

EUROPEAN PHYSICAL SOCIETY



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MARSEILLE



Lepton Flavour Universality tests and determination of $|V_{us}|$ using the HFLAV tau branching fractions fit

for the HFLAV Tau group,

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HFLAV Tau branching fraction fit

inputs		examples
171	measurements of τ branching fractions & branching ratios	$\mathcal{B}(\tau^- \rightarrow \pi^- \nu_\tau), \quad \frac{\mathcal{B}(\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau)}{\mathcal{B}(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)}$
1	nuisance fit parameter measurement (new feature)	$\mathcal{B}(a_1^- \rightarrow \pi^- \gamma) = 0.0021 \pm 0.0008$ [Phys.Rept. 421 (2005) 191]
91	constraints	$\mathcal{B}_{\tau \rightarrow \mu \nu \bar{\nu}} / \tau \rightarrow e \nu \bar{\nu} = \mathcal{B}_{\tau \rightarrow \mu \nu \bar{\nu}} / \mathcal{B}_{\tau \rightarrow e \nu \bar{\nu}}$ tau-decay final states related through different η, ω decay modes
1	uncertainty scale factor	5.44 scale factor for inconsistent BABAR/ Belle $\mathcal{B}(\tau^- \rightarrow K^- K^- K^+ \nu_\tau)$
external uncertain parameters: correct for updated values, account from induced uncertainties' correlations		

$$\chi^2 = \sum_{ijkl} (m_i - M_{ik} q_k) \left(V^{-1} \right)_{ij} (m_j - M_{jl} q_l) + \sum_r \frac{(n_r - p_r)^2}{\sigma_{n_r}^2}$$

137 fit parameters 1 nuisance fit parameter (new feature)
covariance matrix of fit parameters and nuisance fit parameter
 $\chi^2/\text{d.o.f.} = 138/125$ $P(\chi^2) = 20.2\%$
unitarity residual $\mathcal{B}_{\text{ur}} = 1 - \mathcal{B}_{\text{all}} = 0.0007 \pm 0.0011$

- ▶ HFLAV unitarity-constrained tau BR fit used for PDG tau BR “fit” averages
 - ▶ use of nuisance fit parameter removes some modifications of HFLAV tau BR fit for PDG “fit” averages

Lepton universality tests with tau compared with other measurements

2013 [A.Pich, Precision Tau Physics (2014)]

$ g_\mu/g_e $	$\Gamma_{\tau \rightarrow \mu}/\Gamma_{\tau \rightarrow e}$ 1.0018(14)	$\Gamma_{\pi \rightarrow \mu}/\Gamma_{\pi \rightarrow e}$ 1.0021(16)	$\Gamma_{K \rightarrow \mu}/\Gamma_{K \rightarrow e}$ 0.9978(20)	$\Gamma_{K \rightarrow \pi\mu}/\Gamma_{K \rightarrow \pi e}$ 1.0010(25)	$\Gamma_{W \rightarrow \mu}/\Gamma_{W \rightarrow e}$ 0.996(10)
$ g_\tau/g_\mu $	$\Gamma_{\tau \rightarrow e}/\Gamma_{\mu \rightarrow e}$ 1.0011(15)	$\Gamma_{\tau \rightarrow \pi}/\Gamma_{\pi \rightarrow \mu}$ 0.9962(27)	$\Gamma_{\tau \rightarrow K}/\Gamma_{K \rightarrow \mu}$ 0.9858(70)		$\Gamma_{W \rightarrow \tau}/\Gamma_{W \rightarrow \mu}$ 1.034(13)
$ g_\tau/g_e $	$\Gamma_{\tau \rightarrow \mu}/\Gamma_{\mu \rightarrow e}$ 1.0030(15)				$\Gamma_{W \rightarrow \tau}/\Gamma_{W \rightarrow e}$ 1.031(13)

2024 [V.Cirigliano *et al.*, 2022] [HFLAV 2023 report] [PDG 2024]

$ g_\mu/g_e $	$\Gamma_{\tau \rightarrow \mu}/\Gamma_{\tau \rightarrow e}$ 1.002(11)	$\Gamma_{\pi \rightarrow \mu}/\Gamma_{\pi \rightarrow e}$ 1.0010(9)	$\Gamma_{K \rightarrow \mu}/\Gamma_{K \rightarrow e}$ 0.9978(18)	$\Gamma_{K \rightarrow \pi\mu}/\Gamma_{K \rightarrow \pi e}$ 1.0009(18)	$\Gamma_{W \rightarrow \mu}/\Gamma_{W \rightarrow e}$ 1.001(3)
$ g_\tau/g_\mu $	$\Gamma_{\tau \rightarrow e}/\Gamma_{\mu \rightarrow e}$ 1.0016(14)	$\Gamma_{\tau \rightarrow \pi}/\Gamma_{\pi \rightarrow \mu}$ 0.9958(38)	$\Gamma_{\tau \rightarrow K}/\Gamma_{K \rightarrow \mu}$ 0.9856(75)		$\Gamma_{W \rightarrow \tau}/\Gamma_{W \rightarrow \mu}$ 1.007(10)
$ g_\tau/g_e $	$\Gamma_{\tau \rightarrow \mu}/\Gamma_{\mu \rightarrow e}$ 1.0018(14)				$\Gamma_{W \rightarrow \tau}/\Gamma_{W \rightarrow e}$ 1.008(10)

Lepton universality tests with tau leptonic branching fractions

$$\Gamma[\mathcal{L} \rightarrow \nu_{\mathcal{L}} \ell \bar{\nu}_{\ell}(\gamma)] = \Gamma_{\mathcal{L}\ell} = \Gamma_{\mathcal{L}} \mathcal{B}_{\mathcal{L}\ell} = \frac{\mathcal{B}_{\mathcal{L}\ell}}{\tau_{\mathcal{L}}} = \frac{G_{\mathcal{L}} G_{\ell} m_{\mathcal{L}}^5}{192\pi^3} f^{\mathcal{L}\ell} (1 + \delta R_{\gamma}^{\mathcal{L}\ell}) (1 + \delta R_W^{\mathcal{L}\ell})$$

$$G_{\mathcal{L}} = \frac{g_{\mathcal{L}}^2}{4\sqrt{2}M_W^2}; \quad f^{\mathcal{L}\ell} = 1 - 8\rho + 8\rho^3 - \rho^4 - 12\rho^2 \ln \rho \quad \text{with} \quad \rho = \rho^{\mathcal{L}\ell} = m_{\ell}^2/m_{\mathcal{L}}^2$$

$$\delta R_{\gamma}^{\mathcal{L}\ell} = \frac{\alpha(m_{\mathcal{L}})}{2\pi} \left(\frac{25}{4} - \pi^2 \right) \quad \text{QED radiative correction, 1st order}$$

$$\delta R_W^{\mathcal{L}\ell} = \frac{3}{5} \frac{m_{\mathcal{L}}^2}{M_W^2} + \frac{9}{5} \frac{m_{\ell}^2}{M_W^2} \quad \text{EW radiative correction}$$

[Marciano, 1988], [Pich, Precision Tau Physics, 2014]

$$\delta R_{\gamma}^{\mathcal{L}\ell} = H_1^{\mathcal{L}\ell} \frac{\alpha(m_{\mathcal{L}})}{\pi} + H_2^{\mathcal{L}\ell} \frac{\alpha^2(m_{\mathcal{L}})}{\pi^3} + H_3^{\mathcal{L}\ell} \frac{\alpha^3(m_{\mathcal{L}})}{\pi^3} \quad \text{QED radiative correction, 3rd order}$$

$$H_1^{\mathcal{L}\ell} = \frac{25}{8} - \frac{\pi^2}{2} - \left(9 + 4\pi^2 + 12 \ln \rho \right) \rho + 16\pi^2 \rho^{\frac{3}{2}} \quad \text{with} \quad \rho = \rho^{\mathcal{L}\ell}$$

$$H_2^{\mathcal{L}\ell} = \frac{156815}{5184} - \frac{518}{81} \pi^2 - \frac{895}{36} \zeta(3) + \frac{67}{720} \pi^4 + \frac{53}{6} \pi^2 \ln 2 + H_{2,\text{had}}^{\mathcal{L}\ell} - \frac{5}{4} \pi^2 \sqrt{\rho} \quad \text{with} \quad \rho = \rho^{\mathcal{L}\ell}$$

$$H_{2,\text{had}}^{\mu e} = -0.042 \pm 0.002 \quad H_{2,\text{had}}^{\tau\mu} = ? \quad H_{2,\text{had}}^{\tau e} = ? \quad \text{missing values set to zero in the following}$$

$$H_3^{\mu e} = -15.3 \pm 2.3 \quad H_3^{\tau\mu} = ? \quad H_3^{\tau e} = ? \quad \text{missing values set to zero in the following}$$

[J.Erler & A.Freitas, PDG EW review, 2022]

[using modified minimal subtraction scheme ($\overline{\text{MS}}$)]

2nd order QED radiative corrections for leptonic-tau-decays LFU tests

motivation

- ▶ estimate uncertainty on QED radiative corrections, usually neglected
- ▶ check whether present uncertainties are fine for future prospects at Belle II and FCC-ee

$$\delta R_{\gamma}^{\mathcal{L}\ell} = H_1^{\mathcal{L}\ell} \frac{\alpha(m_{\mathcal{L}})}{\pi} + H_2^{\mathcal{L}\ell} \frac{\alpha^2(m_{\mathcal{L}})}{\pi^3} + H_3^{\mathcal{L}\ell} \frac{\alpha^3(m_{\mathcal{L}})}{\pi^3} \quad \text{QED radiative correction, 3rd order}$$

$$H_1^{\mathcal{L}\ell} = \frac{25}{8} - \frac{\pi^2}{2} - \left(9 + 4\pi^2 + 12 \ln \rho\right) \rho + 16\pi^2 \rho^{\frac{3}{2}} \quad \text{with} \quad \rho = \rho^{\mathcal{L}\ell}$$

$$H_2^{\mathcal{L}\ell} = \frac{156815}{5184} - \frac{518}{81}\pi^2 - \frac{895}{36}\zeta(3) + \frac{67}{720}\pi^4 + \frac{53}{6}\pi^2 \ln 2 + H_{2,\text{had}}^{\mathcal{L}\ell} - \frac{5}{4}\pi^2 \sqrt{\rho} \quad \text{with} \quad \rho = \rho^{\mathcal{L}\ell}$$

$$H_{2,\text{had}}^{\mu e} = -0.042 \pm 0.002 \quad H_{2,\text{had}}^{\tau} = ? \quad H_{2,\text{had}}^{\tau e} = ? \quad \text{missing values set to zero in the following}$$

$$H_3^{\mu e} = -15.3 \pm 2.3 \quad H_3^{\tau} = ? \quad H_3^{\tau e} = ? \quad \text{missing values set to zero in the following}$$

[J.Erler & A.Freitas, PDG EW review, 2022] [using modified minimal subtraction scheme ($\overline{\text{MS}}$)]

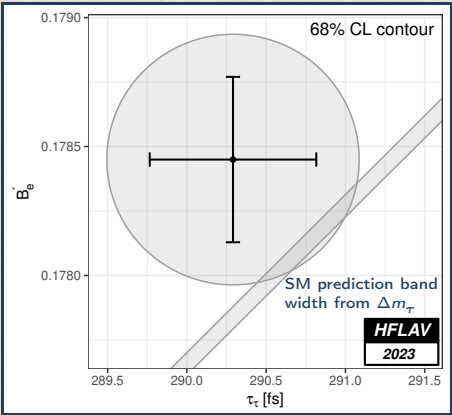
- ▶ for tau leptonic partial widths use up to 2nd order QED radiative correction formula without $H_{2,\text{had}}^{\tau(\mu,e)}$, $H_3^{\tau(\mu,e)}$
- ▶ tau LFU tests uncertainties up to 2.4 ppm if missing terms $H_{2,\text{had}}^{\tau(\mu,e)} \sim 20 \times H_{2,\text{had}}^{\mu e}$, $H_3^{\tau(\mu,e)} \sim 20 \times H_3^{\mu e}$
- ▶ tau mass uncertainty on phase space factors contributes uncertainty up to 1.3 ppm
- ▶ $\alpha(m_{\tau})$ computed using [F.Jegerlehner 2023 code alphaQEDr23.f], uncertainty up to 0.3 ppm on tau LFU tests

Lepton universality tests with tau leptonic branching fractions

	QED rad. corr.	
	1st order	2nd order
$\left(\frac{g_\tau}{g_\mu}\right) = \sqrt{\frac{\mathcal{B}_{\tau e} \tau_\mu m_\mu^5 f^{\mu e} (1 + \delta R_\gamma^{\mu e}) (1 + \delta R_W^{\mu e})}{\mathcal{B}_{\mu e} \tau_\tau m_\tau^5 f^{\tau e} (1 + \delta R_\gamma^{\tau e}) (1 + \delta R_W^{\tau e})}}$	1.0016 ± 0.0014	1.0016 ± 0.0014
$\left(\frac{g_\tau}{g_e}\right) = \sqrt{\frac{\mathcal{B}_{\tau\mu} \tau_\mu m_\mu^5 f^{\mu e} (1 + \delta R_\gamma^{\mu e}) (1 + \delta R_W^{\mu e})}{\mathcal{B}_{\mu e} \tau_\tau m_\tau^5 f^{\tau\mu} (1 + \delta R_\gamma^{\tau\mu}) (1 + \delta R_W^{\tau\mu})}}$	1.0018 ± 0.0014	1.0017 ± 0.0014
$\left(\frac{g_\mu}{g_e}\right) = \sqrt{\frac{\mathcal{B}_{\tau\mu} f^{\tau e} (1 + \delta R_\gamma^{\tau e}) (1 + \delta R_W^{\tau e})}{\mathcal{B}_{\tau e} f^{\tau\mu} (1 + \delta R_\gamma^{\tau\mu}) (1 + \delta R_W^{\tau\mu})}}$	1.0002 ± 0.0011	1.0001 ± 0.0011

- ▶ tau LFU tests limited to 1100–1400 ppm from branching fractions and tau lifetime measurements
- ▶ FCC-ee can improve by a factor of 10 (statistically 1000)
- ▶ expect sub-ppm uncertainty from further terms of EW radiative corrections
- ▶ desirable order-of-magnitude estimate of QED rad. corr. missing terms for tau LFU tests

Canonical tau lepton universality test plot



[HFLAV 2023 report]

$(g_\tau/g_{e\mu}) = 1.0017 \pm 0.0013$
 $[g_{e\mu} = g_e = g_\mu \text{ assuming } g_e = g_\mu]$

$\Delta(g_\tau/g_{e\mu})$ contributions		
input	Δ input	$\Delta(g_\tau/g_{e\mu})$
$B'_{\tau \rightarrow e}$	0.180%	0.090%
τ_τ	0.181%	0.090%
m_τ	0.005%	0.012%
total		0.128%

► $B'(\tau \rightarrow e\bar{\nu}\nu) = \text{average of } \begin{cases} B(\tau \rightarrow e\bar{\nu}\nu) \\ B(\tau \rightarrow \mu\bar{\nu}\nu) \cdot \frac{f^{\tau e} R_\gamma^{\tau e} R_W^{\tau e}}{f^{\tau \mu} R_\gamma^{\tau \mu} R_W^{\tau \mu}} \end{cases}$

► $\frac{B'(\tau \rightarrow e\bar{\nu}\nu)\tau_\mu}{B(\mu \rightarrow e\bar{\nu}\nu)\tau_\tau} = \frac{g_\tau^2}{g_{e\mu}^2} \frac{m_\tau^5 f^{\tau e} R_\gamma^{\tau e} R_W^{\tau e}}{m_\mu^5 f^{\mu e} R_\gamma^{\mu e} R_W^{\mu e}}$

► $\left(\frac{g_\tau}{g_{e\mu}}\right)^2 = \frac{B'(\tau \rightarrow e\bar{\nu}\nu)}{B(\mu \rightarrow e\bar{\nu}\nu)} \frac{\tau_\mu}{\tau_\tau} \frac{m_\mu^5}{m_\tau^5} \frac{f^{\mu e} R_\gamma^{\mu e} R_W^{\mu e}}{f^{\tau e} R_\gamma^{\tau e} R_W^{\tau e}}$

best measurements

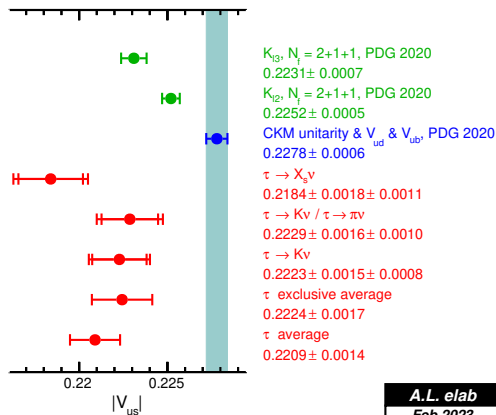
$B'_{\tau \rightarrow e}$	ALEPH
τ_τ	Belle
m_τ	Belle II

$|V_{us}|$ from tau measurements

Cabibbo angle anomaly

- ▶ 2018: CMK 1st row unitarity OK
- ▶ 2019: $\Delta_{\text{CKM}} > 3\sigma$ unitarity violation
 - ▶ dispersive calculation of Δ_R^V inner or universal electroweak radiative corrections (RC) to super-allowed nuclear beta decays
Seng, Gorchtein & Ramsey-Musolf, Phys. Rev. D 100, 013001 (2019)
 - ▶ $\sim 2\times$ more precise
 - ▶ significant shift
- ▶ 2020: revision & inflation of $|V_{ud}|$ th. syst. unc.
- ▶ 2020-2023: PDG reviews quote $\Delta_{\text{CKM}} \sim 3\sigma$
- ▶ 2023: short review [Cirigliano *et al.*, 2023]
- ▶ 2024: PDG review quotes $\Delta_{\text{CKM}} \sim 2\sigma$

$|V_{us}|$ after [Seng *et al.* 2019]



$|V_{us}|$ calculation with “inclusive” tau branching fractions

$|V_{us}|_{\tau\text{-OPE}}$ from $\mathcal{B}(\tau \rightarrow X_s \nu)$ and OPE

$$\triangleright |V_{us}|_{\tau\text{-OPE}} = \sqrt{R_s / \left[\frac{R_{ud}}{|V_{ud}|^2} - \delta R_{\tau, \text{SU3 breaking}} \right]}$$

 $\tau\text{-OPE}$

- $\delta R_{\tau, \text{SU3 breaking}}$ computed with Operator Product Expansion (OPE) perturbative QCD, + m_s W.A.
- E. Gamiz *et al.*, JHEP 01 (2003) 060, PRL 94 (2005) 011803, Nucl. Phys. Proc. Suppl. 169 (2007) 85, PoS KAON (2008) 008
- no strong isospin breaking correction

$|V_{us}|_{\tau\text{-latt}}$ from $\mathcal{B}(\tau \rightarrow X_s \nu)$ and lattice QCD

$$\triangleright |V_{us}|_{\tau\text{-latt}} = \sqrt{\left(\frac{|V_{us}|^2}{R_s} \right)_{\text{latt-incl}} \cdot R_s}$$

 $\tau\text{-latt}$

- $(|V_{us}|^2 / R_s)$ computed with lattice QCD [PRL 132 \(2024\) 26, 261901](#)
- no strong isospin breaking correction

$$\triangleright R_s = \mathcal{B}(\tau \rightarrow X_s \nu) / \mathcal{B}(\tau \rightarrow e \bar{\nu} \nu), \quad R_{ud} = \mathcal{B}(\tau \rightarrow X_{ud} \nu) / \mathcal{B}(\tau \rightarrow e \bar{\nu} \nu)$$

$|V_{us}|$ calculation with “exclusive” tau branching fractions

$$\Gamma(\tau \rightarrow \pi \bar{K} \nu_\tau [\gamma]) = \frac{G_F^2 m_\tau^5}{96 \pi^3} C_K^2 S_{EW}^\tau \left(|V_{us}| f^{K\pi} \right)^2 I_K^\tau \left(1 + \delta_{EM}^{K\tau} + \delta_{SU(2)}^{K\pi} \right)^2$$

 $\tau \rightarrow K \pi$

- C_K Clebsch-Gordan coefficients, $= 1$ for K^0 , $= 1/2$ for K^-
- I_K^ℓ = phase-space form factor integral
- Antonelli *et al.*, JHEP 1310 (2013) 070 (2013)
- (not updated for HFLAV 2023 report)

$$\frac{\mathcal{B}(\tau^- \rightarrow \mathcal{B}(\tau^- \rightarrow K^- \nu_\tau))}{\mathcal{B}(\tau^- \rightarrow \mathcal{B}(\tau^- \rightarrow \pi^- \nu_\tau))} = \left(\frac{f^{K\pm}}{f^{\pi\pm}} \right)^2 \frac{|V_{us}|_{\tau K/\pi}^2}{|V_{ud}|^2} \frac{(m_\tau^2 - m_K^2)^2}{(m_\tau^2 - m_\pi^2)^2} (1 + \delta R_{\tau K/\tau\pi})$$

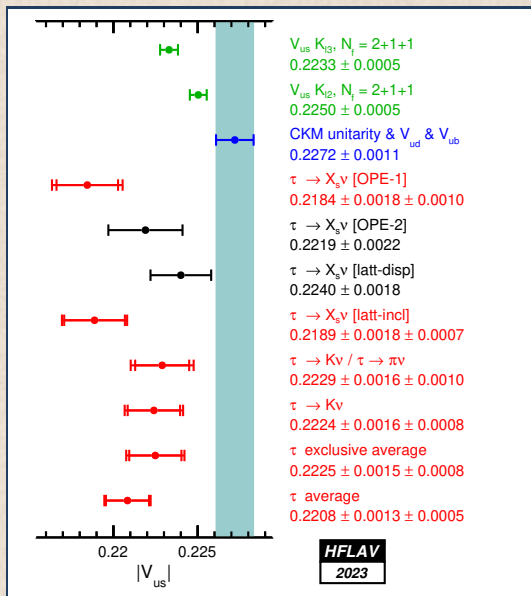
 $\tau \rightarrow K / \tau \rightarrow \pi$

- note that $\Delta \delta R_{\tau K/\tau\pi} = \Delta \delta R_{\tau\pi} \oplus \Delta \delta R_{\tau K}$ [PRD 104 9 (2021) L091502]

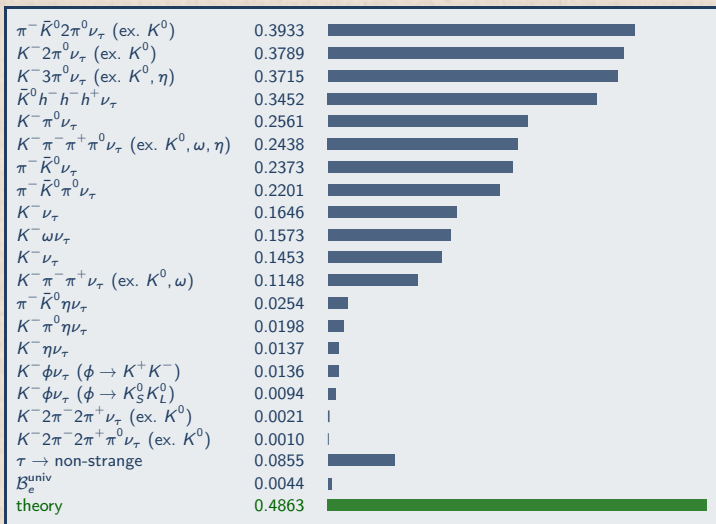
$$\mathcal{B}(\tau^- \rightarrow K^- \nu_\tau) = \frac{1}{16\pi} \left(\frac{G_F}{\hbar^3 c^3} \right)^2 |V_{us}|_{\tau K}^2 f_{K^\pm}^2 \frac{\tau_\tau}{\hbar} m_\tau^3 c^3 \left(1 - \frac{m_K^2}{m_\tau^2} \right)^2 S_{EW}^\tau (1 + \delta R_{\tau K})$$

 $\tau \rightarrow K$

$|V_{us}|$ from tau measurements [HFLAV 2023 report]



$|V_{us}|_{\tau\text{-OPE}}$ uncertainty budget [%] ($|V_{us}|_{\tau\text{-latt}}$ is comparable)



- ▶ Belle II could improve measurements on large multiplicity strang tau decays
- ▶ lattice QCD may plausibly reduce the theory uncertainty

Uncertainty budget of $|V_{us}|$ with exclusive tau measurements

input	Δinput [%]	$\Delta V_{us} _{\tau K/\pi}$ [%]	$\Delta V_{us} _{\tau K}$ [%]
$\mathcal{B}(\tau^- \rightarrow K^- \nu_\tau)$		0.69	0.69
$\mathcal{B}(\tau^- \rightarrow \pi^- \nu_\tau)$		0.24	
$f_{K^\pm}/f_{\pi^\pm}$		0.16	
$f_{K^\pm}$			0.19
S_{EW}^τ			0.02
$\delta R_{\tau K}$		0.28	0.29
$\delta R_{\tau\pi}$		0.28	
$ V_{ud} $		0.03	
tau lifetime			0.09

- ▶ Belle II might improve tau branching fractions
- ▶ if $\Delta\mathcal{B}(\tau^- \rightarrow \pi^- \nu_\tau)$ reduced by ≥ 3 , radiative corrections will be limiting
- ▶ lattice QCD may provide more precise radiative corrections in the future

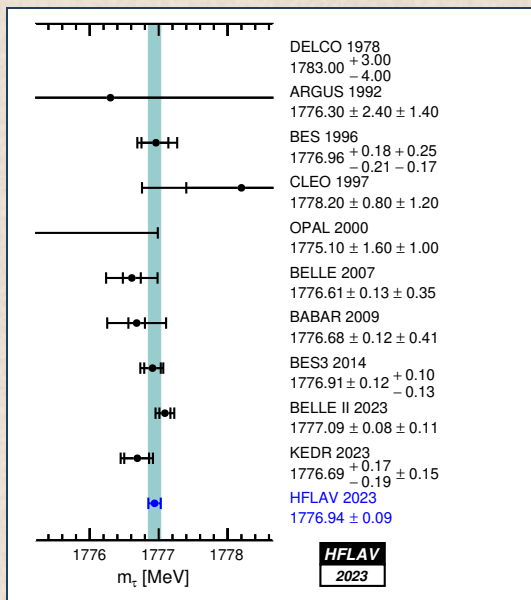
Summary

- ▶ improved HFLAV tau BR fit for HFLAV 2023 report (accepted by PRD, [arXiv:2411.18639](https://arxiv.org/abs/2411.18639) [[hep-ex](#)])
- ▶ HFLAV tau BR fit results used to compute tau-related Lepton Flavour Universality tests
 - ▶ LO $\rightarrow$ NLO QED radiative corrections (first presented here, not included in HFLAV 2023 report)
- ▶ HFLAV tau BR fit results used to compute $|V_{us}|$
 - ▶ added new $|V_{us}|$ calculation relying on recent lattice QCD calculation

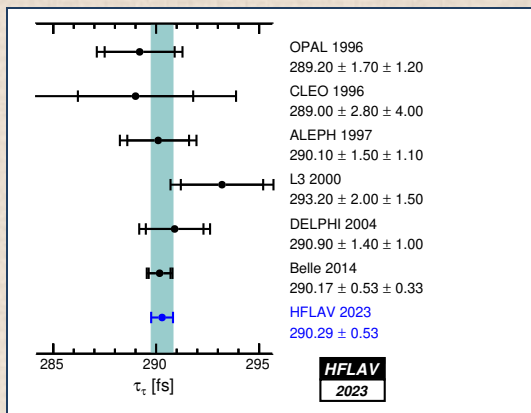
End

Backup slides

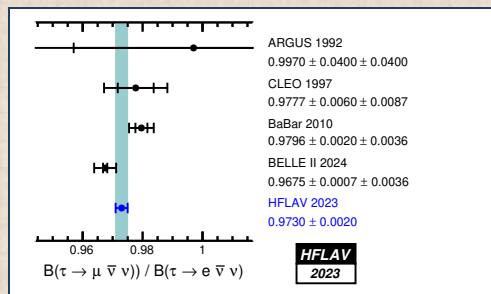
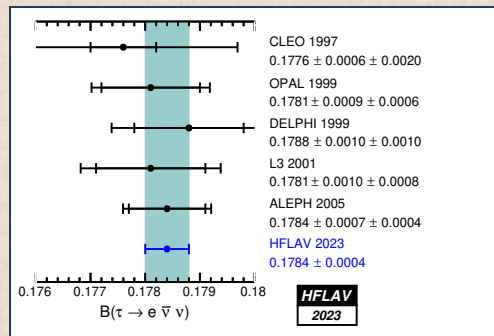
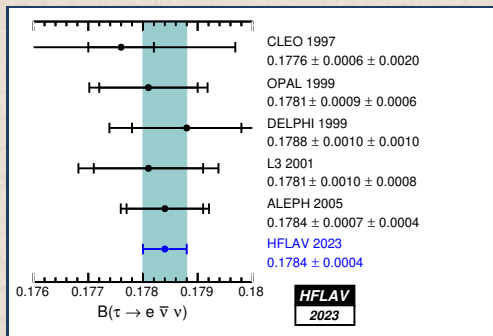
HFLAV Tau mass measurements



HFLAV Tau lifetime measurements



HFLAV Tau branching fractions measurements



$|V_{us}|$ determinations from kaons

$$\Gamma(K \rightarrow \pi \ell \bar{\nu}_\ell [\gamma]) = \frac{G_F^2 m_K^5}{192 \pi^3} C_K^2 S_{EW}^K \left(|V_{us}| f_+^{K\pi}(0) \right)^2 I_K^\ell \left(1 + \delta_{EM}^{K\ell} + \delta_{SU(2)}^{K\pi} \right)^2$$

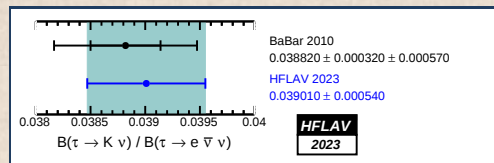
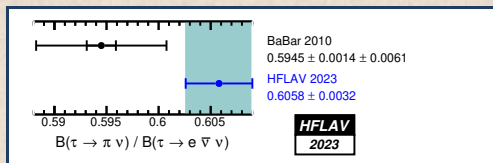
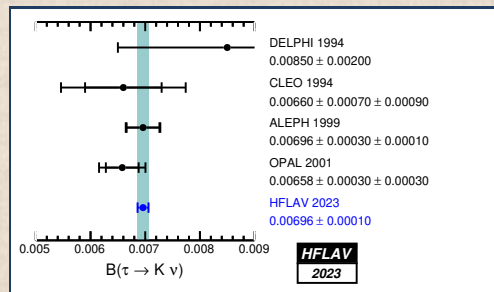
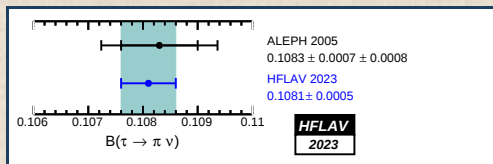
 $K_{\ell 3}$

- ▶ C_K Clebsch-Gordan coefficients, $= 1$ for K^0 , $= 1/2$ for K^-
- ▶ I_K^ℓ = phase-space form factor integral

$$\frac{\Gamma(K^- \rightarrow \ell^- \bar{\nu}_\ell)}{\Gamma(\pi^- \rightarrow \ell^- \bar{\nu}_\ell)} = \frac{|V_{us}|^2}{|V_{ud}|^2} \left(\frac{f^{K^\pm}}{f^{\pi^\pm}} \right)^2 \frac{m_K (1 - m_\ell^2/m_K^2)^2}{m_\pi (1 - m_\ell^2/m_\pi^2)^2} (1 + \delta_{EM})$$

 $K_{\ell 2}$

$\mathcal{B}(\tau^- \rightarrow \pi^- \nu_\tau)$ & $\mathcal{B}(\tau^- \rightarrow K^- \nu_\tau)$ measurements



μ/e universality from $\mathcal{B}(\pi \rightarrow e\nu(\gamma))/\mathcal{B}(\pi \rightarrow \mu\nu(\gamma))$

- ▶ $R_{e/\mu}^{\pi, \text{SM}} = 1.23524(15) \cdot 10^{-4}$
- ▶ $R_{e/\mu}^{\pi, \text{PIENU(2015)}} = 1.2344(23)_{\text{stat}}(19)_{\text{syst}} \cdot 10^{-4}$ PRL 115 (2015) 071801
- ▶ $R_{e/\mu}^{\pi, \text{EXP}} = 1.2327(23) \cdot 10^{-4}$ PDG 2024
- ▶ on-going PEN experiment at PSI also aims to measure $R_{e/\mu}^{\pi}$
- ▶ future: PIONEER (PSI) uncertainty goal equal to SM prediction precision