

Local basis truncation for Quantum Field Theories

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EPS-HEP 09.07.2025



Motivation

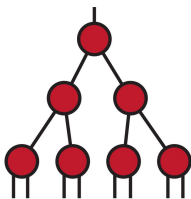
- Tensor Network (TN) methods: Extension to perturbation theory and Monte Carlo

Pros:

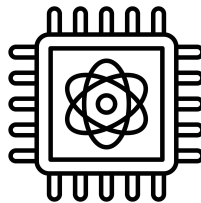
State of the art in 1d
non-perturbative and sign-problem free
non-equilibrium dynamics is accessible

Cons:

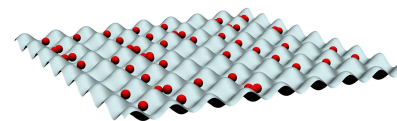
Efficient for low entanglement states
Hard in $d > 1$



Tensor Network



Quantum Computing



Quantum Simulations

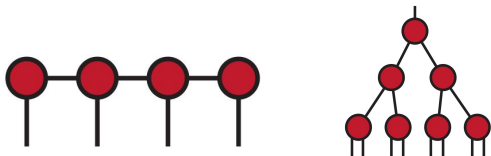


Toolkit:

Quantum many body wavefunction:

$$|\psi\rangle = \sum_{\{i_1, \dots, i_n\}} C_{i_1, \dots, i_n} |i_1, \dots, i_n\rangle$$

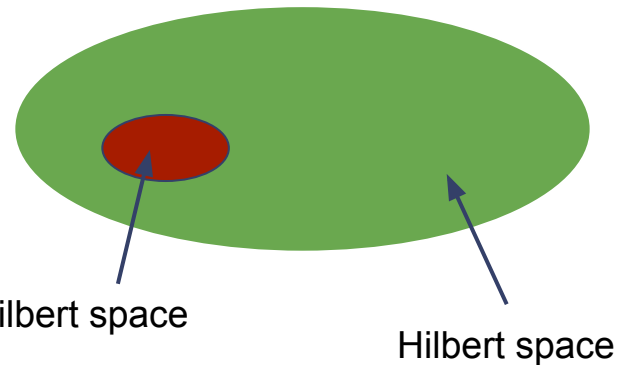
Ansätze:



Computational cost, example:

$$\mathcal{O}(d \cdot L \cdot BD^3)$$

$\sim d^N$



Why does it work?

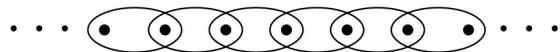
Many physically relevant states obey entanglement area laws.

Optimal local basis truncation of lattice quantum many-body systems

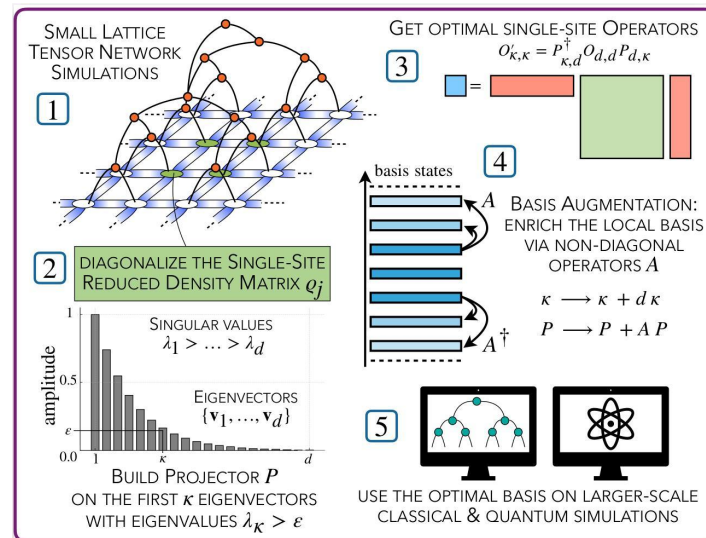


Method:

- Estimate reduced density ρ_i
- MF-theory, TN, Semi-definit programming



- Define projection P_i
- Effective Hamiltonian $O_i^{\text{eff}} = P^\dagger O_i P$

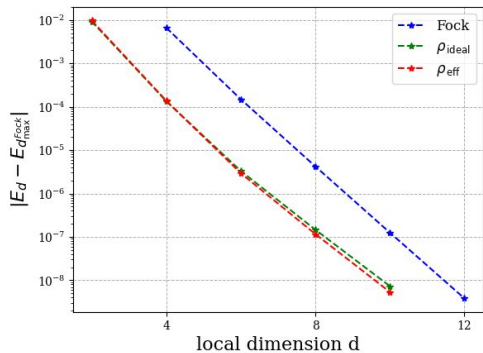




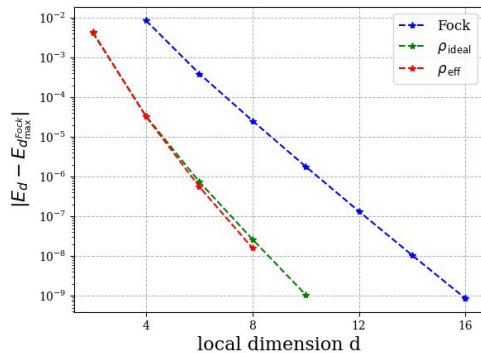
sG-model:

$$aH = \sum_{j \in \Lambda} \pi_j^2 + \frac{1}{2} [(\varphi_{i+1} - \varphi_j)^2 - \frac{M^2}{\beta^2} \cos(\beta \varphi_j)]$$

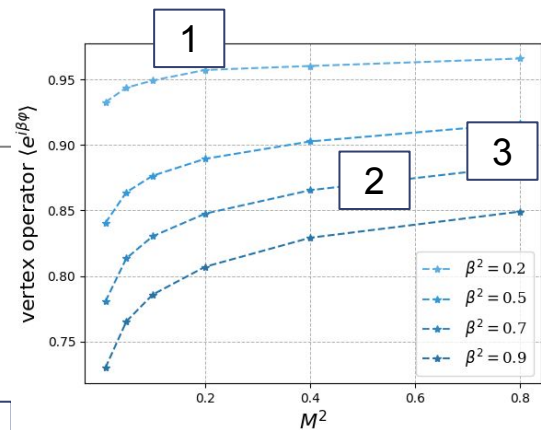
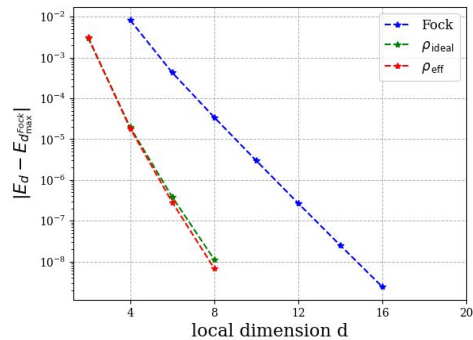
1



2



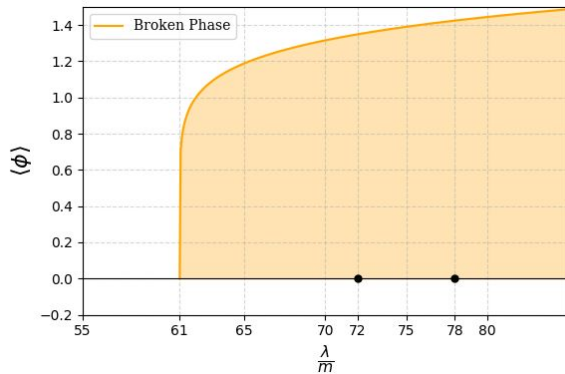
3



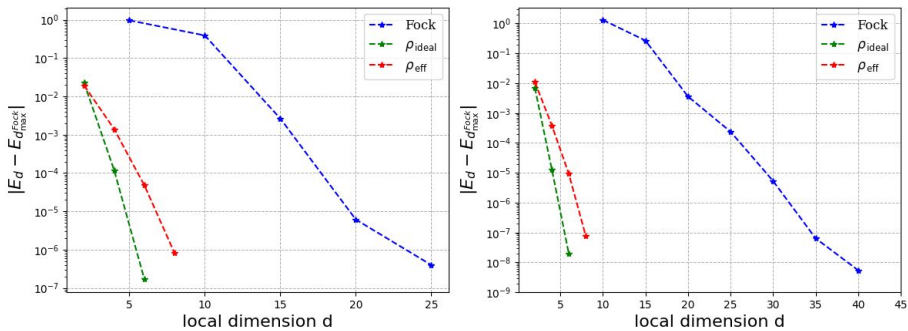


φ^4 -model:

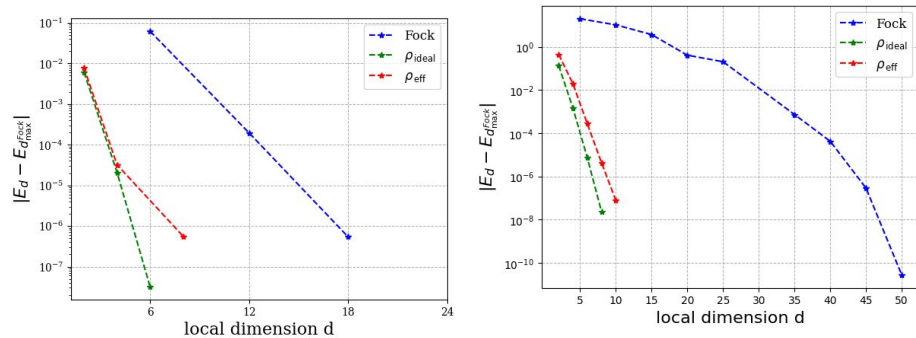
$$aH = \sum_{j \in \Lambda} \left[\frac{1}{2} \pi_j^2 + \frac{(\varphi_{j+1} - \varphi_j)^2}{2} + \frac{a^2 m^2}{2} \varphi_j^2 + \frac{a^2 \lambda}{4!} \varphi_j^4 \right]$$



1D:



2D:

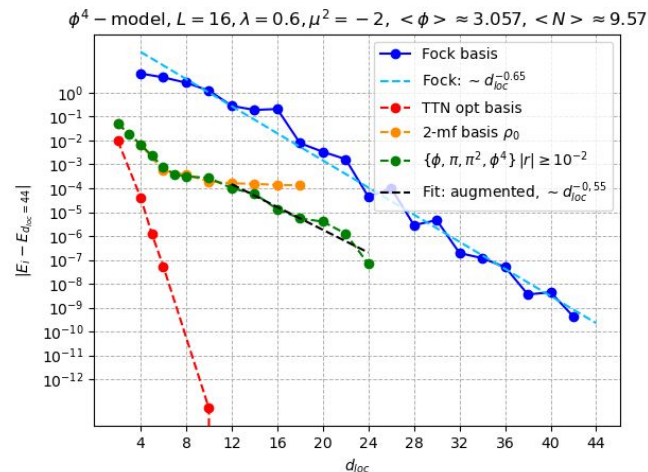




Basis augmentation:

- Estimate reduced density ρ_i
 - A priori, we don't know how good ρ_i is
 - Construct again the projector P_i
- extent basis by acting with local operators

$$|v'_i - \sum_l \langle v'_i | v_l \rangle v_l| > r$$



Lattice gauge theory dressed sites



Credits: Marco Rigobello

DressusG



Hamiltonian LGT:

$$H = \underbrace{\frac{i}{2} \sum_{x,k} \psi_x U_{x,k} \psi_{x+k}^\dagger + \text{h.c.}}_{\text{fermion hopping}} + \underbrace{m \sum_x (-1)^x \psi_x^\dagger \psi_x}_{\text{fermion mass}} + \underbrace{\frac{g^2}{2} \sum_{x,k} E_{x,k}^2 + \frac{1}{g^2} \sum_{\square} P}_{\text{electric+magnetic energy}}$$

Gauss law :

$$G_x^a |\psi_{\text{phys}}\rangle = \left[\sum_{k=1}^D \left(L_{x,k}^a - R_{x,k}^a \right) - Q_x^a \right] |\psi_{\text{phys}}\rangle = 0 \quad \forall x, a$$



Truncation of gauge-dof:

- The change of basis from group element basis $\{|g\rangle\}$ to the representation basis $\{|jmn\rangle\}$ is given by

$$\langle g|jmn\rangle = \sqrt{\frac{\dim(j)}{\text{Vol}(G)}} [\pi_j(g)]_{mn}$$

- Naturally the parallel transporter can then be decomposed as:

$$\langle j'm'n'|U_{M,N}^J|jmn\rangle = \sqrt{\frac{\dim(j)}{\dim(j')}} \sum_{\alpha} \overline{C_{jmJM}^{j'n',\alpha}} C_{jnJN}^{j'm',\alpha}$$



Encoding for TN*:

- Discard un-physical states: $\mathcal{H} = \bigotimes_x \mathcal{H}_x$
- To preserve the tensor product structure, we “double” symmetry content :

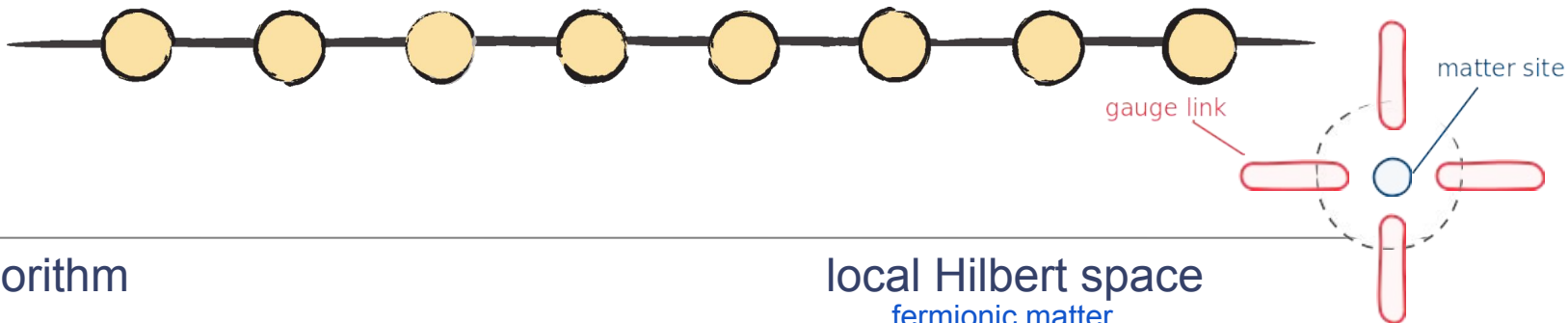
$$[G_x, G_y] = 0$$

$$|j, m_L, m_R\rangle \longrightarrow |j, m_L\rangle \otimes |j, m_R\rangle \quad (\text{in truncated irrep-basis})$$

← split in semi-links

- Fuse site and surrounding semi-links and find singlets
- Impose abelian constraint on semi-links

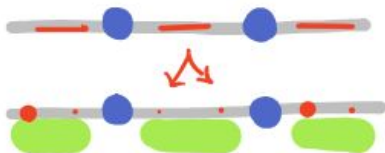
*Silvi et al. PRD 100 074512 (2019)



Algorithm

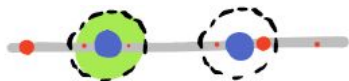
1

split links



2

combine site + semi-links



3

select colour singlets (**Gauss law**)

local Hilbert space

fermionic matter

single component

$$\{\psi_x, \psi_y^\dagger\} = \delta_{xy}$$

bosonic gauge link

$$|g\rangle \quad \forall g \in \mathbf{SU}(3)$$

infinite dimensional

$$\text{Peter-Weyl } \mathcal{H} = \bigoplus_j (j \otimes j^*)$$

$$|j, m, \bar{m}\rangle \quad \forall \text{ irrep } j, \quad m \in j, \quad \bar{m} \in j^*$$

$$\text{algebra } [E^a, U^\dagger] = iT^a U^\dagger$$

$$U |j, m, \bar{m}\rangle = \text{Clebsch-Gordan}$$

$$E^2 |j, m, \bar{m}\rangle = C_2(j) \text{ Casimir}$$



Input:
 (1+d)D
 background charge
 matter content
 gauge group
 truncation scheme

What gauge groups?

$Z_N, U(1), SU(N)$

Non-simple groups:

$U(1) \otimes SU(2) \otimes SU(3)$

```
links: 2
gauge:
  G:
    group: SU 3
    num-C2-levels: 1
matter:
  u:
    irreps:
      G: 1 0 0
  d:
    irreps:
      G: 1 0 0
```



Output:
 Operators

(gauge invariant)
 (abelian constraint)
 (local interaction)

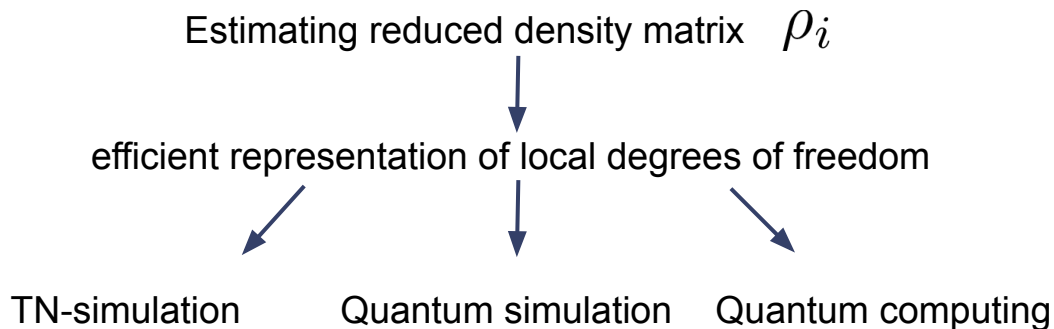
It's modular! You can add your favourite gauge group or truncation.

ℓ	d					
	(2 + 1)-dimensions			(3 + 1)-dimensions		
	U(1)	SU(2)	SU(3)	U(1)	SU(2)	SU(3)
1	35	30	164	267	178	3096
2	165	168	752	3437	3670	52476
3	455	600	3738	18487	35280	813438
4	969	1650	19878	64953	214958	17490134
5	1771	3822	43698	177155	967466	69232482
6	2925	7840	82128	408421	3509062	228461186
7	4495	14688	212496	835311	10828494	1245755754
8	6545	25650	333538	1561841	29473038	2782999996

Magnifico, Cataldi, Rigobello, PM et al.
 arXiv: 2407.03058



Summary:



Dressuag package: Efficient representation of local degrees of freedom of gauge theories for TN-simulations



Collaborators:

Giovanni Cataldi

Marco Rigobello

Mattia Morgavi

Pietro Silvi

Simone Montangero

Thank you!

This work was done under the financial support from the European Union's Horizon 2020 Research and Innovation Programme under the Marie Skłodowska-Curie Grant Agreement No. 101034267.