

Bridging Experiments, Narrowing Uncertainties: When DUNE Meets Hyper-K to unveil insights into 2-3 Oscillation Sector

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Based on [JHEP 10 \(2024\) 243](#)

July 08, 2025

Contribution ID:- 491

EPS-HEP 2025 - Marseille, France



Contribution ID:- 491 EPS-HEP 2025 - Ma

Outline

① Introduction and motivation

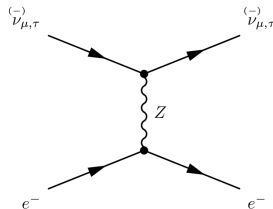
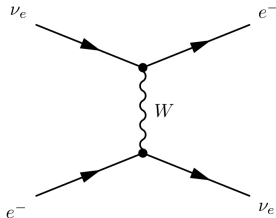
- ▶ Status of current oscillation parameters
- ▶ Deviation of θ_{23} from maximal mixing
- ▶ Role of appearance ($\nu_\mu \rightarrow \nu_e$) and disappearance ($\nu_\mu \rightarrow \nu_\mu$) oscillation channels in probing deviation from maximal mixing

② Regarding this work

- ▶ Sensitivity of DUNE and Hyper-K in determining deviation from maximal θ_{23} with variable exposure
- ▶ Precision measurements in atmospheric oscillation parameters - θ_{23} and Δm_{31}^2
- ▶ Wrong θ_{23} octant exclusion
- ▶ Allowed regions in $(\sin^2 \theta_{23} - \Delta m_{31}^2)$ and $(\sin^2 \theta_{23} - \delta_{CP})$ plane.

③ Summary and Conclusions.

Matter Effect

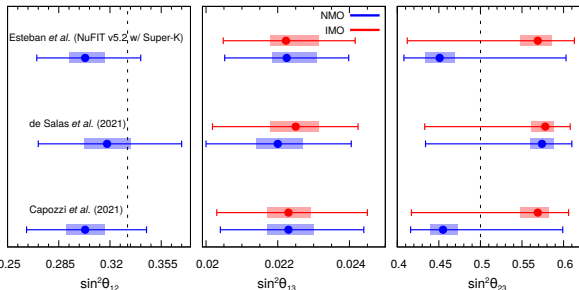


- Neutrinos interact with matter by coherent forward elastic scattering.
- Charge current interaction of ν_e with electrons creates an extra effective matter term for ν_e , i.e.,

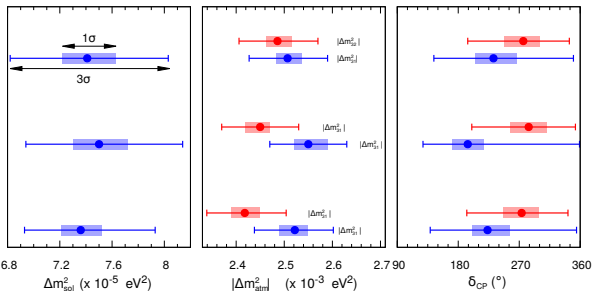
$$A = 2\sqrt{2}G_F N_e E.$$
- Matter term changes sign when we switch from neutrino to anti-neutrino mode.
- Matter term modifies oscillation probability differently depending on the sign of Δm^2 .
- The Hamiltonian corresponding to interaction with matter via CC-interaction is,

$$H = U \left[\frac{1}{2E} \text{diag}(m_1^2, m_2^2, m_3^2) \right] U^\dagger + \text{diag}(V_{CC}, 0, 0)$$

Present global-fit scenario in 3ν -paradigm



- 3σ (1σ) range of ν oscillation parameters, Esteban *et al.* www.nu-fit.org, de Salas *et al.* [arXiv: 2006.11237](https://arxiv.org/abs/2006.11237), and Capozzi *et al.* [arXiv: 2107.00532](https://arxiv.org/abs/2107.00532) in NMO and IMO.



Parameters	Relative 1σ error
$\sin^2 \theta_{12}$	4.5%
Δm_{21}^2	2.3%
$\sin^2 \theta_{23}$	6.7%
Δm_{31}^2	1.1%
$\sin^2 \theta_{13}$	3%
δ_{CP}	16%

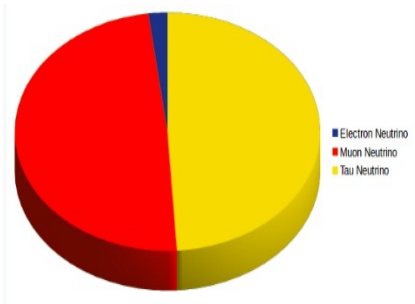
- The two most uncertain parameters are θ_{23} and δ_{CP} .
- θ_{23} is in LO (HO) for NMO (IMO) by Esteban *et al.* and Capozzi *et al.*

Deviation of θ_{23} from maximal mixing

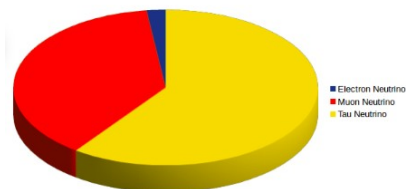
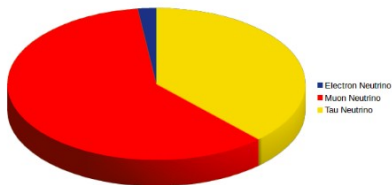
- $\mu \rightarrow \tau$ symmetry

$$\begin{aligned} |\nu_2\rangle &= \cos \theta_{23} |\nu_\mu\rangle + \sin \theta_{23} |\nu_\tau\rangle \\ |\nu_3\rangle &= -\sin \theta_{23} |\nu_\mu\rangle + \cos \theta_{23} |\nu_\tau\rangle \end{aligned}$$

- If $\theta_{23} = 45^\circ$, i.e for MM, ν_2 and ν_3 have equal contributions of ν_μ and ν_τ .



Deviation of θ_{23} from maximal mixing



ν Mixing Model	θ_{23}	θ_{13}	θ_{12}
Tri-bimaximal	45°	0°	35°
Bi-maximal	45°	0°	45°
Tri-bimaximal Cabibbo	45°	8.54°	35°
Bi-large	39°	12.12°	39°
Bi-trimaximal	36.23°	12.18°	36.23°

- Deviation from maximal mixing of θ_{23} indicates the exclusion of several theoretical neutrino mixing models.

Considered values of neutrino oscillation parameters in our work

Parameters	Best-fit	1σ range	3σ range	
$\Delta m_{21}^2/10^{-5} \text{ eV}^2$	7.36	7.21-7.52	6.93-7.93	Fixed
$\sin^2 \theta_{12}/10^{-1}$	3.03	2.90-3.16	2.63-3.45	Fixed
$\sin^2 \theta_{13}/10^{-2}$	2.23	2.17-2.30	2.04-2.44	Fixed
$\sin^2 \theta_{23}/10^{-1}$	4.55	4.40-4.73	4.16-5.99	Free
$\Delta m_{31}^2/10^{-3} \text{ eV}^2$	2.522	2.490-2.545	2.436-2.605	Free
$\delta_{\text{CP}}/^\circ$	223	200-256	139-355	Free

Roles of different channels in our study

- The appearance probability for $\nu_\mu \rightarrow \nu_e$ channel

$$P(\nu_\mu \rightarrow \nu_e) \approx \sin^2 \theta_{23} \sin^2(2\theta_{13}) \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2}$$

$$- \alpha \sin(2\theta_{13}) \zeta \sin \delta_{CP} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2}$$

$$+ \alpha \sin(2\theta_{13}) \zeta \cos \delta_{CP} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2}$$

- The disappearance probability for $\nu_\mu \rightarrow \nu_\mu$ channel

$$P(\nu_\mu \rightarrow \nu_\mu) \approx 1 - \sin^2(2\theta_{23}) \sin^2(\Delta)$$

$$+ (\alpha \Delta) c_{12}^2 \sin^2(2\theta_{23}) \sin(2\Delta)$$

$$- 2\alpha \zeta \cos(\delta_{CP}) \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(\hat{A} - 1)\Delta]}{(\hat{A} - 1)}$$

$$+ \frac{2}{(\hat{A} - 1)} \alpha \zeta \cos(2\theta_{23}) \cos(\delta_{CP}) \sin(\Delta) [\hat{A} \sin(\Delta) - \frac{\sin(\hat{A}\Delta)}{\hat{A}} \cos((\hat{A} - 1)\Delta)]$$

$$\alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \approx 0.033$$

$$\Delta = \Delta m_{31}^2 L / 4E$$

$$\zeta = \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{23}$$

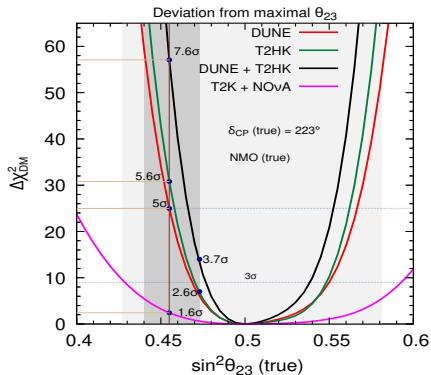
Appearance channel helps in θ_{23} octant exclusion.
Disappearance channel helps in the precision of θ_{23} .

$$\hat{A} = \pm \frac{2\sqrt{2}G_F N_e E}{\Delta m_{31}^2}$$

Salient features of DUNE and Hyper-K

Features	DUNE	Hyper-K
Baseline length	1300 km (Larger matter effect)	295 km (Smaller matter effect)
Detector Mass	40 kt (Smaller statistics)	187 kt (Larger statistics)
Detection technique	LArTPC	Water Cherenkov
Beam type	Wide-band, on-axis	Narrow-band, off-axis (2.5°)
Beam Power	1.2 MW	1.3 MW
Run time	5 yrs ν + 5 yrs $\bar{\nu}$	2.5 yrs ν + 7.5 yrs $\bar{\nu}$
P.O.T/year	1.1×10^{21}	2.7×10^{21}
Syst. Uncertainty in App. (Disapp.) channel	2% (5%)	5% (3%)

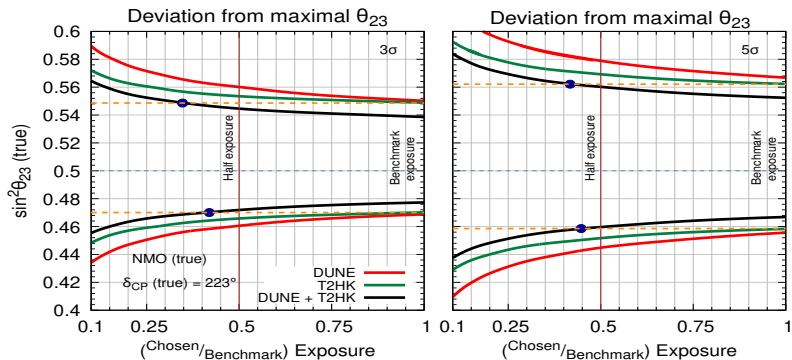
Deviation from maximal θ_{23}



$$\Delta\chi^2_{\text{DM}} = \min_{\delta_{\text{CP}}, \Delta m_{31}^2} \left\{ \chi^2 \left(\sin^2 \theta_{23}^{\text{test}} = 0.5 \right) - \chi^2 \left(\sin^2 \theta_{23}^{\text{true}} \in [0.4, 0.6] \right) \right\}$$

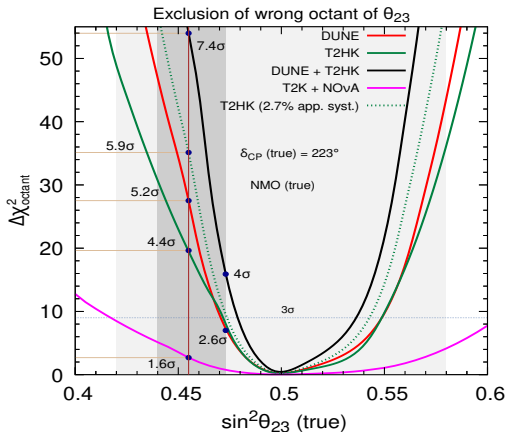
- In Nature, if true $\sin^2 \theta_{23}$ attains the lower value of the current 1σ uncertainty (0.473), only DUNE+Hyper-K can achieve 3σ sensitivity of non-maximal θ_{23} with the present benchmark values.

Deviation from maximal θ_{23} as the function of exposure



- The range of true values of $\sin^2 \theta_{23}$ that can be differentiated from MM choices, by DUNE + Hyper-K with just ~ 0.4 of their nominal exposures, cannot be achieved by either of them individually even at their respective full exposures.
- At lower exposure, Hyper-K always performs better than DUNE irrespective of the values of θ_{23} . At nominal exposure, they perform almost in same way.

Potential of exclusion of the wrong octant of θ_{23}



- At lower confidence, Hyper-K wins due to larger statistics whereas, at higher confidence DUNE wins due to lesser systematics in appearance channel.
- The combined setup of DUNE and Hyper-K boosts their individual performances to exclude the wrong octant solution.

Definition of $\Delta\chi^2_{\text{octant}}$

- For $\sin^2 \theta_{23} (\text{true}) < 0.5$ (LO),

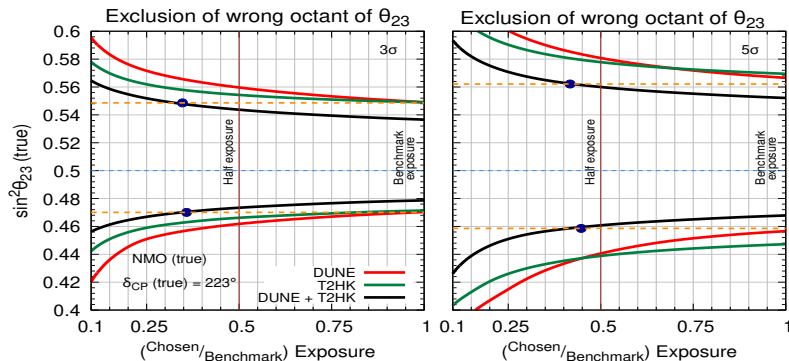
$$\Delta\chi^2_{\text{octant}} = \min_{(\vec{\lambda})} \{ \chi^2 (\sin^2 \theta_{23}^{\text{true}} = [0.4, 0.5]) - \chi^2 (\sin^2 \theta_{23}^{\text{test}} = (0.5, 0.6]) \}$$

- For $\sin^2 \theta_{23} (\text{true}) > 0.5$ (HO),

$$\Delta\chi^2_{\text{octant}} = \min_{(\vec{\lambda})} \{ \chi^2 (\sin^2 \theta_{23}^{\text{true}} = (0.5, 0.6]) - \chi^2 (\sin^2 \theta_{23}^{\text{test}} = [0.4, 0.5]) \}$$

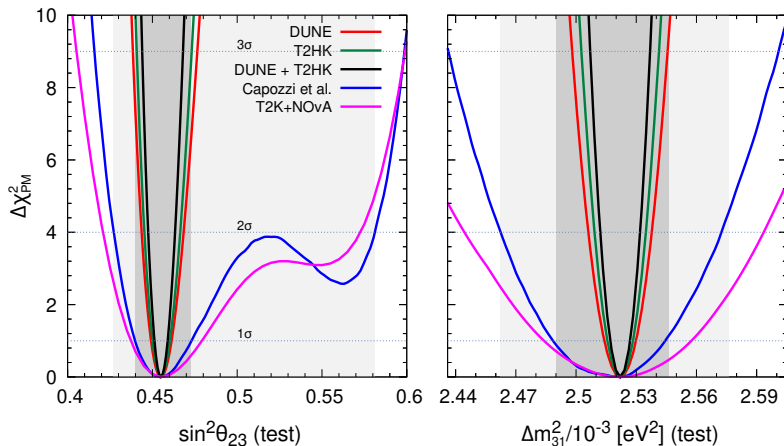
where, $\lambda = \delta_{\text{CP}}$, Δm_{31}^2 is the marginalized parameters.

Efficacy of DUNE and Hyper-K in octant exclusion at various exposures



- With just 0.25 times of the benchmark exposure of the individual experiments, the combined set up can exclude the wrong octant for more than half of the currently allowed $\sin^2 \theta_{23}$.

Precision measurement of θ_{23} and Δm_{31}^2



- The combination of DUNE and Hyper-K outperforms their performances in isolation to the precision measurement of θ_{23} and Δm_{31}^2 .

Definition of Relative 1σ precision and $\Delta\chi_{PM}^2$

The relative 1σ precision in the measurement of oscillation parameters ζ is estimated as follows:

$$p(\zeta) = \frac{\zeta^{\max} - \zeta^{\min}}{6.0 \times \zeta^{\text{true}}} \times 100\%.$$

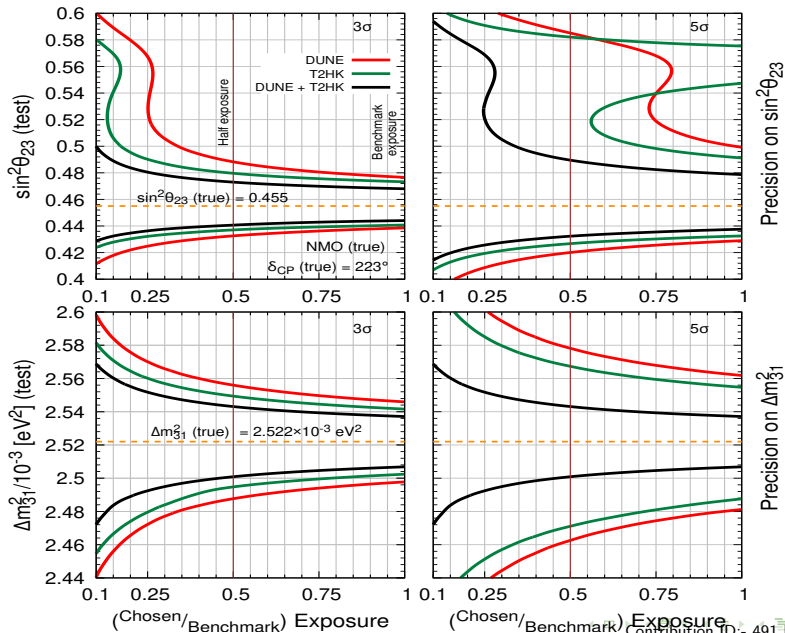
ζ^{\max} and ζ^{\min} are the allowed 3σ upper and lower bounds, respectively.

$$\Delta\chi_{PM, \sin^2 \theta_{23}}^2 = \min_{(\delta_{CP}, \Delta m_{31}^2)} \left\{ \chi^2 \left(\sin^2 \theta_{23}^{\text{true}} \in [0.4, 0.6] \right) - \chi^2 \left(\sin^2 \theta_{23}^{\text{test}} = [0.455] \right) \right\},$$

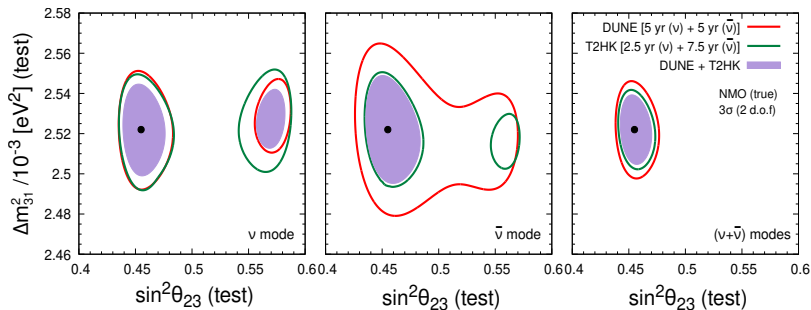
$$\Delta\chi_{PM, \Delta m_{31}^2}^2 = \min_{(\delta_{CP}, \sin^2 \theta_{23})} \left\{ \chi^2 \left(\sin^2 \theta_{23}^{\text{true}} \in [2.4, 2.6] \times 10^{-3} \right) - \chi^2 \left(\sin^2 \theta_{23}^{\text{test}} = 2.522 \times 10^{-3} \right) \right\},$$

Parameter	Relative 1σ precision (%)					
	HK	DUNE	HK+DUNE	T2K+NO ν A	Capozzi <i>et al.</i>	JUNO
$\sin^2 \theta_{23}$	1.18	1.40	0.88	7.10	6.72	—
Δm_{31}^2	0.25	0.31	0.20	0.99	1.09	0.2

Performance of measuring precision at various exposures

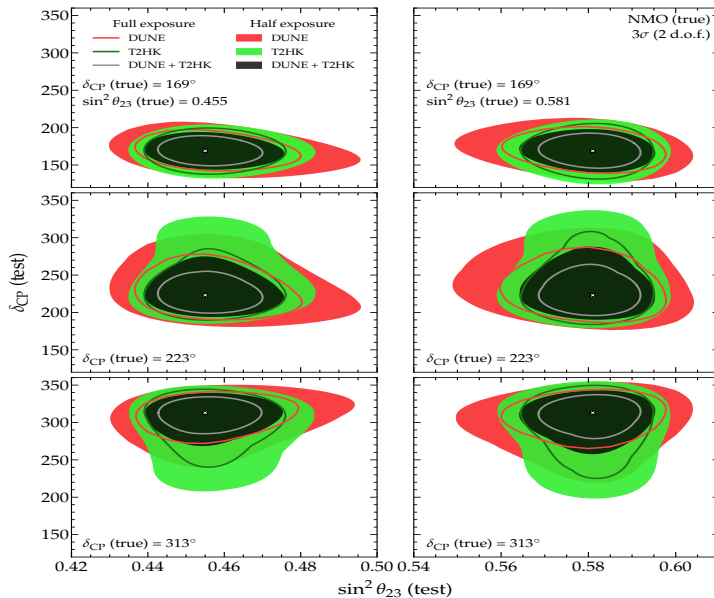


Allowed regions in $\sin^2 \theta_{23} - \Delta m_{31}^2$ plane

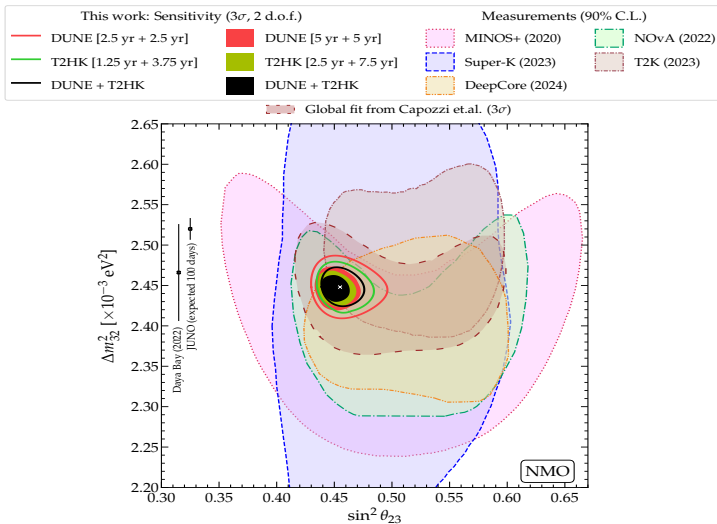


- The combination of DUNE and Hyper-K can exclude the HO only in antineutrino mode at 3σ C.L. breaking $\sin^2 \theta_{23} - \delta_{\text{CP}}$ degeneracy due to higher $\bar{\nu}$ statistics in Hyper-K. So, majority of the appearance events are free from fake (matter-induced) CP-phase.
- HO can be ruled out when both ν and $\bar{\nu}$ modes are considered together.

Allowed regions in $\sin^2 \theta_{23} - \delta_{\text{CP}}$ plane



Allowed regions in $\sin^2 \theta_{23} - \Delta m_{32}^2$ plane given by the other experiments

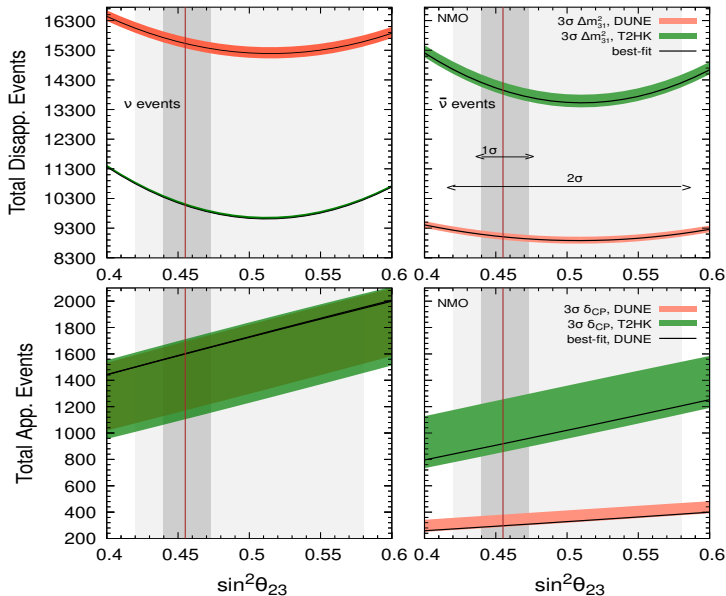


Summary

- Ongoing long-baseline and atmospheric experiments (e.g.- T2K, NO ν A, MINOS+, Super-K etc.) strongly suggest deviation from MM of θ_{23} .
- DUNE has large matter effect so is expected to measure Δm_{31}^2 precisely. But the larger matter effect induces fake CP-asymmetry which is negligible in Hyper-K.
- The disappearance statistics of Hyper-K is larger. So, Hyper-K is expected to measure θ_{23} precisely. But the combined setup improves the present achievable precision of $\sin^2 \theta_{23}$ and Δm_{31}^2 by a factor of 7 and 5, respectively.
- Combination of DUNE and Hyper-K outperforms their isolated performances to establish non-maximal θ_{23} . Furthermore, the range of true values of $\sin^2 \theta_{23}$ that can be differentiated from MM choices, by DUNE + Hyper-K with just ~ 0.4 of their nominal exposures, cannot be achieved by either of them individually even at their respective full exposures.
- With only 0.25 times of the benchmark exposure of the standalone experiments, the combination of DUNE and Hyper-K can exclude the wrong octant of θ_{23} for more than half of the currently allowed $\sin^2 \theta_{23}$.
- DUNE+Hyper-K can exclude the wrong octant solution with only antineutrino mode due to the complementarity between DUNE and Hyper-K.

Thank
You!

Total event rates in DUNE and Hyper-K



Total number of events in DUNE and Hyper-K

Channel		LO ($\sin^2 \theta_{23} = 0.455$)		MM ($\sin^2 \theta_{23} = 0.5$)		HO ($\sin^2 \theta_{23} = 0.599$)	
App.	ν	DUNE	Hyper-K	DUNE	Hyper-K	DUNE	Hyper-K
	$\bar{\nu}$	1601 [1586]	1598 [1588]	1729 [1712]	1725 [1713]	2004 [1983]	1996 [1981]
Disapp.	ν	297 [187]	919 [755]	328 [209]	1021 [844]	399 [260]	1251 [1044]
	$\bar{\nu}$	15529 [14286]	10064 [9487]	15209 [13974]	9628 [9057]	15857 [14597]	10661 [10074]
Disapp.	ν	9008 [4433]	13949 [8985]	8884 [4333]	13541 [8643]	9252 [4648]	14613 [9553]
	$\bar{\nu}$						

Table: Total (Signal + Background) appearance and disappearance event rates in DUNE and Hyper-K assuming 480 kt·MW·years and 2431 kt·MW·years of exposure, respectively. Events in parenthesis does not include the effect of wrong-sign contamination. The events are simulated by General Long Baseline Experiment Simulator (GLoBES).

- Contribution of wrong sign events is more in $\bar{\nu}$ mode than ν due to the cross-section suppression.
- Initially pions or kaons are produced due to pp or pn collision. Positive charged mesons are abundant than the negative one. Hence, contamination of ν in $\bar{\nu}$ beam is more.

Definition of χ^2 used

$$\chi^2 = \min_{(\vec{\zeta}_s, \vec{\zeta}_b)} \left\{ 2 \sum_{i=1}^n (\tilde{y}_i - x_i - x_i \ln \frac{\tilde{y}_i}{x_i}) + \zeta_s^2 + \zeta_b^2 \right\},$$

where, n is the total number of bins and

$$\tilde{y}_i(\{\omega, \alpha_{e\mu}\}, \{\zeta_s, \zeta_b\}) = N_i^{th}(\{\omega, \alpha_{e\mu}\})[1 + \pi^s \zeta_s] + N_i^b(\{\omega, \alpha_{e\mu}\})[1 + \pi^b \zeta_b]$$

where,

- $N_i^{th}(\{\omega, \alpha_{e\mu}\})$ = Predicted no. of events in i -th bin for a set of osc. params. ω and for a given value of $\alpha_{e\mu}$
- $N_i^b(\{\omega, \alpha_{e\mu}\})$ = No. of background events in the i -th bin where CC background depends on ω and $\alpha_{e\mu}$ but NC does not
- π^s, π^b = Systematic errors in signal and background
- ζ_s, ζ_b = 'Pulls' due to systematic errors in signal and background respectively
- $x_i = N_i^{ex} + N_i^b$ (where, N_i^{ex} = No. of observed CC signal events in the i -th bin, N_i^b = Same for the background)

<https://arxiv.org/pdf/1509.03517.pdf>

Contribution ID:- 491 EPS-HEP 2025 - M

Effect of Systematics of DUNE in probing non-maximal θ_{23}

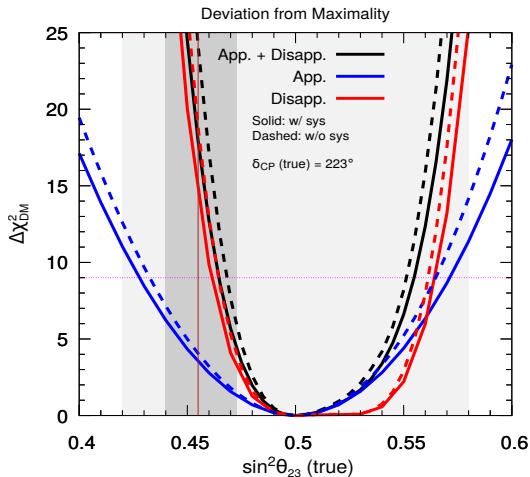


Table of systematics in DUNE

True $\sin^2 \theta_{23}$	Channels	2%, 5%	0%, 0%	5%, 5%	5%, 10%	10%, 10%
0.455 (Best-fit)	App.+Disapp. App. Disapp.	17.64 3.52 14.31	24.13 4.05 18.79	16.88 2.33 14.31	16.74 2.33 14.16	15.42 1.05 14.16
0.473 (1σ upper bound)	App.+Disapp. App. Disapp.	4.28 1.27 2.99	5.72 1.47 3.88	3.88 0.84 2.99	3.84 0.84 2.97	3.42 0.38 2.97

Neutrino mixing in three-flavor oscillation

- Flavor and mass eigen-states are linearly combined as

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad (1)$$

- The three flavor neutrino oscillation is given by

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{j < i} \text{Re}(U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j}) \sin^2(1.27 \Delta m_{ij}^2 L/E) \\ - 2 \sum_{j < i} \text{Im}(U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j}) \sin(2.54 \Delta m_{ij}^2 L/E)$$

where, $\Delta m_{ij}^2 = m_i^2 - m_j^2$ (in eV^2), L is the baseline length (in km), and E is the energy of the neutrino (in GeV).

Parametrization of PMNS matrix

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

- PMNS (Pontecorvo-Maki-Nakagawa-Sakata) matrix is parametrized by three independent mixing angles (θ_{23} , θ_{13} , and θ_{12}), two independent mass-squared difference [Solar mass-squared difference (Δm_{21}^2), Atmospheric mass-squared difference (Δm_{31}^2)] and one δ_{CP} phase.
- For the non-zero value of θ_{13} , we have got 3ν -paradigm in neutrino oscillation.
- The (3×3) matrices in the red, green, and blue color are called "1-2 sector" or Solar Sector, "1-3 sector" or Reactor Sector, and "2-3 sector" or Atmospheric Sector.