

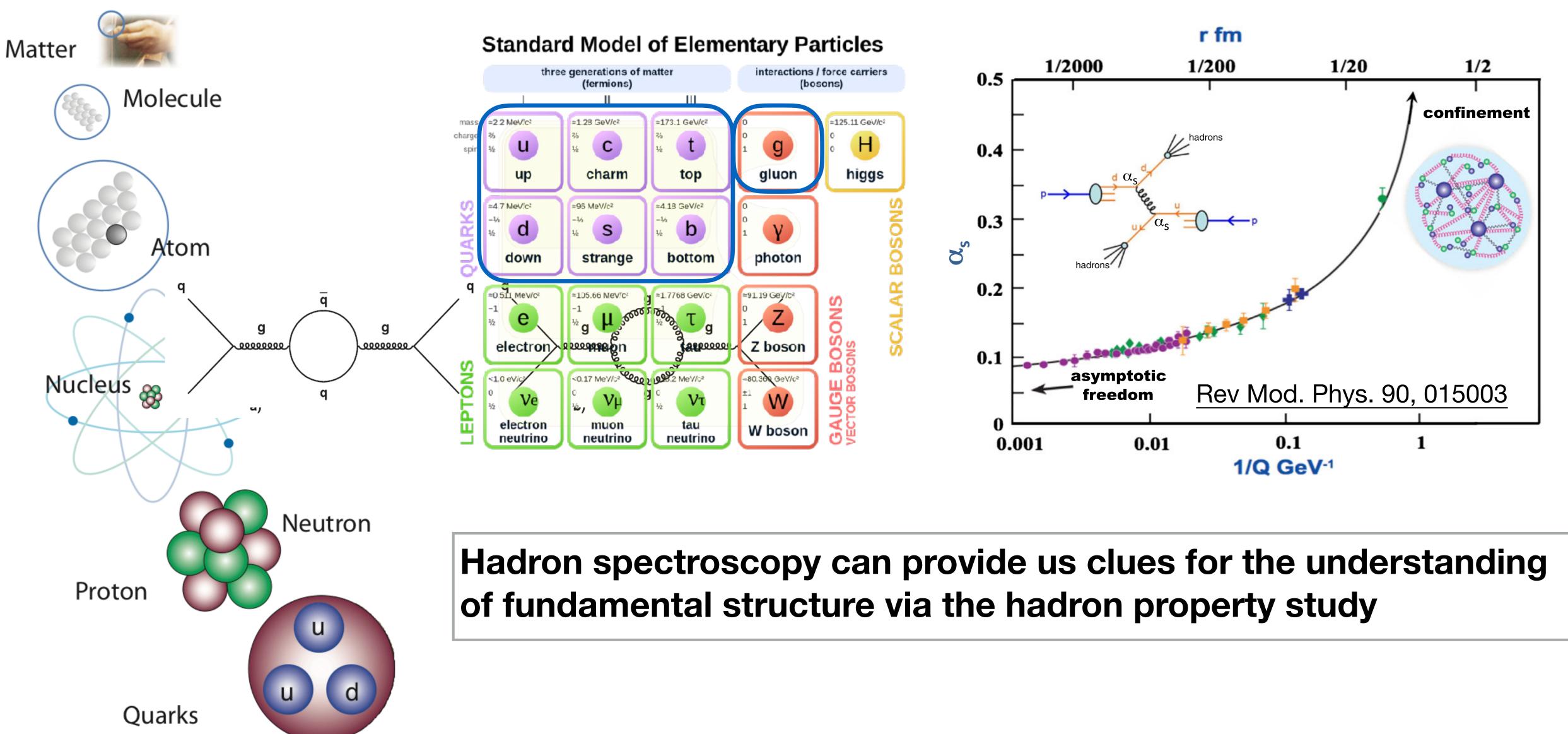
# **Recent Highlight on light hadrons at BESIII**

EPS-HEP CONFERENCE, 07-11 July, 2025, MARSEILLE, FRANCE

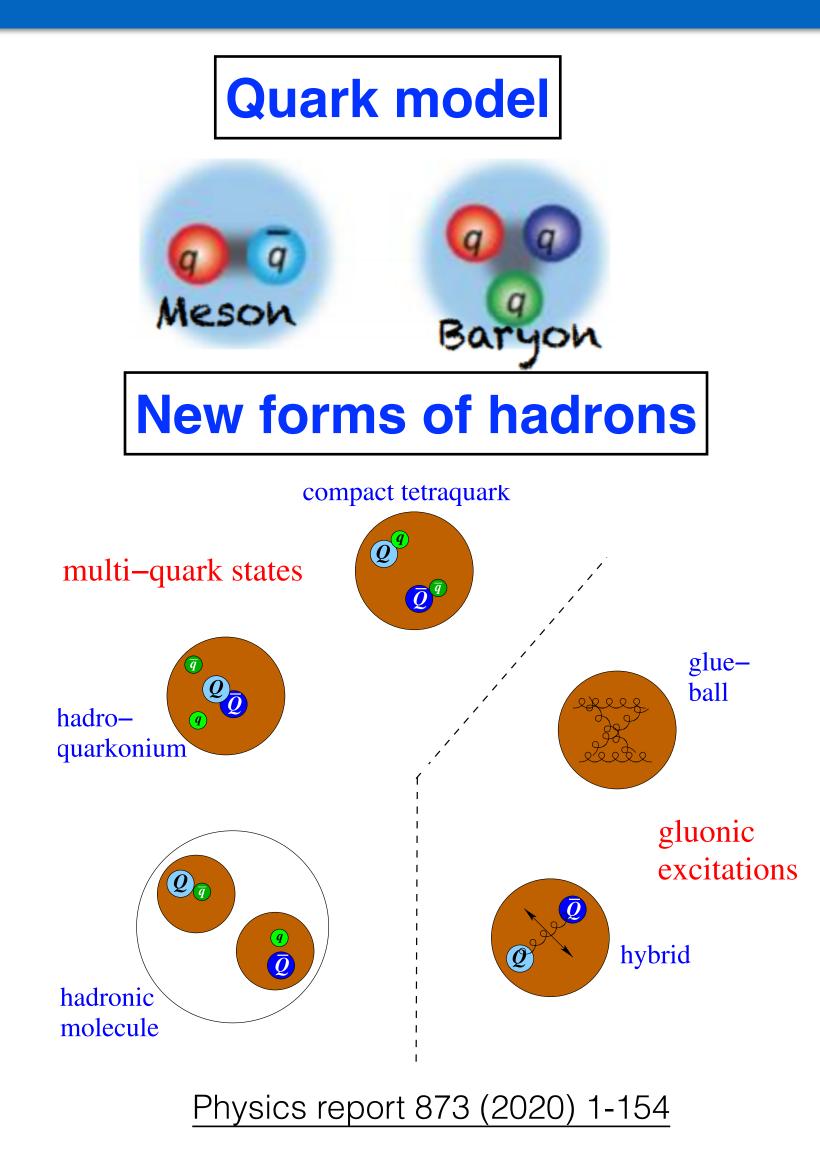
- **Yanping Huang** (On behalf of the BESIII Collaboration)
- Institute of High Energy Physics, CAS

# Fundamental Structure of Matters









#### Quark model (QM)

- Identify hadrons as compound objects consisting of quarks and antiquarks
- Dynamics description inside hadrons
- - Multi-quark: quark number >= 4
  - Hybrid state: the mixture of quark and gluon
  - **Glueball:** composed of gluons +
- Identification from QM: challenging
  - **Exotic quantum states** +
  - Crypto exotic with particular properties

### New form of hadrons:





## **Beijing Electron Positron Collider (BEPCII)**

### World unique e+e- accelerator in charm physics energy region



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#### **BESIII detector**



### **2004: Construction**

- Double rings
- Beam energy:
  - 1.0 2.3 (2.45)GeV
- Designed luminosity:
  - 1×10<sup>33</sup> cm<sup>-2</sup> s<sup>-1</sup>

2008: test run

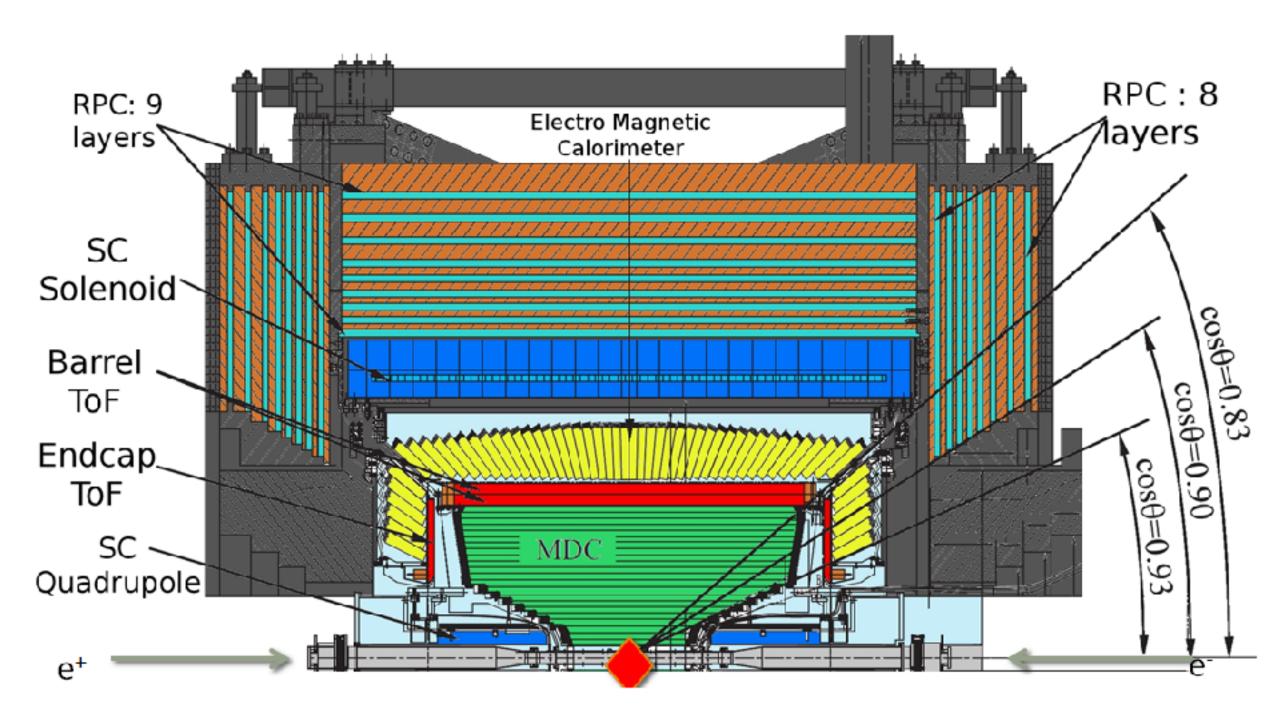
2009-now: BESIII physics runs



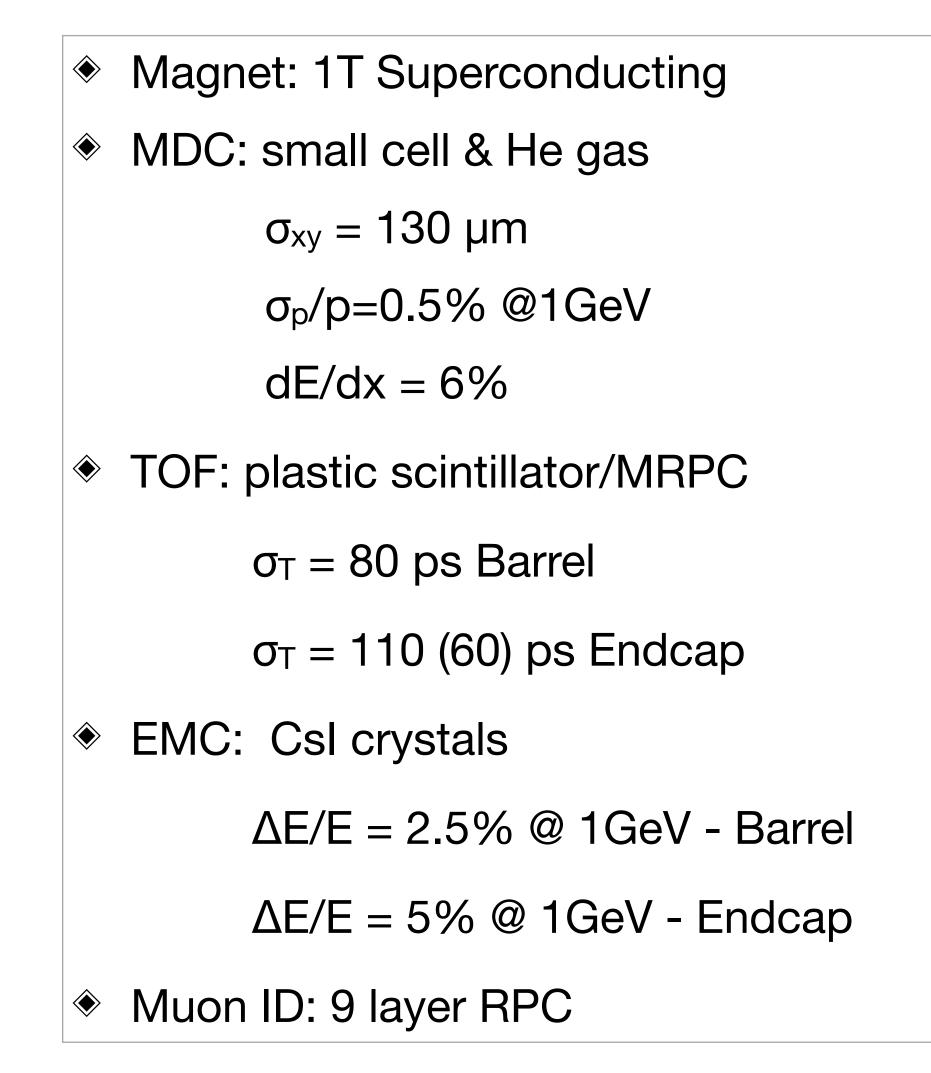


# **BESII detector**

### **Designed for neutral and charged particle with excellent resolution, PID, and large coverage**



Total weight 730 ton, ~40,000 readout channel Data rate: 5kHz, 50Mb/s



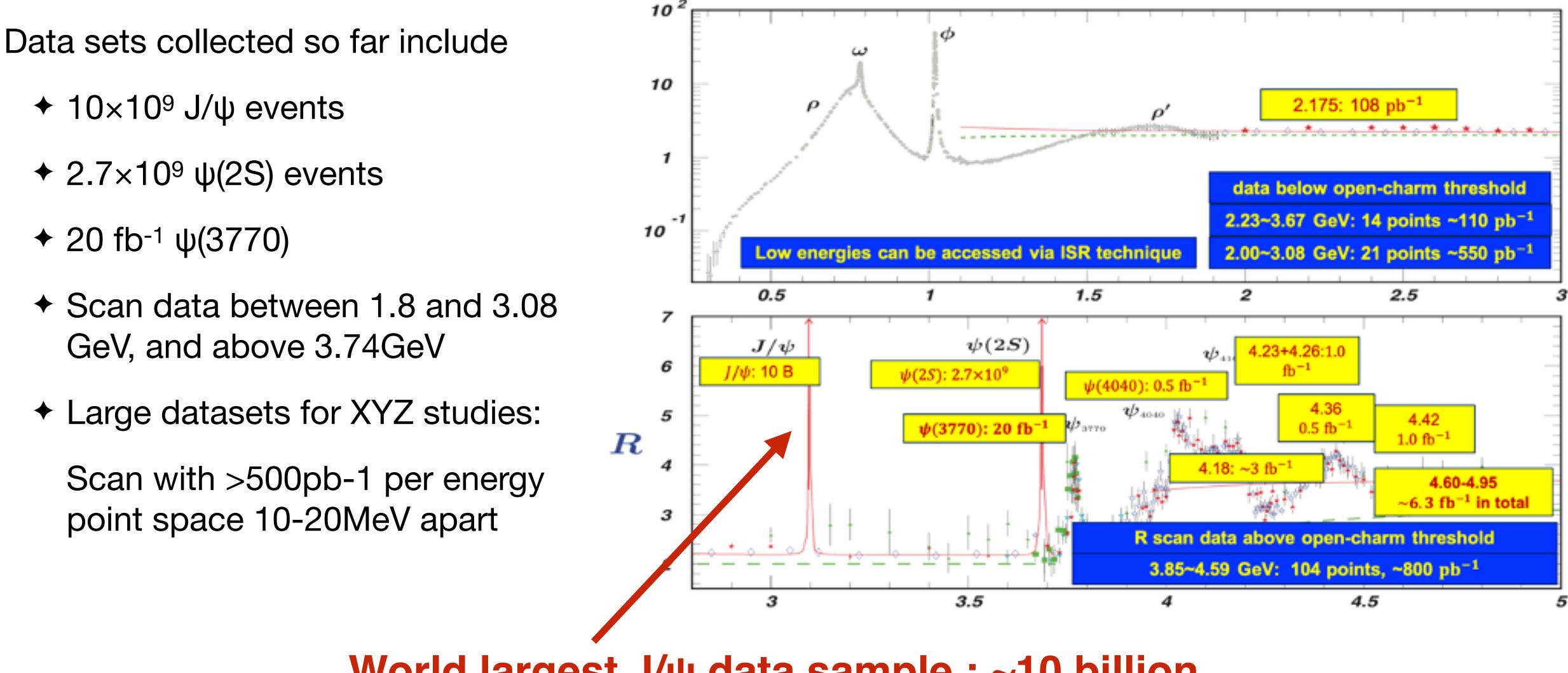
### Has been in full operation since 2008, all sub-detectors are in a very good status!







## **BESIII Data samples**



#### Totally about 50fb<sup>-1</sup> integrated luminosity

### World largest J/ψ data sample : ~10 billion

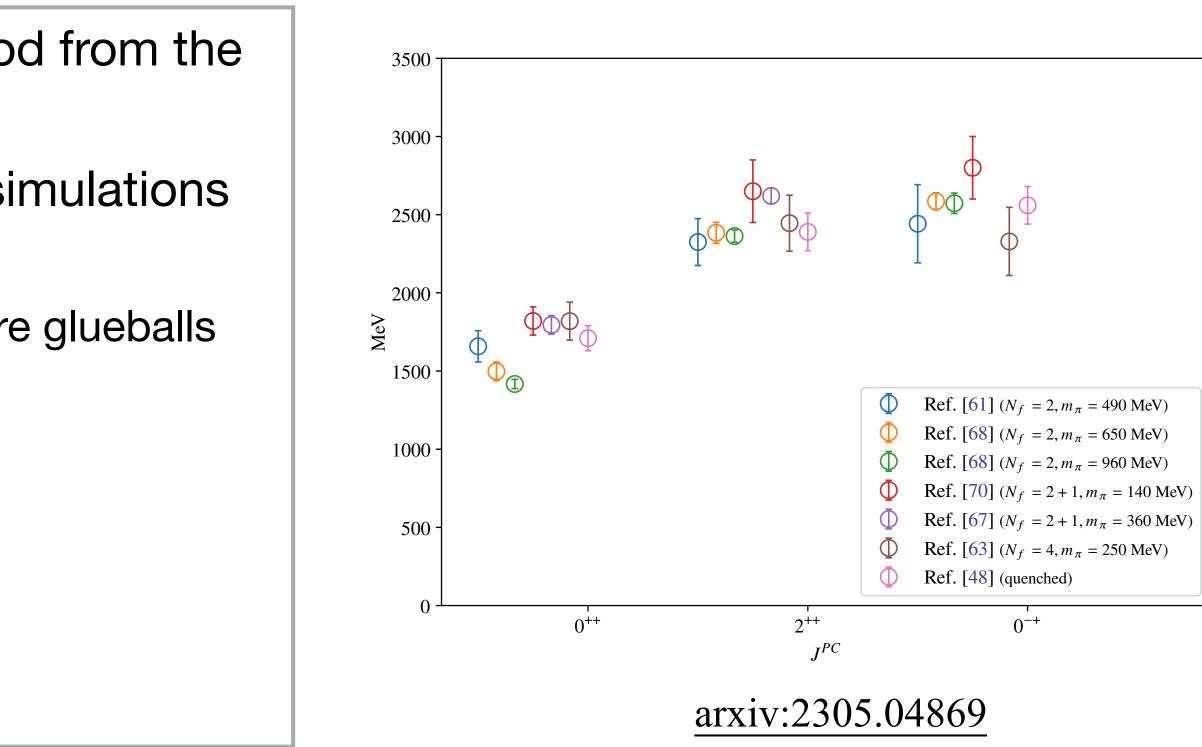




- - Glueballs to QCD is just as important as Higgs Boson to EW
- Lattice QCD (LQCD) is a non-perturbative method from the first principles in theory.
- **Different lattice QCD groups** (including lattice simulations with dynamical quarks)
  - Predictions on masses and production rates of pure glueballs
  - Consistent results and expected to be reliable.
- Lattice QCD predictions on glueball masses:
  - **0++ ground state:** 1.5 1.7 GeV/c<sup>2</sup>
  - ◆ 2++ ground state: 2.3 2.4GeV/c<sup>2</sup>
  - ◆ 0-+ ground state: 2.3 2.6GeV/c<sup>2</sup>

## Glueballs

The basic theory for strong interactions is quantum chromodynamics (QCD) Gluon self-interaction: prediction of non-Abelian Gauge SU(3) QCD theory + Glueballs are unique particles formed with force carriers via self-interactions







- Many experiments searched for glueballs over the past 4 decades
- Many historical glueball candidates, but with some difficulties/controversies.
  - ◆ Scalar Glueball candidate (0++): f<sub>0</sub>(1500), f<sub>0</sub>(1710)
  - + Tensor Glueball candidate (2++): f<sub>2</sub>(2340)
  - Pseudoscalar Glueball candidate (0-+): η(1405)

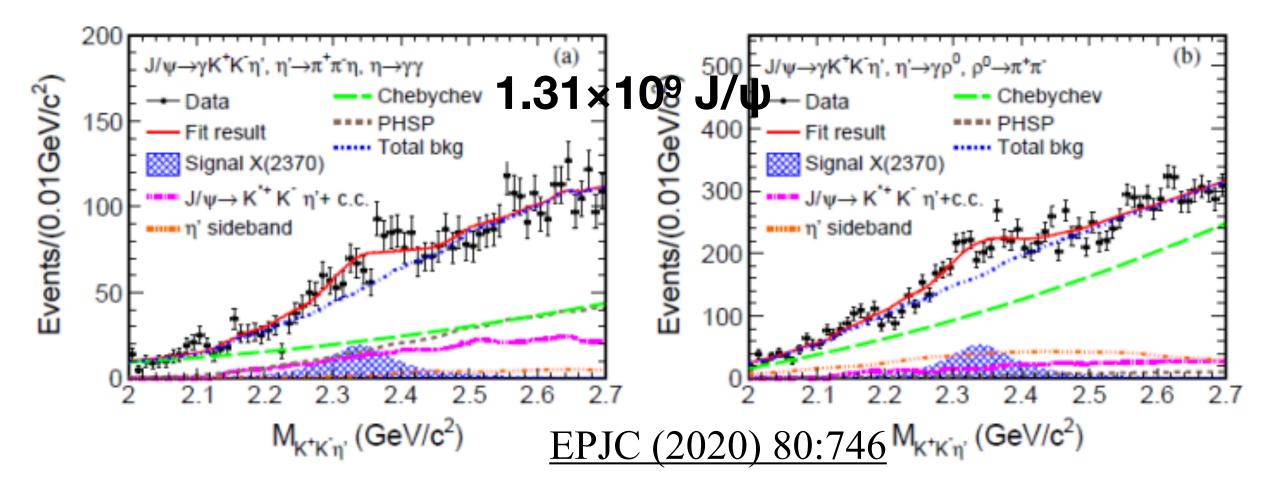




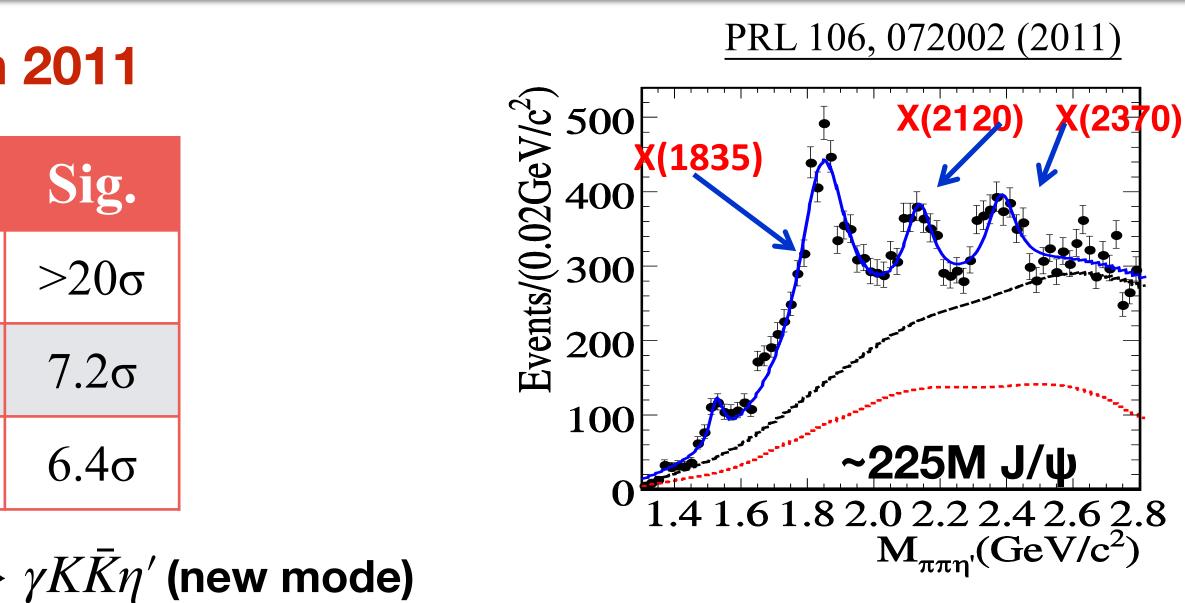
#### **Discovered by BESIII** in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ in 2011

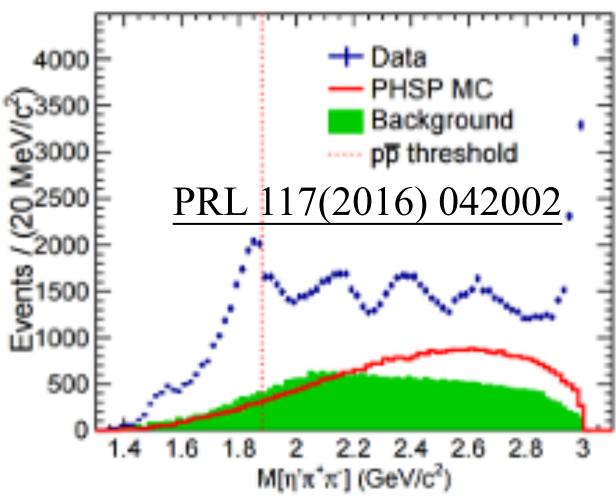
	M(MeV/c <sup>2</sup> )	$\Gamma(MeV/c^2)$	
X(1835)	1836.5±3.0+5.6-2.1	190.1±9.0+38-36	-
X(2120)	2122.4±6.7 <sup>+4.7</sup> -2.7	$83 \pm 16^{+31}$ -11	
X(2370)	$2376.3 \pm 8.7^{+3.2}_{-4.3}$	83±17+44-6	

Confirmed by BESIII in  $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$  and  $J/\psi \rightarrow \gamma K \bar{K} \eta'$  (new mode) ۲

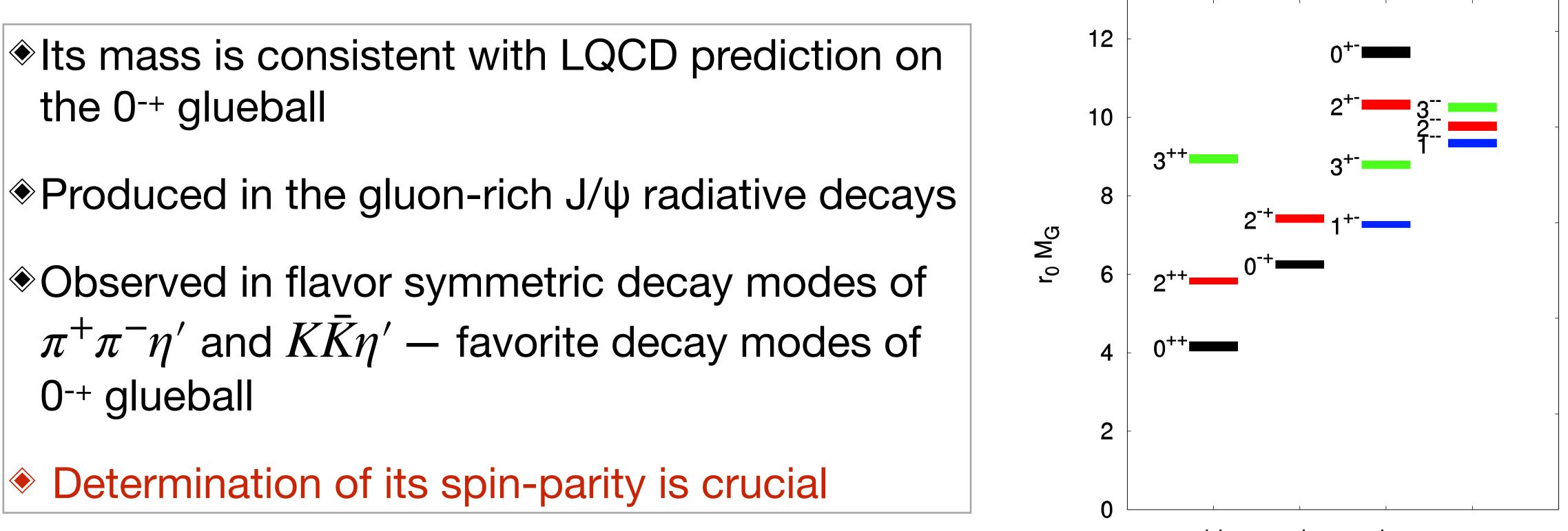


# X(2370)







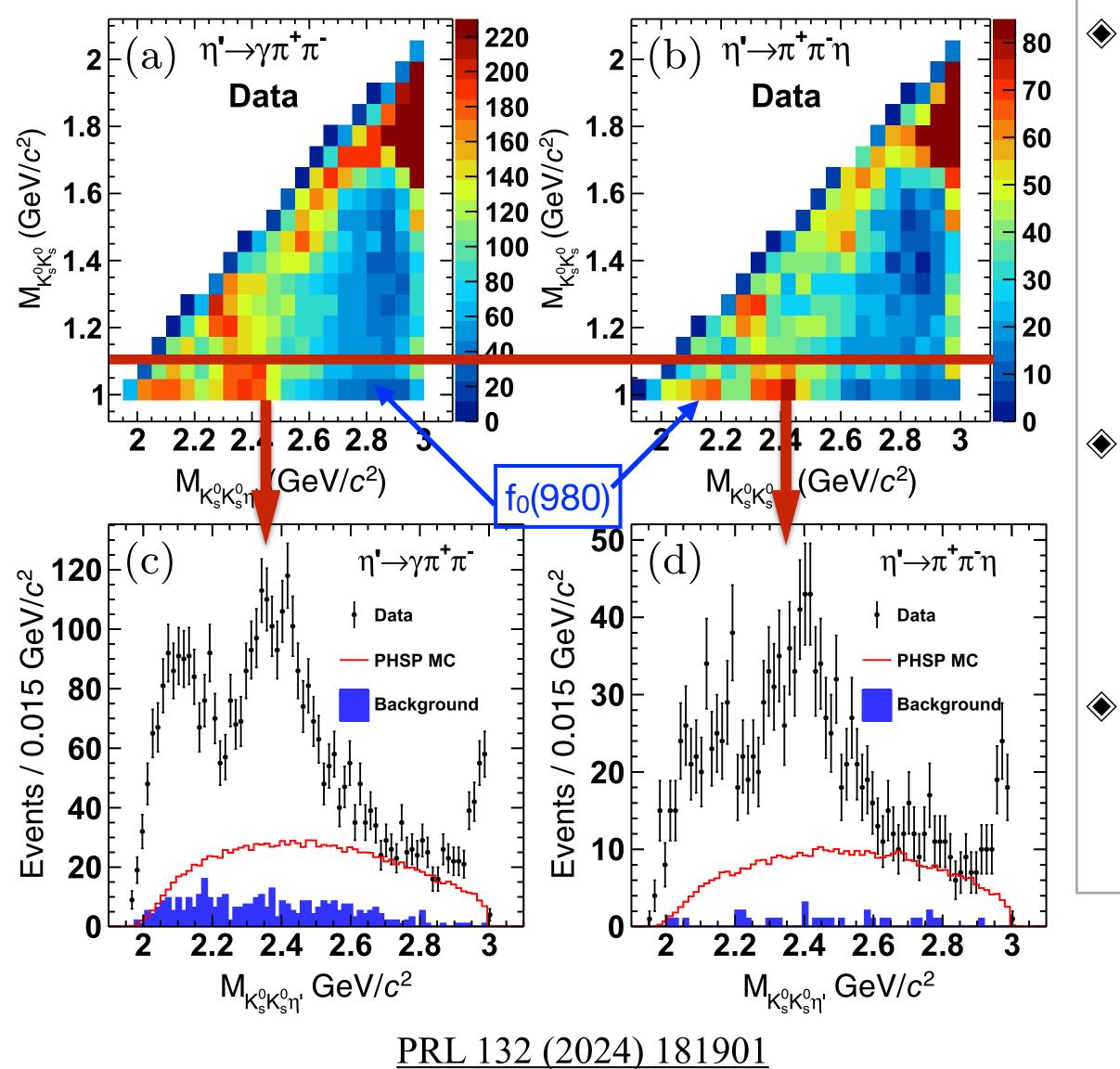


## X(2370) - good candidate of 0<sup>-+</sup> glueball



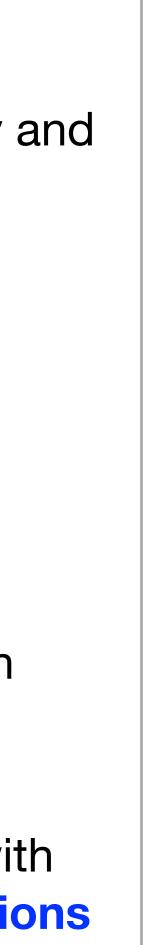


## Spin-Parity determination of the X(2370) in $J/\psi \rightarrow \gamma K^0_s K^0_s \eta^2$



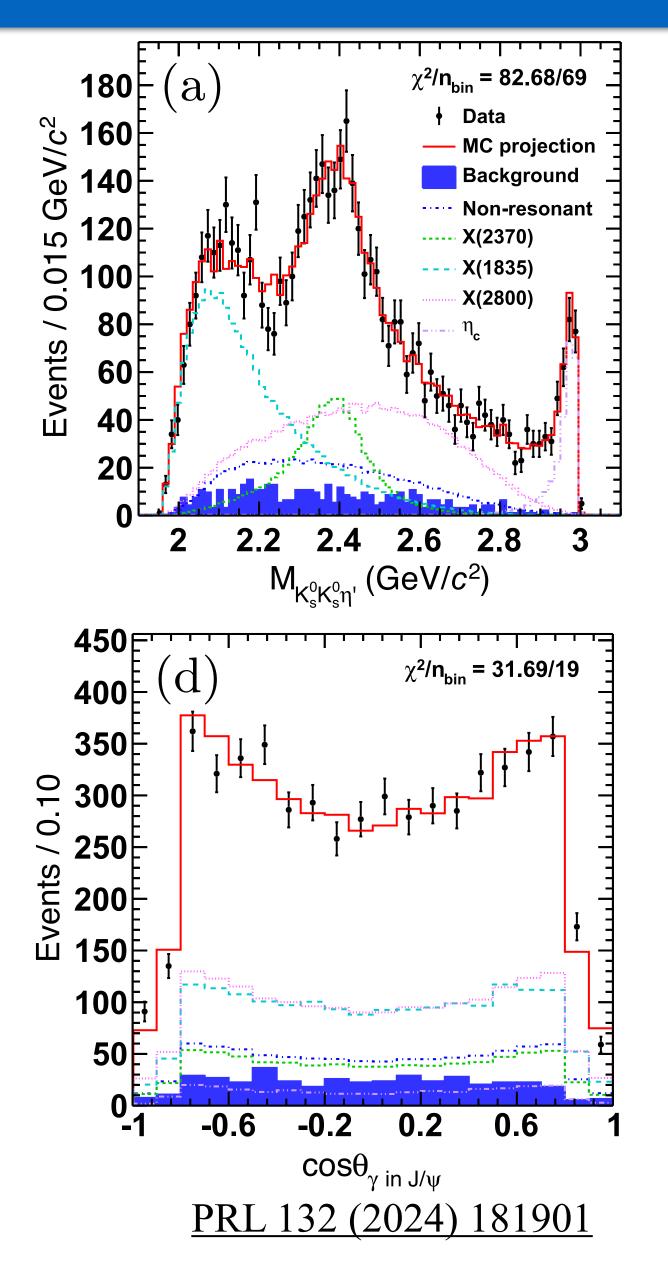
- Analysis advantage of  $J/\psi \rightarrow \gamma K^{0}_{s}K^{0}_{s}\eta'$ :
- Almost background free channel (exchange symmetry and C-parity conservation)
- + 10billion  $J/\psi$  data
- Very good BESIII detector performance
- Similar structures in  $\eta' \rightarrow \pi^+\pi^-\eta / \gamma\pi^+\pi^-$  modes:
  - Evident f<sub>0</sub>(980) in K<sup>0</sup><sub>s</sub>K<sup>0</sup><sub>s</sub> mass threshold
  - + Clear signal of X(1835), X(2370),  $\eta_c$  with f<sub>0</sub>(980) selection
  - Best PWA fit can well describe the data:
  - + Spin-parity of the X(2370) is determined to be 0-+ with significance larger than 9.8σ w.r.t. other J<sup>pc</sup> assumptions







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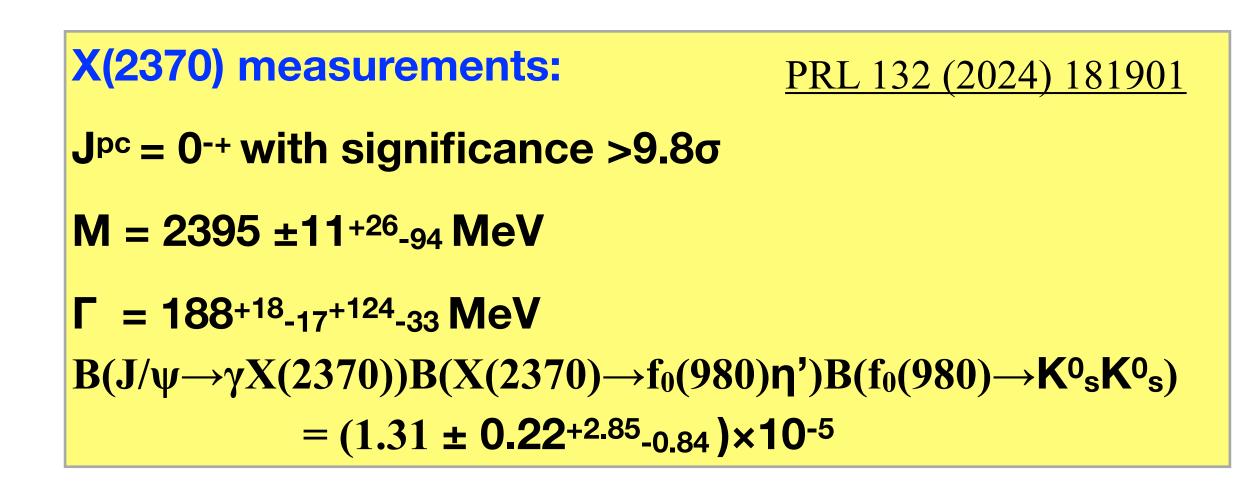
- Analysis advantage of  $J/\psi \rightarrow \gamma K^{0}_{s}K^{0}_{s}\eta'$ :
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  - Clear signal of X(1835), X(2370),  $\eta_c$  with f<sub>0</sub>(980) selection
- Best PWA fit can well describe the data:
  - Spin-parity of the X(2370) is determined to be 0<sup>-+</sup> with significance larger than 9.8 w.r.t. other J<sup>pc</sup> assumptions



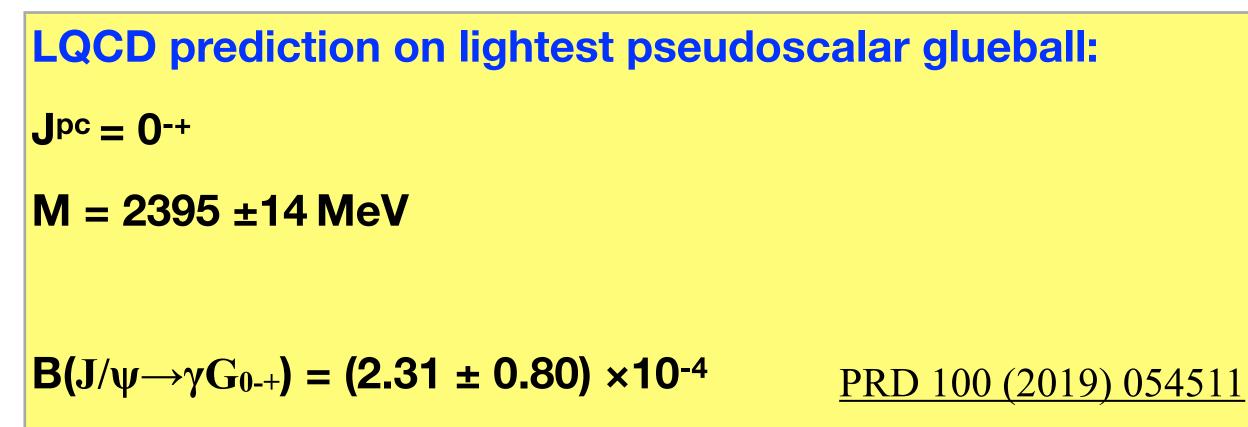




## **Compared with LQCD prediction on Lightest 0-+ Glueball**



- - + The spin-parity of the X(2370) is determined to be 0<sup>-+</sup> for the first time
  - Mass is in a good agreement with LQCD predictions
  - (assuming ~5% decay rate,  $B(J/\psi \rightarrow \gamma X(2370)) = (10.7^{+22.8} 7) \times 10^{-4})$



The measurements are in a good agreement with the predictions on lightest pseudoscalar glueball

+ The estimation on B(J/ $\psi \rightarrow \gamma X(2370)$ ) and prediction on B(J/ $\psi \rightarrow \gamma G_{0-+}$ ) are consistent within errors







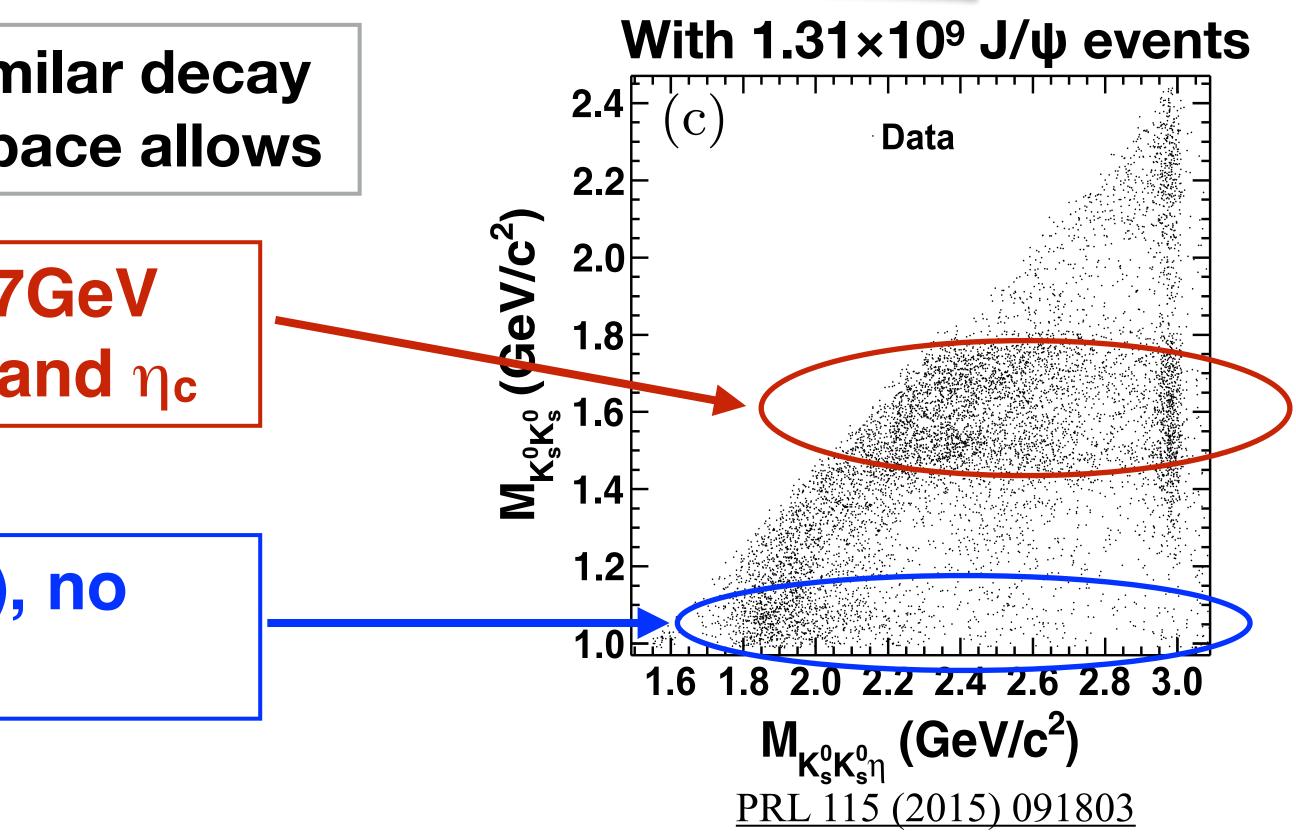
Observation and Spin-Parity Determination of the X(1835) in  $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta$ 

Qualitatively, we can clearly observe: similar decay patterns of the X(2370) and  $\eta_c$  if phase space allows

In the upper KK mass band of 1.5-1.7GeV range, clear signals of both X(2370) and  $\eta_c$ 

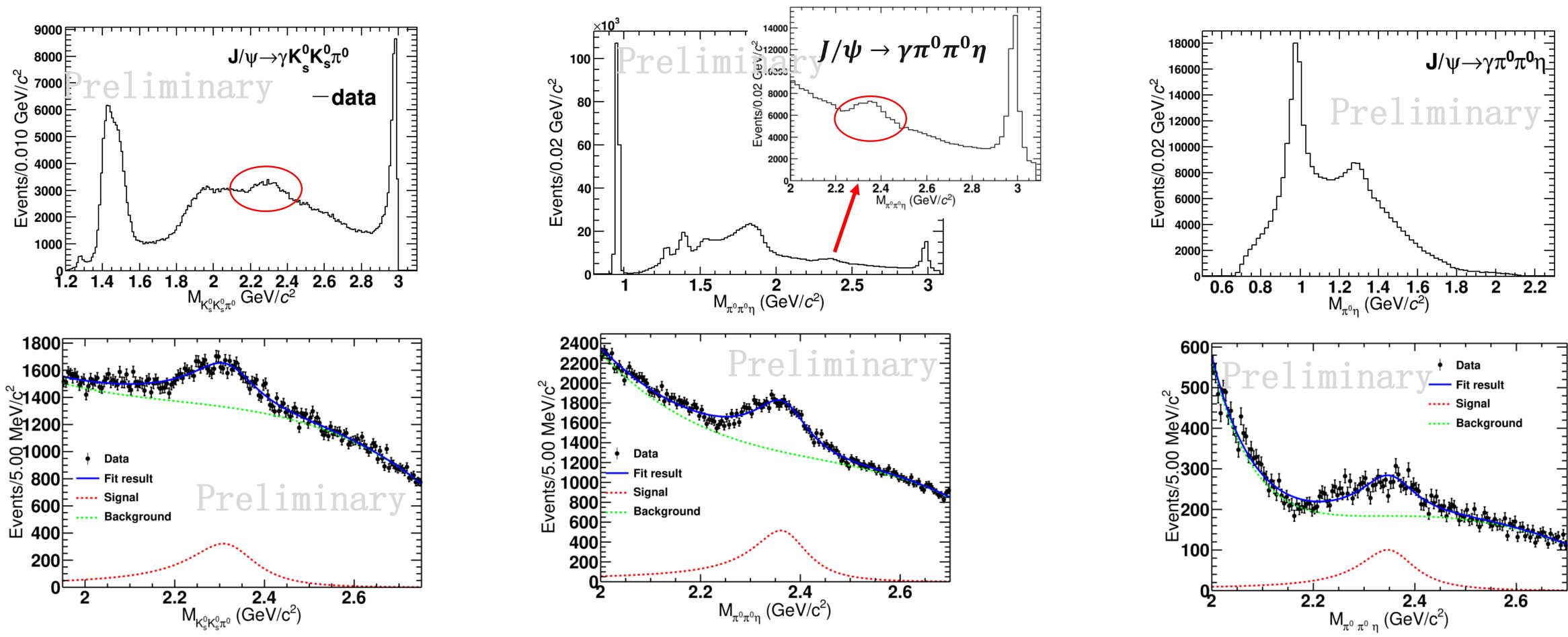
In the lower KK mass band of f<sub>0</sub>(980), no **X(2370), nor** η<sub>c</sub>

## X(2370) in $J/\psi \rightarrow \gamma K^0_s K^0_s \eta$





## **Observation of new decay modes of the X(2370)**



First observation of  $X(2370) \to K_s^0 K_s^0 \pi^0$ ,  $X(2370) \to \pi^0 \pi^0 \eta$  and  $X(2370) \to a(980)\pi^0$ with significances >>  $5\sigma$  and accompanied with  $\eta_c$ 







### **Observation of the X(2370) in the 5 golden decay modes**

- The glueball decays could be the analogy to Charmonium decays since they all decay via gluons (OZI) suppression) [PLB 380 189(1996), Commu. Theor. Phys. 23.373 (1995)]
  - + e.g. the 0<sup>-+</sup> glueball could have similar decays of  $\eta_c$

 $J/\psi \eta_c$ 

5 major  $\eta_c$  decay modes (from PDG) -5 "Golden" modes in 0<sup>-+</sup> glueball traditional searches

#### Decays involving hadronic resonances

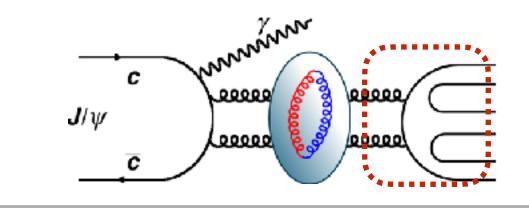
 $\eta'(958)\pi\pi$  $1.87 \pm 0.26)$  %  $\eta'(958) K\overline{K}$  $1.61\pm0.25)\%$ 

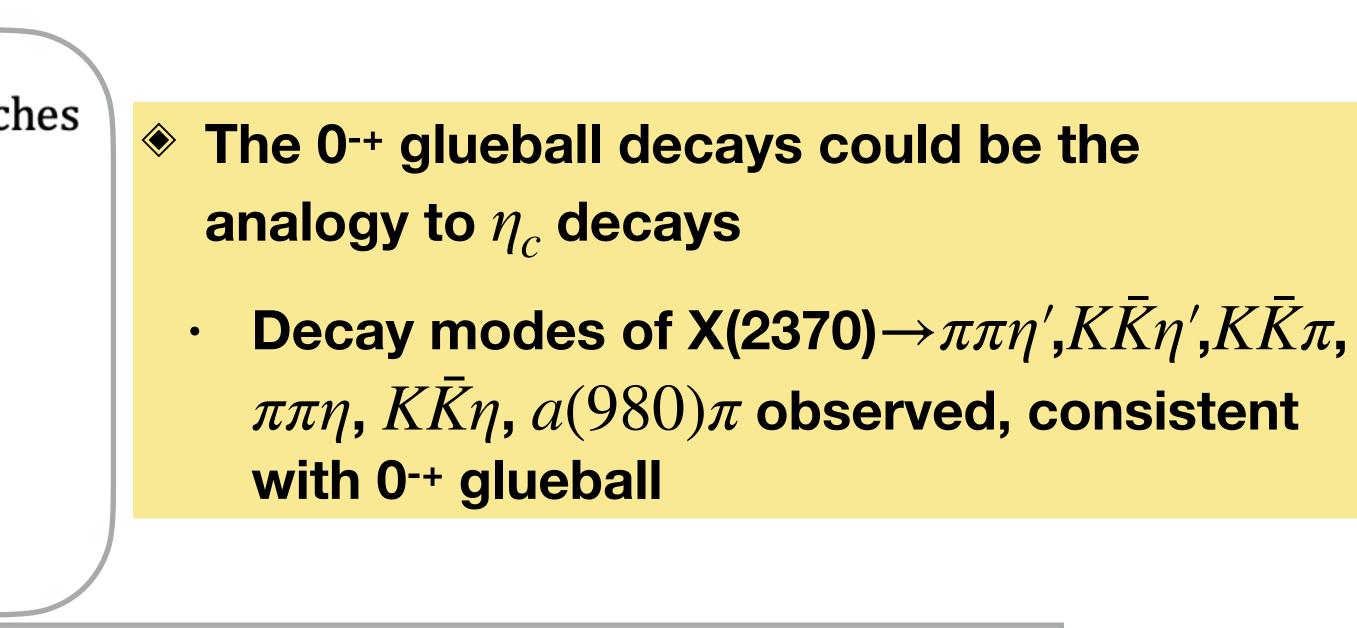
 $\Gamma_1$ 

 $\Gamma_2$ 

#### Decays into stable hadrons

 $K\overline{K}\pi$  $(7.0 \pm 0.4)\%$ I 34  $KK\eta$  $1.32\pm0.15)$  %  $(1.7 \pm 0.5)\%$ 



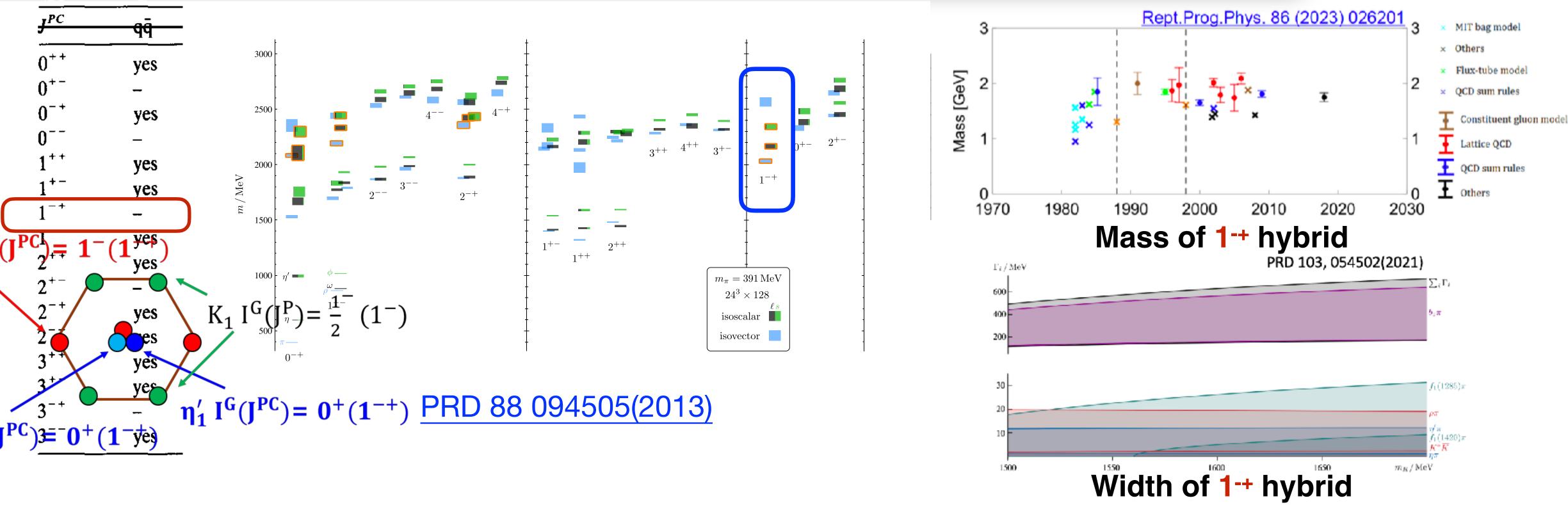


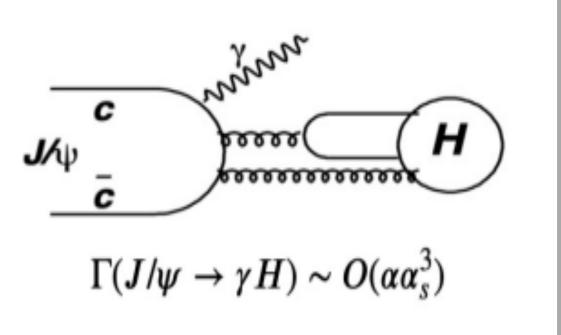
Such high similarity between the X(2370) and  $\eta_c$  decay modes strongly supports the glueball interpretation of the X(2370)





## Exotic 1-+ state





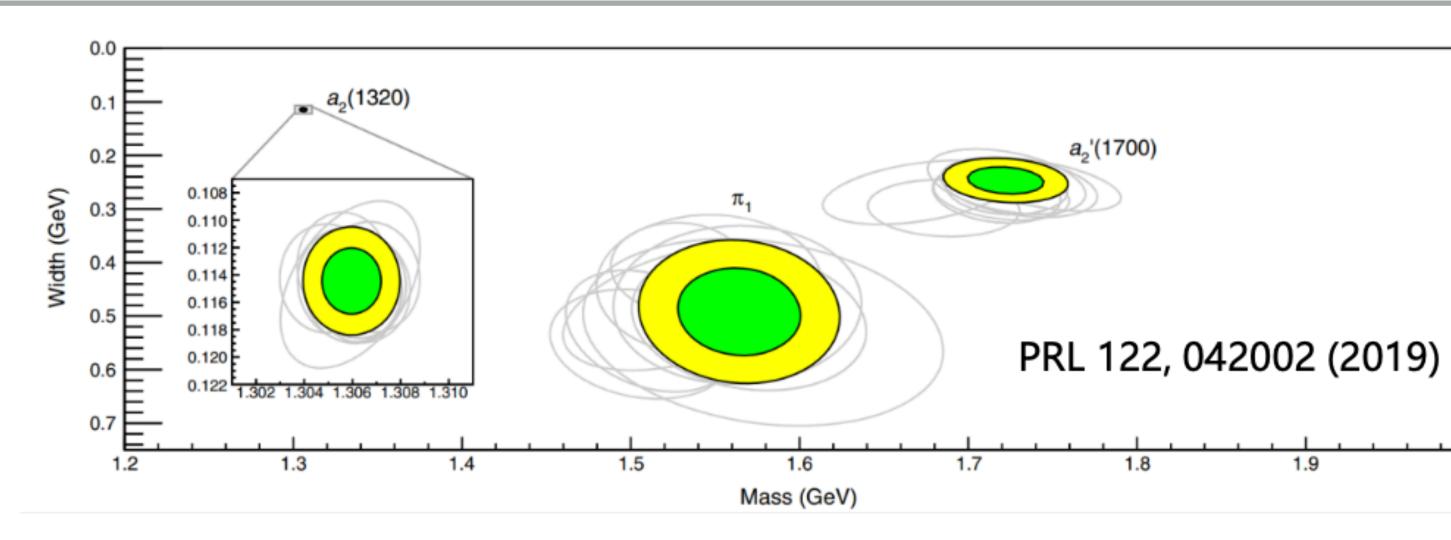
Spin-exotic state of 1<sup>-+</sup>: forbidden in conventional quark model Exotic state 1-+ provide an unique way for hybrid search: LQCD predicts the lightest nonet of 1-+ hybrids: 1.7 - 2.1GeV Can be produced in the gluon-rich charmonium decays



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- Over 3 decades, experimental evidence for 3 candidates with 1-+ state:
  - ✦ All 1<sup>-+</sup> iso-vectors
  - $\pi_1(1400)$  : seen in  $\eta\pi$
  - $\pi_1(1600)$  : seen in  $\rho \pi$ ,  $\eta' \pi$ ,  $b_1 \pi$ ,  $f_1 \pi$
  - +  $\pi_1(2015)$ : seen in  $b_1\pi$  and  $f_1\pi$
- Some claims are controversial
- $\pi_1(1400)$  and  $\pi_1(1600)$  can be one pole

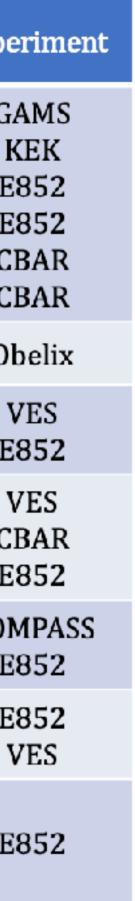


## Spin-exotic mesons

1.9

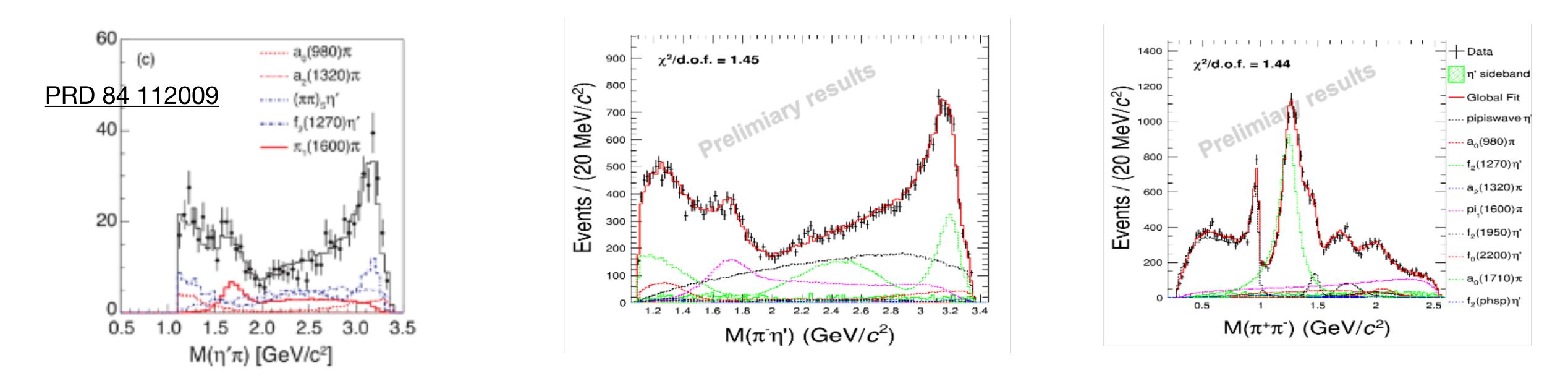
2.0

ates with 1-+ st	ate:				
			Decay mode	Reaction	Exper
		π <sub>1</sub> (1400)	ηπ	$\pi^{-}p \rightarrow \pi^{-}\eta p$ $\pi^{-}p \rightarrow \pi^{0}\eta n$ $\pi^{-}p \rightarrow \pi^{-}\eta p$ $\pi^{-}p \rightarrow \pi^{0}\eta n$ $\bar{p}n \rightarrow \pi^{-}\pi^{0}\eta$ $\bar{p}p \rightarrow \pi^{0}\pi^{0}\eta$	GA KI E8 E8 CB CB
			$ ho\pi$	$ar{p}p  ightarrow 2\pi^+ 2\pi^-$	Ob
				$\pi^{-}Be \rightarrow \eta' \pi^{-} \pi^{0}Be$ $\pi^{-}p \rightarrow \pi^{-} \eta' p$	V E8
		π <sub>1</sub> (1600)	$b_1\pi$	$\pi^{-}Be  ightarrow \omega\pi^{-}\pi^{0}Be$ $\bar{p}p  ightarrow \omega\pi^{+}\pi^{-}\pi^{0}$ $\pi^{-}p  ightarrow \omega\pi^{-}\pi^{0}p$	V) CB E8
			ρπ	$\pi^{-}Pb \rightarrow \pi^{+}\pi^{-}\pi^{-}X$ $\pi^{-}p \rightarrow \pi^{+}\pi^{-}\pi^{-}p$	COM E8
			$f_1\pi$	$\pi^- p  ightarrow p\eta \pi^+ \pi^- \pi^- \pi^- \pi^- A$ $\pi^- A  ightarrow \eta \pi^+ \pi^- \pi^- A$	E8 V
	π <sub>1</sub> (2015)	$f_1\pi$	$\pi^- p \rightarrow \omega \pi^- \pi^0 p$	EC	
		$\pi_1(2015)$	$f_1\pi$ $b_1\pi$	$\pi^- p \to p \eta \pi^+ \pi^- \pi^-$	E8
042002 (2019)					





## Observation of Exotic 1<sup>-+</sup> Isovector state $\pi(1600)$



- PWA in  $\psi' \to \gamma \chi_{c1}(\chi_{c1} \to \pi^+ \pi^- \eta')$  with higher  $\psi'$  data sample @ BESIII:

+ First observation of Exotic 1<sup>-+</sup> Isovector state  $\pi(1600)$  with a significance >10 $\sigma$  better than other  $J^{PC}$ assumption

+ The significance of phase motion is also greater than  $10\sigma$ 

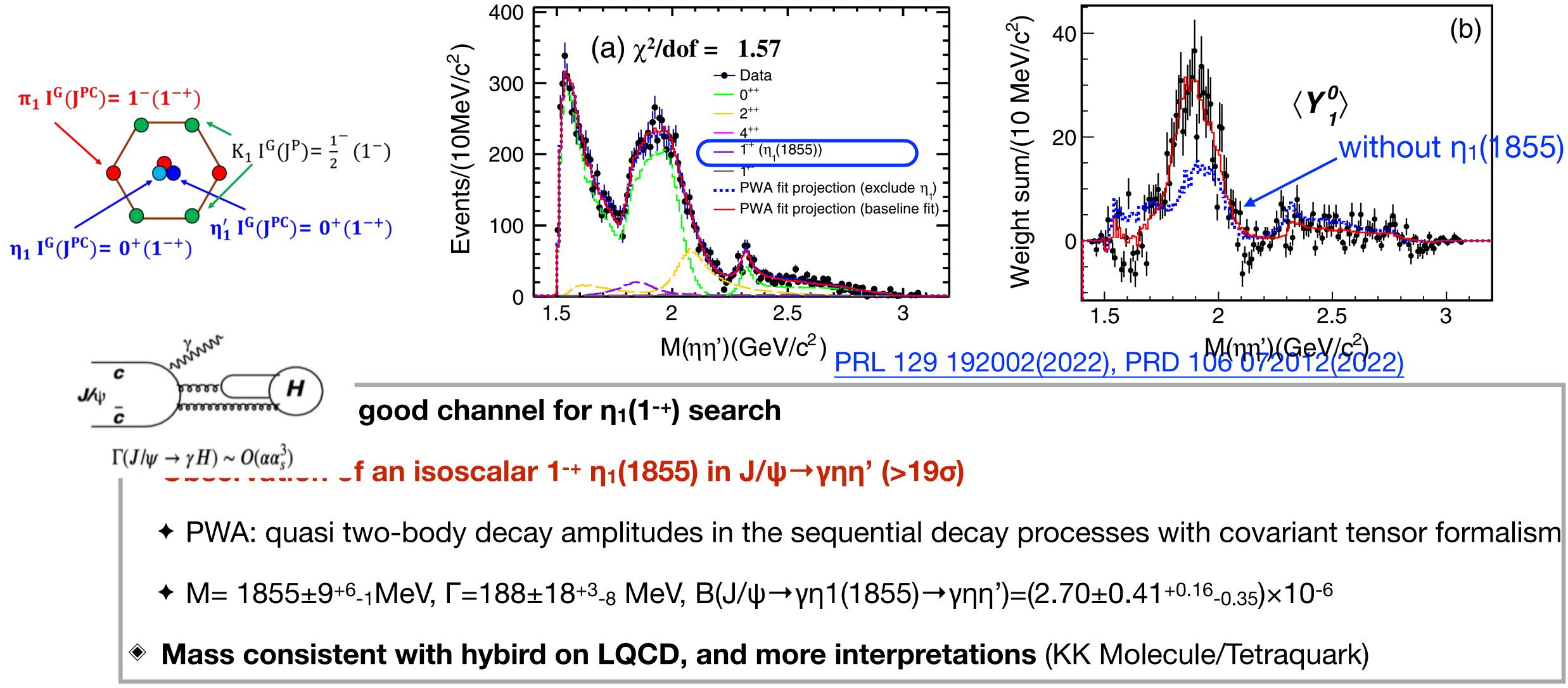
• CLEO-c results: evidence of an exotic P-wave  $\eta'\pi$  amplitude with  $4\sigma$  and but no significant phase motion





## **Observation of An Exotic 1-+ Isoscalar** $\eta_1(1855)$

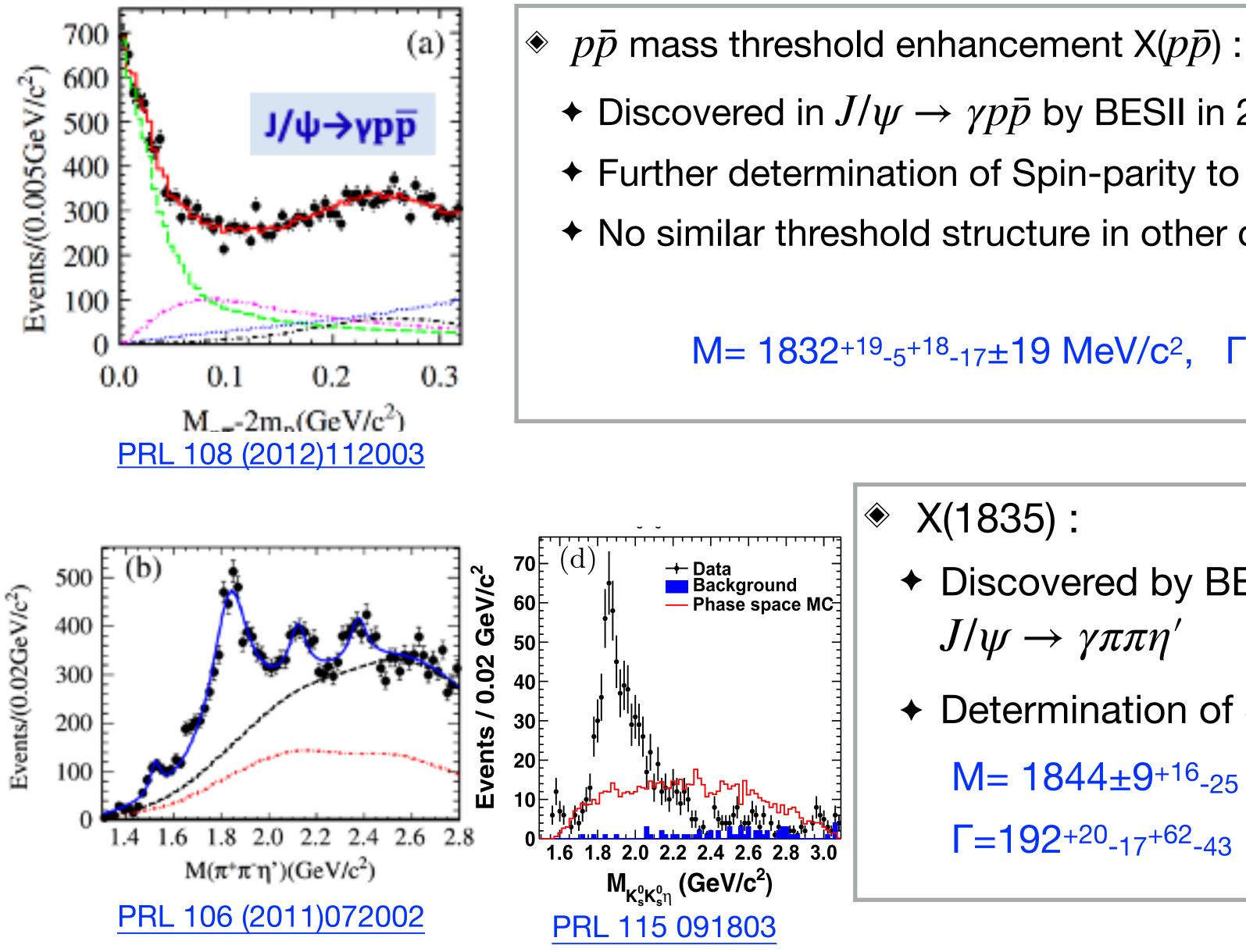
### Isoscalar (1-+) is critical to establish the nonet hybrid multiplet: partners for the Isovector (1-+)







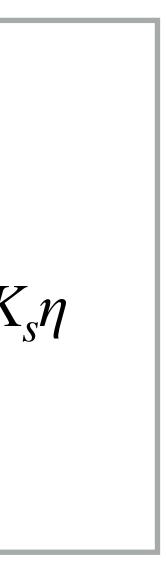
# **Observation of** $X(p\bar{p})$ **and** X(1835)



- + Discovered in  $J/\psi \rightarrow \gamma p \bar{p}$  by BESII in 2003 and confirmed by BESIII and CLEO-c Further determination of Spin-parity to be 0++
- + No similar threshold structure in other channels  $\rightarrow$  It can not be pure FSI effect
  - $M = 1832^{+19}_{-5}^{+18}_{-17} \pm 19 \text{ MeV/c}^2$ ,  $\Gamma = 13 \pm 19 \text{ MeV/c}^2$  (<76 MeV/c<sup>2</sup>@90% C.L.)

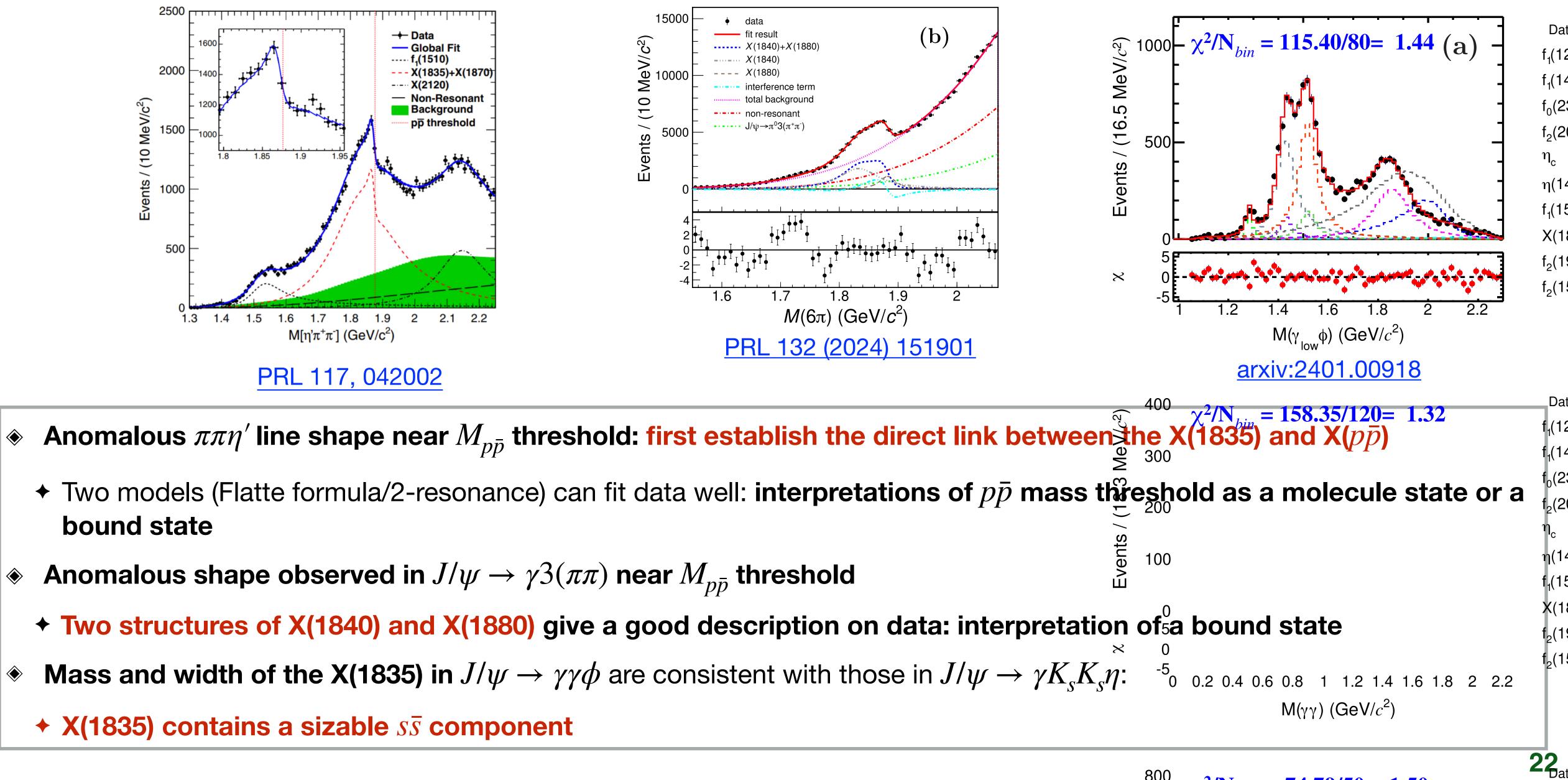
- ♦ X(1835) :
  - Discovered by BESII and confirmed by BESIII in  $J/\psi \rightarrow \gamma \pi \pi \eta'$
  - + Determination of Spin-parity to be 0<sup>-+</sup> in  $J/\psi \rightarrow \gamma K_s K_s \eta$ 
    - $M = 1844 \pm 9^{+16} 25 MeV/c^{2}$
    - $\Gamma = 192^{+20}_{-17}^{+62}_{-43} \text{ MeV/c}^2$



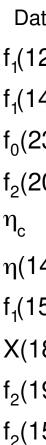




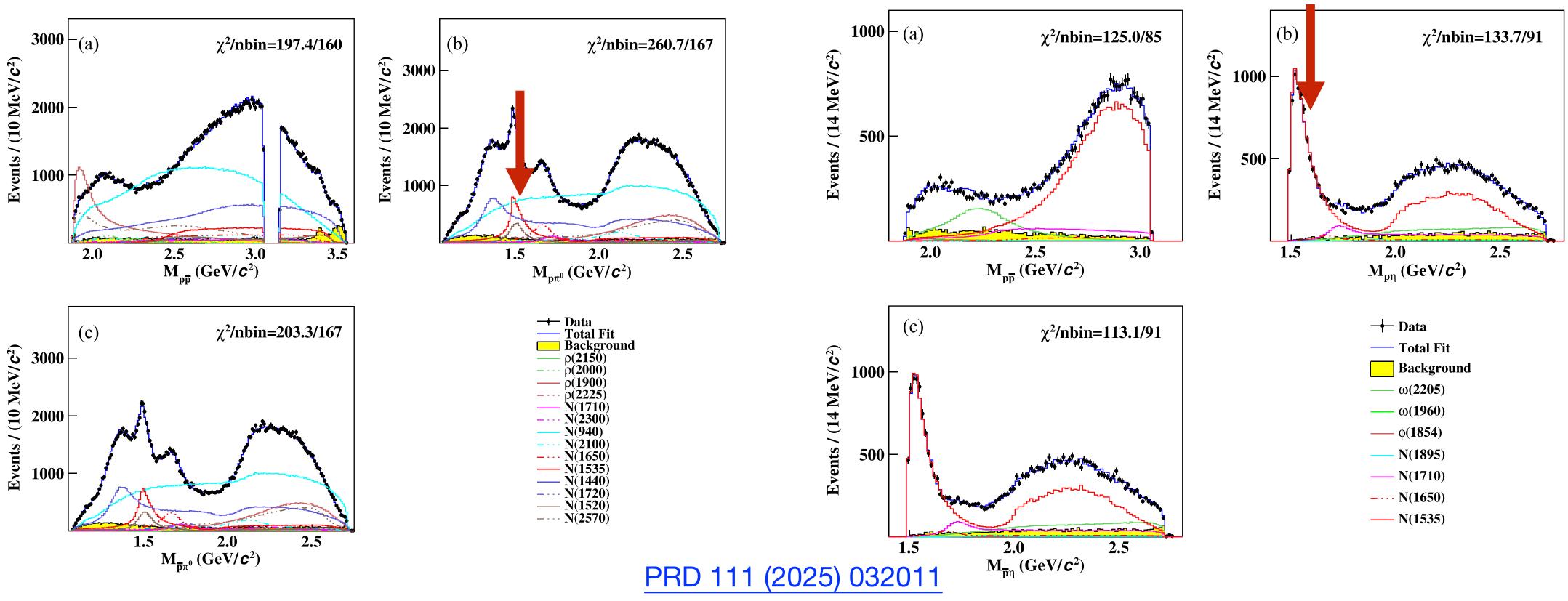
### $M(6\pi)$ (GeV/ $c^2$ ) Direct link between the $X(p\bar{p})$ and X(1835)



- - bound state
- ۲
- ۲
  - + X(1835) contains a sizable  $s\bar{s}$  component



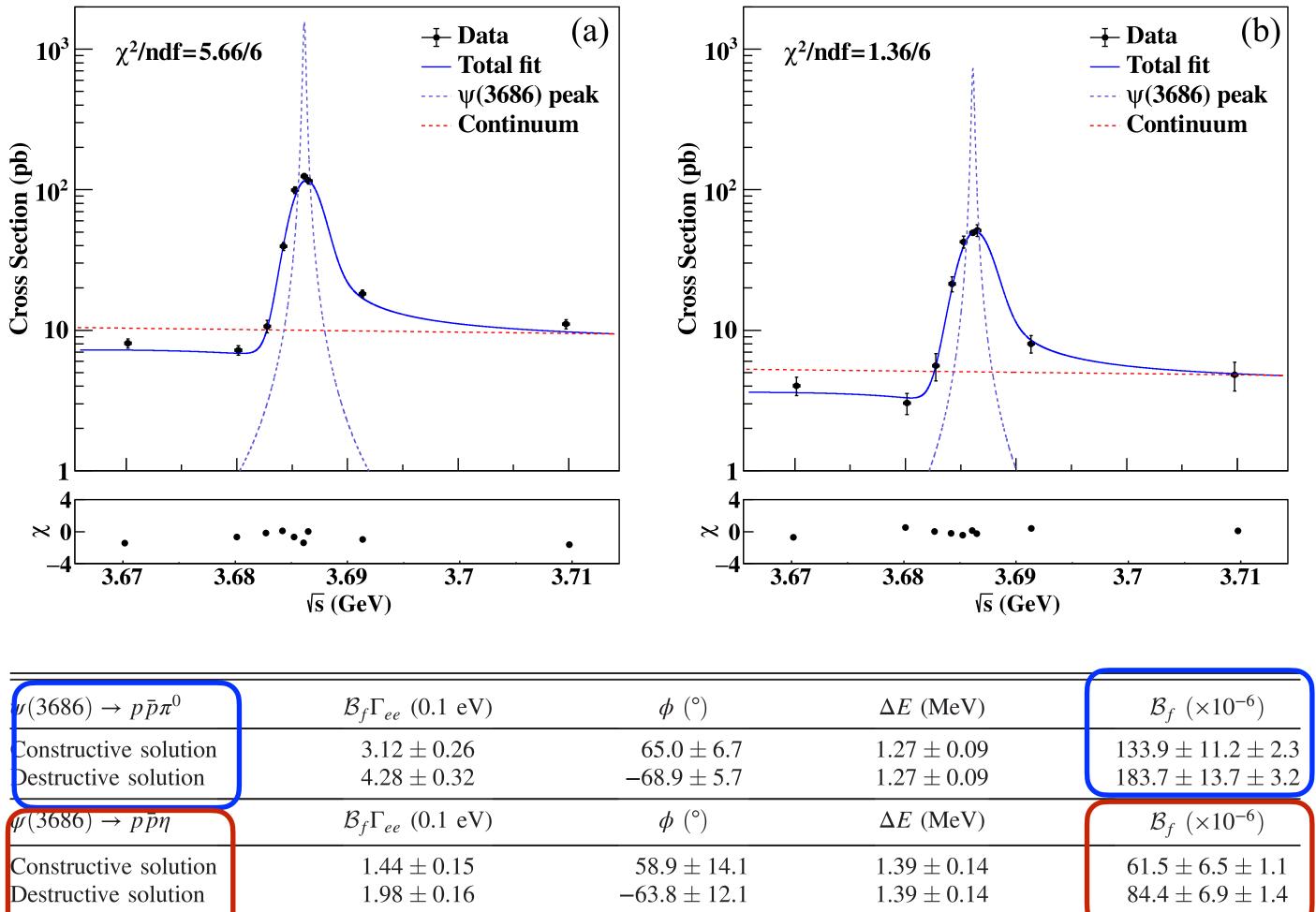
# $\psi(3686) \rightarrow p\bar{p}\pi^0 \text{ and } \psi(3686) \rightarrow p\bar{p}\eta$



Data can be well described by the well-established N\* states Ratio of the N(1535) decays is measured, which suggests a strong  $s\bar{s}$  components ٨  $\Gamma_{N(1535)\to N\eta} / \Gamma_{N(1535)\to N\pi} = 0.99 \pm 0.05 \pm 0.19$ 







Constructive solution Destructive solution	$\begin{array}{c} 3.12 \pm 0.26 \\ 4.28 \pm 0.32 \end{array}$	$65.0 \pm 6.7 \\ -68.9 \pm 5.7$	$\begin{array}{c} 1.27 \pm 0.09 \\ 1.27 \pm 0.09 \end{array}$	
$\psi(3686) \rightarrow p \bar{p} \eta$	$\mathcal{B}_{f}\Gamma_{ee}~(0.1~{ m eV})$	$\phi$ (°)	$\Delta E$ (MeV)	(
Constructive solution Destructive solution	$\begin{array}{c} 1.44 \pm 0.15 \\ 1.98 \pm 0.16 \end{array}$	$58.9 \pm 14.1$ -63.8 ± 12.1	$1.39 \pm 0.14 \\ 1.39 \pm 0.14$	

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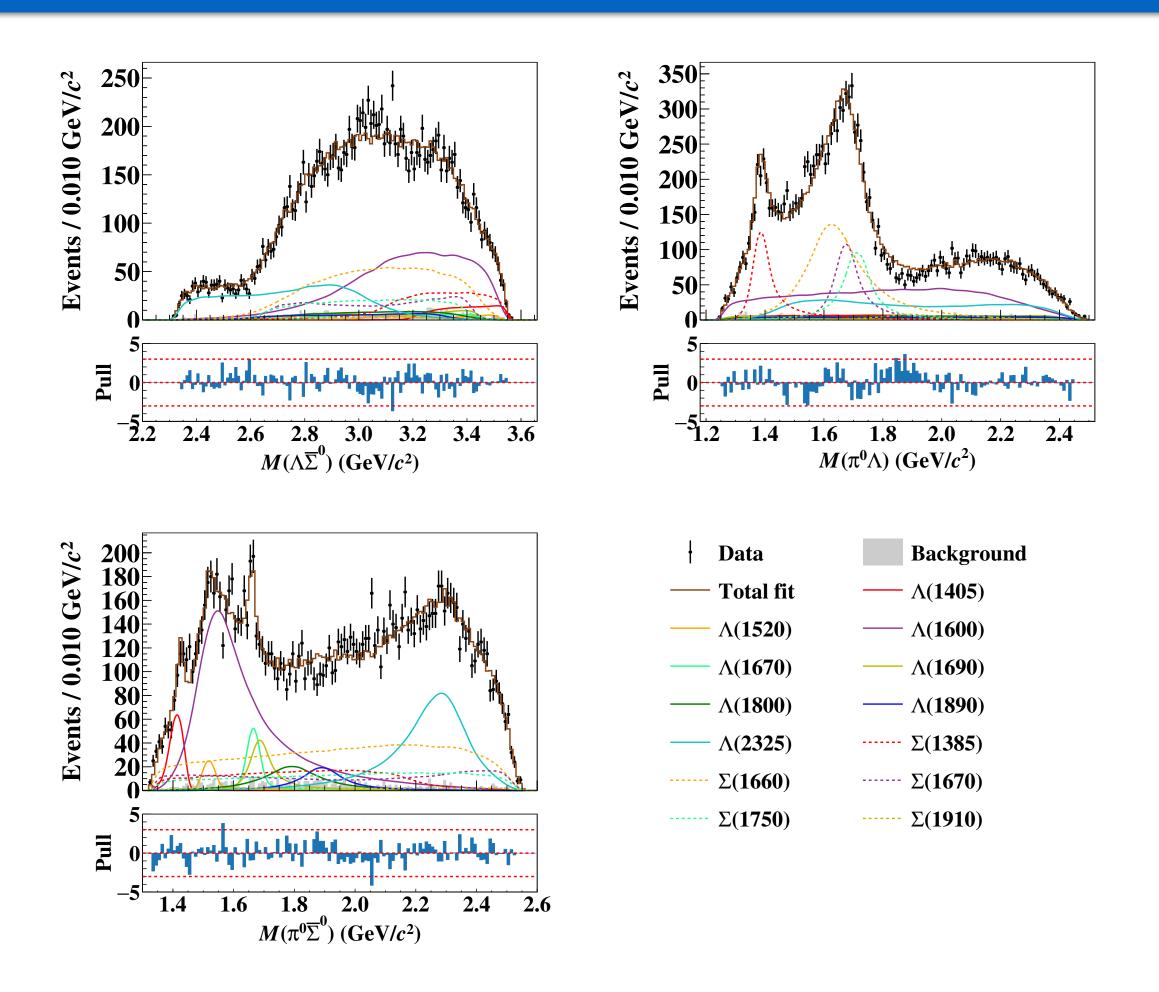
# $\psi(3686) \rightarrow p\bar{p}\pi^0 \text{ and } \psi(3686) \rightarrow p\bar{p}\eta$

 $B(\psi(3686) \to p\bar{p}\pi^0)/B(J/\psi \to p\bar{p}\pi^0)$ 

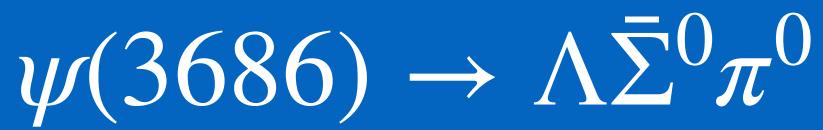
- +  $(11.3 \pm 1.2)$ % (constructive interference)
- +  $(15.4 \pm 1.6)$ % (destructive interference)
- $B(\psi(3686) \rightarrow p\bar{p}\eta)/B(J/\psi \rightarrow p\bar{p}\eta)$ 
  - +  $(3.1 \pm 0.4)$ % (constructive interference)
  - +  $(4.2 \pm 0.4)$  % (destructive interference)
  - Violated from the "12%" rule







 $\Lambda(2325)$  is necessary to better describe data with its spin-parity  $J^P = 3/2^{-1}$ 



 $q_i$ 

#### $\Lambda(1405)$ Faltte-like model

$$R(m) = \frac{1}{m_0^2 - m^2 - im_0 \left[\sum_i g_i \frac{q_i}{m} \times \frac{m_0}{|q_{i0}|} \times \frac{|q_i|^{2l_i}}{|q_{i0}|^{2l_i}} B_{l_i}^{\prime 2}(|q_i|, |q_{i0}|, d)\right]},$$

$$= \begin{cases} \frac{\sqrt{(m^2 - (m_{i,1} + m_{i,2})^2)(m^2 - (m_{i,1} - m_{i,2})^2)}}{2m} & (m^2 - (m_{i,1} + m_{i,2})^2)(m^2 - (m_{i,1} - m_{i,2})^2) \ge 0 \\ (m^2 - (m_{i,1} + m_{i,2})^2)(m^2 - (m_{i,1} - m_{i,2})^2) = 0 \end{cases}$$

#### $\Lambda(1405)$ chiral dynamics model (two poles)

$$\begin{split} R(m) &= \frac{1}{|I - VG|}, \\ V_{ij}(m) &= -C_{ij} \frac{1}{4f^2} (2m - M_i - M_j) \sqrt{\frac{M_i - E_i}{M_i}} \sqrt{\frac{M_j - E_j}{M_j}}, \\ G_k(m;\mu) &= \frac{1}{(4\pi)^2} \bigg\{ a_k(\mu) + \ln \frac{M_k^2}{\mu} + \frac{m_k^2 - M_k^2 + m^2}{2m^2} \ln \frac{m_k^2}{M_k^2} \\ &\quad + \frac{q_k}{m} \big[ \ln(m^2 - (M_k^2 - m_k^2) + 2q_km) + \ln(m^2 + (M_k^2 - m_k^2) + 2q_km) \\ &\quad - \ln(-m^2 + (M_k^2 - m_k^2) + 2q_km) - \ln(-m^2 - (M_k^2 - m_k^2) + 2q_km) \big] \bigg\}, \end{split}$$

#### Due to limited statistics, no separation power for the two $\Lambda(1405)$ models (11 $\sigma$ )





- A set of interesting and important results from the light hadron spectroscopy achieved: •
- **Discovery of a glueball-like particle: X(2370)** 
  - + Strong correlation between the X(1835) and  $M_{p\bar{p}}$  threshold enhancement. A molecule state or a bound state?
  - + Observation of An Exotic 1<sup>-+</sup> Isoscalar state  $\eta_1(1855)$  and Isovector state  $\pi(1600)$

**+** ...

With the more data, the more extensive and intensive investigation are ongoing, looking forward to new results in the near future.





