

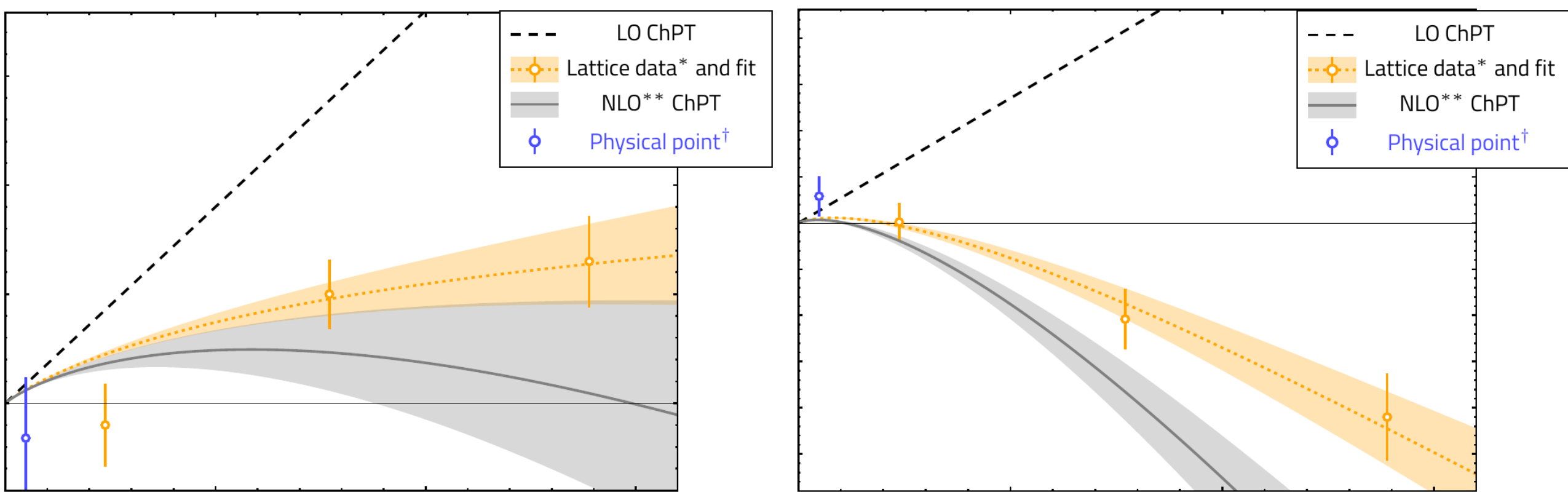
SIX-POINT CHPT AMPLITUDE FOR PIONS AND KAONS

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Abstract

We compute all one-loop six-point scattering amplitudes with external pions and kaons within the framework of three-flavor Chiral Perturbation Theory (ChPT). We use distinct pion, kaon and η meson masses, but maintain the isospin limit, thus neglecting mass differences between charged and neutral particles (as this is a QED-dominated effect). Rather than separate amplitudes for each channel, we write generic amplitudes using the flavor Lie algebra, generalizing the concept of stripped amplitudes.

The goal: generalize this



* Blanton, T. D. et al., *Interactions of two and three mesons including higher partial waves from lattice QCD*, JHEP **10** (2021) 023

** JB, MS et al., *The isospin-3 three-particle K-matrix at NLO in ChPT*, JHEP **05** (2023) 187

† Dawid, S. et al., *Two- and three-meson scattering amplitudes with physical quark masses from lattice QCD*, arXiv:2502.17976

2+1-flavor ChPT

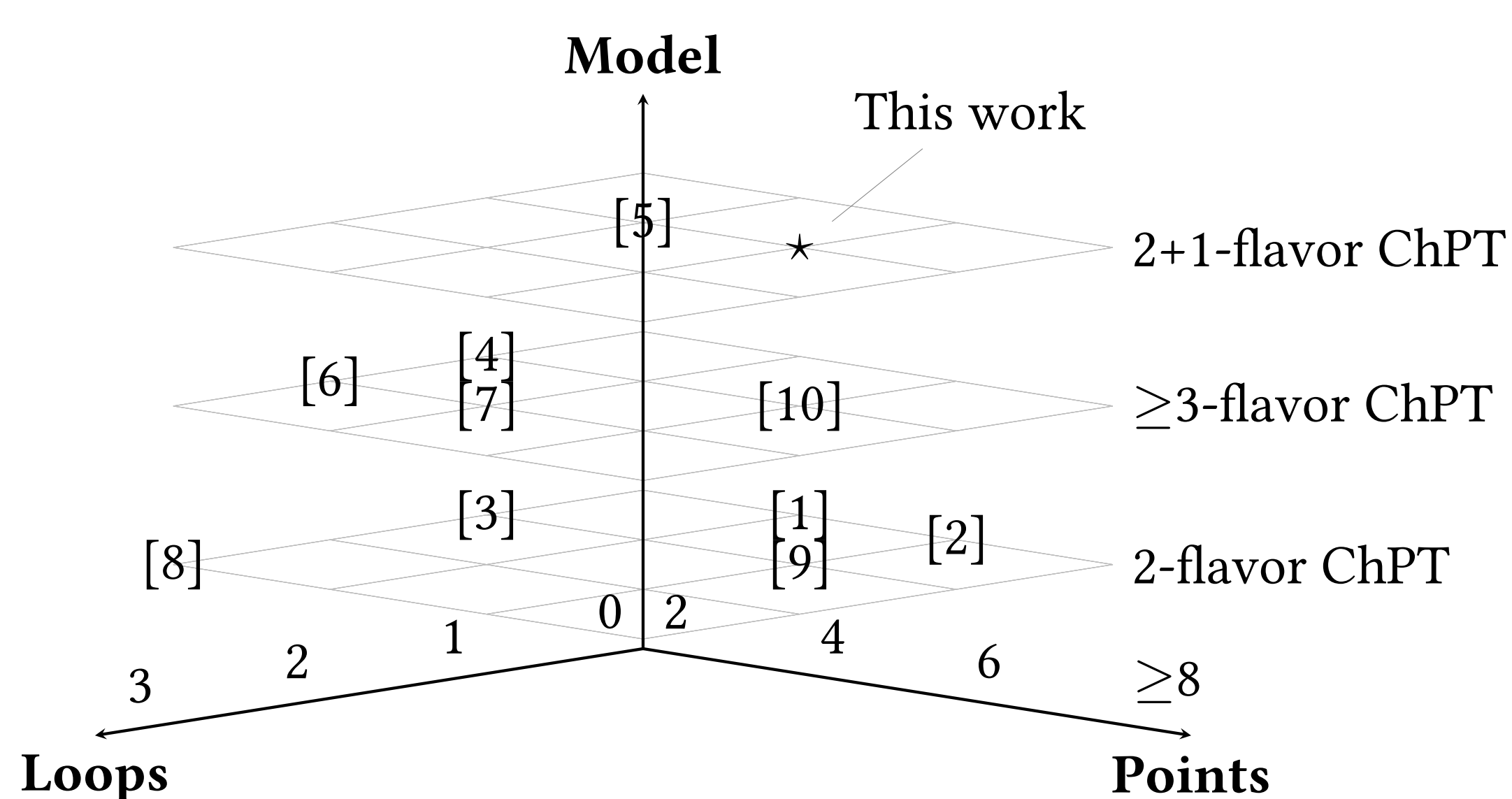
The low-energy effective field theory of QCD with $m_u = m_d < m_s$ and no heavier quarks

$SU(3)_R \times SU(3)_L$ chiral symmetry of massless QCD \rightarrow $\begin{cases} \text{spontaneously broken by quark condensate} & \Rightarrow \text{non-linear sigma model} \\ \text{explicitly broken by quark masses} & \Rightarrow \text{ChPT with meson mass terms} \end{cases}$

$$\mathcal{L}_{\text{ChPT}} = \langle \partial_\mu U \partial^\mu U^\dagger + (U + U^\dagger) \chi \rangle + (\text{counterterms})$$

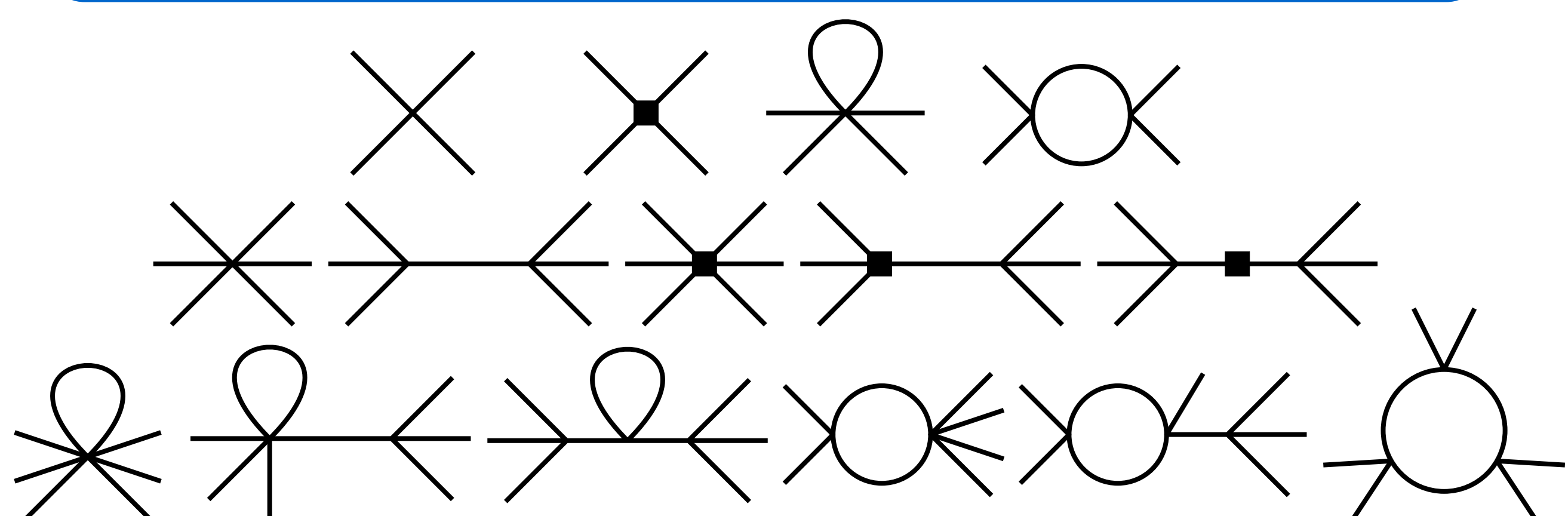
$$U \sim \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & K^0 \\ K^- & \bar{K}^0 & -2\frac{\eta}{\sqrt{6}} \end{pmatrix}, \quad \chi \sim \begin{pmatrix} m_u \\ m_d \\ m_s \end{pmatrix}$$

(incomplete) ChPT amplitude history

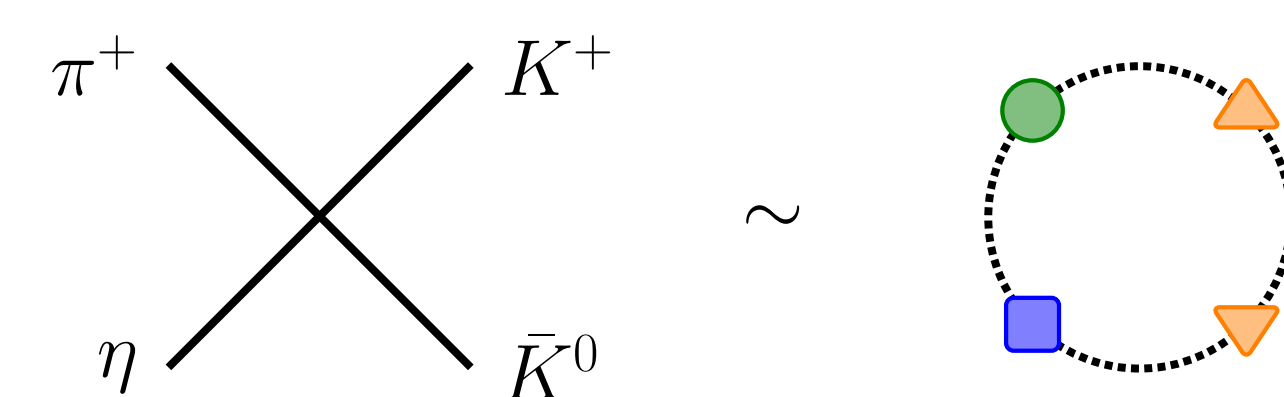


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Feynman diagrams



Strangeness-conscious Lie algebra



..... = index contraction
 = pion (any)
 = η meson
 = kaon (any)
 = kaon ($S = +1$)
 = kaon ($S = -1$)

Removal of η 's

$$\begin{aligned} \text{pion} - \eta &= \text{pion} - \eta = \frac{\text{pion}}{\sqrt{6}} \\ \eta - \text{pion} &= \frac{\text{pion}}{\sqrt{6}} \\ \text{kaon} - \eta &= \frac{\text{kaon}}{\sqrt{6}} \\ \eta - \text{kaon} &= -2 \frac{\text{kaon}}{\sqrt{6}} \\ \text{pion} - \eta &= \frac{\text{pion}}{3} - \frac{\eta}{\sqrt{6}} \end{aligned}$$

..... = 1,2,3,... pions = 1,3,5,... kaons

Strangeness rules

$$\begin{aligned} \text{pion} - \text{pion} &= \text{pion} - \text{pion} = \text{pion} - \text{pion} = \text{pion} - \text{pion} = 0 \\ \text{pion} - \text{kaon} &= \text{pion} - \text{kaon} = \text{pion} - \text{kaon} = \text{pion} - \text{kaon} = 0 \end{aligned}$$

Cayley–Hamilton theorem

$$\text{pion} - \text{pion} + (a \leftrightarrow b) = \text{pion} - \text{pion} = \delta_b^a \quad \text{pion} - \text{pion} = \epsilon^{abc}$$

The multi-mass propagator

N-flavor Fierz identity

$$\begin{aligned} \text{pion} - \text{pion} &= P(q^2) \left[\text{pion} - \text{pion} - \frac{1}{N} \text{pion} - \text{pion} \right] \\ P(q^2) &= \frac{1}{q^2 - M^2} \end{aligned}$$

2+1-flavor Fierz identity

$$\begin{aligned} \sum_a \text{pion} - \text{pion} &= \Pi_0(q^2) \left[\text{pion} - \text{pion} - \frac{1}{3} \text{pion} - \text{pion} \right] \\ &+ \Pi_1(q^2) \left[\text{pion} - \text{pion} + \text{pion} - \text{pion} - \frac{2}{3} (\text{pion} - \text{pion} + \text{pion} - \text{pion}) \right] \\ &+ \Pi_2(q^2) \left[\text{pion} - \text{pion} \right] \\ \Pi_0 &= \frac{P_\pi + 2P_K}{3} \quad \Pi_1 = \frac{2(P_\pi - P_K)}{\sqrt{6}} \quad \Pi_2 = P_\eta + \frac{P_\pi - 4P_K}{3} \end{aligned}$$

The triangle integral

$$\begin{aligned} \text{pion} - \text{pion} &= \int \frac{d^d l}{(2\pi)^d} \frac{(\text{numerator})}{[(l - q_1)^2 - M_a^2][(l + q_2)^2 - M_b^2][l^2 - M_c^2]} \end{aligned}$$

Numerator up to rank 3, master integral reduction not practical
 Large number of symmetries under simultaneous permutation of q_1, q_2, q_3 and M_a, M_b, M_c