

Precise determination of the strong coupling from hadronic τ decays including $\tau \rightarrow \pi^- \pi^0 \nu_\tau$ from Belle

Diogo Boito

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With: A. Eiben, M. Golterman, K. Maltman, L. Mansur, and S. Peris

[2502.08147](#) DB, Eiben, Golterman, Maltman, Mansur, and Peris, *Phys. Rev. D* 111 (2025) 7, 074010

[2012.10440](#) DB, Golterman, Maltman, Peris, Rodrigues, and Schaaf, *Phys. Rev. D* 103 (2021) 3, 034028

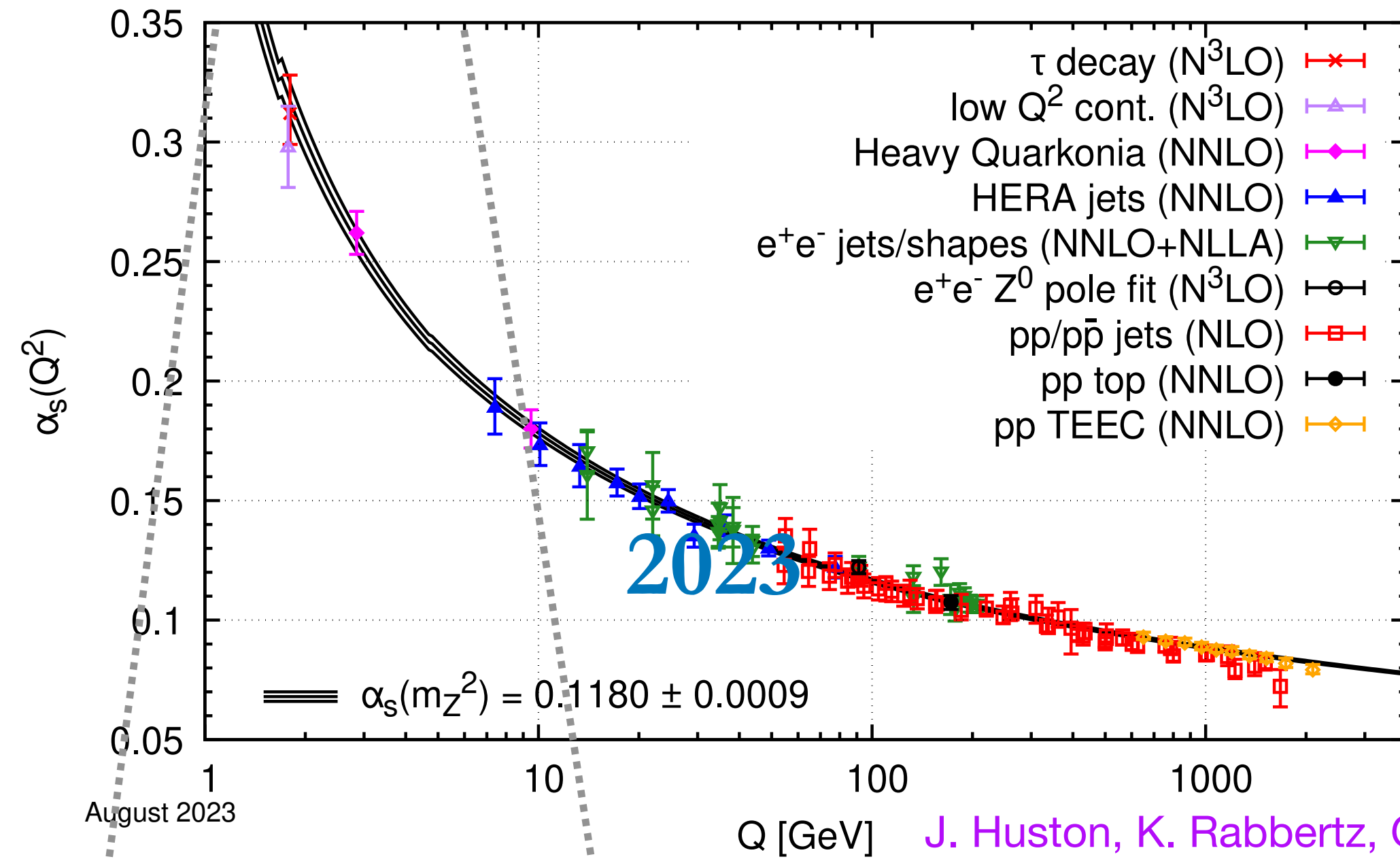


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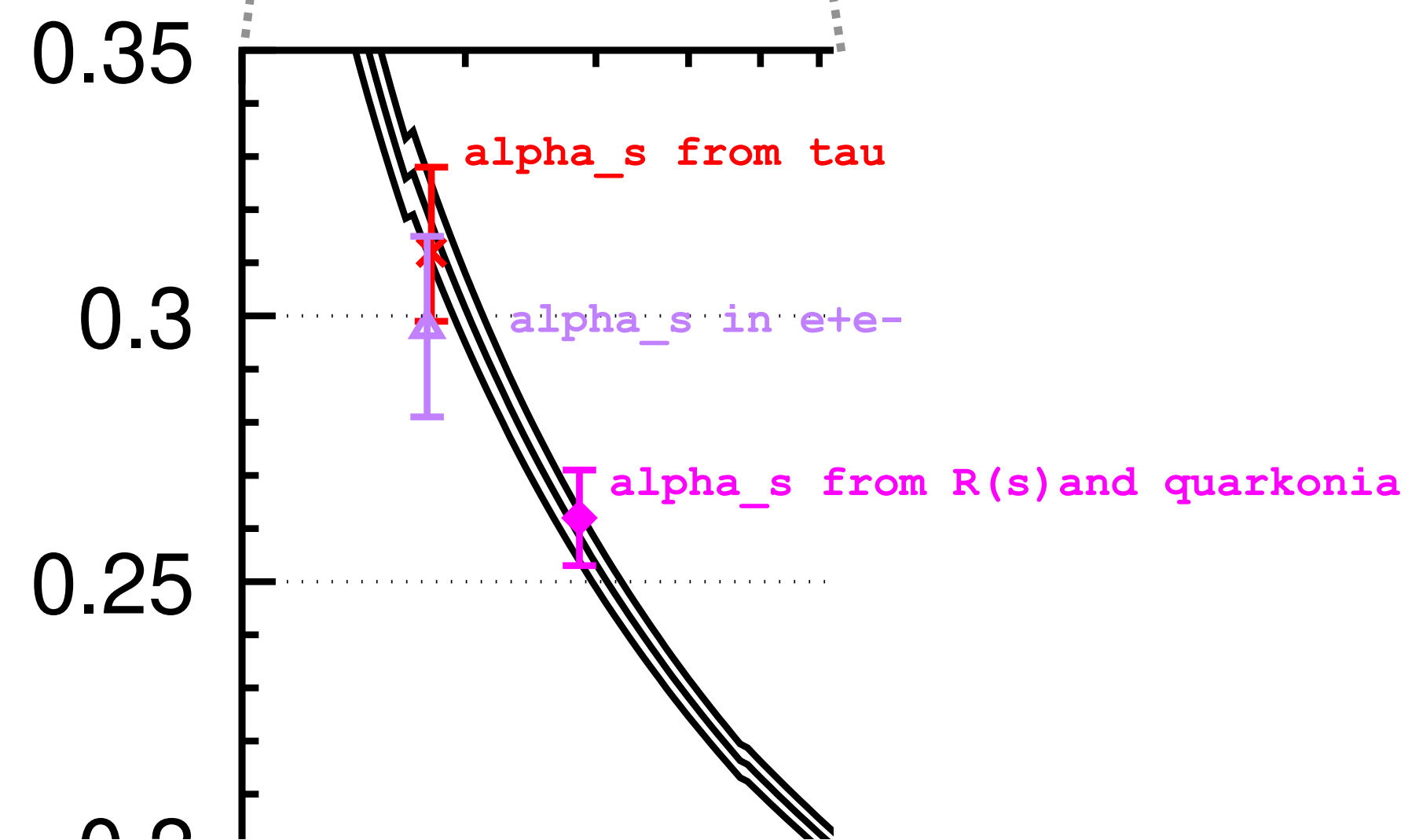
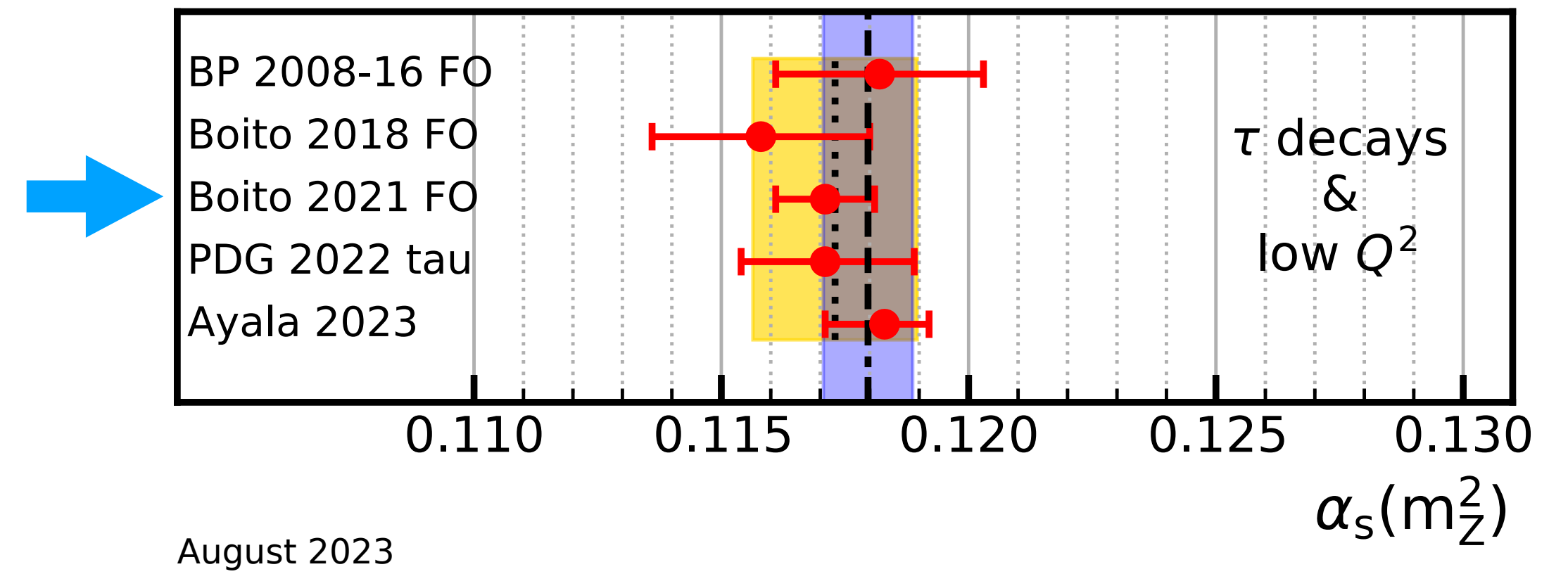


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Strong coupling from hadronic tau decays

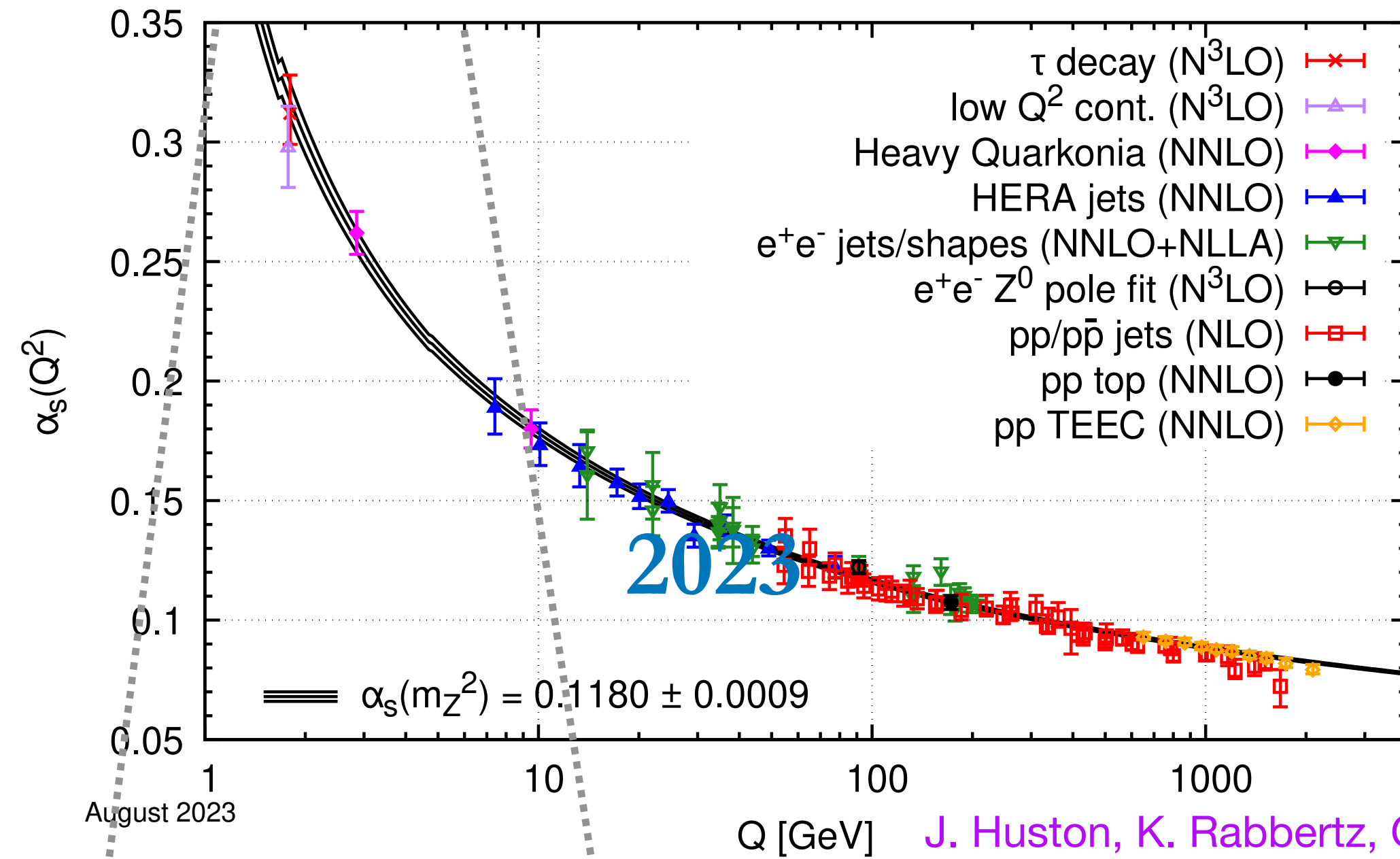


J. Huston, K. Rabbertz, G. Zanderighi,
Particle Data Group, (2023)

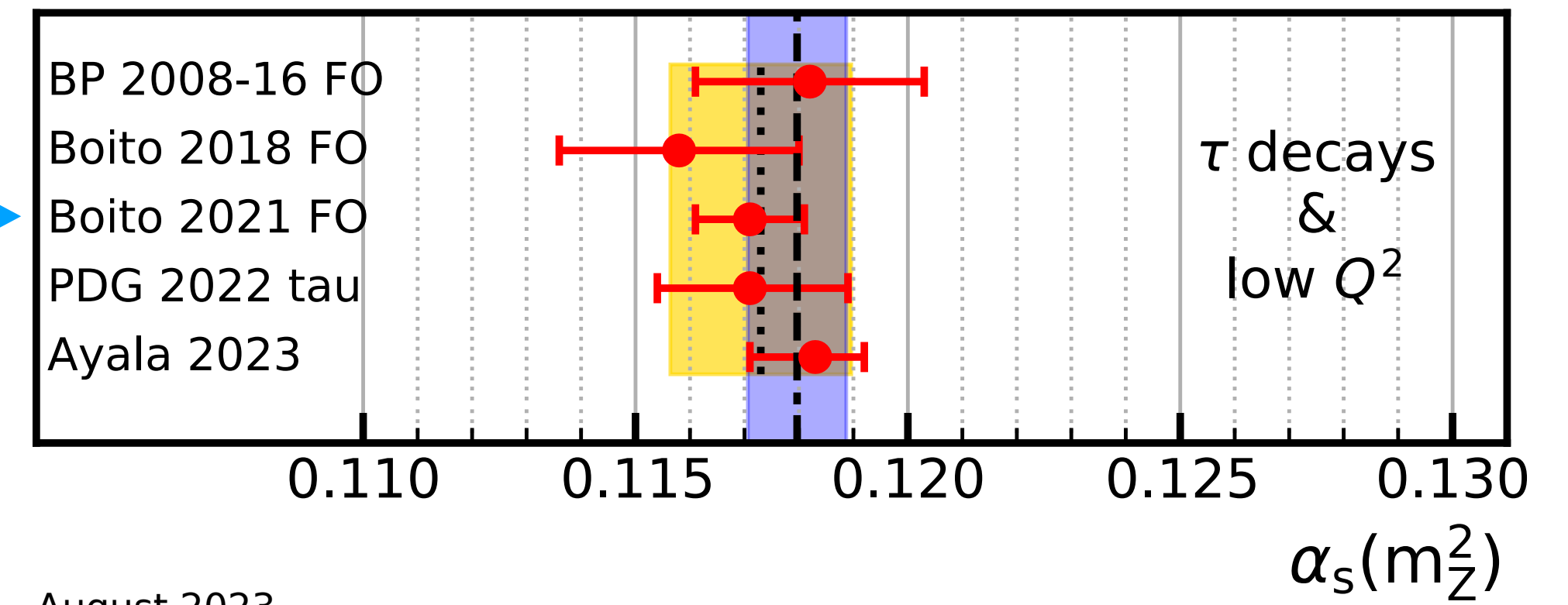
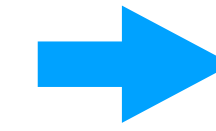


- Important test of asymptotic freedom
- Larger coupling at lower energies: more sensitivity to QCD corrections, less precision required from experimental data

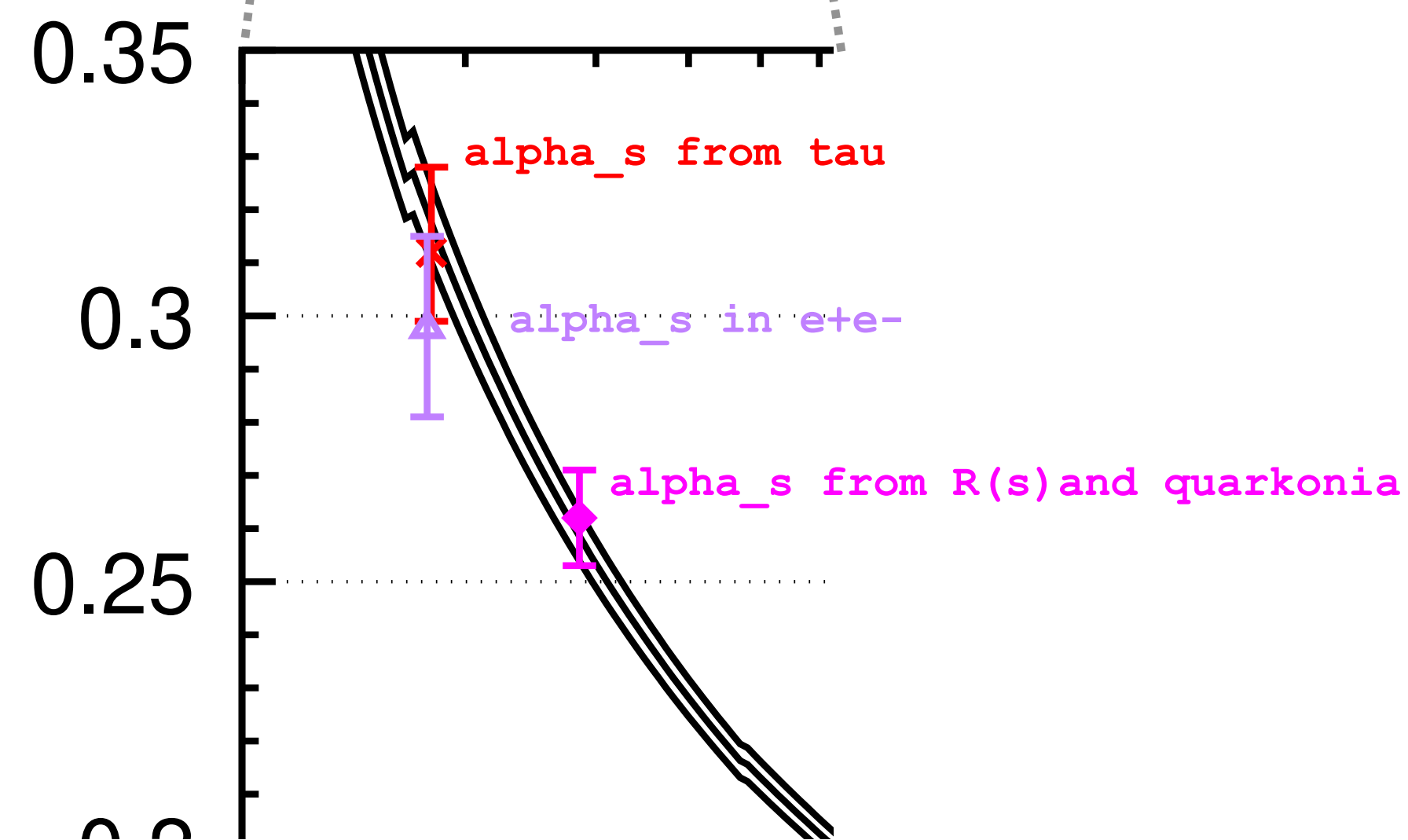
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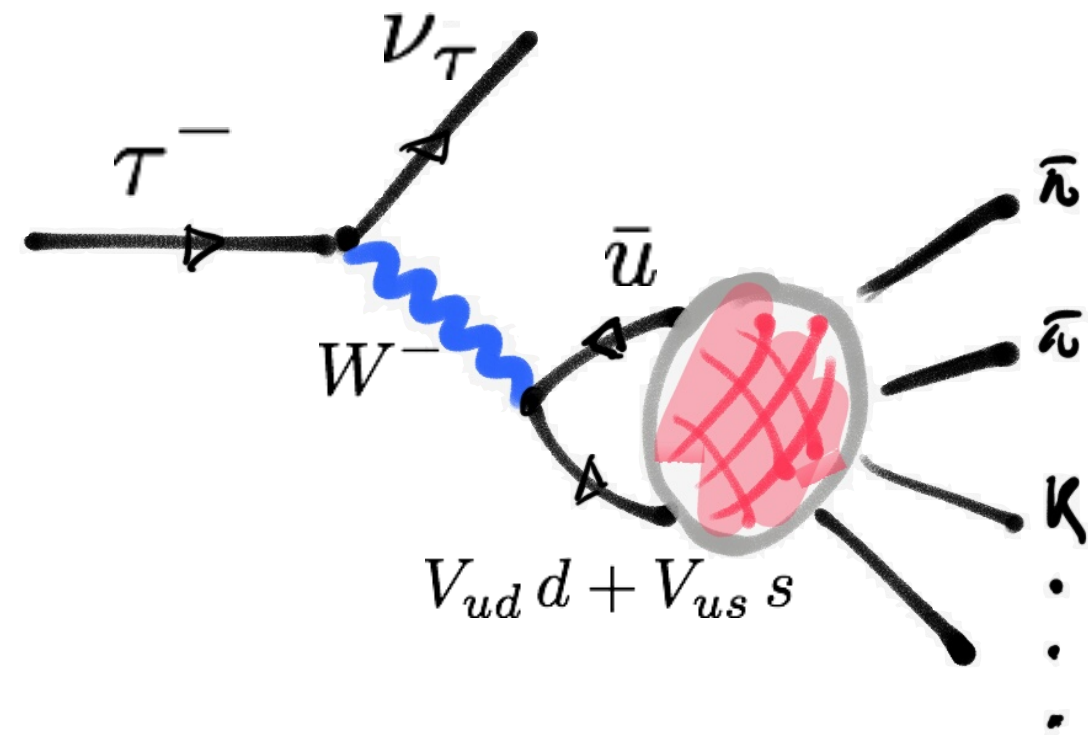


August 2023



- Important test of asymptotic freedom
- Larger coupling at lower energies: more sensitivity to QCD corrections, less precision required from experimental data
- Larger non-perturbative contributions (OPE, Duality Violations)
- Potential convergence issues in perturbation theory are enhanced (renormalons)

Strong coupling from hadronic tau decays



Massless (V & A) correlators

$$\Pi_{\mu\nu}(q) = i \int d^4x e^{iqx} \langle 0 | T \{ J_\mu(x) J_\nu(0)^\dagger \} | 0 \rangle$$

$$\mathcal{J}_\mu = \bar{u} \gamma_\mu (\gamma_5) d$$

Braaten, Narison, and Pich '92

Sum rules (using Cauchy's theorem)

$$\frac{1}{s_0} \int_0^{s_0} ds w(s) \underbrace{\frac{1}{\pi} \text{Im} \Pi(s)}_{\text{experiment}} = - \frac{1}{2\pi i s_0} \oint_{|z|=s_0} dz w(z) \underbrace{\Pi(z)}_{\text{theory}}$$

- Here we focus on the light-quark current. Tiny mass corrections can then be neglected, $m_u, m_d \sim 0$. No net strangeness in final states.
- We extract the coupling from the vector-isovector channel (to be discussed in the next slides)
- Experimental results should be **fully inclusive**

Theory overview

$$\frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi(z) \underset{\text{theory}}{=} \frac{1}{4\pi^2} S_{\text{EW}} \left(\delta_w^{\text{tree}} + \delta_w^{(0)} + \delta_{\text{OPE},w}^{D\geq 4} + \delta_{\text{DV}_s}^{(w)} + \delta_{\text{EW}}^{(w)} \right)$$

Perturbation theory

$$\delta_w^{(0)} = \sum_{n=1}^{\infty} \left(\frac{\alpha_s(\mu^2)}{\pi} \right)^n \sum_{k=1}^n k c_{n,k} \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} W(x) \ln^{k-1}(-s_0 x / \mu^2)$$

$\uparrow$
 $W(x) = 2 \int_x^1 dz w(z)$

$$R_\tau = \frac{\Gamma(\tau \rightarrow (\text{had}) + \nu_\tau)}{\Gamma(\tau \rightarrow e \bar{\nu}_e \nu_\tau)} = 3.634(8) \quad \text{Banerjee et al. (HFLAV), 2411.18639}$$

$$\approx N_c (1 + \delta_{w_\tau}^{(0)}) = 3(1 + 0.2) \sim 3.6$$

Perturbative series for the decay width

Gorishnii, Kataev, Larin '91
Surguladze&Samuel '91

Baikov, Chetyrkin, Kühn '08

$$\delta_{w_\tau}^{(0)} = \alpha_s^1 + \alpha_s^2 + \alpha_s^3 + \alpha_s^4 = 0.1001 + 0.0521 + 0.0264 + 0.0127 = 0.1914$$

(fixed order perturbation theory for R_τ , $\alpha_s(m_\tau^2) = 0.3144$)

pQCD correction is ~20%

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Estimate of fifth order coefficient included

$$c_{5,1} = 283 \pm 140 \quad \begin{array}{l} \text{Beneke and Jamin '08} \\ \text{DB, Masjuan, Oliani '18} \\ \text{Caprini '19} \end{array}$$

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$$\delta_{\text{OPE}}^{D \geq 4} = c_w^{(4)} \frac{C_4}{Q^4} + c_w^{(6)} \frac{C_6}{Q^6} + c_w^{(8)} \frac{C_8}{Q^8} + \dots$$

C_D coefficients are parameters of the fit

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Duality Violations

$$\rho_{\text{DV}}(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s)$$

Ansatz based on widely accepted assumptions about QCD: asymptotic Regge behavior and large- N_c . Main expected corrections: logarithmic and powers of $1/s$.

DB, Caprini, Golterman, Maltman, Peris, PRD '18

Theory: Fixed Order vs Contour Improved

- For a long time the question of renormalization scale fixing led to an ambiguity
- Fixed Order Perturbation Theory (FOPT, fixed scale) leads to smaller values for the coupling than Contour Improved Perturbation Theory (CIPT, running scale on the integration contour)

	α_s^1	α_s^2	α_s^3	α_s^4	
$\delta_{\text{FO}}^{(0)}$	$= 0.1001$	$+ 0.0521$	$+ 0.0264$	$+ 0.0127$	$= 0.1914$
$\delta_{\text{CI}}^{(0)}$	$= 0.1353$	$+ 0.0255$	$+ 0.0100$	$+ 0.0070$	$= 0.1779$

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- **Issue now resolved:** CIPT is intrinsically different and incompatible with the usual form of the OPE condensates [Hoang & Regner '20, '21 \(see also Golterman, Maltman, Peris, '23; Gracia, Hoang, and Mateu '23\)](#)

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- The problem with CIPT can be fixed by an adequate subtraction of the leading IR renormalon [Benitez-Rathgeb, DB, Hoang, Jamin '22](#)

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We will not quote results from CIPT in this talk

Analysis strategy

Any analytical weight function can be used in the FESRs (normally polynomials are used)

$$\frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi(z) = \frac{1}{4\pi^2} S_{\text{EW}} \left(\delta_{\text{tree}}^w + \delta_w^{(0)} + \delta_{\text{OPE},w}^{D \geq 4} + \delta_{\text{DV}s}^{(w)} + \delta_{\text{EW}}^{(w)} \right)$$

Desired properties from the choice of weights:

1. Good perturbative behavior
2. Small condensate contributions
3. Suppression of DVs

Suppression of DVs comes with the price of additional (unknown) higher dim. contributions from the OPE: **see-saw effect**

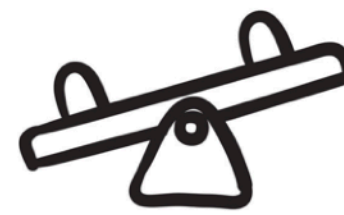
DV strategy (this work)

DB, Eiben, Golterman, Maltman, Mansur, and Peris, 2502.08147.

DB, M. Golterman, K. Maltman, S. Peris, M. V. Rodrigues and W. Schaaf, 2012.10440

- Accept some DVs (oscillations evident in data)
- Strongly suppress contamination on the OPE side
- DVs have to be parametrized with a model:

$$\rho_{\text{DV}}(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s)$$



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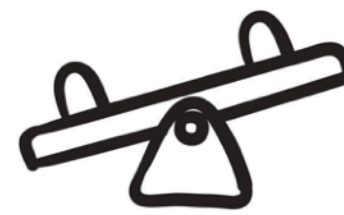
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Truncated OPE strategy

A Pich, A. Rodriguez-Sanchez 1605.06830
Davier, Höcker, Malaescu, Yuan, Zhang 1312.1501

- Suppress DVs by using $w(z)$ with zeros at $z=0$ (pinching) but need to ignore the higher order contributions on the OPE side (too many parameters)
- Contrary to previous understanding, α_s obtained from **perturbative QCD, disregarding all non-perturbative contributions**

see DB, M. Golterman, K. Maltman, S. Peris '24

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Choice of weights

$$w_0(y) = 1$$

$$w_2(y) = 1 - y^2$$

$$w_3(y) = (1 - y)^2(1 + 2y)$$

$$w_4(y) = (1 - y^2)^2$$

Tiny condensate contributions, sensitive to DVs

$D = 6$

$D = 6$ and 8 : **Tau kinematical Moment** (R_τ)

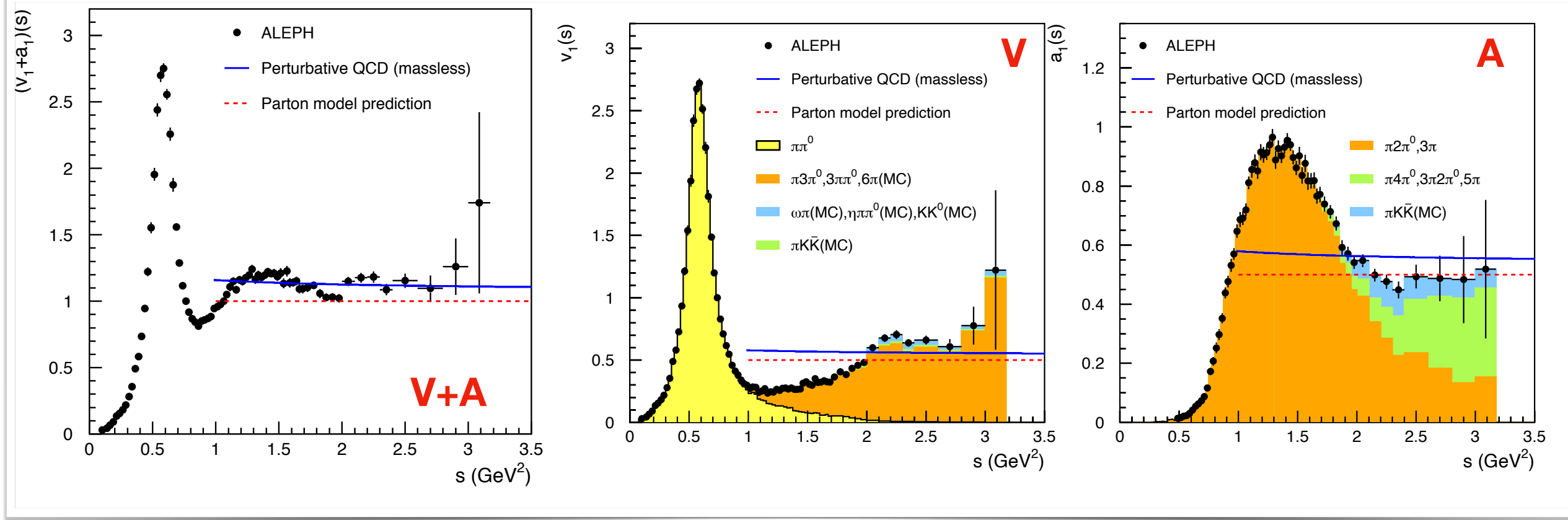
$D = 6$ and 10

DB, Cata, Golterman, Jamin, Maltman '11,
Beneke, DB, Jamin '12,
DB, M. Golterman, K. Maltman, S. Peris '16
DB F Olani '20
DB, Golterman, Maltman, and Peris '24

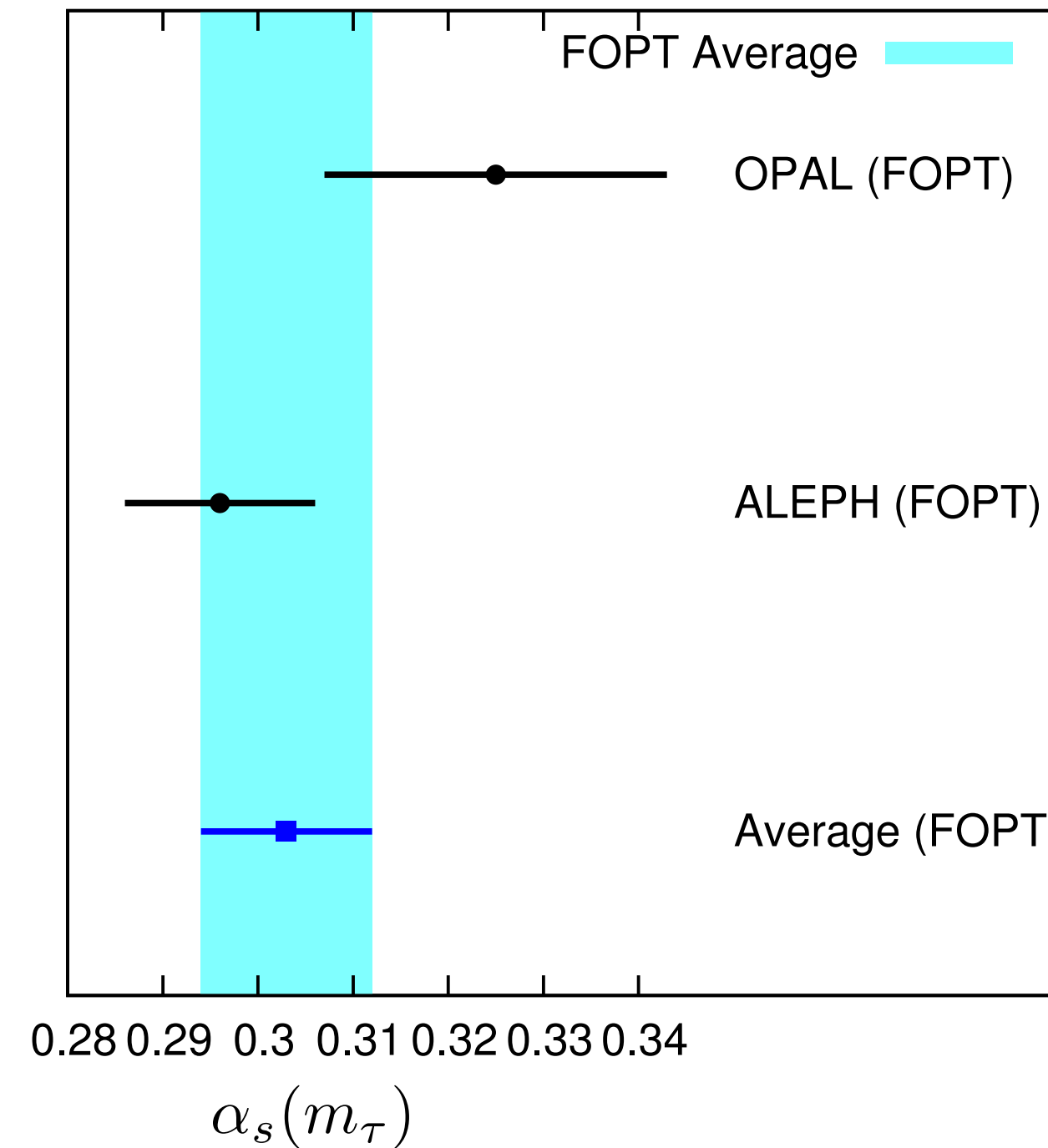
Main results are based on the simplest weight function $w_0(y) = 1$

Experimental data

Davier et al [ALEPH] '14



The two data sets led to compatible values for $\alpha_s(m_\tau)$ for the DV strategy

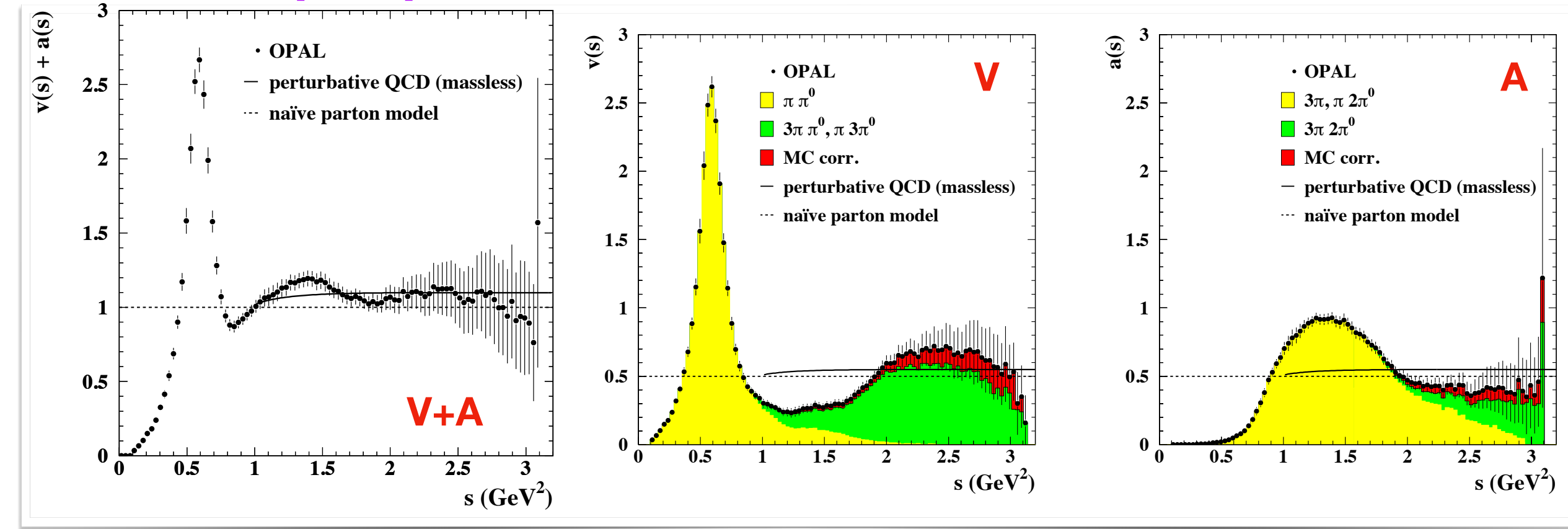


DB, Golterman, Jamin, Mahdavi, Maltman, Osborne, and Peris, '12

DB, Golterman, Maltman, Osborne, and Peris, '15

Ackerstaff et al [OPAL] '98

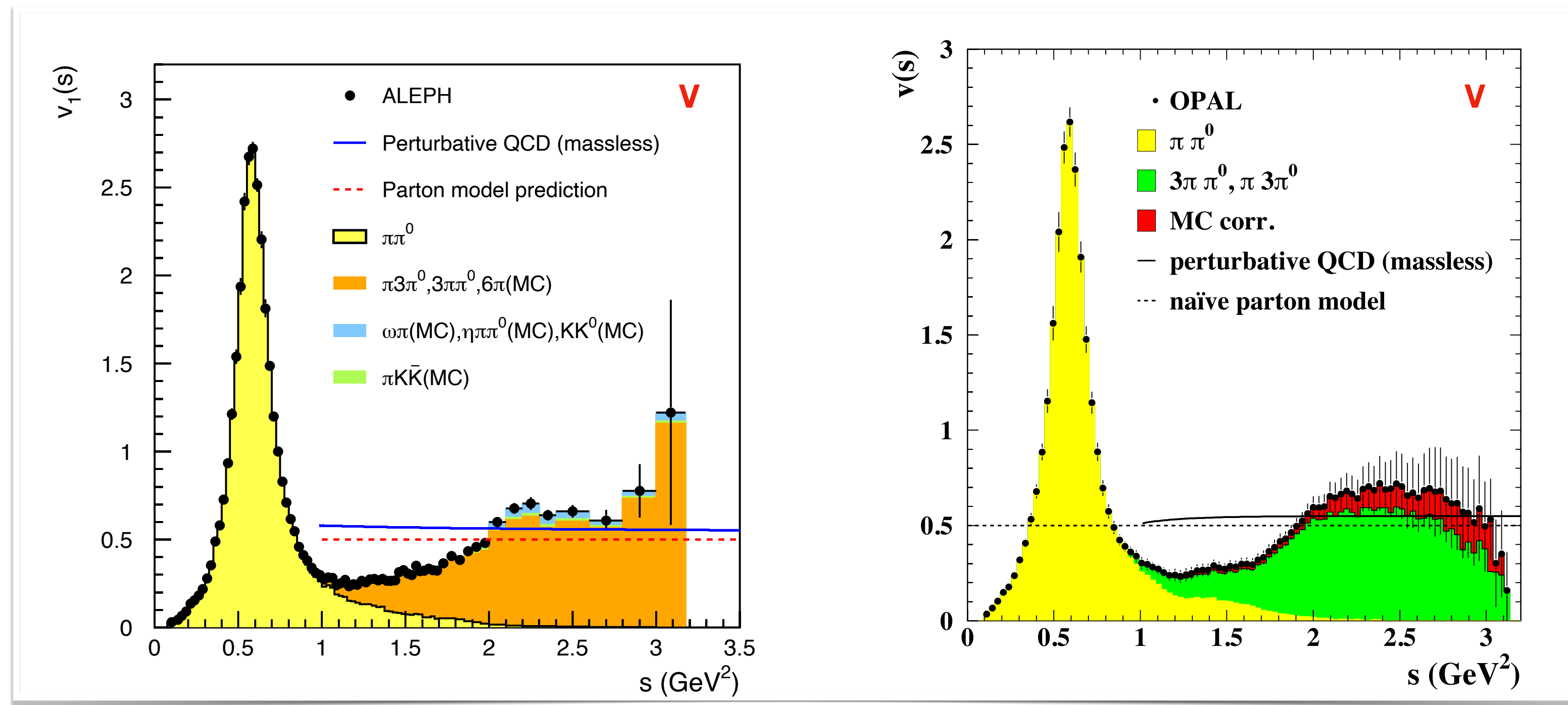
updated for recent values of branching fractions Boito et al '11, '21



- Can we combine the data sets (as done, e.g., for $g-2$ of the muon)?
- Are they (locally) compatible?
- Can we improve on the LEP-based data sets with recent experimental results?

Anatomy of the data sets

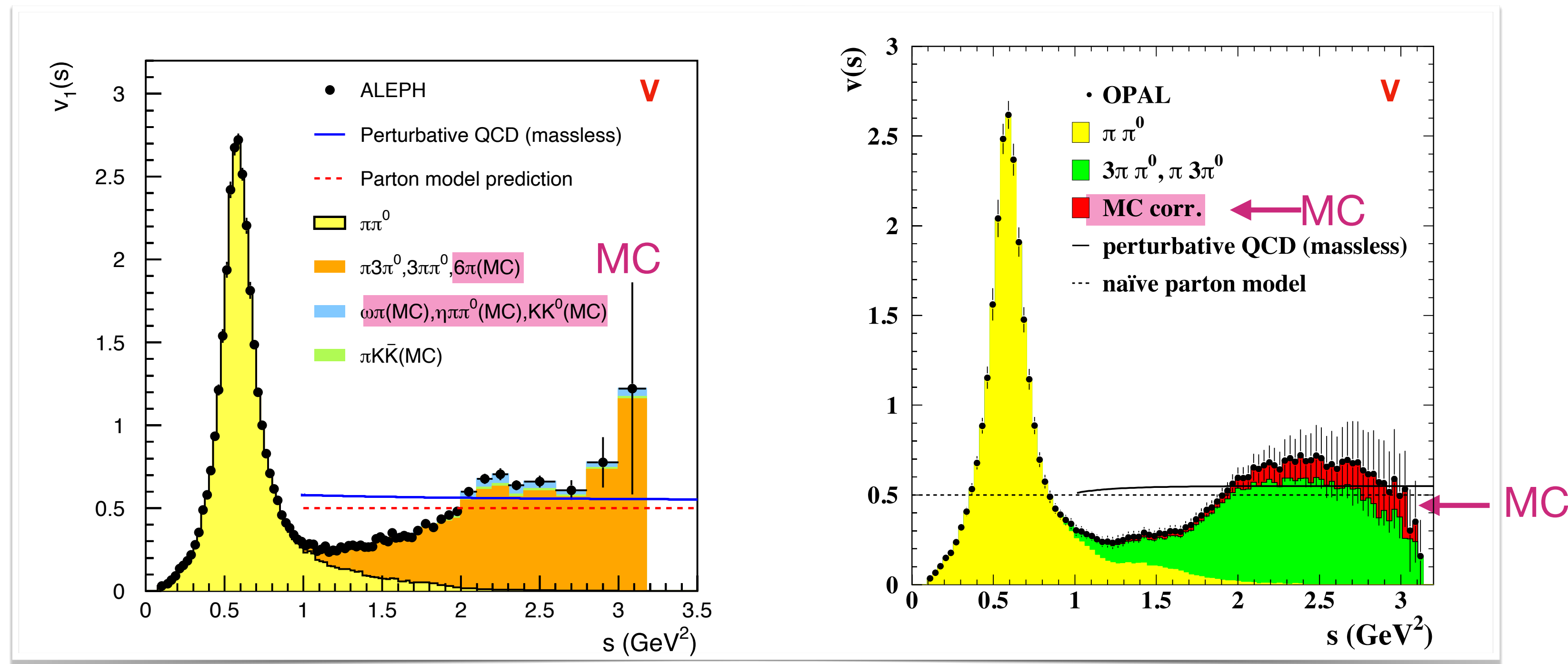
- V channel dominated by $\tau \rightarrow 2\pi + \nu_\tau$ and $\tau \rightarrow 4\pi + \nu_\tau$
- Residual channels are subdominant (but very important for α_s !)
- ALEPH and OPAL used **Monte Carlo input** for several of the subdominant (residual) modes (**not real data**)



Recently measured channels in e^+e^- can be used to improve the vector channel using conserved vector current (CVC) for residual modes in regions where isospin breaking is small

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Improved vector-isovector spectral function

- We introduced in 2020 a method to combine existing data and improve on previous spectral functions

DB, M. Golterman, K. Maltman, S. Peris, M. V. Rodrigues and W. Schaaf, 2012.10440

- Combined $2\pi + 4\pi$ data from ALEPH and OPAL
- Results based **solely** on **experimental data**: 7 residual channels obtained from recent e^+e^- data (**no need for MC inputs**)
- All results updated for recent branching ratio measurements



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DB, Eiben, Golterman, Maltman, Mansur, and Peris, 2502.08147

- **New data combination algorithm** to deal with strong correlations in 4π spectra from ALEPH and OPAL

- Based on the algorithm of Keshavarzi, Nomura, Teubner '18, '19 (used in the context of $g-2$ of the muon)
- Statistical treatment of non- χ^2 fits from Bruno & Sommer, '23

- Inclusion of high statistics Belle spectrum for $\tau \rightarrow \pi^- \pi^0 \nu_\tau$ **for the first time** in inclusive hadronic tau decays

- Results including the CLEO spectrum for $\tau \rightarrow \pi^- \pi^0 \nu_\tau$ also given

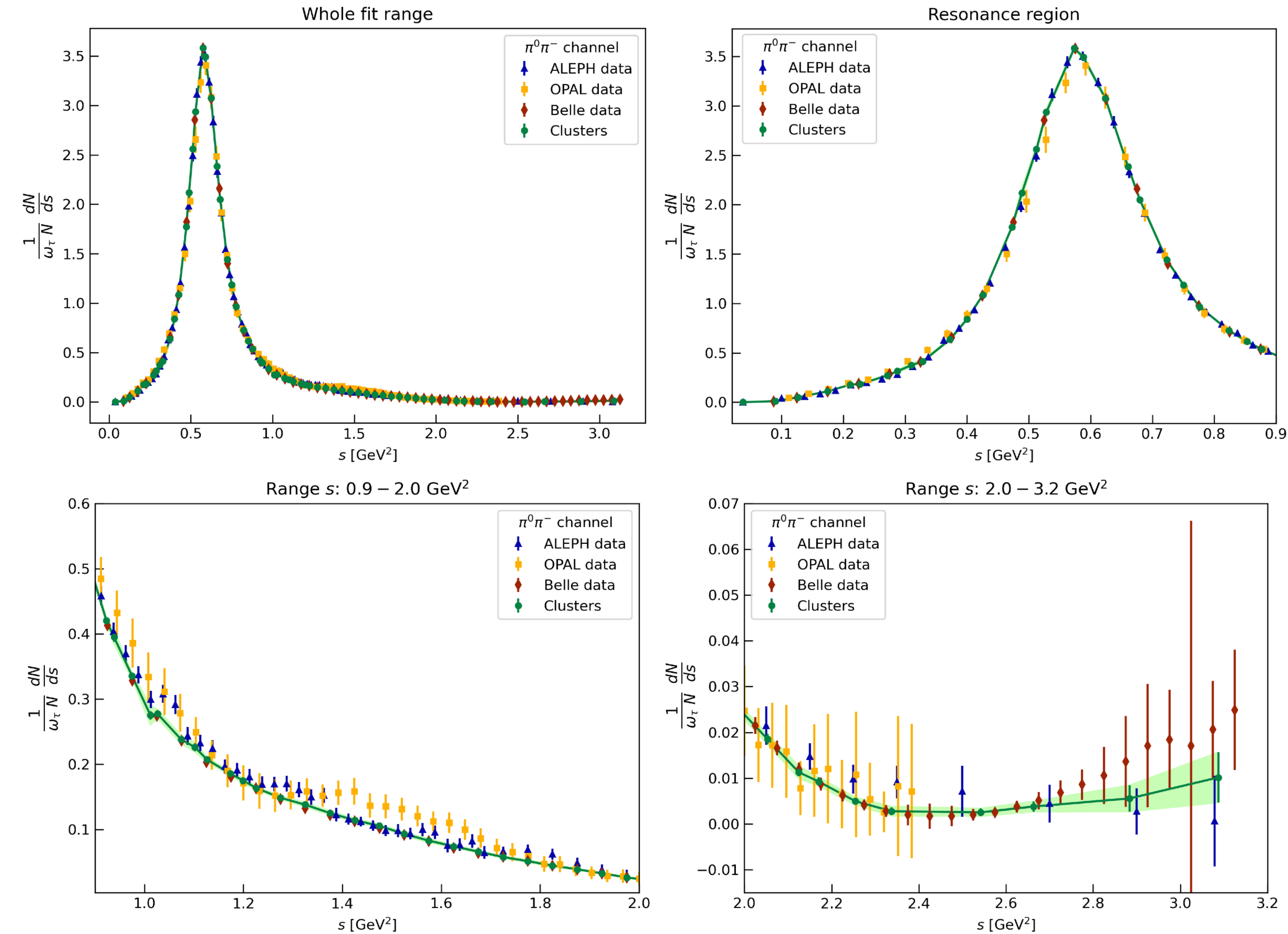
(but not enough information about correlations exist to include CLEO in the final results)

- New HFLAV branching ratios used Banerjee et al. (HFLAV) '24

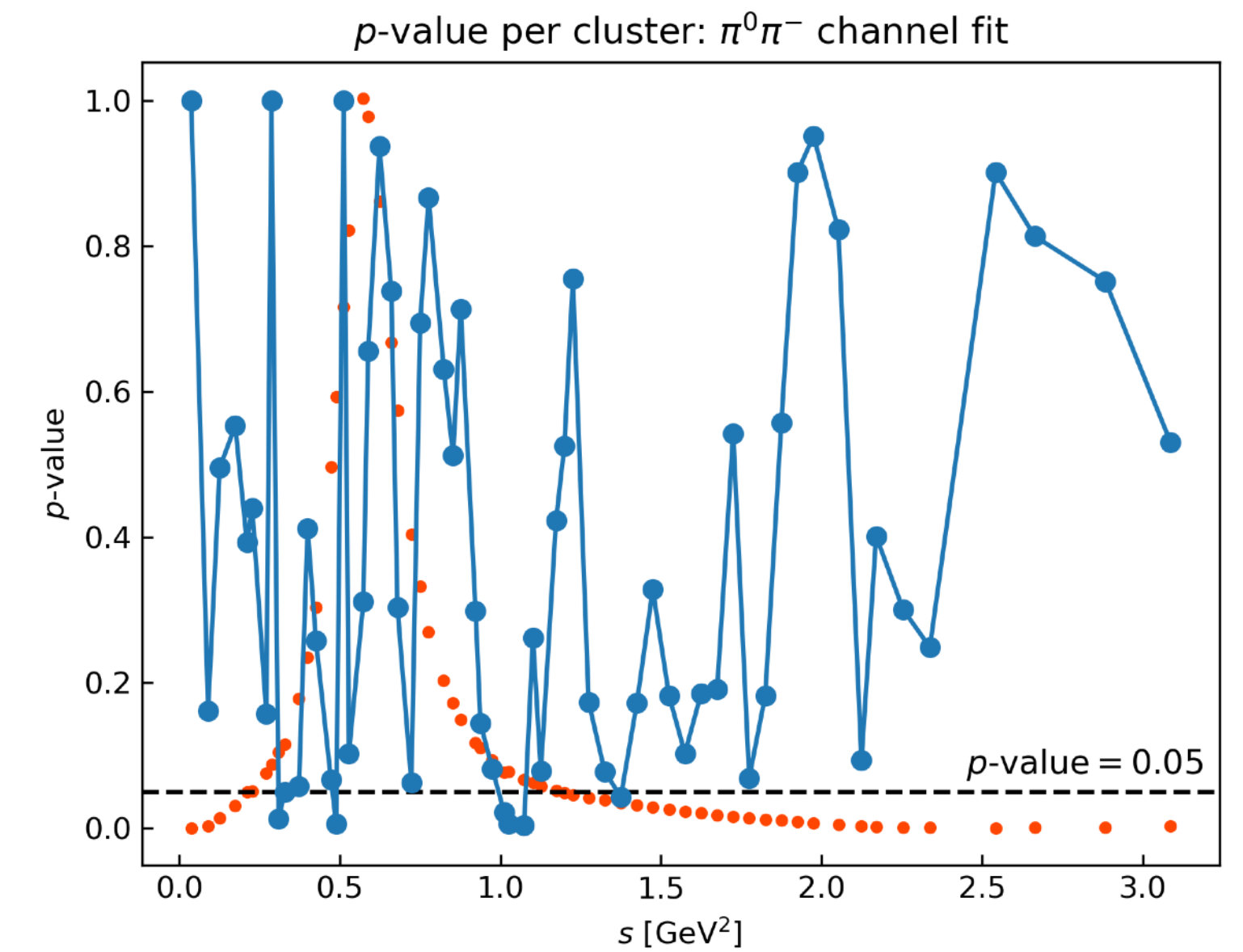
- Algorithm to avoid d'Agostini bias in multiplication by global factors (branching fractions) Ball et al. (NNPDF) '10

Improved vector-isovector spectral function: 2π channel

Combination of 2π spectra (combined results in green)

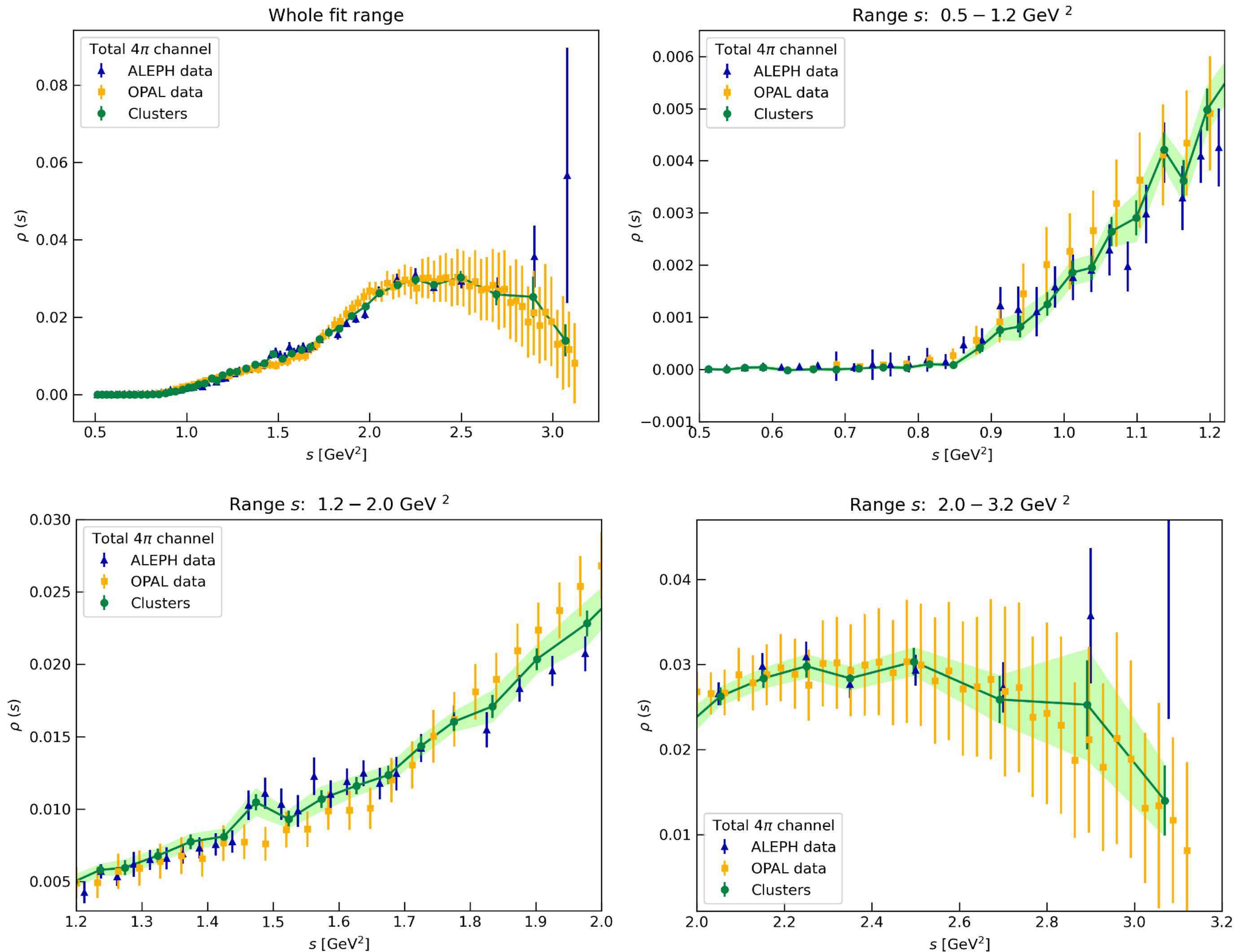


Good χ^2 both locally and globally
Effects of error inflation are minimal



Improved vector-isovector spectral function: 4π channels

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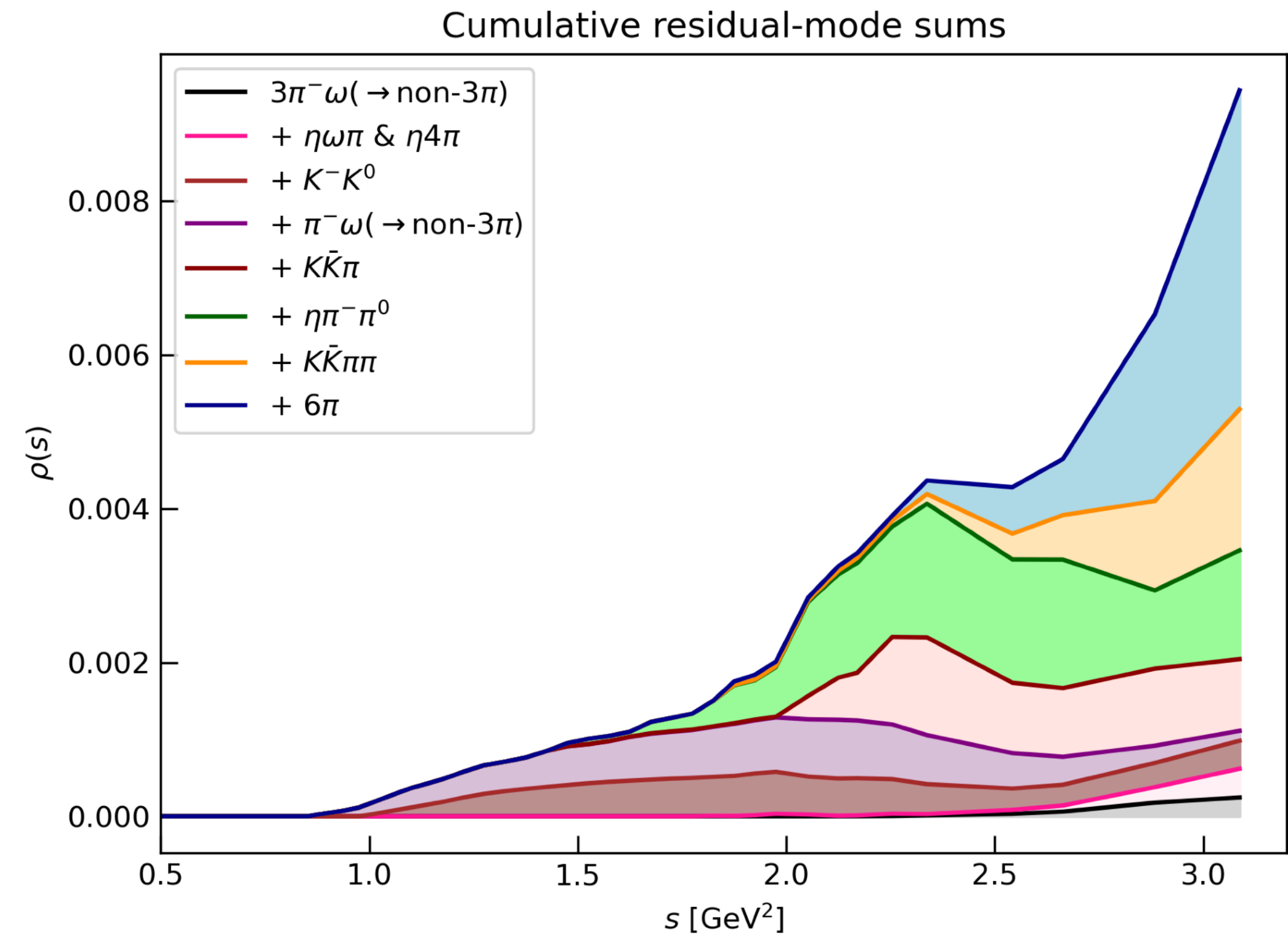
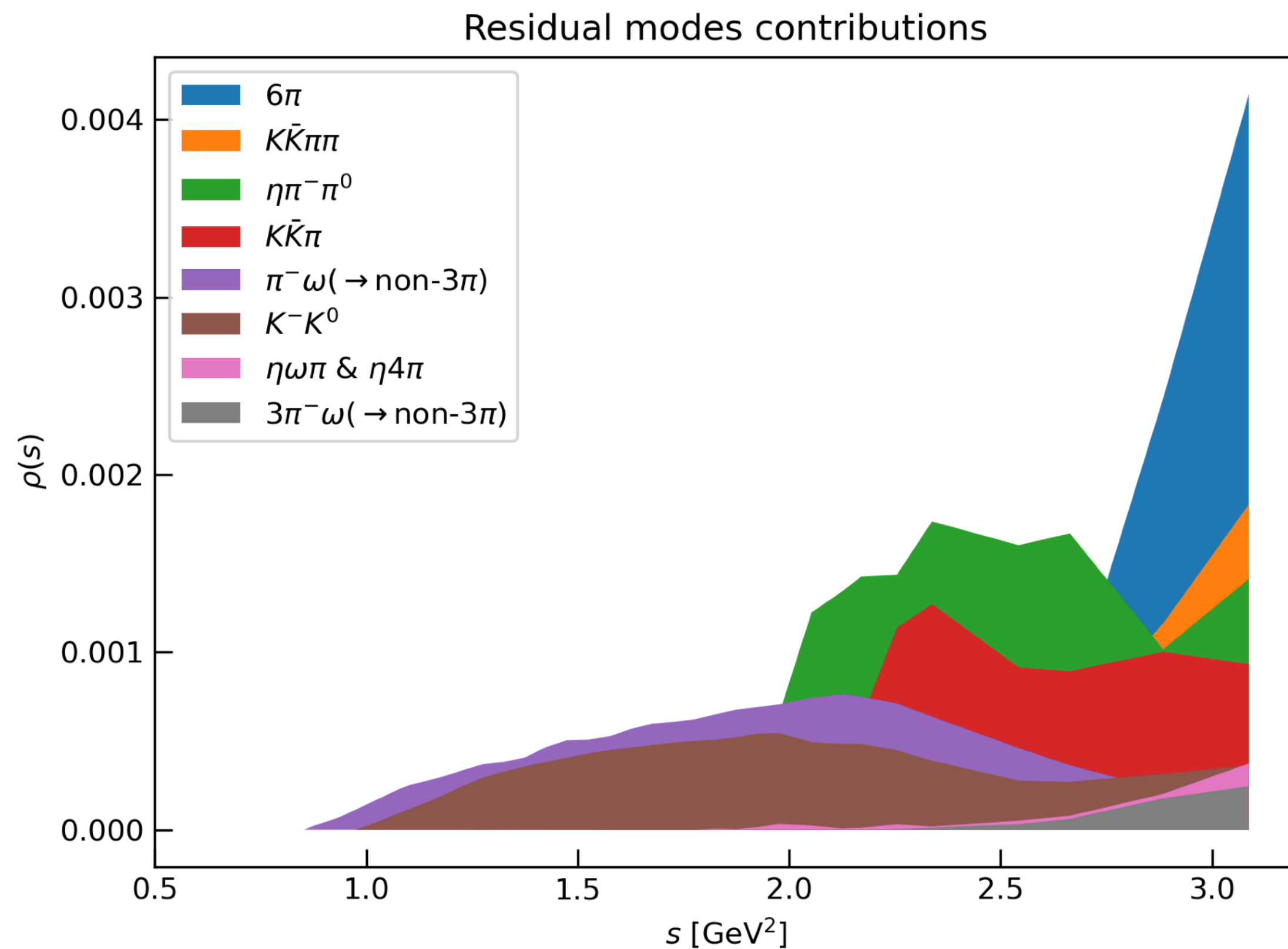


- Correlation matrices for 4π exclusive channels are ill-behaved (negative eigenvalues)
- Adding up the 4π channels softens the problem, but strong correlations in OPAL results lead to poor fits
- The main issue are the strong, but poorly known, correlations in the $\pi^-3\pi^0$ OPAL channel
- Perform fits where the strong correlations in OPAL $\pi^-3\pi^0$ channel are not included in the minimization, but are still used in the error propagation
- General p -value analysis for fits of this type well understood. Implemented with a Monte Carlo (see Bruno & Sommer, '23)

Residual channels

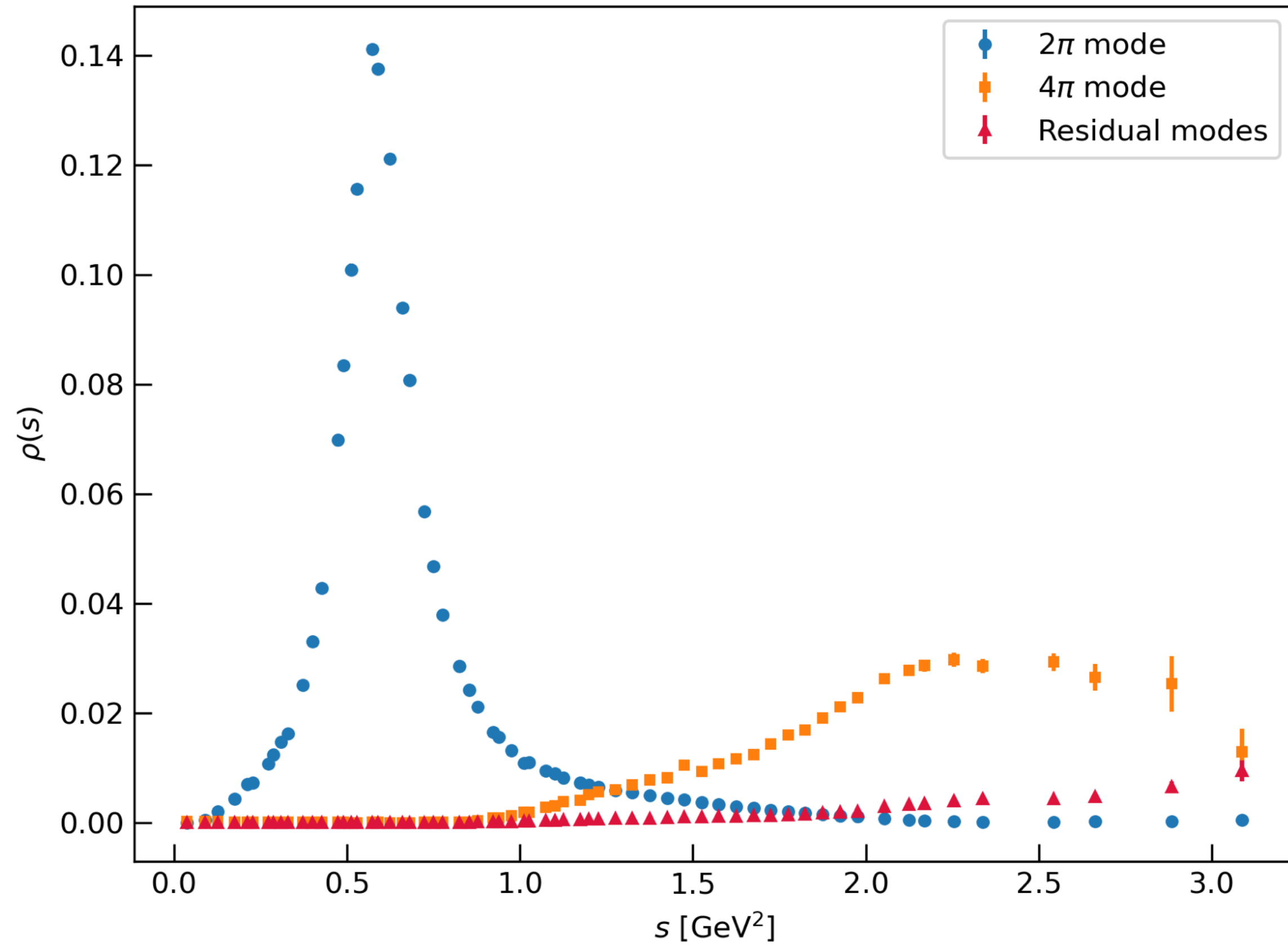
- 7 residual channels extracted from e^+e^- data + BaBar data for $\tau \rightarrow KK_S\nu_\tau$
- Dramatic improvement in errors for higher multiplicity modes (near end point)
- Isospin breaking corrections on already small residual modes (1–2%) would not affect final results

No Monte Carlo input

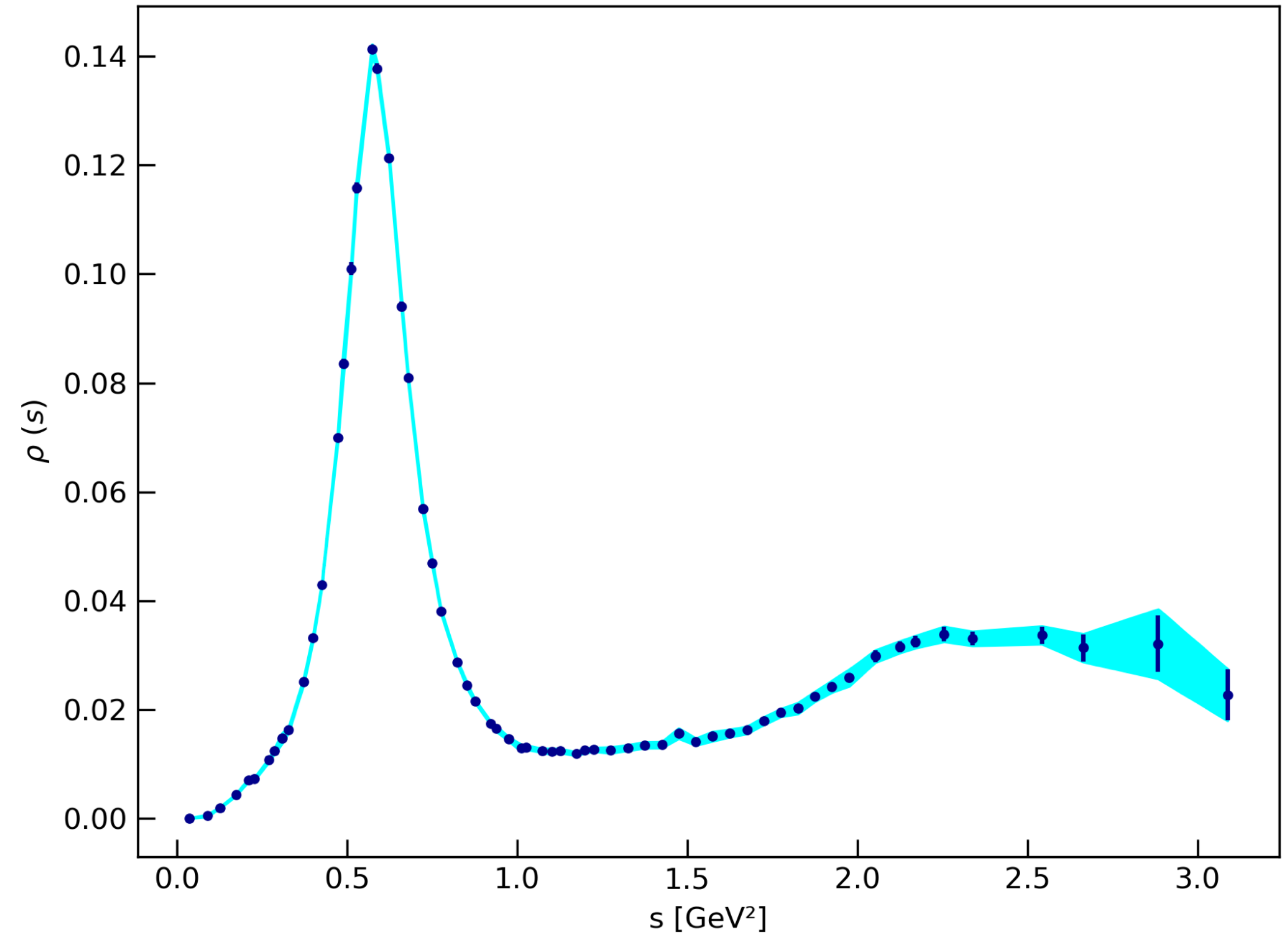


Final spectral function

Contributions



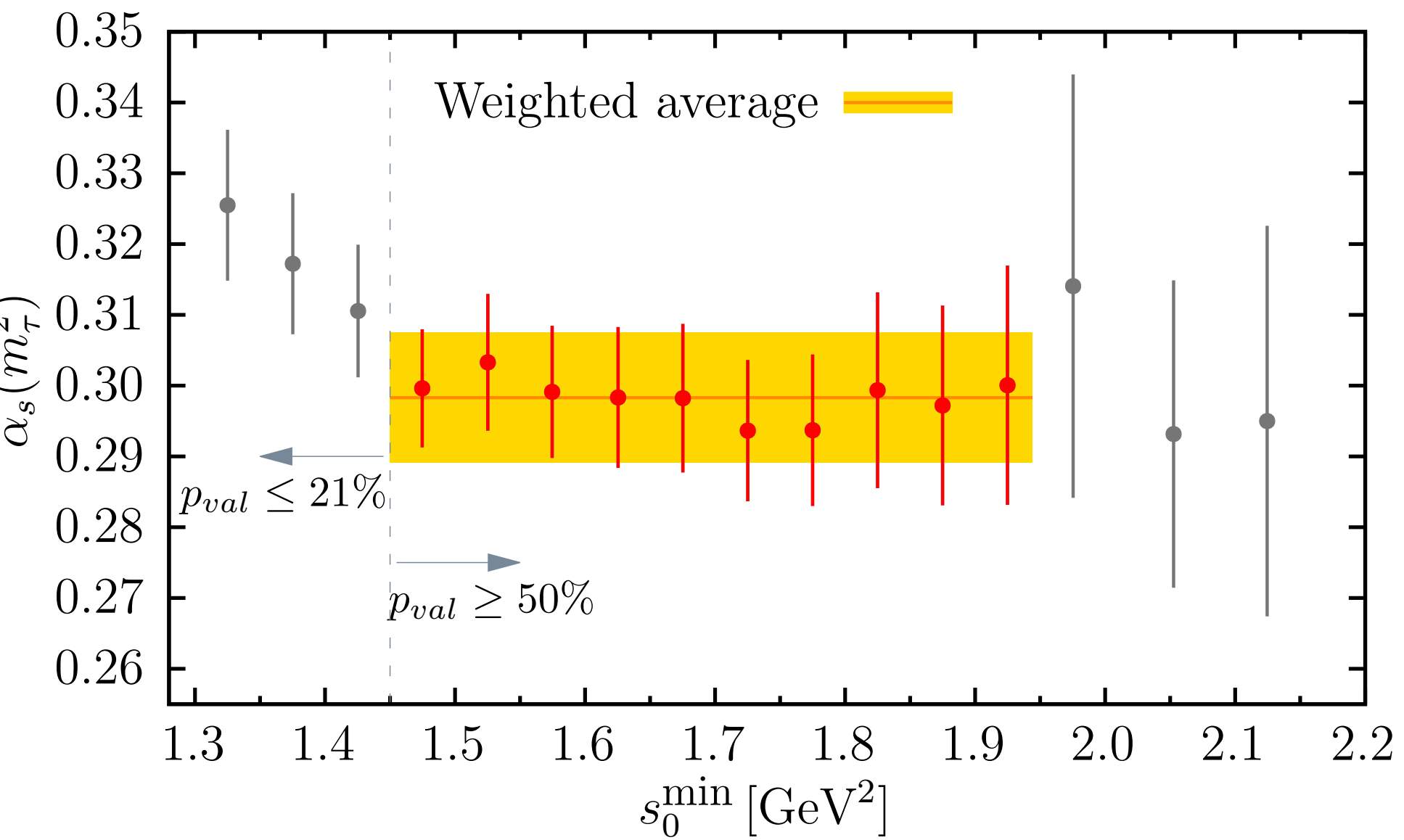
Total vector-isovector spectral function



DB, Eiben, Golterman, Maltman, Mansur, and Peris, 2502.08147

Strong coupling from the new spectral function

- Several fits, single moments or in combination
- Many fit windows: $[s_{\min}, m_{\tau}^2]$
- Consistency between different fits (α_s , condensates, DV params.)



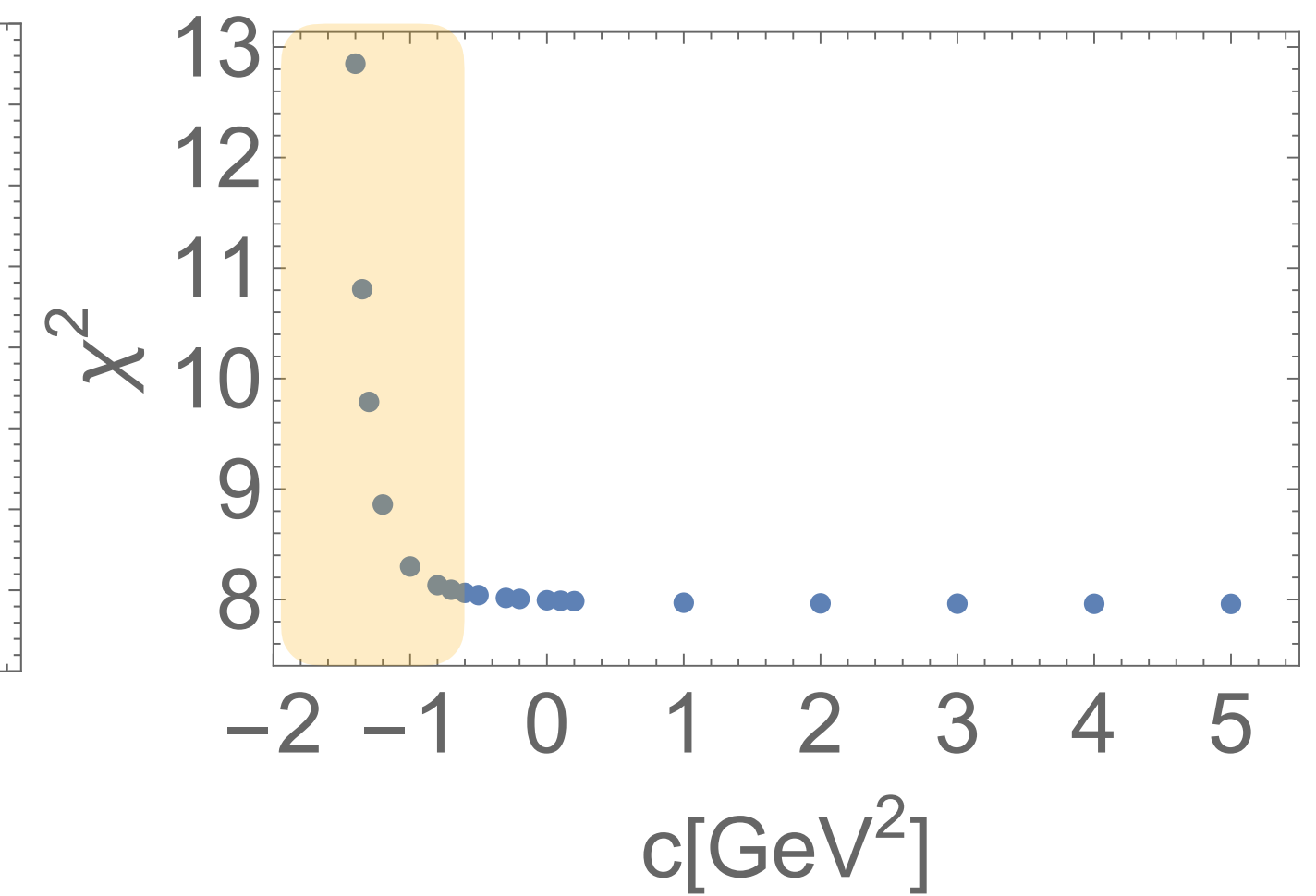
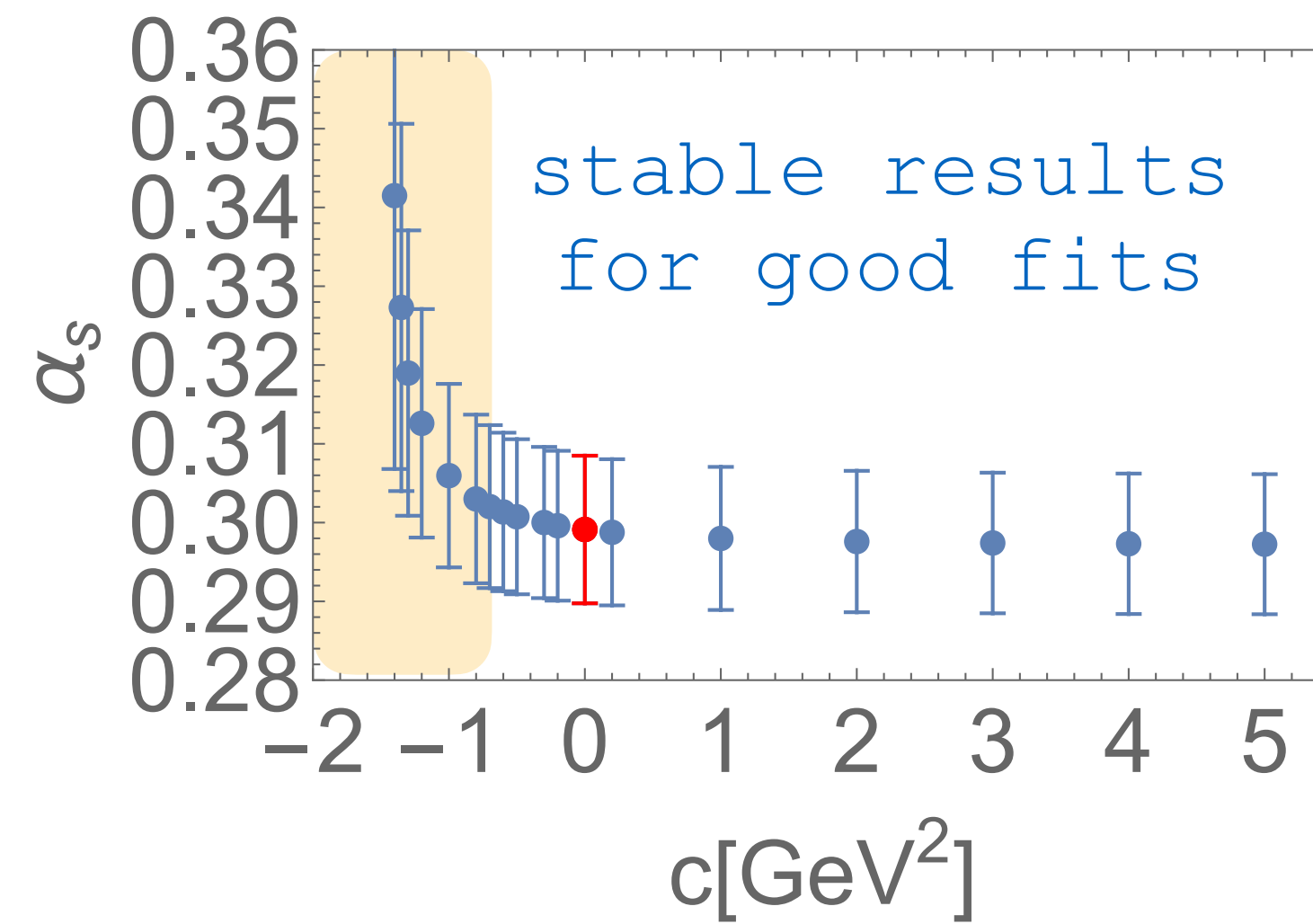
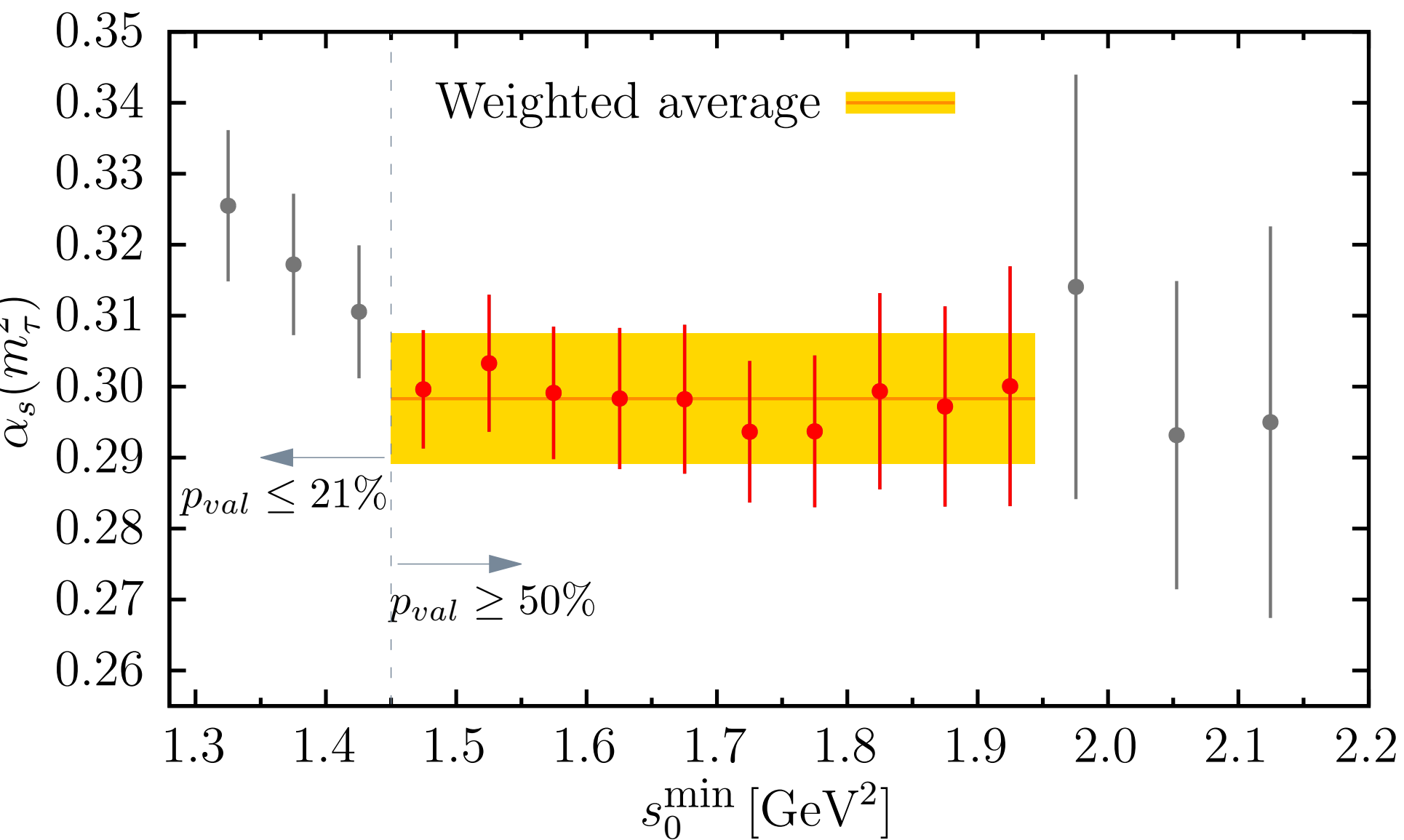
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- **Test of possible correction to the DV ansatz**

DB, Caprini, Golterman, Maltman, Peris, PRD '18

Pich & Rodriguez-Sanchez '22

$$\rho^{\text{DV}}(s) = \left(1 + \frac{c}{s}\right) e^{-\delta - \gamma s} \sin(\alpha + \beta s)$$



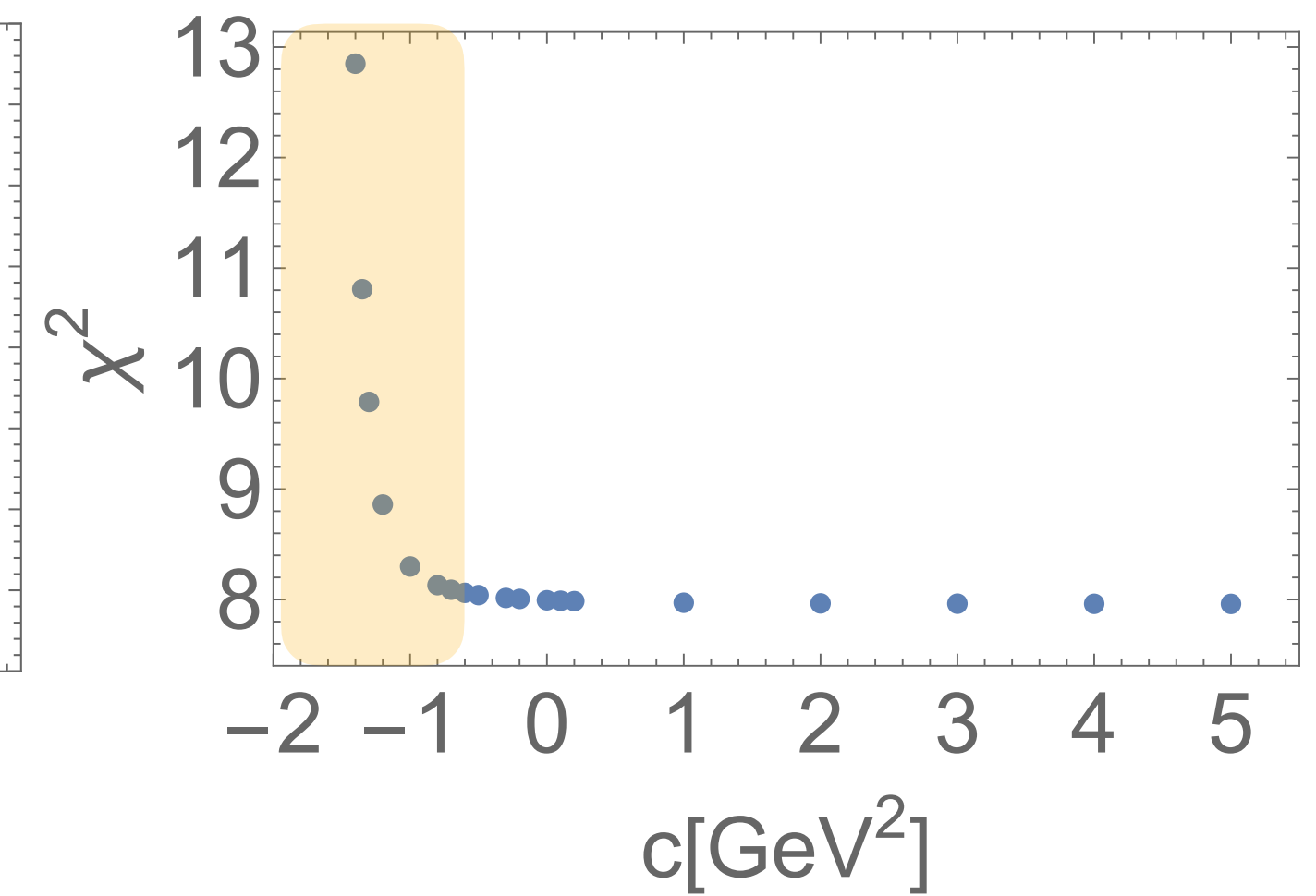
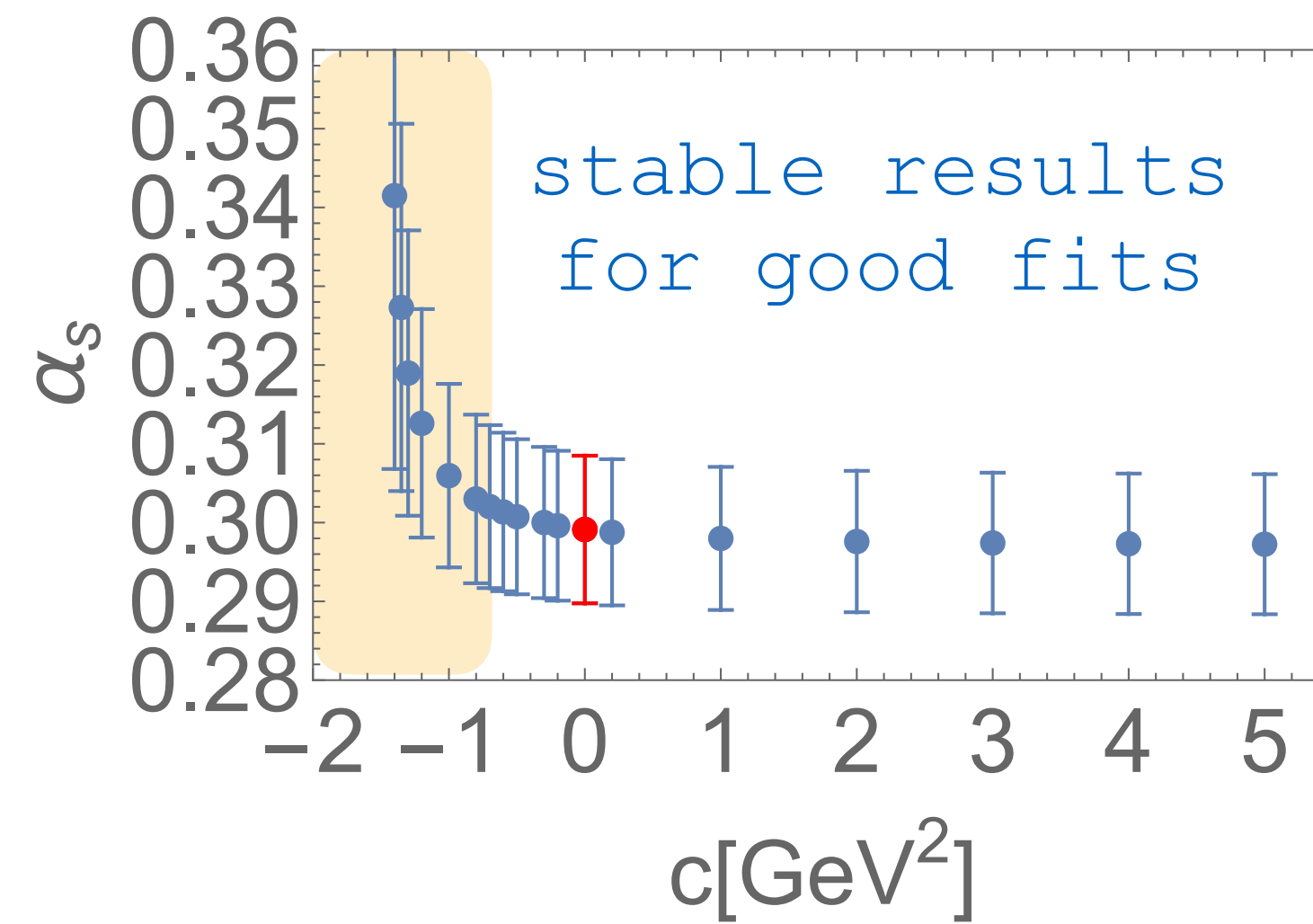
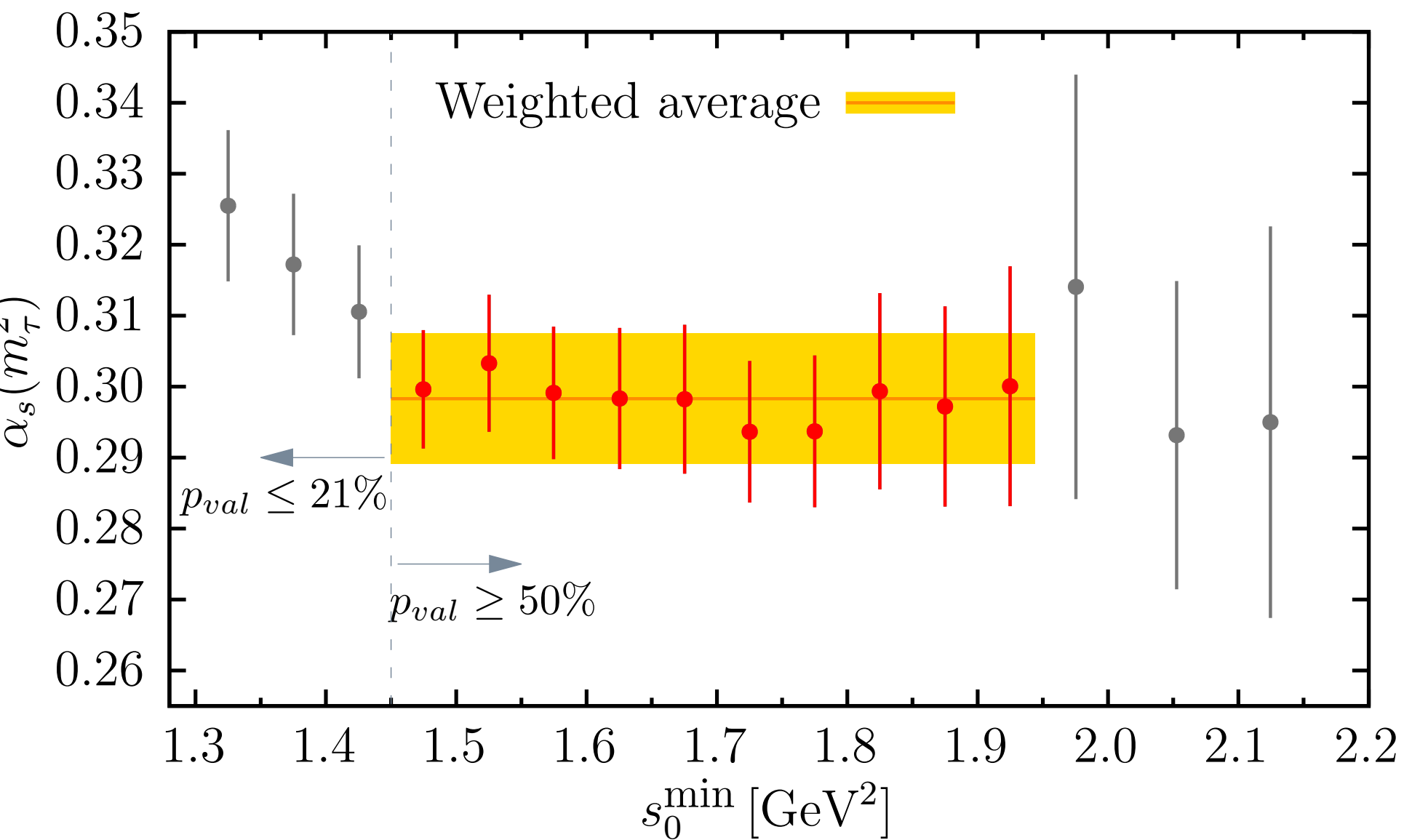
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DB, Eiben, Golterman, Maltman, Mansur, and Peris, 2502.08147

$$\begin{aligned} \alpha_s(m_\tau^2) &= 0.2983 \pm 0.0092_{\text{stat}} \pm 0.0026_{\text{fit}} \pm 0.0022_{\text{pert}} \pm 0.0025_{\text{DVs}} \\ &= 0.2983 \pm 0.0101 \quad (n_f = 3) \end{aligned}$$

$$\alpha_s(m_Z^2) = 0.1159 \pm 0.0014 \quad (n_f = 5)$$

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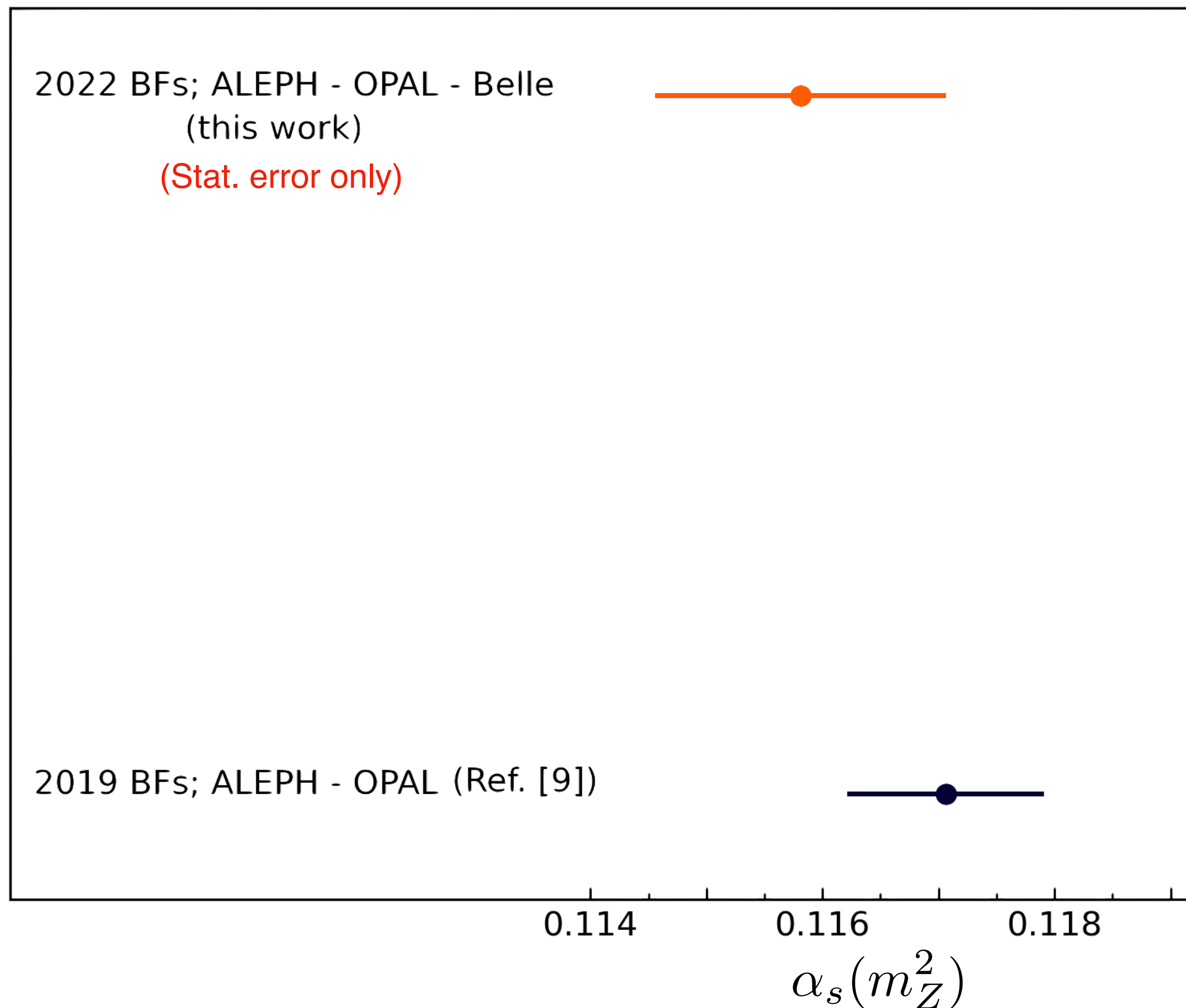
Our previous result from 2020 was significantly larger

$$\begin{aligned}
 \alpha_s(m_\tau^2) &= 0.3077 \pm 0.0065_{\text{stat}} \pm 0.0038_{\text{pert}} \\
 &= 0.3077 \pm 0.0075 .
 \end{aligned}$$

DB, Eiben, Golterman, Maltman,
 Mansur, and Peris, '25

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DB, M. Golterman, K. Maltman, S. Peris,
 M. V. Rodrigues and W. Schaaf, '20



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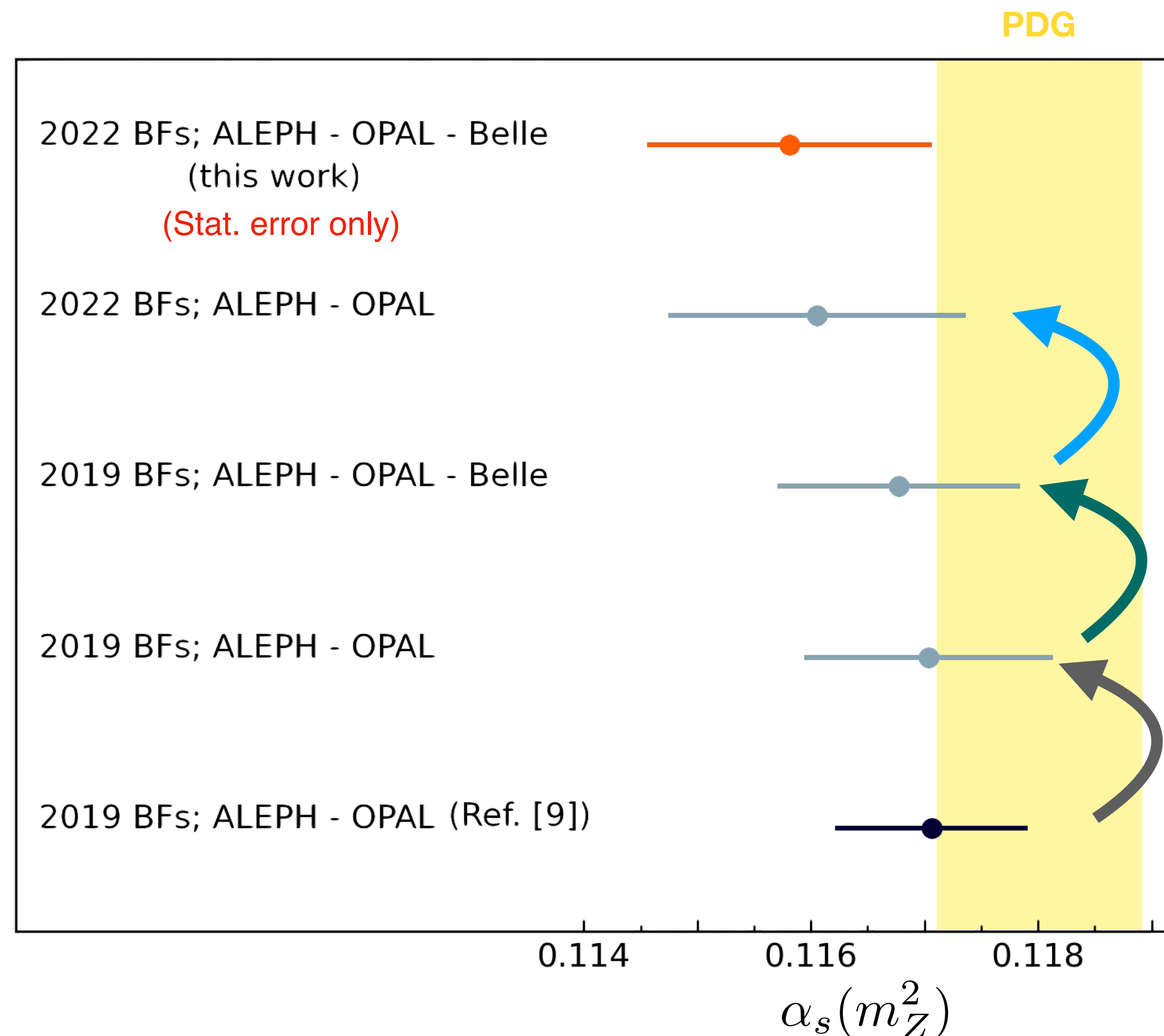
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DB, M. Golterman, K. Maltman, S. Peris,
 M. V. Rodrigues and W. Schaaf, '20



Final result with Belle, new BFs and new algorithm

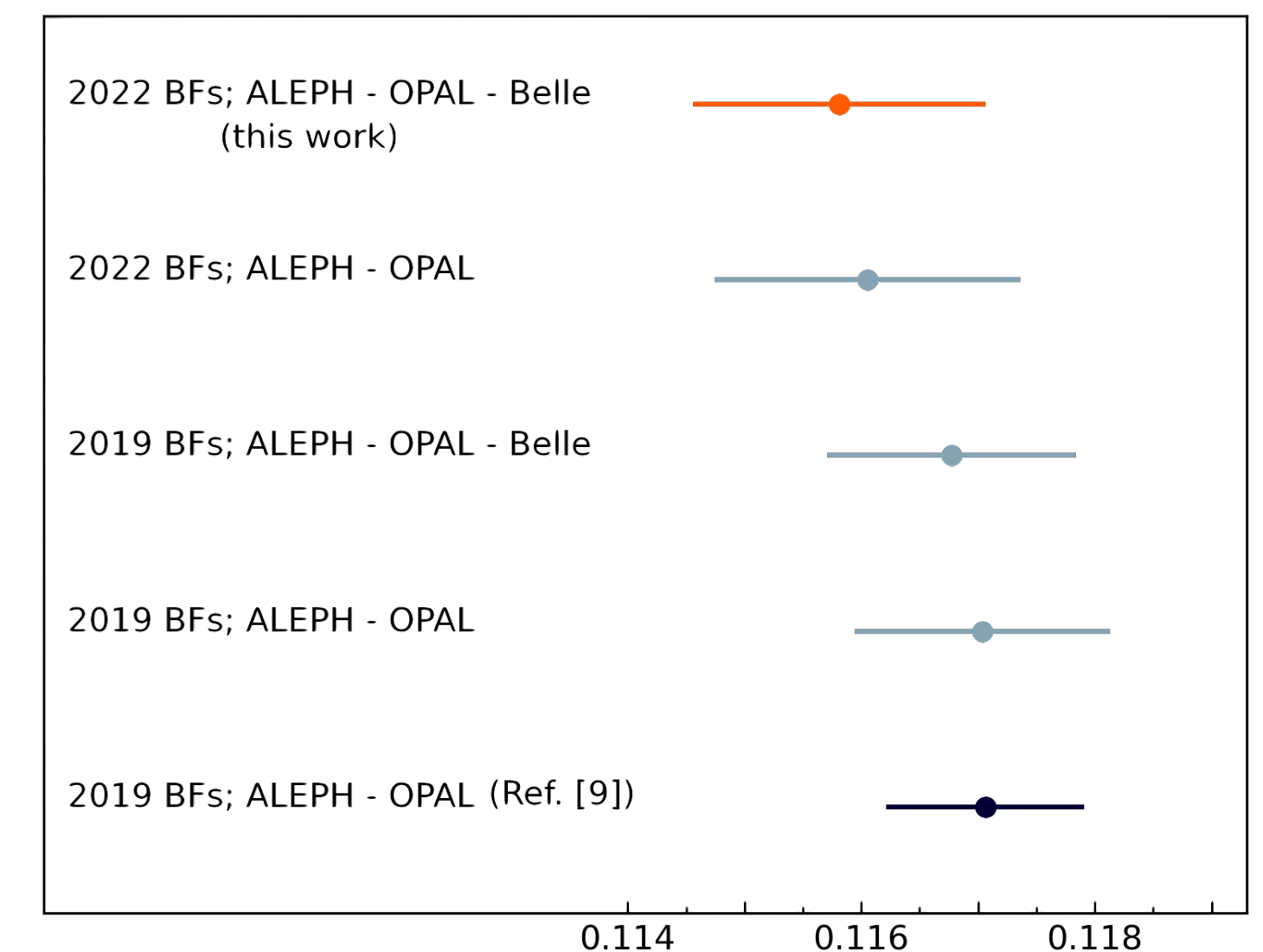
Largest shift due to new values of the 4π branching fractions!

Inclusion of the Belle spectrum leads to a small shift towards smaller values.

Changes in our data combination algorithm lead to slightly larger errors but no shift

Conclusions

- Vector channel in hadronic tau decays is special: e^+e^- data + CVC allows for improvement near kin. end point
- New vector-isovector spectral function purely based on exp. data, **no MC input needed**
- Inclusion, for the first time, of high-statistics Belle 2π spectrum in an inclusive hadronic tau decay analysis
- Our analysis can immediately incorporate any new spectrum for 2π or 4π tau decay channels (Belle II?) 
- Improvements of this type not possible for the axial channel
- Final vector spectral function is competitive
- Final value for the strong coupling lower than before mainly due to changes in exp. BF's



Conclusions

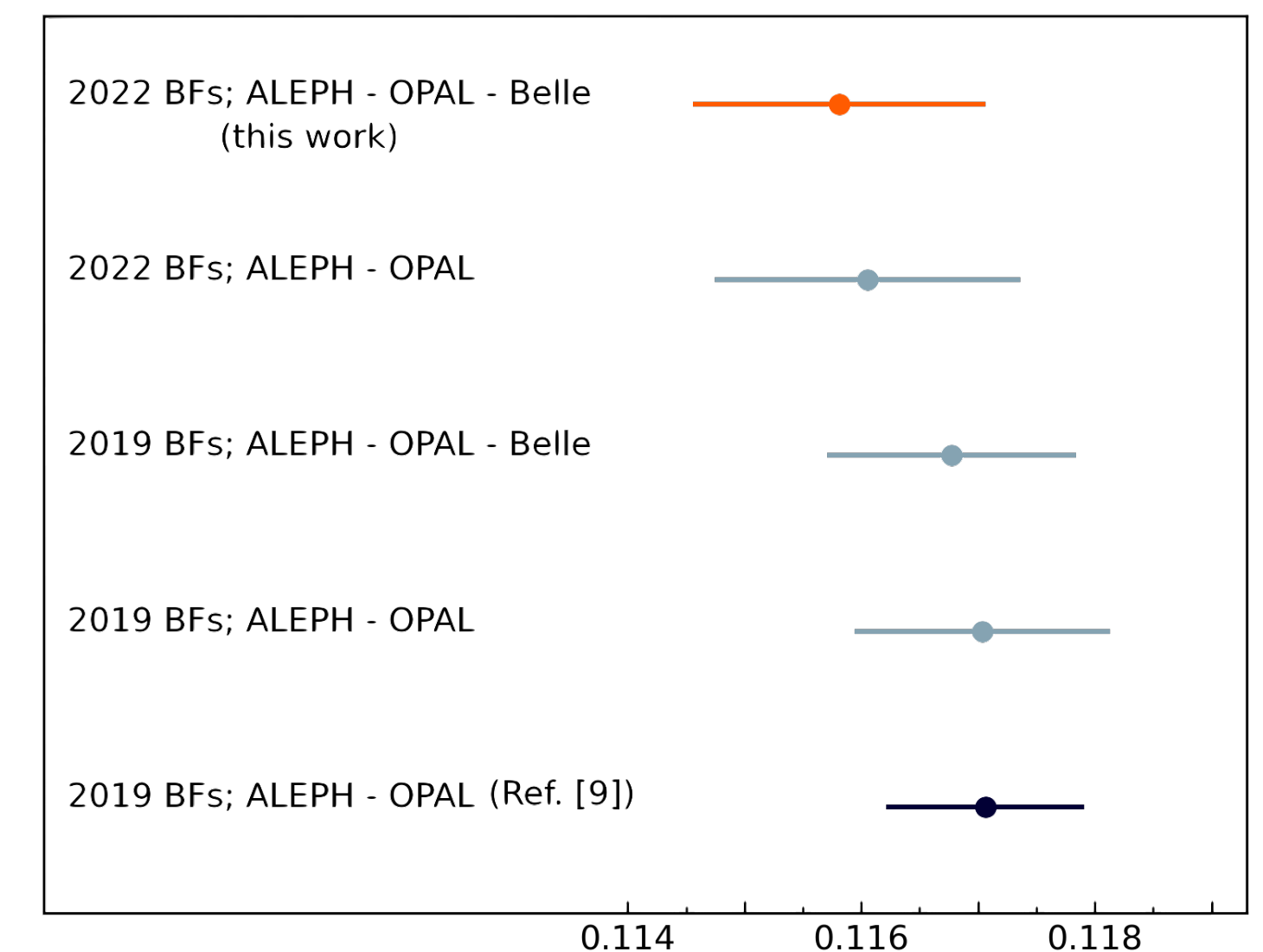
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$$\alpha_s(m_\tau^2) = 0.2983 \pm 0.0092_{\text{stat}} \pm 0.0026_{\text{fit}} \pm 0.0022_{\text{pert}} \pm 0.0025_{\text{DVs}}$$

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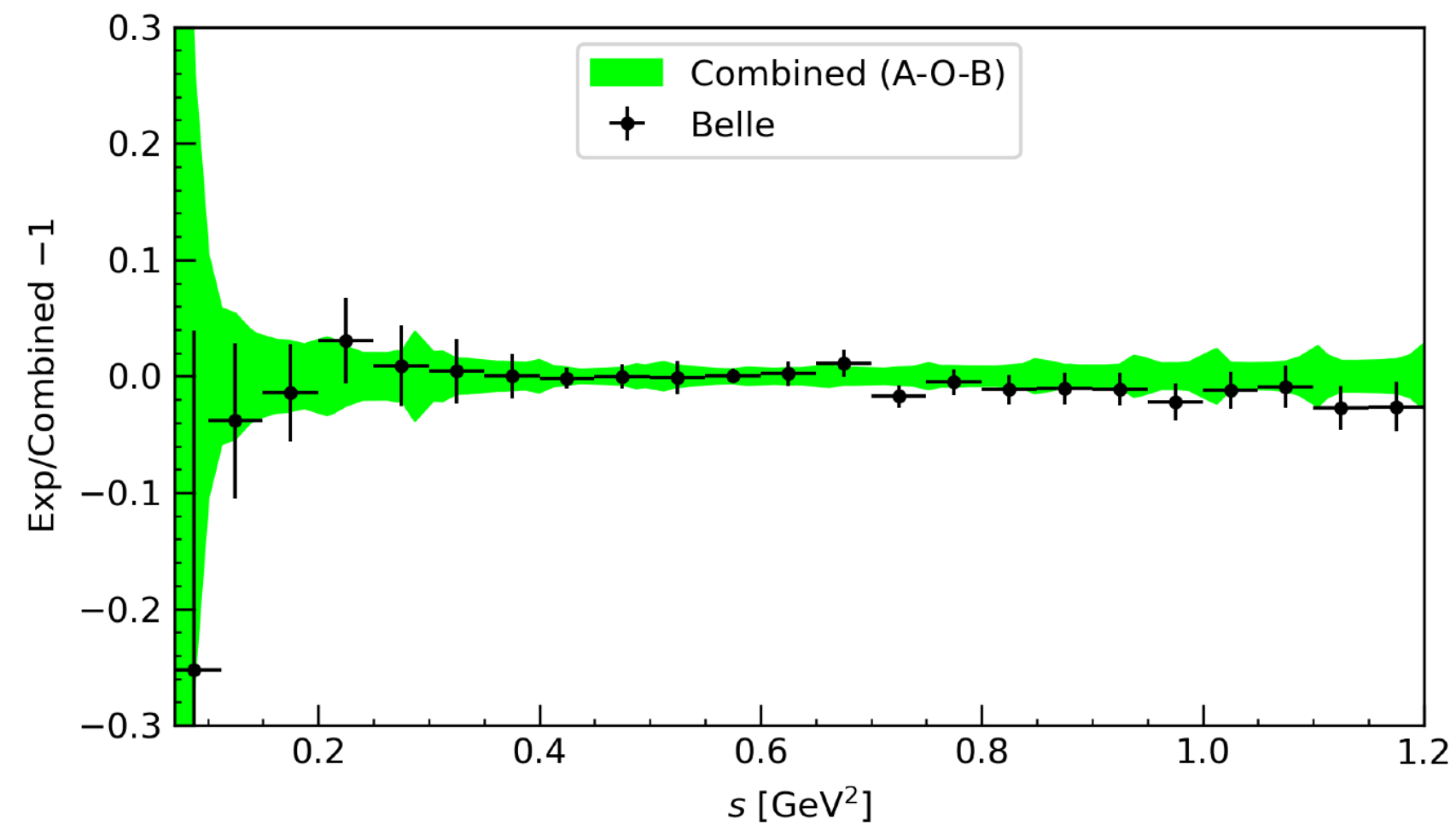
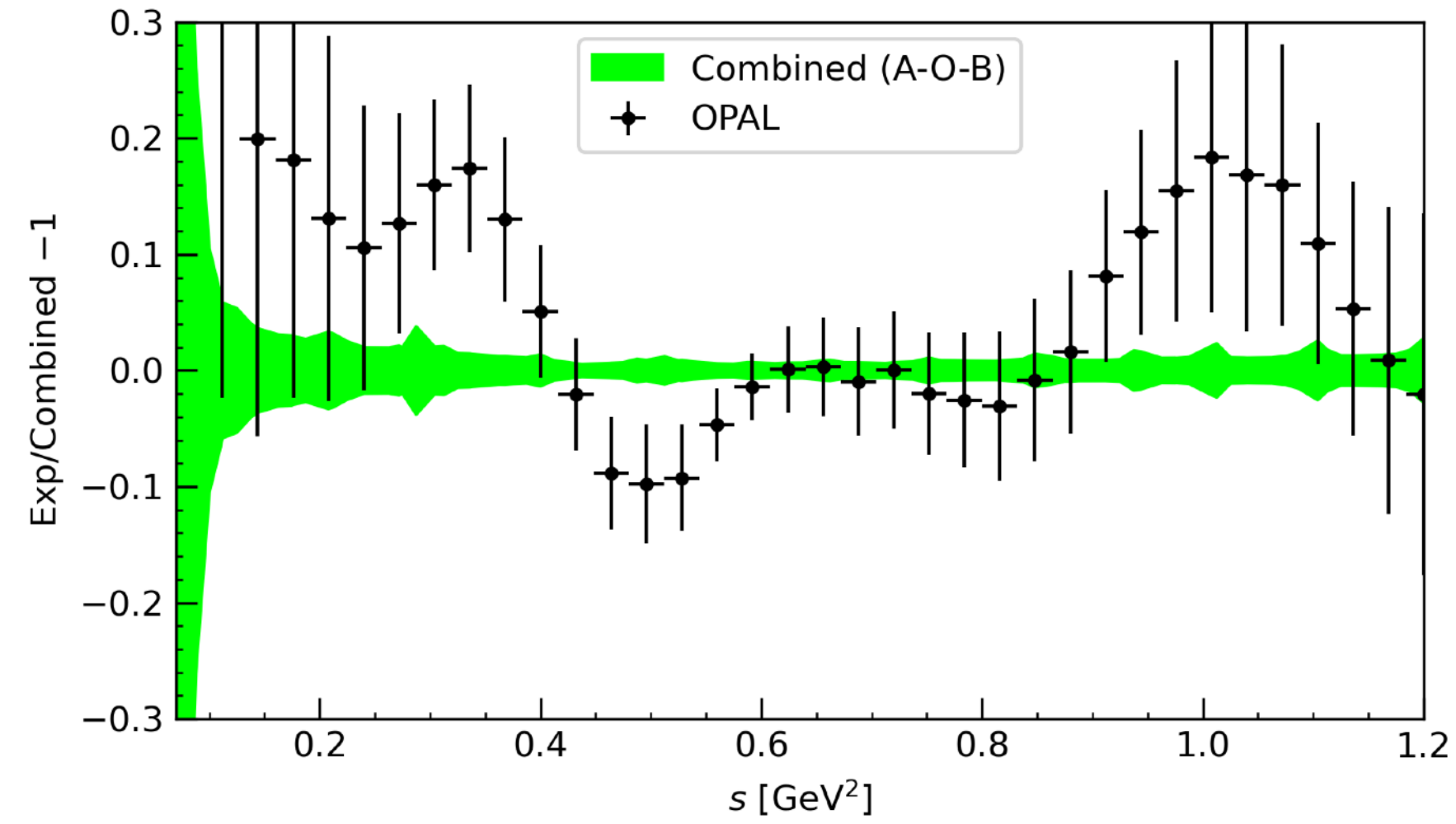
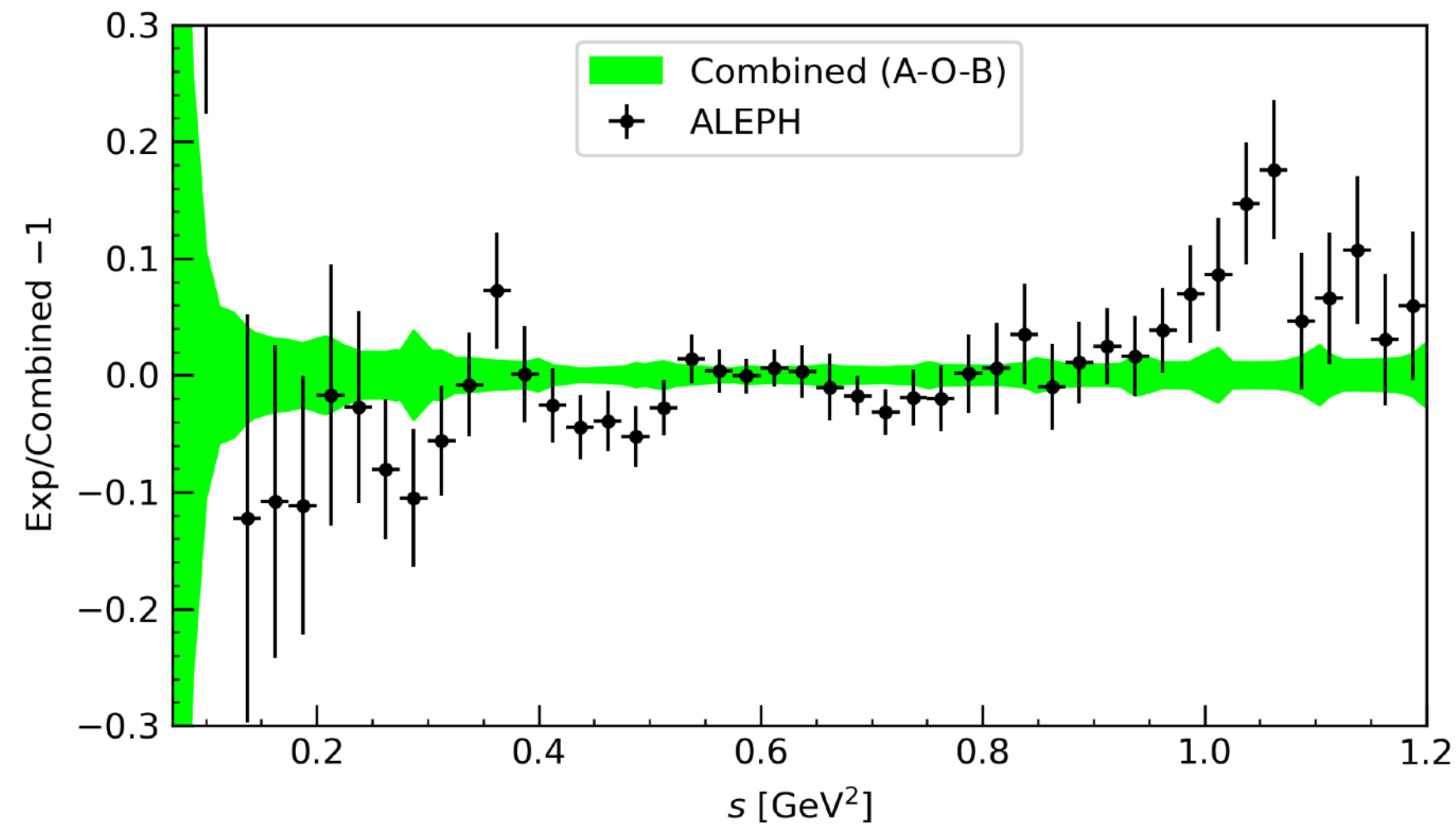
DB, Eiben, Golterman, Maltman, Mansur, and Peris, 2502.08147



Extra

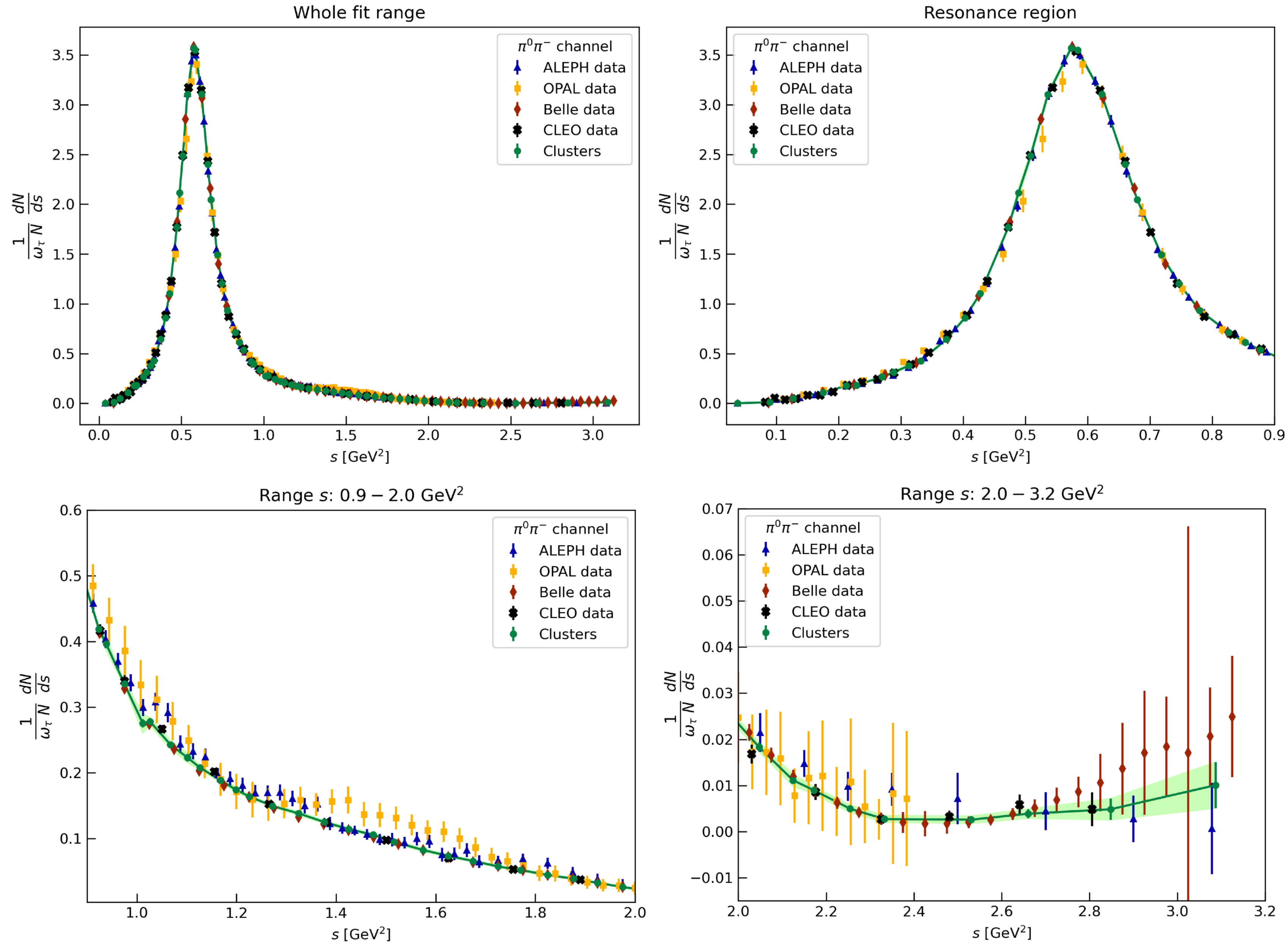
Improved vector-isovector spectral function: 2π channel

Combination of 2π spectra (combined results in green, A-O-B = ALEPH, OPAL, and Belle)

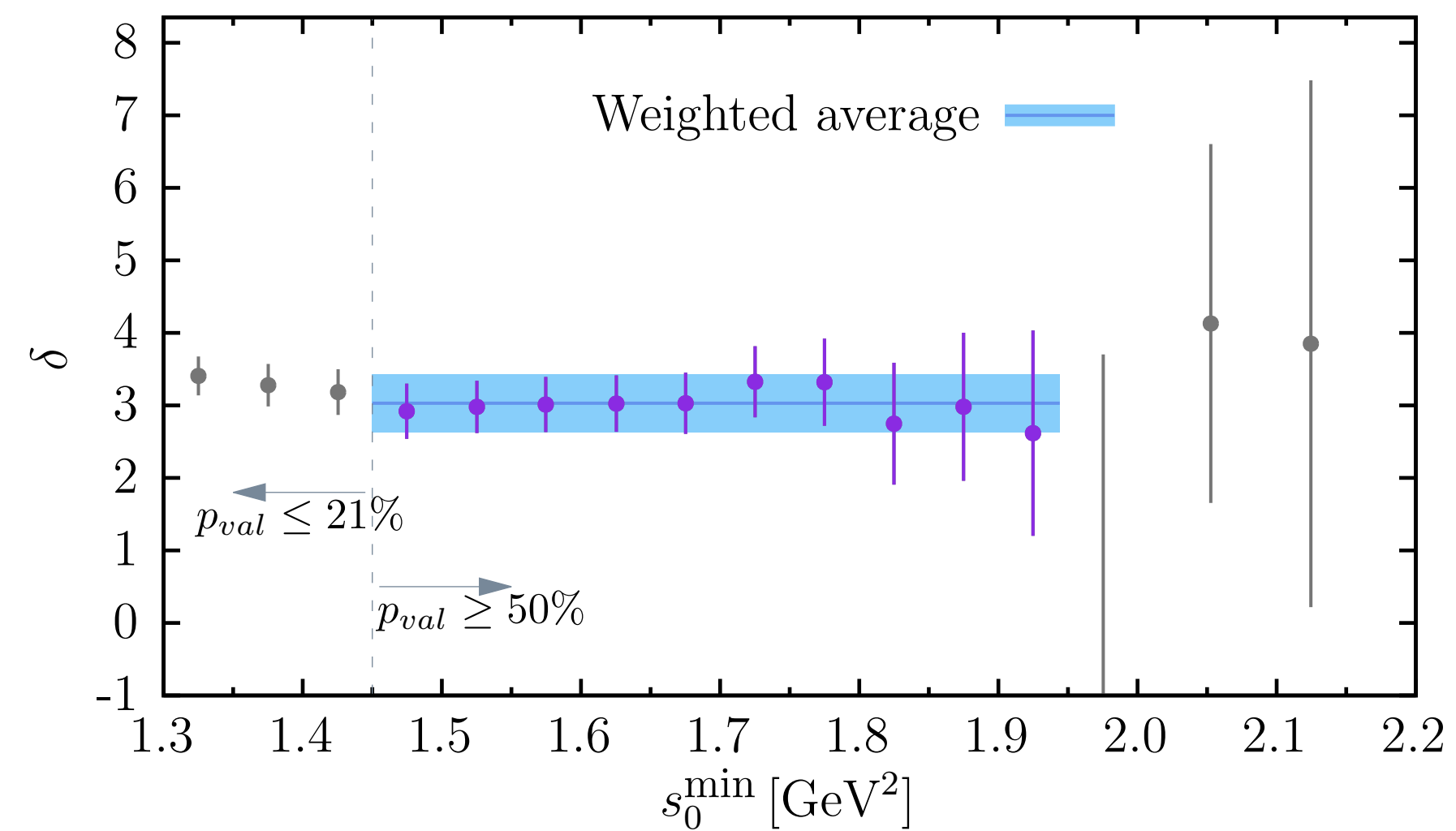
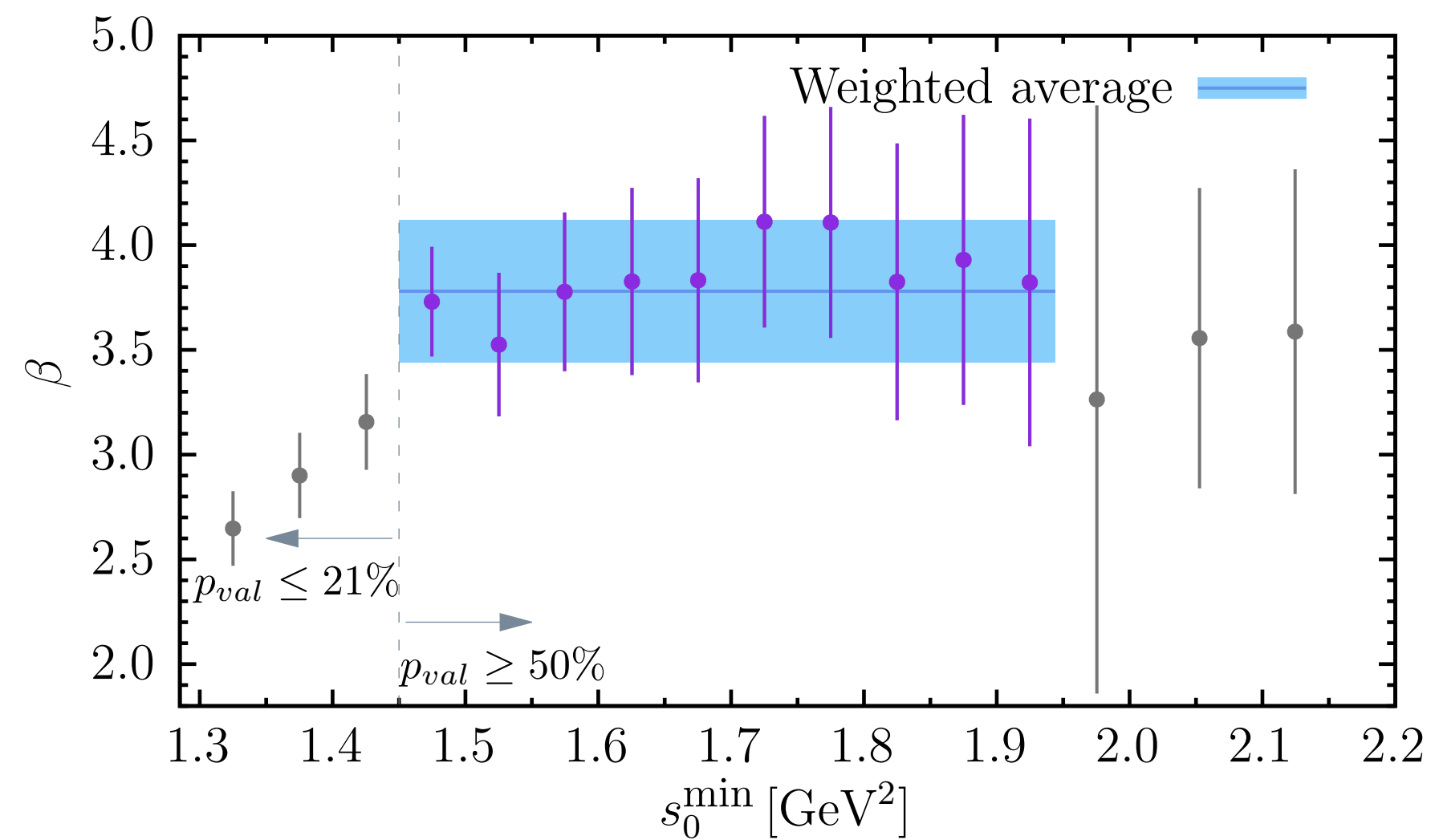
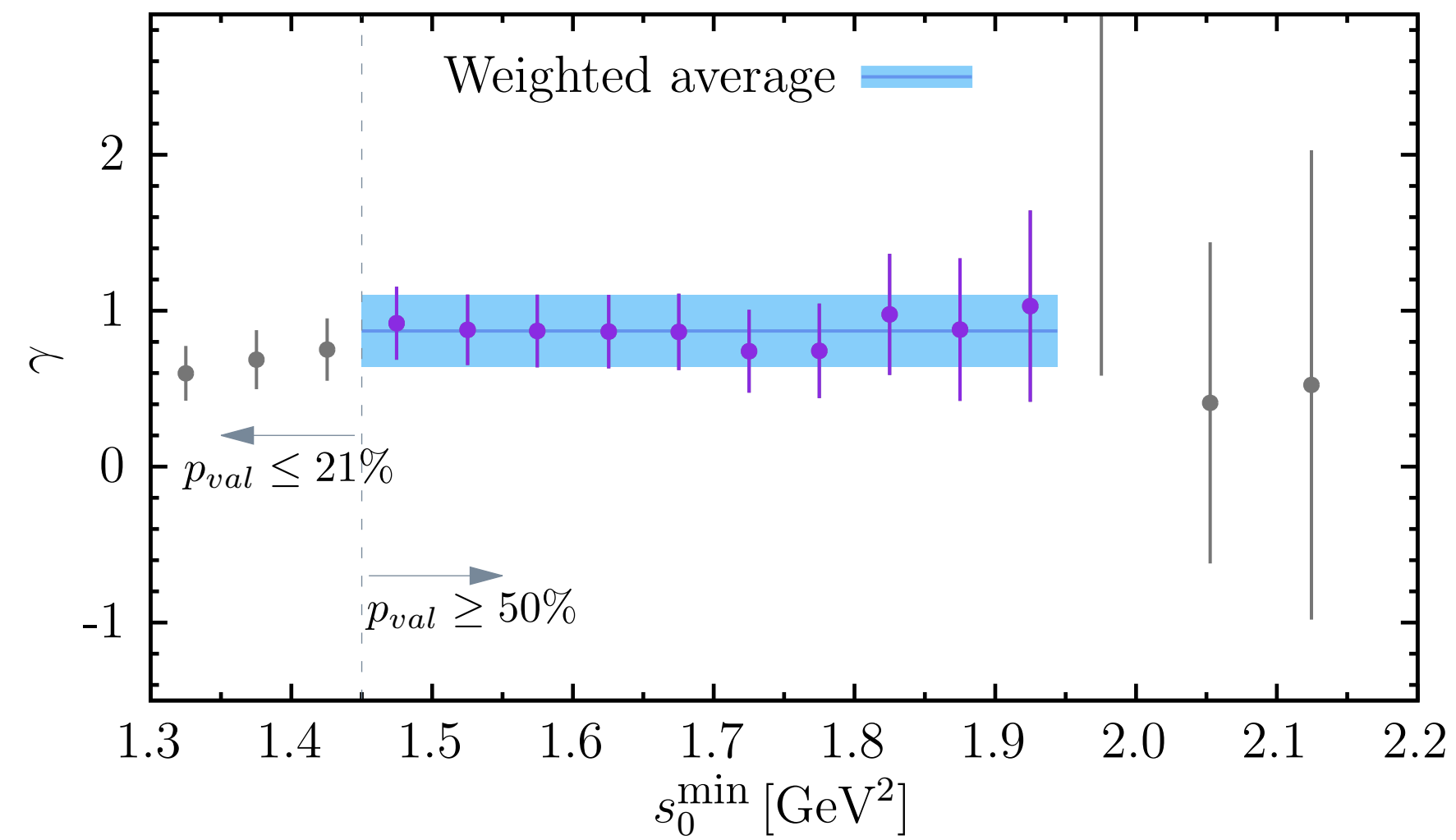
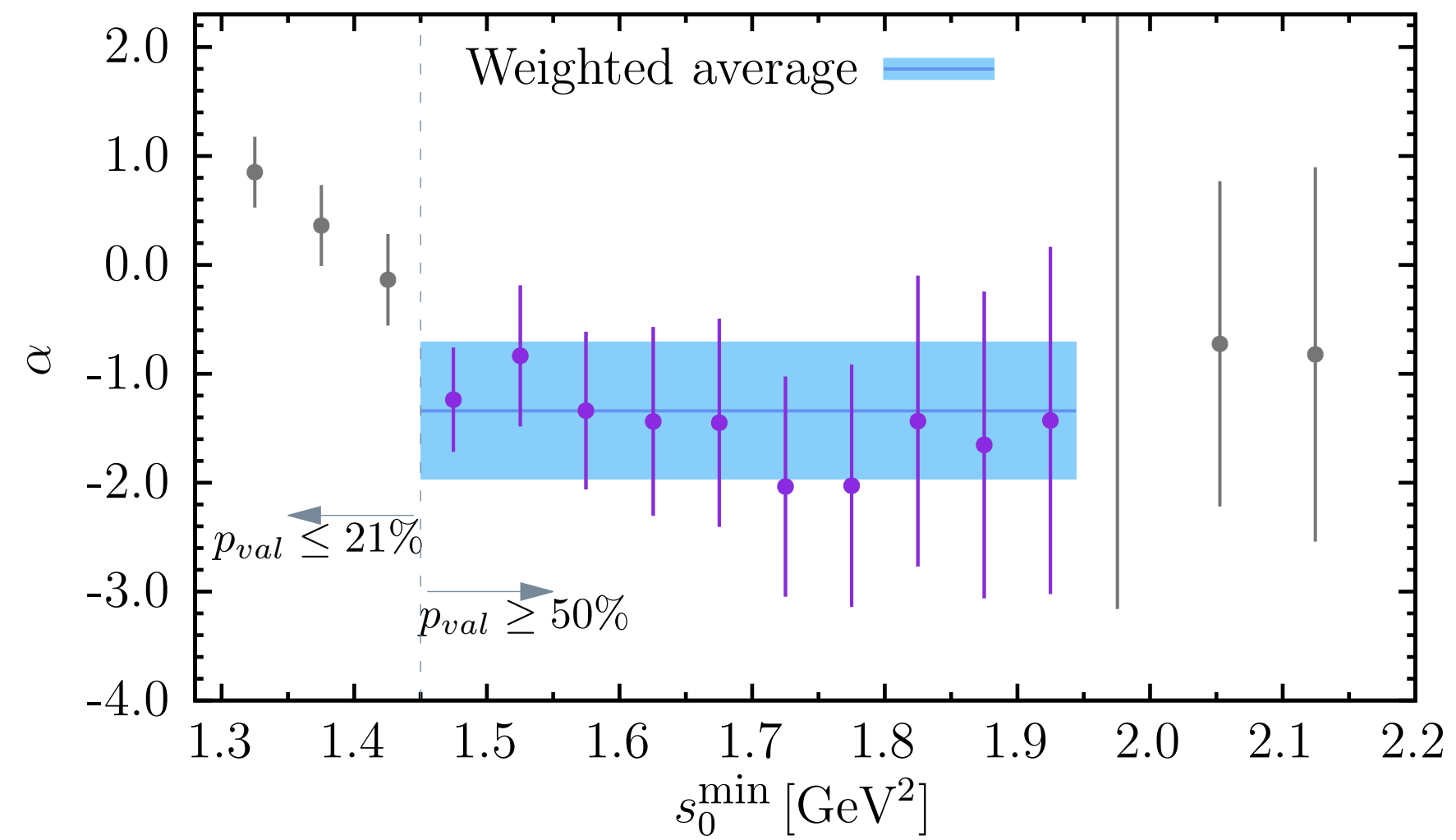


2π data combination dominated by Belle, but Belle spectrum systematically lower for large invariant masses

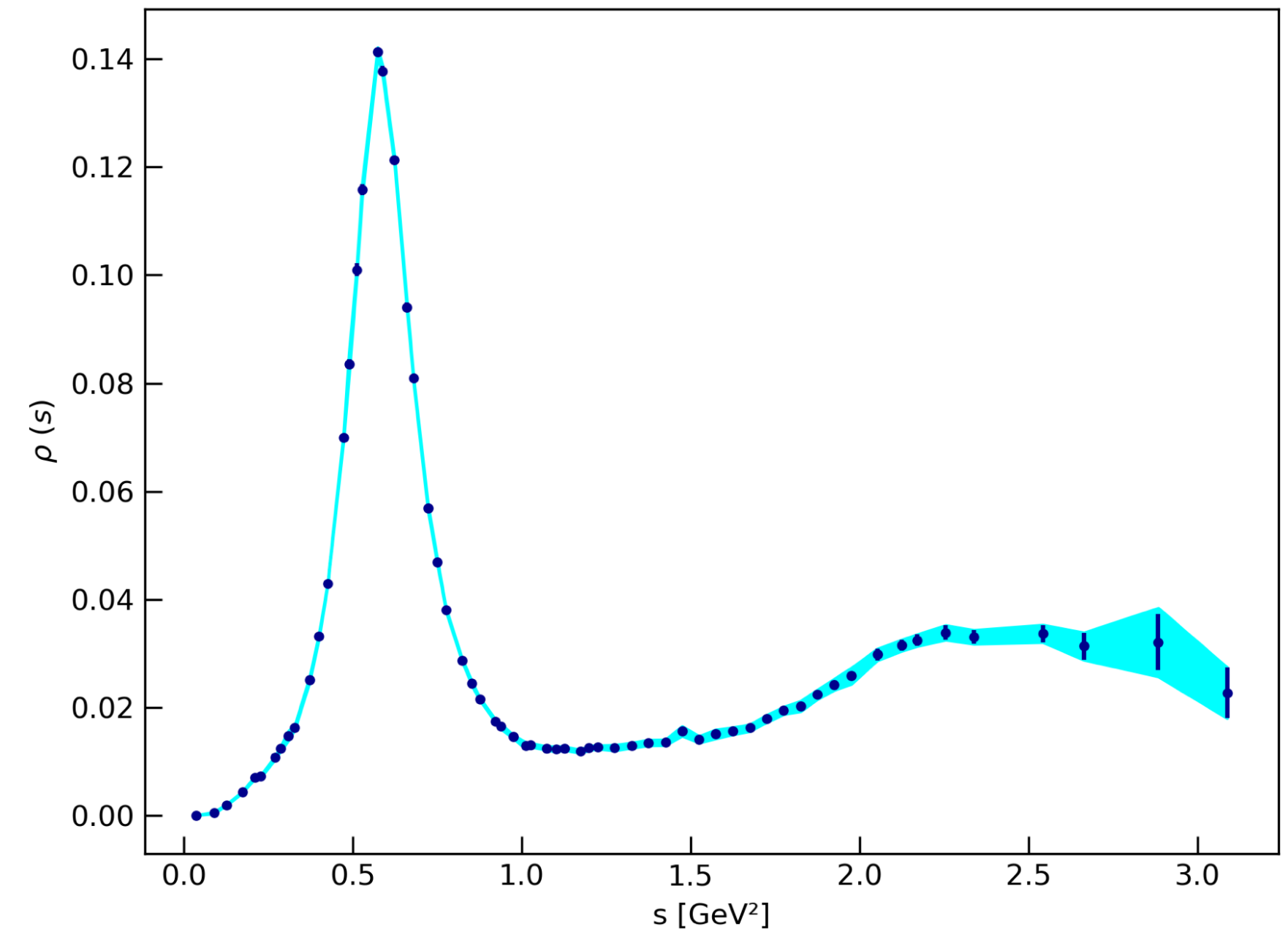
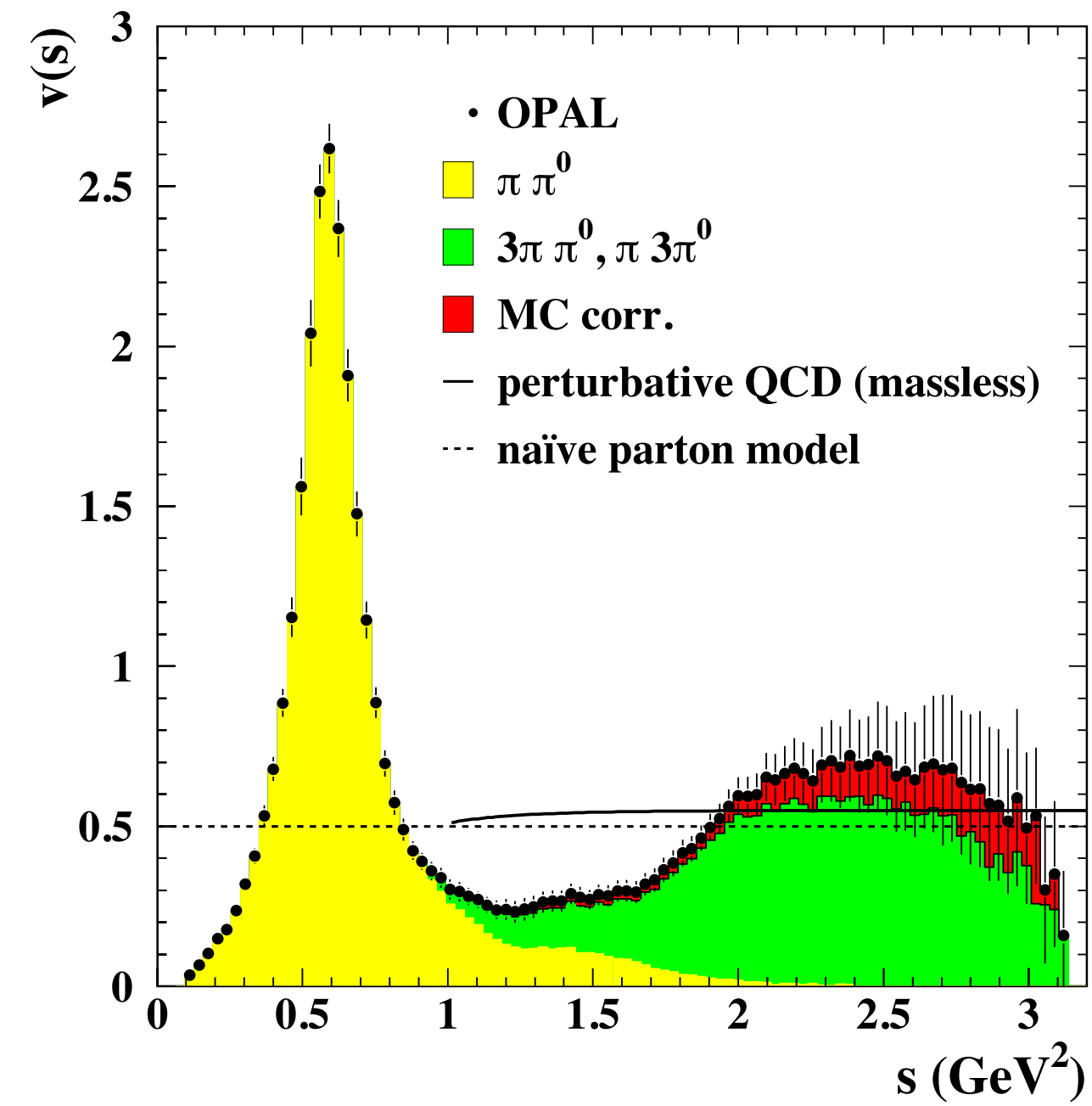
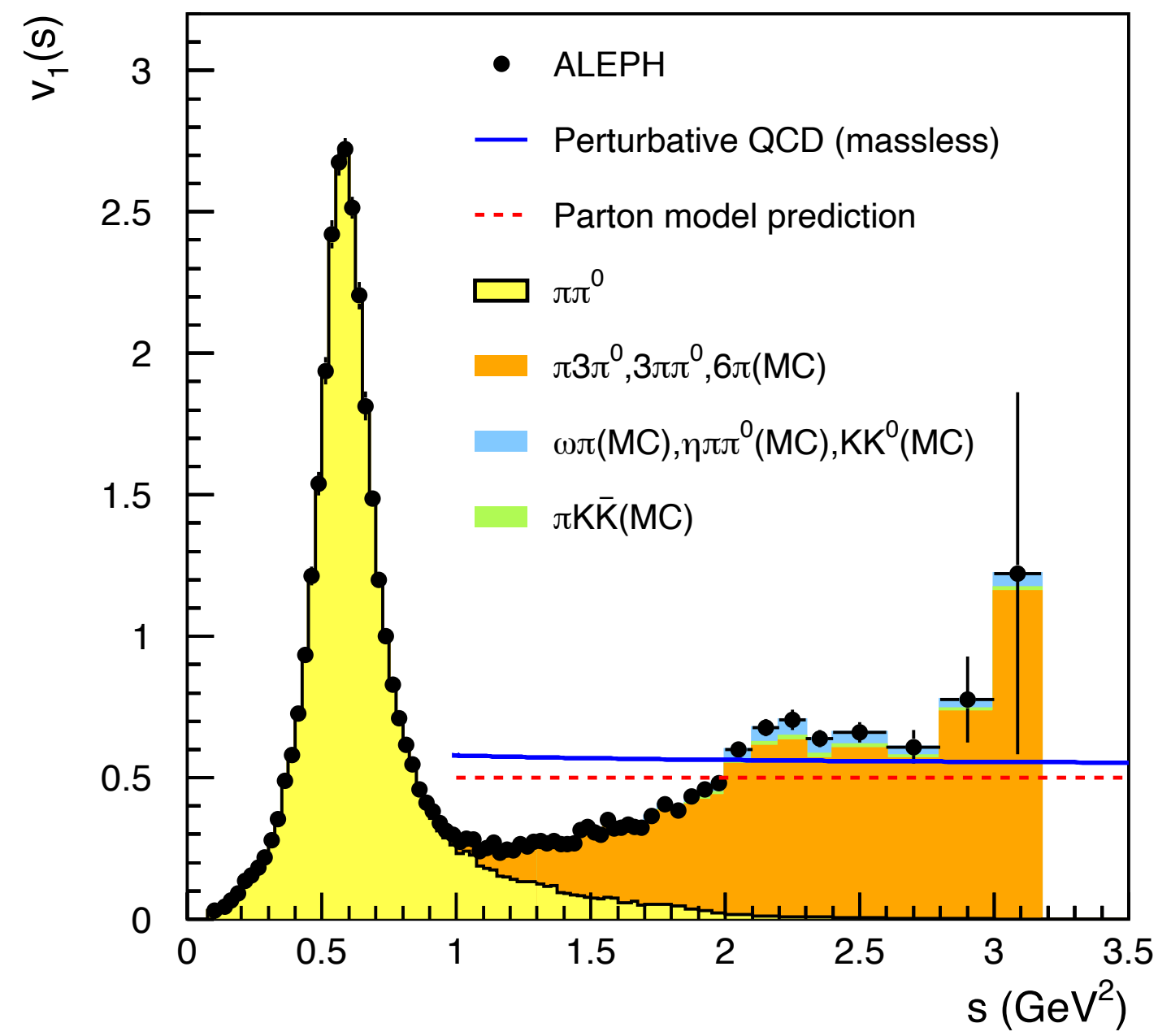
inclusion of CLEO 2π data



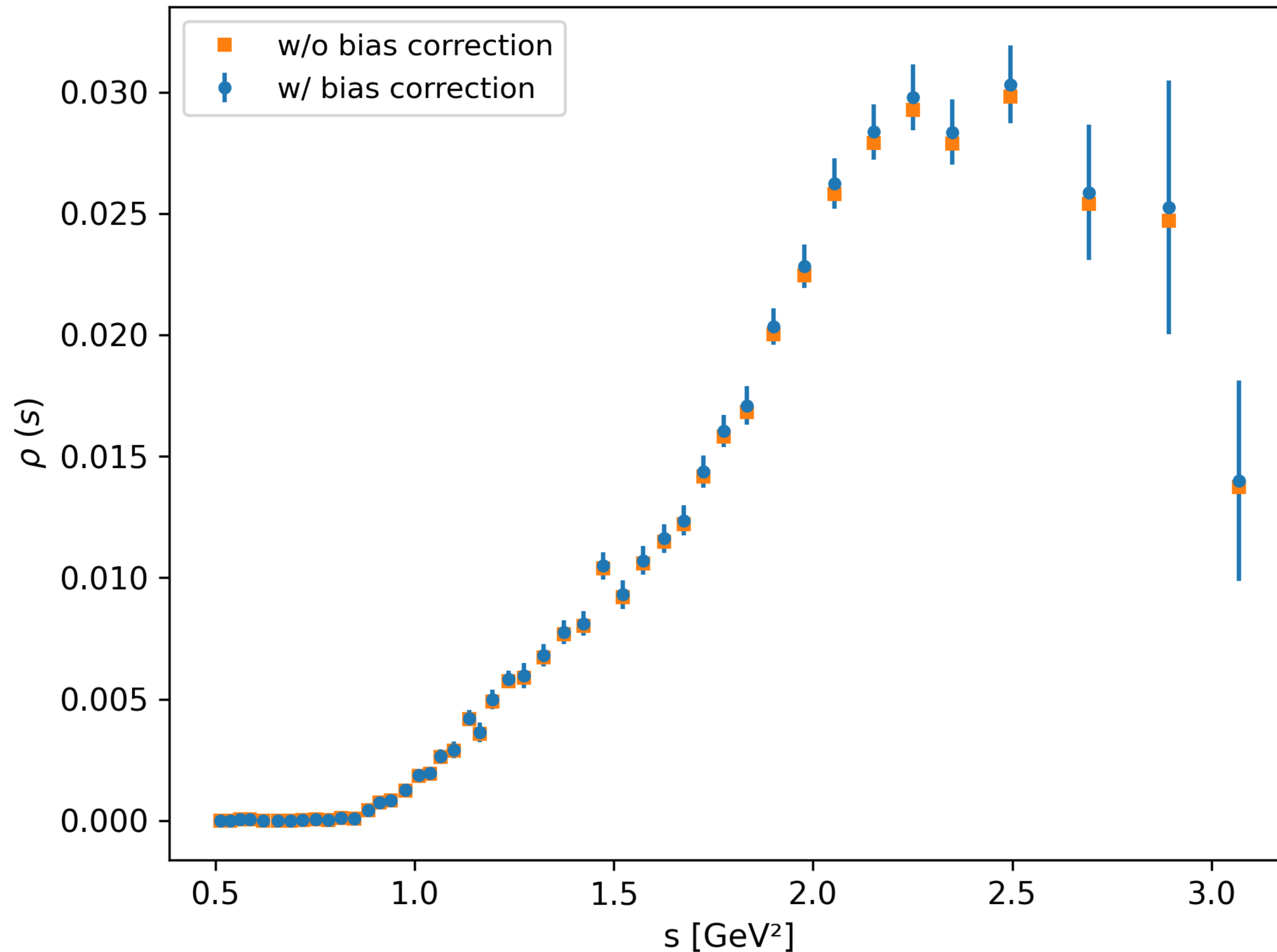
Duality Violation parameters



Vector-isovector spectral functions



d'Agostini bias correction



- We correct for the d'Agostini bias that arise from the multiplication by the Branching Fractions.
- This done with (an approximate version of) NNPDF's iterative algorithm
 Ball et al. (NNPDF) '10
- The effect turns out to be small