

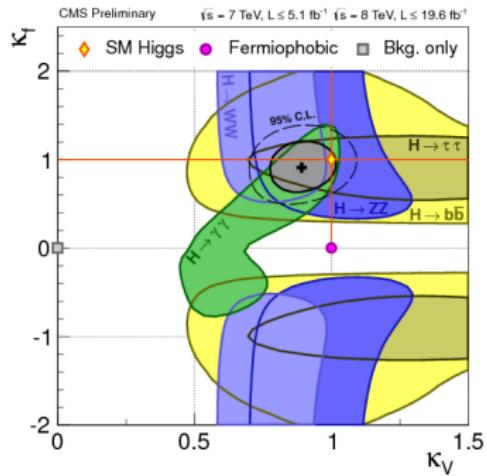
Enhanced Charged Higgs Signal at the LHC

Chung Kao

Homer Dodge Department of Physics and Astronomy
University of Oklahoma

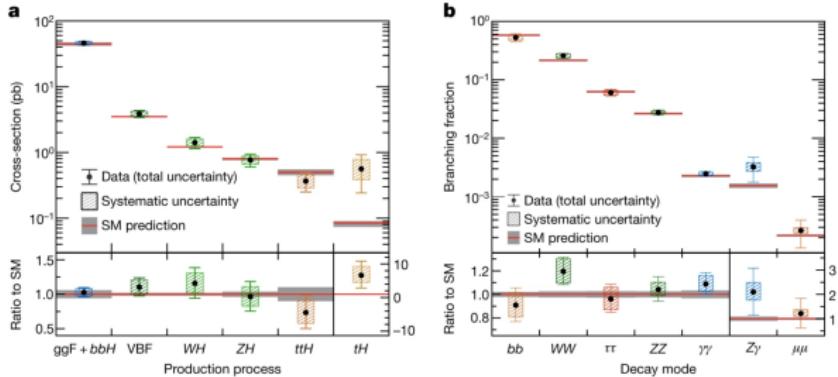
July 09, 2025

Where do we stand?



- A CP-even Higgs boson ($m_h = 125 \text{ GeV}$) was discovered in 2012.
- Where is new physics? Is there a second Higgs doublet?
- This 125 GeV Higgs boson looks almost like the Standard Higgs in the decoupling limit or the alignment limit.

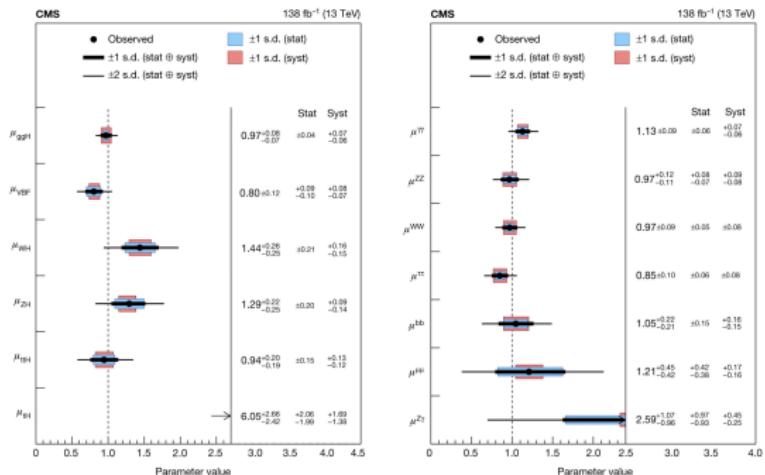
Where do we stand?



a. Cross-sections for different Higgs boson production processes are measured assuming standard model (SM) values for the decay branching fractions. **b.** Branching fractions for different Higgs boson decay modes are measured assuming SM values for the production cross-sections. The lower panels show the ratios of the measured values to their SM predictions. The vertical bar on each point denotes the 68% confidence interval. The p value for compatibility of the measurement and the SM prediction is 65% for **a** and 56% for **b**. Data are from ATLAS Run 2.

The ATLAS Collaboration. A detailed map of Higgs boson interactions by the ATLAS experiment ten years after the discovery. *Nature* 607, 52–59 (2022).

Where do we stand?



Signal-strength parameters extracted for various production modes μ_b , assuming $\mathcal{B}^f = (\mathcal{B}^f)_{\text{SM}}$ (left), and decay channels μ_f , assuming $\sigma_i = (\sigma_i)_{\text{SM}}$ (right). The thick and thin black lines indicate the 1-s.d. and 2-s.d. confidence intervals, respectively, with the systematic (syst) and statistical (stat) components of the 1-s.d. interval indicated by the red and blue bands, respectively. The vertical dashed line at unity represents the values of μ_b and μ_f in the SM. The covariance matrices of the fitted signal-strength parameters are shown in Extended Data Fig. 5. The P values with respect to the SM prediction are 3.1% and 30.1% for the left plot and the right plot, respectively. The P value corresponds to the probability that a result deviates as much, or more, from the SM prediction as the observed one.

The CMS Collaboration. A portrait of the Higgs boson by the CMS experiment ten years after the discovery. Nature 607, 60–68 (2022).

General Two Higgs Doublet Models

The Yukawa Lagrangian in a general two Higgs doublet model:

$$\begin{aligned} \mathcal{L}_Y = & \frac{-1}{\sqrt{2}} \sum_{F=U,D,L} \bar{F} \left\{ \left[\kappa^F s_{\beta-\alpha} + \rho^F c_{\beta-\alpha} \right] h^0 + \left[\kappa^F c_{\beta-\alpha} - \rho^F s_{\beta-\alpha} \right] H^0 \right. \\ & - i \operatorname{sgn}(Q_F) \rho^F A^0 \Big\} P_R F - \bar{U} \left[V \rho^D P_R - \rho^{U\dagger} V P_L \right] D H^+ \\ & - \bar{\nu} \left[\rho^L P_R \right] L H^+ + \text{H.c.} \end{aligned}$$

- $P_{L,R} \equiv (1 \mp \gamma_5)/2$, $c_{\beta-\alpha} \equiv \cos(\beta - \alpha)$, $s_{\beta-\alpha} \equiv \sin(\beta - \alpha)$,
- α is the mixing angle between neutral Higgs scalars,
- $\tan \beta \equiv v_2/v_1$ is the ratio of the vacuum expectation values of the two Higgs doublets,
- Q_F is the fermion charge,
- κ matrices are diagonal and fixed by fermion masses to $\kappa^F = \sqrt{2}m_F/v$ with $v \approx 246$ GeV,
- ρ matrices contain both diagonal and off-diagonal elements with free parameters.

2HDMC

- A computer code that evaluates decay widths for h^0, H^0, A^0 , and H^\pm including higher-order QCD corrections, theoretical constraints and much more, by D. Eriksson, J. Rathsman, and O. Stål¹.
- For neutral Higgs bosons, QCD corrections are evaluated with $m_f(M_\phi)$ for $h^0, H^0, A^0 \rightarrow f\bar{f}$.
- FCNH decays are evaluated with $\rho_{qq'}$ such as $H^0 \rightarrow tc$ and $A^0 \rightarrow tc$.
- For charged Higgs decays into quarks, QCD corrections are calculated with $\rho_{qq'}$ scaled by running mass $m_q(M_{H^\pm})$, such as $H^\pm \rightarrow bc, tb$.

¹ <https://2hdmc.hepforge.org/>;

2HDMC - Two-Higgs-Doublet Model Calculator, D. Eriksson, J. Rathsman, O. Stål Comput. Phys. Commun. 181:189-205 (2010); Comput. Phys. Commun. 181:833-834 (2010); [arXiv:0902.0851]

MadGraph with 2HDM

- *The general Two-Higgs-Doublet Model*, by Claude Duhr, Michel Herquet, and Celine Degrande².
- 2HDM_UFO: leading order.
- 2HDM_NLO: QCD corrections with virtual and real gluons.
- For neutral Higgs bosons, to compare with 2HDMC for Higgs decays, we need 2HDM_NLO and scale $m(\text{pole})$ to $m(Q = M_H)$.

² <https://feynrules.irmp.ucl.ac.be/wiki/2HDM>;

Automatic evaluation of UV and R2 terms for beyond the Standard Model Lagrangians: a proof-of-principle, Celine Degrande, Comput. Phys. Commun. 197 (2015) 239-262.

Yukawa Interactions of Charged Higgs Boson

To study the interaction of $H^+ \bar{c} b$, we follow the Lagrangian ³:

$$\mathcal{L}_Y = -\bar{U} \left[V \rho^D P_R - \rho^{U\dagger} V P_L \right] D H^+ - \bar{\nu} \left[\rho^L P_R \right] L H^+ + \text{H.c.} \quad (1)$$

where $P_{L,R} \equiv (1 \mp \gamma_5)/2$. For simplicity, let us choose ρ as real parameters. The coupling of $H^+ \bar{c} b$ can be simplified as

$$\begin{aligned} \mathcal{L}_{H^+ \bar{c} b} &= -\bar{c} (V_{cd} \rho_{db} + V_{cs} \rho_{sb} + V_{cb} \rho_{bb}) P_R b H^+ \\ &\quad + (\rho_{uc}^* V_{ub} + \rho_{cc}^* V_{cb} + \rho_{tc}^* V_{tb}) P_L b H^+ + \text{H.c.} \\ &\simeq + \bar{c} [\rho_{tc}^* V_{tb} P_L] b H^+ \text{H.c.} \end{aligned}$$

The coupling $\lambda_{H^+ \bar{c} b} \propto \rho_{tc} V_{tb}$ is not suppressed by V_{cb} .

We propose that the process $c g \rightarrow b H^+$ followed by $H^+ \rightarrow \bar{b} c$ might be a promising discovery channel for the charged Higgs boson (H^+).

³S. Davidson and H. E. Haber, Phys. Rev. D **72** (2005), 035004;
F. Mahmoudi and O. Stal, Phys. Rev. D **81** (2010), 035016.

$$cg \rightarrow bH^\pm \rightarrow bbc$$

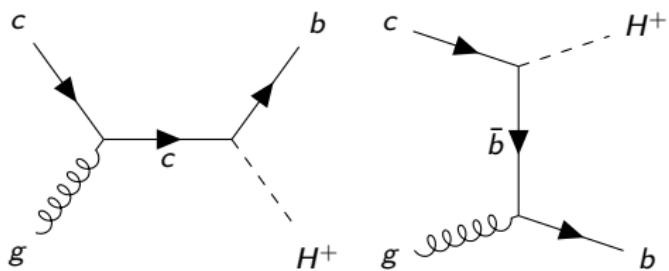
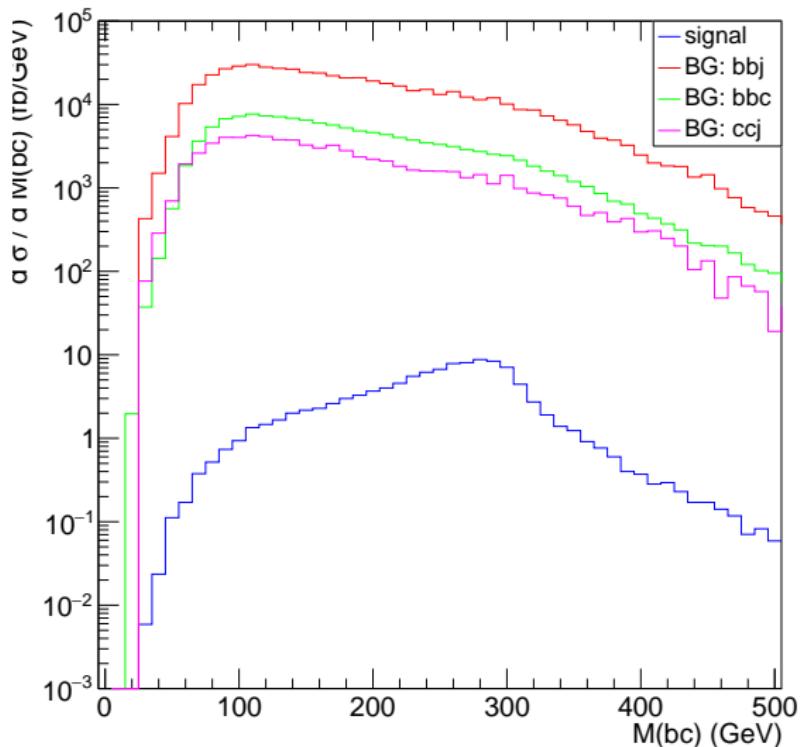


Figure: Leading-order Feynman diagrams for $cg \rightarrow bH^+$.

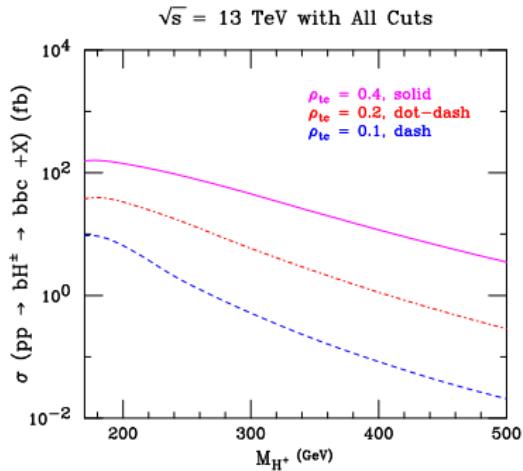
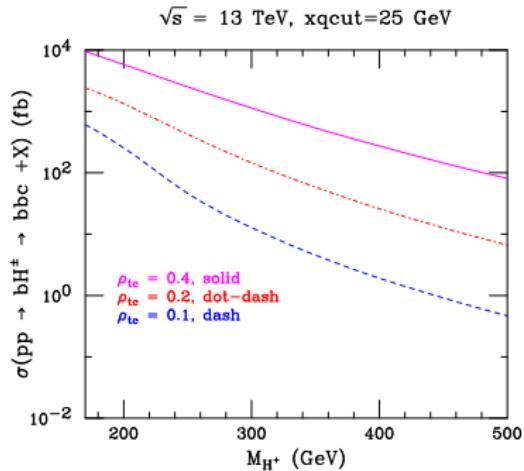
D. K. Ghosh, W. S. Hou and T. Modak, Phys. Rev. Lett. **125** (2020)
no.22, 221801.

Invariant Mass Distributions

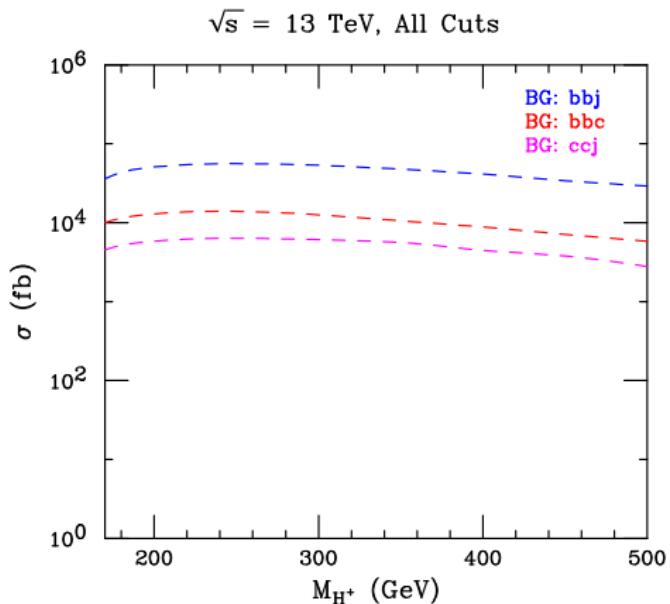
$\sqrt{s} = 13 \text{ TeV}, M(bc)$



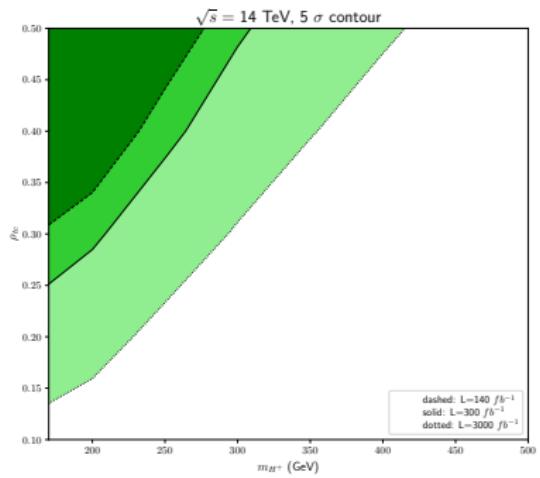
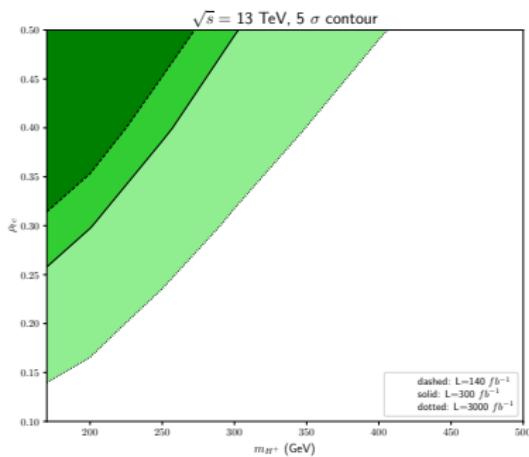
Cross Section versus M_{H^+}



Physics Background



5 σ discovery contours for $c g \rightarrow b H^\pm$

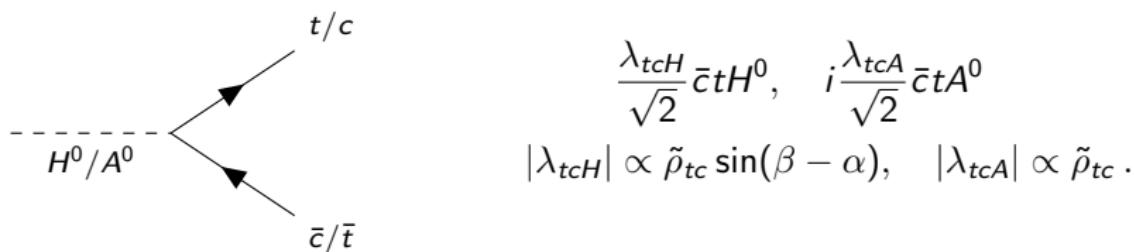


Summary for Enhanced Charged Higgs Signal

- In a general two Higgs doublet model, $H^\pm \rightarrow bc$ does not have V_{cb} suppression.
- The cross section of $cg \rightarrow bH^\pm \rightarrow bbc$ in G2HDM can be significantly larger than that in models with Type II Yukawa couplings with $1/V_{cb}^2$ enhancement.
- This discovery channel offers great promise to search for BSM new physics.
- Its cross section is proportional to ρ_{tc}^4 and might provide a good opportunity to measure ρ_{tc} .

Charming Higgs Bosons with Top Quarks

- With strong constraint $\mathcal{B}(t \rightarrow ch^0) \leq 3.7 \times 10^{-4}$, placed by ATLAS Collaboration⁴ and CMS Collaboration⁵
 - $\lambda_{tch} \leq 0.036$, with $\lambda_{tch} \propto 1.92 \times \sqrt{\mathcal{B}(t \rightarrow ch^0)}$
 - $|\lambda_{tch}| \propto \tilde{\rho}_{tc} \cos(\beta - \alpha)$, where $\tilde{\rho}_{tc} \approx \sqrt{(|\rho_{tc}|^2 + |\rho_{ct}|^2)/2}$
- $H^0/A^0 \rightarrow t\bar{c} + \bar{t}c$



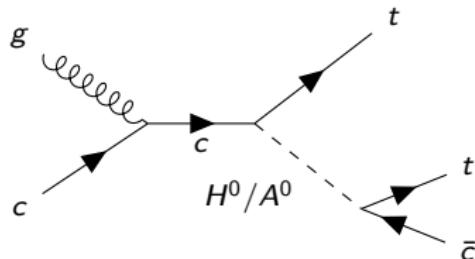
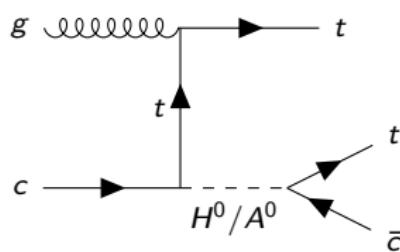
- In the alignment (decoupling) limit, $\cos(\beta - \alpha) \rightarrow 0$, $\sin(\beta - \alpha) \approx 1$, $|\lambda_{tcH}|$ and $|\lambda_{tcA}|$ are not suppressed

⁴ G. Aad et al. [ATLAS], Eur. Phys. J. C **84**, no.7, 757 (2024).

⁵ A. Hayrapetyan et al. [CMS], [arXiv:2407.15172 [hep-ex]]

Charming Higgs Bosons with Top Quarks

- $pp \rightarrow t\phi^0 \rightarrow t(tc)$, where $\phi^0 = H^0$ or A^0



- All Top: $ttc \rightarrow bjjcbl\nu$
 - dominant physics background: (ttj) $pp \rightarrow t\bar{t}j \rightarrow b\bar{b}jjj\ell\nu$
- Same Sign Top: $ttc \rightarrow bbcl^\pm\ell^\pm\nu\nu$ where $\ell = e$ or μ
 - dominant physics background: (ttw) $pp \rightarrow t\bar{t}w^\pm \rightarrow b\bar{b}jj\ell^\pm\ell^\pm\nu$
 - other background: (ttz) $pp \rightarrow t\bar{t}z \rightarrow b\bar{b}jj\ell^\pm\ell^\pm\ell^\mp$

Decays of Heavy Neutral Higgs Bosons

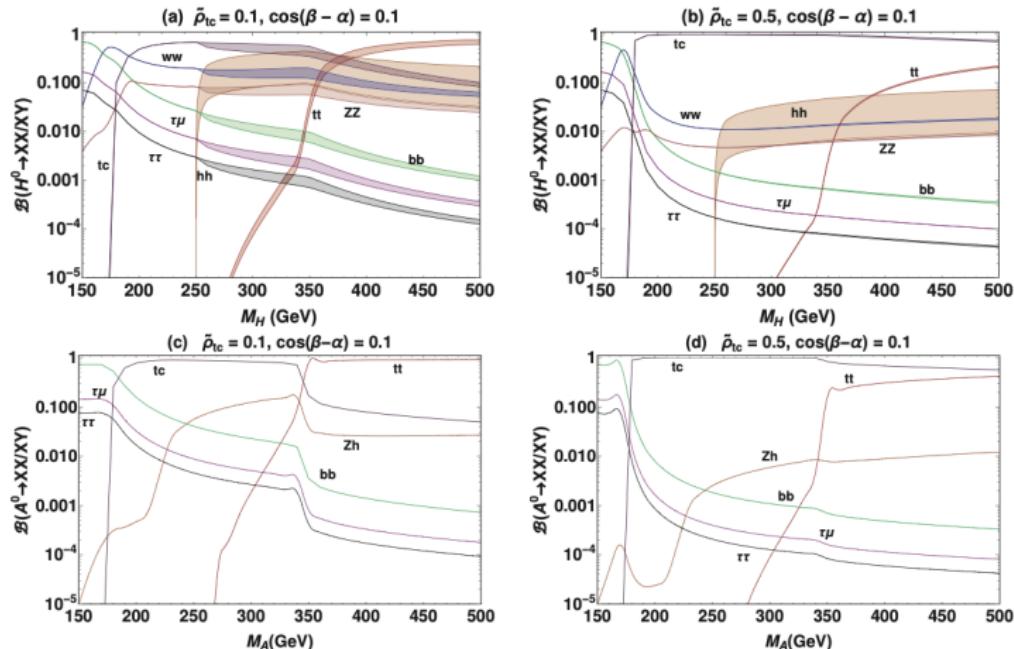
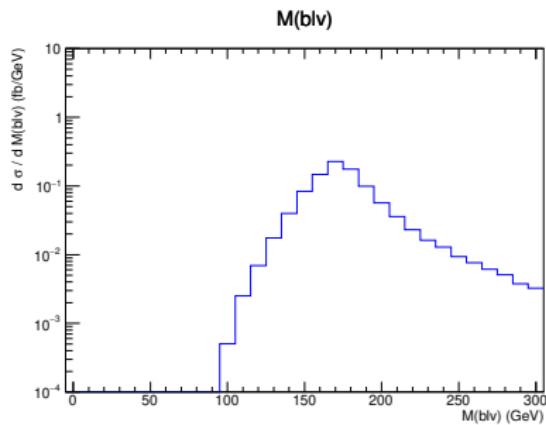
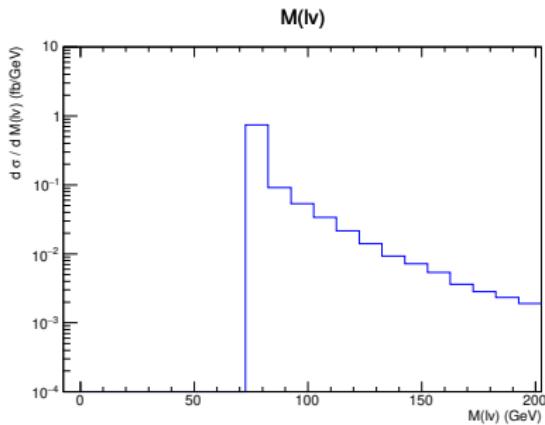


FIG. 1: Major two body decays of H^0 vs M_H for (a) $\tilde{\rho}_{tc} = 0.1$, and (b) $\tilde{\rho}_{tc} = 0.5$, with $\tilde{\rho}_{\tau\mu} = 0.01$, $\lambda_5 = 0$ and $1 \leq \tan\beta \leq 10$. The analogous case for A^0 is given in (c) and (d).

All Top final state ($b jj cb \ell \nu$)

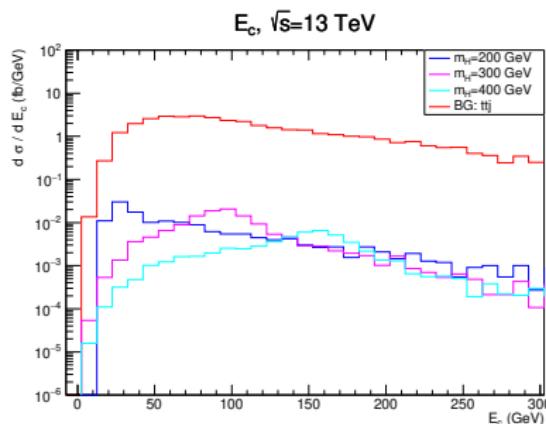
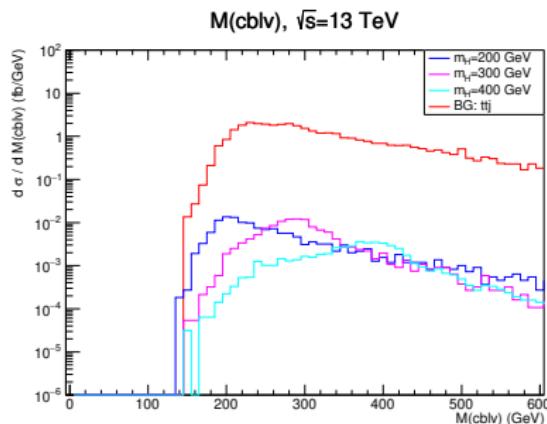
- Basic requirements

- (i) five jets including two b -jet and one c -jet with $P_T(b, c, j) \geq 25$ GeV, $|\eta(b, c, j)| \leq 2.5$
- (iii) a lepton with $p_T(\ell) \geq 20$ GeV, and $|\eta(\ell)| \leq 2.5$,
- (iv) $\Delta R(cj, jj, bj, bc, \ell j, \ell b, \ell c) \geq 0.4$,
- (v) missing transverse energy $\text{MET} = \cancel{E}_T \geq 25$ GeV,
- Reconstruct P_ν with m_w and m_t ($(P_\nu + P_\ell)^2 = m_W^2$, $(P_\nu + P_\ell + P_b)^2 = m_t^2$).



All Top final state ($b jj cb \ell \nu$)

- Reconstruct m_H with $c-jet$.
- Energy of the charm quark (E_c) in the rest frame of the Higgs Boson.
The peak E_c^* is given by $E_c^* = \frac{m_H}{2} \left[1 + \frac{m_c^2}{m_H^2} - \frac{m_t^2}{m_H^2} \right]^{1/2}$.
 $E_c^*(200) = 25.4$ GeV, $E_c^*(300) = 100.3$ GeV, $E_c^*(400) = 162.7$ GeV.

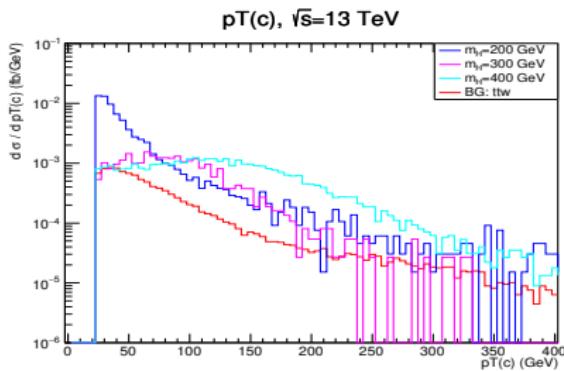


⁶T. Han, J. Jiang and M. Sher, Phys. Lett. B **516**, 337-344 (2001)

Same Sign Top final state^{7, 8}($b\bar{b}c\ell^{\pm}\ell^{\pm}\nu\nu$)

- Basic requirements

- (i) three jets including two b -jet and one c -jet with $P_T(b, c) \geq 25$ GeV, $|\eta(b, c, j)| \leq 2.5$
- (ii) two lepton with $p_T(\ell) \geq 20$ GeV, and $|\eta(\ell)| \leq 2.5$,
- (iii) $\Delta R(bc, \ell b, \ell c, \ell\ell) \geq 0.4$,
- (iv) missing transverse energy $\text{MET} = \cancel{E}_T \geq 30$ GeV,



- For $m_H > 400$ GeV, $P_T(c)$ cuts can reduce background

⁷ W. Hou, M. Kohda, T. Modak, Phys.Lett. B **786** (2018) 212-216

⁸ CMS Collaboration, Phys. Lett. B **850** (2024) 138478

Statistical Significance

N_{SS} is calculated by:

$$N_{SS} = \sqrt{2 \times (N_S + N_B) \ln(1 + N_S/N_B) - 2 \times N_S}.$$

Here N_S and N_B are number of signal and background events, where $N_S = \sigma_s \times L$, $N_B = \sigma_b \times L$, and L is the luminosity of the LHC.⁹

⁹ G. Cowan,K. Cranmer,E. Gross and O. Vitells, Eur. Phys. Jour. C **71** (2011) 1554. ▶ ⟲ ⟳ ⟴ ⟵ ⟶ ⟷ ⟸ ⟹ ⟺ ⟻ ⟼ ⟽ ⟾ ⟷ ⟸ ⟹ ⟺ ⟻ ⟼ ⟽ ⟾

Cross section & N_{SS}

- All Top state at $\sqrt{s} = 13$ TeV:

$\tilde{\rho}_{tc} = 0.4$	Signal	ttj	$N_{SS}(L=1000 \text{ fb}^{-1})$
$m_H = 200$ GeV	0.2866 fb	14.80 fb	2.35
$m_H = 300$ GeV	0.3545 fb	31.62 fb	1.99
$m_H = 400$ GeV	0.1211 fb	18.17 fb	0.90

- Same Sign Top final state at $\sqrt{s} = 13$ TeV:

$\tilde{\rho}_{tc} = 0.4$	Signal	ttw	ttz	$N_{SS}(L=1000 \text{ fb}^{-1})$
$m_H = 200$ GeV	0.2027 fb	0.0339 fb	0.0015 fb	22.41
$m_H = 300$ GeV	0.1791 fb	0.0339 fb	0.0015 fb	20.36
$m_H = 400$ GeV	0.0695 fb	0.0339 fb	0.0015 fb	9.43

5 σ Discovery Contours

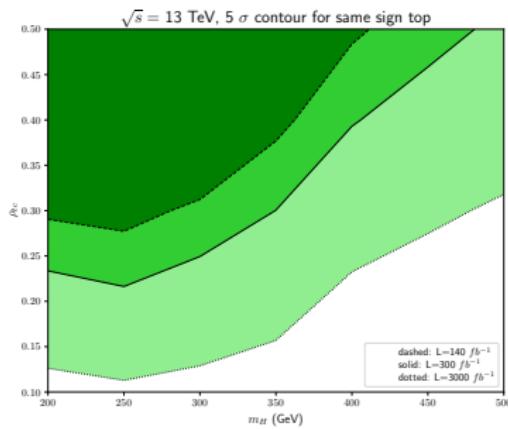
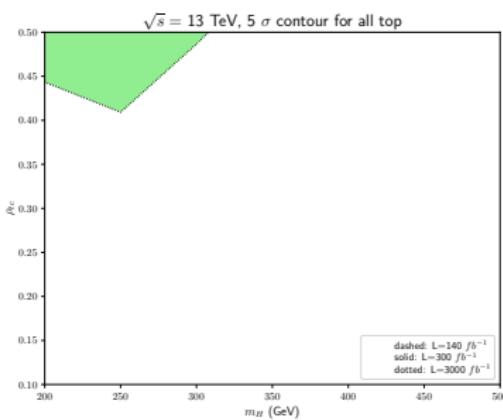


Figure: Discovery contours of $pp \rightarrow tH^0 \rightarrow ttc + X$ for All Tops (left) and Same Sign Tops (right).

3 σ Discovery Contour

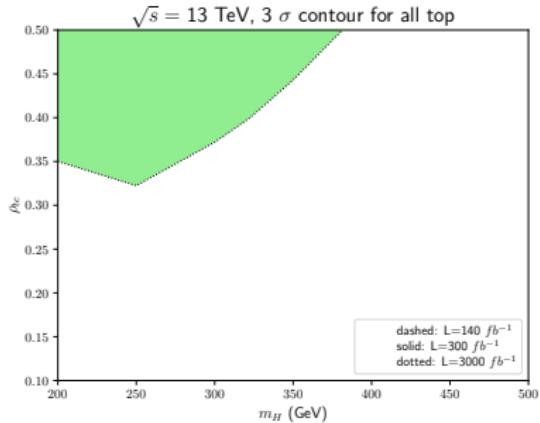
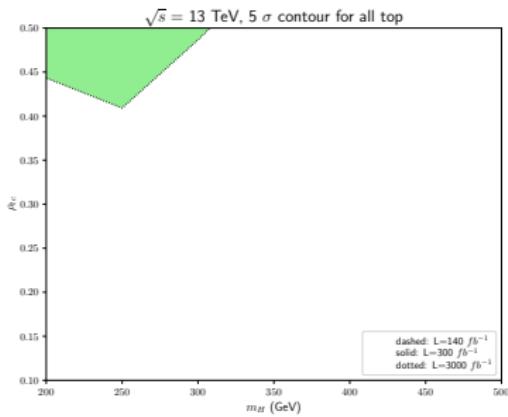
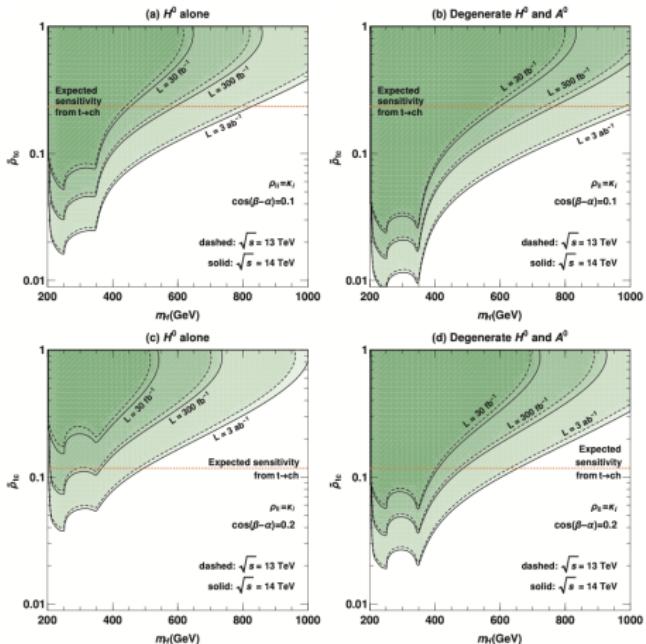


Figure: Discovery contours of $pp \rightarrow tH^0 \rightarrow ttc + X$ for all tops with 5 σ (left) and 3 σ (right).

Discovery Potential of $pp \rightarrow H^0 \rightarrow tc + X$



Flavor Changing Heavy Higgs Interactions at the LHC, B. Altunkaynak, W. S. Hou, C. Kao, M. Kohda and B. McCoy, Phys. Lett. B **751** (2015), 135–142.

Conclusions

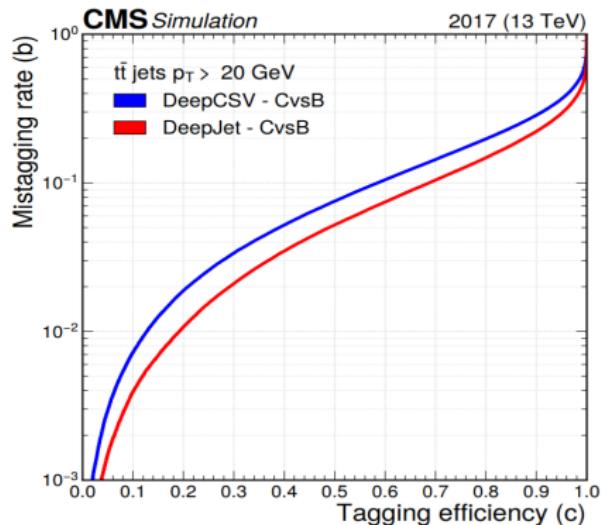
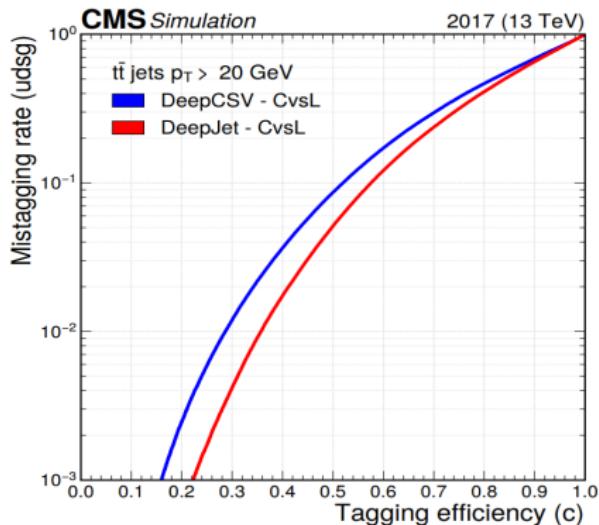
- It is of great interest to investigate the link between the most massive particle (top) and the mass giver (Higgs).
- In the decoupling limit, $\cos(\beta - \alpha) \rightarrow 0$ and $\sin(\beta - \alpha) \approx 1$.
- It is a win-win strategy to search for the FCNH top decay $t \rightarrow ch^0$ and the heavy Higgs decay $H^0, A^0 \rightarrow tc$ since $\lambda_{tch} \propto \tilde{\rho}_{tc} \cos(\beta - \alpha)$ and $\lambda_{Htc} \propto \tilde{\rho}_{tc} \sin(\beta - \alpha)$.
- Same sign top final state is almost background free. Combining with all top final state can help us improve the selection strategy and enhance the discovery potential for $pp \rightarrow tH^0 \rightarrow ttc + X$ at the LHC.
- The charged Higgs boson does not have CKM suppression for λ_{Hcb} . Thus $cg \rightarrow bH^\pm \rightarrow bbc$ is very promising to search for new physics.

Collaborators

- *Enhanced Charged Higgs Signal at the LHC:*
with Chenyu Fang, George W.-S. Hou, and Med Krab
- *Charming Higgs Bosons with Top Quarks at the LHC:*
with Nicolas Bost, Chenyu Fang, and Phillip Gutierrez

Backup: C-Tagging

CMS c-tagging efficiency and missing tagging efficiencies ¹⁰.



¹⁰ A. Tumasyan *et al.* [CMS], JINST **17**, no.03, P03014 (2022)