

The role of Flavour in global SMEFT fits





Universidad de <mark>Granada</mark>

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Based on arXiv: 2507.06191 [hep-ph], in collaboration with

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Disclaimer(s)

- I am not a Flavour expert (not even close to...)
- This talk could be seen as the second part of the talk yesterday in the Joint "Top and Electroweak physics" +
 - There will be some overlap to set the stage and introduce the tools but different focus on the discussion
- With few exceptions^{*}, flavour is one of the sectors that receives less attention in global SMEFT combinations of different types of data sets
 - Not because it is not relevant but because it difficult to treat!
 - As long as one departs from trivial flavour assumptions and/or goes beyond leading-order studies, flavour measurements enter the game
- The purpose of this talk is to illustrate this, with a global study combining EW+Higgs+Top+Flavour in the SMEFT, floating simultaneously all the SMEFT parameters and the SM parameters (+ including RGE effects, matching SMEFT/LEFT,...)

*See e.g. R. Aoude et al., JHEP 12 (2020) 113, R. Bartocci et al. JHEP 05 (2024) 074, S. Bruggisser et al. JHEP 02 (2023) 225, L. Allwicher et al. JHEP 03 (2024) 049, C. Grunwald et al., JHEP 11 (2023) 110









SMEFT Lagrangian:

$$\mathcal{L}_{ ext{Eff}} = \sum_{d=4}^{\infty} rac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{ ext{SM}} + rac{1}{\Lambda} \mathcal{L}_5 + rac{1}{\Lambda^2} \mathcal{L}_6 + \cdots \ \mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \qquad [\mathcal{O}_i] = d \longrightarrow \left(rac{q}{\Lambda}
ight)^{d-4} \ q = v, E < \Lambda \qquad \mathcal{L}_w(?)$$

$$\begin{array}{c} \begin{array}{c} \begin{array}{c} 4 \ \text{Fermion} \\ \end{array} \\ & \mathcal{O}_{lq}^{(1)} \ \mathcal{O}_{lq}^{(3)} \\ & \mathcal{O}_{lq}^{(1)} \ \mathcal{O}_{lq}^{(3)} \\ & \mathcal{O}_{lu} \ \mathcal{O}_{qe} \ \mathcal{O}_{eu} \\ & \mathcal{O}_{lu} \ \mathcal{O}_{qe} \ \mathcal{O}_{eu} \\ & \mathcal{O}_{ud} \ \mathcal{O}_{ud}^{(1)} \\ & \mathcal{O}_{le} \\ & \mathcal{O}_{le} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{ll} \ \mathcal{O}_{ee} \\ & \mathcal{O}_{ld} \ \mathcal{O}_{ed} \\ & \mathcal{O}_{lequ} \ \mathcal{O}_{qu}^{(1)} \ \mathcal{O}_{qu}^{(1)} \\ & \mathcal{O}_{qu}^{(1)} \ \mathcal{O}_{qd}^{(1)} \\ & \mathcal{O}_{lequ}^{(1)} \ \mathcal{O}_{lequ}^{(3)} \\ & \mathcal{O}_{ledq} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{ledq} \\ & \mathcal{O}_{quqd}^{(1)} \mathcal{O}_{quqd}^{(8)} \\ & \mathcal{O}_{quqd} \\ \end{array} \\ \end{array}$$

With some minimal assumptions about the UV, the IR effects of new physics can be parameterized via the

What is not Flavour in the SMEFT? To dimension 6 (2499 operators assuming B and L conservation)

2F-Scalar ${\cal O}_{\phi l}^{(3)}$ **1**) **5**1 $\mathcal{O}_{u\phi}$ $\mathcal{O}_{e\phi}$ \mathcal{O}_{ϕ} $\mathcal{O}_{\phi e}$ ${\cal O}_{d\phi}$ $\mathcal{O}_{\phi \Box}$ $:\mathcal{O}_{\phi D}$ Bosonic 2F-Dipoles **F-Vector** ${\cal O}_{\phi ilde B} \cdot$ ${\cal O}_{\phi B}$ $\mathcal{O}_{eB} \,\, \mathcal{O}_{eW}$ $\begin{array}{c} \mathcal{O}_{u} & \mathcal{O}_{\phi q}^{(3)} & \mathcal{O}_{uB} & \mathcal{O}_{uW} & \mathcal{O}_{uG} \\ \phi u & \mathcal{O}_{\phi d} & \mathcal{O}_{dB} & \mathcal{O}_{dW} & \mathcal{O}_{dG} \end{array}$ $\mathcal{O}_{\phi W}$ $\phi ilde W$. $\mathcal{O}_{\phi WB} \stackrel{\mathsf{CP}}{\longrightarrow} \mathcal{O}_{\phi \tilde{W}B} \stackrel{\varphi}{\longrightarrow} \mathcal{O}_{\phi \tilde{G}} \stackrel{\varphi$ ${\cal O}_{\phi ud}$ The role of Flavour in Global SMEFT fits

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SMEFT Lagrangian:

$$\mathcal{L}_{ ext{Eff}} = \sum_{d=4}^{\infty} rac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{ ext{SM}} + rac{1}{\Lambda} \mathcal{L}_5 + rac{1}{\Lambda^2} \mathcal{L}_6 + \cdots + \mathcal{L}_{ ext{UV}}(2)$$
 $\mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \qquad [\mathcal{O}_i] = d \longrightarrow \left(rac{q}{\Lambda}
ight)^{d-4} = v, E < \Lambda$



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$$U(2)^5 = U(2)_{q_L} imes U(2)_{u_R} imes U(2)_{d_R} imes U(2)_{l_L} imes U(2)_{e_R}$$

Even under these assumptions, Flavour can play an important role in SMEFT analysis and must be carefully implemented in fitting tools



Common choices motivated by "realistic" BSM scenarios

Assume New Physics is flavour blind and respects a U(3)⁵ flavour symmetry Assume New Physics respects the approximate U(2) quark flavour symmetries of the SM \Rightarrow No new sources of flavour mixing but separate 3rd and light generations

	· · · 2F-Scalar	$\mathcal{O}_{\mathcal{O}} = \mathcal{O}_{\mathcal{O}}$
s!	$\mathcal{O}_{u\phi}$	$\mathcal{O}_{W} \stackrel{\mathcal{O}_{\tilde{G}}}{\longleftarrow} \mathcal{O}_{\tilde{W}} \stackrel{\mathcal{O}_{\tilde{G}}}{\longrightarrow} \mathcal{O}_{\tilde{W}} \stackrel{\mathcal{O}_{\tilde{W}}}{\longrightarrow} \mathcal{O}_{\tilde$
changing	${\cal O}_{e\phi} {egin{array}{c} {\cal O}_{u\phi} \ {\cal O}_{d\phi} \end{array}}$	\mathcal{O}_{ϕ} \mathbf{CP} \mathcal{O}_{ϕ}
onstraints ptions to bounds	$\mathcal{O}_{eB} \mathcal{O}_{eW}$ $\mathcal{O}_{uB} \mathcal{O}_{uW} \mathcal{O}_{uG}$	$\begin{array}{c c} \mathcal{O}_{\phi D} \\ \mathcal{O}_{\phi B} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{\phi B} \\ \mathcal{O}_{\phi \tilde{W}} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{\phi \tilde{W}} \\ \mathcal{O}_{\phi \tilde{W}} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{\phi \tilde{W}} \\ \mathcal{O}_{\phi \tilde{W}} \\ \end{array} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{\phi \tilde{W}} \\ \mathcal{O}_{\phi \tilde{W}B} \\ \end{array} \\ \begin{array}{c} \mathcal{O}_{\phi \tilde{W}} \\ \mathcal{O}_{\phi \tilde{W}B} \\ \end{array} \\ \end{array}$
$\mathcal{O}_{\phi ud} = \mathcal{O}_{\phi d}$	${\cal O}_{dB} {\cal O}_{dW} {\cal O}_{dG}$	$\mathcal{O}_{\phi G} \mathcal{O}_{\phi \tilde{G}}$
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The Fitting Framework







under GPL on GitHub)

https://github.com/silvest/HEPfit

Main Reference: JB et al., <u>Eur. Phys. J. C (2020) 80:456</u>, <u>arXiv: 1910.14012 [hep-ph]</u>



Original code already containing a base SMEFT class with a setup for EW/Higgs LO studies

➡ Massive upgrades in the work presented here

General High Energy Physics fitting tool to combine indirect and direct searches of new physics (available)







- The **SMEFT class** in **HEP**fit :
- Implementation of full dimension-6 SMEFT basis:
 - Warsaw basis: All 2499 operators
 - Restrictions assuming different flavour assumptions available
 - U(3)⁵ flavour symmetry
 - U(2)⁵ flavour symmetry: both in the "UP" and "DOWN" bases
- Calculations in both " α " and "M_W" scheme for most observables
- RGE evolution included via **RGESolver** S. Di Noi, L. Silvestrini, Eur. Phys.J.C 83 (2023) 3, 200
 - Multiple possibilities: Exact integration / Matrix Evolution (much faster)
 - Possibility of RGE to multiple scales
 - **Careful:** RGE available only at LO (1-loop). Running between similar scales < TH unc.
- NLO SMEFT finite terms available for several of the most precise observables

Careful: Consistent NLO study requires 2-loop RGE. Not available in literature (yet)





















- U(2)⁵ in the chosen basis $C_i(\Lambda)$ $Y_q^{ij}(\Lambda)$

















RGE

SMEFT

2499



The Global Fit Setup **Combining EW/Higgs/Top/Flavour**





















The Global SMEFT fit <u>SMEET framowork</u> SMEFT $O_{\rm SM}$ State-of-the-art for SM predictions in most precise observables Higgs (EWPO, Flavour) SM parameters (EW and flavour) floated in the fit, together with all the Wilson Coefficients $\Lambda_h(B)$ Jorge de Blas - U. of Granada

Technical details of the fits and comparisons Jorge de Blas - U. of Granada June 23, 2025





Electroweak Observables

- **Electroweak Precision Observables:**
 - \blacktriangleright Z-pole (LEP/SLD): Γ_Z , A_f , A_{FB}^f , R_f , ...
 - W properties (LEP2/Tevatron/LHC):
 - Higgs and Top properties (Tevatron/LHC): $M_H,\ m_t$
 - Tests of lepton universality from Tevatron & LHC
- LEP2 observables
 - Di-Boson: $e^+e^-
 ightarrow W^+W^-$ Berthier et al., 1606.06693 [hep-ph]
 - $e^+e^- \rightarrow f\bar{f}$: leptonic cross sections and asymmetries, hadronic cross section
- Drell-Yan at LHC: $pp
 ightarrow \ell^+ \ell^-, \ \ell
 u$
 - **Differential distributions**
 - Implemented from HighPT code

L. Allwicher et al., 2207.10756, 2207.10714 [hep-ph]







Higgs Boson Observables

- ATLAS+CMS 8 TeV combination for single strengths: $\mu_{ij} = \frac{\sigma_i \times BR_j}{(\sigma_i \times BR_j)_{\rm SM}}$
- ATLAS and CMS I3 TeV results (I39 fb⁻¹)
 - STXS Stage 1.2 binning
- Including full information on all available channels (production and decay)









		Top Obse	rvable	S
	Process	Observable	\sqrt{s}	$\int \mathcal{L}$
	$p\bar{p} \rightarrow t\bar{t}$	${\rm d}A_{\rm FB}^{t\bar{t}}/{\rm d}m_{t\bar{t}}$	1.96 TeV	$9.7 { m ~fb^{-1}}$
nology Constraint	$pp \rightarrow t\bar{t}$	$\sigma_{t\bar{t}}^{13\text{TeV}}/\sigma_{t\bar{t}}^{8\text{TeV}} \\ \sigma_{t\bar{t}}^{8\text{TeV}}/\sigma_{t\bar{t}}^{7\text{TeV}} \\ \sigma_{t\bar{t}} \\ \sigma_{t\bar{t}} \\ d\sigma_{t\bar{t}}/dm_{t\bar{t}} \\ (d\sigma_{t\bar{t}}/dm_{t\bar{t}})/\sigma_{t\bar{t}} \\ dA_{C}/dm_{t\bar{t}} \\ \end{array}$	13 & 8 TeV 8 & 7 TeV 13 TeV 13 TeV 13 TeV 13 TeV 13 TeV	$\begin{array}{c} 20 \& 36 \ \mathrm{fb}^{-1} \\ 20 \& 5 \ \mathrm{fb}^{-1} \\ 36/139 \ \mathrm{fb}^{-1} \\ 36 \ \mathrm{fb}^{-1} \\ 36/137 \ \mathrm{fb}^{-1} \\ 140 \ \mathrm{fb}^{-1} \end{array}$
	$pp \to t\bar{t}Z$	$d\sigma/dp_T^Z$	$13 { m TeV}$	$77.5/140 \text{ fb}^{-1}$
	$pp \to t\bar{t}\gamma$	$d\sigma/dp_T^\gamma$	$13 { m TeV}$	$140 {\rm ~fb^{-1}}$
Тор	$pp \to t\bar{t}W$	$\sigma_{ttW^{\pm}} \ \sigma_{ttW^{+}}/\sigma_{ttW^{-}}$	$13 { m ~TeV}$	$140 {\rm ~fb^{-1}}$
EW	$t \to Wb$	F_0, F_L	8 TeV 13 TeV	20 fb^{-1} 140 fb^{-1}
$\mathcal{L}_{\mathrm{EFT}}$	$pp \to tW$	σ	7 TeV 8 TeV 13 TeV	$\begin{array}{c} 4.6 \ \& \ 1.5 \ \mathrm{fb}^{-1} \\ 20 \ \mathrm{fb}^{-1} \\ 3.2/140 \ \mathrm{fb}^{-1} \end{array}$
	$pp \to t\bar{b} \text{ (s-ch)}$	σ	$8 { m TeV}$ 13 TeV	20 fb^{-1} 140 fb ⁻¹
ur	$pp \rightarrow tq \text{ (t-ch)}$	σ	7 TeV 8 TeV 13 TeV	$\begin{array}{r} 4.6 \ \& \ 1.5 \ \mathrm{fb^{-1}} \\ 20 \ \mathrm{fb^{-1}} \\ 36/140 \ \mathrm{fb^{-1}} \end{array}$
	$pp \rightarrow t\gamma q$	σ	$13 { m TeV}$	$140/36 \text{ fb}^{-1}$
	$pp \rightarrow tZq$	σ	$13 { m TeV}$	$140 {\rm ~fb^{-1}}$
	$pp \to t\bar{t}b\bar{b}$	σ	$13 { m TeV}$	$36 {\rm ~fb^{-1}}$
	$pp \rightarrow t\bar{t}t\bar{t}$	σ	$13 { m TeV}$	$140 {\rm ~fb^{-1}}$







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		Parameter	Value
ор) \		F_{B_s} (GeV)	0.2301 ± 0.0012
. /		$\begin{vmatrix} I B_s & (GCV) \\ F_{B_s} / F_{B_d} \end{vmatrix}$	$\begin{array}{c} 0.2301 \pm 0.0012 \\ 1.208 \pm 0.005 \end{array}$
		$\begin{vmatrix} I B_s / I B_d \\ B_{B_s} (4.2 \text{GeV}) \end{vmatrix}$	0.888 ± 0.040
		$\begin{vmatrix} B_{B_s} (B_{2} \oplus V) \\ B_{B_s} / B_{B_d} \end{vmatrix}$	1.015 ± 0.021
	Computed in	$B_{B_{s},4}(4.2{\rm GeV})$	0.98 ± 0.08
	Lattice QCD	$B_{B_s,5}(4.2 {\rm GeV})$	1.66 ± 0.13
d P		$B_{B_{d},4}(4.2{\rm GeV})$	0.99 ± 0.08
		$B_{B_{d},5}(4.2{\rm GeV})$	1.58 ± 0.18
		$B_K(2{ m GeV})$	0.552 ± 0.012
		$B_{K,4}(2{ m GeV})$	0.904 ± 0.053
		$B_{K,5}(2{ m GeV})$	0.618 ± 0.114
		$\phi_{\varepsilon_K}(^{\circ})$	43.51 ± 0.05
		$\overline{\kappa}_{\varepsilon_K}$	0.97 ± 0.02
	Negligible effect	$(\Delta M_K)^{\rm SM} \ (\rm ns^{-1})$	8.8 ± 3.6
LEF	expected from	$B_{D,1}(3\mathrm{GeV})$	0.765 ± 0.025
	-	$B_{D,4}(3\mathrm{GeV})$	0.98 ± 0.06
	SMEFT	$B_{D,5}(3{ m GeV})$	1.05 ± 0.09
Flavour	(within current	$F_D (GeV)$	0.2120 ± 0.0007
	uncertainties)	$\left (\Delta M_D)^{\rm SM} (\rm ps^{-1}) \right $	0.005 ± 0.005
		$f_K (GeV)$	0.15611 ± 0.00021
		$\int f_K / f_{\pi}$	1.1966 ± 0.0018
		$\delta R_{\pi}^{\rm phys}$	0.0153 ± 0.0019
		$\delta P_{c,u}$	0.04 ± 0.02









The role of Flavour in combined EW+Higgs+Top+Flavour studies







Bo nds controlled mystly by EW/Higgs observables



Not surprisingly, in the U(3)⁵ -symmetric limit contributions to flavour observables are very suppressed

Picture changes dramatically when we relax the hypotheses to U(2)5











SMEFT fit results: U(3)⁵

Global fit: comparison of different choices of Λ



Strong correlations between coefficients significantly relax the bounds but many operators can still be constrained within the perturbative regime (especially for Λ =3 TeV)





The role of Flavour in SMEFT fits: U(2)⁵







The role of Flavour in SMEFT fits: U(2)⁵

Full 10 TeV



(Very small contributions from RH operators to Flavour either in the UP or DOWN bases)







The role of Flavour in SMEFT fits: U(2)⁵

Full 10 TeV Full 3 TeV



sets results: Impact of data ndividual













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Strong dependence on choice of "basis"

Alignment of New Physics with the DOWN sector in the U(2)⁵ limit \Rightarrow Strong suppresion to flavour-changing b-quark processes \Rightarrow Systematic relaxation of flavour bounds (Still providing the leading constraint in dipole interactions)





















SMEFT fit results: U(2)⁵









SMEFT fit results: U(2)⁵





- Despite the very different hypotheses and prominent role of flavour measurements in the U(2)⁵ fit, these assumptions still provide a good flavour "protection"
 - $C_{\varphi q}^{(1)}$, $C_{\varphi q}^{(3)}$ \rightarrow Modify NC quark interactions **Comparable results in U(3)**⁵ and U(2)⁵ limits (both in individual and global fits) **Controlled by EW/Top observables**
 - Relevant for building new physics models not (too) far from the EW scale!











Summary and Conclusions

- SMEFT:

 - - U(3)⁵ flavour symmetry (41 operators)
 - U(2)⁵ flavour symmetry (124 operators)
 - \checkmark Including prior information to ensure the EFT is studied within its perturbative regime
- The role of flavour in the global fit
 - \checkmark Strong dependence of the result on flavour assumptions
 - new interactions...
 - symmetry
 - \checkmark U(2)⁵ symmetry provides a good flavour protection \Rightarrow Important for EW scale BSM models!

In this study, we have presented a consistent combination of EW/Higgs/Top/Flavour constraints in the dimension-6

✓ Including variations of the SMEFT Wilson coefficients and all the SM parameters (inputs + TH uncert.)

✓ Including RGE evolution both in the SMEFT and LEFT starting from a full basis of SMEFT effects in the UV:

Around 200 parameters in the fit !

 \checkmark U(2)⁵ : beyond the trivial U(3)⁵ limit, flavour measurements play a leading role in setting very strong limits in

...though results depend on the direction of the 3rd family in flavour space chosen to define the $U(2)^{5}$



