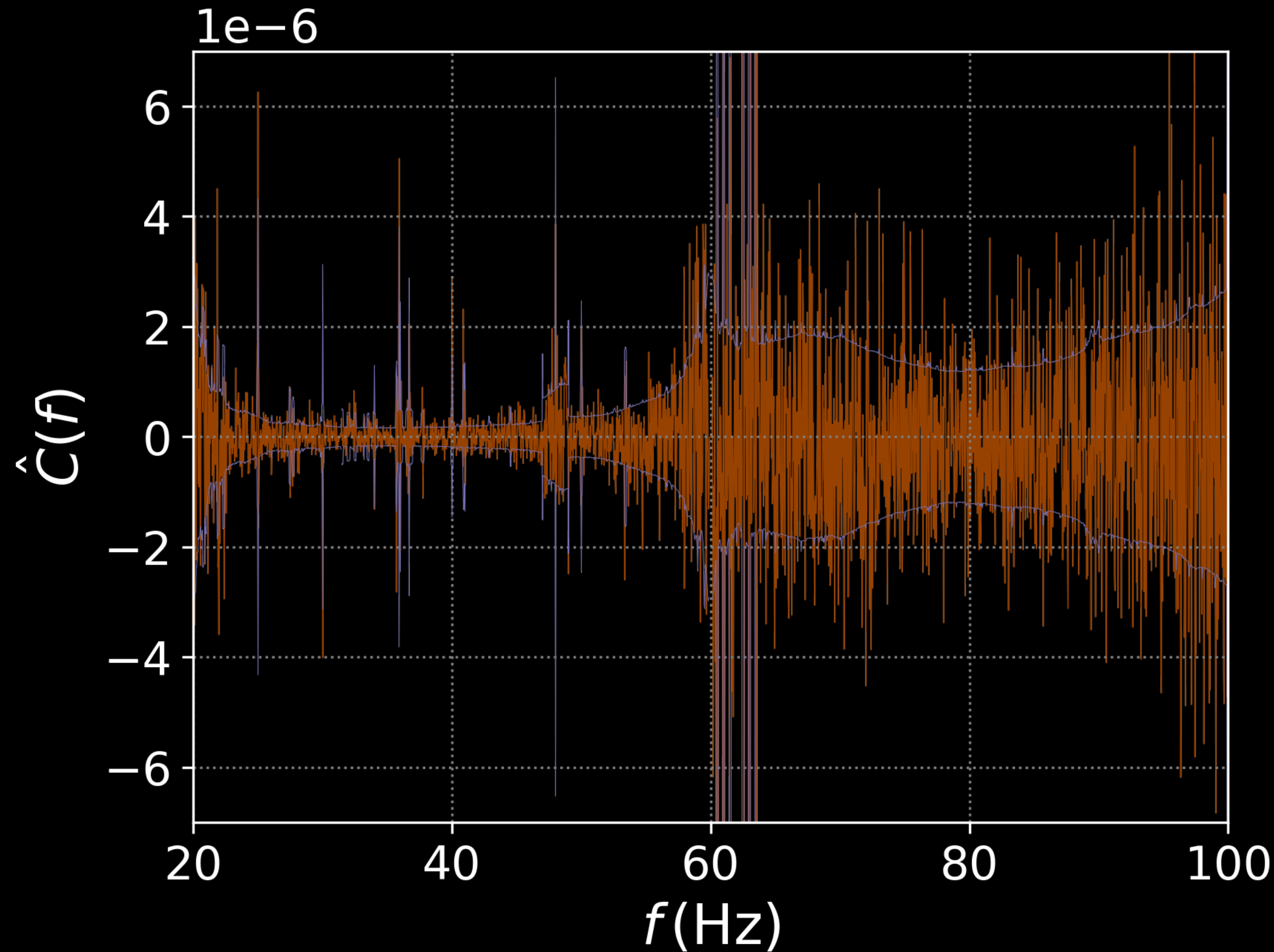
Variance:

$$\sigma_{IJ}^2(f) \approx \frac{1}{2T \Delta f} \frac{P_I(f) P_J(f)}{\Gamma_{IJ}^2(f) S_0^2(f)}$$

Signal to noise ratio:

$$\text{SNR} = \frac{3H_0^2 \sqrt{T}}{10\pi^2} \left(\int_{-\infty}^{\infty} df \frac{\Omega_{\text{GW}}^2(|f|) \Gamma_{12}^2(|f|)}{|f|^6 P_1(|f|) P_2(|f|)} \right)^{1/2}$$

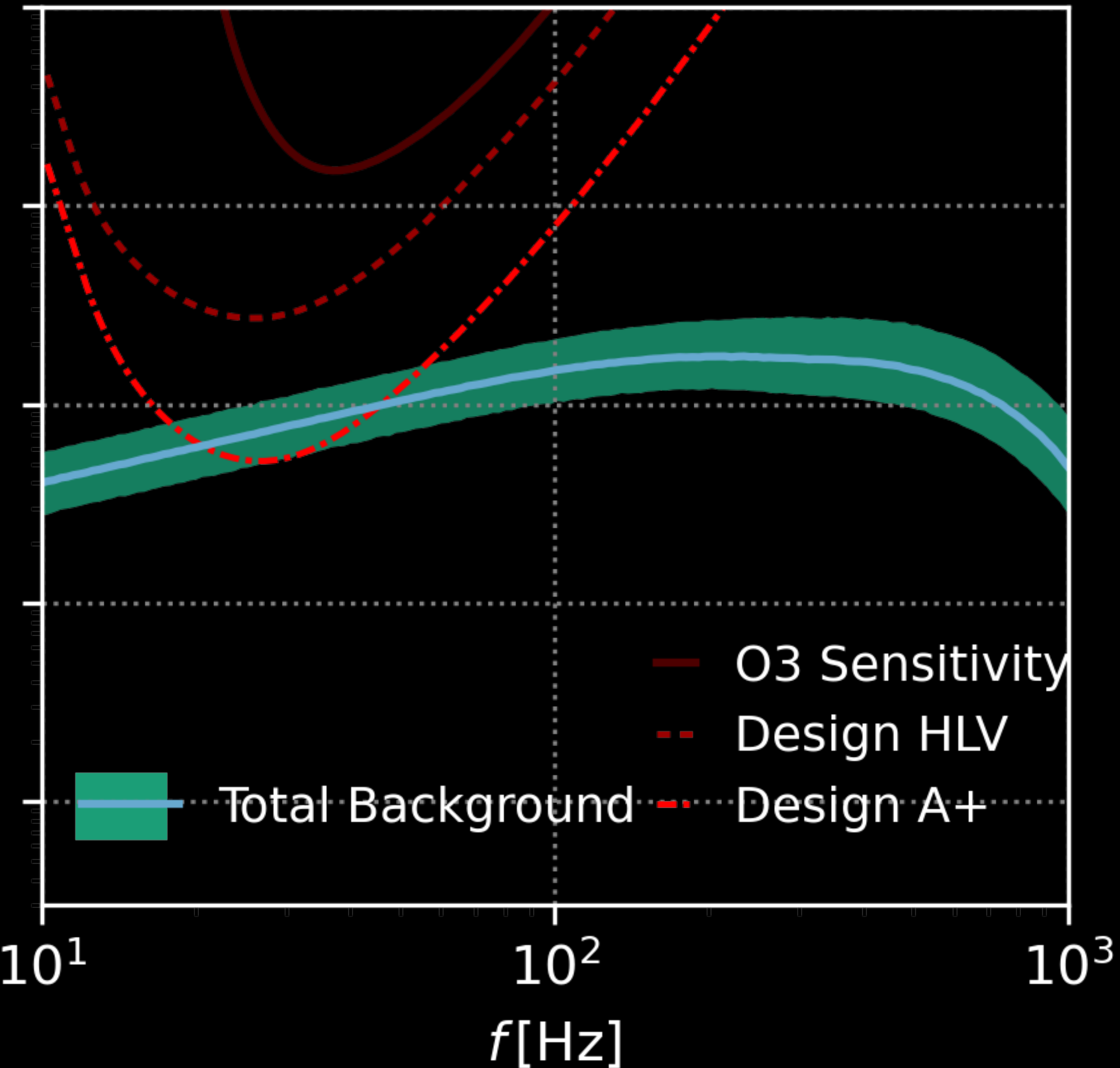
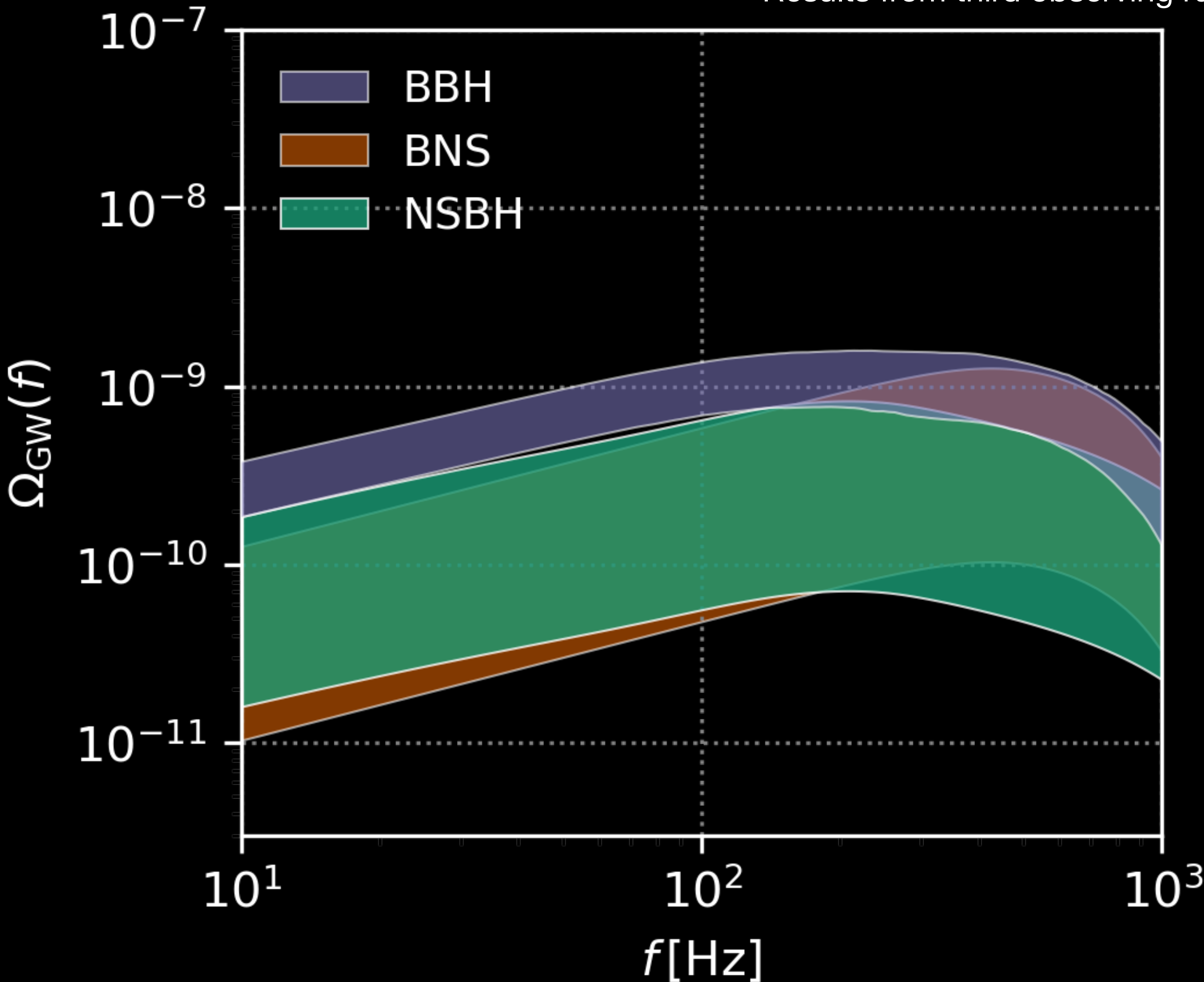
The observed cross-correlation spectra combining data from all three observing runs of LIGO-Virgo-KAGRA detectors: The spectrum is consistent with expectations from uncorrelated, Gaussian noise.



$$\Omega_{\text{CBC}}(f) \propto \int_0^\infty dz \, R(z) \frac{1}{(1+z)E(z)} f_s \left(\frac{dE_{\text{gw}}}{df_s} \right)$$

$$\Omega_{\text{GW}}(25\text{Hz}) = 6.9 \times 10^{-10}$$

Results from third observing run of LIGO-Virgo-KAGRA collaboration.



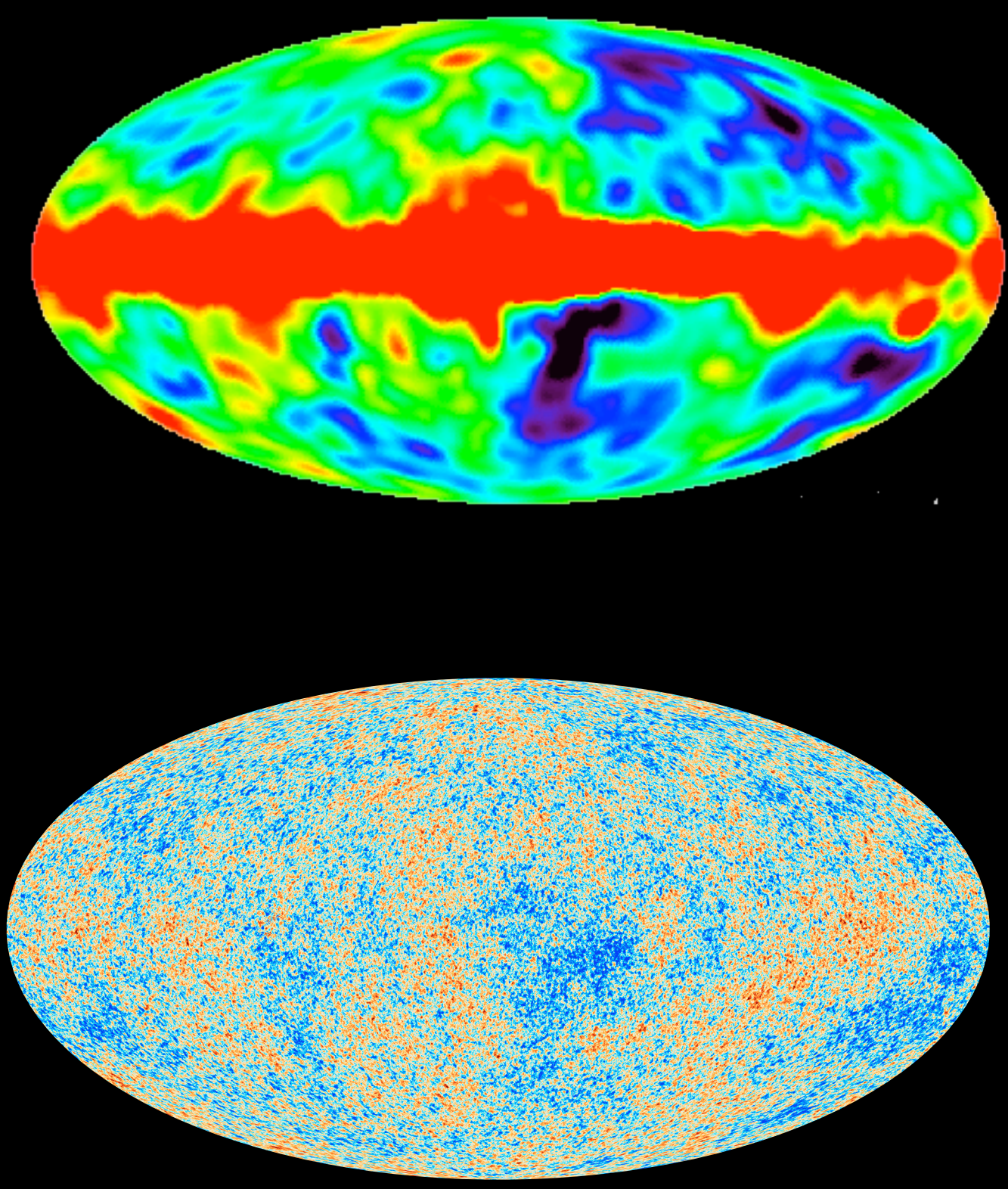
WE ARE NOW AT:

— 1965: Penzias & Wilson

— 1992: COBE

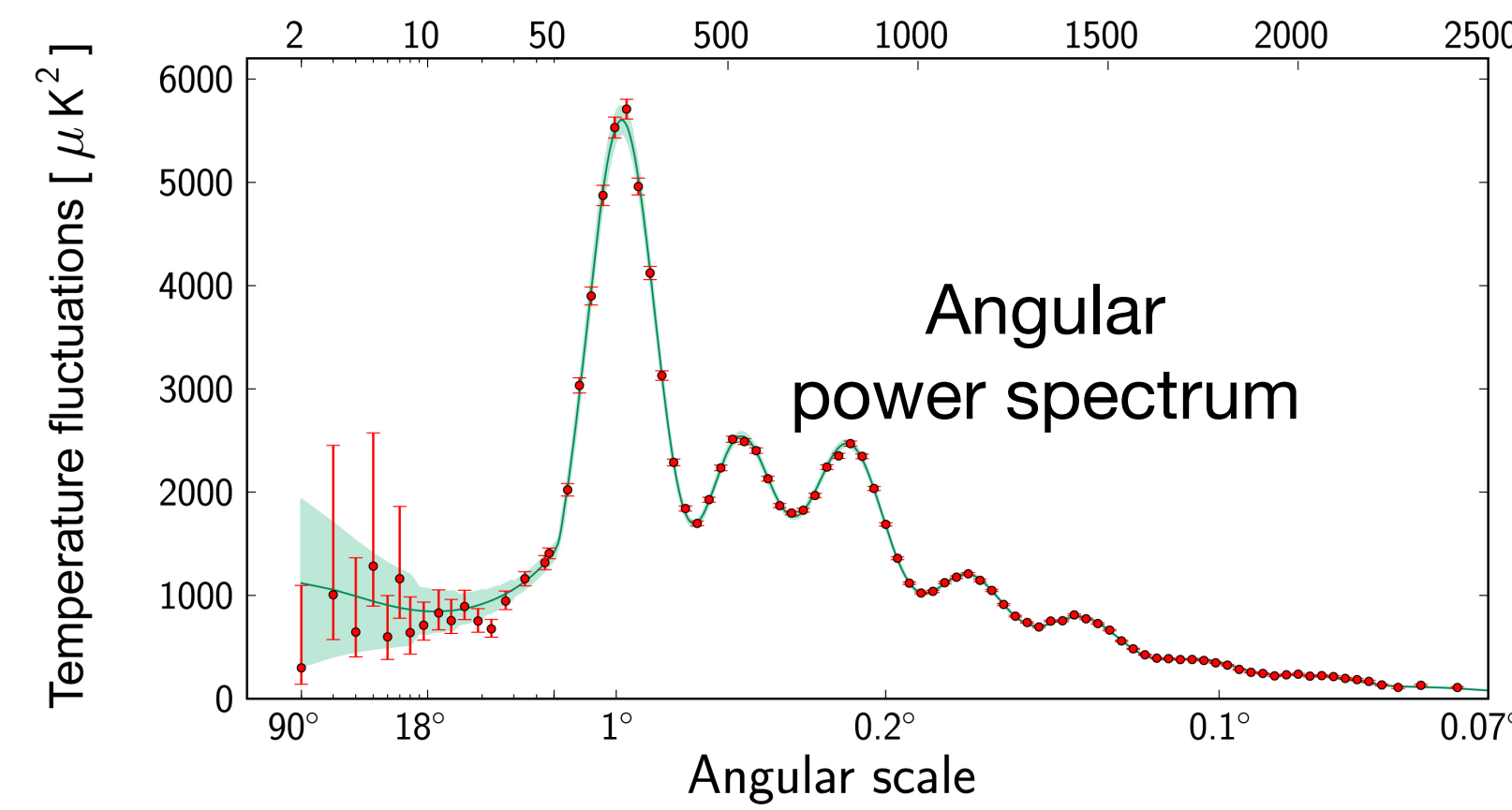
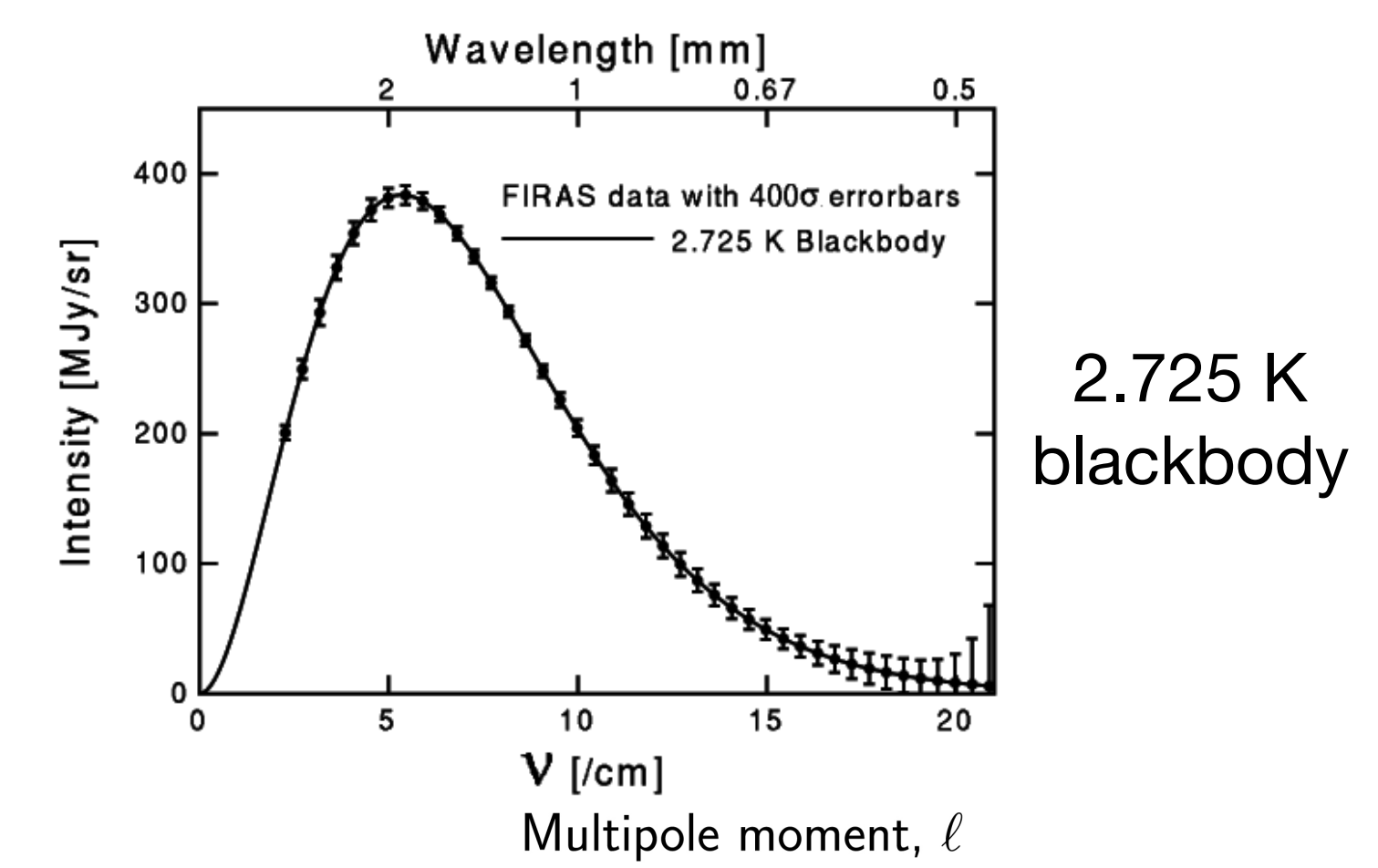
— 2013: WMAP, Planck

— 2025: **We are yet to detect the isotropic component of the GWB!**

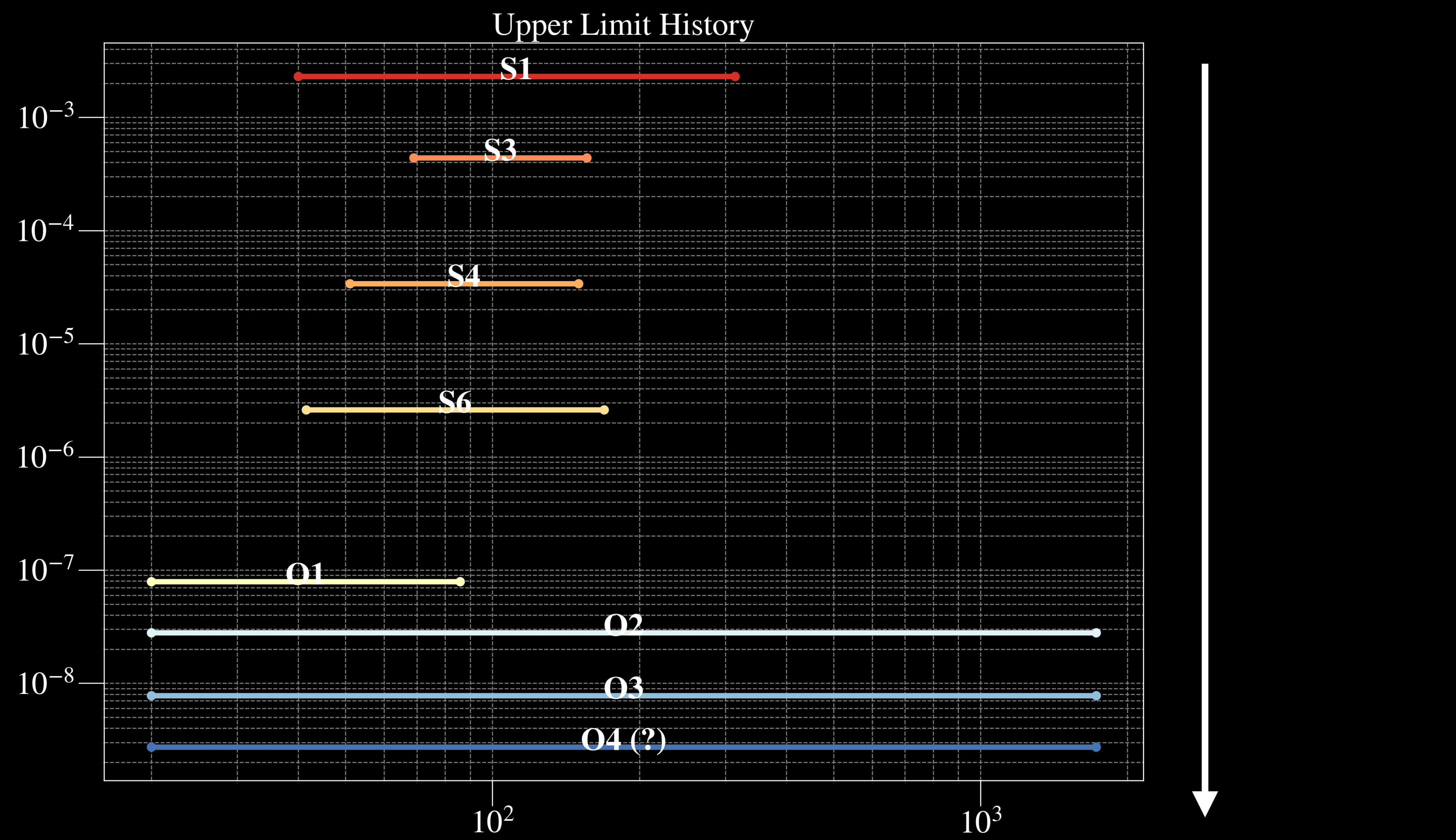


A MEASUREMENT OF EXCESS ANTENNA TEMPERATURE
AT 4080 Mc/s

Measurements of the effective zenith noise temperature of the 20-foot horn-reflector antenna (Crawford, Hogg, and Hunt 1961) at the Crawford Hill Laboratory, Holmdel, New Jersey, at 4080 Mc/s have yielded a value about 3.5° K higher than expected. This excess temperature is, within the limits of our observations, isotropic, unpolarized, and free from seasonal variations (July, 1964–April, 1965). A possible explanation for the observed excess noise temperature is the one given by Dicke, Peebles, Roll, and Wilkinson (1965) in a companion letter in this issue.

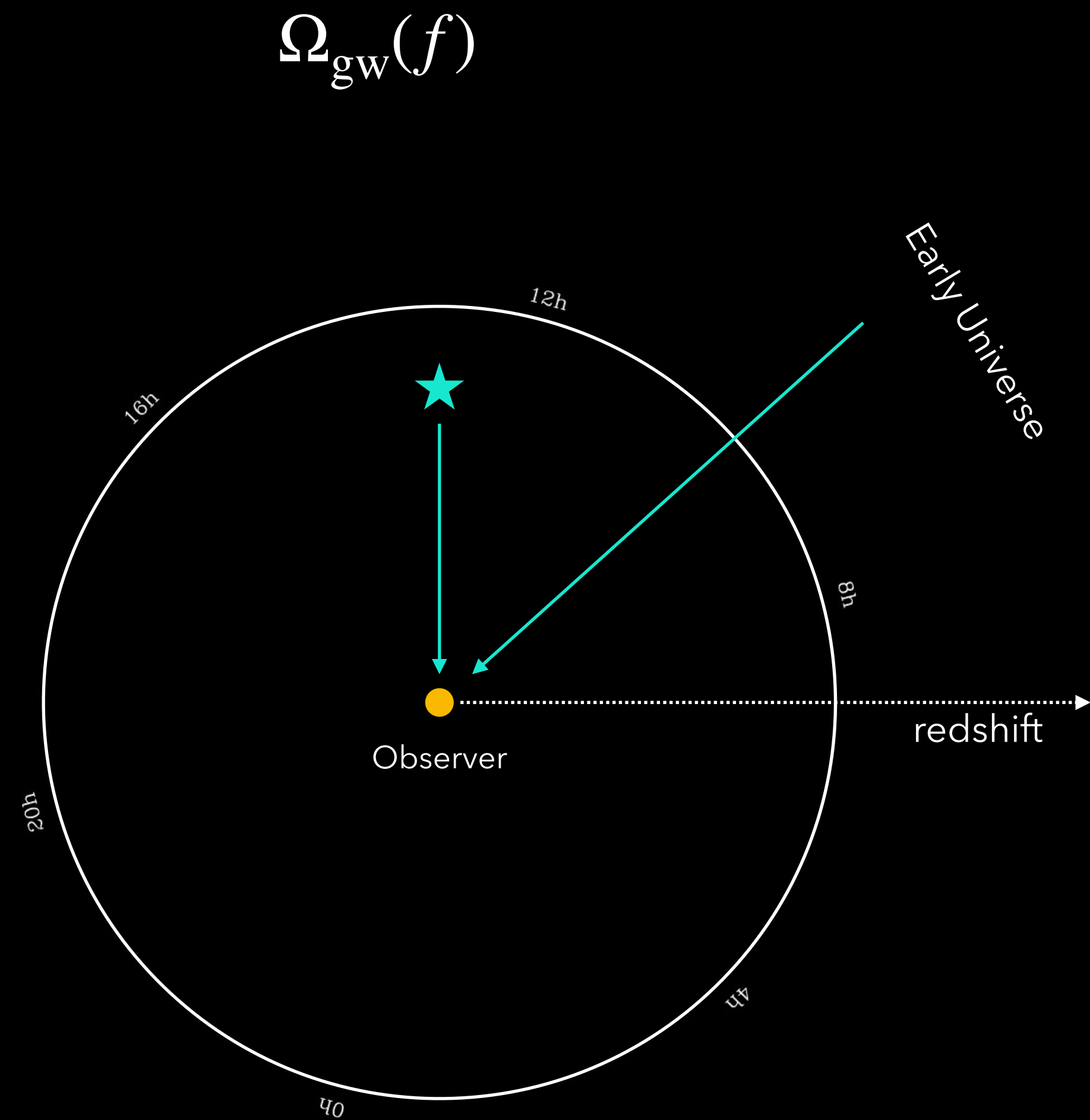


WE ARE NOW AT:



We are reaching there...

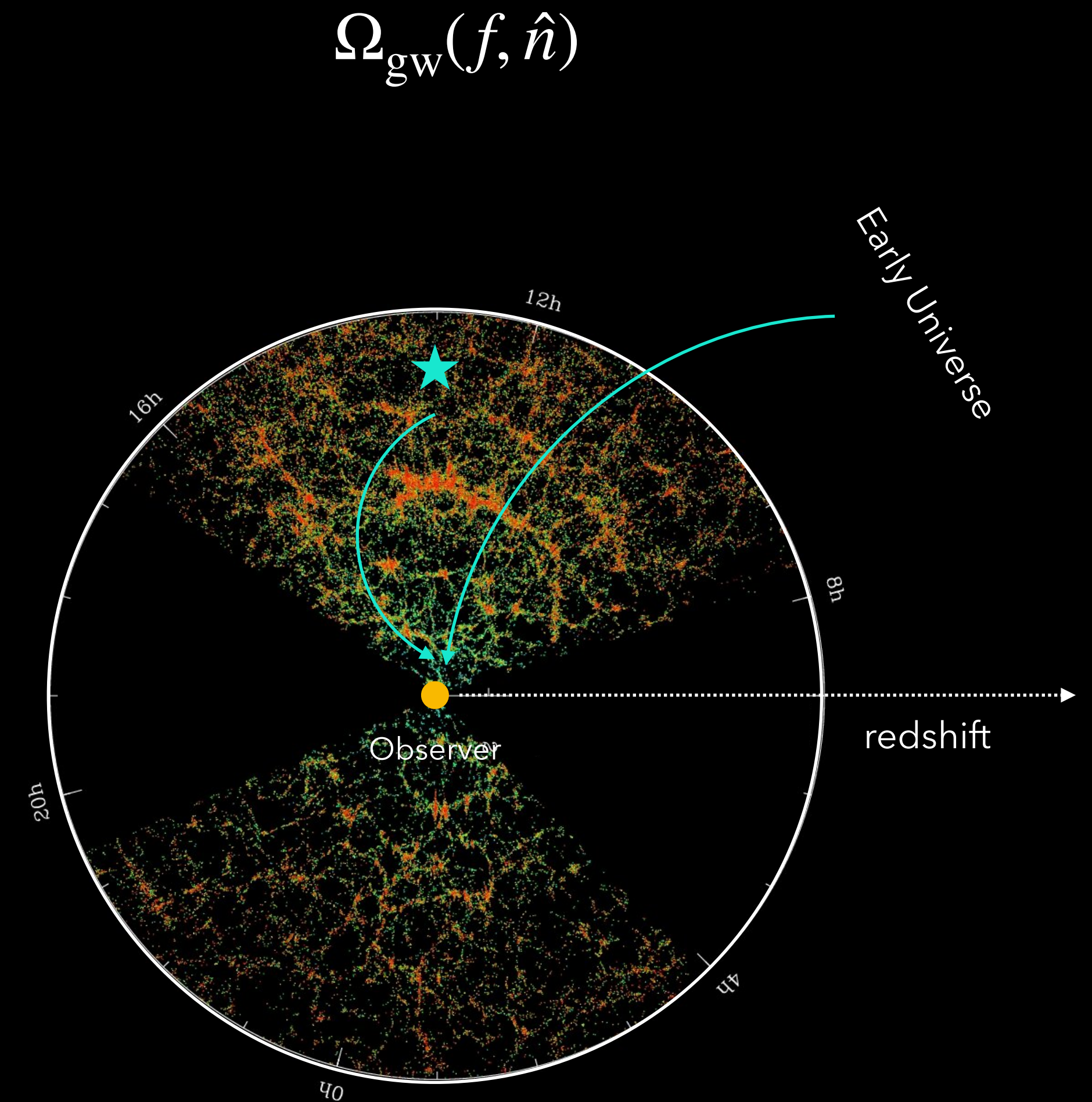
ISOTROPIC GWB



- sources isotropically distributed
- propagation along straight line
- no anisotropy in received flux

ANISOTROPIC GWB

- Anisotropic distribution of the emitting sources.
- Due to propagation: as gravitational-wave propagate, they accumulate line-of-sight effects, crossing different matter density fields which are inhomogeneously distributed in the Universe.

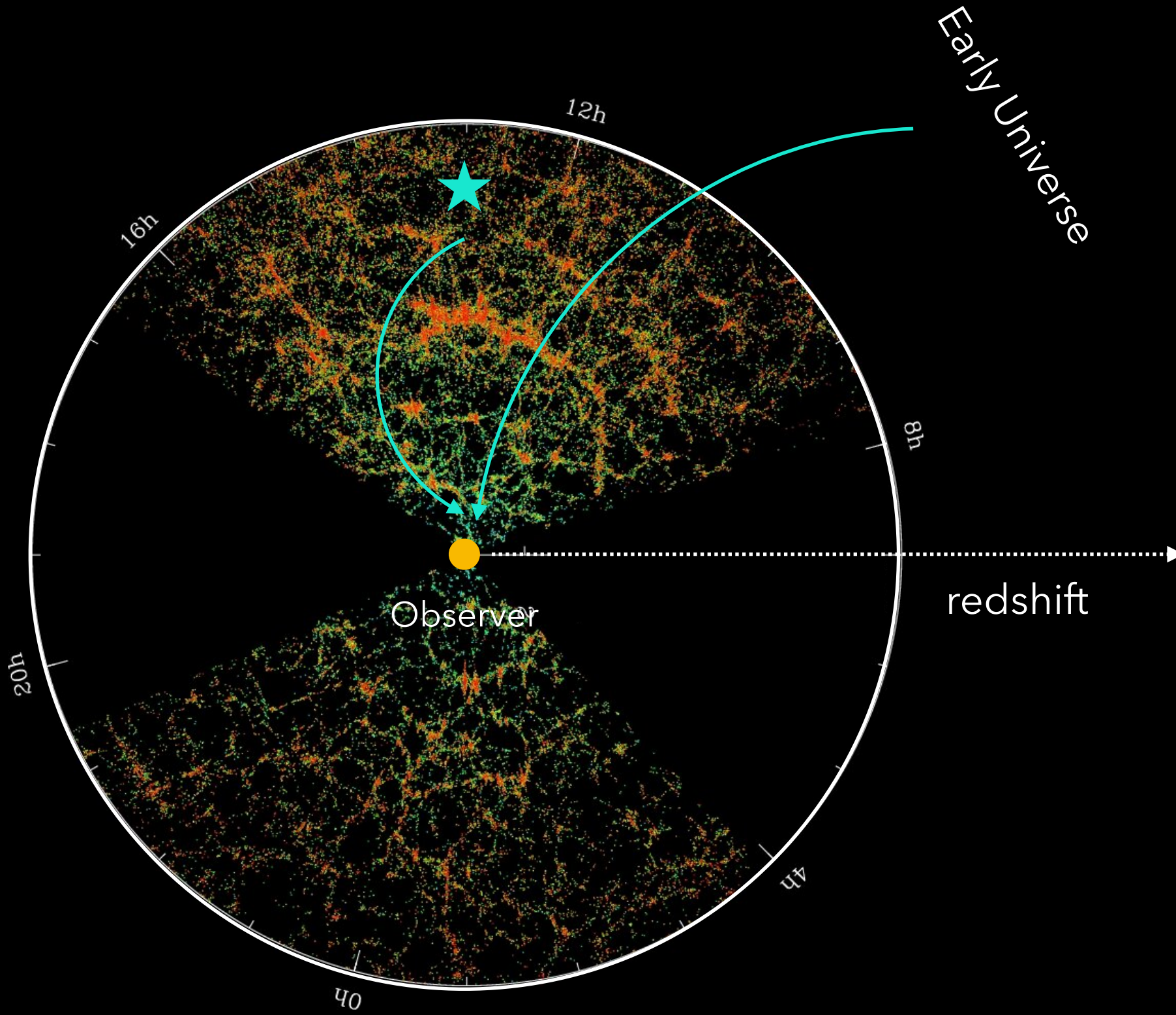
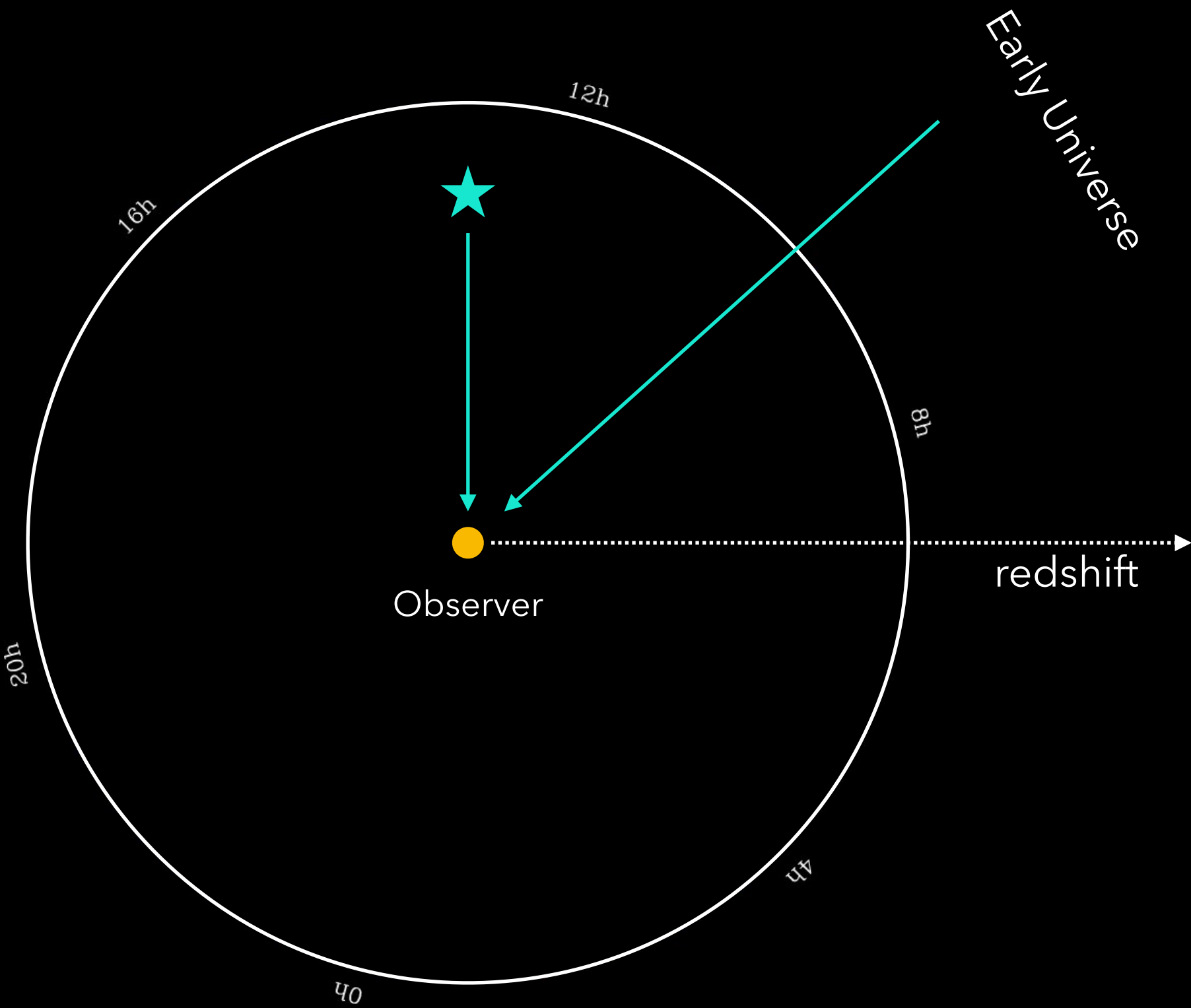


ANISOTROPIC GWB

$$\Omega_{\text{gw}}(f)$$

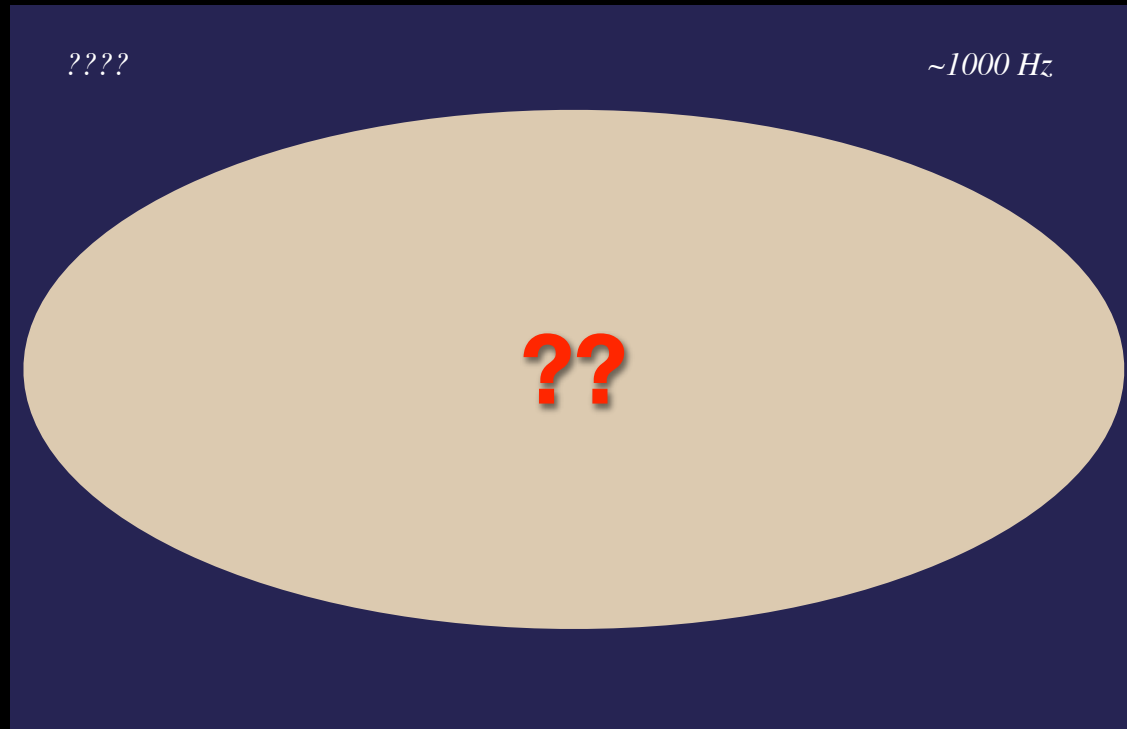
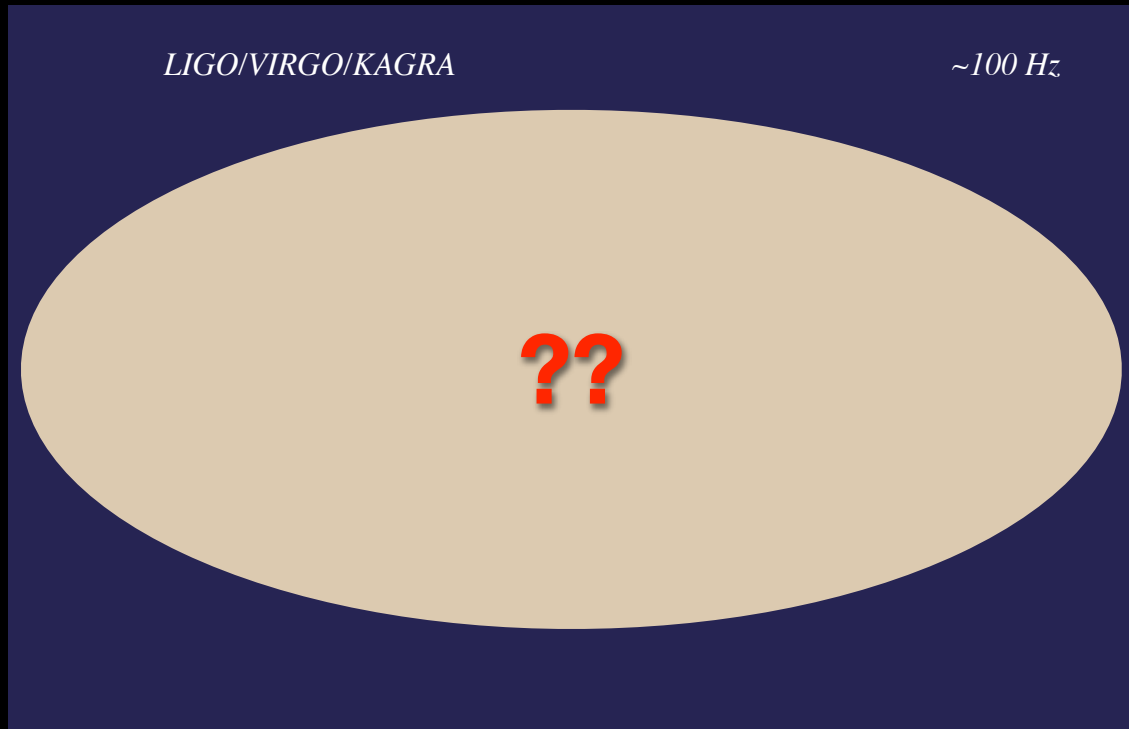
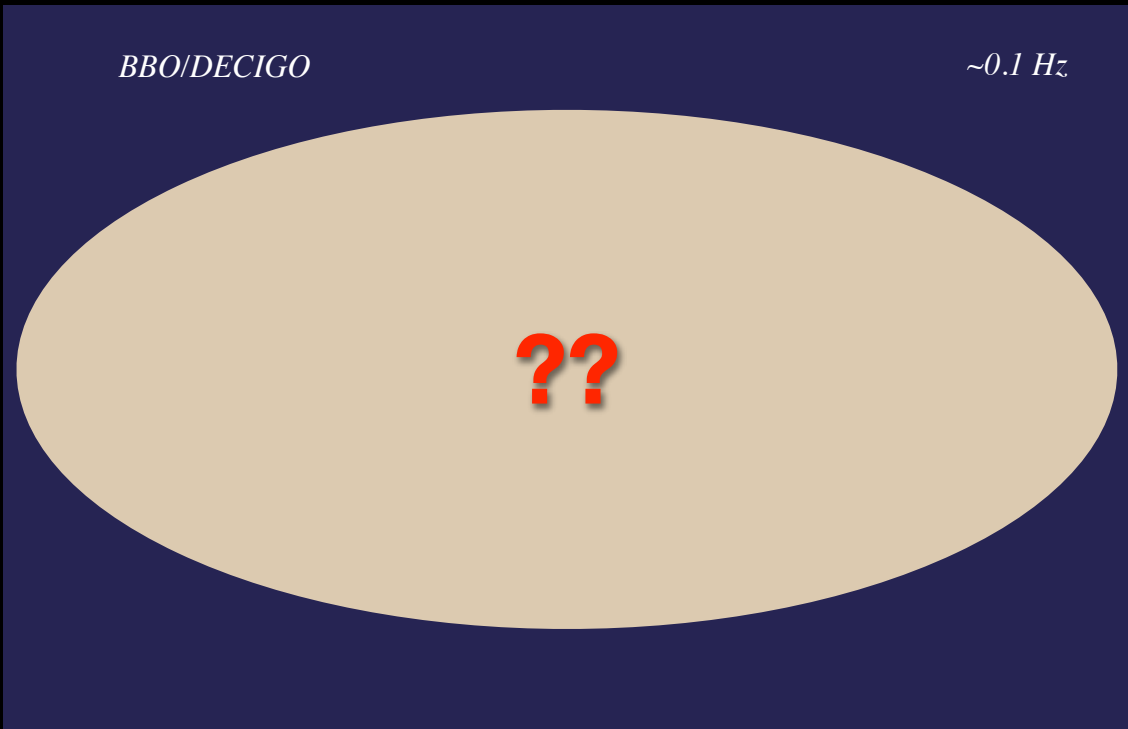
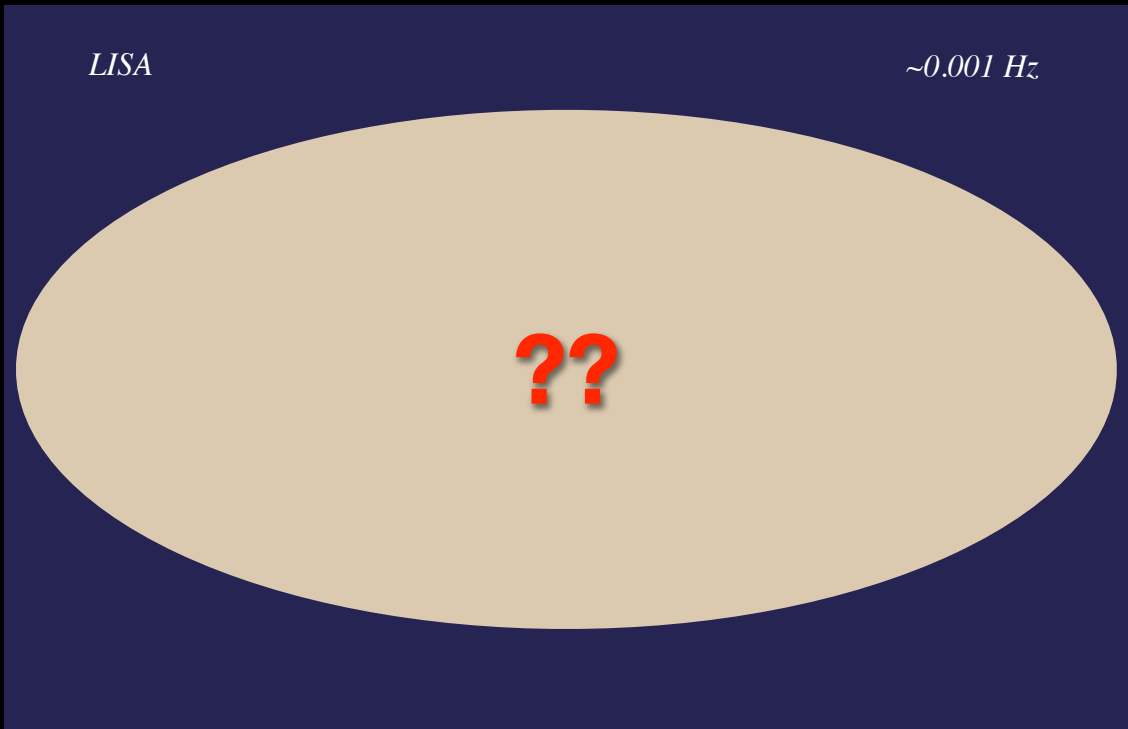
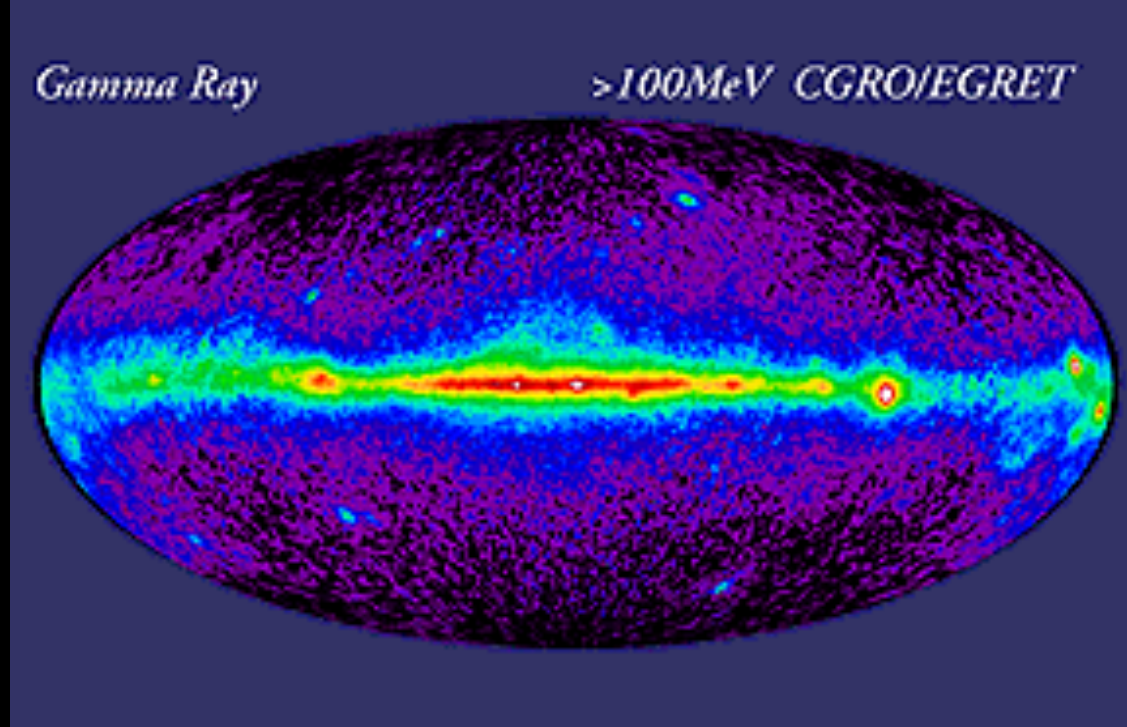
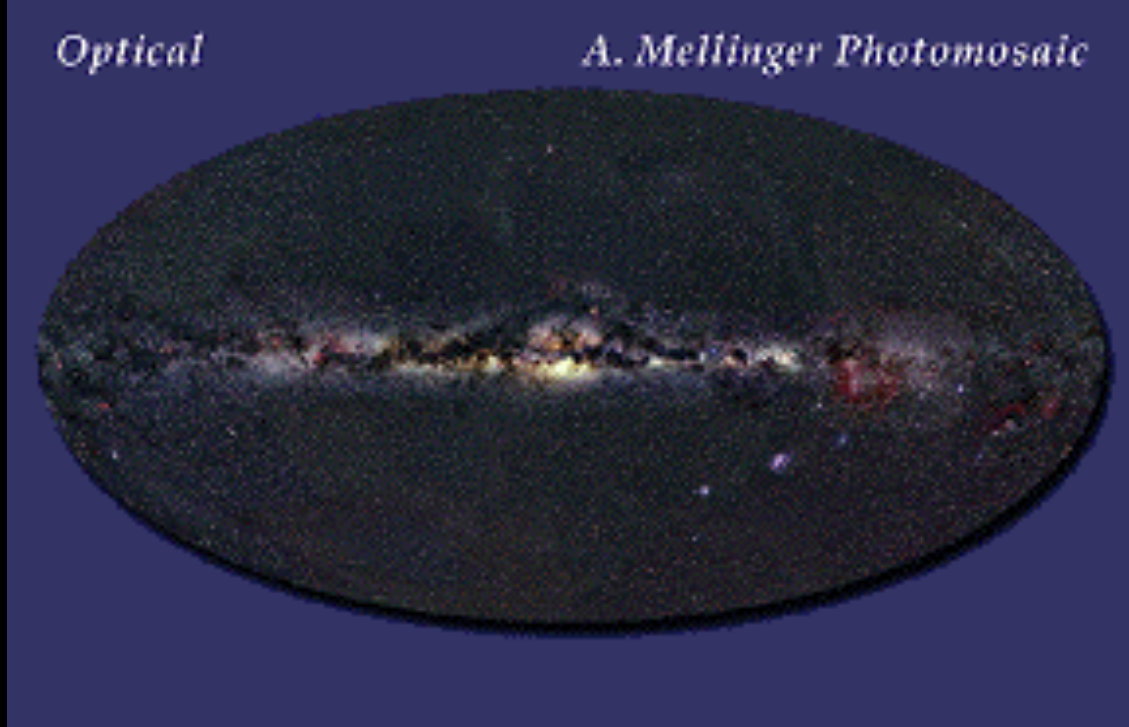
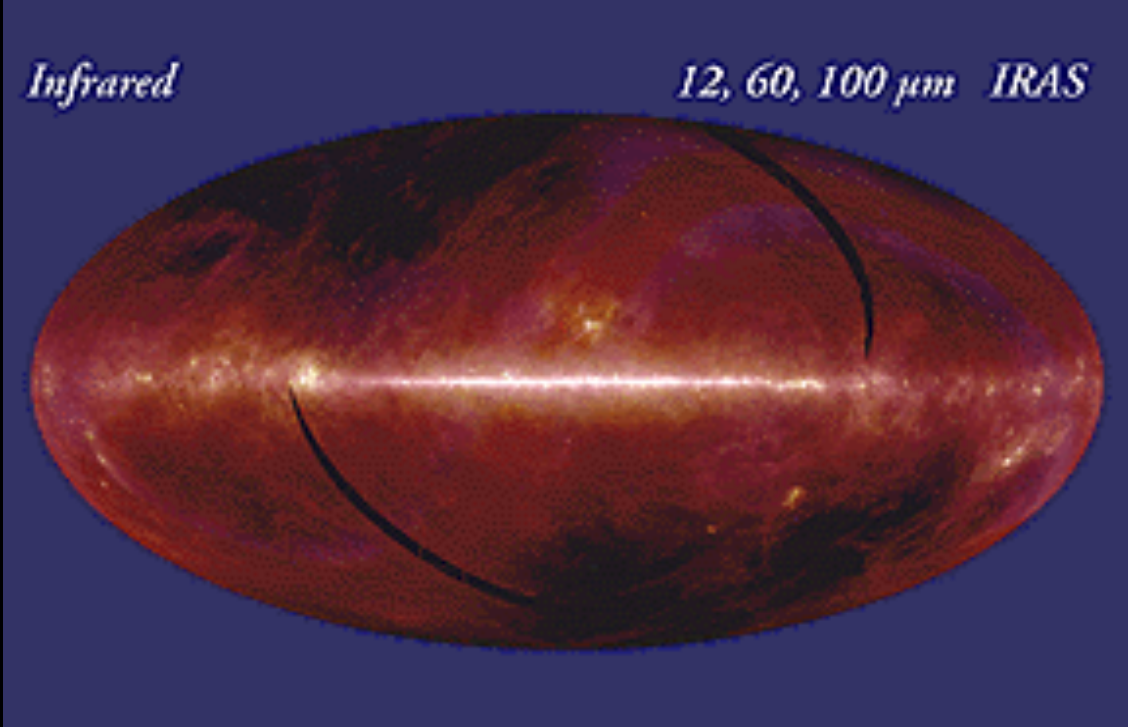
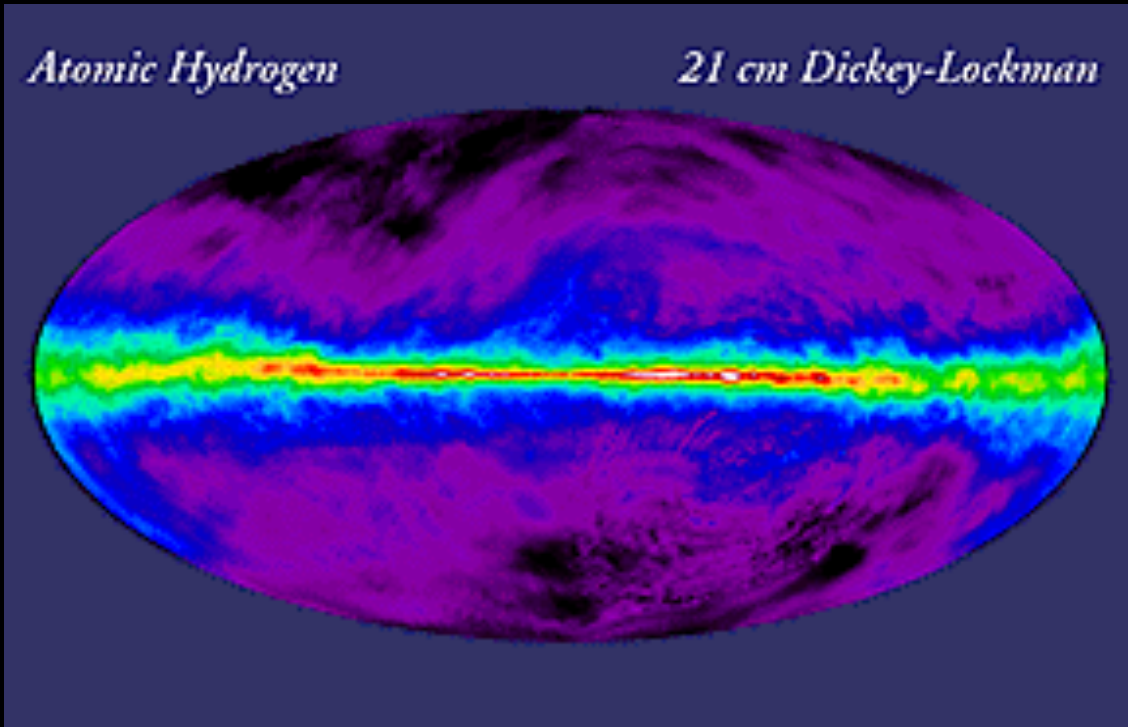
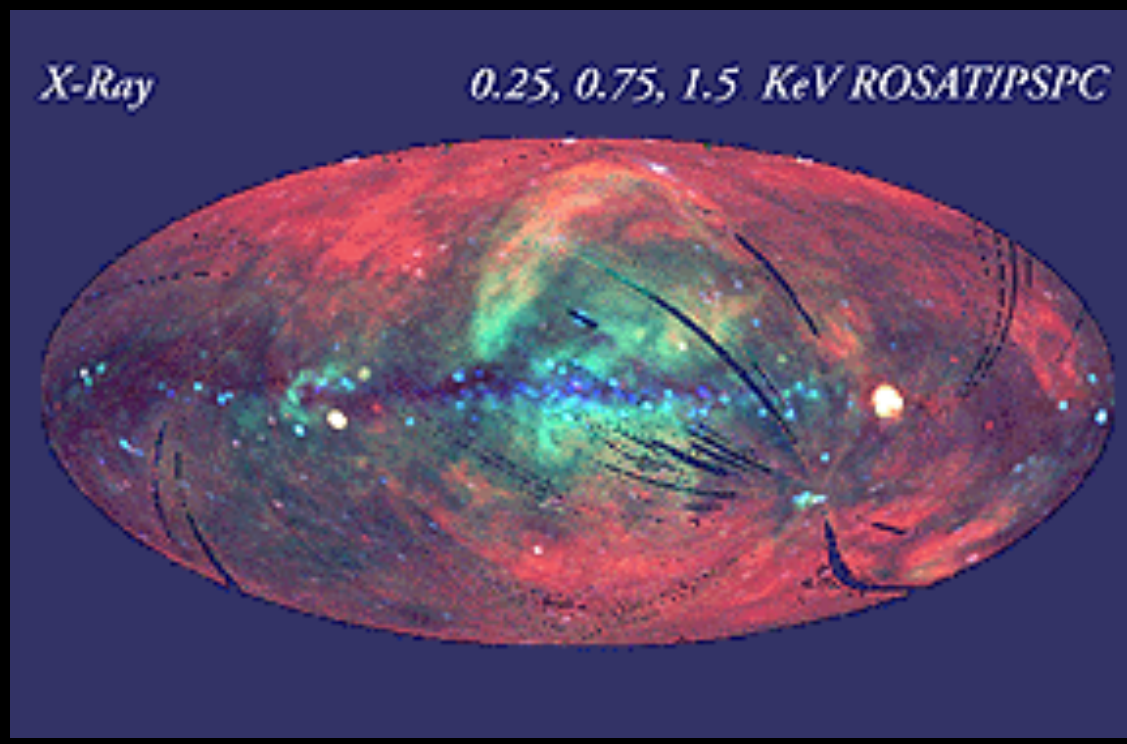
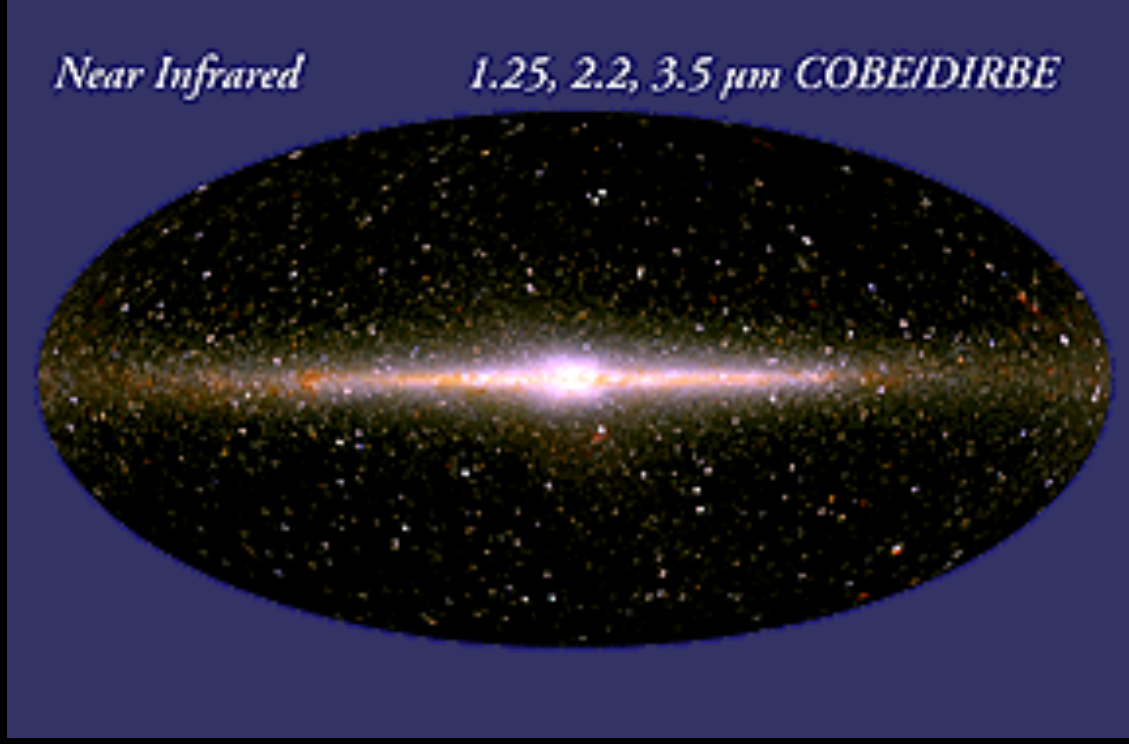
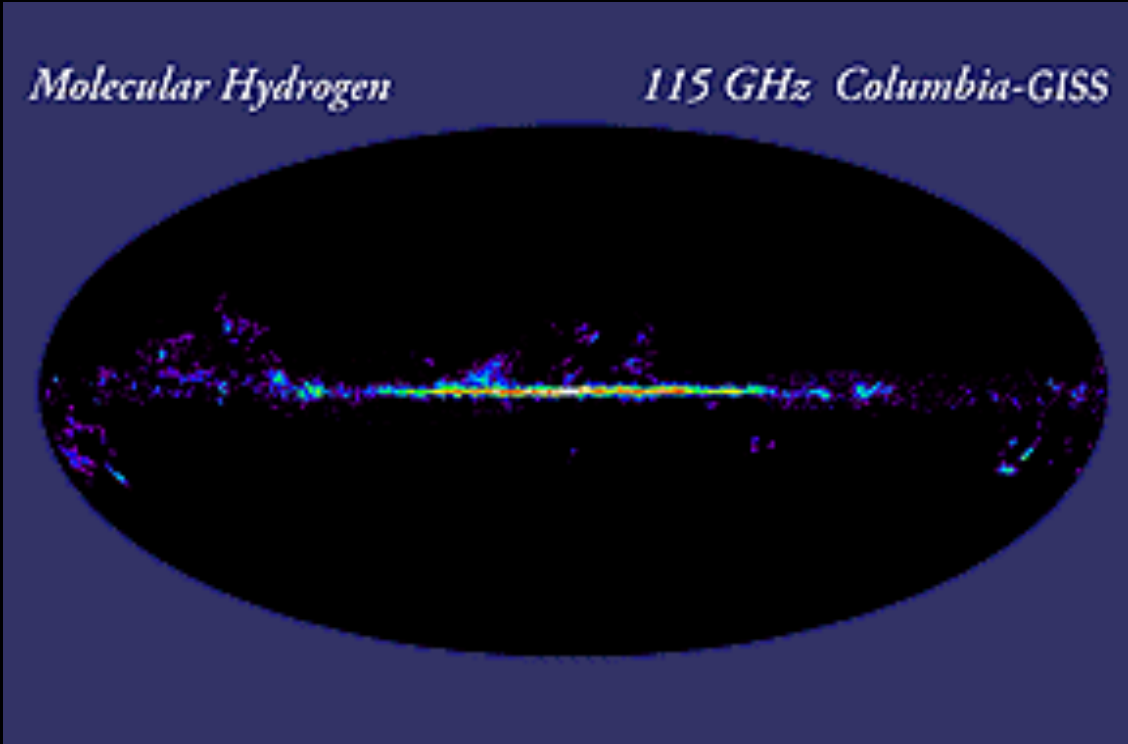
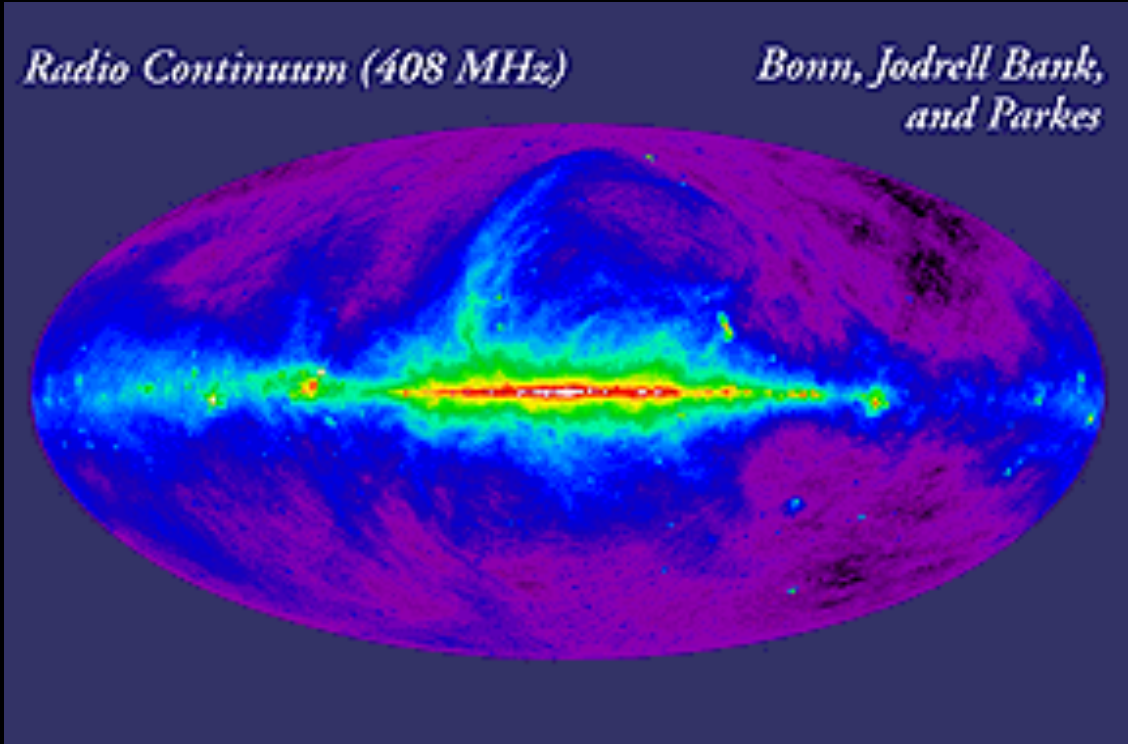


$$\Omega_{\text{gw}}(f, \hat{n})$$



HOW DO WE MAP THE GWB SKY?

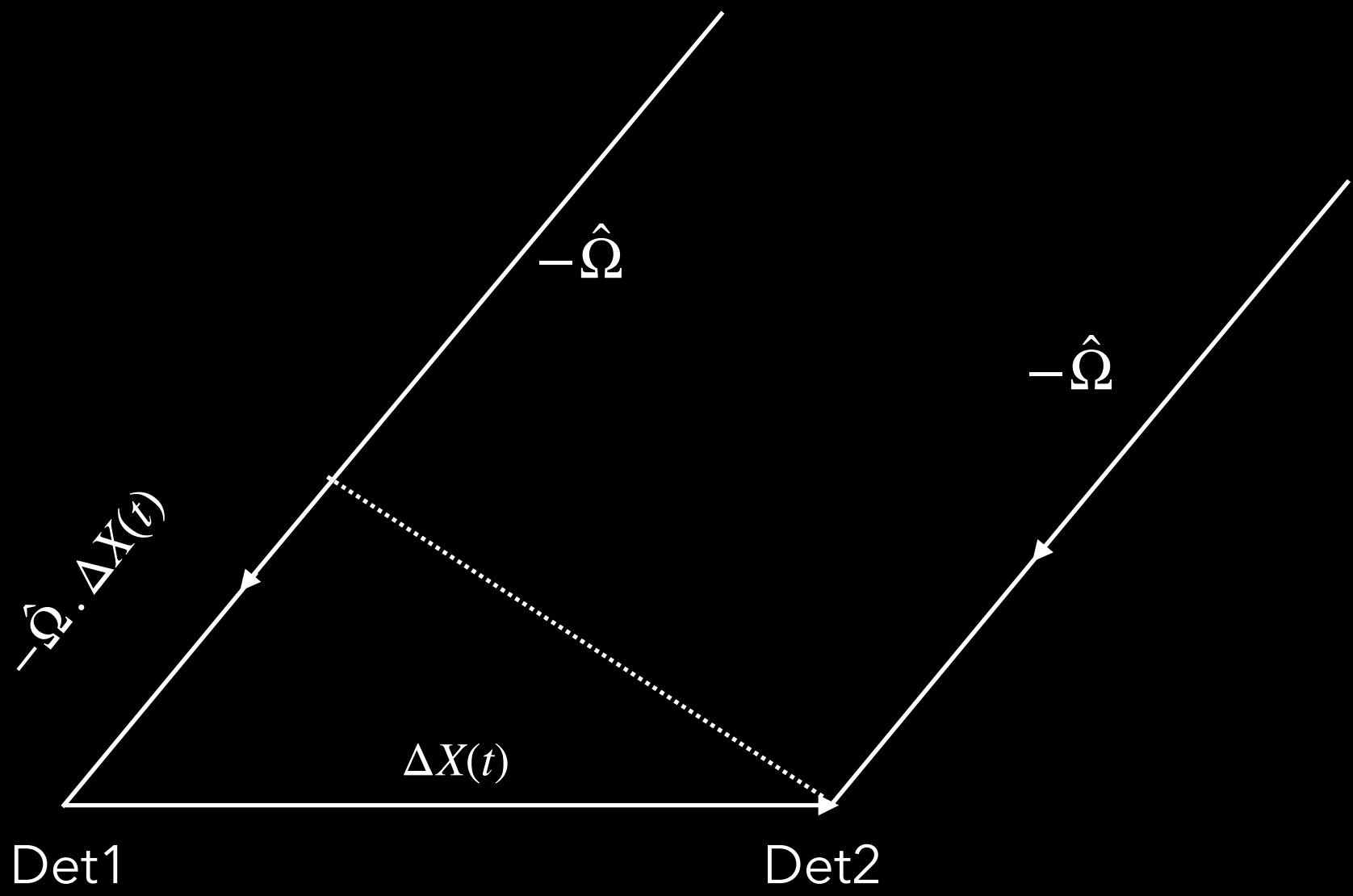
http://mwmw.gsfc.nasa.gov/mmw_allsky.html



HOW DO WE MAP THE GWB SKY?

**Cross-correlation is essentially
a one-dimensional map of the sky**

How do we extract the anisotropic GWB signal?
How to solve for its directionality?
How do we effectively map this signal on the sky?



- The time delay between two detectors
- Rotation of the earth.

GWB energy density $\Omega_{\text{gw}}(f, \hat{n}) \equiv \frac{f}{\rho_c} \frac{d\rho_{\text{GW}}}{df} = \frac{2\pi^2}{3H_0^2} f^3 \mathcal{P}(f, \hat{n})$

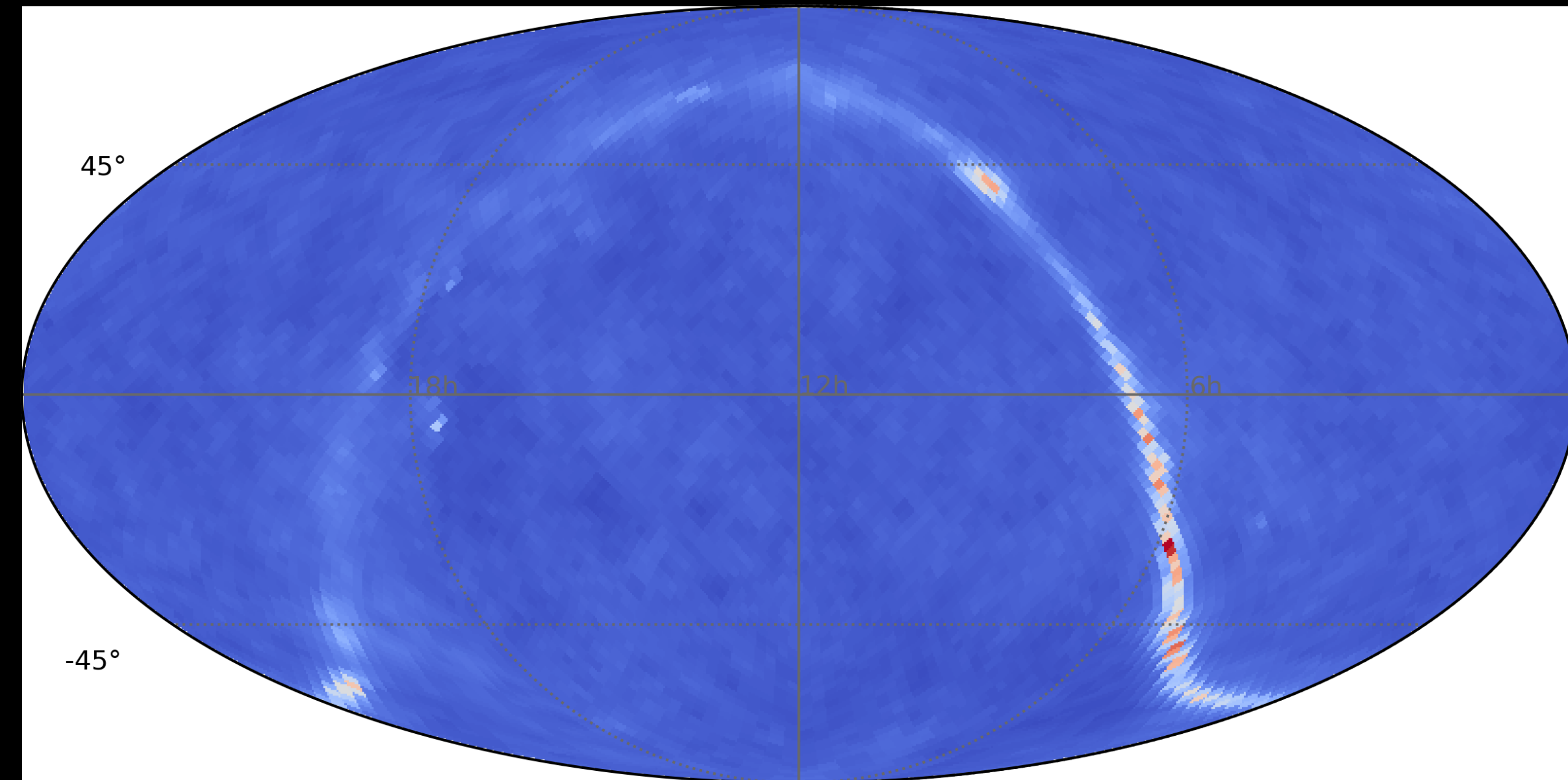
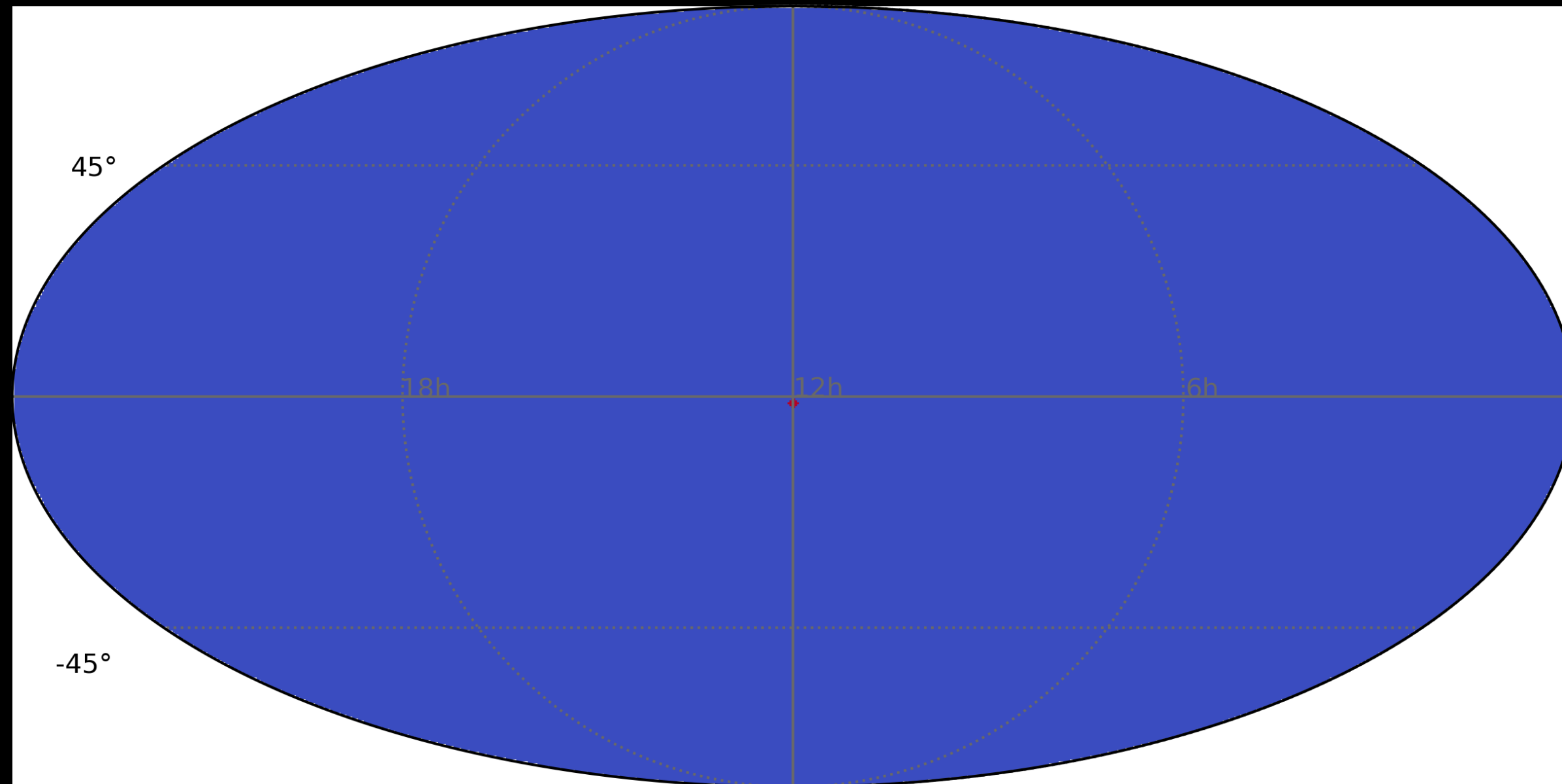
Cross-correlation is essentially a one-dimensional map of the sky.

Anisotropy can be expanded in pixel or spherical harmonic basis

$$\mathcal{P}(f, \hat{n}) = \sum_p \mathcal{P}_p(f) e_p(\hat{n})$$

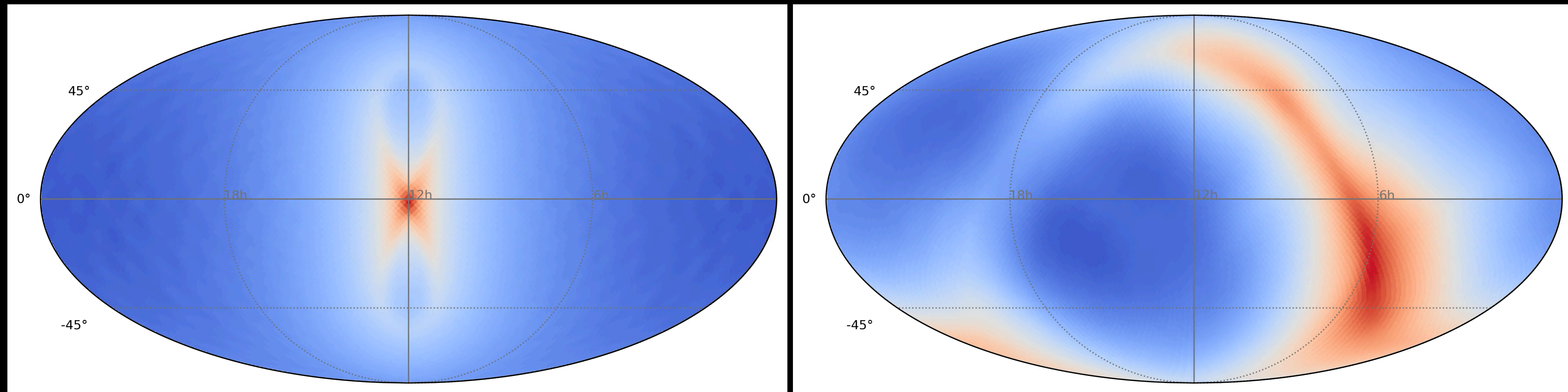
ANISOTROPIC SEARCH

Cross-correlation is essentially a one-dimensional map of the sky.



ANISOTROPIC SEARCH

Cross-correlation is essentially a one-dimensional map of the sky.



- Broadband: point sources with different power-law spectra.
- Narrowband: point sources having narrow GW frequency band (SN 1987A, ScoX-1, GC)
- Spherical harmonics search: Extended or diffuse sources - measure angular power spectra

All-sky BBR Results			Max SNR (% <i>p</i> -value)				Upper limit ranges (10^{-8})	
α	Ω_{GW}	$H(f)$	HL(O3)	HV(O3)	LV(O3)	O1 + O2 + O3 (HLV)	O1 + O2 + O3 (HLV)	O1 + O2 (HL)
0	Constant	$\propto f^{-3}$	2.3 (66)	3.4 (24)	3.1 (51)	2.6 (23)	1.7–7.6	4.4–21
2/3	$\propto f^{2/3}$	$\propto f^{-7/3}$	2.5 (59)	3.7 (14)	3.1 (62)	2.7 (24)	0.85–4.1	2.3–12
3	$\propto f^3$	Constant	3.7 (32)	3.6 (47)	4.1 (12)	3.6 (20)	0.013–0.11	0.046–0.32

- Broadband: point sources with different power-law spectra.
- Narrowband: point sources having narrow GW frequency band (SN 1987A, ScoX-1, GC)
- Spherical harmonics search: Extended or diffuse sources - measure angular power spectra

Narrow band Radiometer Results					
Direction	Max SNR	p-value (%)	Frequency (Hz) (± 0.016 Hz)	Best upper limit (10^{-25})	Frequency band (Hz)
Scorpius X-1	4.1	65.7	630.31	2.1	189.31–190.31
SN 1987A	4.9	1.8	414.0	1.7	185.13–186.13
Galactic Center	4.1	62.3	927.25	2.1	202.56–203.56

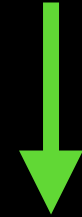
- Broadband: point sources with different power-law spectra.
- Narrowband: point sources having narrow GW frequency band (SN 1987A, ScoX-1, GC)
- Spherical harmonics search: Extended or diffuse sources - measure angular power spectra

SHD Results								
α	Ω_{GW}	$H(f)$	Max SNR (% p -value)				Upper limit range (10^{-9})	
			HL(O3)	HV(O3)	LV(O3)	O1 + O2 + O3 (HLV)	O1 + O2 + O3 (HLV)	O1 + O2 (HL)
0	Constant	$\propto f^{-3}$	1.6 (78)	2.1 (40)	1.5 (83)	2.2 (43)	3.2–9.3	7.8–29
2/3	$\propto f^{2/3}$	$\propto f^{-7/3}$	3.0 (13)	3.9 (0.98)	1.9 (82)	2.9 (18)	2.4–9.3	6.4–25
3	$\propto f^3$	Constant	3.9 (12)	4.0 (10)	3.9 (11)	3.2 (60)	0.57–3.4	1.9–11

PyStoch : fast HEALPix based GWB mapmaking

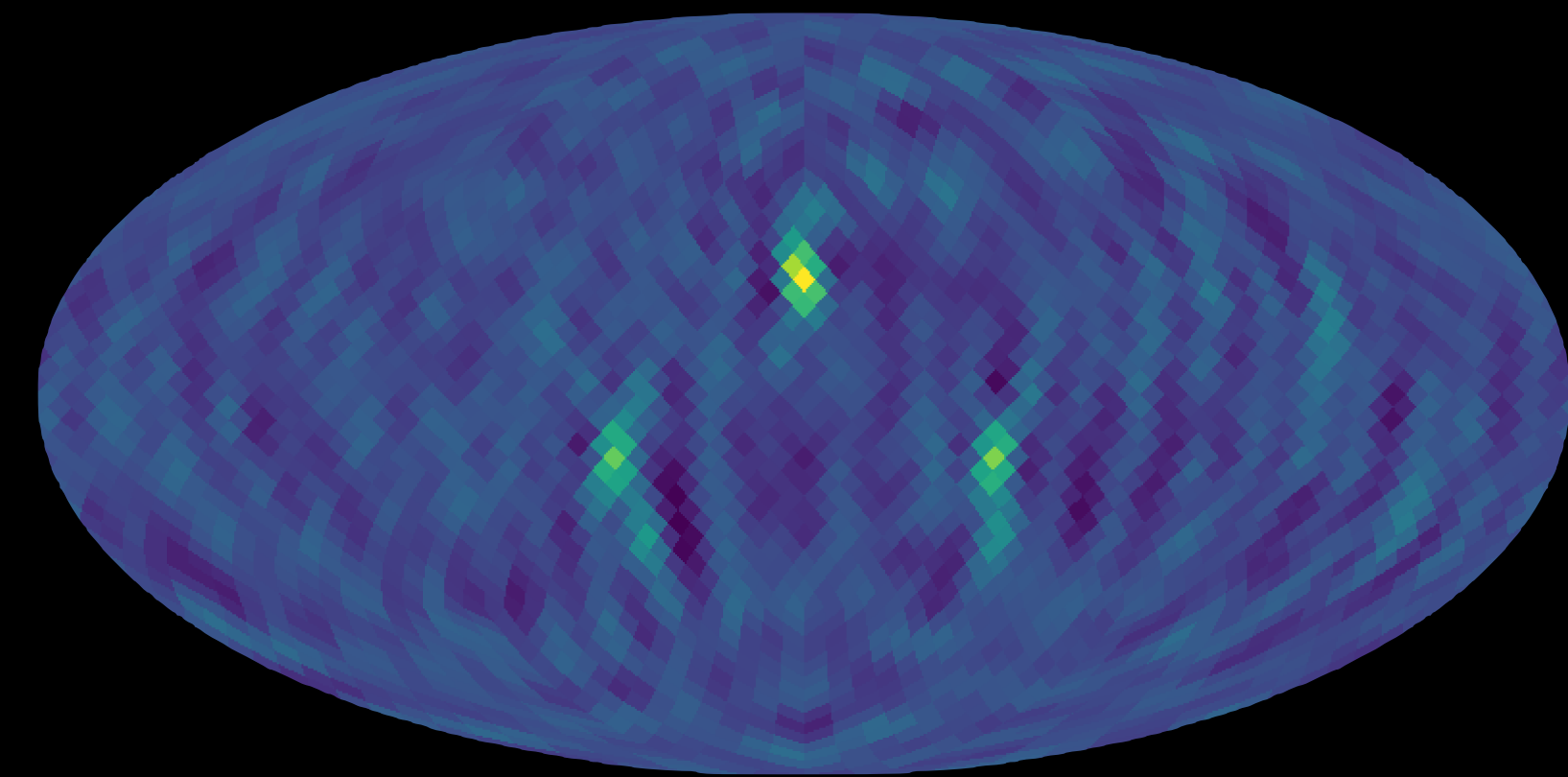
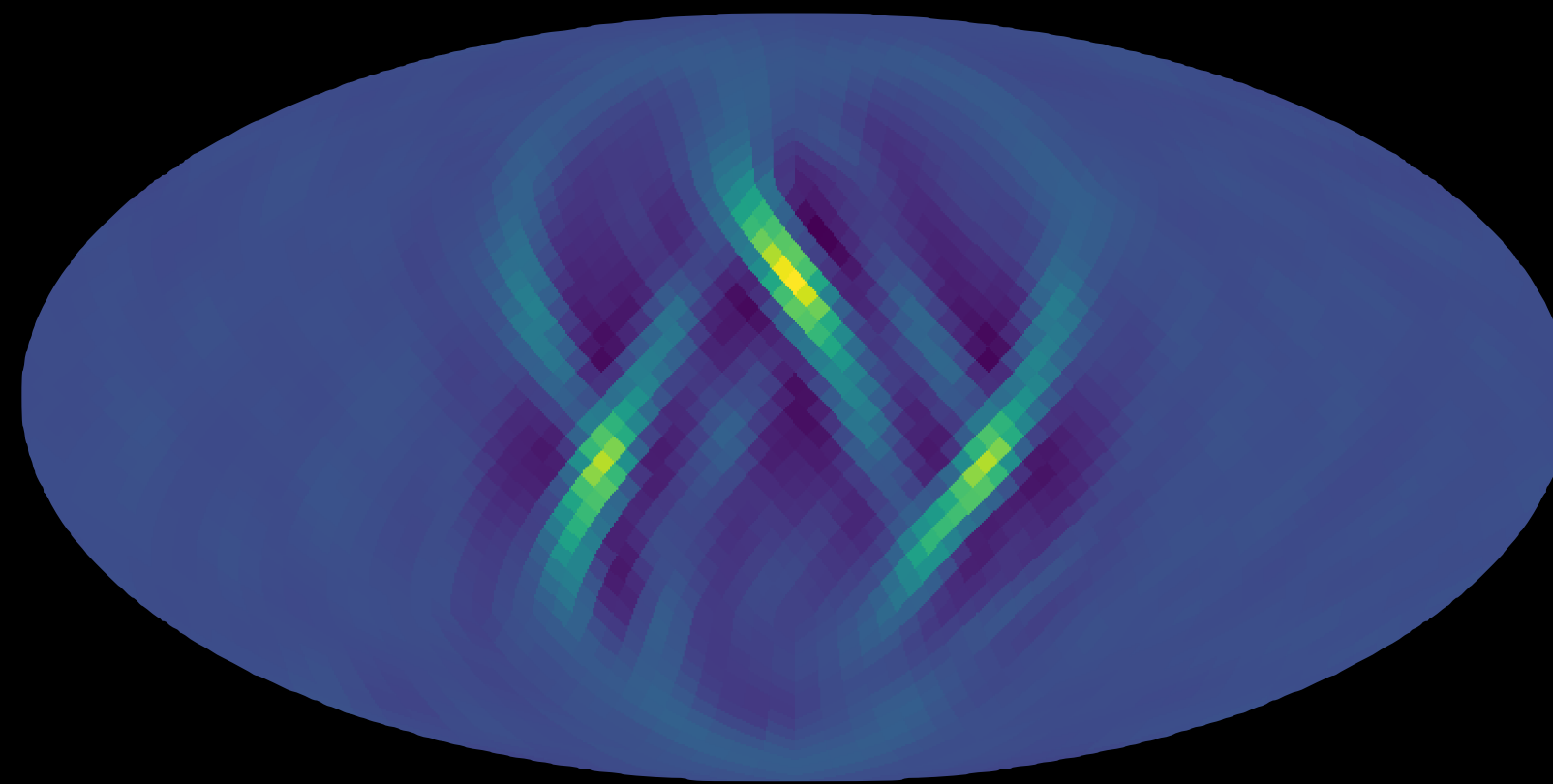
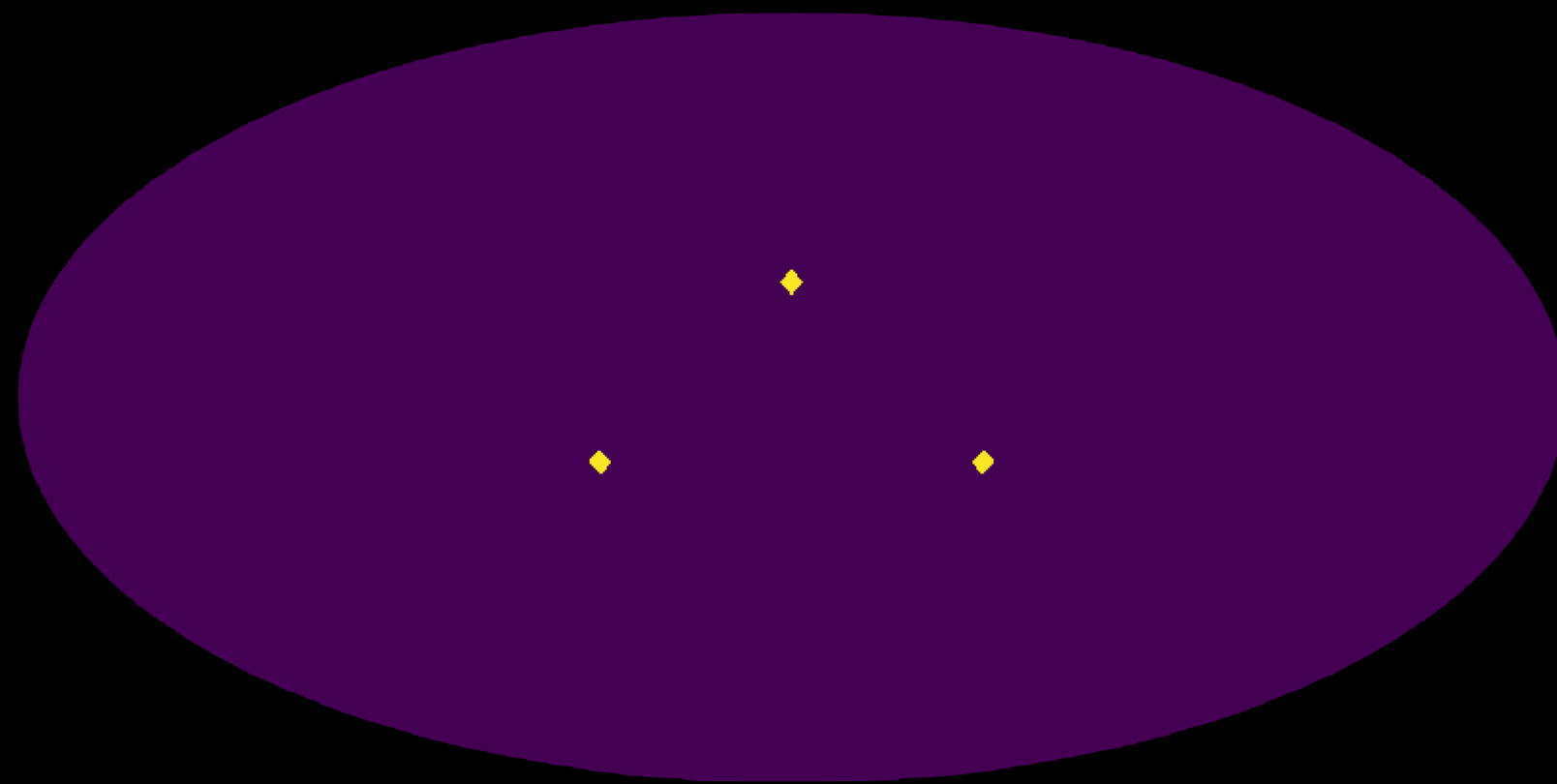


perform the whole analysis on a laptop in a few minutes

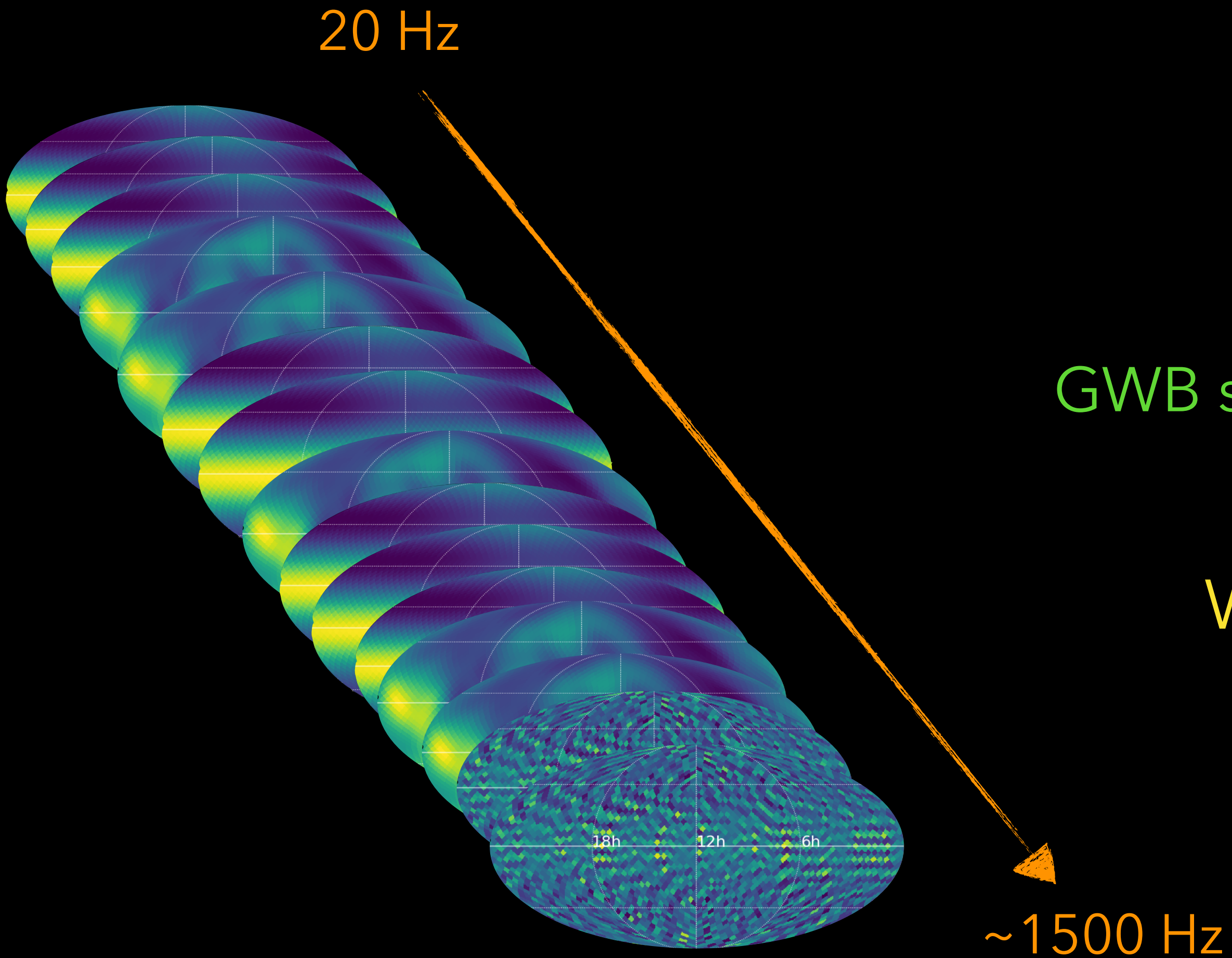


Produces the narrowband maps as an intermediate result

so separate search for different frequency spectra becomes redundant



Now we have all the ingredients to perform an all-sky, all-frequency search, which assumes **no** specific power-law model for the GWB



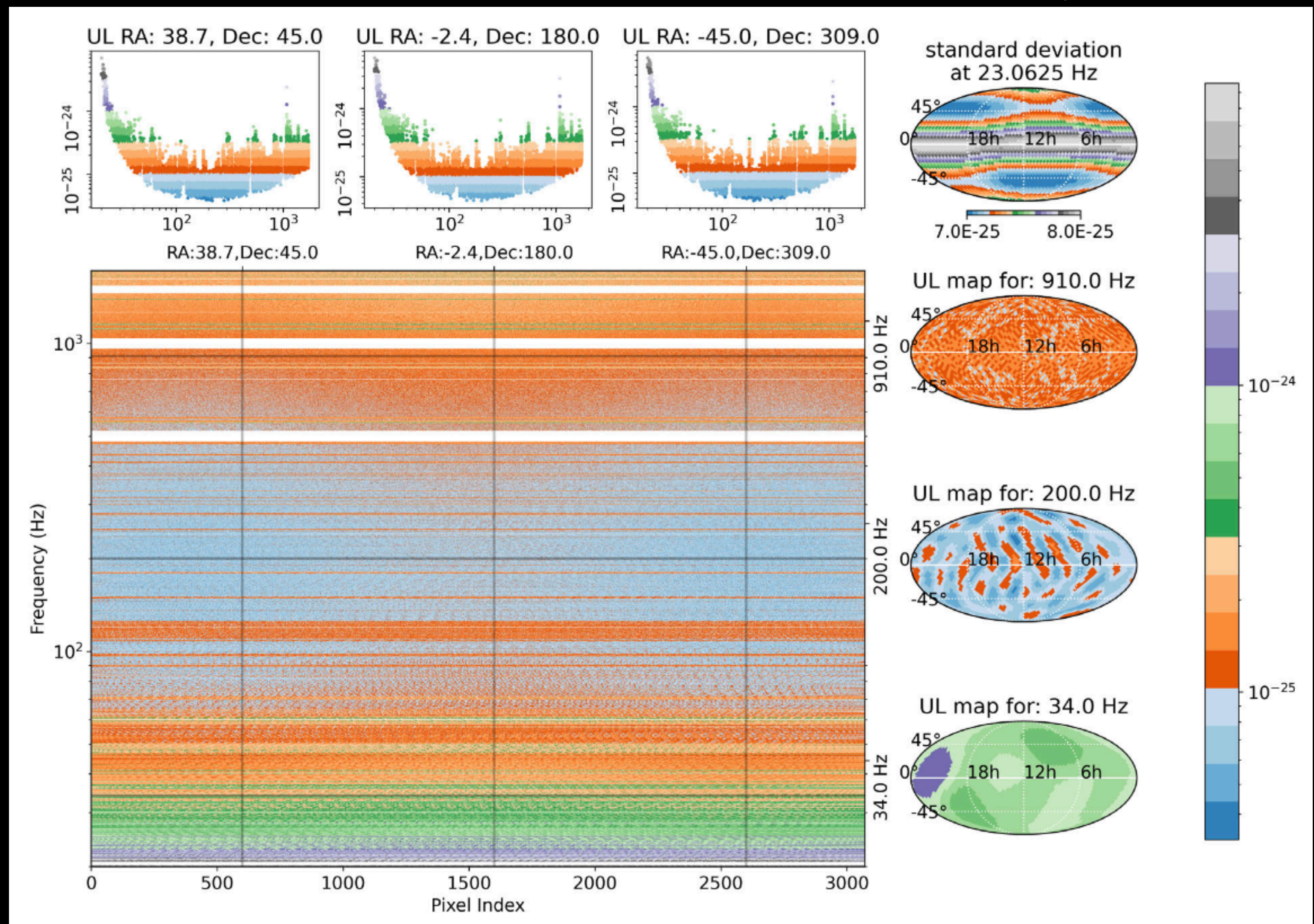
GWB sky maps at every frequency bins

We presented the **first atlas** of GWB sky

Phys.Rev.D 105 (2022) 10, 102001

Given no detection, we set the all-sky all-frequency upper limits on the GWB effective strain*:

$$h(f, \hat{n}) = \sqrt{\mathcal{P}(f, \hat{n}) df}$$



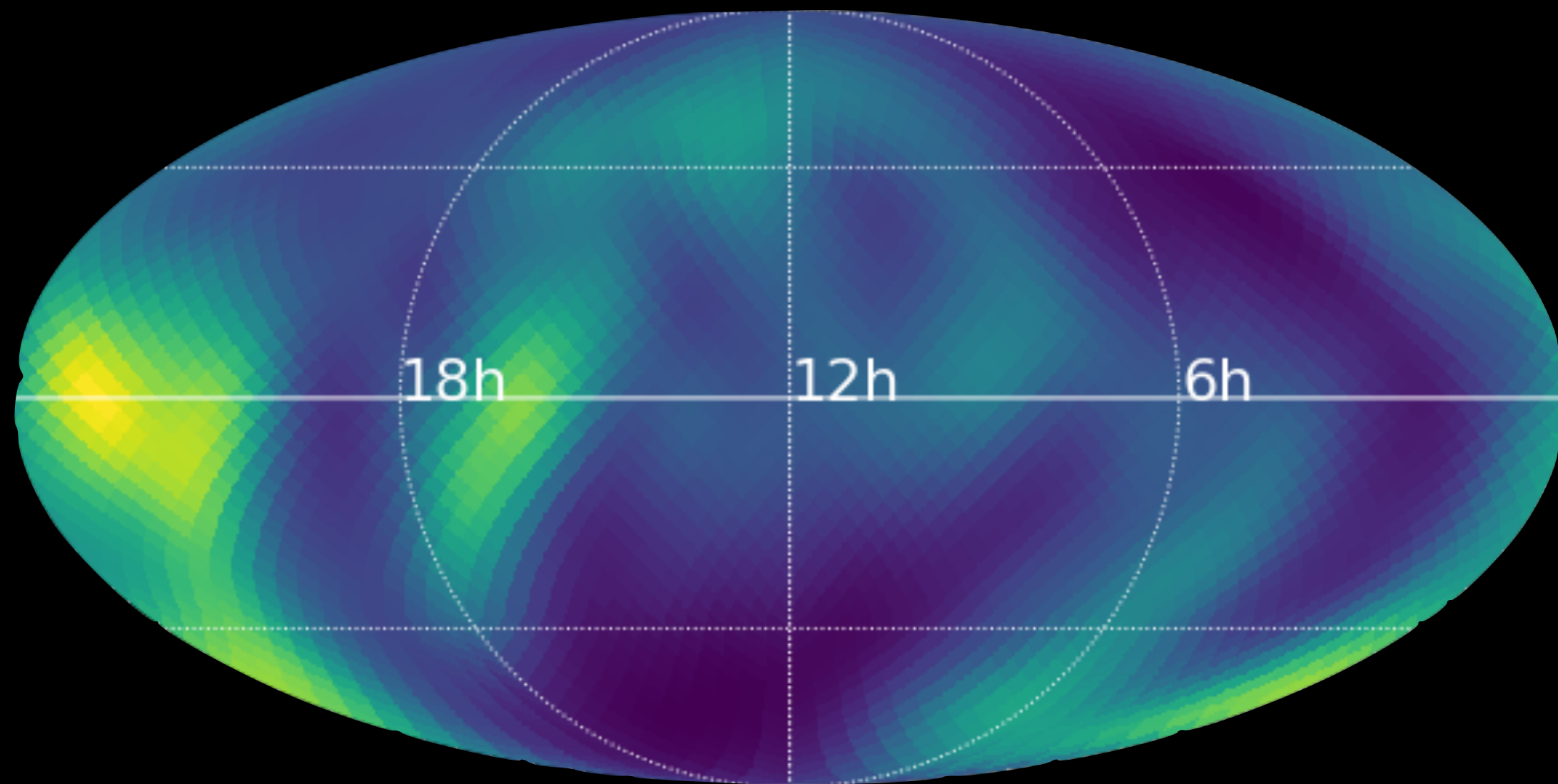
- The **colour bar** here denotes the range of **upper limit variations**.
- The **vertical** cross-section in this diagram shows the **frequency-dependent upper limit in a particular direction**.
- The **Horizontal** cross-sections form a **map of upper limits in a particular frequency**.
- **Notched frequencies** in a baseline appear as horizontal **white bands** in the plot.

*circular polarisation without Doppler correction

Assume a power law and combine these narrowband maps to obtain the 'usual' broadband results

$$\hat{\mathcal{P}}(\hat{\mathbf{n}}) = \frac{\sum_f \hat{\mathcal{P}}(f, \hat{\mathbf{n}}) \sigma_{\hat{\mathbf{n}}}^{-2}(f) H(f)}{\sum_f \sigma_{\hat{\mathbf{n}}}^{-2}(f) H^2(f)}$$

$$\sigma_{\hat{\mathbf{n}}} = \left[\sum_f \sigma_{\hat{\mathbf{n}}}^{-2}(f) H^2(f) \right]^{-1/2}$$

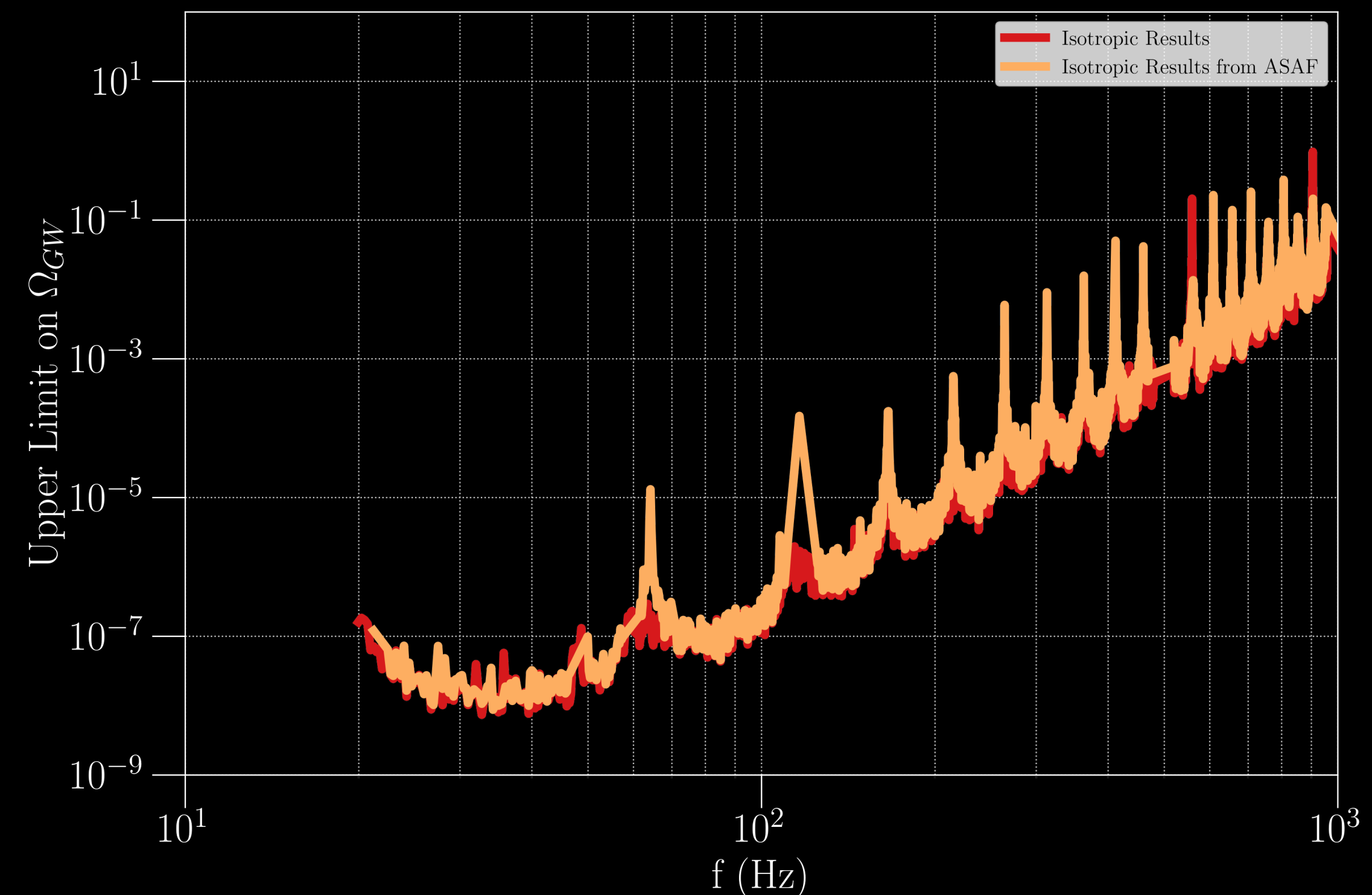


$\alpha = 0$, spectral shape

Assume a power law and sum over all the directions of these narrowband maps to obtain the 'usual' isotropic results

$$\hat{\mathcal{P}}_{\text{iso}}(f) \sigma_{\text{iso}}^{-2}(f) = \frac{5}{4\pi} \int d\hat{\mathbf{n}} \hat{\mathcal{P}}(f, \hat{\mathbf{n}}) \sigma_{\hat{\mathbf{n}}}^{-2}(f)$$

$$\sigma_{\text{iso}}^{-2}(f) = \left(\frac{5}{4\pi} \right)^2 \int d\hat{\mathbf{n}} \int d\hat{\mathbf{n}}' \Gamma_{\hat{\mathbf{n}}, \hat{\mathbf{n}}'}(f)$$

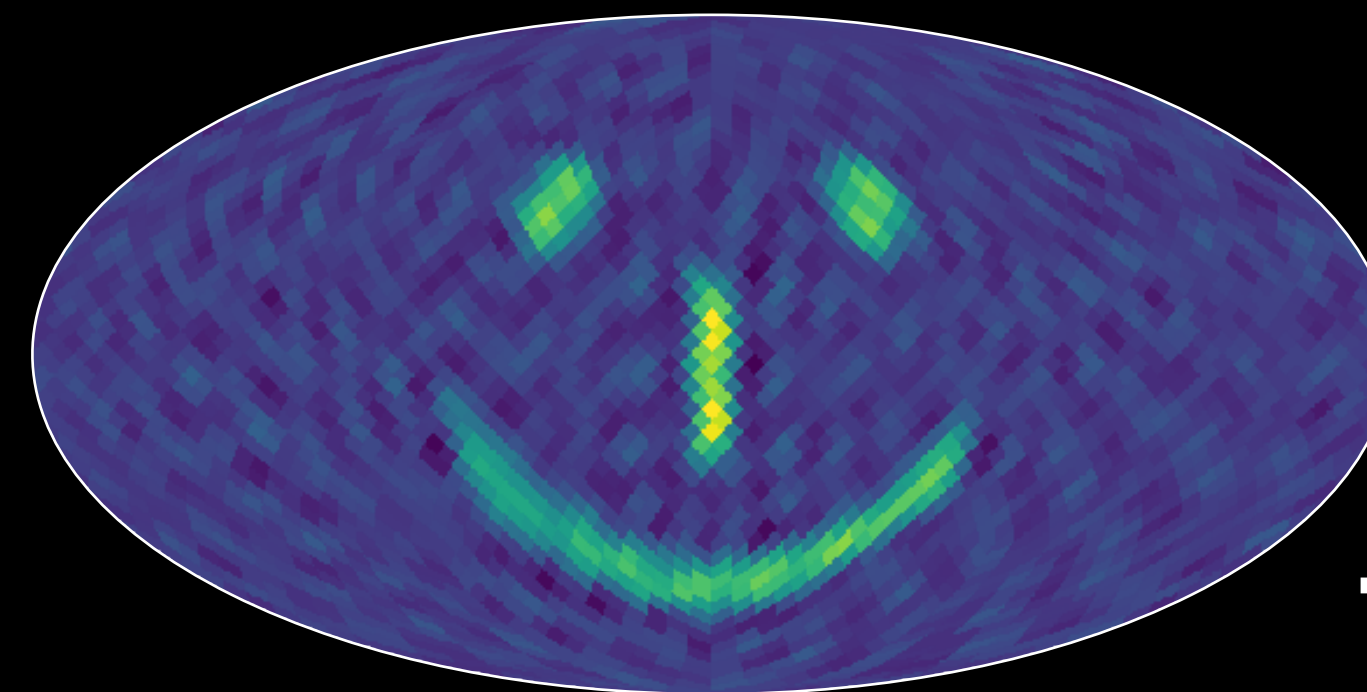


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 - Weak-signal limit, polarized background, non-gaussianity, non-stationarity...
- Plenty more work to do!
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Thank you!

Detectors in **different locations** and with **different orientations** respond differently to a passing GW.

Overlap function encodes reduction in sensitivity of a cross-correlation analysis due to separation and misalignment of the detectors.



$$\gamma_{ft,p}^{IJ} = \sum_A F_I^A(\hat{\Omega}, t) F_J^A(\hat{\Omega}, t) e^{2\pi i f \hat{\Omega} \cdot \Delta \mathbf{x}_J(t)/c}$$

OVERLAP REDUCTION FUNCTION

