



# Conventional and Exotic Heavy Quarkonium Production and Spectroscopy in ATLAS

Dvij Chaitanya Mankad

(On Behalf of the ATLAS Collaboration)

Department of Particle Physics & Astrophysics,  
Weizmann Institute of Science, Israel

EPS-HEP 2025, Marseille | July 07-11, 2025

<https://indico.in2p3.fr/event/33627/contributions/155127/>

# Introduction

- Flavor Physics at ATLAS
  - Measurements of Production Cross-Sections & Spectroscopy in This Talk!
  - Measurements of Weak Decays
- Despite being a general purpose detector, competitive at flavor physics...
  - Large statistics
  - ~Full coverage → phase space complements LHCb
  - Good muon performance
  - Often constrained by trigger, however, constantly optimizing :)

# Introduction

- Flavor Physics at ATLAS

- Measurements of Production Cross-Sections & Spectroscopy in This Talk!

- **Conventional Quarkonia:** Differential Production Cross-Section of  $J/\psi$  and  $\psi(2S)$

- **Exotic Quarkonium:** Search for Di-Charmonium Resonances in  $4\mu$  Final States

- Measurements of Weak Decays

- Despite being a general purpose detector, competitive at flavor physics...

- Large statistics,  $\sim$ full coverage  $\rightarrow$  phase space complements LHCb, good muon performance, often constrained by trigger, however, constantly optimizing! :)

# Introduction

## □ Flavor Physics at ATLAS

- Measurements of Production Cross-Sections & Spectroscopy **in This Talk!**

- **Conventional Quarkonia:** Differential Production Cross-Section of  $J/\psi$  and  $\psi(2S)$

- **Exotic Quarkonium:** Search for Di-Charmonium Resonances in  $4\mu$  Final States

- More results in the talk by Adam Barton (9th July, 17:40)

- Measurements of Weak Decays

- See, the talk by Adam Barton (9th July, 17:40)



In the Flavor Session  
in Touze

## □ Despite being a general purpose detector, competitive at flavor physics...

- Large statistics,  $\sim$ full coverage  $\rightarrow$  phase space complements LHCb, good muon performance, **often constrained by trigger, however, constantly optimizing! :)**

# Differential Production Cross-Section of $J/\psi$ and $\psi(2S)$

arXiv:2309.17177

Eur. Phys. J. C 84 (2024) 169

# Motivation

- **Two sources of quarkonia production:**
  - **Prompt:** Coming from short-lived QCD processes.
  - **Non-prompt:** Coming from decays of  $b$ -hadrons.
- **Understanding quarkonium production in hadronic collisions still incomplete:**
  - Perturbative QCD can describe the non-prompt production well but not prompt production.
  - NRQCD approach to build a universal library of LDMEs has achieved mixed success.
  - Color Evaporation Model is simpler in terms of parameters but faces its own problem in describing data.
- **This analysis:**
  - Provides experimental data in **previously unmeasured kinematic range** of quarkonia production!
  - Can help theoretical models with **qualitatively new information**.



# Experimental Strategy

## □ Kinematic range extension:

□  $J/\psi$  :  $p_T < 100 \text{ GeV} \rightarrow p_T < 360 \text{ GeV}$

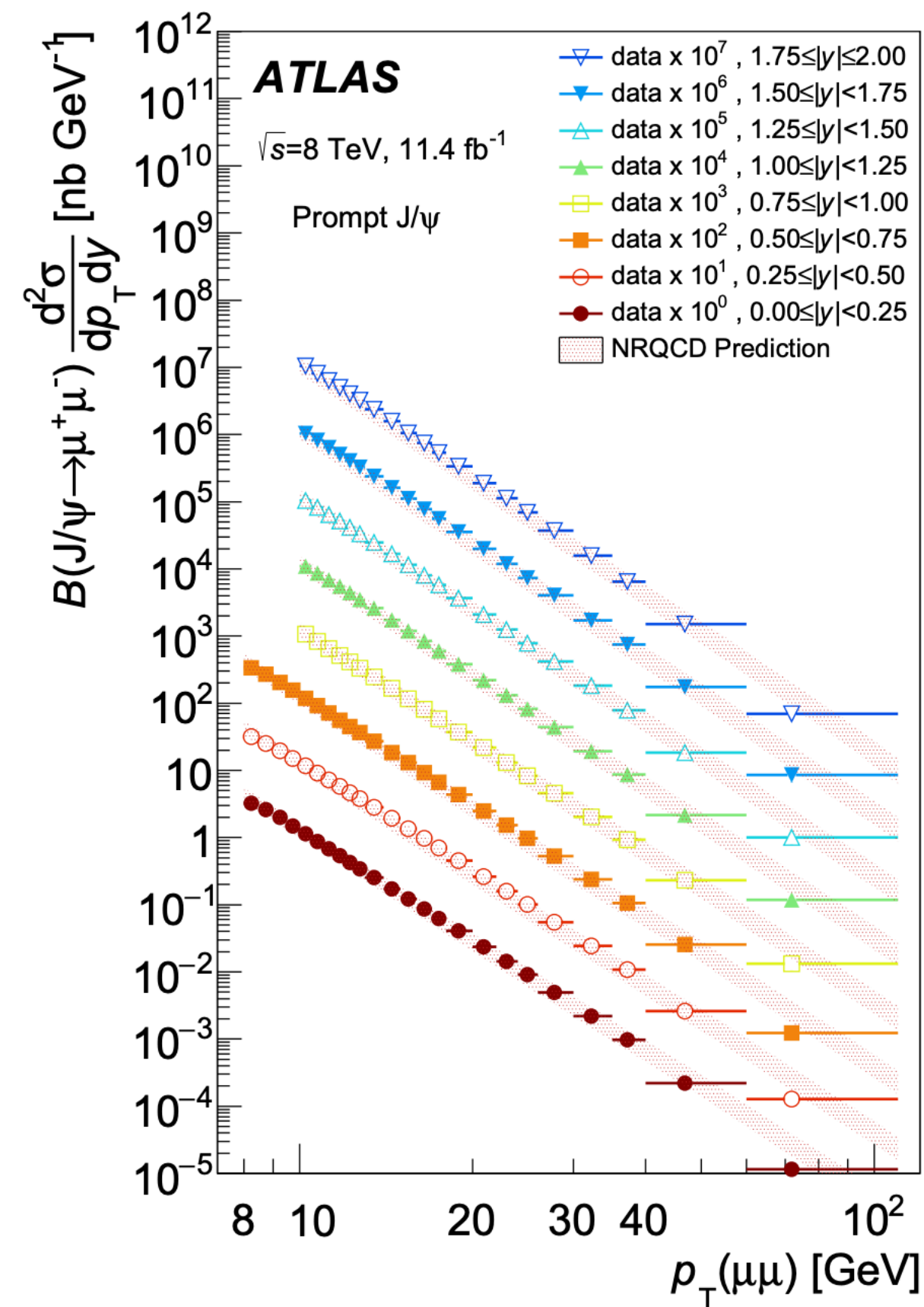
□  $\psi(2S)$  :  $p_T < 100 \text{ GeV} \rightarrow p_T < 140 \text{ GeV}$

## □ Possible due to updated trigger-strategy!

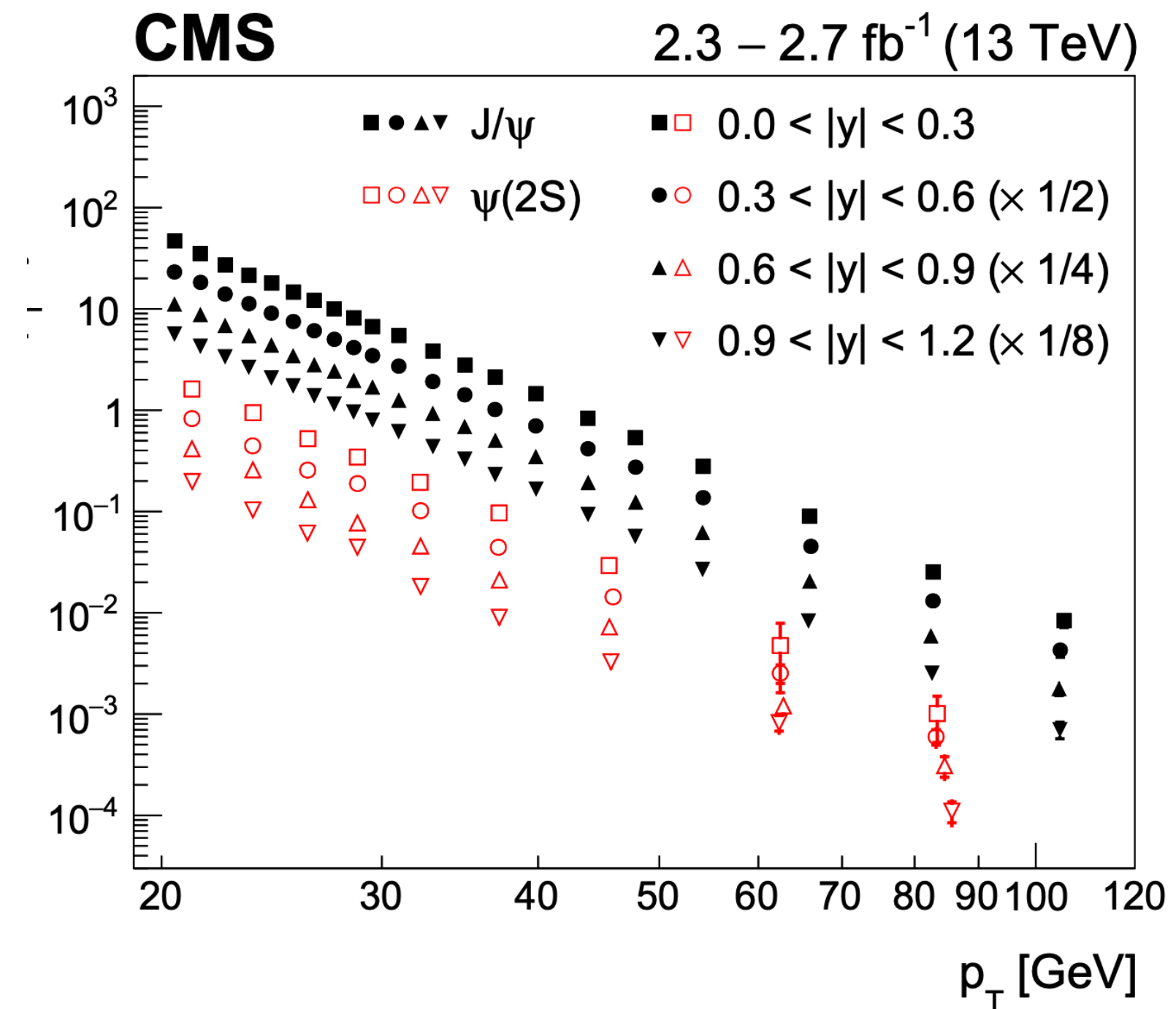
□ Angular resolution of dimuon triggers not sufficient to resolve highly boosted muons coming from charmonia with  $p_T > 100 \text{ GeV}$ .

□ Instead, trigger on a single muon trigger with high muon  $p_T$  threshold ( $\sim 50 \text{ GeV}$ )!

□ Dimuon triggers still used to cover lower charmonia  $p_T$  phase space.



1512.03657  
**(ATLAS Run1)**



1710.11002  
**(CMS Run2)**

# Analysis Strategy

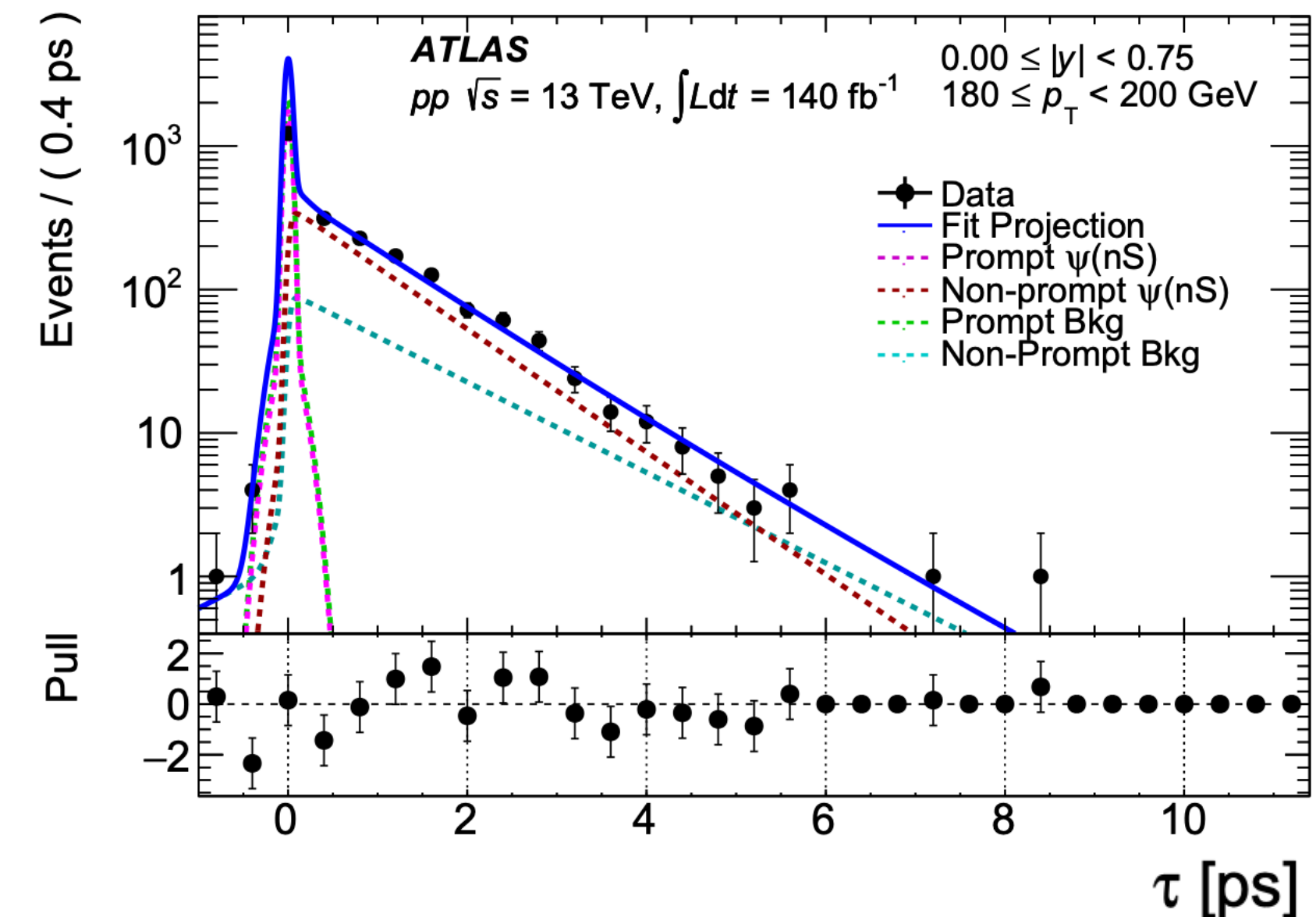
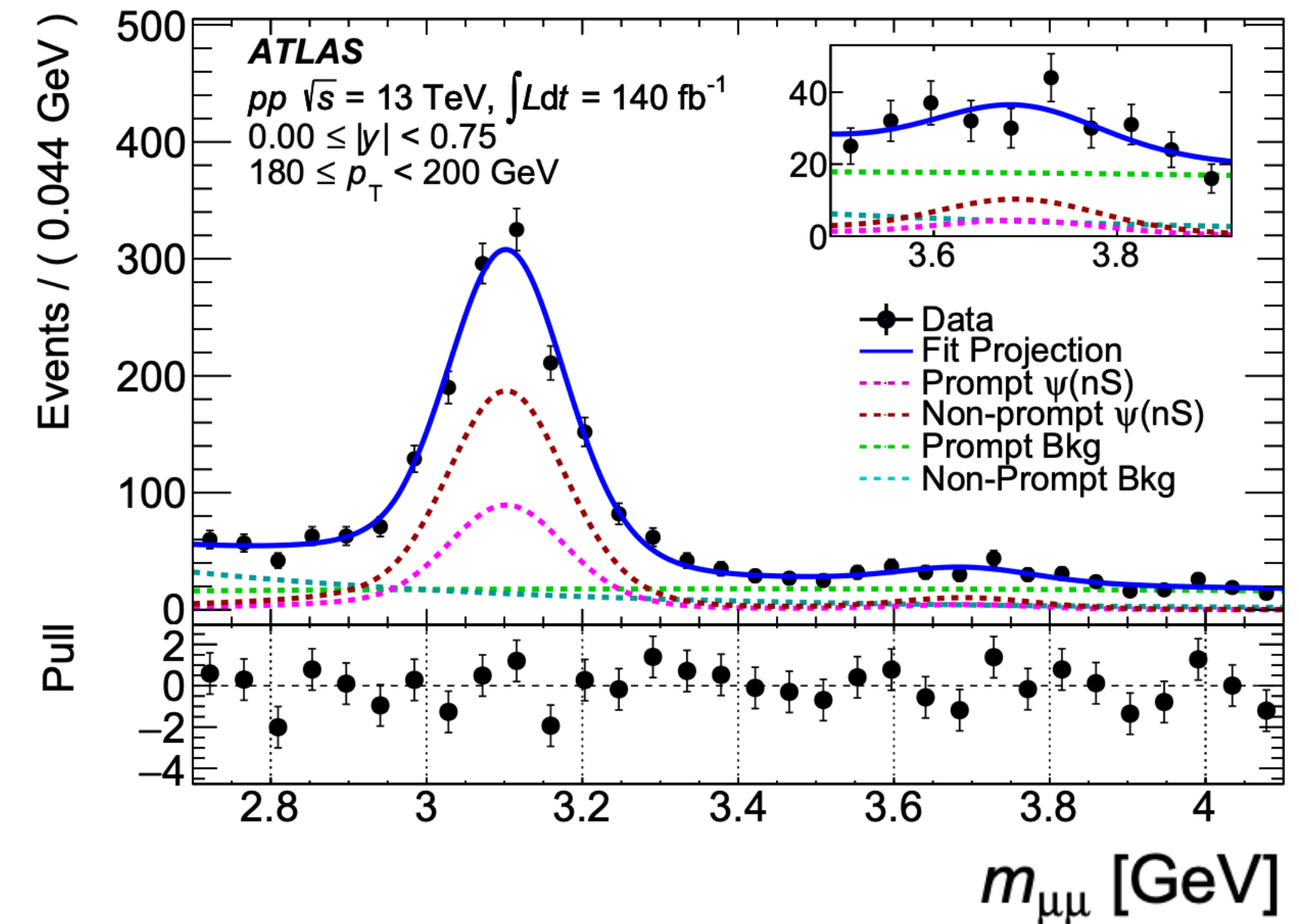
- Prompt and non-prompt contributions separated using the reconstructed pseudo proper lifetime

$$\tau = \frac{m_{\mu\mu} L_{xy}}{p_T c}.$$

- Measured signal yield extracted from **simultaneous fits to reconstructed mass and pseudo proper lifetime** — in bins of  $p_T$  and  $y$ .
- Differential cross-section measured from the measured yield as:

$$\frac{d^2\sigma^{\text{P,NP}}(pp \rightarrow \psi)}{dp_T dy} = \frac{1}{\mathcal{B}(\psi \rightarrow \mu\mu) \int \mathcal{L} dt} \cdot \frac{1}{\mathcal{A}(\psi) \epsilon_{\text{trig}} \epsilon_{\text{reco}} \text{SF}_{\text{trig}} \text{SF}_{\text{reco}}} \cdot \frac{N_{\psi}^{\text{P,NP}}}{\Delta p_T \Delta y}$$

- Fraction of non-prompt production  $F_{\psi}^{\text{NP}}$  and  $\psi(2S)$ -to- $J/\psi$  production ratios  $R_{\psi(2S)}^{\text{P,NP}}$  also **extracted** with partial cancellations of uncertainties!





# Analysis Strategy

- Prompt and non-prompt contributions separated using the reconstructed pseudo proper lifetime

$$\tau = \frac{m_{\mu\mu}}{p_T} \frac{L_{xy}}{c}.$$

- Measured signal yield extracted from **simultaneous fits to reconstructed mass and pseudo proper lifetime** — in bins of  $p_T$  and  $y$ .
- Differential cross-section measured from the measured yield as:

$$\frac{d^2\sigma^{\text{P,NP}}(pp \rightarrow \psi)}{dp_T dy} = \frac{1}{\mathcal{B}(\psi \rightarrow \mu\mu) \int \mathcal{L} dt} \cdot \frac{1}{\mathcal{A}(\psi) \epsilon_{\text{trig}} \epsilon_{\text{reco}} \text{SF}_{\text{trig}} \text{SF}_{\text{reco}}}.$$

- Fraction of non-prompt production  $F_{\psi}^{\text{NP}}$  and  $\psi(2S)$ -to- $J/\psi$  production ratios  $R_{\psi(2S)}^{\text{P,NP}}$  also **extracted** with partial cancellations of uncertainties!

**Non-Prompt-Production Ratio**  
 $R_{J/\psi, \psi(2S)}^{\text{NP}} = F_{J/\psi, \psi(2S)}^{\text{NP}} / (1 - F_{J/\psi, \psi(2S)}^{\text{NP}})$   
**Partial Cancellation of Systematics**

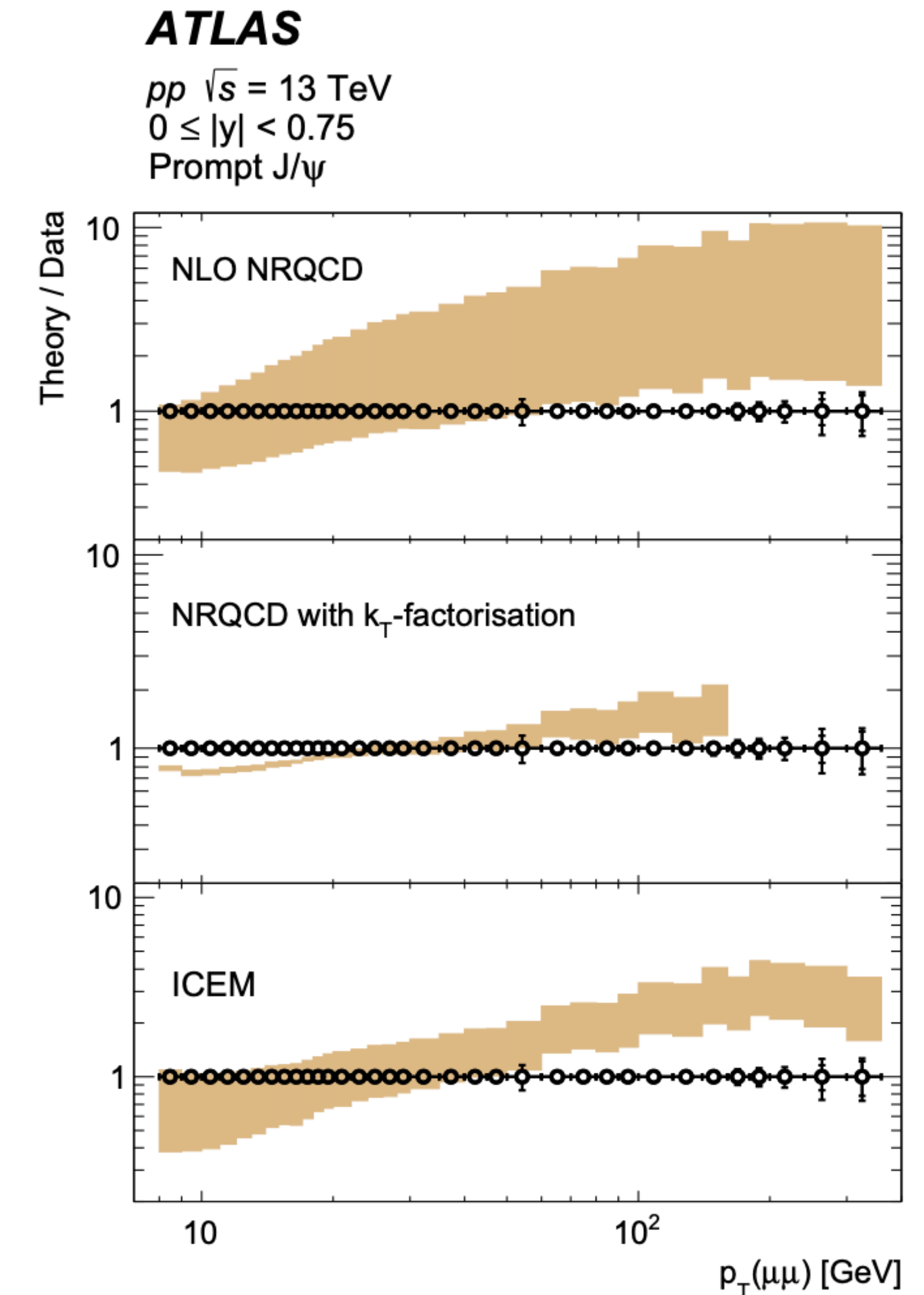
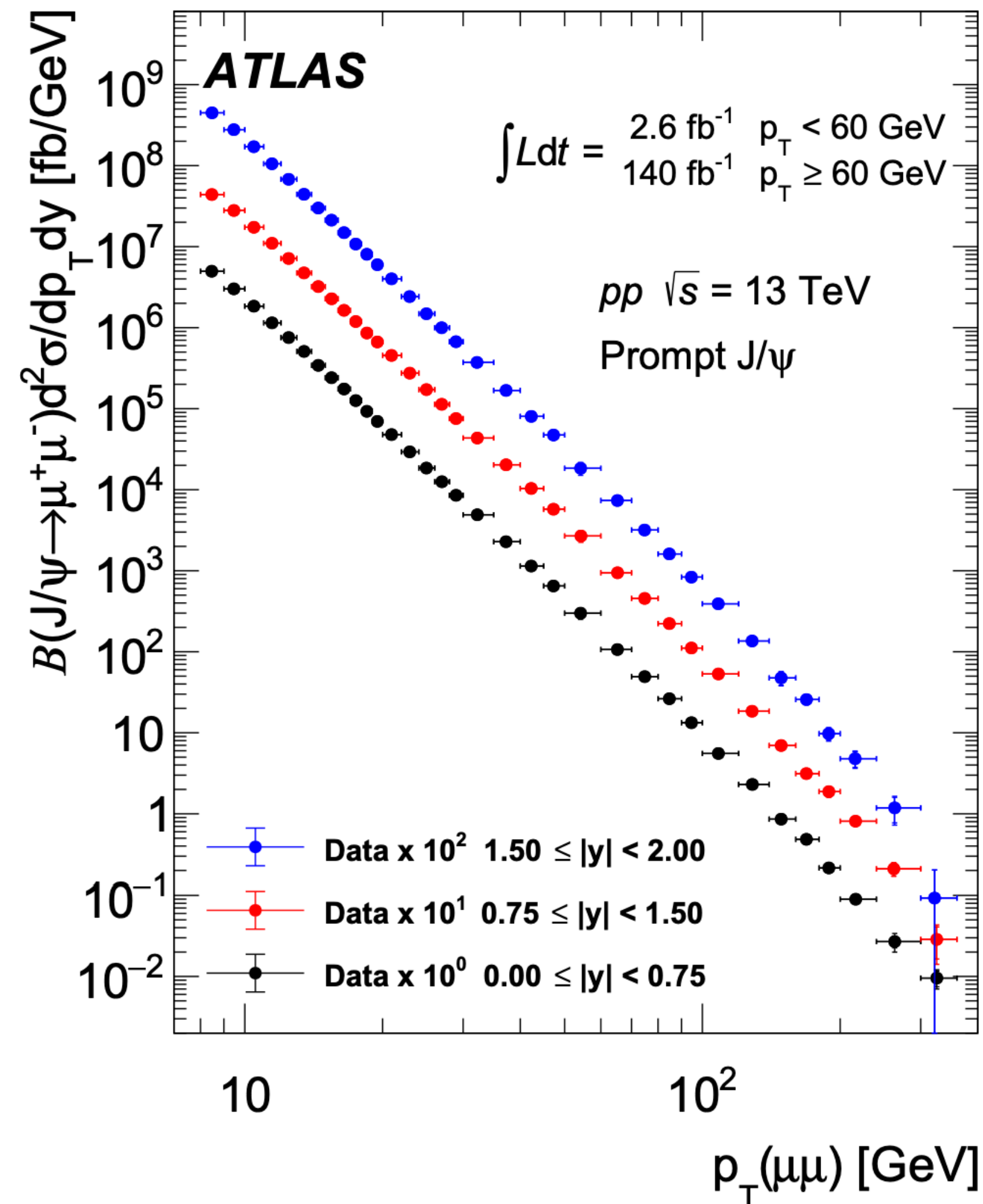
$\frac{d^2}{dp_T dy} \sigma_{J/\psi}^{\text{NP}}$	$\frac{d^2}{dp_T dy} \sigma_{\psi(2S)}^{\text{NP}}$
$\frac{d^2}{dp_T dy} \sigma_{J/\psi}^{\text{P}}$	$\frac{d^2}{dp_T dy} \sigma_{\psi(2S)}^{\text{P}}$

**Partial Cancellation of Systematics**

$\psi(2S)$ -Production Ratio  $R_{\psi(2S)}^{\text{P, NP}}$

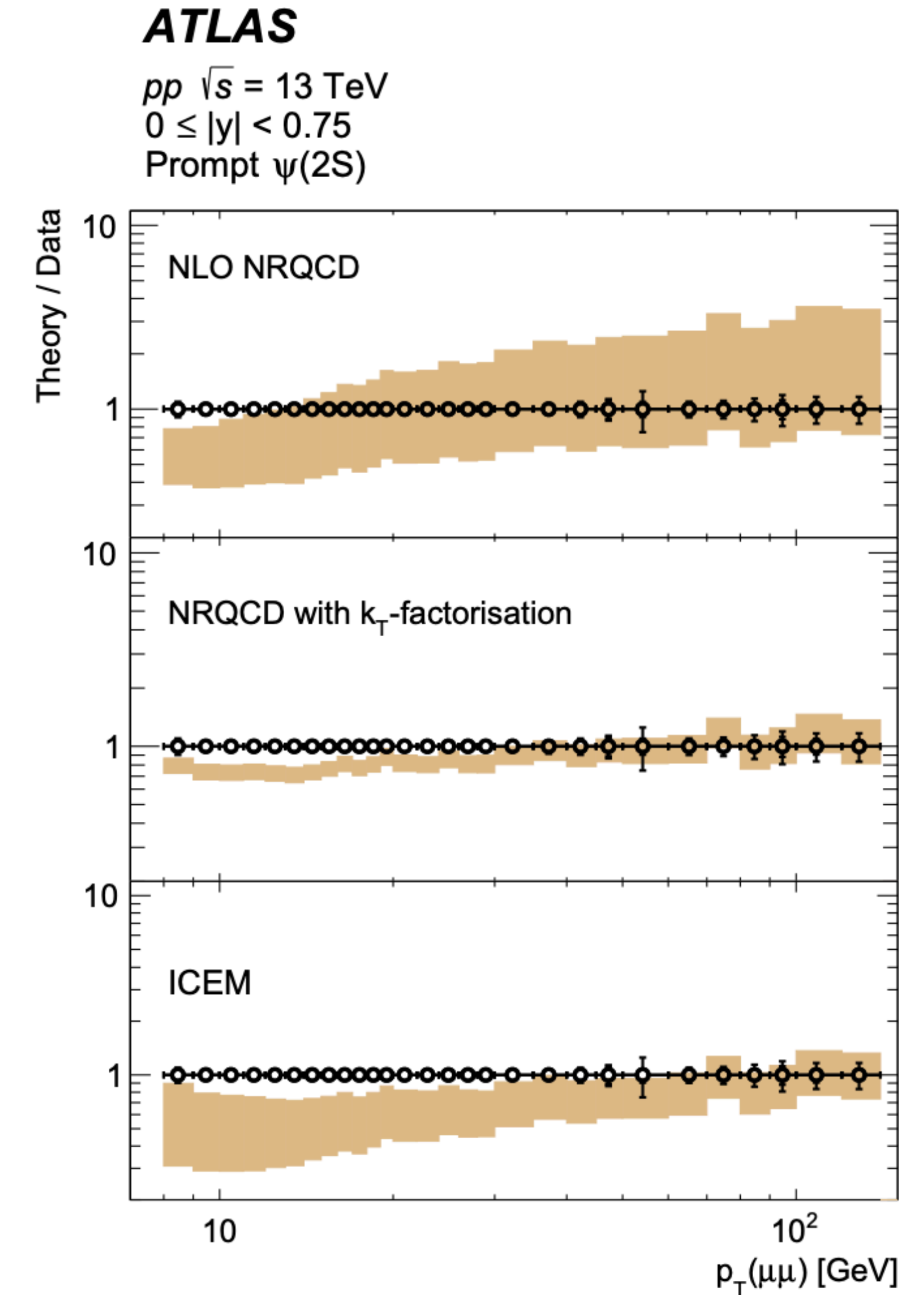
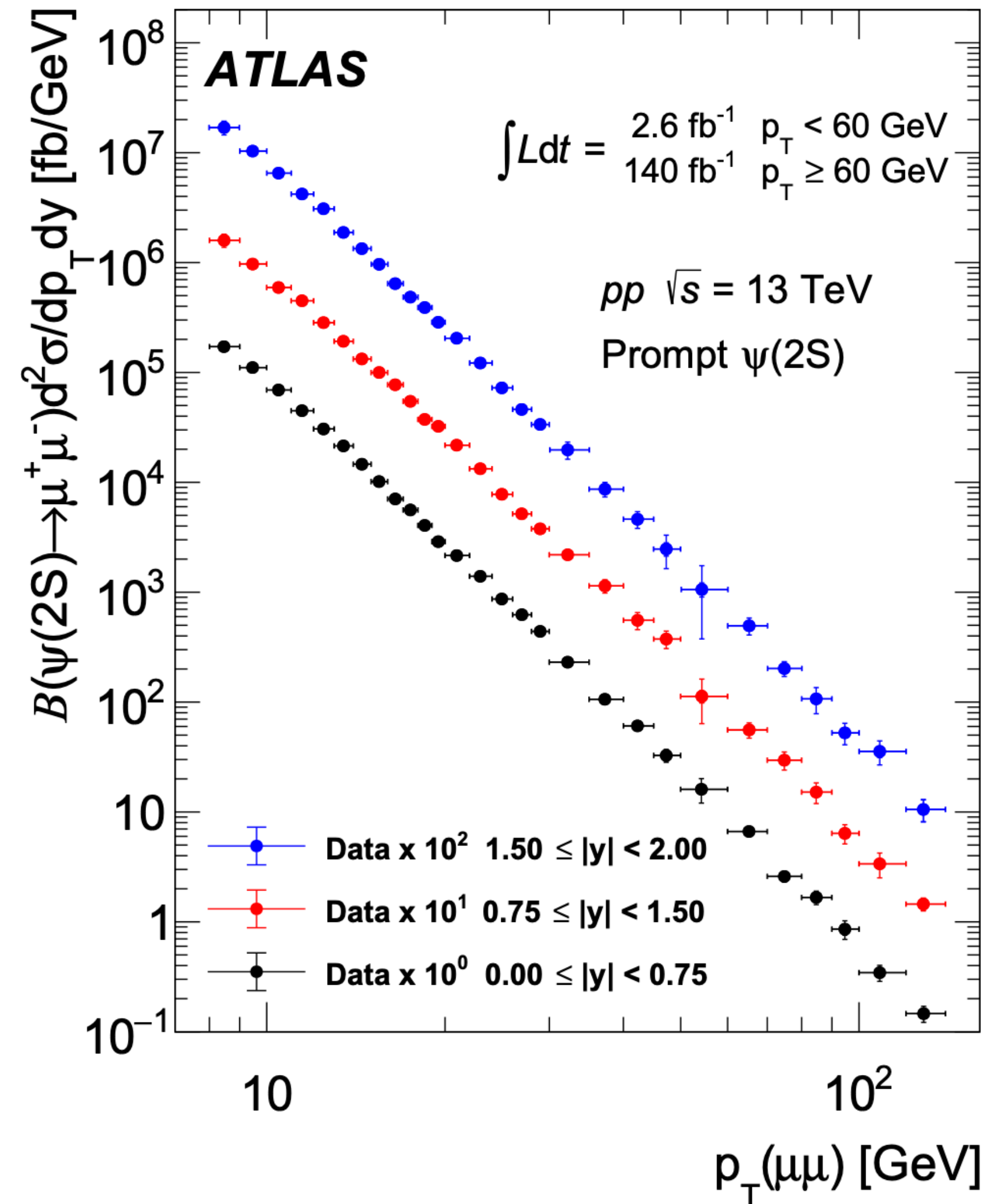
# Results: Prompt $J/\psi$ Production

- NLO NRQCD and ICEM seem to overestimate the high- $p_T$  production.
- NRQCD with  $k_T$ -factorization seems to underestimate the low- $p_T$  production.



# Results: Prompt $\psi(2S)$ Production

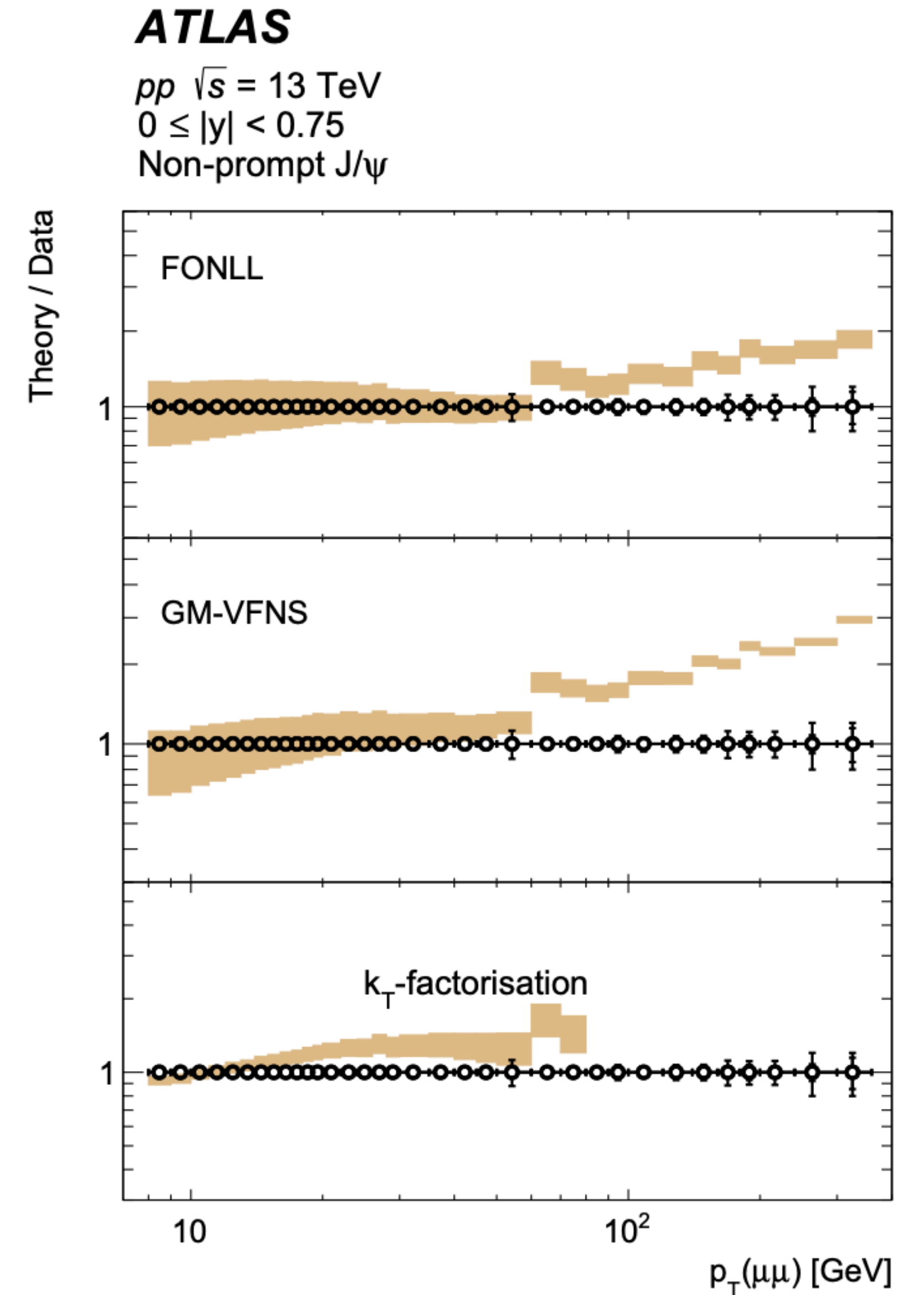
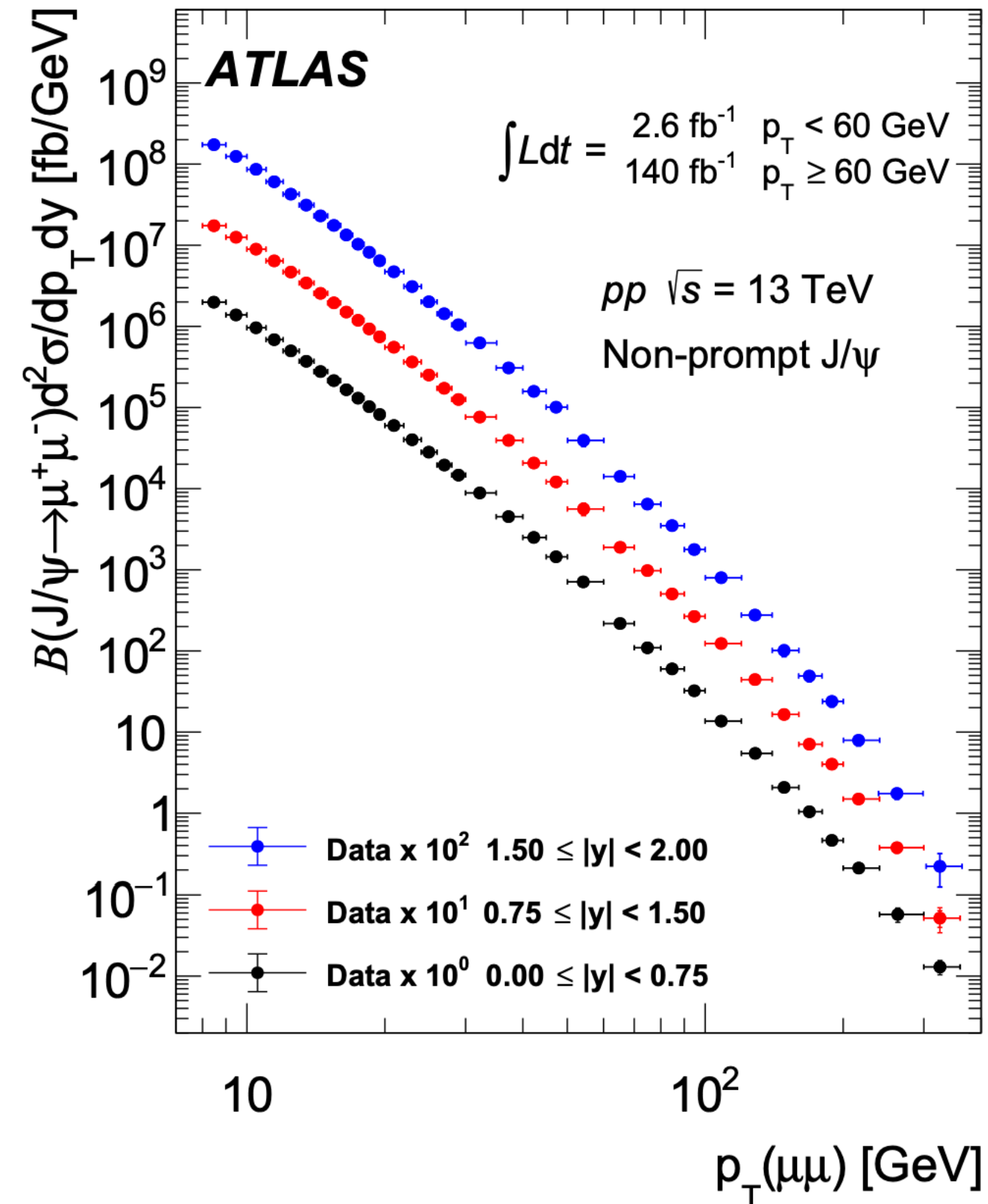
- NLO NRQCD seems to over-estimation of high- $p_T$  production.
- NRQCD with  $k_T$ -factorization seems to underestimate the  $\psi(2S)$  production.





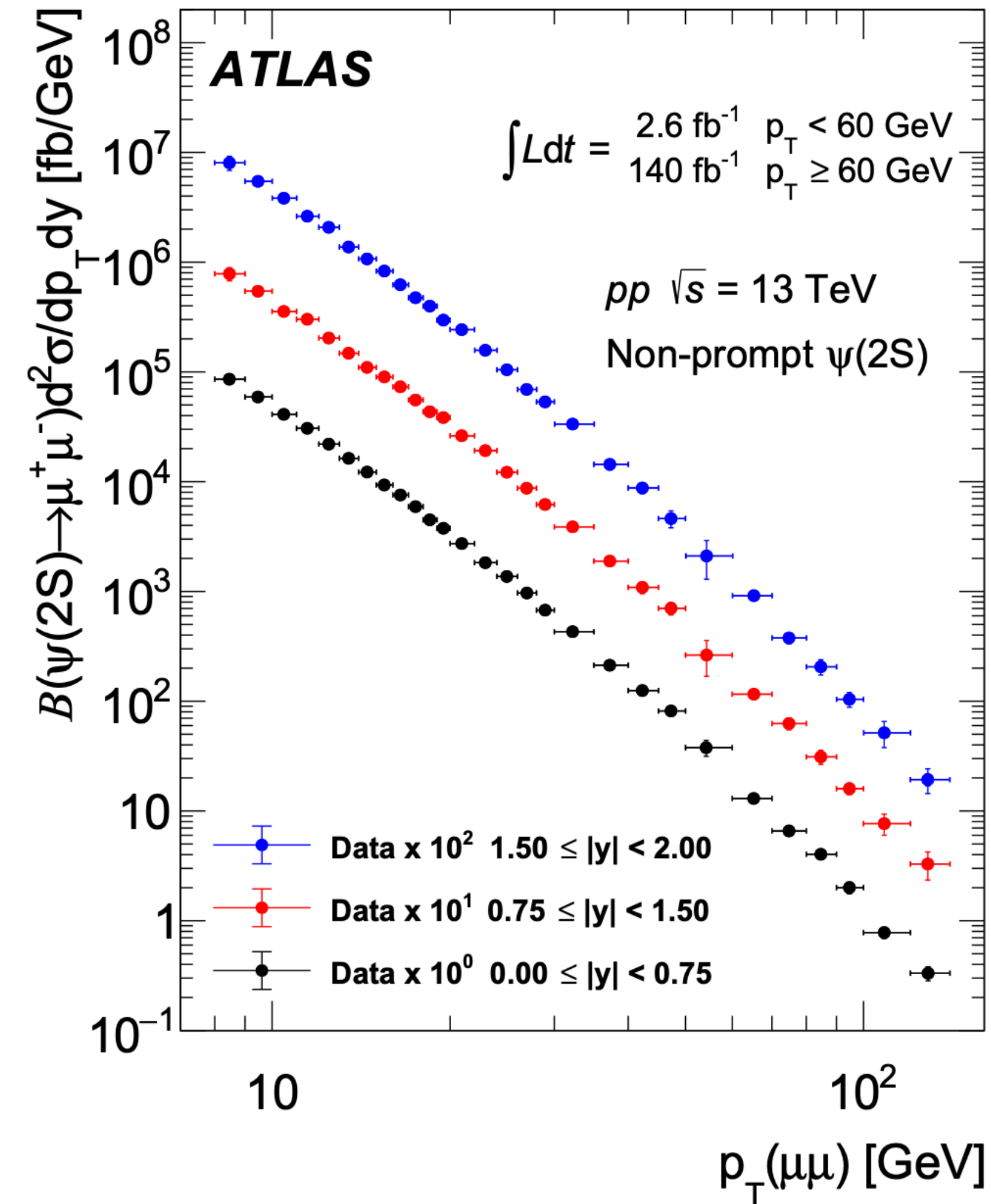
# Results: Non-Prompt $J/\psi$ Production

- FONLL and GM-VFNS seem to overestimate the high- $p_T$  production.
- $k_T$ -factorization limited in availability at high- $p_T$ .



# Results: Non-Prompt $\psi(2S)$ Production

- $k_T$ -factorization seems to underestimate the production at low- $p_T$ .

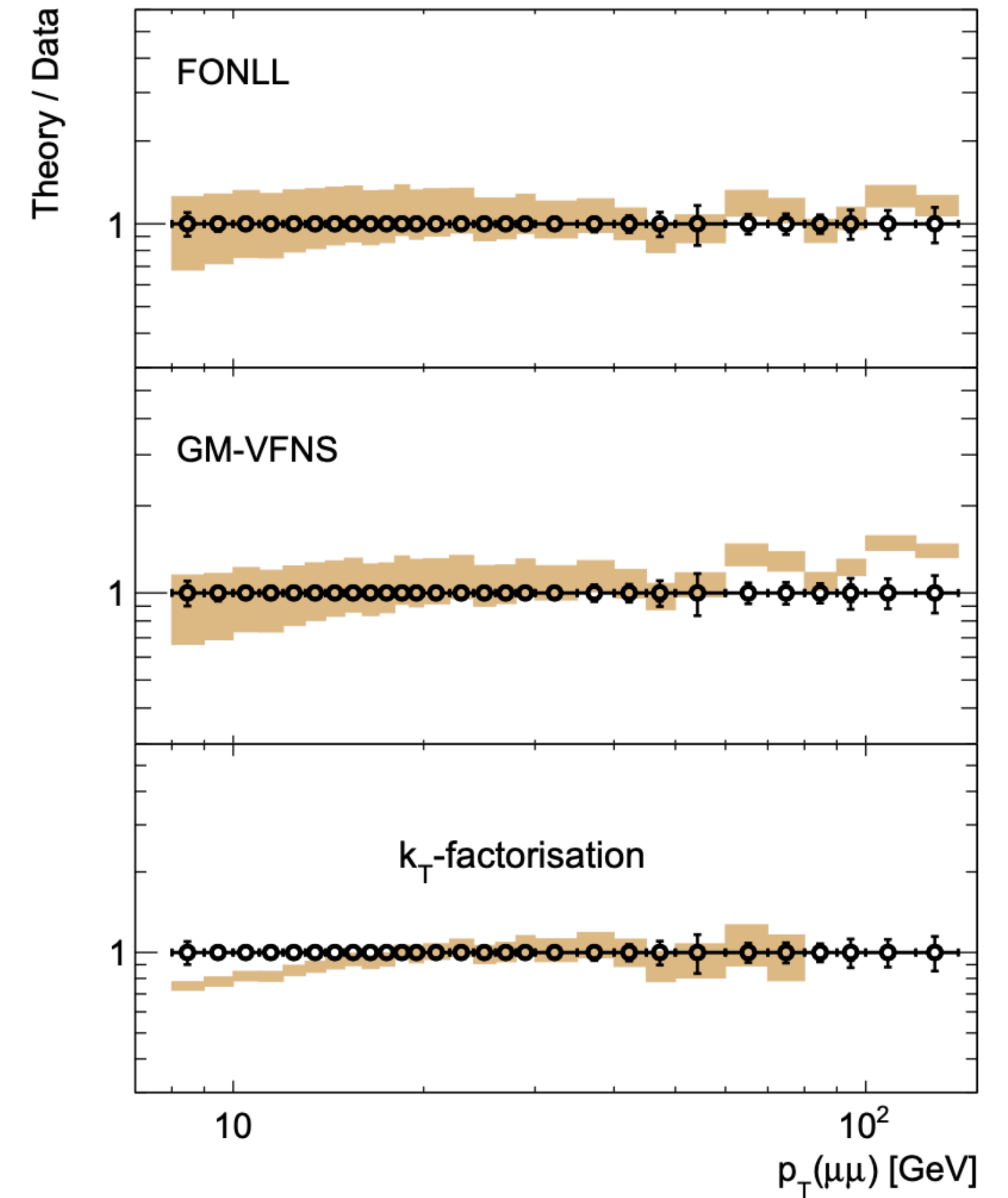


**ATLAS**

$pp \sqrt{s} = 13 \text{ TeV}$

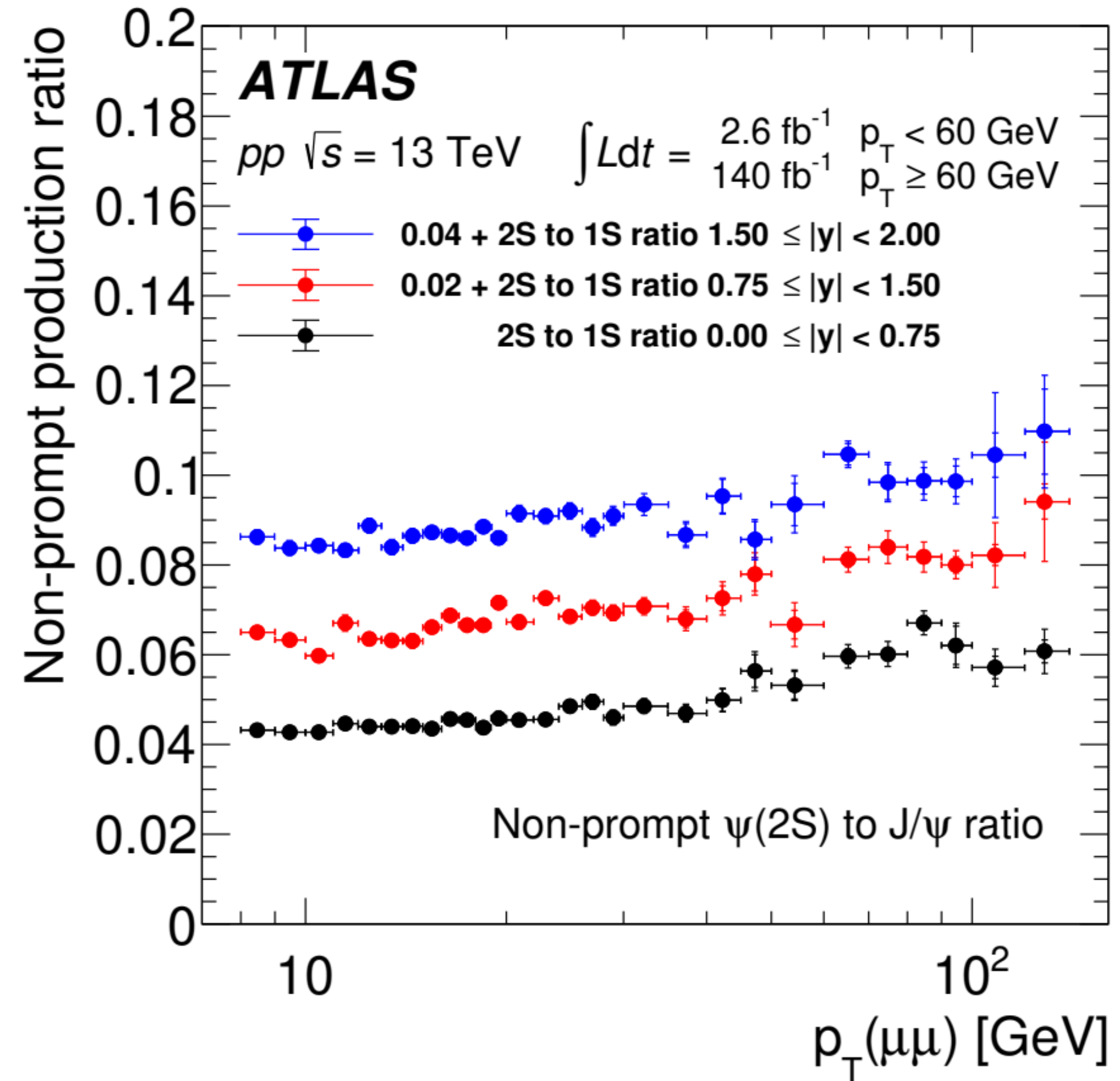
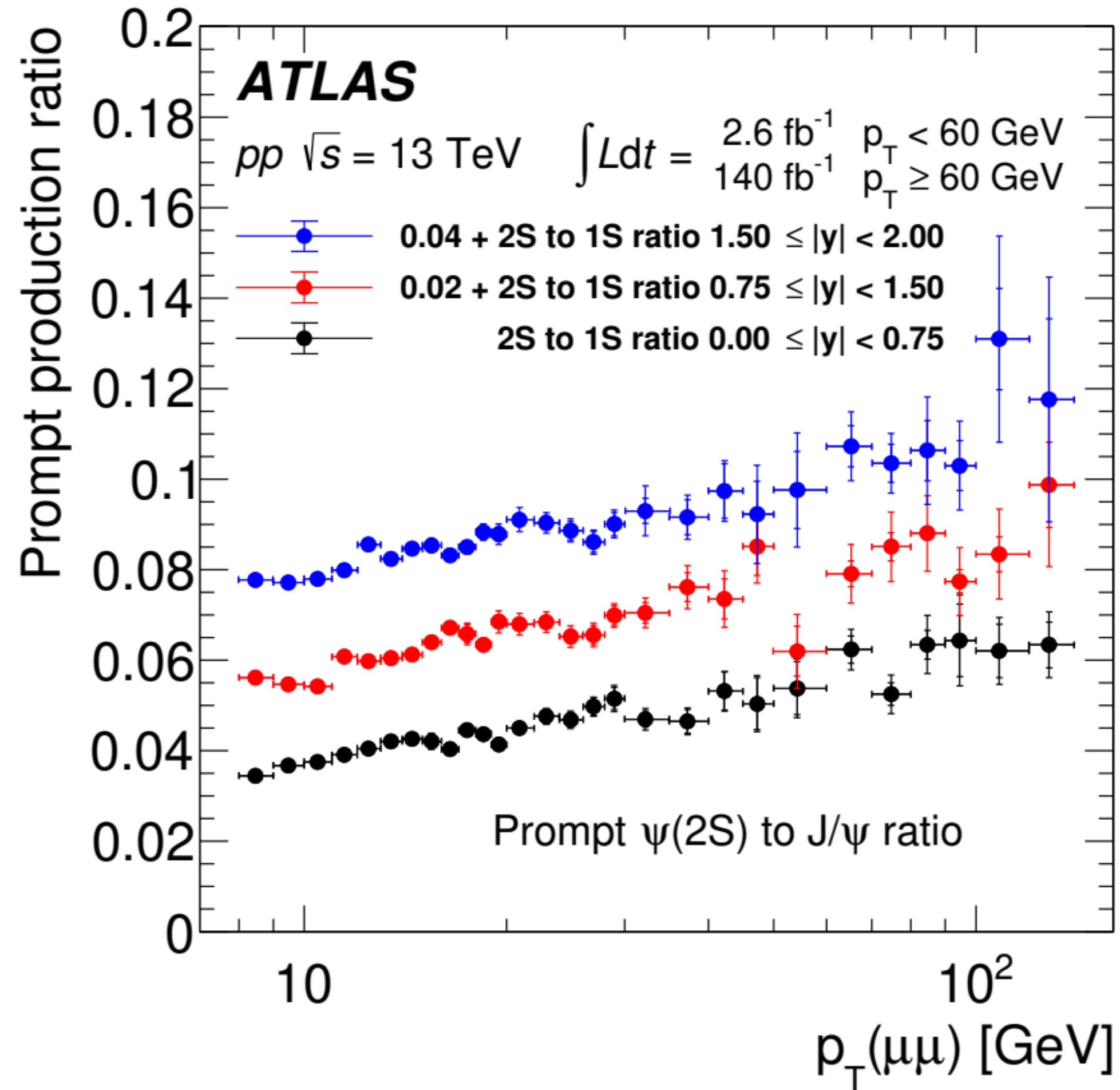
$0 \leq |y| < 0.75$

Non-prompt  $\psi(2S)$



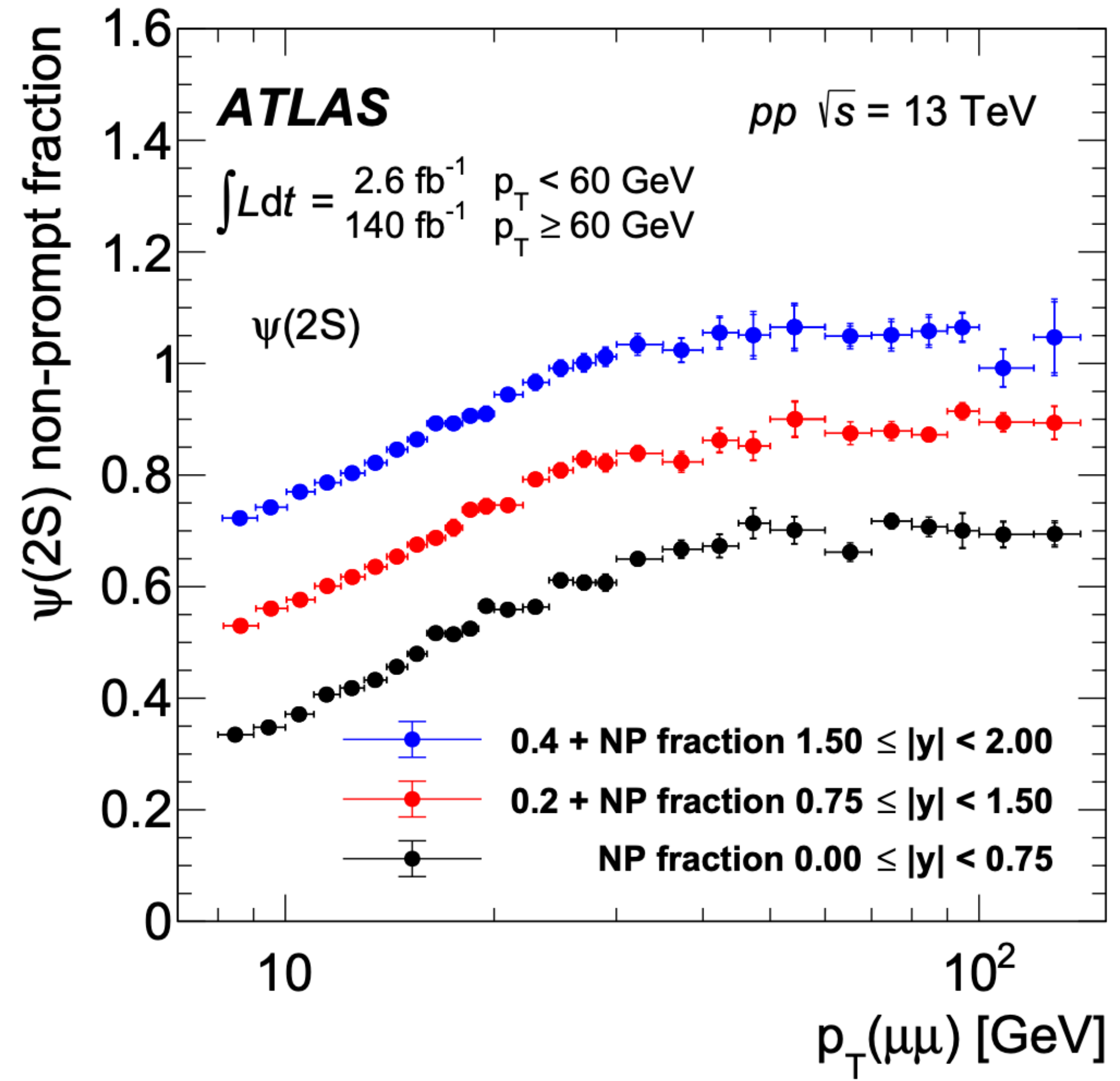
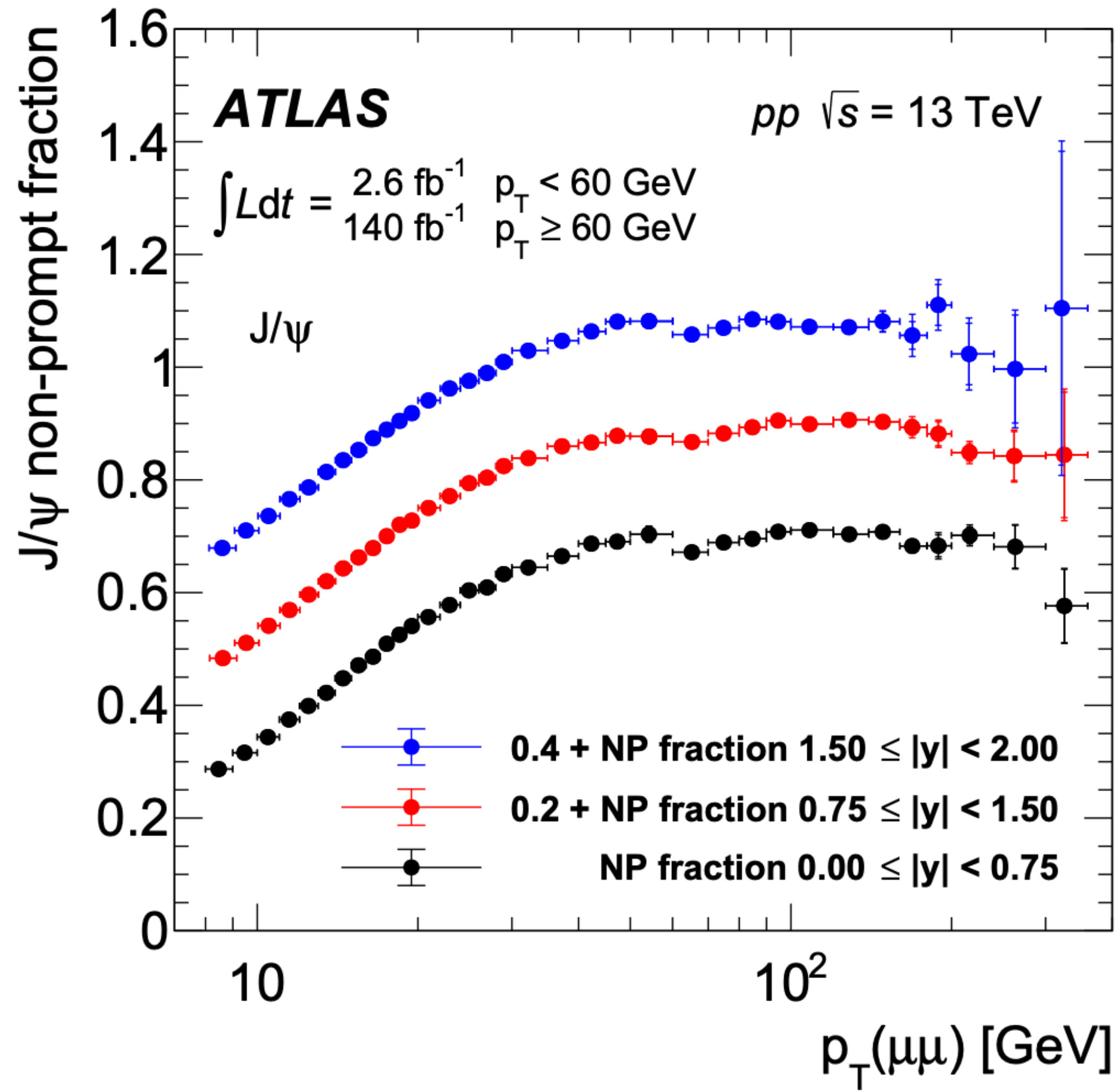


# Results: $\psi(2S)$ -to- $J/\psi$ Production Ratios



# Results: Non-Prompt Production Fractions

2309.17177



# Search for Di-Charmonium Resonances in $4\mu$ Final States

arXiv:2304.08962

Phys. Rev. Lett. 131 (2023) 151902

# Motivation

## □ More broadly:

- Color-confinement allows for exotic bound states of quarks other than baryons and meson:  $qq\bar{q}\bar{q}$ ,  $qqqq\bar{q}$ .

- BSM also predicts resonances in di-quarkonia spectrum.

## □ LHCb observed a narrow $X(6900)$ structure in $m(J/\psi J/\psi)$ in 2020!

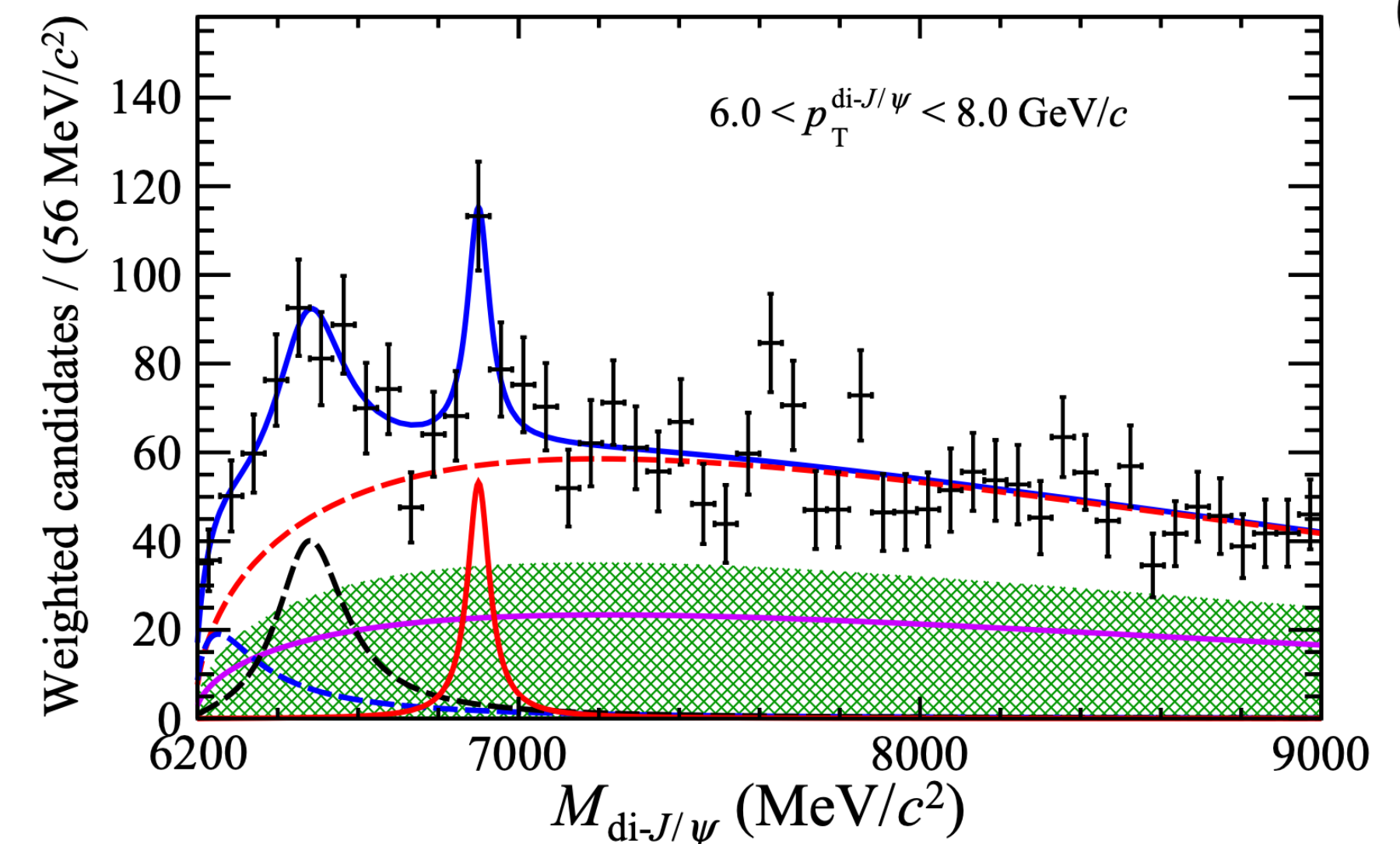
2006.16957

- Consistent with a charming tetraquark  $T_{cc\bar{c}\bar{c}}$ .

## □ This analysis:

- Corroborate the LHCb discovery in **a quite different phase space.**
- Make sense of the **additional enhancement near di- $J/\psi$  mass threshold** that LHCb observed.
- Search for di-charmonium excesses in  $J/\psi + \psi(2S)$  **channel!**

**LHCb** 2006.16957





# Analysis Strategy

## □ Data

- Full Run2 dataset with  $L = 140 \text{ fb}^{-1}$  of  $pp$  collision data at  $\sqrt{s} = 13 \text{ TeV}$
- Trigger on low- $p_T$  dimuon or trimuon triggers requiring an oppositely charged muon pair.

## □ Signal: Search for $4\mu$ final state produced via

- $J/\psi + J/\psi$  Channel
- $J/\psi + \psi(2S)$  Channel

## □ Background:

- **Prompt Di-Charmonia: SPS and DPS events.**  
Modeled with MC corrected with data in mass sidebands.
- **Non-Prompt Di-Charmonia:  $b$ -Hadron Decays:  $b\bar{b} \rightarrow J/\psi(\rightarrow \mu^+\mu^-)J/\psi(\rightarrow \mu^+\mu^-)X$**   
Modeled with MC corrected with data in sidebands of vertexing quality and/or displacement of  $\mu^+\mu^-$  vertices.
- **One Charmonium + Non-Resonant  $\mu^+\mu^-$ : Mostly due to fake muons.**  
Purely data-driven fake estimation.

## □ Event Selection

- $4\mu$  vertexing followed by two  $\mu^+\mu^-$  vertexings with  $J/\psi$  or  $\psi(2S)$  mass-constraints in events with two pairs of  $\mu^+\mu^-$ .
- Background rejection via stringent cuts on the quality of the  $4\mu$  vertex fit and on the angle b/w the charmonia candidates.

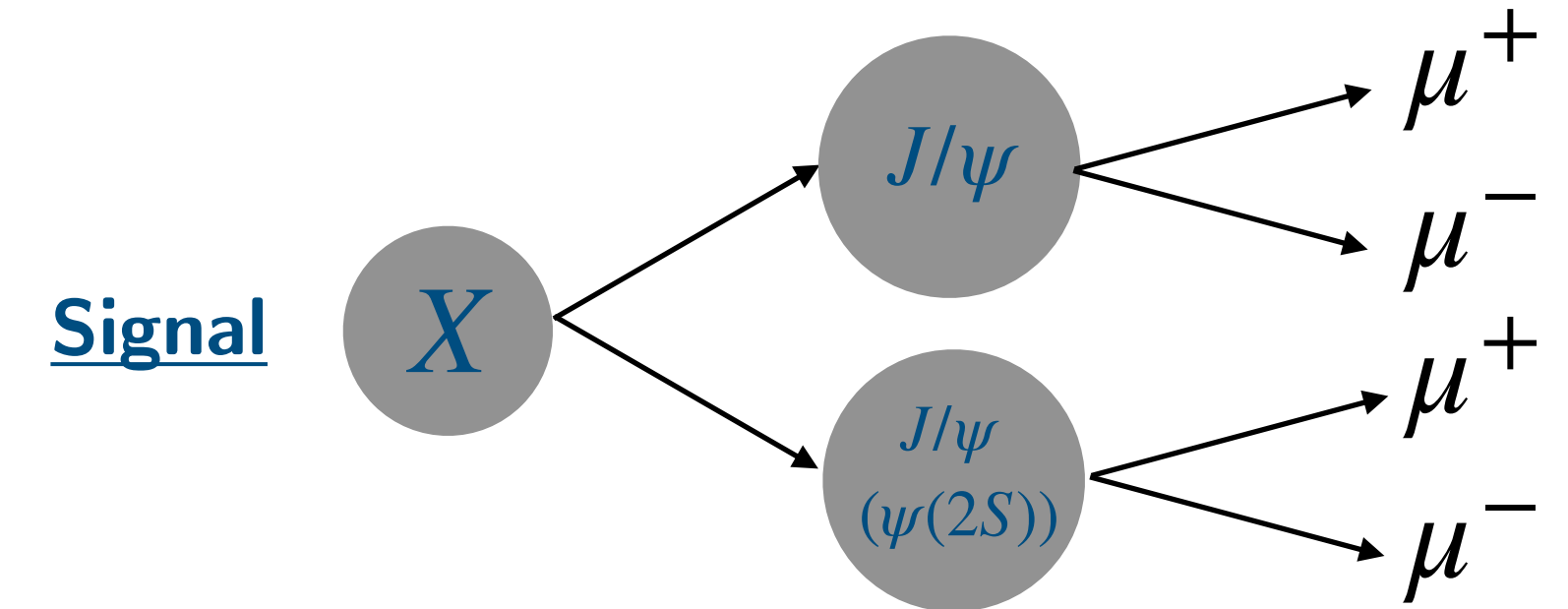
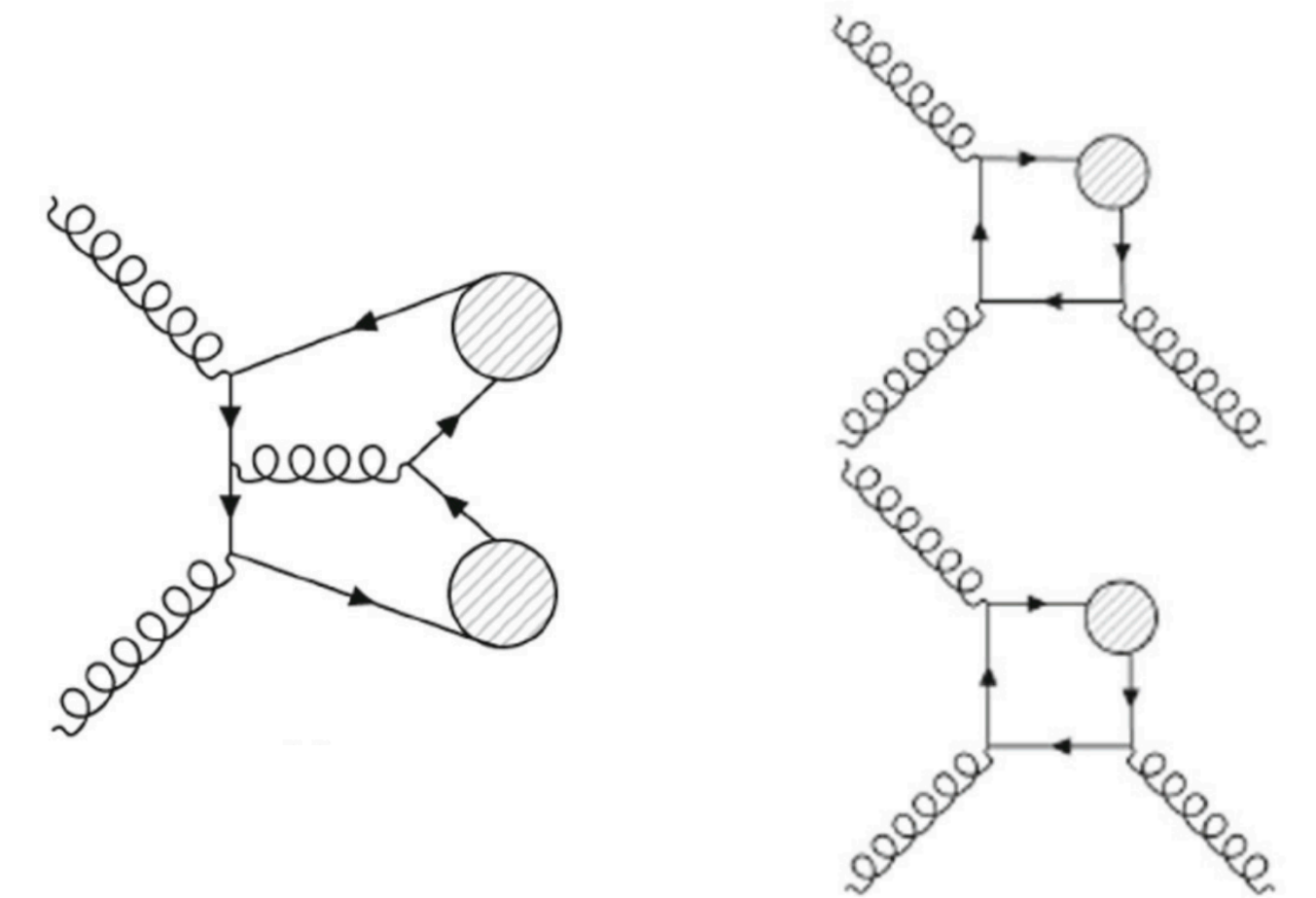


Illustration Credit: Yue Xu ([LHCP2025](#))



SPS Background

DPS Background



# Signal Modeling for the Fits

## □ Signal Modeling for Fits

### □ $J/\psi + J/\psi$ Channel

- Feiddown from  $J/\psi + \psi(2S)$  channel signal fit taken as background

- **Model A:** 3 interfering BW resonances

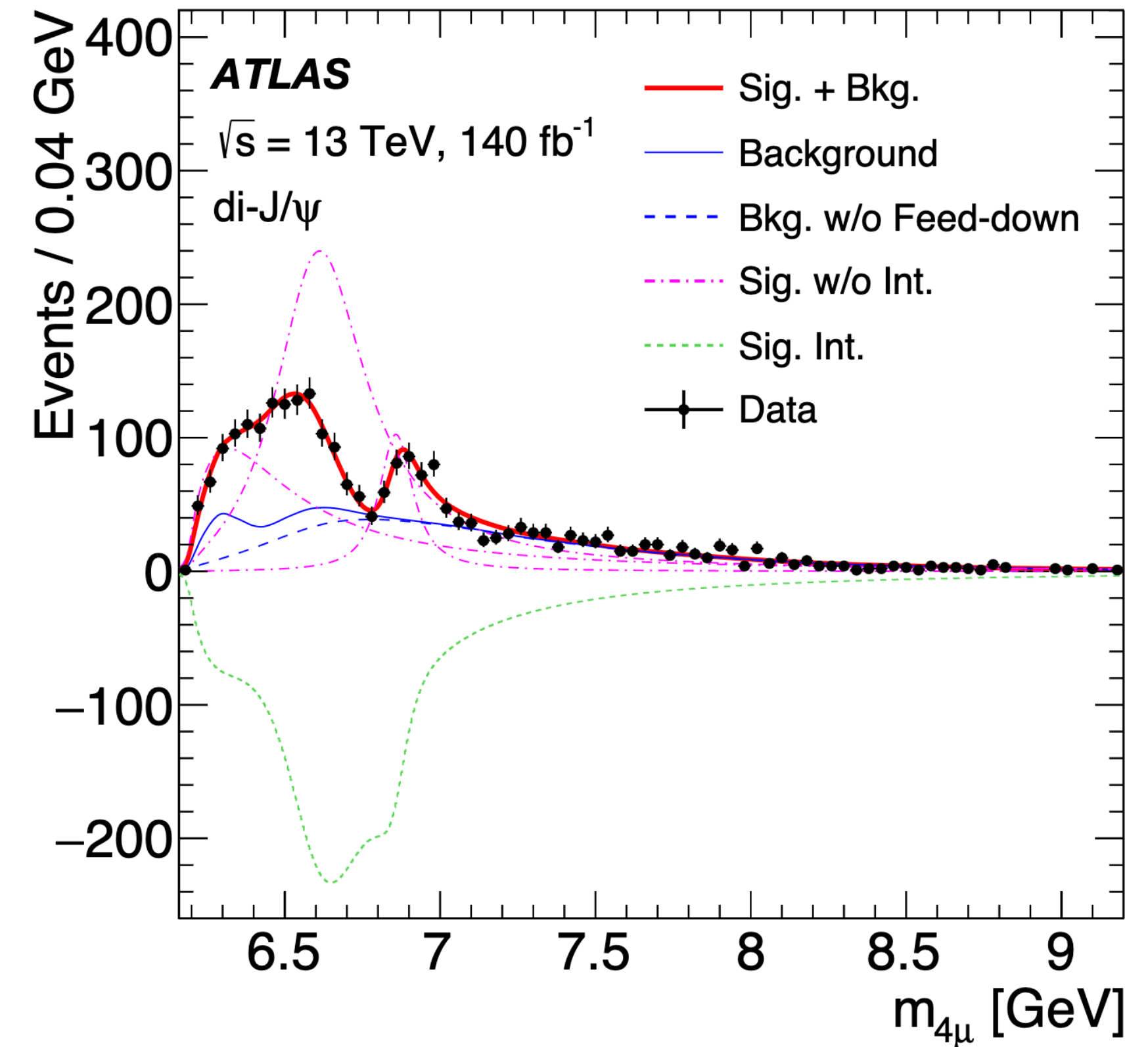
- **Model B:** 2 BW where the lower resonance interferes with the SPS background

### □ $J/\psi + \psi(2S)$ Channel

- **Model  $\alpha$ :** 3 interfering BW resonances of Model A + standalone 4th BW resonance

- **Model  $\beta$ :** Only 1 BW resonance

- Other models considered but excluded in favor of these with the most promising of the excluded ones used to calculate systematics associated with the models.



# Results: $J/\psi + J/\psi$ Channel

- Significance for both models far exceeds  $5\sigma$

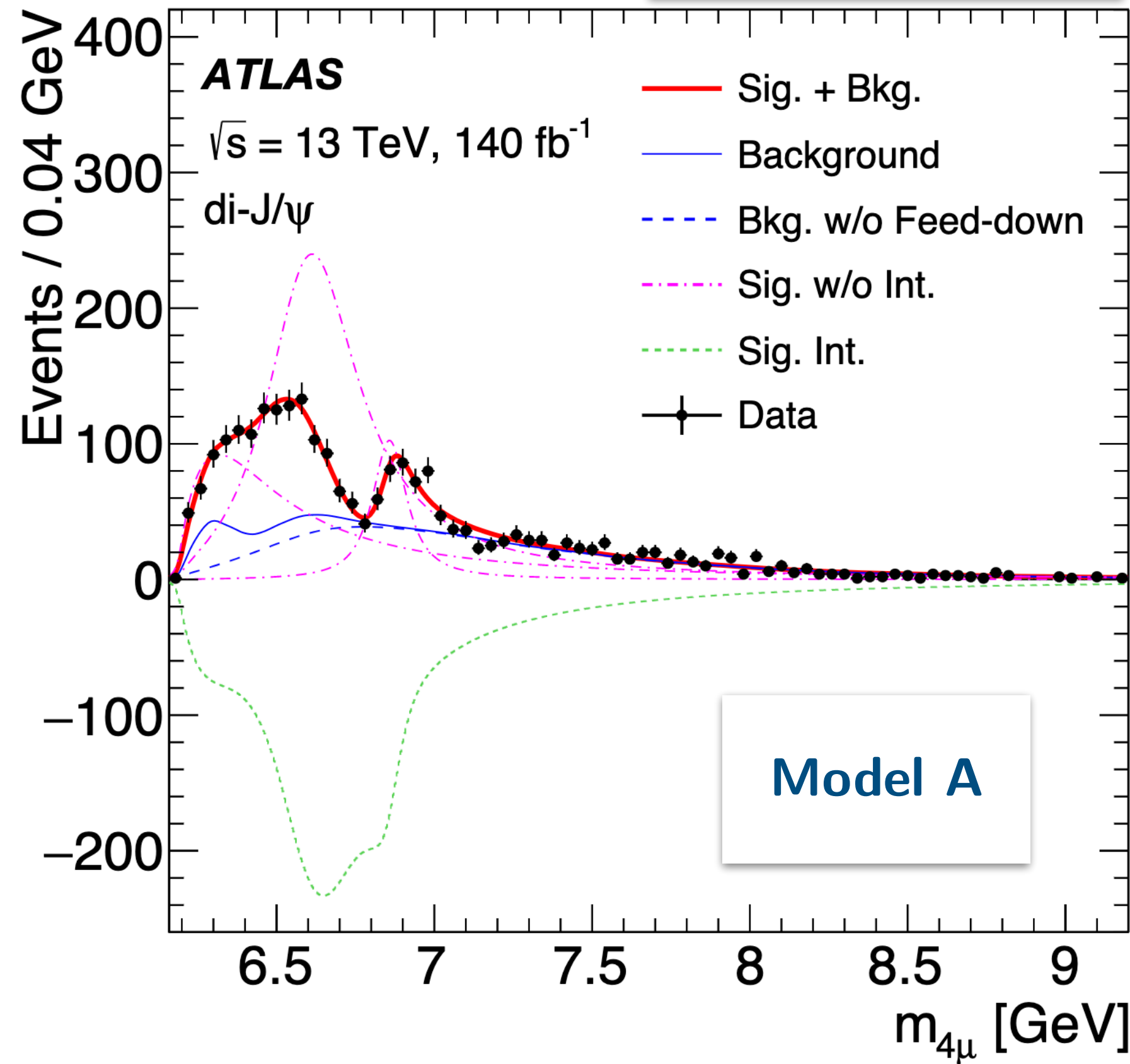
- The mass of the  $m_2$  resonance consistent with the LHCb mass as well as with the CMS search now: [2306.07164](#)

- The broad structure at lower mass can still be from other effects such as feeddown from higher dicharmonium resonances.

$$T_{cc\bar{c}\bar{c}} \rightarrow \chi_{cJ'} \chi_{cJ'} \rightarrow J/\psi J/\psi \gamma \gamma$$

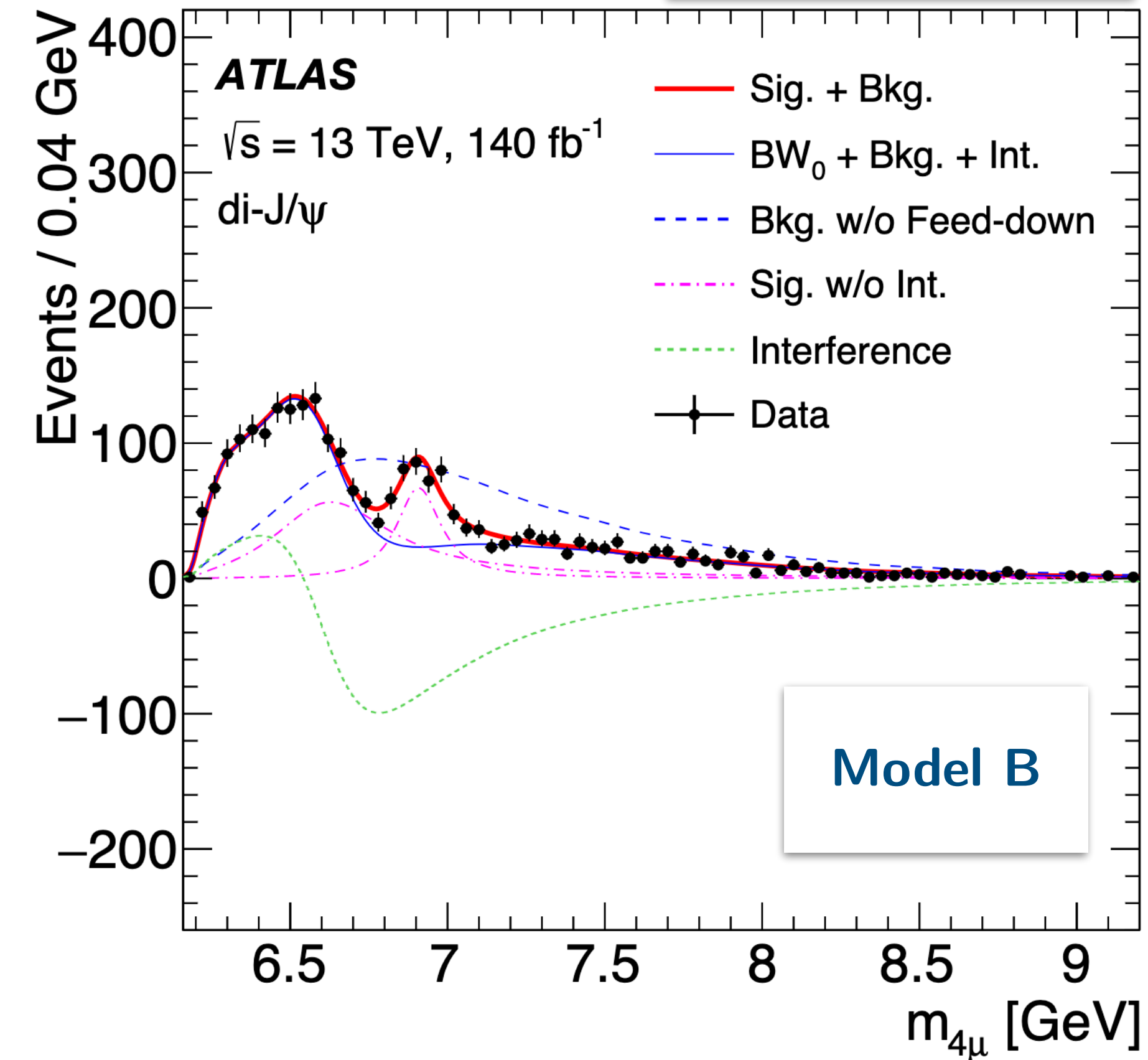
All mutually interfering

di- $J/\psi$	model A
$m_0$	$6.41 \pm 0.08^{+0.08}_{-0.03}$
$\Gamma_0$	$0.59 \pm 0.35^{+0.12}_{-0.20}$
$m_1$	$6.63 \pm 0.05^{+0.08}_{-0.01}$
$\Gamma_1$	$0.35 \pm 0.11^{+0.11}_{-0.04}$
$m_2$	$6.86 \pm 0.03^{+0.01}_{-0.02}$
$\Gamma_2$	$0.11 \pm 0.05^{+0.02}_{-0.01}$
$\Delta s/s$	$\pm 5.1\%^{+8.1\%}_{-8.9\%}$



Interfering with SPS background

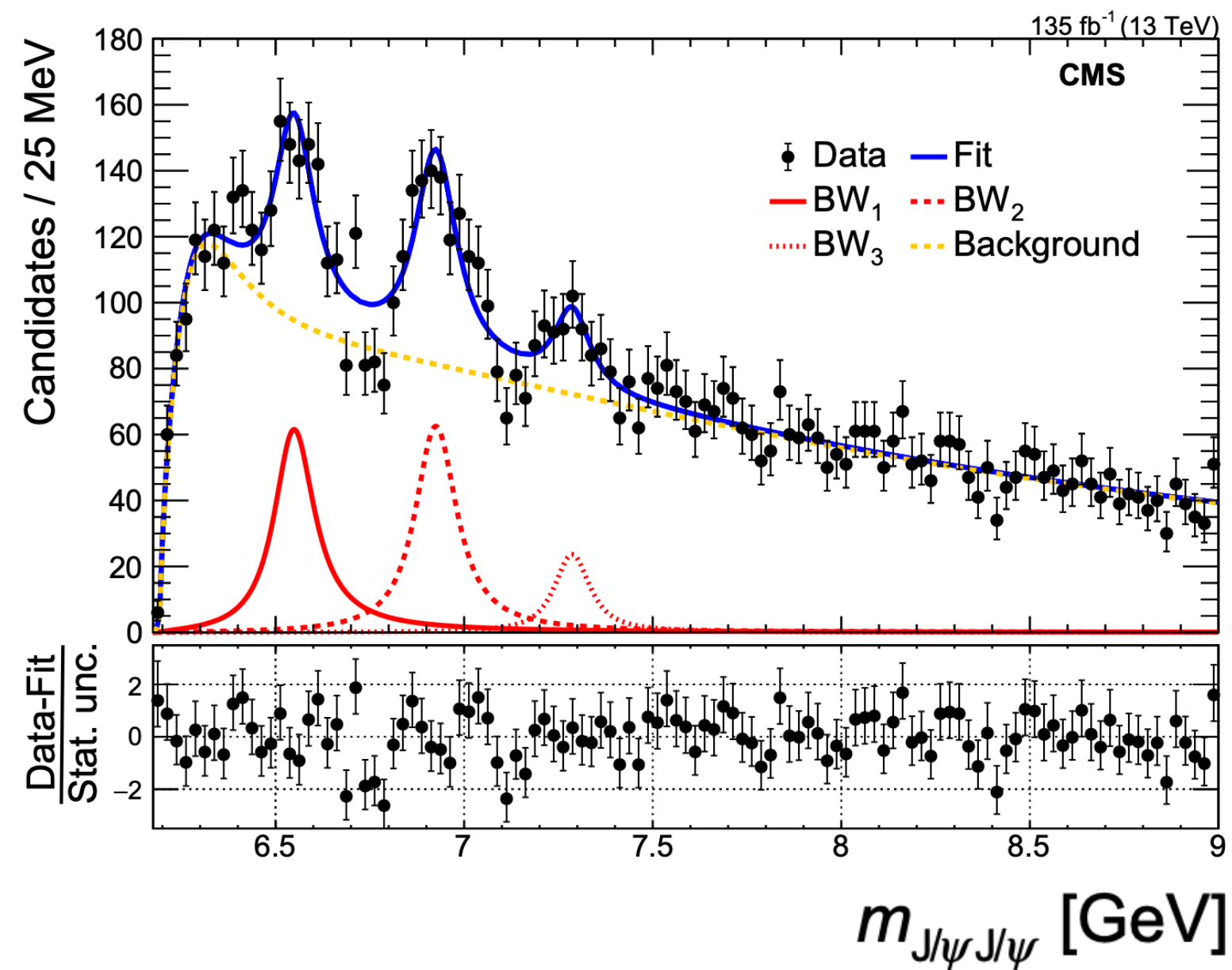
di- $J/\psi$	model B
$m_0$	$6.65 \pm 0.02^{+0.03}_{-0.02}$
$\Gamma_0$	$0.44 \pm 0.05^{+0.06}_{-0.05}$
$m_1$	—
$\Gamma_1$	—
$m_2$	$6.91 \pm 0.01 \pm 0.01$
$\Gamma_2$	$0.15 \pm 0.03 \pm 0.01$
$\Delta s/s$	—



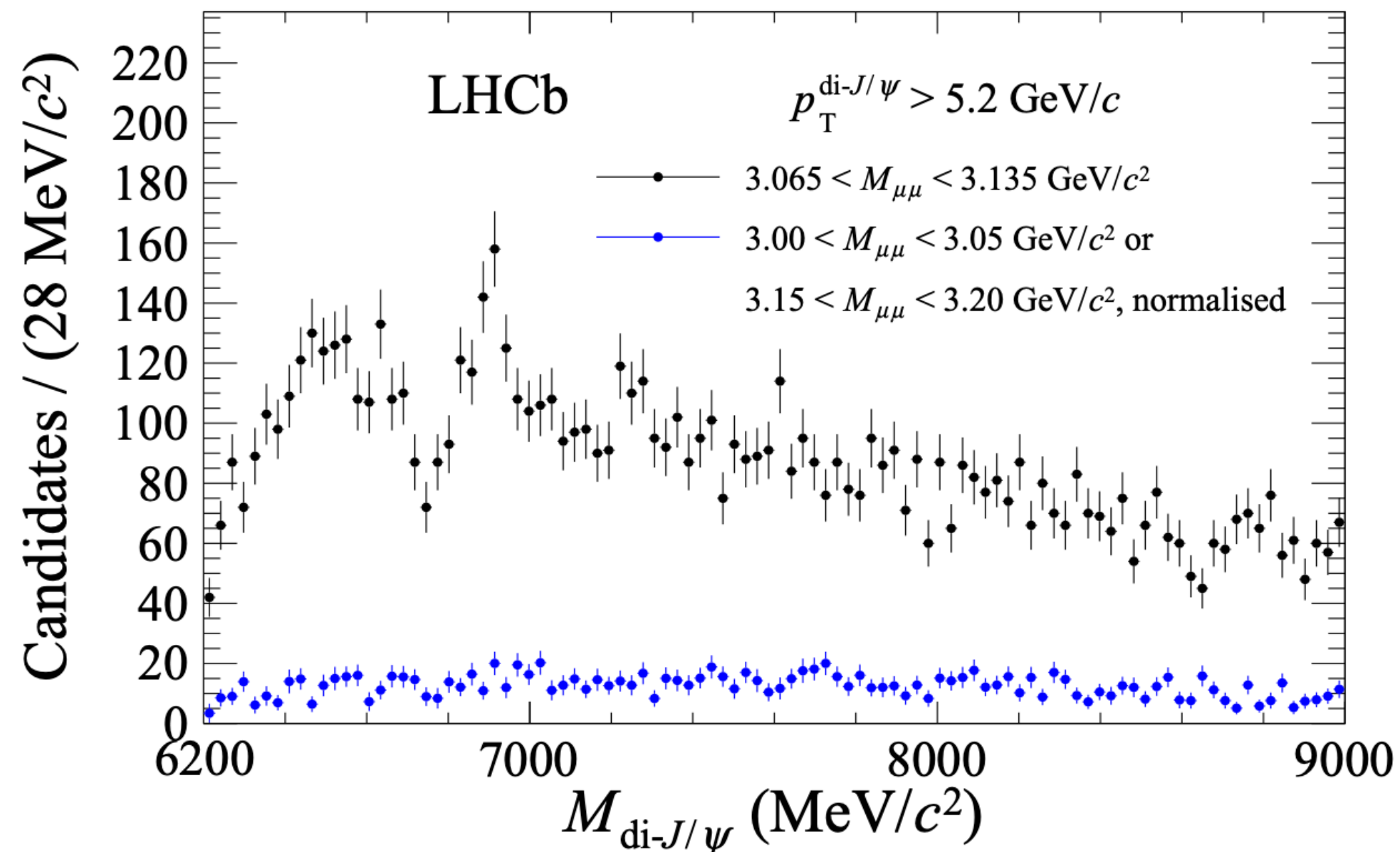
# Brief Recap of Context for $J/\psi + \psi(2S)$ Channel

- $X(6900)$  is just above  $J/\psi + \psi(2S)$  mass threshold!
- So, we might expect an excess at the lower-end of  $J/\psi + \psi(2S)$  spectrum.
- Is there something going on also at  $\sim 7.2$  GeV in the dicharmonia spectrum?
- Technical: To fit the feed-down background for  $J/\psi + J/\psi$  channel.

CMS 2306.07164



LHCb 2006.16957





# Results: $J/\psi + \psi(2S)$ Channel

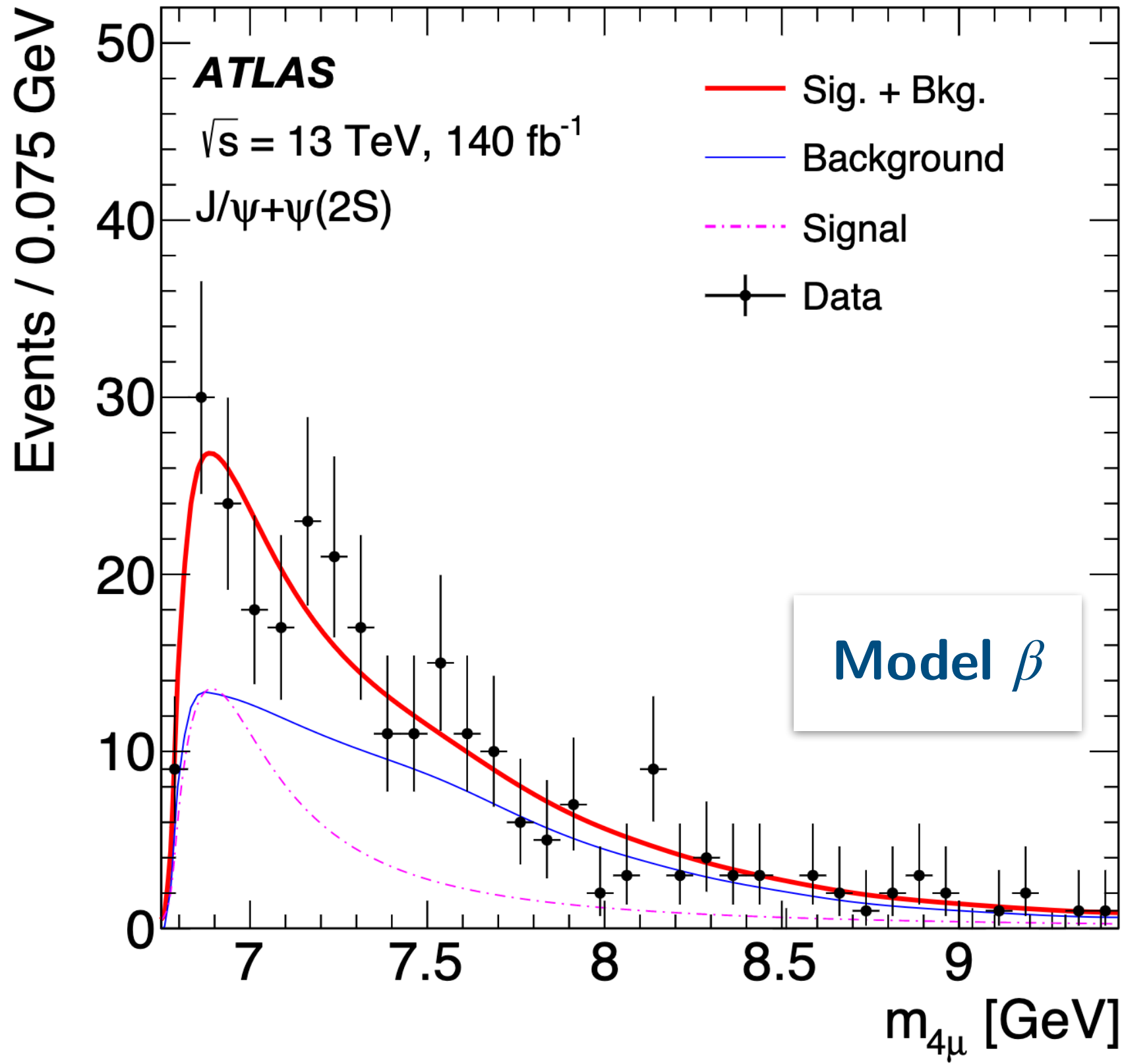
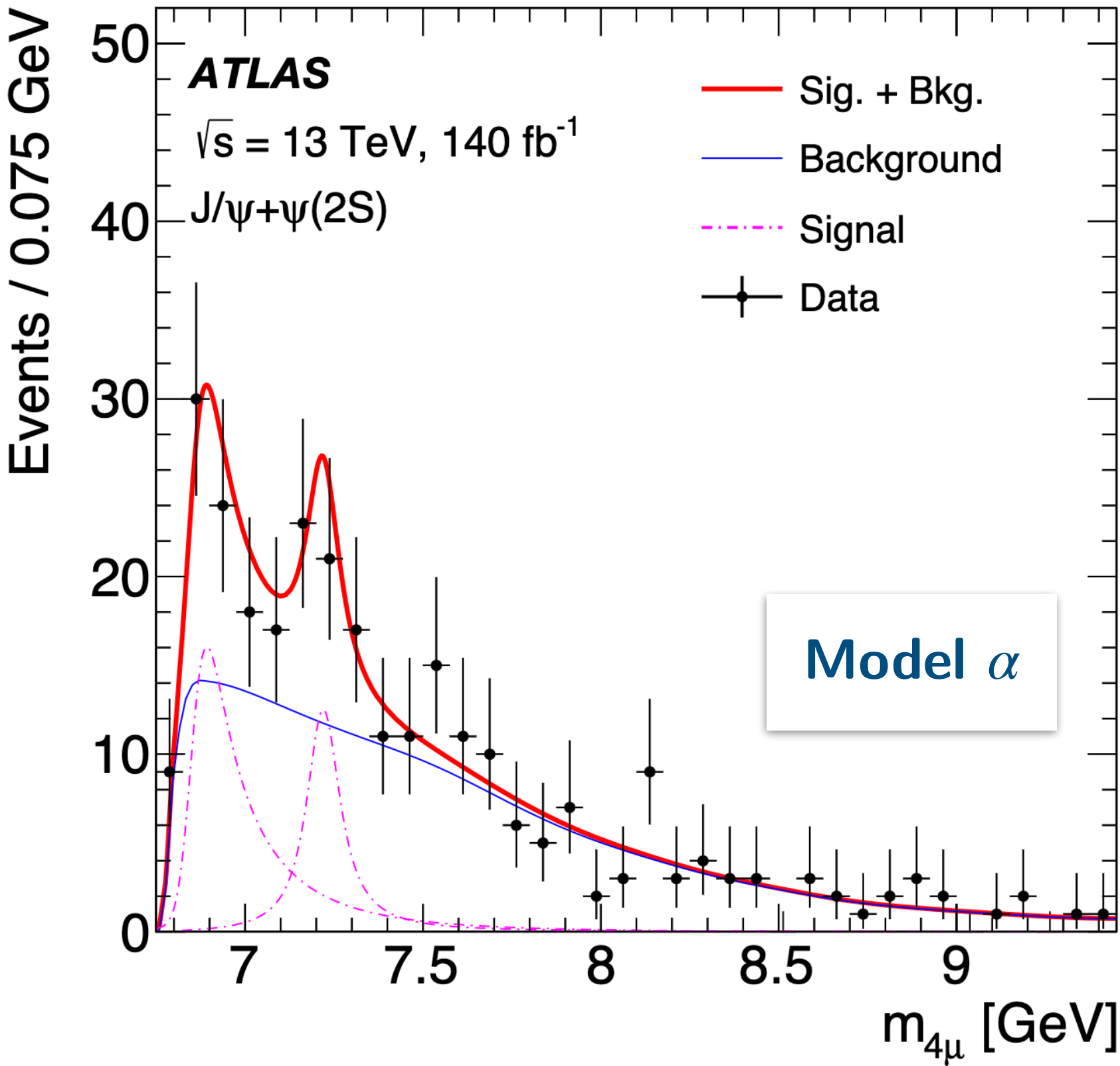
- Model  $\alpha$ :  $4.7\sigma$ 
  - Significance for the second resonance alone at  $3\sigma$ .
- Model  $\beta$ :  $4.3\sigma$

$J/\psi+\psi(2S)$	model $\alpha$
$m_3$	$7.22 \pm 0.03^{+0.01}_{-0.04}$
$\Gamma_3$	$0.09 \pm 0.06^{+0.06}_{-0.05}$
$\Delta s/s$	$\pm 21\%^{+25\%}_{-15\%}$

$J/\psi+\psi(2S)$	model $\beta$
$m_3$	$6.96 \pm 0.05 \pm 0.03$
$\Gamma_3$	$0.51 \pm 0.17^{+0.11}_{-0.10}$
$\Delta s/s$	$\pm 20\% \pm 12\%$

All mutually interfering

di- $J/\psi$	model A
$m_0$	$6.41 \pm 0.08^{+0.08}_{-0.03}$
$\Gamma_0$	$0.59 \pm 0.35^{+0.12}_{-0.20}$
$m_1$	$6.63 \pm 0.05^{+0.08}_{-0.01}$
$\Gamma_1$	$0.35 \pm 0.11^{+0.11}_{-0.04}$
$m_2$	$6.86 \pm 0.03^{+0.01}_{-0.02}$
$\Gamma_2$	$0.11 \pm 0.05^{+0.02}_{-0.01}$
$\Delta s/s$	$\pm 5.1\%^{+8.1\%}_{-8.9\%}$



# Outlook: Heavy Quarkonium Production and Spectroscopy ATLAS in This Talk and Beyond!

- Differential Production Cross-Section of  $J/\psi$  and  $\psi(2S)$ 
  - Extended up to  $p_T^{J/\psi} = 360$  GeV. A lot to work with for theoretical models!
- Search for Di-Charmonium Resonances in  $4\mu$  Final States
  - Corroborated the LHCb discovery of  $X(6900)$  in di- $J/\psi$ .
  - Look out for more work on the evidence of excess in  $J/\psi + \psi(2S)$ .
  - More work also needed to understand the low-mass broad structure observed, also by LHCb.

ATLAS Public Results



**Thank You!**

# Back-Up

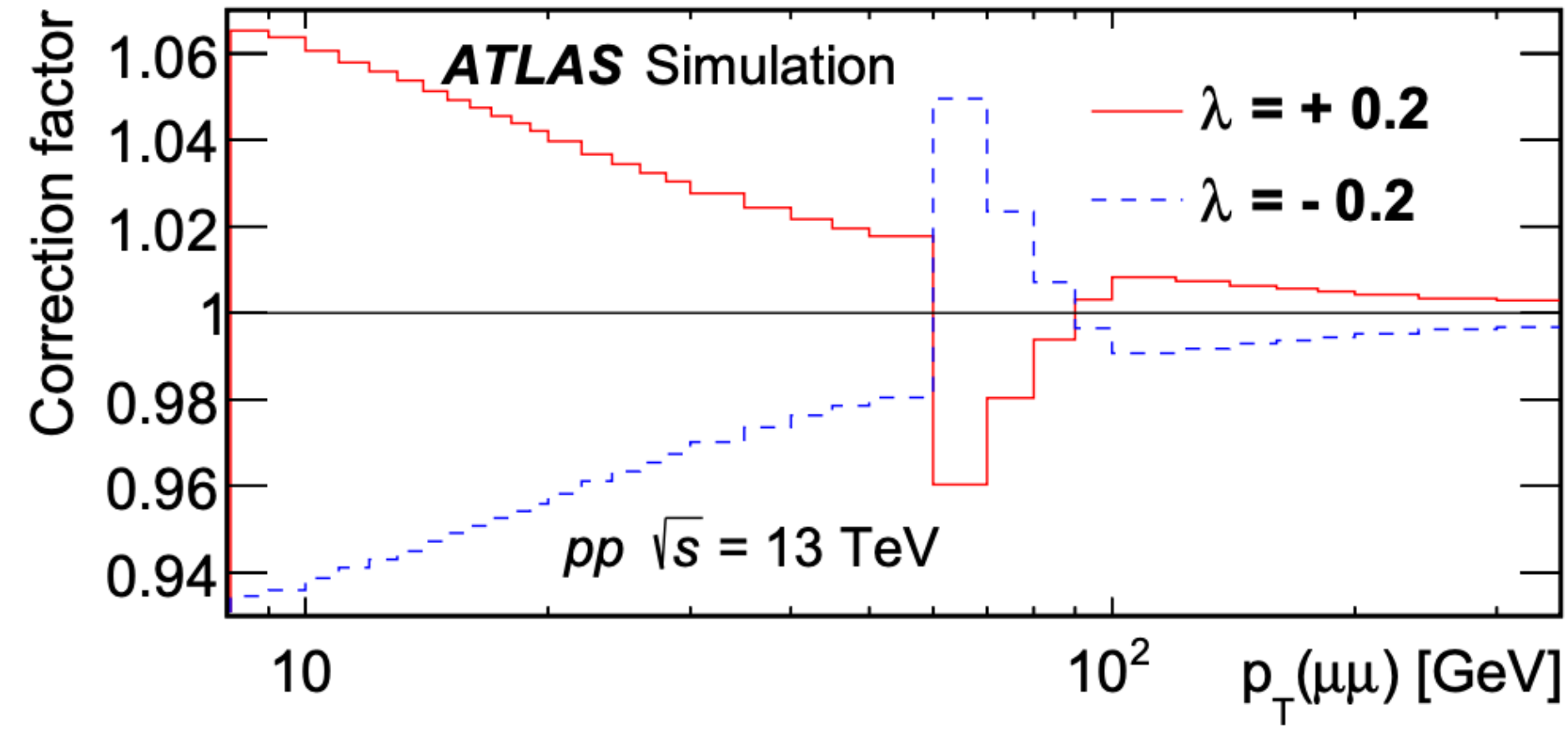
# Differential Production Cross-Section of $J/\psi$ and $\psi(2S)$

arXiv:2309.17177

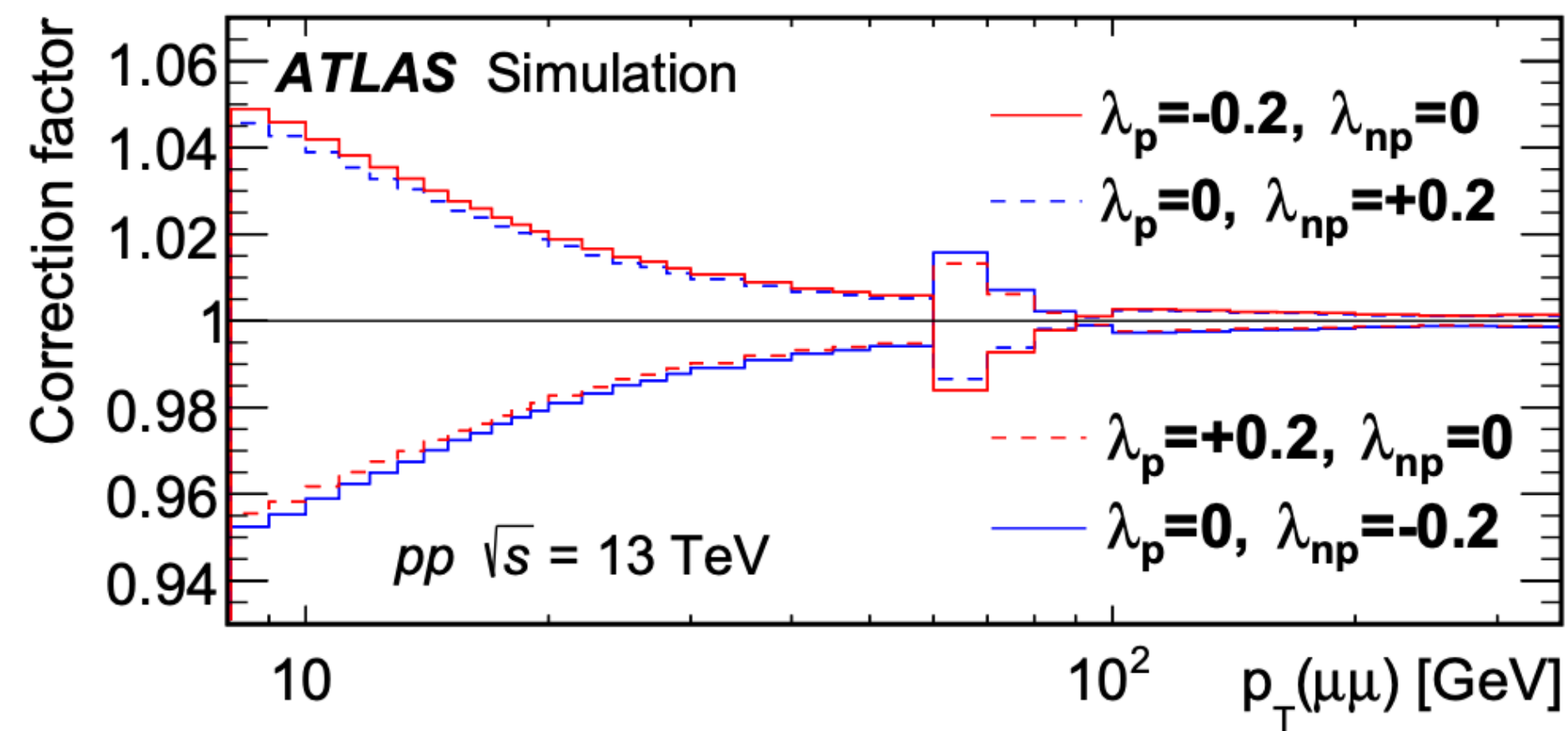
Eur. Phys. J. C 84 (2024) 169

# The Spin-Alignment Correction Factors

The switch in trigger strategy kicks in...  
at 60 GeV

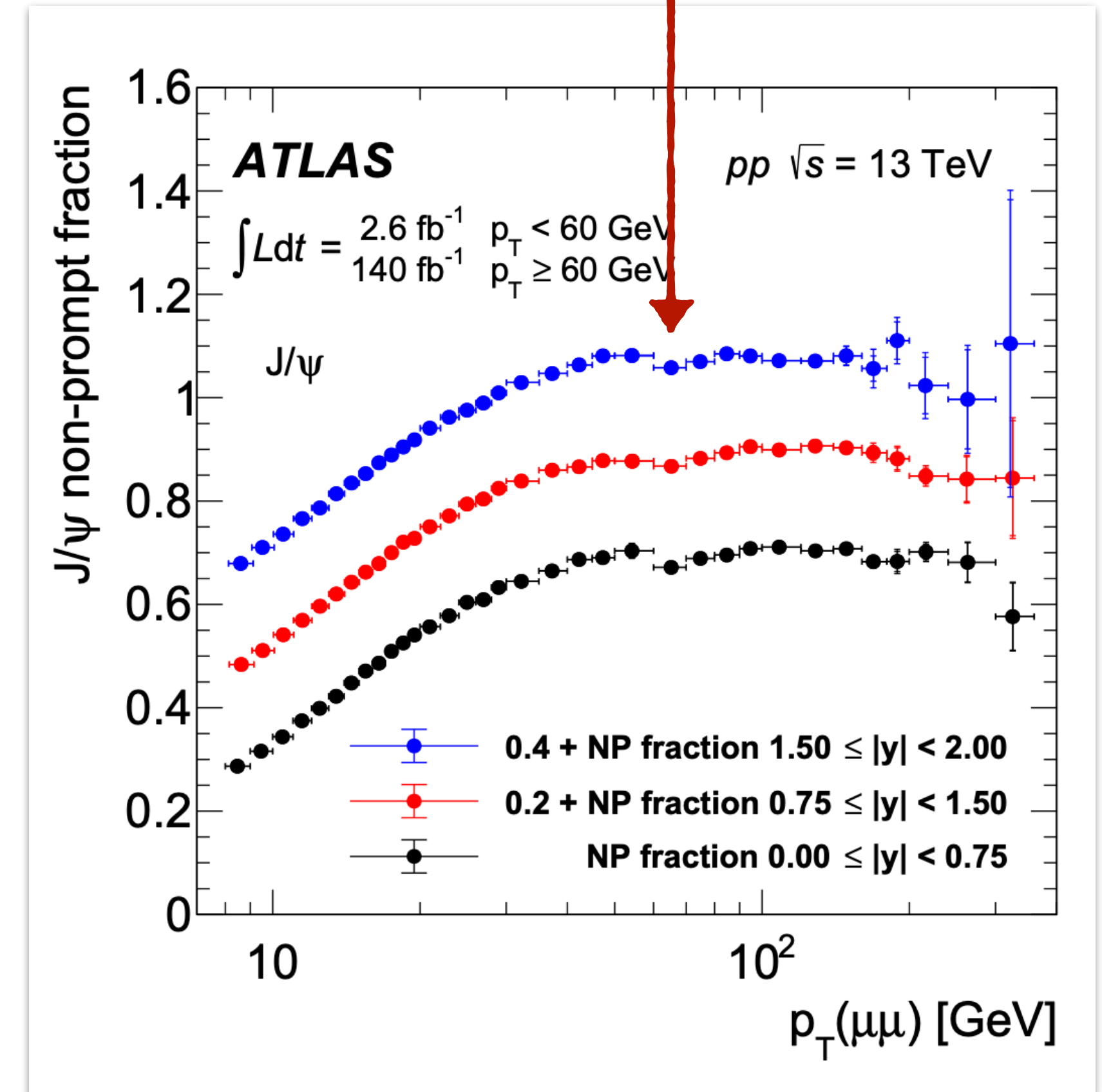


Prompt  $J/\psi$



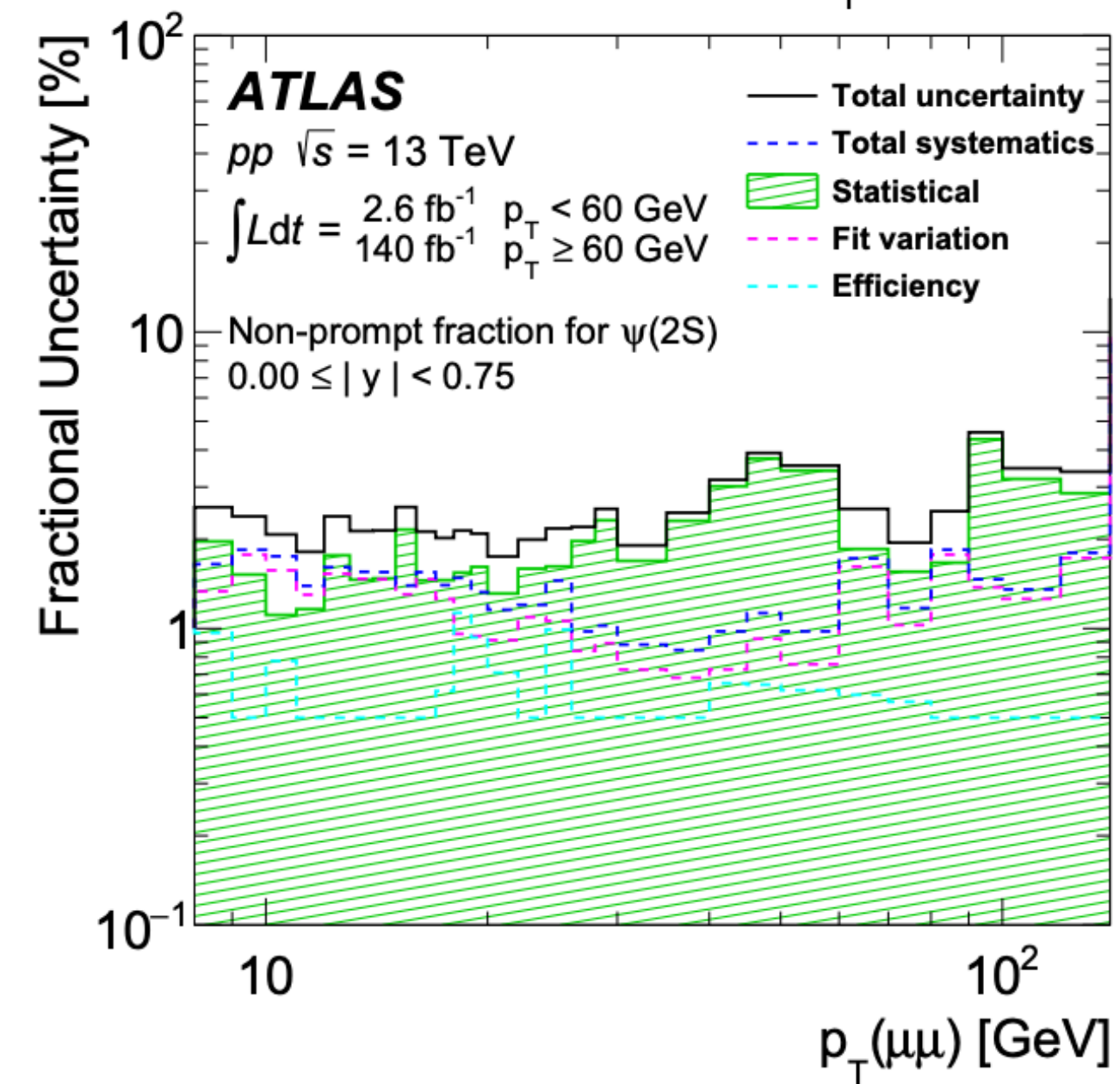
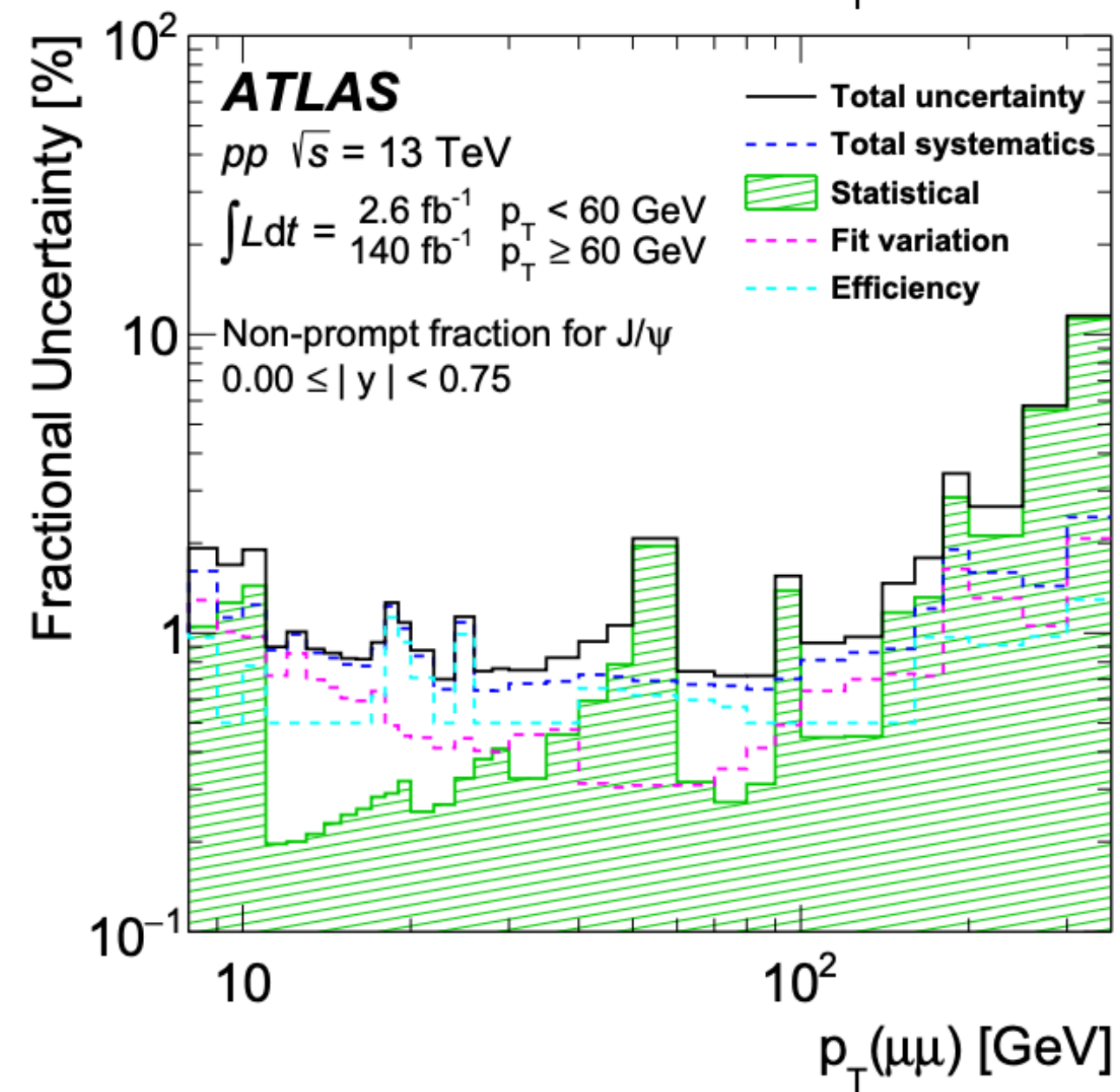
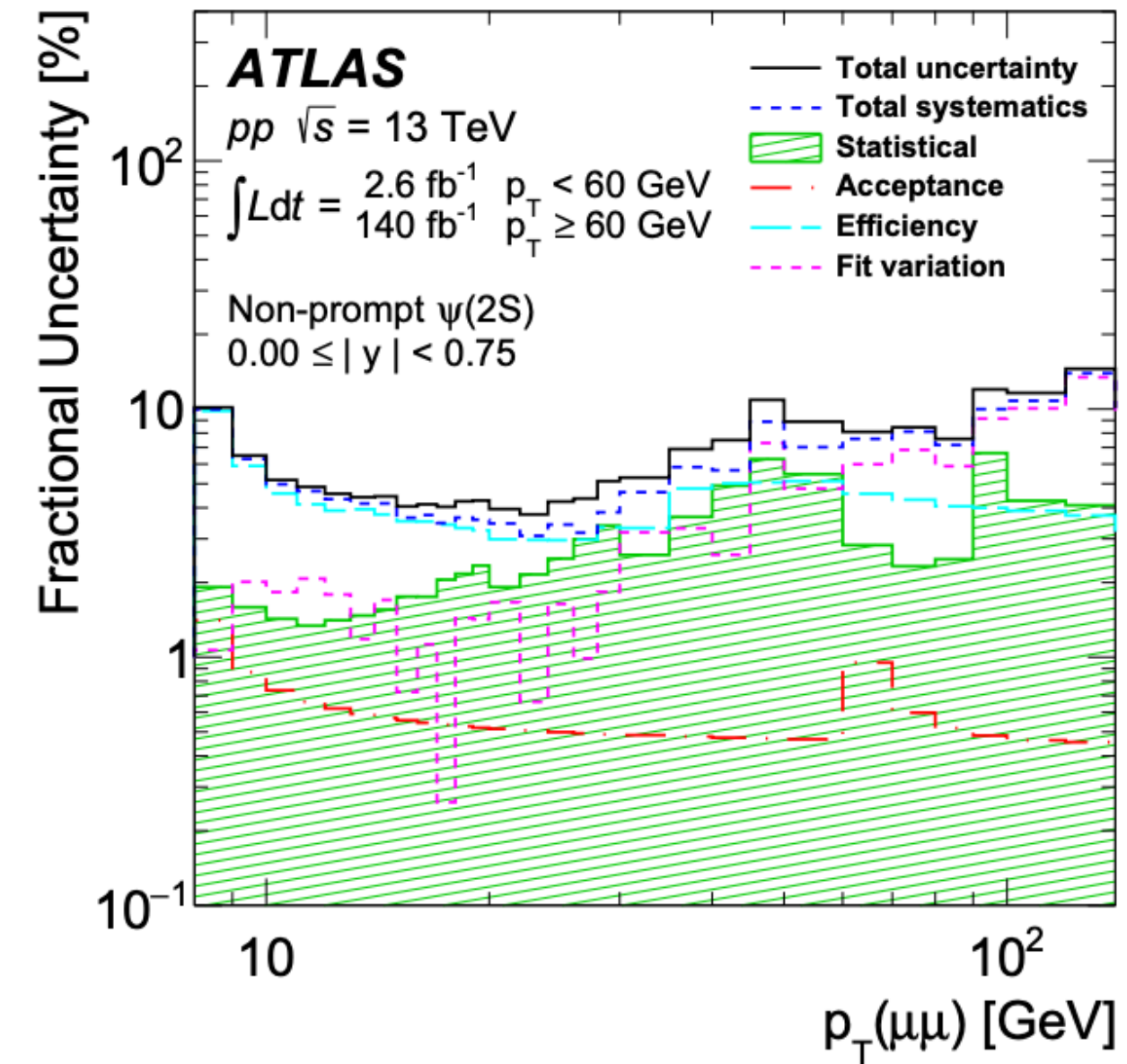
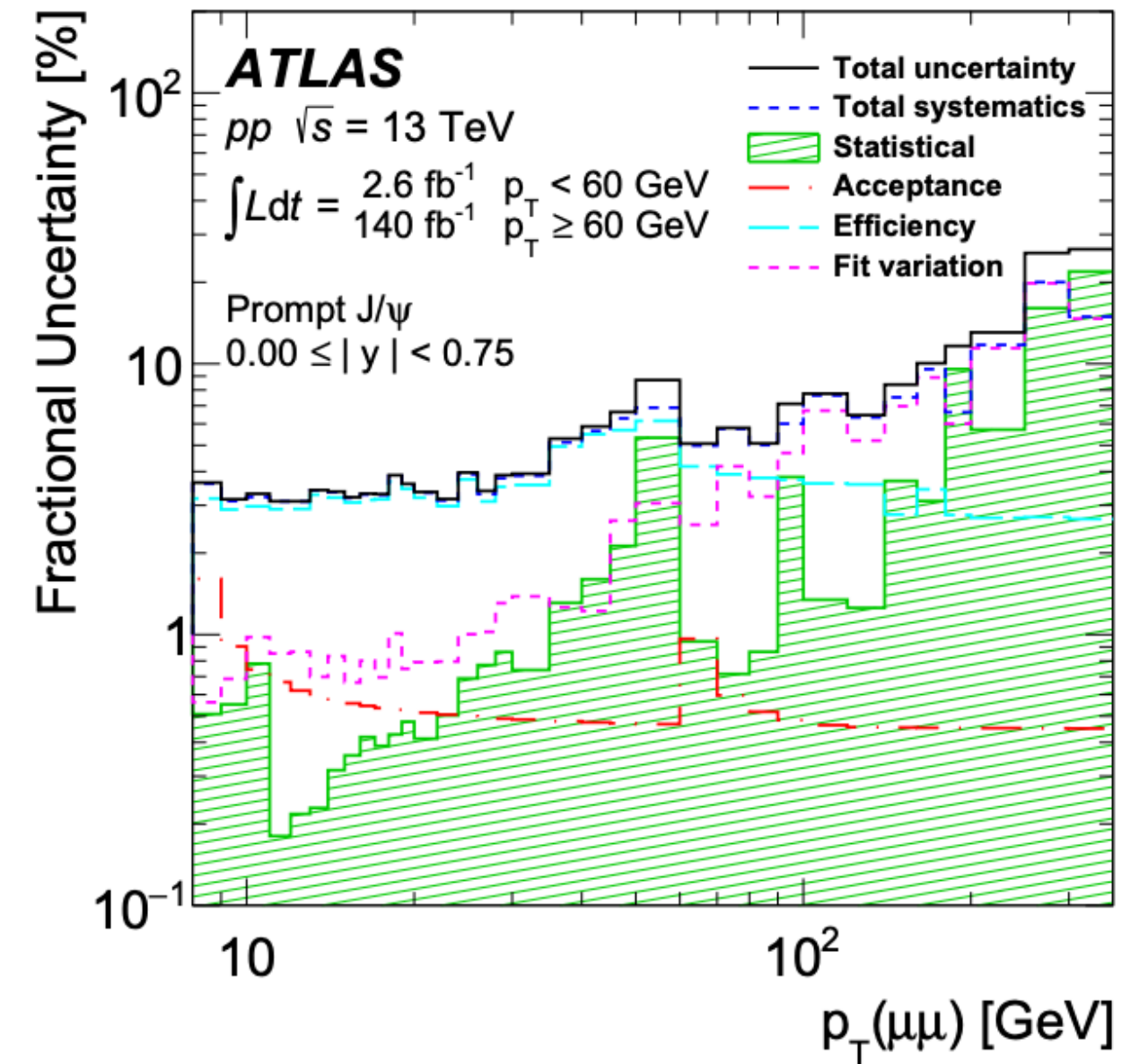
Non-Prompt  $J/\psi$

Most dominant corrections to acceptance under the isotropic assumption, based on the best available data and theory, are taken as systematics.



# Systematics

2309.17177





# Fit Model

Each Signal/Background Component

$$P_i(m_{\mu\mu}, \tau) = f_i(m_{\mu\mu}) \cdot [h_i(\tau) \otimes R(\tau)]$$

$$F(m_{\mu\mu}, \tau) = \sum_{i=1}^7 \kappa_i P_i(m_{\mu\mu}, \tau)$$

Total PDF

$$R(\tau) = \omega_A G_A(\tau) + \omega_B G_B(\tau) + (1 - \omega_A - \omega_B) G_C(\tau)$$

Lifetime Resolution

† Some correlation between mass and lifetime resolution is modeled (next slide.)

i	Type	P/NP	$f_i(m)$	$h_i(\tau)$
1†	$J/\psi$	P	$\omega_0 G_1(m) + (1 - \omega_0)[\omega_1 CB(m) + (1 - \omega_1)G_2(m)]$	$\delta(\tau)$
2	$J/\psi$	NP	$\omega_0 G_1(m) + (1 - \omega_0)[\omega_1 CB(m) + (1 - \omega_1)G_2(m)]$	$\omega_2 E_1(\tau) + (1 - \omega_2)E_1(b\tau)$
3	$\psi(2S)$	P	$\omega_0 G_1(\beta m) + (1 - \omega_0)[\omega_1 CB(\beta m) + (1 - \omega_1)G_2(\beta m)]$	$\delta(\tau)$
4	$\psi(2S)$	NP	$\omega_0 G_1(\beta m) + (1 - \omega_0)[\omega_1 CB(\beta m) + (1 - \omega_1)G_2(\beta m)]$	$E_2(\tau)$
5	Bkg	P	$P$	$\delta(\tau)$
6	Bkg	NP	$E_3(m)$	$E_4(\tau)$
7	Bkg	NP	$E_5(m)$	$E_6( \tau )$

# Fit Model

Each Signal/Background Component

$$P_i(m_{\mu\mu}, \tau) = f_i(m_{\mu\mu}) \cdot [h_i(\tau) \otimes R(\tau)]$$

$$F(m_{\mu\mu}, \tau) = \sum_{i=1}^7 \kappa_i P_i(m_{\mu\mu}, \tau)$$

Total PDF

$$R(\tau) = \omega_A G_A(\tau) + \omega_B G_B(\tau) + (1 - \omega_A - \omega_B) G_C(\tau)$$

Lifetime Resolution

i	Type	P/NP	$f_i(m)$	$h_i(\tau)$
1†	$J/\psi$	P	$\omega_0 G_1(m) + (1 - \omega_0)[\omega_1 CB(m) + (1 - \omega_1)G_2(m)]$	$\delta(\tau)$

† In the  $i = 1$  term, describing the prompt  $J/\psi$  peak, the product of the narrowest Gaussian term in  $f_1(m_{\mu\mu})$  and the narrowest Gaussian term in  $R(\tau)$  was replaced by a bivariate Gaussian function in  $m_{\mu\mu}$  and  $\tau$  with a correlation coefficient  $\rho = 0.3$ ,

$$\omega_0 G_1(m_{\mu\mu}; \sigma_1) \cdot f_A G_A(\tau; \sigma_A) \mapsto \omega_0 f_A G_{BV}(m_{\mu\mu}, \tau; \sigma_1, \sigma_A, \rho), \quad (4)$$

to take into account the observed correlation between the measured values of these quantities. The effect of this correlation was found to be negligible for other terms.