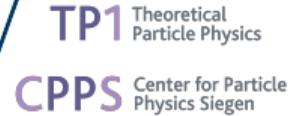
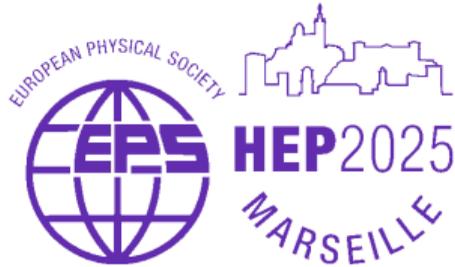


Symmetric Mass Generation

Anna Hasenfratz and **Oliver Witzel**
(for the Lattice Strong Dynamics collaboration)



EPS-HEP 2025
Marseille, France · July 08, 2025



Two paradigms to generate mass:

Chiral Symmetry Breaking (χ SM)

- ▶ **UV limit:** N_f massless flavors gauged with $SU(N_c)$: $SU(N_f)_V \times SU(N_f)_A \times U(1)_V \times U(1)_A$
- ▶ **IR limit:** Spontaneous symmetry breaking
 - Fermions are confined
 - Chiral symmetry breaks
 - Massless Goldstone bosons
 - Bilinear condensate $\langle \psi \bar{\psi} \rangle \neq 0$
 - Non-Goldstone states are gapped
- ▶ If a system has 't Hooft anomalies, massless Goldstones are necessary to satisfy the anomaly matching condition

Symmetric Mass Generation (SMG)

- ▶ Generate mass without symmetry breaking
- ▶ **IR limit:** Symmetric mass generation
 - Fermions are confined
 - No symmetry breaking
 - No massless Goldstones bosons
 - Bilinear condensate $\langle \psi \bar{\psi} \rangle = 0$
 - Expect $\langle \psi_1 \bar{\psi}_1 \psi_2 \bar{\psi}_2 \rangle \neq 0$ and/or instantons
 - All bound states are gapped
- ▶ If all 't Hooft anomalies are canceled SMG is possible

Two paradigms to generate mass:

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Symmetric Mass Generation (SMG)

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▶ **UV limit:** N_f massless flavors gauged with $SU(N_c)$: $SU(N_f)_V \times SU(N_f) \times U(1)_A$

- ▶ **IR limit:** Symmetric
 - Fermions are
 - No symm
 - No ma
 - Bili
 - An
- ▶ Instantons
- ▶ If all 't Hooft anomalies are canceled
SMG is possible

Traditionally:
Confinement cannot happen w/o χ SB

Recently:
Consider generalized symmetries

[García-Etxebarria, Montero JHEP 08 (2019) 003]

How to cancel 't Hooft anomalies (qualitative example by Cenke Xu for illustration)

- ▶ Vector N_f massless flavors gauged with $SU(N_c)$: $SU(N_f)_V \times SU(N_f)_A \times U(1)_V \times U(1)_A$
- ▶ $U(1)_A$: $\psi \rightarrow e^{i\alpha\gamma_5}\psi$, $\bar{\psi} \rightarrow \bar{\psi}e^{i\alpha\gamma_5}$ is explicitly broken by the anomaly

$$Z \rightarrow \exp\{i\alpha N_c N_f Q\} Z, \quad Q = \int d^4x F\bar{F}$$

- ▶ Spin- Z_4 : $\psi_L \rightarrow i\psi_L$, $\psi_R \rightarrow -i\psi_R$, $(\phi \rightarrow -\phi, A_\mu \rightarrow A_\mu)$ is a symmetry of the action

→ The measure changes: $Z \rightarrow \exp\{i\frac{2\pi}{4} N_c N_f Q\} Z$

→ Not anomalous, if $N_f N_c Q \bmod 4 = 0$

- ▶ If the system breaks $U(1)_A \rightarrow$ Spin- Z_4 an SMG phase is possible

→ Breaking could be explicit (add 4-fermion interaction or lattice discretization) or spontaneous

Rigorous examples [García-Etxebarria, Montero JHEP 08 (2019) 003]

▶ **2-D**: Eight Weyl (massless Majorana) fermions cancel all anomalies (chiral and time reversal)

→ Add 4-fermion interaction \Rightarrow SMG phase is predicted [Fidkowski, Kitaev PRB 81 (2010) 134509]

[Review by Wang, You Symmetry 14 (2022) 7]

▶ **4-D**: 16 Weyl fermions or 8 massless Dirac fermions to cancel all 't Hooft anomalies

[You, BenTov, Xu, arXiv:1402.4151] [Butt et al. PRD 104 (2021) 094504]

→ Counting in the UV: $N_f \times N_c \bmod 8 = 0$

→ Possible 4-D candidate models:

- SU(2) gauge + $N_f = 4$ fundamental flavors
- SU(3) gauge + $N_f = 8$ fundamental flavors
- SU(4) gauge + $N_f = 2, 4, 6, \dots$ fundamental flavors

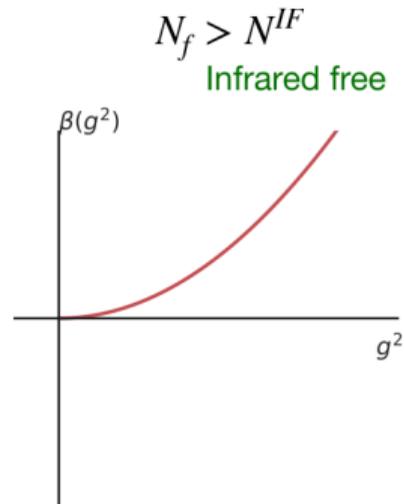
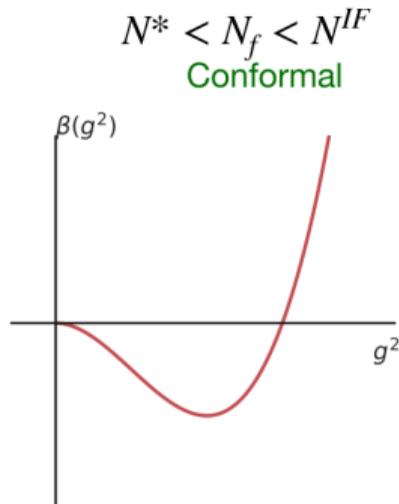
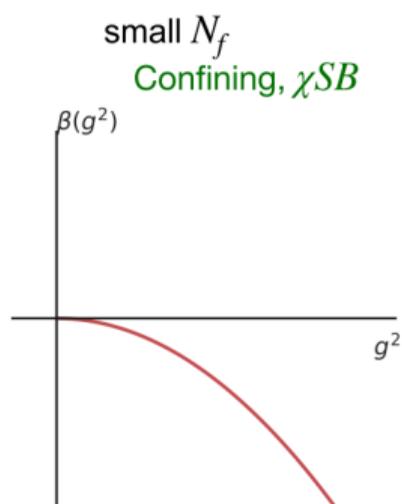
▶ Anomaly cancellation is **necessary but not sufficient** condition

→ Does SMG exist in any of these systems? Where?

→ Does it exist in the continuum limit at a strongly coupled fixed point?

Phases of gauge-fermion systems

- ▶ $SU(N_c)$ gauge + N_f (fundamental) massless flavors



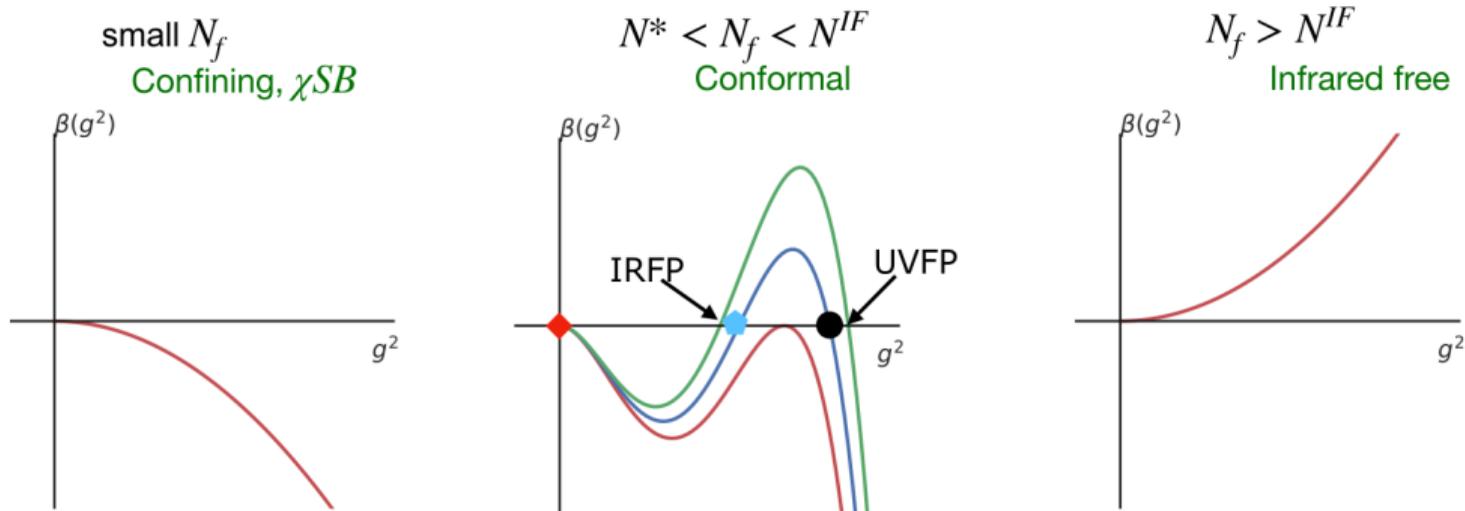
- ▶ RG β function $\beta(g^2) = -\mu^2 \frac{dg^2}{d\mu^2}$ characterizes the phases

→ Known up to 5-loop in \overline{MS} [Baikov, Chetyrkin, Kühn PRL118(2017)082002]

- ▶ **Perturbatively** emerges the IR fixed point at $g_0^2 = \infty$ for $N_f = N^*$ and moves to $g_0^2 = 0$ for $N_f \rightarrow N^{IF}$

Phases of gauge-fermion systems

- ▶ $SU(N_c)$ gauge + N_f (fundamental) massless flavors



- ▶ **Nonperturbatively** the IR fixed point could emerge at finite g_*^2 if $\beta(g^2) \sim (\delta - \delta_*) - (g - g_*)^2$

→ “Conformality lost” at IR-UV fixed point merger

[Kaplan et al. PRD80 (2009) 125005] [Vecchi PRD82 (2010) 045013] [Gorbenko et al. JHEP10 (2018) 108]

→ **UVFP**: phase transition, new phase, new relevant operator

Numerically investigated examples in 4-D

▶ Similar results for two systems using $m_f = 0$

→ SU(3) gauge + $N_f = 8$ [Hasenfratz PRD 106 (2022) 014513] [Hasenfratz, OW (LSD) PoS Lattice 2024 146]

→ SU(2) gauge + $N_f = 4$ [Butt et al. PRL 134 (2025) 031602]

▶ Phase structure

strong coupling

weak coupling



▶ SMG phase

▶ Chirally symmetric

▶ Confining, gapped spectrum

continuous

phase transition



continuum limit exist
RG β function

▶ Appears conformal

▶ Chirally symmetric

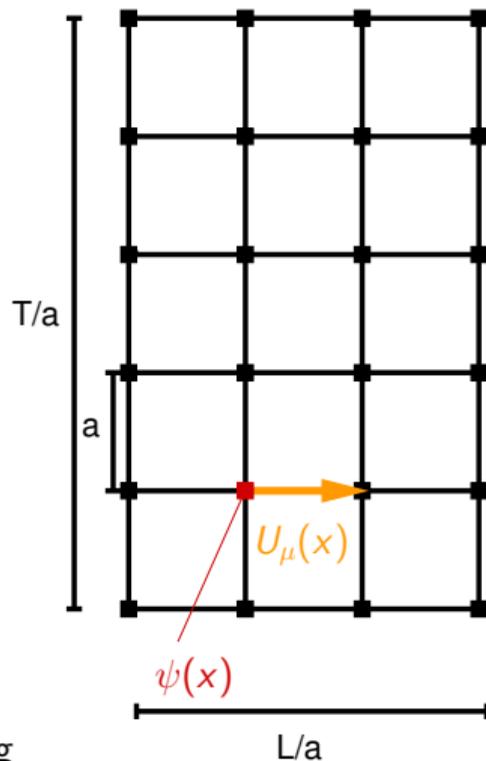
▶ Conformal hyperscaling

Numerical investigations: dynamical LFT simulations

- ▶ **Wick-rotate** to Euclidean time $t \rightarrow i\tau$
- ▶ **Discretize space-time** introducing the lattice spacing a
- ▶ **Restrict volume** to a hypercube of finite extent $(L/a)^3 \times T/a$
- ▶ Formulate problem in terms of **Feynman's path integral**

$$\langle \mathcal{O} \rangle_E = \frac{1}{Z} \int \mathcal{D}[\psi, \bar{\psi}] \mathcal{D}[U] \mathcal{O}[\psi, \bar{\psi}, U] e^{-S_E[\psi, \bar{\psi}, U]}$$

- ▶ The lattice regularizes the theory
 - **Finite lattice spacing a** : UV regulator
 - **Finite volume of length L/a** : IR regulator
- ▶ Stochastic interpretation of the path integral
 - ↪ **Markov chain Monte Carlo** simulations w/ importance sampling
 - Requires High Performance Computing (HPC)



Simulations

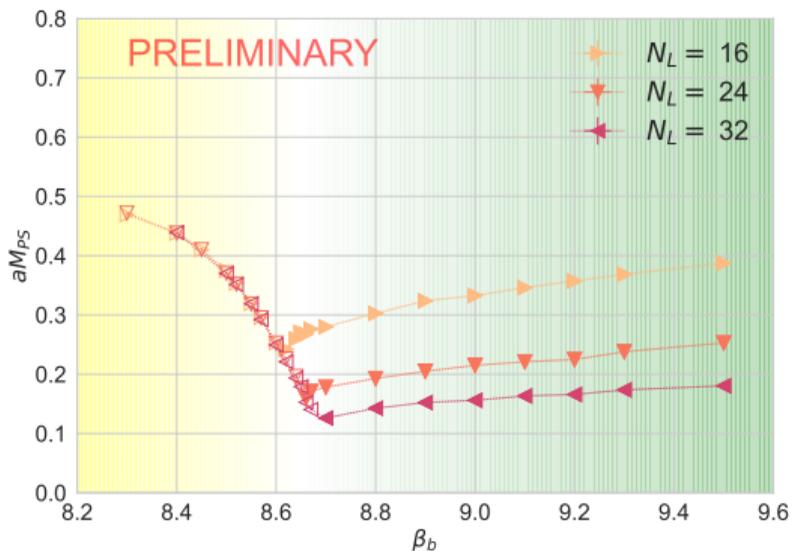
- ▶ nHYP-smearred staggered fermions \rightsquigarrow Dirac-Kähler fermions
- ▶ Additional Pauli-Villars fields
- ▶ Plaquette gauge action with adjoint term

- ▶ Volume $16^3 \times 32$, $24^3 \times 64$, $32^3 \times 64$, $48^3 \times 96$
- ▶ $\beta_b = 8.10, 8.20, 8.30, 8.40, 8.50, 8.60, 8.70, 8.80, 8.90, 9.10, 9.20, 9.30, 9.50$
 - Simulations cross from weak coupling to symmetric mass generation (SMG) phase
[Hasenfratz PRD 106 (2022) 014513]

- ▶ HMC: QEX (Osborn, Jin, Peterson)
- ▶ Spectrum: MILC (DeTar et al., . . . , Schaich, Hasenfratz)
- ▶ Gauge flow: QLUA (Pochinsky et al.)

Pseudoscalar correlator (would-be Goldstone pion)

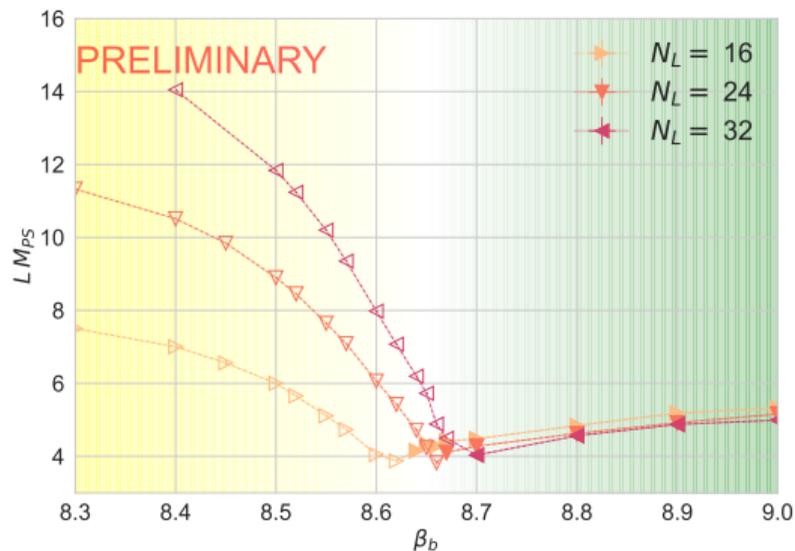
► aM_{PS}



volume independent
“gapped phase”

$1/L$ volume scaling
conformal phase

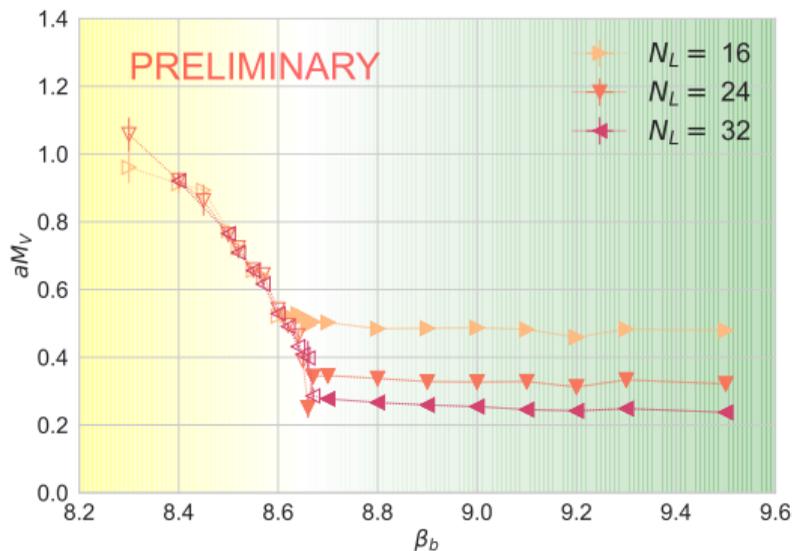
► LM_{PS}



► Conformal scaling for $\beta_b > 8.6$
 $M_{PS}L \approx \text{const}$

Vector correlator

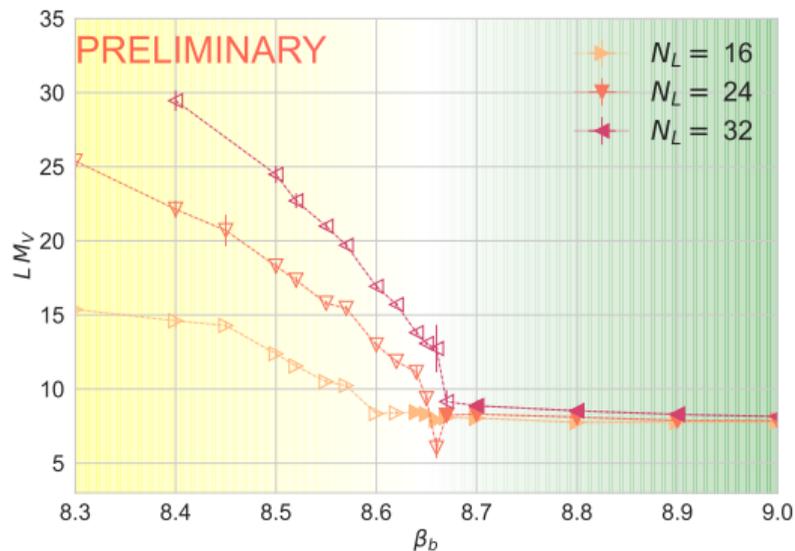
► aM_{VT}



volume independent
“gapped phase”

$1/L$ volume scaling
conformal phase

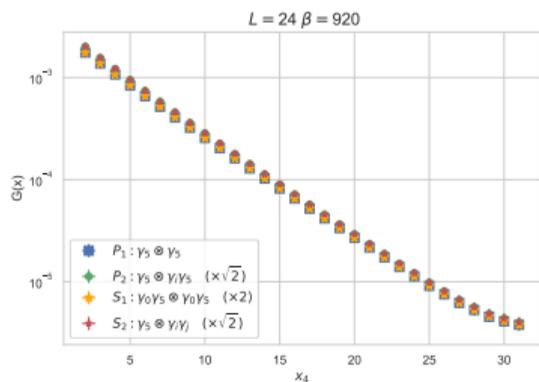
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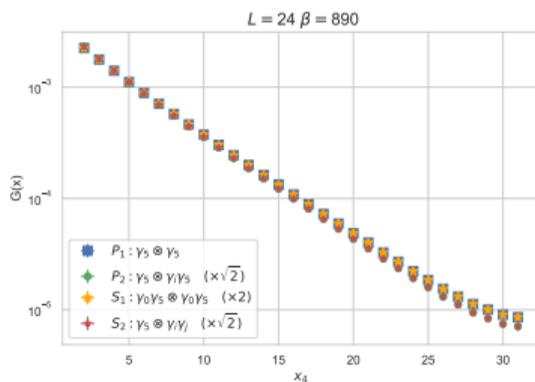
Pseudoscalar correlator (would-be Goldstone pion)

- ▶ Weak coupling
 $\beta_b = 9.20$ (conformal)



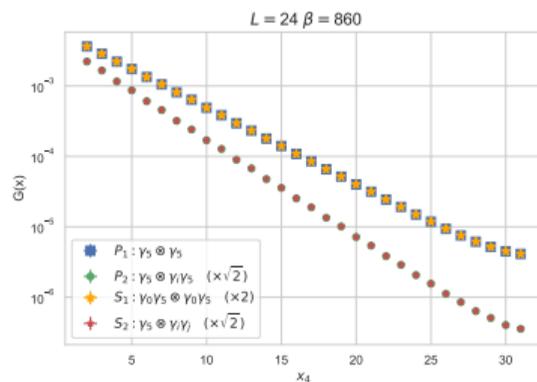
- ▶ Degenerate pseudoscalar and scalar (parity doubling)
- ▶ No taste splitting

- ▶ $\beta_b = 8.90$



- ▶ Degenerate pseudoscalar and scalar (parity doubling)

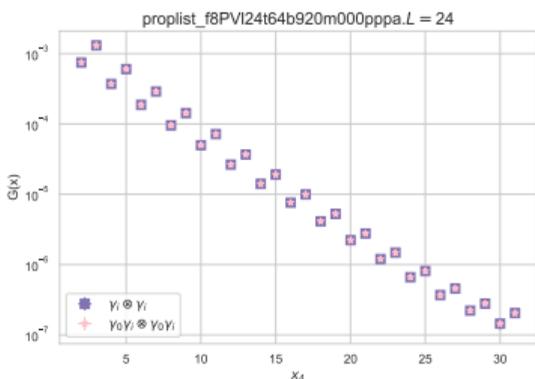
- ▶ Strong coupling
 $\beta_b = 8.60$ (SMG)



- ▶ Degenerate pseudoscalar and scalar (parity doubling)
- ▶ No chiral symmetry breaking

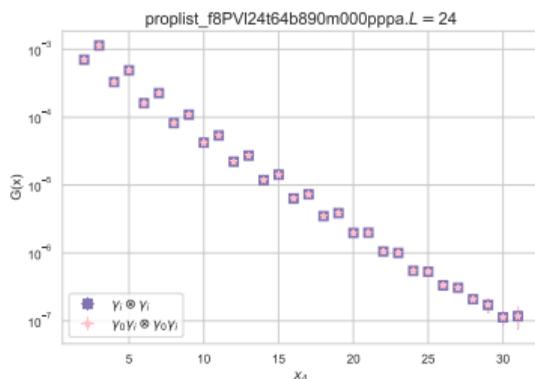
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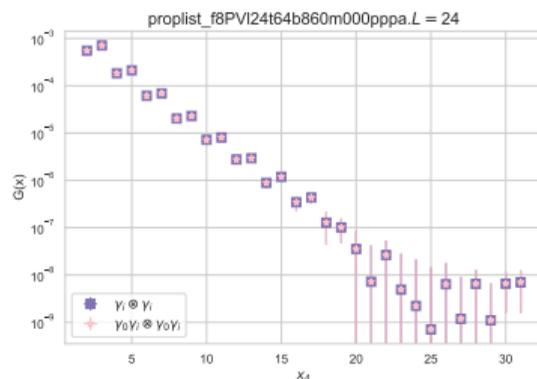
- ▶ Degenerate vector and axial (parity doubling)

- ▶ $\beta_b = 8.90$



- ▶ Degenerate vector and axial (parity doubling)

- ▶ Strong coupling
 $\beta_b = 8.60$ (SMG)

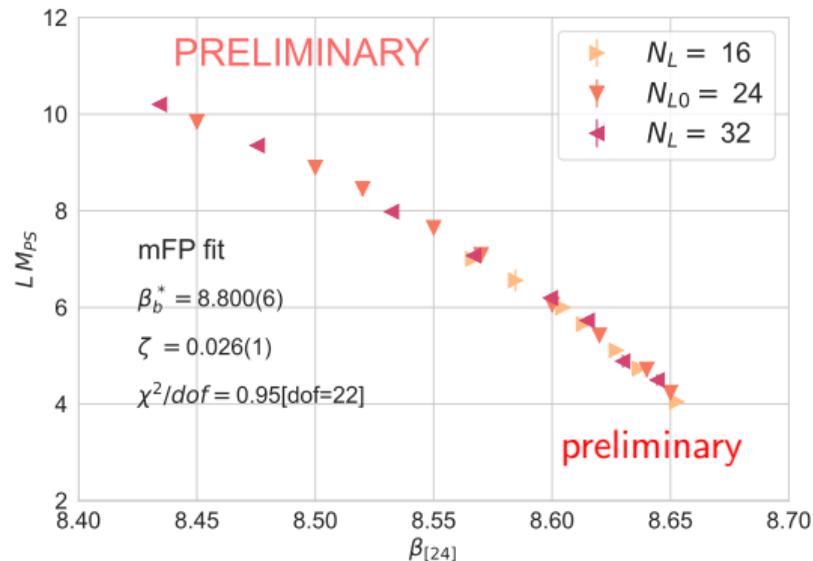
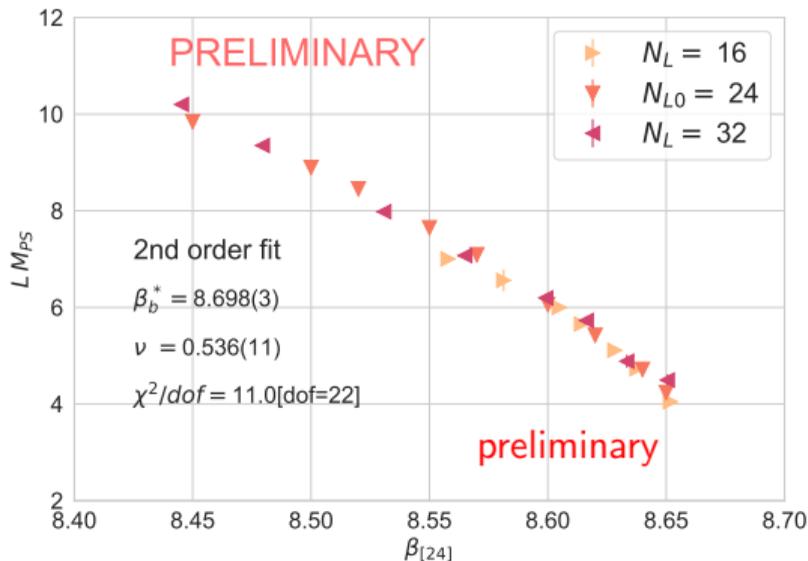


- ▶ Degenerate vector and axial (parity doubling)
- ▶ No chiral symmetry breaking

Order of the phase transition

- ▶ Finite size scaling/curve collapse analysis near the critical point $g \rightarrow g^*$
[Hasenfratz PRD 106 (2022) 014513]
 - \mathcal{O} is a dimensionless operator: $\mathcal{O}(g, L) = f(L/\xi)$
 - ξ correlation length at g
 - $f(x = L/\xi)$ unique curve, independent of L
- ▶ Second order scaling: $\xi \propto |g - g_*|^{-\nu}$
 - First order scaling is similar and has $\nu = 1/d = 0.25$
- ▶ Merged fixed point (BKT) scaling: $\xi \propto \exp\{\zeta/|g - g_*|\}$ if $\beta(g^2) \sim (g^2 - g_*^2)^2$
- ▶ Obtain exponents and g^* by standard curve-collapse analysis

Outlook: finite size scaling analysis



► Finite size scaling

$$2^{\text{nd}} \text{ order: } L \cdot m(\beta; L) = f_{2\text{nd}}^{(c)} \left((\beta/\beta_c - 1) L^{1/\nu} \right)$$

$$\text{BKT: } L \cdot m(\beta; L) = f_{\text{BKT}}^{(c)} \left(L \exp(-\zeta |\beta/\beta_c - 1|^{-\nu}) \right)$$

Summary

- ▶ SMG is possible in 4-D, if all 't Hooft anomalies are canceled
- ▶ Large scale simulations to scrutinize $SU(3)$ with $N_f = 8$ fundamental flavors at strong coupling using additional Pauli-Villars fields
 - Preliminary findings support picture of a conformal and an SMG phase
 - No chiral symmetry breaking
 - Indications for a new non-trivially interacting UVFP in 4-D
 - Implies: $SU(3) + N_f = 8$ fundamental flavors is likely the onset of the conformal window
 - Maybe: merged fixed point with marginal operator like QCD but not asymptotically free
 - Are staggered (Dirac-Kähler) fermions special?
- ▶ Ongoing $48^3 \times 96$ volume simulations



T14, Fri 9:50
Salle Joliette