

A precise determination of α_s from the heavy jet mass distribution



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de Física Fundamental y Matemáticas

Taskforce A. Bhattacharya, M. Benitez, A.H. Hoang, M.D.
Schwartz, I.W. Stewart and X. Zhang based on 2502.12253



July 7th, 2025

EPS-HEP 2025, Marseille

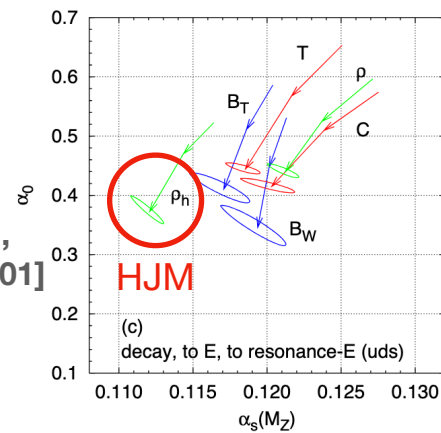
Outline

- A story on the strong coupling
- Event shapes and heavy jet mass
- Theoretical considerations
- Fitting procedure and results
- Summary

α_s from Event Shapes

[See also prior work of Catani, Trentadue, Turnock, Webber 1993]

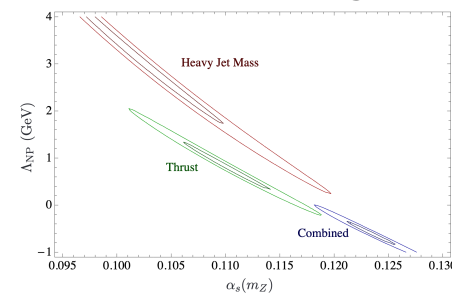
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Secondly fits for the heavy-jet mass (a very non-inclusive variable) lead to values for α_s which are about 10% smaller than for inclusive variables like the thrust or the mean jet mass. This needs to be understood. It could be due to a difference in the behaviour of the perturbation series at higher orders.

[See also Dissertori et al 2007]

N³LL resummation with NNLO matching



Event Shape	$\alpha_s(m_Z)$	Λ_{NP} (GeV)	$\chi^2/\text{d.o.f.}$
Thrust	0.1101	0.821	66.9/47
Heavy Jet Mass	0.1017	3.17	60.4/43
Combined	0.1236	-0.621	453/92

[Chien, Schwartz 2010]

These studies used, in parts, only
25% of data bins:
 $0.08 < \rho < 0.18$

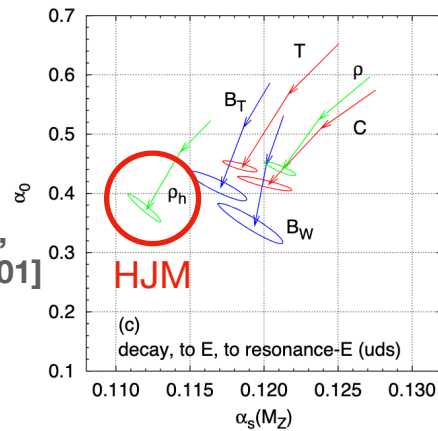
2001-2010

Prior to SCET, only NLL resummation possible

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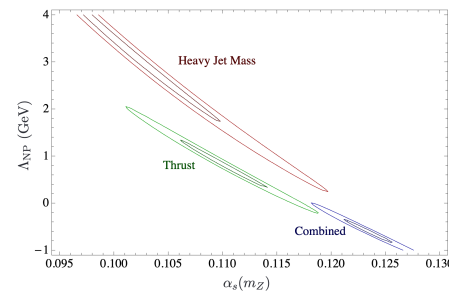
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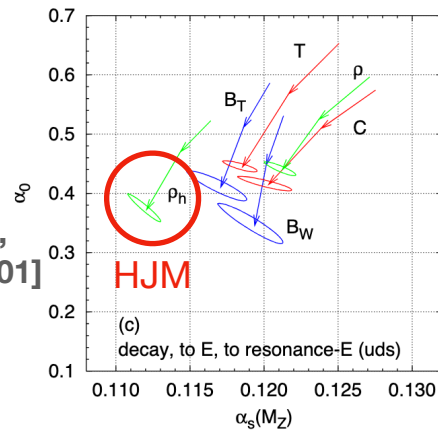
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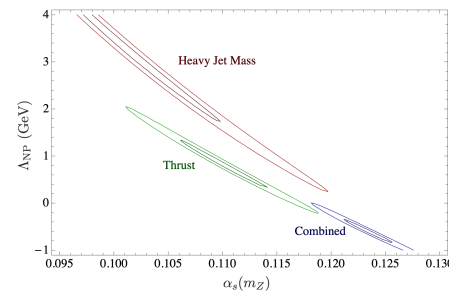
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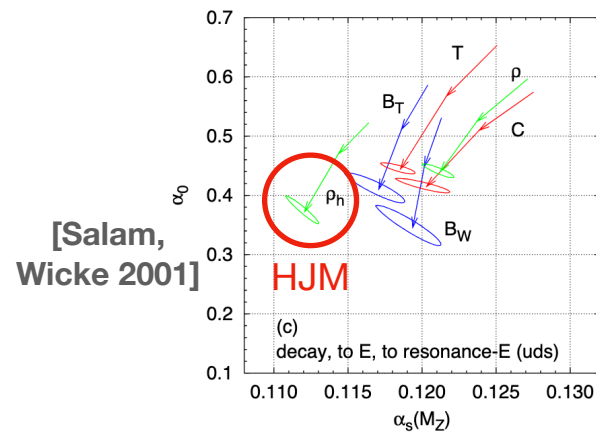
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At that time, fixed-order $\mathcal{O}(\alpha_s^3)$ results became available

Impossible to have consistent results for HJM!

α_s from Event Shapes

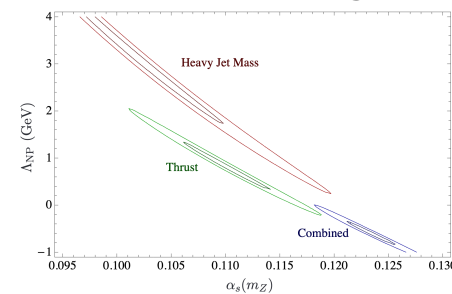
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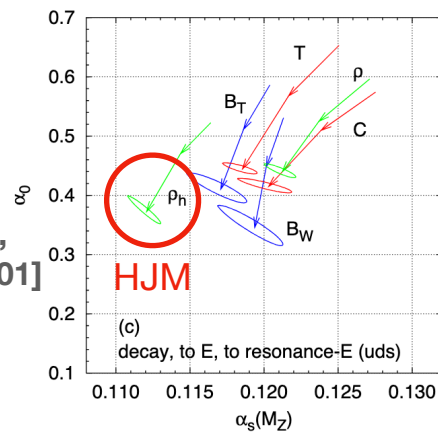
[Hoang, Kolodrubetz, Mateu, Stewart 2015]

very precise value, but well below the world average

α_s from Event Shapes

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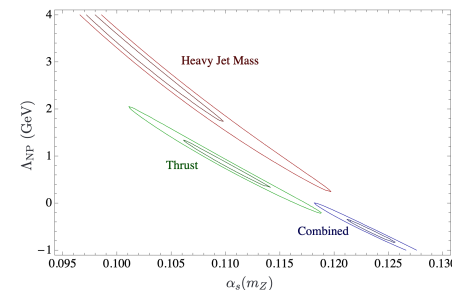
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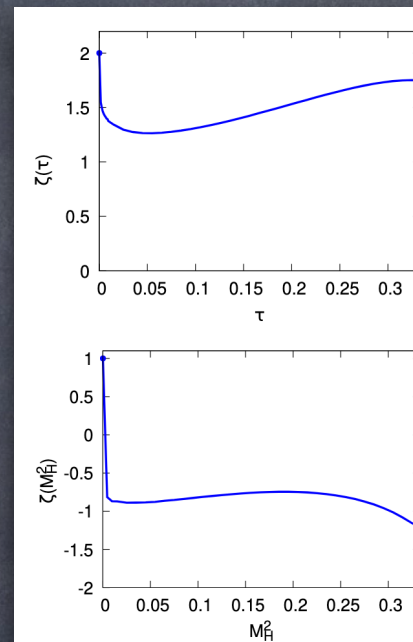


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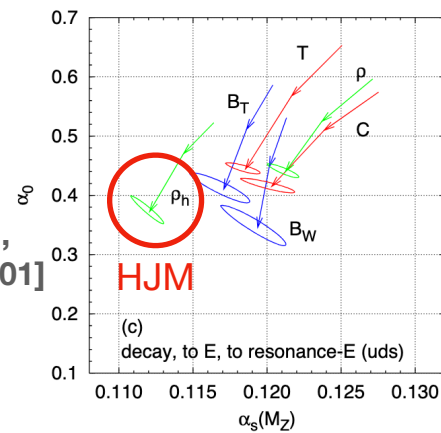
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renormalon-based "computation" of those

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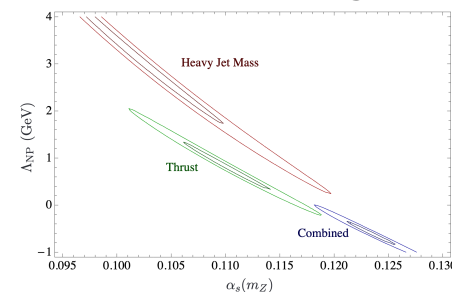
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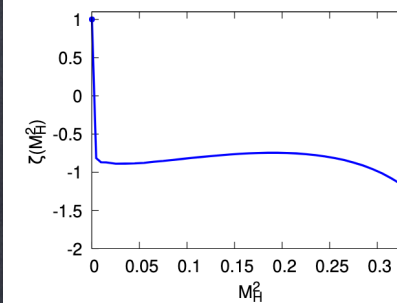
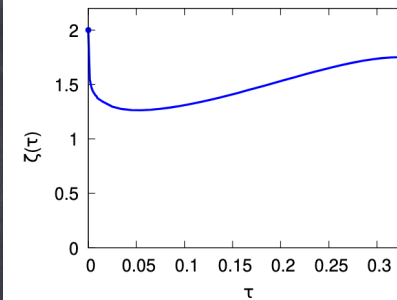


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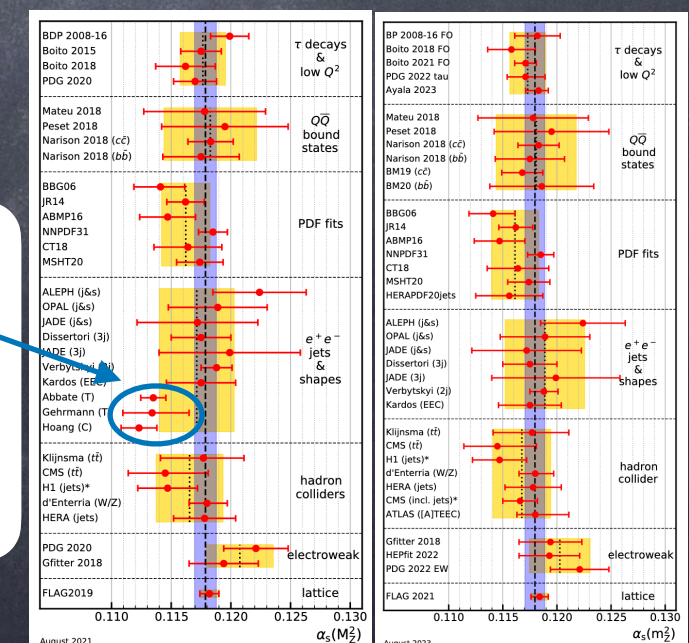
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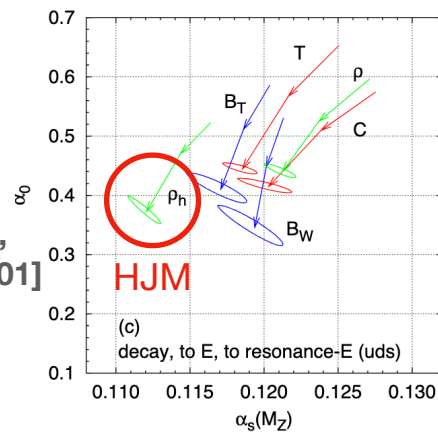
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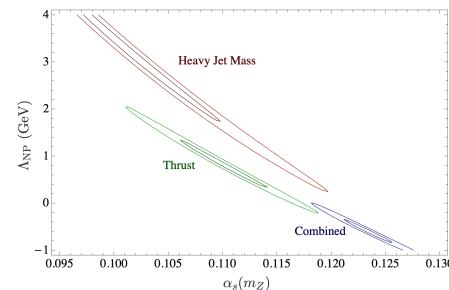
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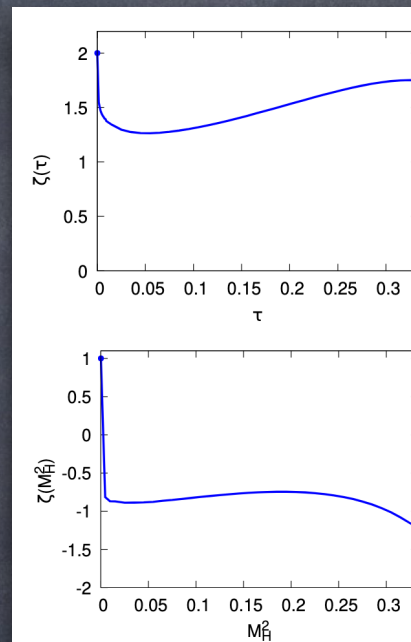


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2024

On Determining $\alpha_s(m_Z)$ from Dijets in e^+e^- Thrust

Miguel A. Benitez^a, André H. Hoang^b, Vicent Mateu^a, Iain W. Stewart^{b,c} and Gherardo Vita^d

pure dijet analysis still yields small α_s

power corrections were not the issue!

2001-2010

2010-2015

2023

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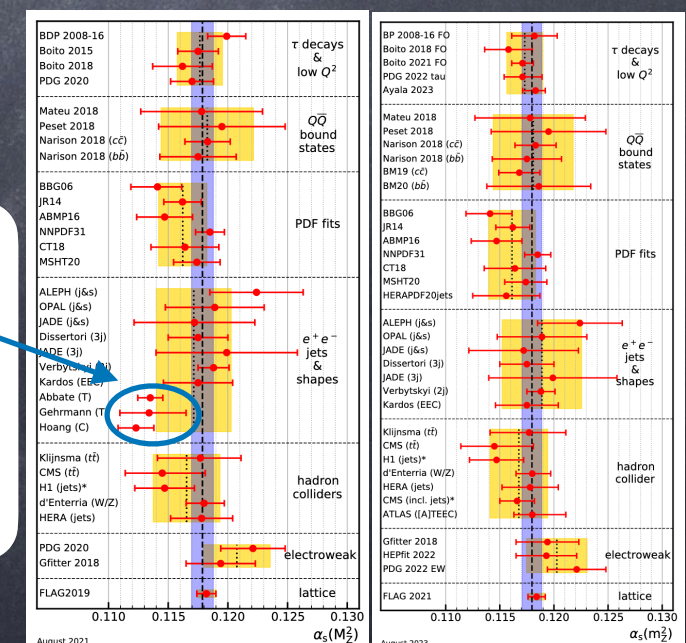
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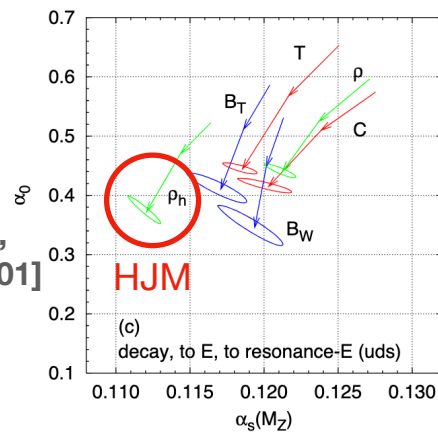
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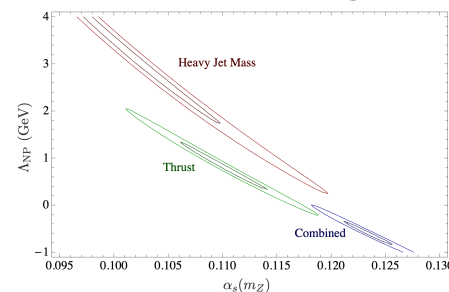
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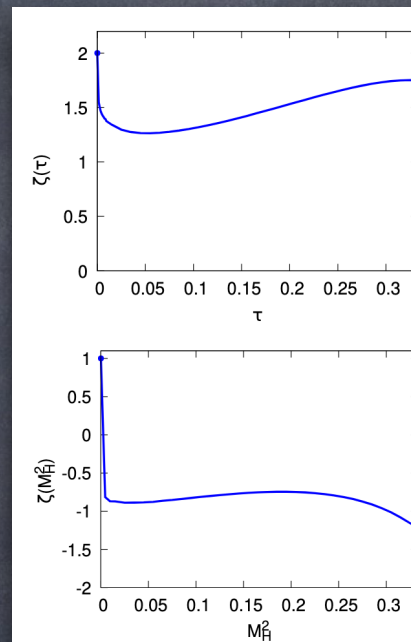


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A Precise Determination of α_s from the Heavy Jet Mass Distribution

Miguel A. Benitez¹, Arindam Bhattacharya², André H. Hoang³, Vicent Mateu¹, Matthew D. Schwartz², Iain W. Stewart^{3,4} and Xiaoyuan Zhang²

This talk

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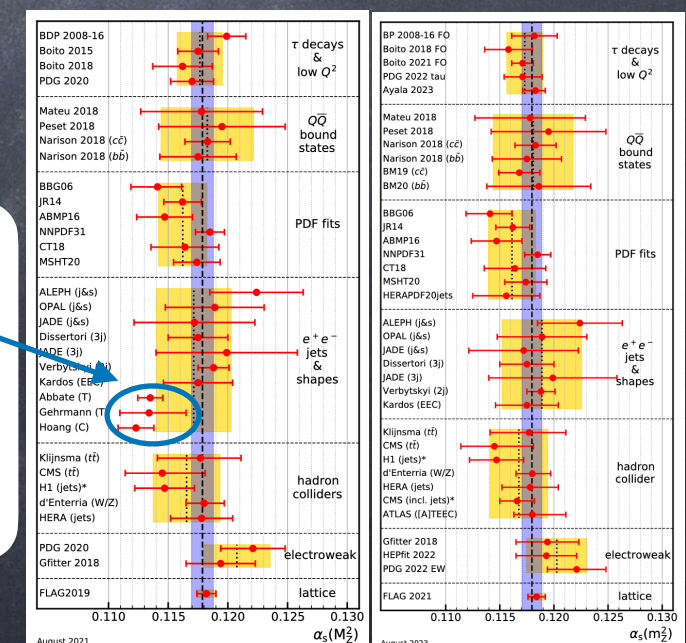
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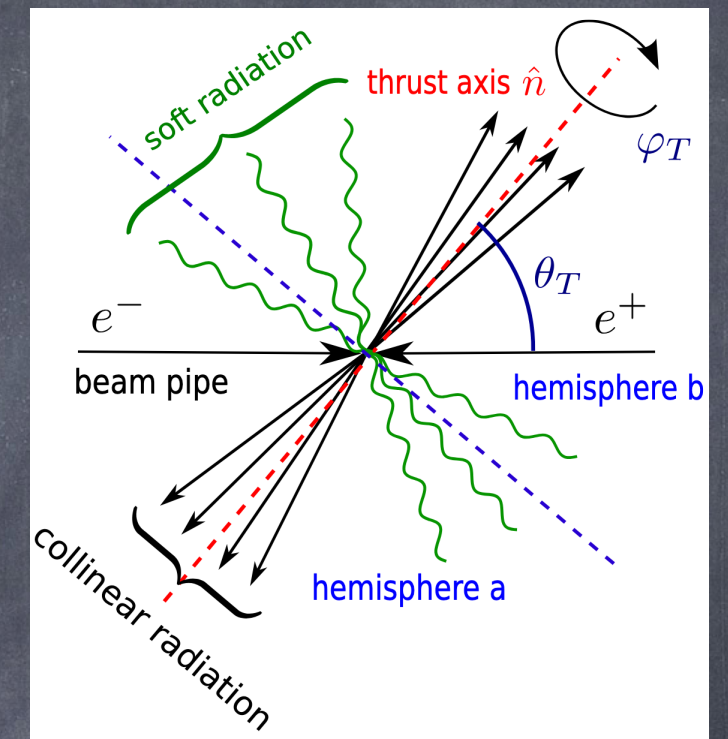
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The HJM distribution

Event shapes describe geometric properties of final-state particles

Anatomy of dijet event



The HJM distribution

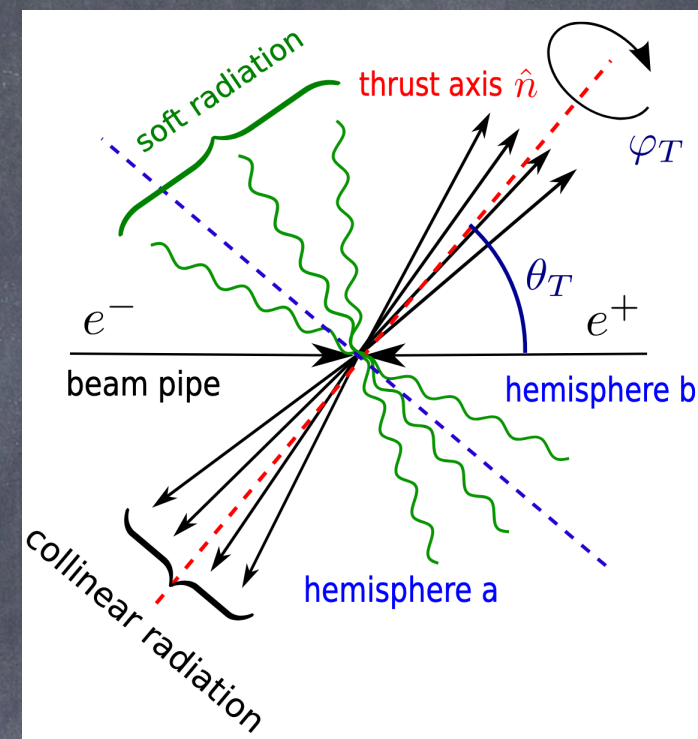
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thrust axis defined as vector that maximises

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plane normal to $\hat{n}$ defines two hemispheres

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hemispheres
invariant masses

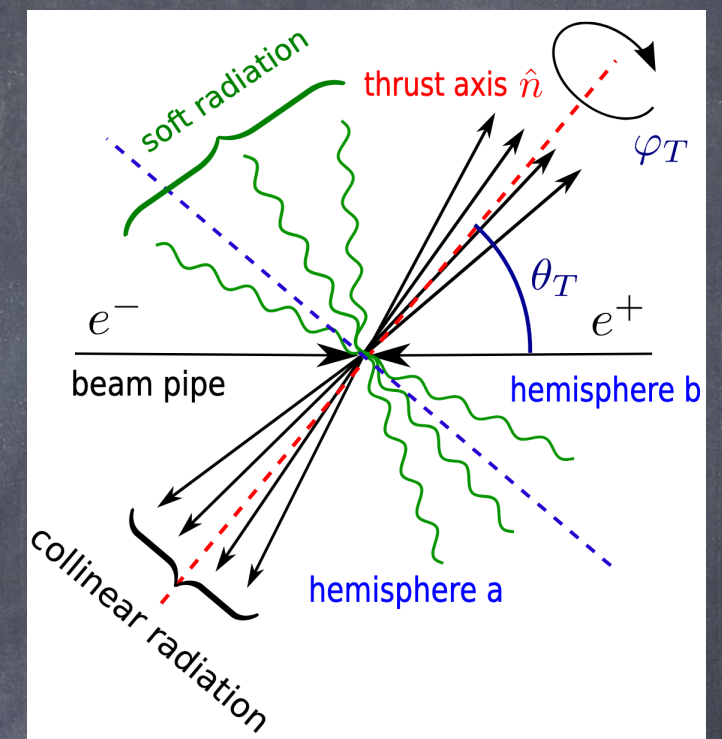
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heavy jet mass $\rho = \frac{1}{Q^2} \max(s_a, s_b)$

[Clavelli 1979] [Chandramohan, Clavelli 1981] [Clavelli, Wyler 1981]

$Q = \sqrt{(p_{e+} + p_{e-})^2}$ c.o.m. energy

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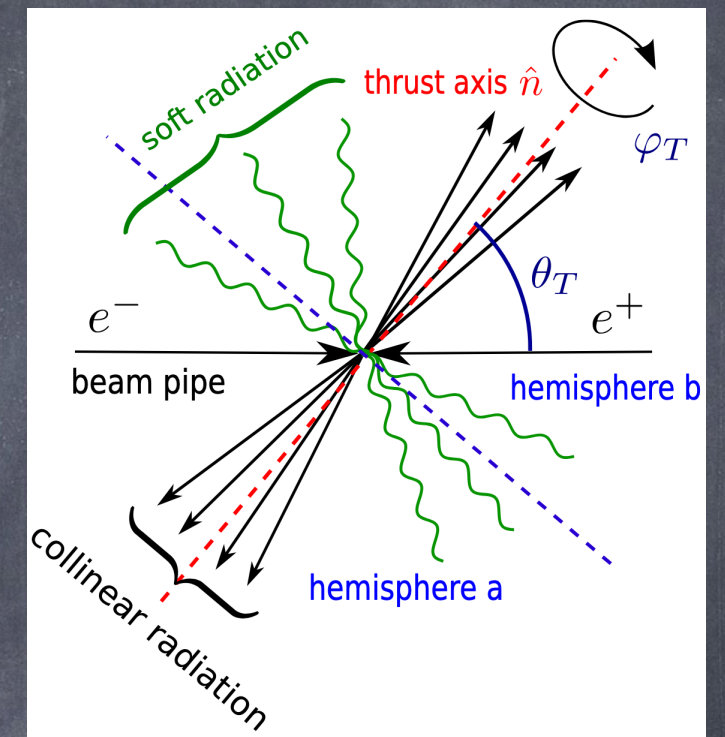
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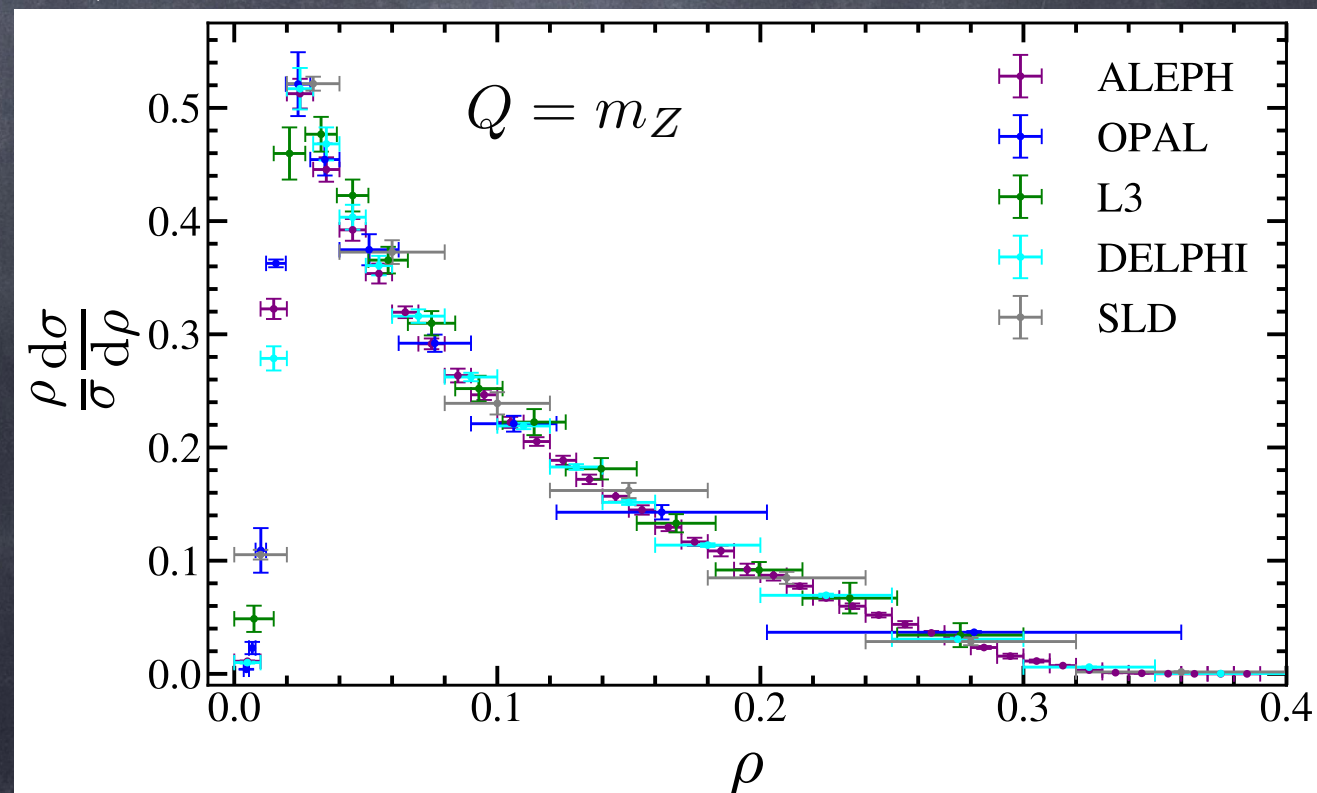
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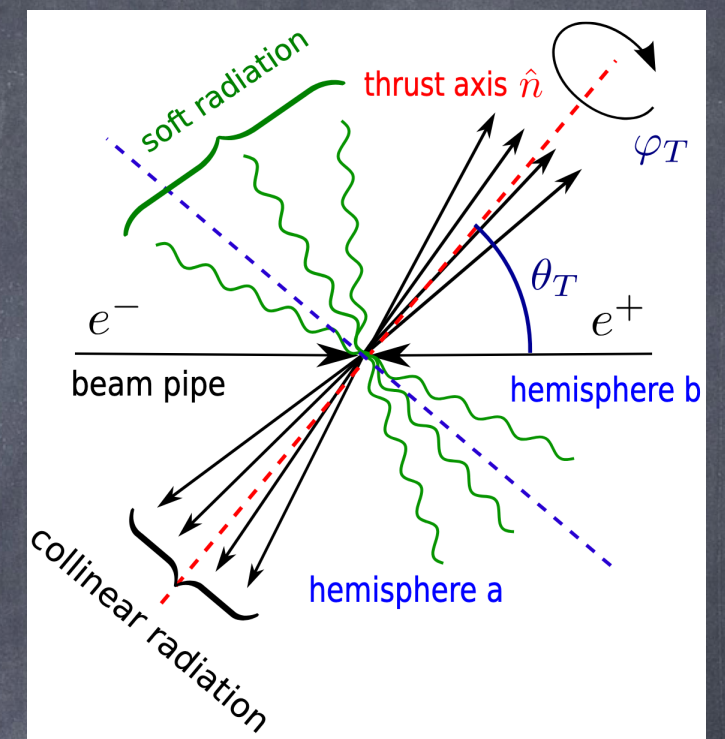
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Peak region

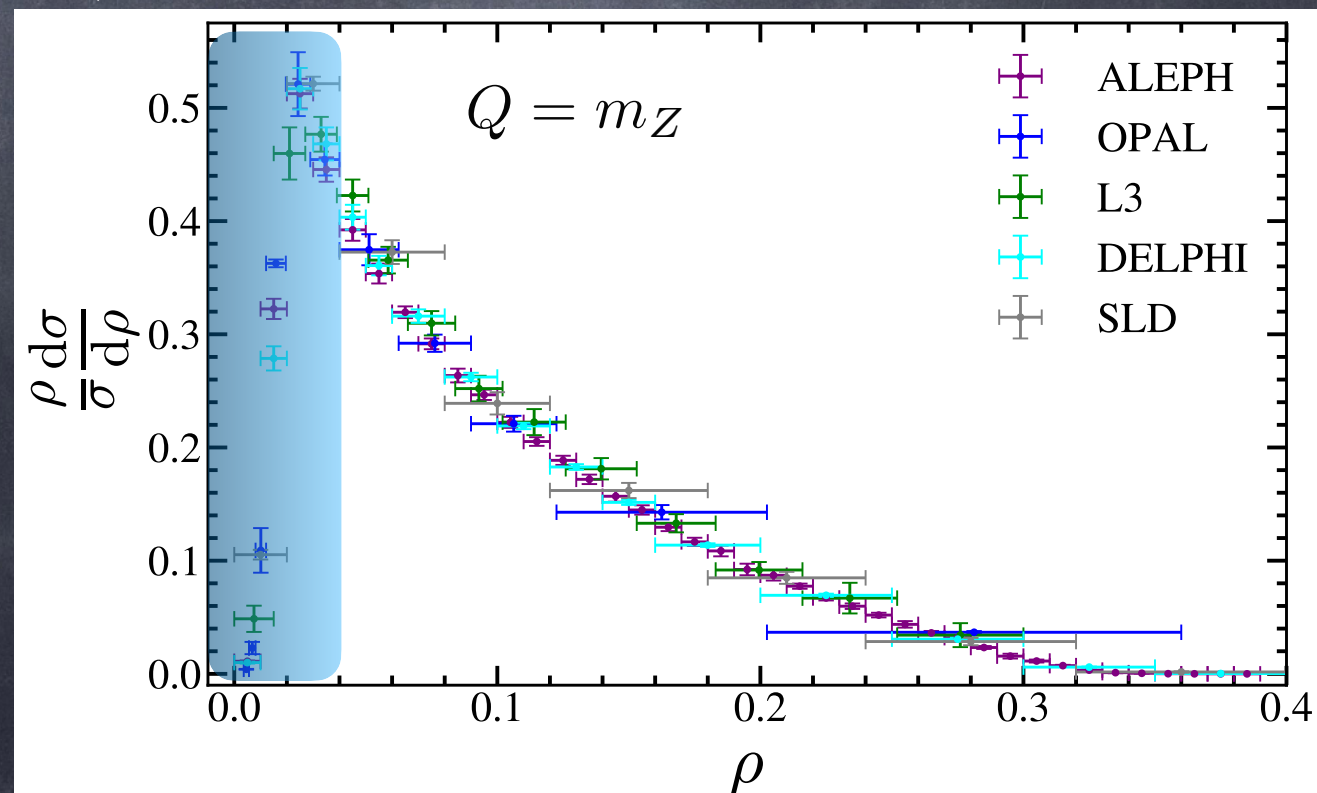
$$\frac{s_i}{Q} \sim \Lambda_{\text{QCD}}$$

Heavily affected by
non-perturbative
dynamics

Anatomy of dijet event



experimental distribution



The HJM distribution

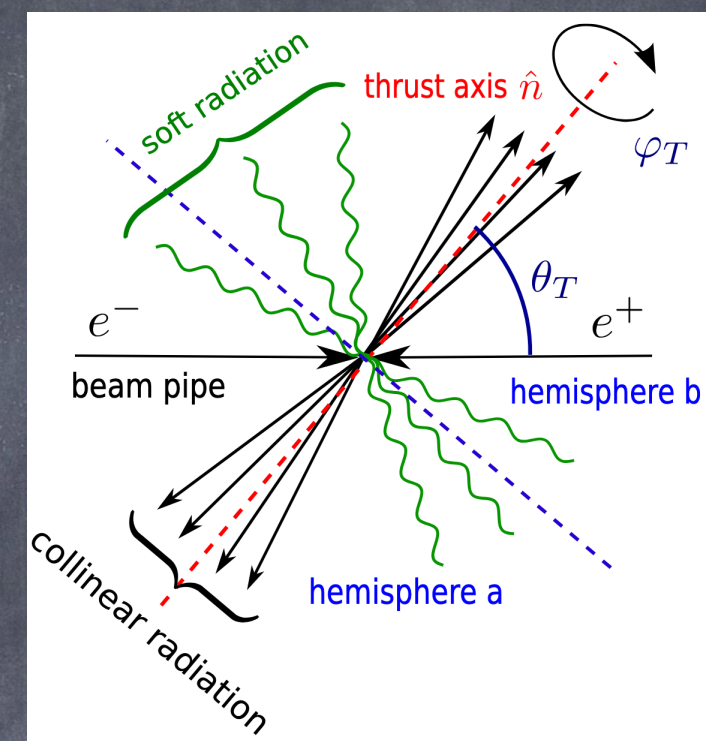
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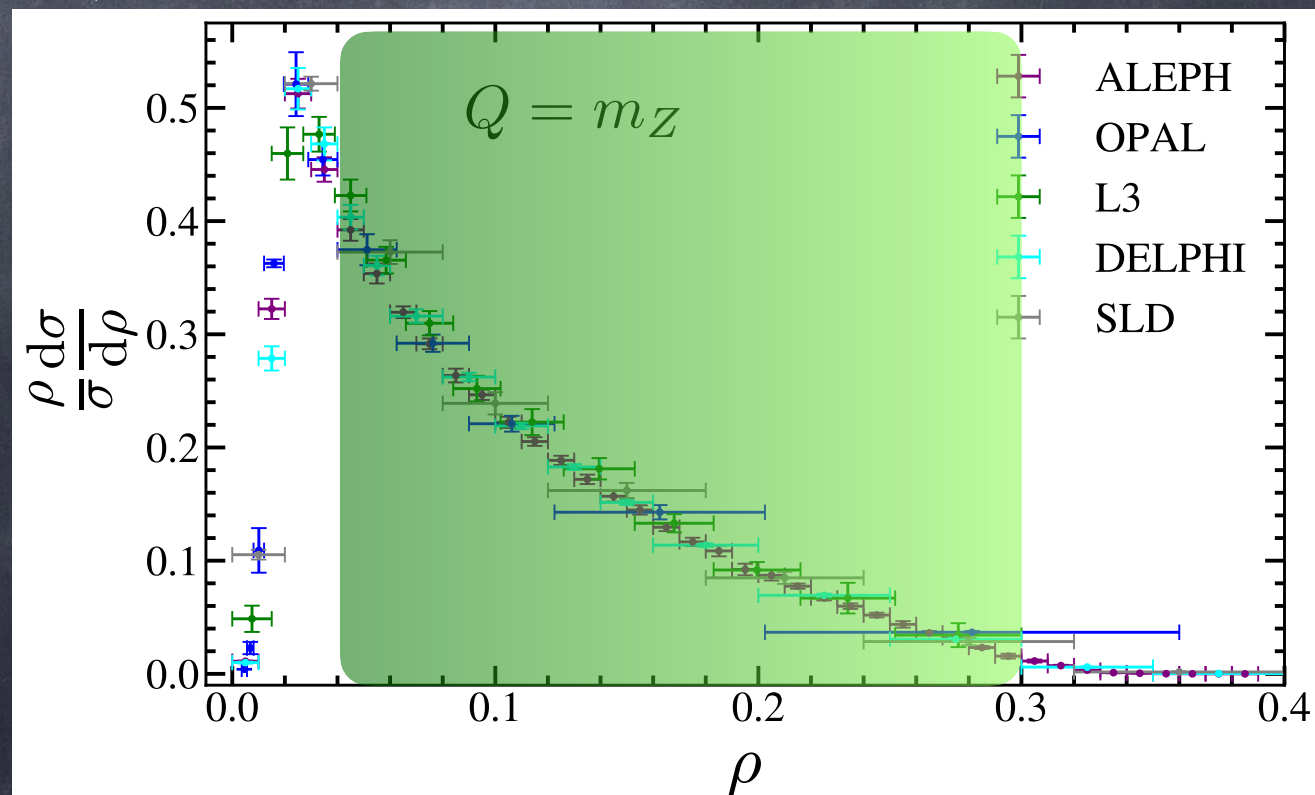
Tail region

$$\Lambda_{\text{QCD}} \ll \frac{s_i}{Q} \ll Q$$

Soft and collinear radiation dominate

clear separation between perturbative and non-perturbative modes

experimental distribution



The HJM distribution

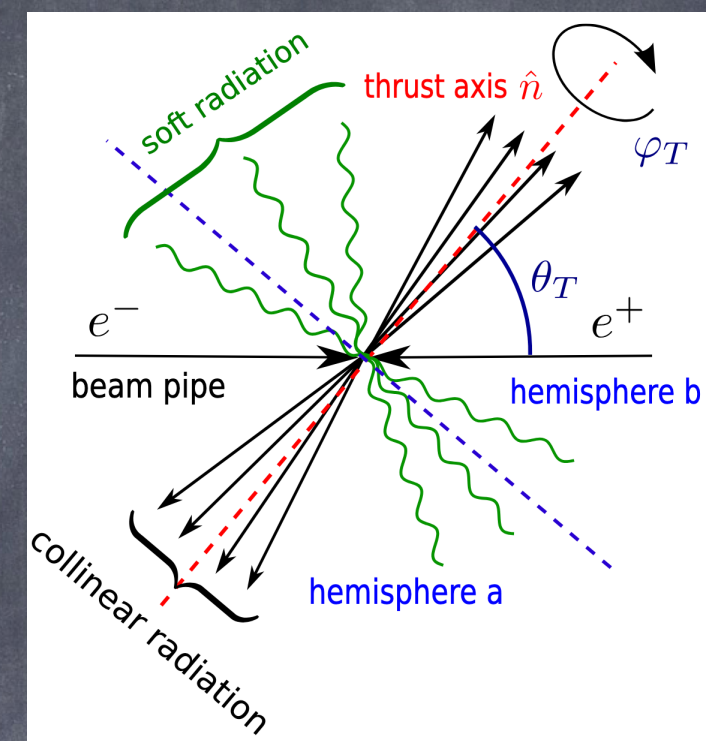
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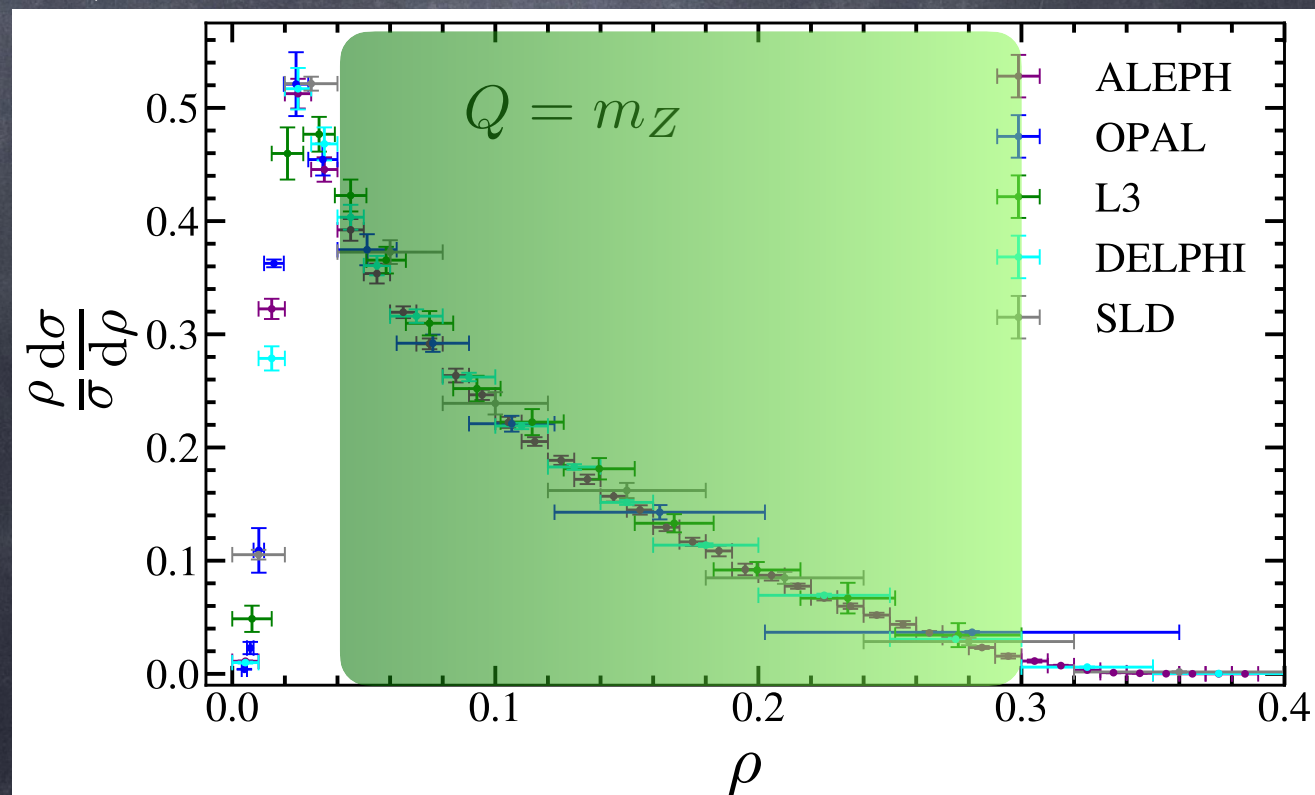
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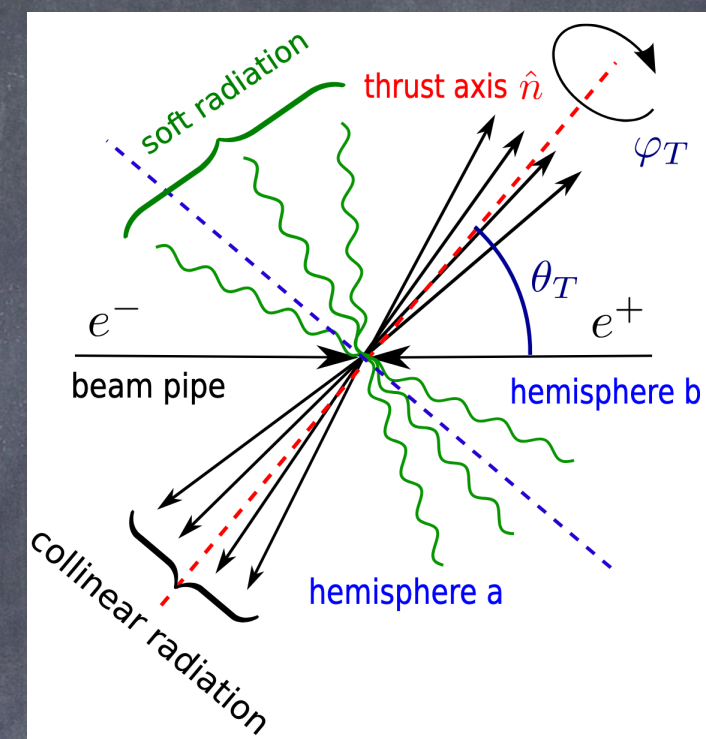
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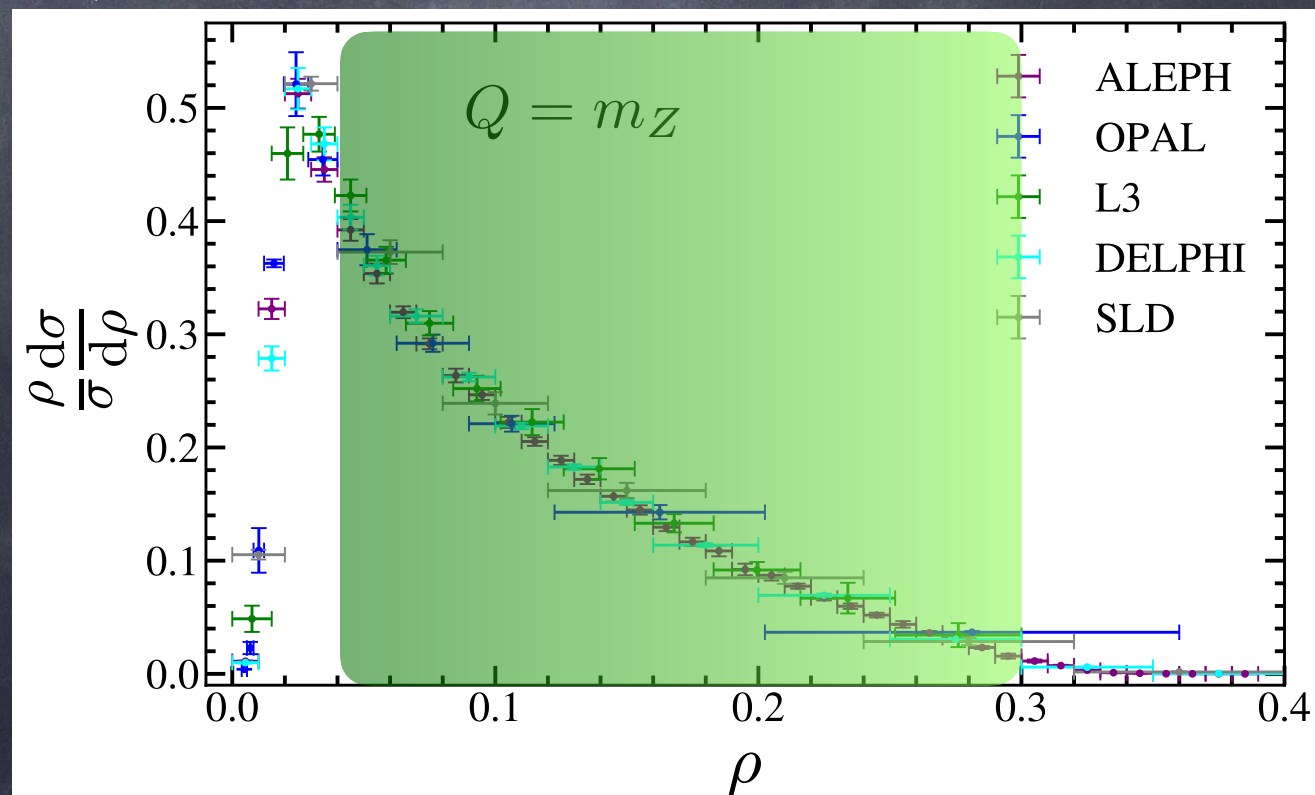
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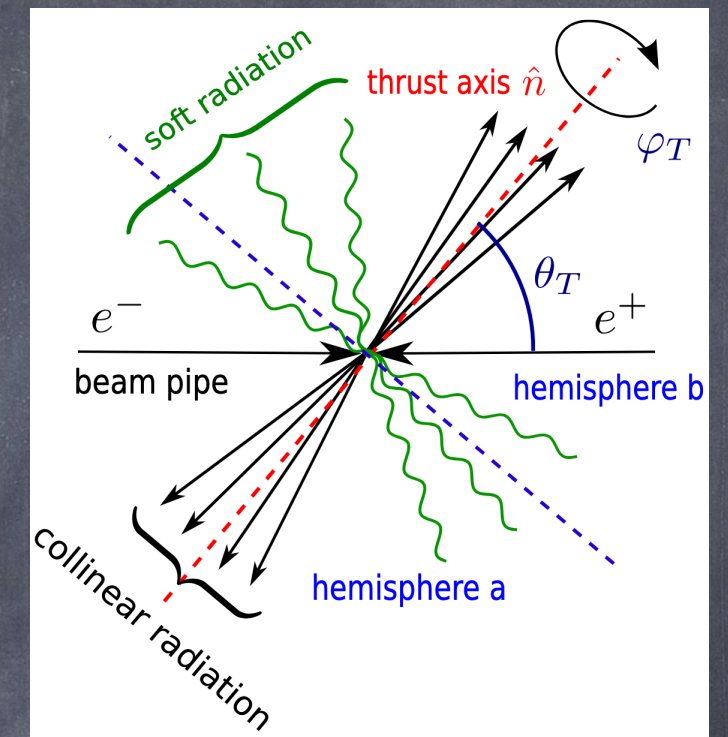
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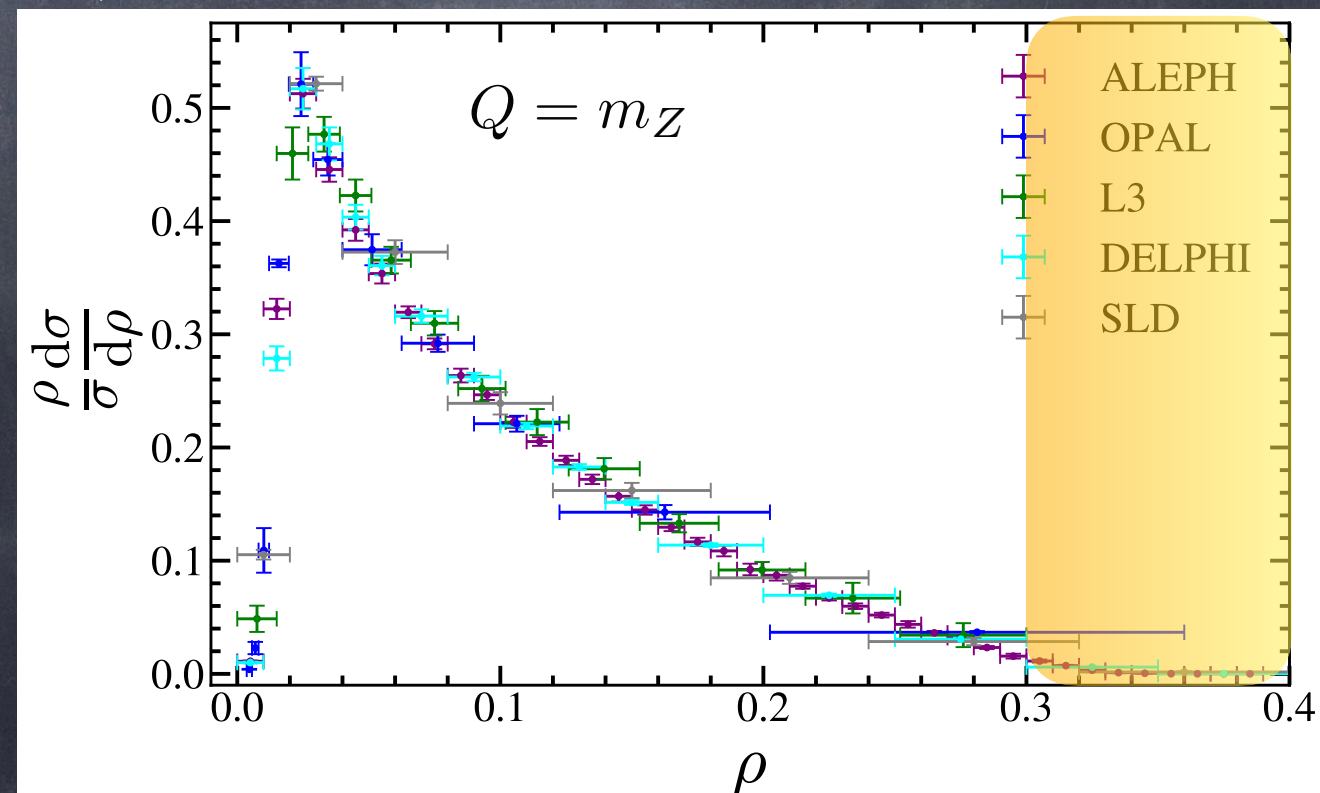


Far-tail region

$$s_i \sim Q^2$$

No clear distinction between three scales
right description: FO
small hadronization corrections

experimental distribution



The HJM distribution

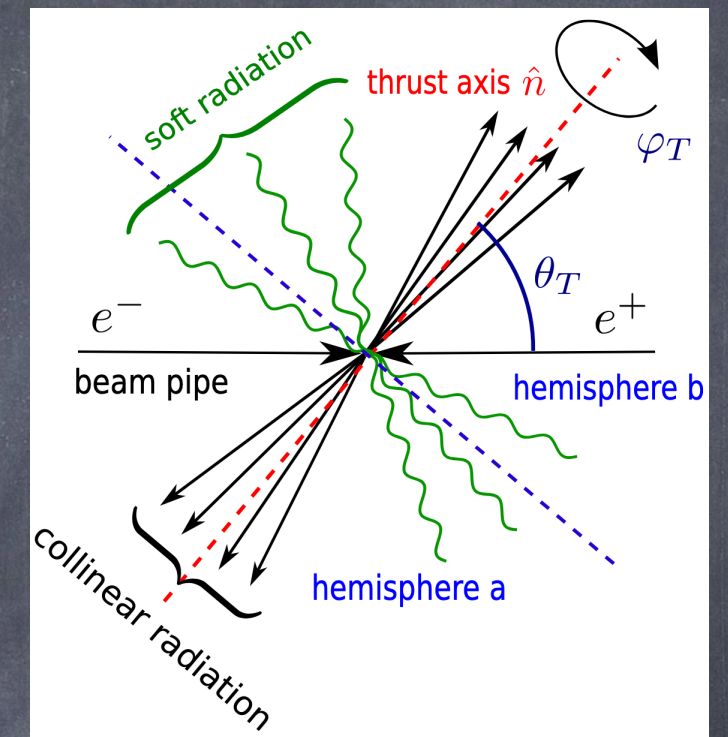
Event shapes describe geometric properties of final-state particles

thrust axis defined as vector that maximises

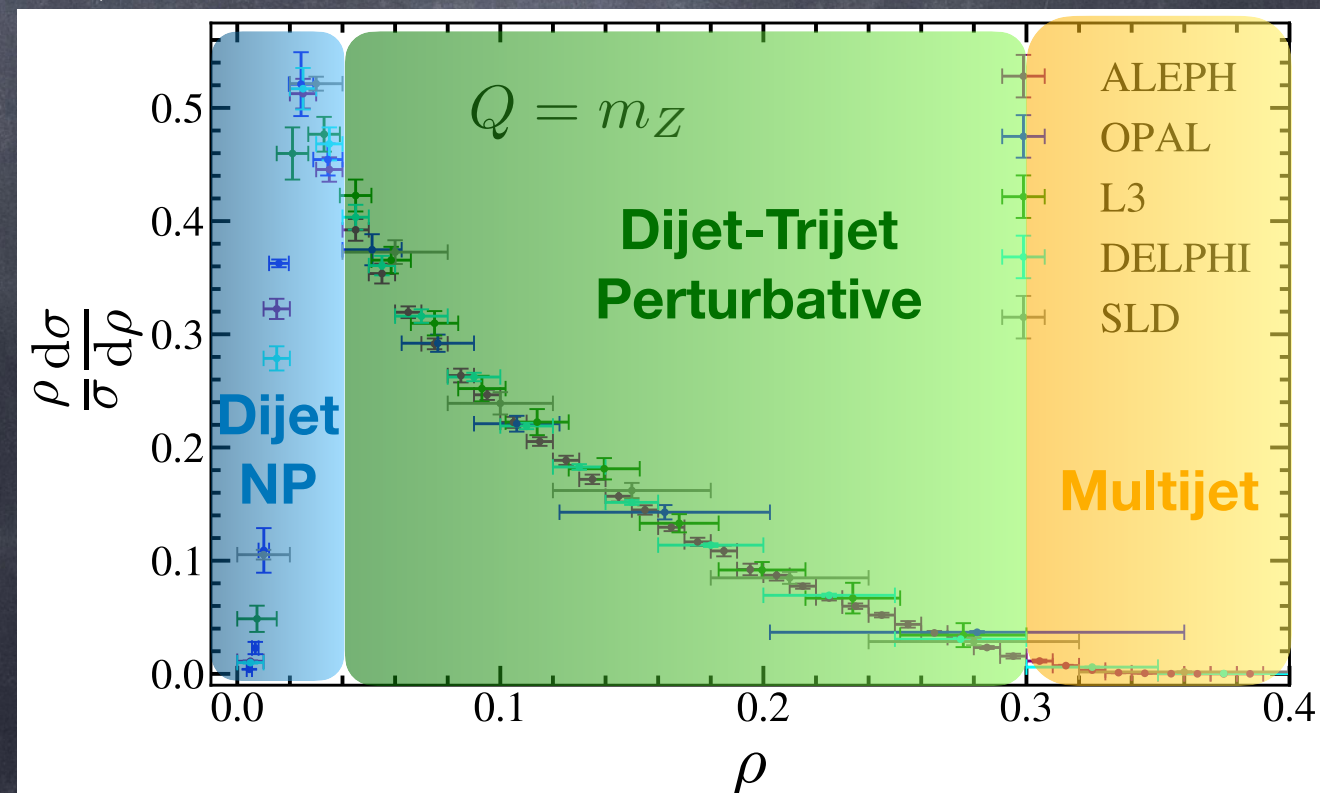
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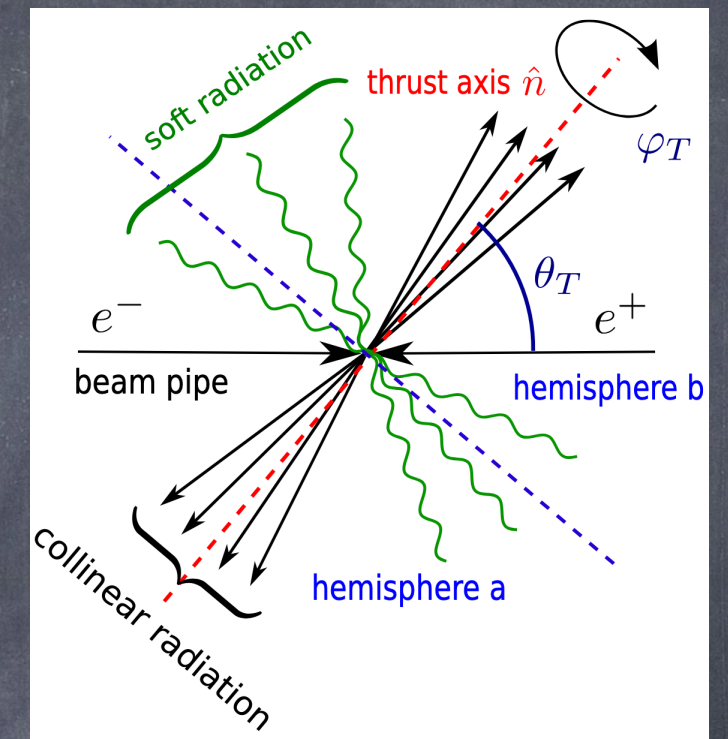
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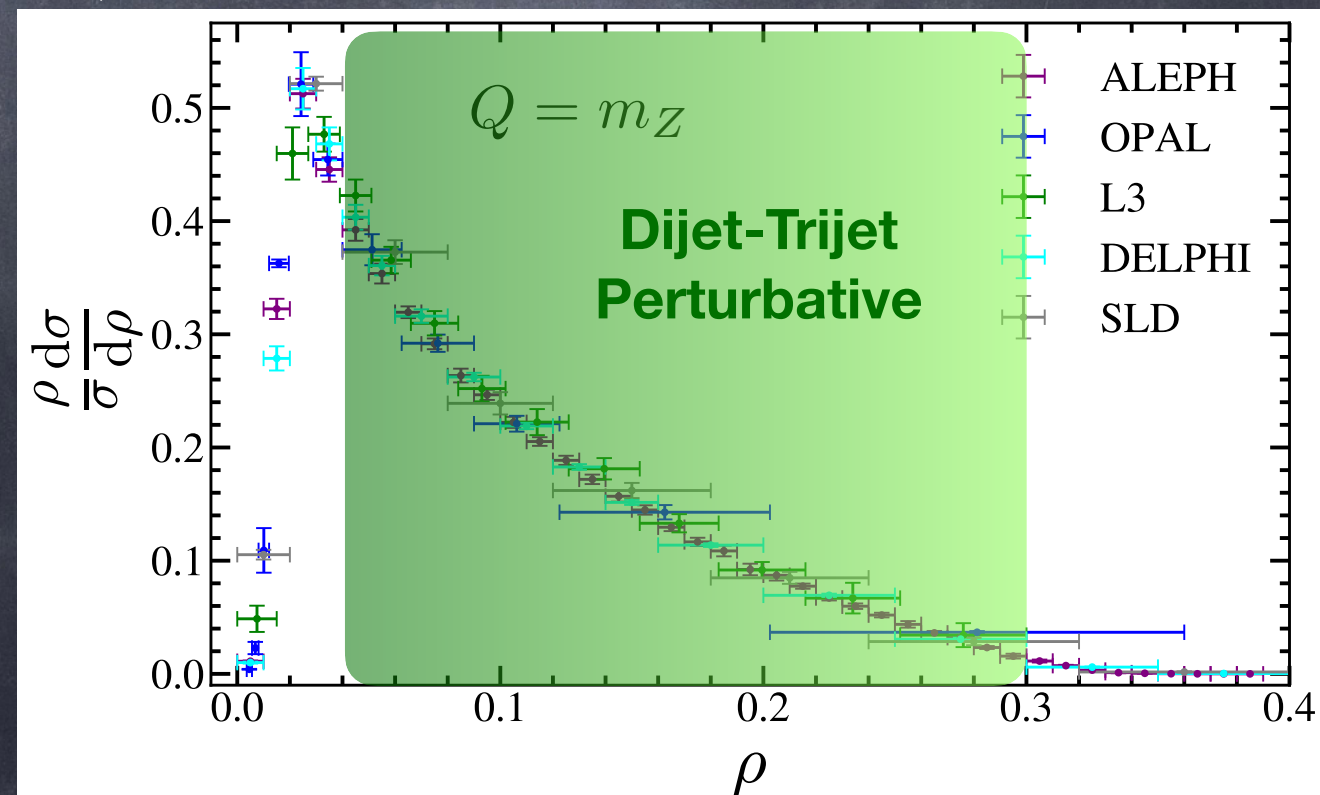
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Anatomy of dijet event



experimental distribution



Ideal place for fits

HJM distribution — dijet

Dijet resummation @ N³LL: [A.H. Hoang, VM, I.W. Stewart & M.D. Schwartz, 2025] along with first-principles power corrections

HJM obtained as marginalisation of dihemsphere mass distro

$$\frac{d\sigma}{d\rho} = 2Q^2 \int_0^{Q^2\rho} ds \frac{d^2\sigma}{ds_1 ds_2} (s_1 = Q^2\rho, s_2 = s) \equiv 2Q^2 \Xi(Q^2\rho, Q^2\rho)$$

semi-cumulative

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Hadronic cross section obtained as 2D convolution:

$$\frac{d\sigma}{d\rho} = 2Q^2 \int dk_1 dk_2 \hat{\Xi}(Q^2\rho - Qk_1, Q^2\rho - Qk_2) F(k_1, k_2)$$

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partonic semi-cumulative

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partonic semi-cumulative 2D shape function

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partonic semi-cumulative 2D shape function

In tail region gives raise to non-trivial OPE in terms of

$$\Omega_{ij} = \Omega_{ji} = \int dk_1 dk_2 k_1^i k_2^j F(k_1, k_2)$$

Leading power: $\frac{d\sigma_{\text{dij}}}{d\rho} = \frac{d\hat{\sigma}_{\text{dij}}}{d\rho} \left(\rho - \frac{\Omega_{10}}{Q} \right)$

HJM distribution – dijet factorization

Factorisation of singular partonic distribution is also 2D

$$\frac{1}{\sigma_0} \frac{d^2 \hat{\sigma}_{\text{sing}}}{ds_1 ds_2} = H(Q, \mu) \int d\ell_1 d\ell_2 J(s_1 - Q \ell_1 - Q \bar{\Delta}(R, \mu), \mu) \\ \times J(s_2 - Q \ell_2 - Q \bar{\Delta}(R, \mu), \mu) e^{\delta(R, \mu) \left(\frac{\partial}{\partial \ell_1} + \frac{\partial}{\partial \ell_2} \right)} \hat{S}(\ell_1, \ell_2, \mu)$$

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Involves one hard factor, two jet functions and one soft function

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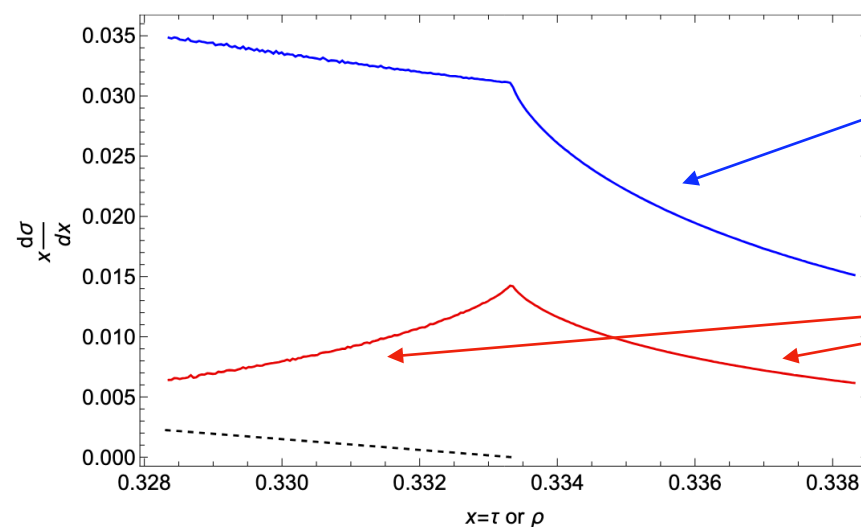
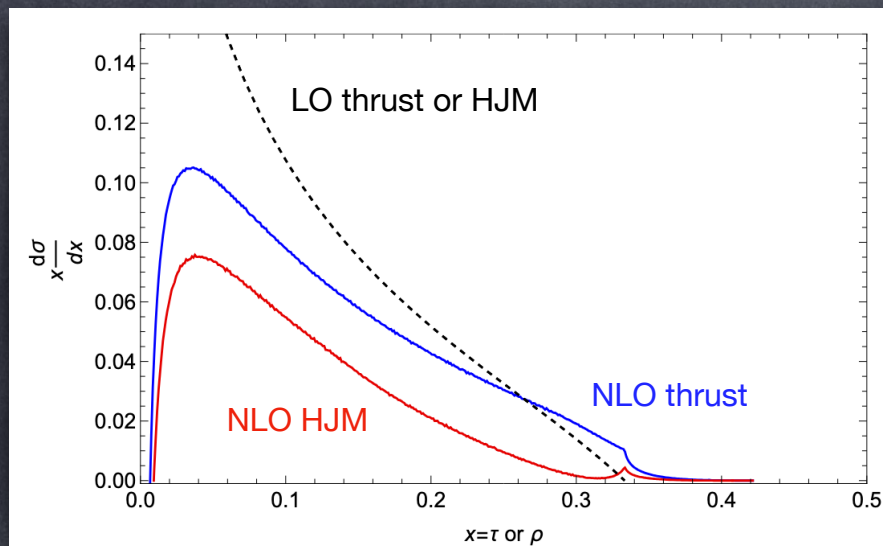
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RGE between different matrix elements sums up large logs

Non-global logs in soft function makes this non-trivial

HJM distribution – dijet + shoulder resummation

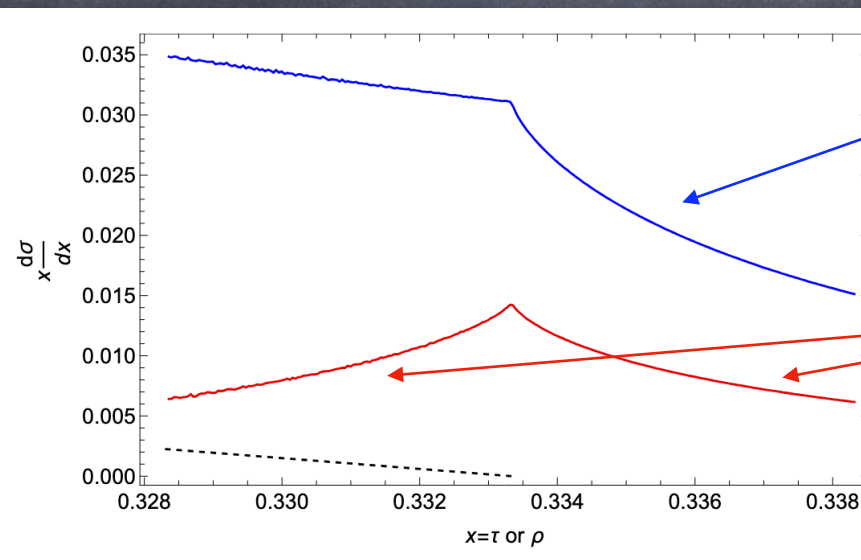
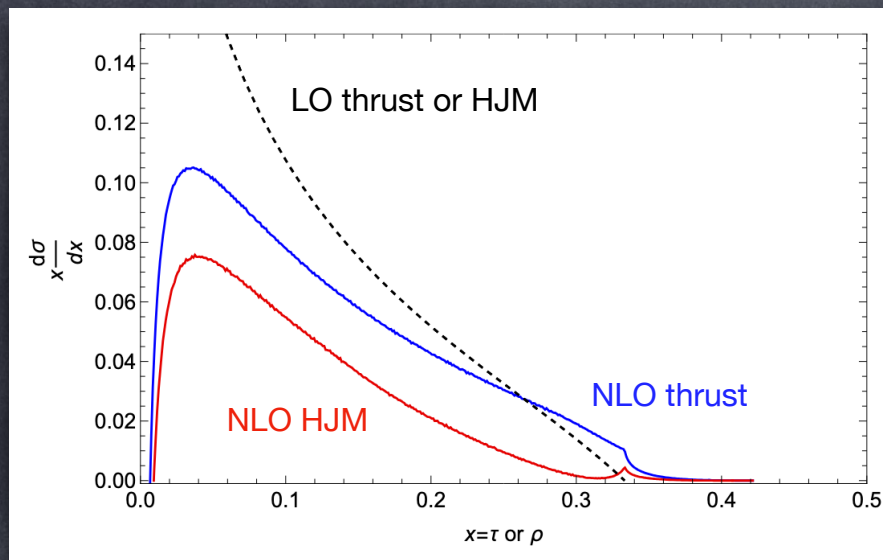


Thrust has **right** shoulder only

HJM has **left** and **right** shoulders

[Bhattacharya, Schwartz, Zhang 2022]

HJM distribution – dijet + shoulder resummation



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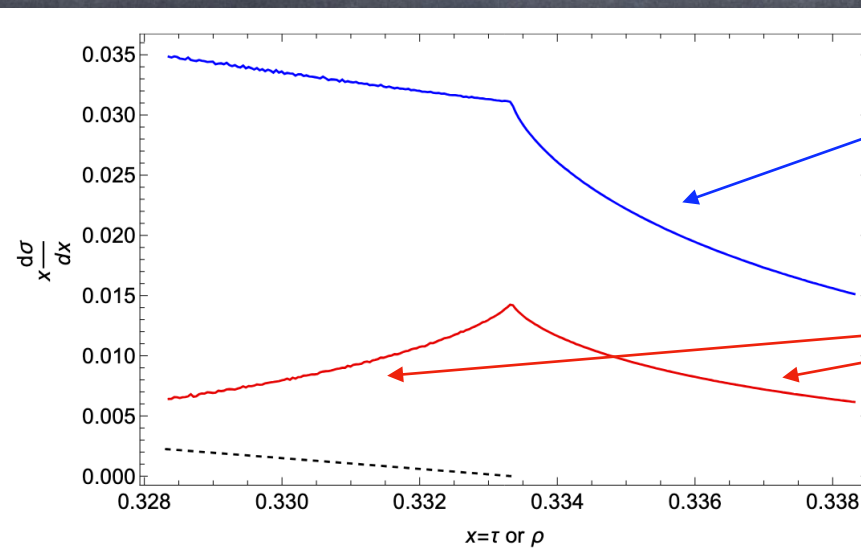
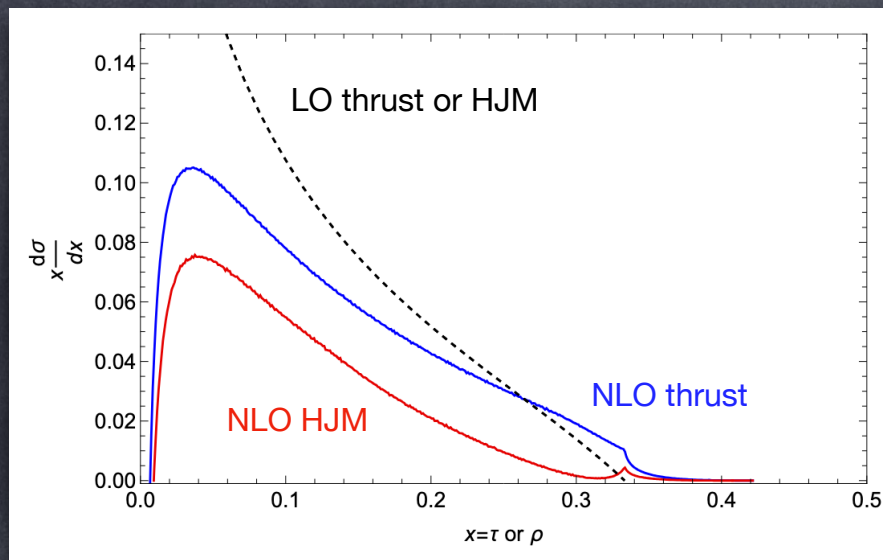
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Around symmetric trijet limit $\rho \rightarrow 1/3$, distribution factorizes as

$$d\hat{\sigma}_{\text{sh}} = H_{\text{sh}} \times J_1 \otimes J_2 \otimes J_3 \otimes S_{1,2,3}$$

[Bhattacharya, Michel, Schwartz, Stewart, Zhang 2023]

HJM distribution – dijet + shoulder resummation



[Bhattacharya, Schwartz, Zhang 2022]

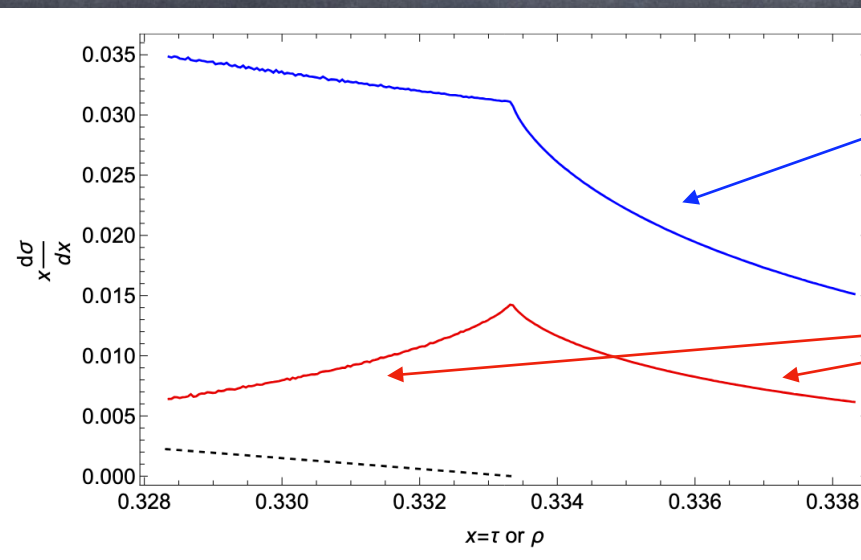
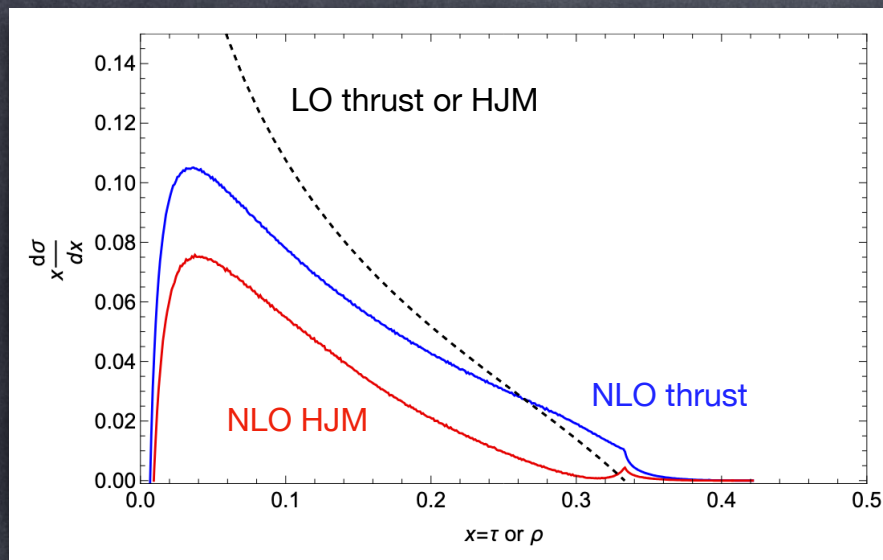
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Matching between dijet, FO & shoulder writing full cross section as

$$d\hat{\sigma} = [d\hat{\sigma}_{\text{dij}} - d\hat{\sigma}_{\text{dij}}^{\text{sing}}] + [d\hat{\sigma}_{\text{sh}} - d\hat{\sigma}_{\text{sh}}^{\text{sing}}] + d\sigma_{\text{FO}}$$

HJM distribution – dijet + shoulder resummation



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Parametrize 3-jet power corrections as

$$\frac{d\sigma_{\text{sh}}}{d\rho} = \frac{d\hat{\sigma}_{\text{sh}}}{d\rho} \left(\rho - \frac{\Theta_1}{Q} \right)$$

Fit procedure

Use χ^2 function including theoretical and experimental uncertainties

Experiment

$35 \text{ GeV} < Q < 207 \text{ GeV}$
(700 datapoints)

Minimal Overlap Model for
systematic uncertainties

$$\sigma_{ij}^{\text{exp}} = \delta_{ij} (\Delta_i^{\text{stat}})^2 + \delta_{D_i D_j} \min(\Delta_i^{\text{sys}}, \Delta_j^{\text{sys}})^2$$

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Theory

pert. uncertainties assessed with
renormalization scale variation
not Gaussian + highly correlated

flat random scan: 5000 sets of 17
parameters, yielding theory &
uncertainty for each data-point x_i

$$\bar{x}_i = \frac{x_i^{\text{max}} + x_i^{\text{min}}}{2} \quad \Delta_i^{\text{theo}} = \frac{x_i^{\text{max}} - x_i^{\text{min}}}{2}$$

correlation $r_{ij}^{\text{theo}} = \frac{\langle (x_i - \bar{x}_i)(x_j - \bar{x}_j) \rangle}{\sqrt{\langle (x_i - \bar{x}_i)^2 \rangle} \sqrt{\langle (x_j - \bar{x}_j)^2 \rangle}}$

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total covariance matrix

$$\sigma_{ij}^{\text{tot}} = \sigma_{ij}^{\text{exp}} + \Delta_i^{\text{theo}} \Delta_j^{\text{theo}} r_{ij}^{\text{theo}}$$

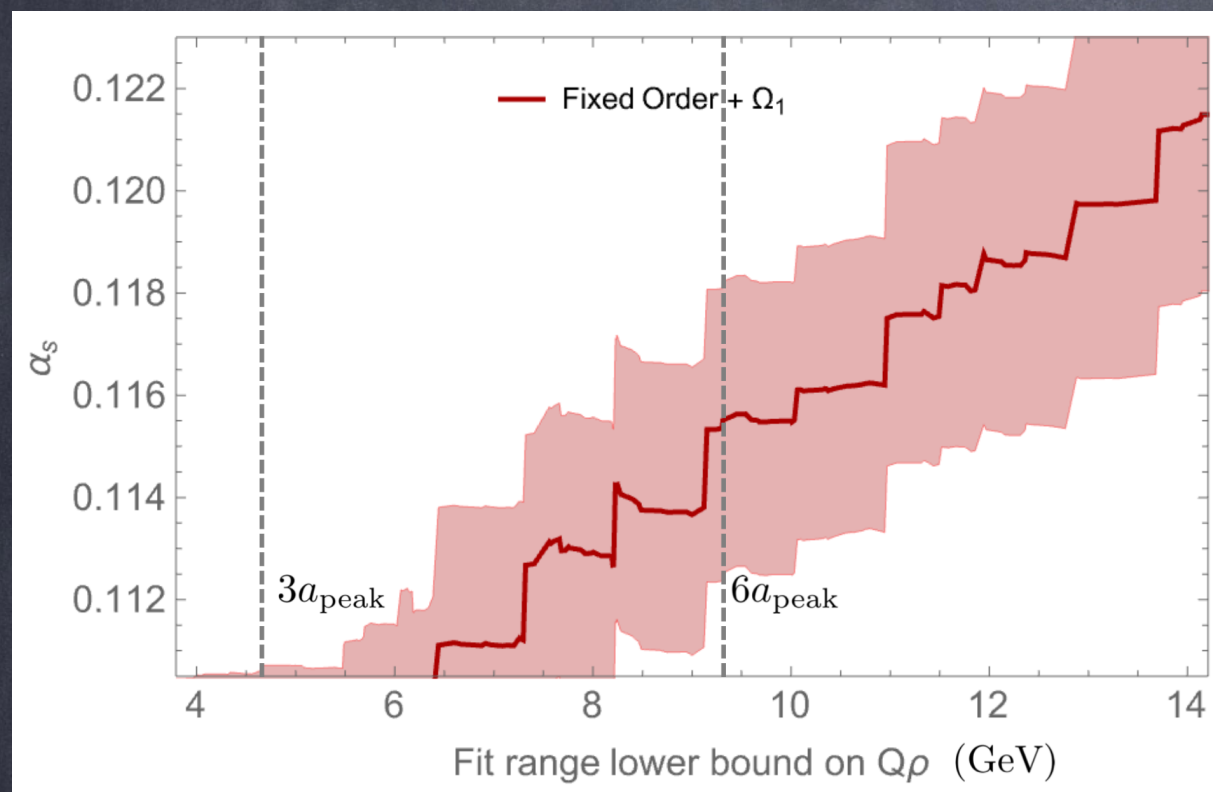


$$\chi^2 = \sum_{i,j=1}^{N_{\text{bins}}} (\bar{x}_i - x_i^{\text{exp}}) (\bar{x}_j - x_j^{\text{exp}}) (\sigma_{\text{tot}}^{-1})_{ij}$$

Fit results — Fixed Order

Results for α_s using fit range $\frac{a}{Q} < \rho < 0.3$ for different a

$a_{\text{peak}} \equiv \text{peak position}$



high-sensitivity to fit range

large fit uncertainty

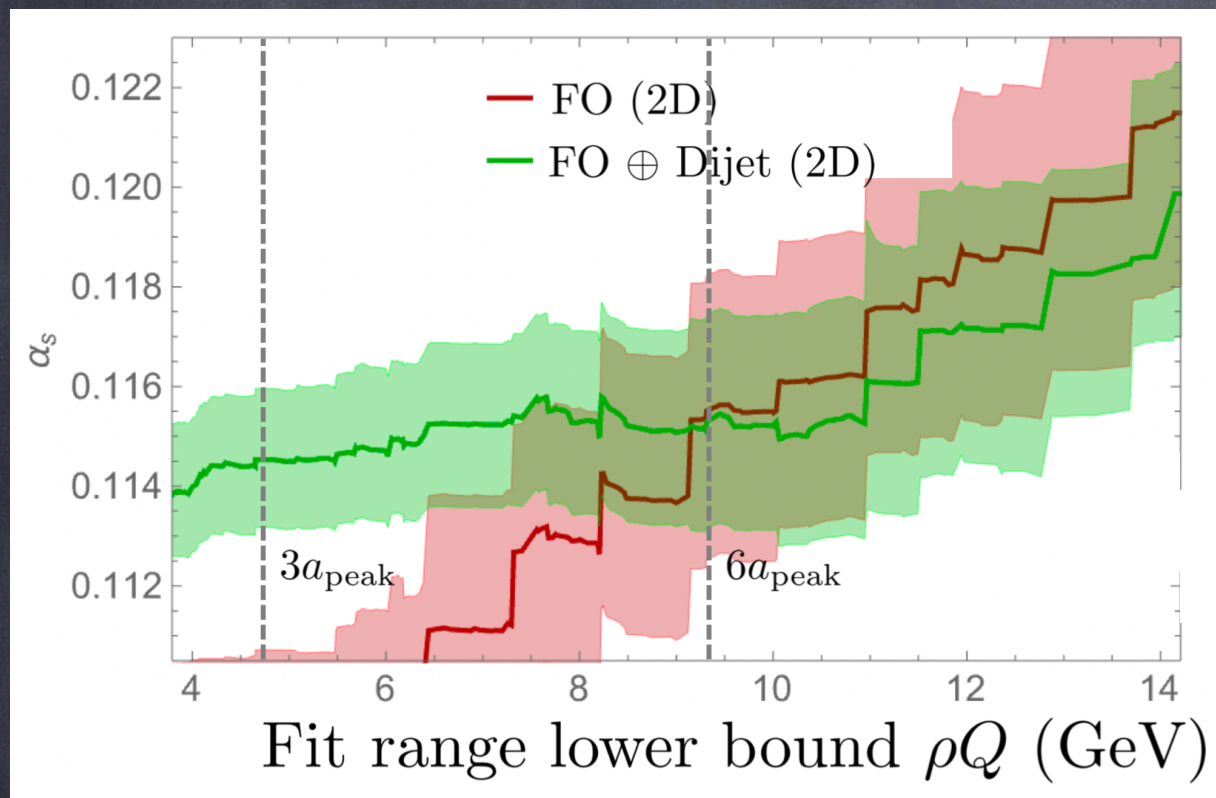
pick whichever value you like!

Model	$\alpha_s(m_Z)$	th+exp	Ω_1^p	Θ_1	fit range	χ^2/dof	Ω_1^p [GeV]	Θ_1 [GeV]
Fixed Order 2D	0.1166 ± 0.0034	± 0.0014	± 0.0027	—	± 0.0015	1.108	0.06 ± 0.13	—

Fit results – dijet resummation

Results for α_s using fit range $\frac{a}{Q} < \rho < 0.3$ for different a

$a_{\text{peak}} \equiv \text{peak position}$



low-sensitivity to fit range

small fit uncertainty

$a \in [3a_{\text{peak}}, 6a_{\text{peak}}]$

data prefers $\Omega_{10} > 0$

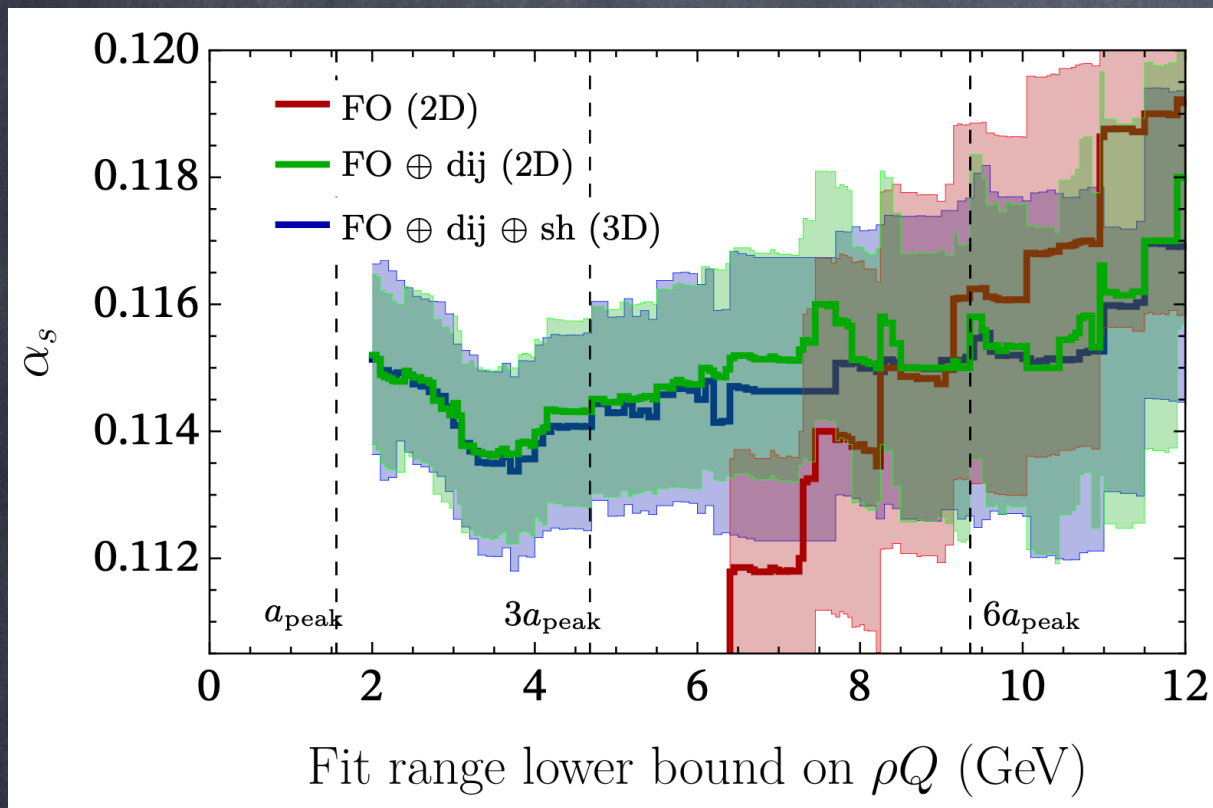
as expected by theory

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Fit results – dijet + shoulder resummation

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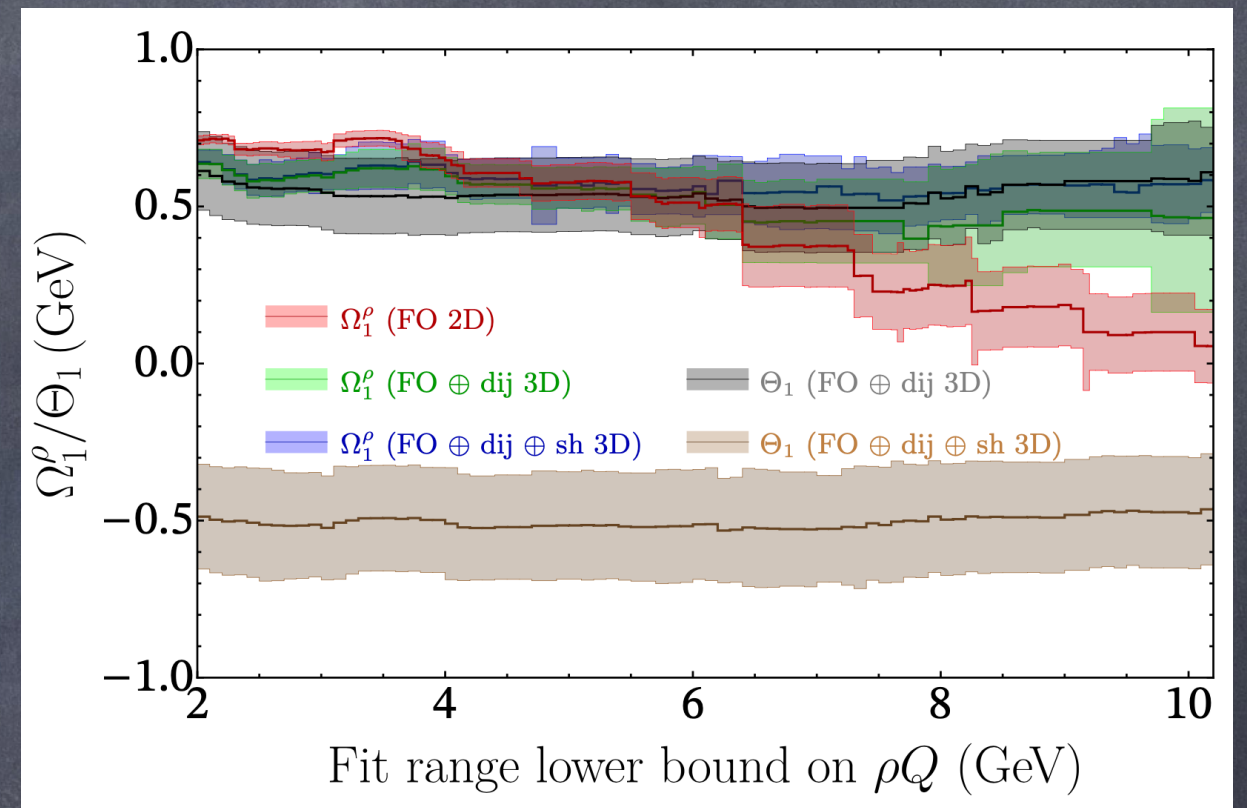
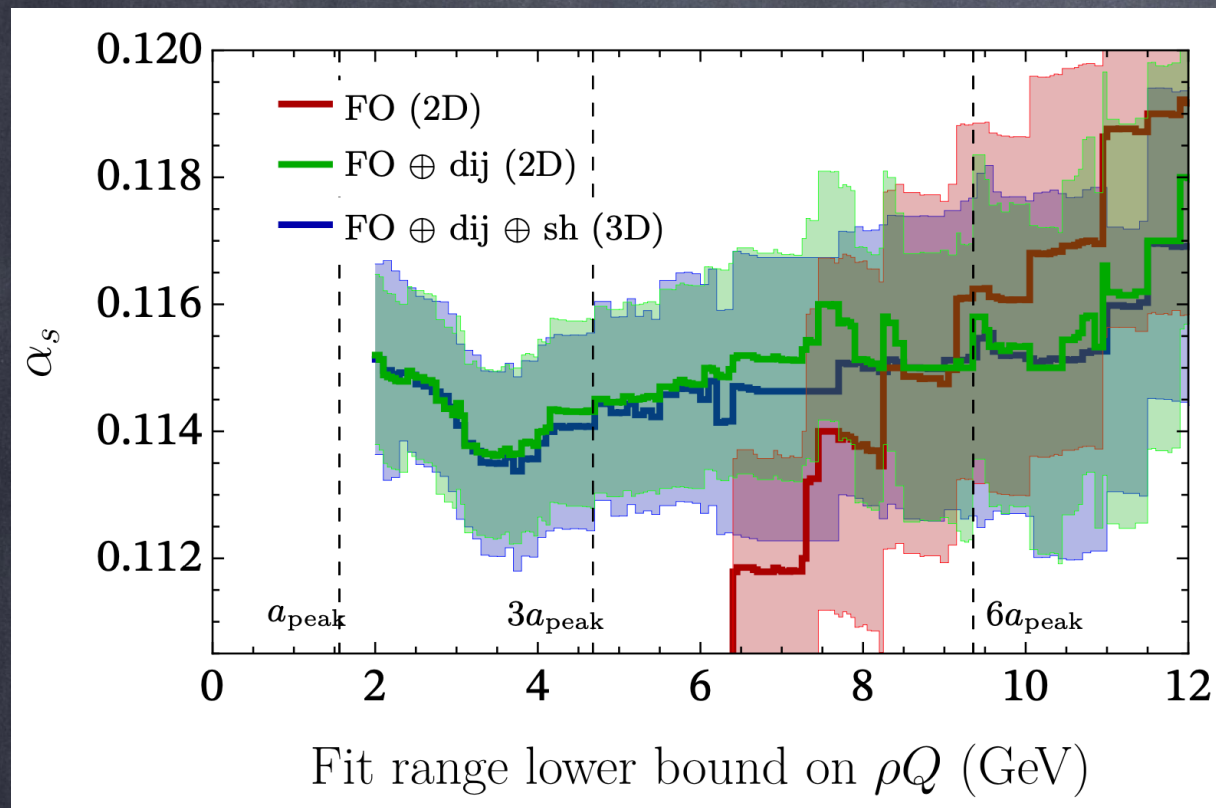
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FO + dijet 3D	0.1156 ± 0.0024	± 0.0010	± 0.0021	± 0.0004	± 0.0007	1.052	0.52 ± 0.08	0.53 ± 0.13
FO + dijet + shoulder 3D	0.1145 ± 0.0020	± 0.0009	± 0.0018	± 0.0001	± 0.0003	1.043	0.57 ± 0.09	-0.50 ± 0.17

Fit results – dijet + shoulder resummation

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negative 3-jet power correction **only with shoulder resummation!**

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as expected
by theory

Summary

- Innovations of our analysis
 - Improved dijet/OPE and trijet/shoulder region
 - Inclusion of theory correlations in fits
- Dijet resummation crucial for robust results
- Shoulder resummation crucial for $\Theta_1 < 0$

final results

$$\chi^2/\text{dof} = 1.04$$

$$\alpha_s(m_Z) = 0.1145^{+0.0021}_{-0.0019}$$

$$\Omega_{10} = 0.57 \pm 0.09 \text{ GeV}$$

$$\Theta_1 = -0.50 \pm 0.17 \text{ GeV}$$

compatible with thrust and C-parameter

comparison
to data

